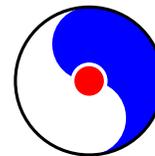


# Muon $g-2$ in Standard Model

Taku Izubuchi  
(RBC&UKQCD collaboration)

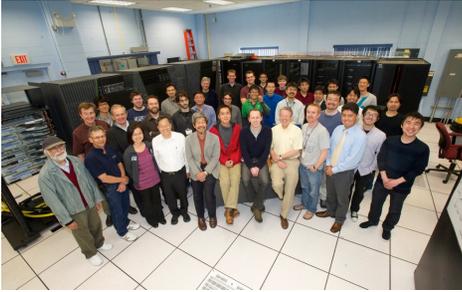


**RIKEN BNL**  
Research Center

2023-08-08 YITP String and Field , Kyoto

# References

- g-2 Hadronic Vacuum Polarization (HVP)  
RBC/UKQCD 2023 <https://arxiv.org/pdf/2301.08696.pdf>  
Phys. Rev. Lett. 121 (2018) 022003
- g-2 Hadronic Light-by-Light (HLbL)  
RBC/UKQCD 2023 <https://arxiv.org/pdf/2304.04423.pdf>  
Phys. Rev. Lett 124 (2020) 13, 132002  
Phys. Rev. D96 (2017) 034515  
Phys. Rev. Lett. 118 (2017) 022005
- talks at 5<sup>th</sup> Plenary Workshop Muon g-2 Theory Initiative  
<https://indico.ph.ed.ac.uk/event/112/timetable/#20220905>  
Tom Blum, Christoph Lehner, Mattia Bruno,  
BMWc (Laurent Lellouch) FNAL/HPQCD (Ethan Neil)  
ETMc (Giuseppe Gagliardi) chiQCD (Gen Wang)
- talks at Radiative correction and MC tools for low-energy hadronic cross sections in e+e- collisions,  
<https://indico.psi.ch/event/13708/timetable/#20230607.detailed>  
Martin Hoferichter , Fedor Ignatov
- talks at Lattice 2023 <https://indico.fnal.gov/event/57249/>  
A. Keshavarzi, S. Kuberski, C. Lehner,



# Collaborators / Machines

g-2 DWF  
HVP & HLbL

- Group 1: Regensburg, Christoph Lehner
- Group 2: UKQCD: Vera Gulpers et al.
- Group 3: Millan: Mattia Bruno + Davide Giusti (Regensburg)
- Group 4: Connecticut: Tom Blum
- Group 5: BNL: Taku Izubuchi  
Peter Boyle , Chulwoo Jung, Chris Kelly,  
Aaron Meyer, Nobuyuki Matsumoto

DWFQCD  
Global fit

- Group 1 Yong-Chull Jung, Norman Christ, Bob Mawhinney (CU), Chris Kelly (BNL)
- Group2 Christoph Lehner

tau input for  
g-2 HVP &  
HVP GEVP

Mattia Bruno (Milano)  
Aaron Meyer (BNL)

Christoph Lehner (Regensburg)  
Taku Izubuchi (BNL & RBRC)

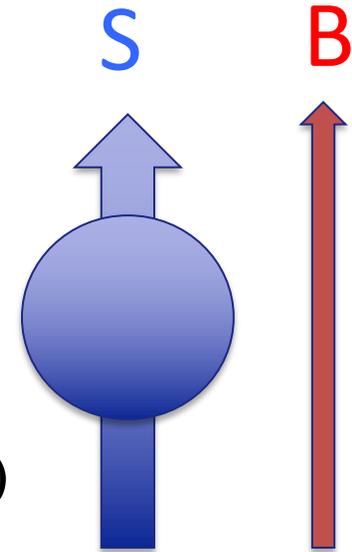
Part of related calculation are done by resources from  
USQCD (DOE), HPCI(Japan), RIKEN, XSEDE, ANL BG/Q Mira (DOE, ALCC), Edinburgh BG/Q,  
Crasher (DOE)  
BNL BG/Q, RIKEN BG/Q and Cluster (RICC, HOKUSAI)  
Support from RIKEN, JSPS, US DOE, and BNL

# Anomalous magnetic moment

- Fermion's energy in the external magnetic field:

$$V(x) = -\vec{\mu}_l \cdot \vec{B} \quad \vec{\mu}_l = g_l \frac{e}{2m_l} \vec{S}_l$$

- Magnetic moment Lande g-factor tree level value **2**
- 1928 P.A.M. Dirac “Quantum Theory of Electron”  
Dirac equation (relativity, minimal gauge interaction)

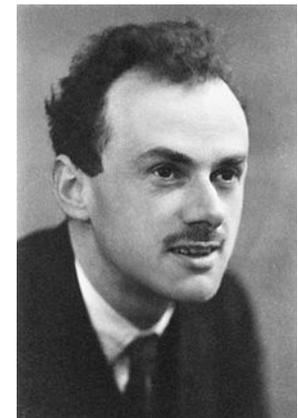


$$i[\partial_\mu - ieA_\mu(x)]\gamma^\mu\psi(x) = m\psi(x)$$

- Non-relativistic and weak constant magnetic field limits of the Dirac equation :

$$-i\hbar\frac{\partial\psi}{\partial t} = \left[ \frac{\nabla^2}{2m} + \frac{e}{2m} \left( \vec{L} + 2\vec{S} \right) \cdot \vec{B} \right] \psi$$

$$g_l = 2 \quad (\text{for Dirac Fermion } l = e, \mu, \tau, \dots)$$

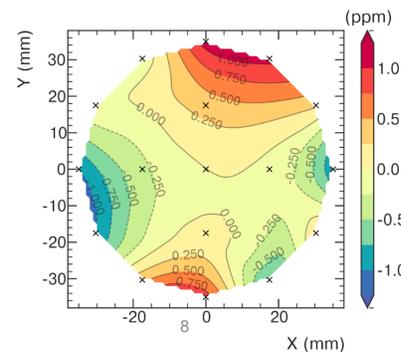
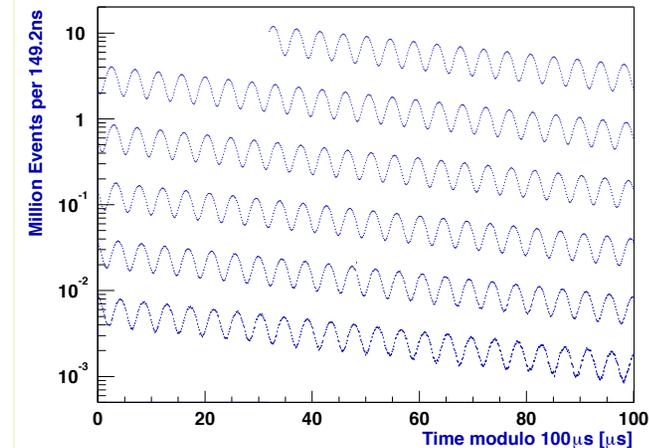
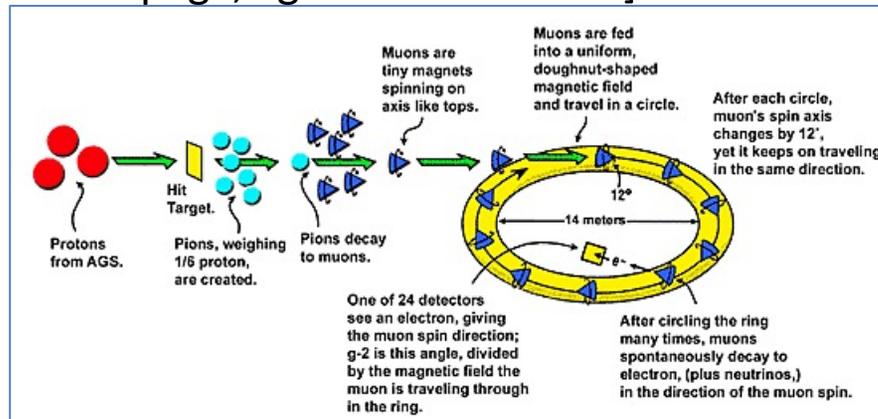


# The Muon g-2 experiments BNL E821 (-2004)

- measure precession of muon spin very accurately

$$N(t) = N_0(E) \exp\left(-t/\gamma\tau_\mu\right) [1 + A(E) \sin(\omega_a t + \phi(E))]$$

[ BNL web page, g-2 collaboration ]

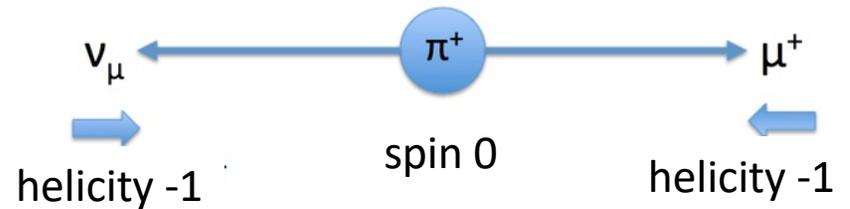


$$a_\mu \propto \frac{\omega_a}{\langle \omega'_p \rangle \times M_\mu}$$

[ Aoyama, LAT15] [ A.Keshavarzi, LAT23]

# Recipe of a g-2 measurement

1. Prepare a polarized muon beam from P-violating pion decay

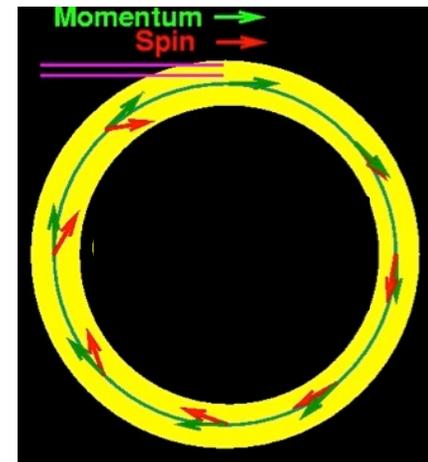


1. Store in a magnetic field (let muon spin precessed)

$$\vec{\omega} = -\frac{e}{m} \left[ a_\mu \vec{B} - \left( a_\mu - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{\beta} \times \vec{E}}{c} + \frac{\eta}{2} \left( \vec{\beta} \times \vec{B} + \frac{\vec{E}}{c} \right) \right]$$

Magic momentum,  $\gamma=30$  ( $p=3$  GeV/c),

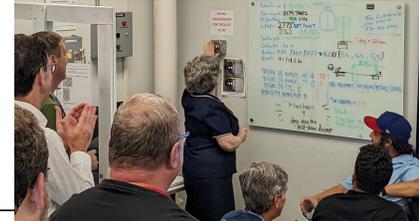
2. Measure positron from P-violating muon decay



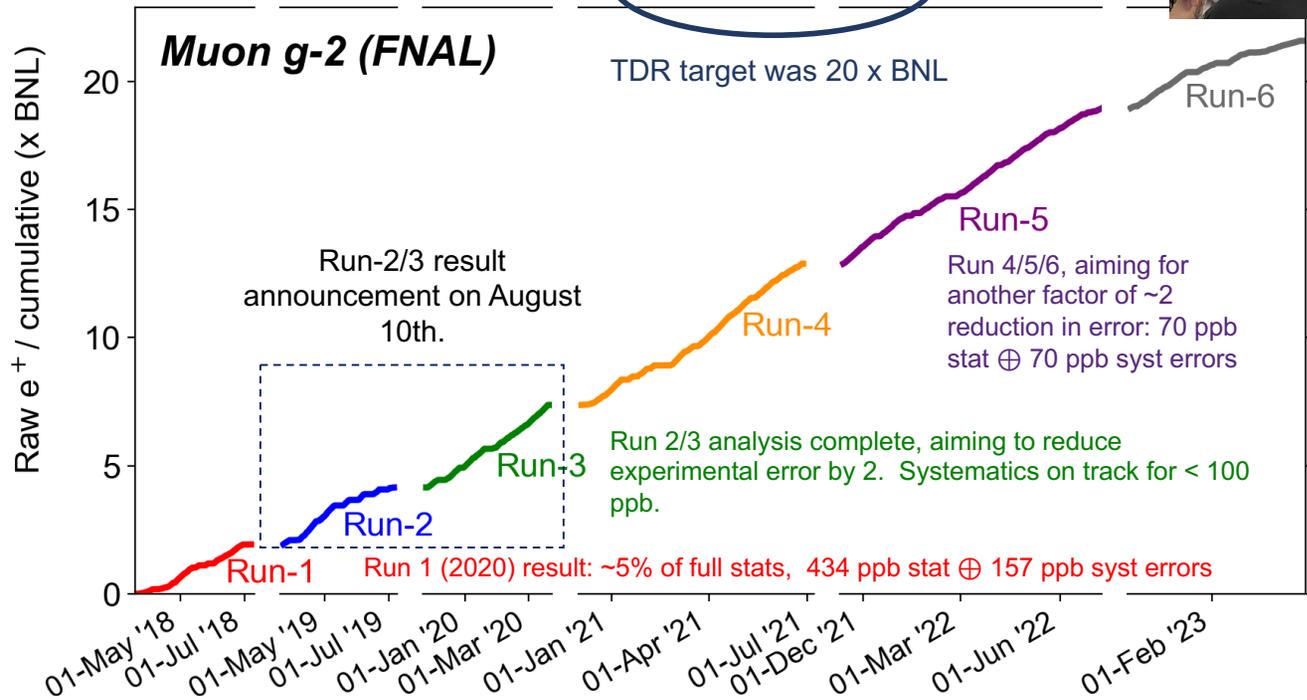
[ Slide from T. Mibe, L. Roberts ]

# Various corrections, error budget [A. Keshavarzi, LAT23]

## The full data-set



Last update: 2023-07-10 10:26; Total = 21.90 (xBNL)

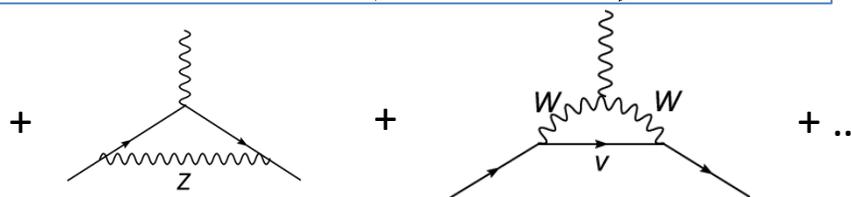
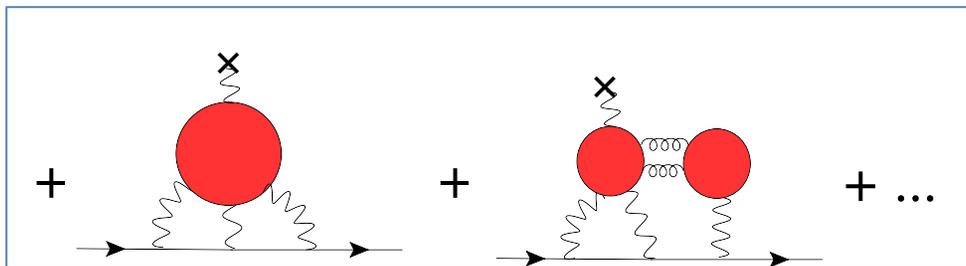
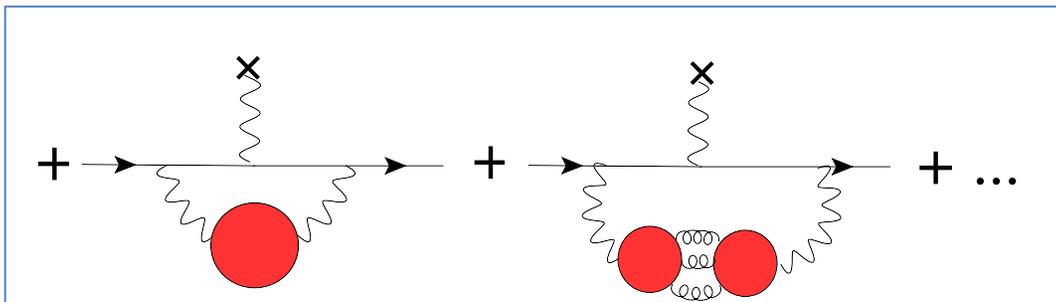
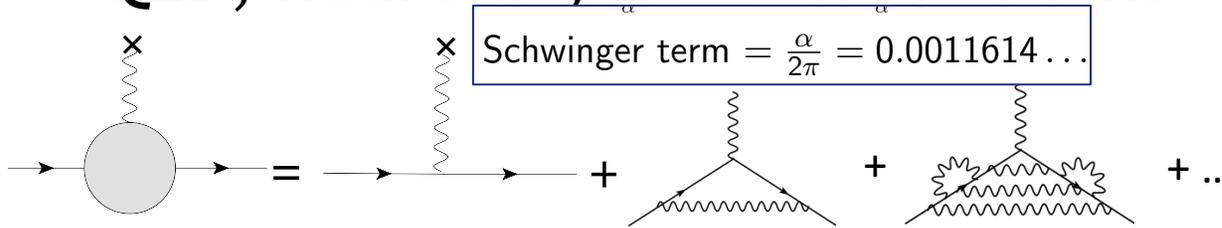


# SM Theory

$$\gamma^\mu \rightarrow \Gamma^\mu(q) = \left( \gamma^\mu F_1(q^2) + \frac{i \sigma^{\mu\nu} q_\nu}{2m} F_2(q^2) \right)$$



## ■ QED, hadronic, EW contributions



QED (5-loop)

Aoyama Hayakawa,

Kinoshita, Nio

PRL109,111808 (2012)

Hadronic vacuum  
polarization (HVP)

Hadronic light-by-light  
(HLbL)

Electroweak (EW)

