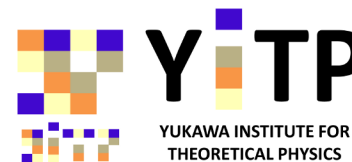

Neutron Stars and Hypernuclei

Akira Ohnishi (YITP, Kyoto U.)

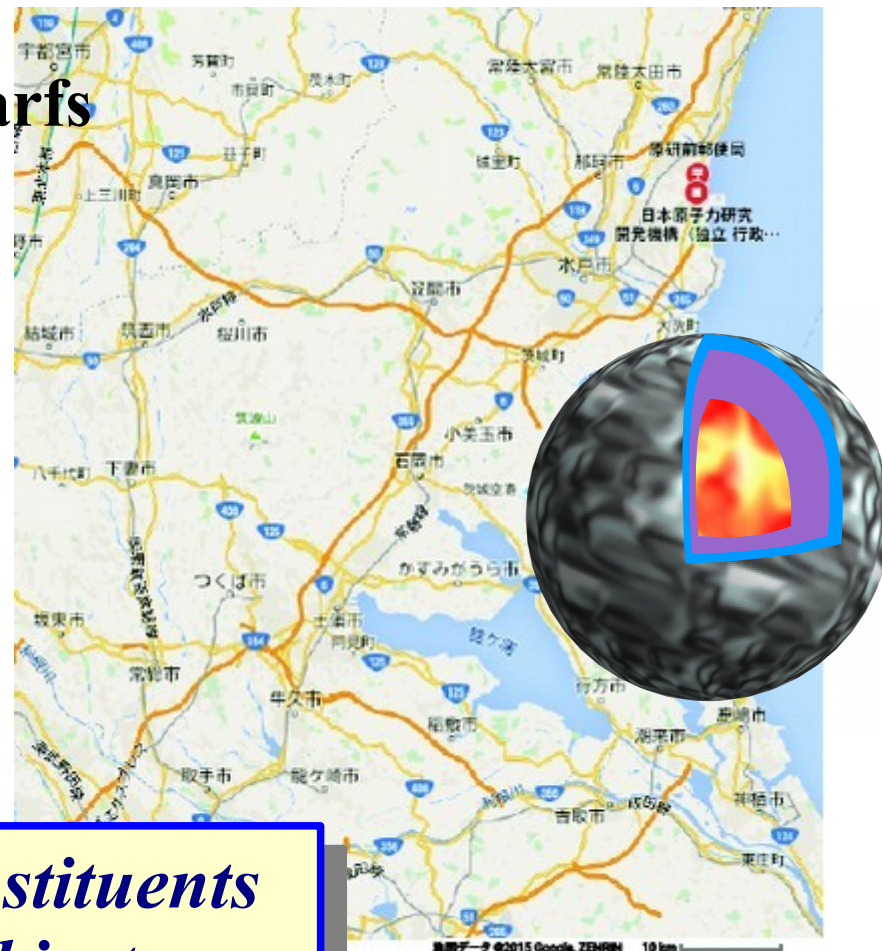
JAEA-ASRC-TPI Theory Lecture Series
June 2, 2015, JAEA

- Introduction
- Part I : Basics of Neutron Star Physics
- Part II : Hypernuclear Physics & Neutron Stars
- Summary



Basic properties of neutron stars

- Mass: $M = (1-2) M_{\odot}$ ($M \sim 1.4 M_{\odot}$)
- Radius: $5 \text{ km} < R < 20 \text{ km}$ ($R \sim 10 \text{ km}$)
- Supported by Nuclear Pressure
c.f. Electron pressure for white dwarfs
- Cold enough
($T \sim 10^6 \text{ K} \sim 100 \text{ eV}$)
compared with
neutron Fermi energy.
- Various constituents
(conjectured)
 $n, p, e, \mu, Y, \bar{K}, \pi, q, g, qq, \dots$

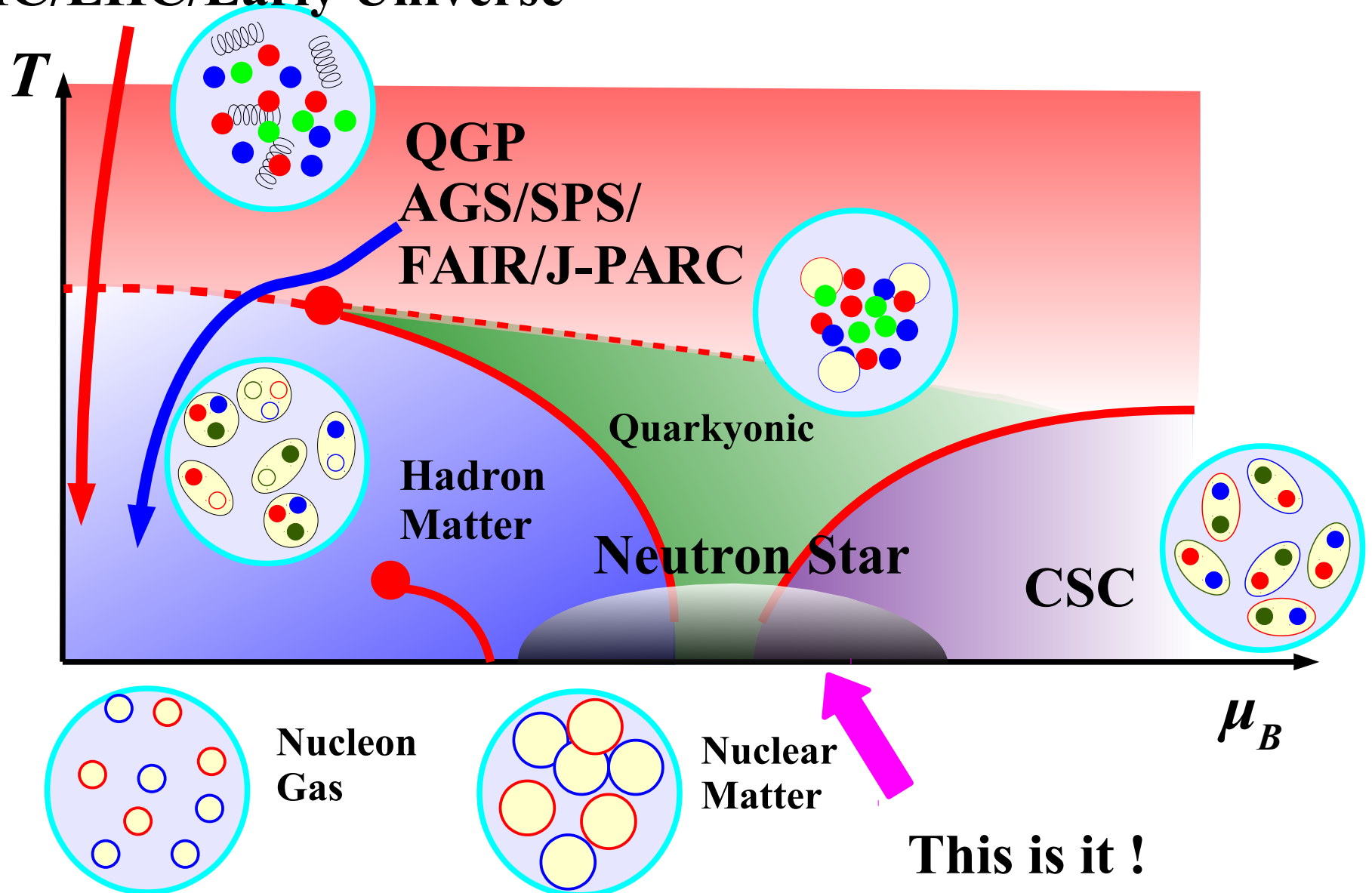


*Wide density range \rightarrow various constituents
NS = high-energy astrophysical objects
and laboratories of dense matter.*

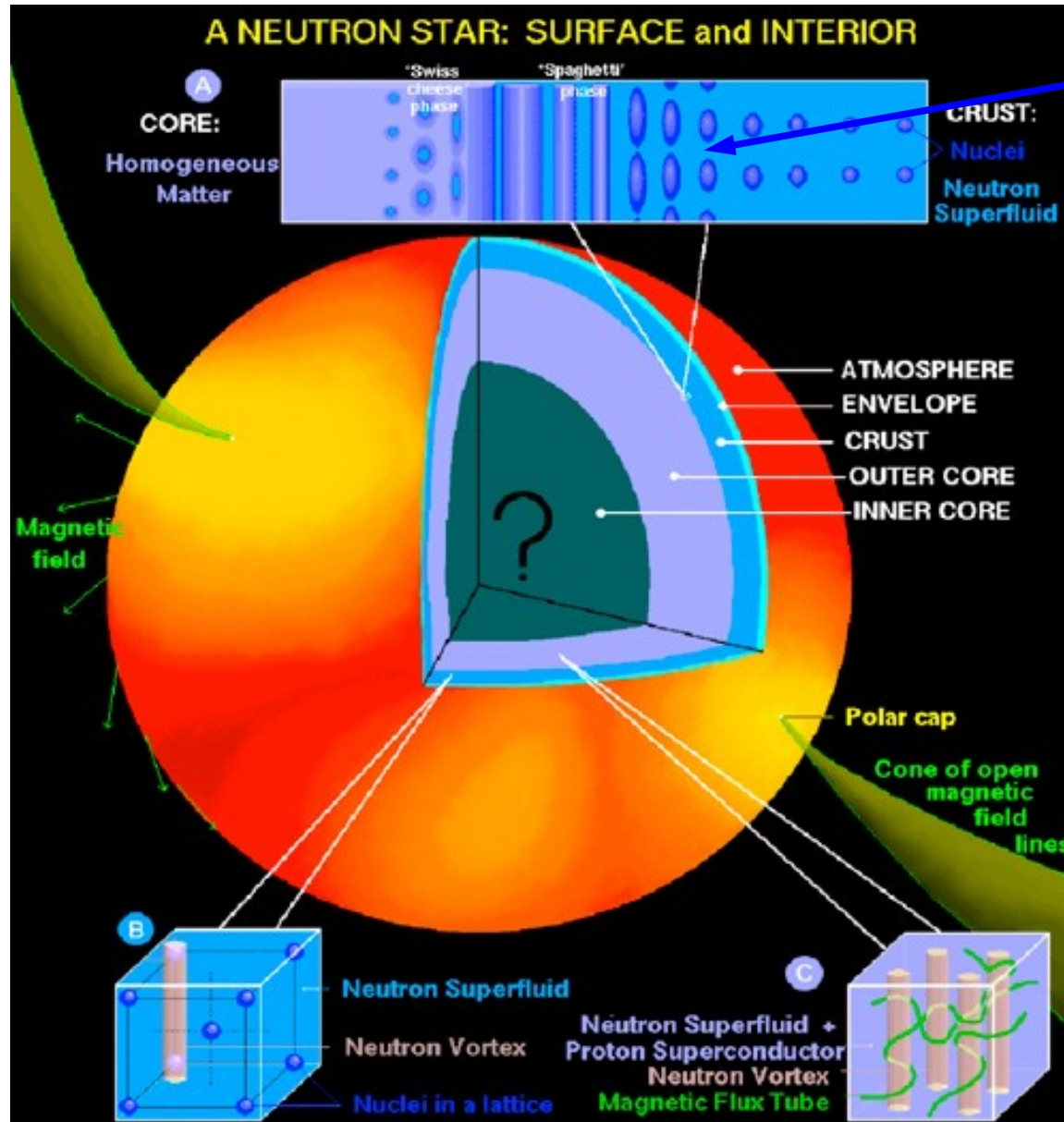
google & zenrin

QCD Phase Diagram

RHIC/LHC/Early Universe



Inside Neutron Stars



For pasta nuclei,
ask Maruyama-san

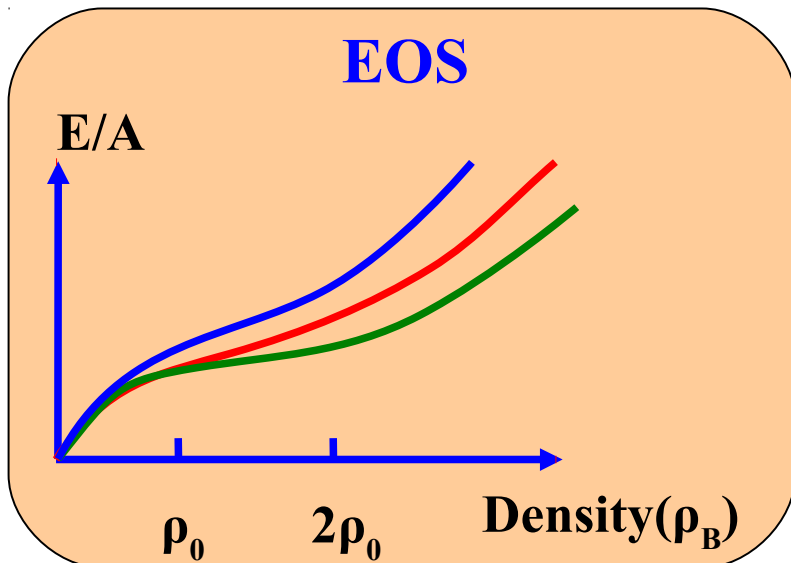
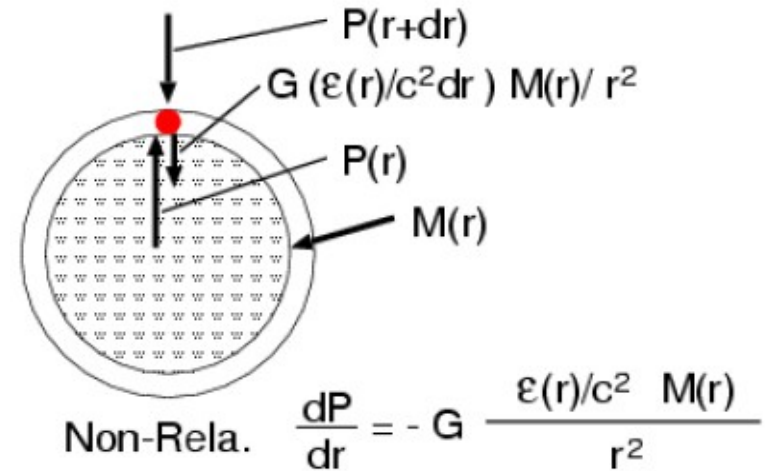
Dany Page

M-R curve and EOS

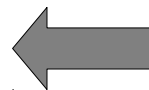
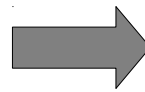
- M-R curve and NS matter EOS has 1 to 1 correspondence
 - TOV(Tolman-Oppenheimer-Volkoff) equation =GR Hydrostatic Eq.

$$\frac{dP}{dr} = -G \frac{(\epsilon/c^2 + P/c^2)(M + 4\pi r^3 P/c^2)}{r^2(1 - 2GM/rc^2)}$$

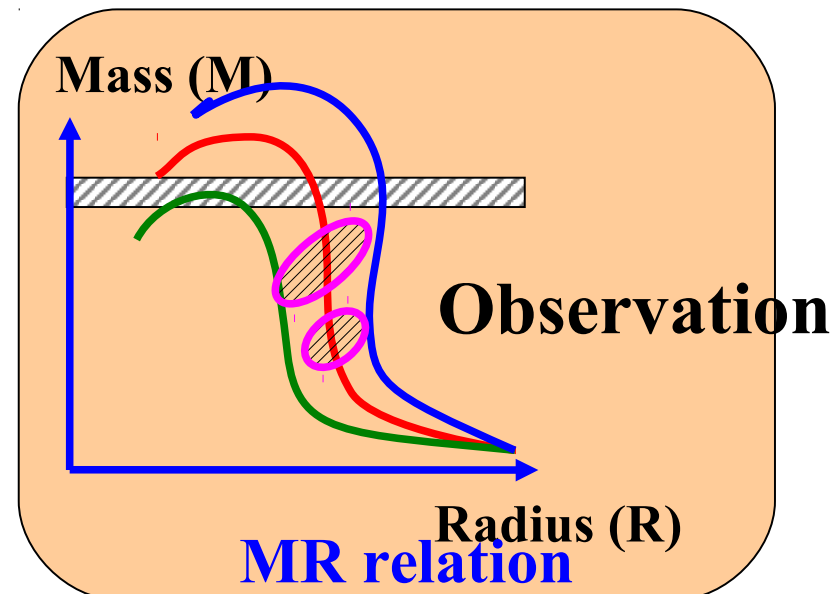
$$\frac{dM}{dr} = 4\pi r^2 \epsilon/c^2, \quad P = P(\epsilon) \quad (\text{EOS})$$



prediction

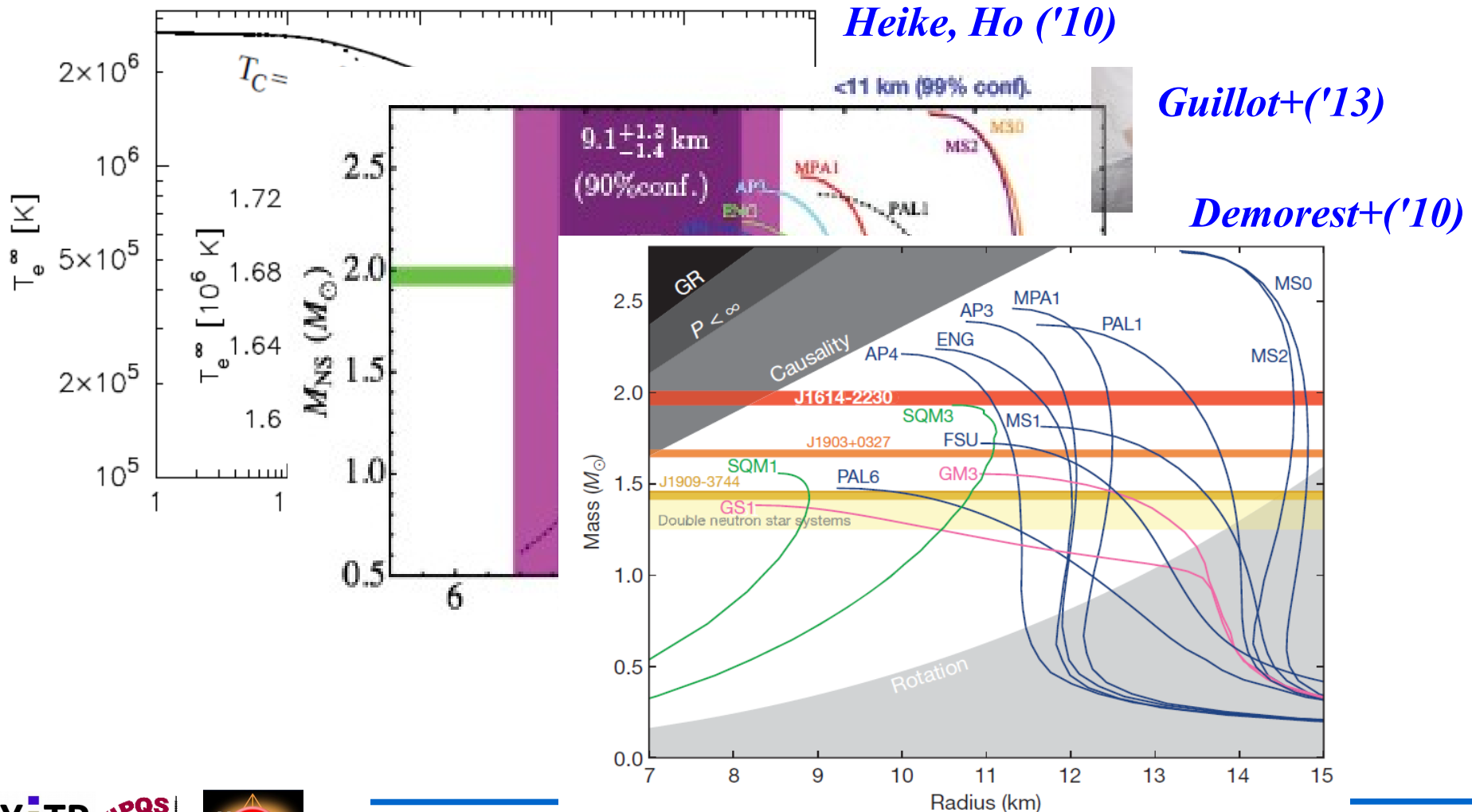


Judge



Puzzles of NS

- Magnetar, NS oscillation,
- Rapid NS cooling puzzle (CasA cools too fast ?)
- Compact NS problem (9 km NS ?)
- Massive NS puzzle ($2 M_{\odot}$ NS ?)



Nuclear Superfluidity and Cooling Curve

■ Surface T measurement and Cooling curve

- Stable superfluid → Gap → Suppression of ν emission

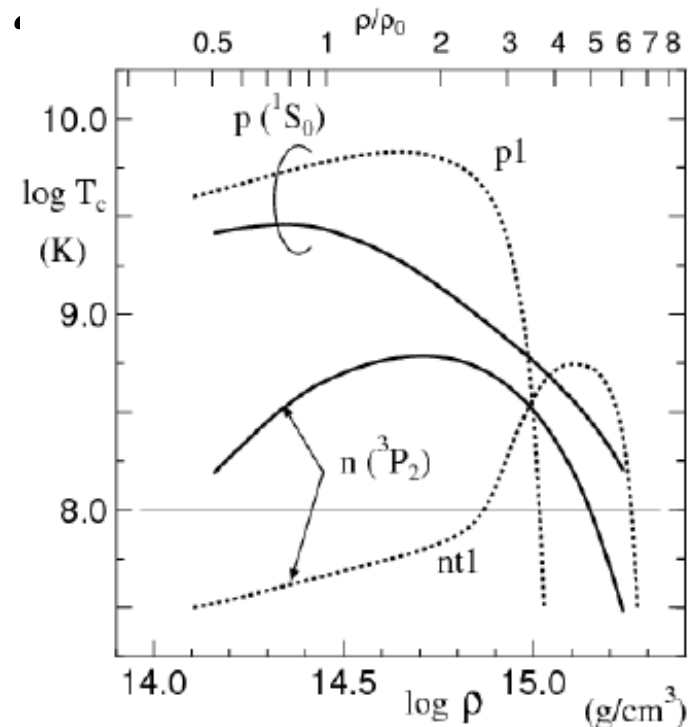
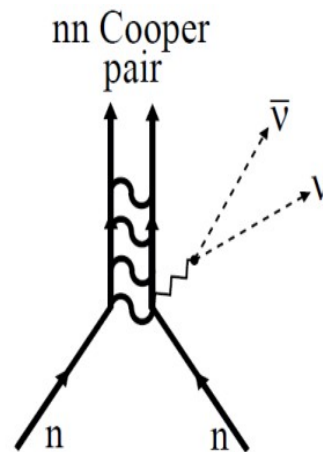
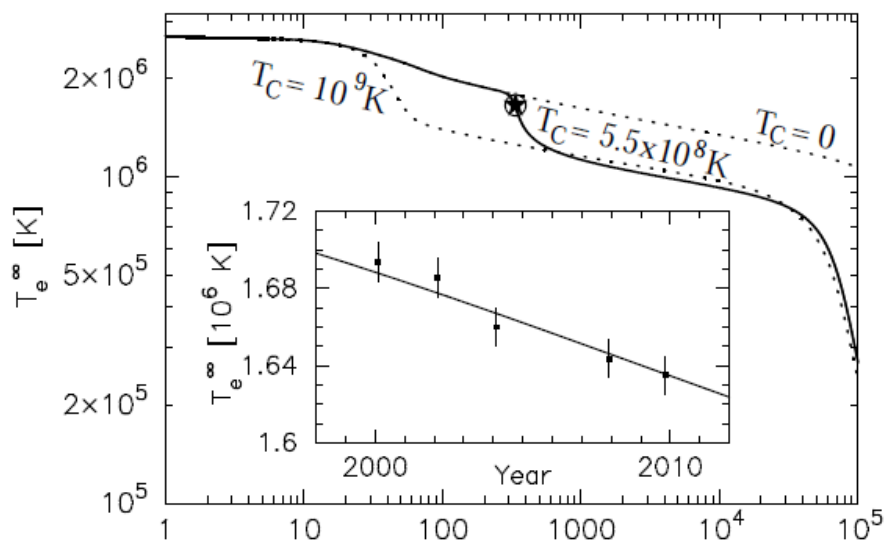
- Onset of superfluidity → Rapid cooling

- Precise T and Cooling rate measurement in Cas A

Heinke, Ho, ApJ 719('10) L167 [arXiv:1007.4719]

Page et al., PRL 106 ('11) 081101 [arXiv:1011.6142]

■ Can we predict the pairing cap around $5\rho_0$?



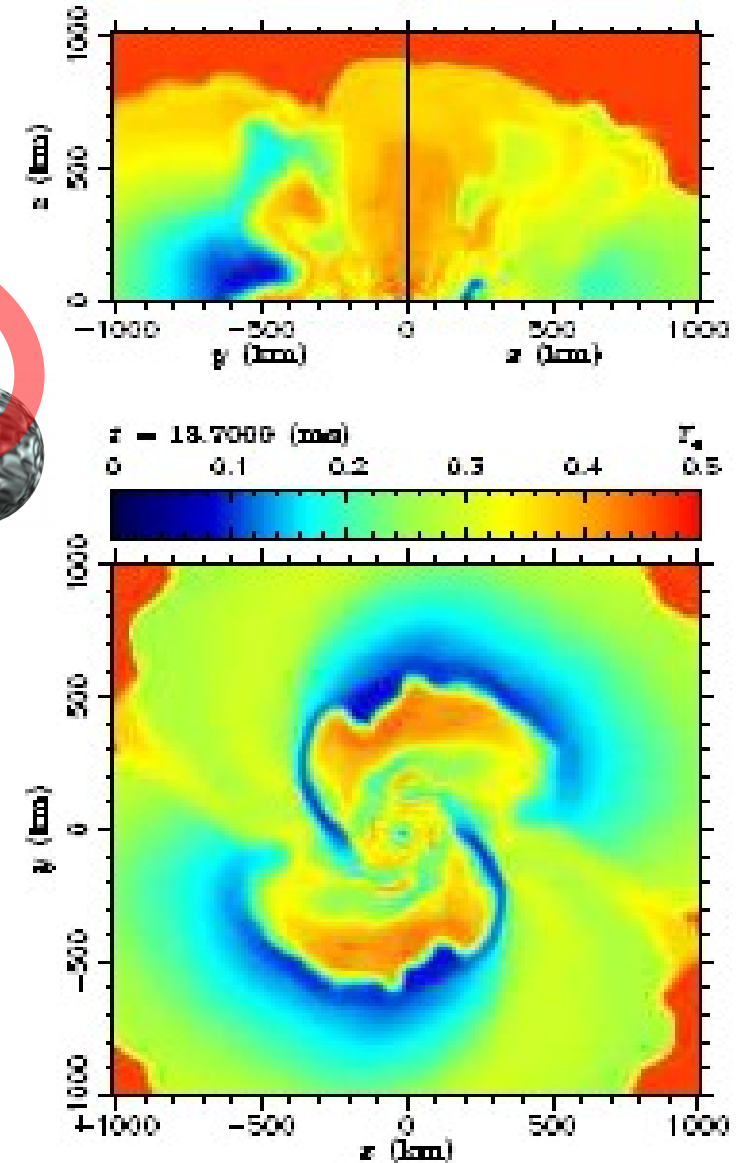
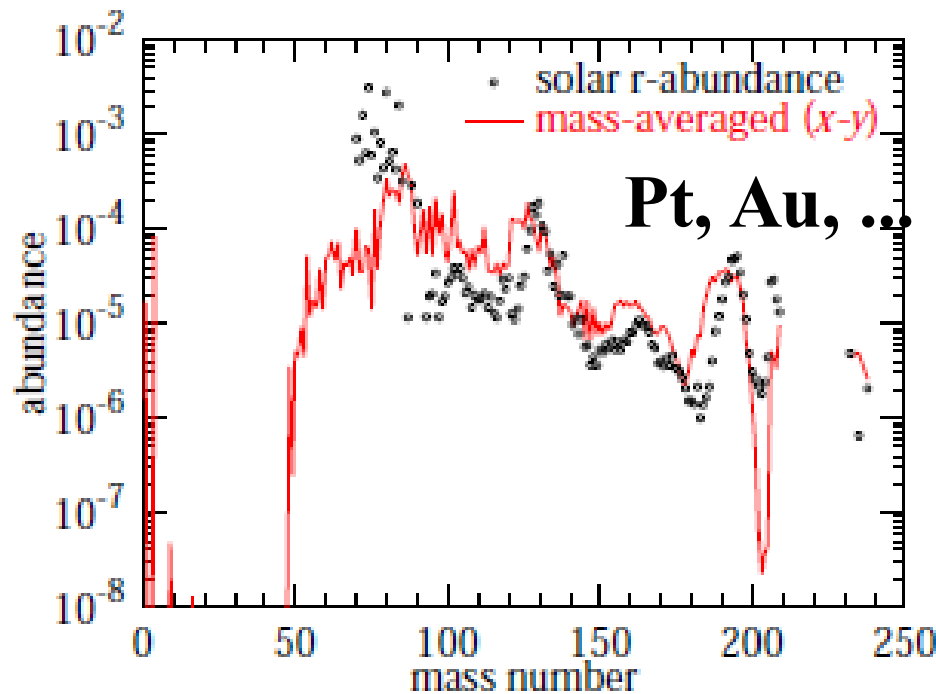
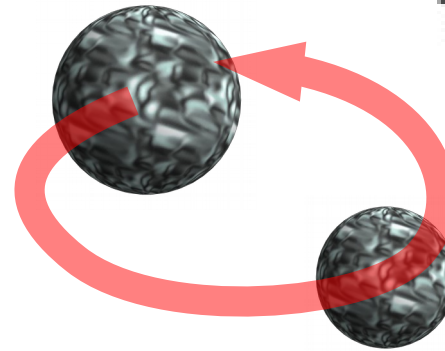
Age [yrs] *Page et al., 2011*

Takatsuka

Binary Neutron Star Mergers and Nucleosynthesis

- New possibility of r-process nucleosynthesis

- Element ratio from binary NS merger is found to reproduce Solar abundance.



Wanajo, Sekiguchi ('14)

Origin of Strong Magnetic Field

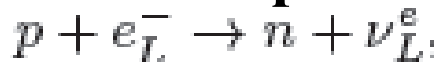
■ Magnetic field in NS $B = 10^{12} - 10^{15}$ G

- How can we make strong B ?
- How can we keep strong B ?
- Fossil, Dynamo, Ferromagnetism, ...

■ A new idea: Chiral Plasma Instability

AO, N. Yamamoto, arXiv:1402.4760

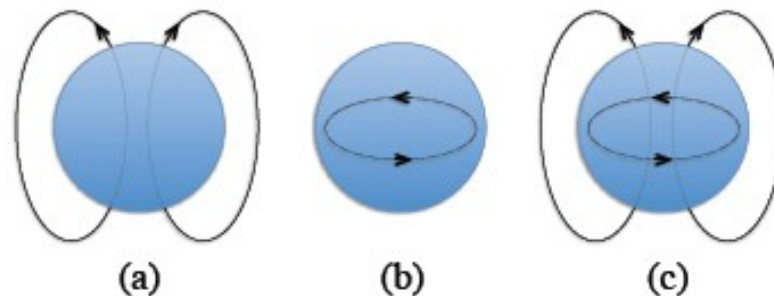
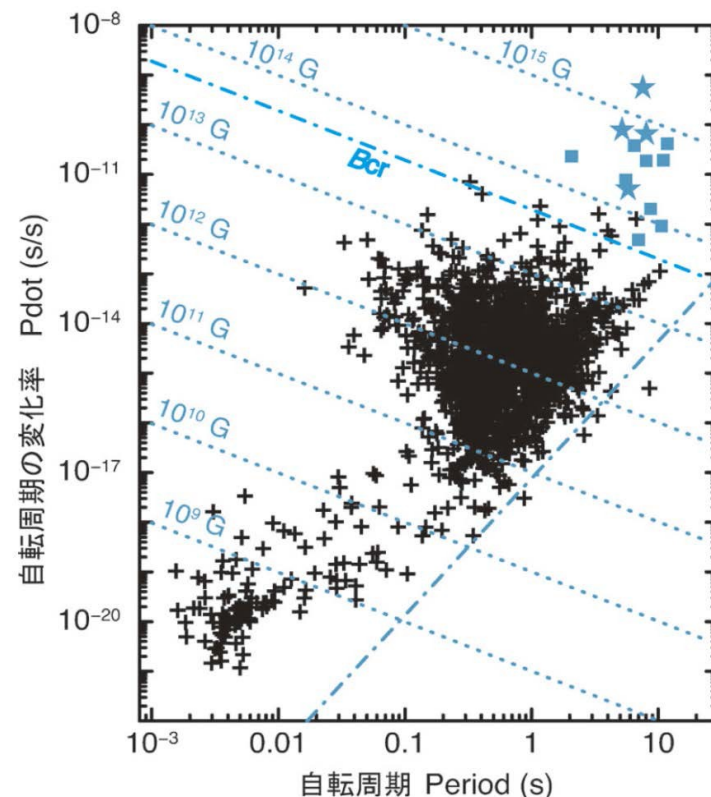
- Left-handed electrons are eaten in electron capture \rightarrow chiral chem. pot.



- Chiral plasma instability: N_5 is converted to magnetic helicity

$$j_z = \frac{2\alpha}{\pi} \mu_5 B_z, \quad \frac{d}{dt} \left(N_5 + \frac{\alpha}{\pi} \mathcal{H} \right) = 0, \quad N_5 = \int dx n_5$$

- Finite magnetic helicity makes magnetic field stable. $\mathcal{H} = \int dx A \cdot B$



See also, D. Grabowska, D. B. Kaplan, S. Reddy, PRD('15)085035

NS matter Grant-in-Aid Study in Japan(2012-)

High ρ (Group A)
head: Tamura, Takahashi

Hypernuclei, Kaonic nuclei
YN & YY int.,
Eff. Interaction
(Heavy-ion collisions)



PI: H. Tamura

Hyperons, mesons, quarks

Asym. nuclear matter
+elec.+ μ

Nuclei+neutron gas+elec.

Nuclei + elec.

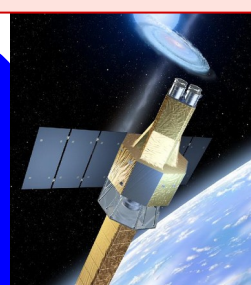
Low ρ (Group B)
head: Murakami,
Nakamura, Horikoshi

Sym. E, Pairing gap,
BEC-BEC cross over,
Cold atom, Unitary gas

NS Obs. (Group C)
head: Takahashi

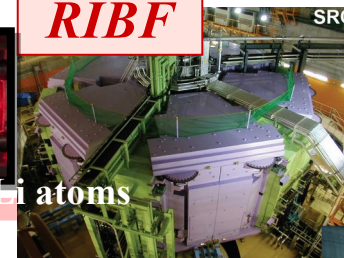
Radius, Mass,
Temp. (Cooling),
Star quake, Pasta

ASTRO-H



Theory (Group D)
head: Ohnishi

RIBF



US: UNEDF, ICNT, FRIB, RHIC, NICER...

Europe: CompStar, EMMI, FAIR, GANIL, LOFT, ...



Accelerators and Satellites for Neutron Star Physics

GANIL

FAIR

NICER

LOFT

J-PARC

RHIC

FRIB

LHC

ASTRO-H

Neutron Star Matter

RIBF

Contents

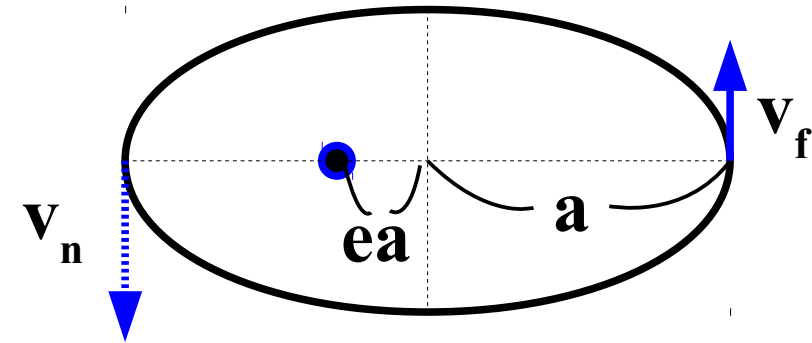
- **Introduction**
- **Part I : Basics of Neutron Star Physics**
 - **Neutron star mass & radius observations**
 - **Nuclear matter EOS and neutron stars**
 - **Massive Neutron Star Puzzle**
- **Part II : Hypernuclear Physics & Neutron Stars**
 - **Hypernuclear Physics : Implications from Experiments**
 - **What is necessary to solve massive NS puzzle ?**
 - **Recent Attempts toward the massive NS puzzle**
- **Summary**

Mass & Radius Measurements of Neutron Stars

Neutron Star Observables: Mass (1)

■ Please remember Kepler motion basics

- major axis= a , eccentricity= e ,
reduced mass= m , total mass= M



$$E/m = \frac{1}{2} v_f^2 - \frac{GM}{a(1+e)} = \frac{1}{2} v_n^2 - \frac{GM}{a(1-e)}$$

$$L = m v_f a(1+e) = m v_n a(1-e)$$

$$\rightarrow v_f^2 = \frac{GM}{a} \frac{1-e}{1+e}, L = 2m \frac{dS}{dt} = m \sqrt{GMa(1-e^2)}$$

$$\rightarrow P = S / (dS/dt) = 2\pi a^2 \sqrt{1-e^2} / \sqrt{GMa(1-e^2)} = 2\pi a^{3/2} / \sqrt{GM}$$

Neutron Star Observables: Mass (2)

Binary stars

- inclination angle = i
- Doppler shift (Pulse timing change) is given by the radial velocity (視線速度)

$$K = v \sin i$$

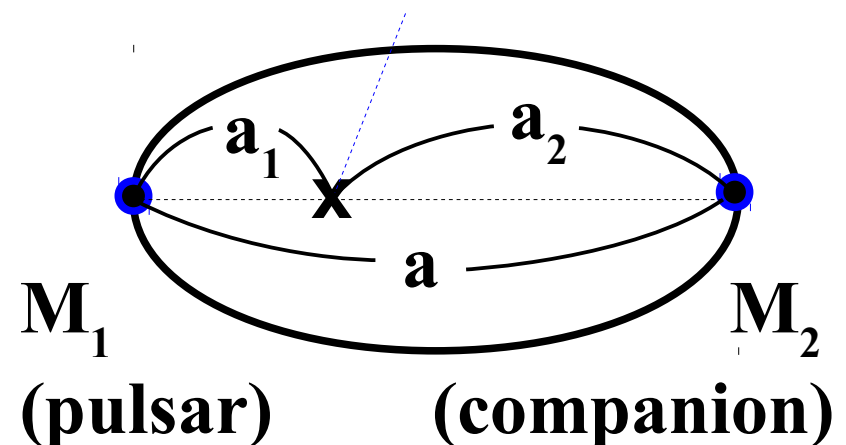
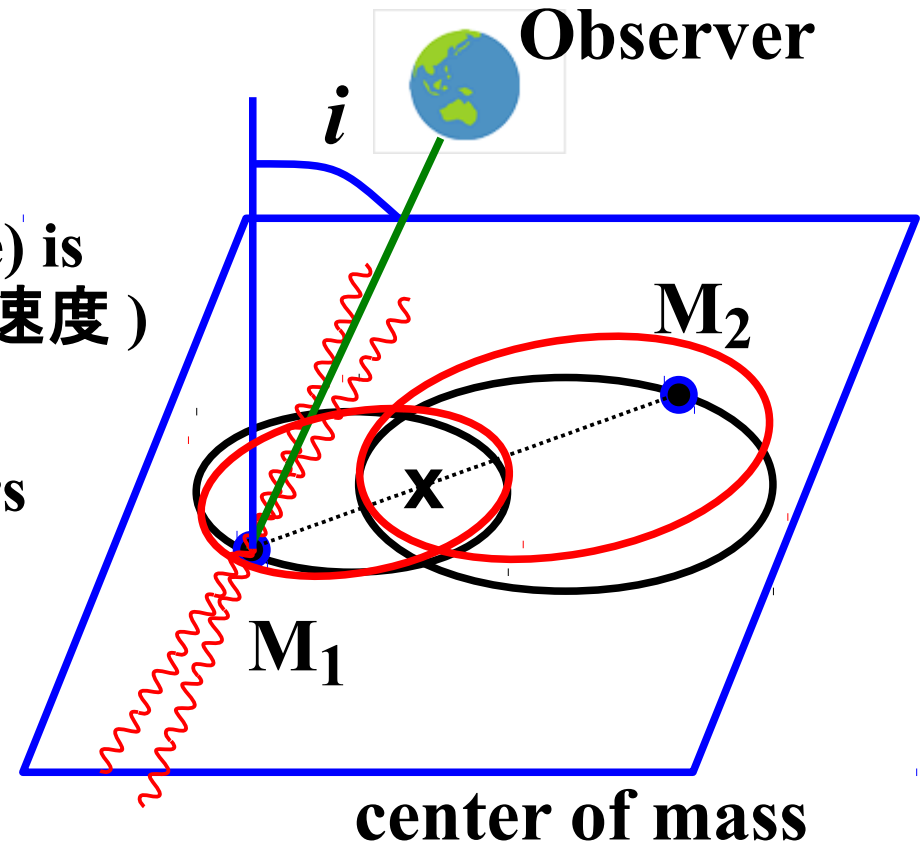
- Radial velocity \rightarrow orbit parameters
- Mass function (observable)

$$f \equiv \frac{(M_2 \sin i)^3}{M^2} = \frac{4\pi^2 (a_1 \sin i)^3}{G} P^{-2}$$

$$= \frac{K^3 P (1 - e^2)^{3/2}}{2\pi G}$$

$$(K = v \sin i, M = M_1 + M_2)$$

- and GR effects ...



