

Tensor renormalization group approach to the four-dimensional lattice gauge theories

Shinichiro Akiyama ^{a), b)}

^{a)} Center for Computational Sciences, University of Tsukuba

^{b)} Endowed Chair for Quantum Software, University of Tokyo

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Research motivation

QFT on a lattice

- ✓ Quantum field theory (QFT) is a fundamental tool to describe various physical phenomena
 - We would like to solve the QFT to understand the physics and make a theoretical prediction. **One of the ways is to solve the path integral of the QFT**
- ✓ Once we consider the QFT on a lattice, we can regard the path integral just as a multiple integral

$$\int \prod_{x \in \mathbb{R}^d} d\phi(x) e^{-S[\phi]} \quad \rightarrow \quad \int \prod_{n \in \Lambda_d} d\phi(n) e^{-S_{\text{lat}}[\phi]}$$

- ✓ The QFT on a lattice provides us with a mathematically rigorous starting point and **we can investigate it based on the procedure of statistical physics**

Standard numerical approach for QFT on a lattice

✓ Monte Carlo (MC) simulation

- The MC is based on the probabilistic interpretation of the given Boltzmann weight $e^{-S[\phi]}$
- Lattice QCD is one of the most successful applications of the MC

✓ There are several difficulties in the MC simulation, though

- We encounter the **sign problem** when $e^{-S[\phi]}$ takes negative or complex value
- Bosons are easily dealt with the MC, but **fermions** are not
In the path integral formalism, fermions are described by the Grassmann numbers, which obey the anti-commutation $\psi\phi + \phi\psi = 0$

Different approach?

- ✓ There are many systems suffering from the sign problem
 - QCD at finite density
- ✓ Many unrevealed physics in such systems
 - Thermodynamic limit (or zero-temperature limit) is almost inaccessible w/ the standard MC approach
- ✓ We need different numerical methods which can give us insights for these systems
 - **How about tensor network?**

A quick overview of higher-dimensional TRG

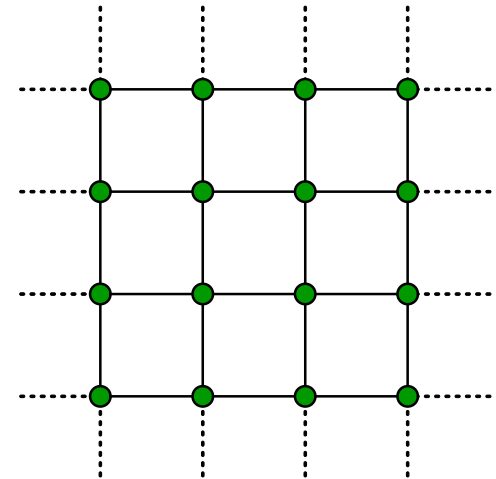
Tensor network & Lattice field theory

- ✓ **A method to investigate quantum many-body system expressing an objective function as a tensor contraction (= tensor network)**
 - Orús, APS Physics 1(2019)538-550
 - Bañuls-Cichy, Rep. Prog. Phys. 83(2020)024401
 - Meurice-Sakai-Unmuth-Yockey, Rev. Mod. Phys. 94(2022)025005
 - Okunishi-Nishino-Ueda, J. Phys. Soc. Jap. 91(2022)062001

- ✓ **TN method provides us with various ways to investigate lattice QFT**
 - **Hamiltonian formalism**
 - Describe a state vector as a TN, which is **variationally optimized**
 - Cf. **DMRG, TEBD**
 - White, PRL69(1992)2863-2866, White, PRB48(1993)10345-10356
 - Vidal, PRL91(2003)147902, Vidal, PRL98(2007)070201
 - Cf. Various talks in this workshop**
 - **Lagrangian formalism**
 - Describe a path integral as a TN, which is **approximately contracted**
 - Cf. **TRG, TNR, loop-TNR, GILT**
 - Levin-Nave, PRL99(2007)120601
 - Evenbly-Vidal, PRL115(2015)180405, Evenbly, PRB95(2017)045117
 - Yang-Gu-Wen, PRL118(2017)110504
 - Hauru-Delcamp-Mizera, PRB97(2018)045111

Pros and Cons

- ✓ Tensor renormalization group (TRG) approximately contracts a given TN based on the idea of real-space renormalization group
 - No sign problem
 - Thermodynamic limit
 - Grassmann variables
 - Path integral
 - **Higher dimension than $d = 2$**
- ✓ **Higher-dimensional TN computation is challenging**
 - Further algorithmic development is necessary
 - Improvement of the TRG based on the removal of short-range correlations
 - Lessons form other TN methods such as TTN, PEPS, isoTNS, etc



Cf. L. Vanderstraeten, 9/26

Cf. F. Pollmann, 9/27

Status of (3+1)D TN calculations

Hamiltonian formalism	Lagrangian formalism
<ul style="list-style-type: none"> • QED at finite density Magnifico+ 	<ul style="list-style-type: none"> • Ising model SA+ • Staggered fermion w/ strongly coupled U(N) Milde+ • Complex ϕ^4 theory at finite density SA+ • Nambu—Jona-Lasinio model at finite density SA+ • Real ϕ^4 theory SA+ • \mathbb{Z}_2 & \mathbb{Z}_3 gauge-Higgs at finite density SA-Kuramashi

- ✓ So far, the (3+1)D TN calculations have been driven by the Lagrangian formalism w/ the TRG approach
- ✓ Development of **parallel computing method** specialized for individual algorithms to reduce their execution time per process
[SA+, PoS\(LATTICE2019\)138](#)
[Yamashita-Sakurai, CPC278\(2022\)108423](#)
- ✓ **Application of ML techniques and/or GPU** is a recent hot topic
[Liao+, PRX9\(2019\)031041](#)
[R. G. Jha-Samlodia, CPC\(2023\)108941](#)

Procedure of TRG approach

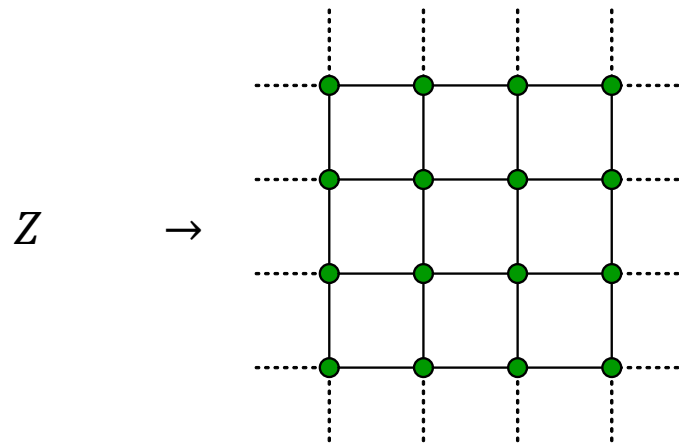
1) Represent the path integral as a tensor network

- Lattice QFTs can be easily represented by TN

Cf. Y. Meurice, 9/15

Meurice-Sakai-Unmuth-Yockey, Rev. Mod. Phys. 94(2022)025005

Meurice, "Quantum Field Theory, A quantum computation approach"



2) Take contractions approximately

- Various algorithms are proposed to achieve this mission
- In 2D, we can also use other schemes to take contractions approximately

Cf. iTEBD for 2D classical Ising model: Orús-Vidal, PRB78(2008)155117

Higher-order TRG (HOTRG)

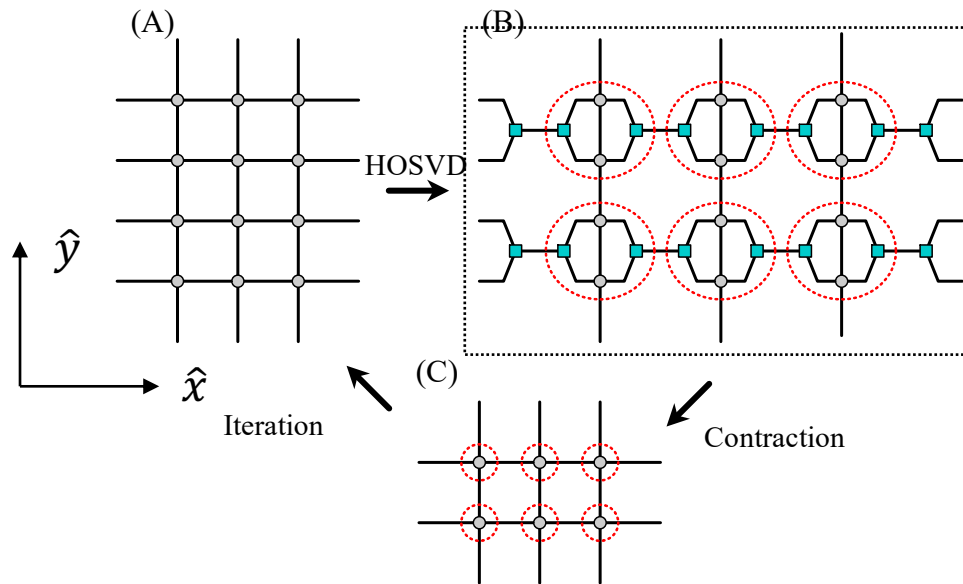
Xie-Chen-Qin-Zhu-Yang-Xiang, PRB86(2012)045139

