Phenomenology of neutrinos for current and future experiments

Tokyo Metropolitan University

Osamu Yasuda

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1. Introduction

Framework of 3 flavor ν oscillation

Mixing matrix

$$\begin{pmatrix}\nu_e\\ \nu_\mu\\ \nu_\tau\end{pmatrix}=\begin{pmatrix}U_{e1} & U_{e2} & U_{e3}\\ U_{\mu1} & U_{\mu2} & U_{\mu3}\\ U_{\tau1} & U_{\tau2} & U_{\tau3}\end{pmatrix}\begin{pmatrix}\nu_1\\ \nu_2\\ \nu_3\end{pmatrix}$$

Functions of mixing angles θ_{12}, θ_{23}, θ_{13}, and CP phase δ

Information we have obtained so far:

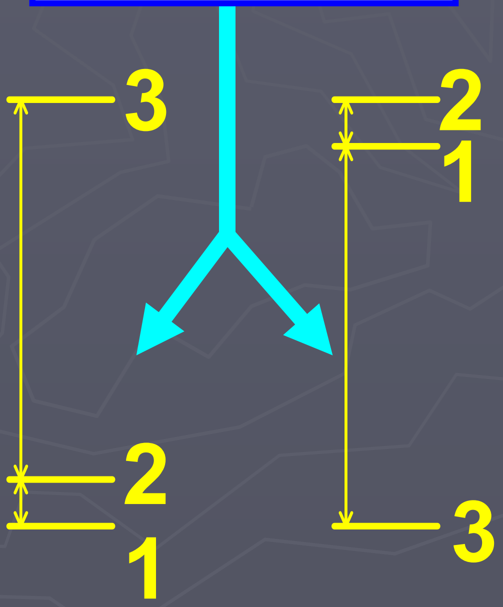
ν_{solar}+KamLAND (reactor) → $\theta_{12}\cong\frac{\pi}{6},\ \Delta m^2_{21}\cong 8\times10^{-5}\,\mathrm{eV}^2$

ν_{atm}+K2K,MINOS(accelerators) → $\theta_{23}\cong\frac{\pi}{4},\ |\Delta m^2_{32}|\cong 2.5\times10^{-3}\,\mathrm{eV}^2$

CHOOZ (reactor) → $|\theta_{13}|\lesssim\sqrt{0.15}/2$

$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \cong \begin{pmatrix} c_{12} & s_{12} & \varepsilon \\ -s_{12}/\sqrt{2} & c_{12}/\sqrt{2} & 1/\sqrt{2} \\ s_{12}/\sqrt{2} & -c_{12}/\sqrt{2} & 1/\sqrt{2} \end{pmatrix}$$

- **Both mass hierarchies are allowed**



normal hierarchy

$\Delta m^2_{32} > 0$

inverted hierarchy

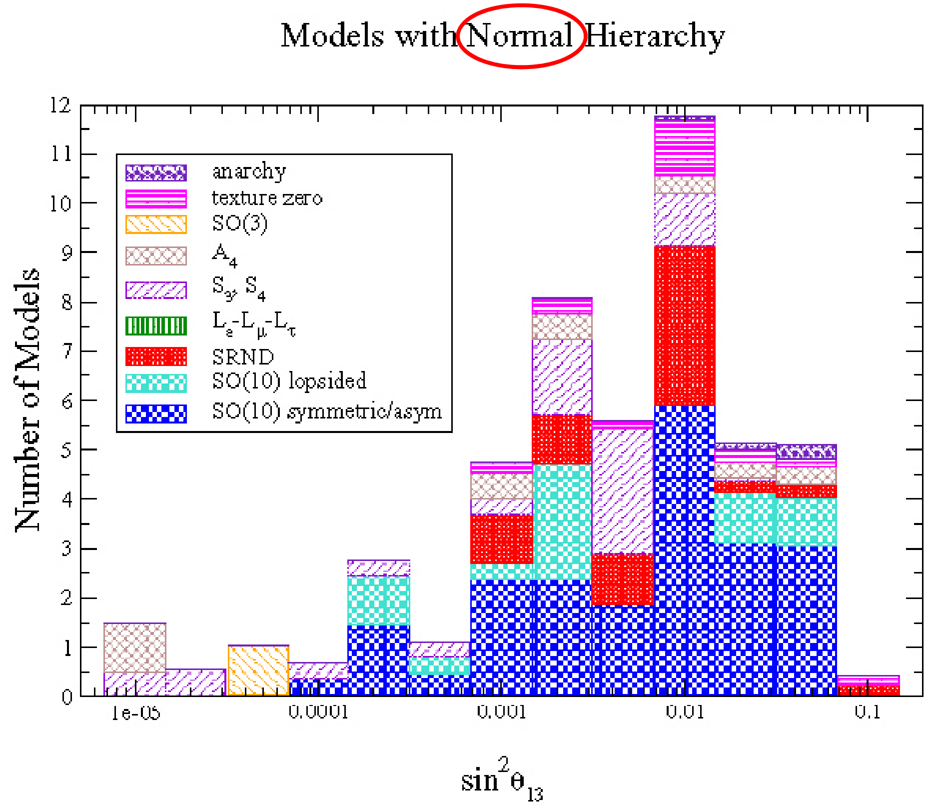
$\Delta m^2_{32} < 0$

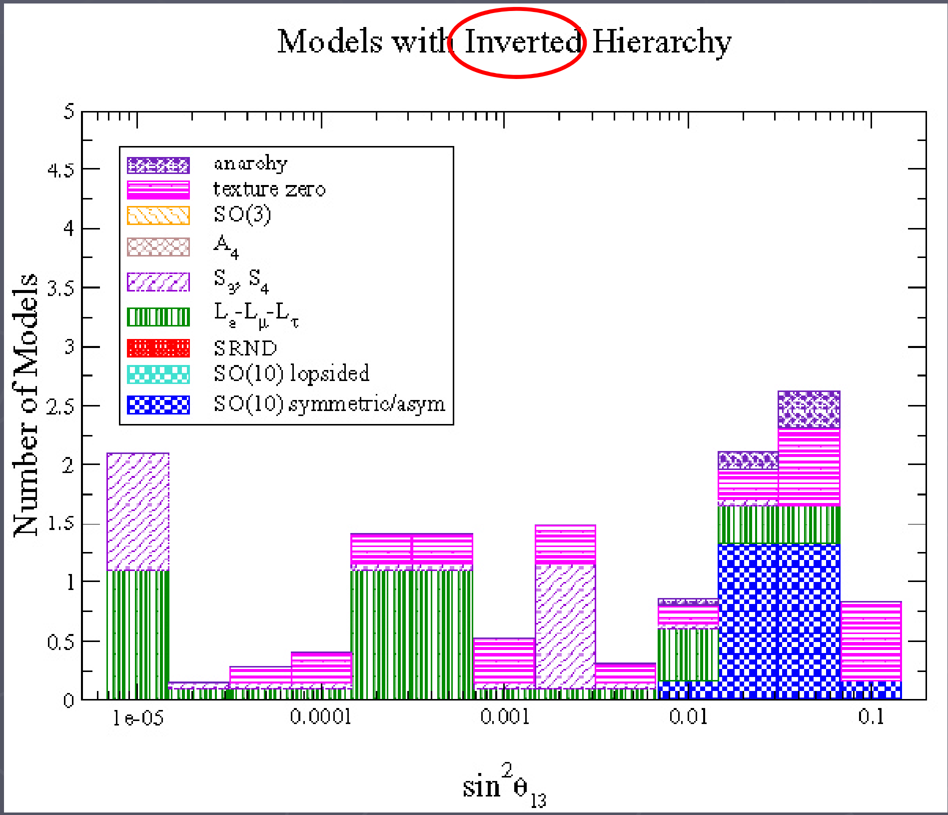
- θ_{13} :only upper bound is known
- δ :undetermined

Next task is to measure θ_{13}, sign(Δm^2_{31}) and δ.

