

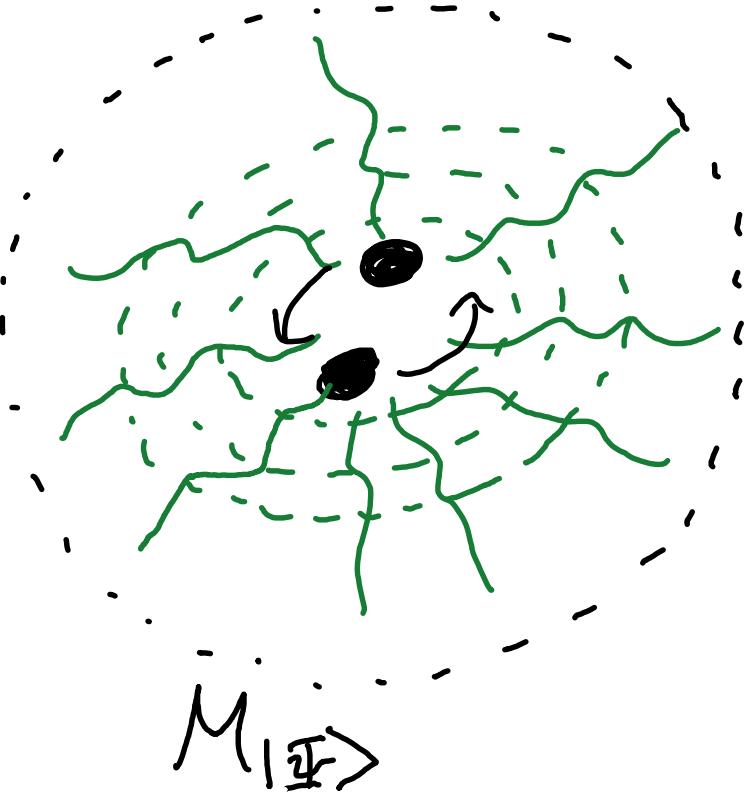
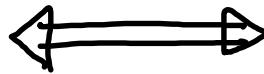
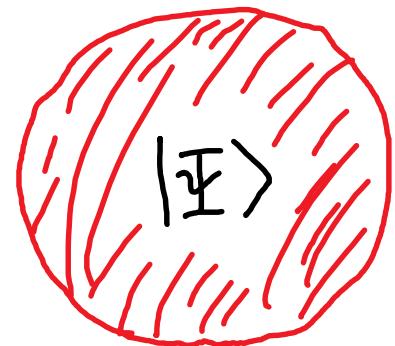
# SPACETIME FROM "BITS"

Mark Van Raamsdonk, UBC

YITP workshop, June 2019

based on: 1808.01197

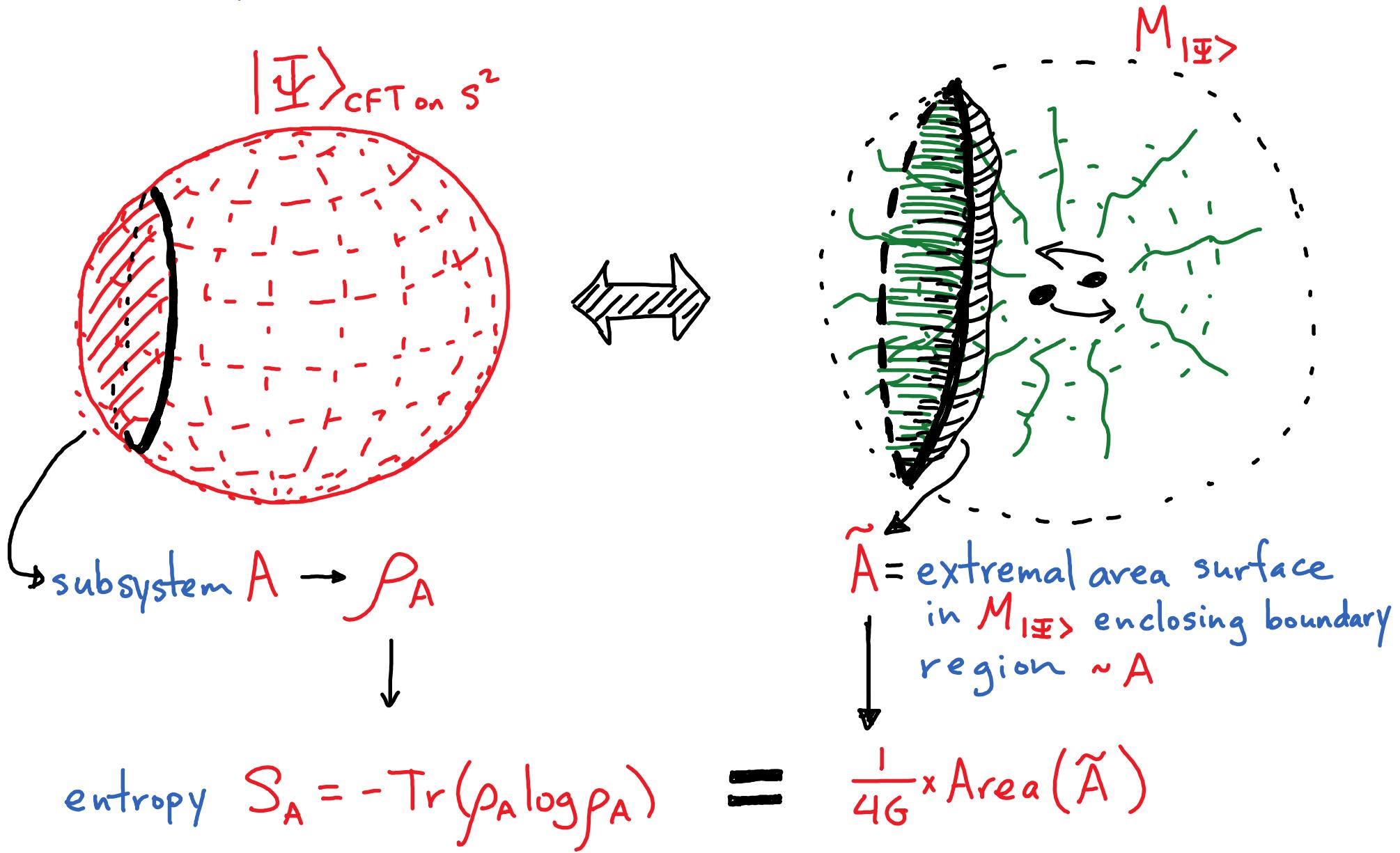
THE AdS/CFT correspondence: spacetime physics +  
gravitation dynamics encoded in a non-gravitational  
quantum system



$$i\hbar \frac{d|\Psi\rangle}{dt} = H |\Psi\rangle$$

$$\dot{M}_{|\Psi\rangle}$$

The Ryu-Takayanagi formula: spacetime geometry is directly related to entanglement structure:

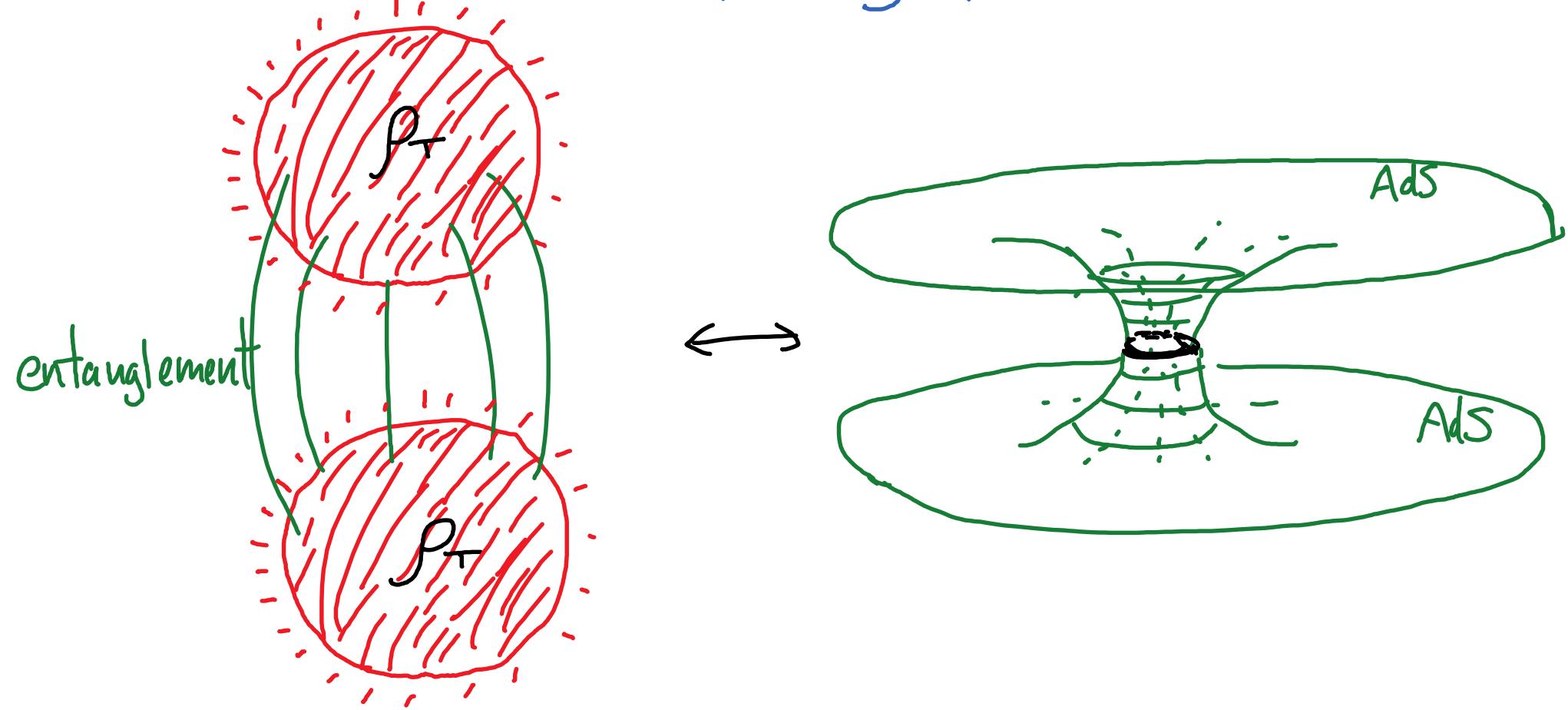


Suggests: Spacetime "emerges" from entanglement in underlying degrees of freedom.

Qualitative picture: removing the entanglement should destroy spacetime

BUT: all QFT states are highly entangled; would cost infinite energy to remove

Inspirational example (Maldacena): Entangling 2 copies of CFT  
can connect the corresponding spacetimes



$$|\Psi\rangle = \sum_i e^{-\beta E_i / z} |E_i\rangle \otimes |E_i\rangle$$

"Thermofield  
Double State"

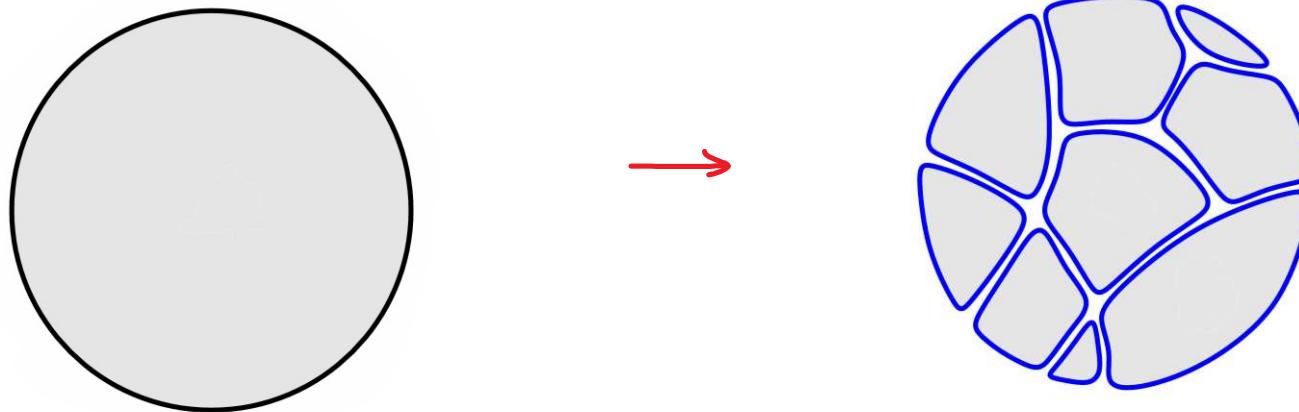
Main question for this talk: can we build up  
any spacetime by entangling discrete parts?

Basic idea: replace continuous CFT with collection of non-interacting systems (CFTs with boundary)

e.g. 1D



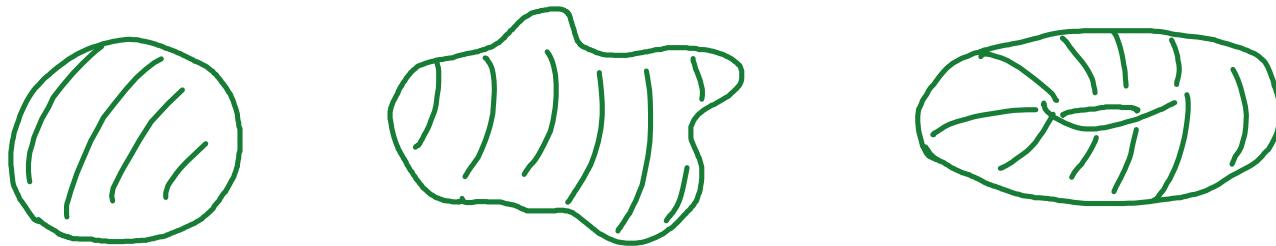
e.g. 2D



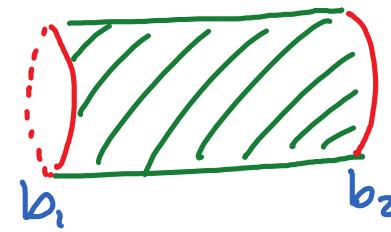
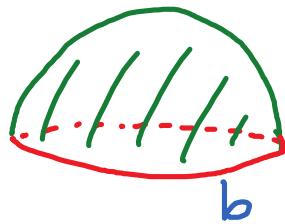
- approximate state  $|\Psi\rangle$  of original theory encoding some spacetime  $M_\Psi$  with entangled state of "discretized" system.
- argue that new state still encodes connected spacetime closely related to  $M_\Psi$

# CFTs with boundary.

- Given a CFT, can define it on various spacetime backgrounds:



- Can also define it on spaces with boundary if we choose boundary conditions for each boundary component



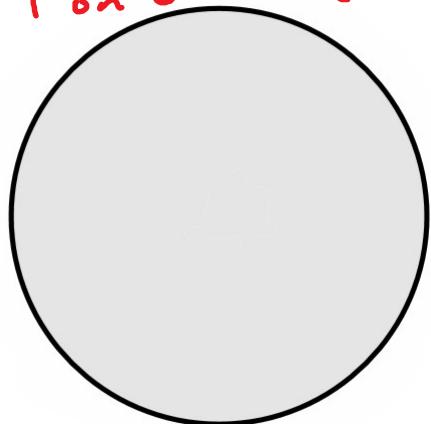
- example: for free scalar theory, can choose Neumann or Dirichlet boundary conditions.

Examples:

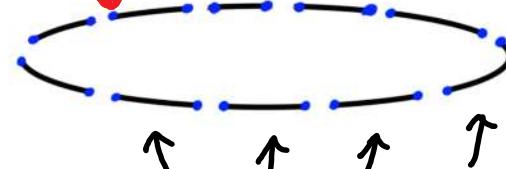
CFT on  $S^1 \times$  time



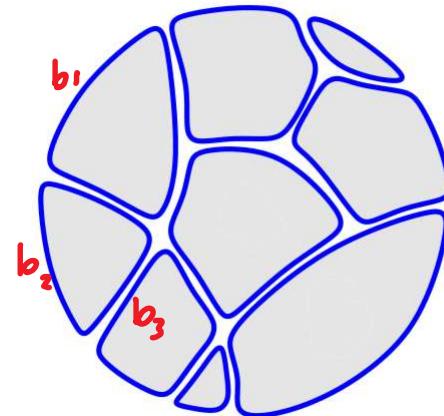
CFT on  $S^2 \times$  time



CFT on interval with boundary conditions  $b_1, b_2$



collection of CFTs on geometries w. disk topology

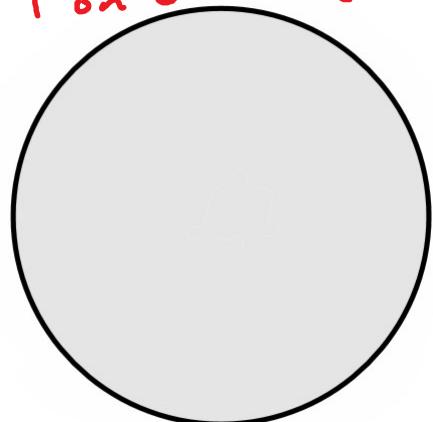


Examples:

CFT on  $S^1 \times \text{time}$



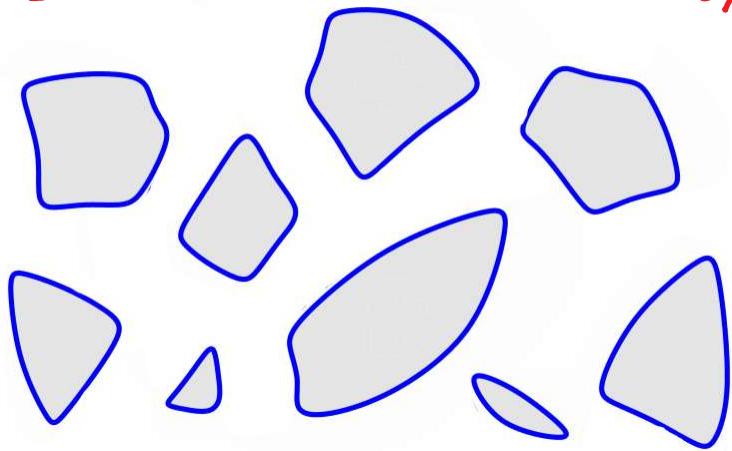
CFT on  $S^2 \times \text{time}$



collection of CFTs on interval



Collection of CFTs on  
geometries with disk topology:



\* no intrinsic spatial arrangement \*

Next step: how do we approximate CFT states by states of the multipart systems?

# Geometrical states via Euclidean path integral

- For "holographic" CFT, only some states encode dual spacetime with nice semiclassical description
- A very large class of these "geometrical" states may be defined via Euclidean path integral.

Standard example: vacuum state

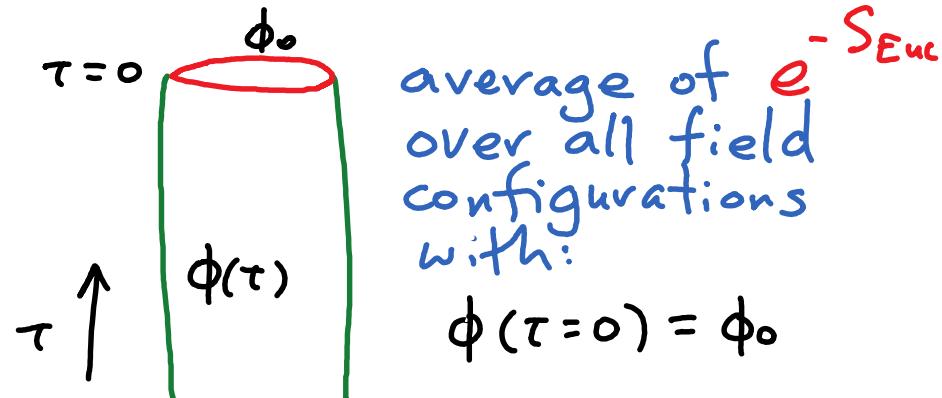
$$|vac\rangle = \lim_{\beta \rightarrow \infty} e^{-\beta H} |\Psi_0\rangle \quad \Rightarrow \quad \text{arbitrary state}$$

with overlaps to vacuum

Euclidean action

$$\Psi_{vac}[\phi_0] = \int_{\tau < 0} [d\phi(\tau)] e^{-S_{Euc}}$$

$\phi(0) = \phi_0$



Harmonic oscillator example:

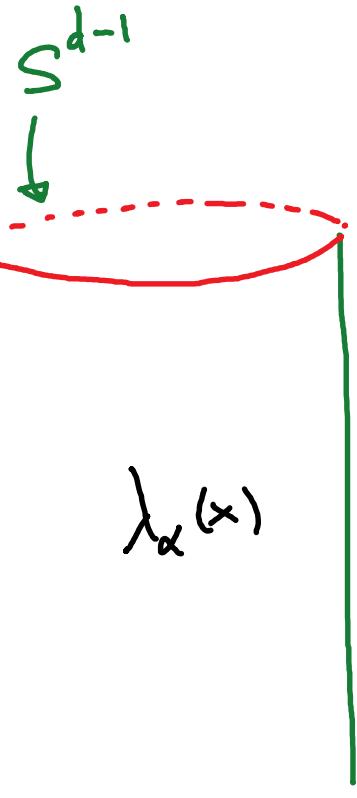
$$\langle x_0 | \text{vac} \rangle = N \cdot \int_{\tau < 0}^{x(0) = x_0} [dx(\tau)] e^{- \int d\tau \left\{ \frac{1}{2} m \dot{x}^2 + \frac{1}{2} m \omega^2 x^2 \right\}}$$

More general states:

$$\langle x_0 | \Psi_\lambda \rangle = N \cdot \int_{\tau < 0}^{x(0) = x_0} [dx(\tau)] e^{- \int d\tau \left\{ \frac{1}{2} m \dot{x}^2 + \frac{1}{2} m \omega^2 x^2 + \lambda(\tau) x(\tau) \right\}}$$

↑  
source  
for  
 $x$  operator

# AdS/CFT:



wavefunction  
for CFT on  $S^{d-1} \times \text{time}$

$$\langle \phi_0 | \Psi_\lambda \rangle = \frac{1}{\sqrt{Z_\lambda}} \int_{\phi(\tau) = \phi_0} [d\phi(\tau)] e^{-S_{\text{Euc}} - \int \lambda_\alpha^{(x)} O_\alpha(x)}$$

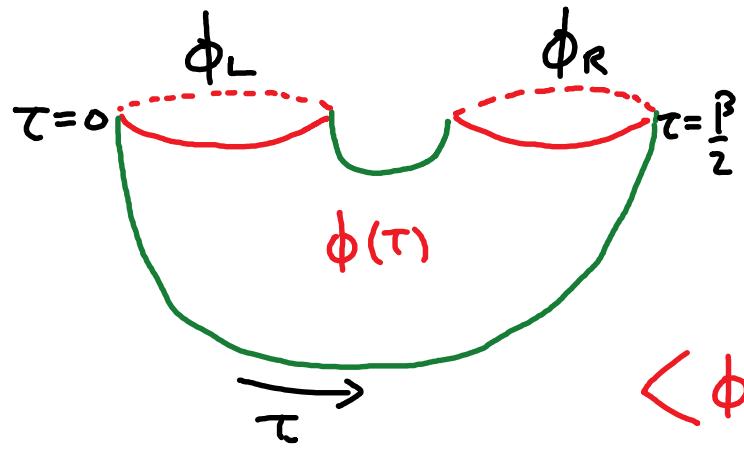
These states have gravity duals with  
a good classical description.  
(will describe dual geometry later)

Entangling via path integrals:

Recall: thermofield double state dual to 2-sided black hole

$$|\Psi_\beta\rangle = \sum_i e^{-\beta E_i/z} |E_i\rangle \otimes |\bar{E}_i\rangle$$

Can define using a Euclidean path integral:



$$\langle \phi_L, \phi_R | \Psi_\beta \rangle \sim \int_{\phi(0)=\phi_L}^{\phi(\beta/2)=\phi_R} [d\phi(\tau)] e^{-S_{\text{EUC}}}$$

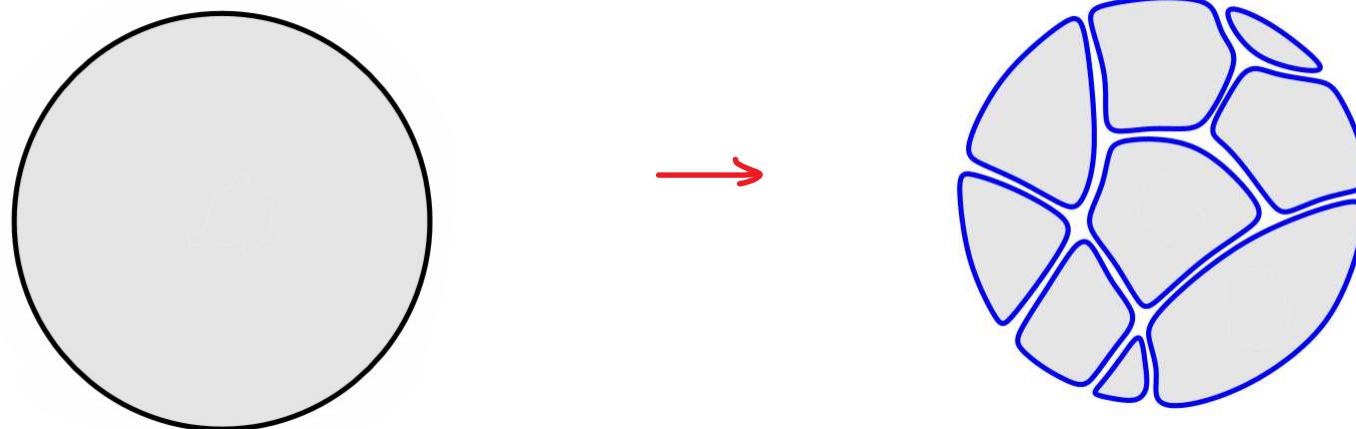
Connected path-integral geometry defines entangled state of non-interacting CFTs.

