

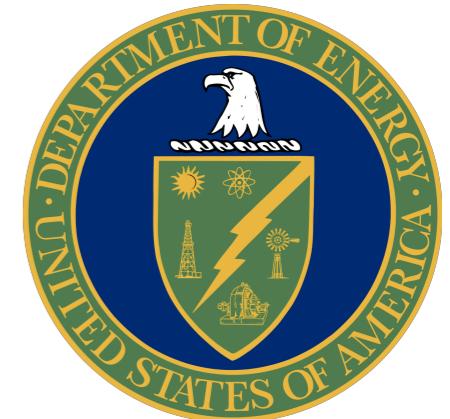
Reflected entropy in AdS/CFT



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Based on: SD & Thomas Faulkner 1905.00577



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Motivation

- Football Field theory encoded as (local geometric objects in?) dual gravity theory

$$S(A) = -\text{Tr}(\rho_A \ln \rho_A) = \frac{\text{Area}(m_A)}{4G_N} + \mathcal{O}(G_N^0)$$

[Ryu, Takayanagi `08]

[Faulkner, Lewkowycz, Maldacena `13]

[Lewkowycz, Maldacena, `13]

- Football HRT surface (time dep.) = Entanglement entropy [Hubeny, Rangamani, Takayanagi `07]

Cross-section of Wheeler-De Witt patch = Complexity? [Stanford, Susskind`14]

Max. vol time slice = Information metric [Miyaji, Numasawa, Shiba, Takayanagi, Watanabe`15]

Entanglement wedge = density matrix [Czech, Karczmarek, Nogueira, van Raamsdonk `12]

Minimal cross-section of entanglement wedge = ? Masamichi, Koji's talks

- Tennis ball Entropy of purification: $E_W(A : B) = E_P(A : B)$

[Takayanagi, Umemoto, `17]

- Tennis ball Reflected entropy: $E_W(A : B) = \frac{1}{2}S_R(A : B)$

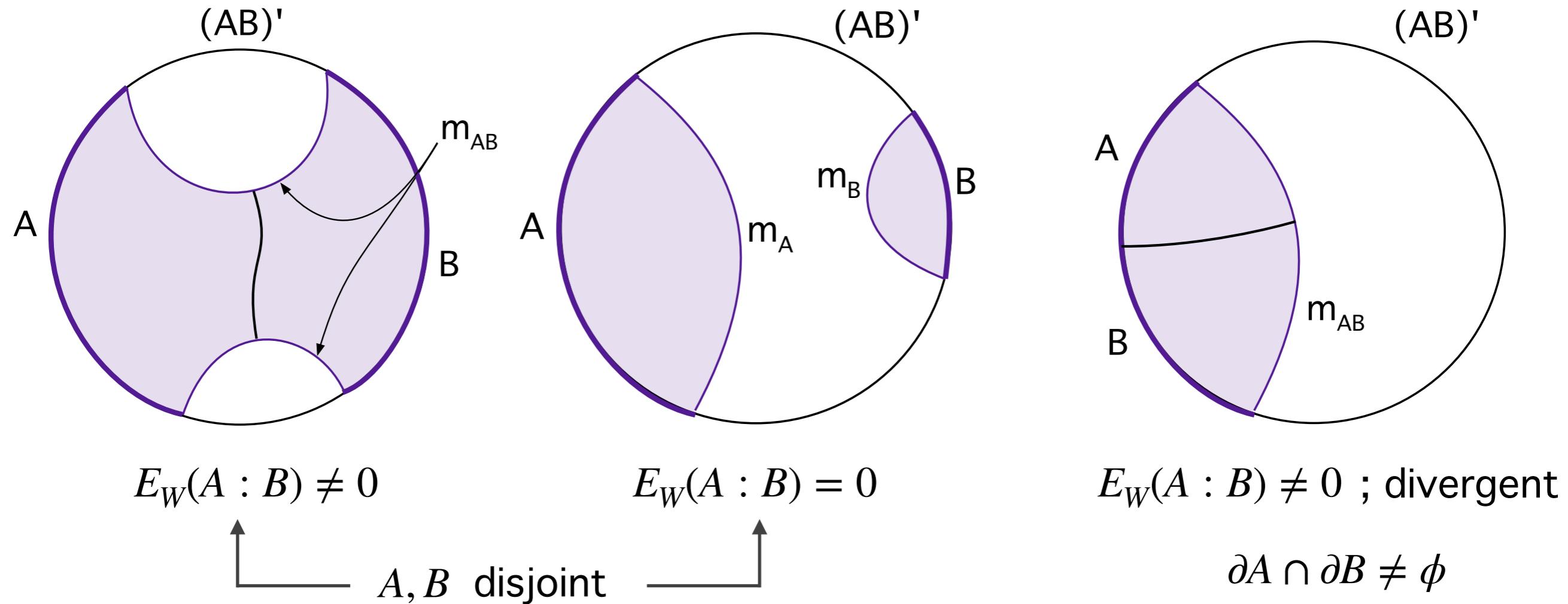
[SD, Faulkner, `19]

Slicing the entanglement wedge

- Set of bulk points spacelike separated from m_{AB} and on the side of AB .

[Czech, Karczmarek, Nogueira, van Raamsdonk `12]

[Headrick, Hubeny, Lawrence, Rangamani `14]



