

Physics of Hadronic Matter
Probed by High Energy Nuclear Collisions

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**Hierarchies in Nuclear Physics
and Physics of Hadronic Matter**

Theoretical Preparation:

Microscopic Transport Model and Statistical Models

Fragment Formation and Liquid-Gas Phase Transition

Roles of Strangeness in Dense Matter

High Energy Heavy-Ion Collision

and Hot/Dense Hadronic Matter

Summary

Hierarchies in Nuclear Physics

Quarks and Gluons (QCD)

Nucleons and Hadrons (NN Interaction, Effective Lagrangian, ...)

Finite Nuclei (Effective NN Interaction, Model Space, ...)

Nuclear/Hadronic Matter

Nuclear Physics = Physics of Four (or Three) Hierarchies

Physics of Nuclear/Hadronic Matter

*** Two-Fold Structure:**

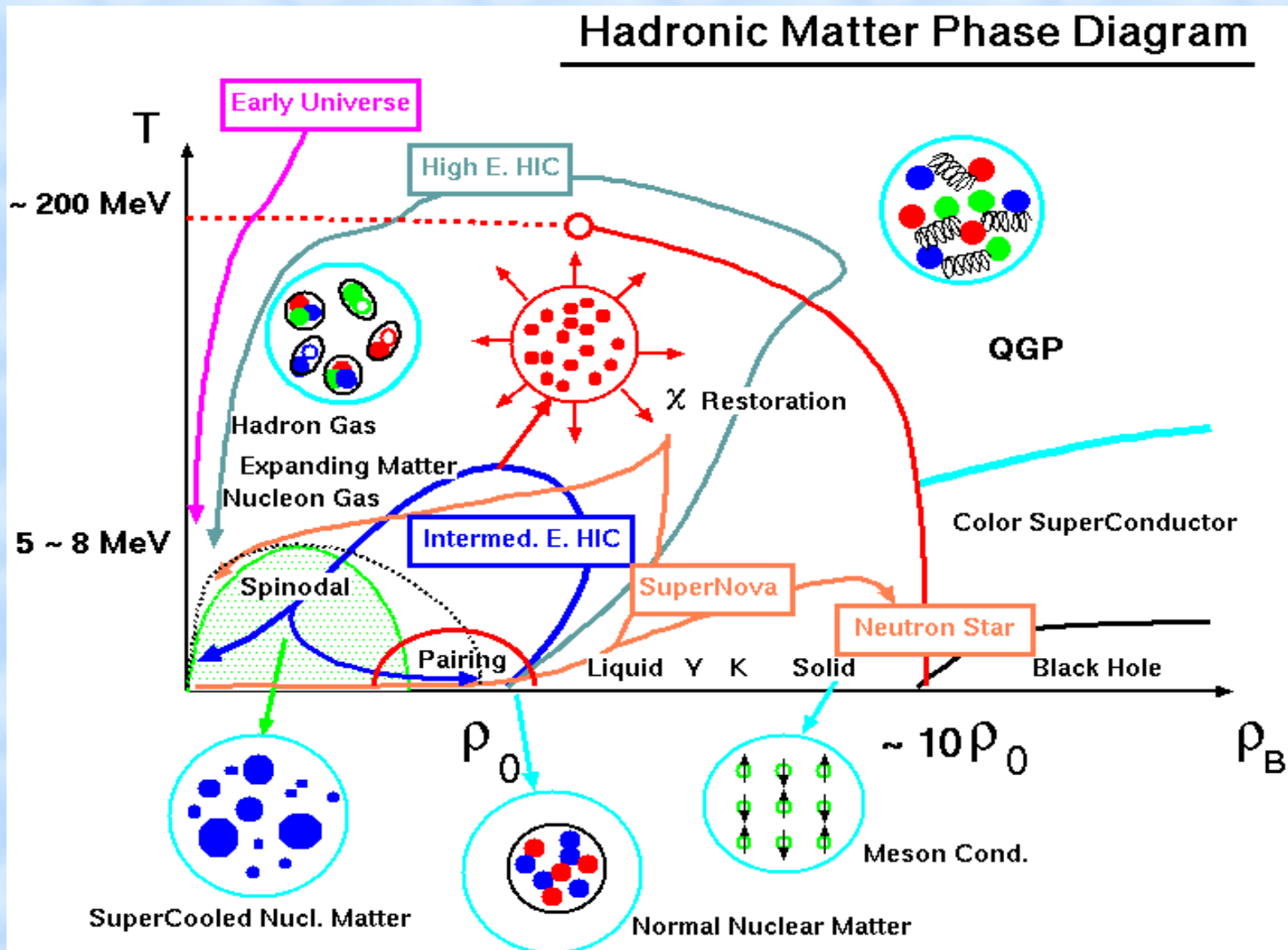
Quark & Gluon \leftrightarrow Nucleon/Hadron \leftrightarrow Nucleus

*** Relation to Astrophysical Objects/Phenomena:**

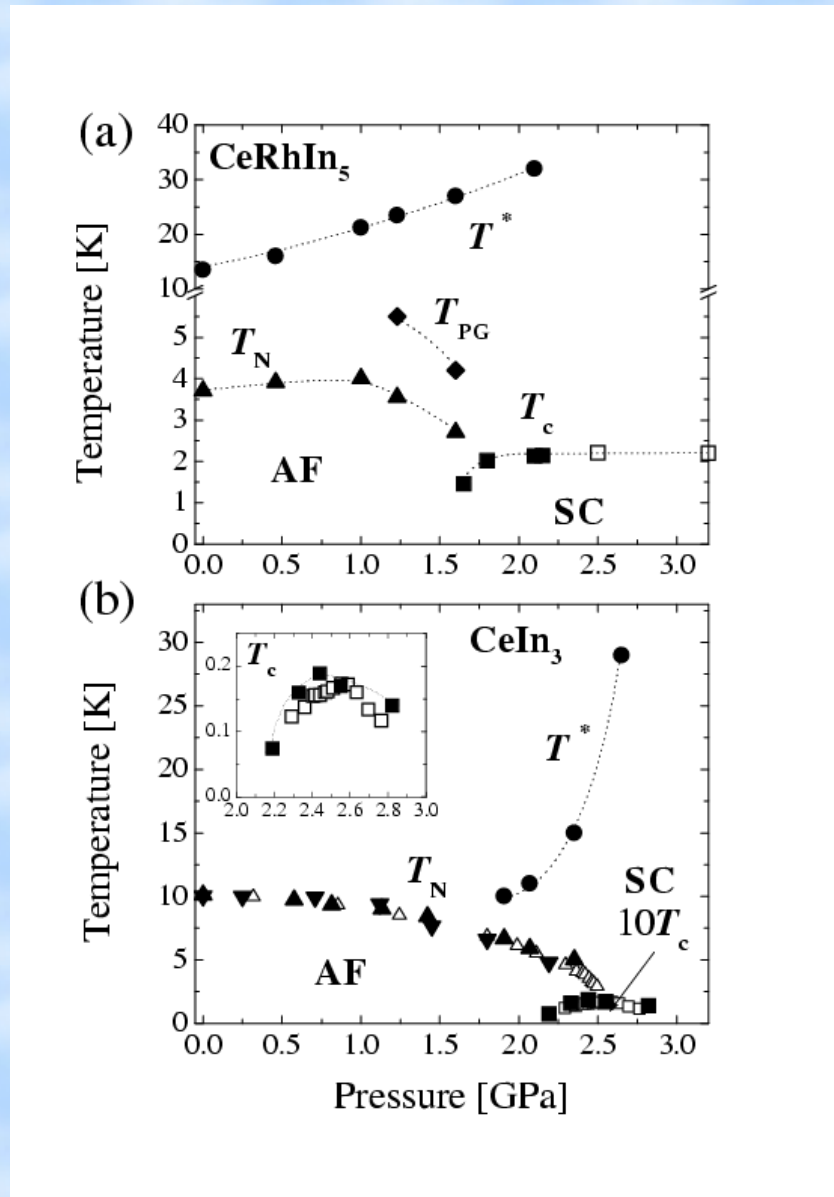
Early Universe, Compact Objects

*** Similarities to Superconductor in Solid State**

Hadronic Matter Phase Diagram



Phase Diagram of Superconductor $CeRhIn_5$ and $CeIn_3$

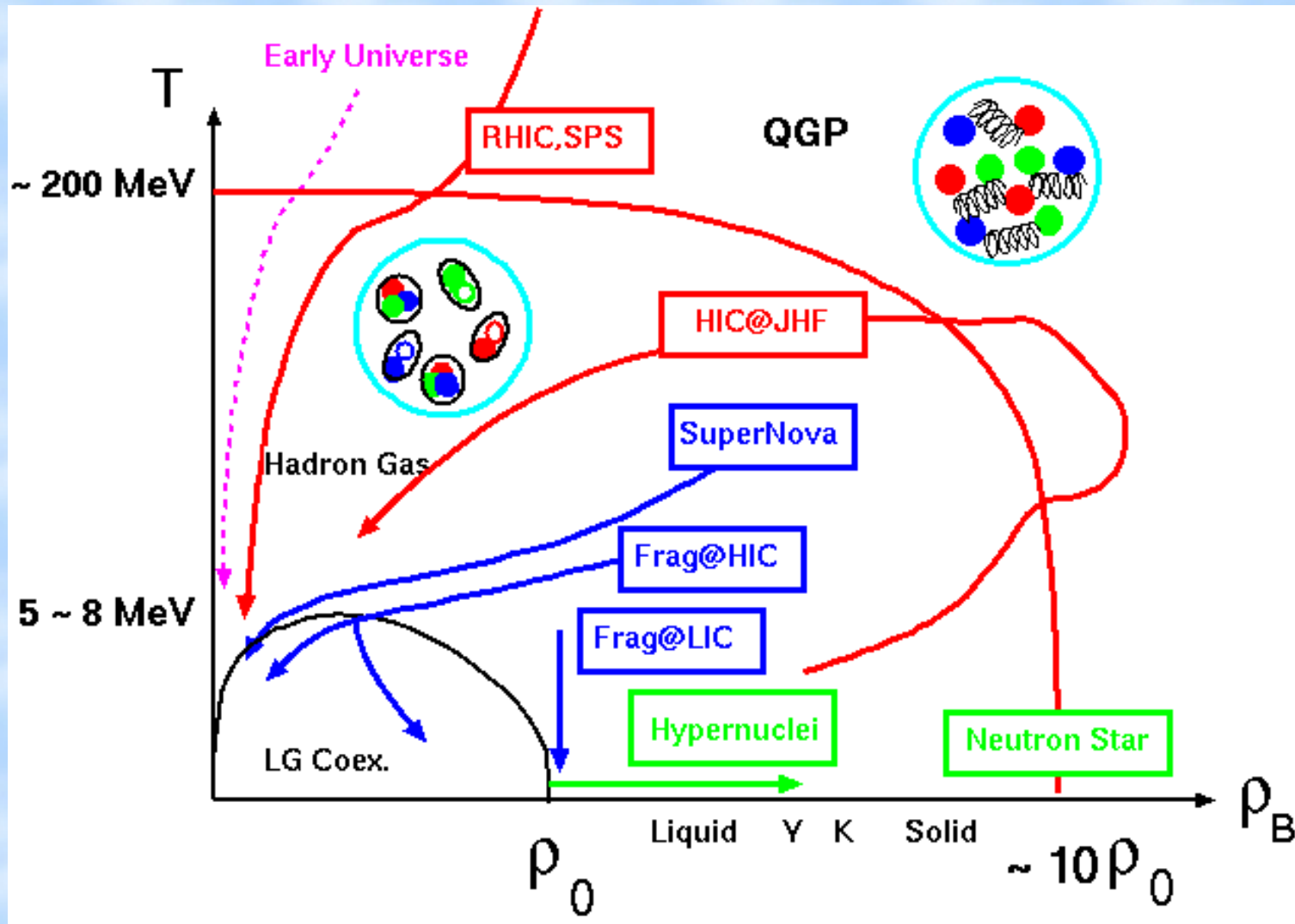


(Kawasaki et al., cond -mat/0110620.)

