A Consistent Modeling of Neutrino-driven Wind with Accretion Flow onto a Proto-Neutron Stars and its Implications for 56 Ni Production

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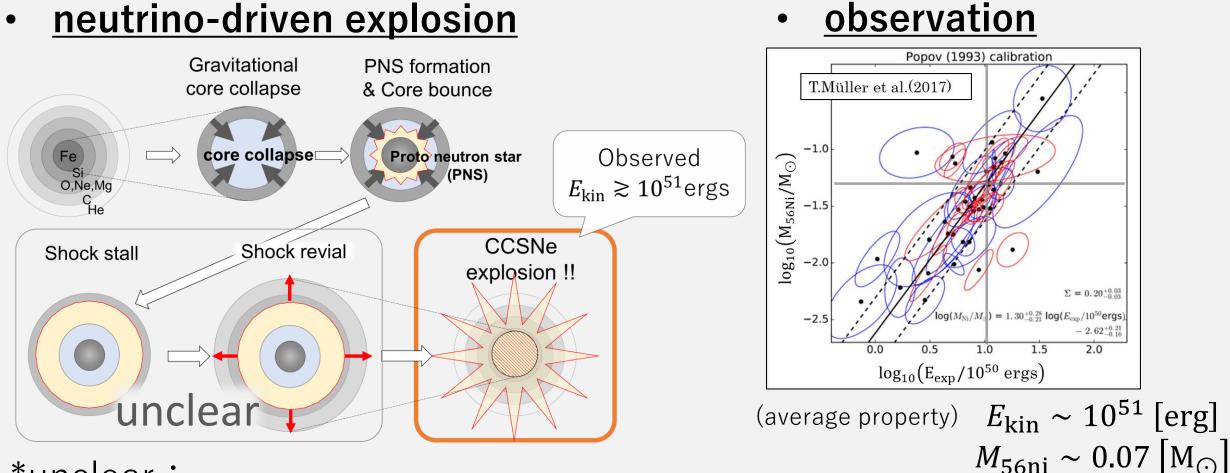
<u>Collaborator</u>: Yudai Suwa (Univ. of Tokyo & YITP, Kyoto Univ.) Keiichi Maeda (Kyoto Univ.)

#### Reference :

- "A Consistent Modeling of Neutrino-driven Wind with Accretion Flow onto a Protoneutron Star and its Implications for 56Ni Production"
   Sawada & Suwa (2021), ApJ, 908, 6 (arxiv. 2010.05615)
- "Nucleosynthesis Constraints on the Energy Growth Timescale of a Core-collapse Supernova Explosion"

Sawada & Maeda (2019), ApJ, 886, 47 (arxiv.1910.06972)

## Explosion mechanism of Core-Collapse SNe



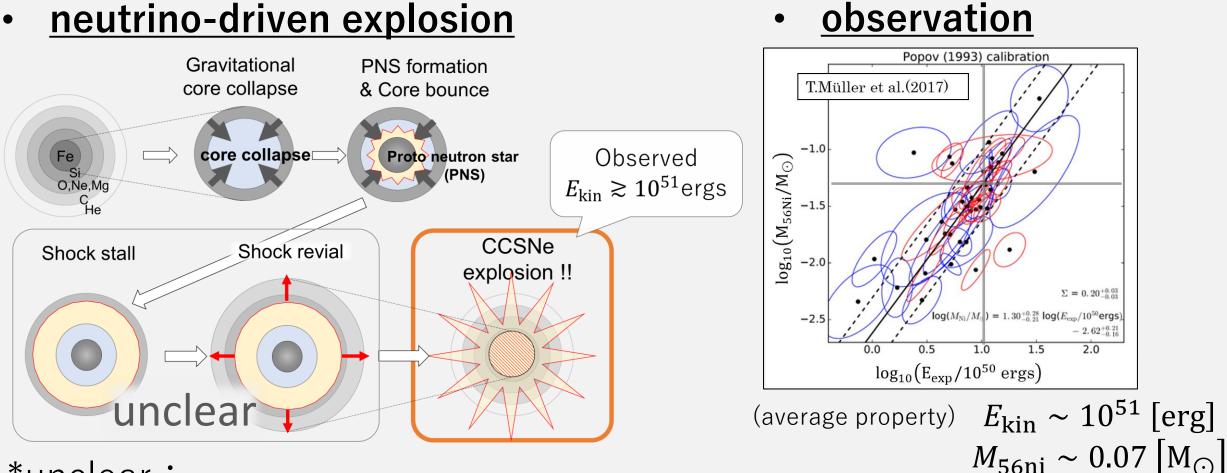
\*unclear:

- 1). Can reach to  $10^{51}$  [erg] ?
- 2). Can synthesize a sufficient amount of 56Ni?