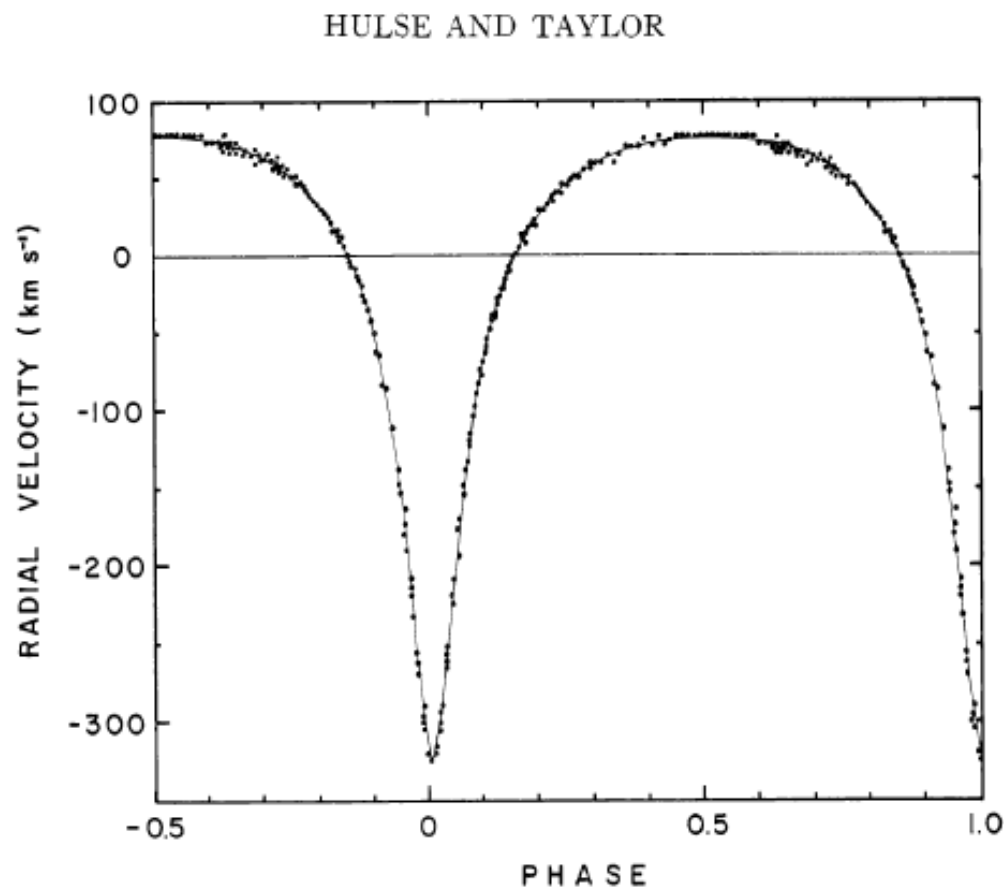
Hulse-Taylor Pulsar (PSR 1913+16)

- Precisely (and firstly) measured neutron star binary (1993 Nobel prize to Hulse & Taylor)
- Radial velocity \rightarrow $P, e, a_1 \sin i \rightarrow$ Mass function

TABLE 2

ELEMENTS OF THE ORBIT

$K_1 = 199 \pm 5 \text{ km s}^{-1}$
$P_b = 27908 \pm 7 \text{ s}$
$e = 0.615 \pm 0.010$
$\omega = 179^\circ \pm 1^\circ$
$T = \text{JD } 2,442,321.433 \pm 0.002$
$a_1 \sin i = 1.00 \pm 0.02 R_\odot$
$f(m) = 0.13 \pm 0.01 M_\odot$



Hulse-Taylor ('75)

More on Hulse-Taylor Pulsar (PSR 1913+16)

General Relativistic Effects

Perihelion shift (近日点移動)

$$\dot{\omega} = 3 \left(\frac{2\pi}{P} \right)^{5/3} \frac{(GM)^{2/3}}{(1-e^2)c^2}$$

Einstein delay

$$\Delta_E = \gamma \sin u$$

(u = eccentric anomaly)

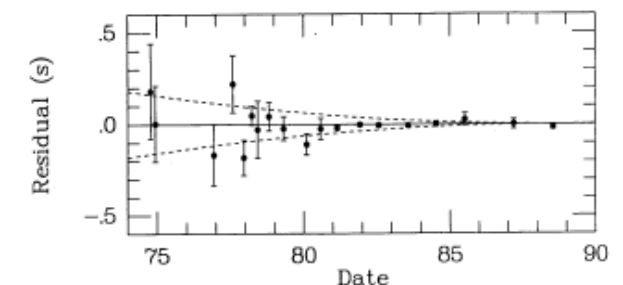
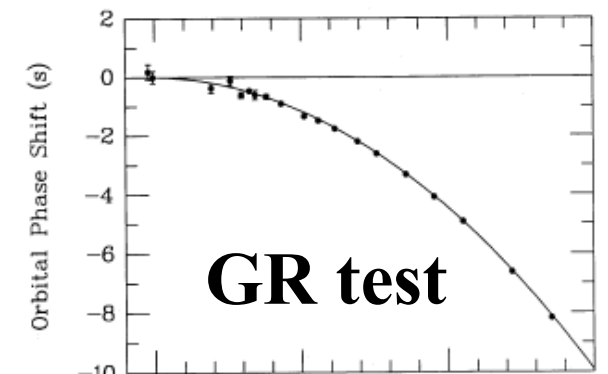
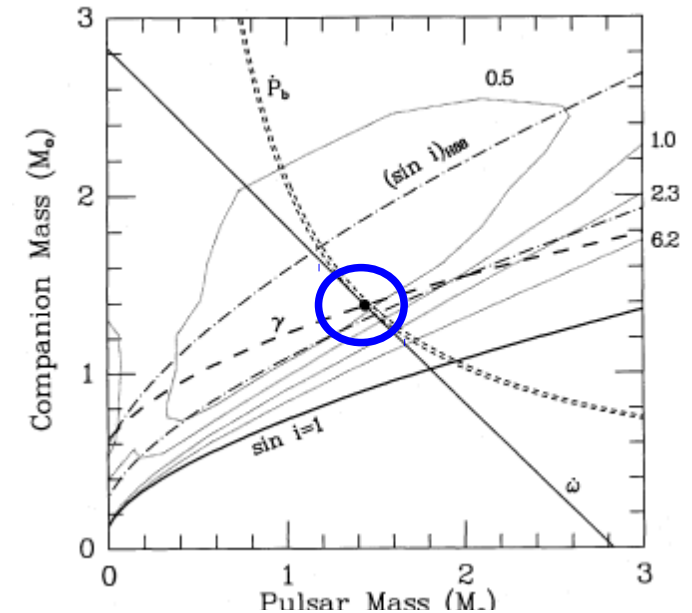
$$\gamma = \frac{eP_b G m_2 (m_1 + 2m_2)}{2\pi c^2 a_R M} \quad \frac{a_R^3}{P_b^2} = \frac{GM}{4\pi^2} \left[1 + \left(\frac{m_1 m_2}{M^2} - 9 \right) \frac{GM}{2a_R c^2} \right]^2$$

Two observable

→ Precise measurement of m_1 and m_2 .

$$m_1 = 1.442 \pm 0.003 M_{\text{sun}}$$

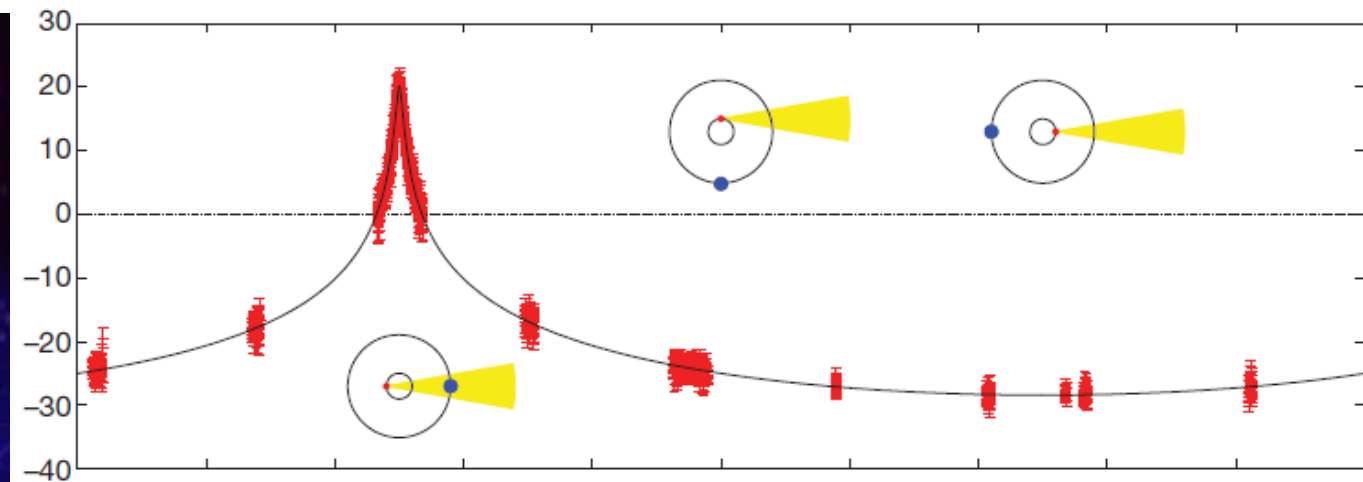
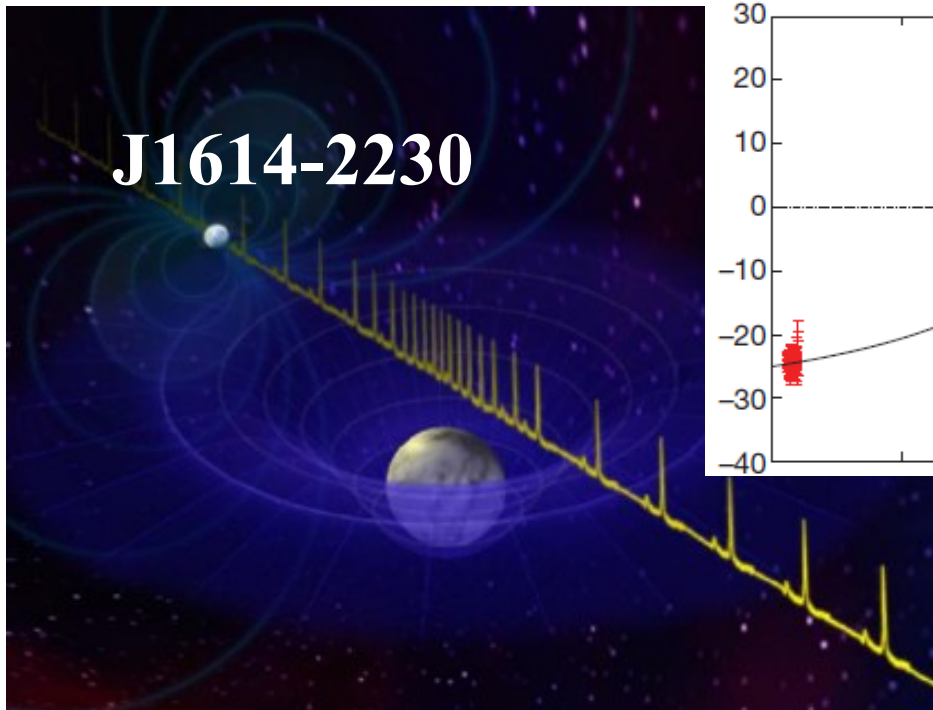
$$m_2 = 1.386 \pm 0.003 M_{\text{sun}}$$



Taylor, Weisenberg ('89)

Massive Neutron Star

- General Relativity Effects on Time Delay
 - Einstein delay : varying grav. red shift
 - Shapiro delay : companion's grav. field
- A massive neutron star (J1614-2230)
 - $M = 1.97 \pm 0.04 M_{\odot}$ is obtained using the Shapiro delay
Demorest et al. (2010)



$$\Delta_S = -2m \left[\ln \frac{r}{a} + \ln (1 - \sin i \sin \phi) \right]$$

Demorest et al., Nature 467 (2010) 1081.

Neutron Star Masses

- NS masses in NS binaries can be measured precisely by using some of GR effects.

- Perihelion shift+Einstein delay

$$\rightarrow M = 1.442 \pm 0.003 M_{\odot}$$

(Hulse-Taylor pulsar)

Taylor, Weisenberg ('89)

- Shapiro delay

$$\rightarrow M = 1.97 \pm 0.04 M_{\odot}$$

Demorest et al. ('10)

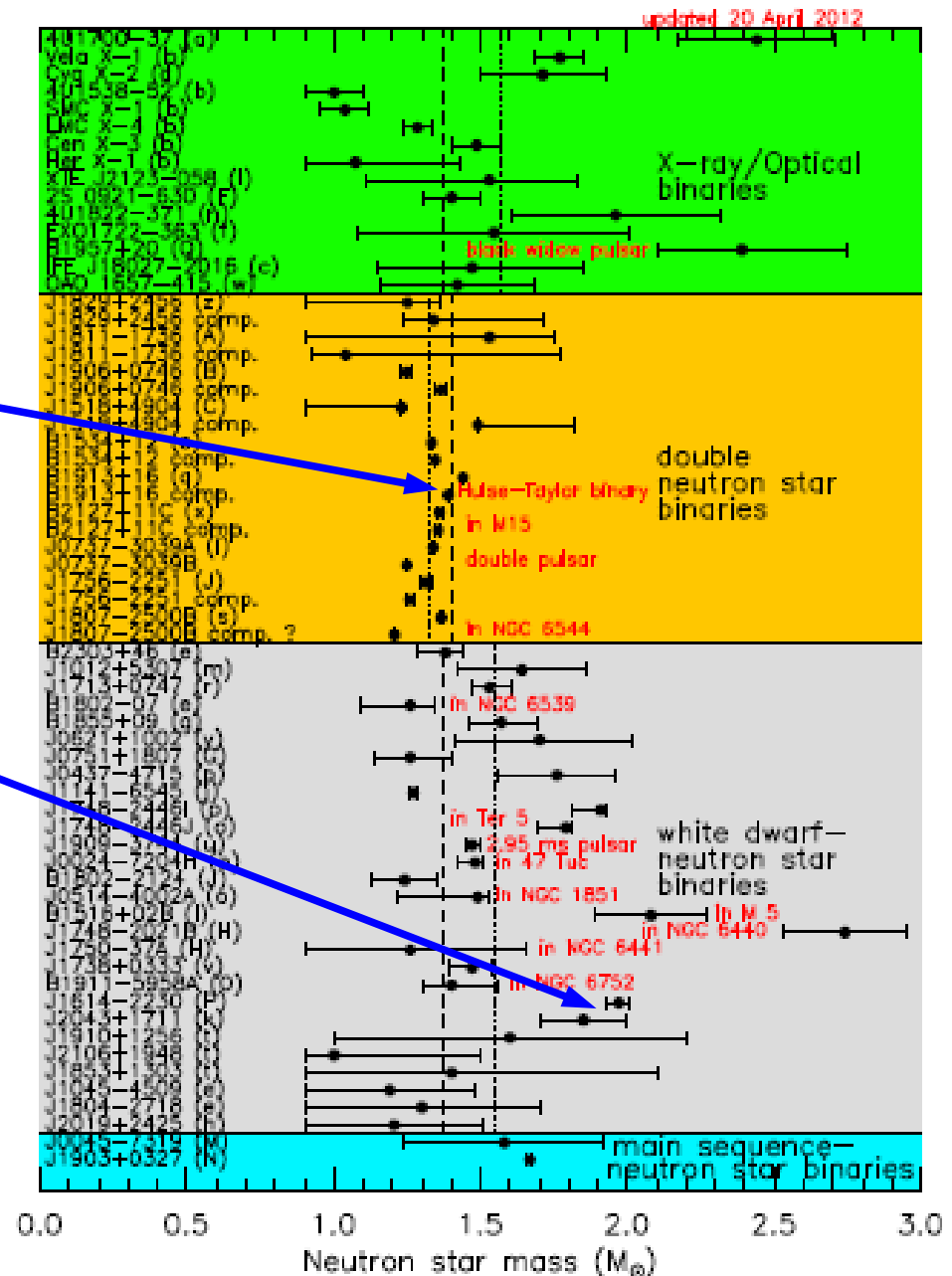
- Another obs.: $M = 2.01 \pm 0.04 M_{\odot}$

Antoniadis et al. ('13)

Neutron Star Mass

$$M = (1-2) M_{\odot}$$

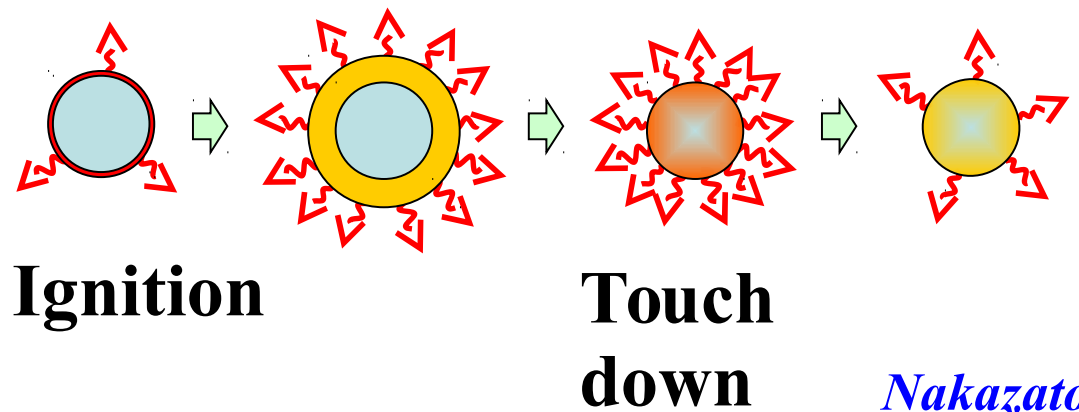
Canonical value = 1.4 M_{\odot}



Lattimer (2013)

Neutron Star Radius

- How can we measure 10 km radius of a star with 10-100 thousands light year distance from us ?
 - Size of galaxy $\sim 3 \times 10^{14}$ km (~ 10 kpc $\sim 3 \times 10^4$ light year)
→ Model analysis is necessary !
- X-ray burster
 - Mass accretion from companion occasionally induces explosive hydrogen / helium burning.
 - High temperature → NS becomes bright !
 - Three methods to measure NS radius



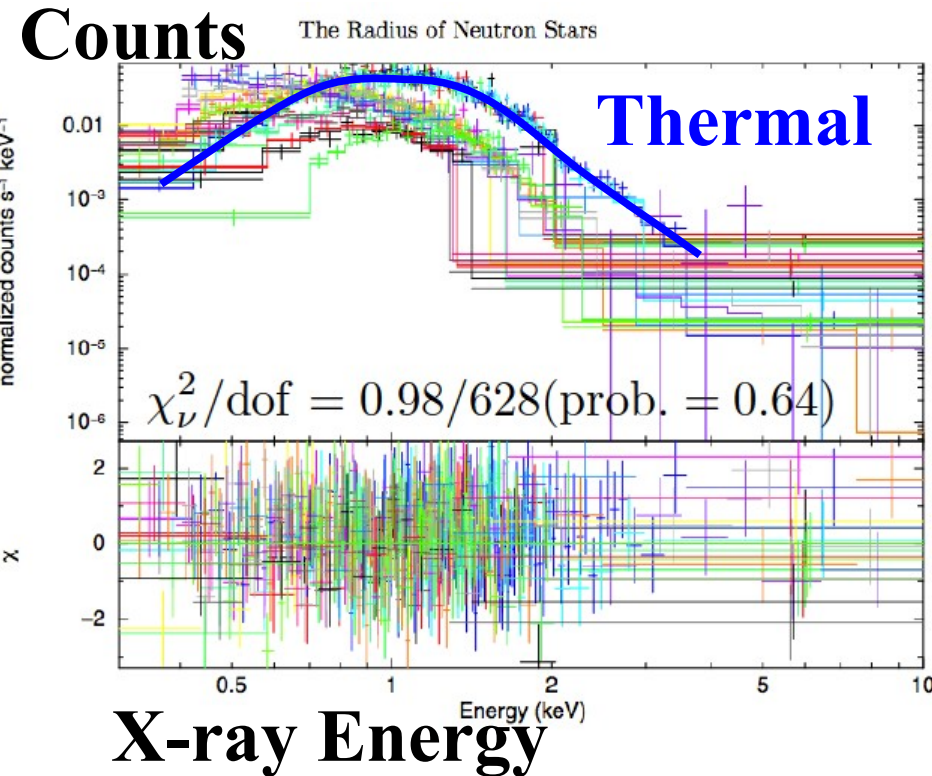
NS Radius Measurement (1)

■ Surface emission

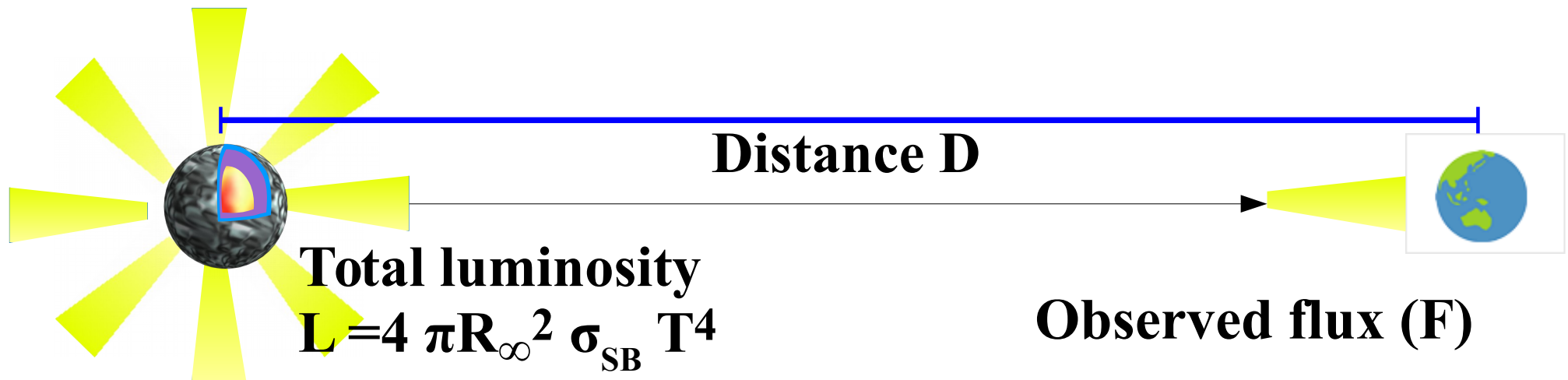
- Stefan-Boltzmann law is assumed
→ NS radius is obtained from Flux, Temperature, and Distance measurement.

$$L = 4 \pi R_{\infty}^2 \sigma_{\text{SB}} T^4, \quad F = \frac{L}{4 \pi D^2}$$

$$\rightarrow R = \sqrt{\frac{F D^2}{\sigma_{\text{SB}} T^4} \left(1 - \frac{2 G M}{R c^2} \right)^{-1/2}}$$



Guillot et al. (2013)

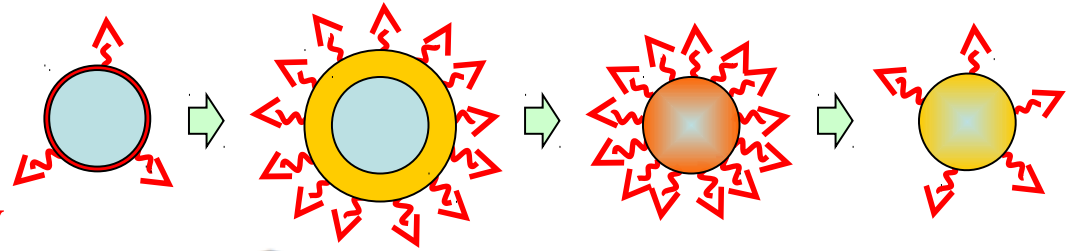


NS Radius Measurement (2)

Eddington Limit

Eddington Limit

radiation pressure = gravity



$$\frac{4\pi r^2 \sigma_{\text{SB}} T^4}{4\pi r^2 c} \cdot N_e \cdot \sigma_{\text{T}}$$

$$= \frac{GM}{r^2} \cdot N_N \cdot m_N$$

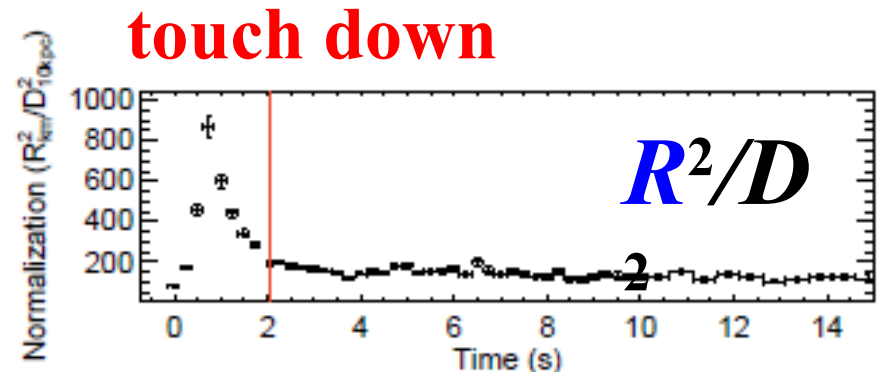
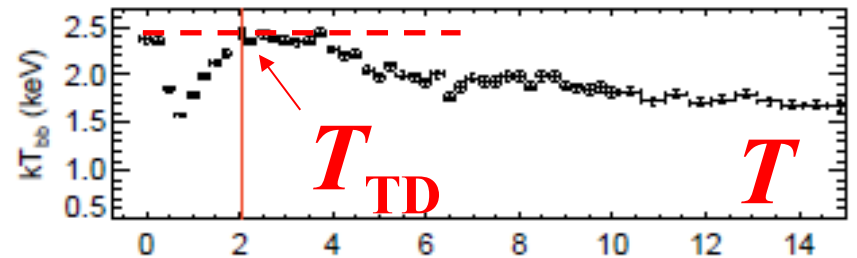
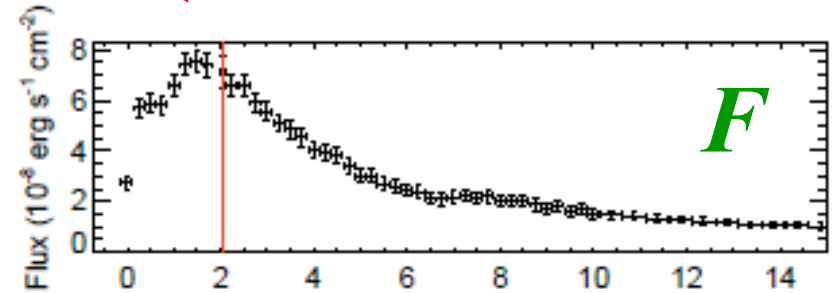
$$\rightarrow R_{\infty}^2 = \frac{2GMcm_N}{\sigma_{\text{T}}\sigma_{\text{SB}}T^4} \frac{N_N}{N_e}$$

Eddington limit is assumed to be achieved at “touch down”.

Electron-nucleon ratio

$$N_e/N_N = (1+X)/2$$

(X=1 for hydrogen atmosphere
X=0 for light elements)



Guver et al., ApJ 747 (2012) 47

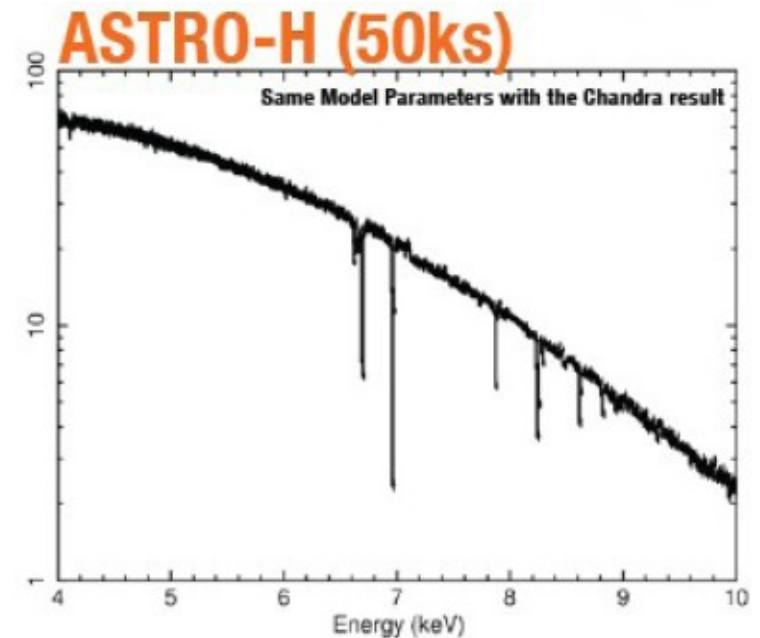
NS Radius Measurement (3)

Red Shift

- Neutron Star surface is expected to contain Irons.
- Absorption lines should be red shifted.
→ Almost direct observation of M/R.

$$E_{\text{obs}} = E_{\text{surf}} \sqrt{1 - \frac{2GM}{Rc^2}}$$

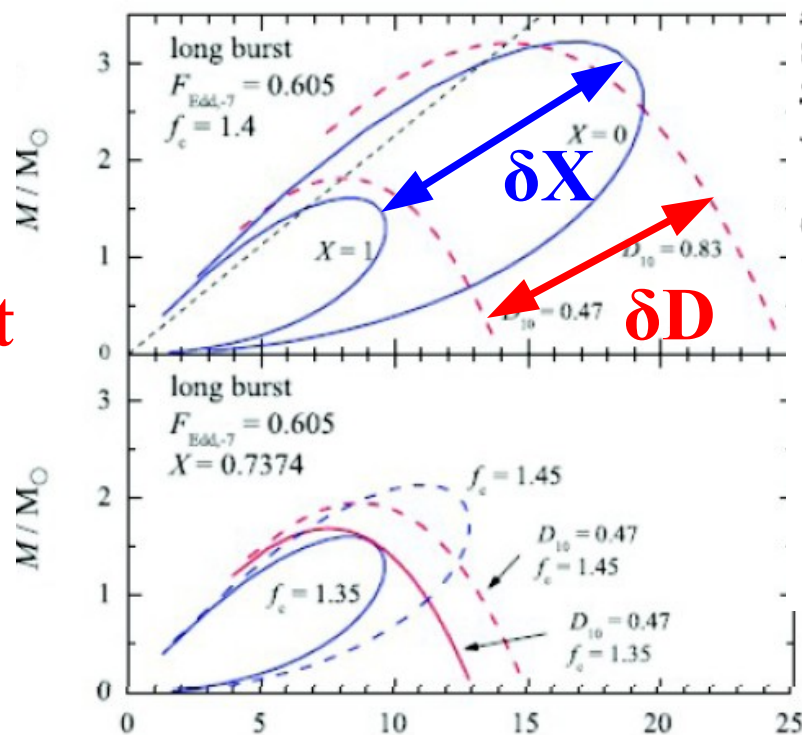
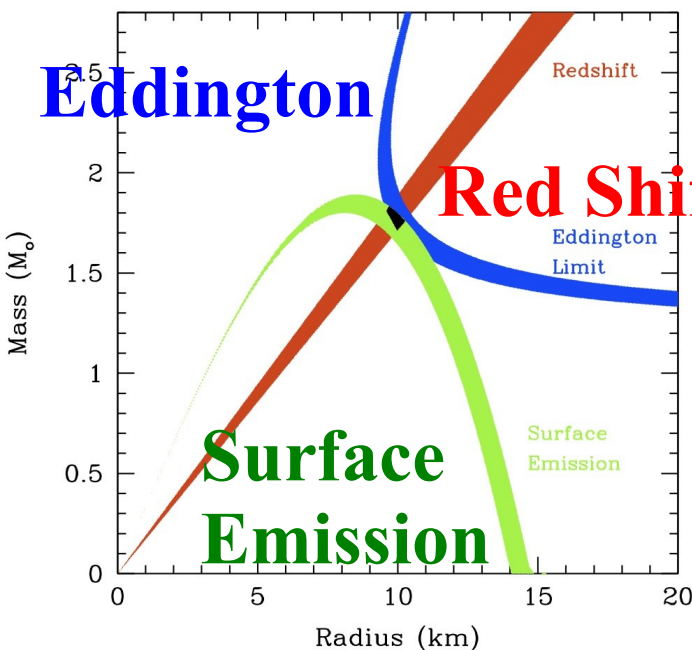
- ASTRO-H will measure Iron absorption line from NS, and determine M/R with 1 % accuracy !



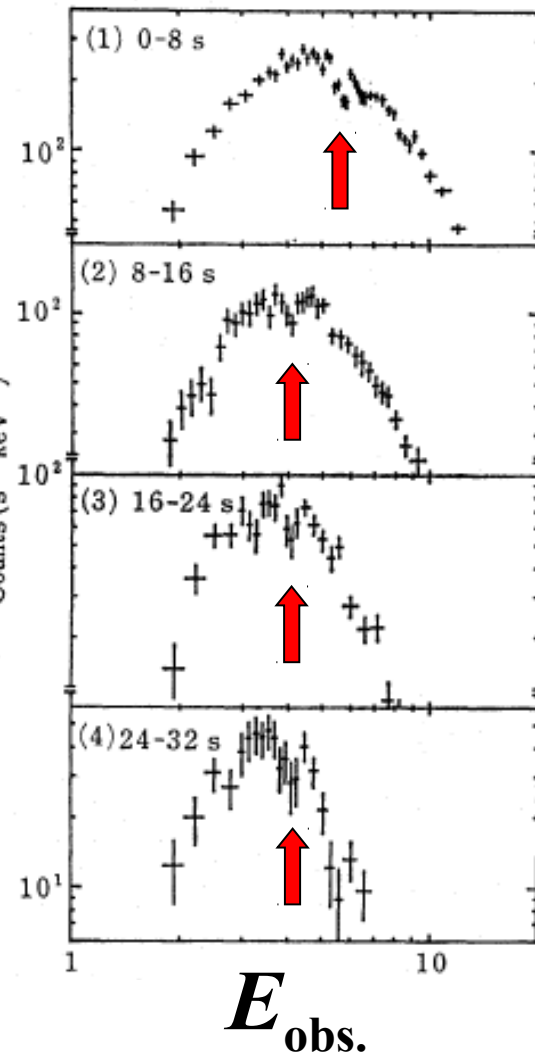
ASTRO-H simulation

Neutron Star Radius

- Do three methods give consistent (M, R) ?
 - Surface emission & Eddington limit have large error bars from Distance & Composition uncertainty.
 - Red shift of discrete lines have not been observed unambiguously.



4U 1724-307, Suleimanov et al., *ApJ*742('11),122

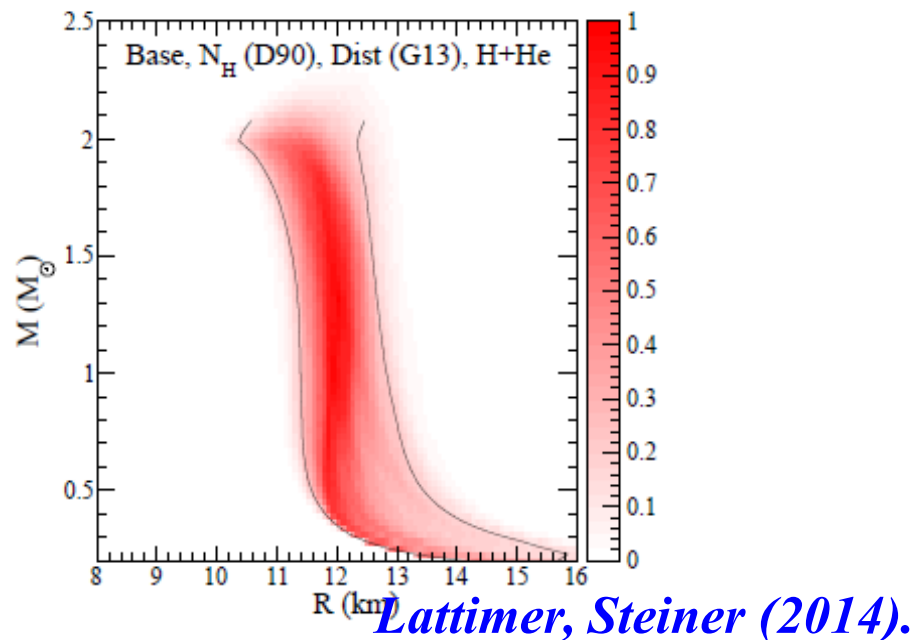
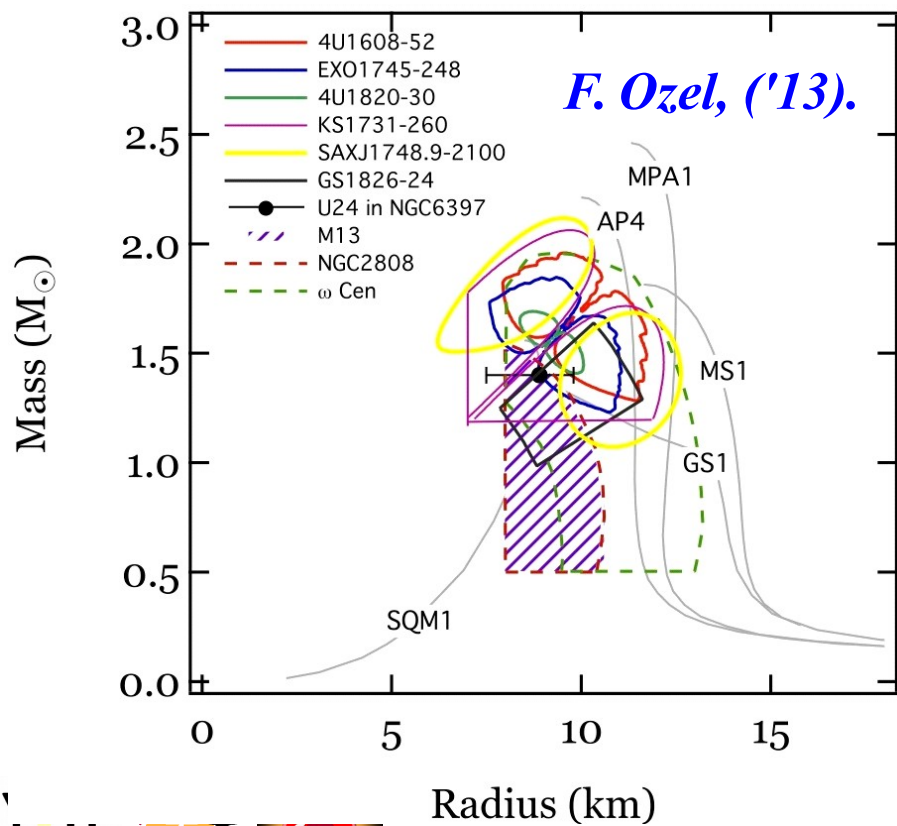
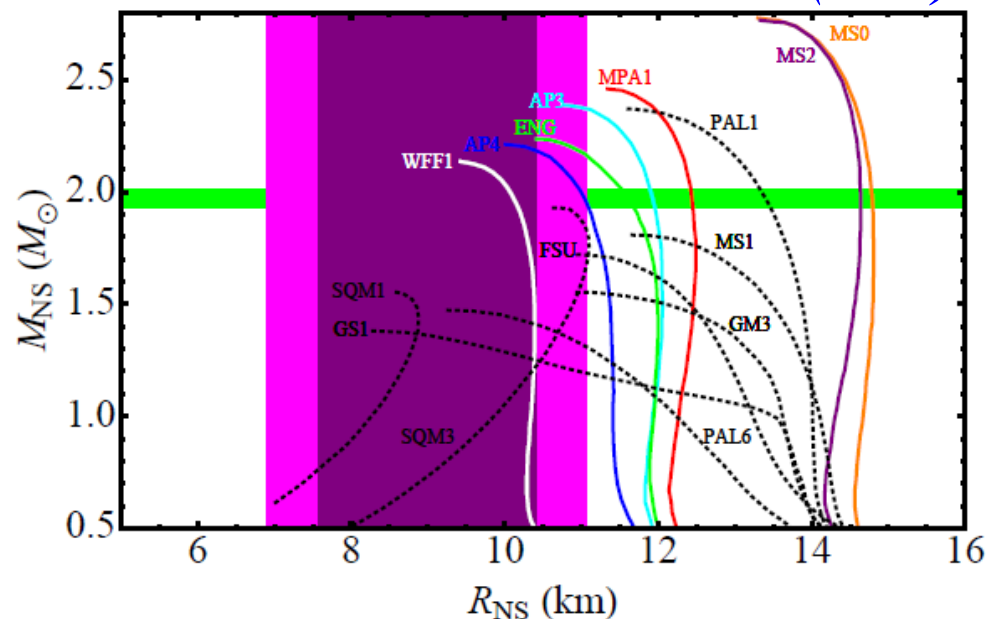


Waki et al., *PASJ*36('84)819

Compact NS puzzle

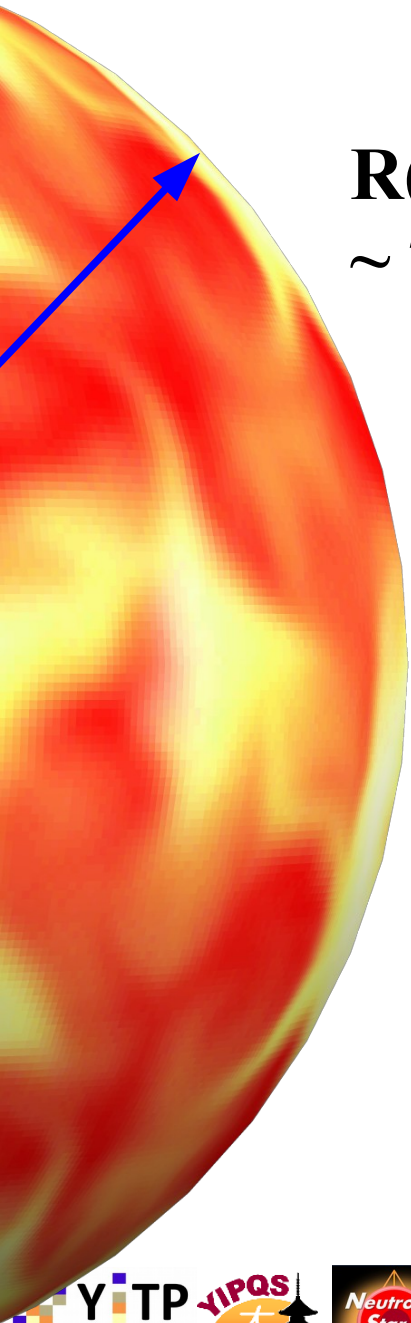
- Some analyses suggest smaller R_{NS} than nucl. phys. predictions.
- Some make objections.
 Suleimanov+, $R_{1.4} > 13.9$ km
 Lattimer+, $R_{1.4} = 12 \pm 1.4$ km

Guillot et al. (2013)

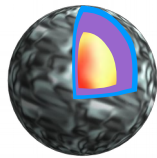


Lattimer, Steiner (2014).

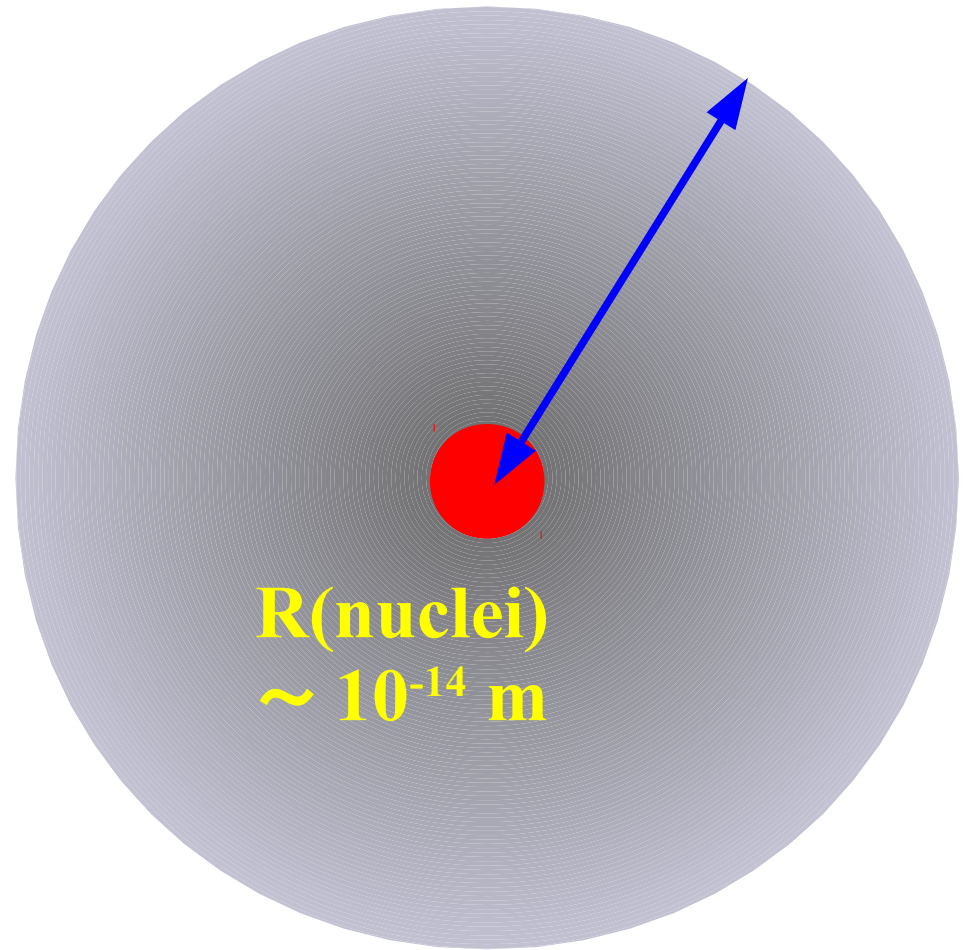
Neutron Star Density



R(Sun)
 $\sim 700,000$ km



R(NS) ~ 10 km
M(NS) $\sim 1.4 M_{\odot}$



R(atom)
 $\sim 10^{-10}$ m

R(nuclei)
 $\sim 10^{-14}$ m

Very High Density !

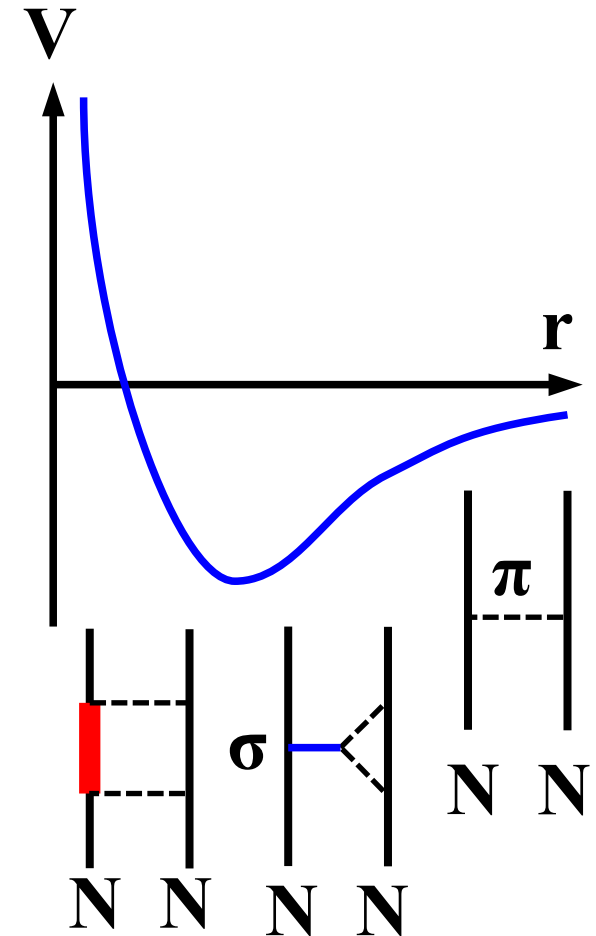
$$m_N \rho(\text{NS}) \sim (2-7) \times 10^{14} \text{ g / cm}^3 \sim (1-3) m_N \rho_0$$

Neutron Stars are supported by Nuclear Force !

