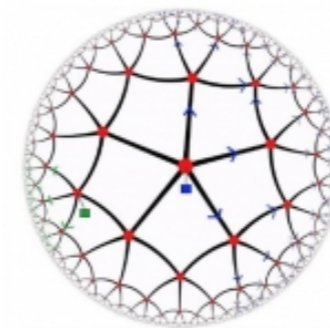
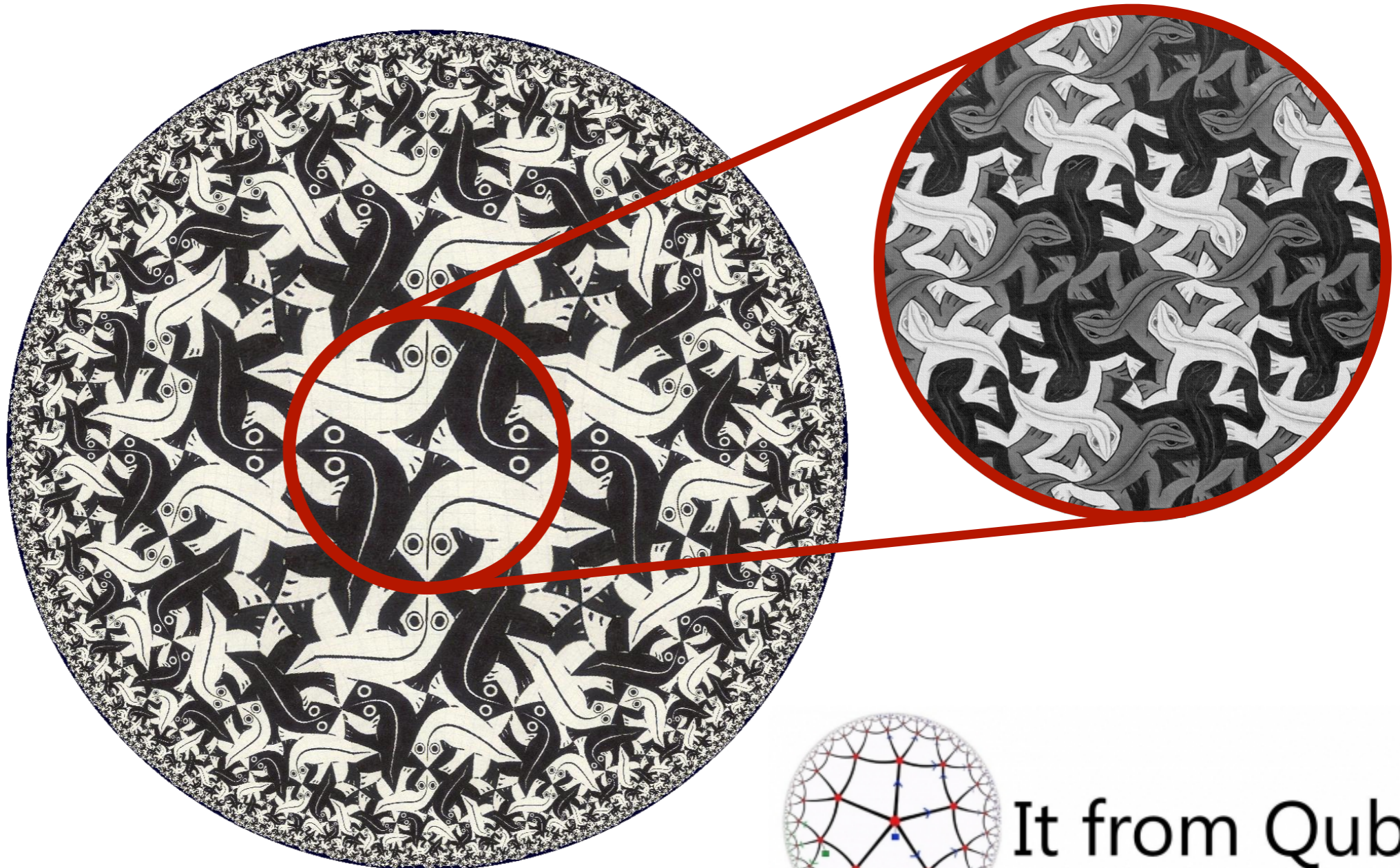




# Flat space physics from AdS/CFT

Eliot Hijano

Based on arxiv:1905.02729



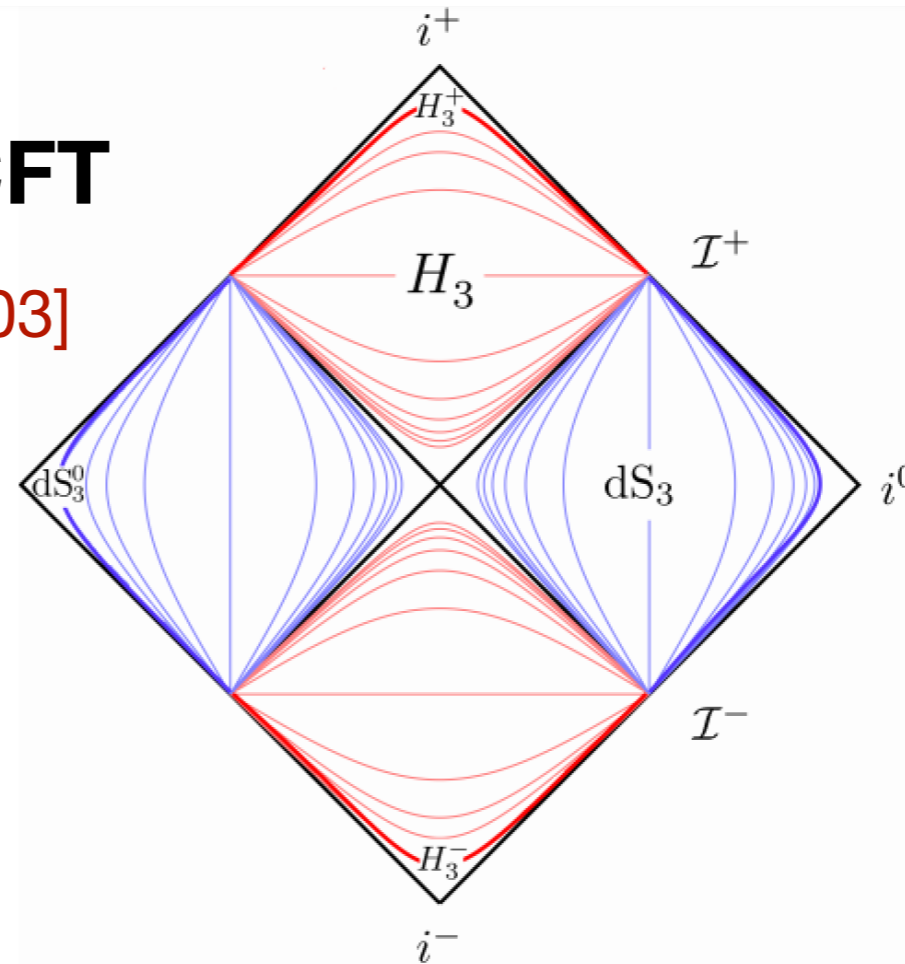
**It from Qubit**  
Simons Collaboration on  
Quantum Fields, Gravity and Information

