

# 多軌道を持つ引力フェルミ原子気体に おける相分離

小林恵太, 太田幸宏, 奥村雅彦, 山田進,  
町田昌彦

# Introduction: ultra cold fermi gas

---

## Controllability to unusual degree of freedom

configuration of trap potential (harmonic trap, optical lattice )  
multi-component  
population imbalance (spin imbalance)  
interaction strength (repulsive to attractive)

## Superfluidity of two component fermi gas

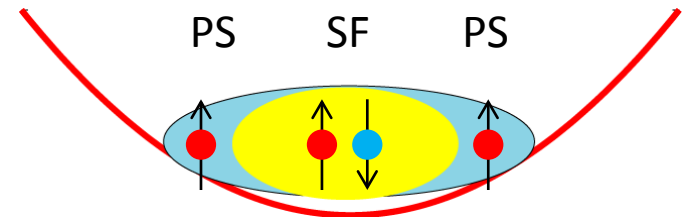
BCS-BEC crossover

FFLO state

Phase separation of polarized component in trap potential

A. E. Feiguin and F. Heidrich-Meisner Phys. Rev. B. **76**, 220508(R) (2007).

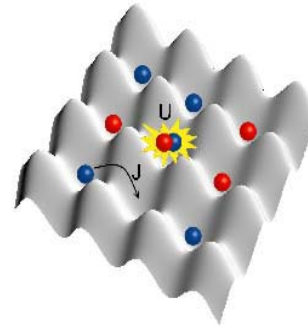
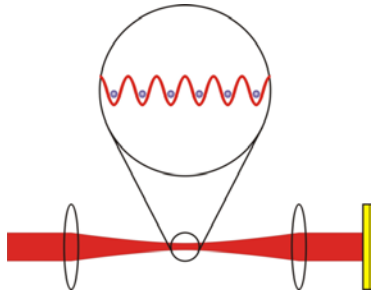
Masaki Tezuka, and Masahiko Ueda, Phys. Rev. Lett. 100, 110403 (2008).



Two component fermi gas in harmonic trap potential

# Introduction: optical lattice

Optical lattice: periodic potential created by laser beams



## multi-band(p-band)

Experimentally atom gas with p-band is achieved

Various phenomena are expected  
(Orbital order, magnetization with high spin etc....)

Phys. Rev. Lett. 99, 200405  
(2007)

Nature. Phys, Vol7, 147, (2011)

Phys. Rev. B 77, 174431 (2008)

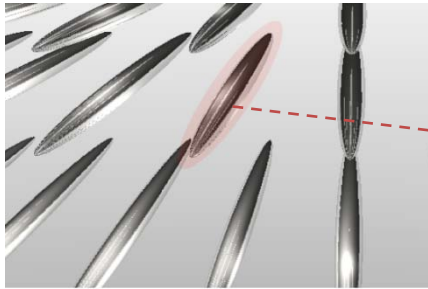
Phys. Rev. Lett, 100, 200406 (2008)

Multi band Superfluidity

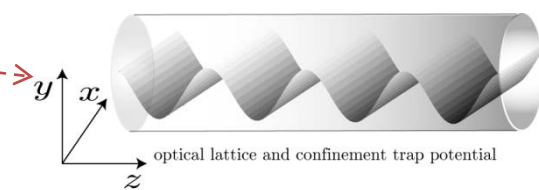
Multi-band effects of attractive fermi gas in 1D optical lattice

# contents

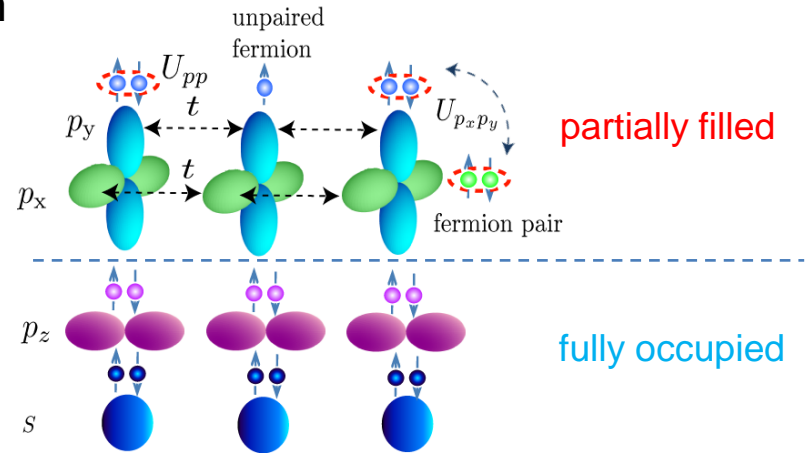
## 1D p-orbitals (two degenerate multi-band) system



1-D optical lattices



optical lattice and confinement trap potential



## Approach

Density matrix renormalization study  
Effective Hamiltonian in strong coupling limit

## Results

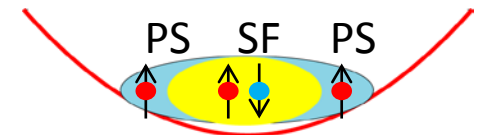
### Haldane insulator phase

Phys. Rev. Lett. 97, 260401 (2006), Phys. Rev. B. 81, 020408(R) (2010).