✓ Applicable to any d -dimensional lattice

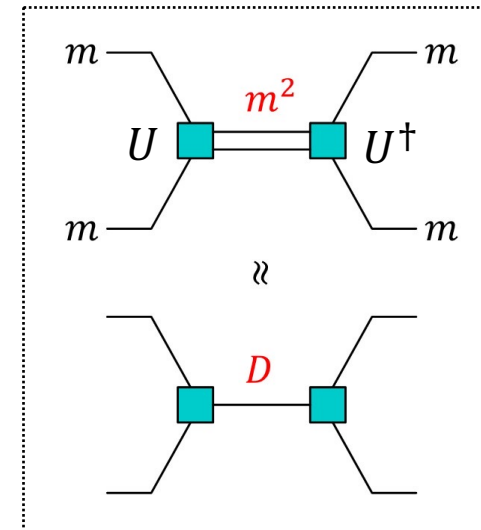
✓ # of tensors are reduced to half

Iterating this CG n times, we can approximately contract 2^n tensors

Memory	Complexity
$O(D^{2d})$	$O(D^{4d-1})$



Sequential coarse-graining along with each direction

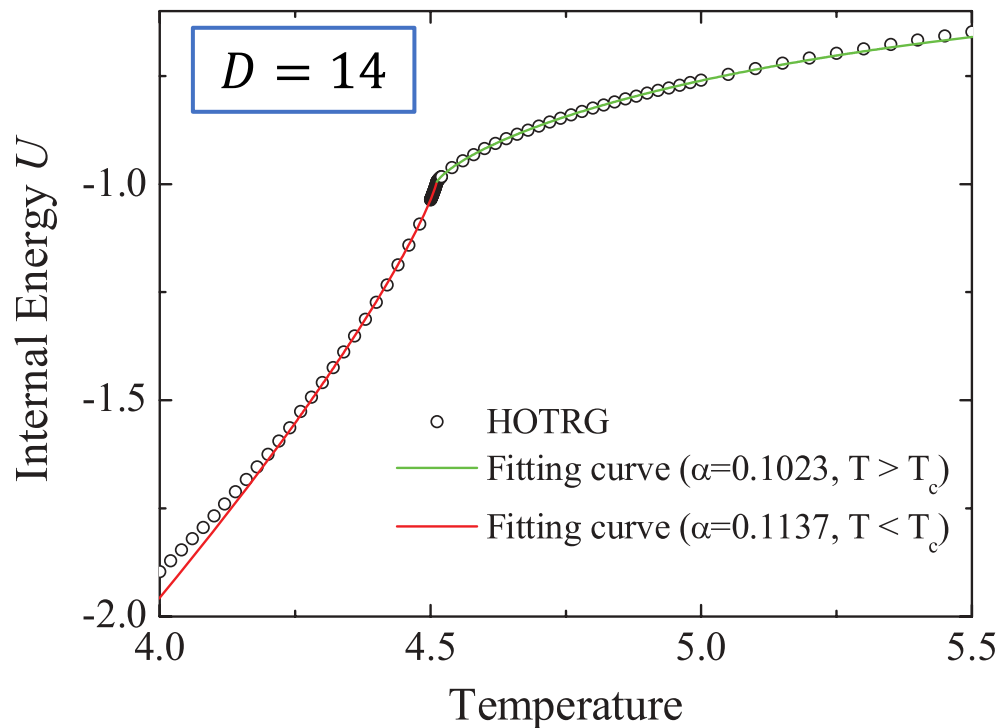


D : bond dimension

Example: 3D Ising model w/ HOTRG

Xie-Chen-Qin-Zhu-Yang-Xiang, PRB86(2012)045139

- ✓ Critical point is precisely located with relatively small bond dimension



Critical point

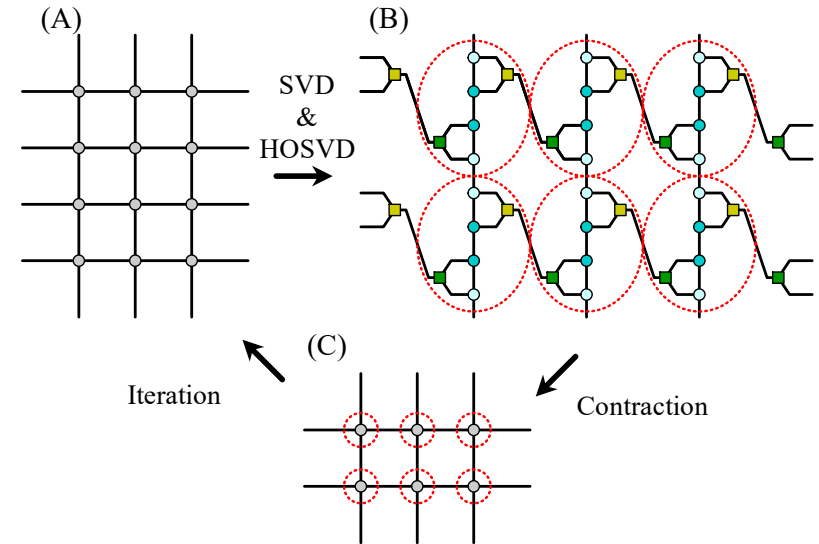
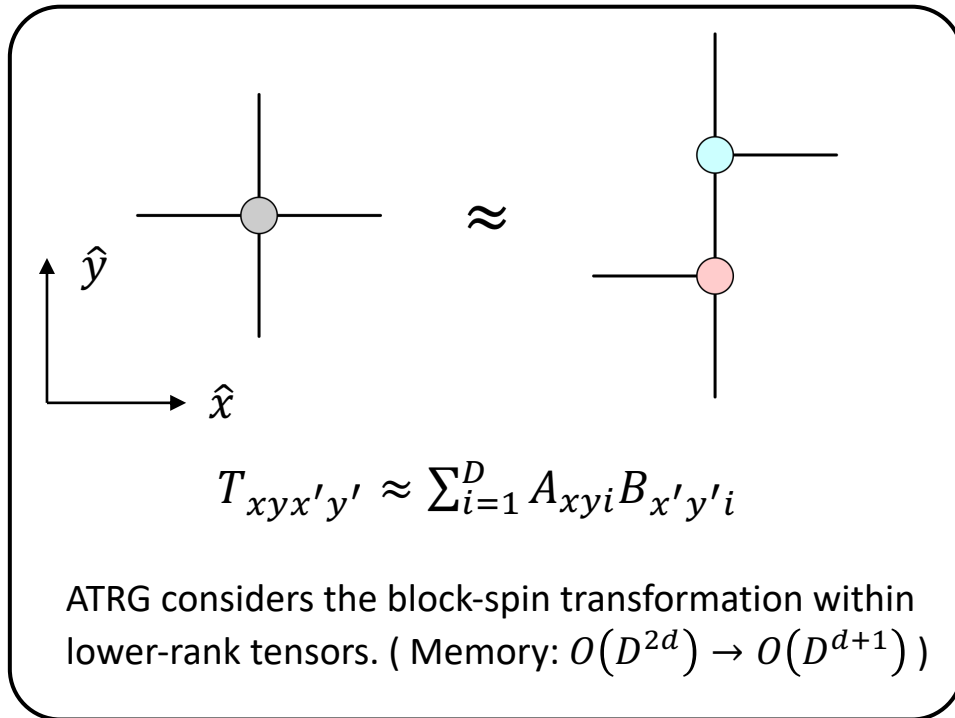
Method	T_c
HOTRG ($D = 16$, from U)	4.511544
HOTRG ($D = 16$, from M)	4.511546
Monte Carlo ³⁷	4.511523
Monte Carlo ³⁸	4.511525
Monte Carlo ³⁹	4.511516
Monte Carlo ³⁵	4.511528
Series expansion ⁴⁰	4.511536
CTMRG ¹²	4.5788
TPVA ¹³	4.5704
CTMRG ¹⁴	4.5393
TPVA ¹⁶	4.554
Algebraic variation ⁴¹	4.547

Anisotropic TRG (ATRG)

