

原子核基礎論B

京大基研 大西 明

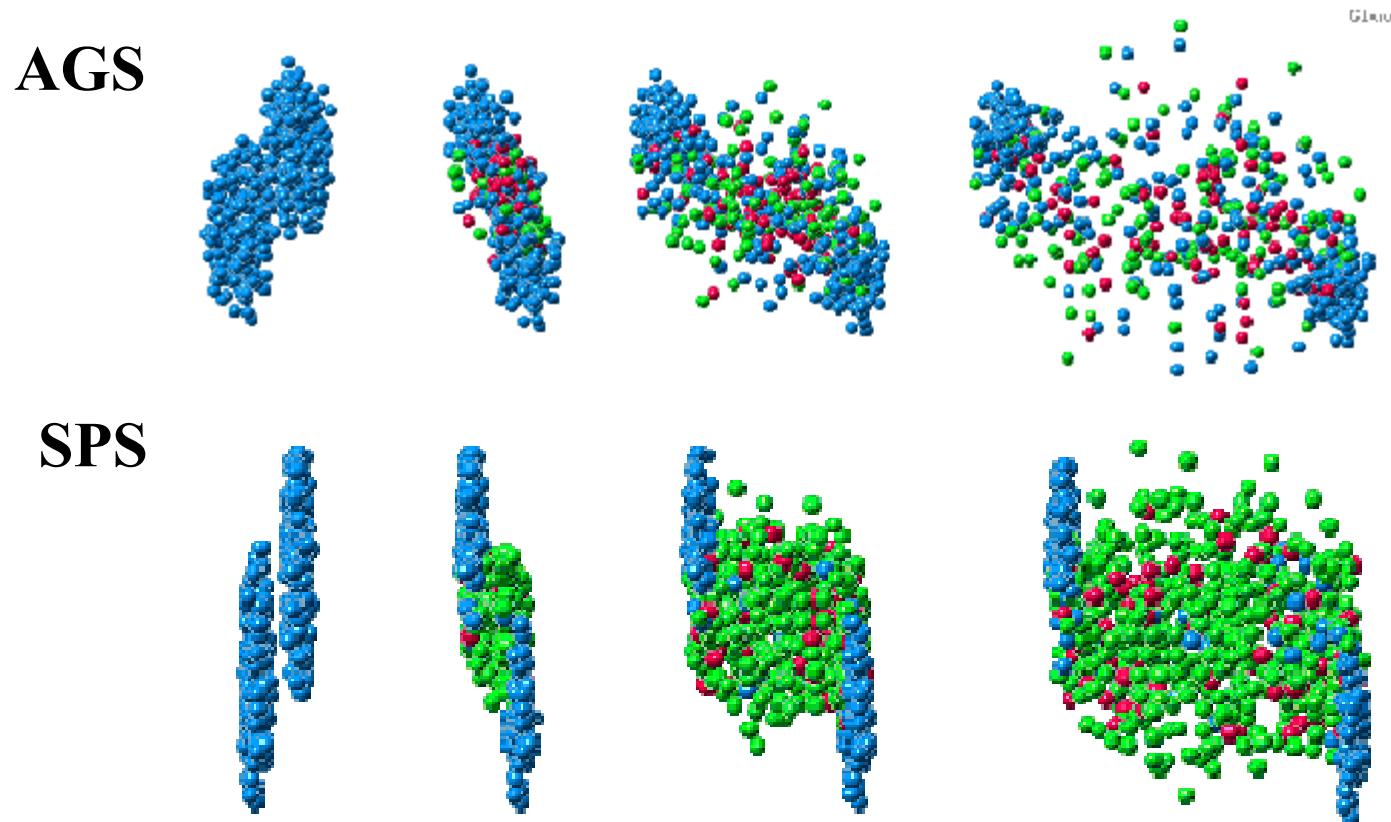
Akira Ohnishi (YITP, Kyoto U.)

1. 核力・特に非中心力や3体力(1回)
2. 原子核構造を記述するための種々の模型の最近の進展(2回)
3. 最近の中性子過剰核の物理の最近の進展(2回)
4. 原子核構造における異なる状態の混合や競合(2回) 板垣
5. 高温・高密度核物質概観(1回)(高エネルギー重イオン衝突、コンパクト天体现象)
→ 前期の Sec. 3 と重なりが大きいのでスキップ
6. 有限温度・密度における場の理論入門(2回)
7. QCD 有効模型における相転移と相図(2回)
8. 有限温度・密度格子 QCD と符号問題(1回) 大西
9. 高エネルギー重イオン衝突における輸送理論(1回)

Nuclear Transport Models and Heavy-Ion Collision

Heavy-Ion Collisions at $E_{\text{inc}} \sim (1\text{-}100) \text{ A GeV}$

- Study of Hot and Dense Hadronic Matter
→ Particle Yield, Collective Dynamics (Flow), EOS,

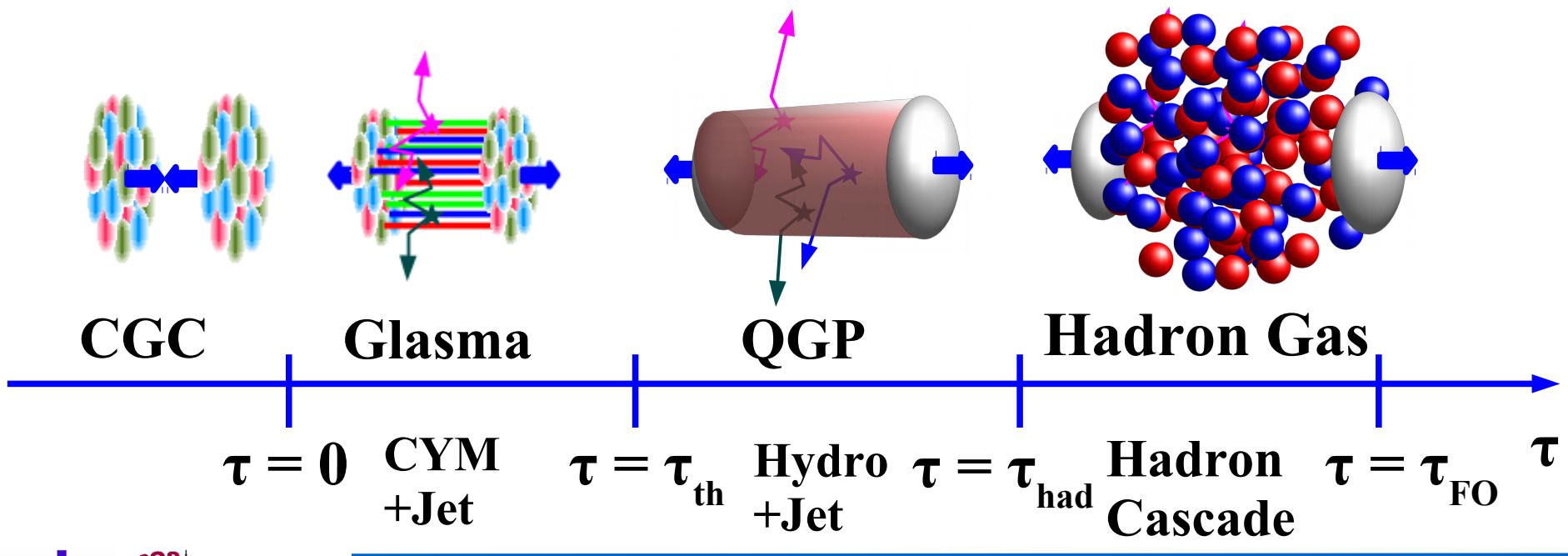


JAMming on the Web, linked from <http://www.jcprg.org/>

RHIC における2つの驚き (cont.)

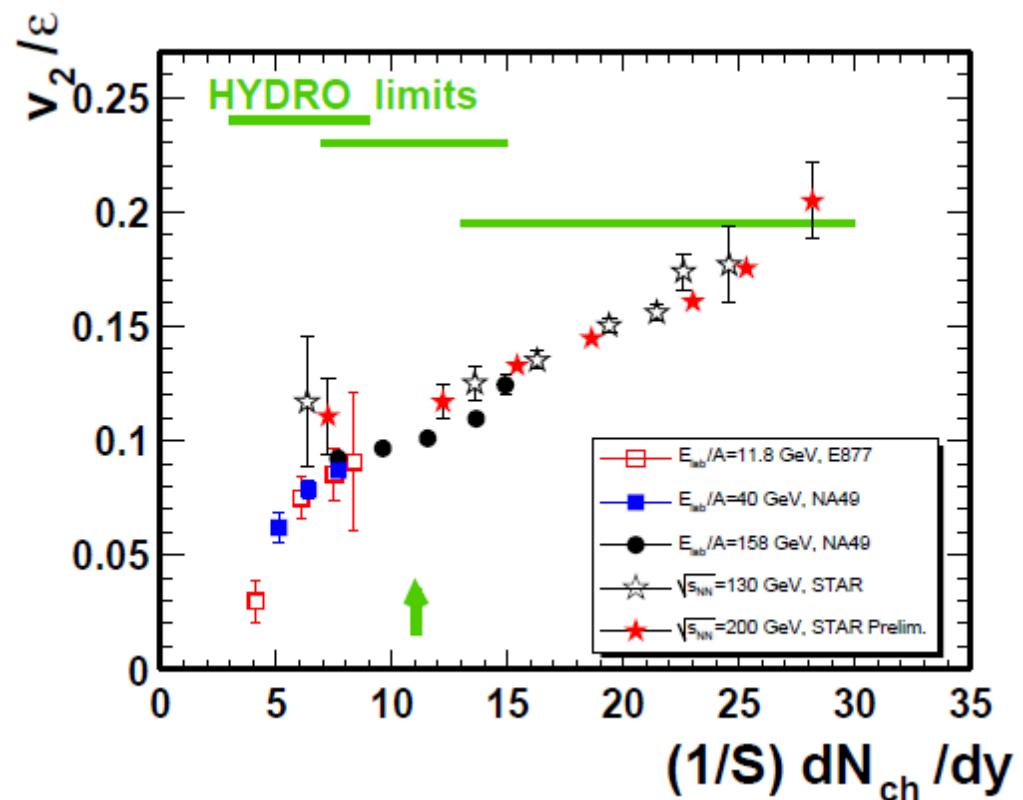
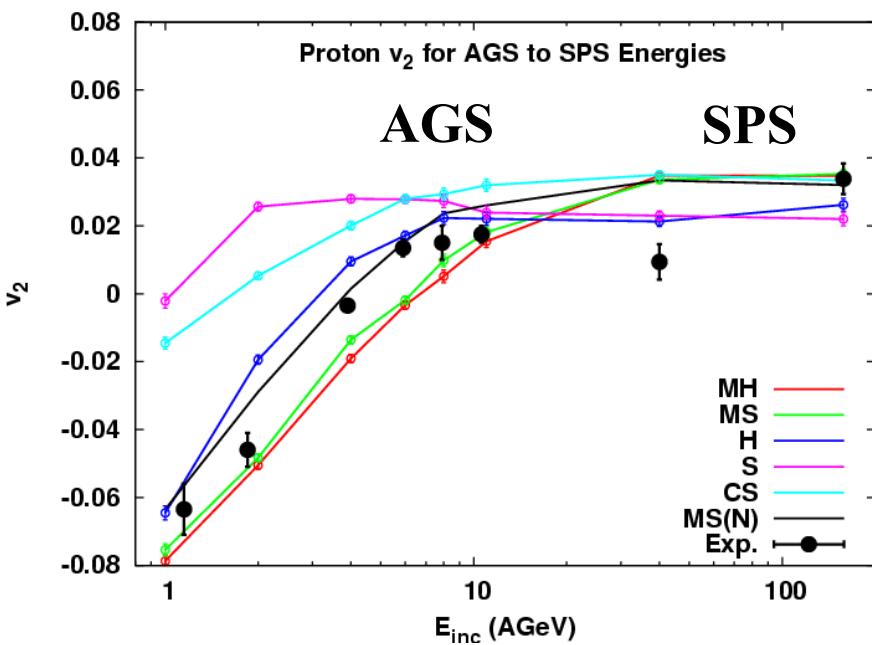
■ 2つの驚き (2): 早い熱平衡化

- 摂動論的 QCD の予言 (2-5 fm/c) に比べて有意に早い時刻 (0.6-1 fm/c) で熱化が起こり、流体力学的時間発展が進む。
→ なぜ早い？
- 高エネルギー重イオン衝突の初期条件
= グラズマ (古典ヤンミルズ場が主要)



Hydrodynamics vs Transport

- $\sqrt{s_{NN}} < 20 \text{ GeV} \rightarrow$ Transport model calculation seems to explain v_2 data.
- RHIC (& LHC) \rightarrow Hydrodynamics is successful.



M. Isse, AO, N. Otuka, P.K. Sahu, Y. Nara,
PRC72 ('05) 064908 [nucl-th/0502058]

U. W. Heinz, AIP Conf. Proc. 739 ('05) 163
[nucl-th/0407067]

HIC Transport Models: Five Major Origins

■ *Nuclear Mean Field Dynamics*

- Basic Element of Low Energy Nuclear Physics
- TDHF → Vlasov → BUU

■ *NN two-body (residual) interaction*

- Main Source of Particle Production
- Boltzmann equation → Intranuclear Cascade Models

■ *Partonic Interaction and String Decay*

- Main Source of high pT Particles at Collider Energies
- JETSET + (previous) PYTHIA (Lund model) → (new) PYTHIA

■ *Relativistic Hydrodynamics*

- Most Successful Picture at RHIC

■ *Classical Field Dynamics*

- Classical Gluon Field = Initial condition of Hydro. at Collider Energies

Nuclear Mean Field Dynamics

TDHF and Vlasov Equation

- Time-Dependent Mean Field Theory (e.g., TDHF) $i\hbar \frac{\partial \phi_i}{\partial t} = h\phi_i$
- Density Matrix

$$\rho(\mathbf{r}, \mathbf{r}') = \sum_i^{\text{Occ}} \phi_i(\mathbf{r}) \phi_i^*(\mathbf{r}') \quad \rightarrow \rho_W = f \text{ (phase space dist.)}$$

- TDHF for Density Matrix

$$i\hbar \frac{\partial \rho}{\partial t} = [h, \rho] \quad \rightarrow \frac{\partial f}{\partial t} = \{h_W, f\} + \mathcal{O}(\hbar^2)$$

$(AB)_W = A_W \exp(i\hbar\Lambda) B_W$, $\Lambda = \nabla'_r \cdot \nabla_p - \nabla'_p \cdot \nabla_r$ (∇' acts on the left)

- Wigner Transformation and Wigner-Kirkwood Expansion
(Ref.: Ring-Schuck)

$$A_W(\mathbf{r}, \mathbf{p}) = \int d^3s \exp(-i\mathbf{p} \cdot \mathbf{s}/\hbar) \langle \mathbf{r} + \mathbf{s}/2 | A | \mathbf{r} - \mathbf{s}/2 \rangle$$

$$(AB)_W = A_W \exp(i\hbar\Lambda) B_W \quad \Lambda \equiv \nabla'_r \cdot \nabla_p - \nabla'_p \cdot \nabla_r \quad (\nabla' \text{ acts on the left})$$

$$[A, B]_W = 2iA_W \sin(\hbar\Lambda/2)B_W = i\hbar \{A_W, B_W\}_{P.B.} + O(\hbar^3)$$

Test Particle Method

■ Vlasov Equation

$$\frac{\partial f}{\partial t} - \{ h_W, f \}_{P.B.} = \frac{\partial f}{\partial t} + v \nabla_r f - \nabla U \cdot \nabla_p f = 0$$

■ Classical Hamiltonian

$$h_W(r, p) = \frac{p^2}{2m} + U(r, p)$$

■ Test Particle Method (C. Y. Wong, 1982)

$$f(r, p) = \frac{1}{N_0} \sum_i^{AN_0} \delta(r - r_i) \delta(p - p_i) \rightarrow \frac{dr_i}{dt} = \nabla_p h_w, \quad \frac{dp_i}{dt} = -\nabla_r h_w,$$

*Mean Field Evolution can be simulated
by Classical Test Particles
→ Opened a possibility to Simulate High Energy HIC
including Two-Body Collisions in Cascade*

BUU (Boltzmann-Uehling-Uhlenbeck) Equation

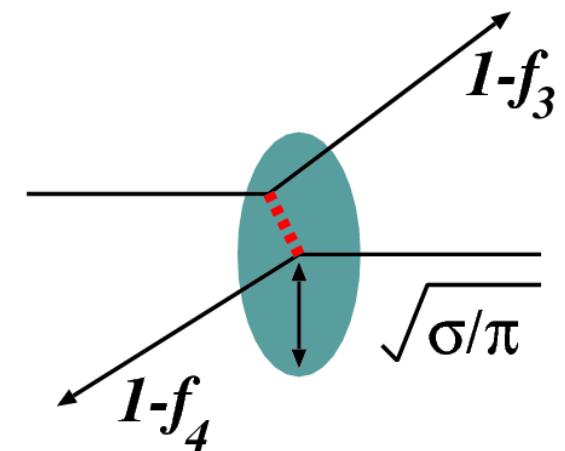
■ BUU Equation (Bertsch and Das Gupta, Phys. Rept. 160(88), 190)

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \nabla_r f - \nabla U \cdot \nabla_p f = I_{coll}[f]$$

$$I_{coll}[f] = -\frac{1}{2} \int \frac{d^3 p_2 d\Omega}{(2\pi\hbar)^3} V_{12} \frac{d\sigma}{d\Omega} \times [f f_2 (1-f_3)(1-f_4) - f_3 f_4 (1-f)(1-f_2)]$$

