

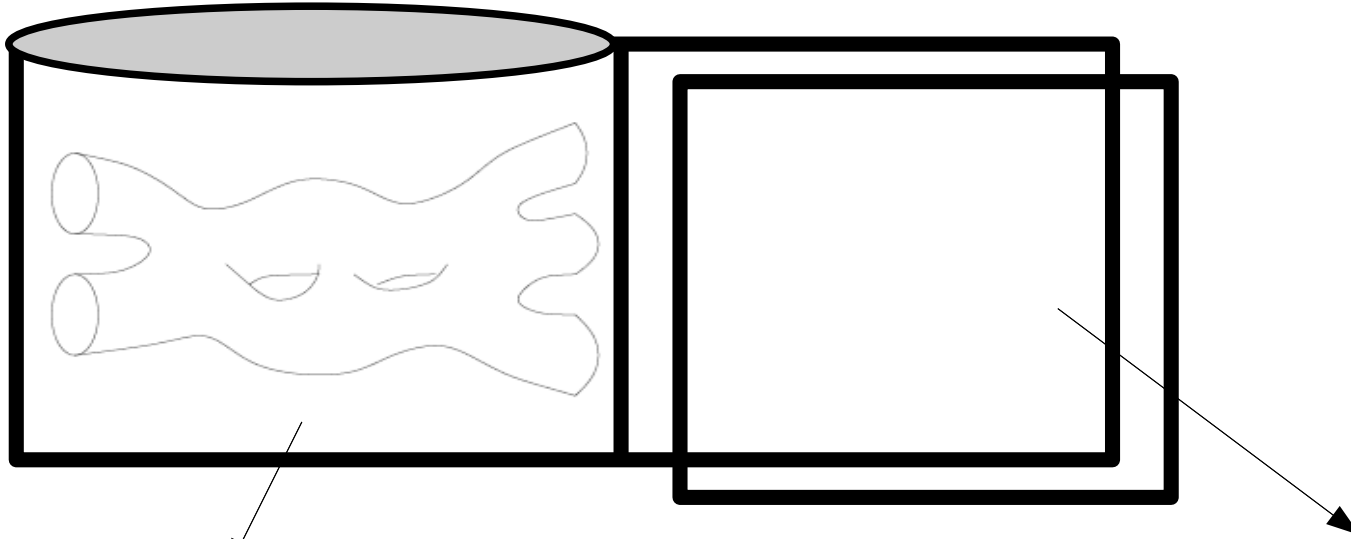
# **Seiberg-Witten prepotential from Liouville classical conformal block**

戴大盛 理研川合

Based on [1008.4332](#)

**JHEP 1010 107 (2010)**

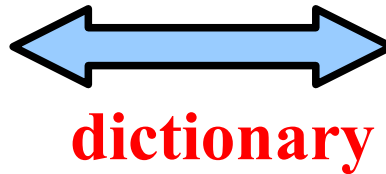
# Review: AGT conjecture



**LFT on Riemann surface  $C$**

**4D  $N=2$   $SU(2)$  SCFT  
(Coulomb phase)**

**Information of  $C$**



**properties of 4D QFT**

**Ex**

**Complex moduli of  $C$**

**UV gauge coupling**

**.... and so on**

# AGT dictionary

Given any  $C$ , one has certain SCFT !!

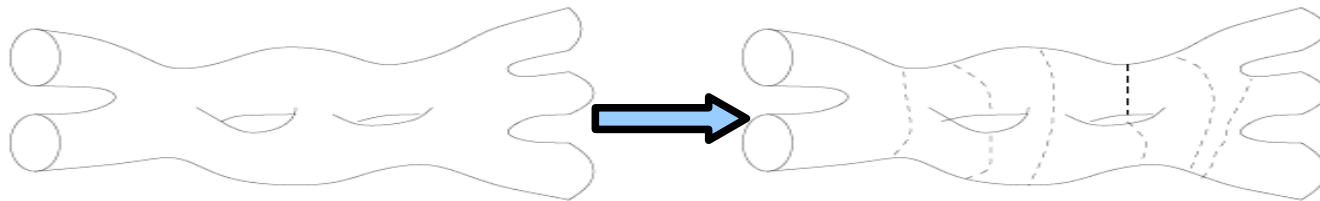
(I) # of gauge coupling (or group) =  $\chi=(3g-3+n)$