Adachi-Okubo-Todo, PRB102(2020)054432

- ✓ Applicable to any d -dimensional lattice
- ✓ More economic than the HOTRG

Memory	Complexity
$O(D^{d+1})$	$O(D^{2d+1})$

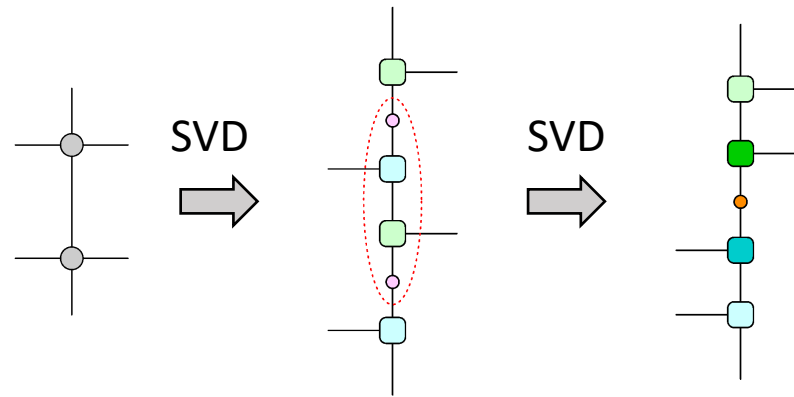


of tensors are reduced to half

Canonical form in ATRG

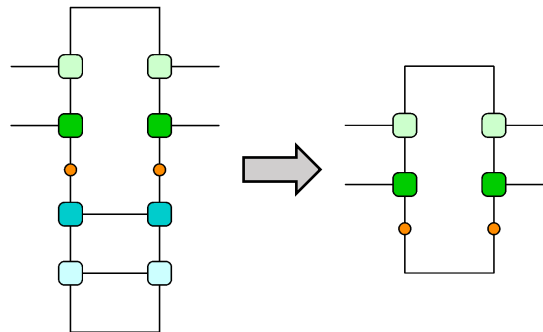
✓ ATRG converts two adjacent tensors into a **canonical form**

- Canonical form is an important idea in MPS [Schollwöck, Annals of Physics 326\(2011\)96-192](#)
Cf. [F. Pollmann, 9/27](#)



✓ “Reduced density matrix” is simplified thanks to the canonical form

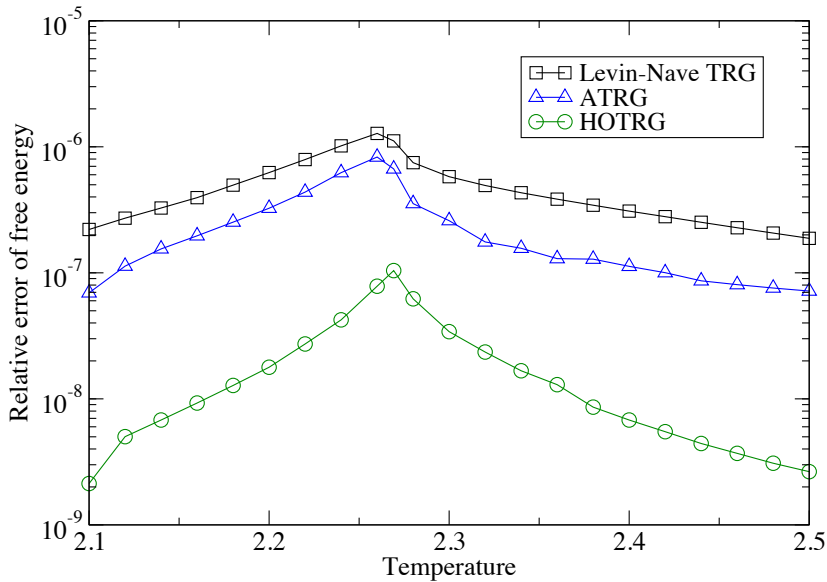
- Highly helpful in practical computations



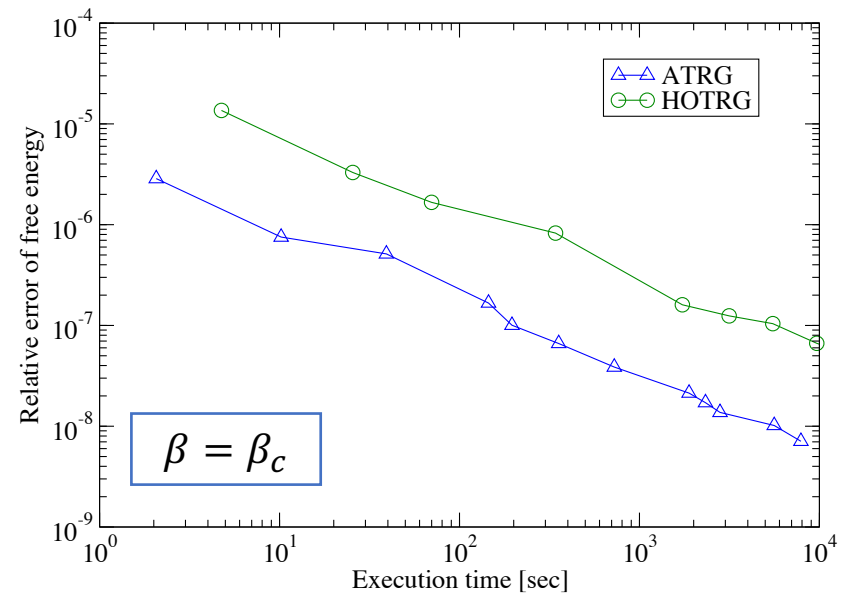
Benchmarking w/ 2D Ising model

- ✓ HOTRG & ATRG improve the accuracy of the original (LN-)TRG at the same D
The exact solution is well reproduced
- ✓ ATRG shows better performance than the HOTRG at the same execution time

Comparison of three types of TRG
w/ $D = 24$



Relative error vs execution time



Grassmann TRG approach

Gu-Verstraete-Wen, arXiv:1004.2563

- ✓ Any TRG algorithm can be applied for fermions
Fermionic path integral can be expressed as a tensor network generated by **Grassmann tensors**

$$\mathcal{J}_{\eta_1 \eta_2 \eta_3 \dots} = \sum_{i_1, i_2, i_3, \dots} T^{i_1 i_2 i_3 \dots} \eta_1^{i_1} \eta_2^{i_2} \eta_3^{i_3} \dots$$

Gu, PRB88(2013)115139

Shimizu-Kuramashi, PRD90(2014)014508

Takeda-Yoshimura, PTEP2015(2015)043B01

Meurice, PoS LATTICE2018(2018)231

Bao's thesis, PhD, Uwaterloo

SA-Kadoh, JHEP10(2021)188

- ✓ A clear correspondence btw tensors and Grassmann tensors

	Tensor	Grassmann tensor
index	integer	Grassmann number
contraction	$\sum_i \dots$	$\int \int d\bar{\eta} d\eta e^{-\bar{\eta}\eta} \dots$

$$e^{A\bar{\psi}_n \psi_{n+\mu}} = \left(\int \int d\bar{\eta}_n d\eta_n e^{-\bar{\eta}_n \eta_n} \right) \exp[-\sqrt{A}\bar{\psi}_n \eta_n + \sqrt{A}\bar{\eta}_n \psi_{n+\mu}]$$

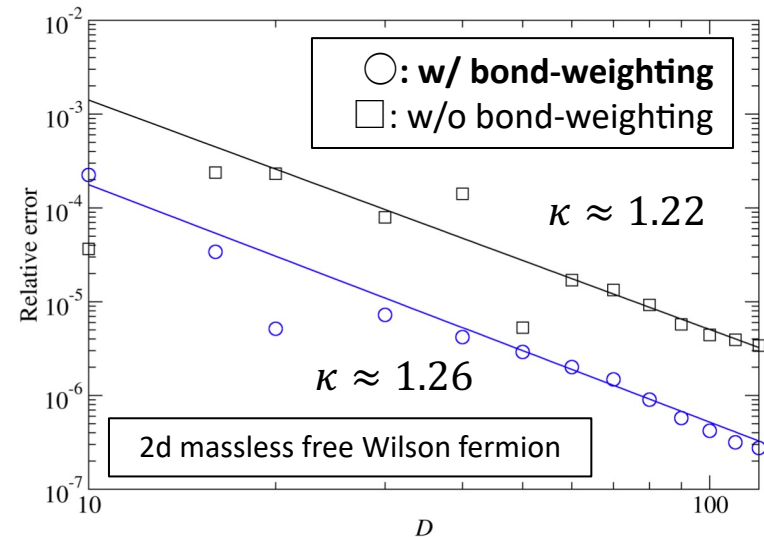
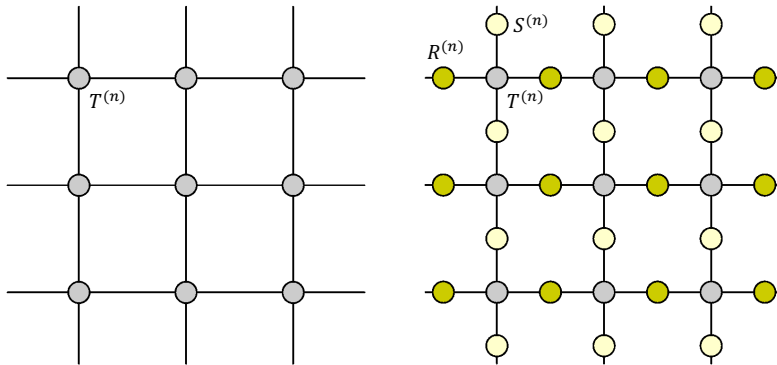
Several public codes of Grassmann TRG

