

# Prompt radio-emitting outflows in tidal disruption events from accretion disk winds

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# Tidal Disruption Events

When a supermassive black hole destroys a star

Jet or outflow  
(radio)

Synchrotron emission from the outflow  
encountering the circumnuclear medium

Unbound debris  
stream (radio?)

Supermassive black  
hole (X-ray)

Accretion disk (X-  
ray/optical)

Bound stellar debris  
(optical)

Outflow ejection mechanism has been hotly  
debated

- Debris related
  - Unbound debris stream (e.g. Krolik+2016)
  - Collision induced outflows (e.g. Lu&Bonnerot2020)
- Accretion-induced
  - Jets (on or off axis) (e.g. van Velzen+2016)
  - Spherical winds (e.g. Alexander+2016)

# Tidal Disruption Events

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Unbound debris

## Outflow ejection mechanism is hotly debated

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  - Collision induced outflows
- Accretion-induced
  - Jets
  - Spherical winds

Not all TDEs are detected in the radio. Approx. 50% of TDEs are radio-detected

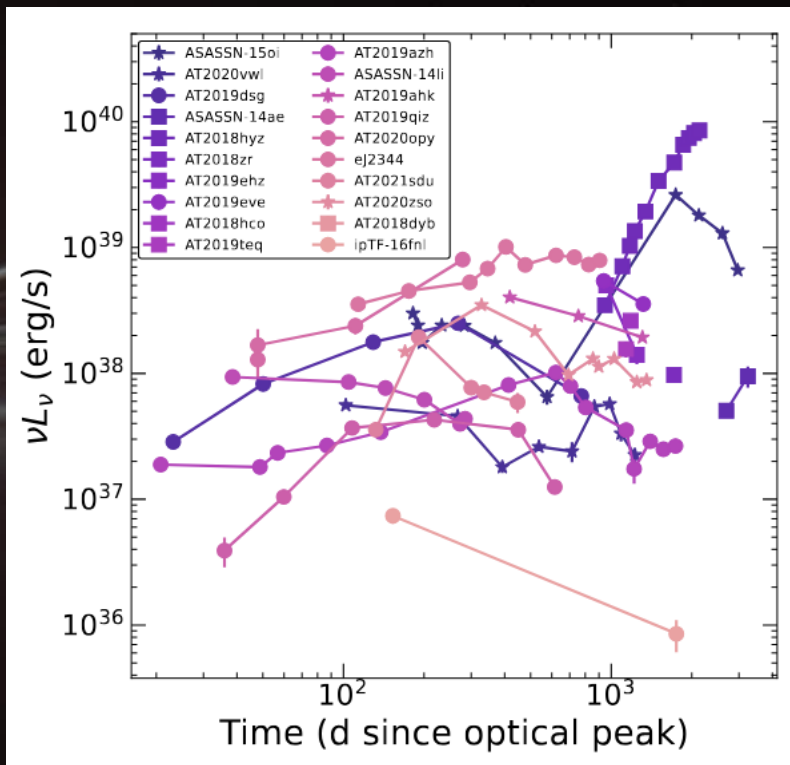
Open questions:

- What produces the radio emission in TDEs?
  - What produces the late-time flares in TDEs?
- How do SMBHs launch jets/outflows?

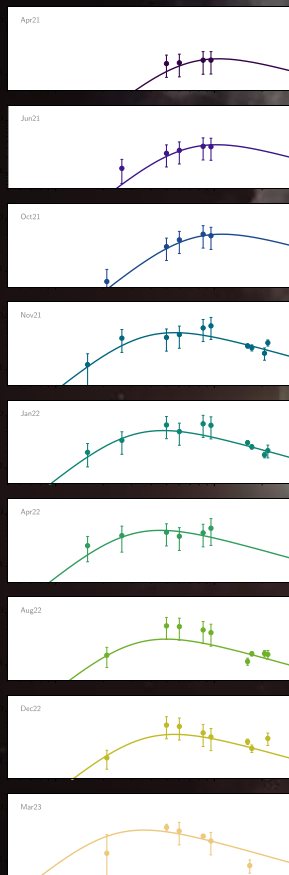
# Radio emission from TDEs

Extremely diverse properties (as we heard from Kate!):

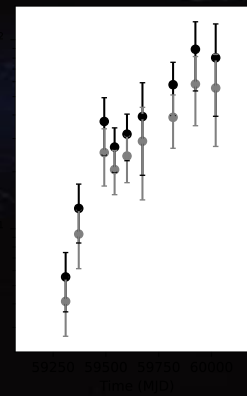
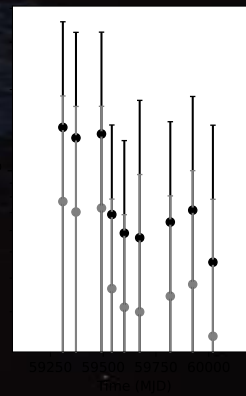
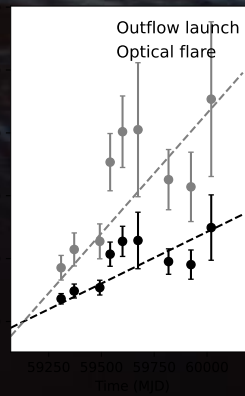
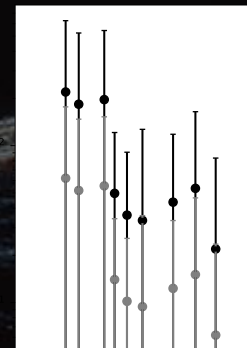
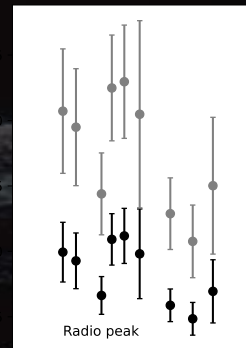
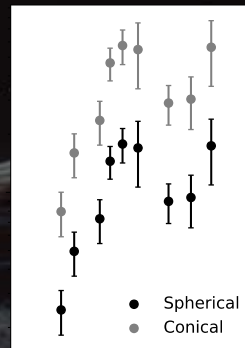
- Double flares
- Rising radio emission >1000s days
- Very early radio emission in some