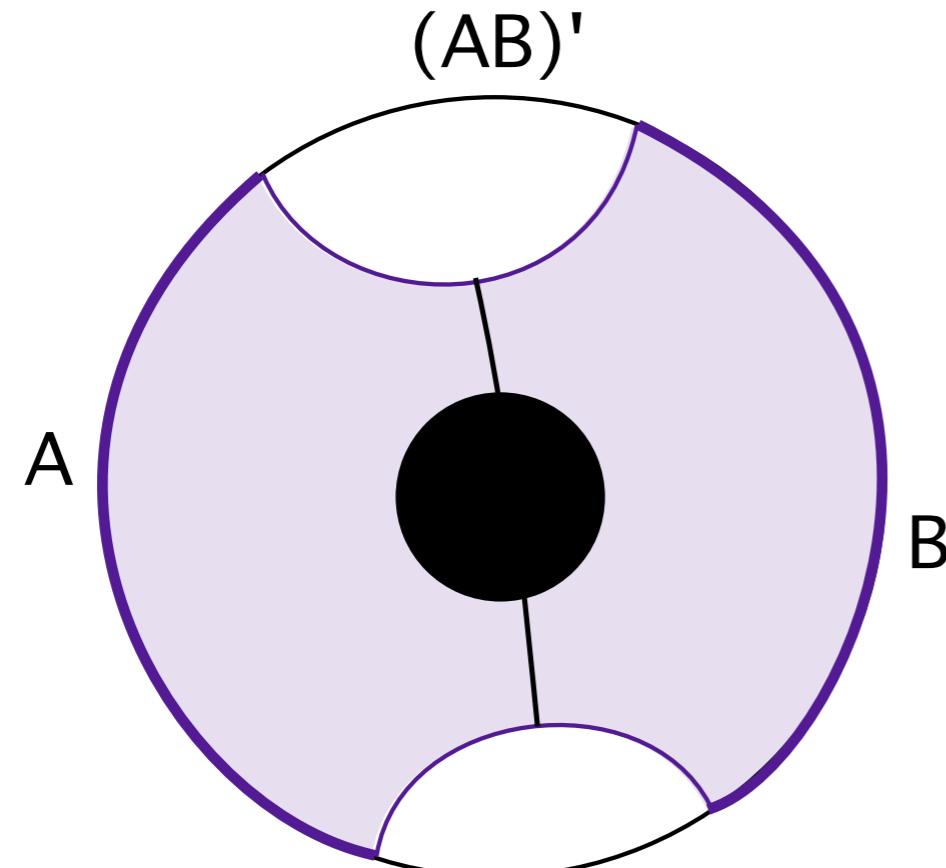
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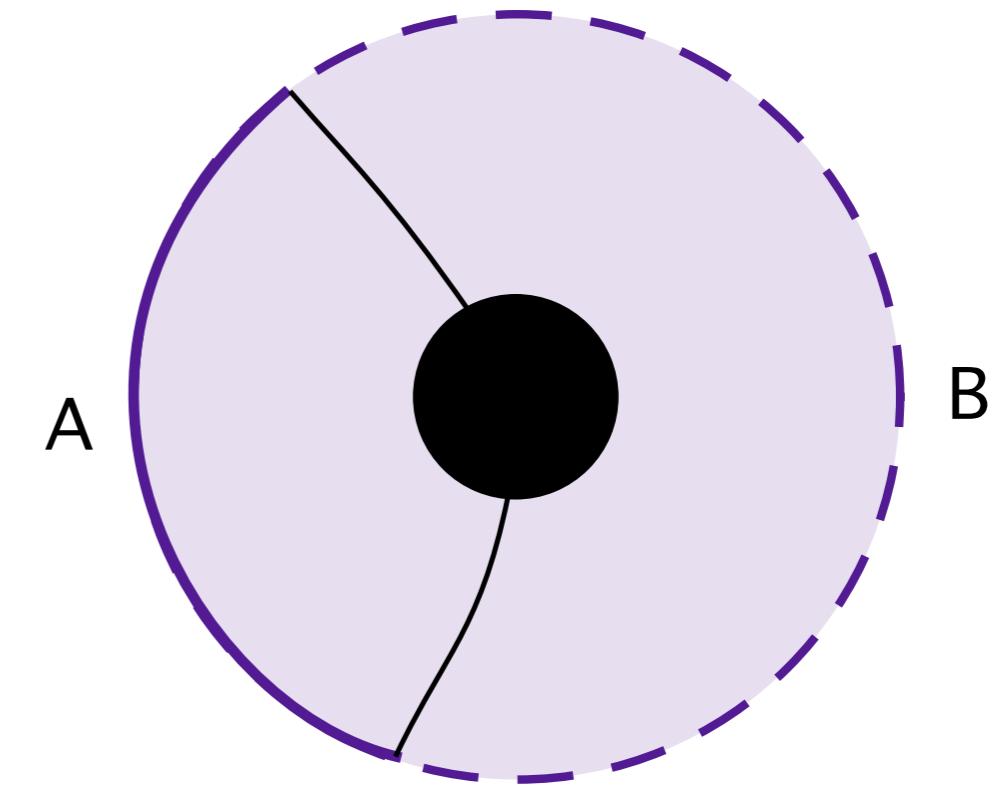
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$(AB)' \neq 0 ; E_W(A : B)$ finite



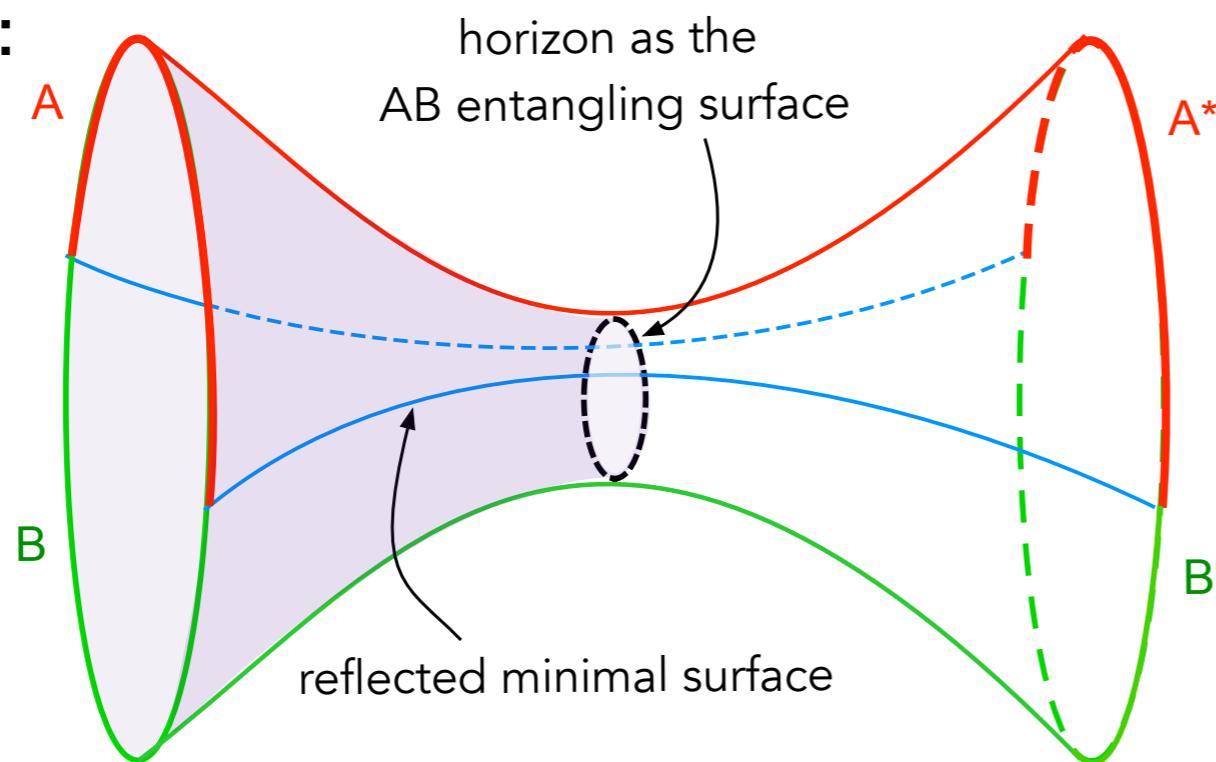
$(AB)' = 0 ; E_W(A : B)$ divergent

Thermal state of a CFT; bulk dual is an AdS black hole.

Warm-up example

$$E_W(A : B) = \frac{1}{2} S(AA^*)_{TFD} \equiv \frac{1}{2} S_R(A : B)$$

Example:



$$\rho_{AB} = e^{-\beta H}$$

$$|TFD\rangle = \sum_i e^{-\beta E_i/2} |E_i\rangle_L \otimes |E_i\rangle_R$$

[Maldacena '01]

Mixed state is a thermal state in AB; bulk dual to $|TFD\rangle$ is eternal black hole.