### Luther-Emery phase

Phys. Rev. Lett. 33, 589 (1974).

Strong phase separation of polarized component without trap potential



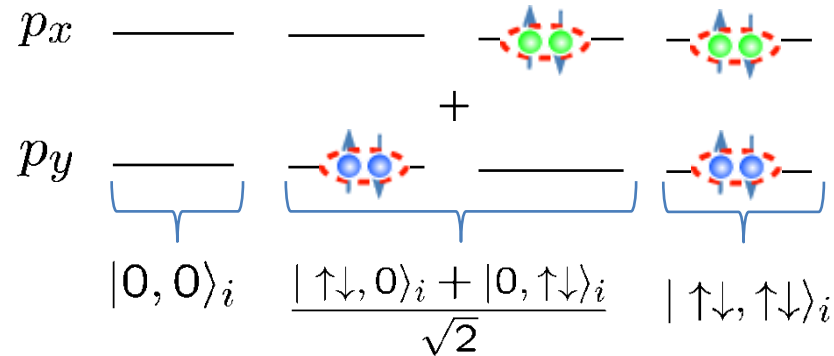
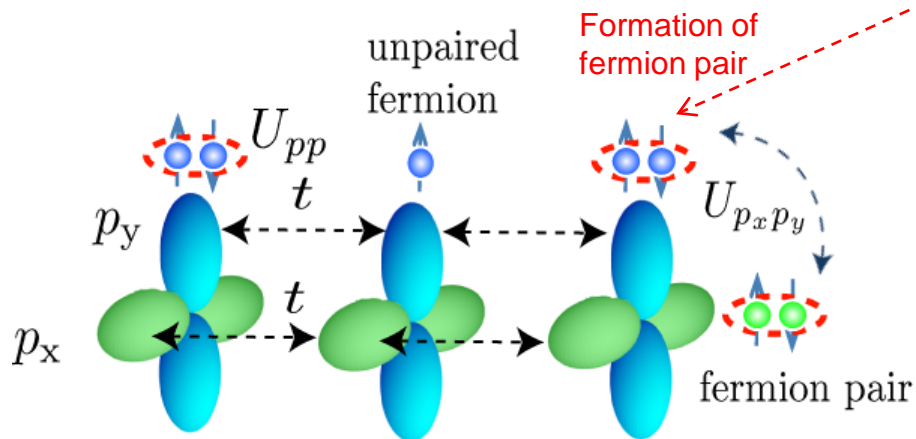
Phase separation in harmonic trap potential (single band)

# Model and Formulation

p-orbitals tight binding Hamiltonian with  $U_{pp}, U_{pxpy} < 0$

$$H - \mu N = - \sum_{p,\sigma} \sum_{\langle i,j \rangle} t c_{p,\sigma,i}^\dagger c_{p,\sigma,j} - \sum_{p,\sigma,i} \mu n_{p,\sigma,i} + \sum_i \left[ \sum_p U_{pp} n_{p,\uparrow,i} n_{p,\downarrow,i} \right. \\ \left. + \sum_{p \neq p'} U_{pp'} \left\{ c_{p,\uparrow,i}^\dagger c_{p',\downarrow,i}^\dagger c_{p,\downarrow,i} c_{p',\uparrow,i} + n_{p,\uparrow,i} n_{p',\downarrow,i} + c_{p,\uparrow,i}^\dagger c_{p,\downarrow,i}^\dagger c_{p',\downarrow,i} c_{p',\uparrow,i} \right\} \right]$$

negative pair hopping term



pseudo spin 1 states

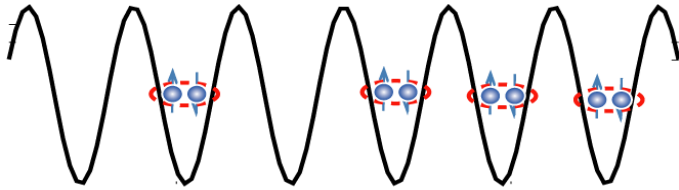
$$t = - \int dz w_{p,i+1} \left[ \frac{-\hbar^2 \nabla^2}{2M} + V_{\text{opt}} \right] w_{p,i}$$

$$U_{pp'} = g \int d^3x w_{p,i}^2 w_{p',i}^2 \quad \text{Wannier function (p-orbital)}$$

