

Going beyond two
- multipartite entanglement and multi inner product -

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YKIS conference “Foundations and Developments of Quantum
Information Theory” 2023/09/22

Introduction

- One of the most fundamental concepts in quantum mechanics is the inner product: $\langle \psi | \phi \rangle$. Represents a transition amplitude.
- It'd also be useful to think it in terms of entanglement:

$$\begin{aligned} (\langle \psi | \otimes \langle \phi |^*) | \Omega \rangle &= \langle \psi | \phi \rangle \\ | \Omega \rangle &= \text{maximally entangled state} \end{aligned}$$

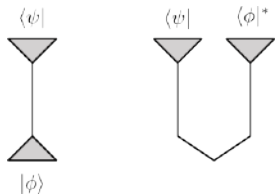
- E.g.

$$|Bell\rangle = (1/\sqrt{2})(| \uparrow \uparrow \rangle + | \downarrow \downarrow \rangle),$$



$$|TFD\rangle = \sqrt{Z}^{-1} \sum_n e^{-\beta E_n/2} |E_n\rangle_L |E_n\rangle_R$$

- Maybe useful to consider in terms of diagrams of diagrams



- Or in terms of geometry;
“ER=EPR” [Maldacena-Susskind (13)]

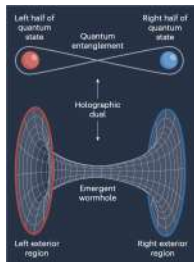
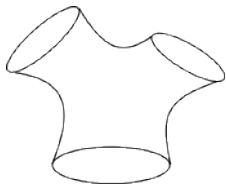


Image: [Brown-Susskind].

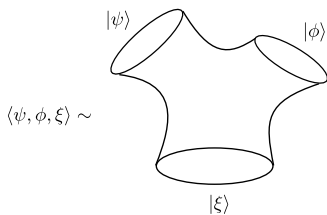
Multiboundary and multipartite entanglement



$$|V\rangle = \sum C_{ijk} |i\rangle |j\rangle |k\rangle$$

E.g., [Balasubramanian-Hayden-Maloney-Marolf-Ross, "Multiboundary Wormholes and Holographic Entanglement," (14)].

... and multi inner product



We go beyond bipartite entanglement and bi inner product.

- (1+1)d critical spin chain (CFT) with [Yuhan Liu and Yijian Zou \(22\)](#)
- (1+1)d gapped system (higher Berry phase) with [Shuhei Ohyama \(23\)](#)
- Tripartite entanglement in (2+1)d topologically-ordered states with [Yuhan Liu, Ramanjit Sohal, Jonah Kudler-Flam, Yuya Kusuki \(21-23\)](#)

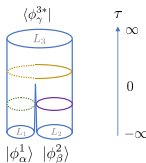
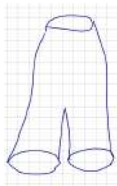
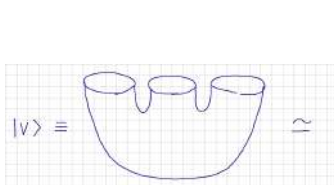
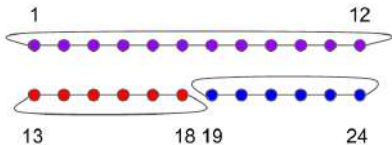
Triple inner product in critical spin chain

- Lattice quantum many-body systems, e.g., $H = -\sum_i X_i X_{i+1} - \sum_i Z_i$,
 $H = \sum (X_i X_{i+1} + Y_i Y_{i+1} + \Delta Z_i Z_{i+1})$.



