

# Signatures of sterile $\nu$ mixing in high energy cosmic $\nu$ flux

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Work in collaboration with  
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01-07-2008 @ nufact08

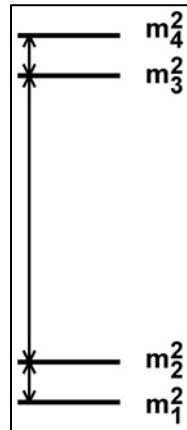
# 1. Four flavor $\nu$ oscillation

## 1.1 Schemes with LSND

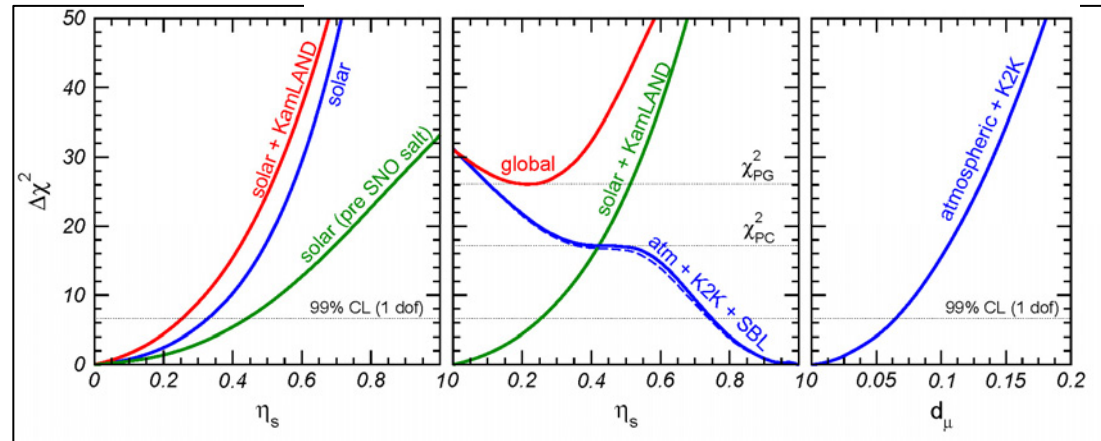
### (1) (2+2)-scheme

excluded ( $\sim 4.9 \sigma$ )  
because it  
contradicts with

$\nu_{\text{atm}}$  or  $\nu_{\text{solar}}$

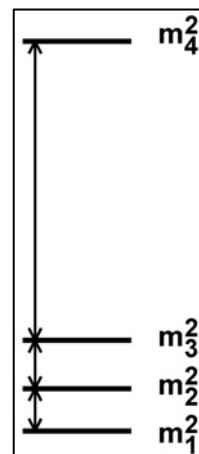


Maltoni et al, hep-ph/0405172

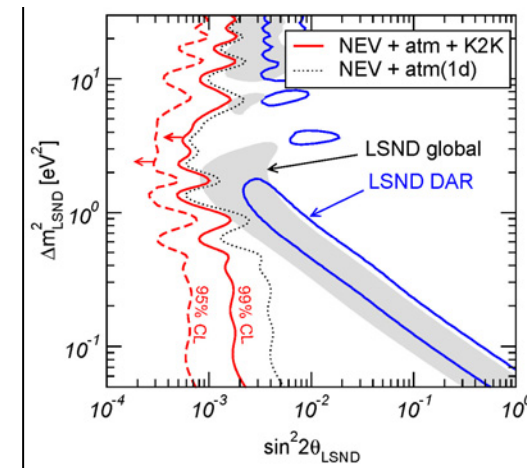


### (2) (3+1)-scheme

Strongly disfavored  
( $\sim 3.2 \sigma$  L) because of the  
tension between **LSND**  
and **Bugey+CDHSW**  
(+other negative results)



