

SATELLITE SYMPOSIUM
NEW ERA IN NEUTRINO PHYSICS

T.M.U., June 1998

THREE-FLAVOR MIXING
and
TRIANGLE GRAPHS

E. LISI, INFN Bari, Italy

"...The purpose of the meeting is twofold:

1) To have an extensive discussion on some restricted topics like 3-flavor analyses of various neutrino physics phenomenology, new experimental projects which will start in a few years, etc.

2) to bring young people in Japan to the meeting and tell them how exciting is the field, what is the present understanding of the neutrino physics, and what are the basics to understand them..."

- Purpose of this seminar (pedagogical):
- **UNDERSTANDING THE BASICS OF 3-FLAVOR ANALYSES**

● TWO-FLAVOR PARAMETER SPACE :

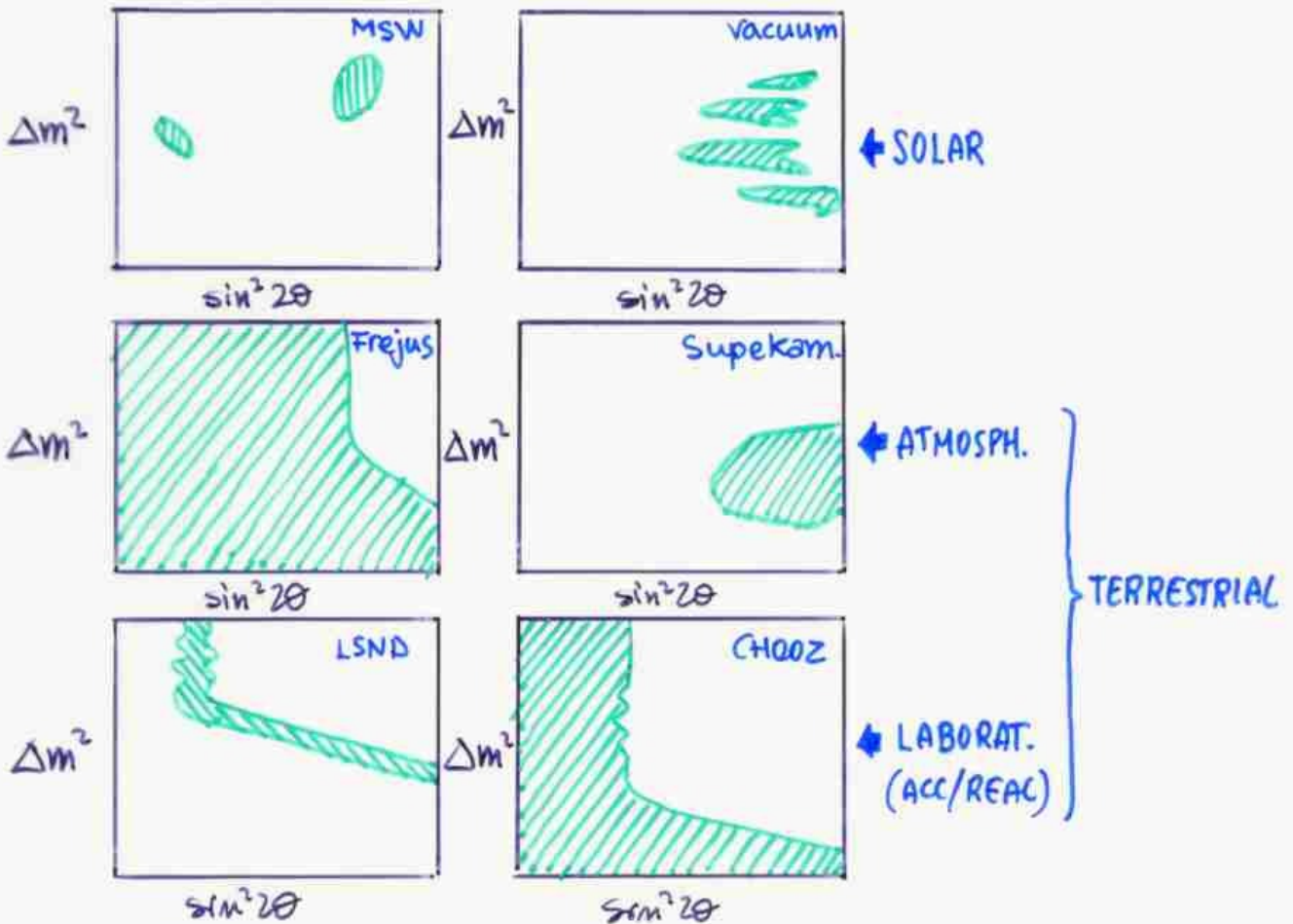
2-dimensional
 $(\Delta m^2, \sin^2 2\theta)$



$$\begin{pmatrix} \nu_\alpha \\ \nu_\beta \end{pmatrix} = \begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$$

$\alpha, \beta = e, \mu, \tau$
 $\theta \in [0, \pi/4]$

● USUAL GRAPHICAL REPRESENTATIONS :



... etc.

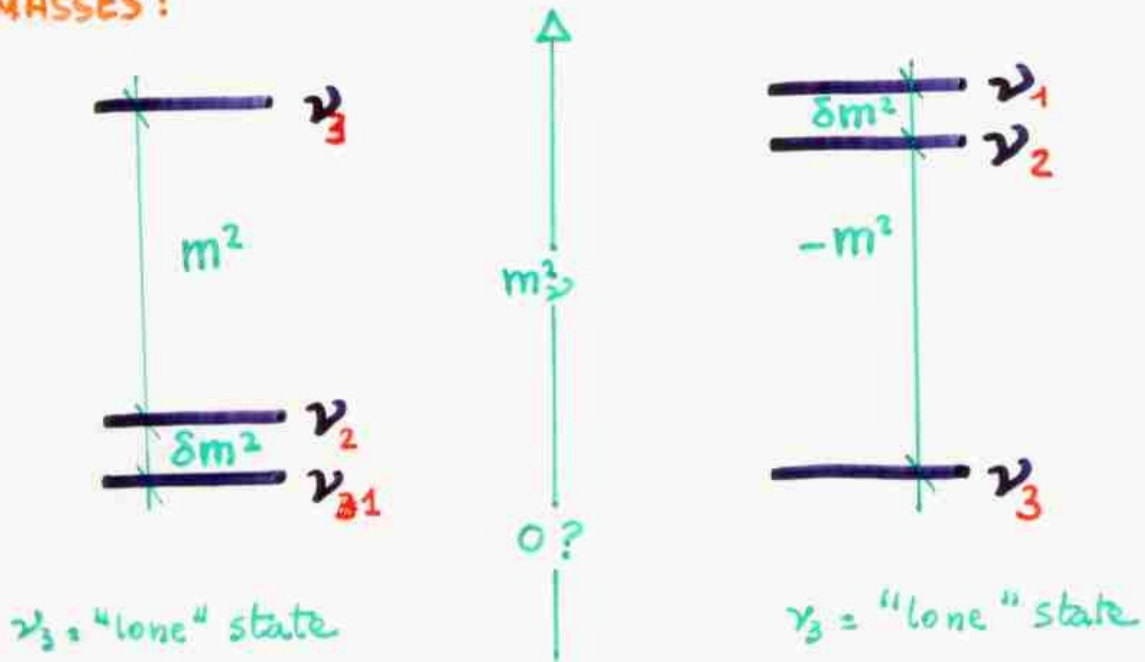
• THREE-FLAVOR PARAMETER SPACE

6-dimensional

$$(\delta m^2, m^2, \omega, \psi, \varphi, \delta_{CP})$$

$\delta m^2 \leq m^2$ by convention
 $\omega, \varphi, \psi \in [0, \pi/2]$

• MASSES:



• MIXING:

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_\psi & s_\psi \\ 0 & -s_\psi & c_\psi \end{pmatrix} \begin{pmatrix} c_\varphi & 0 & s_\varphi \\ 0 & 1 & 0 \\ -s_\varphi & 0 & c_\varphi \end{pmatrix} \begin{pmatrix} c_\omega & s_\omega & 0 \\ -s_\omega & c_\omega & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} \otimes \delta_{CP}$$

- PROBLEM: Find useful representations of this huge parameter space

NO EASY solution to this problem in the most general cases

However, parameter space reduction is possible under the reasonable hypothesis

$$\delta m^2 \ll m^2$$

motivated by the current ν phenomenology:

$$\delta m^2 \lesssim 10^{-4} \text{ eV}^2 \quad (\text{solar neutrinos})$$

$$m^2 \gtrsim 10^{-3} \text{ eV}^2 \quad (\text{terrestrial neutrinos})$$



$\delta m^2/m^2 \lesssim 1/10$, and often much smaller

- 1st welcome simplification:

CP violation effects vanish as $\delta m^2/m^2 \rightarrow 0$
(Cabibbo, 1978)

→ $\delta_{CP} \sim$ unobservable

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

↙ mixing matrix $U_{\alpha i}$ real

$\nu/\bar{\nu}$ oscillations indistinguishable (in vacuum)

PARAMETER SPACE REDUCTION FOR SOLAR NEUTRINOS

Only $P(\nu_e \rightarrow \nu_e)$ is probed by γ_0
(no μ, τ appearance expts!)

Hence, rotations in the (ν_μ, ν_τ) subspace
do not change the physics. In particular,
redefine ν_μ and ν_τ :

$$\begin{pmatrix} \nu'_\mu \\ \nu'_\tau \end{pmatrix} = \begin{pmatrix} c\psi & -s\psi \\ s\psi & c\psi \end{pmatrix} \begin{pmatrix} \nu_\mu \\ \nu_\tau \end{pmatrix}$$

so that

$$\begin{pmatrix} \nu_e \\ \nu'_\mu \\ \nu'_\tau \end{pmatrix} = \begin{pmatrix} c\varphi & 0 & s\varphi \\ 0 & 1 & 0 \\ -s\varphi & 0 & c\varphi \end{pmatrix} \begin{pmatrix} c\omega & s\omega & 0 \\ -s\omega & c\omega & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

and only φ and ω are observable.

In the $m^2 \gg \delta m^2$ limit, m^2 -driven oscillations
are averaged away



SOLAR NEUTRINOS PROBE

$$\left(\delta m^2, \omega, \varphi \right)$$

IN THE LIMIT $\delta m^2/m^2 \rightarrow 0$

PARAMETER SPACE REDUCTION FOR TERRESTRIAL NEUTRINOS

In the limit $\delta m^2 \rightarrow 0$, terrestrial (lab. or atmosph.) experiments become insensitive to rotations in the subspace spanned by the quasi-degenerate doublet (ν_1, ν_2) . In particular, redefine

$$\begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix} = \begin{pmatrix} c\omega & -s\omega \\ s\omega & c\omega \end{pmatrix} \begin{pmatrix} \nu'_1 \\ \nu'_2 \end{pmatrix}$$

so that

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c\psi & s\psi \\ 0 & -s\psi & c\psi \end{pmatrix} \begin{pmatrix} c\varphi & 0 & s\varphi \\ 0 & 1 & 0 \\ -s\varphi & 0 & c\varphi \end{pmatrix} \begin{pmatrix} \nu'_1 \\ \nu'_2 \\ \nu'_3 \end{pmatrix}$$

and only ψ and φ are observable.