- Triple wave function overlap $A_{\alpha\beta\gamma}$

$$A_{\alpha\beta\gamma} = \langle \phi_\gamma | (|\phi_\alpha\rangle |\phi_\beta\rangle) \quad [= \langle V | (|\phi_\alpha^1\rangle |\phi_\beta^2\rangle |\phi_\gamma^3\rangle^*)]$$



Numerical approaches to critical “spin” chains

- Exact diagonalization, DMRG, Monte Carlo, etc.
- A quantum spin chain may be at conformal critical point \Rightarrow How do we identify CFT?
 - Central charge \Leftarrow entanglement entropy scaling
 - Dimensions \Leftarrow numerical exact diagonalization
- These may not be enough to nail down full information of CFT
- $A_{\alpha\beta\gamma}$ can be used to extract OPE coefficients (and more) of (1+1)d lattice quantum systems at conformal critical point [Zou-Vidal (21), Liu-Zou-SR (22)]

$$\phi_{\alpha}(z, \bar{z})\phi_{\beta}(w, \bar{w}) = \sum_{\gamma} |z - w|^{-\Delta_{\alpha} - \Delta_{\beta} + \Delta_{\gamma}} C_{\alpha\beta\gamma} \phi_{\gamma}(z, \bar{z}) + \dots$$

Wavefunction overlap and OPE

- Wave function overlap can be mapped on to a 3pt function on a plane:

[Zou-Vidal (21)]

$$\begin{aligned}\frac{A_{\alpha\beta\gamma}}{A_{1111}} &= \langle \phi_{\alpha}^1(-\infty) \phi_{\beta}^2(-\infty) \phi_{\gamma}^3(+\infty) \rangle_{\text{Pants}} \\ &= \text{Jacobian} \cdot \langle \phi_{\alpha}^1(w_1) \phi_{\beta}^2(w_2) \phi_{\gamma}^3(w_3) \rangle_{\mathbb{C}} \propto C_{\alpha\beta\gamma}\end{aligned}$$

- Finite size corrections [Liu-Zou-SR (22)] :

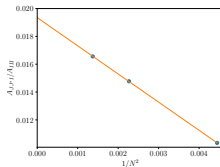
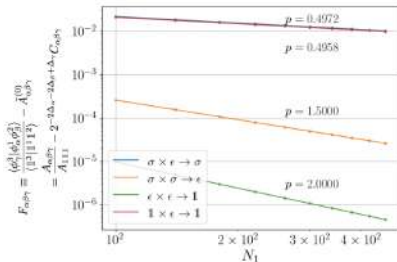
$$A_{\alpha\beta\gamma} = \sum_{\delta, \hat{\chi}=0,1} a_{(\delta, \hat{\chi})} 2^{-2\Delta_{\alpha}-2\Delta_{\beta}+2\Delta_{\gamma}} C_{\alpha\beta\gamma} L^{-\Delta_{(\delta, \hat{\chi})}}$$

$\Delta_{(\delta, \hat{\chi})}$: operator content of the \mathbb{Z}_2 orbifold theory

E.g. $A_{1111} \propto L^{-c/8}$

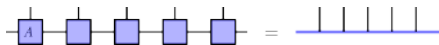
Applications

- Critical Ising chain, XXZ model
- Interacting anyon chain for the Haagerup fusion category ($c \sim 2$)
[\[Huang-Lin-Ohmori-Tachikawa-Tezuka \(21\); Vanhove-Lootens-Damme-Wolf-Osborne-Haegeman-Verstraete \(21\)\]](#)
 Tested a conjectured CFT
- “Building block” for tensor network [\[Ueda-Yamazaki \(23\)\]](#)



Gapped states in $(1+1)d$ and (higher) Berry phase

- Short-range entangled or invertible states in $(1+1)d$
- Is there Berry phase associated with the triple inner product?
- We will use matrix product states (MPS)



[Previous and related works on higher Berry phase:

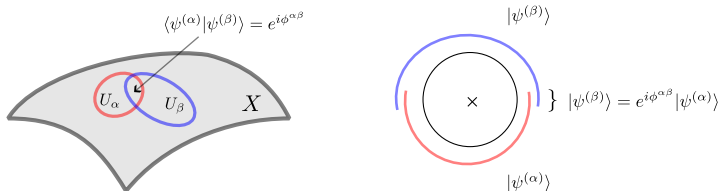
Kitaev (2019); Cordova-Freed-Lam-Seiberg (19); Kapustin-Spodyneiko (20); Kapustin-Sopenko (22);

