

ブラックホール時空で運動する 弦のカオスの普遍性

棚橋典大 [九州大学 マス・フォア・インダストリ研究所]

based on

Chaos of Wilson Loop from String Motion near Black Hole Horizon

橋本幸士、村田佳樹、棚橋典大 [arXiv:1803.06756]

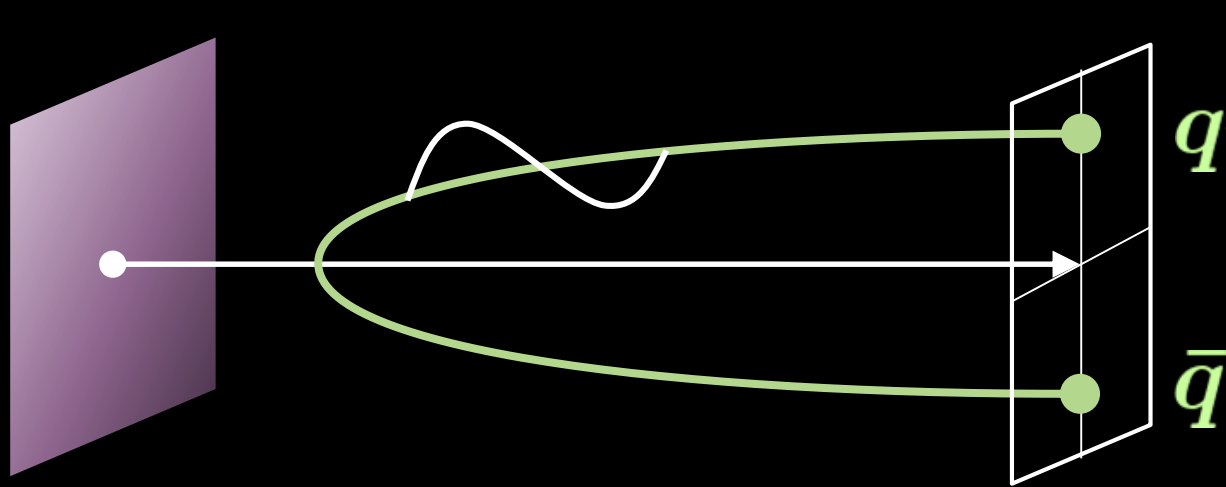
Universality in Chaos of Particle Motion near Black Hole Horizon

橋本幸士、棚橋典大 [arXiv:1610.06070]

Setup:

Fundamental String moving near **AdS BH horizon**

= Dynamics of “quark-anti quark pair” at finite temperature



Maldacena '98
Rey & Yee '98

Results:

- ✓ **String motion become chaotic** due to BH gravity
⇔ **Force between the quarks** becomes chaotic when $T \neq 0$.
- ✓ **Lyapunov exponent λ** is smaller than **surface gravity κ**

$$\lambda \leq \kappa = 2\pi T / \hbar$$

[Maldacena-Shenker-Stanford '15]

CONTENTS

1. AdS string & perturbative motion
2. Nonlinear time evolution & chaos
3. Summary

CLASSICAL CHAOS

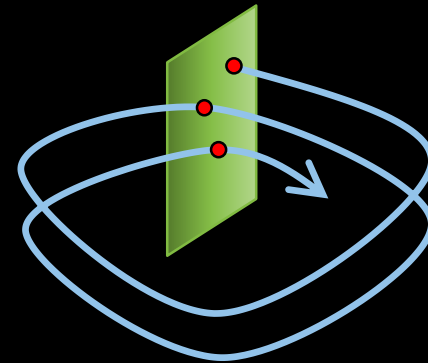
◆ Diagnostics of classical chaos

• Poincaré plot

= Section of orbits in phase space

non-chaotic → regular shaped plot

chaotic → scattered plot



• Lyapunov exponent λ = Separation growth rate of nearby orbits

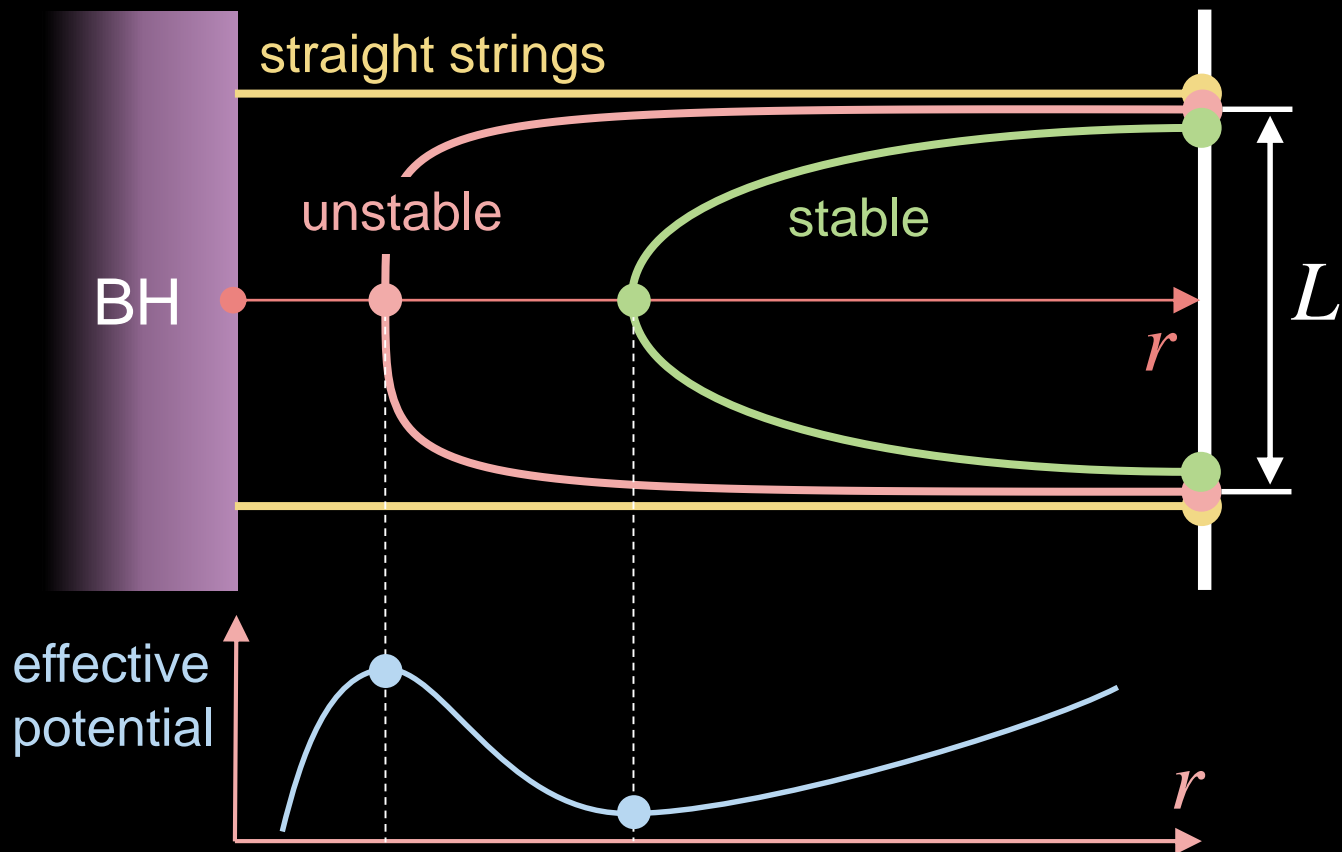
$d(0)$

$d(t) \sim d(0) \exp(\lambda t)$

✓ We will focus on the string motion for a while.

