

From zero flavors to twelve: the opening of the conformal window

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YITP Workshop
QCD phase diagram and lattice QCD

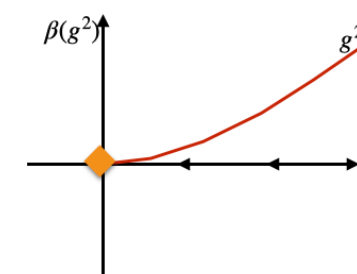
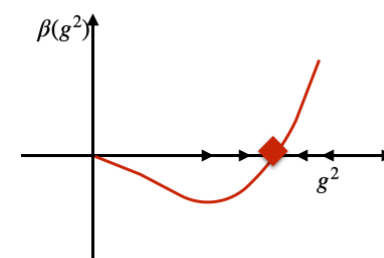
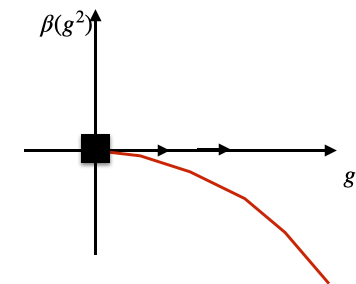
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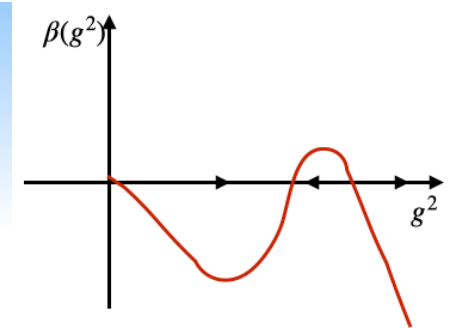
View of Kyoto

The phase diagram as the function of N_f

- ▶ $N_f = 0$ YM is confining (zero temperature)
- ▶ QCD with 2 light flavors is chirally broken, confining (zero temperature)
- ▶ As N_f increases a phase transition to a conformal phase occurs ($N_f^* \approx 8 - 10$ for SU(3) color, fundamental fermions)
 - the conformal phase is characterized by an infrared fixed point
 - there might be a UV fixed point with a new relevant operator (4-fermion?)
- ▶ At $N_f = 16.5$ (SU(3) color) asymptotic freedom is lost
 - the system could be “trivial”
 - or a UV-safe fixed point might emerge (again, new relevant operator)

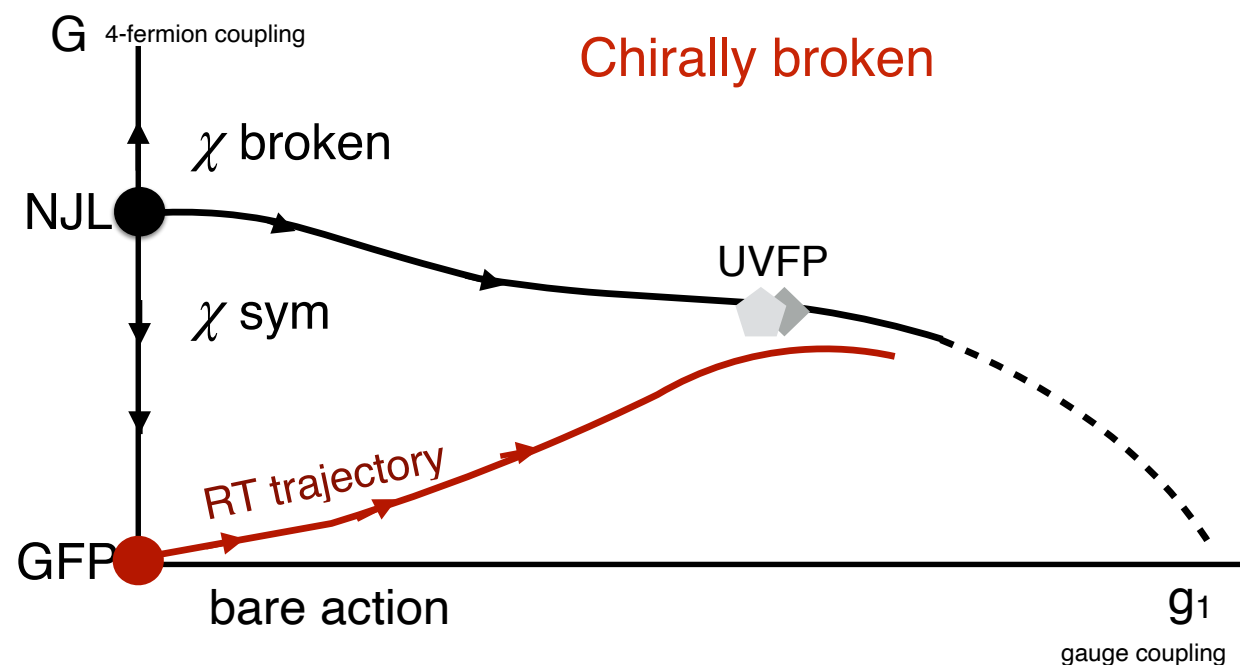
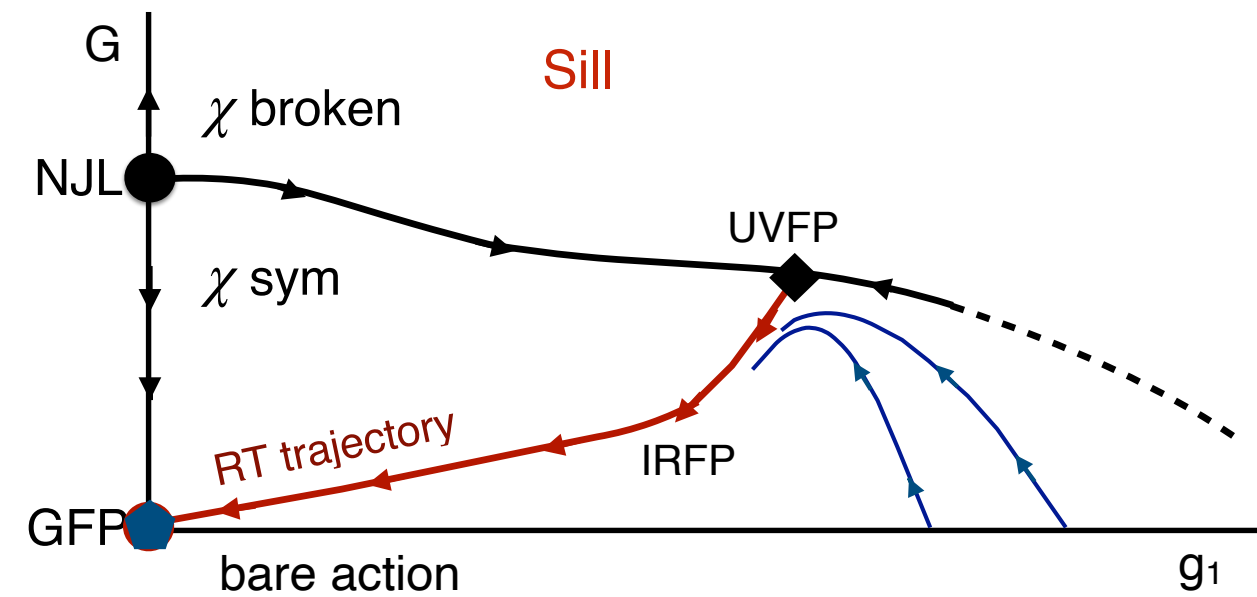
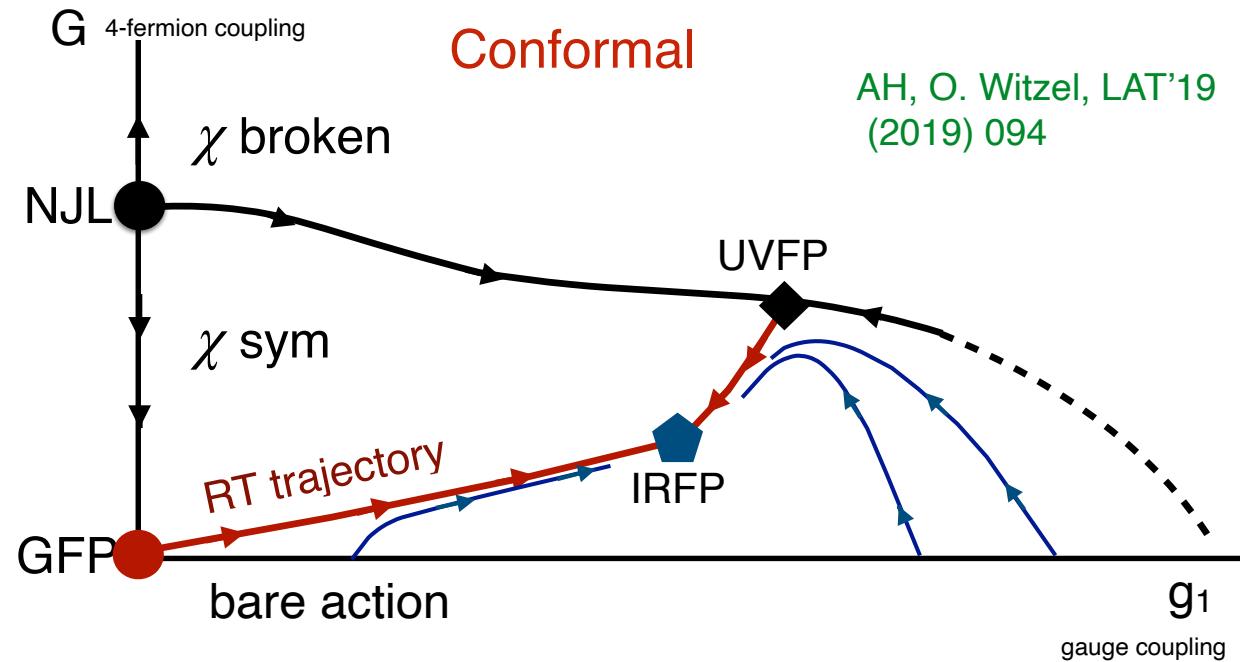


Phase structure of (near-)conformal systems



Conjectured phase diagram in the extended parameter space

Kaplan et al, *Phys.Rev.D* 80 (2009) 125005
 Gorbenko et al, *JHEP* 10 (2018) 108



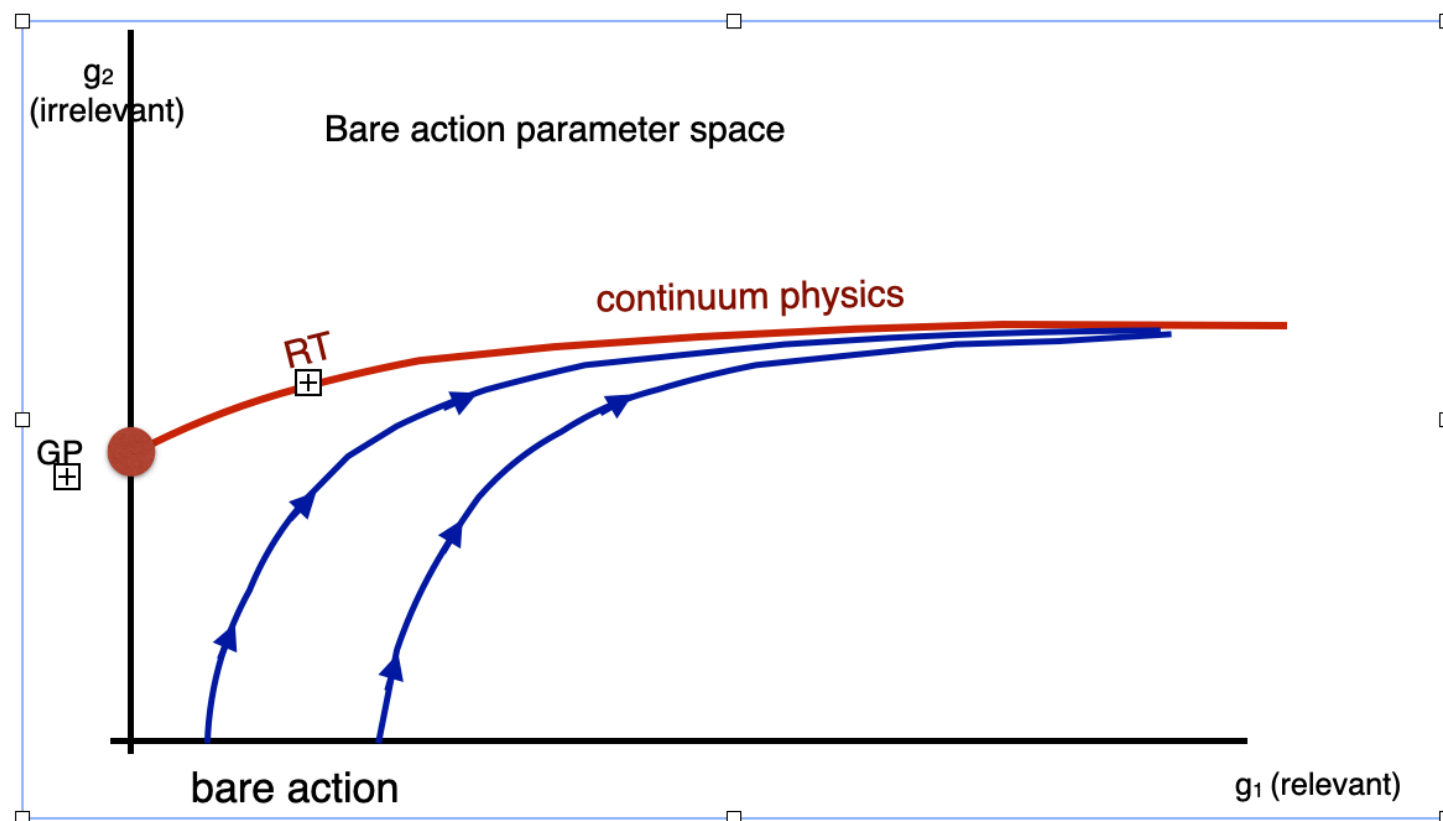
Goal is to determine the RG flow along the renormalized trajectory:

- this is a 1-parameter function: RG β fn
- can be done without NJL action

The continuous β function on the lattice

In Wilsonian RG language:

- ◆ **RG transformation:** determines the FP and its renormalized trajectory (**RT**) (continuum physics)
- ◆ **Lattice action:** starting point of RG flow
 - along the **RT** RG flows from different bare couplings overlap, and describe continuum physics
- ◆ **Operator:** g_{GF}^2 should correspond to scaling operator along the **RT**
 - $g_{GF}^2 \propto \mu^{-4} \langle E \rangle$ (no anomalous dimension)



gauge-fermion only

The continuous β function from gradient flow

GF is a smoothing transformation (Luscher JHEP 2010, 071)

It can be used to define a real-space Wilsonian RG transformation: $\mu \propto 1/\sqrt{8t}$,

$$\beta(g_{GF}) = -t \frac{dg_{GF}^2}{dt}$$

A. Carosso, AH, E. Neil,
PRL 121,201601 (2018)

Sonoda, H., Suzuki,
H. PTEP,023B05 (2021)

Lattice details:

The RG picture is valid

- in infinite volume
- in $am = 0$ chiral limit

The continuum limit is $t/a^2 \rightarrow \infty$

- the flows approach the **RT**
- the correct scaling operator is projected out
- Remaining cut-off effects are removed by $a^2/t \rightarrow 0$

AH, O. Witzel, *Phys.Rev.D* 101 (2020) 3

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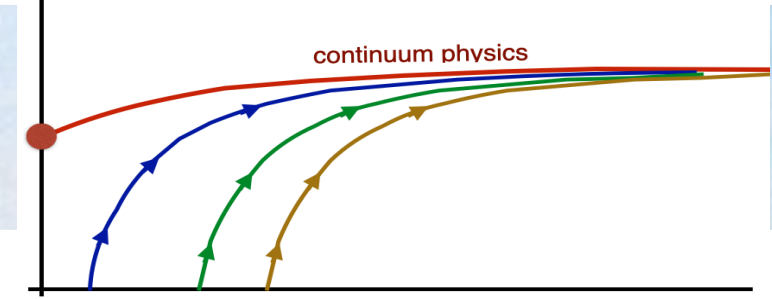
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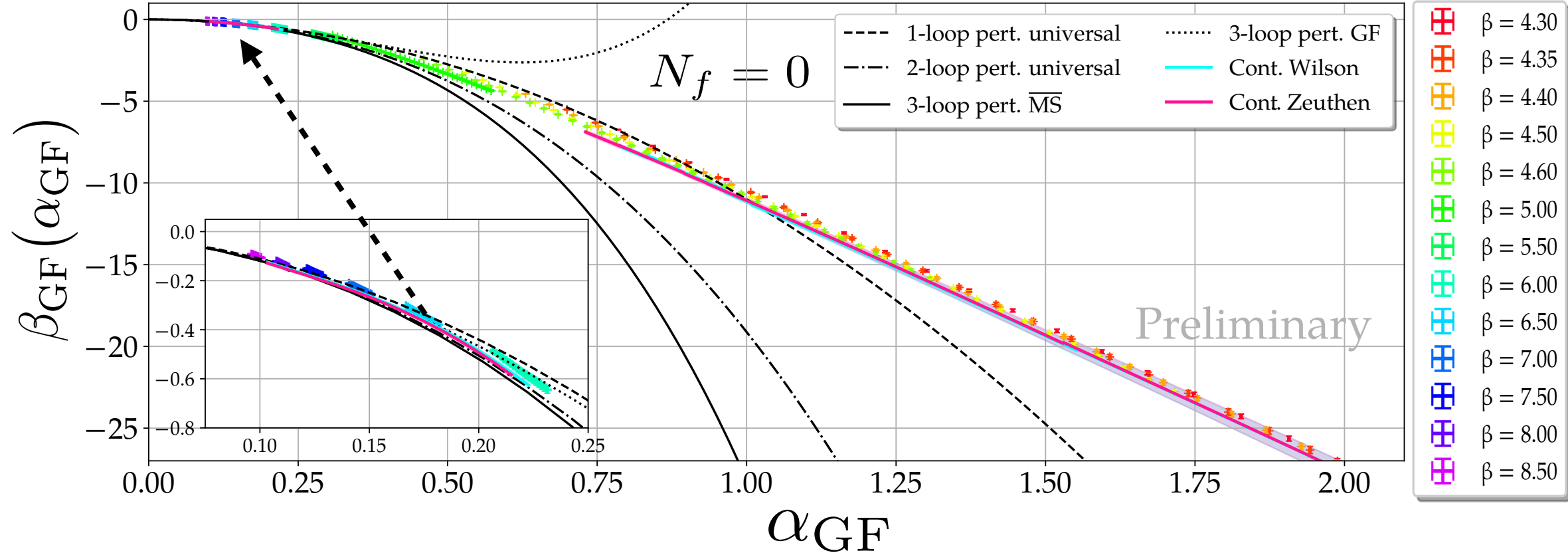
AH, O. Witzel, *Phys.Rev.D* 101 (2020) 3

