

Neutrino oscillation phenomenology

Tokyo Metropolitan University

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**INO-KEK meeting
27 November 2007@KEK**

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1. Introduction

Framework of 3 flavor ν oscillation

Mixing matrix

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

Functions of mixing angles θ_{12} , θ_{23} , θ_{13} , and CP phase δ

Information we have obtained so far:

ν_{solar} + KamLAND (reactor)

$$\theta_{12} \cong \frac{\pi}{6}, \Delta m_{21}^2 \cong 8 \times 10^{-5} \text{ eV}^2$$

ν_{atm} + K2K, MINOS (accelerators)

$$\theta_{23} \cong \frac{\pi}{4}, |\Delta m_{32}^2| \cong 2.5 \times 10^{-3} \text{ eV}^2$$

CHOOZ (reactor)

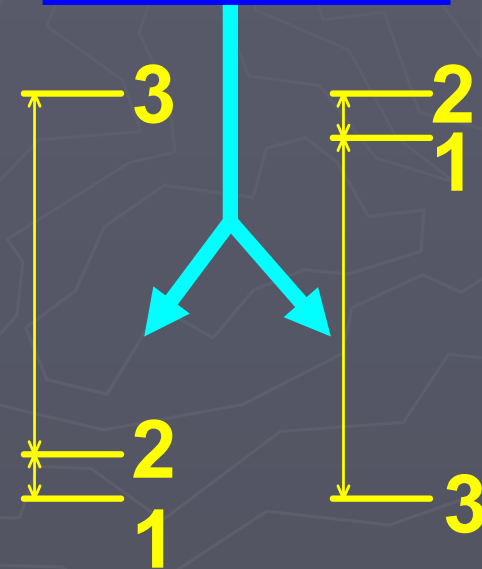
$$|\theta_{13}| \leq \sqrt{0.15}/2$$

$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \cong \begin{pmatrix} \mathbf{c}_{12} & \mathbf{s}_{12} & \boldsymbol{\varepsilon} \\ -\mathbf{s}_{12}/\sqrt{2} & \mathbf{c}_{12}/\sqrt{2} & 1/\sqrt{2} \\ \mathbf{s}_{12}/\sqrt{2} & -\mathbf{c}_{12}/\sqrt{2} & 1/\sqrt{2} \end{pmatrix}$$

- θ_{13} : only upper bound is known
- δ : undetermined

Next task is to measure θ_{13} , $\text{sign}(\Delta m^2_{31})$ and δ .

- Both mass hierarchies are allowed



normal hierarchy

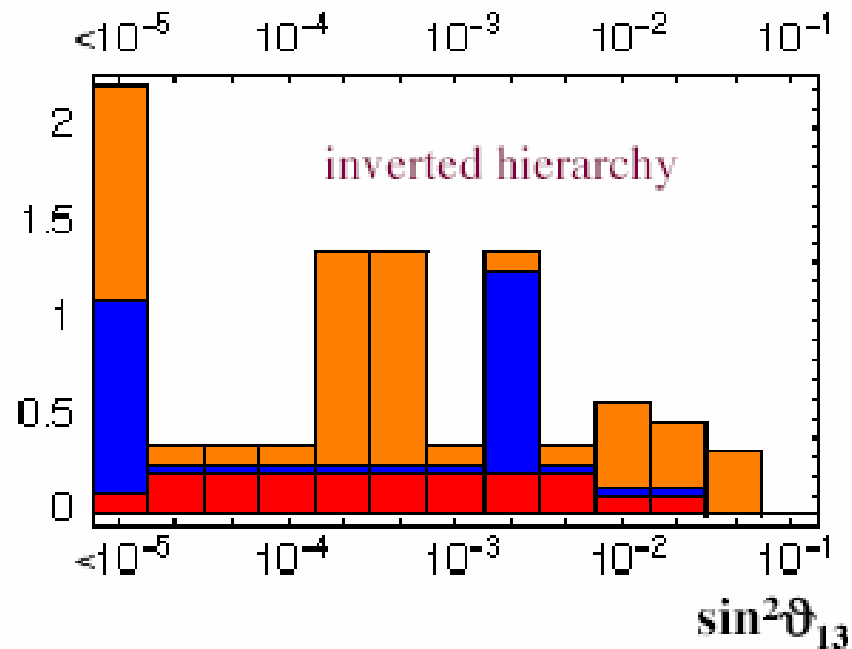
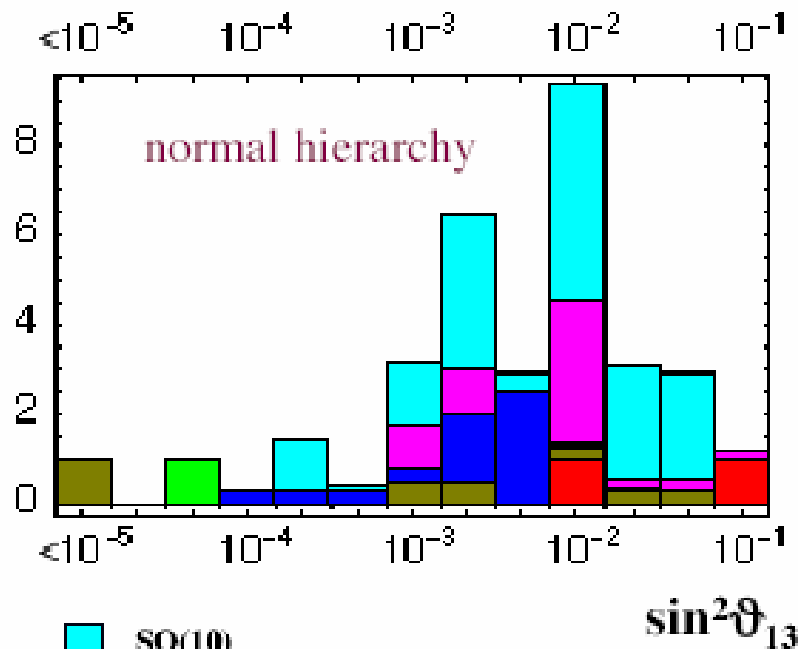
$$\Delta m^2_{32} > 0$$

inverted hierarchy

$$\Delta m^2_{32} < 0$$

Theoretical prediction for θ_{13}

Chen: ISS 3rd plenary ('06) @ RAL



- SO(10)
- SRND
- $L_e - L_\mu - L_\tau$
- S3, S4
- A4
- SO(3)
- Texture

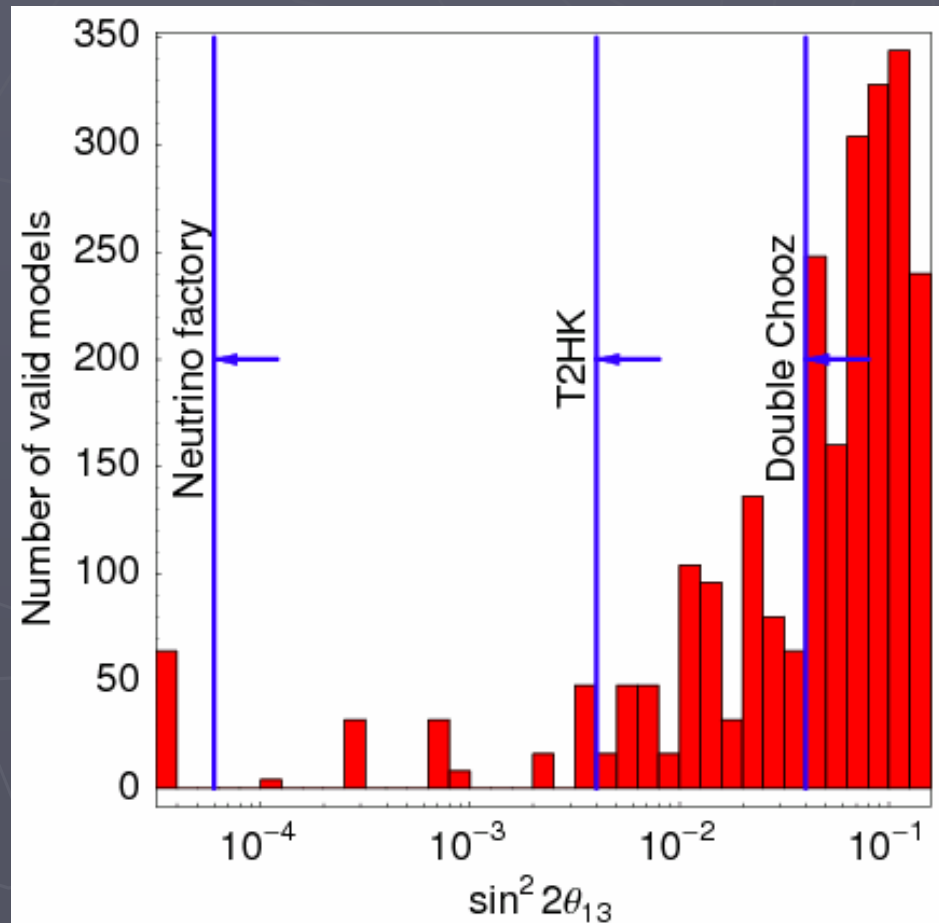
C. Albright & M.-C.C. (2006)

Winter: nufact07 @ Okayama

Systematic generation of ν mass matrices by extended QLC

- ▶ Parameter space analysis based on realizations
- ▶ Large θ_{13} preferred
- ▶ Compared to the GUT literature:
Some realizations with very small $\sin^2 2\theta_{13} \sim 3.3 \cdot 10^{-5}$

(Plentinger, Seidl, Winter,
hep-ph/0612169)



All kinds of values of θ_{13} are predicted by theory, and it doesn't look like illuminating.