■ Incorporated Physics in BUU

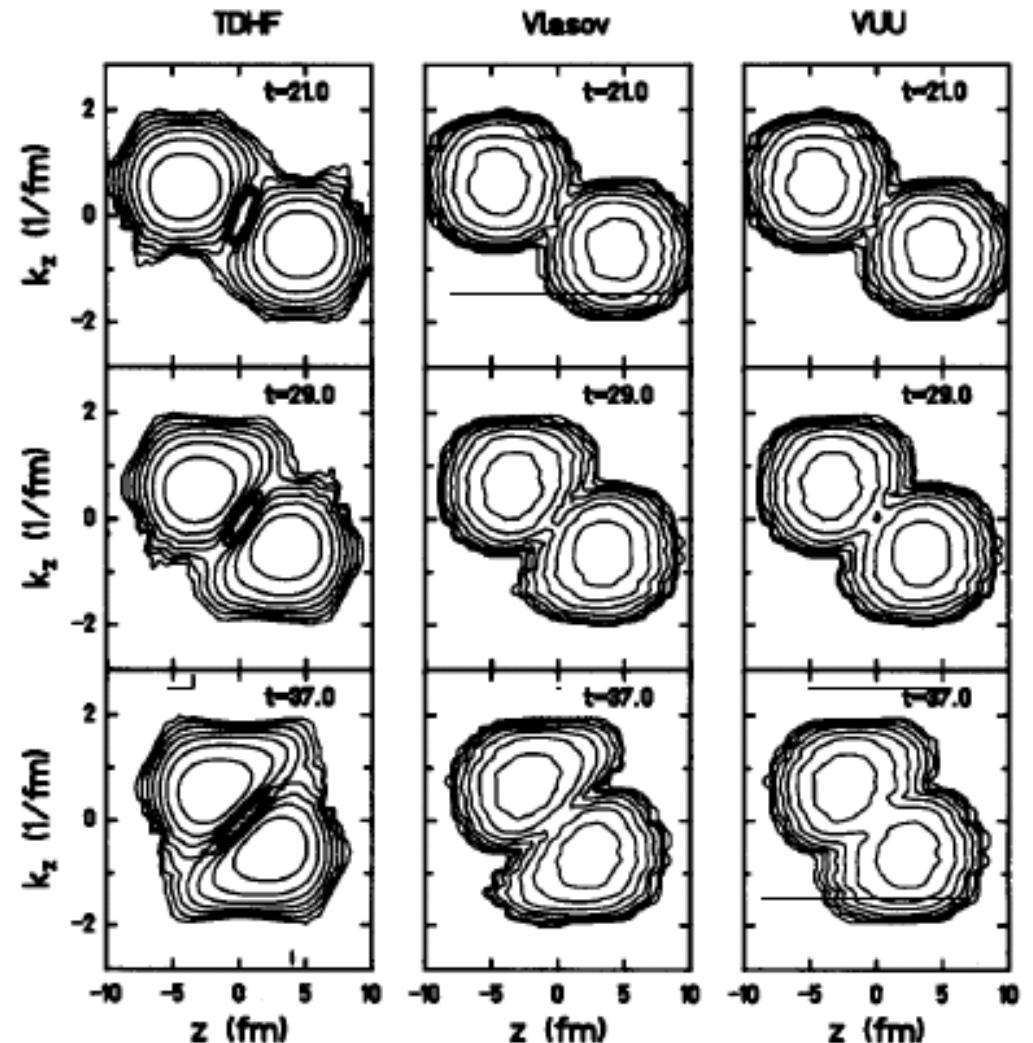
- Mean Field Evolution
- (Incoherent) Two-Body Collisions
- Pauli Blocking in Two-Body Collisions



○ *One-Body Observables (Particle Spectra, Collective Flow,..)*
✗ *Event-by-Event Fluctuation (Fragment, Intermittency,..)*

Comarison of TDHF, Vlasov and BUU(VUU)

- Ca+Ca, 40 A MeV
(Cassing-Metag-Mosel-Niita, Phys. Rep. 188 (1990) 363).



Collision Term and Particle Production

Baryon-Baryon and Meson-Baryon Collisions

■ NN collision mechanism

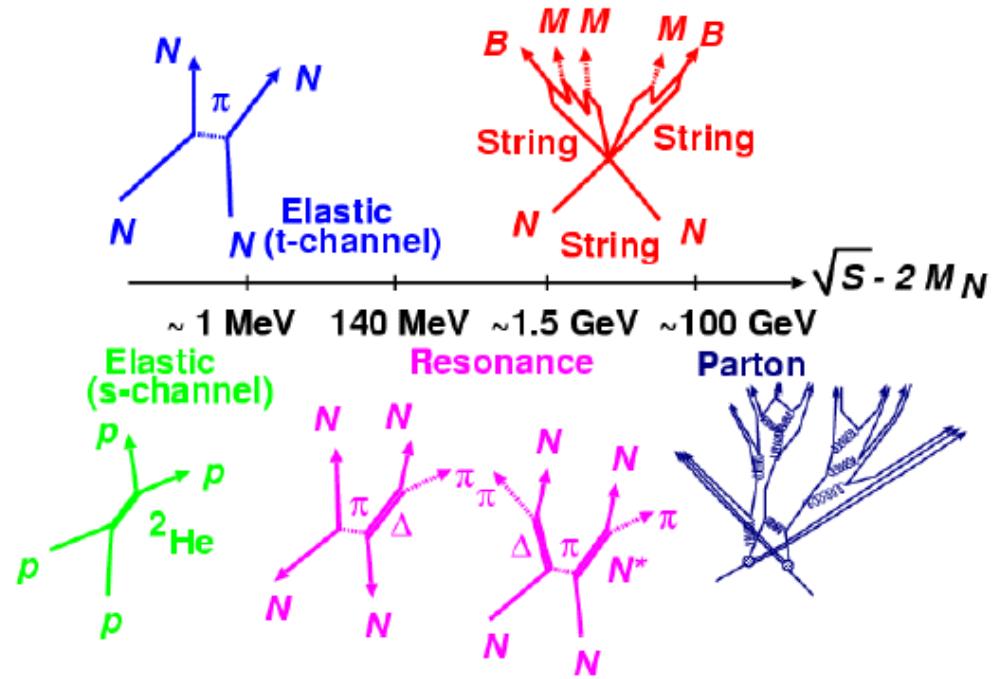
Elastic

→ Resonance

→ String

→ Jet

Energy Dependence of NN Reaction Mechanism

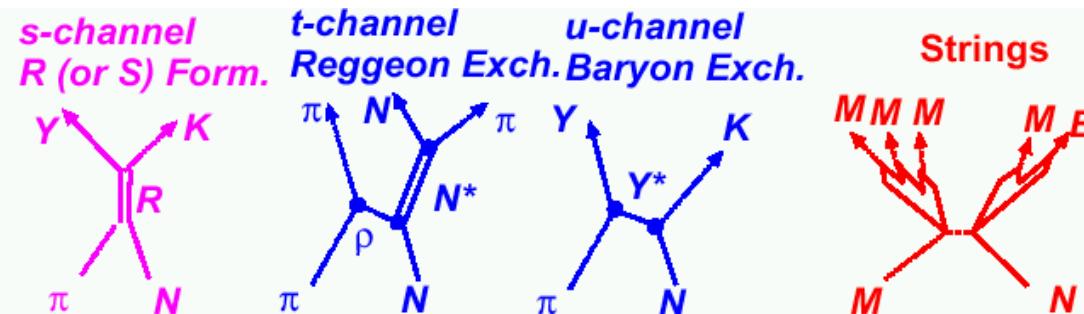


■ Meson-Nucleon Collision

→ s-channel Resonance

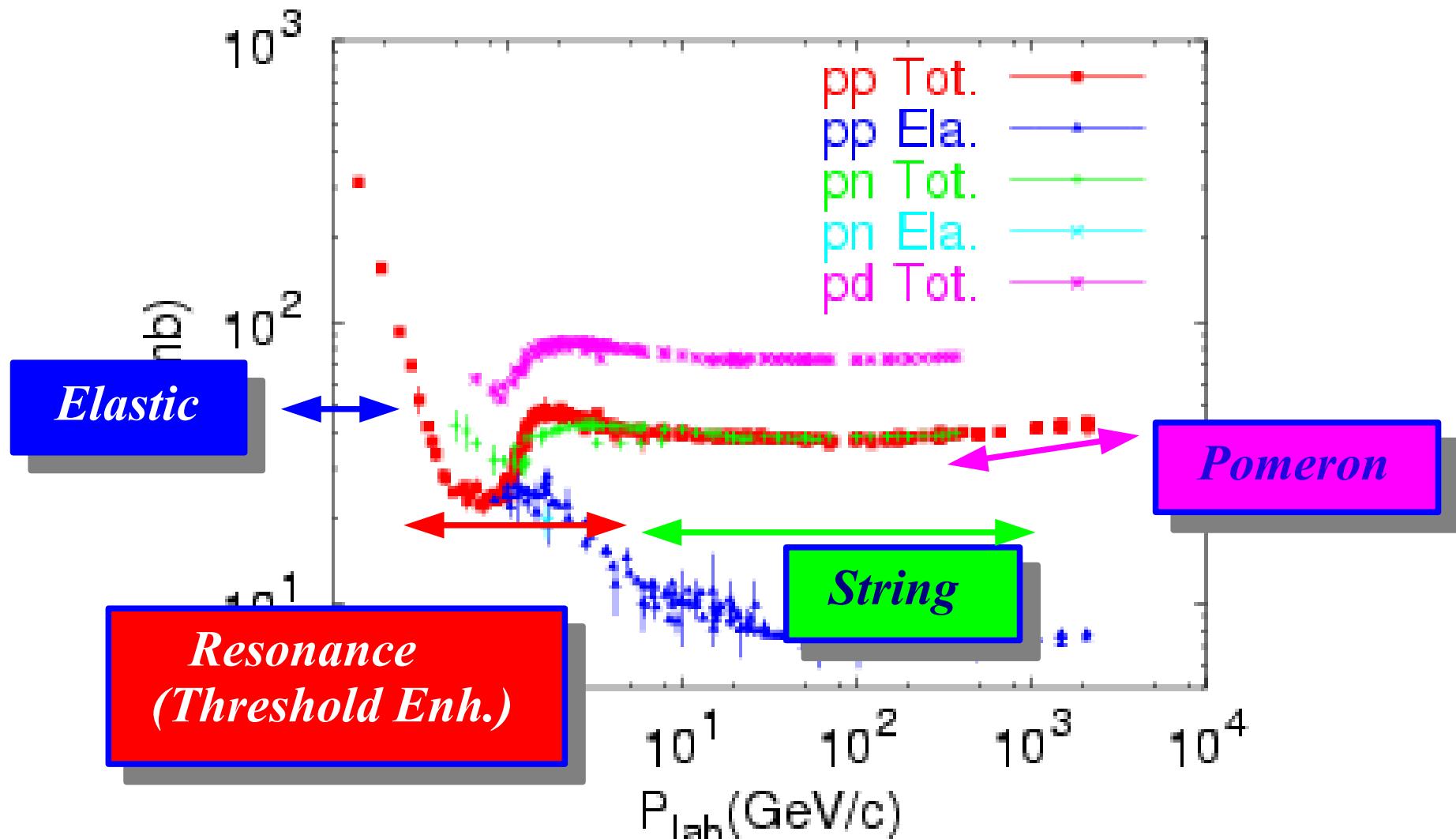
→ t-(u-) channel Res.

→ String formation

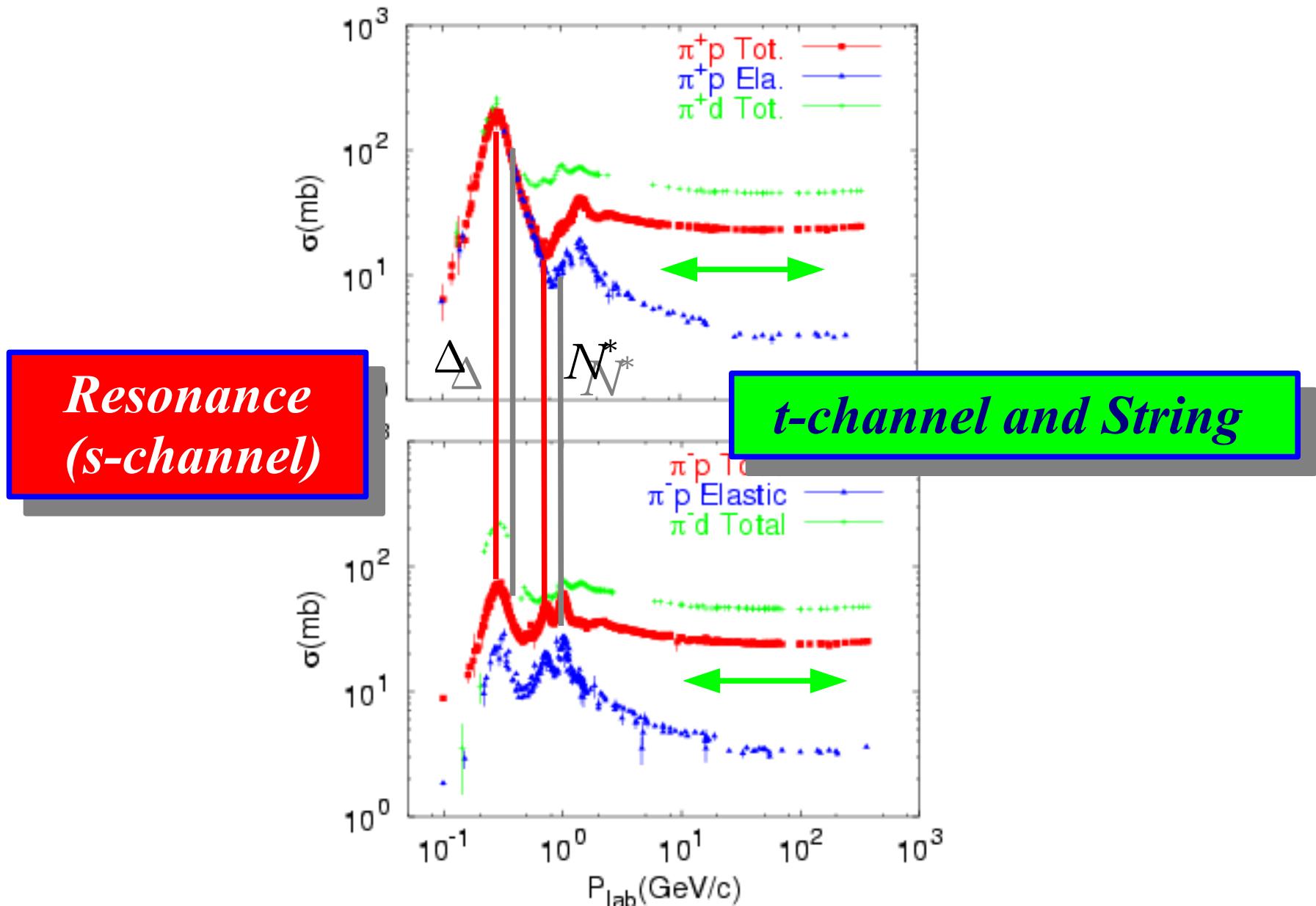


NN Cross Sections

From Particle Data Group



Meson-Baryon Cross Section



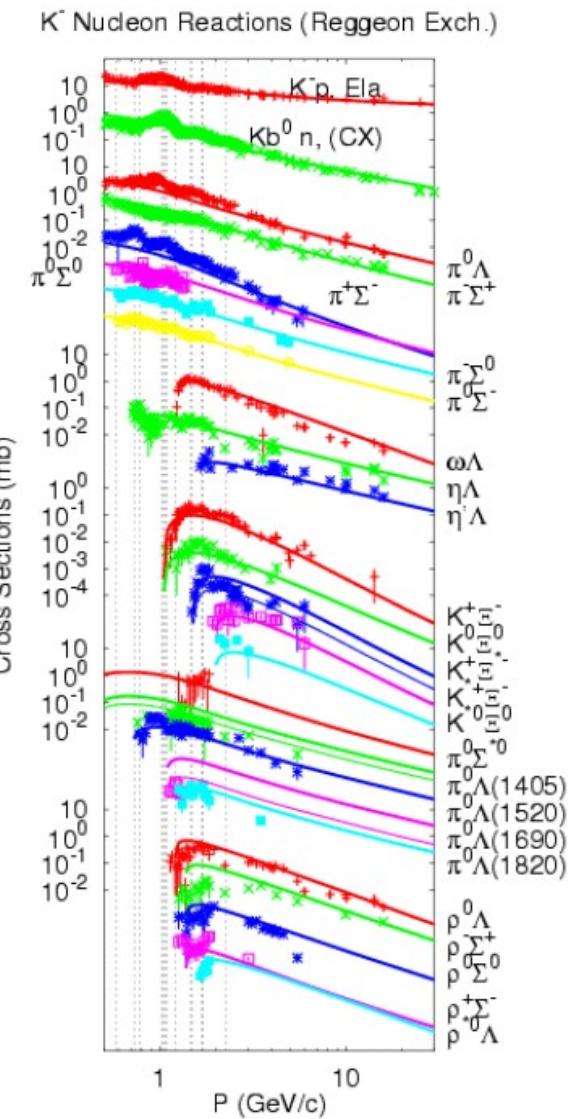
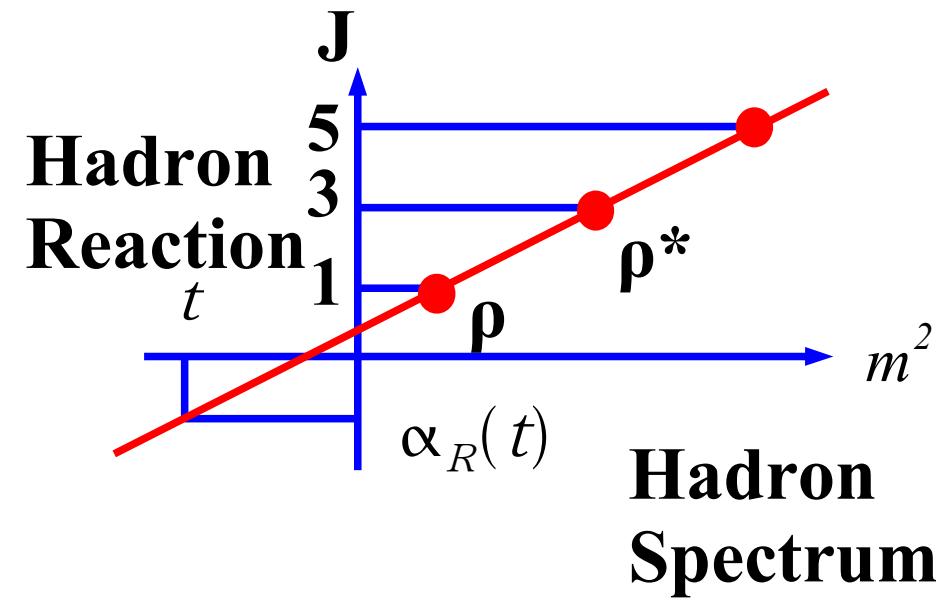
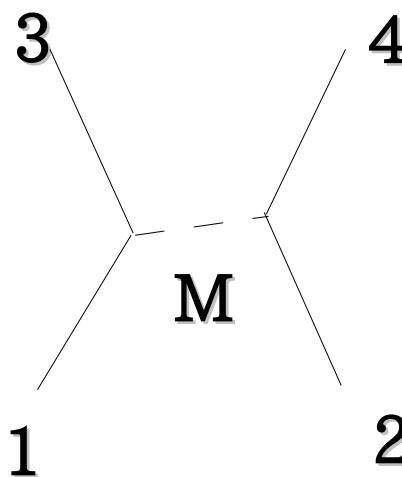
Reggeon Exchange