(II) flavor mass = weight of insertion

(III) Coulomb parameter = internal momentum

(IV) SW curve = double cover of  $C$

**Ex**



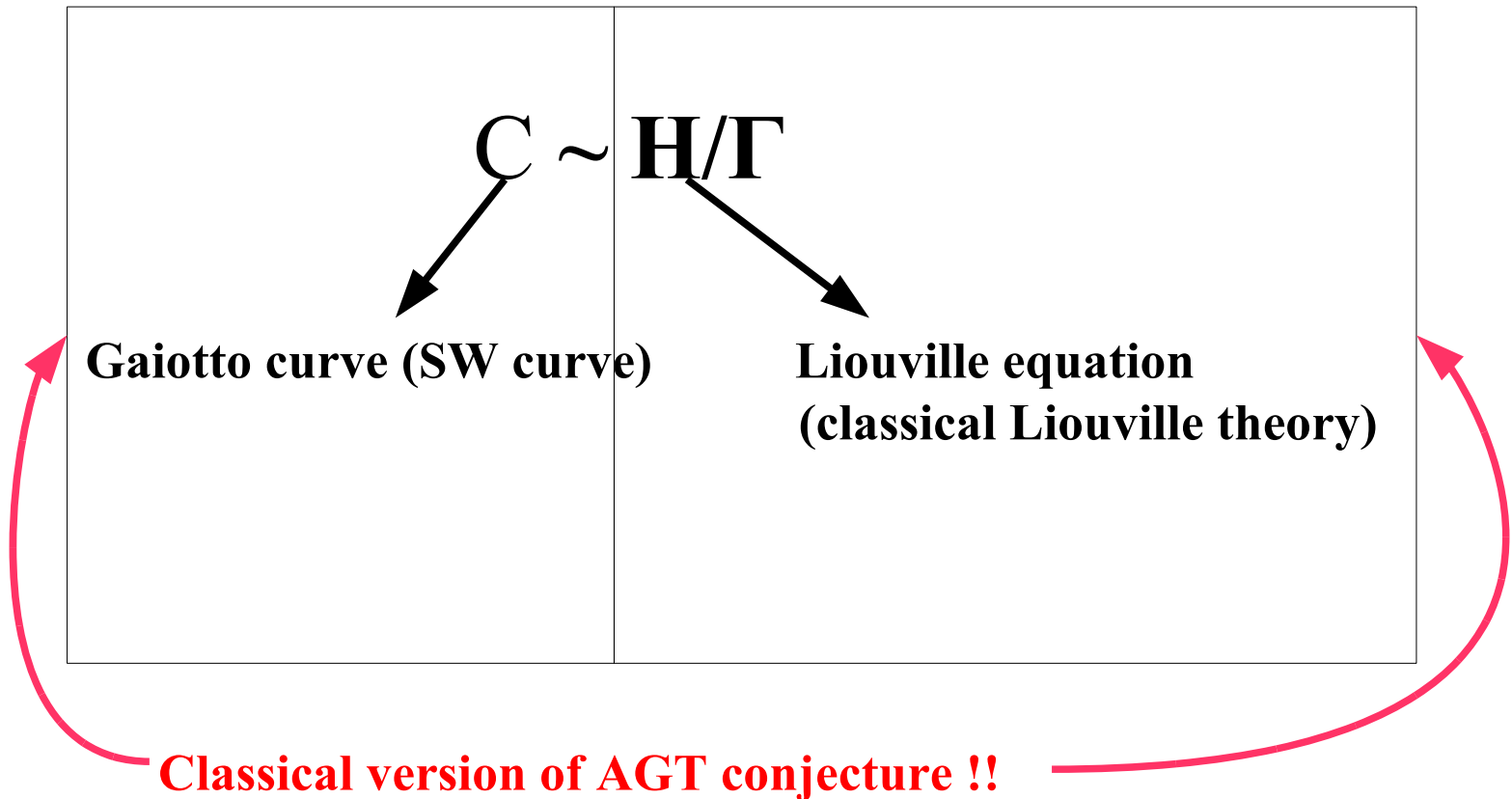
$\chi=(3g-3+n)=8$ , # of Coulomb parameters is eight!!

Ways of bootstrap are many...

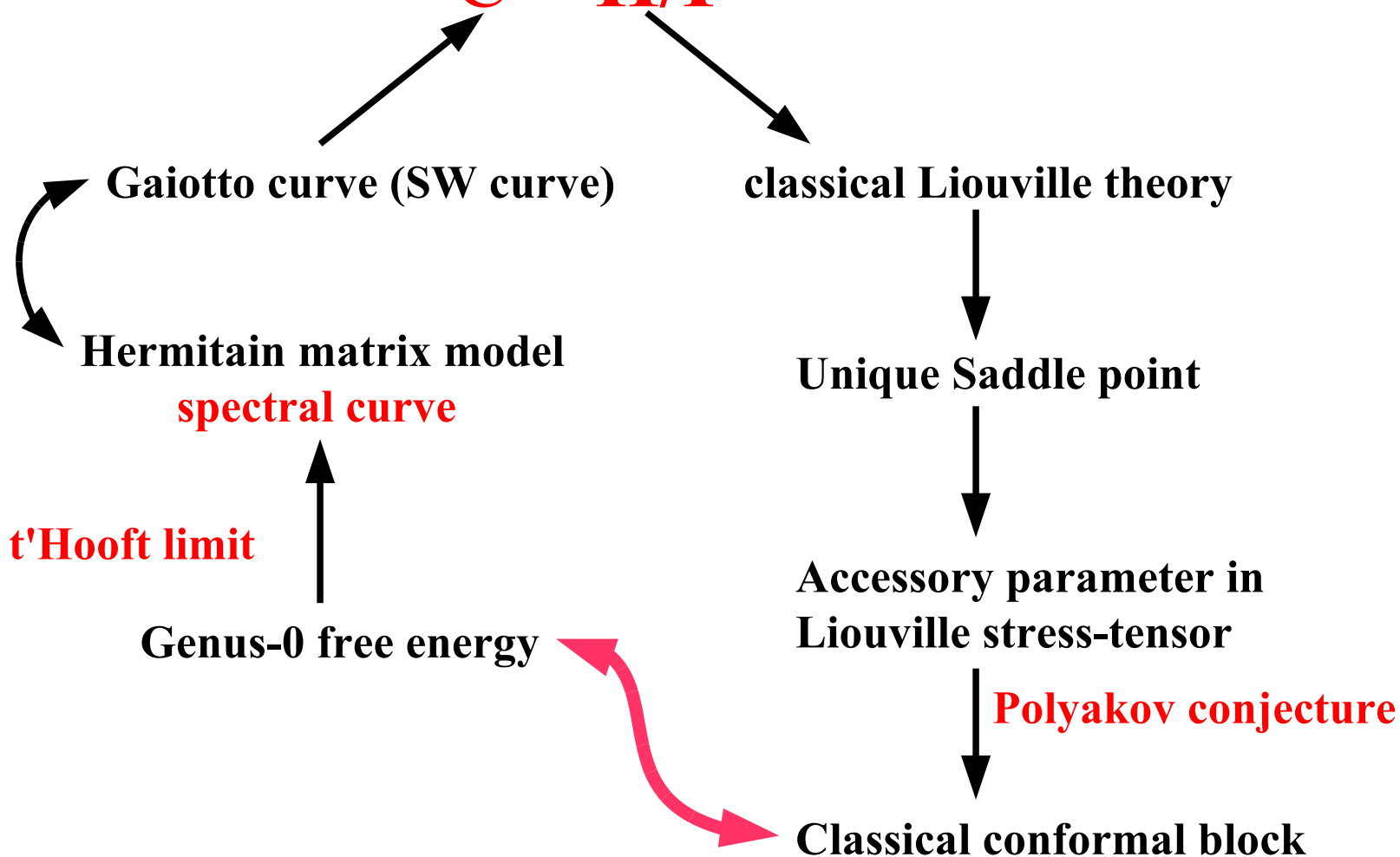
# Main idea is ...

