

高密度物質と中性子星の物理

Physics of Neutron Star Matter

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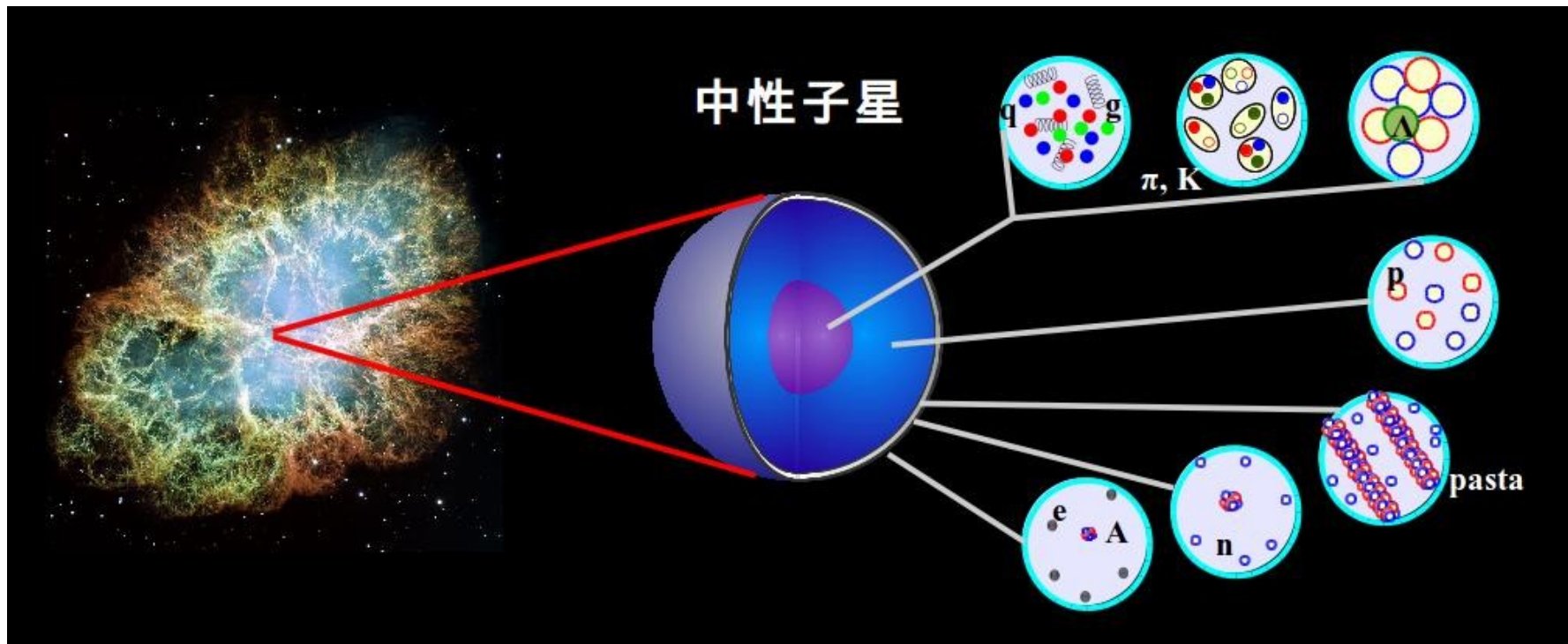
1. 中性子星の基本的性質
 2. 状態方程式を記述する理論模型
 3. 対称エネルギーと非対称核物質の状態方程式
 4. ハイパー核物理と高密度核物質の状態方程式
- 中性子星におけるエキゾチック自由度
→ YYN3 体力：重イオン衝突とハイパー核から中性子星へ

九州大学集中講義 7/8-10



Neutron Star Matter

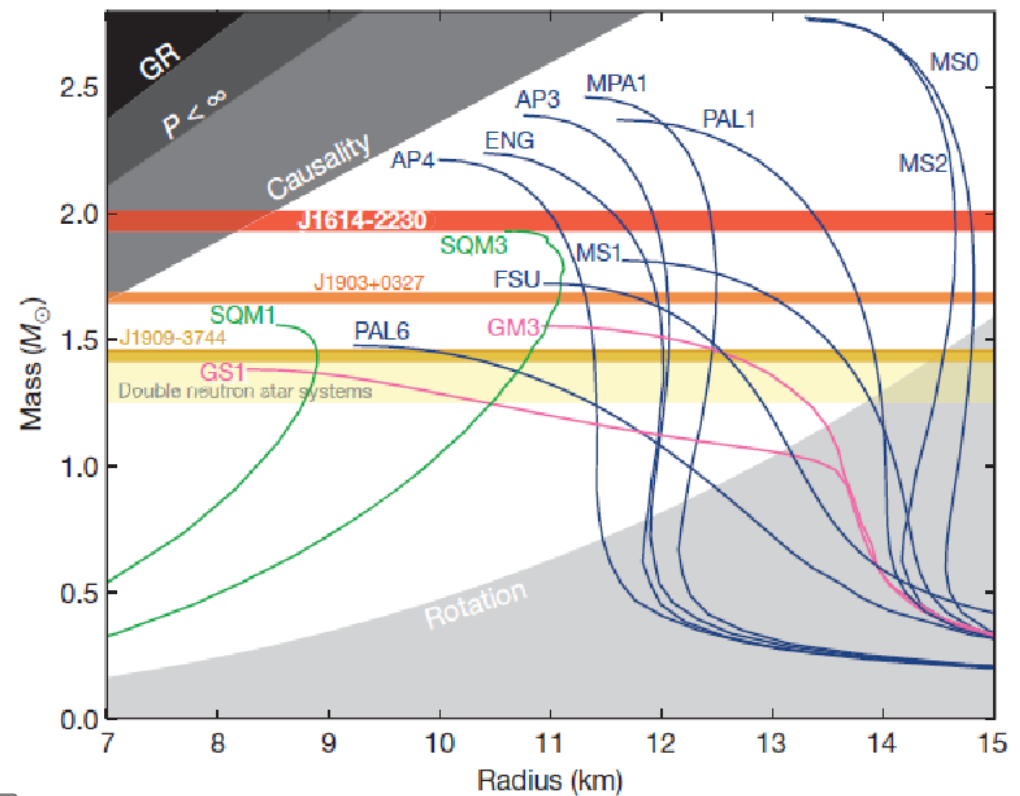
- Cold, dense, charge neutral
- Constituents
 - $n, p, e, \mu, Y, \pi, K, q, \text{di-quark}, \dots$



Can we determine int. btw constituents ?

Massive Neutron Stars

- $M_{\text{NS}} = 1.97 \pm 0.04 \text{ Msun}$ measured using Shapiro delay (GR effect).
- EOSs w/ strange hadrons are “ruled out”, while Lab. exp. suggest their existence in NS.



Massive Neutron Star Puzzle

Demorest et al., Nature 467 (2010) 1081.

NS matter EOS with hyperons

■ “Ruled-out” hyperonic EOS = Naive RMF

Glendenning, Moszkowski ('91)

- SU(6) coupling ($\sim$ quark counting) $g_{\sigma\Lambda} = 2/3 g_{\sigma N}$
- No $\bar{s}s$ mesons

■ Proposed prescription after 2 Msun NS

- Modify coupling constant from SU(6) value

Weisenborn et al., ('11); Tsubakihara, AO, Harada ('13)

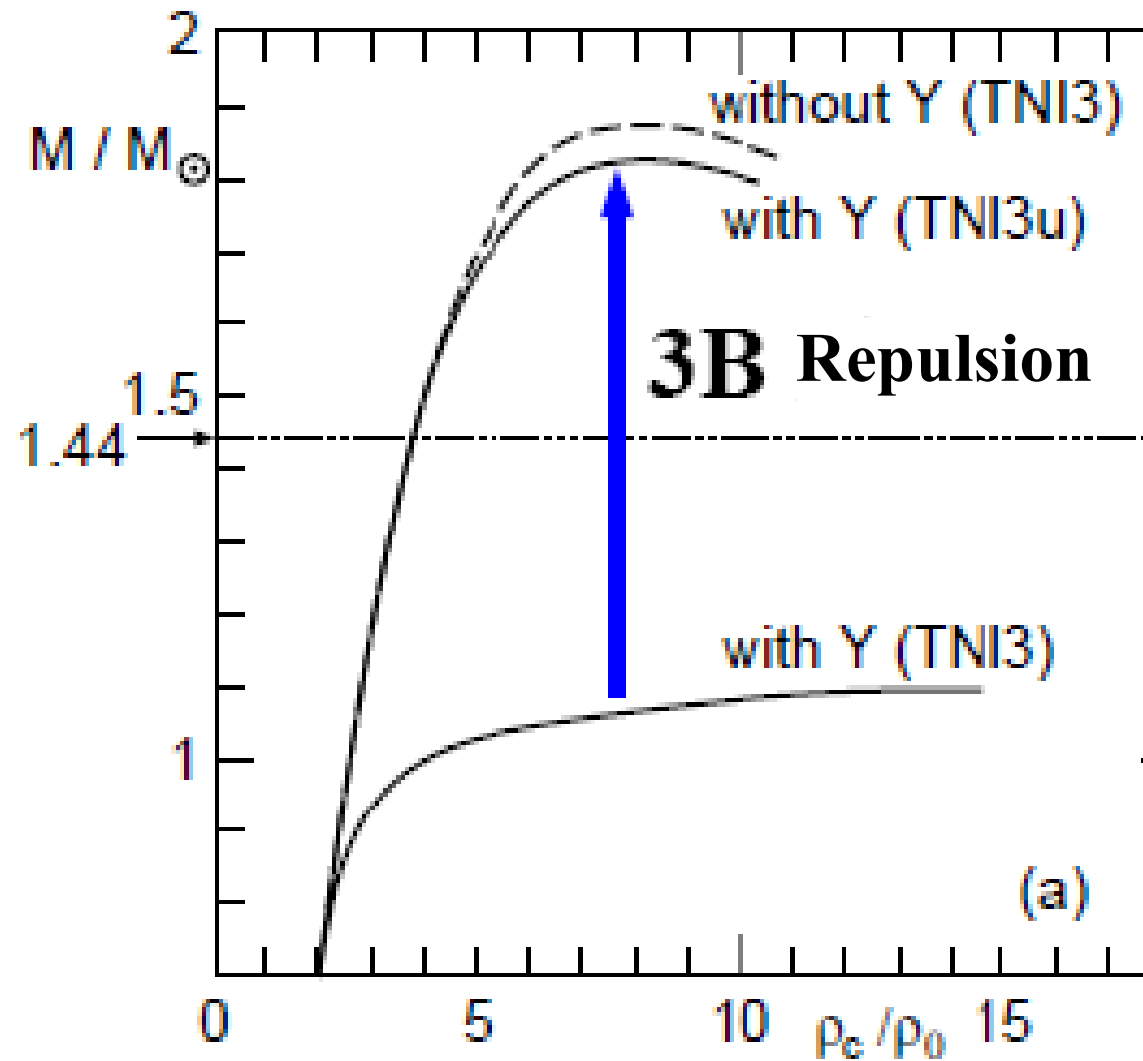
- Introducing three-body repulsion

Bednarek, et al.('11); Miyatsu, Yamamuro, Nakazato ('13)

