



Reconstructing Neutron Star Masses from Metal-poor Stars

Erika M. Holmbeck

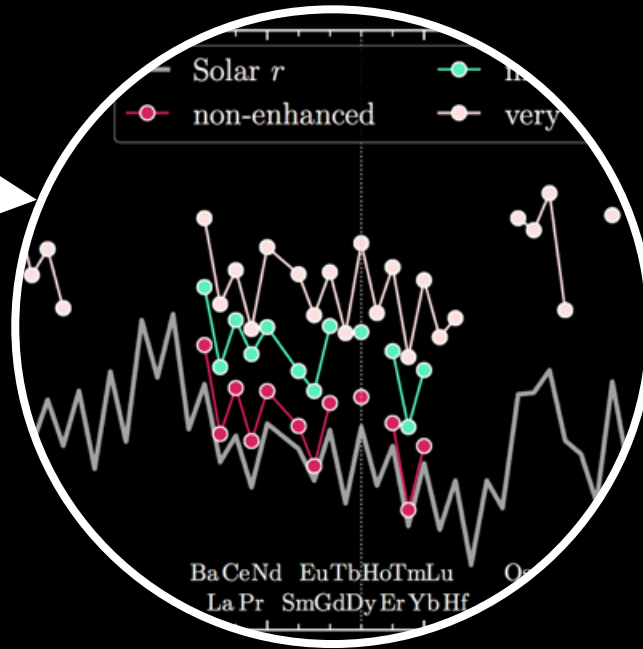
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Holmbeck, Frebel, McLaughlin, et al. 2021, ApJ, 909, 21

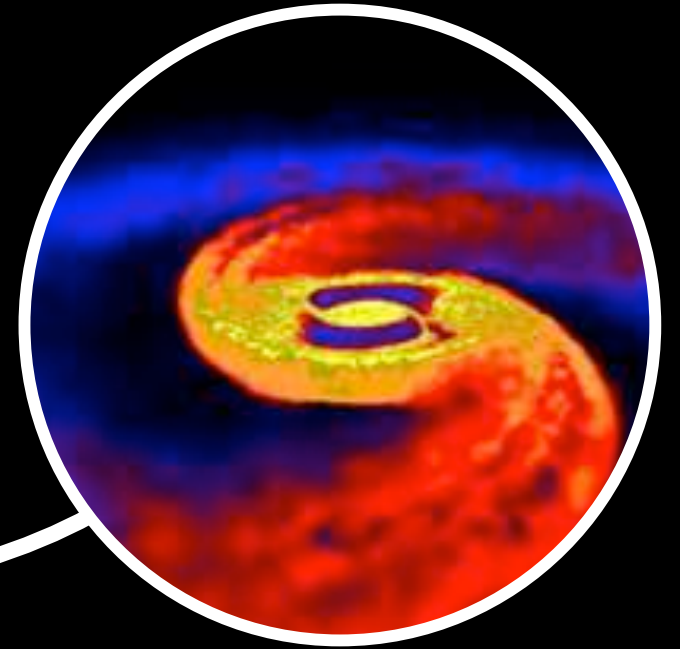
Star-forming gas



Enriches the inter-
stellar medium

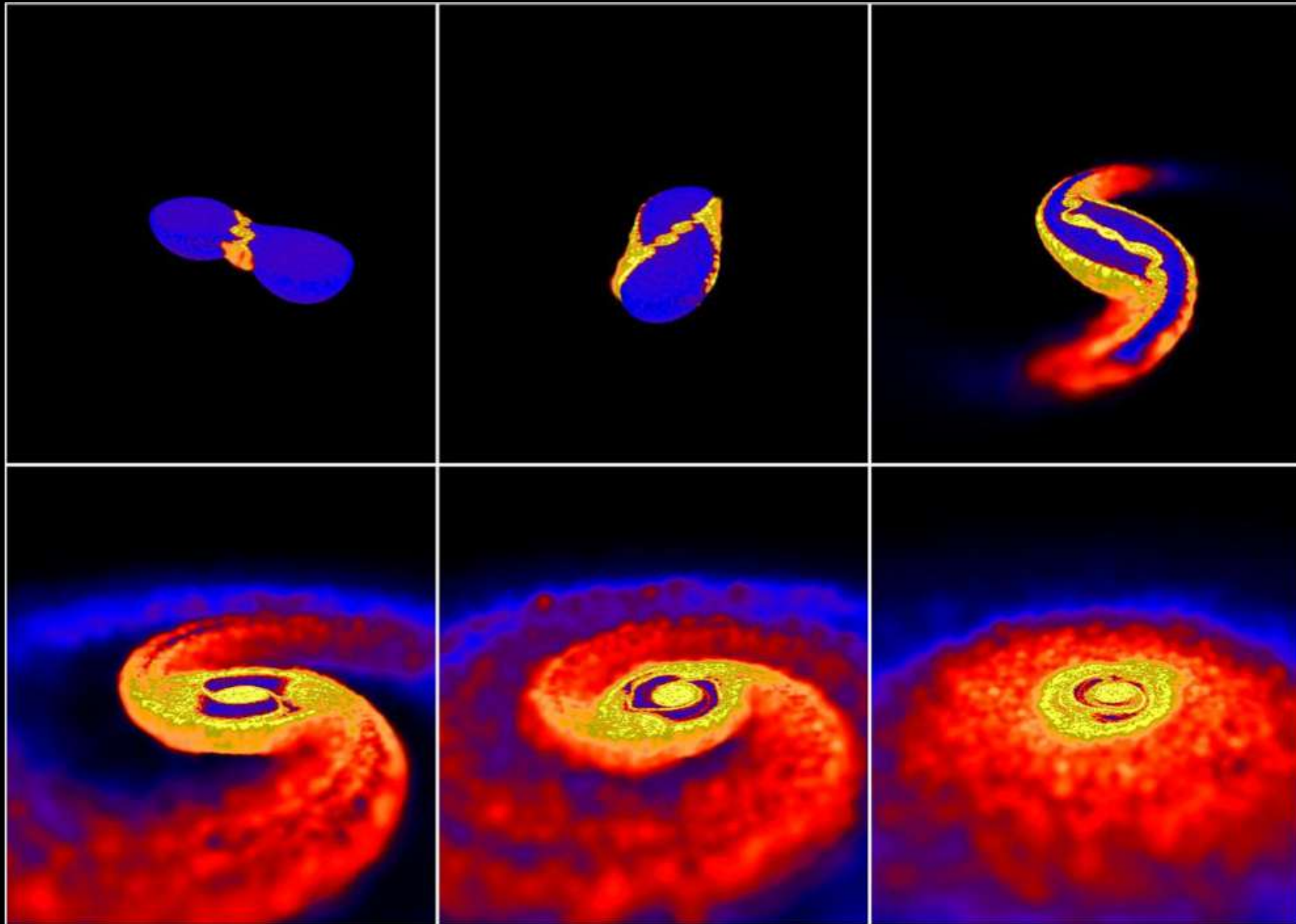


Stars with *r*-process
signatures

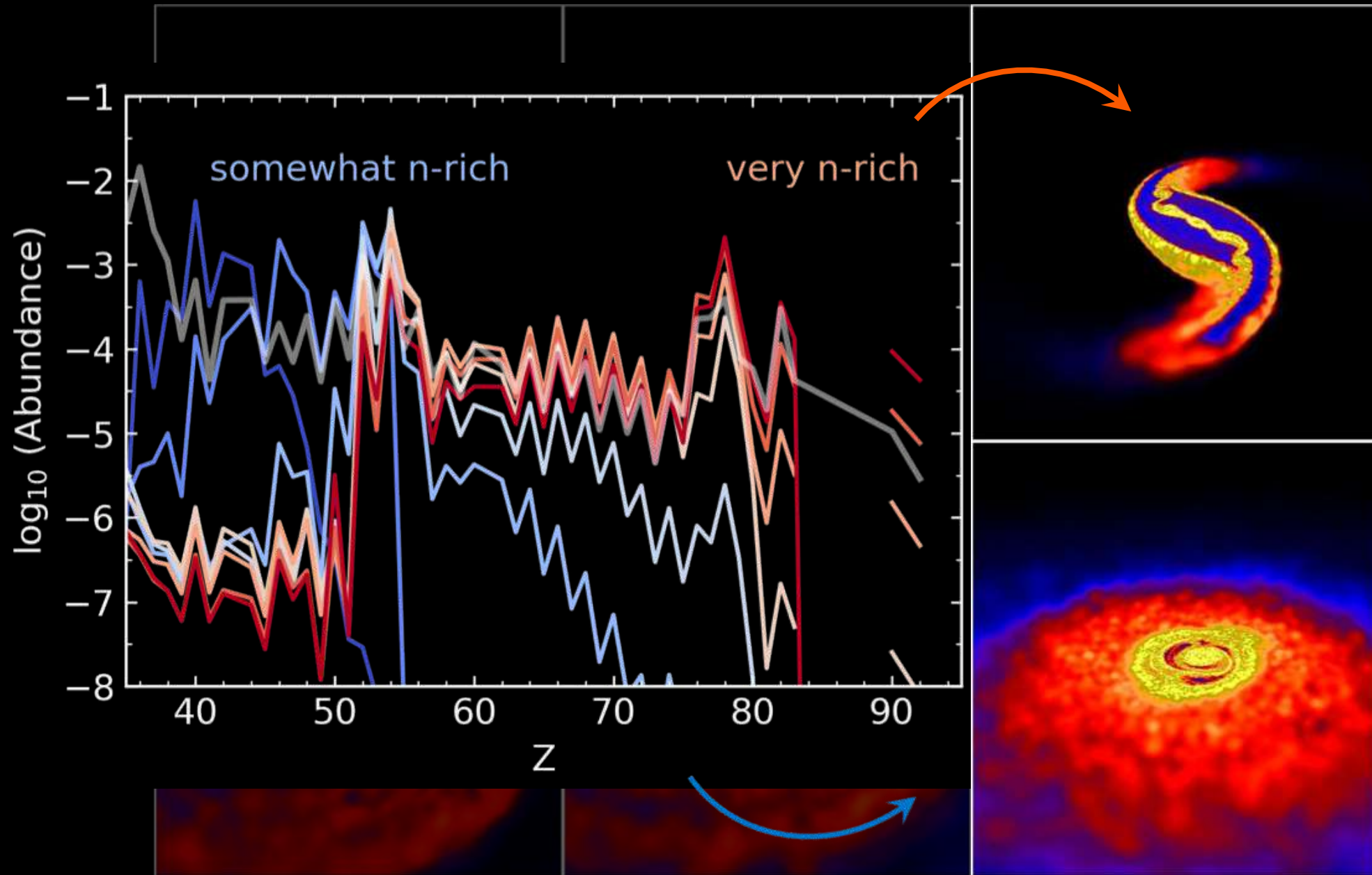


Merger ejects *r*-process material

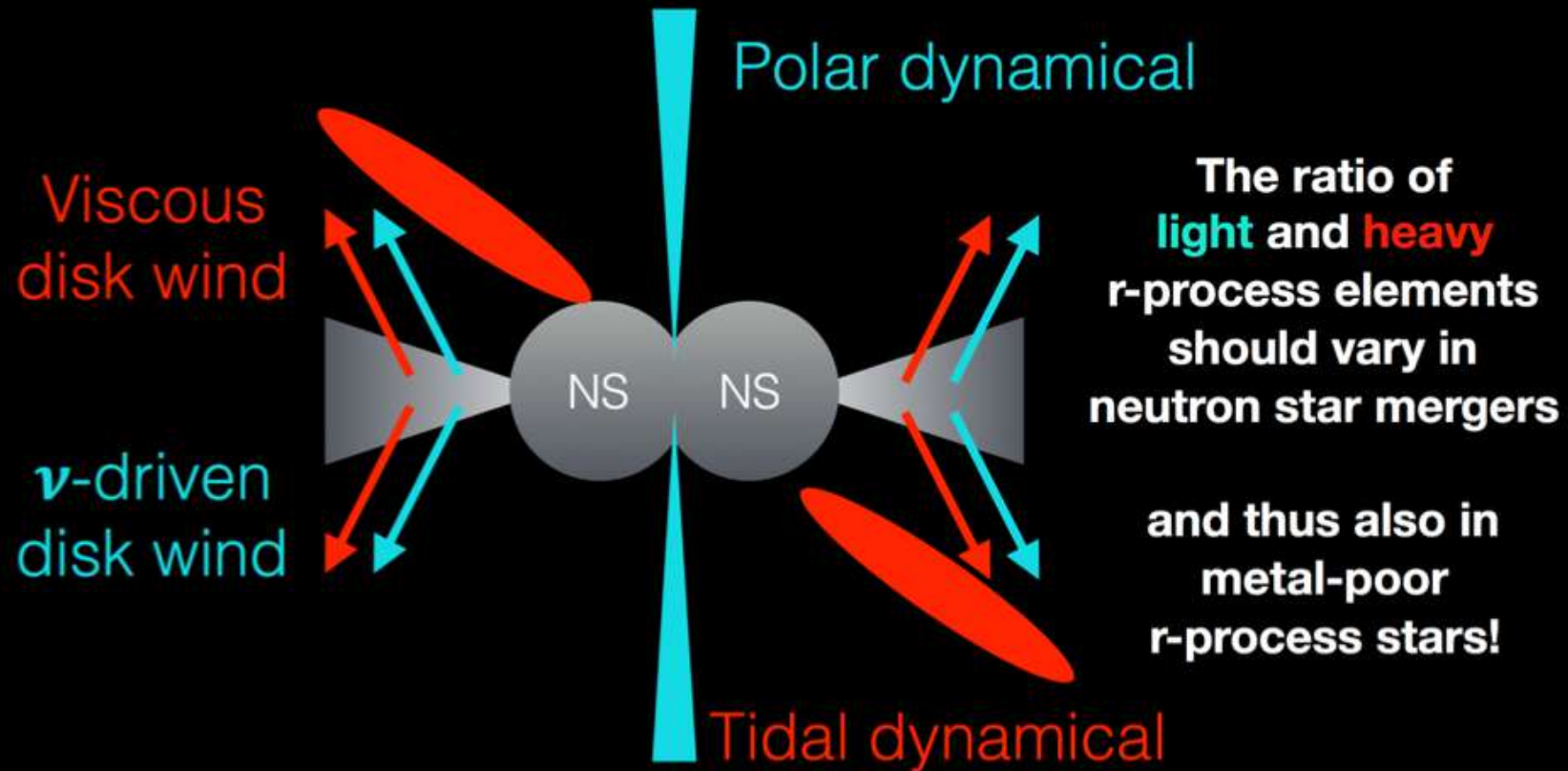
NSMs eject material through primarily two mechanisms