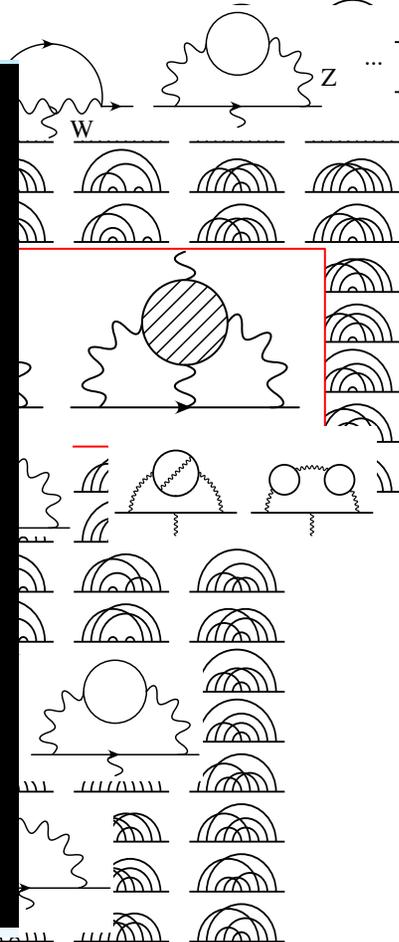
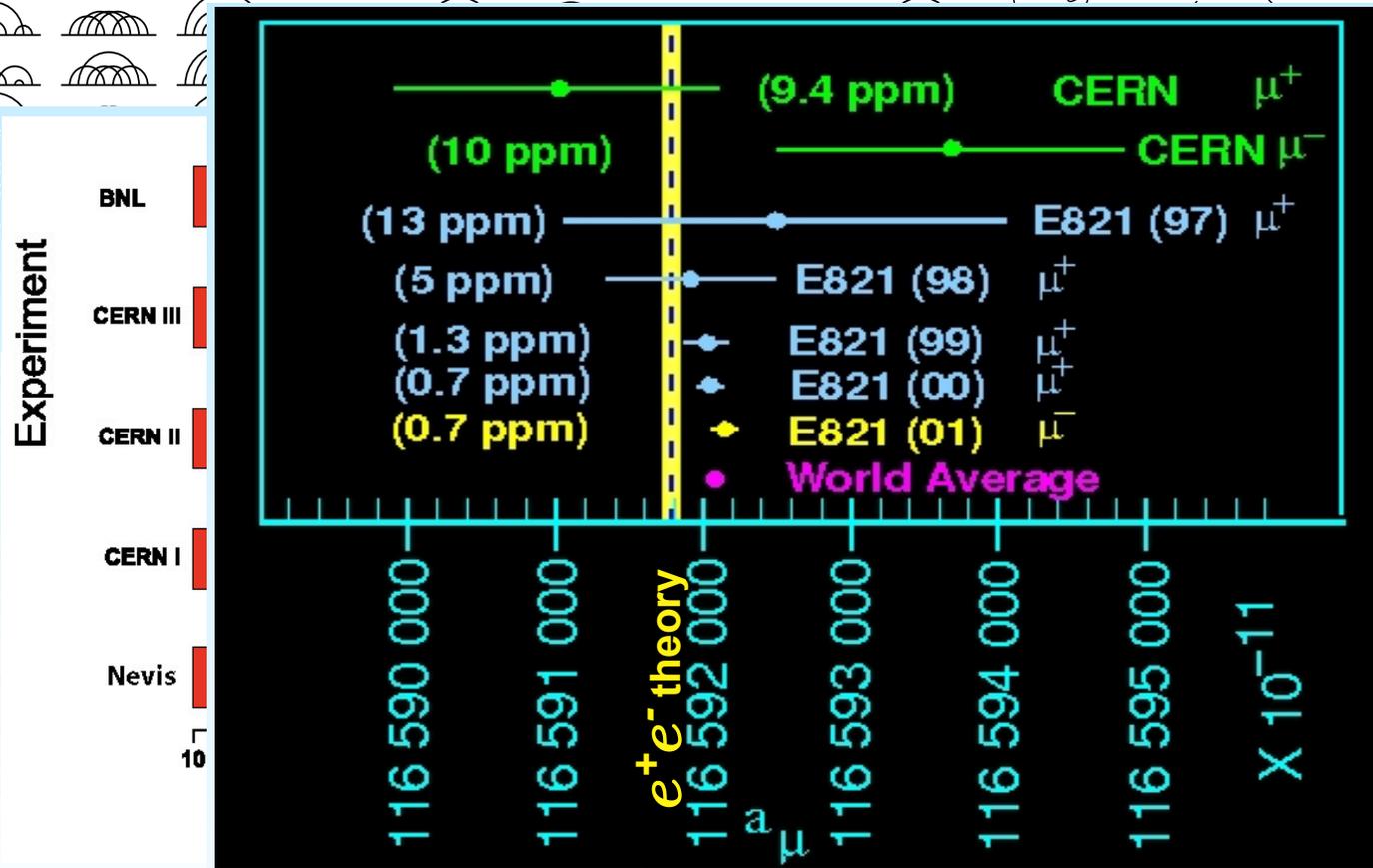
Knecht et al 02

Czarnecki et al. 02

.....

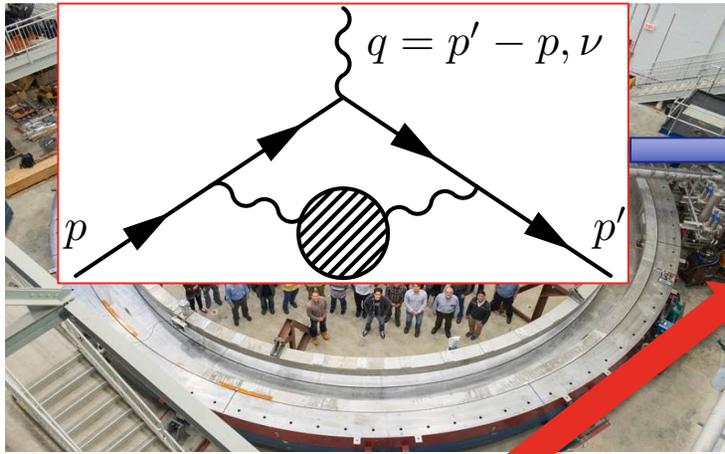
# Experiments history

Lee Roberts



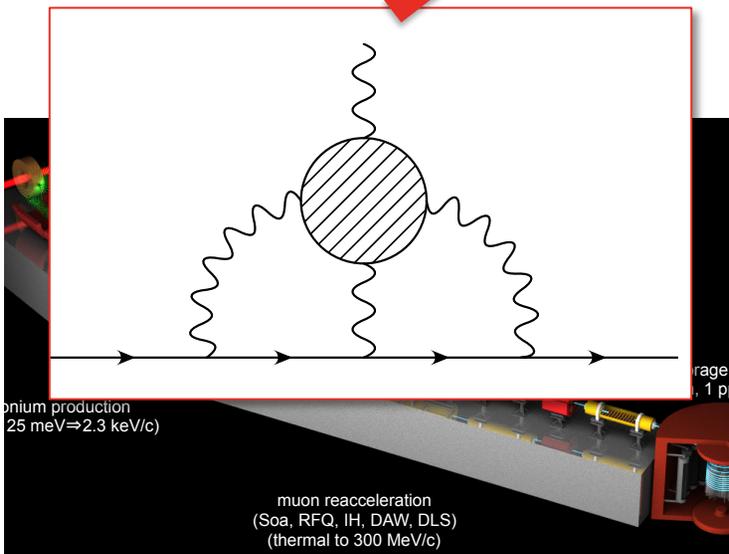
# muon anomalous magnetic moment

BNL g-2 till 2004 :  $\sim 3.7 \sigma$  larger than SM prediction



Contribution	Value $\times 10^{10}$	Uncertainty $\times 10^{10}$
QED (5 loops)	11 658 471.895	0.008
EW	15.4	0.1
<b>HVP LO</b>	692.3	<b>4.2</b>
HVP NLO	-9.84	0.06
HVP NNLO	1.24	0.01
<b>Hadronic light-by-light</b>	10.5	<b>2.6</b>
Total SM prediction	11 659 181.5	4.9
BNL E821 result	11 659 209.1	6.3
FNAL E989/J-PARC E34 goal		$\approx 1.6$

$$a_{\mu}^{\text{EXP}} - a_{\mu}^{\text{SM}} = 27.4 \underbrace{(2.7)}_{\text{HVP}} \underbrace{(2.6)}_{\text{HLbL}} \underbrace{(0.1)}_{\text{other}} \underbrace{(6.3)}_{\text{EXP}} \times 10^{-10}$$



FNAL E989 (2017-)

2021-04: announces  $\sim$ BNL level error,  $4.6 \times 10^{-10}$

**2023-08-10**: Run-2, Run-3 results

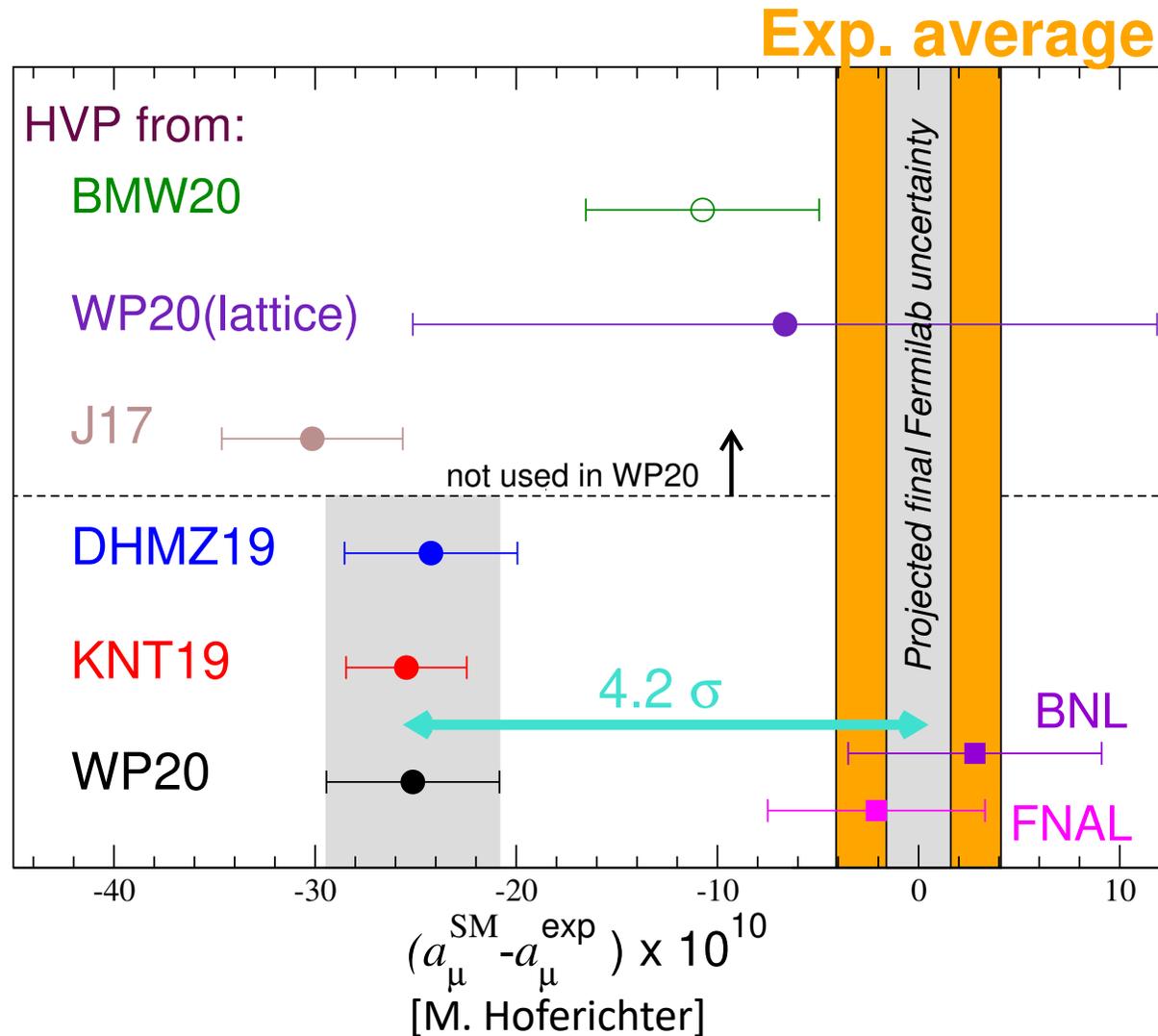
All 6 Run completed, x22 more statistics than BNL aiming for error  $1.6 \times 10^{-10}$  0.14ppm

J-PARC E34 (IMPORTANT different systematics !)

ultra-cold muon beam

0.37 ppm then 0.1 ppm, also EDM

# April 2020 status muon g-2 HVP



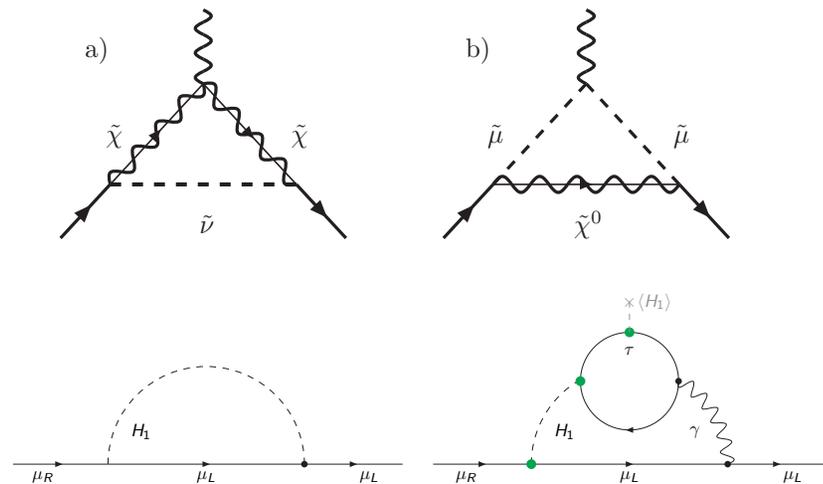
# G-2 from BSM sources

- Typical new particle contribute g-2  
 $g-2 \sim C (m_\mu / m_{NP})^2$
- To explain current discrepancy

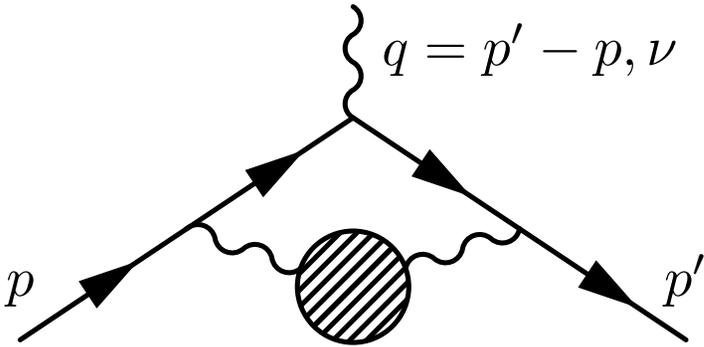
$C$	1	$\frac{\alpha}{\pi}$	$(\frac{\alpha}{\pi})^2$
$M_{NP}$	$2.0^{+0.4}_{-0.3}$ TeV	$100^{+21}_{-13}$ GeV	$5^{+1}_{-1}$ GeV

[A. Nyfler ]

- SUSY (scalar-lepton )
- 2 Higgs doublet models  
Type-X, ....
- Dark photons  
from kinematical mixings  
 $\varepsilon F_{\mu\nu} F'_{\mu\nu}$

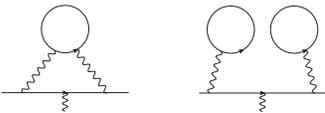


# Hadronic Vacuum Polarization (HVP) contribution to $g-2$

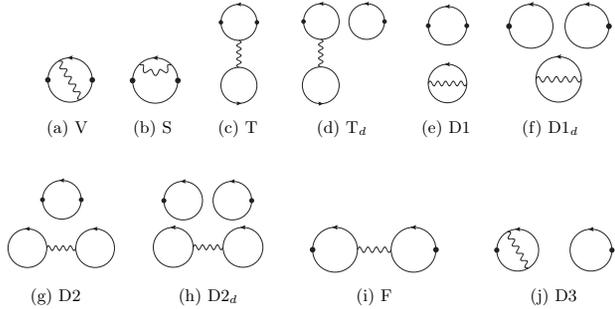


Quark & anti-quark contribution

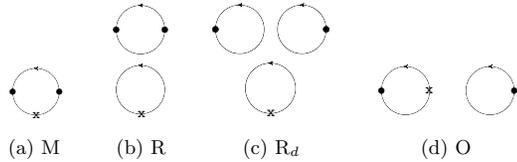
Isospin limit



QED corrections



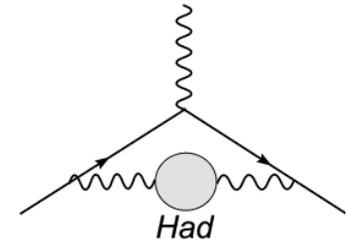
Strong isospin breaking



# Leading order of hadronic contribution (HVP)

- Hadronic vacuum polarization (HVP)

$$V_\mu \quad \text{[diagram: photon with hadronic loop]} \quad V_\nu = (q^2 g_{\mu\nu} - q_\mu q_\nu) \Pi_V(q^2)$$



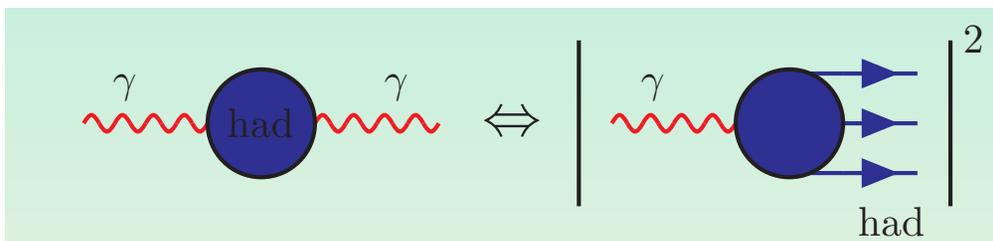
quark's EM current :  $V_\mu = \sum_f Q_f \bar{f} \gamma_\mu f$

- Unitarity, Optical Theorem

$$\text{Im}\Pi_V(s) = \frac{s}{4\pi\alpha} \sigma_{\text{tot}}(e^+ e^- \rightarrow X)$$

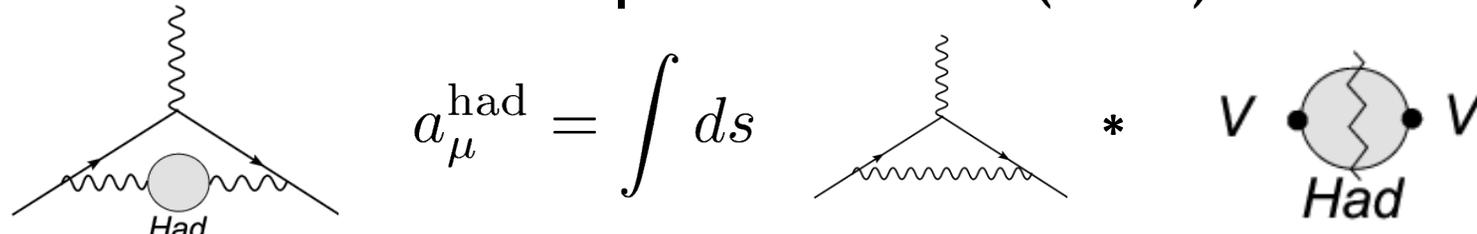
- Analyticity

$$\Pi_V(s) - \Pi_V(0) = \frac{k^2}{\pi} \int_{4m_\pi^2}^{\infty} ds \frac{\text{Im}\Pi_V(s)}{s(s - k^2 - i\epsilon)}$$