✓ Grassmann bond-weighted TRG <https://github.com/akiyama-es/Grassmann-BTRG> by SA

- Originally proposed for spin models
- Works also well for fermions

Adachi-Okubo-Todo, PRB105(2022)L060402

SA, JHEP11(2022)030



$$\delta f \sim D^{2\kappa} \text{ w/ } \kappa = 1.344 \text{ when } c = 1$$

Tagliacozzo+, PRB78(2008)024410

Pollmann+, PRL102(2009)255701

✓ GrassmannTN <https://github.com/ayosprakob/grassmanntn> by A. Yosprakob

- Python package for Grassmann TRG computations

Yosprakob, arXiv:2309.07557

TRG study of (3+1)D \mathbb{Z}_2 & \mathbb{Z}_3 gauge-Higgs models

First application of TRG to (3+1)D LGT

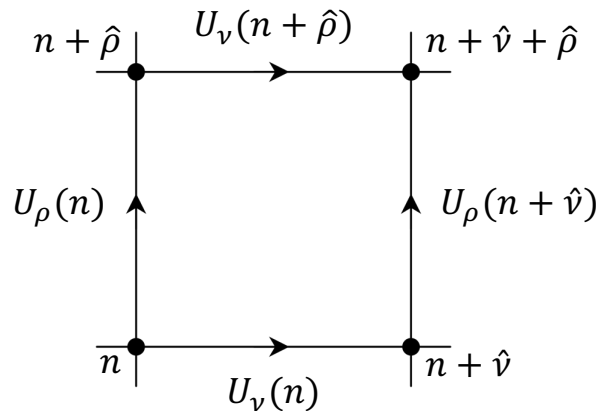
SA-Kuramashi, JHEP05(2022)102

SA-Kuramashi, arXiv:2304.07934 (To appear in JHEP)

\mathbb{Z}_n gauge-Higgs model in the unitary gauge

$$S = -\beta \sum_n \sum_{\nu > \rho} \text{Re}[U_\nu(n) U_\rho(n + \hat{\nu}) U_\nu^*(n + \hat{\rho}) U_\rho^*(n)]$$

$$- \eta \sum_n \sum_\nu [e^{\mu \delta_{\nu,4}} \sigma^*(n) U_\nu(n) \sigma(n + \hat{\nu}) + e^{-\mu \delta_{\nu,4}} \sigma^*(n) U_\nu^*(n - \hat{\nu}) \sigma(n - \hat{\nu})]$$



$U_\nu(n) (\in \mathbb{Z}_n)$: link variable living on edges
 $\sigma(n) (\in \mathbb{Z}_n)$: matter field living on sites
 β : inverse gauge coupling
 η : spin-spin coupling
 μ : chemical potential

✓ Unitary gauge fixing: $\sigma^*(n) U_\nu(n) \sigma(n + \hat{\nu}) \mapsto U_\nu(n)$

$$S = -\beta \sum_n \sum_{\nu > \rho} \text{Re}[U_\nu(n) U_\rho(n + \hat{\nu}) U_\nu^*(n + \hat{\rho}) U_\rho^*(n)]$$

$$- 2\eta \sum_n \sum_\nu [\cosh(\mu \delta_{\nu,4}) \text{Re} U_\nu(n) + i \sinh(\mu \delta_{\nu,4}) \text{Im} U_\nu(n)]$$

Motivation of studying \mathbb{Z}_n gauge-Higgs model

✓ The simplest lattice gauge theory coupling to a matter field

- A good target to see whether the TRG is efficient for the (3+1)D lattice gauge theory or not

✓ The model possesses the critical endpoint (CEP)

- QCD at finite temperature and density also has the CEP
Can we use the TRG to specify the precise location of CEP?

✓ We can consider the model at finite density

- We can investigate how the CEP moves by finite chemical potential

The \mathbb{Z}_2 model is free from the sign problem but the \mathbb{Z}_3 model is not

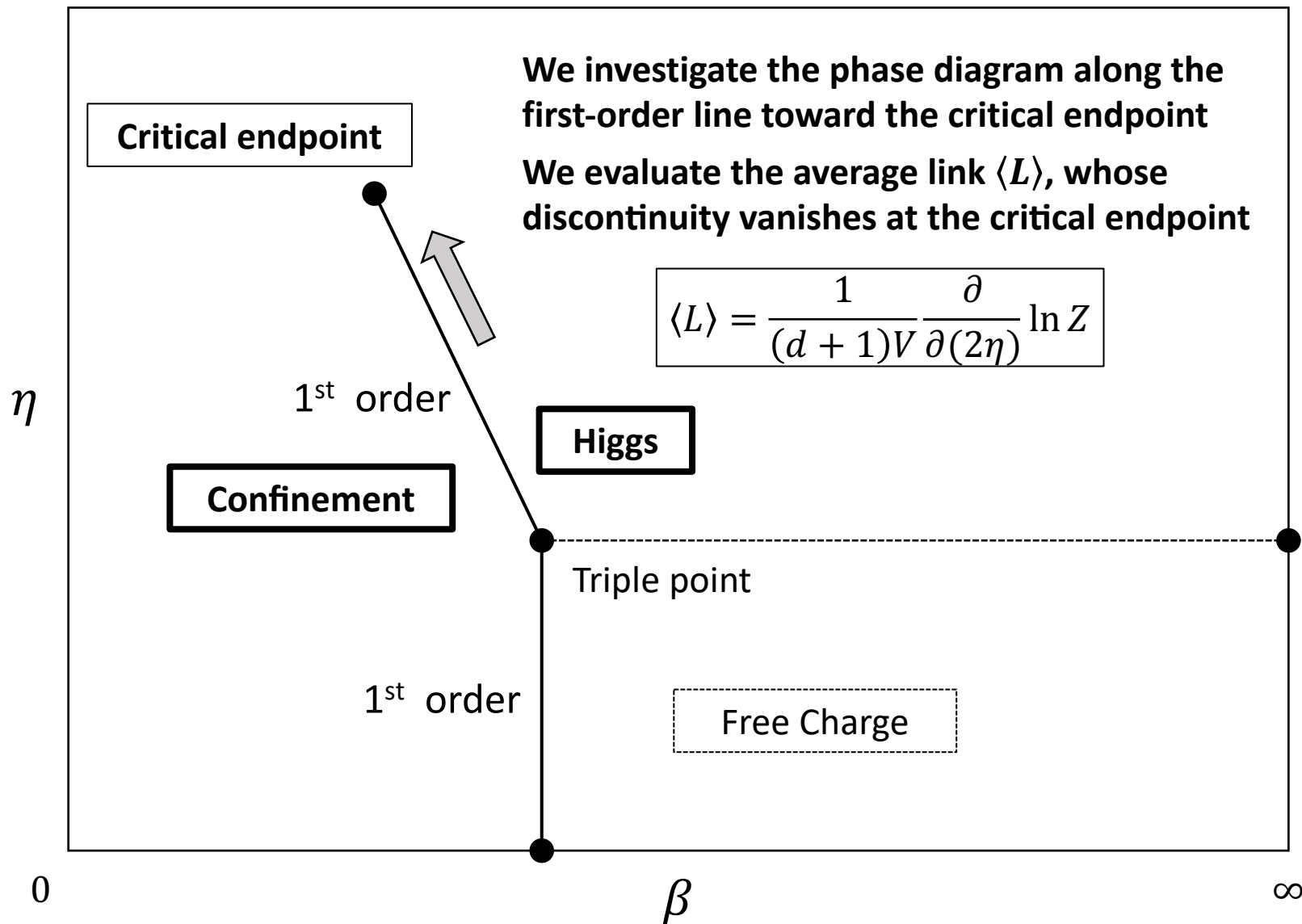
Cf. TRG studies of gauge-Higgs models in 2D

Unmuth–Yockey+, PRD98(2018)094511, Bazavov+, PRD99(2019)114507, Butt+, PRD101(2020)094509