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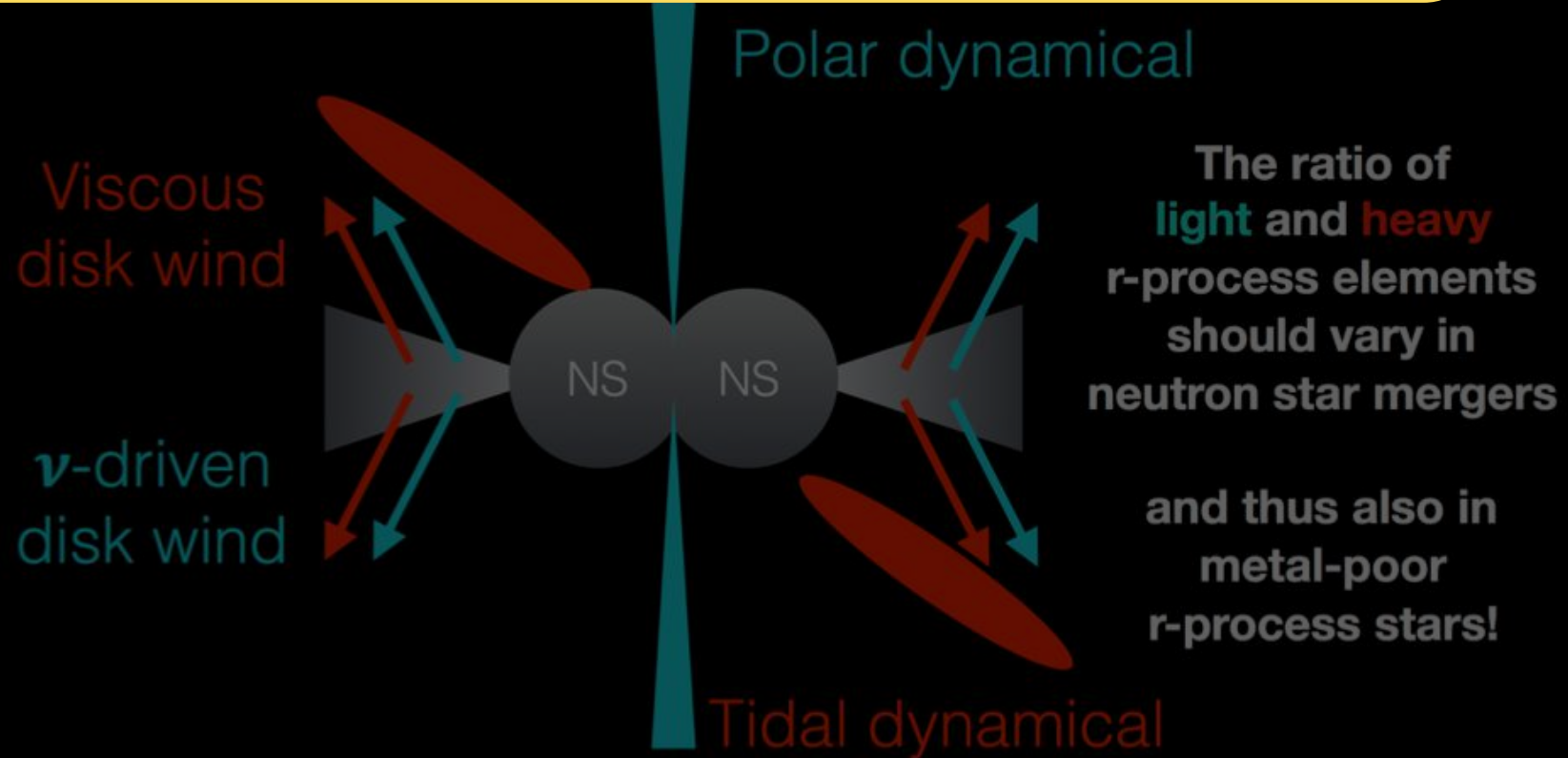
NSMs should not make identical r-process yields



Blue: high Y_e , neutron-poor, **light** 1st-peak r-process

Red: low Y_e , neutron-rich, **heavy** 2nd+3rd peak r-process

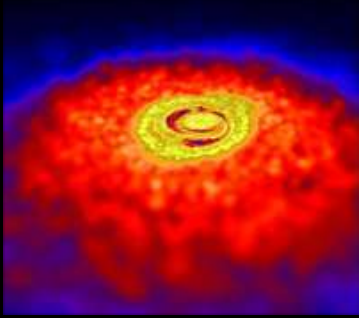
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Blue: high Y_e , neutron-poor, **light** 1st-peak r-process

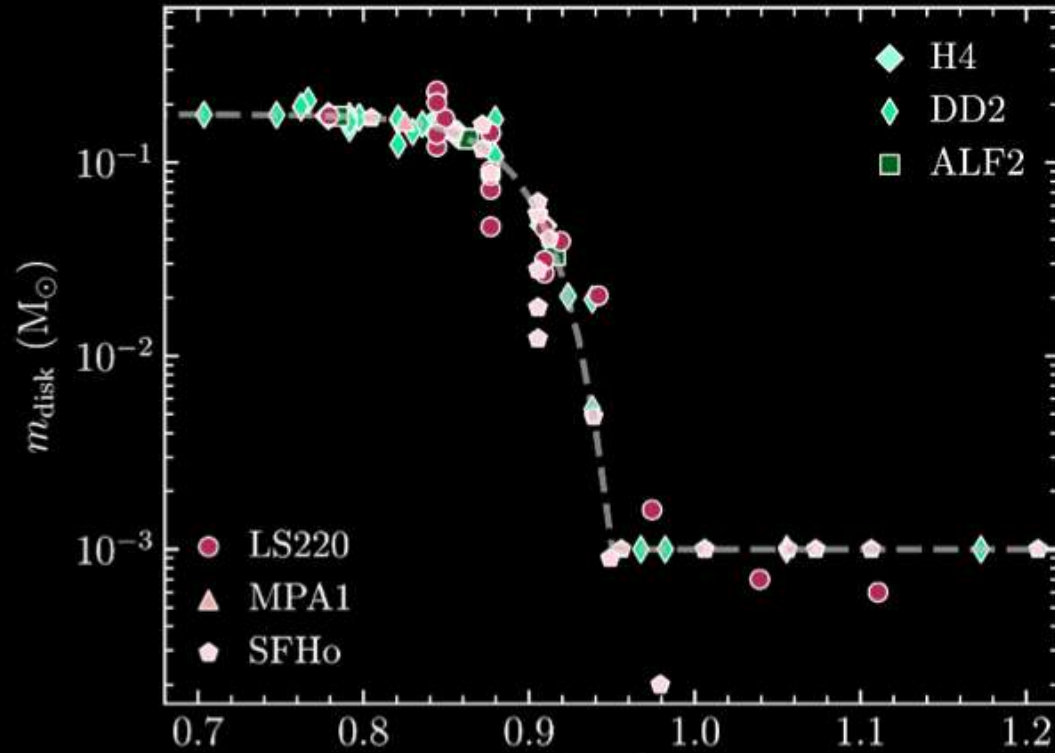
Red: low Y_e , neutron-rich, **heavy** 2nd+3rd peak r-process

The mass ejected by the different outflow mechanisms depend on the BNS properties

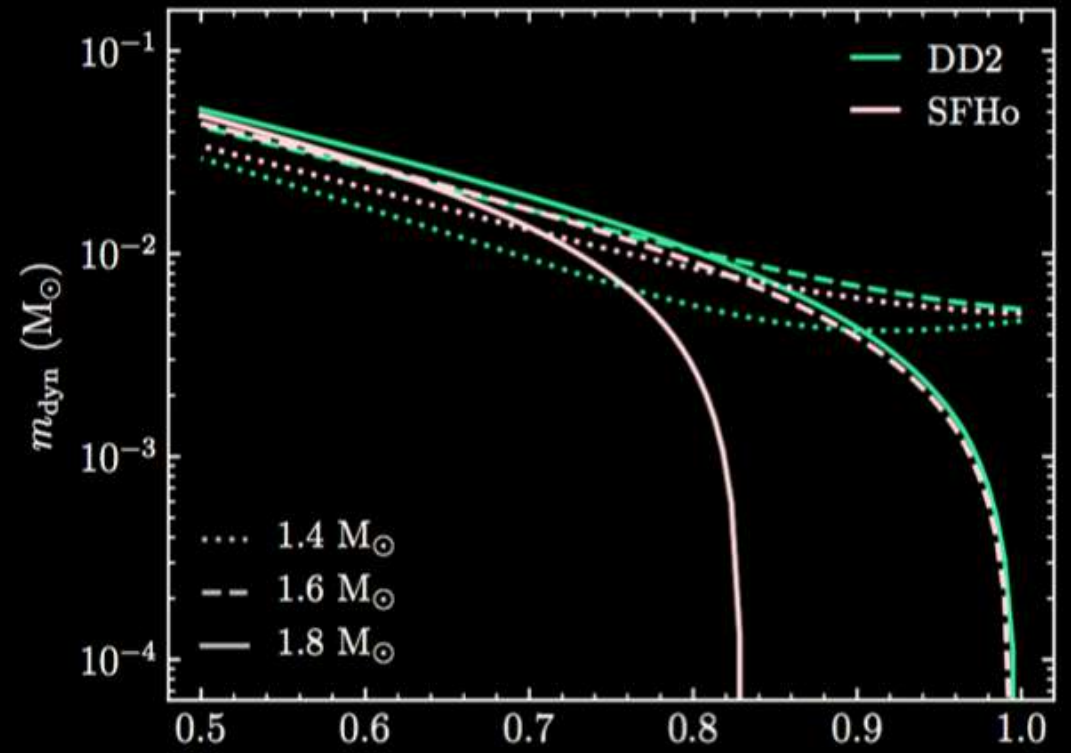


$$\log_{10}(m_{\text{disk}}) = \max \left\{ -3, a \left(1 + b \tanh \left[\frac{c - \frac{M_{\text{tot}}}{M_{\text{thr}}}}{d} \right] \right) \right\},$$

$$\frac{m_{\text{dyn}}}{10^{-3} M_{\odot}} = \left[\frac{a}{c_1} + b \left(\frac{M_2}{M_1} \right)^r + c c_1 \right] M_1 + [1 \leftrightarrow 2]$$

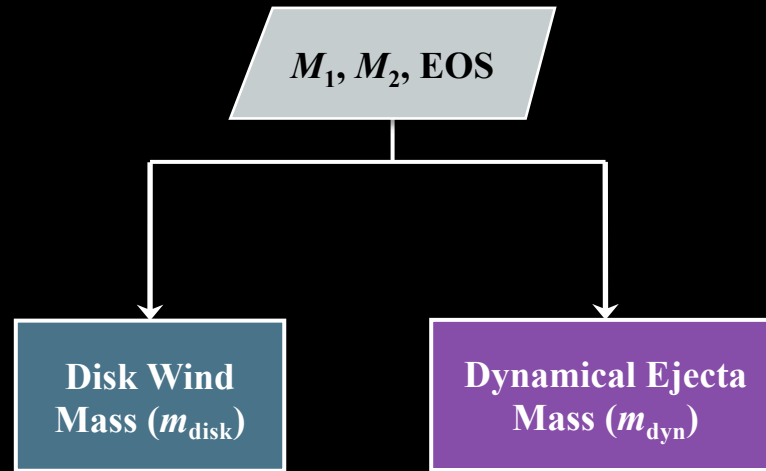


more total mass \longrightarrow

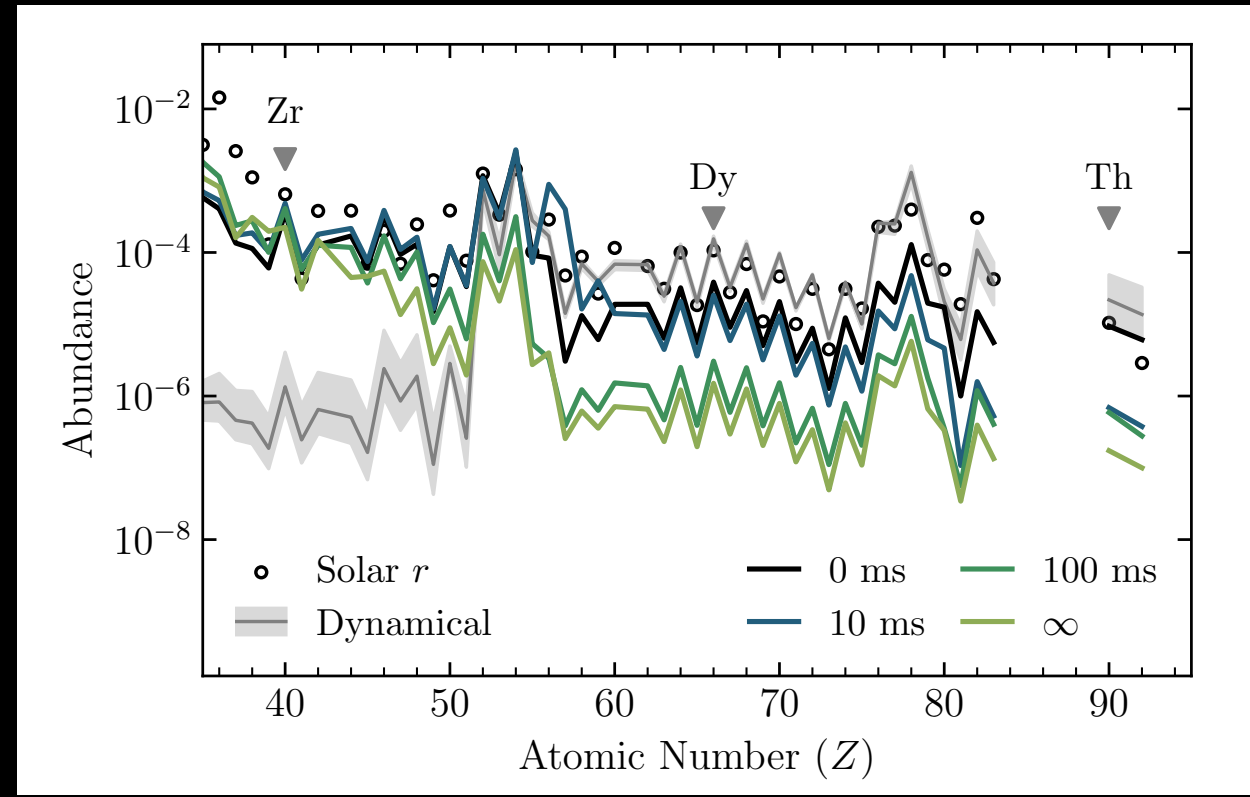
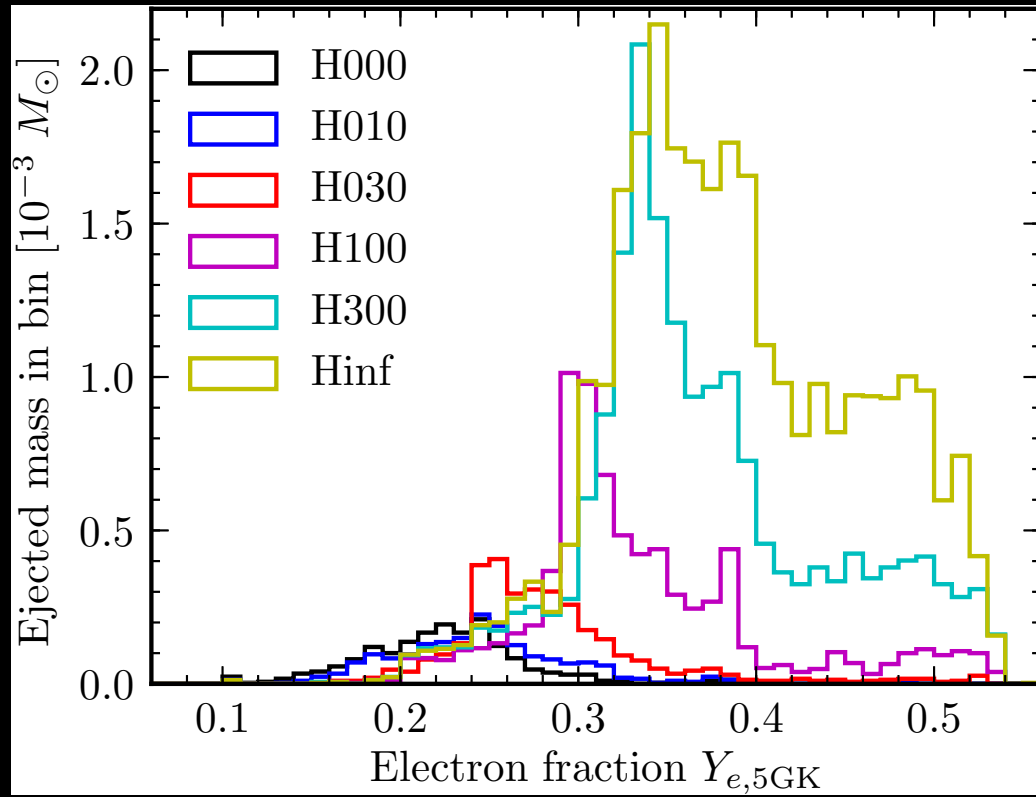


