

# 高密度物質と中性子星の物理 *Physics of Neutron Star Matter*

京大基研 大西 明

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- 中性子星の基本的性質
- 状態方程式を記述する理論模型
- 対称エネルギーと非対称核物質の状態方程式
- ハイパー核物理と高密度核物質の状態方程式
- 中性子星におけるエキゾチック自由度
- Supplementary Contents
  - 実験・観測・理論で解き明かす中性子星物質状態方程式
  - 重イオン衝突とハイパー核から中性子星へ

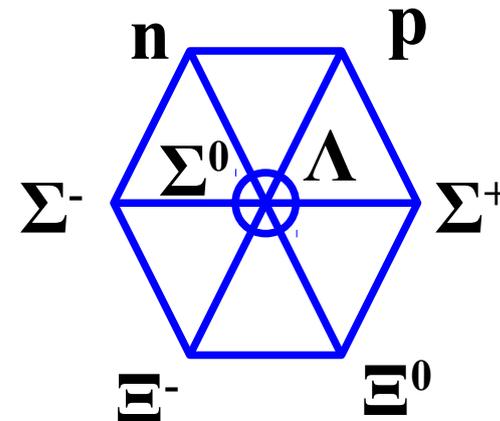
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# Hyperons (Baryons with Strangeness)

- Ground state baryon  $SU(3)_f$  octet ( $J^\pi=1/2^+$ )

Baryon	M(Mev)	S	Comp.
n	940	0	udd
p	938	0	uud
$\Lambda$	1116	-1	$(uds-dus)/\sqrt{2}$
$\Sigma^+$	1189	-1	uus
$\Sigma^0$	1193	-1	$(uds+dus)/\sqrt{2}$
$\Sigma^-$	1197	-1	dds
$\Xi^0$	1315	-2	uss
$\Xi^-$	1321	-2	dss



# $SU(3)_f$ transformation

- Fundamental triplet  $(u,d,s)^T = \mathbf{q} \rightarrow \mathbf{q}' = \mathbf{U} \mathbf{q}$  ( $\mathbf{U} \in SU(3)$ )
- Diquark  $\mathbf{D}_i = \varepsilon_{ijk} \mathbf{q}_j \mathbf{q}_k \rightarrow \mathbf{D}' = \mathbf{D} \mathbf{U}^+$
- Baryon octet  $\mathbf{B}_{ij} = \mathbf{D}_j \mathbf{q}_i \rightarrow \mathbf{B}' = \mathbf{U} \mathbf{B} \mathbf{U}^+$

$$\begin{pmatrix} [ds]u & [su]u & [ud]u \\ [ds]d & [su]d & [ud]d \\ [ds]s & [su]s & [ud]s \end{pmatrix} = \begin{pmatrix} \frac{\Lambda}{\sqrt{6}} + \frac{\Sigma^0}{\sqrt{2}} & \Sigma^+ & p \\ \Sigma^- & \frac{\Lambda}{\sqrt{6}} - \frac{\Sigma^0}{\sqrt{2}} & n \\ \Xi^- & \Xi^0 & -\frac{2\Lambda}{\sqrt{6}} \end{pmatrix}$$

# $SU(3)_f$ transformation

- Fundamental triplet  $(u,d,s)^T = \mathbf{q} \rightarrow \mathbf{q}' = \mathbf{U} \mathbf{q}$  ( $\mathbf{U} \in SU(3)$ )
- Anti-quark  $\bar{\mathbf{q}} \rightarrow \bar{\mathbf{q}}' = \bar{\mathbf{q}} \mathbf{U}^+$
- Meson octet  $\mathbf{M}_{ij} = \bar{\mathbf{q}}_j \mathbf{q}_i \rightarrow \mathbf{M}' = \mathbf{U} \mathbf{M} \mathbf{U}^+$

$$\begin{pmatrix} \bar{u}u & \bar{d}u & \bar{s}u \\ \bar{u}d & \bar{d}d & \bar{s}d \\ \bar{u}s & \bar{d}s & \bar{s}s \end{pmatrix} = \begin{pmatrix} \frac{\eta}{\sqrt{6}} + \frac{\pi^0}{\sqrt{2}} & \pi^+ & K^+ \\ \pi^- & \frac{\eta}{\sqrt{6}} - \frac{\pi^0}{\sqrt{2}} & K^0 \\ K^- & \bar{K}^0 & -\frac{2\eta}{\sqrt{6}} \end{pmatrix} = P$$

$$S = \begin{pmatrix} \frac{\sigma}{\sqrt{2}} + \frac{a_0}{\sqrt{2}} & a_0^+ & \kappa^+ \\ a_0^- & \frac{\sigma}{\sqrt{2}} - \frac{a_0}{\sqrt{2}} & \kappa^0 \\ \kappa^- & \bar{\kappa}^0 & \xi \end{pmatrix} \quad V = \begin{pmatrix} \frac{\omega}{\sqrt{2}} + \frac{\rho^0}{\sqrt{2}} & \rho^+ & K^{*+} \\ \rho^- & \frac{\omega}{\sqrt{2}} - \frac{\rho^0}{\sqrt{2}} & K^{*0} \\ K^{*-} & \bar{K}^{*0} & \phi \end{pmatrix}$$

# $SU(3)_f$ invariant coupling

## ■ Baryon-Meson coupling

$$\begin{aligned}\mathcal{L}_{BV} &= \sqrt{2}\{g_s \text{tr}(M_v) \text{tr}(\bar{B}B) + g_D \text{tr}(\bar{B}\{M_v, B\}) + g_F \text{tr}(\bar{B}[M_v, B])\} \\ &= \sqrt{2}\{g_s \text{tr}(M_v) \text{tr}(\bar{B}B) + g_1 \text{tr}(\bar{B}M_v B) + g_2 \text{tr}(B\bar{B}M_v)\}\end{aligned}$$

## ■ Assumption

- BM coupling is  $SU(3)$  invariant
- $N$  does not couple with  $\bar{s}s$  vector meson

$$g_{\omega\Lambda} = \frac{5}{6}g_{\omega N} - \frac{1}{2}g_{\rho N}, \quad g_{\phi\Lambda} = \frac{\sqrt{2}}{6}(g_{\omega N} + 3g_{\rho N})$$

## ■ Further simplification: $g_{\rho N} = g_{\omega N}/3$ (quark counting)

$$g_{\omega N} = g_v, \quad g_{\rho N} = g_v/3, \quad g_{\omega\Lambda} = 2g_v/3, \quad g_{\phi\Lambda} = \sqrt{2}g_v/3$$

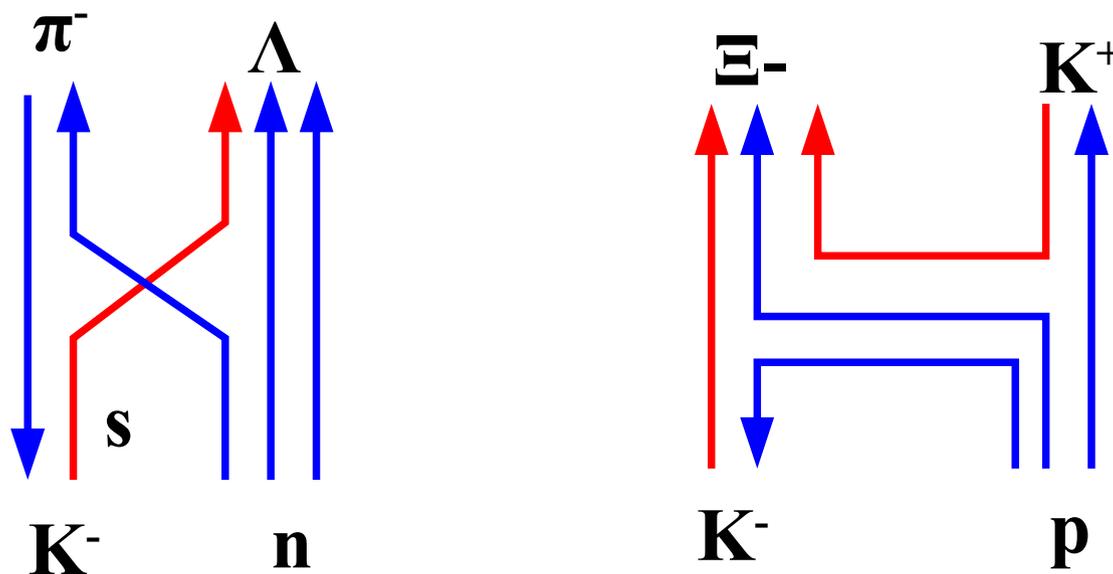
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# *Hypernuclear Formation*

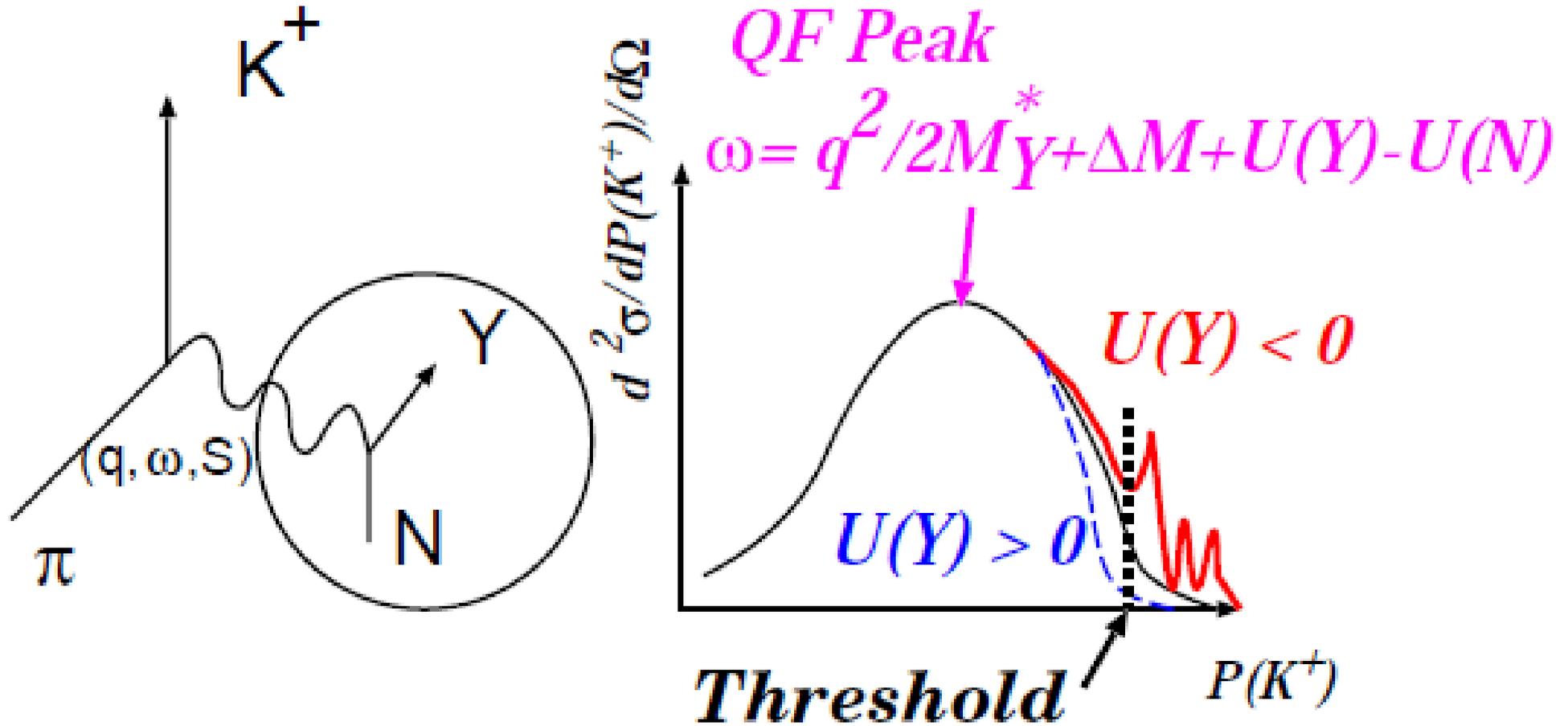
# Hypernuclear formation

- $(K^-, \pi^-)$ ,  $(\pi^-, K^+)$ , and  $(K^-, K^+)$  reactions on nuclei  $\rightarrow$  Hypernuclei

Reaction	Elementary Processes	
	Main Process	Other Processes
$(K^-, \pi^-)$	$K^- n \rightarrow \pi^- \Lambda$ ,	$K^- n \rightarrow \pi^- \Sigma^0, K^- p \rightarrow \pi^- \Sigma^+$
$(K^-, \pi^+)$	$K^- p \rightarrow \pi^+ \Sigma^-$ ,	$K^- pp \rightarrow \pi^+ \Lambda n$ (n-rich hypernuclear formation)
$(\pi^+, K^+)$	$\pi^+ n \rightarrow K^+ \Lambda$ ,	$\pi^+ n \rightarrow K^+ \Sigma^0, \pi^+ p \rightarrow K^+ \Sigma^+$
$(\pi^-, K^+)$	$\pi^- p \rightarrow K^+ \Sigma^-$ ,	$\pi^- pp \rightarrow K^+ \Lambda n$ (n-rich hypernuclear formation)
$(K^-, K^+)$	$K^- p \rightarrow K^+ \Xi^-$ ,	$K^- pp \rightarrow K^+ \Lambda \Lambda$

