`Nickel Mass Problem' in the CCSN Explosion Mechanism, and Neutrino-Driven Wind Model as a Solution to it.

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Reference :

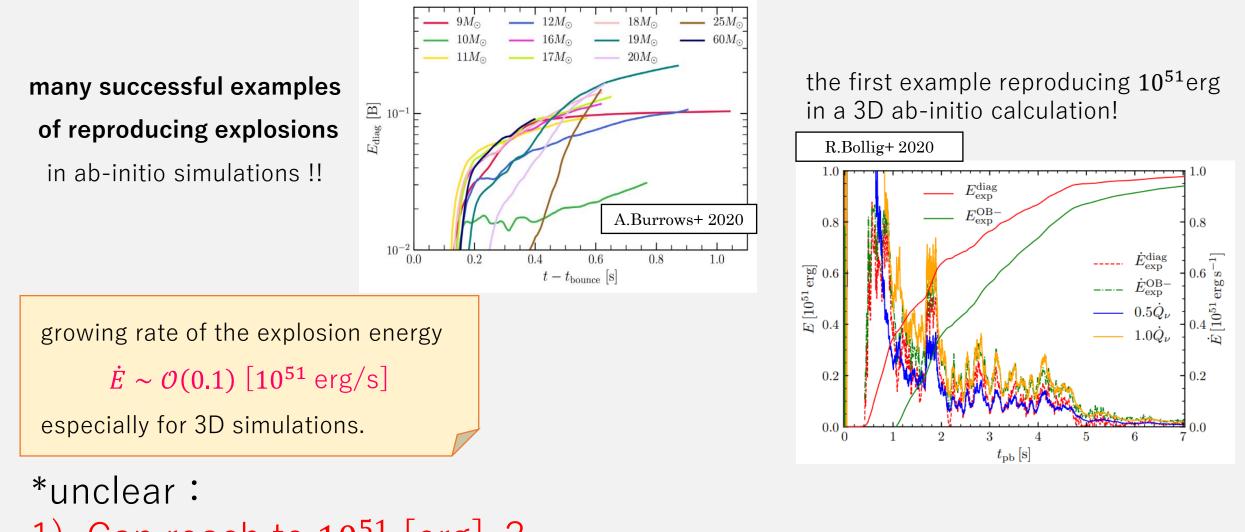
- "A Consistent Modeling of Neutrino-driven Wind with Accretion Flow onto a Protoneutron Star and its Implications for 56Ni Production"
 Sawada & Suwa (2021), ApJ, 908, 6 (arxiv. 2010.05615)
- "Nucleosynthesis Constraints on the Energy Growth Timescale of a Core-collapse Supernova Explosion"

Sawada & Maeda (2019), ApJ, 886, 47 (arxiv.1910.06972)

Explosion mechanism of Core-Collapse SNe

observation neutrino-driven explosion Popov (1993) calibration Gravitational **PNS** formation T.Müller et al.(2017) core collapse & Core bounce -1.0core collapse Proto neutron star Observed $\log_{10}(M_{56Ni}/M_{\odot})$ (PNS) $E_{\rm kin} \gtrsim 10^{51} {\rm ergs}$ O,Ne,Mg C He -1.5 **CCSNe** Shock stall Shock revial -2.0explosion !! $\Sigma = 0.20^{+0.03}_{-0.03}$ 0.0 0.5 1.0 1.5 2.0 $\log_{10}(E_{exp}/10^{50} \text{ ergs})$ unclear (average property) $E_{\rm kin} \sim 10^{51} \, [{\rm erg}]$ $M_{56ni} \sim 0.07 \,[M_{\odot}]$ *unclear: 1). Can reach to 10^{51} [erg] ? my interest 2). Can synthesize a sufficient amount of 56Ni?