# Leading order of hadronic contribution (HVP)

## ■ Hadronic vacuum polarization (HVP)

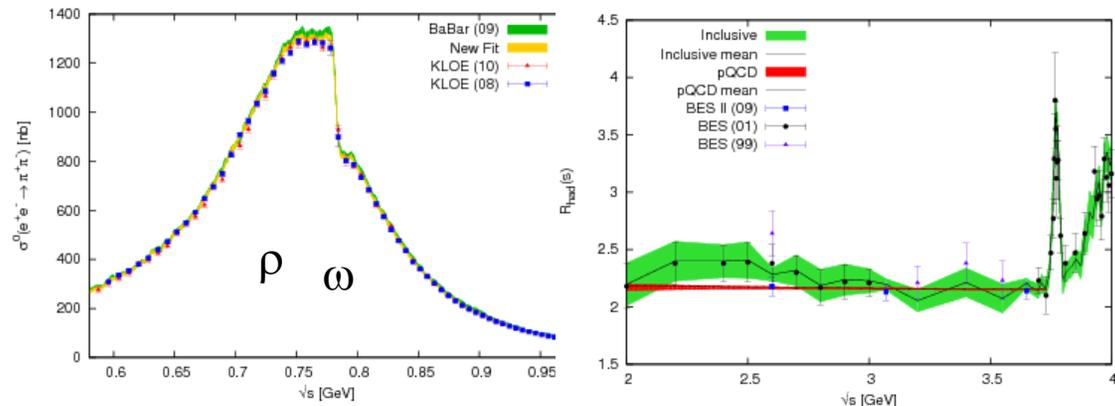


$$a_{\mu}^{\text{had}} = \int ds$$

$$= \frac{\alpha}{\pi^2} \int_{m_{\pi}^2}^{\infty} \frac{ds}{s} \text{Im}\Pi(s) K(s) \quad K(s) = \int_0^1 dx \frac{x^2(1-x)}{x^2 + (s/m_{\mu}^2)(1-x)}$$

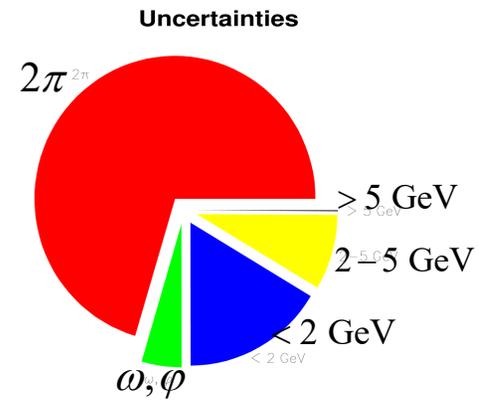
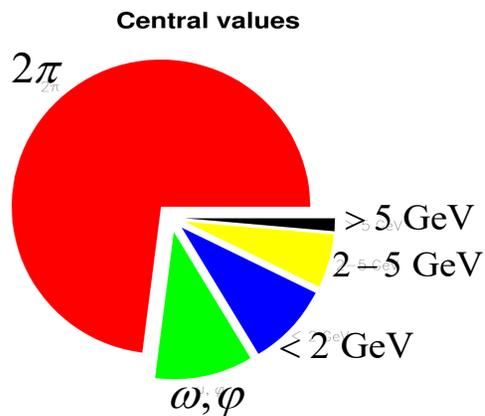
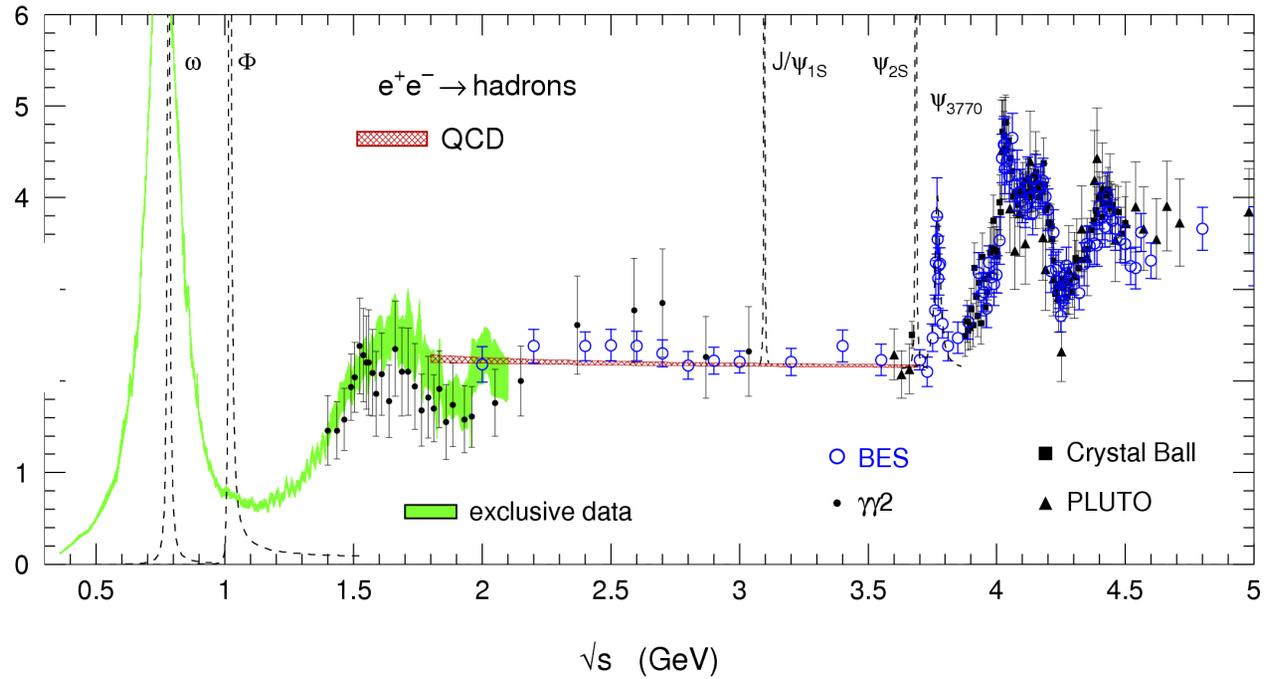
$$= \frac{1}{3} \left(\frac{\alpha}{\pi}\right)^2 \left[ \int_{m_{\pi}^2}^{s_{\text{cut}}} ds \frac{K(s)}{s} R_{\text{had}}^{\text{data}}(s) + \int_{s_{\text{cut}}}^{\infty} ds \frac{K(s)}{s} R_{\text{had}}^{\text{pQCD}}(s) \right]$$

$$R(s) = \sigma_{had}^0(s) / \left( \frac{4\pi\alpha^2}{3s} \right)$$



# g-2 from R-ratio experiments

$$R(s) = \sigma_{had}^0(s) / \left( \frac{4\pi\alpha^2}{3s} \right)$$





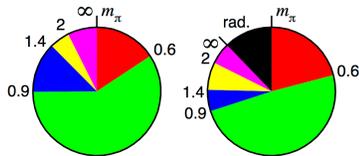
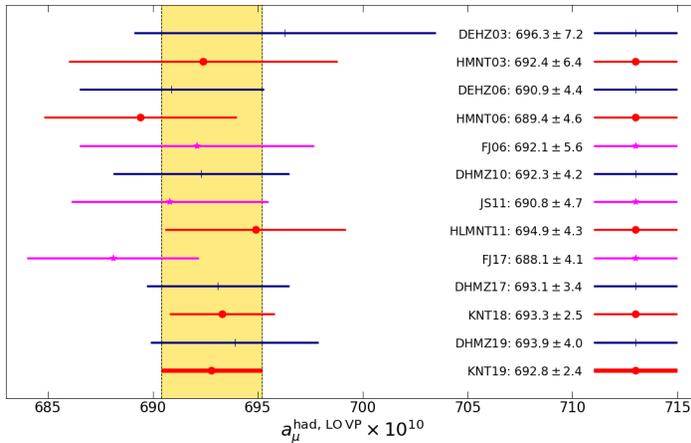
# Comparisons and the 2021 WP result

$$a_\mu^{\text{had, LOVP}} = 693.84 \pm 1.19_{\text{stat}} \pm 1.96_{\text{sys}} \pm 0.22_{\text{vp}} \pm 0.71_{\text{fsr}}$$

$$= 692.78 \pm 2.42_{\text{tot}}$$

KNT19, Phys.Rev.D 97 (2018) 114025, Phys.Rev.D 101 (2020) 014029.

Phys.Rept. 887 (2020) 1-166.



➤ Clear  $\pi^+\pi^-$  dominance

➤ Precision better than 0.4% (uncertainties include all available correlations and  $\chi^2$  inflation)

Detailed comparisons by-channel and energy range between direct integration results:

	DHMZ19	KNT19	Difference
$\pi^+\pi^-$	507.85(0.83)(3.23)(0.55)	504.23(1.90)	3.62
$\pi^+\pi^-\pi^0$	46.21(0.40)(1.10)(0.86)	46.63(94)	-0.42
$\pi^+\pi^-\pi^+\pi^-$	13.68(0.03)(0.27)(0.14)	13.99(19)	-0.31
$\pi^+\pi^-\pi^0\pi^0$	18.03(0.06)(0.48)(0.26)	18.15(74)	-0.12
$K^+K^-$	23.08(0.20)(0.33)(0.21)	23.00(22)	0.08
$K_S K_L$	12.82(0.06)(0.18)(0.15)	13.04(19)	-0.22
$\pi^0\gamma$	4.41(0.06)(0.04)(0.07)	4.58(10)	-0.17
Sum of the above	626.08(0.95)(3.48)(1.47)	623.62(2.27)	2.46
[1.8, 3.7] GeV (without $c\bar{c}$ )	33.45(71)	34.45(56)	-1.00
$J/\psi, \psi(2S)$	7.76(12)	7.84(19)	-0.08
[3.7, $\infty$ ] GeV	17.15(31)	16.95(19)	0.20
Total $a_\mu^{\text{HVP, LO}}$	694.0(1.0)(3.5)(1.6)(0.1) $_{\psi(0.7)}$ $_{\text{DV+QCD}}$	692.8(2.4)	1.2

+ evaluations using unitarity & analyticity constraints for  $\pi\pi$  and  $\pi\pi\pi$  channels [CHS 2018, HHKS 2019]

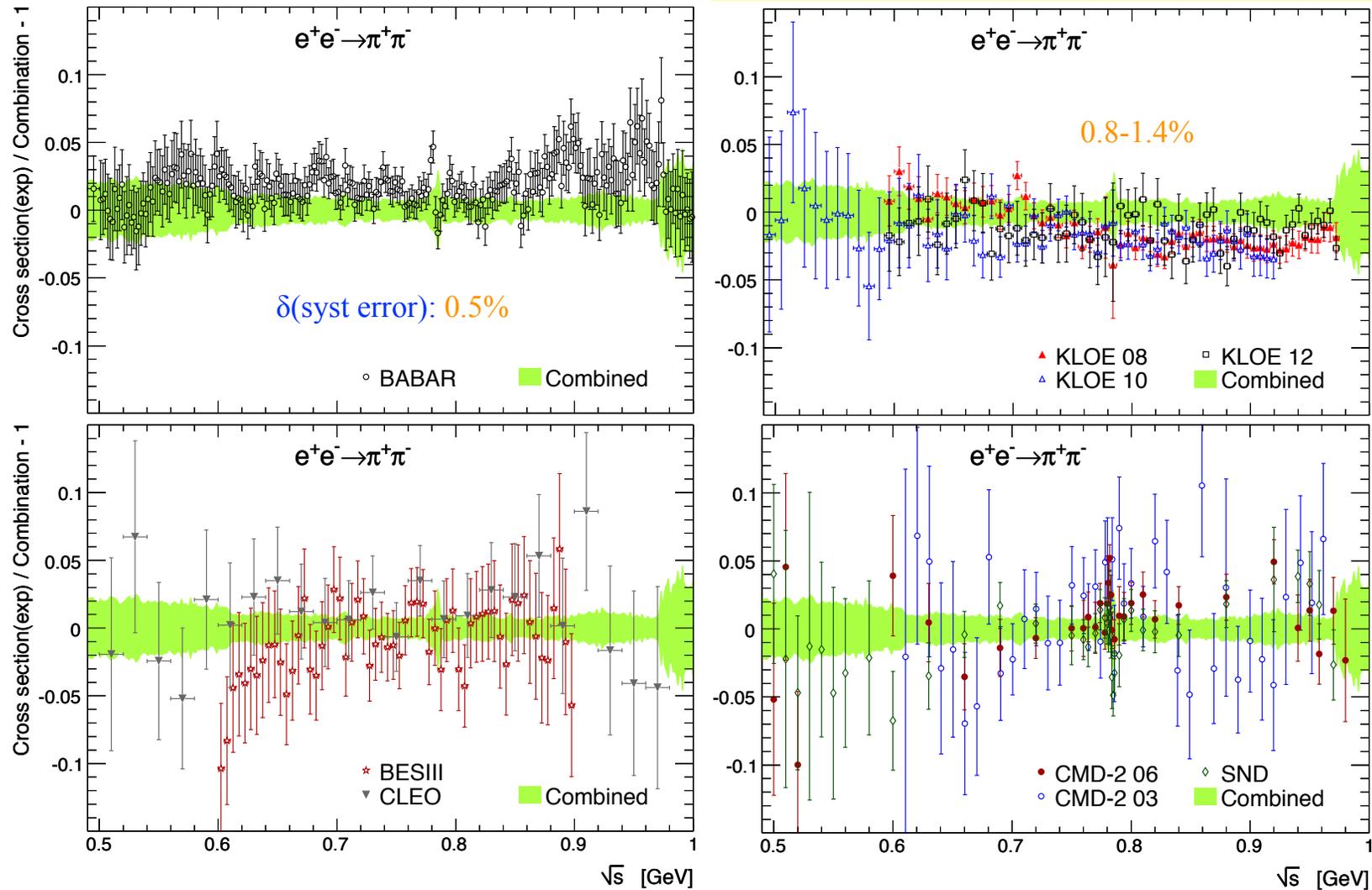
Conservative merging to obtain a realistic assessment of the underlying uncertainties:

- Account for differences in results from the same experimental inputs.
- Include correlations between systematic errors

$$\Rightarrow a_\mu^{\text{HVP, LO}} = 693.1 (4.0) \times 10^{-10}$$

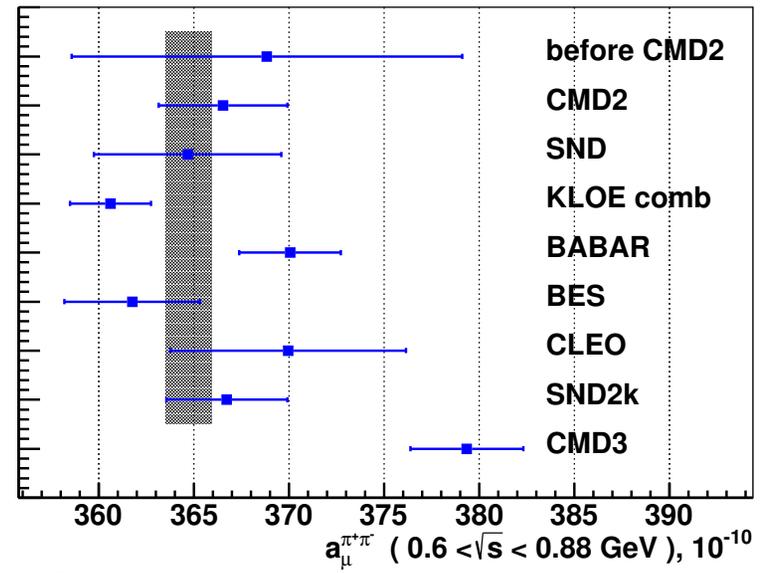
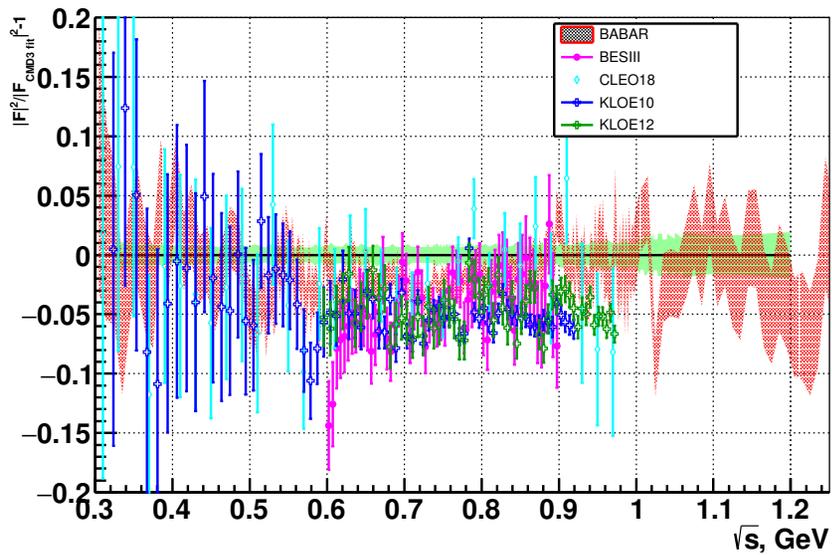
25

# The Dominant $\pi^+\pi^-$ Channel (2)



BABAR & KLOE dominates 0.6-0.9 GeV  $\pi\pi$  data,  
 Has a large discrepancy between BABAR & KLOE -> inflate error (dominant)

# 2023 CMD-3 $e^+ e^-$ to $\pi^+ \pi^-$



CMD-3, 2302.08834

[https://indico.psi.ch/event/13708/contributions/43296/attachments/25270/46331/pipiFinal\\_7June2023\\_ZurichRadcorMC.pdf](https://indico.psi.ch/event/13708/contributions/43296/attachments/25270/46331/pipiFinal_7June2023_ZurichRadcorMC.pdf)

# [ A. Keshavarzi, LAT23 ]



## CMD-3 compared to KNT19

In collaboration with Genessa Benton, Diogo Boito, Maarten Golterman, Kim Maltman & Santi Peris.

CMD-3 [F. Ignatov et al, arXiv:2302.08834]

To be able to compare CMD-3 with KNT19 data combination:

- Data published as pion form factor,  $|F_\pi|^2$ .
- Must subtract vacuum polarisation effects using Fedor Ignatov's VP correction update.
- Must include final-state-radiation effects.
- Put data on fine, common binning.