# Approaches to flat holography

## Uplifting A(dS)/CFT

[de Boer, Solodukhin '03]

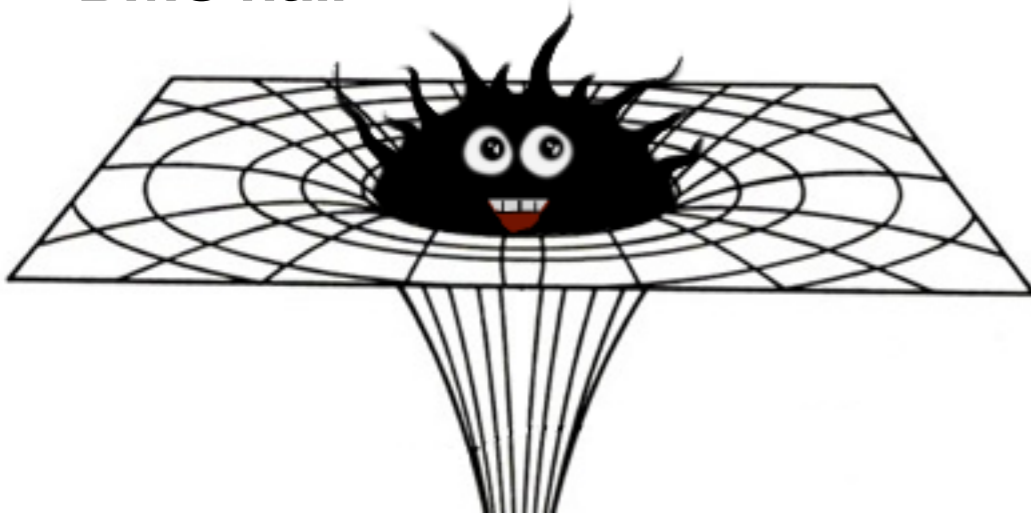
[Ball et al '19]



## Direct approach

[Hawking, Perry, Strominger '16]

## BMS hair



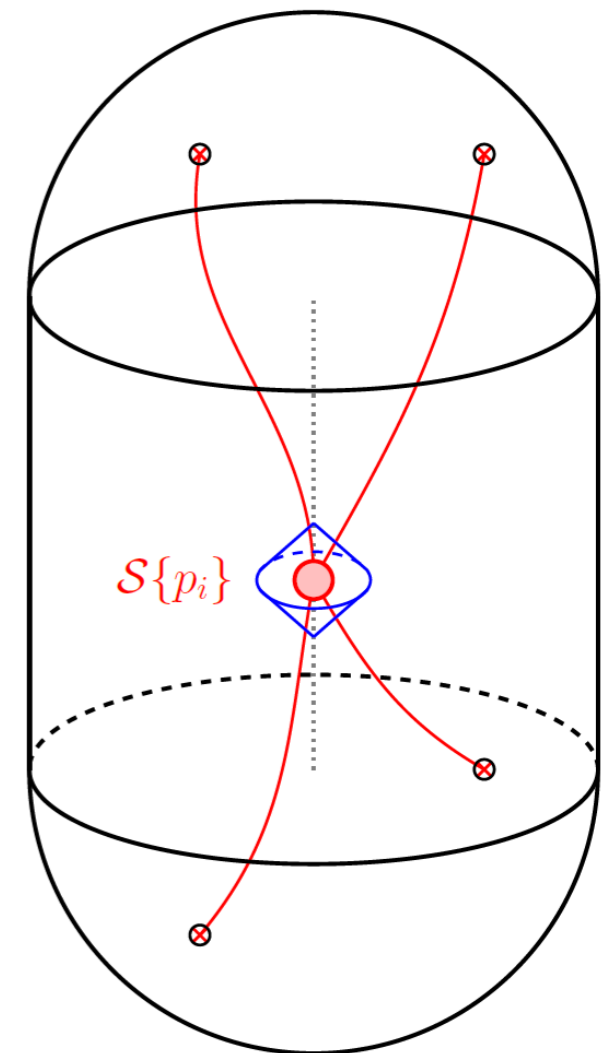
## Indirect approach

[Penedones '10]