Re-express  $C$  owing to Poincare, Koebe and Klein  
(Uniformization theorem, about 100 years ago)

- Uniform punctured Riemann surface by upper half-plane  $\mathbf{H}$  (universal cover)
- Endow  $C \sim \mathbf{H}/\Gamma$  with hyperbolic metric w/ discrete Fuchsian group  $\Gamma$



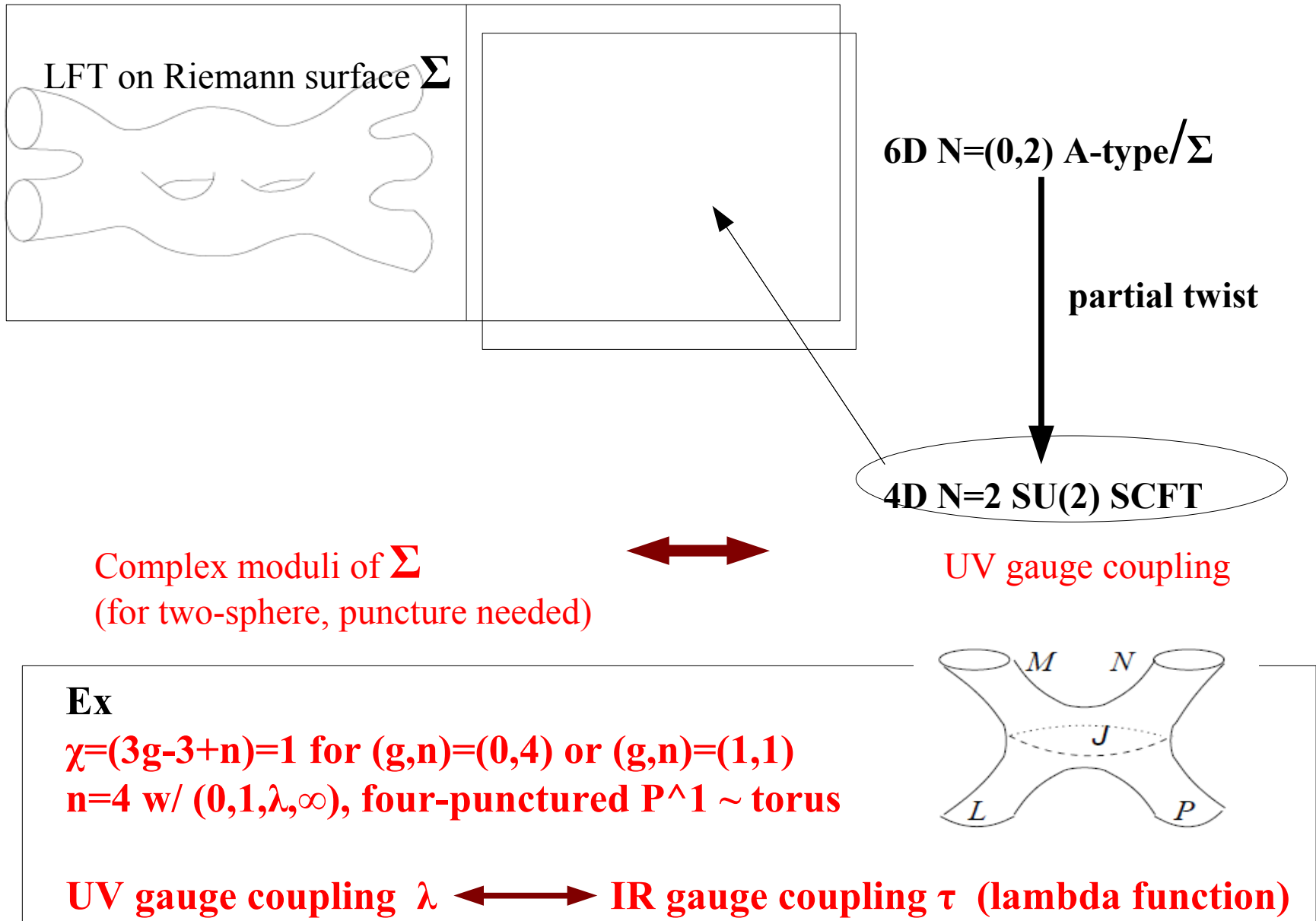
$$C \sim H/\Gamma$$



## Plan of talk

- Introduction—very brief review of AGT relation (2pp)
- What is Uniformization (均一化) (6pp)
- Big picture (2pp)
- Example I, II & III (4pp)  
(extract SW prepotential from classical conformal block)
- Summary (1pp)

