

Ultra light axion dark matter and small scale problems

Kohei Hayashi (JSPS fellow)
ICRR, The University of Tokyo

Contents

1. Introduction

1.1 Cold dark matter theory and dwarf spheroidal galaxies

1.2 Small scale problems in Λ CDM models

2. Ultralight axion dark matter as a solution to small scale problems

3. Non-sphericity of ultralight axion dark matter halos in the Galactic dwarf satellites

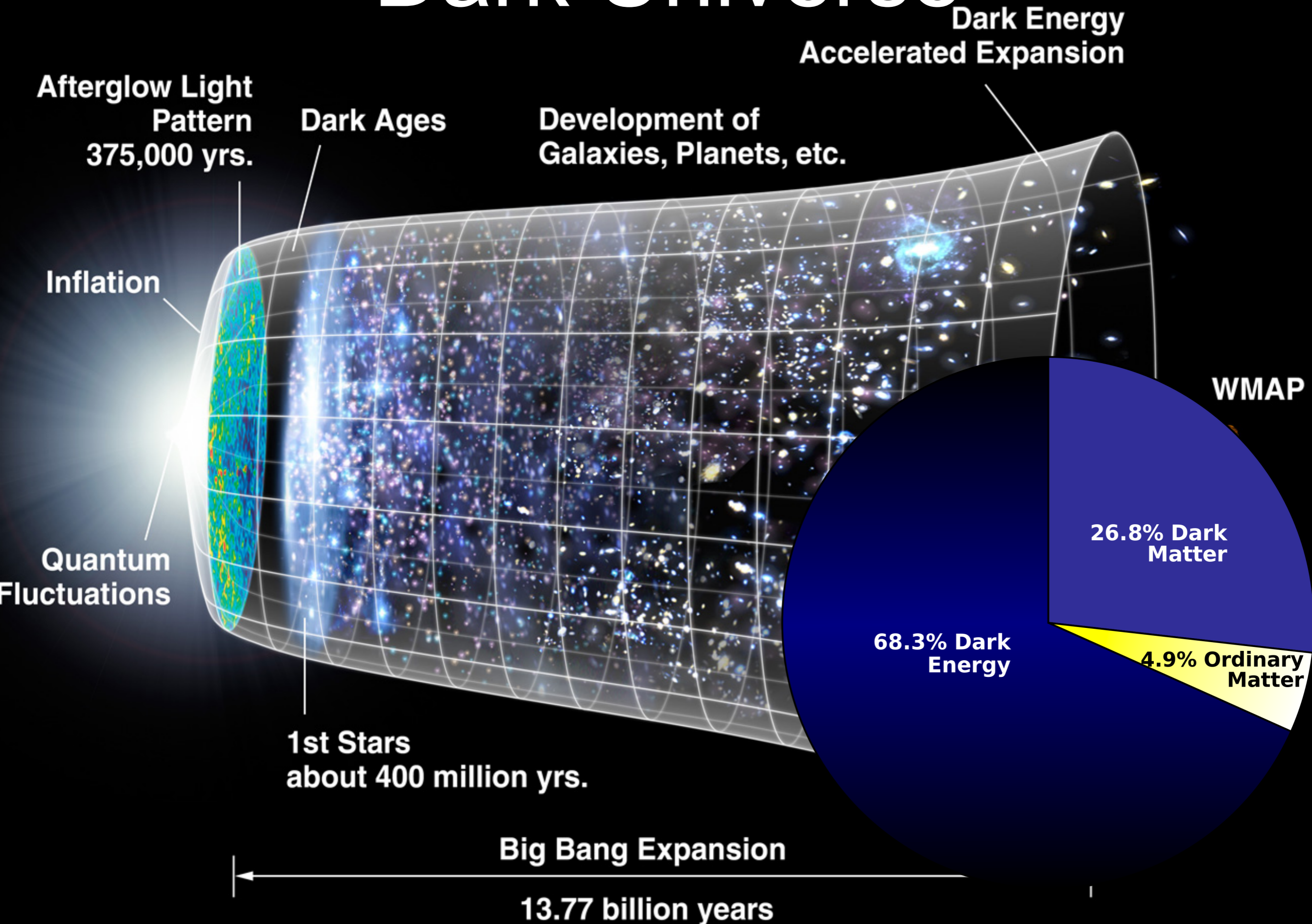
4. Revisit core-cusp problem

5. Summary

Introduction

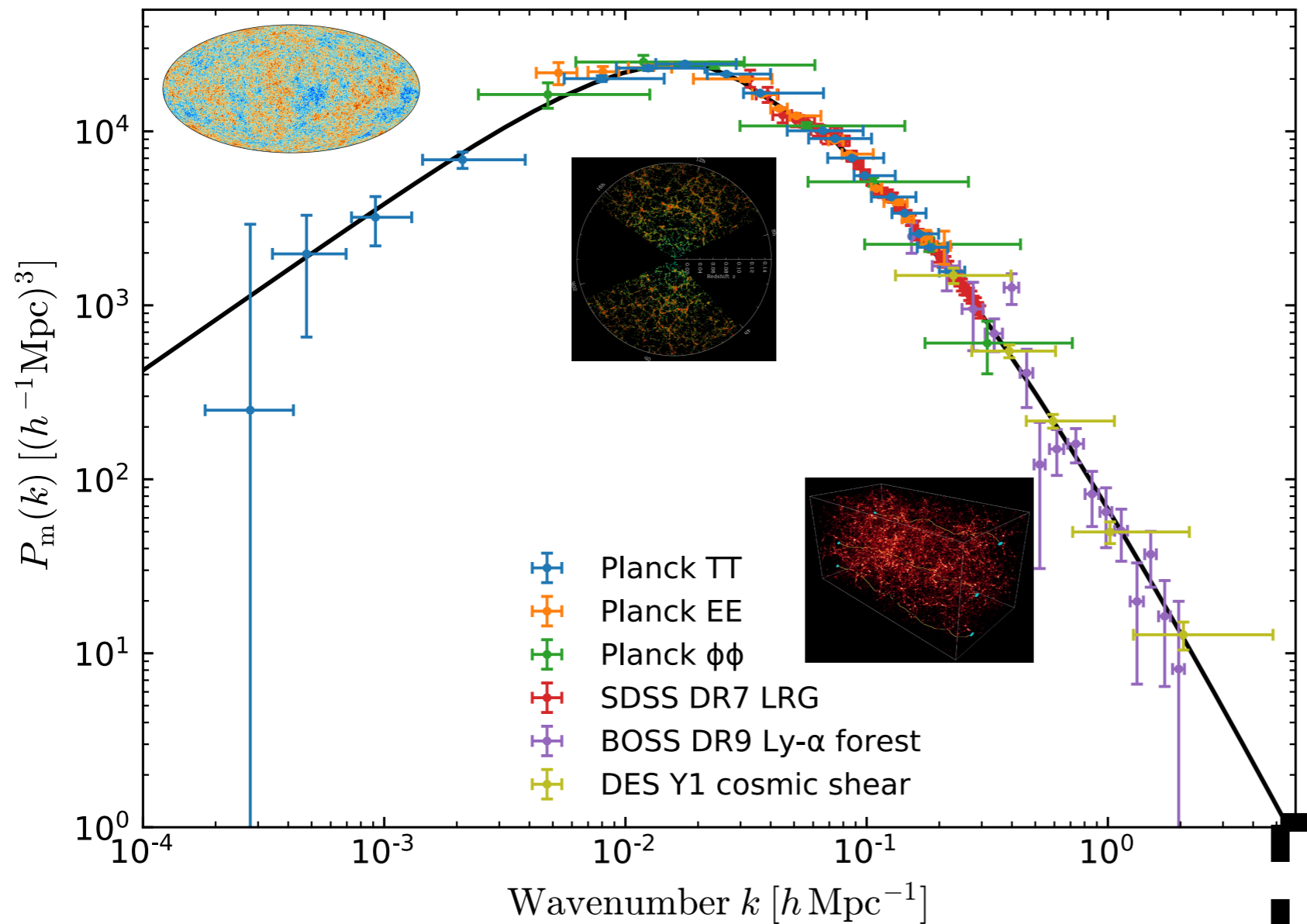
**Cold dark matter theory
and
dwarf spheroidal galaxies**

Dark Universe



Λ Cold Dark Matter model

ESA and the Planck Collaboration

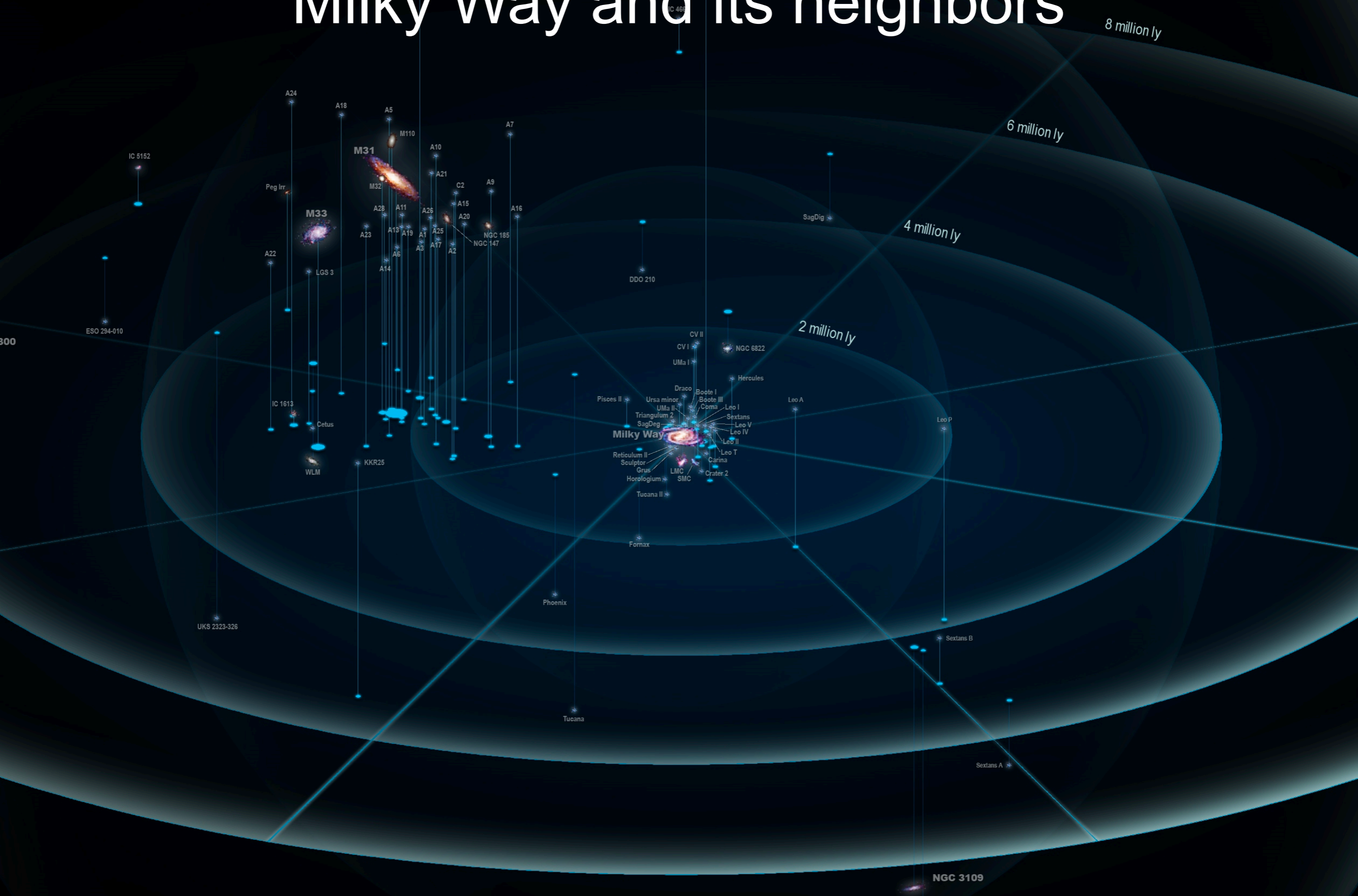


Large scale ($> 1 \text{Mpc}$)
 \Rightarrow Remarkable success!