In the $m^2 \gg \delta m^2$ limit, δm^2 -driven oscillations are frozen (in vacuum)



TERRESTRIAL NEUTRINOS PROBE

$$(m^2, \psi, \varphi)$$

IN THE LIMIT $\delta m^2/m^2 \rightarrow 0$

SUMMARY:

At zeroth order in $\delta m^2/m^2$,

- SOLAR NEUTRINOS PROBE
- TERRESTRIAL NEUTRINOS PROBE

$$\begin{array}{l} \delta m^2, \omega, \varphi \\ m^2, \psi, \varphi \end{array}$$

or, equivalently:

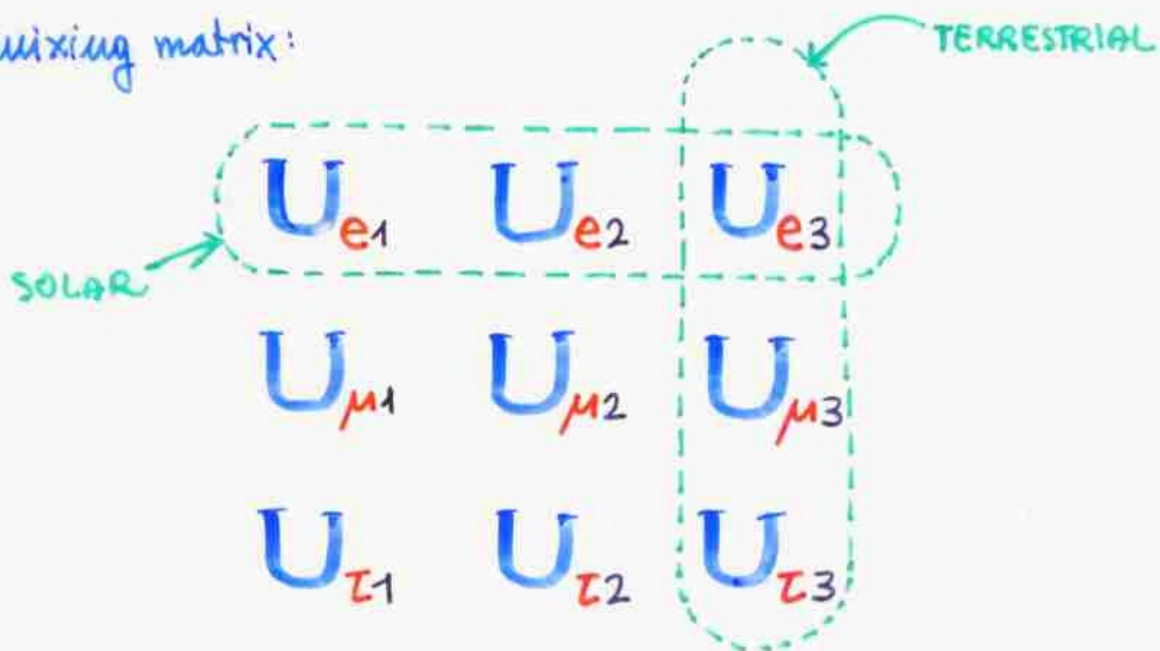
- Solar neutrinos probe $(\delta m^2, U_{e1}, U_{e2}, U_{e3})$, i.e. the mass composition of ν_e :

$$\begin{aligned} \nu_e &= U_{e1} \nu_1 + U_{e2} \nu_2 + U_{e3} \nu_3 \\ &= C_\varphi (C_\omega \nu_1 + S_\omega \nu_2) + S_\varphi \nu_3 \end{aligned}$$

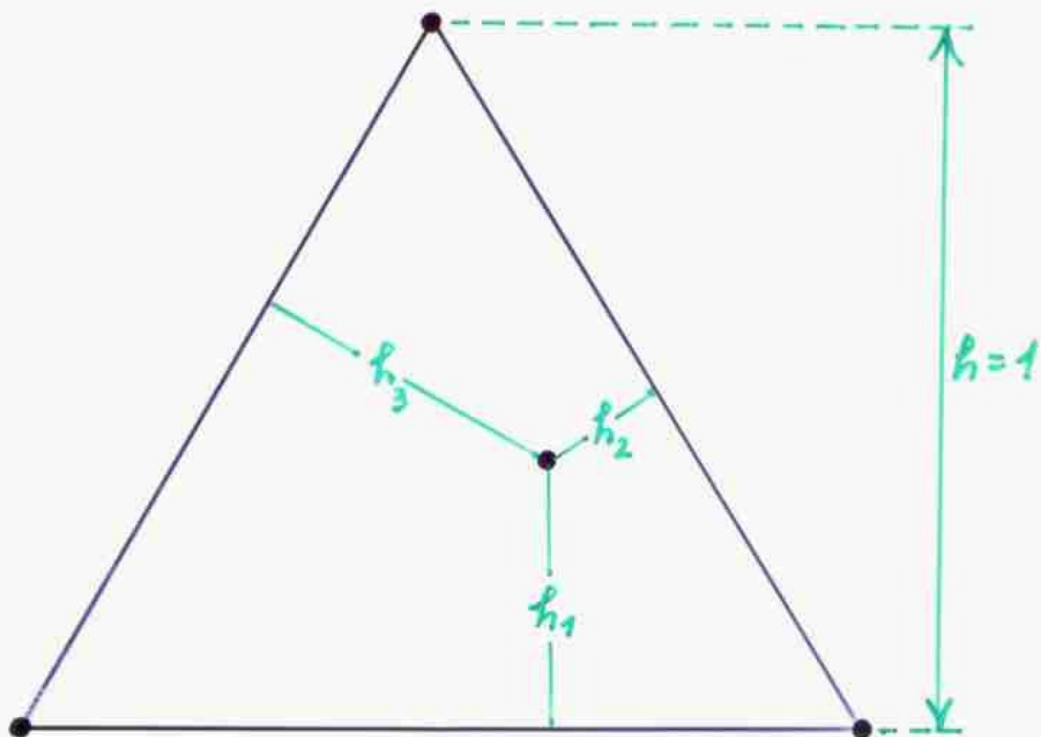
- Terrestrial neutrinos probe $(m^2, U_{e3}, U_{\mu 3}, U_{\tau 3})$, i.e. the flavor composition of ν_3 :

$$\begin{aligned} \nu_3 &= U_{e3} \nu_e + U_{\mu 3} \nu_\mu + U_{\tau 3} \nu_\tau \\ &= C_\varphi (C_\psi \nu_\tau + S_\psi \nu_\mu) + S_\varphi \nu_e \end{aligned}$$

mixing matrix:



Unitarity ($U_{e3}^2 + U_{\mu 3}^2 + U_{\tau 3}^2 = 1$ or $U_{e1}^2 + U_{e2}^2 + U_{e3}^2 = 1$)
can be implemented through
TRIANGLE GRAPHS :

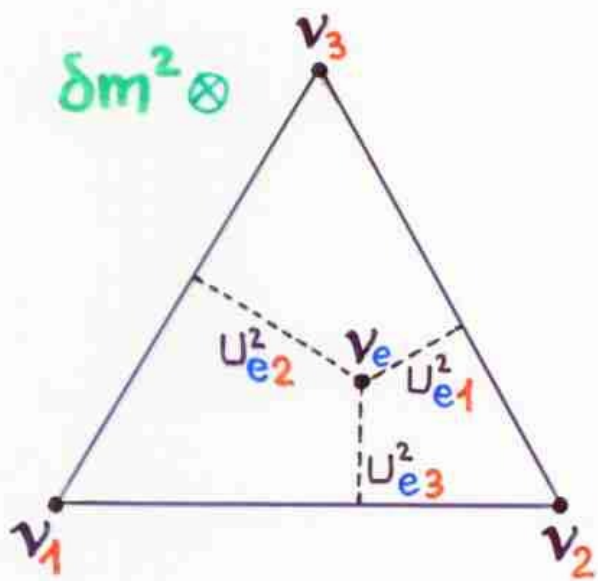


$$h_1 + h_2 + h_3 \equiv h = 1$$

Unitarity suggests the following representations:

SOLAR ν

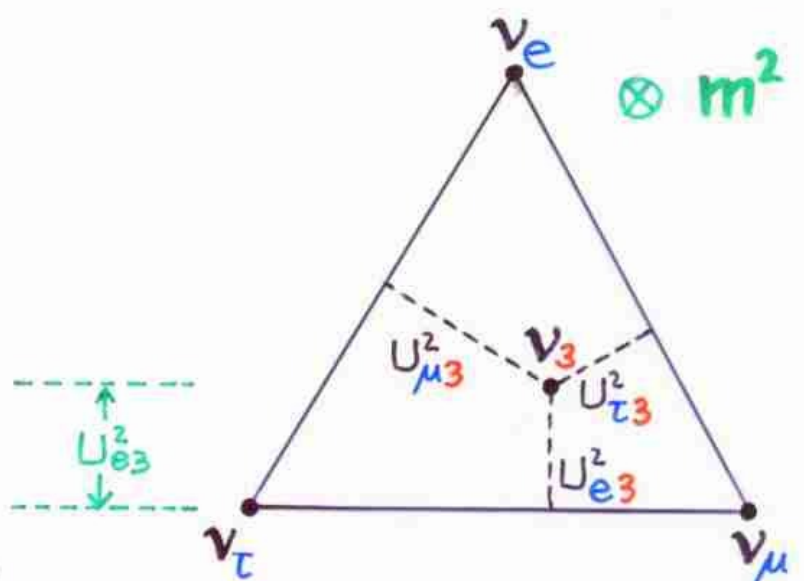
parameter space:



$$U_{e1}^2 + U_{e2}^2 + U_{e3}^2 = 1$$

TERRESTRIAL ν

parameter space:

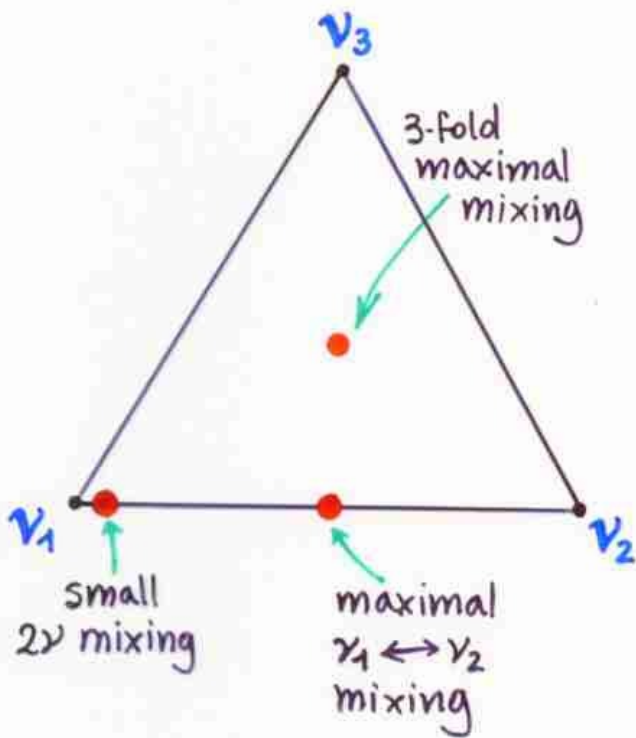


$$U_{e3}^2 + U_{\mu 3}^2 + U_{\tau 3}^2 = 1$$

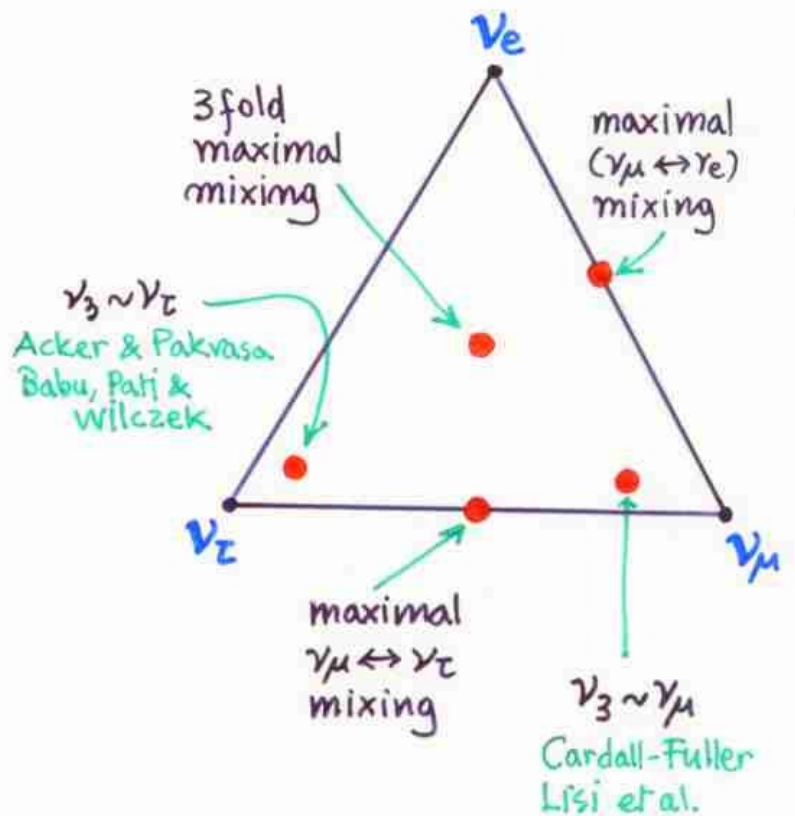
(Lisi, Fogli, Montanino, Scioscia)