Maltoni et al, hep-ph/0405172



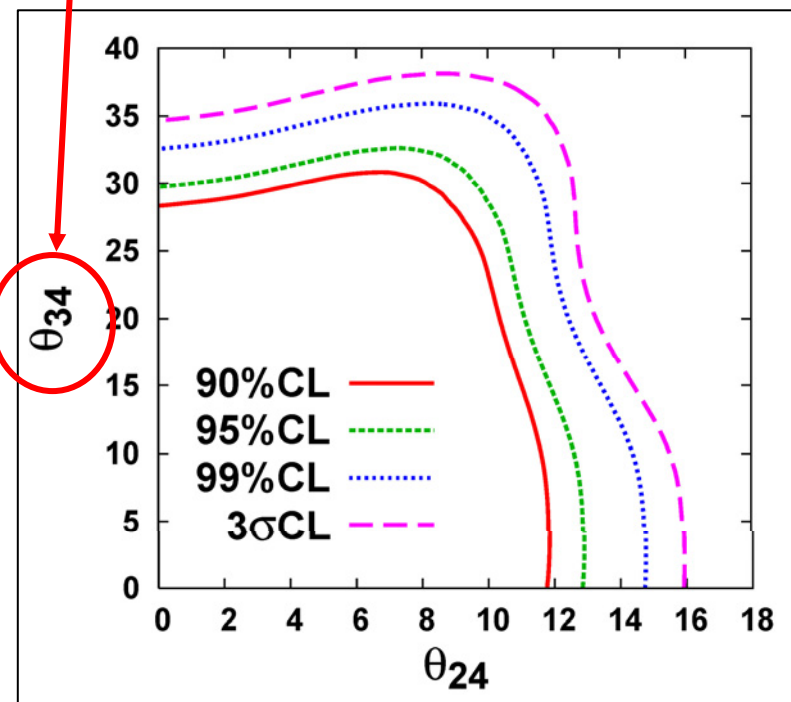
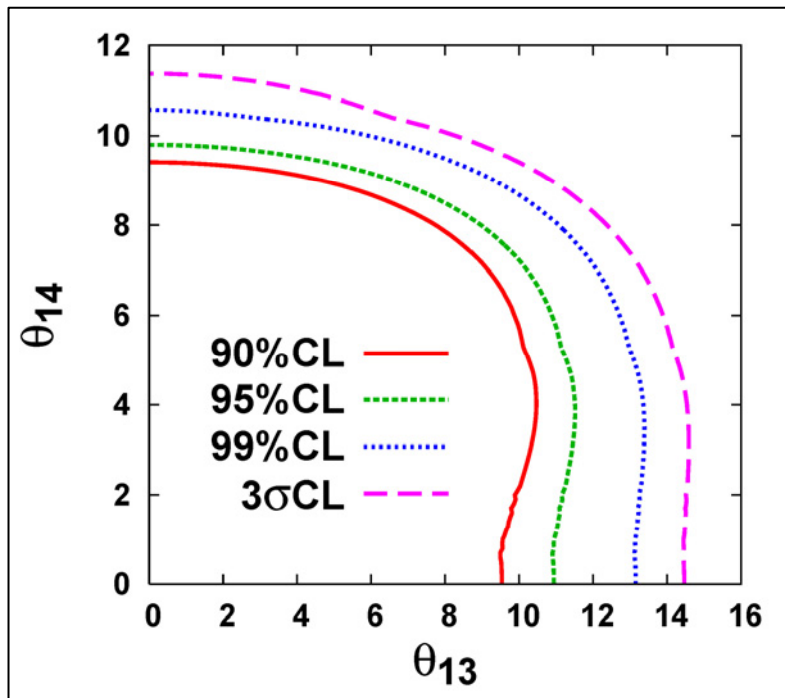
# 1.2 (3+1)-scheme without LSND

Donini-Maltoni-Meloni-Migliozzi-Terranova  
arXiv:0704.0388v2

$$U = R_{34}(\theta_{34}) R_{24}(\theta_{24}) R_{23}(\theta_{23}, \delta_3) R_{14}(\theta_{14}) R_{13}(\theta_{13}, \delta_2) R_{12}(\theta_{12}, \delta_1)$$