# Introduction—very brief review of AGT relation



# AGT dictionary

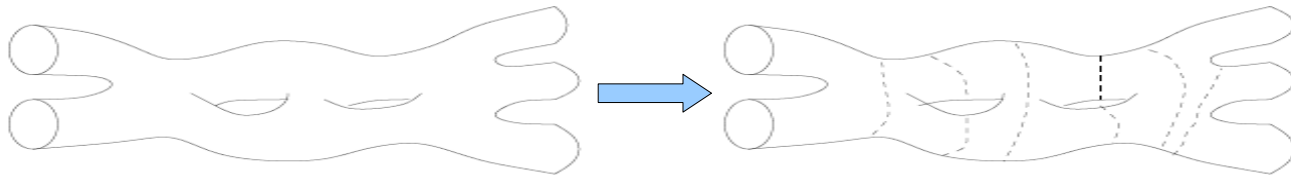
(I) # of gauge coupling (or group) =  $\chi=(3g-3+n)$

(II) flavor mass = weight of insertion

(III) Coulomb parameter = internal momentum

(IV) SW curve =  $\Sigma$

**Ex**



**$\chi=(3g-3+n)=8$ , # of Coulomb parameters is eight!!**

**Ways of bootstrap are many...**



# What is Uniformization (均一化)

- Uniform punctured Riemann surface by upper half-plane (universal cover)
- Endow  $\Sigma \sim \mathbf{H}/\mathbf{T}$  with hyperbolic metric
  - conformal factor  $\varphi$
  - negative constant curvature → Liouville equation
  - saddle point (e.o.m.) of Liouville field theory (LFT)
  - classical regime of LFT

**Dawn(曙) of the bridge in between LFT & SW curve!!**

**Without relying on AGT conjecture!!**

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From conformal factor  $\varphi$  we can have

→ stress-tensor T of LFT :  $T_L(z) = Q\partial_z^2\phi - (\partial_z\phi)^2$  /  $T(z) \equiv \frac{1}{2}\partial_z^2\varphi_{cl} - \frac{1}{4}(\partial_z\varphi_{cl})^2$

→ 2<sup>nd</sup> order ODE of Fuchsian type  $(\partial_z^2 + T(z))\Psi = 0$  on  $\Sigma$

→ determine T by either **monodromy** on  $\Sigma$  or **Polyakov conjecture (Ward identity)**

**Polyakov conjecture:**

$$c_i = -\frac{\partial\mathcal{S}[\varphi_{cl}(\delta_i, z_i)]}{\partial z_i}, \quad i \neq (0, 1, \infty).$$



$$c_2(x) = -\frac{\partial\mathcal{S}_{cl}[\xi_i; x]}{\partial x} = \left( \frac{\partial}{\partial x} f_\delta \left[ \begin{array}{cc} \delta_3 & \delta_2 \\ \delta_4 & \delta_1 \end{array} \right] (x) \right)_{p=p_s(x)}$$

**Basically, it is Ward identity of T (next page)**

**Make possible a direct connection to classical conformal block f**

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## Ward identity of T

$$\langle T(z) X_\alpha \rangle = \sum_{i=1}^n \left( \frac{\Delta_{\alpha_i}}{(z - z_i)^2} + \frac{\partial z_i}{z - z_i} \right) \langle X_\alpha \rangle \quad ; \quad X_\alpha = V_{\alpha_1}(z_1) \cdots V_{\alpha_n}(z_n).$$

$$T = \sum_{i=1}^{n-1} \left( \frac{\Delta_{\alpha_i}}{(z - z_i)^2} + \frac{c_i}{z - z_i} \right)$$

For four-punctured sphere

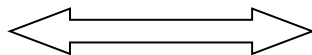
$$c_2(x) = -\frac{\partial \mathcal{S}_{cl}[\xi_i; x]}{\partial x} = \left( \frac{\partial}{\partial x} f_\delta \begin{bmatrix} \delta_3 & \delta_2 \\ \delta_4 & \delta_1 \end{bmatrix} (x) \right)_{p=p_s(x)}$$

Accessory parameter is determined by Polyakov conjecture,  
proved by mathematicians [Leon Takhtajan et al]

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$$c_2(x) = -\frac{\partial \mathcal{S}_{cl}[\xi_i; x]}{\partial x} = \left( \frac{\partial}{\partial x} f_\delta \begin{bmatrix} \delta_3 & \delta_2 \\ \delta_4 & \delta_1 \end{bmatrix} (x) \right)_{p=p_s(x)}$$

classical geometry  $\Sigma$



classical conformal block

**Reason: factorization of LFT action as  $b$  goes to 0**

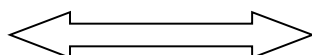
$$\mathcal{S}_{cl}[\xi_1, \dots, \xi_4; x] = \mathcal{S}^{(3)}(\delta_4, \delta_3, \delta_s) + \mathcal{S}^{(3)}(\delta_s, \delta_2, \delta_1) - f_{\delta_s, \delta_i}(x) - f_{\delta_s, \delta_i}(\bar{x}).$$

Ultimately, we find  $f$ =SW prepotential (instanton part)

**Reason: analogous to Hermitian matrix model at large- $N$**

$$\langle T_M(z) \rangle \rightarrow \mathcal{W}'(z)^2 + f(z) \quad c_2(x) = \frac{\partial}{\partial x} \hbar^{-2} \mathcal{F}_0 \quad Z = \frac{1}{\text{vol}U(N)} \int_{N \times N} dM \exp\left(\frac{1}{\hbar} \text{Tr} \mathcal{W}(M)\right) = \exp\left(\sum_{g \geq 0} \hbar^{2g-2} \mathcal{F}_g\right)$$

Gaiotto curve



Nekrasov partition function

?