The duality proposal

- For ANY mixed state ρ_{AB} ,

$$E_W(A : B) = \frac{1}{2} S(AA^\star) \sqrt{\rho_{AB}} \equiv \frac{1}{2} S_R(A : B)$$

$\sqrt{\rho_{AB}}$: canonical purification

- $\rho_{AB}^{1/2} \in \mathcal{H}_A \otimes \mathcal{H}_B \longleftrightarrow |\sqrt{\rho_{AB}}\rangle \in \underbrace{\text{End}(\mathcal{H}_A) \otimes \text{End}(\mathcal{H}_B)}$

- space of linear maps/matrices acting on \mathcal{H}_A
 - forms a Hilbert space, $\langle \sigma_A | \sigma'_A \rangle = \text{Tr}_A \sigma_A^\dagger \sigma'_A$
 - $\text{End}(\mathcal{H}_A) \cong \mathcal{H}_A \otimes \mathcal{H}_A^\star$
- $\Rightarrow |\sqrt{\rho_{AB}}\rangle \in (\mathcal{H}_A \otimes \mathcal{H}_A^\star) \otimes (\mathcal{H}_B \otimes \mathcal{H}_B^\star)$

$$\rho_{AB} = e^{-\beta H}$$

$$\sum_i e^{-\beta E_i/2} |E_i\rangle_L \otimes |E_i\rangle_R$$

such that $\text{Tr}_{A^\star B^\star} |\sqrt{\rho_{AB}}\rangle \langle \sqrt{\rho_{AB}}| = \rho_{AB}$

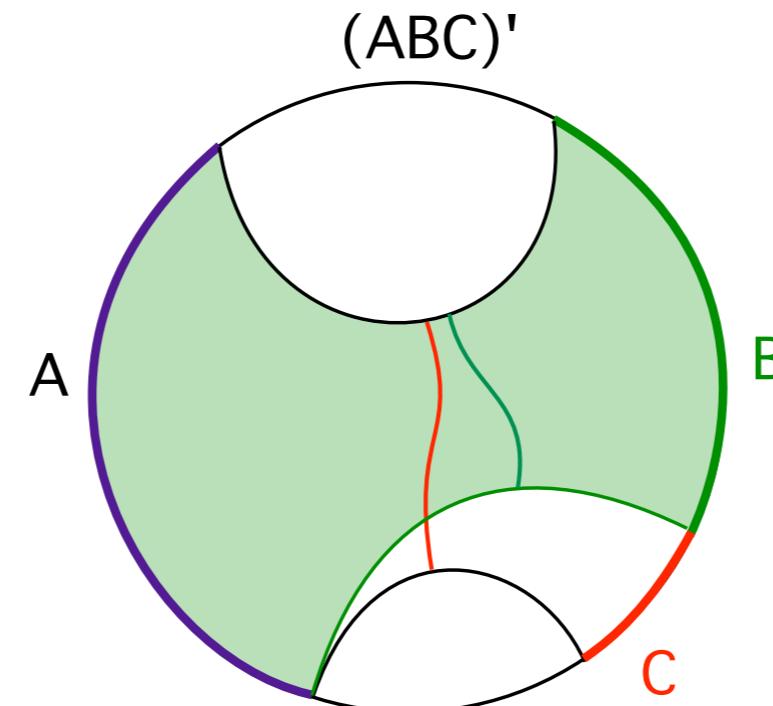
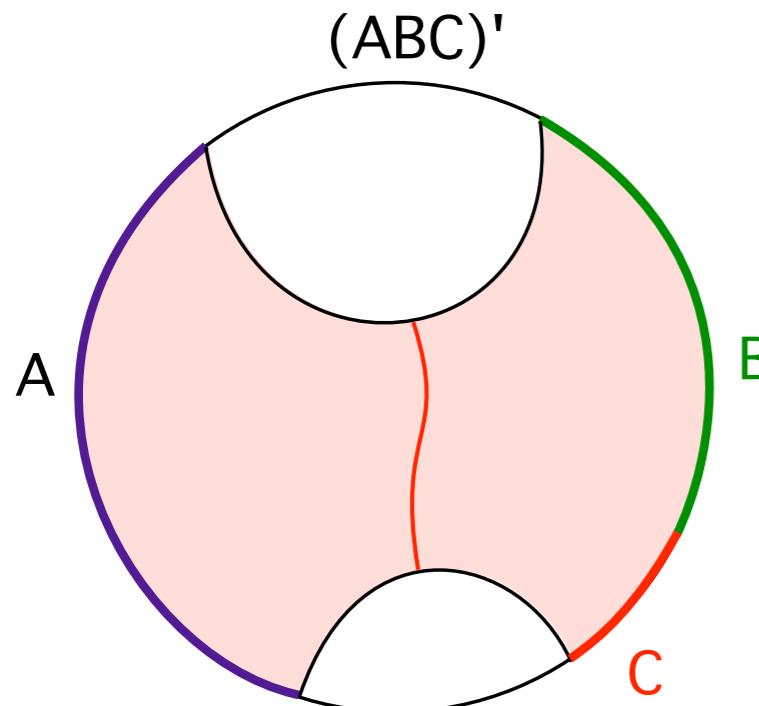
Reflected entropy: properties

$$S_R(A : B) = S(AA^\star)_{\sqrt{\rho_{AB}}}$$

- ⚽ If ρ_{AB} is pure: $S_R(A : B) = S(A) + S(A^\star) = 2S(A)$
 - ⚽ If ρ_{AB} is factorized: $\sqrt{\rho_{AB}} = \sqrt{\rho_A} \otimes \sqrt{\rho_B} \implies S_R(A : B) = 0$
 - ⚽ Measures entanglement and classical correlations.
- $$0 \leq I(A : B) \leq S_R(A : B) \leq 2 \min [S(A), S(B)]$$
- ⚽ E wedge inequality: $S_R(A : BC) \geq I(A : B) + I(A : C)$ [Hayden, Headrick, Maloney `11]
 - ⚽ Monotonicity (?): $S_R(A : BC) \geq S_R(A : B)$

Reflected entropy: properties

$$S_R(A : B) = S(AA^\star)_{\sqrt{\rho_{AB}}}$$



Expected from

- Ent wedge nesting
- $E_W(A : BC) \geq E_W(A : B)$

[Chen, Dong, Lewkowycz, Qi `18]

[Faulkner, Li, Wang `17]