(Barger and Cline (Benjamin, 1969), H. Sorge, PRC (1995), RQMD2.1)

- Regge Trajectory $J = \alpha_R(t) \sim \alpha_R(0) + \alpha'_R(0)t$
- 2 to 2 Cross Section

$$\frac{d\sigma}{d\Omega} = \frac{p_f}{64\pi s p_i} |M(s, t)|^2$$

$$M(s, t) \sim \sum_R \frac{(p_i p_f)^J}{t - M_R} \sim F(t) \exp[\alpha_R(t) \log(s/s_0)]$$

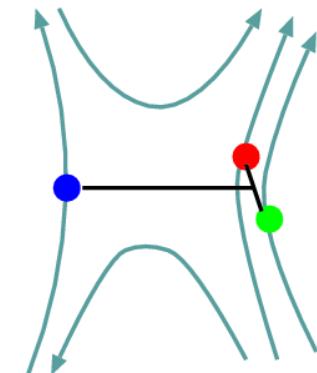


String formation and decay

- What does the regge trajectory suggest ?
→ Existence of (color- or hadron-)String !

$$M = 2 \int_0^R \frac{\kappa dr}{\sqrt{1 - (r/R)^2}} = \pi \kappa R , \quad J = 2 \int_0^R r \times \frac{\kappa dr}{\sqrt{1 - (r/R)^2}} \frac{r}{R} = \frac{\pi \kappa R^2}{2} \pi$$

$$\rightarrow J = \frac{M^2}{2\pi\kappa}$$

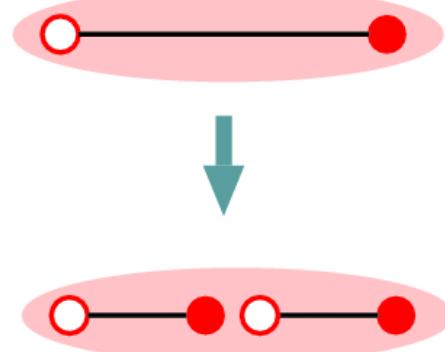


- String Tension

$$\frac{1}{2\pi\kappa} = \alpha'_R(0) \approx 0.9 \text{ GeV}^{-2} \rightarrow \kappa \approx 1 \text{ GeV/fm}$$

$$\text{string} = s + d + g + \dots$$

- String decay
- Extended String
→ Large E stored
→ q qbar pair creation (Schwinger mech.)



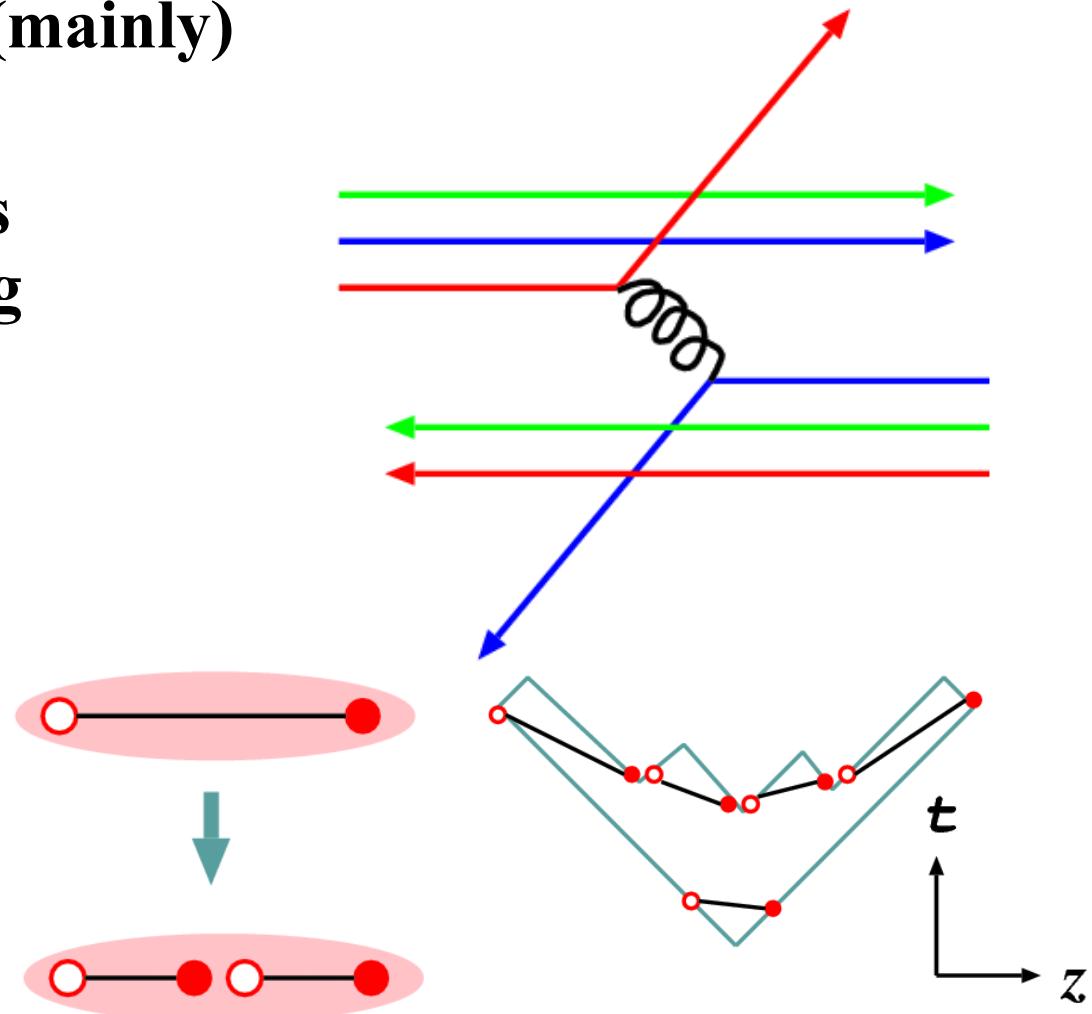
String = Coherent superposition of hadron resonances with various J

Jet Production

- Elastic Scattering of Partons (mainly) with One Gluon Exch.
- Color Exch. between Hadrons
 - Complex color flux starting from leading partons
 - many hadron production
 - Jet production

PYTHIA

- Event Generator of High Energy Reactions
 - Jet production +String decay for QCD processes



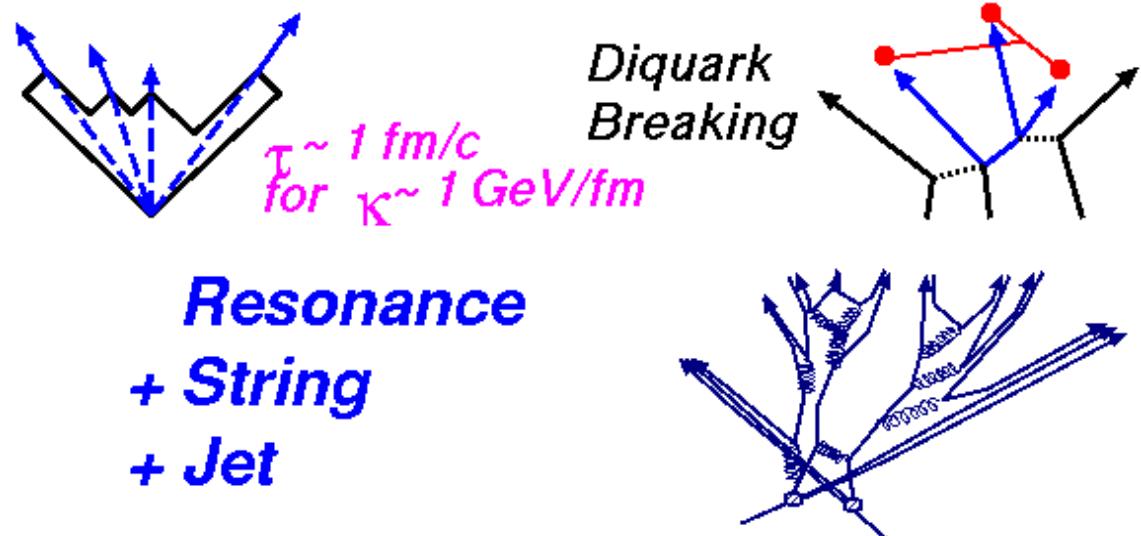
(T. Sjöstrand et al., Comput. Phys. Commun. 135 (2001), 238.)

JAM (*Jet AA Microscopic transport model*)

Nara, Otuka, AO, Niita, Chiba, Phys. Rev. C61 (2000), 024901.

■ Hadron-String Cascade with Jet production

- hh collision with Res. up to $m < 2$ GeV (3.5 GeV) for M (B)
- String excitation and decay
- String-Hadron collisions are simulated by hh collisions in the formation time.
- jet production is incl. using PYTHIA
- Secondary partonic int.:
 NOT incl.
- Color transparency:
 NOT taken care of



Relativistic Hydrodynamics

Relativistic Hydrodynamics

■ EOM: Conservation Laws

$\partial_\mu T^{\mu\nu} = 0$ Energy Momentum Conservation

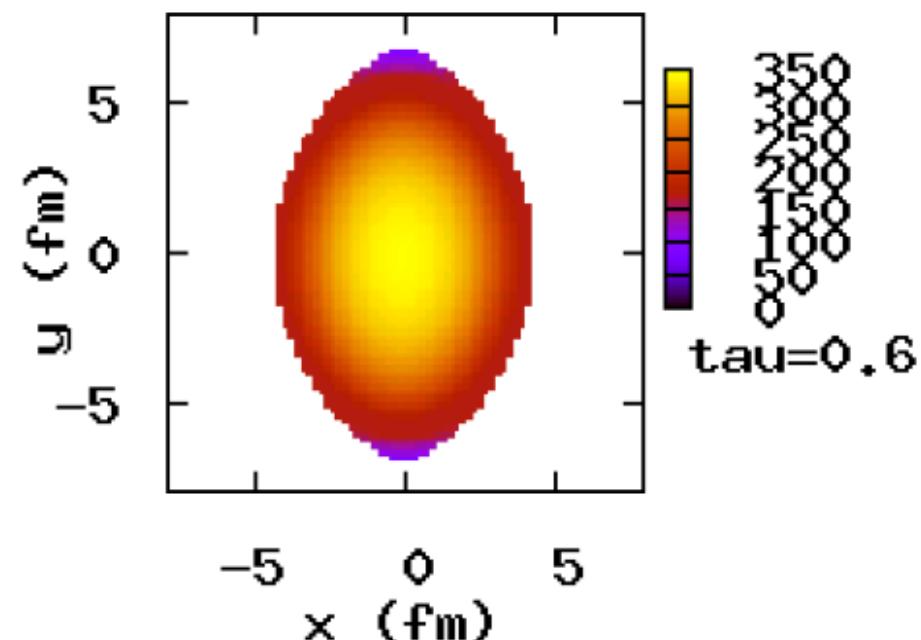
$\partial_\mu (n_i u^\mu) = 0$ Conservation of Charge (Baryon, Strangeness, ...)

$$T^{\mu\nu} = (e + P)u^\mu u^\nu - Pg^{\mu\nu}$$

e : energy density, P : pressure,

u^μ : four velocity $\gamma(1, v)$,

n_i : number density

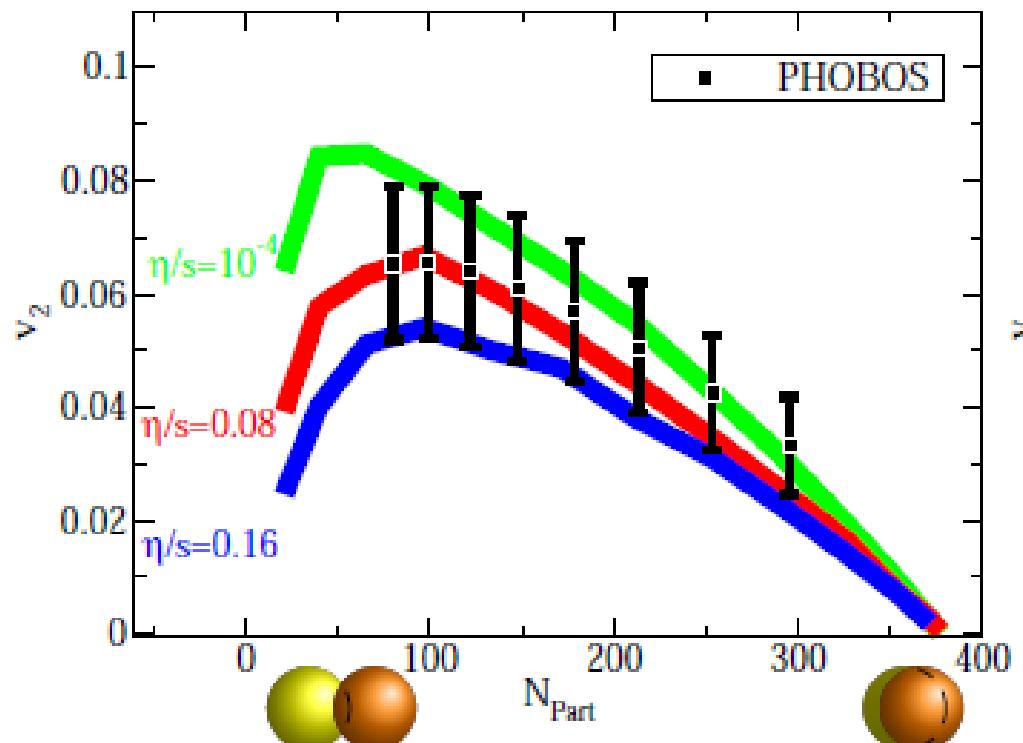


T. Hirano, Y. Nara, NPA743, 305 (2004)

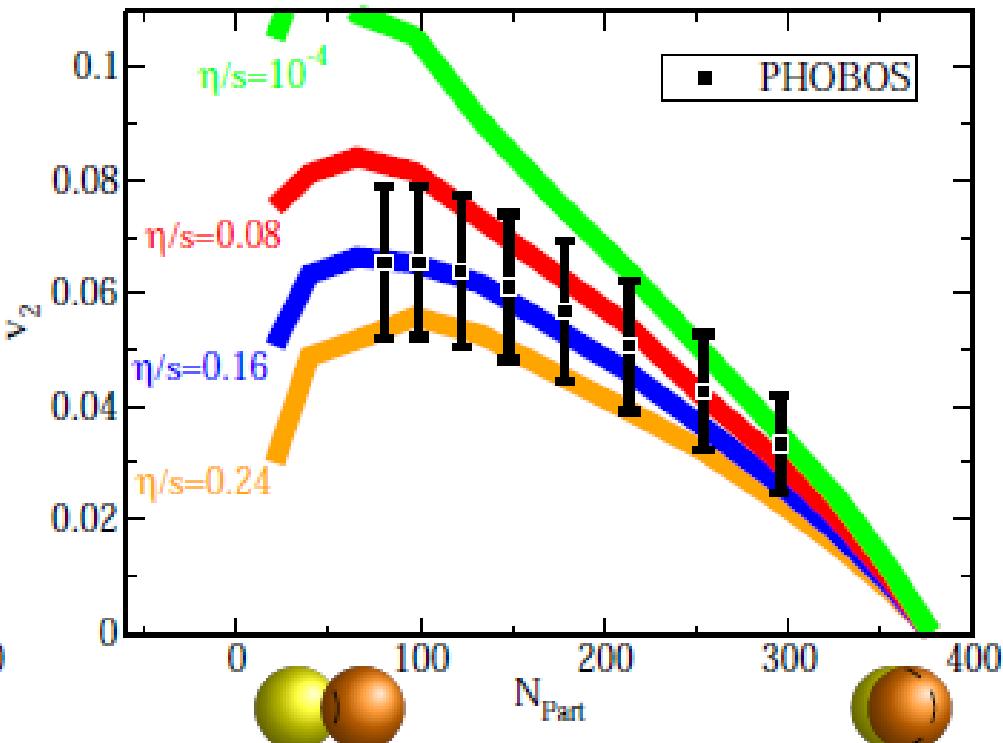
T. Hirano, K. Tsuda, PRC 66, 054905(2002)

