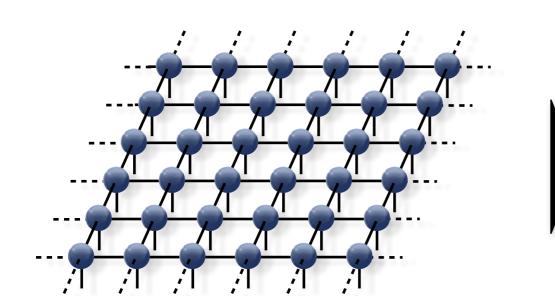
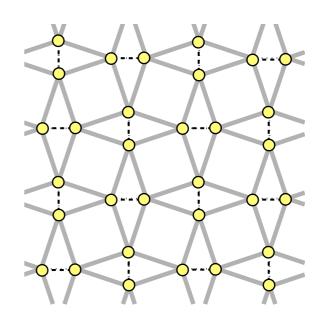
# Progress in understanding the quantum phases of SrCu<sub>2</sub>(BO<sub>3</sub>)<sub>2</sub>

Philippe Corboz, Institute for Theoretical Physics, University of Amsterdam





**iPEPS**2D tensor network ansatz

Shastry-Sutherland model







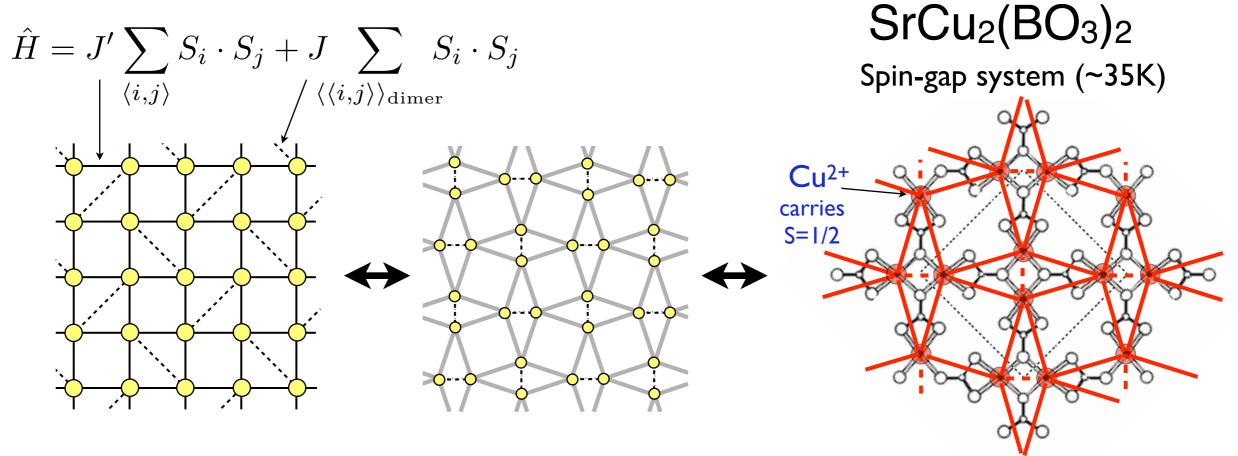




#### Outline

- ▶ Part I: SCBO under pressure at finite temperature
  - ◆ Critical point at finite temperature, analogous to the critical point of water
    J. L. Jiménez, S. P. G. Crone, E. Fogh, M. E. Zayed, R. Lortz, E. Pomjakushina,
    K. Conder, A. M. Läuchli, L. Weber, S. Wessel, A. Honecker, B. Normand, C. Rüegg,
    PC, H. M. Rønnow, and F. Mila, Nature 592, 370 (2021).
- ▶ Part II: SCBO under extreme conditions of field & pressure
  - New type of 1/5 plateau and supersolid phases in the (p, H) phase diagram
     Z. Shi, S. Dissanayake, PC, W. Steinhardt, D. Graf, D. M. Silevitch, H. A. Dabkowska,
     T. F. Rosenbaum, F. Mila, S. Haravifard, Nat Commun 13, 1 (2022)
- ▶ Part III: SCBO up to the saturation magnetic field
  - Close agreement between iPEPS and experiments & insights into ultrasound velocity
     T. Nomura, PC, A. Miyata, S. Zherlitsyn, Y. Ishii, Y. Kohama, Y. Matsuda, A. Ikeda,
     C. Zhong, H. Kageyama, F. Mila, arXiv:2209.07652
- ▶ Part IV: Role of interlayer coupling in SCBO
  - ◆ Phase diagram of the SSM with interlayer coupling with new iPEPS approach
     P. Vlaar, PC, arxiv:2208.06423; in preparation

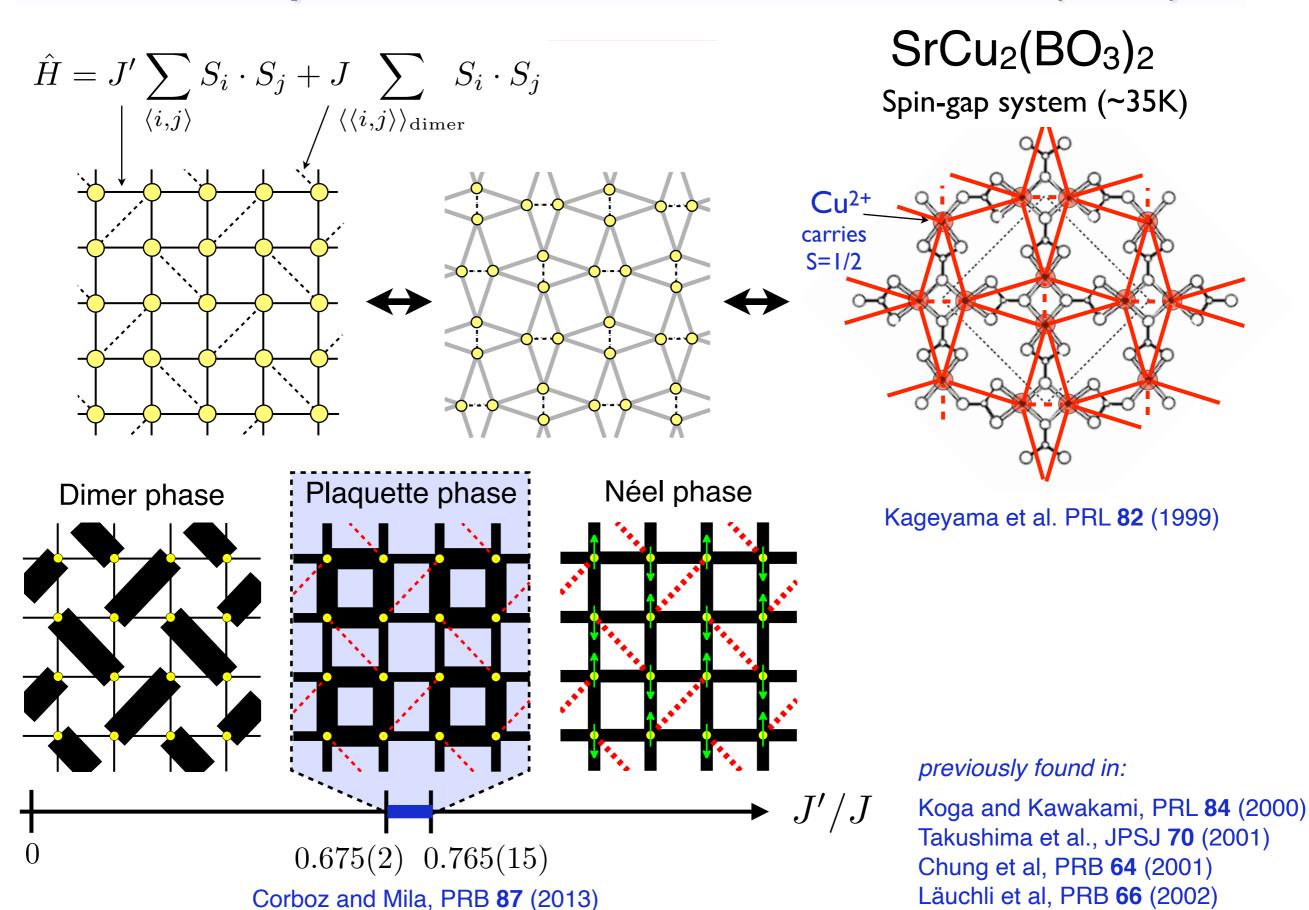
# The Shastry-Sutherland model and SrCu<sub>2</sub>(BO<sub>3</sub>)<sub>2</sub>



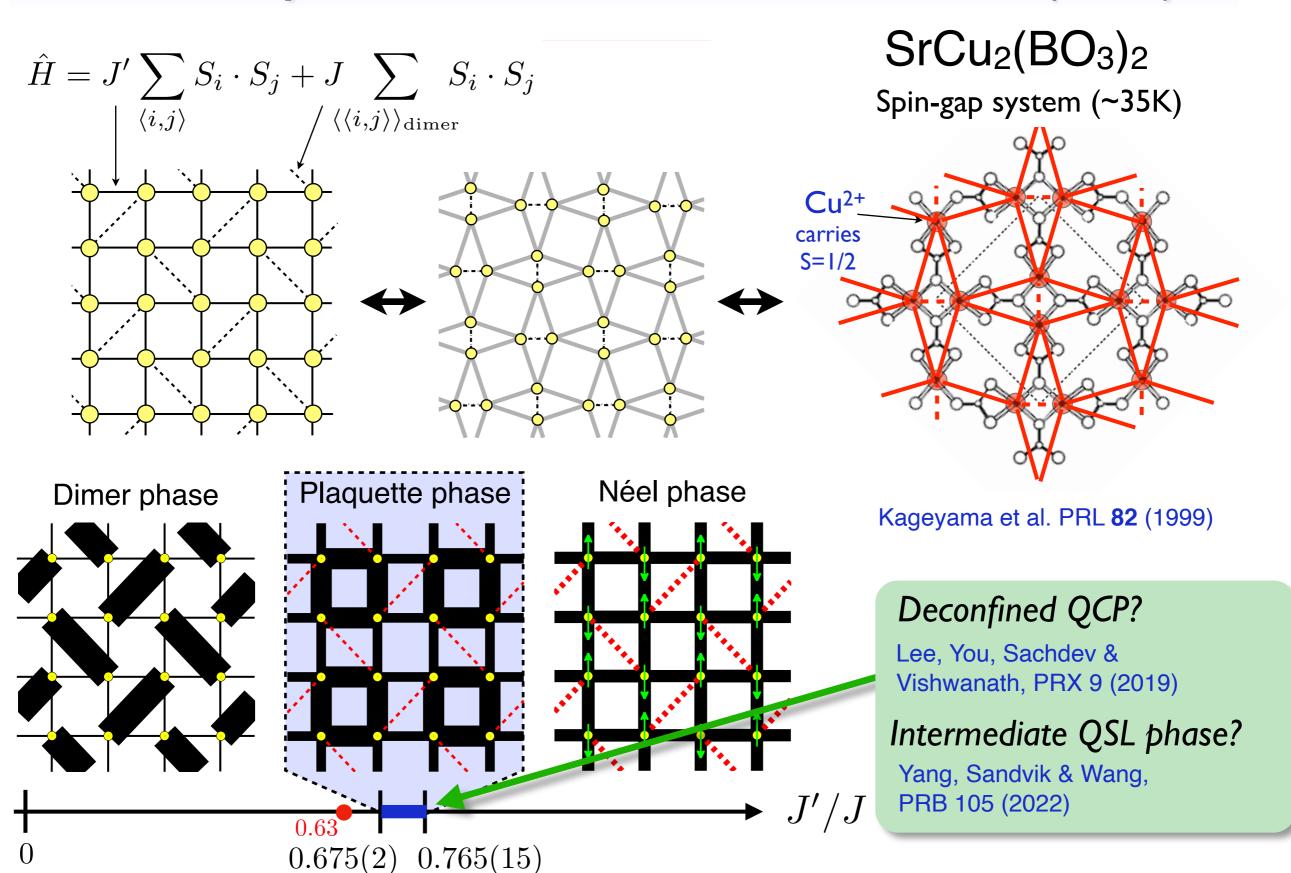
Shastry & Sutherland, Physica B+C 108 (1981).