- Crossover transition to quark matter

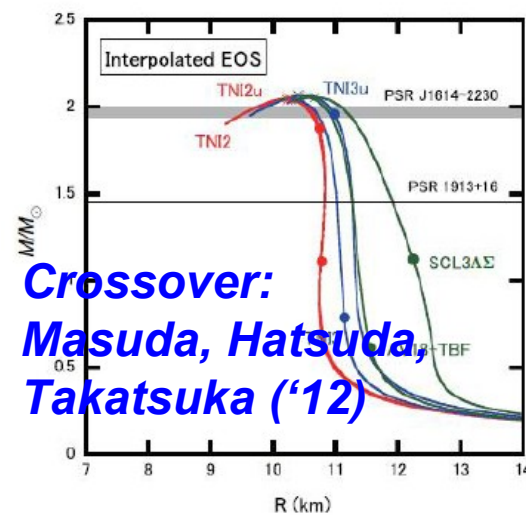
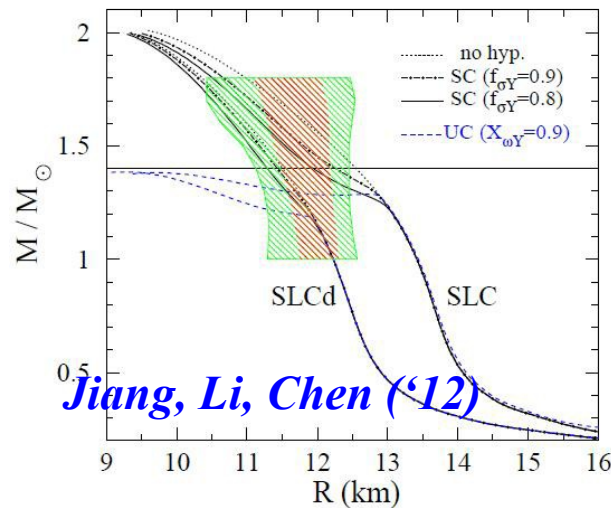
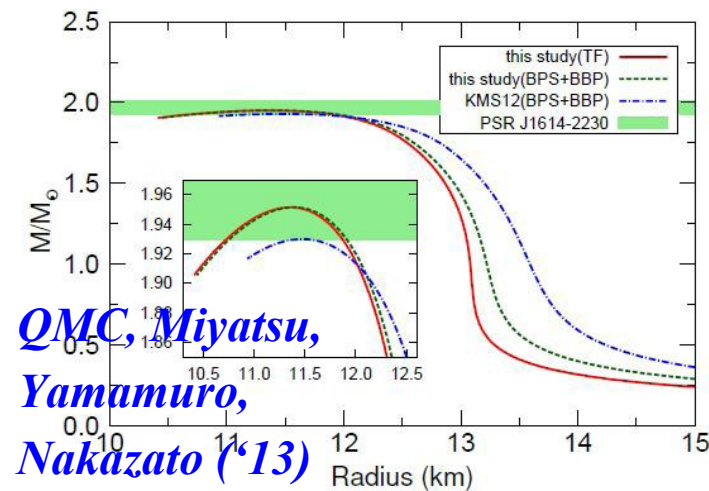
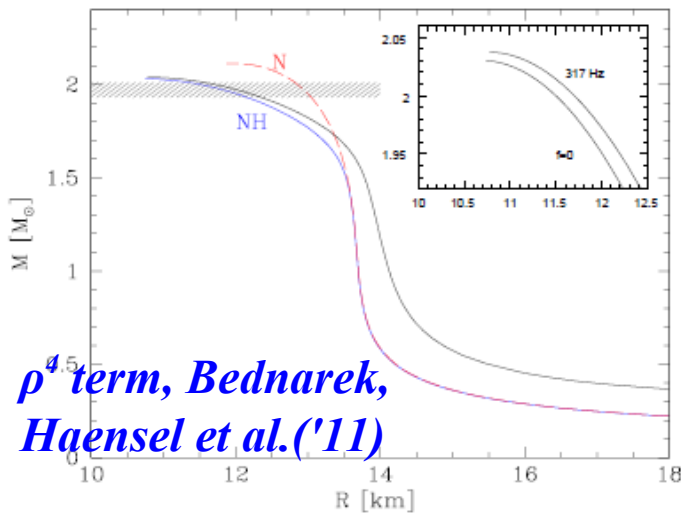
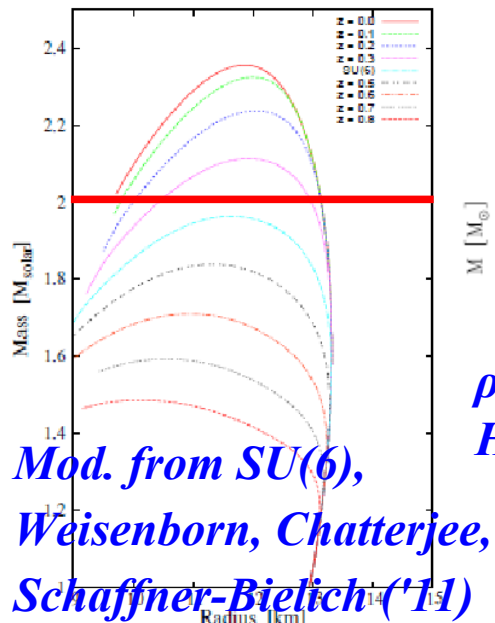
Masuda, Hatsuda, Takatsuka ('12)

Three Baryon Repulsion

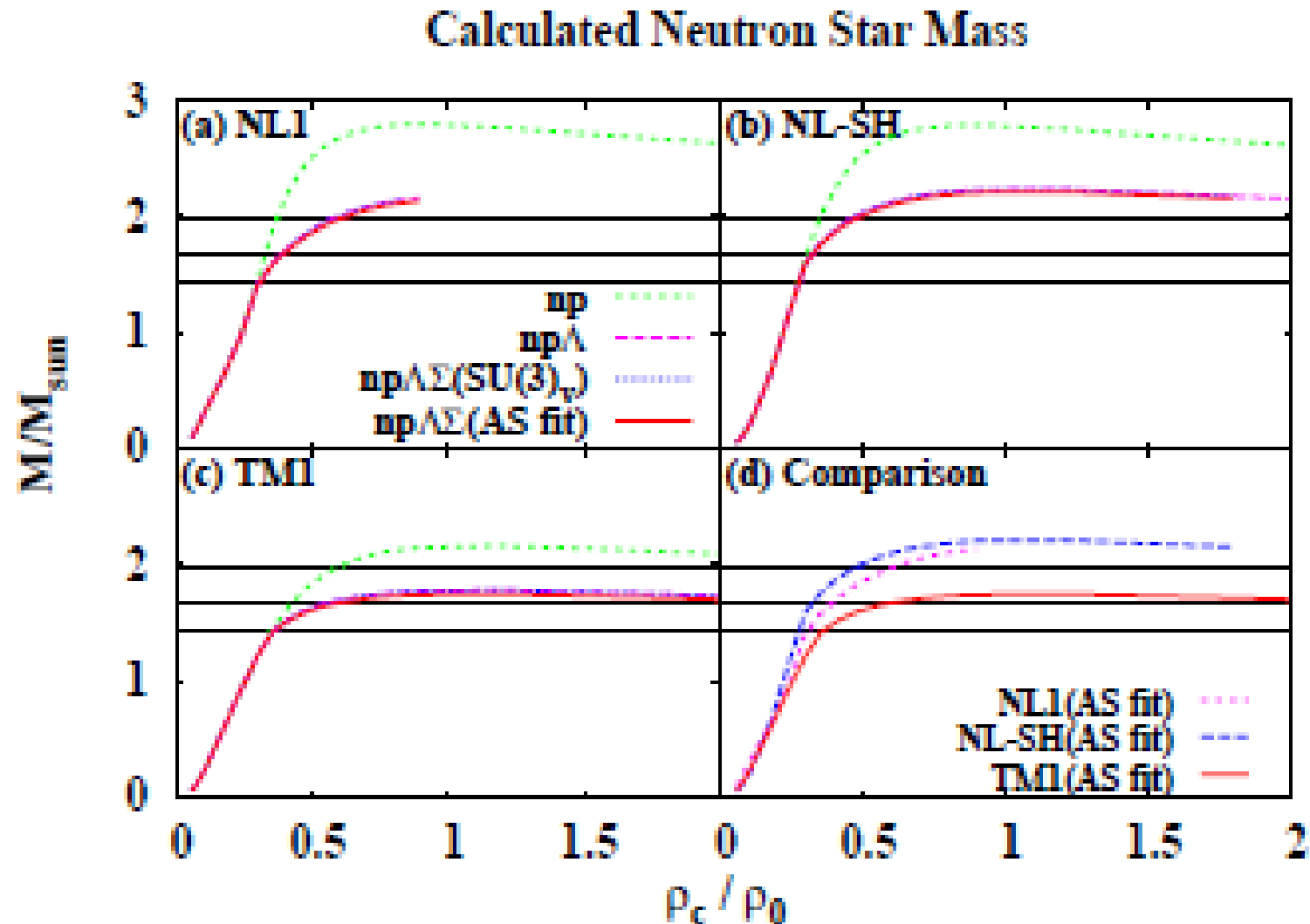


Nishizaki, Takatsuka, Yamamoto ('02)

NS matter EOS with hyperons



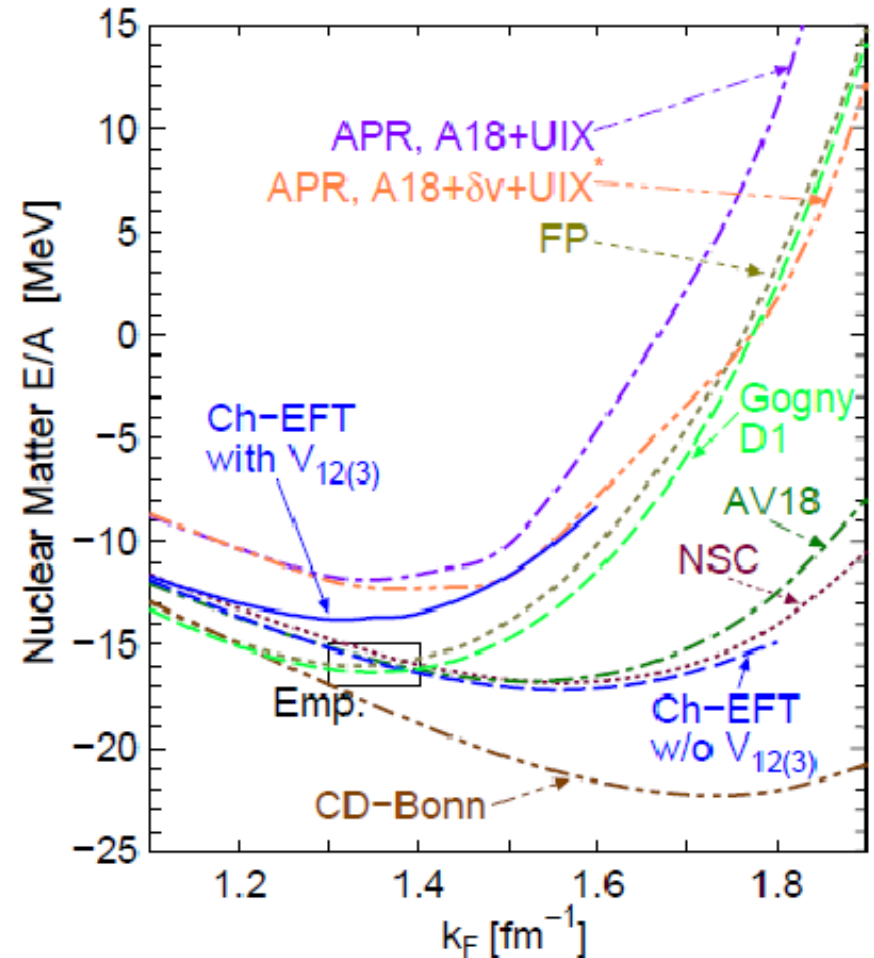
NS matter EOS with hyperons



*Non-linear term dep. +Atomic shift fit:
Tsubakihara, AO, Harada, arXiv:14020979*

Ch-EFT EOS

- Phen. models need inputs from
Experimental Data and/or Microscopic (Ab initio) Calc.
- Recent Ch-EFT EOS is promising !
NN (N3LO)+3NF(N2LO)
Kohno ('13)



M. Kohno, PRC 88 ('13) 064005

“Universal” mechanism of “Three-body” repulsion

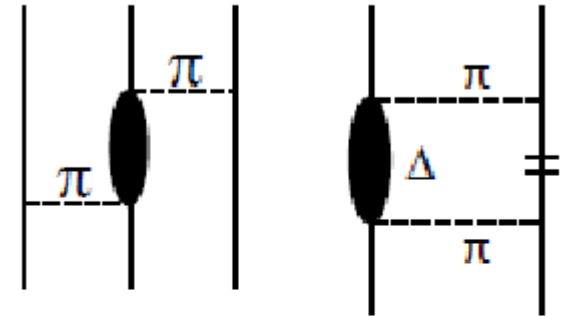
- “Universal” 3-body repulsion is necessary to support NS.
Nishizaki, Takatsuka, Yamamoto (‘02)
- Mechanism of “Universal” Three-Baryon Repulsion.
 - “ σ ”-exchange $\sim$ two pion exch. w/ res.
 - Large attraction from two pion exchange is suppressed by the Pauli blocking in the intermediate stage.
Kohn ('13)

“Universal” TBR

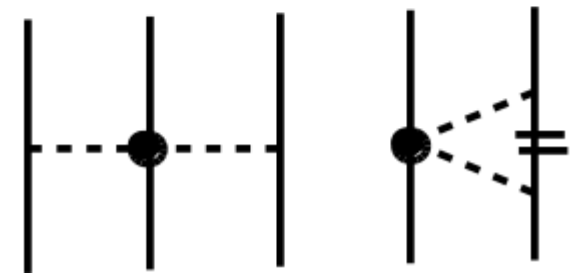
- *Coupling to Res. (hidden DOF)*
- *Reduced “ σ ” exch. pot. ?*

How about YNN or YYN ?