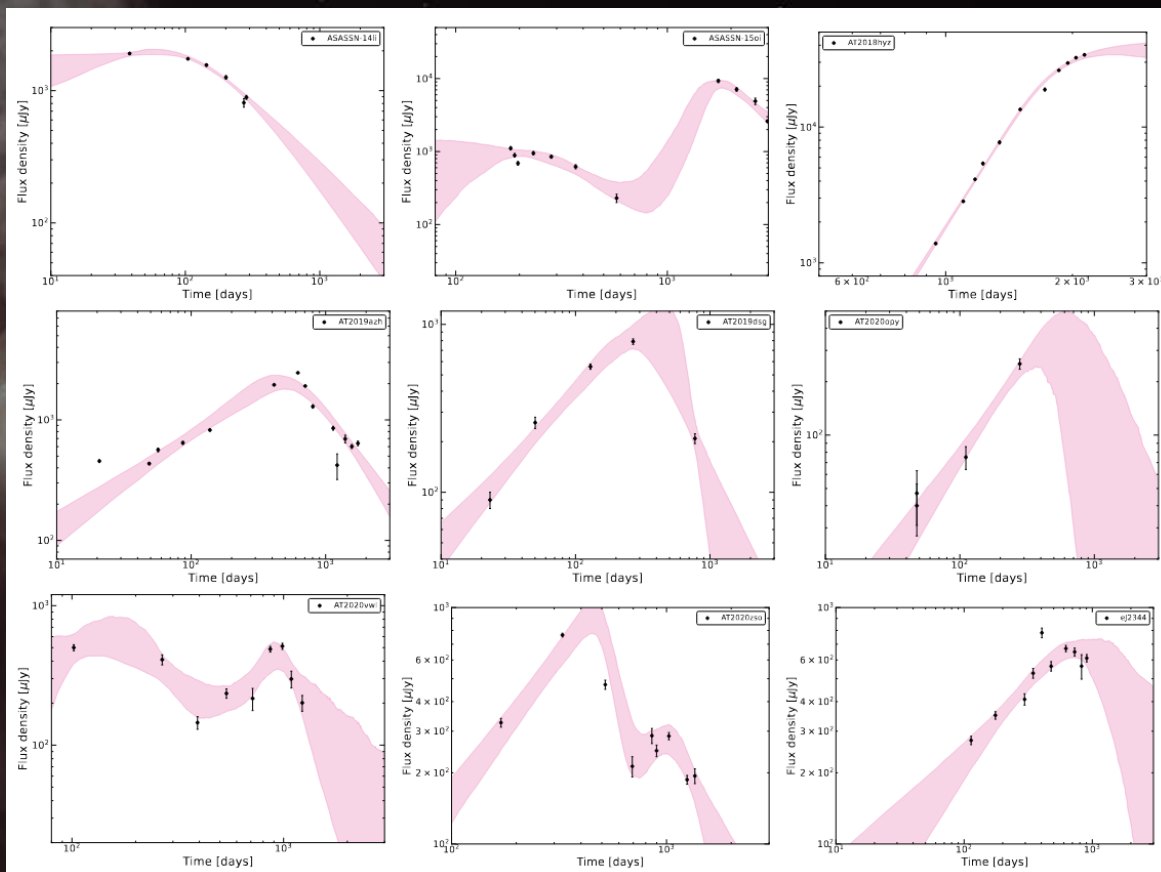
# Broadband spectral monitoring constrains physical outflow properties



ATCA observations of eJ2344 (Goodwin+2024)



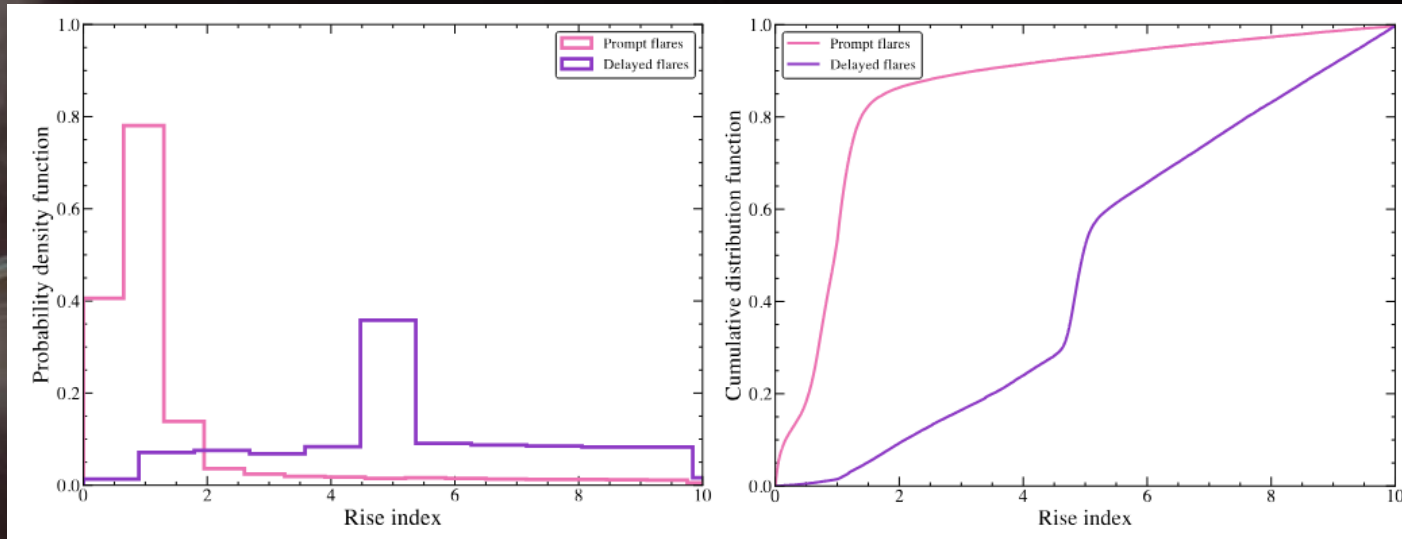
# Lightcurve fitting reveals very different behaviour for prompt vs delayed flares



