# Resonant Magnetic Field Generation from axions

with

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### **Cosmological Magnetic Fields**



### Various Magnetogenesis Models

- Primordial Magnetic Fields
  - Inflationary
  - Phase transition

B. Ratra 1992, Martin & Yokoyama 2008, Durrer+ 2011, ...

Vachaspati 1992, Durrer and Nerenov+2013, ..

- Post Inflationary
  - Onset of LSS formation
  - Biermann batteries

Galactic winds : Bertone et al 2006, .. Relativistic outflows from AGNs : Rees 1987, ..

Subramaniam 2018

#### Setups

$$\mathcal{L} = -\frac{1}{2}\partial_{\mu}\phi\partial^{\mu}\phi - V(\phi) - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{\alpha}{4}\frac{\phi}{f}F_{\mu\nu}\tilde{F}^{\mu\nu}$$

Chern-Simons Coupling

Dominant

- Inflationary (Monodromy)
- Reheating

# Subdominant

- Spectator
- Post Inflationary

### **General equations**

#### Klein-Gordon EOM

$$\phi'' + 2\mathcal{H}\phi' + a^2 V_{,\phi} = a^2 \frac{\alpha}{f} \langle \mathbf{E} \cdot \mathbf{B} \rangle$$

#### Gauge field EOM

$$\mathcal{A}_{\pm}^{\prime\prime}(\eta,\,k) + \left(k^2 \mp \frac{\alpha}{f} \frac{\phi^{\prime}}{\mathcal{H}} k \mathcal{H}\right) \mathcal{A}_{\pm}(\eta,\,k) = 0$$

Backreaction

$$\langle \mathbf{E} \cdot \mathbf{B} \rangle = -\frac{1}{8\pi^2} \int \mathrm{d}\ln k \left(\frac{k}{a}\right)^4 \left[\frac{1}{k} \frac{\mathrm{d}}{\mathrm{d}\eta} \left(|\sqrt{2k}\mathcal{A}_+|^2\right) - \frac{1}{k} \frac{\mathrm{d}}{\mathrm{d}\eta} \left(|\sqrt{2k}\mathcal{A}_+|^2\right)\right]$$

### Monodromy

$$V(\phi) = \mu^2 \phi^2 + \Lambda^4 \cos \frac{\phi}{f}$$

$$\varepsilon \equiv \frac{1}{2 \mathrm{Mpl}^2} \left(\frac{\phi'}{\mathcal{H}}\right)^2 \qquad |\xi| = \sqrt{2\varepsilon} \alpha \frac{\mathrm{Mpl}}{f}$$

 $\overline{}$ 



 $\alpha \left( f/\mathrm{Mpl} \right) : 10$ 

### Monodromy





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### Reheating

$$V(\phi) = \frac{(mf)^2}{2} \left[ 1 - \frac{1}{(1 + \tilde{\phi}^2/p)^p} \right] \equiv (mf)^2 \tilde{V}(\tilde{\phi})$$

### **Pure Natural Potential**



### Reheating

Background

 $V(\phi) = \frac{(mf)^2}{2} \left[ 1 - \frac{1}{(1 + \tilde{\phi}^2/p)^p} \right] \equiv (mf)^2 \tilde{V}(\tilde{\phi})$  $\left(\frac{H}{m}\right)^2 = \frac{1}{6} \left(\frac{f}{\mathrm{Mpl}}\right)^2 \left| \left(\frac{\mathrm{d}\tilde{\phi}}{\mathrm{d}\tilde{t}}\right)^2 + 2\tilde{V} \right|$  $\frac{\mathrm{d}^2 \tilde{\phi}}{\mathrm{d}\tilde{t}^2} + 3\frac{H}{m}\frac{\mathrm{d}\tilde{\phi}}{\mathrm{d}\tilde{t}} + \tilde{V}_{,\tilde{\phi}} = 0$ 



$$\frac{\mathrm{d}\mathcal{A}_h}{\mathrm{d}\tilde{t}^2} + \frac{H}{m}\frac{\mathrm{d}\mathcal{A}_h}{\mathrm{d}\tilde{t}} + \omega_h^2\mathcal{A}_h = 0 \qquad \omega_h^2 \equiv \frac{k^2}{a^2m^2} - h\alpha\frac{k}{ma}\left(\frac{\mathrm{d}\tilde{\phi}}{\mathrm{d}\tilde{t}}\right)$$



### Reheating

#### **Generated Spectra**



## Tachyonic Resonance

Evolution of spectrum generated by a spectator (Next section)