Kageyama et al. PRL 82 (1999)

# The Shastry-Sutherland model and SrCu<sub>2</sub>(BO<sub>3</sub>)<sub>2</sub>



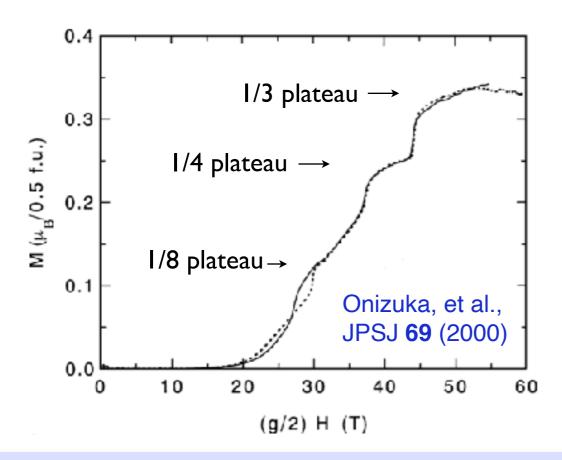
# The Shastry-Sutherland model and SrCu<sub>2</sub>(BO<sub>3</sub>)<sub>2</sub>



Corboz and Mila, PRB 87 (2013)

### Magnetization plateaus

SrCu<sub>2</sub>(BO<sub>3</sub>)<sub>2</sub> in a magnetic field exhibits several magnetization plateaus

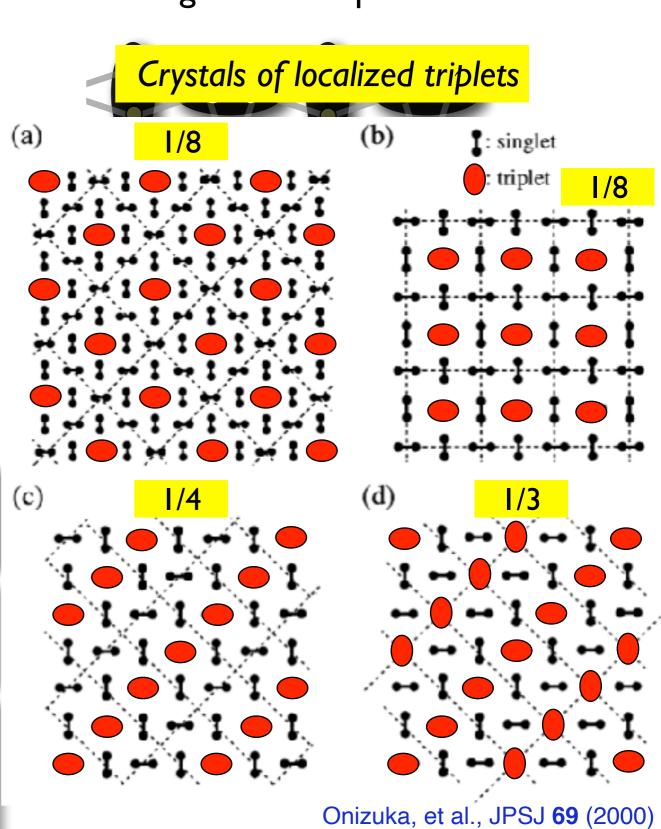


The SSM has almost localized triplet excitations [Miyahara&Ueda'99, Kageyama et al. '00]

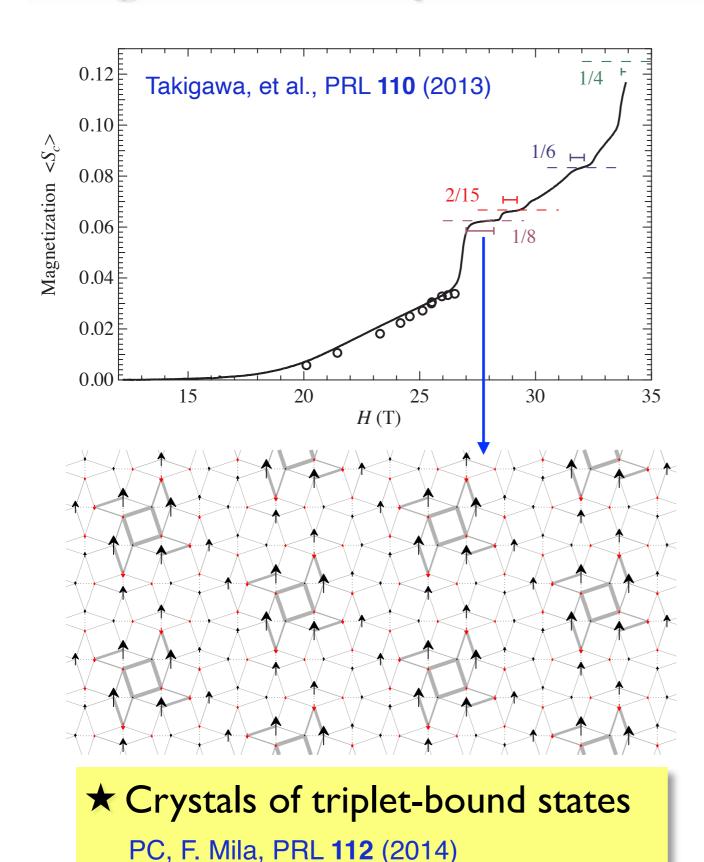
Triplets repel each other (on the mean-field level)

#### Common assumption:

magnetization plateaus correspond to crystals of localized triplets!



### Magnetization plateaus below the 1/4 plateau



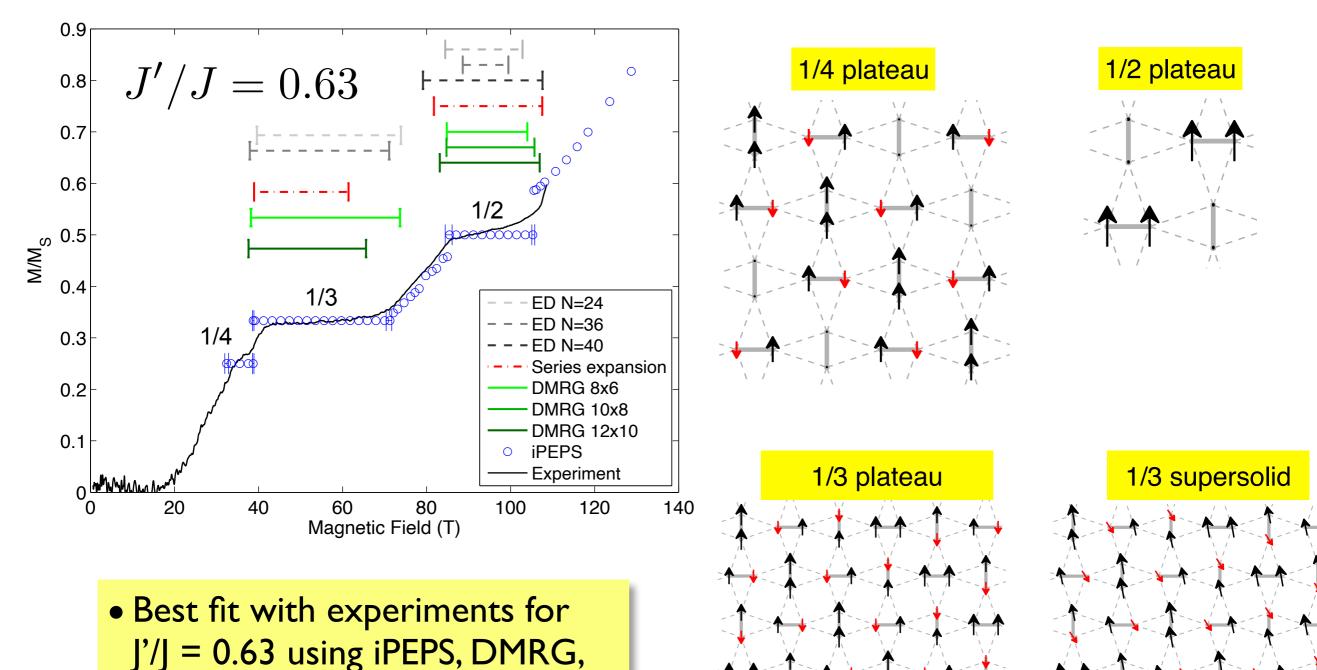
Many experimental / theoretical studies

Kageyama et al, PRL **82** (1999) Onizuka et al, JPSJ **69** (2000) Kageyama et al, PRL **84** (2000) Kodama et al, Science 298 (2002) Takigawa et al, Physica 27 (2004) Levy et al, EPL **81** (2008) Sebastian et al, PNAS **105** (2008) Isaev et al, PRL 103 (2009) Jaime et al, PNAS **109** (2012) Takigawa et al, PRL **110** (2013) Matsuda et al, PRL 111 (2013) Miyahara and K. Ueda, PRL 82 (1999) Momoi and Totsuka, PRB 61 (2000) Momoi and Totsuka, PRB 62 (2000) Fukumoto and Oguchi, JPSJ 69 (2000) Fukumoto, JPSJ **70** (2001) Miyahara and Ueda, JPCM 15 (2003) Miyahara, Becca and Mila, PRB 68 (2003) Dorier, Schmidt, and Mila, PRL 101 (2008) Abendschein & Capponi, PRL 101 (2008) Takigawa et al, JPSJ 79 (2010) Nemec et al, PRB **86** (2012) Matsuda et al., PRL 111 (2013)

...

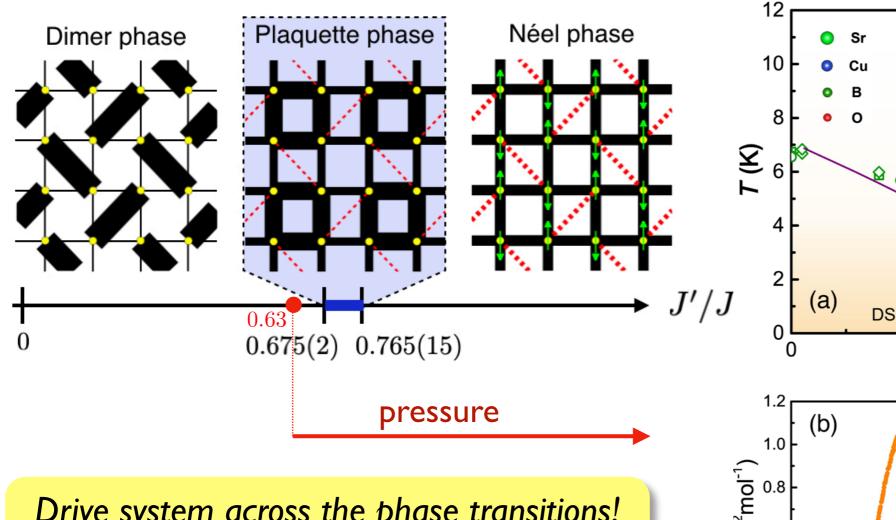
### Magnetization plateaus at high fields

Matsuda, Abe, Takeyama, Kageyama, PC, Honecker, Manmana, Foltin, Schmidt & Mila, PRL 111 (2013)



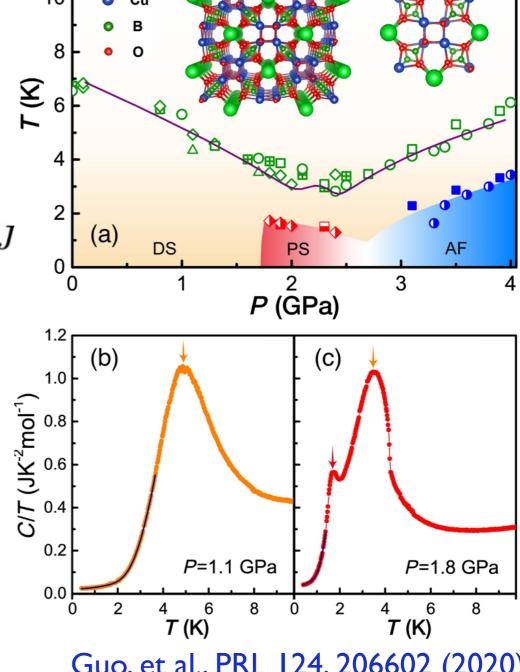
- ED, series expansion
- Supersolid phases at high fields

# SrCu<sub>2</sub>(BO<sub>3</sub>)<sub>2</sub> under pressure



Drive system across the phase transitions!