→ Theory is not yet developed enough to say something on mass & mixing of quarks & leptons.

Reference hep-ex/0402041	$\sin \theta_{13}$	$\sin^2 2\theta_{13}$
<i>SO(10)</i>		
Goh, Mohapatra, Ng [40]	0.18	0.13
<i>Orbifold SO(10)</i>		
Asaka, Buchmüller, Covi [41]	0.1	0.04
<i>SO(10) + flavor symmetry</i>		
Babu, Pati, Wilczek [42]	$5.5 \cdot 10^{-4}$	$1.2 \cdot 10^{-6}$
Blazek, Raby, Tobe [43]	0.05	0.01
Kitano, Mimura [44]	0.22	0.18
Albright, Barr [45]	0.014	$7.8 \cdot 10^{-4}$
Maekawa [46]	0.22	0.18
Ross, Velasco-Sevilla [47]	0.07	0.02
Chen, Mahanthappa [48]	0.15	0.09
Raby [49]	0.1	0.04
<i>SO(10) + texture</i>		
Buchmüller, Wyler [50]	0.1	0.04
Bando, Obara [51]	0.01 .. 0.06	$4 \cdot 10^{-4}$.. 0.01
<i>Flavor symmetries</i>		
Grimus, Lavoura [52, 53]	0	0
Grimus, Lavoura [52]	0.3	0.3
Babu, Ma, Valle [54]	0.14	0.08
Kuchimanchi, Mohapatra [55]	0.08 .. 0.4	0.03 .. 0.5
Ohlsson, Seidl [56]	0.07 .. 0.14	0.02 .. 0.08
King, Ross [57]	0.2	0.15
<i>Textures</i>		
Honda, Kaneko, Tanimoto [58]	0.08 .. 0.20	0.03 .. 0.15
Lebed, Martin [59]	0.1	0.04
Bando, Kaneko, Obara, Tanimoto [60]	0.01 .. 0.05	$4 \cdot 10^{-4}$.. 0.01
Ibarra, Ross [61]	0.2	0.15
<i>3 × 2 see-saw</i>		
Appelquist, Piai, Shrock [62, 63]	0.05	0.01
Frampton, Glashow, Yanagida [64]	0.1	0.04
Mei, Xing [65] (normal hierarchy)	0.07	0.02
(inverted hierarchy)	> 0.006	> $1.6 \cdot 10^{-4}$
<i>Anarchy</i>		
de Gouvêa, Murayama [66]	> 0.1	> 0.04
<i>Renormalization group enhancement</i>		
Mohapatra, Parida, Rajasekaran [67]	0.08 .. 0.1	0.03 .. 0.04

2. Future LBL (Long BaseLine experiments)

Most realistic way to measure θ_{13} , $\text{sign}(\Delta m^2_{31})$ and δ is long base line experiments by **accelerators or reactors**.

→ **Matter effect** contributes in LBL in most cases

- Two points to be taken into account for precise measurements:

(1) Correlation of errors

(2) Parameter degeneracy

(1) Correlation of errors

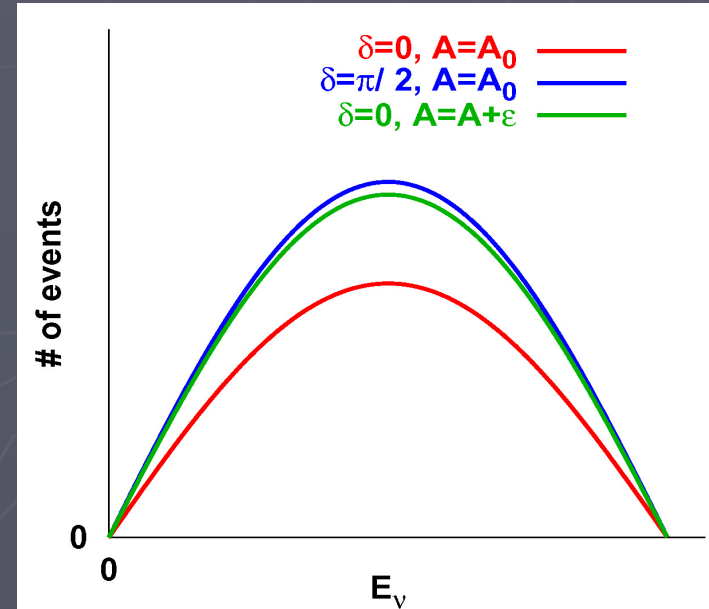
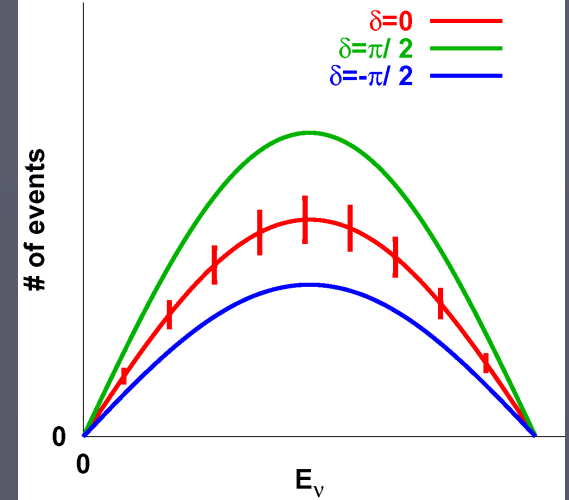
If uncertainties of other parameters (such as density of matter $\rho \propto A$) mimic the dependence on δ , then we cannot determine δ (correlation of errors)

→ We have take into account the uncertainties of other parameters to reject “ $\delta=0$ ”

$$\Delta\chi^2 = \sum_j \frac{[N_j(\delta) - N_j(\delta = 0)]^2}{\sigma_j^2}$$



$$\Delta\chi^2 = \min_{\theta_{k\ell}, \Delta m_{k\ell}^2, \bar{A}} \sum_j \frac{[N_j(\delta; \theta_{k\ell}, \Delta m_{k\ell}^2, \bar{A}) - N_j(\delta = 0; \theta_{k\ell}, \Delta m_{k\ell}^2, \bar{A})]^2}{\sigma_j^2}$$



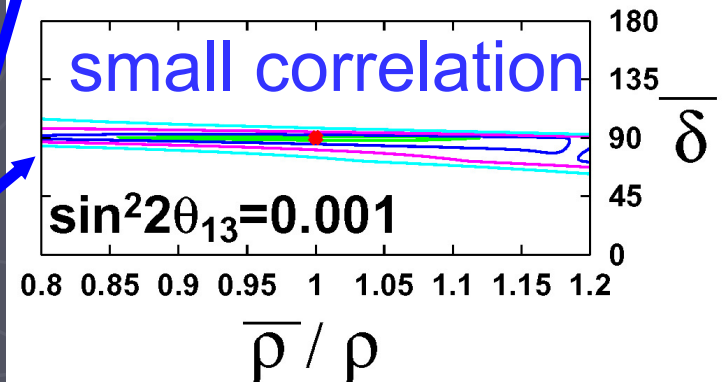
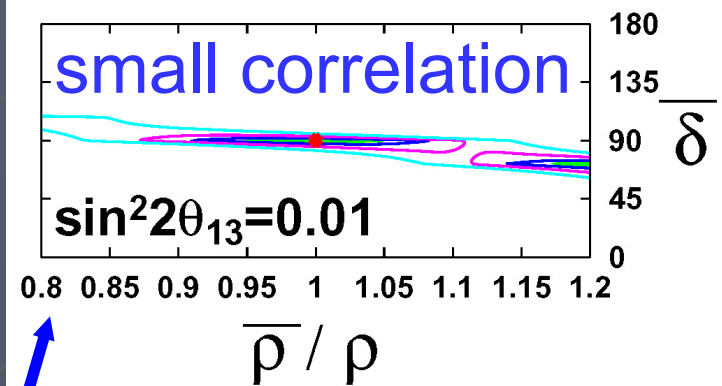
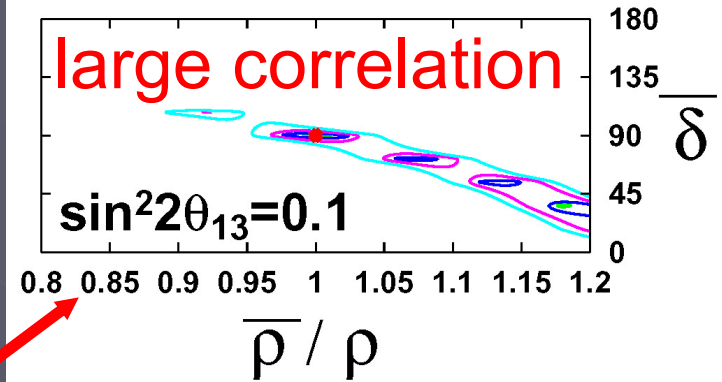
(Example) There is correlation of errors at a neutrino factory:

