

Extremal Surfaces in Asymptotic AdS Black Hole and Holographic Entanglement Entropy

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Based on T. Eguchi, and CH, [arXiv:15xx.xxxxx [hep-th]] [work in progress](#)

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Introduction:

A review of $\text{AdS}_{d+2}/\text{CFT}_{d+1}$ correspondence, entanglement entropy, and black hole entropy

AdS_{d+2} (bulk space-time)

$\xleftarrow{\text{holographic dual}}$ **CFT_{d+1}** (boundary)

- holographic entanglement entropy:

[Ryu-Takayanagi '06]

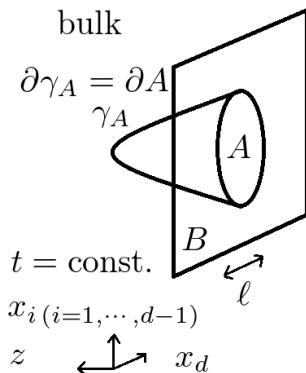
$$S_{A_{\text{HEE}}} = \frac{\text{Area}(\gamma_A)}{4G_N},$$

γ_A : minimal surface in bulk space-time

cf. black hole entropy: $S_{\text{BH}} = \frac{A}{4G_N}$,

e.g. AdS_{d+2} boundary topology : $\mathbb{R}^{1,d}$

bulk



- entanglement entropy:

$$S_{A_{\text{EE}}} = -\text{Tr}_{\mathcal{H}_A} \rho_A \log \rho_A,$$

$$= - \lim_{n \rightarrow 1} \frac{\partial}{\partial n} \text{Tr}_{\mathcal{H}_A} (\rho_A)^n,$$

(replica trick)

e.g. $n = 3$ replica trick in CFT₂

$$t_E = \infty$$

A trapezoid with vertices labeled u and v at the bottom. The top edge is horizontal and the bottom edge is slanted downwards to the right.

conformal mapping

$$\phi^{(2)}$$

\mathbb{R}^2

$$\phi^{(3)}$$

$$\mathbb{S}^1 \times \mathbb{R}^1$$

$$\mathbb{R}^1 \times \mathbb{S}^1$$

T²

Results of holographic entanglement entropy as extremal (and/or minimal) surface:

-- In the case of BTZ black hole

- rotating BTZ black hole metric:

$$\begin{aligned} ds_{\text{BTZ black hole}}^2 &= -f(r)dt^2 + f(r)^{-1}dr^2 + r^2 \left(d\phi - \frac{J}{2r^2}dt \right)^2, \\ &= -\frac{(r^2 - r_+^2)(r^2 - r_-^2)}{R_{\text{AdS}}^2 r^2} dt^2 + \frac{r^2 R_{\text{AdS}}^2}{(r^2 - r_+^2)(r^2 - r_-^2)} dr^2 + r^2 \left(d\phi + \frac{r_+ r_-}{R_{\text{AdS}} r^2} dt \right)^2, \\ f(r) &= \frac{r^2}{R_{\text{AdS}}^2} - M + \left(\frac{J}{2r} \right)^2, \\ M &= \frac{r_+^2 + r_-^2}{R_{\text{AdS}}^2} > 0, \quad J = \frac{2r_+ r_-}{R_{\text{AdS}}}, \quad M \geq \frac{J}{R_{\text{AdS}}}, \end{aligned}$$

-- non-rotating (t=const. and J=0) holographic entanglement entropy:

[Hubeny-Rangamani-Takayanagi '07]

$$S_A = \frac{\text{Area}(\gamma_A)}{4G_N^{(3)}} = \frac{c}{3} \log \left(\frac{\beta}{\pi\epsilon} \sinh \left| \frac{2\pi\Delta\phi_A}{\beta} \right| \right),$$

$$c = \frac{2R_{\text{AdS}}}{3G_N^{(3)}}, \quad \beta = \frac{2\pi}{\sqrt{M}},$$

Results of holographic entanglement entropy as extremal (and/or minimal) surface:

-- In the case of BTZ black hole

(detail)