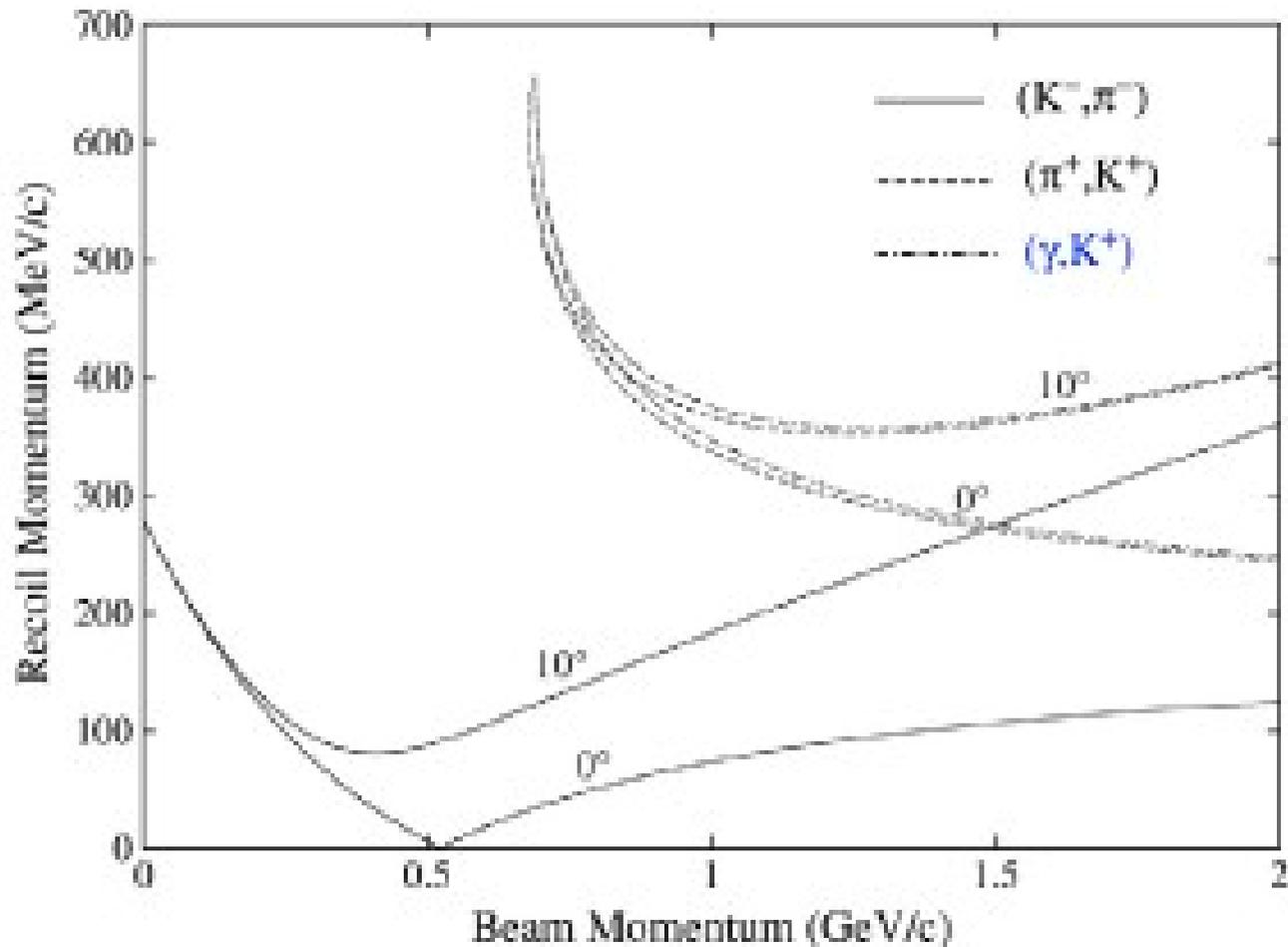


# Hypernuclear formation



# Hypernuclear formation

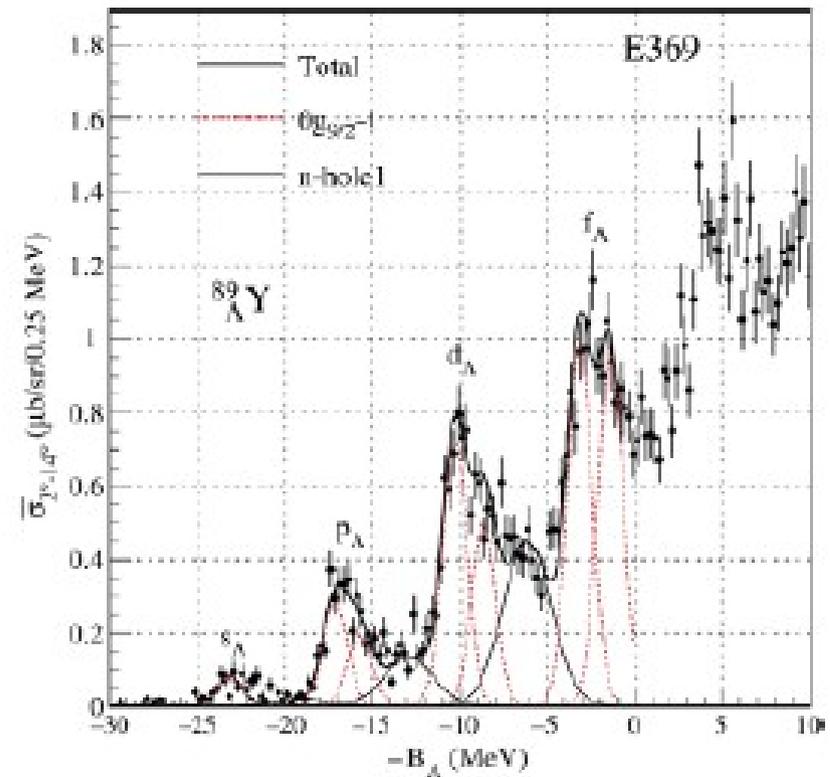
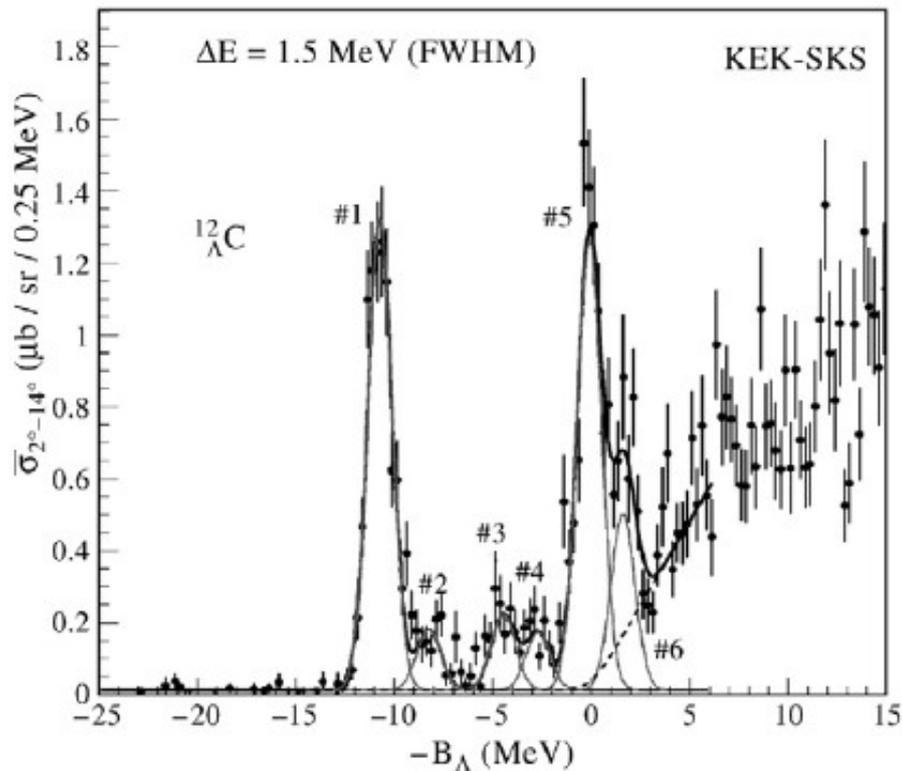
- $(K^-, \pi^-)$ :  $Q > 0$ , Small momentum transfer  $\rightarrow$  substitutional reaction
- $(\pi^-, K^+)$ :  $Q < 0$ , Momentum transfer  $\sim 300 \text{ MeV}/c \sim k_F$



# $\Lambda$ hypernuclear formation

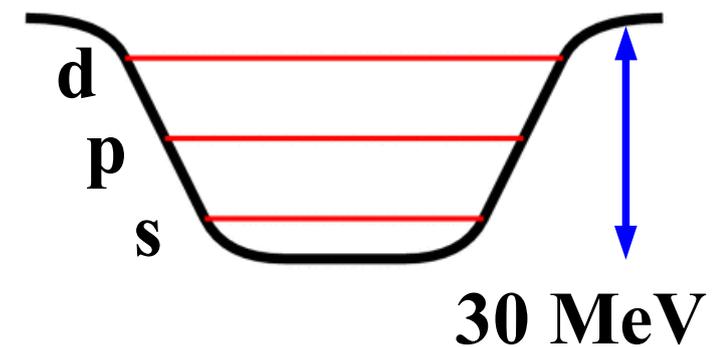
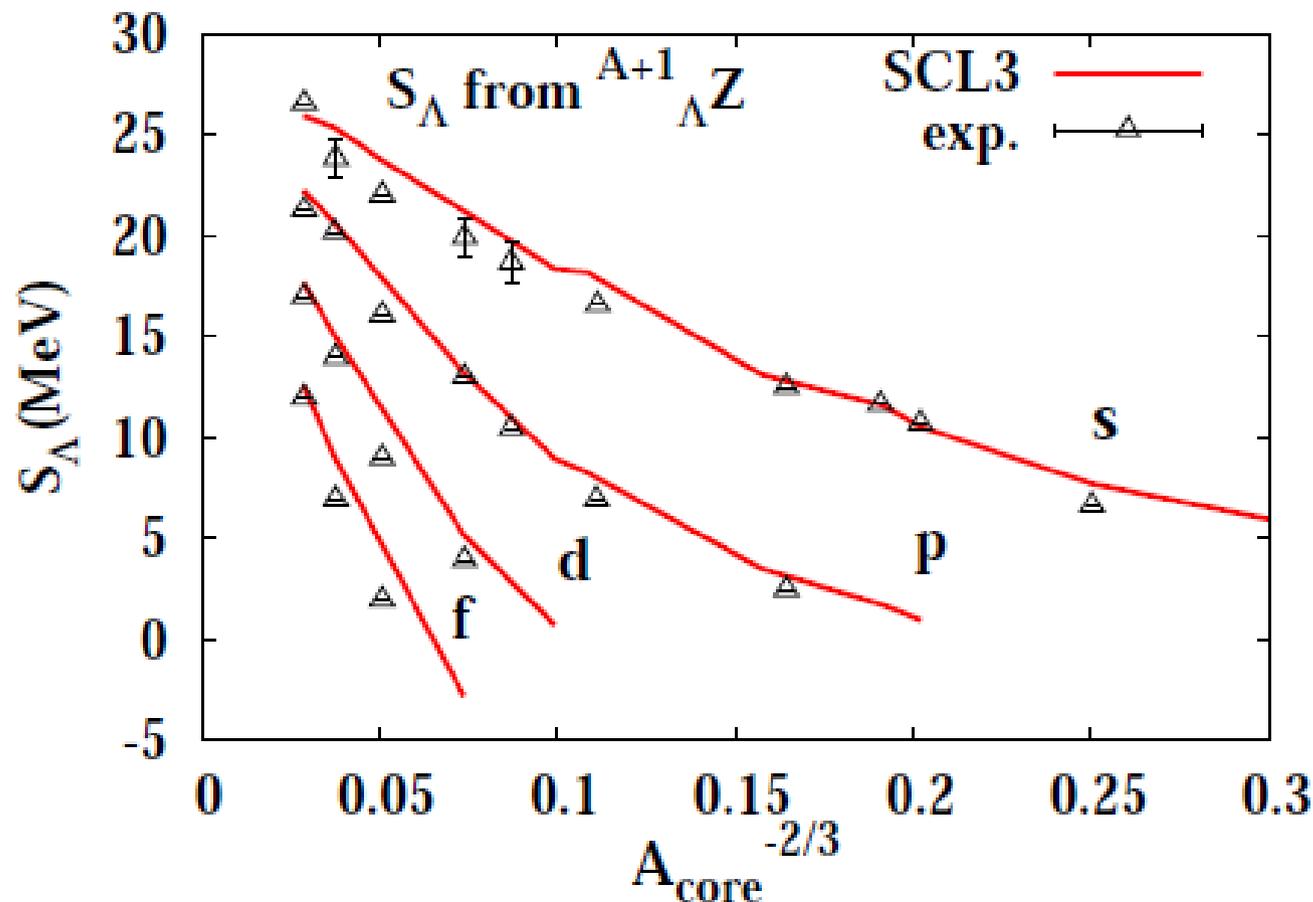
## ■ $(\pi^+, K^+)$ reactions on nuclei

- $q \sim k_F \rightarrow$  various s.p. states of  $\Lambda$  are populated



# Single particles states of $\Lambda$ in nuclei

- Single particle potential depth of  $\Lambda$  is around -30 MeV
  - s, p, d, f, ... states are clearly seen
  - $A_{\text{core}}^{-2/3} \propto R^{-2} \propto \text{K.E. of } \Lambda$



# $\Sigma$ production in nuclei

- Only one bound state  ${}^4_{\Sigma}\text{He}$  (Too light !)  
→ Continuum (Quasi-Free) Spectroscopy is necessary
- Cont. Spec. Theory = Distorted Wave Impulse Approx. (DWIA)

$$\frac{d^2 \sigma}{dE_K d\Omega_K} = \beta \left( \frac{d\sigma}{d\Omega} \right)_{N\pi \rightarrow KY}^{Elem.} S(E, q)$$

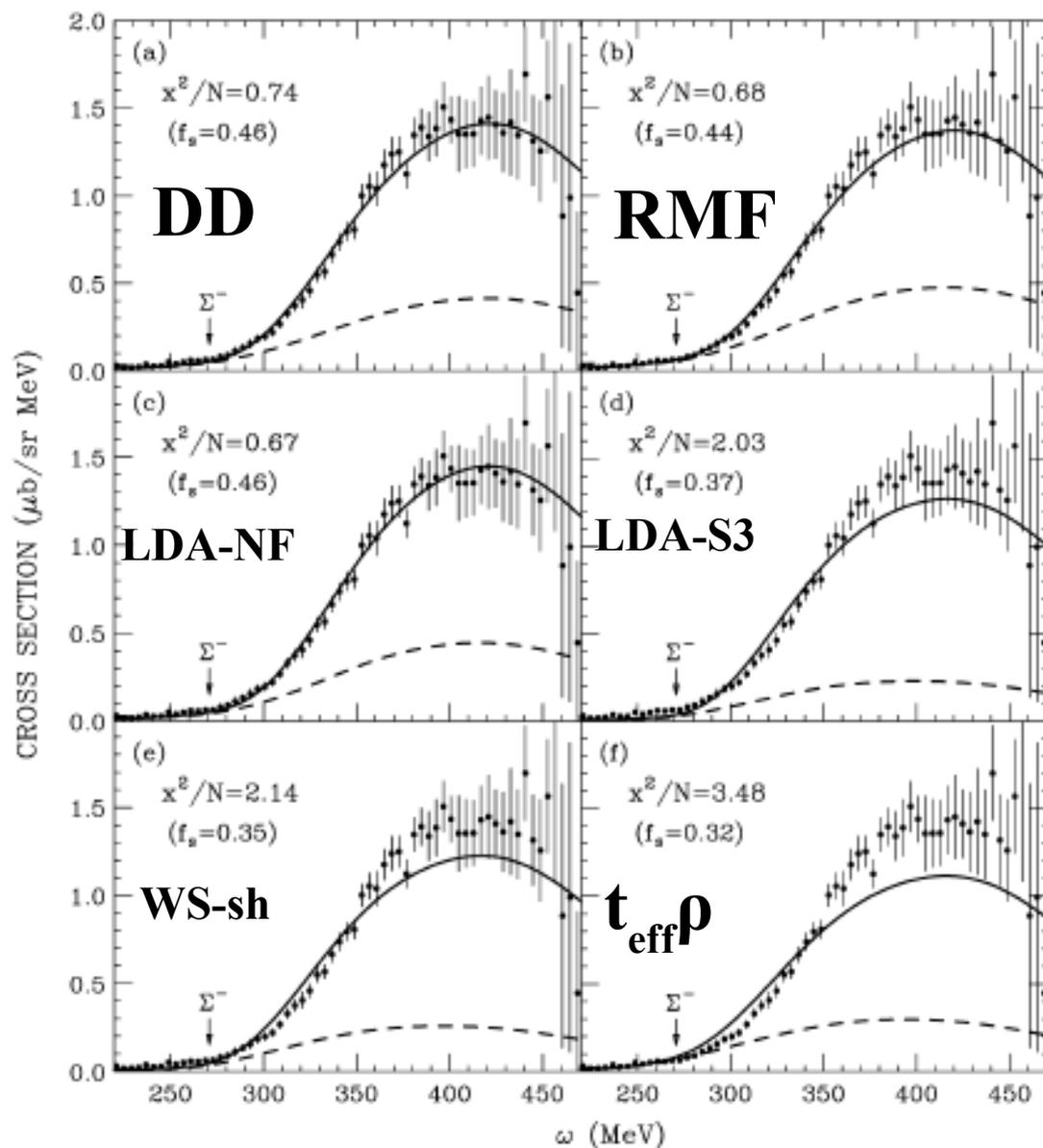
Kinematical Factor

Elem. Cross Sec.