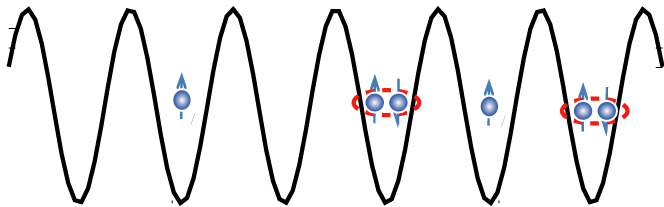
# DMRG study

---

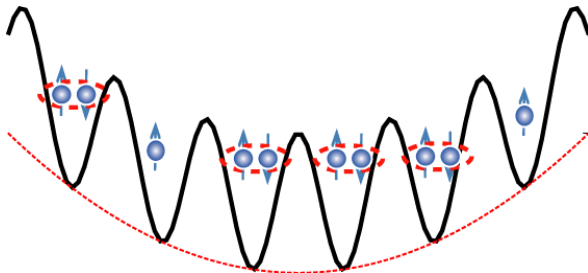
- Without spin imbalance



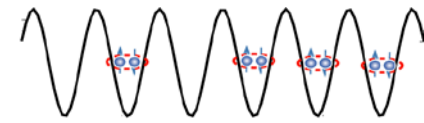
- With spin imbalance



- Effect of harmonic trap potential



# Result (without spin imbalance)



Open boundary condition

Total lattice site number:  $L = 100$

Filling rate:  $\tilde{n} = \sum_{p,i,\sigma} n_{p,i,\sigma} / 2L$  ( $0 \leq \tilde{n} \leq 2$ )

Inter-band interaction:  $U_{pp}/t = -10$

Intra-band interaction:  $U_{p_x p_y}/t = 4U_{pp}/9$

Spin imbalance:  $P = \sum_{p,i} (n_{p,\uparrow,i} - n_{p,\downarrow,i}) = 0$

Effective Hamiltonian (large U limit)

1-D pseudo spin 1 XXZ model

$$H_{\text{eff}} = J_{\text{ex}} \sum_{\langle i,j \rangle} \left( \rho_i^{(z)} \rho_j^{(z)} - \rho_i^{(+)} \rho_j^{(-)} / 2 \right) - \sum_i 2\bar{\mu} \rho_i^z$$

pseudo spin 1 operator

$$\rho_i^{(+)} = \sum_p c_{p,\uparrow,i}^\dagger c_{p,\downarrow,i}^\dagger$$

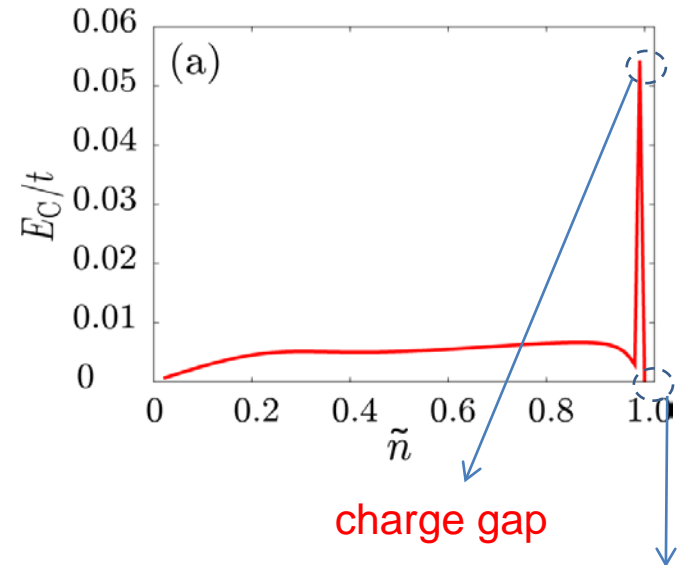
$$\rho_i^{(-)} = [\rho_i^{(+)}]^\dagger$$

$$\rho_p^{(z)} = \sum_\sigma (\sum_i n_{p,\sigma,i} - 1) / 2$$

$$\bar{\mu} = \mu + \frac{|U_{pp}| + |U_{p_x p_y}|}{2}$$

$$J_{\text{ex}} = \frac{2t^2}{|U_{pp}| + |U_{p_x p_y}|}$$

charge gap:  $E_C \equiv E(N+ \uparrow\downarrow) + E(N- \uparrow\downarrow) - 2E(N)$