η/s FROM FLOW (HISTORICAL)

“Glauber” initial conditions



“CGC” initial conditions



(ML & Romatschke, Phys. Rev. C78 (2008) 034915)

- Best extraction of η/s by comparing viscous hydro to flow data
- Largest uncertainty from unknown initial condition

低エネルギーでのハドロン輸送模型の成功
+高エネルギーでの流体模型の成功
→ QGP 生成の始まりは $\sqrt{s_{NN}} = 40 \text{ GeV}$ 程度か？

$\sqrt{s}_{NN} = (5-20) \text{ GeV}$ で QCD 相転移は見えるか？

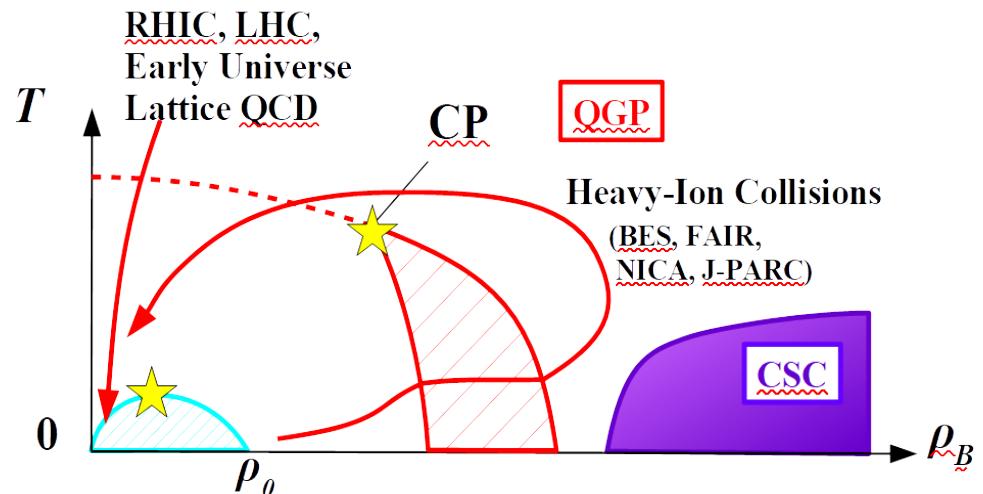
高エネルギー重イオン衝突の物理

■ クオーク・グルーオン・プラズマ (QGP) の発見・物性

- 流体模型とジェットによる記述の成功 → 精緻化へ
(RHIC の最高エネルギー、LHC)

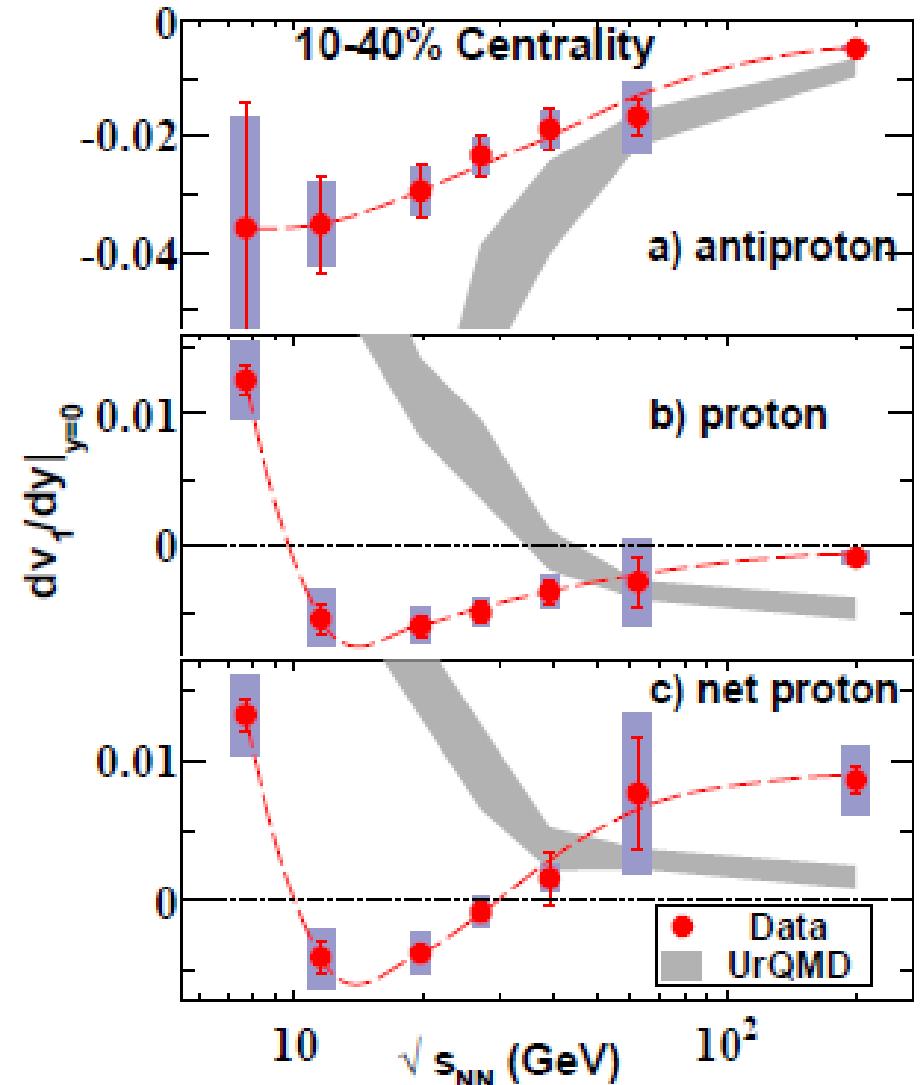
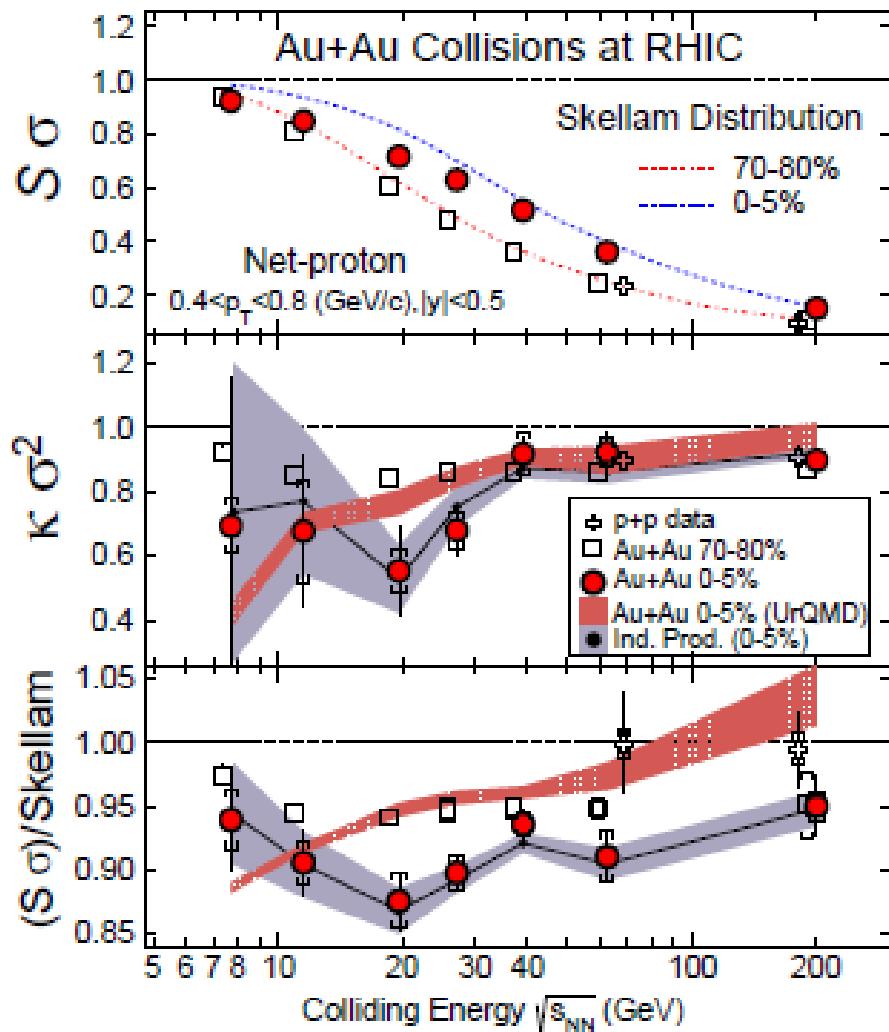
■ 高密度領域における状態方程式 (EOS) ・ QCD 相転移

- $\sqrt{s_{NN}} = (5-20) \text{ GeV}$ の衝突エネルギー領域 → $\rho = (5-10)\rho_0$
- 高密度での EOS は決められるか？
- 高密度での相転移は 1 次か？クロスオーバーか？



QCD 一次相転移は見えたか？

- 衝突エネルギー関数として非単調性が見えている ($\kappa\sigma^2$, dv_1/dy)



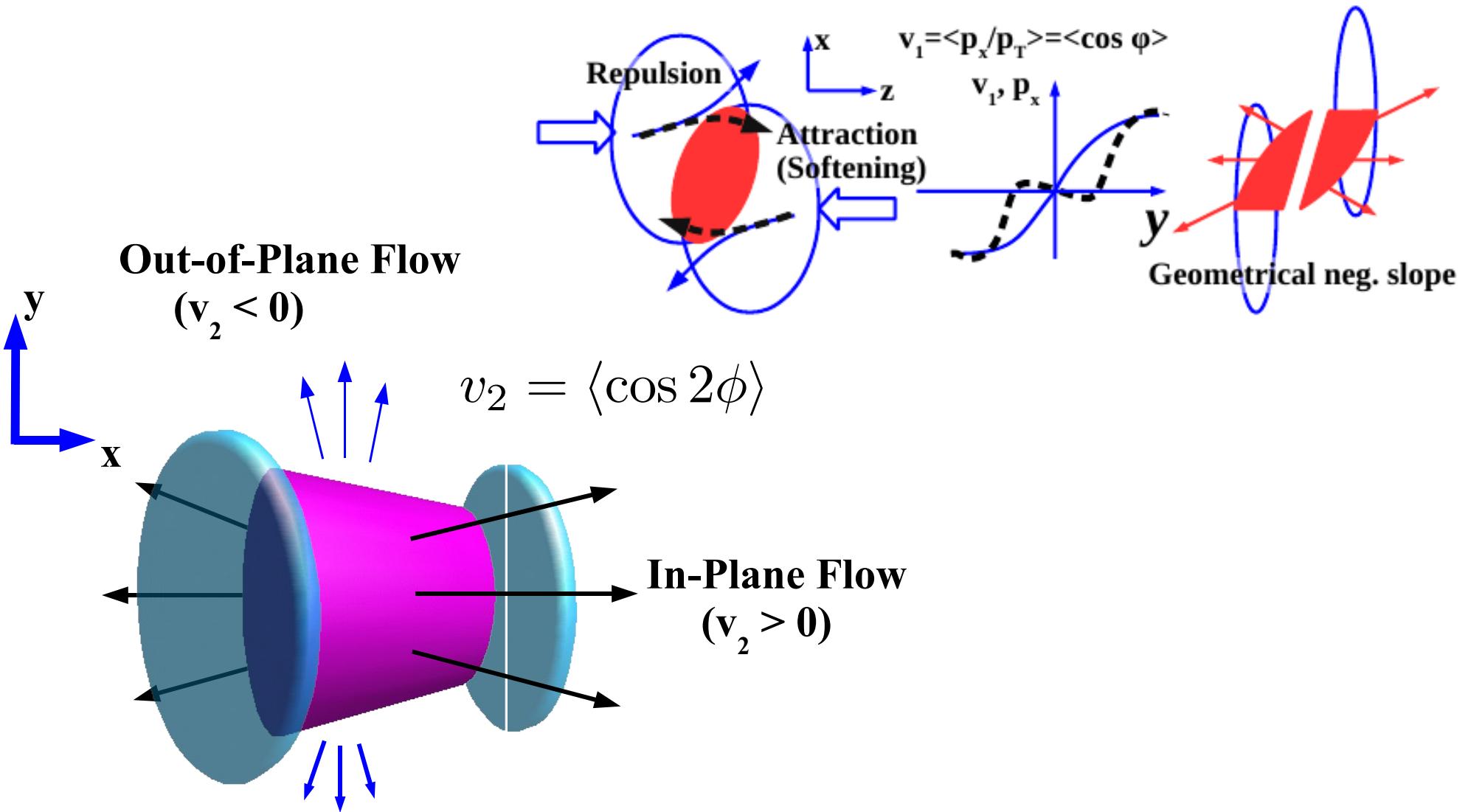
STAR Collab., PRL 112('14)032302.

STAR Collab., PRL 112('14)162301.

集団フロー

- Directed flow ($v_1, \langle p_x \rangle$), Elliptic flow (v_2)

→ 衝突初期に作られ、高密度の状態方程式 (EOS) に敏感



高エネルギー重イオン衝突の物理

■ クオーク・グルーオン・プラズマ (QGP) の発見・物性

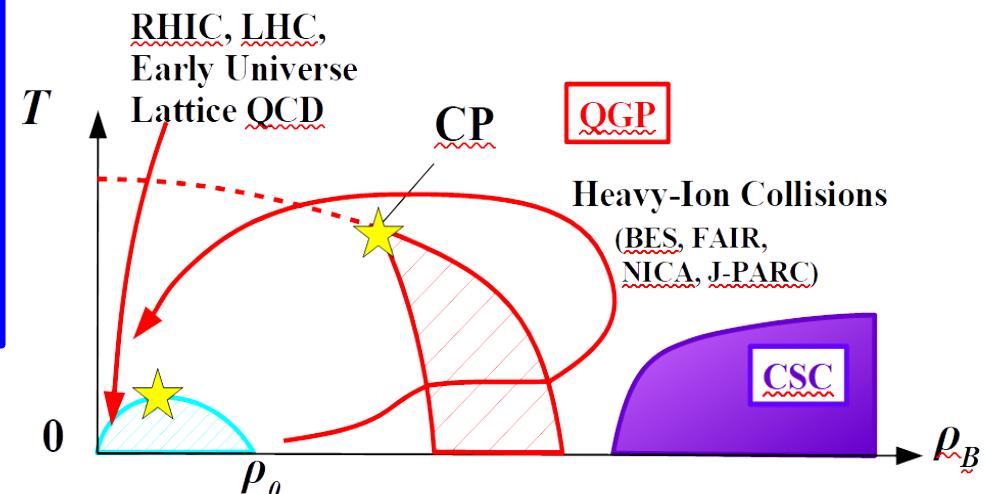
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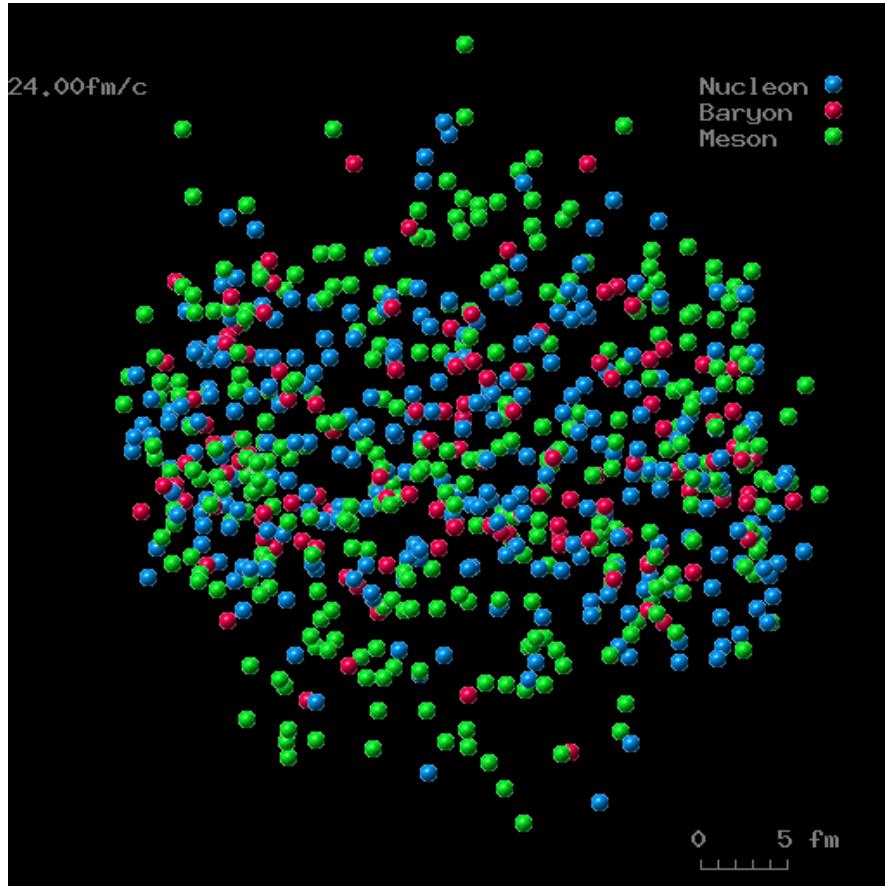
■ ハドロン輸送模型に基づく研究