# Lightcurve fitting reveals very different behaviour

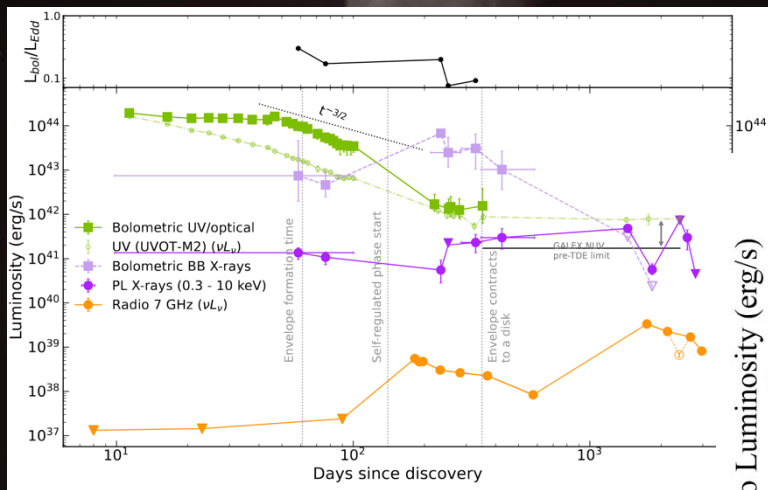
Prompt radio emission  
rises with index 1-2

Delayed flares typically  
rise much more steeply

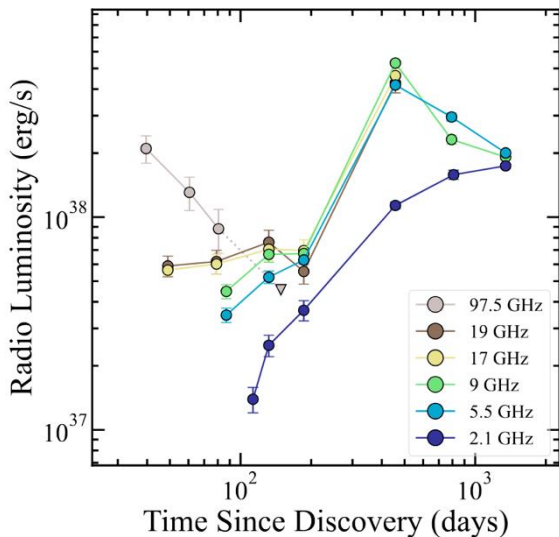


# The case for multiple ejection mechanisms

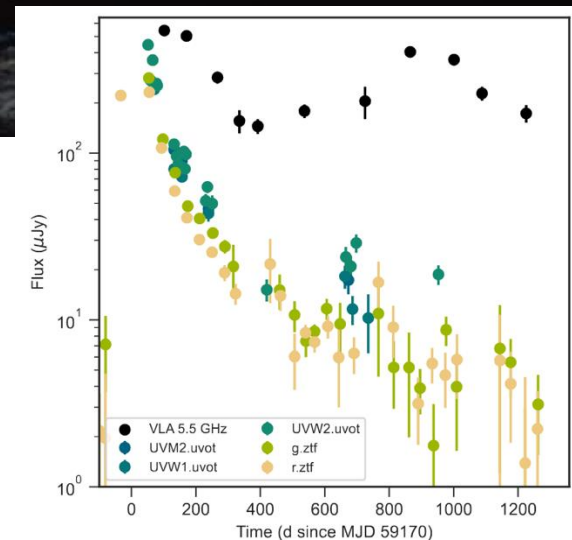
ASASSN-15oi (Hajela+2025)



ASASSN-19bt  
(Christy+2024)

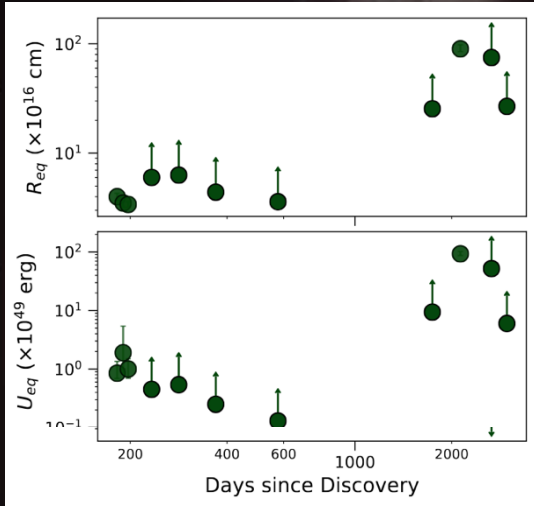


AT2020vwl  
(Goodwin+2025)

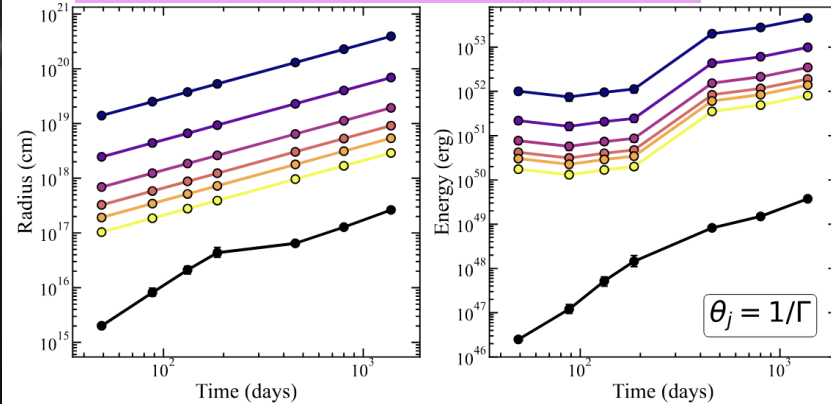


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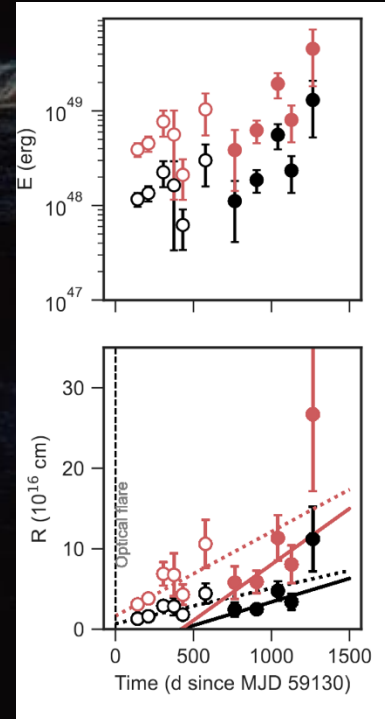
ASASSN-15oi (Hajela+2025)