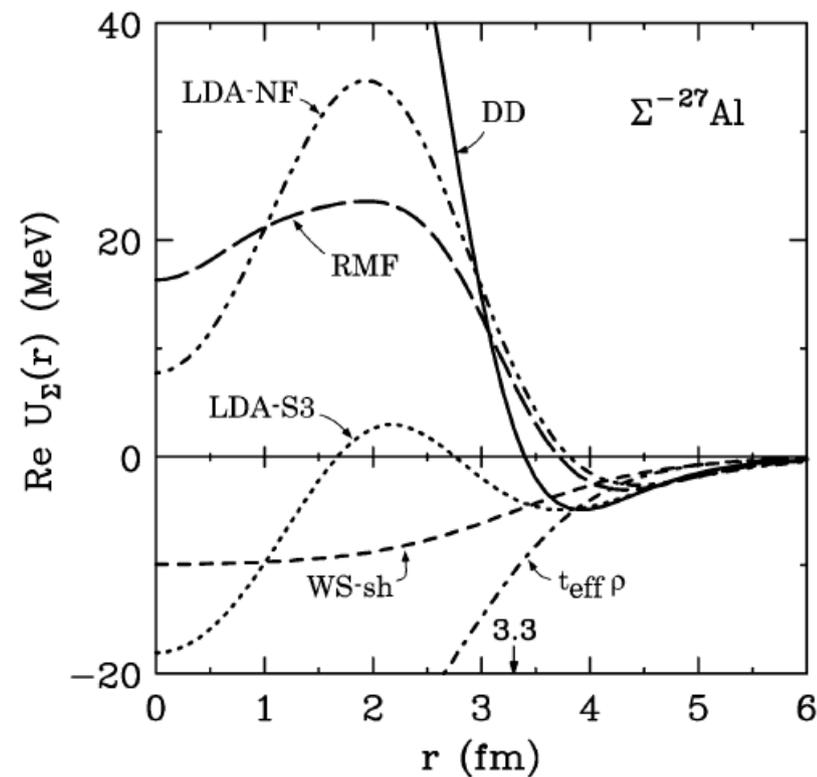
Strength Func.

- Large ( $\omega, q$ ) range → Important to respect **On-Shell Kinematics**
- Another way:  $\Sigma^-$  atomic shift
  - Atomic shift of  $\Sigma^-$  with O, Mg, Al, S, Si, W, Pb core are measured
- $\Sigma$  potential in nuclei
  - Isoscalar part: 15-35 MeV repulsion
  - Isovector part: 20-30 % of SU(3) value

# $\Sigma$ production in nuclei

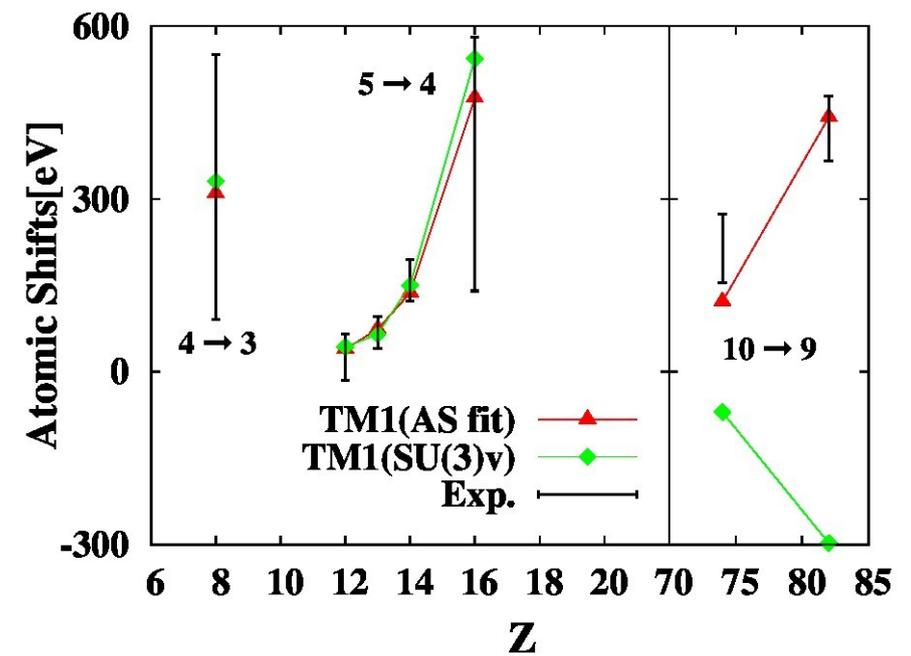
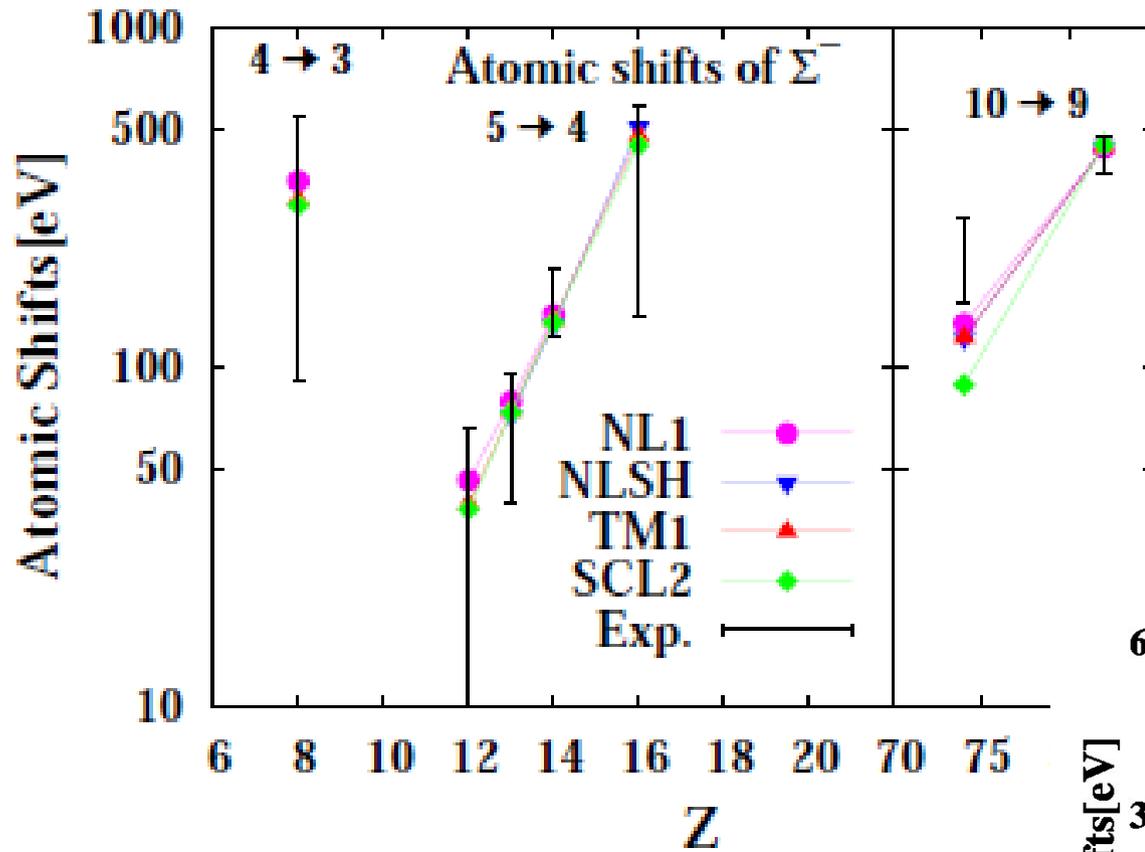


$^{28}\text{Si}(\pi^-, \text{K}^+)$



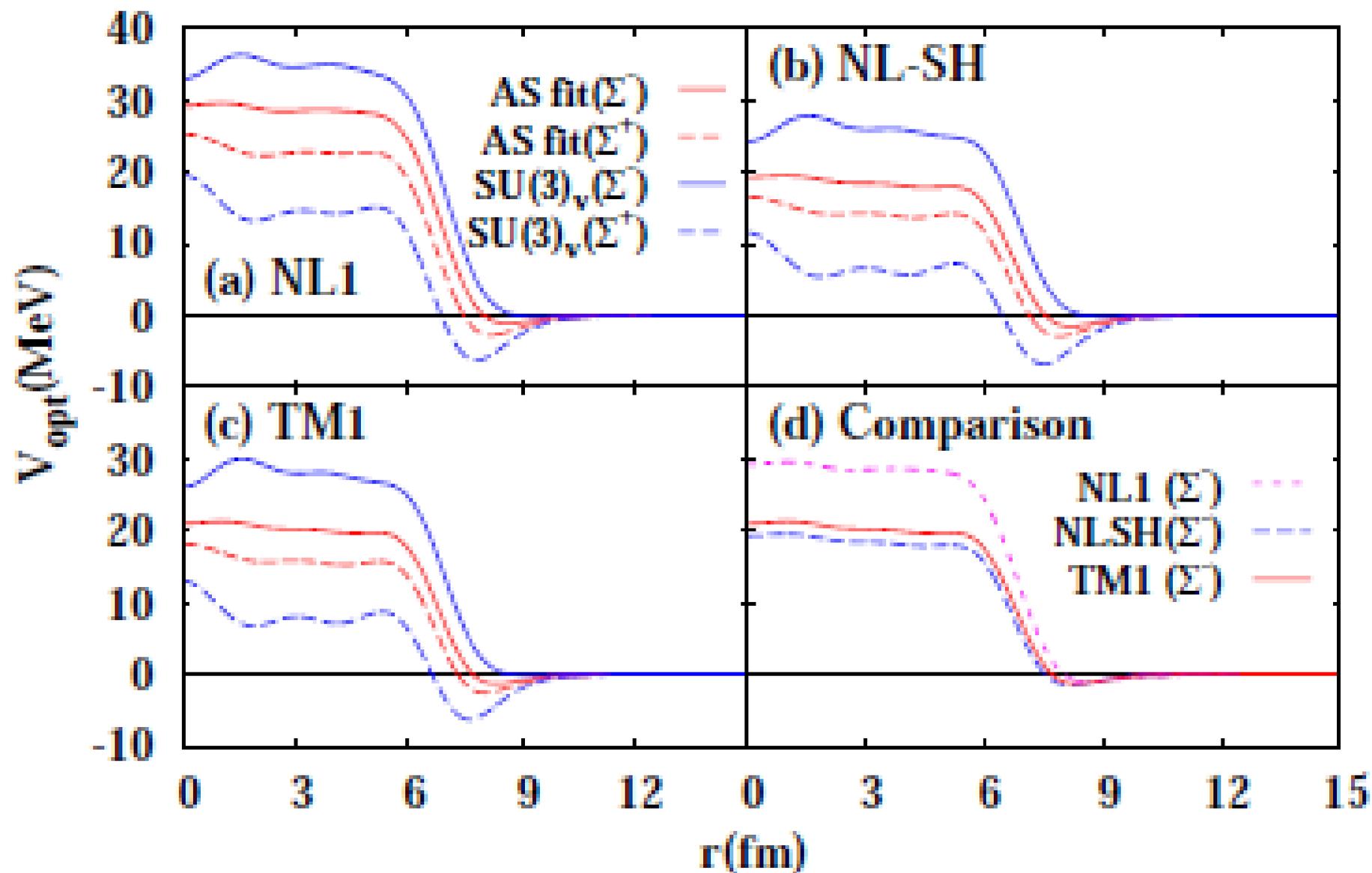
Harada, Hirabayashi ('05)

