中性子星物理入門 Introduction to Physics of Neutron Stars

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原子核三者若手夏の学校 2015 8/17-22, 2015, ホテルたつき, 蒲郡

- Introducton
- Basics of Neutron Star Physics
- Neutron Star Matter EOS
- Massive NS Puzzle
- Summary







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概要

中性子星は密度、構成要素ともにバラエティに富む多体問題の 宝庫であり、近年の実験・観測の進展により、実験データから示 唆される相互作用の性質と観測データをつき合わせて中性子星 物質状態方程式を定量的に議論できる時代を迎えつつある。 この三者共通講義では、まず中性子星の大まかな性質を概観 した後、近年大きな問題となっている重い中性子星パズル・コン パクトな中性子星パズル・中性子星の冷却・中性子星の強い磁 場などについて解説する。次に状態方程式を理解する上で基本と なる理論の枠組みを解説し、理論・実験・観測による最近の取り 組みを紹介する。



Contents

- Introduction
- Neutron star basics
 - NS mass: Kepler motion, Mass function, and GR effects
 - NS radius: Stephan-Boltzmann, Eddington limit, Red shift
 - A little on NS cooling and magnetic field
- Nuclear matter and neutron star matter EOS
 - Tolman-Oppenheimer-Volkoff (TOV) equation
 - Saturation Point, Incompressibility, and Symmetry Energy
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Crab Nebula SN1054 (e.g. Meigetsu-ki, Teika Fujiwara) Crab pulsar (PSR J0534+2200), discovered in 1968.





Basic properties of neutron stars

- **Mass:** $M = (1-2) M_{\odot} (M \sim 1.4 M_{\odot})$
- Radius: 5 km < R < 20 km (R ~ 10 km)</p>
- Supported by Nuclear Pressure c.f. Electron pressure for white dwarfs
- Cold enough $(T \sim 10^{6} \text{ K} \sim 100 \text{ eV})$ compared with neutron Fermi energy.

чикама ихітиля года уни Курово Калана Кал

Various constituents (conjectured) n, p, e, μ, Y, K, π, q, g, qq,



5

Wide density range \rightarrow various constituents NS = high-energy astrophysical objects and laboratories of dense matter.

Inside Neutron Stars



Dany Page



QCD Phase Diagram





M-R curve and EOS

- M-R curve and NS matter EOS has 1 to 1 correspondence
 - TOV(Tolman-Oppenheimer-Volkoff) equation =GR Hydrostatic Eq.





Puzzles of NS

- Magnetar, NS oscillation,
- Rapid NS cooling puzzle (CasA cools too fast ?)
- Compact NS problem (9 km NS ?)
- Massive NS puzzle (2 M_{\odot} NS ?)



Gravitational Collapse of Massive Star



Binary Neutron Star Mergers and Nucleosynthesis





Dynamical Black Hole Formation

- Gravitational collapse of heavy (e.g. 40 M_o) progenitor would lead to BH formation.
 - Shock stalls, and heating by v is not enough to take over strong accretion. → failed supernova
 - v emission time ~ (1-2) sec w/o exotic matter.
 - emission time is shortened by exotic dof (quarks, hyperons, pions).



Sumiyoshi, Yamada, Suzuki, Sumiyoshi,Ishizuka, AO, Yamada, Nakazato, Sumiyoshi, Chiba, PRL 97('06)091101. Suzuki, ApJL 690('09)43. Yamada, PRD77('08)103006



Binary Neutron Star Merger

T ~ 40 MeV, $\rho_B \sim 10^{15}$ g/cm³ ~ 4 ρ_0 ($\rho_0 \sim 2.5 \text{ x } 10^{14}$ g/cm³), Ye ~ 0.1



Courtesy of K. Kiuchi Data are from Y. Sekiguchi, K. Kiuchi, K. Kyotoku, M. Shibata, PRD91('15)064059.



Physics Opportunities in Neutron Stars

- Equation of state of dense matter
 - Laboratory of exotic constituents
 - Laboratory of QCD phase transition at high density
- Equation of state of isospin asymmetric matter
 - Symmetry energy connect laboratory exp. and astronomical obs.
 - Baryon superfluidity above nuclear density
 - Realization of unitary gas, which can be simulated by cold atoms
- Compact astrophysical objects, whose structure is yet unknown
 - Challenge to measure mass, radius, temperature, magnetic field, ...
- Promising site of gravitational wave source
- Promising site of r-process nucleosynthesis
- Examination of general relativity
- Neutrino emission determines the cooling of NSs.