A word on theory: Theoretical prediction for θ_{13}

Albright, Chen: Phys.Rev.D74:113006,2006





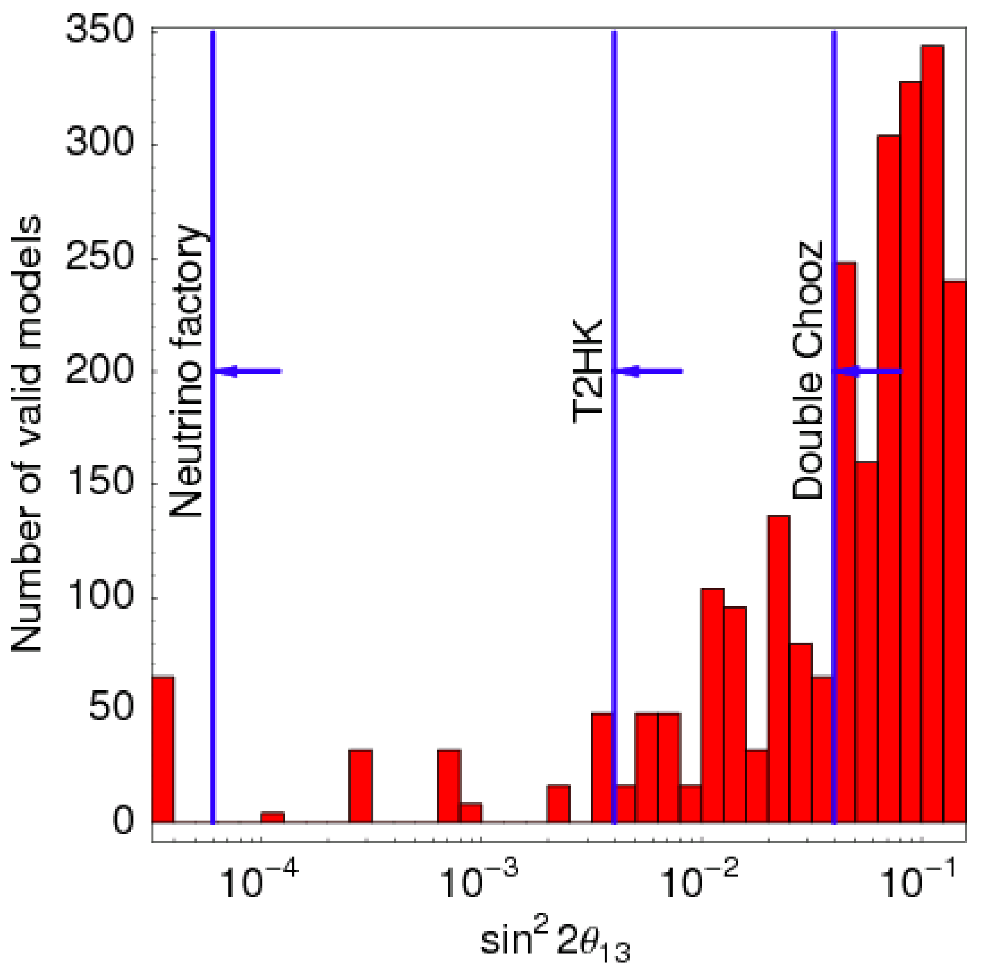
Depending on the mass hierarchy, the predictions differ.

Plentinger, Seidl, Winter: Nucl.Phys.B791:60-92,2008

Systematic generation of ν mass matrices by extended QLC

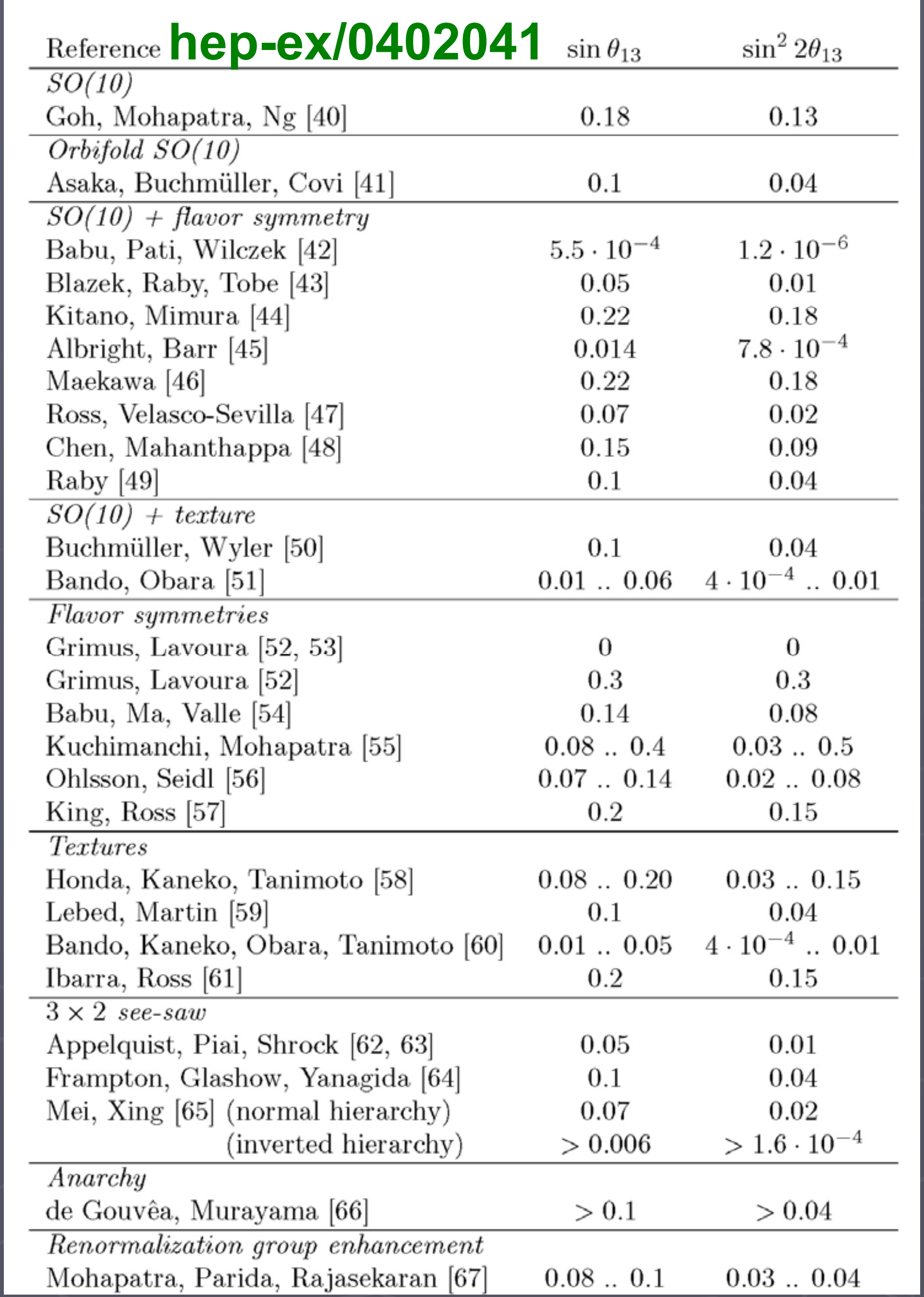
1981 textures !!!

- Parameter space analysis based on realizations
- Large θ_{13} preferred
- Compared to the GUT literature:
  Some realizations with very small $\sin^2 2\theta_{13}$ ~3.3 10^{-5}



All kinds of values of θ_{13} are predicted by theory, and it doesn't look like illuminating.

→ Theory is not yet developed enough to say something on mass & mixing of quarks & leptons.



hep-ex/0402041

Reference	$\sin\theta_{13}$	$\sin^2 2\theta_{13}$
SO(10)		
Goh, Mohapatra, Ng [40]	0.18	0.13
Orbifold SO(10)		
Asaka, Buchmüller, Covi [41]	0.1	0.04
SO(10) + flavor symmetry		
Babu, Pati, Wilczek [42]	$5.5 \cdot 10^{-4}$	$1.2 \cdot 10^{-6}$
Blazek, Raby, Tobe [43]	0.05	0.01
Kitano, Mimura [44]	0.22	0.18
Albright, Barr [45]	0.014	$7.8 \cdot 10^{-4}$
Maekawa [46]	0.22	0.18
Ross, Velasco-Sevilla [47]	0.07	0.02
Chen, Mahanthappa [48]	0.15	0.09
Raby [49]	0.1	0.04
SO(10) + texture		
Buchmüller, Wyler [50]	0.1	0.04
Bando, Obara [51]	0.01 .. 0.06	$4 \cdot 10^{-4}$ .. 0.01
Flavor symmetries		
Grimus, Lavoura [52, 53]	0	0
Grimus, Lavoura [52]	0.3	0.3
Babu, Ma, Valle [54]	0.14	0.08
Kuchimanchi, Mohapatra [55]	0.08 .. 0.4	0.03 .. 0.5
Ohlsson, Seidl [56]	0.07 .. 0.14	0.02 .. 0.08
King, Ross [57]	0.2	0.15
Textures		
Honda, Kaneko, Tanimoto [58]	0.08 .. 0.20	0.03 .. 0.15
Lebed, Martin [59]	0.1	0.04
Bando, Kaneko, Obara, Tanimoto [60]	0.01 .. 0.05	$4 \cdot 10^{-4}$ .. 0.01
Ibarra, Ross [61]	0.2	0.15
3×2 see-saw		
Appelquist, Piai, Shrock [62, 63]	0.05	0.01
Frampton, Glashow, Yanagida [64]	0.1	0.04
Mei, Xing [65] (normal hierarchy)	0.07	0.02
(inverted hierarchy)	> 0.006	$> 1.6 \cdot 10^{-4}$
Anarchy		
de Gouvêa, Murayama [66]	> 0.1	> 0.04
Renormalization group enhancement		
Mohapatra, Parida, Rajasekaran [67]	0.08 .. 0.1	0.03 .. 0.04