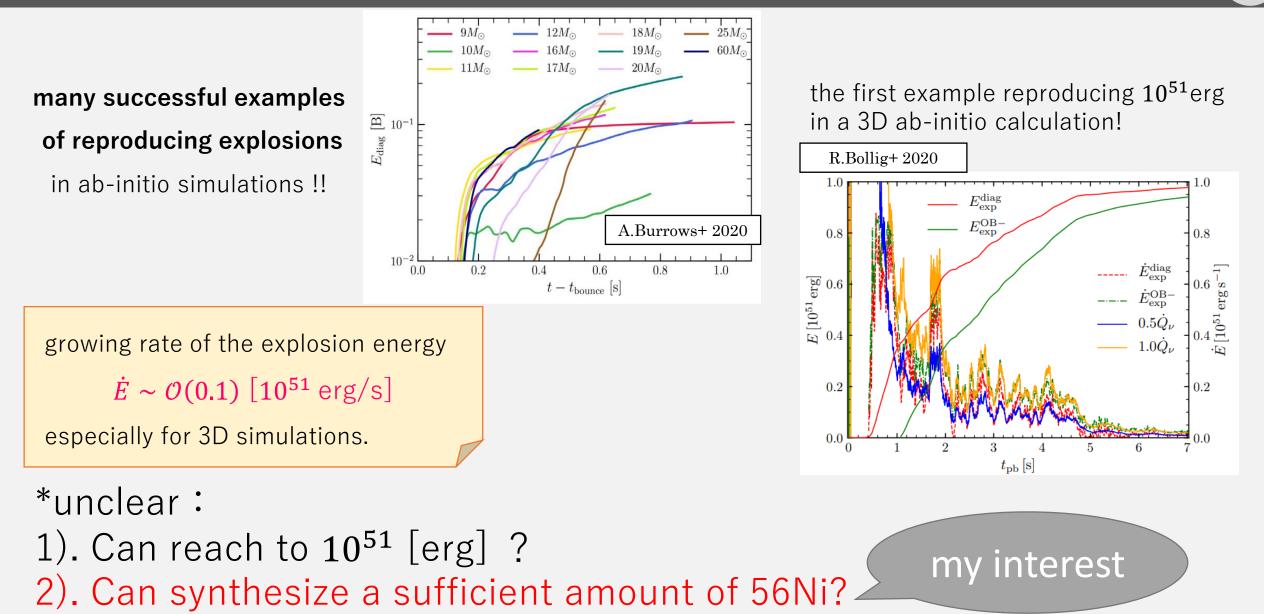
Current results from the ab-initio calculation



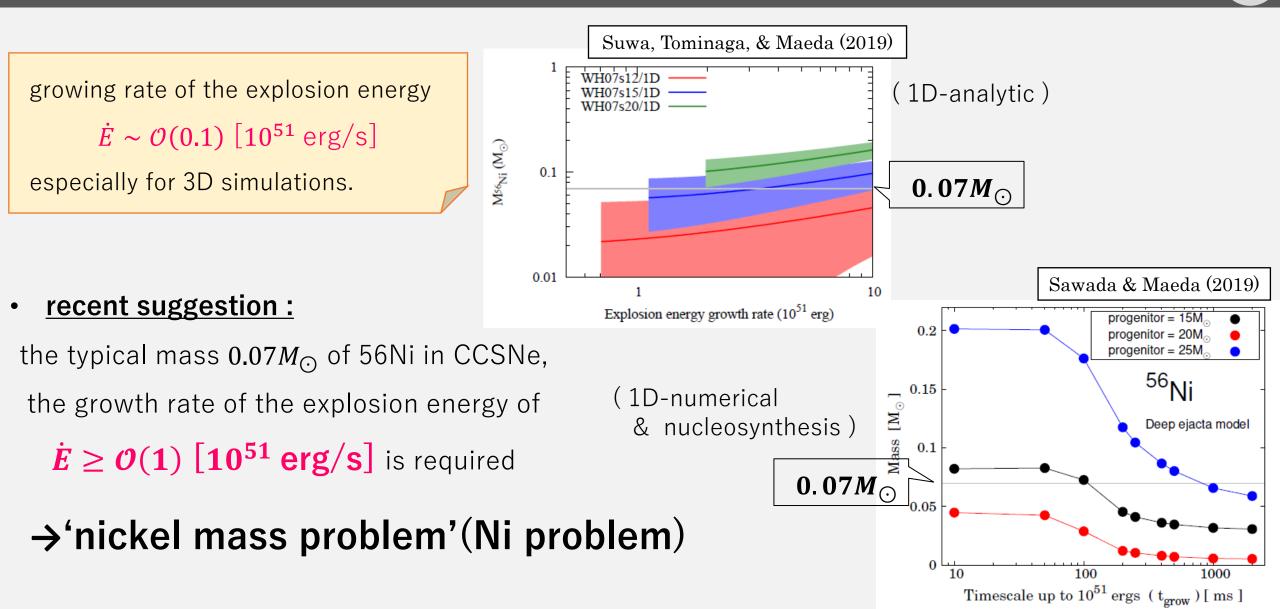
- 1). Can reach to 10^{51} [erg] ?
- 2). Can synthesize a sufficient amount of 56Ni?