Haldane phase  
spin gap  
edge 1/2 spin

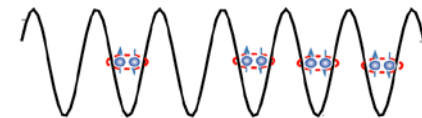
Phys. Lett. A 93, 464 (1983)  
Phys. Rev. Lett. 59, 799 (1987)

gapless  
excitation

Haldane insulator phase

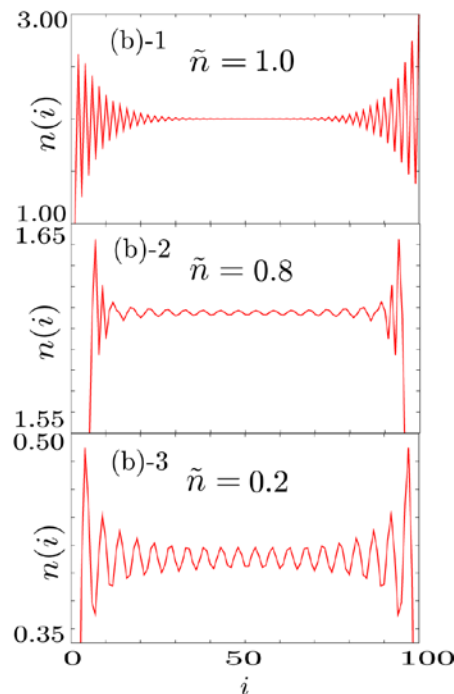
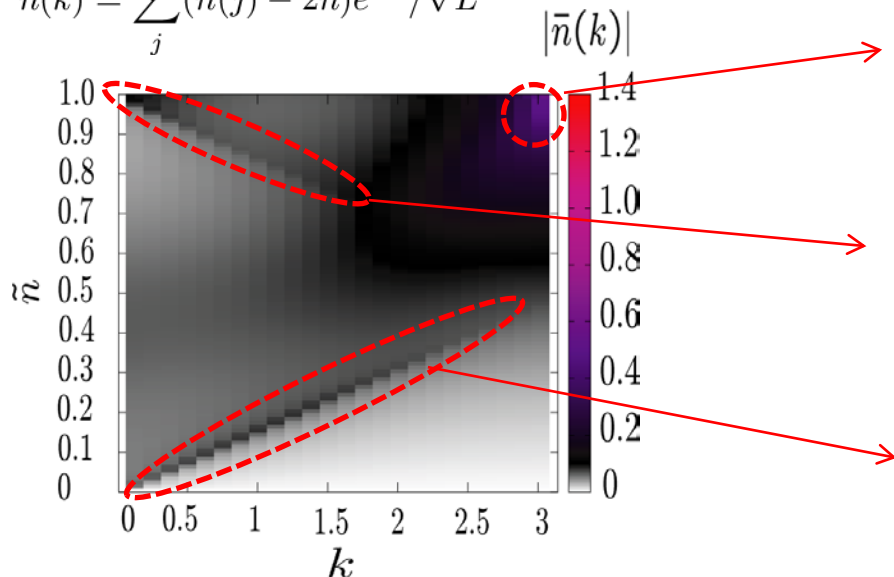
pseudo spin gap  
→ charge gap  
edge pseudo 1/2 spin  
→ gapless excitation at the edges

# Result (without spin imbalance)



Absolute value of Fourier transformation of density fluctuation

$$\tilde{n}(k) = \sum_j (n(j) - 2\tilde{n}) e^{ikj} / \sqrt{L}$$



**strong peak**  
edge state: staggered CDW  
with exponential decay

**weak peak**  
CDW

Periodicity of CDW

$$H_{\text{eff}} = J_{\text{ex}} \sum_{\langle i,j \rangle} \left( \rho_i^{(z)} \rho_j^{(z)} - \rho_i^{(+)} \rho_j^{(-)} / 2 \right) - \sum_i 2\bar{\mu} \rho_i^z$$

$$\rho_i^{(+)} = \sum_p c_{p,\uparrow,i}^\dagger c_{p,\downarrow,i}^\dagger \quad \rho_i^{(-)} = [\rho_i^{(+)}]^\dagger \quad \rho_p^{(z)} = \sum_\sigma (\sum_i n_{p,\sigma,i} - 1) / 2$$



$$2k_F = 2\pi \left( 1 - \left| \sum_i \rho_i^{(z)} \right| \right) = 2\pi \tilde{n}$$

Fermi vector in equivalent spinless fermion

Spin 1 Heisenberg model with magnetic field  
[Gabor Fath, Phys. Rev. B **68**, 134445, (2003).]