Monotonicity (?):

$$S_R(A : BC) \geq S_R(A : B)$$

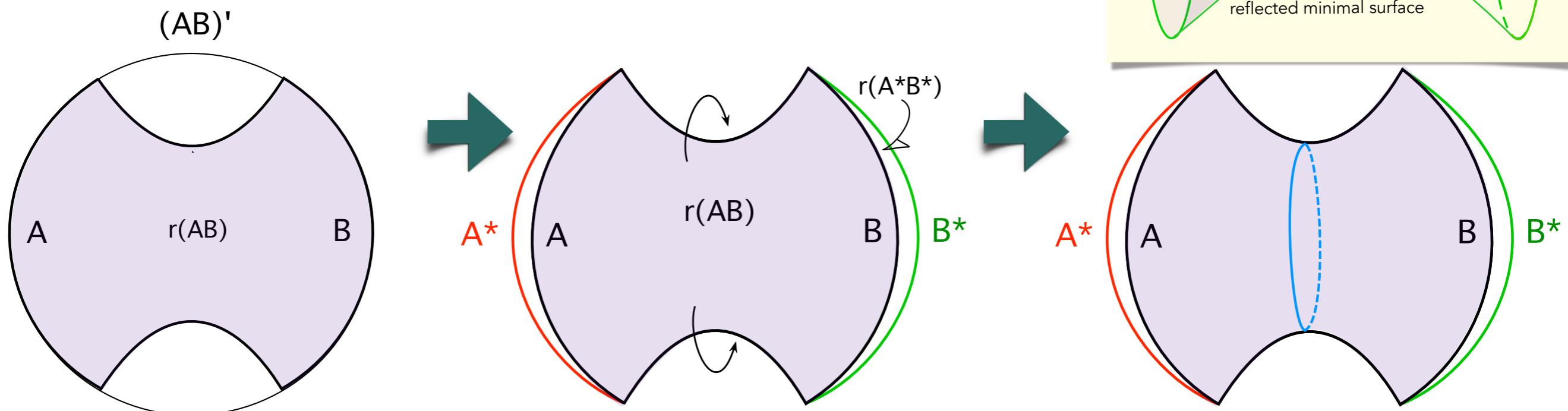
Rènyi version ✓

$$n \rightarrow 1 ?$$

Reflected minimal surfaces

$$S_R(A : B) = \frac{\text{Ar}(m_{AA^*})}{4G_N} + \dots$$

m_{AA^*} : Bulk surfaces that compute the reflected entropy

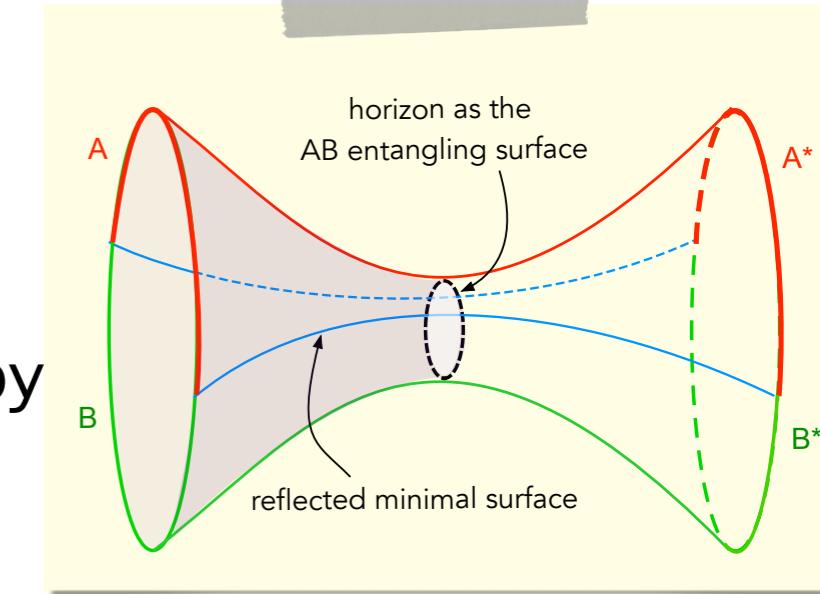


$$\partial A \cap \partial B = \phi$$

Hilbert space split as
A, B and $(AB)'$

$r(AB) \cup r(A^*B^*)$
(glued along m_{AB})
dual to $\sqrt{\rho_{AB}}$

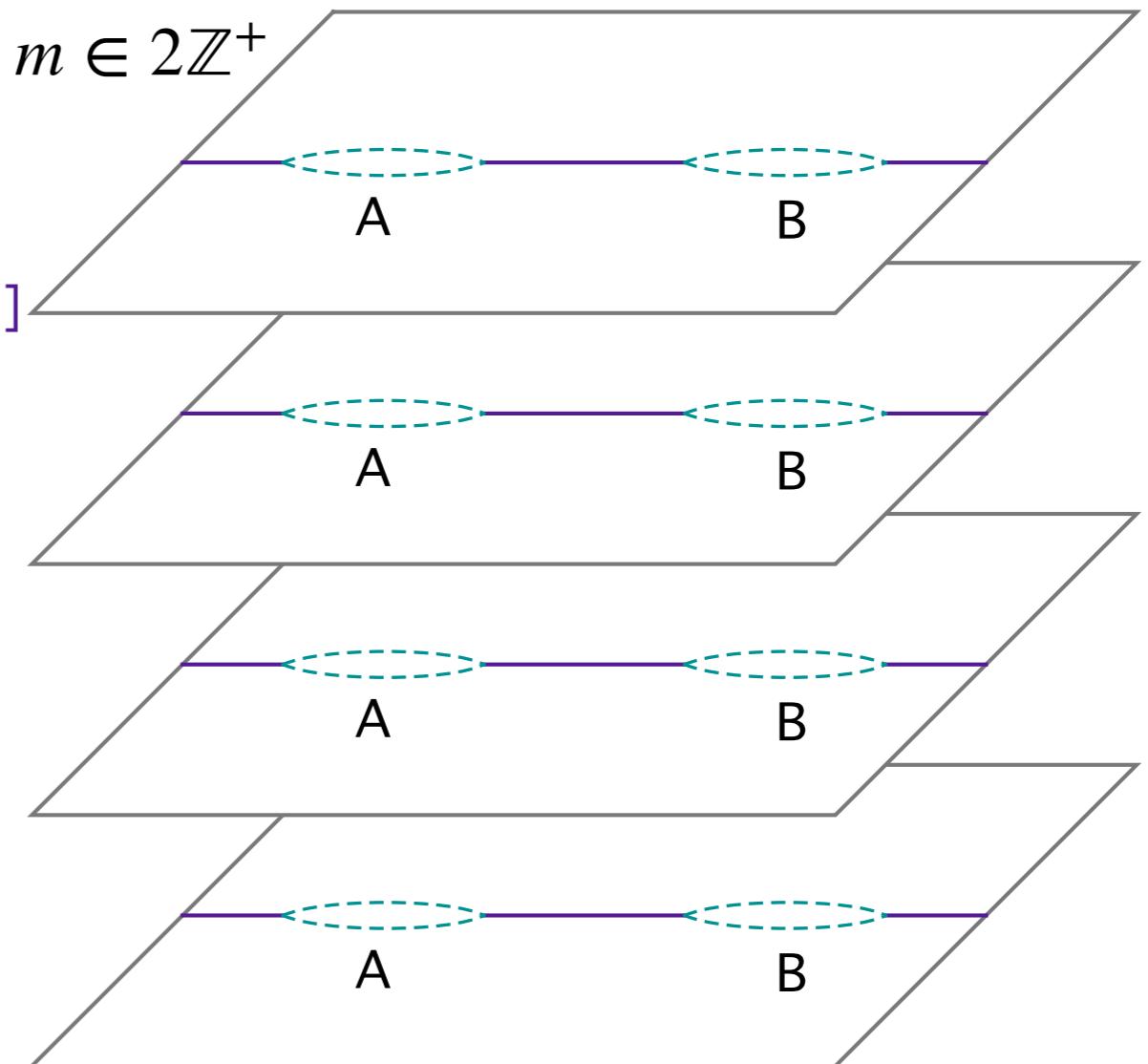
m_{AA^*} is non-contractible
cycle in glued spacetime



Replica trick for reflected entropy

Construct canonical purification $|\rho_{AB}^{1/2}\rangle$, and the corresponding bulk dual.