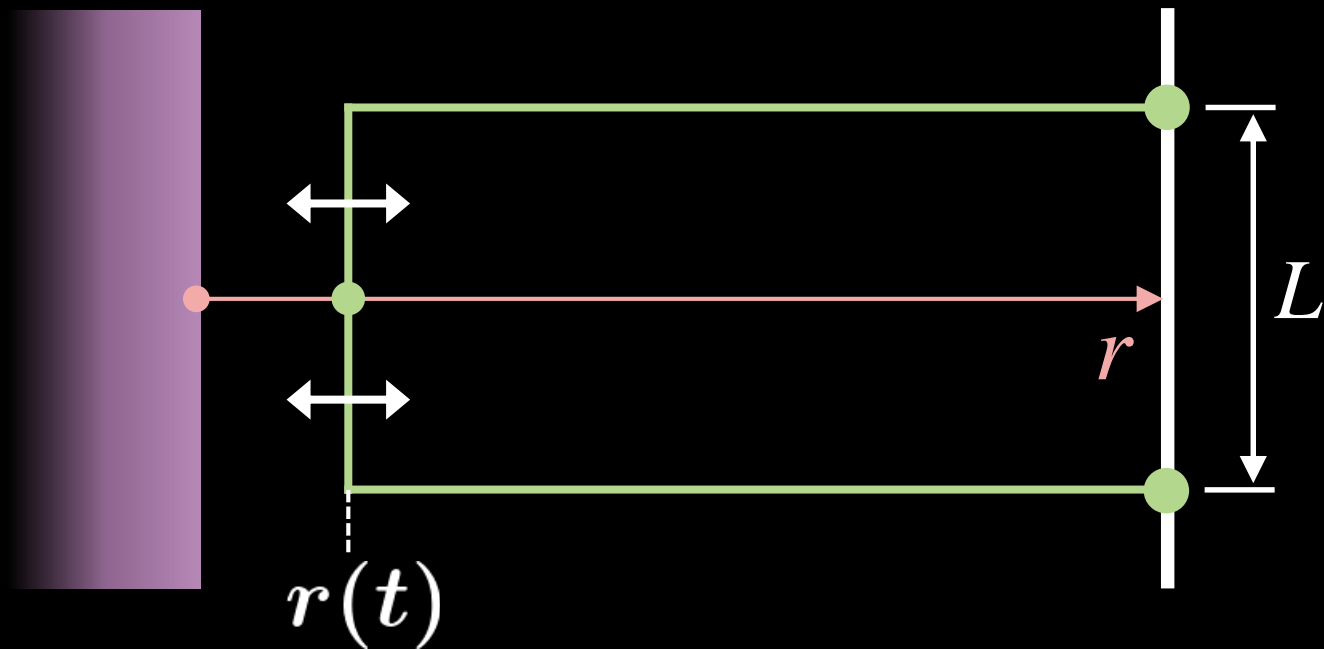
STRING IN AdS

Three shapes of static Nambu-Goto string in AdS



STRING IN AdS

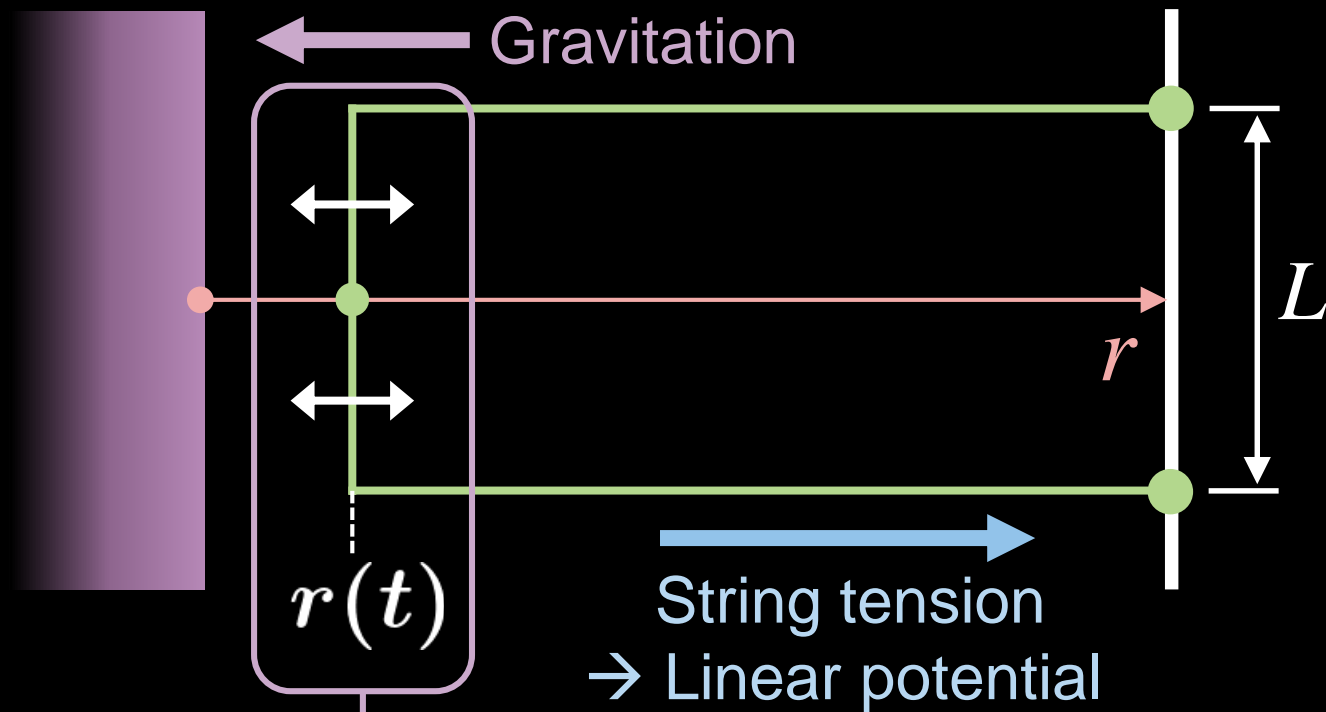
“Square-shaped string” approximation



$$\mathcal{L} \simeq -L \sqrt{r^4(t) f(r(t)) - \frac{\dot{r}^2(t)}{f(r(t))}} + 2(r(t) - r_H) \left[f(r) = 1 - \frac{r_H^4}{r^4} \right]$$

STRING IN AdS

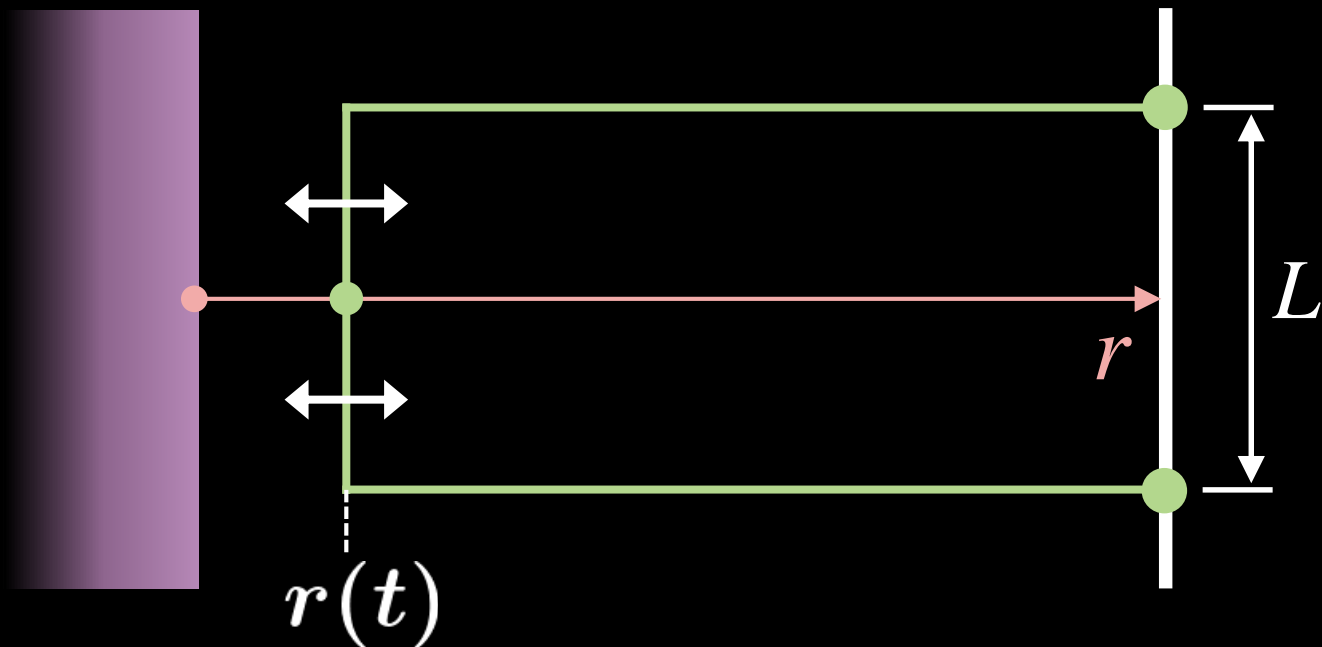
“Square-shaped string” approximation



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STRING IN AdS

“Square-shaped string” approximation

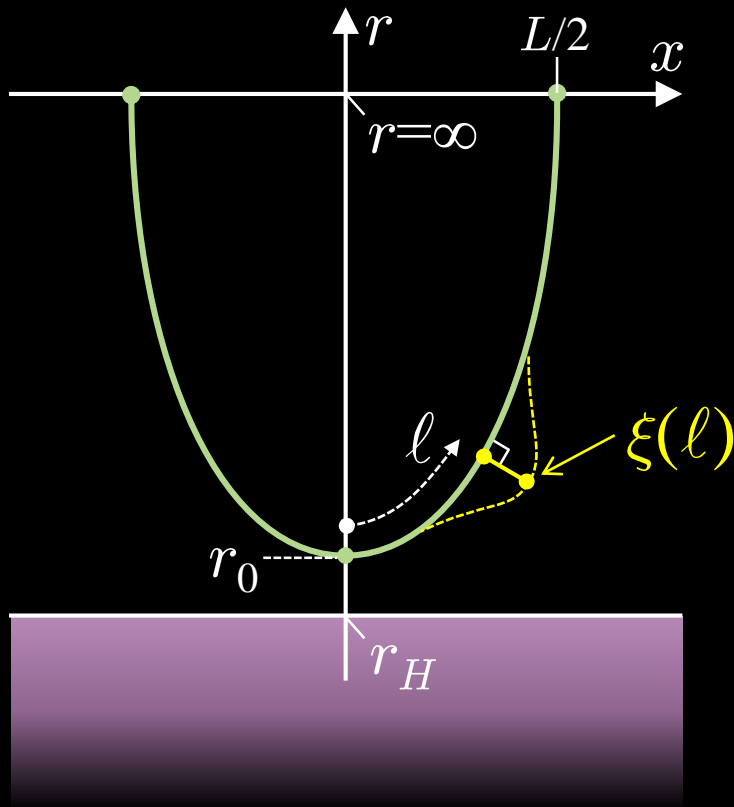


When the string is on the potential maximum near horizon,

$$\mathcal{L} \simeq \frac{1}{2r_H^5 L^2} \left[\dot{r}^2 + \kappa^2 (r(t) - r_*)^2 \right] \Rightarrow r(t) \sim r_* + e^{\kappa t}$$

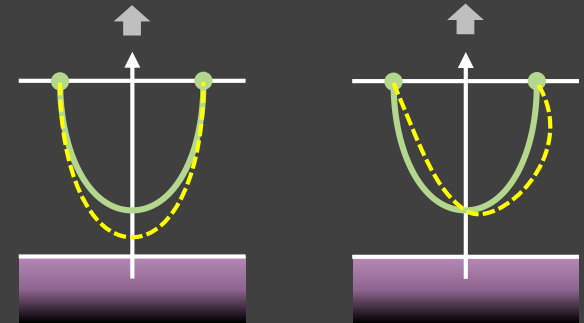
$$\Rightarrow \lambda \lesssim \kappa = 2\pi T / \hbar$$

Perturbative string motion



Truncate $\xi(l)$ up to lowest 2 modes

$$\xi(t, l) = c_0(t)e_0(l) + c_1(t)e_1(l)$$



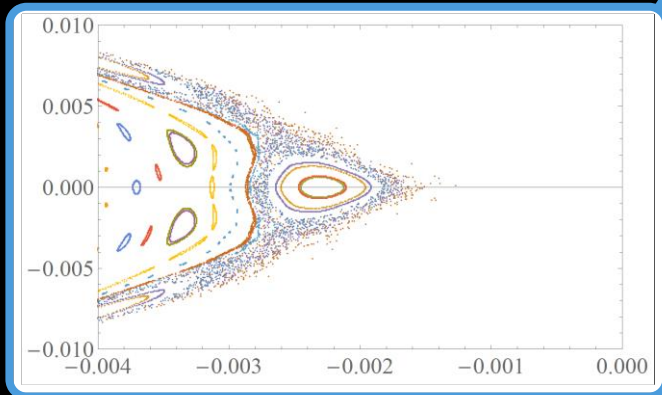
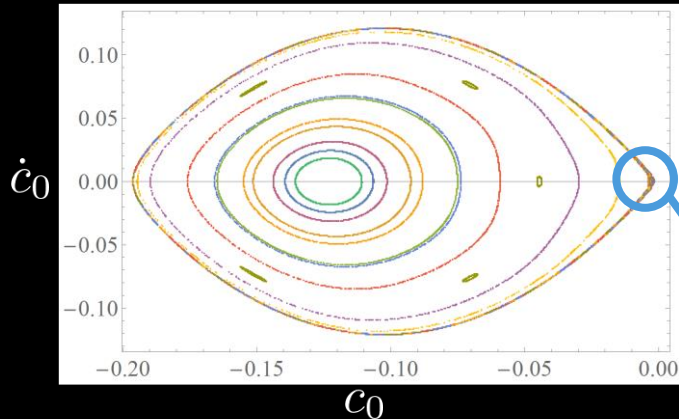
Then expand \mathcal{L} up to 3rd order in $c_i(t)$

ex.) For $r_0 = 1.1 r_H$,

$$\begin{aligned} \frac{\mathcal{L}}{\mathcal{T}} = & \sum_{n=0,1} (\dot{c}_n^2 - \omega_n^2 c_n^2) + 7.11c_0^3 + 35.3c_0c_1^2 \\ & + 4.66c_0\dot{c}_0^2 + 1.32c_0\dot{c}_1^2 - 7.57\dot{c}_0c_1\dot{c}_1 \\ & (\omega_0^2 = -1.40, \omega_1^2 = 7.57) \end{aligned}$$