INTERESTING SUBCASES :

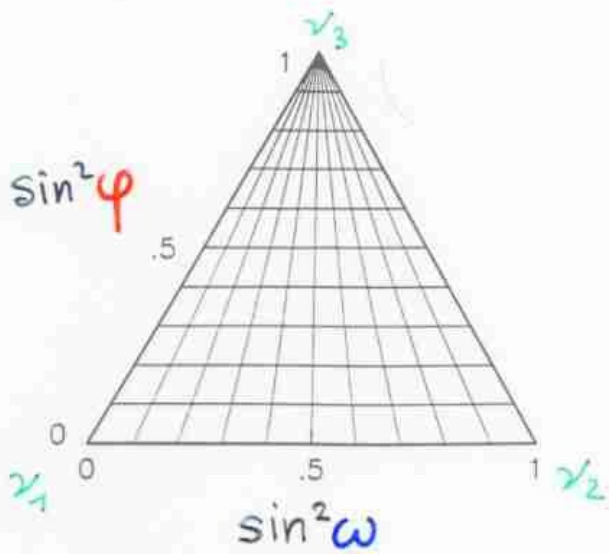
SOLAR ν



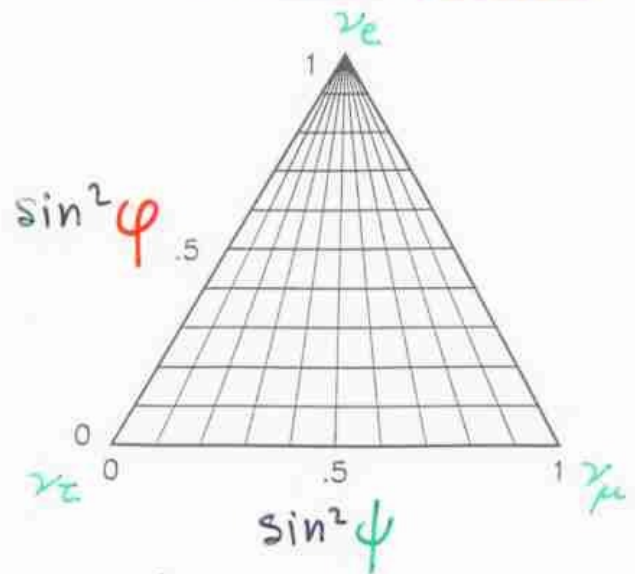
TERRESTRIAL ν



SOLAR

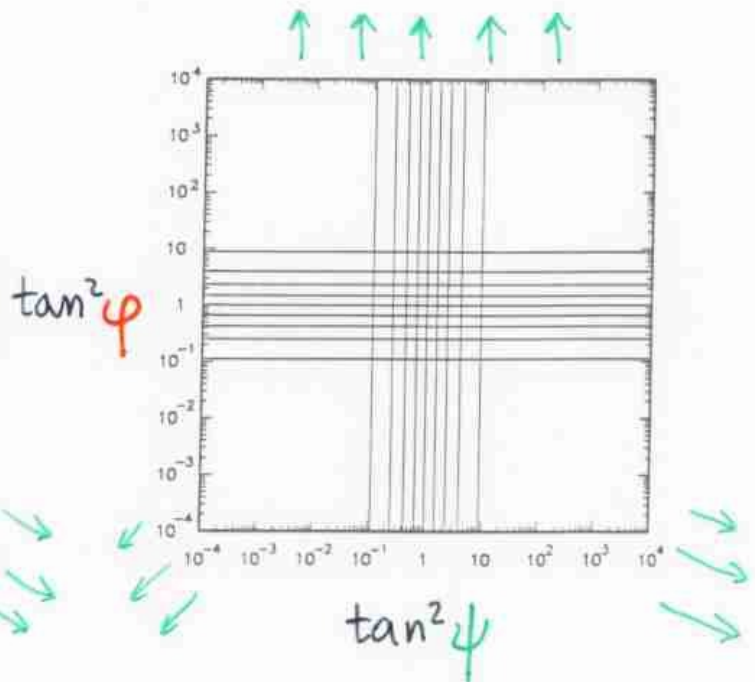
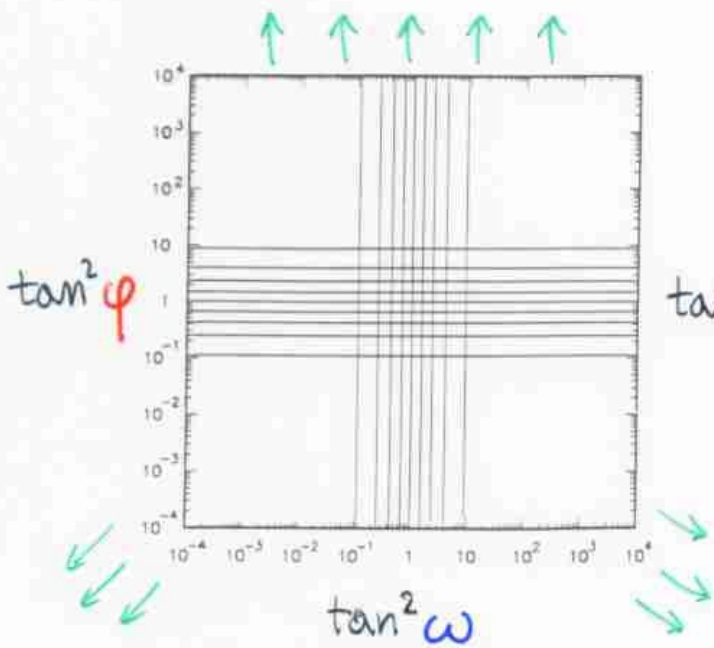


TERRESTRIAL



δm^2 fixed

m^2 fixed



BILOGARITHMIC MIXING/MIXING REPRESENTATION

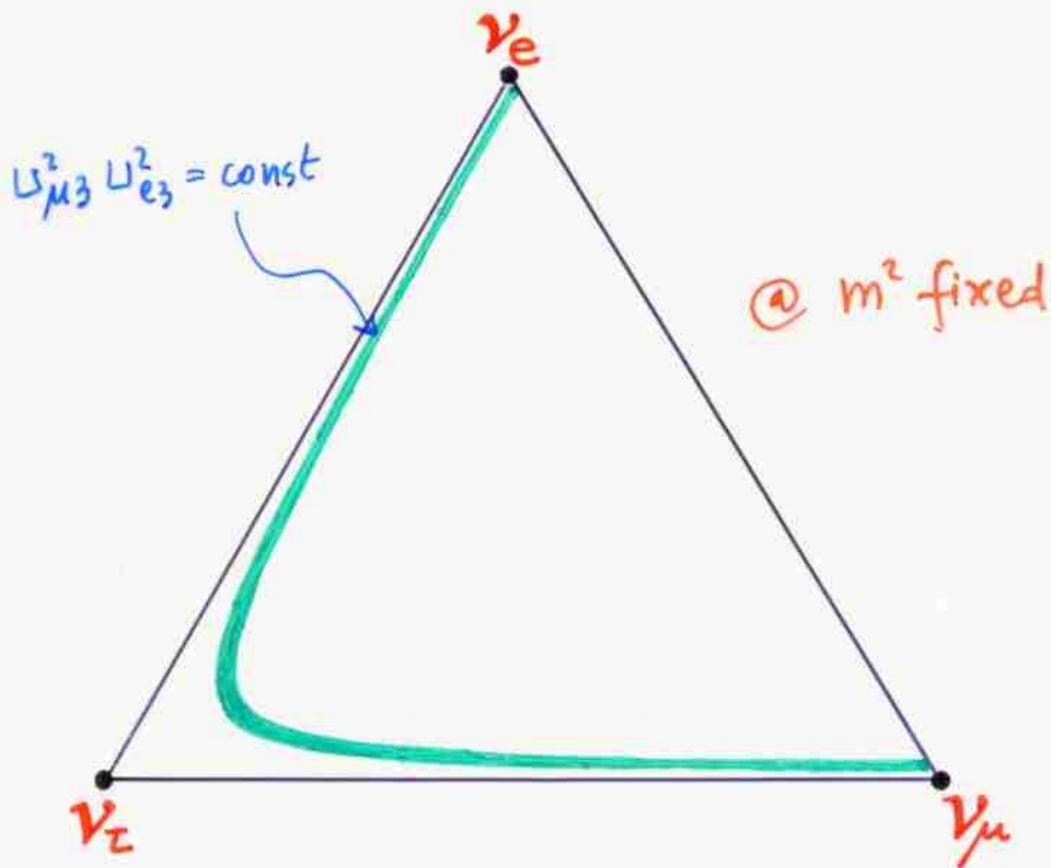
1st EXAMPLE : LSND

$$P_{\text{LSND}} = P(\nu_{\mu} \leftrightarrow \nu_e) \neq 0$$

$$32: P_{\text{LSND}} = 4 U_{\mu 3}^2 U_{e 3}^2 \sin^2 \left(1.27 \frac{m^2 L}{E} \right) \quad \left(\frac{\text{eV}^2 \text{ km}}{\text{GeV}} \right)$$