- **Step scaling function** requires that the volume is the only scale; valid only in the deconfined regime
- **Continuous β function** requires infinite volume extrapolation but it is correct even in the confining regime

The continuous β function, $N_f = 0$

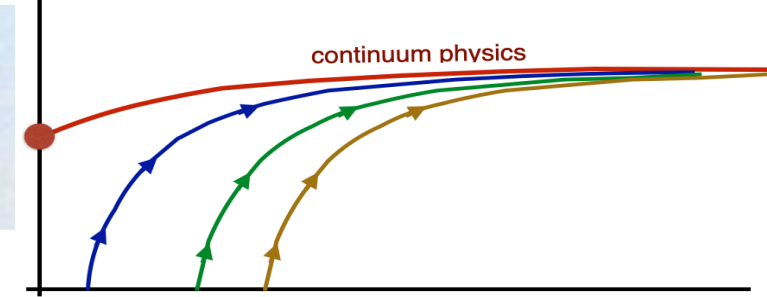


C. Peterson et al 2109.09720 (Lat'21)

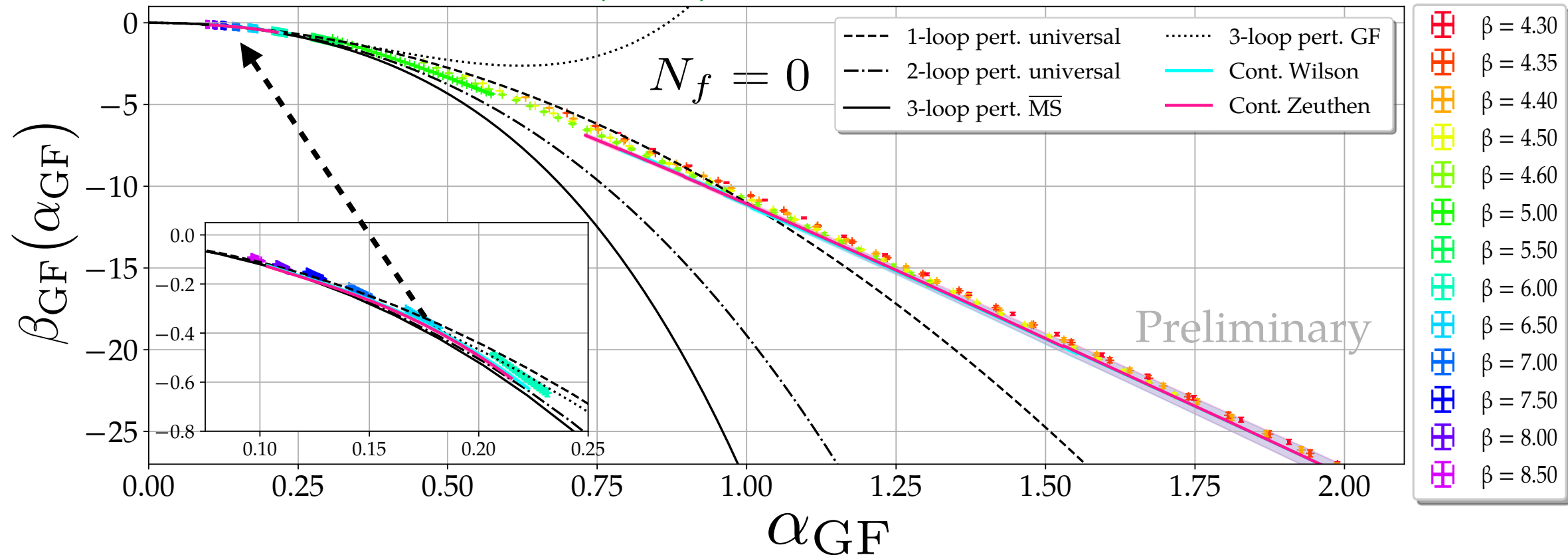


- ◆ Different bare couplings overlap, form a unique curve :
 - flow reached the RT
 - cutoff effects are small

The continuous β function, $N_f = 0$



C. Peterson et al 2109.09720 (Lat'21)



Nakamura, Schierholz
2106.11369

◆ In the confining regime $\beta(g^2) \propto g^2$: non-perturbative; Topology?

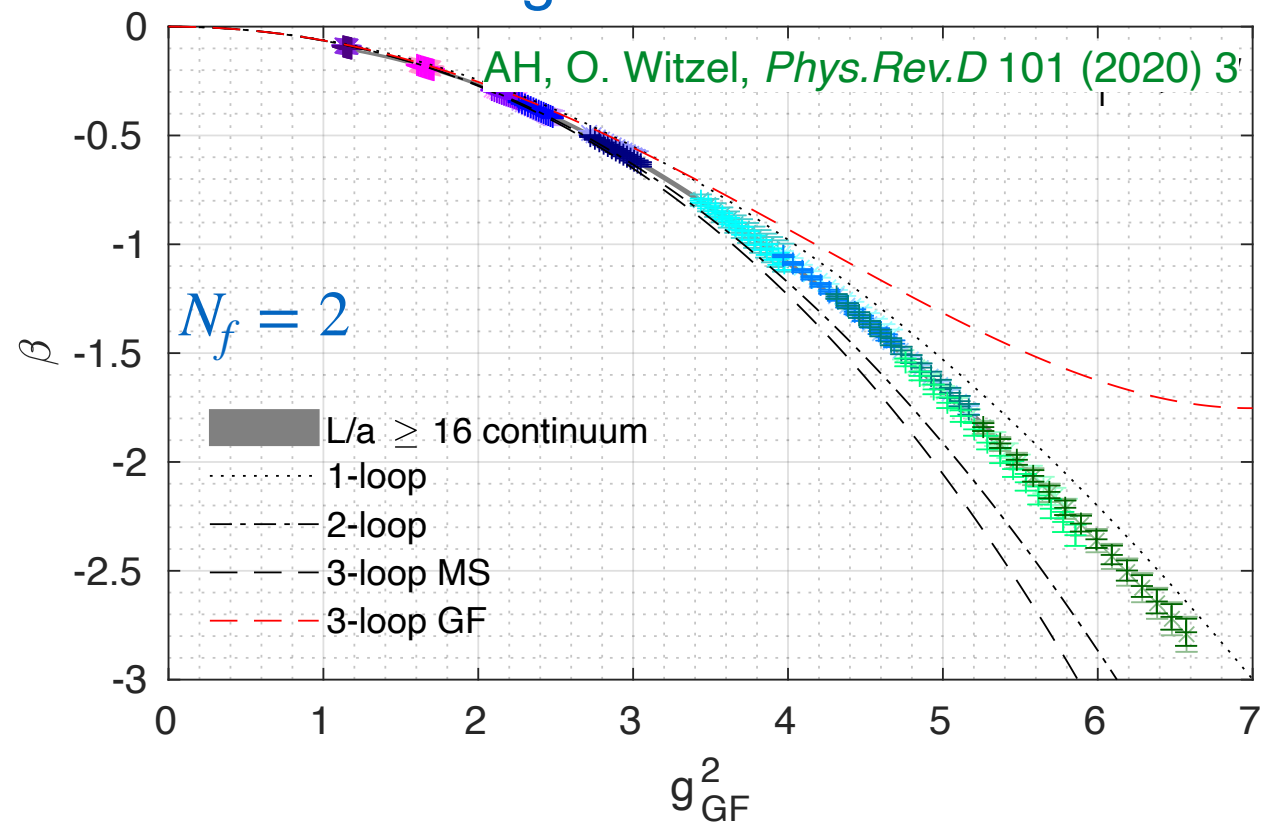
$$N_T = N_I + N_A \text{ instantons} \rightarrow \text{vacuum energy density: } \langle E \rangle = \langle E \rangle_{N_T=0} + N_T * S_I/V$$

$$\text{GF coupling: } g_{GF}^2 = g_{GF,0}^2 + ct^2 N_T/V$$

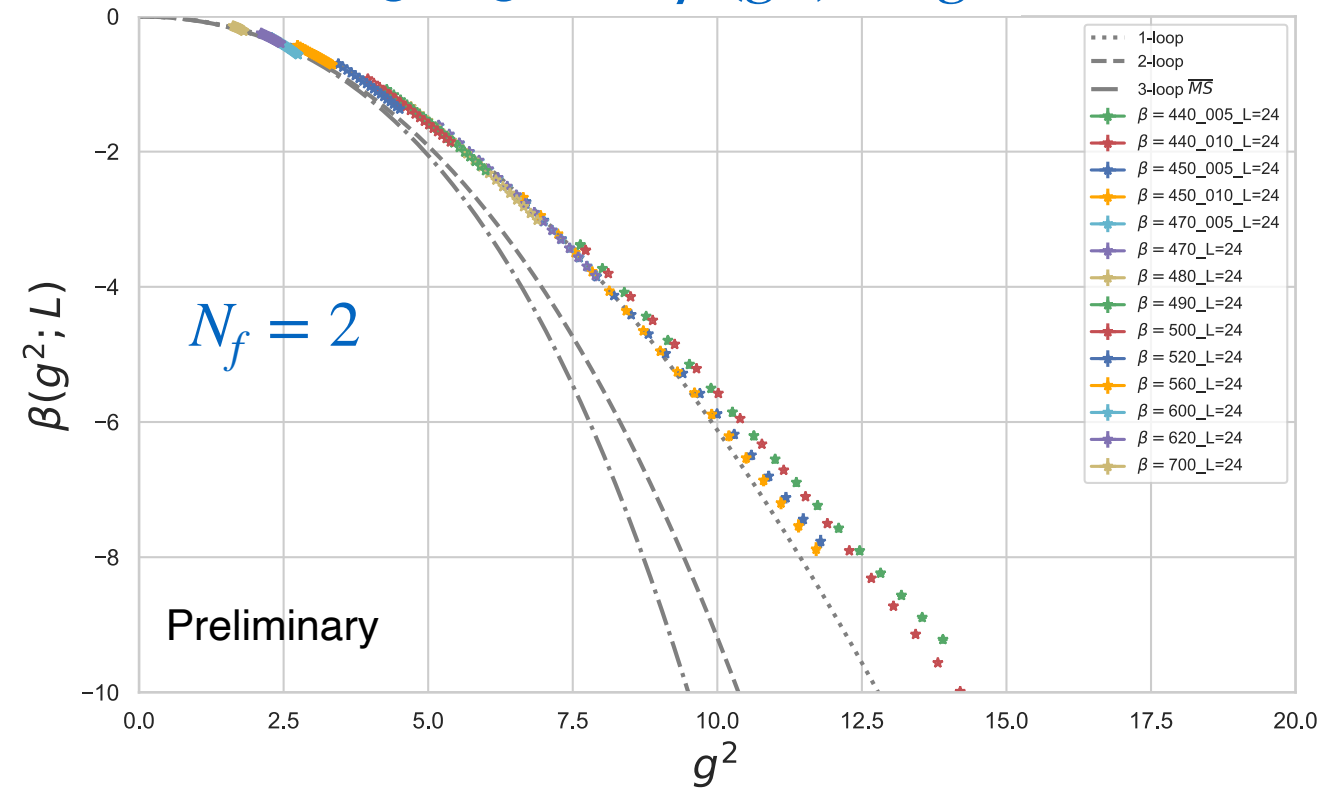
If $\lim_{V \rightarrow \infty} N_T/V = \text{finite}$, topology dominates and $\beta(g^2) \propto g^2$