# $\Sigma^-$ atomic shift



# $\Sigma$ potential in nuclei

## Optical $\Sigma$ potentials in Pb



# $\Xi$ hypernuclear formation

- Missing mass spectroscopy

BNL E885  $^{12}\text{C}(\text{K}^-, \text{K}^+)$

*Fukuda et al. PRC58('98),1306;*

*Khaustov et al. PRC61('00), 054603.*

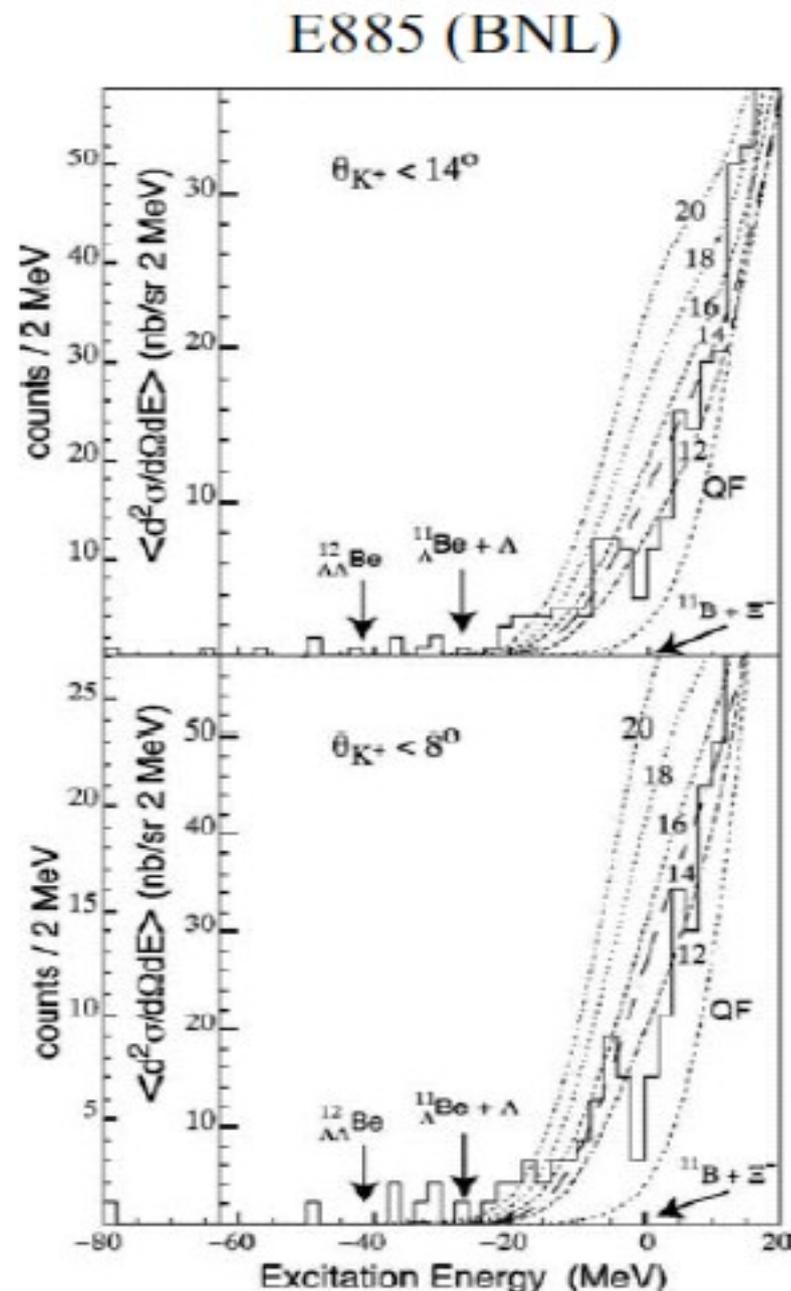
- No clear bound states found

- Twin hypernuclear formation

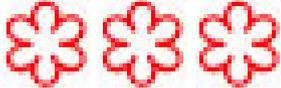
*Aoki et al. PLB355('95),45.*

- Potential depth

$$U_{\Xi} \sim -14 \text{ MeV}$$



# “Stars” of Hyperon Potentials (A la Michelin)

- $U_{\Lambda}(\rho_0) \sim -30 \text{ MeV}$  
  - *Bound State Spectroscopy + Continuum Spectroscopy*
- $U_{\Sigma}(\rho_0) > +15 \text{ MeV}$  
  - Continuum (Quasi-Free) spectroscopy
  - Atomic shift data (attractive at surface) should be respected.
- $U_{\Xi}(\rho_0) \sim -14 \text{ MeV}$  
  - No confirmed bound state, No atomic data, High mom. transf., ....  $\rightarrow$  Small Potential Deps.
  - Continuum low-res. spectrum shape  $\rightarrow -14 \text{ MeV}$



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# *Hyperons in Neutron Stars*

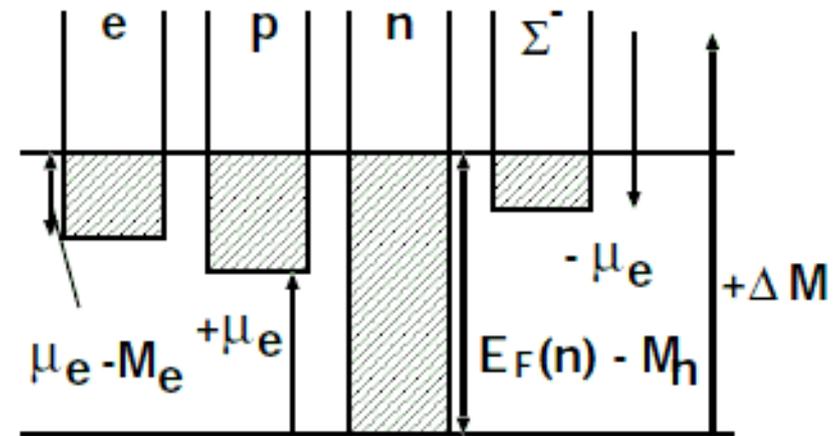
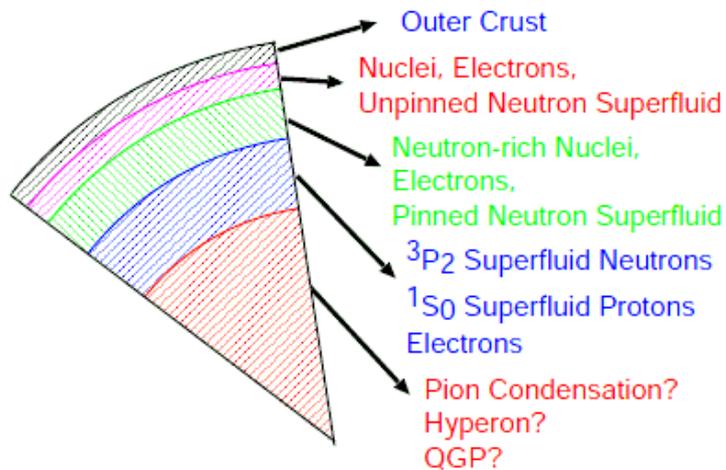
# Hyperons in Dense Matter

## ■ What appears at high density ?

- Nucleon superfluid ( $^3S_1$ ,  $^3P_2$ ), Pion condensation, Kaon condensation, Baryon Rich QGP, Color SuperConductor (CSC), Quarkyonic Matter, ....

### ● Hyperons

Tsuruta, Cameron (66); Langer, Rosen (70); Pandharipande (71); Itoh(75); Glendenning; Weber, Weigel; Sugahara, Toki; Schaffner, Mishustin; Balberg, Gal; Baldo et al.; Vidana et al.; Nishizaki, Yamamoto, Takatsuka; Kohno, Fujiwara et al.; Sahu, Ohnishi; Ishizuka, Ohnishi, Sumiyoshi, Yamada; ...

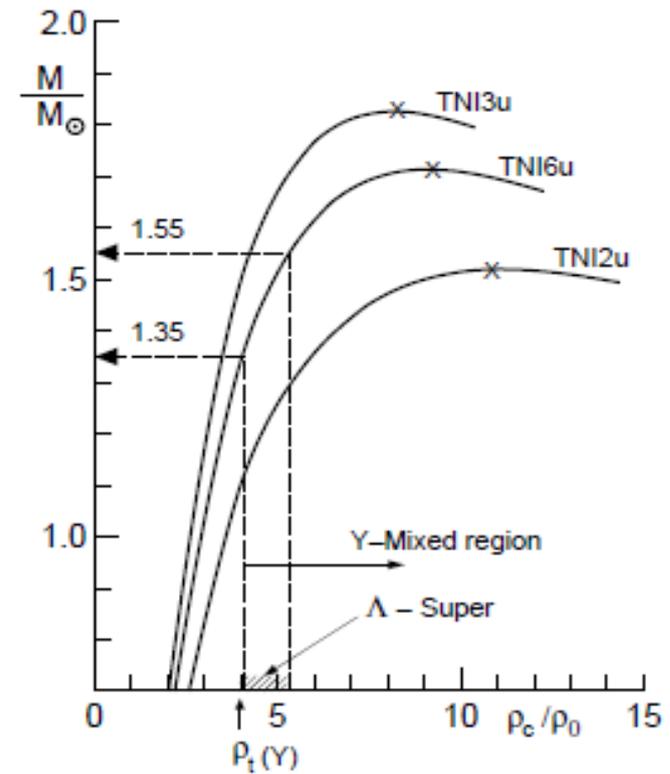
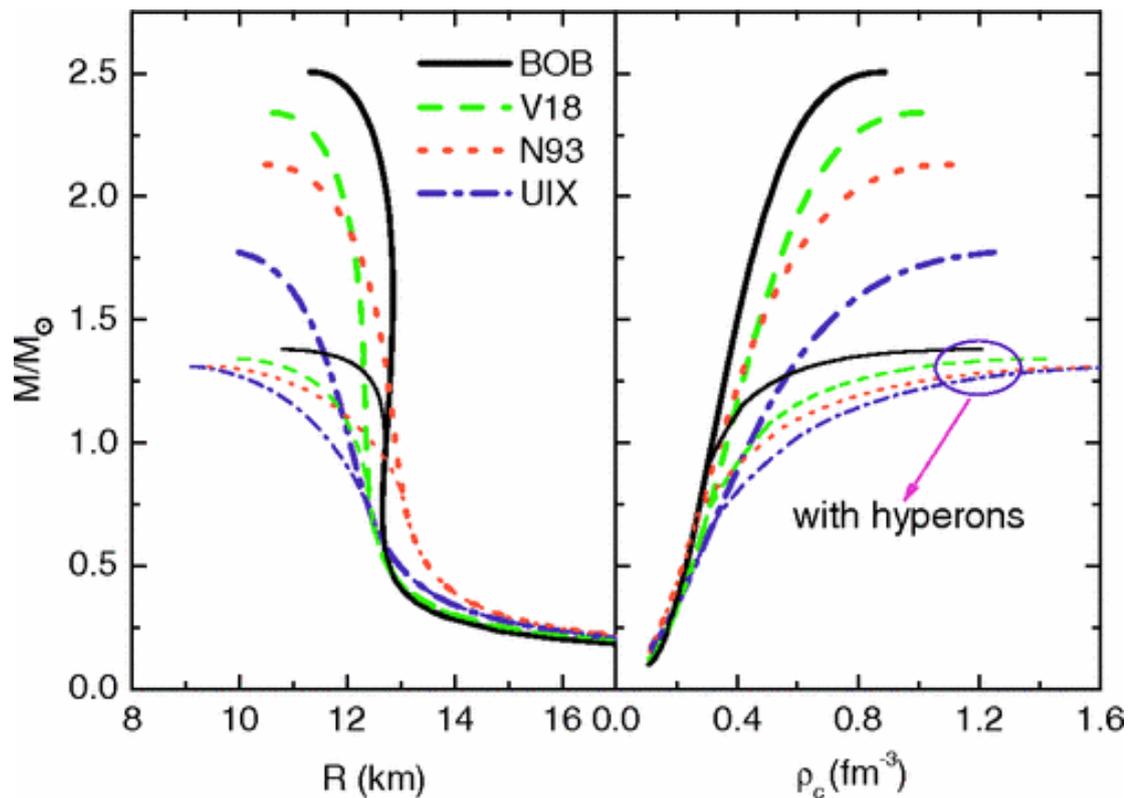


***Nobody says “Hyperons cannot appear in neutron star core” !***

$Y$  appears when  $\mu_B = E_F(n) + U(n) \geq M(Y) + U(Y) + Q_Y \mu_e$

# Bruckner-Hartree-Fock theory with Hyperons

- Microscopic G-matrix calculation with realistic NN, YN potential and microscopic (or phen.) 3N force (or 3B force).
  - Interaction dep. (V18, N93, ...) is large → Need finite nuclear info.  
*E.Hiyama, T.Motoba, Y.Yamamoto, M.Kamimura / M.Tamura et al.*
  - NS collapses with hyperons w/o 3BF.



*Z.H.Li, H.-J.Schulze, PRC78('08),028801.*

*S. Nishizaki, T. Takatsuka,  
Y. Yamamoto, PTP108('02)703.*

# RMF with Hyperons (Single $\Lambda$ hypernuclei)

## ■ RMF for $\Lambda$ hypernuclei

$x \sim 1/3$ : R. Brockmann, W. Weise, PLB69('77)167; J. Boguta and S. Bohrman, PLB102('81)93.

$x \sim 2/3$ : N. K. Glendenning, PRC23('81)2757, PLB114('82)392;

Tensor: Y. Sugahara, H. Toki, PTP92('94)803; H. Shen, F. Yang, H. Toki, PTP115('06)325;  
J. Mares, B. K. Jennings, PRC49('94)2472.

$\rho$ -dep. coupling: H. Lenske, Lect. Notes Phys. 641('04)147; C. M. Keil, F. Hofmann, H. Lenske, PRC 61('00)064309.

SU(3) or SU(6) ( $\zeta, \varphi$ ): J. Schaffner, C. B. Dover, A. Gal, C. Greiner, H. Stoecker, PRL71('93)1328;  
Schaffner et al., Ann.Phys.235('94)35; J. Schaffner, I. N. Mishustin, PRC 53('96)1416.

Chiral SU(3) RMF: K. Tsubakihara, H. Maekawa, H. Matsumiya, AO, PRC81('10)065206.

- Sep. E. of  $\Lambda$  is well fitted by  $U_{\Lambda} \sim -30 \text{ MeV} \sim 2/3 U_N$

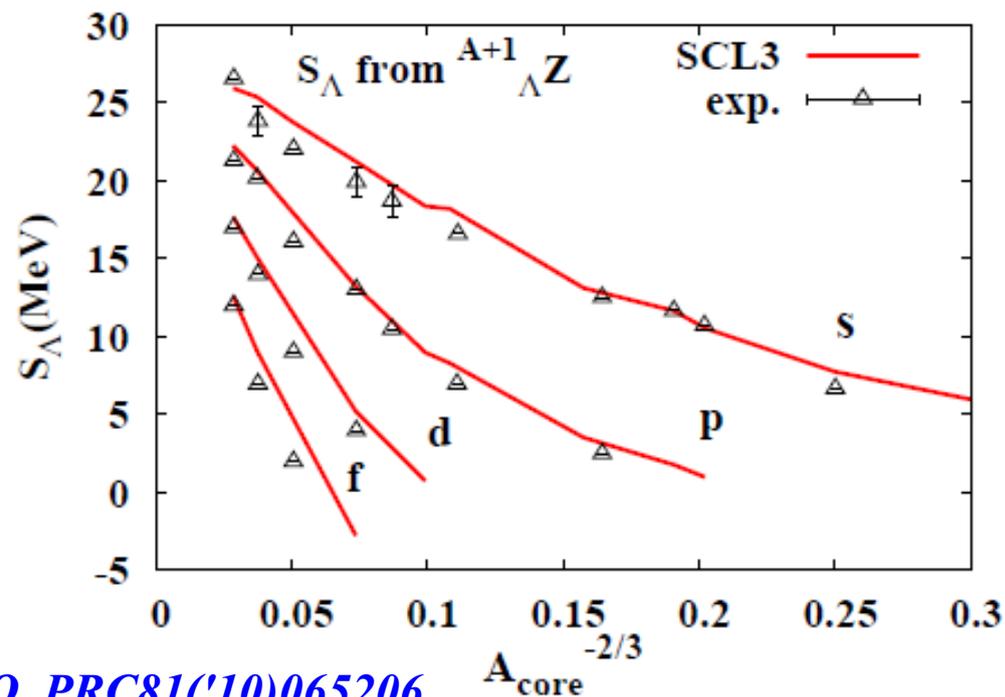
- Coupling with mesons

$$x_M = g_{M\Lambda} / g_{MN}$$

quark counting:  $x_{\sigma} \sim 2/3$

$\pi$  exchanges:  $x_{\sigma} \sim 1/3$

→ Which is true ?



K. Tsubakihara, H. Maekawa, H. Matsumiya, AO, PRC81('10)065206.

# Hyperon Composition in Dense Matter

■ Hyperon start to emerge at  $(2-3)\rho_0$  in Neutron Star Matter !

■ Hyperon composition in NS is sensitive to Hyperon potential.

