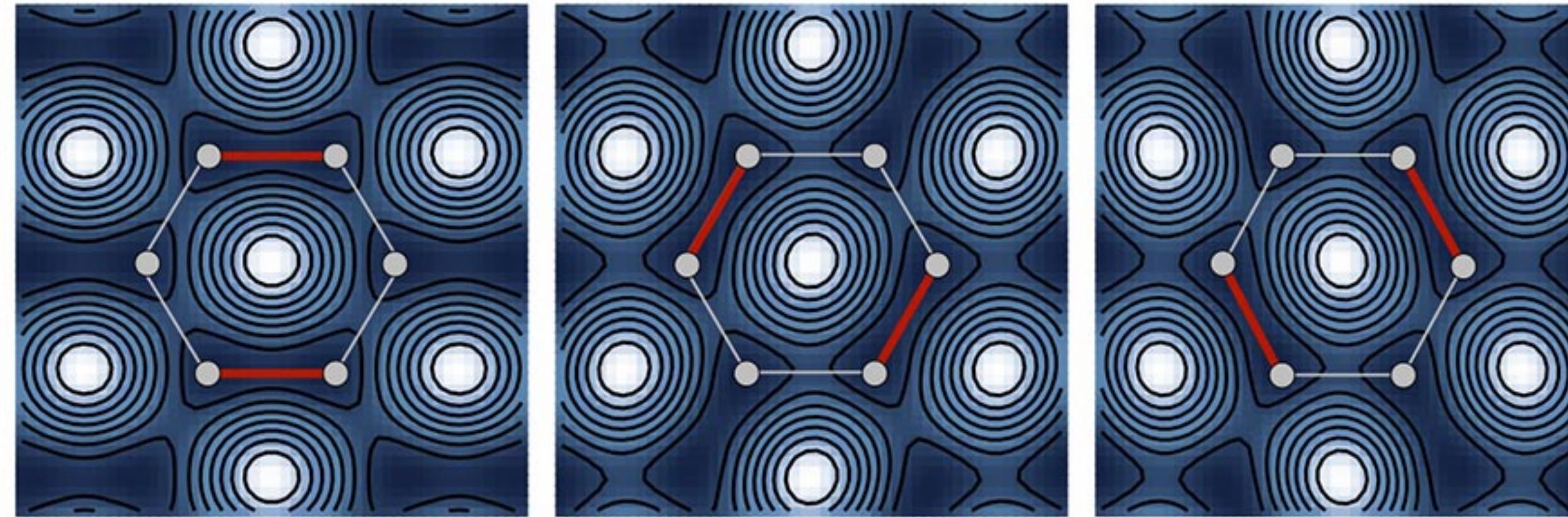
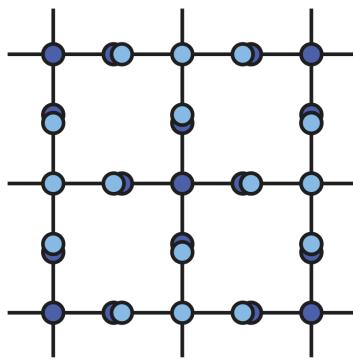


Anomalous Floquet topological systems with ultracold atoms



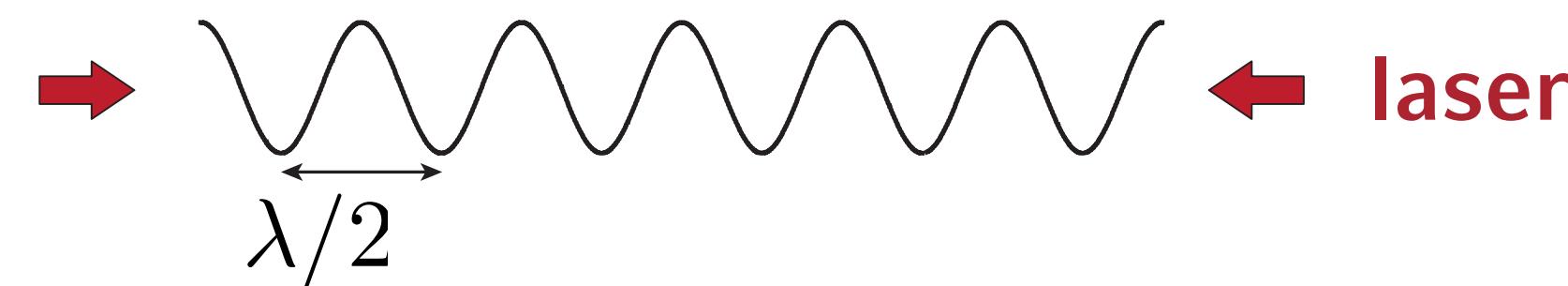
Monika Aidelsburger
Ludwig-Maximilians Universität München
Munich Center for Quantum Science & Technology



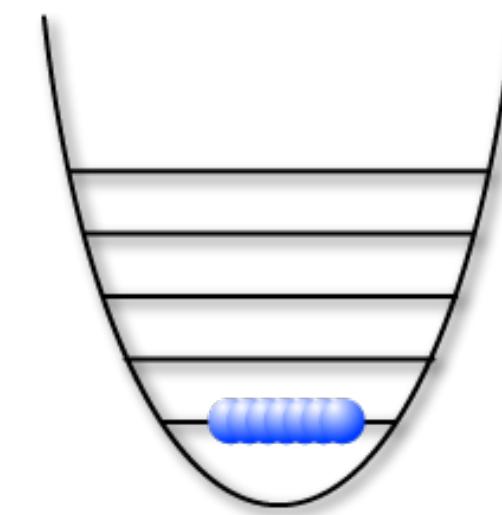
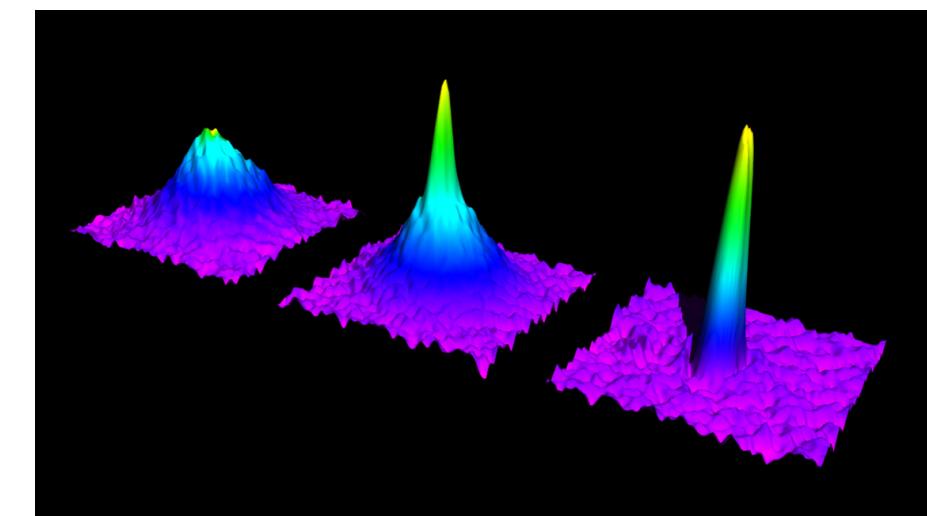
Ultracold atoms in optical lattices

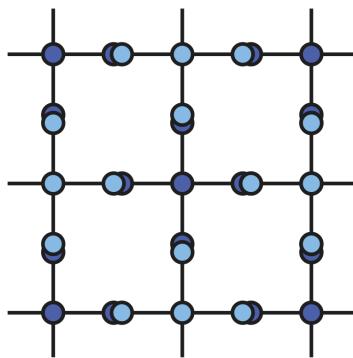
Quantum simulation with ultracold atoms in optical lattices:

- Atoms confined in periodic potentials



Bose-Einstein Condensate:

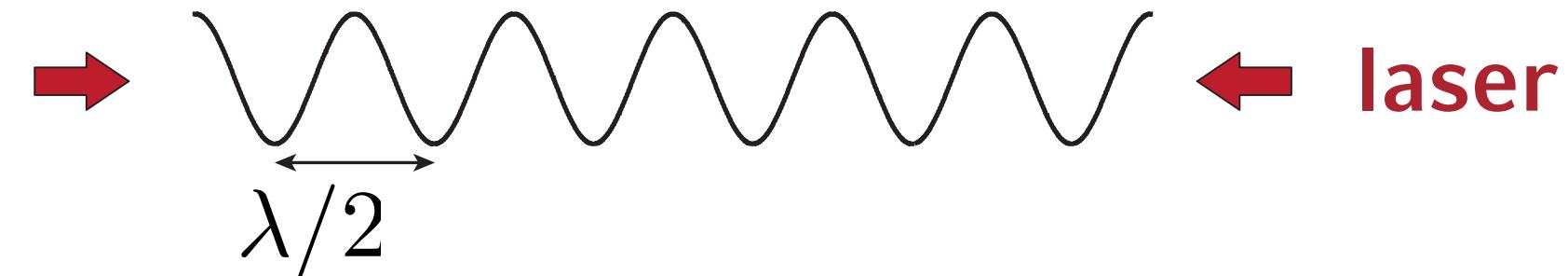




Ultracold atoms in optical lattices

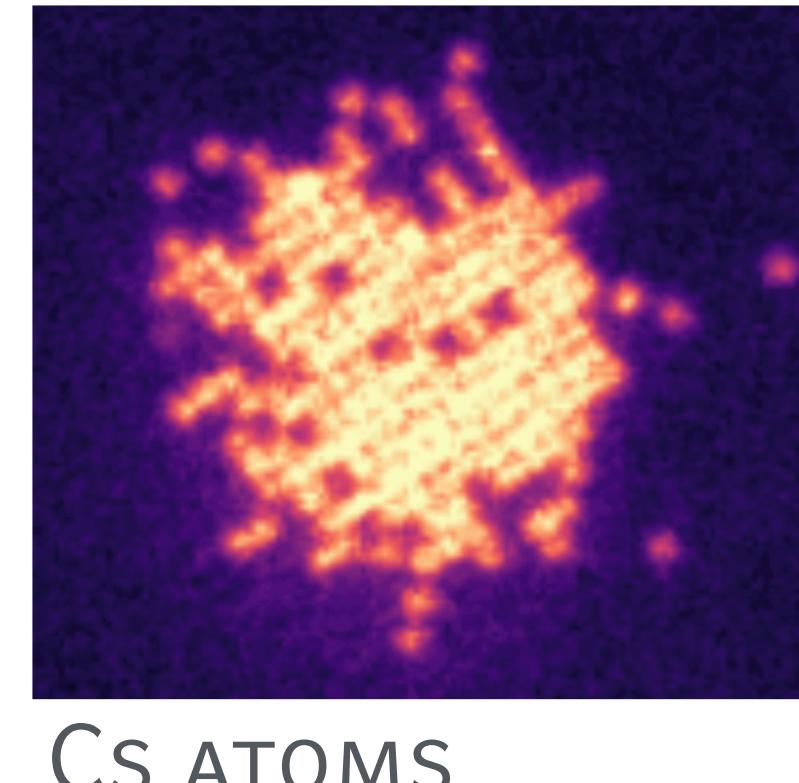
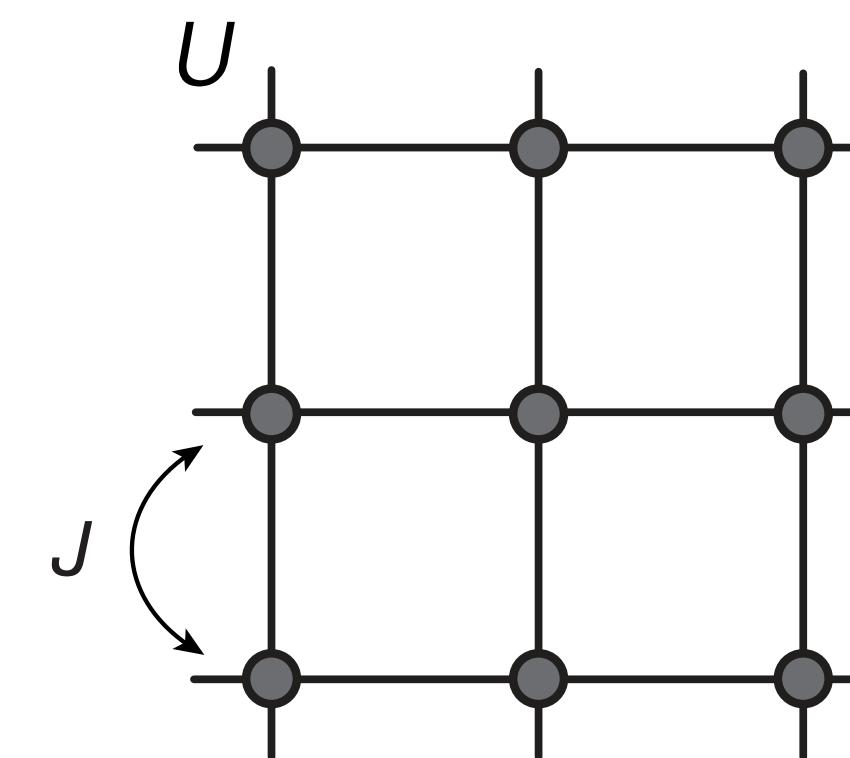
Quantum simulation with ultracold atoms in optical lattices:

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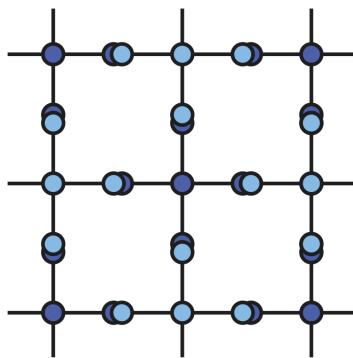


- Quantum simulation of Hubbard models

- Access to local observables using quantum gas microscopes

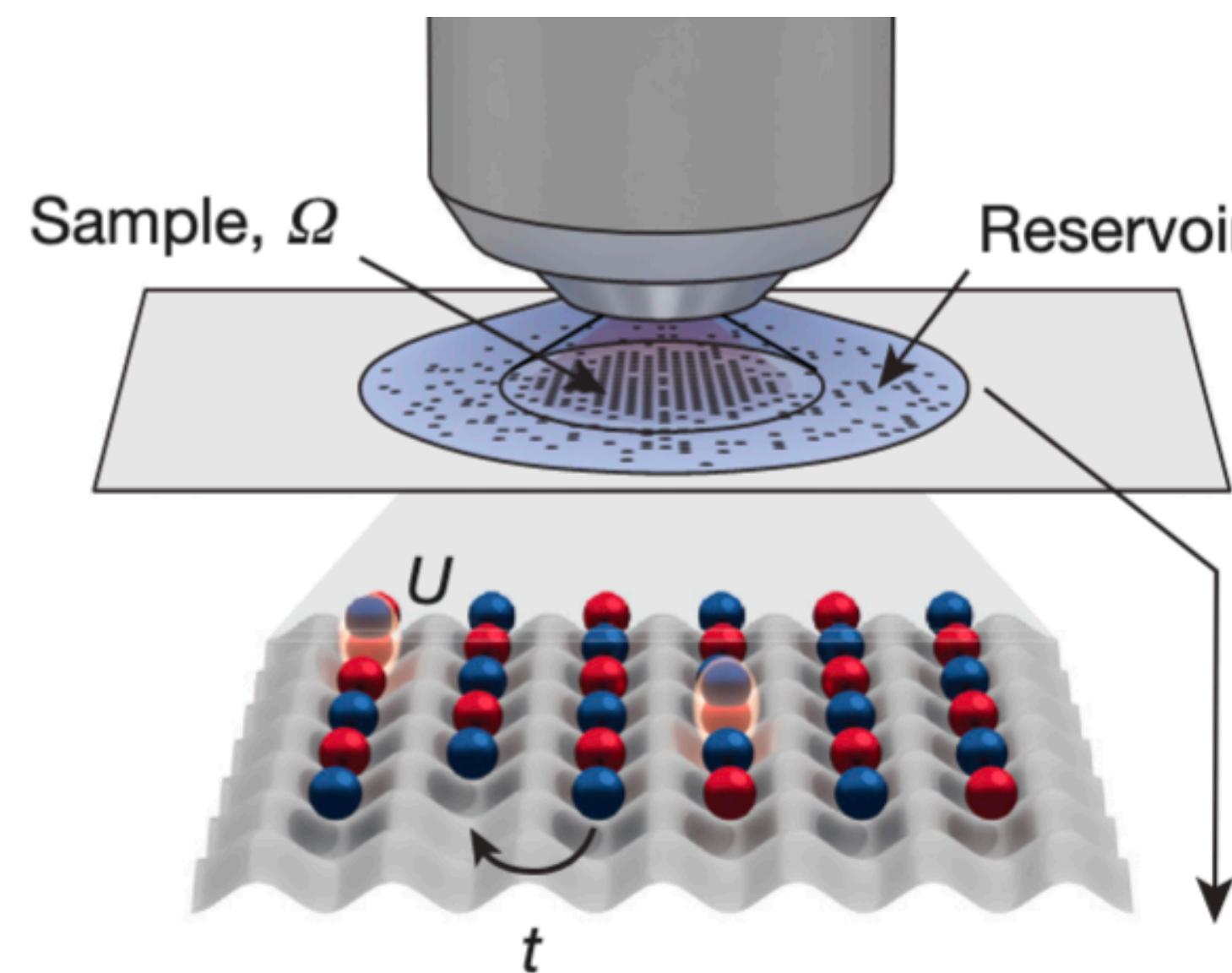


$$\hat{H} = -J \sum_{\langle i,j \rangle} \hat{a}_i^\dagger \hat{a}_j + \frac{U}{2} \sum_i \hat{n}_i (\hat{n}_i - 1)$$



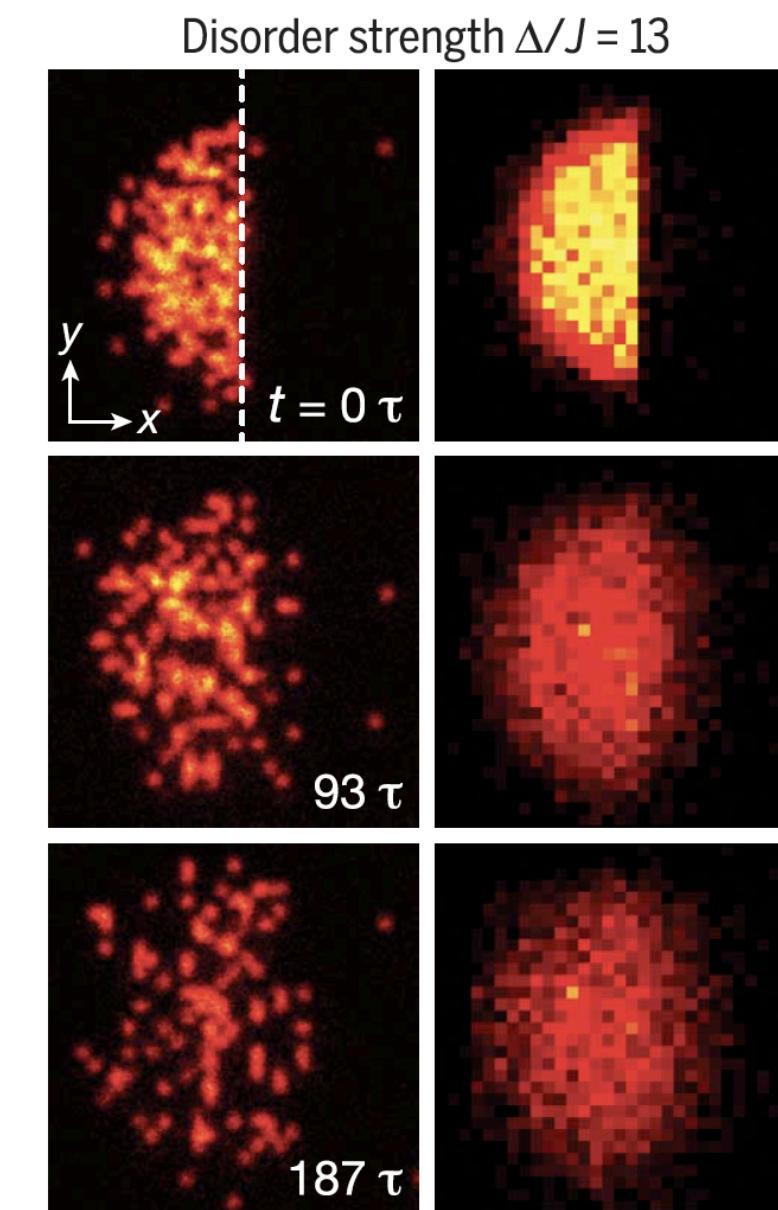
Ultracold atoms in optical lattices

Anti-ferromagnetic correlations in the Fermi-Hubbard model:

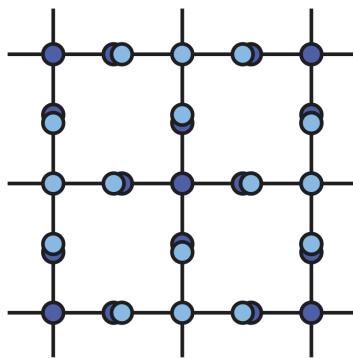


A. MARUZENKO, ... M. GREINER, NATURE (2017)

Thermalization of isolated quantum-many body systems:

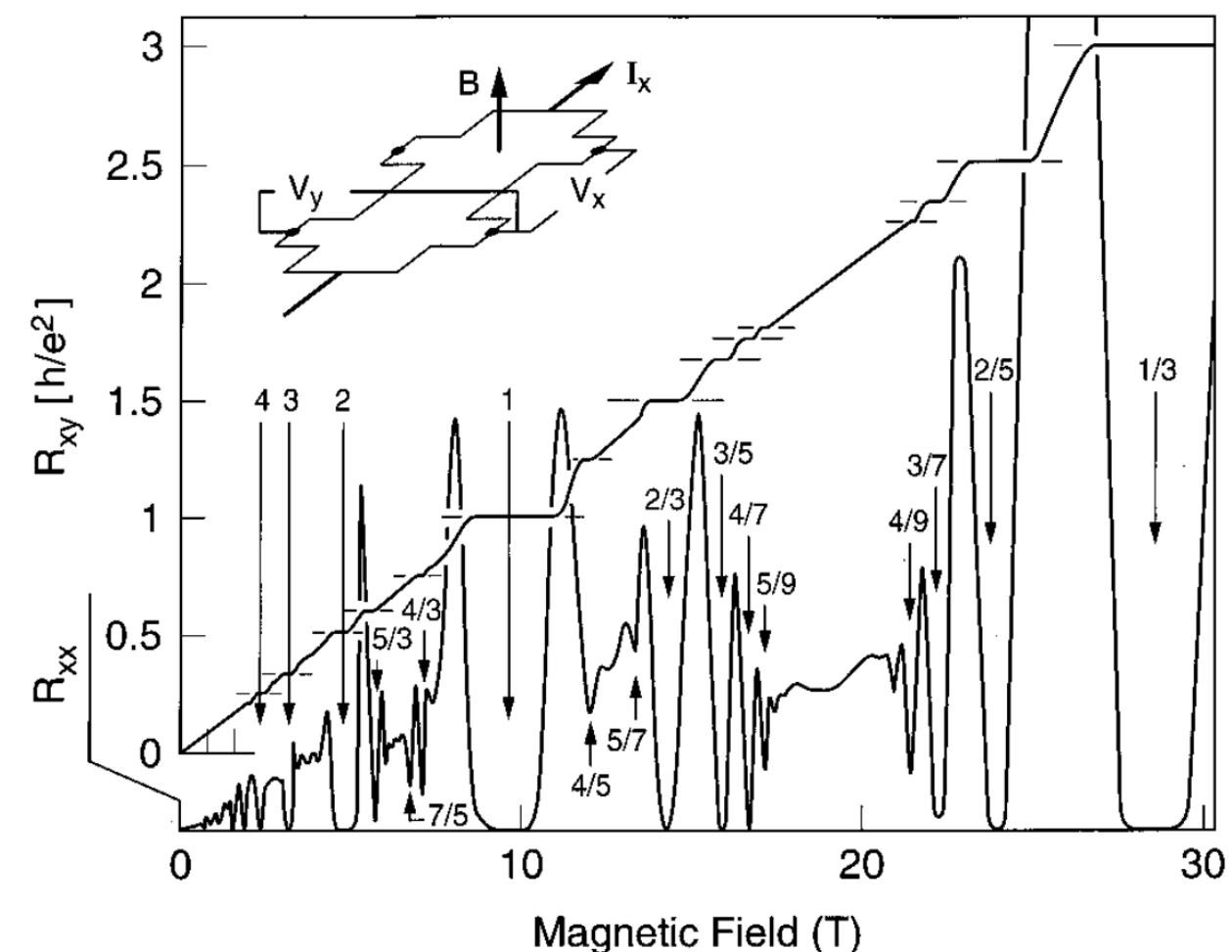


J. Y. CHOI, ..., I. BLOCH, SCIENCE 352, 1547 (2016)



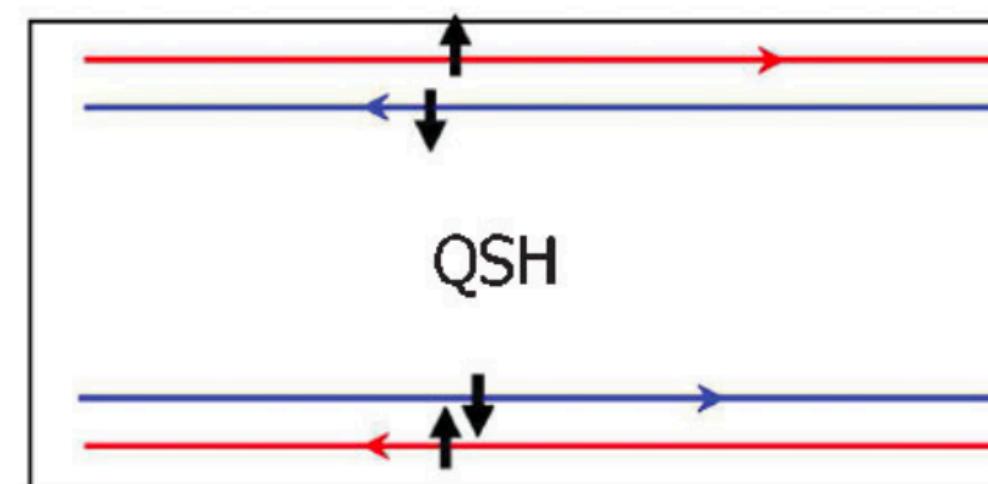
Topological phases of matter

Integer & fractional quantum Hall insulators



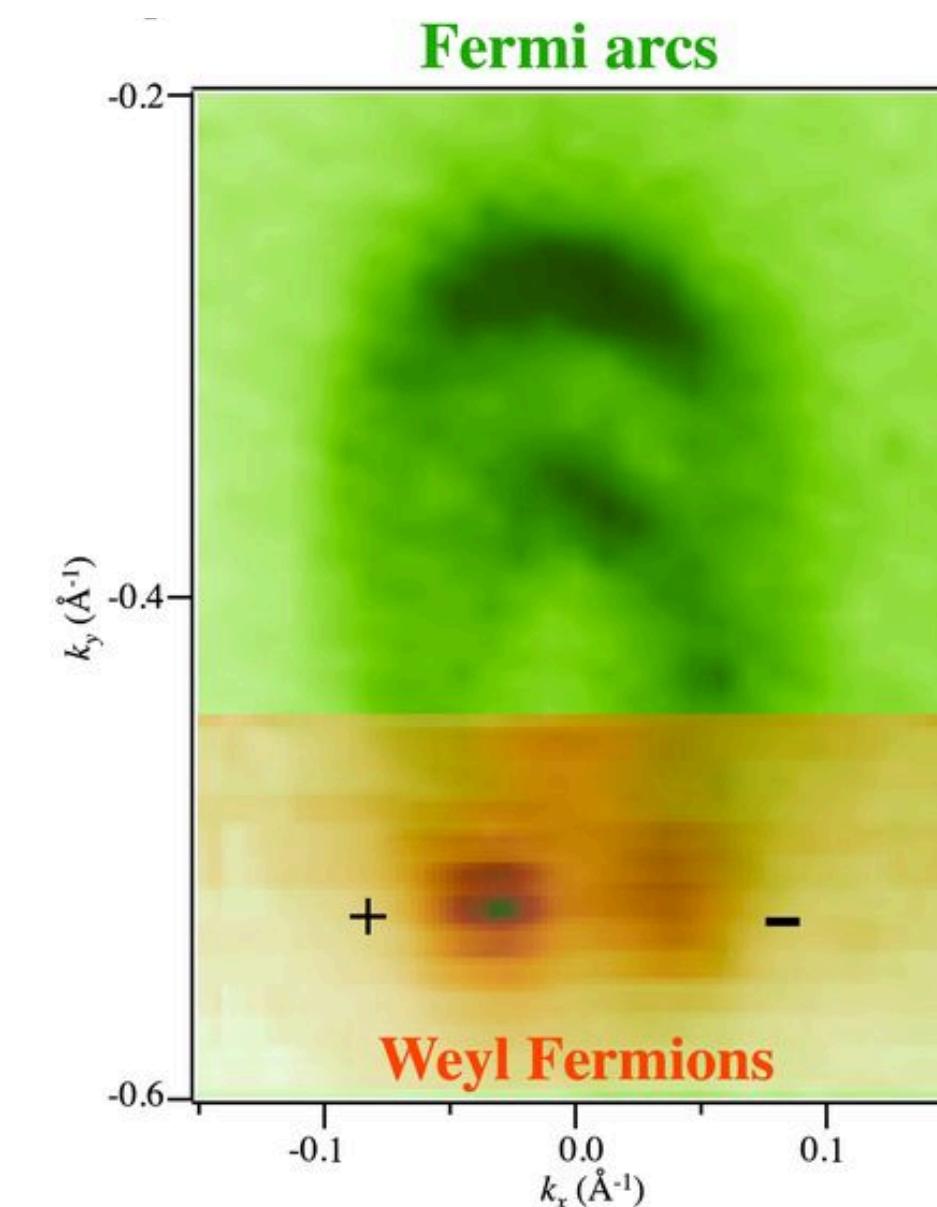
K. KLITZING, REV. MOD. PHYS. (1986)
STORMER ET AL., REV. MOD. PHYS. (1999)