Sawada & Suwa (ApJ, 2021) arxiv. 2010.05615

Current results from the ab-initio calculation



Sawada & Suwa (ApJ, 2021) arxiv. 2010.05615



Sawada & Suwa (ApJ, 2021) arxiv. 2010.05615

(1) radiation dominant & isothermal @post-shock region,

(2) adiabatic/constant vel. expansion($r_{\text{shock}} = v_{\text{shock}} \cdot t$),

$$E_{\rm exp} = (aT^4) \times \left(\frac{4\pi}{3}r_{\rm shock}^3\right) \quad \Rightarrow T_{\rm peak} \propto t^{-3/4}$$

Suwa, Tominaga, & Maeda (2019)

Sawada & Maeda (2019)

 \checkmark When the explosion is sufficiently instantaneous, then

Deposit $E_{exp} = 10^{51} erg$ from the initial setup.

 \checkmark When taking into account the growth timescale

 $E_{\exp}(t) = \frac{10^{51} \text{erg}}{t_{\text{grow}}} \cdot t$

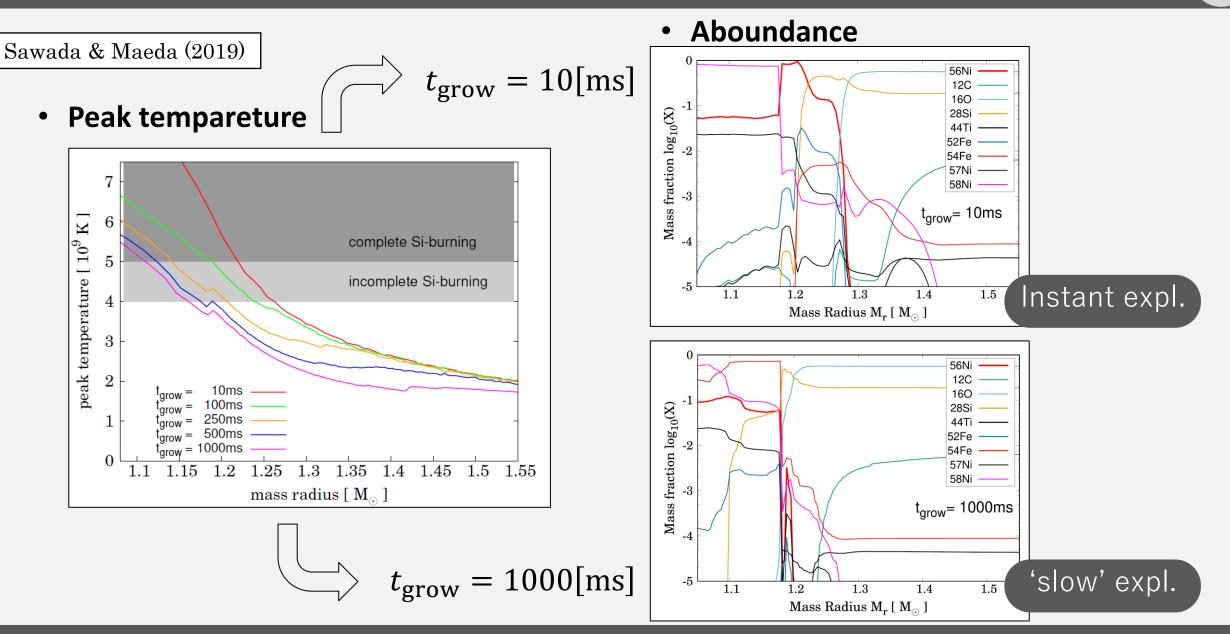
of the explosion energy, then

e.g., Woosley+ 2002

this work

Woosley+02 $t_{grow} = 100 ms$ 12 $t_{grow} = 200 ms$ Instantaneous _10 $t_{grow} = 400 ms$ explosion 10^{9} $t_{grow} = 1000 ms$ 8 T_9 [K / 6 "slow" explosion 0 0.10.30.20.4 0.50 (or Radius) time [sec]