The continuous β function, $N_f = 2$ DWF, $m_f = 0$

Deconfined regime

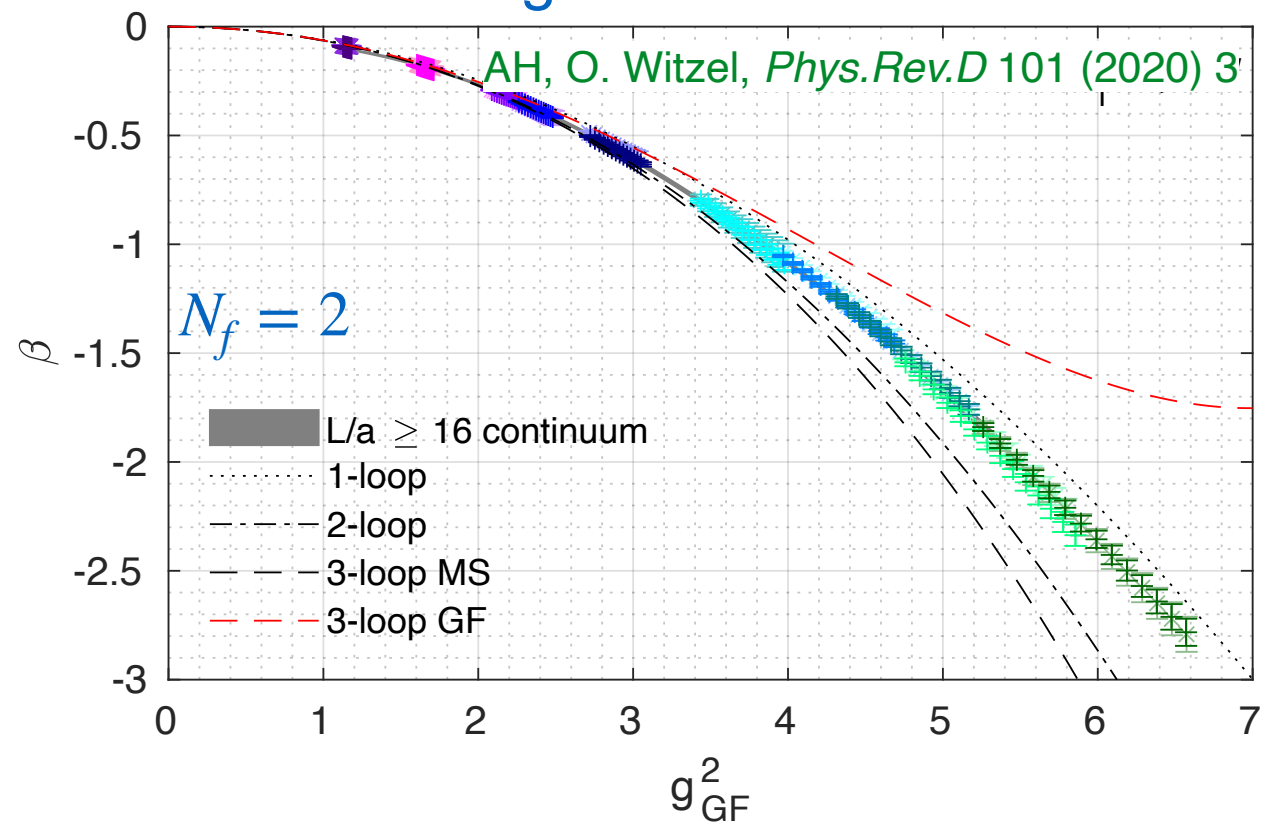


Confining regime : $\beta(g^2) \propto g^2$

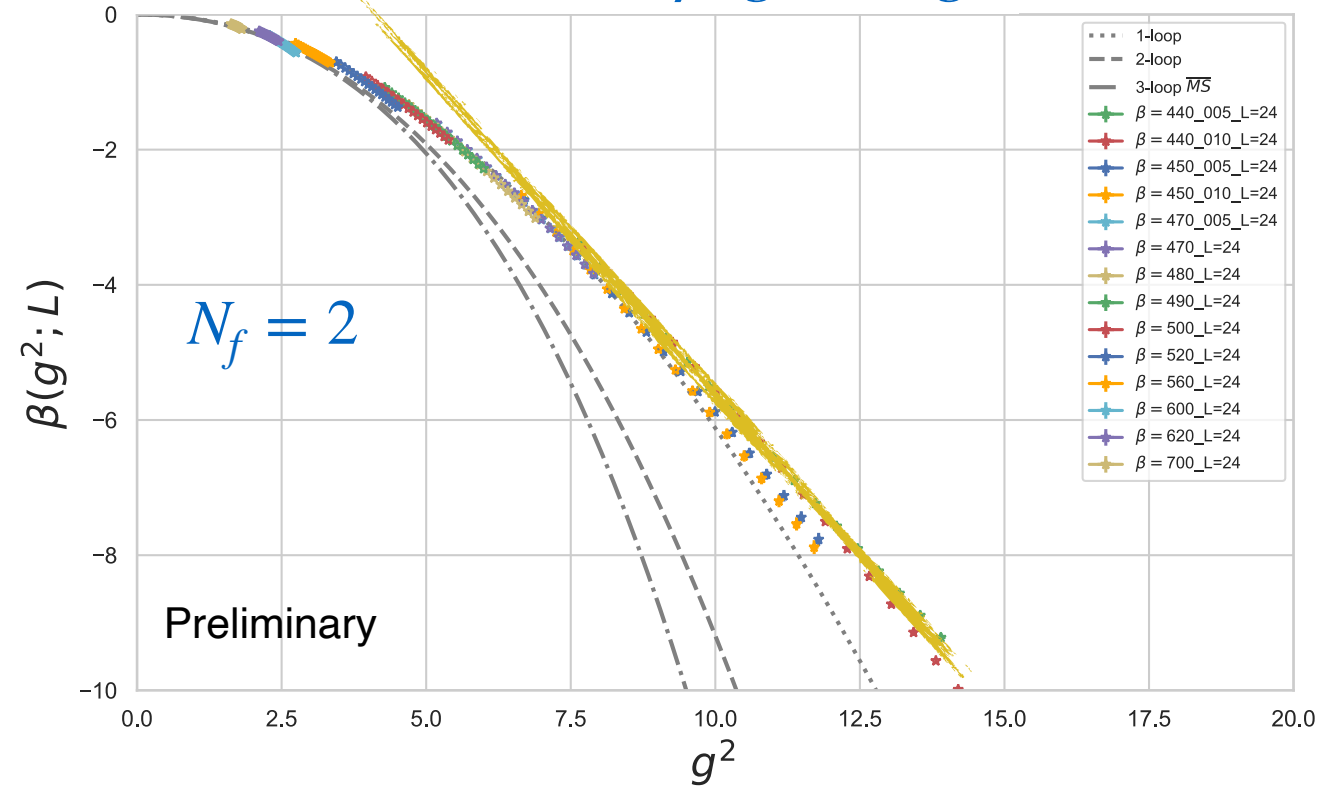


The continuous β function, $N_f = 2$ DWF, $m_f = 0$

Deconfined regime

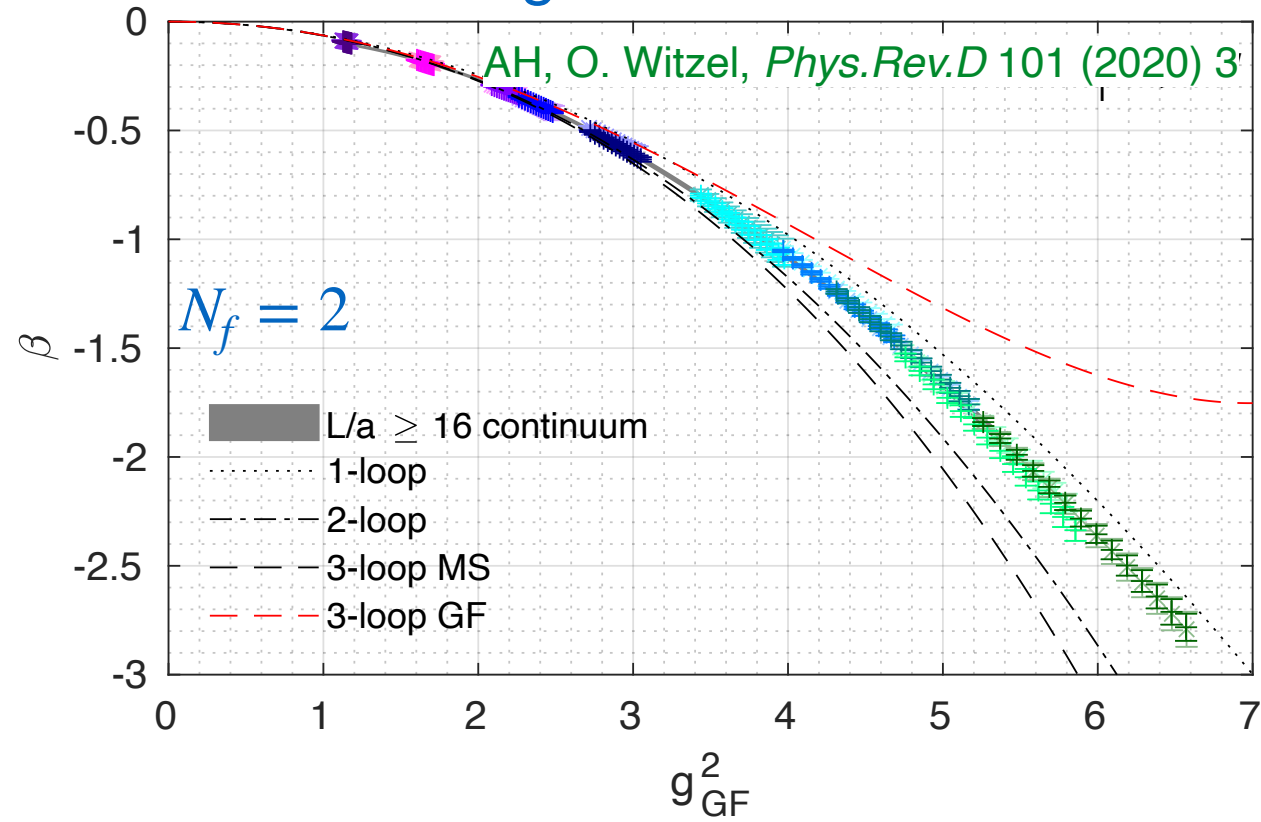


Confining regime : $\beta(g^2) \propto g^2$

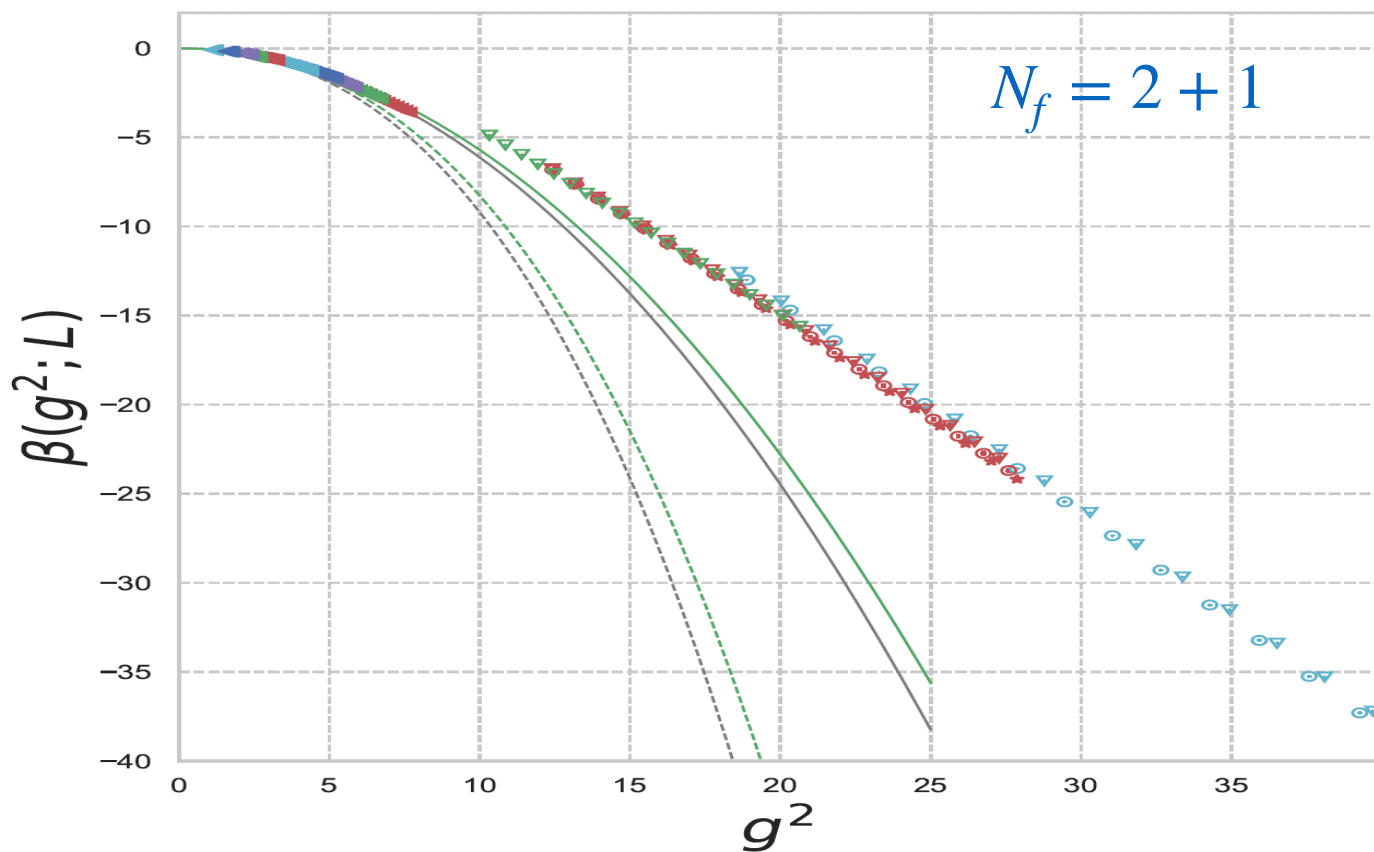
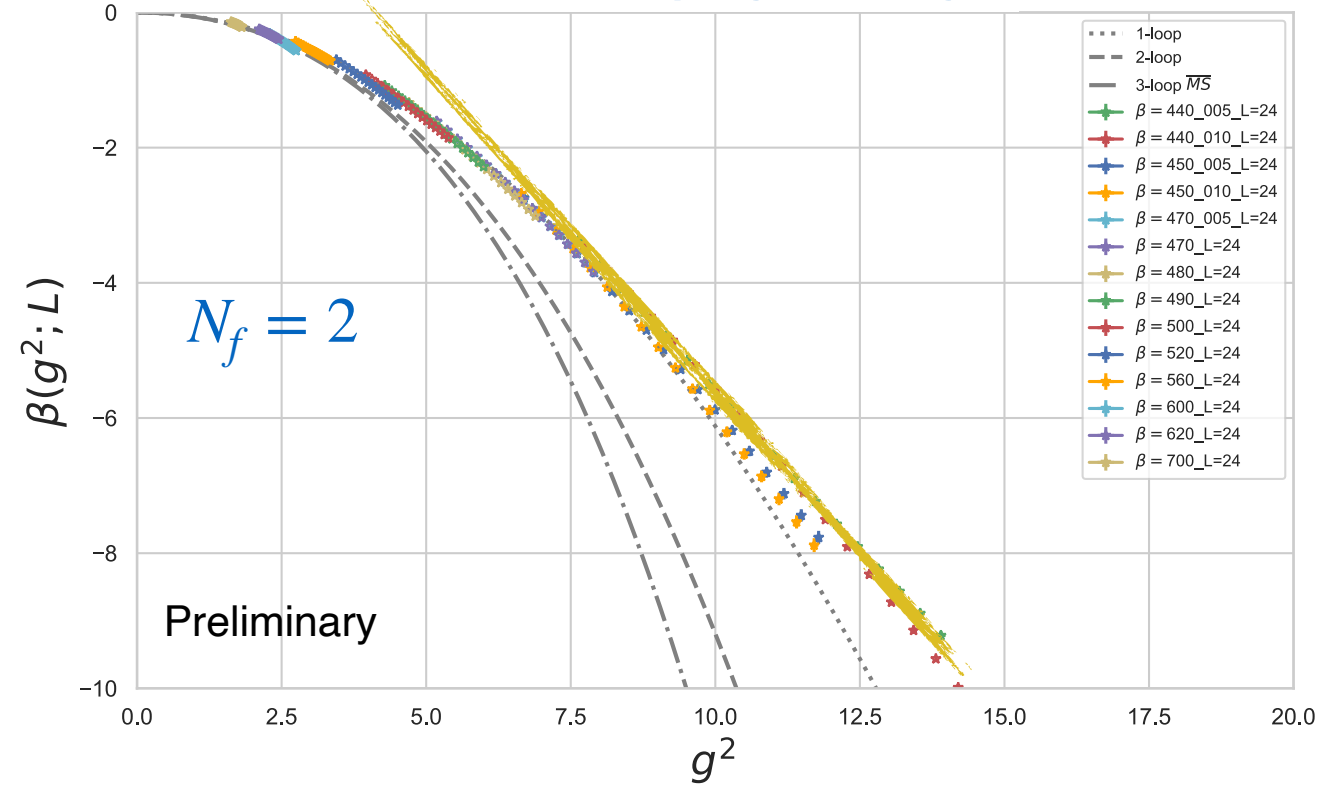


The continuous β function, $N_f = 2, 2 + 1$ DWF, $m_f = 0$

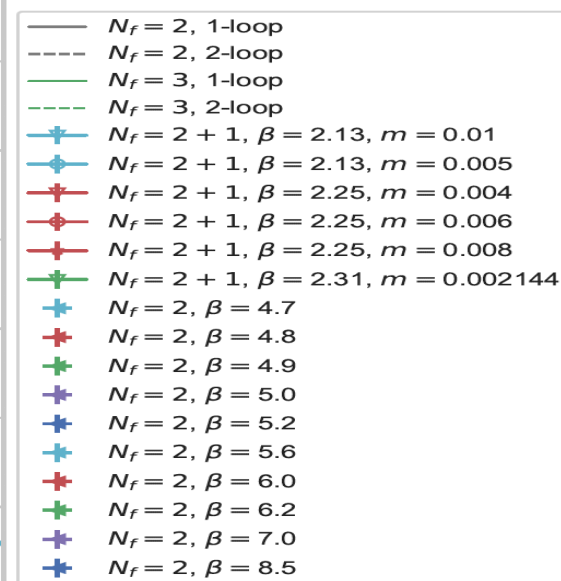
Deconfined regime



Confining regime : $\beta(g^2) \propto g^2$



C. Monahan, Lat'21



Confining regime :
 $\beta(g^2) \propto g^2$ for
 $N_f = 0, 2, 2 + 1$

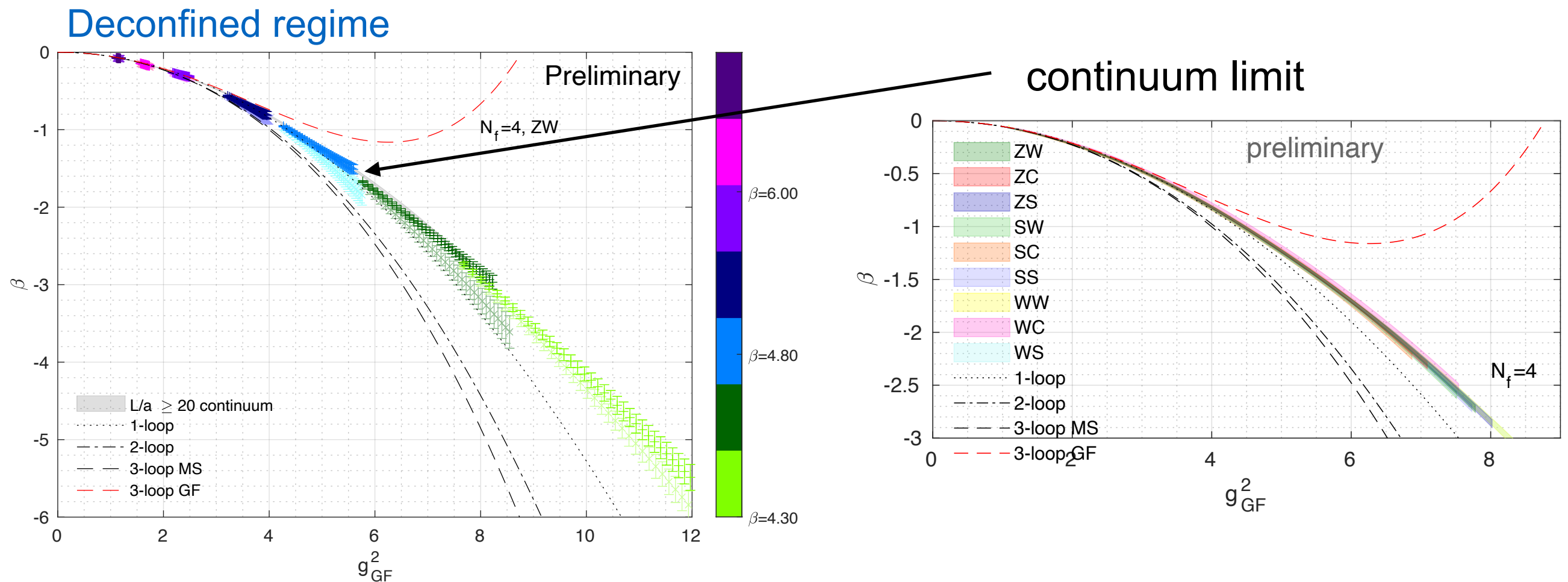
The continuous β function, $N_f = 4$ DWF, $m_f = 0$