\longleftarrow more asymmetric

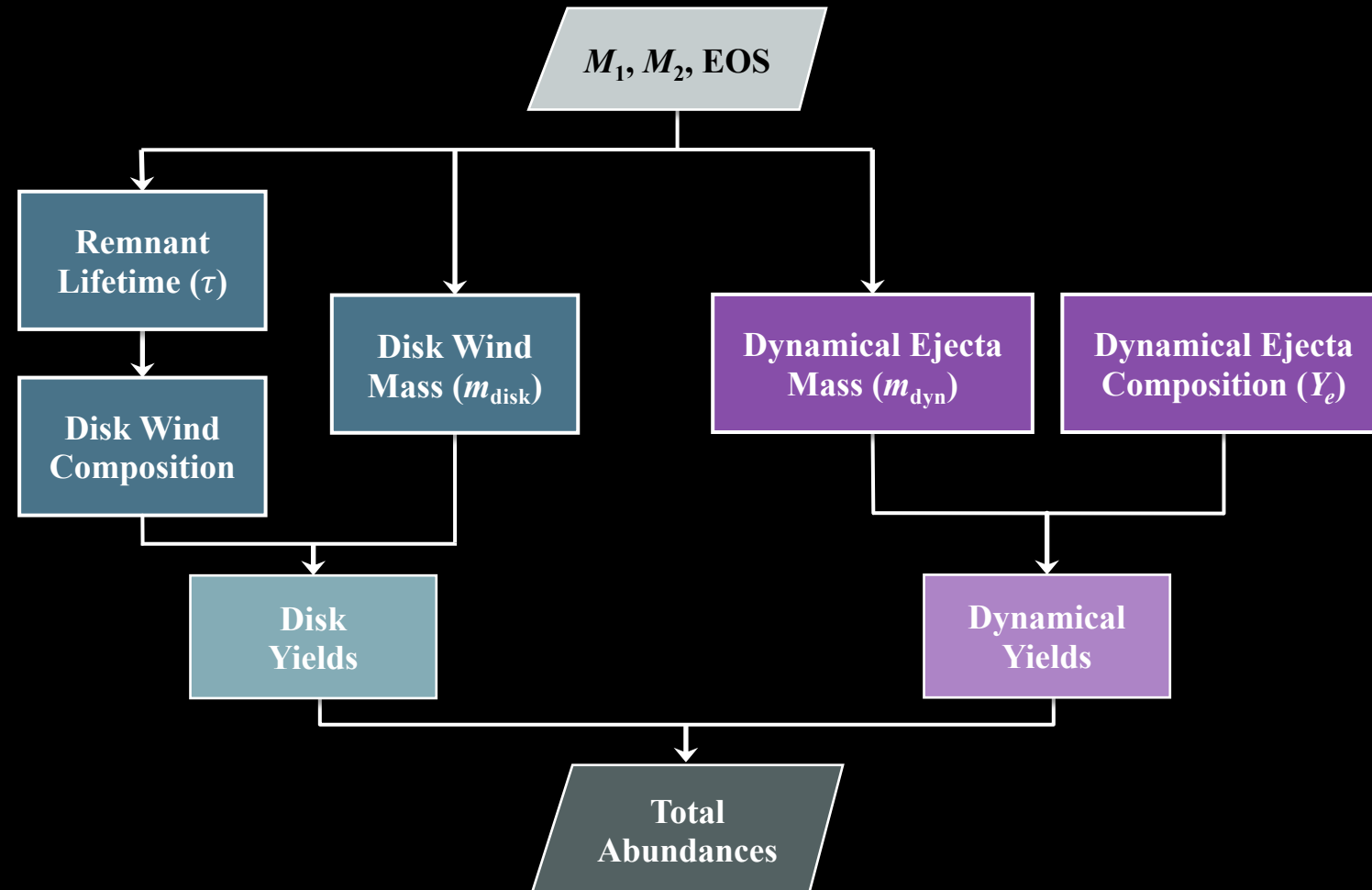
From a selection of M_1 , M_2 , and EOS, we can find the total mass ejected by the two components

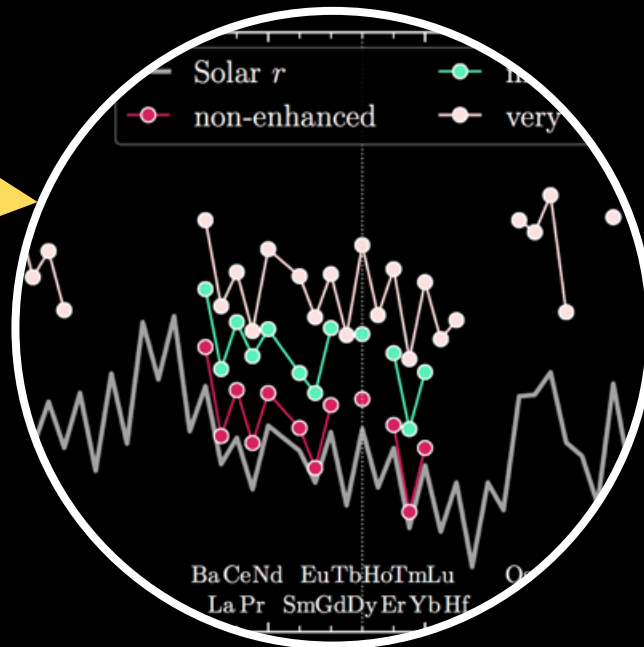
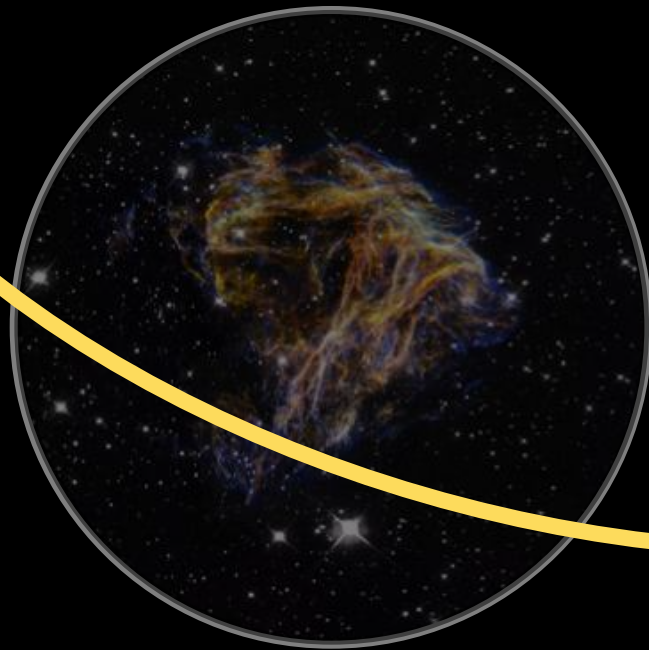


The elemental composition of the ejecta also depend on the binary properties

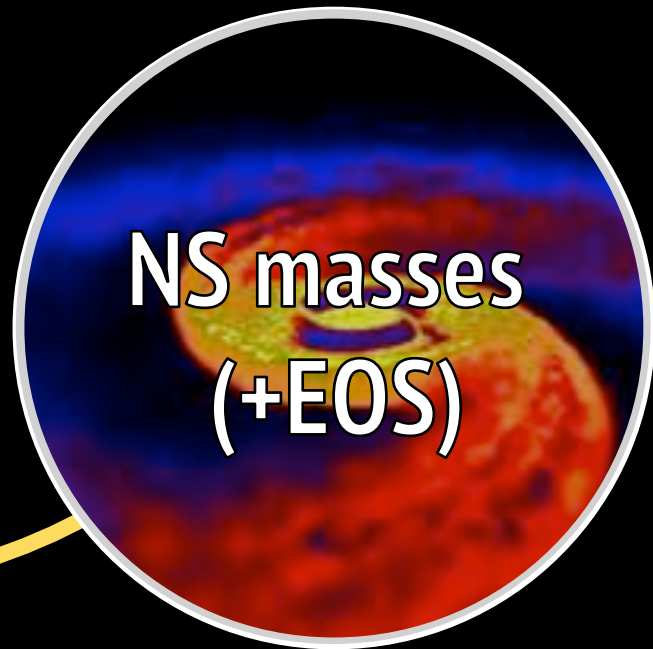


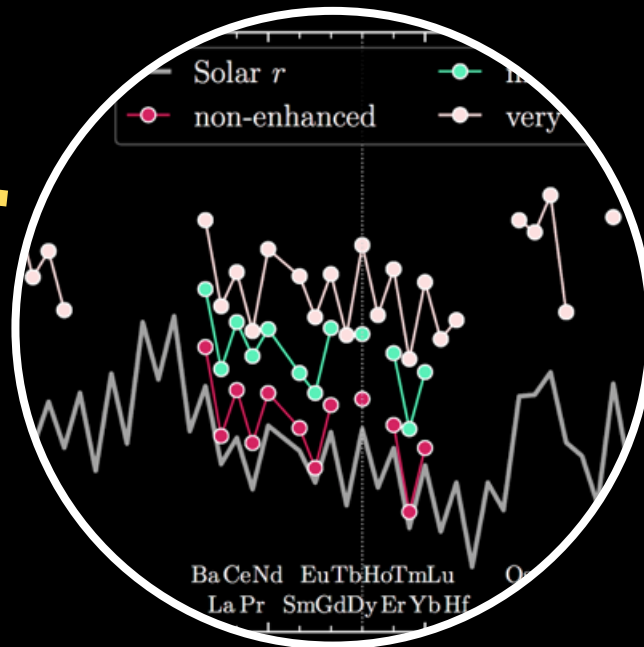
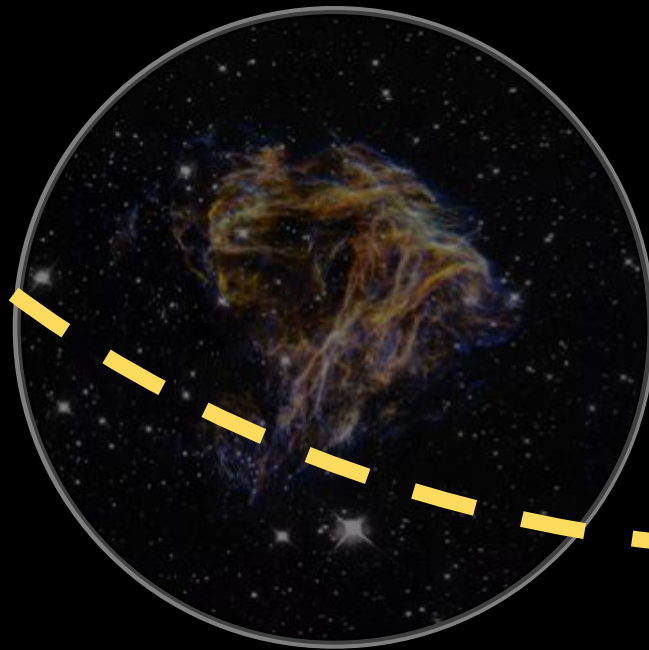
From a selection of M_1 , M_2 , and EOS, we can find the total r -process yields from an NSM event



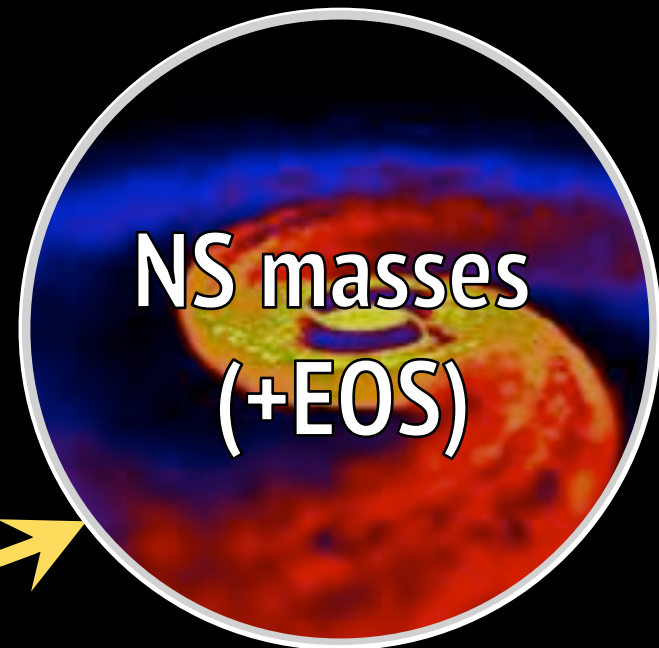


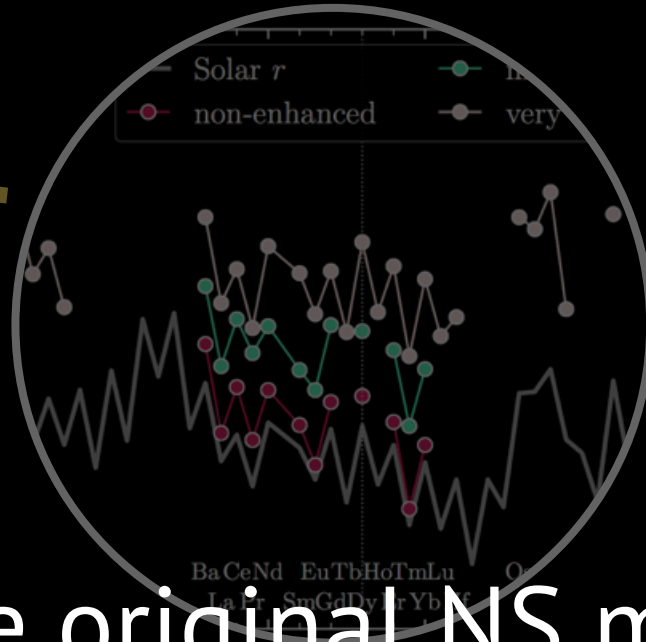
Stars with r -process signatures





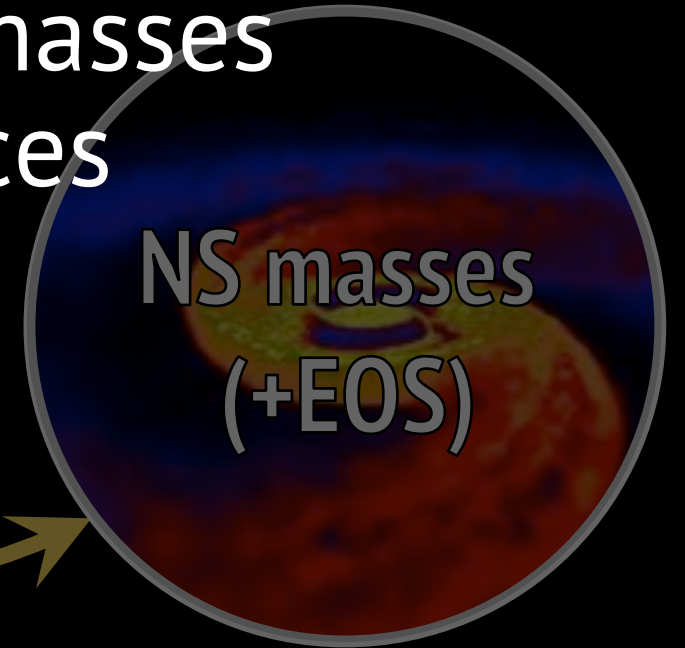
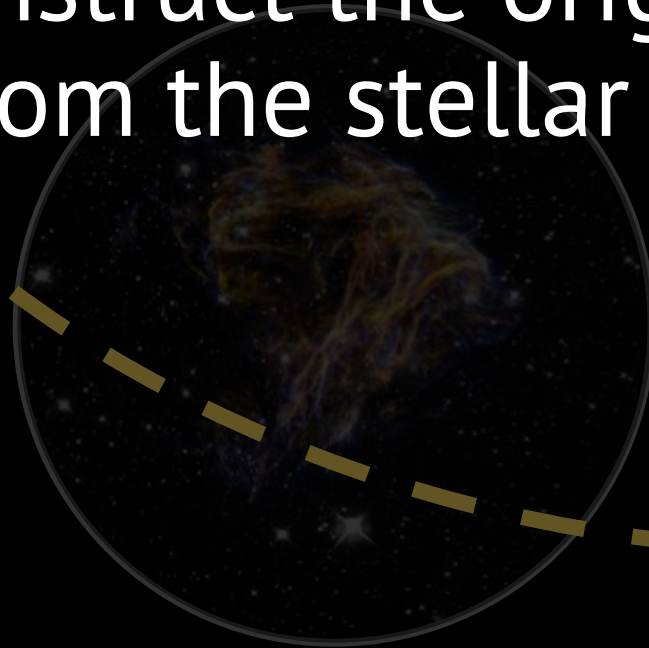
Stars with r -process signatures



