

# Volume of Moduli Space of Vortices and Localization Formula

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# Introduction

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Topic:

- ❖ Calculation of the volume of the moduli space of BPS solitons

Moduli space?

- ❖ Kähler or hyper-Kähler quotient space

{Solutions of F & D-term constraints} / {Gauge symmetry}

BPS soliton

- ❖ Abelian/non-Abelian vortex on a *compact* Riemann surface with genus  $h$  ( $\Sigma_h$ )

# Introduction

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BPS equations ( $G=U(N_c)$  and  $N_f$  flavors):

$$\mu_r \equiv F - \frac{g^2}{2}(c - HH^\dagger)\omega = 0$$

$$\mu_{\bar{z}} \equiv \mathcal{D}_{\bar{z}}H = 0$$

$$\mu_z \equiv \mathcal{D}_z H^\dagger = 0$$

where  $g$ : gauge coupling,  $c$ : FI parameter,  $\omega$ : Kähler 2-form on  $\Sigma_h$ ,  
 $H: N_c \times N_f$  matrix.

$$\mathcal{M}_k \equiv \frac{\{\text{Solutions of } \mu_r = \mu_z = \mu_{\bar{z}} = 0 \text{ with } \frac{1}{2\pi} \int F = k\}}{U(N_c)}$$

Volume of  $\mathcal{M}_k$ ?

# Introduction

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The volume of moduli space of the BPS solitons relates to

- ❖ Non-perturbative corrections in supersymmetric gauge theory (Nekrasov's formula)
- ❖ Thermodynamical partition function of the BPS solitons (Manton et al.)

$$Z = \frac{1}{\hbar^{2k}} \int d^k p d^k x e^{-\frac{1}{2T} g^{ij} p_i p_j} = \left( \frac{2\pi^2 T}{\hbar^2} \right)^k \text{Vol}(\mathcal{M}_k)$$

The volume of moduli space of the BPS solitons (supersymmetric systems) is a key to understand the non-perturbative dynamics and dualities in gauge/string theory

# Method

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Straightforwardly,

BPS eqs.  $\Rightarrow$  BPS solutions  $\Rightarrow$  effective action  $\Rightarrow$  metric  $\Rightarrow$  volume

Instead, we define the field theoretical partition function to obtain the volume of the moduli space [Moore-Nekrasov-Shatashvili (1997)]:

$$\mathcal{Z}_k^{N_c, N_f}(\Sigma_h) = \int \mathcal{D}\Phi \mathcal{D}^2 A \mathcal{D}^2 \lambda \mathcal{D}^2 H \mathcal{D}^2 \psi \mathcal{D}^2 Y \mathcal{D}^2 \chi e^{-S_0 - S_1}$$

where

$$S_0 = \int_{\Sigma_h} \text{Tr} \left[ i\Phi \left\{ F - \frac{g^2}{2} (c - HH^\dagger) \omega \right\} + \frac{1}{2} \lambda \wedge \lambda + \frac{g^2}{2} \psi \psi^\dagger \omega \right]$$

$$S_1 = \int_{\Sigma_h} d^2 z \text{Tr} [g^{z\bar{z}} (Y_z Y_{\bar{z}} + i\Phi \chi_z \chi_{\bar{z}})] + \dots$$

# Method

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This is essentially a constrained system on the moduli space of the vortex

$$\begin{aligned} \mathcal{Z}_k^{N_c, N_f} &= \int \mathcal{D}^2 A \mathcal{D}^2 H \delta(\mu_r) J_r \delta(\mu_z) J_z \delta(\mu_{\bar{z}}) J_{\bar{z}} \cdots \\ &= \text{Vol}(\mathcal{M}_k^{N_c, N_f}) \end{aligned}$$

We can perform the path integral, which reduces to residue integrals over zero modes of Lagrange multiplier field  $\Phi$

$$\begin{aligned} \mathcal{Z}_k^{N_c, N_f}(\Sigma_h) &= \sum_{\sum_a k_a = k} (-1)^\sigma \int \prod_a \frac{d\phi_a}{2\pi} \frac{\prod_a (1 + \frac{N_f}{2\pi i \phi_a})^h \prod_{a < b} (i\phi_a - i\phi_b)^{2-2h}}{\prod_a (i\phi_a)^{N_f(1-h+k_a)}} \\ &\quad \times e^{2\pi i \sum_a \phi_a (\frac{g^2 c}{4\pi} \mathcal{A} - k_a)} \end{aligned}$$