- ハドロン輸送模型 JAM
- EOS の導入とフロー
- 負の側方フローと EOS の軟化

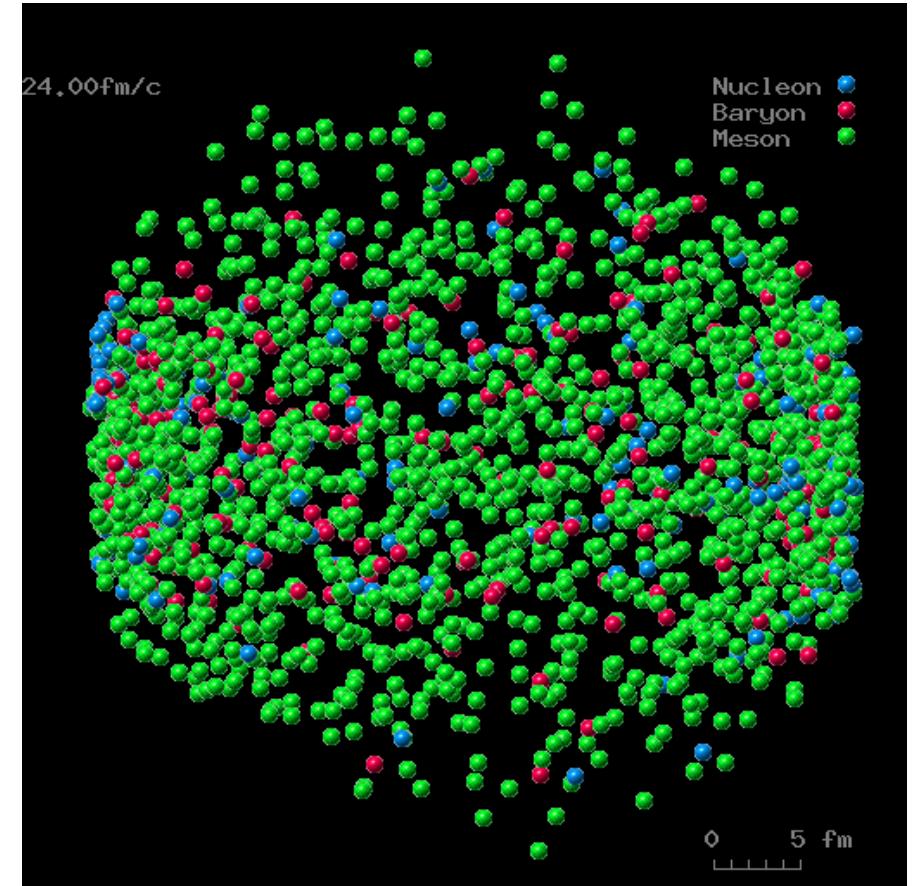


How do heavy-ion collisions look like ?

Au+Au, 10.6 A GeV



Pb+Pb, 158 A GeV



$$\sqrt{s_{NN}} \sim 5 \text{ GeV}$$

$$\sqrt{s_{NN}} \sim 20 \text{ GeV}$$

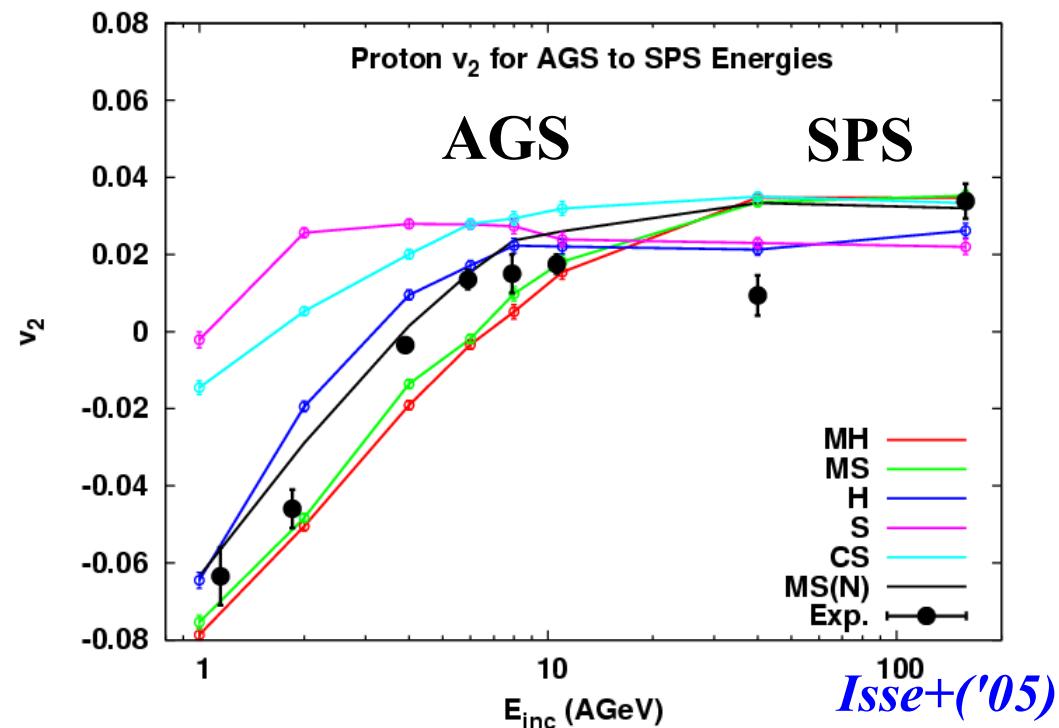
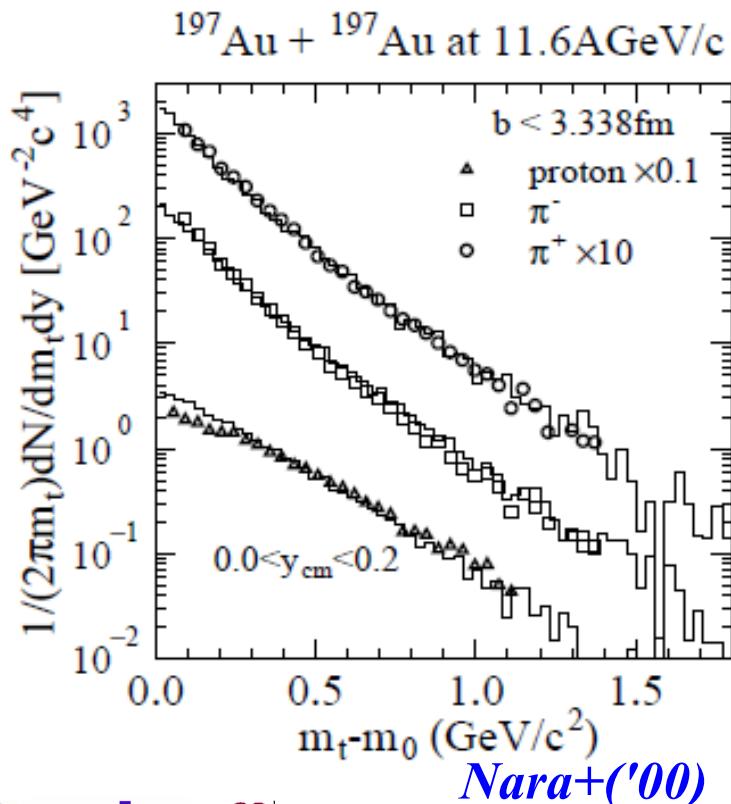
JAMming on the Web <http://www.jcprg.org/jow/>

ハドロン輸送模型 JAM

*Y.Nara, N.Otuka, AO, K.Niita, S.Chiba, PRC61('00), 024901
M.Isse, AO, N.Otuka, P.K.Sahu, Y.Nara, PRC72 ('05)064908*

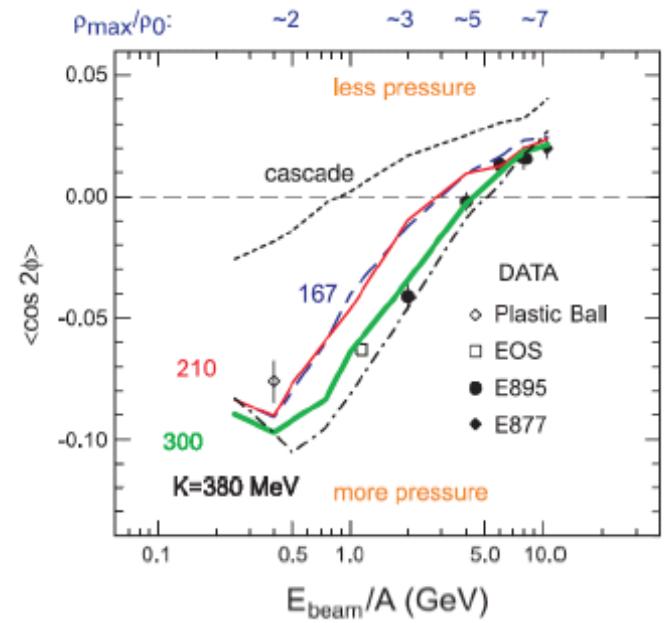
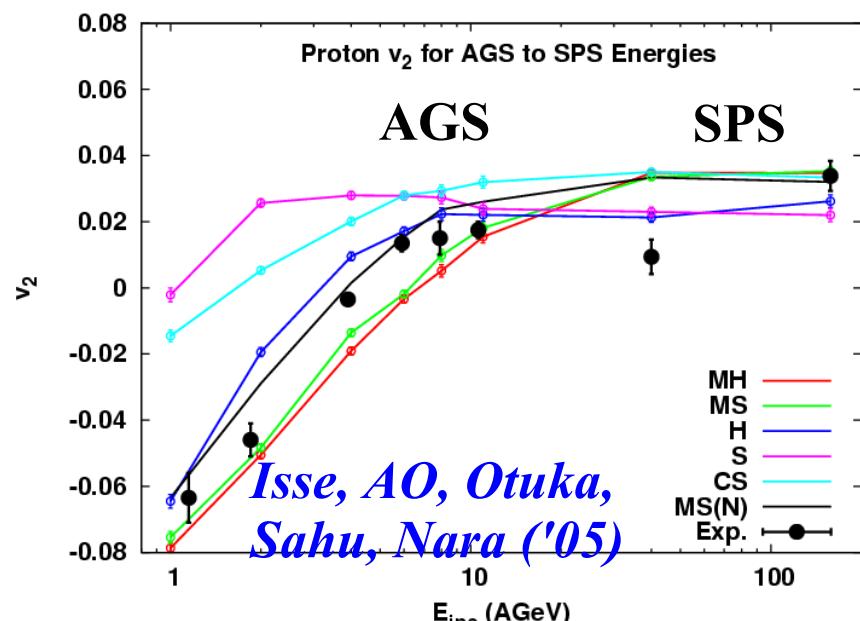
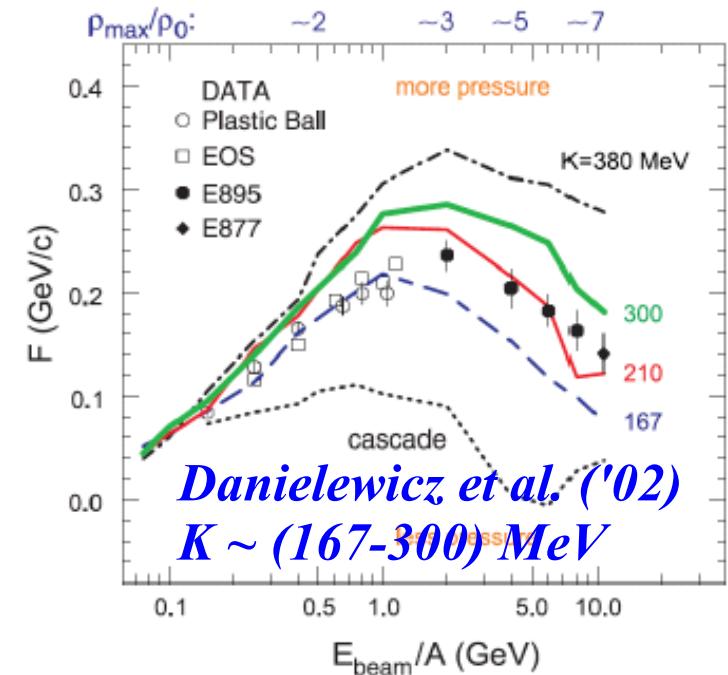
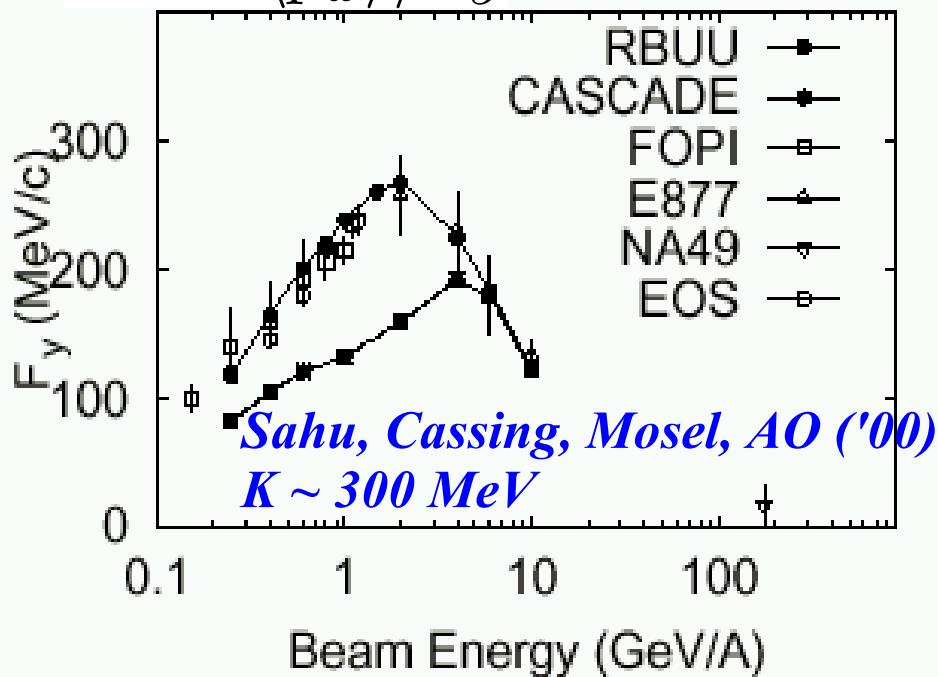
■ Event Generator: Jet AA Microscopic transport model (JAM)

- 多くの自由度・素過程を取り入れた輸送模型
- 平均場を導入して $E_{\text{inc}} = (1-158) \text{ AGeV}$ でのフローをほぼ説明
- しかし $\sqrt{s_{\text{NN}}} = 11.5 \text{ GeV}$ での負のフローは説明できない



集団フローから状態方程式へ

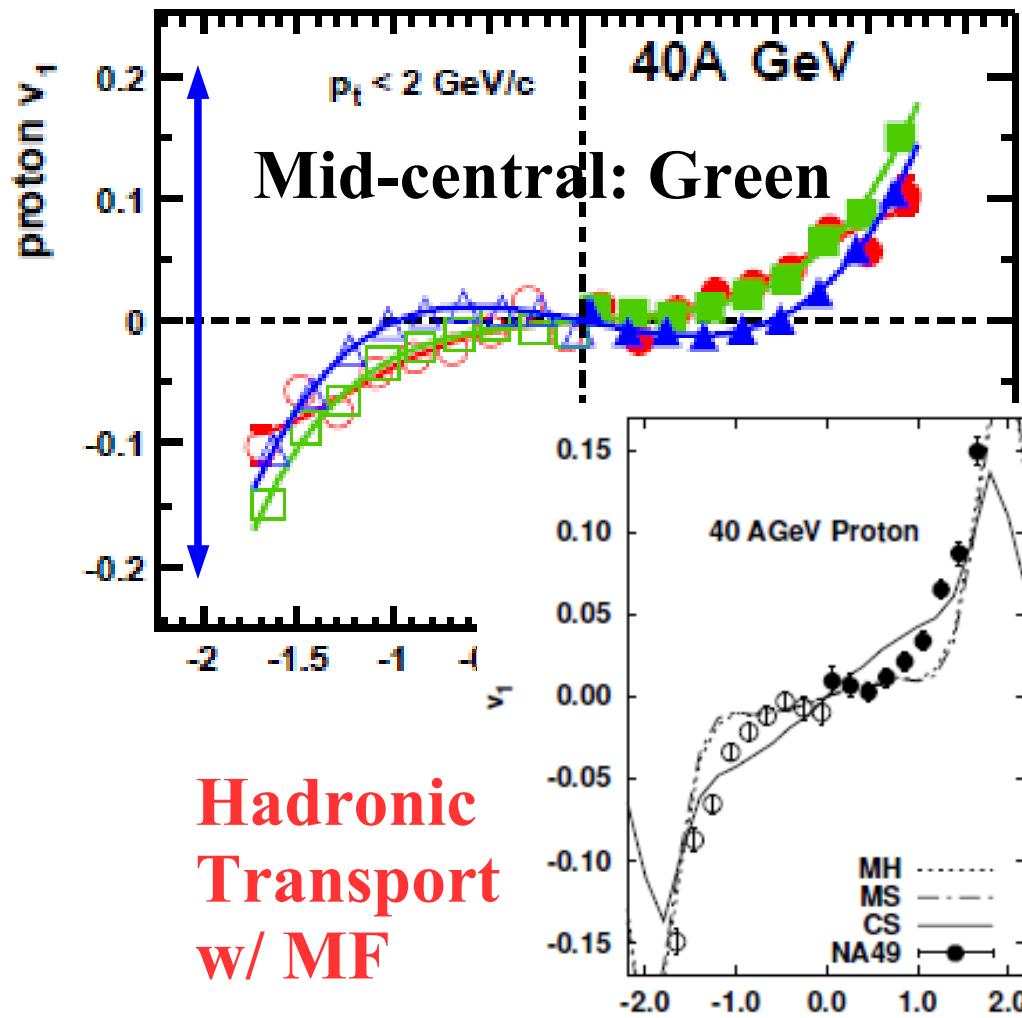
$$F = d\langle p_x \rangle / dy$$



SPS(NA49) vs RHIC(STAR)