➤ Second flare associated with energy injection



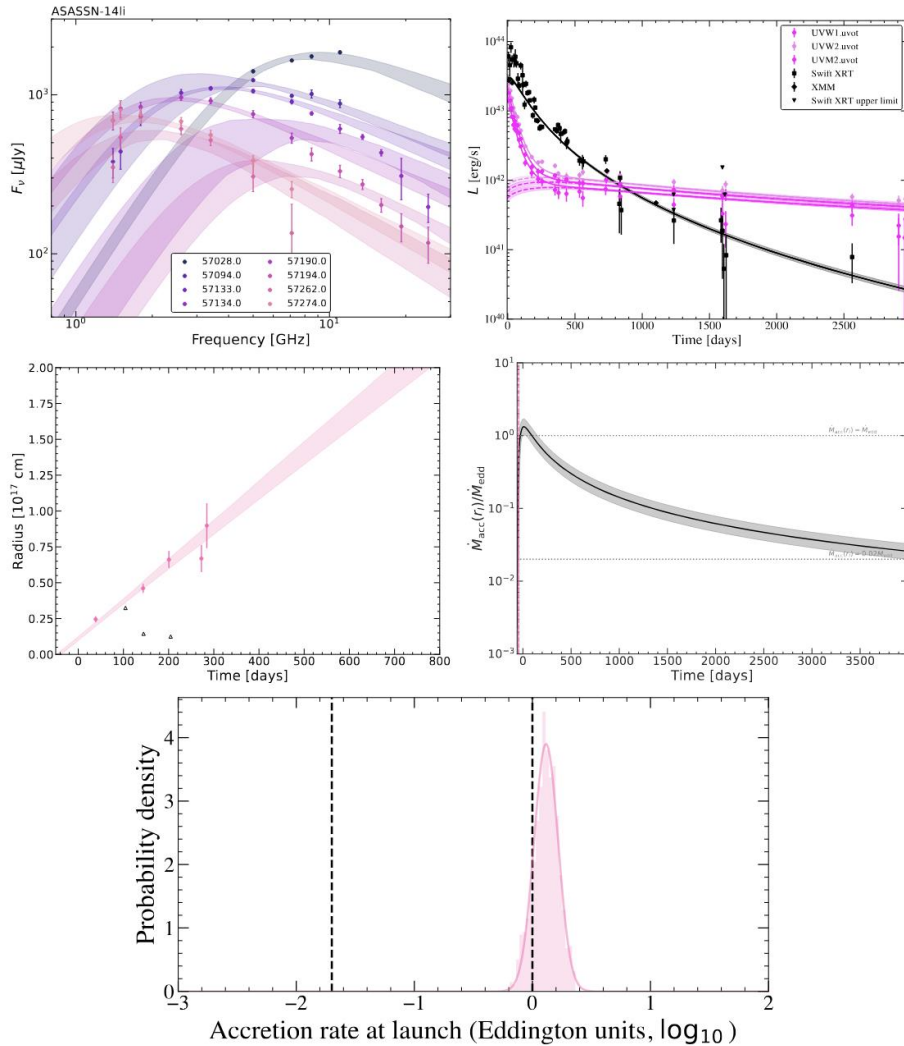
AT2020vwl (Goodwin+2025)





**Is there a link between the outflow launch time and the state of the accretion disk at this time?**

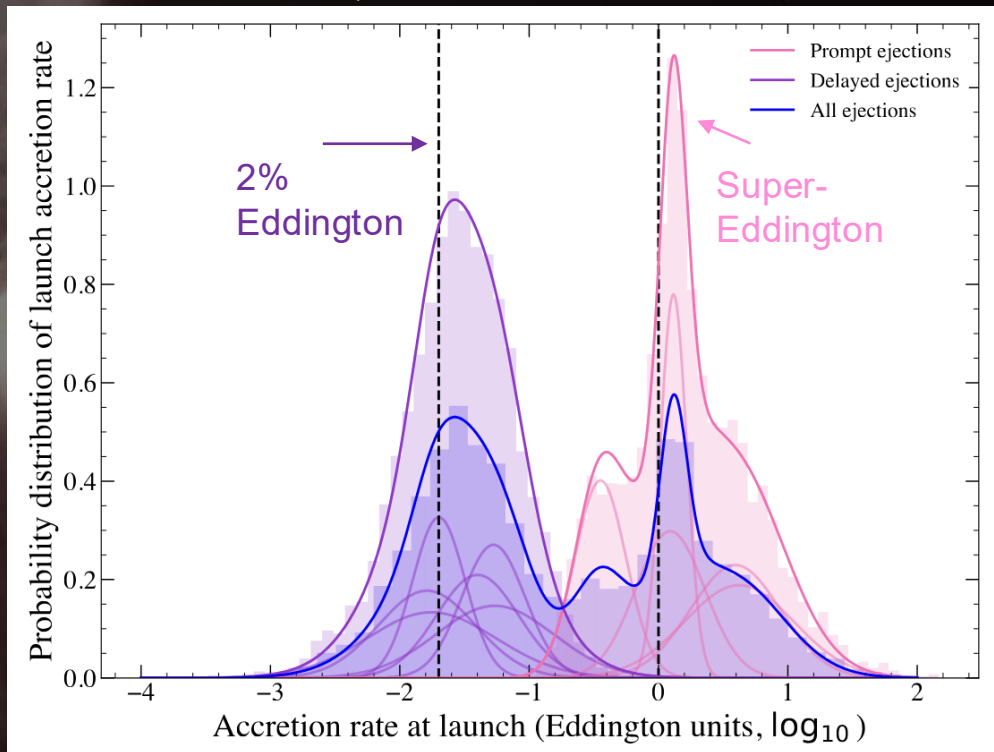
# Accretion disk / outflow connections



Model a sample of TDEs to extract the disk accretion rate at the time of outflow launch