Recall: we wanted to replace CFT with collection of non-interacting systems

e.g. 1D



e.g. 2D

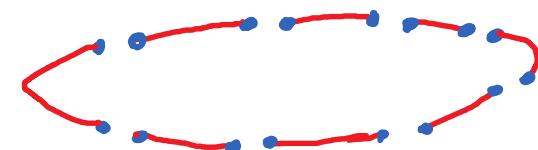


We are now ready to describe how to approximate the CFT states with entangled states of the new system

Old system:



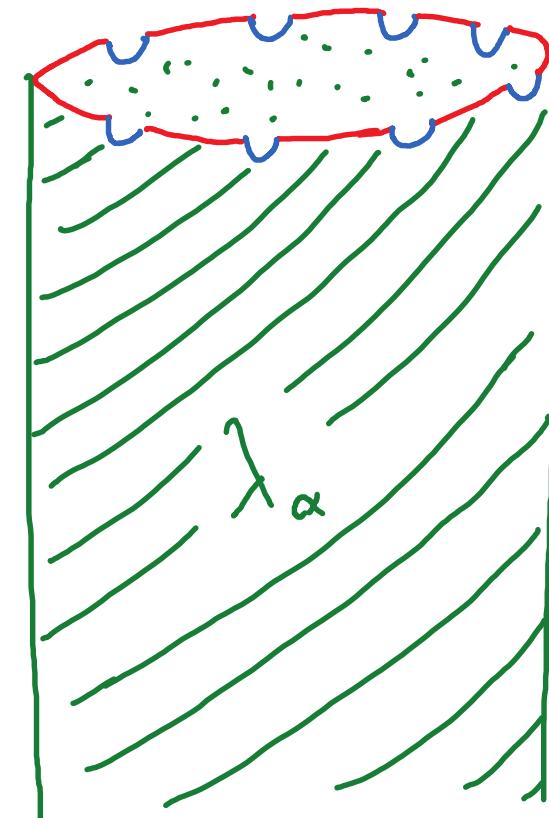
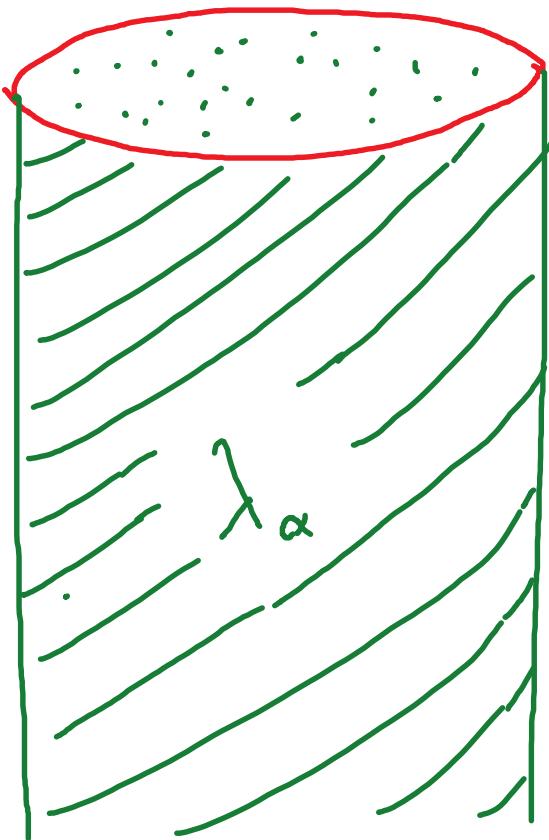
New system:



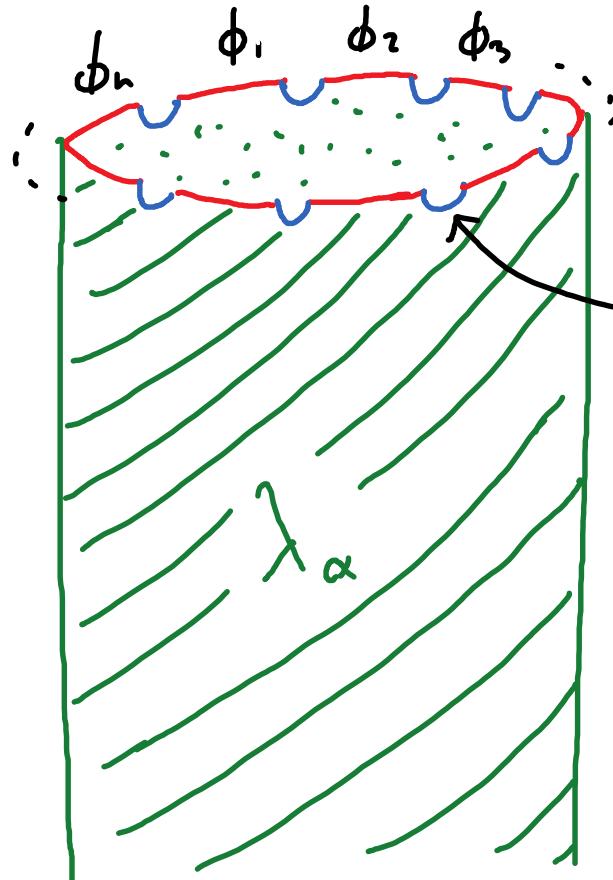
Main idea:

Given CFT state  
defined via path integral:

Define a new path integral  
by introducing boundary  
components:



Explicit expression:



$$\langle \phi_1, \dots, \phi_n | \Psi \rangle$$

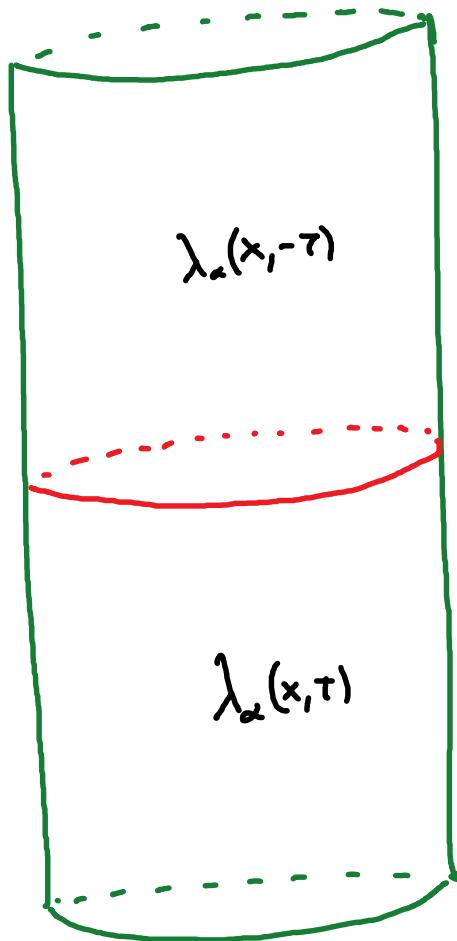
$$\sim \int [d\phi] e^{-S_{\text{Euc}} - \int \lambda_\alpha \partial_\alpha}$$

Need to  
impose  
boundary  
conditions  
here

Defines state of  $n$   
non-interacting CFTs w. boundary.