- Average density of NS $\sim (1-3) \rho_0$, Max. density $\sim (5-10) \rho_0$
→ Supported by Nuclear Force
c.f. White Dwarfs are supported by electron pressure.

- Nuclear Force

- Long-range part: π exchange
Yukawa (1935)
- Medium-range attraction:
2 π exchange, σ exchange,
- Short-range repulsion:
Vector meson exchange,
Pauli blocking btw. quarks
Gluon exchange
Tamagaki; Oka, Yazaki;
Aoki, Hatsuda, Ishii



Neutron Star Matter EOS

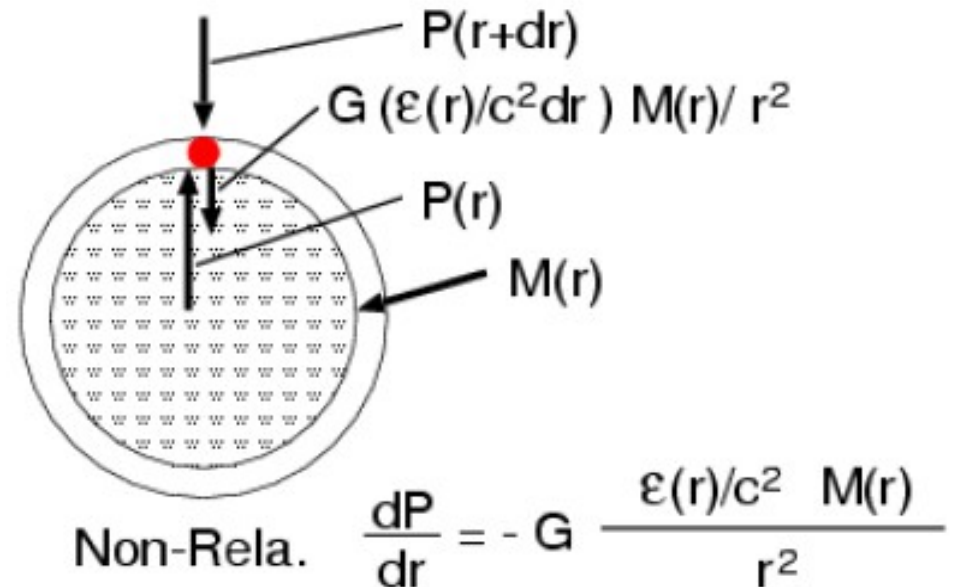
TOV equation

- **General Relativistic Hydrostatic Equation**
= **TOV(Tolman-Oppenheimer-Volkoff) equation**

$$\frac{dP}{dr} = -G \frac{(\epsilon/c^2 + P/c^2)(M + 4\pi r^3 P/c^2)}{r^2(1 - 2GM/rc^2)}$$

$$\frac{dM}{dr} = 4\pi r^2 \epsilon/c^2, \quad P = P(\epsilon) \quad (\text{EOS})$$

- **Spherical and non-rotating.**
- **3 Variables ($\epsilon(r)$, $P(r)$, $M(r)$),
3 Equations.**
- **Initial cond. $\epsilon(r=0)$
Solve TOV until $P=0$**

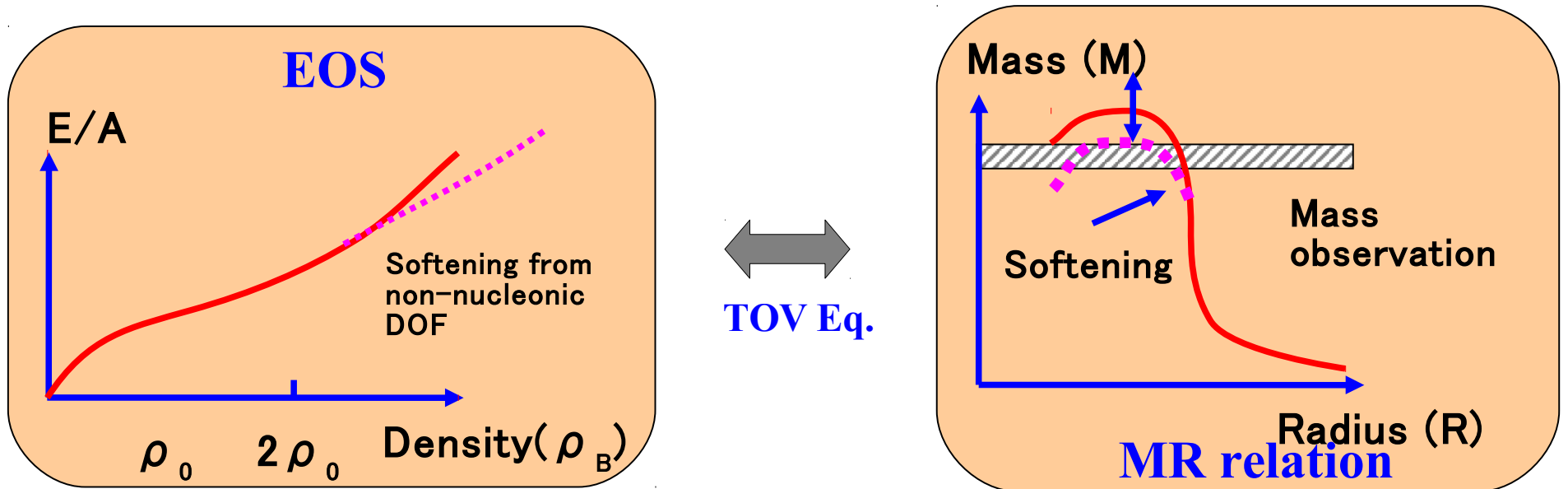


M-R Relation and EOS

- Solving TOV eq. starting from the “initial” condition, $\varepsilon(r=0) = \varepsilon_c = \text{given}$ until the “boundary” condition $P(r)=0$ is satisfied.
→ M and R are the functions of $\varepsilon(r=0)$ and functionals of EOS, $P=P(\varepsilon)$.

$$M = M(\varepsilon_c)[P(\varepsilon)] \quad , \quad R = R(\varepsilon_c)[P(\varepsilon)]$$

→ M-R curve and NS matter EOS : 1 to 1 correspondence



- Bethe-Weizsacker mass formula

Nuclear binding energy is roughly given by Liquid drop.

Nuclear size measurement $\rightarrow R = r_0 A^{1/3}$

$$B(A, Z) = \underbrace{a_v A}_{\text{Volume}} - \underbrace{a_s A^{2/3}}_{\text{Surface}} - \underbrace{a_C \frac{Z^2}{A^{1/3}}}_{\text{Coulomb}} - \underbrace{a_a \frac{(N - Z)^2}{A}}_{\text{Symmetry}} + \underbrace{a_p \frac{\delta_p}{A^\gamma}}_{\text{Pairing}}$$

Volume	Surface	Coulomb	Symmetry	Pairing
$A \propto \frac{4\pi}{3} R^3$	$A^{2/3} \propto 4\pi R^2$	$\propto \frac{Q^2}{R}$		

- Ignore Coulomb, consider $A \rightarrow \infty$,

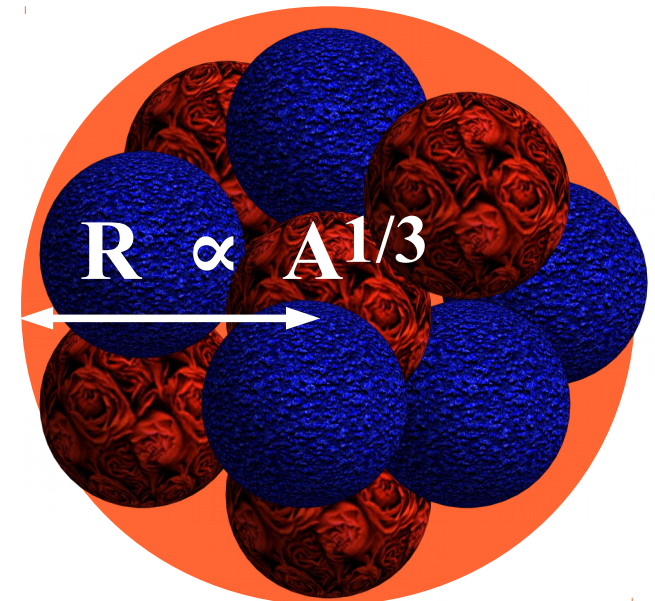
$$B/A = a_v(\rho) - a_a(\rho) \delta^2, \quad \delta = (N - Z)/A$$

$$a_v \simeq 16 \text{ MeV}$$

$$a_a \simeq 23 \text{ MeV} \quad (a_a(\text{vol}) \simeq 30 \text{ MeV})$$

Coef. may depend on the number density ρ

\rightarrow Nuclear Matter EOS



Neutron Star Matter EOS

Energy per nucleon in nuclear matter

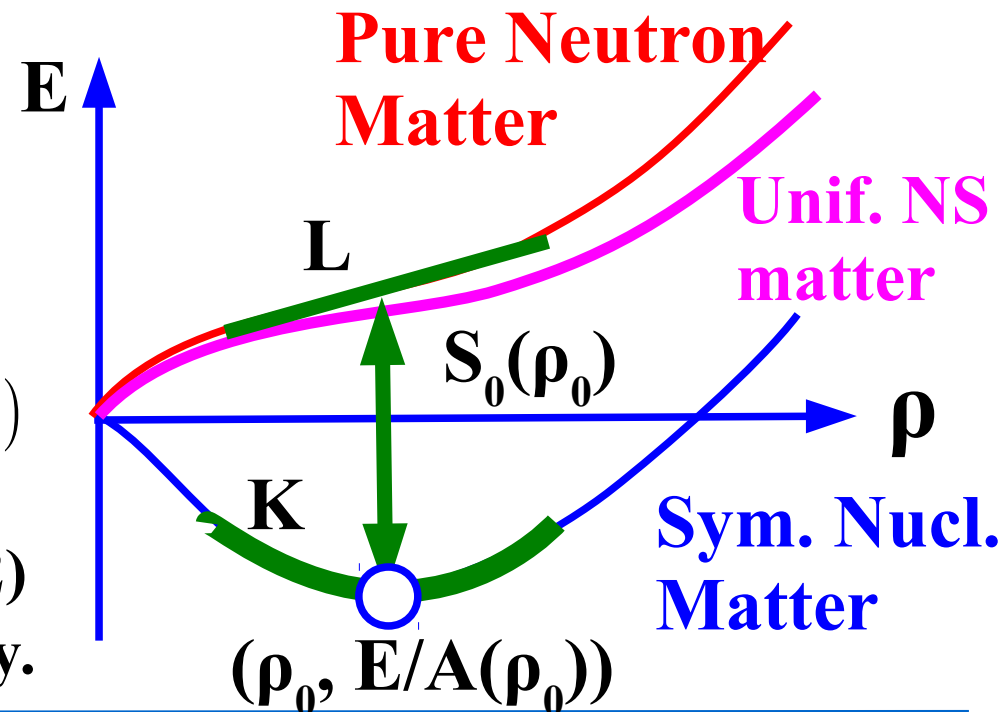
$$E_{\text{NM}}(\rho, \delta) = E_{\text{SNM}}(\rho) + S(\rho)\delta^2, \quad \delta = (N - Z)/A$$

$$E_{\text{SNM}}(\rho) \simeq E_0 + \frac{K(\rho - \rho_0)^2}{18\rho_0^2}, \quad S(\rho) = S_0 + \frac{L(\rho - \rho_0)}{3\rho_0}$$

- Saturation point $(\rho_0, E_0) \sim (0.16 \text{ fm}^{-3}, -16 \text{ MeV})$
- Symmetry energy parameters $(S_0 (=J), L) \sim (30 \text{ MeV}, 70 \text{ MeV})$
- Incompressibility $K \sim 230 \text{ MeV}$

Uniform neutron star matter

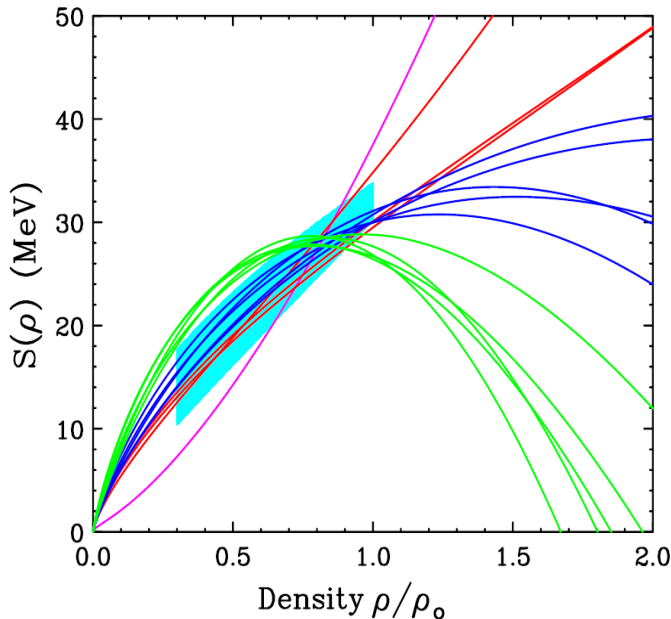
- Constituents at low density = proton, neutron and electron
- Charge neutrality $\rightarrow \rho(\text{elec.}) = \rho(p) \quad (\rho_e = \rho_p = \rho(1 - \delta)/2)$
 δ is optimized to minimize energy.



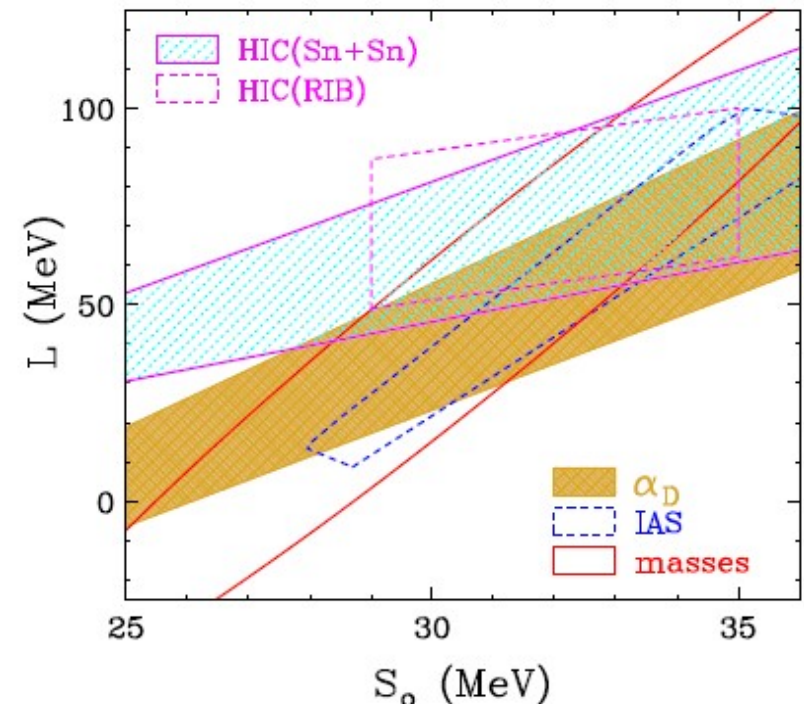
Symmetry Energy

- Symmetry Energy has been extracted from various observations.
 - Mass formula, Isobaric Analog State, Pygmy Dipole Resonance, Isospin Diffusion, Neutron Skin thickness, Dipole Polarizability, Asteroseismology

Recent recommended value
 $S_0 = 30-35 \text{ MeV}, L = 50-90 \text{ MeV}$
Is it enough for NS radii ?



M.B. Tsang et al.
(NuSYM2011),
PRC 86 ('12)015803.



C.J.Horowitz, E.F.Brown, Y.Kim,
W.G.Lynch, R.Michaels, A. Ono, J.
Piekarewicz, M. B. Tsang, H.H.Wolter
(NuSYM13), JPG41('14) 093001

Simple parametrized EOS

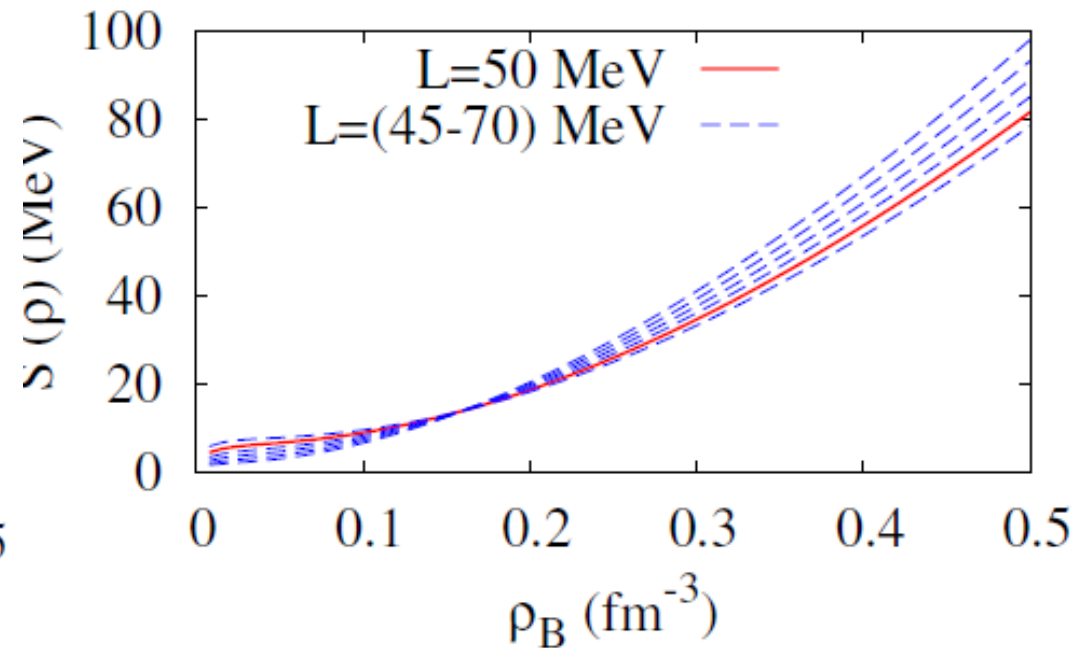
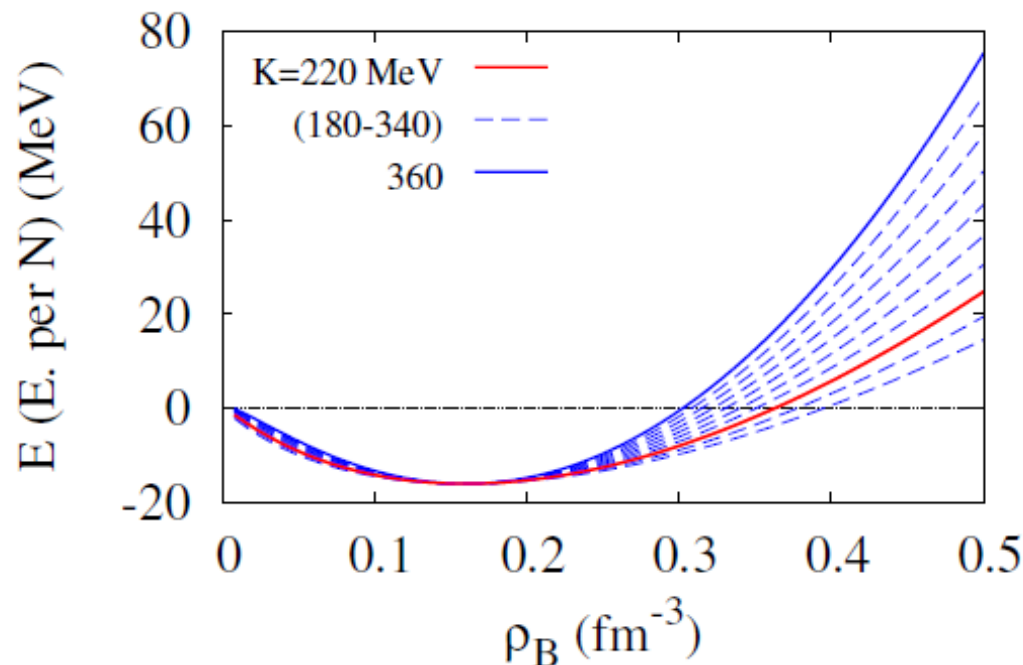
■ Skyrme int. motivated parameterization

$$E_{\text{SNM}} = \frac{3}{5} E_F(\rho) + \frac{\alpha}{2} \left(\frac{\rho}{\rho_0} \right) + \frac{\beta}{2 + \gamma} \left(\frac{\rho}{\rho_0} \right)^{1+\gamma}$$

$$S(\rho) = \frac{1}{3} E_F(\rho) + \left[S_0 - \frac{1}{3} E_F(\rho_0) \right] \left(\frac{\rho}{\rho_0} \right)^{\gamma_{\text{sym}}}$$

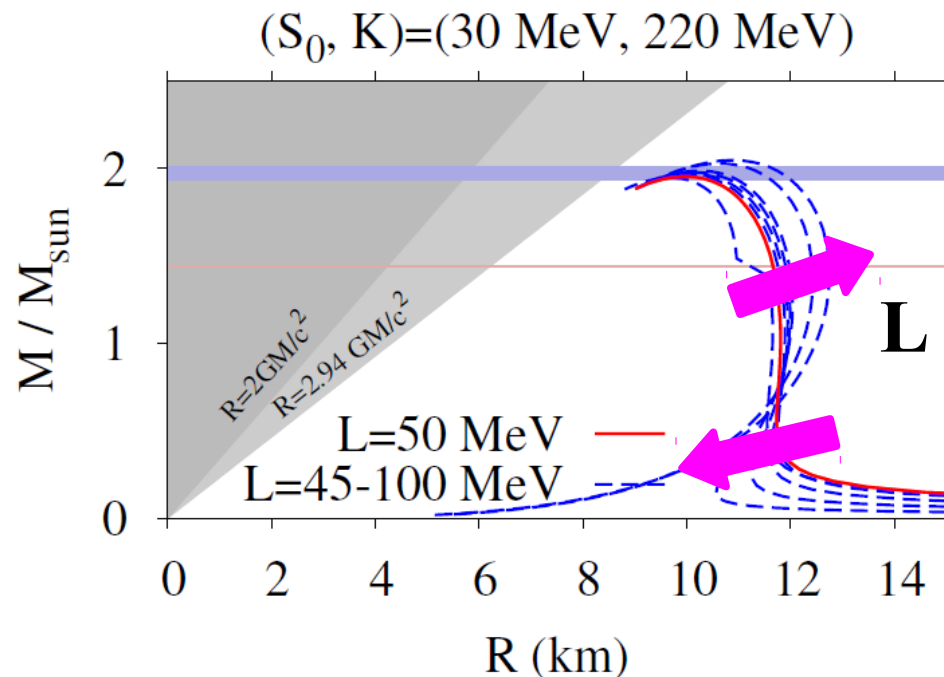
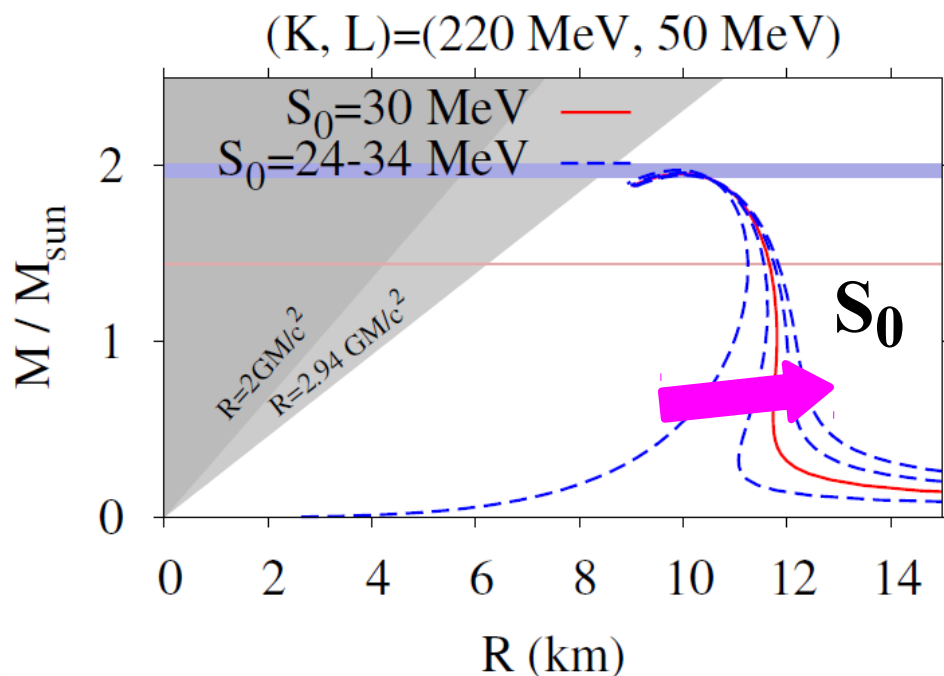
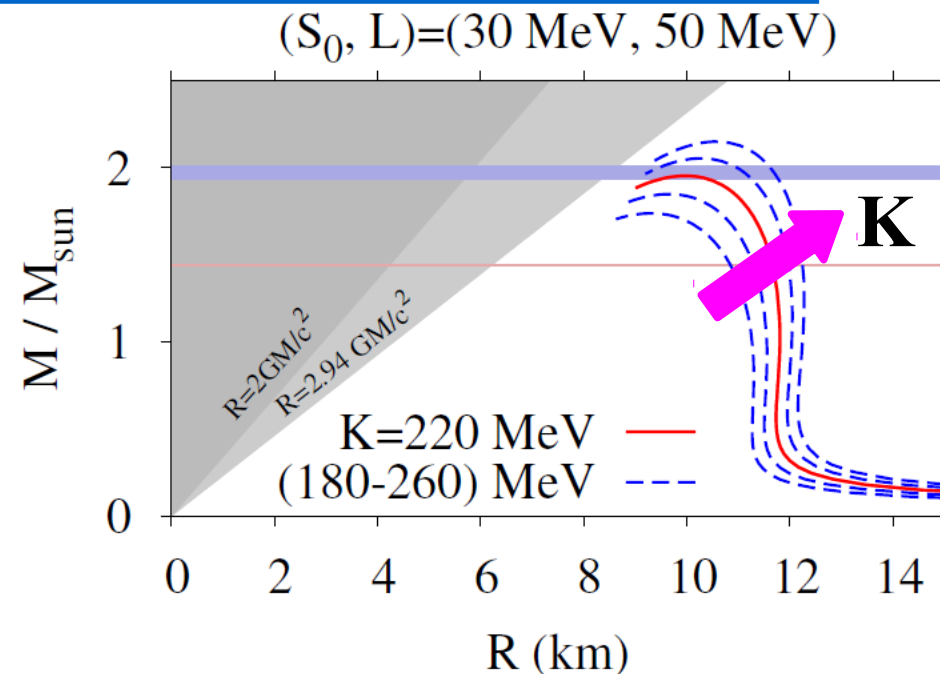
● $\rho_0, E/A(\rho_0), K \rightarrow \alpha, \beta, \gamma, L \rightarrow \gamma_{\text{sym}}$

$K=220 \text{ MeV}, S_0=30 \text{ MeV}$



Simple parametrized EOS

- Larger $K \rightarrow M \uparrow, R \uparrow$
- Larger $S_0 \rightarrow R \downarrow$ at small M
- Larger L
 $\rightarrow R \uparrow(\downarrow)$ at large (small) M



Theories/Models for Nuclear Matter EOS

■ Mean Field from Effective Int. ~ Nuclear Density Functionals

● Skyrme Hartree-Fock

- ◆ Non.-Rel., Zero Range, Two-body + Three-body (or ρ -dep. two-body)

$$\frac{E}{A} = \left\langle \frac{\mathbf{p}^2}{2m^*} \right\rangle + V(\rho, \delta), \quad V \simeq \frac{\alpha}{2} \frac{\rho}{\rho_0} + \frac{\alpha' \delta}{2} \frac{\rho}{\rho_0} + \frac{\beta}{1 + \gamma} \left(\frac{\rho}{\rho_0} \right)^\gamma + \dots$$

● Relativistic Mean Field

- ◆ Relativistic, Meson-Baryon coupling, Meson self-energies

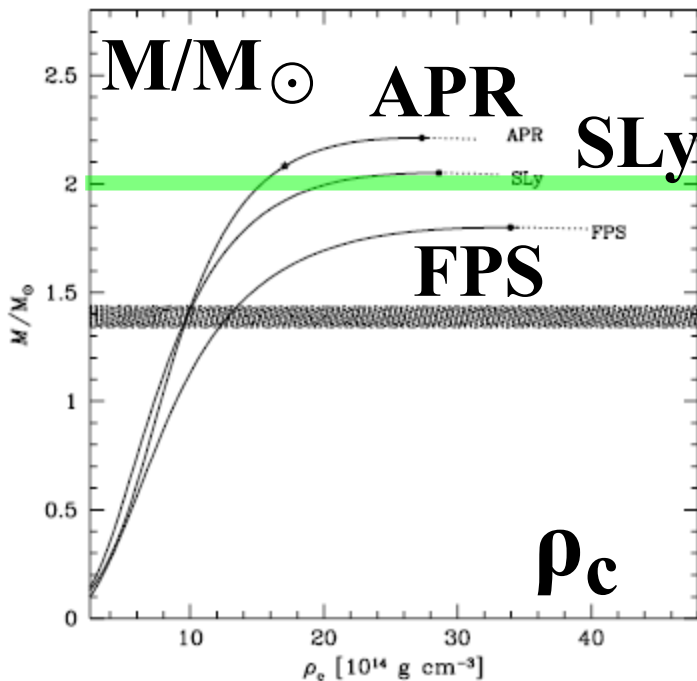
$$\frac{E}{A} = \left\langle \sqrt{\mathbf{p}^2 (M - g_\sigma \sigma)^2} \right\rangle + g_\omega \omega + \frac{1}{\rho_B} \left[\frac{1}{2} m_\sigma \sigma^2 - \frac{1}{2} m_\omega \omega^2 + \dots \right]$$

■ Microscopic (ab initio) Approaches (starting from bare NN int.)

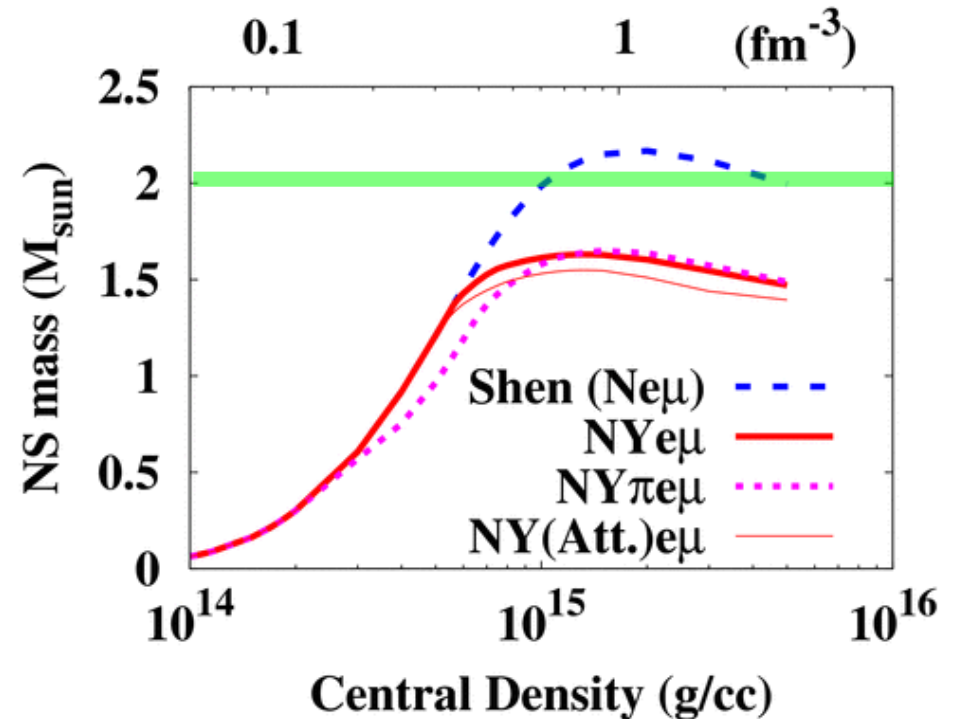
- Variational calculation
- Quantum Monte-Carlo
- Bruckner Theory (G-matrix)

Mean Field models

- Fit parameters to nuclear properties (B.E., radius, ...)
 - predict neutron star (M,R).
- Non-Rel. treatment with SLy (std. parametrization), FPS (impr.)
 - $M_{\max} \sim (1.8-2.0) M_{\odot}$
- Rel. MF (TM1) → $M_{\max} \sim 2.2 M_{\odot}$



F. Douchin, P. Haensel.
Astron. Astrophys. 380('01)151.



Ishizuka, AO, Tsubakihara, Sumiyoshi,
Yamada, J. Phys. G35(08),085201
c.f. H. Shen+('09) → n, p, Λ EOS

Variational Calculation

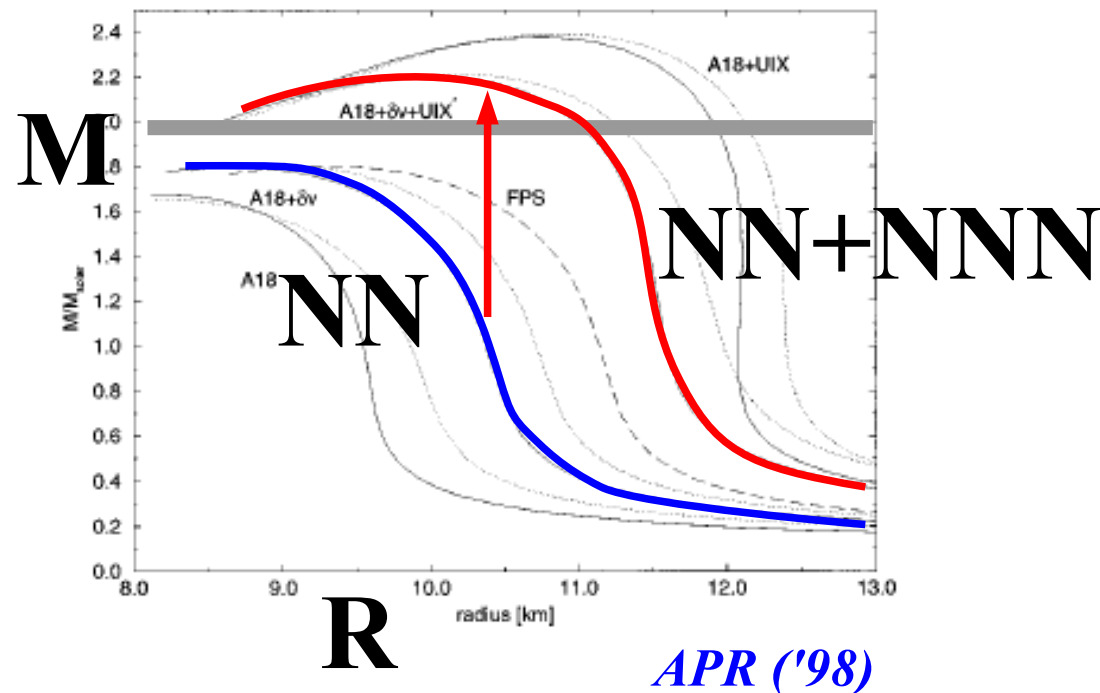
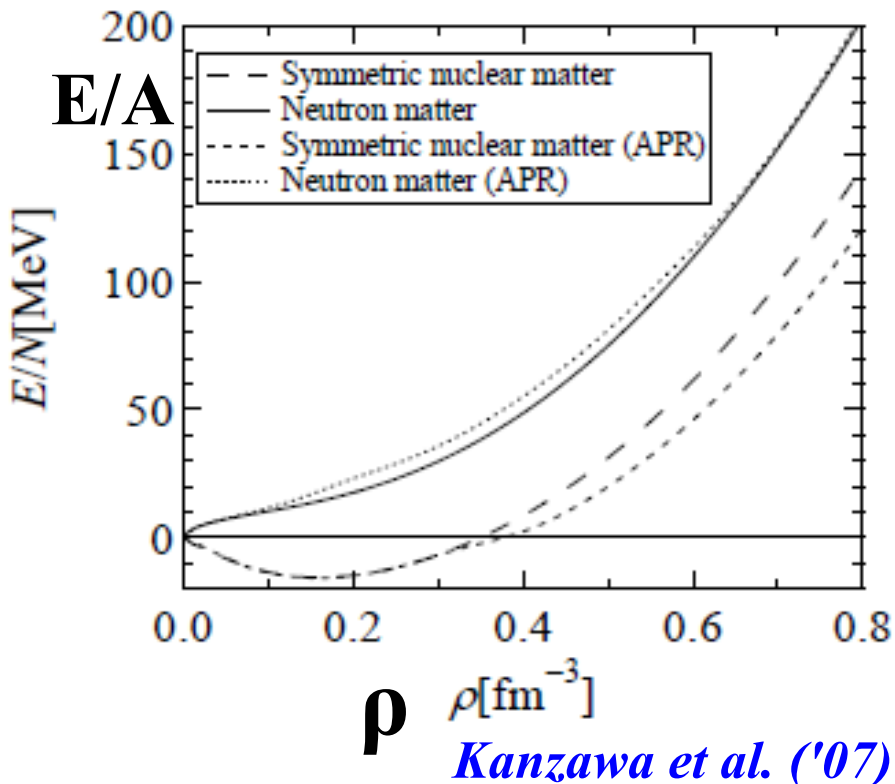
Variational Calculation starting from bare nuclear force

B. Friedman, V.R. Pandharipande, NPA361('81)502;

A. Akmal, V.R. Pandharipande, D.G. Ravenhall, PRC58('98)1804;

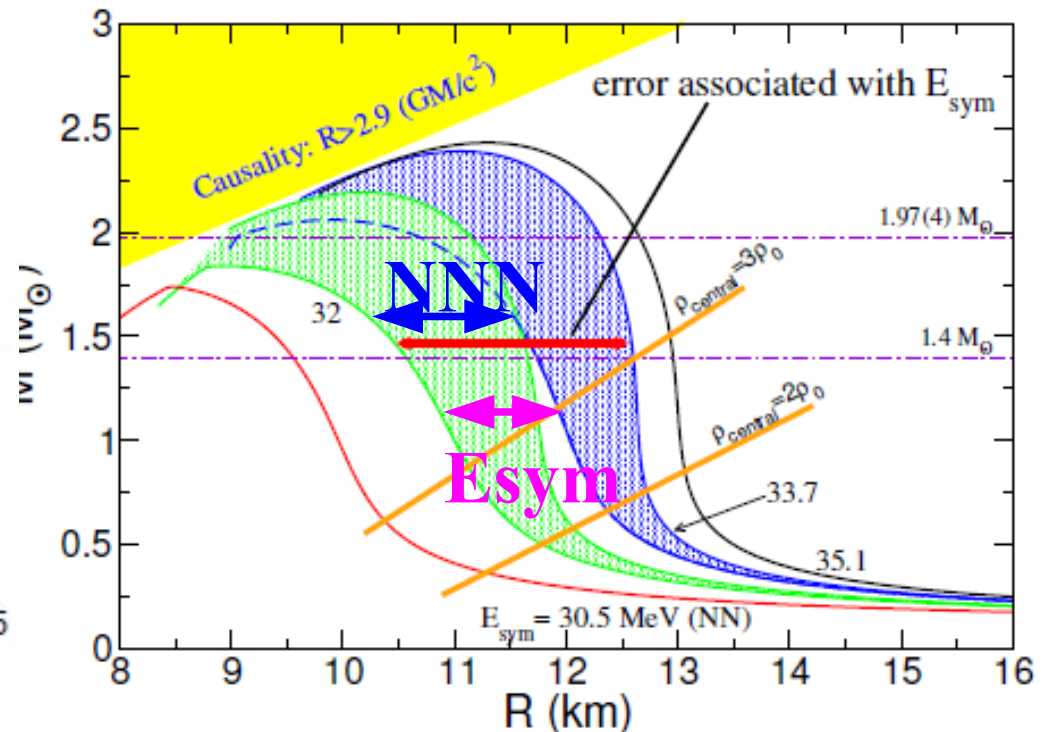
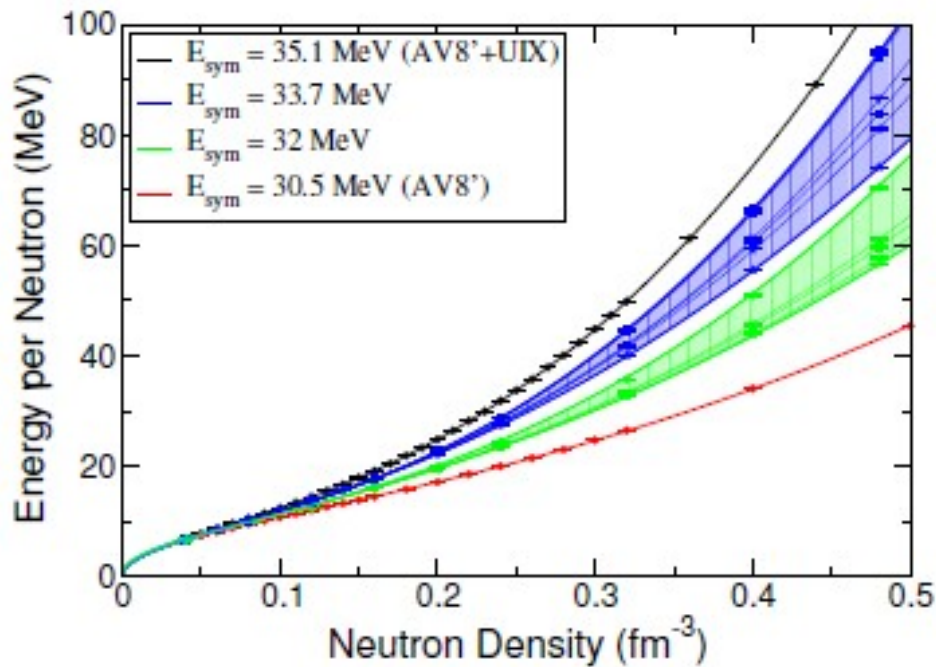
H. Kanzawa, K. Oyamatsu, K. Sumiyoshi, M. Takano, NPA791 ('07) 232.

Argonne v18(v14) + Rel. corr. + Three Nucleon Int.



Quantum Monte-Carlo calc.

- Auxiliary Field Diffusion Monte-Carlo (AFDMC) calc.
 - Hubbard-Stratonovich transf. + MC integral over aux. fields.
 - 3n force parameters are tuned to fit finite nuclei.
 - 2 MeV Difference in E_{sym} results in 1.5 km (15 %) diff. in R_{NS} .



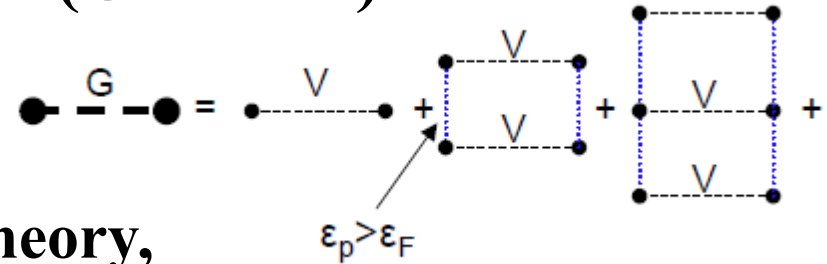
Gandolfi, Carlson, Reddy, PRC 032801, 85 (2012).