Approaches for Hadronic Matter Study

Constructive Way: Quarks and Gluons → Hadrons → Matter

My Individual Problems: Phenomenology of Each Region



Theoretical Preparation

Mean Field Evolution

TDHF – Vlasov – BUU, AMD

Collision Term

Resonance, String and Jet

Statistical Model

Statistical Model of Hadrons and Fragments

TDHF and Wigner Transformation

Mean Field Theory (e.g., TDHF)

$$i \hbar \frac{\partial \phi}{\partial t} = h \phi, \quad \rho(r, r') = \sum_i^{Occ} \phi_i(r) \phi_i^*(r'), \quad i \hbar \frac{\partial \rho}{\partial t} = [h, \rho]$$

Wigner Transformation and Wigner-Kirkwood Expansion

(Ref.: Ring-Schuck)

$$O_W(r, p) \equiv \int d^3 s \exp(-i p \cdot s) \langle r + s/2 | O | r - s/2 \rangle$$

$$(AB)_W = A_W \exp(i \hbar \Lambda) B_W$$

$$\Lambda \equiv \nabla'_r \cdot \nabla_p - \nabla'_p \cdot \nabla_r \quad (\nabla' \text{ acts on the left})$$

$$[A, B]_W = 2i A_W \sin(\hbar \Lambda/2) B_W = i \hbar \{A_W, B_W\}_{P.B.} + O(\hbar^3)$$

Vlasov Equation and Test Particle Method

Wigner Transform of TDHF

$$\left(i \hbar \frac{\partial \rho}{\partial t} = [h, \rho] \right)_W \rightarrow \frac{\partial f}{\partial t} = \{ h_W, f \}_{P.B.} + O(\hbar^2)$$

Winger Function: $f(r,p)$... Phase Space Density

Classical Hamiltonian: $h_W(r, p) = \frac{p^2}{2m} + U(r, p)$

Vlasov Equation

$$\frac{\partial f}{\partial t} + v \cdot \nabla_r f - \nabla U \cdot \nabla_p f = 0$$

Test Particle Method (C. Y. Wong, 1982)

$$f(r, p) = \sum_i^{AN_0} \delta(r - r_i) \delta(p - p_i) \rightarrow \frac{dr_i}{dt} = \nabla_p h, \quad \frac{dp_i}{dt} = -\nabla_r h,$$

**Mean Field Evolution can be simulated
by Classical Test Particles**

BUU (Boltzmann-Uehling-Uhlenbeck) Equation

(Bertsch and Das Gupta, Phys. Rept. 160(88), 190)

$$\frac{\partial f}{\partial t} + v \cdot \nabla_r f - \nabla U \cdot \nabla_p f = I_{coll}[f]$$

$$I_{coll}[f] = -\frac{1}{2} \int \frac{d^3 p_2 d\Omega}{(2\pi\hbar)^3} v_{12} \frac{d\sigma}{d\Omega} \\ \times [f f_2 (1-f_3)(1-f_4) - f_3 f_4 (1-f)(1-f_2)]$$

Incorporated Physics in BUU

Mean Field Evolution

(Incoherent) Two-Body Collisions

Pauli Blocking in Two-Body Collisions

- One-Body Observables (Particle Spectra, Collective Flow, ..)**
- ✗ Event-by-Event Fluctuation (Fragment, Intermittency, ...)**

AMD (Antisymmetrized Molecular Dynamics)

Wave Function

$$|\Psi\rangle = A \prod |\psi_i\rangle, \quad \psi_i = \phi(r; Z_i) \chi(\sigma, \tau),$$

$$\phi(r; Z) = \left(\frac{2\nu}{\pi}\right)^{3/4} \exp\left(-\nu(r - Z/\sqrt{\nu})^2 + Z^2/2\right)$$

$$\propto \exp\left(-\nu(r - D)^2 + iK \cdot (r - D)/\hbar\right) \left(Z = \sqrt{\nu} D + \frac{i}{2\hbar\sqrt{\nu}} K \right)$$

Equation of Motion

$$L = \frac{\langle \Psi | i\hbar \partial / \partial t - H | \Psi \rangle}{\langle \Psi | \Psi \rangle}, \quad \frac{d}{dt} \frac{\partial L}{\partial (d\bar{Z}_i/dt)} - \frac{\partial L}{\partial \bar{Z}_i} = 0$$

$$\rightarrow i\hbar C_{i\alpha, j\beta} \frac{dZ_i}{dt} = \frac{\partial H}{\partial \bar{Z}_i}$$

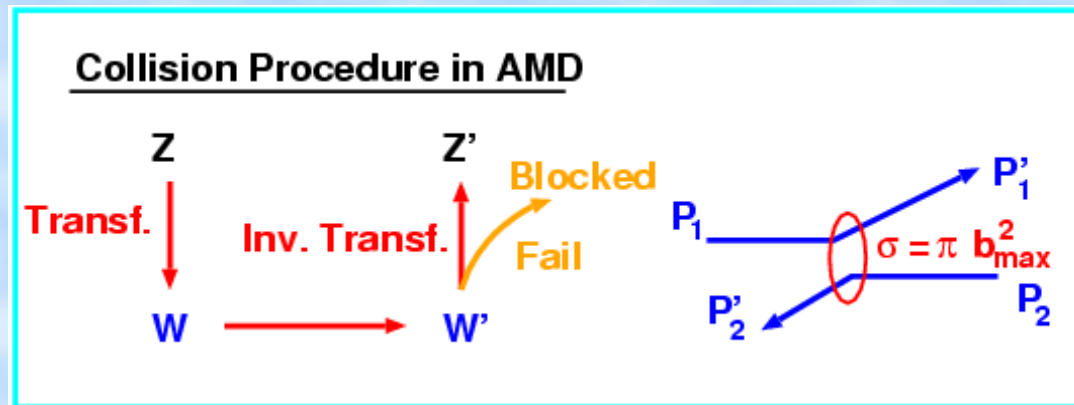
$$\text{No Antisymmetrization} \quad C = \delta \rightarrow \frac{dD_i}{dt} = \frac{\partial H}{\partial K_i}, \quad \frac{dK_i}{dt} = -\frac{\partial H}{\partial D_i}$$

Collision Term in AMD

Approximate Canonical Variables

$$W_i = \sqrt{Q_{ij}} Z_j = \sqrt{\nabla} R_i + \frac{i}{\sqrt{\nabla} \hbar} P_i, \quad Q_{ij} \equiv B_{ij} B_{ij}^{-1}, \quad B_{ij} = \langle \psi_i | \psi_j \rangle$$

Example $\langle L \rangle = \sum_{ij} B_{ji}^{-1} B_{ij} \frac{1}{i} \bar{Z}_i \times Z_j = \sum_i \bar{W}_i \times W_i$



Collision Term = "Canonical" Variable + Classical Analogy

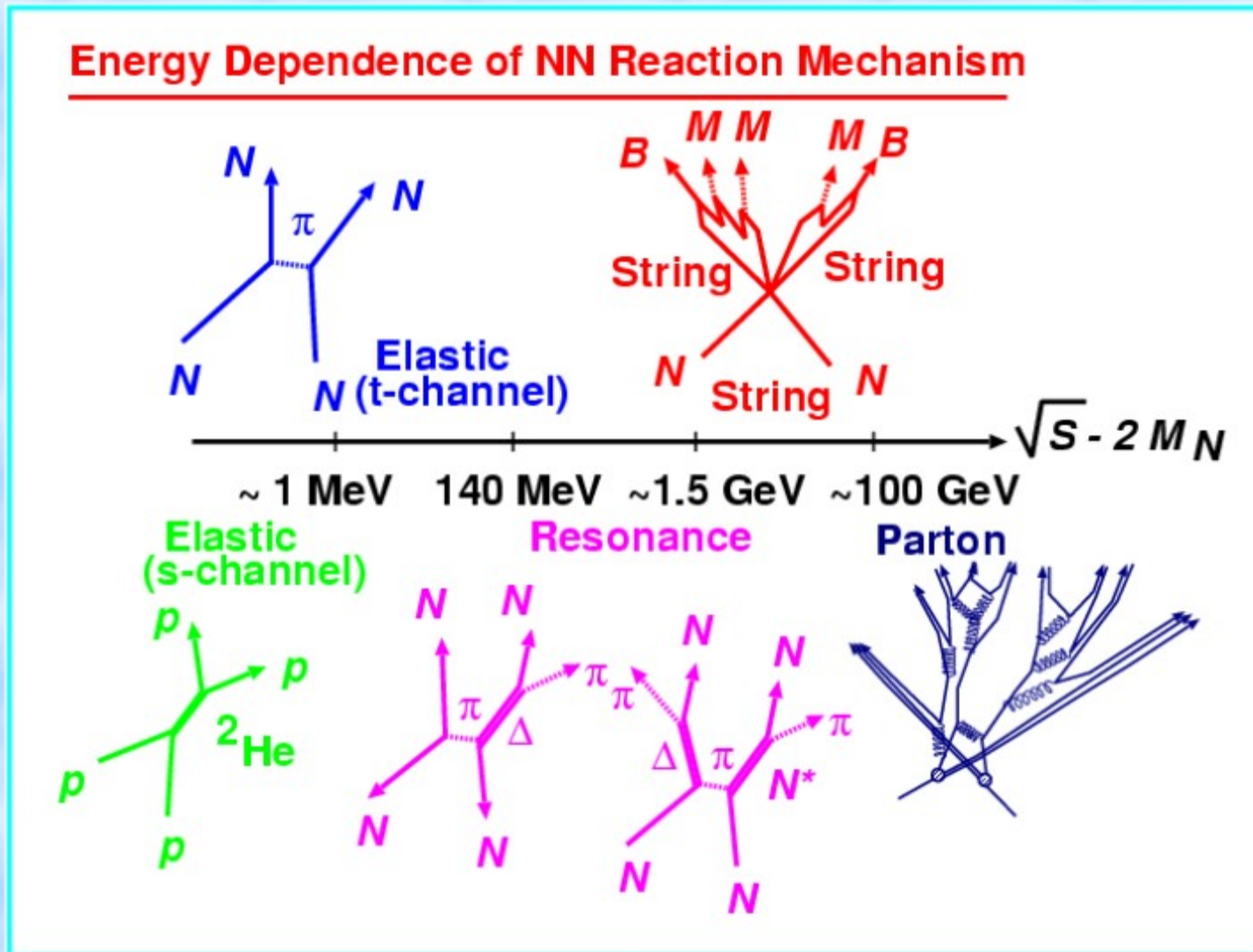
Physics included in AMD

Time Evolution of Anti-Symmetrized Wave Function
Two-Body Collisions with Pauli Blocking
Event-by-Event Fluctuation

Points to be Improved in AMD

Wave Packet Dynamics → Not an Eigen State of Energy, $J\pi$
Initial and Final Fragment State → Not Quantized
Two Body Collisions → Classical Analogue (Not Derived)

