Mesoscopic transport with ultracold atomic gases

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- Small sample (conduction channel) attached to macroscopic reservoirs
- Electric current induced by external bias voltage







C. Rossler et al., APL 93, 071107 (2008)

Fractional charge / measurement in FQHE



R. de-Picciotto et al., Nature **389**, 162 (1997); L. Saminadyar et al., PRL **79**, 2526 (1997).

Mesoscopic system w/ cold atoms



Two-terminal setup realized by Esslinger's group at ETH

2.0

- Atomic current (charge neutral)
- Current induced by biases on thermodynamic quantities (chemical potential, temperature)







~500 µm





Why cold atoms?





Control of quantum statistics



S. Eckel et al., PRA 93, 063619 (2016).

(AEA) Transport in strongly-interacting Fermi gases A



(IEA) Transport in strongly-interacting Fermi gases



(AEA) Mesoscopic system + dissipation (atom loss)





Experimental results were interpreted with a phenomenological non-Hermitian Landauer-Büttiker analysis

(AEA) Mesoscopic system + dissipation (atom loss)



Current formula of lossy point contact

SU, arXiv:2206.09088 (PRA in press)

$$I = \int \frac{d\omega}{2\pi} \Big[\mathcal{T}(\omega) + \frac{\mathcal{L}(\omega)}{2} \Big] [n_L(\omega) - n_R(\omega)] \qquad \begin{array}{l} \mathcal{T}(\omega) : \text{transmittance} \\ \mathcal{L}(\omega) : \text{loss probability} \end{array}$$

- Obtained with an analysis based on Keldysh+ Lindblad formalism or three-terminal Landauer-Büttiker analysis
- Consistent with the non-Hermitian Landauer-Büttiker analysis

(AFA) Strongly-interacting Fermi gas + dissipation



0.00

0.01

0.03

 $\Delta \mu / \Delta$

0.02

0.04

0.05

0.06



Synthetic dimensions



Time, au (ms

 $\frac{2}{t_s/t_x}$

0.3

0.4

0.5

±∓

 $\frac{2}{t_s/t_x}$



(Mesoscopic transport via magnetic impurity



$$\mathcal{H} = \int d^3r \Big[\sum_{\sigma=\uparrow,\downarrow} \psi_{\sigma}^{\dagger} \Big\{ -\frac{\hbar^2 \nabla^2}{2m} + V(\mathbf{r}) \Big\} \psi_{\sigma} - g \psi_{\uparrow}^{\dagger} \psi_{\downarrow}^{\dagger} \psi_{\downarrow} \psi_{\uparrow} \Big] + \sum_{\sigma} V_{\sigma} \psi_{\sigma}^{\dagger} \psi_{\sigma}(\mathbf{0}) \Big]$$

$$\mathcal{H} = \int d^3r \Big[\sum_{\alpha=1,2} \psi_{\alpha}^{\dagger} \Big\{ -\frac{\hbar^2 \nabla^2}{2m} + V(\mathbf{r}) \Big\} \psi_{\alpha} - g \psi_{\uparrow}^{\dagger} \psi_{\downarrow}^{\dagger} \psi_{2} \psi_{1} \Big] + \sum_{\sigma} V_{\sigma} \psi_{\sigma}^{\dagger} \psi_{\sigma}(\mathbf{0}) \Big]$$

$$\mathcal{H} = \int d^3r \Big[\sum_{\alpha=1,2} \psi_{\alpha}^{\dagger} \Big\{ -\frac{\hbar^2 \nabla^2}{2m} + V(\mathbf{r}) \Big\} \psi_{\alpha} - g \psi_{\uparrow}^{\dagger} \psi_{2}^{\dagger} \psi_{2} \psi_{1} \Big] + \sum_{\alpha,\beta=1,2} \psi^{\dagger} V_{\alpha\beta} \psi_{\beta}(\mathbf{0}) \Big]$$

$$V_{12(21)} \neq 0 \quad \text{if } V_{\uparrow} \neq V_{\downarrow}$$



Landauer-Büttiker formula $I = \int \frac{d\epsilon}{h} \mathcal{T}(\epsilon) [f_1(\epsilon) - f_2(\epsilon)]$