Bruckner-Hartree-Fock

- Effective interaction from bare NN int. (G-matrix).

$$g(E) = V + V \frac{Q}{E - H_0} g(E)$$

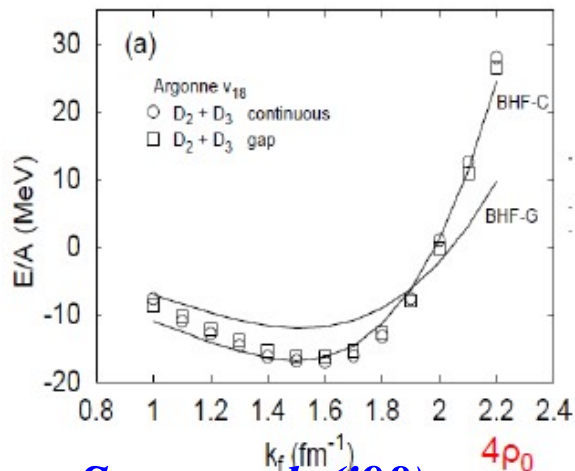
Pauli



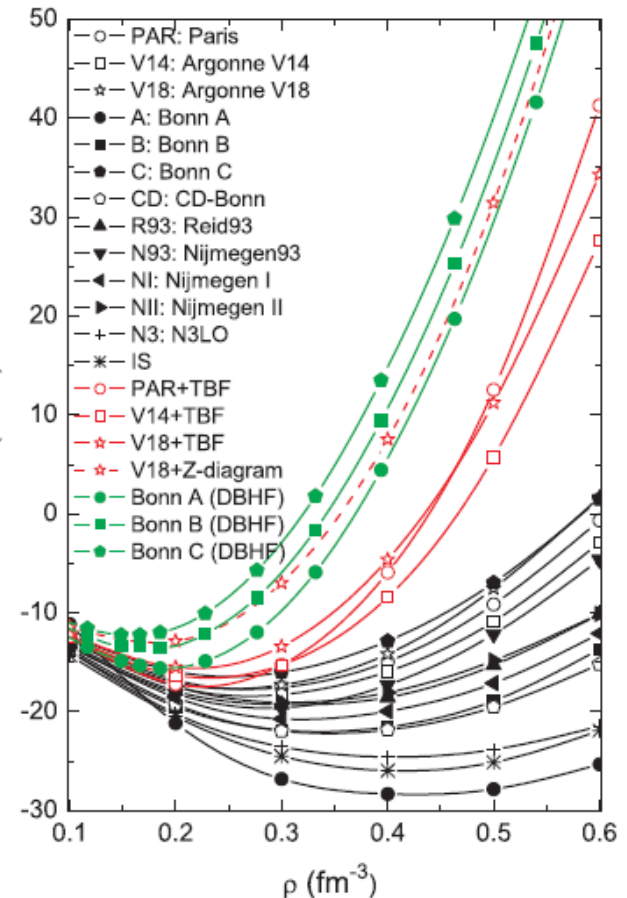
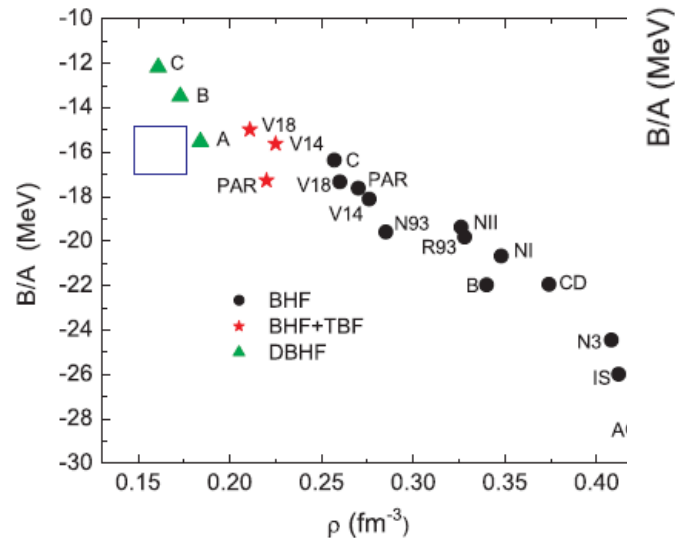
- G-matrix = Lowest order Bruckner theory, but next-to-leading terms give small effects at $\rho < 4 \rho_0$.

Song, Baldo, Giansiracusa, Lombardo ('98)

- Need 3-body force to reproduce saturation point.



Song et al. ('98)



Z.H.Li, U. Lombardo, H.-J. Schulze, W. Zuo, L. W. Chen, H. R. Ma, PRC74('06)047304.

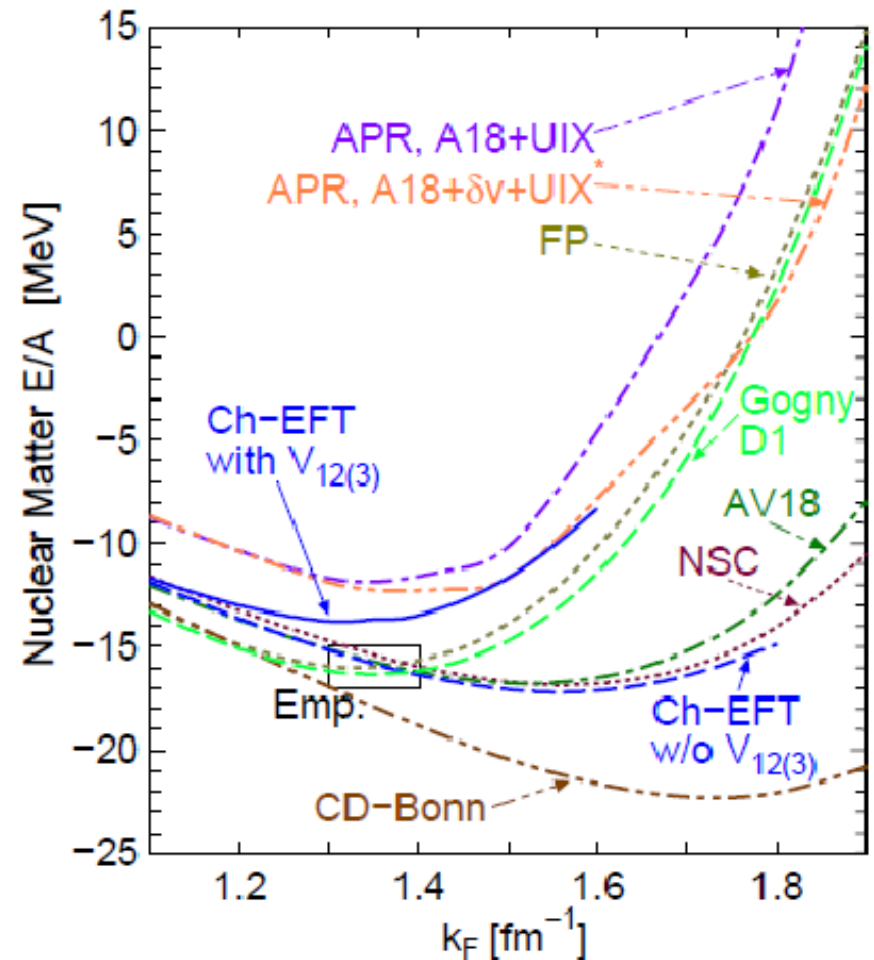
BHF with Ch-EFT & Lattice NN force

- Bruckner-HF calc. with NN (N3LO)+3NF(N2LO) interactions from Chiral Effective Field Theory *M.Kohno ('13)*

- Ch-EFT = Eff. Field Theory with the same symmetry as QCD
Weinberg; Gasser, Leutwyler ('84)
→ Systematically gives NN & NNN interaction terms.
Epelbaum, Gockle, Meissner ('05)

- Bruckner HF calc. with NN int. from Lattice QCD.
Inoue et al. (HAL QCD Coll.), PRL111 ('13)112503

- Not yet reliable but promising !



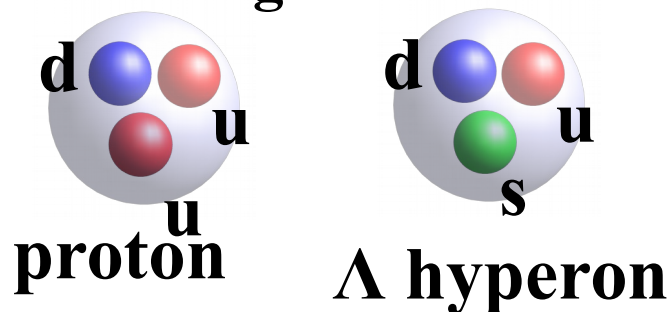
M. Kohno, PRC88('13)064005

Massive Neutron Star puzzle

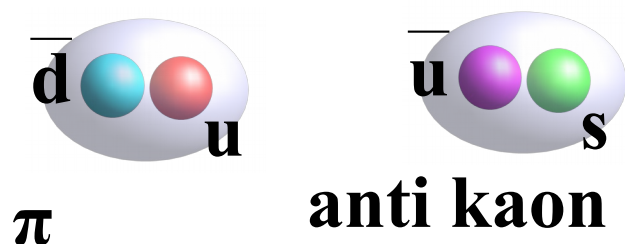
Neutron star – Is it made of neutrons ?

■ Possibilities of various constituents in neutron star core

• Strange Hadrons

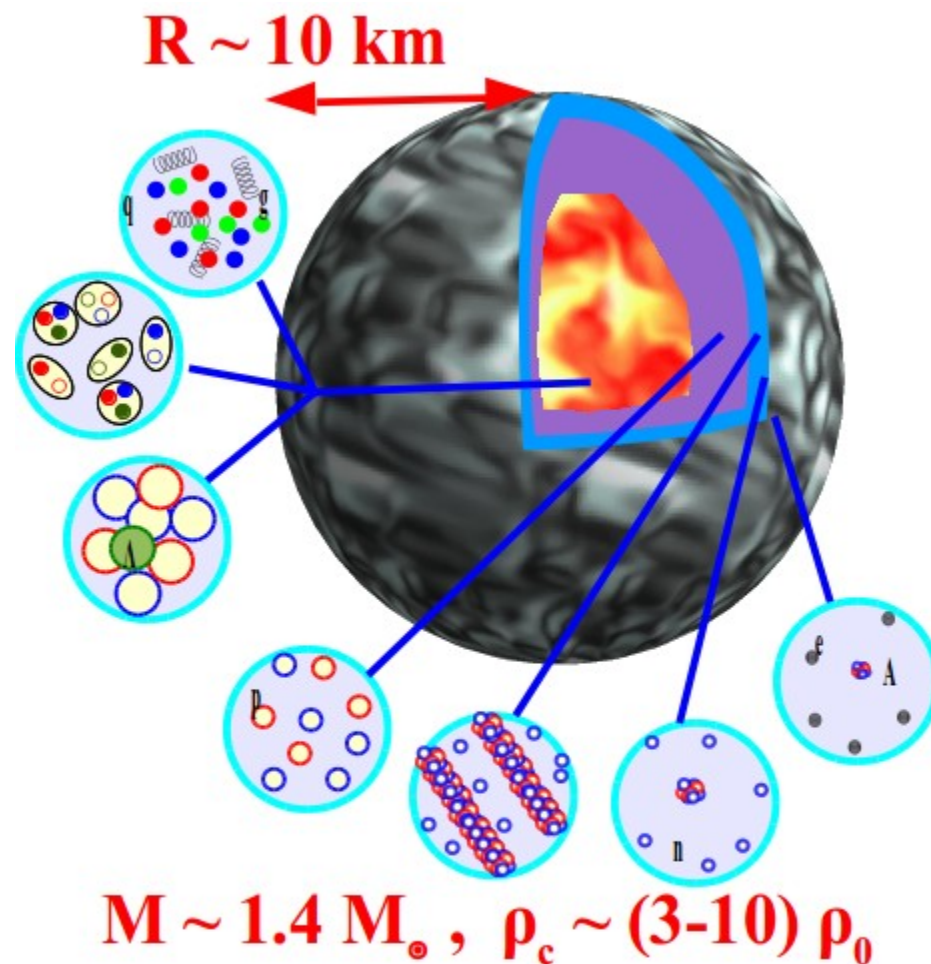


• Meson condensate (K, π)



• Quark matter

• Quark pair condensate (Color superconductor)



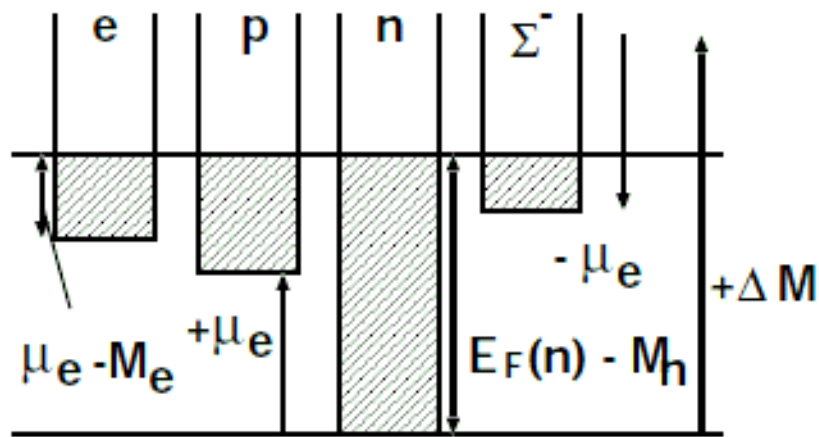
Hyperons in Dense Matter

■ What appears at high density ?

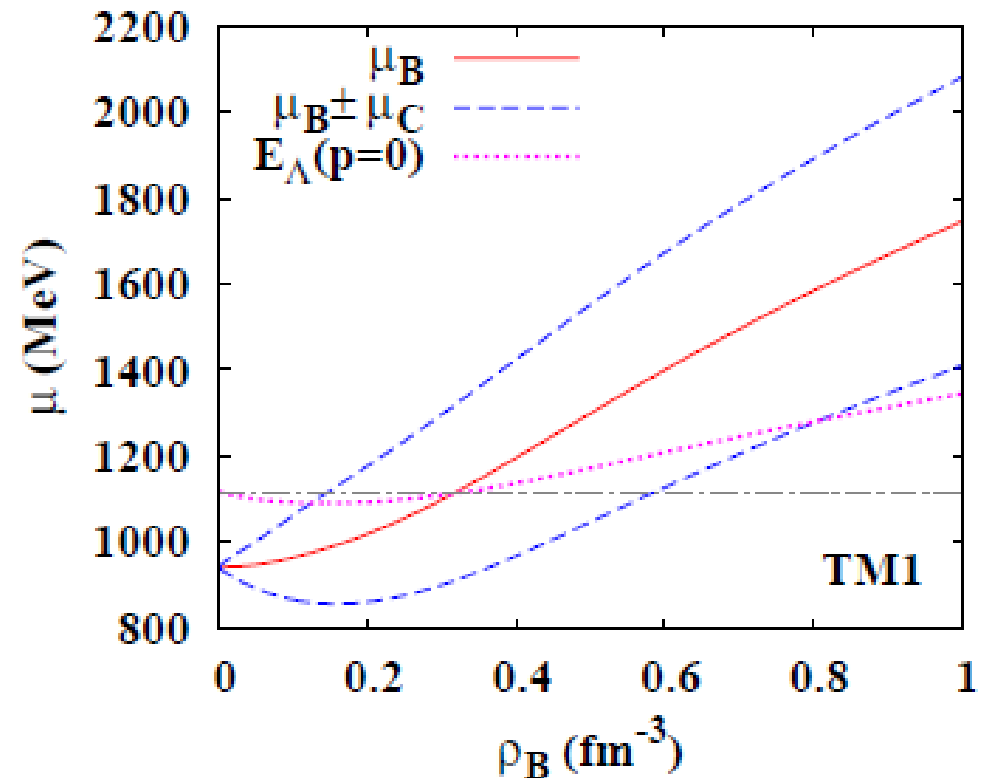
- Nucleon superfluid (3S_1 , 3P_2), Pion condensation, Kaon condensation, Baryon Rich QGP, Color SuperConductor (CSC), Quarkyonic Matter, ...

● Hyperons

Tsuruta, Cameron (66); Langer, Rosen (70); Pandharipande (71); Itoh(75); Glendenning; Weber, Weigel; Sugahara, Toki; Schaffner, Mishustin; Balberg, Gal; Baldo et al.; Vidana et al.; Nishizaki, Yamamoto, Takatsuka; Kohno, Fujiwara et al.; Sahu, Ohnishi; Ishizuka, Ohnishi, Sumiyoshi, Yamada; ...



*Chemical potential overtakes Λ mass
→ appearance of Λ*



Neutron Star Masses

- NS masses in NS binaries can be measured precisely by using some of GR effects via doppler shifts.

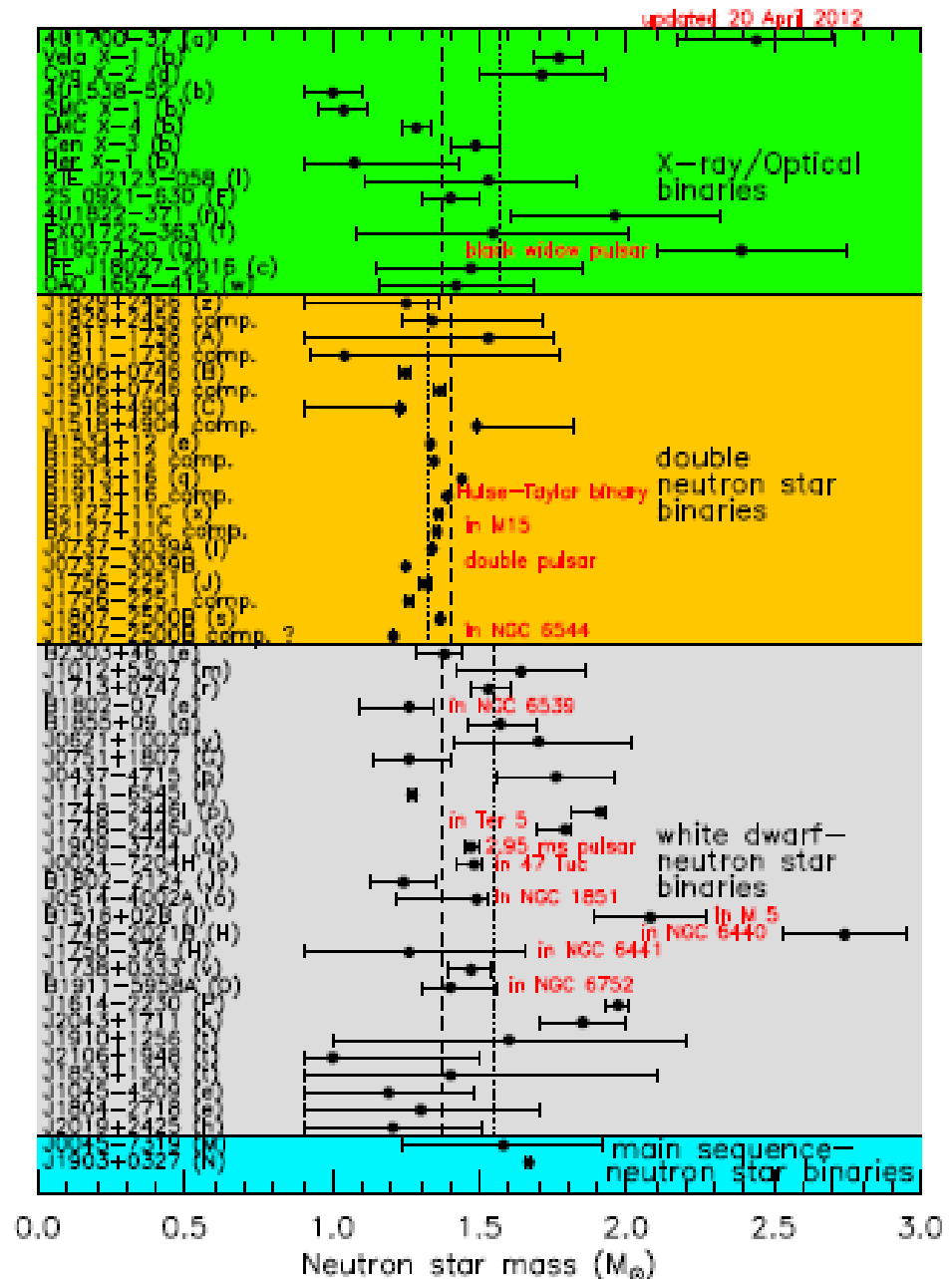
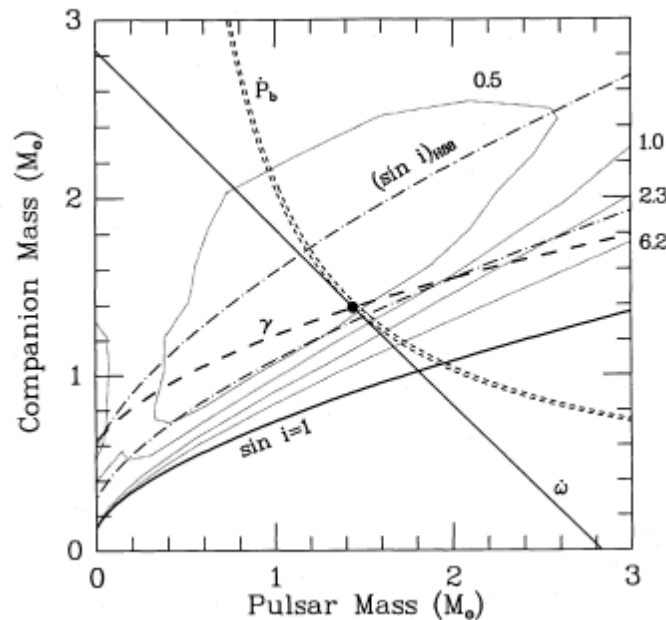
- Perihelion shift+Einstein delay

$$\rightarrow M = 1.442 \pm 0.003 M_{\odot}$$

(Hulse-Taylor pulsar)

Taylor, Weisenberg ('89)

- Many NSs have $M \sim 1.4 M_{\odot}$.



Lattimer (2013)

Massive Neutron Star Puzzle

- Observation of massive neutron stars ($M \sim 2 M_{\odot}$)

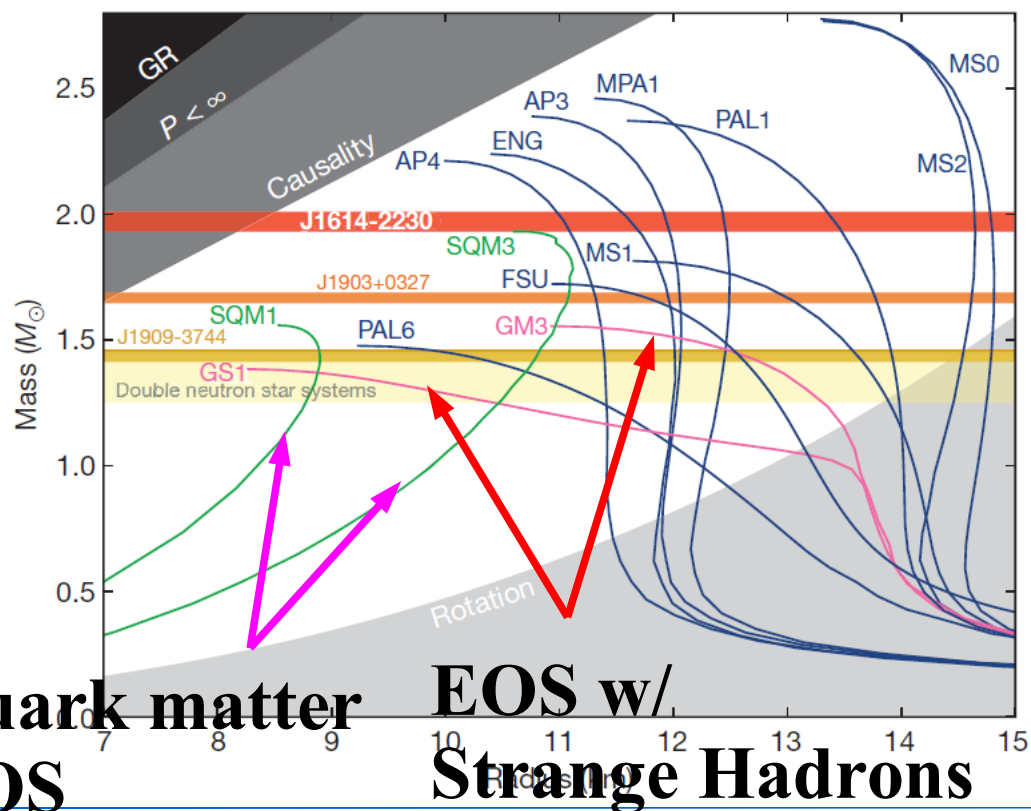
- PSR J1614-2230 (NS-WD binary), $1.97 \pm 0.04 M_{\odot}$

Demorest et al., Nature 467('10)1081 (Oct.28, 2010).

”Kinematical” measurement (Shapiro delay, GR)
+ large inclination angle

- PSR J0348+0432 (NS-WS binary), $2.01 \pm 0.04 M_{\odot}$

Antoniadis et al., Science 340('13)

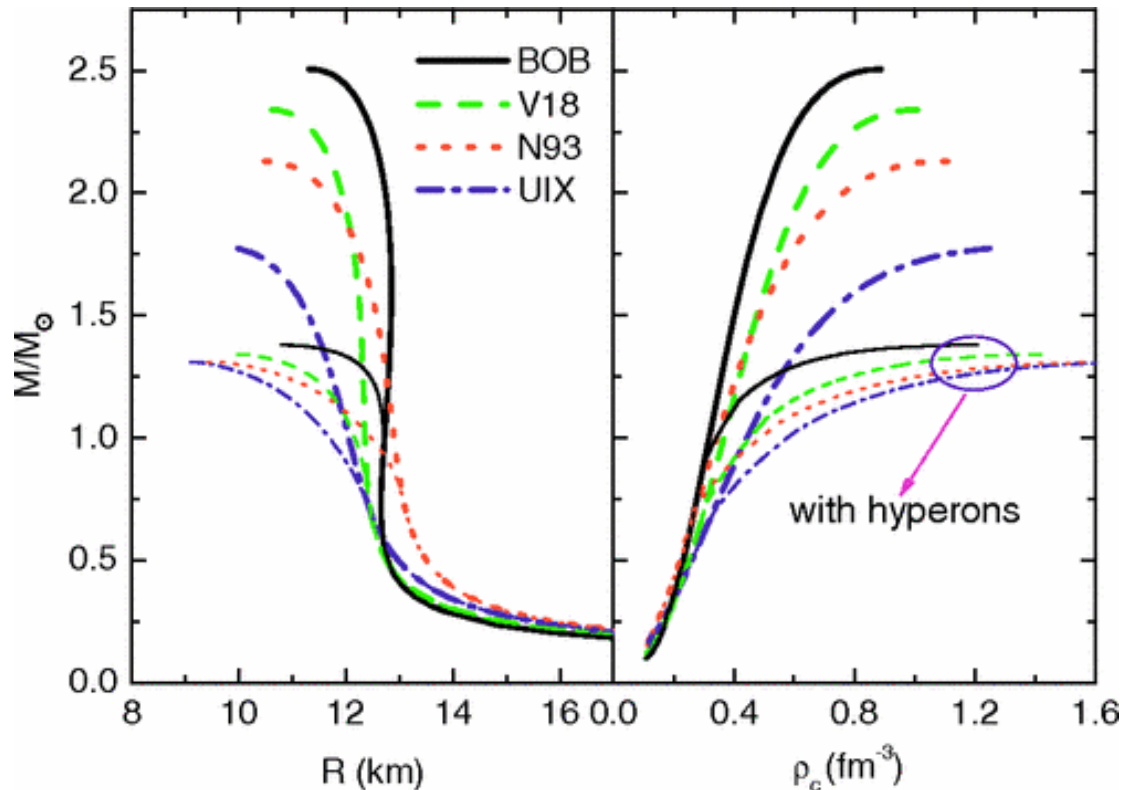


No Exotics in NS ?

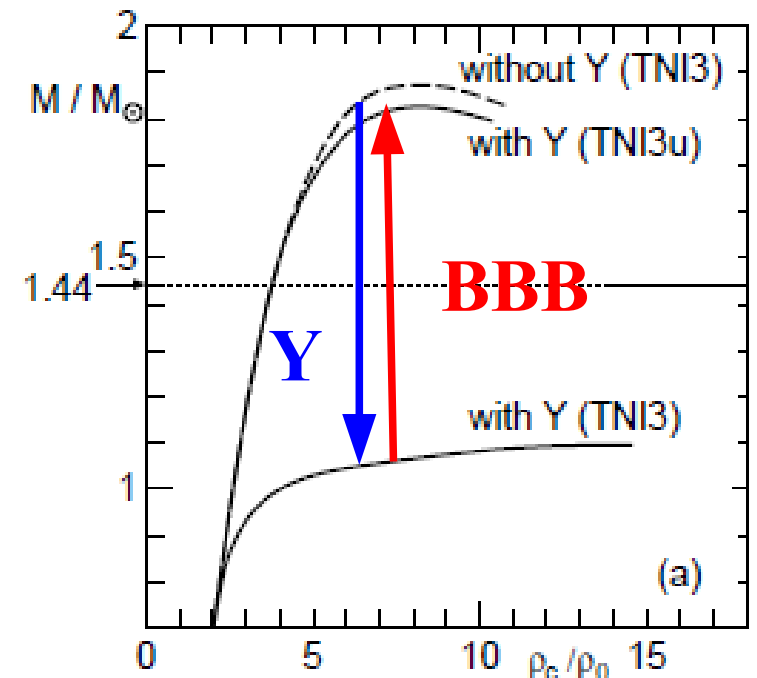
Quark matter EOS EOS w/ Strange Hadrons

Bruckner-Hartree-Fock theory with Hyperons

- Microscopic G-matrix calculation with realistic NN, YN potential and microscopic (or phen.) 3N force (or 3B force).
 - Interaction dep. (V18, N93, ...) is large → Need finite nuclear info.
E.Hiyama, T.Motoba, Y.Yamamoto, M.Kamimura / M.Tamura et al.
 - NS collapses with hyperons w/o 3BF.



Z.H.Li, H.-J.Schulze, PRC78('08),028801.



S. Nishizaki, T. Takatsuka, Y. Yamamoto, PTP108('02)703.

RMF with Hyperons (Single Λ hypernuclei)

■ RMF for Λ hypernuclei

$x \sim 1/3$: R. Brockmann, W. Weise, PLB69('77)167; J. Boguta and S. Bohrman, PLB102('81)93.

$x \sim 2/3$: N. K. Glendenning, PRC23('81)2757, PLB114('82)392;

Tensor: Y. Sugahara, H. Toki, PTP92('94)803; H. Shen, F. Yang, H. Toki, PTP115('06)325;
J. Mares, B. K. Jennings, PRC49('94)2472.

ρ -dep. coupling: H. Lenske, Lect. Notes Phys. 641('04)147; C. M. Keil, F. Hofmann, H. Lenske, PRC 61('00)064309.

$SU(3)$ or $SU(6)$ (ζ, φ): J. Schaffner, C. B. Dover, A. Gal, C. Greiner, H. Stoecker, PRL71('93)1328;
Schaffner et al., Ann.Phys.235('94)35; J. Schaffner, I. N. Mishustin, PRC 53('96)1416.

Chiral $SU(3)$ RMF: K. Tsubakihara, H. Maekawa, H. Matsumiya, AO, PRC81('10)065206.

- Sep. E. of Λ is well fitted by $U_{\Lambda} \sim -30 \text{ MeV} \sim 2/3 U_N$

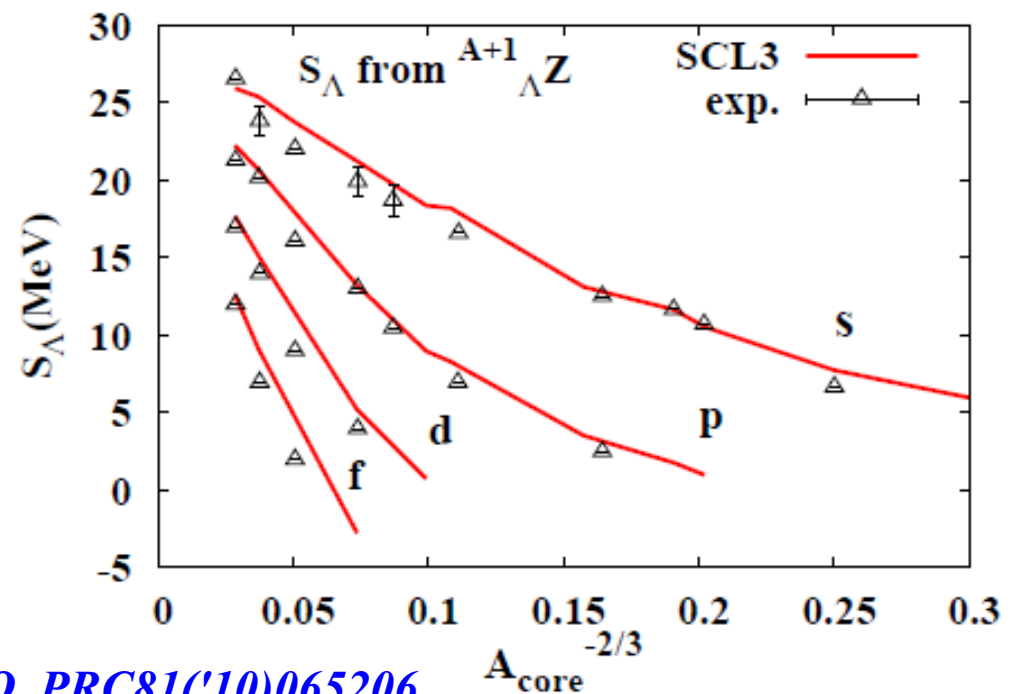
- Coupling with mesons

$$x_M = g_{M\Lambda} / g_{MN}$$

quark counting: $x_{\sigma} \sim 2/3$

π exchanges: $x_{\sigma} \sim 1/3$

→ Which is true ?



K. Tsubakihara, H. Maekawa, H. Matsumiya, AO, PRC81('10)065206.

Hyperon Composition in Dense Matter

■ Hyperon start to emerge at $(2-3)\rho_0$ in Neutron Star Matter !

■ Hyperon composition in NS is sensitive to Hyperon potential.

- $U_\Lambda \sim -30$ MeV: Well-known

- $U_\Sigma \sim -(12-15)$ MeV

(K^-, K^+) reaction, twin hypernuclei

P. Khaustov et al. (E885), PRC61('00)054603;

S. Aoki et al., PLB355('95)45.

- $U_\Sigma \sim -30$ MeV (Old conjecture)

→ Σ - appears prior to Λ

- $U_\Sigma > 0$ (repulsive) → No Σ in NS

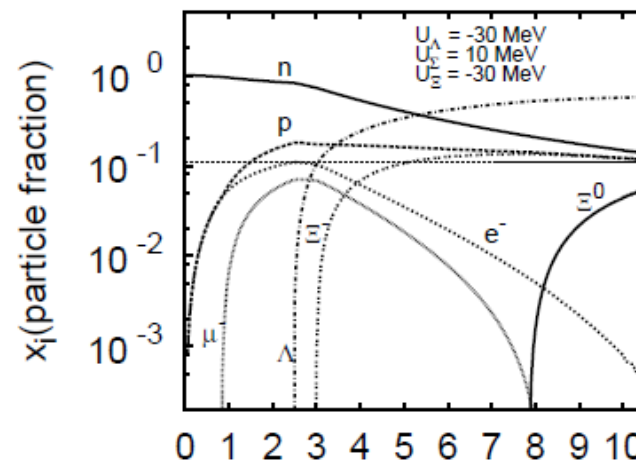
Σ atom (phen. fit), QF prod.

S. Balberg, A. Gal, NPA625('97)435;

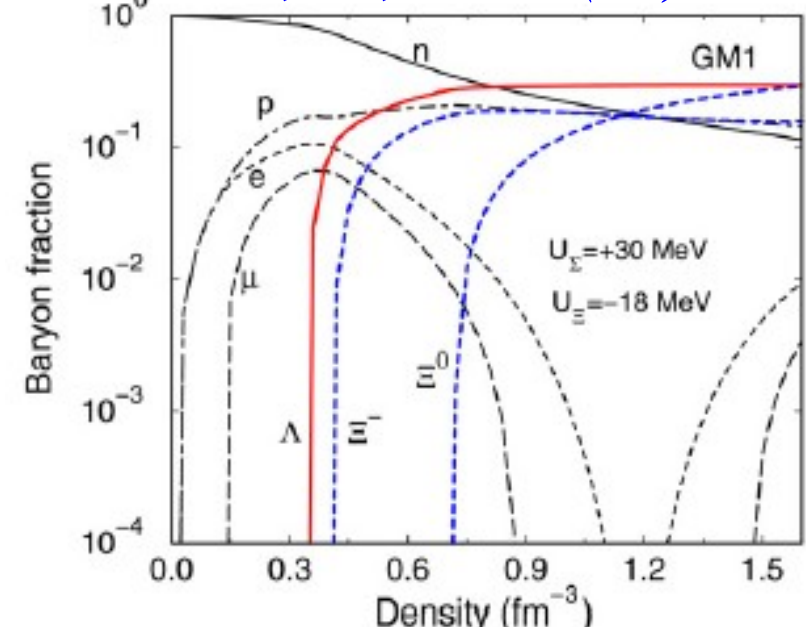
H. Noumi et al., PRL89('02)072301;

T. Harada, Y. Hirabayashi, NPA759('05)143;

M. Kohno et al. PRC74('06)064613.



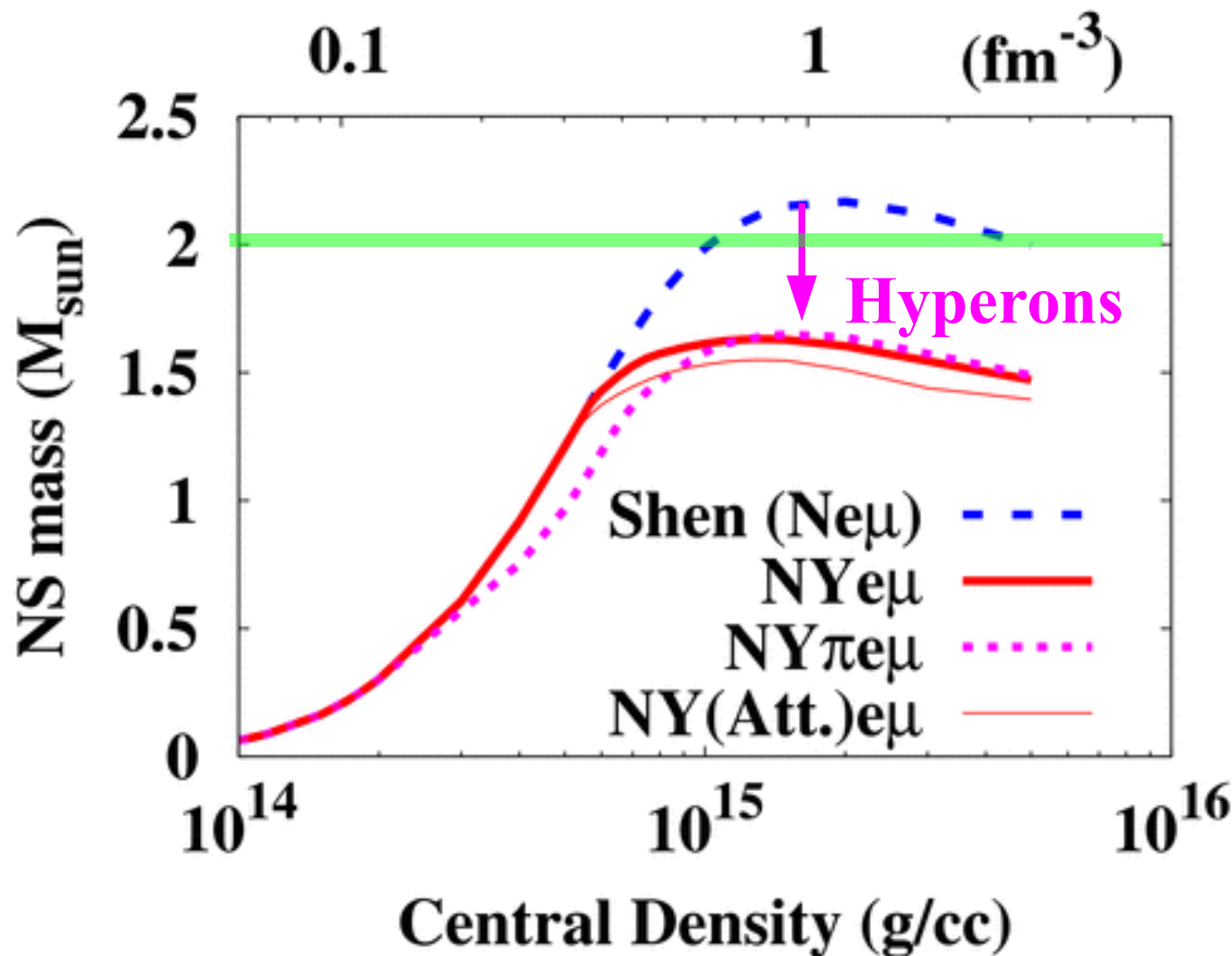
P.K.Sahu, AO, NPA691('01)439c



J. Schaffner-Bielich, NPA804('08)309.

NS M-R Relation in RMF with Hyperons

Ishizuka, AO, Tsubakihara, Sumiyoshi, Yamada, *J. Phys. G35(08),085201*



c.f. H.Shen+('09) → n, p, Λ EOS

Summary of Part I

- Neutron star can be regarded as a gigantic nucleus
 - $M \sim 1.4 M_{\odot}$, $R \sim 10$ km \rightarrow Density $\sim (1-3) \rho_0$
- Some of NS masses have been obtained precisely.
 - Kepler orbit + General Relativity corrections $\rightarrow m_1, m_2$
 - Two NSs are found to have $M \sim 2 M_{\odot}$ recently.
PSR J1614-2230, PSR J0348+0432
- NS radii are more difficult to measure, and model dependent.
 - Three methods: Surface emission, Eddington limit, Redshift.
 - Conservative estimate: $R_{\text{NS}} = (8-15)$ km.
- Nuclear matter EOSs have been studied in various ways.
 - Microscopic calc. (starting from bare NN force),
Phen. model calc. (using effective NN force).
 - Many of them with hyperons predict NS max. mass $< 2 M_{\odot}$.

*Thank you for your attention !
See you 15(?) minutes later.*

Contents

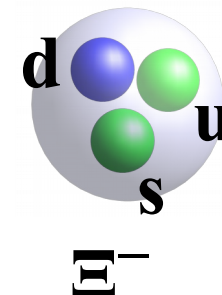
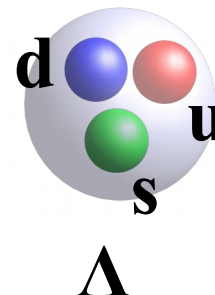
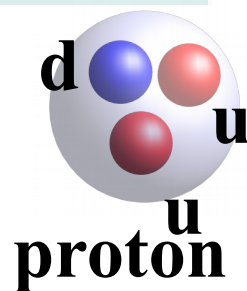
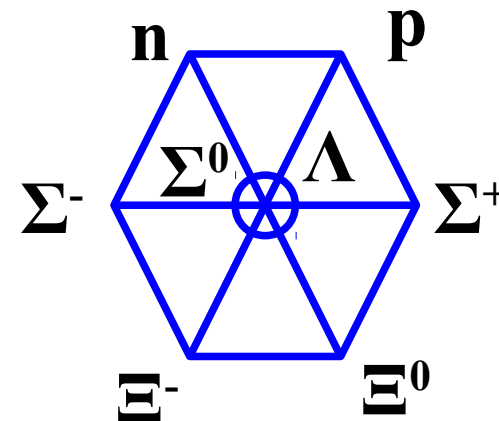
- **Introduction**
- **Part I : Basics of Neutron Star Physics**
 - **Neutron star mass & radius observations**
 - **Nuclear matter EOS and neutron stars**
 - **Massive Neutron Star Puzzle**
- **Part II : Hypernuclear Physics & Neutron Stars**
 - **Hypernuclear Physics : Implications from Experiments**
 - **What is necessary to solve massive NS puzzle ?**
 - **Recent Attempts toward the massive NS puzzle**
- **Summary**

Hypernuclear Physics Implications from Experiments

Hyperons (Baryons with Strangeness)

- Ground state baryon $SU(3)_f$ octet ($J^\pi=1/2^+$)

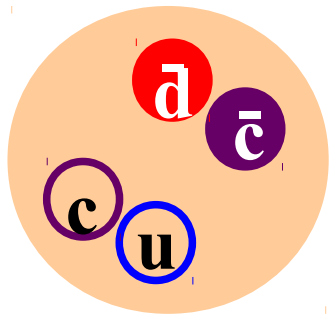
Baryon	M(Mev)	S	Comp.
n	940	0	udd
p	938	0	uud
Λ	1116	-1	$(uds-dus)/\sqrt{2}$
Σ^+	1189	-1	uus
Σ^0	1193	-1	$(uds+dus)/\sqrt{2}$
Σ^-	1197	-1	dds
Ξ^0	1315	-2	uss
Ξ^-	1321	-2	dss



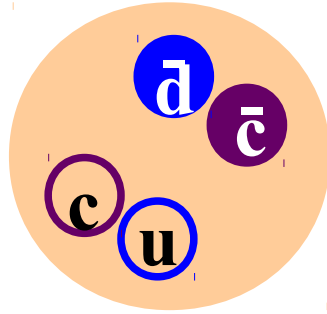
Exotic Hadrons

Exotic hadrons

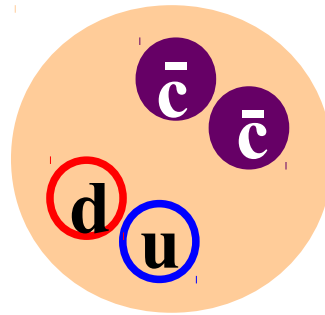
→ X, Y, Z, Θ^+ , Discovered/Proposed at LEPs, Belle, BaBar,...