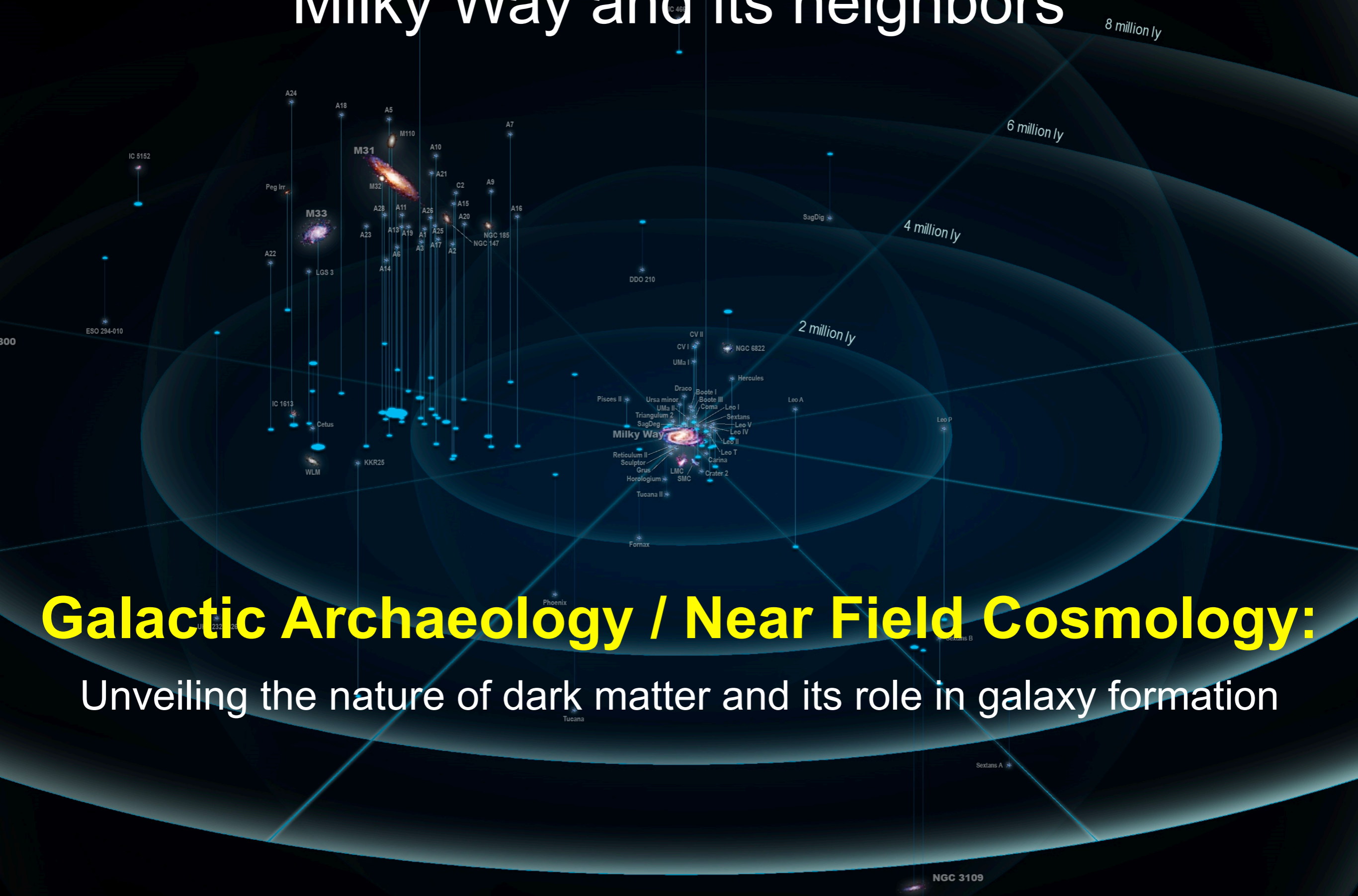
Small scale ($< 1 \text{Mpc}$)
 \Rightarrow What's going on?



The structures on small scales (<1Mpc): Milky Way and its neighbors



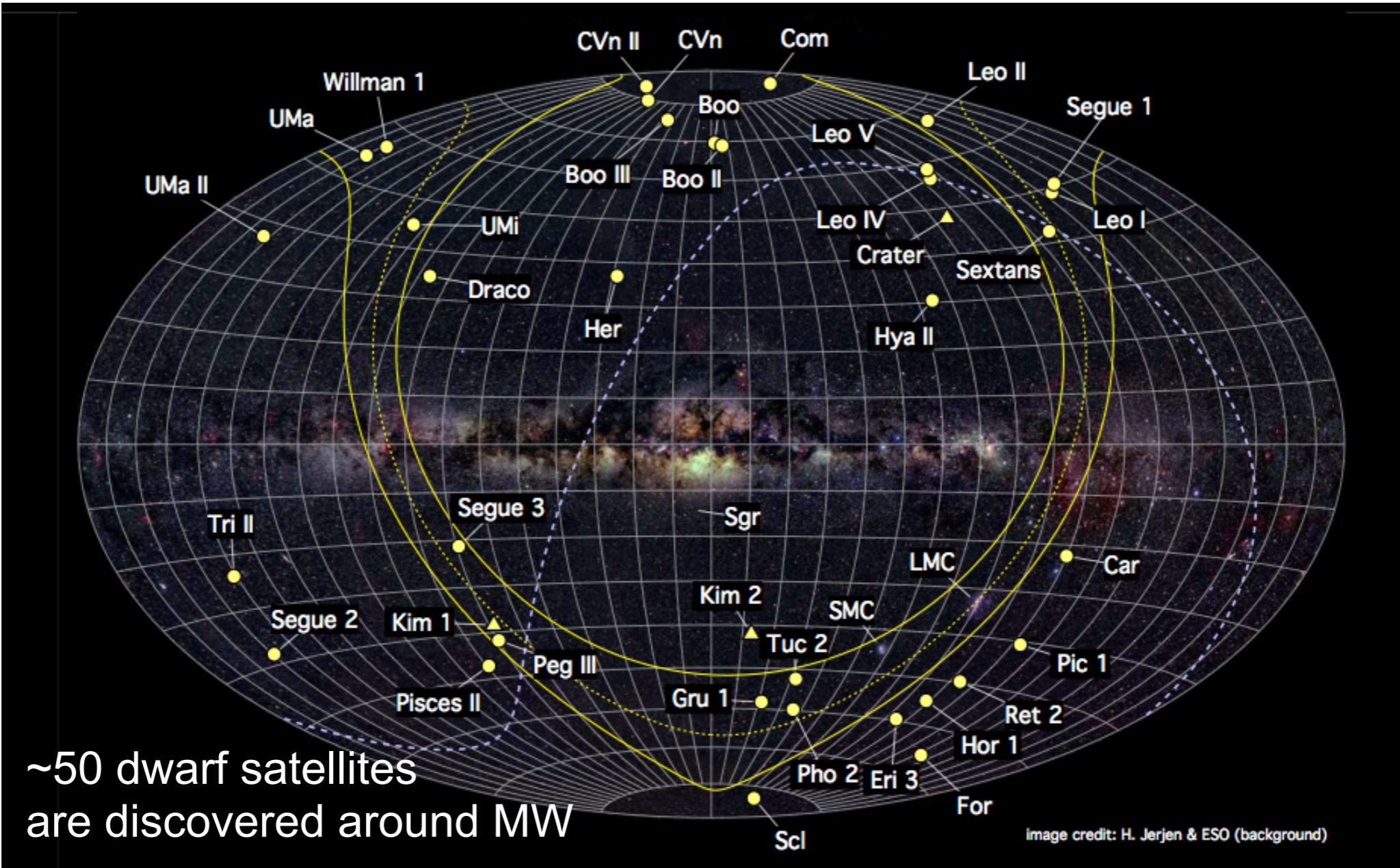
The structures on small scales (<1Mpc): Milky Way and its neighbors



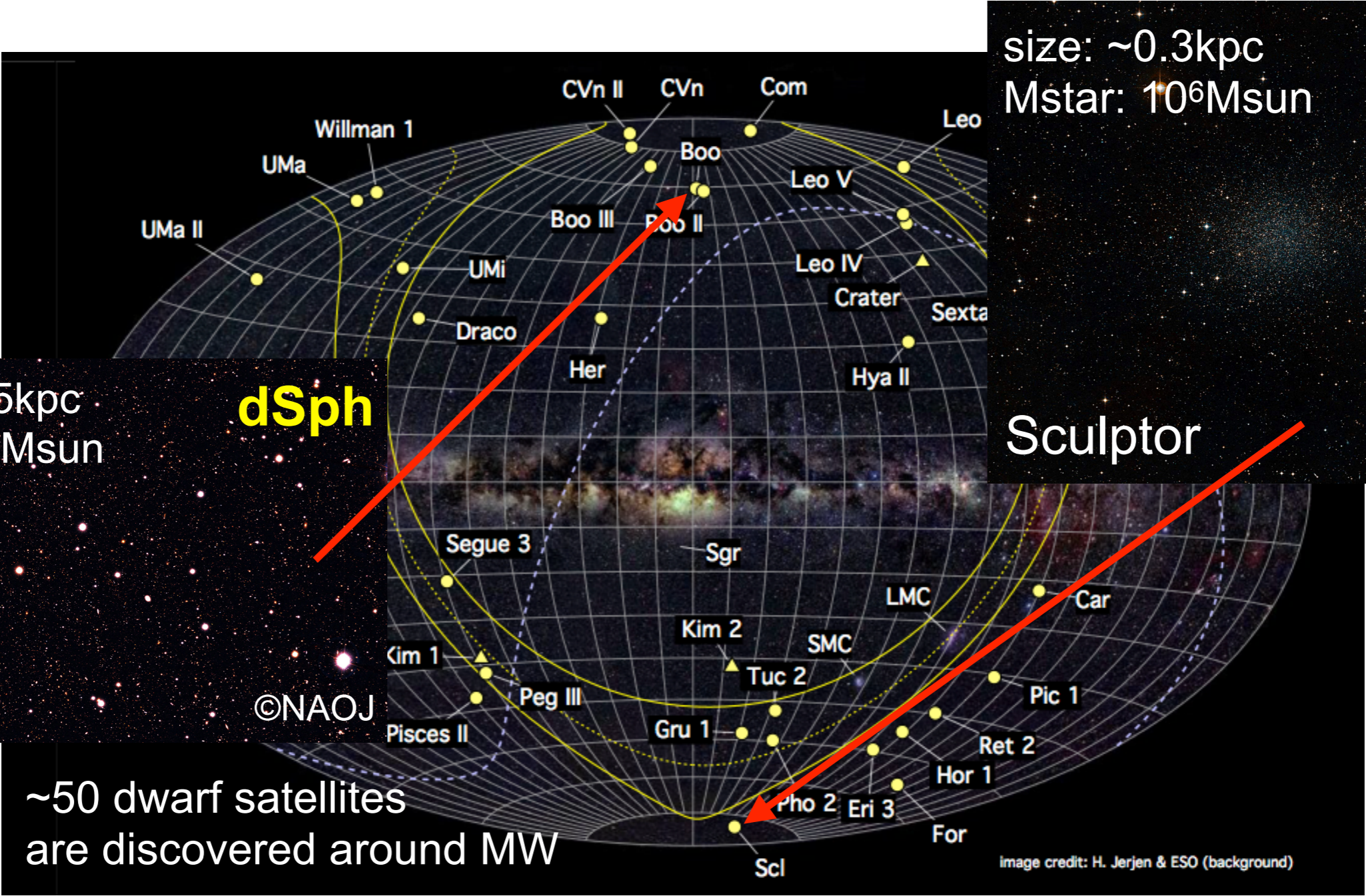
Galactic Archaeology / Near Field Cosmology:

Unveiling the nature of dark matter and its role in galaxy formation

Dwarf Spheroidal Galaxy (dSph): basic properties



Dwarf Spheroidal Galaxy (dSph): basic properties



size: ~0.25kpc
Mstar: 10⁴Msun

dSph

size: ~0.3kpc
Mstar: 10⁶Msun

dSph

Sculptor

Bootes I

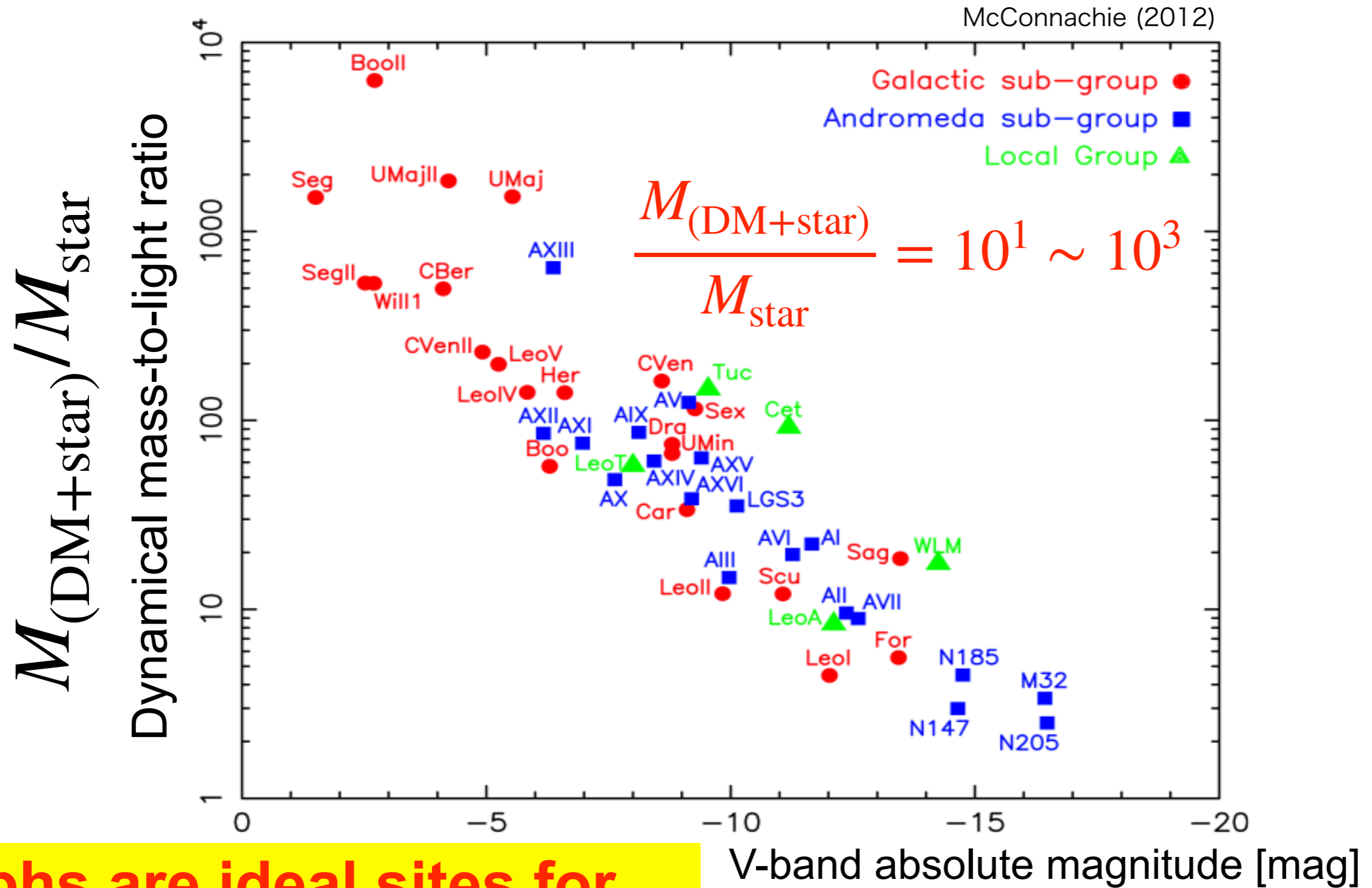
~50 dwarf satellites
are discovered around MW

©NAOJ

image credit: H. Jerjen & ESO (background)

❖ no gas, no current SF
❖ smallest and oldest galaxy ➡ **Fossil records of galaxy formation**

dSphs: dark-matter dominated system



DSphs are ideal sites for studying the nature of DM!

Small scale problems in Λ CDM models

Small-scale challenges to Λ CDM paradigm

Definition of “small scales”

$$M_{\text{virial}} < 10^{11} M_{\odot}, k > 3 \text{Mpc}^{-1}, r < 1 \text{Mpc}$$

$$\implies r_{\text{vir}} < 150 \text{kpc}, V_{\text{virial}} < 50 \text{km/s}$$

i.e., galaxy and dwarf-galaxy scales

Bullock &
Boylan-Kolchin (2017)

◆ **Missing satellite problem (Moore+99, Klypin+99)**

- Overabundance of dark subhalos

◆ **Core-cusp problem (de Blok 2002, Gilmore+ 07)**

- Cuspy central density in CDM halos vs. cores in observed galaxies

◆ **Too-big-to-fail problem (Boylan-Kolchin+ 11)**

- Most massive subhalos are more concentrated than observed luminous satellites

✦ **the other problems (Pawlowski+ 12, KH & Chiba 12)**

(satellite planes, shapes of dark halo, etc...)

Small-scale challenges to Λ CDM paradigm

Definition of “small scales”

$$M_{\text{virial}} < 10^{11} M_{\odot}, k > 3 \text{Mpc}^{-1}, r < 1 \text{Mpc}$$

$$\implies r_{\text{vir}} < 150 \text{kpc}, V_{\text{virial}} < 50 \text{km/s}$$

i.e., galaxy and dwarf-galaxy scales

Bullock &
Boylan-Kolchin (2017)

◆ **Missing satellite problem (Moore+99, Klypin+99)**

- Overabundance of dark subhalos

◆ **Core-cusp problem (de Blok 2002, Gilmore+ 07)**

- Cuspy central density in CDM halos vs. cores in observed galaxies

◆ **Too-big-to-fail problem (Boylan-Kolchin+ 11)**

- Most massive subhalos are more concentrated than observed luminous satellites

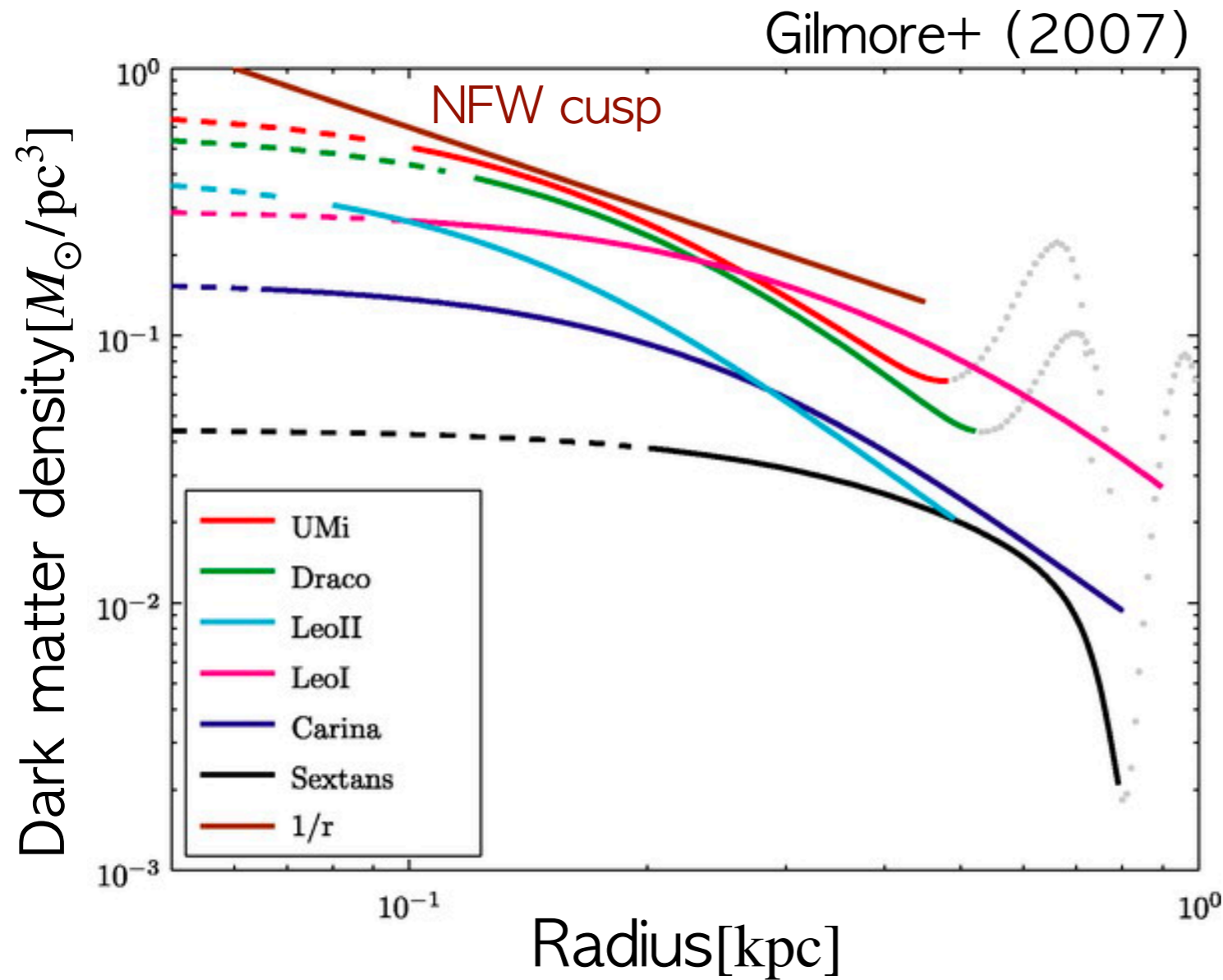
✦ **the other problems (Pawlowski+ 12, KH & Chiba 12)**

(satellite planes, shapes of dark halo, etc...)

Solution to core-cusp problem

◆ Core-cusp problem

- Cuspy central density in CDM halos vs. cores in observed galaxies



Possible solutions:

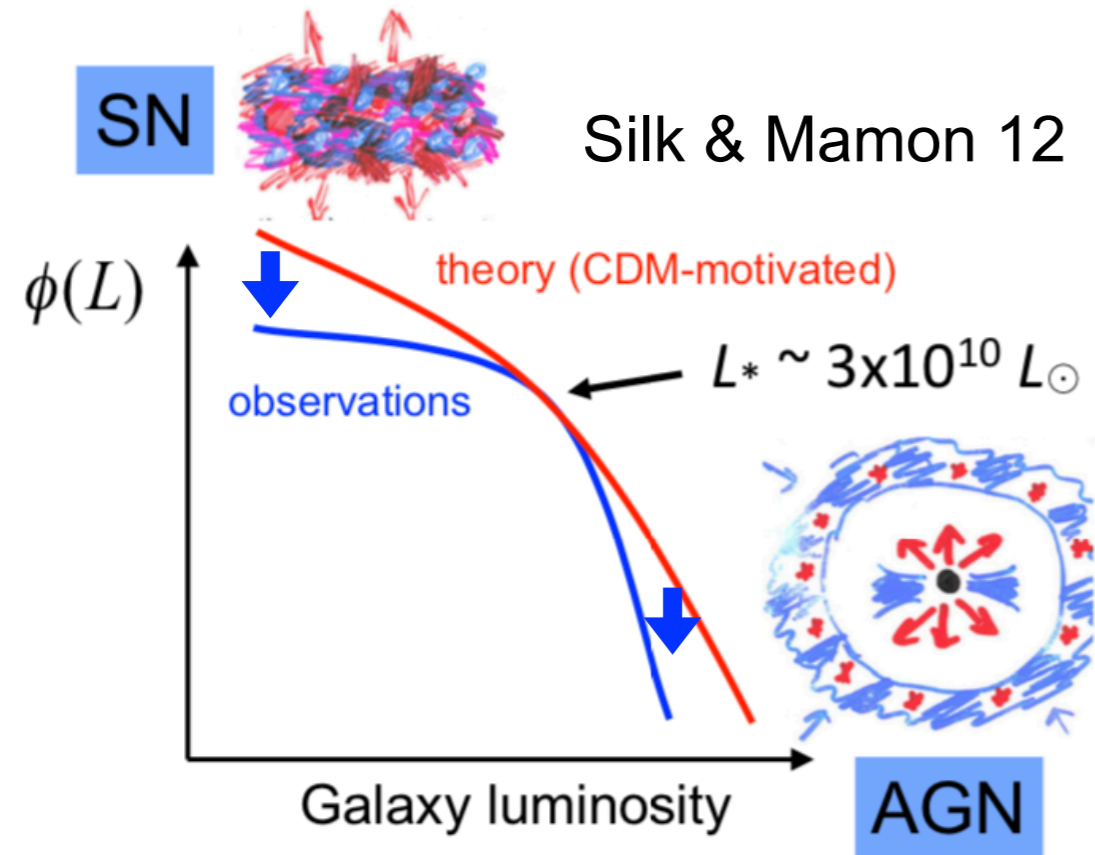
- **Baryonic feedback**
Stellar feedbacks such as SNe can transform central cusp into cored dark matter profiles.
- **Alternative dark matter models**
The other dark matter models motivated by particle physics (SIDM, SIMP, Axion..) can create a cored density profiles without relying on any baryon effects.