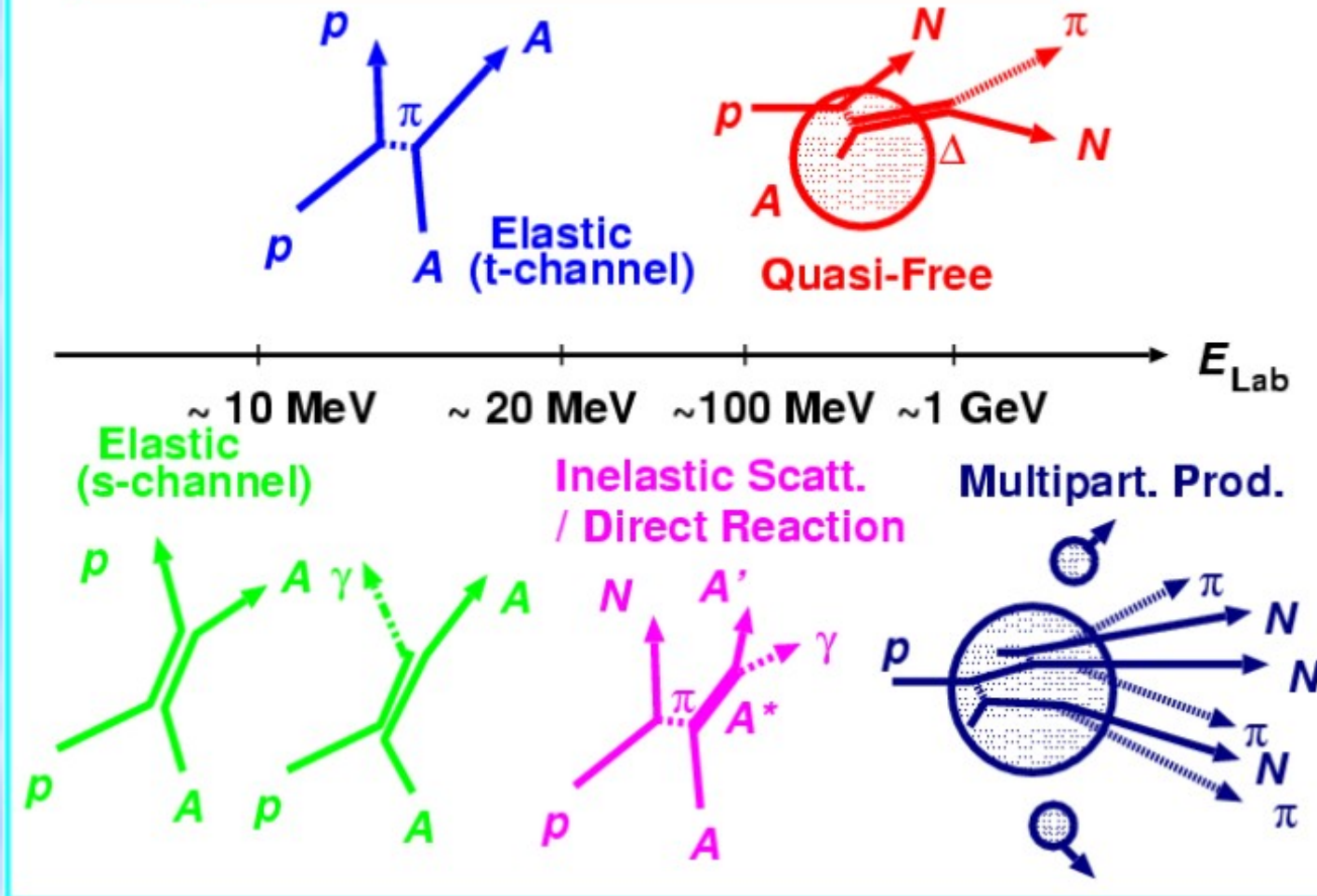
Rough Reaction Mechanism of NN Collisions



Elastic \rightarrow Resonance \rightarrow String \rightarrow Jet

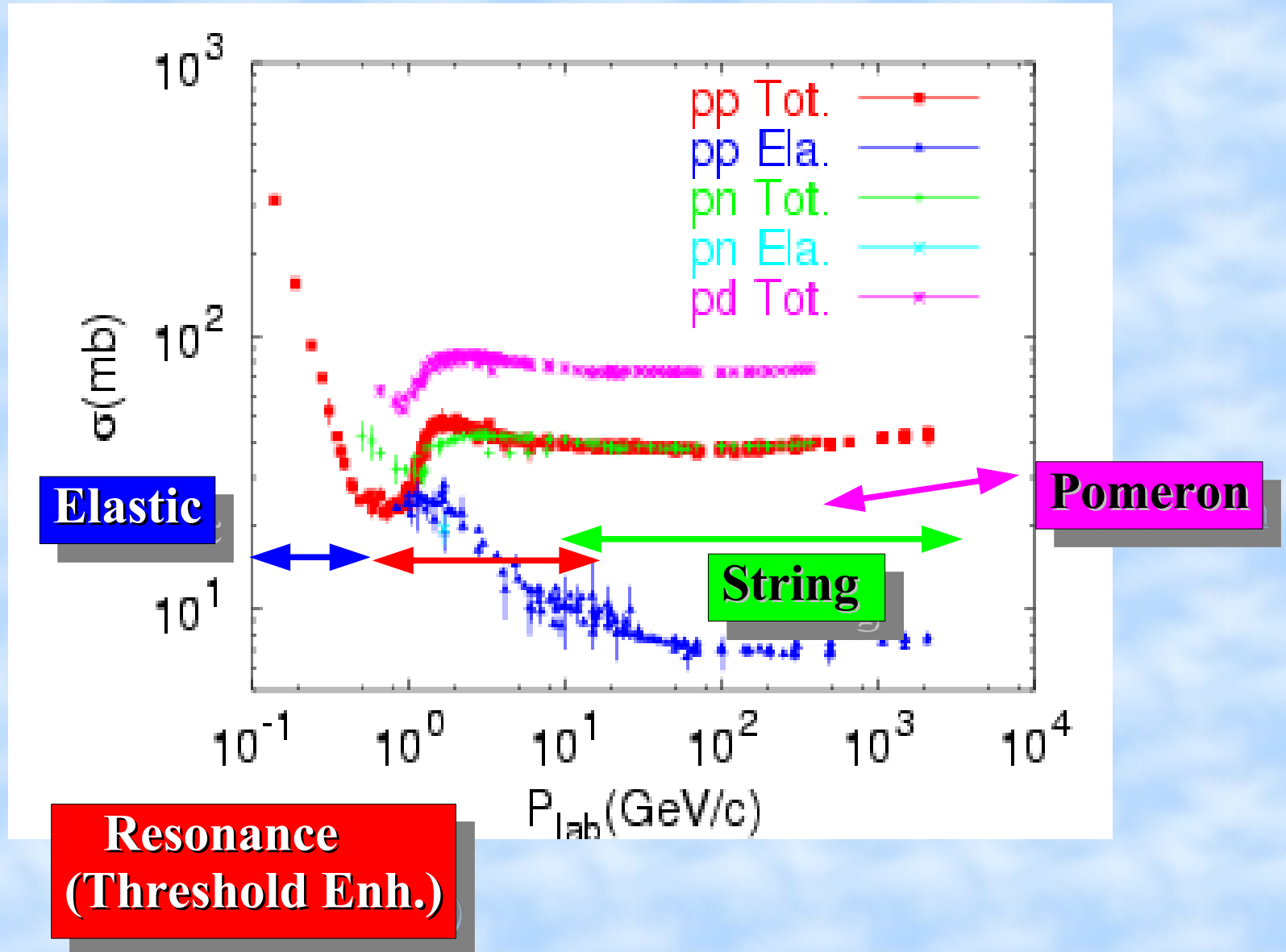
Rough Reaction Mechanism of NA Collisions

Energy Dependence of pA Reaction Mechanism



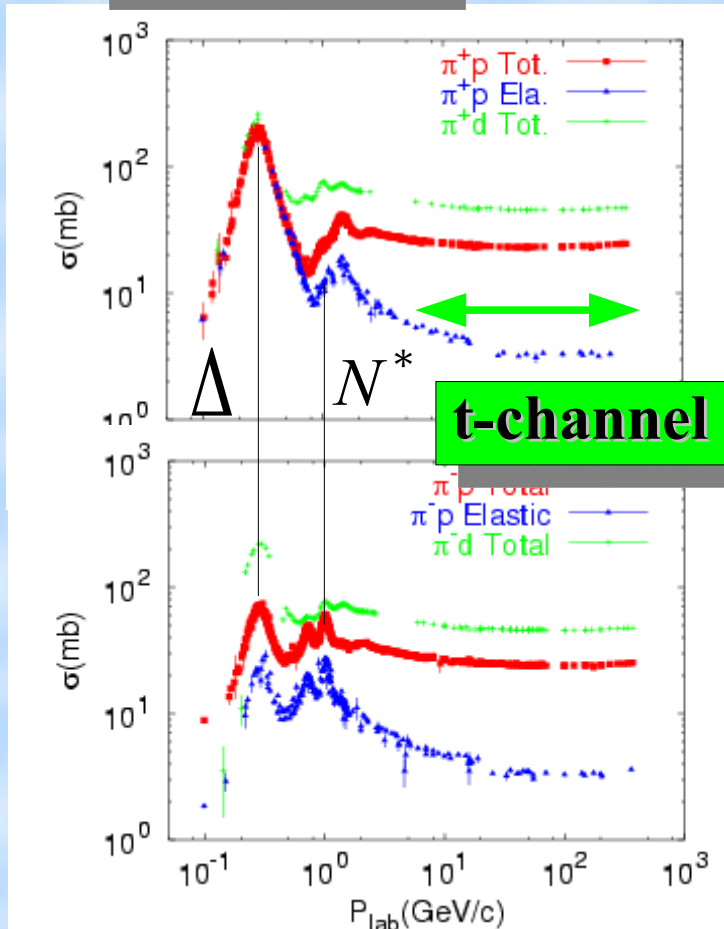
NN Cross Sections

From Particle Data Group

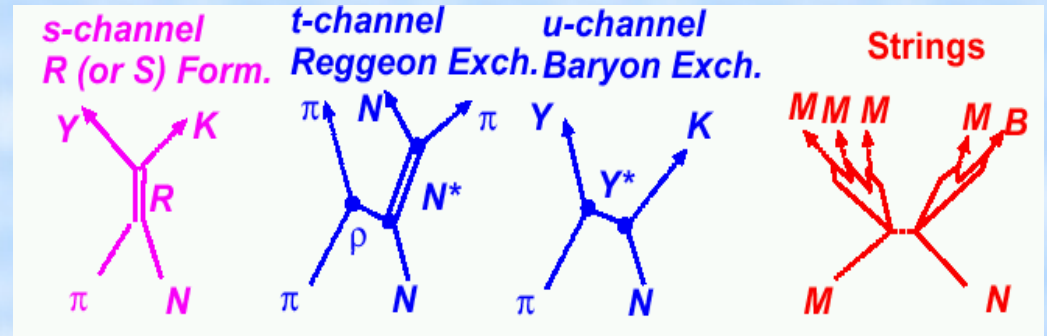


Meson-Baryon Cross Section

**Resonance
(s-channel)**



t-channel and String



**Strong Resonance Formation at Low Energy
 → Smooth Behavior at High Energy**

Statistical Model of Hadrons and Fragments

Grand Canonical Statistical Ensemble of Constituents

$$\begin{aligned} N_i &= \int d\Gamma_i f_i(E_i - \mu_i, T) , \\ d\Gamma_i &\equiv \frac{g_i d^3 r_i d^3 p_i}{(2\pi)^3} , \\ f_i &\equiv \frac{1}{\exp [(E_i - \mu_i)/T] \mp 1} , \end{aligned}$$

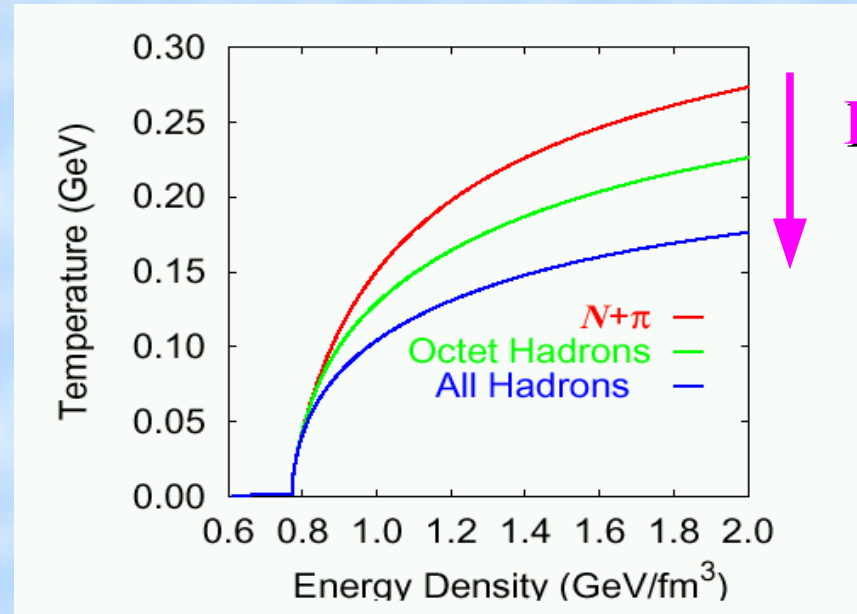
Thermodynamical Functions

$$\begin{aligned} \Omega &= - \sum_i \int d\Gamma_i f_i \frac{p_i^2}{3E_i} = -PV , \\ S &= \frac{1}{T} \sum_i \int d\Gamma_i f_i \left(E_i - \mu_i + \frac{p_i^2}{3E_i} \right) , \\ P &= -\Omega/V \\ N_i &= \int d\Gamma_i f_i , \\ E &= \sum_i \int d\Gamma_i f_i E_i . \end{aligned}$$