Sawada & Suwa (ApJ, 2021) arxiv. 2010.05615

observation

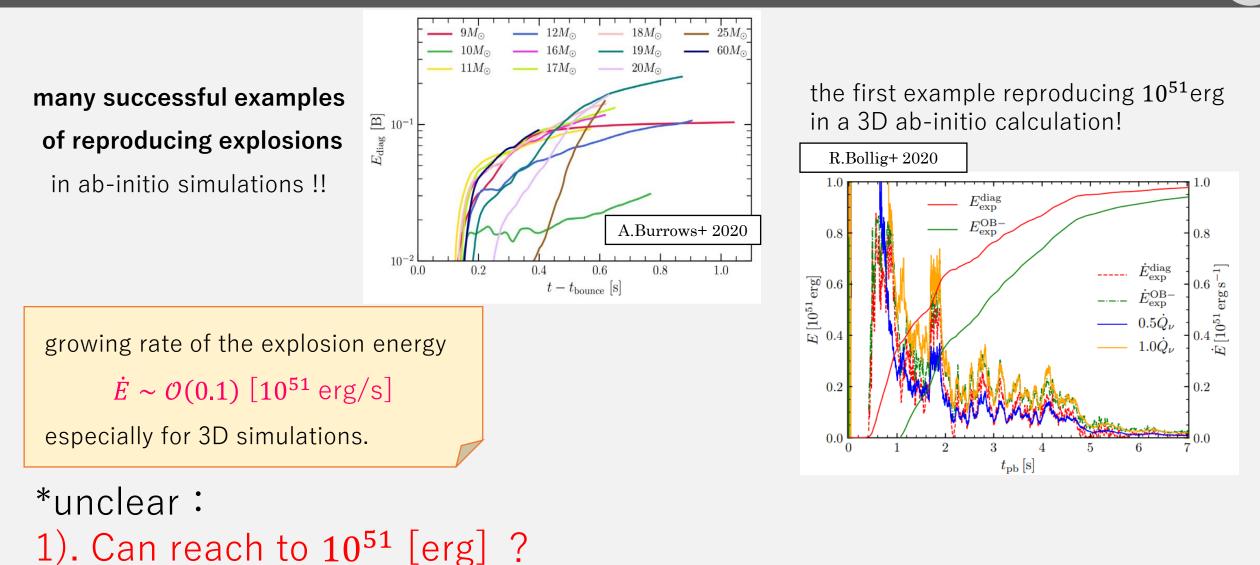
## Explosion mechanism of Core-Collapse SNe



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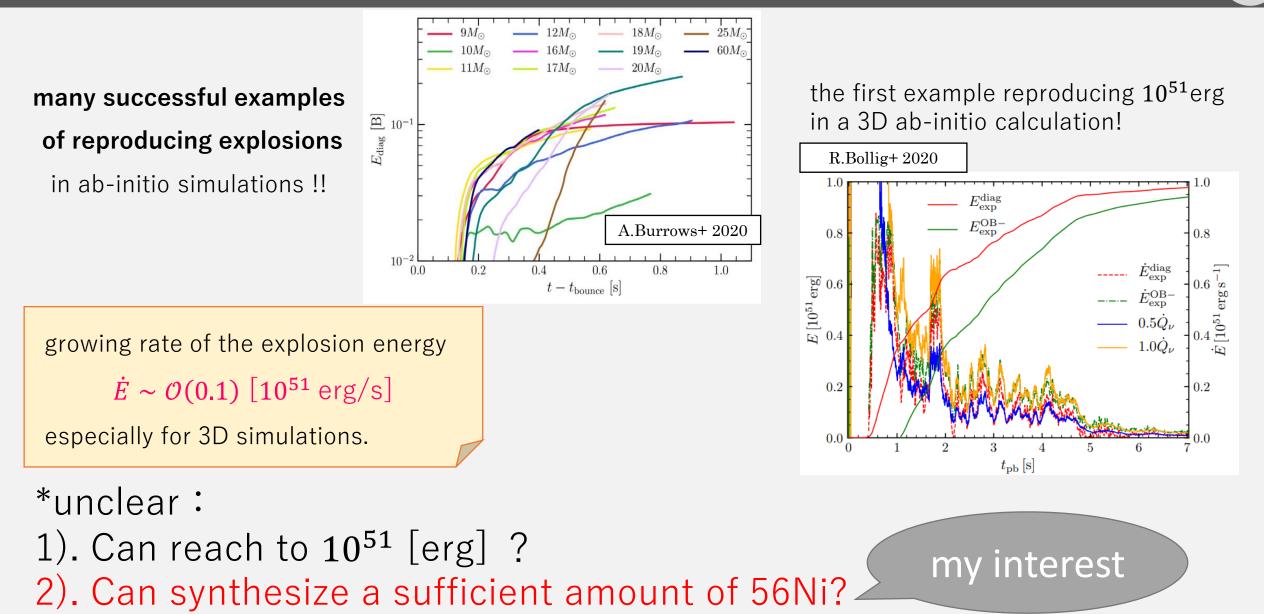
# Current results from the ab-initio calculation



