Directed flow of Λ from heavy-ion collisions and hyperon puzzle of neutron stars

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- Introduction Hyperon puzzle
- **U**_{Λ} from chiral EFT
- Directed flow of Λ
- Summary

Y.Nara, A. Jinno, K. Murase, AO, in prep.



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Hyperon Puzzle of Neutron Stars

- Observation of massive neutron stars rules out hyperonic EOS ?
 - Attractive $U_{\Lambda}(\rho)$ causes hyperon mixing in NS at (2-4) ρ_0 , softens the EOS, and reduces $M_{max} = (1.3-1.6) M_{\odot}$
- Proposed solutions
 - Three-body ANN repulsion \rightarrow repulsive U_A(ρ) at high density
 - Transition to quark matter before Λ appears
 - General relativity → Modified gravity



Repulsive $U_{\Lambda}(\rho)$ at high density in chiral EFT

- Chiral effective field theory (chiral EFT) may cause repulsive Λ potential at high densities *Gerstung, Kaiser, Weise (2001.10563), Kohno (1802.05388)*
- **Yet unknown parameters are tuned to support 2** M_{\odot} **neutron stars.**
 - → Repulsion at high densities needs to be verified e.g. in heavy-ion collisions.



Directed flow of protons

- Directed flow has been utilized to study EOS
 - $v_1 = \langle \cos \phi \rangle$ (directed flow), $\langle p_x \rangle$ (side flow)

E.g. Sahu, Cassing, Mosel, AO (nucl-th/9907002), Snellings+(nucl-ex/9908001)

- Non-monotonic colliding energy dep. of proton v1 slope has been a puzzle STAR (1401.3043), Nara+(JAM, '16,'17), Ivanov+ (3FD, 1601.03902), Konchakovski+ (PHSD, 1404.2765)
 - An explanation is found. *Nara+(2109.07594)* **Compression & expansion stages** cause different slope
 - \rightarrow non-monotonicity.











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Since the directed flow of protons can be described by a single EOS, U_{Λ} may be constrained by Λ flow.

We study Λ flow using U_{Λ} from chiral EFT. This is complementary to precision hypernuclear spectroscopy.







Chiral EFT (decuplet saturation model)





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Momentum dependence

- 3BF from chiral EFT comes from higher-order diagrams → Larger number of propagators and derivatives
 - → Momentum dependence is needed (also from BHF calc.)
- U_Λ(k, ρ) from chiral EFT seems to have stronger mom. dep. than quark model YN potential.





U_{Λ} from Chiral EFT

Chiral EFT with 3BF and hyperons

Gerstung+(2001.10563)(GKW, decuplet saturation model), Kohno (1802.05388)

ρ-dep. potential using Fermi mom. expansion *Tews+(1611.07133)* + momentum dep. fitted to *Kohno('18)*

$$U_{\Lambda}(\rho, k) = a \frac{\rho}{\rho_0} + b \left(\frac{\rho}{\rho_0}\right)^{4/3} + c \left(\frac{\rho}{\rho_0}\right)^{5/3} + \sum_n \frac{C_n}{\rho_0} \int \frac{d\mathbf{k}'}{(2\pi)^3} \frac{f(\mathbf{r}, \mathbf{k}')}{1 + (\mathbf{k} - \mathbf{k}')^2 / \mu_n^2}$$

Range of fit







Directed flow of Λ

Directed flow of Λ is expected to be smaller than v₁(p) from the compression+tilted expansion mechanism, but data show v₁(Λ) ~ v₁(p) *STAR*, *PRL120* ('18),062301 (1708.07132) → Stronger repulsion for Λ at high densities ?



Transport model calculations using an event generator JAM2.1.

- Collision time scheme is updated X.-L. Zhao+(2001.10140)
- Quark potential in the leading hadron during the formation time is switched on. The gaussian width (simulating interaction range) is tuned to explain proton flow.



Comparison with data

- $v_1(\Lambda)$ at (3-27) GeV *STAR (PRL, (1708.07132); PLB (2108.00908)*).
- Cascade (w/o potential effects) does not explain the data.
- Chiral EFT U_{Λ} (GKW3) explains the data well.
- Strong mom.-dep. (MD1) suppresses v₁ at large |y|.





Comparison with data



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v₁ slope

- **U**_{Λ} from chiral EFT describes the v₁ slope of Λ (dv₁(Λ)/dy) well.
 - Repulsive potential at high densities may realize a large v₁ value at low colliding energy.
 - dv₁/dy at 3 GeV is slightly underestimated (also for protons).



MS2: mom.-dep. soft (K=210 MeV) pot. for N 2 GKW3: chiral EFT w/2+3-body force, no mom. dep. GKW3+MD1/MD2:

chiral EFT w/ 2+3-body force, with mom. dep.