Integrals are localized at poles (Localization formula)

# Results

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$N_c=N_f=1$  (Abrikosov–Nielsen–Olesen vortex)


$$\mathcal{Z}_k^{1,1}(\Sigma_h) = (2\pi)^{k-h} \sum_{j=0}^h \frac{h!}{j!(k-j)!(h-j)!} \left( \frac{g^2 c}{4\pi} \mathcal{A} - k \right)^{k-j}$$

We find

$$\mathcal{A} \geq \frac{4\pi}{g^2 c} k \quad \text{Bradlow limit}$$

For  $h=0$  (sphere)

$$\mathcal{Z}_k(S^2) = \frac{(2\pi)^k}{k!} \left( \frac{g^2 c}{4\pi} \mathcal{A} - k \right)^k \xrightarrow{\mathcal{A} \rightarrow \infty} \text{Vol}((S^2)^k / \mathcal{S}_k) \sim \frac{\mathcal{A}^k}{k!}$$

 Eq of state:  $P \left( \mathcal{A} - \frac{4\pi}{g^2 c} k \right) = kT$

# Results

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$N_c=2$ ,  $N_f$  (non-Abelian *semi-local* vortex) on the sphere

$$\mathcal{Z}_0^{2,N_f}(S^2) = \frac{2!}{(N_f - 1)!(N_f - 2)!} (2\pi \tilde{\mathcal{A}})^{2(N_f - 2)}$$

$$\begin{aligned} \mathcal{Z}_1^{2,N_f}(S^2) &= \frac{(2\pi)^{3N_f - 4}}{(2N_f - 1)(N_f - 1)!(2N_f - 3)!} \tilde{\mathcal{A}}^{N_f - 3} (\tilde{\mathcal{A}} - 1)^{2N_f - 3} \\ &\quad \times \left( (N_f - 2)\tilde{\mathcal{A}}^2 + 2(N_f + 1)\tilde{\mathcal{A}} + (N_f - 2) \right) \end{aligned}$$

$$\begin{aligned} \mathcal{Z}_2^{2,N_f}(S^2) &= 2(2\pi)^{4N_f - 4} \left[ \frac{-2}{(N_f - 1)!(3N_f - 1)!} \tilde{\mathcal{A}}^{N_f - 3} (\tilde{\mathcal{A}} - 2)^{3N_f - 3} \right. \\ &\quad \times \left( (2N_f^2 - 2N_f + 1)\tilde{\mathcal{A}}^2 + 2(2N_f + 1)(N_f - 1)\tilde{\mathcal{A}} + 2(N_f - 1)(N_f - 2) \right) \\ &\quad \left. + \frac{1}{(2N_f - 1)!(2N_f - 2)!} (\tilde{\mathcal{A}} - 1)^{4N_f - 4} \right] \end{aligned}$$

where  $\tilde{\mathcal{A}} \equiv \frac{g^2 c}{4\pi} \mathcal{A}$



# Comments

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On the sphere,  $k=0$  (perturbative part) gives the volume of the vacuum moduli space

$$\mathcal{Z}_0^{N_c, N_f}(S^2) = N_c! \times \text{Vol}(G_{N_c, N_f}) \tilde{\mathcal{A}}^{N_c(N_f - N_c)}$$

Vol. of Grassmannian

The case of  $N_c=N_f$  (non-Abelian *local* vortex) is rather special

$$\mathcal{Z}_0^{2,2}(S^2) = 2$$

$$\mathcal{Z}_1^{2,2}(S^2) = 2 \times (2\pi)^2 (\tilde{\mathcal{A}} - 1)$$

$$\mathcal{Z}_2^{2,2}(S^2) = 2 \times \frac{(2\pi)^4}{2!} \left( \tilde{\mathcal{A}}^2 - \frac{20}{6} \tilde{\mathcal{A}} + \frac{17}{6} \right)$$

In the limit of  $\mathcal{A} \rightarrow \infty$ , the divergence comes from the position moduli only

# Conclusion

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- ❖ We evaluate the volume of the moduli space of BPS vortices on the Riemann surface by using the localization formula
- ❖ We can see not only the volume itself but also the geometrical structure of the moduli space
- ❖ Using the deconstruction, we can also apply the localization method to the solitons in the Chern-Simons-Higgs system, etc.
- ❖ The volume of the moduli space (aka localization formula) is a key to understand the dualities between gauge theories, string theories, matrix models, and more!