Caveats:

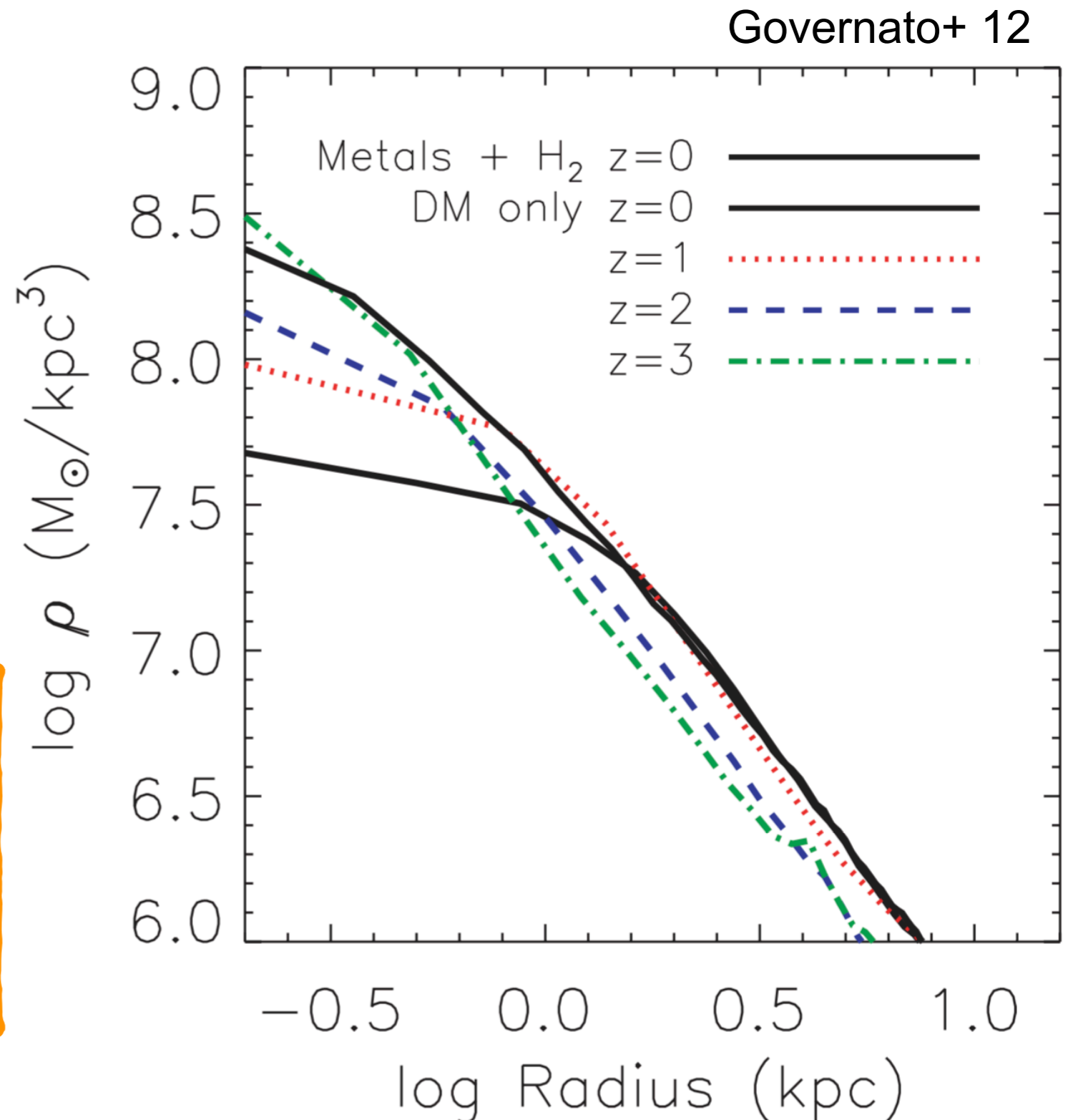
- **Incomplete observed data**
- **Uncertainties on dynamical analysis**

Solution to core-cusp problem

- Baryonic feedback



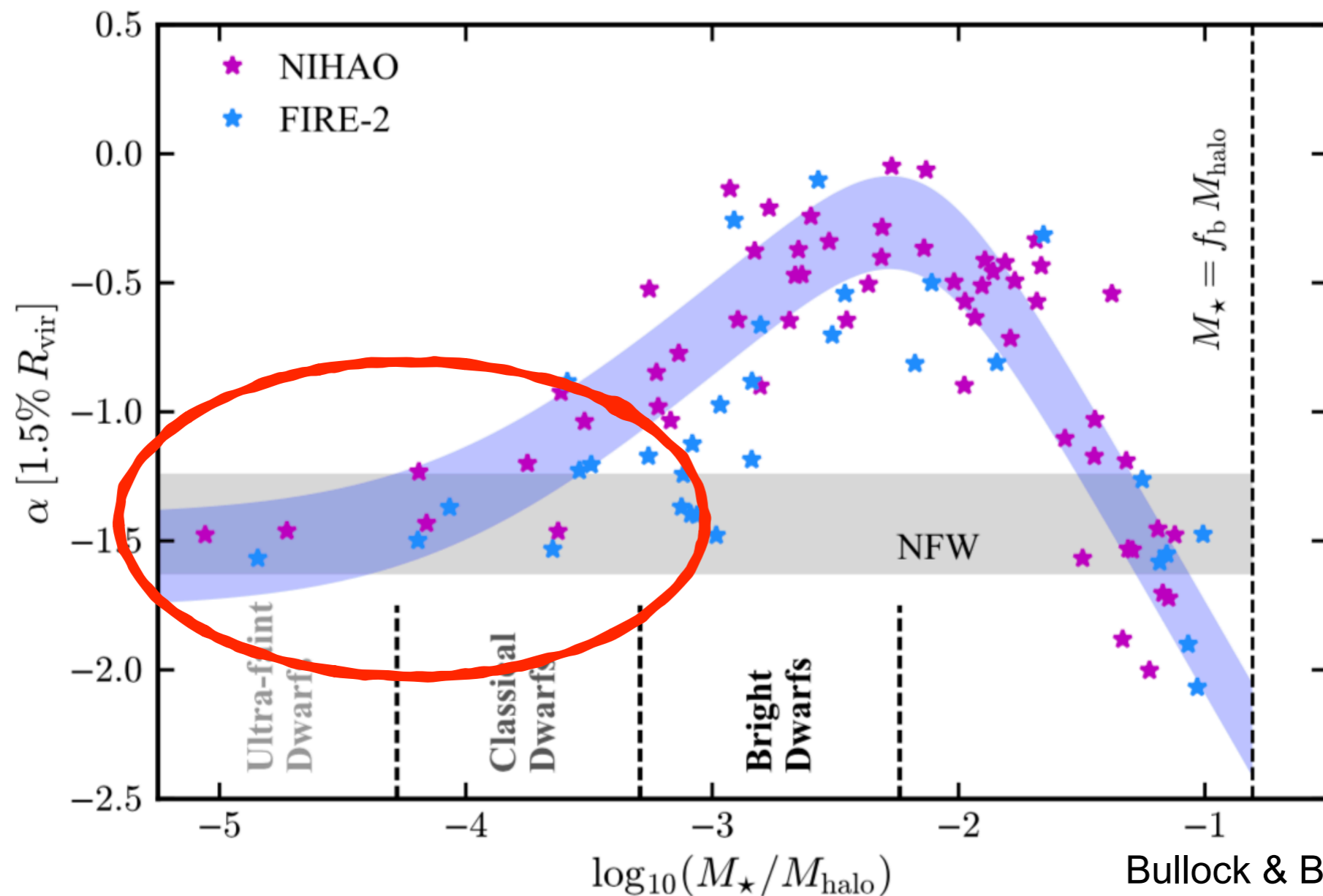
- Supernova feedback is dominated in less-massive galaxies.
- Numerical simulations have predicted that this feedback process can transform central cusp into cored dark matter profiles.



Solution to core-cusp problem

- Baryonic feedback

Recent high resolution simulations argue that **the impact of baryonic feedback is negligible on classical and ultra faint dwarf galaxy mass scales.**

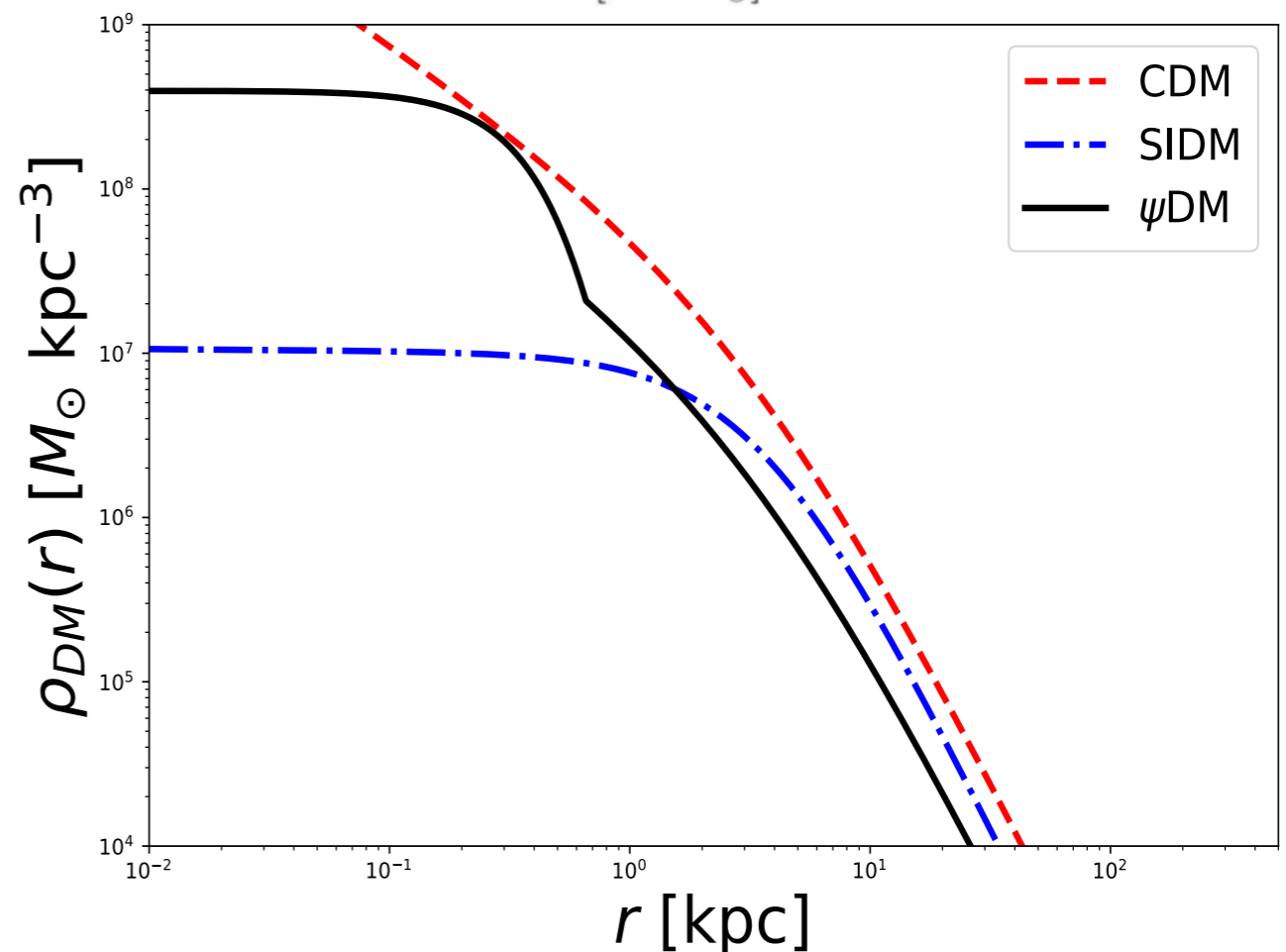
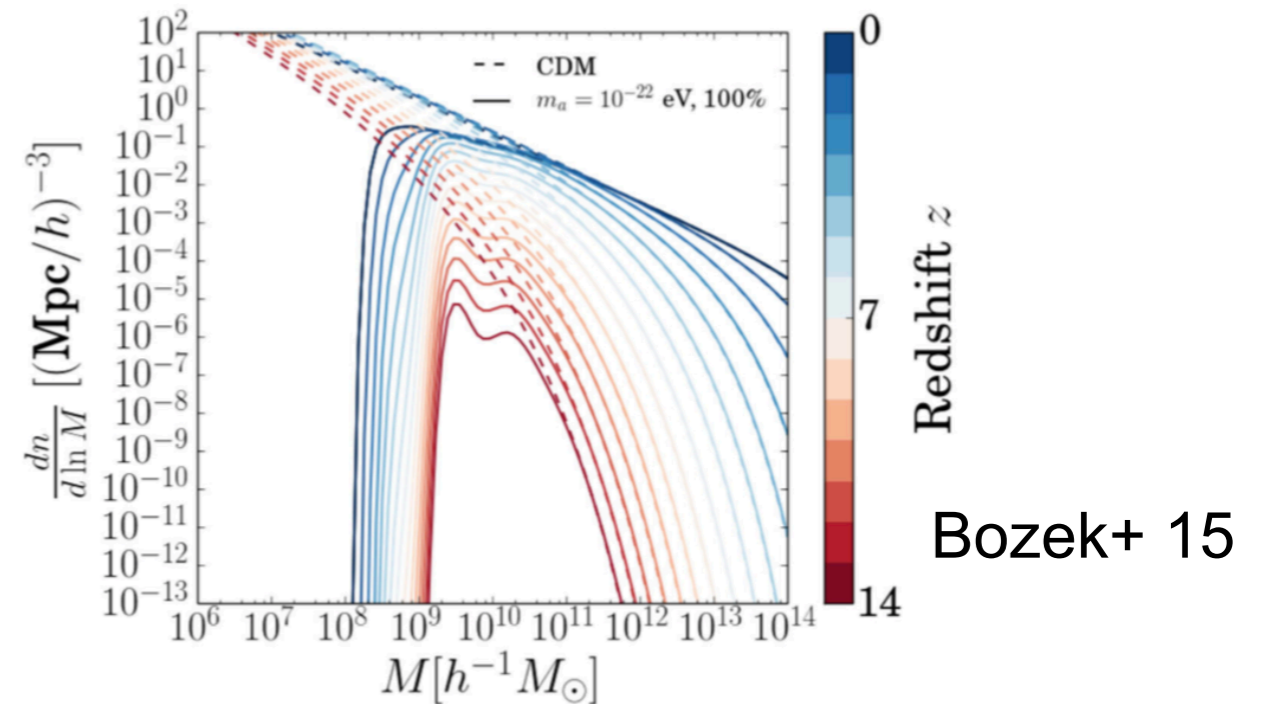


Bullock & Boylan-Kolchin (2017)

Solution to core-cusp problem

- Alternative DM models (WDM, SIDM, SIMP, Ultra light axion,...)