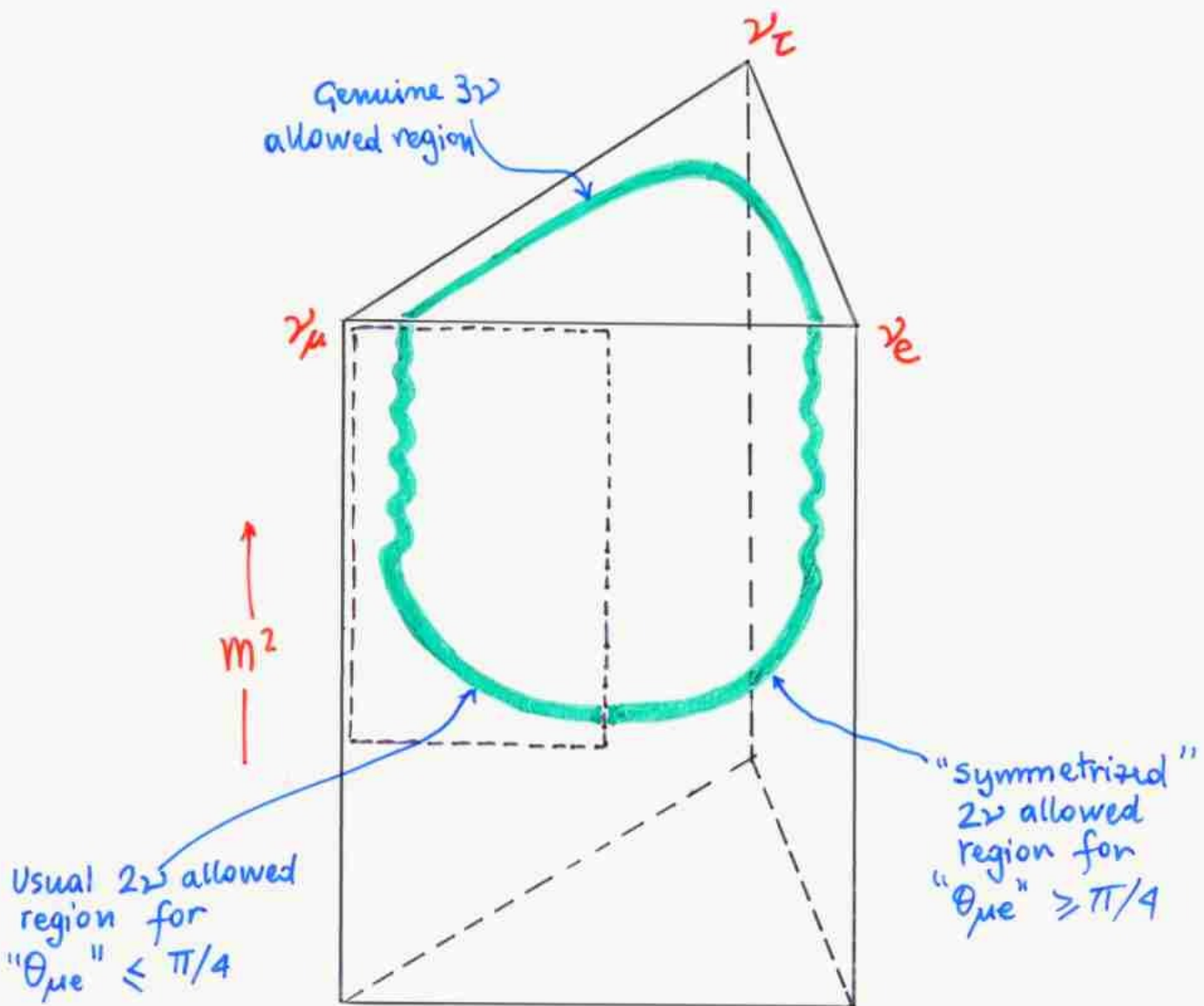
- Region allowed by LSND at fixed m^2
(bounded by $P_{\text{LSND}} = \text{const.}$ curves)



$$" \sin^2 2\theta_{e\mu} " \equiv 4 U_{\mu 3}^2 U_{e 3}^2 = 4 s_{\psi}^2 c_{\psi}^2 s_{\psi}^2$$

VIEW OF THE LSND ALLOWED REGION IN THE $(m^2, \omega_{\alpha 3}^2)$ SPACE

■ ALLOWED "SHELL"



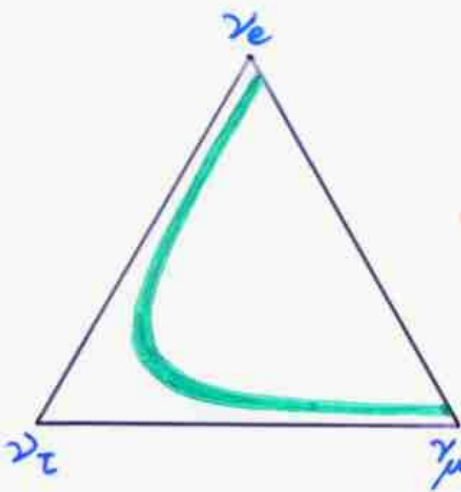
LSND

+

negative
acc/react
searches

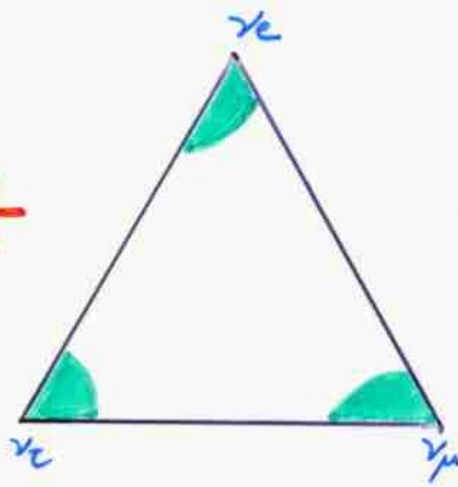
=

$\nu_3 \sim \nu_e$
or $\nu_3 \sim \nu_\mu$



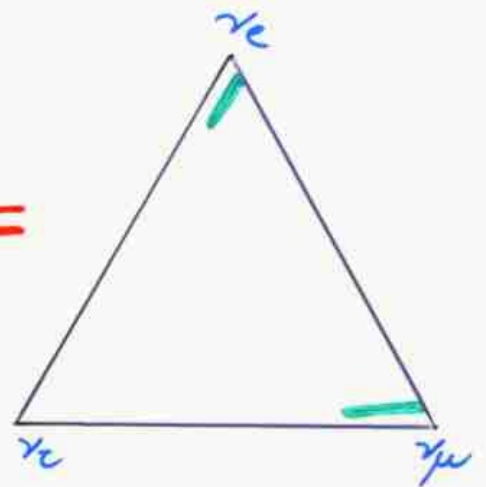
LSND allowed region:
 $U_{e3}^2 U_{\mu 3}^2 \sim \text{const}$

+



Regions allowed
by negative
searches:
 $\nu_3 \sim \nu_\alpha$ ($\alpha = e, \mu, \tau$)
to suppress
oscillations

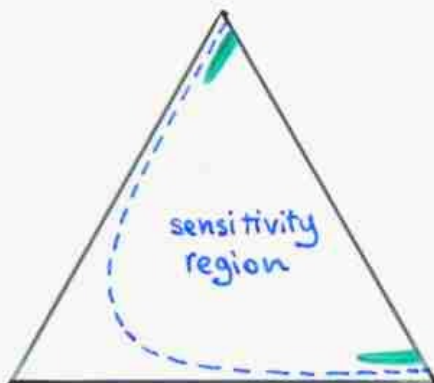
=



Combination:
 ν_3 can be close
to either ν_e or
 ν_μ but not
to ν_τ (at 99% CL)

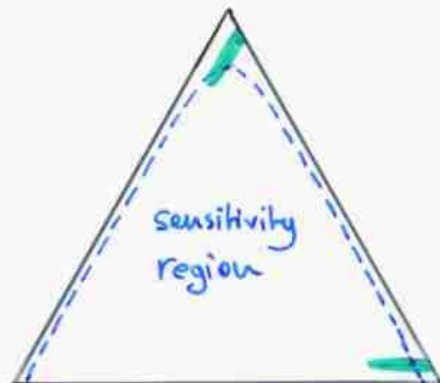
POSSIBLE TESTS AT ACCELERATORS:

$\nu_\mu \leftrightarrow \nu_e$ channel



KARMEN
BOONE
CERN expt. ?

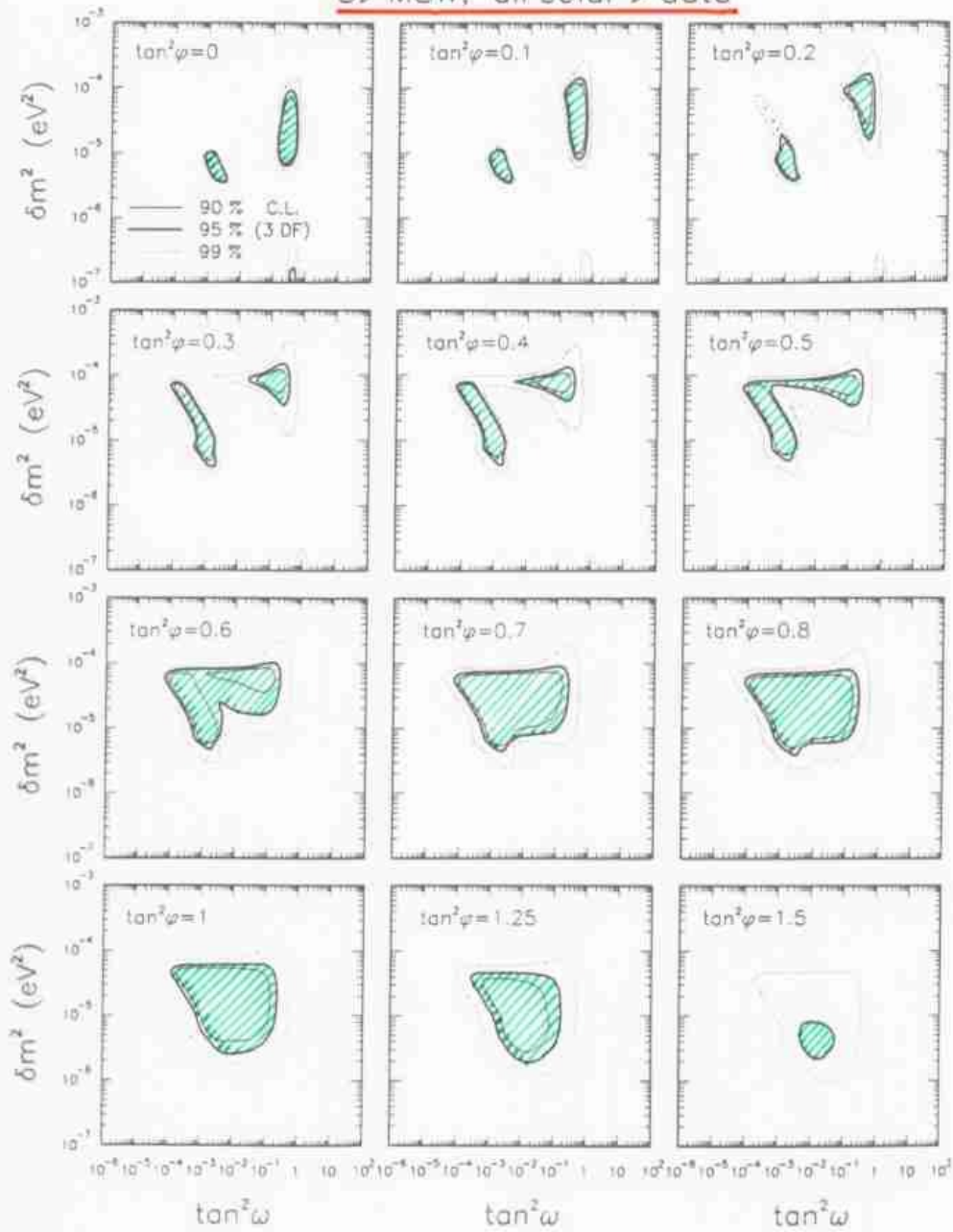
$\nu_\mu \leftrightarrow \nu_\tau$ channel