Constituents can either be Hadrons (high-T) or Fragments (low-T)

Hadronic Caloric Curve and Hagedorn Gas Behavior

Hadronic Caloric Curve (Otuka, Thesis)



Larger DOF
→ Lower T

Hagedorn Gas Behavior:

Exponentially Growing Level Density → Limiting Temperature

$$\rho(m) \rightarrow am^{-\frac{5}{2}} \exp\left(\frac{m}{T_0}\right) (\text{GeV}^{-1}).$$

$$\begin{aligned} Z(V, T) &= \sum_{(k)} \exp\left(\frac{1}{T} \sum_i \sum_{\tau} \epsilon_{i,\tau} \nu_{i,\tau}\right) - 1 \\ &= \exp\left[\frac{VT}{2\pi^2} \sum_{n=1}^{\infty} \frac{1}{n^2} \int_0^{\infty} dm \rho(m, n) m^2 K_2\left(\frac{nm}{T}\right)\right] - 1 \\ &\equiv \frac{VT^3}{(2\pi)^{\frac{3}{2}}} \int_0^{\infty} dm m^{\frac{3}{2}} \rho(m) Q\left(\frac{m}{T}\right), \end{aligned}$$

Simple Model of QCD Phase Transition

Massless Particles at Zero Chem. Pot.

$$\Omega = -\frac{\pi^2 V}{90} \left(\sum_B g_B + \frac{7}{8} \sum_F g_F \right) T^4,$$

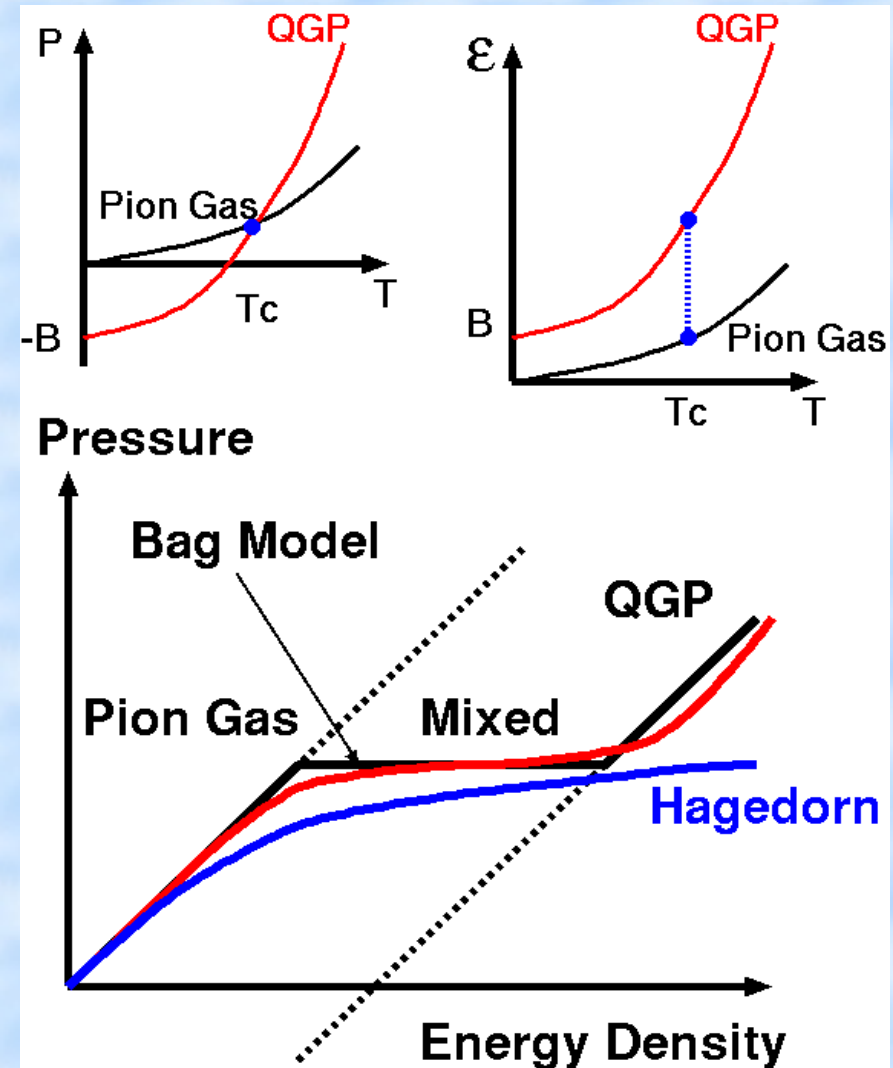
Massless (Free) Pion Gas

$$P_\pi = \frac{\pi^2}{30} T^4, \quad \epsilon_\pi = \frac{\pi^2}{10} T^4$$

QGP with Finite Bag Constant

$$P_{QGP} = \frac{37 \pi^2}{90} T^4 - B$$

$$\epsilon_{QGP} = \frac{37 \pi^2}{30} T^4 + B$$



$$DOF = 2(\text{spin}) \times 2(q, \bar{q}) \times 3(\text{color}) \times 2(\text{flavor}) \times \frac{7}{8}(\text{Fermion}) + 2(\text{spin}) \times 8(\text{color})$$

Statistical Model of Nuclear Fragments

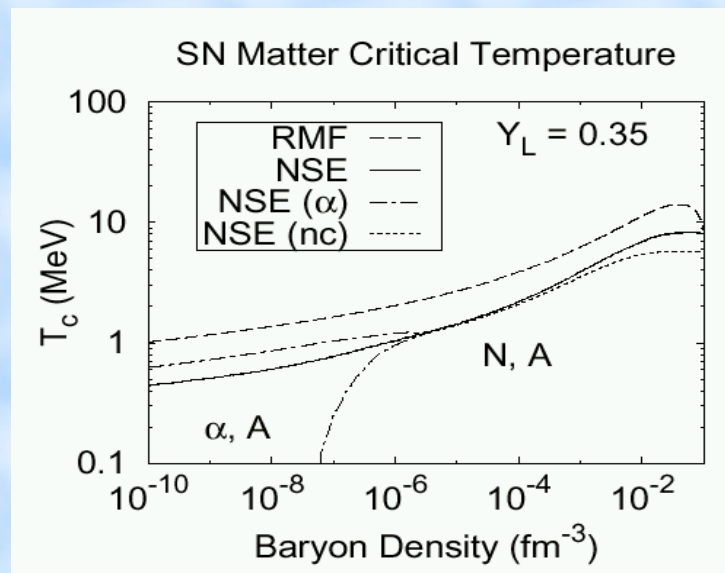
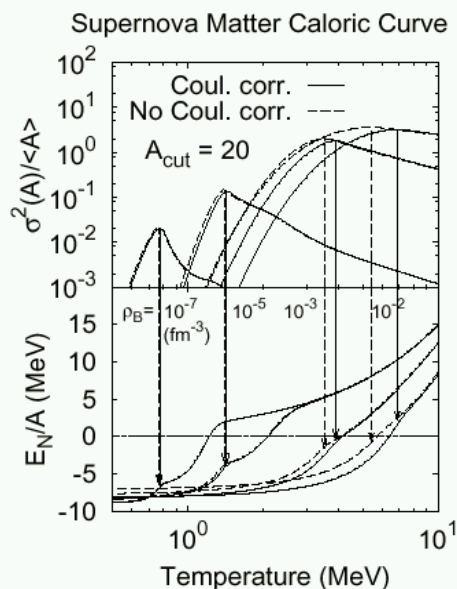
Nuclear Level Density (taken from Randrup & Fai)

$$\zeta_f(T) = \sum_i g_f^{(i)} \exp(-E_f^{*(i)}/T)$$

$$\simeq g_f^{(g.s.)} + \frac{c_1}{A_f^{5/3}} \int_0^\infty dE^* e^{-E^*/T} \exp(2\sqrt{a_f E^*}),$$

$$a_f = \frac{A_f}{8} (1 - c_2 A_f^{-1/3}) \text{ (MeV}^{-1}\text{)}, \quad c_1 = 0.2 \text{ (MeV}^{-1}\text{)}, \quad c_2 = 0.8$$

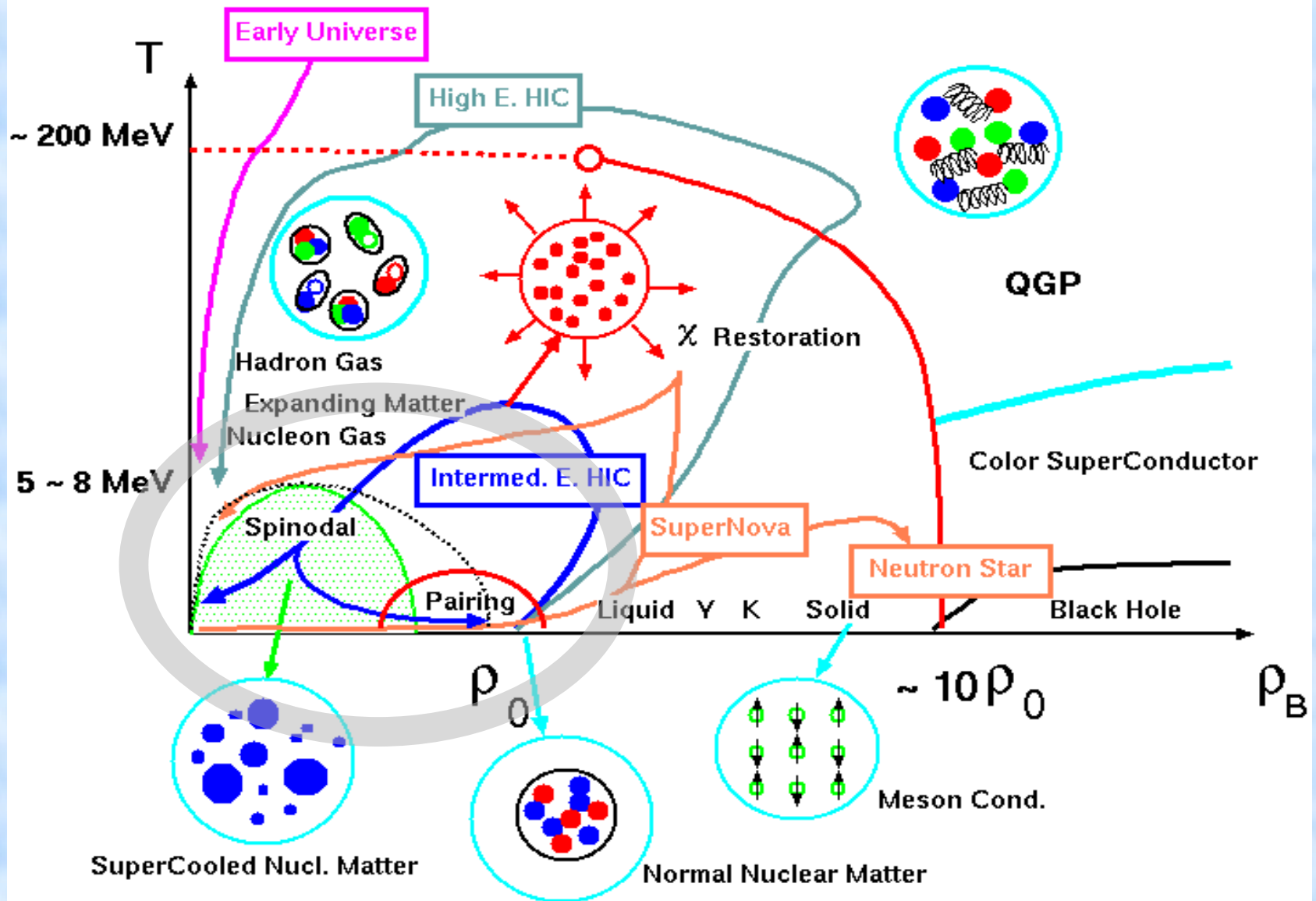
$$\zeta \propto \exp(2\sqrt{a E^*}) \rightarrow E^* = a T^2$$



Ishizuka, AO, Sumiyoshi
(nucl-th/0208020)

Part I:
Fragment Formation
and Nuclear Liquid Gas Phase Transition

Hadronic Matter Phase Diagram



Contents

Introduction — Liquid-Gas Phase Transition

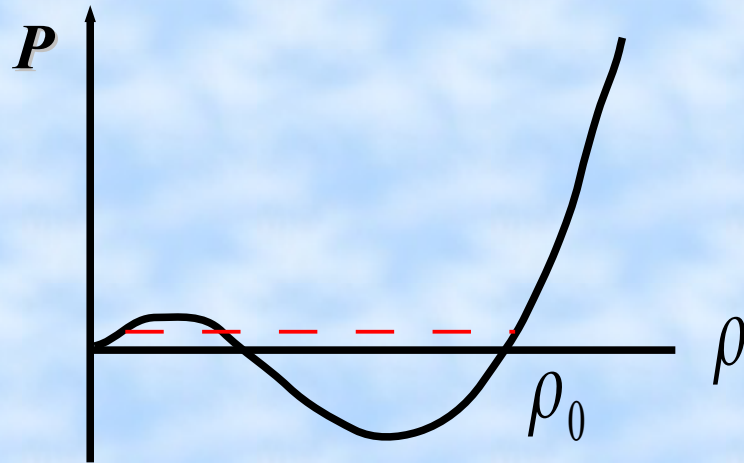
Sideward Peak of IMF Emission in High Energy pA Reaction

Phase Transition of Supernova Matter

Nuclear Int. Van der Waals Int.