Topological insulators in 2D & 3D



M. KÖNIG ET AL., SCIENCE (2007)
A. ROTH ET AL., SCIENCE (2009)

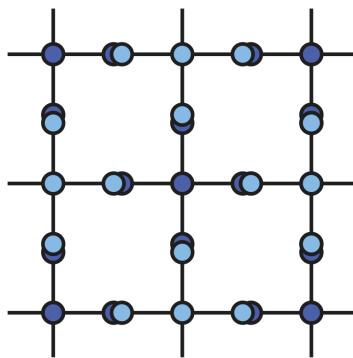
Weyl semimetals



L. LU ET AL., SCIENCE (2015)
S.-Y. XU ET AL., SCIENCE (2015)

...

How to engineer
topological phases?



Floquet engineering

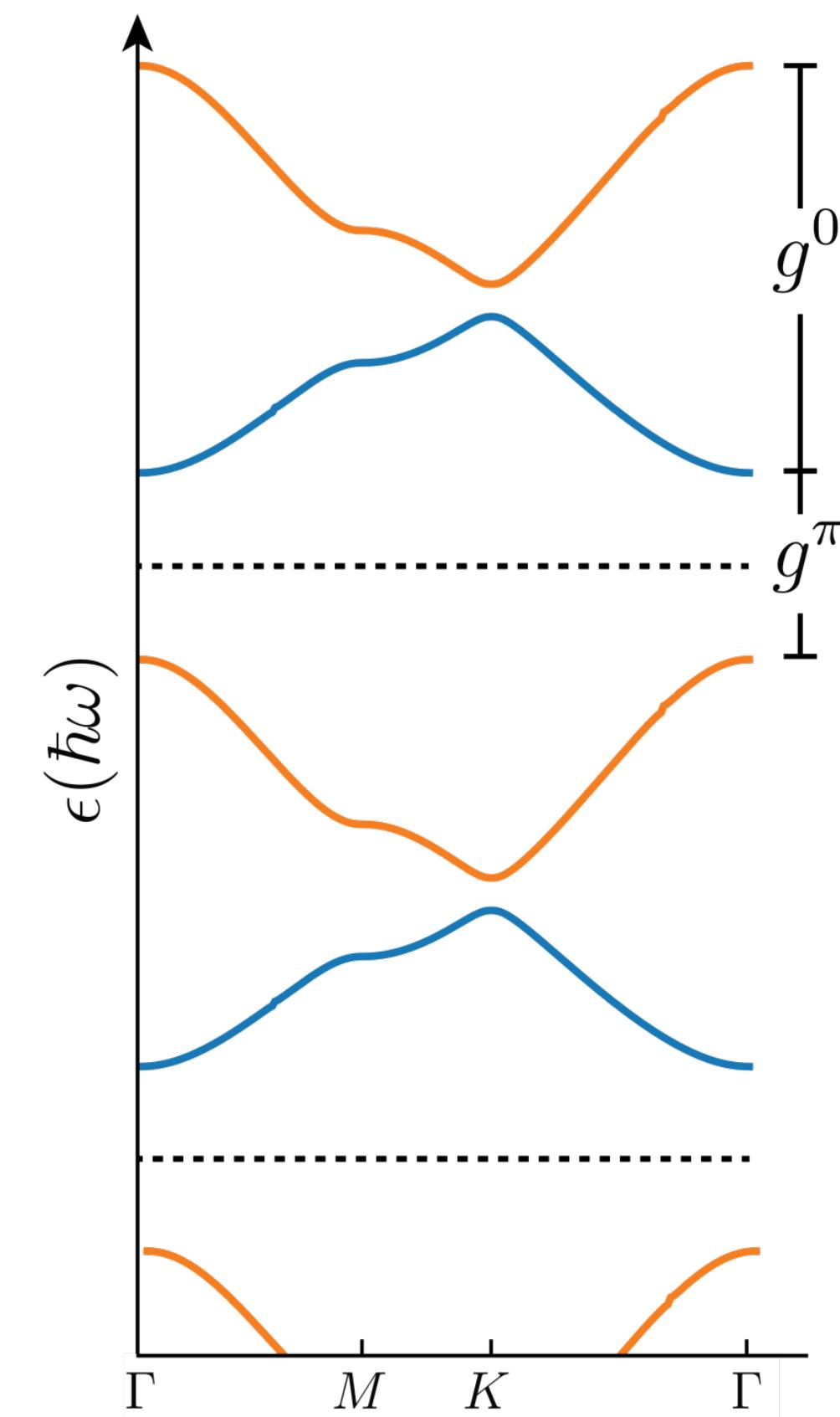
- Time-periodic driven Hamiltonian

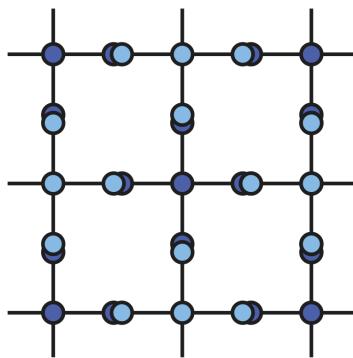
$$\hat{H}(t) = \hat{H}(t + T)$$

- Stroboscopic time evolution governed by effective Floquet Hamiltonian \hat{H}^F

$$\hat{U}(T, 0) = \exp\left(-\frac{i}{\hbar}T\hat{H}^F\right)$$

Engineer \hat{H}_F with topological properties!

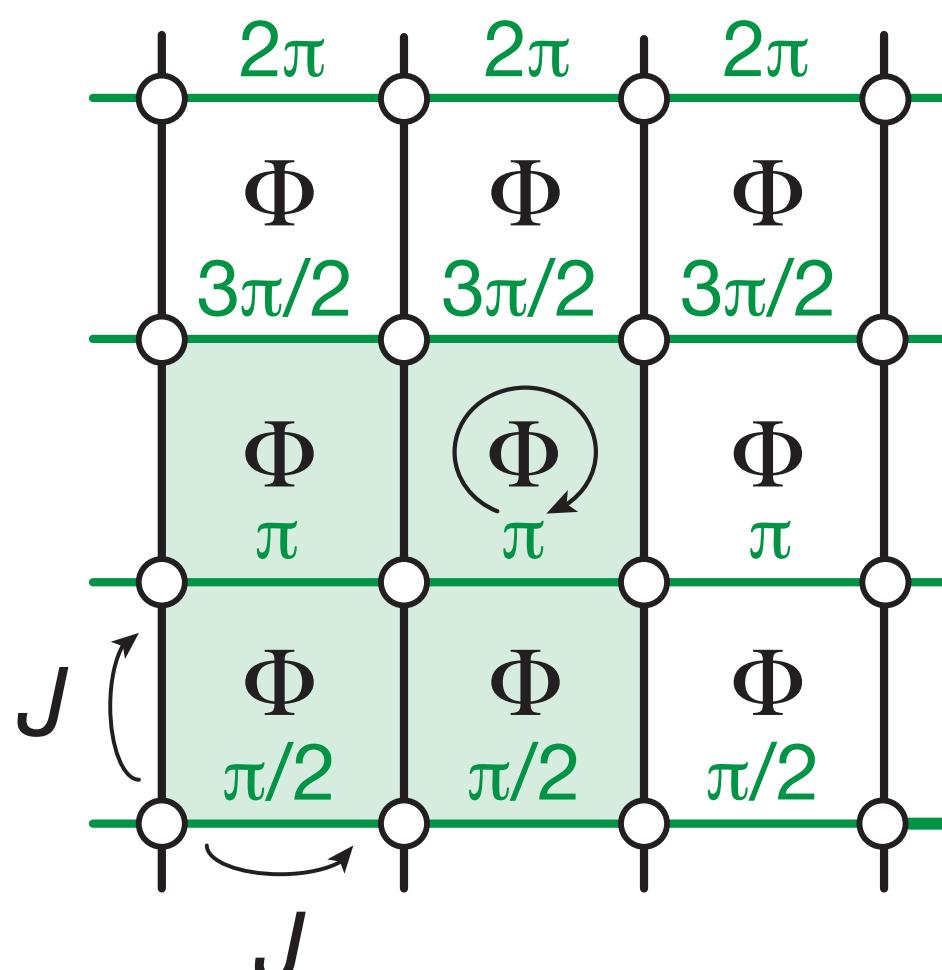




Topological lattice models

Hofstadter model

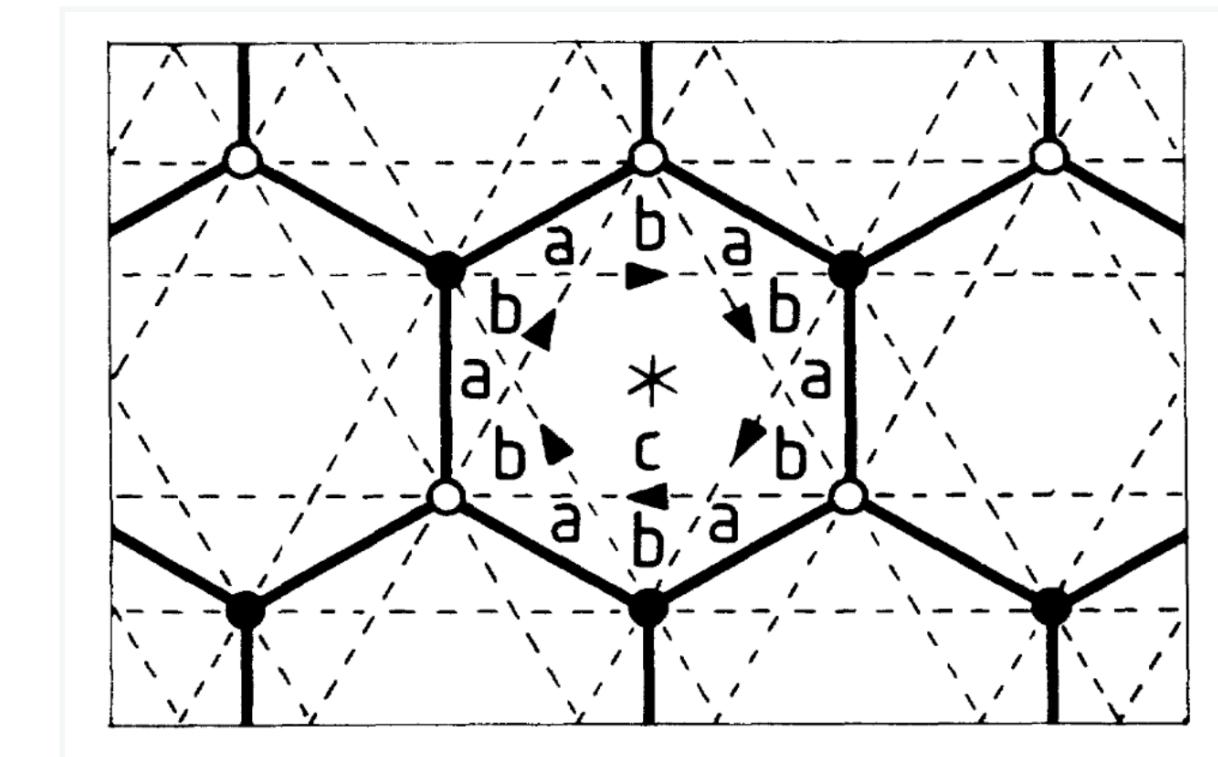
HARPER, PROC. PHYS. SOC., SECT.A **68**, 874 (1955)
 AZBEL, Zh. EKSP. TEOR. FIZ. **46**, 929 (1964)
 HOFSTADTER, PRB **14**, 2239 (1976)



$$\hat{H} = -J \sum_{m,n} \left(e^{in\Phi} \hat{a}_{m+1,n}^\dagger \hat{a}_{m,n} + \hat{a}_{m,n+1}^\dagger \hat{a}_{m,n} + \text{h.c.} \right)$$

Haldane model

HALDANE, PRL **61**, 2015 (1988)

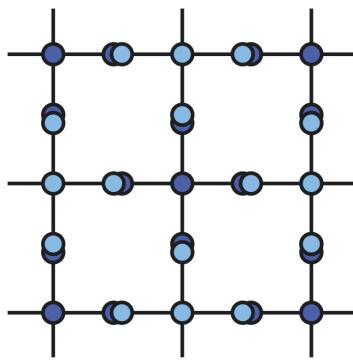


$$\hat{H} = \sum_{\langle ij \rangle} t_{ij} \hat{c}_i^\dagger \hat{c}_j + \sum_{\langle\langle ij \rangle\rangle} e^{i\Phi_{ij}} t'_{ij} \hat{c}_i^\dagger \hat{c}_j + \Delta_{AB} \sum_{i \in A} \hat{c}_i^\dagger \hat{c}_i$$

MA ET AL., PRL (2013); H. MIYAKE ET AL., PRL (2013)

E. M. TAI ET AL., NATURE (2017)

G. JOTZU ET AL., NATURE (2014) ; TARNOWSKI ET AL., NAT. COMM. (2019)



Topological invariants

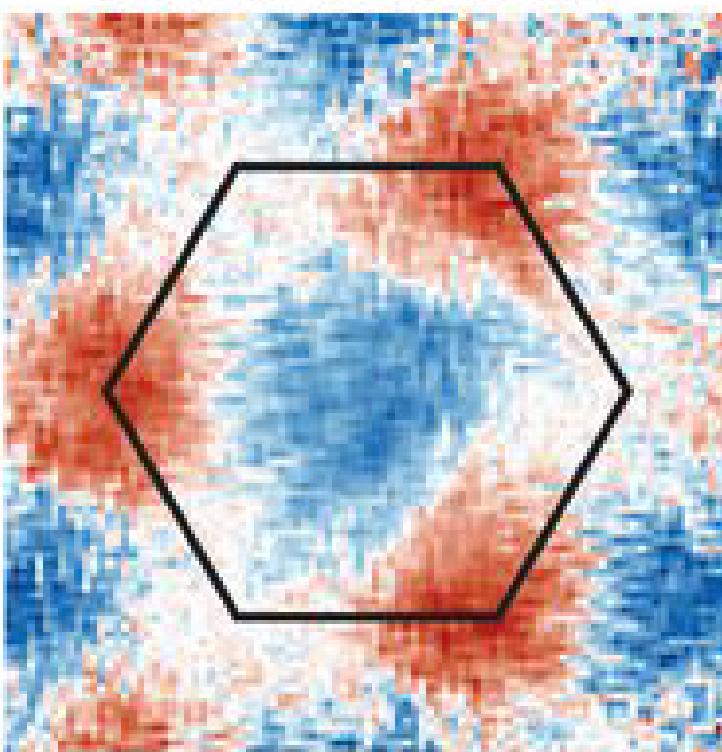
Chern number:

$$\mathcal{C}^\mu = \frac{1}{2\pi} \int_{\text{BZ}} \Omega^\mu d^2q$$

$|u_\mu(\mathbf{q})\rangle$: periodic Bloch function
 μ : band index

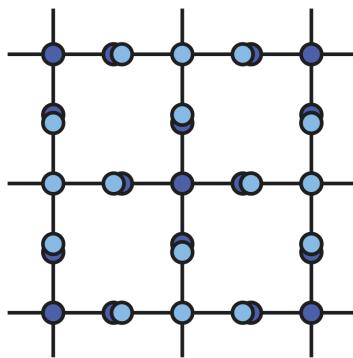
Berry curvature:

$$\Omega_\mu = i \left(\langle \partial_{q_x} u_\mu | \partial_{q_y} u_\mu \rangle - \langle \partial_{q_y} u_\mu | \partial_{q_x} u_\mu \rangle \right)$$



WEITENBERG/SENGSTOCK

M. ATALA, ET AL., NAT. PHYS. (2013); L. DUCA ET AL., SCIENCE (2015)
G. JOTZU ET AL., NATURE (2014); M. A. ET AL., NATURE PHYS. (2015)
N. FLÄSCHNER, SCIENCE (2016); T. LI, SCIENCE (2016)
TARNOWSKI ET AL., NAT. COMM. (2019);
L. ASTERIA ET AL., NAT. PHYS. (2019);
B. REM ET AL., NAT. PHYS. (2019);

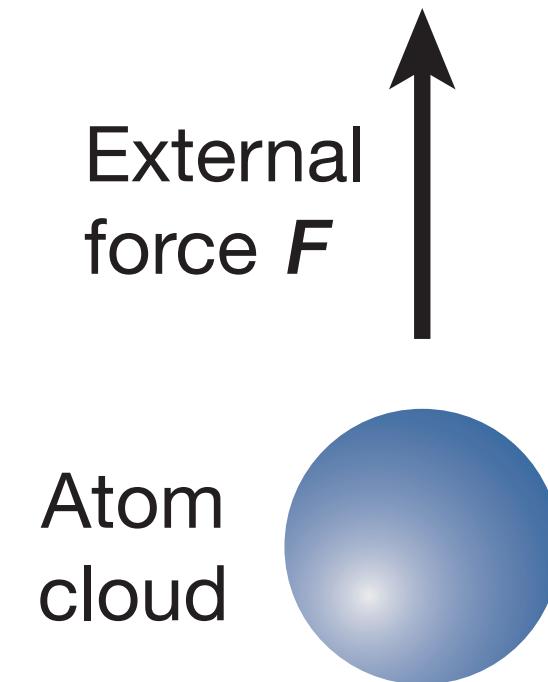


Transport measurements

Semiclassical dynamics

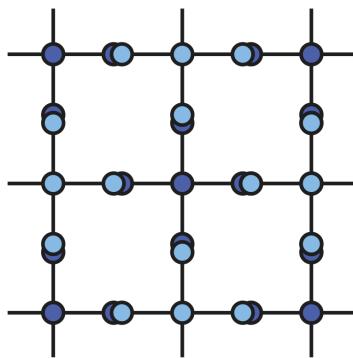
- *Bloch oscillations* captured by *band velocity*

$$\mathbf{v}_\mu^{\text{band}} = \frac{1}{\hbar} \partial_{\mathbf{q}} E_\mu \quad \mu: \text{band index}$$



- *Anomalous transverse velocity*
proportional to *Berry curv.* Ω_μ

$$\mathbf{v}_\mu^x(\mathbf{q}) = -\frac{F}{\hbar} \Omega_\mu(\mathbf{q})$$



Transport measurements

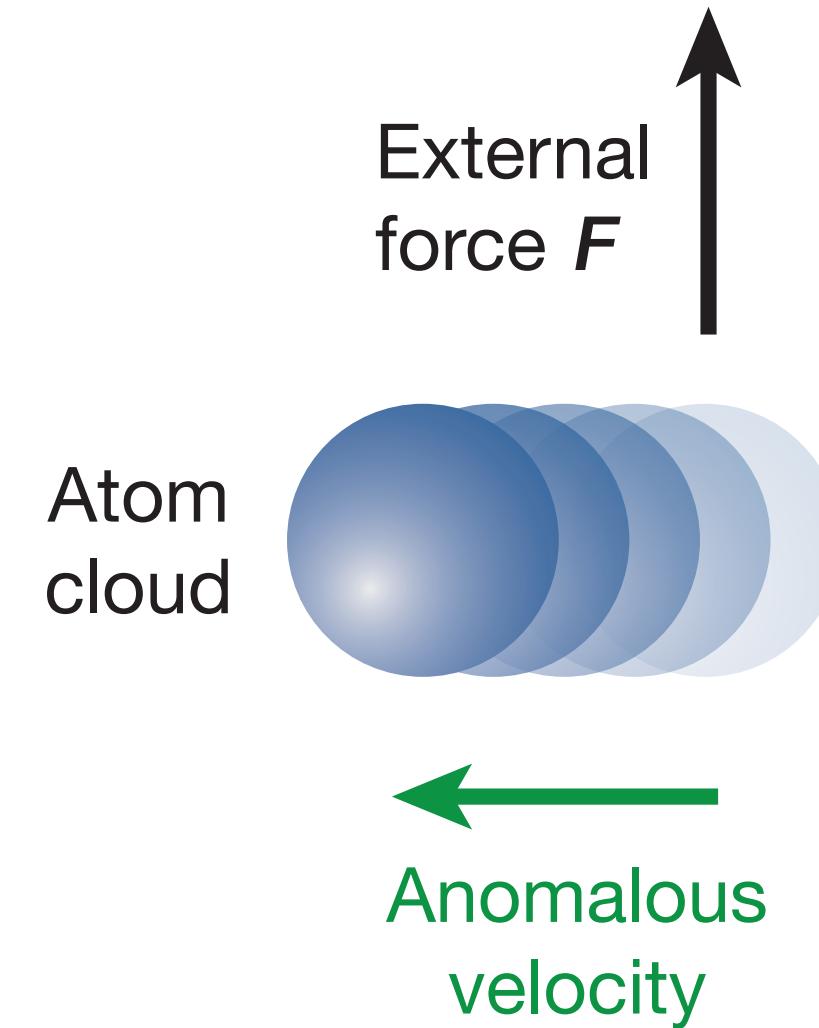
Semiclassical dynamics

- *Bloch oscillations* captured by *band velocity*

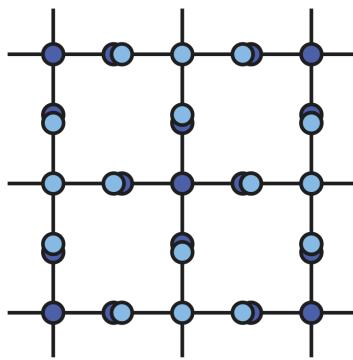
$$\mathbf{v}_\mu^{\text{band}} = \frac{1}{\hbar} \partial_{\mathbf{q}} E_\mu \quad \mu: \text{band index}$$

- *Anomalous transverse velocity*
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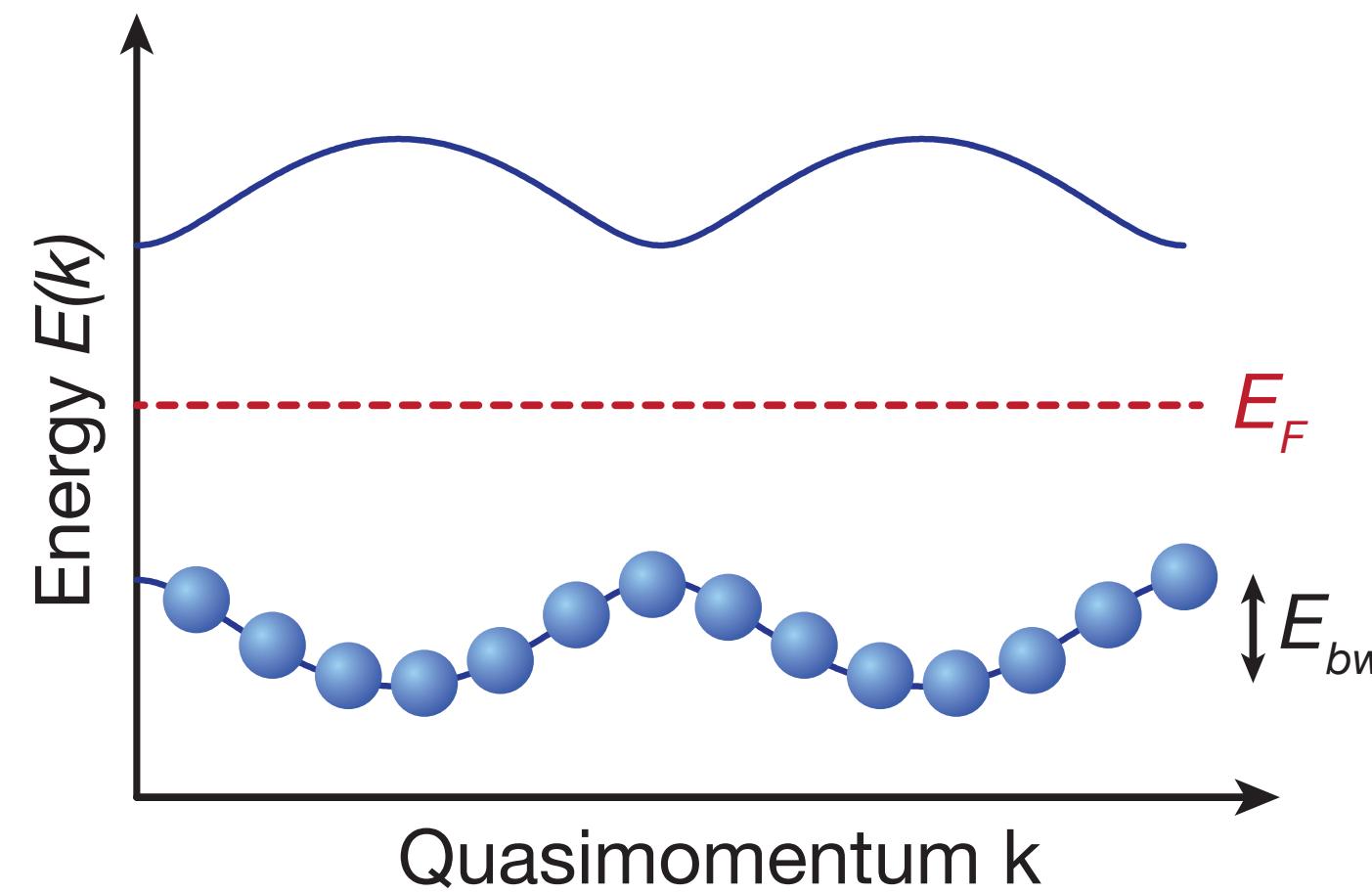
$$\mathbf{v}_\mu^x(\mathbf{q}) = -\frac{F}{\hbar} \Omega_\mu(\mathbf{q})$$



