

# Path-Integral Optimization & Complexity in CFT

(+an attempt in BCFT)

Kento Watanabe (U.Tokyo)

Based on

- Phys.Rev.Lett. 119 (2017) no.7, 071602
  - JHEP 1711 (2017) 097
- w/ Pawel Caputa, Nilay Kundu,  
Masamichi Miyaji and Tadashi Takayanagi

(+ work just started w/ Yoshiki Sato(IPMU))

In this talk, I will review

## Path-Integral Optimization & the “Complexity” in CFT

[Caputa-Kundu-Miyaji-Takayanagi-KW'17] 2 years ago...

Maybe useful appetizer to listen [Michal's talk](#) on complexity  
[Masamichi's talk](#) on EoP

Then, I will try to show you our recent attempt in BCFT...

Preliminary observation  
just started w/ [Yoshiki Sato](#)

# AdS/CFT & Quantum Information

Modern Perspective of AdS/CFT  $\longrightarrow$  A “Geometrization” of Quantum States

Quantum Entanglement



Emergent Geometry

AdS/CFT  
Tensor Network

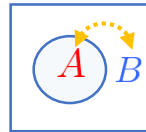
“Gravity”  
“Spacetime”

Ex: Entanglement Entropy (EE)

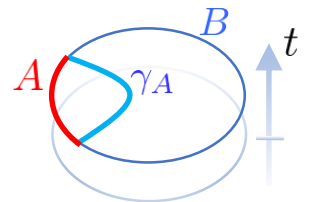


Min. Area of codim.-2 Surface

$$S_A = -\text{Tr} \rho_A \log \rho_A$$

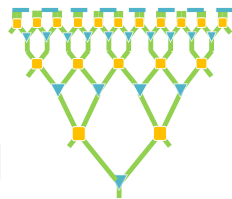


$$\min_{\gamma_A} \left[ \frac{\text{Area}(\gamma_A)}{4G_N} \right] \quad [\text{Ryu-Takayanagi 06}]$$



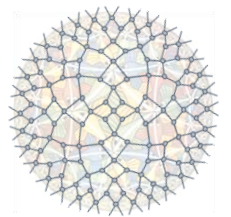
MERA network  
(for CFT vacuum)

[Vidal 05,06]



Hyperbolic geometry (a time slice of AdS)

[Swingle 09]



But still  
mysterious  
mechanism...

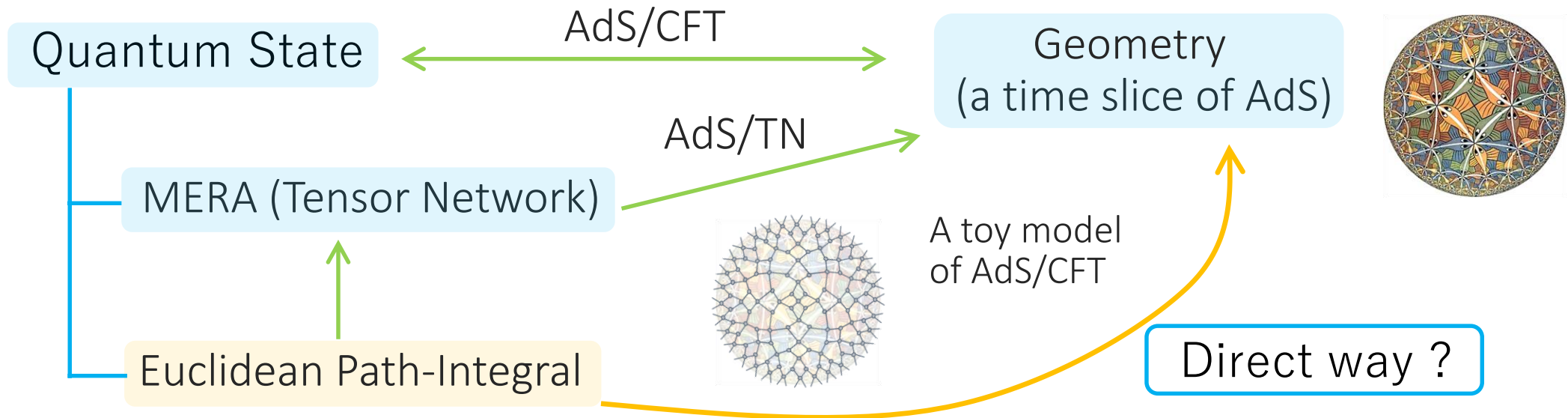


More direct way to dual geometries from quantum states?  
Entanglement is NOT enough... Better probe? Complexity?

# Motivation & Proposal (1)

AdS/CFT → 20 years old!  
but still mysterious...

Direct or Systematic Way to Get Information about Dual Geometries?