Hsin-Kapustin-Thorngren (20); Artymowicz-Kapustin-Sopenko (20); Choi-Ohmori (22); Shiozaki (21);

Wen-Qi-Beaudry-Moreno-Pflaum-Spiegel-Vishwanath-Hermele (21); Beaudry-Hermele-Moreno-Pflaum-Qi-Spiegel

(23); Ohyama-Shiozaki-Sato (22); Ohyama-Terashima-Shiozaki (23)]

Regular Berry phase for single-particle wavefunctions



- Parameterized state $|\psi(x)\rangle$, $x = (x^1, x^2, \dots) \in X$.
- Berry connection $\langle \psi(x) | \psi(x + dx) \rangle = \mathcal{A}_\mu(x) dx^\mu$, “associated with” the gauge transformation $|\psi(x)\rangle \rightarrow e^{i\phi(x)} |\psi(x)\rangle$
- Mathematical framework: complex line bundle;
Classified by Chern class, $H^2(X, \mathbb{Z})$
- Wu-Yang description of magnetic monopole
At the double intersection $U_\alpha \cap U_\beta$, $\langle \phi^{(\alpha)} | \phi^{(\beta)} \rangle = e^{i\phi^{\alpha\beta}}$, which defines $[e^{i\phi^{\alpha\beta}}] \in H^1(X, U(1)) \simeq H^2(X, \mathbb{Z})$.

Thouless pump

[Thouless (1983)]

Topological Thouless Pumping of Ultracold Fermions

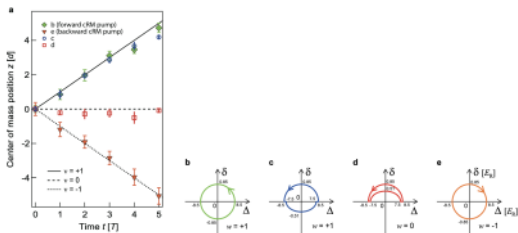
Shuta Nakajima^{1,*}, Takafumi Tomita¹, Shintaro Taie¹, Tomohiro Ichinose¹,

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[Nakajima et al (16)]

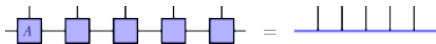


$$H = \sum_i \left[-(J + \delta) f_i^\dagger d_i - (J - \delta) f_i^\dagger d_{i+1} + h.c. + \Delta (f_i^\dagger f_i - d_i^\dagger d_i) \right]$$

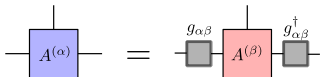
$$\nu = \frac{1}{2\pi} \int_0^T dt \int_{-\pi}^{\pi} dk \Omega(k, t) = \text{integer}$$

MPS and MPS gauge transformation

- Parameterized injective MPS $\{A^i(x)\}$ over some parameter space X .



- MPS gauge transformation, $A \rightarrow gAg^\dagger$: [Perez-Garcia-Verstraete-Wolf-Cirac]

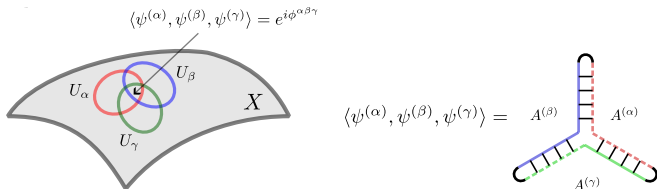


where $g_{\alpha\beta} \in PU(\text{bond dim})$.

- Plays a crucial role in the classification of SPT phases in $(1+1)d$
[Pollmann-Turner-Berg-Oshikawa (10), Chen-Gu-Wen (11), Schuch-Perez-Garcia-Cirac(11)]
- Is there “Berry phase” associated with the MPS gauge transformation?