Next step: understand dual geometries

# AdS/CFT recipe:

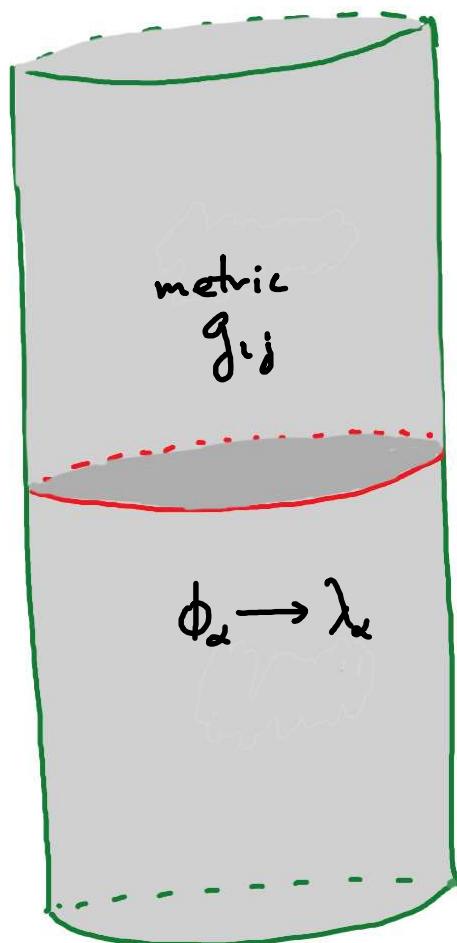


Step 1: Find Euclidean asymptotically  
AdS space with:

boundary geometry = geometry of path  
integral for  $\bar{Z}_\lambda$

boundary conditions for fields  
determined by sources  $\lambda_\alpha$

AdS/CFT recipe:

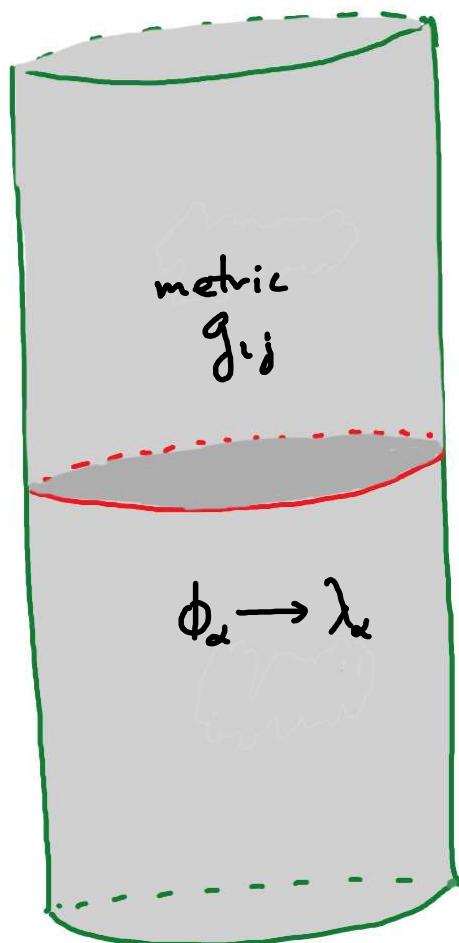


Step 1: Find Euclidean asymptotically  
AdS space with:

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AdS/CFT recipe:



Step 1: Find Euclidean asymptotically AdS space with:

boundary geometry = geometry of path integral for  $\bar{Z}_\lambda$

boundary conditions for fields determined by sources  $\lambda_\alpha$

Step 2: Take slice at  $\tau=0$  surface (shaded). This provides initial data for Lorentzian evolution.

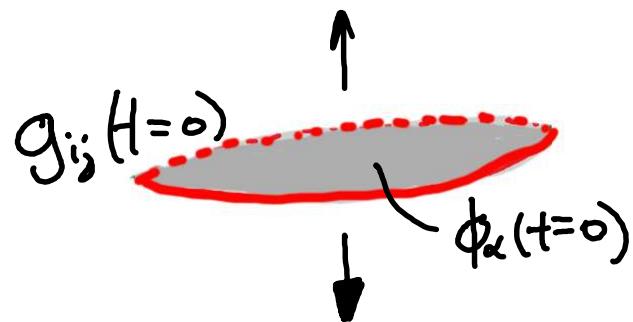
i.e.: evolve this forward & backward in time to find spacetime dual to CFT state.

AdS/CFT recipe:

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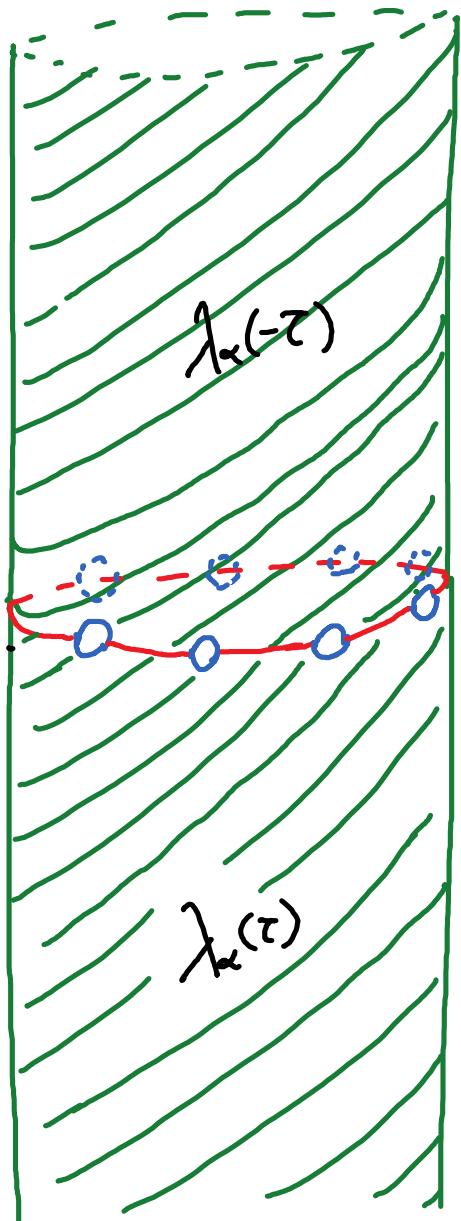


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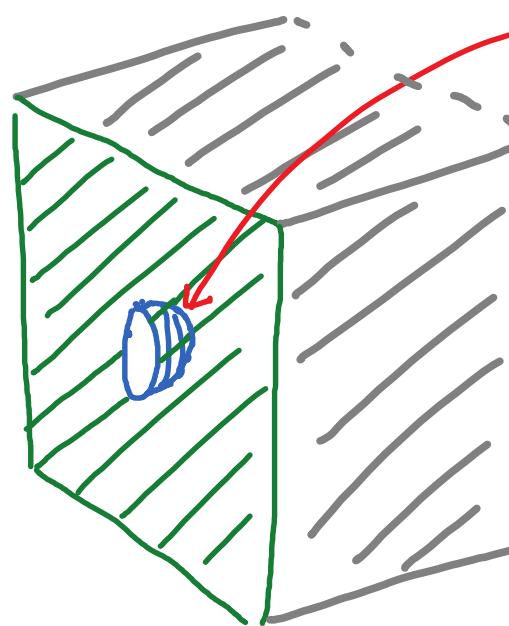
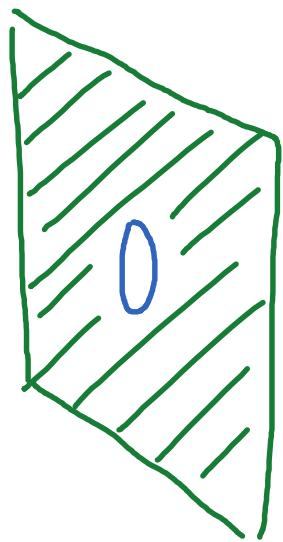
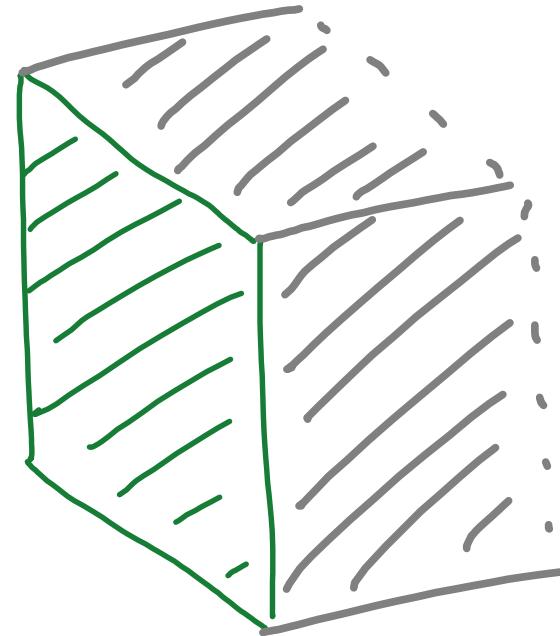
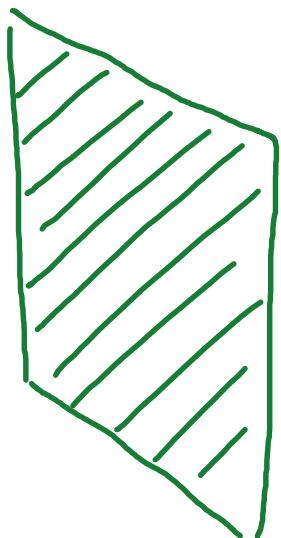
i.e.: evolve this forward & backward in time to find spacetime dual to CFT state.