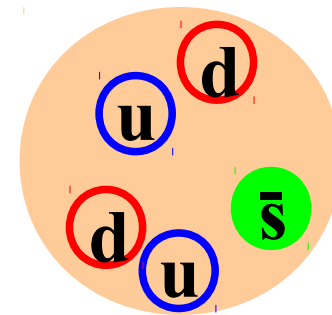
Z(4430)



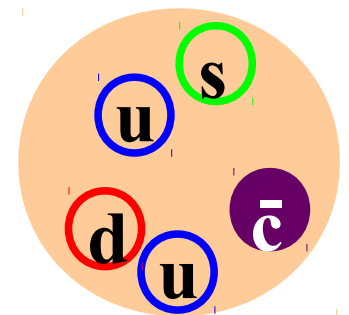
X(3872)



T_{cc}



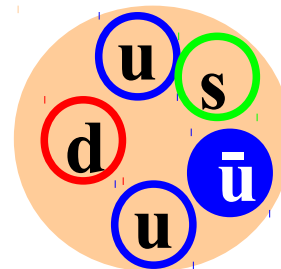
Θ^+



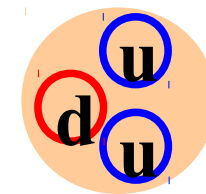
Θ^+_{cs}

Various pictures

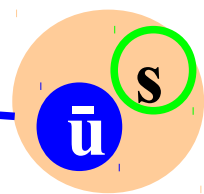
- Di-quark component
- Hadronic molecule
- QQ couples with QQ $q\bar{q}$



$uuds_{\bar{u}}$



p



K^-

$\Lambda(1405)$

$SU(3)_f$ transformation

- Fundamental triplet $(u,d,s)^T = \mathbf{q} \rightarrow \mathbf{q}' = \mathbf{U} \mathbf{q}$ ($\mathbf{U} \in SU(3)$)
- Diquark $\mathbf{D}_i = \varepsilon_{ijk} \mathbf{q}_j \mathbf{q}_k \rightarrow \mathbf{D}' = \mathbf{D} \mathbf{U}^+$
- Baryon octet $\mathbf{B}_{ij} = \mathbf{D}_j \mathbf{q}_i \rightarrow \mathbf{B}' = \mathbf{U} \mathbf{B} \mathbf{U}^+$

$$\begin{pmatrix} [ds]u & [su]u & [ud]u \\ [ds]d & [su]d & [ud]d \\ [ds]s & [su]s & [ud]s \end{pmatrix} = \begin{pmatrix} \frac{\Lambda}{\sqrt{6}} + \frac{\Sigma^0}{\sqrt{2}} & \Sigma^+ & p \\ \Sigma^- & \frac{\Lambda}{\sqrt{6}} - \frac{\Sigma^0}{\sqrt{2}} & n \\ \Xi^- & \Xi^0 & -\frac{2\Lambda}{\sqrt{6}} \end{pmatrix}$$

$SU(3)_f$ transformation

- Fundamental triplet $(u,d,s)^T = \mathbf{q} \rightarrow \mathbf{q}' = \mathbf{U} \mathbf{q}$ ($\mathbf{U} \in SU(3)$)
- Anti-quark $\bar{\mathbf{q}} \rightarrow \bar{\mathbf{q}}' = \bar{\mathbf{q}} \mathbf{U}^+$
- Meson octet $\mathbf{M}_{ij} = \bar{\mathbf{q}}_j \mathbf{q}_i \rightarrow \mathbf{M}' = \mathbf{U} \mathbf{M} \mathbf{U}^+$

$$\begin{pmatrix} \bar{u}u & \bar{d}u & \bar{s}u \\ \bar{u}d & \bar{d}d & \bar{s}d \\ \bar{u}s & \bar{d}s & \bar{s}s \end{pmatrix} = \begin{pmatrix} \frac{\eta}{\sqrt{6}} + \frac{\pi^0}{\sqrt{2}} & \pi^+ & K^+ \\ \pi^- & \frac{\eta}{\sqrt{6}} - \frac{\pi^0}{\sqrt{2}} & K^0 \\ K^- & \bar{K}^0 & -\frac{2\eta}{\sqrt{6}} \end{pmatrix} = P$$

$$S = \begin{pmatrix} \frac{\sigma}{\sqrt{2}} + \frac{a_0}{\sqrt{2}} & a_0^+ & \kappa^+ \\ a_0^- & \frac{\sigma}{\sqrt{2}} - \frac{a_0}{\sqrt{2}} & \kappa^0 \\ \kappa^- & \bar{\kappa}^0 & \xi \end{pmatrix} \quad V = \begin{pmatrix} \frac{\omega}{\sqrt{2}} + \frac{\rho^0}{\sqrt{2}} & \rho^+ & K^{*+} \\ \rho^- & \frac{\omega}{\sqrt{2}} - \frac{\rho^0}{\sqrt{2}} & K^{*0} \\ K^{*-} & \bar{K}^{*0} & \varphi \end{pmatrix}$$

$SU(3)_f$ invariant coupling

■ Baryon-Meson coupling

$$\begin{aligned}\mathcal{L}_{BV} &= \sqrt{2}\{g_s \text{tr}(M_v) \text{tr}(\bar{B}B) + g_D \text{tr}(\bar{B}\{M_v, B\}) + g_F \text{tr}(\bar{B}[M_v, B])\} \\ &= \sqrt{2}\{g_s \text{tr}(M_v) \text{tr}(\bar{B}B) + g_1 \text{tr}(\bar{B}M_v B) + g_2 \text{tr}(BBM_v)\}\end{aligned}$$

■ Assumption

- BM coupling is $SU(3)$ invariant
- N does not couple with $\bar{s}s$ vector meson

$$g_{\omega\Lambda} = \frac{5}{6}g_{\omega N} - \frac{1}{2}g_{\rho N}, \quad g_{\phi\Lambda} = \frac{\sqrt{2}}{6}(g_{\omega N} + 3g_{\rho N})$$

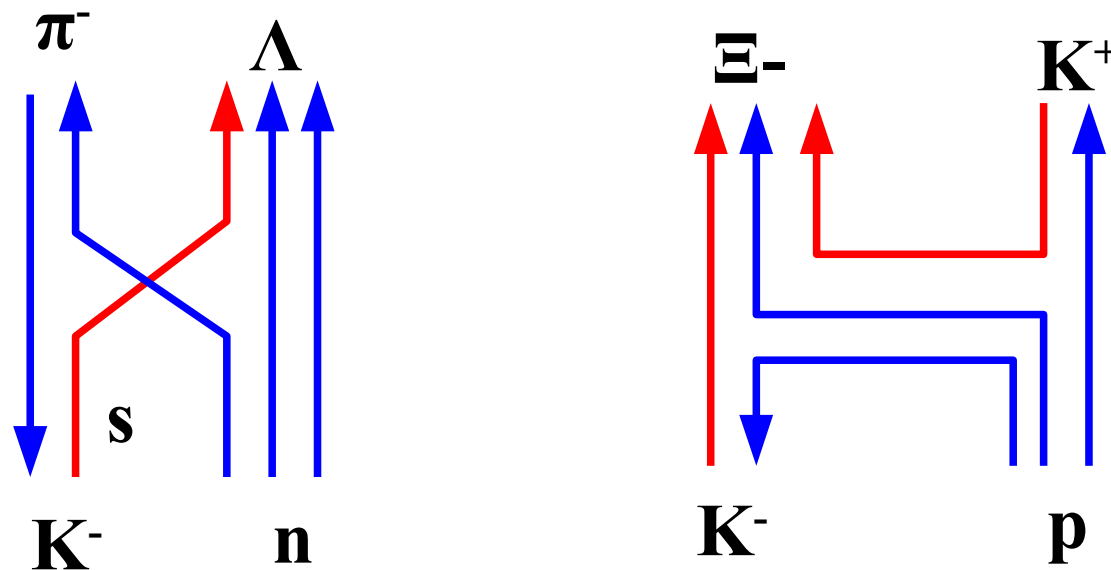
■ Further simplification: $g_{\rho N} = g_{\omega N}/3$ (quark counting)

$$g_{\omega N} = g_v, \quad g_{\rho N} = g_v/3, \quad g_{\omega\Lambda} = 2g_v/3, \quad g_{\phi\Lambda} = \sqrt{2}g_v/3$$

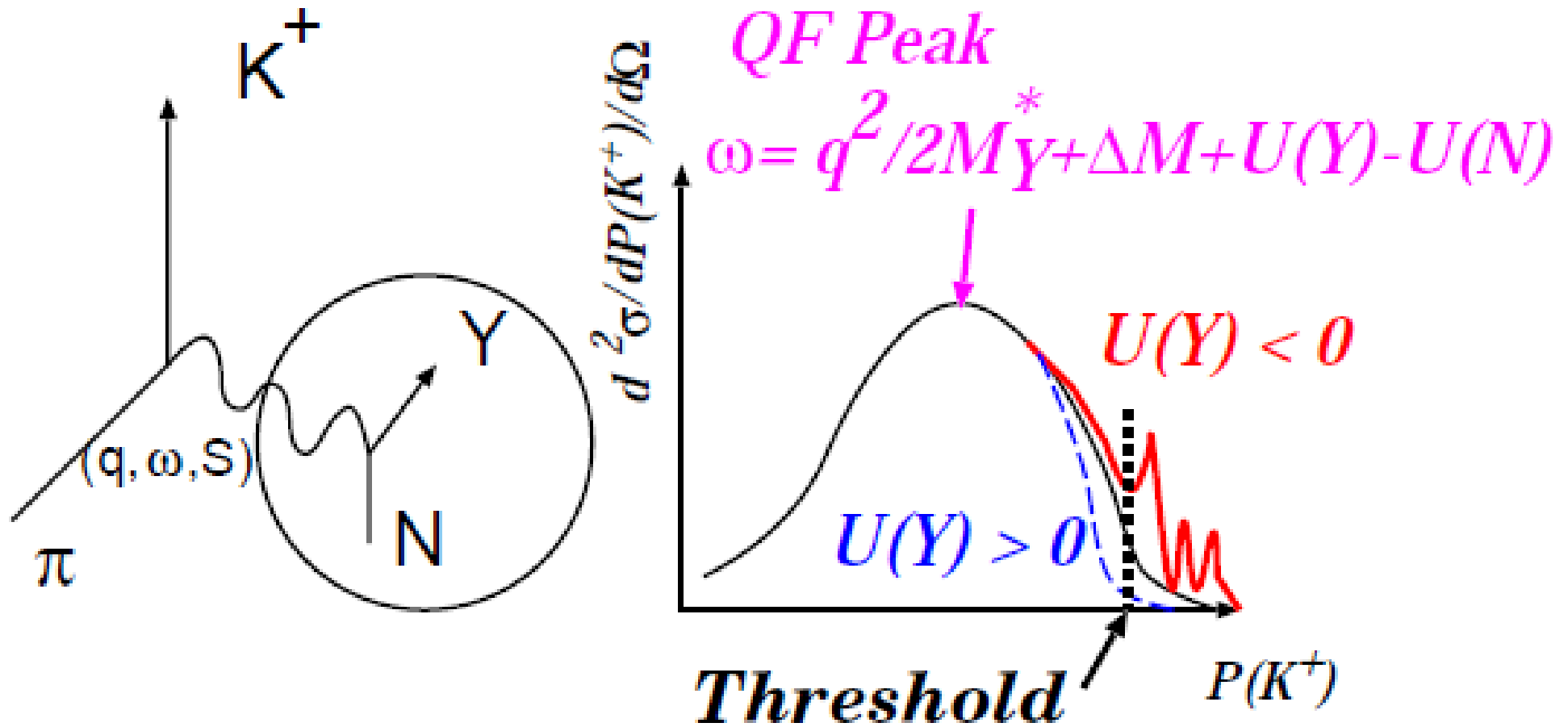
Hypernuclear formation

- (K^-, π^-) , (π^-, K^+) , and (K^-, K^+) reactions on nuclei \rightarrow Hypernuclei

Reaction	Elementary Processes	
	Main Process	Other Processes
(K^-, π^-)	$K^- n \rightarrow \pi^- \Lambda$,	$K^- n \rightarrow \pi^- \Sigma^0, K^- p \rightarrow \pi^- \Sigma^+$
(K^-, π^+)	$K^- p \rightarrow \pi^+ \Sigma^-$,	$K^- pp \rightarrow \pi^+ \Lambda n$ (n-rich hypernuclear formation)
(π^+, K^+)	$\pi^+ n \rightarrow K^+ \Lambda$,	$\pi^+ n \rightarrow K^+ \Sigma^0, \pi^+ p \rightarrow K^+ \Sigma^+$
(π^-, K^+)	$\pi^- p \rightarrow K^+ \Sigma^-$,	$\pi^- pp \rightarrow K^+ \Lambda n$ (n-rich hypernuclear formation)
(K^-, K^+)	$K^- p \rightarrow K^+ \Xi^-$,	$K^- pp \rightarrow K^+ \Lambda \Lambda$

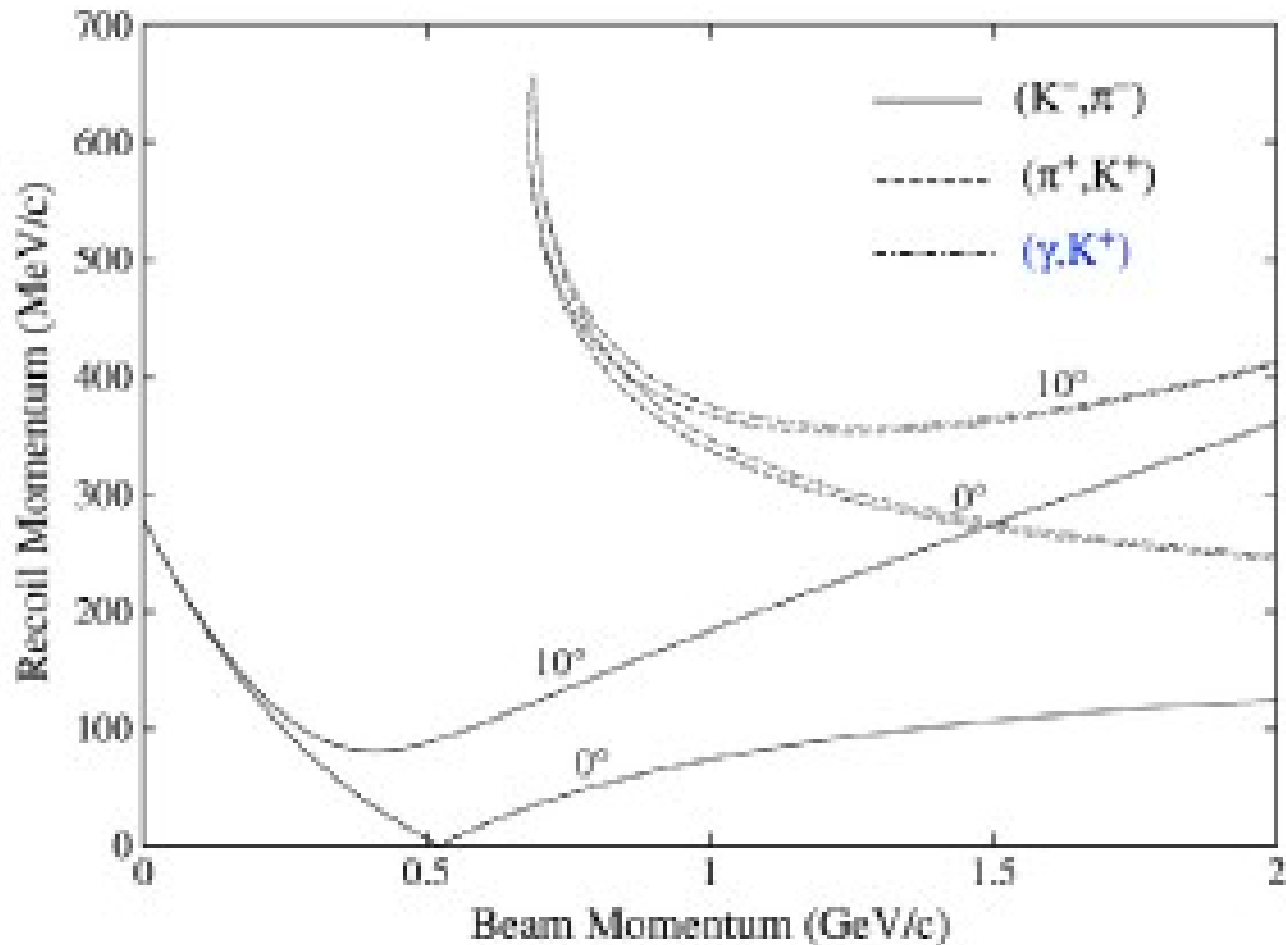


Hypernuclear formation



Hypernuclear formation

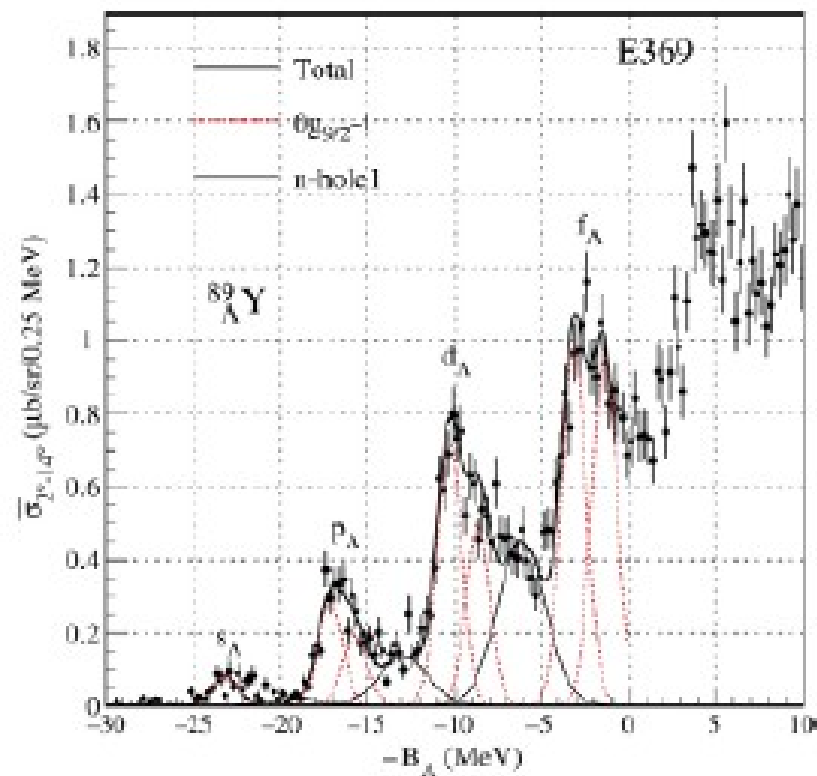
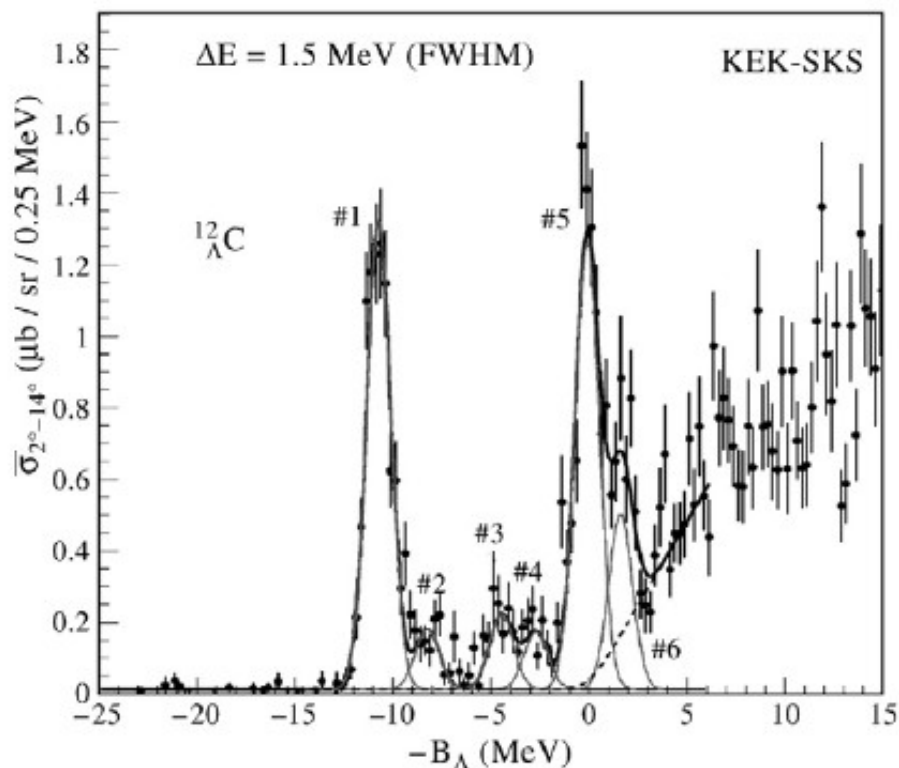
- (K^-, π^-) : $Q > 0$, Small momentum transfer \rightarrow substitutional reaction
- (π, K^+) : $Q < 0$, Momentum transfer $\sim 300 \text{ MeV}/c \sim k_F$



Λ hypernuclear formation

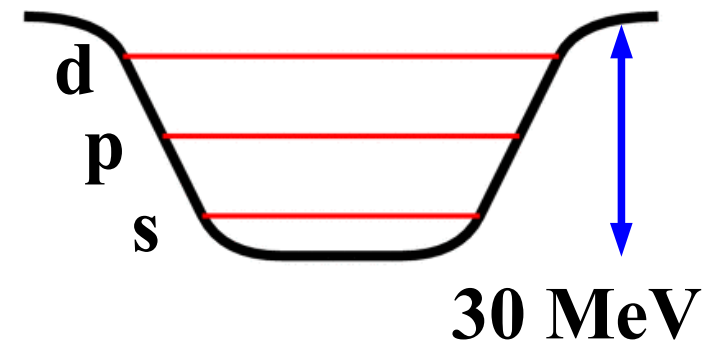
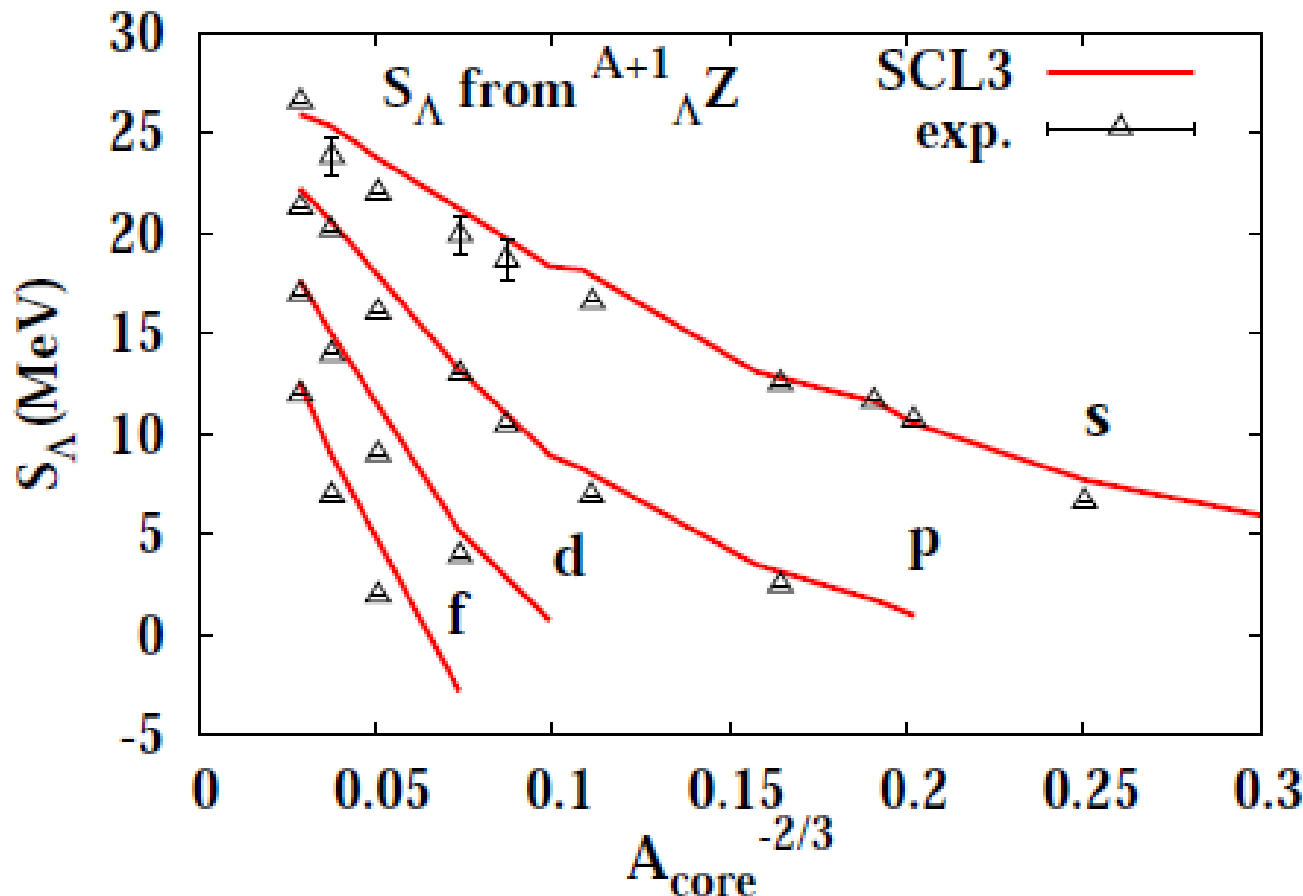
■ (π^+, K^+) reactions on nuclei

- $q \sim k_F \rightarrow$ various s.p. states of Λ are populated



Single particles states of Λ in nuclei

- Single particle potential depth of Λ is around -30 MeV
 - s, p, d, f, ... states are clearly seen
 - $A_{\text{core}}^{-2/3} \propto R^{-2} \propto \text{K.E. of } \Lambda$



Σ production in nuclei

- Only one bound state $^4_{\Sigma}\text{He}$ (Too light !)
→ Continuum (Quasi-Free) Spectroscopy is necessary
- Cont. Spec. Theory = Distorted Wave Impulse Approx. (DWIA)

$$\frac{d^2 \sigma}{dE_K d\Omega_K} = \beta \left(\frac{d\sigma}{d\Omega} \right)_{N\pi \rightarrow KY}^{Elem.} S(E, q)$$

Kinematical Factor

Elem. Cross Sec.

Strength Func.

- Large (ω, q) range → Important to respect **On-Shell Kinematics**

- Another way: Σ^- atomic shift

- Atomic shift of Σ^- with O, Mg, Al, S, Si, W, Pb core are measured

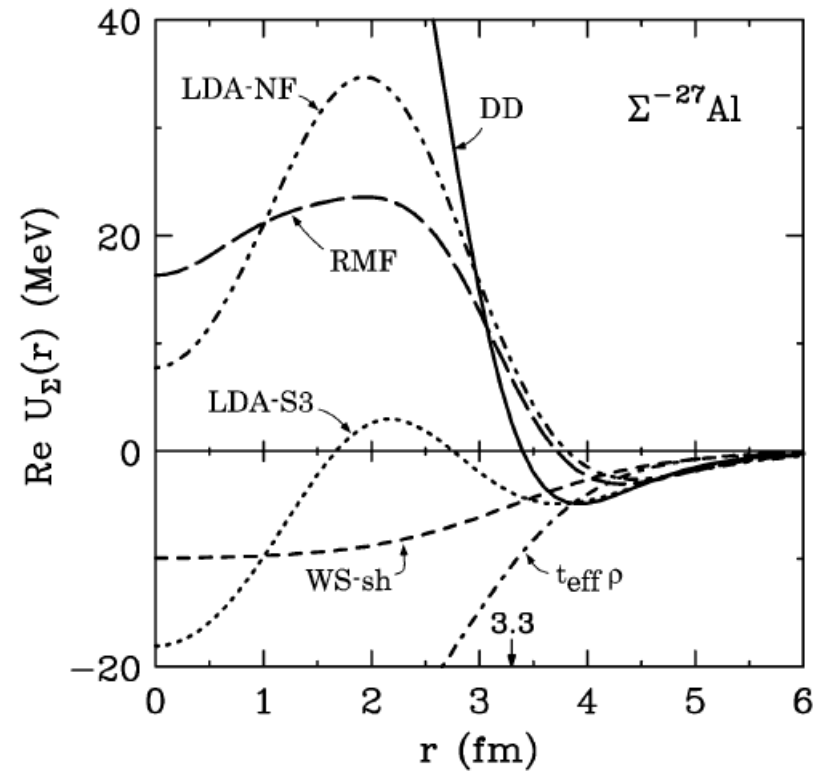
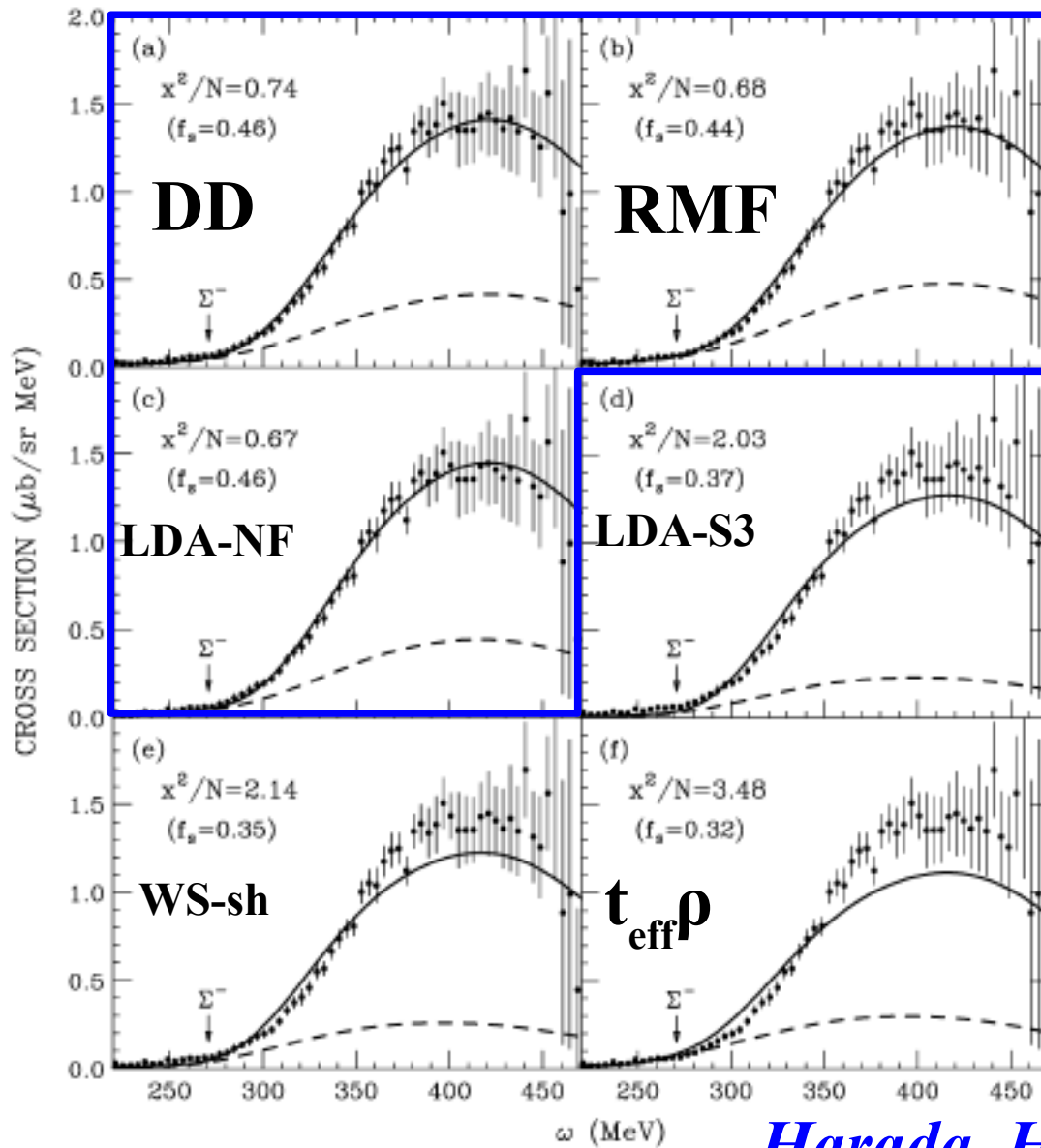
- Σ potential in nuclei

- Isoscalar part: 15-35 MeV repulsion
- Isovector part: 20-30 % of SU(3) value

Σ production in nuclei

$\chi^2/\text{DOF} < 1$

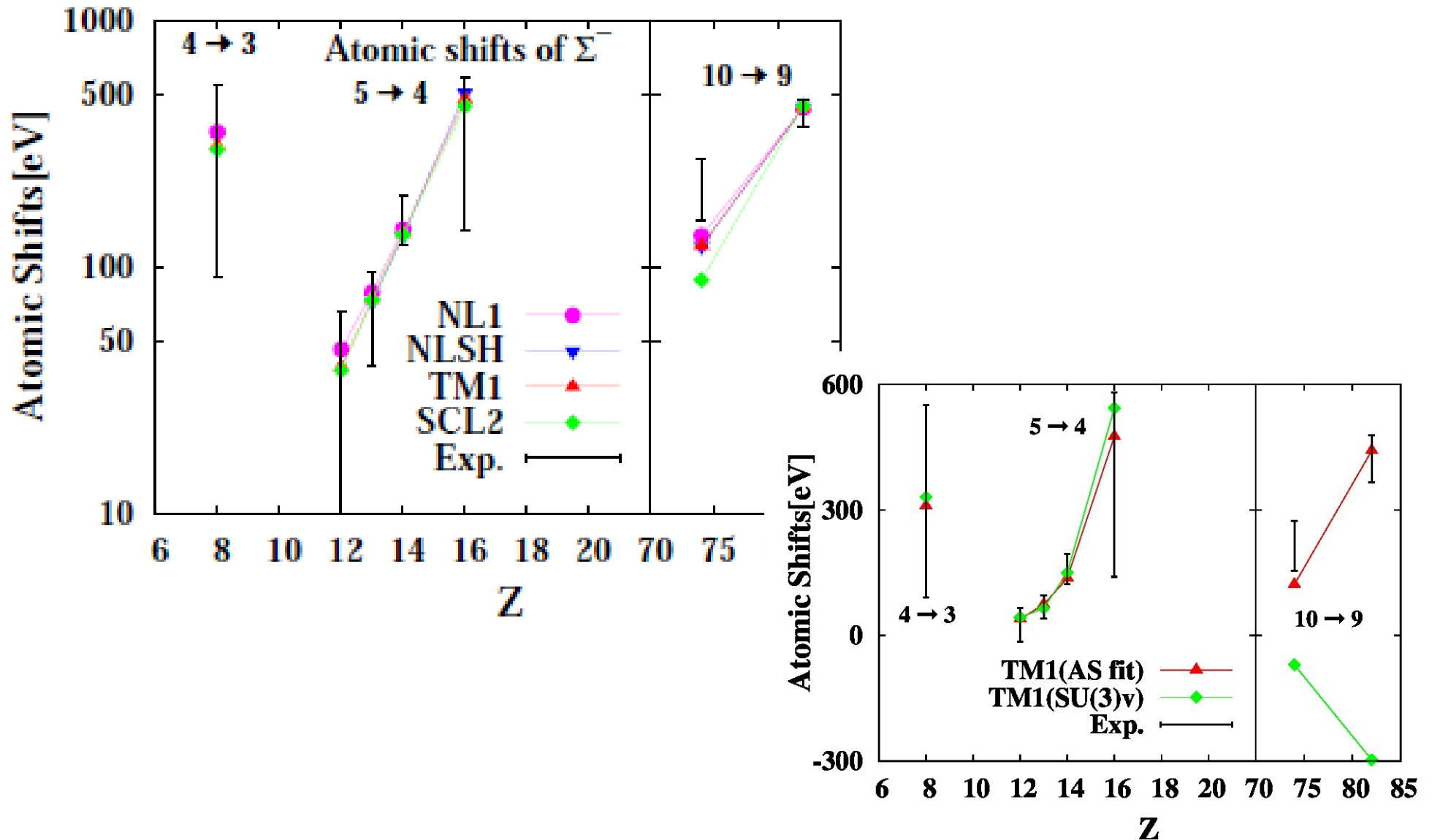
$^{28}\text{Si}(\pi^-, \text{K}^+)$



Harada, Hirabayashi ('05)

Data: Noumi et al. ('02); Saha et al. ('04)

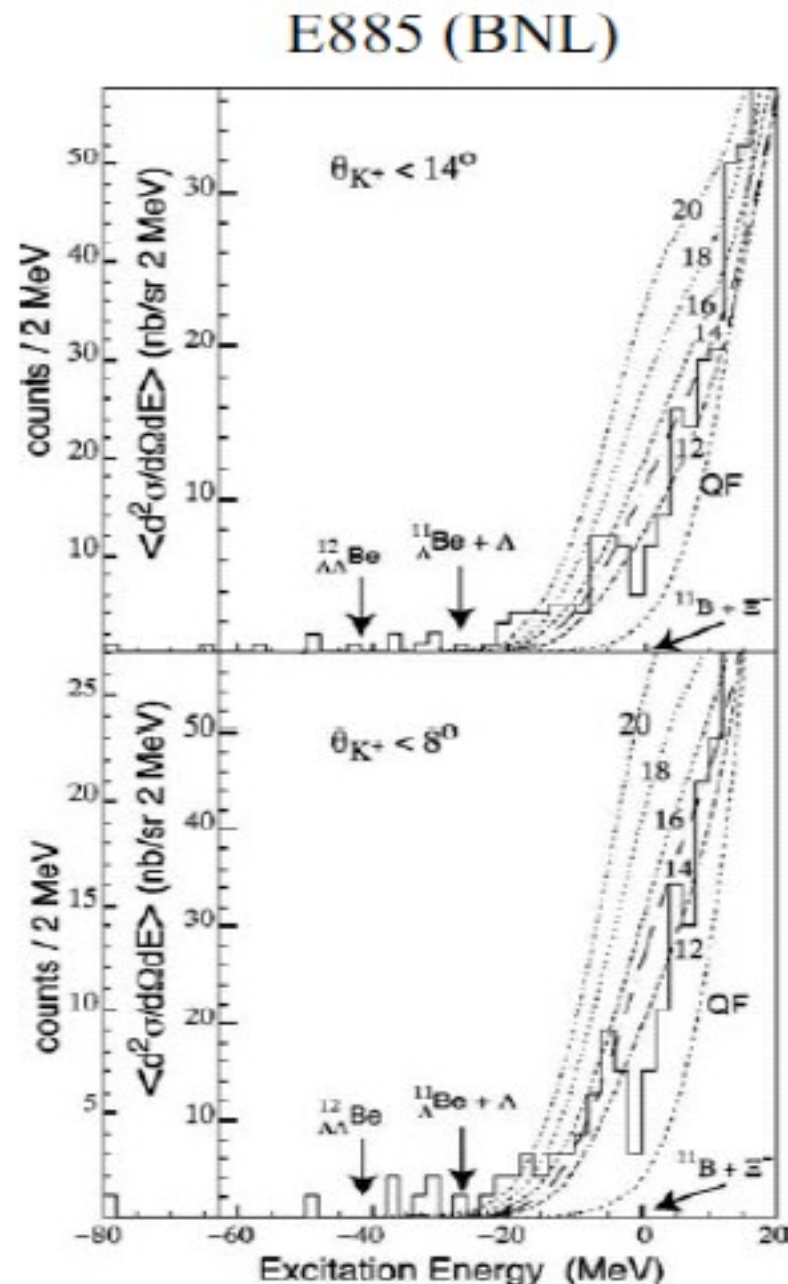
Σ^- atomic shift



Compiled data: Mares, Friedman, Gal, Jennings ('95)
 Calc.: Tsubakihara, Harada, AO: arXiv:1402.0979

Ξ hypernuclear formation

- Missing mass spectroscopy
BNL E885 $^{12}\text{C}(\text{K}^-, \text{K}^+)$
Fukuda et al. PRC58('98),1306;
Khaustov et al. PRC61('00), 054603.
 - No clear bound states found
- Twin hypernuclear formation
Aoki et al. PLB355('95),45.
- Potential depth
 $U_{\Xi} \sim -14 \text{ MeV}$



Where is the $S=-2$ dibaryon ($uuddss$) “H” ?

- Jaffe's prediction (1977)
 - 80 MeV below $\Lambda\Lambda$
 - (strong attraction from color mag. int.)

- Double hypernuclei ${}_{\Lambda\Lambda}{}^6\text{He}$ (Nagara)
 - No deeply bound “H”

- Resonance or Bound “H” ?

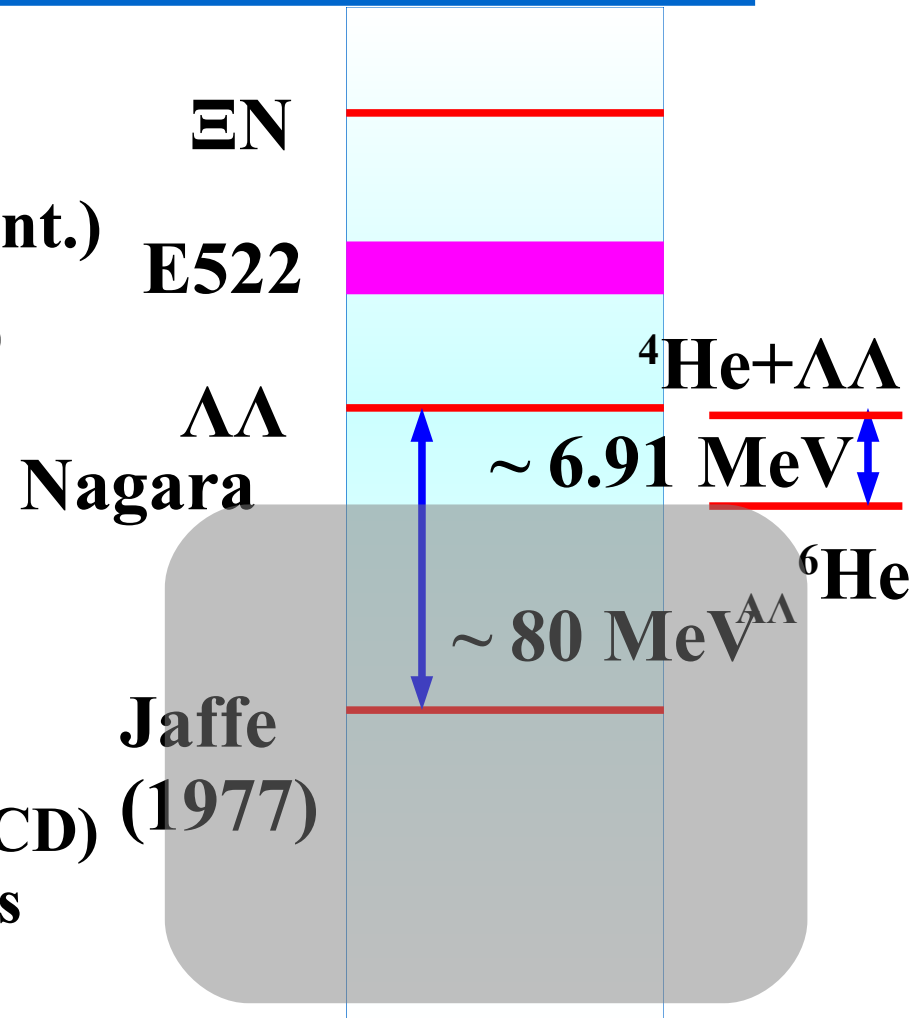
- KEK-E522 (Yoon et al., ('07))
 - “bump” at $E_{\Lambda\Lambda} \sim 15$ MeV

- Lattice QCD (HAL QCD & NPLQCD) (1977)
 - bound H at large ud quark mass

- How about HIC ?