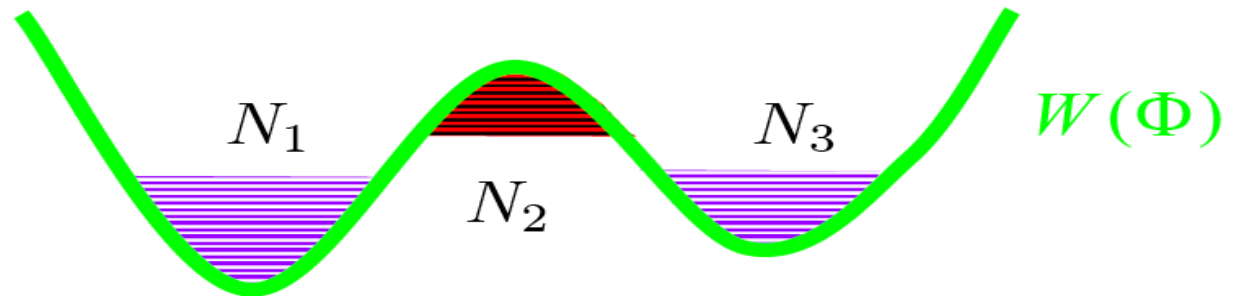
# Degression to Hermitian Matrix Model

$$Z = \int_{N \times N} d\Phi e^{\text{Tr} W(\Phi)/\hbar}$$

't Hooft limit

$$N \rightarrow \infty, \hbar \rightarrow 0, \hbar N = t \text{ constant}$$

Filling fractions  $t_i = \hbar N_i$

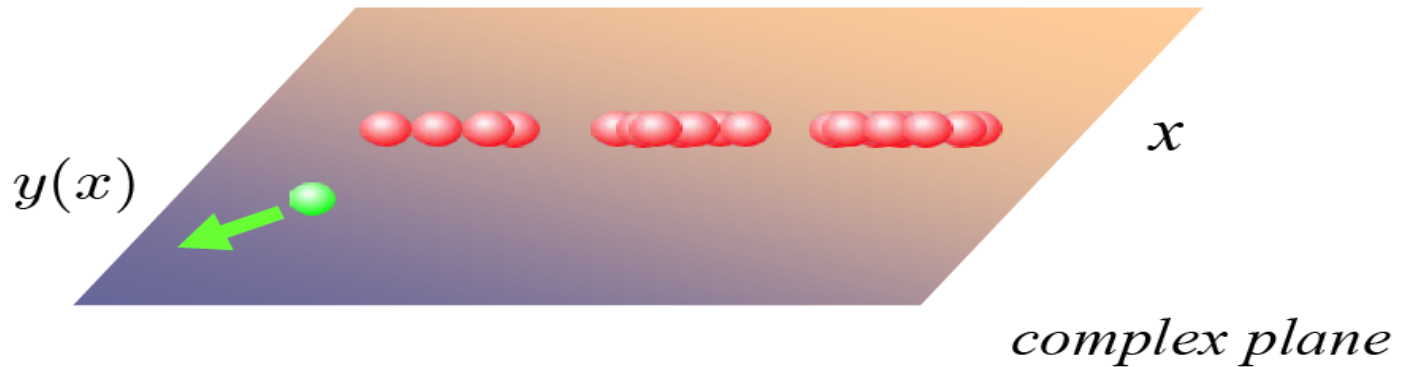


classical geometry  $\Sigma$

$$Z = \int d^N \lambda \prod_{i < j} (\lambda_i - \lambda_j)^2 \exp \sum_i W(\lambda_i) / \hbar$$