2. Future accelerator and reactor experiments

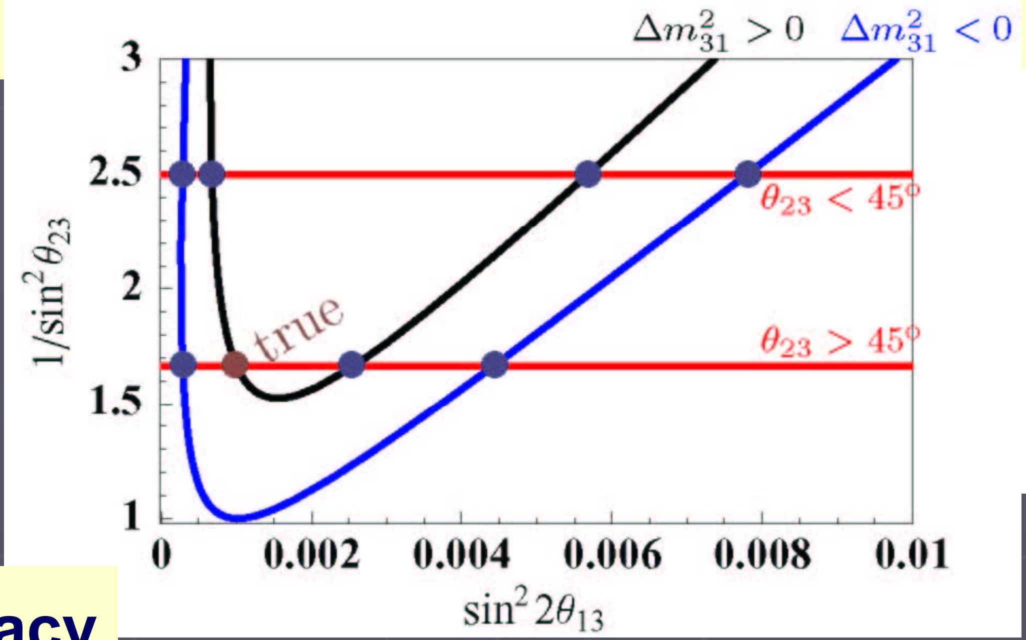
Most realistic way to measure θ_{13} , sign(Δm^2_{31}) and δ is long base line experiments by accelerators or reactors.

● One issue has to be taken into account for precise measurements:

Parameter degeneracy

Parameter degeneracy

Even if we know $P(\nu_\mu \to \nu_e)$ and $P(\overline{\nu_\mu} \to \overline{\nu_e})$ in a long baseline accelerator experiments with approximately monoenergetic neutrino beam, precise determination of θ_{13}, sign(Δm^2_{31}) and δ is difficult because of the 8-fold parameter degeneracy.



- intrinsic (δ, θ_{13}) degeneracy
- $\Delta m^2_{31} \Leftrightarrow -\Delta m^2_{31}$ degeneracy
- $\theta_{23} \Leftrightarrow \pi/2 - \theta_{23}$ degeneracy

To solve parameter degeneracy, combine the following:

(A) LBL measurement at $|\Delta m^2_{31}|L/4E = \pi/2$

$\rightarrow$ hyperbola shrinks to a straight line

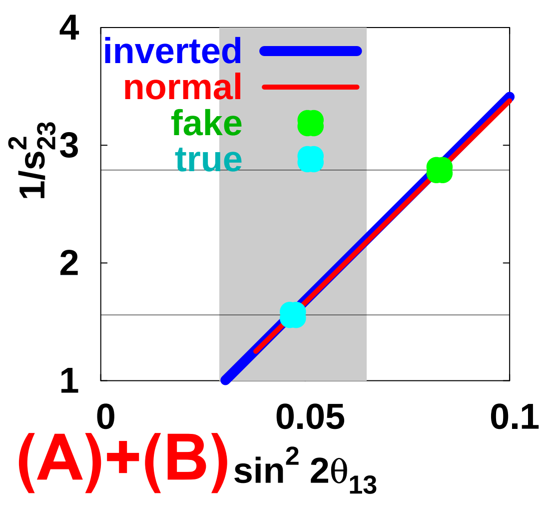
(B) reactor measurement of θ_{13} $\bar{\nu}_e \rightarrow \bar{\nu}_e$

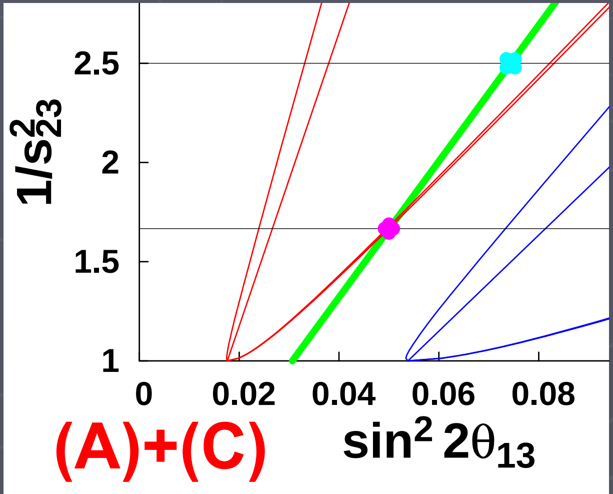
$\rightarrow$ depends only on θ_{13}

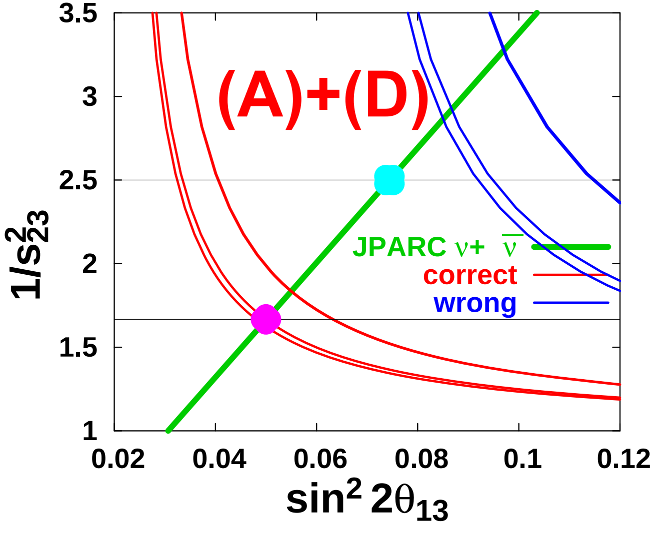
(C) LBL measurement of $\nu_\mu \rightarrow \nu_e$ (or $\nu_e \rightarrow \nu_\mu$)

with different L/E

(D) measurement of $\nu_e \rightarrow \nu_\tau$







Future accelerator LBL experiments

To perform precise measurements of θ_{13} and δ, one has to have a lot of numbers of events to improve statistical errors.

→We need **high intensity beam**

Candidates for high intensity beam in the future:

● (conventional) superbeam

$$\begin{cases} \pi^+ \to \mu^+ + \nu_\mu \\ \pi^- \to \mu^- + \bar{\nu}_\mu \end{cases}$$

$\nu_\mu \to \nu_e$

$\bar{\nu}_\mu \to \bar{\nu}_e$

● neutrino factory

μ in a storage ring

$$\begin{cases} \mu^+ \to e^+ + \nu_e + \bar{\nu}_\mu \\ \mu^- \to e^- + \bar{\nu}_e + \nu_\mu \end{cases}$$

$\nu_e \to \nu_\mu$

$\bar{\nu}_e \to \bar{\nu}_\mu$

● beta beam

RI in a storage ring

$$\begin{cases} {}^6_2\mathrm{He} \to {}^6_3\mathrm{Li} + e^- + \bar{\nu}_e \\ {}^{18}_{10}\mathrm{Ne} \to {}^{18}_{9}\mathrm{F} + e^+ + \nu_e \end{cases}$$

$\bar{\nu}_e \to \bar{\nu}_\mu$

$\nu_e \to \nu_\mu$

Future LBL exp. (under construction / proposed)

- superbeam

Kobayashi's talk

T2K phase I (2009-, 0.75MW, E~1GeV, L=295km)

T2K phase II (4MW+HK, E~1GeV, L=295km)

T2KK (JAERI→HK&Korea, E~1GeV, L=295km&1000km)

NOvA (FNAL→ Ash River (MN), E~2GeV, L=810km)

VBLNO (BNL→Homestake, E~2GeV, L>2500km)

SPL (CERN→Frejus, E~0.25GeV, L=130km)

- neutrino factory (E_ν<50GeV, L~3000km)
- beta beam (E_ν=0.5-1.5GeV, L~130km)

Future reactor experiments (E~4MeV, L~2km)