**Our Proposal** [Caputa-Kundu-Miyaji-Takayanagi-KW'17]

“Optimization” of Euclidean Path-Integral for Wave Functional in CFTs

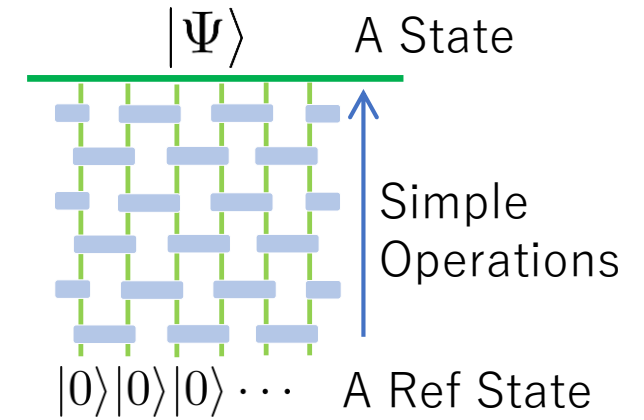
In 2d CFT  
→

“Minimization” of the Liouville Action of the Back-ground Metric

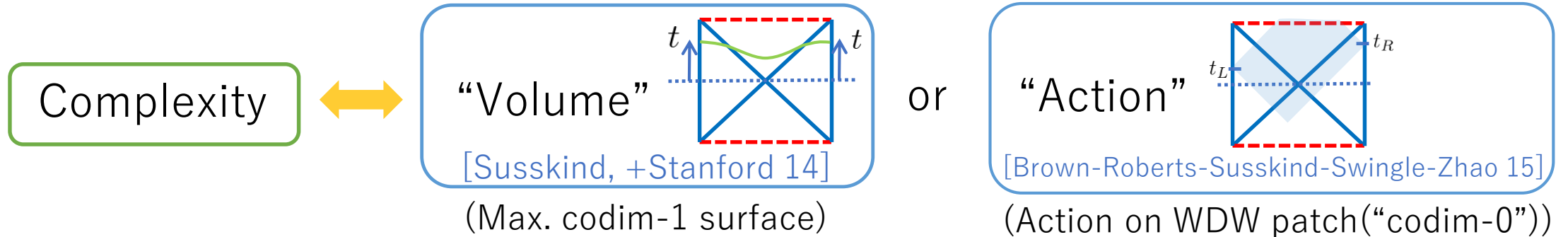
# Motivation & Proposal (2)

CFT Analogue of “Complexity of Quantum State” ?

$$C_{\Psi} = \min[ \#(\text{Operations}) ]$$



- Holographic Complexity  $\rightarrow$  A new probe for dual spacetime beyond HEE



- CFT analogue?  $\rightarrow$  No Definition of the complexity in CFTs so far...

A few developing attempts...  
(including ours)

[Susskind, + collaborators...] [Chapman-Heller-Marrochio-Pastawski 17]  
 [Jefferson-Myers 17] [Yang, + collaborators 17] **Michal's talk?**  
 [Caputa-Magan 18]...etc

# Motivation & Proposal (2)

CFT Analogue of “Complexity of Quantum State” ?

**Our Proposal** [Caputa-Kundu-Miyaji-Takayanagi-KW'17]

Complexity of States in CFTs

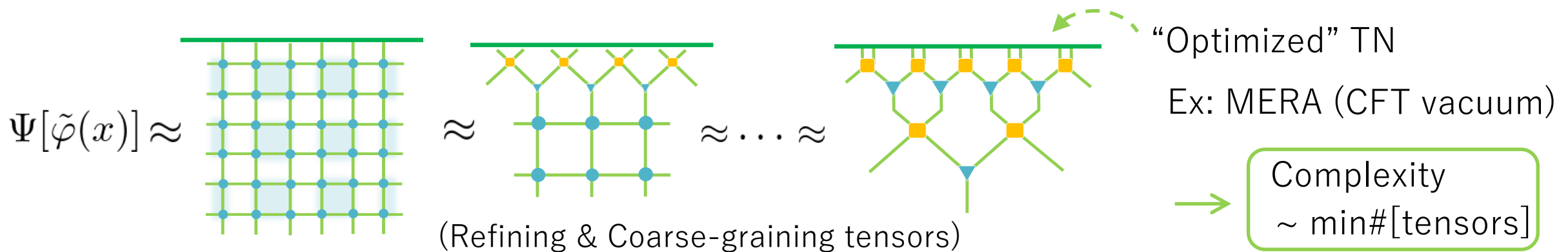


“Optimized Action”

In 2d CFT,

“Liouville Action”

**Hint** Tensor Network Renormalization (TNR) Procedure  $\longrightarrow$  “Optimization”



# Optimization of Euclidean Path-Integrals & Complexity in CFTs

[Caputa-Kundu-Miyaji-Takayanagi-KW'17]

More detail...