Perturbative string motion

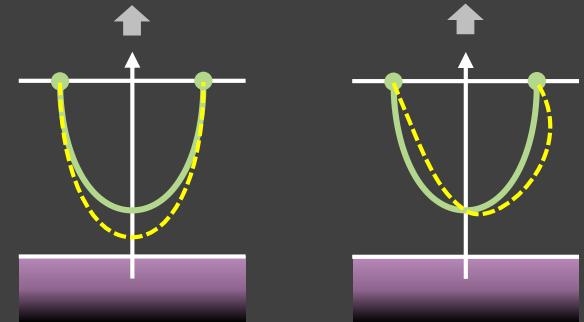
Poincaré plot



- Chaotic (scattered plot) in near-horizon region
- $\lambda \sim 0.04 \times 2\pi T / \hbar$

Truncate $\xi(\ell)$ up to lowest 2 modes

$$\xi(t, \ell) = c_0(t)e_0(\ell) + c_1(t)e_1(\ell)$$



Then expand \mathcal{L} up to 3rd order in $c_i(t)$

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Chaos Bound for AdS String

Nonlinear dynamics of string

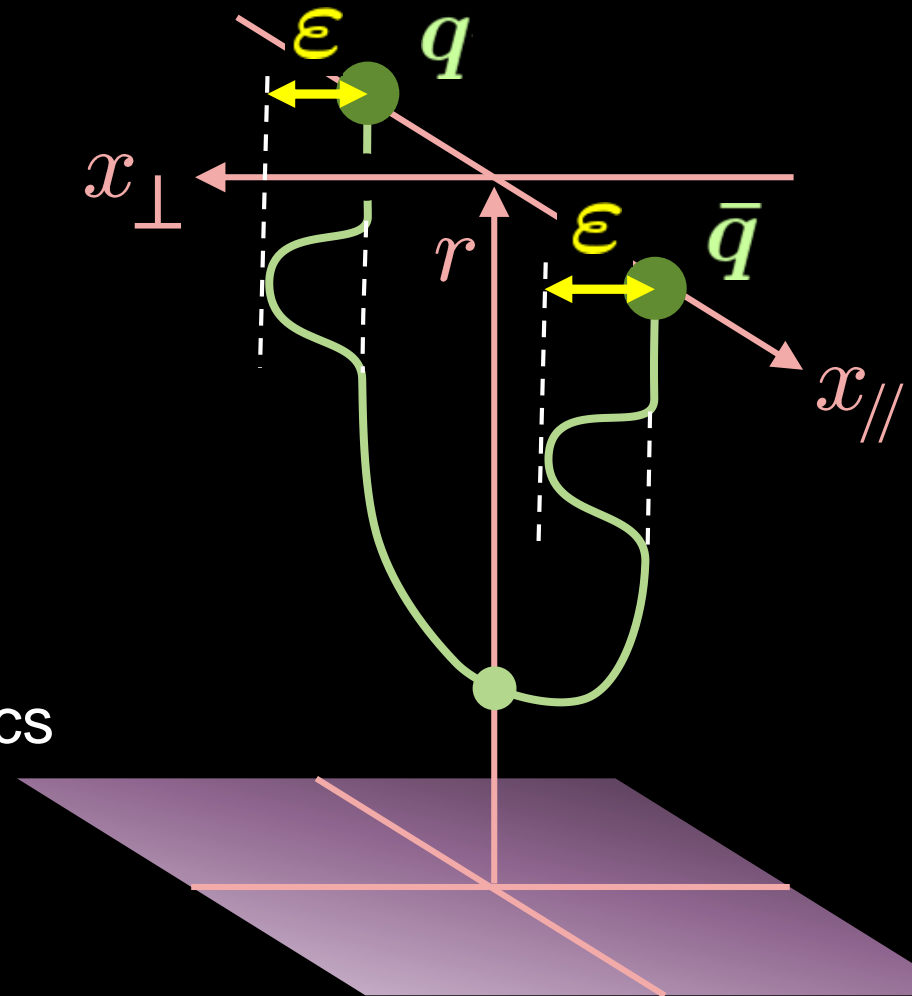
Shake the string end points
by amplitude ϵ

→ Nonlinear string motion



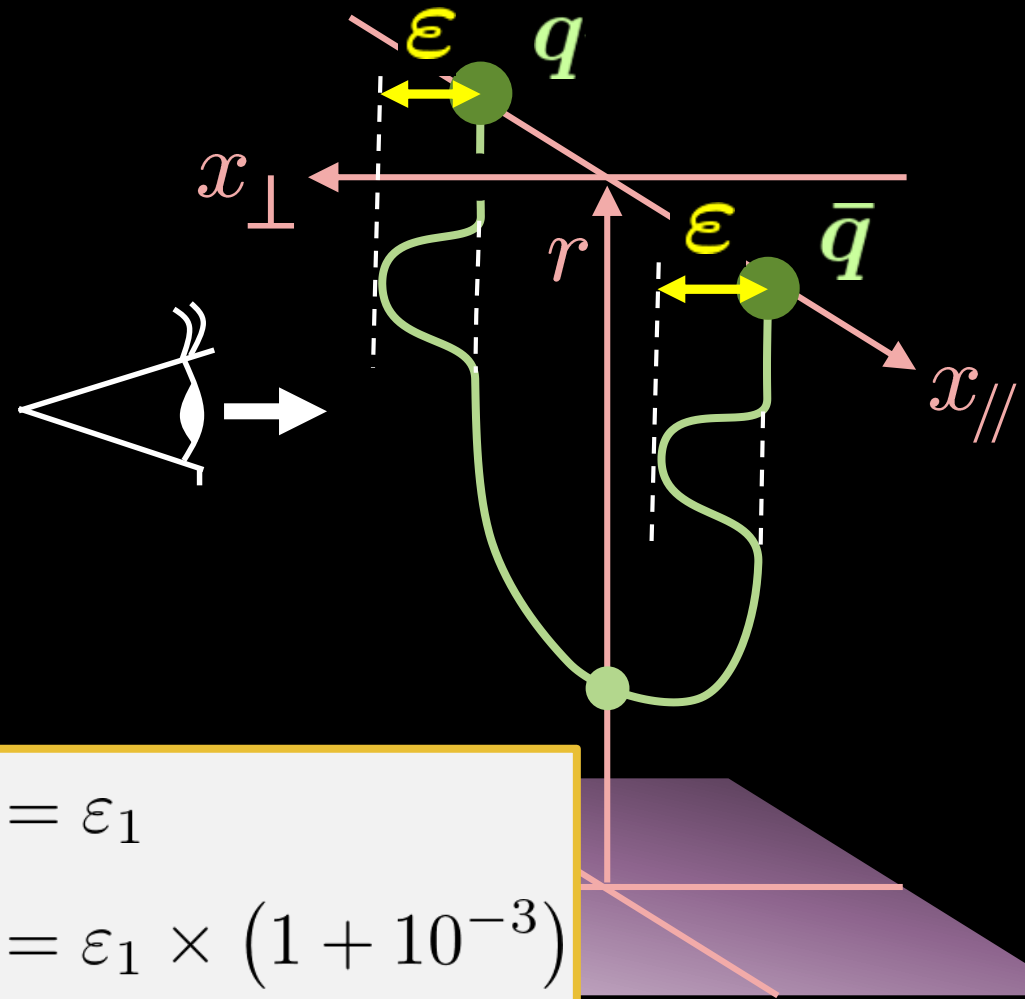
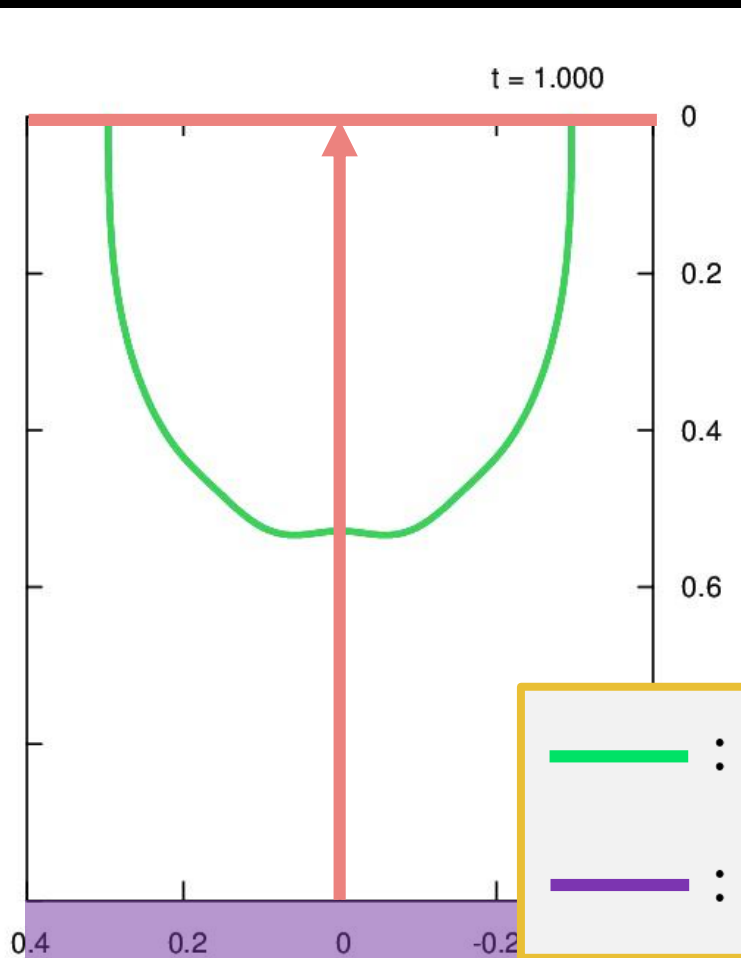
Shake the quark— $\bar{\text{quark}}$ pair