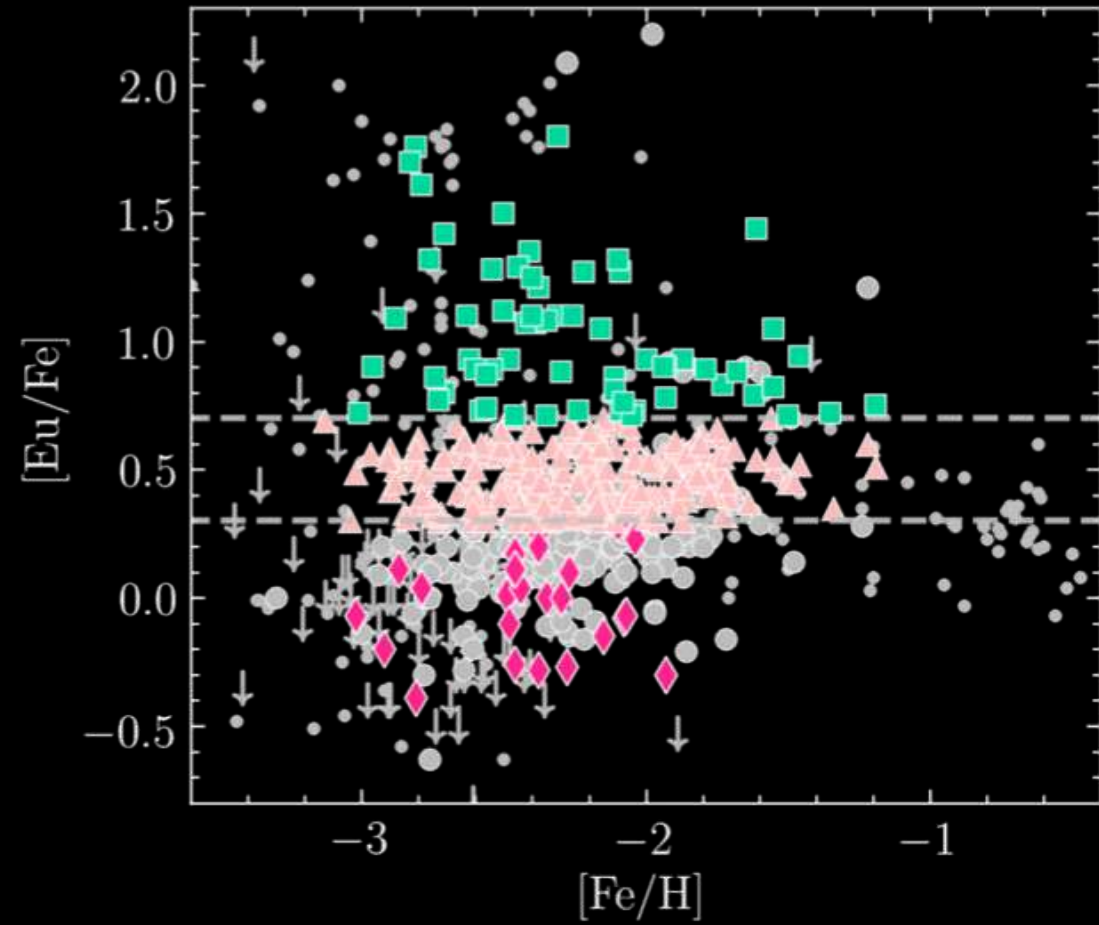
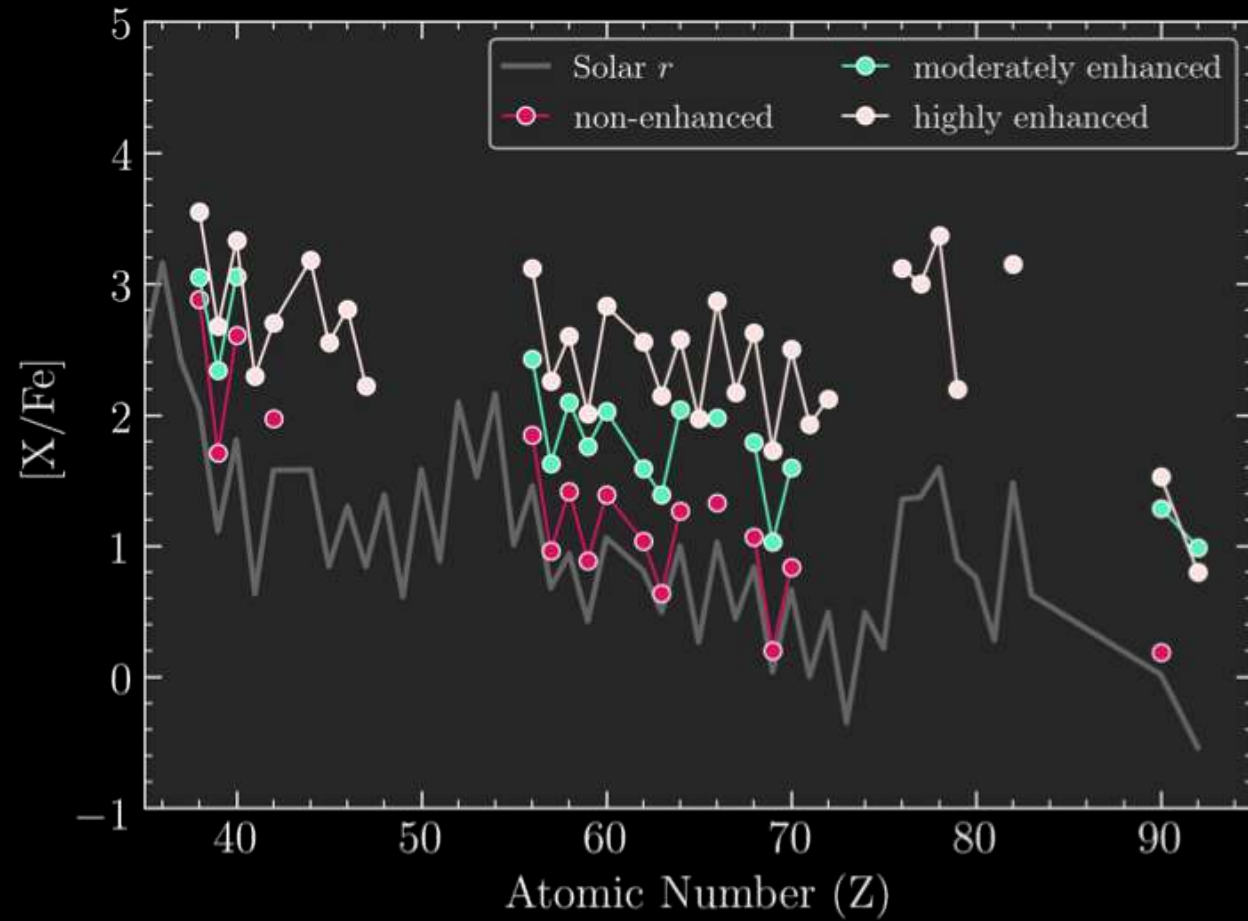


Stars with r -process signatures

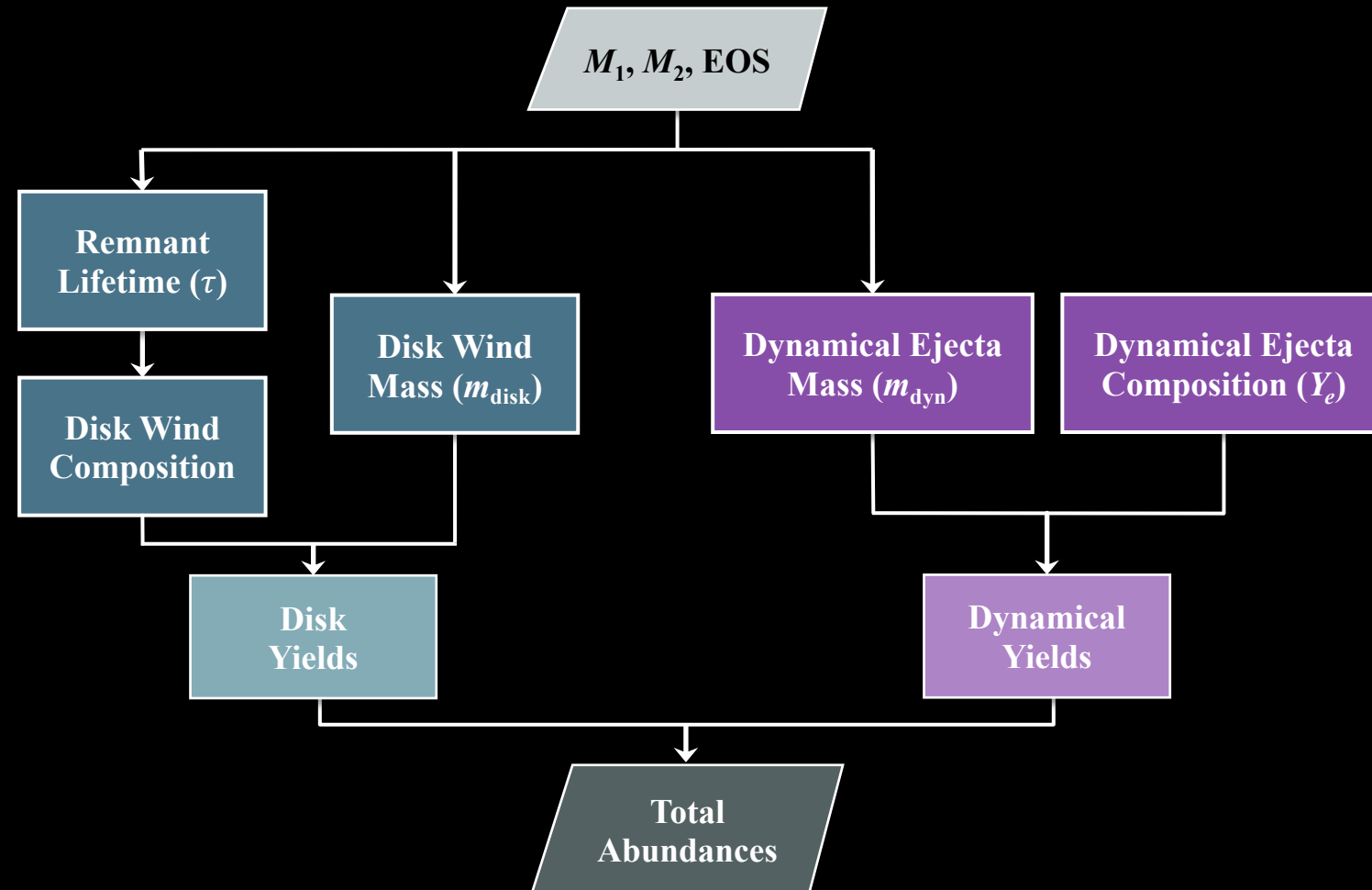
Reconstruct the original NS masses from the stellar abundances



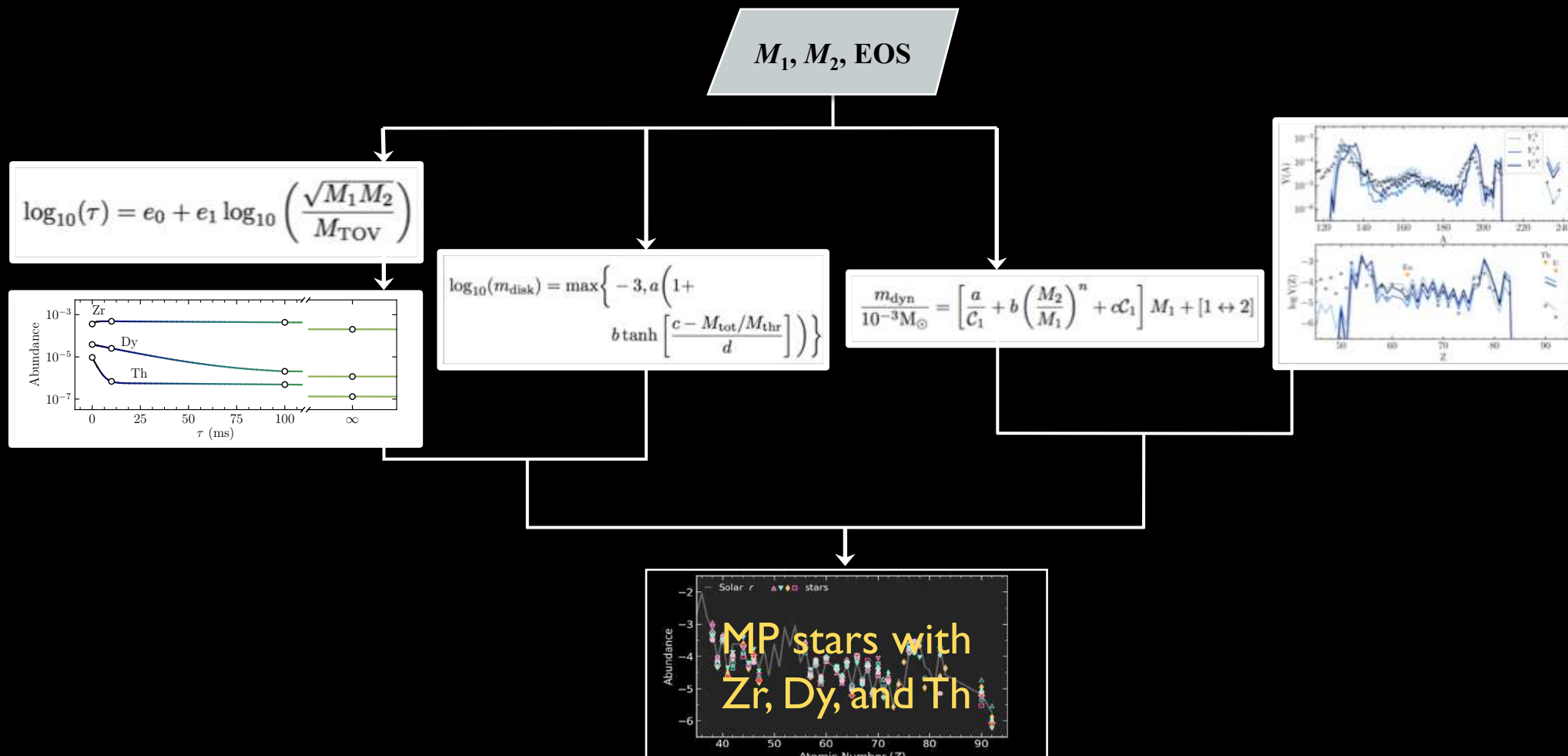
We need to consider the “non-enhanced” stars too