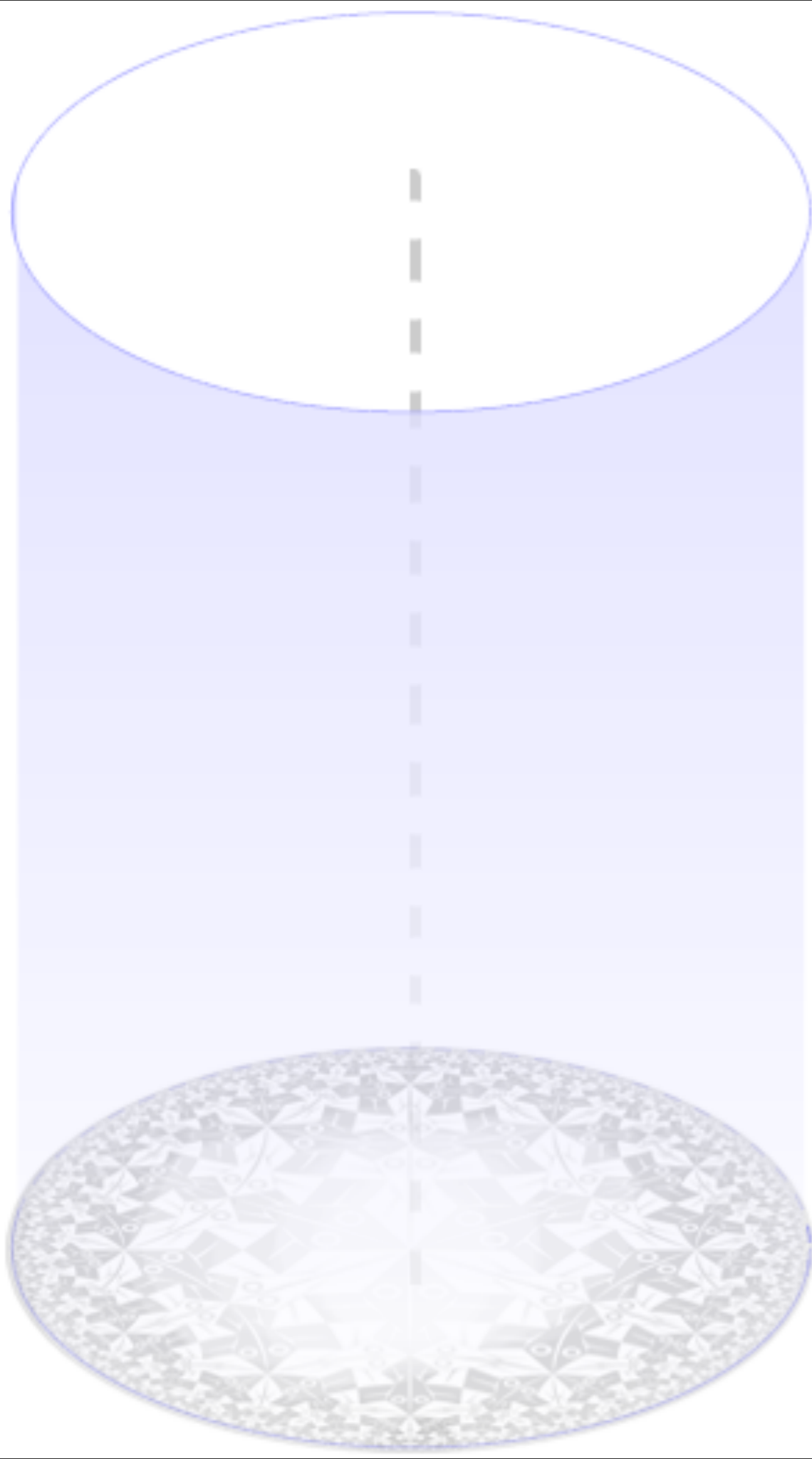
[Fitzpatrick, Kaplan '11]

[Paulos et al '17]

[EH, '19]