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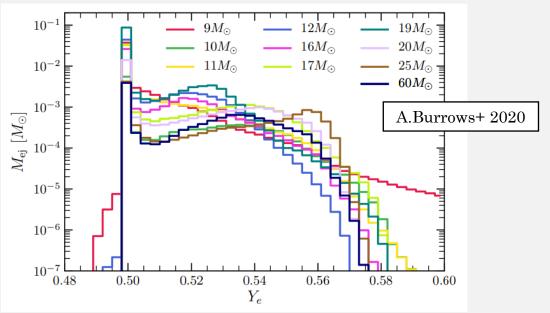


Sawada & Suwa (ApJ, 2021) arxiv. 2010.05615

• <u>NOTE :</u>

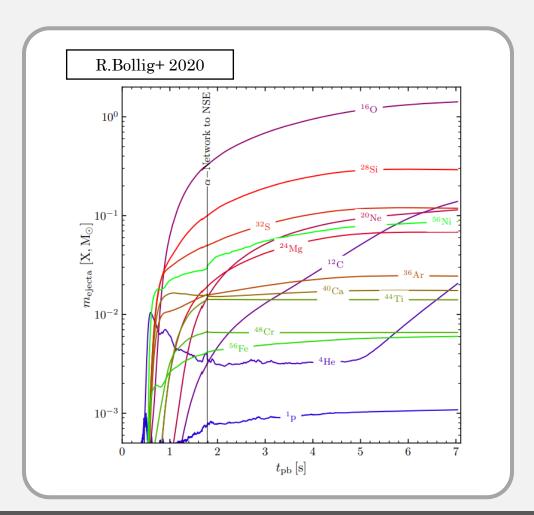
some models in ab-initio simulations have succeeded

in producing the typical mass $0.07 M_{\odot}$ of 56Ni in CCSNe,

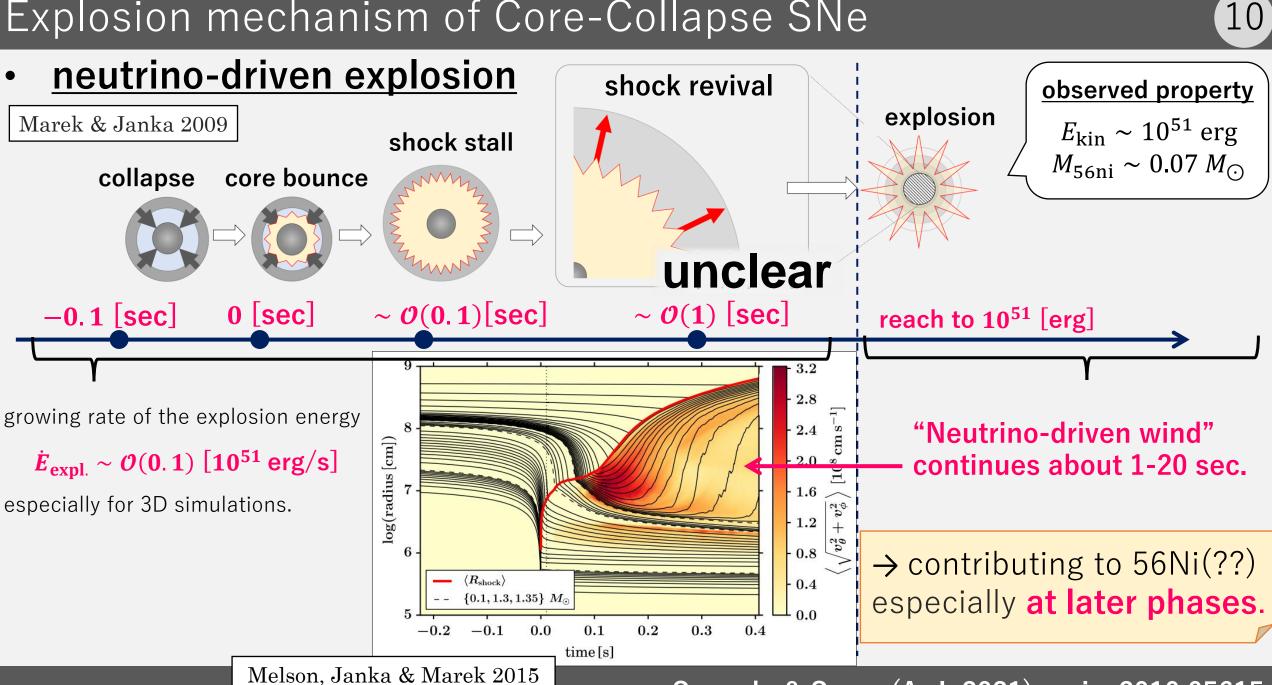


But 'nickel mass problem'(Ni problem)