K. Ono et al., Nat. Commun. 12, 6724 (2021).

Two-orbital lattice system with ¹⁷³Yb

¹S₀ atoms: itinerant fermions

³P₀ atom: localized impurity

Spin-dependent potential can be tuned with the orbital Feshbach resonance

G. Pagano et al., PRL **115**, 265301 (2015); M. Hofer et al., PRL **115**, 265302 (2015).







(AEA) Mesoscopic transport of Bose superfluid



Tunneling Hamiltonian formalism



$$H = H_L + H_R + H_T$$

$$I = -\dot{N}_L = i[N_L, H_T]$$

$$H_T = \sum_{\mathbf{k}, \mathbf{p}} \left(e^{-i\Delta\mu\tau} t_{\mathbf{k}, \mathbf{p}} b_{\mathbf{k}, L}^{\dagger} b_{\mathbf{p}, R} + h.c. \right)$$
absence of the momentum conservation

There must be the conversion process between condensation and normal elements.

Linear response theory: F. Meier & W. Zwerger PRA **64** 033610 (2001). Beyond linear response effect: SU & J.P. Brantut, PRR **2**, 023284 (2020); SU, PRR **2**, 023340 (2020).

Experiment: G. Del Pace et al., PRL **126**, 055301 (2021).







Editors' Suggestion

Asymmetry and nonlinearity of current-bias characteristics in superfluid–normal-state junctions of weakly interacting Bose gases

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Normal Superfluid bosons bosons

Normal bosons: Hartree-Fock theory Superfluid bosons: Bogoliubov theory



- Asymmetry arises from the conversion process between condensation and normal elements, and the bosonic Andreev reflection.
- Quasiparticle current shows a symmetric response and is suppressed with decreasing T.

Strongly interacting Fermi gas in a cavity

V. Helson et al., PRR 4, 133199 (2022)

$$H = H_{\text{atom}} + H_c + H_{\text{int}}$$
$$H_c = \Delta a^{\dagger} a$$
$$H_{\text{int}} = \Omega a^{\dagger} a \int d^3 r n(\mathbf{r}) \cos^2 \mathbf{k}_c \cdot \mathbf{r}$$

$$\mathbf{a}$$

$$\mathbf{b}$$

• Photon measurement reflects density-density correlation of atoms $\int S(q, w) + S(-q, w)$

$$\chi^{R}(\mathbf{q},\omega=0) = -\int d\omega \left[\frac{S(\mathbf{q},\omega) + S(-\mathbf{q},\omega)}{\omega}\right]$$

compressibility sum rule

 Agreement with a theory with the operator product expansion







Two-terminal transport of Fermi gases

Nonlinear current-bias characteristics D. Husmann et al., Science 350, 1498 (2015).

Breakdown of conductance quantization

SU and M. Ueda, PRL **118**, 105303 (2017).

Particle loss effect in mesoscopic transport

SU, arXiv:2206.09088 M.-Z. Huang et al., arXiv:2210.03371

• Transport with synthetic junctions

Realization of two and three terminal transport

S. Nakada et al., PRA 102, 031302(R) (2020). K. Ono et al., Nat. Commun. **12,** 6724 (2021).

Transport of bosons

Asymmetry and nonlinearity in SN junction SU, PRA **106**, L011303 (2022). Compressibility sum rule via optical cavity

V. Helson et al., PRR **4**, 133199 (2022).

Collaborators

Transport of Fermi gases









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Transport with synthetic dimensions



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Strongly-interacting Fermi gas inside a cavity



Brantut group @ EPFL J.P. Brnatut, V. Helson, T. Zwettler



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