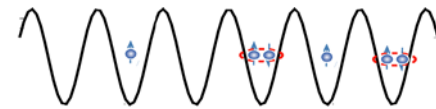
**Luther-Emery phase**

Superconducting pair correlation  $\langle \rho_i^{(+)} \rho_0^{(-)} \rangle$   
Density-density correlation  $\langle \rho_i^{(z)} \rho_0^{(z)} \rangle$

developed with power-low decay



# Result (with spin imbalance)



Spin imbalance:

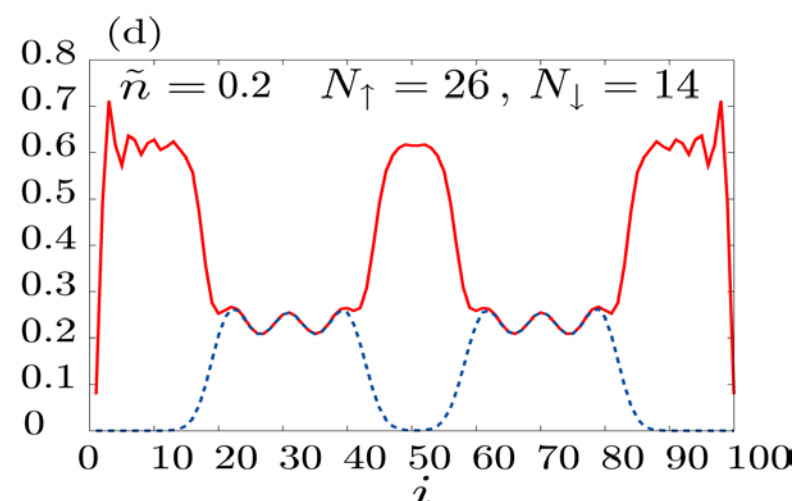
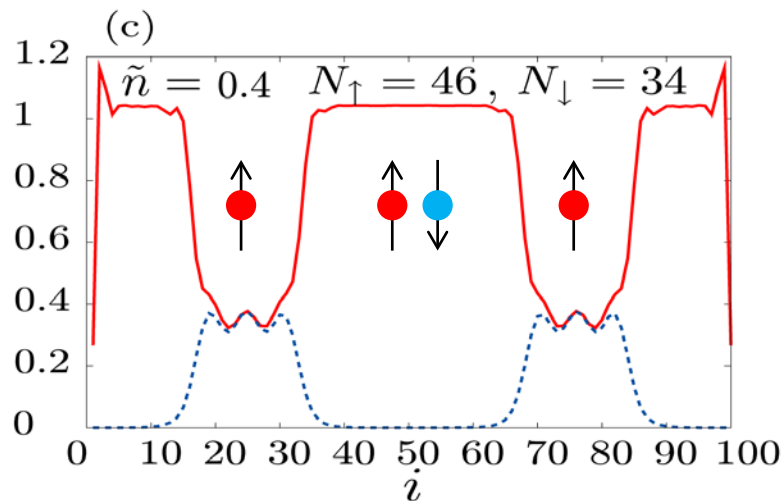
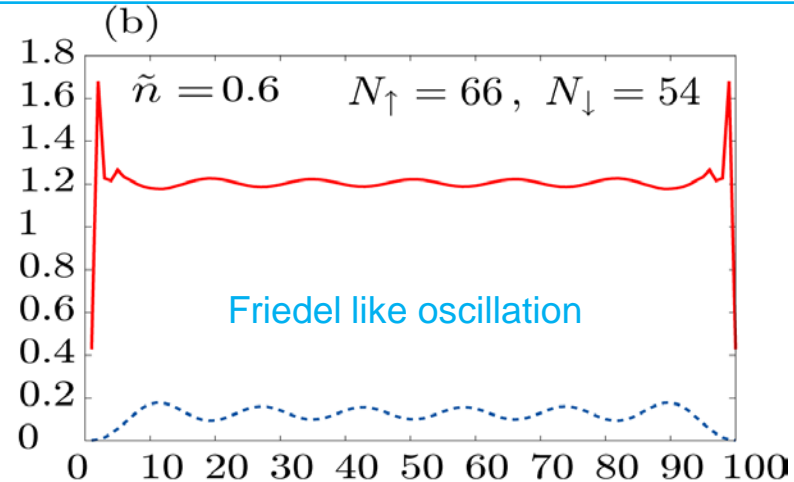
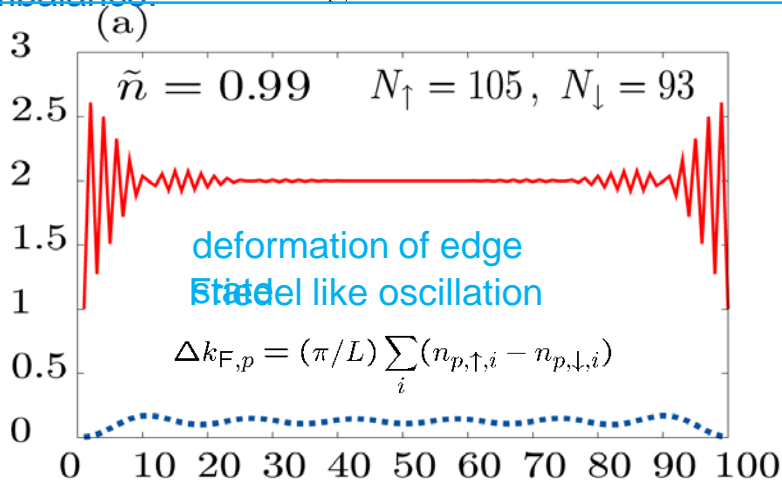
$$P = \sum_{p,i} (n_{p,\uparrow,i} - n_{p,\downarrow,i}) = 12$$