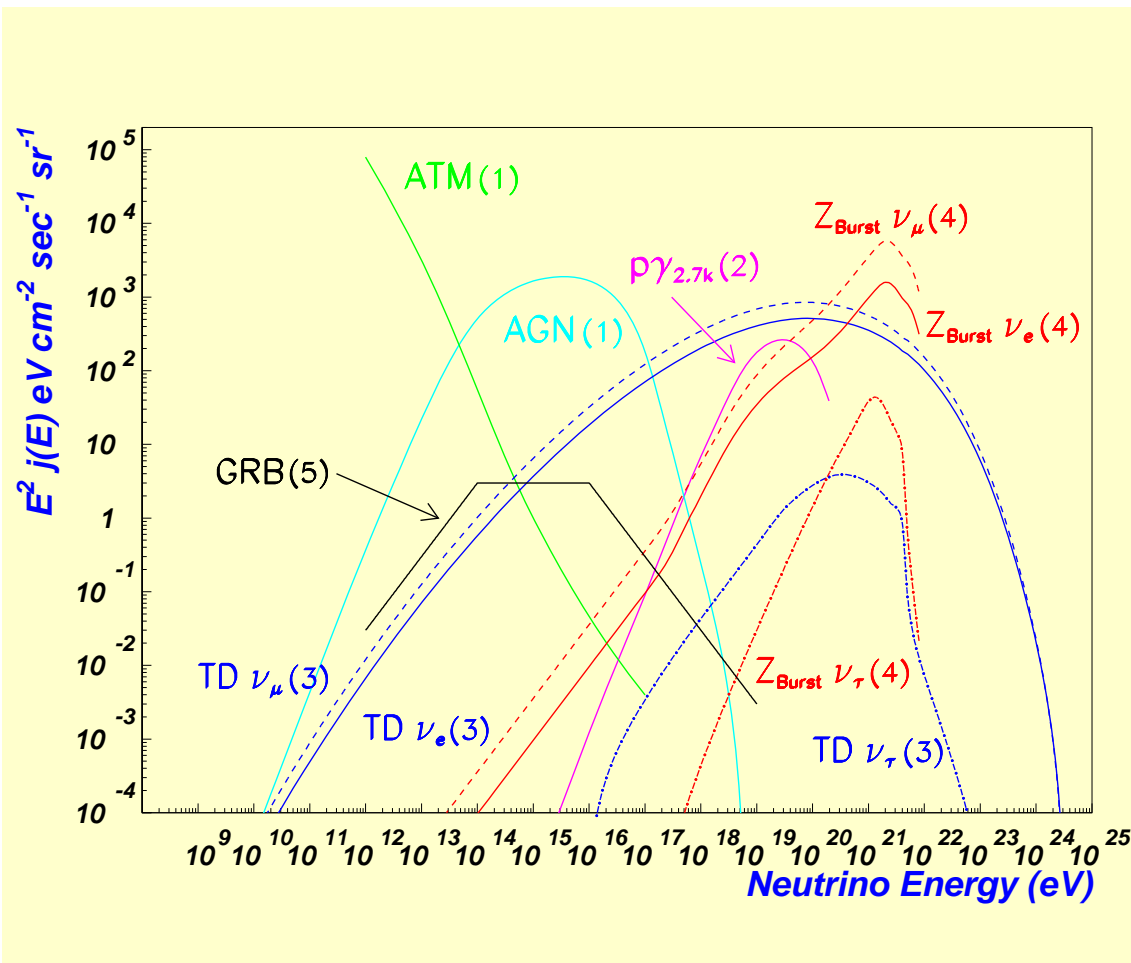
Constraints by all the negative results give the allowed region

$\theta_{34}$  : could be relatively large



## 2. Effects of $\nu$ oscillation on high energy cosmic $\nu$

### 2.1 Flux of high energy cosmic $\nu$



● Active Galactic Nuclei (AGN) & Gamma Ray Burst (GRB) are speculated to produce high energy  $\nu$

●  $E > 1 \text{ TeV}$ : small BG from  $\nu_{\text{atm}}$

## 2.2 Flavor ratio of $\nu$ flux for $N_\nu=3$

In standard  $N_\nu=3$ , when  $L \rightarrow \infty$

Learned, Pakvasa '95

### ● Oscillation for $L \rightarrow \infty$ :

$$P(\nu_\alpha \rightarrow \nu_\beta) \cong \sum_j |U_{\alpha j}|^2 |U_{\beta j}|^2$$

**CHOOZ+  $\nu_{\text{atm}}$  :  $|\theta_{13}| \ll 1$**

**$\nu_{\text{atm}}$  :  $|\pi/4 - \theta_{23}| \ll 1$**

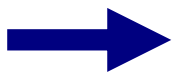
### ● Initial flux:

**Just like in  $\nu_{\text{atm}}$ , the source of  $\nu$  is  $\pi$  decay**

$$F^0(\nu_e) : F^0(\nu_\mu) : F^0(\nu_\tau) \\ \cong 1 : 2 : 0$$

### ● Observed flux:

$$F(\nu_e) : F(\nu_\mu) : F(\nu_\tau) \\ \cong 1 : 1 : 1$$



## 2.3 Triangle representation of flux

Athar-Jezabek-OY '00

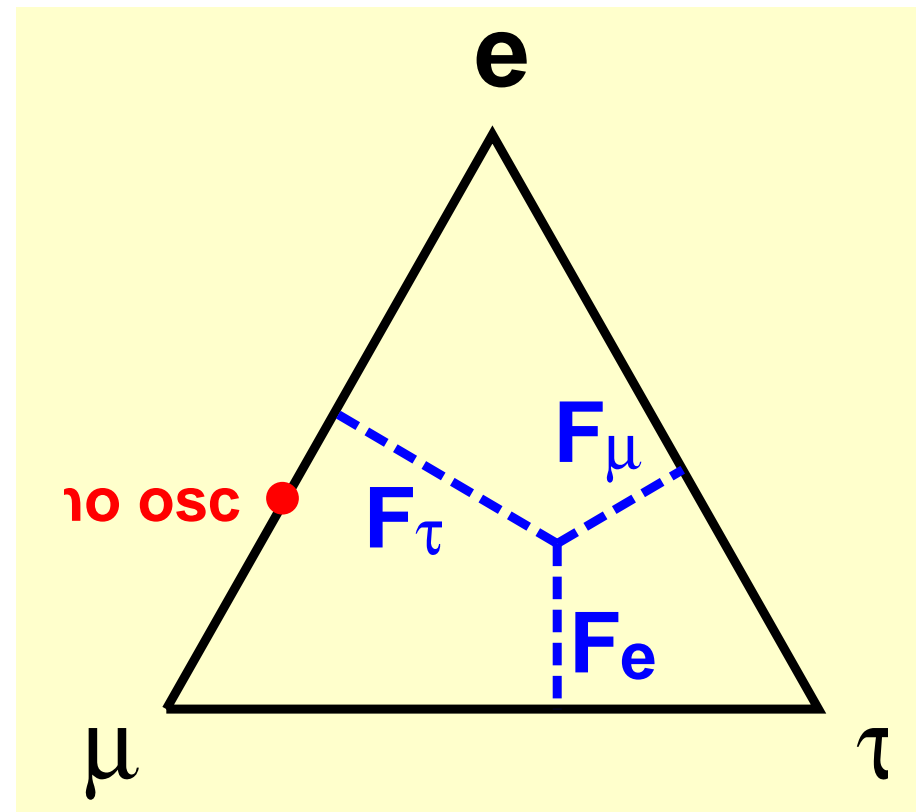
Precise normalization is not known

→ The ratio of different flavors is important quantity to observe

● The case for sterile  $\nu$

The normalized ratio of active flavors is useful:

$$\tilde{F}(\nu_\alpha) \equiv \frac{F(\nu_\alpha)}{F(\nu_e) + F(\nu_\mu) + F(\nu_\tau)}$$