Prepare condensate at momentum \mathbf{q}
and *probe Berry curvature locally!*



Hall deflection

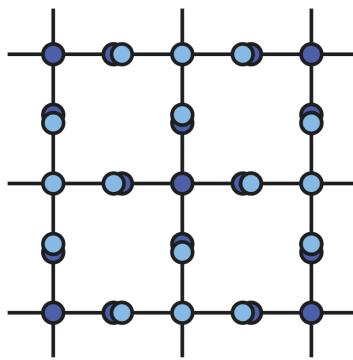


Contribution from band velocity vanishes

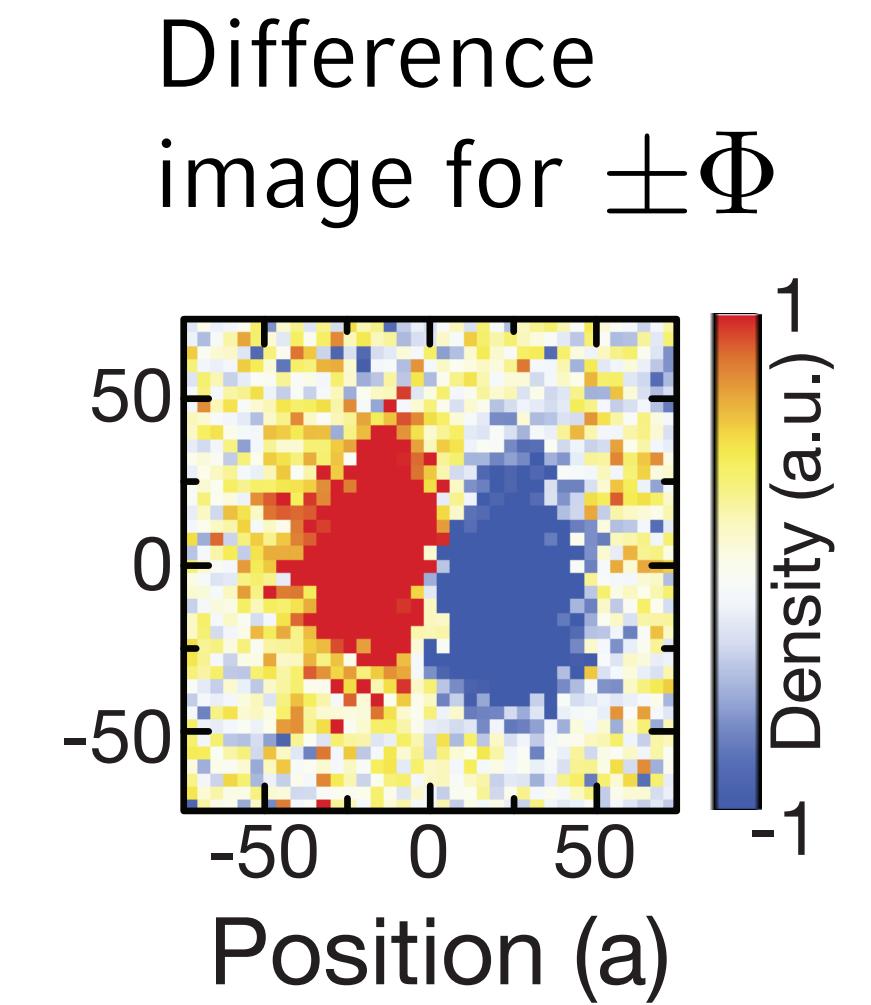
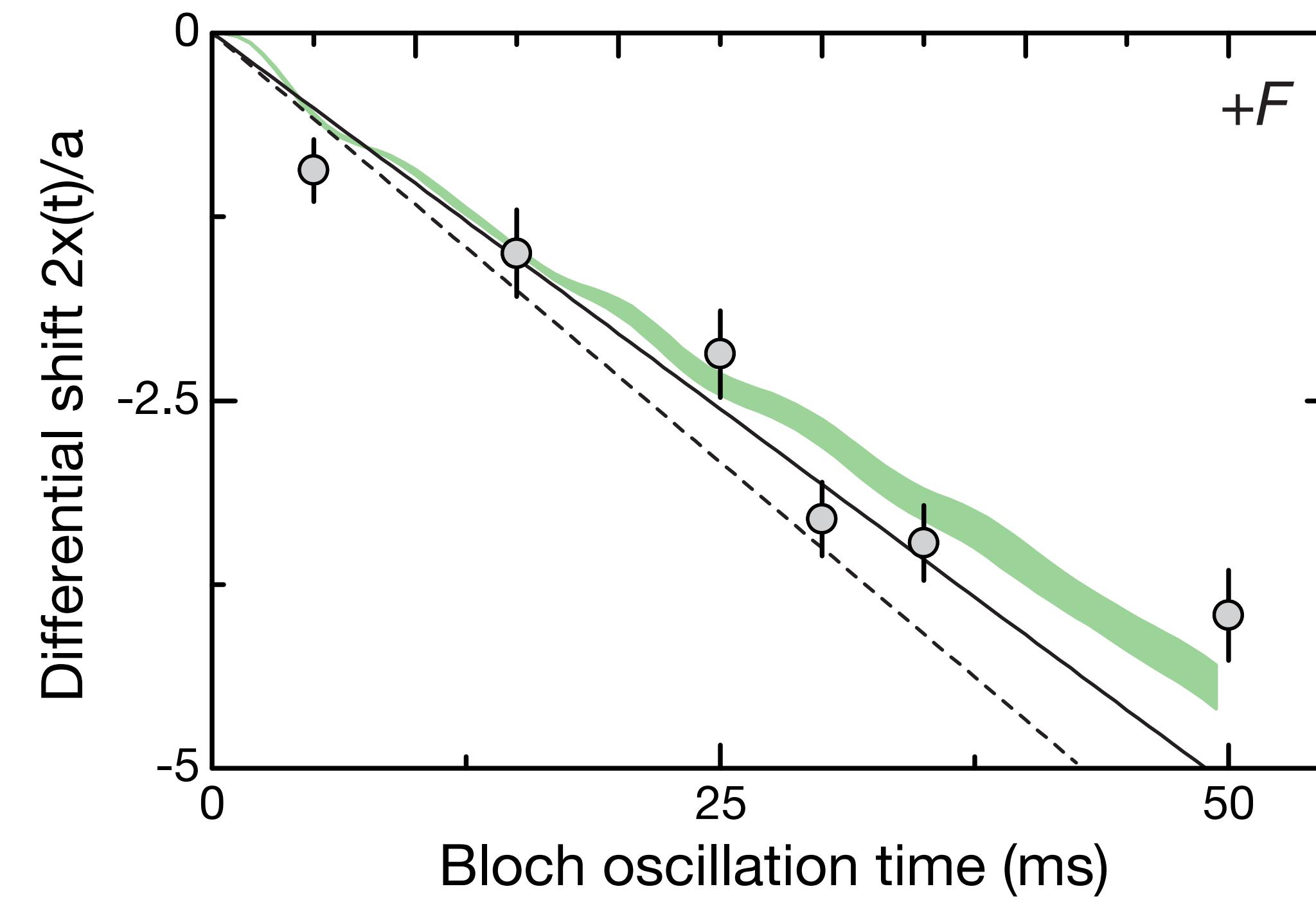
$$\int \partial E_\mu / \partial k_{x,y} d^2k = 0$$

Transverse center-of-mass motion (μ th band):

$$x_\mu(t) = -\frac{4a^2 F}{h} \nu_\mu t$$

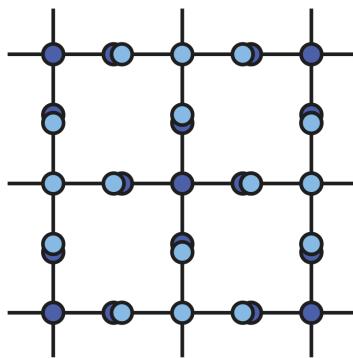


Hall deflection



$$\nu_{\text{exp}} = 0.99(5)$$

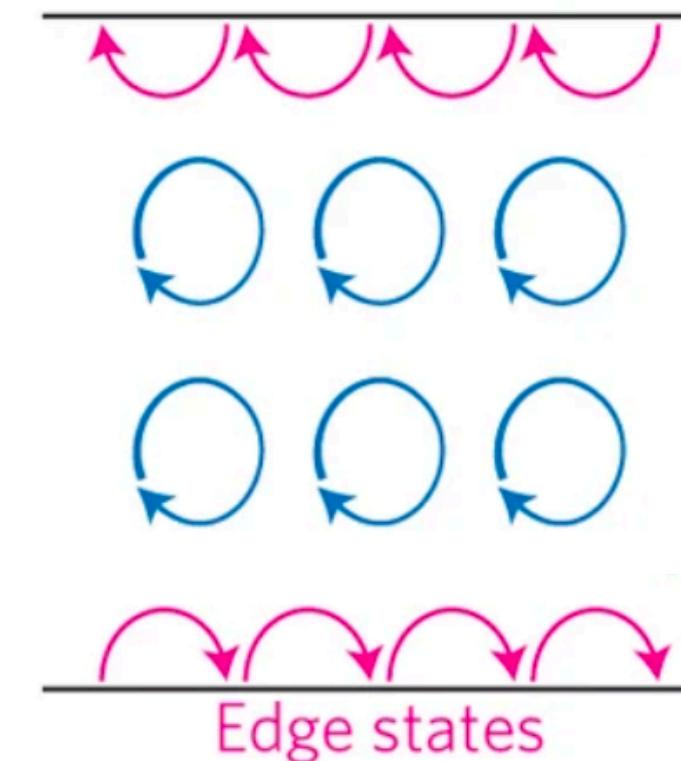
Can we go beyond emulating
static topological systems?



Anomalous Floquet phases

Floquet engineering:

- High-frequency limit:
→ known **bulk-edge correspondence**
- Chern number uniquely determines
net edge modes in the gap

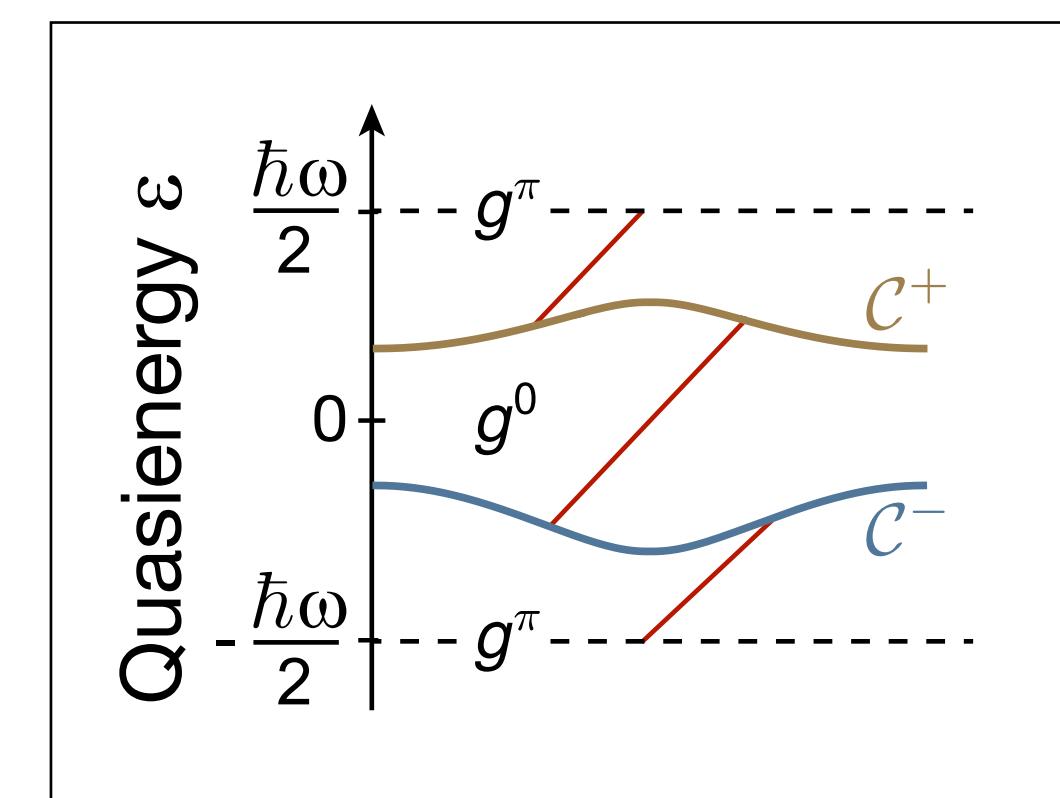


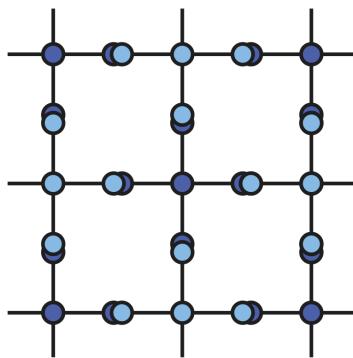
Beyond high-frequency limit:

T. KITAGAWA ET AL., PRB 82, 235114 (2010)

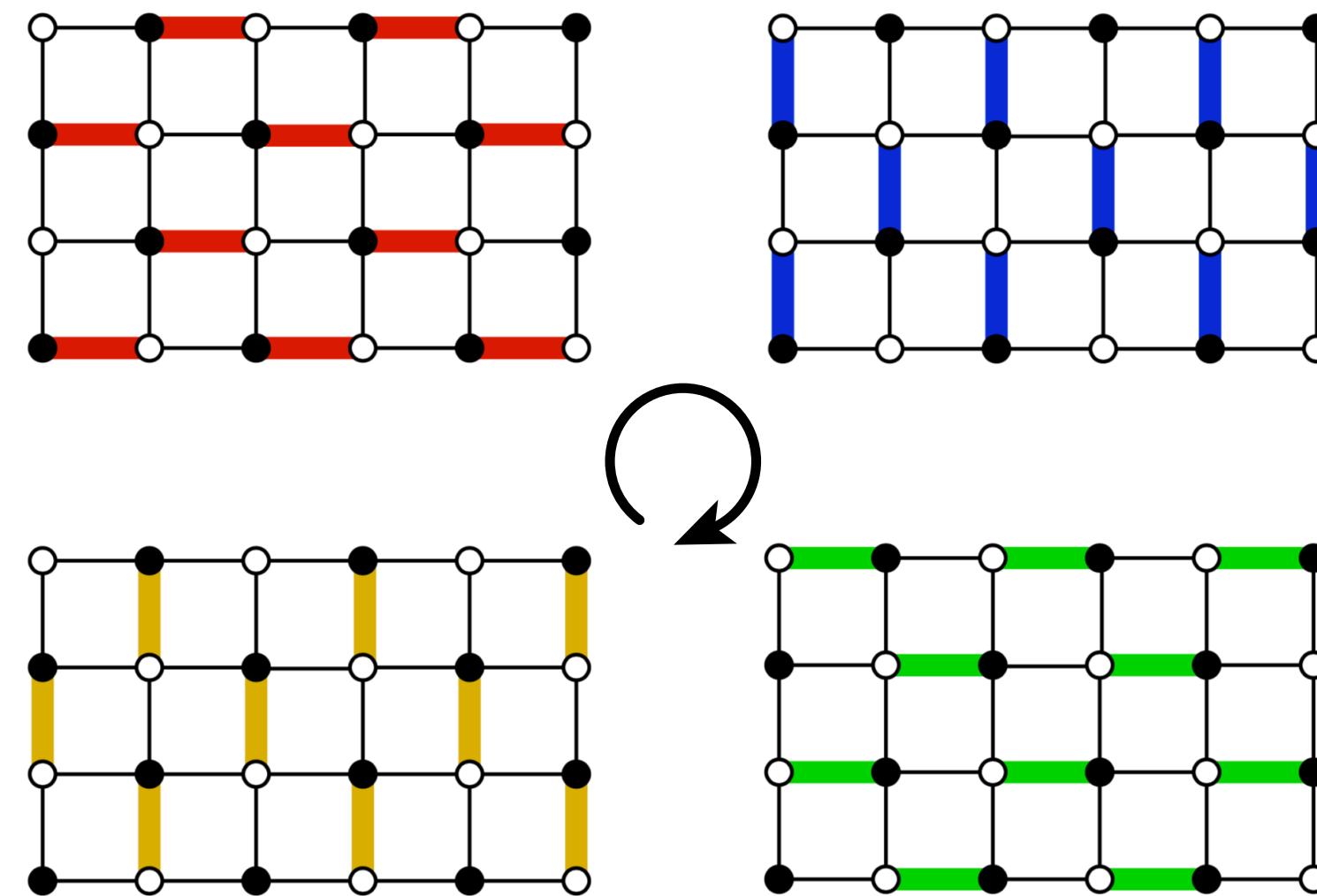
M. RUDNER ET AL., PRX 3, 031005 (2013)

- Edge modes, but **zero** Chern numbers!
- Robust edge transport \longleftrightarrow localized bulk!

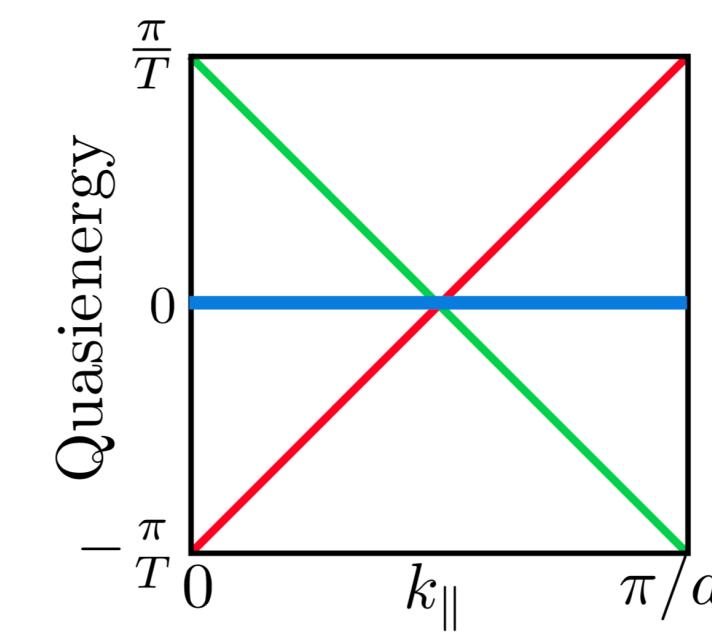
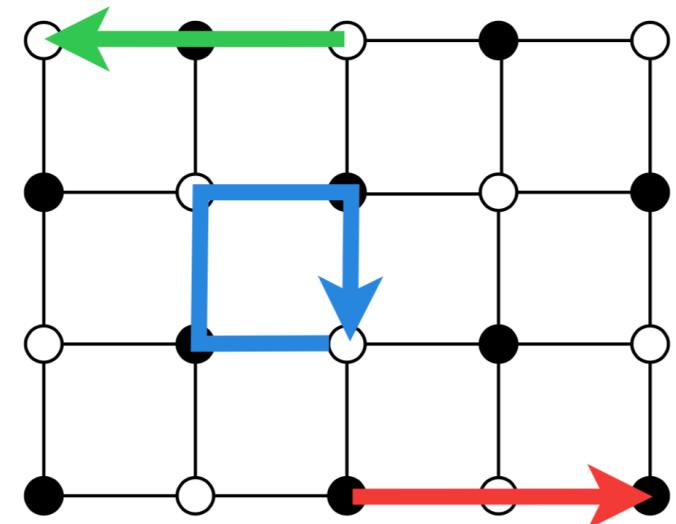




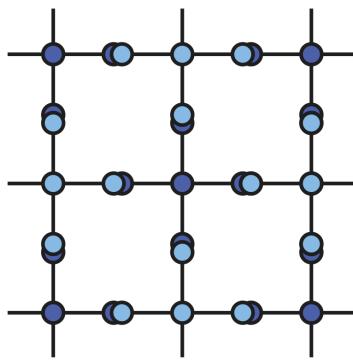
Stepwise modulated square lattice



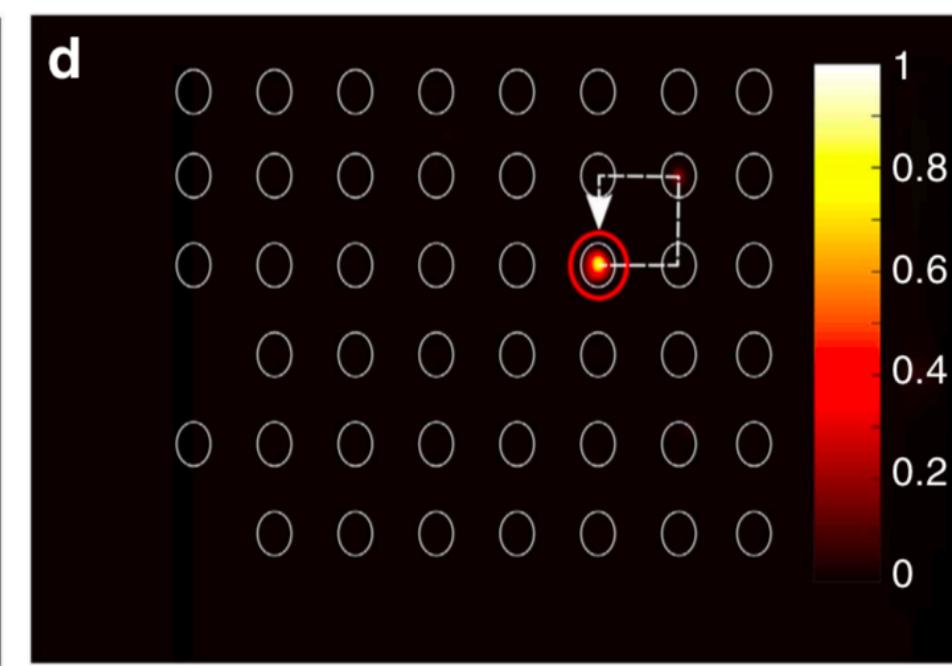
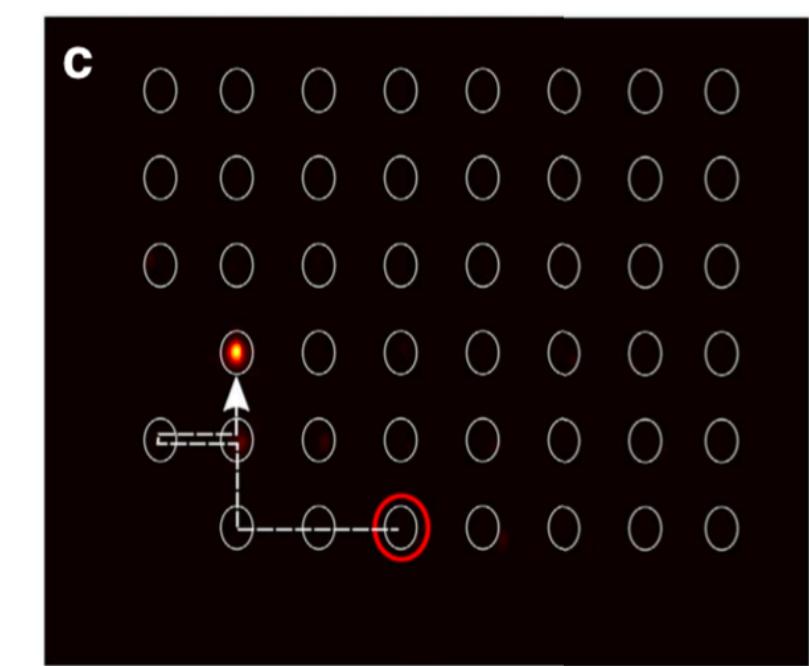
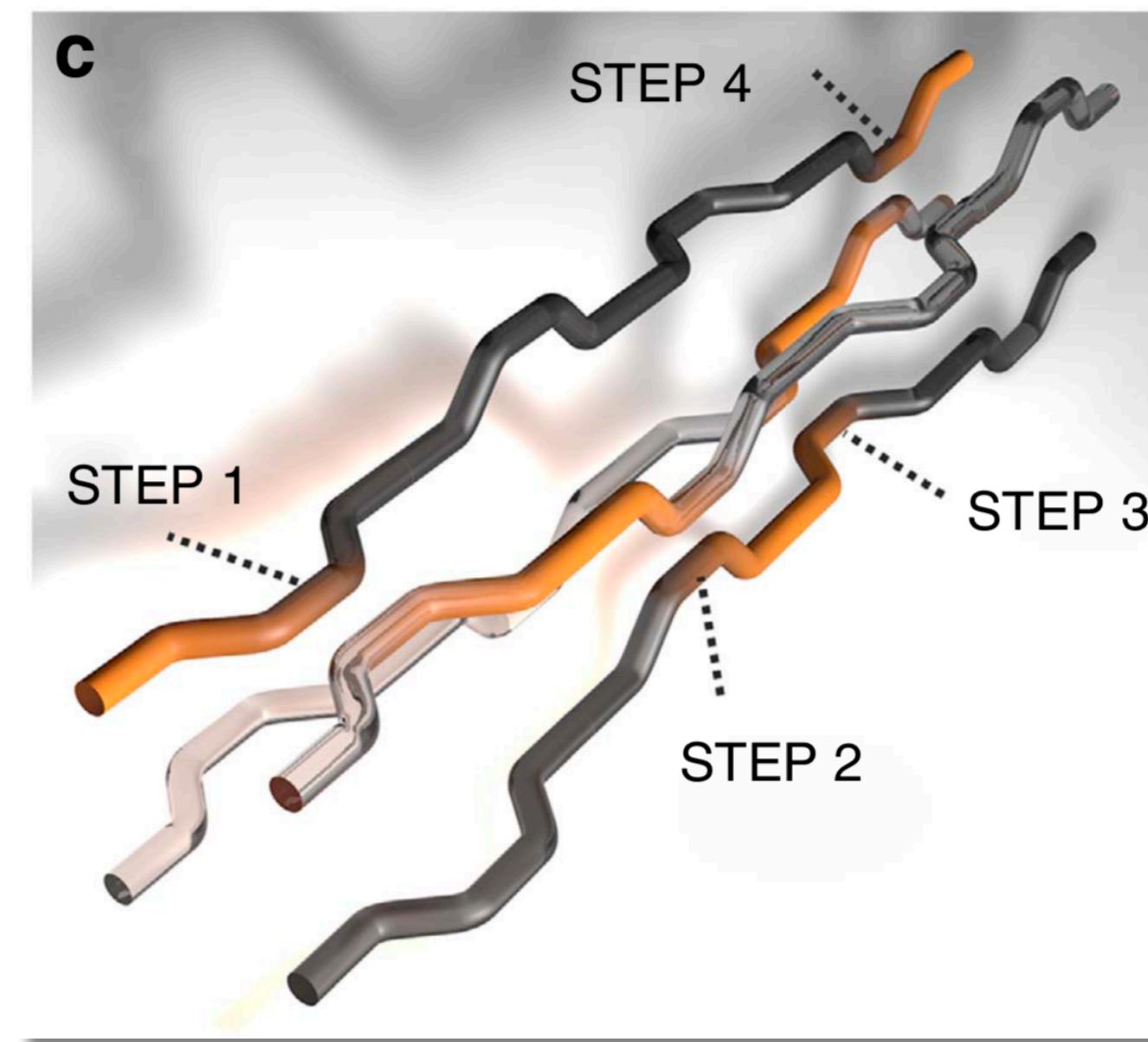
→ two degenerate flat bands
at $\varepsilon=0$, with *trivial topology*!