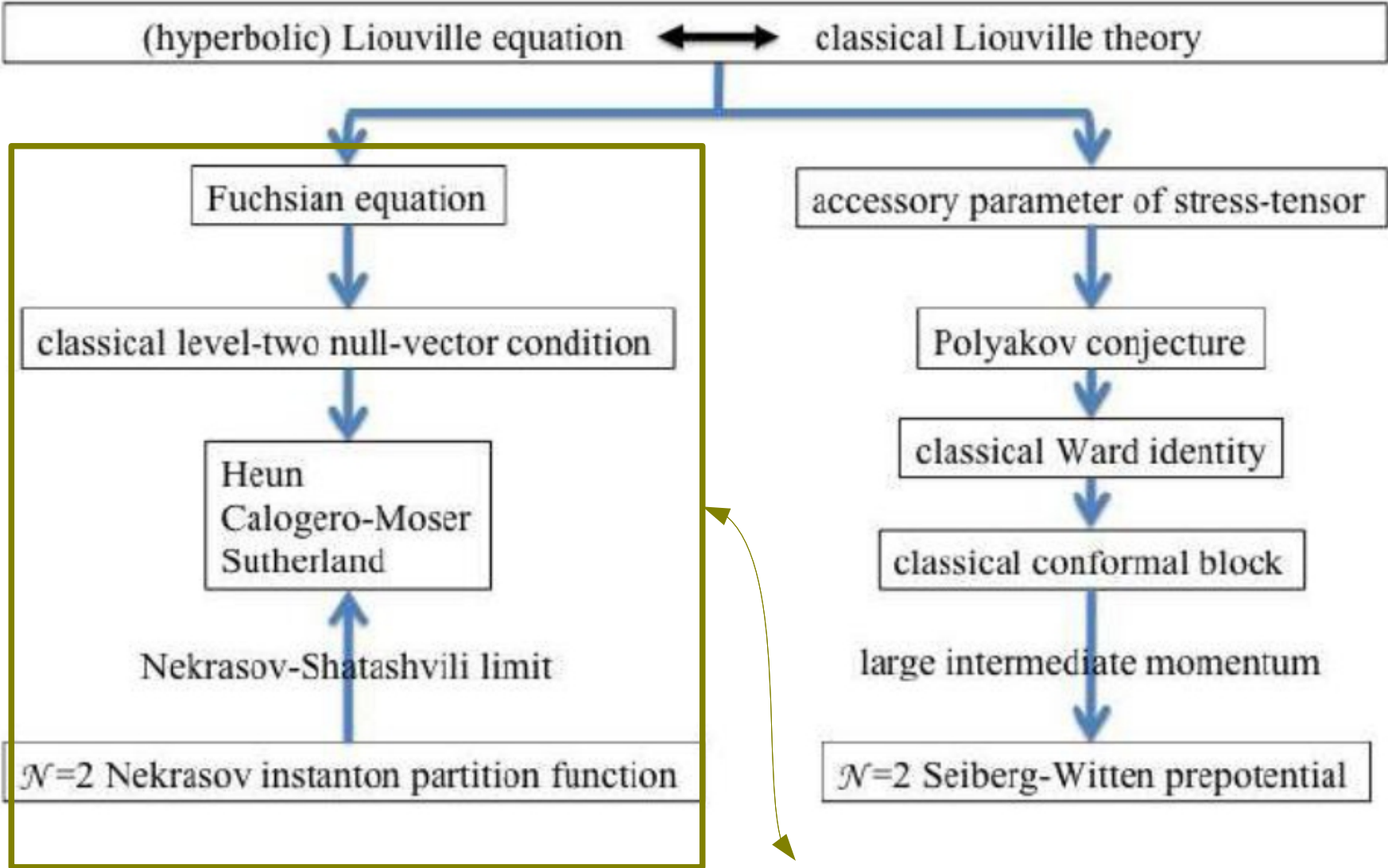
Effective force/resolvent

$$y(x) = \frac{dS_{eff}}{dx} = W'(x) - 2\hbar \text{Tr} \frac{1}{x - \Phi}$$



classical geometry  $\Sigma$

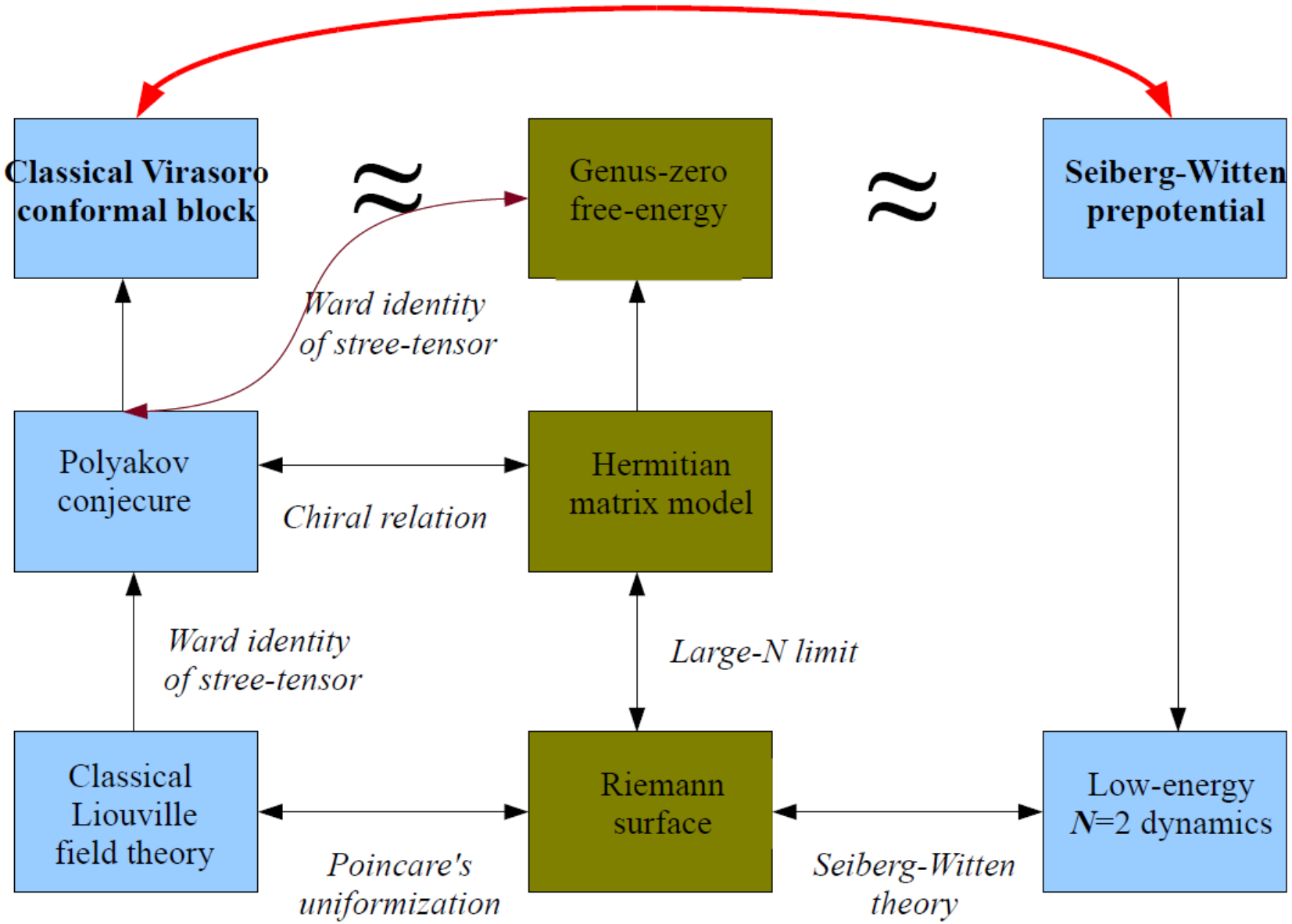
# uniformization



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# Why classical Virasoro conformal block reproduces Seiberg-Witten prepotential ?

Unexpected or not ...





# Example I, II & III

(extract SW prepotential from classical conformal block)

(I) All  $m_i = 0$

(II)  $m_1 = m_2 = \frac{\xi}{4}$  and  $m_3 = m_4 = 0$

(III) Arbitrary four flavor masses

bare flavor masses and weights assigned to punctures are related by

$$\begin{cases} \xi_1 = m_1 + m_2 + \frac{1}{2}, & \xi_2 = -m_1 + m_2 + \frac{1}{2}, \\ \xi_3 = m_3 + m_4 + \frac{1}{2}, & \xi_4 = -m_3 + m_4 + \frac{1}{2}, \end{cases}$$

# Example (I)

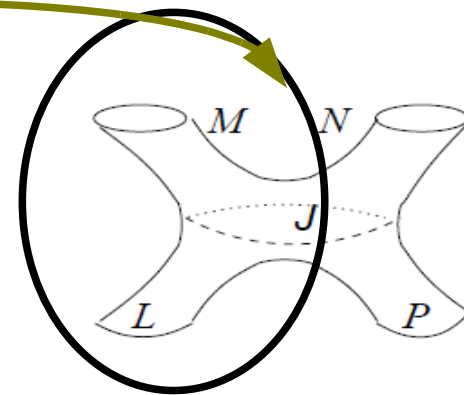
$$f_{\delta_s, \delta_i = \frac{1}{4}}(x)$$

$$= (p_s^2 - \frac{1}{4}) \log x + (p_s^2 + \frac{1}{4}) \frac{x}{2} + (\frac{13p_s^2}{16} + \frac{9}{32} + \frac{1}{256p_s^2 + 256}) \frac{x^2}{4} + \dots$$