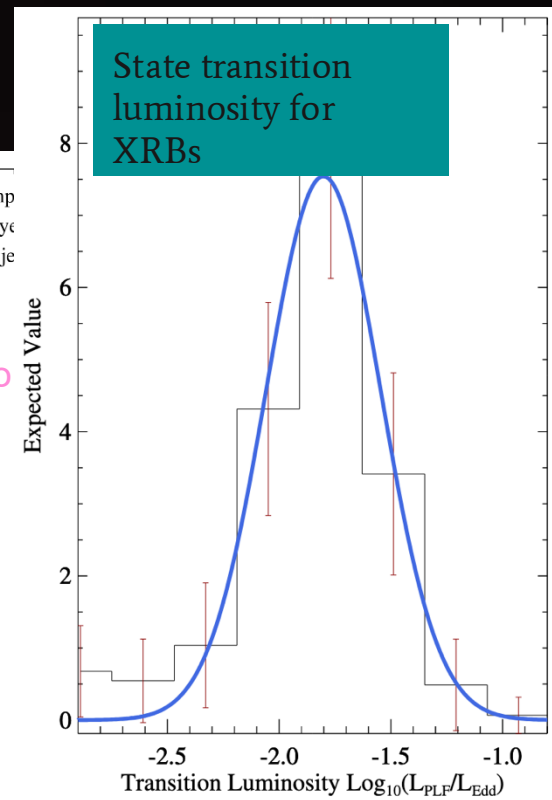
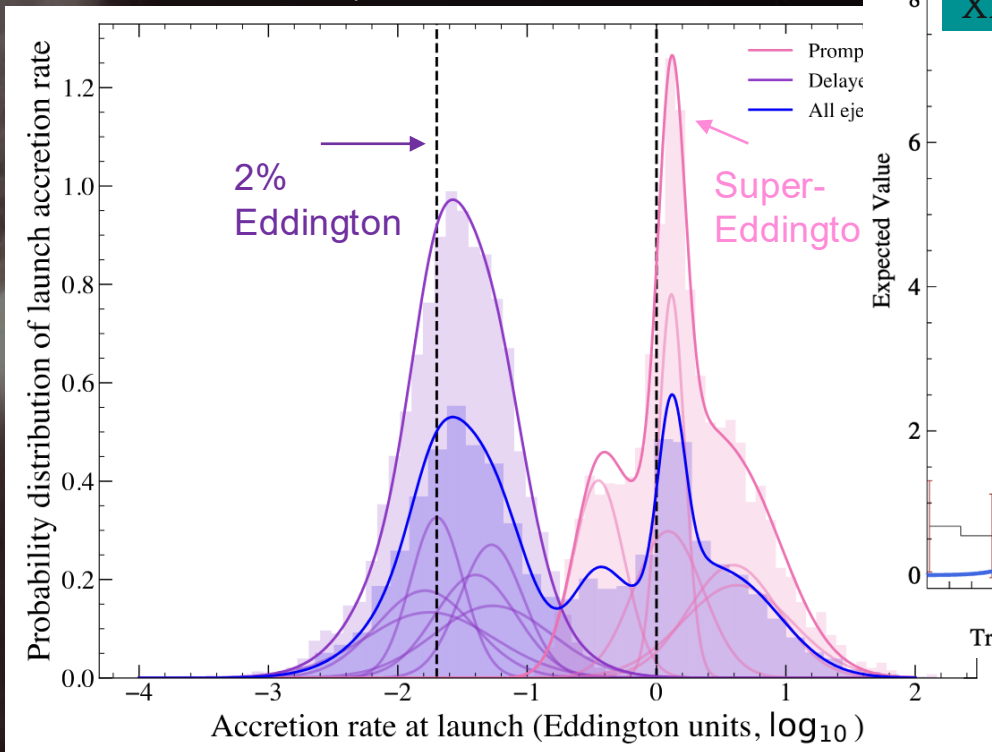
# Accretion rate at outflow launch is a bimodal distribution

Goodwin & Mummery (2026)



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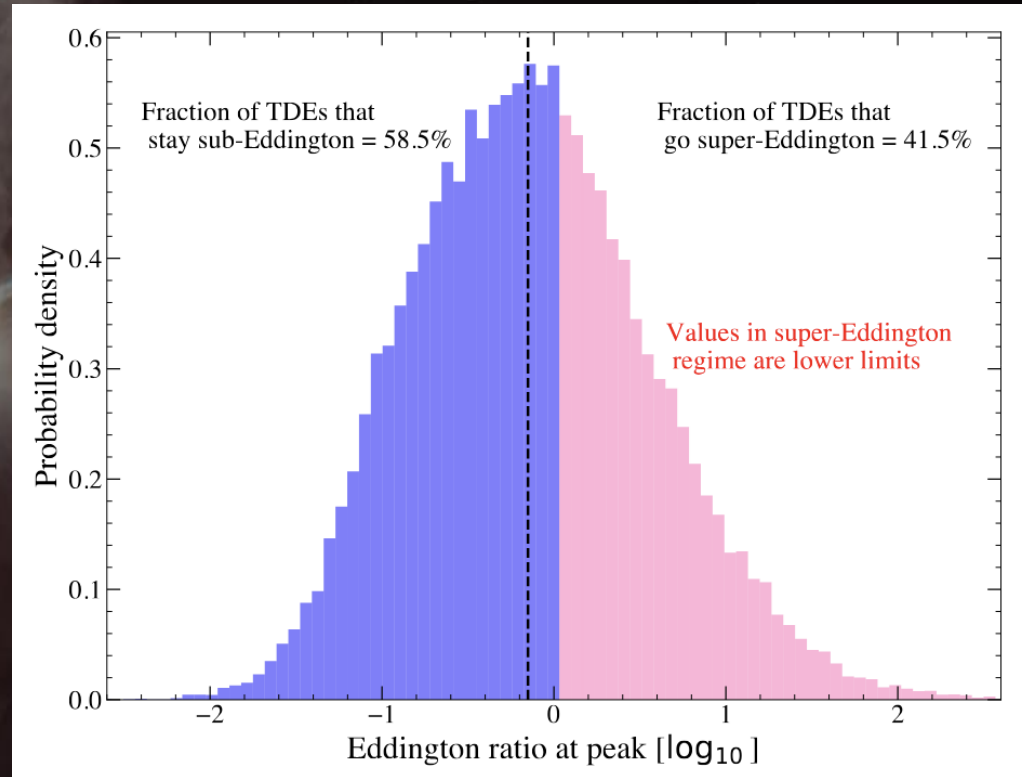
Goodwin & Mummery (2026)