$$E_{\mu} = 50\text{GeV}, L = 3000\text{km}$$

correlation of errors in ρ and δ is serious for $\sin^2 2\theta_{13} \sim 0.1$

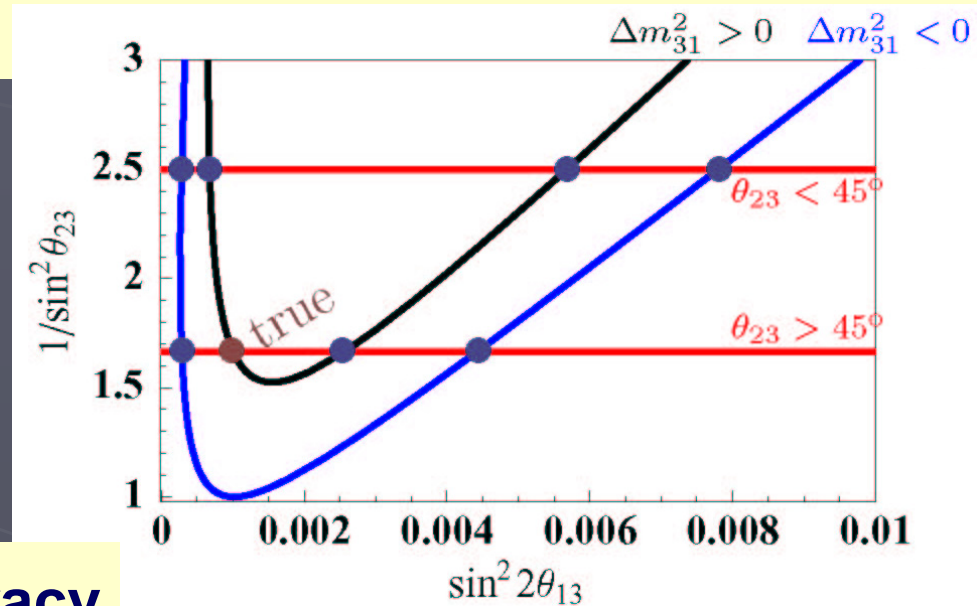
→ sensitivity to $\sin^2 2\theta_{13}$ of a ν factory is poor for $\sin^2 2\theta_{13} \sim 0.1$

correlation of errors in ρ and δ is not serious for $\sin^2 2\theta_{13} < 0.01$



(2) Parameter degeneracy

Even if we know $P(\nu_\mu \rightarrow \nu_e)$ and $P(\overline{\nu}_\mu \rightarrow \overline{\nu}_e)$ in a long baseline accelerator experiments with approximately monoenergetic neutrino beam, precise determination of θ_{13} , $\text{sign}(\Delta m_{31}^2)$ and δ is difficult because of the 8-fold parameter degeneracy.



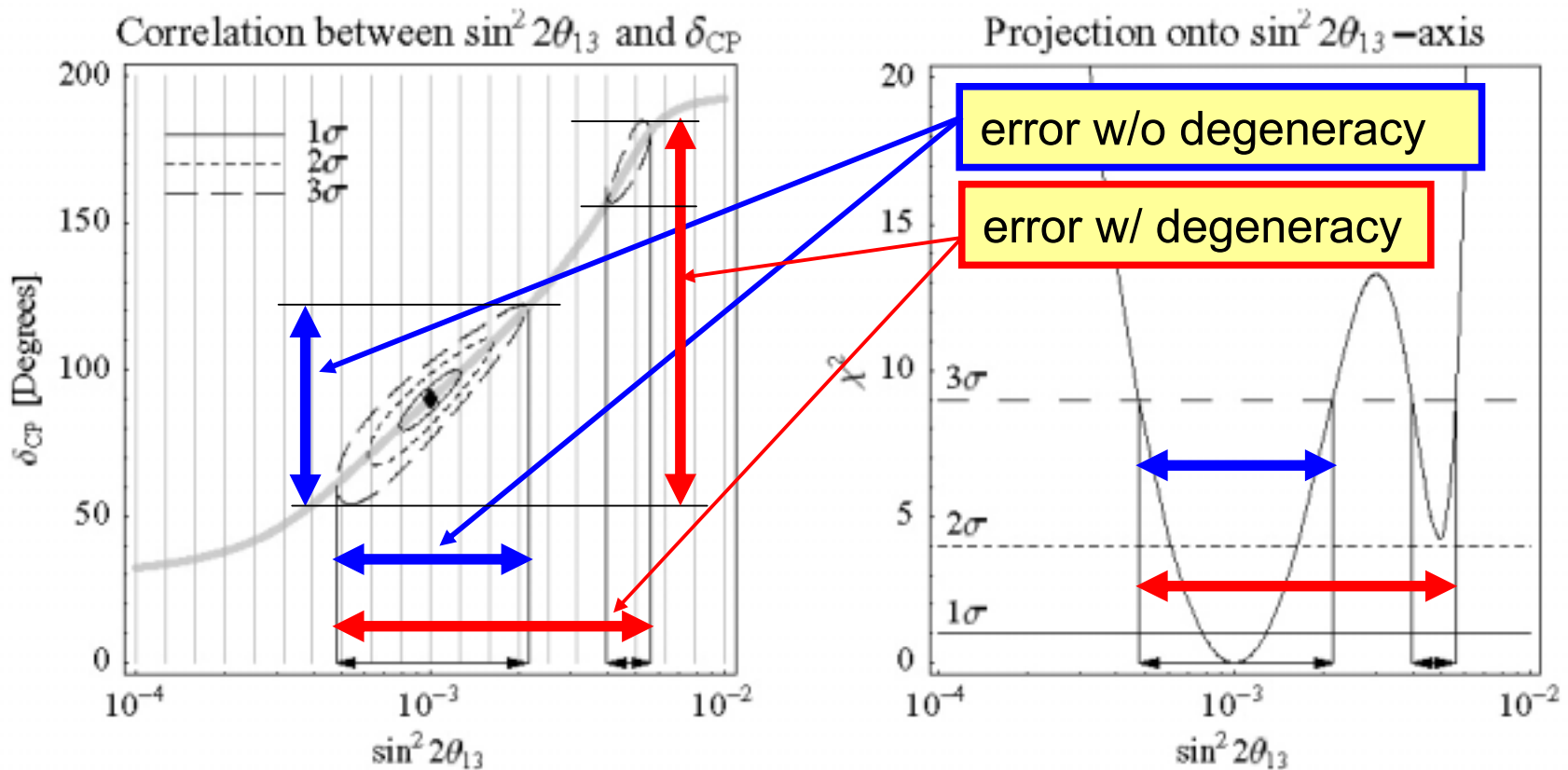
● intrinsic (δ, θ_{13}) degeneracy

● $\Delta m_{31}^2 \leftrightarrow -\Delta m_{31}^2$ degeneracy

● $\theta_{23} \leftrightarrow \pi/2 - \theta_{23}$ degeneracy

If parameter degeneracy exists, then the errors of the parameters become unnecessarily large.

Resolution of parameter degeneracy is important.



Future LBL experiments

To perform precise measurements of θ_{13} and δ , one has to have a lot of numbers of events to improve statistical errors.