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 $\sqrt{s_{NN}}$ (GeV)



Summary

- The directed flow (v₁) of Λ from HIC is studied by using the Λ potential from chiral EFT with 2+3-body interactions, which can support 2 solar mass neutron stars.
 - A flow is found to be sensitive to U_{Λ} !
 - U_{Λ} from chiral EFT with no or weak momentum dependence is consistent with the directed flow data from heavy-ion collisions. [Similar results for $< p_x > at \sqrt{s_{NN}} = 3.0$ GeV, D.C. Zhang+ (2107.00277)]
 - Momentum dependence of U_{Λ} needs to be evaluated carefully.
 - The forward and backward v_1 values seem to be more sensitive to the Λ potential at high densities and/or high momentum.
 - v₁ of p and Λ are sensitive to the details of the transport scheme, such as the potential during the formation time and the collision time choice.
- **How can we pin down U_{\Lambda} at high densities further ?**
 - More data at 3-10 GeV with wider rapidity coverage.
 - Elliptic flow (v₂) and other observables
 - A-nucleus scattering (Emulsion or Femtoscopy) \rightarrow mom. dep.
 - Hypernuclear spectroscopy



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Thank you for your attention !



Effects of mom.-dep. potential

- Momentum dep. of U_Λ enhances
 both compression / expansion effects on v₁ slope.
 - More repulsion for finite p particles in the compression stage.
 - Stronger reduction of v1 from repulsion partly from spectators (?).

0.05

0.04

، 0.03 ح

0.02

0.01

0.00



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Approaches to repulsive U_{Λ} at high densities

- **Top-down approach: Lattice QCD calculation at high densities**
 - The sign problem needs to be solved.
- Bottom-up approach: Precision experiments of hypernuclei
 - Comparison of data and calculated results with (reliable) AN potential together with (uncertain) ANN potential
 - Extrapolation to high densities ($\rho \sim \rho_0$ to $\rho = (2-5)\rho_0$) is necessary.
- Phenomenological approach: Hyperons in heavy-ion collisions
 - High-density matter is created during HIC.
 - A flow using theoretical U_{Λ} can be compared with data.
 - Non-equilibrium effects need to be taken care of.

We study Λ flow in HIC using U_{Λ} from chiral EFT (Phenomenology based on first-principles results)





Past tries



(There was a mistake...)



Directed flow of protons



Semi-Classical Nuclear Transport Theories

- Wigner(-Weyl) transform of TDHF = Vlasov equation
 - Wigner transform of density matrix=Wigner fn. (phase space dist.)
 - Wigner transform of commutator ~ i $\hbar \times$ Poisson bracket

$$i\hbar \frac{d\rho}{dt} = [h, \rho] \to \frac{\partial f}{\partial t} + \boldsymbol{v} \cdot \boldsymbol{\nabla} f - \boldsymbol{\nabla} U \cdot \boldsymbol{\nabla}_p f = 0$$
$$[f = \rho_W, [A, B]_W = i\hbar \{A_W, B_W\}_{PB} + \mathcal{O}(\hbar^2)]$$

• Test particle solution of the Vlasov equation \rightarrow Classical EOM

$$f(\boldsymbol{r}, \boldsymbol{p}) = \frac{(2\pi)^3}{N} \sum_{i=1,NA} \delta(\boldsymbol{r} - \boldsymbol{r}_i) \delta(\boldsymbol{p} - \boldsymbol{p}_i)$$
$$\rightarrow \frac{d\boldsymbol{r}_i}{dt} = \frac{\partial h}{\partial \boldsymbol{p}}\Big|_{\boldsymbol{p}=\boldsymbol{p}_i} = \frac{\boldsymbol{p}}{m} + \frac{\partial U}{\partial \boldsymbol{p}}\Big|_{\boldsymbol{p}=\boldsymbol{p}_i}, \ \frac{d\boldsymbol{p}_i}{dt} = -\frac{\partial U}{\partial \boldsymbol{r}}\Big|_{\boldsymbol{r}=\boldsymbol{r}_i}$$

- Relativistic Quantum Molecular Dynamics
 - Transport model applicable to high energies Sorge, Stoecker, Greiner ('89); Maruyama et al. ('96)
 - Stronger potential effects are necessary → Vector potential Nara et al. ('20), Nara, AO ('21)
 - Stochastic collisions are also included



Can we access EOS by using flows ?

- EOS from Flow is a Notorious problem!
 - Momentum-dependent potential can simulate stiff EOS, and then we cannot extract stiffness. (1980s ~)
 - Directed flow value depends on the details of the theoretical treatment.
- A New (?) Hope (Episode IV)
 - After fixing momentum-dependent pot. from pA scattering data and explaining v₁ data, EOS dependence of v₂ (elliptic flow) remains ! (Global analysis of multiple observables will help.)
- **How about** Λ ?





Nara, AO (PRC'('22), 2109.07594)