- Topology not determined by properties of \hat{H}_{eff} !
- Full time-dependence is important
- Non-trivial winding of quasi-energy spectrum



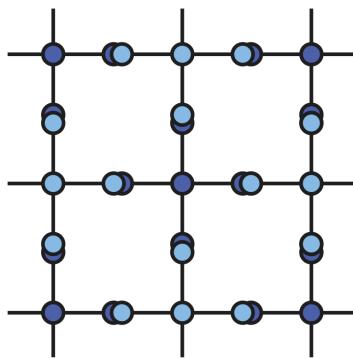
Coupled waveguide arrays



→ Robust edge transport

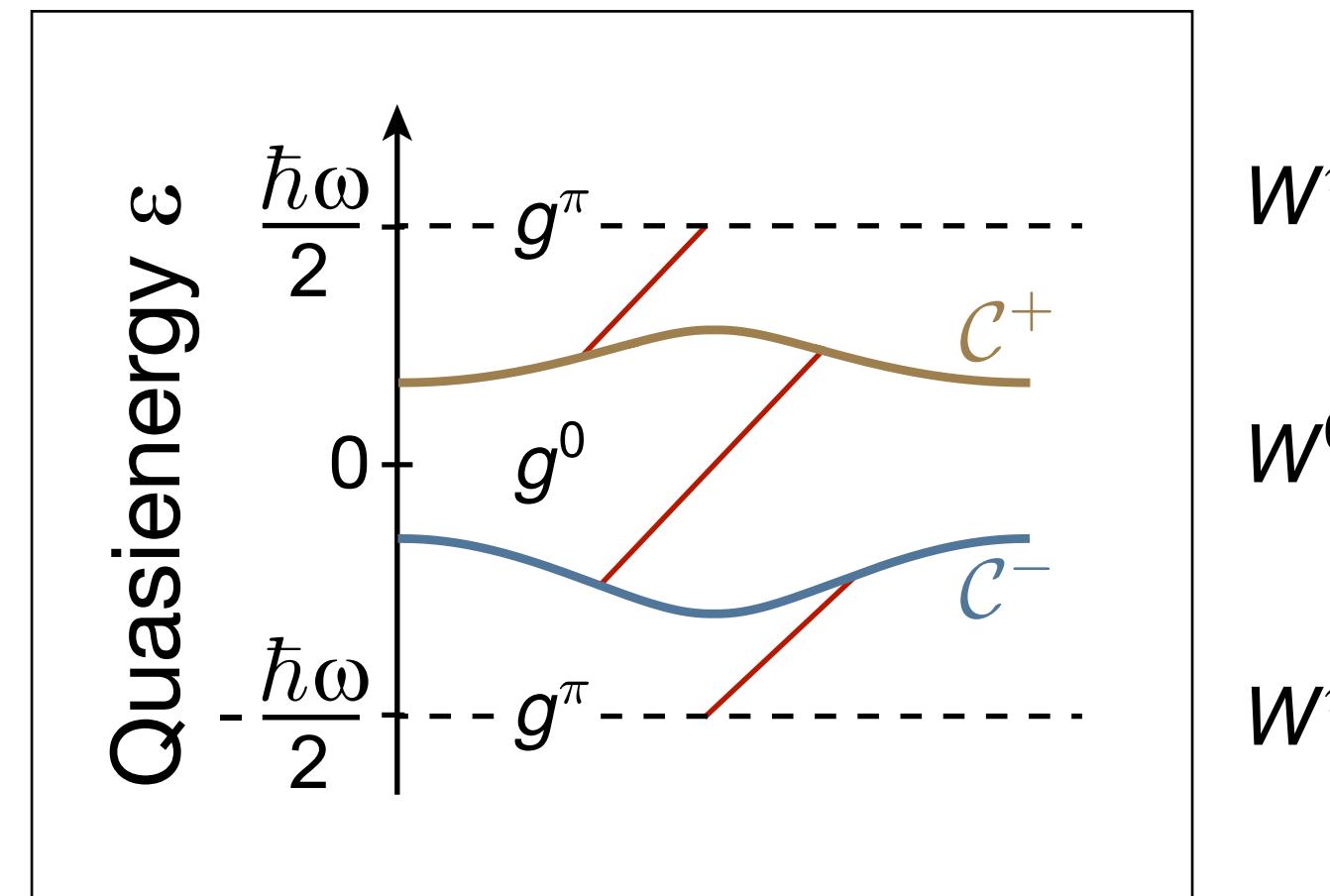
Cold atoms:

- Reveal *bulk and edge-state physics*
- Study *many-body phases*



Topological Floquet phases

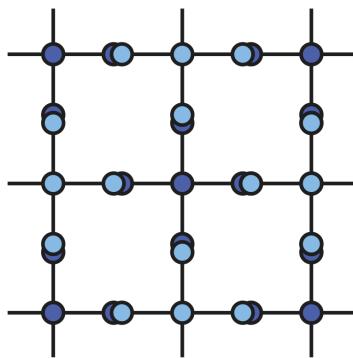
How to characterize the topological properties?



*Winding number counts
edge modes in gap!*

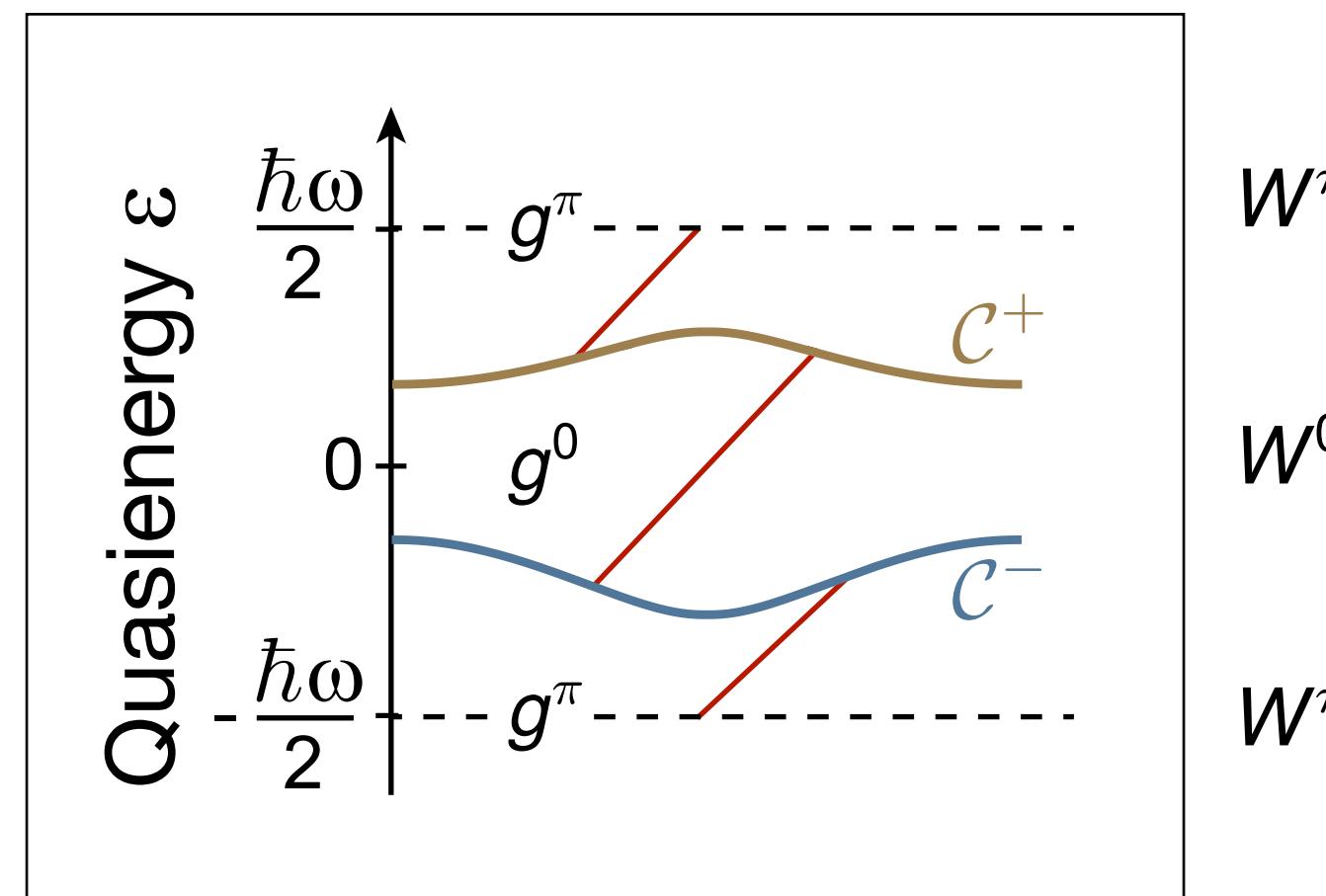
$$C^\pm = \mp(W^0 - W^\pi)$$

$$W[U] = \frac{1}{8\pi^2} \int dt dq_x dq_y \times \text{Tr}(U^{-1} \partial_t U [U^{-1} \partial_{q_x} U, U^{-1} \partial_{q_y} U])$$



Topological Floquet phases

How to characterize the topological properties?

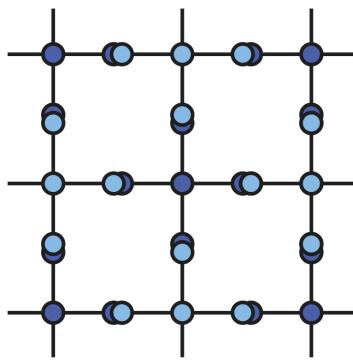


*Winding number counts
edge modes in gap!*

$$C^\pm = \mp(W^0 - W^\pi)$$

Complete set of invariants given by winding numbers!

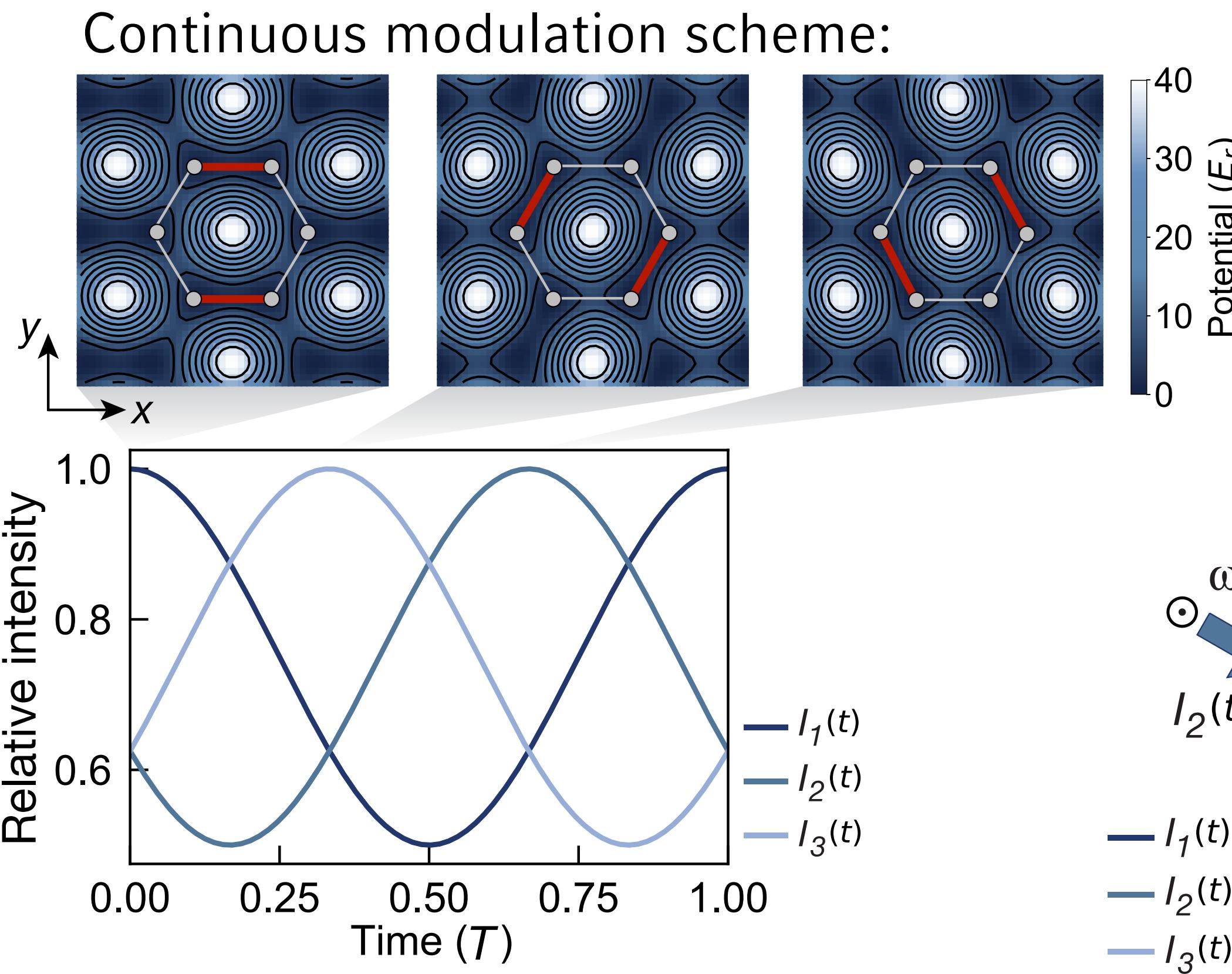
Experimental realization



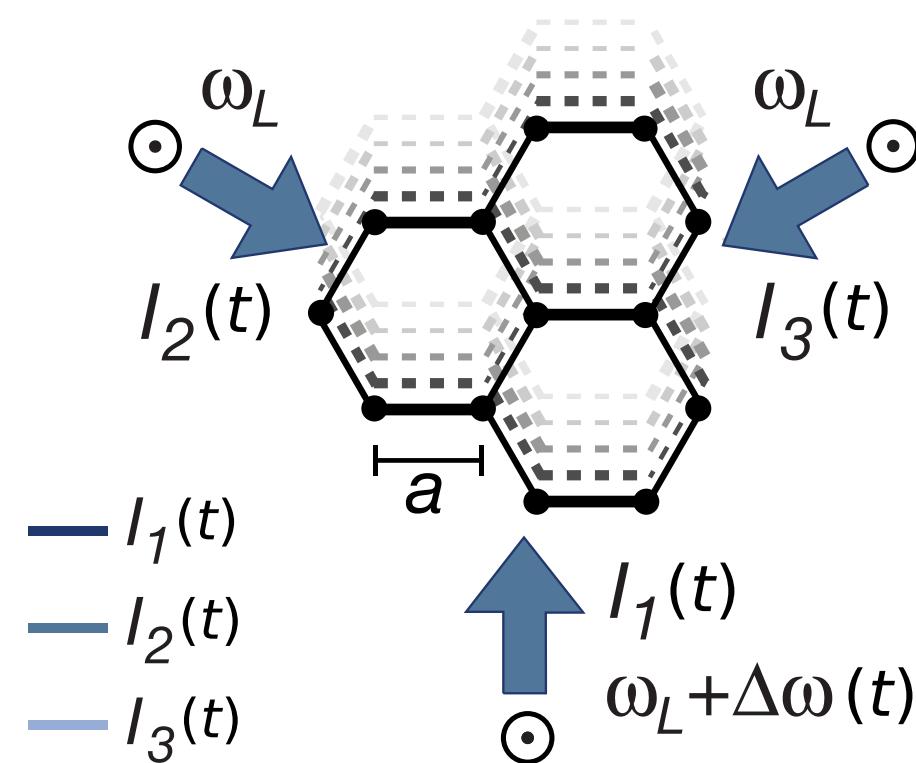
Periodically-modulated hexagonal lattice

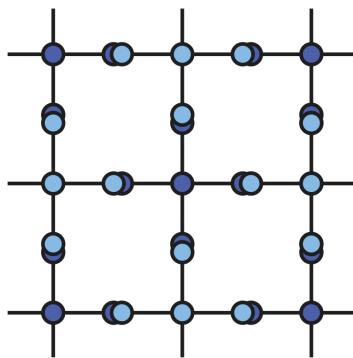
Modulate NN
tunnelings
between λJ and J

T. KITAGAWA ET AL.,
PRB 82, 235114 (2010)

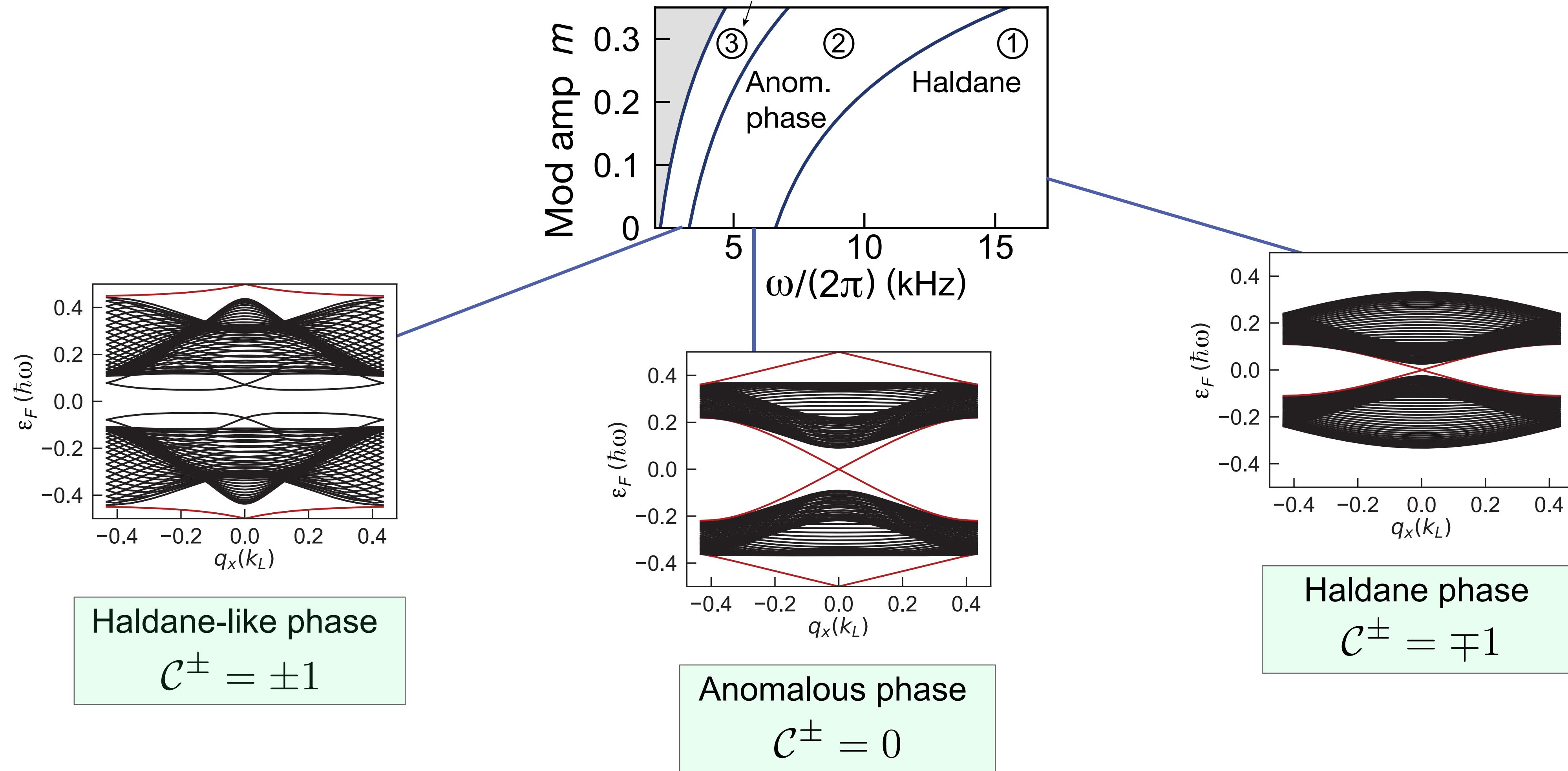


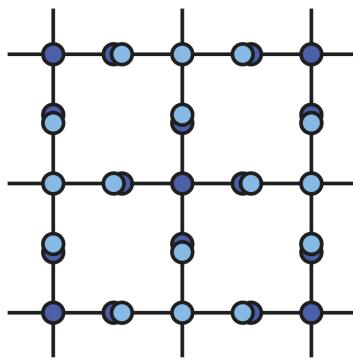
Here: Bosonic ^{39}K
 $\lambda = 737\text{nm}$, $a = 284\text{nm}$