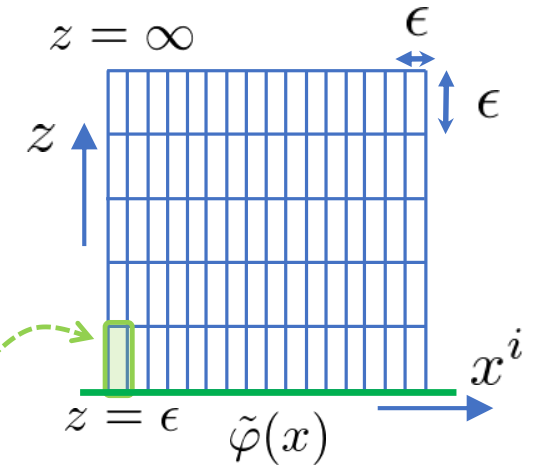
# Basic Rules for Our “Optimization” Procedure

( CFT Analogue of TNR Procedure )

- Discretization of Euclidean path-integral  
(Regularization)

$$\Psi[\tilde{\varphi}(x)] = \int \left( \prod_x \prod_{\epsilon < z < \infty} D\varphi(z, x) \right) e^{-S_{CFT}(\varphi)} \cdot \prod_x \delta(\varphi(\epsilon, x) - \tilde{\varphi}(x))$$

→ Metric (one cell = unit area)  $ds^2 = \epsilon^{-2} \cdot (dz^2 + dx^i dx^i)$

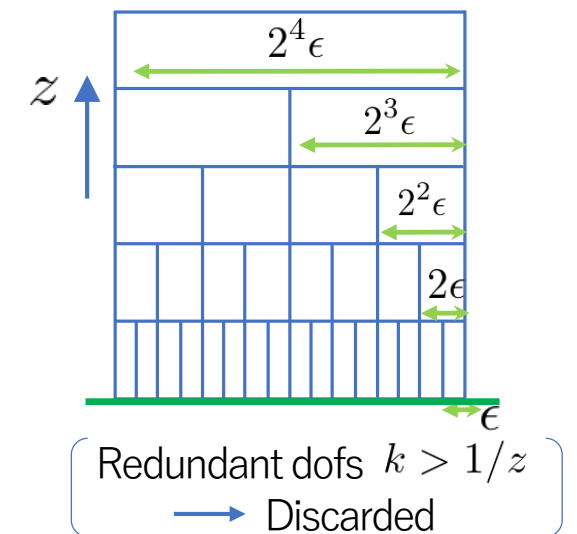


- “Optimization” of the path-integral

→ Changing the geometry of the lattice regularization

→ Modifying the back-ground metric for the path-integral  
with fixing the UV bdy condition

$$g_{zz}(z = \epsilon, x) = \epsilon^{-2}, \quad g_{ij}(z = \epsilon, x) = \delta_{ij} \cdot \epsilon^{-2} \quad g_{iz}(z =, x) = 0$$





# Basic Rules for Our “Optimization” Procedure

( CFT Analogue of TNR Procedure )

- After optimization, reproduce the correct wave functional up to a normalization (Ansatz)

$$\longrightarrow \Psi_{g_{ab}}[\tilde{\varphi}(x)] = N_{(g,\delta)} \cdot \Psi_{\delta_{ab}}[\tilde{\varphi}(x)]$$

Estimate redundant dofs  $\longrightarrow$  Minimize this!!

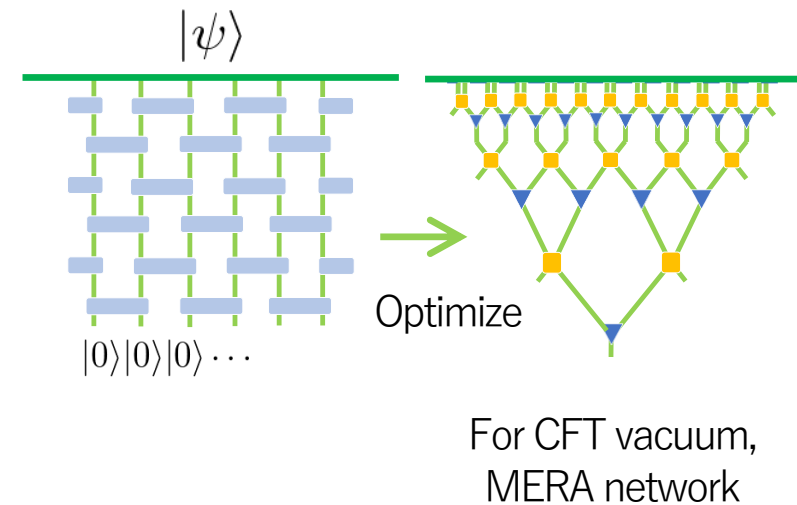
- “Optimization” of the path-integral

$\longrightarrow$  Minimizing “# of lattice points”  
(= “# of tensors in a TN”)

