

Symmetry-resolved entanglement in AdS and BCFT

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Charged Moments in W_3 Higher Spin Holography

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Symmetry-resolved entanglement in AdS_3/CFT_2 coupled to $U(1)$ Chern-Simons theory



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Symmetry-resolved entanglement for excited states and two entangling intervals in AdS_3/CFT_2

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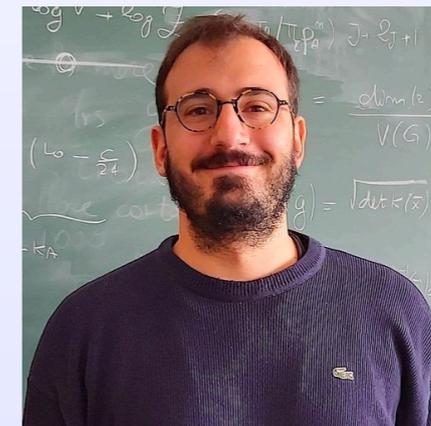
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On the boundary conformal field theory approach to symmetry-resolved entanglement

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Symmetry Resolved Entanglement

Entanglement entropy in each charge sector, e.g. U(1)

charge operator Q

$$Q = Q_A \oplus Q_B.$$

eigenstate of Q : $[\rho, Q] = 0.$



$$[\rho_A, Q_A] = 0.$$

Block decomposition:

$$\rho_A = \oplus_q \rho_A(q)$$

Symmetry Resolved Renyi and Entanglement Entropy:

$$S_n(q) = \frac{1}{1-n} \log \text{Tr} \left(\frac{\rho_A(q)}{P_A(q)} \right)^n$$

$$P_A(q) = \frac{\text{Tr} \rho_A(q)}{\text{Tr} \rho_A} = \text{Tr} \rho_A(q)$$

$$S_1(q) = \lim_{n \rightarrow 1} S_n(q) = -\text{Tr} \left(\frac{\rho_A(q)}{P_A(q)} \log \frac{\rho_A(q)}{P_A(q)} \right)$$

AdS₃ dual to U(1)_k Kac-Moody CFT

3D Einstein-Hilbert gravity

$$S_g = \frac{1}{16\pi G_3} \int d^3x \sqrt{g} \left(R + \frac{2}{L^2} \right)$$

L... Curvature Radius of AdS₃ space-time

To suppress quantum gravity effects: $c = \frac{3L}{2G_3} \gg 1$

U(1)_k Chern-Simons theory

$$S_{CS} = \frac{ik}{8\pi} \int A \wedge dA$$

k... Chern-Simons level

Asymptotic symmetry analysis:

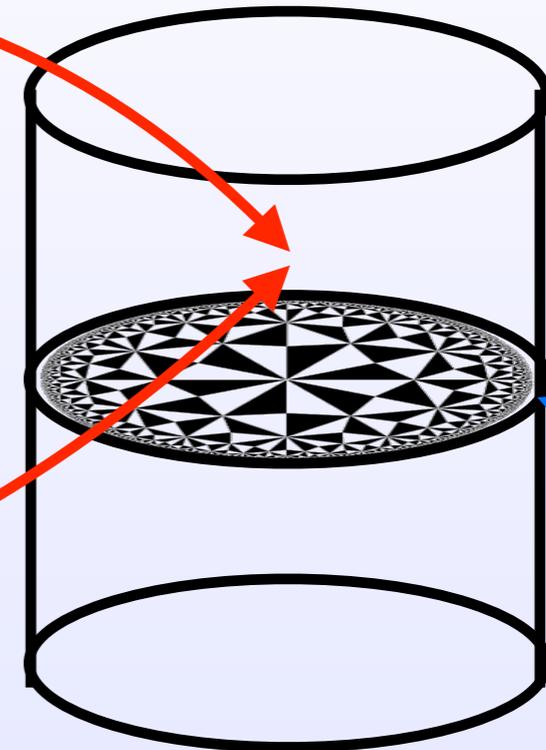
Boundary theory has U(1)_k Kac-Moody symmetry