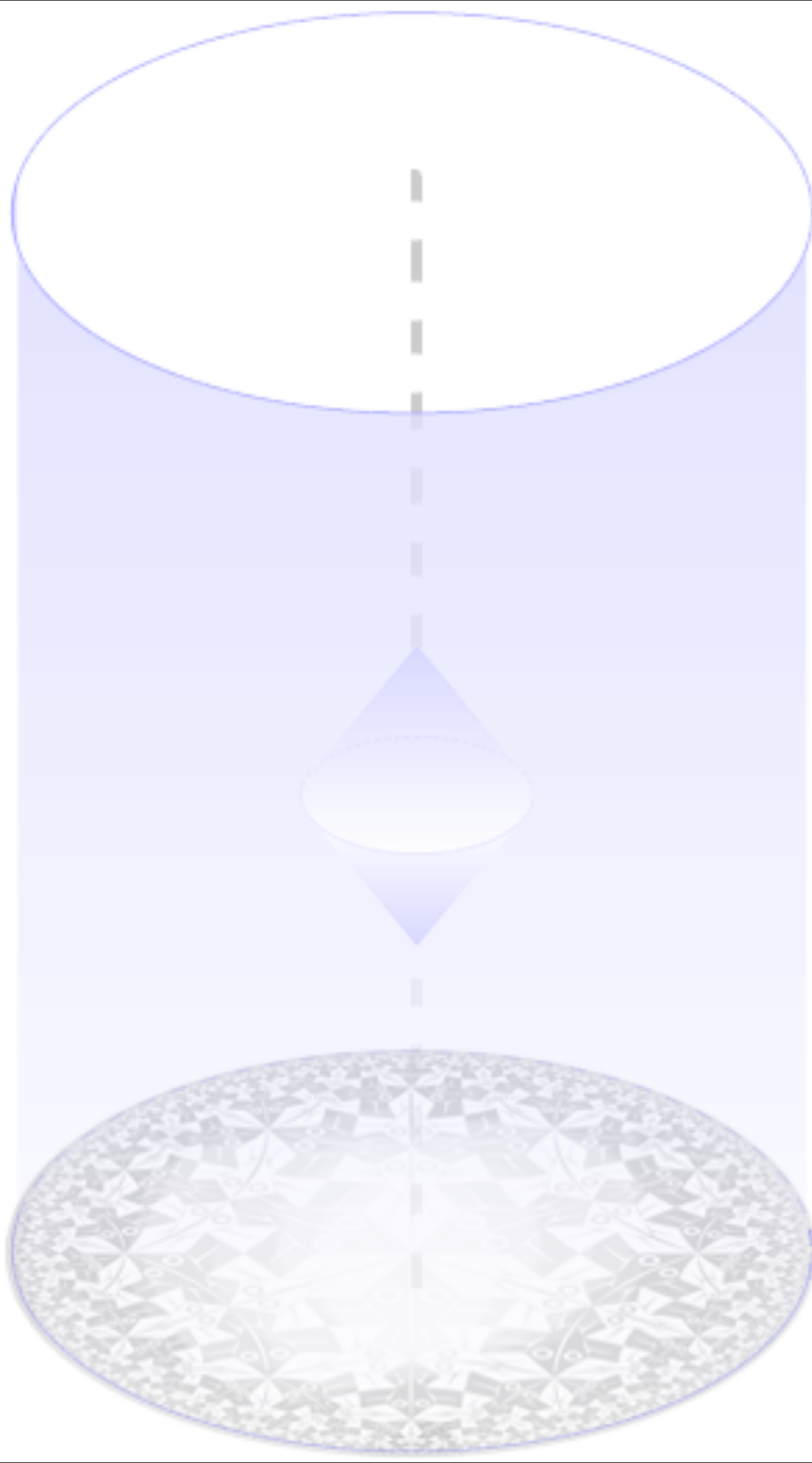


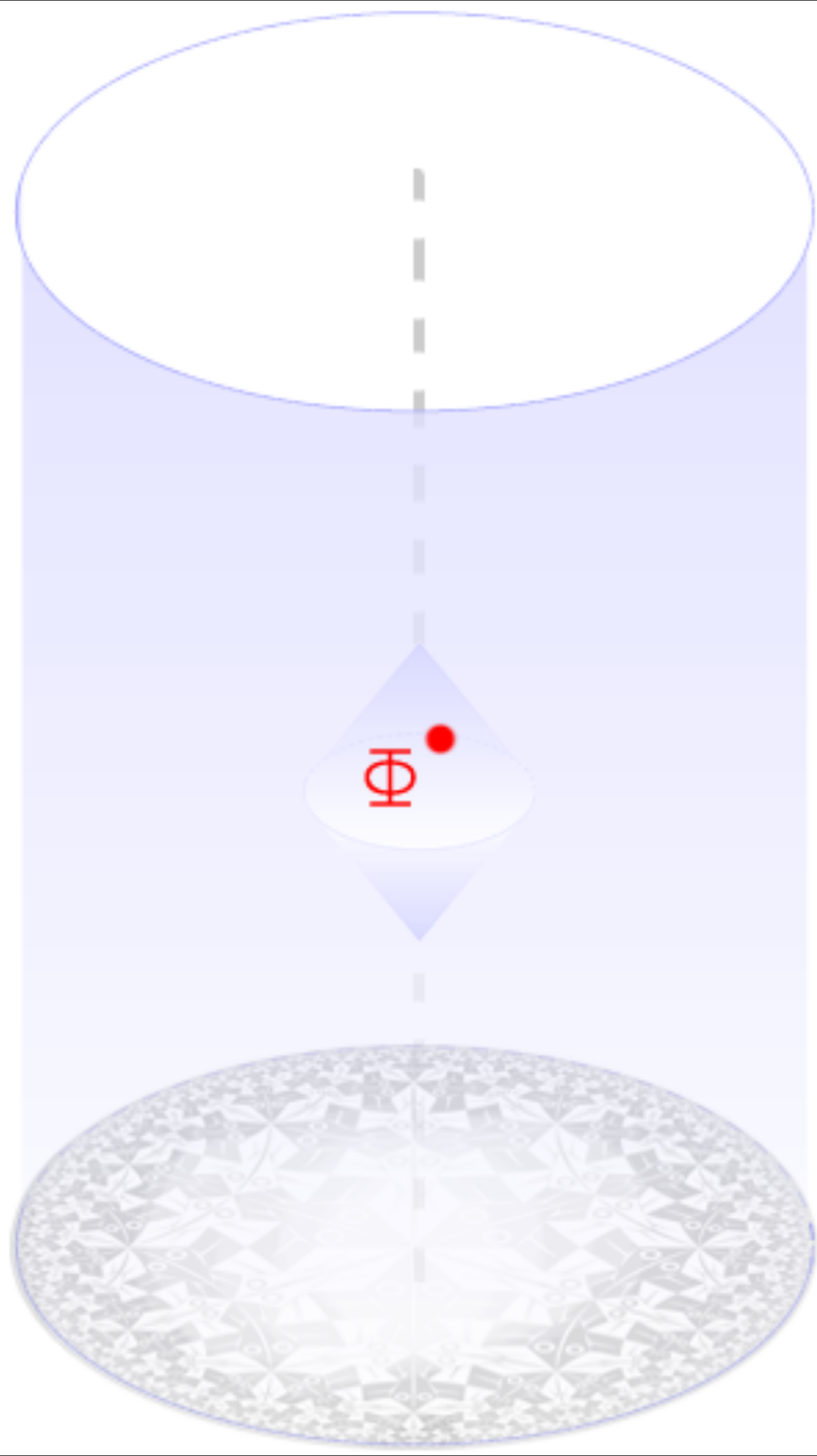
Start with Global AdS



Start with Global AdS

Define a scattering region at the center

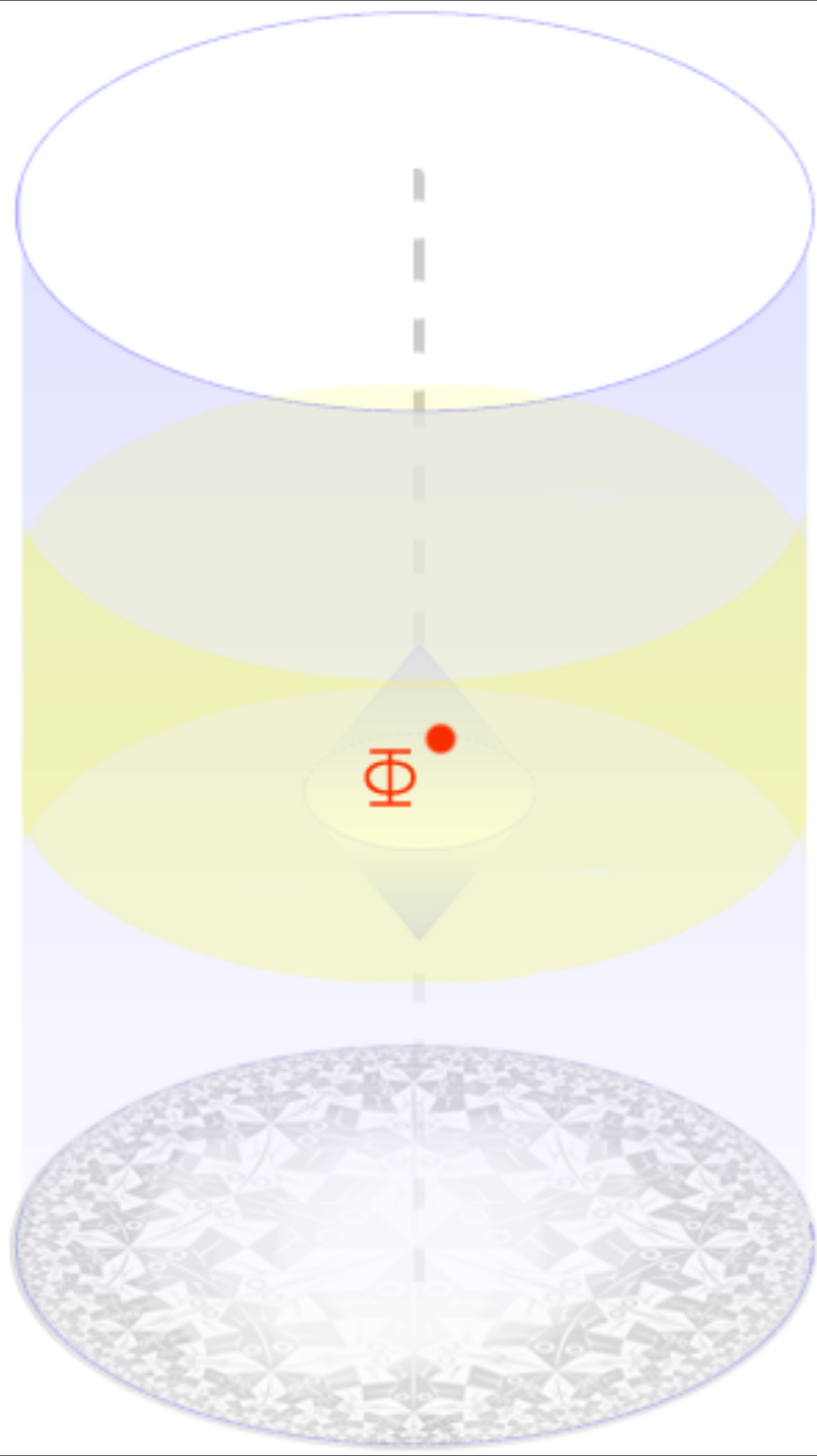




Start with Global AdS

Define a scattering region at the center