→ Nonlinear flux tube dynamics



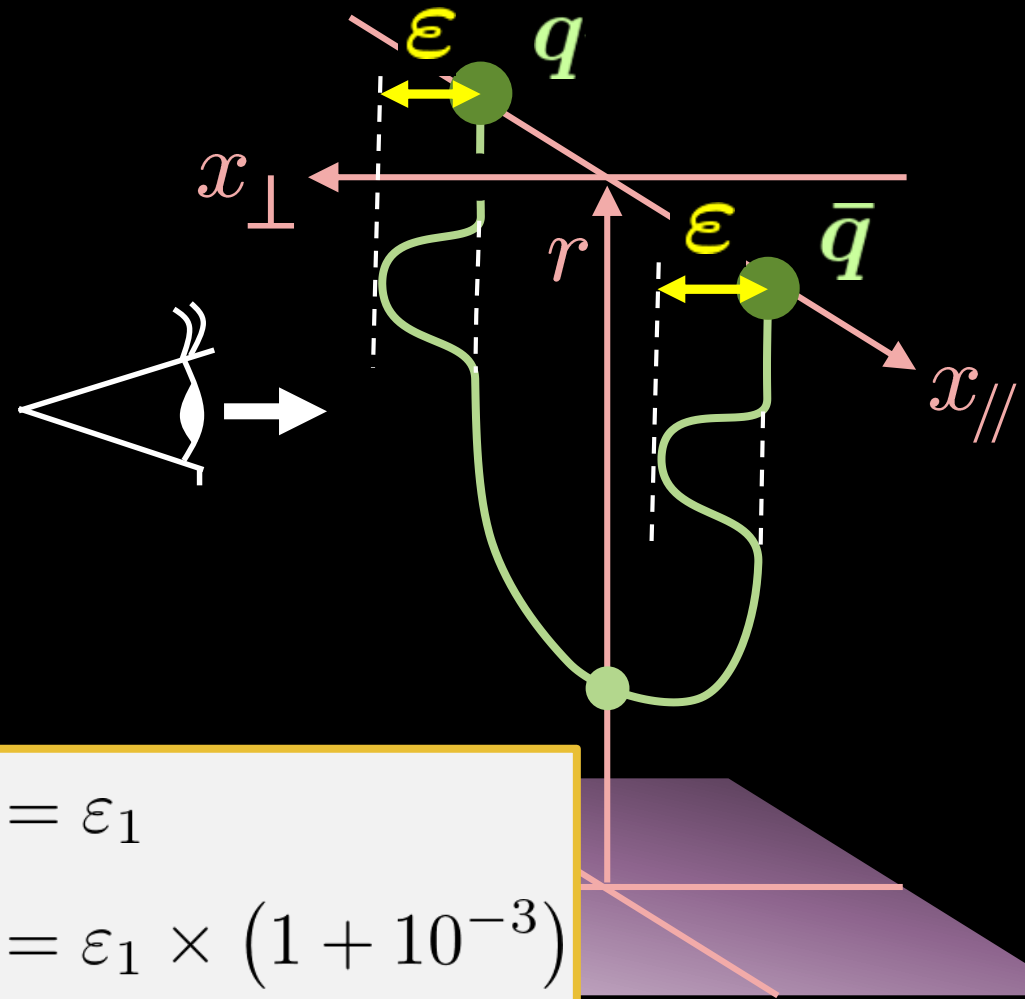
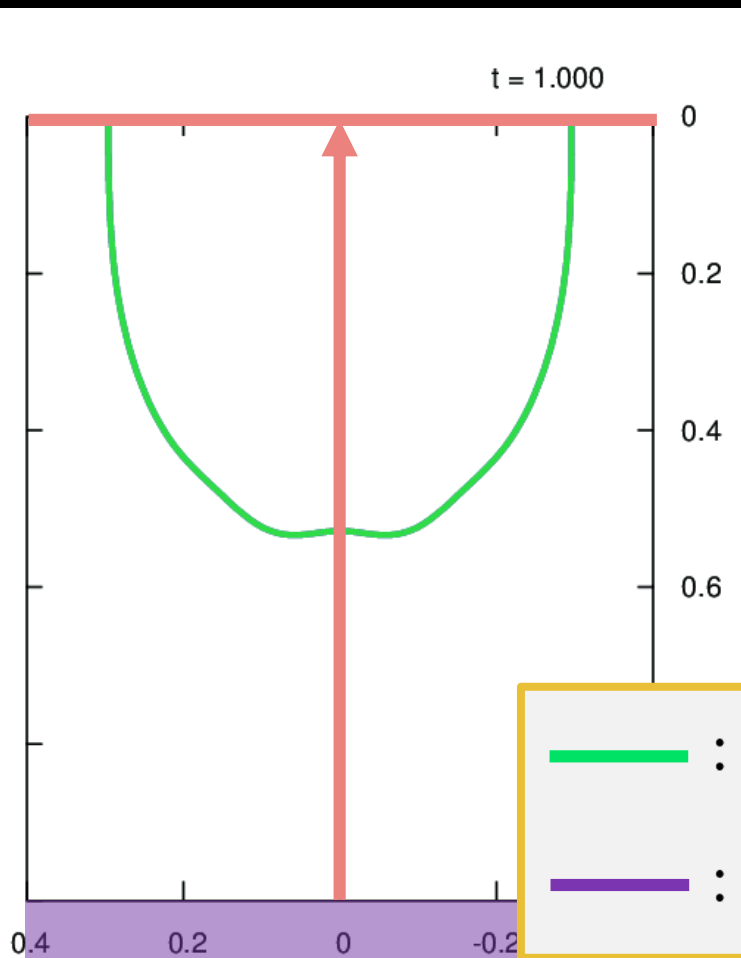
Chaos Bound for AdS String

Nonlinear dynamics of string for slightly different ϵ



Chaos Bound for AdS String

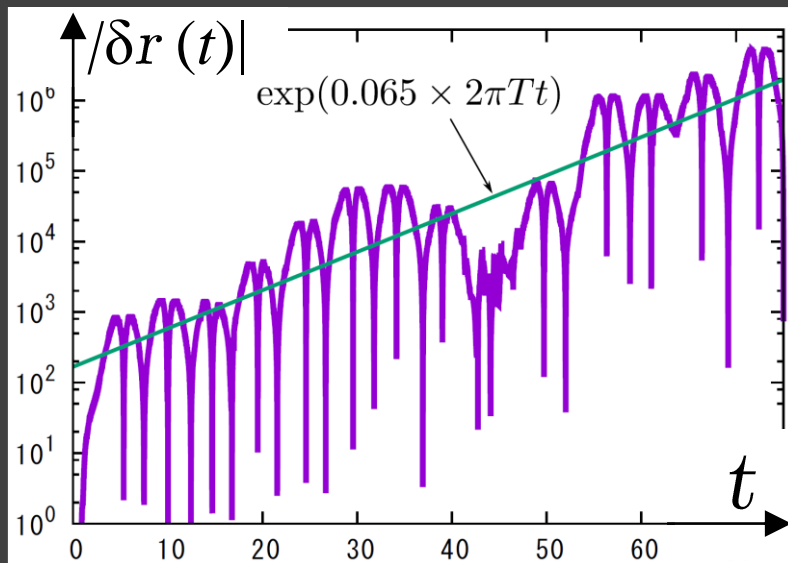
Nonlinear dynamics of string for slightly different ϵ



Chaos Bound for AdS String

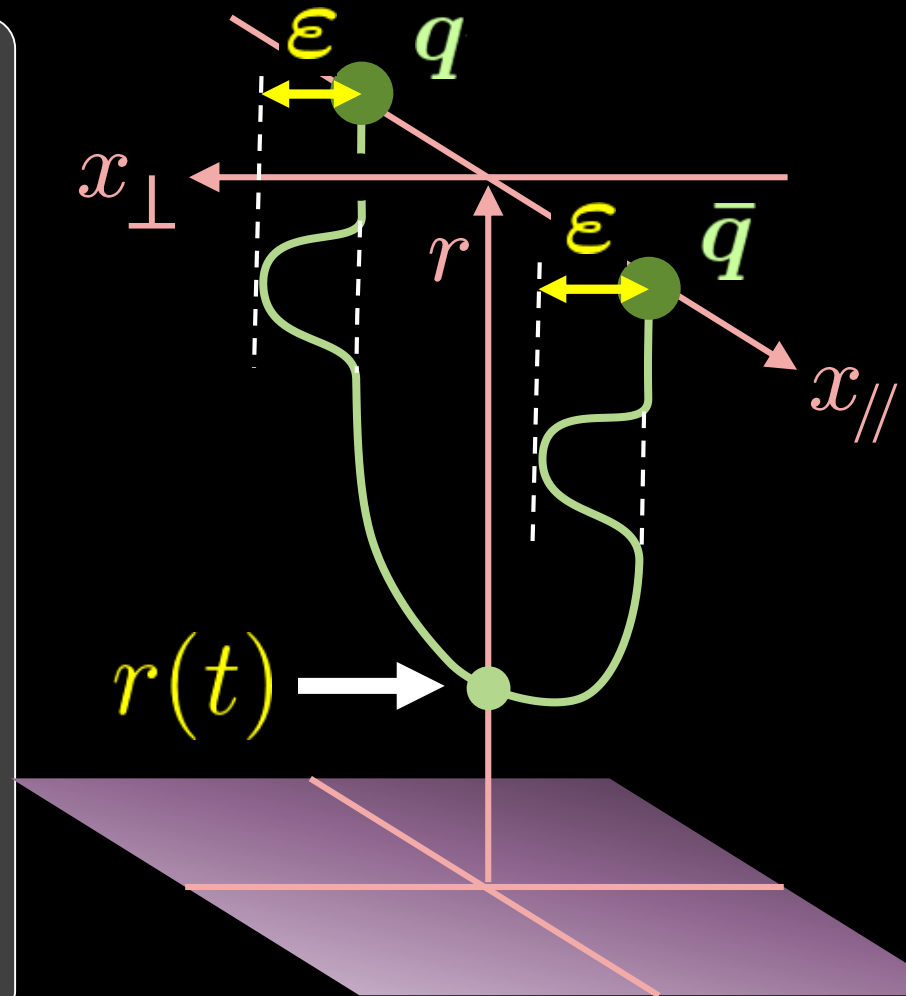
Nonlinear dynamics of string for slightly different ϵ