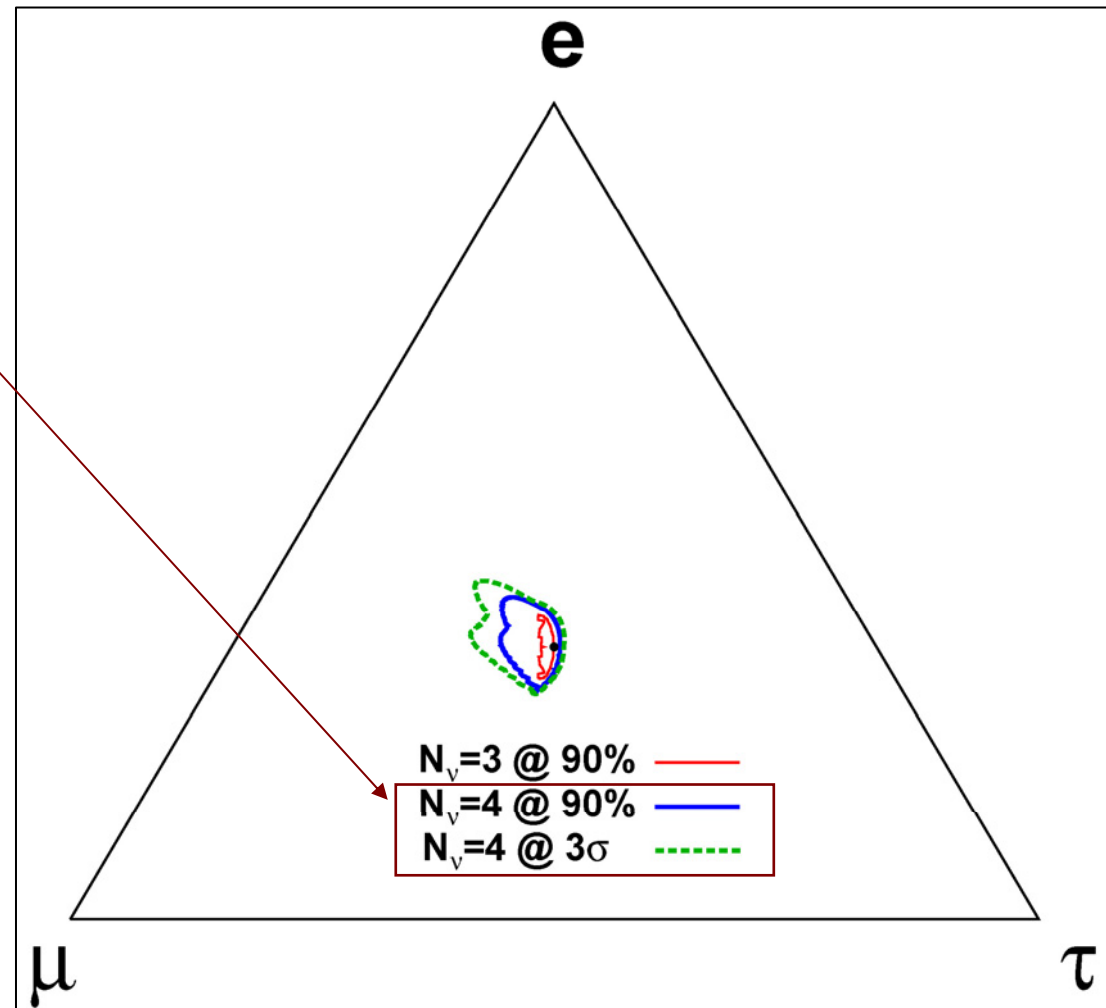
## 2.4 Flavor ratio of $\nu$ flux for (3+1)- scheme

Donini-OY '07

(3+1)-scheme w/o LSND gives the prediction which could be distinguished from  $N_\nu=3$  case



In principle, (3+1)-  
scheme could be  
distinguished  
from the three  
flavor case



# 2.5 Theoretical uncertainties of original $\nu$ flux

Lipari-Lusignoli-Meloni Phys.Rev.D75:123005,2007

For illustrations, they discussed  $\nu$  flux from GRB using Waxman-Bahcall

- Proton energy spectrum

$$N_p(E_p) \propto E_p^{-\alpha}$$

- Photon number density

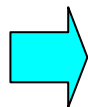
$$n_\gamma(\epsilon) \propto \begin{cases} (\epsilon/\epsilon_b)^{-\beta_1} & \text{for } \epsilon \leq \epsilon_b, \\ (\epsilon/\epsilon_b)^{-\beta_2} & \text{for } \epsilon_b < \epsilon < \epsilon_{\max} \\ 0 & \text{for } \epsilon \geq \epsilon_{\max}, \end{cases}$$