- Path integral description for $|\rho_{AB}^{m/2}\rangle$, $m \in 2\mathbb{Z}^+$
- by “splitting” usual replica trick for calculating $\text{Tr } \rho_{AB}^m$ [Calabrese, Cardy, ‘05]

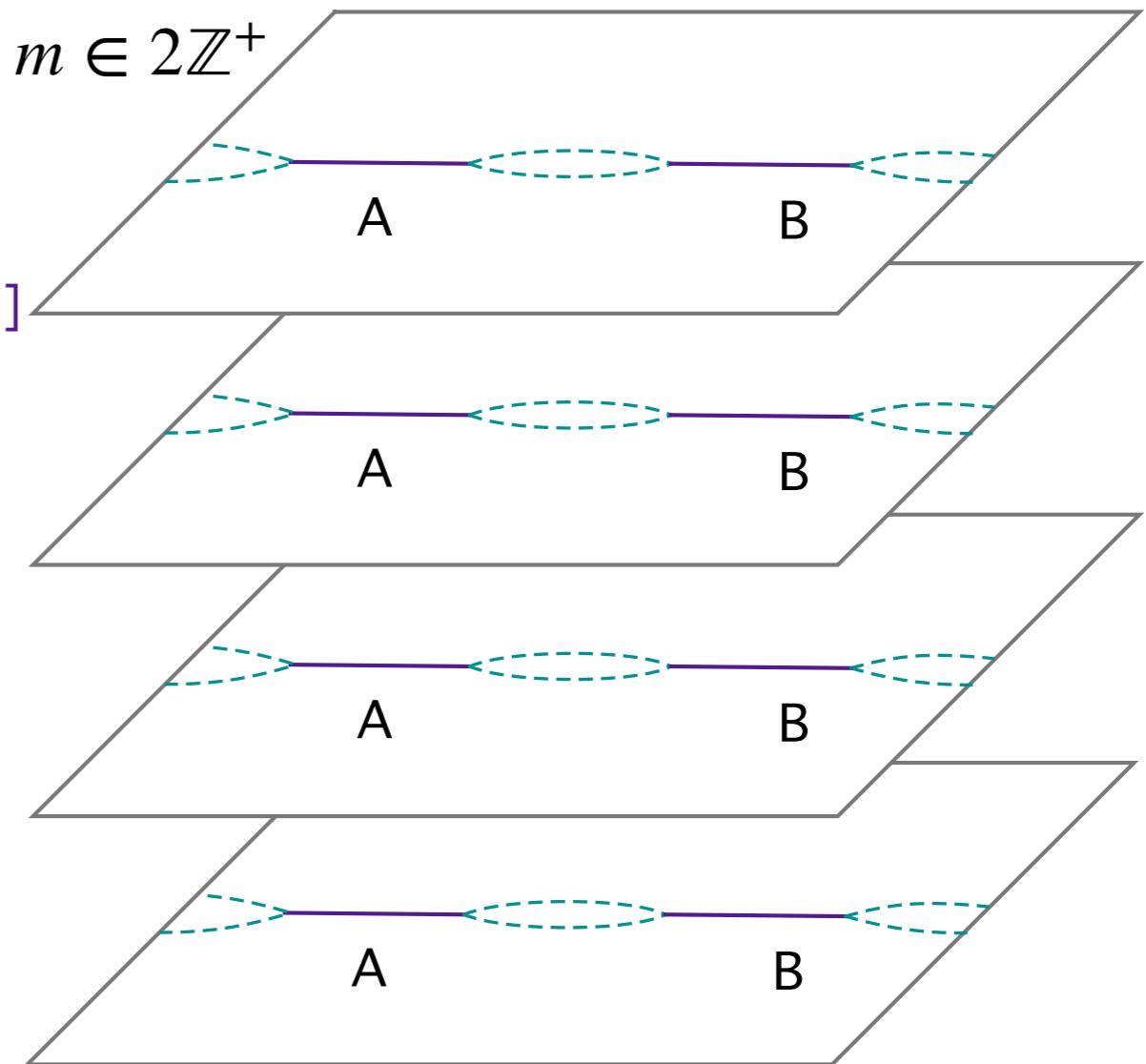


$$m = 4$$

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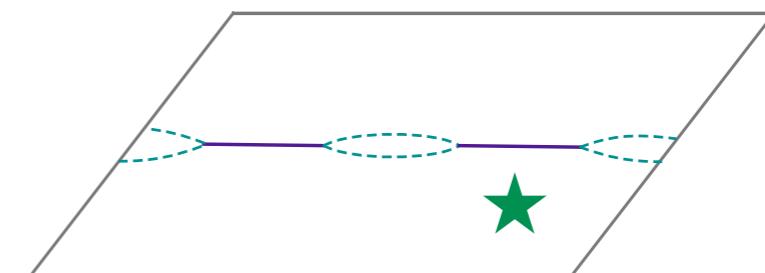
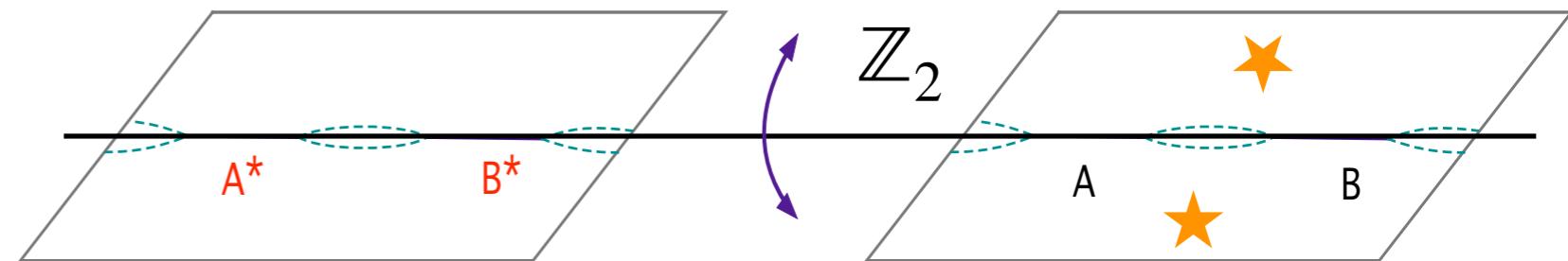
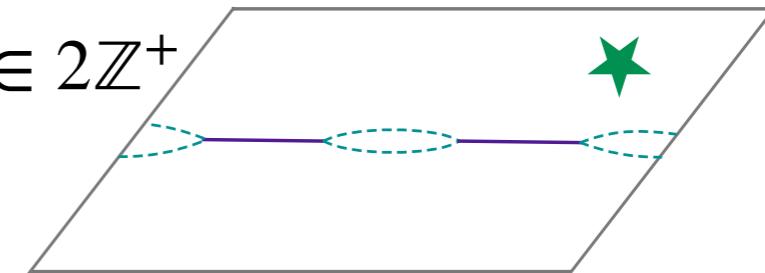