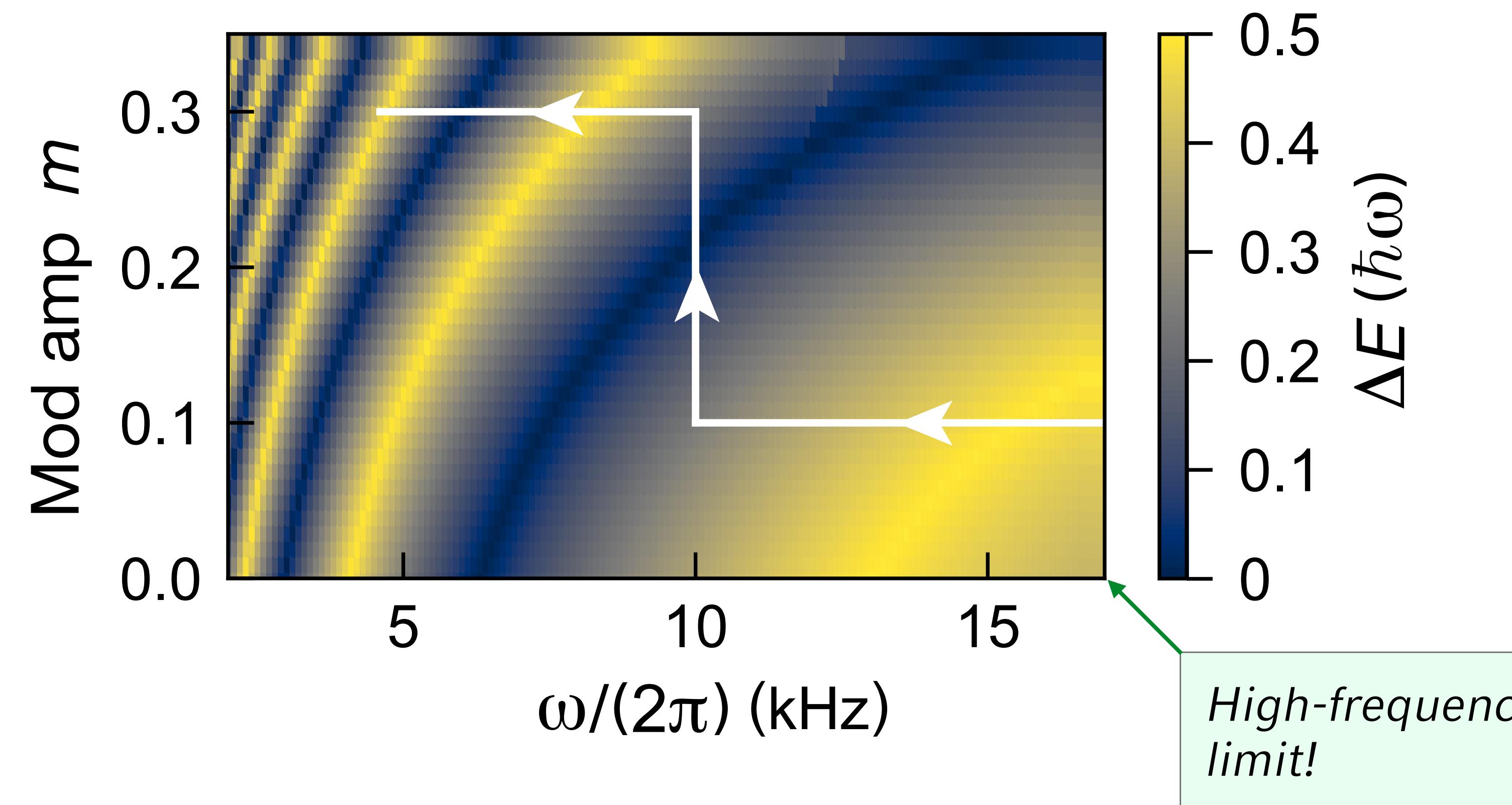


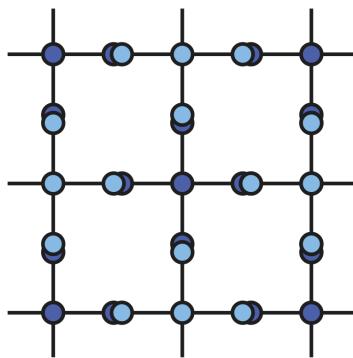
Phase diagram



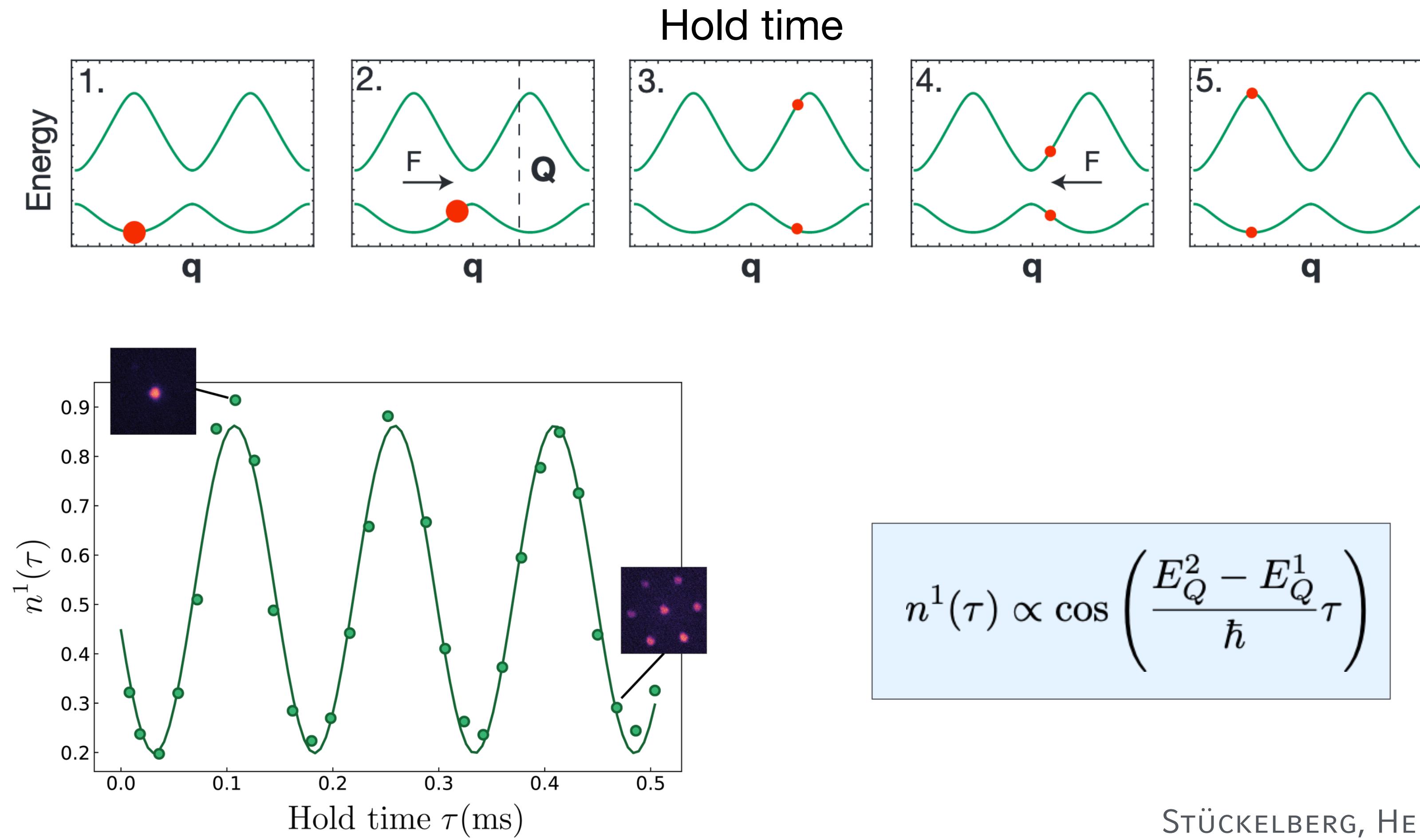


Probing the phase diagram

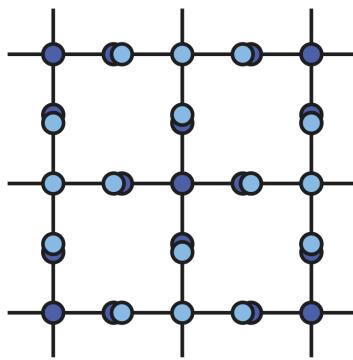




Stückelberg oscillations

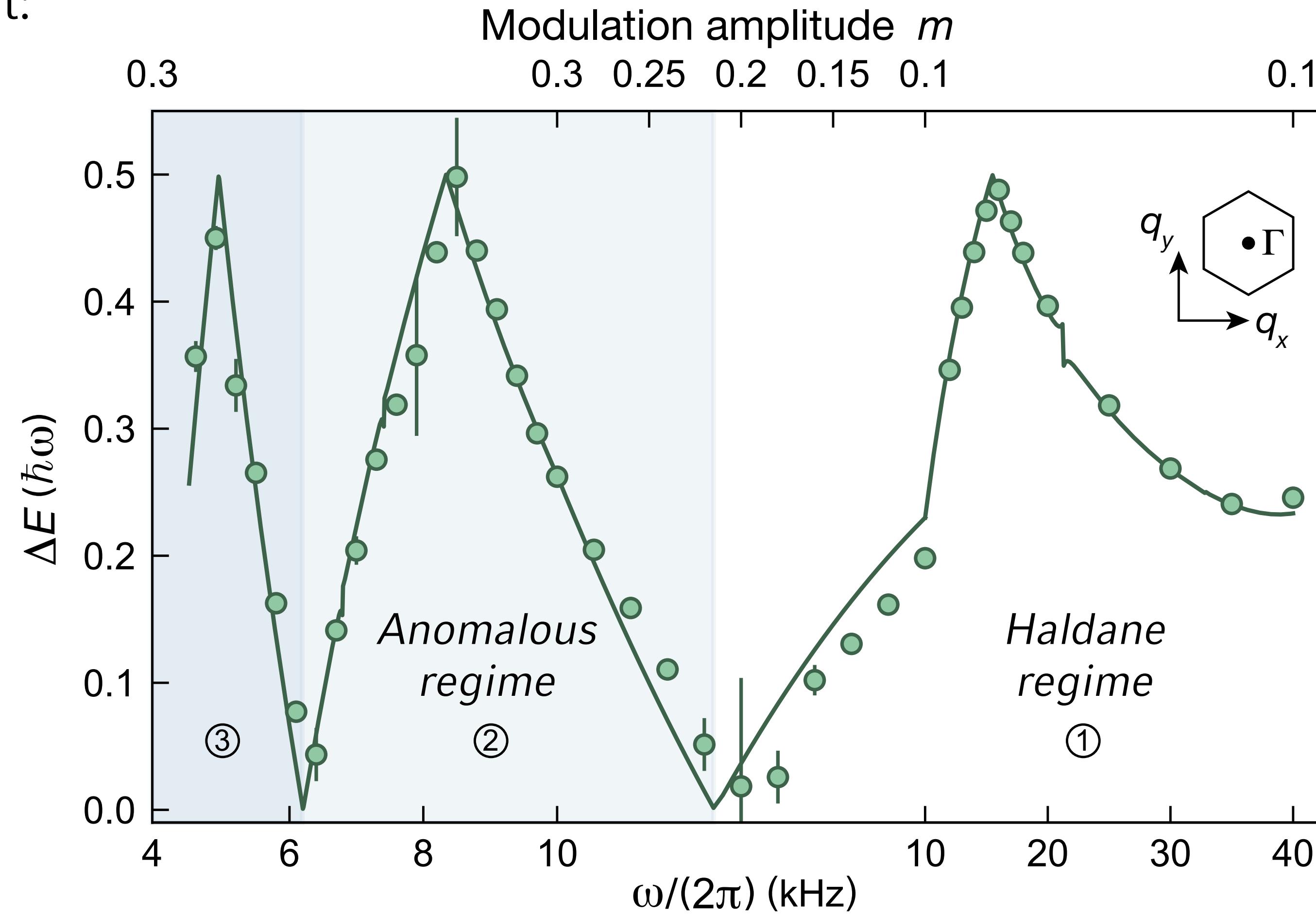


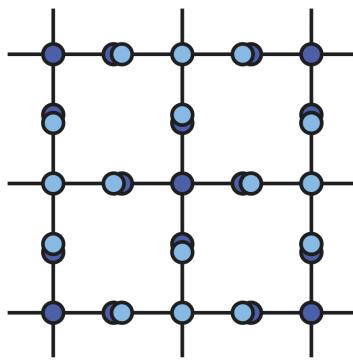
STÜCKELBERG, HELV. PHYS. ACTA 5, 369 (1932)
SHEVCHENKO ET AL., PHYS. REP. 492, 1 (2010)
ZANESINI ET AL., PRA 82, 065601, (2010), WEITZ PRL 105, 215301 (2010)



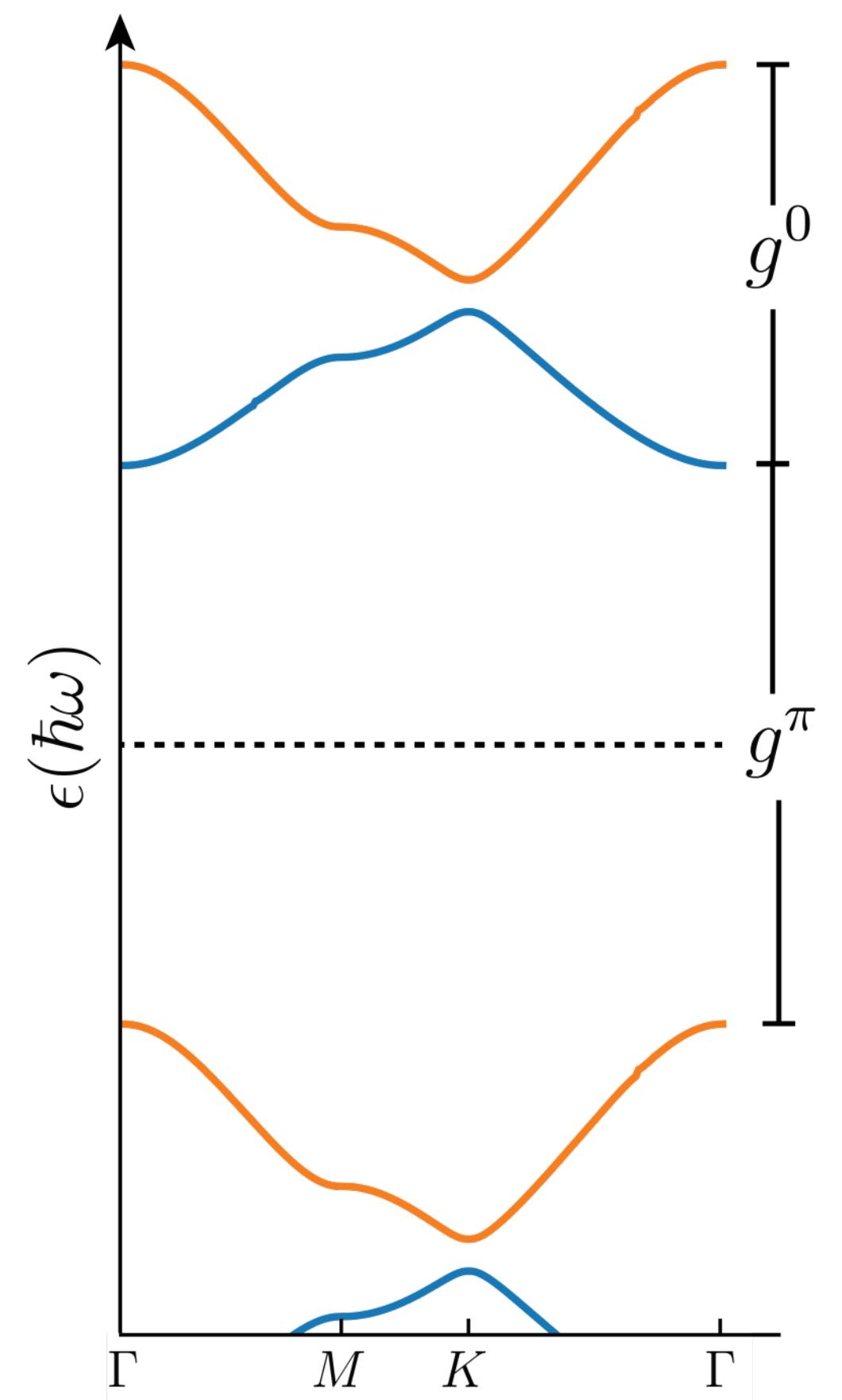
Identify phase transition via gap closing

Γ -point:

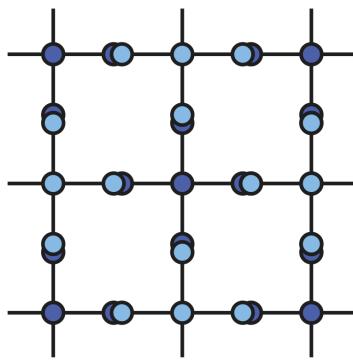




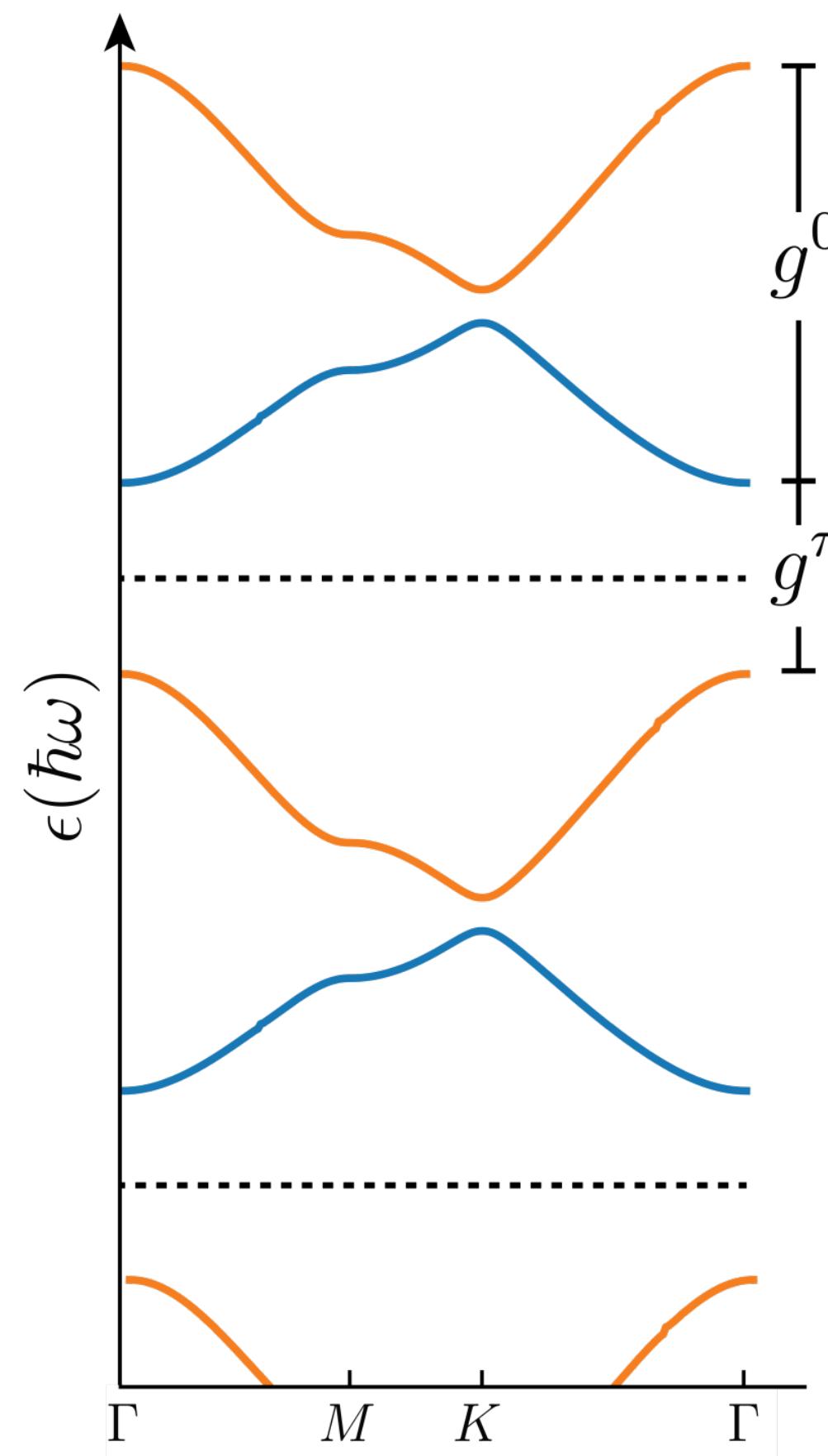
Identify which gap closes



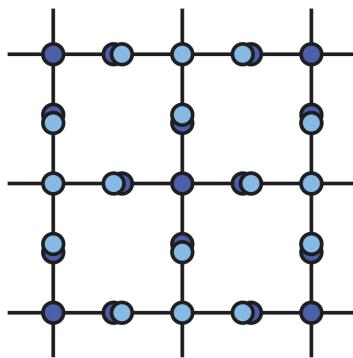
- We always measure the *smaller energy gap*: $\min(g^0, g^\pi)$
- In the *high-frequency limit* we probe g^0



Identify which gap closes

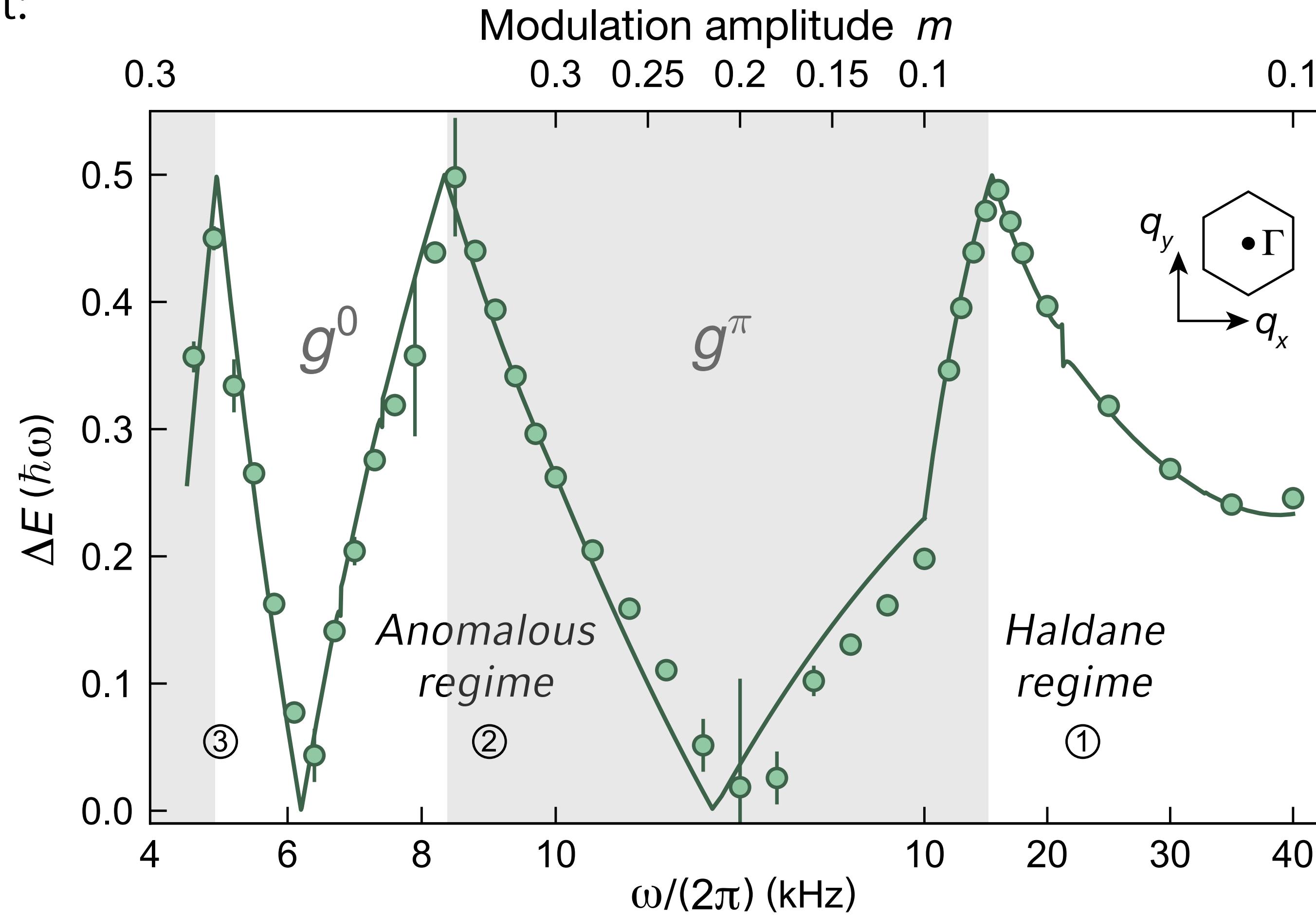


- We always measure the *smaller energy gap*: $\min(g^0, g^\pi)$
- In the *high-frequency limit* we probe g^0
- Transition, when
$$g^0 = g^\pi = \hbar\omega/2$$

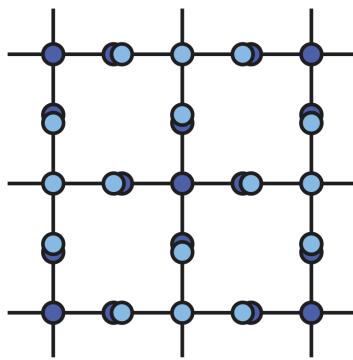


Identify which gap closes

Γ -point:

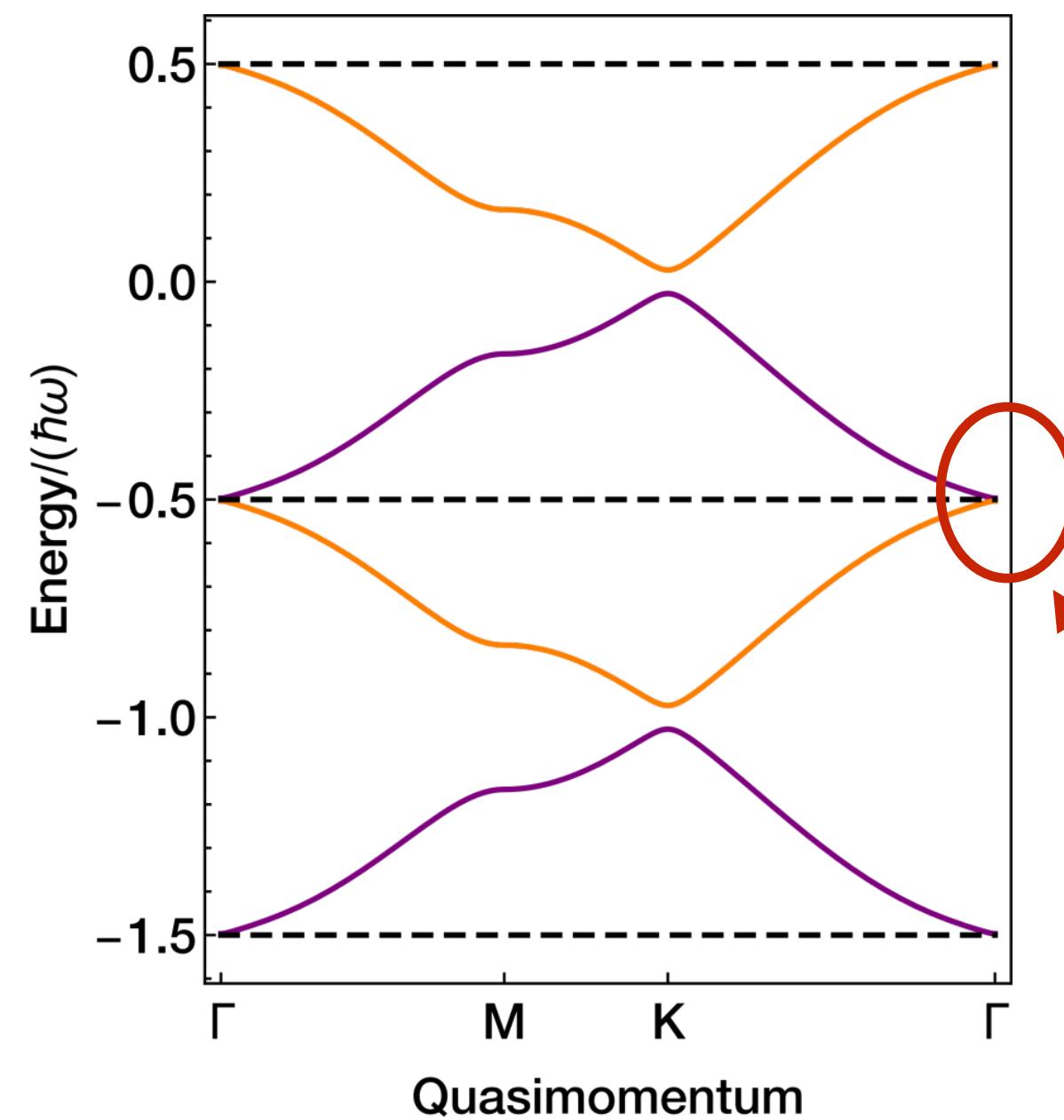


Determine bulk geometric properties

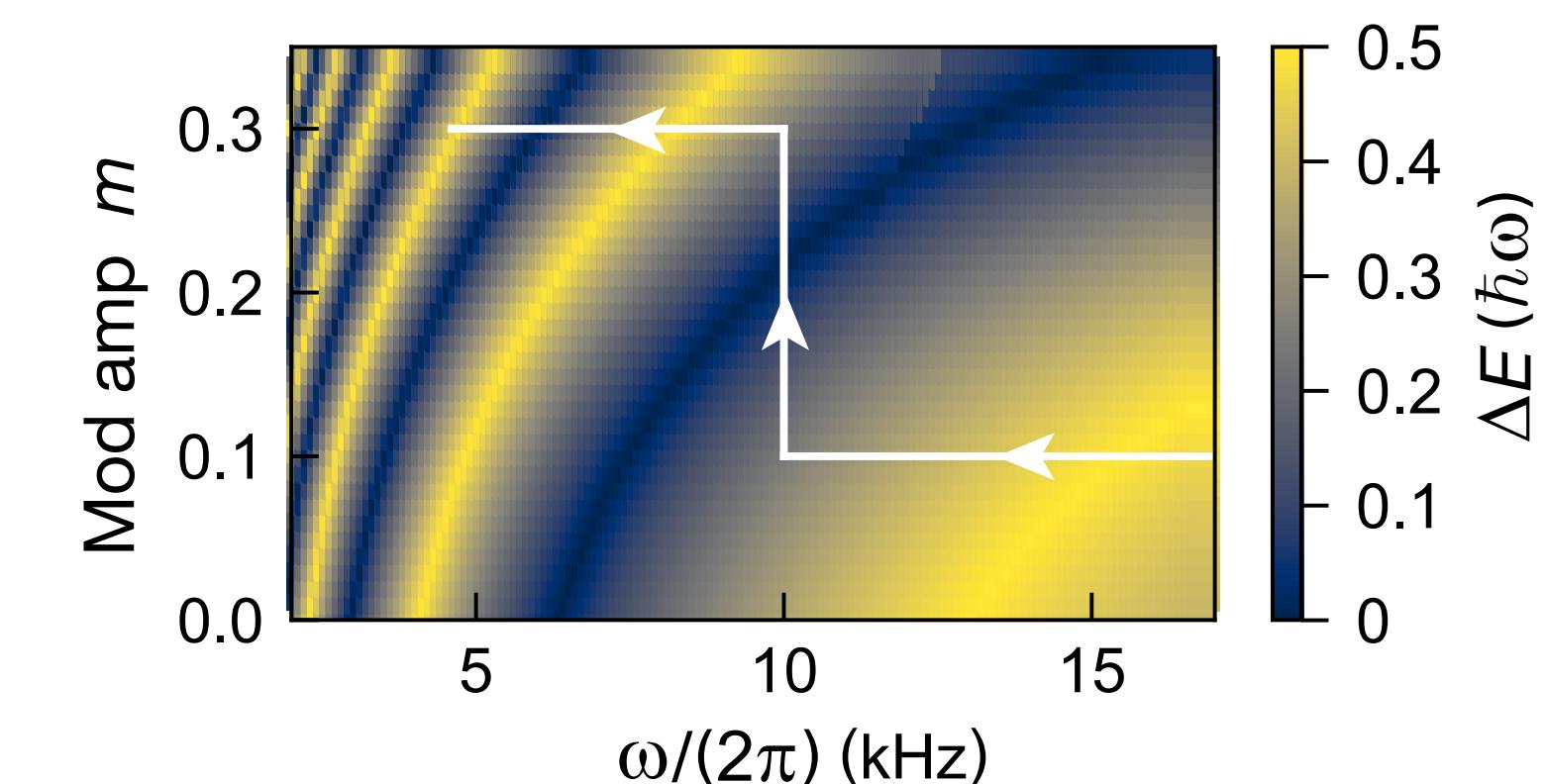


Topology of gap-closing point

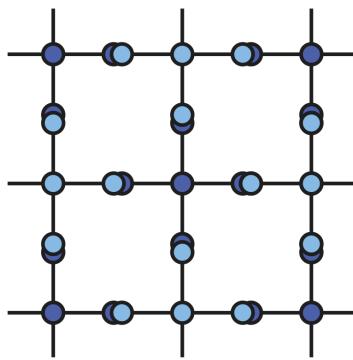
Band-touching singularities



- Consider 3D-parameter space: (q_x, q_y, λ)

