

Potential Analysis in Holographic Schwinger Effect

Yoshiki Sato (Kyoto Univ.)

In collaboration with Kentaroh Yoshida (Kyoto Univ.)

Abstract : We consider the Schwinger effect from the viewpoint of quark-anti quark potential. The potential analysis gives a strong support for Semenoff-Zaremba's proposal from another perspective.

★What is the Schwinger effect ? [Schwinger, PR 82(1951) 664]

- Pair creations of electron and positron in an external electric field.

Pair creations of particle and anti-particles in an external field are said ubiquitously.

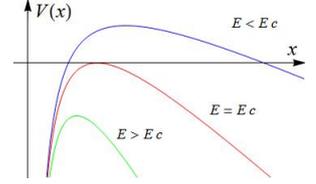
- Non-perturbative phenomenon

The Schwinger effect is a tunneling process with the potential barrier $V(x) = 2m - eEx - \frac{\alpha_s}{x}$.

- Critical electric field

The potential barrier decreases gradually as E becomes large. When $E=E_c$, the potential barrier vanishes.

➡ E_c is a critical electric field. However, this is not confirmed in QED.



★Pair creation in N=4 SYM and Potential Analysis [YS, K. Yoshida, JHEP 08 (2013) 002 [arXiv:1304.7917]]

- To argue the Schwinger effect, a U(1) gauge field should be introduced. It is necessary to spontaneously break SU(N+1) to SU(N) x U(1).

- Quark : SU(N) fundamental rep., U(1) charge: g_{YM}

The potential barrier vanishes when $E \simeq 0.70 \frac{2\pi m^2}{\sqrt{\lambda}}$ with $\alpha_s = 4\pi^2 \sqrt{\lambda} / \Gamma(1/4)^4$

The critical electric field is not the same as the one of DBI result. $\left[E_{DBI} = \frac{2\pi m^2}{\sqrt{\lambda}}, S_{DBI} = -T_{D3} \frac{r_0^4}{L^4} \int d^4x \sqrt{1 - \frac{(2\pi\alpha')^2 L^4 E^2}{r_0^4}} \right]$

We revisit this issue.

The probe is put near the horizon!

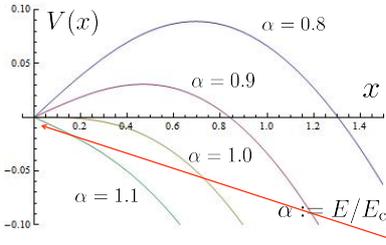
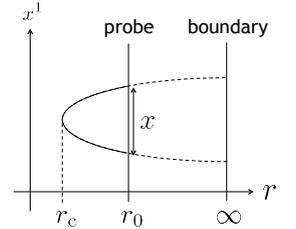
cf. [Semenoff-Zaremba, PRL 107 (2011) 171601]

The metric of AdS₅ x S⁵:

$$ds^2 = \frac{r^2}{L^2} \eta_{\mu\nu} dx^\mu dx^\nu + \frac{L^2}{r^2} dr^2 + L^2 d\Omega_5^2$$

The distance between quark and anti quark:

$$x = \frac{2L^2}{r_0 a} \int_1^{1/a} \frac{dy}{y^2 \sqrt{y^4 - 1}} \quad \left(a := \frac{r_c}{r_0} \right)$$



The Coulomb potential + the static energy:

$$V_{CP+SE} = 2T_F r_0 \int_1^{1/a} dy \frac{ay^2}{\sqrt{y^4 - 1}}$$

The potential barrier vanishes when $\alpha = 1.0$!!

Also analytically shown.

Finite!!



The theory is not an usual gauge theory but a non-linear QED.

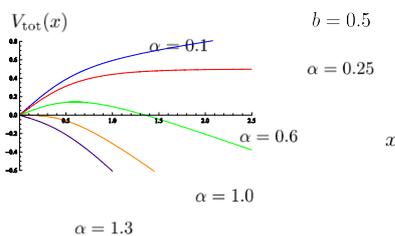
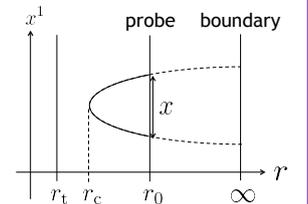
★Potential Analysis in confining background [YS, K. Yoshida, arXiv:1306.5512]

The metric (AdS-Soliton): [G. T. Horowitz, R. C. Myers, PRD 59 (1998) 026005]

$$ds^2 = \frac{r^2}{L^2} (\eta_{\mu\nu} dx^\mu dx^\nu + f(r)(dx^3)^2) + \frac{L^2}{r^2} f(r)^{-1} dr^2 + L^2 d\Omega_5^2, \quad f(r) = 1 - \frac{r_+^4}{r^4}$$

The distance between quark and anti quark: $x = \frac{2L^2}{r_0 a} \int_1^{1/a} \frac{dy}{\sqrt{(y^4 - 1)(y^4 - (b/a)^4)}} \quad \left(a := \frac{r_c}{r_0}, b := \frac{r_+}{r_0} \right)$

The potential energy + the static energy: $V_{PE+SE} = 2T_F r_0 \int_1^{1/a} dy \frac{ay^4}{\sqrt{(y^4 - 1)(y^4 - (b/a)^4)}}$



Two critical electric fields.

- (1) The orange line ($\alpha = 1.0$):

The potential barrier vanishes. This behaviour is the same as the previous case.

- (2) The red line ($\alpha = b^2$):

The potential barrier becomes flat around $x \rightarrow \infty$.

The electric field is balanced with the confining string tension $T_F \left(\frac{r_0}{L} \right)^2 b^2$.