→ We need **high intensity** beam

Candidates for high intensity beam in the future:

- (conventional) superbeam

{	$\pi^+ \rightarrow \mu^+ + \nu_\mu$	$\nu_\mu \rightarrow \nu_e$
	$\pi^- \rightarrow \mu^- + \bar{\nu}_\mu$	$\bar{\nu}_\mu \rightarrow \bar{\nu}_e$

- neutrino factory

{	$\mu^+ \rightarrow e^+ + \nu_e + \bar{\nu}_\mu$	$\nu_e \rightarrow \nu_\mu$
	$\mu^- \rightarrow e^- + \bar{\nu}_e + \nu_\mu$	$\bar{\nu}_e \rightarrow \bar{\nu}_\mu$

μ in a storage ring

- beta beam

{	${}^6_2\text{He} \rightarrow {}^6_3\text{Li} + e^- + \bar{\nu}_e$	$\bar{\nu}_e \rightarrow \bar{\nu}_\mu$
	${}^{18}_{10}\text{Ne} \rightarrow {}^{18}_9\text{F} + e^+ + \nu_e$	$\nu_e \rightarrow \nu_\mu$

RI in a storage ring

Future LBL exp. (**under construction** / proposed)

● **superbeam**

T2K phase I (2009-, 0.75MW, $E \sim 1\text{GeV}$, $L=295\text{km}$)

T2K phase II (4MW+HK, $E \sim 1\text{GeV}$, $L=295\text{km}$)

T2KK (JAERI \rightarrow HK&Korea, $E \sim 1\text{GeV}$, $L=295\text{km} \& 1000\text{km}$)

NOvA (FNAL \rightarrow Ash River (MN), $E \sim 2\text{GeV}$, $L=810\text{km}$)

VBLNO (BNL \rightarrow DUSEL*, $E \sim 2\text{GeV}$, $L > 2500\text{km}$)

(*Deep Underground Science and Engineering Laboratory:
Homestake(SD), Icicle Creek(WA), San Jacinto(CA), Soudan(MN),
Kimballton(VA), Henderson(CO))

SPL (CERN \rightarrow Frejus, $E \sim 0.25\text{GeV}$, $L=130\text{km}$)

● **neutrino factory** ($E_\nu < 50\text{GeV}$, $L \sim 3000\text{km}$)

● **beta beam** ($E_\nu = 0.5-1.5\text{GeV}$, $L \sim 130\text{km}$)

Proposed reactor experiments ($E \sim 4\text{MeV}$, $L \sim 2\text{km}$)

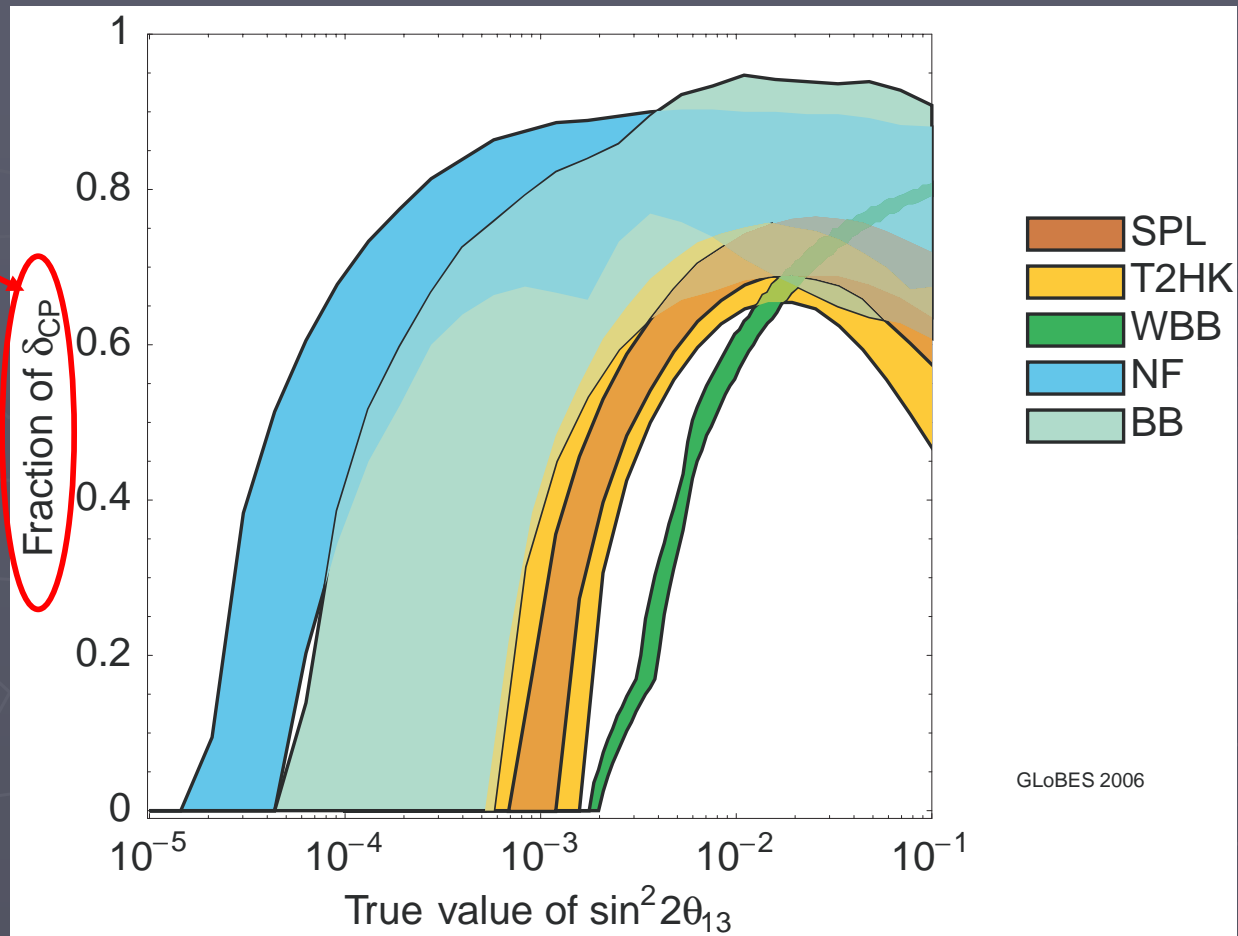
Double CHOOZ (France) , **Daya Bay** (China),
Reno (Korea), **Angra** (Brazil)

sensitivity to the CP phase δ of future experiments

error of δ

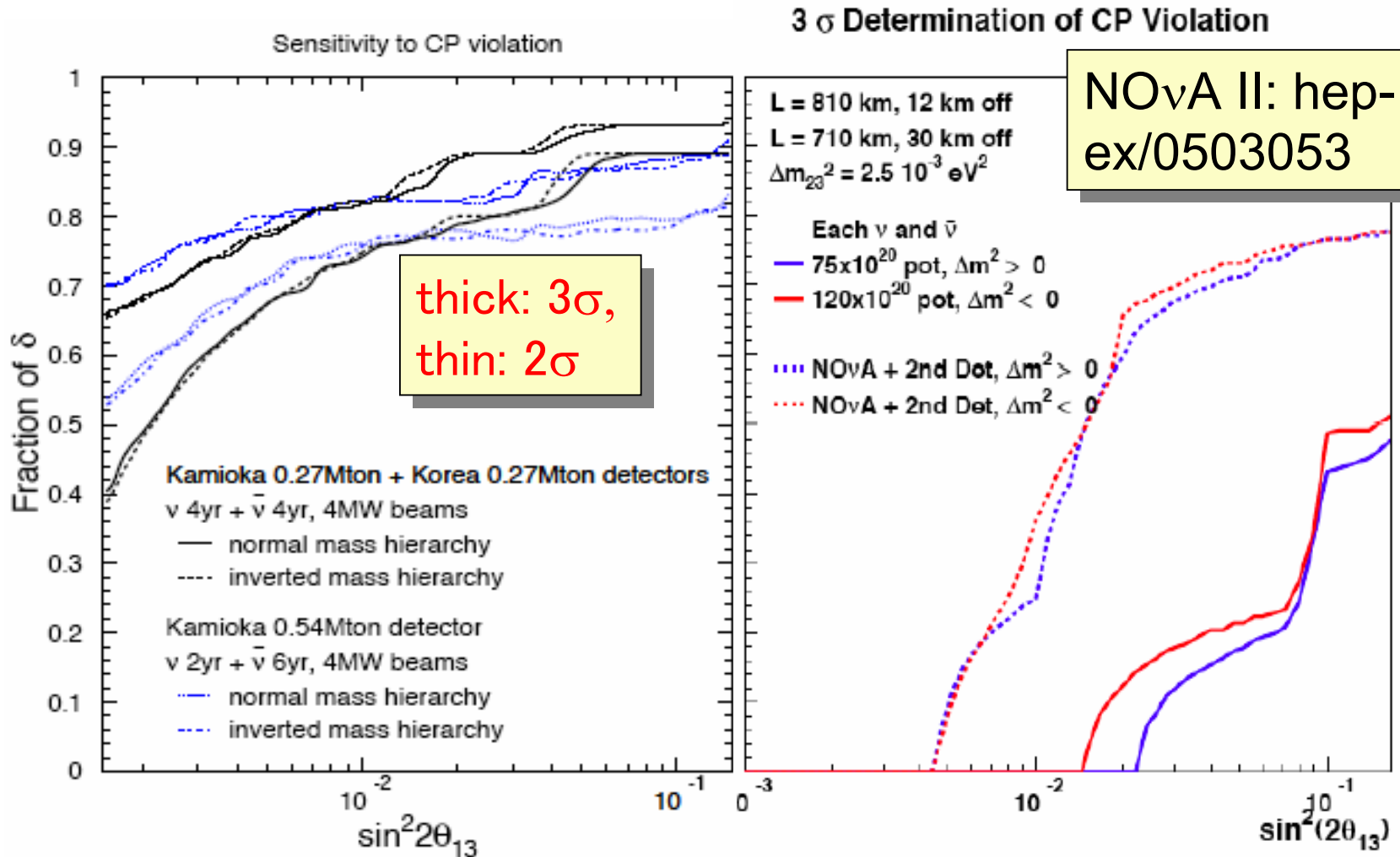
T2HK = Tokai to
HyperKamiokande

Fraction of δ_{CP}



arXiv:0710.4947 [hep-ph]