- Suppress the matter power spectrum on small scales.
- Dark matter cannot be concentrated on smaller spatial scales.
- **Create a cored dark matter density profiles** without relying on any baryon physics.



Ultra light Axion Dark Matter as a solution to small scale problems

Ultralight axion dark matter (ULADM)

- The lightest particle among dark matter candidates ($m_\psi \sim 10^{-22}$ eV)
- Create a core (\sim kpc) comes from quantum pressure

$$r_{\text{core}} \sim \lambda_{\text{dB}} \equiv \frac{h}{m_\psi v}$$

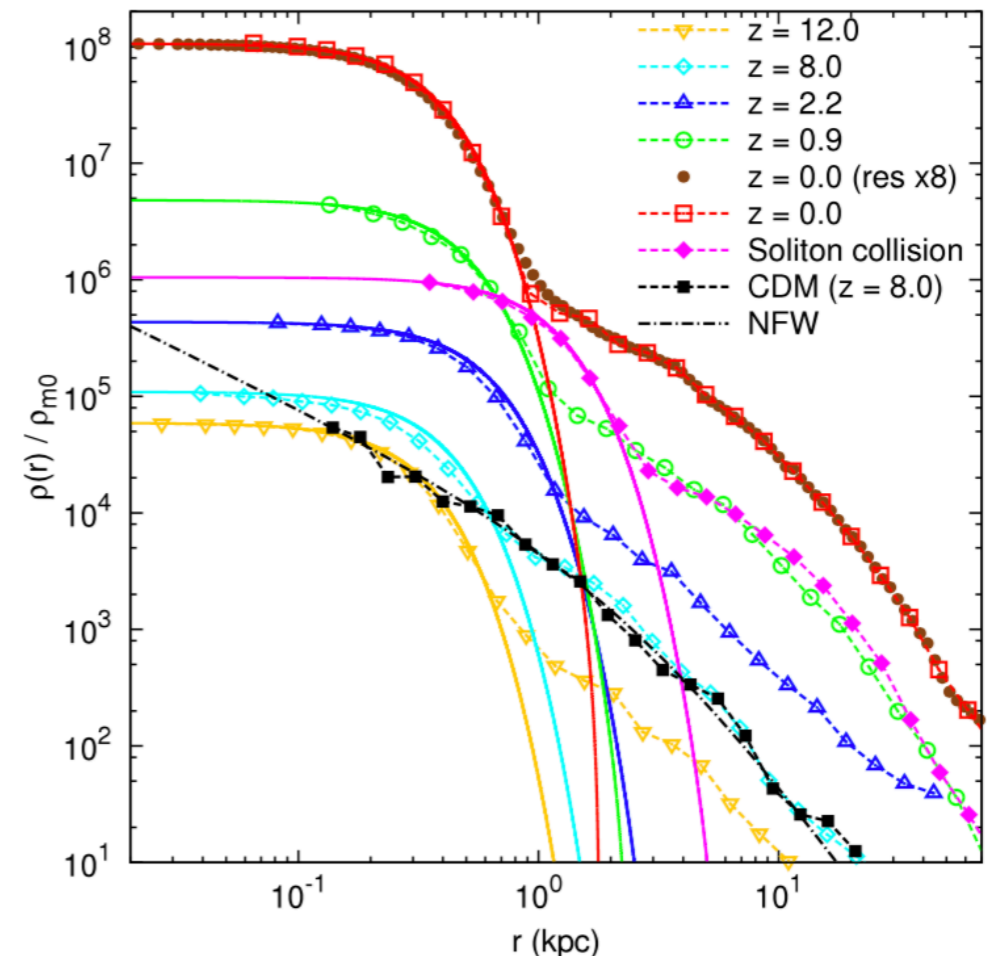
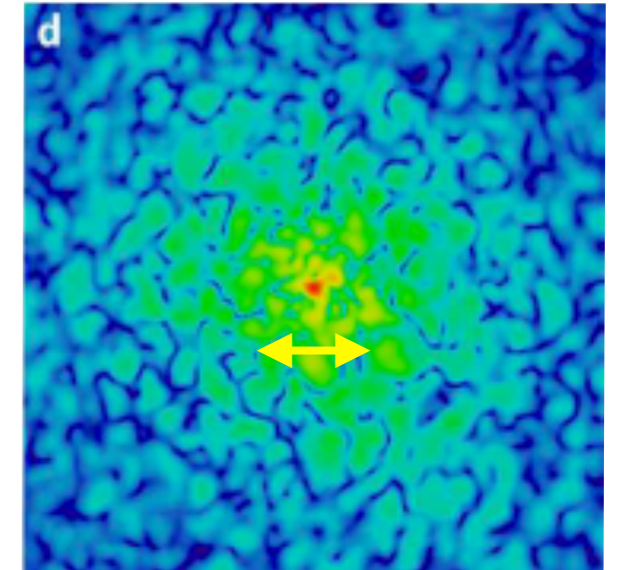
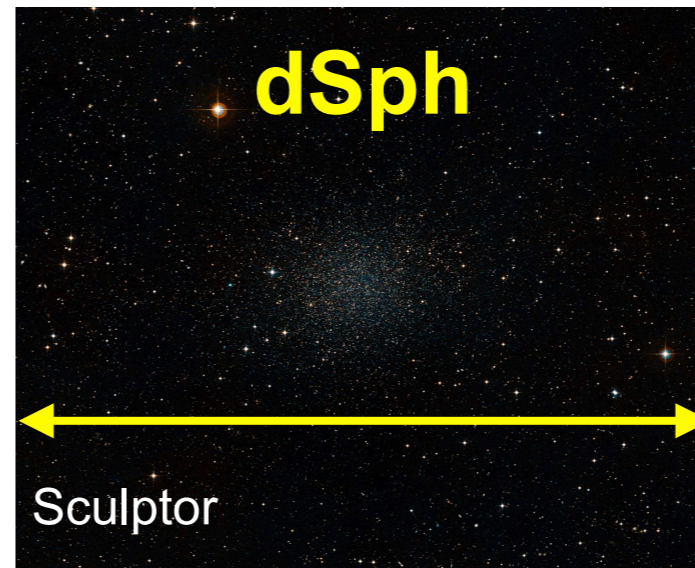
- central soliton core + outer NFW DM profile

Soliton-core dark matter density profile

$$\rho_{\text{soliton}}(r) = \frac{\rho_c}{[1 + 0.091(r/r_c)^2]^8}$$

$$\rho_c = 1.9 \times 10^{12} \left(\frac{m_\psi}{10^{-23} \text{ eV}} \right)^{-2} \left(\frac{r_c}{\text{pc}} \right)^{-4} [M_\odot \text{ pc}^{-3}]$$

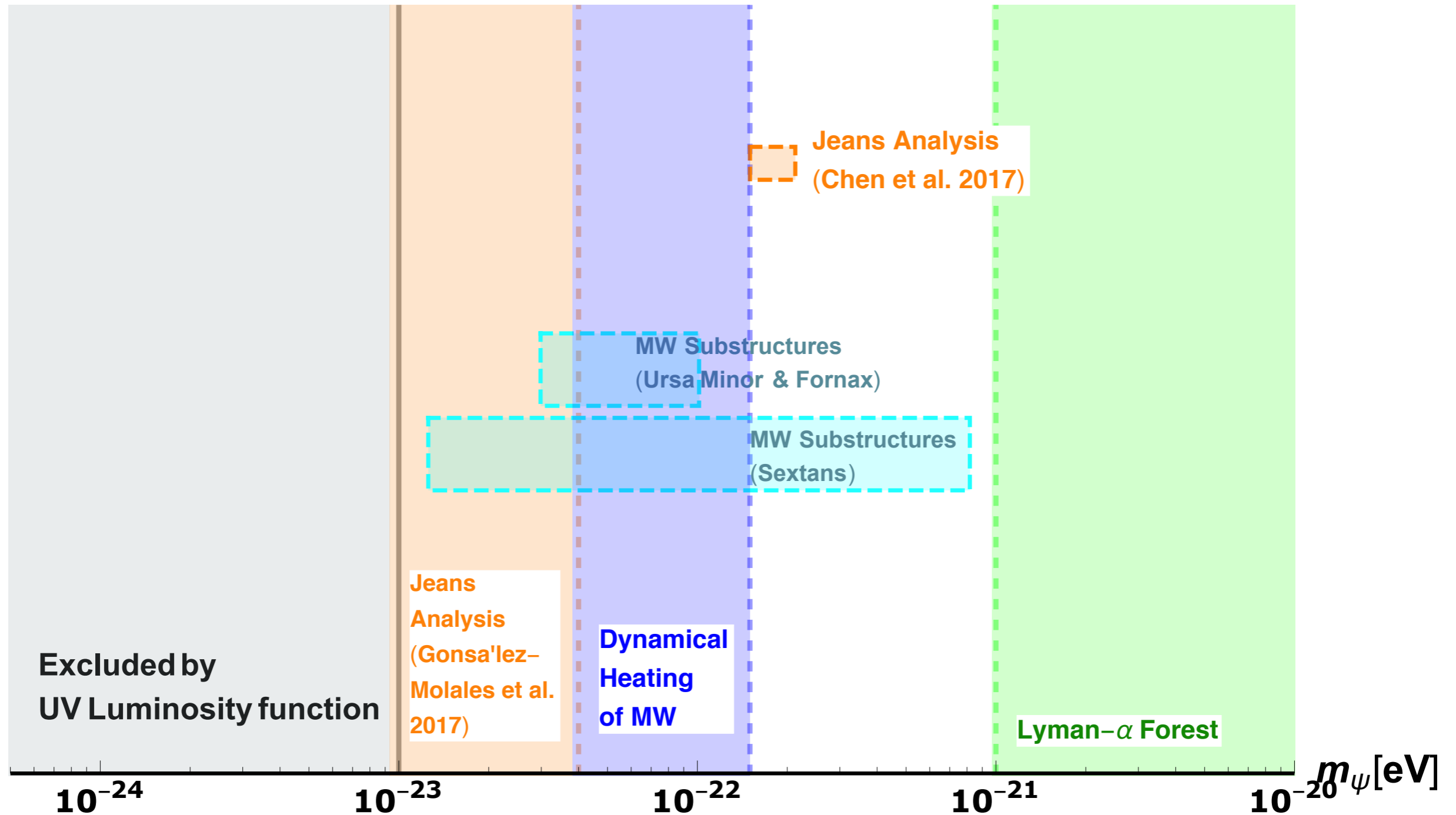
Schive et al. (2014)



Ultralight axion dark matter (ULADM)

- Current constraints on particle mass of ULADM

KH & Obata 19



Constraining particle mass of ULADM

Soliton-core dark matter density profile

$$\rho_{\text{soliton}}(r) = \frac{\rho_c}{[1 + 0.091(r/r_c)^2]^8}$$

$$\rho_c = 1.9 \times 10^{12} \left(\frac{m_\psi}{10^{-23} \text{ eV}} \right)^{-2} \left(\frac{r_c}{\text{pc}} \right)^{-4} [M_\odot \text{ pc}^{-3}]$$

Dynamical models for stars
ex) Spherical Jeans equation

$\sigma_{\text{l.o.s}}$ (Theory)