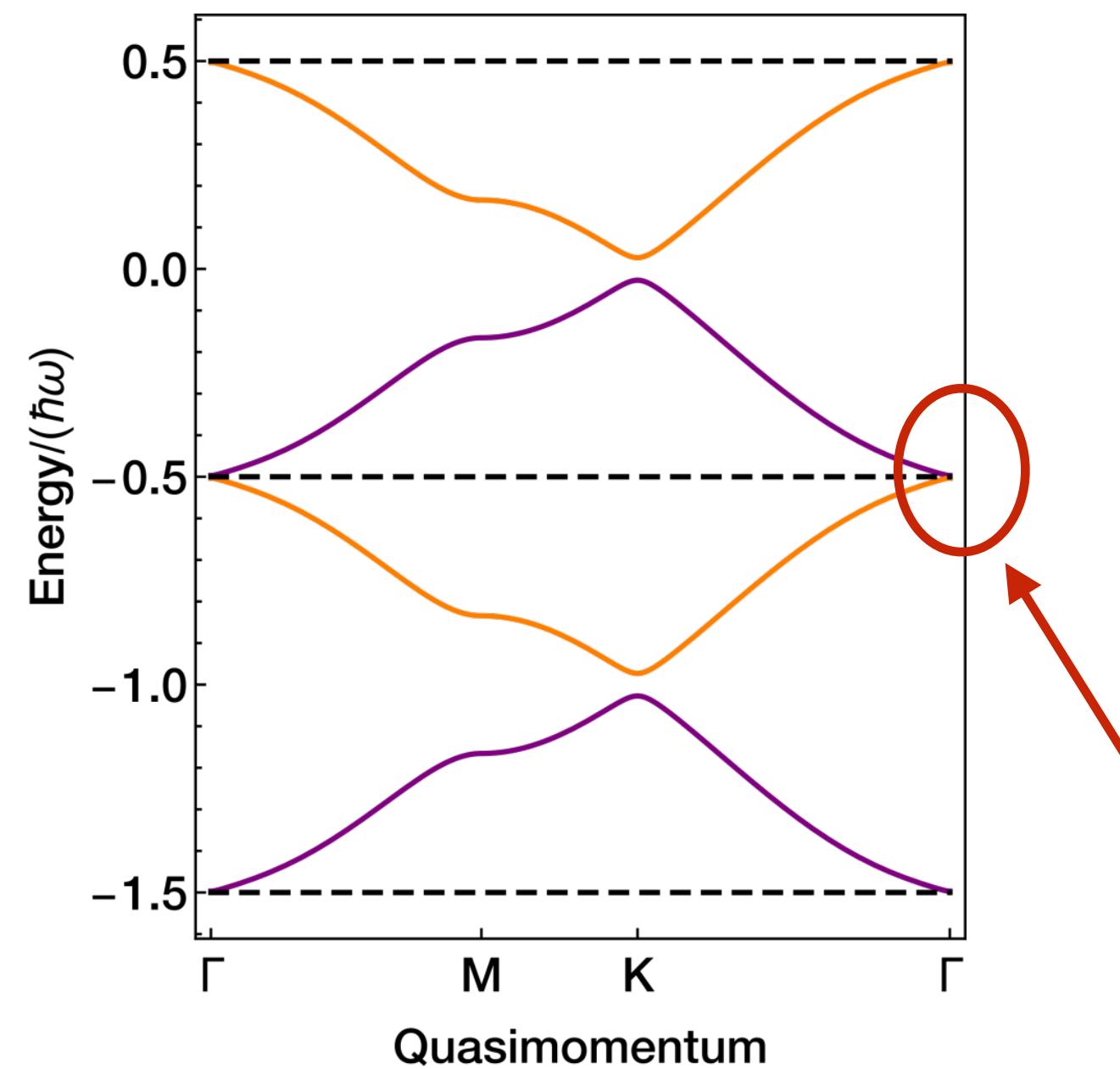


Topological charge $Q_s = \pm 1$



Topology of gap-closing point

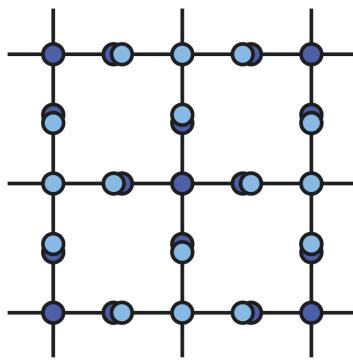
Band-touching singularities



- Consider 3D-parameter space: (q_x, q_y, λ)
- Topological charge describes change in Winding number

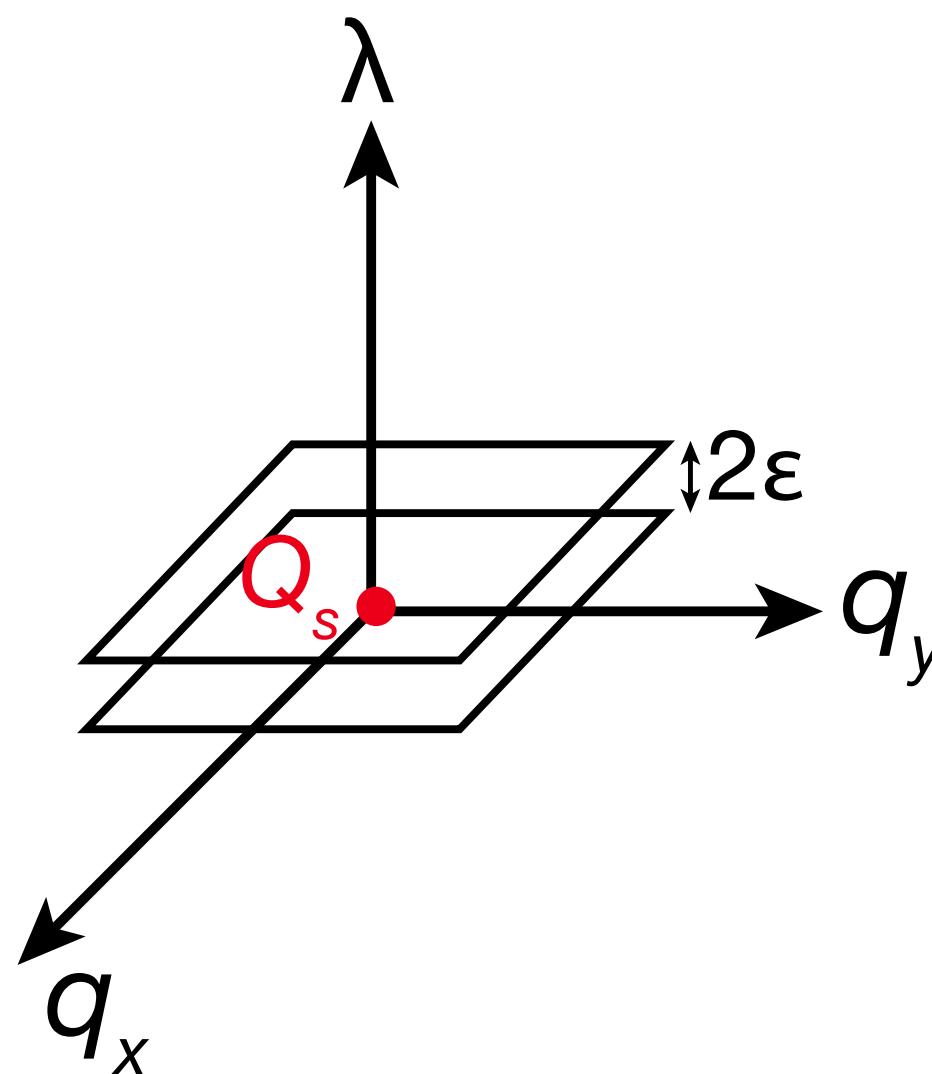
$$W_{\lambda_s + \varepsilon}^j = W_{\lambda_s - \varepsilon}^j + Q_s^j$$

Topological charge $Q_s = \pm 1$

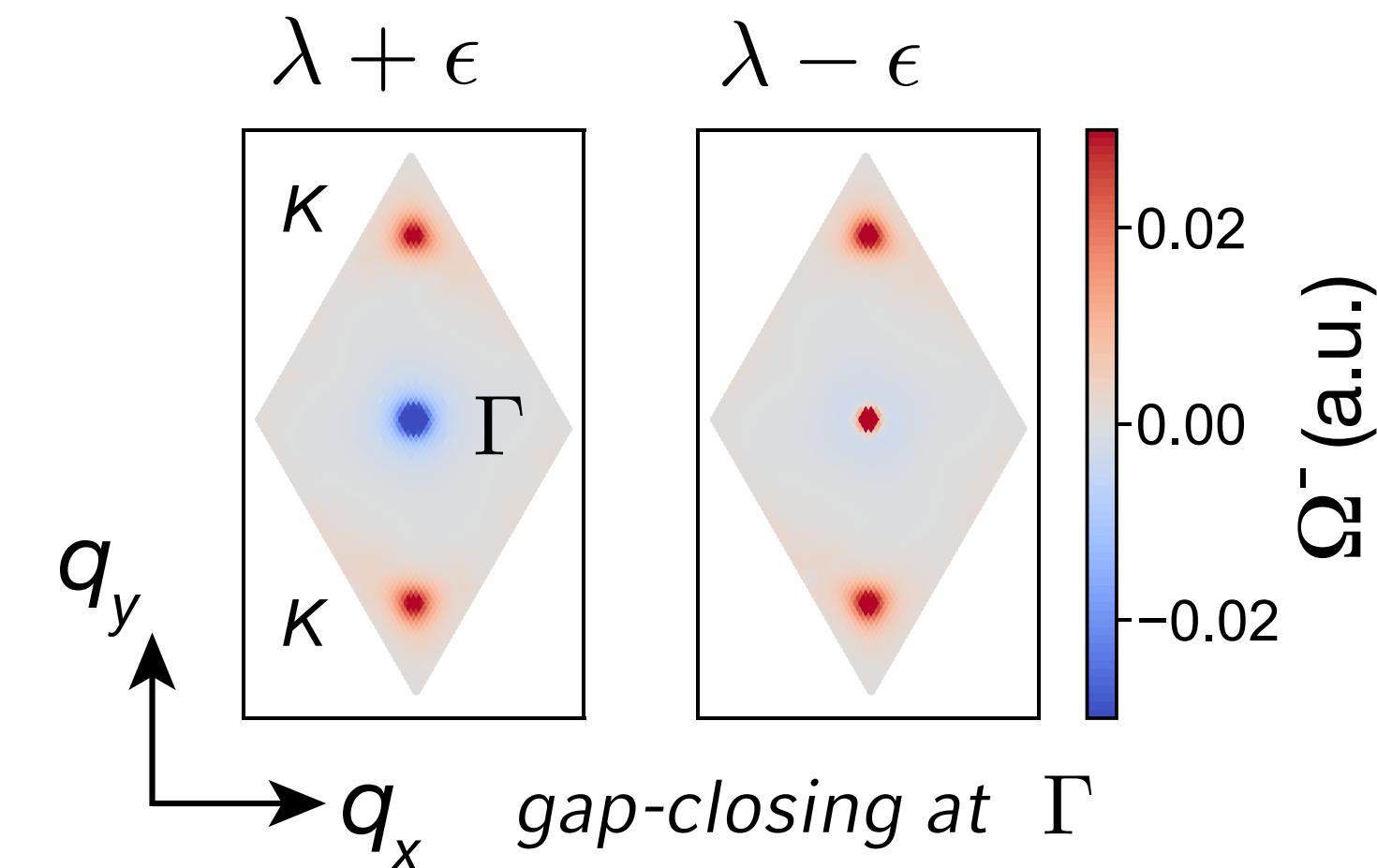


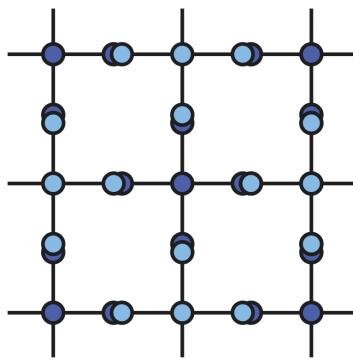
Determine topological charge

Measure Berry flux associated with Q_s



Measure *Berry-flux locally shortly before and after* the gap-closing point!

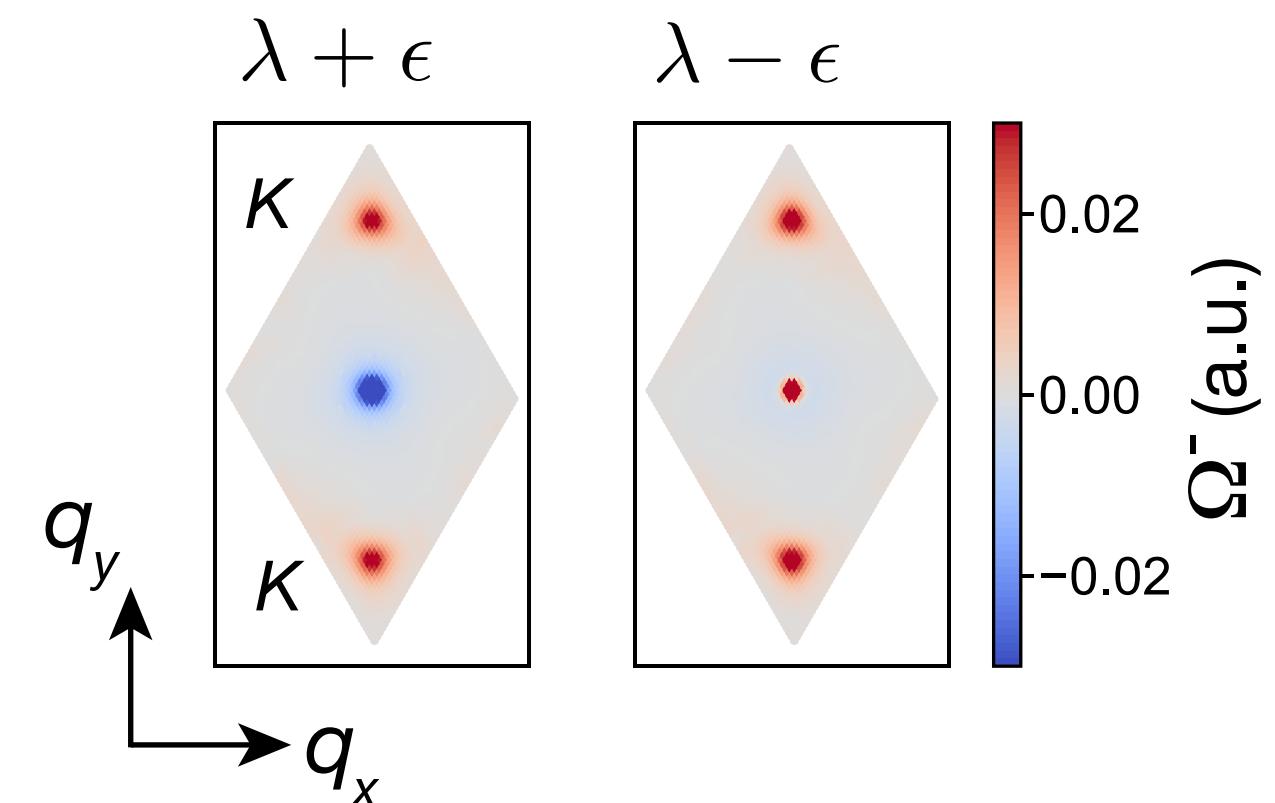
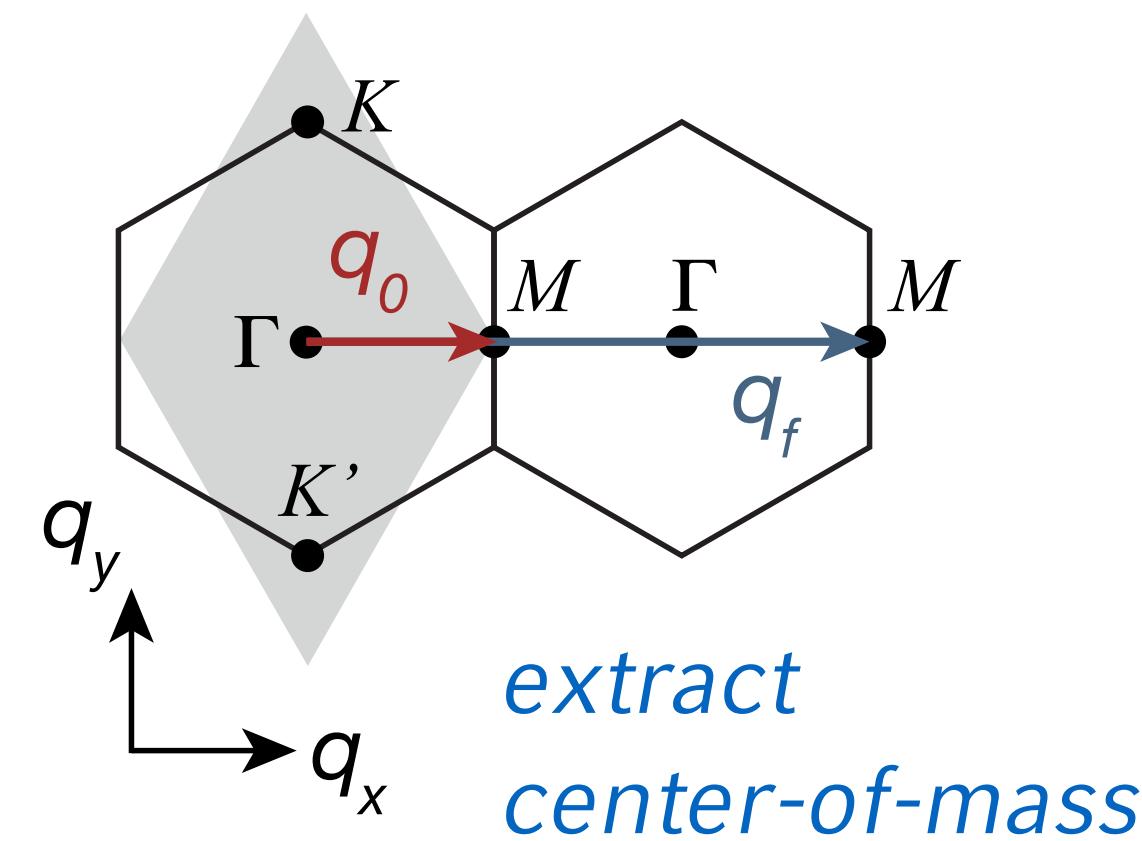




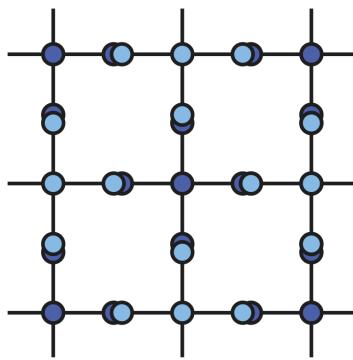
Local Hall deflection

Anomalous transverse velocity
proportional to *Berry curv.* Ω_μ

$$\mathbf{v}_\mu^x(\mathbf{q}) = -\frac{F}{\hbar} \Omega_\mu(\mathbf{q})$$

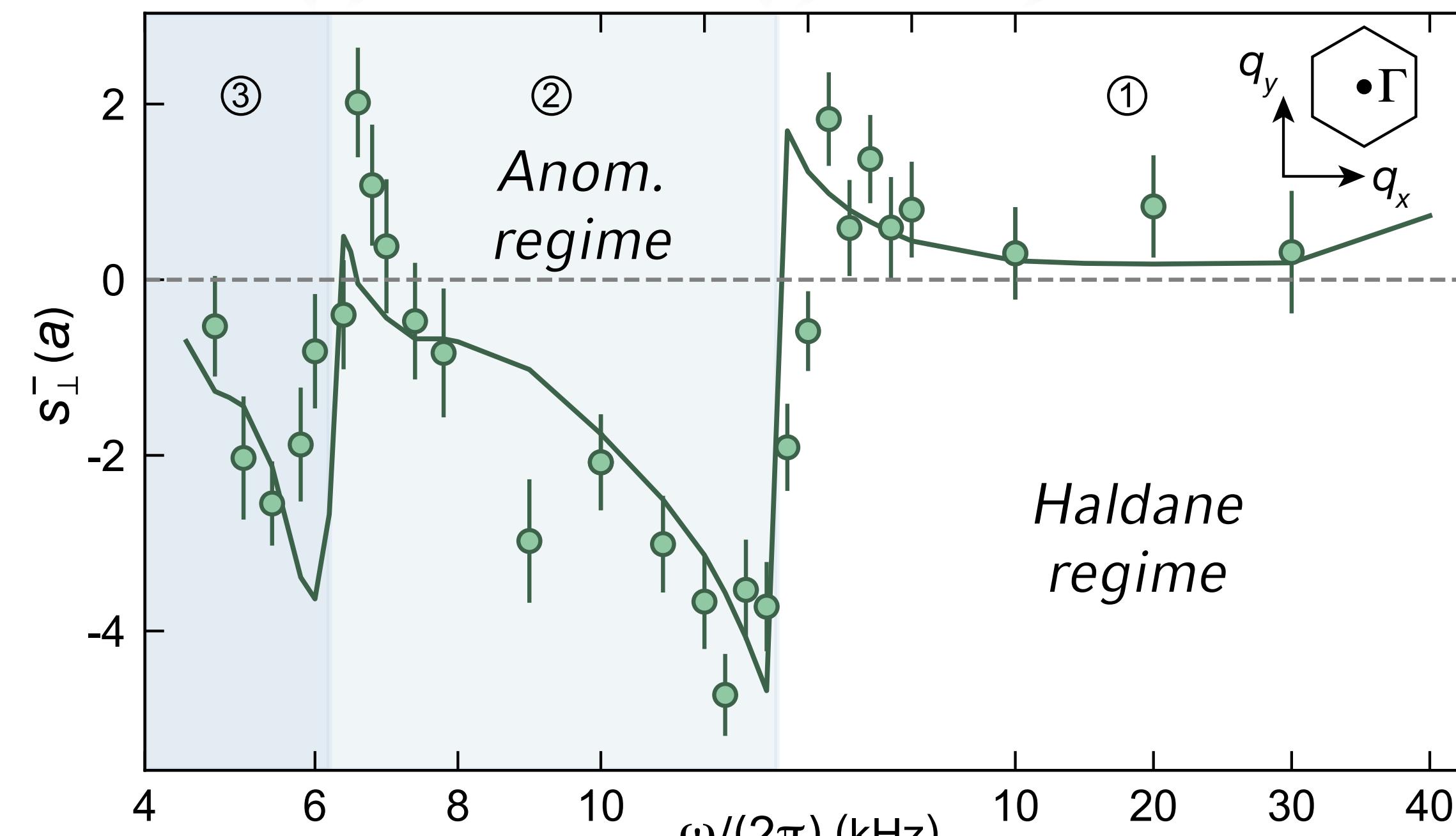
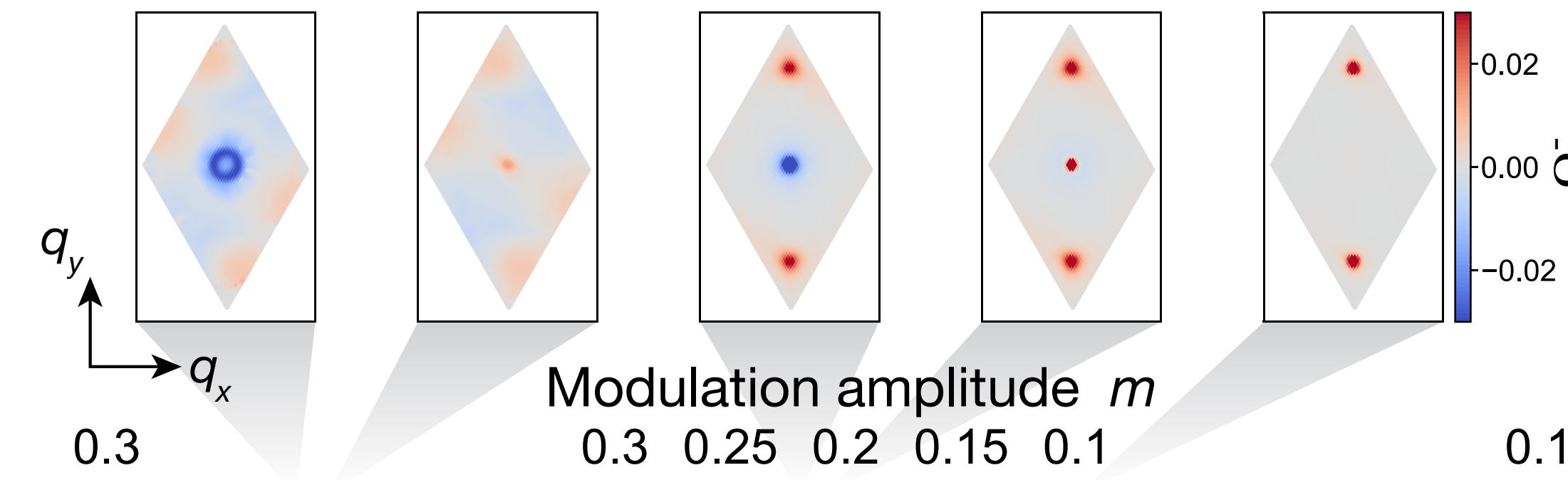


$$Q_s^0 = \text{sgn} (\Delta s_\perp^-(\mathbf{q}_s))$$

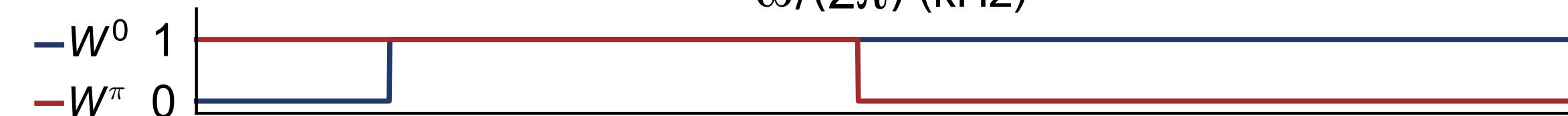


Berry curvature from local Hall deflection

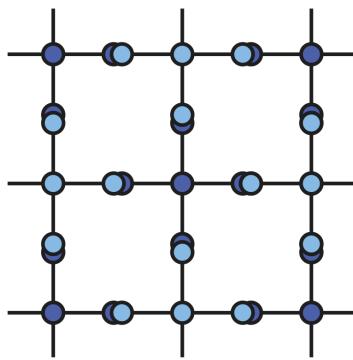
Γ -point:



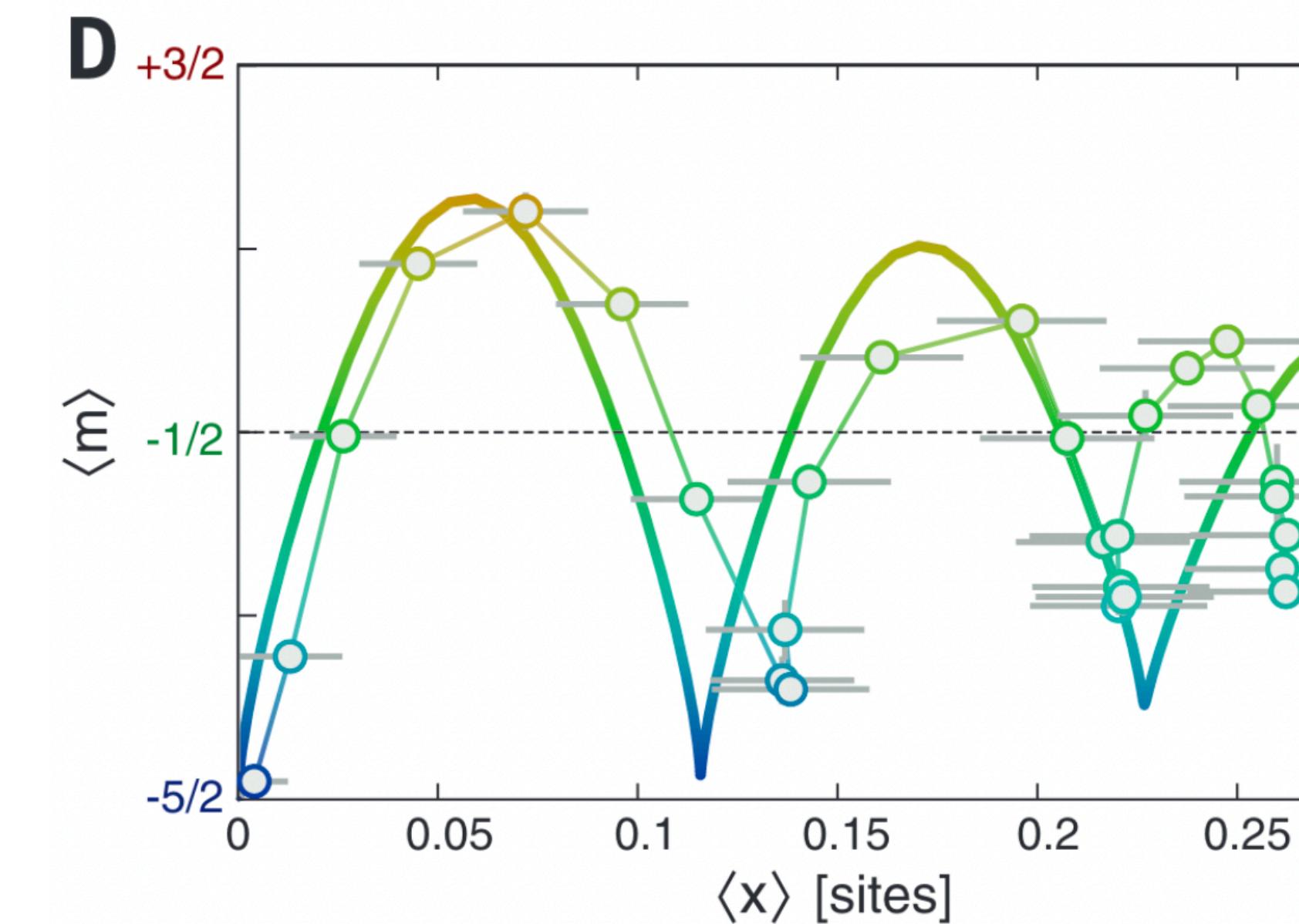
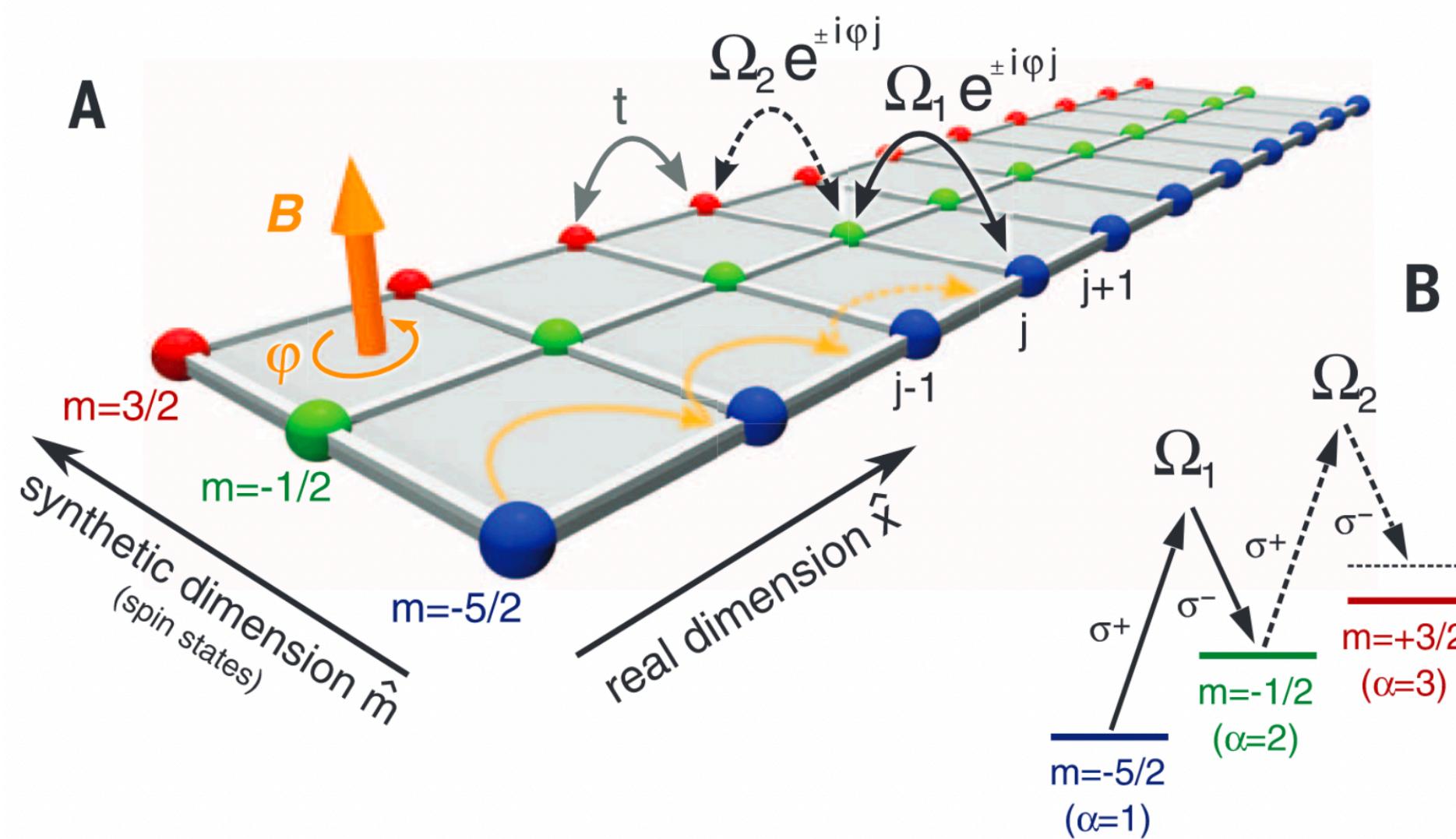
Complete set of topological invariants!



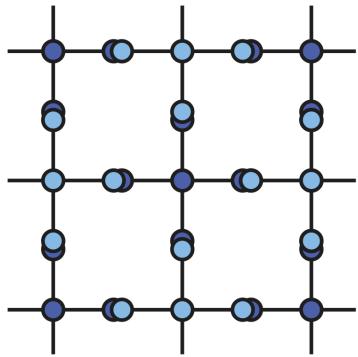
Can we observe topological
edge modes?



Synthetic dimensions

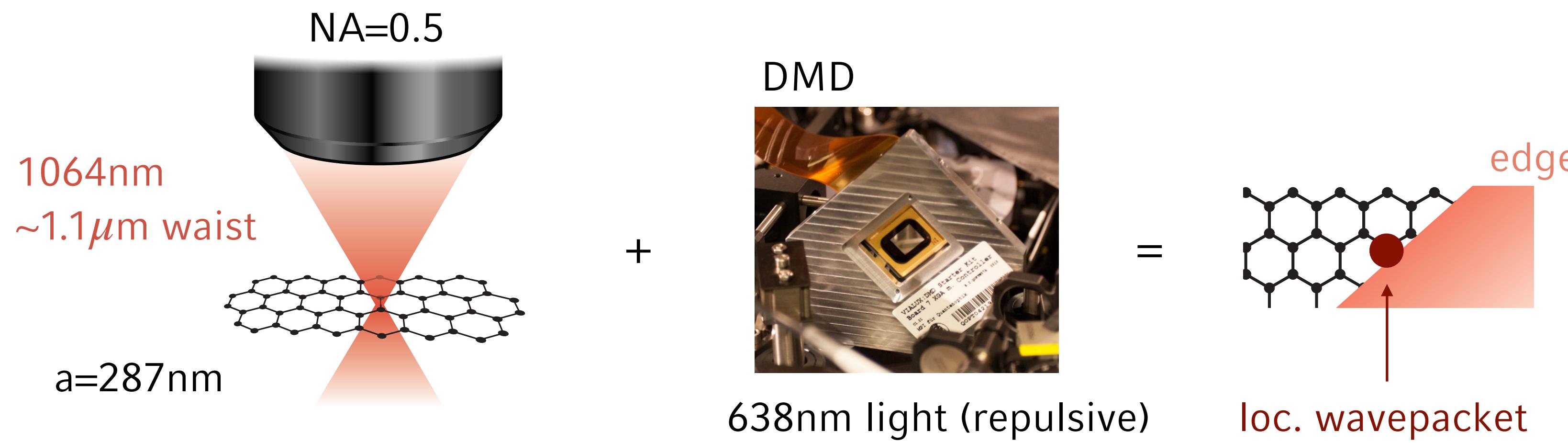


MANCINI, ..., FALLANI, SCIENCE 349 (2015)
 STUHL, ..., SPIELMAN, SCIENCE 349 (2015)



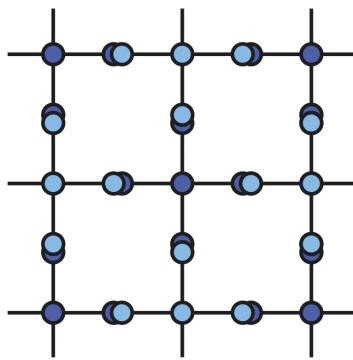
Edge dynamics

Realizing a sharp edge:



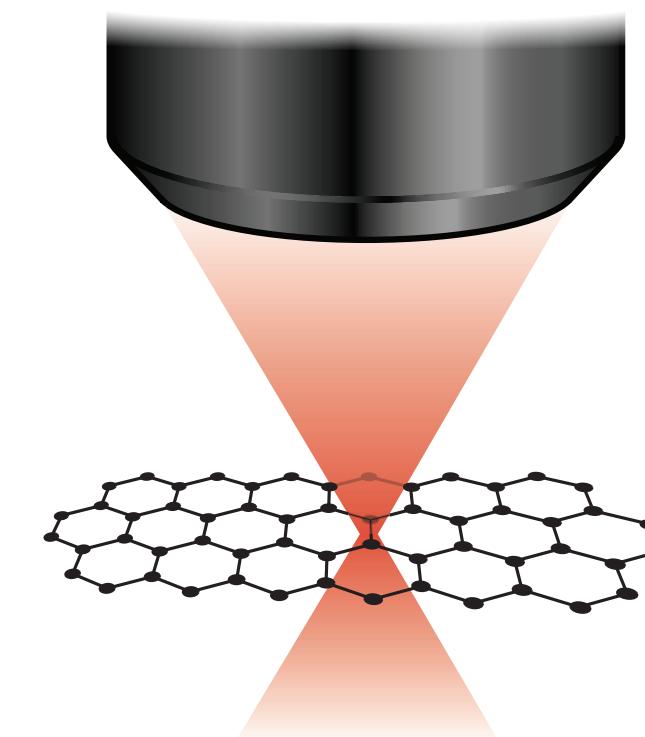
Width of the edge: 2-3 lattice sites!

Preliminary!

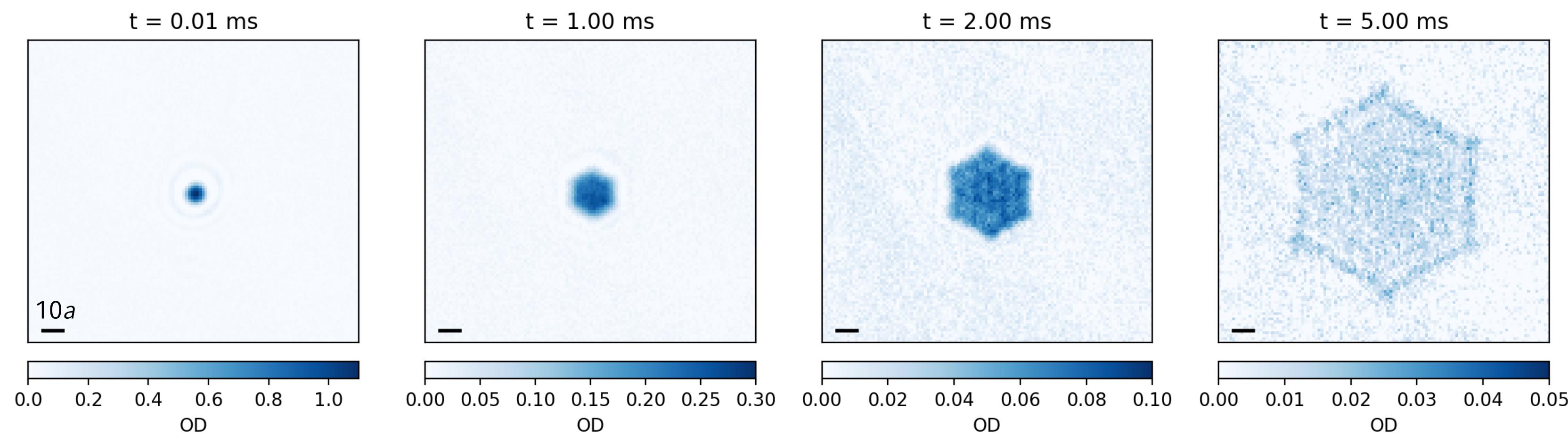


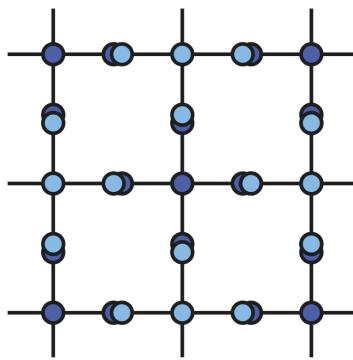
Bulk expansion

Lattice depth: 6Er; $\lambda=745\text{nm}$;
 $\omega_r = 2\pi \times 17\text{Hz}$; $\omega_z = 2\pi \times 250\text{Hz}$

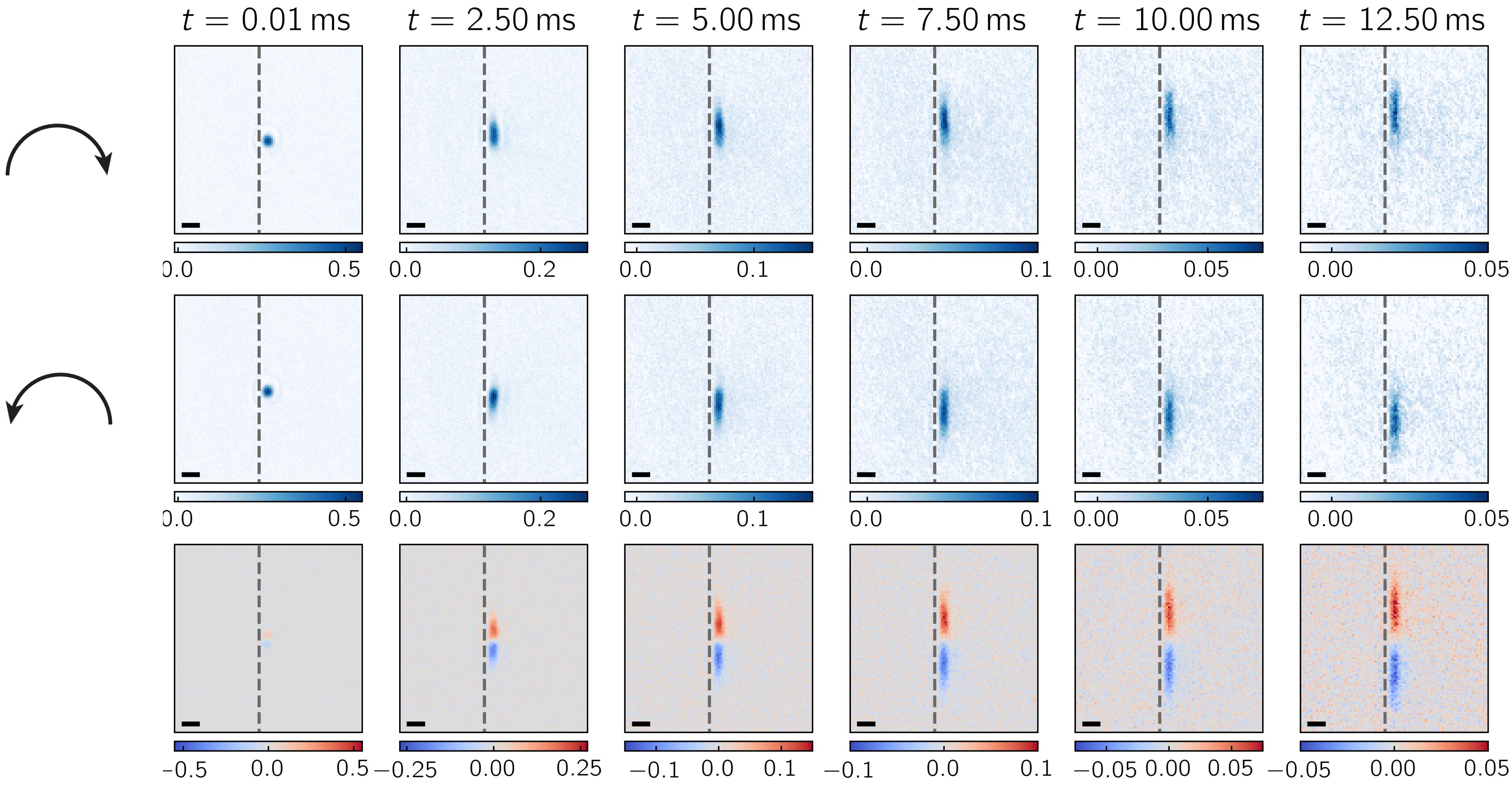


$$\mathbf{v}_\mu^{\text{band}} = \frac{1}{\hbar} \partial_{\mathbf{q}} E_\mu$$



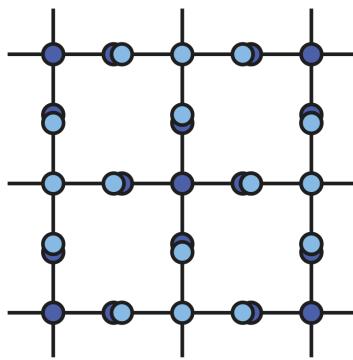


Edge dynamics in anomalous regime

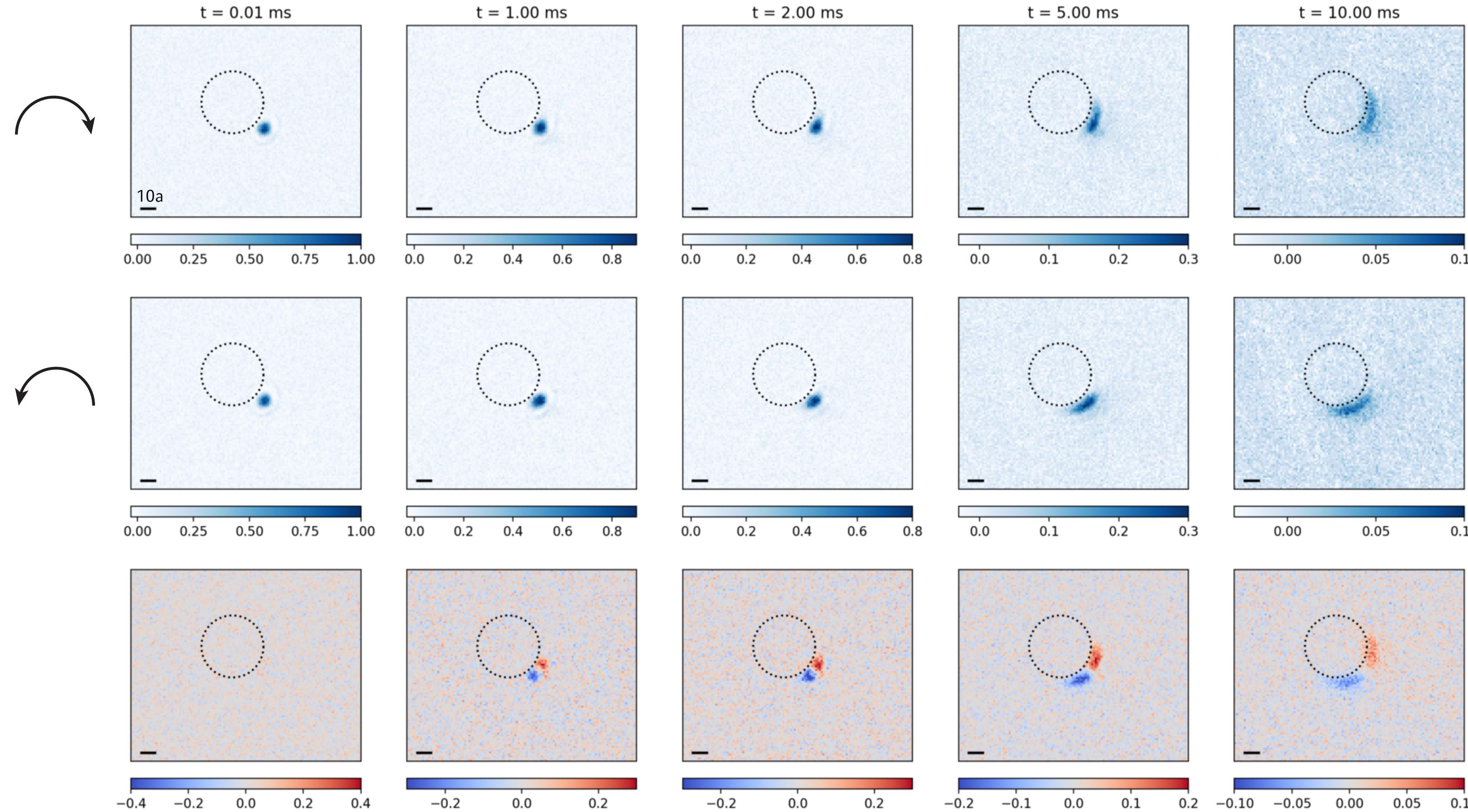


f=7kHz, m=0.25

Preliminary!