density profile

$$n_{\uparrow}(i) + n_{\downarrow}(i)$$

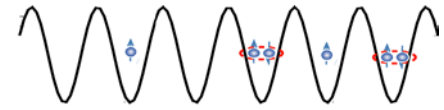
spin density profile

$$n_{\uparrow}(i) - n_{\downarrow}(i)$$



Strong phase separation of polarized component without trap potential

# Result (with spin imbalance)



Effective Hamiltonian with spin imbalance (large U limit)

$$H_{\text{eff}} = P \left[ - \sum_{p,\sigma} \sum_{\langle i,j \rangle} \underline{tc_{p,\sigma,i}^\dagger c_{p,\sigma,j}} - 2\bar{\mu} \sum_i \rho_i^{(z)} - \frac{J_{\text{ex}}}{2} \sum_{\langle i,j \rangle} \rho_i^{(+)} \rho_j^{(-)} + J_{\text{ex}} \sum_{\langle i,j \rangle} \rho_i^{(z)} \rho_j^{(z)} \right] P$$

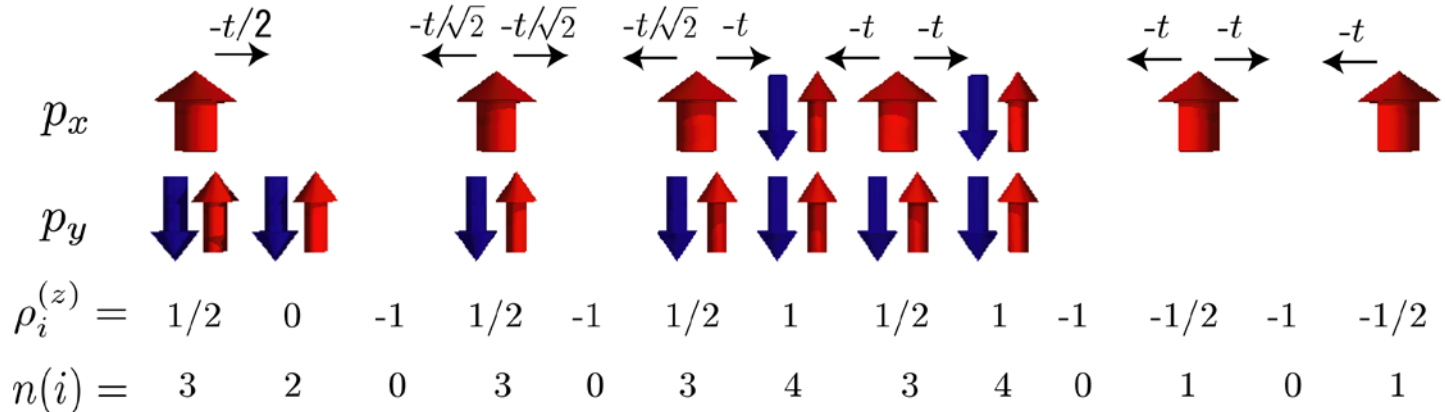
$$- \sum_{pp'} P \left[ J_{\text{ex}} \left( n_{p,\uparrow,i}^{(z)} + \frac{1}{2} \right) \left( n_{p',\uparrow,j}^{(z)} + \frac{1}{2} \right) \right] P$$

$$\bar{\mu} = \mu + \frac{|U_{pp}| + |U_{p_r p_y}|}{2} \quad J_{\text{ex}} = \frac{2t^2}{|U_{pp}| + |U_{p_r p_y}|}$$