Repeat this for modified path integral

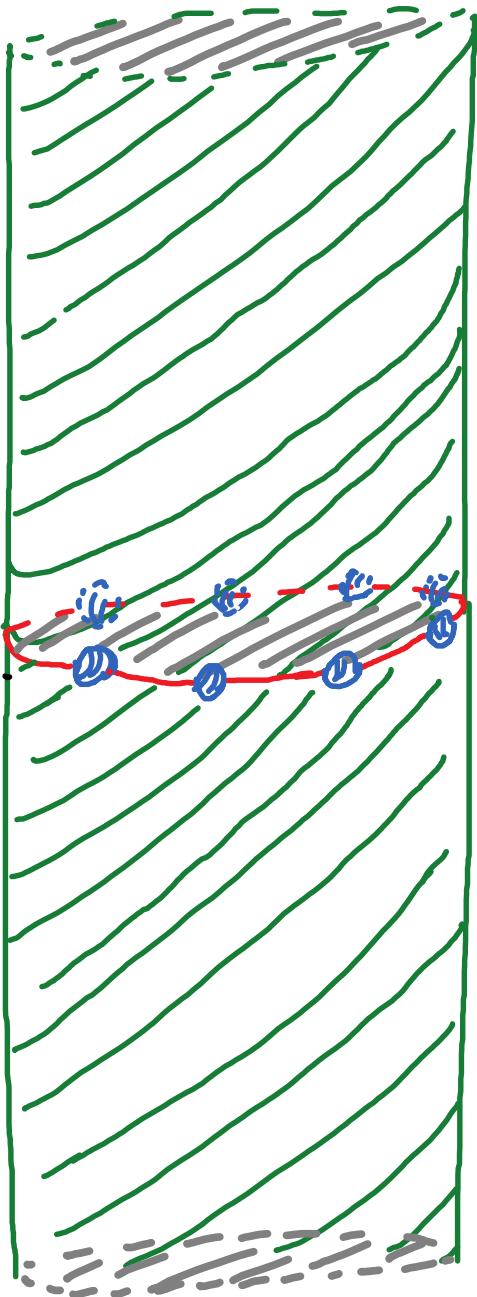
What is the effect of the extra boundaries?



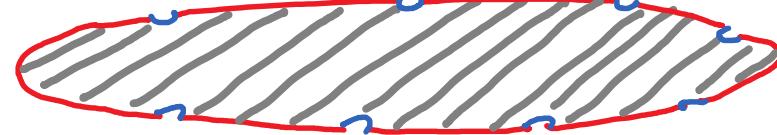
Takayanagi: model by "end of the world" brane



spacetime  
ends on  
this surface

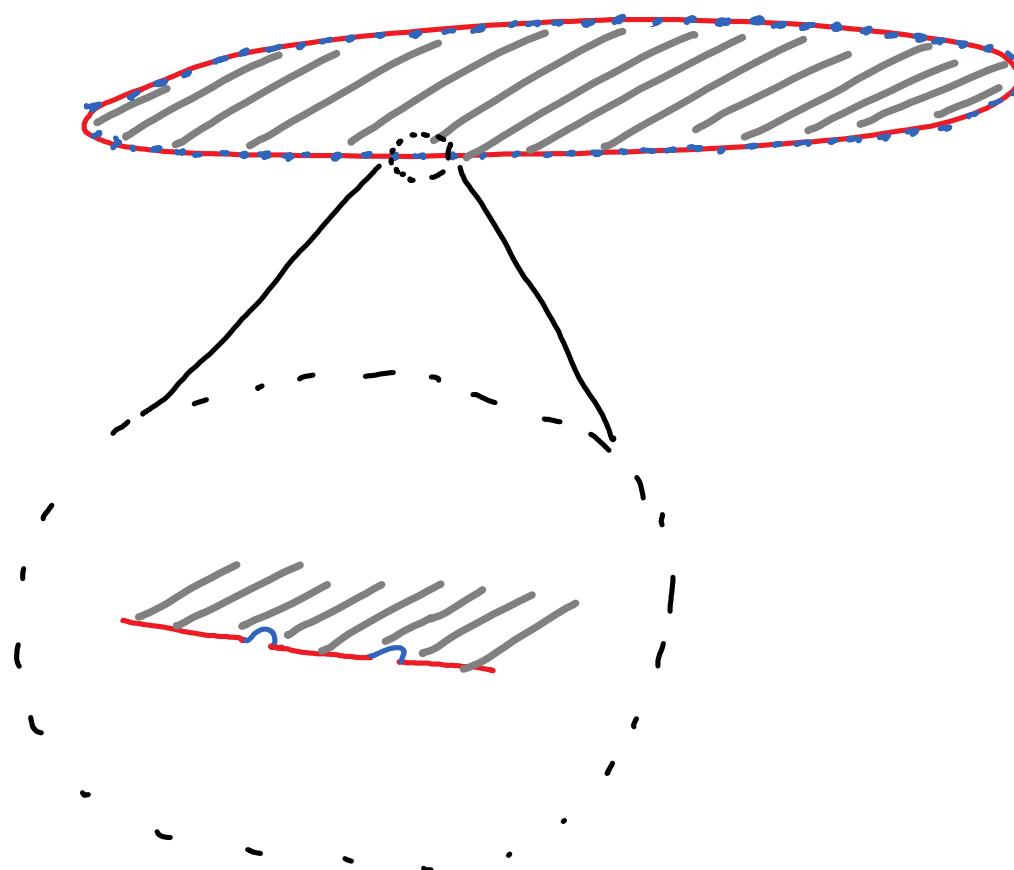


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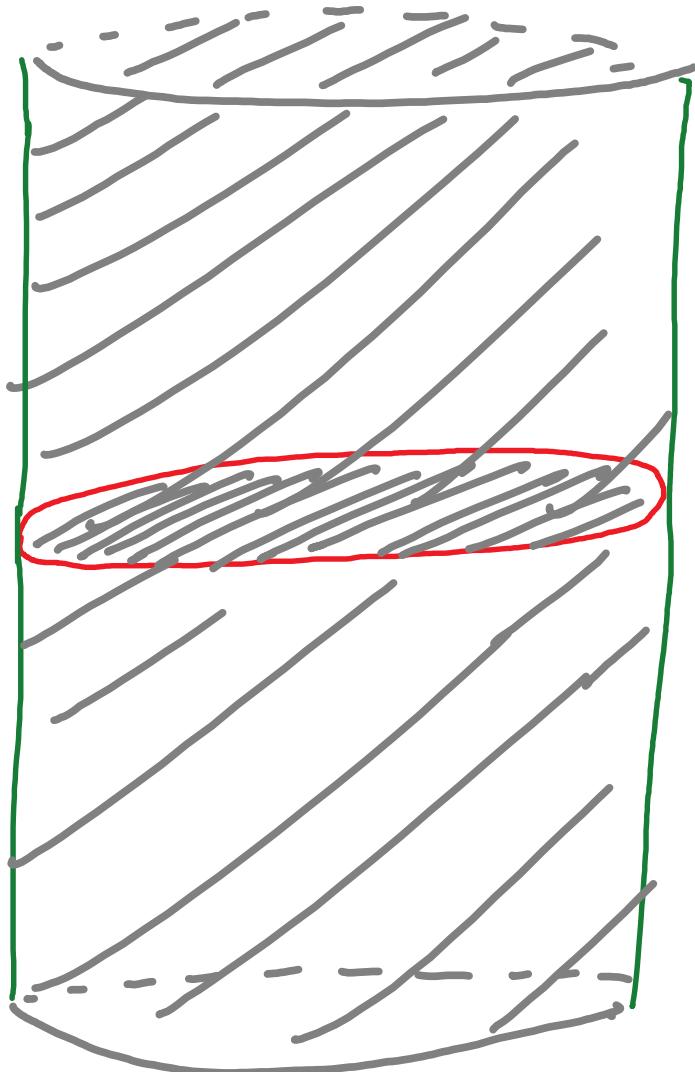
For very small boundary components, geometry of space at initial time differs significantly from original case only near these boundary points.