科研費新学術領域の複数がコンパクト天体に関連

- 重力波天体 領域代表:中村卓(京大)
- 地下素核研究 領域代表:井上邦雄 (東北大
- 中性子星核物質 領域代表:田村裕和(東北大)
- ニュートリノフロンティア 領域代表:中家剛(京大)

Joint symposium by three innovative areas: Gravitational Wave Source / Underground Particle-Nuclear Research / Neutron Star Matter "Universe and Astronomical Objects

Uncovered by Multi-Fold Approach"

新学術3領域(重力波天体・地下素核研究・中性子星核物質) 合同シンポジウム

「多面的アプローチで解きあかす宇宙と天体」

July 24(Fri.) 13:00 ~ 25(Sat.) 15:30, 2015 Aoba Science Hall (2F in "Science Complex C"), Graduate School of Science, Tohoku University

2015年7月24日(金)13:00 ~ 25日(土)15:30 東北大学理学研究科 青葉サイエンスホール(理学合同C棟2F)

http://lambda.phys.tohoku.ac.jp/nstar/symposium/three-areas

力波天体の多様な観測による宇宙物理学の新展開 ew development in astrophysics.through multimessenger observations of gravitational wave sources

宇宙の歴史をひもとく地下素粒子原子核研究 文部科学省研究費補助金新学術領域領域番号2603(平成26年~30年度)



atter in neutron stars investigated by experiments and astronomical observation

NS matter Grant-in-Aid Study in Japan(2012-)



Accelerators and Satellites for Neutron Star Physics





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Mass & Radius Measurements of Neutron Stars



Neutron Star Observables: Mass (1)

- Please remember Kepler motion basics
 - major axis=a, eccentricity=e, reduced mass=m, total mass=M



$$E/m = \frac{1}{2}v_f^2 - \frac{GM}{a(1+e)} = \frac{1}{2}v_n^2 - \frac{GM}{a(1-e)}$$

$$L = mv_f a(1+e) = mv_n a(1-e)$$

$$\Rightarrow v_f^2 = \frac{GM}{a} \frac{1-e}{1+e}, L = 2m\frac{dS}{dt} = m\sqrt{GMa(1-e^2)}$$

$$\Rightarrow P = S/(dS/dt) = 2\pi a^2 \sqrt{1-e^2}/\sqrt{GMa(1-e^2)} = 2\pi a^{3/2}/\sqrt{GM}$$



Neutron Star Observables: Mass (2)

- Binary stars
 - inclination angle = i
 - Doppler shift (Pulse timing change) is given by the radial velocity (視線速度)
 K = v sin i

 - Mass function (observable)

$$f = \frac{(M_2 \sin i)^3}{M^2} = \frac{4\pi^2 (a_1 \sin i)^3}{G} P^2$$
$$= \frac{K^3 P (1 - e^2)^{3/2}}{2\pi G}$$
$$(K = v \sin i, M = M_1 + M_2)$$

and GR effects ...





Hulse-Taylor Pulsar (PSR 1913+16)

Precisely (and firstly) measured neutron star binary (1993 Nobel prize to Hulse & Taylor)

a Radial velocity \rightarrow P, e, K \rightarrow Mass function





More on Hulse-Taylor Pulsar (PSR 1913+16)

- General Relativistic Effects
 - Perihelion shift (近日点移動)

$$\dot{\omega} = 3 \left(\frac{2\pi}{P} \right)^{5/3} \frac{(GM)^{2/3}}{(1-e^2)c^2}$$

Einstein delay

$$\Delta_E = \gamma \sin u$$

(u=eccentric anomaly)

$$v = \frac{eP_b Gm_2(m_1 + 2m_2)}{2\pi c^2 a_{\rm R} M} \qquad \frac{a_{\rm R}^3}{P_b^2} = \frac{GM}{4\pi^2} \left[1 + \left(\frac{m_1 m_2}{M^2} - 9\right) \frac{GM}{2a_{\rm R} c^2} \right]^2$$

Two observable

 \rightarrow Precise measurement of m₁ and m₂.

$$m_1 = 1.442 \pm 0.003 M_{sun}$$

 $m_2 = 1.386 \pm 0.003 M_{sun}$





Massive Neutron Star

- General Relativity Effects on Time Delay
 - Einstein delay : varying grav. red shift
 - Shapiro delay : companion's grav. field
- A massive neutron star (J1614-2230)
 - M = 1.97 \pm 0.04 M_{\odot} is obtained using the Shapiro delay Demorest et al. (2010)





Neutron Star Masses

- NS masses in NS binaries can be measured precisely by using some of GR effects.
 - Perihelion shift+Einstein delay
 → M = 1.442 ± 0.003 M_☉
 (Hulse-Taylor pulsar)
 Taylor, Weisenberg ('89)
 - Shapiro delay $\rightarrow M = 1.97 \pm 0.04 M_{\odot}$

Demorest et al. ('10)

• Another obs.: $M = 2.01 \pm 0.04 M_{\odot}$ Antoniadis et al. ('13)

Neutron Star Mass $M = (1-2) M_{\odot}$ Canonical value = 1.4 M_{\odot}





Neutron Star Radius

- How can we measure 10 km radius of a star with 10-100 thousands light year distance from us ?
 - Size of galaxy ~ 3×10^{14} km (~ $10 \text{ kpc} \sim 3 \times 10^{4}$ light year)
 - \rightarrow Model analysis is necessary !
- X-ray burster
 - Mass accretion from companion occasionally induces explosive hydrogen / helium burning.
 - High temperature \rightarrow NS becomes bright !
 - Three methods to measure NS radius





NS Radius Measurement (1)

- Surface emission
 - Stefan-Boltzmann law is assumed
 → NS radius is obtained
 from Flux, Temperature,
 and Distance measurement.

$$L = 4 \pi R_{\infty}^{2} \sigma_{\rm SB} T^{4} , \quad F = \frac{L}{4 \pi D^{2}}$$

$$\Rightarrow R = \sqrt{\frac{F D^{2}}{\sigma_{\rm SB} T^{4}}} \left(1 - \frac{2 G M}{R c^{2}}\right)^{-1/2}$$