-- non-rotating ($t=\text{const.}$ and $J=0$) holographic entanglement entropy:
 [Hubeny-Rangamani-Takayanagi '07]

space-like geodesic equation:

$$\frac{dr}{d\phi} = \pm \frac{1}{R_{\text{AdS}}} r \sqrt{r^2 - r_+^2} \sqrt{\frac{r^2}{r^{*2}} - 1},$$

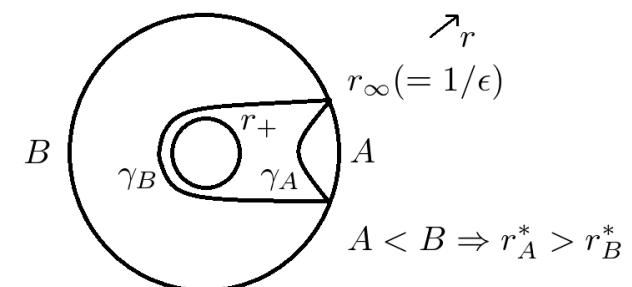
$$2\Delta\phi_A = \pm R_{\text{AdS}} \int_{r_\infty}^{r_A^*} \frac{dr}{r \sqrt{r^2 - r_+^2} \sqrt{\frac{r^2}{r_A^{*2}} - 1}},$$

$$2\Delta\phi_B = \pm R_{\text{AdS}} \int_{r_\infty}^{r_B^*} \frac{dr}{r \sqrt{r^2 - r_+^2} \sqrt{\frac{r^2}{r_B^{*2}} - 1}},$$

$$S_A = \frac{\text{Area}(\gamma_A)}{4G_N^{(3)}} = \frac{c}{3} \log \left(\frac{\beta}{\pi\epsilon} \sinh \left| \frac{2\pi\Delta\phi_A}{\beta} \right| \right),$$

$$S_B = \frac{\text{Area}(\gamma_B)}{4G_N^{(3)}} = \frac{c}{3} \log \left(\frac{\beta}{\pi\epsilon} \sinh \left| \frac{2\pi\Delta\phi_B}{\beta} \right| \right),$$

e.g. minimal surfaces in non-rotating BTZ black hole



$$\Delta\phi_A + \Delta\phi_B = \pi, \gamma_A \neq \gamma_B \Rightarrow S_{A_{\text{HEE}}} \neq S_{B_{\text{HEE}}}$$

$$(2+1)\text{dim.} \Rightarrow S_{\text{BH}} \neq |S_{B_{\text{HEE}}} - S_{A_{\text{HEE}}}|$$

Results of holographic entanglement entropy as extremal (and/or minimal) surface:

-- In the case of Schwarzschild-AdS₄ black hole

- Schwarzschild-AdS₄ black hole metric (AdS boundary topology R^{1,d}):

$$ds_{\text{Schwarzschild-AdS}_4 \text{ BH}}^2 = R_{\text{AdS}}^2 \frac{-f(z)dt^2 + \frac{1}{f(z)}dz^2 + (dx^2 + dy^2)}{z^2},$$

$$f(z) = 1 - \left(\frac{z}{z_H}\right)^3,$$

cf. AdS₄ space-time in Poincare AdS₄ coordinate:

$$ds_{\text{Poincare AdS}_4}^2 = R_{\text{AdS}}^2 \frac{-dt^2 + dz^2 + (dx^2 + dy^2)}{z^2},$$

-- static (t=const. and translational invariant in y direction) holographic entanglement entropy:

$$S_{A_{\text{HEE}}} = \frac{\text{Area}(\gamma_A)}{4G_N} = \frac{R_{\text{AdS}}^2 L}{4G_N} \int_{\epsilon_{\text{UV}}}^{z^*} \frac{dz}{z^2 \sqrt{1 - \left(\frac{z}{z_H}\right)^3} \sqrt{-1 + \left(\frac{z^*}{z}\right)^4}} \left(\frac{z^*}{z}\right)^2 \sim \mathcal{O}(\epsilon_{\text{UV}}^{-1}),$$

Area low

Results of holographic entanglement entropy as extremal (and/or minimal) surface:

-- In the case of Schwarzschild- AdS_4 black hole

(detail)