Physical Picture

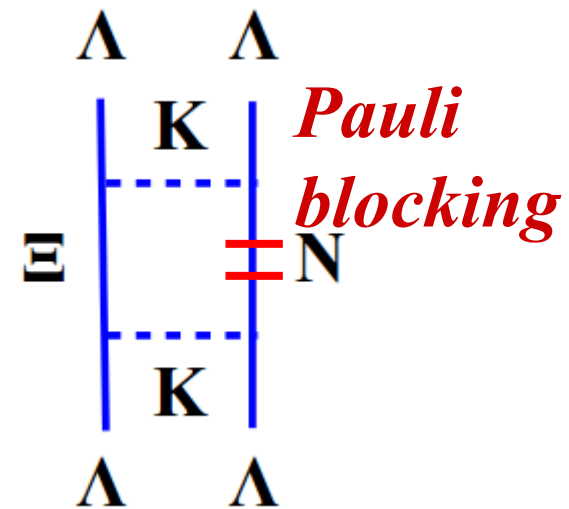


χ EFT



$\Lambda\Lambda$ interaction in vacuum and in nuclear medium

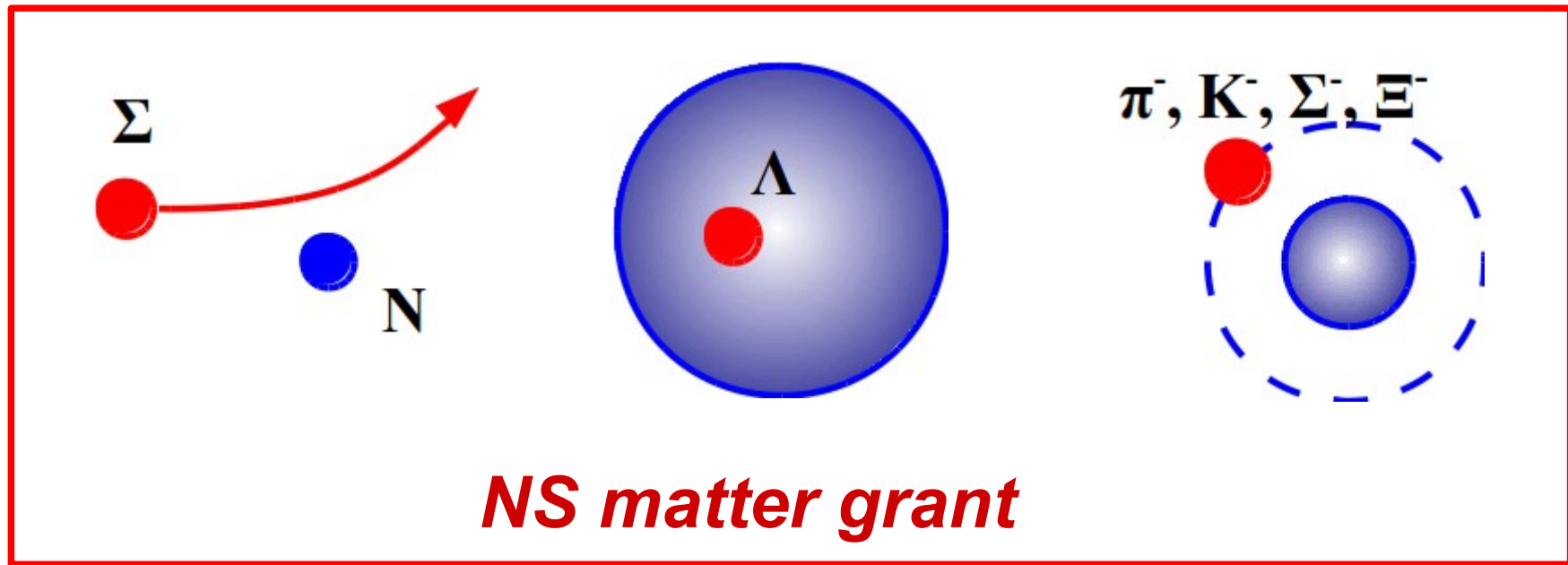
- Vacuum $\Lambda\Lambda$ interaction may be theoretically accessible
Lattice QCD calc. HAL QCD ('11) & NPLQCD ('11)
- In-medium $\Lambda\Lambda$ interaction may be experimentally accessible
 - Bond energy of ${}^6_{\Lambda\Lambda}\text{He}$
 $\Delta B_{\Lambda\Lambda} = 1.0 \text{ MeV} \rightarrow 0.6 \text{ MeV}$
 - $a_0(\text{Nagara fit}) = -0.575 \text{ fm}$ ($\Delta B_{\Lambda\Lambda} = 1.0 \text{ MeV}$)
Hiyama et al. ('02)
- Difference of vacuum & in-medium $\Lambda\Lambda$ int. would inform us $\Lambda\Lambda\text{N}$ int. effects.
 - $\Lambda\Lambda$ - ΞN couples in vacuum
 - Coupling is suppressed in ${}^6_{\Lambda\Lambda}\text{He}$



Is there Any way to access vacuum $\Lambda\Lambda$ int. experimentally ?

Interactions btw short-lived hadrons

- Scattering, Nuclear bound state, Atomic shift



- Exotic hadron spectroscopy

- Correlations from heavy-ion collisions

- STAR data of $\Lambda\Lambda$ correlation

This Lecture

■ Introduction

- Massive NS and NS matter EOS
- Strangeness in NS matter
- “Universal” Three-Baryon Repulsion

■ Constraint on $\Lambda\Lambda$ interaction from HIC

- $\Lambda\Lambda$ correlation in heavy-ion collisions
- Constraint on $\Lambda\Lambda$ interaction from HIC data

■ Discussion

- Can we see the difference btw vacuum and in-medium $\Lambda\Lambda$ interaction ?
- Do we see H in $\Lambda\Lambda$ correlation ?

■ Summary

Where is the $S=-2$ dibaryon ($uuddss$) “H” ?

- Jaffe's prediction (1977)
→ 80 MeV below $\Lambda\Lambda$
(strong attraction from color mag. int.)

- Double hypernuclei $_{\Lambda\Lambda}^6\text{He}$ (Nagara)
→ No deeply bound “H”

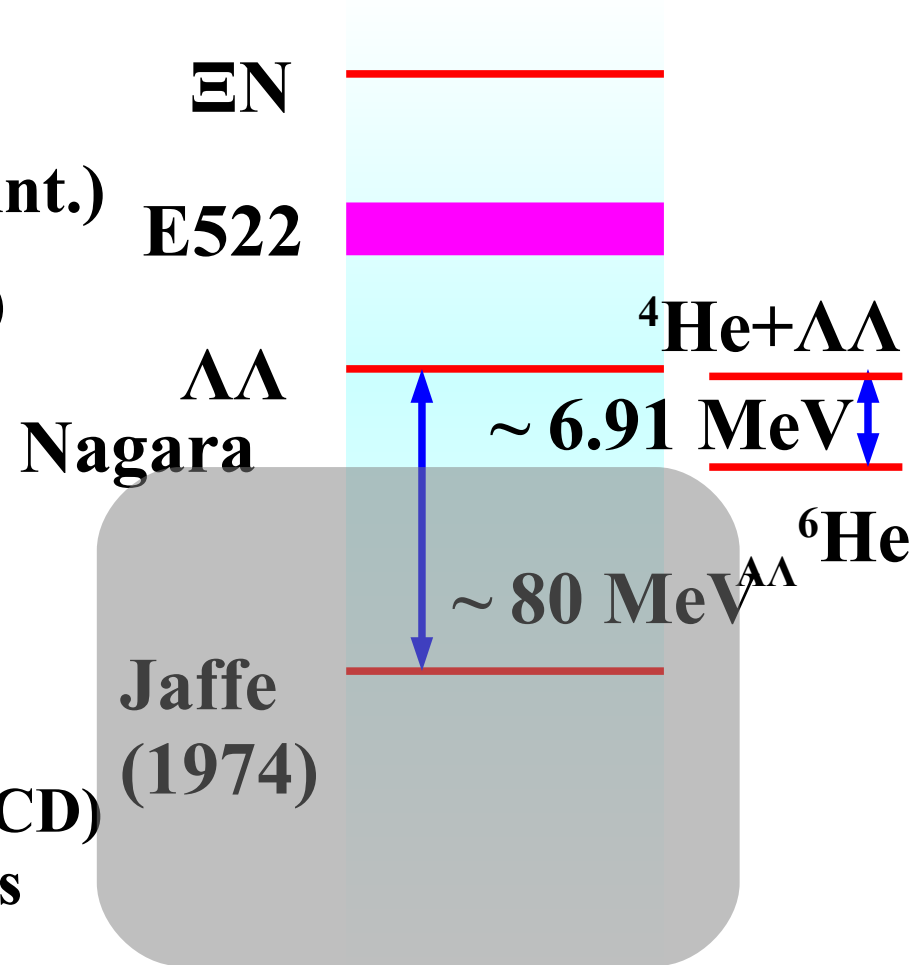
- Resonance or Bound “H” ?

- KEK-E522 (Yoon et al., ('07))
→ “bump” at $E_{\Lambda\Lambda} \sim 15 \text{ MeV}$
- Lattice QCD (HAL QCD & NPLQCD)
→ bound H at large ud quark mass

- How about HIC ?

- RHIC & LHC = Hadron Factory including Exotics
- “H” would be formed as frequently as stat. model predicts.

Cho, Furumoto, Hyodo, Jido, Ko, Lee, Nielsen, AO, Sekihara, Yasui, Yazaki
(*ExHIC Collab.*), *PRL('11)212001; arXiv:t:1107.1302*



Nagara event

■ $\Lambda\Lambda$ ${}^6\text{He}$ hypernuclei

Takahashi et al., PRL87('01)212502
(KEK-E373 experiment)

Lambpha

$$m({}_{\Lambda\Lambda}^6\text{He}) = 5951.82 \pm 0.54 \text{ MeV}$$

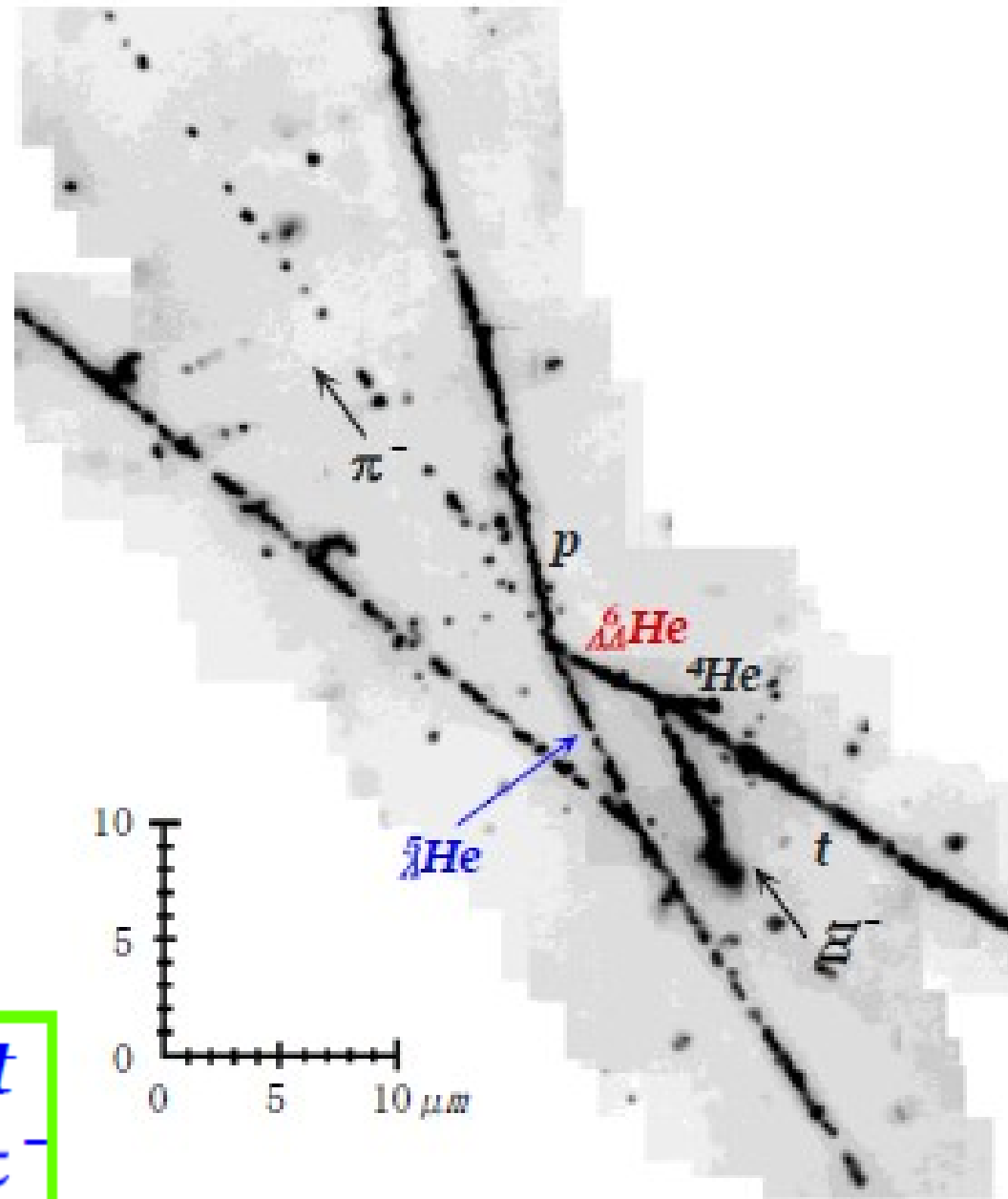
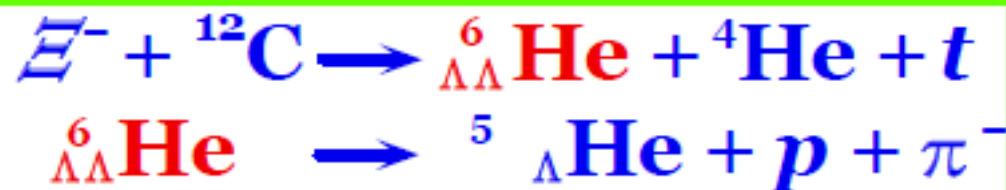
$$B_{\Lambda\Lambda} = 7.25 \pm 0.19_{-0.11}^{+0.18} \text{ MeV}$$

$$\Delta B_{\Lambda\Lambda} = 1.01 \pm 0.20_{-0.11}^{+0.18} \text{ MeV}$$

(assumed $B_{\Xi^-} = 0.13 \text{ MeV}$)