- RHIC & LHC = Hadron Factory including Exotics
- “H” would be formed as frequently as stat. model predicts.

*Cho, Furumoto, Hyodo, Jido, Ko, Lee, Nielsen, AO, Sekihara, Yasui, Yazaki
(ExHIC Collab.), PRL('11)212001; arXiv:t:1107.1302*



Nagara event

■ ${}_{\Lambda\Lambda}{}^6\text{He}$ hypernuclei

Takahashi et al., PRL87('01)212502

(KEK-E373 experiment)

Lambpha

$$m({}_{\Lambda\Lambda}{}^6\text{He}) = 5951.82 \pm 0.54 \text{ MeV}$$

$$B_{\Lambda\Lambda} = 7.25 \pm 0.19_{-0.11}^{+0.18} \text{ MeV}$$

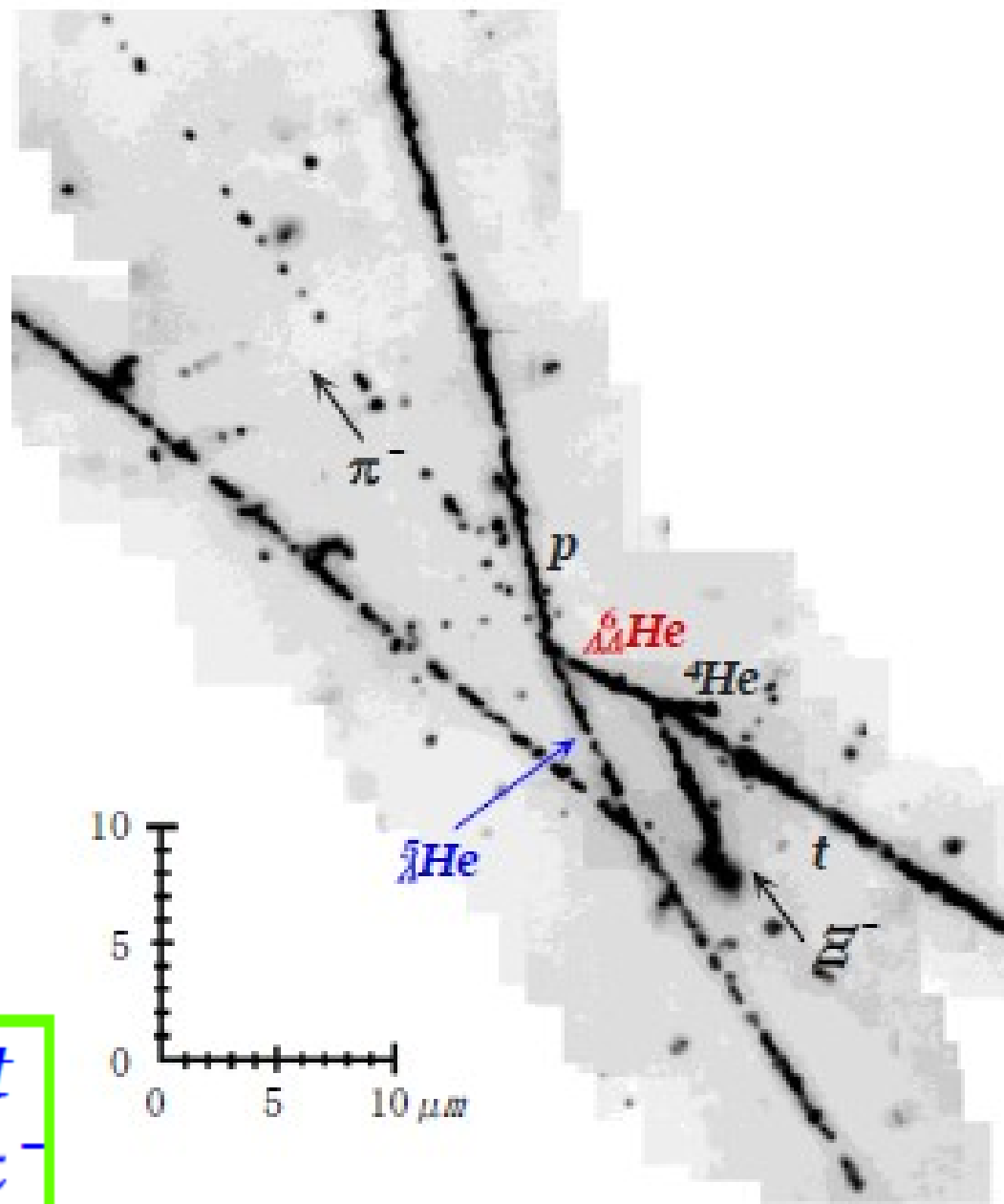
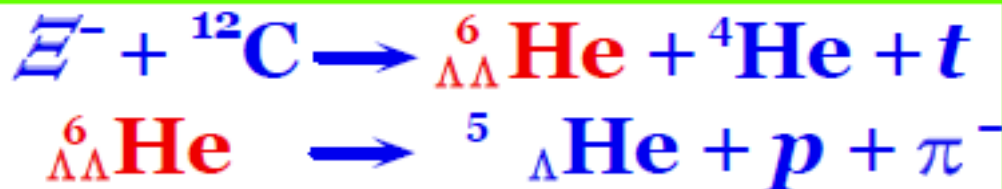
$$\Delta B_{\Lambda\Lambda} = 1.01 \pm 0.20_{-0.11}^{+0.18} \text{ MeV}$$

(assumed $B_{\Xi^-} = 0.13 \text{ MeV}$)

$$\rightarrow B_{\Lambda\Lambda} = 6.91 \text{ MeV}$$

(PDG modified(updated)

Ξ^- mass)



$\Lambda\Lambda$ interaction models

■ Boson exchange potentials

- Nijmegen potentials: various versions

Rijken et al., ('77-'10)

Hard core: Nijmegen model D & F (ND, NF)

Soft core: Nijmegen soft core '89 & '97 (NSC89, NSC97)

Extended soft core: ESC08

- Ehime potential: would be too attractive.

Ueda et al., ('98)

Ehime fits old double Λ hypernucl. data, $\Delta B_{\Lambda\Lambda} = 4$ MeV

■ Quark cluster model

- fss2

Fujiwara, Kohno, Nakamoto, Suzuki ('01)

Short range repulsion from quark Pauli blocking & OGE

Core is softer due to non-locality

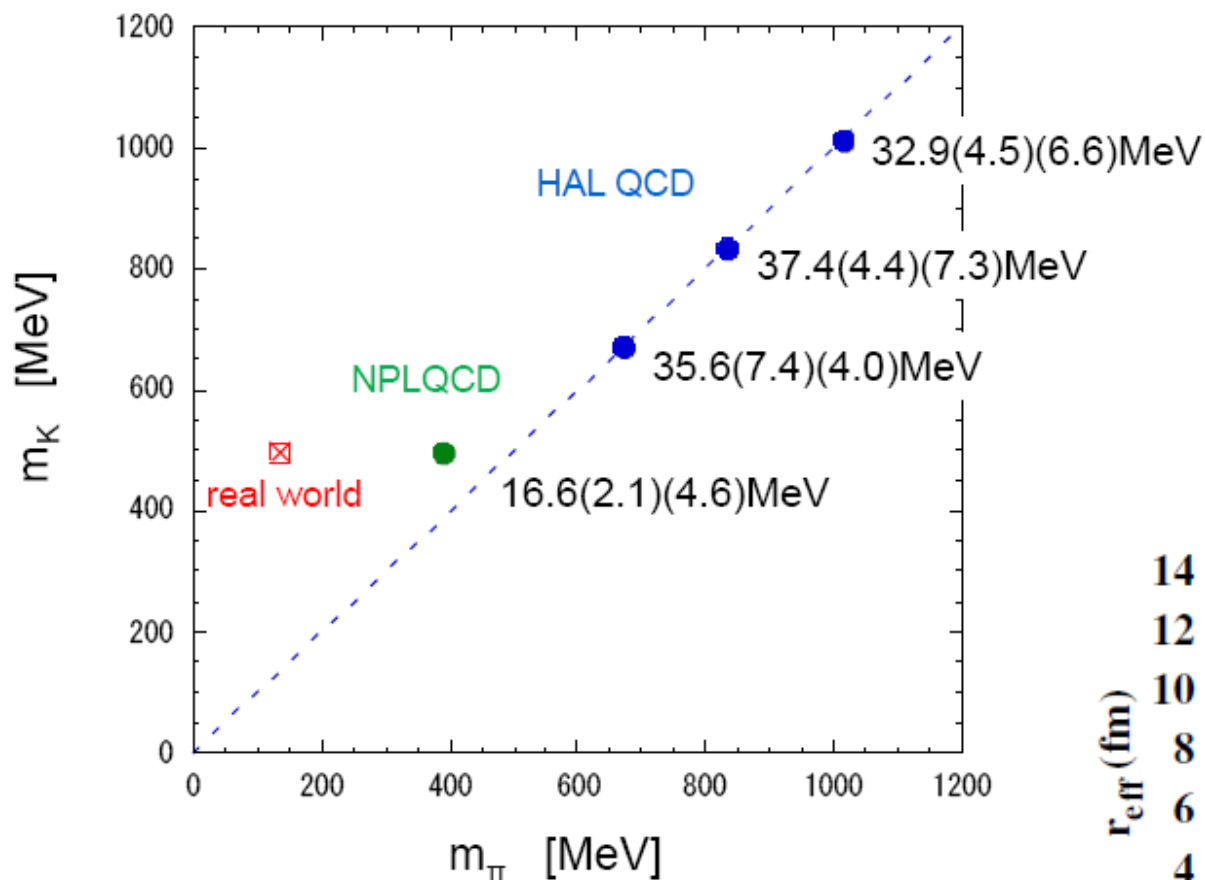
■ Modified Nijmegen potentials fitting Nagara data.

Filikhin, Gal ('02), Hiyama et al.('02)

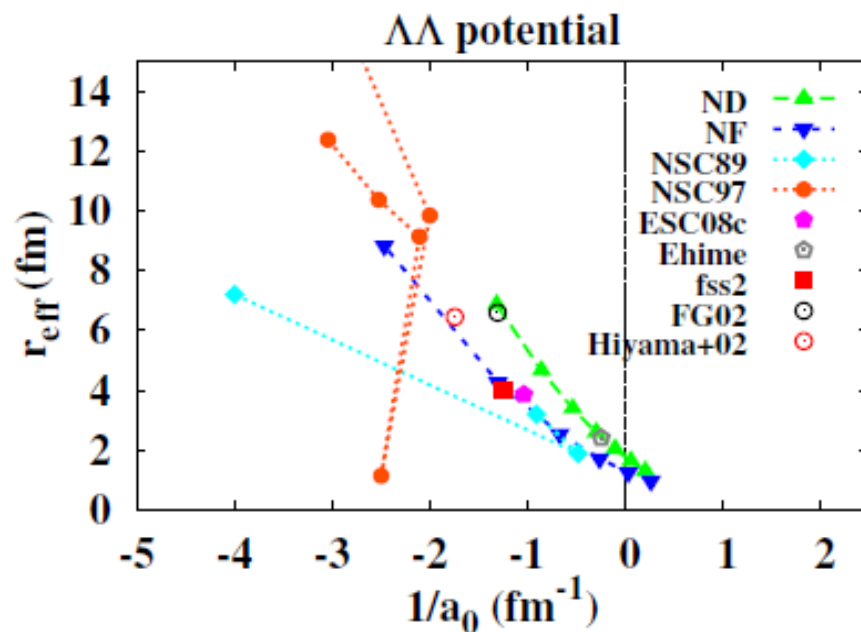
- Potential Fitting Nagara data $\Delta B_{\Lambda\Lambda} = 1.0$ MeV

Lattice QCD predicts bound “H”

“H” bounds with heavy π ($M_\pi > 400$ MeV)



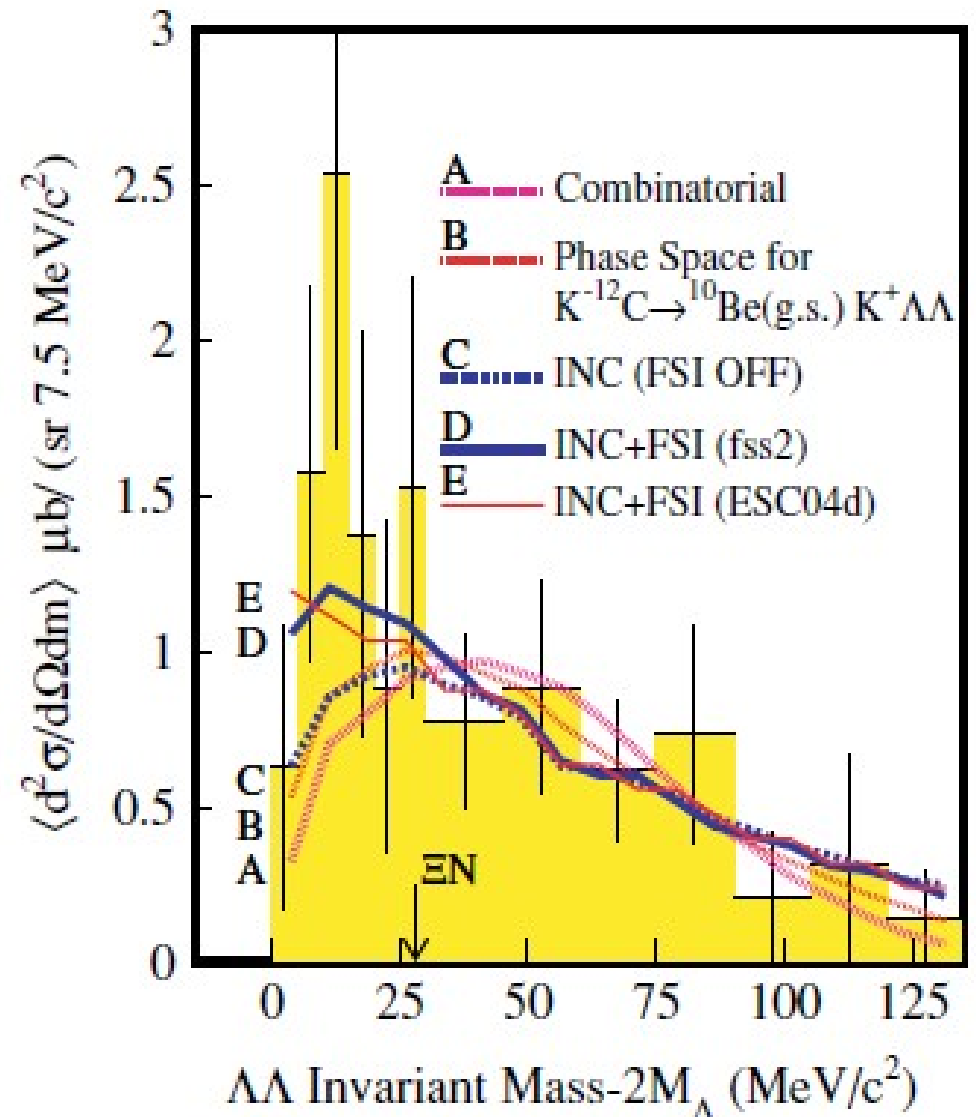
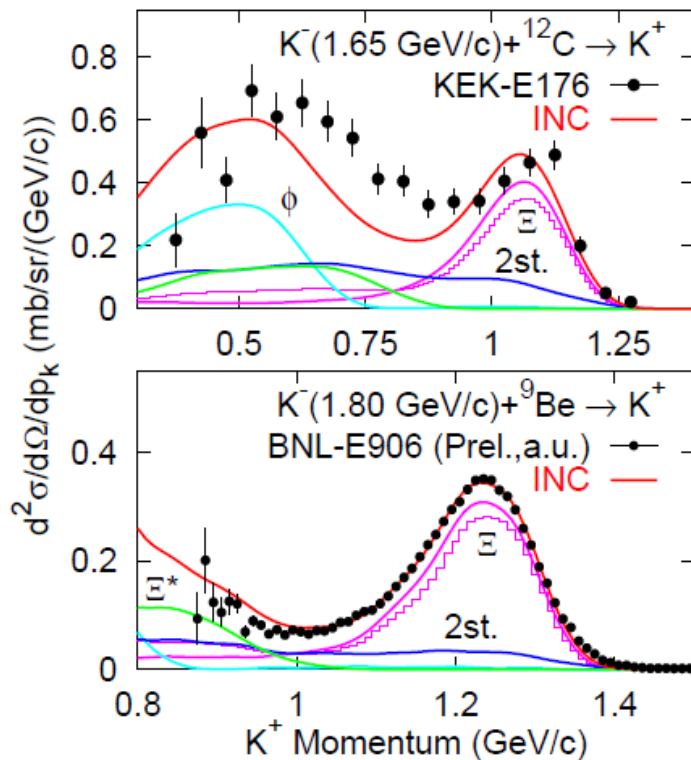
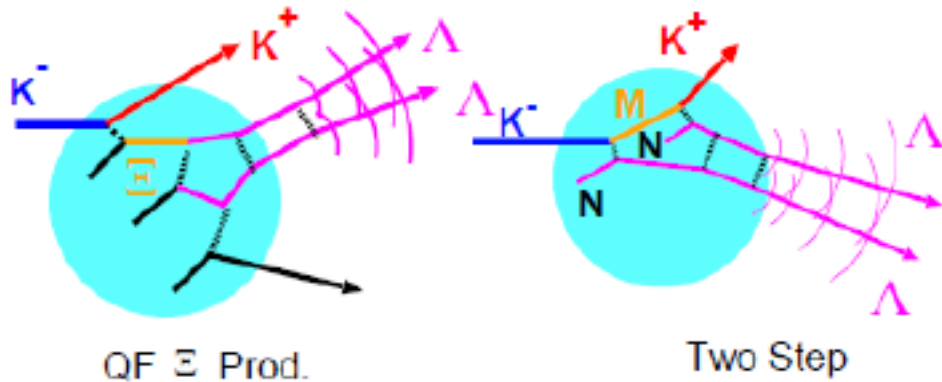
NPLQCD Collab., PRL 106 (2011) 162001;
HAL QCD Collab., PRL 106 (2011) 162002



Models


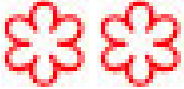


$\Lambda\Lambda$ correlation from $(K^-, K^+ \Lambda\Lambda)$ reaction

- Enhancement at $\sim 2 M(\Lambda) + 10$ MeV,



C.J. Yoon, ..., (KEK-E522), AO, PRC75 (2007) 022201(R)
J. K. Ahn et al. (KEK-E224).

“Stars” of Hyperon Potentials (A la Michelin)

- $U_{\Lambda}(\rho_0) \sim -30 \text{ MeV}$ 
 - *Bound State Spectroscopy + Continuum Spectroscopy*
- $U_{\Sigma}(\rho_0) > +15 \text{ MeV}$ 
 - Continuum (Quasi-Free) spectroscopy
 - Atomic shift data (attractive at surface) should be respected.
- $U_{\Xi}(\rho_0) \sim -14 \text{ MeV}$ 
 - No confirmed bound state, No atomic data, High mom. transf., \rightarrow Small Potential Deps.
 - Continuum low-res. spectrum shape $\rightarrow -14 \text{ MeV}$
- $V_{\Lambda\Lambda}$: Weakly attractive. 

But these potentials lead to collapse of massive NS



*Toward the Solution
of Massive Neutron Star puzzle*

Massive Neutron Star Puzzle

- Observation of massive neutron stars ($M \sim 2 M_{\odot}$)

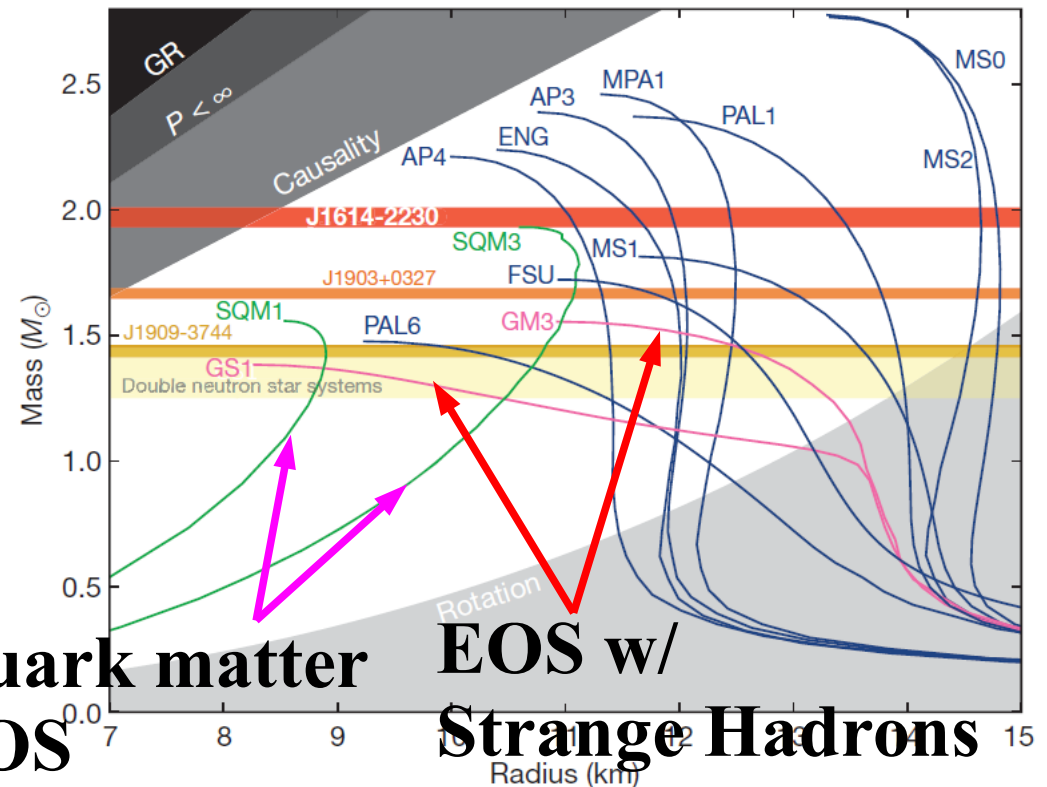
- PSR J1614-2230 (NS-WD binary), $1.97 \pm 0.04 M_{\odot}$

Demorest et al., Nature 467('10)1081 (Oct.28, 2010).

”Kinematical” measurement (Shapiro delay, GR)
+ large inclination angle

- PSR J0348+0432 (NS-WS binary), $2.01 \pm 0.04 M_{\odot}$

Antoniadis et al., Science 340('13)1233232.



No Exotics in NS ?

Possible Solutions to Massive NS puzzle

■ Proposed “Solutions” of Massive NS puzzle

● Modification of YN interaction

*Weisenborn, Chatterjee, Schaffner-Bielich ('11); Jiang, Li, Chen ('12);
Tsubakihara, AO ('13)*

● Introducing BBB repulsion

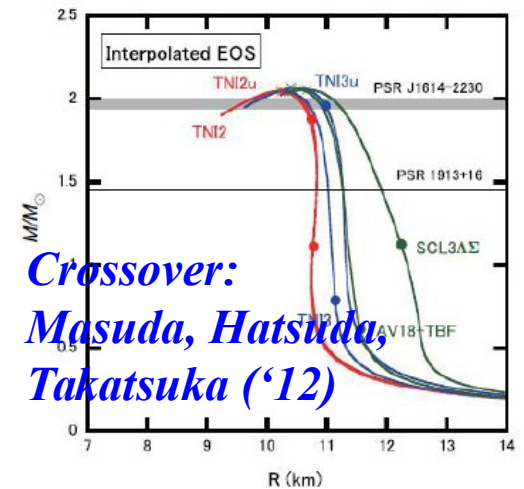
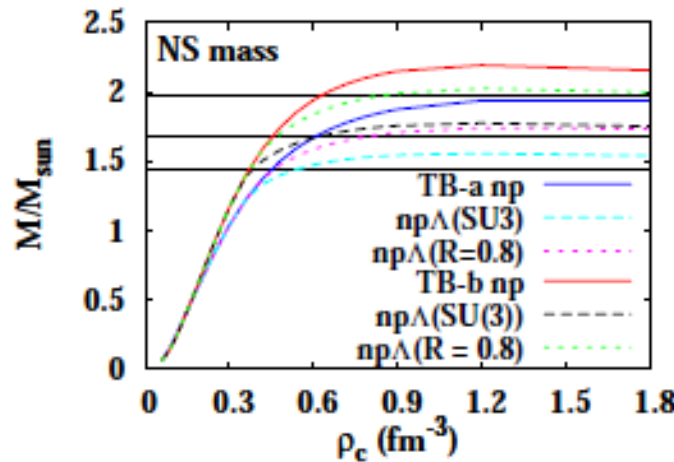
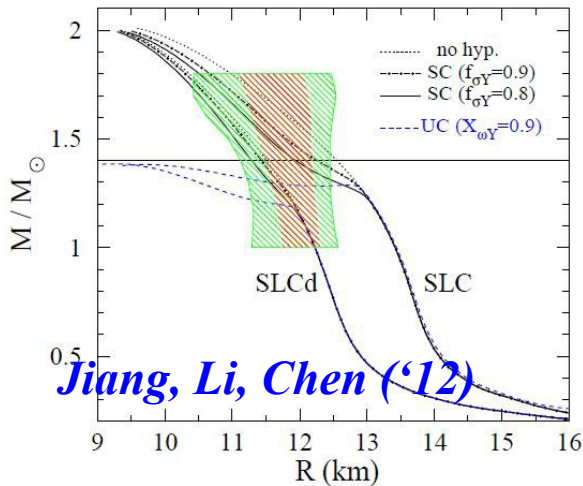
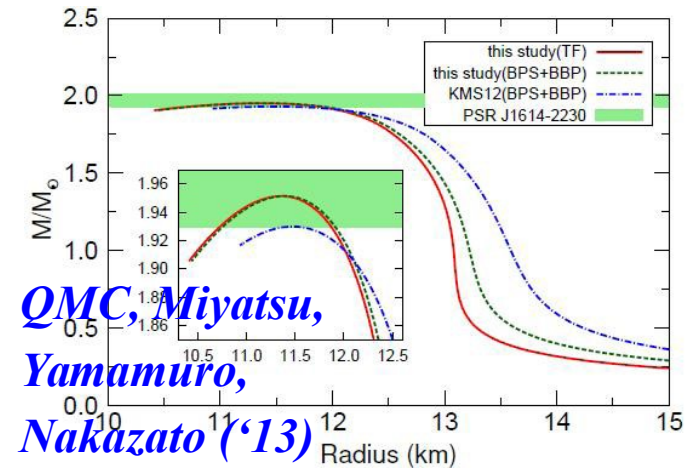
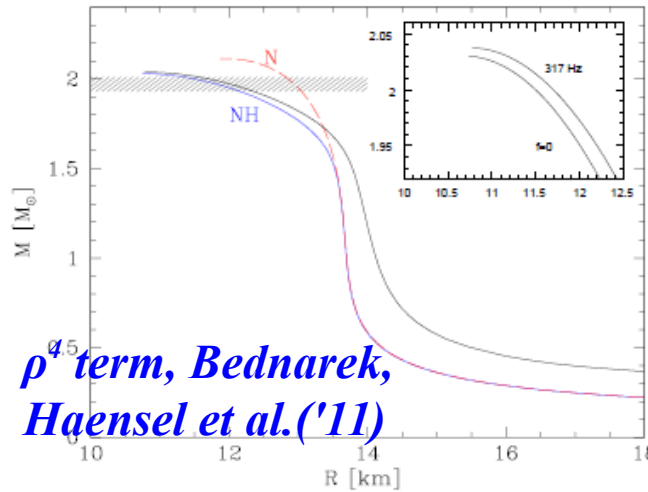
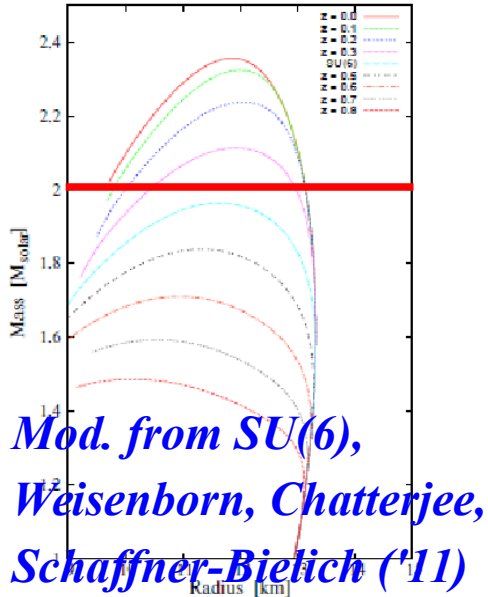
*Bednarek, Haensel et al.('11); Miyatsu, Yamamuro, Nakazato ('13);
Tsubakihara, this session.*

● Early crossover transition to quark matter

Masuda, Hatsuda, Takatsuka ('12)

● Choose Stiff EOS for nuclear matter *Tsubakihara, AO ('14)*

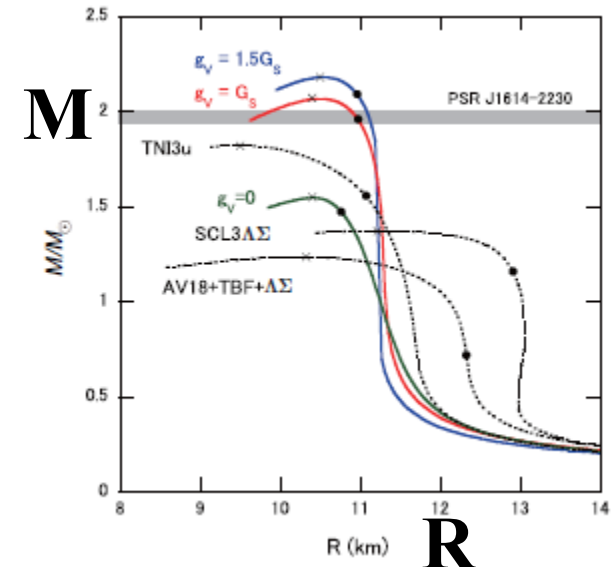
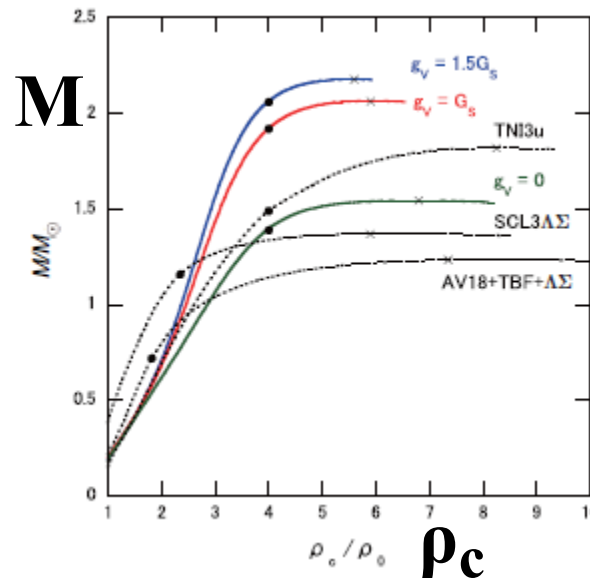
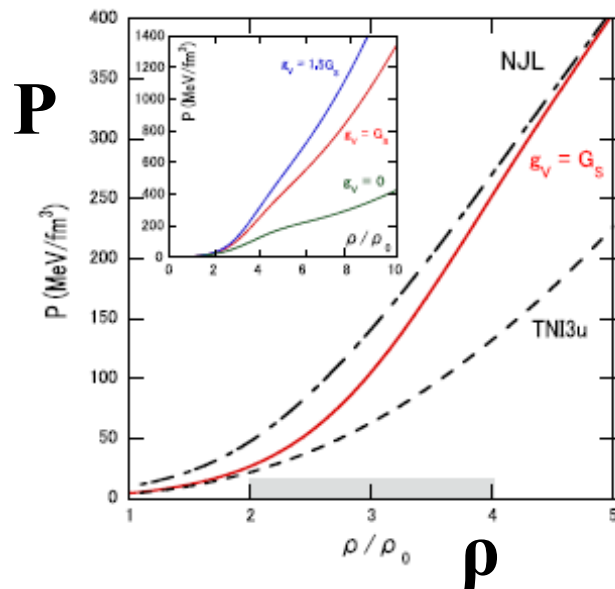
NS matter EOS with hyperons



*These are phenomenological “solutions”.
How can we examine them ?*

Early crossover transition to quark matter

- A possible solution of the massive NS puzzle = quark matter
K. Masuda, T. Hatsuda, T. Takatsuka, ApJ764('13)12
 - With large vector qq coupling, finite density QCD phase transition can be crossover.
 - Crossover transition → EOS can be stiffer and can support $2 M_{\odot}$
 - Transition density = $(2-4) \rho_0$.
- Is it consistent with heavy-ion collisions ?



Masuda, Hatsuda, Takatsuka ('13)

A. Ohnishi @ JAEA, June 2, 2015 79

Possible Solutions to Massive NS puzzle

■ Proposed “Solutions” of Massive NS puzzle

- **Modification of YN interaction**

Weisenborn, Chatterjee, Schaffner-Bielich ('11); Jiang, Li, Chen ('12); Tsubakihara, AO ('13)

- **Introducing BBB repulsion**

Bednarek, Haensel et al.('11); Miyatsu, Yamamuro, Nakazato ('13).

- **Early crossover transition to quark matter**

Masuda, Hatsuda, Takatsuka ('12)

- **Choose Stiff EOS for nuclear matter** *Tsubakihara, AO ('14)*

■ What is necessary to solve the massive NS puzzle ?

- **EOS of nucleon matter need to be precisely settled.**

- **Yet un-explored YN & YY interactions**

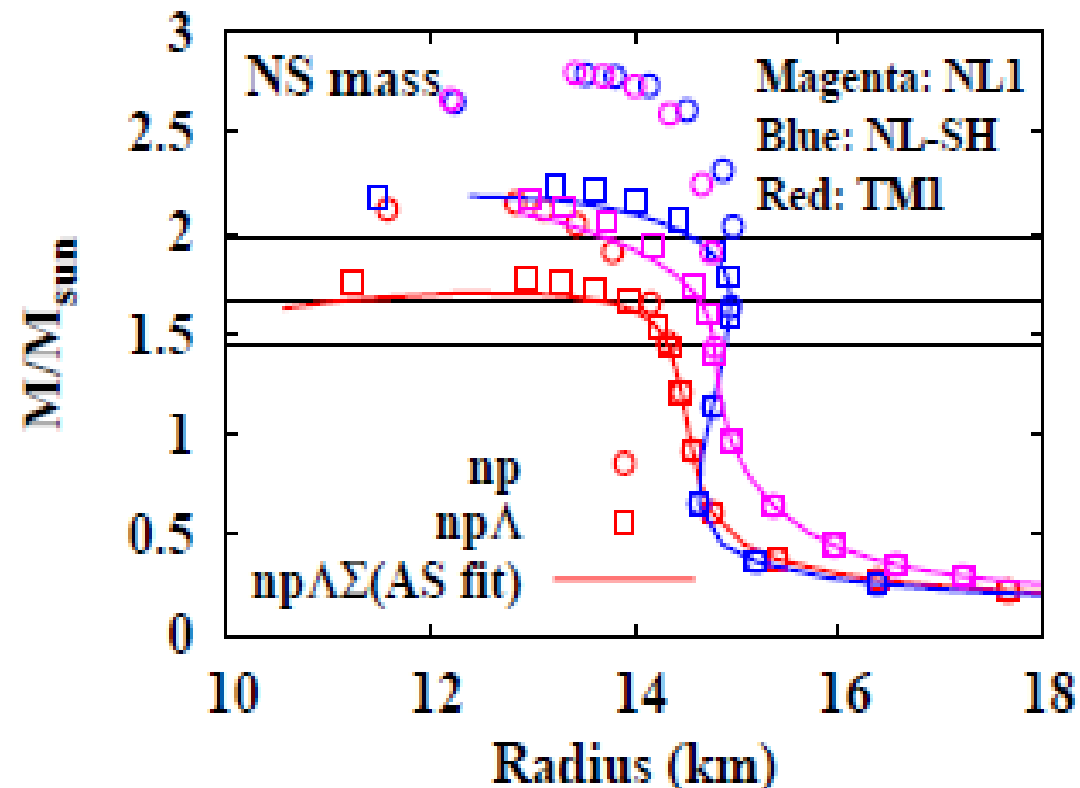
- **Three-body interaction including hyperons (YNN, YYN, YYY) and its effects on EOS**

- **Finding onset density of quark matter**

Massive Neutron Stars with Hyperons

Tsubakihara, Harada, AO, arXiv:1402.0979

- Ruled-out EOS with hyperons = GM3
Glendenning & Moszkowski (1991)
- We did NOTHING special and find $2 M_{\odot}$ NS can be supported.
 - “Typical” RMF for nucl. matter
NL1, NL-SH, TM1
*Reinhardt et al. ('86);
Sharma, Nagarajan, Ring ('93);
Sugahara, Toki ('94).*
 - ss mesons are introduced
 - Hypernuclear data
 Λ , $\Lambda\Lambda$ hypernuclei
 Σ atomic shifts
SU(3) relation to isoscalar
-vector couplings



Yet Un-explored ΞN & ΛN Interactions

■ ΞN interaction

- Indirect evidence of $U_{\Xi}(\rho_0) \sim -14$ MeV

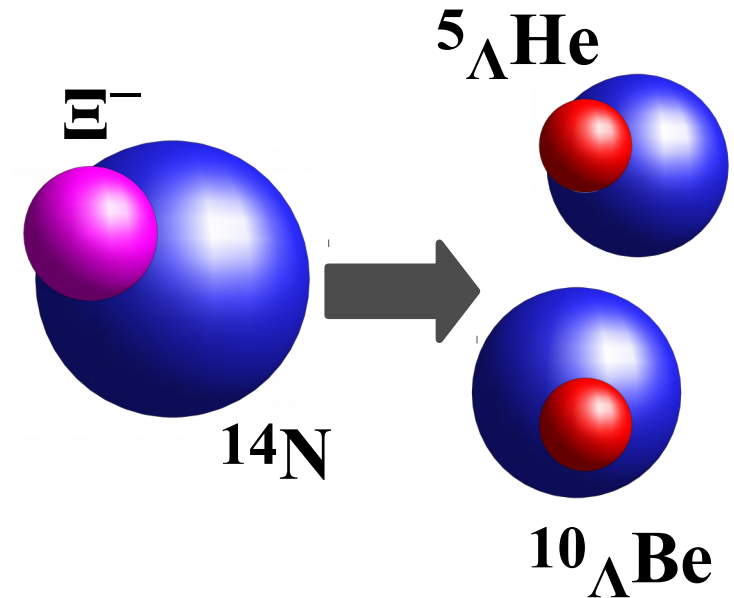
- First evidence of deeply bound

Ξ hypernucleus $\Xi^- + {}^{14}\text{N}$

K. Nakazawa et al., PTEP 2015, 033D02.

$B(\Xi^-) = 4.38 \pm 0.25$ MeV (g.s.)

or 1.11 ± 0.25 MeV (exc.)



■ ΛN interaction

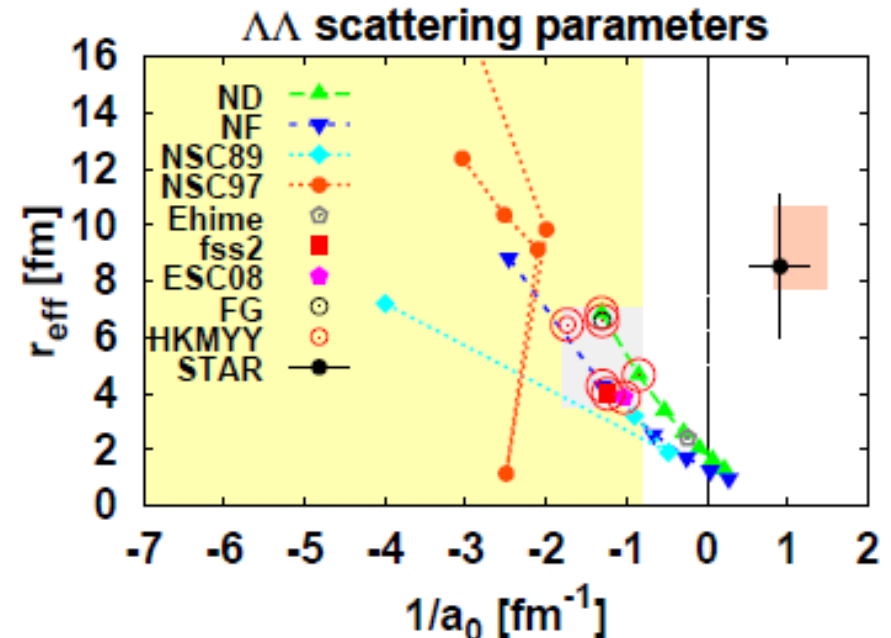
- Double hypernuclear bond energy

$\Delta B_{\Lambda\Lambda}({}^6_{\Lambda\Lambda}\text{He}) = 0.67 \pm 0.16$ MeV

- ΛN int. from two-particle corr. from heavy-ion collisions

$-1.25 \text{ fm} < a_0(\Lambda N) < 0$

■ Isospin dependent part of ΣN



K. Morita, T. Furumoto, AO, PRC91('15),024916

A. Ohnishi @ JAEA, June 2, 2015 82

Hadron-Hadron correlation in HIC

- **Correlation function formula** *Bauer, Gelbke, Pratt ('92); Lednicky ('09).*

$$C(q) = \int d\mathbf{x}_{12} \underbrace{S(\mathbf{x}_{12})}_{\text{Source}} \underbrace{|\Psi(\mathbf{x}_{12})|^2}_{\text{wave fn.}}$$

- **Free boson + Gaussian source**
= Hanbury-Brown & Twiss effect

$$C(q) = 1 + \exp(-4q^2 R^2)$$

- **Free fermion + Gaussian source**

$$C(q) = 1 - \frac{1}{2} \exp(-4q^2 R^2)$$

- **Correlation fn. has info. both on source and w.f. (~ int.)**

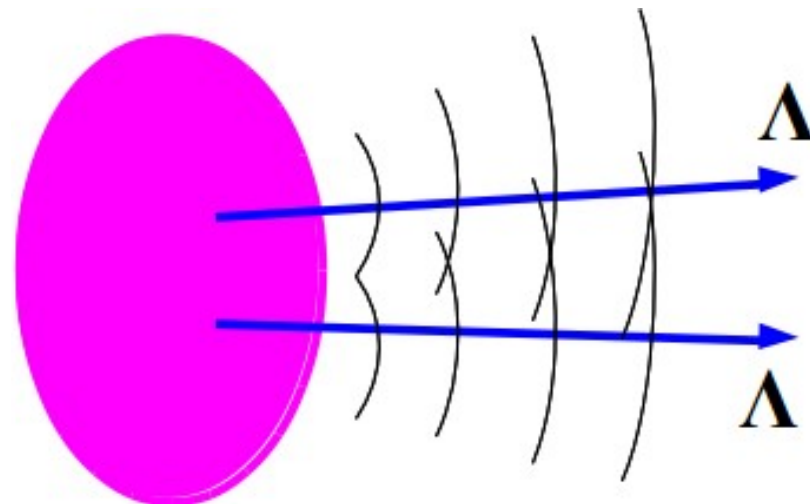
- **$\Lambda\Lambda$ correlation measurement**

- **(K⁻,K⁺) reaction** *C.J.Yoon et al. (KEK-E522)('07); J.K.Ahn et al. (KEK-E224); AO, Hirata, Nara, Shinmura, Akaishi ('01).*

- **Heavy-ion collisions**

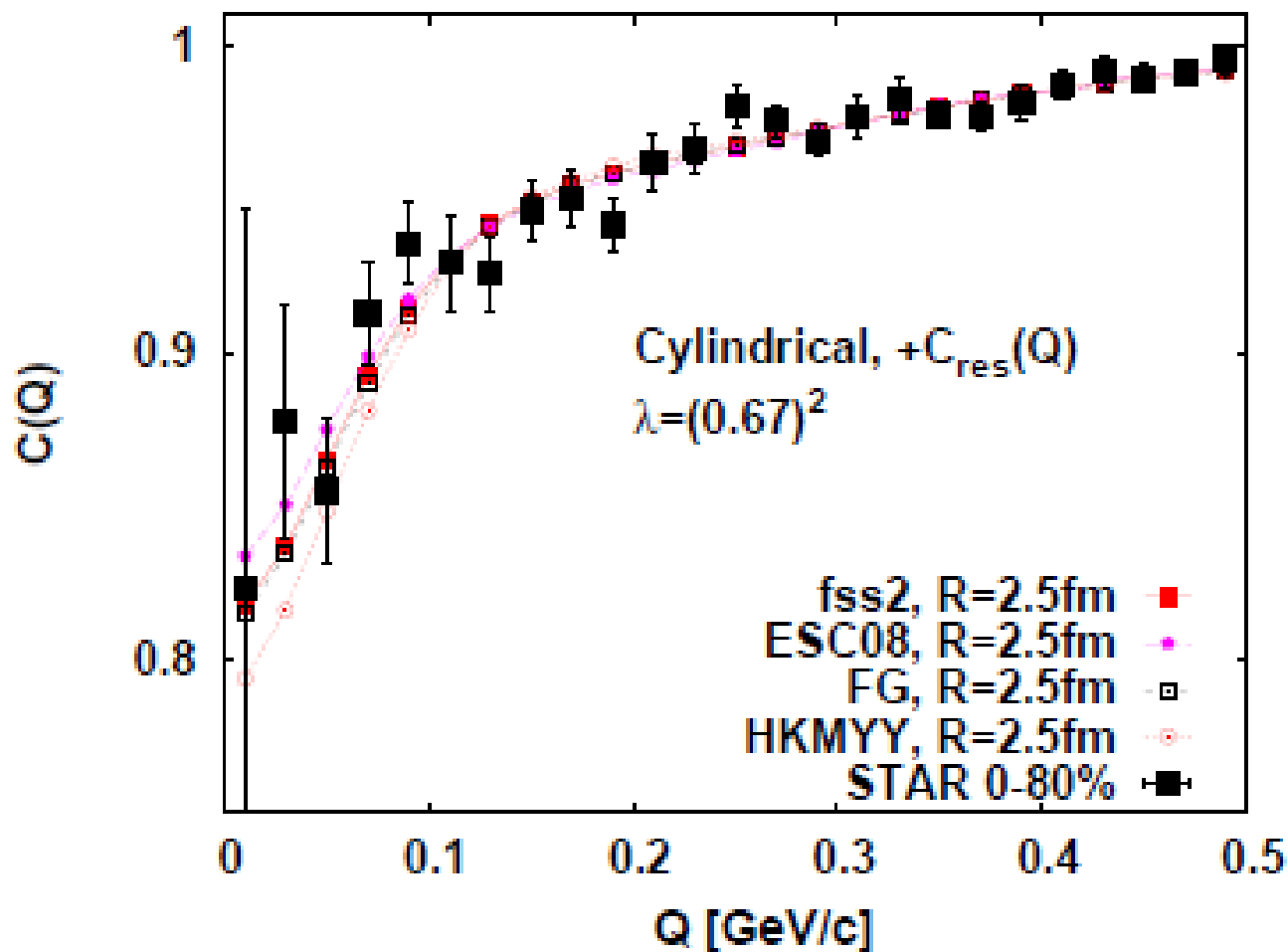
STAR collab. arXiv:1408.4360;

C. Greiner, B. Muller ('89); AO, Hirata, Nara, Shinmura, Akaishi ('01).



$\Lambda\Lambda$ correlation and favored $\Lambda\Lambda$ interaction

$\Lambda\Lambda$ correlation with long. and transverse flow effects, Σ^0 feed down, and unknown long tail effects



*K.Morita, T.Furumoto, AO, PRC91('15)024916 [arXiv:1408.6682]
Data: Adamczyk et al. (STAR Collaboration), PRL 114 ('15) 022301.*

Do we see $\Lambda\Lambda$ interaction ?