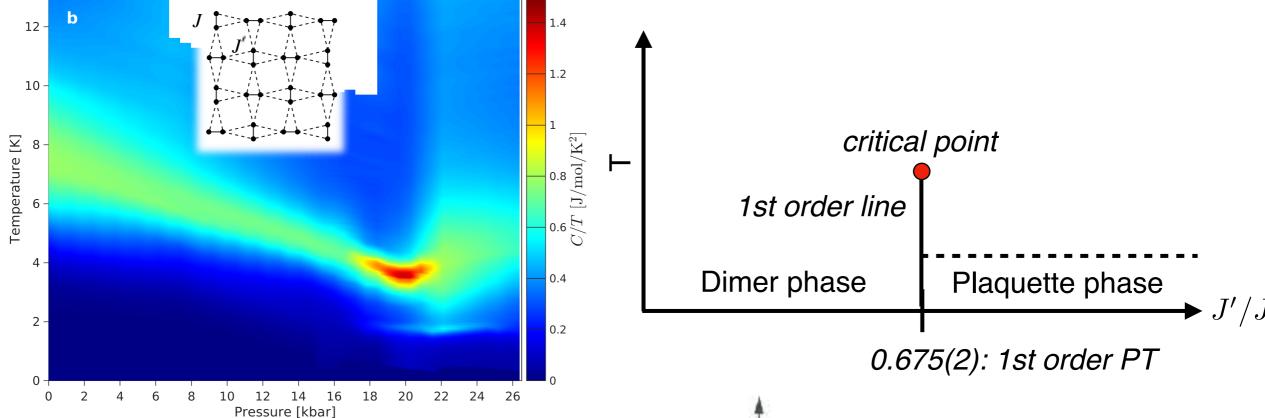
Waki, et al. J. Phys. Soc. Jpn. 76, 073710 (2007) Haravifard, et al. Nat. Commun. 7, 11956 (2016) Zayed, et al., Nat. Phys. 13, 962 (2017) Sakurai, et al., J. Phys. Soc. Jpn. 87, 033701 (2018) Guo, et al., PRL 124, 206602 (2020) Bettler, et al., Phys. Rev. Research 2, 012010 (2020)



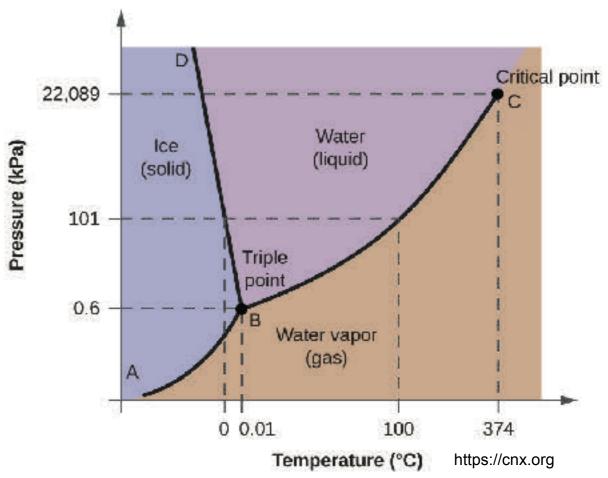
Guo, et al., PRL 124, 206602 (2020)

# Specific heat data (group of H. M. Rønnow)

Jiménez, Crone, et al., Nature 592, 370 (2021)

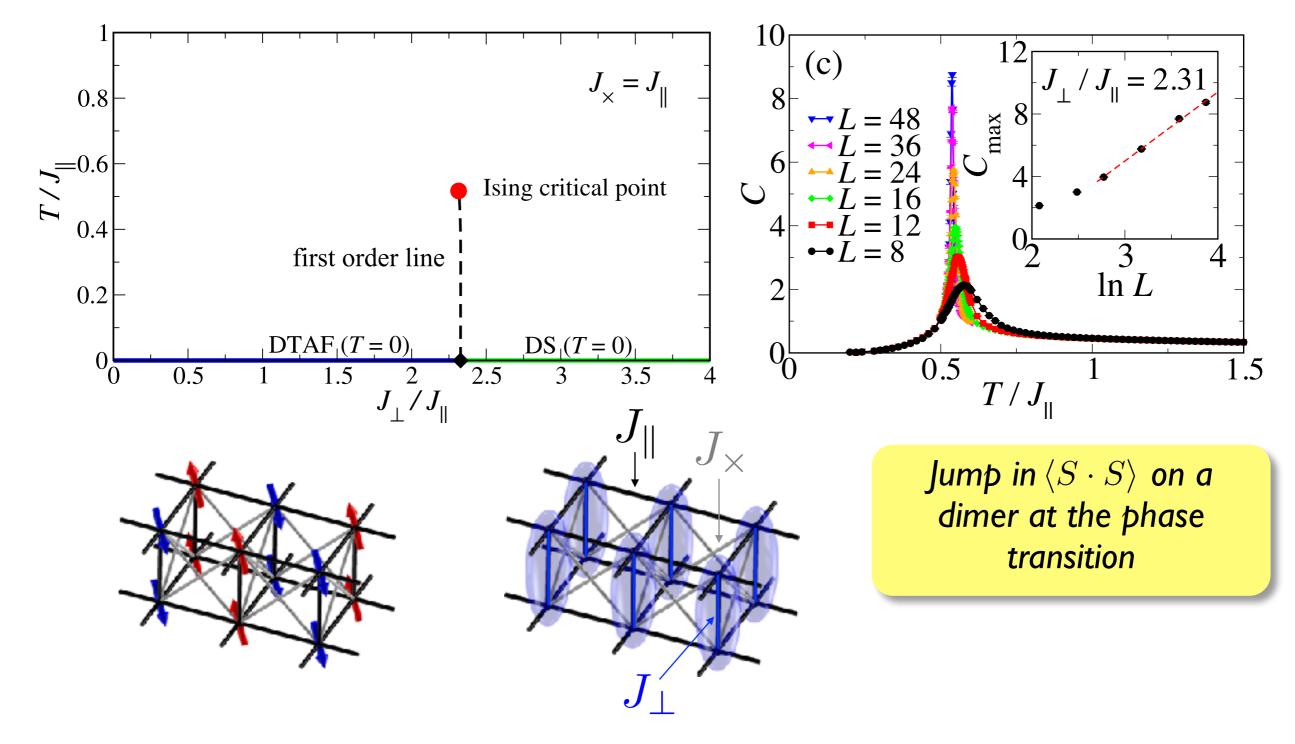


Can we reproduce this with iPEPS?



#### Inspiration: fully frustrated Heisenberg bilayer model

Stapmanns, PC, Mila, Honecker, Normand, and Wessel, PRL 121 (2018)



Dimer triplets

Dimer singlets

#### MPS & PEPS



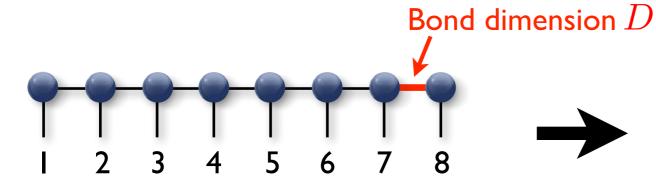
#### **MPS**

Matrix-product state



#### PEPS (TPS)

projected entangled-pair state (tensor product state)

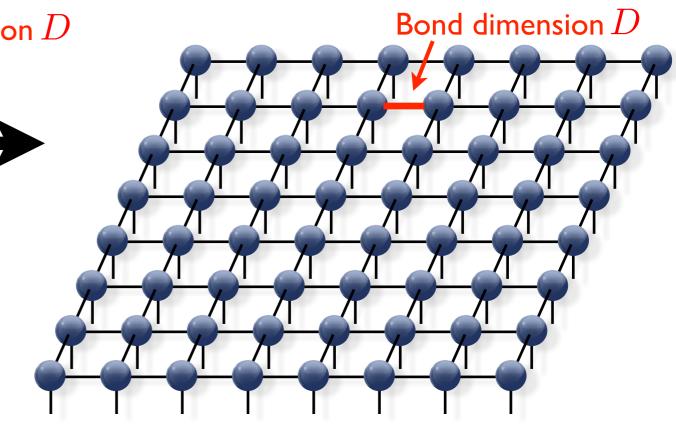


Physical indices (lattices sites)

S. R. White, PRL 69, 2863 (1992)

Fannes et al., CMP 144, 443 (1992)

Östlund, Rommer, PRL 75, 3537 (1995)



F. Verstraete, J. I. Cirac, cond-mat/0407066 Nishio, Maeshima, Gendiar, Nishino, cond-mat/0401115

# Infinite PEPS (iPEPS)



#### **IMPS**

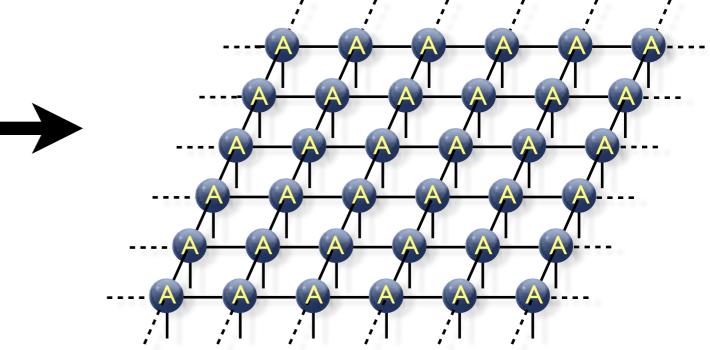
infinite matrix-product state





#### **iPEPS**

infinite projected entangled-pair state



Jordan, Orus, Vidal, Verstraete, Cirac, PRL (2008)

★ Work directly in the thermodynamic limit:
No finite size and boundary effects!

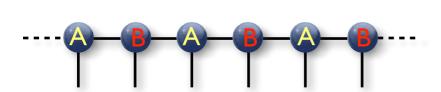
## Infinite PEPS (iPEPS)



#### **IMPS**

infinite matrix-product state

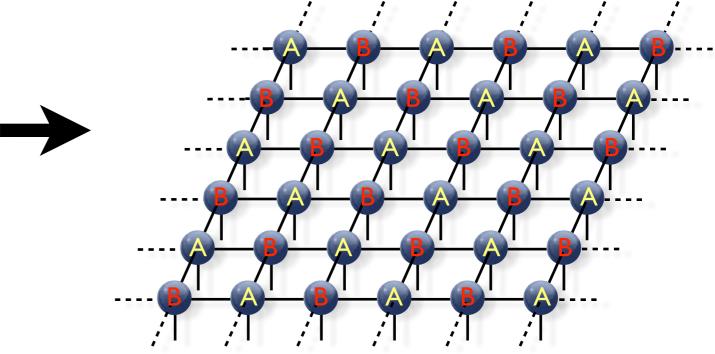






#### **iPEPS**

infinite projected entangled-pair state



Jordan, Orus, Vidal, Verstraete, Cirac, PRL (2008)

★ Work directly in the thermodynamic limit:
No finite size and boundary effects!