T2KK vs. NO ν A; CP





International scoping study of a future Neutrino Factory and super-beam facility

June 2005 ~ Aug. 2006

<http://www.hep.ph.ic.ac.uk/iss/>

- Evaluate the physics case for the facility
- Study options for the accelerator complex and neutrino detection systems

- ◆ Physics Group **Y. Nagashima**
- ◆ Detector Group **A. Blondel**
- ◆ Accelerator Group **M. Zisman**

- Theory Subgroup **S.F. King**
- Phenomenology Subgroup **OY**
- Experiment Subgroup **K. Long**
- Muon Subgroup **L. Roberts**

➤ **Theory Subgroup**

Model building for neutrino mass & mixing

➤ **Phenomenology Subgroup**

Deviation from SM with massive neutrinos

➤ **Experiment Subgroup**

Estimation of sensitivity and resolution of degeneracy for ν factories and β beams

**Final report of Physics Group :
arXiv:0710.4947 [hep-ph]**



International Design Study of the **Neutrino Factory**

2007 ~

<http://www.hep.ph.ic.ac.uk/ids/>

- ◆ Accelerator Group
- ◆ Detector Group
- ◆ **Physics and Performance Evaluation Group**

Tentative plan of PPEG

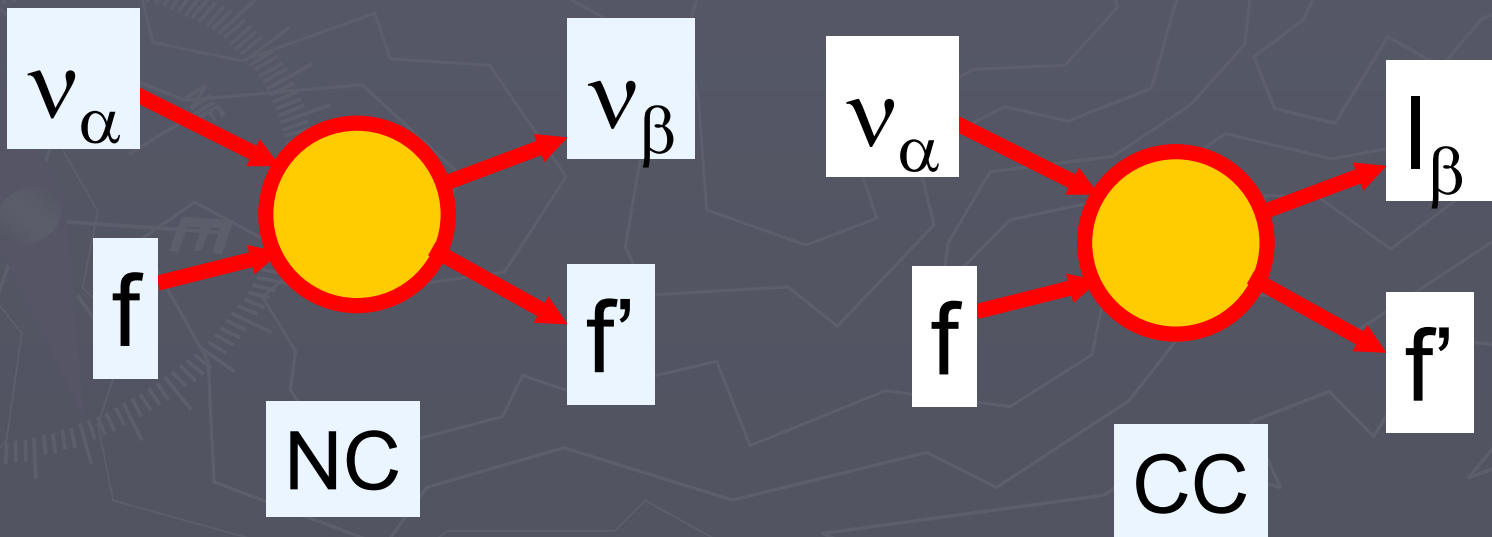
- provide sensitivity estimates as requested
- track new developments like low E NF
- understand optimization in the context of non-standard physics
- establish a physics case for all values of θ_{13}
- keep track of competitors to the NF
- near detector, both at a SB and NF
- status of cross-sections
- muon physics

3. Deviation from standard 3 flavor framework

- (1) New physics (NP) (exotic interactions)**
- (2) violation of unitarity (like at a B factory)**
- (3) Sterile neutrinos**

(1) New physics (NP) (exotic interactions)

Flavors are not necessarily conserved:



Effects of New Physics on ν oscillations

$$\mathcal{A}_{SM} \equiv A \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$A \equiv \sqrt{2}G_F N_e$$

$$P(\nu_\alpha \rightarrow \nu_\beta) = \left| \left[U^d \tilde{U} \exp \left\{ -i \text{diag}(\tilde{E}_j) L \right\} \tilde{U}^{-1} (U^s)^{-1} \right]_{\beta\alpha} \right|^2$$

with

$$U \text{diag}(E_j) U^{-1} + \mathcal{A}_{SM} + \mathcal{A}_{NP} \equiv \tilde{U} \text{diag}(\tilde{E}_j) \tilde{U}^{-1}$$

source

$$|(U^s-1)_{\alpha\beta}| < \mathcal{O}(10^{-2})$$

$$\begin{pmatrix} \nu_e^s \\ \nu_\mu^s \\ \nu_\tau^s \end{pmatrix} = U^s U_{MNS} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

propagation

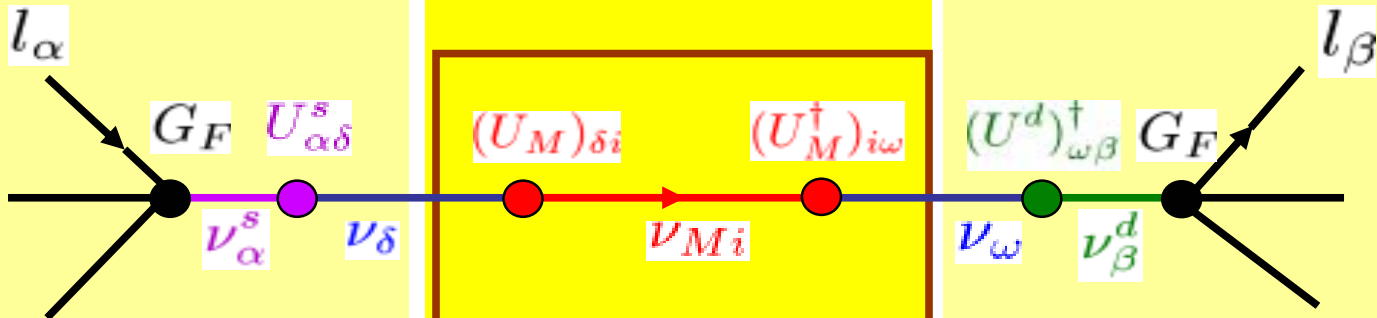
$$\epsilon_{ee}, \epsilon_{e\tau}, \epsilon_{\tau\tau} \sim \mathcal{O}(1)$$

$$A \begin{pmatrix} \epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu}^* & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau}^* & \epsilon_{\mu\tau}^* & \epsilon_{\tau\tau} \end{pmatrix}$$

detection

$$|(U^d-1)_{\alpha\beta}| < \mathcal{O}(10^{-2})$$

$$\begin{pmatrix} \nu_e^d \\ \nu_\mu^d \\ \nu_\tau^d \end{pmatrix} = U^d U_{MNS} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$



(i) Effects of New Physics at source and detector

Deviation from the standard form is small:

Grossman (PLB359:141,1995)

$$|(\mathbf{U}^s - 1)_{\alpha\beta}| < \mathcal{O}(10^{-2}), \quad |(\mathbf{U}^d - 1)_{\alpha\beta}| < \mathcal{O}(10^{-2})$$

(ii) New Physics effects in propagation

1. Constraints from various ν experiments:

Davidson et al (JHEP 0303:011,2003)

$$\left(\begin{array}{ccc} -3 \lesssim \epsilon_{ee} \lesssim 2 & |\epsilon_{e\mu}| \lesssim 0.5 & |\epsilon_{e\tau}| \lesssim 1.5 \\ |\epsilon_{e\mu}| \lesssim 0.5 & |\epsilon_{\mu\mu}| \lesssim 0.05 & |\epsilon_{\mu\tau}| \lesssim 0.15 \\ |\epsilon_{e\tau}| \lesssim 1.5 & |\epsilon_{\mu\tau}| \lesssim 0.15 & |\epsilon_{\tau\tau}| \lesssim 6 \end{array} \right)$$