→ unclear whether we can reproduce sufficient 56Ni amount <u>as a canonical nature.</u>



Explosion mechanism of Core-Collapse SNe



aim and content of our work

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problem

1D (e.g., Wanajo2013) :

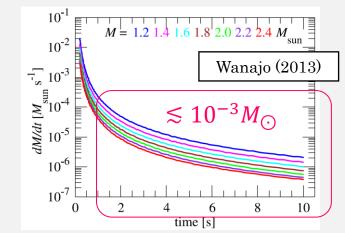
it is already known that 1D wind simulation could not solve Ni problem. (\rightarrow multi-D, especially energy injection by accretion, is important.)

• multi-D (e.g.,Wanajo+2018) :

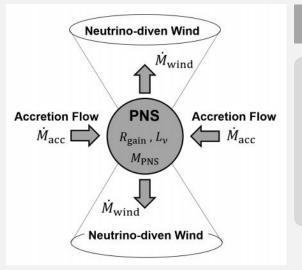
it may solve the Ni problem .

 \rightarrow However, calculation time is limited.

Can it be solved if we follow it to the late explosion stage?

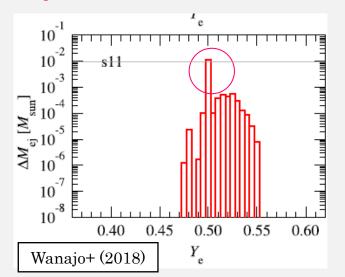


 \gtrsim 0.01 M_{\odot} ejection is required for Ni problem



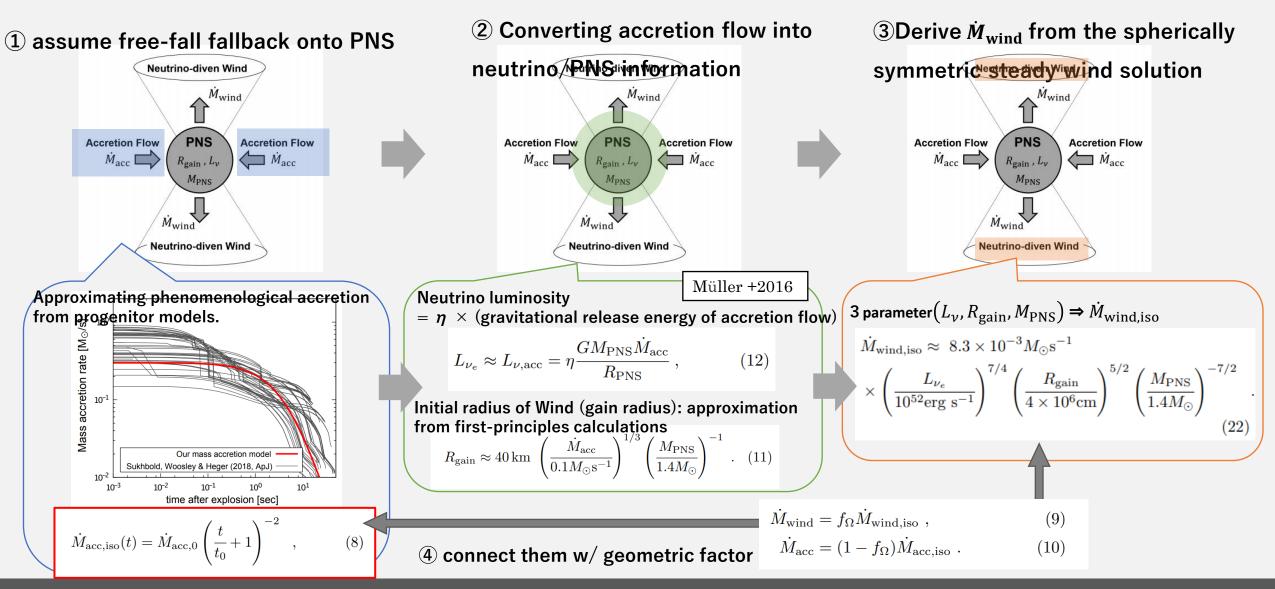
our work

- 1. Build a consistent model of the neutrino-driven wind with an accretion flow onto a PNS.
- 2. Investigate the possibility that neutrino-driven wind can solve the "Ni-problem"