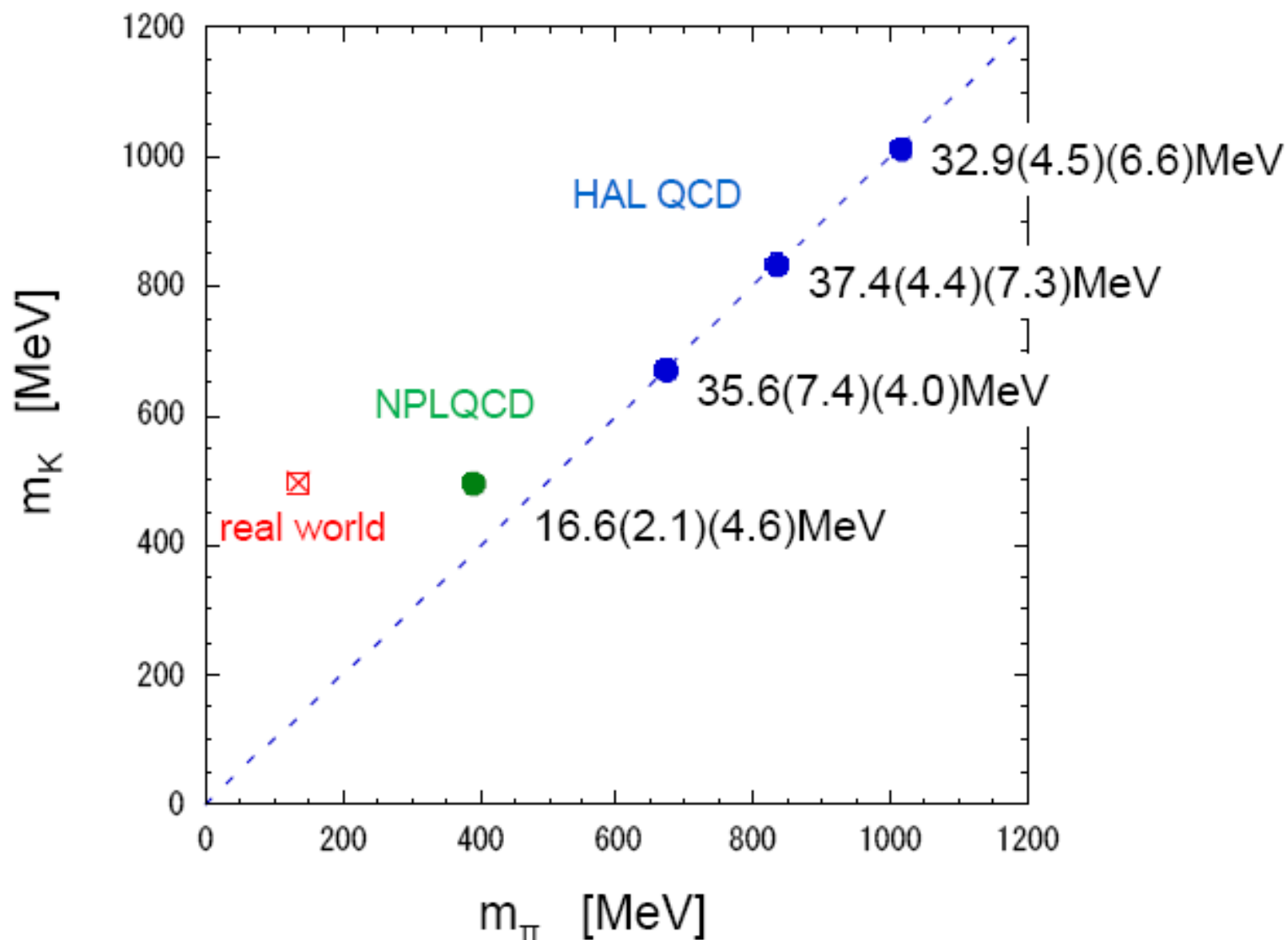
$$\rightarrow B_{\Lambda\Lambda} = 6.91 \text{ MeV}$$

(PDG modified(updated)
 Ξ^- mass)



Lattice QCD predicts bound “H”

- “H” bounds with heavy π ($M_\pi > 400$ MeV)



NPLQCD Collab., PRL 106 (2011) 162001; HAL QCD Collab., PRL 106 (2011) 162002

Hadron-Hadron correlation in HIC

- Hanbury-Brown and Twiss Effects for free bosons

$$C(q) = \int d\mathbf{x} \frac{\exp(-\mathbf{x}^2/4R^2)}{(4\pi R^2)^{-3/2}} \underbrace{|\sqrt{2} \cos \mathbf{q} \cdot \mathbf{x}|^2}_{|\Psi|^2} = 1 + \exp(-4q^2 R^2)$$

- Correlation func. $\sim \int \text{Source } \mathbf{x} |\Psi|^2$
→ If source is known, corr. fn. tells us w.f. or interaction.

Bauer, Gelbke, Pratt ('92); Lednicky ('09).

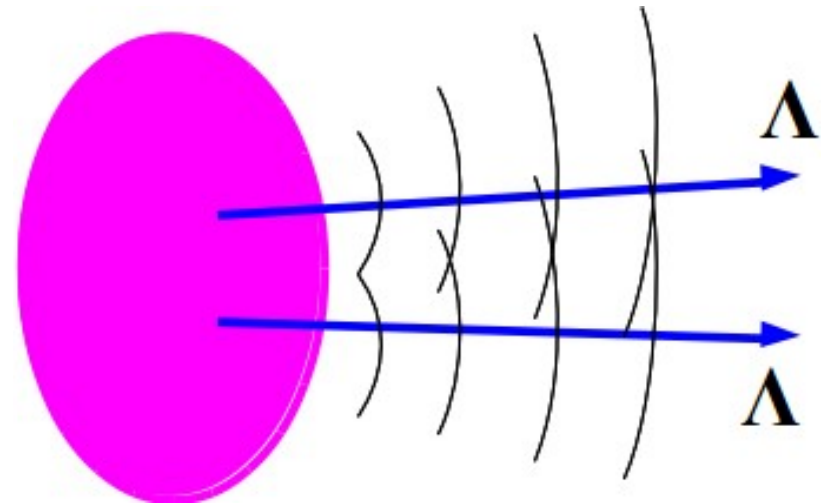
- $\Lambda\Lambda$ correlation is measured in (K-,K+) reaction

C.J.Yoon et al. (KEK-E522)('07); J.K.Ahn et al. (KEK-E224)

- STAR measured $\Lambda\Lambda$ correlation at RHIC

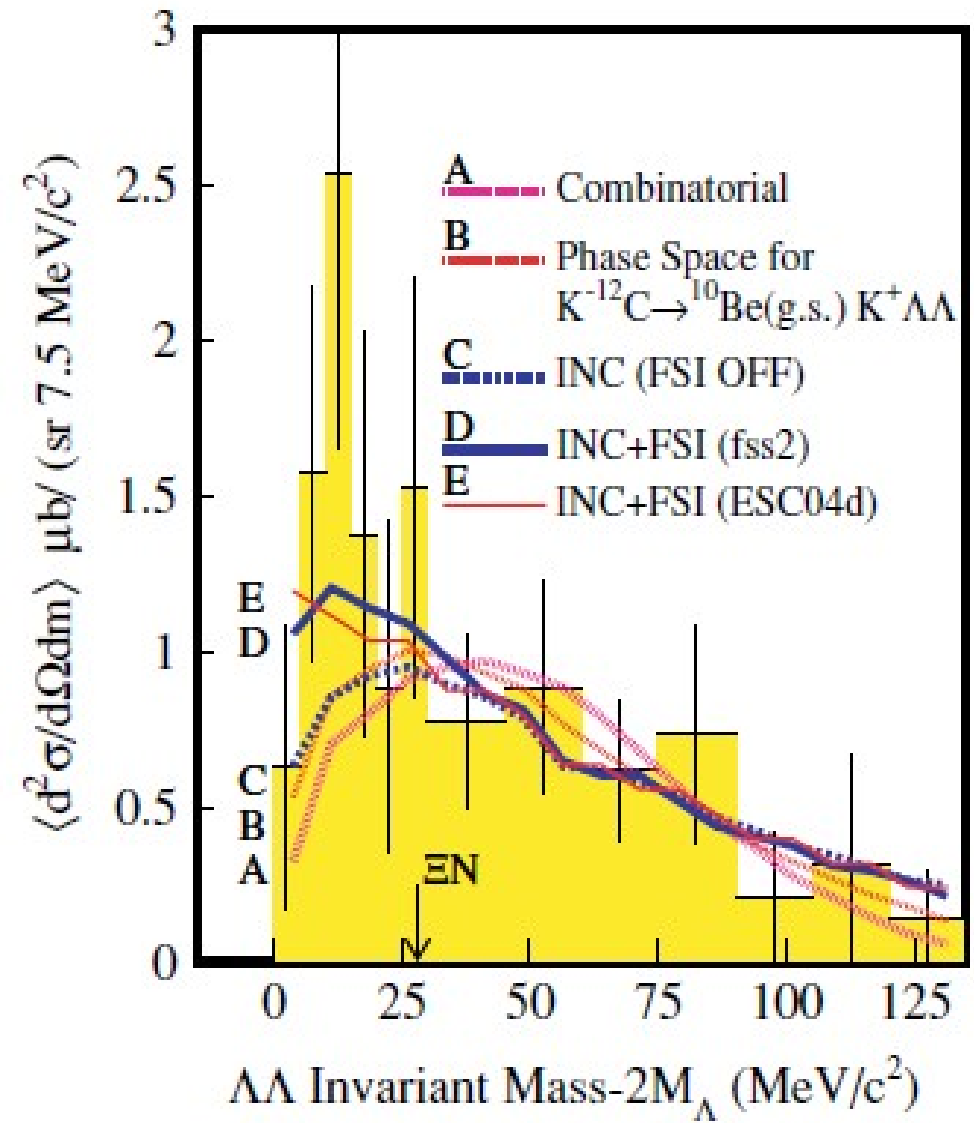
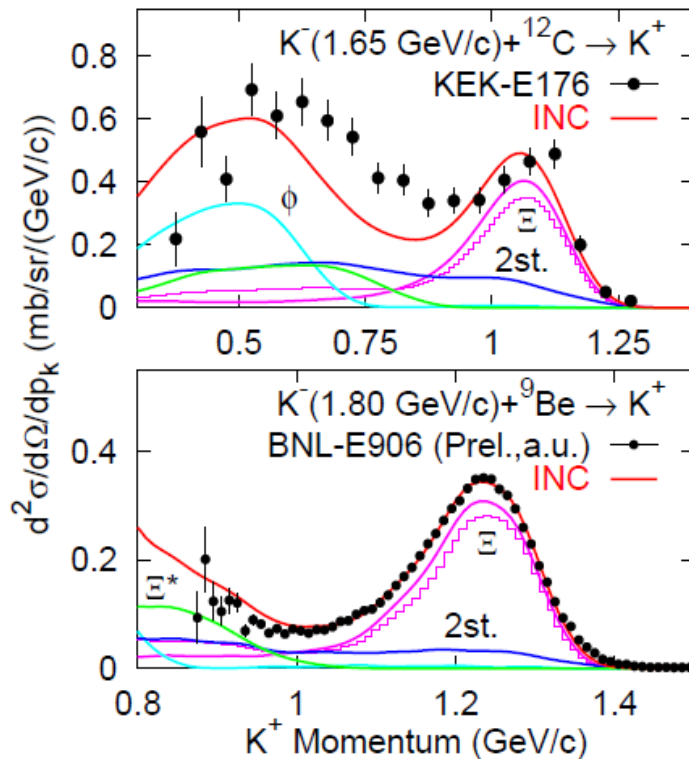
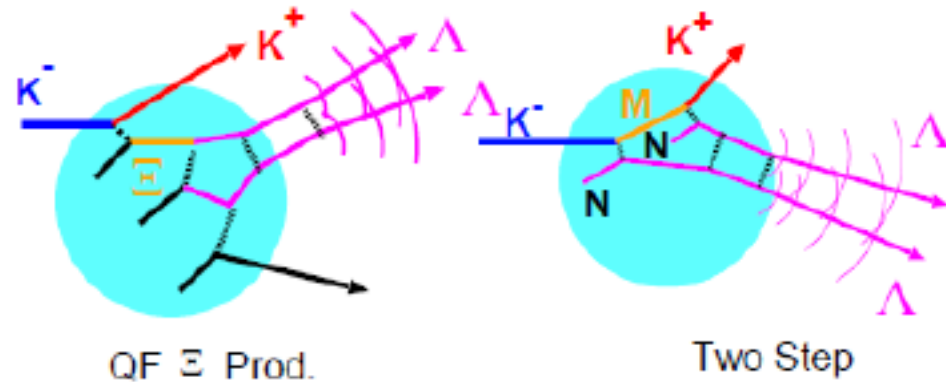
N. Shah et al.('12)