- $U_\Lambda \sim -30$  MeV: Well-known

- $U_\Sigma \sim -(12-15)$  MeV

( $K^-, K^+$ ) reaction, twin hypernuclei

*P. Khaustov et al. (E885), PRC61('00)054603;*

*S. Aoki et al., PLB355('95)45.*

- $U_\Sigma \sim -30$  MeV (Old conjecture)

→  $\Sigma$ - appears prior to  $\Lambda$

- $U_\Sigma > 0$  (repulsive) → No  $\Sigma$  in NS

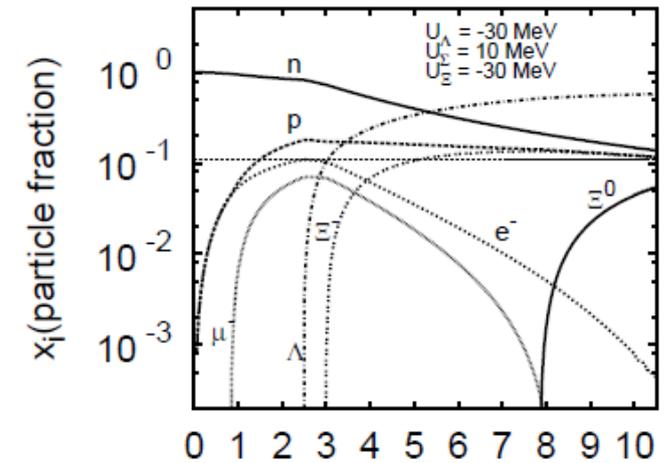
$\Sigma$  atom (phen. fit), QF prod.

*S. Balberg, A. Gal, NPA625('97)435;*

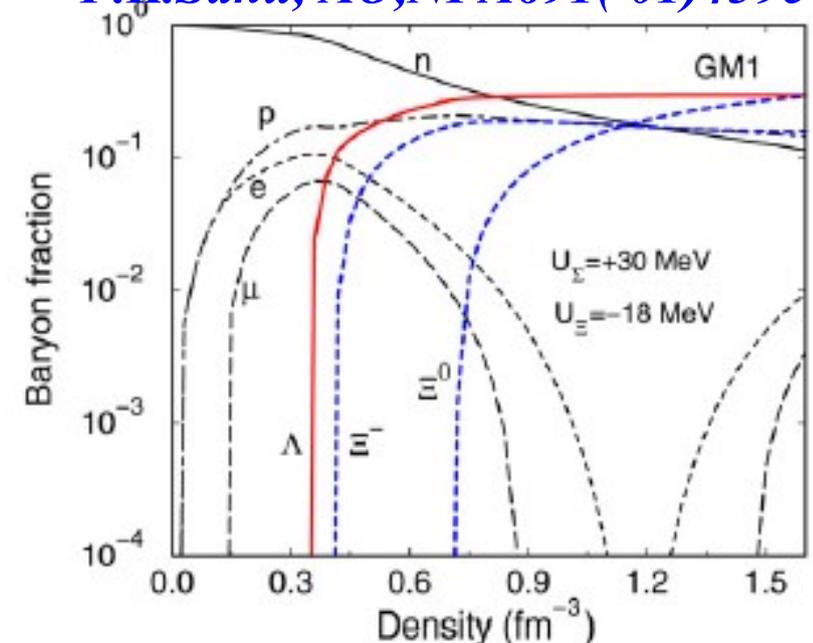
*H. Noumi et al., PRL89('02)072301;*

*T. Harada, Y. Hirabayashi, NPA759('05)143;*

*M. Kohno et al. PRC74('06)064613.*



*P.K.Sahu, AO, NPA691('01)439c*



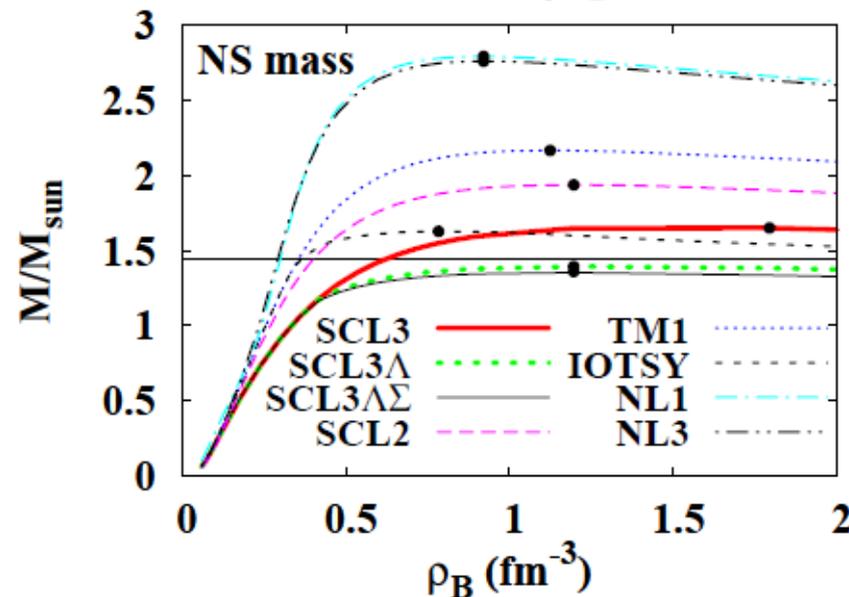
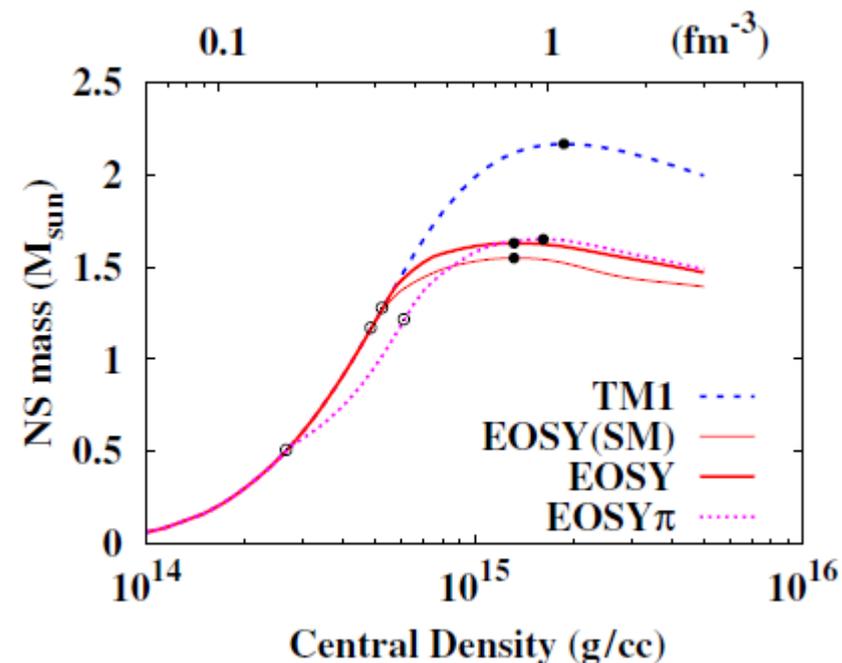
*J. Schaffner-Bielich, NPA804('08)309.*

# Neutron Star Mass

- Large fraction of hyperons softens EOS at  $\rho_B > (0.3-0.4) \text{ fm}^{-3}$ 
  - NS star max. mass red.  $\sim 1 M_{\text{sun}}$ .
  - RMF generally predicts stiff EOS at high density. (Scalar attraction saturation, or Z-graph in NR view.)
  - Some of RMF with Y do not support  $1.44 M_{\text{sun}}$ .
- Additional Repulsion at high  $\rho$  ?
  - Vector mass mod.
    - stronger repulsion at high  $\rho$ .
- Another term such as  $\text{NN}\omega\sigma$ .

*C. Ishizuka, AO, K. Tsubakihara, K. Sumiyoshi, S. Yamada, JPG35('08)085201.*

*K. Tsubakihara, H. Maekawa, H. Matsumiya, AO, PRC81('10)065206.*



# RMF with hyperons (1)

Ishizuka, AO, Tsubakihara, Sumiyoshi, Yamada, *J. Phys. G35(08),085201*

## ■ Lagrangian including hyperons

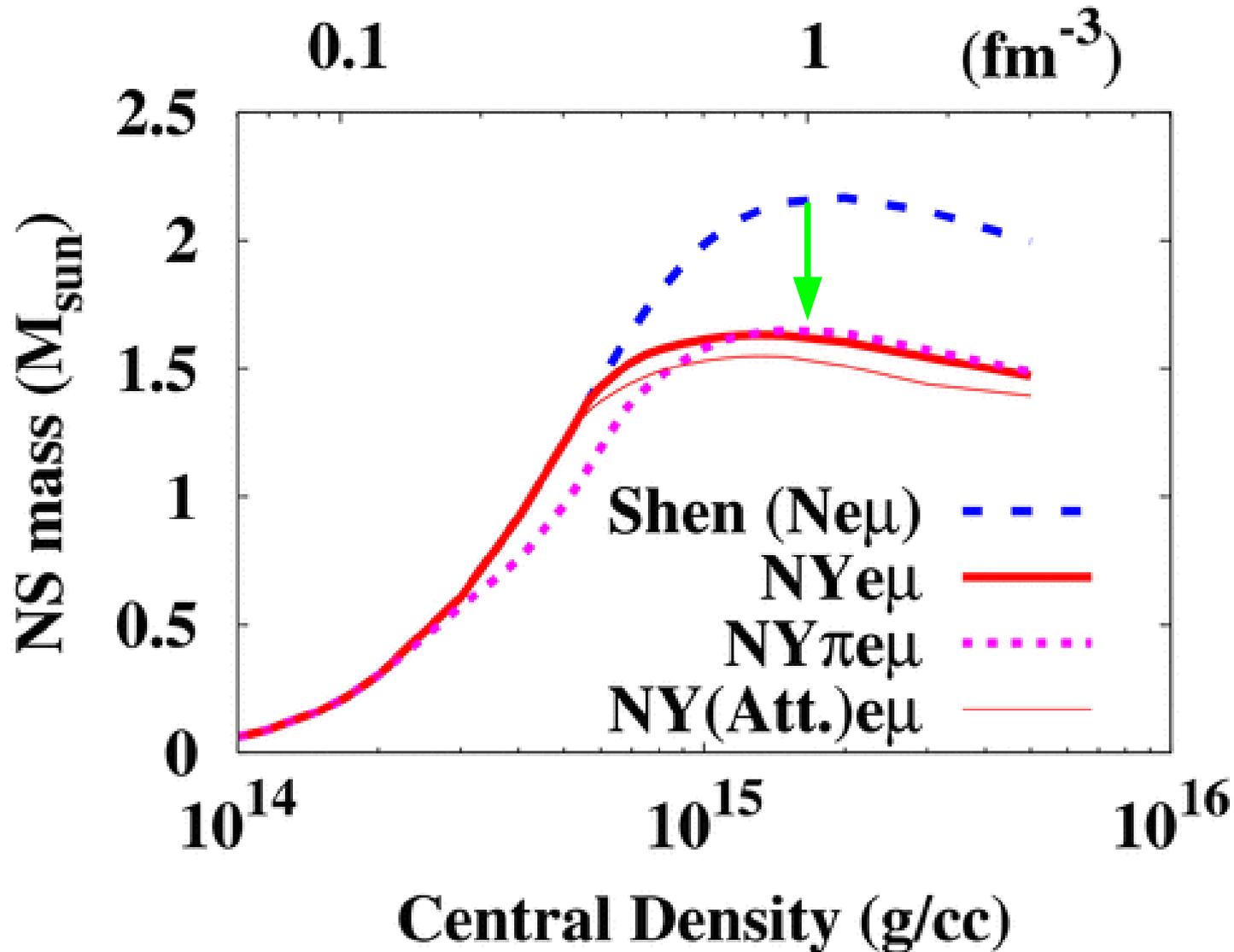
$$\begin{aligned} \mathcal{L} = & \sum_B \bar{\Psi}_B (i\cancel{\partial} - M_B) \Psi_B + \frac{1}{2} \partial^\mu \sigma \partial_\mu \sigma - U_\sigma(\sigma) \\ & - \frac{1}{4} \omega^{\mu\nu} \omega_{\mu\nu} + \frac{1}{2} m_\omega^2 \omega^\mu \omega_\mu - \frac{1}{4} \vec{R}^{\mu\nu} \cdot \vec{R}_{\mu\nu} + \frac{1}{2} m_\rho^2 \vec{R}^\mu \cdot \vec{R}_\mu \\ & - \sum_B \bar{\Psi}_B \left( g_{\sigma B} \sigma + g_{\omega B} \cancel{\partial} + g_{\rho B} \vec{R} \cdot \vec{t}_B \right) \Psi_B + \frac{1}{4} c_\omega (\omega^\mu \omega_\mu)^2 + \mathcal{L}^{YY}, \\ U_\sigma(\sigma) = & \frac{1}{2} m_\sigma^2 \sigma^2 + \frac{g_3}{3} \sigma^3 + \frac{g_4}{4} \sigma^4, \\ \mathcal{L}^{YY} = & \frac{1}{2} \partial_\nu \zeta \partial^\nu \zeta - \frac{1}{2} m_\zeta^2 \zeta^2 - \frac{1}{4} \phi_{\mu\nu} \phi^{\mu\nu} + \frac{1}{2} m_\phi^2 \phi_\mu \phi^\mu \\ & - \sum_B \bar{\Psi}_B (g_{\zeta B} \zeta + g_{\phi B} \gamma^\mu \phi_\mu) \Psi_B, \end{aligned}$$

- Nucleon part = TM1
- Vector coupling = Quark counting (SU(6))
- Scalar couplings are chosen to reproduce hyperon potentials in NM

$$U_\Sigma^{(N)}(\rho_0) \simeq +30 \text{ MeV}, \quad U_\Xi^{(N)}(\rho_0) \simeq -15 \text{ MeV}.$$