NS Radius Measurement (2)

- Eddington Limit
 - Eddington Limit radiation pressure = gravity

$$\frac{4\pi r^2 \sigma_{\rm SB} T^4}{4\pi r^2 c} \cdot N_e \cdot \sigma_{\rm T}$$
$$= \frac{GM}{r^2} \cdot N_N \cdot m_N$$
$$\rightarrow R_{\infty}^2 = \frac{2GMcm_N}{\sigma_{\rm T}\sigma_{\rm SB} T^4} \frac{N_N}{N_e}$$

- Eddington limit is assumed to be achieved at "touch down".
- Electron-nucleon ratio N_e/N_N=(1+X)/2 (X=1 for hydrogen atmosphere X=0 for light elements)





NS Radius Measurement (3)

Red Shift

- Neutron Star surface is expected to contain Irons.
- Absorption lines should be red shifted.
 → Almost direct observation of M/R.

$$E_{\rm obs} = E_{\rm surf} \sqrt{1 - \frac{2GM}{Rc^2}}$$

 ASTRO-H will measure Iron absorption line from NS, and determine M/R with 1 % accuracy !





ASTRO-H simulation



Neutron Star Radius

(1) 0-8 s

(2) 8-16 s

 10^{2}

- Do three methods give consistent (M, R) ?
 - Surface emission & Eddington limit have large error bars from Distance & Composition uncertainty.
 - Red shift of discrete lines have not been observed unambiguously.



Compact NS puzzle

Some analyses suggest smaller Guillot et al. (2013) **R_{NS}** than nucl. phys. predictions. 2.5 MPA1 PAL1 Some make objections. WFF1 ([©]W) ^{2.0} ^{(©}W) ^{SN}W ^{1.5} Suleimanov+, $R_{14} > 13.9$ km MS1 Lattimer+, $R_{14} = 12 \pm 1.4$ km 1.00.5 3.0F 4U1608-52 10 12 6 8 14 16 EX01745-248 *F. Ozel, ('13).* $R_{\rm NS}$ (km) 4U1820-30 <S1731-260 2.5SAXJ1748.9-2100 MPA1 Base, N₁₁ (D90), Dist (G13), H+He GS1826-24 0.9 U24 in NGC6397 AP4 M13 2.0 0.8 NGC2808 Mass (M_☉) ω Cen 0.7 1.5 0.6 MS1 (^oM (M_o) 0.5 1.0 GS1 0.4 0.3 0.5 0.2 SQM1 0.5 0.1 0.0 ± 13 14 15 16 11 12 5 10 15 0 ^R (kn Lattimer, Steiner (2014). Radius (km) 31 A. Ohnishi (a) YONUPA, Aug. 17, 2015 Matt

Neutron Star Density



Neutron Stars are supported by Nuclear Force !

- Average density of NS ~ (1-3) ρ_0 , Max. density ~ (5-10) ρ_0
 - → Supported by Nuclear Force
 - c.f. White Dwarfs are supported by electron pressure.
- Nuclear Force
 - Long-range part: π exchange Yukawa (1935)
 - Medium-range attraction:
 2 π exchange, σ exchange,
 Nambu, Jona-Lasinio (1961)
 - Short-range repulsion: Vector meson exchange, Pauli blocking btw. quarks Gluon exchange

Neudatchin, Smirnov, Tamagaki; Oka, Yazaki; Aoki, Hatsuda, Ishii









A little on NS cooling & Magnetic Field



Neutron Star Cooling

- Direct URCA process
 - $n \rightarrow p + e^- + \bar{\nu}_e \ , \ \ e^- + p \rightarrow n + \nu_e$
 - Dominant at high T (T>10⁹ K)
 - Suppressed at low T (T < 10⁹ K)
- Modified URCA process

 $n+n \rightarrow n+p+e^- + \bar{\nu}_e \ , \ \ n+p+e^- \rightarrow n+n+\nu_e$

- Standard" cooling process of young NS (t < 10⁴ yrs, T > 10⁸ K)
- Non-standard cooling processes
 - Y-URCA

 $Y \rightarrow N + e^- + \bar{\nu}_e \ , \ \ e^- + Y \rightarrow N + \nu_e$

• π cooling

 $\pi^- + n \rightarrow n + e^- + \bar{\nu}_e$, $n + e^- \rightarrow n + \pi^- + \nu_e$

quark beta decay

$$d \to u + e^- + \bar{\nu}_e$$
, $u + e^- \to d + \nu_e$

Casino de Urca @ Rio

Direct URCA suppression

- **D-URCA is suppressed at Y**_p < 0.11
 - Equilibrium condition: $\mu_n = \mu_p + \mu_e$ $\frac{P_F^2(n)}{2M_n} + M_n + U_n = \frac{P_F^2(p)}{2M_p} + M_p + U_p + P_F(e)$
 - Charge neutrality: P_F(p)=P_F(e)
 - Momentum conservation for zero momentum v emission

$$P_F(n) = 2P_F(p) \rightarrow Y_p = Z/(N+Z) = 1/9 = 0.12$$

• Y-DURCA and q-DURCA is free from suppression M-URCA is slow

$$\Gamma = \frac{(2\pi)^4}{\hbar V} \int \delta(E_f - E_i) \delta^3(\mathbf{p}_f - \mathbf{p}_i) |H_{fi}|^2 f_1 f_2 (1 - f_1') (1 - f_p) (1 - f_e) \prod_i V \frac{d^3 p_i}{(2\pi)^3}$$
$$L_{\nu}^{\text{mURCA}} = C \frac{M}{M_{\odot}} \left(\frac{\rho_0}{\rho}\right)^{1/3} \left(\frac{T}{10^9 \text{ K}}\right)^8 \quad (C \simeq (0.8 - 5) \times 10^{39} \text{erg/s})$$
Shapiro texbook