FIT

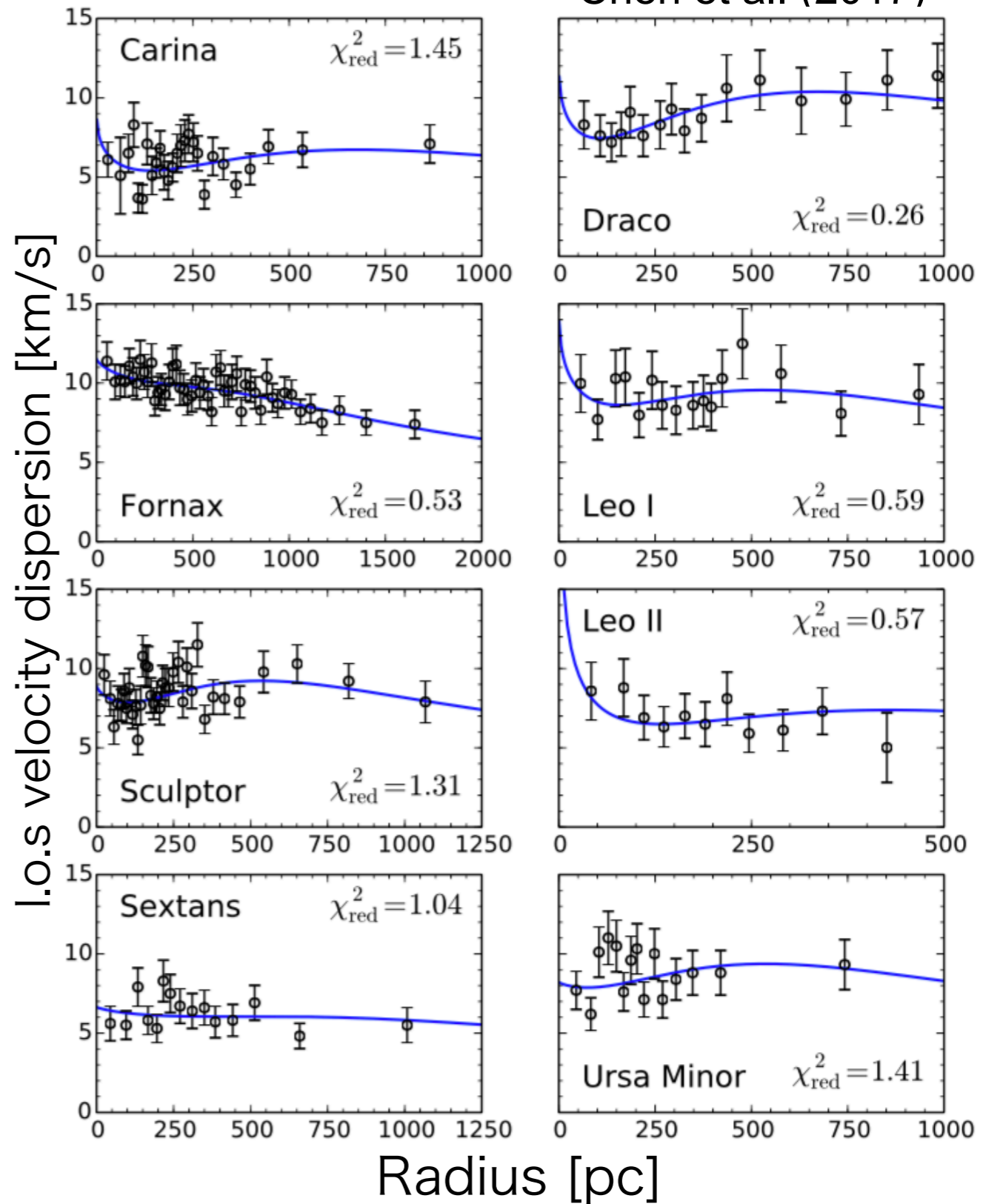
$\sigma_{\text{l.o.s}}$ (observed)

$$\sigma_{\text{l.o.s}}^2(R) = \frac{2}{I(R)} \int_R^\infty \left(1 - \beta_a \frac{R^2}{r^2} \right) \frac{\nu(r) v_r^2}{\sqrt{r^2 - R^2}} dr$$

- Current constraint from Spherical mass models

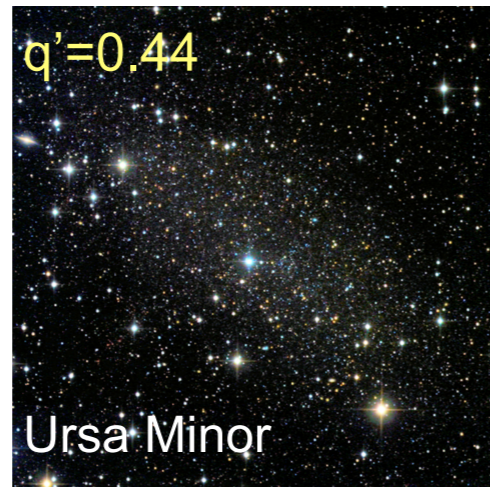
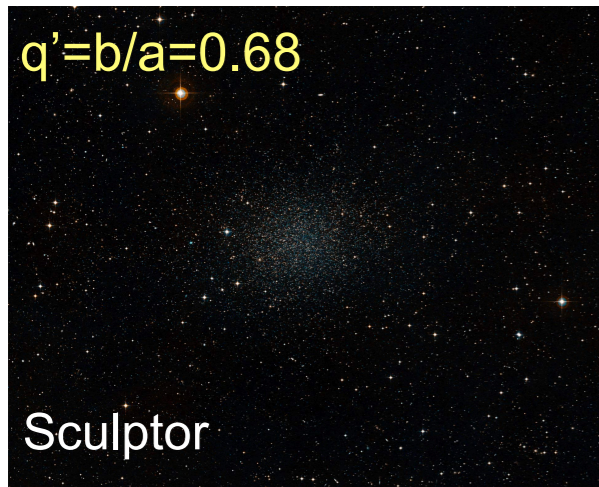
$$m_\psi = 1.79^{+0.35}_{-0.33} \times 10^{-22} \text{ eV } (2\sigma)$$

Chen et al. (2017)

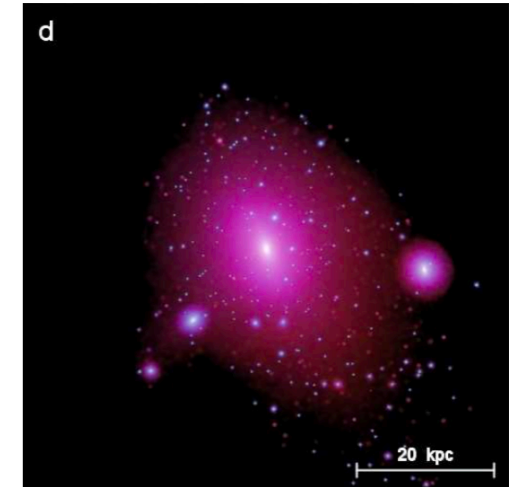
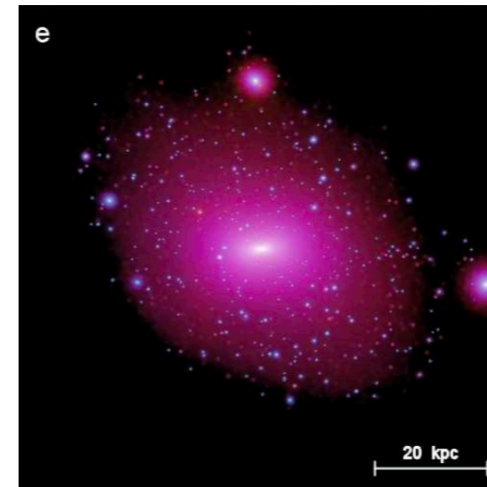


Major systematic uncertainty: Spherical Symmetry

1. Observed dSphs are **NOT** spherical shape

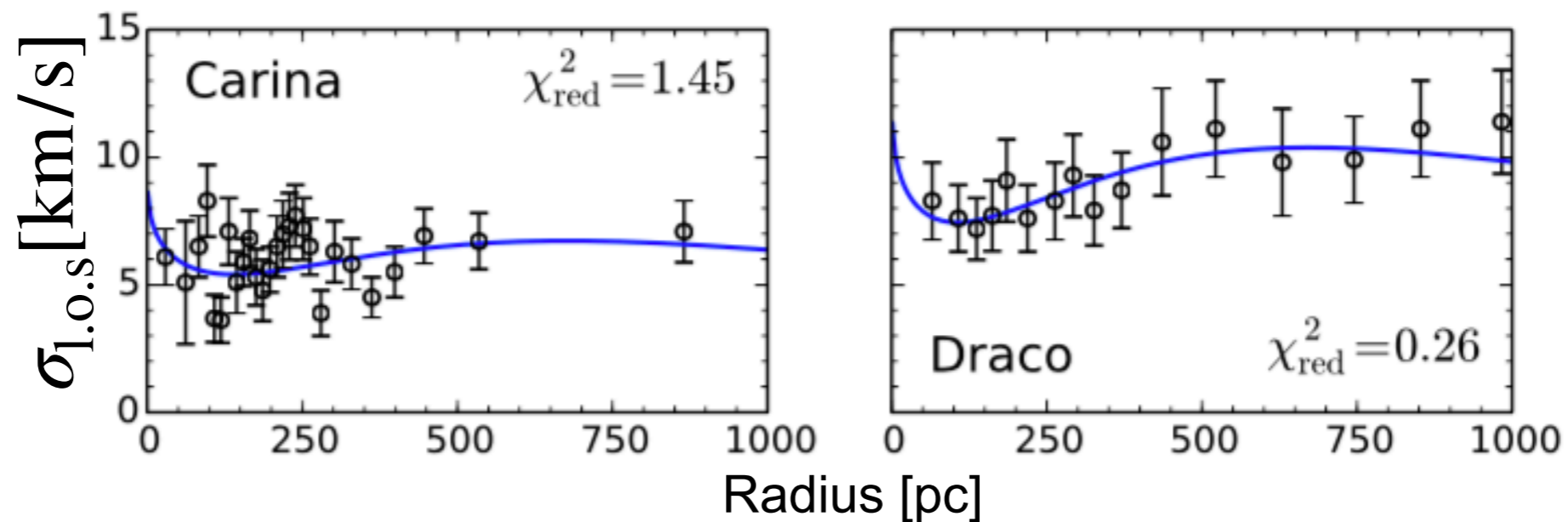


2. DM models predict **NON-spherical** DM halo



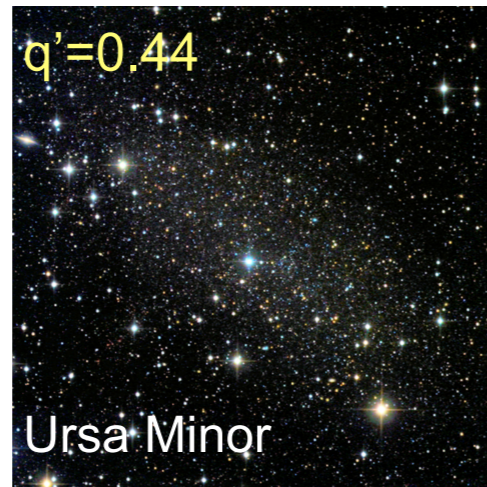
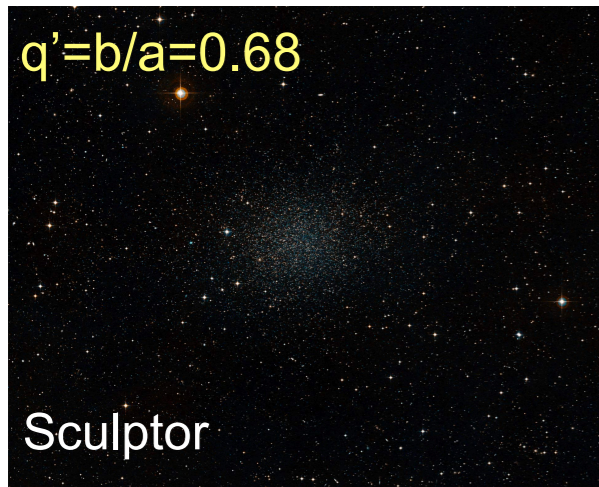
credit: Aquarius project

3. 1D spatial information

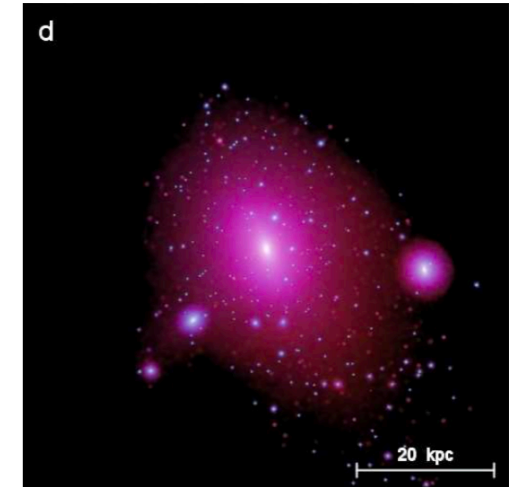
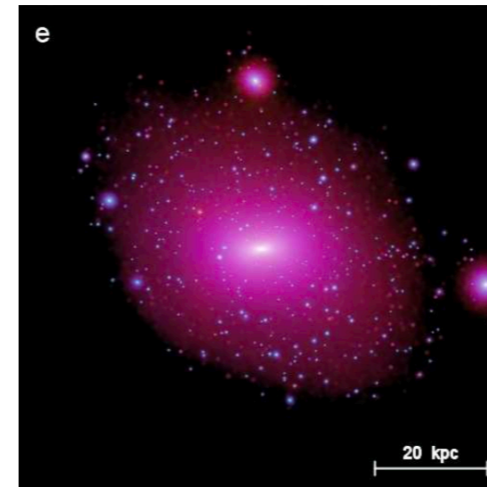