String position difference $\delta r(t)$



$$|\delta r(t)| \sim e^{\lambda t},$$

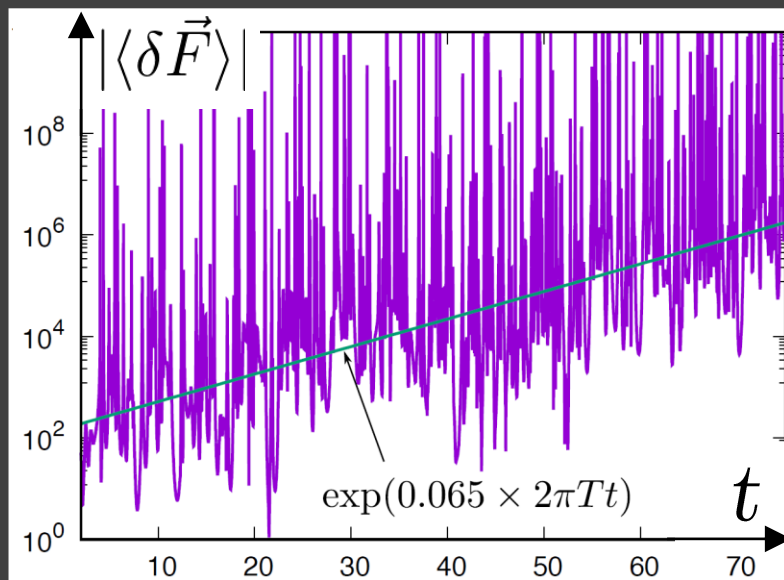
$$\lambda \sim 0.065 \times 2\pi T / \hbar$$



Chaos Bound for AdS String

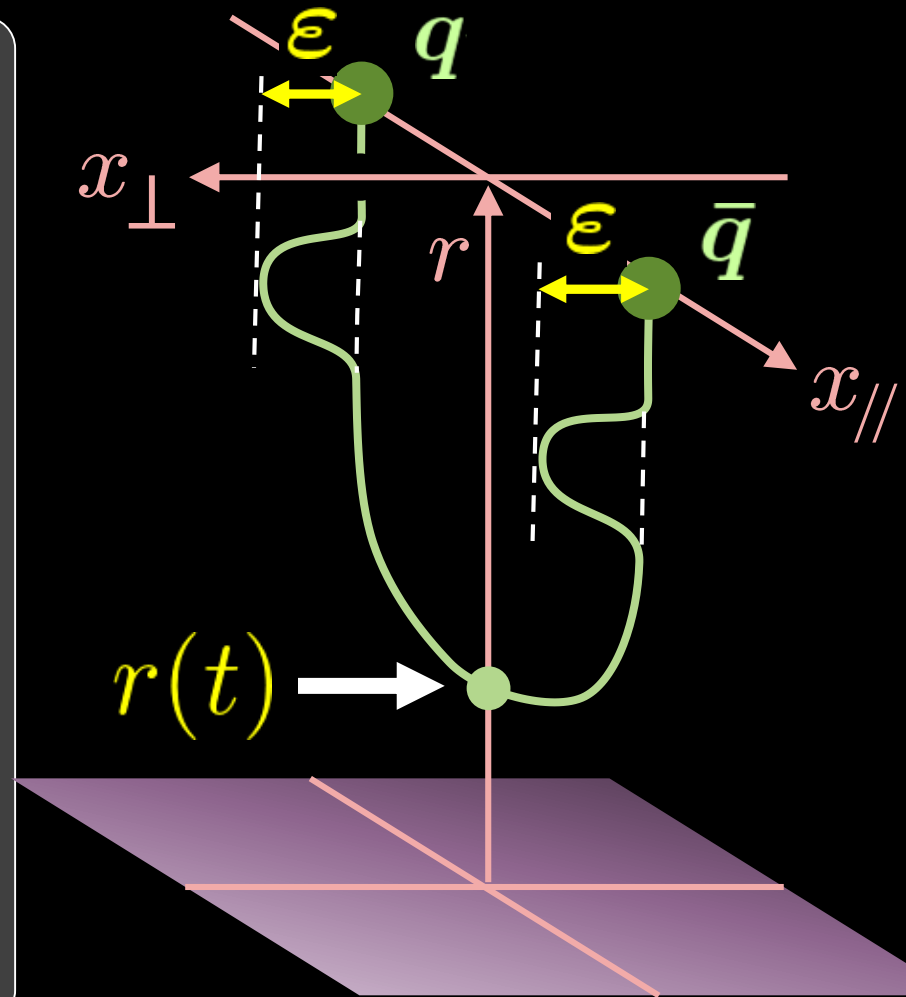
Nonlinear dynamics of string for slightly different ϵ

Inter-quark force $\langle \vec{F}(t) \rangle \propto \partial_r^3 \vec{X}(t, r)$



$$|\langle \delta \vec{F} \rangle| \sim e^{\lambda t},$$

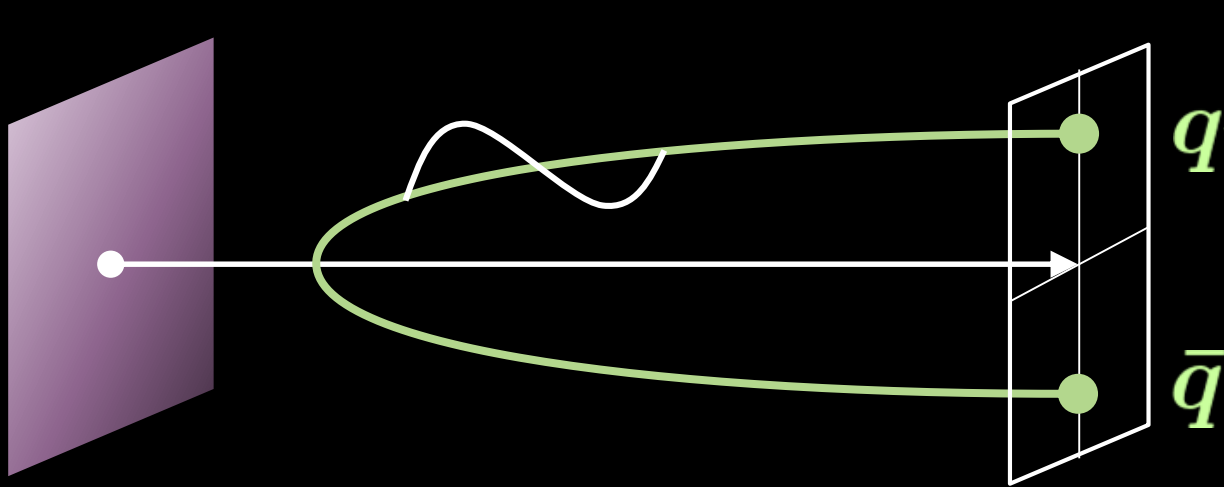
$$\lambda \sim 0.065 \times 2\pi T / \hbar$$



Setup:

Fundamental String moving near **AdS BH horizon**

= Dynamics of “quark-anti quark pair” at finite temperature



Maldacena '98
Rey & Yee '98

Results:

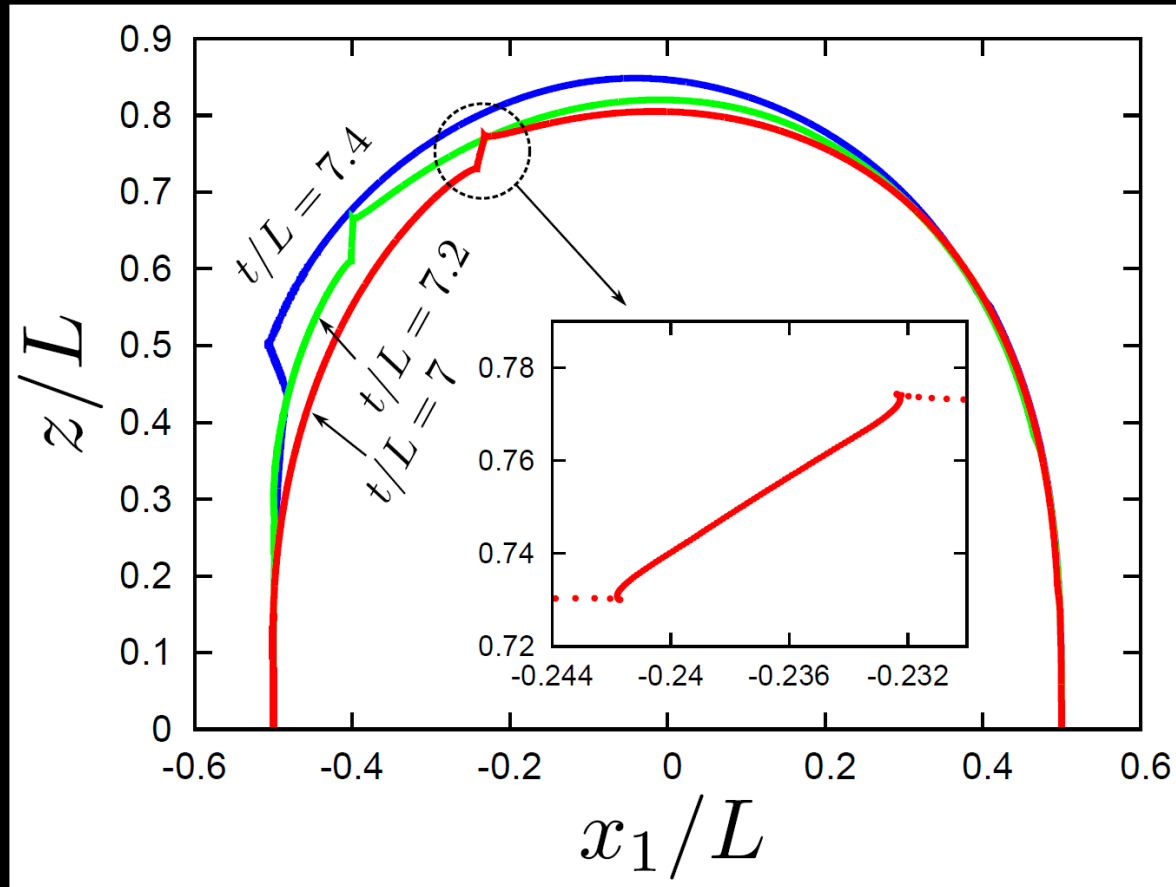
- ✓ **String motion become chaotic** due to BH gravity
⇔ **Force between the quarks** becomes chaotic when $T \neq 0$.