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n

 $Y_p < 1/9$

P(p) P(e)

P(n)

P(p) P(e)

P(n)
Neutron Star Cooling (cont.)

- Many of neutron star temperature observations are consistent with "standard" modified URCA cooling (with some heating).
- Some require faster cooling. Need some exotics.
- Exotic cooling is too fast if there is no suppression mechanism. Superfluidity is a promising candidate.



, EVOLUTION OF HYPERON-MIXED NEUTRON STARS



S. Tsuruta, Grossmann Medalist, 2015



Nuclear Superfluidity and Cooling Curve

- Surface T measurement and Cooling curve
 - Stable superfluid \rightarrow Gap \rightarrow Suppression of v emission

 - Precise T and Cooling rate measurement in Cas A Heinke, Ho, ApJ 719('10) L167 [arXiv:1007.4719] Page et al., PRL 106 ('11) 081101 [arXiv:1011.6142]



Magnetic Field

- Magnetic Dipole Model (cf. Shapiro, Teukolsky)
 - Magnetic Dipole Moment

$$|\mathbf{m}| = \frac{1}{2} B_p R^3 ,$$

$$\dot{E} = -\frac{2}{3c^3} |\ddot{\mathbf{m}}|^2 = -\frac{B_p^2 R^6 \Omega^4 \sin^2 \alpha}{6c^3}$$

Rotation Energy of NS

$$\begin{split} E = &\frac{1}{2}I\Omega^2 , \quad \dot{E} = I\Omega\dot{\Omega} , \\ T \equiv &-\left(\frac{\Omega}{\dot{\Omega}}\right)_0 = \frac{6Ic^3}{B_p^2 R^6 \sin^2 \alpha \Omega_0^2} , \\ \text{age} : t \simeq T/2 \end{split}$$

- **Magnetic field in NS** $B = 10^{12} 10^{15} G$
 - From P and dP/dt, we can guess B and t (age) of NS



Ho, Klus, Coe, Andersson ('13)





Origin of Strong Magnetic Field

- How can we make strong B ? cf. H. C. Spruit, AIP Conf.Proc.983('08)391.
 - Fossil field hypothesis (化石磁場) (flux conservation)
 - Dynamo process in progenitor star evolution
 - Ferromagnetism e.g. Yoshiike, Nishiyama, Tatsumi ('15)
- How can we keep strong B ?
 - Dipole magnetic field is not stable Flowers, Ruderman ('77)
 - Finite magnetic helicity $\mathcal{H} = \int dx A \cdot B$

makes magnetic field stable. Prendergast ('56); AO, N. Yamamoto, arXiv:1402.4760; D. Grabowska, D. B. Kaplan, S. Reddy, PRD('15)085035.







Flowers, Ruderman ('77)



Chiral Plasma Instability ?

Chiral Plasma Instability

AO, N. Yamamoto, arXiv:1402.4760

 ■ Left-handed electrons are eaten in electron capture → chiral chem. pot.

$$p + e_L^- \rightarrow n + \nu_{L^+}^e$$

 Chiral plasma instability: N₅ is converted to magnetic helicity Akamatsu, Yamamoto ('13, '14)

$$j_z = \frac{2\alpha}{\pi} \mu_5 B_z, \qquad \frac{d}{dt} \left(N_5 + \frac{\alpha}{\pi} \mathcal{H} \right) = 0, \quad N_5 = \int d\boldsymbol{x} \, n_5$$

Finite magnetic helicity makes magnetic field stable.

$$\mathcal{H} = \int dx A \cdot B$$

Electron Mass may kill the instability.

D. Grabowska, D. B. Kaplan, S. Reddy, PRD('15)085035



(a)

(b)

(c)

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TOV equation

General Relativistic Hydrostatic Equation
 = TOV(Tolman-Oppenheimer-Volkoff) equation

$$\frac{dP}{dr} = -G \frac{(\varepsilon/c^2 + P/c^2)(M + 4\pi r^3 P/c^2)}{r^2(1 - 2GM/rc^2)}$$
$$\frac{dM}{dr} = 4\pi r^2 \varepsilon/c^2, \ P = P(\varepsilon) \ (EOS)$$

- Spherical and non-rotating.
- 3 Variables (ε(r), P(r), M(r)),
 3 Equations.
- Initial cond. ε(r=0)
 Solve TOV until P=0





M-R Relation and EOS

Solving TOV eq.

starting from the "initial" condition, $\varepsilon(r=0) = \varepsilon_c =$ given until the "boundary" condition P(r)=0 is satisfied.