2. Constraints from atmospheric neutrinos:

Friedland-Lunardini (Phys.Rev.D72:053009,2005)

$\epsilon_{ee}, \epsilon_{e\tau}, \epsilon_{\tau\tau} \sim \mathbf{O(1)}$ are consistent with ν_{atm} data, provided

$$\epsilon_{\tau\tau} \sim \frac{|\epsilon_{e\tau}|^2}{1 + \epsilon_{ee}}$$

$$0 \leq |\epsilon_{e\tau}| \lesssim 1 + \epsilon_{ee}$$

$$-1 \lesssim \epsilon_{ee} \lesssim 1.5$$

In general:

NP effects at production and at detection becomes important when L is smaller

$$|\epsilon_{\alpha\beta}| < O(10^{-2})$$

$$P(\nu_{\alpha} \rightarrow \nu_{\beta}) \rightarrow \left| \left[U^d (U^s)^{-1} \right]_{\beta\alpha} \right|^2$$

i.e., no BG from osc.
in the limit of $L \rightarrow 0$



Experiments with a shorter baseline are advantageous

NP effects in propagation becomes important when L is larger

$$\epsilon_{\alpha\beta} \text{ can be of } O(1)$$

because $AL \epsilon_{\alpha\beta} \sim \epsilon_{\alpha\beta} (L/2000\text{km})$



Experiments with a longer baseline (e.g., **INO**) are advantageous

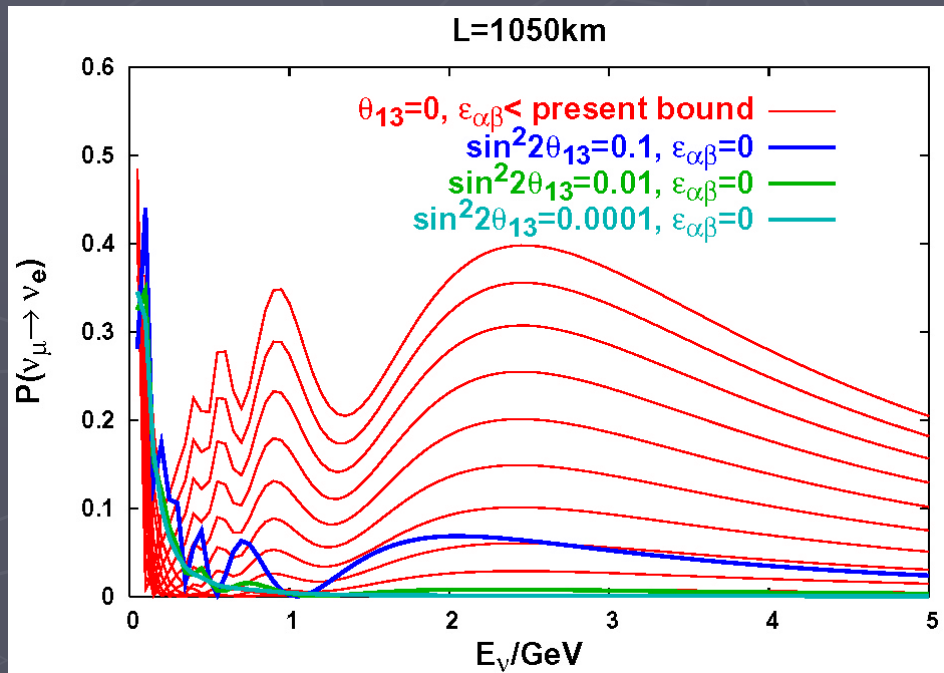
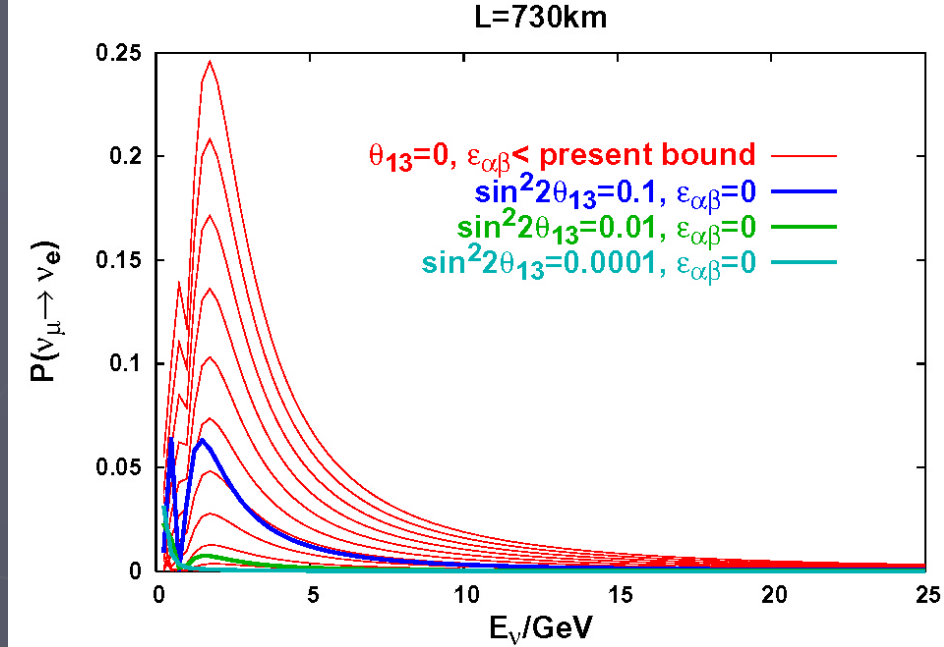
MINOS (ν_e appearance)

$$\nu_{\mu} \rightarrow \nu_e$$

T2KK (ν_e appearance)

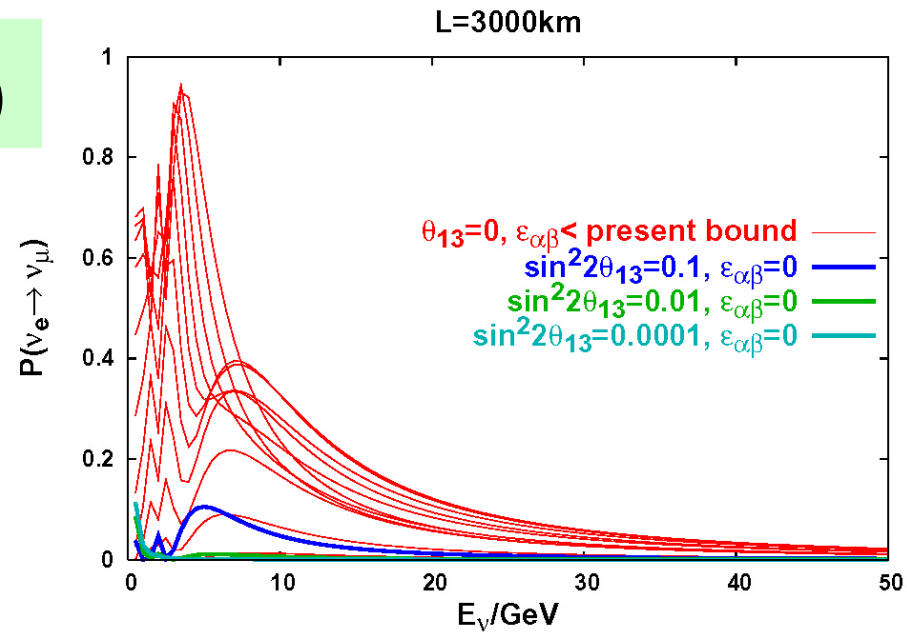
$$\nu_{\mu} \rightarrow \nu_e$$

NK,HS,OY, hep-ph/0606013



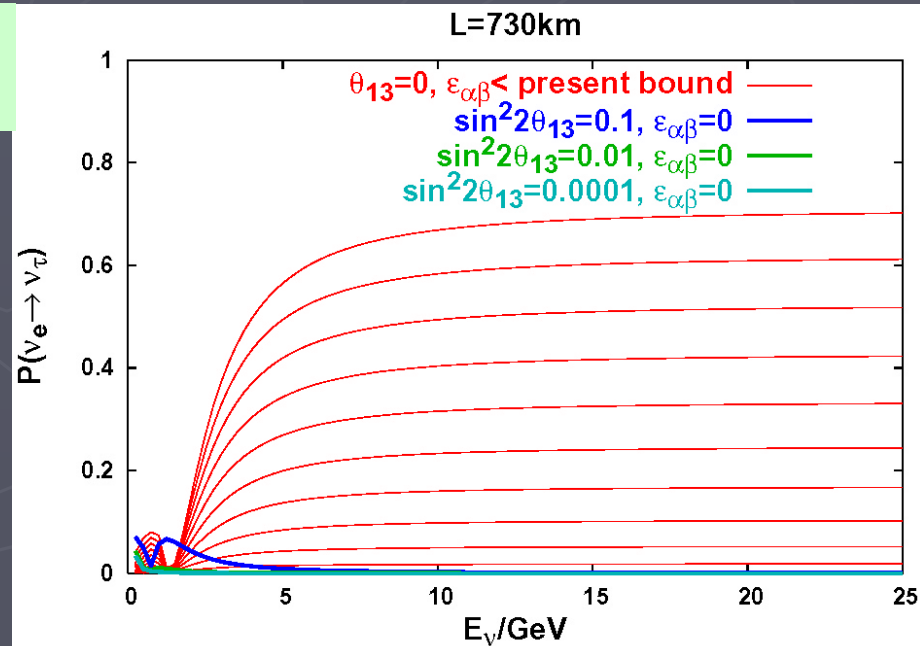
ν factory (golden channel)