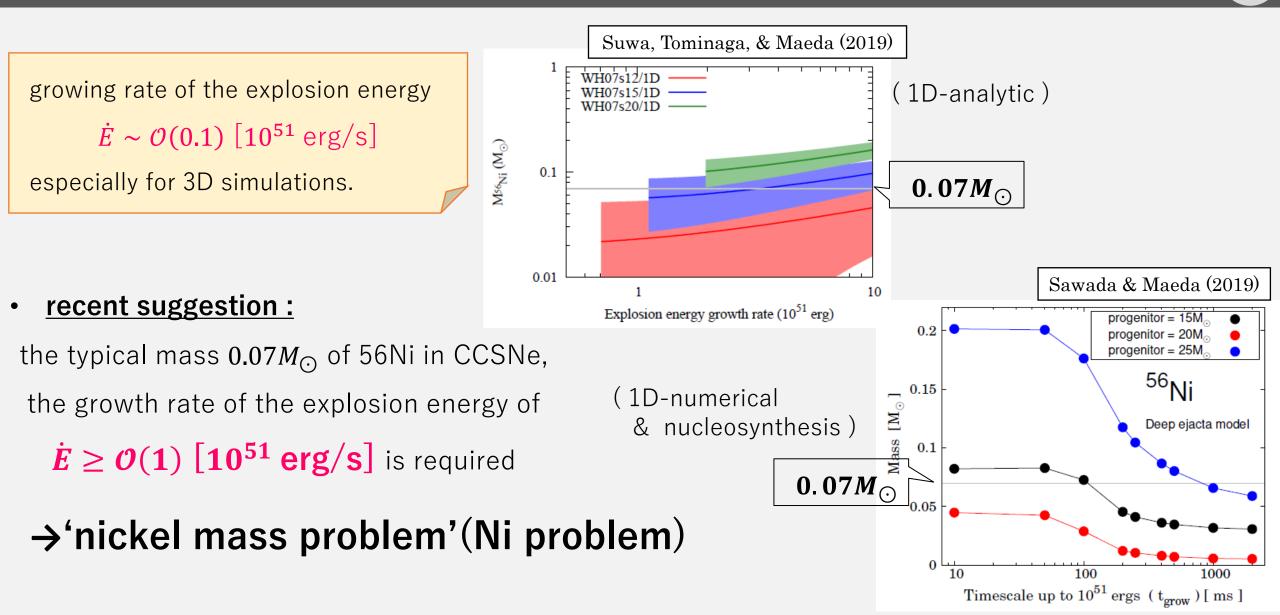
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Sawada & Suwa (ApJ, 2021) arxiv. 2010.05615

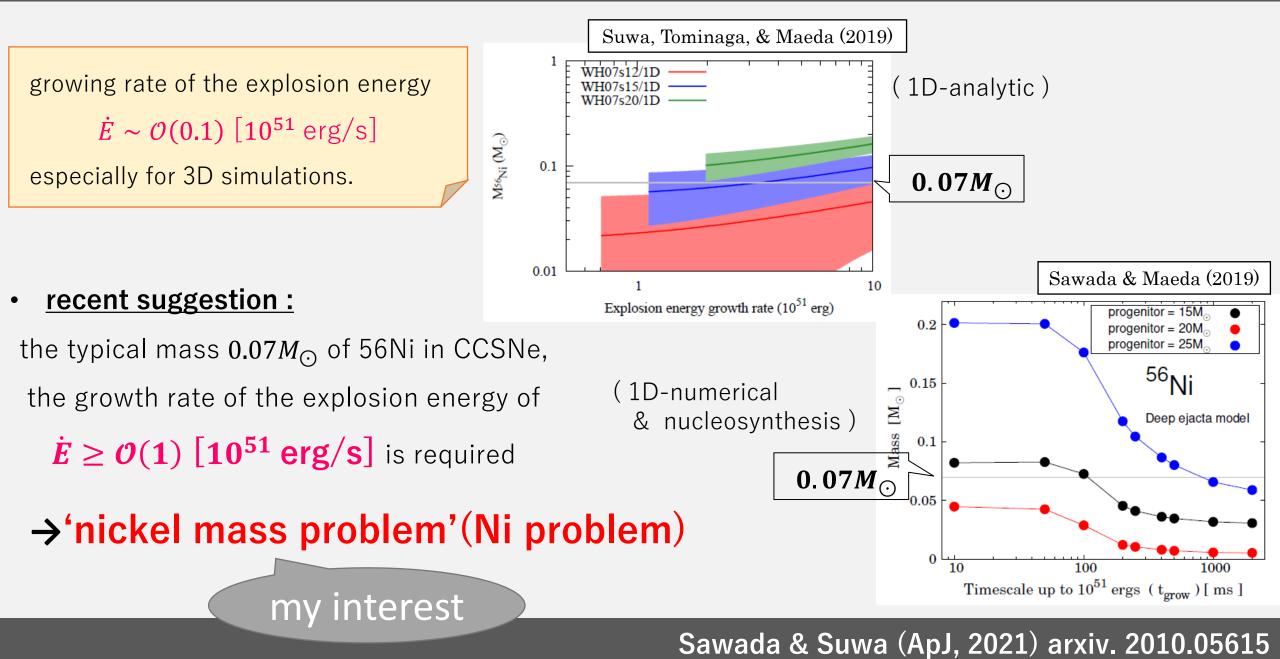
## Current results from the ab-initio calculation



Sawada & Suwa (ApJ, 2021) arxiv. 2010.05615



### Sawada & Suwa (ApJ, 2021) arxiv. 2010.05615



(1) radiation dominant & isothermal @post-shock region,

(2) adiabatic/constant vel. expansion( $r_{\text{shock}} = v_{\text{shock}} \cdot t$ ),

$$E_{\rm exp} = (aT^4) \times \left(\frac{4\pi}{3}r_{\rm shock}^3\right) , \implies T_{\rm peak} \propto t^{-3/4}$$

Suwa, Tominaga, & Maeda (2019) Sawada & Maeda (2019)

5

14Woosley+02  $t_{grow} = 100 ms$ 12 $t_{grow} = 200 ms$ 10 $t_{grow} = 400 ms$  $10^{9}$  $t_{grow} = 1000 ms$ 8  $T_9 [K/$ 6  $\mathbf{2}$ 0 0.10.20.30.4 0.50 (or Radius) time [sec]

### Sawada & Suwa (ApJ, 2021) arxiv. 2010.05615



Deposit 
$$E_{exp} = 10^{51} erg$$
 from the initial setup.  
e.g., Woosley+ 2002