 $\rightarrow M \text{ and } R \text{ are the functions of } \epsilon(r=0)$ and functionals of EOS, $P=P(\epsilon)$.

$$M = M(\varepsilon_c)[P(\varepsilon)]$$
, $R = R(\varepsilon_c)[P(\varepsilon)]$

→ M-R curve and NS matter EOS : 1 to 1 correspondence





A. Ohnishi @ YONUPA, Aug.17, 2015 45

Nuclear Mass

Bethe-Weizsacker mass formula Nuclear binding energy is roughly given by Liquid drop. Nuclear size measurement $\rightarrow R = r_0 A^{1/3}$

$$B(A,Z) = a_v A - a_s A^{2/3} - a_C \frac{Z^2}{A^{1/3}} - a_a \frac{(N-Z)^2}{A} + a_p \frac{\delta_p}{A^2}$$

Volume Surface Coulomb Symmetry

 $A \propto \frac{4\pi}{3} R^3 \qquad A^{2/3} \propto 4\pi R^2 \qquad \propto \frac{Q^2}{R}$ **Ignore Coulomb, consider A** $\rightarrow \infty$,

$$B/A = a_v(\rho) - a_a(\rho)\delta^2$$
, $\delta = (N - Z)/A$
 $a_v \simeq 16 \text{ MeV}$
 $a_a \simeq 23 \text{ MeV}(a_a(\text{vol}) \simeq 30 \text{ MeV})$

Coef. may depend on the number density $\rho \rightarrow$ Nuclear Matter EOS



 $\mathbf{R} \propto \mathbf{A} \frac{1}{3}$

Pairing

Neutron Star Matter EOS

Energy per nucleon in nuclear matter

$$E_{\rm NM}(\rho, \delta) = E_{\rm SNM}(\rho) + S(\rho)\delta^2 , \quad \delta = (N - Z)/A$$
$$E_{\rm SNM}(\rho) \simeq E_0 + \frac{K(\rho - \rho_0)^2}{18\rho_0^2} , \quad S(\rho) = S_0 + \frac{L(\rho - \rho_0)}{3\rho_0}$$

- Saturation point $(\rho_0, E_0) \sim (0.16 \text{ fm}^{-3}, -16 \text{ MeV})$
- Symmetry energy parameters (S $_0$ (=J), L) ~ (30 MeV, 70 MeV)
- Incompressibility K ~ 230 MeV
- Uniform neutron star matter
 - Constituents at low density
 = proton, neutron and electron

$$E_{\rm NSM}(\rho) = E_{\rm NM}(\rho, \delta) + E_e(\rho_e = \rho_p)$$

• Charge neutrality $\rightarrow \rho(\text{elec.}) = \rho(p) (\rho_e = \rho_p = \rho(1 - \delta)/2)$ δ is optimized to minimize energy.







Symmetry Energy

- Symmetry Energy has been extracted from various observations.
 - Mass formula, Isobaric Analog State, Pygmy Dipole Resonance, Isospin Diffusion, Neutron Skin thickness, Dipole Polarizability, Asteroseismology



Simple parametrized EOS

Skyrme int. motivated parameterization

$$E_{\rm SNM} = \frac{3}{5} E_F(\rho) + \frac{\alpha}{2} \left(\frac{\rho}{\rho_0}\right) + \frac{\beta}{2+\gamma} \left(\frac{\rho}{\rho_0}\right)^{1+\gamma}$$
$$S(\rho) = \frac{1}{3} E_F(\rho) + \left[S_0 - \frac{1}{3} E_F(\rho_0)\right] \left(\frac{\rho}{\rho_0}\right)^{\gamma_{\rm sym}}$$

• $(\rho_0, E/A(\rho_0), K) \rightarrow (\alpha, \beta, \gamma), L \rightarrow \gamma_{sym}$ K=220 MeV, S₀=30 MeV





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Simple parametrized EOS





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Theories/Models for Nuclear Matter EOS

- Mean Field from Effective Int. ~ Nuclear Density Functionals
 - Skyrme Hartree-Fock

Non.-Rel.,Zero Range, Two-body + Three-body (or ρ-dep. two-body)

$$\frac{E}{A} = \left\langle \frac{\mathbf{p}^2}{2m^*} \right\rangle + V(\rho, \delta) , \quad V \simeq \frac{\alpha}{2} \frac{\rho}{\rho_0} + \frac{\alpha' \delta}{2} \frac{\rho}{\rho_0} + \frac{\beta}{1+\gamma} \left(\frac{\rho}{\rho_0}\right)^{\gamma} + \dots$$

- Relativistic Mean Field
 - Relativistic, Meson-Baryon coupling, Meson self-energies

$$\frac{E}{A} = \left\langle \sqrt{\mathbf{p}^2 + (M - g_\sigma \sigma)^2} \right\rangle + g_\omega \omega + \frac{1}{\rho_B} \left[\frac{1}{2} m_\sigma^2 \sigma^2 - \frac{1}{2} m_\omega^2 \omega^2 + \dots \right]$$

- Microscopic (ab initio) Approaches (starting from bare NN int.)
 - Variational calculation
 - Quantum Monte-Carlo
 - Bruckner Theory (G-matrix)



Mean Field models

- Fit parameters to nuclear properties (B.E., radius, ...) → predict neutron star (M,R).
 - In Non-Rel. treatment with SLy (std. parametrization), FPS (impr.) → Mmax ~ (1.8-2.0) M_☉
 - Rel. MF (TM1) \rightarrow Mmax \sim 2.2 M $_{\odot}$



Variational Calculation

Variational Calculation starting from bare nuclear force

B. Friedman, V.R. Pandharipande, NPA361('81)502; A. Akmal, V.R.Pandharipande, D.G. Ravenhall, PRC58('98)1804; H. Kanzawa, K. Oyamatsu, K. Sumiyoshi, M. Takano, NPA791 ('07) 232.