Consider a **local bulk operator** in the scattering region



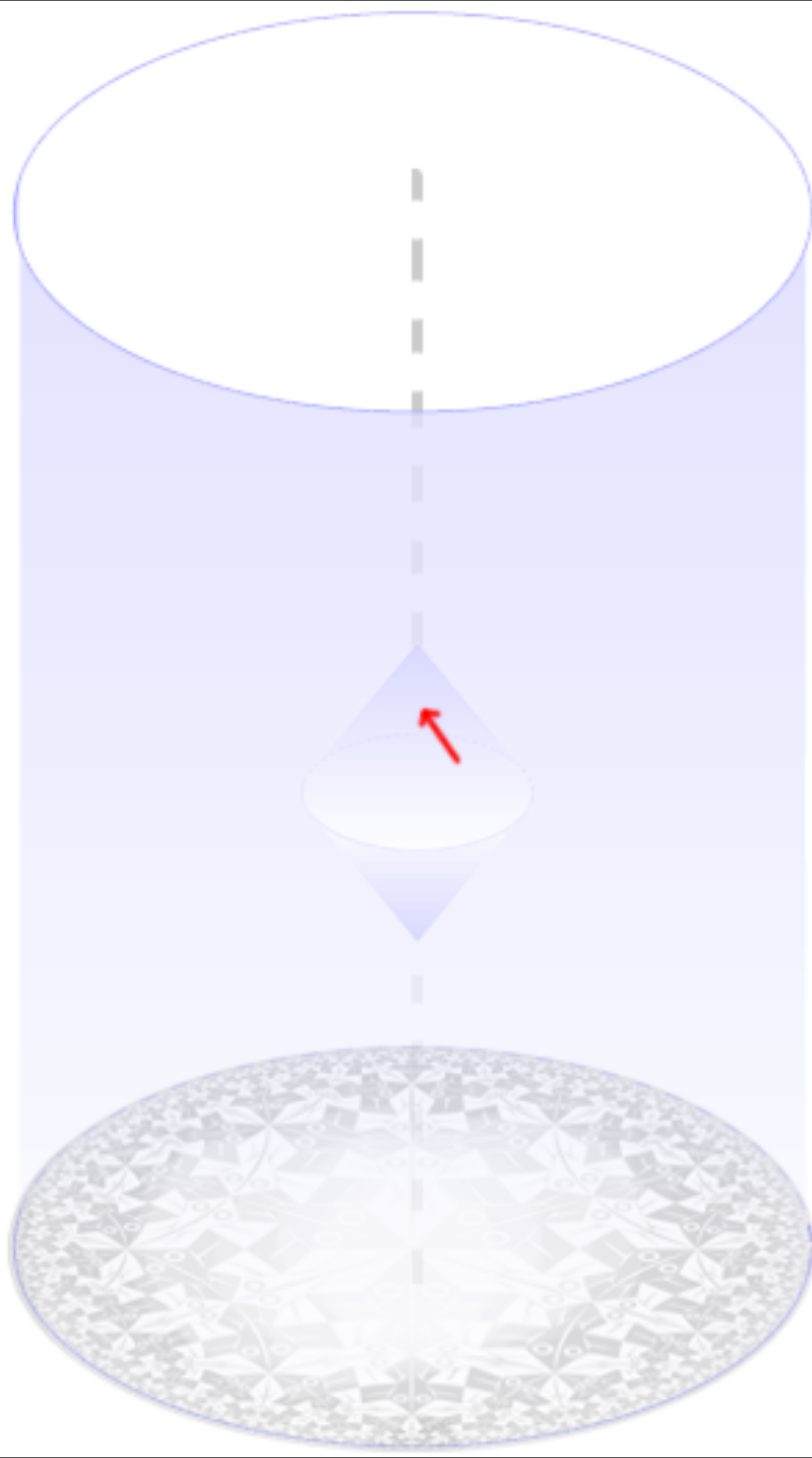
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Reconstruct the operator at the conformal boundary (HKLL)

$$\Phi(X) = \int_{\partial\text{AdS}_{d+1}} d^d x K_{\text{HKLL}}(X, x) \mathcal{O}(x).$$