$\longrightarrow$  Our Conjecture for Complexity in CFT

$$N_{opt} \approx \exp(\text{Complexity})$$

$\longrightarrow$  Min[#(Operations)]



# 2d CFT case

- In 2d CFT, we can diagonalize the general back-ground metric

Weyl Scaling  $ds^2 = e^{2\phi(z,x)}(dz^2 + dx^2)$  with  $e^{2\phi(z=\epsilon,x)} = \frac{1}{\epsilon^2}$

- The change of the measure is characterized by the Liouville action

$$[D\varphi]_{g_{ab}=e^{2\phi}\delta_{ab}} = e^{S_L[\phi]-S_L[0]} \cdot [D\varphi]_{g_{ab}=\delta_{ab}}$$

$$S_L[\phi] = \frac{c}{24\pi} \int_{-\infty}^{\infty} dx \int_{\epsilon}^{\infty} dz [(\partial_x \phi)^2 + (\partial_z \phi)^2 + \mu e^{2\phi} + R_0 \phi]$$

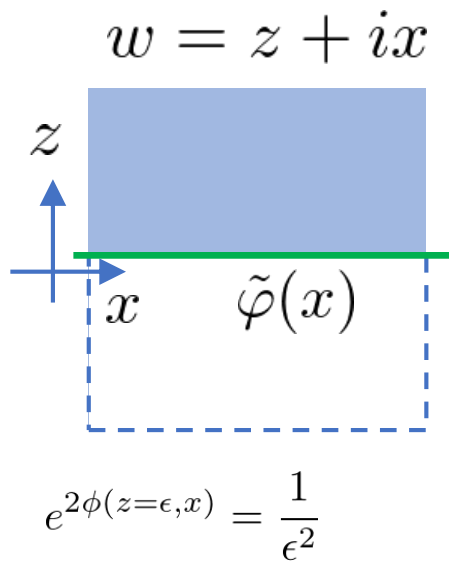
UV regularization  $\dashrightarrow$  # of unitaries in TN

Conformal Anomaly  $\dashrightarrow$  # of isometries [Czech 17]

$$\longrightarrow \Psi_{g_{ab}=e^{2\phi}\delta_{ab}}[\tilde{\varphi}(x)] = e^{S_L[\phi]-S_L[0]} \cdot \Psi_{g_{ab}=\delta_{ab}}[\tilde{\varphi}(x)]$$

Minimize this !!  $\min_{\phi} [S_L(\phi) - S_L(0)] \approx (\text{Complexity})$

# Vacuum on Plane (Poincare AdS<sub>3</sub>)



$$\text{EOM: } \partial_w \partial_{\bar{w}} \phi = \frac{\mu}{4} e^{2\phi} \longrightarrow e^{2\phi} = \frac{4}{\mu} \cdot \frac{A'(w)B'(\bar{w})}{(1 - A(w)B(\bar{w}))^2}$$

$$\text{Especially, } A(w) = w, \quad B(\bar{w}) = -1/\bar{w}$$

a time slice of Poincare AdS<sub>3</sub>

$$\longrightarrow e^{2\phi} = \frac{4}{\mu} \cdot \frac{1}{(w + \bar{w})^2} = \frac{1}{\mu} \cdot \frac{1}{z^2} \longrightarrow ds^2 = \frac{dx^2 + dz^2}{z^2} : H_2$$

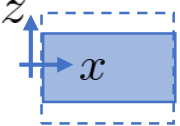
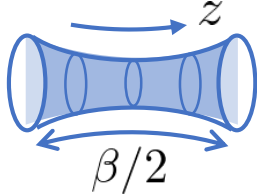
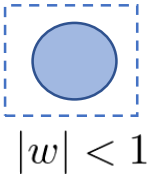
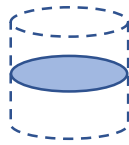
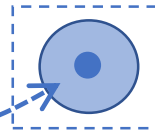
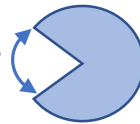
This solution clearly minimizes the Liouville action :  $e^{2\phi(z=\infty, x)} = 0 \quad L = \int dx$

$$S_L = \frac{c}{24\pi} \int dx dz \left[ \underbrace{(\partial_x \phi)^2}_{\text{dashed}} + \underbrace{(\partial_z \phi + \sqrt{\mu} e^\phi)^2}_{\text{dashed}} \right] - \frac{c}{12\pi} \int dx \left[ \sqrt{\mu} e^\phi \right]_{z=\epsilon}^{z=\infty} \geq \frac{c\sqrt{\mu}L}{12\pi\epsilon}$$