*Let's try to constrain
 $\Lambda\Lambda$ interaction !*



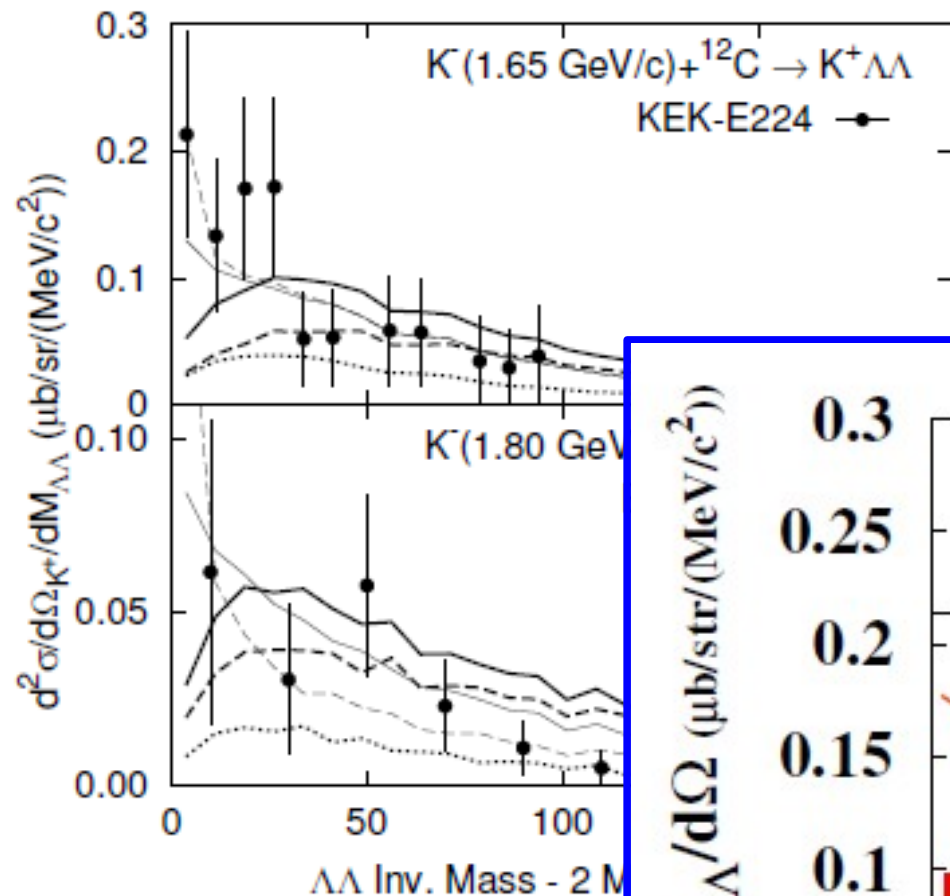
$\Lambda\Lambda$ correlation from $(K^-, K^+ \Lambda\Lambda)$ reaction

■ Enhancement at $\sim 2 M(\Lambda) + 10$ MeV,

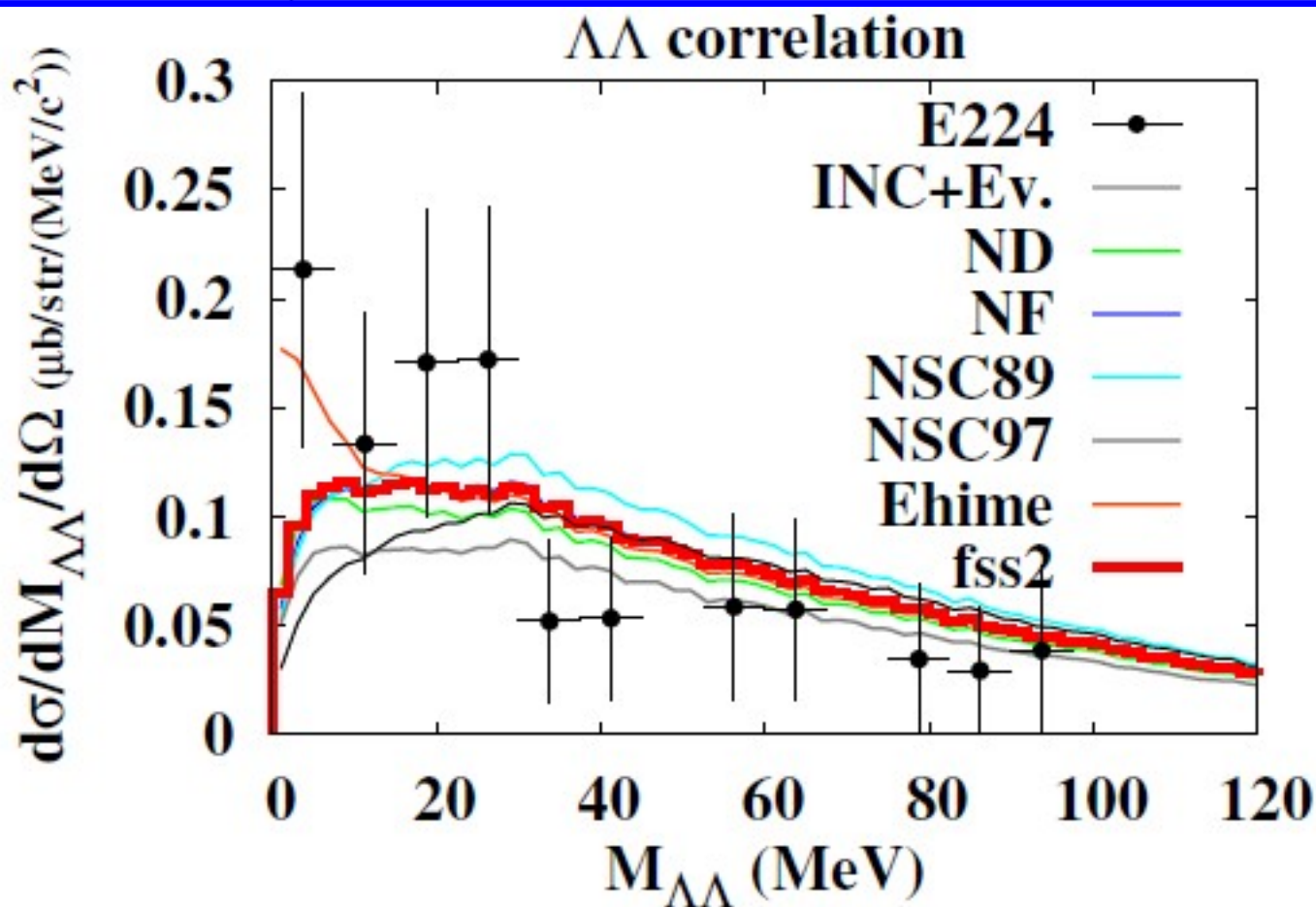


C.J.Yoon, ..., (KEK-E522), AO, PRC75 (2007) 022201(R)
J. K. Ahn et al. (KEK-E224).

$\Lambda\Lambda$ Invariant Mass Spectrum



$\Lambda\Lambda$ int. constrained in HIC
are consistent with $\Lambda\Lambda$ inv.
mass spectrum in $(K^-, K^+ \Lambda\Lambda)$



AO, Hirata, Nara,
Shinmura, Akaishi,
NPA684(2001), 595c

$\Lambda\Lambda$ correlation in HIC

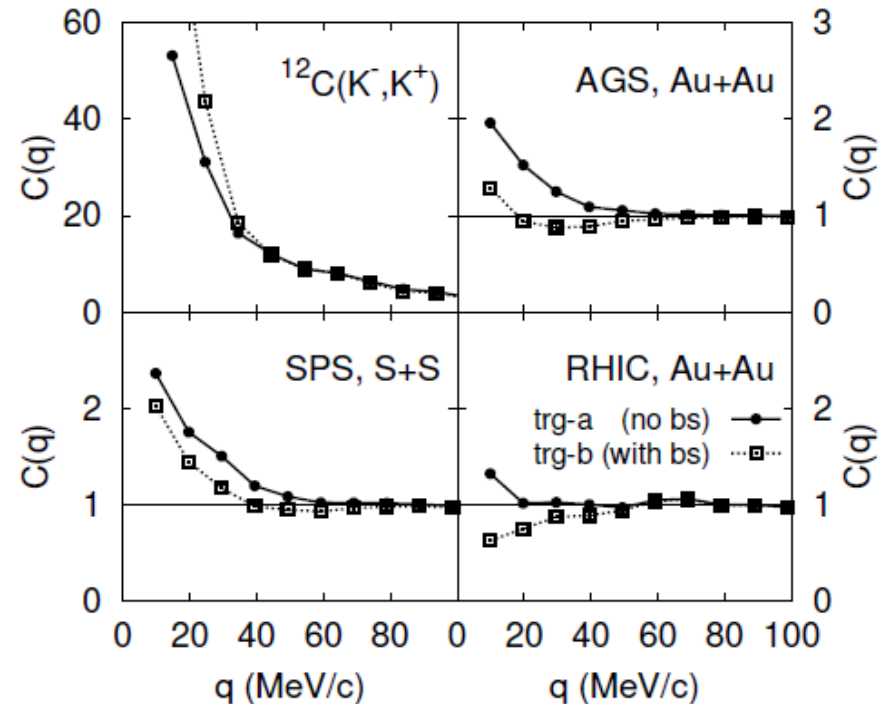
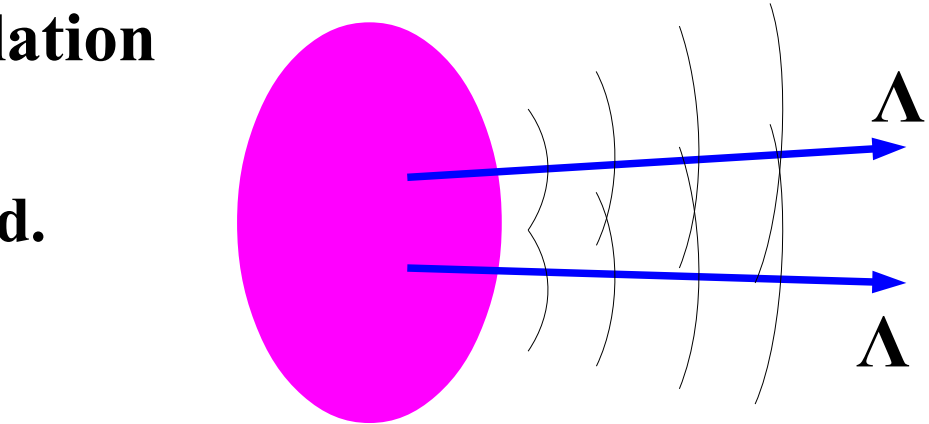
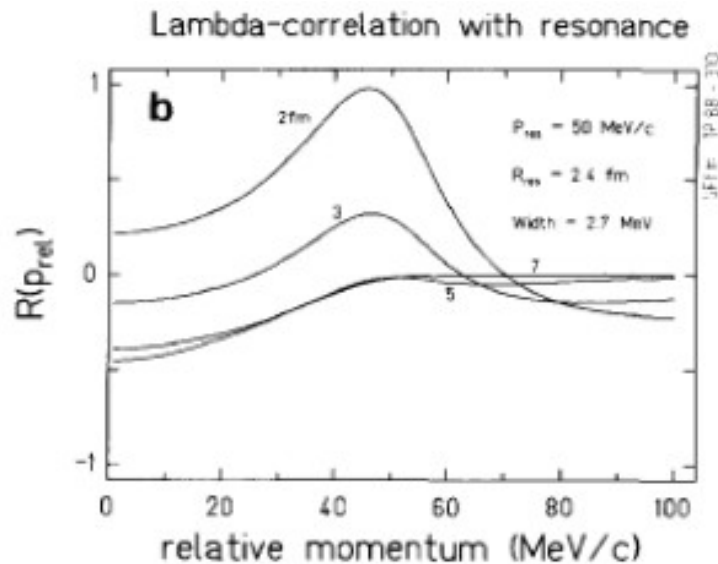
■ Merit of HIC to measure $\Lambda\Lambda$ correlation

- Source is “Simple and Clean” !
T, μ , flow, size, ... are well-analyzed.