→ LG Phase Transition is expected.

RMF



Recent Experimental Progress

Two indep. exp. on two indep. Observables show the Existence of First Order L.-G. Phase Transition.

Nuclear Caloric Curve

**J. Pochadzalla et al.
Phys. Rev. Lett. 75 (1995) 1040.
(GSI-ALLADIN collab.)**

Boiling Temperature is Clearly Seen

**Fragment Yields are assumed
to follow Equilibrium Statistics**

$$Y_f \propto g_f \exp((B_f + Z \mu_p + N \mu_n)/T)$$

$$\rightarrow \frac{Y(^4\text{He})/Y(^3\text{He})}{Y(^7\text{Li})/Y(^6\text{Li})} \propto \exp(\Delta B/T)$$

Negative Heat Capacity

Negative Heat Capacity
→ *First Order*

T and E^* are determined
from *Fragment Multiplicity*
and *Kinetic Energy*
based on Theoretical Model

M. D Agostino et al.
(MSU Exp./INFN-IN2P3 Collab.)
PLB 473 (2000) 219.

For theoretical understanding, we need models,

Which describes both of Reactions and Statistics

T in Exp. Fragment Yield in H.I. Collisions + Equil. Assumption

With Quantum Statistical nature

$$E \propto T^2 \quad \text{in nuclei at low T.}$$

Molecular Dynamics with Quantum Fluctuation

AO & Randrup,

Nucl. Phys. A565(1993), 474; Phys. Rev. Lett. 75(1995), 596;

Ann. Phys. 253 (1997), 279; Phys. Lett. B394(1997), 260;

Phys. Rev. A55(1997), 3315R)

Hirata, Nara, AO, Harada, Randrup,

Prog. Theor. Phys. 102 (1999), 89.

Ono-Horiuchi (AMD-V, Phys. Rev. C53 (1996), 2958)

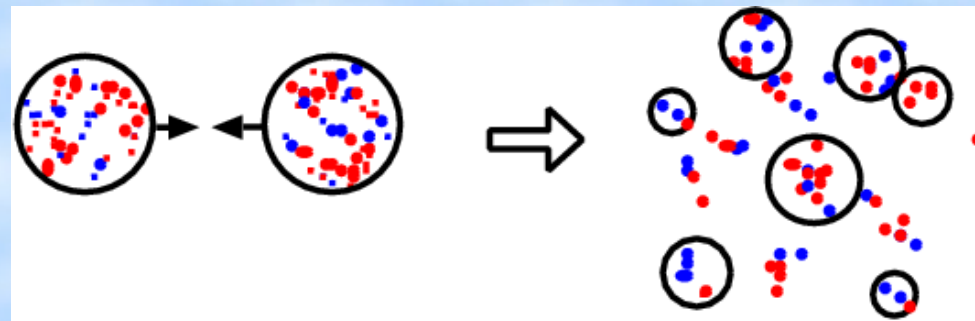
Sugawa-Horiuchi-Ono (AMD-MF, Phys. Rev. C60 (1999) 064607)

Wave Packet Statistics

Mass Distribution

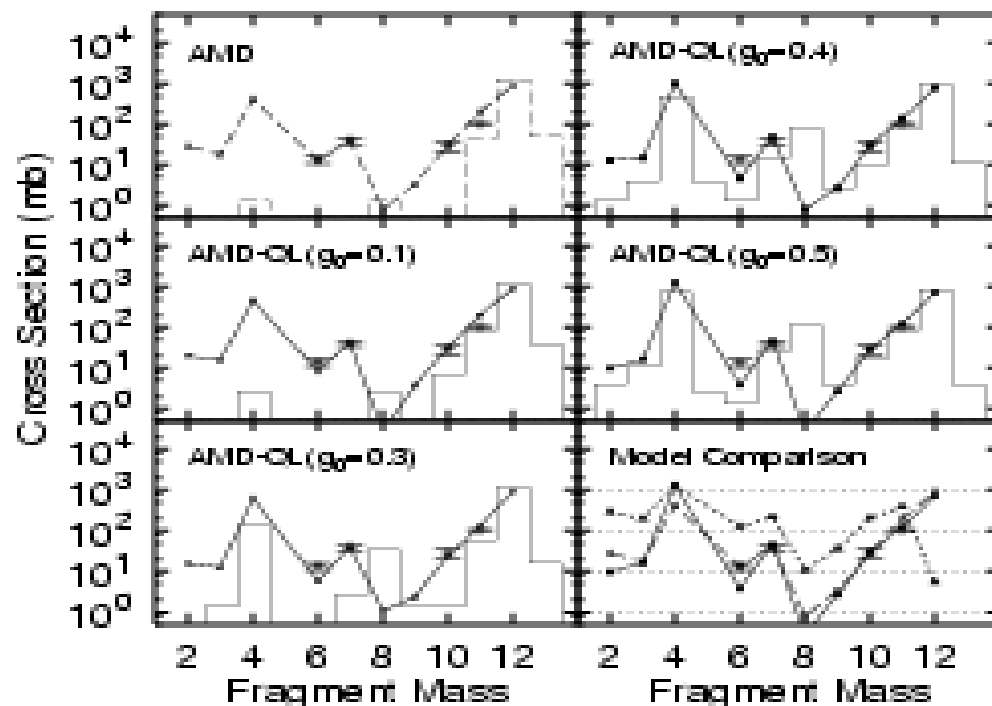
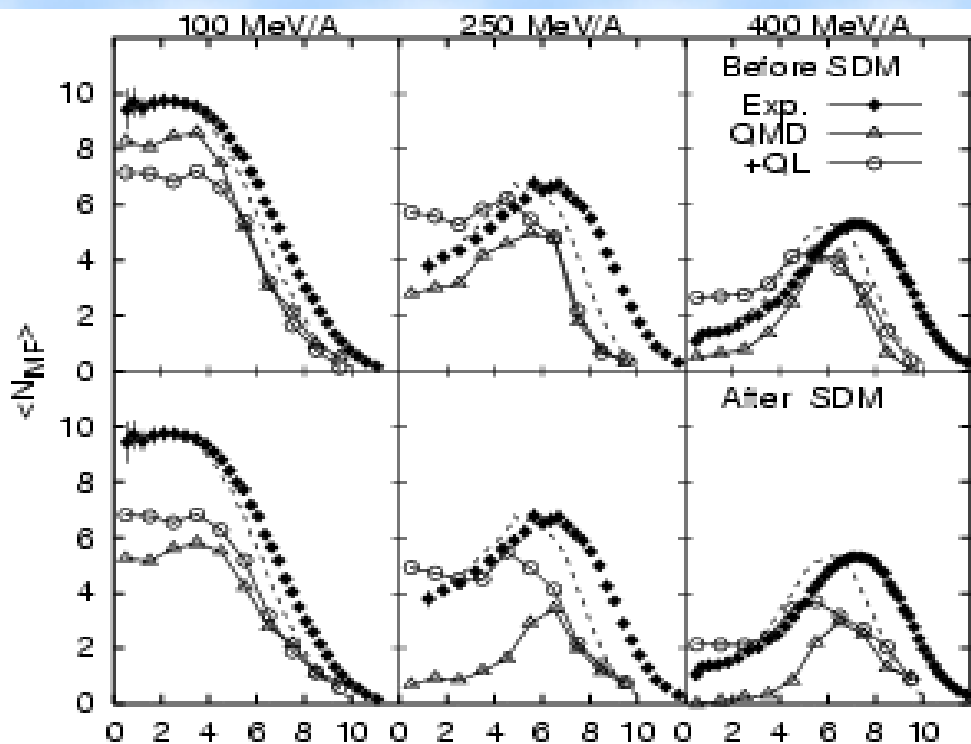
Caloric Curve

Wave Packet Dynamics



Au + Au Collisions

p + 12 C Collisions



(Hirata et al. PTP102(99),89)

What is Understood ?

LG Phase Transition is of First Order (Exp.).

**It can be understood in Microscopic MD qualitatively,
e.g. Fragment Yield.**

What is NOT Understood ?

**Direct Relation between Fragment Formation
and the Properties of Nuclear Matter**

Are Fragments Produced through LG Phase Transition ?

“Initial” Condition of Fragmentation

At Which T and ρ Fragments are Formed ?

Is Equilibrium Reached in Heavy-Ion Collisions ?



Simpler Cases: pA Reaction & Supernova Explosion !

***Shape effects
on high-energy proton induced
nuclear fragmentation***

**Shuji Yamaguchi
Akira Ohnishi**

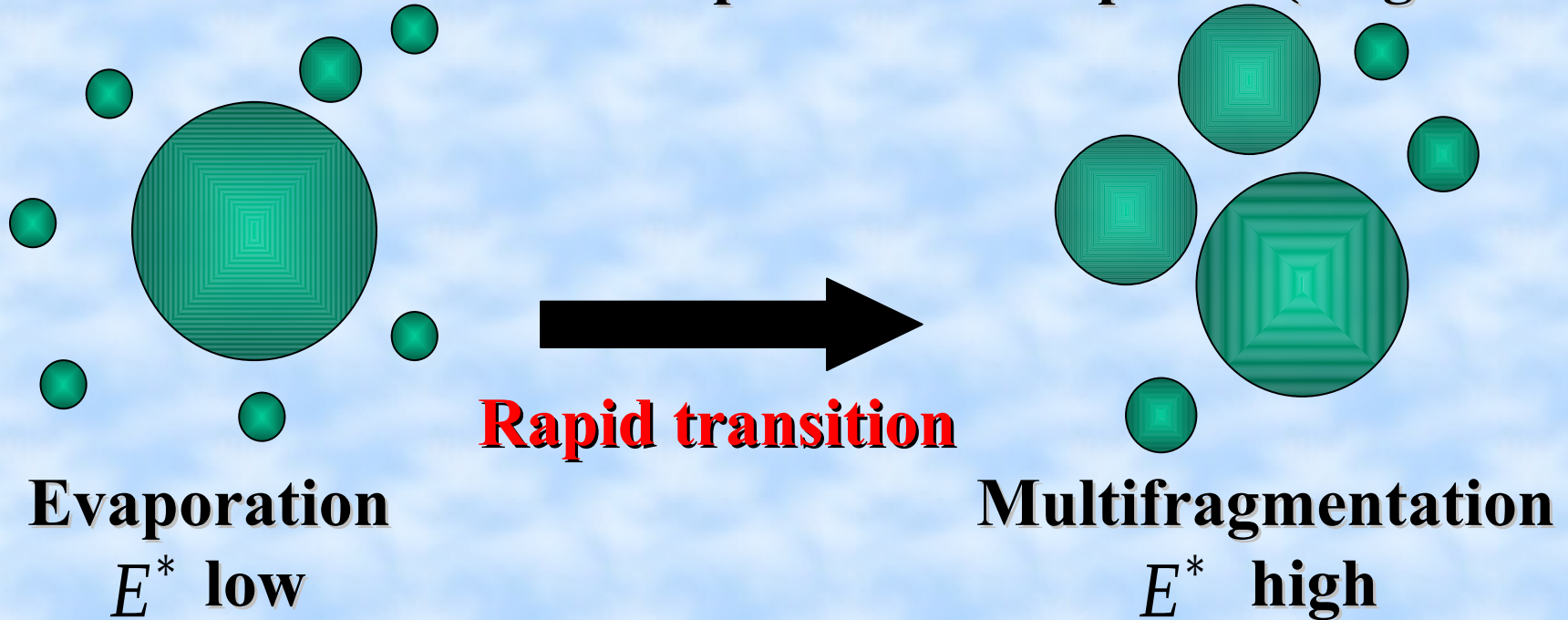
Hokkaido Univ.

