Symmetry energy in dense matter and its relation to phase boundary

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- Introduction
- QCD phase diagram in asymmetric nuclear matter and black hole formation
- Symmetry energy dependence of dense matter EOS
- Summary



**QCD** Phase diagram and Nuclear Matter EOS

Phase diagram and EOS

= Two important aspects of Nuclear Matter

- Dense nuclear matter has rich physics
  - → Many-body theory, Exotic compositions, CEP, Astrophysical applications, ...



Nuclear matter EOS

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= Subjects in Nuclear, Quark-Hadron, Particle, Astro, and Condensed Matter Physics !

## Symmetry Energy

Recent data suggest that EOS becomes softer in asymmetric nuclear matter.

$$K = K_{sym} + K_{asy} \delta^2$$
,  $K_{asy} \sim -550 \,\text{MeV}$   
 $E_{sym} \simeq 31.6 \,(\rho / \rho_0)^{1.05}$ 

- Isoscalar Giant Monopole Resonance (ISGMR) of Sn isotopes
  - ISGMR in Isotope chain (<sup>112</sup>Sn ~ <sup>124</sup>Sn) is systematically studied.





T. Li, U. Garg, et al., PRC81('10), 034309.



### Symmetry Energy





### **Neutron Star Composition**





**Red Shift** 

**MR curve** 

R

Μ

Ohnishi, Sym. E., Nov.10-12, 2011, YITP, Kyoto, Japan 5

ρ

 $\rho_{B}$ 

Sym. E.

## $1.97 \pm 0.04 M_{\odot}$ Neutron Star



signature. We calculate the pulsar mass to be  $(1.97 \pm 0.04)M_{\odot}$ , which rules out almost all currently proposed<sup>2-5</sup> hyperon or boson condensate equations of state ( $M_{\odot}$ , solar mass). Quark matter can support a star this massive only if the quarks are strongly interacting and are therefore not 'free' quarks<sup>12</sup>.

## Symmetry Energy at High Density

- **Symmetric matter pressure is zero at**  $\rho_0$ .
  - $\rightarrow$  A large part of pressure in neutron star matter at ~  $\rho_0$  comes from symmetry energy.
- Symmetry energy helps to enhance electron chemical potential.
  - → How does it affect to pion condensation and transition to quark matter ?
- In this work, we study the role of symmetry energy at high density from the comparison of

the hadronic EOS (Relativistic Mean Field; RMF) and

chiral effective model (NJL, PNJL, PQM, ...) results.



Let me advertize our recent work on ...





## Quark Matter in Compact Stars



- Probable in Neutron Star (Cold, dense, asymmetric) E.g. N. Glendenning, "Compact Stars"; F. Weber, Prog. Part.Nucl.Phys.54('05)193
- Suggestions in Supernovae (Warm, mildely dense (~1.8 ρ<sub>0</sub>))
   *T. Hatsuda*, MPLA2('87)805; I. Sagert et al., PRL102 ('09) 081101.
- May be in Neutron star merger (Very dense)
- How about Dynamical black hole formation

   → Very Hot (~90 MeV), dense(~5ρ₀), Asym. (Y<sub>p</sub>~ (0.1-0.3))
   M. Liebendorfer et al., ApJS 150('04)263; K. Sumiyoshi, et al., PRL97('06) 091101;
   K.Sumiyoshi, C.Ishizuka, AO, S.Yamada, H.Suzuki, ApJL690('09),L43



#### **Dynamical Black Hole Formation**

Gravitational collapse of heavy (e.g. 40  $M_{\odot}$ ) progenitor  $\rightarrow$  BH

• v heating is not enough  $\rightarrow$  failed supernova

radius [km]



#### Model Details

BH formation calculation

Sumiyoshi, Yamada, Suzuki, Chiba, PRL 97('06)091101.

- v radiation 1D (spherical) Hydrodynamics
- v transport is calculated exactly by solving the Boltzmann eq.
- $\blacklozenge$  Gravitational collapse of 40  $M_{\odot}$  star
- Initial condition: WW95 S.E.Woosley, T.A.Weaver, ApJS 101 ('95) 181
- Shen EOS (npeµ)
- QCD effective models
  - NJL, PNJL, PNJL with 8 quark int., PQM
  - N<sub>f</sub>=2
  - Vector coupling  $\rightarrow G_v/G_s$  (g<sub>v</sub>/g<sub>s</sub> in PQM)=0, 0.2

**Chiral Effective Models** 

- Approaches to Phase Diagram
  - Lattice QCD: Reliable at small μ (μ << T), but has the sign problem at large μ</li>
  - Chiral Effective models: NJL, PNJL, PQM Nambu, Jona-Lasinio ('61), Fukushima('03), Ratti, Thaler, Weise ('06),

B.J.Schafer, Pawlowski, Wambach ('07); Skokov, Friman, E.Nakano, Redlich('10) Spontaneous breaking & restoration of chiral symmetry Polyakov loop extension → Deconf. transitions