### **Growth Rate**

 $y(z) \propto e^{\mu z} \Phi(z) \qquad \Phi(z+T)$ 



### **Growth rate**





## **Energy Density**

 $f/{\rm Mpl}: 10^{-2}$   $f/{\rm Mpl}: 5 \times 10^{-3}$ 



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### Helicity



#### Spectator

Kitajima, Soda and Yuko, 2018 GW emission in R.D Universe

Setup  

$$H/m = (a/a_{H=m})^{-\epsilon} \quad V(\phi) = \frac{(mf)^2}{2} \left[ 1 - \frac{1}{(1 + \tilde{\phi}^2/p)^p} \right] \equiv (mf)^2 \tilde{V}(\tilde{\phi})$$



### Spectator

### Helicity



### Spectator



### **Post Reheating generation**

#### Subdominant axion field oscillating during Radiation Dominated

$$\frac{\mathrm{d}^2 \mathcal{A}_h}{\mathrm{d}\tilde{t}^2} + \left(\frac{H}{m} + \frac{4\pi\sigma}{m}\right)\frac{\mathrm{d}\mathcal{A}_h}{\mathrm{d}\tilde{t}} + \omega_h^2 \mathcal{A}_h = 0 \qquad \qquad \frac{H}{m} = \frac{1}{2mt}$$

$$\sigma \begin{cases} 10 \left(\frac{T}{1 \text{GeV}}\right) \text{GeV } T \gg m_e = 0.511 \text{MeV} & \text{Baym and} \\ \text{Heiselberg, 1997} \end{cases}$$
$$3.2 \times 10^4 \text{GeV} \left(\frac{T}{\text{GeV}}\right)^{\frac{3}{2}} T \ll m_e & \text{Spitzer, 1953} \end{cases}$$

#### Criteria



### **Post Inflationary Evolution**

Non-Helical

$$B \propto rac{1}{a^2}$$
  $\lambda \propto a$ 

• Helical

Inverse cascade  

$$\begin{array}{l}
 Helicity \\
 conservation
\end{array} \quad \begin{array}{l}
 Helicity \\
 Conservation
\end{array} \quad \begin{array}{l}
 Helicity \\
 Evolution
\end{array}$$

$$\begin{array}{l}
 (\mathfrak{H}) = \sum_{k=1}^{3} \frac{a_{gen}}{k_{gen}} B_{k_{m,gen}}^{2} \quad \frac{a_{rec}}{k_{m,rec}} \approx \frac{B_{k_{m,rec}}}{\rho_{rec}/M_{pl}}$$

$$B_{km,0} \simeq 10^{-8} \mathrm{G} \times \left(\frac{\lambda_{m,0}}{1 \mathrm{Mpc}}\right)$$

A diabatia

### **Order Estimations**

• Axion Monodromy

$$\lambda_{m,0} \ll 10^2 \mathrm{Mpc}\left(\frac{a_{\star}}{a_{\mathrm{R}}}\right) \left(\frac{10^5 \mathrm{GeV}}{T_{\gamma,\mathrm{R}}}\right) \left(\frac{f}{10^{-3} \mathrm{Mpl}}\right)^{\frac{2}{3}} \left(\frac{1}{\Delta \ln k_{\mathrm{rec}}}\right)^{\frac{1}{3}} \dots$$

• Axion Spectator

$$\lambda_{m,0} \ll 4 \mathrm{Mpc}\left(\frac{a_{\star}}{a_R}\right) \left(\frac{10^6 \mathrm{GeV}}{T_{\gamma,\mathrm{R}}}\right) \left(\frac{1}{\alpha}\right)^{\frac{1}{3}} \left(\frac{\rho_{\phi}}{\rho_{\mathrm{inf}}}\right)^{\frac{1}{3}} \dots$$

Reheating

$$\lambda_{m,0} \ll 10^3 \mathrm{Mpc} \left(\frac{10^5 \mathrm{GeV}}{T_{\gamma,\mathrm{R}}}\right) \left(\frac{f}{10^{-3} \mathrm{Mpl}}\right)^{1/3} \left(\frac{1}{\Delta(\ln k)_{\mathrm{rec}}}\right)^{\frac{1}{3}} \dots$$



- Resonant production of 'helical' magnetic field from axions is efficient
- Generation in the plasma-filled Universe is limited
- These processes can saturate the energy density(without taking non-linear processes into account), and can be promising candidates for origin of cosmic magnetic fields