AH, O. Witzel, in preparation

Simulations in the weak coupling **chirally symmetric** regime, 20^4 , 24^4 , 32^4 volumes

Raw data (Zeuthen flow, Wilson op)

- shows significant overlap and minimal finite volume effects
- ➔ sits almost on top of continuum extrapolation



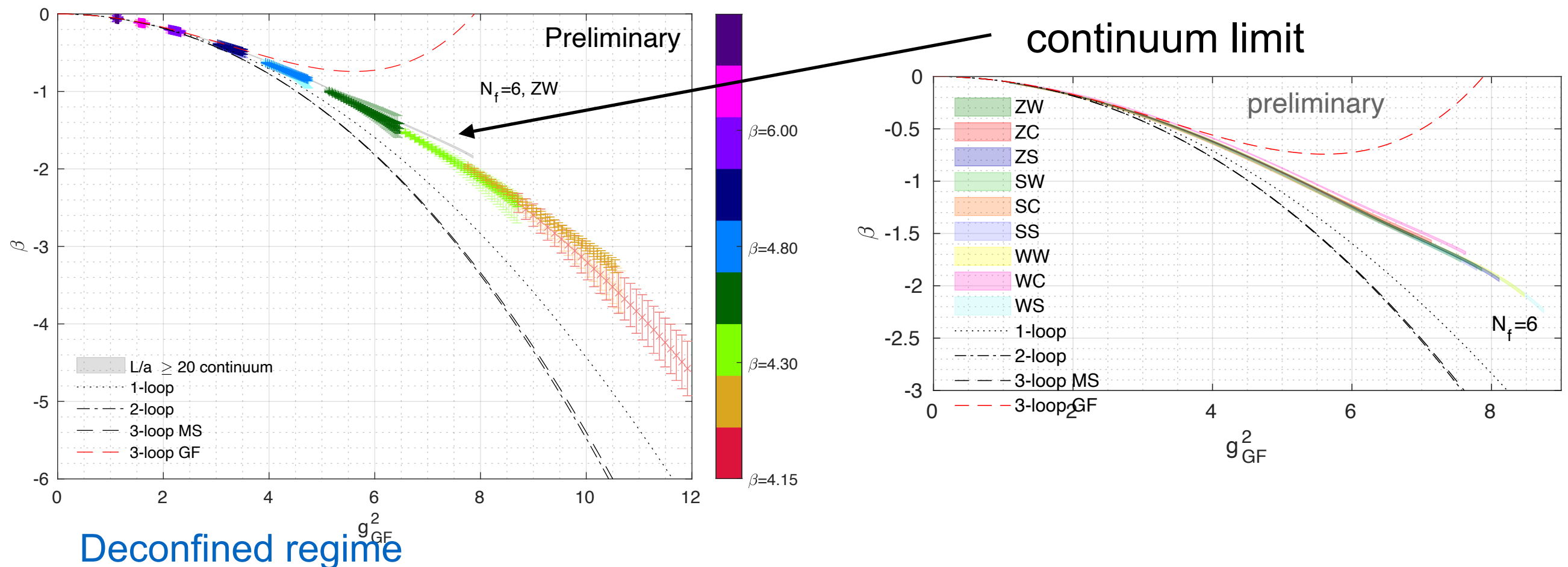
The continuous β function, $N_f = 6$ DWF, $m_f = 0$

AH, O. Witzel, in preparation

Simulations in the weak coupling **chirally symmetric** regime, $20^4, 24^4, 32^4, 40^4$ volumes

Raw data

- shows significant overlap and minimal finite volume effects
- ➔ sits almost on top of continuum extrapolation

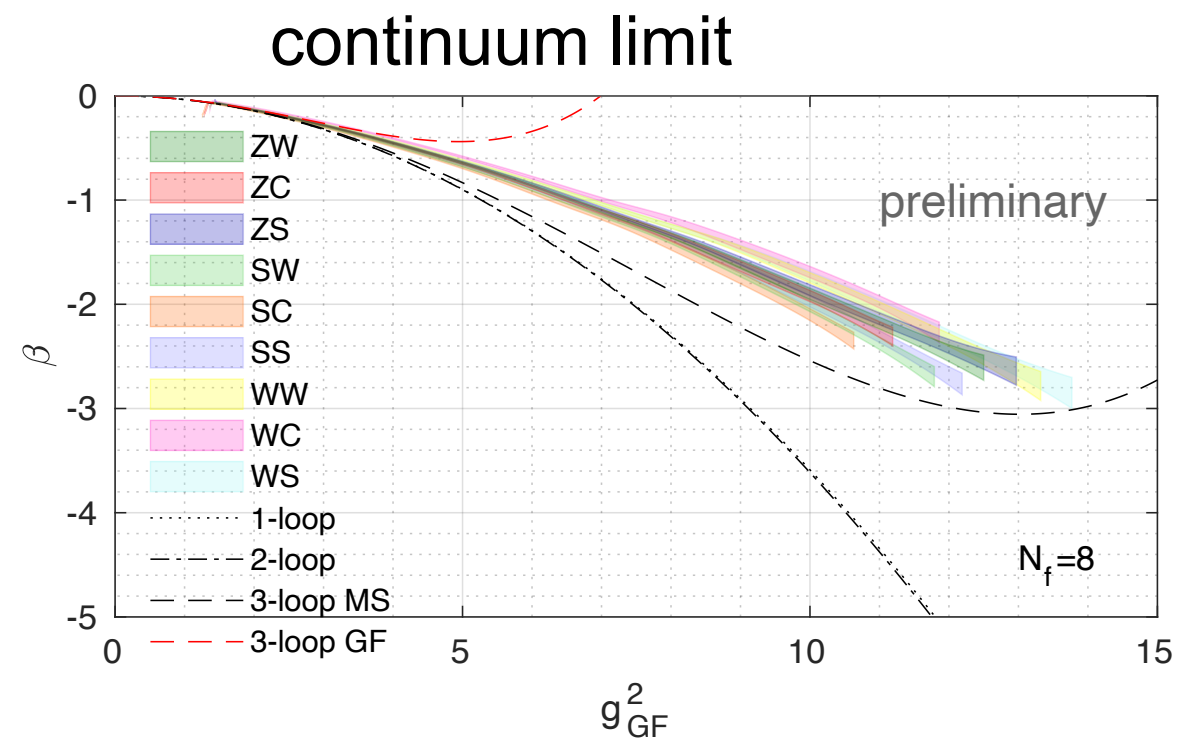
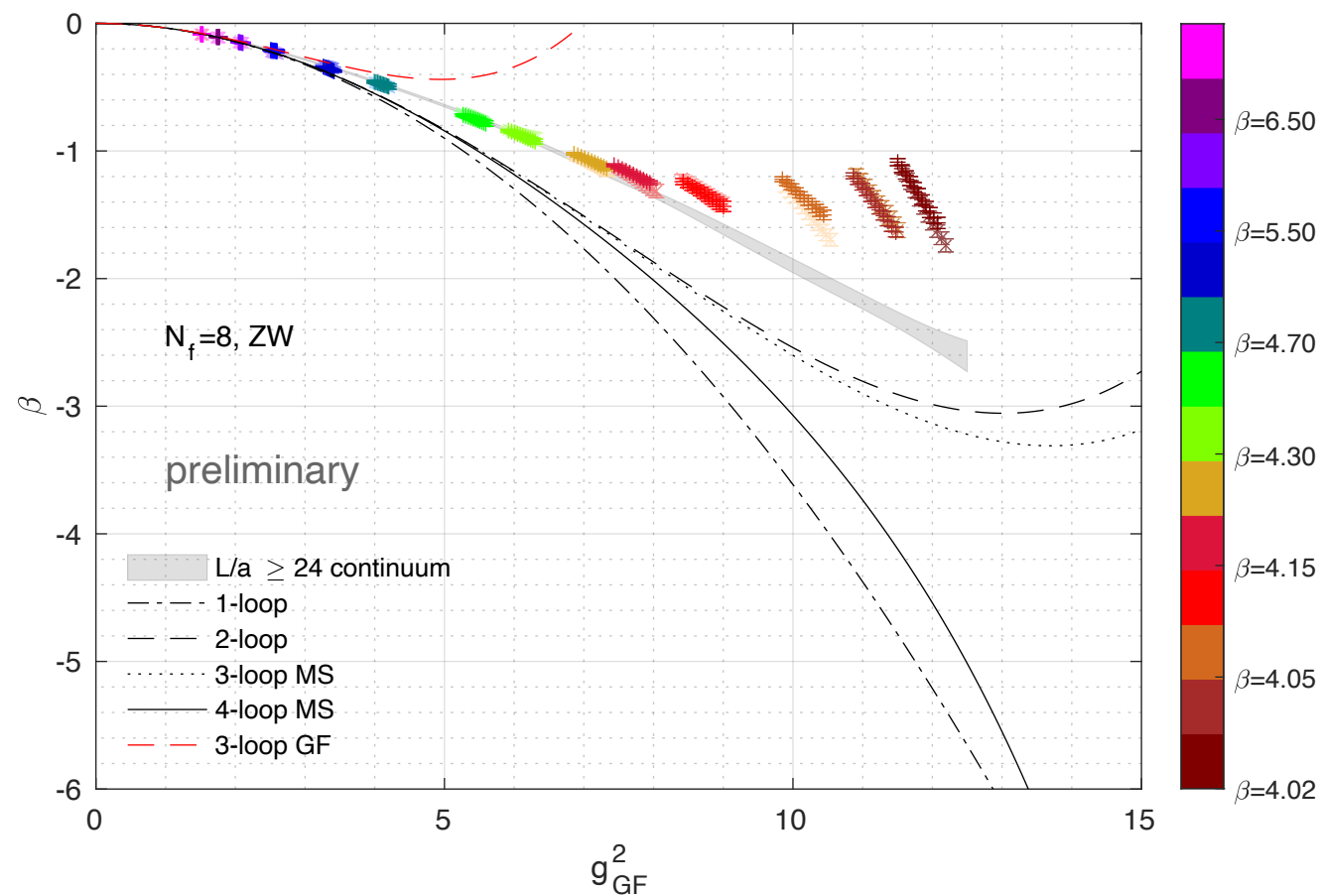


The continuous β function, $N_f = 8$ DWF, $m_f = 0$

AH, O. Witzel, in preparation

Simulations in the weak coupling chirally symmetric regime, $16^4, 24^4, 32^4, 48^4$ volumes
Raw data

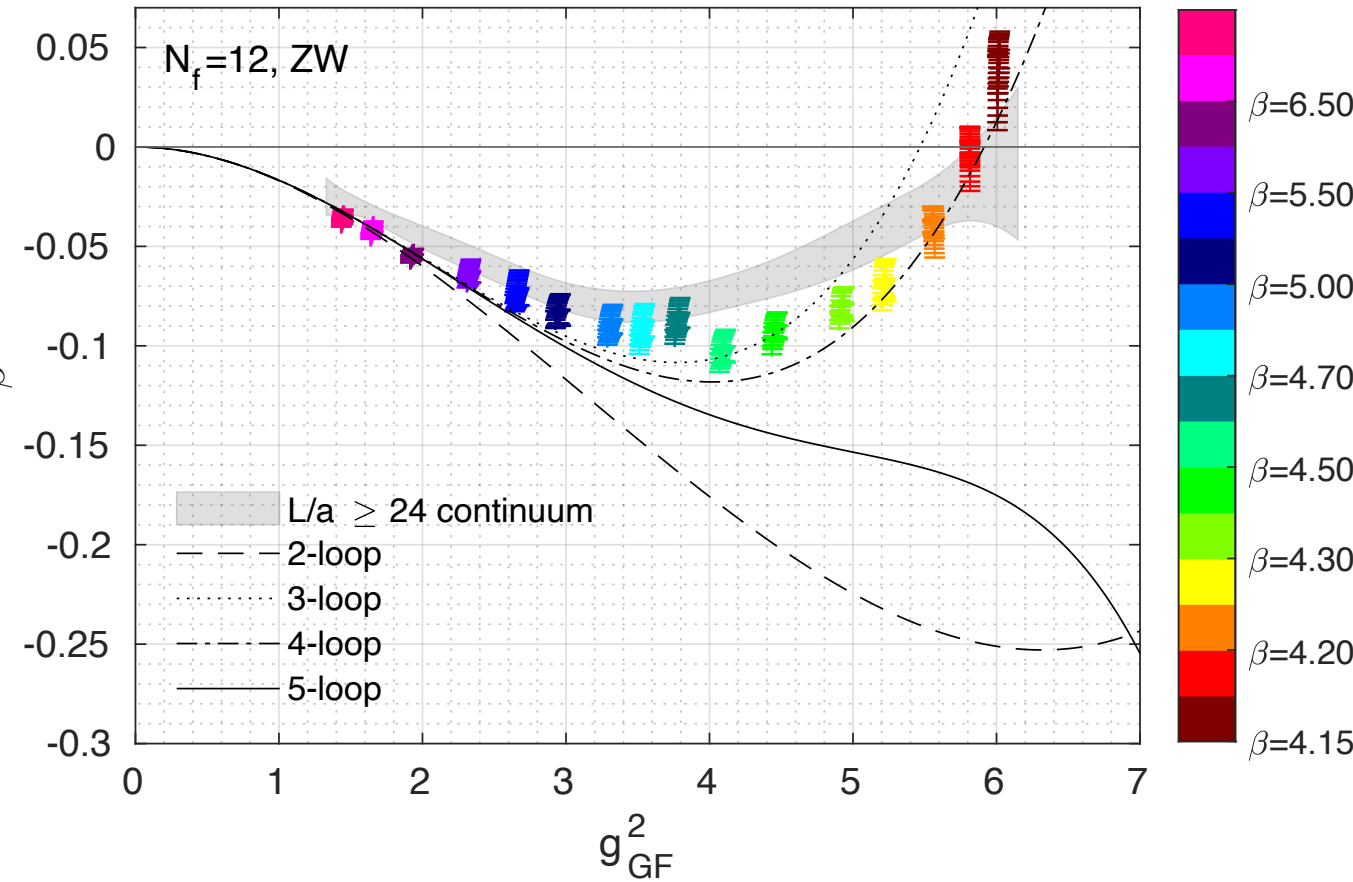
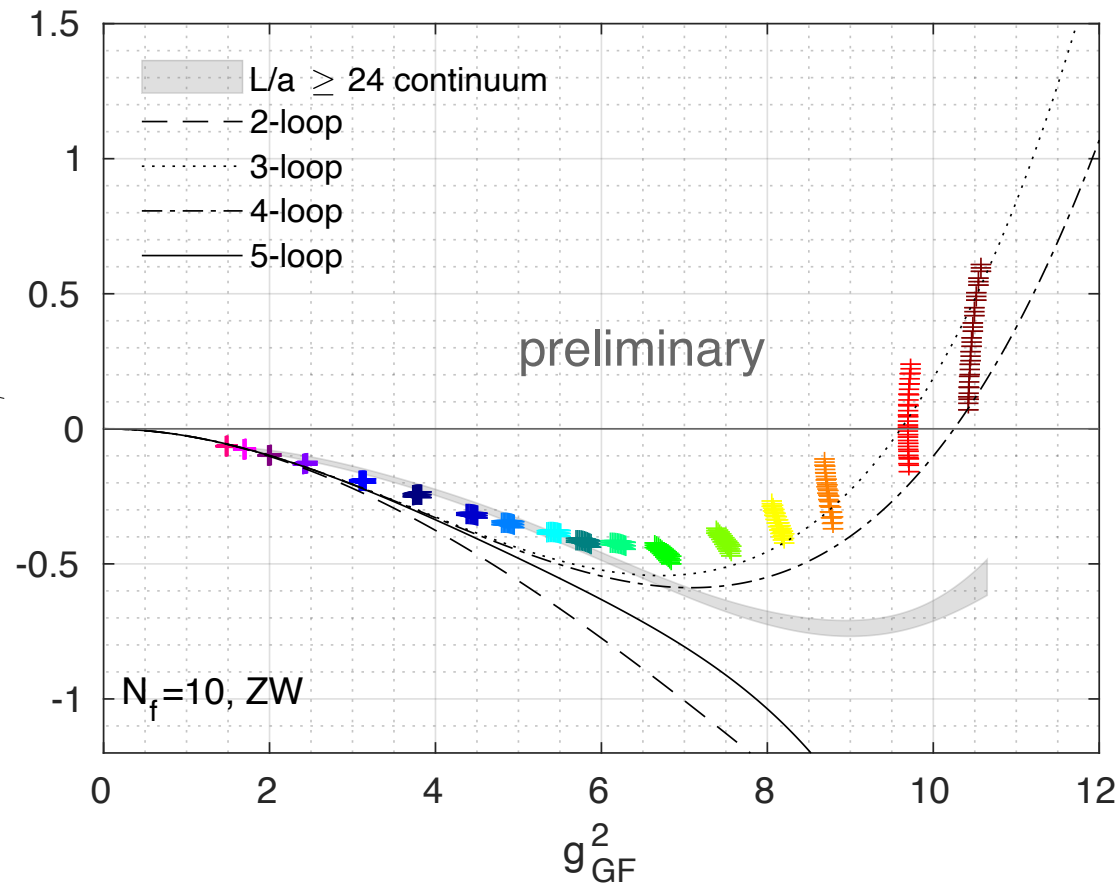
- shows significant overlap and minimal finite volume effects
- ➔ sits almost on top of continuum extrapolation up to $g_{GF}^2 \approx 8$
- at stronger coupling the perturbative guidance breaks down, no overlap, no unique curve traced out by raw data