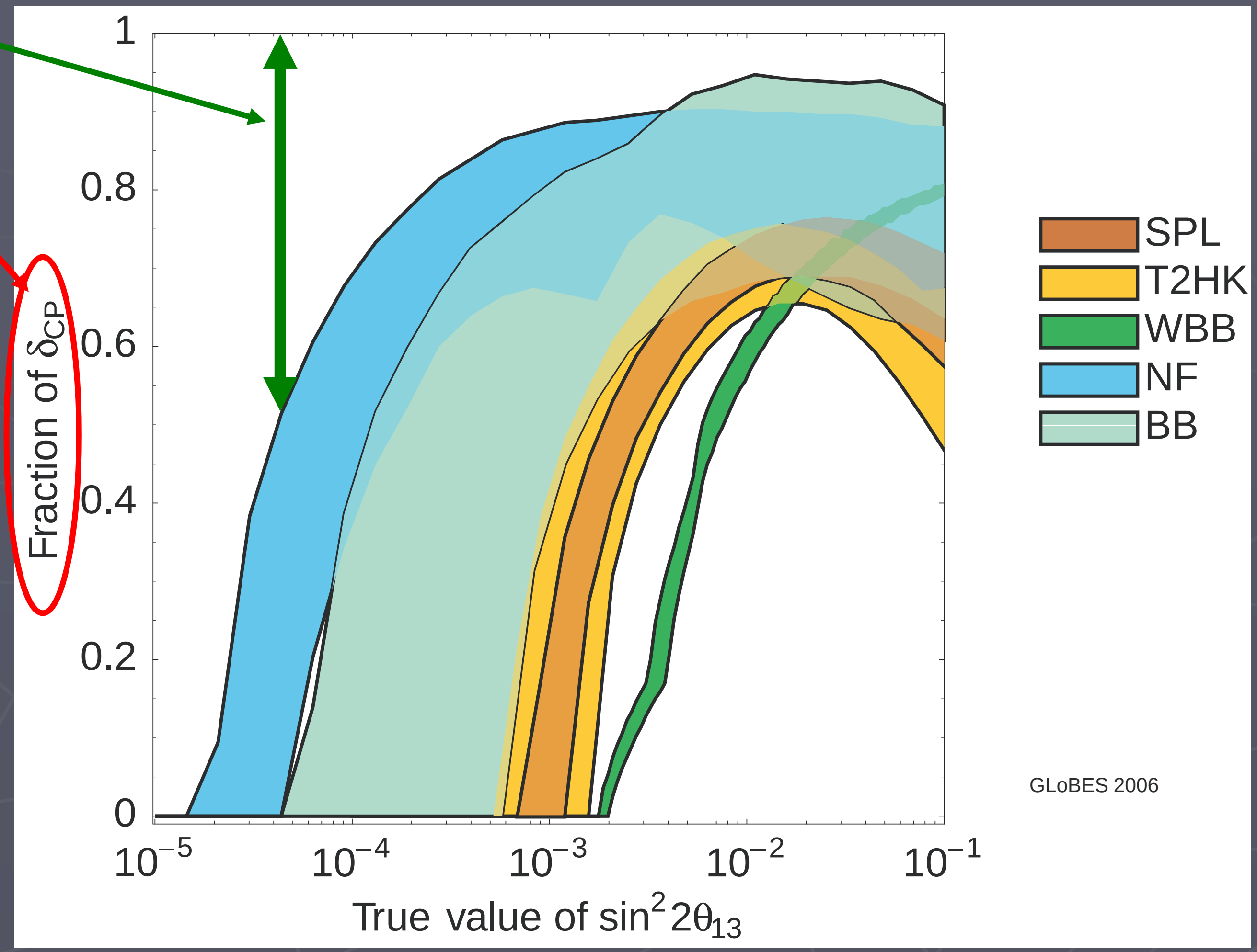
Complementary to accelerator exp.

Double CHOOZ (France) , Daya Bay (China), Reno (Korea), Angra (Brazil)

Kuze's talk

Joo's talk

sensitivity to the CP phase δ of future experiments



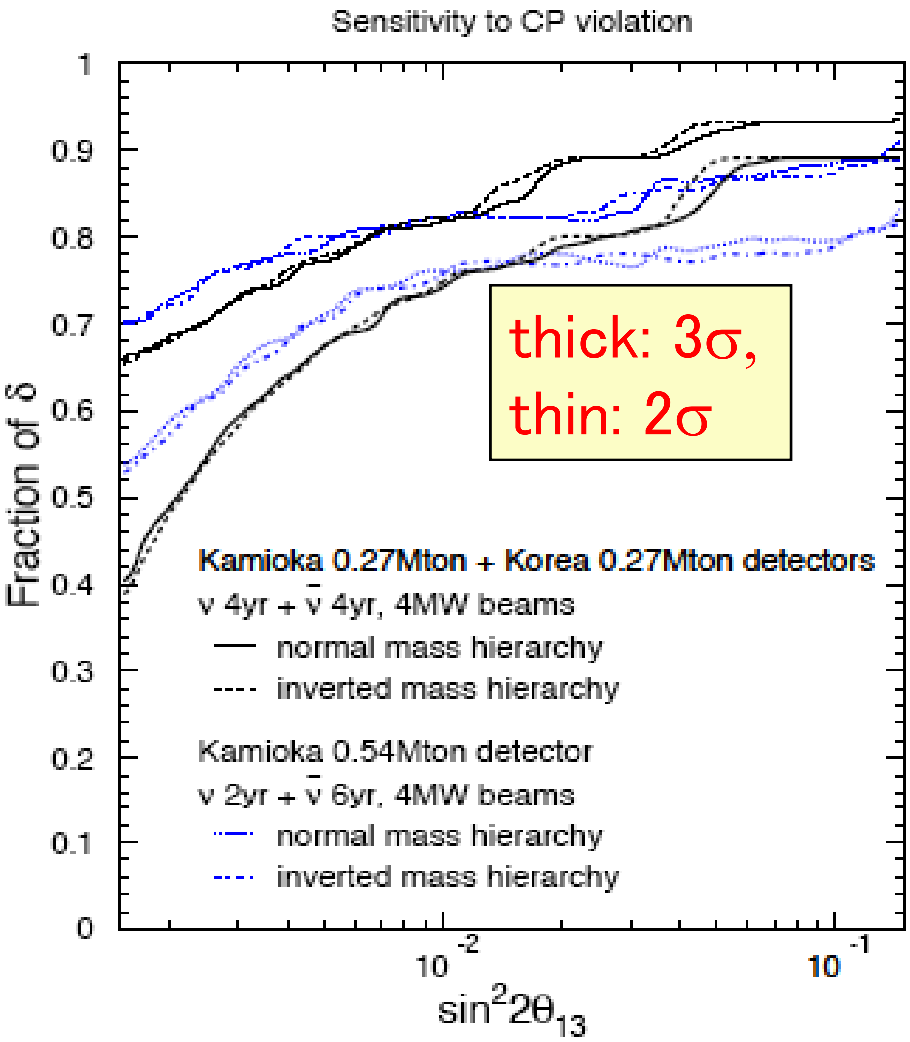
error of δ

1-(error of δ)

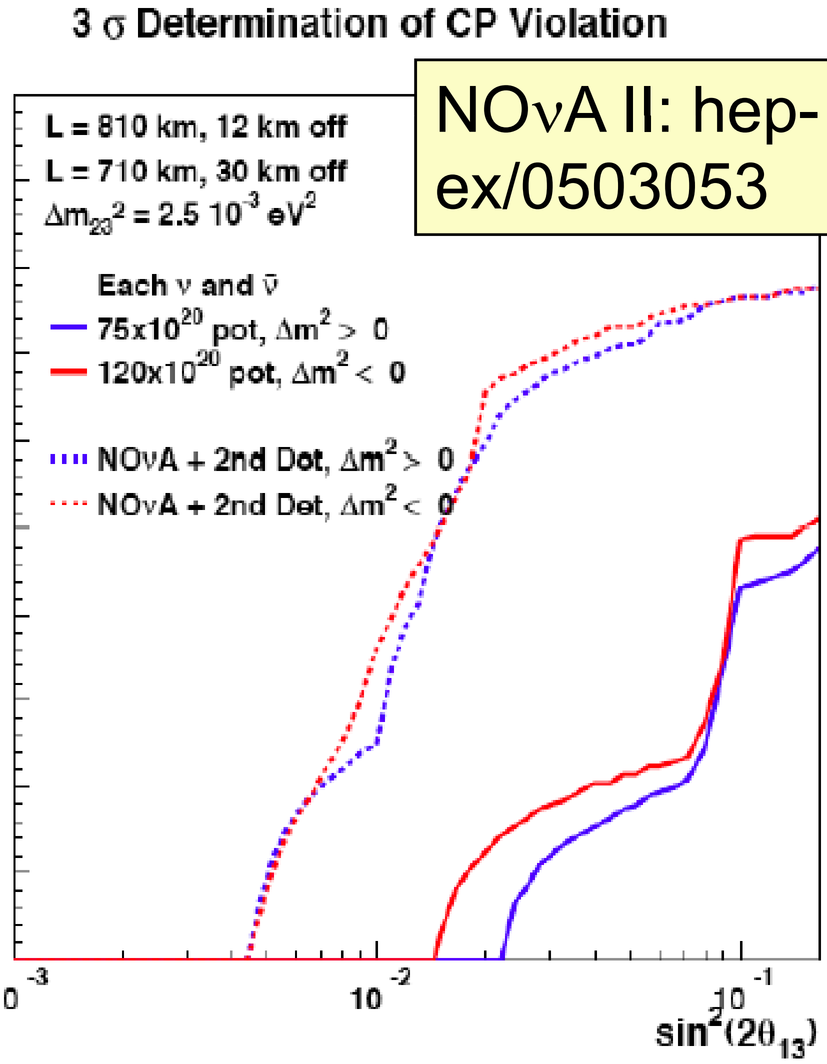
T2HK = Tokai to HyperKamiokande

ISS report: arXiv:0710.4947 [hep-ph]