Hilbert space factorizes into gravitational and U(1) part

P. Kraus [arXiv:hep-th/0609074](https://arxiv.org/abs/hep-th/0609074)

S. Zhao, C. Northe, RM, JHEP 2021 (arXiv: 2012.11274)

K. Weisenberger, S. Zhao, C. Northe, RM, JHEP 2021 (arXiv: 2108.09210)

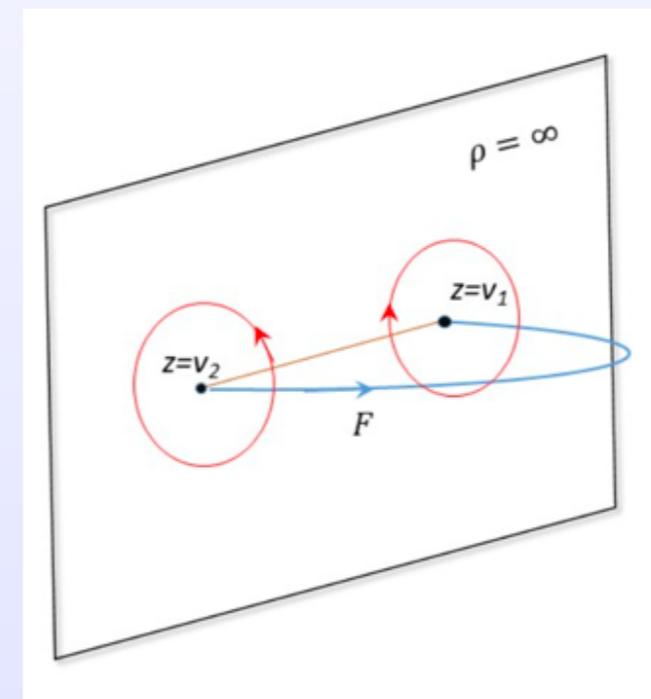
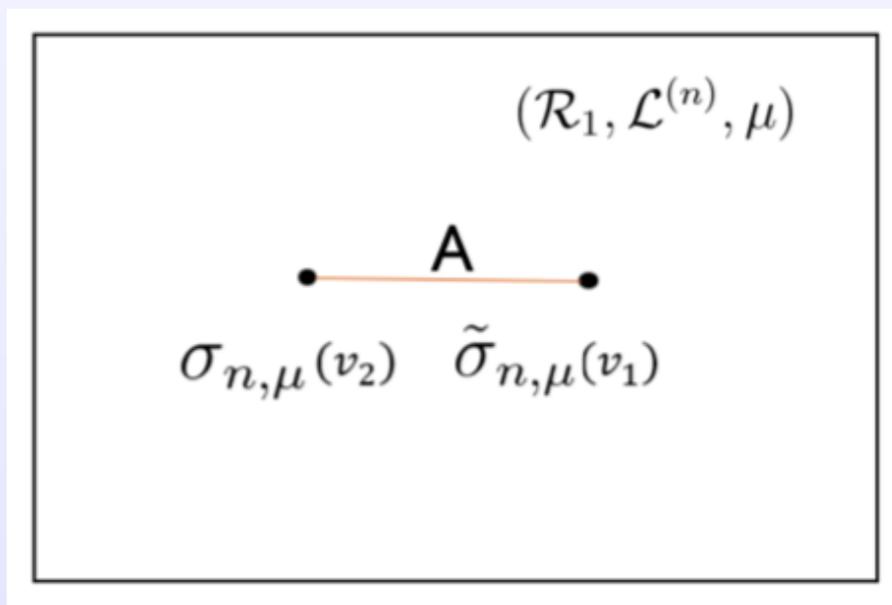


Single Interval SREE in AdS₃/CFT₂

Charged twist operator induces flux:

Sela, Goldstein PRL 2018

Wilson line following the Ryu-Takayanagi geodesic



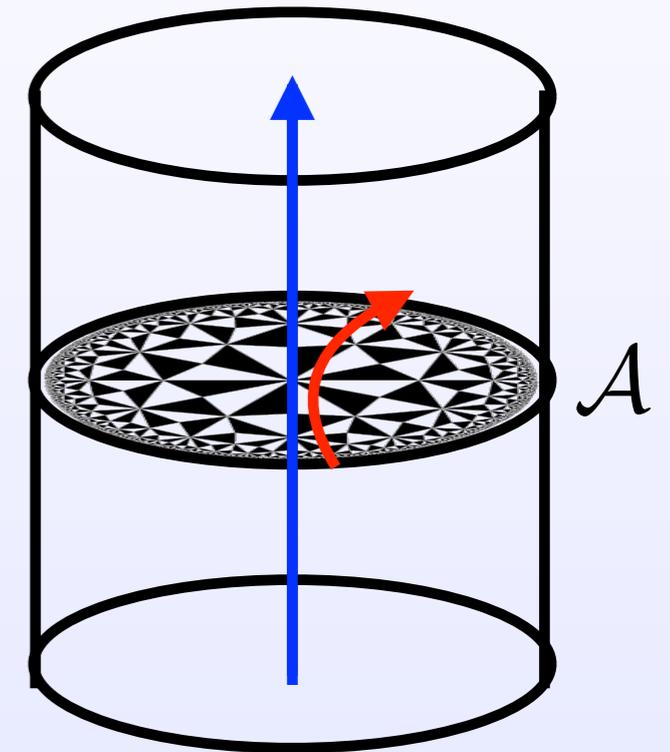
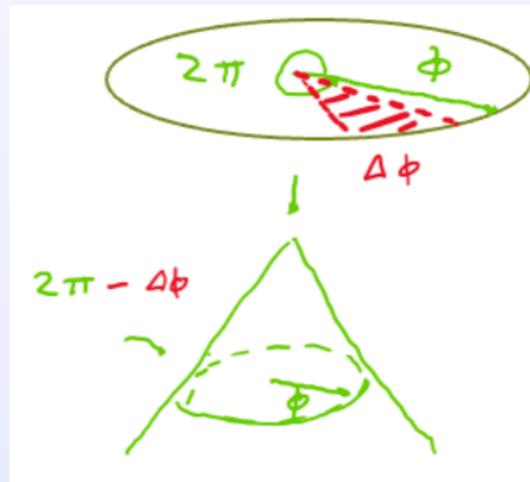
CFT₂ result matches AdS₃ result for $c \gg 1$

$$S_1(q) = \frac{c}{6} \ell - \frac{1}{2} \log \left(\frac{k\ell}{2\pi} \right) \quad \text{with} \quad \ell = 2 \log \frac{|v_1 - v_2|}{\epsilon}$$

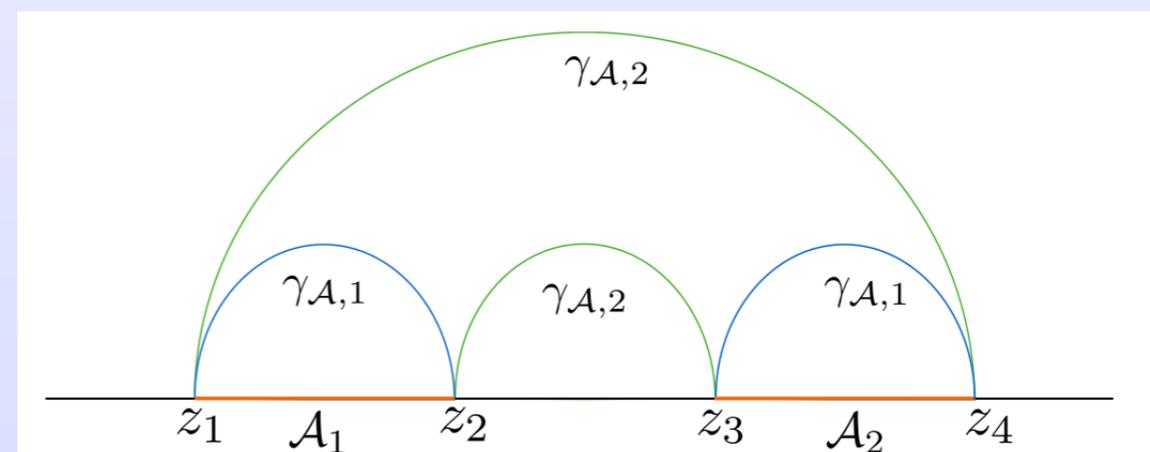
Equipartition of entanglement!

Further Checks

- Single interval with uncharged/charged heavy primary insertions



- Two intervals in the ground state



S. Zhao, C. Northe, RM, JHEP 2021 (arXiv: 2012.11274)

K. Weisenberger, S. Zhao, C. Northe, RM, JHEP 2021 (arXiv: 2108.09210)

Breakdown of Equipartition

SL(3,R) Higher Spin Gravity in 3D: W_3 symmetric CFT
Energy-momentum tensor plus Spin 3 current

$$T(z)W(w) = \frac{3W(w)}{(z-w)^2} + \frac{\partial W(w)}{z-w} + \dots$$

Charged moments for a single interval:

Topological black hole grand canonical partition function

Perturbative result in μ

Kraus+Perlmutter (2011)

Gaberdiel+Hartman+Jin (2012)