~ CHORUS, NOMAD
COSMOS
TOSCA

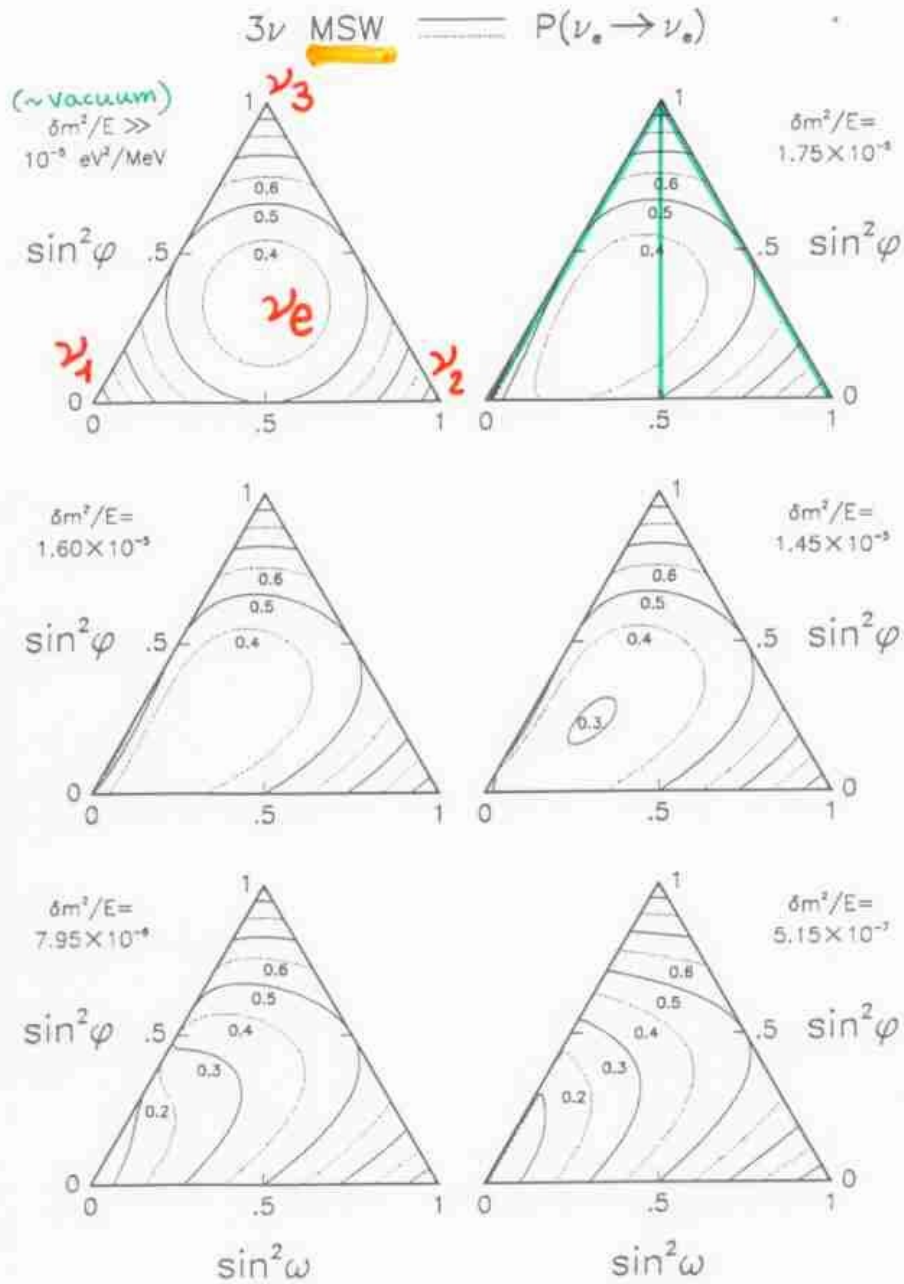


3ν MSW, all solar ν data



$(\delta m^2, \tan^2 \omega)$ @ $\tan^2 \varphi$ fixed

SOLAR ν_s



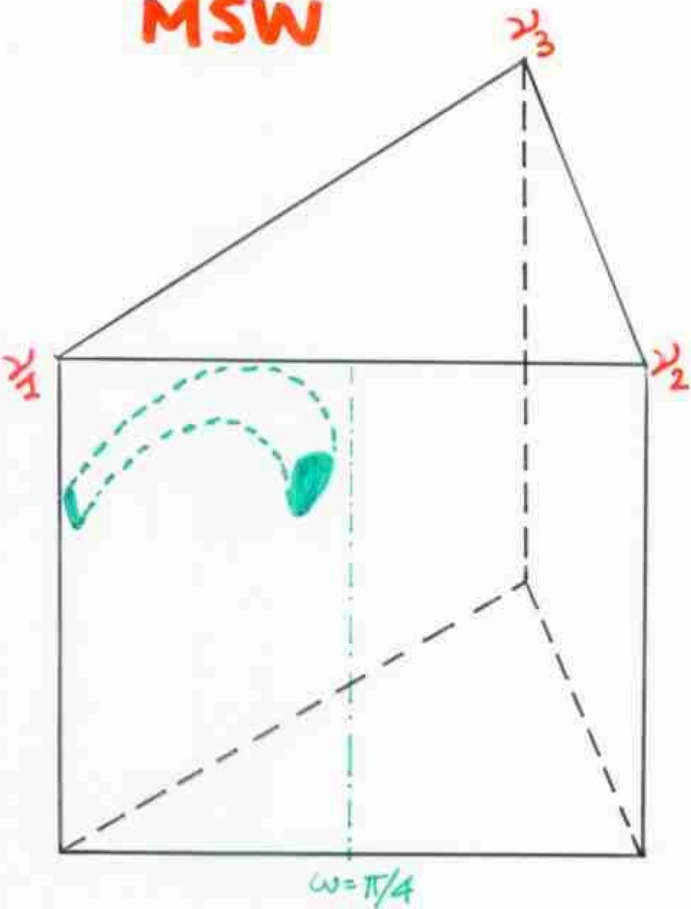
MSW probability contours

Along green lines ($\omega = 0, \pi/4, \pi/2$): $P_{\text{MSW}} = P_{\text{vac}}$

2nd example : Solar ν 's

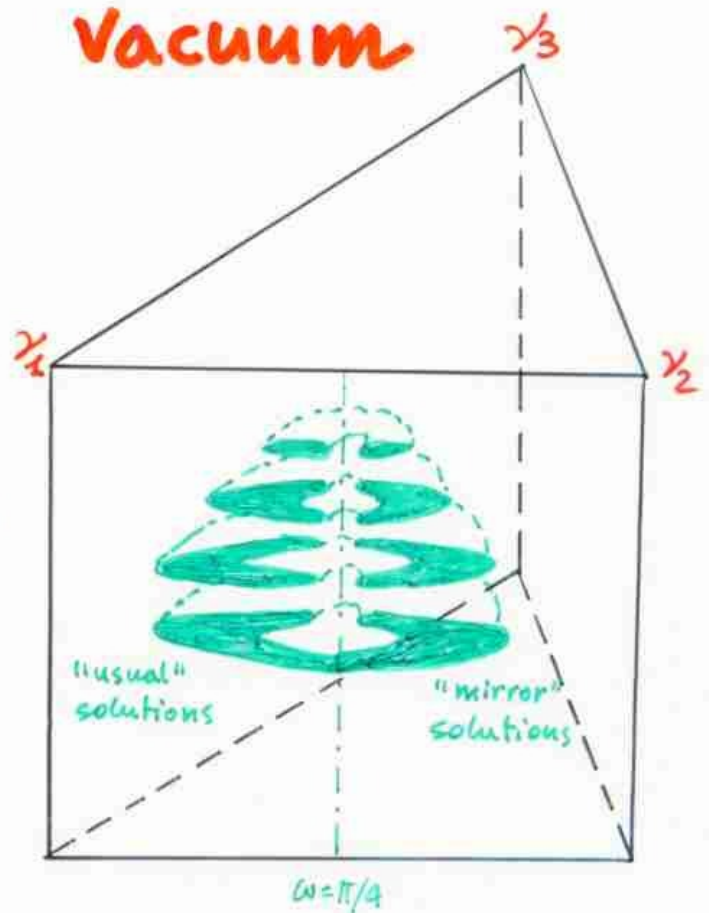
$$\nu_e = U_{e1} \nu_1 + U_{e2} \nu_2 + U_{e3} \nu_3$$

MSW



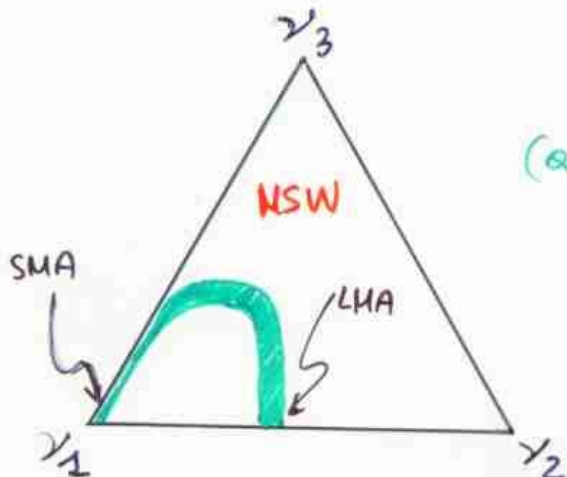
δm^2

Vacuum

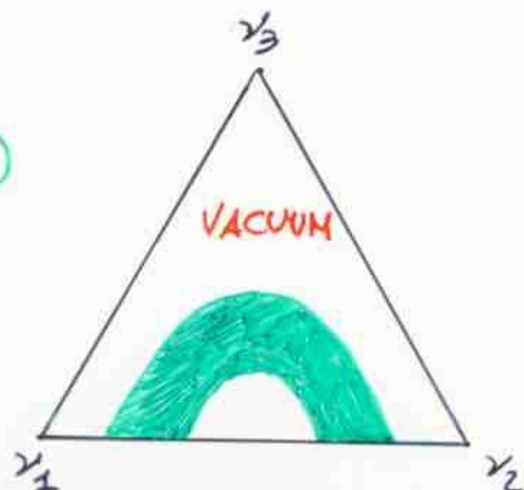


SMALL AND LARGE MIXING
MSW SOLUTIONS ARE
CONTINUOUSLY CONNECTED

VACUUM SOLUTIONS
ARE MIRRORED AND
CONNECTED



(qualitative)



In any case (MSW or vacuum osc.)
large values of ϕ ($\rightarrow \pi/2$) are
excluded :

$$\phi \rightarrow \pi/2 \Leftrightarrow U_{e3}^2 = S_{\phi}^2 \rightarrow 1 \Leftrightarrow \nu_e \rightarrow \nu_3$$

$$\text{and } P_{\text{solar}}(\nu_e \rightarrow \nu_e) \rightarrow 1$$

\rightarrow NO SOLAR ν DEFICIT !!!