Vahdat Motlagh+2018

# How many TDEs go super-Eddington?

Goodwin & Mummery (2026)



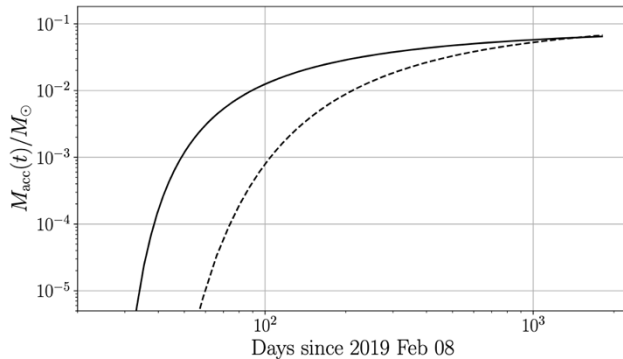
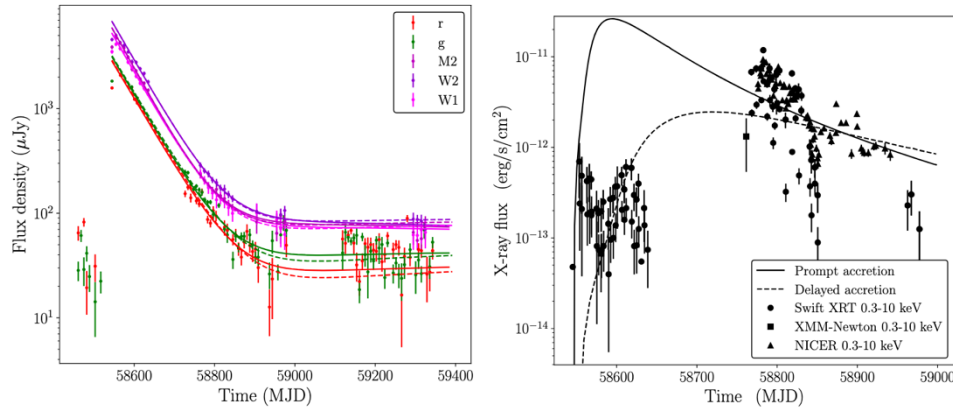
Peak accretion rate depends on:

- Stellar mass
- Viscous time

Whether super-Eddington depends on:

- Black hole mass

# What about X-rays?



X-ray emission alone is a poor tracer of the accretion rate in the disk due to obscuration

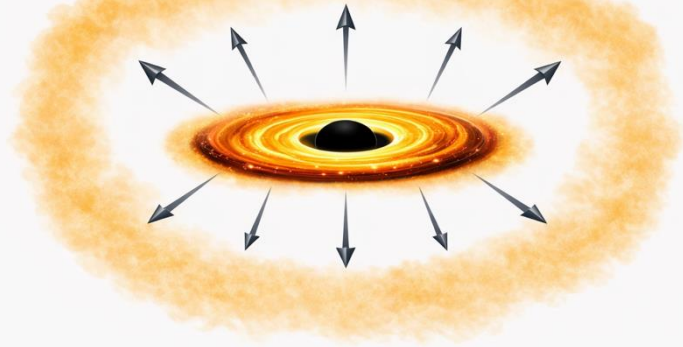
See Guolo+2024

AT2019azh  
(Goodwin+2022)

# Expectations of a disk wind

~spherical disk wind  
launched during super-  
Eddington phase

Shocks CNM produces  
synchrotron emission



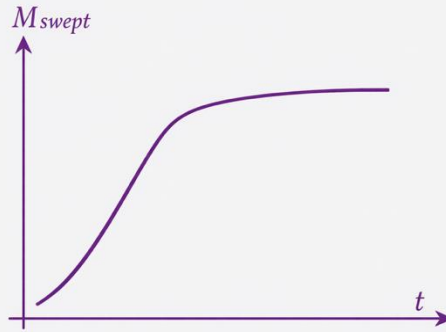
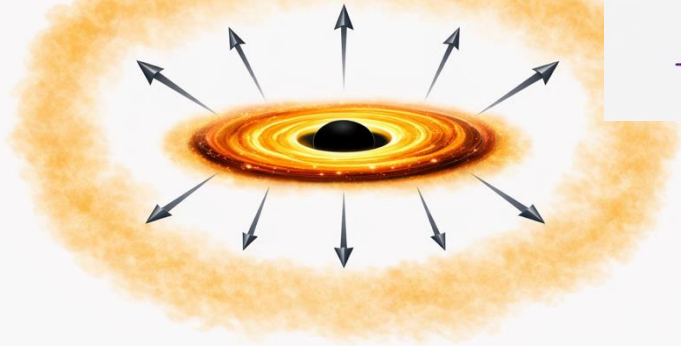
Synchrotron emission  
depends on:

- Mass/energy in outflow
- Density of ambient environment

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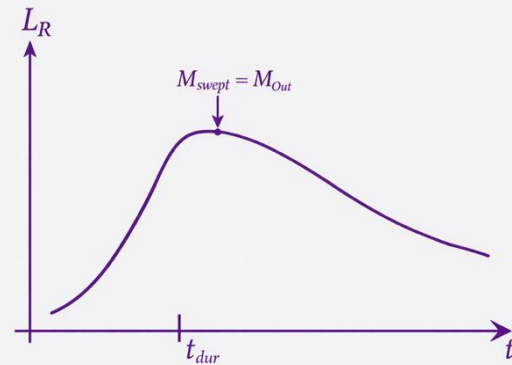
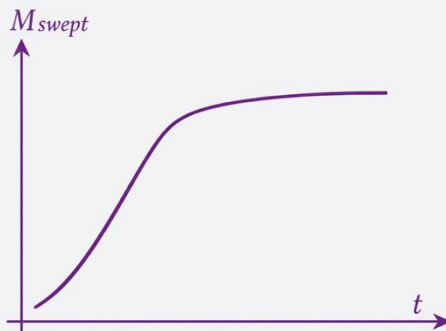
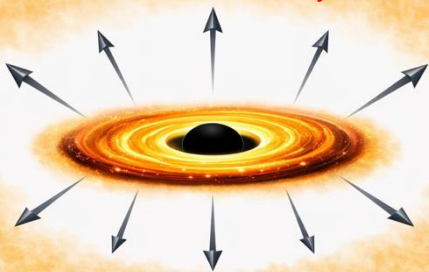
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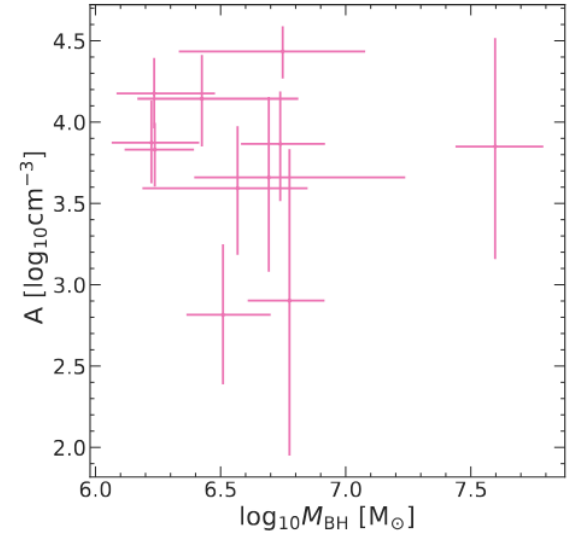
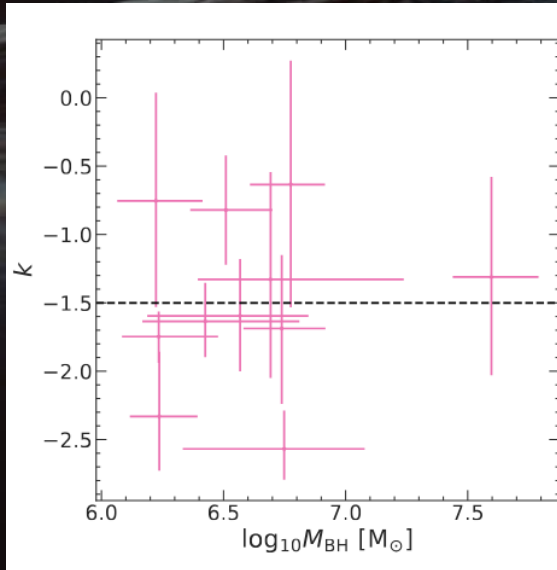
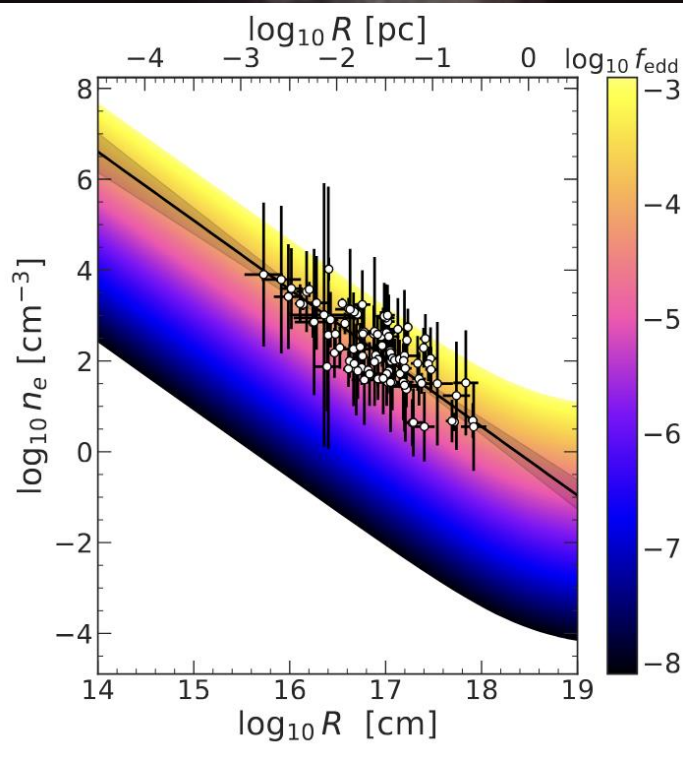
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# Effect of environment

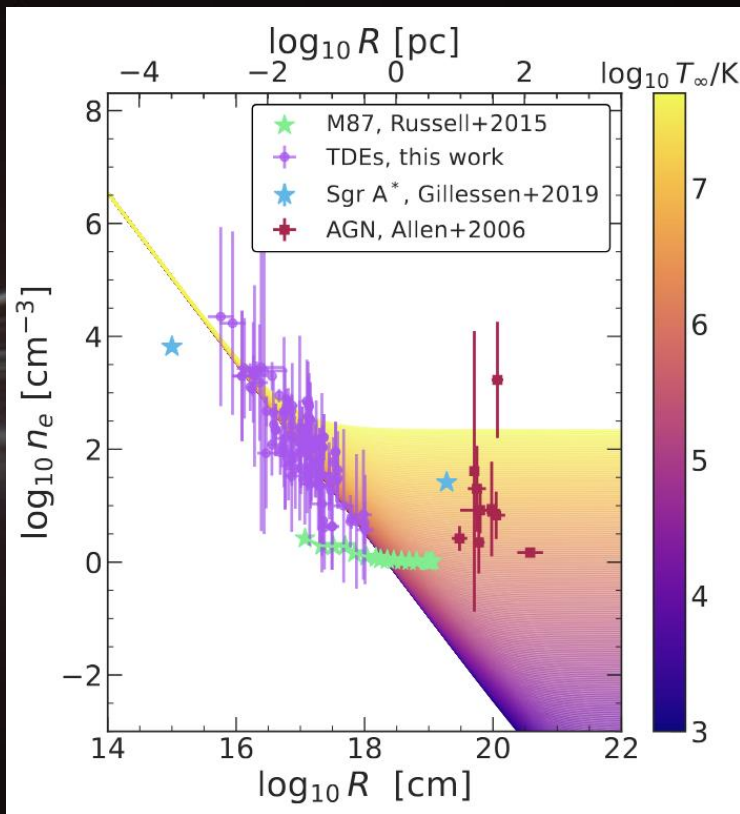
TDE hosts have very similar environment profiles