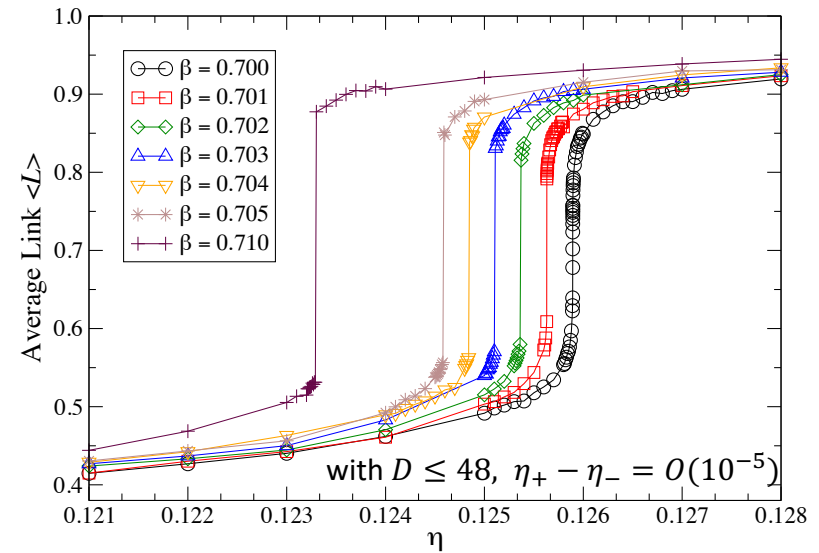
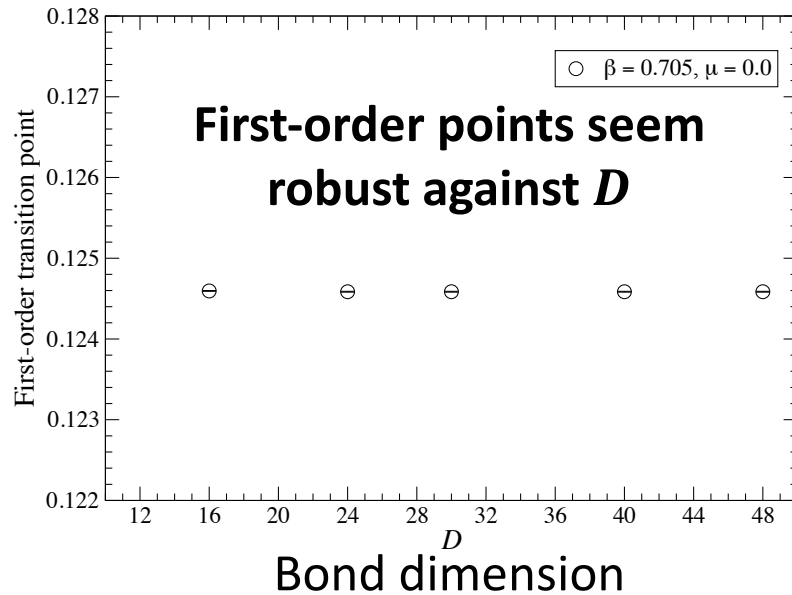
✓ We use the ATRG w/ parallel computation (slicing fundamental tensors)

SA+, PoS(LATTICE2019)138

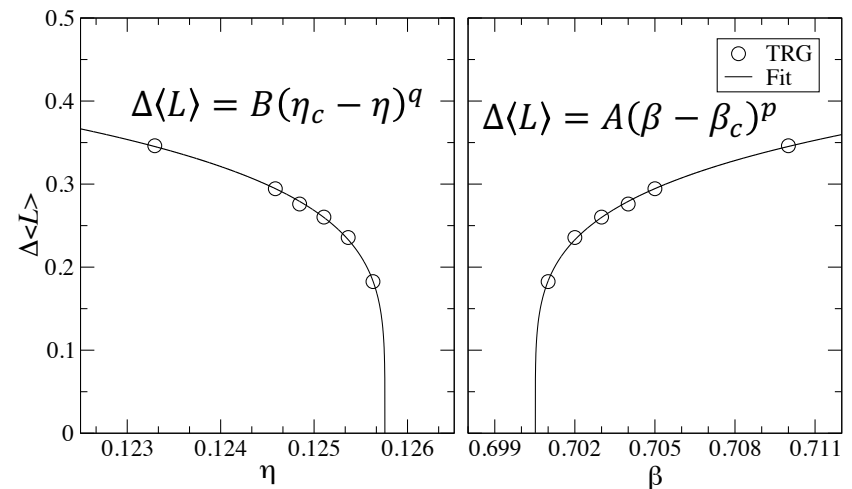
Phase diagram of the (3+1)D model at $\mu = 0$



Study of the (2+1)D model at $\mu = 0$

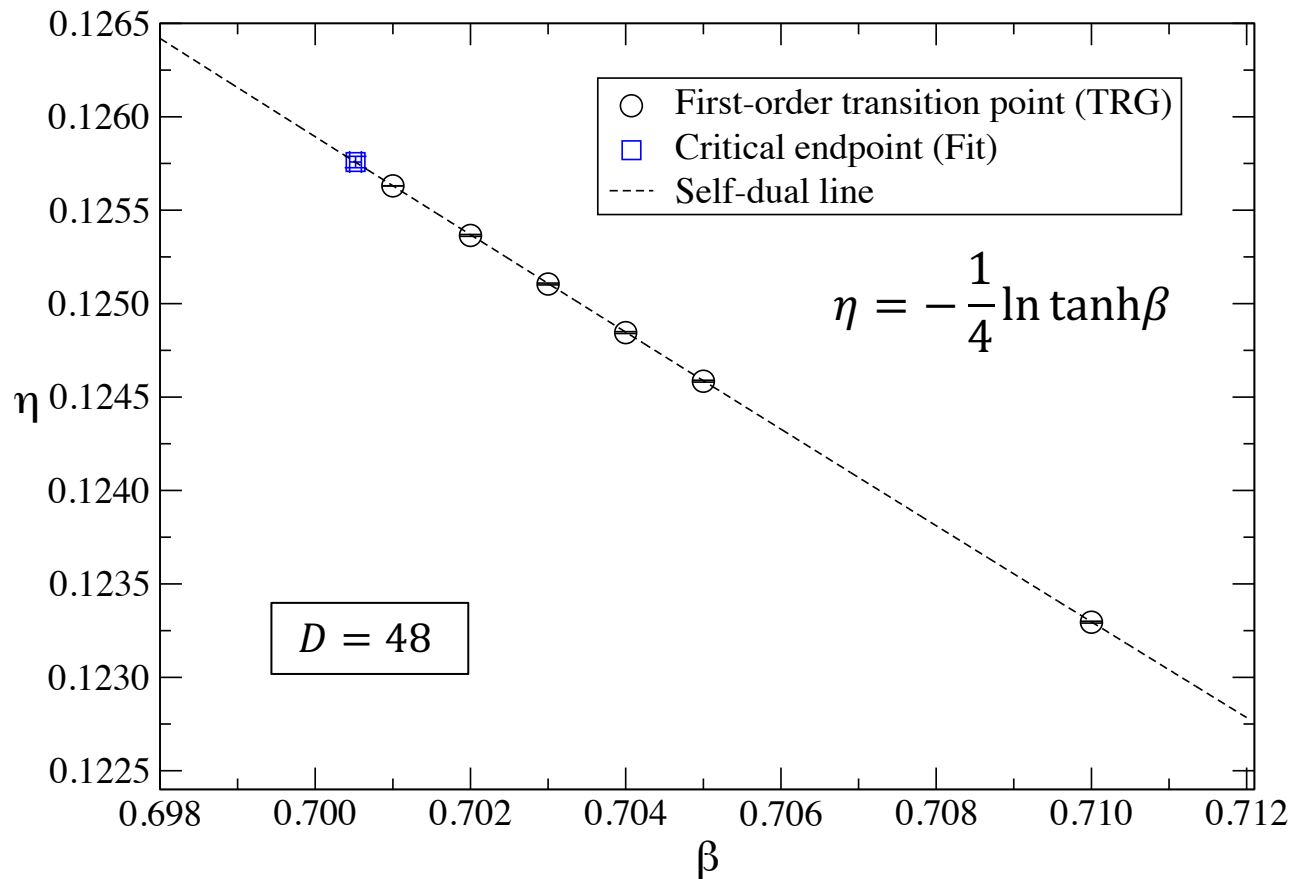


MC Somoza+, PRX11(2021)041008	$\beta_c \approx 0.701$
TRG this work	$\beta_c = 0.70051(7)$