### iPEPS with arbitrary unit cells



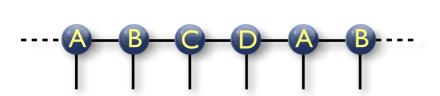
**IMPS** 

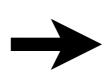
infinite matrix-product state

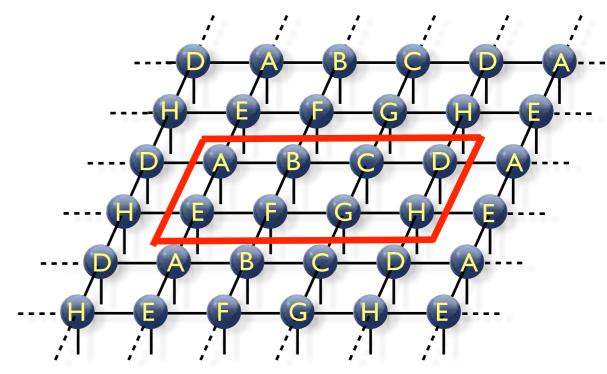


**iPEPS** 

with arbitrary unit cell of tensors







here: 4x2 unit cell

PC, White, Vidal, Troyer, PRB 84 (2011)

★ Run simulations with different unit cell sizes and compare variational energies

### iPEPS ground state simulations

... and many more ...

 Many applications to challenging problems, including frustrated spin, SU(N), and bosonic systems, t-J / Hubbard models, and more, see e.g.:

```
P. Corboz, A. M. Läuchli, K. Penc, M. Troyer and F. Mila, PRL 107 (2011)
S. Dusuel, M. Kamfor, R. Orús, K. P. Schmidt, and J. Vidal, PRL 106, 107203 (2011)
H. H. Zhao, C. Xu, Q. N. Chen, Z. C. Wei, M. P. Qin, G. M. Zhang and T. Xiang, PRB 85 (2012)
P. Corboz, M. Lajkó, A. M. Läuchli, K. Penc and F. Mila, PRX 2 (2012)
P. Corboz and F. Mila, PRB 87 (2013); PRL 112 (2014)
Z.-C. Gu, H.-C. Jiang, D. N. Sheng, H. Yao, L. Balents and X.-G. Wen, PRB 88 (2013)
J. Osorio Iregui, P. Corboz and M. Troyer, PRB 90 (2014)
P. Corboz, T. Rice and M. Troyer, PRL 113 (2014)
T. Picot and D. Poilblanc, PRB 91 (2015)
T. Picot, M. Ziegler, R. Orús and D. Poilblanc, PRB 93 (2016)
P. Nataf, M. Lajkó, P. Corboz, A. M. Läuchli, K. Penc and F. Mila, PRB 93 (2016)
H. Liao, Z. Xie, J. Chen, Z. Liu, H. Xie, R. Huang, B. Normand and T. Xiang, PRL 118 (2017)
B.-X. Zheng, et al., Science 358, 1155 (2017)
I. Niesen and P. Corboz, PRB 95 (2017); SciPost Physics 3, 030 (2017); Rev. B 97, 245146 (2018)
R. Haghshenas, W.-W. Lan, S.-S. Gong, and D. N. Sheng, PRB 97 (2018)
J.-Y. Chen, L. Vanderstraeten, S. Capponi, and D. Poilblanc, PRB 98 (2018)
S. S. Jahromi and R. Orús, PRB 98 (2018)
H.-Y. Lee and N. Kawashima, PRB 97 (2018)
H. Yamaguchi, Y. Sasaki, T. Okubo, et al., PRB 98, 094402 (2018)
R. Haghshenas, S.-S. Gong, and D. N. Sheng, PRB 99, 174423 (2019)
S. S. Chung and P. Corboz, PRB 100 (2019)
B. Ponsioen, S. S. Chung, and P. Corboz, PRB 100 (2019)
C. Boos, S. P. G. Crone, I. A. Niesen, P. Corboz, K. P. Schmidt, and F. Mila, PRB 100 (2019)
Z. Shi, et al, Nature Communications 10, 2439 (2019)
A. Kshetrimayum, C. Balz, B. Lake, and J. Eisert, ArXiv:1904.00028 (2019)
H.-Y. Lee, R. Kaneko, T. Okubo, and N. Kawashima, PRL 123, 087203 (2019).
O. Gauthé, S. Capponi, M. Mambrini, and D. Poilblanc, PRB 101, 205144 (2020).
H.-Y. Lee, R. Kaneko, L. E. Chern, T. Okubo, Y. Yamaji, N. Kawashima, and Y. B. Kim, Nature Communications 11 (2020)
W.-Y. Liu, S.-S. Gong, Y.-B. Li, D. Poilblanc, W.-Q. Chen, and Z.-C. Gu, ArXiv:2009.01821 (2020)
J.-Y. Chen, S. Capponi, A. Wietek, M. Mambrini, N. Schuch, and D. Poilblanc, PRL 125, 017201 (2020)
J. Hasik, D. Poilblanc, and F. Becca, SciPost Physics 10, 012 (2021)
```

#### Finite temperature simulations with iPEPS

▶ Methodological developments (2D):

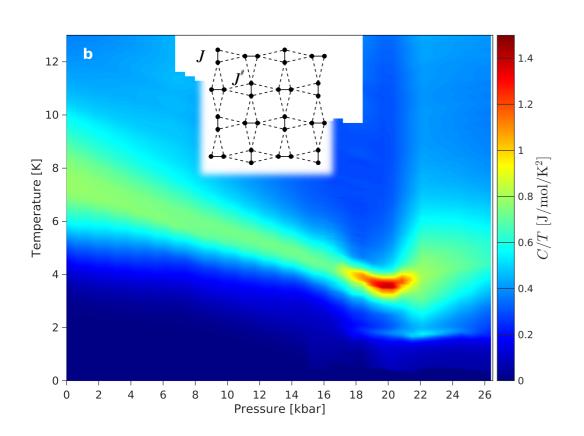
Li et al. PRL 106 (2011); Czarnik et al. PRB 86 (2012); Czarnik & Dziarmaga PRB 90 (2014); Czarnik & Dziarmaga PRB 92 (2015); Czarnik et al. PRB 94 (2016); Dai et al PRB 95 (2017); Kshetrimayum, Rizzi, Eisert, Orus, PRL 122 (2019), P. Czarnik, J. Dziarmaga, PC, PRB 99 (2019), ...

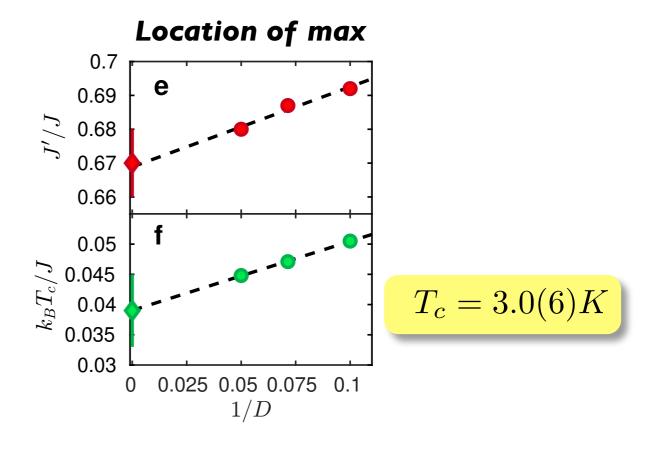
> Density-operator:  $\hat{\rho}=e^{-\beta\hat{H}}\approx$ 

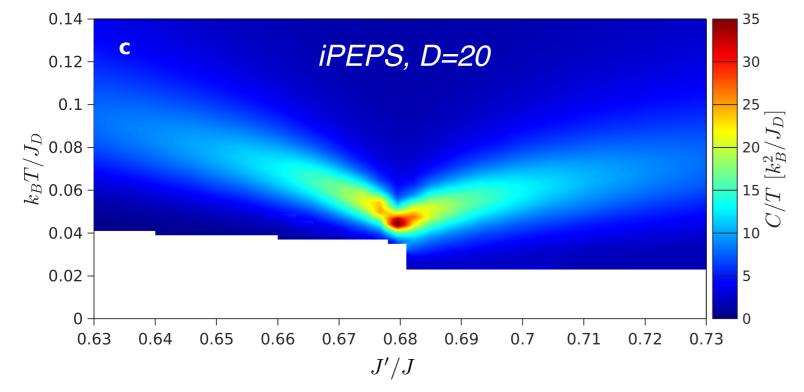
> Symmetric form:  $e^{-\beta \hat{H}/2} \approx \hat{\rho}(\beta) \approx \hat{\rho}(\beta)$ 

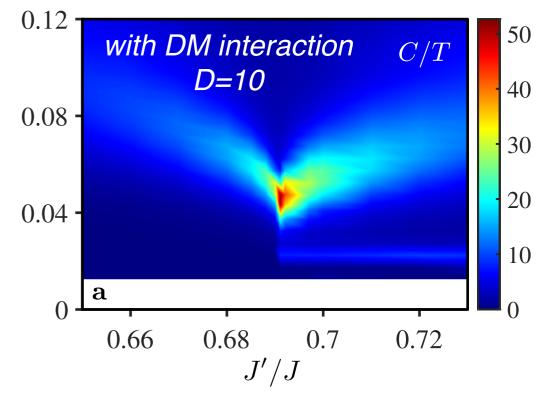
$$\hat{\rho}(\beta) = \hat{\rho}^{\dagger}(\beta)$$
 by construction

### Specific heat: experiments vs iPEPS



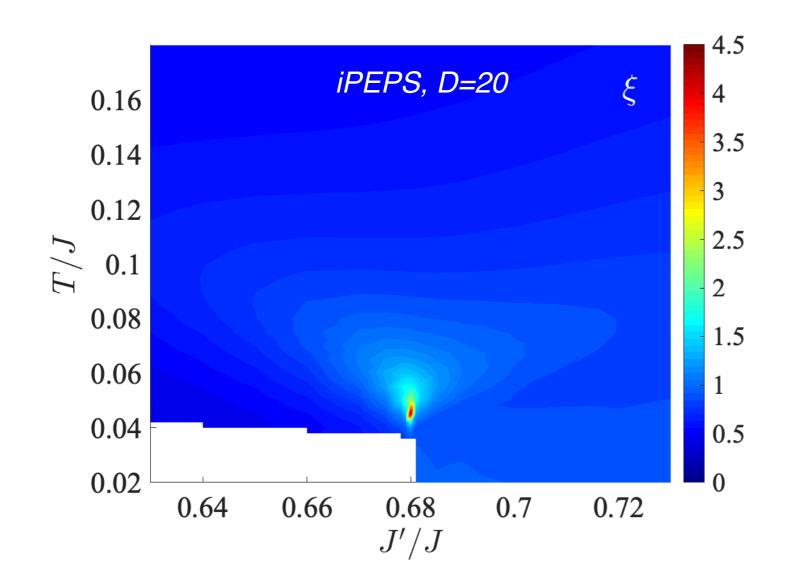




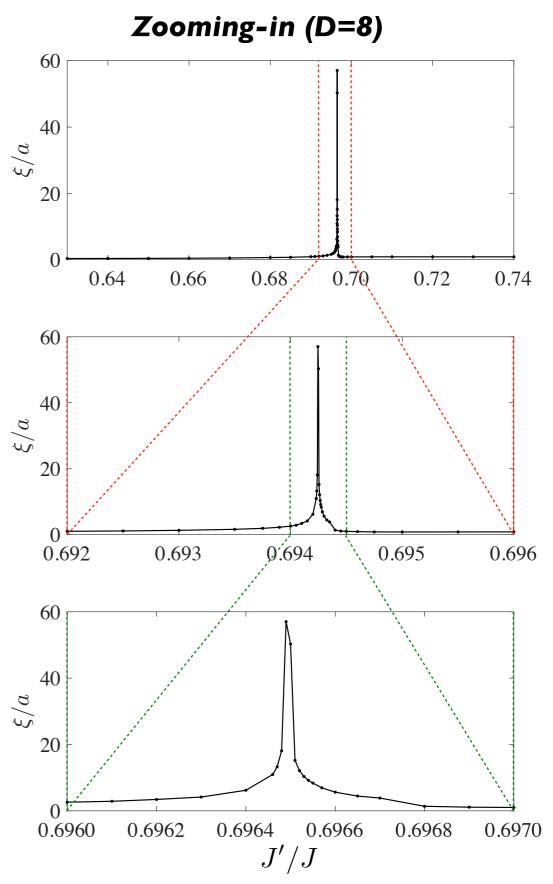


Jiménez, Crone, et al., Nature 592, 370 (2021)