$$\rightarrow a^2 \left( \log x + \frac{x}{2} + \frac{13x^2}{64} + \frac{23x^3}{192} + \dots \right)$$

$$a \equiv \oint dz \sqrt{T(z)} = \frac{\ell(\gamma_{12})}{4\pi b}$$

$$x \equiv \exp(2\pi i \tau_{UV})$$



Perfect agreement, up to a perturbative piece  $-\log 16$

**Example (II)**  $m_1 = m_2 = \frac{\xi}{4}$  and  $m_3 = m_4 = 0$


$$f_{\frac{1}{4}+p^2} \left[ \begin{array}{c} \frac{1}{4} \\ \frac{1}{4} \end{array} \right] (x) = \left( p^2 - \frac{1 - \xi^2}{4} \right) \log x + \left( \frac{1 - \xi^2}{8} + \frac{p^2}{2} \right) x \\ + \left( \frac{9(1 - \xi^2)}{128} + \frac{13p^2}{64} + \frac{(1 - \xi^2)^2}{1024(1 + p^2)} \right) x^2 + \mathcal{O}(x^3)$$

$$\mathcal{F}_{inst}^{SW} = (a^2 - m^2) \log x + (a^2 + m^2) \frac{x}{2} + \left( 13a^2 + 18m^2 + \frac{m^4}{a^2} + \mathcal{O}(a^{-4}) \right) \frac{x^2}{64} + \mathcal{O}(x^3)$$

carrying out  $(a^2, m^2) \rightarrow (p^2, -\frac{\xi^2}{4})$ . up to a perturbative piece