$$\nu_e \rightarrow \nu_\mu$$



ν factory (silver channel)

$$\nu_e \rightarrow \nu_\tau$$



NK,HS,OY, hep-ph/0606013

(2) Unitarity violation due to heavy ν

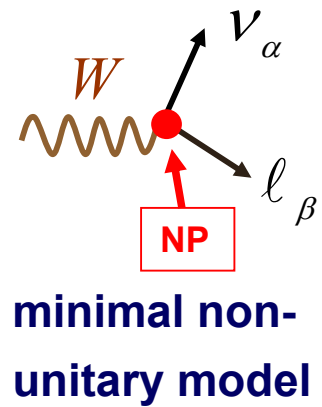
$U \rightarrow N$ (non-unitary)

$NN^\dagger - 1$: deviation from unitarity

mostly from rare decays

$$|NN^\dagger| \approx \begin{pmatrix} 0.994 \pm 0.005 & < 7.1 \cdot 10^{-5} & < 1.6 \cdot 10^{-2} \\ < 7.1 \cdot 10^{-5} & 0.995 \pm 0.005 & < 1.0 \cdot 10^{-2} \\ < 1.6 \cdot 10^{-2} & < 1.0 \cdot 10^{-2} & 0.995 \pm 0.005 \end{pmatrix}$$

90% C.L.



Constraints on unitary violation are strong:

$$|(NN^\dagger - 1)_{\alpha\beta}| < \mathcal{O}(10^{-2})$$

Minimal case: Antusch, Biggio, Fernandez-Martinez, Gavela, Lopez-Pavon, JHEP 0610:084,2006.

Non-minimal case: Abada, Biggio, Bonnet, Gavela, Hambye, arXiv:0707.4058 [hep-ph]

Even stronger constraints \rightarrow more difficult to detect

Unitarity violation is similar to NP effects
at production and at detection →
becomes important when L is smaller

$$P(\mathbf{v}_\alpha \rightarrow \mathbf{v}_\beta) \rightarrow \left| \left[(1+\eta) U \exp[-i \text{diag}(E_j)] U^{-1} (1+\eta)^{-1} \right]_{\beta\alpha} \right|^2$$

$$\eta \equiv (NN^\dagger - 1) / 2$$

$$|\eta_{\alpha\beta}| < O(10^{-2})$$

$$U^s \rightarrow 1 + \eta/2, U^d \rightarrow 1 + \eta/2$$



Experiments with a shorter baseline
are advantageous

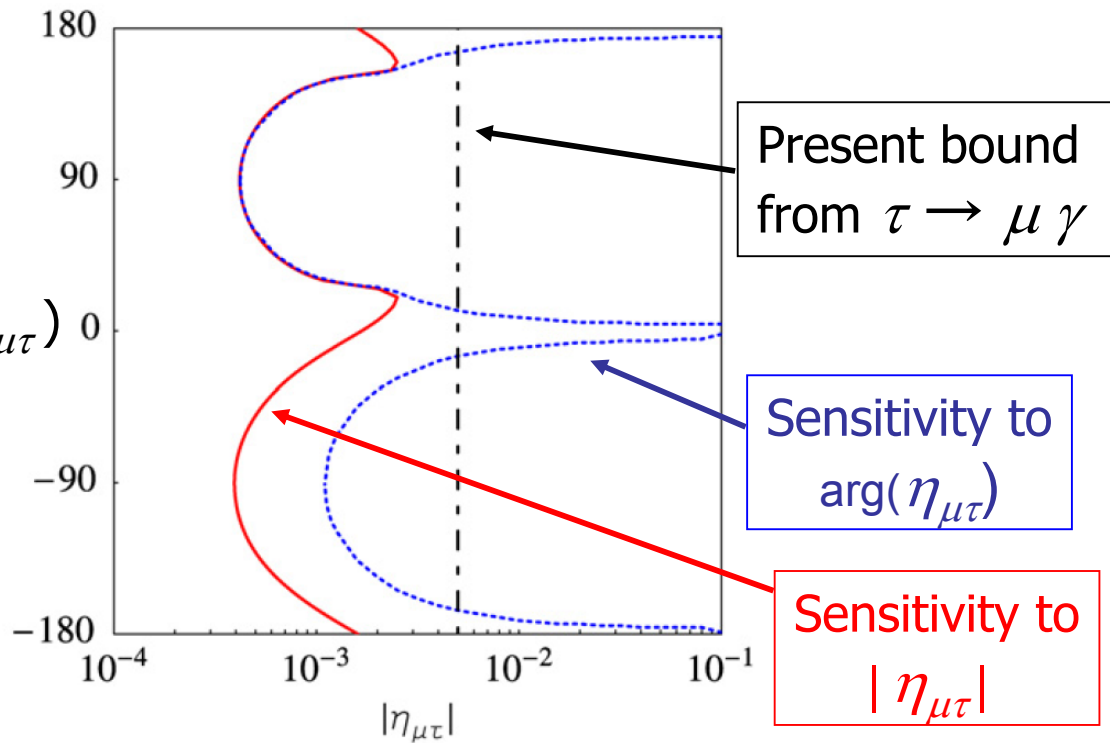
Neutrino factory with a baseline $L=130\text{km}$

Fernandez-Martinez, Gavela, Lopez-Pavon, OY,
Phys.Lett.B649:427-435,2007

Phase of N

$$\nu_{\mu} \leftrightarrow \nu_{\tau}$$

For non-trivial $\arg(\eta_{\mu\tau})$,
one order of magnitude
improvement for $|\eta_{\mu\tau}|$



(3A) Sterile neutrinos assuming LSND

$$\Delta m_{\text{sol}}^2 \sim 10^{-4} \text{ eV}^2$$

$$\Delta m_{\text{atm}}^2 \sim 10^{-3} \text{ eV}^2$$

$$\Delta m_{\text{LSND}}^2 \sim O(1) \text{ eV}^2$$

→ at least one ν_s is required

LSND($\bar{\nu}_\mu \rightarrow \bar{\nu}_e$): affirmative

MiniBOONE($\nu_\mu \rightarrow \nu_e$): negative

difference between ν & anti- ν may offer a promising fit

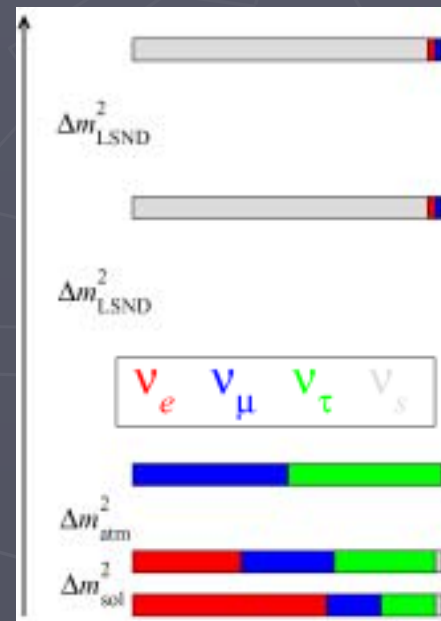
(3+2)-scheme w/ CP phase δ

$$\begin{aligned}
 P_{\nu_\mu \rightarrow \nu_e} &= 4 |U_{e4}|^2 |U_{\mu4}|^2 \sin^2 \phi_{41} \\
 &+ 4 |U_{e5}|^2 |U_{\mu5}|^2 \sin^2 \phi_{51} \\
 &+ 8 |U_{e4} U_{\mu4} U_{e5} U_{\mu5}| \sin \phi_{41} \sin \phi_{51} \cos(\phi_{54} - \delta)
 \end{aligned}$$