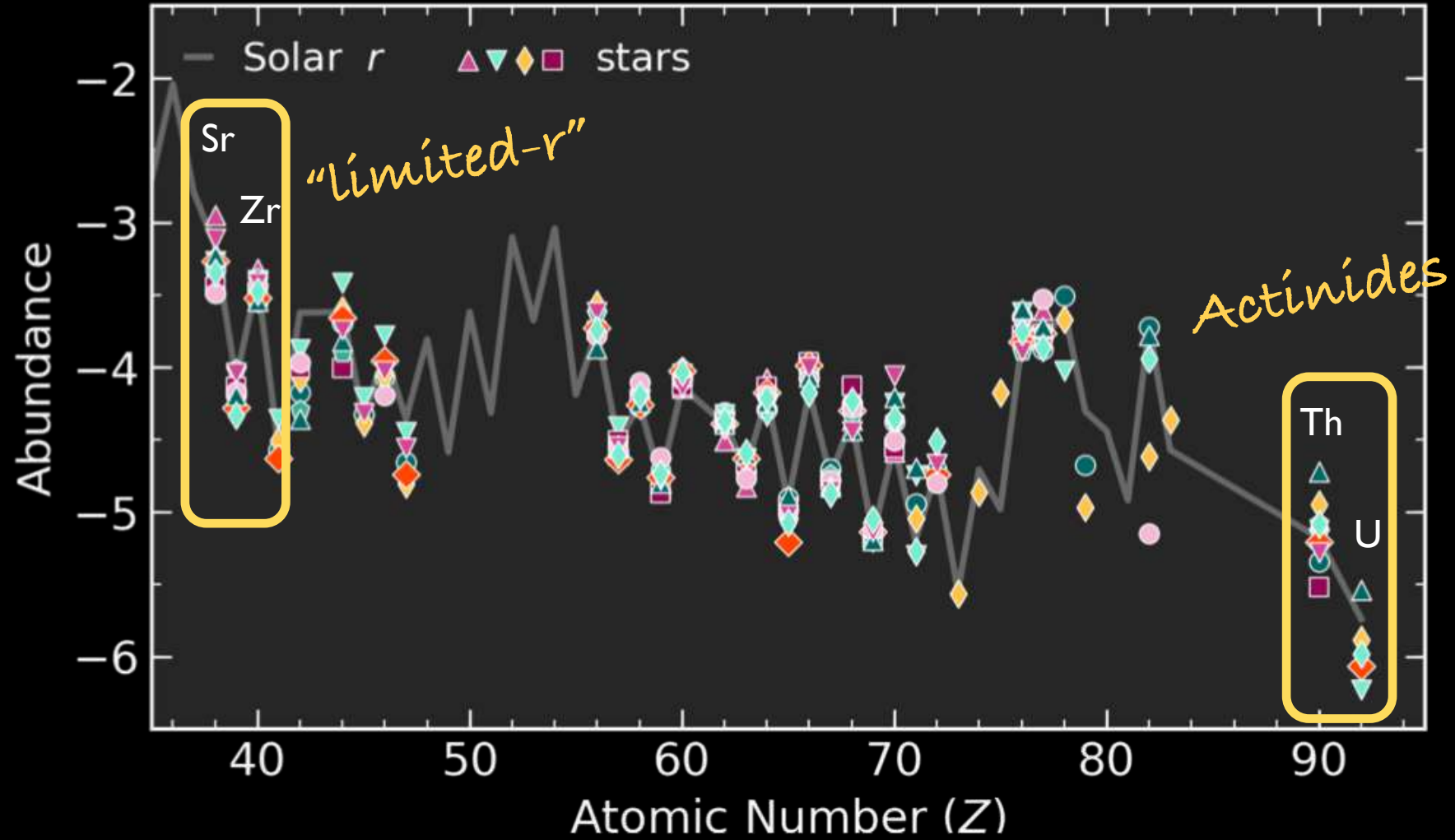
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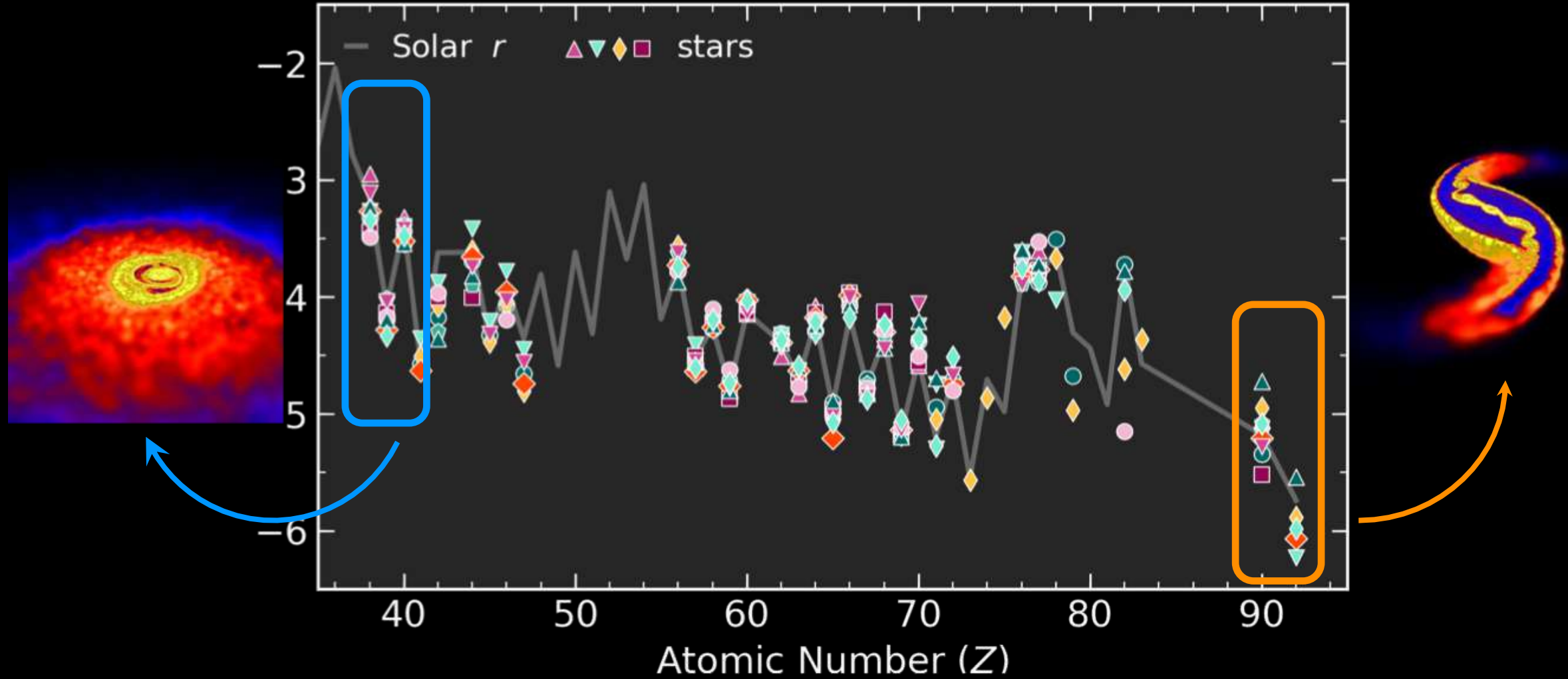
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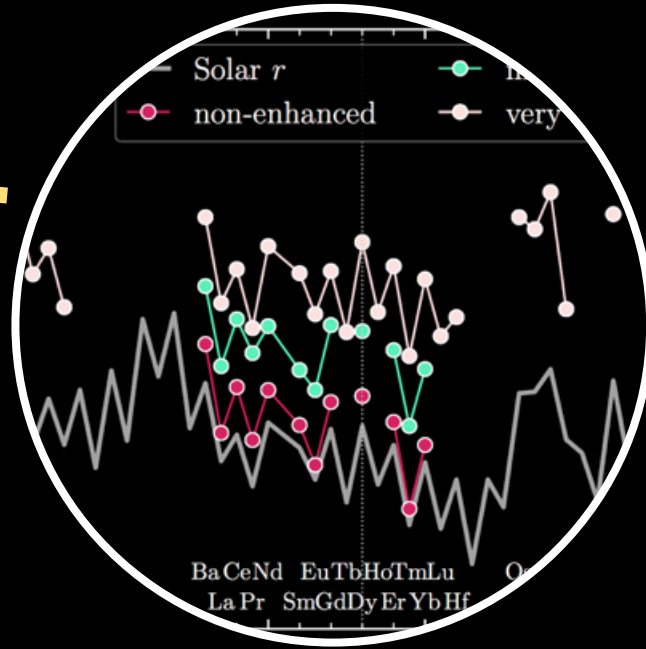


Elemental patterns of the r -stars differ in the lightest and heaviest elements



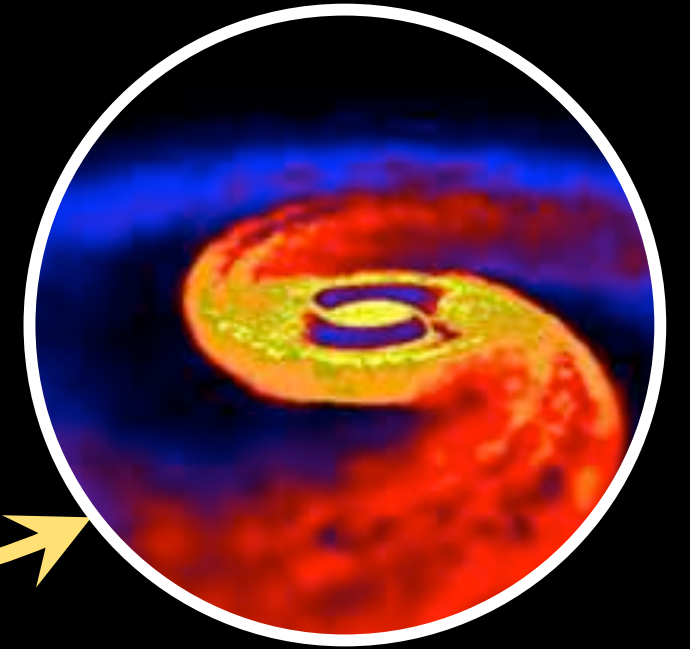
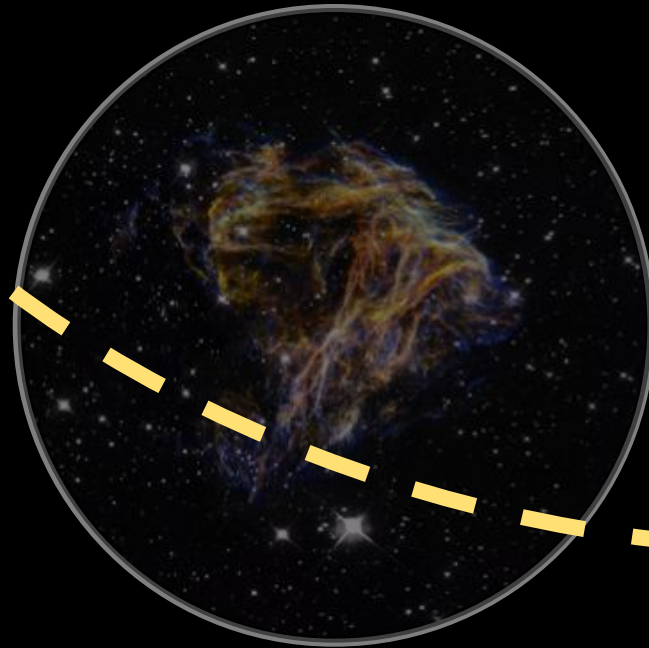
These are the regions where differences
in NS binaries manifest





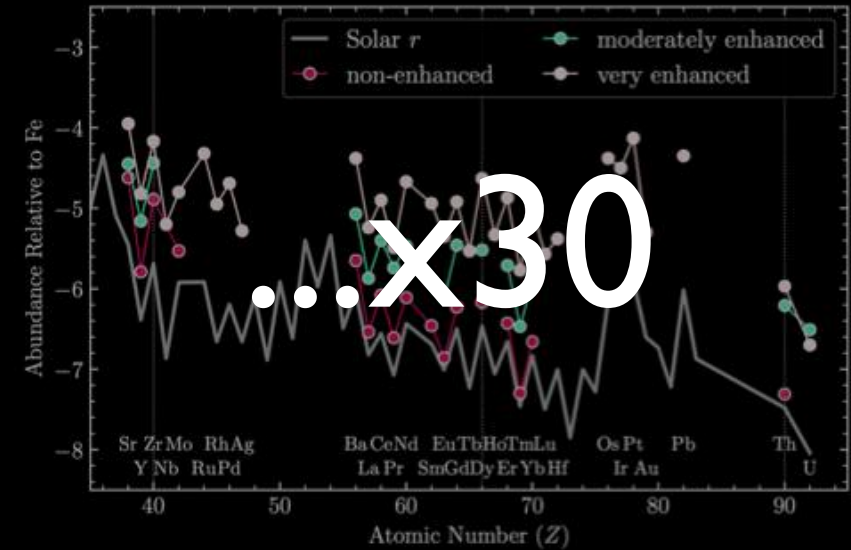
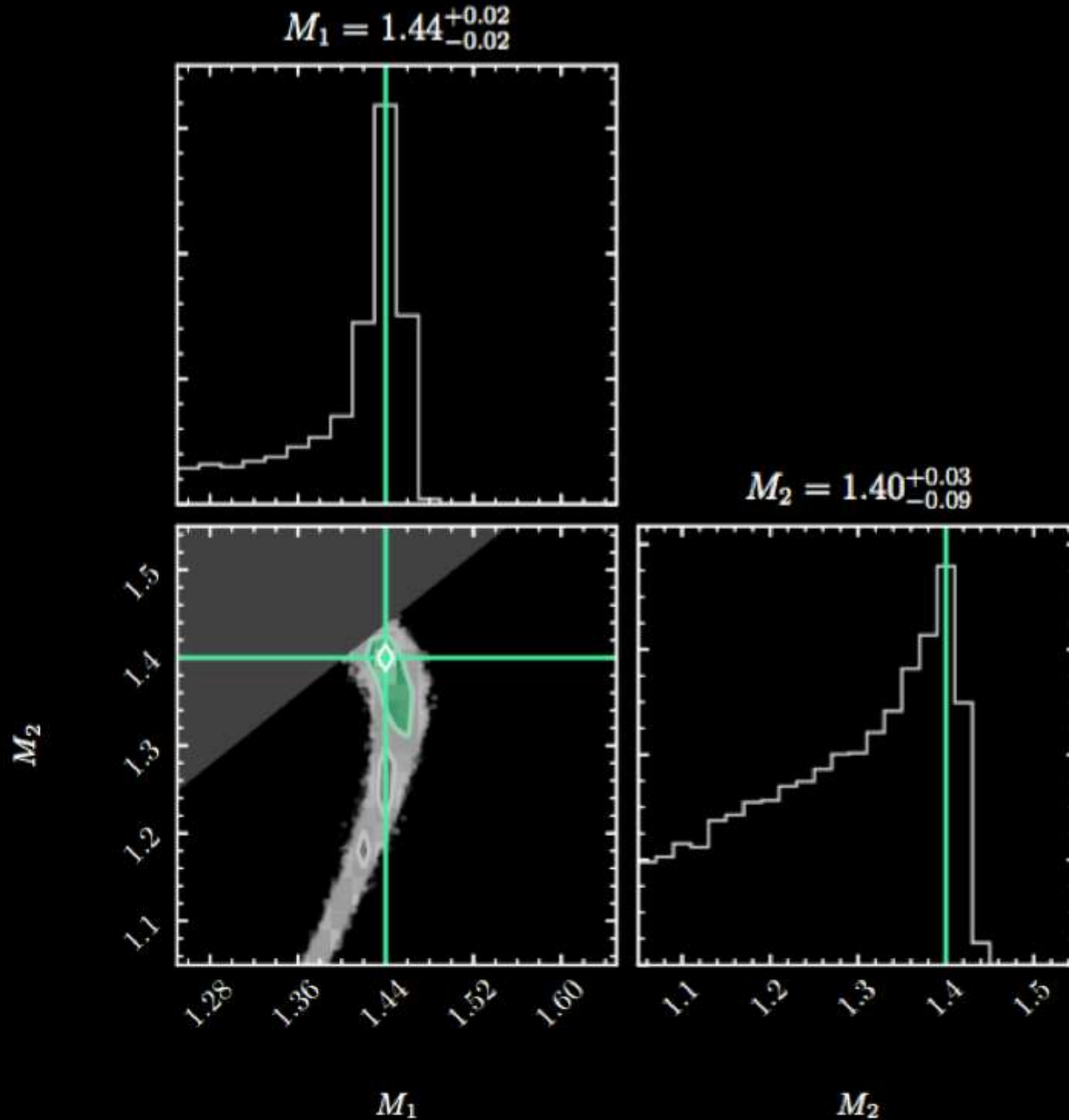
Stars with *r*-process signatures