- Muon energy loss due to synchrotron radiation

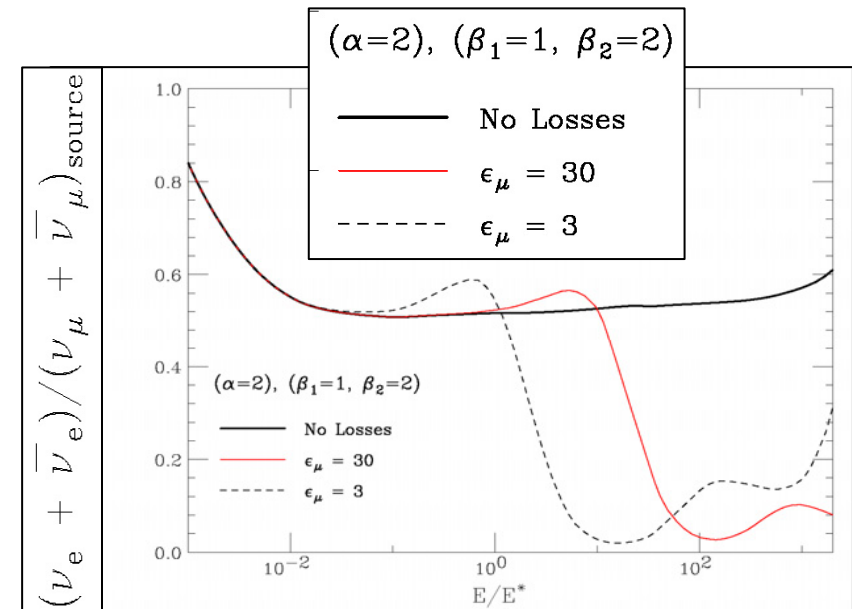
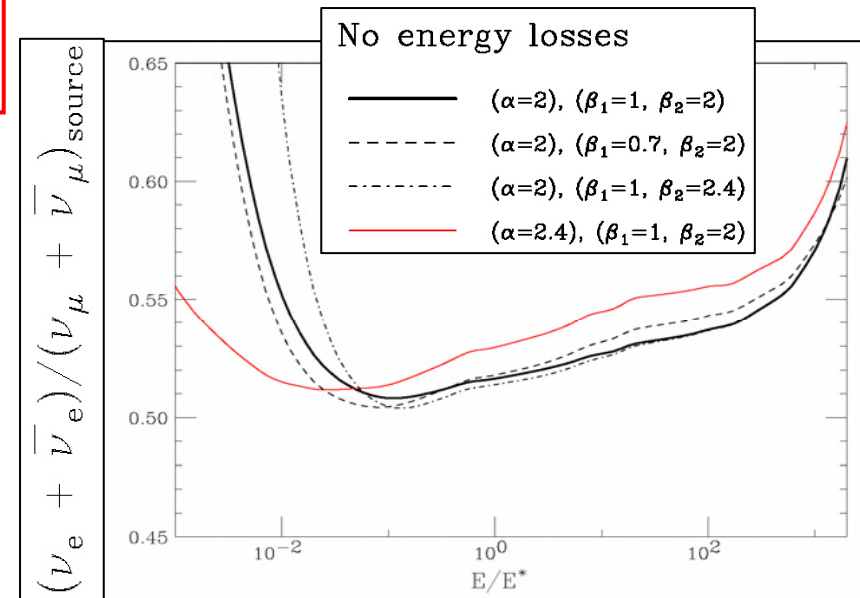
$$\epsilon_\mu = \frac{E_{\text{syn}}^\mu}{E^*} = 8.4 \times 10^4 \left( \frac{\text{Gauss}}{B} \right) \left( \frac{\epsilon_b}{\text{KeV}} \right)$$