Is this a signal of nearby IRFP/UVFP ? (Somewhere at $N_f \gtrsim 8$)

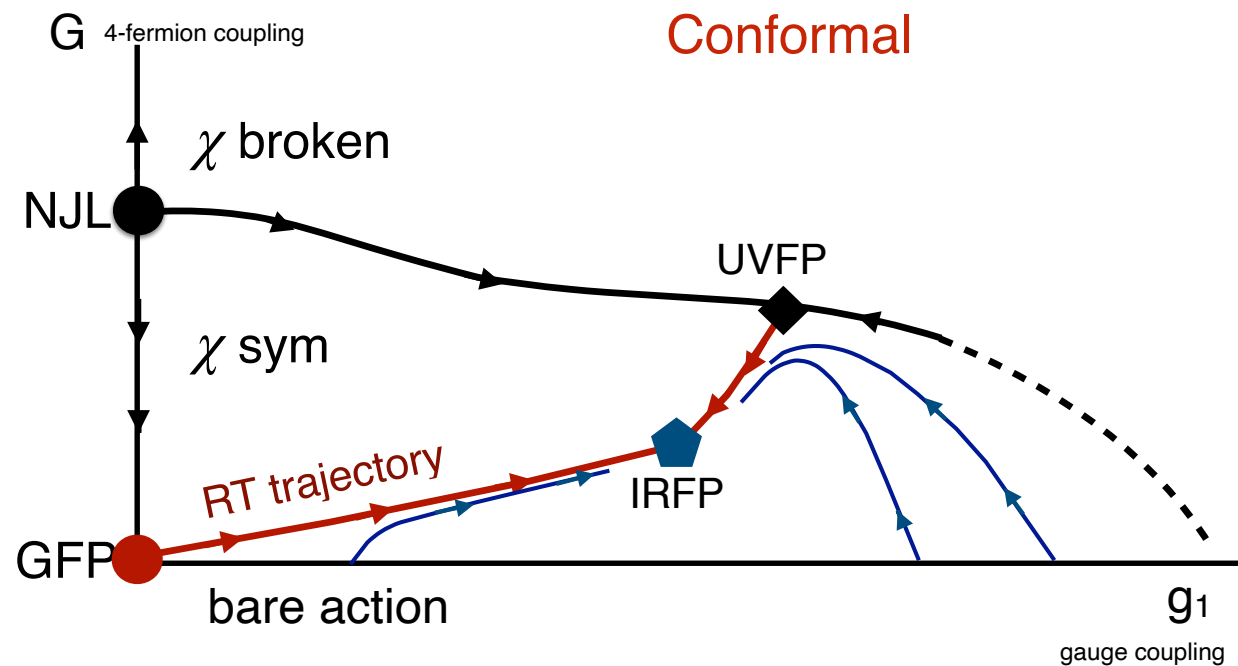
The continuous β function, $N_f = 10$ and 12 DWF

AH,O.Witzel, *Phys.Rev.D* 101 (2020) 11, 114508



- Very slow running
- At stronger couplings there is no overlap, no unique curve traced out by raw data
- Is this a signal of nearby IRFP/UVFP ?

The β function of (near-)conformal systems



Any calculation based on the perturbative Gaussian FP will cover only the regime between GFP and IRFP:

- ▶ $N_f = 0 - 6$ shows slow running but cutoff effects are easy to control
- ▶ $N_f \geq 8$ shows increasing non-perturbative effects
 - perturbatively justified extrapolation in $(a/L)^2$ or (a^2/t) might need to be replaced with $(a/L)^\alpha$, $(a^2/t)^{\alpha/2}$
 - new relevant operator at the UVFP means that a 1-parameter β function might not describe the RG flows

To solve these issues will require a new approach

Most simulations are plagued by (bulk) 1st order transition that prevent investigations at strong coupling

Improved action with Pauli-Villars bosons

AH, T. DeGrand PRD 49 (1994) 466

AH, B. Svetitsky, Y. Shamir, *Phys.Rev.D* 104 (2021) 7

Simulations in the strong coupling are difficult. Many flavor studies are particularly hard due to the induced **effective gauge action** from fermions

$$e^{-S_{ind}[U]} = \int \prod_{i=1}^{N_f} d\psi_i d\bar{\psi}_i e \left[-\sum_{i=1}^{N_f} \bar{\psi}_i (D+m_f) \psi_i \right] = \text{Det}(D + m_f)^{N_f} = e^{N_f \text{Tr} \log(D+m_f)}$$

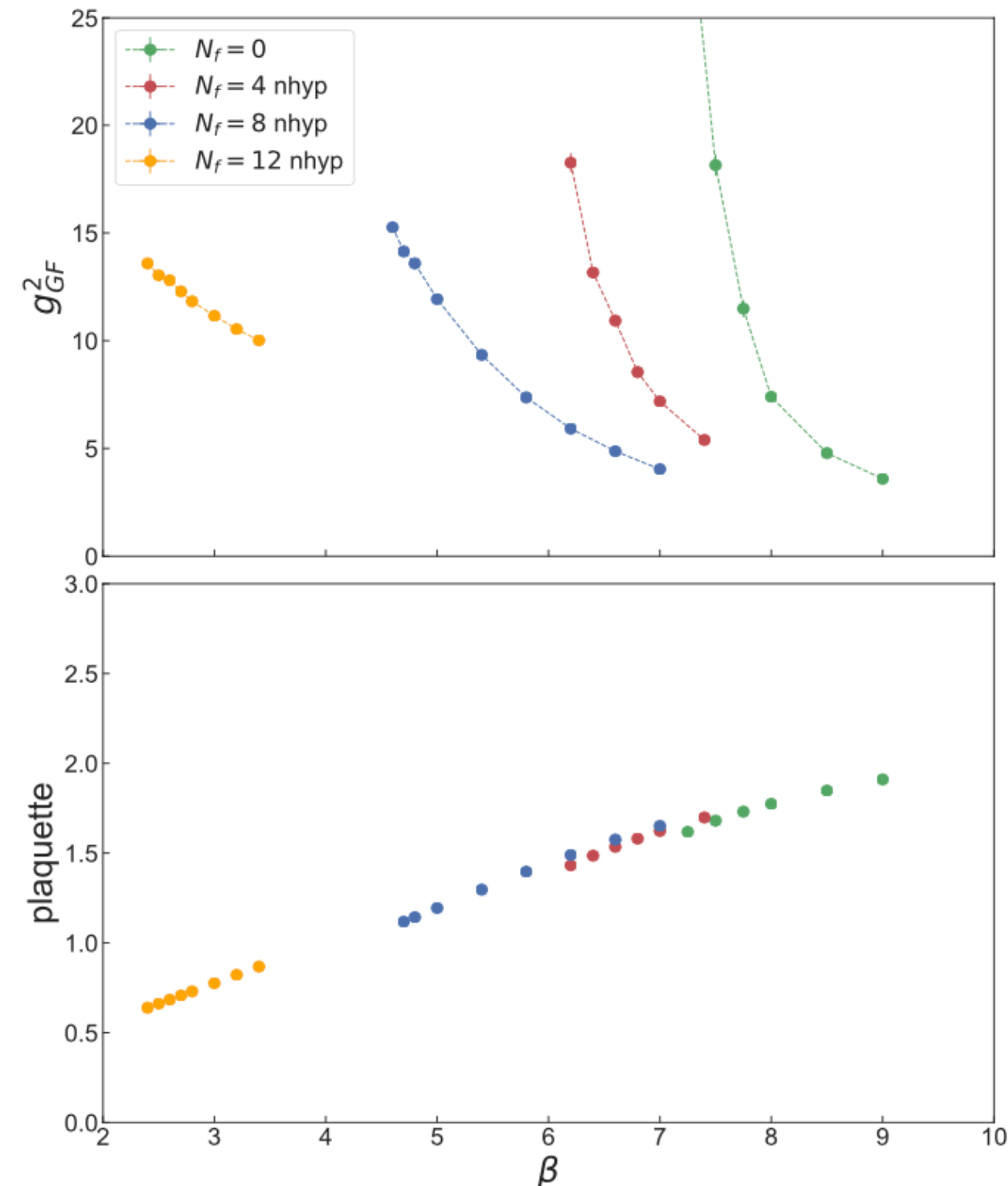
Expanding S_{ind} in $1/m_f$ gives a sum of gauge loops. For staggered fermions

$$S_{ind} = -N_f/4 \sum_{\ell} \frac{(-1)^{\ell/2}}{\ell (2am_f)^\ell} \sum_x \sum_{\mathcal{C}_\ell} \mathcal{E}_{\mathcal{C}_\ell} \text{Tr} U_{\mathcal{C}_\ell}$$

Leading term is plaquette: $\beta_{ind}^{(p)} = \frac{6N_f}{4(2am_f)^4} : \beta_{eff} = \beta_g + \beta_{ind}$

The induced action of fermions

Compare simulations with $N_f = 0, 4, 8,$ and 12 fundamental (staggered) fermions



As N_f grows,

- identical renormalized coupling corresponds to smaller the bare coupling
- the plaquette much rougher

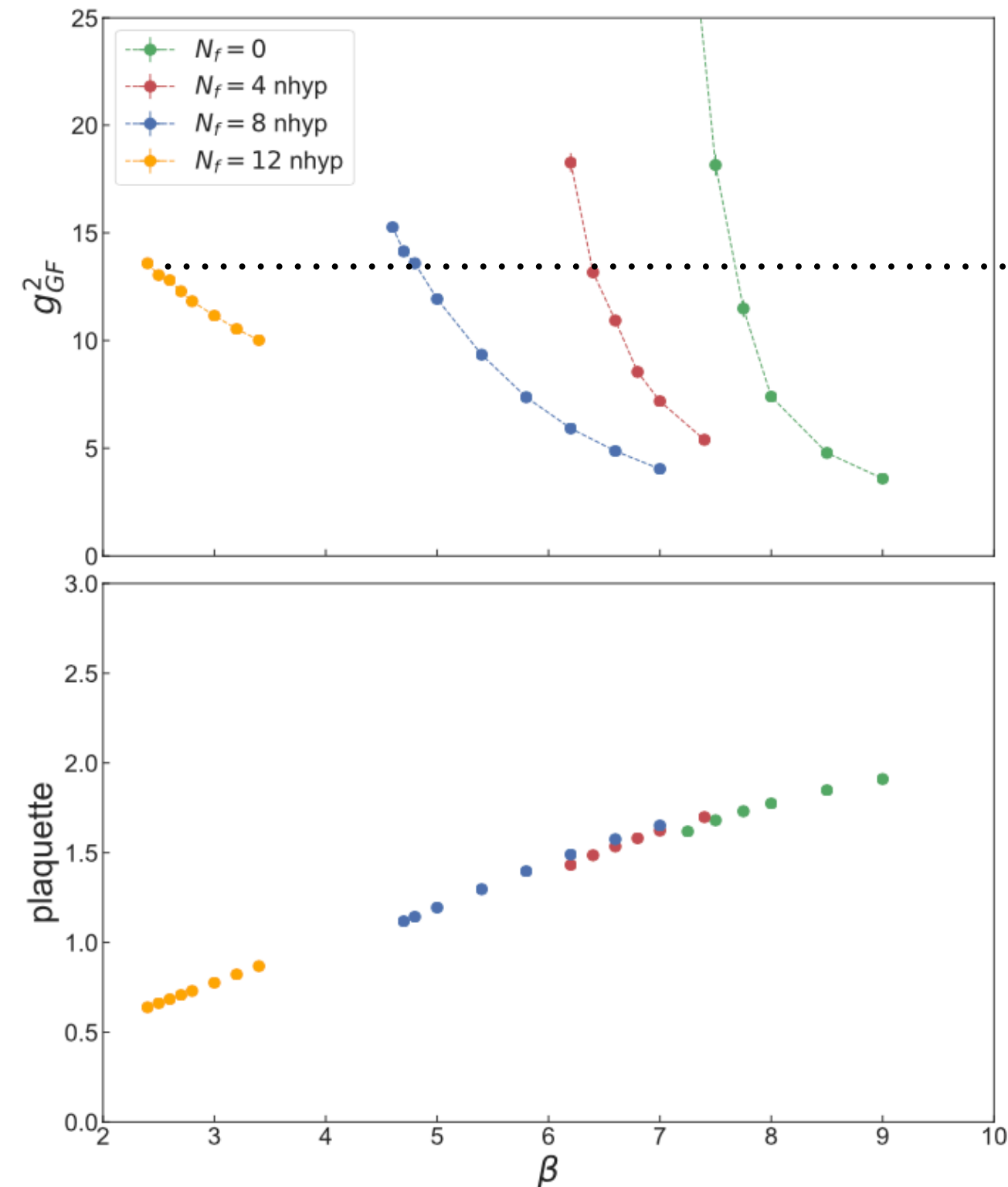
Simulations to be done at strong bare coupling

- lattice artifacts are large
- spurious bulk transitions limit the renormalized coupling range

Both $N_f = 8$ and 12 exhibit 1st order bulk transition that limits the accessible coupling range

The induced action of fermions

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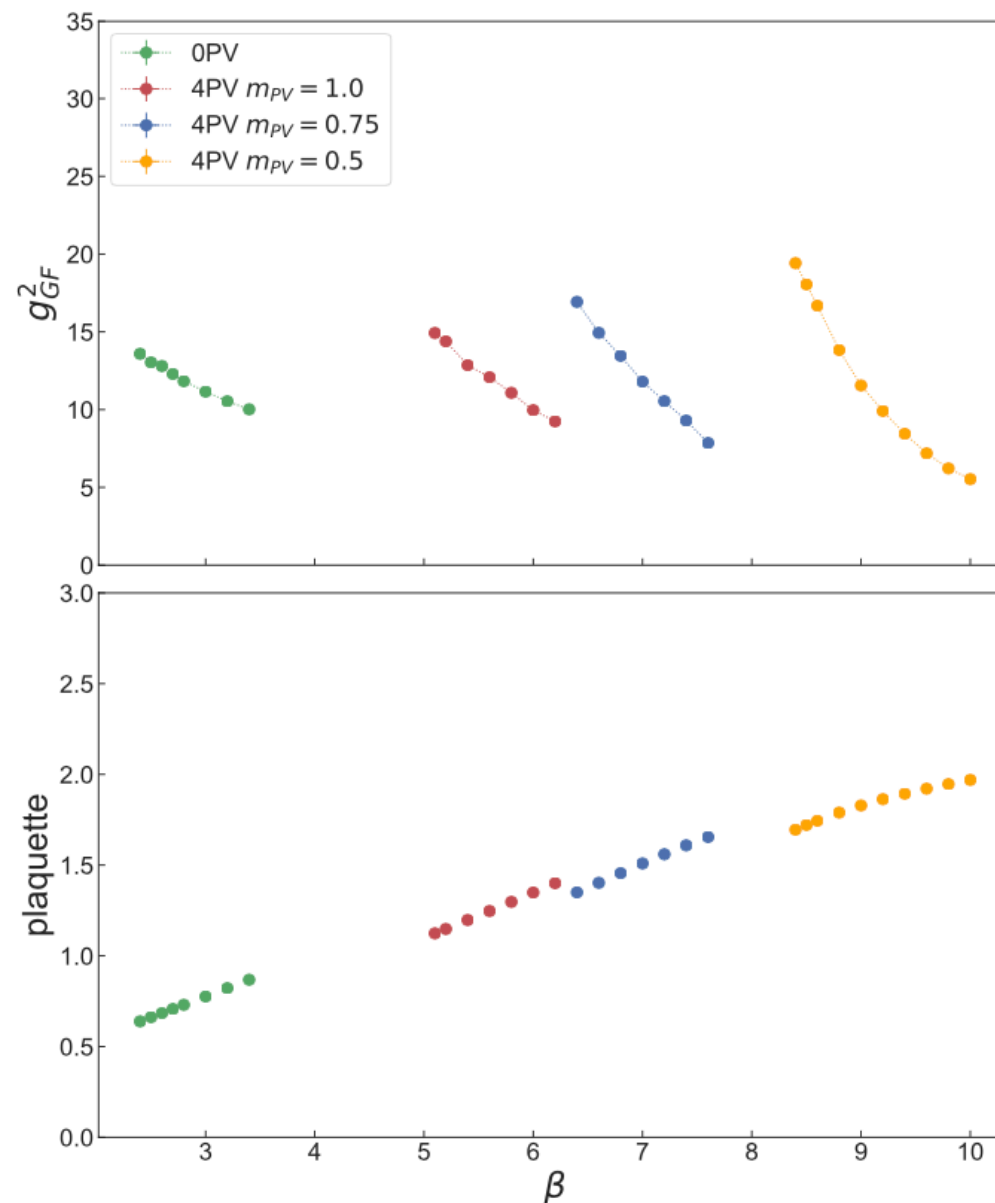
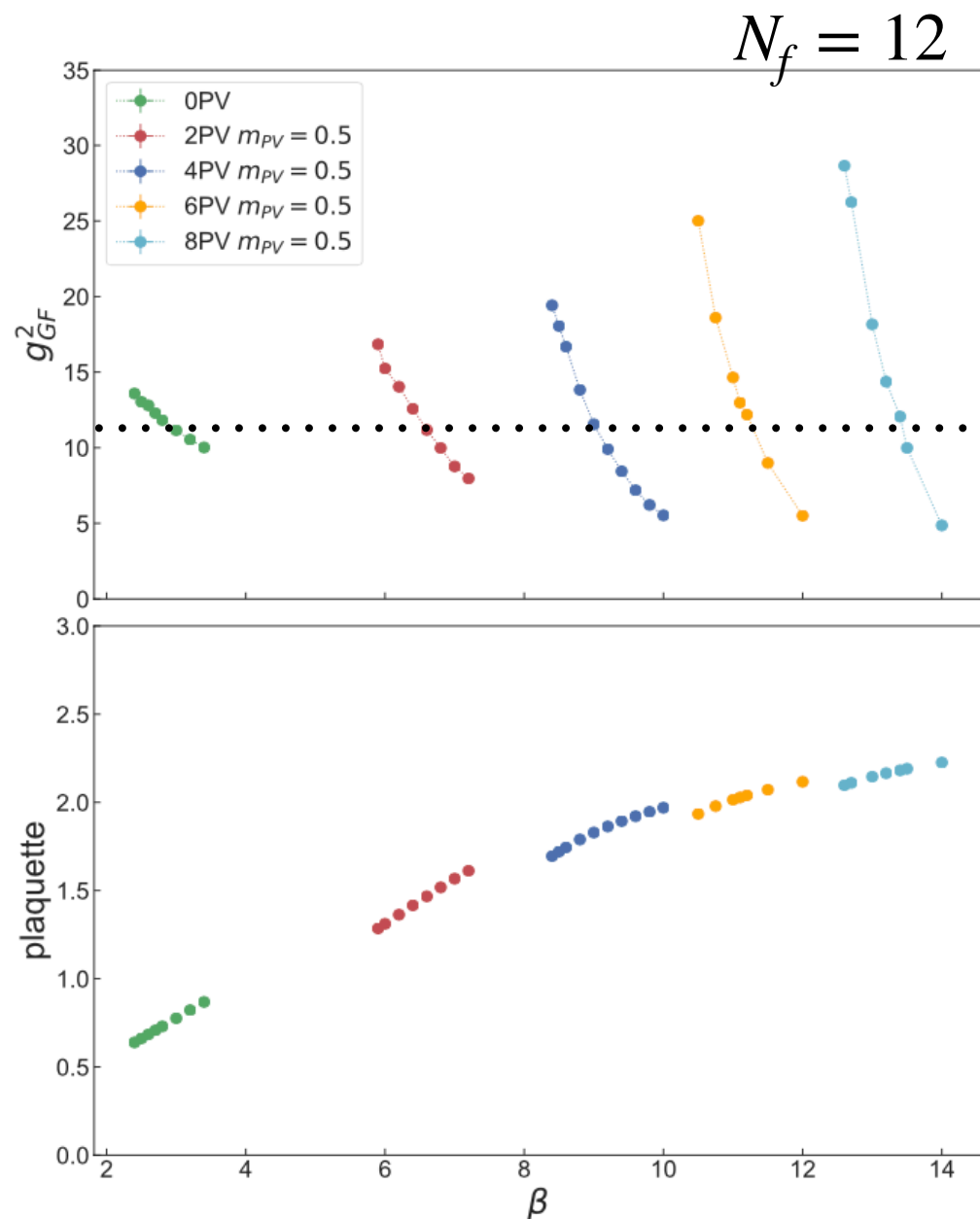
Both $N_f = 8$ and 12 exhibit 1st order bulk transition that limits the accessible coupling range