Main results

[Ohyama-SR (23)]



- Defined and computed triple inner product of three MPSs.
- At triple intersection, $U_\alpha \cap U_\beta \cap U_\gamma$, $\langle \psi^{(\alpha)}, \psi^{(\beta)}, \psi^{(\gamma)} \rangle = e^{i\phi^{\alpha\beta\gamma}}$ where $g_{\alpha\beta}g_{\beta\gamma} = e^{i\phi^{\alpha\beta\gamma}} g_{\alpha\gamma}$.
- $e^{i\phi^{\alpha\beta\gamma}}$ satisfies the cocycle condition, $\delta e^{i\phi^{\alpha\beta\gamma}} = 1$, and defines $[e^{i\phi^{\alpha\beta\gamma}}] \in H^2(X, U(1)) \simeq H^3(X, \mathbb{Z})$
- Mathematical framework: *gerbe* – a generalization of complex line bundle and classified by Dixmier-Douady class, $H^3(X, \mathbb{Z})$ [See also \[Qi-Stephen-Wen-Spiegel-Markus-Pflaum-Beaudry-Hermele \(23\)\]](#)

Some more details

- Transfer matrix and mixed transfer matrix

$$T_\alpha = \begin{array}{c} \boxed{A_\alpha^*} \\ | \\ \boxed{A_\alpha} \end{array} = \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array} \quad T_{\alpha\beta} = \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array}$$

- (Right and left) eigen vector $\Lambda_{\alpha\beta}^{R,L}$:

$$T_{\alpha\beta} \cdot \Lambda_{\alpha\beta}^R = \Lambda_{\alpha\beta}^R \iff \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array} \text{D} = \text{D}$$

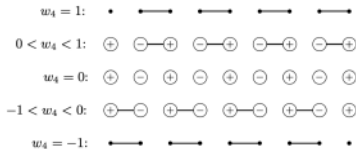
Related to the gauge transformation as $\Lambda_{\alpha\beta}^R = g_{\alpha\beta} \Lambda_\beta^R$, $\Lambda_{\alpha\beta}^L = g_{\alpha\beta} \Lambda_\beta^L$.

- Triple inner product:

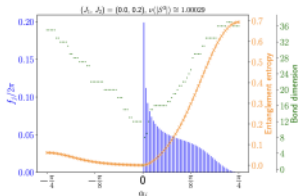
$$\begin{array}{c} A^{(\beta)} \\ | \\ A^{(\alpha)} \\ | \\ A^{(\gamma)} \end{array} = \text{tr} \left(\Lambda_{\beta\alpha}^L \Lambda_{\beta\gamma}^R \Lambda_{\gamma\alpha}^R \right) = \text{tr} \left(\Lambda_\alpha^L \hat{g}_{\alpha\beta} 1_n \hat{g}_{\beta\gamma} 1_n \hat{g}_{\gamma\alpha} \right)$$

“Higher” Thouless pump

- Topological classification of parameterized family of states (“adiabatic process”) Many-body generalization of Thouless pump (“higher Thouless pump”)



- A concrete model parameterized over $X = S^3$: [Wen-Qi-Beaudry-Moreno-Pflaum-Spiegel-Vishwanath-Hermele (21)]
- What about the Berry connection and curvature? Practical method to calculate topological invariant? [Shiozaki-Heinsdorf-Ohyama (23)]



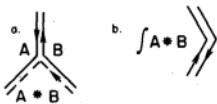
- "Aharonov-Bohm phase" of MPS: 2-form gauge field
- Possible because of "internal structure" of MPS, can go beyond single-particles
- 2-form gauge field couples naturally to 1d extended objects; strings
- No fundamental string in condensed matter physics, but we could have emergent ones.
- More (formal) direct link with string theory (string field theory):

NON-COMMUTATIVE GEOMETRY AND STRING FIELD THEORY

Edward WITTEN*

Joseph Henry Laboratories, Princeton University, Princeton, New Jersey 08540, USA