Moreover, solar ν data indicate
that 2ν fits are better than 3ν

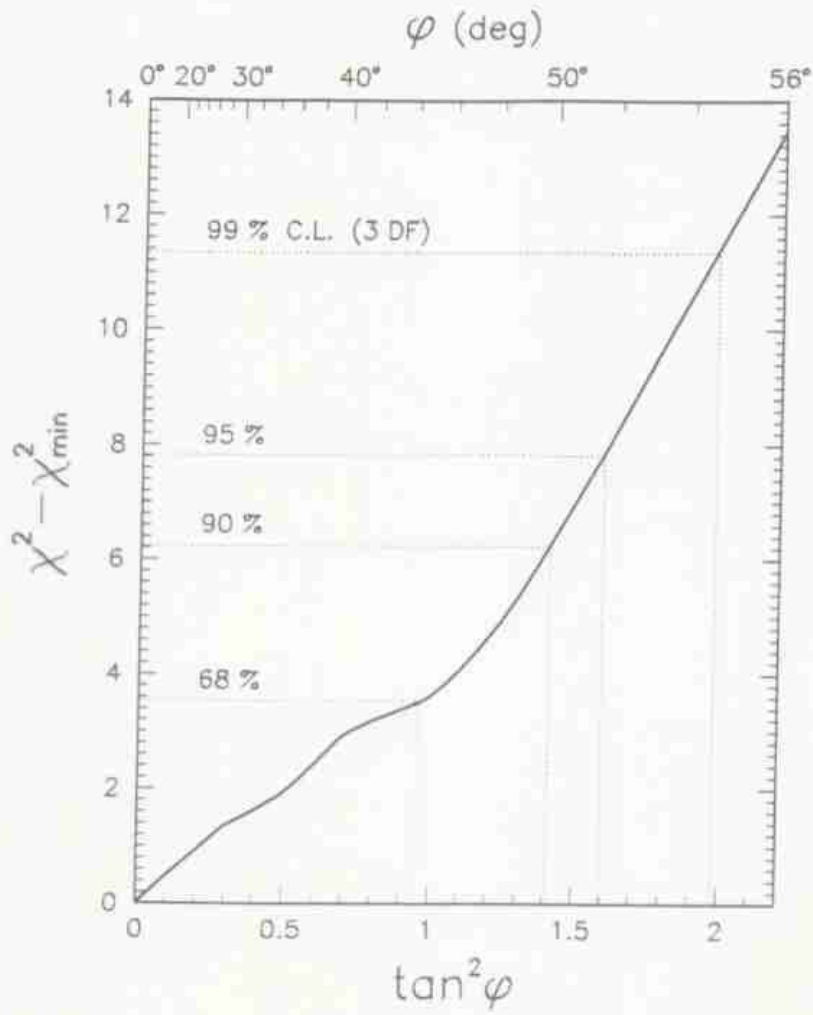
$$\chi^2(2\nu) < \chi^2(3\nu)$$

(MSW: Fogli, E.L., Montanino)
(vacuum: Berezhiani, Rossi)

although up to $\phi \sim 40^\circ$
or 50° is allowed

\uparrow
NONTRIVIAL INDICATION THAT
IS WORTH CHECKING WITH THE
LATEST DATA

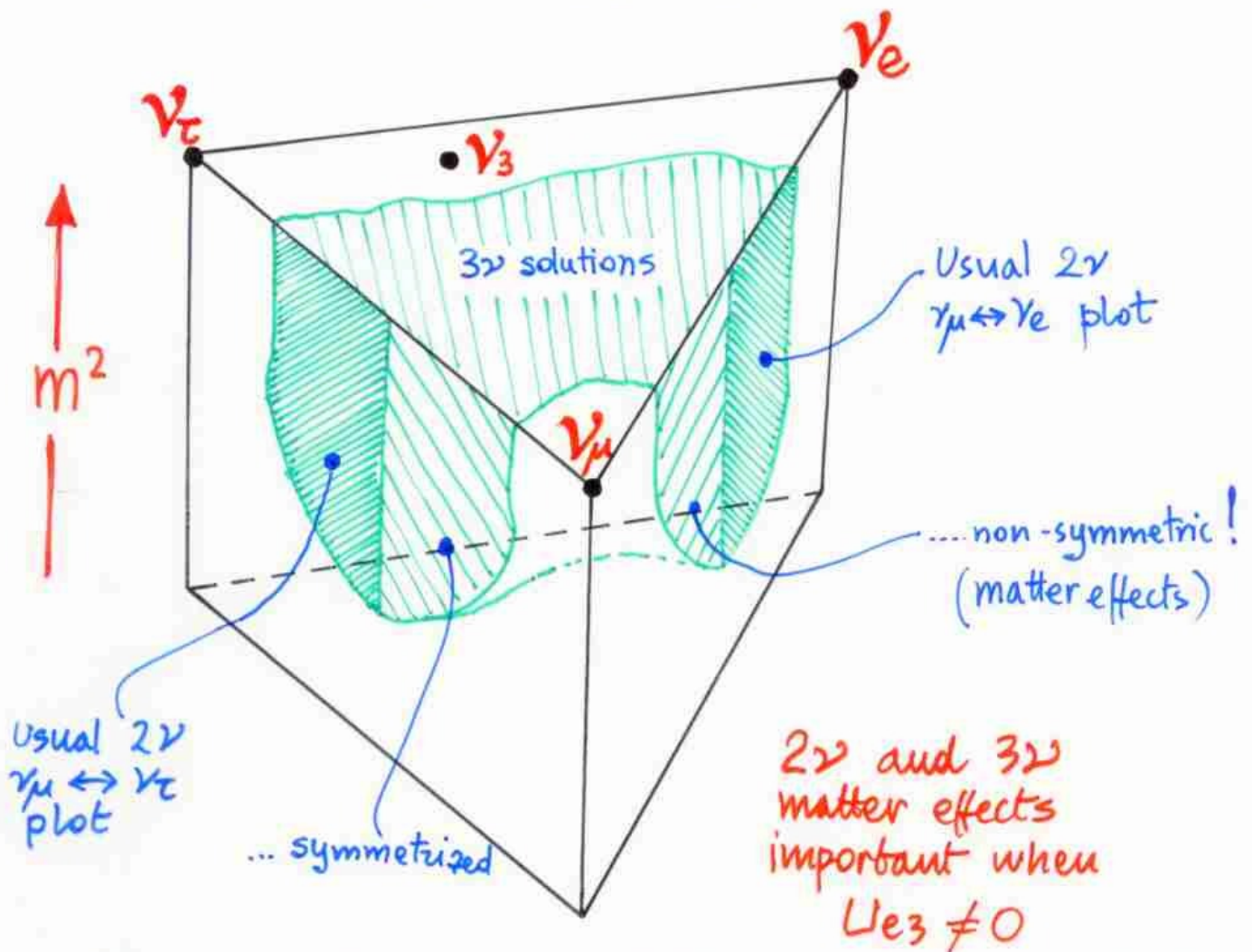
Solar neutrinos,
Pre-Superkamiokande data (MSW)



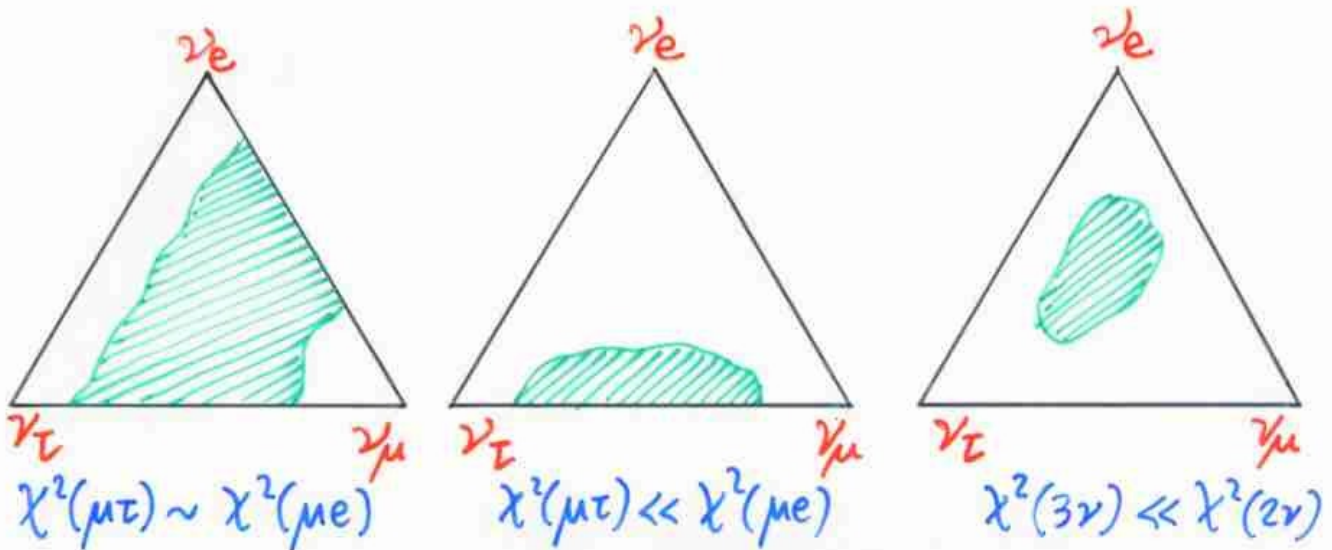
δm^2 and $\tan^2 \omega$ unconstrained

3rd example : Atmospheric ν

Parameter space : $(m^2, U_{e3}^2, U_{\mu 3}^2, U_{\tau 3}^2)$

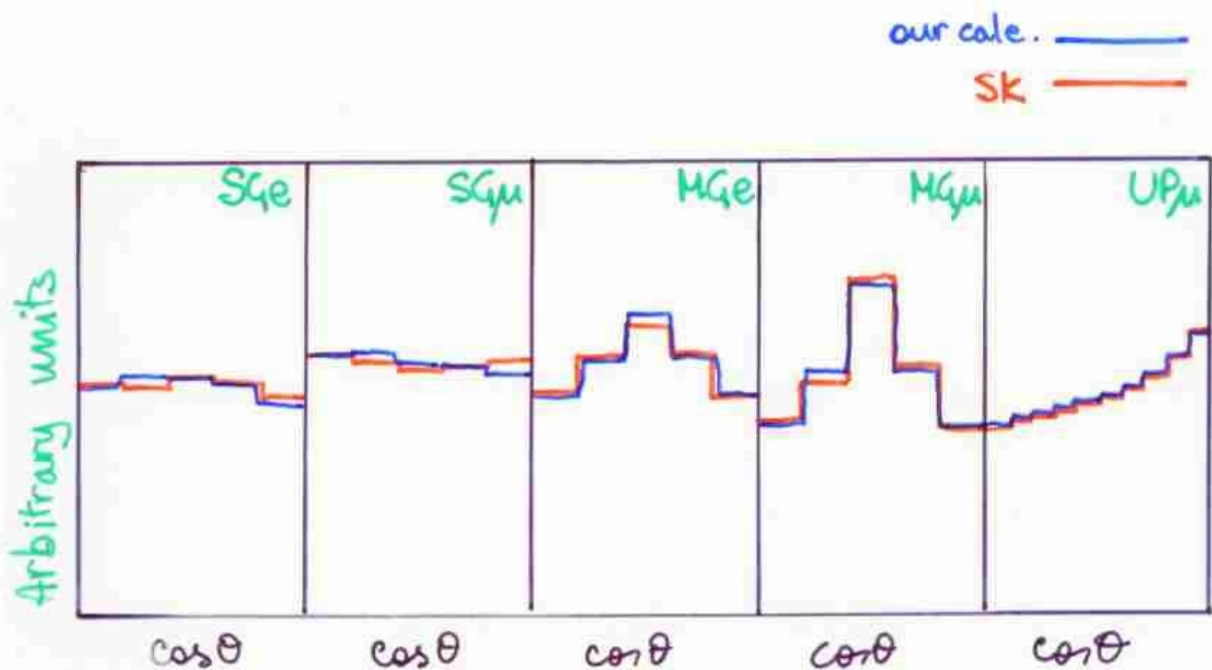


Depending on the DATA and on m^2 , several situations are possible, e.g. :