Major systematic uncertainty: Spherical Symmetry

1. Observed dSphs are **NOT** spherical shape

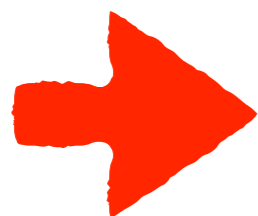
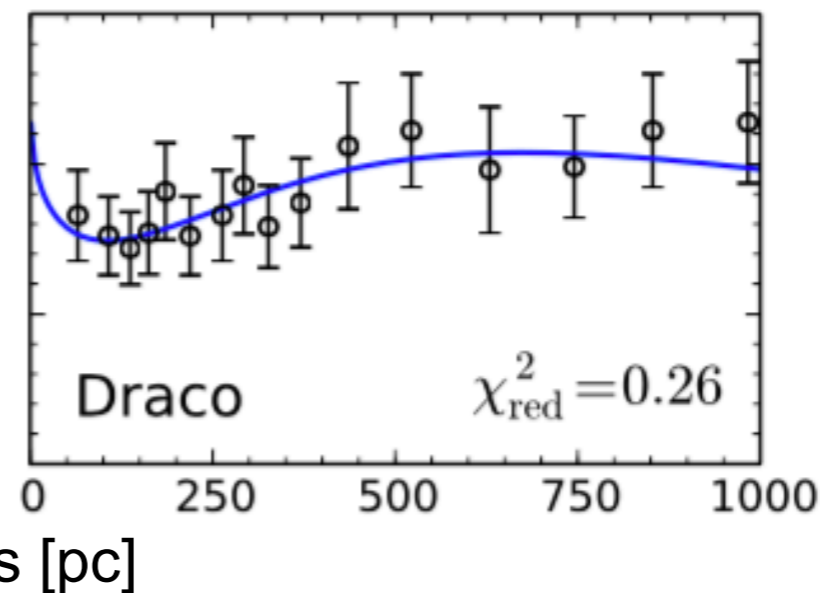
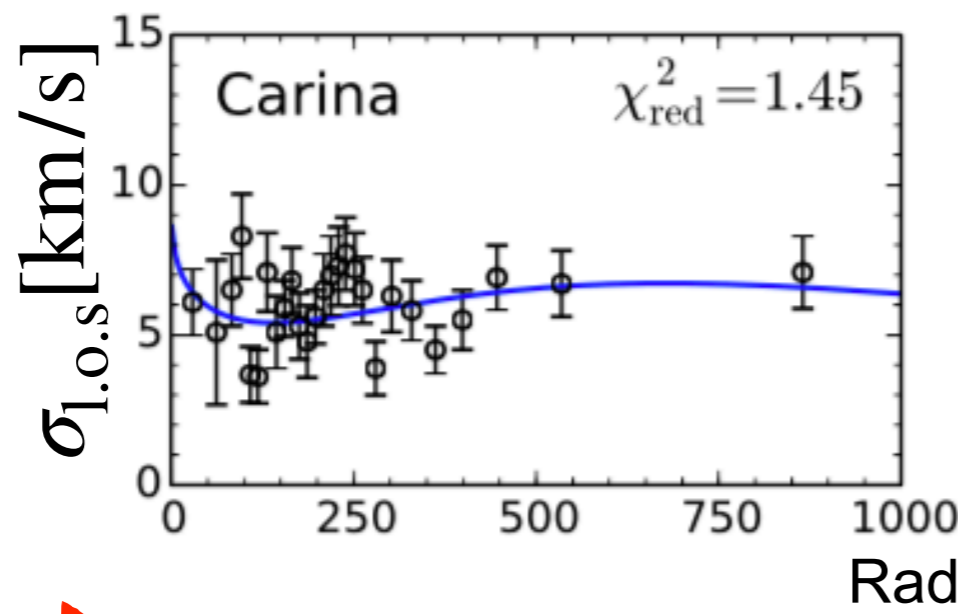


2. DM models predict **NON-spherical** DM halo



credit: Aquarius project

3. 1D spatial information



Non-spherical mass model

Non-sphericity of ultralight axion dark matter halos in the Galactic dwarf satellites

Hayashi and Obata (2019), arXiv: 190203054

Non-spherical dynamical mass models

Unobservable

Non-spherical dark matter density profile

$$\rho_{\text{soliton}}(r) = \frac{\rho_c}{[1 + 0.091(r/r_c)^2]^8}$$

$$\rho_c = 1.9 \times 10^{12} \left(\frac{m_\psi}{10^{-23} \text{ eV}} \right)^{-2} \left(\frac{r_c}{\text{pc}} \right)^{-4} [M_\odot \text{ pc}^{-3}]$$

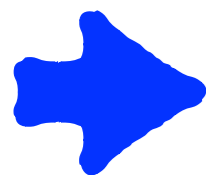
$$r^2 = R^2 + \frac{z^2}{Q^2} \quad \leftarrow \text{DM halo axial ratio}$$

Non-spherical stellar profile

$$\rho_*(r_*) = \frac{3L}{4\pi r_p^3} \left[1 + \frac{r_*^2}{r_p^2} \right]^{-5/2}$$

$$r_*^2 = R^2 + \frac{z^2}{q^2} \quad \leftarrow \text{stellar axial ratio}$$

Axisymmetric Jeans equations

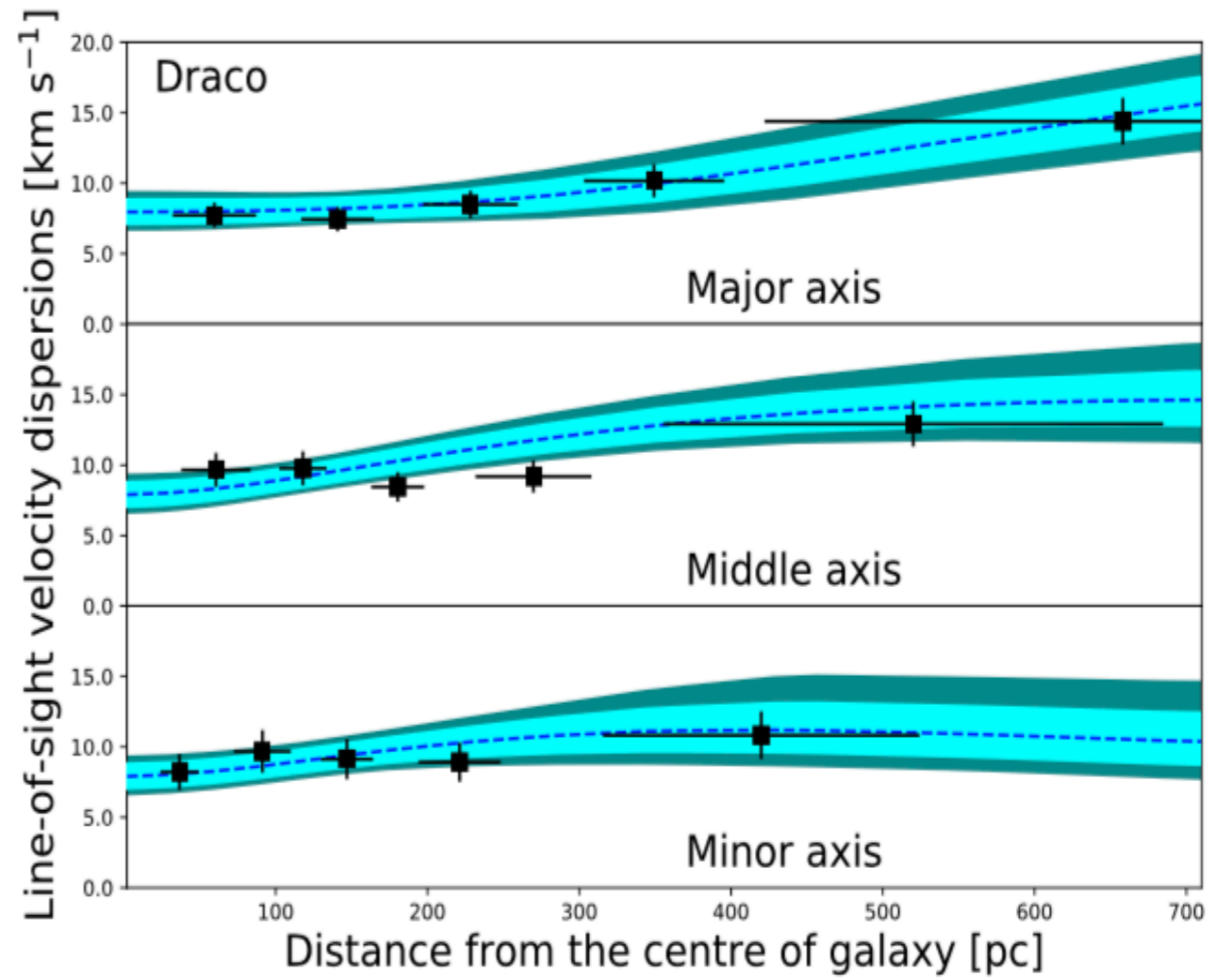

 $(\overline{v_z^2}, \overline{v_\phi^2})$

$\overline{v_R^2}$ is unknown parameter as $\beta_z = 1 - \overline{v_z^2}/\overline{v_R^2}$

$\sigma_{\text{l.o.s}}$ (theory)

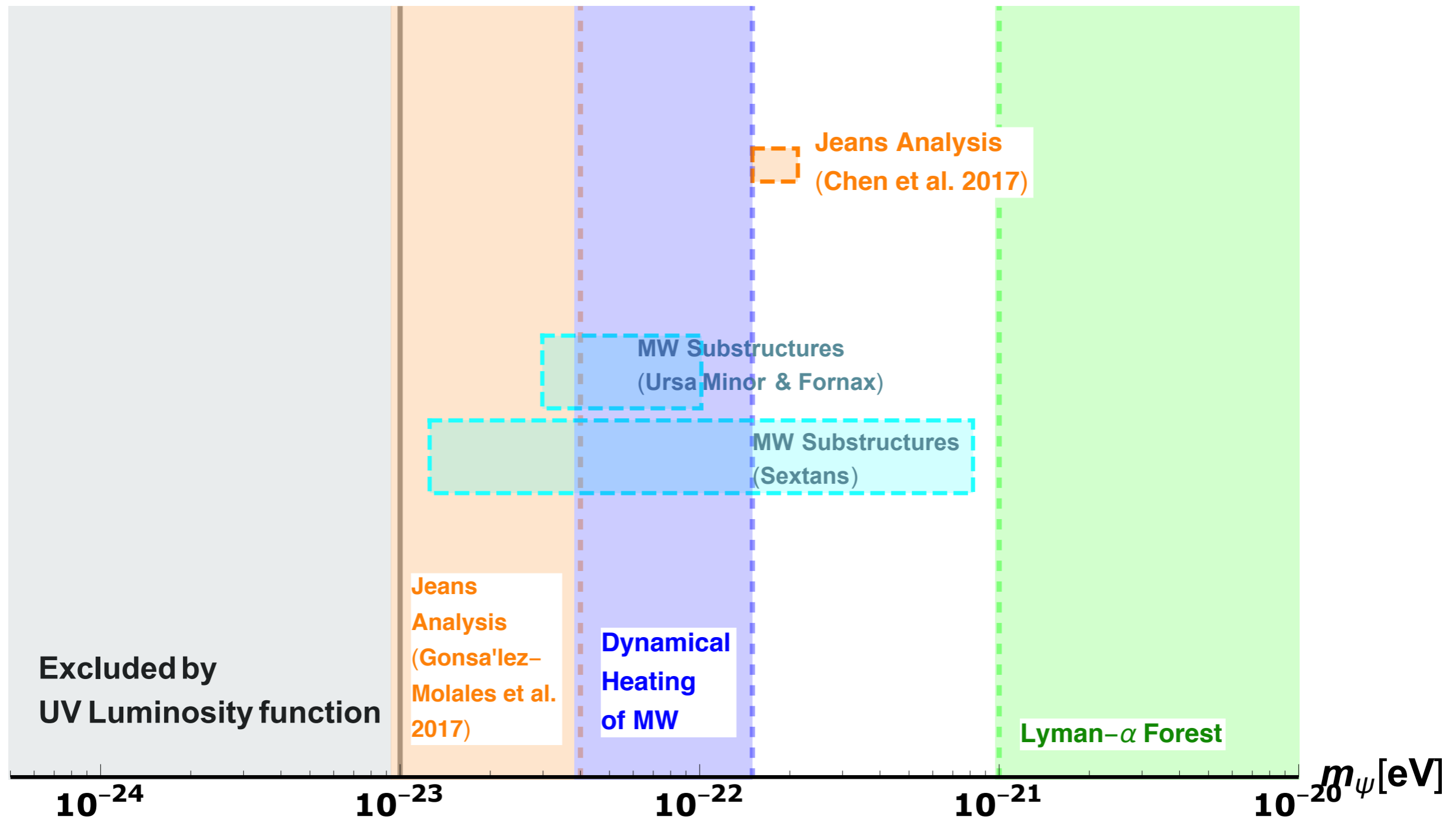
FIT

$\sigma_{\text{l.o.s}}$ (observed)



