Krylov complexity and chaos in quantum mechanics



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Krylov complexity and chaos in quantum mechanics

Classical & quantum chaos in billiard systems

We find significant correlations between indicators of chaos



These correlations are universal for any billiards (stadium, Sinai, ...)
 Subtlety in Krylov complexity value?



Classical chaos in billiard system

- Classical chaos in non-linear deterministic dynamical system
 - Orbits sensitive to initial conditions \rightarrow "information loss"
 - Lyapunov exponent $\lambda > 0 \Rightarrow$ chaotic



Example: Stadium billiard [Bunimovich '74]



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Quantum chaos in billiard system

Schrodinger eq. = linear equation of wave function Quantum effect washes out small scale $(\Delta x \Delta p \ge \hbar)$

 \rightarrow No chaos in quantum regime?

- Probes of chaos in quantum billiard
 - Energy level statistics



Statistics of adjacent energy level spacing

$$s_n = E_{n+1} - E_n$$

• Poisson statistics (→ non-chaotic)

 $P(s) \sim e^{-s}$

• Wigner-Dyson statistics (→ chaotic)

$$P(s) \sim s^{\beta} \exp(-s^2)$$

Out-of-time-order correlators (OTOC) in quantum billiard [Hashimoto+ '17]



Krylov operator complexity

- Yet another indicator of quantum chaos: Krylov complexity
 - Krylov operator complexity for $\mathcal{O}(t) = e^{iHt} \mathcal{O}(0) e^{-iHt}$
 - Krylov state complexity for $|\psi(t)\rangle = e^{-iHt}|\psi(0)\rangle$
 - Also variation of Lanczos coefficients can be an indicator of chaos [Rabinovici+ '21]
- Can we use it to probe quantum chaos in billiard system?
 - Comparison with level statistics
- Relationship with *classical chaos*?
 - Comparison with classical Lyapunov exponent

[Parker+ '18]



Krylov operator/state complexity

Krylov operator complexity

Time evolution of an operator \mathcal{O} for Hamiltonian *H*

$$\mathcal{O}(t) = e^{iHt} \mathcal{O}(0) e^{-iHt} = \sum_{n=0}^{\infty} \frac{(it)^n}{n!} \mathcal{L}^n \mathcal{O}(0) \qquad (\mathcal{L} \equiv [H, \cdot])$$

• Krylov space $\mathcal{H}_{\mathcal{O}} \equiv \{\mathcal{O}, \mathcal{L}\mathcal{O}, \mathcal{L}^2\mathcal{O}, ...\}$

► Introduce an inner product by $(\mathcal{O}_1 | \mathcal{O}_2) \equiv \operatorname{Tr} \left[\mathcal{O}_1^{\dagger} \mathcal{O}_2 \right] \leftarrow \left(\begin{array}{c} \text{cf. Wightman inner product} \\ (\mathcal{O}_1 | \mathcal{O}_2) \simeq \operatorname{Tr} \left[e^{-\frac{\beta H}{2}} \mathcal{O}_1^{\dagger} e^{-\frac{\beta H}{2}} \mathcal{O}_2 \right] \right)$

Switch to orthonormal basis $\mathcal{H}_{\mathcal{O}} = \{\mathcal{O}_0 = \mathcal{O}(0), \mathcal{O}_1, \mathcal{O}_2, ...\}$ by Lanczos algorithm $\mathcal{O}(t) = \sum_{n=0}^{\infty} i^n \varphi_n(t) \mathcal{O}_n$

• Krylov operator complexity $C_{\mathcal{O}}(t) \equiv \sum_{n=0}^{\infty} n |\varphi_n(t)|^2$



Krylov operator/state complexity

Krylov operator complexity

► Lanczos algorithm → Orthonormal basis $\{\mathcal{O}_0, \mathcal{O}_1, \mathcal{O}_2, ...\}$ & Lanczos coeff. b_n

- 1. $b_0 \equiv 0$, $\mathcal{O}_{-1} \equiv 0$
- 2. $\mathcal{O}_0 \equiv \mathcal{O}/\|\mathcal{O}\|$, where $\|\mathcal{O}\| \equiv \sqrt{(\mathcal{O}|\mathcal{O})}$
- 3. For $n \ge 1$: $\mathcal{A}_n = \mathcal{LO}_{n-1} b_{n-1}\mathcal{O}_{n-2}$
- 4. Set $b_n = \|\mathcal{A}_n\|$
- 5. If $b_n = 0$ stop; otherwise set $\mathcal{O}_n = \mathcal{A}_n/b_n$ and go to step 3.