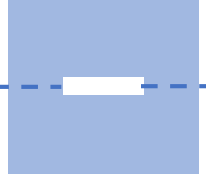
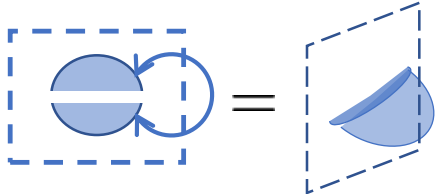

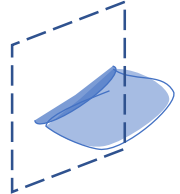
Volume divergence!!

# Other Examples

Work well for simple examples!!

- Finite T or TFD state   $\beta/2$   $\rightarrow$  Wormhole  (A time slice of BTZ BH)
- Vacuum on disk   $|w| < 1$   $\rightarrow$  Hyperbolic disk  $H_2$   (A slice of global AdS3)
- Primary State  $O(w, \bar{w}) \propto e^{-2h\phi}$    $\rightarrow$  Conical Singular Geometry (Match to AdS3/CFT2 for  $h \ll c$ )  
Deficit angle 

Similarly,

- $\rho_A = \text{Tr}_{A^c} |\Psi\rangle\langle\Psi|$    $\rightarrow$  Entanglement Wedge 
- $S_A = -\text{Tr} \rho_A \log \rho_A$    $n \rightarrow 1$   $\rightarrow$  Holographic EE  $S_A \propto \text{Min}[\text{Area}]$  

# Examples : Our “Complexity” in CFTs

## 2d CFT

$$\min_{\phi} [S_L(\phi) - S_L(0)] \approx (\text{Complexity})$$

Volume divergence !!

- Vacuum on plane (Poincare AdS3)  $S_L = \frac{cL}{12\pi\epsilon}$   $L = \int dx$
- Vacuum on disk (Global AdS3)  $S_L = \frac{c}{6} \cdot \left( \frac{1}{\epsilon} - 1 \right)$
- TFD (BTZ BH)  $S_L = \frac{c}{3} \cdot \left( \frac{1}{\epsilon} - \frac{\pi^2}{2\beta} \right)$

With a naïve extension...

## 3d CFT

- (Global AdS4)  $N = \frac{1}{8\pi G_N}$

$$S_L = 4\pi N \cdot \left( \frac{1}{\epsilon^2} + \frac{1}{2} + \log \left( \frac{2}{\epsilon} \right) \right)$$

## 4d CFT

- (Global AdS5)  $N = \frac{3}{16\pi G_N}$

$$S_L = 2\pi^2 N \cdot \left( \frac{2}{3\epsilon^3} + \frac{1}{\epsilon} - \frac{5}{12} \right)$$

- $\left[ \begin{array}{l} \bullet \text{ The volume law leading divergence} \\ \quad \& \text{ The divergence structure} \end{array} \right]$  agree with the holographic complexities !!
- The relative coefficients are different in general...

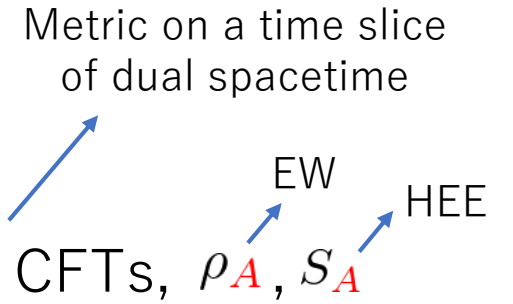
[Chapman-Marrochio-Myers 16]

[Lehner-Poisson-Myers-Sorkin 16]

[Carmi-Myers-Rath 16] [Reynolds-Ross16]

# Summary [Caputa-Kundu-Miyaji-Takayanagi-KW'17]

- A proposal to define complexity of states in CFTs using “Optimization” of Euclidean Path-Integrals
  - Some checks for some states in 2d CFTs,  $\rho_A, S_A$
- “Complexity” = “Optimal Action” = “the Liouville Action” (2d) →  $\min_{\phi}[S_L(\phi) - S_L(0)] \approx (\text{Complexity})$
- (Qualitative) Matching to the Holographic Conjectures!!



## Further Works *Developing in past 2 years!!*

- Higher dims, Time depend., Relation to Holographic Complexity
  - Non-CFT case
  - RG flow
  - Application to EoP → Masamichi's talk
  - ... etc
- • How about CFT w/ defect or bdy ?