Argonne v18(v14) + Rel. corr. + Three Nucleon Int.

Star Mat





Quantum Monte-Carlo calc.

- Auxiliary Field Diffusion Monte-Carlo (AFDMC) calc.
 - Hubbard-Stratonovich transf. + MC integral over aux. fields.
 - In force parameters are tuned to fit finite nuclei.
 - 2 MeV Difference in Esym results in 1.5 km (15 %) diff. in R_{NS}.





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Bruckner-Hartree-Fock

- Effective interaction from bare NN int. (G-matrix).
 - G-matrix = Lowest order Bruckner theory, but next-to-leading terms give small effects at ρ < 4 ρ₀.

Song, Baldo, Giansiracusa, Lombardo ('98)

Need 3-body force to reproduce saturation point.

 $g(E) = V + V \underbrace{\simeq}_{E - H_0}$



Z.H.Li, U. Lombardo, H.-J. Schulze, W. Zuo, L. W. Chen, H. R. Ma, PRC74('06)047304.



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 $\epsilon_{\rm p} > \epsilon_{\rm F}$

- □— V14: Argonne V14 - ☆— V18: Argonne V18 - ●— A: Bonn A

-o-PAR: Paris

B: Bonn B
 C: Bonn C
 CD: CD-Bonn

-▲— R93: Reid93 -▼— N93: Nijmegen93 -◀— NI: Nijmegen I

– ►– NII: Nijmegen II

50

20

BHF with Ch-EFT & Lattice NN force

- Bruckner-HF calc. with NN (N3LO)+3NF(N2LO) interactions from Chiral Effective Field Theory *M.Kohno (*13)*
 - Ch-EFT = Eff. Field Theory with the same symmetry as QCD *Weinberg; Gasser, Leutwyler ('84)*
 - → Systematically gives NN & NNN interaction terms.

Epelbaum, Gockle, Meissner ('05)

- Bruckner HF calc. with NN int. from Lattice QCD. *Inoue et al. (HAL QCD Coll.), PRL111 ('13)112503*
 - Not yet reliable but promising !



M. Kohno, PRC88('13)064005



Contents

- Introduction
- Neutron star basics
 - NS mass: Kepler motion, Mass function, and GR effects
 - NS radius: Stephan-Boltzmann, Eddington limit, Red shift
 - A little on NS cooling and magnetic field
- Nuclear matter and neutron star matter EOS
 - Tolman-Oppenheimer-Volkoff (TOV) equation
 - Saturation Point, Incompressibility, and Symmetry Energy
- Massive neutron star puzzle
 - How can we sustain two-solar-mass NSs ?
 - Proposed mechanisms to sustain massive NSs
 - What is necessary to solve massive NS puzzle ?



Summary





Neutron star – Is it made of neutrons ?

Possibilities of various constituents in neutron star core







Massive Neutron Star

- General Relativity Effects on Time Delay
 - Einstein delay : varying grav. red shift
 - Shapiro delay : companion's grav. field
- A massive neutron star (J1614-2230)
 - M = 1.97 \pm 0.04 M_{\odot} is obtained using the Shapiro delay Demorest et al. (2010)





Massive Neutron Star Puzzle



61 A. Ohnishi (a) YONUPA, Aug. 17, 2015

Hyperons in Dense Matter

- What appears at high density ?
 - Nucleon superfluid (³S₁, ³P₂), Pion condensation, Kaon condensation, Baryon Rich QGP, Color SuperConductor (CSC), Quarkyonic Matter,

Hyperons

Tsuruta, Cameron (66); Langer, Rosen (70); Pandharipande (71); Itoh(75); Glendenning; Weber, Weigel; Sugahara, Toki; Schaffner, Mishustin; Balberg, Gal; Baldo et al.; Vidana et al.; Nishizaki,Yamamoto, Takatsuka; Kohno,Fujiwara et al.; Sahu,Ohnishi; Ishizuka, Ohnishi, Sumiyoshi, Yamada; ...





NS matter EOS with hyperons



These are phenomenological "solutions". How can we examine them ?