✓ When taking into account the growth timescale of the explosion energy, then

$$E_{\exp}(t) = \frac{10^{51} \text{erg}}{t_{\text{grow}}} \cdot t$$
 this work

5

## (1) radiation dominant & isothermal @post-shock region,

(2) adiabatic/constant vel. expansion( $r_{\text{shock}} = v_{\text{shock}} \cdot t$ ),

$$E_{\rm exp} = (aT^4) \times \left(\frac{4\pi}{3}r_{\rm shock}^3\right) \quad \Rightarrow T_{\rm peak} \propto t^{-3/4}$$

Suwa, Tominaga, & Maeda (2019) Sawada & Maeda (2019)

14 Woosley+02  $t_{grow} = 100 ms$ 12 $t_{grow} = 200 ms$ Instantaneous 10 $t_{grow} = 400 ms$ explosion  $10^{9}$  $t_{grow} = 1000 ms$ 8  $T_9$  [ K / 6  $\mathbf{2}$ 0 0.20.3 0.10.4 0.50 (or Radius) time [sec]

Sawada & Suwa (ApJ, 2021) arxiv. 2010.05615

When the explosion is sufficiently instantaneous, then  
Deposit 
$$E_{exp} = 10^{51} erg$$
 from the initial setup.

e.g., Woosley+ 2002

this work

✓ When taking into account the growth timescale of the explosion energy, then 10<sup>51</sup>erg

$$E_{\rm exp}(t) = \frac{10^{-1} {\rm erg}}{t_{\rm grow}} \cdot t$$

(1) radiation dominant & isothermal @post-shock region,

(2) adiabatic/constant vel. expansion( $r_{\text{shock}} = v_{\text{shock}} \cdot t$ ),

$$E_{\rm exp} = (aT^4) \times \left(\frac{4\pi}{3}r_{\rm shock}^3\right) , \implies T_{\rm peak} \propto t^{-3/4}$$

Suwa, Tominaga, & Maeda (2019) Sawada & Maeda (2019)

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14Woosley+02  $t_{grow} = 100 ms$ 12 $t_{grow} = 200 ms$ Instantaneous 10 $t_{grow} = 400 ms$ explosion  $10^{9}$  $t_{grow} = 1000 ms$ 8  $T_9$  [ K / 6 "slow" explosion 0 0.3 0.10.20.4 0.50 (or Radius) time [sec]

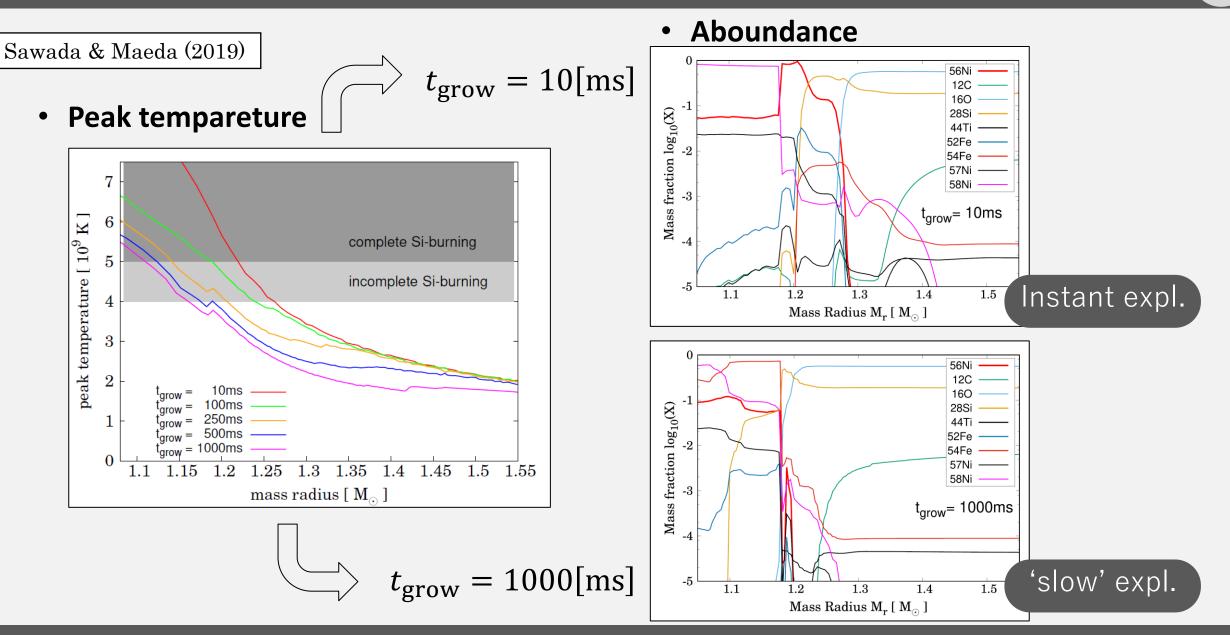
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Deposit 
$$E_{exp} = 10^{51} erg$$
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✓ When taking into account the growth timescale of the explosion energy, then  $E_{exp}(t) = \frac{10^{51} erg}{t_{grow}} \cdot t \quad \text{this work}$ 