$$N_e = AR^k$$



# Revealing the intra-Bondi density profile

TDEs are powerful probes of intra-Bondi radii in galaxies



Goodwin & Mummery (in prep)

# Expectations of a disk wind

- In the viscous limited regime,  $\dot{M}_{\text{peak}}$  depends only on disk mass and viscous time (not BH mass)
- $\dot{M}_{\text{out}}$  depends on disk mass
- Velocities are roughly the same
- $R_{\text{dec}}$  depends on  $\dot{M}_{\text{out}}$  and ambient density (BH mass)
- Radio luminosity (at peak) depends on  $\dot{M}_{\text{out}}$  and B field

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Therefore, we expect:

$$E_{\text{kin},0} \propto M_{\text{disk}},$$

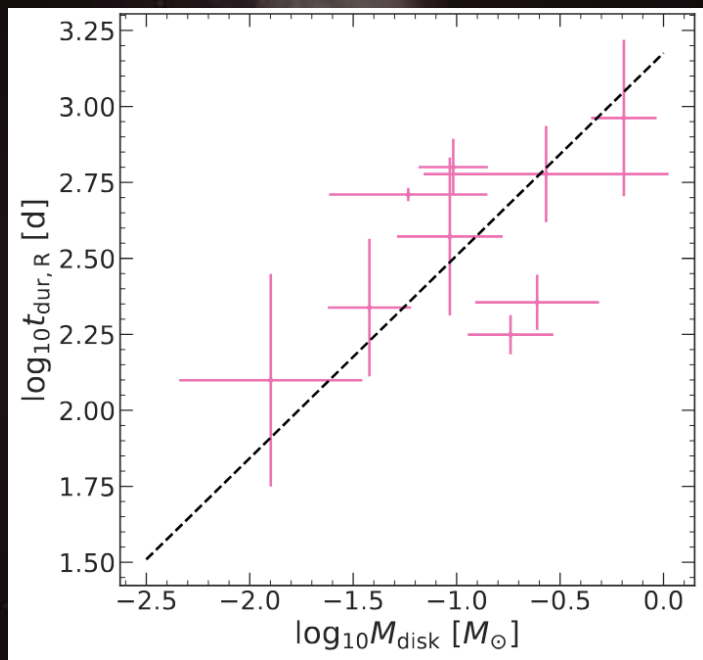
$$R_{\text{dec}} \propto M_{\text{disk}}^{2/3} M_{\bullet}^{-1/3},$$

$$L_{\nu,\text{max}} \propto M_{\text{disk}} (M_{\bullet})^{1/4},$$

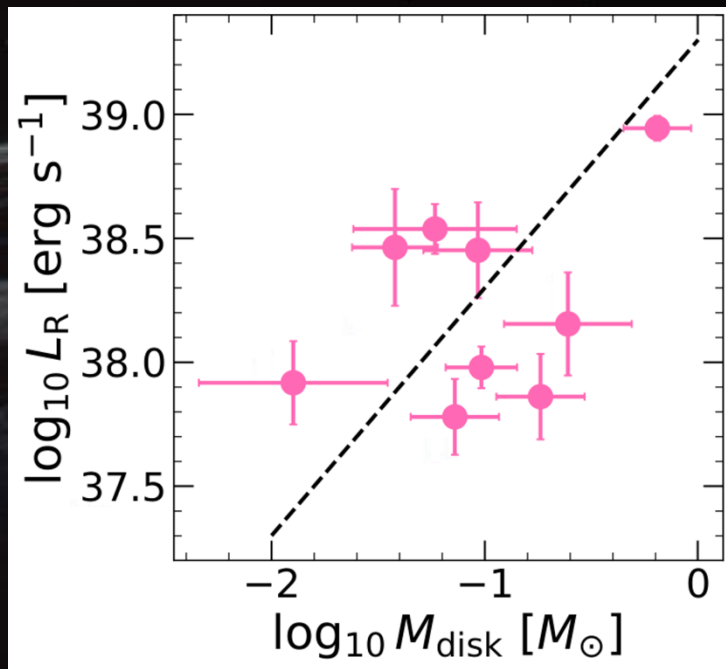
$$t_R \propto M_{\text{disk}}^{2/3} M_{\bullet}^{-1/3}.$$

# Expectations of a disk wind

$$t_{dur} \sim M_{disk}^{2/3}$$



$$L_{R,pk} \sim M_{disk}$$



# Discussion points

- Prompt radio outflows launched by accretion disk winds requires efficient disk circularization
- Observed fraction of TDEs that exhibit prompt radio emission is ~consistent with theoretical prediction of 40%. Implies uncollimated outflows and efficient disk circularization in all events
- Presence of prompt radio emission governed by both disk (stellar) mass and black hole mass
  - Low mass black holes are more likely to go Super-Eddington but launch less luminous outflows

# Summary

- ★ Combing radio and accretion disk modelling in TDEs reveals outflows are launched at two critical accretion rates: during a super-Eddington wind phase and at a critical accretion rate of 2% Eddington
- ★ Approximately ~40% of TDEs should launch a (prompt) super-Eddington disk wind
- ★ With the (small) sample publicly available, the duration and luminosity of prompt radio flares correlates as expected with a disk mass for a disk wind
- ★ Radio emission reveals that TDE hosts have very similar ambient density profiles that are consistent with the expectations of a Bondi sphere
- ★ This model requires that disk formation is prompt in majority of TDEs