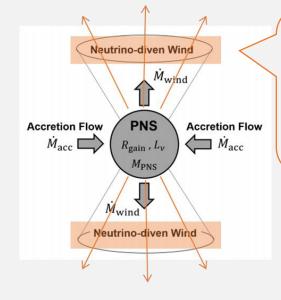
aim and content of our work

our work the neutrino-driven wind with an accretion flow onto a PNS.



Sawada & Suwa (ApJ, 2021) arxiv. 2010.05615

semi-analytic wind model (e.g., Otsuki et al. 2000).



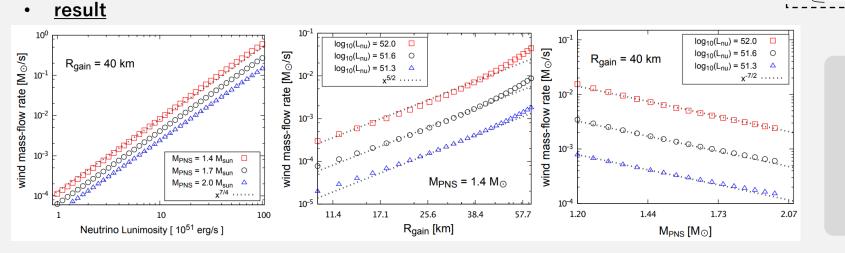
Assume a neutrino-driven wind blowing in the "radial direction" through the low-density region swept by the non-spherically symmetric shock wave

<u>spherically symmetric steady-state wind</u> <u>solution(e.g., Otsuki et al. 2000).</u>

• $(R_{\text{gain}}, L_{\nu}, M_{\text{PNS}})$ \Rightarrow tran-sonic solution $\nu_{\text{tran}} \Leftrightarrow \text{maximum } \dot{M}_{\text{wind}}$ • $v \frac{dv}{dr} = -\frac{1 + (v/c)^2 - (2GM/c^2r) dP}{\rho(1 + \epsilon/c^2) + P/c^2} \frac{dP}{dr} - \frac{GM_{PNS}}{r^2}$, (2) • $\dot{Q} = v \left(\frac{d\epsilon}{dr} - \frac{P}{\rho^2} \frac{d\rho}{dr}\right)$, (3) • Helmholtz EoS (Timmes & Swesty 2000), • Boundary condition: $r = R_{gain} (\dot{Q} \approx 0)$ Given: $\rho_0 = 10^{10} \text{g cm}^{-3} \cdot L_{\nu,51}^{1/2}$, (Fujibayashi+ 2015) • $Y_e = \left[1 + \frac{L_{\nu_e}^n \langle \sigma_{\overline{\nu_e}p} \rangle}{L_{\nu_e}^n \langle \sigma_{\nu_e n} \rangle}\right]^{-1} = 0.5$, (4) (e.g., Bliss+ 2018).

(1)