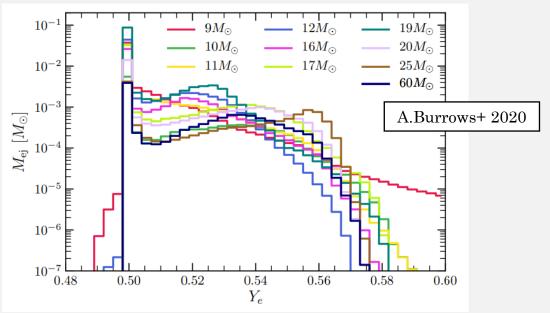


Sawada & Suwa (ApJ, 2021) arxiv. 2010.05615

## • <u>NOTE :</u>

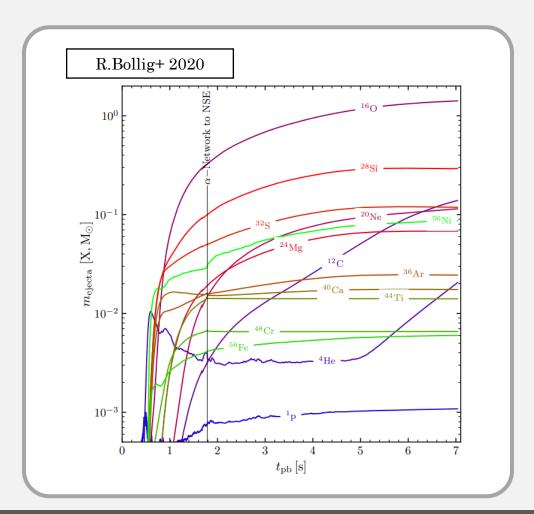
some models in ab-initio simulations have succeeded

in producing the typical mass  $0.07 M_{\odot}$  of 56Ni in CCSNe,

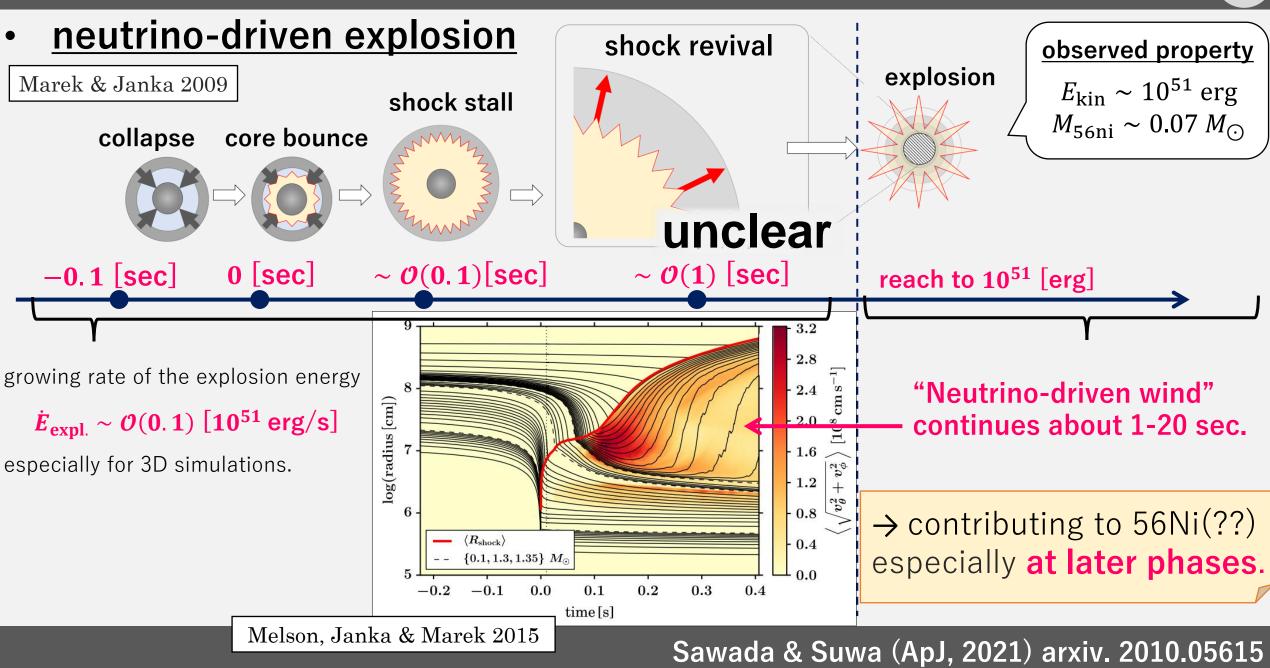


But 'nickel mass problem'(Ni problem)

→ unclear whether we can reproduce sufficient 56Ni amount <u>as a canonical nature.</u>



## Explosion mechanism of Core-Collapse SNe



aim and content of our work

motivation • investigate the potential of the neutrino-driven wind to solve 'Ni problem', especially **at later phases.** 

problem

**1D** (e.g., Wanajo2013) :

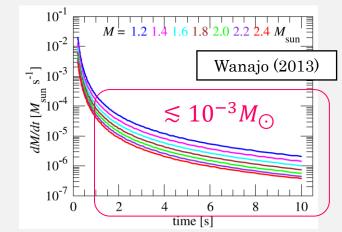
it is already known that 1D wind simulation could not solve Ni problem. ( $\rightarrow$  multi-D, especially energy injection by accretion, is important.)

### • multi-D (e.g.,Wanajo+2018) :

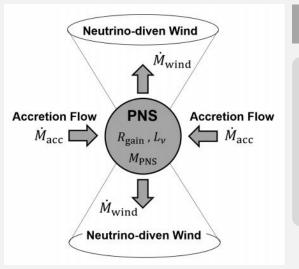
it may solve the Ni problem .

 $\rightarrow$  However, calculation time is limited.

Can it be solved if we follow it to the late explosion stage?