In the full  $2\pi$  data combination range, the KNT19 analysis found:

$$a_\mu^{\pi^+\pi^-}(0.305 \rightarrow 1.937 \text{ GeV}) = (503.46 \pm 1.91) \times 10^{-10}.$$

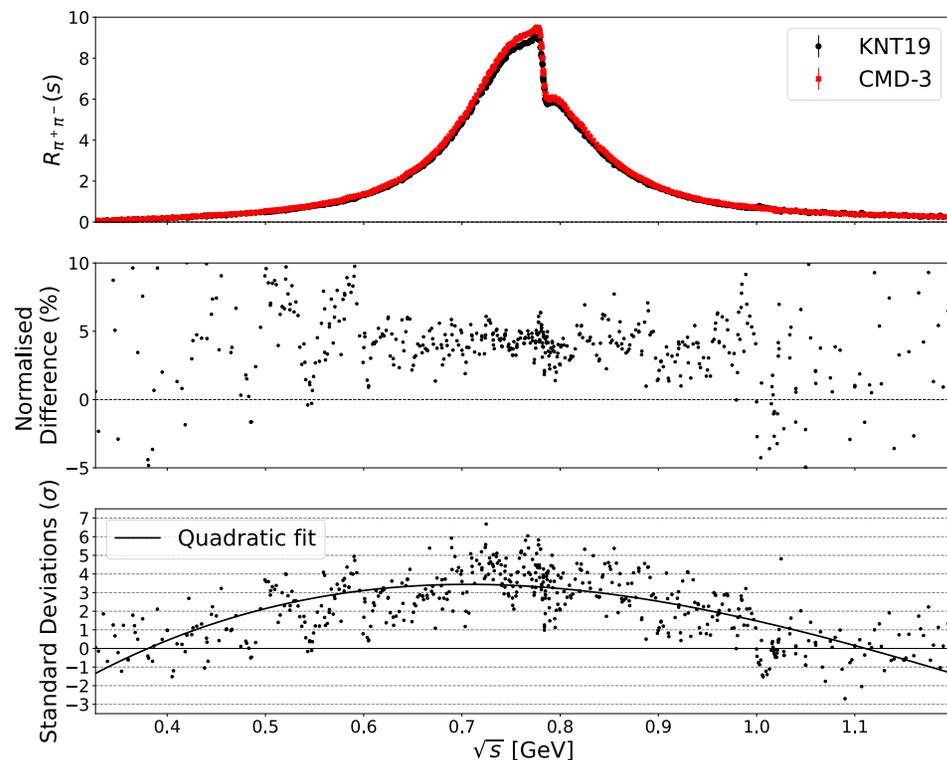
Replacing KNT19  $2\pi$  data in the region  $0.33 \rightarrow 1.20 \text{ GeV}$  with CMD-3 data:

$$a_\mu^{\pi^+\pi^-}(0.305 \rightarrow 1.937 \text{ GeV}) = (525.17 \pm 4.18) \times 10^{-10}.$$

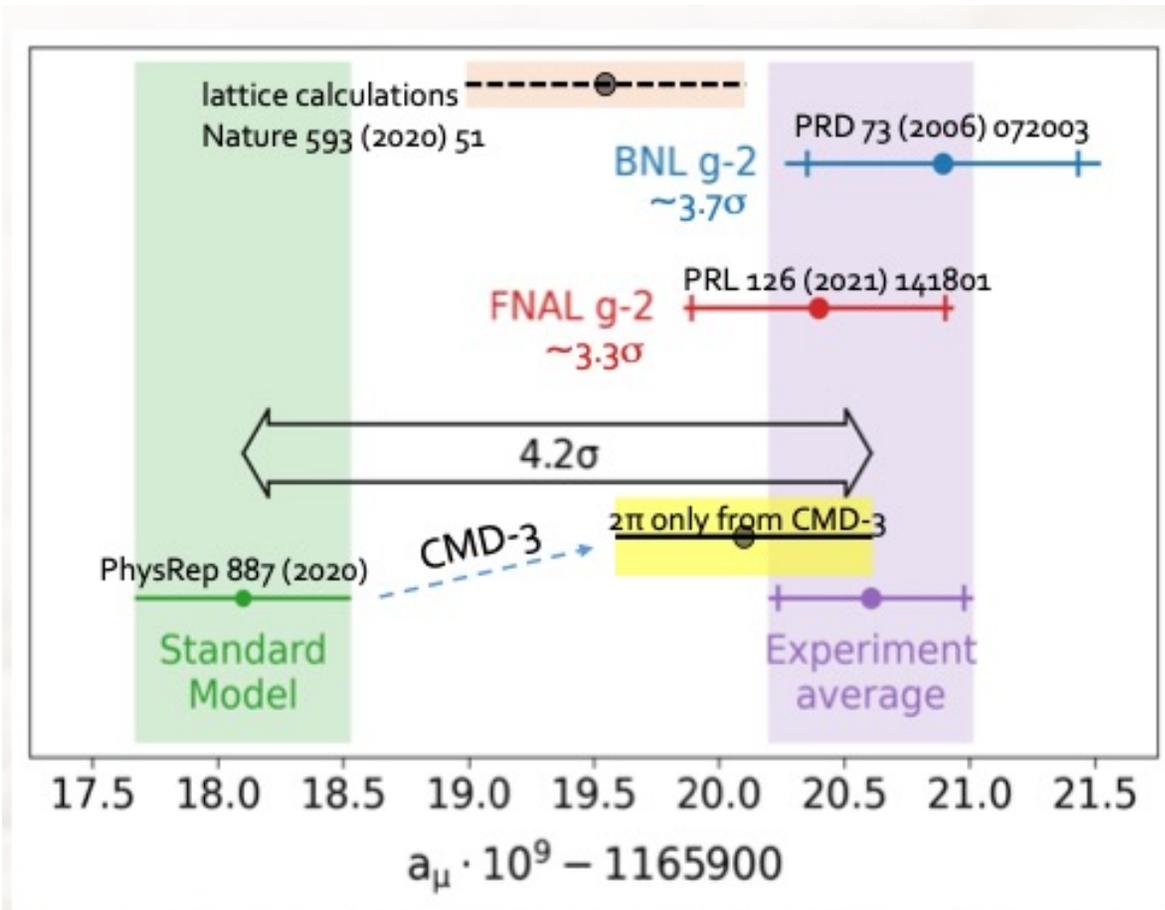
Neglecting possible correlations between e.g. CMD-3 and CMD-2, this results in a difference of:

$$\Delta a_\mu^{\pi^+\pi^-} = (21.71 \pm 4.96) \times 10^{-10} \rightarrow 4.4\sigma,$$

This removes the experiment vs. SM Muon g-2 discrepancy.



27



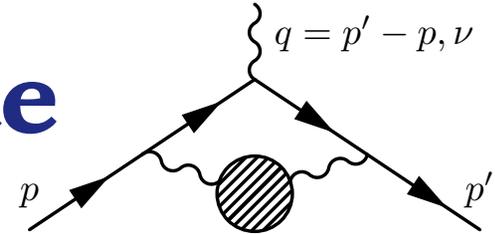
If it will be only CMD-3 than SM will be solved. But CMD-3 is only one now over many other experiments (BaBar, KLOE, BES, CMD-2, SND, ...)

**Unfortunately at the moment, we don't know the reasons of the disagreement between different experiments.**

[https://indico.psi.ch/event/13708/contributions/43296/attachments/25270/46331/pipiFinal\\_7June2023\\_ZurichRadcorMC.pdf](https://indico.psi.ch/event/13708/contributions/43296/attachments/25270/46331/pipiFinal_7June2023_ZurichRadcorMC.pdf)

# g-2 HVP from Lattice

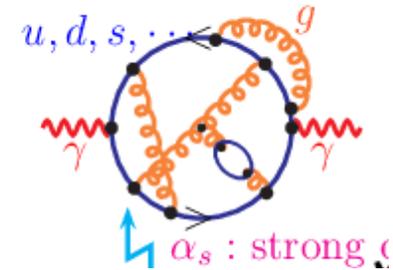
[ T. Blum, 2003 ]



[Bernecker Meyer 2011, Feng et al. 2013]

In Euclidean space-time, project vector 2 pt to zero spacial momentum,  $\vec{p} = 0$  :

$$C(t) = \frac{1}{3} \sum_{x,i} \langle j_i(x) j_i(0) \rangle$$



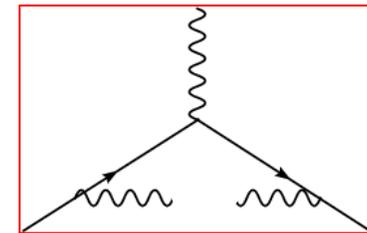
g-2 HVP contribution is

$$w(t) \sim t^4$$

$$a_\mu^{HVP} = \sum_t w(t) C(t)$$

$$w(t) = 2 \int_0^\infty \frac{d\omega}{\omega} f_{\text{QED}}(\omega^2) \left[ \frac{\cos \omega t - 1}{\omega^2} + \frac{t^2}{2} \right]$$

$f_{\text{QED}}(\omega^2)$

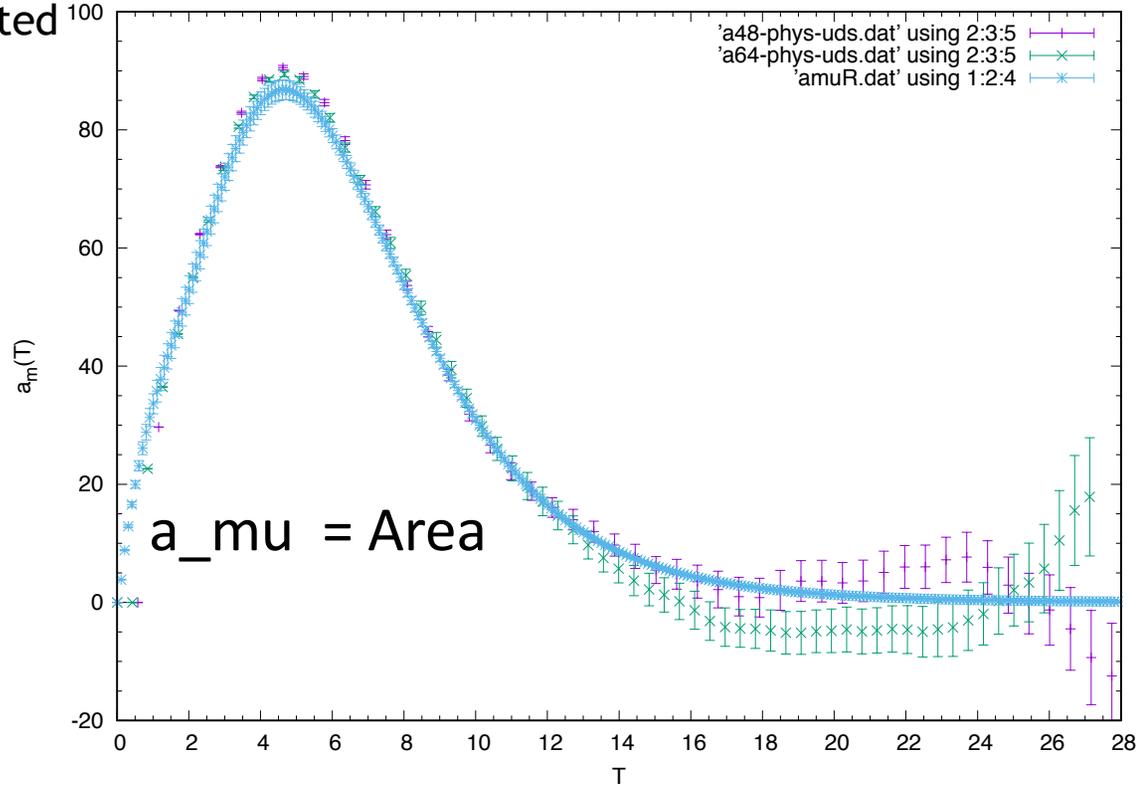


- Subtraction  $\Pi(0)$  is performed.  
Noise/Signal  $\sim e^{(E_{\pi\pi} - m_\pi)t}$ , is improved [Lehner et al. 2015] .

# Comparison of R-ratio and Lattice

## [ F. Jegerlehner alphaQED 2016 ]

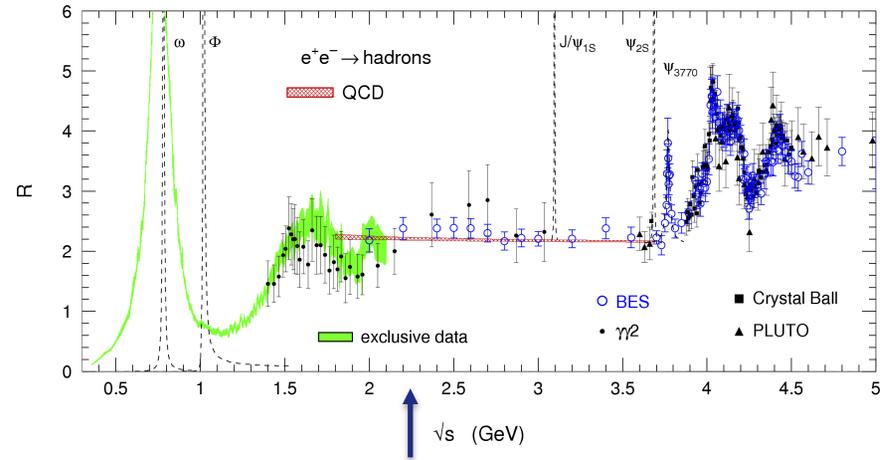
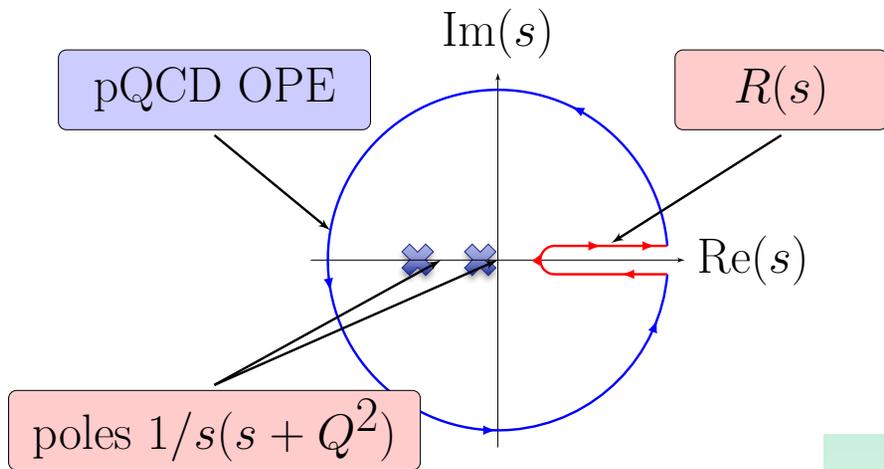
- Covariance matrix among energy bin in R-ratio is not available, assumes 100% correlated



$$a_{\mu}^{HVP} = \sum_t w(t) C(t)$$

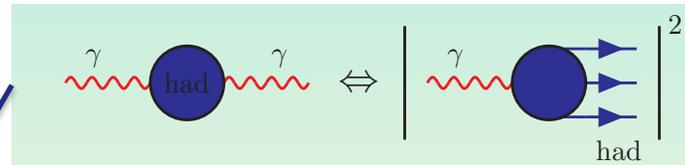
$$w(t) = 2 \int_0^{\infty} \frac{d\omega}{\omega} f_{\text{QED}}(\omega^2) \left[ \frac{\cos \omega t - 1}{\omega^2} + \frac{t^2}{2} \right]$$

$$C(t) = \frac{1}{3} \sum_{x,i} \langle j_i(x) j_i(0) \rangle$$



$$\hat{\Pi}(Q^2) = Q^2 \int_0^\infty ds \frac{R(s)}{s(s+Q^2)}$$

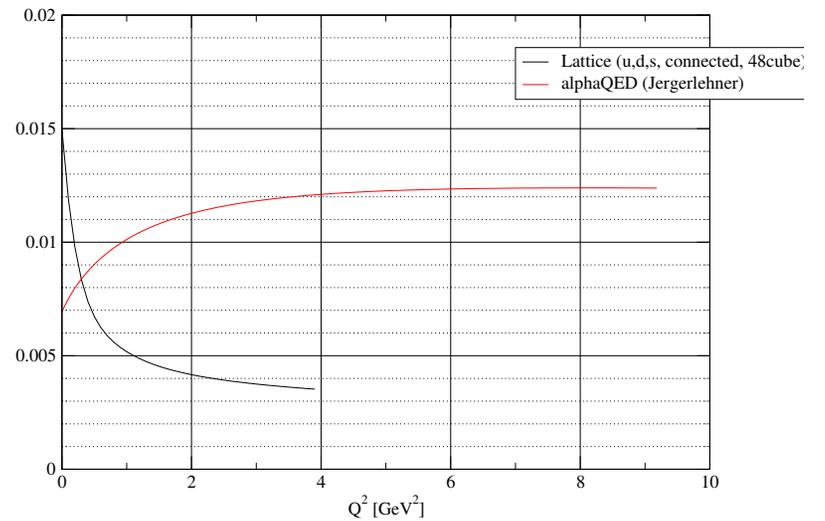
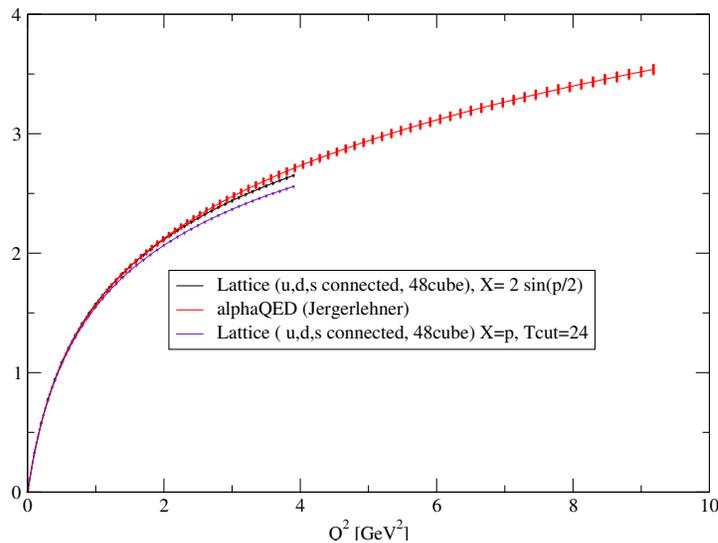
$(1/a = 1.78 \text{ GeV},$



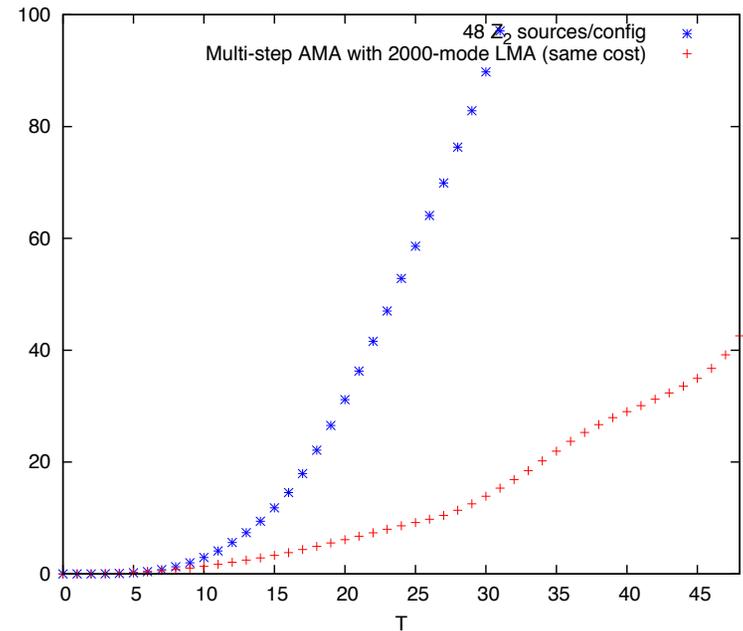
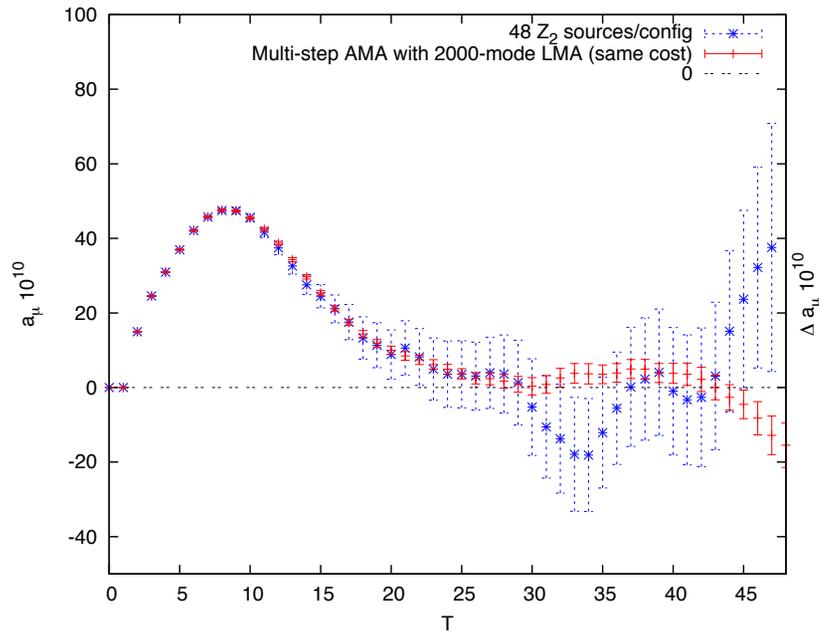
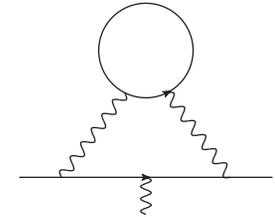
Relative statistical error)