# ハイペロンを含むRMFでの中性子星

Ishizuka, AO, Tsubakihara, Sumiyoshi, Yamada, *J. Phys. G*35(08),085201



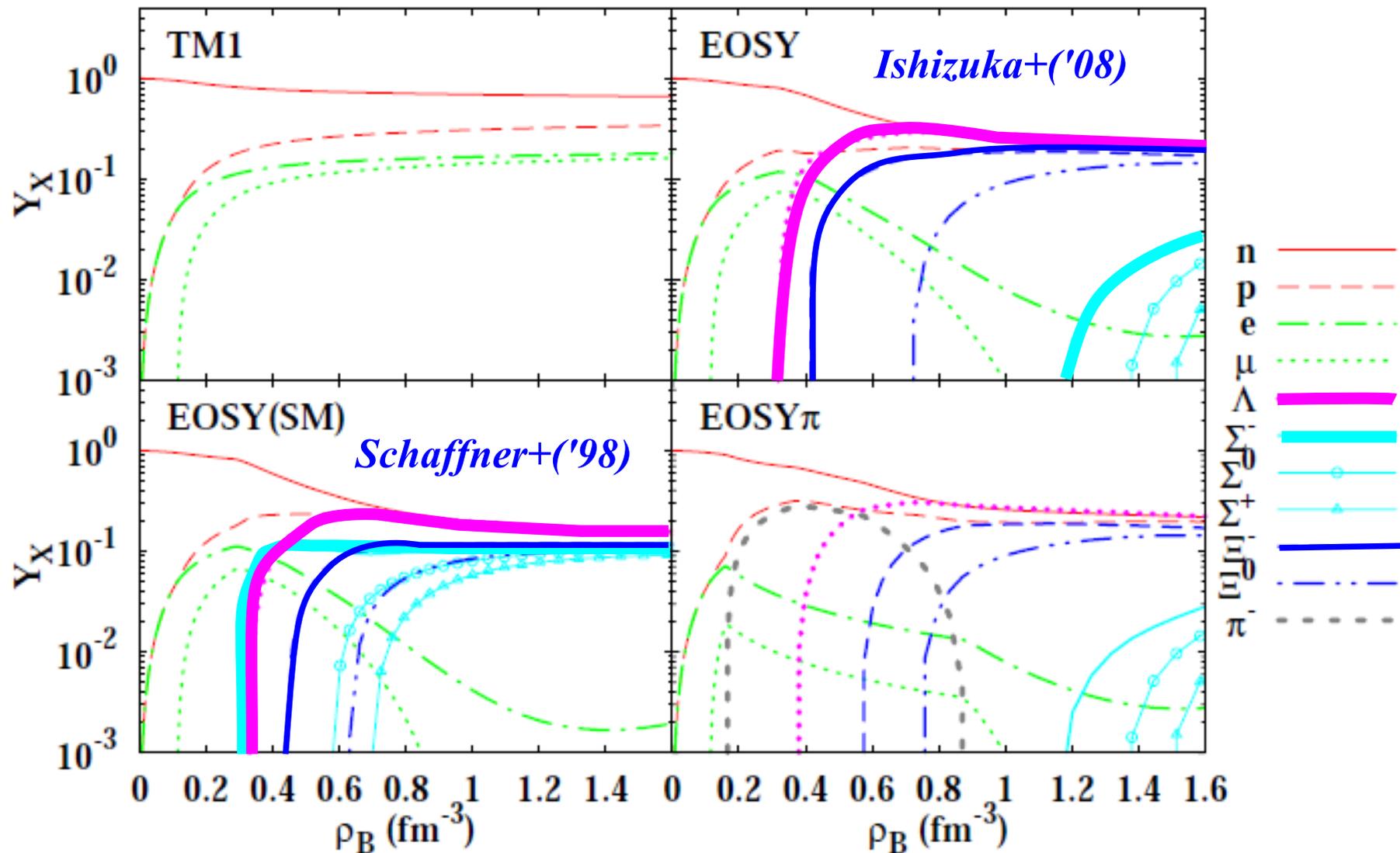
*c.f. H. Shen+('09)  $\rightarrow$  n, p,  $\Lambda$  EOS*

*Ohnishi @ Osaka U., 2014 25*

# ハイペロンを含むRMFでの中性子星

Ishizuka, AO, Tsubakihara, Sumiyoshi, Yamada, *J. Phys. G*35(08),085201

## Neutron Star Matter



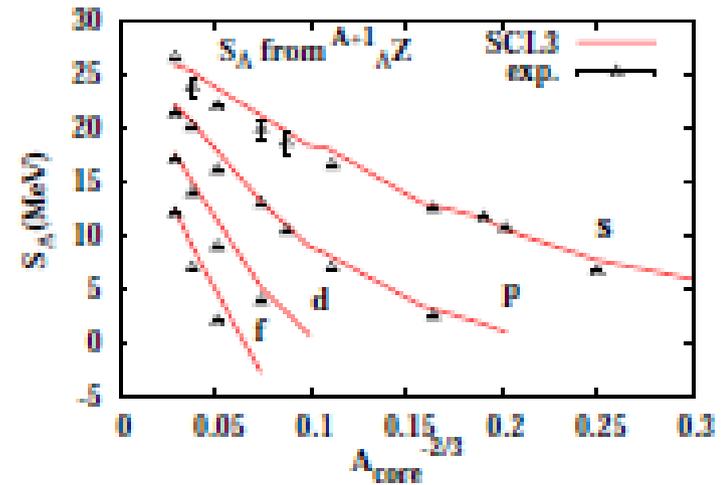
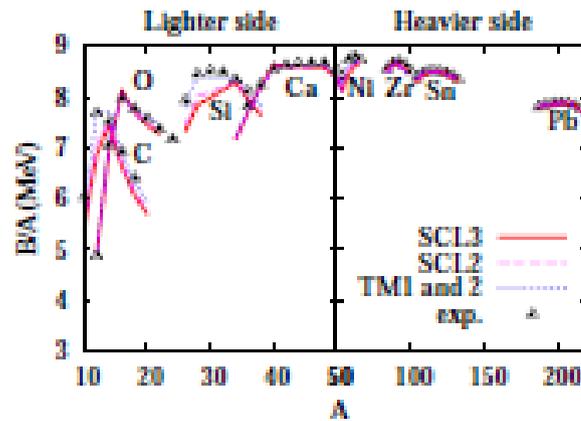
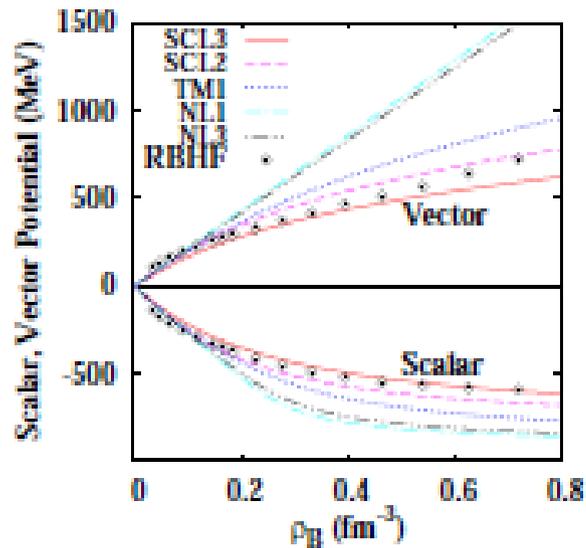
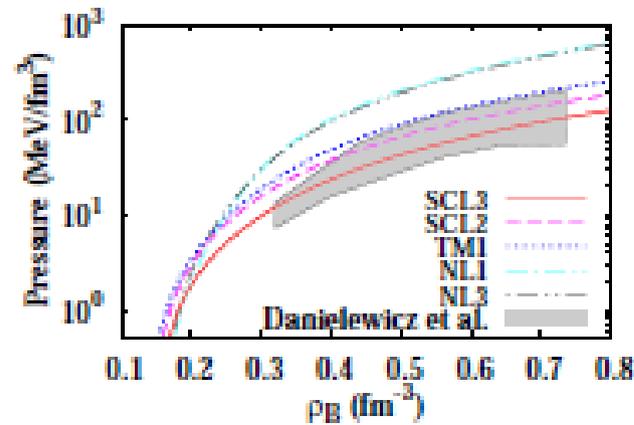
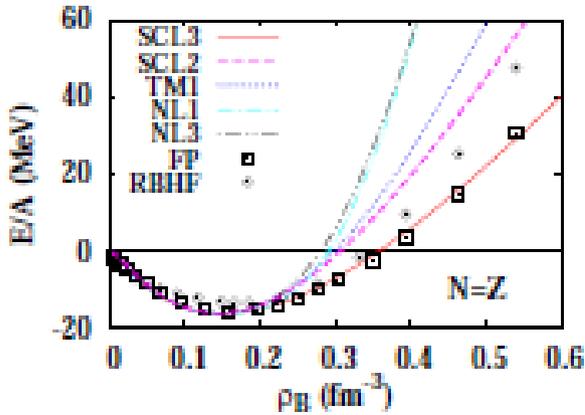
# *RMF with Hyperons (SCL3)*

*K. Tsubakihara, H. Maekawa, H. Matsumiya, A. Ohnishi, PRC81 (2010), 065206.*

- **Fit available data as much as possible**
  - **BE of normal nuclei, EOS (incompressibility), Pressure suggested from HIC, and Vector potential in RBHF.**
  - **Single  $\Lambda$  single particle energies, Double  $\Lambda$  bond energy,  $\Sigma$ - atomic shift,  $\Xi$  potential depth**
- **Non-linear terms**
  - **Logrithmic chiral symmetric  $\sigma$  potential suggested by the strong coupling limit of lattice QCD**
  - **$\omega^4$  term as in TM1**
  - **$U(1)_A$  breaking term (Kobayashi-Masukawa-'t Hooft)**
- **Published on June 22, 2010.**

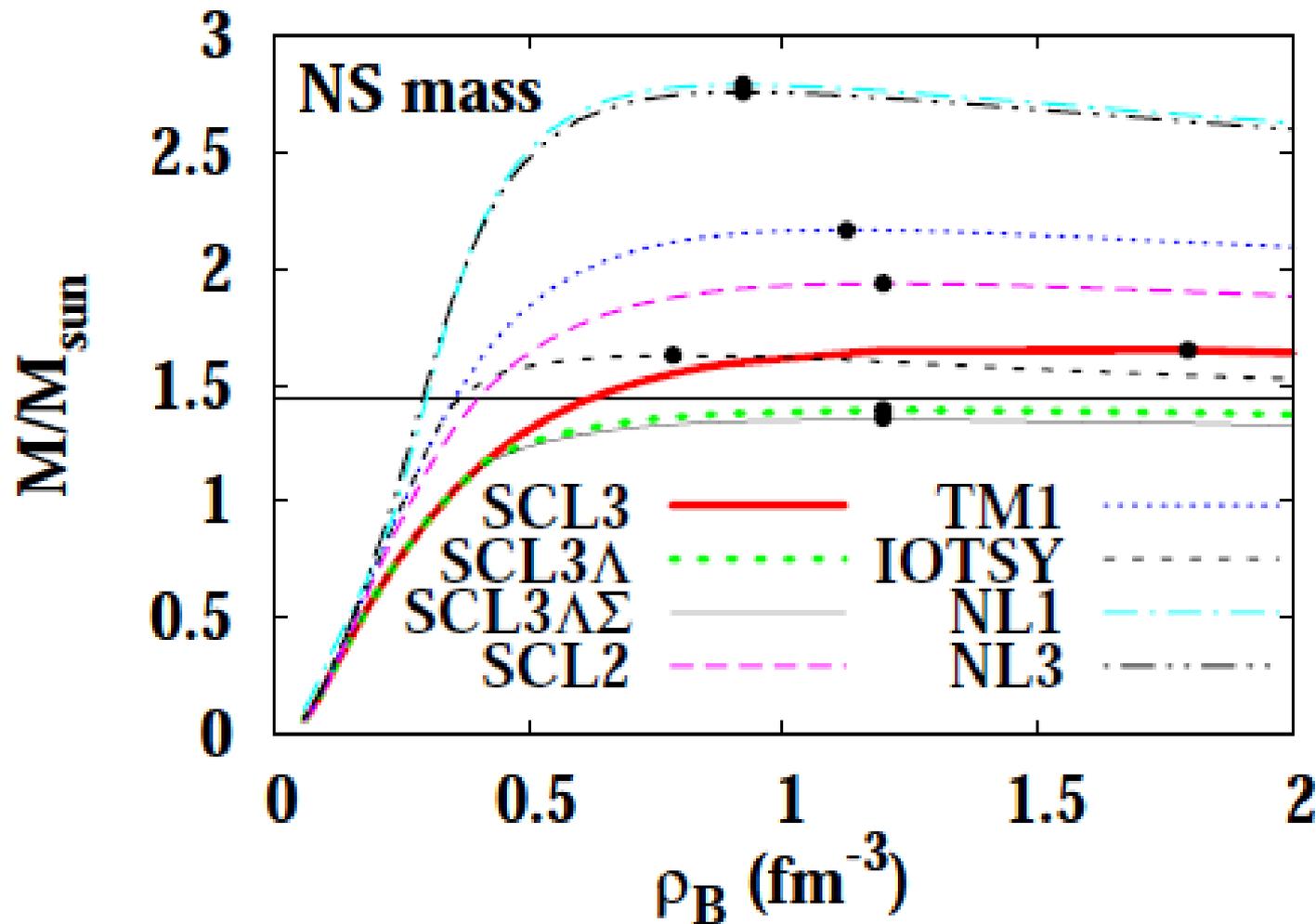
# EOS related data

K. Tsubakihara, H. Maekawa, H. Matsumiya, A. Ohnishi, *PRC81* (2010), 065206.



# Neutron Star Mass

*K. Tsubakihara, H. Maekawa, H. Matsumiya, A. Ohnishi, PRC81 (2010), 065206.*



**$M_{\text{max}} < 1.44 M_{\odot}$  with hyperons**

# *RMF as a phenomenological MODEL !*

- **Baryon one-loop approximation (Hartree approximation) makes RMF a phenomenological model.**

→ We need **DATA** and **AB INITIO** results.

- **Saturation point ( $\rho_0$  and  $E/A(\rho_0)$ ) from mass formula**
- **Nuclear binding energies**
- **$U_v$  and  $U_s$  from DBHF results**
- **$P(\rho_B)$  from heavy-ion data**
- **$\Lambda$  separation energy from single  $\Lambda$  hypernuclear data**
- **$\Lambda\Lambda$  bond energy from double  $\Lambda$  hypernuclear data**
- **$\Sigma$  atomic shift**
- **$\Sigma$  and  $\Xi$  potential depth from quasi-free production data**
- *Pure neutron matter EOS from ab initio calculations (not used here)*
- **Neutron Star Max. Mass  $\sim 1.40 M_\odot$ , a little smaller  $1.44 M_\odot$ .**

*The Judgement Day, Oct. 28, 2010.  $1.97 M_\odot$*

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# *Massive Neutron Star Puzzle*

# Massive Neutron Star Puzzle

## ■ 重い中性子星 (2 倍の太陽質量) の観測

*Demorest et al., Nature 467 (2010) 1081 (Oct.28, 2010).*

PSR J1614-2230 (NS-WD binary),  $1.97 \pm 0.04 M_{\odot}$

- 一般相対性理論 (Shapiro delay) に基づく質量決定

- 幸運な公転面の向き + 美しい観測結果

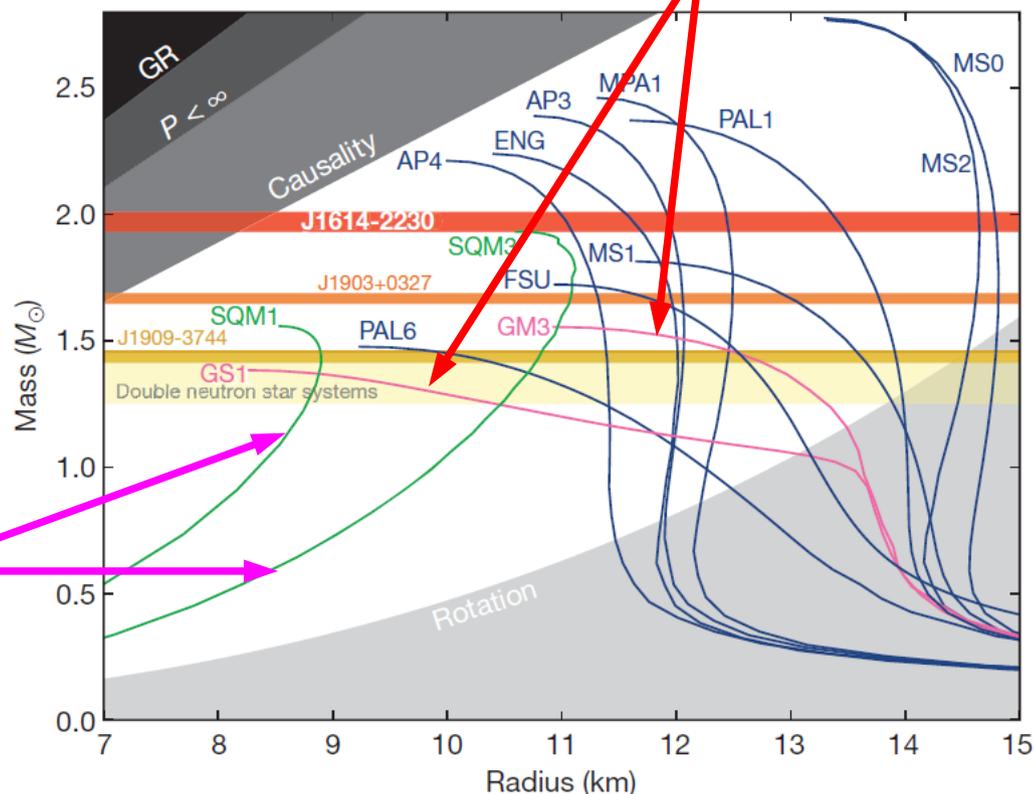
Strange Hadron  
を含む EOS

## ■ 高密度状態方程式 (EOS) に強い制限

- Strange Hadron (ハイペロン・K 中間子) 凝縮を含む EOS は棄却 (?)