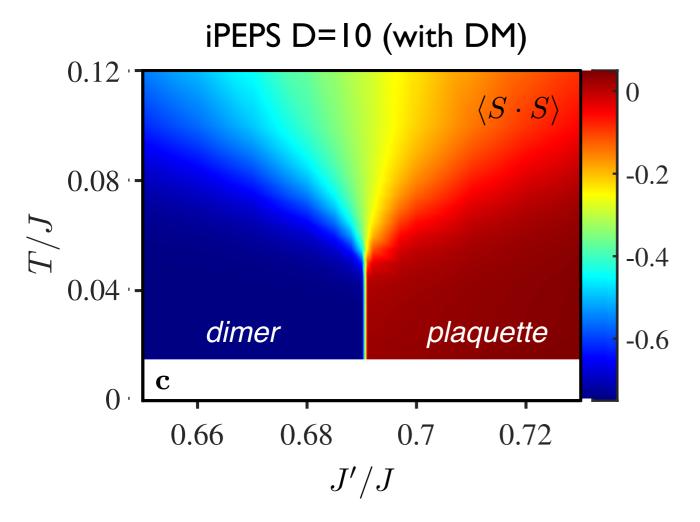
### Correlation length



Diverging correlation length compatible with a critical point

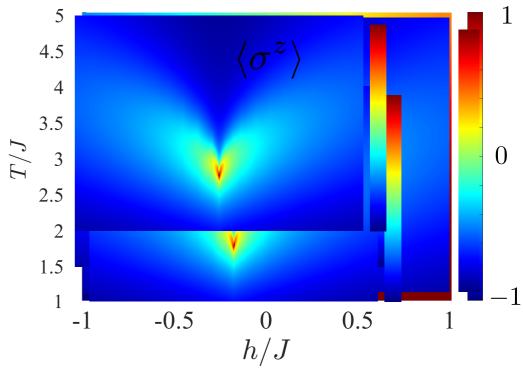


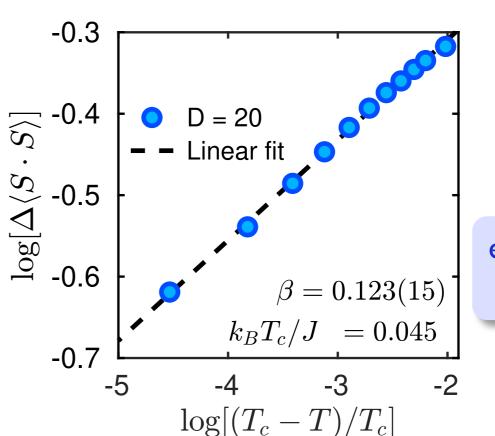
# Jump in $<S \cdot S>$ on dimer



Clear evidence of a first order line with a critical point compatible with the 2D Ising universality class

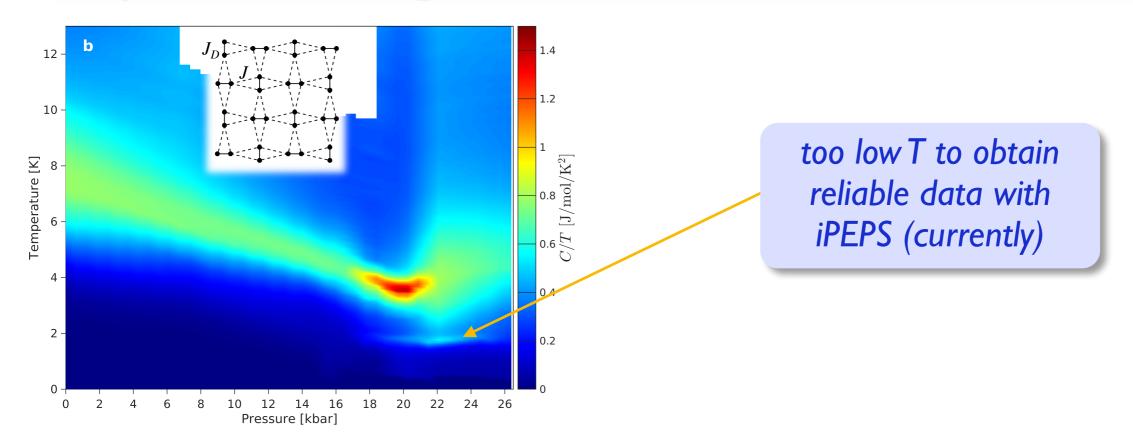


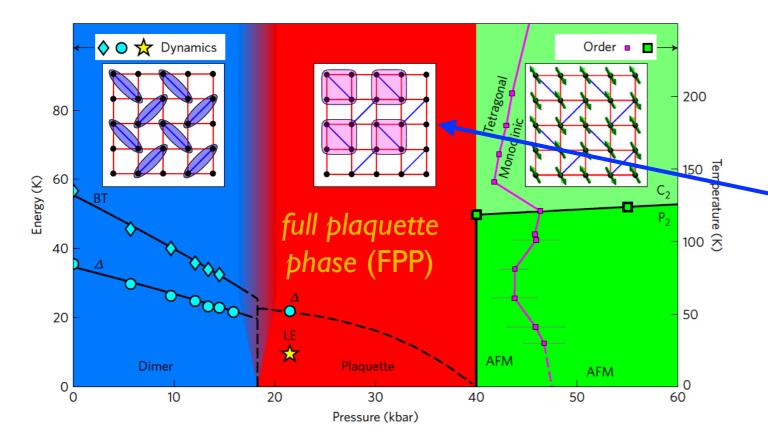




exact value:  $\beta = 0.125$ 

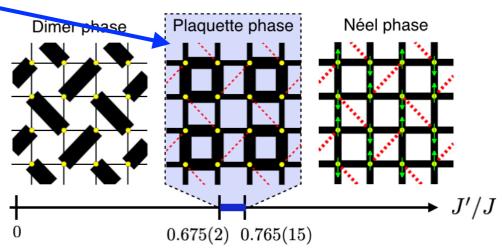
# Open challenges



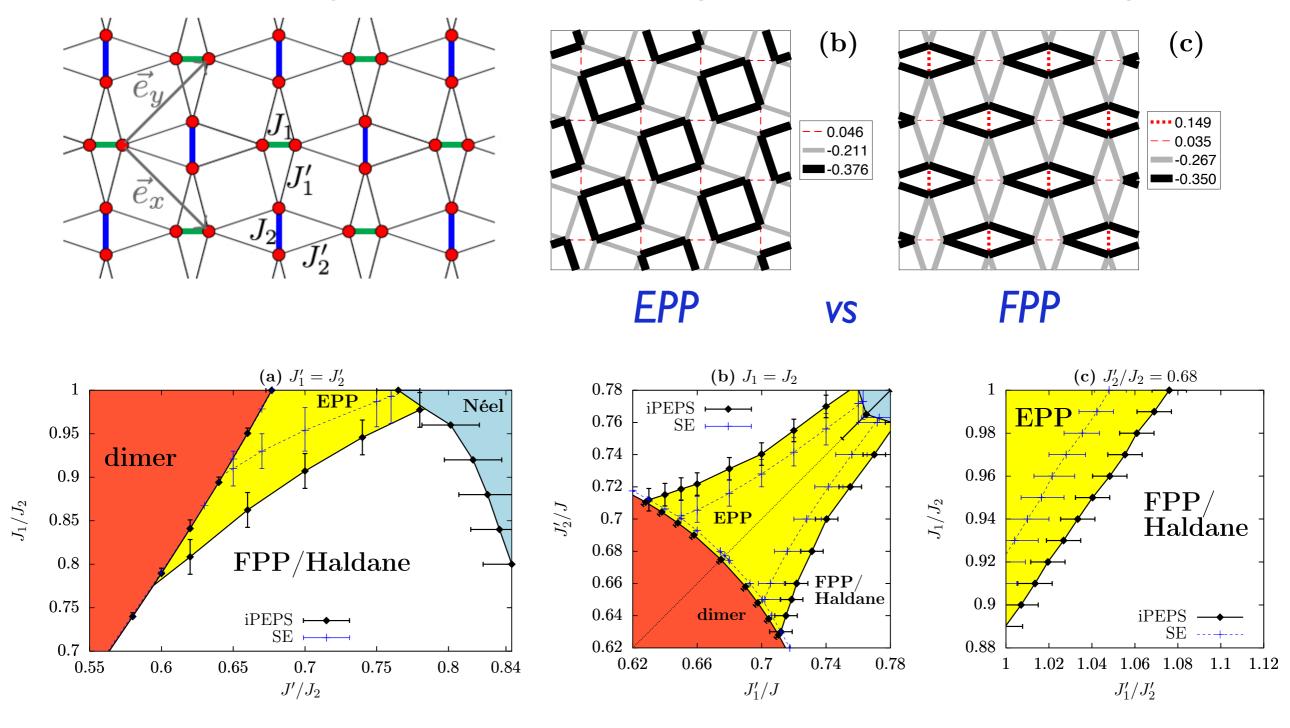


Inelastic neutron scattering: full plaquette phase (FPP), not empty plaquette phase (EPP)

Zayed, et al., Nat. Phys. 13, 962 (2017)



▶ Distorted Shastry-Sutherland model: competition between EPP and FPP phase



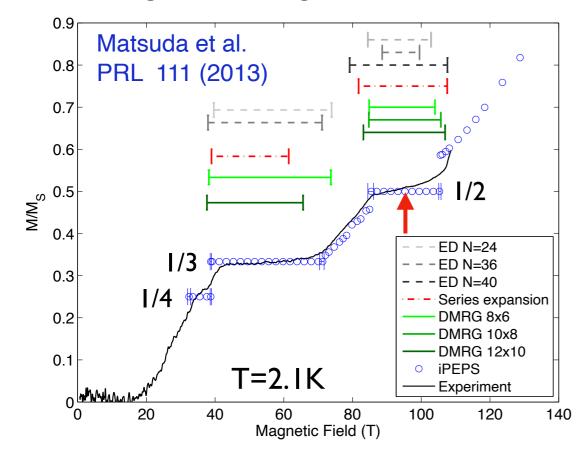
▶ Small deformation leads to FPP phase!

But precise model still unclear...