T2KK vs. NOνA; CP



thick: 3σ, thin: 2σ



NOνA II: hep-ex/0503053

Minakata@2nd T2KK workshop ('06)

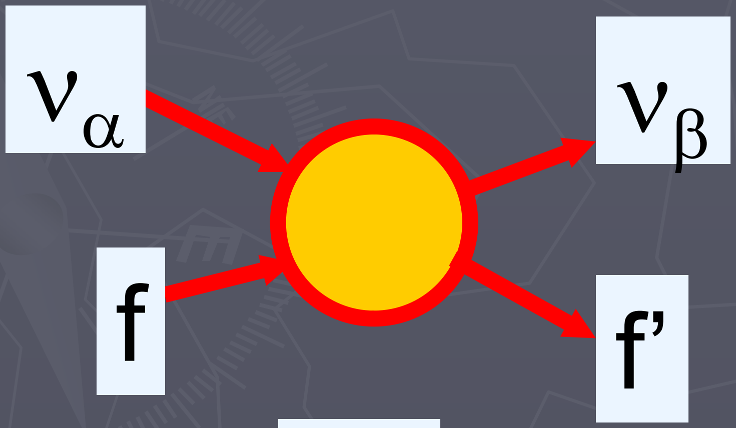
3. Deviation from standard 3 flavor framework

Just like at B factories, high precision measurements of ν oscillation can be used also to probe physics beyond SM by looking at deviation from SM+masssive ν.

Here I will discuss the following topics:
(1) New physics (NP) (exotic interactions)
(2) Violation of unitarity (like at a B factory)
(3) Sterile neutrinos

(1) New physics (NP) (exotic interactions)

Flavors are not necessarily conserved in these interactions: $\alpha=\beta$ & $\alpha\neq\beta$



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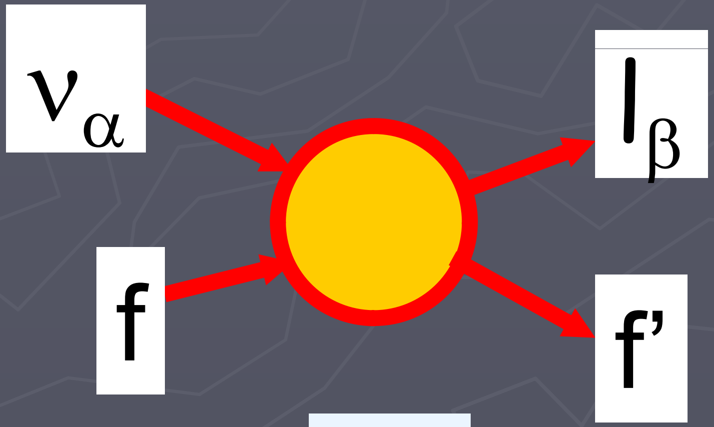


CC

Effects of New Physics on ν oscillations

$A \equiv \sqrt{2}G_F N_e$

$$\mathcal{A}_{SM} \equiv A \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$P(\nu_\alpha \to \nu_\beta) = \left| \left[\boxed{U^d} \tilde{U} \exp\left\{ -i\, \mathrm{diag}(\tilde{E}_j) L \right\} \tilde{U}^{-1} \boxed{(U^s)^{-1}} \right]_{\beta\alpha} \right|^2$$

with

$$U\, \mathrm{diag}(E_j)\, U^{-1} + \mathcal{A}_{SM} \boxed{+ \mathcal{A}_{NP}} \equiv \tilde{U}\, \mathrm{diag}(\tilde{E}_j)\, \tilde{U}^{-1}$$

source

$|(U^s\text{-}1)_{\alpha\beta}| < O(10^{-2})$

$$\begin{pmatrix} \nu_e^S \\ \nu_\mu^S \\ \nu_\tau^S \end{pmatrix} = \boxed{U^s} U_{MNS} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

propagation

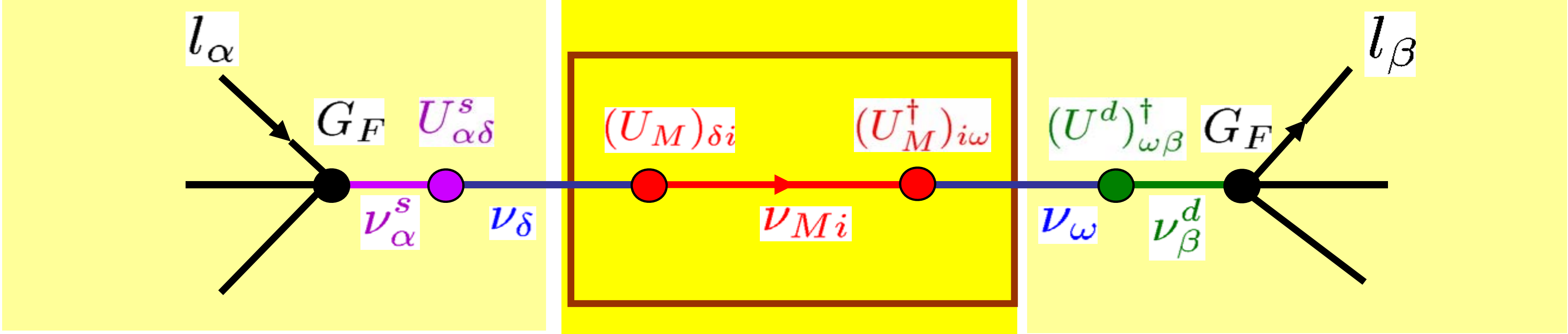
ε_{ee} , $\varepsilon_{e\tau}$, $\varepsilon_{\tau\tau}$ ~O(1)

$$A \begin{pmatrix} \epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu}^* & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau}^* & \epsilon_{\mu\tau}^* & \epsilon_{\tau\tau} \end{pmatrix}$$

detection

$|(U^d\text{-}1)_{\alpha\beta}| < O(10^{-2})$

$$\begin{pmatrix} \nu_e^d \\ \nu_\mu^d \\ \nu_\tau^d \end{pmatrix} = \boxed{U^d} U_{MNS} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$



(i) Effects of New Physics at source and detector

Deviation from the standard form is small:

Grossman (PLB359:141,1995)

$$|(U^{s}-1)_{\alpha\beta}| < O(10^{-2}),\ |(U^{d}-1)_{\alpha\beta}| < O(10^{-2})$$

(ii) New Physics effects in propagation

1. Constraints from various ν experiments:

Davidson et al (JHEP 0303:011,2003)

$$\begin{pmatrix} \boxed{-3 \lesssim \epsilon_{ee} \lesssim 2} & |\epsilon_{e\mu}| \lesssim 0.5 & \boxed{|\epsilon_{e\tau}| \lesssim 1.5} \\ |\epsilon_{e\mu}| \lesssim 0.5 & |\epsilon_{\mu\mu}| \lesssim 0.05 & |\epsilon_{\mu\tau}| \lesssim 0.15 \\ \boxed{|\epsilon_{e\tau}| \lesssim 1.5} & |\epsilon_{\mu\tau}| \lesssim 0.15 & \boxed{|\epsilon_{\tau\tau}| \lesssim 6} \end{pmatrix}$$

2. Constraints from atmospheric neutrinos:

Friedland-Lunardini (Phys.Rev.D72:053009,2005)

ε_{ee} , $\varepsilon_{e\tau}$, $\varepsilon_{\tau\tau}$ ~O(1) are consistent with ν_{atm} data, provided

➡ **Deviation could be large**

$$\boxed{\epsilon_{\tau\tau} \sim \frac{|\epsilon_{e\tau}|^2}{1+\epsilon_{ee}}} \qquad \boxed{0 \leq |\epsilon_{e\tau}| \lesssim 1 + \epsilon_{ee}} \qquad \boxed{-1 \lesssim \epsilon_{ee} \lesssim 1.5}$$

In general:

NP effects at production and at detection becomes important when L is smaller

$|\varepsilon_{\alpha\beta}| < O(10^{-2})$

$$P(\nu_\alpha \to \nu_\beta) \to \left| \left[U^d (U^s)^{-1} \right]_{\beta\,\alpha} \right|^2$$

i.e., no BG from osc. in the limit of $L \to 0$

⇒ Experiments with a shorter baseline are advantageous

NP effects in propagation becomes important when L is larger

$\varepsilon_{\alpha\beta}$ can be of O(1)