Star Matte

Possible Solutions to Massive NS puzzle

- Proposed "Solutions" of Massive NS puzzle
 - Choose Stiff EOS for nuclear matter Tsubakihara, Harada, AO ('14)
 - Modification of YN interaction Weisenborn, Chatterjee, Schaffner-Bielich ('11); Jiang, Li, Chen ('12); Tsubakihara, AO ('13)
 - Introducing BBB repulsion S. Nishizaki, T. Takatsuka, Y. Yamamoto ('02); Bednarek, Haensel et al.('11); Miyatsu, Yamamuro, Nakazato ('13); Tamagaki ('08). Togashi, Hiyama, Takano, Yamamoto; Nakamoto, Suzuki;
 - Early transition to quark matter Masuda, Hatsuda, Takatsuka ('12)
- What is necessary to solve the massive NS puzzle ?
 - EOS of symmetric nuclear matter at high density
 - Symmetry Energy at supra nuclear density.
 - Yet un-explored YN & YY interactions
 - Three-body interaction including hyperons (YNN, YYN, YYY) and its effects on EOS
 - Finding onset density of quark matter



NNN force

- NNN force is necessary to reproduce saturation point and to support massive neutron stars
 - Variational cal. + phen. NNN force

 A. Akmal, V.R.Pandharipande, D.G. Ravenhall, PRC58('98)1804;
 H. Kanzawa, K. Oyamatsu, K. Sumiyoshi, M. Takano, NPA791 ('07) 232.
 - Chiral EFT NN+NNN force M. Kohno, PRC88('13)064005





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NNN force from Lattice QCD

- HAL QCD method for BB int.
 Aoki, Hatsuda, Ishii ('07)
 Nambu-Bethe-Salpeter amplitude ~ w.f.
 → NN force from Sch. Eq.
 - Consistent with Luscher's method in asymptotic region Luscher ('91), NPLQCD Collab. ('06, ππ)
- NNN force T. Doi (HAL QCD Collab.)('12)





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Hyperons & YY interaction

- Hyperons are expected to appear in NS and soften EOS.
 - Hypernuclear data \rightarrow max. NS mass reduction of (0.5-1.0) M_{\odot}
 - Nagara event (ΛΛ nuclei) and heavy-ion collisions (ΛΛ correlation) implies ΛΛ int. is weakly attractive.



BBB force including Hyperons

Repulsive BBB int. incl. Y is necessary to support 2 M_o NS.

"Universal" BBB force

Nishizaki, Takatsuka, Yamamoto ('02), Yamamoto, Furumoto, Yasutake, Rijken('13)

Variational calc. including hyperons

Togashi et al. (in prep.)



Early crossover transition to quark matter

- Early crossover to quark matter → massive NS K. Masuda, T. Hatsuda, T. Takatsuka, ApJ764('13)12
- QCD phase diagram in asymmetric matter AO et al. ('11), Ueda et al. ('13)
 - Disappearance of 1st order phase transition at large isospin chem. pot.



Summary

- 中性子星は「極限状況の物質」物理の宝庫である。
 - 高密度、アイソスピン非対称、超流動、エキゾチックな構成要素
- 中性子星物質状態方程式の研究が活発に行われている。
 - RI加速器施設 (RIBF, FRIB, SPIRAL, RAON, ...)、
 ハドロン加速器 (J-PARC, JLAB, ...)、
 重イオン衝突型加速器 (RHIC, LHC, NICA, FAIR, J-PARC, ...)
 - 人工衛星による観測 (ASTRO-H, LOFT, NICER, ...)
 - 理論研究(量子モンテカルロ、カイラル EFT、格子 QCD、 有効相互作用、...)
- 現在、中性子星にまつわる複数のパズルが存在
 - 重い中性子星パズル、中性子星半径の謎、急速な冷却、 強い磁場の起源、....
 - 重い中性子星パズル:ハイペロンを含む3体力?クォーク物質?



Thank you for your attention !



 $Q: なぜクォーク物質では M \rightarrow 0 で R \rightarrow 0 ?$

- Ans: Self-bound するから。
 - クォーク物質では u:d:s=1:1:1で 電気的に中性。
 - 準安定な密度が存在すると、圧力は0

$$P = \rho_B^2 \frac{\partial (E/A)}{\partial \rho_B}$$

表面で準安定な密度と真空が接触。




Birth, Life and Death of Matter in Our Universe





Chiral EFT NN & NNN force



E. Epelbaum ('09)



中性子物質と冷却原子

- **BEC-BCS crossover and unitary gas**
 - 散乱長 >> 粒子間距離 → EOS は普遍的 (unitary gas)
 $E^{\text{Unitary}} = \xi E^{\text{Free}} \quad \xi \simeq 0.4 (\text{Bertsch parameter})$ nn 間の ¹S₀ 散乱長は長い! (a₀=-18.5 fm)
 - → Drip した中性子ガスは、ほぼ unitary gas (-1/ $k_{F}a_{0} \sim 0.1$)

My question

- 核子あたりの相互作用エネルギー $\infty k_F^2 \propto \rho^{2/3}$ $\frac{V^{\text{Unitary}}}{N} = (\xi - 1) \frac{3}{5} \frac{\hbar^2 k_F^2}{2m} \propto \rho^{2/3}$
- どのようにして EOS(密度汎関数) 取り込むか? (Hartree なら ∞ ρ)
- unitary gas / BEC-BCS crossover は ^{bcs limit} クラスト・原子核の性質に どのような影響を及ぼすか?



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 $+\infty$

中性子星物質の状態方程式

- 変分法による計算結果 Friedman-Pandharipande (1981)
 - 広い密度領域において E_{unit} < E_{FP} < E_{Fermi}
 - 低密度領域でポテンシャルエネ ルギーは ρ^{2/3} と振る舞ってい るか?







What is necessary to solve the massive NS puzzle ?