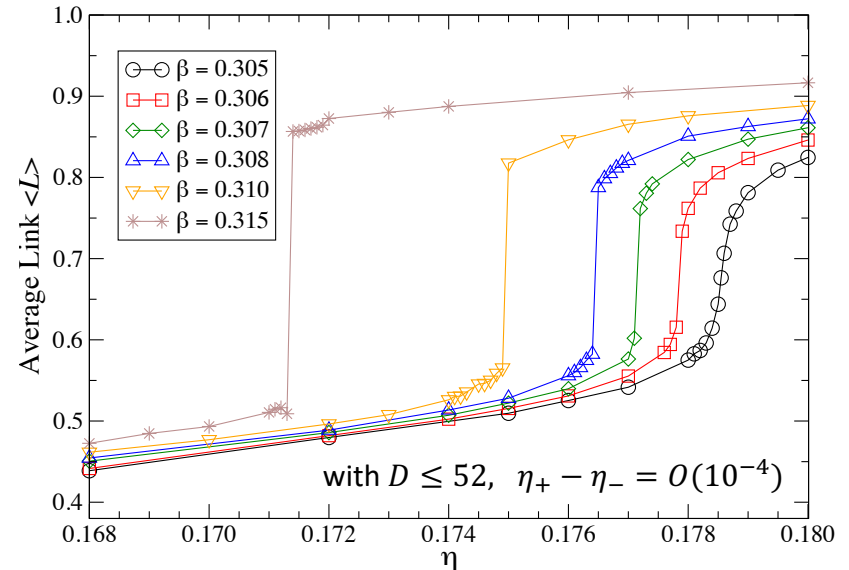
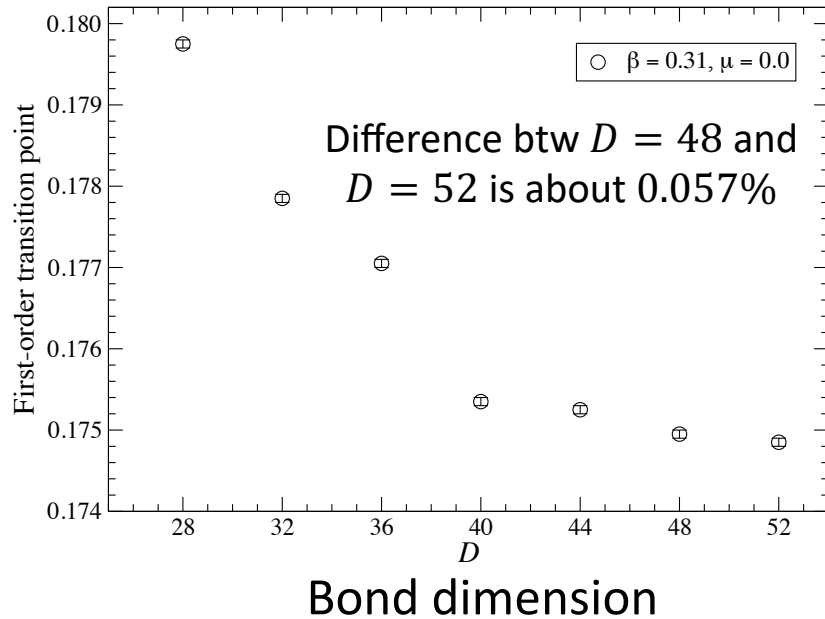


Comparison with the self-dual line

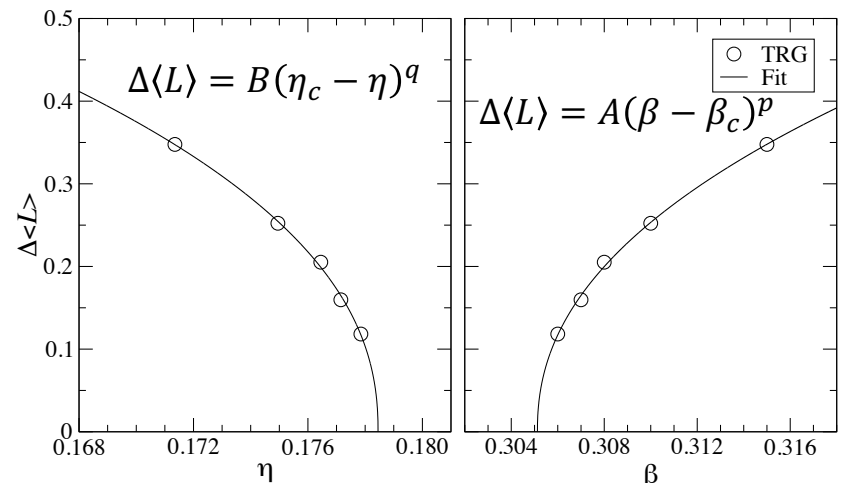
✓ All transition points are well located on the self-dual line



(3+1)D model at vanishing density

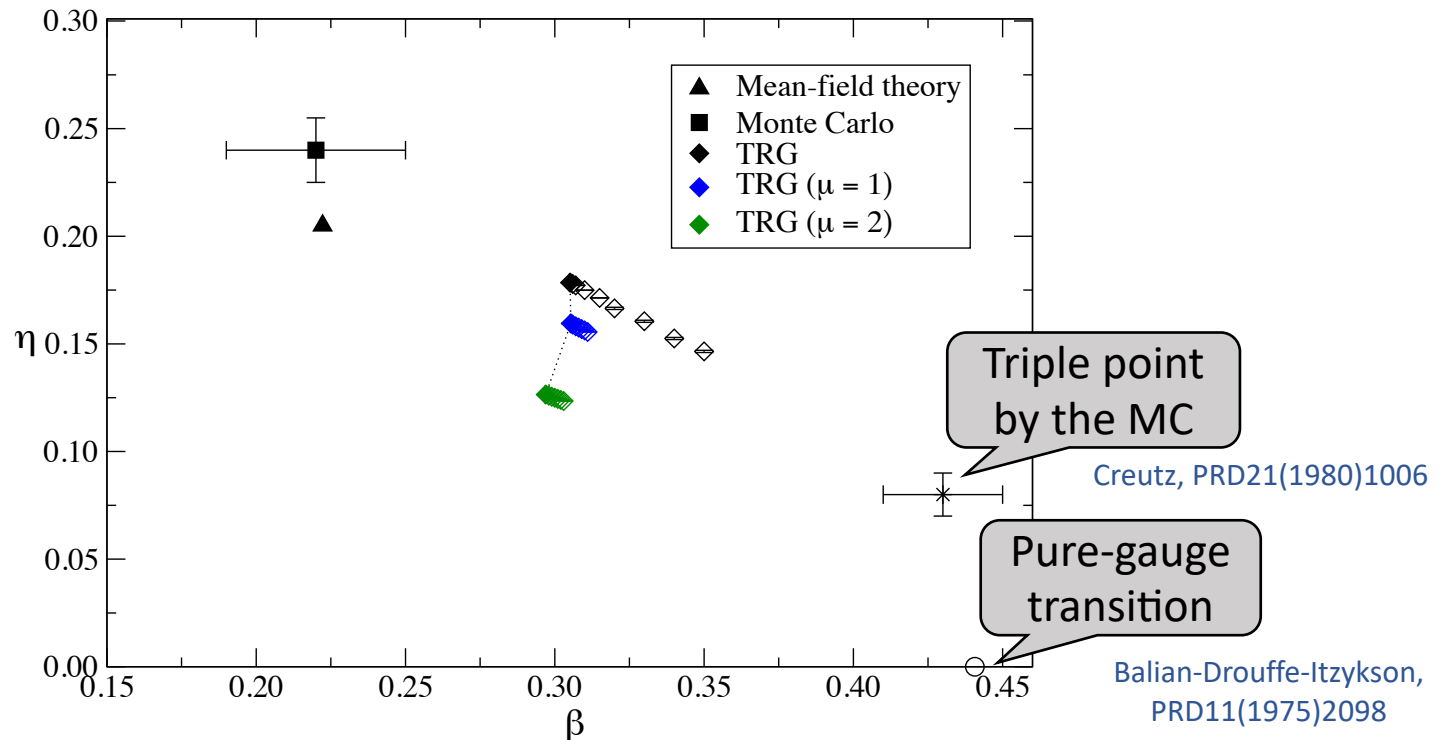


Mean-field Brezin-Drouffe, NPB200(1982)93	$(\beta_c, \eta_c) = (0.22, 0.205)$
MC on $V = 8^4$ Creutz, PRD21(1980)1006	(β_c, η_c) $= (0.22(3), 0.24(2))$
TRG w/ $D = 52$ this work	(β_c, η_c) $= (0.3051(2), 0.1784(2))$



Status of the phase diagram near the CEP

- ✓ It seems that TRG and MC share a similar first-order line at $\mu = 0$
- ✓ Revisiting the CEP by the modern MC simulation should be meaningful



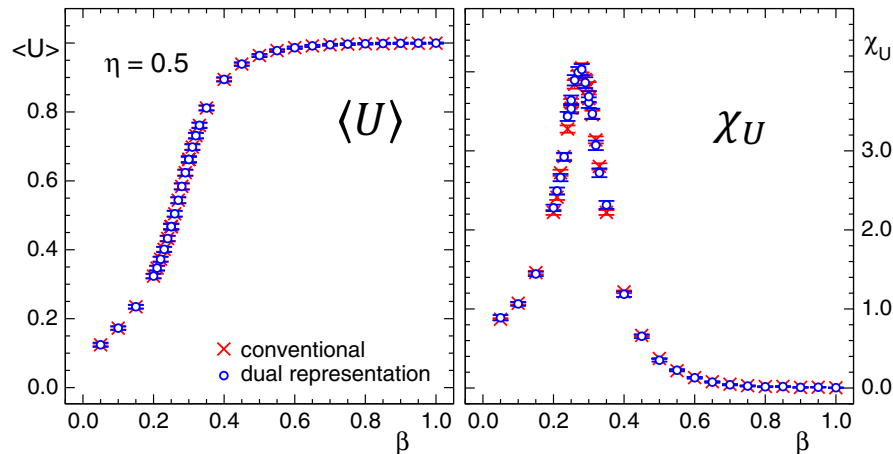
(3+1)D \mathbb{Z}_3 model at vanishing density 1/2

✓ Comparison btw MC and TRG via the average plaquette $\langle U \rangle$ and its susceptibility