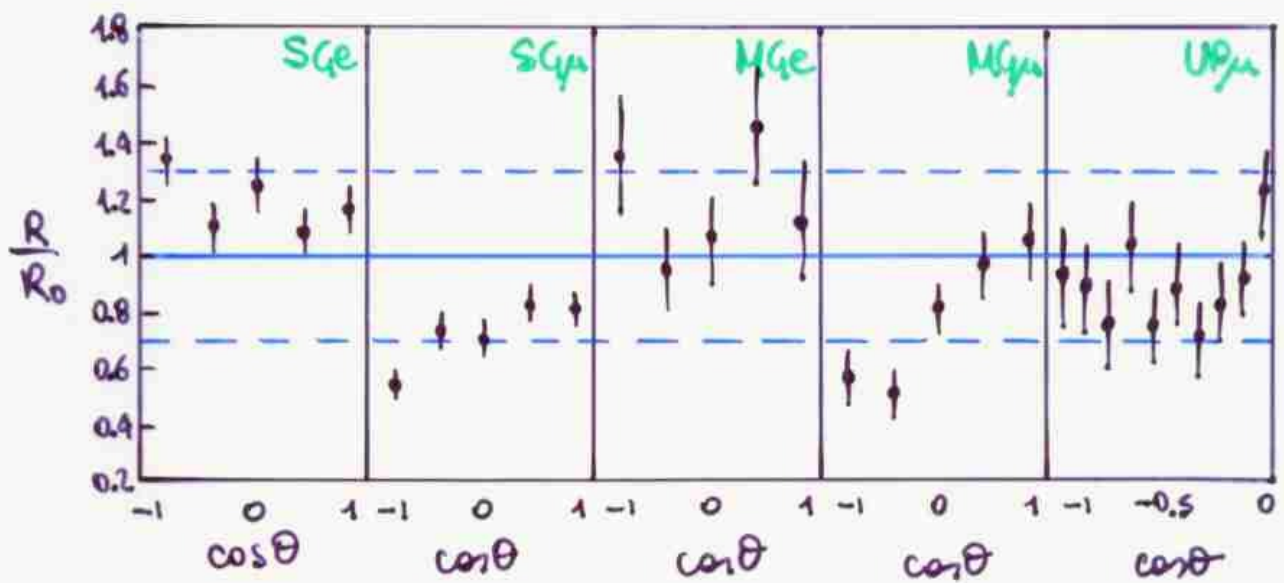
G.L. Fogli, E.L., A. Marrone, G. Scioscia
- Preliminary - (~ 1y work)

NO-OSCILLATION CHECK OF S-K ANGULAR DISTRIBUTIONS



- + checks for oscillation points
- + $\theta_{\nu e}, \theta_{\nu \mu} \neq 0$ at different energies
- +
- + $P(\nu_{\alpha} \rightarrow \nu_{\beta})$ and $P(\bar{\nu}_{\alpha} \rightarrow \bar{\nu}_{\beta})$ calculated numerically for the Earth density profile @ generic $m_1, m_2, m_3, \omega, \varphi, \psi$

SK 400 day

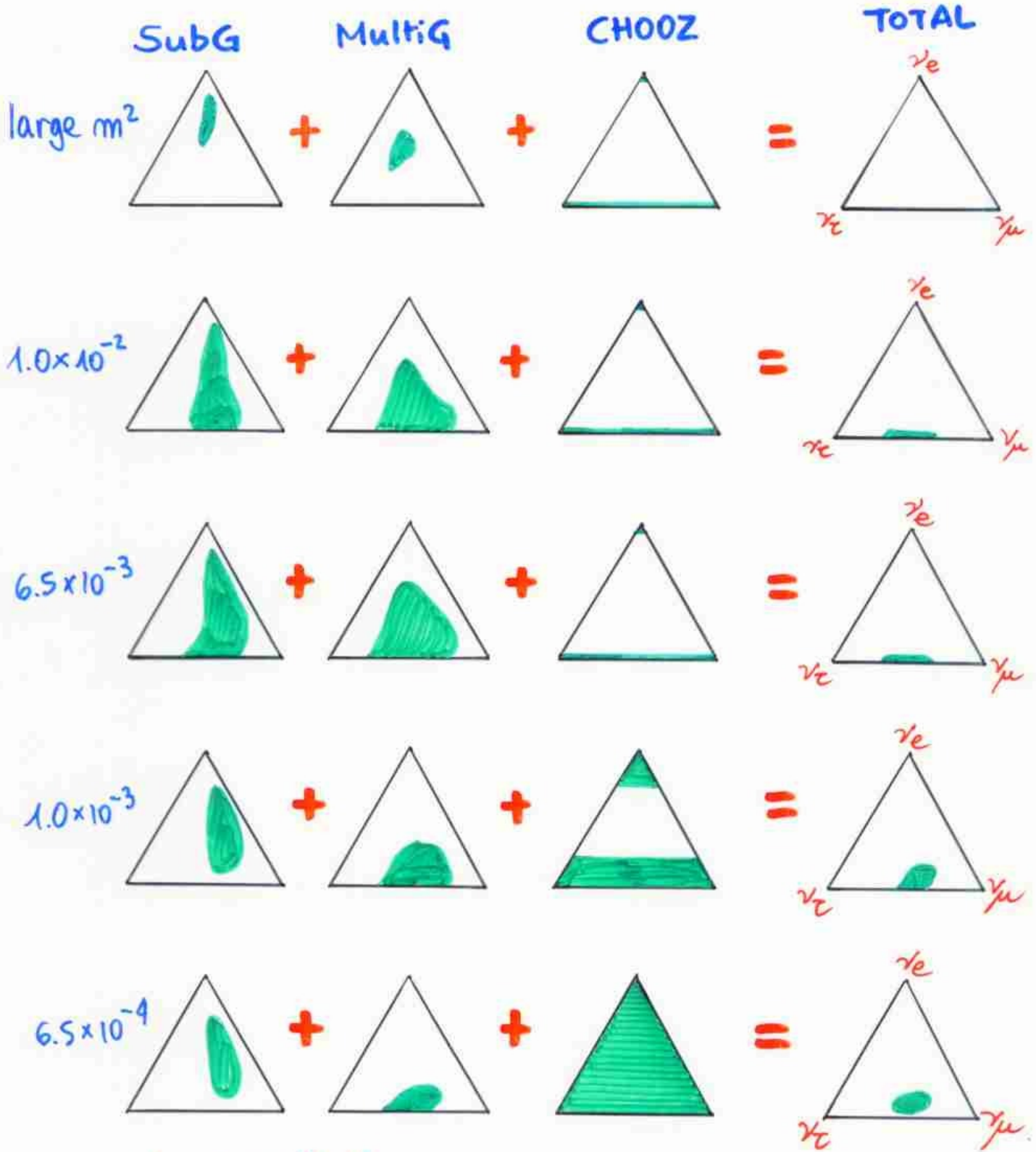


Theory = 1 ± 0.3 in this plot

SuperKamiokande 400 day DATA PRELIMINARY ANALYSIS

(Fogli, E.L., Marrone, Scioccia)

We include generic 3ν matter effects numerically (≥ 1994)



$\chi^2(3\nu) \ll \chi^2(2\nu) \quad ? \quad \text{NOT CLEARLY ESTABLISHED BY 400d SK DATA}$

$L_{e3} \neq 0$ leads to a much richer phenomenology and helps fitting SK data.

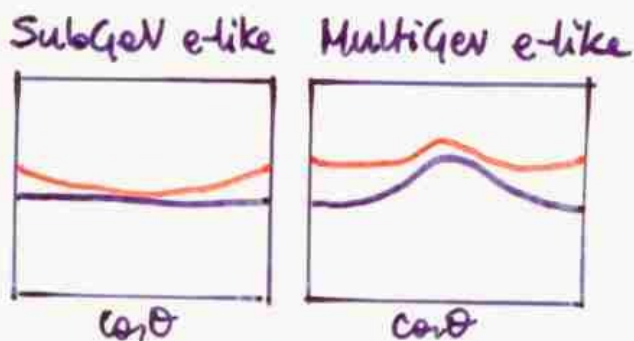
$L_{e3} \neq 0 \rightarrow$ several new effects that do not show the "canonical" L/E dependence

One such effect is due to the fact that

$\frac{\nu_{\mu}^{(-)}}{\nu_{e}^{(-)}}$ increases for $\begin{cases} \text{high energies} \\ \cos\theta \rightarrow \pm 1 \text{ (vertical)} \end{cases}$

\rightarrow THE ATMOSPHERIC NEUTRINO BEAM IS RICHER IN $\nu_{\mu}^{(-)}$'s at high energy and along the vertical
Best place to look for $\nu_{\mu} \rightarrow \nu_{e}$ oscillations!

$L_{e3} \neq 0 \Rightarrow \nu_{\mu} \leftrightarrow \nu_{e}$ open \Rightarrow e-like rate increases
effect larger as E increases and $\cos\theta \rightarrow \pm 1$!!



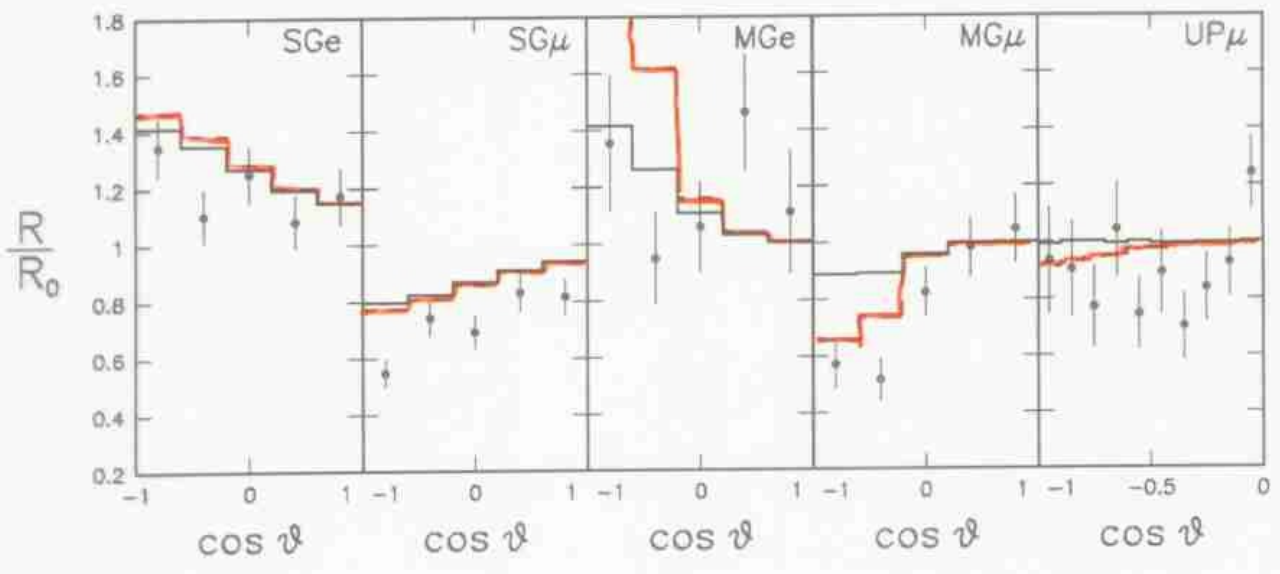
SPECTRAL DEFORMATION NOT DEPENDENT ON L/E !!! (even in vacuum)

PRELIMIN.

Example of 2ν matter effects (not dependent on L/E !)