- 1. Introduction**
- 2. Equilibrium Percolation Model**
- 3. Non Equilibrium Percolation Model**
- 4. Summary**

1. Introduction

Nuclear fragmentation

Excited nucleus breaks up into several pieces (fragments).



This evolution has been discussed in relation to the phase transition of nuclear matter [1].

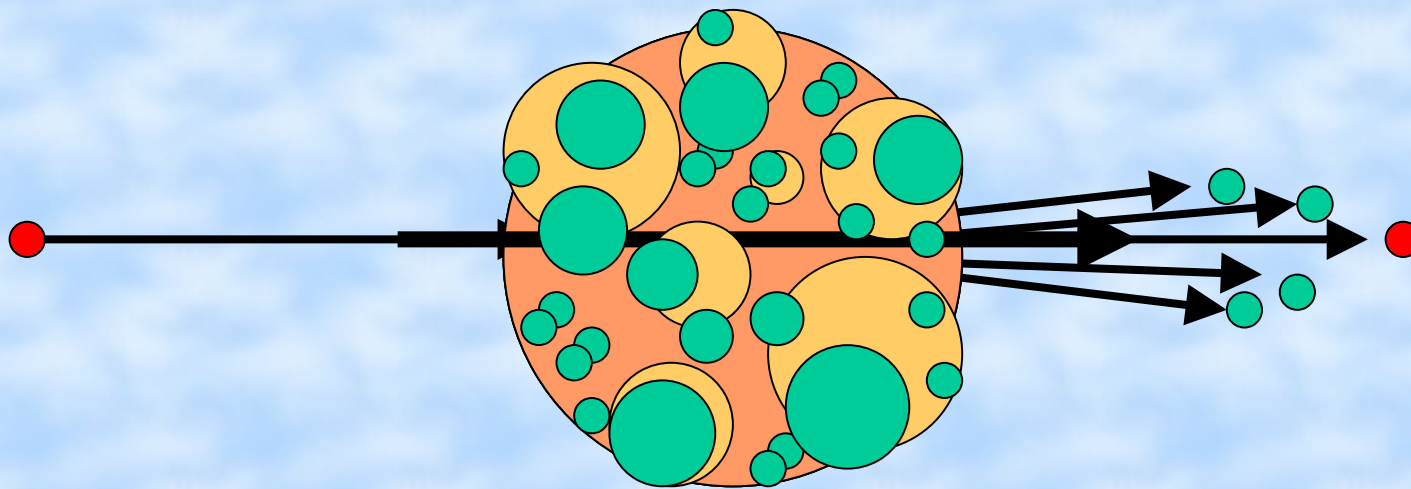
Especially in **Equilibrium Percolation Model [2].**

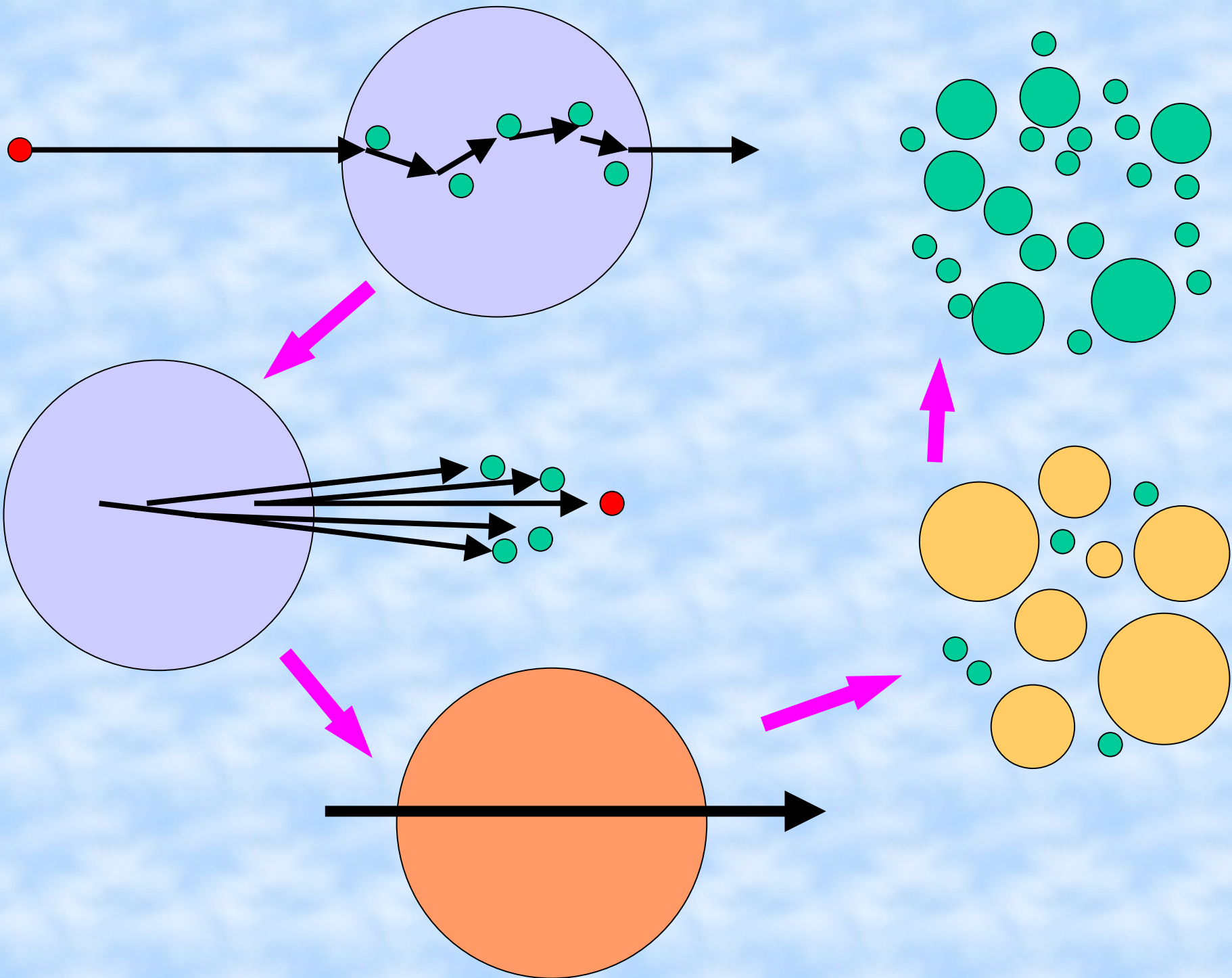
[1] M.K.Berkenbusch P.R.L. 88 (2002) 022701

[2] W.Bauer P.R.C. 38 3 (1988) 1297

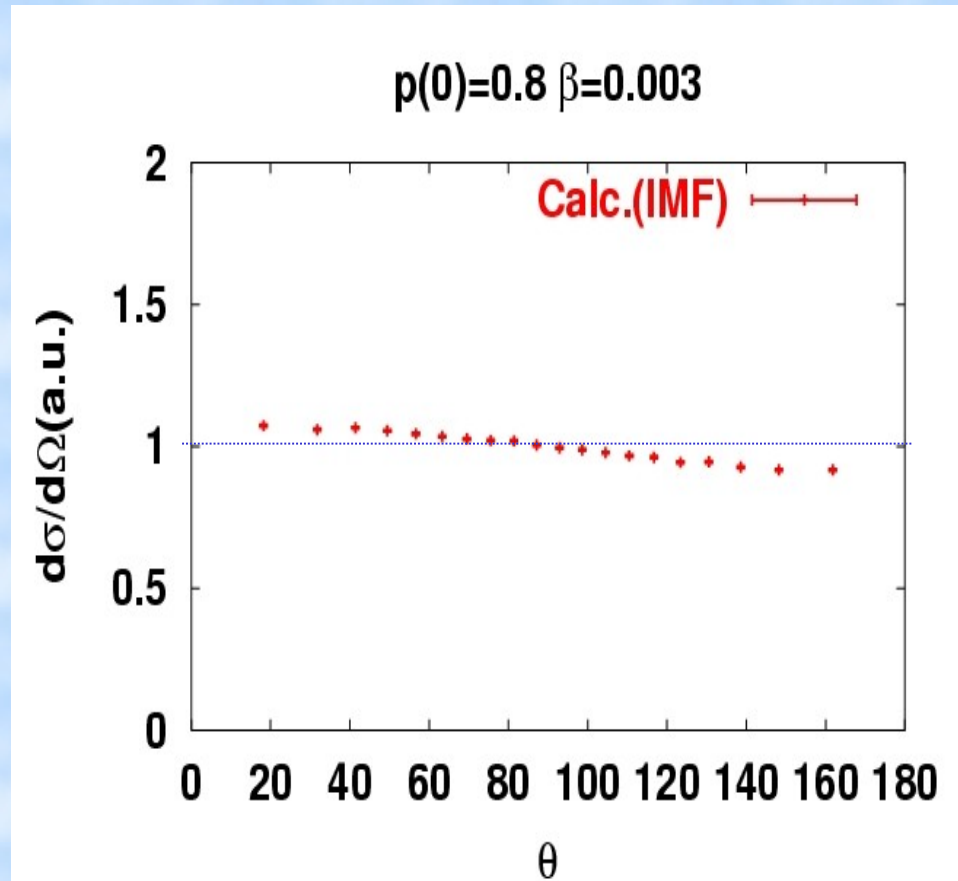
Standard Picture of proton induced Nuclear Fragmentation

**The incident proton passes through the target nucleus.
Residual excited nucleus reaches equilibrium, moving in forward.
The excited nucleus breaks up into fragments.
These fragments evaporate nucleons and alphas or make fission.**





The fragment angular distribution is peaked in forward.



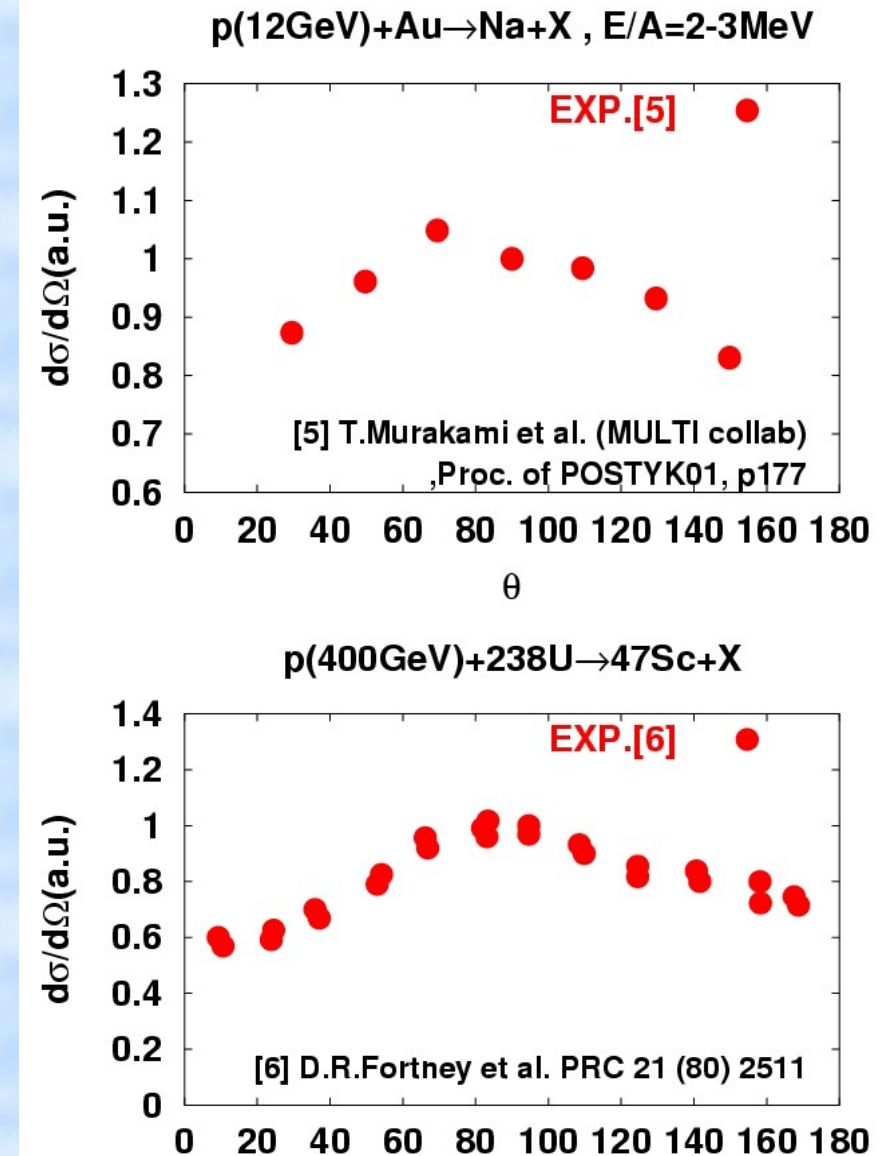
Long-standing problems

1. p(10GeV~)+A

Sideward peak of
IMF angular distribution

2. p(100GeV~)+A

backward enhancement of
IMF angular distribution



These phenomena can not be explained^θ

(IMF: Intermediate Mass Fragment, A=10~50)

Standard picture need to be modified!