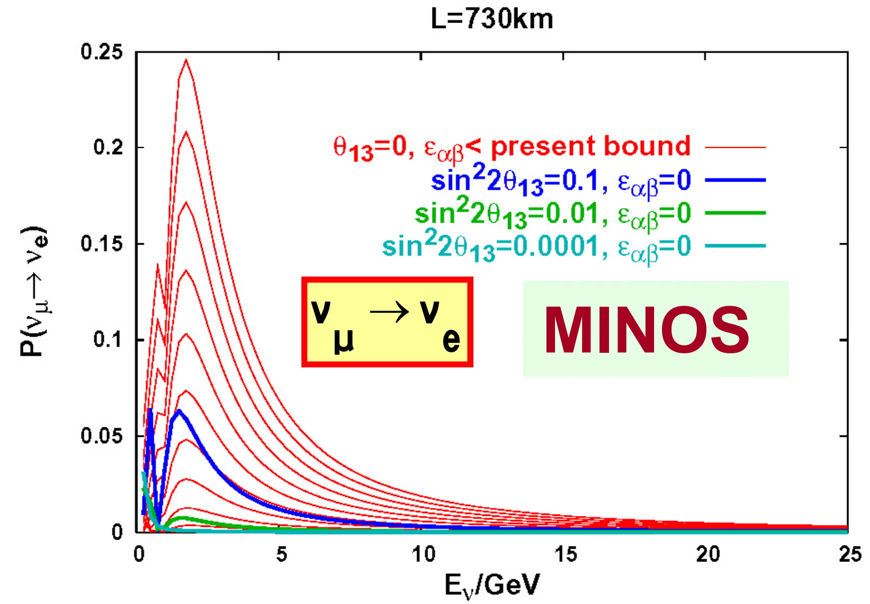
because $A L\, \varepsilon_{\alpha\beta} \sim \varepsilon_{\alpha\beta}\, (L/2000\text{km})$

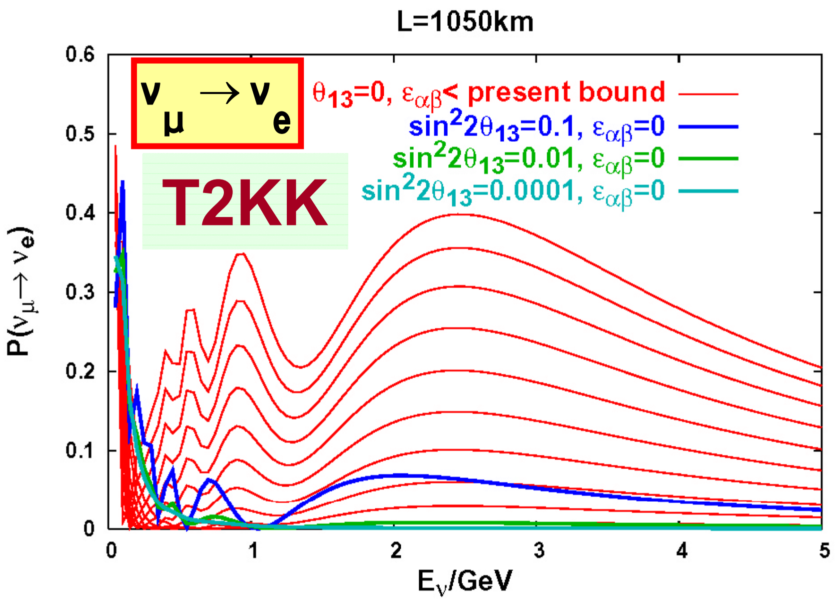
⇒ Experiments with a longer baseline are advantageous to see them;

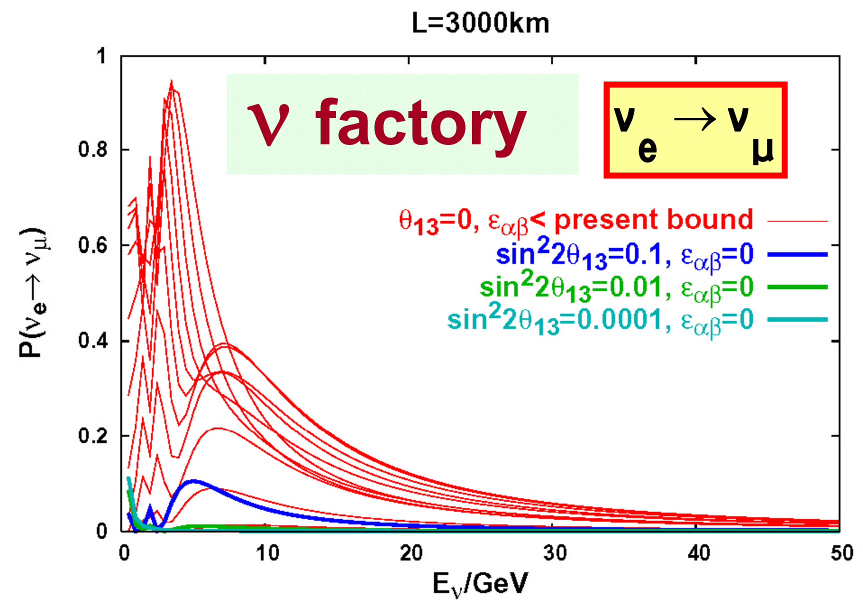
Experiments with very short baselines (e.g., reactors) have no BG from NP effects in propagation

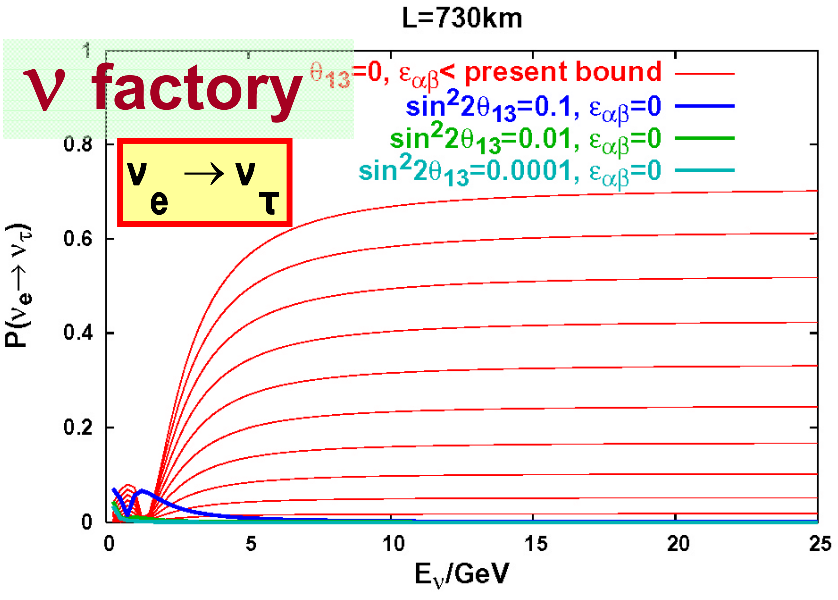
Potentially expected New Physics effects in propagation