### Chiral Effective Models ( $N_f=2$ )

- Lagrangian (PQM, as an example)
  - $L = \overline{q} \Big[ i \gamma^{\mu} \underline{D}_{\mu} g_{\sigma} (\underline{\sigma} + i \gamma_{5} \tau \cdot \pi) \Big] q + \frac{1}{2} \partial^{\mu} \sigma \partial_{\mu} \sigma + \frac{1}{2} \partial^{\mu} \pi \cdot \partial_{\mu} \pi \underbrace{U_{\sigma} (\sigma, \pi)}_{\text{chiral}} \underbrace{U_{\Phi} (\Phi, \overline{\Phi})}_{\text{Polyakov}}$   $F_{\text{eff}} \equiv \Omega / V = U_{\sigma} (\sigma, \pi = 0) + U_{\Phi} (\Phi, \overline{\Phi}) + \underbrace{F_{\text{therm}}}_{\text{therm}} + \underbrace{U_{\text{vac}} (\sigma, \Phi, \overline{\Phi})}_{\text{vac}}$   $particle \text{ exc. } \mathbf{q \text{ zero point}}$ 
    - $U_{\sigma} \sim \phi^4$  theory,  $U_{\phi} \sim -\log$  (Haar Measure)
    - Parameters are fixed by fitting vacuum hadron masses  $(U_{\sigma})$ and lattice data  $(\Phi(T) \rightarrow U_{\Phi})$
  - Vector coupling is not known well  $\rightarrow$  Comparison of  $\mathbf{g}_{\mathbf{v}}/\mathbf{g}_{\mathbf{s}}=0, 0.2$  $L_{V}=-g_{v}\overline{q}\gamma_{\mu}(\omega^{\mu}+\boldsymbol{\tau}\cdot\boldsymbol{R}^{\mu})q-\frac{1}{4}\omega_{\mu\nu}\omega^{\mu\nu}-\frac{1}{4}\boldsymbol{R}_{\mu\nu}\cdot\boldsymbol{R}^{\mu\nu}+\frac{1}{2}m_{\omega}^{2}\omega_{\mu}\omega^{\mu}+\frac{1}{2}m_{\rho}^{2}R_{\mu}R^{\mu}$
  - 8 Fermi interaction *T. Sasaki, Y. Sakai, H. Kouno, M. Yahiro ('10)*  $G_{\sigma 8} \left[ (\bar{q}q)^2 + (\bar{q}i\gamma_5\tau q)^2 \right]^2$



### 3D phase diagram and BH formation

- **Isospin chemical potential**  $\delta \mu = (\mu_d \mu_u)/2 = (\mu_n \mu_p)/2$ 
  - Smaller "Effective" number of flavors  $\rightarrow$  smaller T<sub>CP</sub>
- BH formation process → Quark matter formation & CP sweep
  - Highest  $\mu_B$  just ~ 1300 MeV >  $\mu_c$  (1000-1100 MeV in eff. models)
  - Highest T ~ 90 MeV > T<sub>CP</sub> (at  $\delta\mu$ ~50 MeV)





## Swept Region of Phase Diagram during BH formation

- CP location in Symmetric Matter
  - Lattice QCD μ<sub>CP</sub>=(400-900) MeV
  - Effecitve models
     μ<sub>CP</sub>=(700-1050) MeV
- CP in Asymmetric Matter (E.g. δμ=50 MeV)
  - $T_{CP}$  decreases at finite  $\delta\mu$ .
    - $\rightarrow$  Accessible (T,  $\mu_B$ ) region during BH formation

M.A.Stephanov, Prog.Theor.Phys.Suppl.153 ('04)139; FK02:Z. Fodor, S.D.Katz, JHEP 0203 (2002) 014 LTE:S. Ejiri et al., Prog.Theor.Phys.Suppl. 153 (2004) 118; Can: S. Ejiri, PRD78 (2008) 074507 Stat.:A. Andronic et al., NPA 772('06)167









#### How about Neutron Stars ?

- Neutron Star matter in RMF
  - Solve equilibrium condition at T=0  $\delta\mu = \mu_e/2$ ,  $\rho_c = 0$  (v less, no charge)
  - Various RMFs predict similar δμ values
  - max. δμ ~ 100 MeV
- Phase boundary at T=0 in eff. models
  - First order phase transition disappears at δμ = (60-80) MeV → Possibility of cross over in NS





**Density dependence of Symmetry Energy** 

- RMFs have small ambiguity in Esym. Is it true ?
  - RMF Esym is determined to fit finite nuclear BE, thus reflects average values in the ρ<sub>B</sub> < ρ<sub>0</sub> region.
  - Nuclear effective potential (g-matrix) suggests S-curve behavior of Esym, while RMF gives Esym almost linear in ρ<sub>B</sub>.
    - → RMF may overestimate Esym at high density.







# Effects of Symmetry Energy Change on Phase Transition

- Simple try: Reduce g<sub>ρ</sub> (ρ-N coupling)
   by a factor 0.9.
  - No re-fit of nuclear BE
  - Not connected to low-density (nuclear mixed) EOS (i.e. Do not believe the results.)
- Light neutron star radii sensitively depends on the symmetry energy strength.
  - Pressure at around ρ<sub>0</sub> is dominated by symmetry energy.





## Phase transition with reduced Nuclear Sym. Energy

#### Isospin chemical potential

- Smaller Esym leads to smaller δμ.
- Dependence is not large, but moves in the region of gv/gs=0-0.2.





## **Summary**

- Critical point temperature would decrease at finite isospin chemical potential,  $\delta\mu = (\mu_d \mu_u)/2$ .
  - Quark matter formation and critical point sweep are expected in black hole formation processes.
  - There is also a possibility for the first order phase boundary to disappear in neutron stars because of large δμ.
- Symmetry energy strength at high density is relevant in low-mass neutron star radius (already known) and phase transition.
  - Reduced SymE  $\rightarrow$  Smaller  $\delta\mu$  in NS  $\rightarrow$  Possibility of transition order change.
- There are many subjects to be discussed
  - Construction of Hadron-Quark matter EOS with CP (c.f. J. Steinheimer; D. Blaschke), and its application.
  - S-shape dep. of symmetry energy, 1.97  $M_{\odot}$  neutron star, ...



Thank you for your attention !

#### Collaborators

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