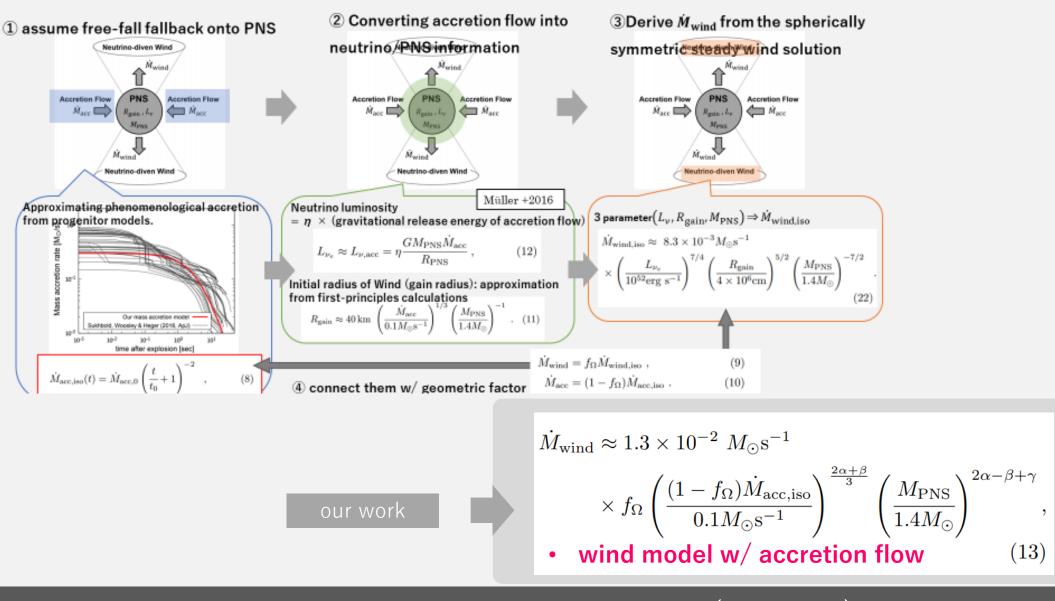
• $\dot{M} = 4\pi r^2 \rho v$,



• spherical wind $\dot{M}_{wind,iso}$

 $\dot{M}_{\rm wind,iso} \approx 8.3 \times 10^{-3} M_{\odot} {\rm s}^{-1} \\ \left(\frac{L_{\nu_e}}{10^{52} {\rm erg \ s}^{-1}}\right)^{\alpha} \left(\frac{R_{\rm gain}}{4 \times 10^6 {\rm cm}}\right)^{\beta} \left(\frac{M_{\rm PNS}}{1.4 M_{\odot}}\right)^{\gamma},$ (5)

ur work the neutrino-driven wind with an accretion flow onto a PNS.



maximum parameter sets

 $M_{\rm PNS,0} \ge 1.4 M_{\odot},$

• total ejected mass of the wind...

$$\begin{split} M_{\rm ej,\infty} &= \int_0^\infty dt \ \dot{M}_{\rm wind} \\ &\approx 4.3 \times 10^{-3} M_\odot s^{-1} f_\Omega (1 - f_\Omega)^2 \left(\frac{\dot{M}_{\rm acc,0}}{0.1 M_\odot s^{-1}}\right)^2 \left(\frac{M_{\rm PNS,0}}{1.4 M_\odot}\right)^{-\frac{5}{2}} \leq 0.067 M_\odot \end{split}$$

• the time evolution of the cumulative ejected mass of the wind...

$$M_{\rm ej}(t_e) = \int_0^{t_e} dt \, \dot{M}_{\rm wind}$$

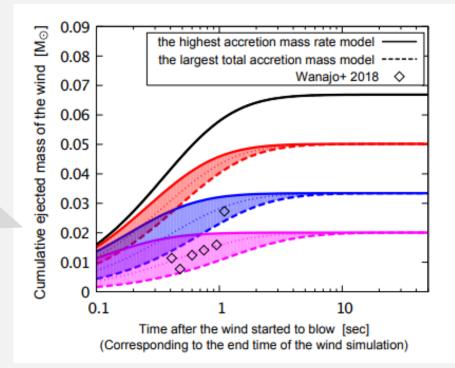
\$\approx 6.4 \times 10^{-4} M_{\overlinessim} \left[1 - \left(\frac{t_0}{t_0 + t_e} \right)^3 \right] \times \left(\frac{t_0}{1s} \right) \left(\frac{\dot{M}_{\rm acc,0}}{0.1M_{\odot}s^{-1}} \right)^2 \left(\frac{M_{\rm PNS,0}}{1.4M_{\odot}} \right)^{-5/2}

- Conclusion.
- 1. the total ejectable is determined within \sim 2 sec from the onset of the explosion.
- 2. the supplementable amount at a late phase (t > 1 sec) remains $M_{ej} < 0.01 M_{\odot}$.

\rightarrow difficult to solve the Ni problem

at the late phase of the explosion by the neutrino-driven wind.