NK,HS,OY, hep-ph/0606013







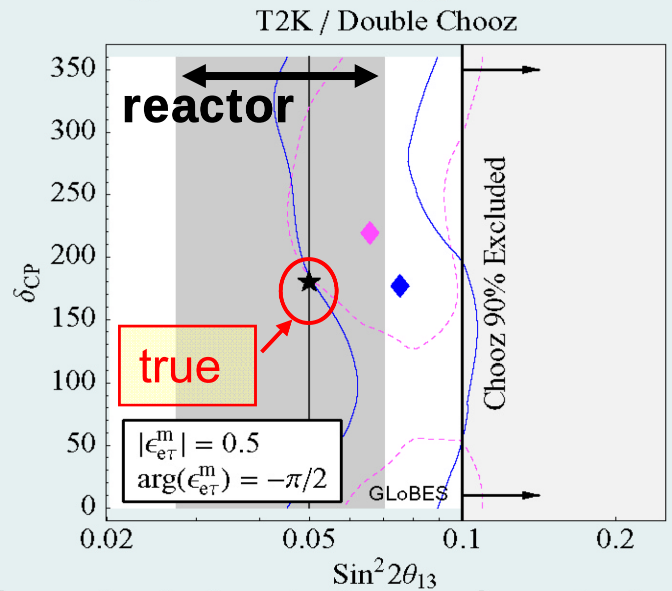


Confusion in the presence of New Physics

Kopp, Lindner, Ota, Sato, PRD77:013007,2008

Mismatch

Best fit points for reactor & T2K don't agree



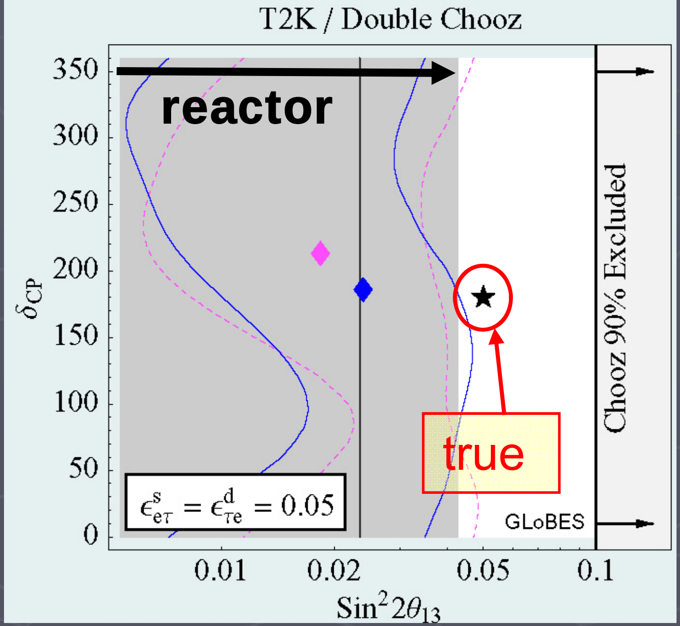
ε^m: NP in propagation

◆ Best fit of T2K with Normal Hierarchy

◆ Best fit of T2K with Inverted Hierarchy

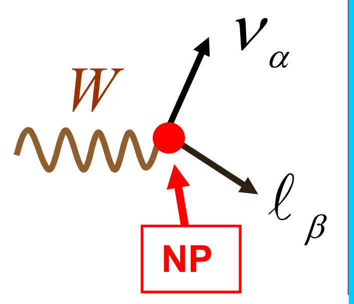
Offset

Best fit points for two agree but miss the true point



$\varepsilon^{s,d}$:NP at source or detection

(2) Unitarity violation due to heavy ν



minimal non-unitary model

$U \rightarrow N$ (non-unitary)

$NN^\dagger - 1$: deviation from unitarity

mostly from rare decays

$$|NN^\dagger| \approx \begin{pmatrix} 0.994 \pm 0.005 & < 7.1 \cdot 10^{-5} & < 1.6 \cdot 10^{-2} \\ < 7.1 \cdot 10^{-5} & 0.995 \pm 0.005 & < 1.0 \cdot 10^{-2} \\ < 1.6 \cdot 10^{-2} & < 1.0 \cdot 10^{-2} & 0.995 \pm 0.005 \end{pmatrix}$$

90% C.L.

Constraints on unitary violation are strong:

$$\boxed{|(NN^\dagger - 1)_{\alpha\beta}| < O(10^{-2})}$$

Minimal case: Antusch, Biggio, Fernandez-Martinez, Gavela, Lopez-Pavon: JHEP 0610:084,2006.

Non-minimal case: Abada, Biggio, Bonnet, Gavela, Hambye: JHEP 0712:061,2007

Even stronger constraints → more difficult to detect

**Unitarity violation is similar to NP effects at production and at detection →
becomes important when L is smaller**

$$P(\nu_\alpha \to \nu_\beta) \to \left| \left[(1+\eta)\, U \exp\left[-i\, \mathrm{diag}\,(E_j)\right] U^{-1} (1+\eta)^{-1} \right]_{\beta\,\alpha} \right|^2$$

$$\eta \equiv (NN^\dagger - 1)/2$$

$$\boxed{|\eta_{\alpha\beta}| < O(10^{-2})}$$

$$U^s \to 1+\eta/2,\ U^d \to 1+\eta/2$$

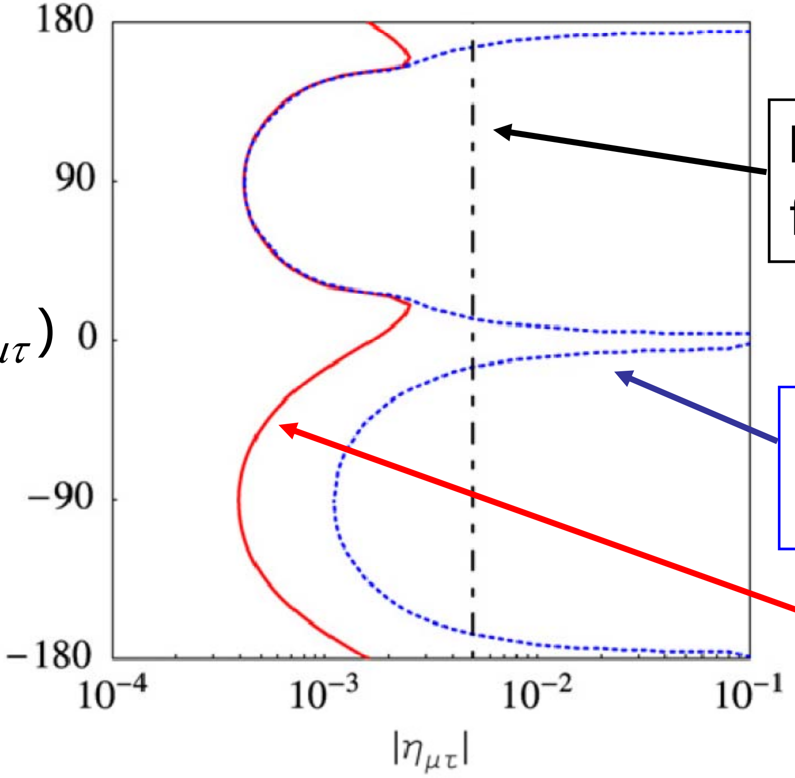
Experiments with a shorter baseline are advantageous

Neutrino factory with a baseline L=130km

Fernandez-Martinez, Gavela, Lopez-Pavon, OY, Phys.Lett.B649:427-435,2007

Phase of N

$\nu_\mu \leftrightarrow \nu_\tau$



Present bound from $\tau \longrightarrow \mu\,\gamma$

Sensitivity to $\arg(\eta_{\mu\tau})$

Sensitivity to $|\eta_{\mu\tau}|$

For non-trivial $\arg(\eta_{\mu\tau})$, one order of magnitude improvement for $\eta_{\mu\tau}$

(3A) Sterile neutrinos assuming LSND

$\Delta m^2_{\text{sol}} \sim 10^{-4}\,\text{eV}^2$ $\Delta m^2_{\text{atm}} \sim 10^{-3}\,\text{eV}^2$ $\Delta m^2_{\text{LSND}} \sim O(1)\,\text{eV}^2$

at least one ν_s is required

LSND$(\bar{\nu}_\mu \to \bar{\nu}_e)$: affirmative

MiniBOONE$(\nu_\mu \to \nu_e)$: negative

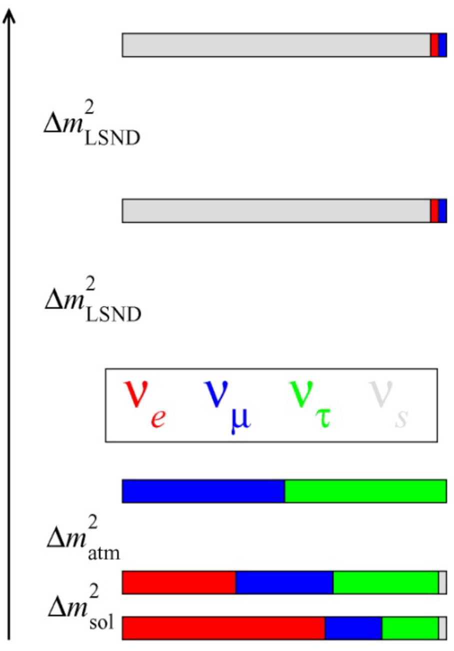
with one more ν_s difference between ν & anti-ν may offer a promising fit

(3+2)-scheme with CP phase δ

$$\begin{aligned} P_{\nu_\mu \to \nu_e} &= 4\,|U_{e4}|^2|U_{\mu 4}|^2 \sin^2 \phi_{41} \\ &+ 4\,|U_{e5}|^2|U_{\mu 5}|^2 \sin^2 \phi_{51} \\ &+ 8\,|U_{e4}U_{\mu 4}U_{e5}U_{\mu 5}| \sin\phi_{41} \sin\phi_{51} \cos(\phi_{54} - \delta) \end{aligned}$$

with the definitions

$$\phi_{ij} \equiv \frac{\Delta m^2_{ij} L}{4E}\,, \qquad \boxed{\delta \equiv \arg\left(U^*_{e4}U_{\mu 4}U_{e5}U^*_{\mu 5}\right)}\,.$$



Schwetz-Mangold@nufact07

$\varphi_{54}^{best} = 1.64\pi$

$$\Delta m^2_{41} = 0.89 \text{ eV}^2$$
$$\Delta m^2_{51} = 6.49 \text{ eV}^2$$

$$\chi^2_{\text{min}} = 94.5/(107-7)$$

Karagiorgi@nufact07

$\chi^2/ndf = 146.7/156$
gof=69%
$\varphi_{54}^{best} = 1.74\pi$

- (3+2) schemes
  - offer the possiblity of CP violation to reconcile LSND and MiniBooNE,
  - but there is tension between appearance and disappearance data (3σ, 4σ for MB300)