Purpose

The sideward peak and the backward enhancement

We study

the non-equilibrium shape effects of nuclear fragmentation,
within Bond Percolation Model.

Fragmentation should occur **BEFORE** equilibrium.

→ Temperature distribution may **NOT BE UNIFORM**.

Thus, we consider two models.

Cylinder model and **Cone** model

2. *Equilibrium Percolation Model* [2]

Show a power law
Position information

Nucleons □ **Sites on a simple cubic lattice in 3D**

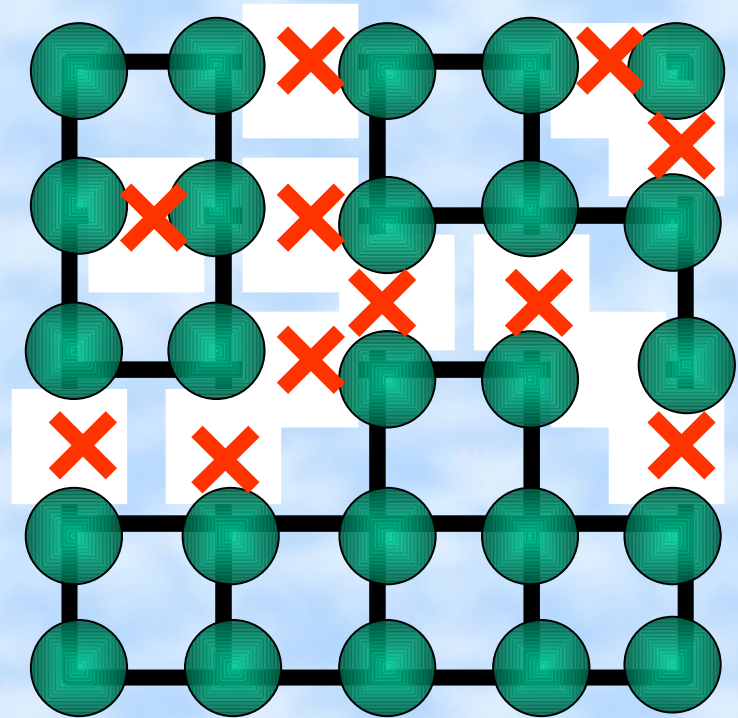
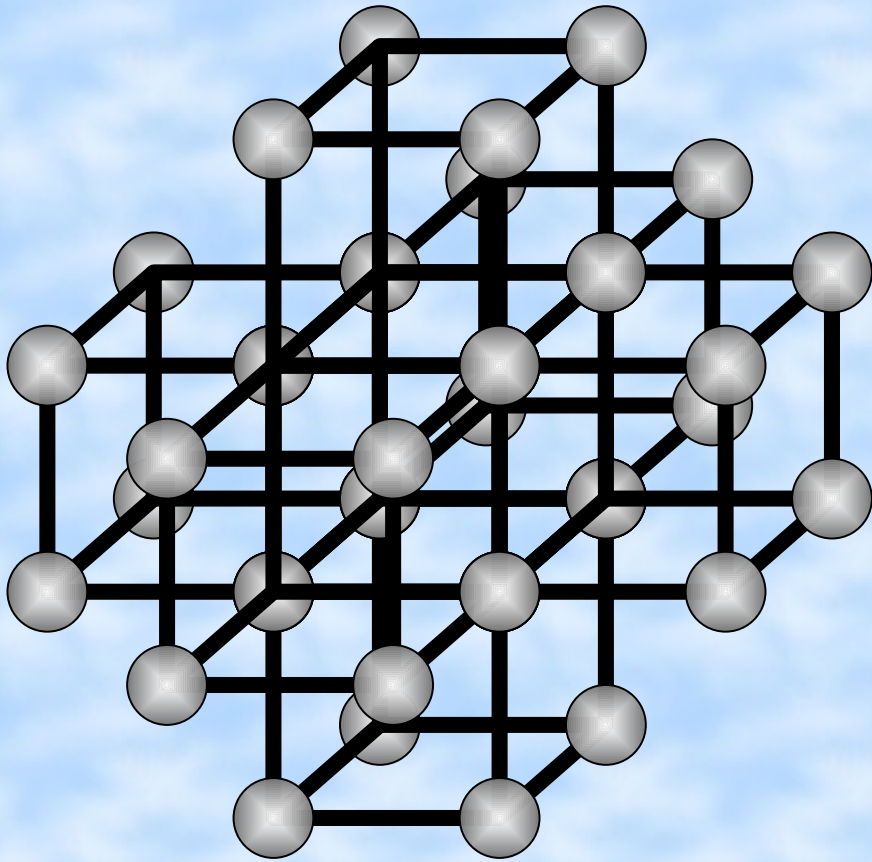
Bond breaking probability P(b)

$$P(b) = P(0) \frac{\int \rho(x=b, 0, z) dz}{\int \rho(0, 0, z) dz} \quad \text{(Eikonal)}$$

Connected sites □ **Fragment**

$$\sigma(A_f) = 2\pi \int_0^{b^{(max)}} b db \frac{N(A_f, b)}{N_{run}(b)}$$

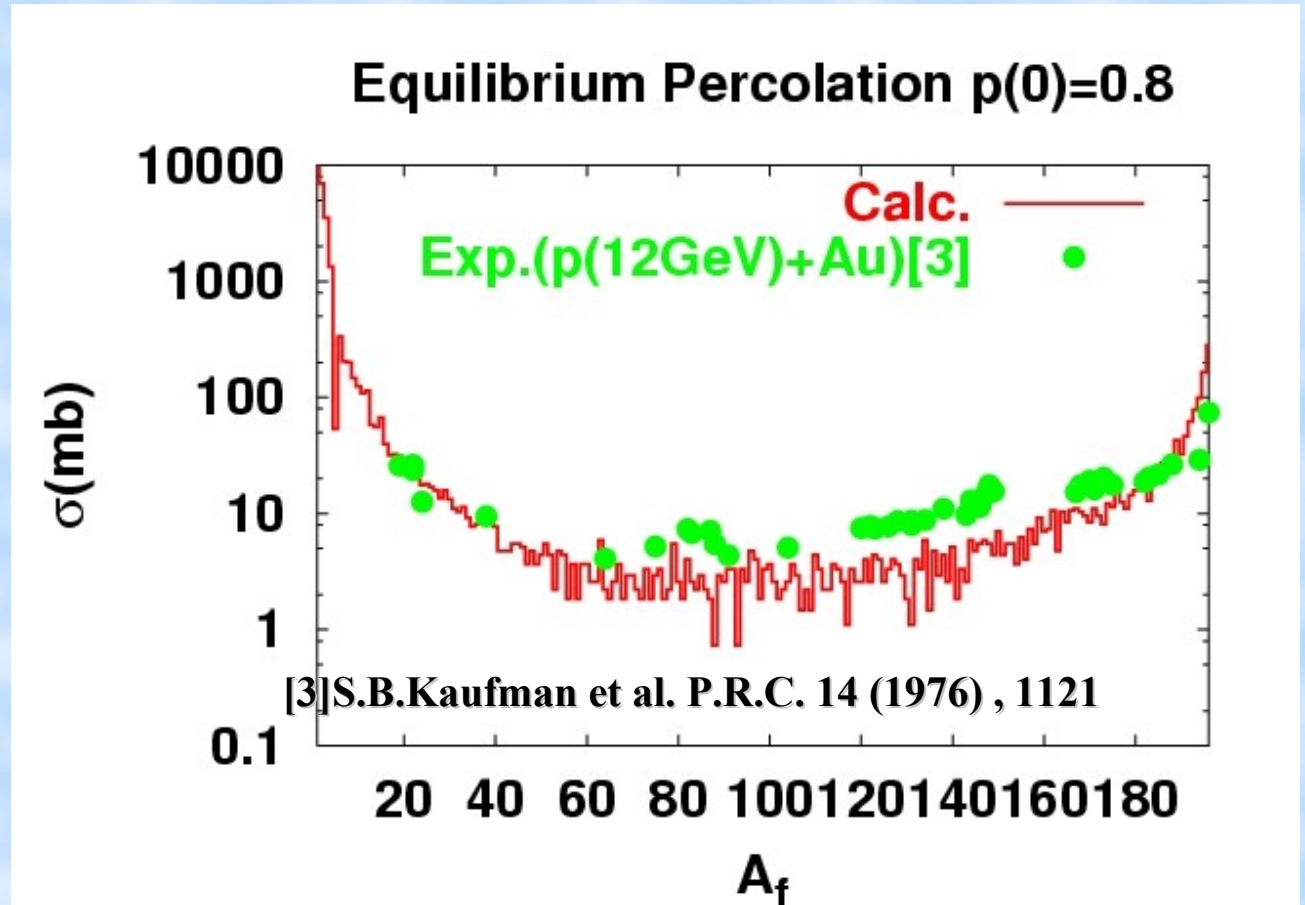
Examples



Fragment Mass Distribution in Equilibrium Percolation Model

Features

U-Shape
A power law



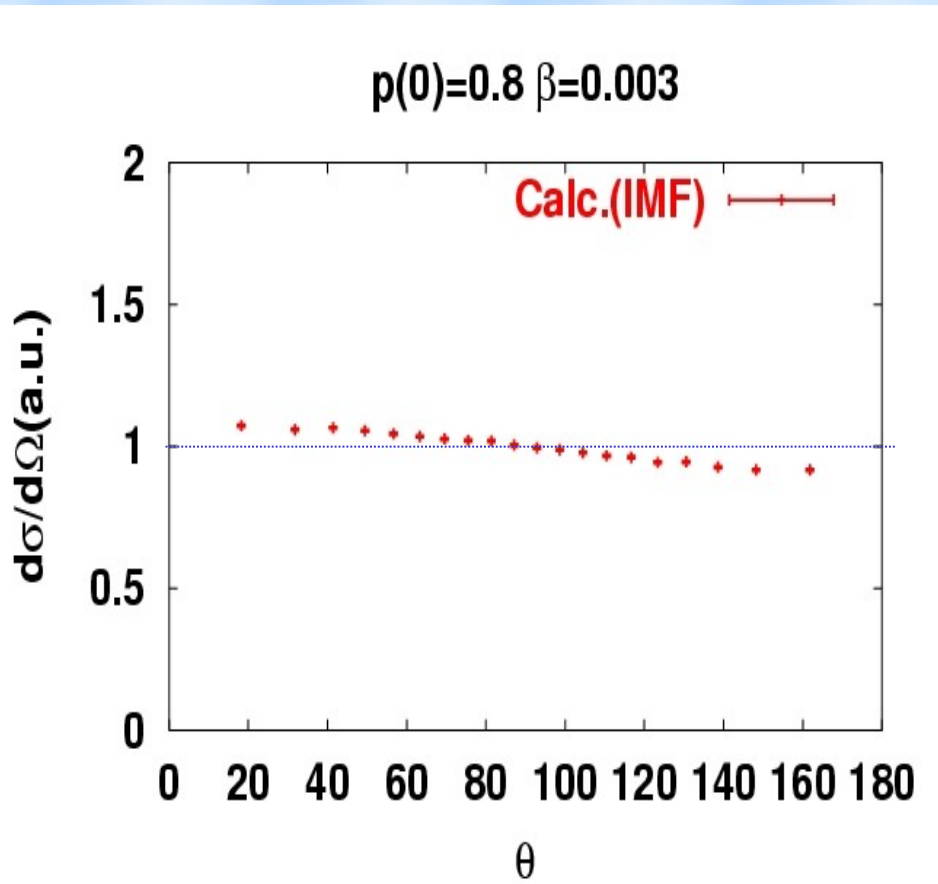
Using SDM(Statistical Decay Model)[4]
to consider evaporation and fission of fragments