- ✓ **Lyapunov exponent λ** is smaller than **surface gravity κ**

$$\lambda \leq \kappa = 2\pi T / \hbar$$

[Maldacena-Shenker-Stanford '15]

◆ Cusp formation on the string

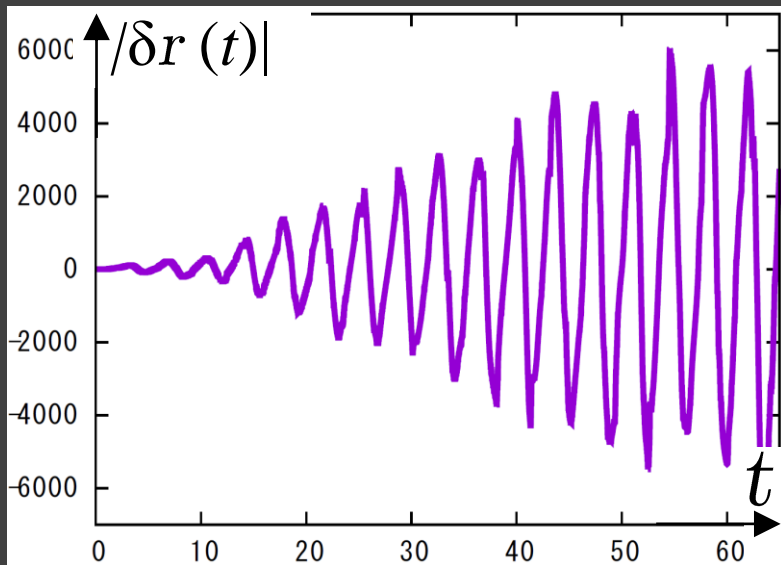


[Ishii, Murata '15]

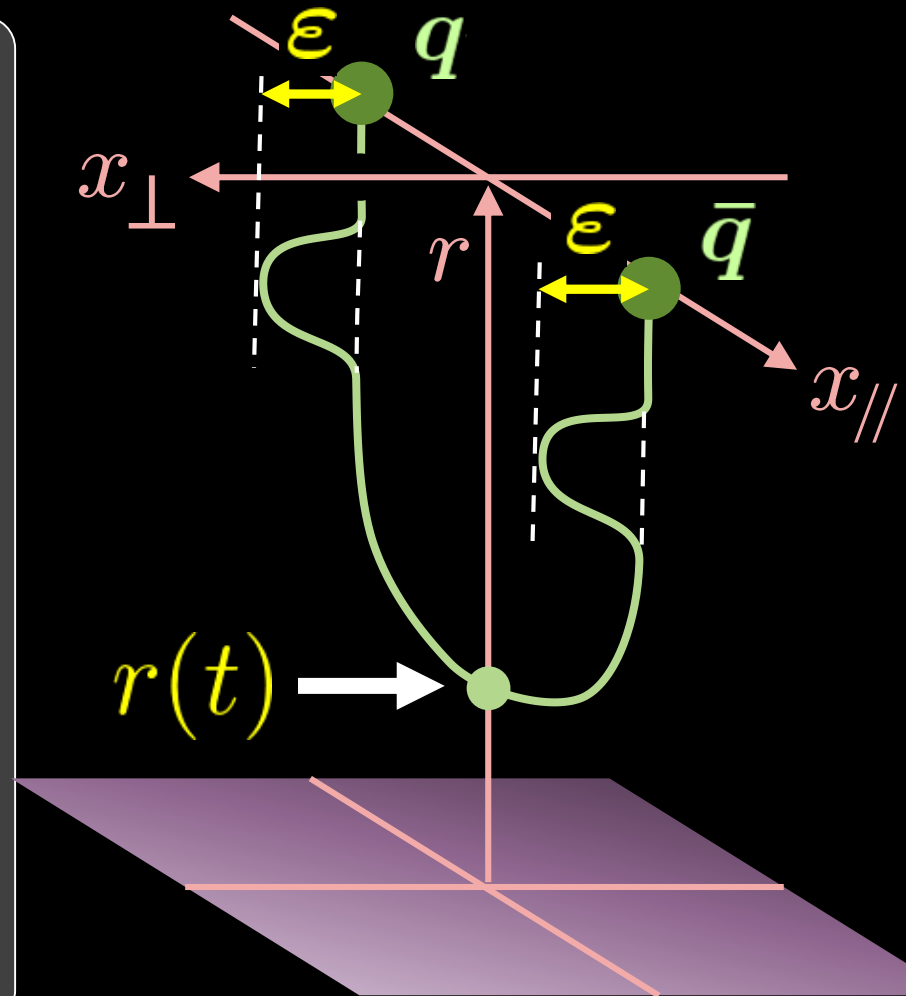
Chaos Bound for AdS String

Nonlinear dynamics of string for slightly different ϵ

When the string is
not close to the horizon,



λ grows just by power-law,
not exponentially.



◆ Motivation:

A bound on chaos in QFT at temperature T :

$$\lambda \leq 2\pi T / \hbar$$

[Maldacena-Shenker-Stanford '15]

◆ Probing the effect of temperature T to chaos in QFT.

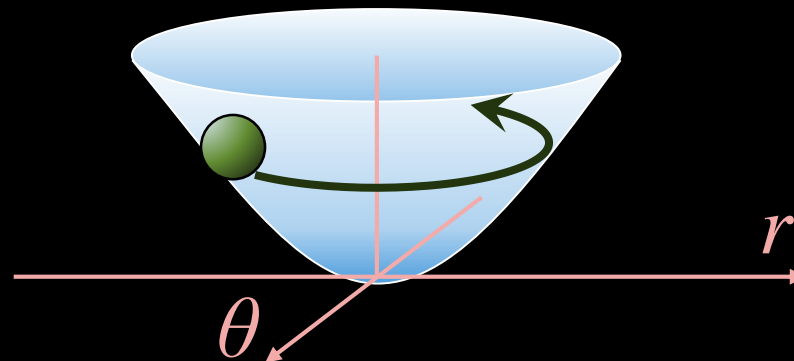
■ We study effect of temperature to chaos in classical gravity.

Use BH surface gravity $\kappa = 2\pi T / \hbar$ instead.

■ To probe effect of κ , we look at trajectories very close to BH.

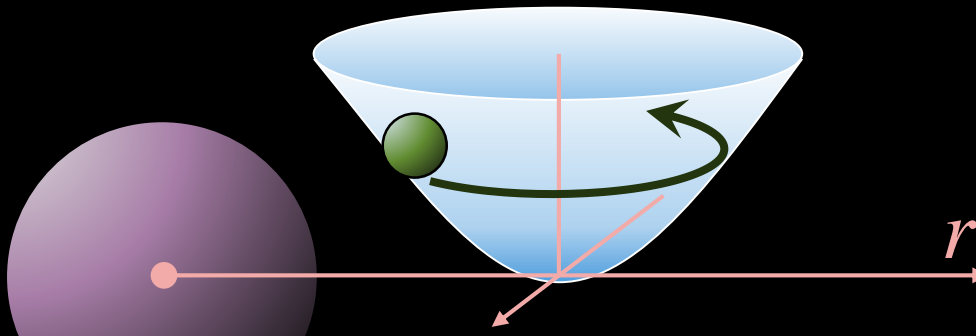
Chaos Bound for Particle near BH

- To realize a particle moving very close to BH horizon,
 1. put a particle in a trapping harmonic potential



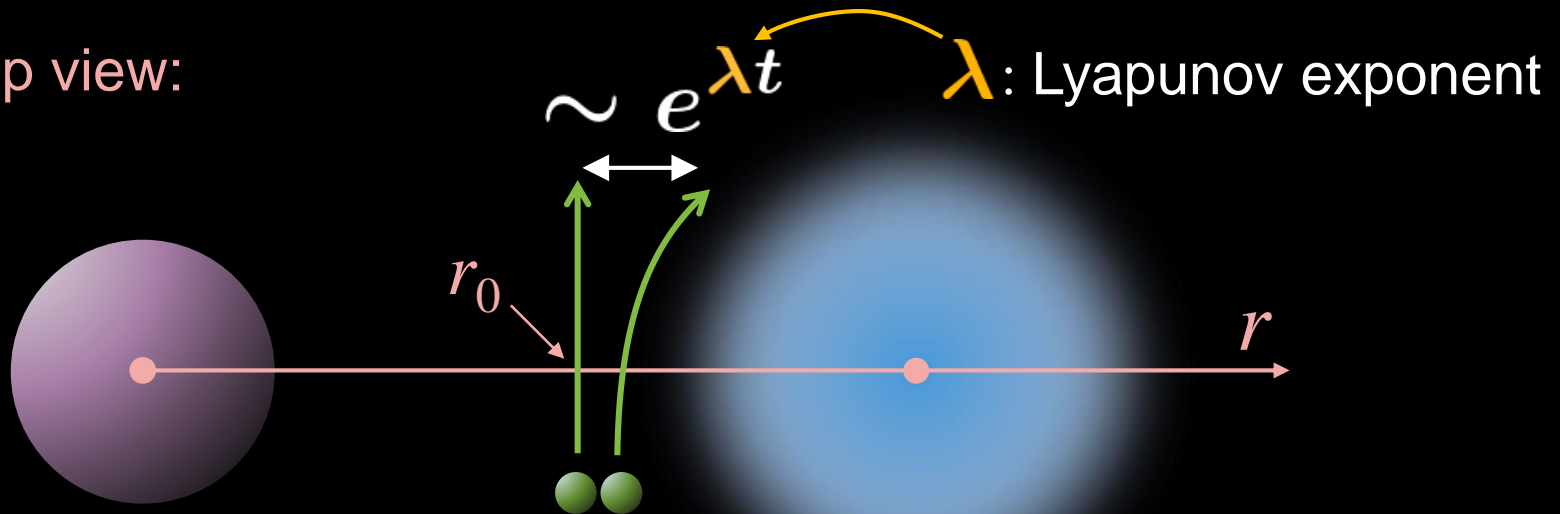
(← no chaos)

2. take it close to a BH horizon

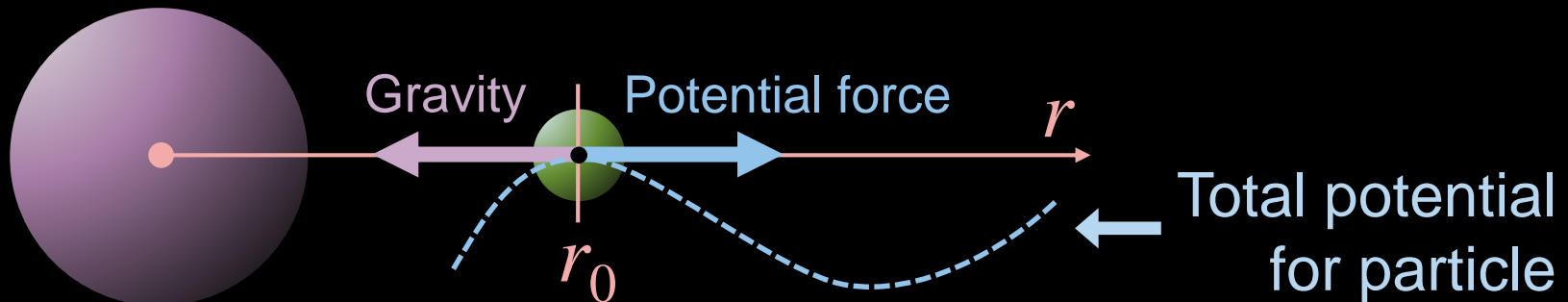


Chaos Bound for Particle near BH

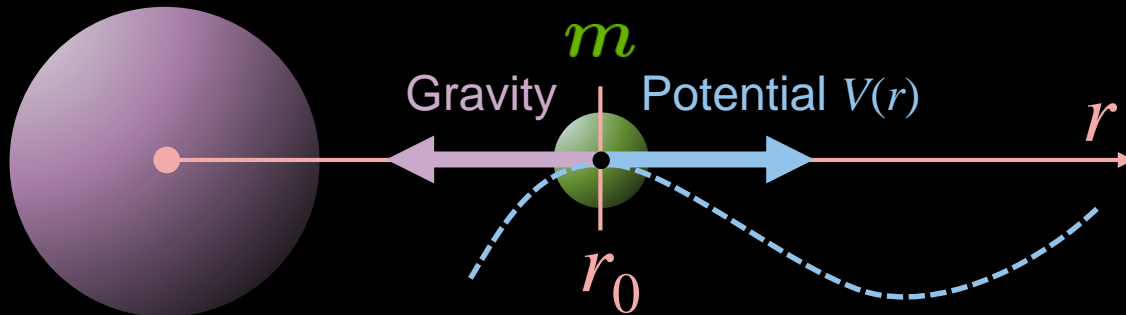
Top view:



2. take it close to a BH horizon



Chaos Bound for Particle near BH



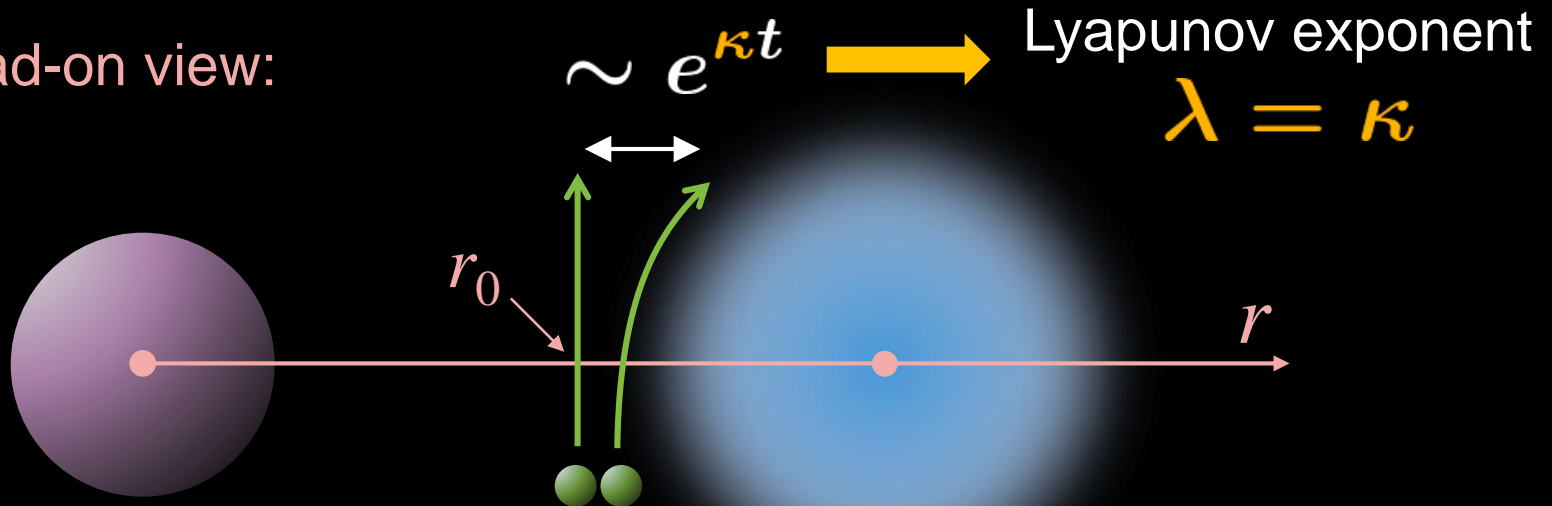
$$\mathcal{L} = -m\sqrt{-g_{\mu\nu}\dot{X}^\mu\dot{X}^\nu} - V(X) \simeq -m\sqrt{f(r) - \frac{\dot{r}^2}{f(r)} - V(r)}$$

- Slow radial velocity near $r = r_0$
- Linear approximation for $V(r) \sim (\text{slope}) \times (r - r_{\text{horizon}})$
- Near-horizon limit $r_0 \rightarrow r_{\text{horizon}}$

$$\mathcal{L} \simeq C(m, \kappa, \text{slope of } V) \times [\dot{r}^2 + \kappa^2 (r - r_0)^2]$$

Chaos Bound for Particle near BH

Head-on view:



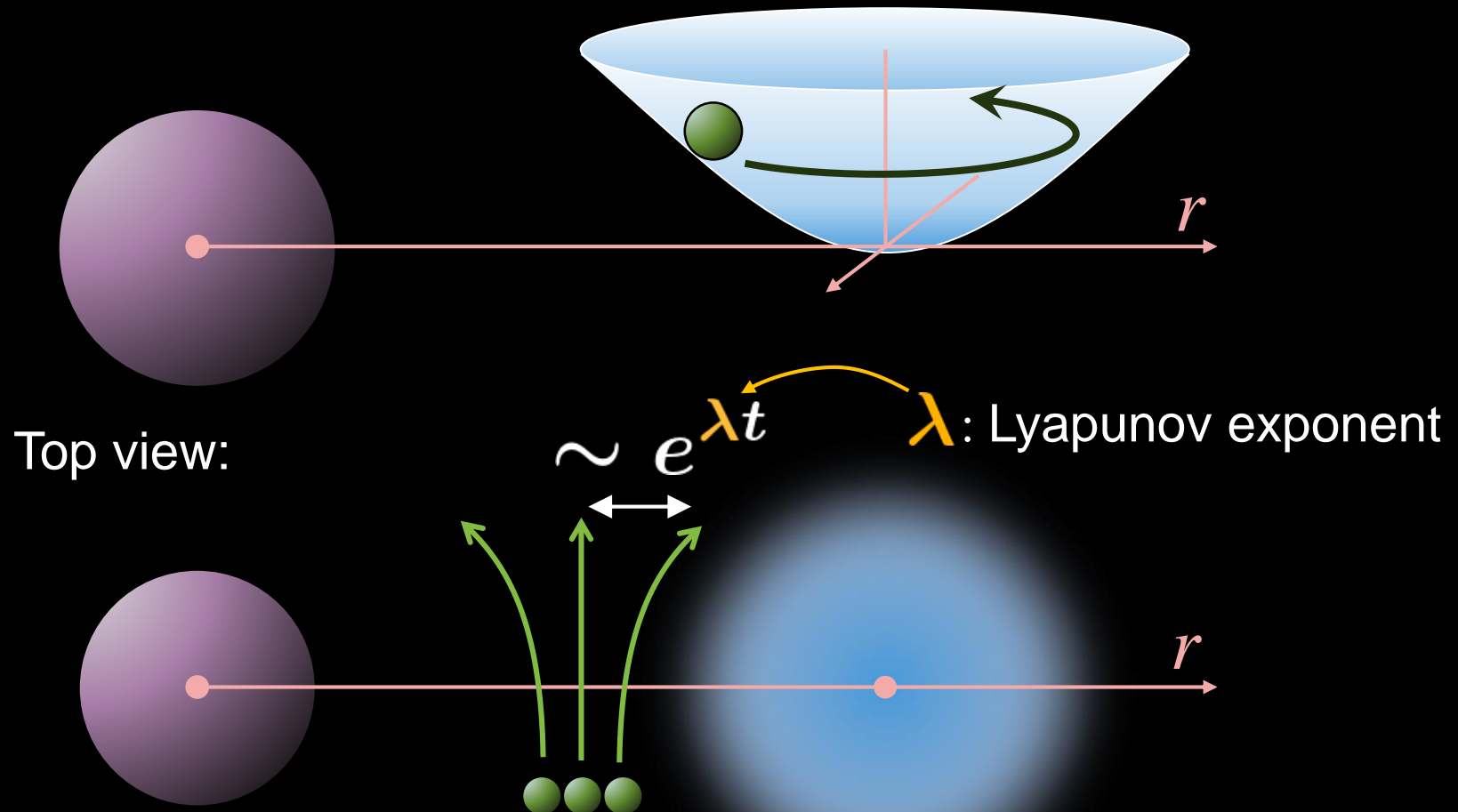
- Linear approximation for $V(r) \sim (\text{slope}) \times (r - r_{\text{horizon}})$
- **Near-horizon limit** $r_0 \rightarrow r_{\text{horizon}}$

$$\mathcal{L} \simeq C(m, \kappa, \text{slope of } V) \times [\dot{r}^2 + \kappa^2 (r - r_0)^2]$$

✓ A generic trajectory obeys $\lambda \leq \kappa = 2\pi T / \hbar$

Chaos Bound for Particle near BH

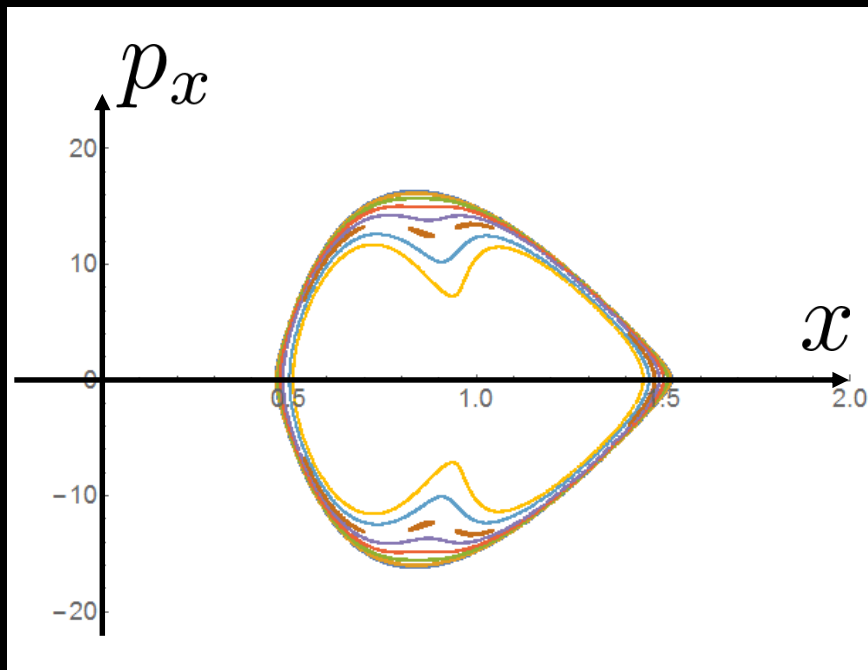
Particle in harmonic potential near a BH