$\hat{\Pi}_{\text{had}}(Q^2)$

Relative Err of  $\hat{\Pi}_{\text{had}}(Q^2)$



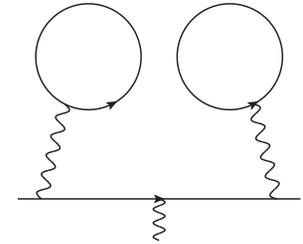
# up/down quark HVP



120 conf ( $a=0.11\text{fm}$ ), 80 conf ( $a=0.086\text{fm}$ ) physical point  $N_f=2+1$  Mobius DWF  
4D full volume LMA with 2,000 eigen vector (of e/o preconditioned zMobius  $D^+D$ )  
EV compression (1/10 memory) using local coherence [ C. Lehner Lat2017 Poster ]  
In addition, 50 sloppy / conf via multi-level AMA  
**more than x 1,000 speed up** compared to simple CG

# disconnected quark loop contribution

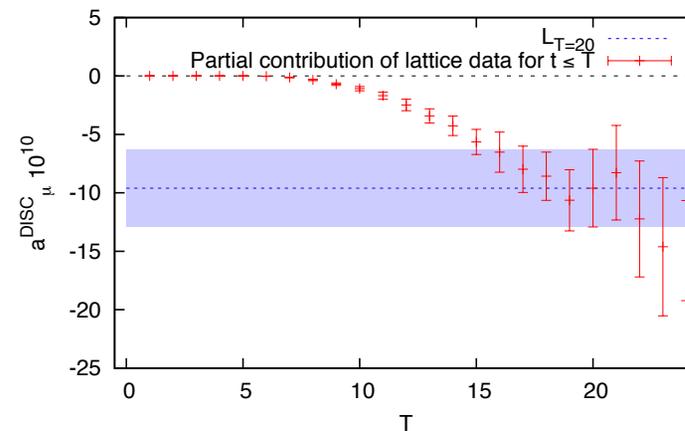
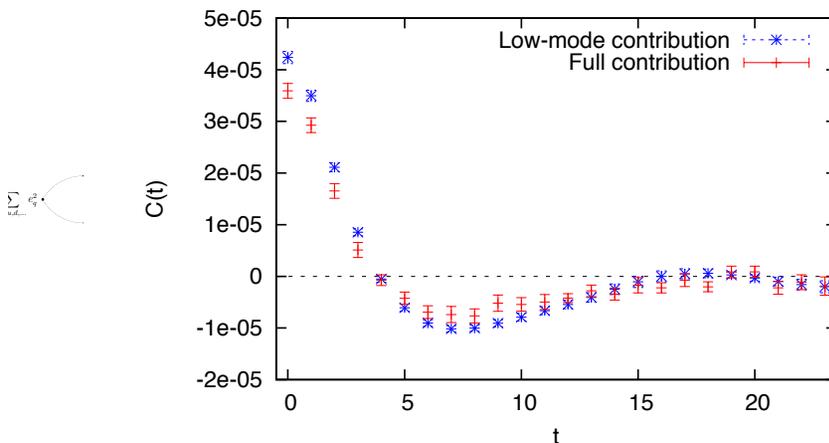
- [ C. Lehner et al. (RBC/UKQCD 2015, arXiv:1512.09054, PRL) ]
- Very challenging calculation due to statistical noise
- Small contribution, vanishes in SU(3) limit,  $Q_u + Q_d + Q_s = 0$
- Use low mode of quark propagator, treat it exactly ( all-to-all propagator with sparse random source )
- First non-zero signal



$$a_{\mu}^{\text{HVP (LO) DISC}} = -9.6(3.3)_{\text{stat}}(2.3)_{\text{sys}} \times 10^{-10}$$

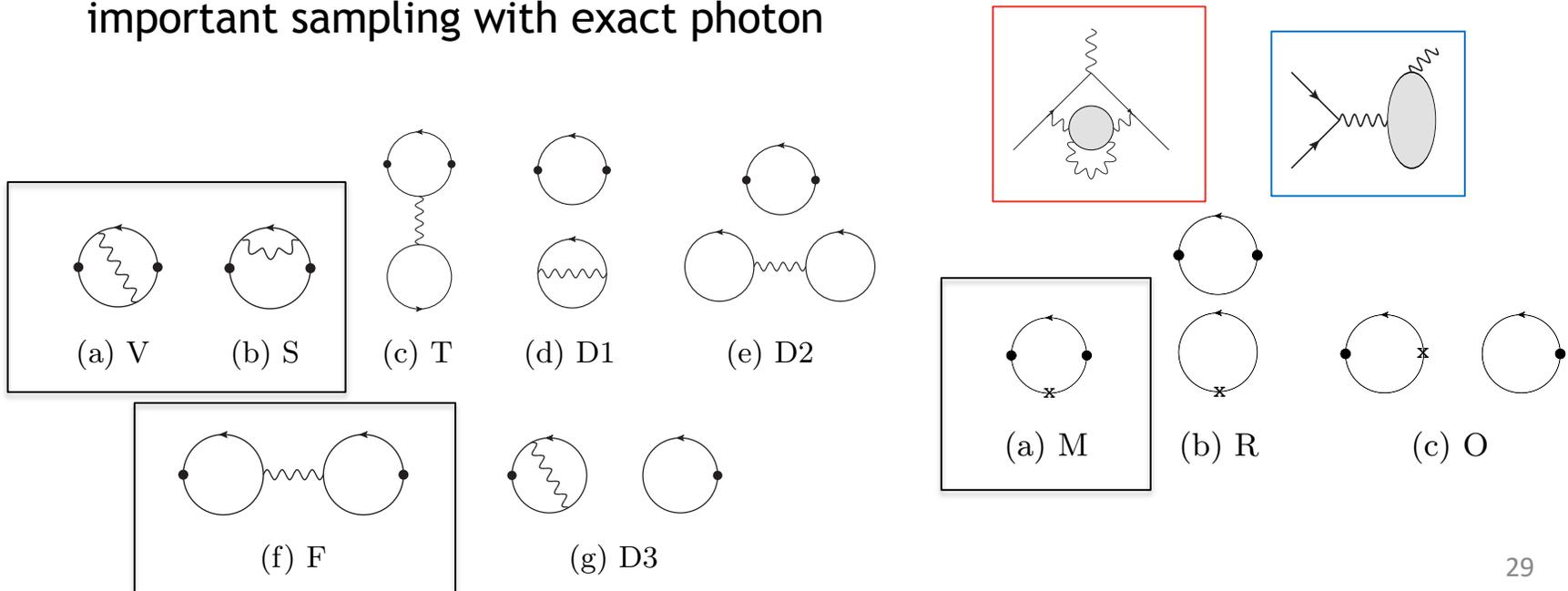
Sensitive to  $m_{\pi}$

crucial to compute at physical mass



# HVP QED+ strong IB corrections

- HVP is computed so far at Iso-symmetric quark mass, needs to compute **isospin breaking** corrections :  $Q_u, Q_d, m_u - m_d \neq 0$
- $u, d, s$  quark mass and lattice spacing are re-tuned using **{charge, neutral} x {pion, kaon}** and ( **Omega** baryon masses )
- For now,  $V, S, F, M$  are computed : assumes EM and IB of sea quark and also shift to lattice spacing is small (correction to disconnected diagram)
- Point-source method : stochastically sample pair of 2 EM vertices a la important sampling with exact photon



# Window values

## Combine R-ratio and Lattice

[ Christoph Lehner et al PRL18 ]

- Divide total  $a_\mu$  into { short, mid, long } distance contributions
- Useful to crosscheck among different lattice group and R-ratios

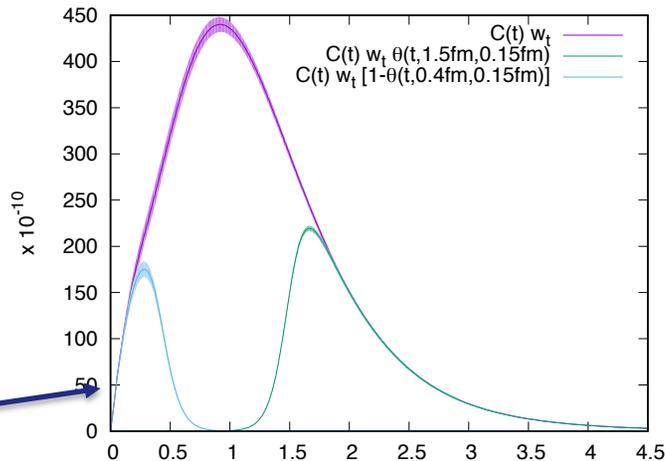
$$\Theta(t, \mu, \sigma) \equiv [1 + \tanh [(t - \mu)/\sigma]]$$

$$a_\mu = \sum_t w_t C(t) \equiv a_\mu^{\text{SD}} + a_\mu^{\text{W}} + a_\mu^{\text{LD}}$$

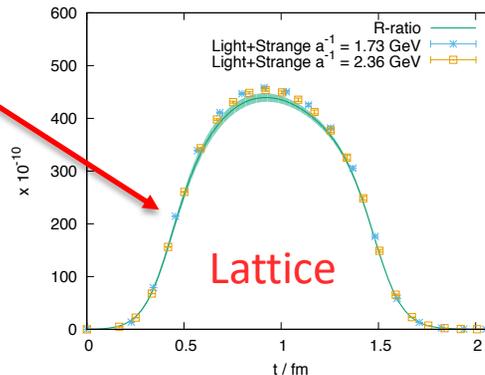
$$a_\mu^{\text{SD}} = \sum_t C(t) w_t [1 - \Theta(t, t_0, \Delta)],$$

$$a_\mu^{\text{W}} = \sum_t C(t) w_t [\Theta(t, t_0, \Delta) - \Theta(t, t_1, \Delta)],$$

$$a_\mu^{\text{LD}} = \sum_t C(t) w_t \Theta(t, t_1, \Delta)$$



+



## Euclidean time correlation from $e^+e^- R(s)$ data

From  $e^+e^- R(s)$  ratio, using dispersive relation, zero-spacial momentum projected Euclidean correlation function  $C(t)$  is obtained

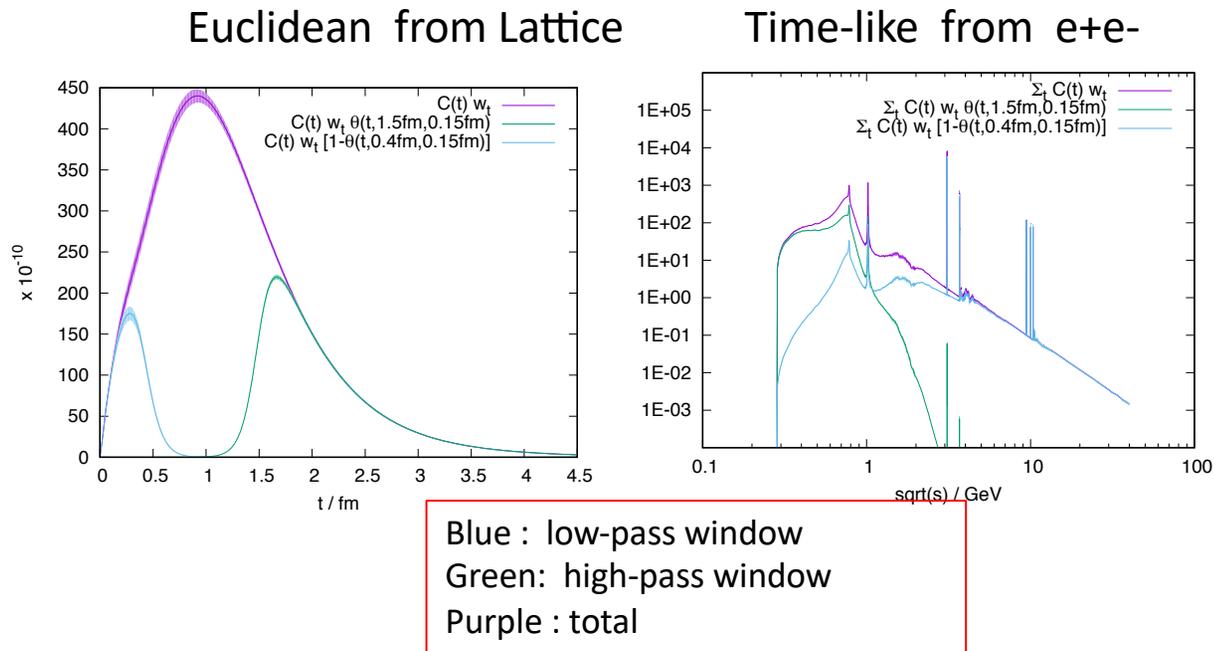
$$\hat{\Pi}(Q^2) = Q^2 \int_0^\infty ds \frac{R(s)}{s(s+Q^2)}$$

Lattice can compute Integral of Inclusive cross sections accurately

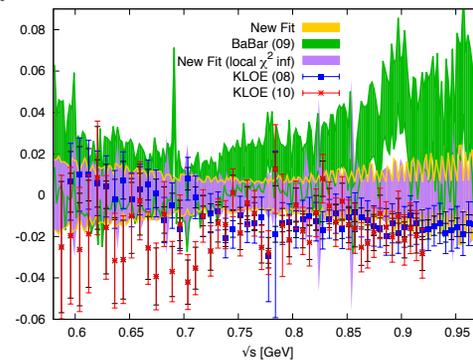
$$C^{\text{R-ratio}}(t) = \frac{1}{12\pi^2} \int_0^\infty \frac{d\omega}{2\pi} \hat{\Pi}(\omega^2) e^{i\omega t} = \frac{1}{12\pi^2} \int_0^\infty ds \sqrt{s} R(s) e^{-\sqrt{s}t}$$

- $C(t)$  or  $w(t)C(t)$  are directly comparable to Lattice results with the proper limits ( $m_q \rightarrow m_q^{\text{phys}}$ ,  $a \rightarrow 0$ ,  $V \rightarrow \infty$ , QED ...)
- Lattice: long distance has large statistical noise, (short distance: discretization error, removed by  $a \rightarrow 0$  and/or pQCD )
- R-ratio : short distance has larger error

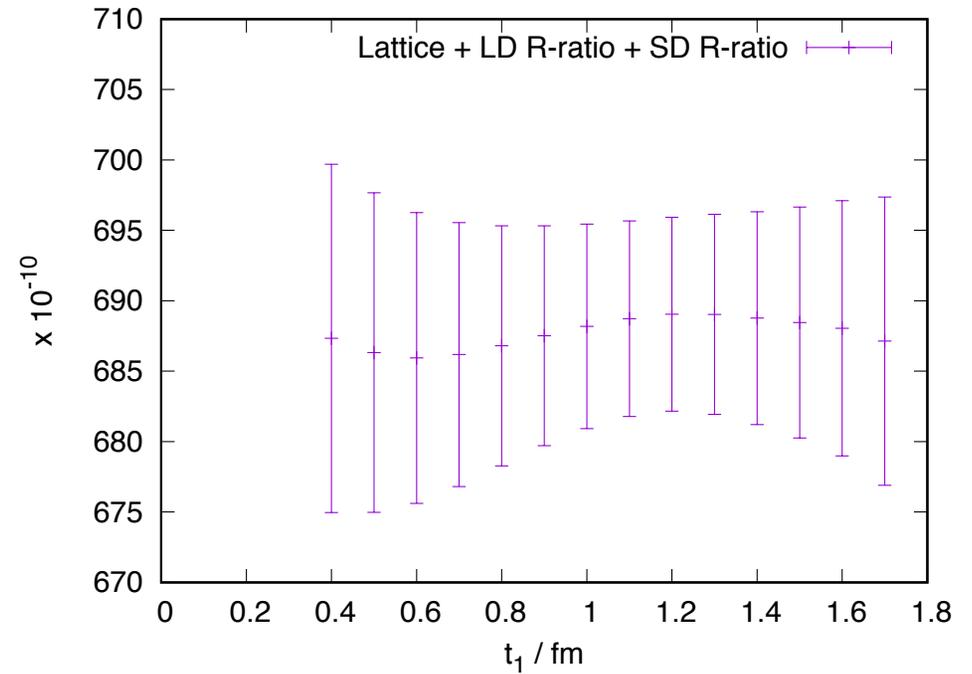
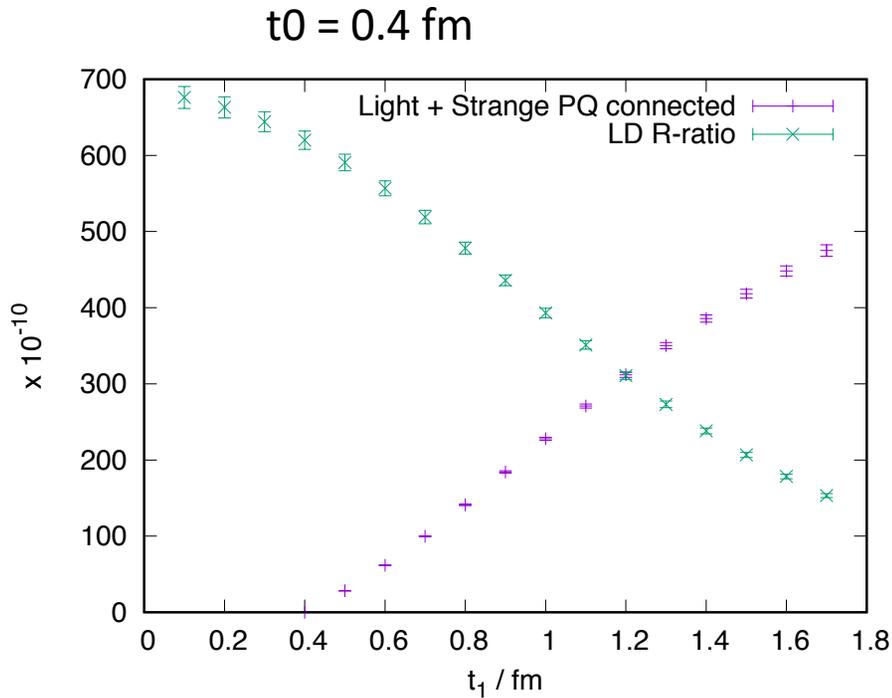
How does this translate to the time-like region?