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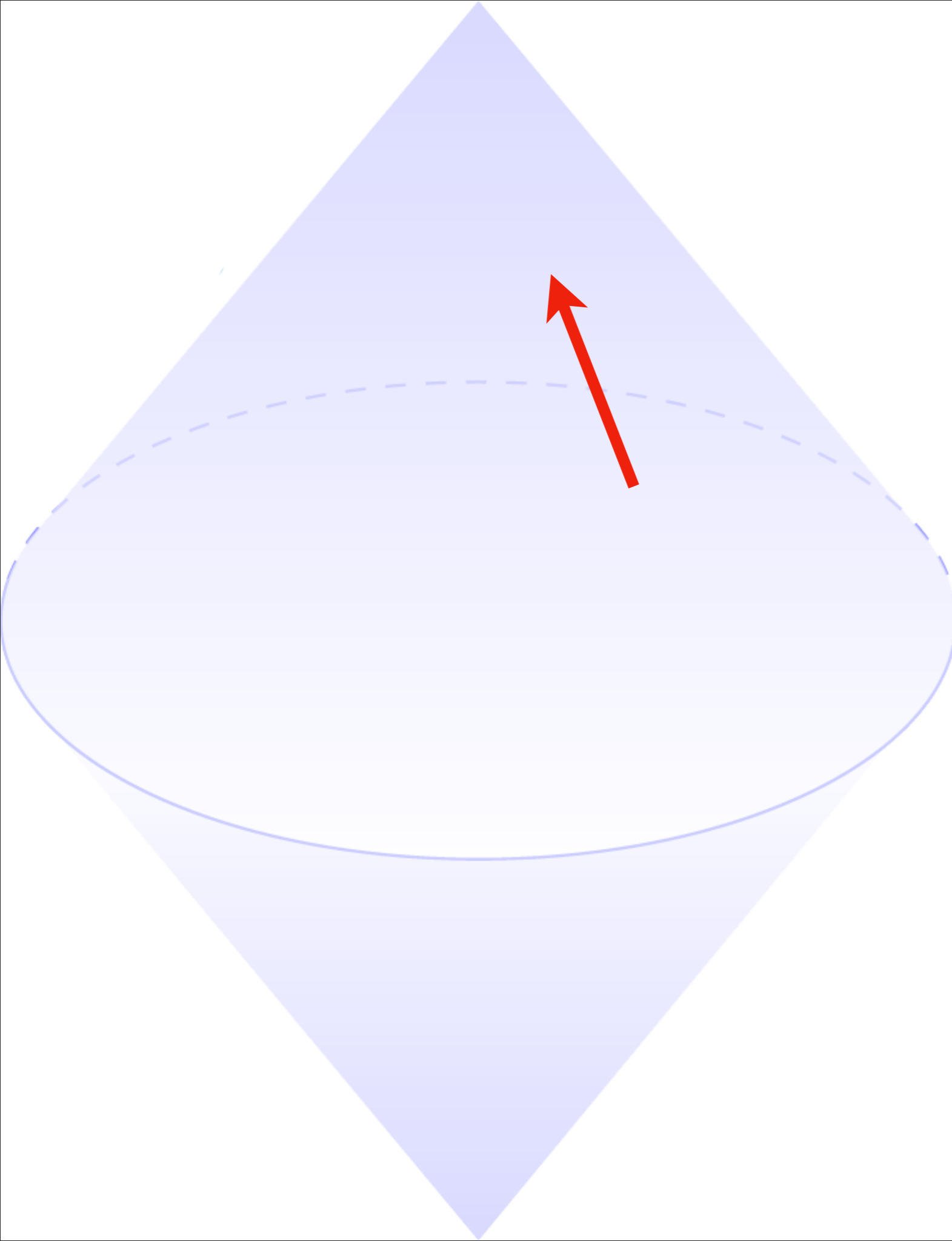
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Fourier transform = operator in momentum space.

$$\Phi(p) = \int d^D X \Phi(X) e^{ip \cdot X}$$



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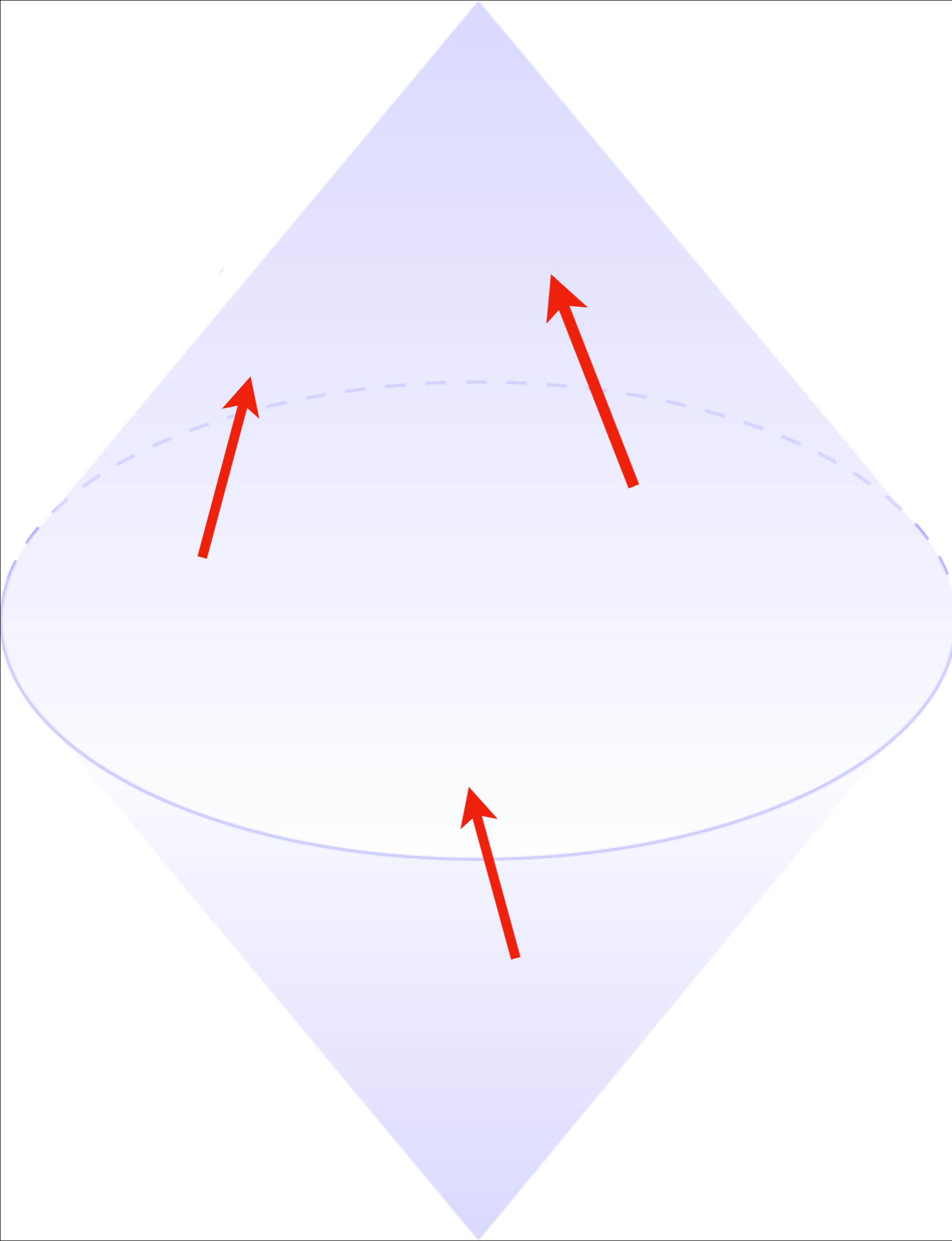
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Large AdS radius limit zooms into the scattering region