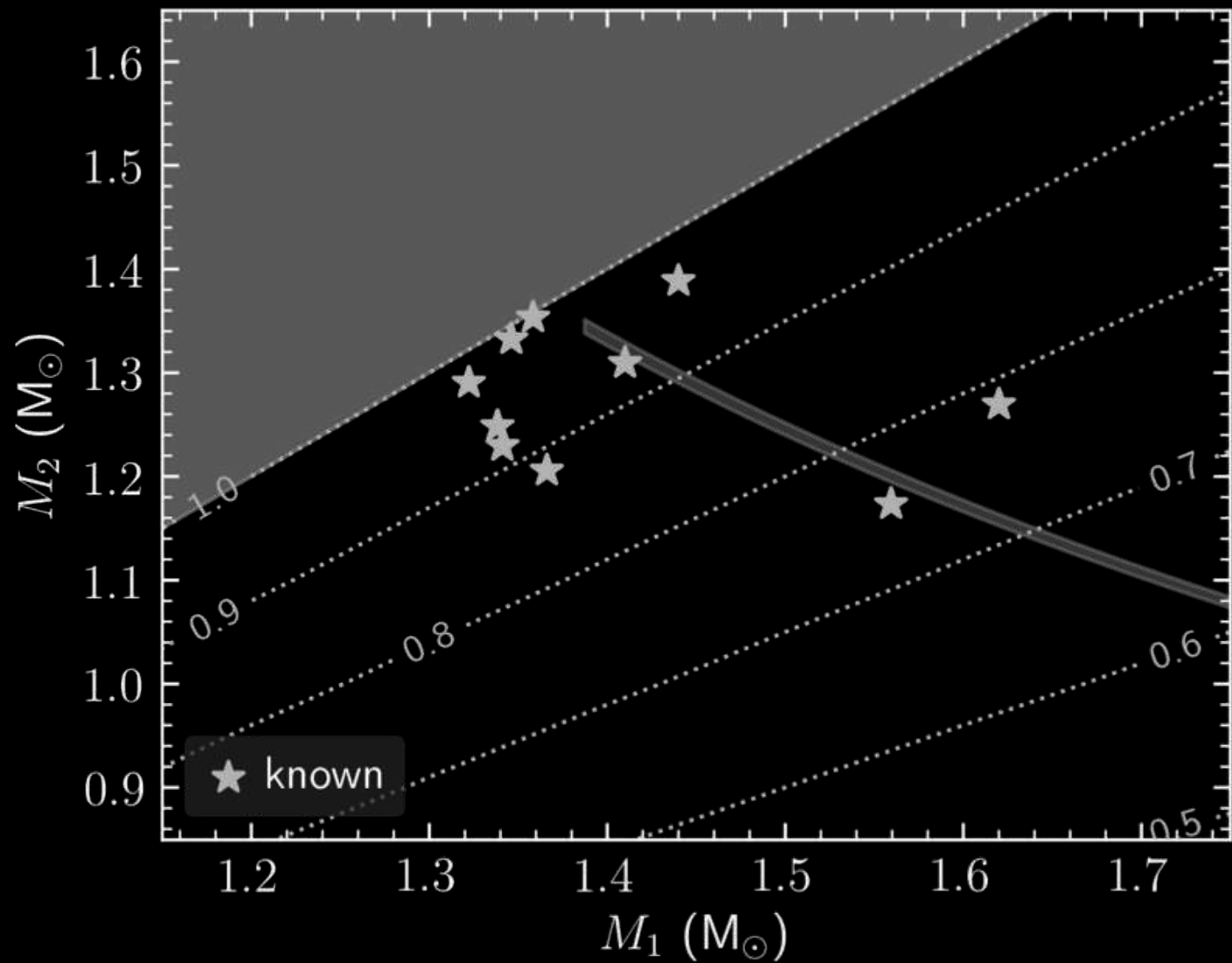
MCMC

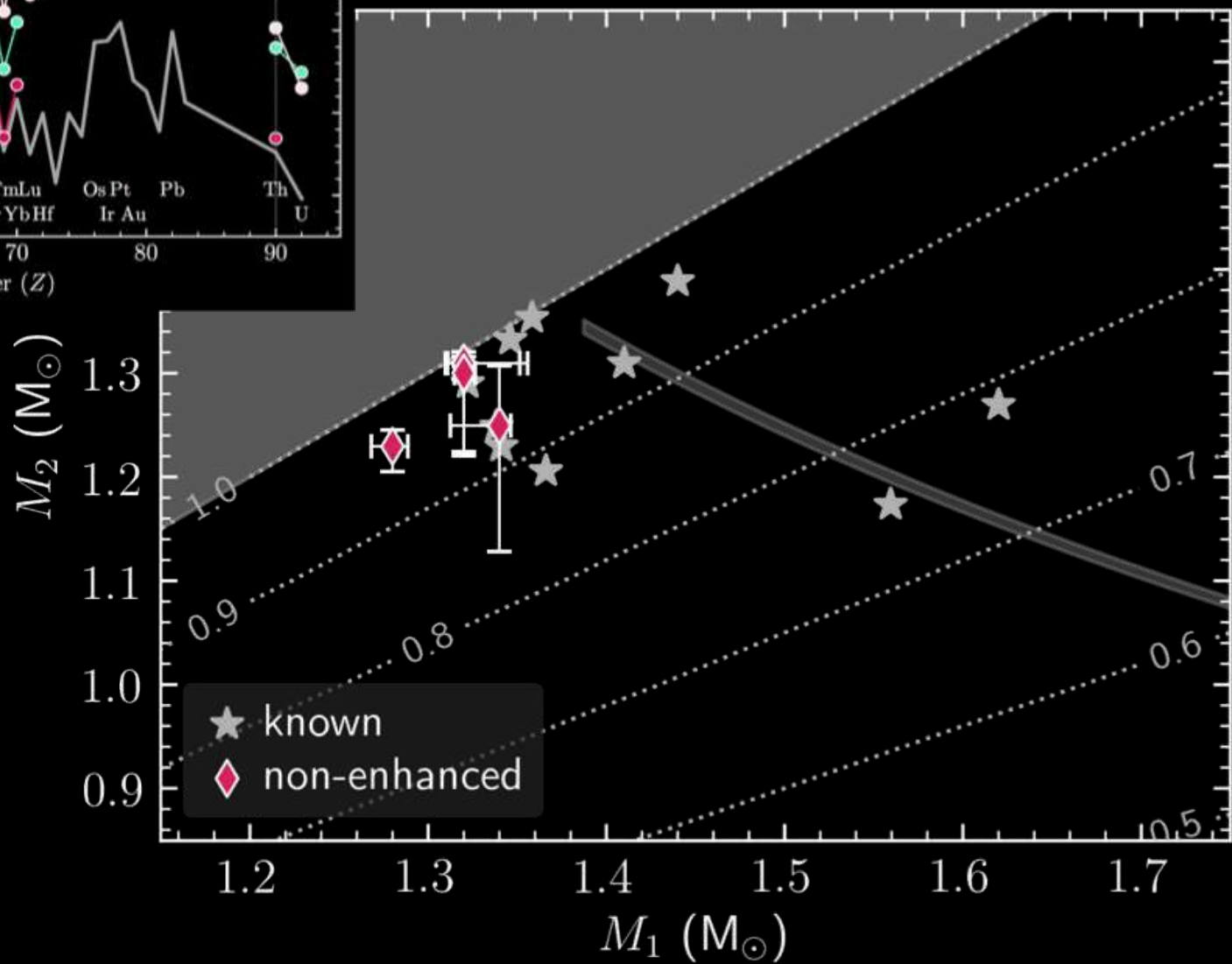
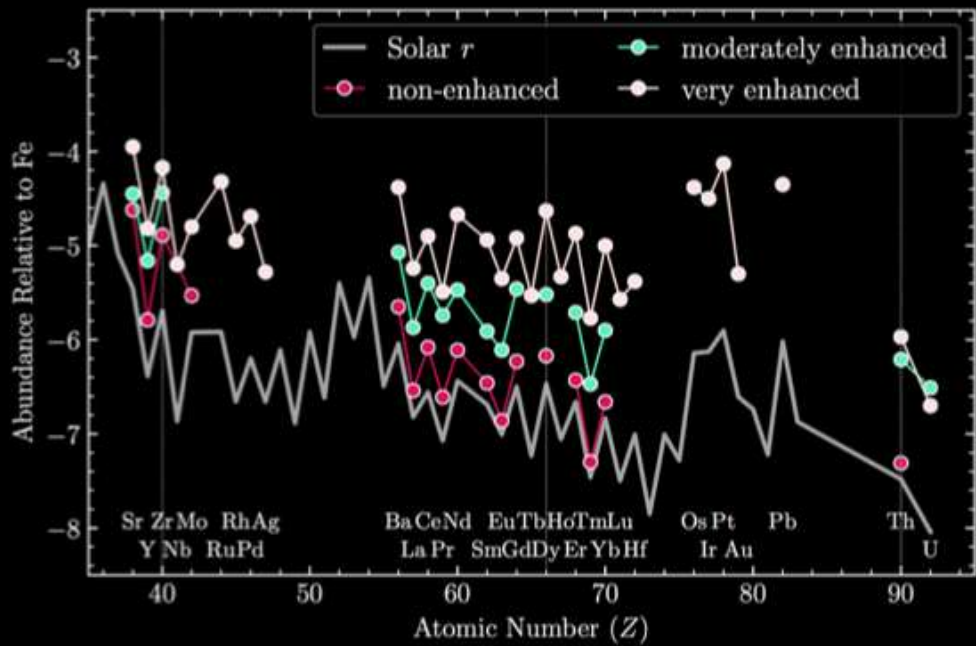


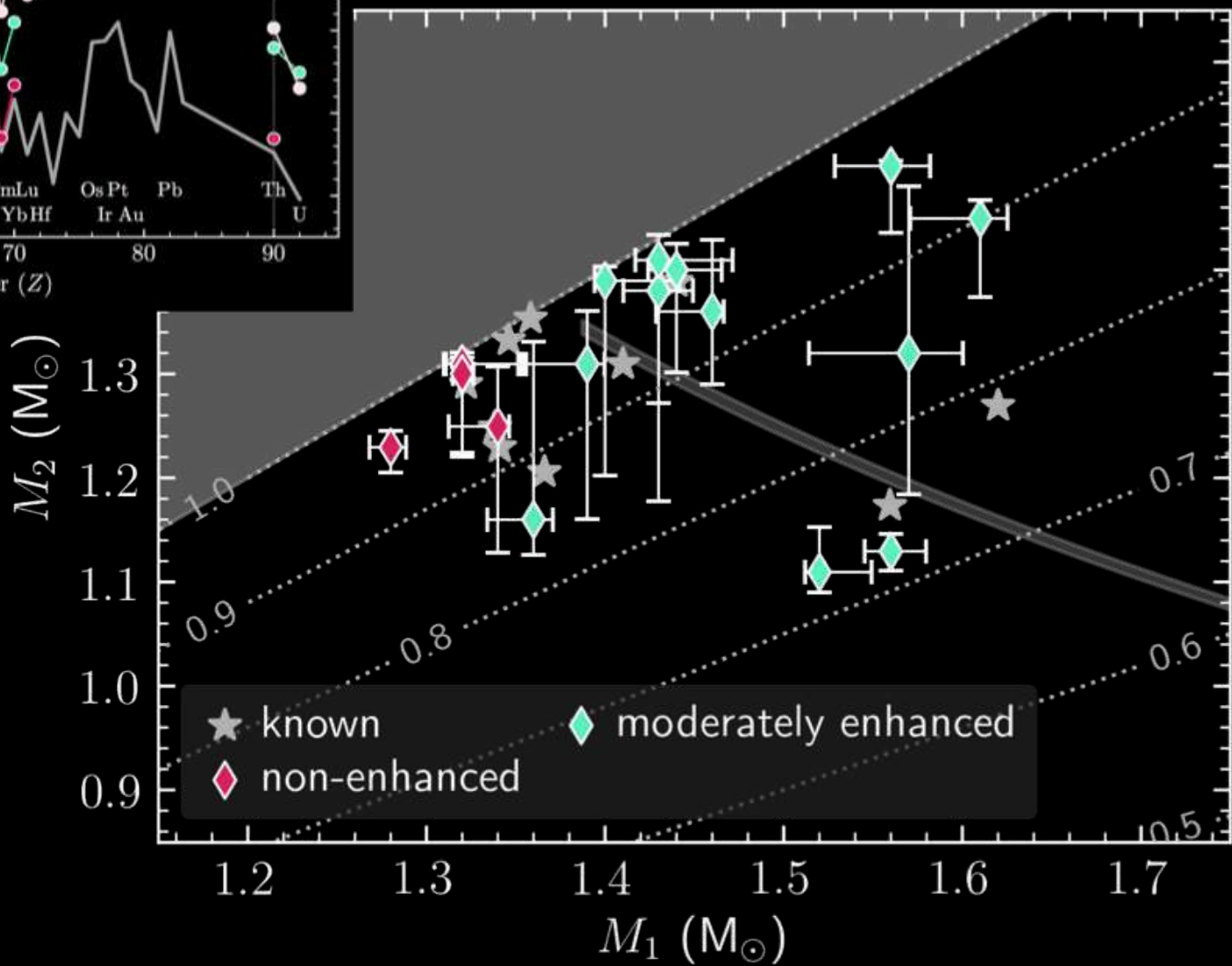
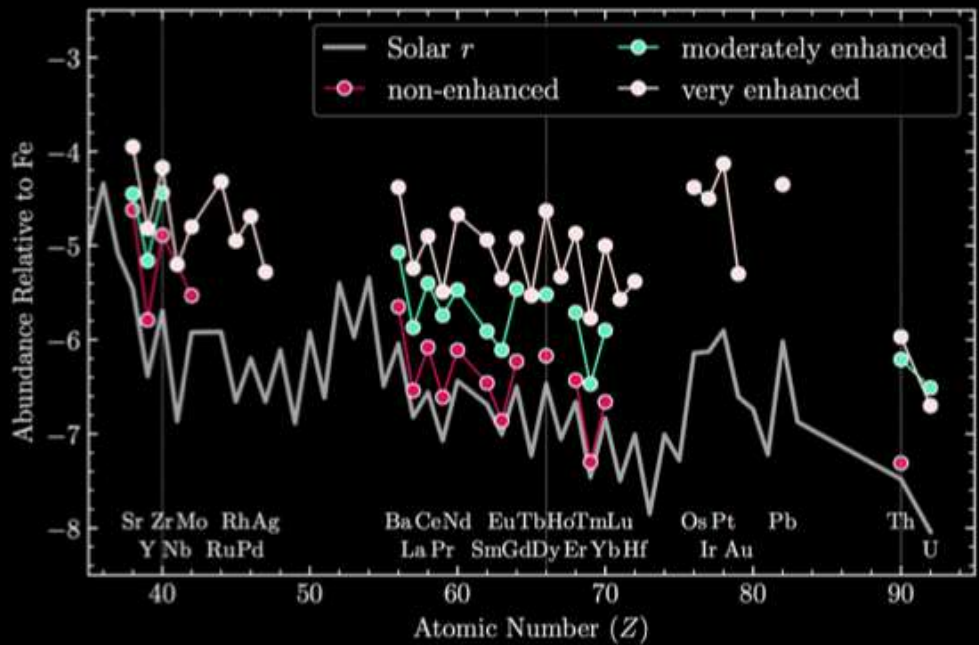
NS masses

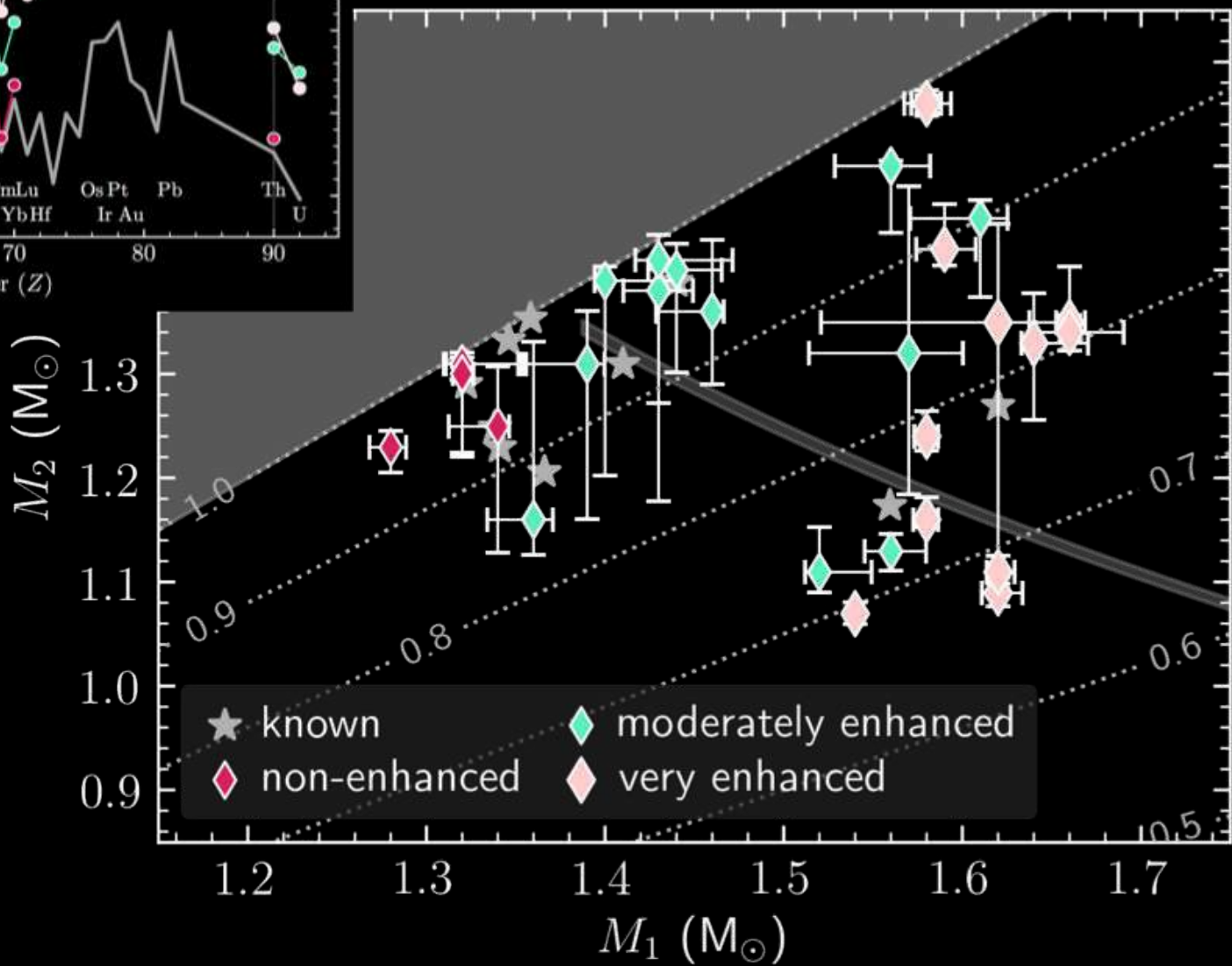
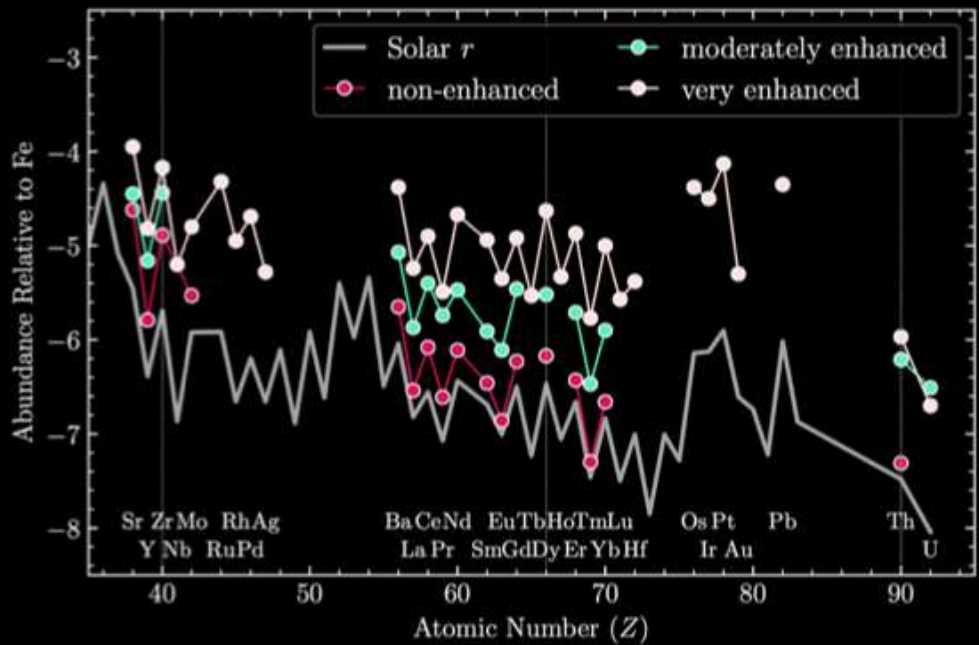
Use an MCMC method to explore the M_1 - M_2 parameter space