$\bar{\nu}_e \rightarrow \bar{\nu}_e$: Bugey

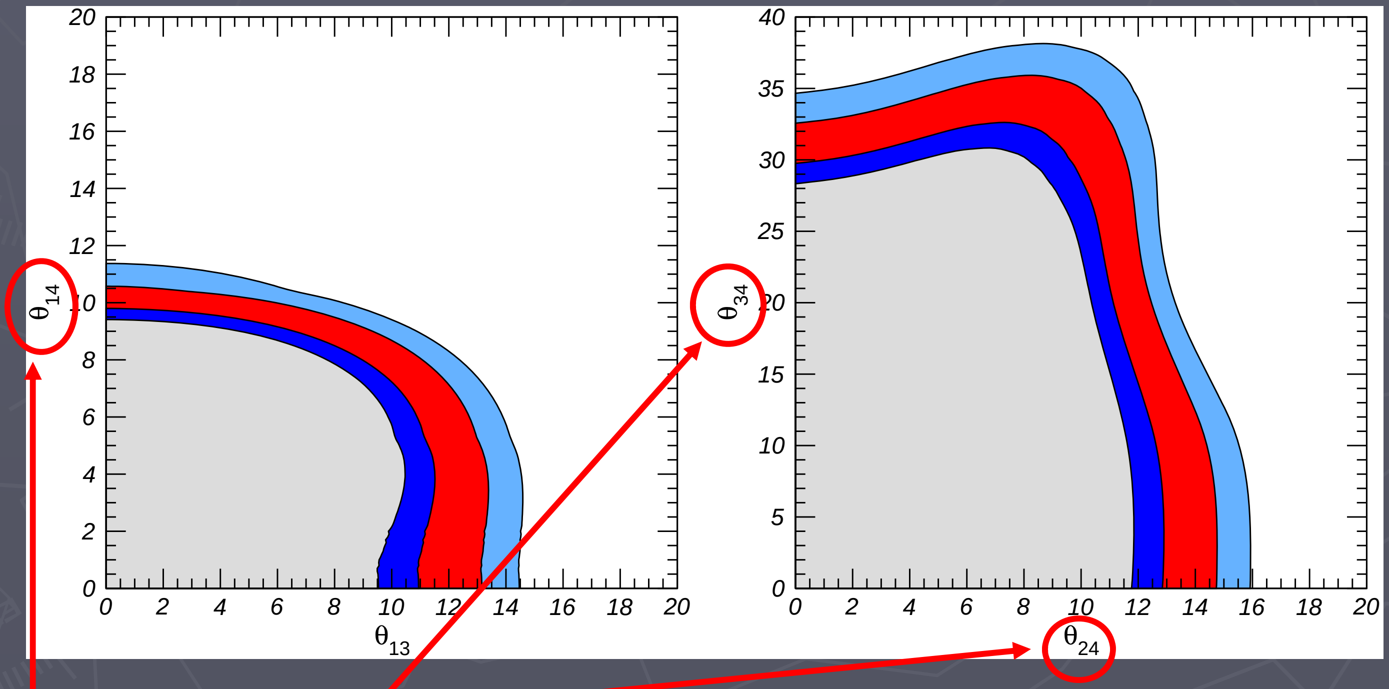
$\nu_\mu \rightarrow \nu_\mu$: CDHSW

So from the global fit, (3+2) is probably not a promising scheme…

(3B) Sterile neutrinos w/o assuming LSND

Without assuming LSND and imposing all the negative constraints one can still consider (3+1)-scheme

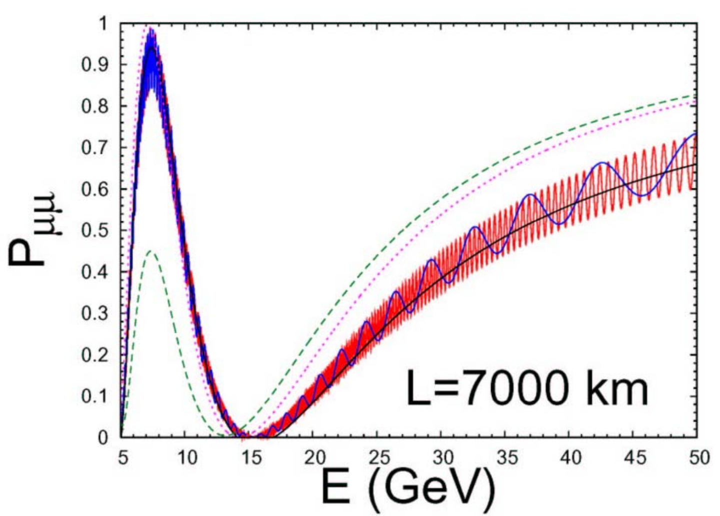
Donini, Maltoni, Meloni, Migliozzi, Terranova, JHEP 0712:013,2007

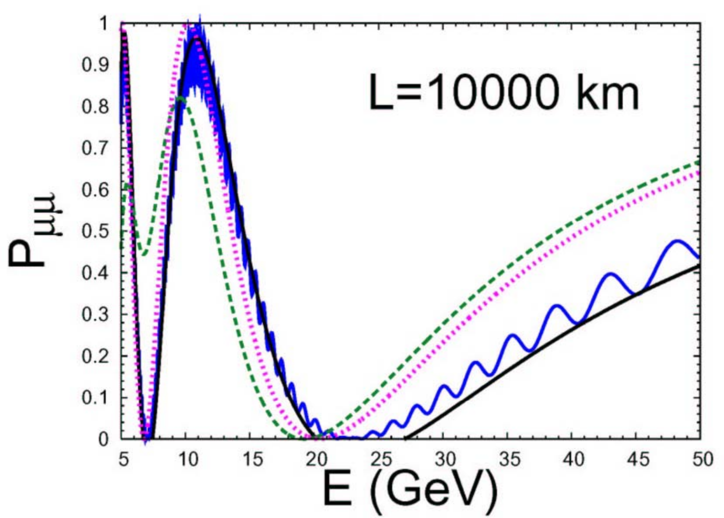


θ_{14} , θ_{24} , θ_{34} : angles which appear only in 4ν scenario

Potentially expected effects of sterile neutrinos at longer baseline

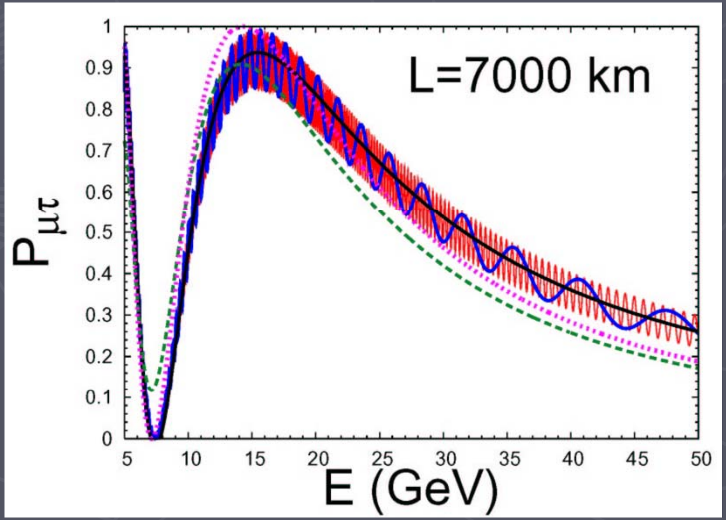
Dighe, Ray: Phys.Rev.D76:113001,2007

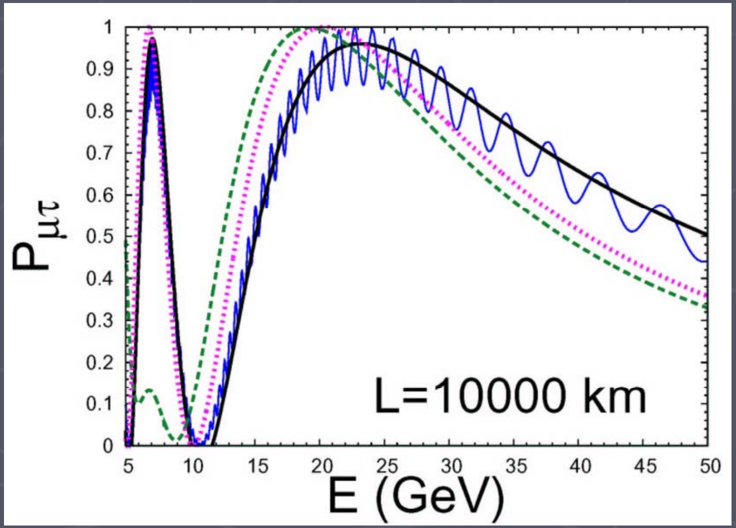




dashed lines: N=3

solid lines: N=4





4. Summary

- A brief review was given on the known parameters in SM+massive ν. Efforts to determine the unknown parameters (θ_{13} , δ , sign(Δm^2_{31})) in the future experiments were described.
- The future neutrino experiments with high precision will be able to see deviation from SM such as non-standard interactions, unitarity violation, sterile neutrinos, etc.
- Beyond SM+massive ν, there are still a lot of things to be worked out: sensitivities to NP, optimization to NP, degeneracy in the presence of NP, etc.
- Accelerator and reactor experiments are complementary not only in SM+massive ν but also in Beyond SM+massive ν.