Projection operator P

pseudo spin 1 state  $\mathcal{H}_{\rho=1} = \{ |0,0\rangle, (|\uparrow\downarrow,0\rangle + |0,\uparrow\downarrow\rangle)/\sqrt{2}, |\uparrow\downarrow,\uparrow\downarrow\rangle \}$   
 unpaired fermion  $\mathcal{H}_\uparrow = \{ |\uparrow,0\rangle, |0,\uparrow\rangle, |\uparrow,\uparrow\downarrow\rangle, |\uparrow\downarrow,\uparrow\rangle, |\uparrow,\uparrow\rangle \}$

hopping process of unpaired fermion



CDW like configuration

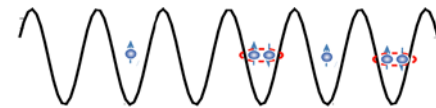
high density

low density

Transfer energy of unpaired fermion is reduced

unpaired fermion tend to form low (or high ) density region

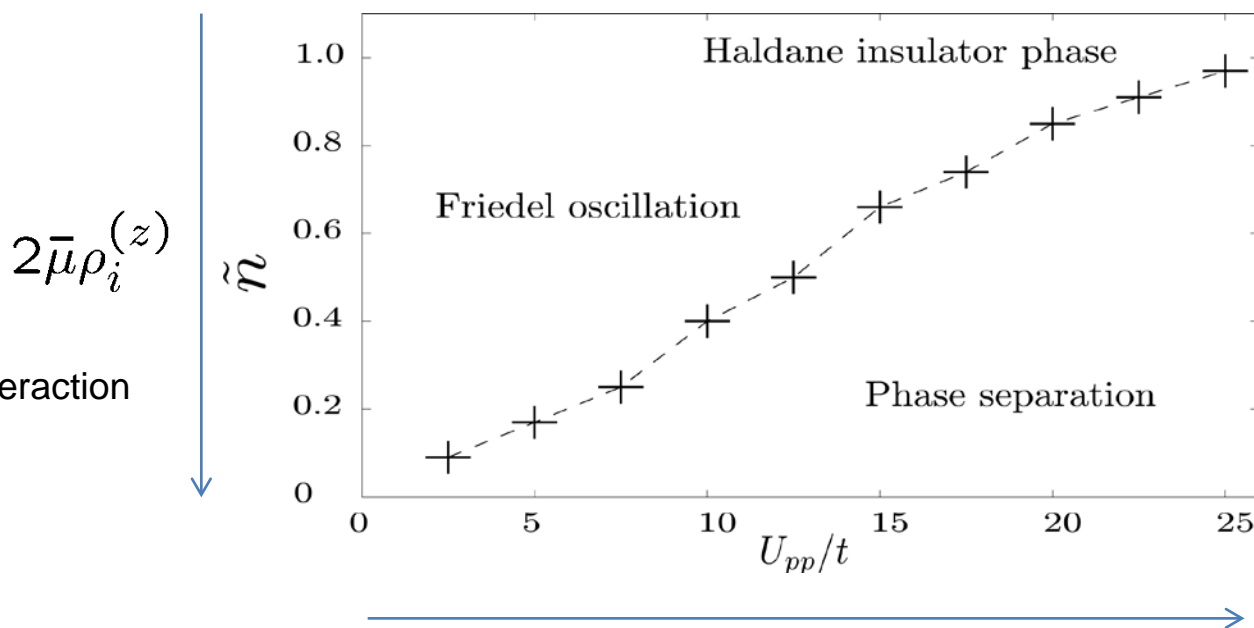
# Result (with spin imbalance)



$$H_{\text{eff}} = P \left[ - \sum_{p,\sigma} \sum_{\langle i,j \rangle} t c_{p,\sigma,i}^\dagger c_{p,\sigma,j} - 2\bar{\mu} \sum_i \rho_i^{(z)} - \frac{J_{\text{ex}}}{2} \sum_{\langle i,j \rangle} \rho_i^{(+)} \rho_j^{(-)} + J_{\text{ex}} \sum_{\langle i,j \rangle} \rho_i^{(z)} \rho_j^{(z)} \right] P$$

$$- \sum_{pp'} P \left[ J_{\text{ex}} \left( n_{p,\uparrow,i}^{(z)} + \frac{1}{2} \right) \left( n_{p',\uparrow,j}^{(z)} + \frac{1}{2} \right) \right] P$$