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- Global \mathbb{Z}_2



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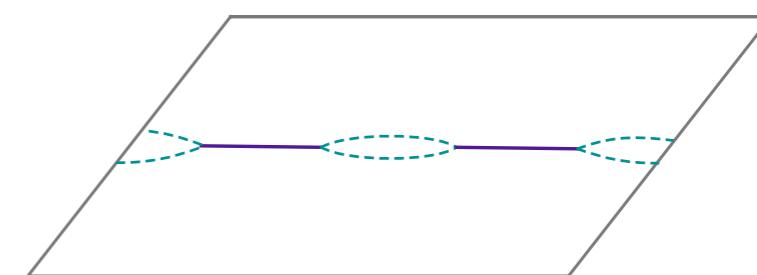
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Can treat the Euclidean P.I. on lower half as $|\rho_{AB}^{m/2}\rangle$

$$\rho_{AA^*BB^*}^{(m)} = \frac{1}{\text{Tr} \rho_{AB}^m} |\rho_{AB}^{m/2}\rangle \langle \rho_{AB}^{m/2}|$$



$m = 4$

Rènyi reflected entropy

- To find the EE of AA^* in the state $\rho_{AA^*BB^*}^{(m)}$, we first trace over BB^* and then compute the Rènyi entropy for replica index $n \in \mathbb{Z}^+$

$$S_n(AA^*)_{\rho_{AB}^{m/2}} = \frac{1}{n-1} \ln \text{Tr} \left(\rho_{AA^*}^{(m)} \right)^n$$

- In 2 dimensions, 4-point function of twist operators

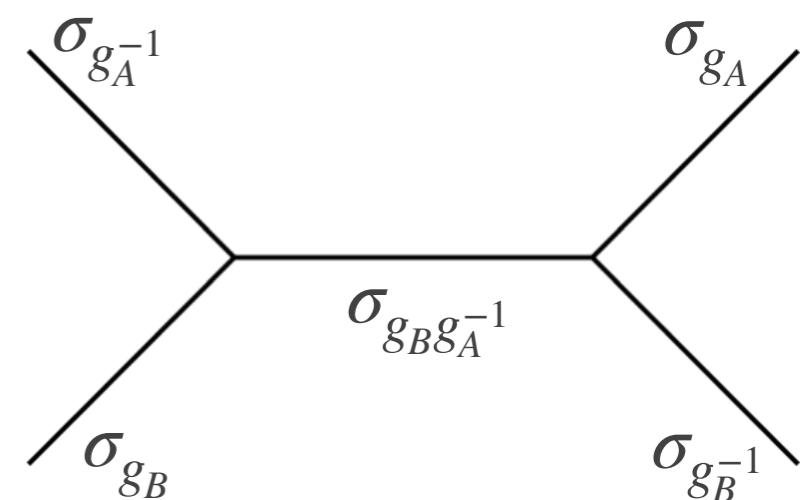
$$S_n(AA^*)_{\rho_{AB}^{m/2}} \propto \ln \left\langle \sigma_{g_A}(a_1) \sigma_{g_A^{-1}}(a_2) \sigma_{g_B}(b_1) \sigma_{g_B^{-1}}(b_2) \right\rangle$$

$$h_{g_A} = h_{g_B} = \frac{nc}{24} \left(m - \frac{1}{m} \right)$$

- For large-c CFTs with sparse spectrum

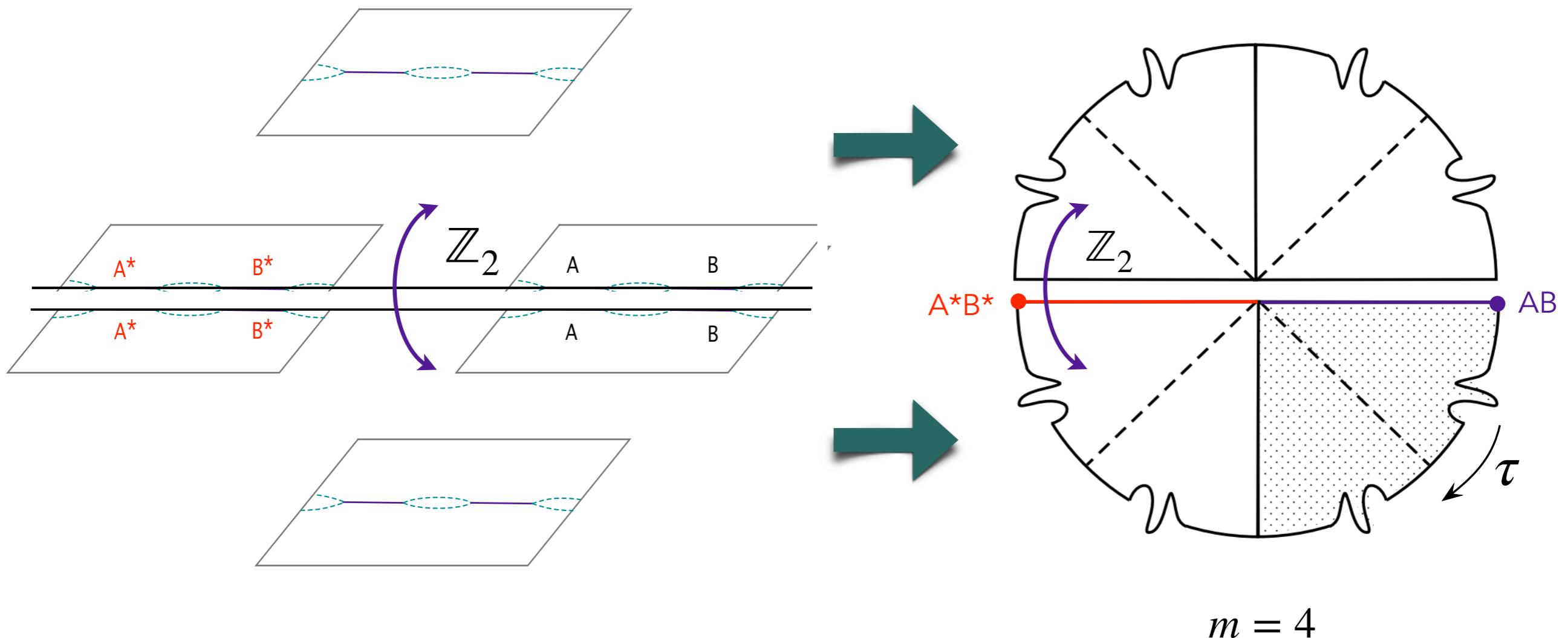
$$\lim_{\{m,n\} \rightarrow 1} S_n(AA^*) \approx \frac{2c}{3} \log \left(\frac{1 + \sqrt{1-x}}{\sqrt{x}} \right)$$