$\nu_\mu \leftrightarrow \nu_e$ with $m^2 = 8.E-4$ *
Size of matter effects

	m^2	U_{e3}^2	$U_{\mu 3}^2$	$U_{\tau 3}^2$
— (mat.)	$8.E-4$	0.50	0.50	0.00
— (vac.)	$8.E-4$	0.50	0.50	0.00



$$\sin^2 2\theta_{vac} \sim 1 \quad \rightarrow \quad \sin^2 2\theta_{mat} < 1$$

(suppression of oscillation)

* Below Chooz bounds

+ GENUINELY 3ν EFFECTS IN MATTER

Under same conditions:

$$\begin{pmatrix} P_{ee} & P_{e\mu} & P_{e\tau} \\ P_{\mu e} & P_{\mu\mu} & P_{\mu\tau} \\ P_{\tau e} & P_{\tau\mu} & P_{\tau\tau} \end{pmatrix}_{\text{matter}} \approx \begin{pmatrix} P_{ee} & P_{e\mu} & P_{e\tau} \\ P_{\mu e} & P_{\mu\mu} & P_{\mu\tau} \\ P_{\tau e} & P_{\tau\mu} & P_{\tau\tau} \end{pmatrix}_{\text{vac}} + \begin{pmatrix} 0 & 0 & 0 \\ 0 & -\delta P & +\delta P \\ 0 & +\delta P & -\delta P \end{pmatrix}$$

$$\delta P = 4 \frac{U_{e3}^2 U_{\mu 3}^2 U_{\tau 3}^2}{(1 - U_{e3}^2)^2} \sin^2 \left(\frac{G_F}{\sqrt{2}} N_e (1 - U_{e3}^2) L \right)$$

$$\approx 2.47 (1 - U_{e3}^2) \frac{N_e}{\text{mol/cm}^3} \cos \theta$$

3ν large phase!



$\delta P = 0$ for pure 2ν oscillations
no L/E dependence!

Within Chooz bounds: no effect on $\left(\frac{U}{D}\right)_e$
few % effect on $\left(\frac{U}{D}\right)_\mu$

Fogli, E.L., Marrone, Montanino
hep-ph/9711421, to appear in PLB

Solar ν

Terrestrial ν

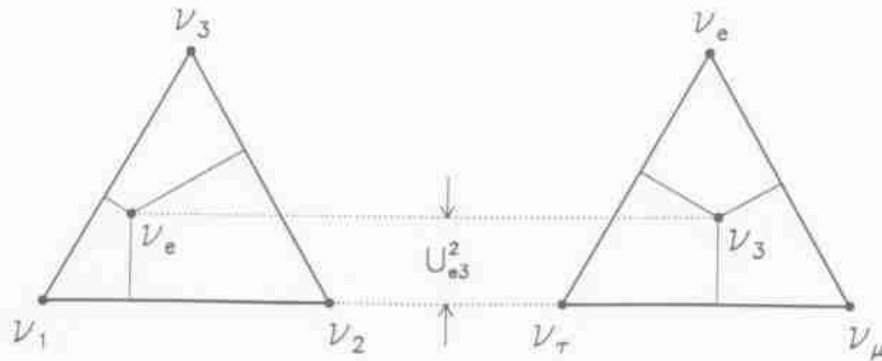
$$\nu_e = U_{e1}\nu_1 + U_{e2}\nu_2 + U_{e3}\nu_3$$

$$\delta m^2 = m_2^2 - m_1^2$$

$$\nu_3 = U_{e3}\nu_e + U_{\mu 3}\nu_\mu + U_{\tau 3}\nu_\tau$$

$$m^2 = m_3^2 - m_{1,2}^2$$

$$(\delta m^2 \ll m^2)$$



ν CKM elements probed:

$$\begin{pmatrix} \underline{e1} & \underline{e2} & \underline{e3} \\ \mu 1 & \mu 2 & \mu 3 \\ \tau 1 & \tau 2 & \tau 3 \end{pmatrix}$$

$$\begin{pmatrix} e1 & e2 & \underline{e3} \\ \mu 1 & \mu 2 & \underline{\mu 3} \\ \tau 1 & \tau 2 & \underline{\tau 3} \end{pmatrix}$$

Parameter space:

$$(\delta m^2, \omega, \varphi)$$

$$(m^2, \psi, \varphi)$$

$$U_{e1}^2 = \cos^2 \varphi \cos^2 \omega$$

$$U_{e2}^2 = \cos^2 \varphi \sin^2 \omega$$

$$U_{e3}^2 = \sin^2 \varphi$$

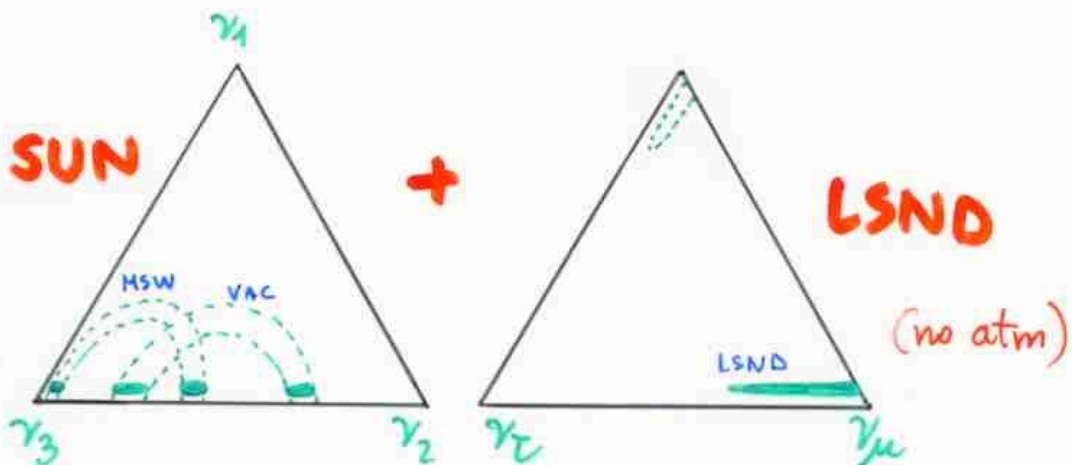
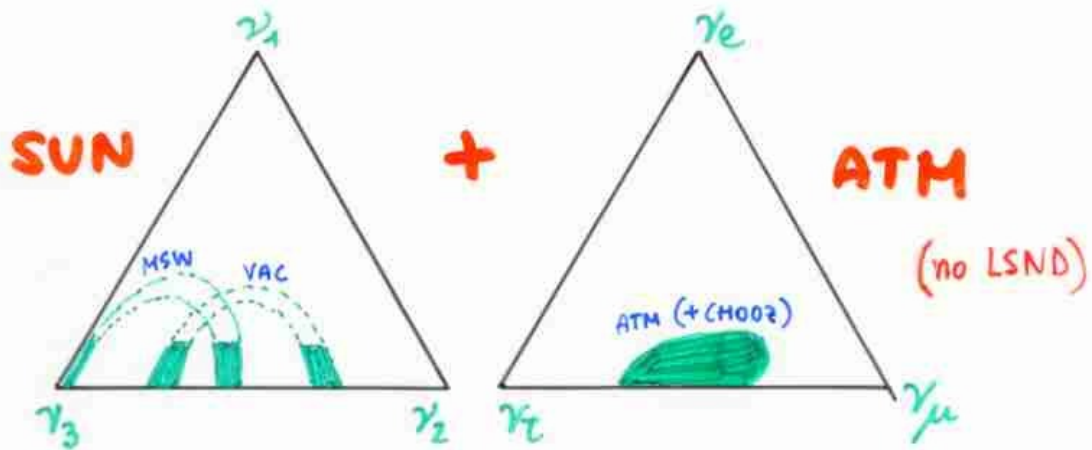
$$U_{\mu 3}^2 = \cos^2 \varphi \sin^2 \psi$$

$$U_{\tau 3}^2 = \cos^2 \varphi \cos^2 \psi$$

$$0 < \omega, \varphi, \psi < \pi/2 ; \delta_{CP} \text{ unobservable}$$

CONCLUSIONS :

ONLY TWO POSSIBLE "SOLUTIONS"
IN THREE-FLAVOR (HIERARCHICAL) SCHEMES :

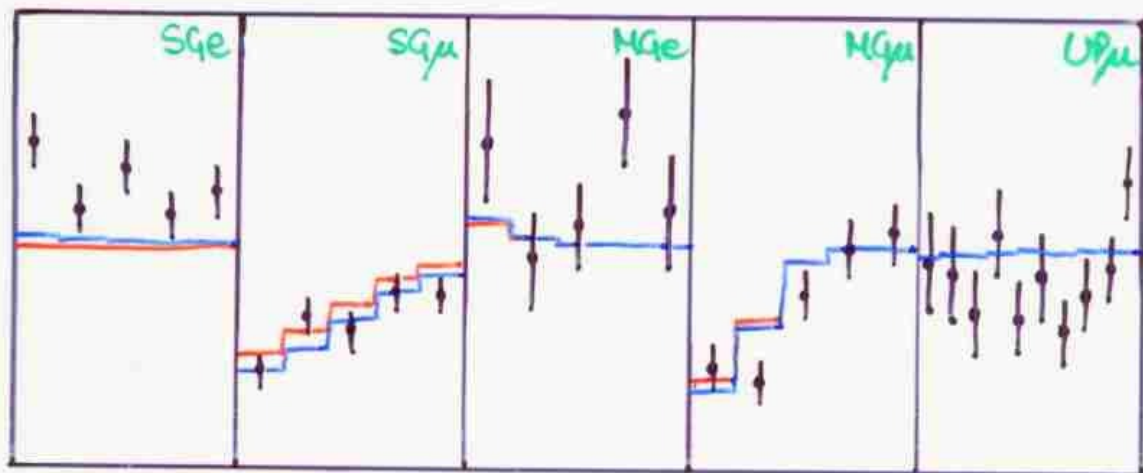
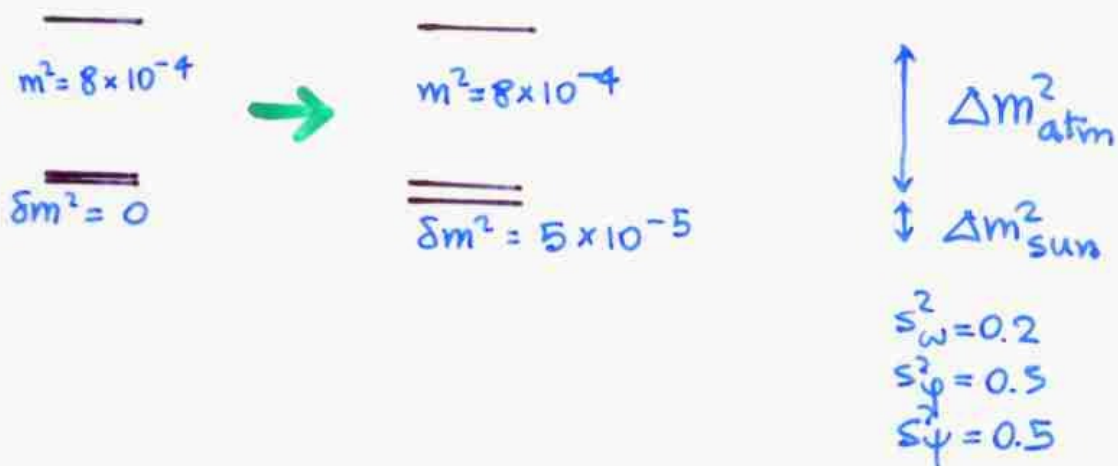


U_{e3} can be nonnegligible in case 1.

Tests : - Prove that U_{e3} is constrained by solar data (or atm. data) only

- Probe U_{e3} with lab. expts in the $\nu_{\mu} \rightarrow \nu_e$ channel

... THE NEXT FRONTIER :
 SOLAR \rightarrow MASS EFFECT
 IN ATMOSPHERIC NEUTRINO
 OSCILLATIONS ?



— $\delta m^2 = 0$
 — $\delta m^2 = 5 \times 10^{-5}$