phase diagram

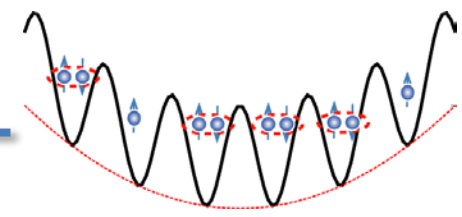


pseudo spin-spin interaction become irrelevant

$$J_{\text{ex}} = \frac{2t^2}{|U_{pp}| + |U_{p_x p_y}|}$$

pseudo spin-spin interaction become irrelevant

# Result (effect of harmonic trap potential)



Multi-band Hamiltonian + harmonic trap potential

$$\sum_{p,\sigma,i} V_{ho}(i) n_{p,\sigma,i}$$

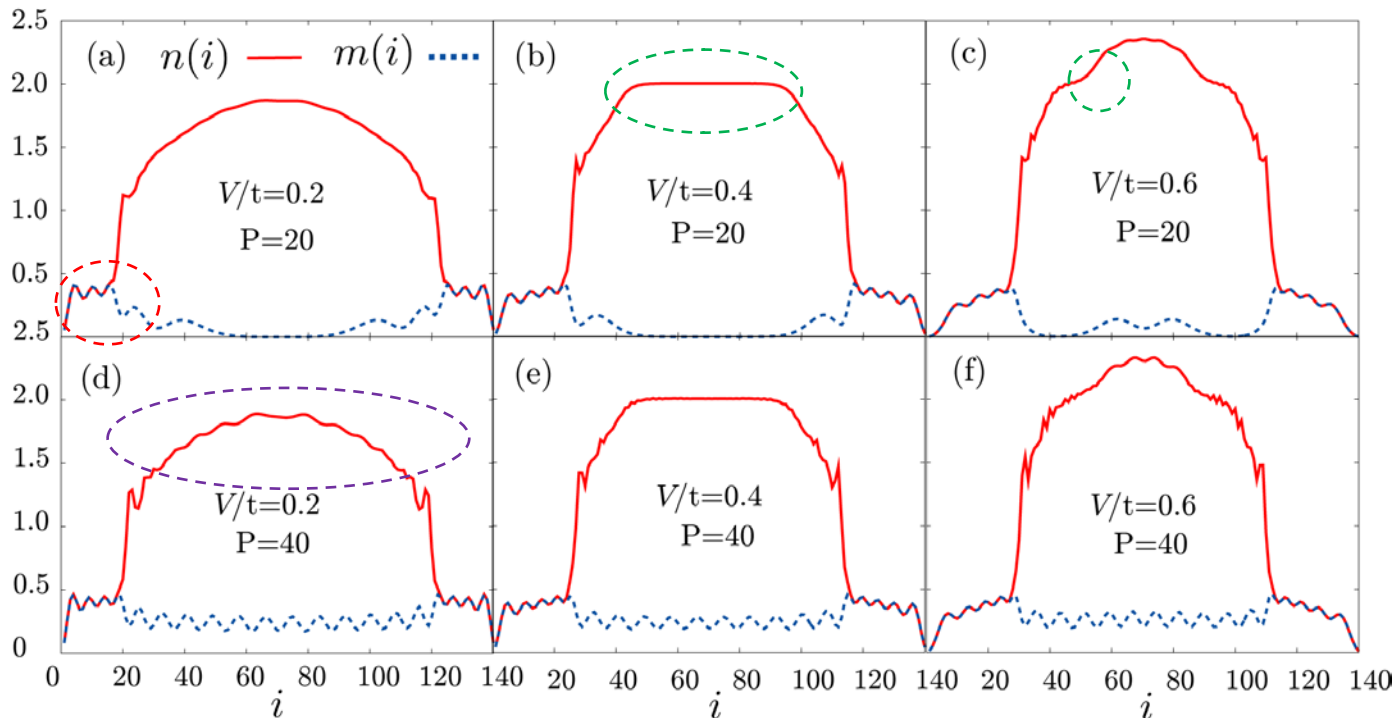
$$V_{ho}(i) = V(2/(L-1))^2 (i - (L+1)/2)^2$$

$$L = 140$$

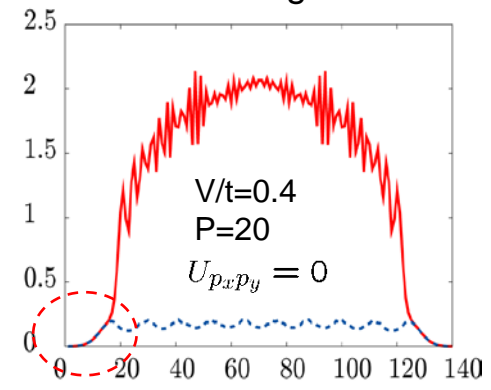
$$U_{pp}/t = -10 \quad U_{pxpy}/t = 4U_{pp}/9$$

$$P = \sum_{p,i} (n_{p,\uparrow,i} - n_{p,\downarrow,i})$$

Result of multi-band



Result of single-band



Strong phase

separation

Mott core like structure: Haldane insulator phase (polarization is excluded)

Friedel like oscillation (polarization and density)

# Summary

---

1D p-orbitals fermi gas with attractive interaction

## Various quantum phase

Haldane insulator phase

Luther-Emery phase

Friedal oscillation of density and spin density profiles

Strong phase separation of polarized component with out trap potential

## Effect of trap potential

Various phase can coexist by controlling filling rate, strength of trap potential and spin imbalance rate.