- Nearly Stat. prod.
→ Many exotics will be produced.

Cho et al.(ExHIC Collab.) ('11)

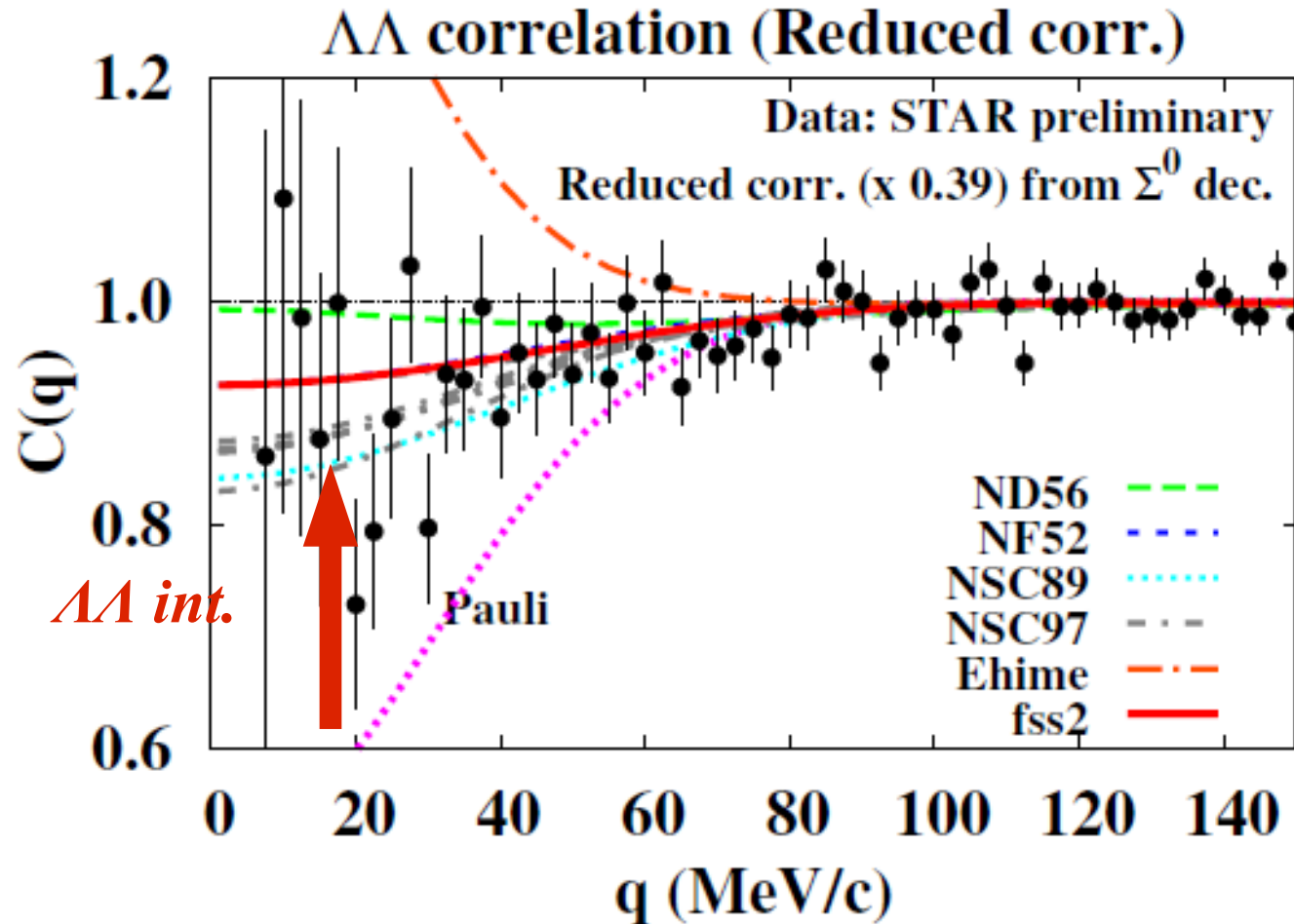
- Discovery of “H” and/or Constraint on $\Lambda\Lambda$ int.



C. Greiner, B. Muller, PLB219('89)199.

AO, Hirata, Nara, Shinmura, Akaishi, NPA670('00)297c

$\Lambda\Lambda$ correlation at RHIC



Data (STAR prelim.): N. Shah et al.('12), Cal.: AO for ExHIC ('13)

$\Lambda\Lambda$ interaction models

■ Boson exchange potentials

- Nijmegen potentials: various versions

Rijken et al., ('77-'10)

Hard core: Nijmegen model D & F (ND, NF)

Soft core: Nijmegen soft core '89 & '97 (NSC89, NSC97)

Extended soft core: ESC08

- Ehime potential: would be too attractive.

Ueda et al., ('98)

Ehime fits old double Λ hypernucl. data, $\Delta B_{\Lambda\Lambda} = 4$ MeV

■ Quark cluster model

- fss2

Fujiwara, Kohno, Nakamoto, Suzuki ('01)

Short range repulsion from quark Pauli blocking & OGE

Core is softer due to non-locality

■ Modified Nijmegen potentials fitting Nagara data.

Filikhin, Gal ('02), Hiyama et al.('02)

- Potential Fitting Nagara data $\Delta B_{\Lambda\Lambda} = 1.0$ MeV

$\Lambda\Lambda$ interaction models

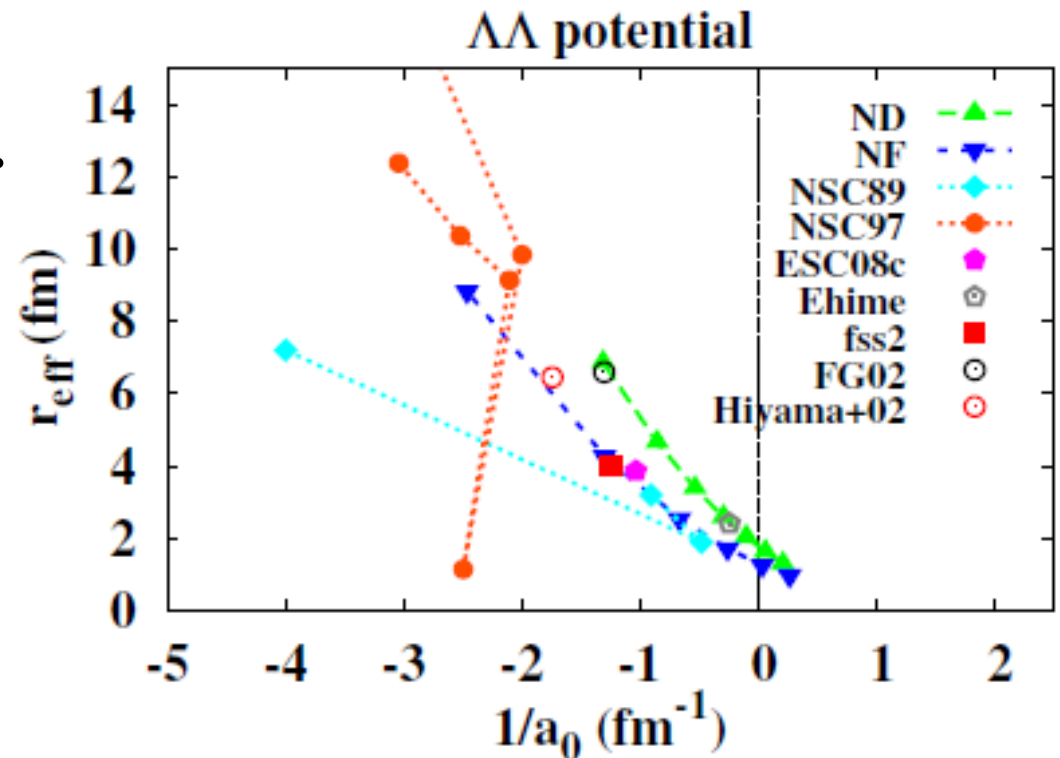
■ Low energy scattering parameters, (a_0, r_{eff})

$$k \cot \delta = -\frac{1}{a_0} + \frac{1}{2} r_{\text{eff}} k^2 + O(k^4)$$

- $a_0 > 0$ (bound region), $a_0 < 0$ (no bound region)

■ Potential parameters

- Hard core radius (ND, NF),
cutoff mass (NSC89),
spin dependence (NSC97a-f).



$\Lambda\Lambda$ correlation function

■ Two particle correlation function

Koonin ('77)

$$C_2(Q, K) = \frac{W_2(k_1, k_2)}{W_1(k_1)W_1(k_2)} = \frac{\int d^4x_1 d^4x_2 S(x_1, K) S(x_2, K) |\Psi_{12}(Q, x_1 - x_2 - (t_2 - t_1)K/m)|^2}{\int d^4x_1 d^4x_2 S(x_1, k_1) S(x_1, k_2)}$$

- $W_1(k)$, $W_2(k_1, k_2)$: 1 & 2 partcl. dist., $S(x, k)$: phase spc. dist.
 $Q=(k_1-k_2)$, $K=(k_1+k_2)/2$
- Wave fn. Ψ (assumption: only the s-wave partial wave is modified.)

$$\Psi_s = \sqrt{2} \left[\cos \mathbf{Q} \cdot \mathbf{r} / 2 + \chi_Q(r) - j_0(Qr/2) \right]$$

$$\Psi_t = \sqrt{2} i \sin \mathbf{Q} \cdot \mathbf{r} / 2$$

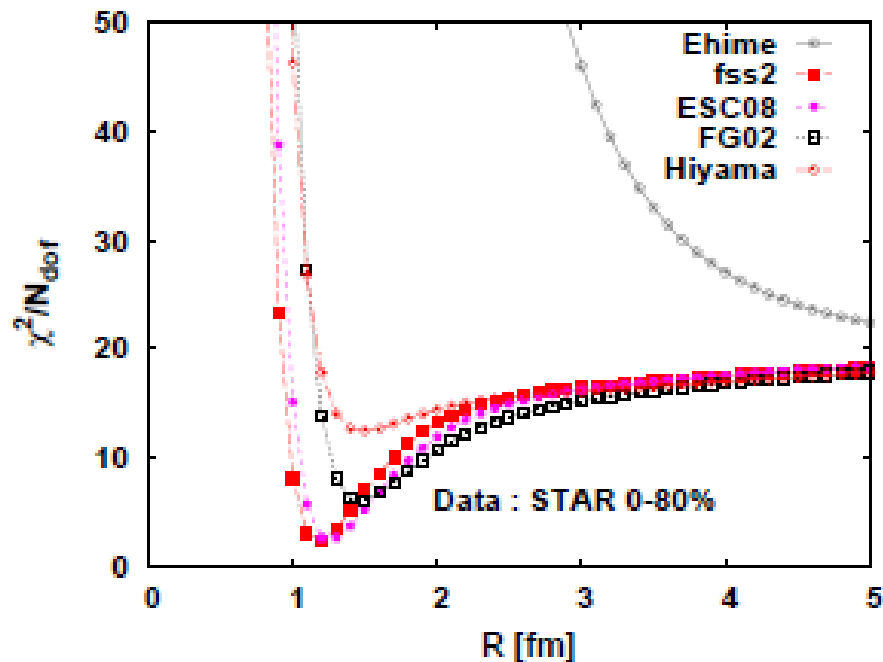
$$|\Psi_{12}|^2 = 1 - \frac{1}{2} \cos \mathbf{Q} \cdot \mathbf{r} + \cos(\mathbf{Q} \cdot \mathbf{r} / 2) \Delta \chi_Q(r) + [\Delta \chi_Q(r)]^2$$

$$\Delta \chi_Q(r) = \chi_Q(r) - j_0(Qr/2)$$

Static Spherical Source

■ Correlation fn. with static, spherical gaussian source

$$C_{\Lambda\Lambda}(Q) \simeq 1 - \frac{1}{2} \exp(-Q^2 R^2) + \frac{1}{2} \int dr S_{12}(r) (|\chi_Q(r)|^2 - |j_0(Qr/2)|^2)$$