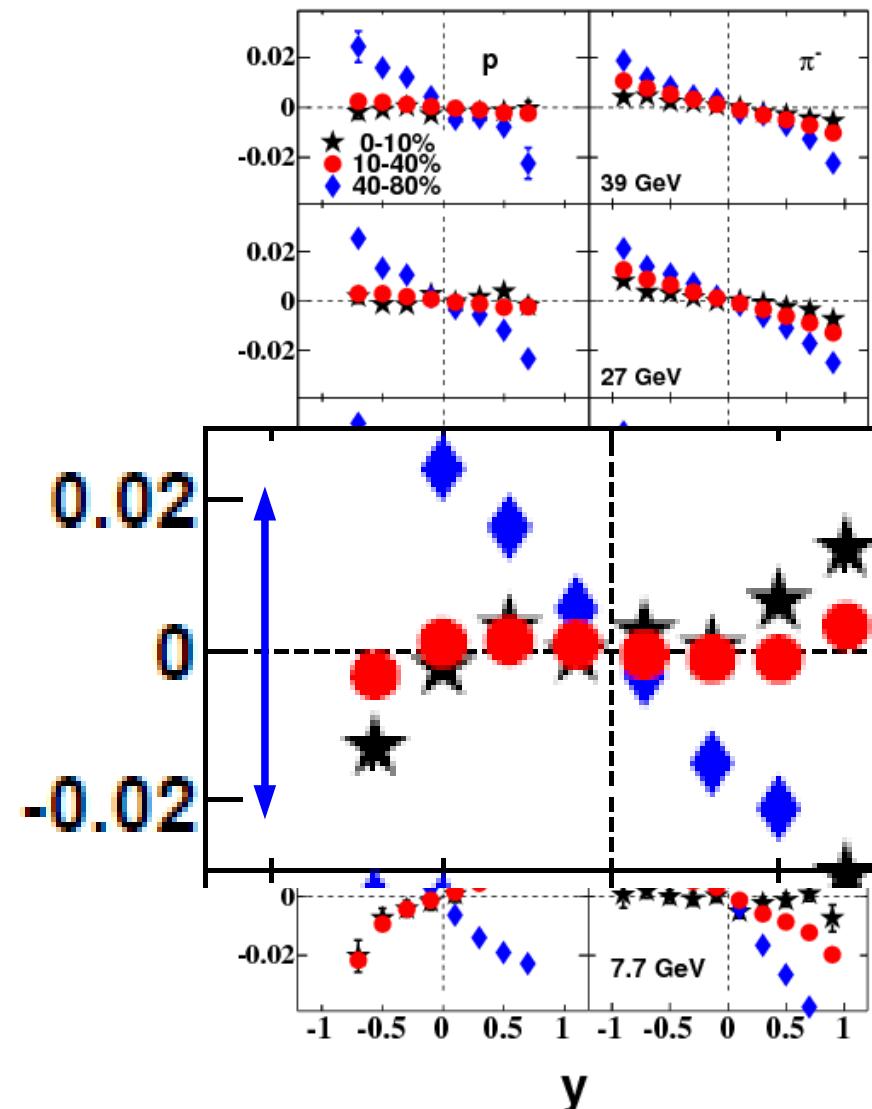
■ SPS (NA49), $\sqrt{s}_{\text{NN}} = 8.9 \text{ GeV}$

C. Alt et al. (NA49), PRC68 ('03) 034903



*M. Isse, A.O., N. Otuka, P.K. Sahu, Y. Nara,
PRC72 ('05) 064908*

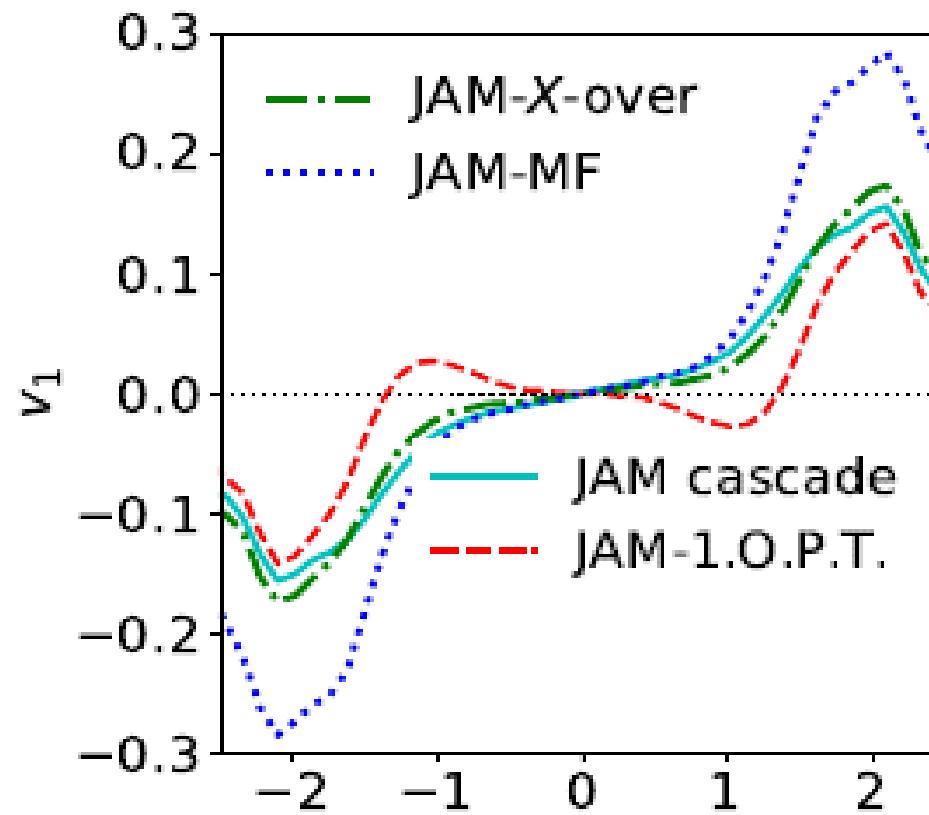
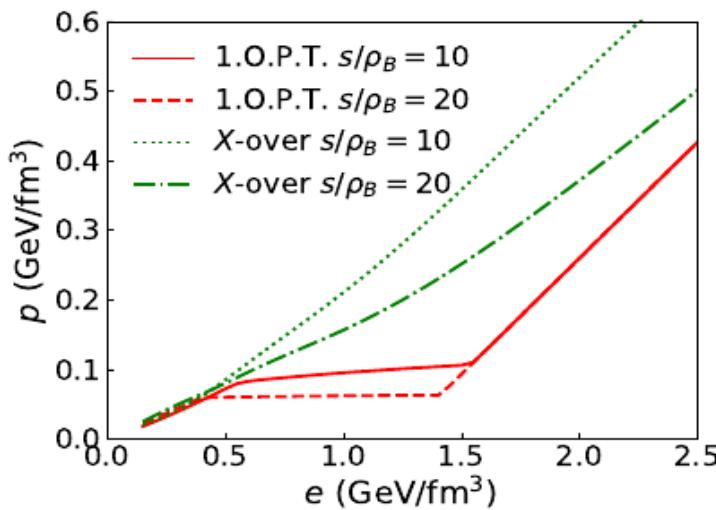
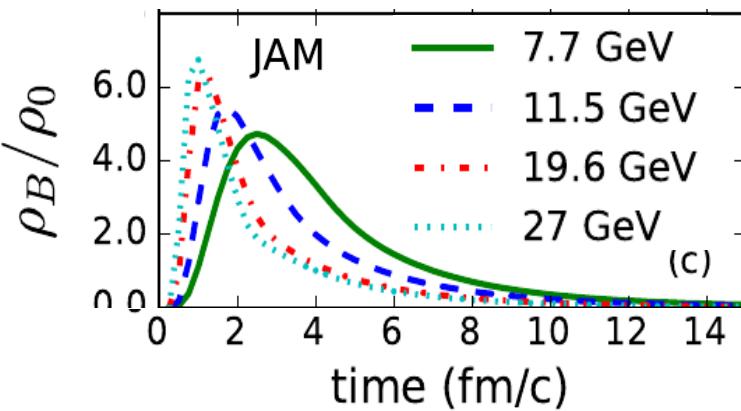
■ RHIC(STAR), 7.7-39 GeV



*L. Adamczyk et al. (STAR),
PRL 112(2014)162301*

負のフローとEOSの軟化

- ビリアル定理を使って任意の EOS を取り込めるように理論を拡張
- $\sqrt{s_{NN}} = 11.5 \text{ GeV}$ で見られる負のフロー ($dv_1/dy < 0$)
 → $(5\text{-}10)\rho_0$ において急激な EOS の軟化あれば説明可能



*Y.Nara, H.Niemi, AO, H.Stoecker, PRC94('16)034906.
 Y. Nara, H. Niemi, AO, J. Steinheimer, X.-F. Luo,
 H. Stoecker, EPJA 54 ('18)18*

ビリアル定理

■ Virial

$$G = \sum_i \mathbf{p}_i \cdot \mathbf{r}_i$$

$$\rightarrow \frac{dG}{dt} = \sum_i \mathbf{p}_i \cdot \mathbf{v}_i - \sum_i \nabla_i U \cdot \mathbf{r}_i + \boxed{\frac{1}{\Delta t} \sum_{\text{collision}} \mathbf{q}_i \cdot (\mathbf{r}_i - \mathbf{r}_j)} = 3VP$$

Kinetic Potential Pressure from Collisions

■ Attractive / Repulsive Orbit Scatterings

- 通常は散乱角はランダム → 衝突項の圧力への影響はゼロ
- Attractive orbits → $\Delta P < 0$ (softening)
- Repulsive orbits → $\Delta P > 0$ (hardening)

■ Boltzmann Eq. simulating a given EOS

$P > P(\epsilon) \rightarrow$ Attractive orbit, $P < P(\epsilon) \rightarrow$ Repulsive orbit

衝突が十分に頻繁であれば、ボルツマン方程式だけでポテンシャル効果をシミュレートできる！

$(5-10)\rho_0$ で QCD 相転移がありそう。
軟化が必要なことから（対称核物質では）
1 次相転移が想定される！

→ 輸送模型と流体力学を組み合わせた
理論開発が求められる。

レポート(5)

- 全部で 5-7 問程度出します。3 問程度以上レポートを出してください。レポート(5) の〆切は 1 月末
- (Report 6) Wigner 変換の性質

$$([A, B])_W = i\hbar \{ A_W, B_W \}_{\text{PB}} + \mathcal{O}(\hbar^3)$$

を示せ。

(Ring-Schuck text の Appendix 参照。面倒です。)

- (Report 7) 分布関数 f を用いてフェルミオン系のエントロピーを定義し、フェルミ統計を含む衝突項による時間発展において、エントロピーが増加することを示せ。

Backups

Collective Flows at AGS and SPS Energies

Collective Flow and EOS: Old Problem ?

- 1970's-1980's: First Suggestions and Measurement
 - Hydrodynamics suggested the Existence of Flow.
 - Strong Collective Flow suggests Hard EOS
- 1980's-1990's: Deeper Discussions in Wider E_{inc} Range
 - Momentum Dep. Pot. can generate Strong Flows.
 - E_{inc} deps. implies the importance of Momentum Deps.
 - Flow Measurement up to AGS Energies.
- 2000's: Extension to SPS and RHIC Energies
 - EOS is determined with Mom. AND Density Dep. Pot. ?

Old but New (Continuing) Problem !

What is Collective Flow ?

(Directed) Flow (dP_X/dY)

Stiffness (Low E)
+ Time Scale (High E)

Elliptic Flow (V_2)

Thermalization
& Pressure Gradient

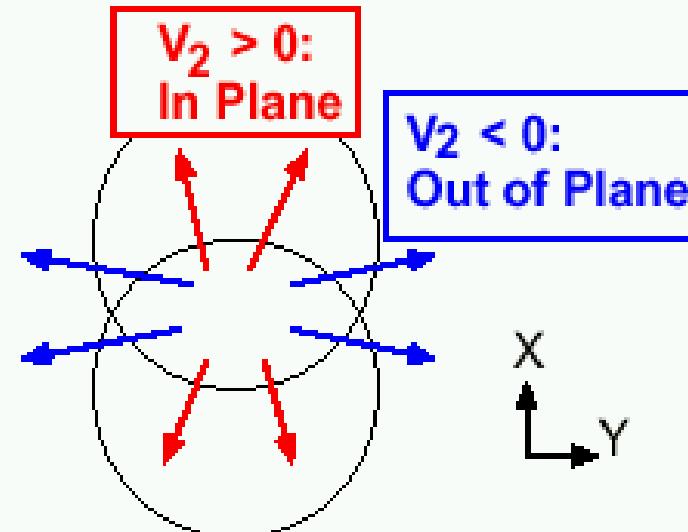
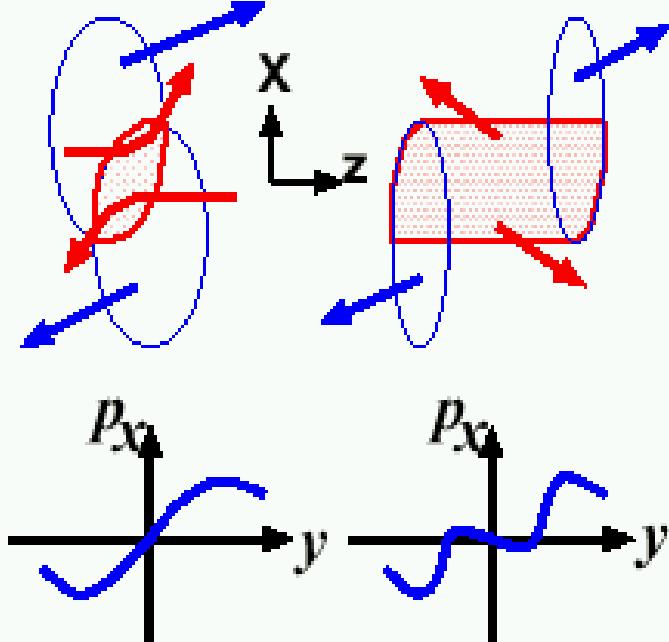
Radial Flow (β_T)