• Krylov complexity $C_{\mathcal{O}}(t) \equiv \sum_{n=0}^{\infty} n |\varphi_n(t)|^2 \approx \text{Distance from initial } \mathcal{O}(0)$



Krylov operator/state complexity

Krylov state complexity

Time evolution of a state $|\psi\rangle$ for Hamiltonian *H*

$$\psi(t)\rangle = e^{-iHt}|\psi(0)\rangle = \sum_{n=0}^{\infty} \frac{(-it)^n}{n!} H^n|\psi(0)\rangle$$

• Krylov space $\mathcal{H}_{|\psi(t)\rangle} \equiv \{|\psi(0)\rangle, H|\psi(0)\rangle, H^2|\psi(0)\rangle, ...\}$

• Orthonormal basis $\mathcal{H}_{|\psi(t)\rangle} = \{|K_0\rangle = |\psi(0)\rangle, |K_1\rangle, |K_2\rangle, ...\}$ by Lanczos algorithm

$$|\psi(t)\rangle = \sum_{n=0}^{\infty} \psi_n(t) |K_n\rangle \rightarrow i \dot{\psi}_n(t) = a_n \psi_n(t) + b_n \psi_{n-1}(t) + b_{n+1} \psi_{n+1}(t)$$

$$\psi_0(t) \psi_1(t) \psi_2(t) \psi_3(t) \psi_4(t)$$

$$\dots$$

$$a_0 b_1 a_1 b_2 a_2 b_3 a_3 b_4 a_4$$

• Krylov state complexity $C_{|\psi(t)\rangle}(t) \equiv \sum_{n=0}^{\infty} n |\psi_n(t)|^2 \approx \text{Distance from initial } |\psi\rangle$



Krylov complexity and chaos in quantum mechanics

Setup: stadium billiard (with unit area & velocity)



- Classical billiard \rightarrow Positive Lyapunov exponent λ
- Quantum billiard
 - Hamiltonian $H = p_x^2 + p_y^2$ with Dirichlet BC at the wall: $\psi(x, y)|_{wall} = 0$
 - Calculate energy level statistics & Krylov operator/state complexity [cf. Finite temperature case: Camargo+ '23]



Classical / quantum chaos in stadium billiard

Classical Lyapunov exponent λ

 λ maximized at $a/R \sim 1$

Energy level statistics

$$\tilde{r}_n \equiv \frac{\min(s_n, s_{n-1})}{\max(s_n, s_{n-1})} \qquad (s_n = E_{n+1} - E_n)$$
$$\langle \tilde{r} \rangle \approx \begin{cases} 0.38629 \quad (\text{Poisson})\\ 0.53590 \quad (\text{GOE}) \end{cases}$$

Poisson (non-chaotic) \rightarrow GOE (chaotic)







Classical / quantum chaos in stadium billiard

- **Krylov complexity for operator** p_x on stadium billiard
 - Solve Schrödinger eq. to obtain eigenvalues E_n & eigenfunctions $\phi_n(x, y)$ $H|n\rangle = E_n|n\rangle, \quad \phi_n(x, y) = \langle x, y|n\rangle \quad (n = 1, 2, ...)$

• Matrix representation of operators x and p_x

$$x_{mn} \equiv \langle m \mid x \mid n \rangle = \int dx \, dy \, \phi_m^*(x, y) \, x \, \phi_n(x, y)$$

$$P_{mn} \equiv \langle m \mid p_x \mid n \rangle = \left\langle m \mid \frac{l}{2} [H, x] \mid n \right\rangle = \frac{l}{2} (E_m - E_n) x_{mn}$$

▶ For numerical calculation, we truncate the energy levels at n = N_{max} = 100
 ▶ Lanczos algorithm → Lanczos coefficients b_n & Krylov complexity C_{px}(t)





a/R

• 0

 10^{4}

Linear growth of b_n at low *n* not obvious (Finite temperature $\rightarrow b_n \leq \pi T \times n$ [Camargo+ '23])



Classical / quantum chaos



Classical / quantum chaos in stadium billiard

• Krylov complexity for operator p_x on stadium billiard Krylov complexity $C_{p_x}(t)$



- No exponential growth of $C_{p_x}(t)$ at early time
- Late-time value correlated with classical λ ? (turn out to be system dependent)





- These correlations are universal for any billiards (stadium, Sinai, ...), for both Krylov operator & state complexity
- Non-universality in time dependence of Krylov complexity?
- Other quantum mechanical systems? [Camargo+'23, ...]
- Implication to holography? [Rabinovici+ '23, ...]