with the definitions

$$\phi_{ij} \equiv \frac{\Delta m_{ij}^2 L}{4E},$$

$$\delta \equiv \arg(U_{e4}^* U_{\mu4} U_{e5} U_{\mu5}^*).$$



Schwetz-Mangold@nufact07

$$\phi_{54}^{best} = 1.64 \pi$$

$$\Delta m_{41}^2 = 0.89 \text{ eV}^2$$

$$\Delta m_{51}^2 = 6.49 \text{ eV}^2$$

$$\chi_{\min}^2 = 94.5 / (107 - 7)$$

Karagiorgi@nufact07

$$\chi^2/ndf = 146.7/156$$

gof=69%

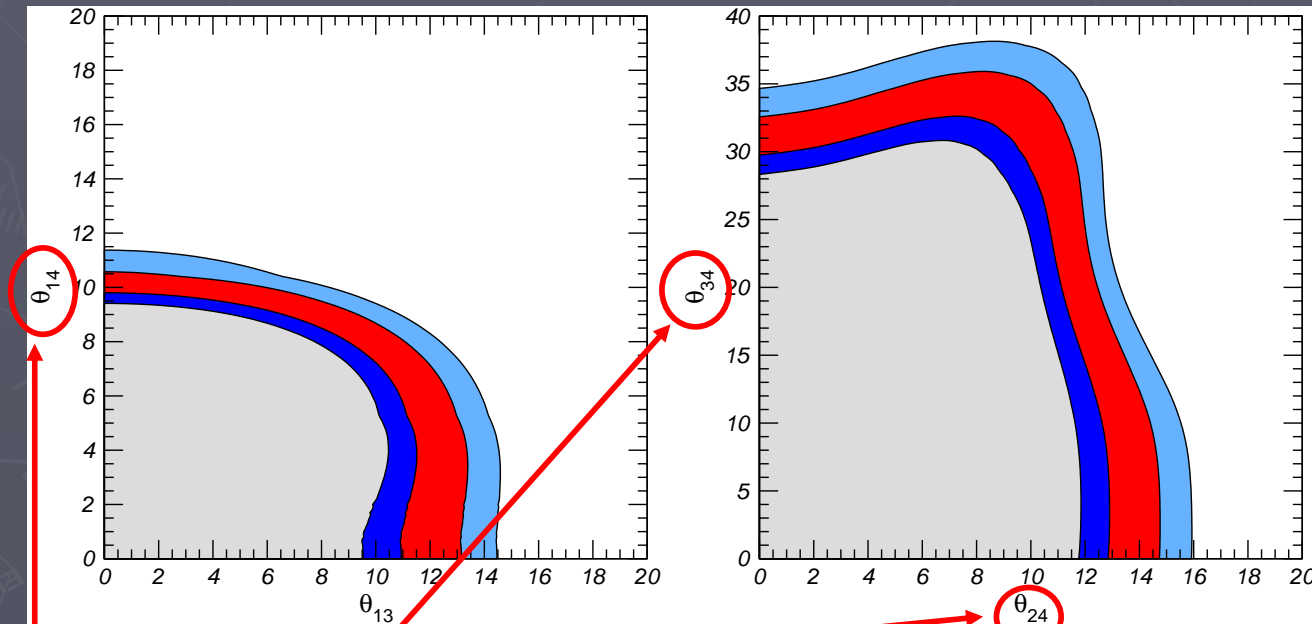
$$\phi_{54}^{best} = 1.74 \pi$$

- (3+2) schemes
 - offer the possibility of CP violation to reconcile LSND and MiniBooNE,
 - but there is tension between appearance and disappearance data (3σ , 4σ for MB300)

(3B) Sterile neutrinos w/o assuming LSND

Without assuming LSND and imposing all the negative constraints one can still consider (3+1)-scheme

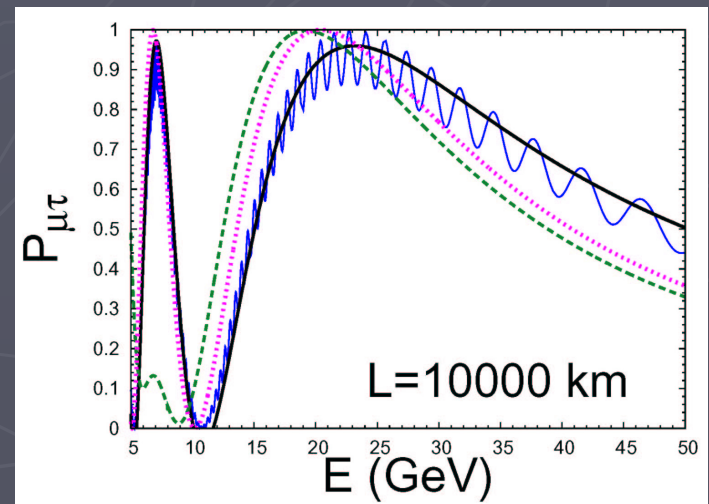
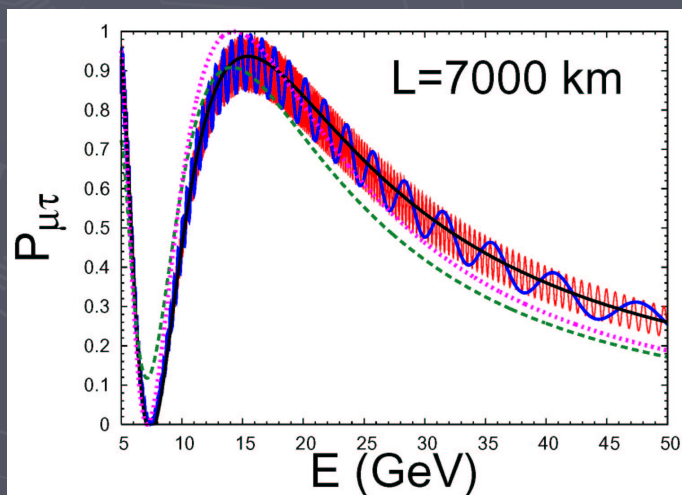
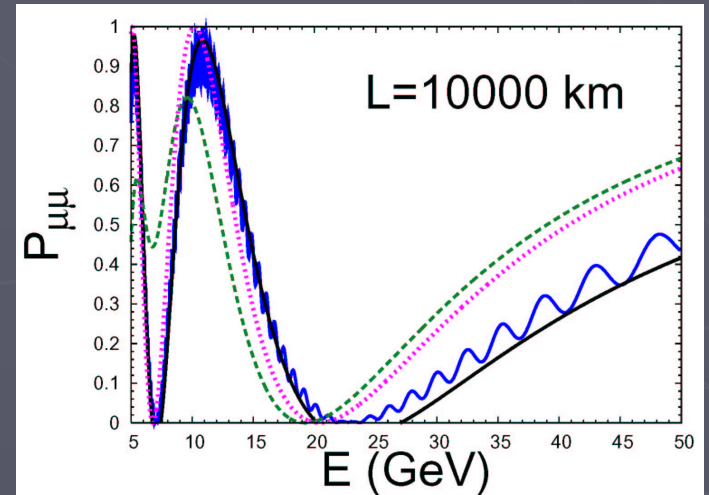
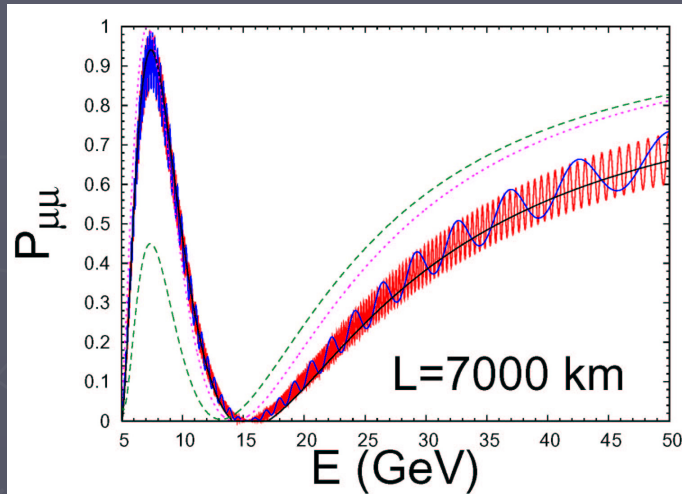
“3+1 sterile neutrinos at the CNGS” Donini, Maltoni, Meloni, Migliozzi, Terranova, arXiv:0704.0388v2 [hep-ph]



θ_{14} , θ_{24} , θ_{34} : angles which appear only in 4 ν scenario

“Signatures of heavy sterile neutrinos at long baseline experiments”

Dighe, Ray, arXiv:0709.0383 [hep-ph]



In either case (3A) or (3B), sterile neutrino oscillations will exhibit enhancement/suppression for ν or anti- ν , and experiments with a longer baseline (such as **INO**) are expected to have good sensitivity to sterile neutrinos.

3. Summary

- A brief review was given on the known parameters in SM+massive ν . Efforts to determine the unknown parameters (θ_{13} , δ , $\text{sign}(\Delta m^2_{31})$) in the future experiments were described.
- The future neutrino experiments with high precision will be able to see deviation from SM such as non-std. interactions, unitarity violation, sterile neutrinos, etc.
- Experiments with longer baselines (such as **INO**) are advantageous to search for NP in propagation and sterile neutrinos.