# Defects distinguish Holographic Proposals?

[Chapman-Ge-Policastro'18]

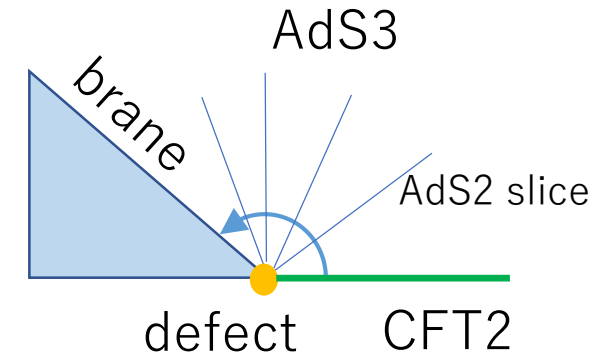
(cf[Flory'17] for CV)

A toy model of AdS3/D(or B)CFT2

[Azeyanagi-Karch-Takayanagi-Thompson'07]

CFT2 w/ a defect (or bdy)  $\longleftrightarrow$  AdS3 w/ a brane on a AdS2 slice

$$I_{\text{bulk}} + I_{\text{brane}} = \frac{1}{16\pi G_N} \int d^3x \sqrt{-g} \left( R - \frac{2}{R_{\text{AdS}}^2} \right) - \lambda \int d^2x \sqrt{-h}$$



[Chapman-Ge-Policastro'18]

$$\begin{aligned} \longrightarrow & \left[ \begin{aligned} C_V - C_{V,\text{vac}} &= \frac{8c}{3} \sinh y^* \cdot \log \left( \frac{2L_B}{\delta} \right) \propto \text{defect brane tension} \\ C_A - C_{A,\text{vac}} &= 0 \quad \text{No defect contribution! (scheme dependent?)} \end{aligned} \right. \end{aligned}$$

$\longrightarrow$  Which is better as the complexity? Need the CFT counterpart...

# Path-Integral Optimization in BCFT (Preliminary...)

(Just started working w/ [Yoshiki Sato](#)...)

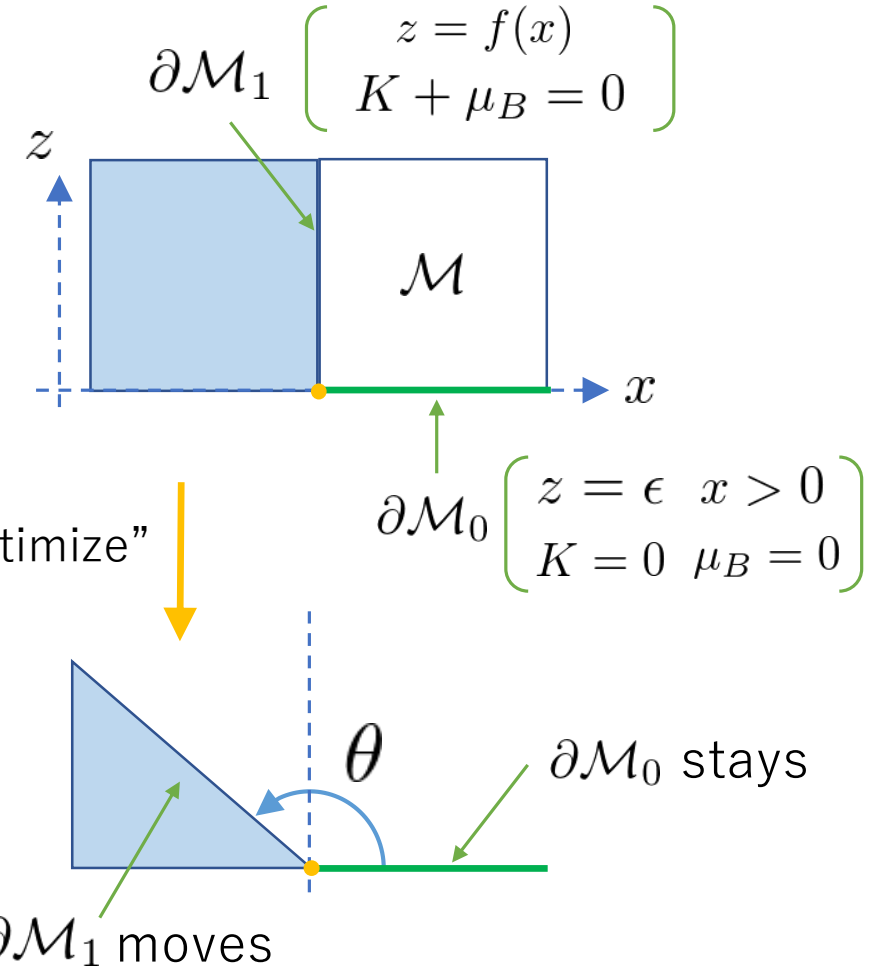
Boundary Liouville field theory

$$S_L[\phi] = \frac{c}{24\pi} \int_{\mathcal{M}} ((\partial\phi)^2 + R\phi + \mu e^{2b\phi}) + \frac{c}{12\pi} \int_{\partial\mathcal{M}} (K\phi + \mu_B e^{b\phi})$$