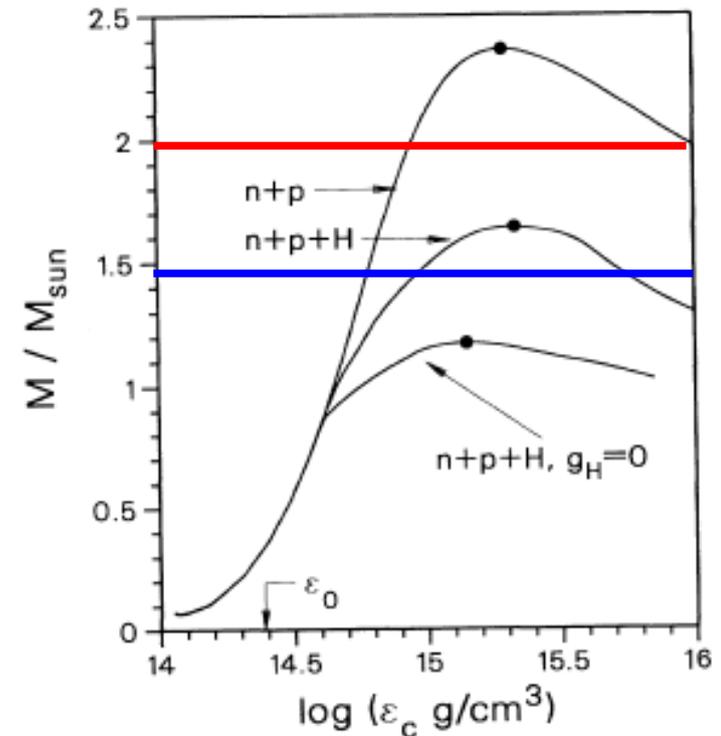
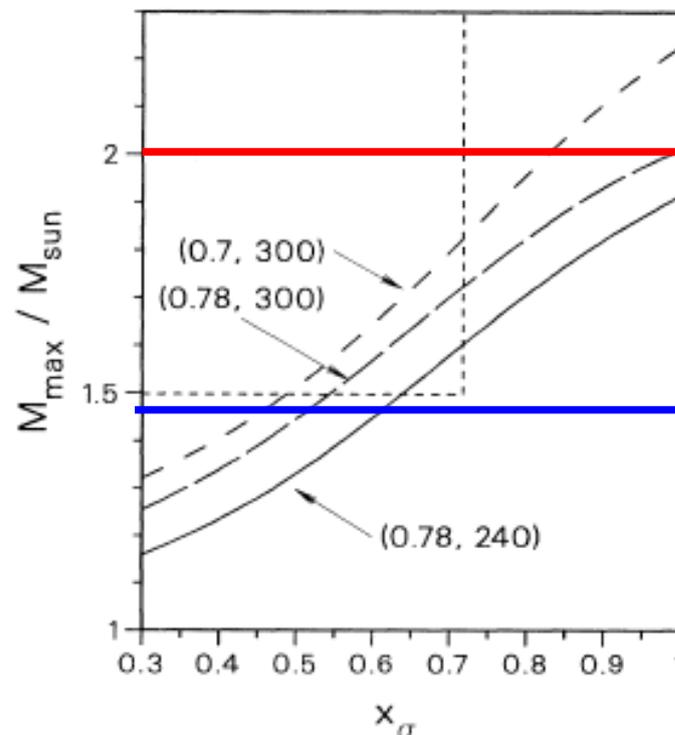
- クォーク物質でも相互作用に制限

クォーク物質  
の EOS



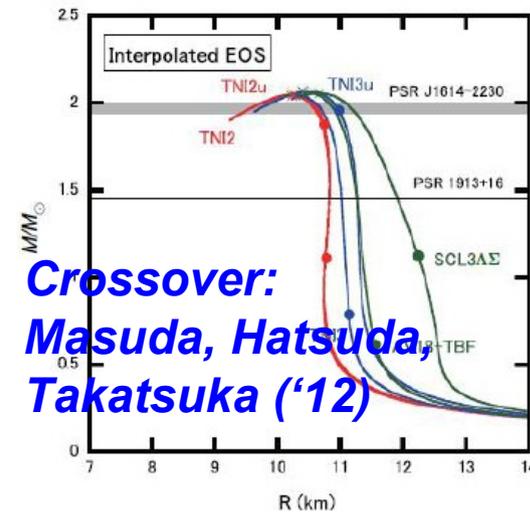
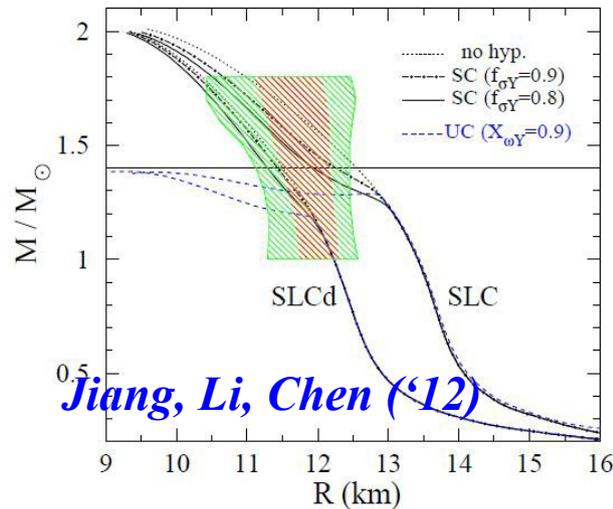
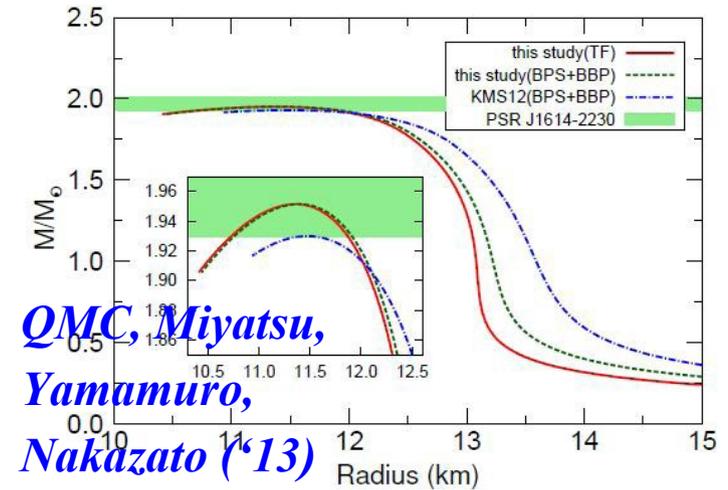
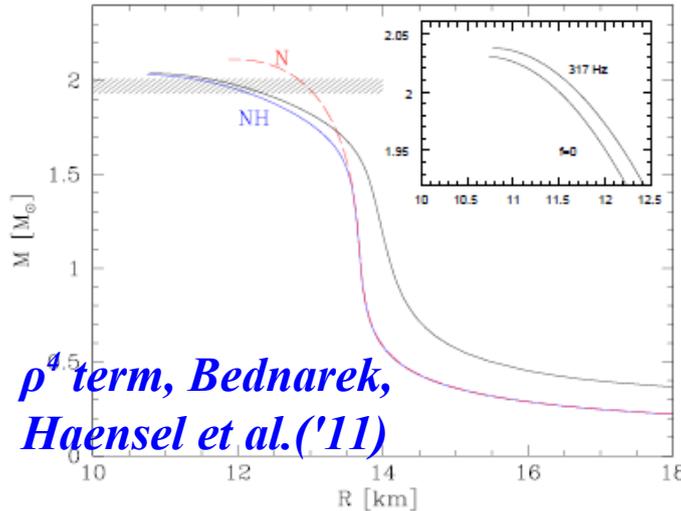
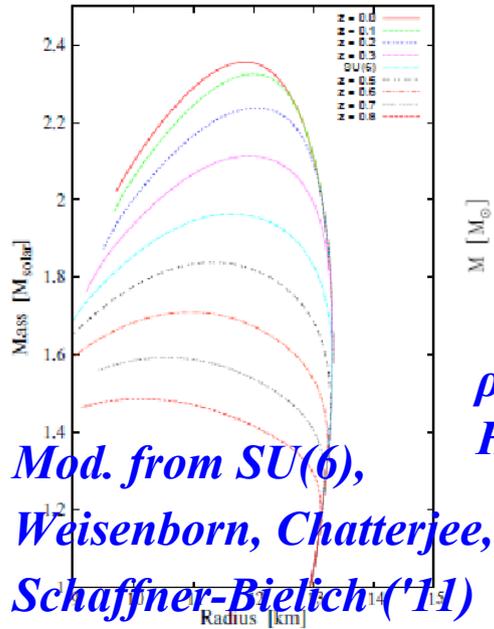
# ハイペロンと重い中性子星

- 「柔らかい EOS は否定されたが、exotic な構成粒子が否定されたわけではない。」 (Lattimer)
  - 否定されたハイペロン・K 中間子を含む EOS
    - = 相対論的平均場理論 (RMF) において結合定数をほぼ SU(6) にしたがって選んだもの ( $g_{\sigma\Lambda} / g_{\sigma N} \sim 2/3$ )



N.K. Glendenning, S.A. Moszkowski, *PRL*67('91)2414

# Massive Neutron Star with Hyperons



# RMF with many-body coupling

## ■ Naive dimensional analysis (NDA) and naturalness

*Manohar, Georgi ('84)*

The vertex is called “natural” if  $C \sim 1$ .

$$L_{\text{int}} \sim (f_\pi \Lambda)^2 \sum_{l,m,n,p} \frac{C_{lmnp}}{m!n!p!} \left( \frac{\bar{\psi} \Gamma \psi}{f_\pi^2 \Lambda} \right)^l \left( \frac{\sigma}{f_\pi} \right)^m \left( \frac{\omega}{f_\pi} \right)^n \left( \frac{R}{f_\pi} \right)^p$$

→ Consistent with the idea that the vertex is generated by loop diagrams under the assumption that the QCD coupling is small.

## ■ FST truncation

*R. J. Furnstahl, B. D. Serot, H. B. Tang, NPA615 ('97)441.*

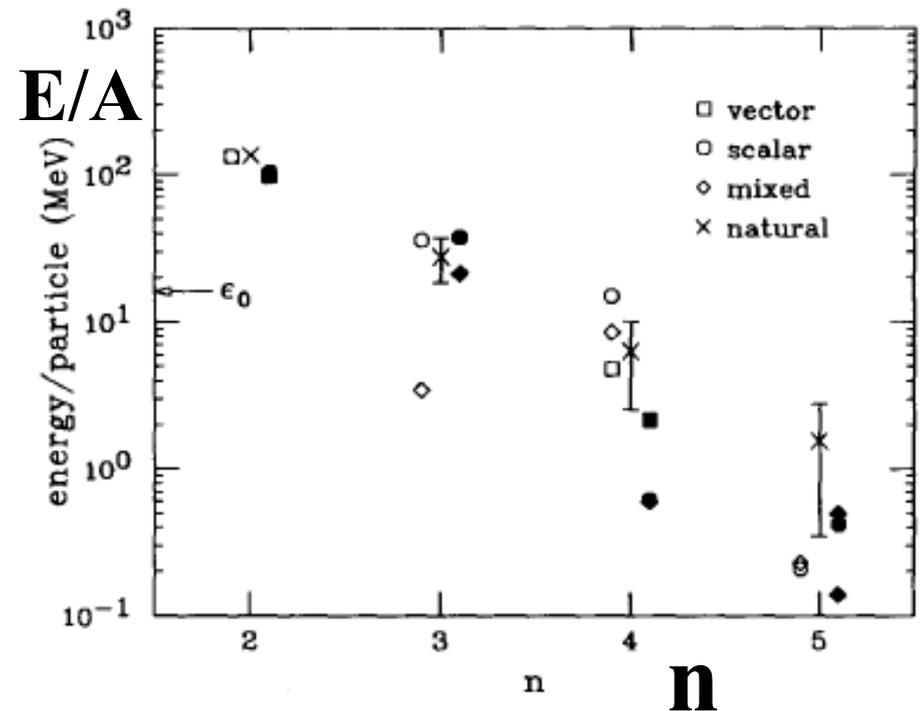
At a given density, we can truncate the Lagrangian by the index

$$\mathbf{n} = \mathbf{B}/2 + \mathbf{M} + \mathbf{D}$$

(B: baryon field, M: Non NG boson, D: derivatives)