Poincaré plot at $y = 0$

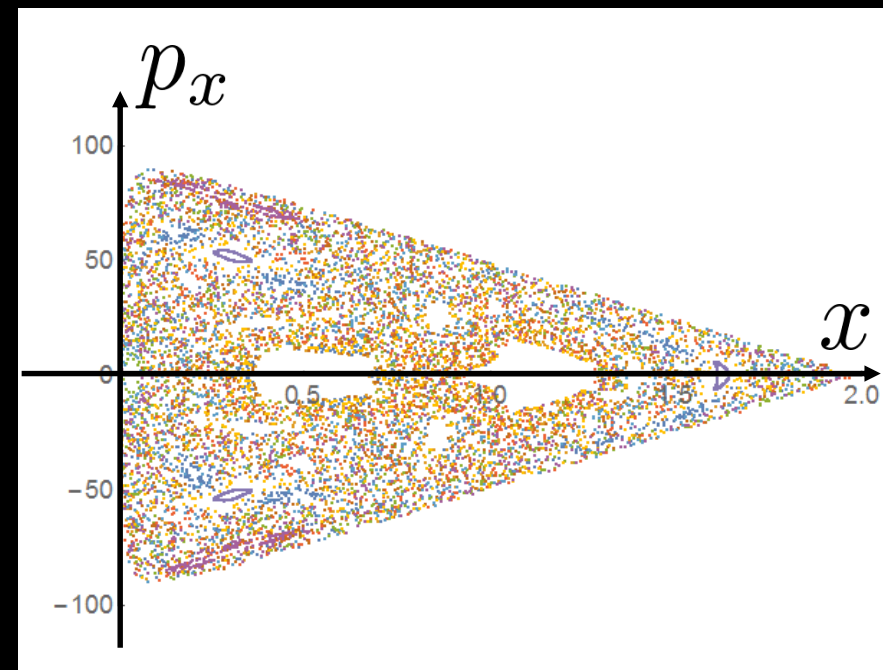
$$\mathcal{L} = -\sqrt{f(x) - \frac{\dot{x}^2}{f(x)} - \dot{y}^2} - \frac{\omega^2}{2} \left[(x - x_c)^2 + y^2 \right] \quad \left[f(x) \equiv 2\kappa x \right]$$

Particle near Potential Minimum



Regular KAM tori, no chaos

Particle near BH Horizon

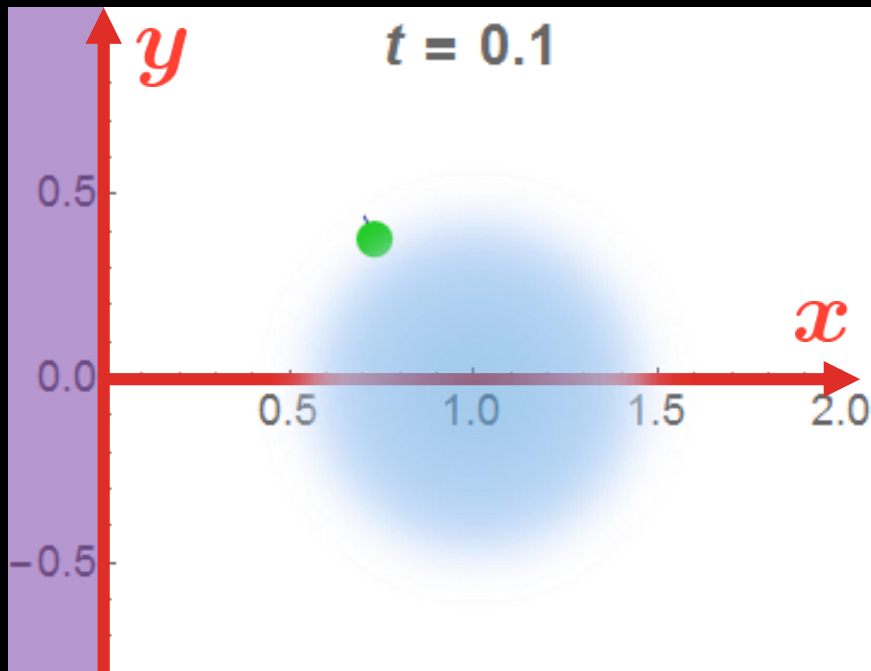


Lyapunov exponent $\lambda \sim 0.2 \kappa$
Satisfies the bound $\lambda \leq \kappa$

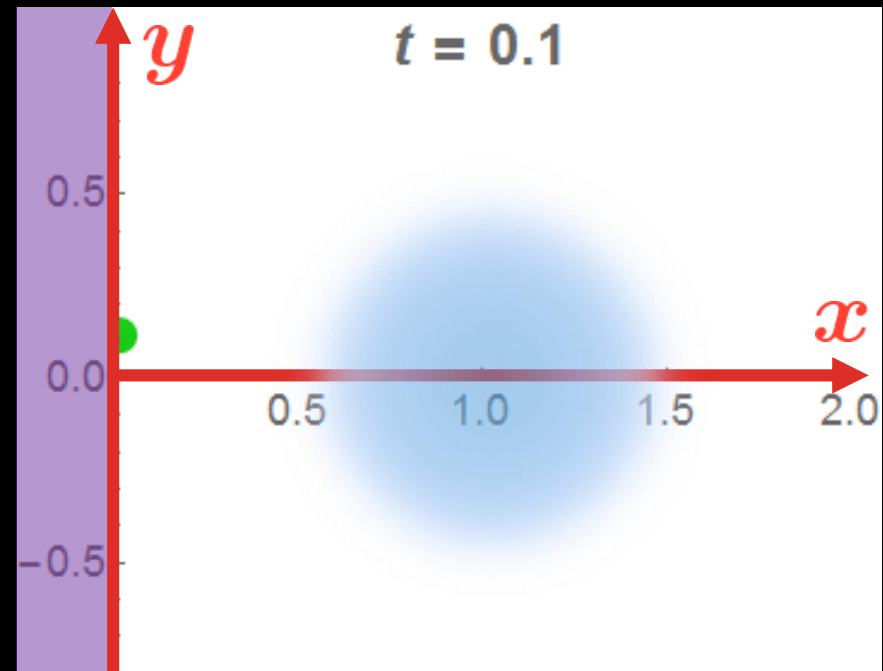
Numerical check

$$\mathcal{L} = -\sqrt{f(x) - \frac{\dot{x}^2}{f(x)} - \dot{y}^2} - \frac{\omega^2}{2} \left[(x - x_c)^2 + y^2 \right] \quad \left[f(x) \equiv 2\kappa x \right]$$

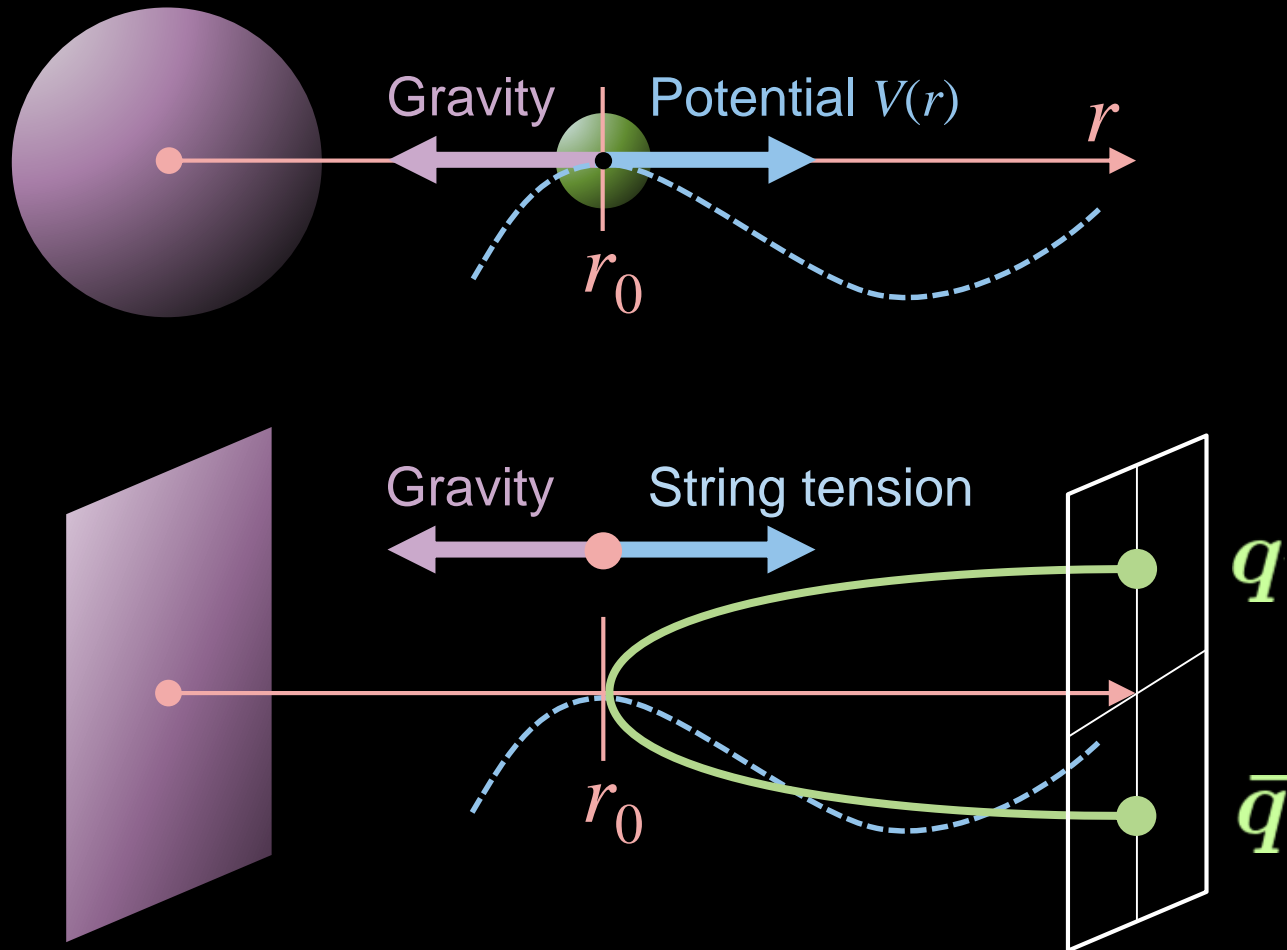
Particle **near Potential Minimum**



Particle **near BH Horizon**



Chaos Bound for AdS String



Fundamental string in AdS = “quark-anti quark pair”

Maldacena '98
Rey & Yee '98