$$E^* \simeq 6.9 \times 10^{13} (\epsilon_b/\text{KeV})^{-1} \text{eV}$$

: Proton threshold energy for inelastic interactions with  $\gamma$



- $e:\mu=1:2$  not necessarily correct
- Energy dependence expected





# 2.6 $e/\mu$ ratio vs $\mu/\tau$ ratio & energy spectrum for (3+1)- scheme

Donini-OY '07

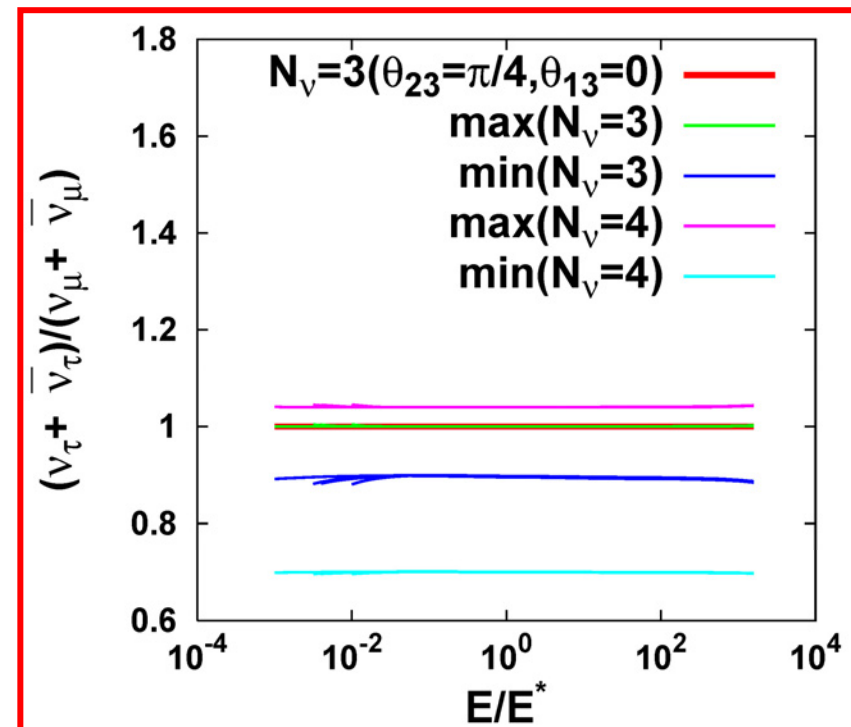
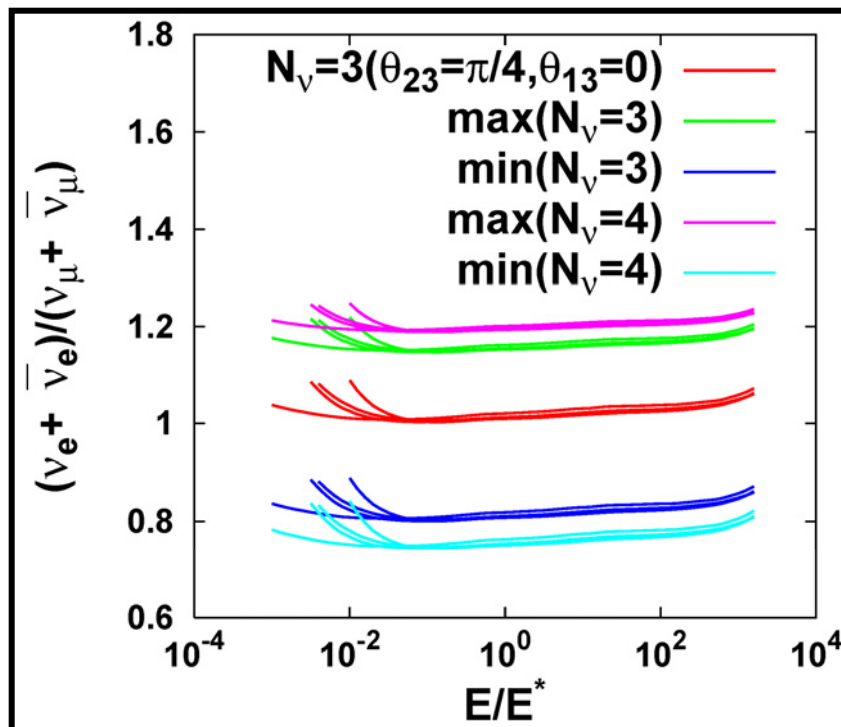
4 curves (energy dependence of  $p, \gamma$ ) corresponding to uncertainties by Lipari et al.

No energy losses

- $(\alpha=2), (\beta_1=1, \beta_2=2)$
- - -  $(\alpha=2), (\beta_1=0.7, \beta_2=2)$
- · - ·  $(\alpha=2), (\beta_1=1, \beta_2=2.4)$
- $(\alpha=2.4), (\beta_1=1, \beta_2=2)$

$\mu/\tau$  ratio is less energy dependent  
(to 1<sup>st</sup> order in small parameters)