Uur mass accretion model 10⁻¹ 10⁻¹ Uur mass accretion model 10⁻² 10⁻³ 10⁻² 10⁻² 10⁻¹ 10⁻¹ 10⁰ 10¹ time after explosion [sec]

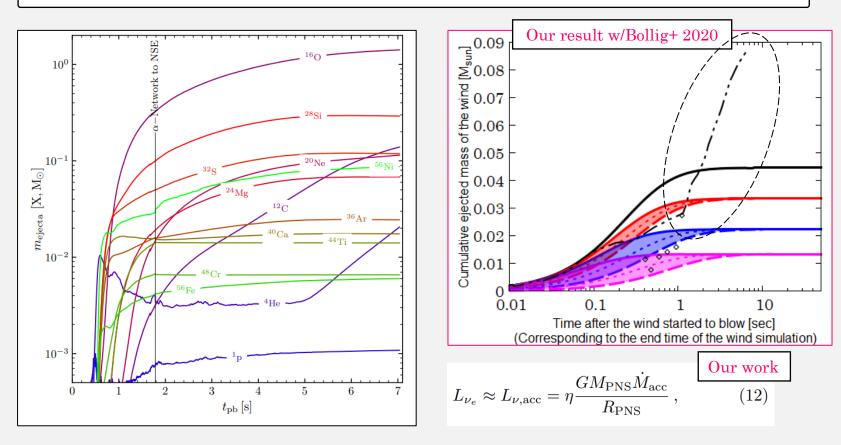


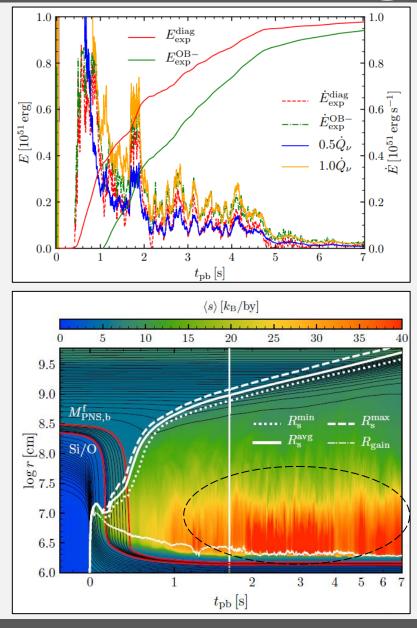
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• **NOTE :** R.Bollig+ 2020

"The converged value of the explosion energy at infinity (with overburden subtracted) is roughly 1B and the ejected 56Ni mass **up to 0.087** solar masses"

"Our final 56Ni mass is therefore an upper limit, and we expect the actual mass to be around **0.05 M**. Nevertheless, it demonstrates that 56Ni masses in the ballpark of those of typical CCSNe can be ejected in 3D neutrino-driven explosions."





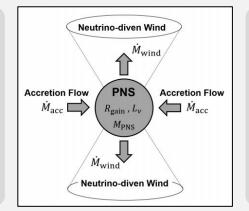
Summary

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1. a model of the neutrino-driven wind with an accretion flow onto a PNS.

spherical wind

$$\dot{M}_{\rm wind,iso} \approx 8.3 \times 10^{-3} M_{\odot} {\rm s}^{-1} \\ \times \left(\frac{L_{\nu_e}}{10^{52} {\rm erg \ s}^{-1}}\right)^{7/4} \left(\frac{R_{\rm gain}}{4 \times 10^6 {\rm cm}}\right)^{5/2} \left(\frac{M_{\rm PNS}}{1.4 M_{\odot}}\right)^{-7/2}$$
(22)



wind model w/ accretion flow

$$\dot{M}_{\text{wind}} \approx 1.3 \times 10^{-2} M_{\odot} \text{s}^{-1} \times f_{\Omega} \left(\frac{(1 - f_{\Omega}) \dot{M}_{\text{acc,iso}}}{0.1 M_{\odot} \text{s}^{-1}} \right)^2 \left(\frac{M_{\text{PNS},0}}{1.4 M_{\odot}} \right)^{-5/2}$$
(23)

2. the possibility that neutrino-driven wind can solve the "Ni-problem"

- 1. the total ejectable is determined within \sim 2 sec from the onset of the explosion.
- 2. the supplementable amount at a late phase (t > 1 sec) remains $M_{ej} < 0.01 M_{\odot}$.

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at the late phase of the explosion by the neutrino-driven wind.