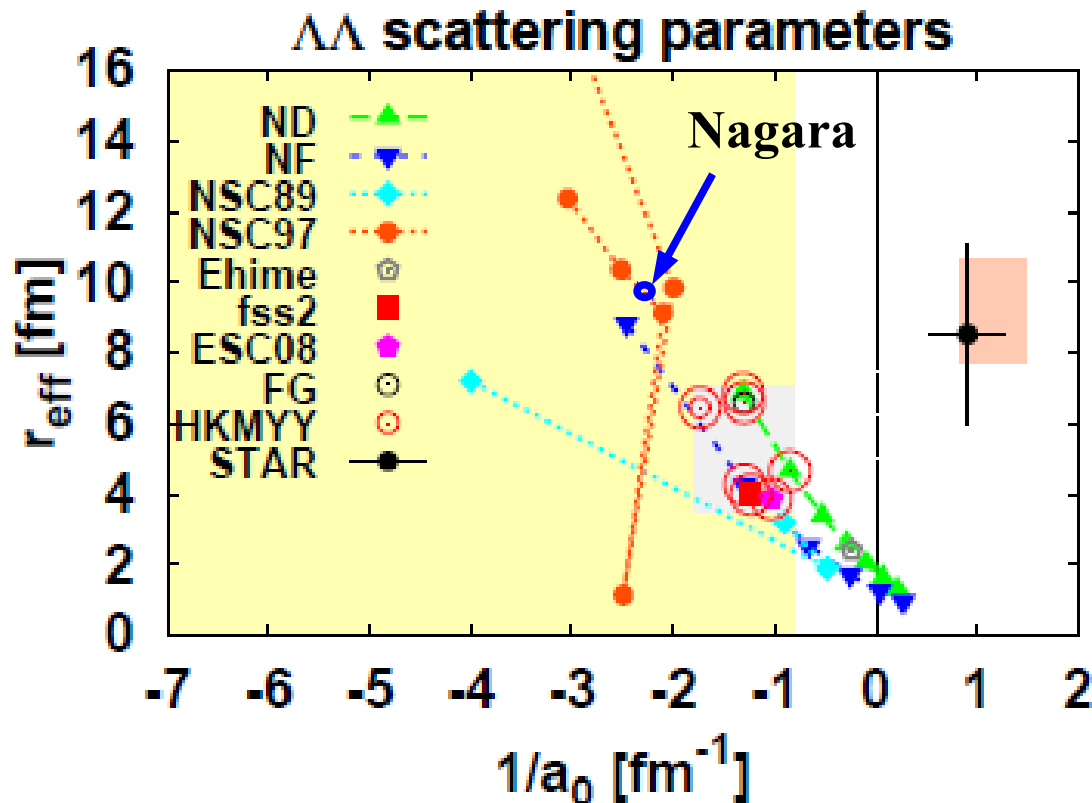
Do we see diff. btw

$V_{\Lambda\Lambda}$ from RHIC (\sim vacuum $\Lambda\Lambda$ int.) and $V_{\Lambda\Lambda}$ from Nagara ?

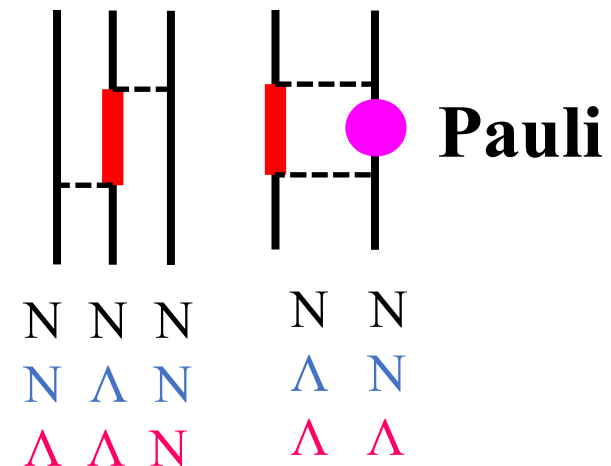
Hiyama et al. ('02); Filikhin, Gal ('02)

Mechanism: Pauli blocking in the intermediate ΞN channel

Kohno ('13) / Myint, Shinmura, Akaishi ('03) / Nishizaki, Takatsuka, Yamamoto('02) / Machleidt.



Morita et al. ('15)



BBB interaction including Hyperons

- BBB int. incl. YNN, YYN and YYY should exist and contribute to EOS.

Nishizaki, Takatsuka, Yamamoto ('02)

- Chiral EFT, Multi-Pomeron exch., Quark Pauli, Lattice 3BF, SJ, ..

Kohno('10); Heidenbauer+('13);

Yamamoto+('14); Nakamoto, Suzuki;

Doi+(HALQCD,'12); Tamagaki('08); ...

- Quant. MC study

D.Lonardonni, S.Gandolfi, F.Pederiva.('13)

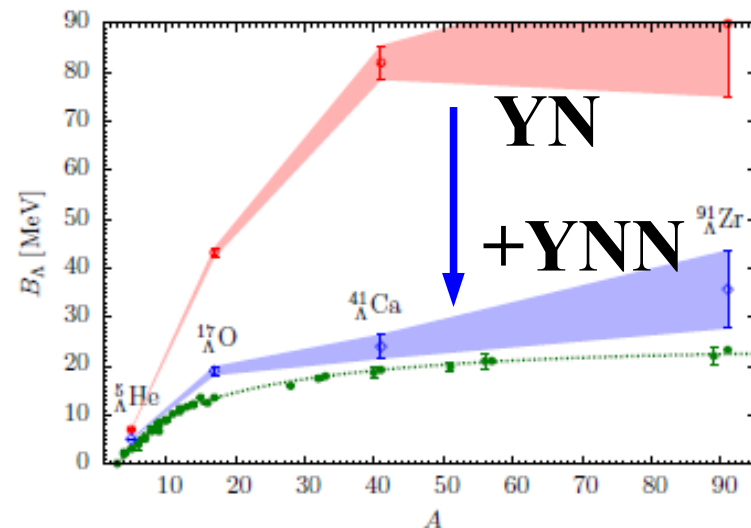
- Quark Meson Coupling

Miyatsu et al.; Thomas (HHIQCD)

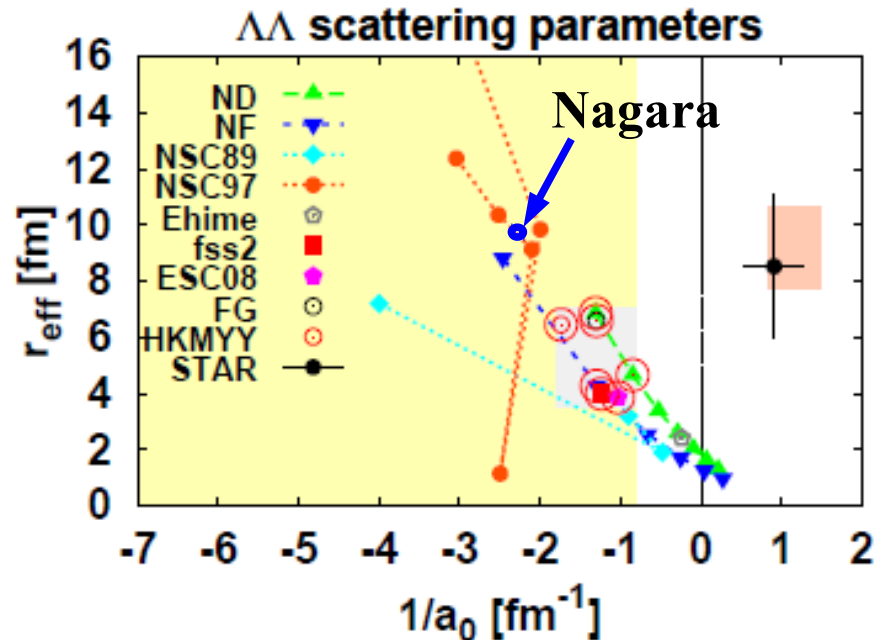
- $\Lambda\Lambda$ *K. Morita, T. Furumoto, AO,*

PRC91('15)024916

- **Caveat: Missing data**
(or data precision is low ...)



Lonardonni et al.('14)



Morita et al. ('15)

NNN force from Lattice QCD

■ HAL QCD method for BB int.

Aoki, Hatsuda, Ishii ('07)

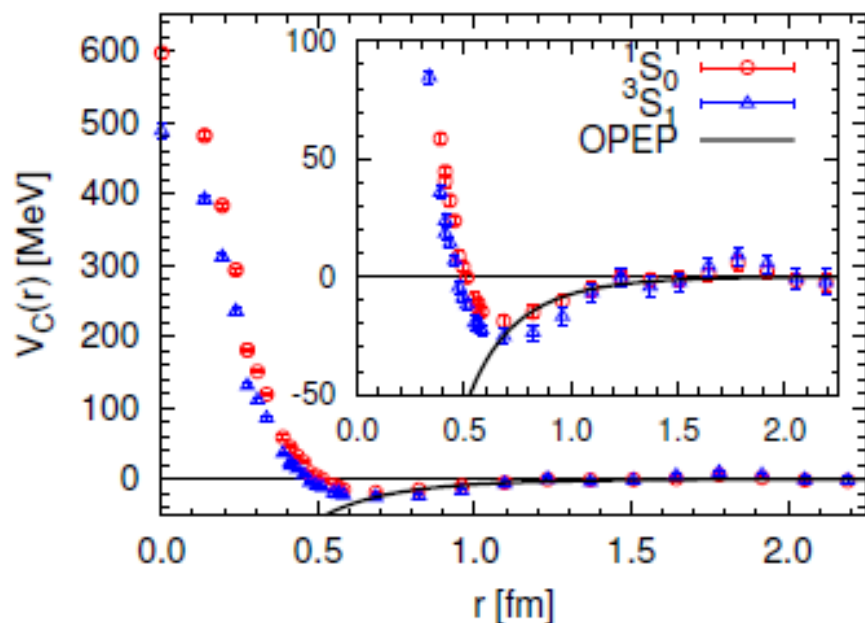
Nambu-Bethe-Salpeter amplitude \sim w.f.

\rightarrow NN force from Sch. Eq.

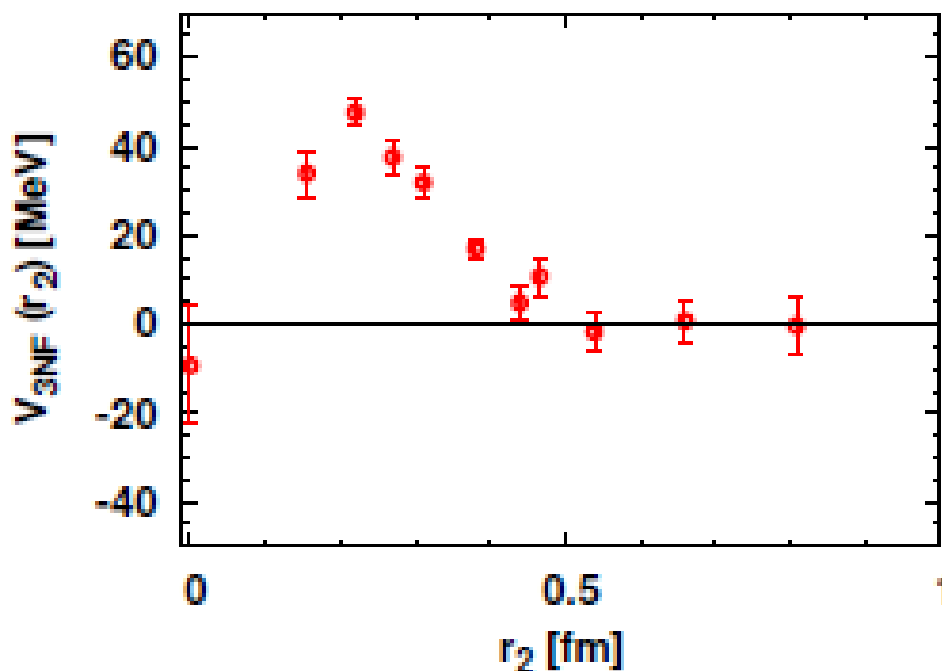
● Consistent with Luscher's method in asymptotic region

Luscher ('91), NPLQCD Collab. ('06, $\pi\pi$)

■ NNN force T. Doi (HAL QCD Collab.)('12)



Aoki, Hatsuda, Ishii ('07)



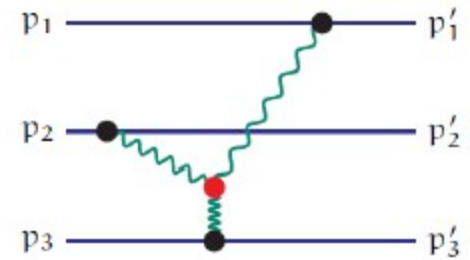
T. Doi et al. (HAL QCD Collab.) ('12)

BBB force incl. hyperons

Triple pomeron vertex model of BBB repulsion

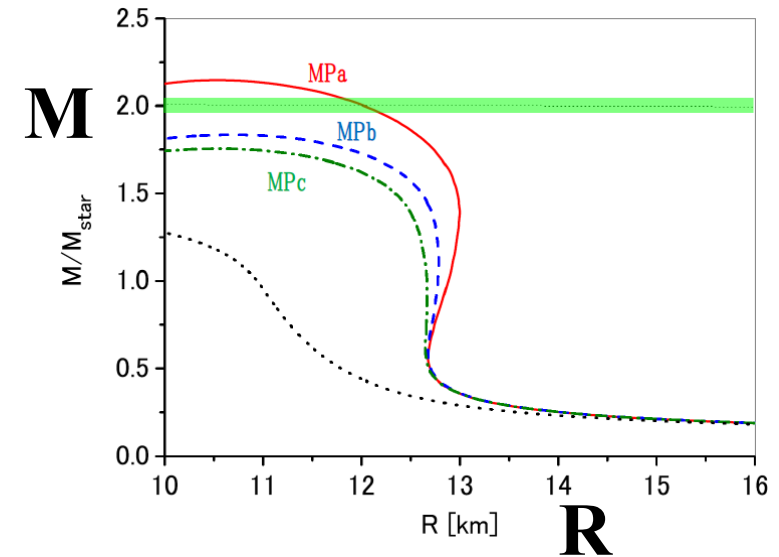
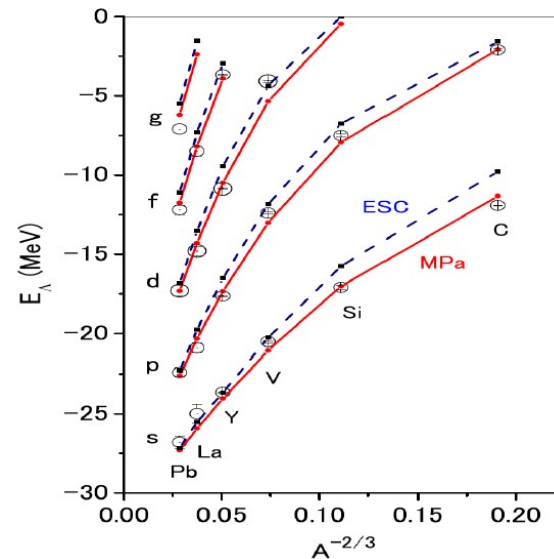
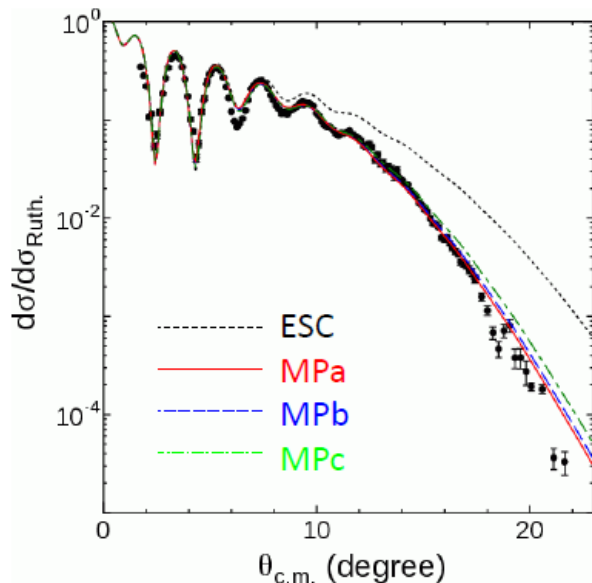
$$\mathcal{L}_{PPP} = g_{3P} \mathcal{M} \sigma_p^3(x)/3!$$

- Pomeron = gluon ladder, flavor blind, induces BB repulsion



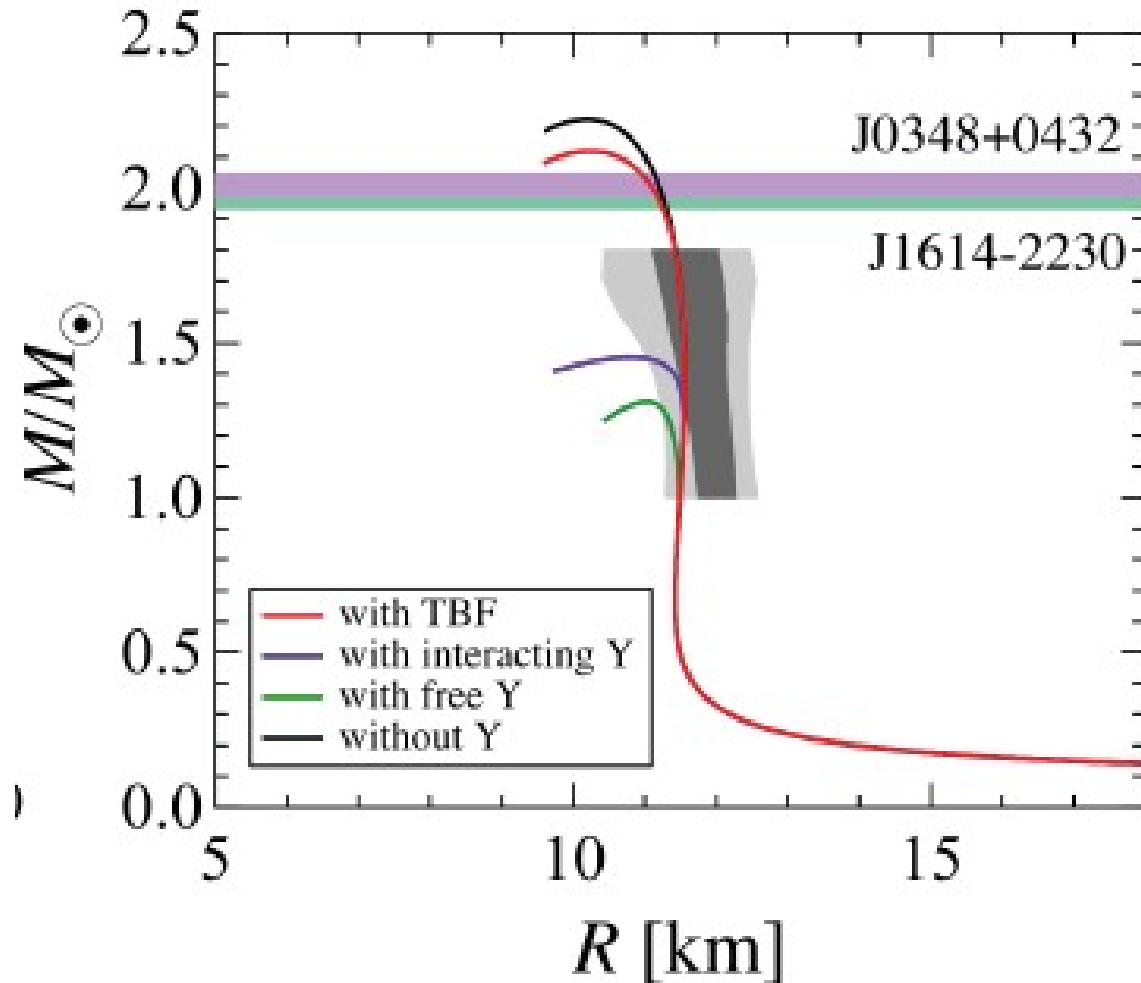
How can we fix BBB force strength ?

- Multi-pomeron coupling strength is determined in AA scattering.
- Multi-baryon force gives better S_Λ , and may support $2 M_\odot$ NS.



Yamamoto, Furumoto, Yasutake, Rijken ('13)

Variational Calculation including Hyperons



Togashi, Hiyama, Takano, Yamamoto

Summary of Part II

- **Strangeness nuclear physics has extracted the depths of hyperons in nuclei, and EOSs based on these potentials fail to support $2M_{\odot}$ NS. (Massive NS puzzle)**
- **Solving the massive NS puzzle is a big challenge in physics.**
 - **All relevant BB (and MB) interactions, Many-body theories, Multi-body interactions, and Transition to quark matter have to be understood properly.**
- **There are several attempts to answer the massive NS puzzle, but not convincing yet.**
 - **How can we justify “model assumptions” ?**
 - **How can we access various YY (and MB) interactions ?**
→ J-PARC experiments, heavy-ion collisions, ...
 - **How can we determine BBB (and BBBB) force ?**
→ Chiral EFT, Lattice QCD, Quark model (Nakamoto), ...
“Model” BBB force + phenomenology (precise S_{Λ} is necessary)

Summary

- **Neutron Star physics is attracting much attention, and many current/future facilities/projects are aiming at solving NS puzzles.**
 - **Radioactive beam facilities** → Sym. E. at $\rho < \rho_0$ and $\rho \sim (2-3) \rho_0$
 - **Hadron machines** → YN and YY int., Hadrons in nuclear matter
 - **Heavy-ion machines** → EOS at high density, hh int.
- **More NS observations are necessary !**
 - **Precise (and assumption free) measurement of R_{NS}**
→ many satellites will be launched soon: **ASTRO-H, NICER, LOFT**
 - **Larger NS mass will further constrain (or kill ?) nuclear physics.**
- **There are more subjects in neutron star physics.**
 - **Cooling, Magnetic field, Crust, Pasta, finite T, ... were not discussed.**
(Ask Maruyama-san on Pasta and Crust.)

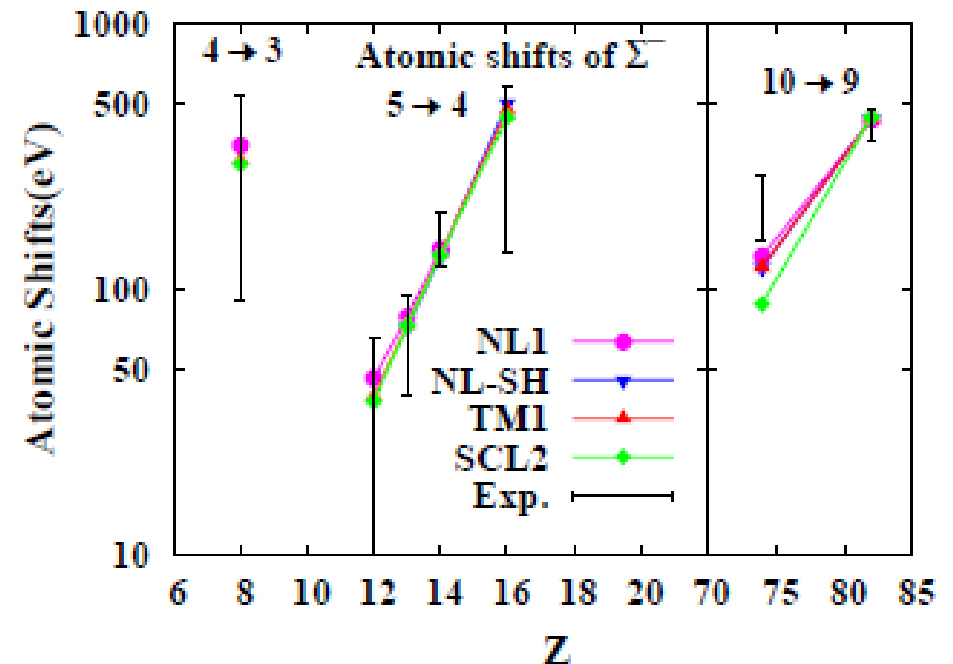
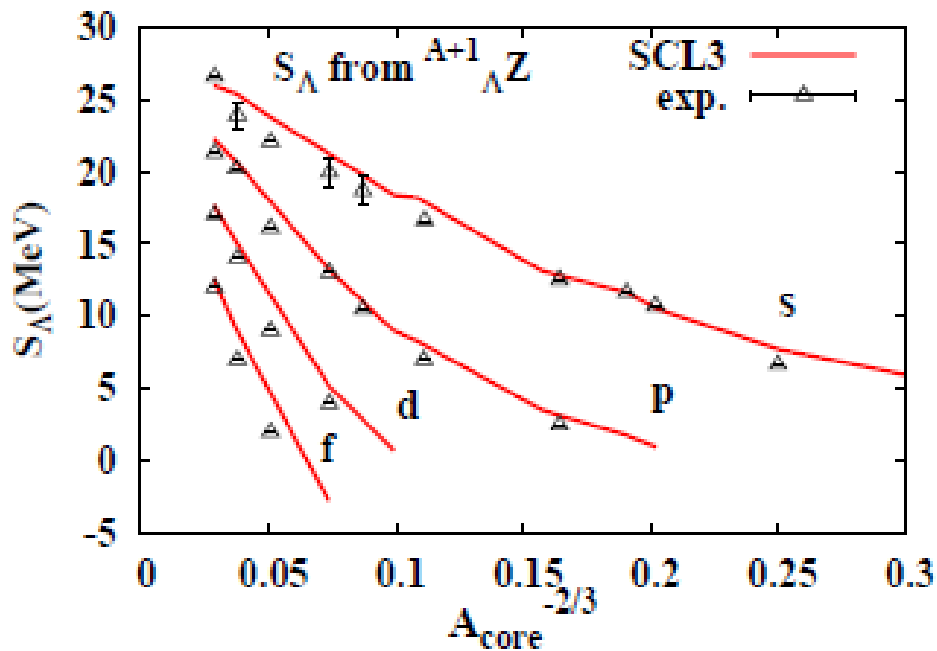
Do I have time ?

Alternative approach

Alternative method

~ “Ab initio” Nucl. Matter EOS + Y phen.

- Fit “Ab initio” EOSs in a phen. model,
- Include hyperons, and explain hypernuclear data.



Tsubakihara et al., PRC81('10)065206

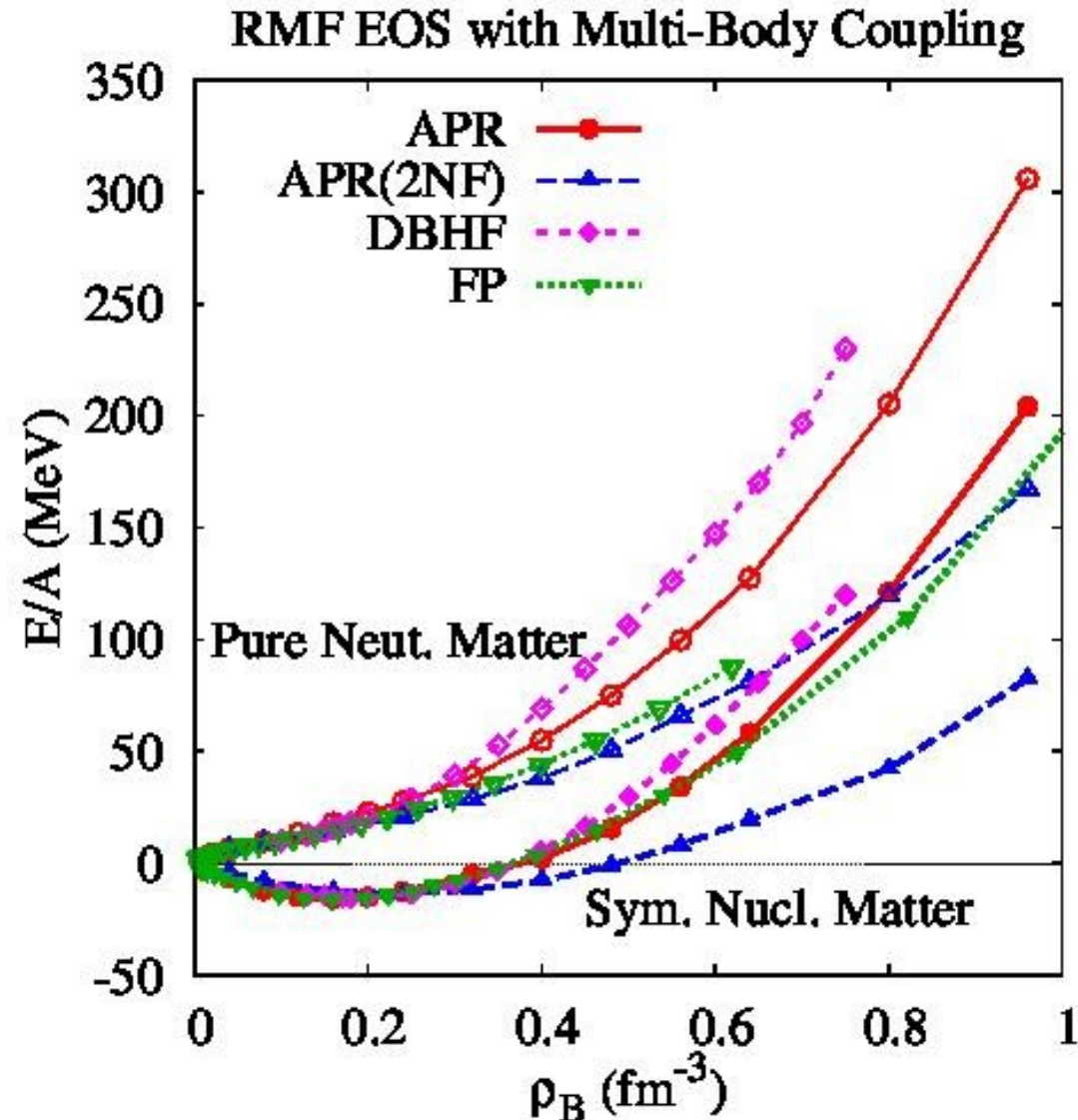
Tsubakihara, Harada, AO, arXiv:1402.0979

We fit ab initio EOS in RMF with multi-body couplings, and introduce hyperons.

“Ab initio” EOS

■ “Ab initio” EOS under consideration

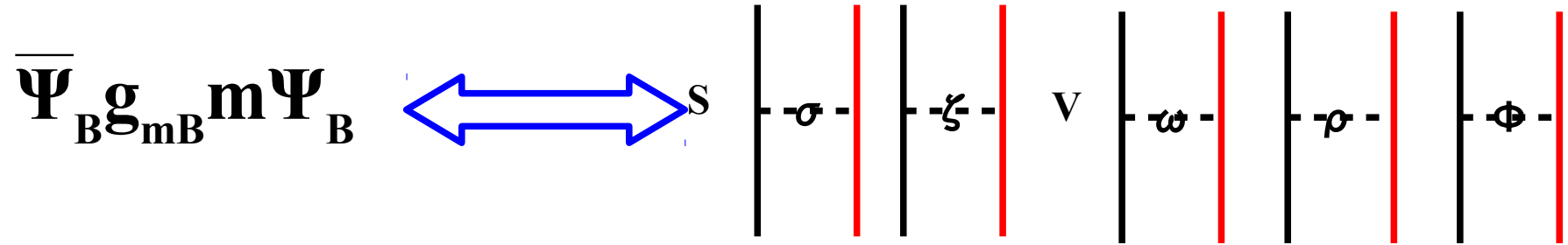
- **FP: Variational calc.**
(Av14+3NF(att.+repl.))
B. Friedman, V.R. Pandharipande, NPA361('81)502.
- **APR: Variational chain summation**
(Av18+rel. corr. ; Av18+ rel. corr.+3NF)
A. Akmal, V.R.Pandharipande, D.G. Ravenhall, PRC58('98)1804.
- **DBHF: Dirac Bruckner approach (Bonn A)**
G. Q. Li, R. Machleidt, R. Brockmann, PRC45('92)2782



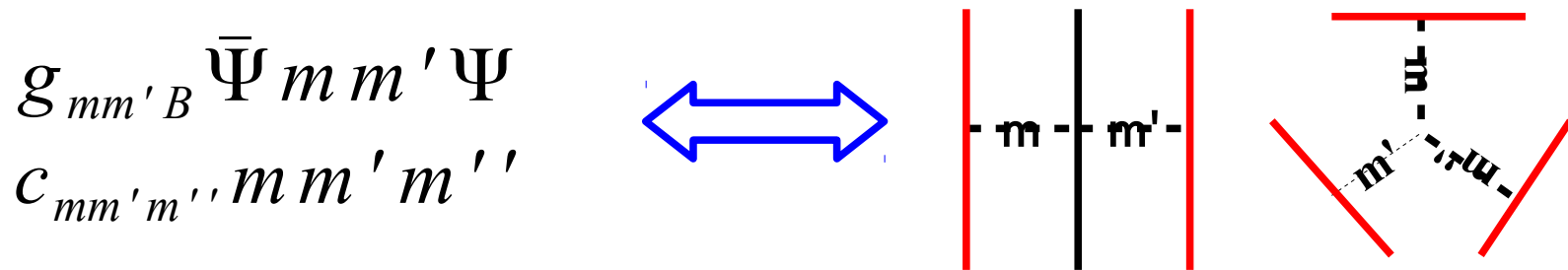
$n=2$ and $n=3$ terms in RMF

Tsubakihara

- $n=B/2+M+D=2$ RMF model (+ effective pot.)
 \rightarrow 2-body interaction (and rel. 3-body corr.)



- $n=3$ model \rightarrow 3-body coupling



Bmm terms are ignored in FST paper (field redefinitions).