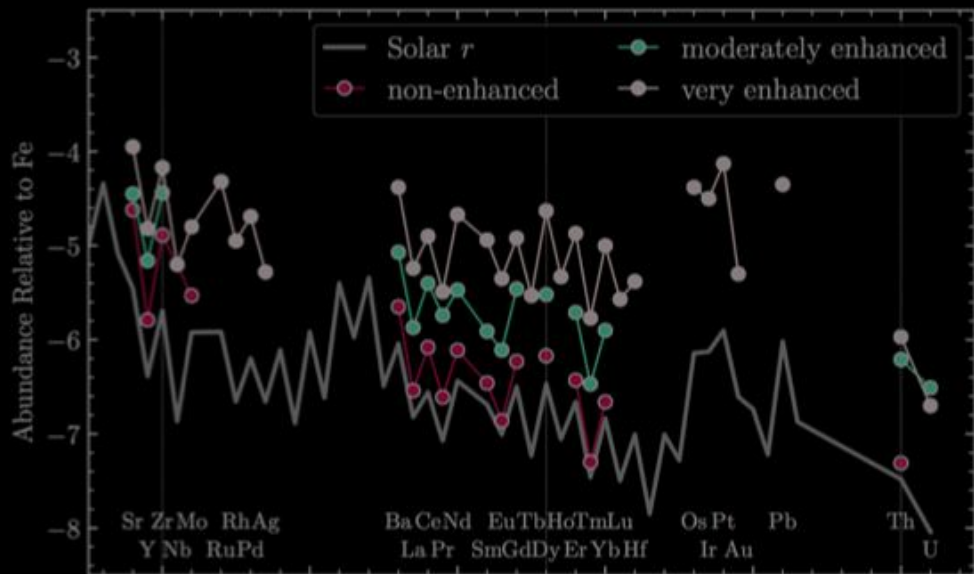




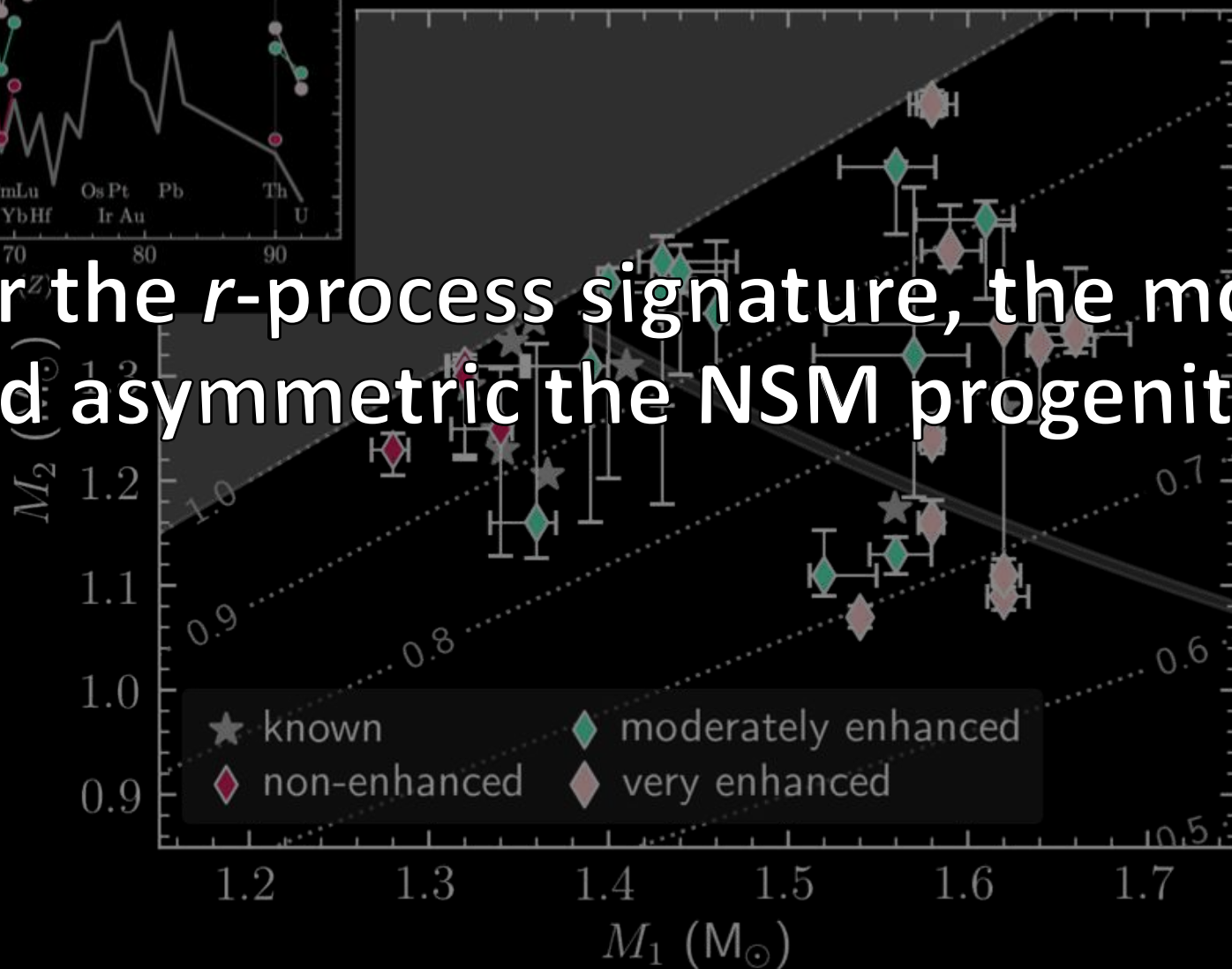




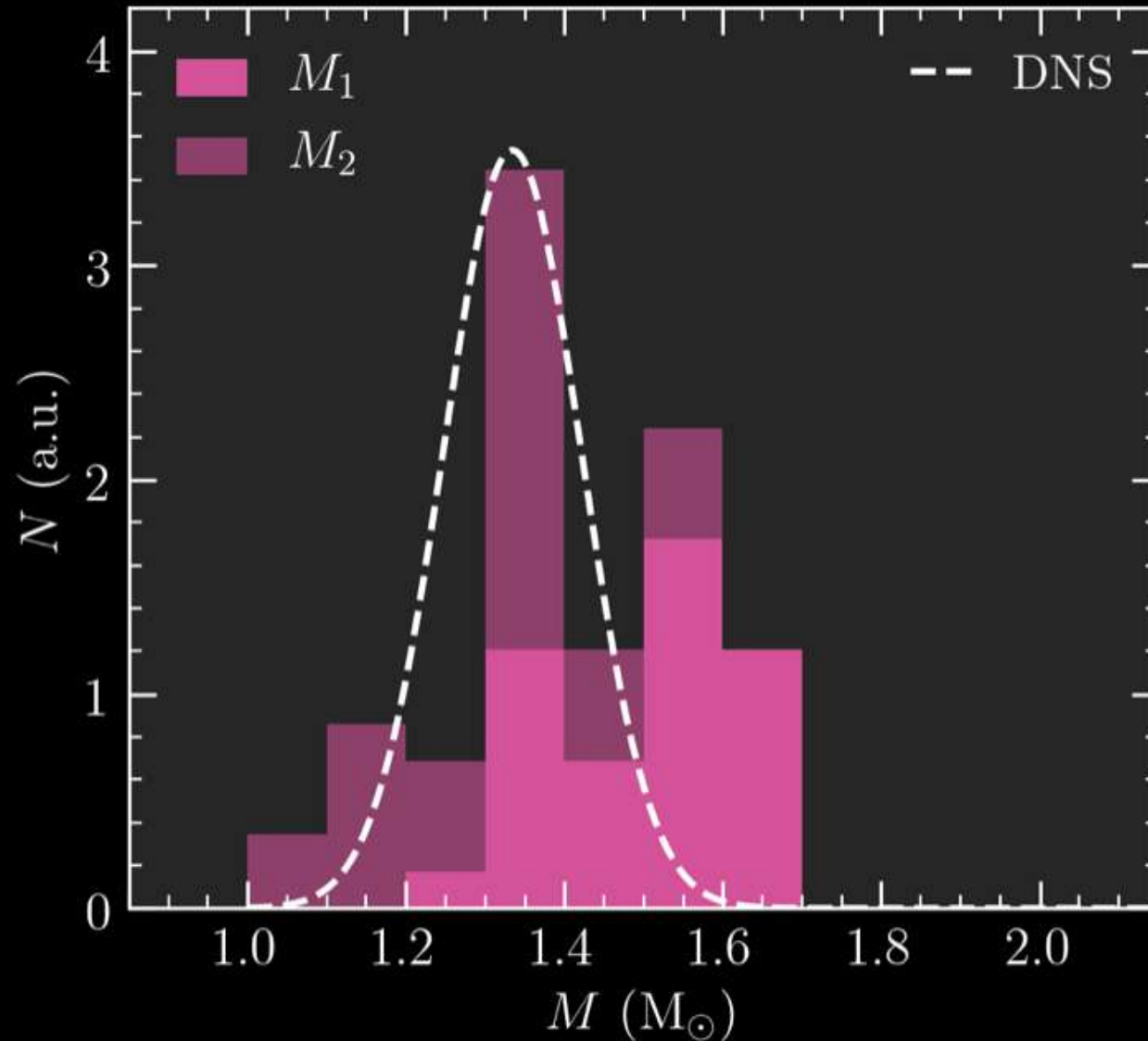




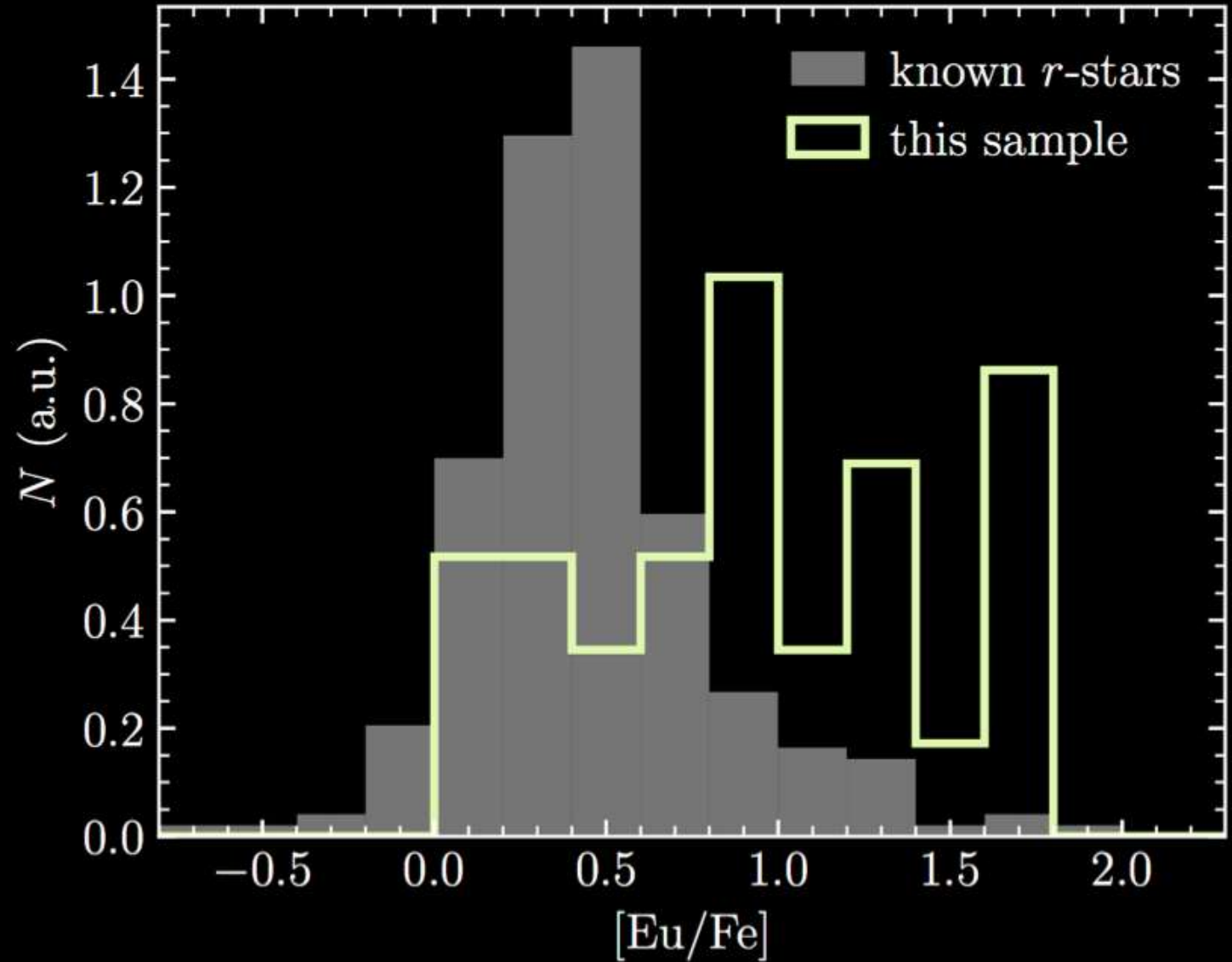
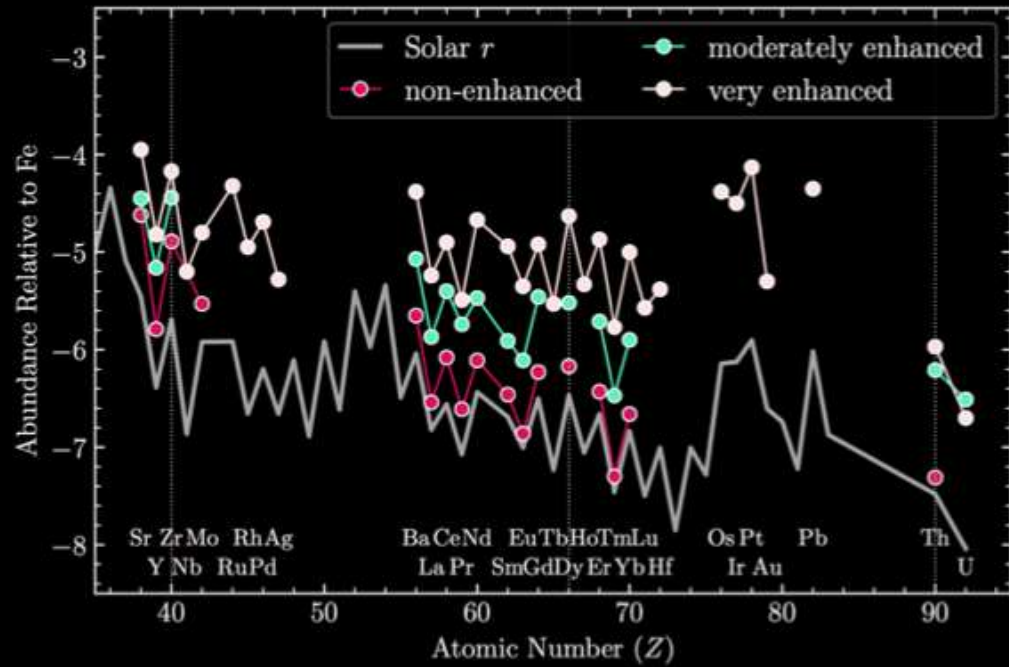
The stronger the r -process signature, the more massive and asymmetric the NSM progenitor