Most of  $\pi\pi$  peak is captured by window from  $t_0 = 0.4$  fm to  $t_1 = 1.5$  fm, so replacing this region with lattice data reduces the dependence on BaBar versus KLOE data sets.



# R-ratio + Lattice



$t_1$  dependence is flat  $\Rightarrow$  a consistency between R-ratio and Lattice

$t_1 = 1.2 \text{ fm}$ , R-ratio : Lattice = 50:50

$t_1 = 1.2 \text{ fm}$  current error (note 100% correlation in R-ratio) is minimum

# 2022/2023 HVP update

- New fine ensemble 96l to check discretization error
- 6 accompanying smaller / heavy pion QCD samples to correct and check various small mistuning and systematic errors, and Nf=2+1+1 ensembles to check sea charm quark effects

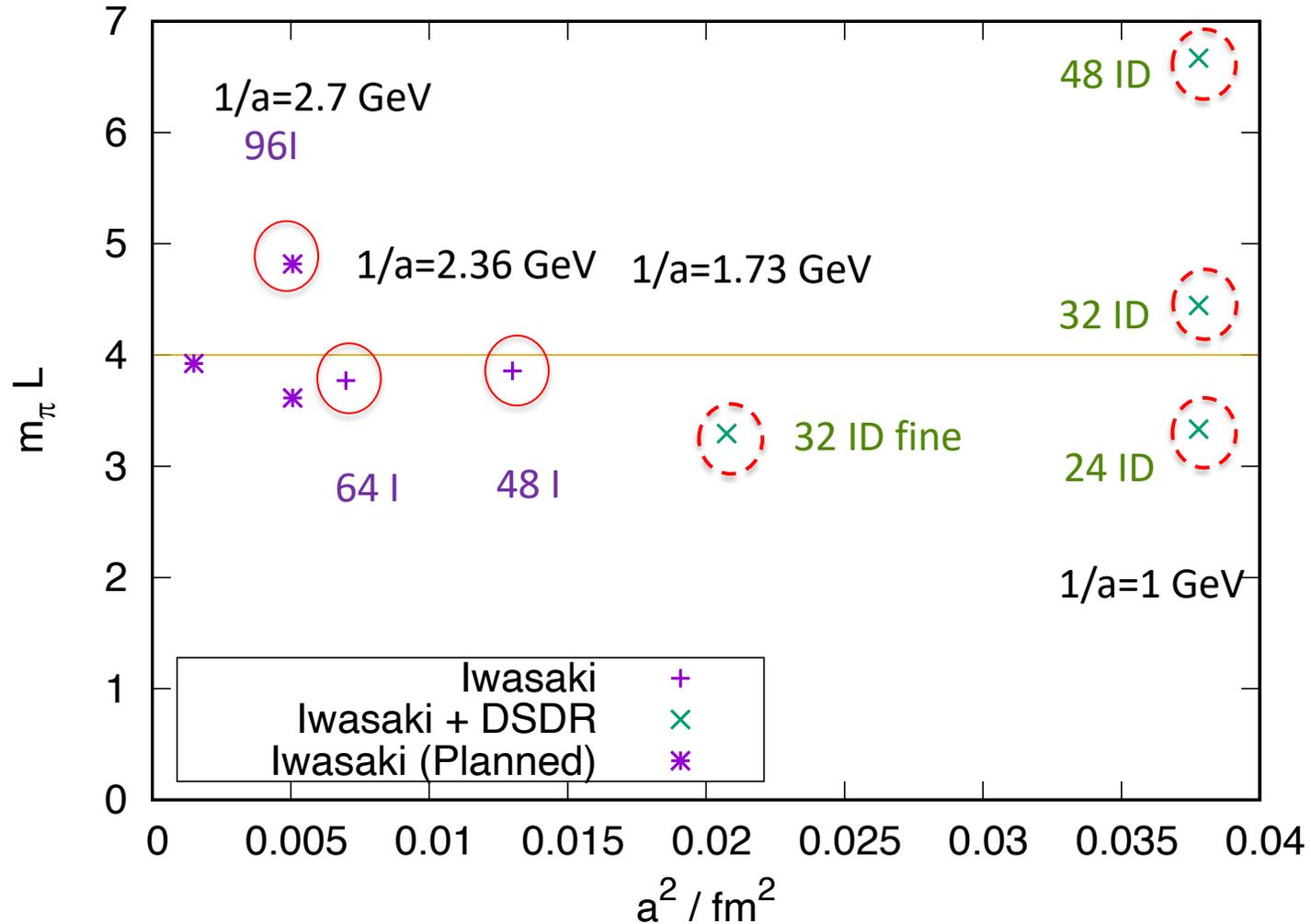
Name	$m_l$	$m_s$	$L_s$	$m_\pi$ /GeV	$m_K$ /GeV	$a^{-1}$ /GeV	$V$	$m_\pi L$	$\beta$
48l	0.00078	0.0362	24	141	500	1.730	$48^3 \times 96$	3.92	2.13
<a href="#">48l ml</a>	0.0025	0.0362	24	210	512	1.730	$32^3 \times 64$	3.89	2.13
<a href="#">48l ms</a>	0.0025	0.05	24	208(2)	600(3)	1.730	$32^3 \times 64$	3.85	2.13
<a href="#">48l ml2 Ls</a>	0.0055	0.0368	32	288	530	1.730	$24^3 \times 48$	4.00	2.13
<a href="#">48l ml2 Ls2</a>	0.002356	0.03366	8	280	530	1.730	$24^3 \times 48$	3.88	2.12
<a href="#">48l ml2</a>	0.0049	0.0362	24	280	530	1.730	$24^3 \times 48$	3.88	2.13
<a href="#">64l ml2</a>	0.00372	0.0257	12	280	530	2.359	$32^3 \times 64$	3.80	2.25

- Blind analysis by 5 groups (HVP) and 2 groups (lattice scale and quark mass)

$$C_b(t) = (b_0 + b_1 a^2 + b_2 a^4) C_0(t)$$

- Among many other **continuum extrapolation** and **Finite Volume** correction were significant and scrutinized

# New $N_f=2+1$ DWF QCD ensemble at physical quark mass



# Lattice EM currents: Two Operators and Three normalization schemes

- $C(t) = \sum_{x,i} \langle V_i(x,t) V_i(0) \rangle$
- Two variants of lattice EM vector current  $V_{i(x,t)}$  are used to check discretization error

$$V_i(x) = \bar{q}(x) \gamma_i q(x) \quad (\text{local current})$$

$$V_i(x) = \overline{\psi(x)} (1 + \gamma_i) U_i(x) \psi(x + \hat{\mu}) + c. c. \quad (\text{conserved current})$$

- EM vector current  $V_i(x,t)$  on lattice is matched to continuum current multiplicatively

$$V_i(x; cont) = Z_V V_i(x; lattice)$$

by matching matrix element of operator to a state

Three variants of states :

Z\_V : 0 momentum single pion state,

Z\_K : 0 momentum single Kaon state

Z\_r : a 0 momentum state specified by the Euclidean distance from another vector operator

# Finite Volume correction estimates

- FV correction by long-distance two pion contributions
  - scalar QED
  - Using pion form factor (Gounaris-Sakurai parametrization) & Lellouche Luscher's FV formula
  - Hansen-Patella FV correction

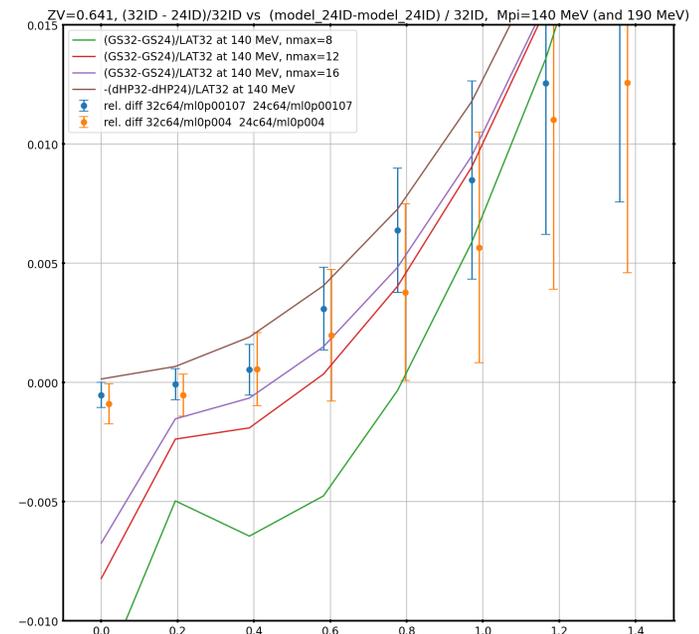
$$a_{\mu}^{\text{HVP,LO}}(T, L) = \text{diagram}$$

$$a_{\mu}^{\text{HVP}}(L = 6.22 \text{ fm}) - a_{\mu}^{\text{HVP}}(L = 4.66 \text{ fm})$$

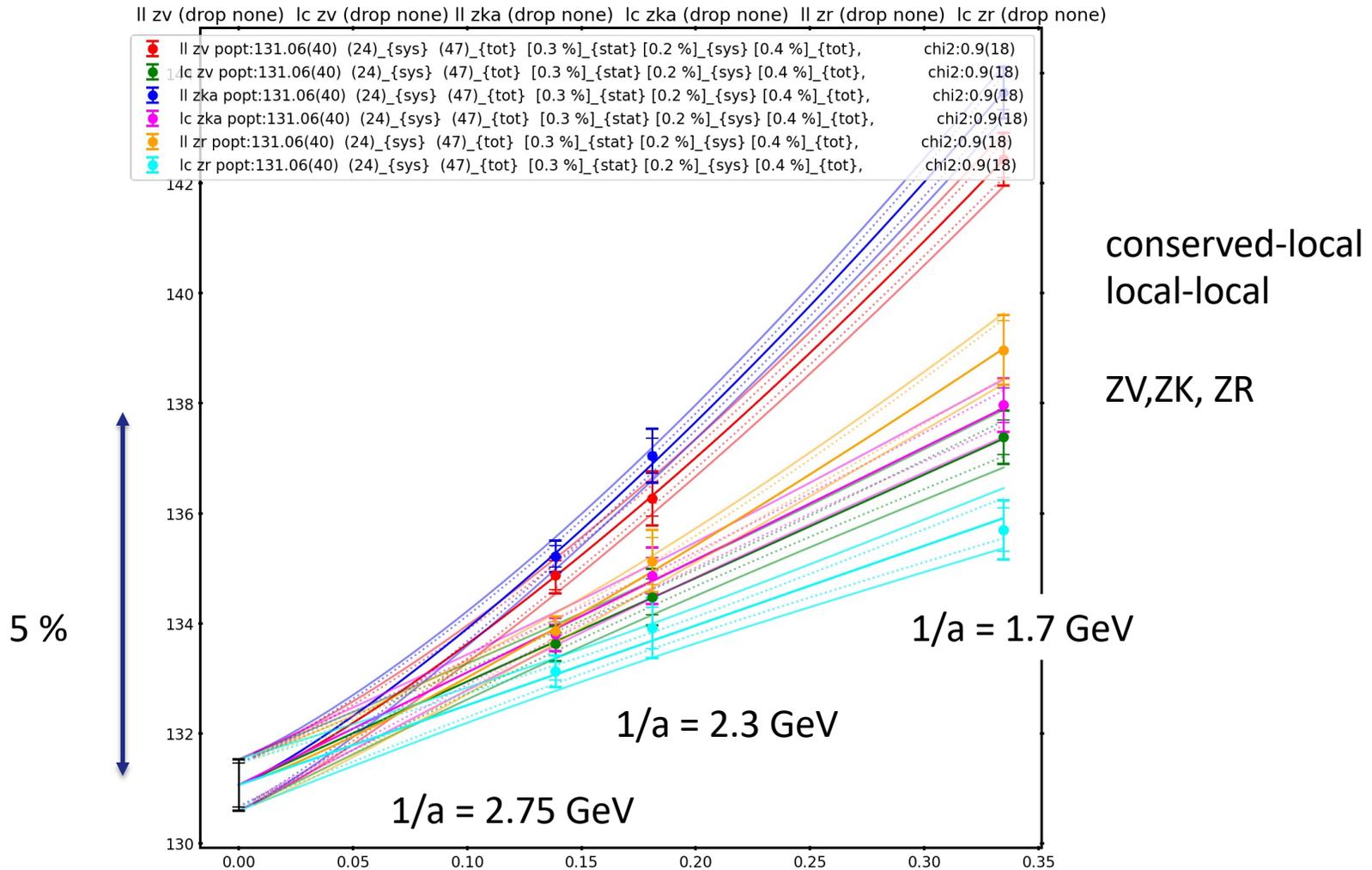
$$= \begin{cases} 12.2 \times 10^{-10} & \text{sQED} \\ 21.6(6.3) \times 10^{-10} & \text{LQCD} \\ 20(3) \times 10^{-10} & \text{GSL} \end{cases}$$

- Revised FV estimation :

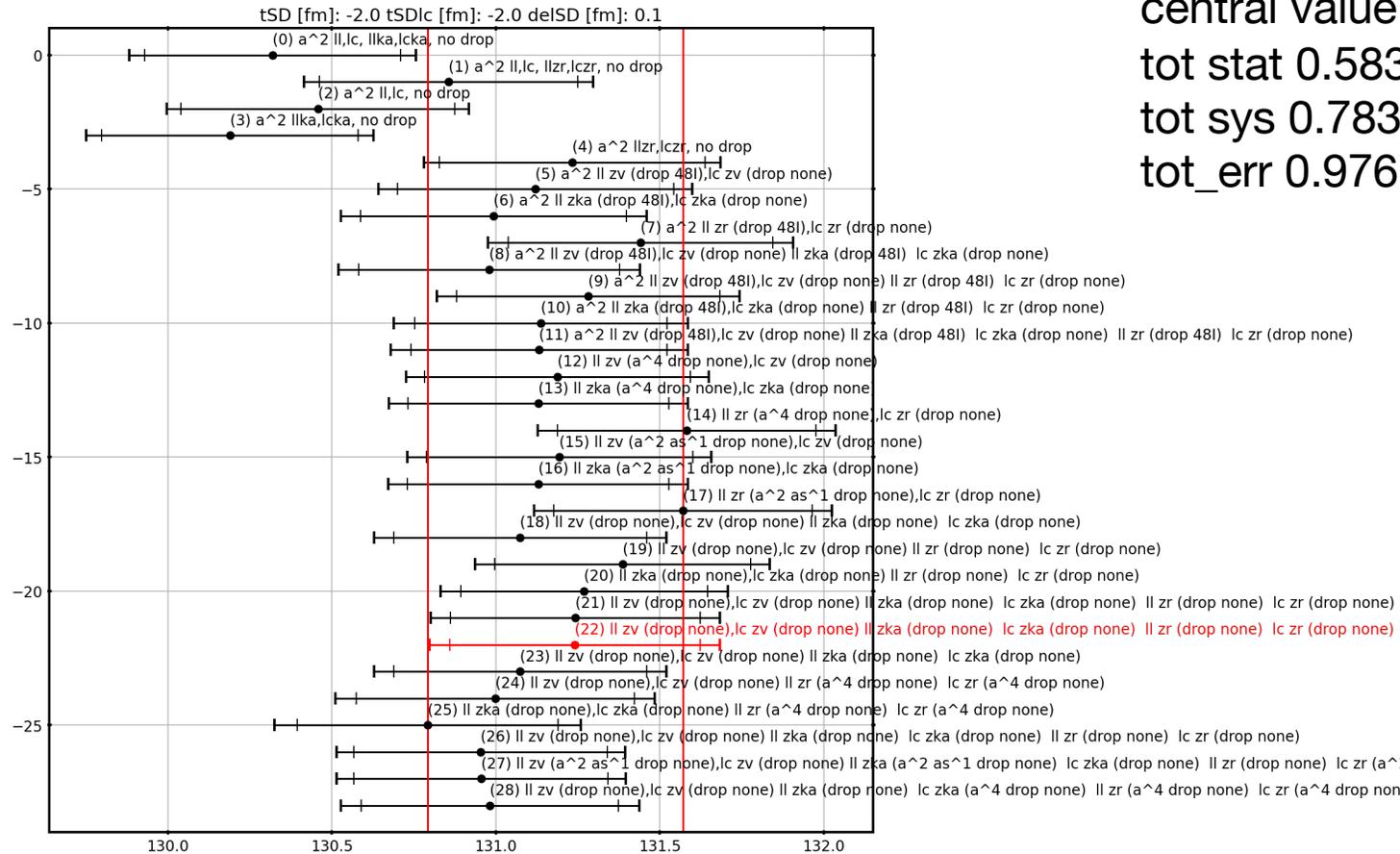
$$a_{\mu}^{\text{HVP}}(L = \infty) - a_{\mu}^{\text{HVP}}(L = 5.47 \text{ fm}) = 22(1) \times 10^{-10}$$



# Continuum limit extrapolation window value [0.4 fm, 1.0 fm]



- Fit forms,  $a^2$ ,  $a^2 + a^4$ ,  $a^2 + a^2 \log(a)$
- Two currents x Three ZV