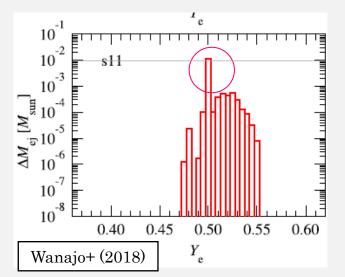


 $\gtrsim 0.01 M_{\odot}$  ejection is required for Ni problem

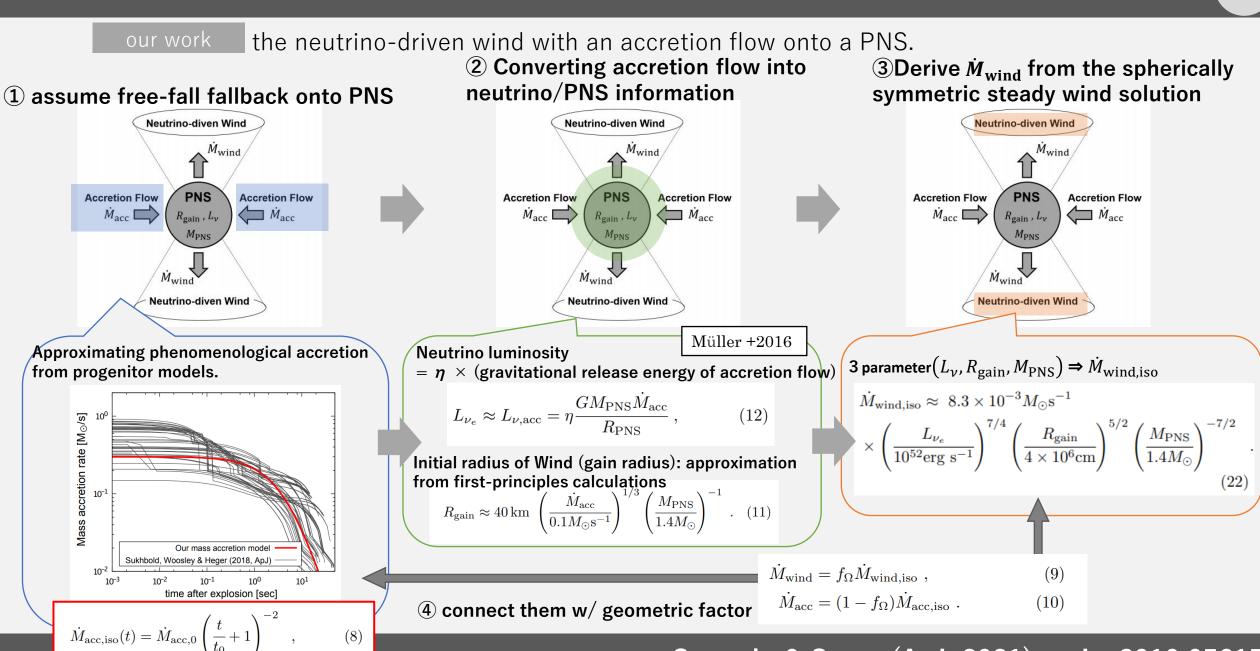


#### our work

- 1. Build a consistent model of the neutrino-driven wind with an accretion flow onto a PNS.
- 2. Investigate the possibility that neutrino-driven wind can solve the "Ni-problem"

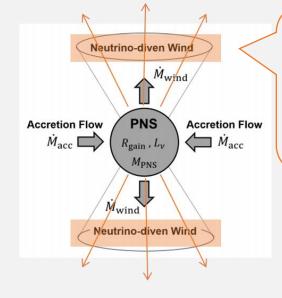


## aim and content of our work



### Sawada & Suwa (ApJ, 2021) arxiv. 2010.05615

# semi-analytic wind model (e.g., Otsuki et al. 2000).



Assume a neutrino-driven wind blowing in the "radial direction" through the low-density region swept by the non-spherically symmetric shock wave

### <u>spherically symmetric steady-state wind</u> <u>solution(e.g., Otsuki et al. 2000).</u>

•  $(R_{\text{gain}}, L_{\nu}, M_{\text{PNS}})$  $\Rightarrow$  tran-sonic solution  $\nu_{\text{tran}} \Leftrightarrow \text{maximum } \dot{M}_{\text{wind}}$  • Helmholtz EoS (Timmes & Swesty 2000), • Boundary condition:  $r = R_{gain} (\dot{Q} \approx 0)$ Given:  $\rho_0 = 10^{10} \text{g cm}^{-3} \cdot L_{\nu,51}^{1/2}$ , (Fujibayashi+ 2015) •  $Y_e = \left[1 + \frac{L_{\overline{\nu}_e}^n \langle \sigma_{\overline{\nu}_e p} \rangle}{L_{\nu_e}^n \langle \sigma_{\nu_e n} \rangle}\right]^{-1} = 0.5, \qquad (4)$ (e.g., Bliss+ 2018).