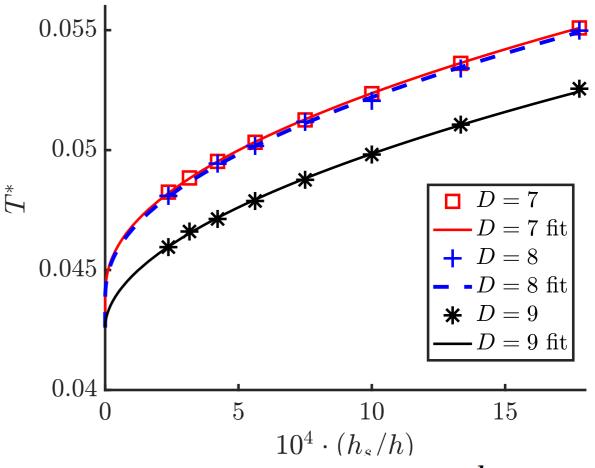
#### Finite T iPEPS study of the m=1/2 plateau in SCBO

P. Czarnik, M. M. Rams, PC, and J. Dziarmaga, PRB 103, 075113 (2021)

#### High-field magnetization data



#### Systematic scaling analysis



$$T^*(h_s, \xi_D) = T_c + ah_s^{1/\tilde{\beta}\delta} + \frac{b}{\xi_D^c} h_s^{(1-c\nu)/\tilde{\beta}\delta}$$

Critical exponents compatible with 2D Ising universality class

$$T_c = 0.043(2)J \approx 3.5(2)K$$

#### Part II: SCBO under extreme conditions of field & pressure

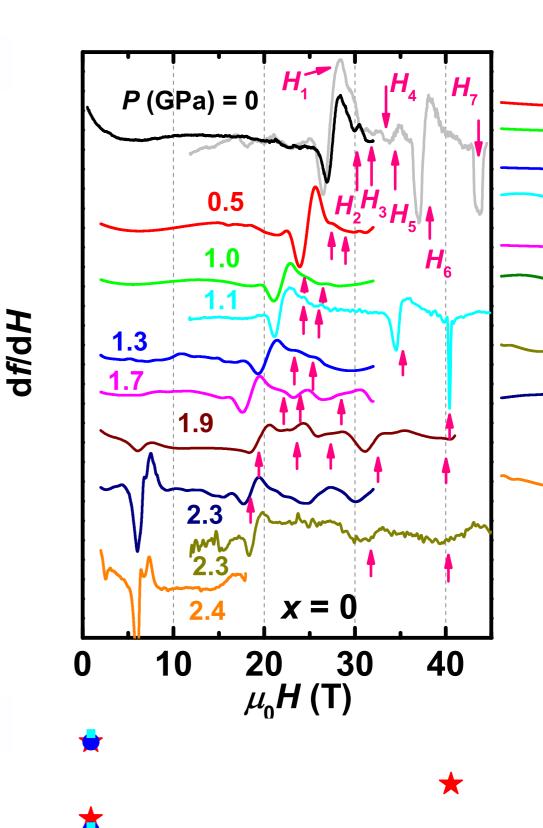
Shi, Dissanayake, PC, William Steinhardt, Graf, Silevitch, Dabkowska, Rosenbaum, Mila, Haravifard, Nat Commun 13, 1 (2022)

Experiments: tunnel diode oscillator (TDO) technique

$$\frac{df}{dH} \propto -\frac{d^2M}{dH^2}$$

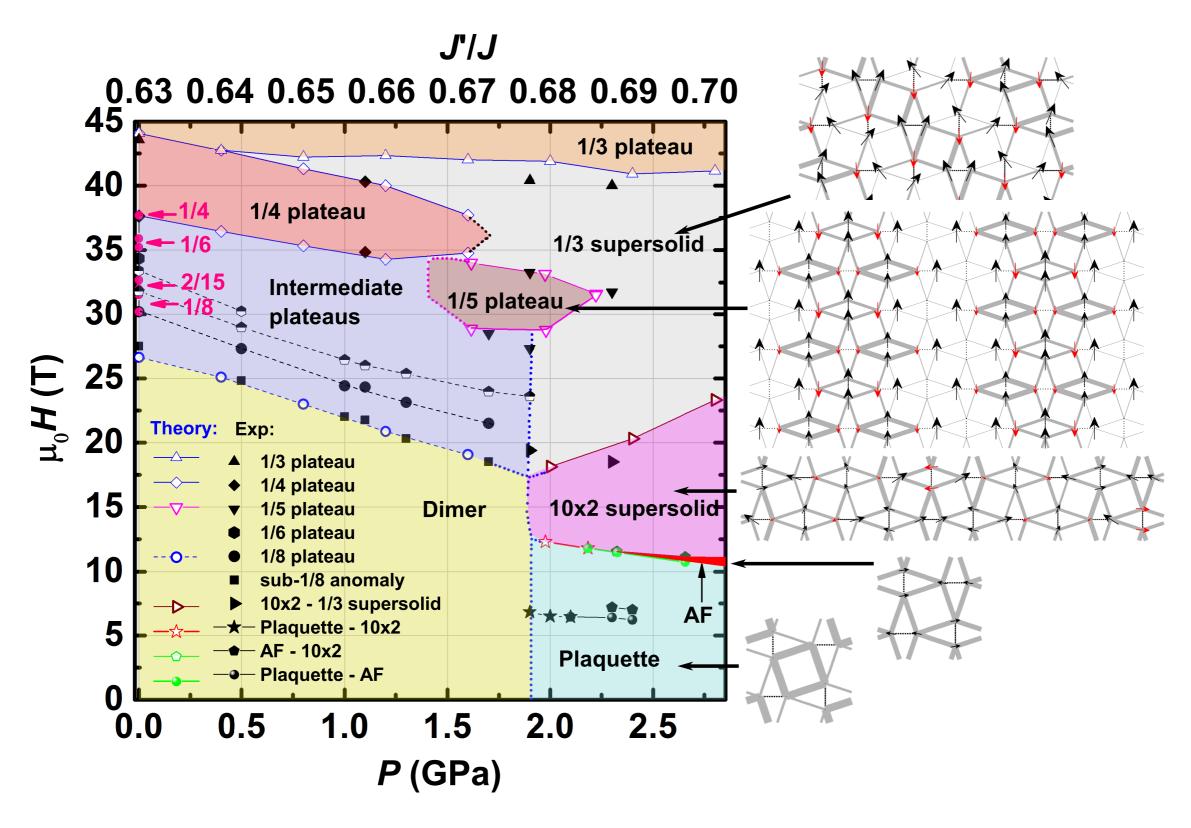
Non-zero df/dH ↔ slope change in M

- ▶ Identify anomalies → phase boundaries
- Compare with iPEPS phase diagram

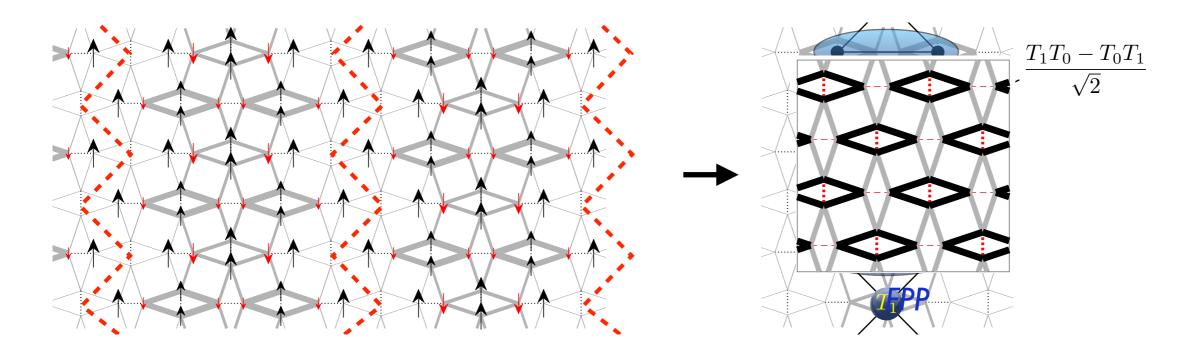


#### SrCu<sub>2</sub>(BO<sub>3</sub>)<sub>2</sub> under extreme conditions of field & pressure

Shi, Dissanayake, PC, William Steinhardt, Graf, Silevitch, Dabkowska, Rosenbaum, Mila, Haravifard, Nat Commun 13, 1 (2022)



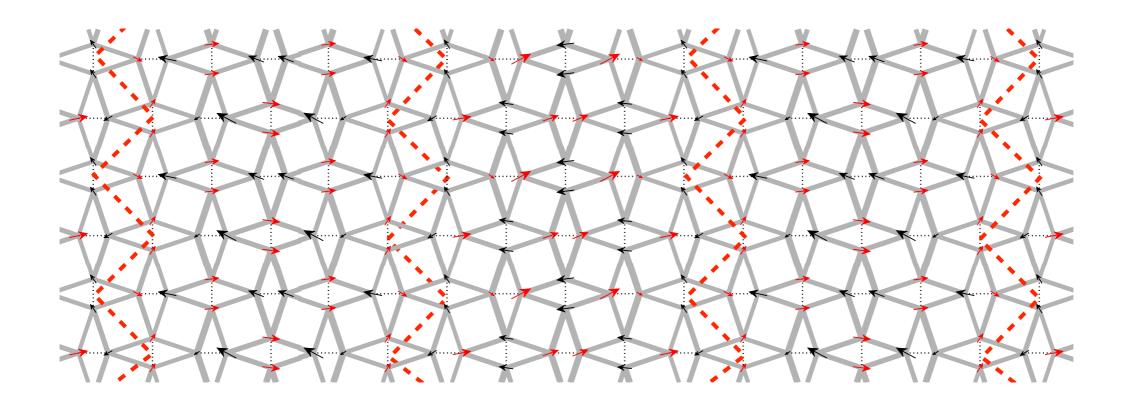
### Nature of the 1/5 plateau



- Vertical stripes separated by dimer singlets (along red dashed lines)
- Strong triplets in the center of each stripe, with a weaker pair of triplets in between
  - ♦ Neither a crystal of triplets, nor a crystal of bound states!
- ▶ Full-plaquette formation: reminiscent of full plaquette phase (FPP)
  - **♦** FPP: triplets on dimers within plaquette and singlets on dimers outside of plaquette, where the triplets form effective S=1 Haldane chains
- ▶ Effective description: S=I diamond chain with m=2/3
- ▶ Also found in a thin SSM tube made of 2 orthogonal dimer chains

Manmana, Picon, Schmidt, and Mila, EPL 94, 67004 (2011)

### Nature of the 10x2 supersolid



- ▶ Descendant of the I/5 plateau state
- Alternating rotation of the spins of successive stripes clockwise or counterclockwise by 90 degrees
- Finite component in the field direction, also on the boundary between stripes

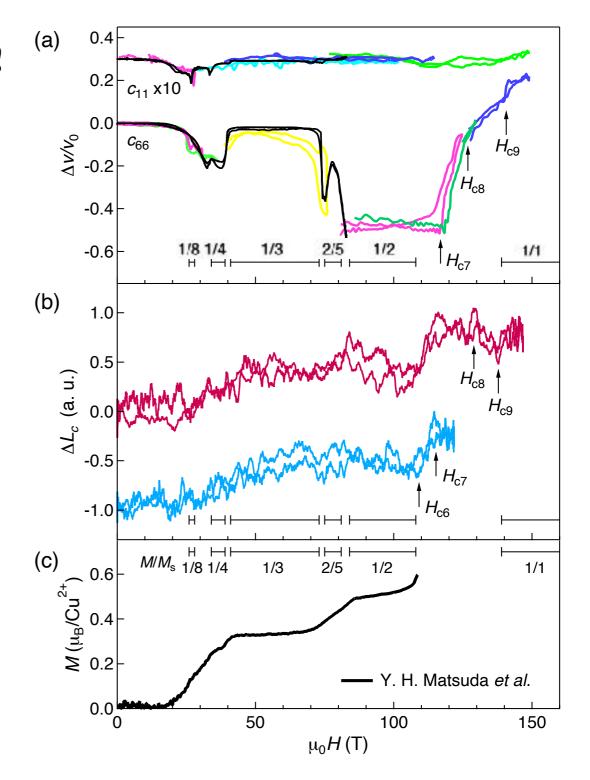
Full plaquette physics appearing at finite magnetic field!