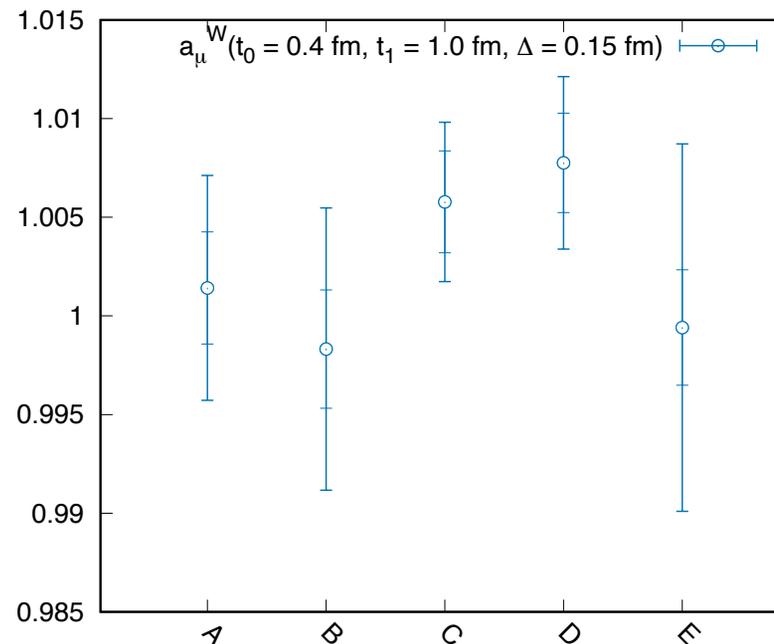


# Blind analysis

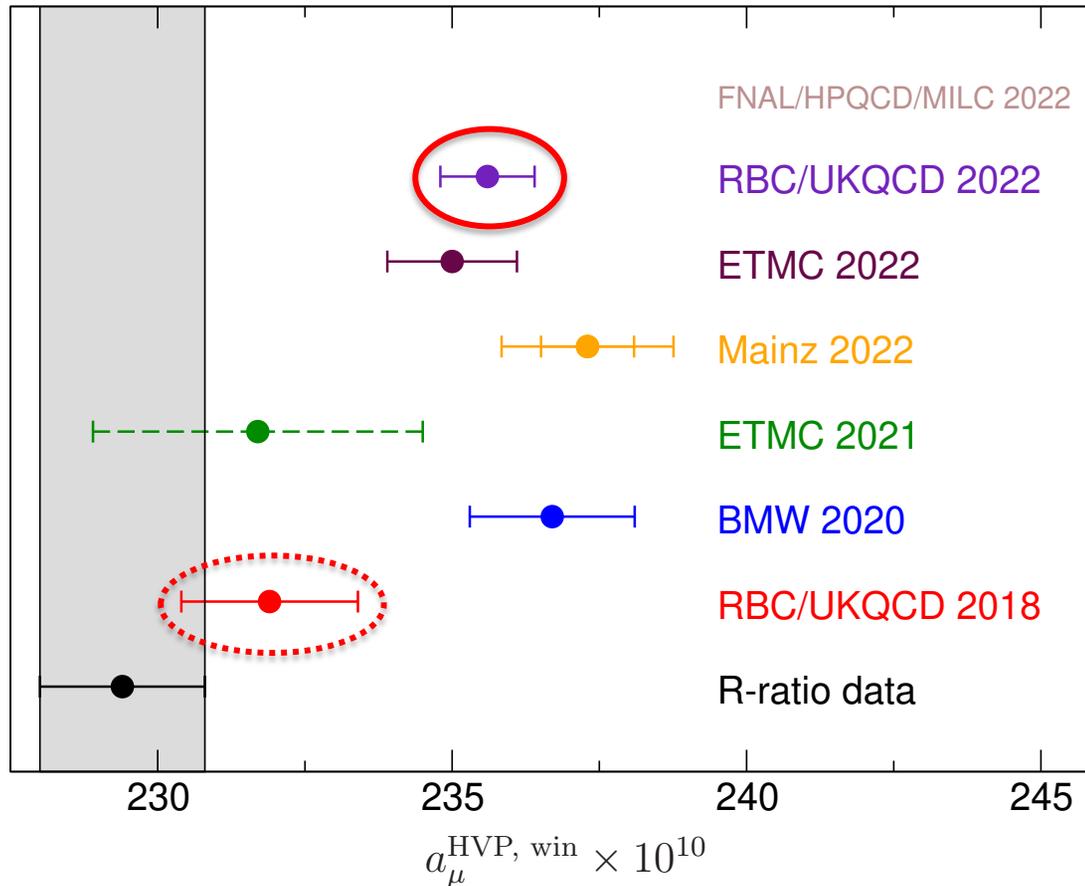
- Blinding procedure :  
lattice spacing dependent blind factor

$$C_b(t) = (b_0 + b_1 a^2 + b_2 a^4) C_0(t)$$

- 5 groups



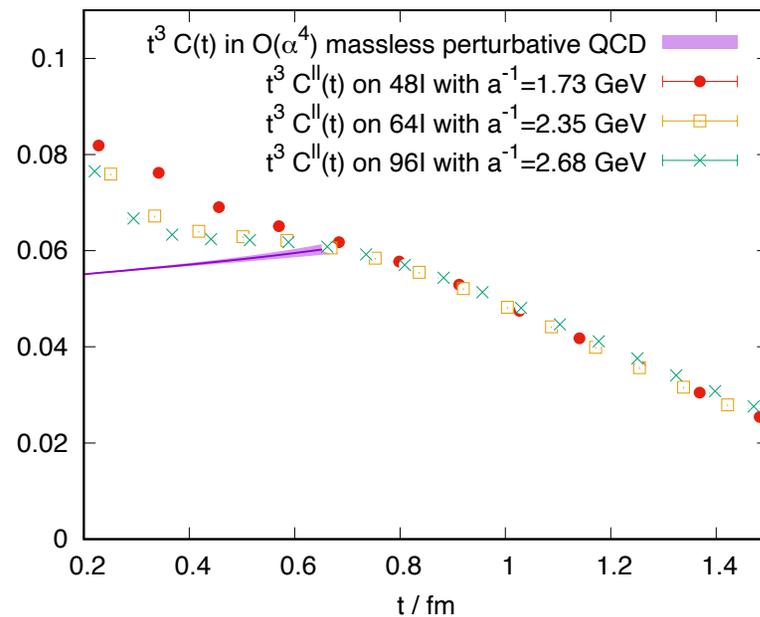
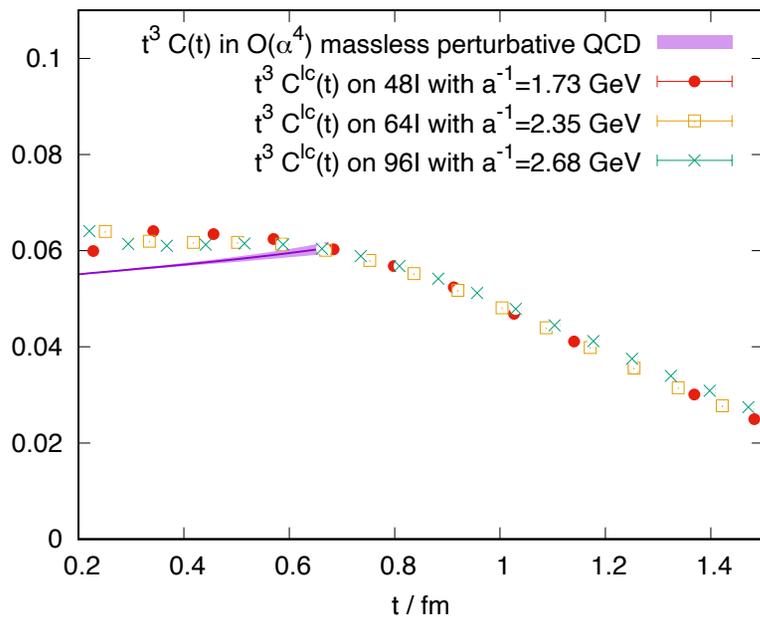
# Mid-Window value of $g-2$ Intermediate energy region



total amu  $\sim 700$ , exp vs theory tension  $\sim 25$   
mid-window value [0.4 fm, 1.0 fm]  $\sim$  rho meson peak

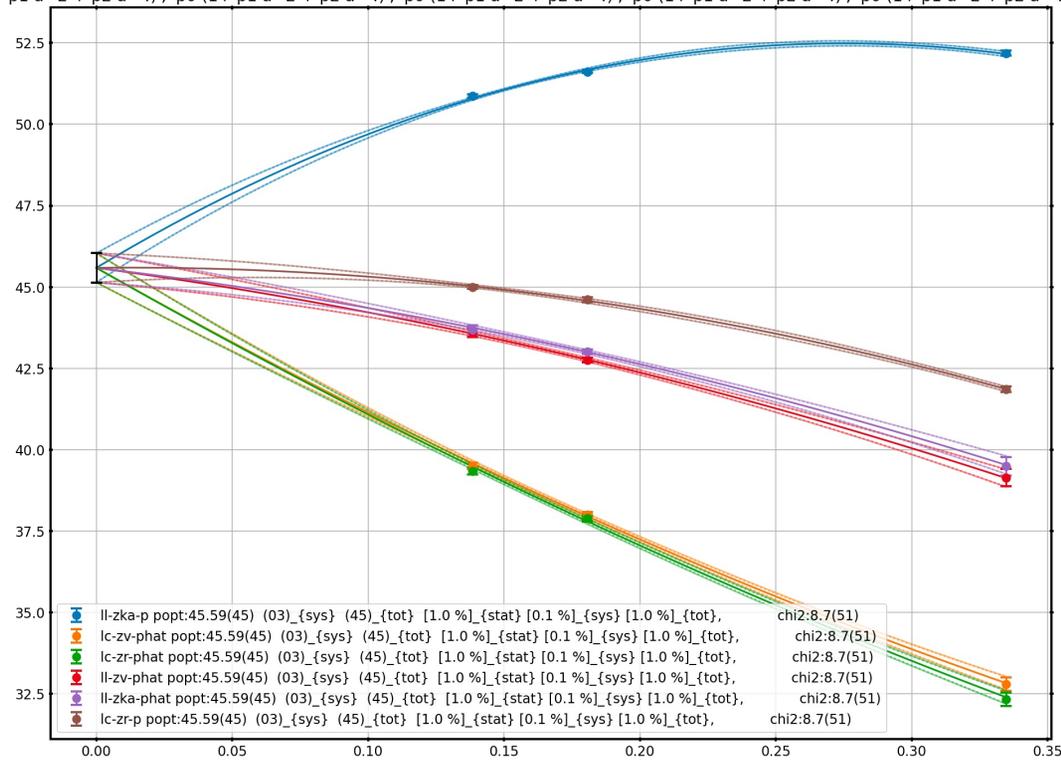
# Short distance window discovered a large discretization error

- $t^3 C(t)$  local-conserved vs local-local
- pQCD, 3 lattice spacings
- Large discretization error at short distance



# Short distance Window (after MF improved tree-level correction)

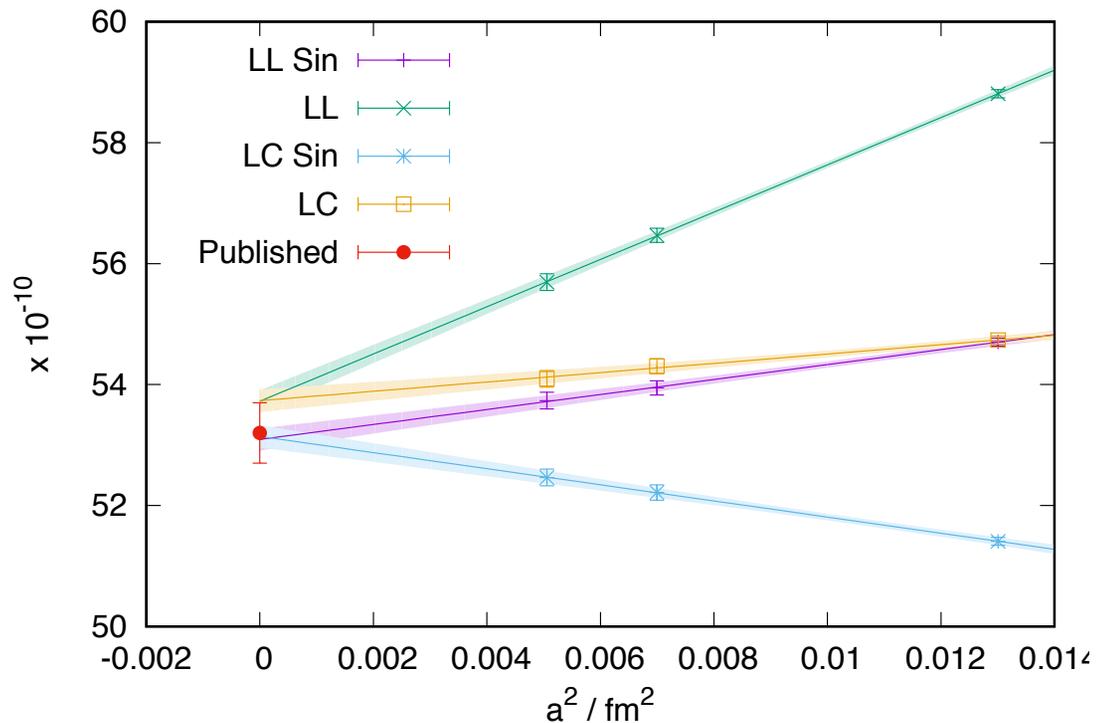
tSD [fm]: -2.0 tSDlc [fm]: -2.0 deISD [fm]: 0.1  
 ['p0 (1+ p1 a^2 + p2 a^4)', 'p0 (1+ p1 a^2 + p2 a^4)'] t0=-2.0 fm, t1=0.4 fm, delta = 0.1 fm



# Strange quark contribution

Add  $a^{-1} = 2.77$  GeV lattice spacing

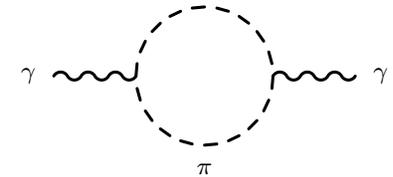
- ▶ Third lattice spacing for strange data ( $a^{-1} = 2.77$  GeV with  $m_\pi = 234$  MeV with sea light-quark mass corrected from global fit):



- ▶ For light quark need new ensemble at physical pion mass. Started run on Summit Machine at Oak Ridge this year ( $a^{-1} = 2.77$  GeV with  $m_\pi = 139$  MeV).

# Reconstruction of HVP for LD from multi-channel Greens function

- Correlation function among N operators  $O_n$ ,  $n=0,1,\dots, N-1$



- Point (or smeared) vector  $O_0 = \sum_x \bar{\psi}(x) \gamma_\mu \psi(x)$ ,  $\mu \in \{1, 2, 3\}$

- 2  $\pi$  operator

$$O_n = \left| \sum_{xyz} \bar{\psi}(x) f(x-z) e^{-i\vec{p}_\pi \cdot \vec{z}} \gamma_5 f(z-y) \psi(y) \right|^2$$

- (4  $\pi$  operator)

$$O_{4\pi} = \left| \sum_{xyz} \bar{\psi}(x) f(x-z) e^{-i\vec{p}_\pi \cdot \vec{z}} \gamma_5 f(z-y) \psi(y) \right|^2 \left| \sum_{xy} \bar{\psi}(x) f(x-y) \gamma_5 \psi(y) \right|^2$$

two pion  
rho-resonance

- NxN correlation function  $\langle O_i(t) O_j(0) \rangle$  (using distillation)

- Solve NxN spectrum  $E_n$  of eigenstates  $|E_n\rangle$  and Overwrap factors  $\langle E_n | O_0 | 0 \rangle$  (GEVP)

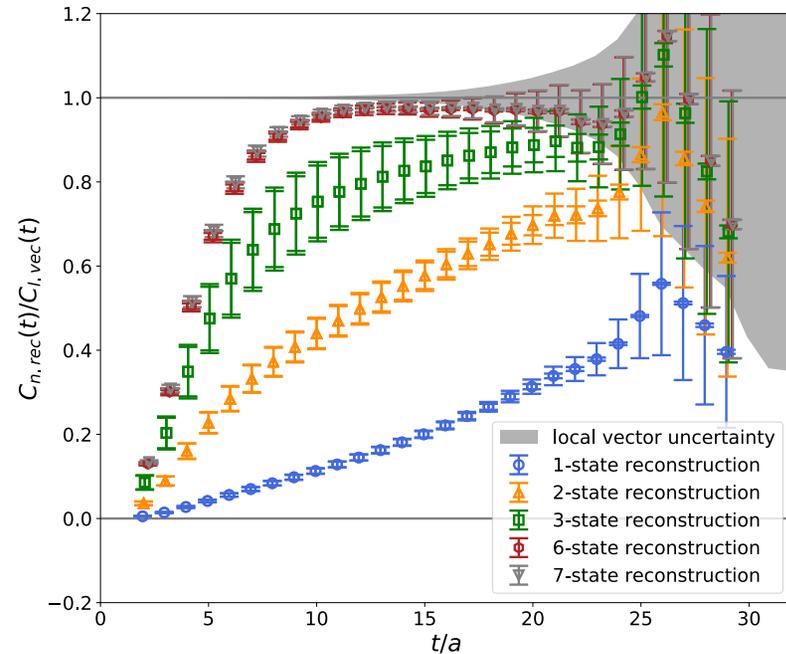
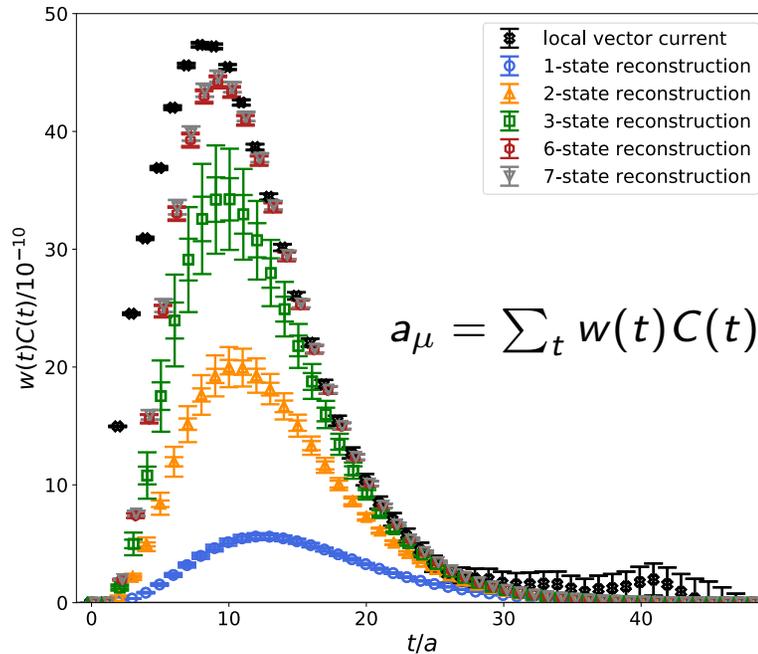
- Reconstruct V-V correlator, and bound contribution from the (N+1)-th states and above

$$\langle O_0(t) O_0^\dagger(0) \rangle = \sum_{n=0}^{N-1} |\langle 0 | O_0 | n \rangle|^2 e^{-E_n t}$$

+ (contributions from  $n \geq N$  states)

# GEVP & Reconstruction of I=1 VV

[ Aaron Meyer ]



Left:  $a_\mu$  integrand, Right: ratio reconstruction/local vector

- ▶ More states  $\implies$  better reconstruction
- ▶ 6 state  $\implies$   $1\sigma$  consistent at  $t \geq 16a \sim 1.7$  fm

# Bounds for $a_\mu$

- Upper & lower bounds from unitarity

$$\tilde{C}(t; t_{\max}, E) = \begin{cases} C(t) & t < t_{\max} \\ C(t_{\max})e^{-E(t-t_{\max})} & t \geq t_{\max} \end{cases}$$