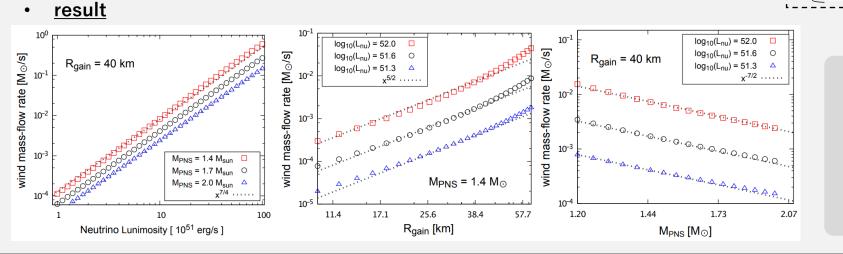
•  $v \frac{dv}{dr} = -\frac{1 + (v/c)^2 - (2GM/c^2r)}{\rho(1 + \epsilon/c^2) + P/c^2} \frac{dP}{dr} - \frac{GM_{\text{PNS}}}{r^2}$ , (2)

(1)

(3)

•  $\dot{M} = 4\pi r^2 \rho v$ ,

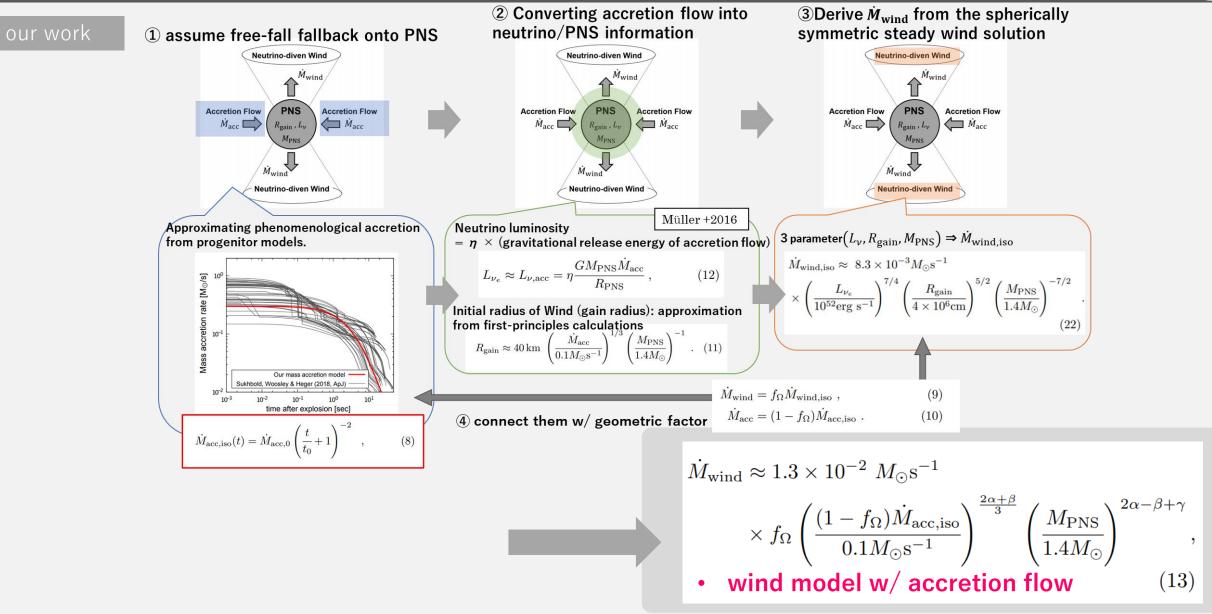
•  $\dot{Q} = v \left( \frac{d\epsilon}{dr} - \frac{P}{\rho^2} \frac{d\rho}{dr} \right) ,$ 



### • spherical wind $\dot{M}_{wind,iso}$

 $\dot{M}_{\rm wind,iso} \approx 8.3 \times 10^{-3} M_{\odot} {\rm s}^{-1} \\ \left(\frac{L_{\nu_e}}{10^{52} {\rm erg \ s}^{-1}}\right)^{\alpha} \left(\frac{R_{\rm gain}}{4 \times 10^6 {\rm cm}}\right)^{\beta} \left(\frac{M_{\rm PNS}}{1.4 M_{\odot}}\right)^{\gamma},$ (5)

## result



maximum parameter sets

 $M_{\rm PNS,0} \ge 1.4 M_{\odot},$ 

total ejected mass of the wind... ٠

 $\dot{M}_{\rm acc,0} < 1.0 M_{\odot} s^{-1}$  $M_{\rm ej,\infty} = \int_0^\infty dt \, \dot{M}_{\rm wind}$ Total accretion mass  $< 0.7 M_{\odot}$  $\approx 4.3 \times 10^{-3} M_{\odot} s^{-1} f_{\Omega} (1 - f_{\Omega})^2 \left(\frac{\dot{M}_{\text{acc},0}}{0.1 M_{\odot} s^{-1}}\right)^2 \left(\frac{M_{\text{PNS},0}}{1.4 M_{\odot}}\right)^{-\frac{3}{2}} \le 0.067 M_{\odot}^{2}$ 

the time evolution of the cumulative ejected mass of the wind... ٠

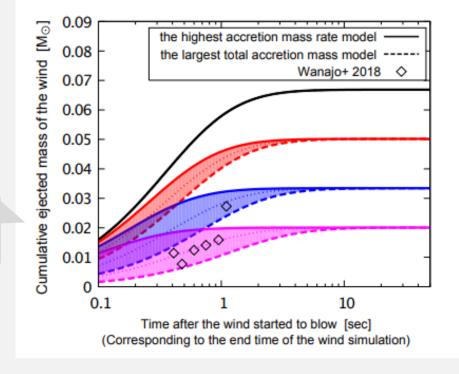
$$M_{\rm ej}(t_e) = \int_0^{t_e} dt \, \dot{M}_{\rm wind}$$
  