Morita, Furumoto, AO (to be submitted)

Geometry & Flow Effects

■ Boost invariant source with flow effects

S. Chapman, P. Scotto, U. Heinz, Heavy Ion Phys. 1, 1 (1995).

$$S(x, k) = \frac{m_T \cosh(y - Y_L)}{(2\pi)^3 \sqrt{2\pi} (\Delta\tau)^2} n_f(u \cdot k, T) \exp \left[-\frac{(\tau - \tau_0)}{2(\Delta\tau)^2} - \frac{x^2 + y^2}{2R^2} \right]$$

● Fluid velocity

$$u^t = \cosh Y_T \cosh Y_L$$

$$u^z = \cosh Y_T \sinh Y_L$$

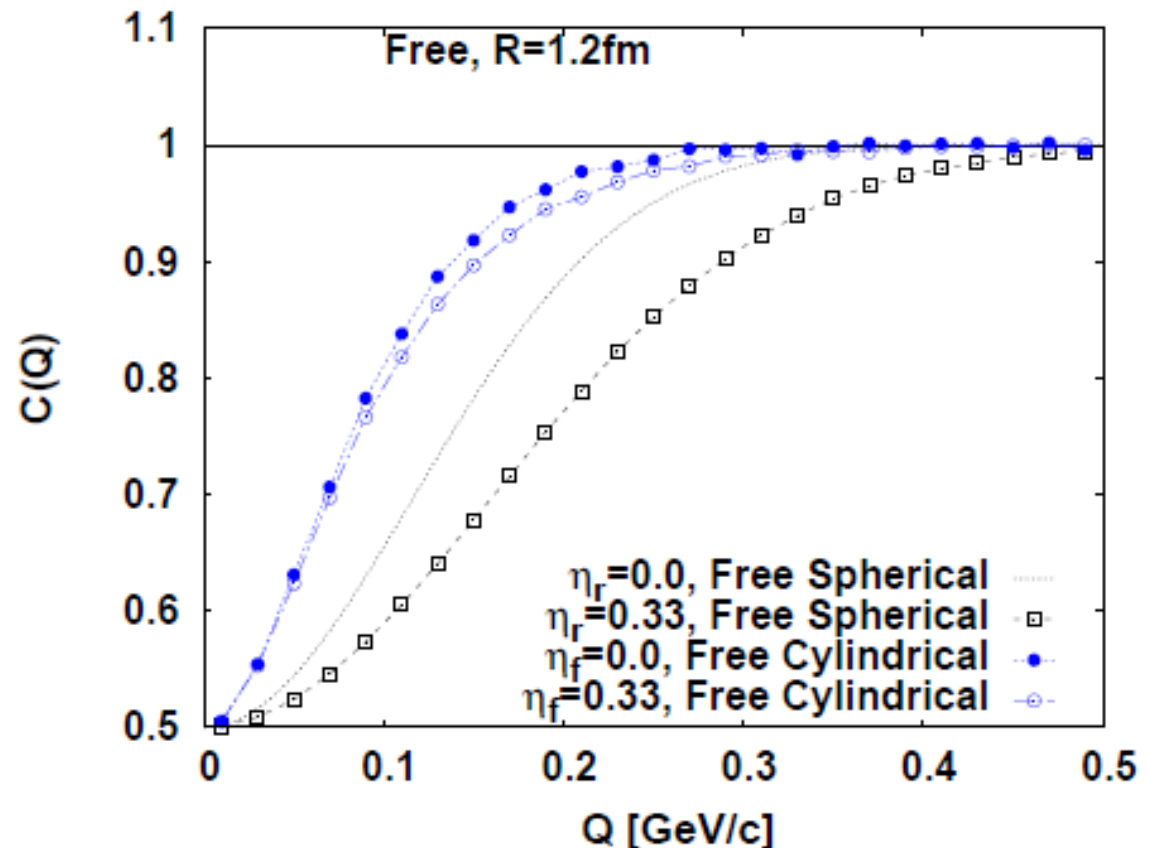
$$u^x = \sinh Y_T \cos \phi$$

$$u^y = \sinh Y_T \sin \phi.$$

$$Y_L = \eta_s = \frac{1}{2} \ln \frac{t + z}{t - z}$$

$$Y_T = \eta_f \frac{r_T}{R}$$

Transverse flow $\eta_f = 0.33$
(from m_T spectrum)

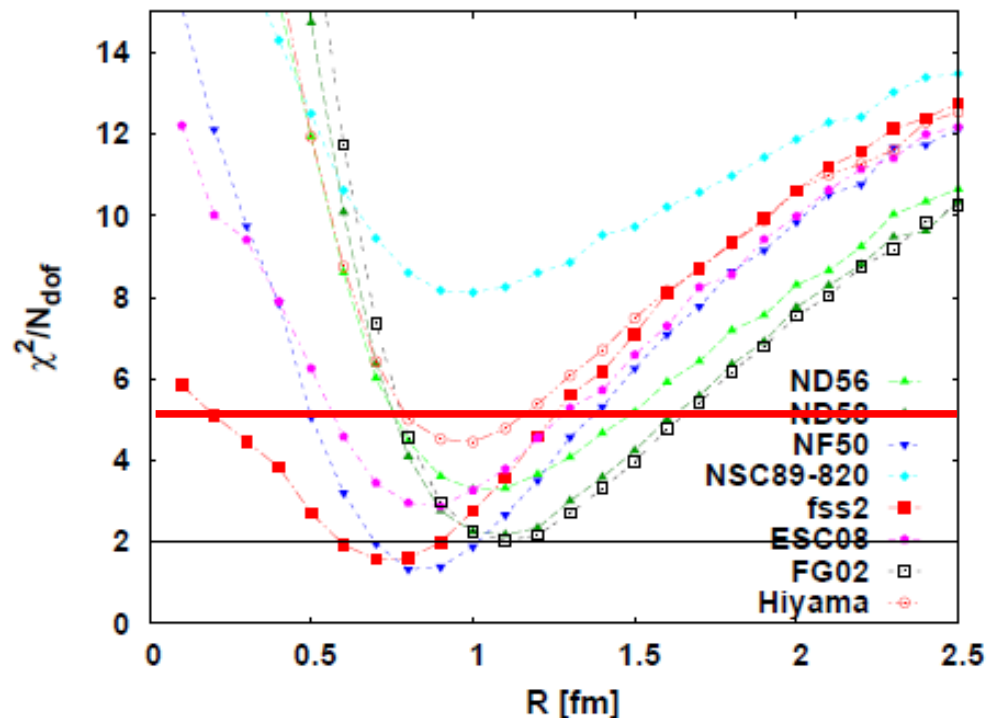


Morita, Furumoto, AO (to be submitted)

$\Lambda\Lambda$ correlation with flow effects

■ Results with flow effects

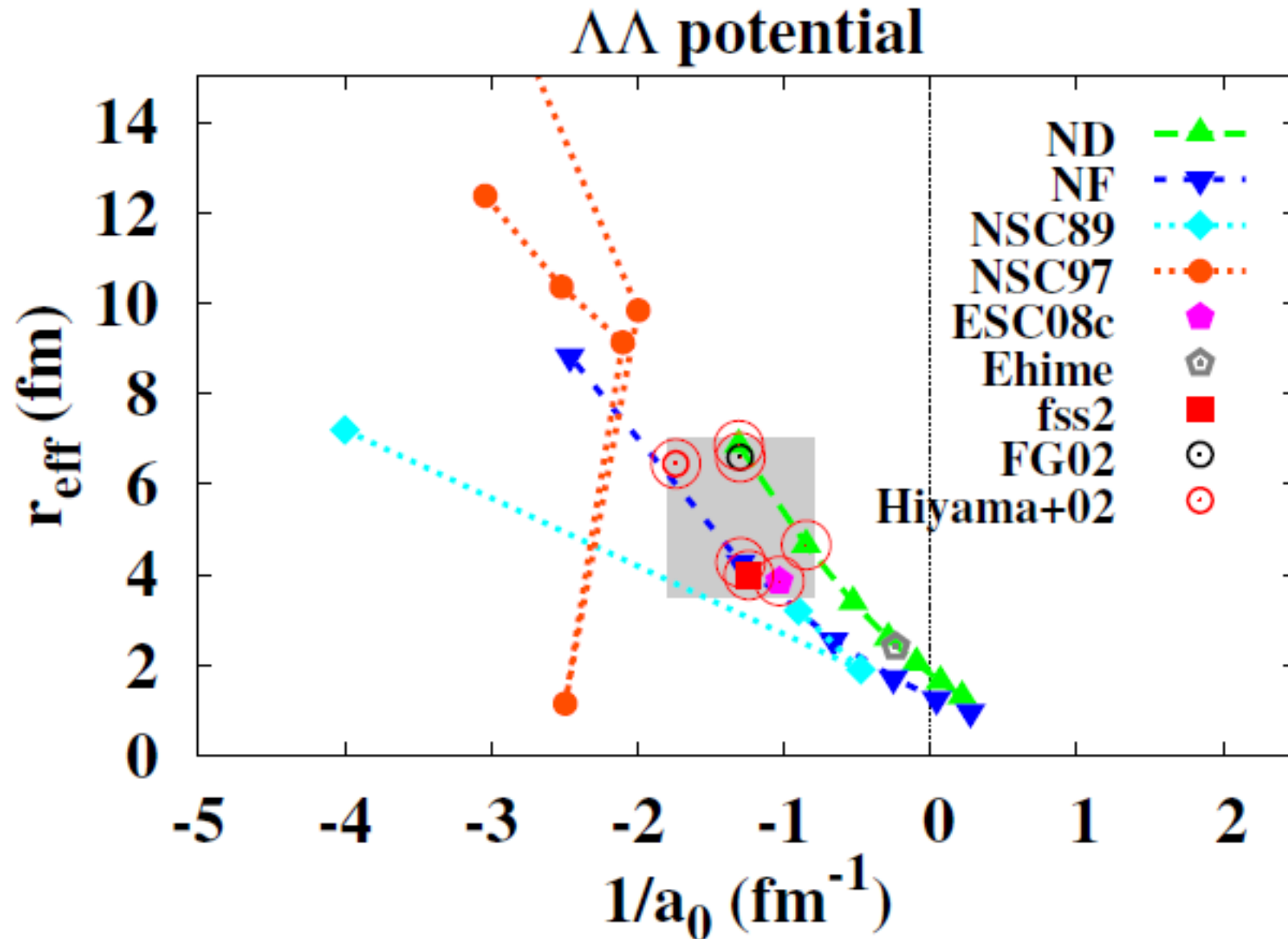
- Optimal transverse source size $R \sim (0.8-1.1)$ fm
- HBT source size is interpreted as the “homogeneity length”, but it is still too small compared with the proton source size, $R_p \sim (2-4)$ fm.



Cylindrical

Morita, Furumoto, AO (to be submitted)

Preferred $\Lambda\Lambda$ interactions



Preferred $\Lambda\Lambda$ interaction parameters

$$-1.8 \text{ fm}^{-1} < 1/a_0 < -0.8 \text{ fm}^{-1}, \quad 3.5 \text{ fm} < r_{\text{eff}} < 7 \text{ fm}$$

■ Introduction

- Massive NS and NS matter EOS
- Strangeness in NS matter
- “Universal” Three-Baryon Repulsion

■ Constraint on $\Lambda\Lambda$ interaction from HIC

- $\Lambda\Lambda$ correlation in heavy-ion collisions
- Constraint on $\Lambda\Lambda$ interaction from HIC data

■ Discussion

- Can we see the difference btw vacuum and in-medium $\Lambda\Lambda$ interaction ?
- Do we see H in $\Lambda\Lambda$ correlation ?

■ Summary

Comparison with In-medium interaction.