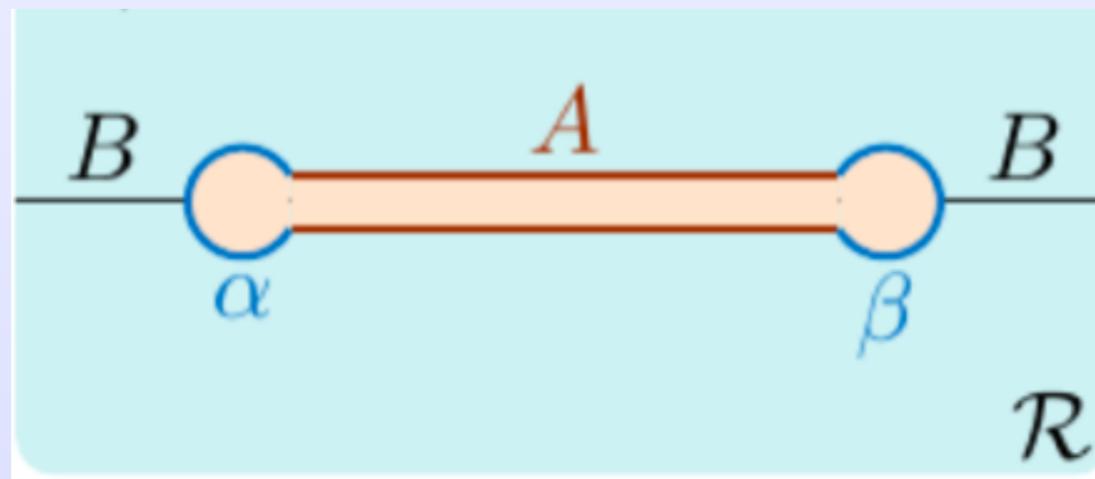
$$\log \text{Tr} \left(e^{-2\pi n \mathcal{H} + 2\pi i \mu Q_{\mathcal{A}}} \right) = \frac{c\ell}{6n} \left(-\frac{1}{3} \frac{\mu^2}{n^4} + \frac{10}{27} \frac{\mu^4}{n^8} + \dots \right)$$

Fourier transformation and taking the replica limit yields breakdown of equipartition in SREE at large c .

Why BCFT?

BCFT corrections to symmetry resolved entanglement

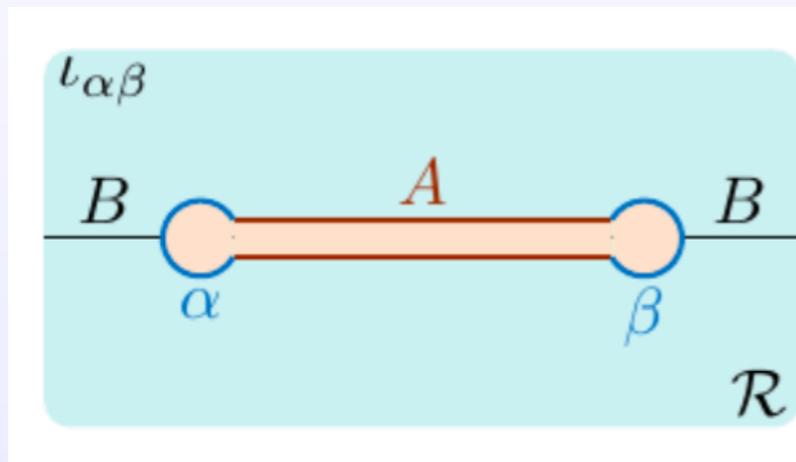
Boundary conditions are important



Exact equipartition of entanglement for free boson to all orders in the UV cutoff expansion

Boundary Conformal Field Theory

Boundary Conditions: Define Hilbert space factorization



Choose boundary conditions which preserve conformal symmetry and $U(1)$

$$T = \bar{T}|_{bdy}$$

$$J = \pm \bar{J}|_{bdy}$$

Dirichlet

Neumann

Example: Compact Free Boson

$$S = \frac{g}{2} \int_{\Sigma} d^2x \partial_{\mu} \varphi \partial^{\mu} \varphi, \quad \varphi = \varphi + 2\pi R$$