Pressure History

$$\epsilon \frac{D\mathbf{V}}{Dt} = -\nabla P$$

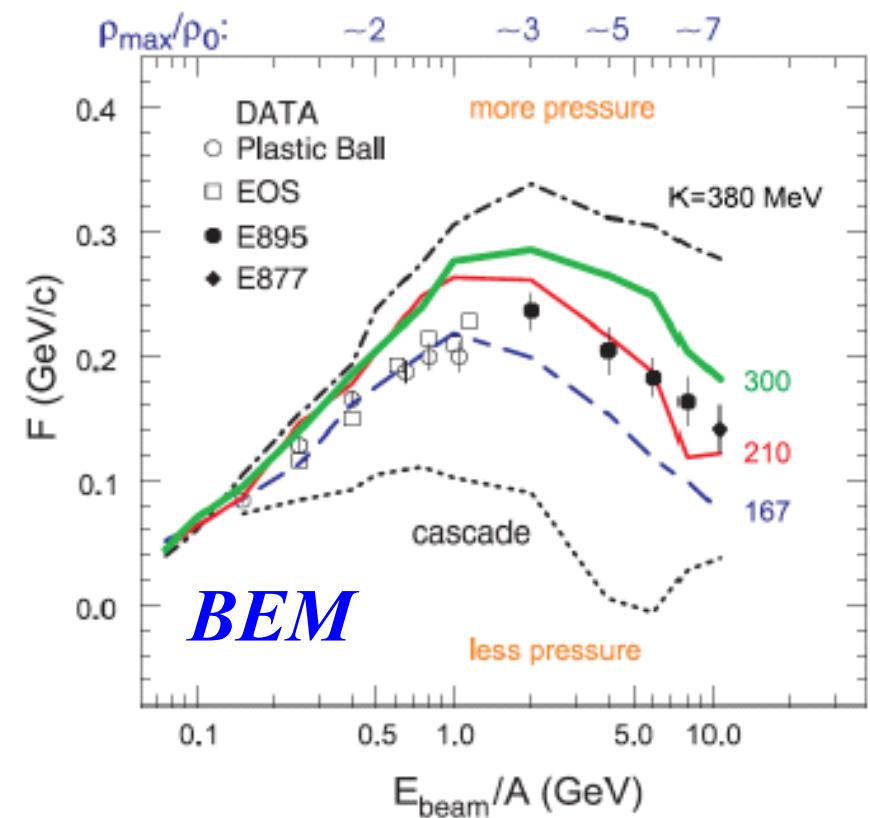
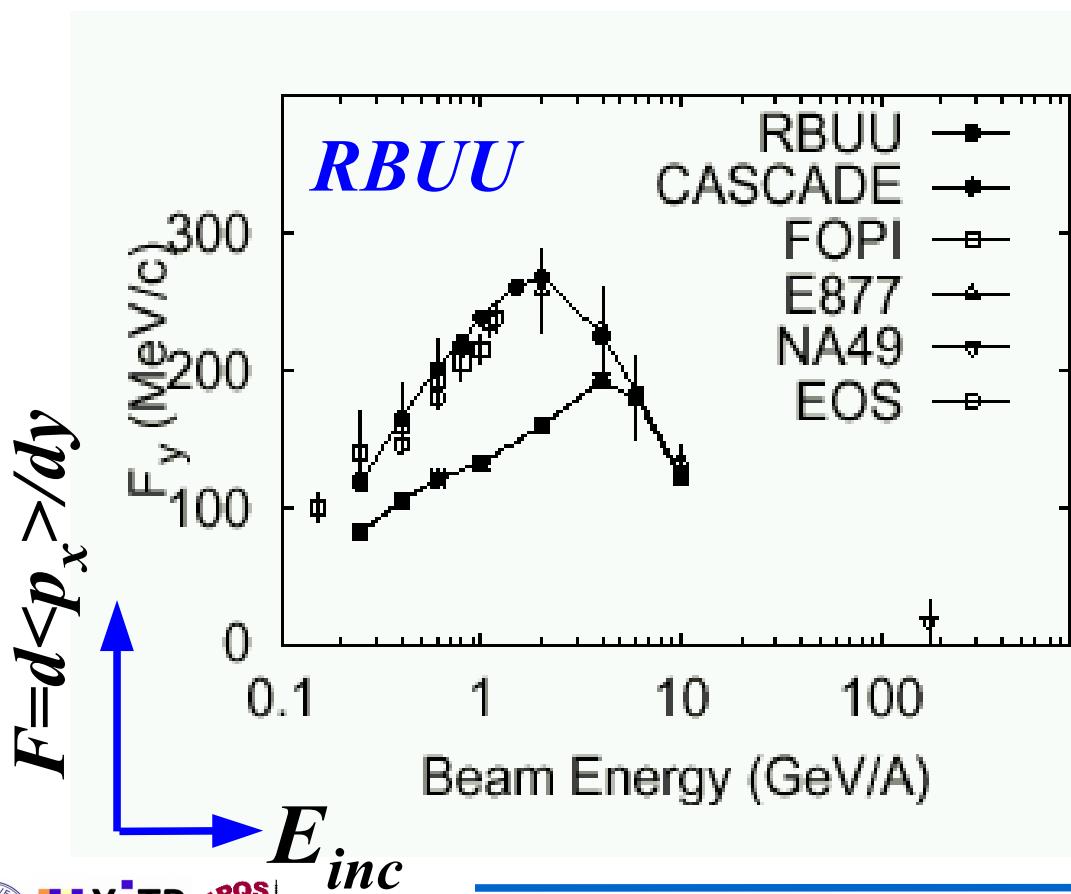
$$\rightarrow V = \int_{path} \frac{-\nabla P dt}{\epsilon}$$

Until AGS Above SPS



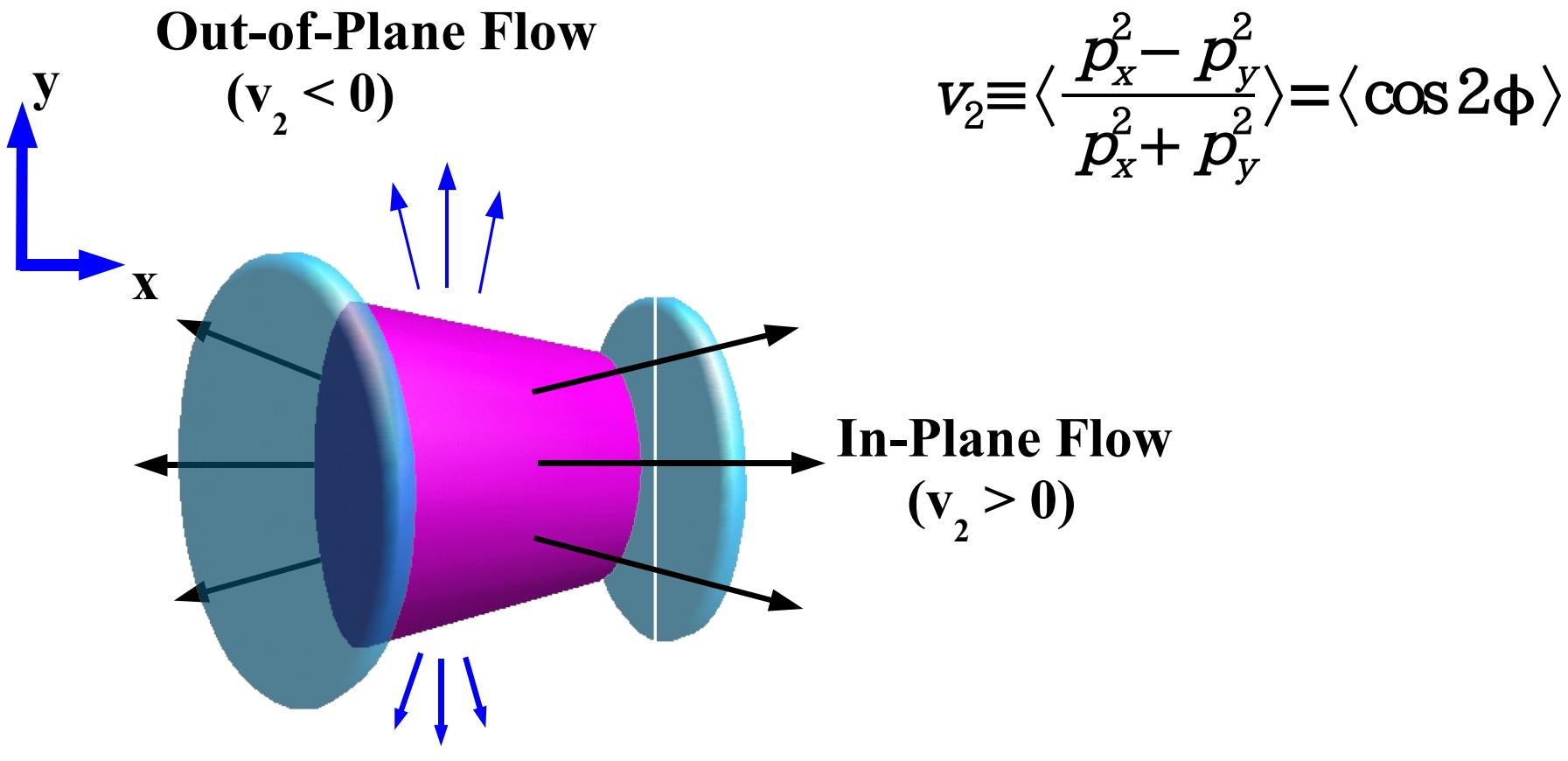
Side Flow at AGS Energies

- Relativistic BUU (RBUU) model: $K \sim 300 \text{ MeV}$
(Sahu, Cassing, Mosel, AO, Nucl. Phys. A672 (2000), 376.)
- Boltzmann Equation Model (BEM): $K=167\sim210 \text{ MeV}$
(P. Danielewicz, R. Lacey, W.G. Lynch, Science 298(2002), 1592.)



Elliptic Flow

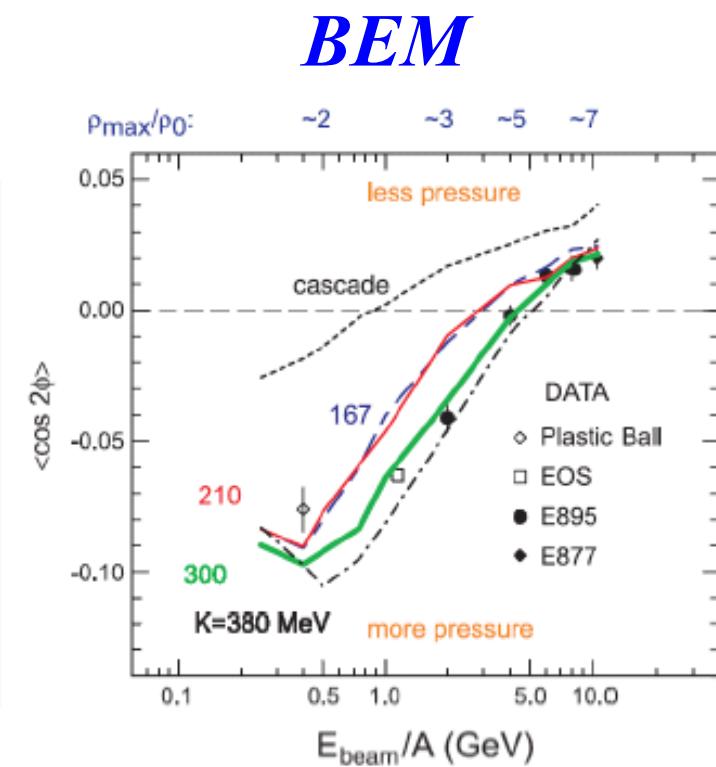
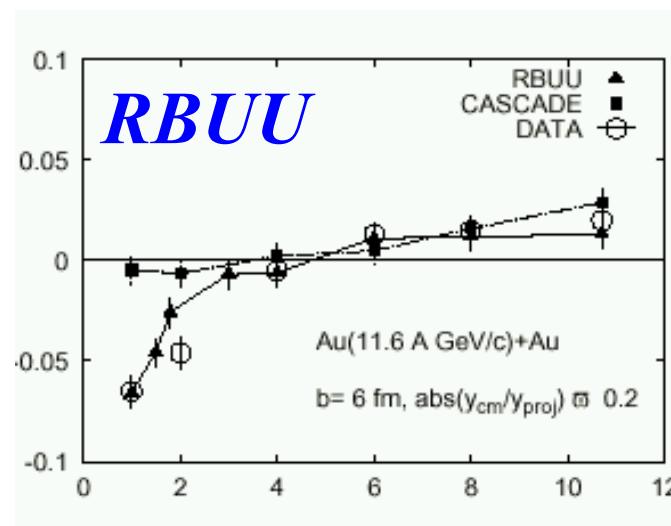
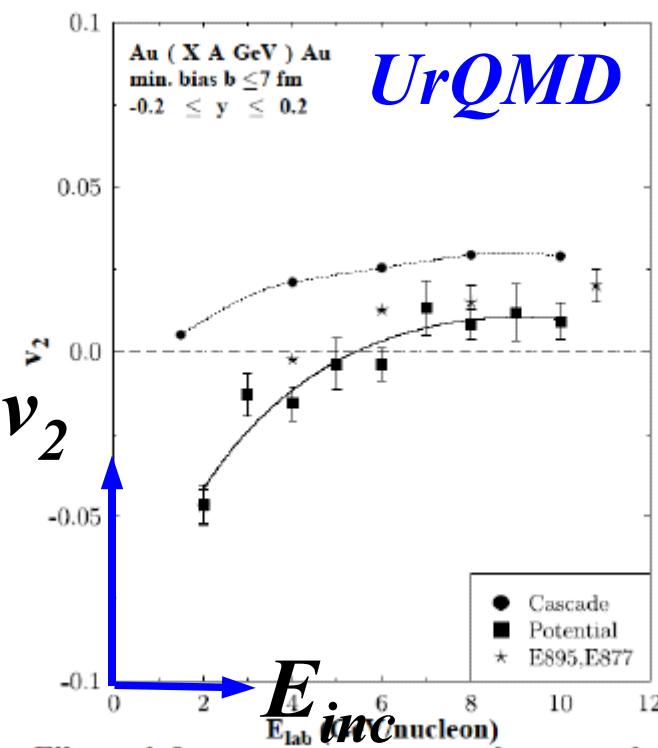
- What is Elliptic Flow ? → Anisotropy in P space
- Hydrodynamical Picture
 - Sensitive to the Pressure Anisotropy in the Early Stage
 - Early Thermalization is Required for Large V2



Elliptic Flow at AGS

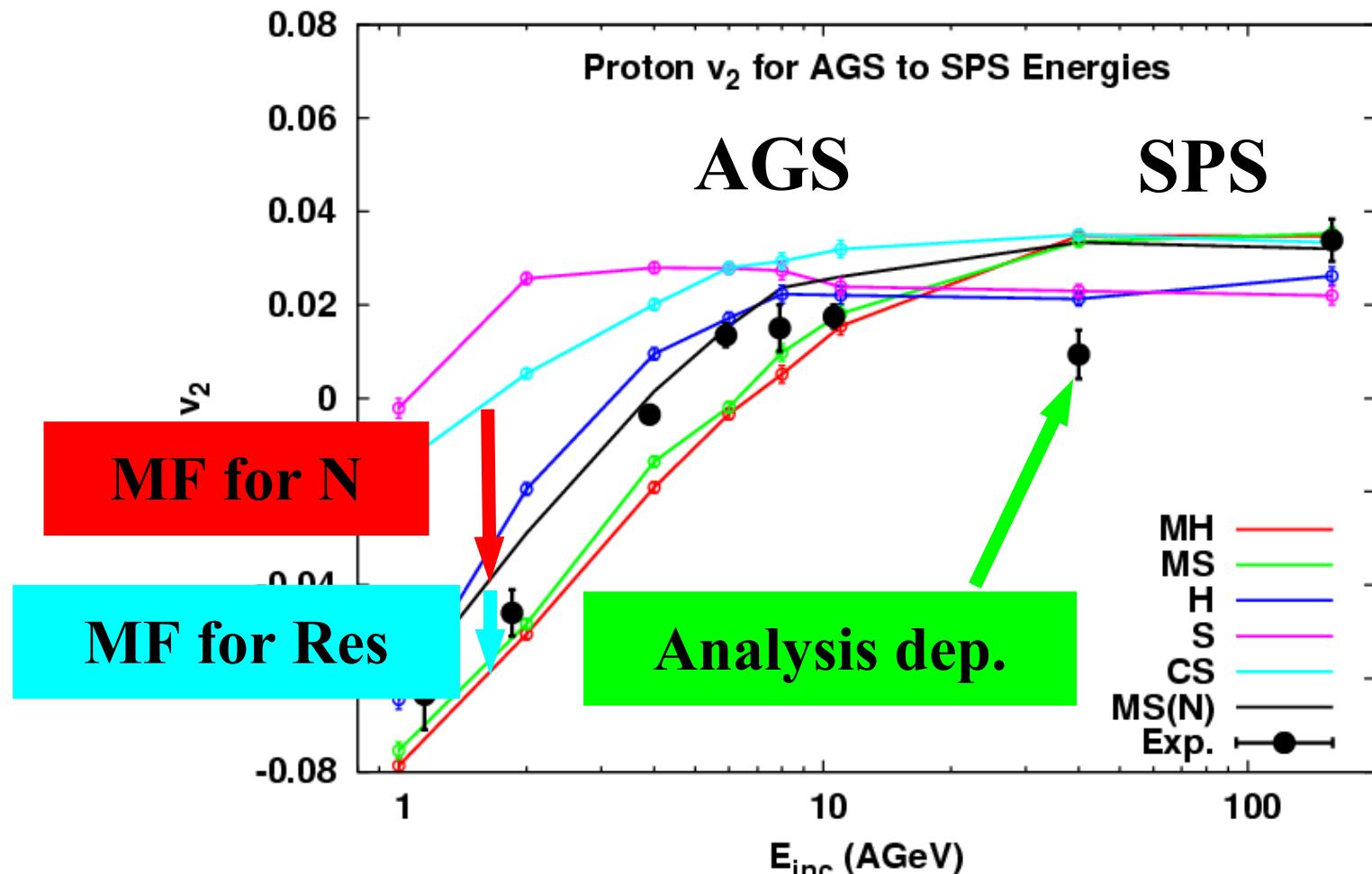
Strong Squeezing Effects at low E (2-4 A GeV)

- UrQMD: Hard EOS (S.Soff et al., nucl-th/9903061)
- RBUU (Sahu-Cassing-Mosel-AO, 2000): $K \sim 300$ MeV
- BEM(Danielewicz2002): $K = 167 \rightarrow 300$ MeV



Elliptic Flow from AGS to SPS

- JAM-MF with p dep. MF explains proton v2 at 1-158 A GeV
 - v2 is not very sensitive to K (incompressibility)
 - Data lies between MS(B) and MS(N)