Received 2 December 1985



String \star product and triple inner product

- Following [Witten (85)], we can introduce “ \star ” product and “integration” for MPSs

$$\Psi_\alpha = \text{---|---|---|---|---|---|---|---|---|---|}$$

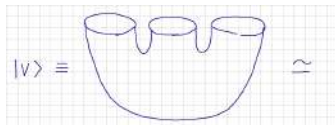
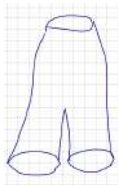
$$\Psi_\alpha * \Psi_\beta = \overbrace{\text{---|---|---|---|}}^{\Psi_\alpha^L} \overbrace{\text{---|---|---|---|}}^{\Psi_\alpha^R} * \overbrace{\text{---|---|---|---|}}^{\Psi_\beta^L} \overbrace{\text{---|---|---|---|}}^{\Psi_\beta^R} = \begin{array}{c} \Psi_\alpha^L \quad \Psi_\beta^R \\ \text{---|---|---|---|} \\ \Psi_\alpha^R \quad \Psi_\beta^L \\ \text{---|---|---|---|} \\ \text{---|---|---|---|} \end{array}$$

$$\int \Psi_\alpha = \int \text{---|---|---|---|---|---|---|---|---|---|} = \text{---|---|---|---|---|---|---|---|---|---|}$$

- Triple inner product is given by “string field theory vertex”

$$\int \Psi_\alpha * \Psi_\beta * \Psi_\gamma = \begin{array}{c} \text{---|---|---|---|---|---|---|---|---|---|} \\ \text{---|---|---|---|---|---|---|---|---|---|} \\ \text{---|---|---|---|---|---|---|---|---|---|} \end{array} = \begin{array}{c} \text{---|---|---|---|---|---|---|---|---|---|} \\ \text{---|---|---|---|---|---|---|---|---|---|} \\ \text{---|---|---|---|---|---|---|---|---|---|} \end{array}$$

Multipartite entanglement?



- $|V\rangle \in \mathcal{H}_{\text{chiralCFT}}^{\otimes 3}$
- What “type” of tripartite states?

$$|\text{GHZ}\rangle = \frac{1}{\sqrt{2}} [|\uparrow\uparrow\uparrow\rangle + |\downarrow\downarrow\downarrow\rangle], \quad |\text{W}\rangle = \frac{1}{\sqrt{3}} [|\uparrow\uparrow\downarrow\rangle + |\uparrow\downarrow\uparrow\rangle + |\downarrow\uparrow\uparrow\rangle]$$

“Three qubits can be entangled in two inequivalent ways” [Dür-Vidal-Cirac (00)]

- Reflected entropy $R_{A:B}$ and Markov gap $h_{A:B} = R_{A:B} - MI_{A:B}$
[Dutta-Faulkner (19)][Akers-Rath (19)][Hayden-Parrikar-Sorce (21)]

Reflected entropy in 2d topological liquids



- By bulk-boundary correspondence, the problem turns out to be related to tripartite entanglement of $(2+1)d$ topological liquids
C.f. Topological entanglement entropy
- Chiral topological liquids with chiral central charge c :

$$h_{A:B} = R_{A:B} - I_{A:B} = \frac{c}{3} \ln 2, \quad c: \text{ central charge.}$$

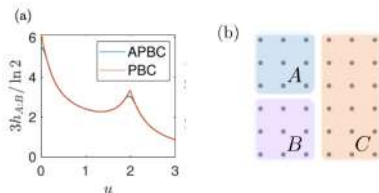
[Liu-Sohal-Kudler-Flam-SR (21)], [Liu-Kusuki-Sohal-Kudler-Flam-SR (23)]

C.f. [Zou-Siva-Soejima-Mong-Zaletel (21)]

- Conjecture [Zou-Siva-Soejima-Mong-Zaletel (21)] :

$$h_{A:B} \leq (c/3) \ln 2$$

where c is the central charge of the total ungappable degrees of freedom



Lattice Chern insulator calculation [Liu et al (21)]

- May have an implication on numerics; non-zero $h_{A:B}$ may be an obstruction to have a finite dim PEPS representation
C.f. diverging correlation length of chiral PEPS

Summary/Outlook

- Introduced generalized inner product (multi-inner product)
- Applications in critical/gapped 1d many-body systems
- Multipartite entanglement of 2d topological liquids

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