\$\approx 6.4 \times 10^{-4} M\_{\overlinessim} \left[ 1 - \left( \frac{t\_0}{t\_0 + t\_e} \right)^3 \right] \times \left( \frac{t\_0}{1s} \right) \left( \frac{\dot{M}\_{\rm acc,0}}{0.1M\_{\odot}s^{-1}} \right)^2 \left( \frac{M\_{\rm PNS,0}}{1.4M\_{\odot}} \right)^{-5/2}

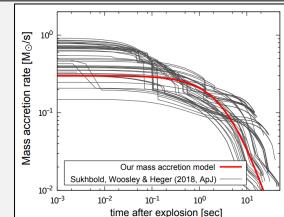
- Conclusion.
- the total ejectable is determined within  $\sim 2$  sec from the onset of the explosion. 1.
- the supplementable amount at a late phase (t > 1 sec) remains  $M_{ei} < 0.01 M_{\odot}$ . 2.

#### $\rightarrow$ difficult to solve the Ni problem

at the late phase of the explosion by the neutrino-driven wind.

Mass accretion rate [ $M_{\odot}$ /s] 10-1 Our mass accretion mode Sukhbold, Woosley & Heger (2018, ApJ 100 time after explosion [sec]





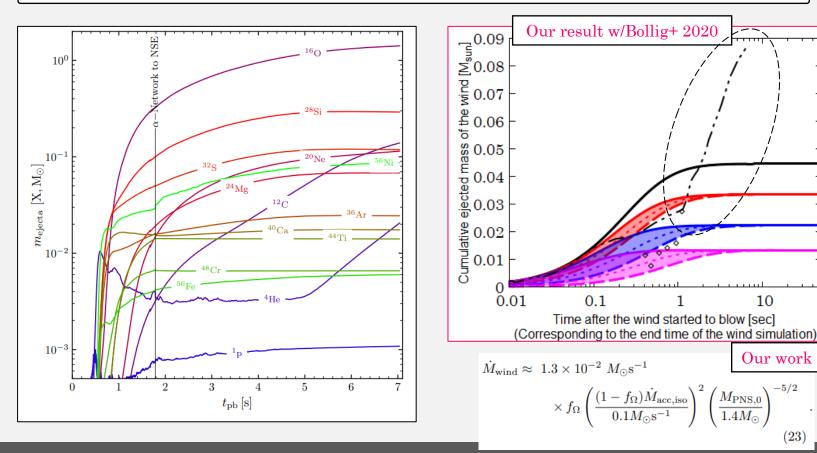
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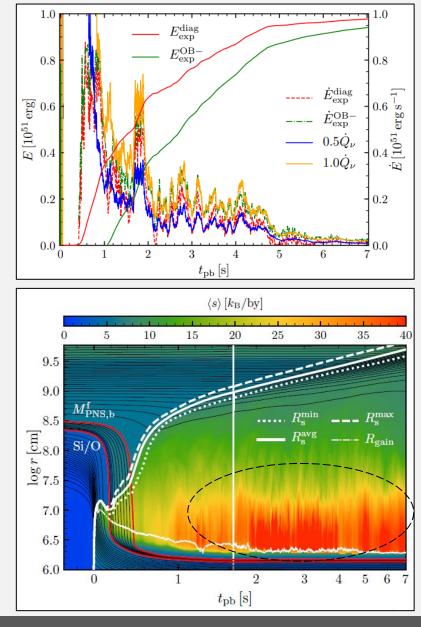
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• **NOTE :** R.Bollig+ 2020

"The converged value of the explosion energy at infinity (with overburden subtracted) is roughly 1B and the ejected 56Ni mass **up to 0.087** solar masses"

"Our final 56Ni mass is therefore an upper limit, and we expect the actual mass to be around **0.05 M**. Nevertheless, it demonstrates that 56Ni masses in the ballpark of those of typical CCSNe can be ejected in 3D neutrino-driven explosions."





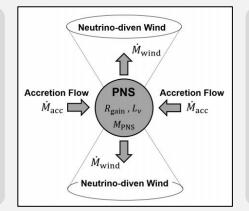
## Summary

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1. a model of the neutrino-driven wind with an accretion flow onto a PNS.

### spherical wind

$$\dot{M}_{\rm wind,iso} \approx 8.3 \times 10^{-3} M_{\odot} {\rm s}^{-1} \\ \times \left(\frac{L_{\nu_e}}{10^{52} {\rm erg \ s}^{-1}}\right)^{7/4} \left(\frac{R_{\rm gain}}{4 \times 10^6 {\rm cm}}\right)^{5/2} \left(\frac{M_{\rm PNS}}{1.4 M_{\odot}}\right)^{-7/2}$$
(22)



### wind model w/ accretion flow

$$\dot{M}_{\text{wind}} \approx 1.3 \times 10^{-2} M_{\odot} \text{s}^{-1} \times f_{\Omega} \left( \frac{(1 - f_{\Omega}) \dot{M}_{\text{acc,iso}}}{0.1 M_{\odot} \text{s}^{-1}} \right)^2 \left( \frac{M_{\text{PNS},0}}{1.4 M_{\odot}} \right)^{-5/2}$$
(23)

2. the possibility that neutrino-driven wind can solve the "Ni-problem"

- 1. the total ejectable is determined within  $\sim$ 2 sec from the onset of the explosion.
- 2. the supplementable amount at a late phase (t > 1 sec) remains  $M_{ej} < 0.01 M_{\odot}$ .

### ightarrow difficult to solve the Ni problem

at the late phase of the explosion by the neutrino-driven wind.