Improved action with Pauli-Villars bosons

Compensate the effect of fermions by adding heavy Pauli-Villars bosons:

- If $m_{PV} \gg m_f, 1/L, \Lambda_{QCD}$ the PV fields decouple and do not change the IR dynamics
- PV bosons induce a negative effective action
- any number of PV bosons can be added (even with different Dirac operator)

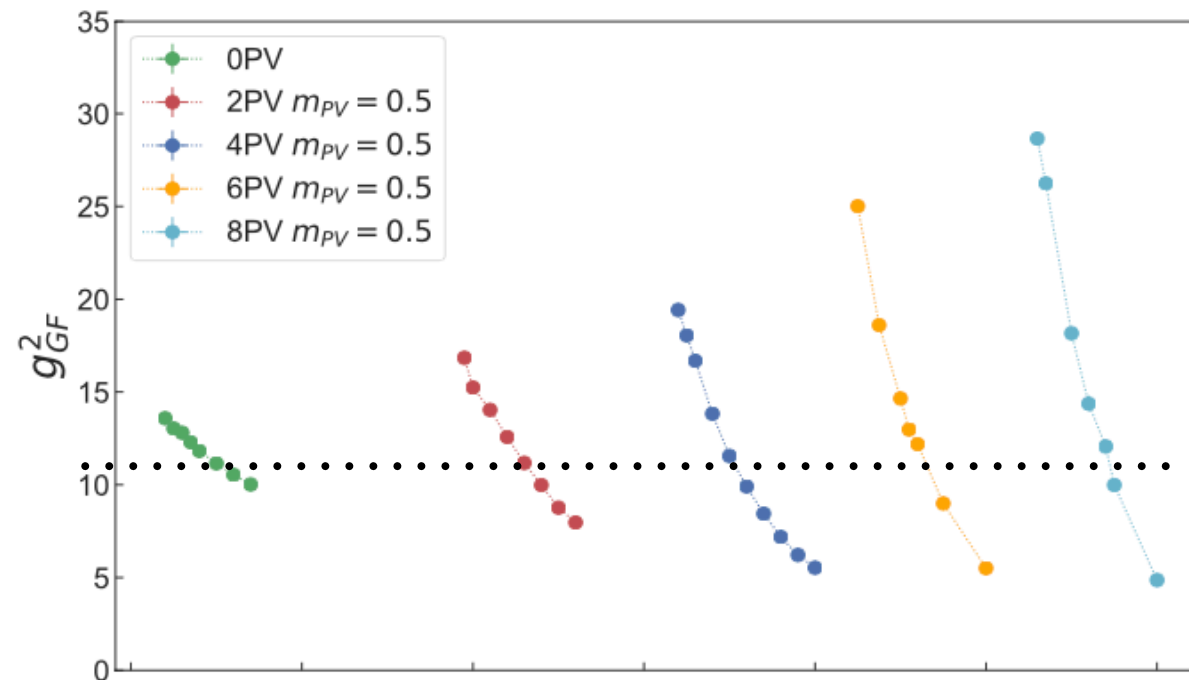
- this is already done in DW simulations



Improved action with Pauli-Villars bosons

Including PV bosons are computationally inexpensive

They likely speed up the simulations while remove cutoff effects and open up the strong coupling parameter space



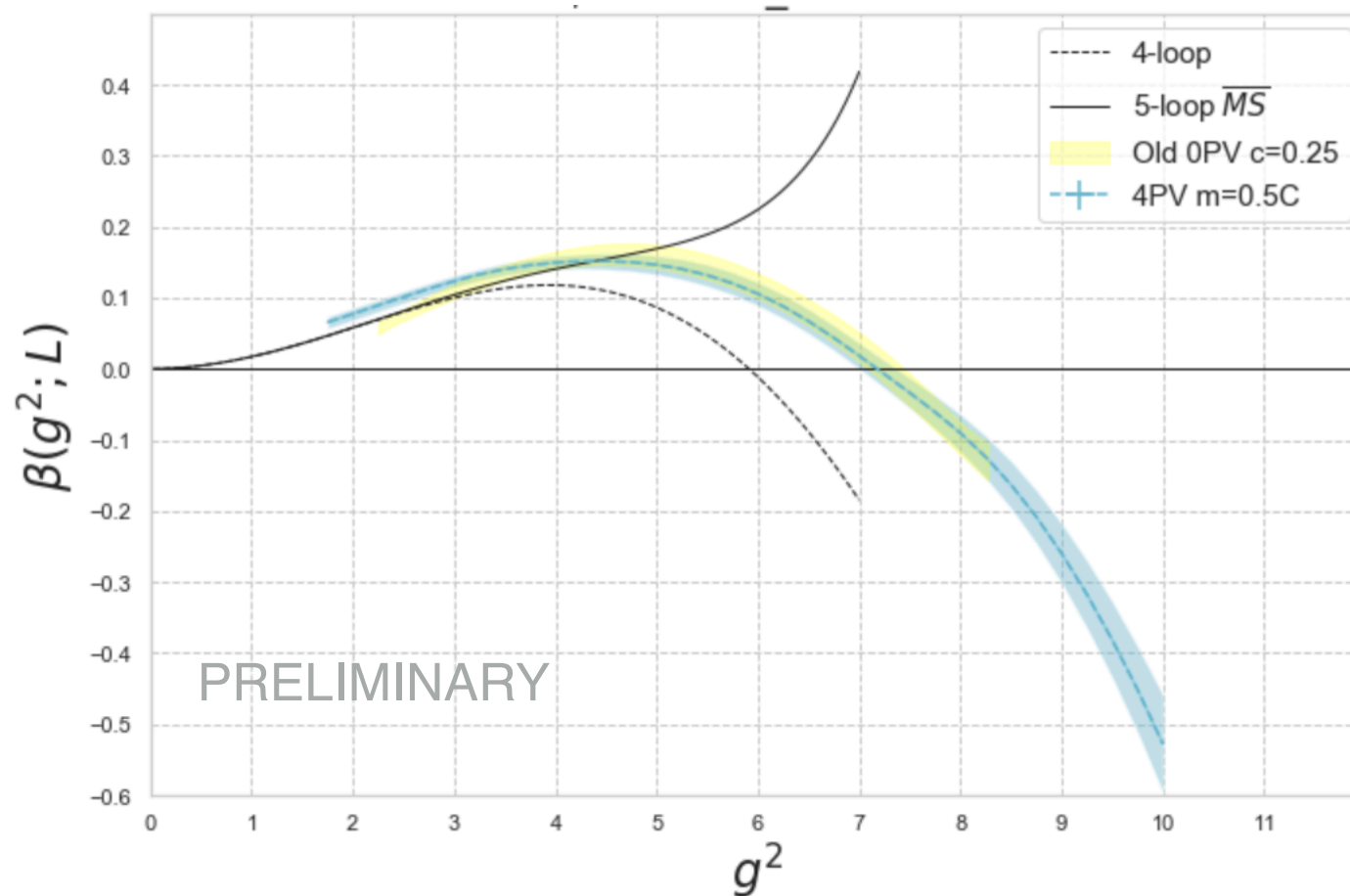
$$N_f = 12$$

N_{PV}	β	g_{GF}^2	N_{step}	N_{CG}	$ \delta H $
0	2.8	10.87(6)	15	1190	0.45
2	6.8	10.89(9)	12	400	0.062
4	9.0	10.58(8)	12	365	0.033
6	11.2	10.87(30)	12	353	0.033
8	13.4	10.87(100)	12	359	0.034

$N_f = 12$ with/wo Pauli-Villars bosons

Predicted step scaling beta functions are identical (staggered)

PV action can explore much stronger couplings



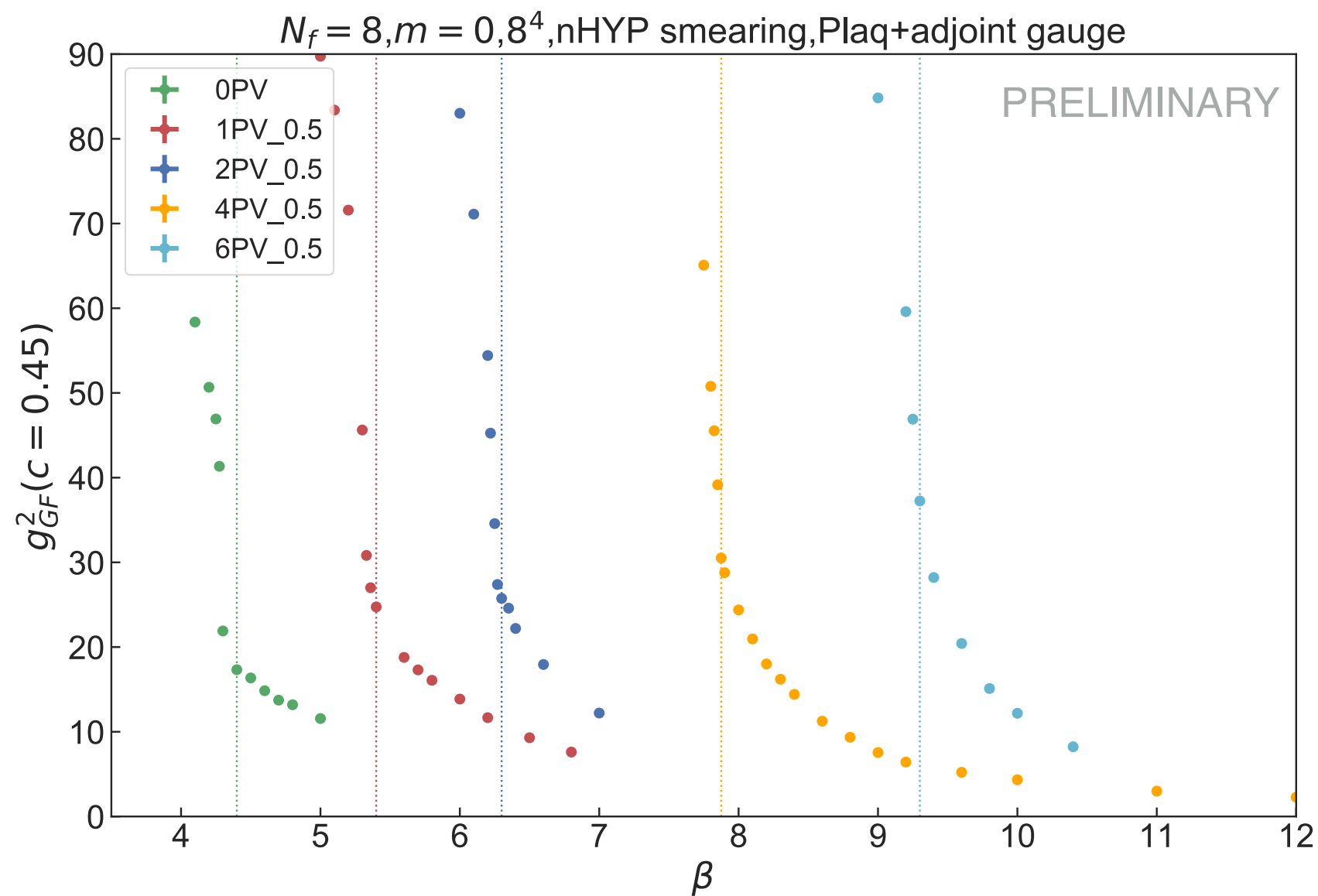
Old:

AH, D. Schaich JHEP 02 (2018) 132

$N_f = 8$ with/wo Pauli-Villars bosons

Expected to be very close to the conformal window

Simulations without PV bosons are limited by a 1st order bulk transition to an S4 phase where staggered shift symmetries are broken



S4 phase remains as PV bosons are added but the transition becomes smooth, appears continuous

Could $N_f = 8$ be the boundary to the conformal window?