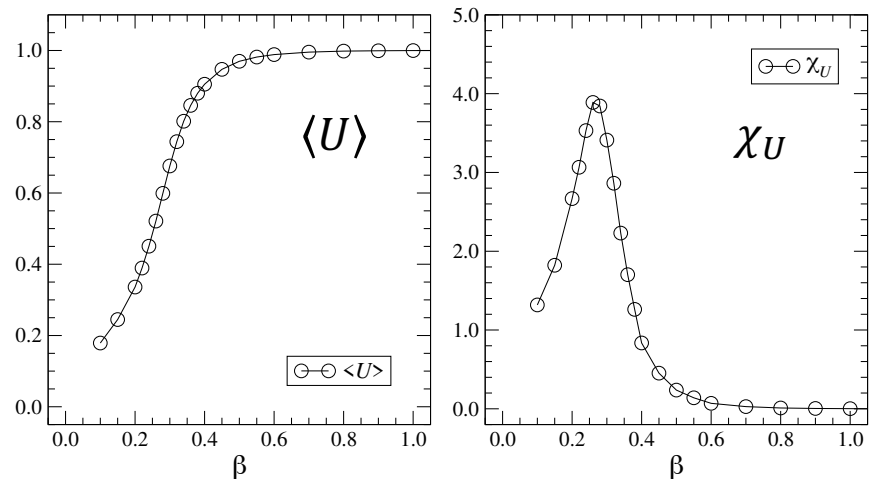
- Good agreement just w/ $D = 45$ at finite- η regime $\langle U \rangle = -\frac{1}{6V} \frac{\partial}{\partial \beta} \ln Z$
- The susceptibility of $\langle U \rangle$ is obtained by numerical difference in case of TRG

Monte Carlo

Gattringer-Schmidt, PRD86(2012)094506



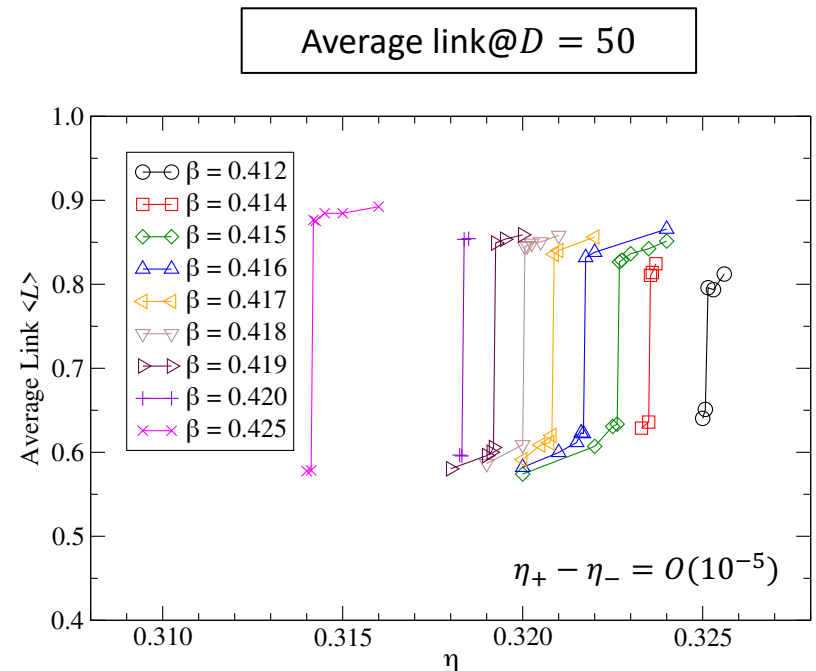
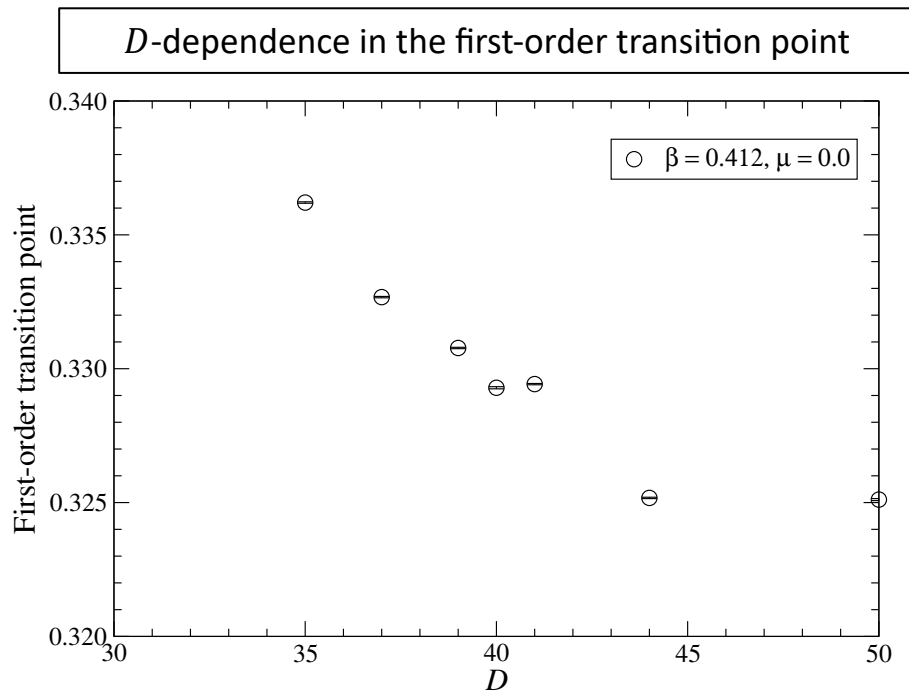
TRG w/ $D = 45$



$\eta = 0.5, \mu = 0$

(3+1)D \mathbb{Z}_3 model at vanishing density 2/2

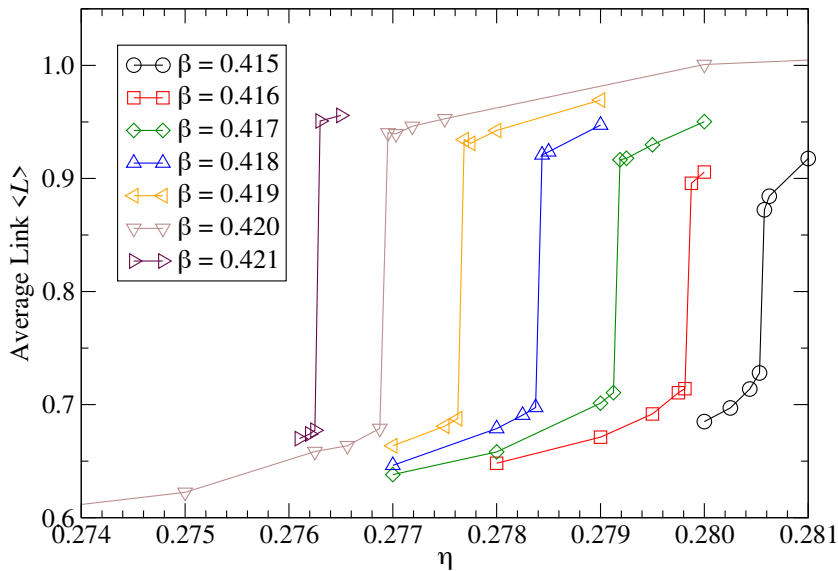
- ✓ Location of the transition point seems converging w.r.t D
 - Relative error btw $D = 44$ and $D = 50$ is **0.019%**
- ✓ $\Delta\langle L \rangle$ becomes smaller when β becomes smaller as expected



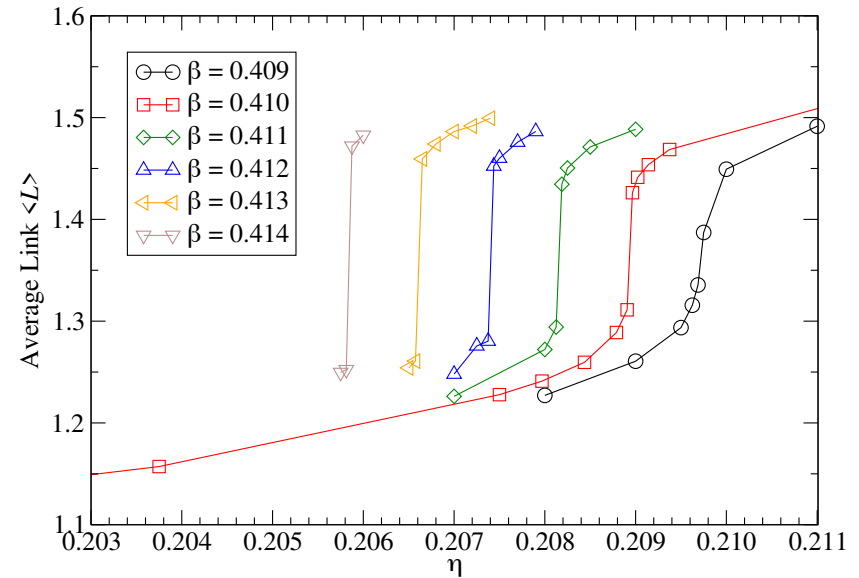
(3+1)D \mathbb{Z}_3 model at finite density