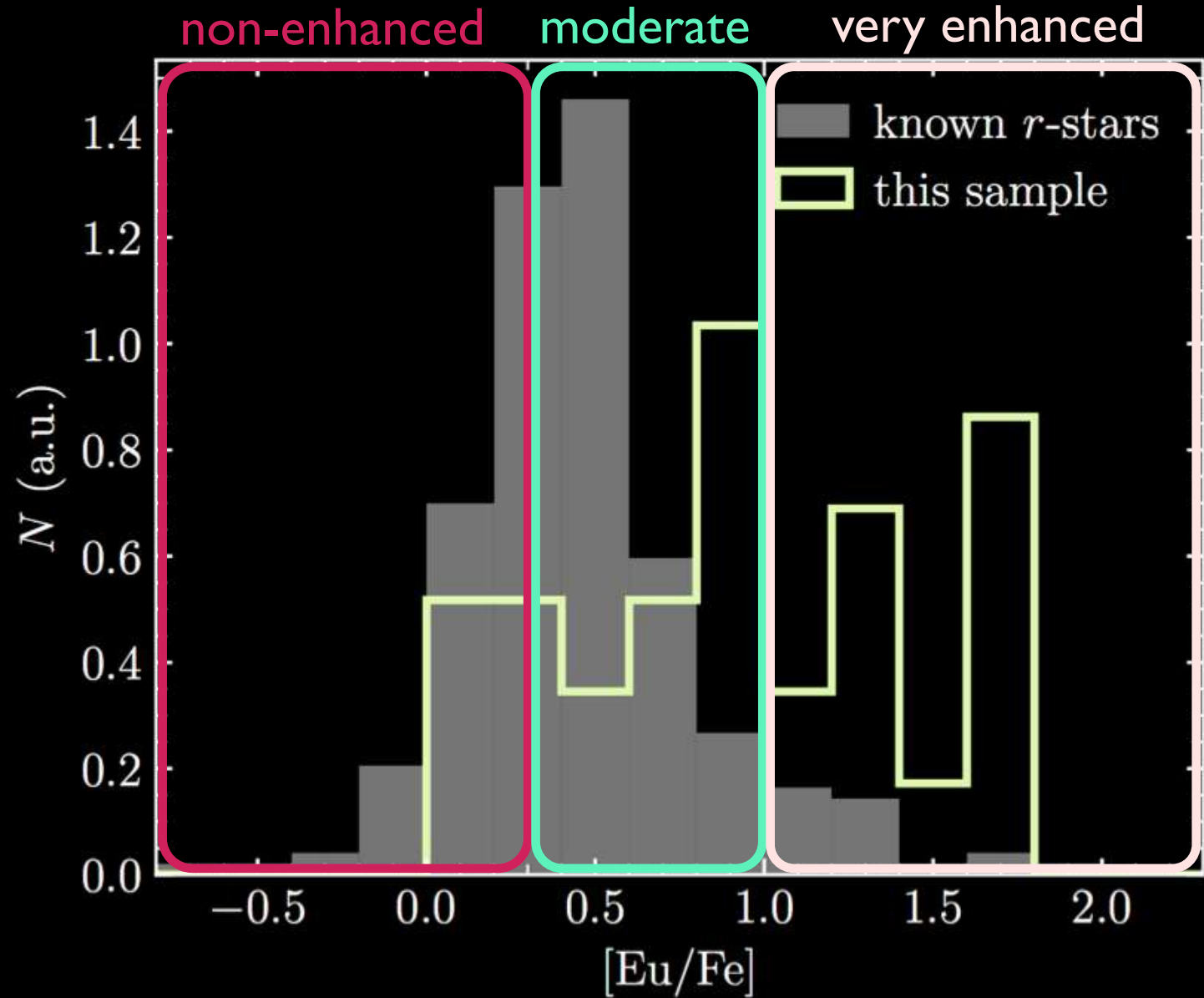
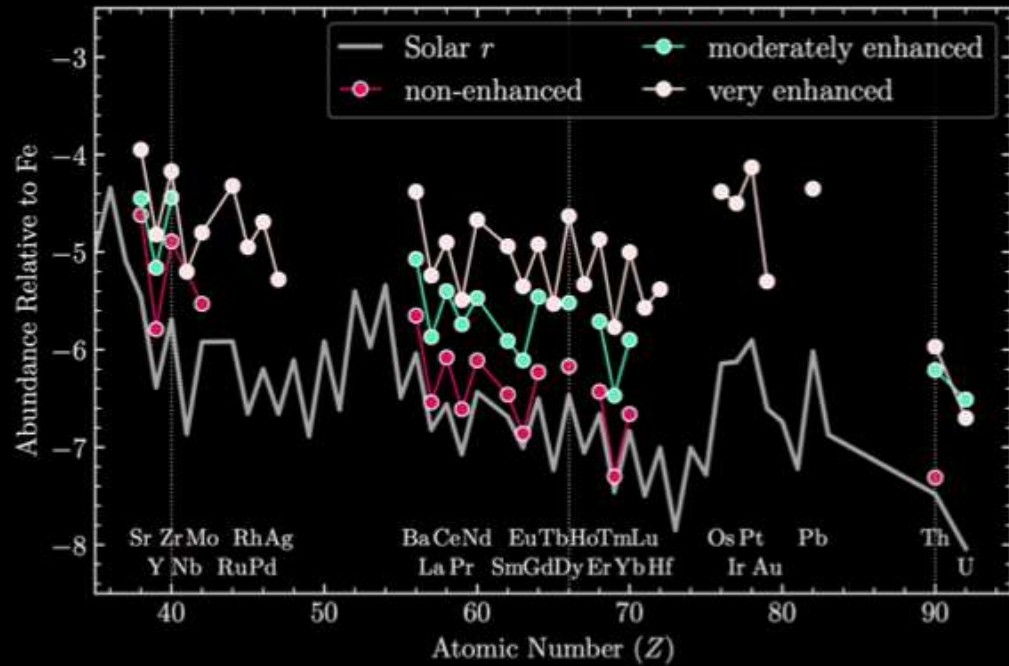
Results require more **higher-mass** DNS systems than are currently found in the Milky Way



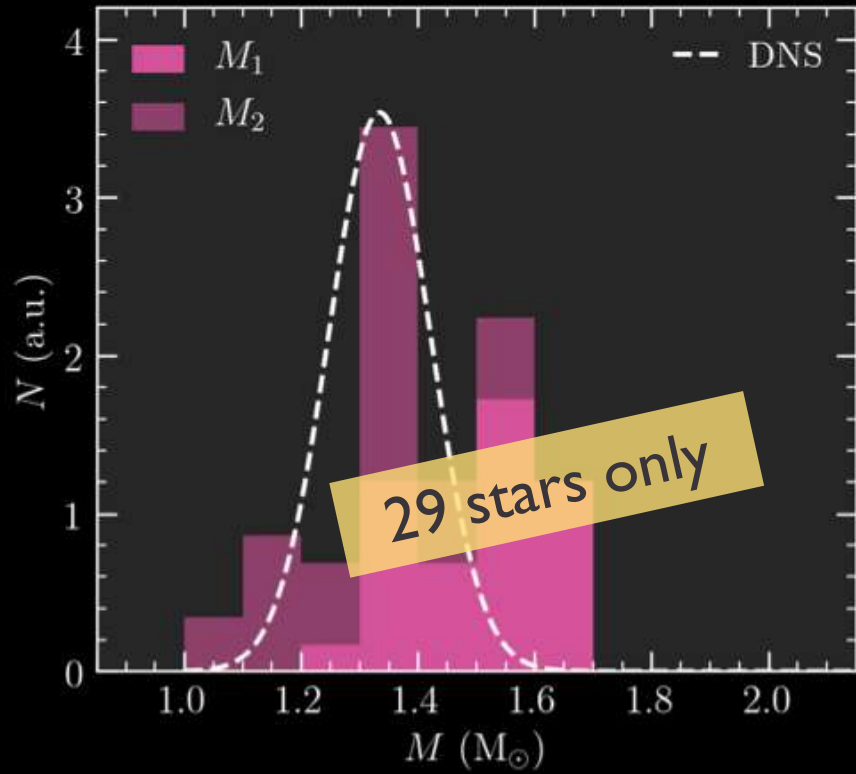
The ~30 stars in our sample are already **more enhanced** than typical *r*-stars



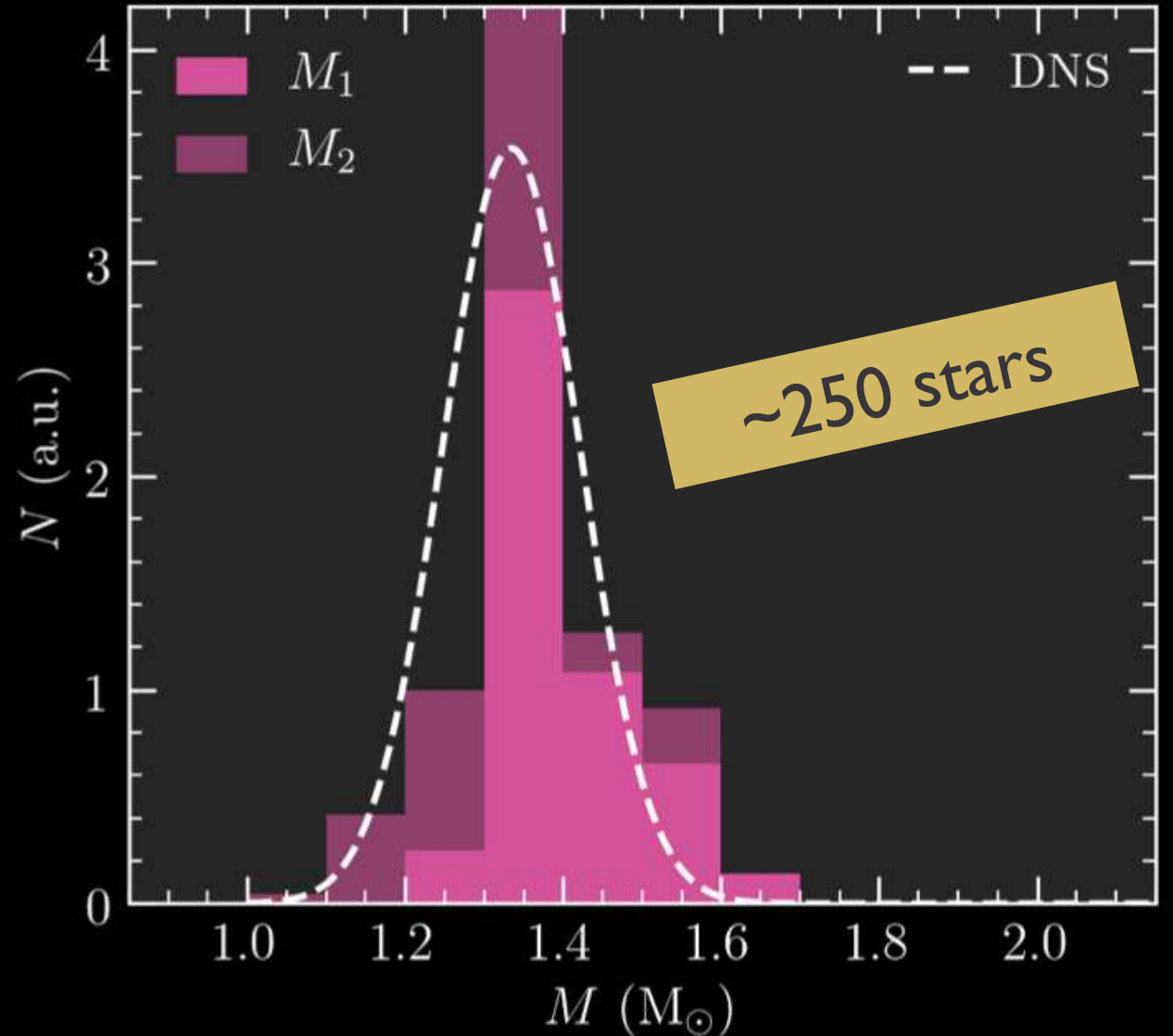
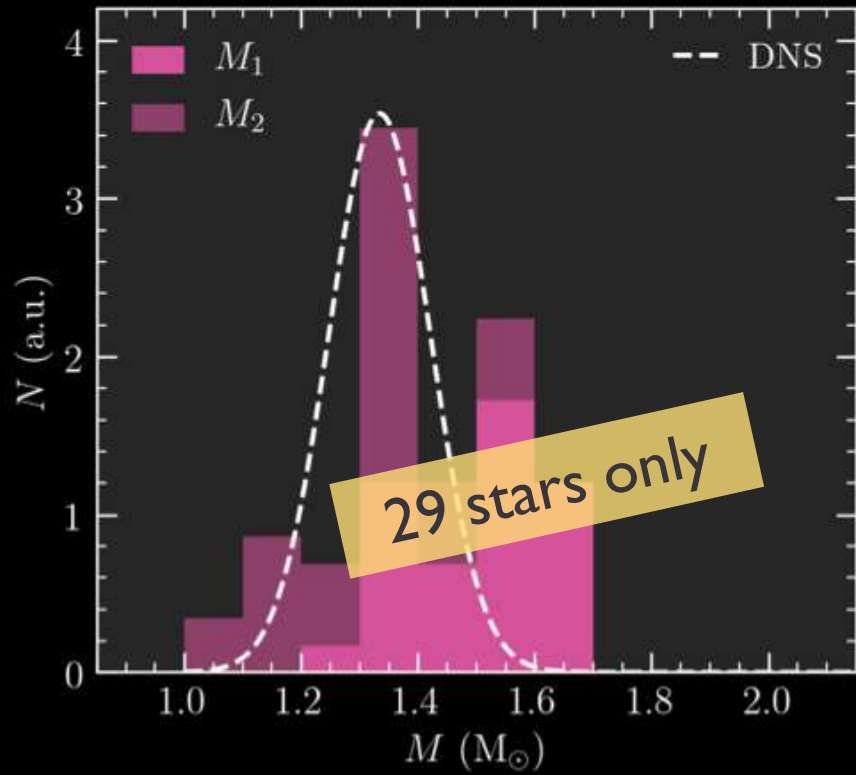
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Assuming that **much of *r*-process material** came from NSMs, those *past* NS-NS systems are similar to *present-day* NS-NSs



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Summary & Implications

Differences in stellar abundances can be used
to study **specific merger conditions**

NSMs could be the **dominant source** of
r-process material in the MW

Inherently assuming that **all MP *r*-stars** have NSM origins

Implies that the **past** DNS population was
similar to the **present-day** distribution

This is not limited to *r*-process or even NSM studies...
Other stellar signatures (e.g., the *i*-process) exist!

Richard O'Shaugnessy (RIT)

Rebecca Surman (ND)

Trevor Sprouse (LANL)

Nicole Vassh (ND)

Matt Mumpower (LANL)

Gail McLaughlin
(NC State)

Terese Hansen (TAMU)

Anna Frebel (MIT)

Ian Roederer (UMich.)

Charli Sakari (SFSU)

Rana Ezzeddine (UFL)

Timothy Beers (ND)

Vini Placco (NOIRlab)

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Rodrigo Fernandez
(U.Alberta)

Xilu Wang (ND)

RIT | Rochester Institute of Technology



JINA-CEE