### Part III: SrCu<sub>2</sub>(BO<sub>3</sub>)<sub>2</sub> up to the saturation field

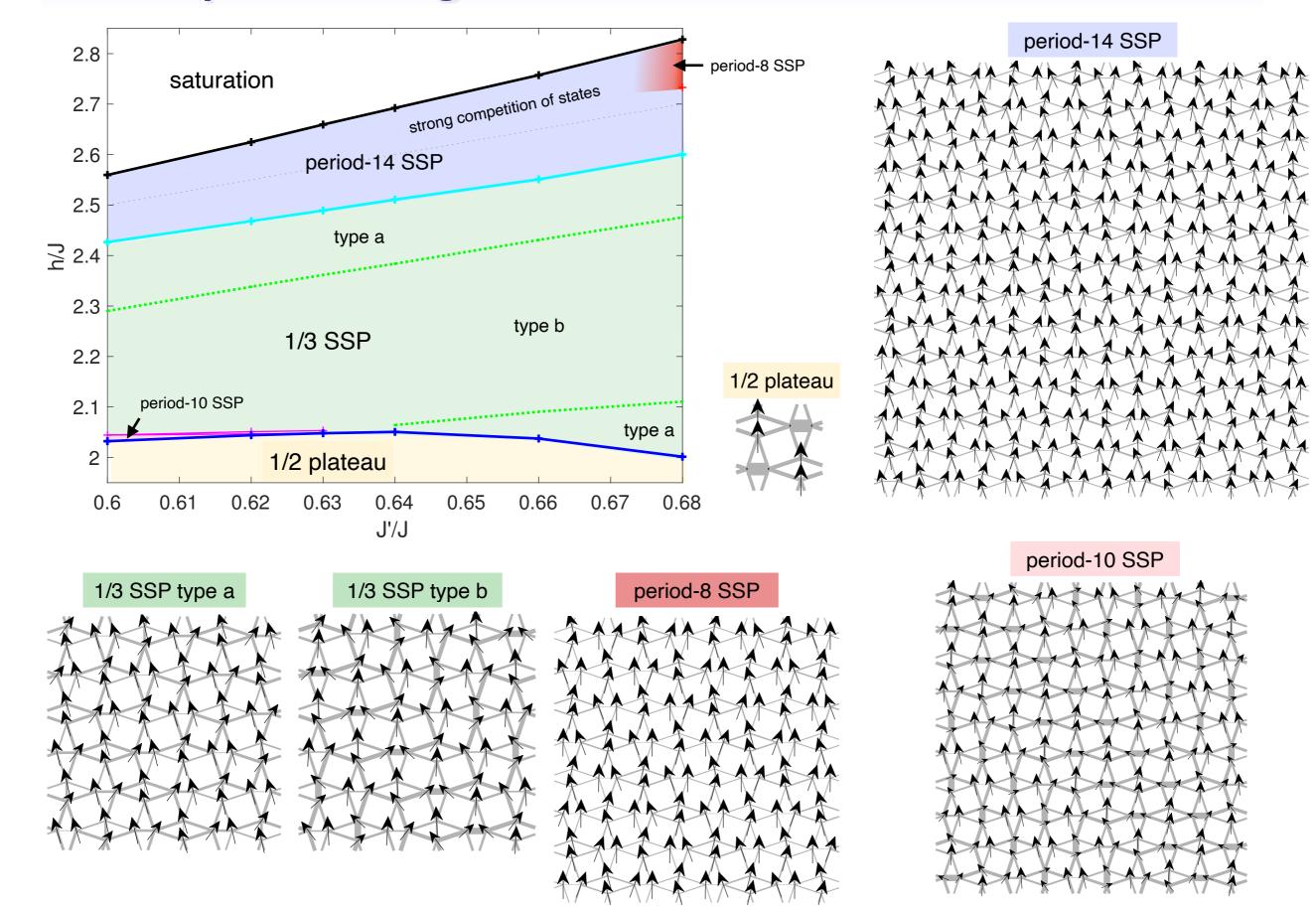
Nomura, PC, Miyata, Zherlitsyn, Ishii, Kohama, Matsuda, Ikeda, Zhong, Kageyama, Mila, arXiv:2209.07652

- ▶ Experiments: ultrahigh fields up to **I50T!**
- Ultrasound velocity & magnetostriction
- ▶ Identified new anomalies above the I/2 plateau at I 16T, I 27T, and I 39T
- ▶ Saturation field: I39T

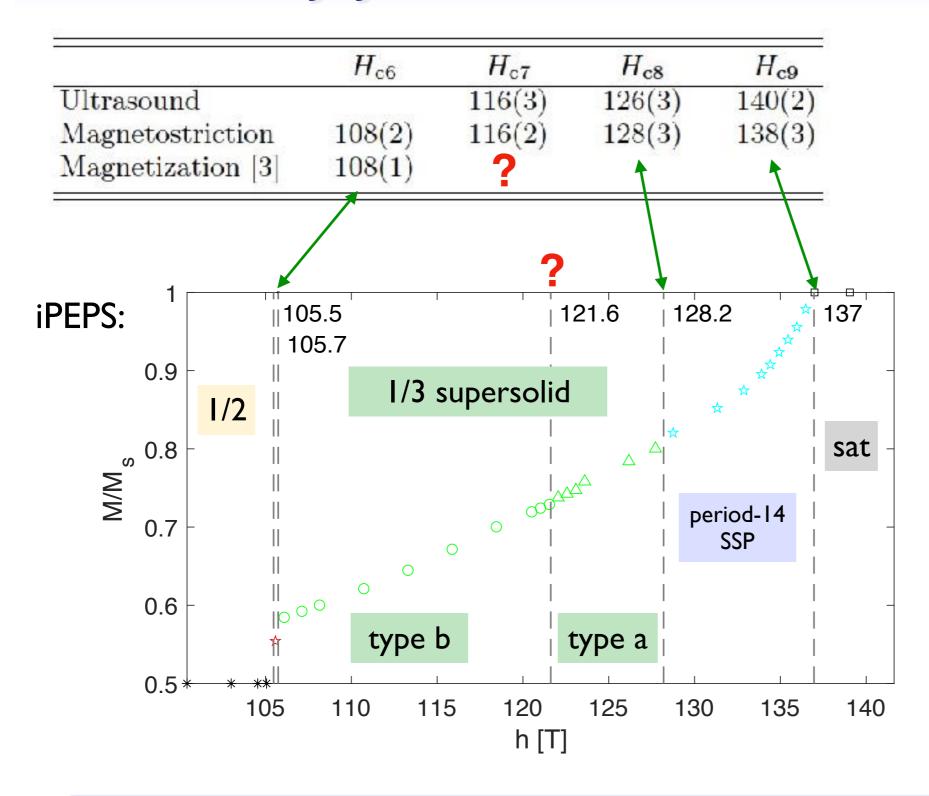
- Comparison with iPEPS results
- Strong decrease of the sound velocity of the c<sub>66</sub> acoustic in the I/2 plateau?



## iPEPS phase diagram

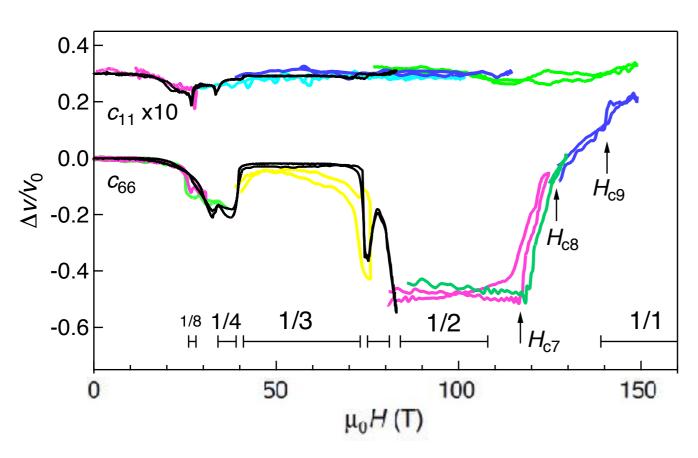


## Results for J'/J=0.63

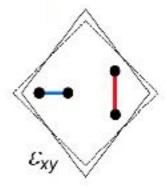


Anomalies  $H_{c6}$ ,  $H_{c8}$ , and  $H_{c9}$  compatible with iPEPS results

### Reduction of ultrasound velocity in 1/2 plateau?

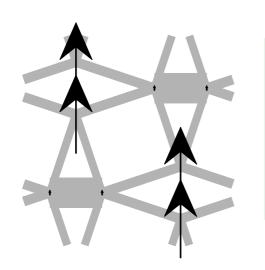


Strain of c66 mode:



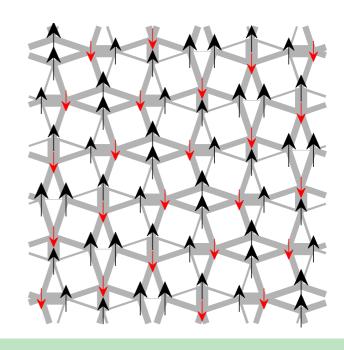
- ▶ Strong reduction in 1/2 plateau
- ▶ Tiny reduction in 1/3 plateau

- ▶ Both E' and E" of the magnetic energy contribute to the elastic constant
- ▶ E' has large magnitude in 1/2, but vanishing for 1/3 plateau

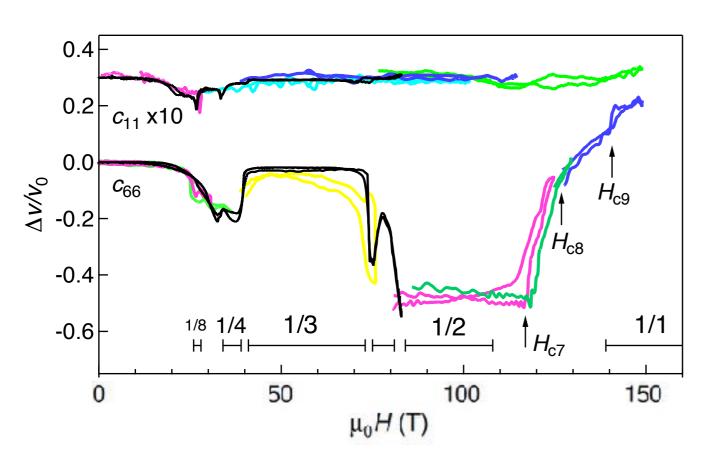


Checkerboard structure with positive (triplet) and negative (singlet) bonds

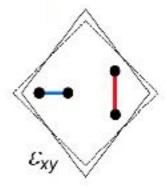
→ contributions to E' add up → large magnitude



## Reduction of ultrasound velocity in 1/2 plateau?



Strain of c66 mode:



- ▶ Strong reduction in 1/2 plateau
- ▶ Tiny reduction in 1/3 plateau

▶ Both E' and E" of the magnetic energy contribute to the elastic constant

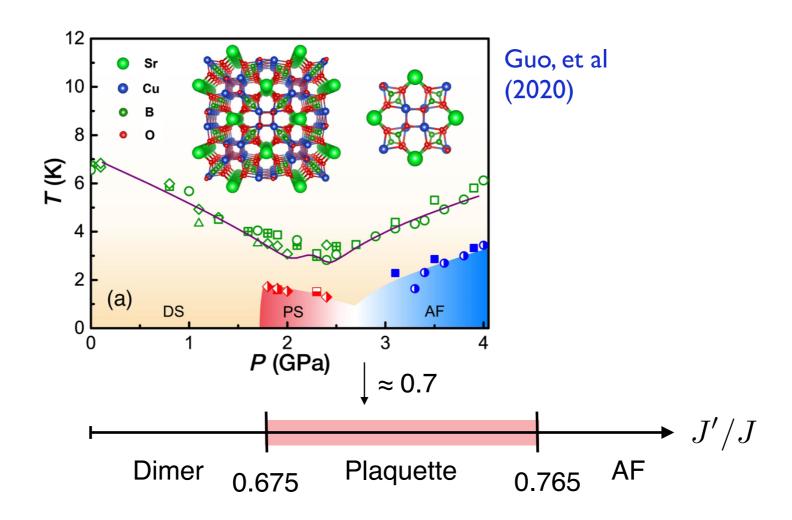
Fit: 
$$\lambda_1 E' + \lambda_2 E'' \sim \Delta c/c$$
,  $c = \rho v^2$ 

Plateau	E'	$E^{\prime\prime}$	$\Delta v/v_0$	$\Delta v/v_0$
			<b>iPEPS</b>	Exp.
1/8	-0.014	-0.44	-0.13	-0.11(4)
1/4	-0.051	-0.31	-0.16	-0.19(3)
1/3	0	-0.18	-0.04	-0.03(2)
1/2	-0.23	-0.034	-0.49	-0.48(7)

Estimates in good agreement with experiment

# Part IV: SSM with interlayer coupling

Extent of the plaquette phase is smaller in experiments than in theory



▶ Pressure model: Shi, et al. (2022)

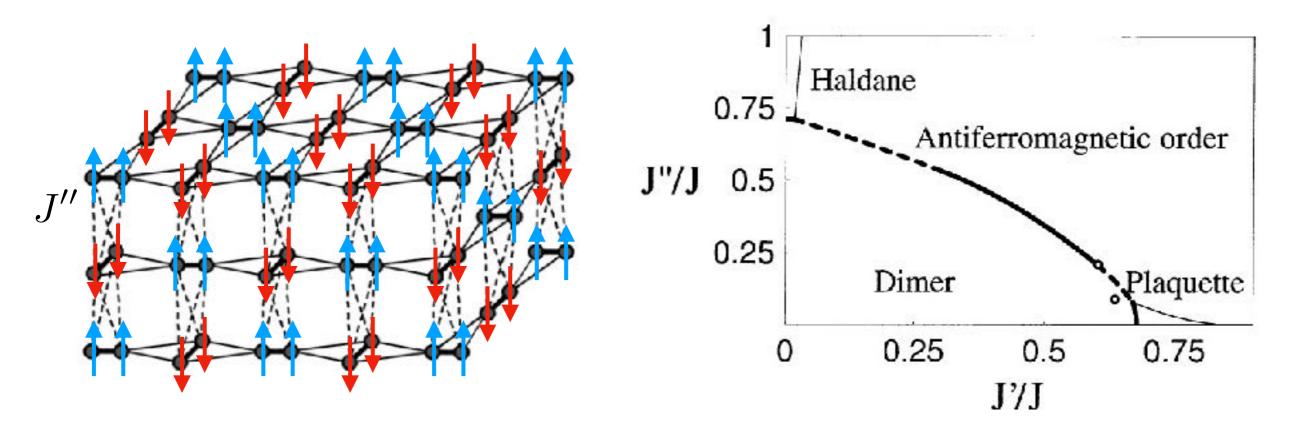
$$p_0 \leftrightarrow J'/J = 0.63$$
  $p_c = 1.8GPa \leftrightarrow J'/J = 0.675$   $J = 81.5K$ 

J'(p), J(p): linear functions. J'(p) changes by 5% between  $p_0$  and  $p_c$  (ESR)

Sakurai, et al., J. Phys. Soc. Jpn. 87, 033701 (2018)

# SSM with interlayer coupling

▶ Series expansion results Koga, J. Phys. Soc. Japan 69 (2000)

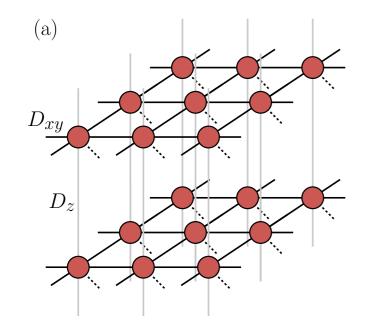


AF phase becomes favored over plaquette phase with increasing J"

- Predicted values for J" (no consensus yet)
  - → J"/J = 0.09 ... 0.21 from fits to susceptibility [Miyahara & Ueda (2000), Knetter et al (2000)]
  - ♦ J"/J ≈ 0.03 from ab-initio calculations [Radtke et al., PNAS 112 (2015)]

### iPEPS for layered systems

Vlaar, PC, arxiv:2208.06423



#### **Ansatz:**

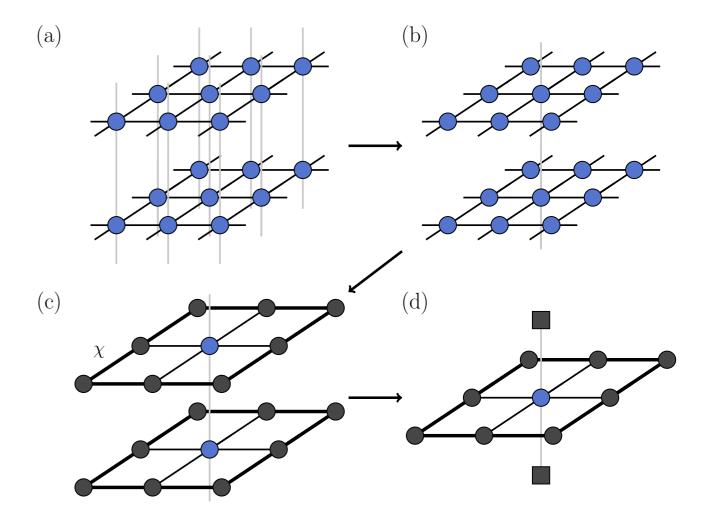
- ▶ 3D tensor network ansatz (coupled iPEPS)
- $D_{xy} > D_z$  for weak interlayer coupling
- ▶  $D_z = I \rightarrow \text{product state of iPEPSs}$



Patrick Vlaar

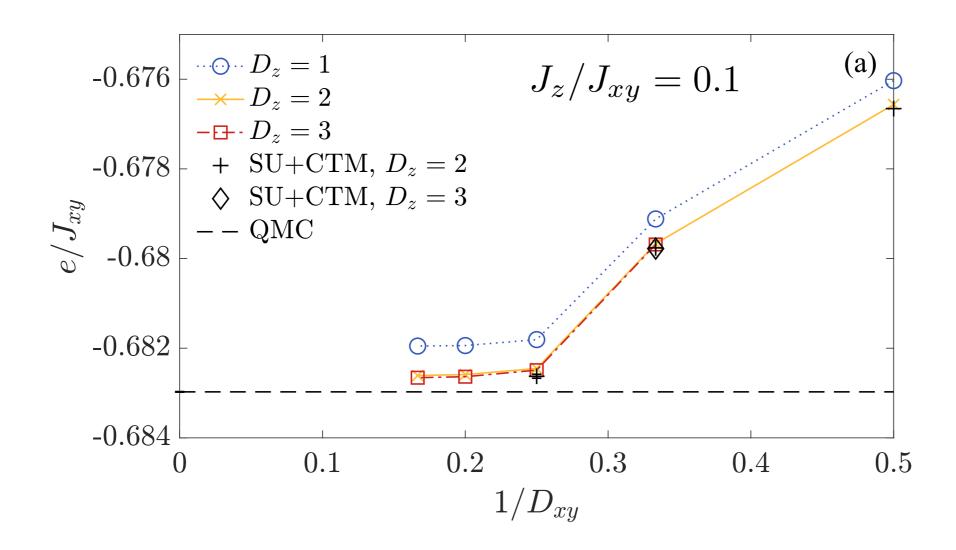
#### **Contraction:**

- $\rightarrow$  D<sub>z</sub> = I : contract individual layers (2D)
- ▶  $D_z > I$ : perform effective decoupling away from center → 2D contraction
- Interlayer correlations beyond mean-field level are included by the  $D_z > 1$  bonds in the center
- Layered corner transfer matrix (LCTM) method



#### Benchmarks for 3D anisotropic Heisenberg model

Vlaar, PC, arxiv:2208.06423

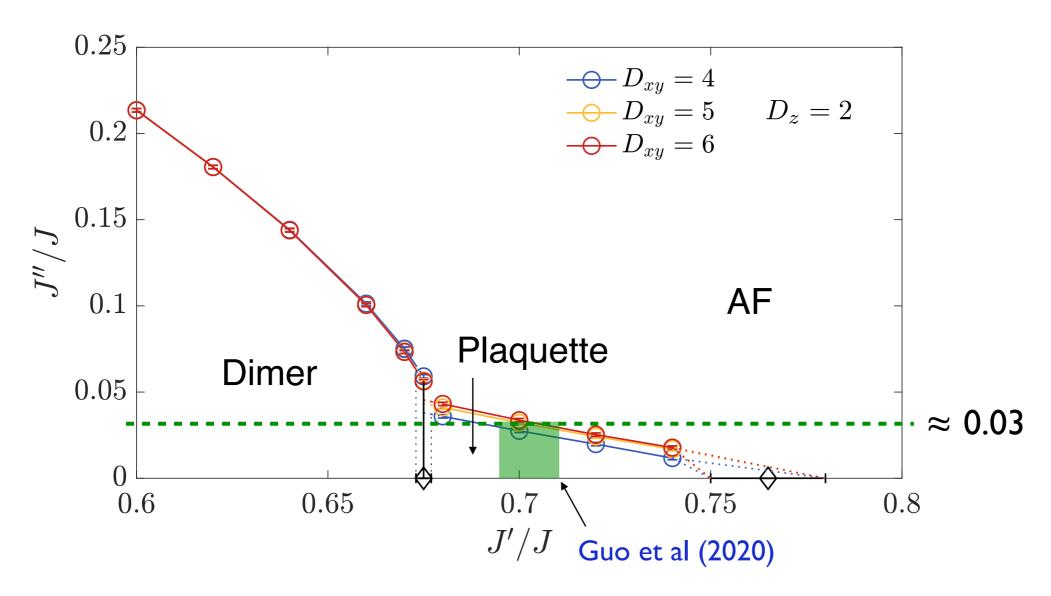


- ▶ Substantial improvement from  $D_z = 1$  to  $D_z = 2$
- Values close to the extrapolated QMC result
- ▶ In agreement with more expensive full 3D contractions

Vlaar & PC, PRB 103, 205137 (2021)

#### Phase diagram: SSM with interlayer coupling

(work in progress)



Estimate for the strength of interlayer coupling: J"/J  $\approx 0.03$ 

LCTM: Promising approach also for other layered systems!

#### Conclusion

- ✓  $SrCu_2(BO_3)_2$  under pressure / in a magnetic field exhibits very rich physics!
- Finite temperature Ising critical point, analogous to critical point of water
- New type of 1/5 plateau and supersolid phases at high pressure & field
- √ Results up to saturation in good agreement with experiments & understanding of the reduction of ultrasound velocity in the 1/2 plateau
- √ Reduction of plaquette phase due to interlayer coupling
- Progress with iPEPS: versatile tool for ground state and finite temperature calculations + extensions to layered systems

#### Thank you for your attention!

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