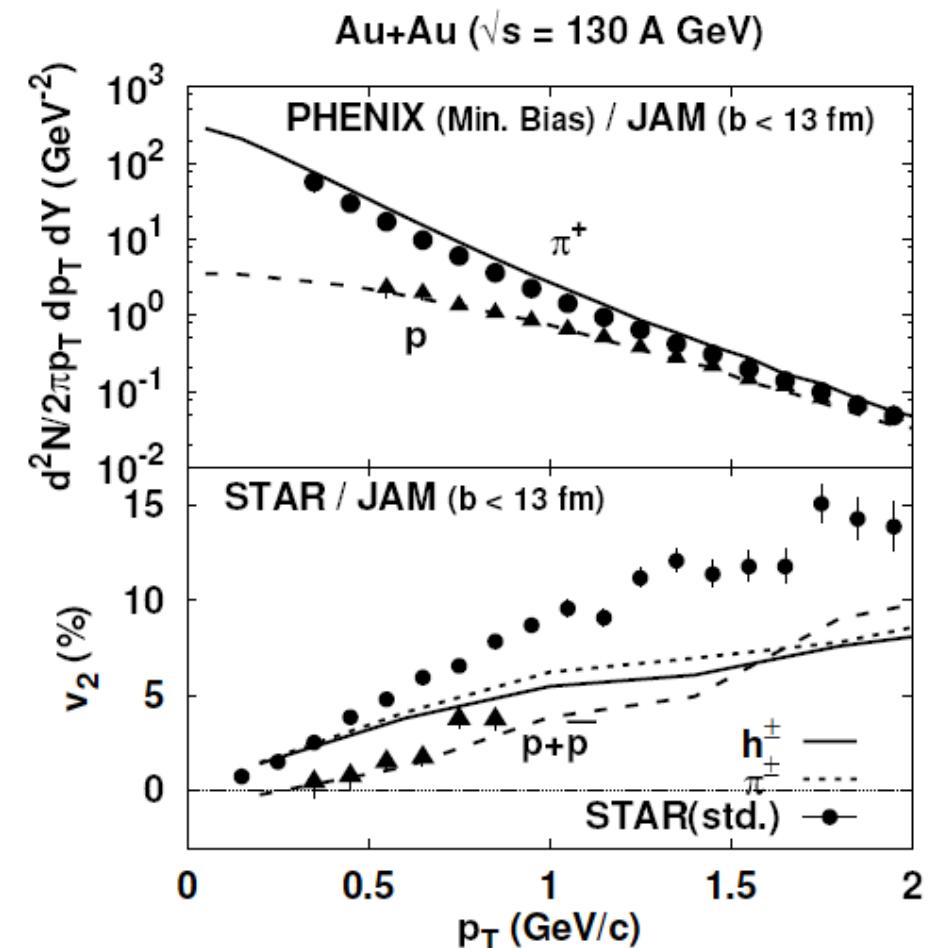
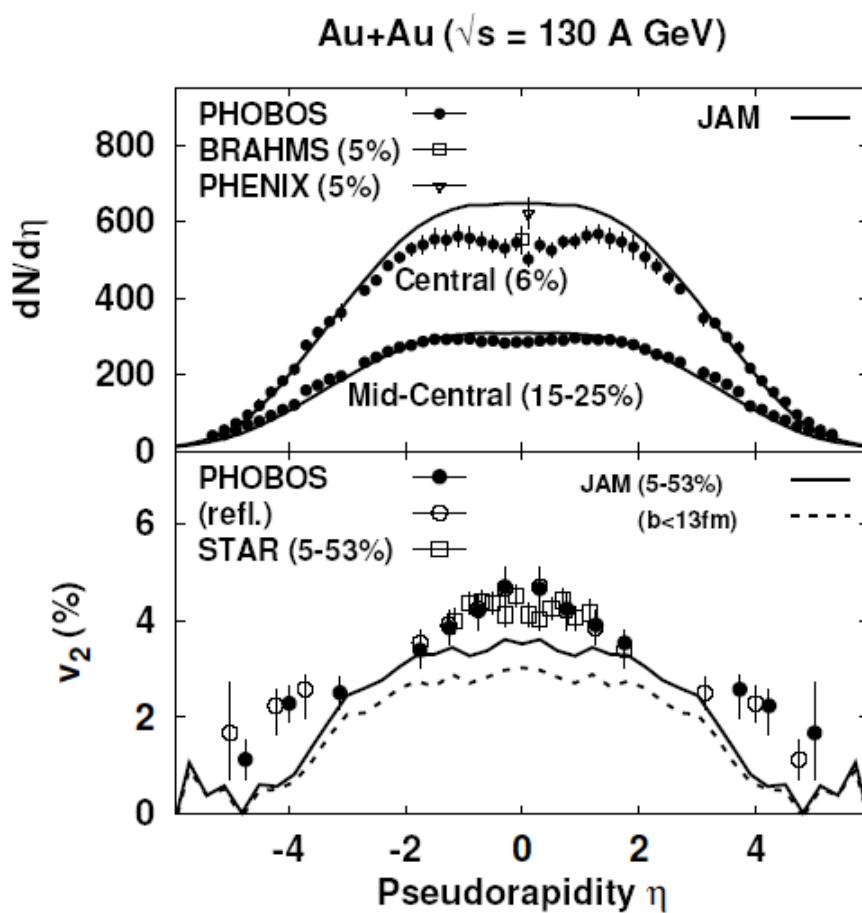


Elliptic Flow @ RHIC

Elliptic Flow in Hadron-String Cascade (I)

■ Hadron-String Cascade (JAM) @ RHIC

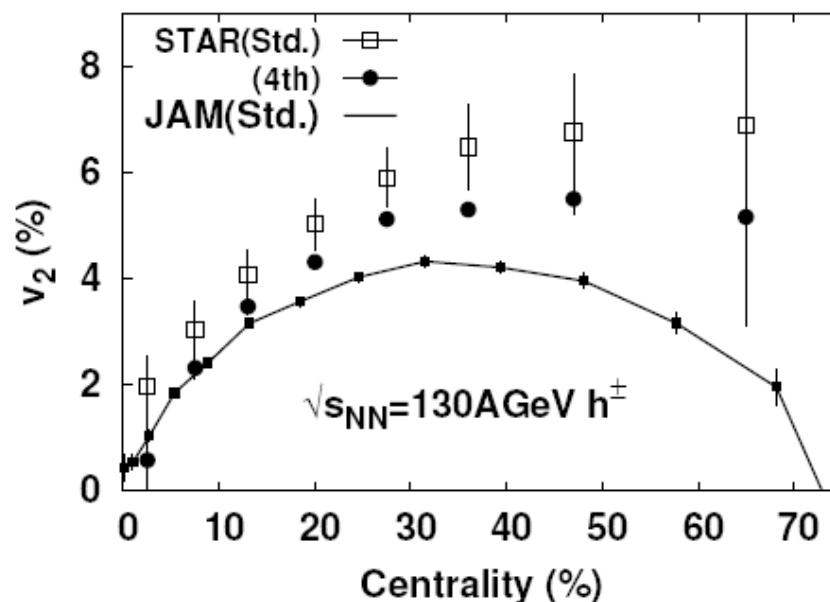
- Hadron Yield is reasonably explained up to 2 GeV/c (10-20 % error)
- v_2 is underestimated (20-30 % (integrated), 50 % ($p_T > 1$ GeV))



Elliptic Flow in Hadron-String Cascade (II)

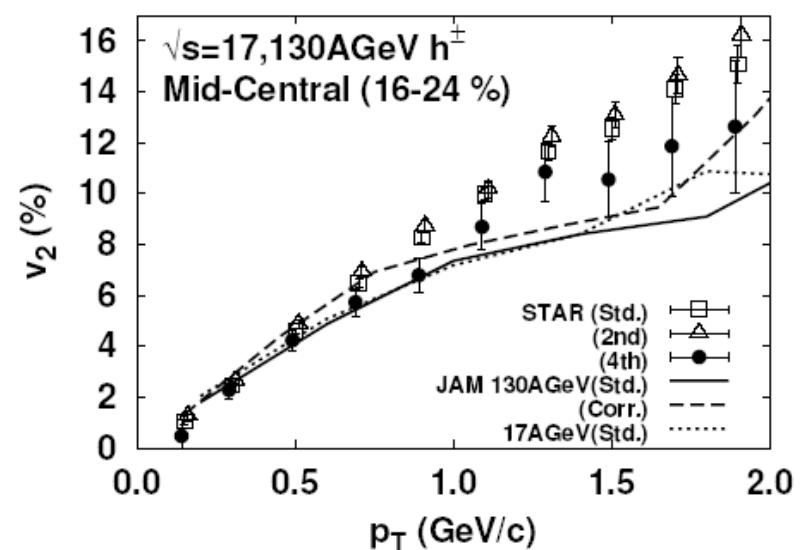
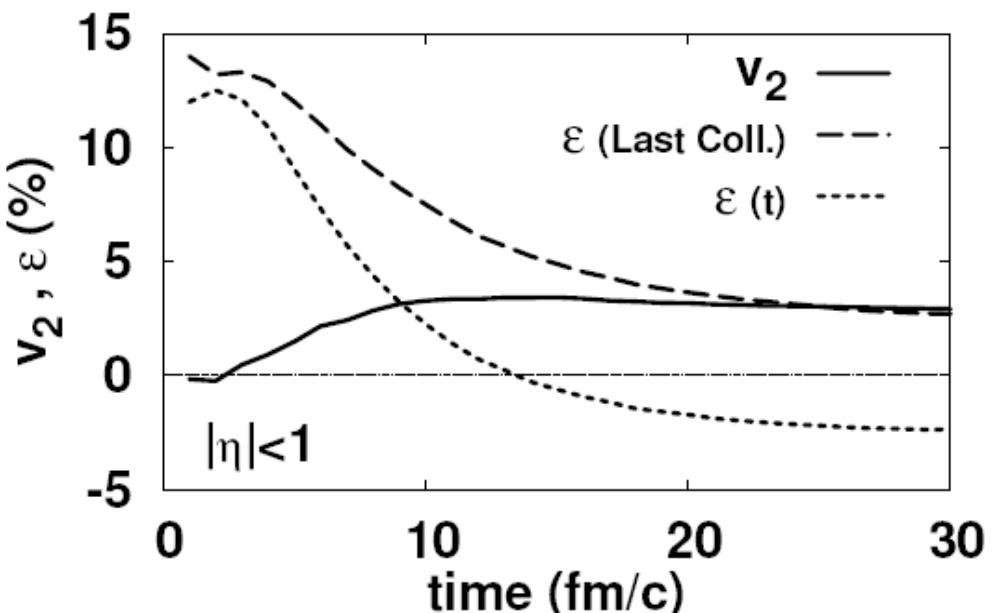
■ Why do we underestimate v_2 in Hadron-String Cascade ?

- v_2 growth time is long (~ 10 fm/c), due to hadron formation time ($\tau \sim 1$ fm/c).
→ much longer than hydro



Sahu-Isse-AO-Otuka-Phatak 2006

Au+Au, $\sqrt{s_{NN}} = 130$ GeV, $b < 13$ fm



Results of Parton Cascade

- Unexpectedly high parton cross sections of $\sigma = 5\text{-}6 \text{ mb}$ have to be assumed in parton cascades in order to reproduce the elliptic flow.

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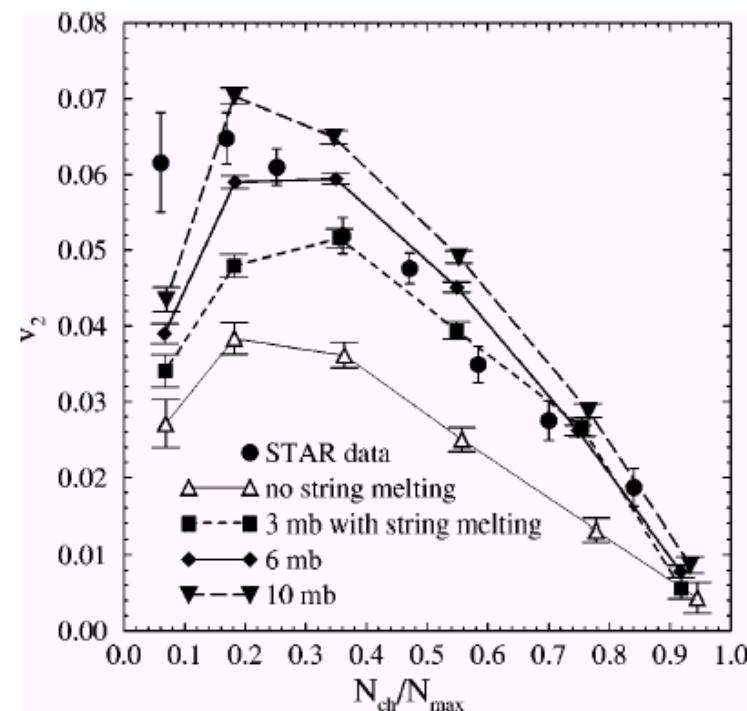


FIG. 3. Impact parameter dependence of elliptic flow at 130 A GeV. The data from the STAR collaboration [7] are shown by filled circles, while the theoretical results for different partonic dynamics are given by curves.

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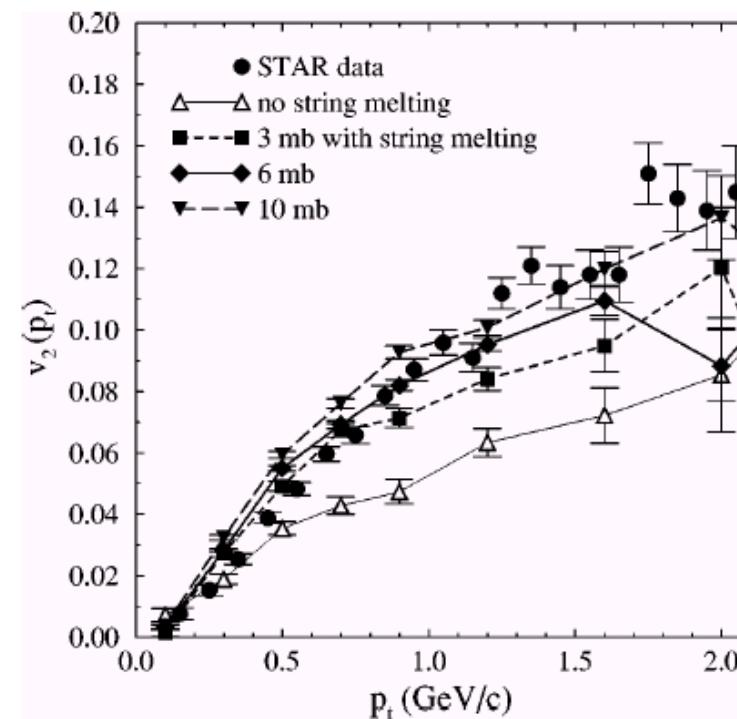
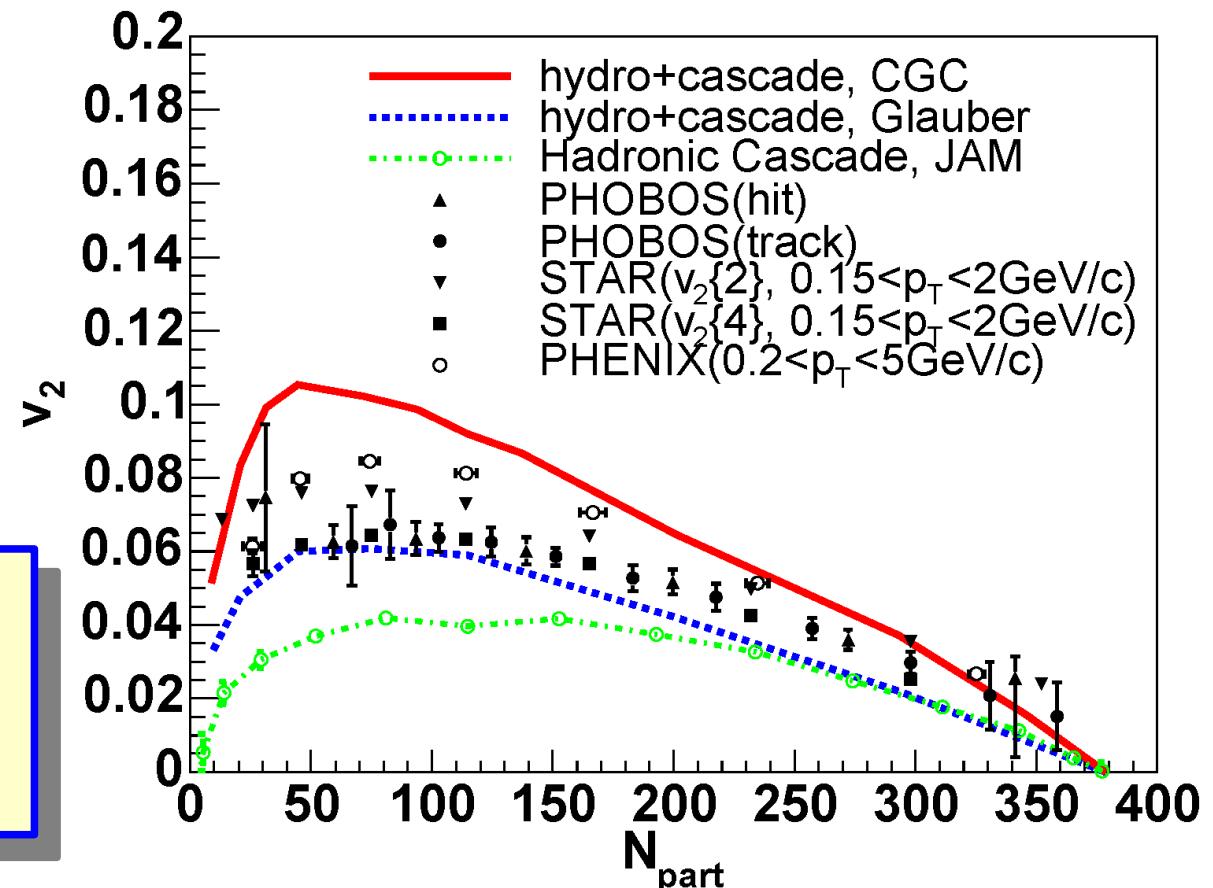


FIG. 4. Transverse momentum dependence of elliptic flow at 130 A GeV. Circles are the STAR data for minimum-bias Au+Au collisions [7], and curves represent the minimum-bias results for charged particles within $\eta \in (-1.3, 1.3)$ from the AMPT model.

Cascade vs Hydro @ RHIC: Au+Au

- Comparison of v_2 as a function of N_{part}
 - Cascade predict smaller v_2 in peripheral collisions
 - Data lies between hydro results with two different initial condition CGC (Color Glass Condensate) and Glauber type initial condition.



*Hydro is better,
CGC may be realized
in central collisions.*