Upper bound:  $E = E_0$ , lowest state in spectrum

Lower bound:  $E = \log\left[\frac{C(t_{\max})}{C(t_{\max}+1)}\right]$

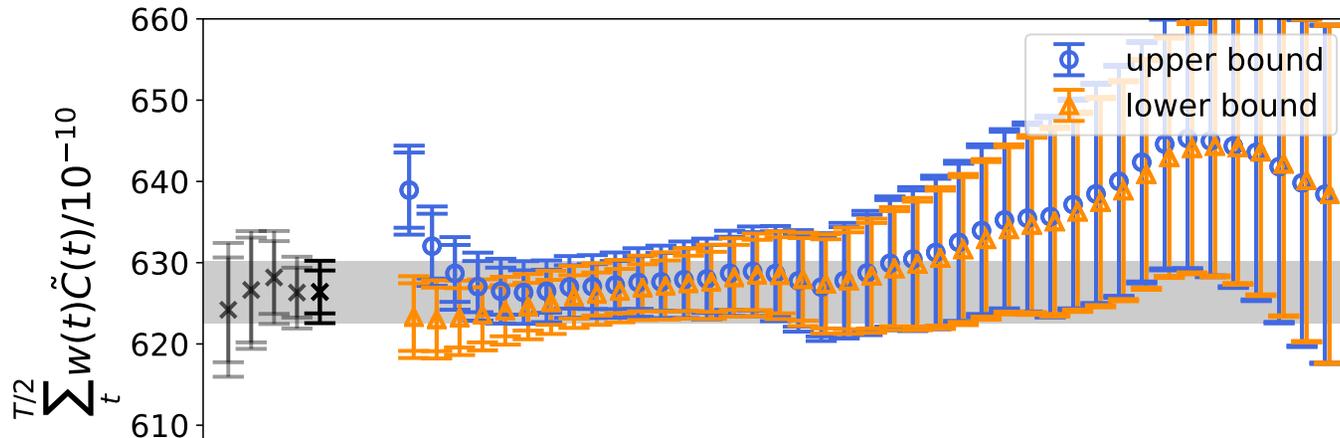
- Also bounds for the  $n$  in  $[N+1, \infty]$  states contribution

Replace  $C(t) \rightarrow C(t) - \sum_n^N |c_n|^2 e^{-E_n t}$

$\implies$  Long distance convergence now  $\propto e^{-E_{N+1} t}$

$\implies$  Smaller overall contribution from neglected states

# GEVP + Bounding Method [A. Meyer]



a factor of 2.5 smaller statistical error by bounding method  
 a factor of 5 smaller statistical error by bounding method + 5 state reconstruction

continuum limit and other systematic studies are on-going (again blinding)

Bounding method  $t_{\max} = 5.0$  fm, 1 state reconstruction:  $a_{\mu}^{\text{HVP}} = 620.0(7.2)$

Bounding method  $t_{\max} = 2.5$  fm, 2 state reconstruction:  $a_{\mu}^{\text{HVP}} = 628.2(5.7)$

Bounding method  $t_{\max} = 2.1$  fm, 5 state reconstruction:  $a_{\mu}^{\text{HVP}} = 626.3(4.4)$

Bounding method  $t_{\max} = 1.7$  fm, 6 state reconstruction:

$$a_{\mu}^{\text{HVP, conn, iso, 48l}} = 626.6(2.7)_{\text{stat}}(0.4)_{Z_V, 48l}(2.6)_{a^{-1}, 48l}(0.5)_{\text{bound}}(0.5)_{\text{exc}}$$

Bounding method gives factor of 2.5 improvement over no bounding method

Improving the bounding method increases gain to factor of 5, including systematics

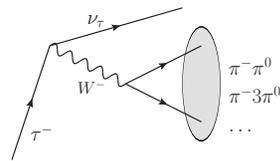
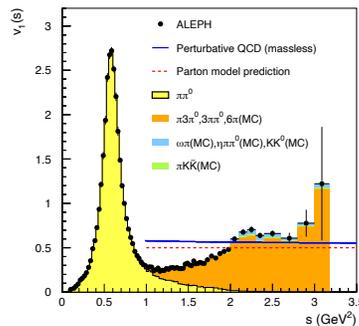
# Use of tau decay data

## [ M. Bruno ]

- Belle II : tau factory
- Isospin corrections from Lattice QCD
- Different systematic errors to cross-check

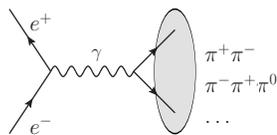


### MOTIVATIONS FOR $\tau$



V - A current

Final states  $I = 1$  charged



EM current

Final states  $I = 0, 1$  neutral

$\tau$  data can improve  $a_\mu[\pi\pi]$   
 $\rightarrow$  72% of total Hadronic LO  
 $\rightarrow$  competitive precision on  $a_\mu^W$



### NEUTRAL VS CHARGED

$$\frac{i}{2}(\bar{u}\gamma_\mu u - \bar{d}\gamma_\mu d), \left[ \begin{matrix} I = 1 \\ I_3 = 0 \end{matrix} \right] \rightarrow j_\mu^{(1,-)} = \frac{i}{\sqrt{2}}(\bar{u}\gamma_\mu d), \left[ \begin{matrix} I = 1 \\ I_3 = -1 \end{matrix} \right]$$

$$\text{Isospin 1 charged correlator } G_{11}^W = \frac{1}{3} \sum_k \int dx \langle j_k^{(1,+)}(x) j_k^{(1,-)}(0) \rangle$$

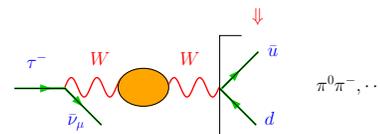
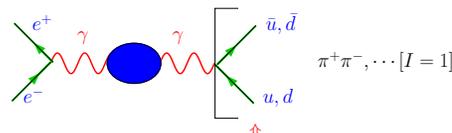
$$\delta G_{11} \equiv G_{11}^\gamma - G_{11}^W \quad [\text{MB et al.' Latt18}]$$

$$= Z_V^4 (4\pi\alpha) \frac{(Q_u - Q_d)^4}{4} \left[ \text{diagram 1} + \text{diagram 2} \right]$$

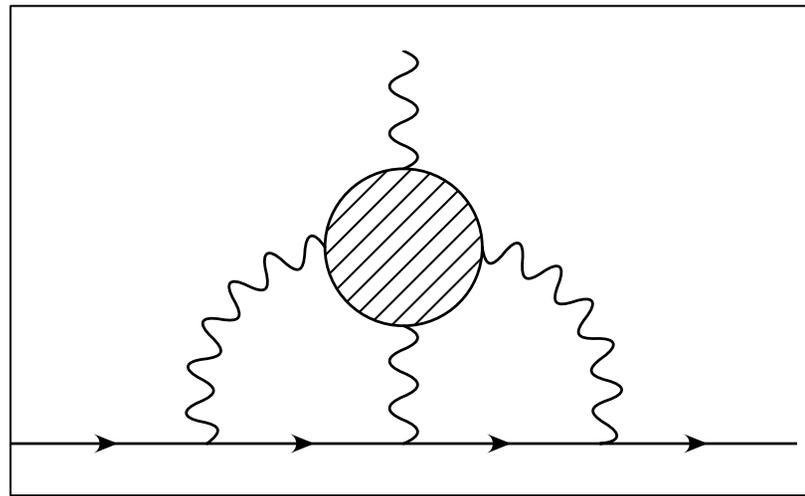
$$G_{01}^\gamma = Z_V^4 \frac{(Q_u^2 - Q_d^2)^2}{2} (4\pi\alpha) \left[ \text{diagram 1} + 2 \times \text{diagram 2} + \text{diagram 3} + \dots \right]$$

$$+ Z_V^2 \frac{Q_u^2 - Q_d^2}{2} (m_u - m_d) \left[ 2 \times \text{diagram 4} + \dots \right]$$

... = subleading diagrams



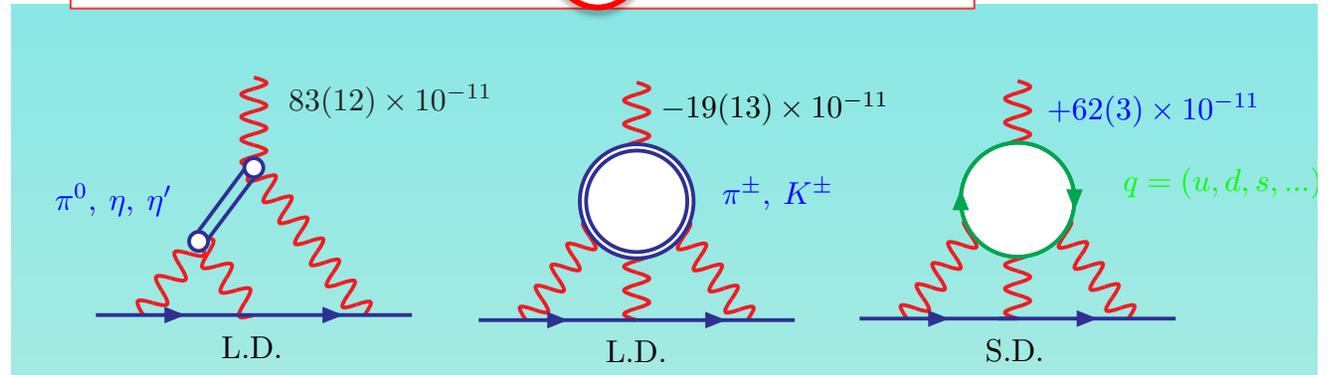
# Hadronic Light-by-Light (HLbL) contributions



# HLbL from Models

- Model estimate with non-perturbative constraints at the chiral / low energy limits using anomaly :  $(9-12) \times 10^{-10}$  with 25-40% uncertainty

$$a_{\mu}^{\text{EXP}} - a_{\mu}^{\text{SM}} = 27.4 \underbrace{(2.7)}_{\text{HVP}} \underbrace{(2.6)}_{\text{HLbL}} \underbrace{(0.1)}_{\text{other}} \underbrace{(6.3)}_{\text{EXP}} \times 10^{-10}$$



F. Jegerlehner ,  $\times 10^{11}$

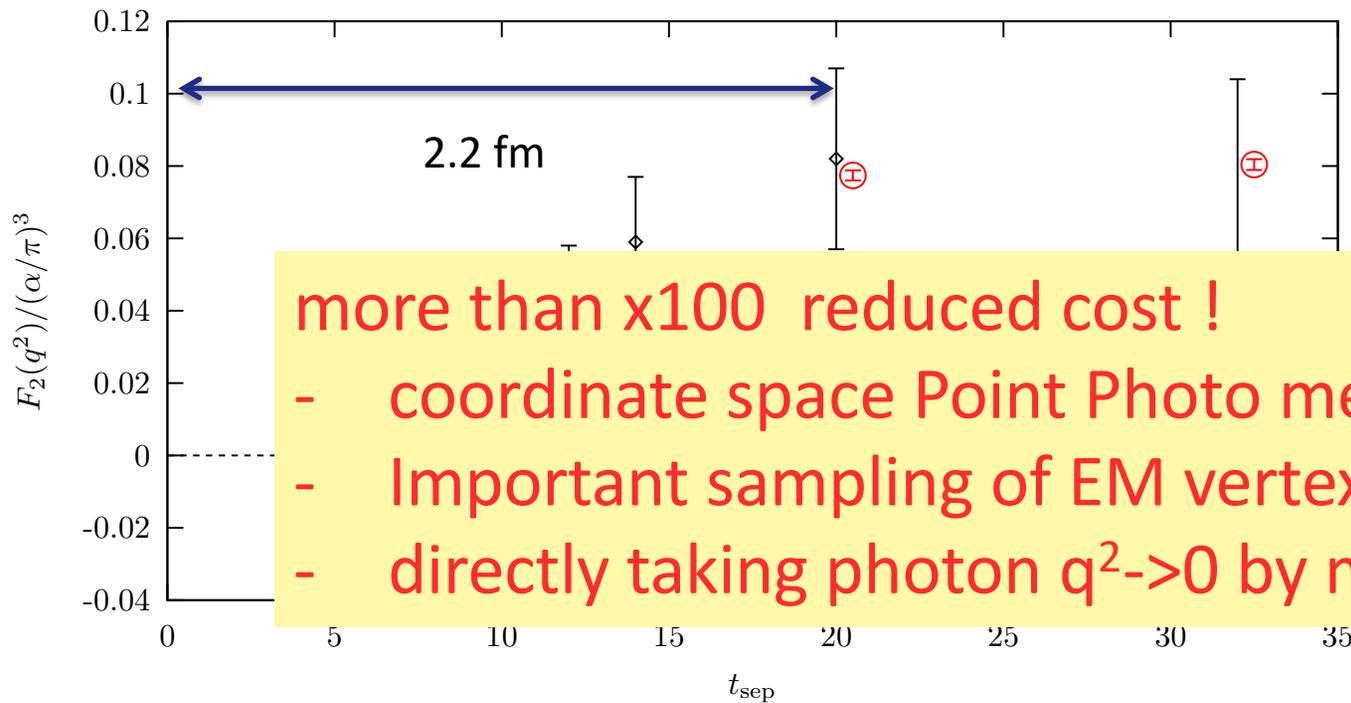
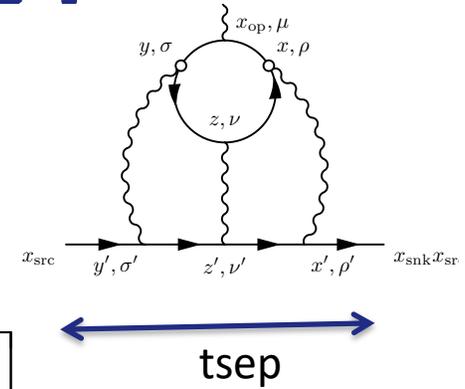
Contribution	BPP	HKS	KN	MV	PdRV	N/JN
$\pi^0, \eta, \eta'$	$85 \pm 13$	$82.7 \pm 6.4$	$83 \pm 12$	$114 \pm 10$	$114 \pm 13$	$99 \pm 16$
$\pi, K$ loops	$-19 \pm 13$	$-4.5 \pm 8.1$	—	$0 \pm 10$	$-19 \pm 19$	$-19 \pm 13$
axial vectors	$2.5 \pm 1.0$	$1.7 \pm 1.7$	—	$22 \pm 5$	$15 \pm 10$	$22 \pm 5$
scalars	$-6.8 \pm 2.0$	—	—	—	$-7 \pm 7$	$-7 \pm 2$
quark loops	$21 \pm 3$	$9.7 \pm 11.1$	—	—	2.3	$21 \pm 3$
total	$83 \pm 32$	$89.6 \pm 15.4$	$80 \pm 40$	$136 \pm 25$	$105 \pm 26$	$116 \pm 39$

# Dramatic Improvement !

## Luchang Jin

$a=0.11$  fm,  $24^3 \times 64$  ( $2.7$  fm) $^3$ ,  
 $m_\pi = 329$  MeV,  $m_\mu \approx 190$  MeV,  $e=1$

$q = 2\pi/L$   $N_{\text{prop}} = 81000$   $\blacklozenge$   
 $q = 0$   $N_{\text{prop}} = 26568$   $\oplus$

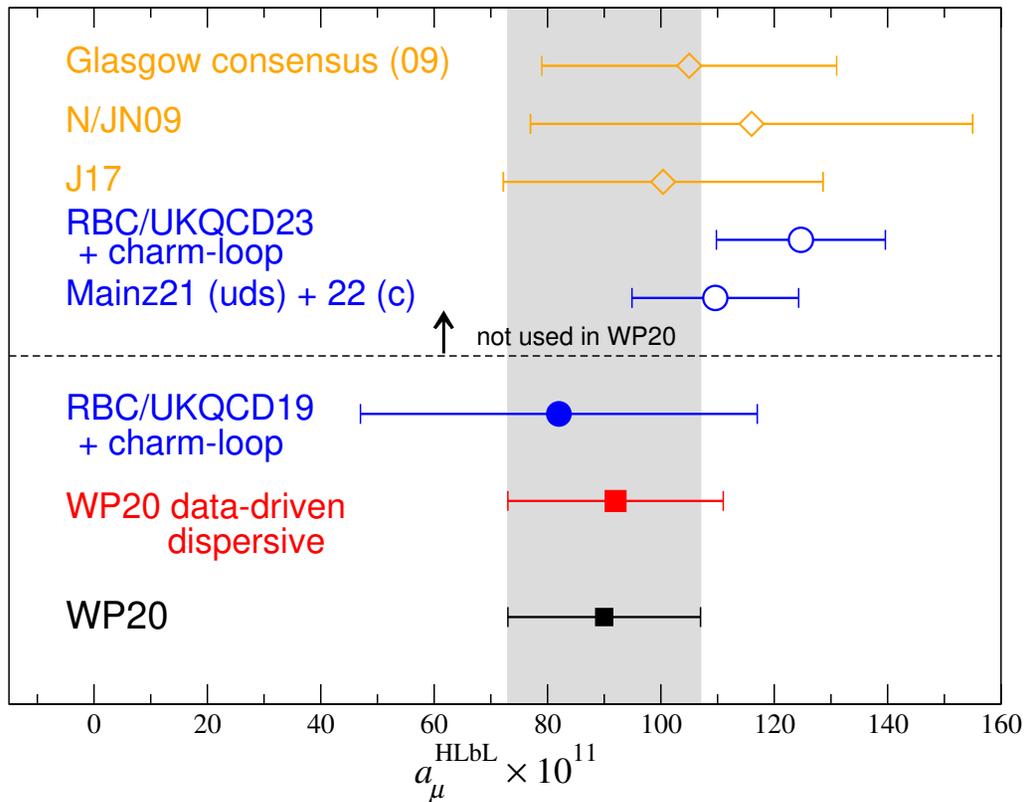


more than x100 reduced cost !

- coordinate space Point Photo method
- Important sampling of EM vertex
- directly taking photon  $q^2 \rightarrow 0$  by moment

Method	$F_2/(\alpha/\pi)^3$	$N_{\text{conf}}$	$N_{\text{prop}}$	$\sqrt{\text{Var}}$
Conserved	0.0825(32)	12	$(118 + 128) \times 2 \times 7$	0.65
Mom.	0.0804(15)	18	$(118 + 128) \times 2 \times 3$	0.24

# 2023 HLbL status



- Lattice QCD Mainz 2021, 2022:

$$a_\mu^{\text{HLbL}}[uds] = 107(15) \times 10^{-11}$$

$$a_\mu^{\text{HLbL}}[c] = 2.8(5) \times 10^{-11}$$

- New result RBC/UKQCD 2023:

$$a_\mu^{\text{HLbL}}[uds] = 122(15) \times 10^{-11}$$

[M. Hoferichter ]

# Summary

- FNAL muon g-2 Run-2, Run-3 results will be announced  
<https://indico.fnal.gov/event/60738/>  
August 10, 10AM central == August 11, 0:00 Japan  
expect a factor of 2 smaller error
- R-ratio data driven approach and muon g-2 has  $4.2 \sigma$  tension,
- New CMD-3 two pion data is significantly larger, may make muon g-2 value consistent with SM, if difference between other R-ratio results will be understood.
- Lattice QCD
- For short distance and mid distance, lattice calculations are now mostly agreeing to each other, a tension with R-ratio may
- For window value [0.4 fm, 1.0 fm], adding new finer ensemble analysis, our 2018 analysis seems to underestimate  $a^4$  discretization error  
 $(a \Lambda)^4 \sim 0.5\%$  for  $\Lambda = 0.4$  fm
- BMW results so far only precise SD, MD, LD, close to g-2
- New full g-2 HVP results including Long Distance contribution soon
- Tau decay experiment input + Lattice IB correction may shed a light
- For HLbL, good agreement between Lattice and phenomenology, but needs another factor of 2 improvement for 10% relative error goal.