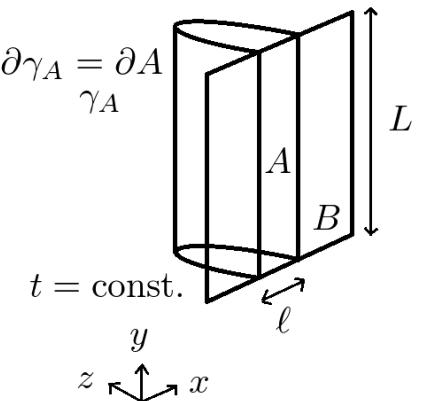
-- static ($t=\text{const.}$ and translational invariant in y direction) **holographic entanglement entropy**:

$$\text{Area}(\gamma_A) = R_{\text{AdS}}^2 L \int_{-\frac{\ell}{2}}^{\frac{\ell}{2}} \frac{dx}{z^2} \sqrt{1 + \frac{1}{f(z)} \left(\frac{dz}{dx} \right)^2},$$

$$0 = \delta \text{Area}(\gamma_A) \quad \Rightarrow \quad \frac{d^2 z}{dx^2} - \frac{1}{2} \frac{f'}{f} \left(\frac{dz}{dx} \right)^2 + \left[f + \left(\frac{dz}{dx} \right)^2 \right] \frac{2}{z} = 0,$$

$$\frac{dz}{dx} = \sqrt{1 - \left(\frac{z}{z_H} \right)^3} \sqrt{-1 + \left(\frac{z^*}{z} \right)^4},$$

e.g. AdS_4 boundary topology : $\mathbb{R}^{1,2}$



$$S_{A_{\text{HEE}}} = \frac{\text{Area}(\gamma_A)}{4G_N} = \frac{R_{\text{AdS}}^2 L}{4G_N} \int_{\epsilon_{\text{UV}}}^{z^*} \frac{dz}{z^2 \sqrt{1 - \left(\frac{z}{z_H} \right)^3} \sqrt{-1 + \left(\frac{z^*}{z} \right)^4}} \left(\frac{z^*}{z} \right)^2 \sim \mathcal{O}(\epsilon_{\text{UV}}^{-1}),$$

Area low

Discussions (form the view point of entanglement entropy and black hole entropy)

v.s. black hole entropy

- Static BTZ black hole:

$$S_{\text{BH}} \neq |S_{B_{\text{HEE}}} - S_{A_{\text{HEE}}}|,$$

This results is coursed by (2+1) dimensional space-time specially.

- Time-dependent AdS_4 Black Hole:
(under investigation)

$$S_{\text{BH}} = |S_{B_{\text{HEE}}} - S_{A_{\text{HEE}}}|,$$

In higher than (2+1) dimensional space-time, above results is expected.

v.s. entanglement entropy

Because there is Hawking radiation from black hole in bulk space-time, total system density operator on boundary interpret as mixed state:

$$\rho_{\text{tot}} \neq |\Psi\rangle\langle\Psi|,$$

Then, generally:

$$S_{A_{\text{EE}}} \neq S_{B_{\text{EE}}},$$

This property was confirmed in holographic entanglement entropy:

$$S_{A_{\text{HEE}}} \neq S_{B_{\text{HEE}}},$$

cf. pure state : $\rho_{\text{tot}} = |\Psi\rangle\langle\Psi|, \Rightarrow S_{A_{\text{EE}}} = S_{B_{\text{EE}}}.$

Summary and future directions

Summary:

- We calculated **holographic entanglement entropy**, which is defined by minimal surface in bulk space-time, in the case of **BTZ black hole** and **Schwarzschild-AdS₄ black hole**, and **reconsidered** these results **from the viewpoint of entanglement entropy and black hole entropy**.
- Holographic entanglement entropy is important physical quantity in a study for AdS_{d+2}/CFT_{d+1} correspondence.

Future directions:

- To manifest degree of freedom on boundary in conformal field theory which is holographic dual with asymptotic AdS black hole in bulk space-time.
- To study in the case of other asymptotic AdS black holes which have not only mass but also charge and/or **angular momentum**.
- **To expand to covariant holographic entanglement entropy.**

References

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