↓ EOM w/ b.c. on  $\partial\mathcal{M}_0$  &  $\partial\mathcal{M}_1$

$$C_L = \frac{1}{2} \cdot C_{L,\text{vac}} + C_{L,\text{bdy}}(\mu_B)$$

Might appear bdy contribution  
(or corner contribution?)



Check the bdy entropy, g-thm...

Work harder! Please discuss!

Thank you !



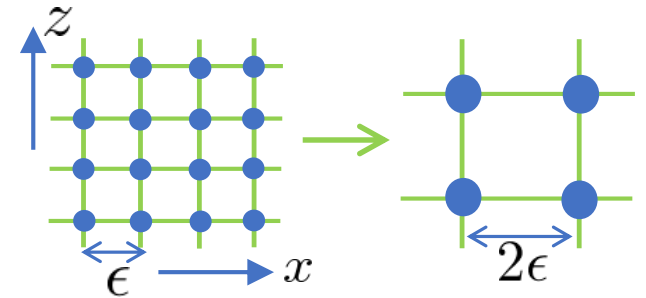


# Naïve Estimation: Liouville Action From TNR

[Czech 17] [Caputa-Kundu-Miyaji-Takayanagi-KW17]

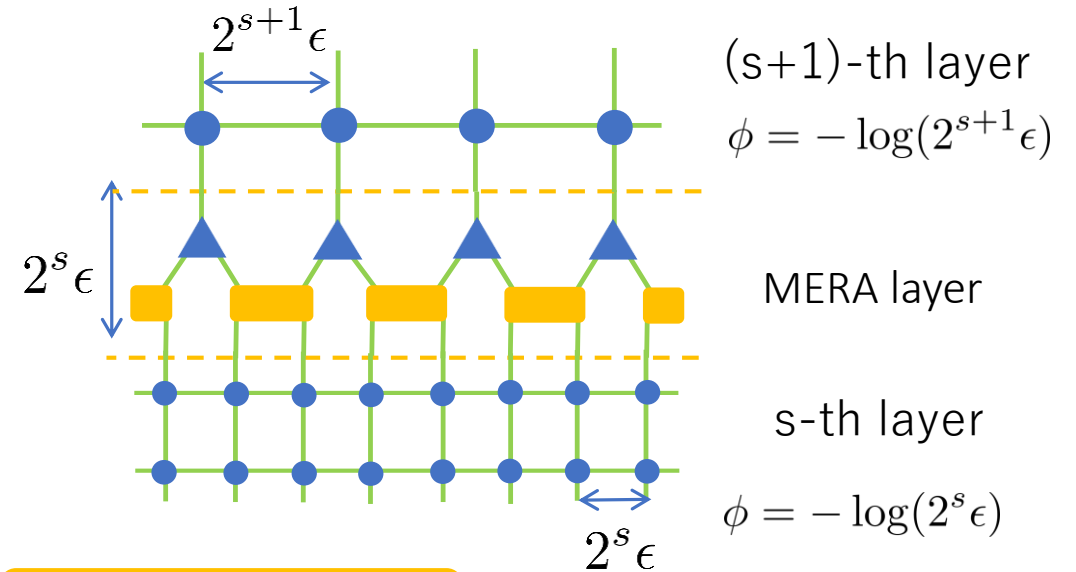
- Suppose each tensor has unit area in the original square lattice

$$\text{---} \bullet \text{---} \frac{dx dz}{\epsilon^2} = e^{2\phi(z,x)} dx dz \xrightarrow{\text{Coarse-graining}} \frac{dx dz}{(2\epsilon)^2}$$



- For the s-th layer of MERA network, per unit cell,

$$\begin{aligned} \text{---} \text{---} \text{---} \text{---} \text{---} &\xrightarrow{\frac{dx dz}{(2^s \epsilon)^2}} e^{2\phi} \\ \text{---} \blacktriangle \text{---} &\xrightarrow{\frac{dx dz}{(2^{s-1} \epsilon)(2^s \epsilon)}} \partial\phi \end{aligned}$$