Same is true even in limit of a very large number of boundary components if size of "gaps" small relative to size of subsystems

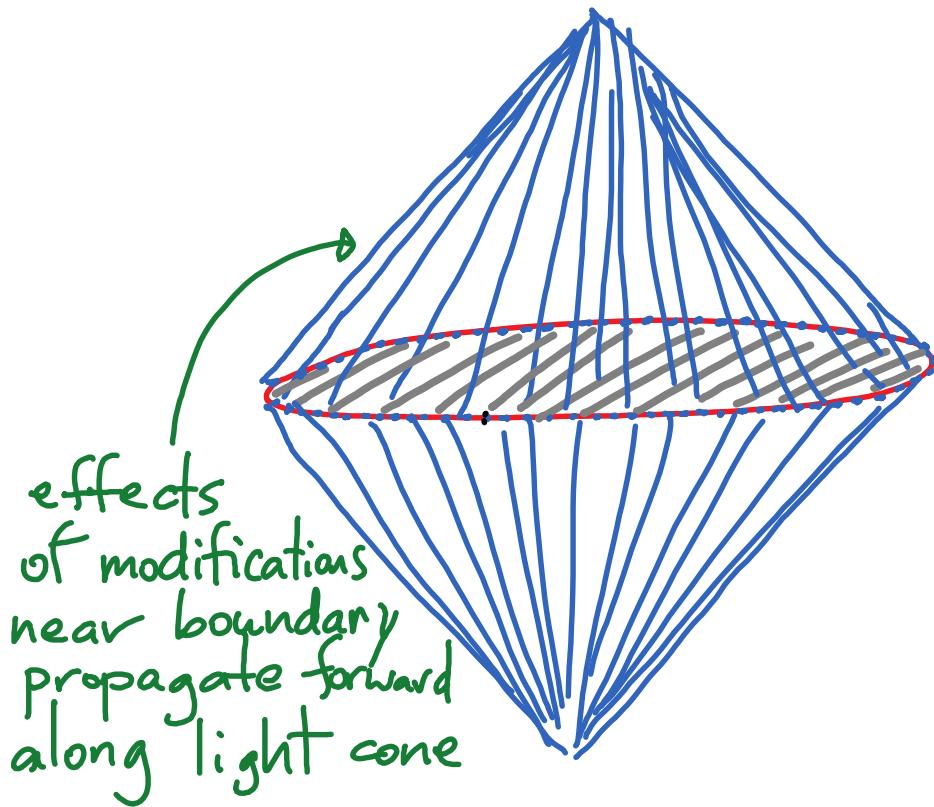


# Effects on Lorentzian geometries

original spacetime:



new spacetime:



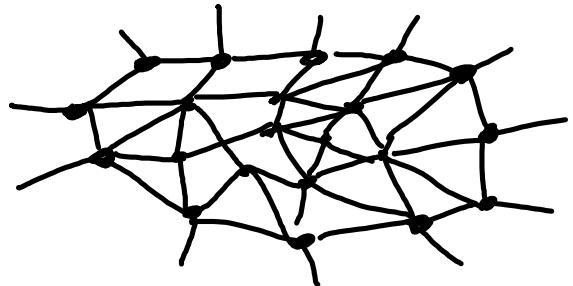
\* Spacetime inside "diamond" is arbitrarily close to previous spacetime \*

Key message: we can describe very general connected spacetimes via multipart entangled states of many non-interacting systems.

- \* Almost no information about spacetime geometry in state of individual "bits"
- \* Removing entanglement manifestly destroys spacetime

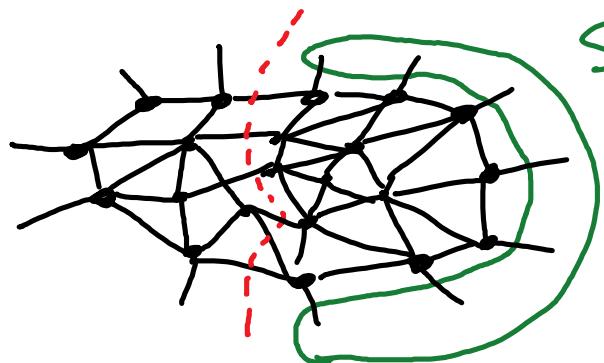
Interesting question: does it matter what the bits are?

## EXTRA: Connection to tensor network models of holography.



Tensor network states based on perfect tensors or random tensors for large bond dimension have been shown to obey a Ryu-Takayanagi formula:

(Pastawski, Yoshida, Harlow, Preskill; Hayden et. al.)



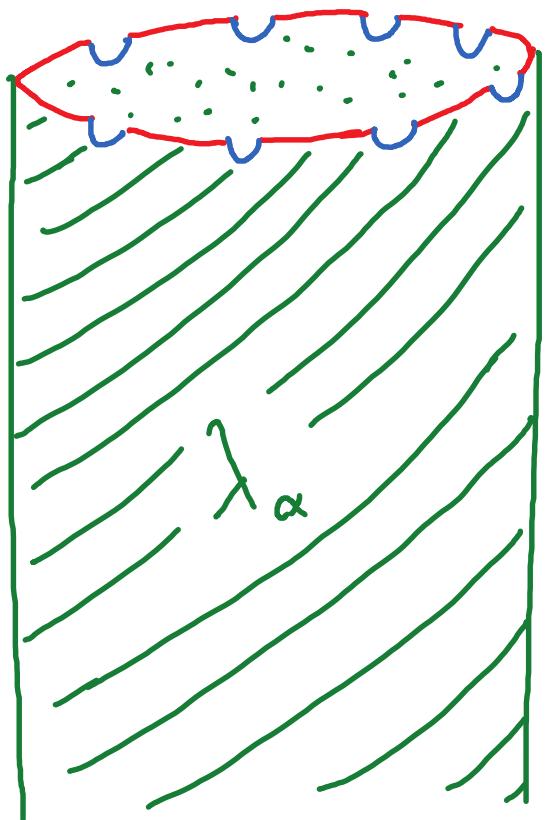
Subsystem A

$$S(A) \sim \min(\# \text{ of bonds cut} \text{ to separate } A)$$

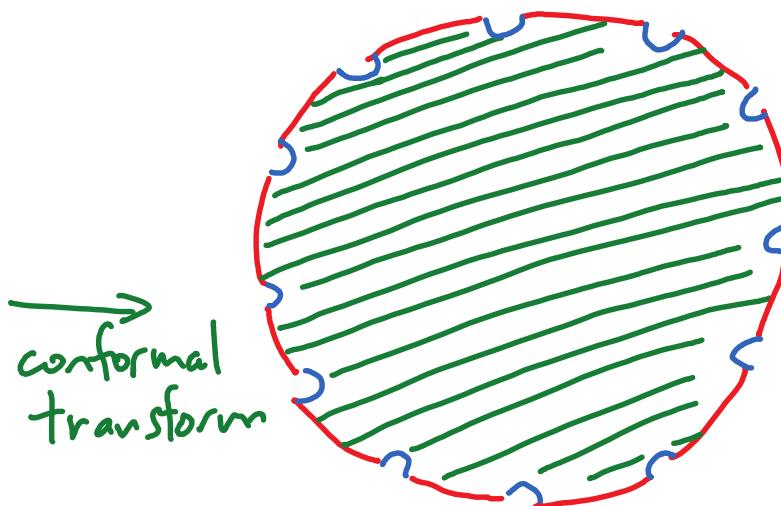
Suggests that tensor network states provide good toy model for holographic states, with network geometry playing role of dual spatial geometry (Swingle)

- Hard to understand if CFT states have this structure since CFT is continuous system.
- With our discrete version, can make a more direct connection to tensor networks:

Start with:

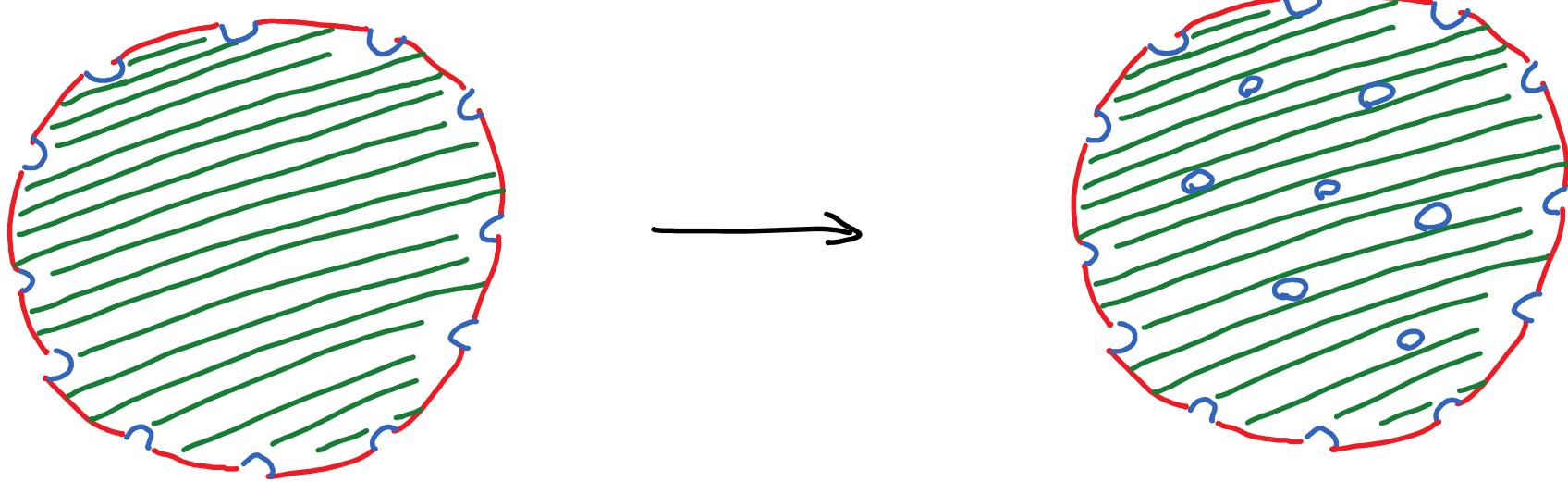


Equivalent to:



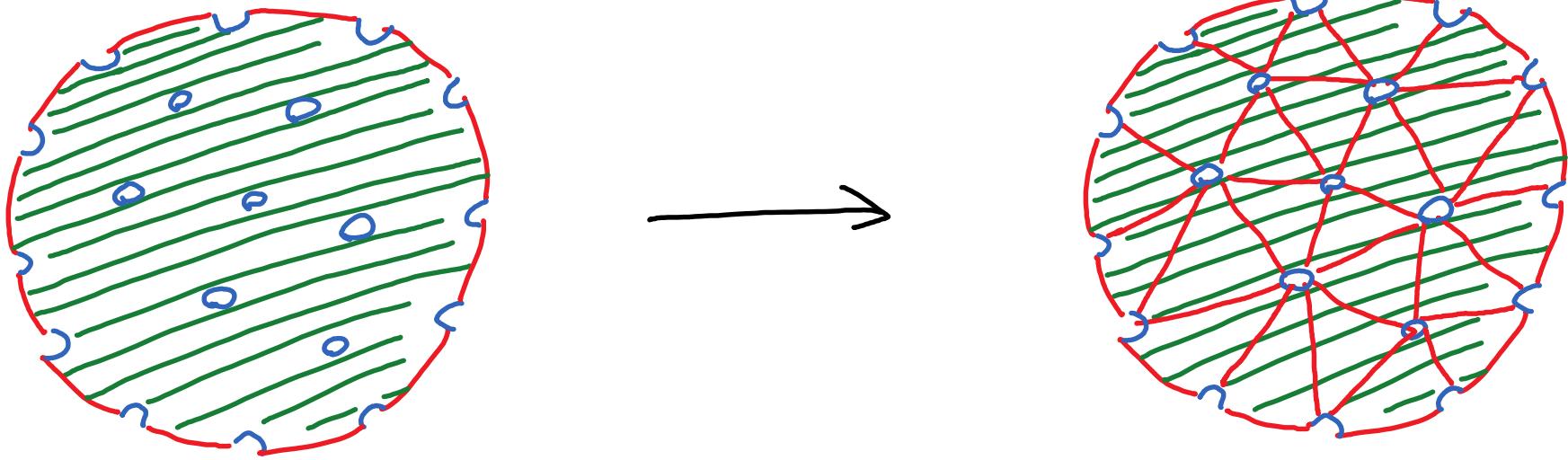
(this is the  
geometry +  
our path  
integral)

Add some extra boundary components:

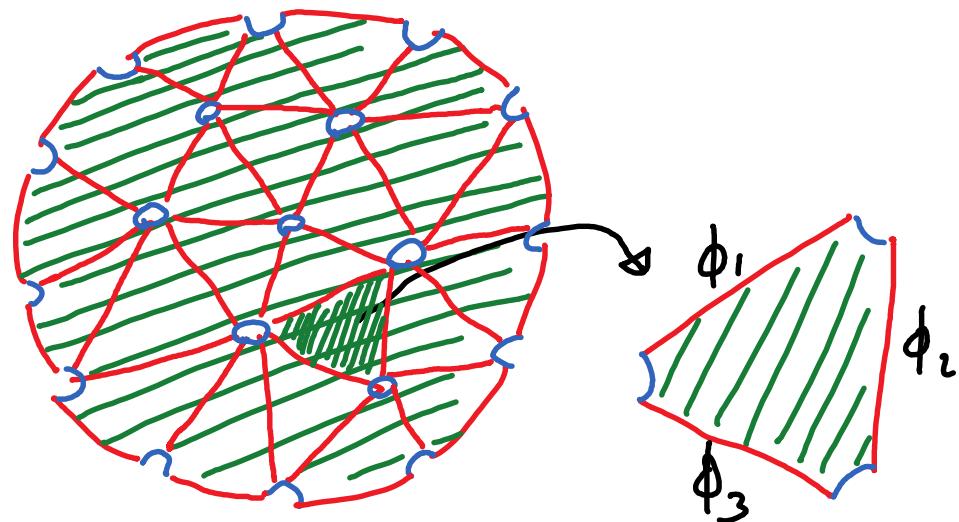


Negligible effect on state in limit of small size

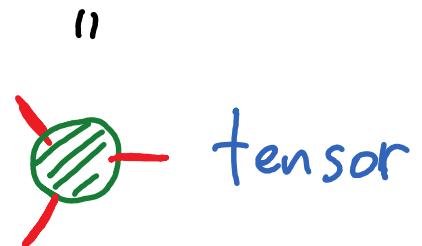
Divide up path integral into pieces:



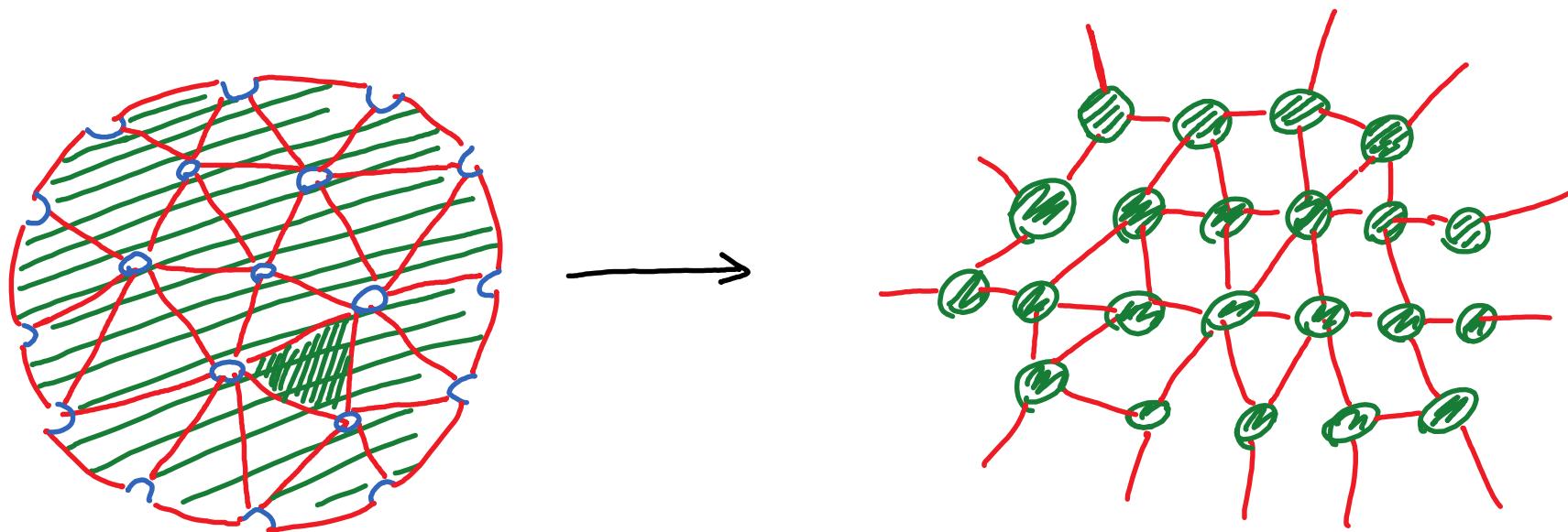
Each piece can be understood as path integral for a state of a small number of CFTs with boundary.



$$\langle \phi_1, \phi_2, \phi_3 | \Psi \rangle = \int [d\phi] e^{-S}$$



Full path integral corresponds to contracting up these states into tensor network (as in PEPS)



- Many possible networks since many choices for where to place extra boundaries+divisions
- Is there a rule that gives a network whose geometry matches geometry of the dual space in gravity theory?  
(c.f. Kaputa, Kundu, Miyagi  
Takayanagi, Watanabe: Path integral optimization)

Thank you!