✓ Again, $\Delta\langle L \rangle$ becomes smaller when β becomes smaller as expected

Average link@ $\mu = 1$



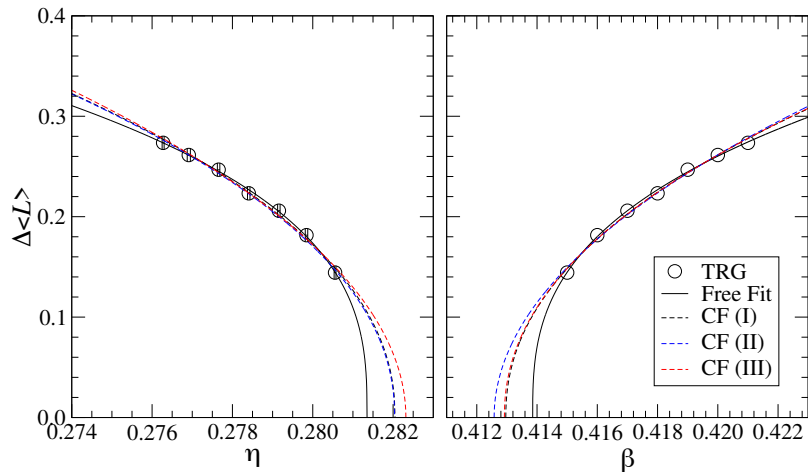
Average link@ $\mu = 2$



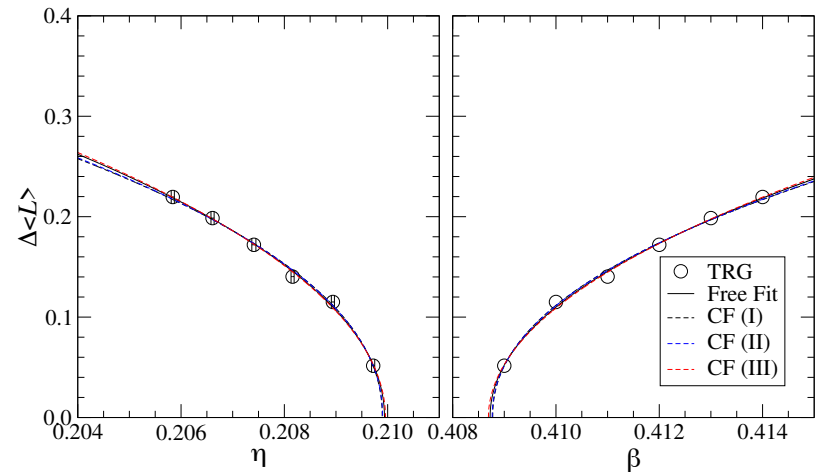
CEP in (3+1)D \mathbb{Z}_3 model at finite μ

- ✓ CEP is determined via the similar fit to the \mathbb{Z}_2 model
 - Fit by $\Delta\langle L \rangle = A(\beta - \beta_c)^p$ and $\Delta\langle L \rangle = B(\eta_c - \eta)^q$
 - According to the mean field theory, $p = q = 0.5$
 - The simultaneous fit among different μ suggests $p = \mathbf{0.46(2)}$, $q = \mathbf{0.46(3)}$

$\mu = 1$



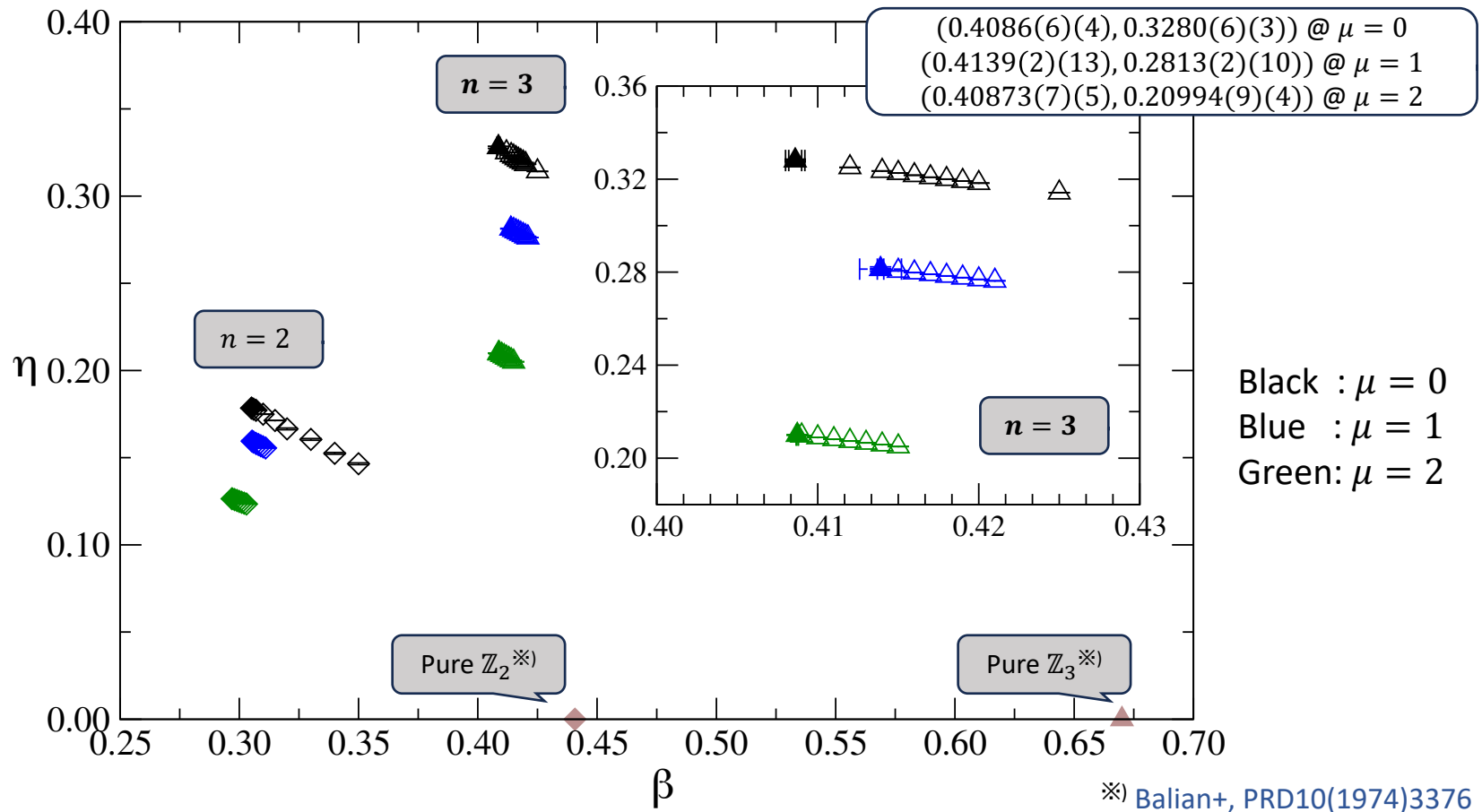
$\mu = 2$



Phase diagram of \mathbb{Z}_n gauge-Higgs model ($n = 2,3$)

✓ n -dependence in the resulting CEP is consistent w/ previous studies

Cf. U(1) gauge-Higgs studies, Baig-Clua, PRD57(1998)3902, Franzki+, PRD57(1998)6625



Summary

- ✓ TRG is a typical TN algorithm, which enables us to perform TN contraction approximately using the idea of RSRG
- ✓ TRG w/ parallel computation has been a good way to investigate higher-dimensional QFT on a thermodynamic lattice
- ✓ Several public codes for Grassmann TRG
- ✓ **The first application of TRG for (3+1)D LGT has been made**
We have obtained the TRG estimates of **CEP in \mathbb{Z}_2 & \mathbb{Z}_3 gauge-Higgs model at finite density**

Future Perspective

✓ A next interesting (challenging) target can be the (3+1)D QED

- Variational approach based on the tree TN for the (3+1)D lattice QED ($L \leq 8$)

Magnifico+, Nature Commun. 12(2021)1

✓ How can we approach $D \rightarrow \infty$?

- This problem should be addressed

Cf. Finite-entanglement scaling:

Tagliacozzo+, PRB78(2008)024410

Pollmann+, PRL102(2009)255701

Cf. L. Vanderstraeten, 9/26

✓ Although TRG is based on Lagrangian formalism, some problems are shared with quantum computations based on Hamiltonian formalism

- TRG may give us insights from the viewpoint of classical computation and vice versa
- How can we deal with higher-dimensional non-abelian gauge theories with TN?

Cf. TRG approach for SU(N) gauge theory Fukuma-Kadoh-Matsumoto, PTEP2021(2021)123B03

Hirasawa+, JHEP12(2021)011

Kuwahara-Tsuchiya, PTEP2022(2022)093B02

- Multi-flavor fermions?

SA, PRD108(2023)034514

Yosprakob-Nishimura-Okunishi, arXiv:2309.01422