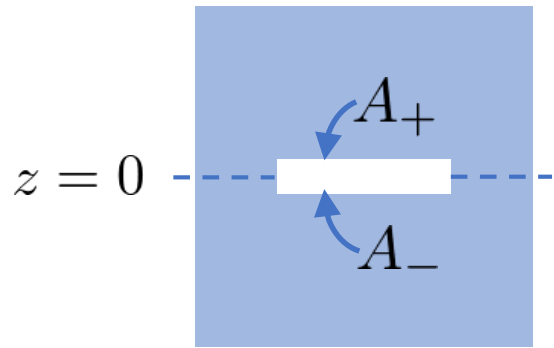


- Total # of tensors in the optimal network

$$\xrightarrow{\text{---}} S_L[\phi] = \frac{c}{24\pi} \int_{-\infty}^{\infty} dx \int_{\epsilon}^{\infty} dz [(\partial\phi)^2 + \mu e^{2\phi}] \sim \text{Complexity !!}$$

# Optimization for Density Matrix

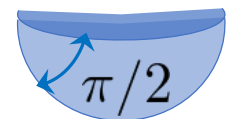
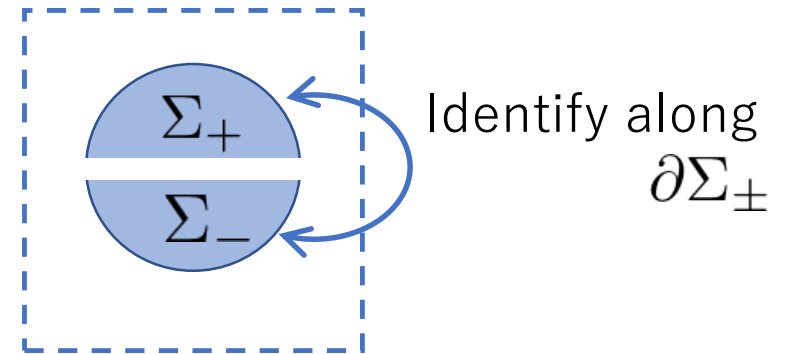
$$\rho_A = \text{Tr}_{A^c} |\Psi\rangle\langle\Psi|$$



Optimize  $S_L$

$$ds^2 = e^{2\phi(z,x)}(dz^2 + dx^2)$$

$$e^\phi|_{A_\pm} = \frac{1}{\epsilon}$$



- The shape of  $\partial\Sigma_\pm$  is given by extremizing

$$S_{Lb} = \frac{c}{12\pi} \int_{\partial\Sigma_\pm} ds [K_0\phi + \mu_B e^\phi] \longrightarrow \mu_B = 0 \text{ (tensionless)}$$

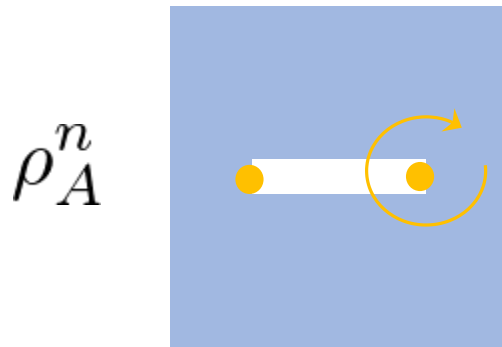
Same as holographic EE !!

- Neumann bdy condition  $\longrightarrow e^\phi K = n^a \partial_a \phi + K_0 = 0$

- Half circle  $x^2 + z^2 = l^2$
- $\longrightarrow$  Entanglement wedge

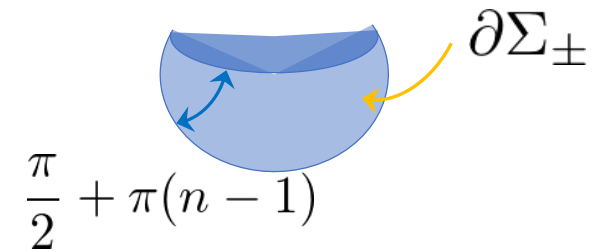
# Entanglement Entropy from Optimization

n-sheeted geometry  
w/ deficit angle  $2\pi(1-n)$



Optimize

$$|n-1| \ll 1$$



- Bdy condition & shape of  $\partial\Sigma_{\pm}$  changes

$$K + \mu_B = 0$$



$$K = \pi(n-1)$$

$$\mu_B = \pi(1-n)$$

Same as holographic EE !!

- $$S_A = -\partial_n \left[ \frac{c\mu_B}{12\pi} \int_{\partial\Sigma_+} ds e^\phi + \frac{c\mu_B}{12\pi} \int_{\partial\Sigma_-} ds e^\phi \right]_{n=1} = \frac{c}{6} \int_{\partial\Sigma_+} ds e^\phi = \frac{c}{3} \log \frac{2l}{\epsilon},$$

$\propto$  Length of the extremal surface