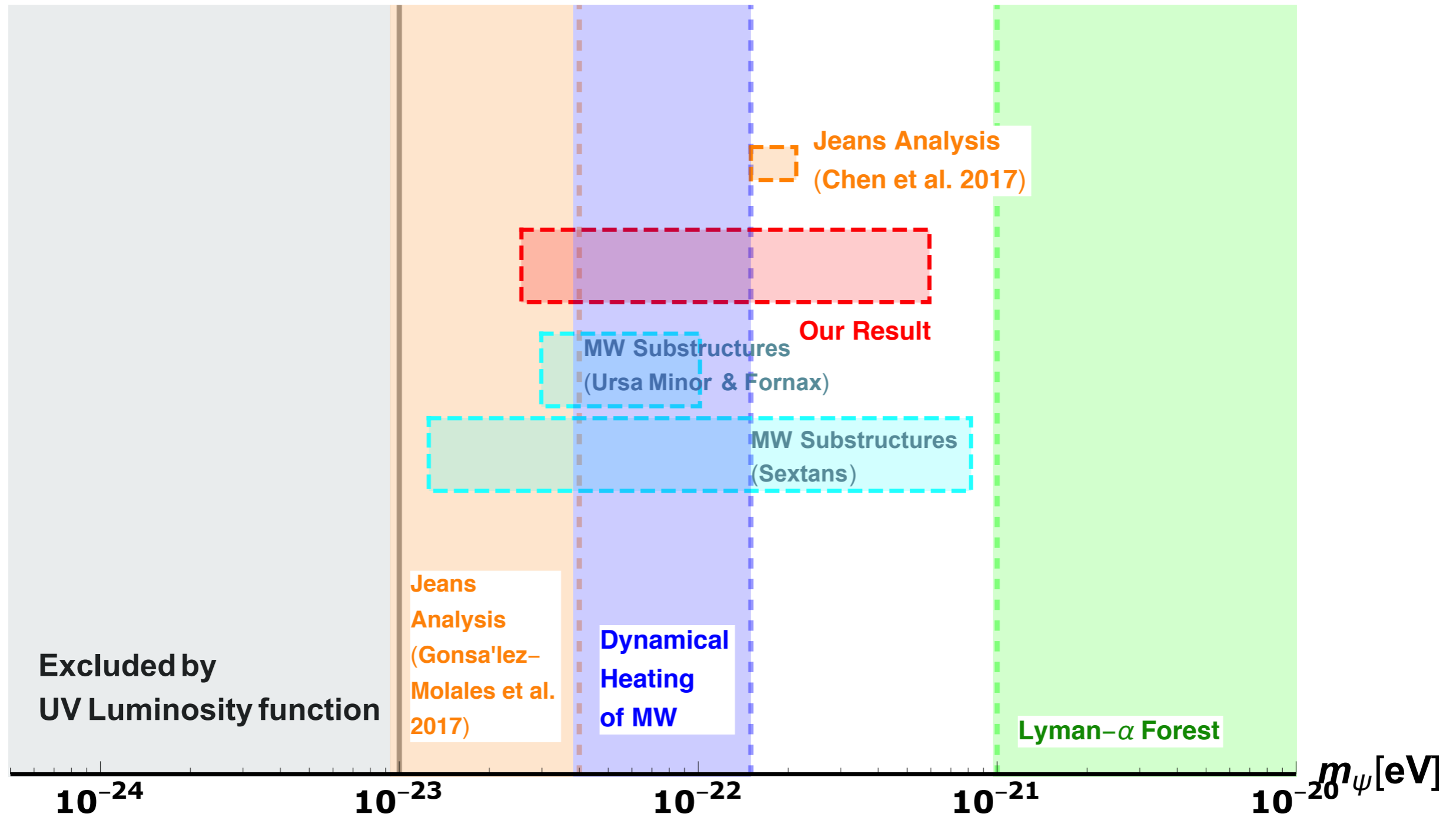
Constraints on ULADM via non-spherical analysis

Hayashi & Obata (2019), 1902.03054



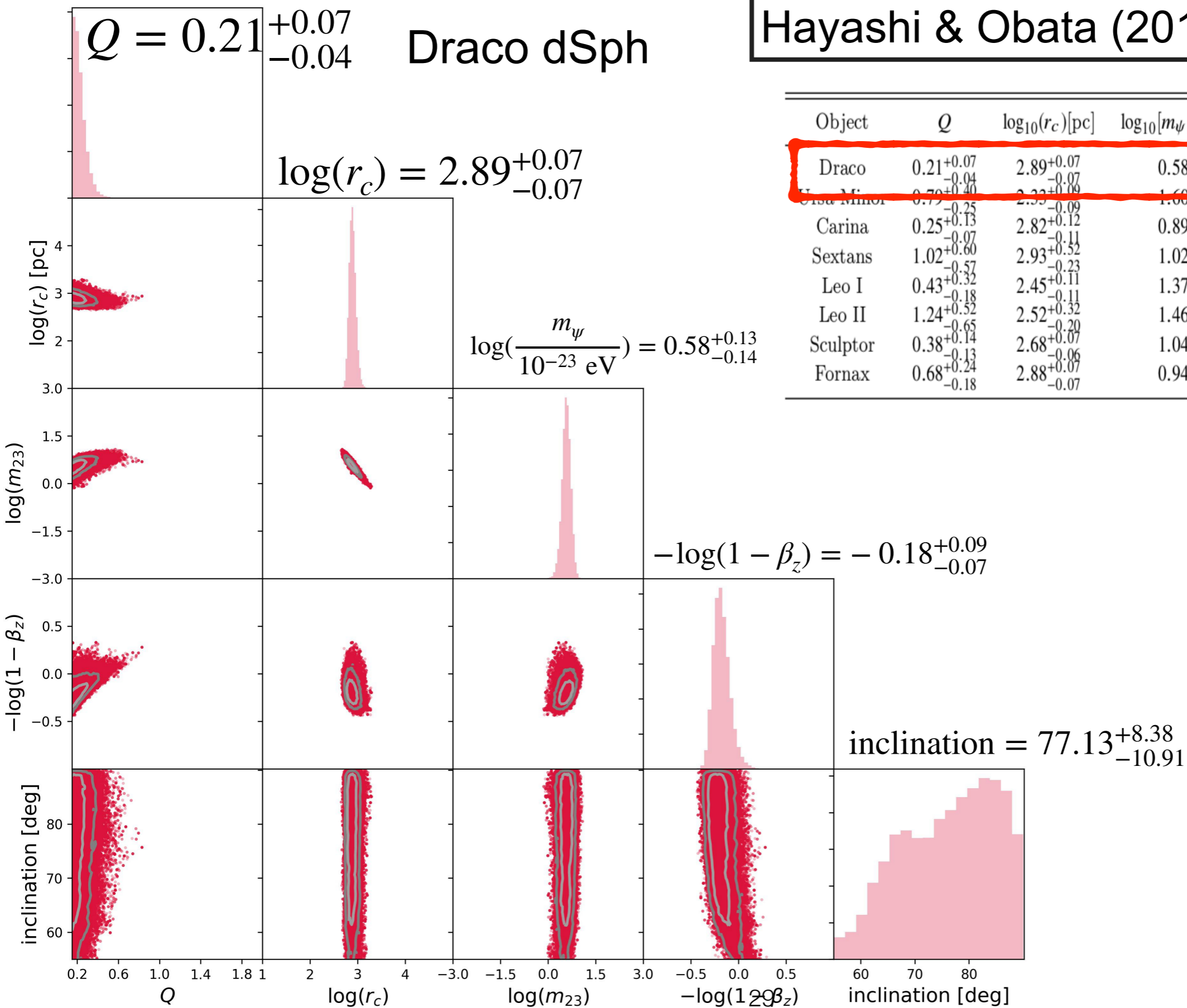
Constraints on ULADM via non-spherical analysis

Hayashi & Obata (2019), 1902.03054



Constraints on ULADM via non-spherical analysis

Hayashi & Obata (2019), 1902.03054



Object	Q	$\log_{10}(r_c)$ [pc]	$\log_{10}[m_\psi/10^{-23}\text{eV}]$	$-\log_{10}(1-\beta_z)$	i [deg]
Draco	$0.21^{+0.07}_{-0.04}$	$2.89^{+0.07}_{-0.07}$	$0.58^{+0.13}_{-0.14}$	$-0.18^{+0.09}_{-0.07}$	$77.13^{+8.38}_{-10.91}$
Ursa Minor	$0.79^{+0.40}_{-0.25}$	$2.33^{+0.09}_{-0.09}$	$1.60^{+0.15}_{-0.16}$	$0.47^{+0.13}_{-0.08}$	$81.15^{+5.64}_{-7.41}$
Carina	$0.25^{+0.13}_{-0.07}$	$2.82^{+0.12}_{-0.11}$	$0.89^{+0.21}_{-0.22}$	$-0.08^{+0.14}_{-0.11}$	$73.73^{+10.77}_{-11.33}$
Sextans	$1.02^{+0.60}_{-0.57}$	$2.93^{+0.52}_{-0.23}$	$1.02^{+0.39}_{-0.89}$	$0.31^{+0.33}_{-0.25}$	$70.34^{+12.95}_{-10.21}$
Leo I	$0.43^{+0.32}_{-0.18}$	$2.45^{+0.11}_{-0.11}$	$1.37^{+0.20}_{-0.20}$	$-0.12^{+0.13}_{-0.09}$	$65.57^{+15.52}_{-14.38}$
Leo II	$1.24^{+0.52}_{-0.65}$	$2.52^{+0.32}_{-0.20}$	$1.46^{+0.33}_{-0.58}$	$0.27^{+0.18}_{-0.22}$	$63.47^{+17.19}_{-14.42}$
Sculptor	$0.38^{+0.14}_{-0.13}$	$2.68^{+0.07}_{-0.06}$	$1.04^{+0.15}_{-0.17}$	$0.06^{+0.19}_{-0.09}$	$64.37^{+14.61}_{-11.55}$
Fornax	$0.68^{+0.24}_{-0.18}$	$2.88^{+0.07}_{-0.07}$	$0.94^{+0.14}_{-0.15}$	$0.17^{+0.11}_{-0.07}$	$69.79^{+11.02}_{-10.74}$

Constraints on ULADM via non-spherical analysis

Hayashi & Obata (2019), 1902.03054

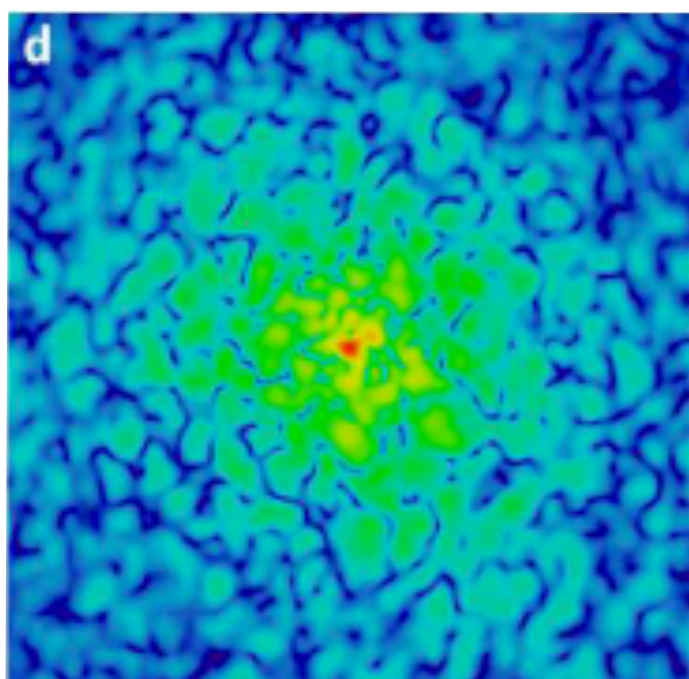
Stellar & DM halo axial ratio of Draco

$$q = (b/a)_{\text{star}} = 0.69$$

$$Q = (b/a)_{\text{DM}} = 0.21$$



- Draco has strongly elongated dark halo, **$Q \sim 0.2$** .
- Draco's ULADM halo is **much more flattened** than N-body predictions and stellar distributions.
- Further understanding of baryonic and DM physics should be needed.



$Q \sim 1.0$

Schive et al. (2014)

Summary

- Λ CDM theory faces the serious challenges on dwarf galaxy scales.
- Ultralight axion dark matter is one of the dark matter candidates, because it can resolve small scale problems.
- The MW dSphs are ideal sites for studying the nature of dark matter because these are DM-dominated systems.
- To obtain realistic limits on DM models, we construct new dynamical modeling with taking into account non-sphericity.
- Our mass models place less stringent constraints on ULADM mass but require unphysically elongated ULADM halos.
- We revisit core-cusp problem and find that the diversity of inner slopes of DM profiles in the classical dwarfs.