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Several insertions = Scattering amplitudes involving multiple particles

## Scattering amplitude:

$$\mathcal{S}\{p_i\} = \lim_{l \rightarrow \infty} l^{\frac{d-3}{2}} \left[ \prod_i C(p_i) \int dt_i e^{\pm i\omega_i t_i} \right] \langle 0 | \mathcal{O}(\tau_1, \Omega_1) \cdots | 0 \rangle$$

$$\tau_i = \pm \frac{\pi}{2} \pm i \cosh^{-1} \frac{\omega_i}{k_i} + \frac{t_i}{l}, \quad \text{and} \quad \Omega_i - \chi_i = \frac{\pi}{2} \mp \frac{\pi}{2}.$$

Reduces to existing formulae in the literature when all particles are either simultaneously massive or massless.

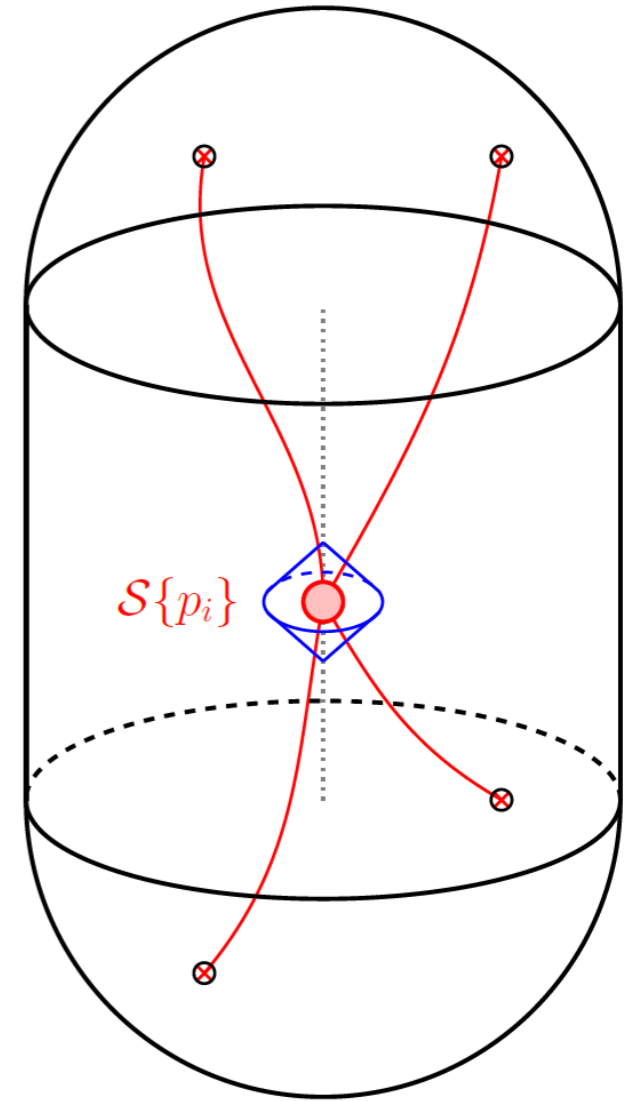
## Some easy examples in $\text{AdS}_{d+1}$ :

$$\mathcal{S}\{p_1, p_2\} \sim \delta^{(d)}(p_1 + p_2),$$

$$\mathcal{S}\{p_1, p_2, p_3\} \sim \delta^{(d)}(p_1 + p_2 + p_3).$$

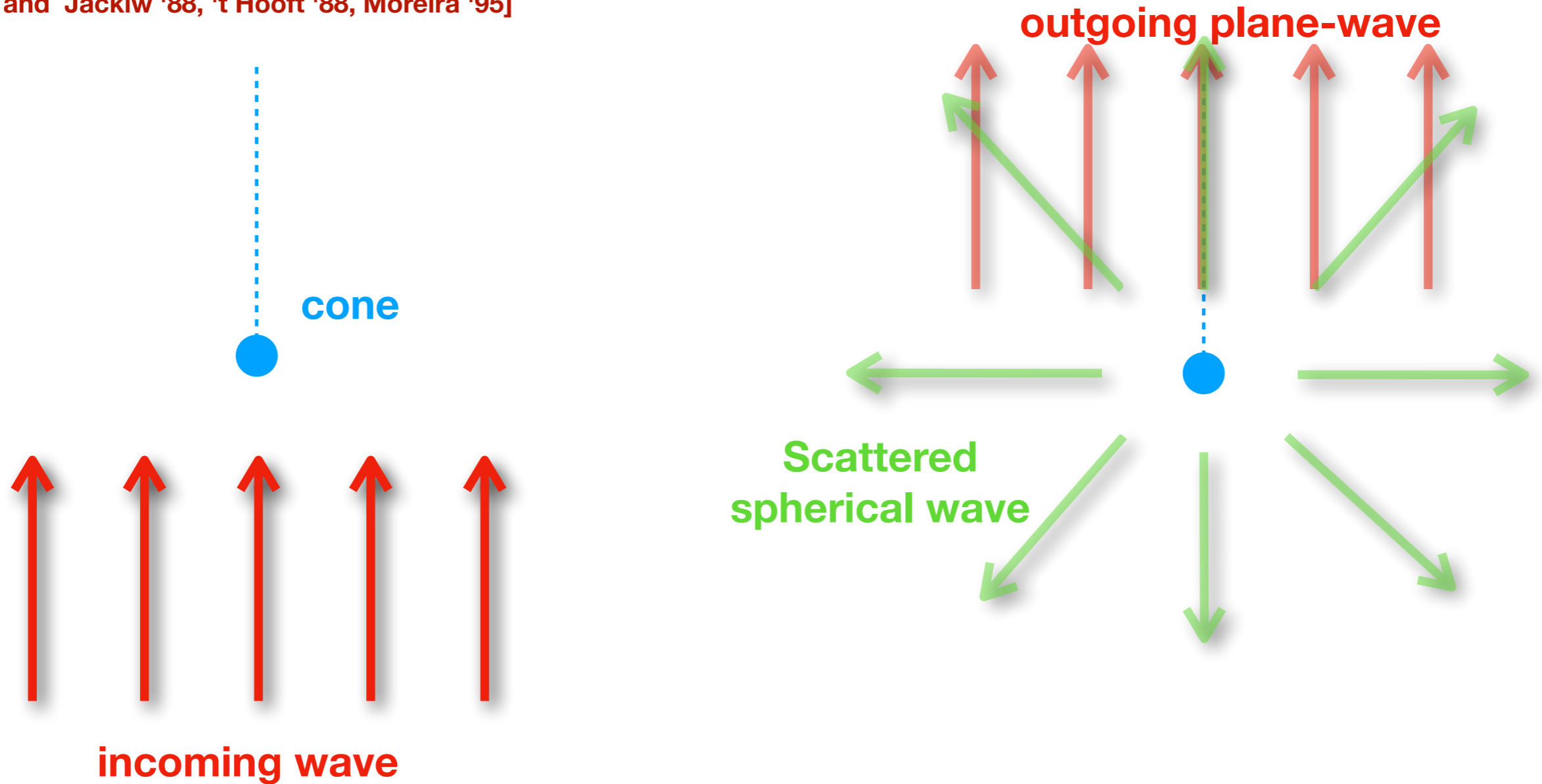
## $\text{BMS}_3$ global block from $\text{CFT}_2$ global block:

$$\mathcal{S}\{p_1, p_2, p_3, p_4\}_{p'^2 = -m'^2} \sim \delta^{(3)}(p_1 + p_2 + p_3 + p_4) \delta(s + m'^2).$$



# Scattering against a cone (D=2+1)

[Deser and Jackiw '88, 't Hooft '88, Moreira '95]



$$\mathcal{S}\{p, p'\} = \delta^{(3)}(p + p') + \delta(\omega + \omega') \frac{\sin \frac{\pi}{\alpha}}{\cos \frac{\pi}{\alpha} - \cos(\phi - \phi')}$$

# Same result from the flat limit of a CFT<sub>2</sub> correlator

**CFT deficit state**

$$|\alpha\rangle_{\text{CFT}_2} = \mathcal{O}_{\Delta=\frac{c}{12}(1-\alpha^2)}|0\rangle_{\text{CFT}_2}$$

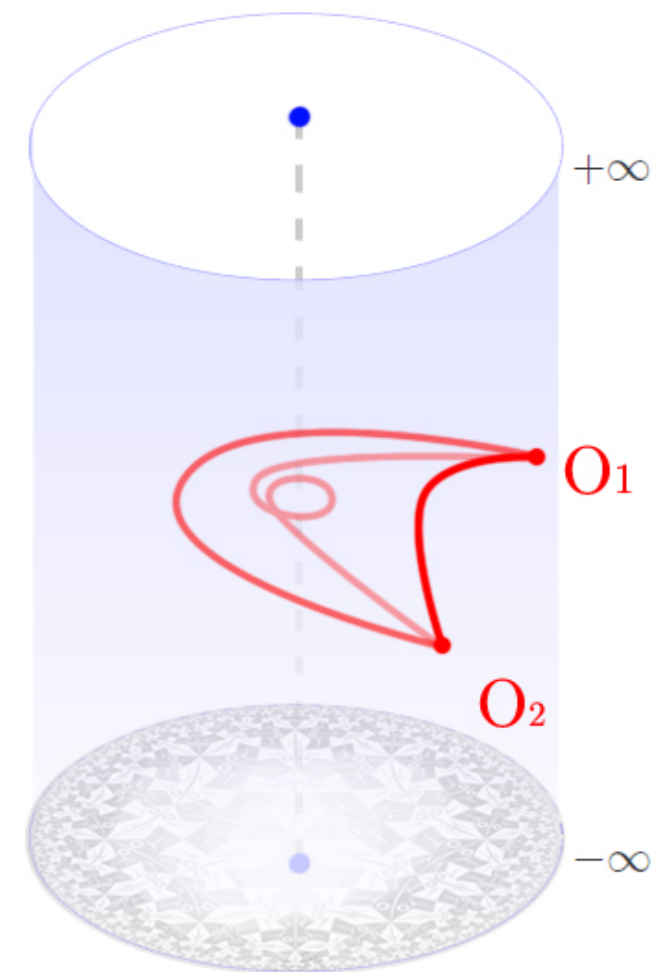
**Dual to a conical deficit  
AdS<sub>3</sub> geometry**

$$ds^2 = \frac{l^2}{\cos^2 \rho} (d\rho^2 - \alpha^2 d\tau^2 + \alpha^2 \sin^2 \rho^2 d\phi^2)$$

$$\mathcal{S}\{p_1, p_2\} = \lim_{l \rightarrow \infty} l^{\frac{d-3}{2}} \left[ \prod_{i=1}^2 C(p_i) \int dt_i e^{-i\omega_i t_i} \right] \langle \alpha | \mathcal{O}(\tau_1, \chi_1) \mathcal{O}(\tau_2, \chi_2) | \alpha \rangle$$

$$\mathcal{S}\{p, p'\} = \delta^{(3)}(p + p') + \delta(\omega + \omega') \frac{\sin \frac{\pi}{\alpha}}{\cos \frac{\pi}{\alpha} - \cos(\phi - \phi')}$$

Non-trivial CFT<sub>2</sub> correlators turn into non-trivial scattering events in asymptotically flat geometries.



# Current/Future Work

- Soft theorems. How do they arise from flat limits of conformal Ward identities? How do soft theorems look like in 2+1 dimensions?
- Exploration of scattering events in 4D black hole backgrounds. What features of conformal correlators imply unitarity of black hole evolution?  
't Hooft S-matrix ansatz:  $S = S_{in} S_{hor} S_{out}$
- Implementation of CFT bootstrap programme. [Paulos et al '17]
- AdS/CFT: Are the divergences of the correlators that turn into S-matrices a further diagnostic of bulk locality?

ありがとうございます!