## Example (III) Arbitrary four flavor masses

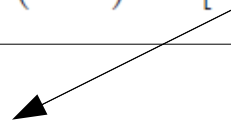
$$\begin{aligned} f_{\delta, \delta_i}(x) &= (\delta - \delta_1 - \delta_2) \log x + \frac{(\delta + \delta_1 - \delta_2)(\delta + \delta_3 - \delta_4)}{2\delta} x + \mathcal{O}(x^2) \\ &\rightarrow (a^2 + 2m_1^2 + 2m_2^2) \log x + \frac{a^4 - 4a^2(m_1m_2 + m_3m_4) + 16m_1m_2m_3m_4}{2a^2} x + \dots \end{aligned}$$



$$(a^2 - m^2) \log x + \frac{(a^4 + 2a^2m^2 + m^4)}{2a^2} x + \dots$$

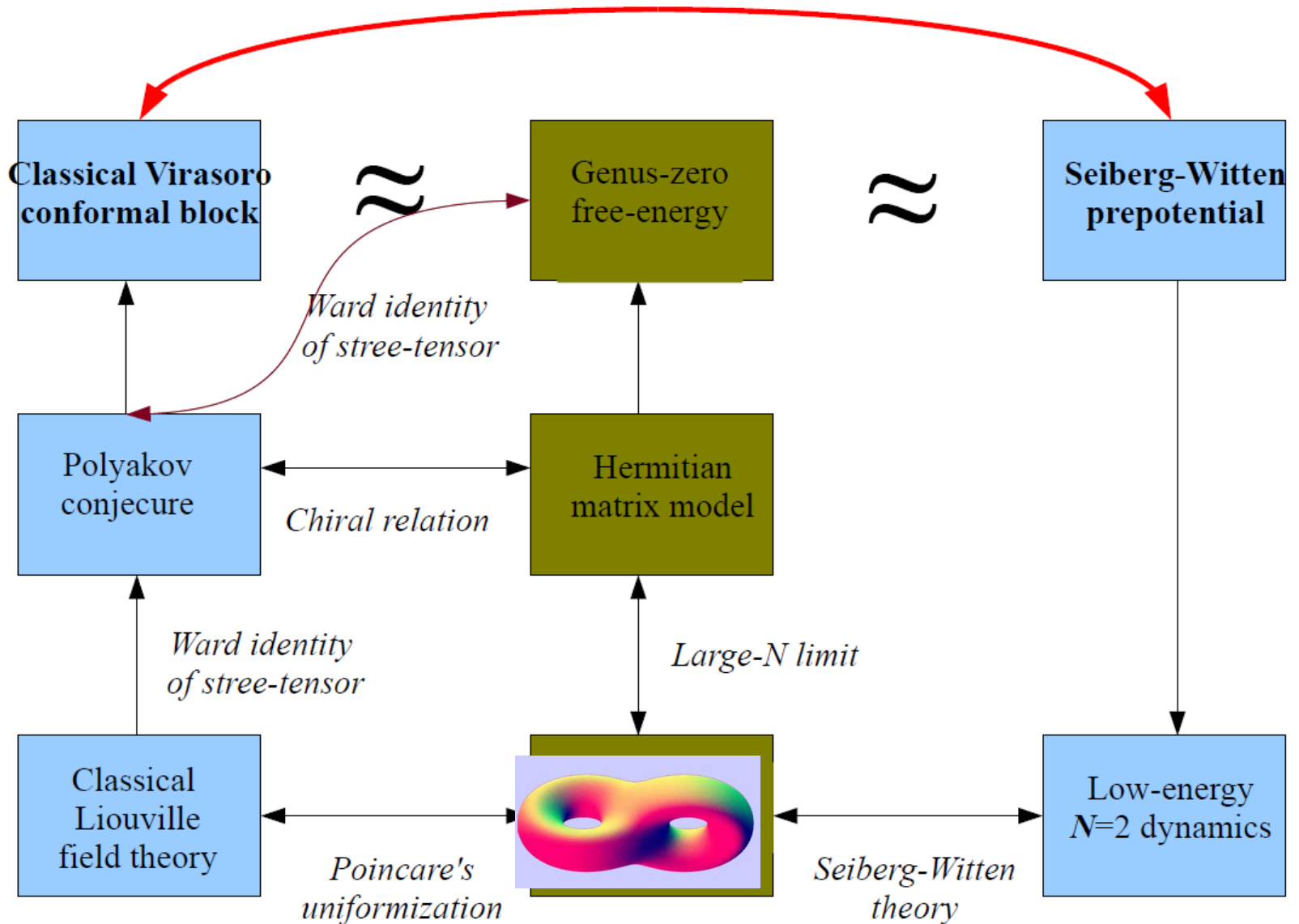
**Especially,**

for  $m_1 = m_2 = m_3 = m_4 = \frac{i}{2}m$  which agrees to (3.48) of [24] up to the  $U(1)$  part

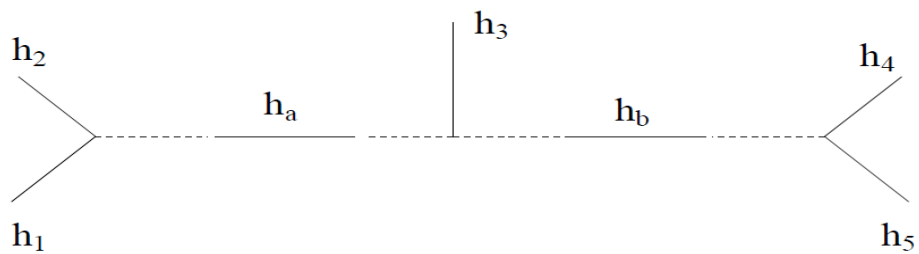
  
Eguchi & Maruyoshi JHEP 1007 (2010) 081  
arXiv:1006.0828

# Why classical Virasoro conformal block reproduces Seiberg-Witten prepotential ?

Unexpected or not ...



**THE END**



$$\mathcal{F}_{g=0}^{5pt} = 1 + \frac{(-h_1 + h_2 + h_a)(h_3 + h_a - h_b)}{2h_a} q_1 + \frac{(-h_4 + h_5 + h_b)(h_3 - h_a + h_b)}{2h_b} q_2 + \dots$$