$$J^{\mu} = \partial^{\mu} \varphi$$

$$\tilde{J}^{\mu} = \epsilon^{\mu\nu} \partial_{\nu} \varphi$$

NN & DD conditions preserve a diagonal Vir x $U(1)$

ND & DN conditions break both $U(1)$,

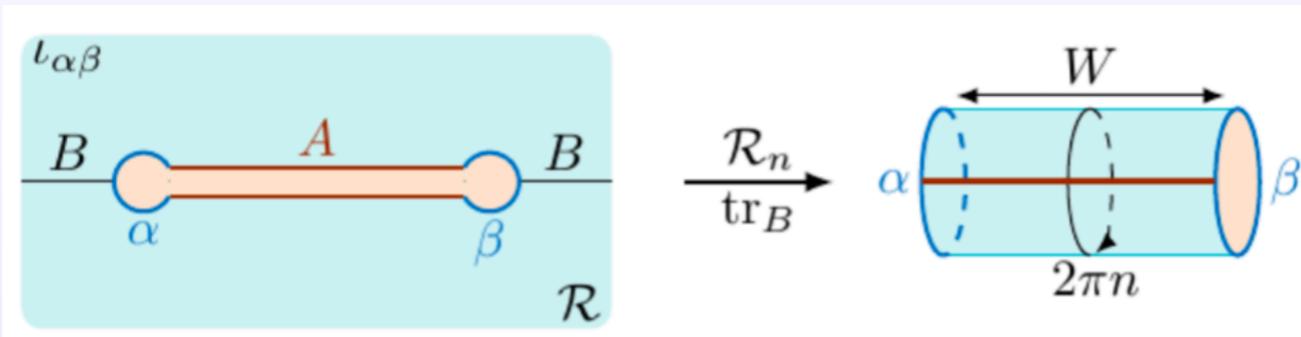
preserve Z_2 & twisted $U(1)$

Ohmori & Tachikawa (2015), Cardy & Tonni (2016)

Di Giulio, RM, Northe, Scheppach, Zhao, SciPost Physics Core, arXiv: 2212.09767

SRE in $U(1)_k$ BCFT

Exact equipartition to all orders in UV cutoff



$$\rho = |0\rangle\langle 0|$$

$$\rho_A = \frac{q^{L_0 - c/24}}{Z_{\alpha\beta}(q)}$$

$$q = e^{-2\pi^2/W}$$

Partition function splits into representations \mathcal{H}_i of chiral algebra with characters $\chi_i(q) = \text{tr}_{\mathcal{H}_i} q^{L_0 - c/24}$

$$Z_{\alpha\beta}(q) = \sum_{i \in \sigma} n_{\alpha\beta}^i \chi_i(q) \approx \chi_{\Omega}(\tilde{q}) \underbrace{\times g_{\alpha} g_{\beta}}_{S_i^{\alpha\beta} = \log[g_{\alpha} g_{\beta}]},$$

$$\chi_Q(q) = \frac{q^{Q^2}}{\eta(q)}$$

Charge-projected partition functions select one representation

$$\mathcal{Z}_n(Q) = \text{tr}_A[\Pi_Q \rho_A^n] = \frac{\chi_Q(q^n)}{(Z_{\pm}(q))^n} \quad p_Q = \frac{\chi_Q(q)}{Z_{\pm}(q)}$$

$$S_n(Q) = \frac{1}{1-n} \log \left[\frac{q^{nQ^2}}{(q^{Q^2})^n} \frac{\eta^n(q)}{\eta(q^n)} \right]$$

$$S_n(Q) = \underbrace{\frac{W}{12} \frac{n+1}{n} - \frac{1}{2} \log \left[\frac{W}{\pi} \right] + \frac{1}{2} \frac{\log n}{1-n}}_{\text{Known terms}} + \frac{1}{1-n} \sum_{k=1}^{\infty} \log \left[\frac{(1 - e^{-2Wk})^n}{1 - e^{-2Wk/n}} \right]$$

Conclusions & Outlook

- New insight into AdS/CFT from quantum information, e.g. Entanglement entropy, **Symmetry resolved entanglement**
- New tests of AdS₃/CFT₂ plus Chern-Simons
- First example of breakdown of equipartition at large c
- BCFT approach yields exact results due to symmetries, U(1) equipartition to all orders for free bosons
- SREE Analysis of further instances of AdS/CFT
- Symmetry resolving other quantum information measures
- Higher dimensions? **P. Bueno, P.A. Cano, A. Murcia, A.R. Sanchez, PRL 2022**
- Other symmetry groups? **Y. Kusuki et. al., arXiv:2309.03287**
- Explore implications on bulk entanglement, in particular AdS/CFT duals of higher order corrections?
- Field theory and Condensed Matter Applications of SRE? **C. Northe, arXiv:2303.07724**