$N_f = 8$ phase structure with 4PV, $am_{PV} = 0.4$

Running coupling g_{GF}^2 at flow time

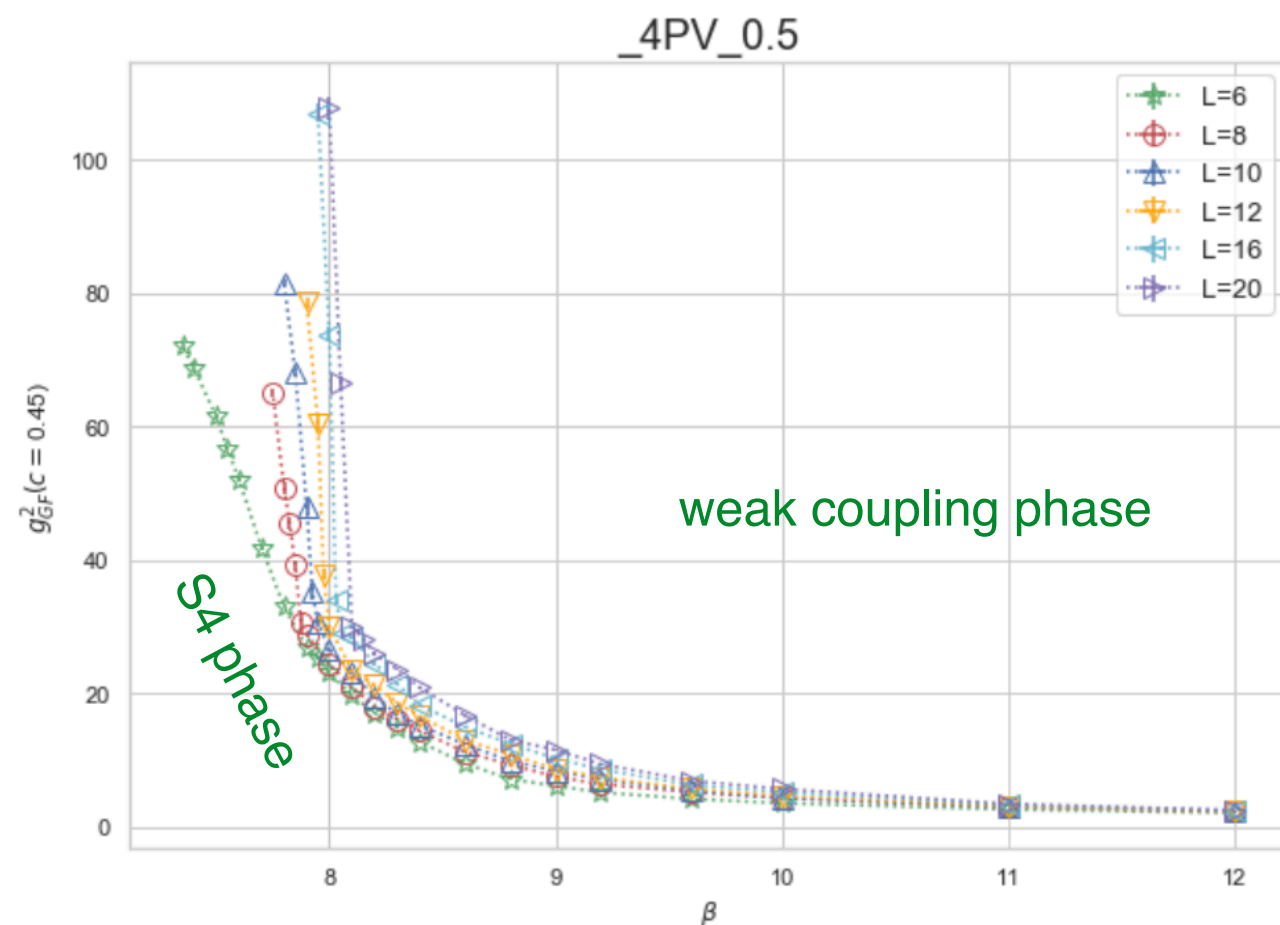
$$\sqrt{8t} = 0.45L, L = 6, \dots, 20$$

Significant finite volume dependence

—> finite size scaling

Topological susceptibility is rising even though simulations are at $m_f = 0$

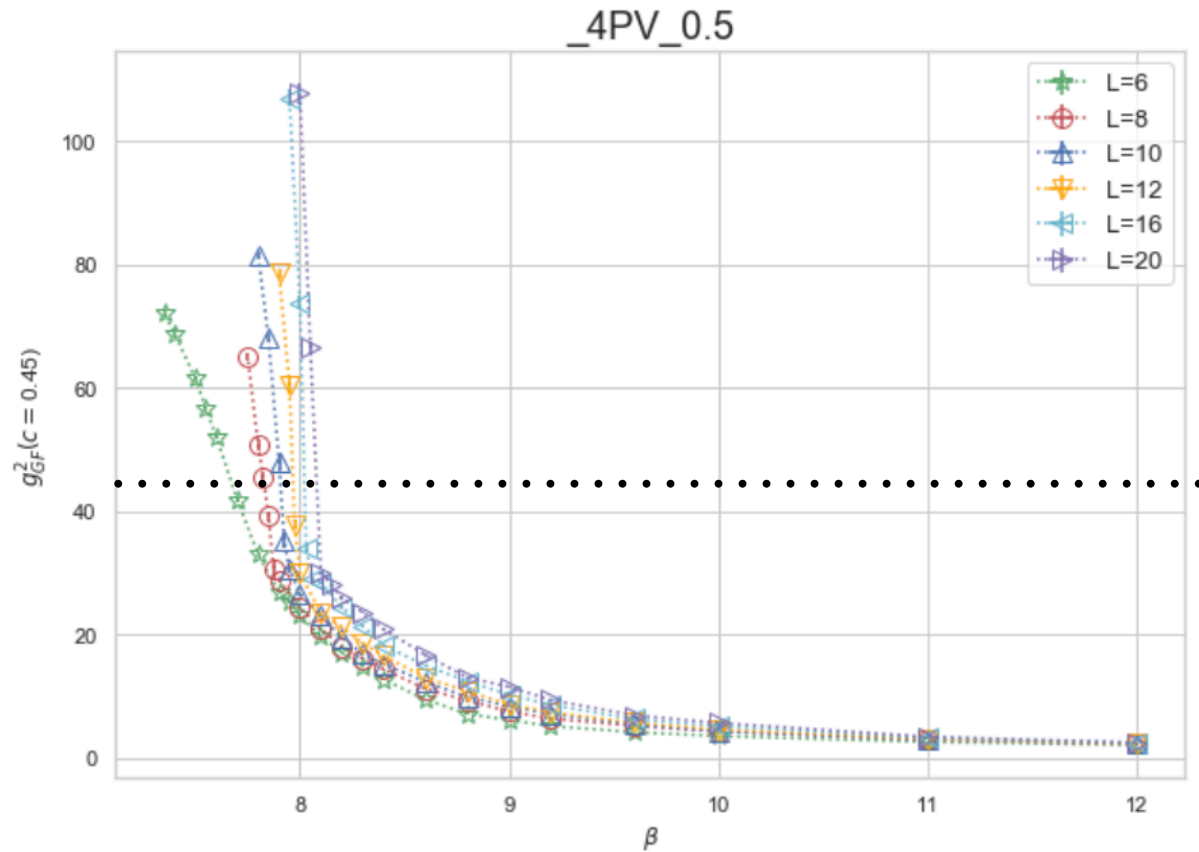
Staggered fermions are Dirac-Kaehler, equivalent to Dirac only at $g_0^2 = 0$



$N_f = 8$: finite size scaling

PRELIMINARY

Running coupling at flow time $\sqrt{8t} = 0.45L$



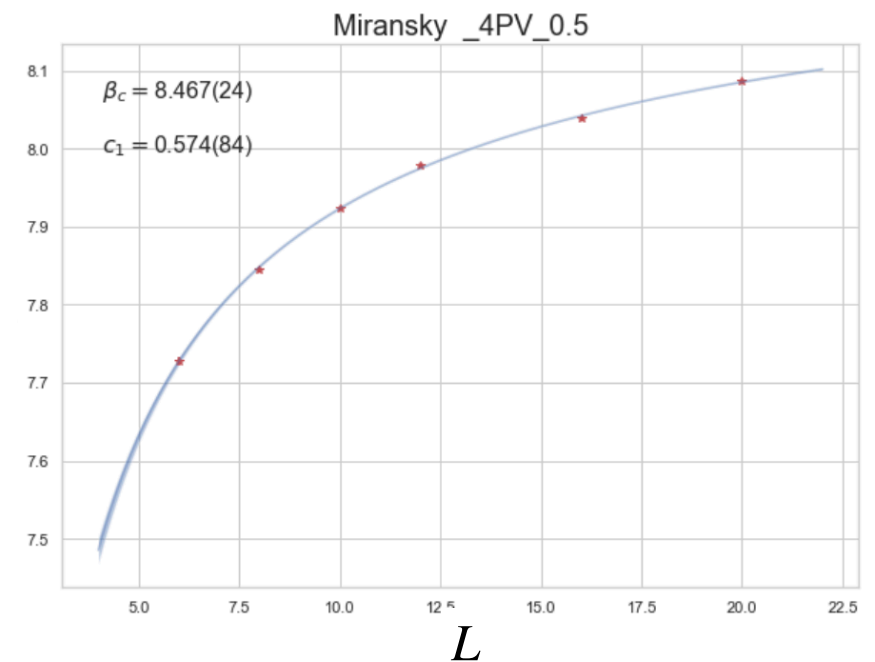
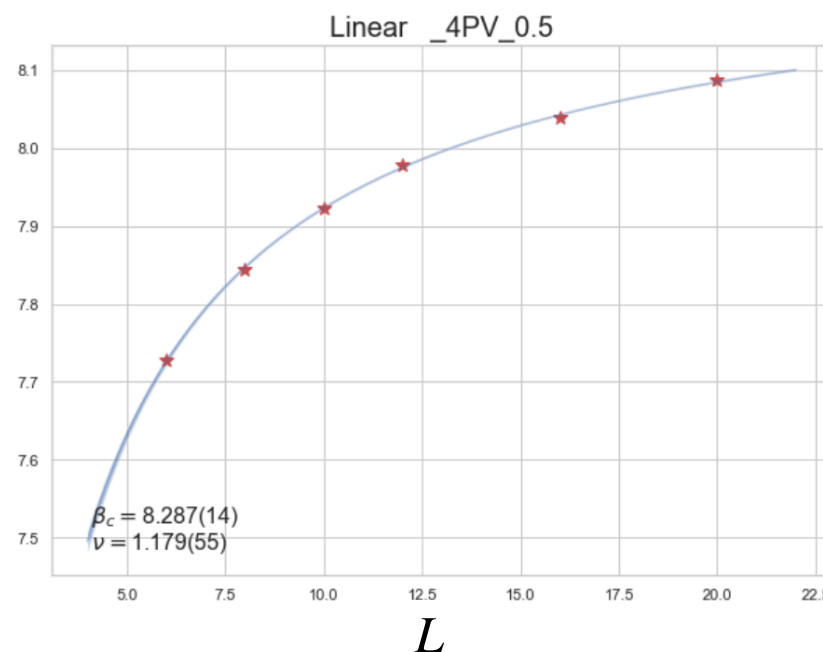
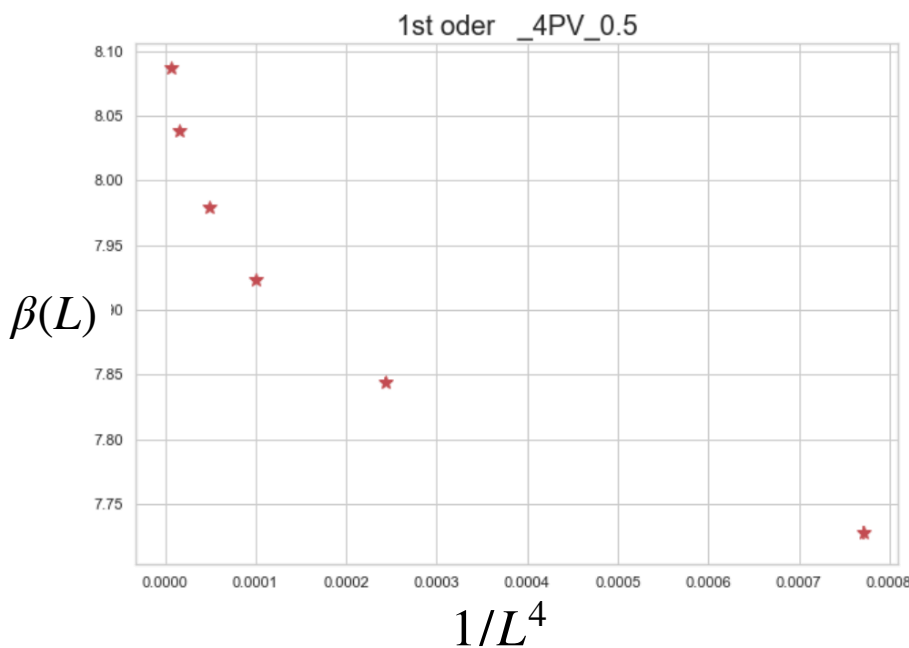
At fixed g_{GF}^2 the coupling scale as

1st order : $\beta(L) - \beta^* \propto L^{-1/4}$

2nd order : $\beta(L) - \beta^* \propto L^{-1/\nu}$

Essential singularity (XY model/
Miransky) : $\beta(L) - \beta^* \propto 1/\log^2(cL)$

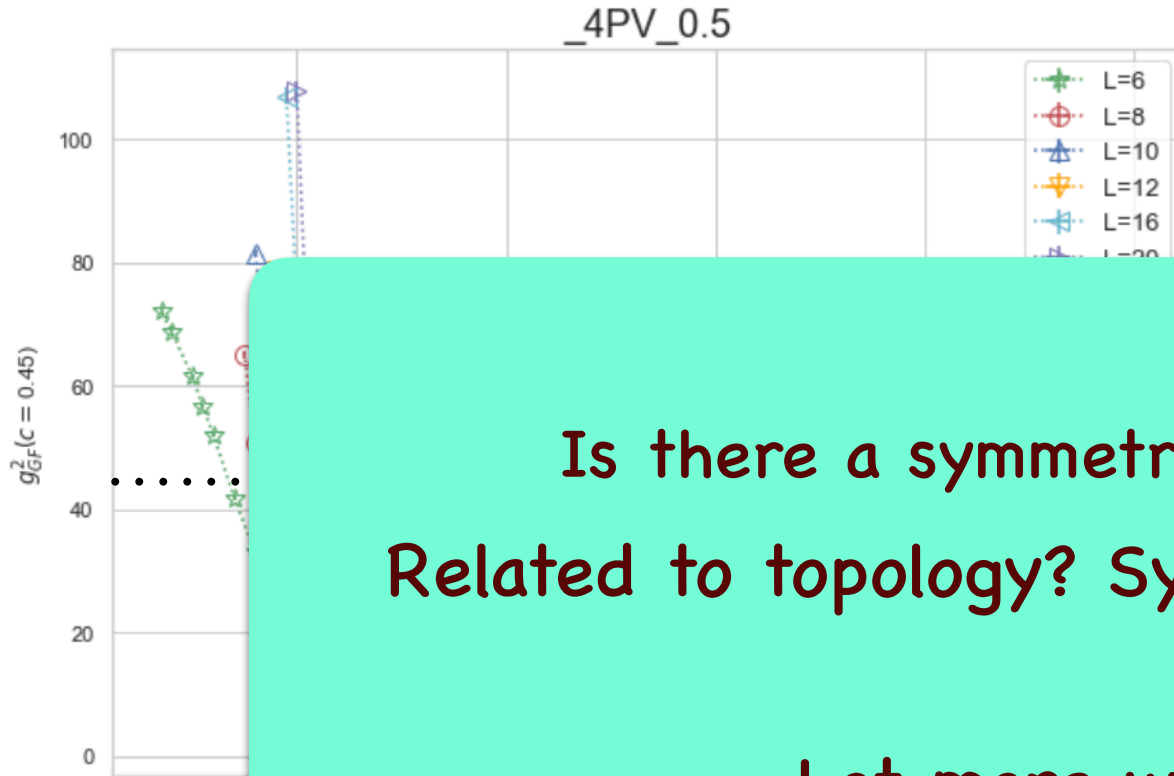
Data are not consistent with 1st order transition



$N_f = 8$: finite size scaling

PRELIMINARY

Running coupling at flow time $\sqrt{8t} = 0.45L$



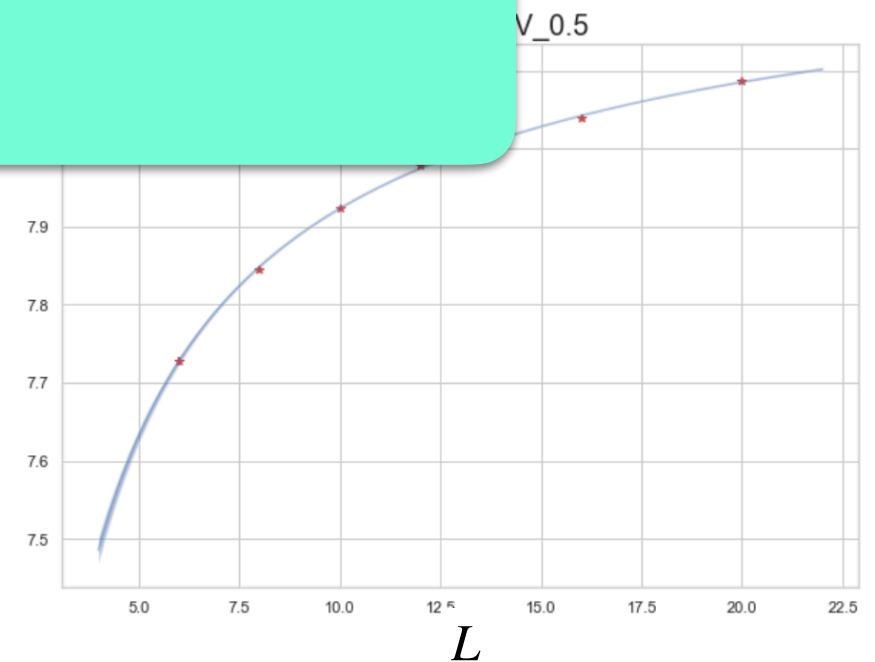
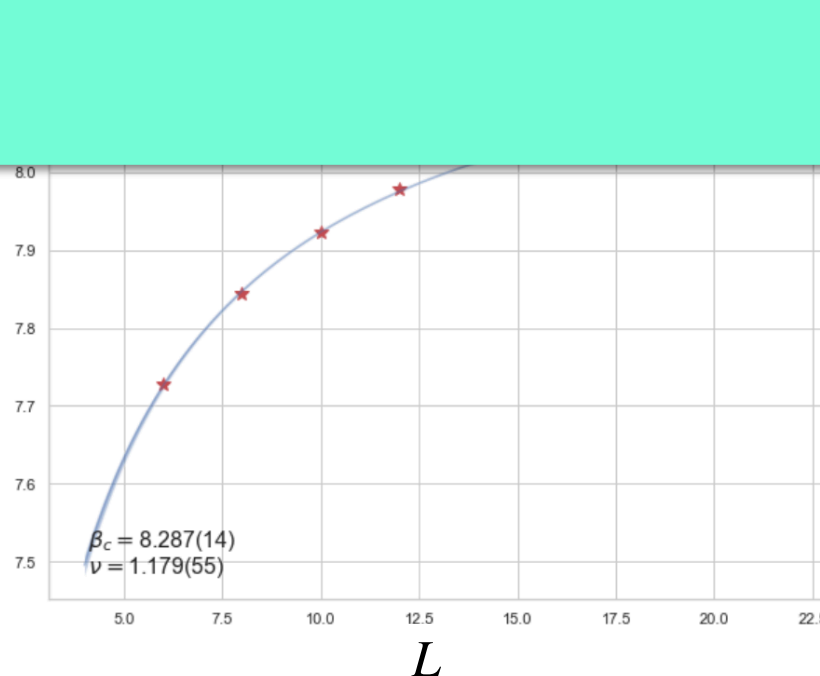
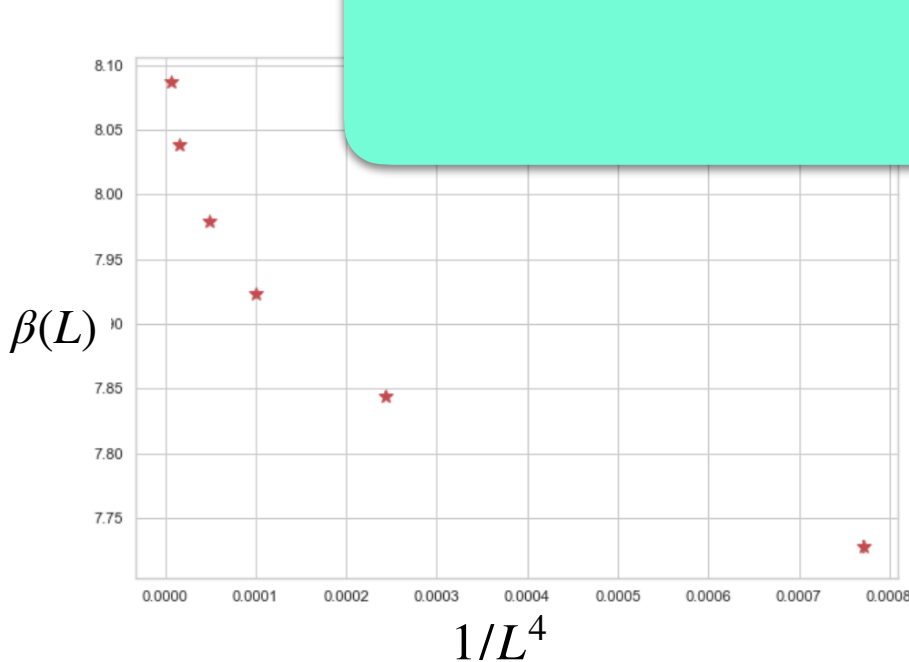
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Is there a symmetry that selects $N_f = 8$?
 Related to topology? Symmetric mass generation?