Fitting “Ab initio” EOS via RMF

■ RMF with multi-body couplings: 15 parameters

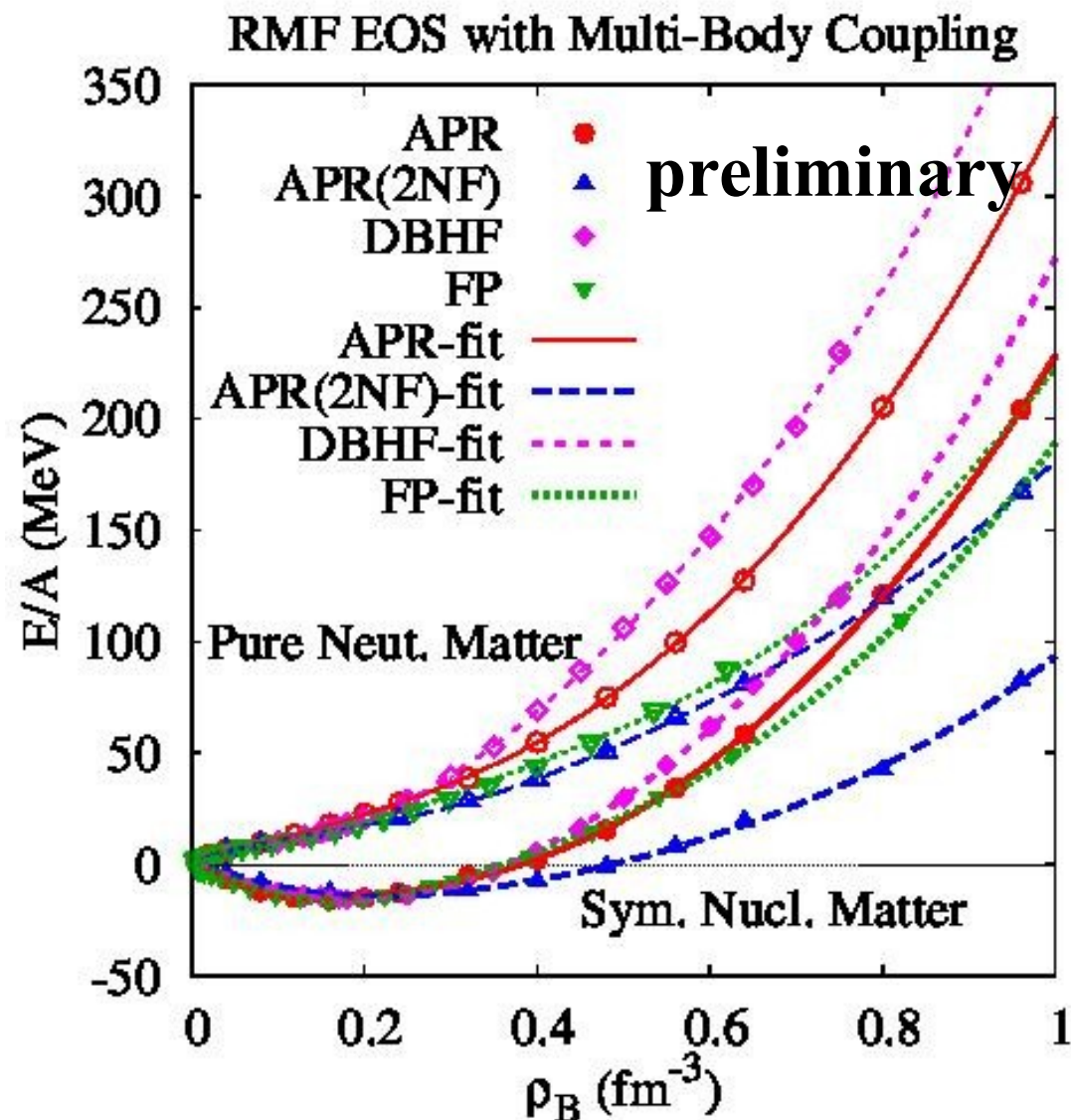
- Working hypothesis
 σ self-energy: SCL2 model

Tsubakihara, AO ('07)

$$M_N \rightarrow 0 @ \sigma \rightarrow f_\pi$$

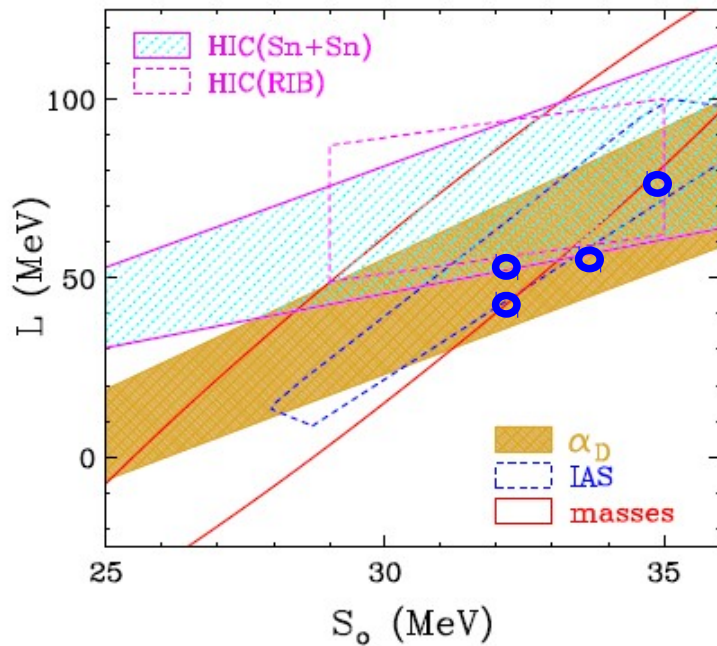
■ Markov Chain Monte-Carlo (MCMC)-like parameter search

- Langevin type shift +Metropolis judge
- Simultaneous fit of SNM and PNM is essential.
- std. dev=0.5-0.7 MeV

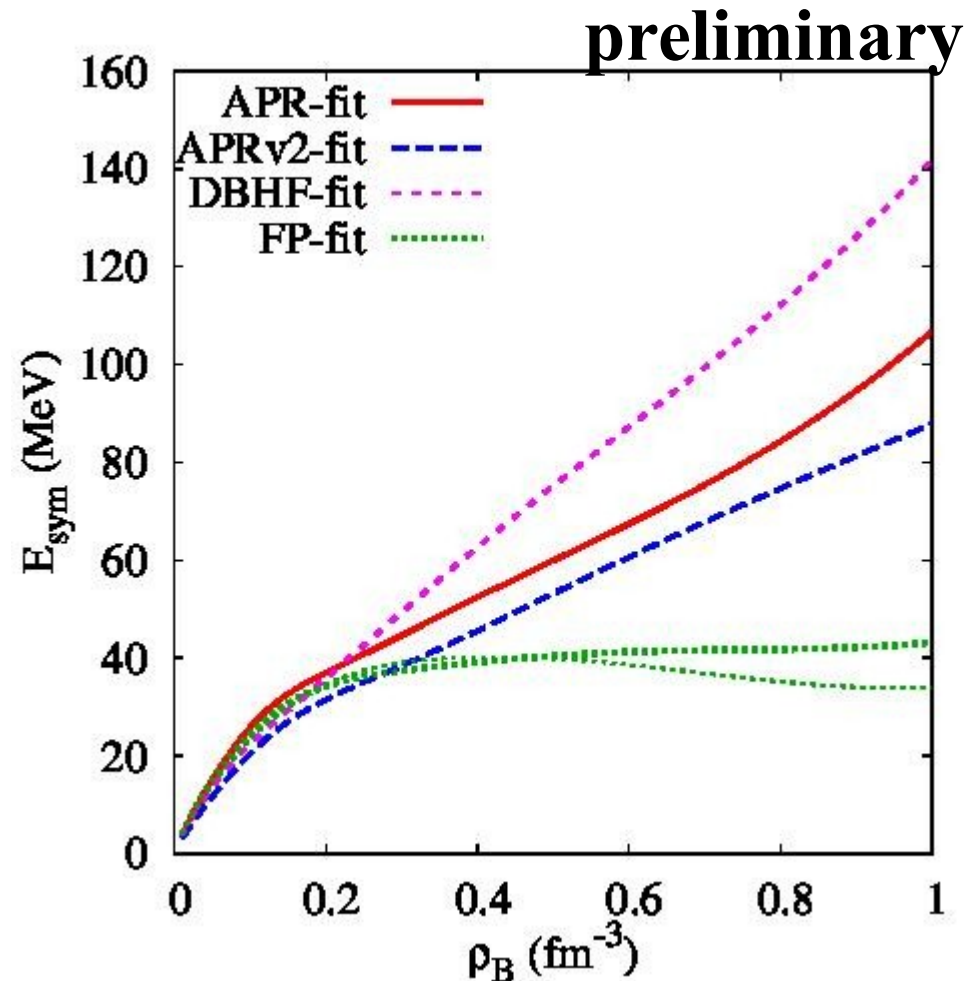


Symmetry Energy

- Symmetry $E. = E(\text{PNM}) - E(\text{SNM})$
 - APR-fit: $(S_0, L) = (32, 47)$ MeV
 - APRv2-fit: $(S_0, L) = (33, 47)$ MeV
 - DBHF-fit: $(S_0, L) = (35, 75)$ MeV
 - FP-fit: $(S_0, L) = (32, 40)$ MeV



Horowitz et al. ('14)



Neutron Star Matter EOS

- Asymmetric Nuclear Matter EOS

$$E_{\text{ANM}}(\rho) = E_{\text{SNM}}(\rho) + \delta^2 S(\rho)$$

β -equilibrium condition \rightarrow NS matter EOS

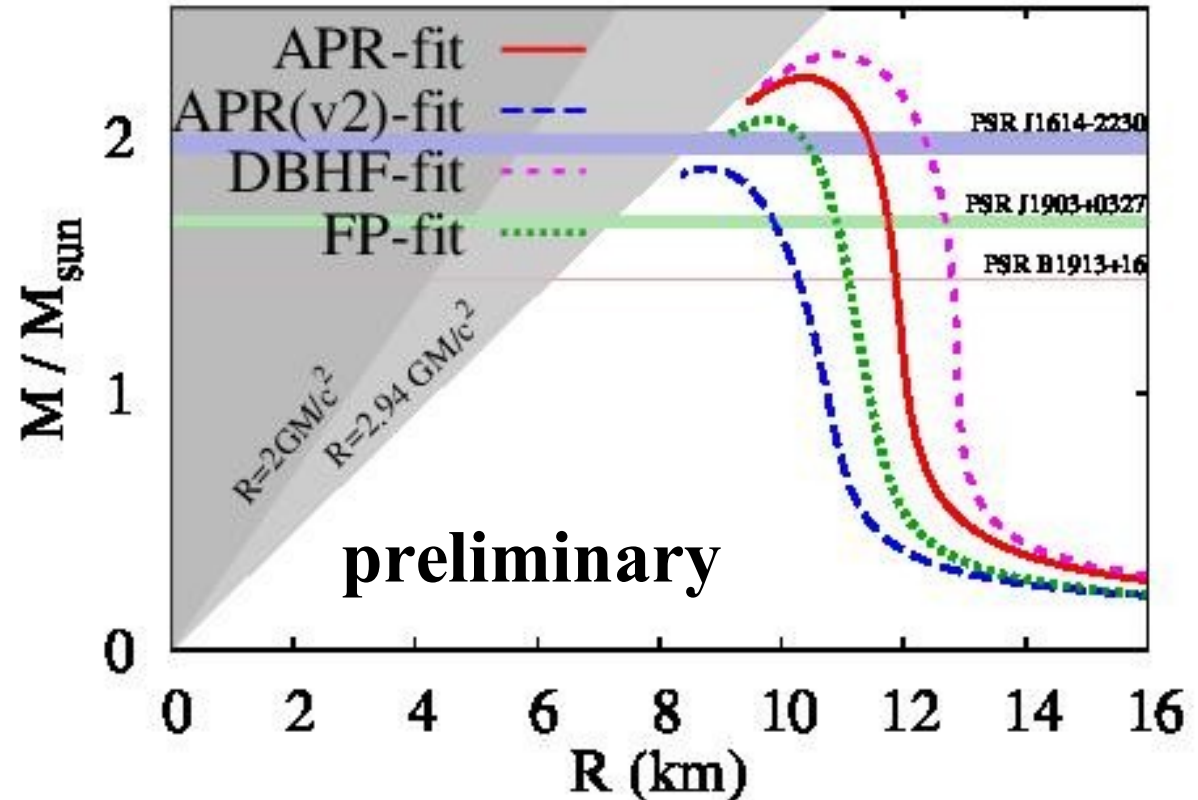
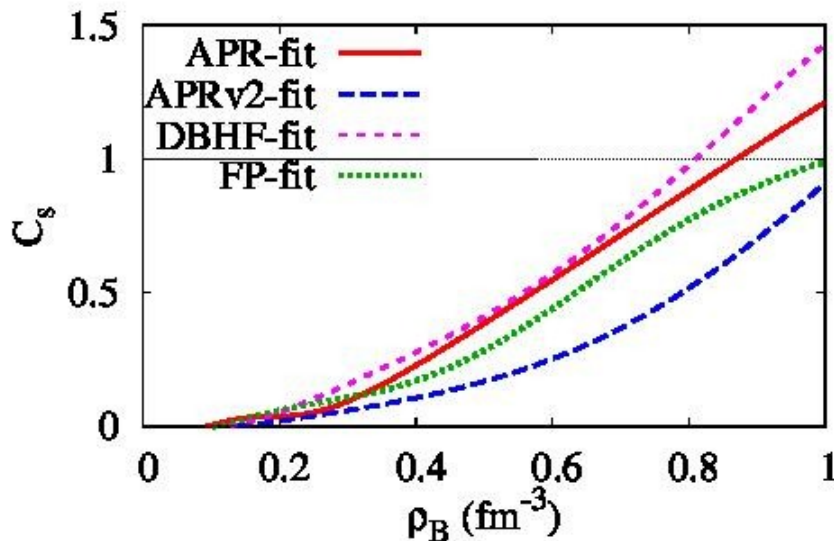
- Max. mass in the fit EOS deviates from the original one by $\sim 0.1 M_{\odot}$.

$$\eta = (KL^2)^{1/3} ?$$

Sotani et al.(2014)

- Caveat:

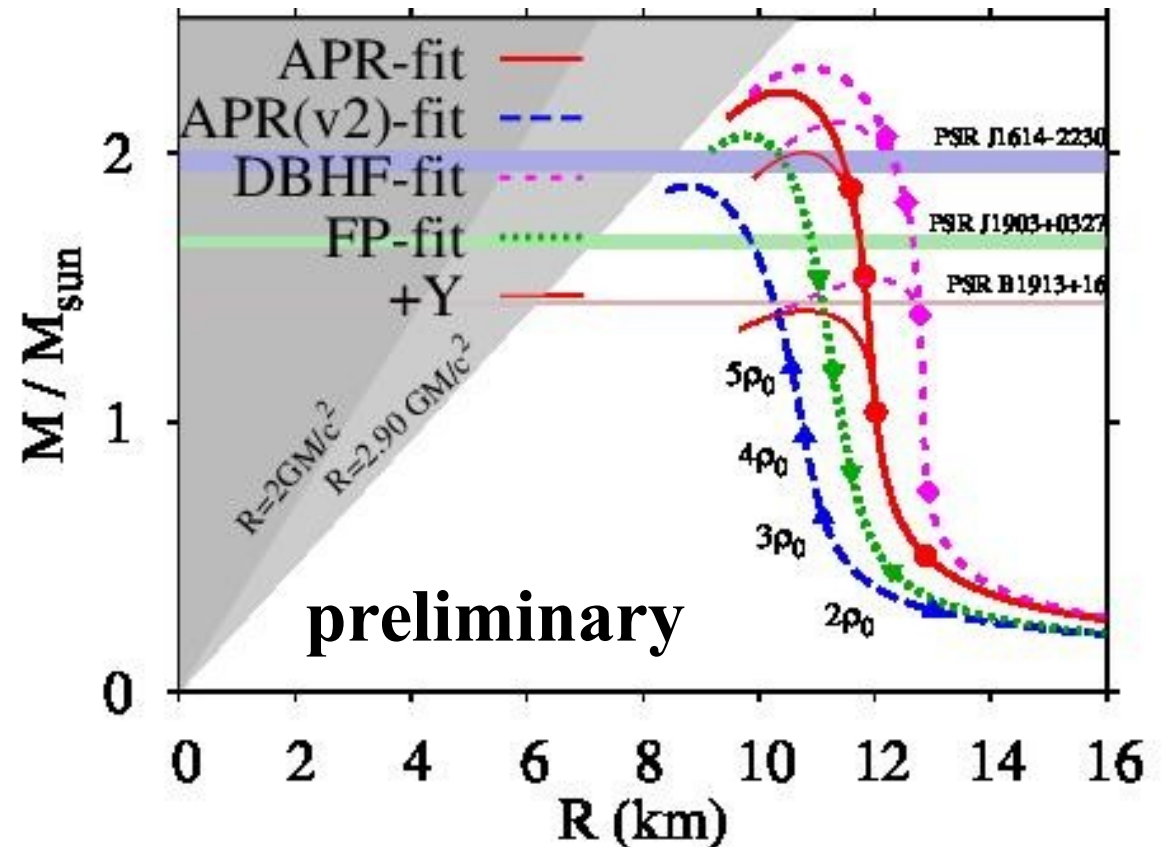
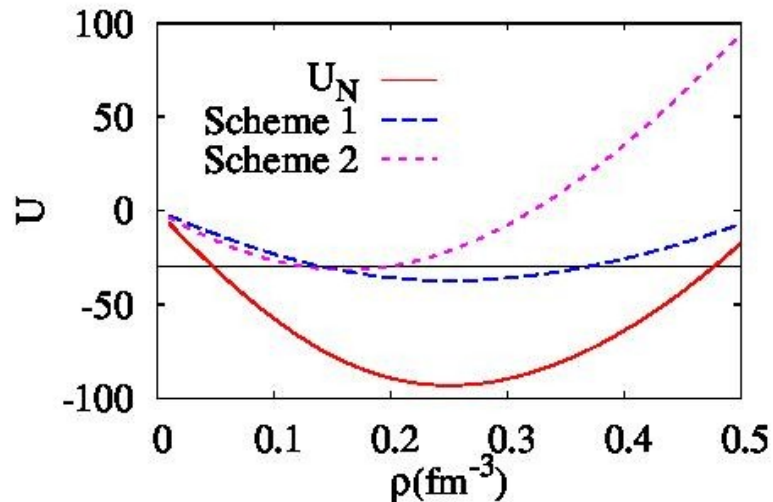
$c_s > c$ at high density



NS matter in “ab initio”-fit + Λ

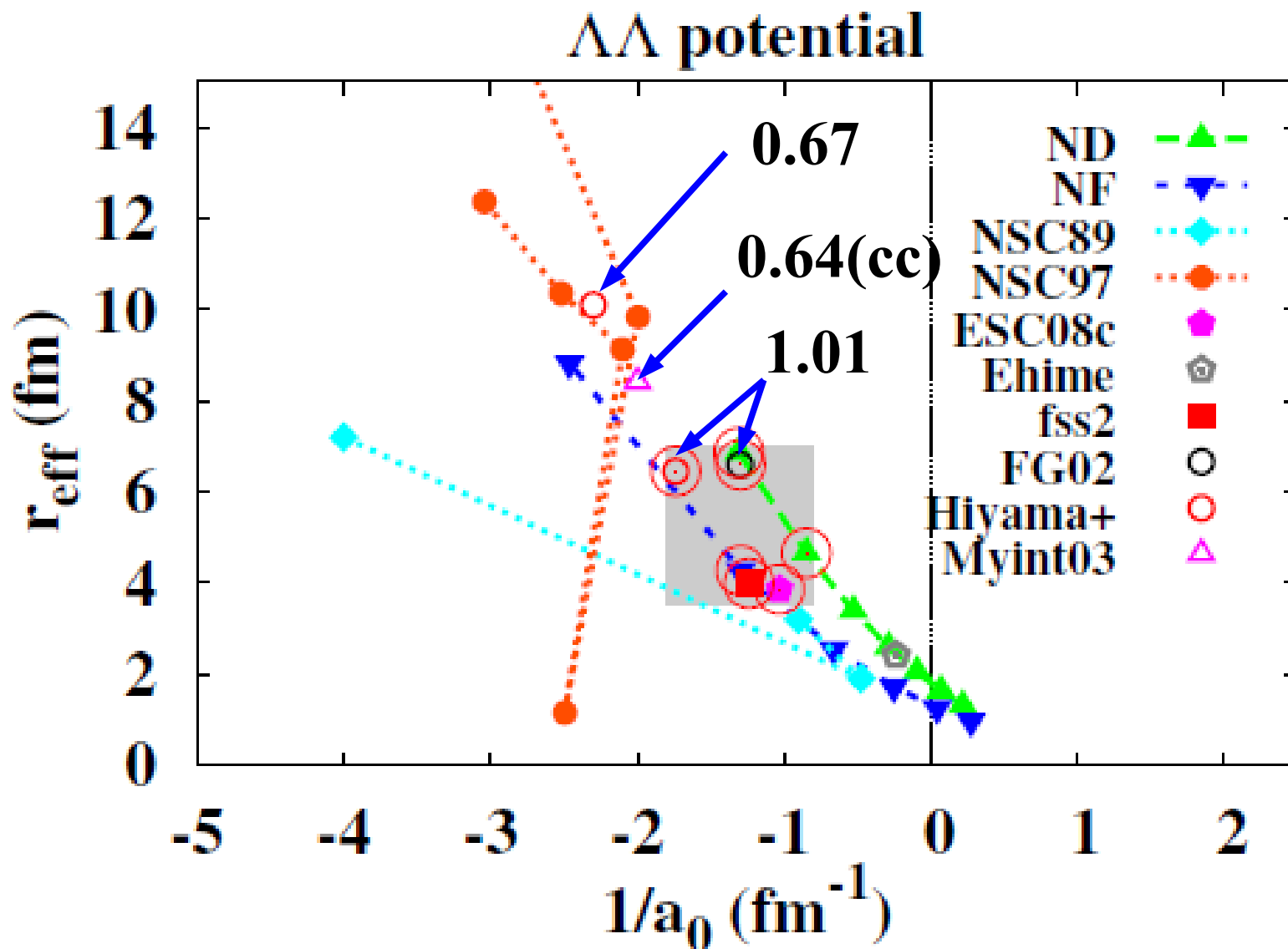
- Λ potential in nuclear matter at $\rho_0 \sim -30$ MeV

- Scheme 1: $U_\Lambda(\rho) = \alpha U_N(\rho)$
- Scheme 2: $U_\Lambda(\rho) = 2/3 U_N^{n=2}(\rho) + \beta U_N^{n>2}(\rho)$



Thank you !

Results with smaller $\Lambda\Lambda$ bond energy



Chiral EFT NN & NNN force

	2N force	3N force	4N force
LO		—	—
NLO		—	—
N ² LO			—
N ³ LO			

E. Epelbaum ('09)

中性子物質と冷却原子

■ BEC-BCS crossover and unitary gas

- 散乱長 \gg 粒子間距離 \rightarrow EOS は普遍的 (unitary gas)

$$E^{\text{Unitary}} = \xi E^{\text{Free}} \quad \xi \simeq 0.4 \text{ (Bertsch parameter)}$$

- nn 間の 1S_0 散乱長は長い! ($a_0 = -18.5 \text{ fm}$)

\rightarrow Dripした中性子ガスは、ほぼ unitary gas ($-1/k_F a_0 \sim 0.1$)

■ My question

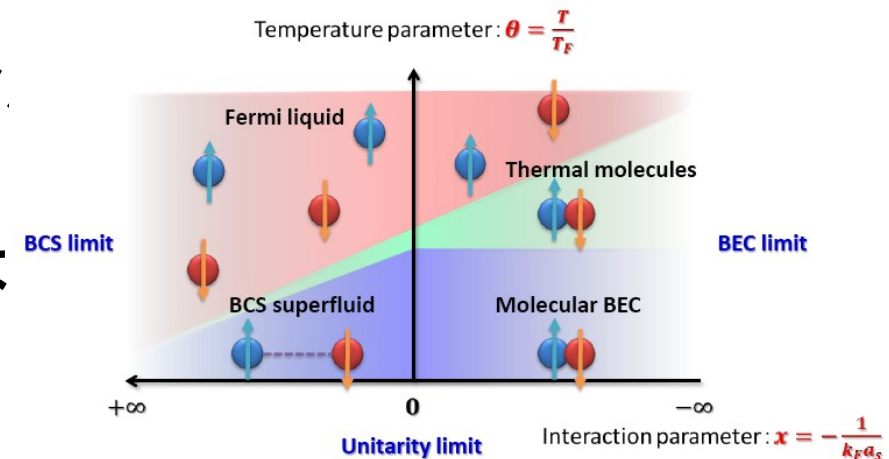
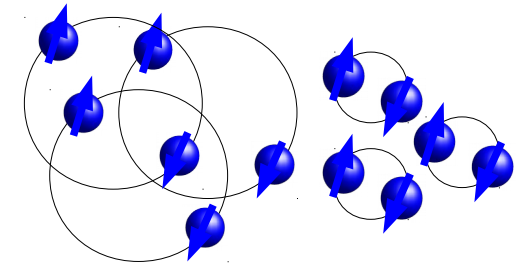
- 核子あたりの相互作用エネルギー

$$\propto k_F^2 \propto \rho^{2/3}$$

$$\frac{V^{\text{Unitary}}}{N} = (\xi - 1) \frac{3}{5} \frac{\hbar^2 k_F^2}{2m} \propto \rho^{2/3}$$

- どのようにして EOS(密度汎関数) | 取り込むか? (Hartree なら $\propto \rho$)

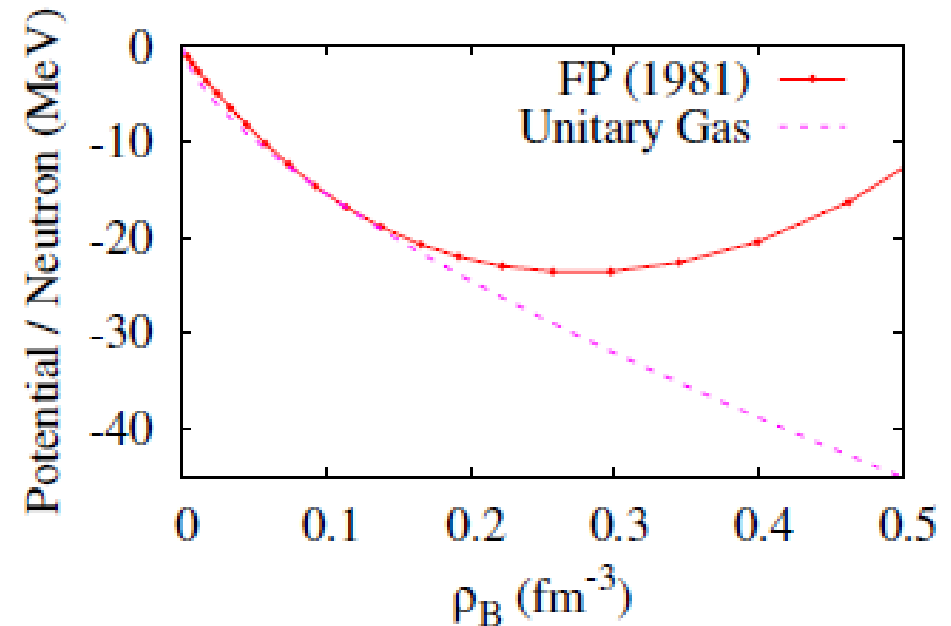
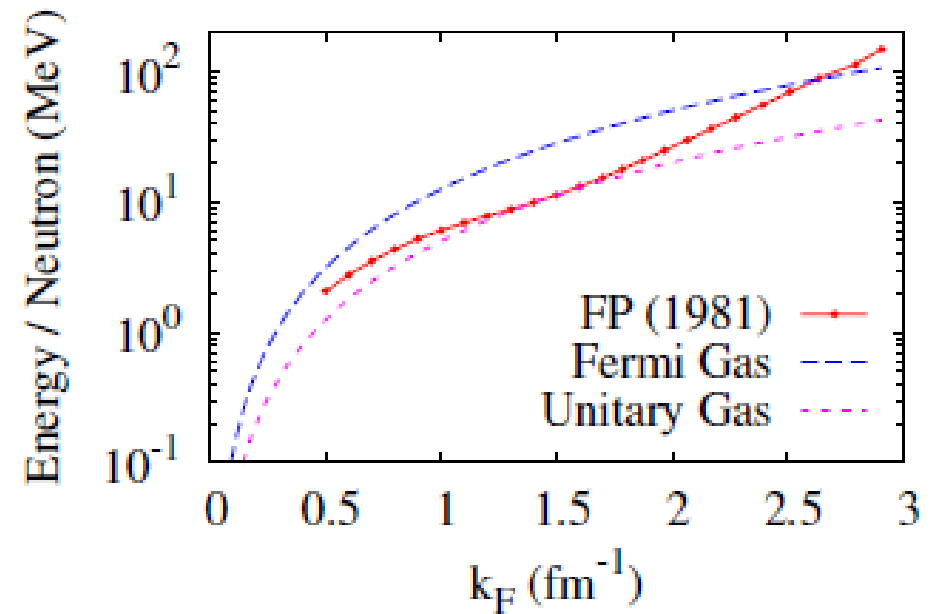
- unitary gas / BEC-BCS crossover は クラスト・原子核の性質に どのような影響を及ぼすか?



中性子星物質の状態方程式

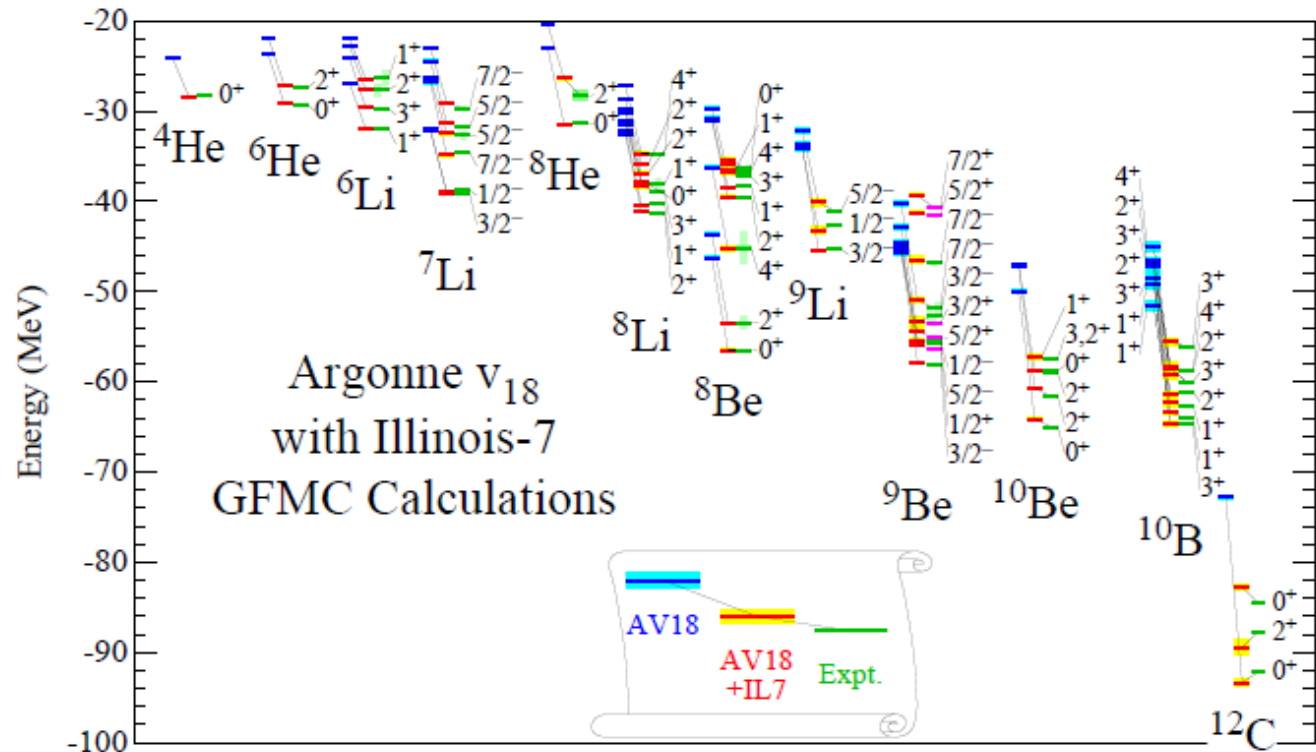
■ 変分法による計算結果 Friedman-Pandharipande (1981)

- 広い密度領域において
 $E_{\text{unit}} < E_{\text{FP}} < E_{\text{Fermi}}$
- 低密度領域でポテンシャルエネルギーは $\rho^{2/3}$ と振る舞っているか？



What is necessary to solve the massive NS puzzle ?

- There are many “model” solutions.
- Ab initio calculation including three-baryon force (3BF)
 - Bare 2NF+Phen. 3NF(UIX, IL2-7) + many-body theory (verified in light nuclei).
 - Chiral EFT (2NF+3NF) + many-body theory
 - Dirac-Bruckner-HF (no 3NF)



J. Carlson et al. ('14)

Relativistic Mean Field with Multi-body couplings

$\sigma\omega\rho$ model +std. non-linear terms + multi-body couplings

$$\mathcal{L}_N = \bar{\psi} (i\gamma^\mu \partial_\mu - M_N - U_s - \gamma^\mu U_\mu) \psi + \mathcal{L}_{\sigma\omega\rho}$$

$$\mathcal{L}_{\sigma\omega\rho} = \frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma - \frac{1}{4} \omega_{\mu\nu} \omega^{\mu\nu} - \frac{1}{4} R_{\mu\nu} \cdot R^{\mu\nu} - \mathcal{V}_{\sigma\omega\rho}$$

$$U_s = -g_\sigma \sigma [1 + r_{\sigma\sigma}(1 - \sigma/f_\pi)] + g_\sigma \omega^\mu \omega_\mu / f_\pi [r_{\omega\omega} + r_{\sigma\omega\omega}(1 - \sigma/f_\pi)]$$

$$U_\mu = g_\omega \omega_\mu [1 - r_{\sigma\omega}\sigma/f_\pi + r_{\omega 3}\omega^\nu \omega_\nu / f_\pi^2]$$

$$+ g_\rho \tau \cdot R_\mu [1 - r_{\sigma\rho}\sigma/f_\pi + r_{\omega\rho}\omega^\nu \omega_\nu / f_\pi^2]$$

$$\mathcal{V}_{\sigma\omega\rho} = \frac{1}{2} m_\sigma^2 \sigma^2 - a_\sigma f \log(\sigma/f_\pi) + \frac{1}{4} c_{\sigma 4} (\sigma^4 - 4f_\pi \sigma^3)$$

$$- \frac{1}{2} m_\omega^2 \omega^\mu \omega_\mu [1 - c_{\sigma\omega}\sigma/f_\pi] - \frac{1}{4} c_{\omega 4} (\omega^\mu \omega_\mu)^2$$

$$- \frac{1}{2} m_\rho^2 R^\mu \cdot R_\mu [1 - c_{\sigma\rho}\sigma/f_\pi + c_{\omega\rho}\omega^\mu \omega_\mu / f_\pi^2] - \frac{1}{4} c_{\rho 4} (R^\mu \cdot R_\mu)^2$$

$$f \log(x) = \log(1-x) + x + \frac{1}{2} x^2 \quad a_\sigma = f_\pi^2 (m_\sigma^2 - m_\pi^2) / 2 - f_\pi^4 c_{\sigma 4}$$

RMF with many-body coupling

■ Naive dimensional analysis (NDA) and naturalness

Manohar, Georgi ('84)

The vertex is called “natural” if $C \sim 1$.

$$L_{\text{int}} \sim (f_\pi \Lambda)^2 \sum_{l,m,n,p} \frac{C_{lmnp}}{m!n!p!} \left(\frac{\bar{\psi} \Gamma \psi}{f_\pi^2 \Lambda} \right)^l \left(\frac{\sigma}{f_\pi} \right)^m \left(\frac{\omega}{f_\pi} \right)^n \left(\frac{R}{f_\pi} \right)^p$$

→ Consistent with the idea that the vertex is generated by loop diagrams under the assumption that the QCD coupling is small.

■ FST truncation

R. J. Furnstahl, B. D. Serot, H. B. Tang, NPA615 ('97)441.

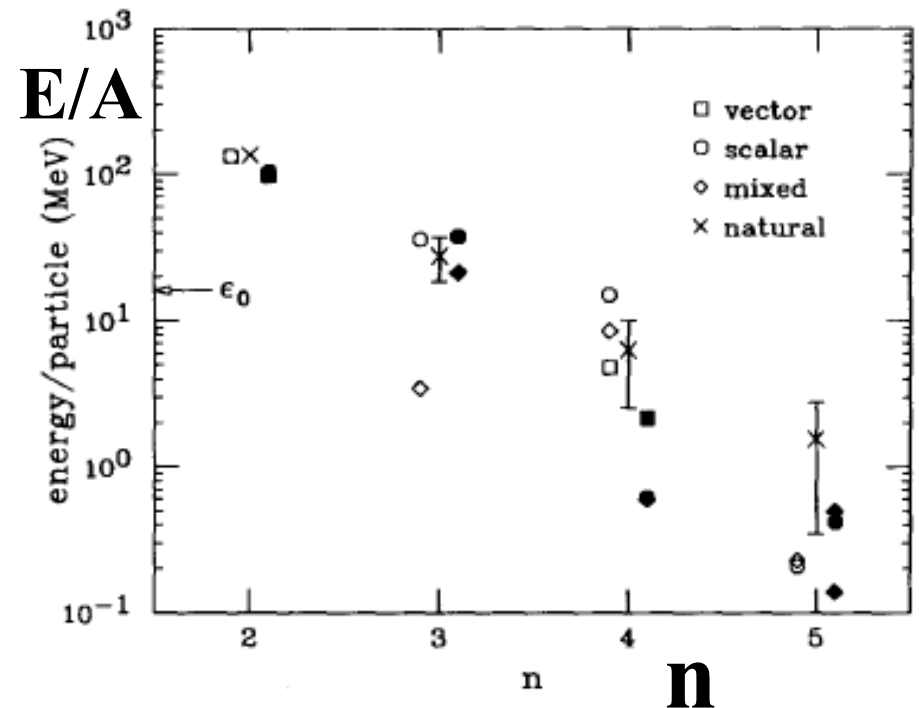
At a given density, we can truncate the Lagrangian by the index

$$n = B/2 + M + D$$

(B: baryon field, M: Non NG boson, D: derivatives)

Naturalness → $V \sim \rho^n/n!$

→ small for large n



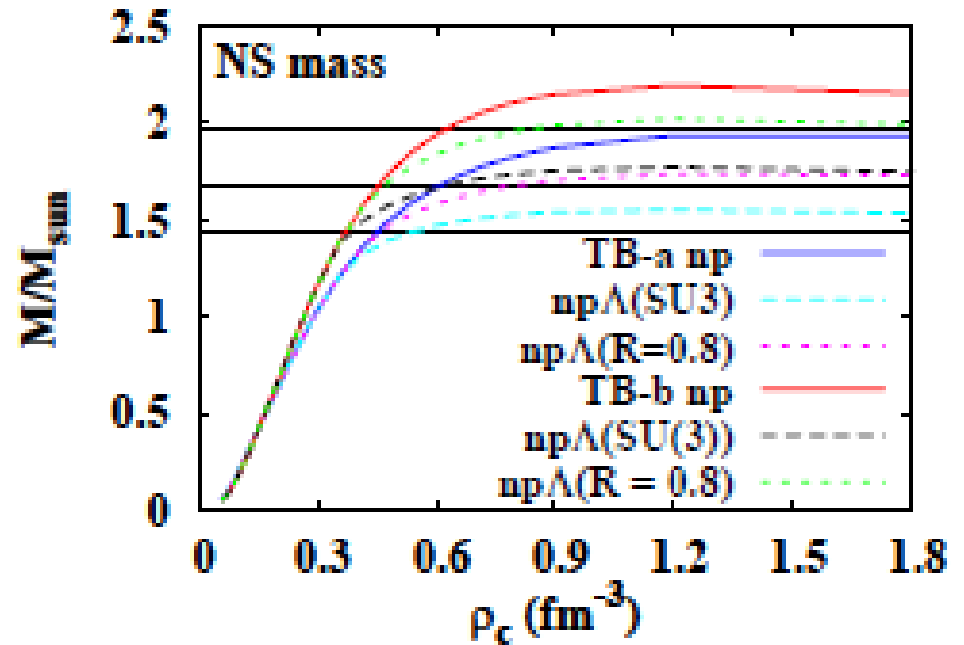
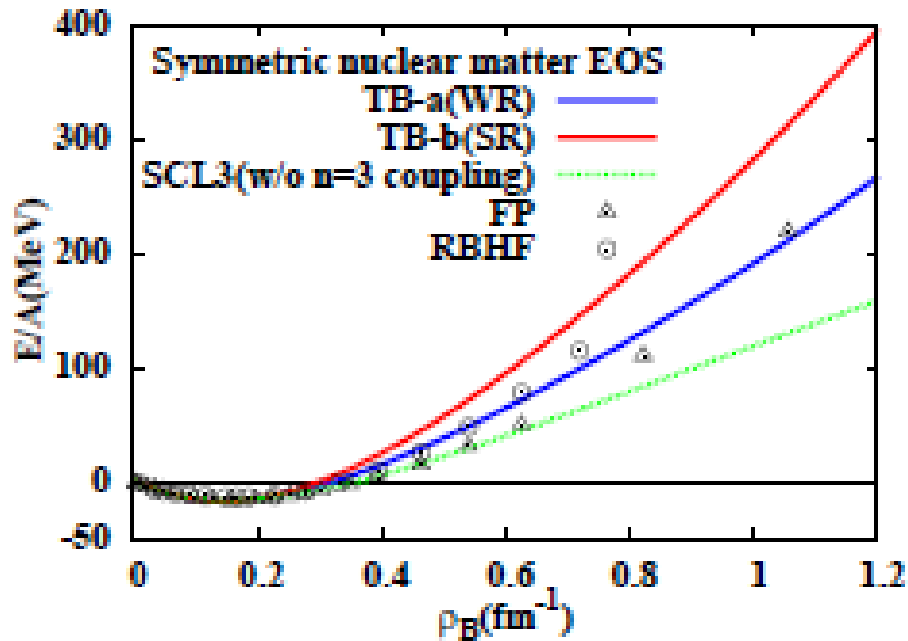
$n=3$ coupling terms

■ RMF with $n=3$ terms

- $n=B/2+M+D$; baryon, meson, derivative

$$\mathcal{L}_{n=3}^{\sigma\omega} = -\frac{1}{f_\pi} \sum_B \bar{\psi}_B \left[g_{\sigma\sigma B} \sigma^2 + g_{\omega\omega B} \omega_\mu \omega^\mu - g_{\sigma\omega B} \sigma \omega_\mu \gamma^\mu \right] \psi_B - c_{\sigma\omega\omega} f_\pi \sigma \omega_\mu \omega^\mu$$

- $g_{\sigma\Lambda} / g_{\sigma N} \sim 0.8 > 2/3 \rightarrow 2 M_\odot$ NS
- Parameter fitting: $(\rho_0, E/A)$, Vector pot. in DBHF, S_0, L, \dots



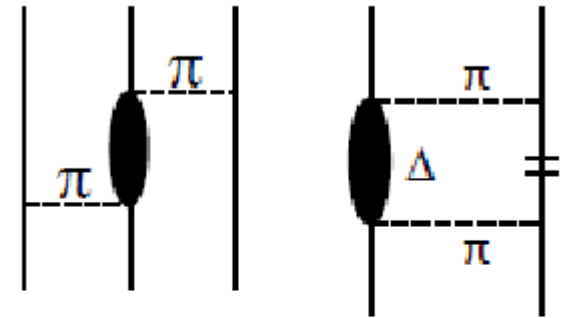
Tsubakihara, AO, NPA914 ('13), 438.

“Universal” mechanism of “Three-body” repulsion

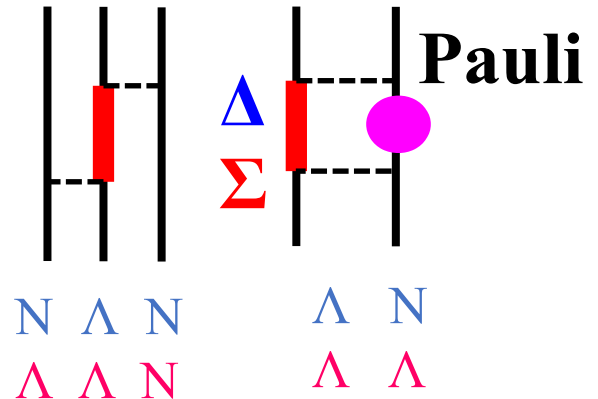
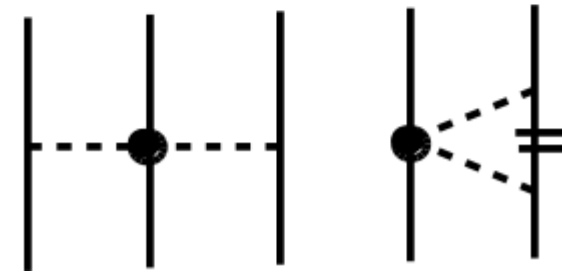
■ Mechanism of “Universal” Three-Baryon Repulsion.

- “ σ ”-exchange \sim two pion exch. w/ res.
- Large attraction from two pion exchange is suppressed by the Pauli blocking in the intermediate stage.

Physical Picture



χ EFT



“Universal” TBR

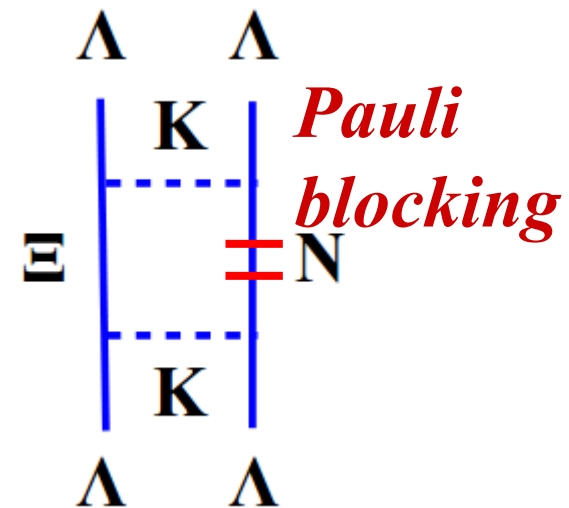
- Coupling to Res. (hidden DOF)
- Reduced “ σ ” exch. pot. ?

How about YNN or YYN ?

$\Lambda\Lambda$ interaction in vacuum and in nuclear medium

- Vacuum $\Lambda\Lambda$ interaction may be theoretically accessible
Lattice QCD calc. HAL QCD ('11) & NPLQCD ('11)
- In-medium $\Lambda\Lambda$ interaction may be experimentally accessible
 - $a_0(\text{Nagara fit}) = -0.575 \text{ fm}, -0.77 \text{ fm}$ ($\Delta B_{\Lambda\Lambda} = 1.0 \text{ MeV}$)
Hiyama et al. ('02), Filikhin, Gal ('02)
 - Bond energy of ${}^6_{\Lambda\Lambda}\text{He}$: $\Delta B_{\Lambda\Lambda} = 1.0 \text{ MeV} \rightarrow 0.6 \text{ MeV}$
Nakazawa, Takahashi ('10)
- Difference of vacuum & in-medium $\Lambda\Lambda$ int. would inform us $\Lambda\Lambda\text{N}$ int. effects.

- $\Lambda\Lambda$ - $\bar{\text{N}}$ couples in vacuum
- Coupling is suppressed in ${}^6_{\Lambda\Lambda}\text{He}$



*Is there Any way to access
“vacuum” $\Lambda\Lambda$ int. experimentally ?*

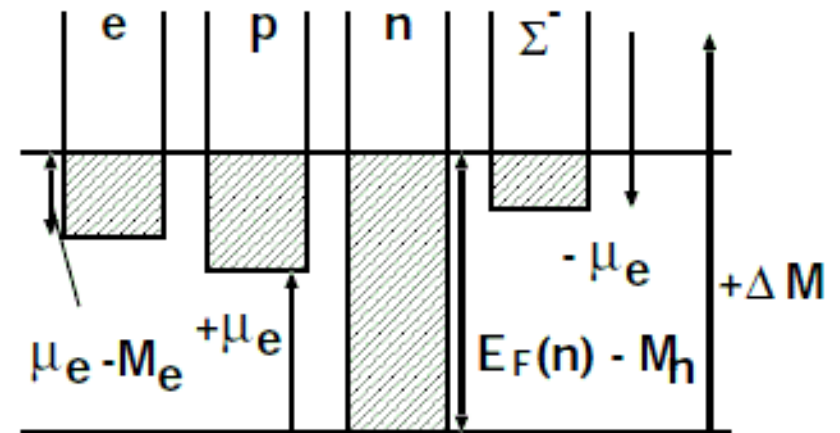
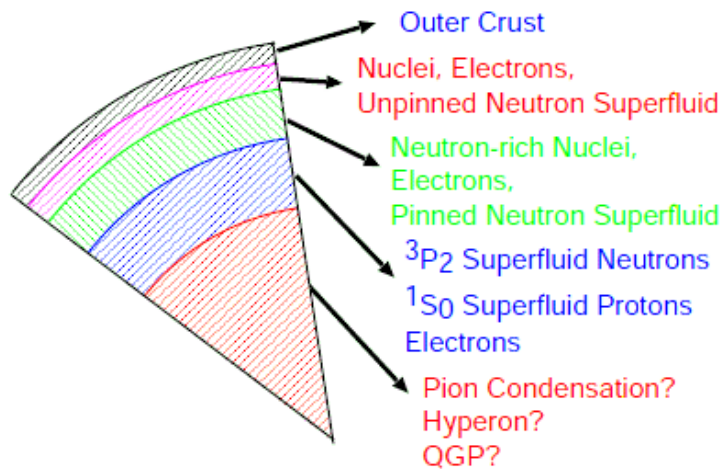
Hyperons in Dense Matter

What appears at high density ?

- Nucleon superfluid (3S_1 , 3P_2), Pion condensation, Kaon condensation, Baryon Rich QGP, Color SuperConductor (CSC), Quarkyonic Matter, ...

Hyperons

Tsuruta, Cameron (66); Langer, Rosen (70); Pandharipande (71); Itoh(75); Glendenning; Weber, Weigel; Sugahara, Toki; Schaffner, Mishustin; Balberg, Gal; Baldo et al.; Vidana et al.; Nishizaki, Yamamoto, Takatsuka; Kohno, Fujiwara et al.; Sahu, Ohnishi; Ishizuka, Ohnishi, Sumiyoshi, Yamada; ...



Nobody says “Hyperons cannot appear in neutron star core” !

Y appears when $\mu_B = E_F(n) + U(n) \geq M(Y) + U(Y) + Q_Y \mu_e$