■ $\Lambda\Lambda$ interactions from Nagara event ($\Delta B_{\Lambda\Lambda} = 1.0$ MeV)

- Hiyama, Kamimura, Motoba, Yamada, Yamamoto ('02)

$$(a_0, r_{\text{eff}}) = (-0.575 \text{ fm}, 6.45 \text{ fm})$$

- Filikhin, Gal ('02)

$$(a_0, r_{\text{eff}}) = (-0.77 \text{ fm}, 6.59 \text{ fm})$$

■ Ξ^- mass is updated by PDG

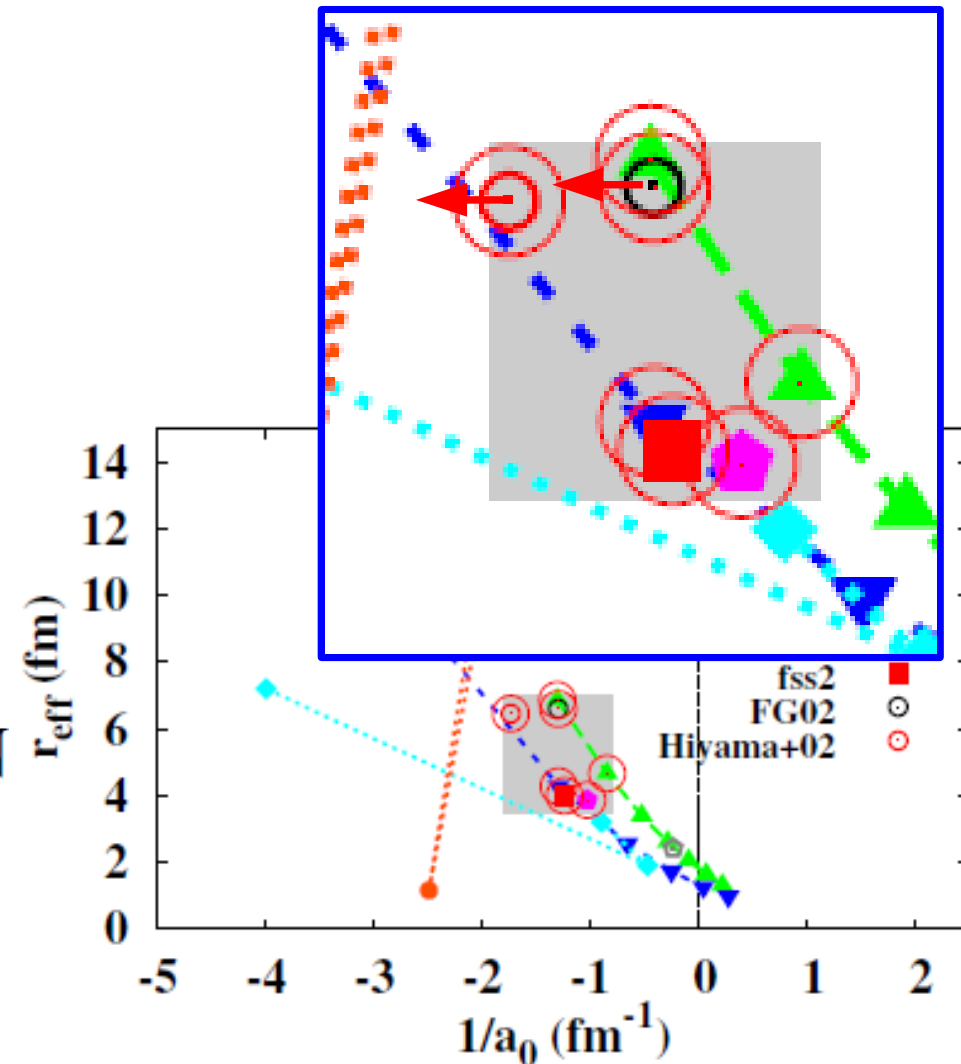
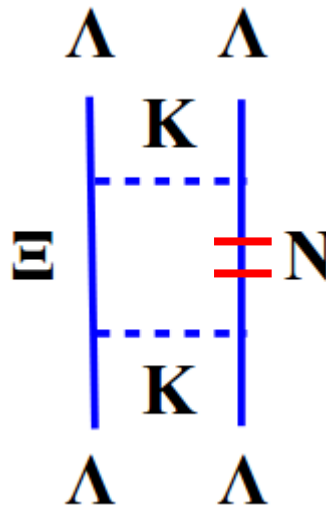
- Bond energy is updated

$$\Delta B_{\Lambda\Lambda} = 0.67 \text{ MeV}$$

→ a_0 will be reduced

by 10-20 %

$$a_0 \sim -(0.5-0.65) \text{ fm}$$



Do we see H as a resonance ?

- Deeply bound H is ruled out by double Λ hypernuclear mass.

$$M_H > 2 M_\Lambda - 6.91 \text{ MeV}$$

- Existence of H as a resonance is not ruled out.

→ Let's try to find it !

- Procedure

- Assume the bump comes from H, and give (E_H, Γ_H) .
- Compare the bump height with statistical model yield.
- If H exists at low E ($E=(1-2) \text{ MeV}$), we can find the signal by reducing the error by a factor of two.

Morita, Furumoto, AO (to be submitted)

Other source of correlation ?

- $\Lambda\Lambda$ correlation would be modified by
 - Feed down effects from Ξ and Σ^0 decay,
 $\Xi^- \rightarrow \pi^- \Lambda$ (detectable)
 $\Sigma^0 \rightarrow \gamma \Lambda$ (will be detectable at LHC (Kwon et al.))
 Σ^0 effects can be taken care of by multiplying 0.41 to (C-1),
and preferred $V(\Lambda\Lambda)$ are similar to the present result.
 - If feed down Λ is included, the correlation is affected by the parent pair interaction.
E.g. pp correlation is significantly affected by $V(p\Lambda)$.
($\Lambda \rightarrow \pi^- p$ and no Coulomb suppression in $p\Lambda$ channel.)

Since there is no Coulomb suppression in $\Lambda\Lambda$ pair, parent pair interaction effects may be less serious than in pp correlation.
- Further investigation is necessary to pin down $\Lambda\Lambda$ interaction more precisely.

Summary of Lecture(?) 5

- We need additional repulsion to solve massive neutron star puzzle related to strangeness hadrons.
 - Need exp. data and ab initio calc.
J-PARC exp. / Lattice BB and BBB int. / Ch-EFT
- We have constrained $\Lambda\Lambda$ low energy scattering parameters using $\Lambda\Lambda$ correlation data from STAR collaboration.
 - Optimal source size & flow parameter are fixed by using correlation and pT spectrum.
 - Preferred scattering parameters are found to be in the range,
 $-1.8 \text{ fm}^{-1} < 1/a_0 < -0.8 \text{ fm}^{-1}$, $3.5 \text{ fm} < r_{\text{eff}} < 7 \text{ fm}$
 - Other mechanisms may need to be taken care of.
- Information on $\Lambda\Lambda N$ may be accessible via correlation in HIC
 - In-medium $\Lambda\Lambda$ interaction seems to be weaker than vacuum interaction.

- 中性子、陽子、電子のみからなる中性子星物質を考える。電子の質量を無視すると、核子あたりのエネルギーは、Lecture 1 で示したように

$$E_{\text{NSM}}(\rho) = E_{\text{SNM}}(\rho) + S(\rho)\delta^2 + \frac{\Delta M}{2}\delta + \frac{3}{8}\hbar k_F(1-\delta)^{4/3}$$

と与えられる。ここで $\Delta M = M_n - M_p$ 、 k_F は同じ密度での対称核物質のフェルミ波数である。非対称度 δ は、核子あたりのエネルギーが最小になるように選ばれる。

- 上の表式を導け。
- 核子あたりのエネルギーが最小となる非対称度 δ を求めよ。
(3次方程式を解くこととなる。 $S(\rho)$, k_F , ΔM は与えられているとしてよい。)
- 今回の講義において、中性子物質の物理の課題の中で各自が興味を持った項目をあげ、その理由を述べよ。

単位認定について

- 次の条件で単位を出します。
 - 6コマの講義中、4コマ以上出席。
 - 6コマの講義中、1コマ以上出席し、レポートを提出。
- レポート問題は Lec.2 後にお知らせしたものです。
1つ目の問題を次の問題に置き換えてもよい。
 - SU(3) 不変な相互作用

$$\begin{aligned}\mathcal{L}_{\text{BV}} &= \sqrt{2}\{g_s \text{tr}(M_v) \text{tr}(\bar{B}B) + g_D \text{tr}(\bar{B}\{M_v, B\}) + g_F \text{tr}(\bar{B}[M_v, B])\} \\ &= \sqrt{2}\{g_s \text{tr}(M_v) \text{tr}(\bar{B}B) + g_1 \text{tr}(\bar{B}M_v B) + g_2 \text{tr}(BBM_v)\}\end{aligned}$$

において、バリオン・メソンの結合定数を g_D, g_F, g_s を用いて表せ。

Thank you

$\Lambda\Lambda$ interaction and correlation

■ $\Lambda\Lambda$ int.

- Nijmegen models
Rijken et al.
- quark model (fss2)
Fujiwara et al. ('01)
- Nagara fit
Filikhin, Gal ('02);
Hiyama et al.('02)

■ Source models

- sph. static source
- cylindrical source w/ flow

K. Morita, AO, T. Furumoto (in prep.)

$\Lambda\Lambda$ correlation in HIC and $\Lambda\Lambda$ interaction

■ Two particle correlation from chaotic source

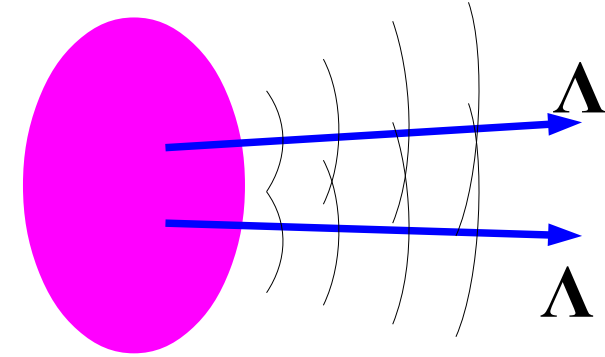
c.f. Bauer, Gelbke, Pratt,

Annu. Rev. Nucl. Part. Sci. 42('92)77.

$$C_{\Lambda\Lambda}(q) = \frac{\int dx_1 dx_2 S(x_1, p+q) S(x_2, p-q) |\psi^{(-)}(x_{12}, q)|^2}{\int dx_1 dx_2 S(x_1, p+q) S(x_2, p-q)}$$

$$\simeq 1 - \frac{1}{2} \exp(-4q^2 R^2) + \frac{1}{2} \int dr S_{12}(r) (|\chi_0(r)|^2 - |j_0(qr)|^2)$$

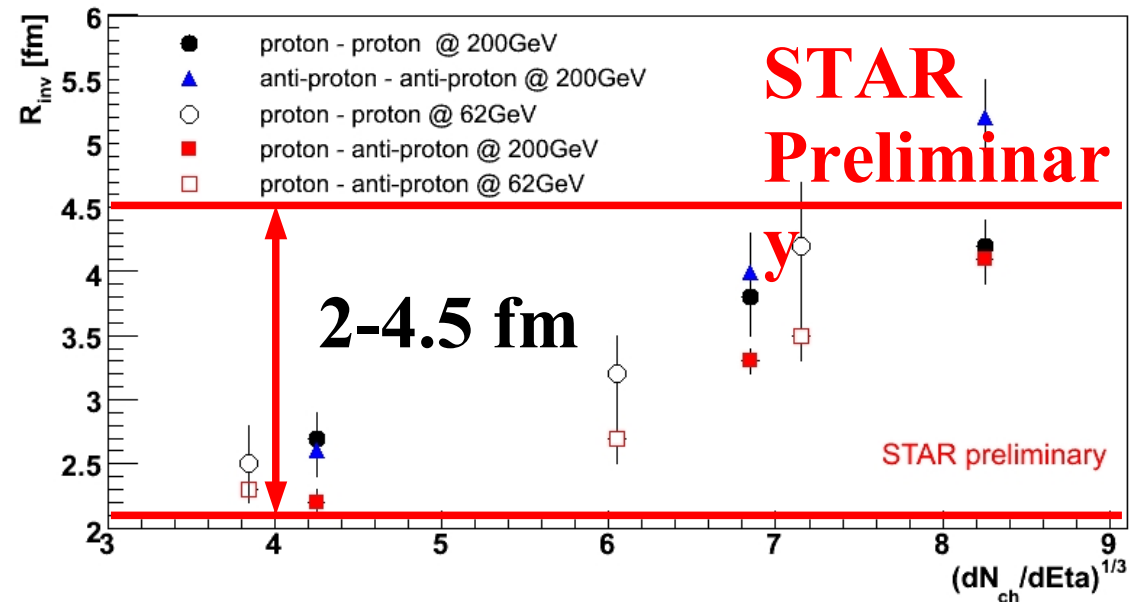
(χ_0 : s-wave wave func., $S_{12}(x) = (2R\sqrt{\pi})^{-3} \exp(-r^2/4R^2)$)



■ Baryon Source size

$R = (2-4.5) \text{ fm}$

● Smaller than π , K source.



$\Lambda\Lambda$ interaction

■ Type of $\Lambda\Lambda$ interaction

- Meson exchange models: Nijmegen model D, F, Soft Core (89, 97)
Nagels, Rijken, de Swart ('77, '79), Maessen, Rijken, de Swart ('89), Rijken, Stoks, Yamamoto ('99)
- Quark cluster model interaction: fss2
Fujiwara, Fujita, Kohno, Nakamoto, Suzuki ('00)
- Phenomenological model: Ehime

■ Two (or three) range gaussian fit results are used in the analysis.