Lot more work is needed!



Summary and Outlook

The equivalence of Wilsonian RG and gradient flow allows a theoretically solid description of the strong coupling regime of lattice models

- ◆ This is particularly important in near-conformal / conformal systems where new fixed points, new relevant operators appear

The continuous β function:

- ◆ In QCD-like chirally broken systems it is well controlled with minimal cutoff effects with improved action/flow/operator in QCD-like chirally broken systems
- ◆ Near the conformal sill strong non-perturbative effects appear :
 - extended actions (4-fermion term?) might be necessary
 - larger volumes (brute-force approach) can extend the accessible range, but will not cover the IRFP or UVFP beyond

PV bosons reduce cutoff effects, open up the parameter space

- ◆ cleaner numerical results
- ◆ possibly revealing new dynamics
- ◆ is $N_f = 8$ special?

EXTRA SLIDES

The continuous β function, $N_f = 2$

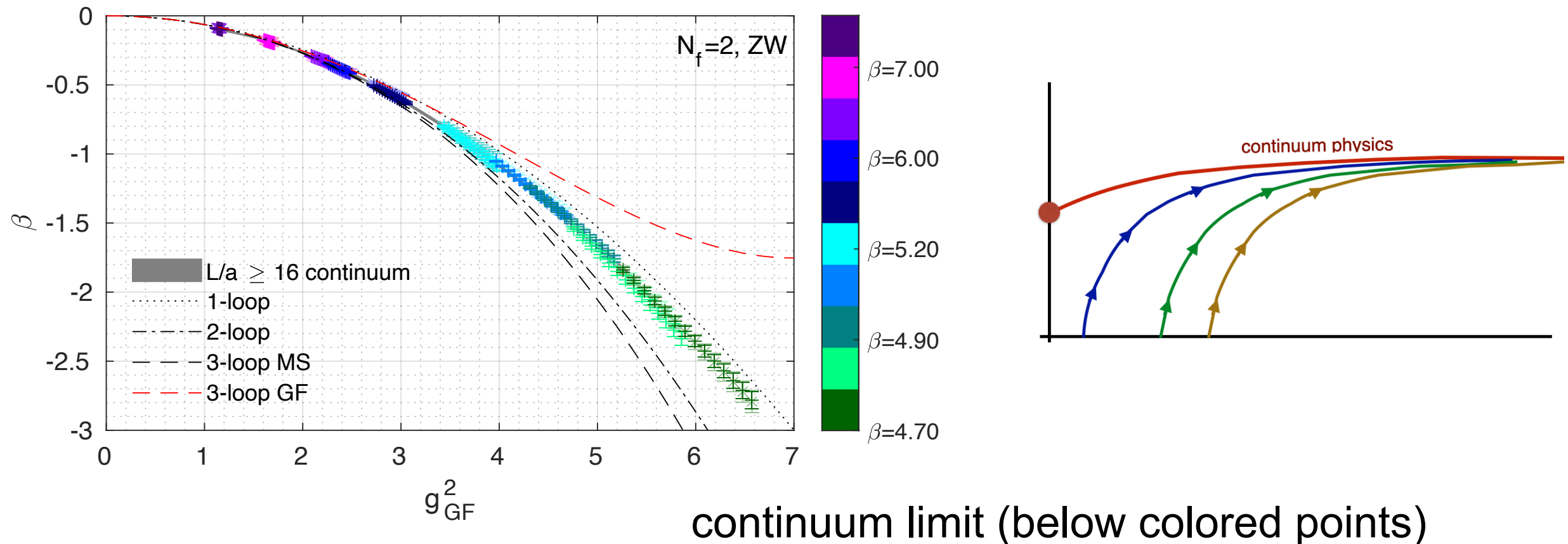
AH, O. Witzel, *Phys.Rev.D* 101 (2020) 3

We use

- Symanzik gauge action, Mobius domain wall fermions (Grid)
- Zeuthen (Z), Wilson(W) and Symanzik(S) flows : optimize to pull the RT close
- Wilson plaquette(W), clover(C) and Symanzik(S) operators : combine to optimize for the scaling operator

ZW raw data

- shows minimal finite volume effects and significant overlap

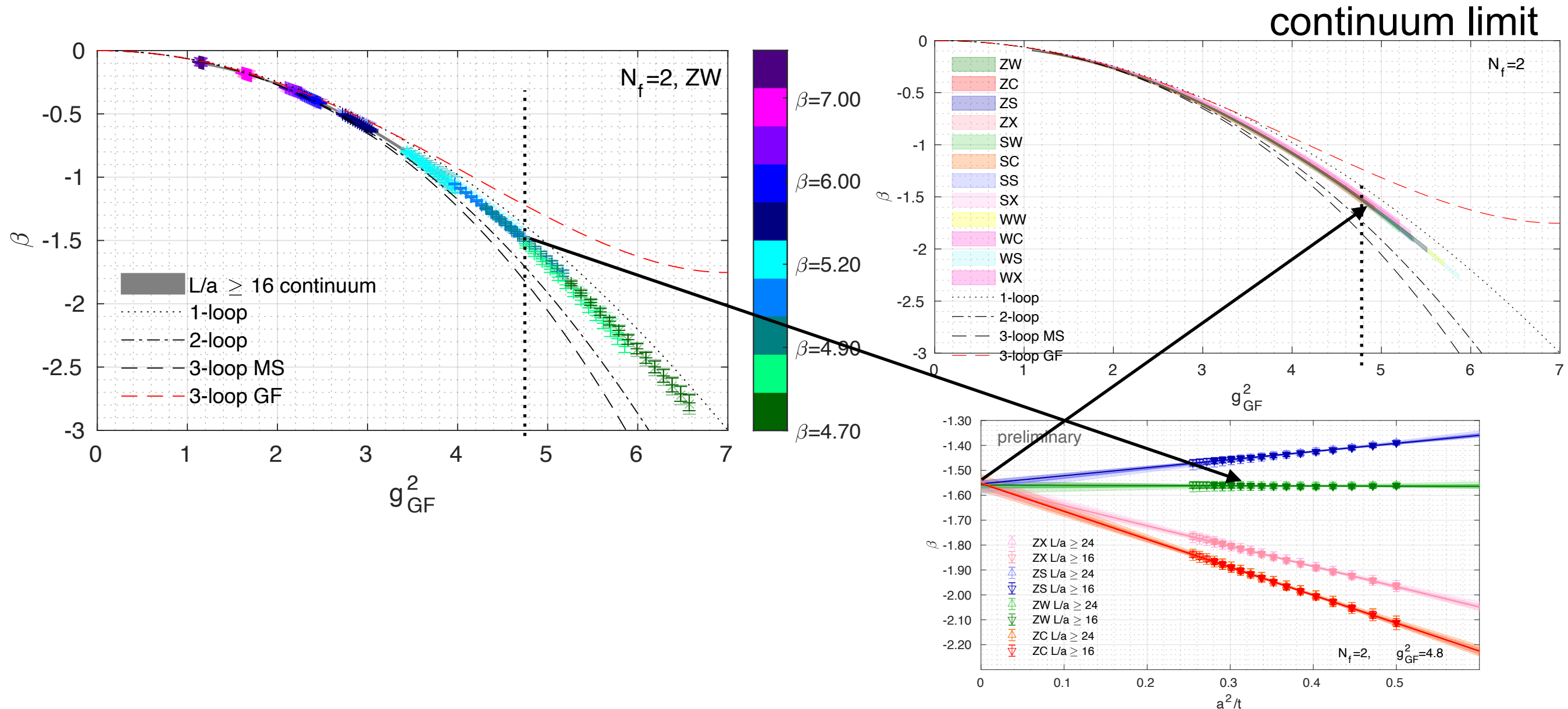


The continuous β function, $N_f = 2$

AH, O. Witzel, *Phys.Rev.D* 101 (2020) 3

Analysis steps for continuum limit:

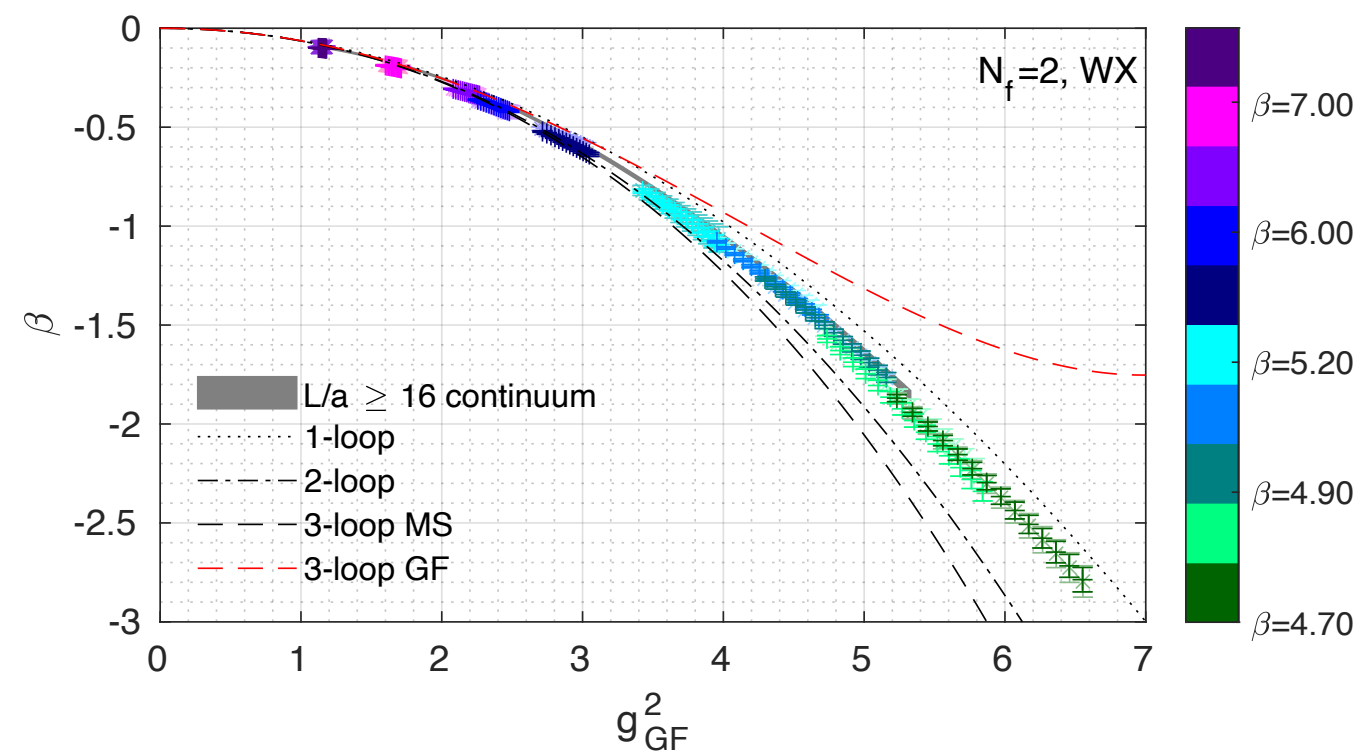
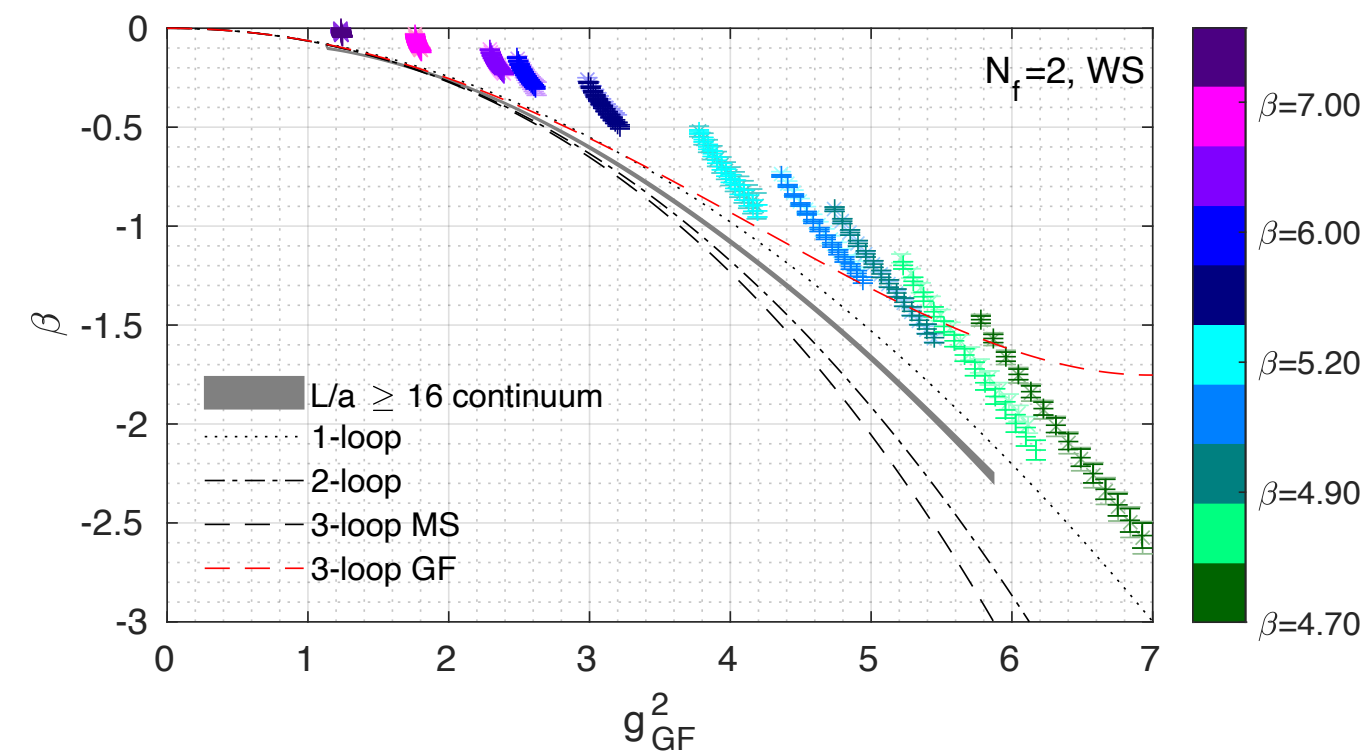
1. Infinite volume extrapolation ($1/L^4$ in the chirally symmetric regime)
2. Infinite flow time extrapolation (a^2/t) at fixed g_{GF}^2



The continuous β function, $N_f = 2$

AH, O. Witzel, *Phys.Rev.D* 101 (2020) 3

- Other flow/operator combos predict the same continuum limit, sometimes with larger cut-off effects (WS)
- It is possible to “optimize” the operator for different flows and find “scaling operators” (e.g. $X = 0.25W + 0.75C$ is optimal for Wilson flow) (WX)



GF β function is closest to 1-loop; consistent with GF 3-loop up to $g_{GF}^2 \approx 2.5$

The continuous β function, $N_f = 0 - 8$

In all cases:

- The β functions runs slower than 1-loop PT
- minimal cutoff effects if Zeuthen flow+Wilson op (or tree level improved coupling)

In the confining regime apparently $\beta(g^2) \propto g^2$