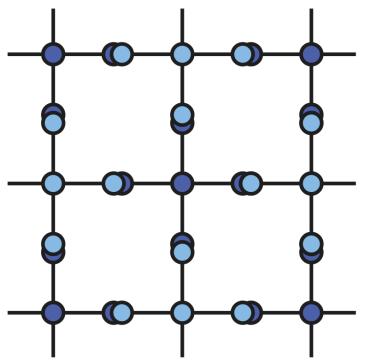


Edge dynamics in anomalous regime

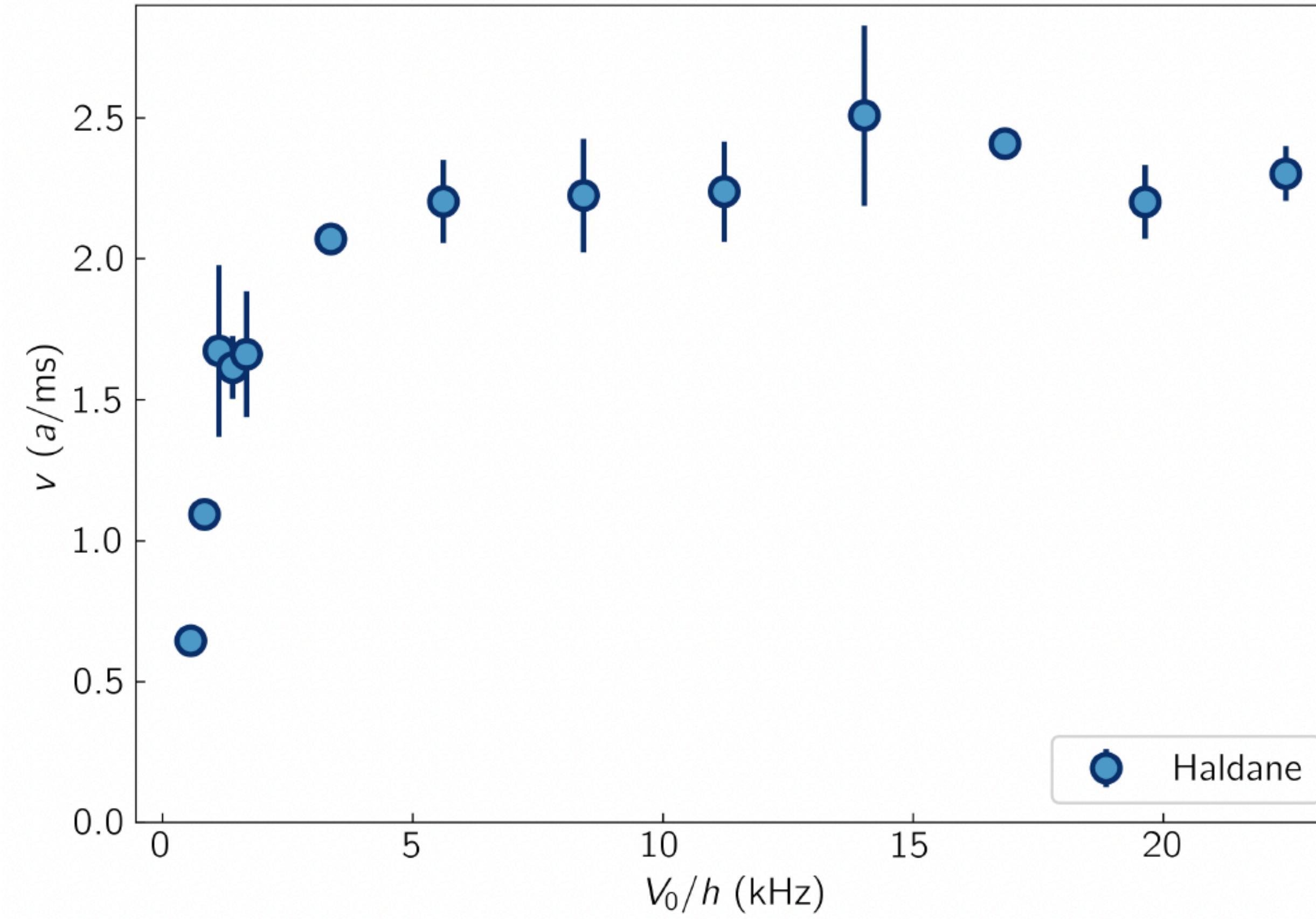


Preliminary!

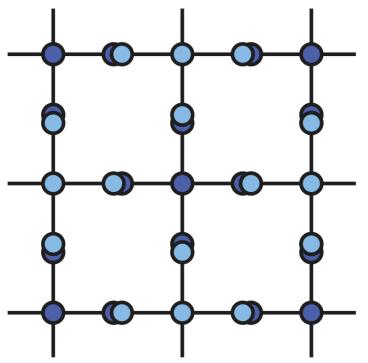
$f=7\text{kHz}, m=0.25$



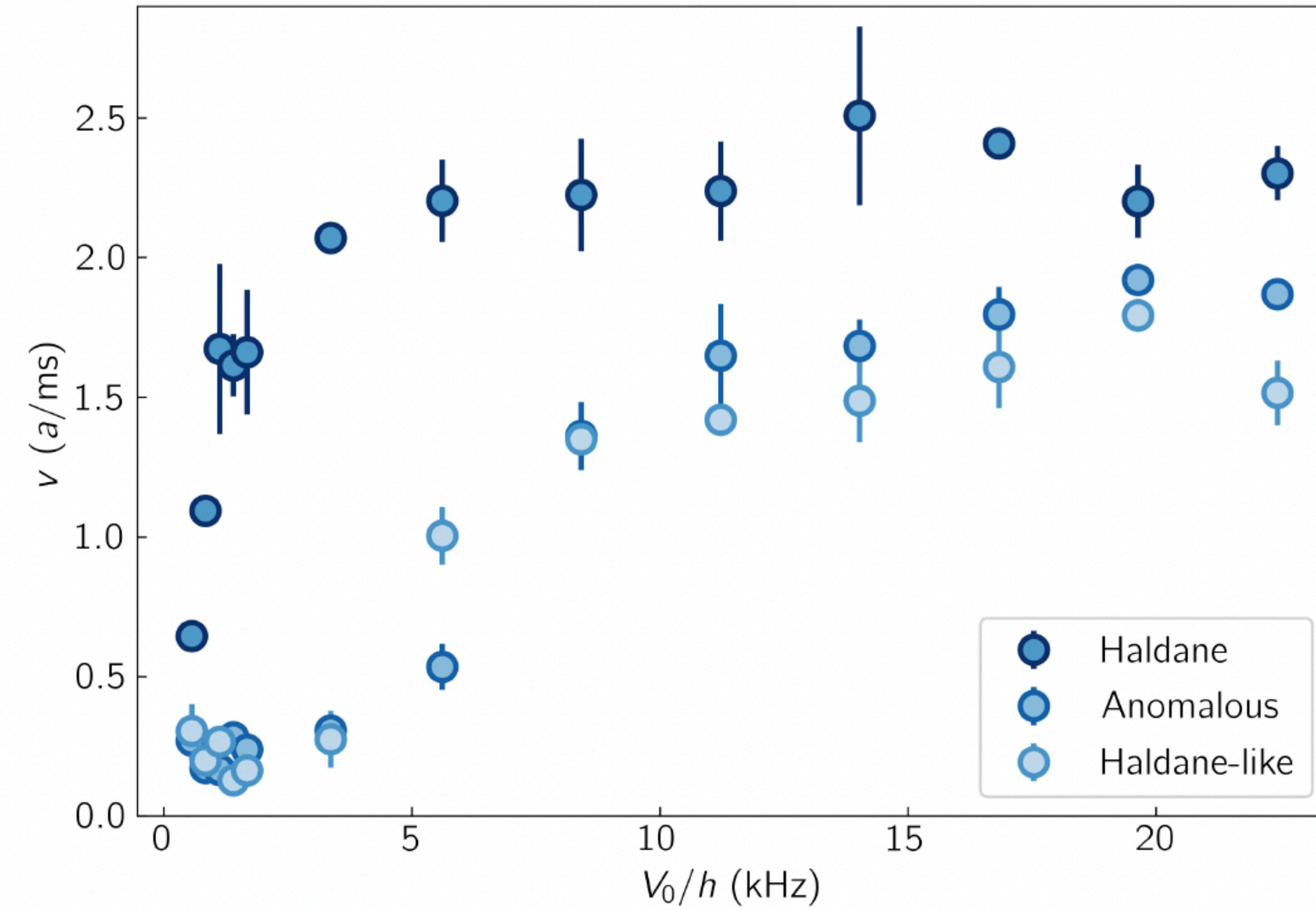
Generating a topological interface



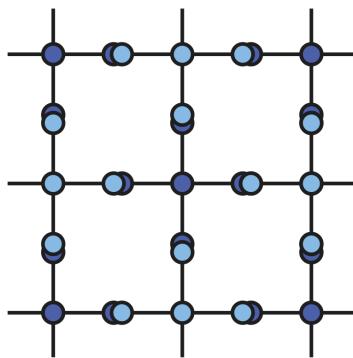
Preliminary!



Generating a topological interface

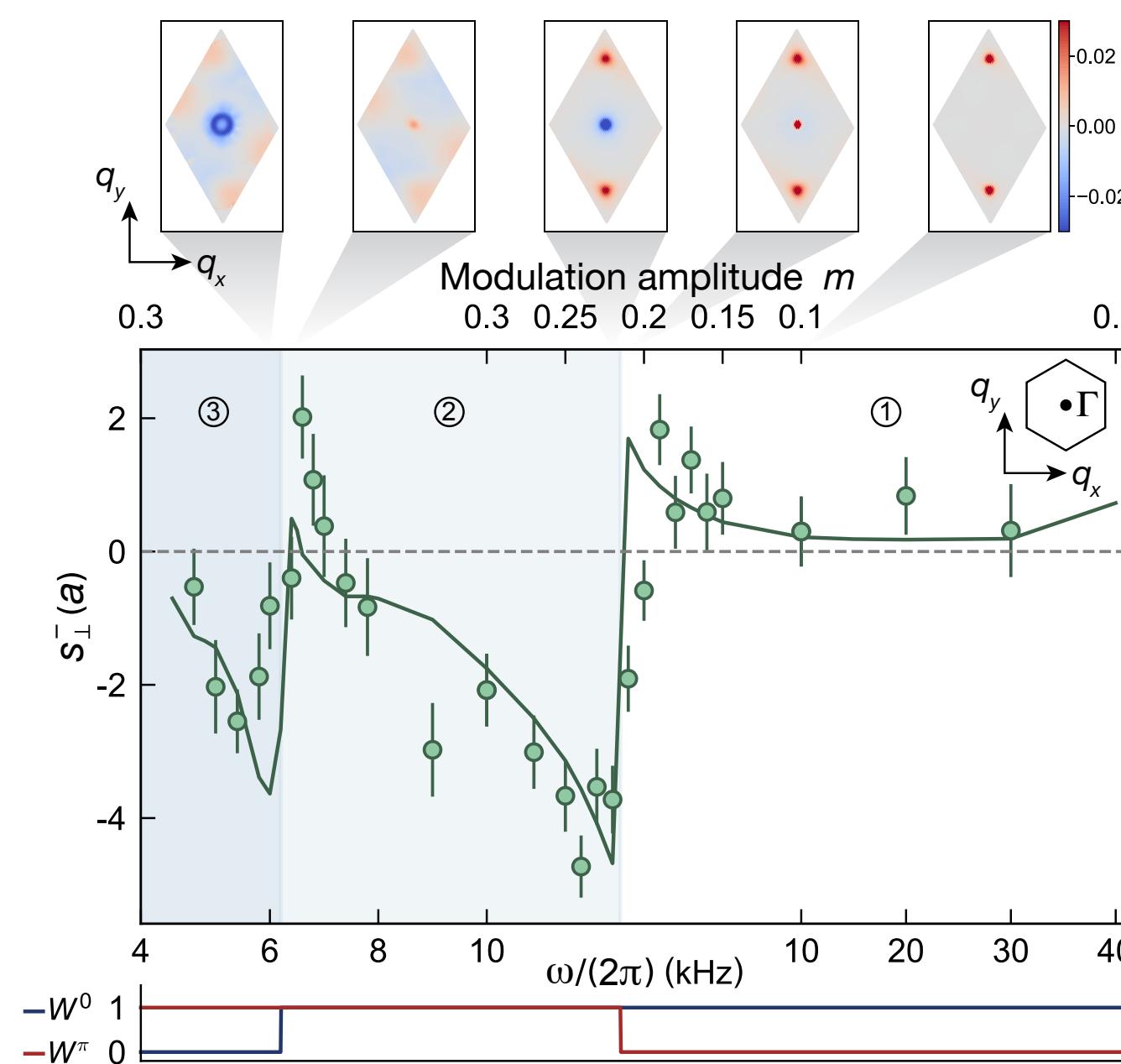


Preliminary!

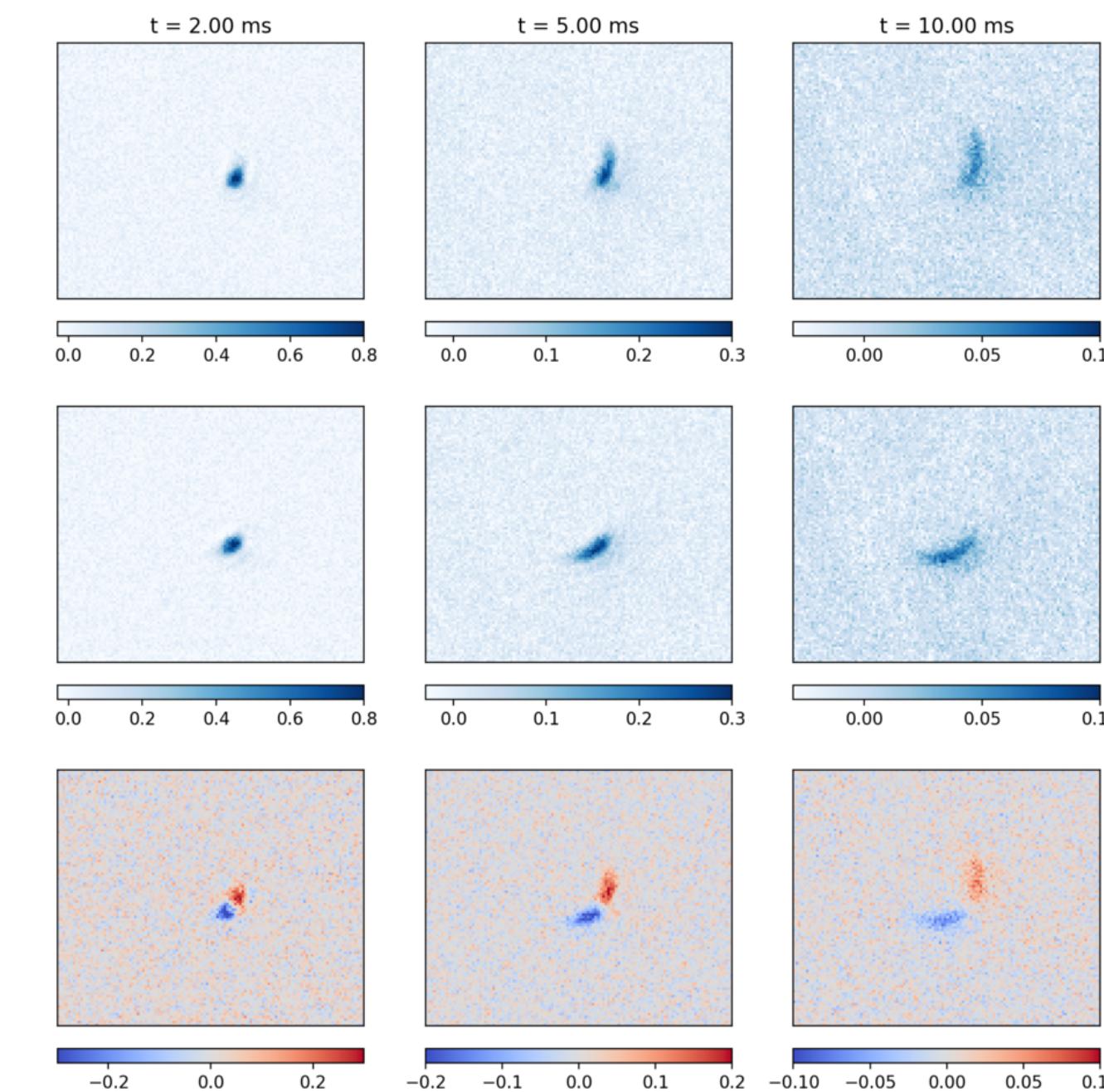


Summary

Bulk



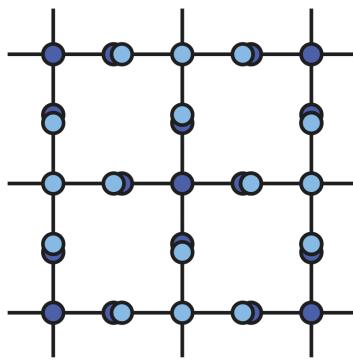
Edge



K. WINTERSPERGER,, MA,
NATURE PHYS. 16, 1058-1063 (2020)

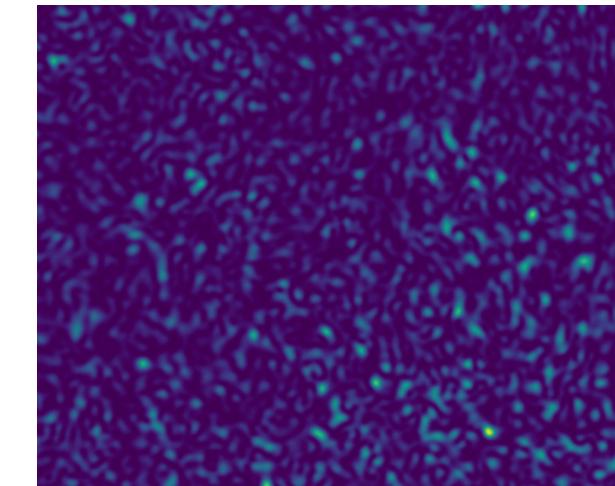
Preliminary!

Outlook



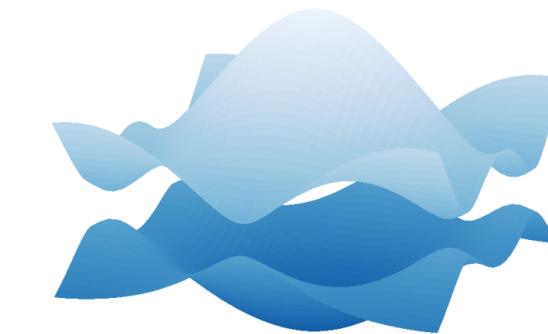
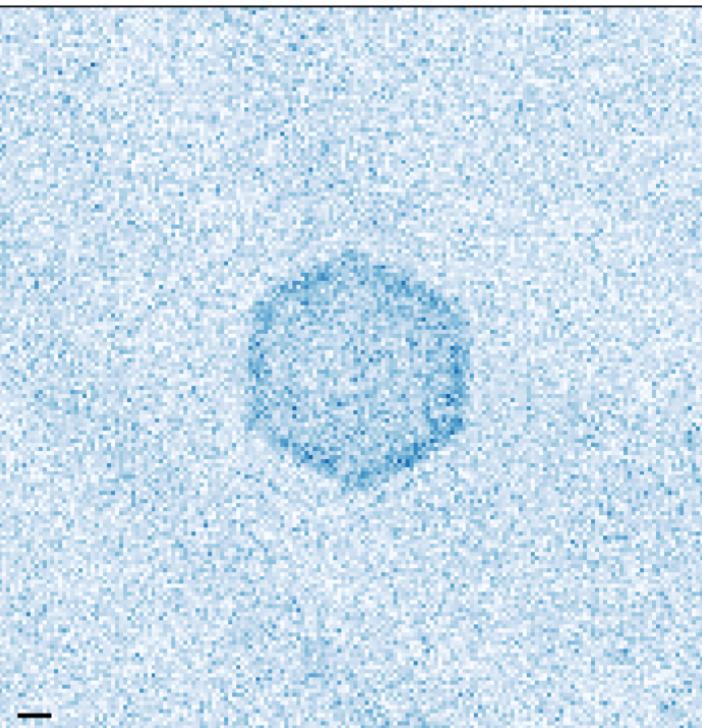
Outlook:

- Interplay of topology & disorder
- Realization of many-body top. phases

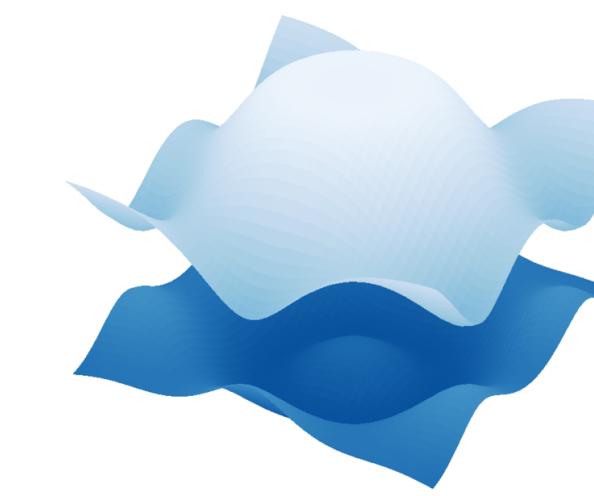
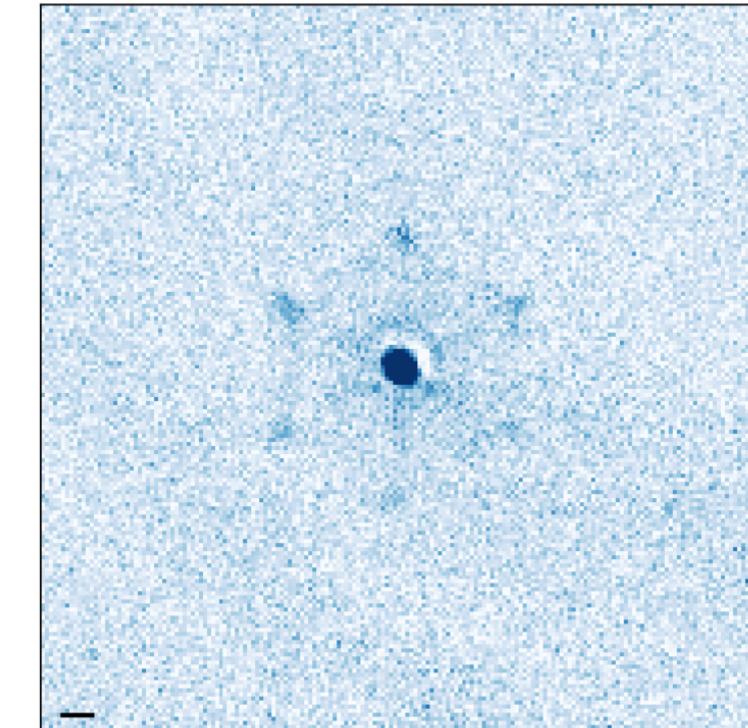


Bulk expansion in the modulated lattice: $v_\mu^{\text{band}} = \frac{1}{\hbar} \partial_{\mathbf{q}} E_\mu$

Haldane:

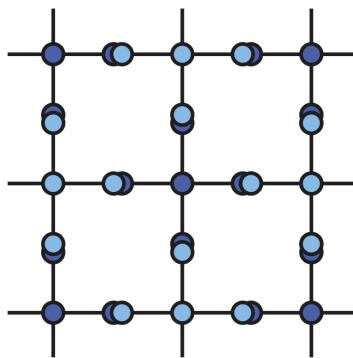


Anomalous:

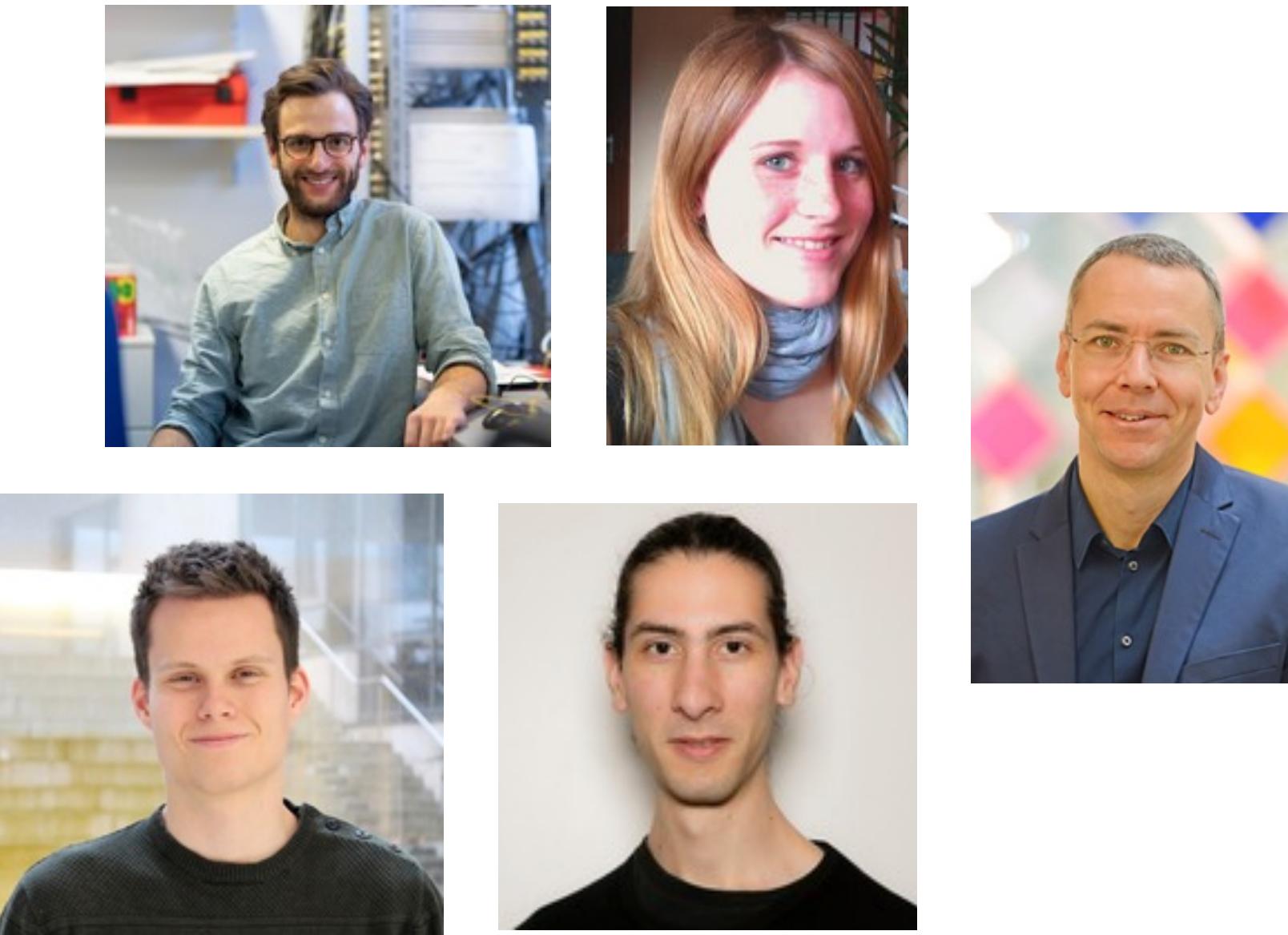


C. A. BRACAMONTES ET AL.
PHYS. REV. LETT. 128 (2022)

Preliminary!



The team



Experiment

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Theory anomalous Floquet top. systems

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- André Eckardt (TU Berlin)
- Nathan Goldman (ULB Brussels)
- Marco Di Liberto (U of Padova)