✓ Cross-section using AdS_3



Establishing $2E_W(A : B) = S_R(A : B)$

- Fill in bulk à la [Lewkowycz, Maldacena `13]

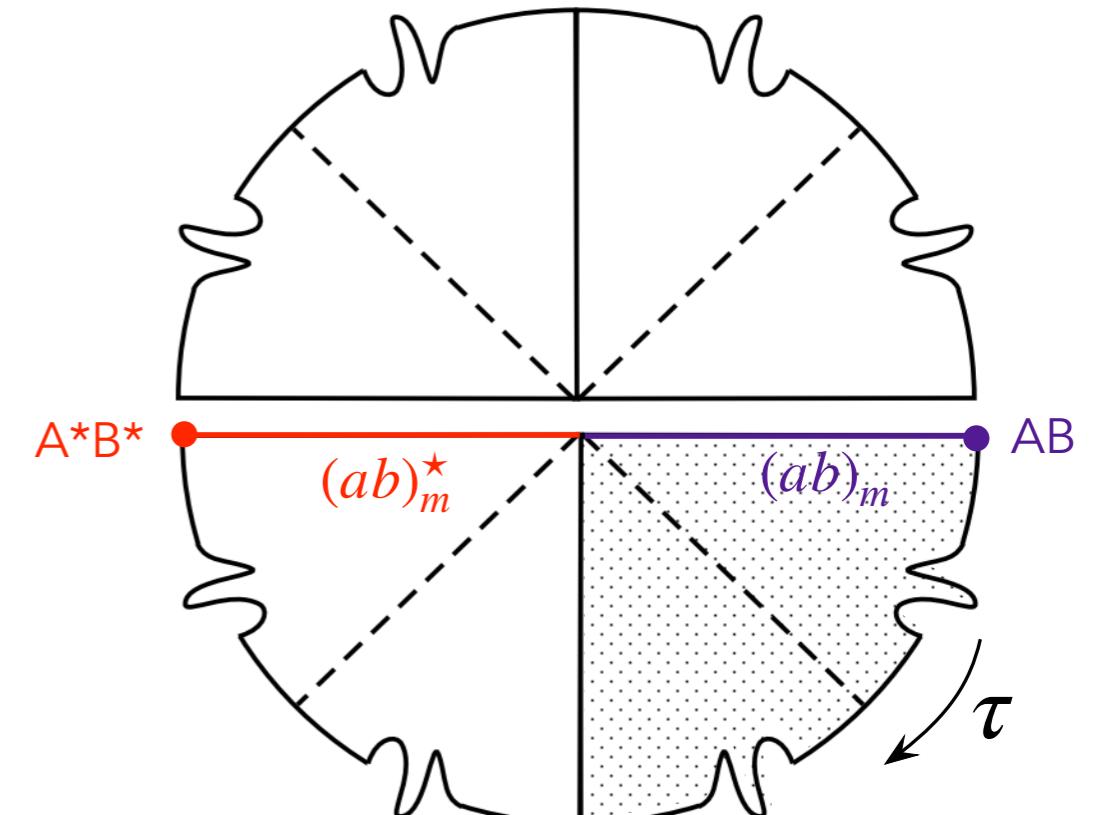


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Gabriel's talk

$(ab)_m$ and $(ab)_m^*$ are entanglement wedges for AB , A^*B^*



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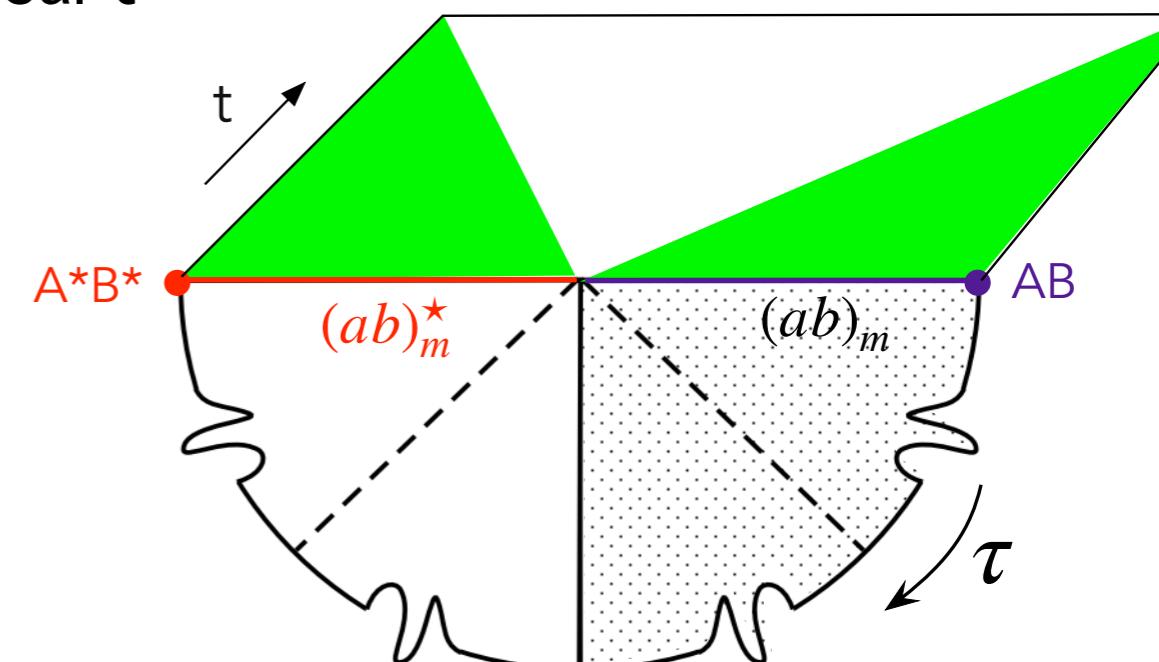
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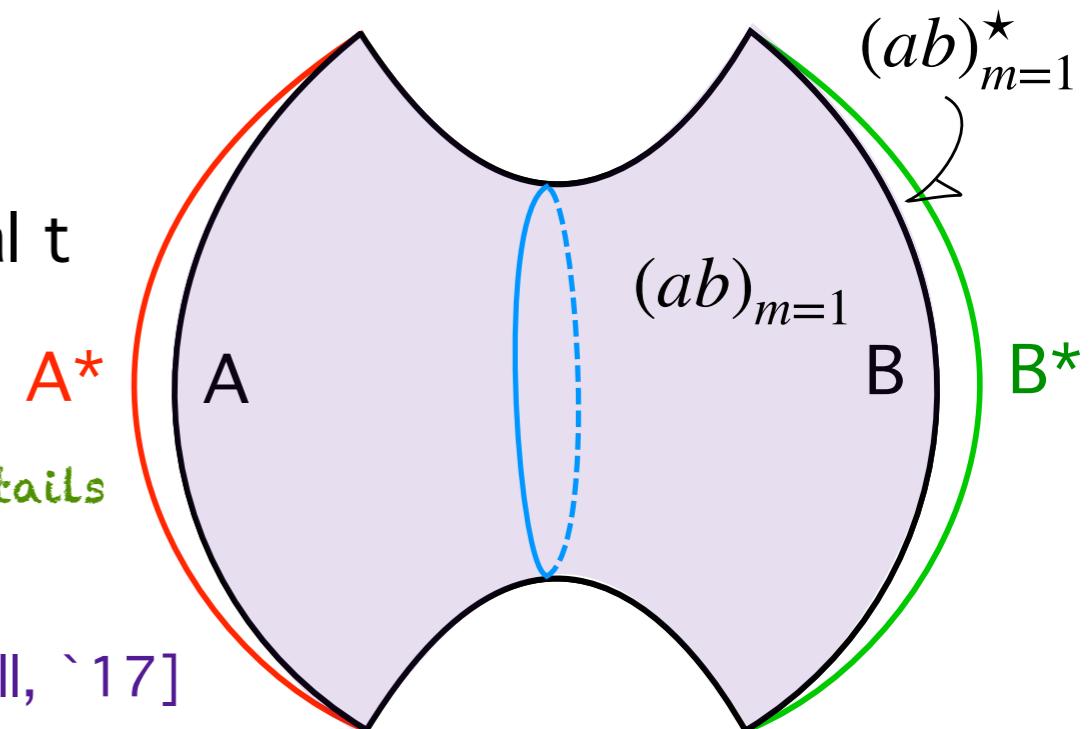
- In the limit $m \rightarrow 1$:

See Arvin's talk for details

- Spacetime gluing prescription across extremal surface [Engelhardt, Wall, `17]

- \mathbb{Z}_m fixed point becomes RT surface for AB

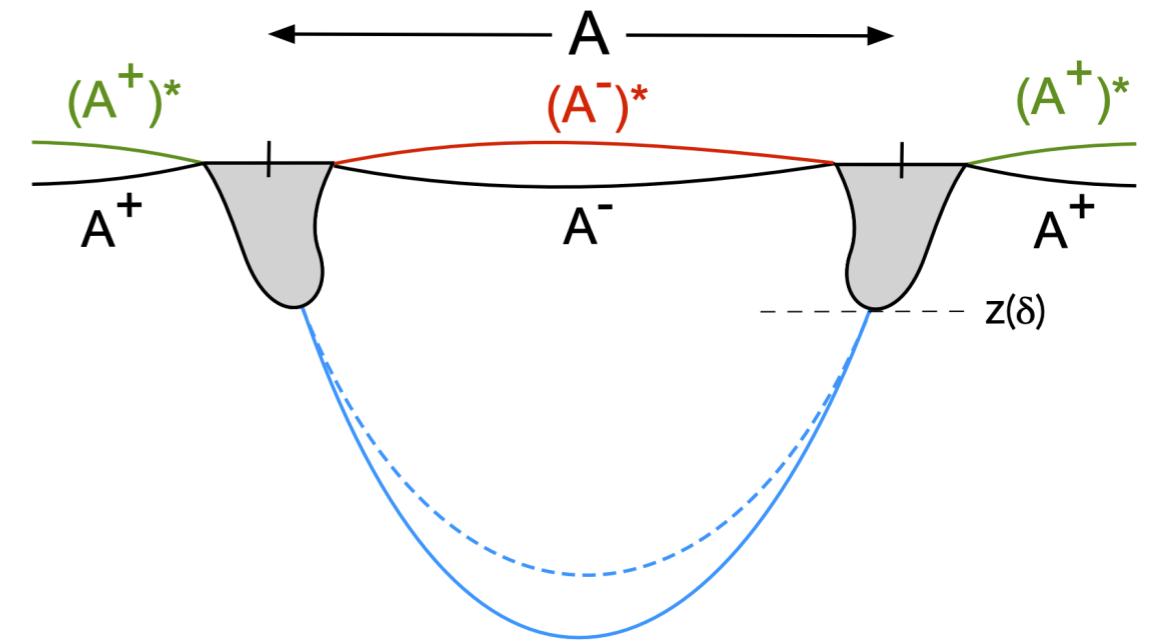
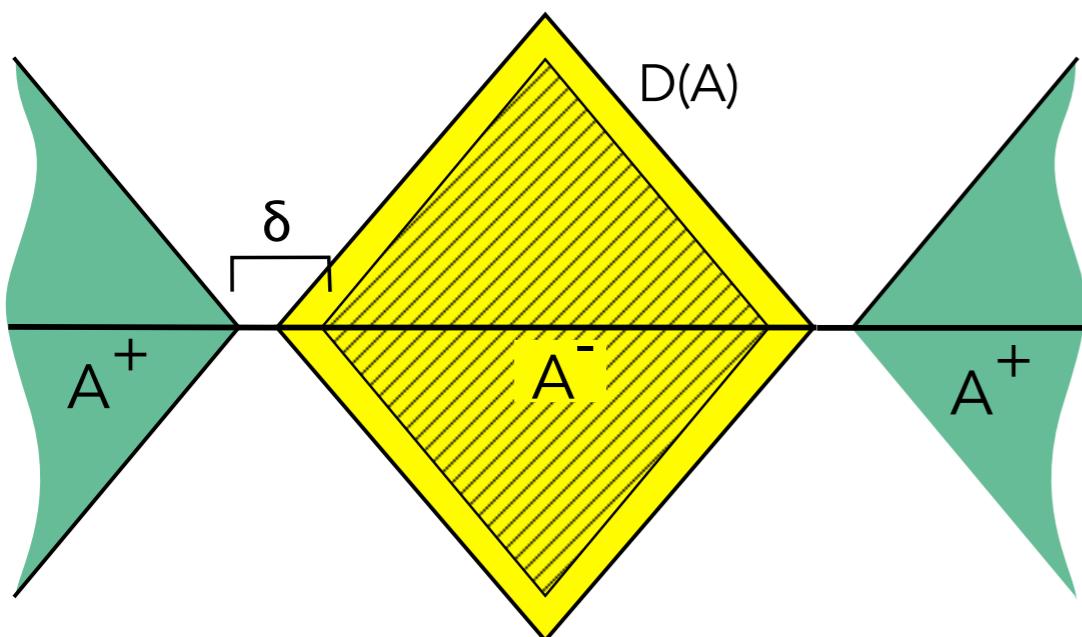
- By construction, $(ab)_{m=1} = r(ab)$, $(ab)_{m=1}^* = r(a^*b^*)$



$S_R(A : B)$ as geometric regulator

- Reflected entropy as natural regulator for EE in QFTs
- Similar to the Mutual Information regulator

[Casini, Huerta, Myers, Yale, '15]



Boundary: $\mathcal{D}(A^-) \subset \mathcal{D}(A) \subset \mathcal{D}((A^+)^c)$

$$\frac{1}{2}I(A^+ : A^-) \leq \frac{1}{2}S_R(A^- : A^+)$$



Bulk: Entanglement wedge for $A^- \cup A^+$

$$z(\delta) \propto \delta \implies \frac{1}{2}S_R(A^- : A^+) \approx S_{EE}^{(\delta)}(A)_\psi$$

[Calabrese, Cardy, '05]

ありがとうございます