$$R_{\tau\mu}^{(4-fam)} = C_{34}^2 = C_{34}^2 R_{\tau\mu}^{(3-fam)}$$

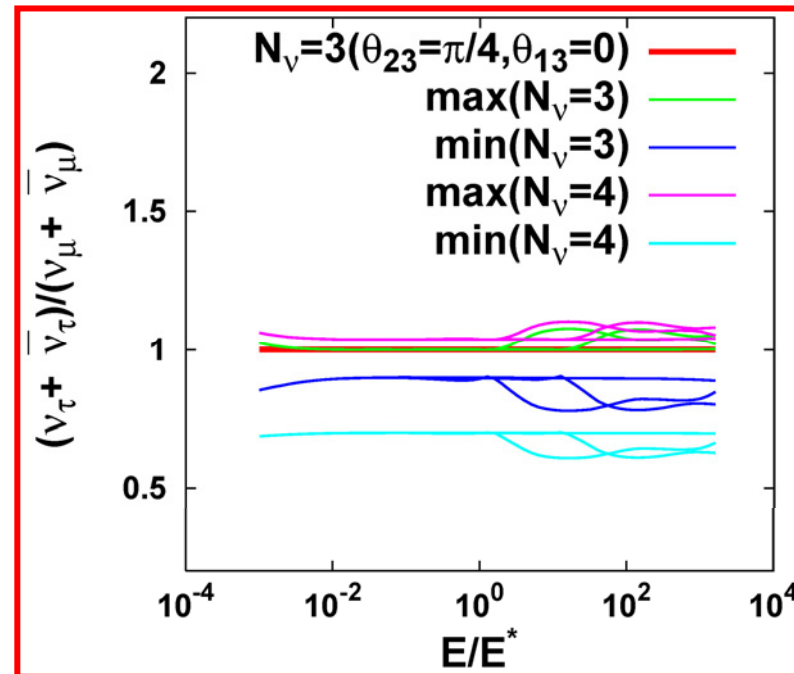
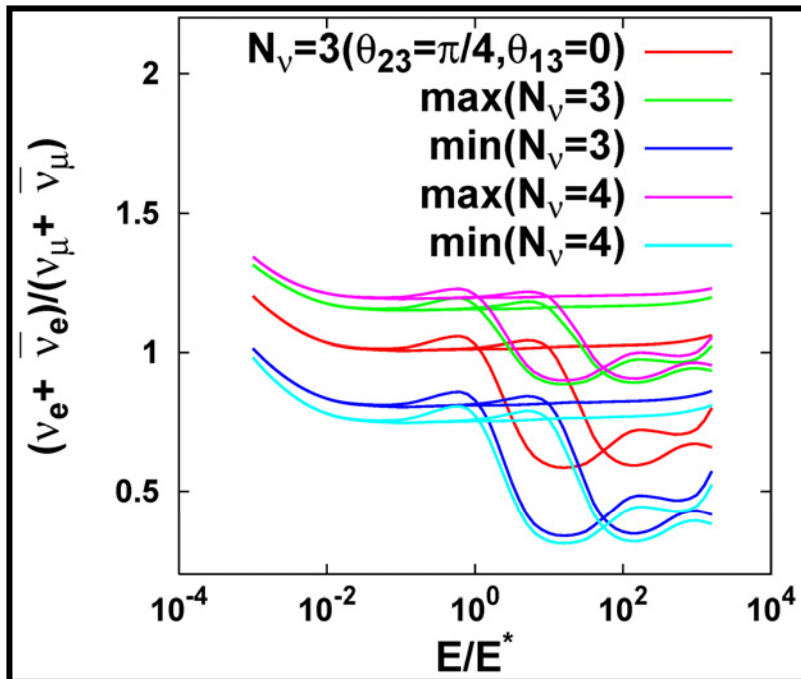


3 curves (muon energy loss)  
corresponding to  
uncertainties by Lipari et al.

$(\alpha=2), (\beta_1=1, \beta_2=2)$

— No Losses  
—  $\epsilon_\mu = 30$   
- - -  $\epsilon_\mu = 3$

Donini-OY '07



- Energy dependence gives us hint on the uncertainties
- $\mu/\tau$  ratio is less energy dependent

$$R_{\tau\mu}^{(4-fam)} = c_{34}^2 = c_{34}^2 R_{\tau\mu}^{(3-fam)}$$

(to 1<sup>st</sup> order in small parameters)

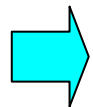
## 2.7 Statistics of expected events

For a typical galactic source, ten years of running at a km<sup>3</sup> water equivalent detector:

$e, \mu$  events  $\sim O(100)$

P. Lipari, astro-ph/0605535

$\tau$  events  $\sim O(30)$



$\mu/\tau$  ratio ( $N_\nu=3$ ) =  $0.30 \mp 0.03(\text{theo}) \mp 0.06(\text{stat})$

$\mu/\tau$  ratio ( $N_\nu=4$ )  $\sim 0.2$

**Statistics is not sufficient at all!**

**A possible way out:**

- to integrate over all the galactic and extragalactic sources, or
- to sum over many GRB's events of similar intensities

## 3. Conclusions

- The **(3+1)-scheme without LSND** constraint predicts flavor ratio of HE cosmic  $\nu$  which could be in principle distinguished from standard case.
- The  **$\mu/\tau$  ratio** suffers relatively less from theoretical uncertainties, and plays an important role to look for signatures of **sterile  $\nu$** .
- Information from **energy spectrum** could be also important to check theoretical uncertainties.
- Statistics from one source is not sufficient to get signatures of sterile  $\nu$ , but if we **sum over data from many sources** then we may be able to say something about **sterile  $\nu$** .