- There are many "model" solutions.
- Ab initio calculation including three-baryon force (3BF)
 - Bare 2NF+Phen. 3NF(UIX, IL2-7) + many-body theory (verified in light nuclei).
 - Chiral EFT (2NF+3NF) + many-body theory



J. Carlson et al. ('14)

"Universal" mechanism of "Three-body" repulsion

- Mechanism of "Universal" Three-Baryon Repulsion.
 - "σ"-exchange ~ two pion exch. w/ res.
 - Large attraction from two pion exchange is suppressed by the Pauli blocking in the intermediate stage.

• Coupling to Res. (hidden DOF)

• Reduced " σ " exch. pot. ?

How about YNN or YYN?





"Universal" TBR

AA interaction in vacuum and in nuclear medium

- Vacuum ΛΛ interaction may be theoretically accessible Lattice QCD calc. HAL QCD ('11) & NPLQCD ('11)
- In-medium ΛΛ interaction may be experimentally accessible
 - a_0 (Nagara fit) = 0.575 fm, -0.77 fm ($\Delta B_{\Lambda\Lambda}$ =1.0 MeV) *Hiyama et al. ('02), Filikhin, Gal ('02)*
 - Bond energy of ${}^{6}_{\Lambda\Lambda}$ He: $\Delta B_{\Lambda\Lambda}$ =1.0 MeV \rightarrow 0.6 MeV Nakazawa, Takahashi ('10)
- Difference of vacuum & in-medium
 ΛΛ int. would inform us ΛΛΝ int. effects.
 - ΛΛ-ΞΝ couples in vacuum
 - Coupling is suppressed in ${}^{6}_{\Lambda\Lambda}$ He

Is there Any way to access "vacuum" AA int. experimentally ?





Exotic Hadrons

Exotic hadrons

 \rightarrow X, Y, Z, Θ^+ , Discovered/Proposed at LEPS, Belle, BaBar,...



- Various pictures
 - Di-quark component
 - Hadronic molecule
 - $Q\overline{Q}$ couples with $Q\overline{Q}$ $q\overline{q}$





Ab initio EOS fit + Hyperons in RMF with multi-body couplings



Alternative approach

- Alternative method
 ~ "Ab initio" Nucl. Matter EOS + Y phen.
 - Fit "Ab initio" EOSs in a phen. model,
 - Include hyperons, and explain hypernuclear data.



Tsubakihara et al., PRC81('10)065206 Tsubakihara, Harada, AO, arXiv:1402.0979

We fit ab initio EOS in RMF with multi-body couplings, and introduce hyperons.

82

"Ab initio" EOS

- "Ab initio" EOS under consideration
 - FP: Variational calc. (Av14+3NF(att.+repl.))
 B. Friedman, V.R. Pandharipande, NPA361('81)502.
 - APR: Variational chain summation (Av18+rel. corr. ; Av18+ rel. corr.+3NF)
 A. Akmal, V.R.Pandharipande, D.G. Ravenhall, PRC58('98)1804.
 - DBHF: Dirac Bruckner approach (Bonn A)
 G. Q. Li, R. Machleidt,
 R. Brockmann,
 PRC45('92)2782





n=2 and n=3 terms in RMF

■ n=B/2+M+D=2 RMF model (+ effective pot.) →2-body interaction (and rel. 3-body corr.)

■ n=3 model → 3-body coupling

$$g_{mm'B}\overline{\Psi}mm'\Psi$$

$$C_{mm'm'}mm'm''$$

Bmm terms are ignored in FST paper (field redefinitions).



Tsubakihara

Fitting "Ab initio" EOS via RMF

RMF with multi-body couplings: 15 parameters

- Working hypothesis σ self-energy: SCL2 model *Tsubakihara, AO ('07)* $M_N \rightarrow 0 @ \sigma \rightarrow f_{\pi}$
- Markov Chain Monte-Carlo (MCMC)-like parameter search
 - Langevin type shift +Metropolis judge
 - Simultaneous fit of SNM and PNM is essential.
 - std. dev=0.5-0.7 MeV





Symmetry Energy

- Symmetry E. = E(PNM)-E(SNM)
 - APR-fit: (S₀, L)=(32, 47) MeV
 - APRv2-fit: (S₀, L)=(33, 47) MeV
 - DBHF-fit: (S₀, L)=(35, 75) MeV
 - FP-fit: (S₀, L)=(32, 40) MeV



Star Mat



Neutron Star Matter EOS

- Asymmetric Nuclear Matter EOS
 E_{ANM}(ρ)=E_{SNM}(ρ)+ δ² S(ρ)
 β-equilibrium condition → NS matter EOS
- Max. mass in the fit EOS deviates from the original one by ~ 0.1 M_☉.





A. Ohnishi @ YONUPA, Aug.17, 2015 87

 $\eta = (KL^2)^{1/3}$?

Sotani et al.(2014)

NS matter in "ab initio"-fit + A

A potential in nuclear matter at $\rho_0 \sim -30$ MeV

- Scheme 1: $U_{\Lambda}(\rho) = \alpha U_{N}(\rho)$
- Scheme 2: $U_{\Lambda}(\rho) = 2/3 U^{n=2}N(\rho) + \beta U^{n>2}N(\rho)$