Naturalness →  $V \sim \rho^n/n!$

→ small for large n

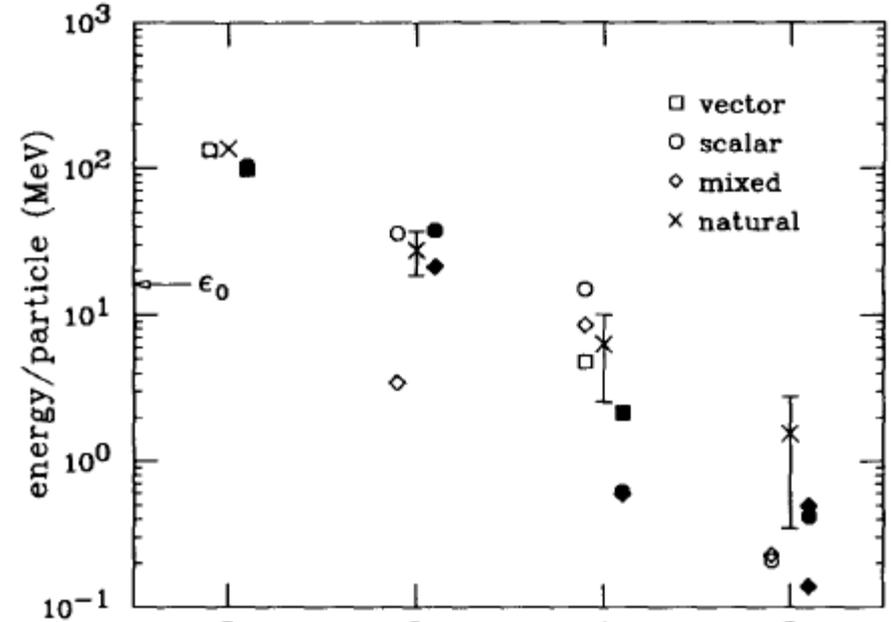
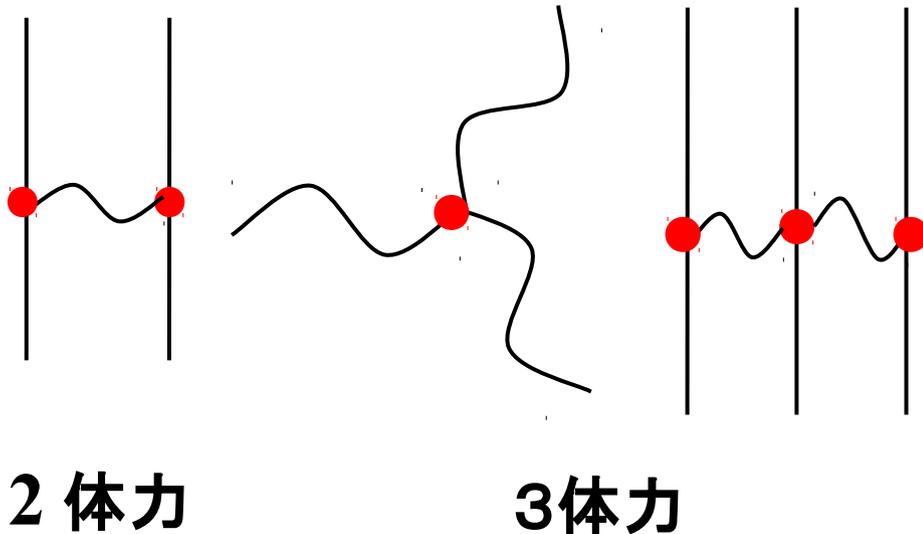


# 3 体力を含む RMF

## ■ 3 体力を含む相互作用項

$$\delta L = -\frac{1}{2} c_{\sigma\omega} \sigma \omega_\mu \omega^\mu - \sum_B \bar{\Psi}_B \left[ g_{\sigma\sigma B} \sigma^2 + g_{\sigma\omega B} \sigma \omega_\mu \gamma^\mu + g_{\omega\omega B} \omega_\mu \omega^\mu \right] \Psi_B$$

- BBMM 項は通常無視 (場の再定義で吸収可能)  
 しかし場の再定義は Naïve dimensional analysis (NDA) の次数を変え、高密度では n 体力が重要な役割を果たす。  
 $n = d + B/2 + M$  (d; 微分結合、B: バリオン場、M: non-NG ボソン)



Furnstahl, Serot, Tang ('97)

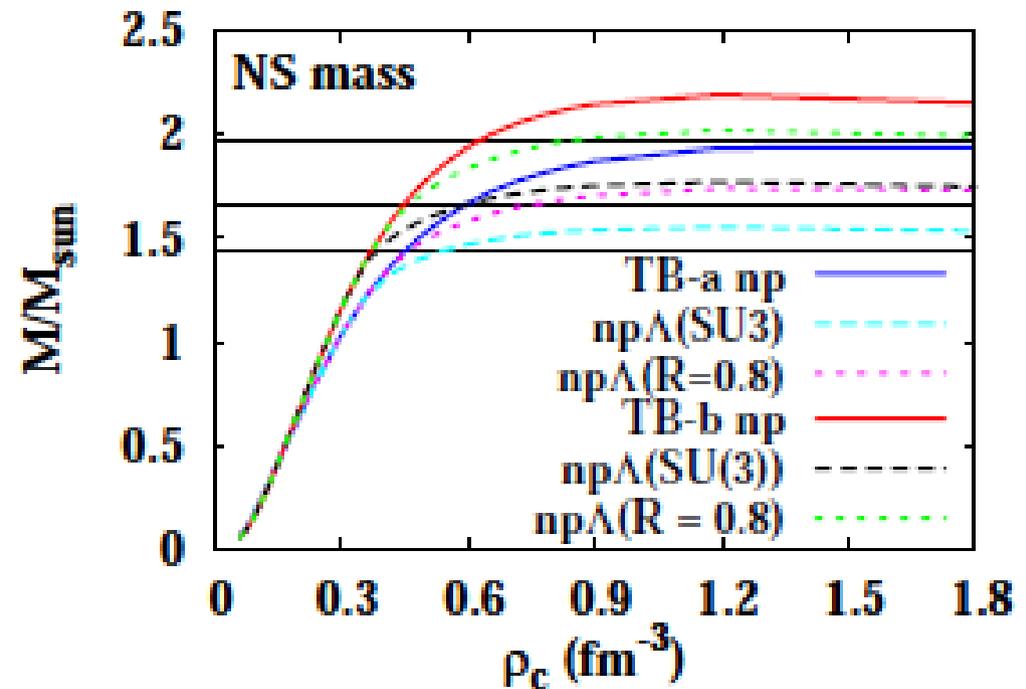
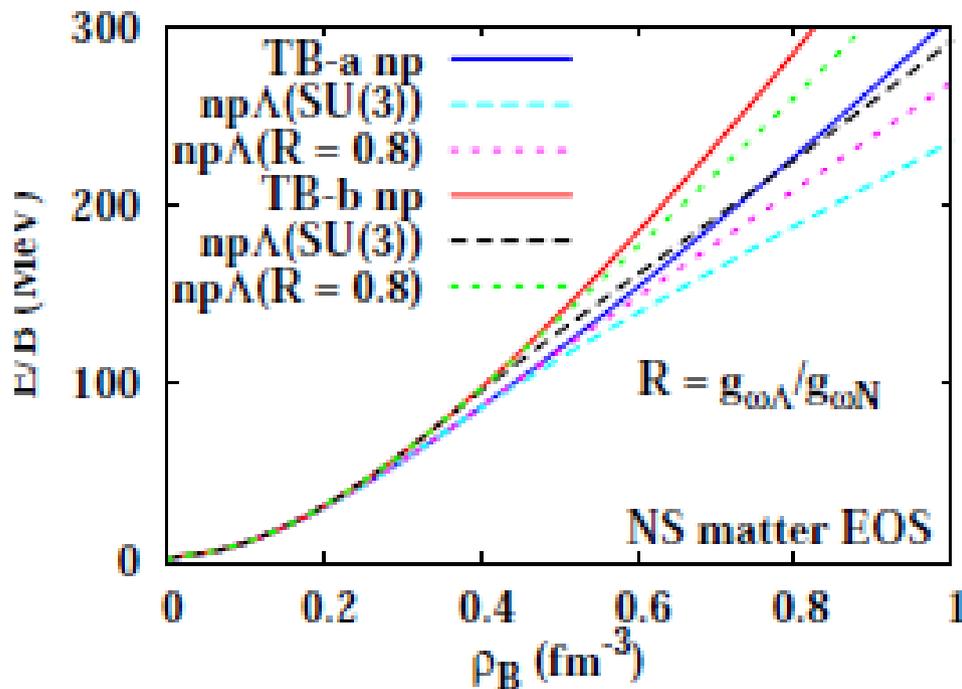
# RMF with 3BF + $SU(3)_f$ “violation”

## ■ Two types of modification

- 3-baryon repulsion → EOS becomes stiff gradually at high density. (Fitting meson mass (E325) and  $U_v$  in RBHF)

- $R = g_{\omega\Lambda} / g_{\omega N} \sim 0.8 (> 2/3 \text{ (SU(3))})$

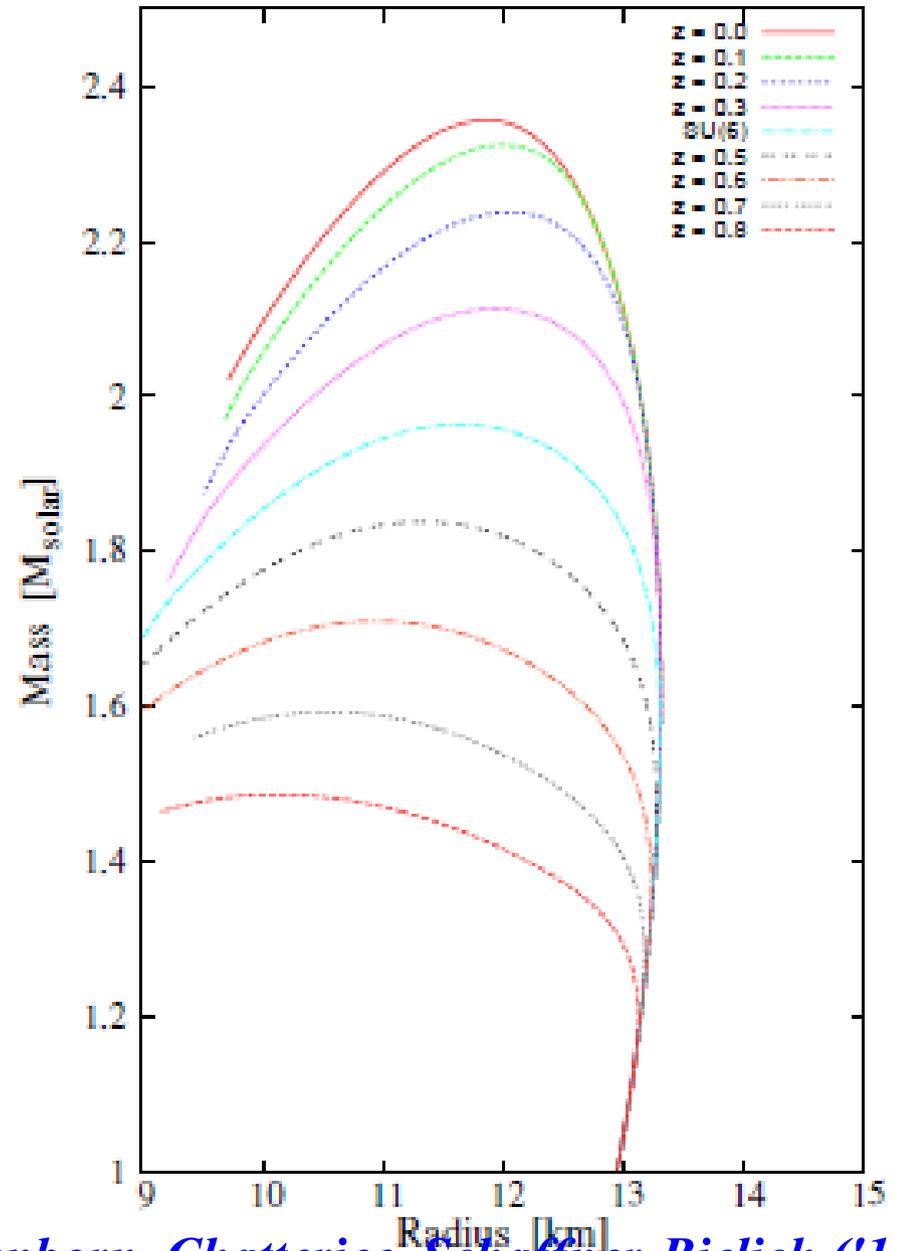
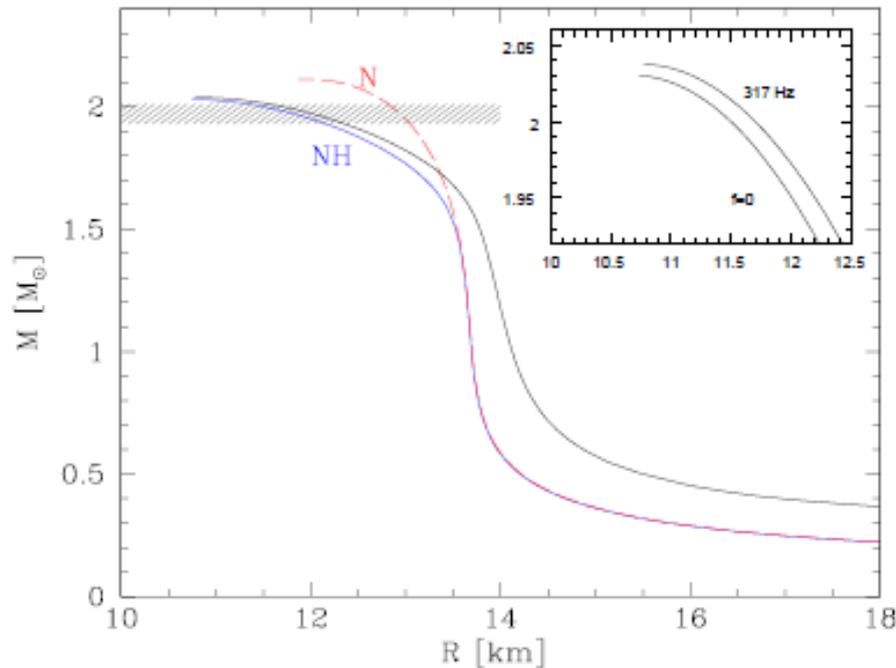
→  $M_{\max} \sim 2.02 M_{\odot}$  with hyperons ( $\sim 1.4 M_{\odot}$  w/o 3BF, violation)



Tsubakihara, AO, NPA914 ('13)438.

# 他の解決方法

- Hidden strange meson の4次
  - $\zeta^4, \zeta^2(\omega^2 + \rho^2)$  を導入
  - $\Lambda$  間の引力を小さく見積もる
  - SU(6) 関係式は保持
- Vector 結合に SU(6) の破れを導入



*Weisenborn, Chatterjee, Schaffner-Bielich ('11)*

*Bednarek, Haensel et al. ('11)*

## Summary of Lecture 4

- ハイパー核物理の進展により、核物質中でのハイペロン・ポテンシャルの深さについて理解が進んでいる。  
 $U_{\Lambda} \sim -30 \text{ MeV}$ ,  $U_{\Sigma} \sim +30 \text{ MeV}$ ,  $U_{\Xi} \sim -15 \text{ MeV}$
- 重い中性子星パズルは現象論的には解決可能。
  - SU(6) or SU(3) の破れ、現象論的 3 体力、Quark-meson coupling、Quark matter への crossover transition、...
- しかしながらこれらの扱いには実験的に確認できていないパラメータ・仮定を含んでおり、第一原理計算・実験データによる検証が必要。

*Hyperons in neutron stars: to be continued*

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*Thank you !*