Excitation energy in SDM

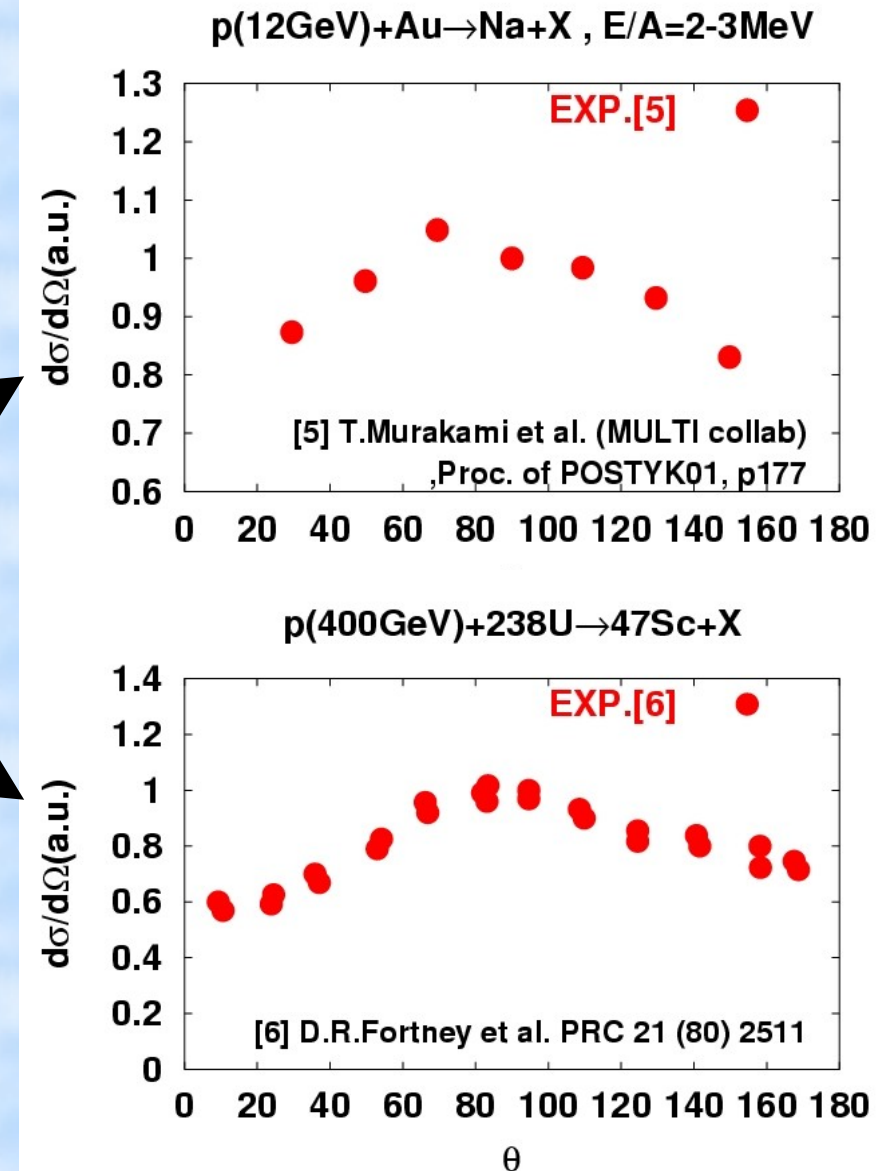
$$E^* = 5.25 \times p(b) \times (\text{All Bonds}) - 5.25 \times (\text{Broken Bonds})$$

[4] K.Niita et al. PRC 52 (95) 2620

However, the observed fragment angular distribution cannot be explained in Equilibrium Percolation Model.



Equilibrium Percolation Model



3. Non Equilibrium Percolation Model

Sideward peak : **Cylinder Model**

Cylindrical hot region[7]

Backward enhancement : **Cone Model**

Conic hot region

Previous works[8,9] :low ρ  **Present work** : high T

[7] Y.Hirata et al.N.P.A707 (2002) 193

[8] T.Maruyama et al.PTP 97 4 (1997) 579

[9] J.Hufner Phys.rep.125(85)129

3.1. Cylinder Model + SDM

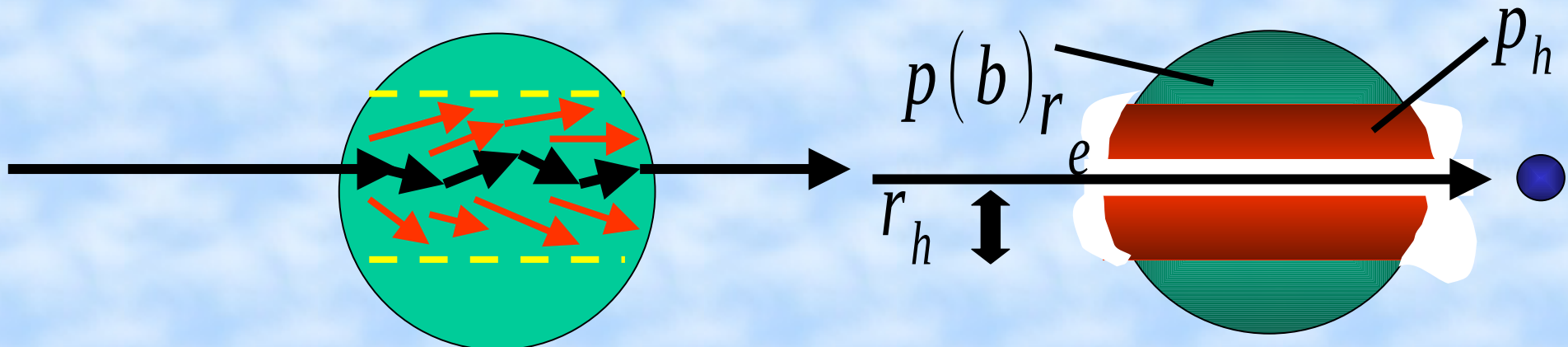
Processes

When nucleus breaks up ,high temperature zone remains.

Isotropic Fermi momentum is distributed to each nucleon.

Coulomb expansion.

Residual nucleus moves in forward direction with a velocity β .



$$\beta = \frac{P_Z(Res)}{M_N A(Res)}$$

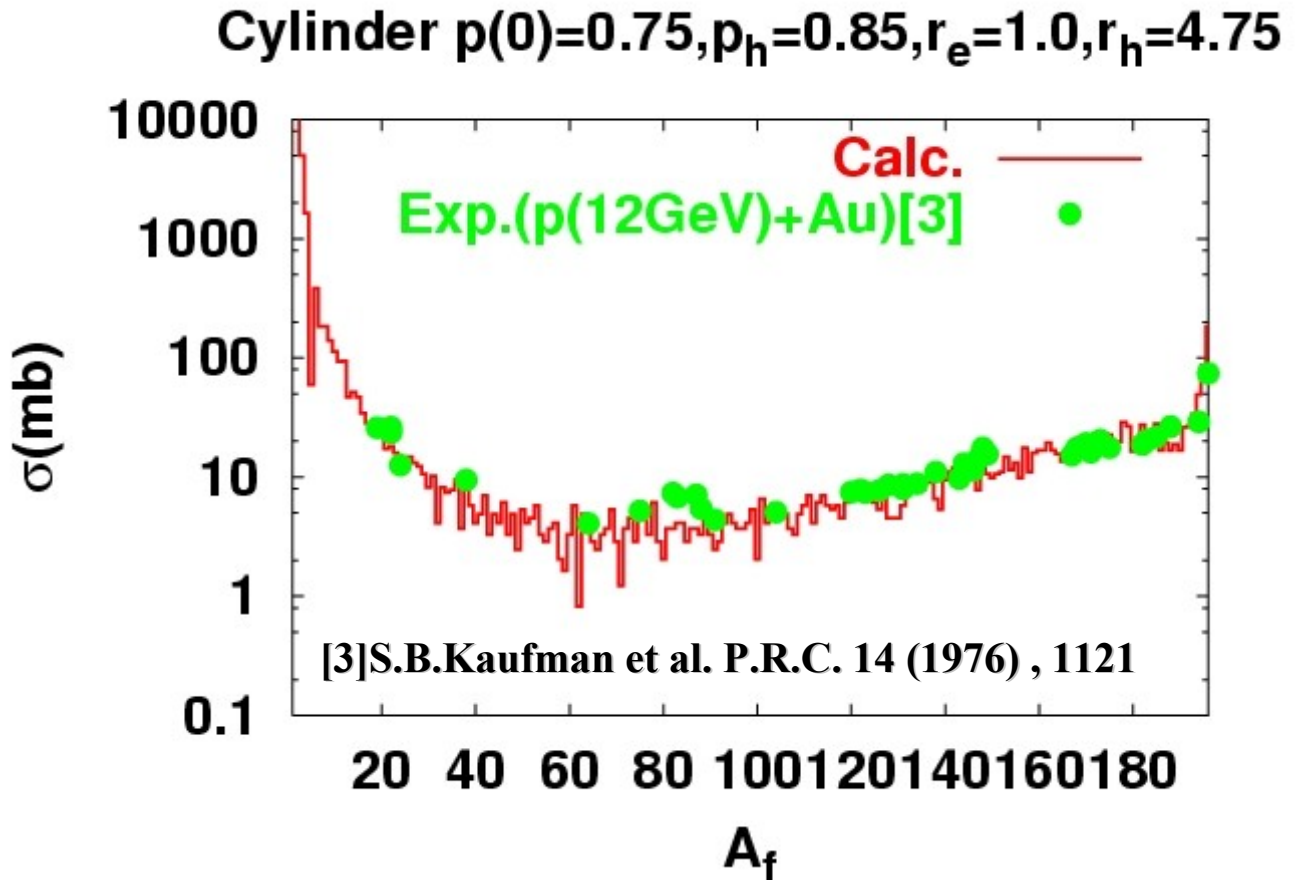
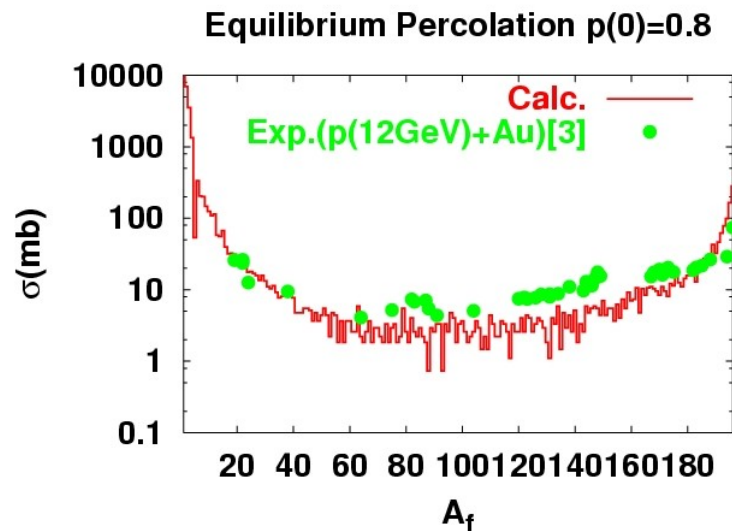
M_N : mass of nucleon

A : mass number of target

P_Z : momentum of z-axis

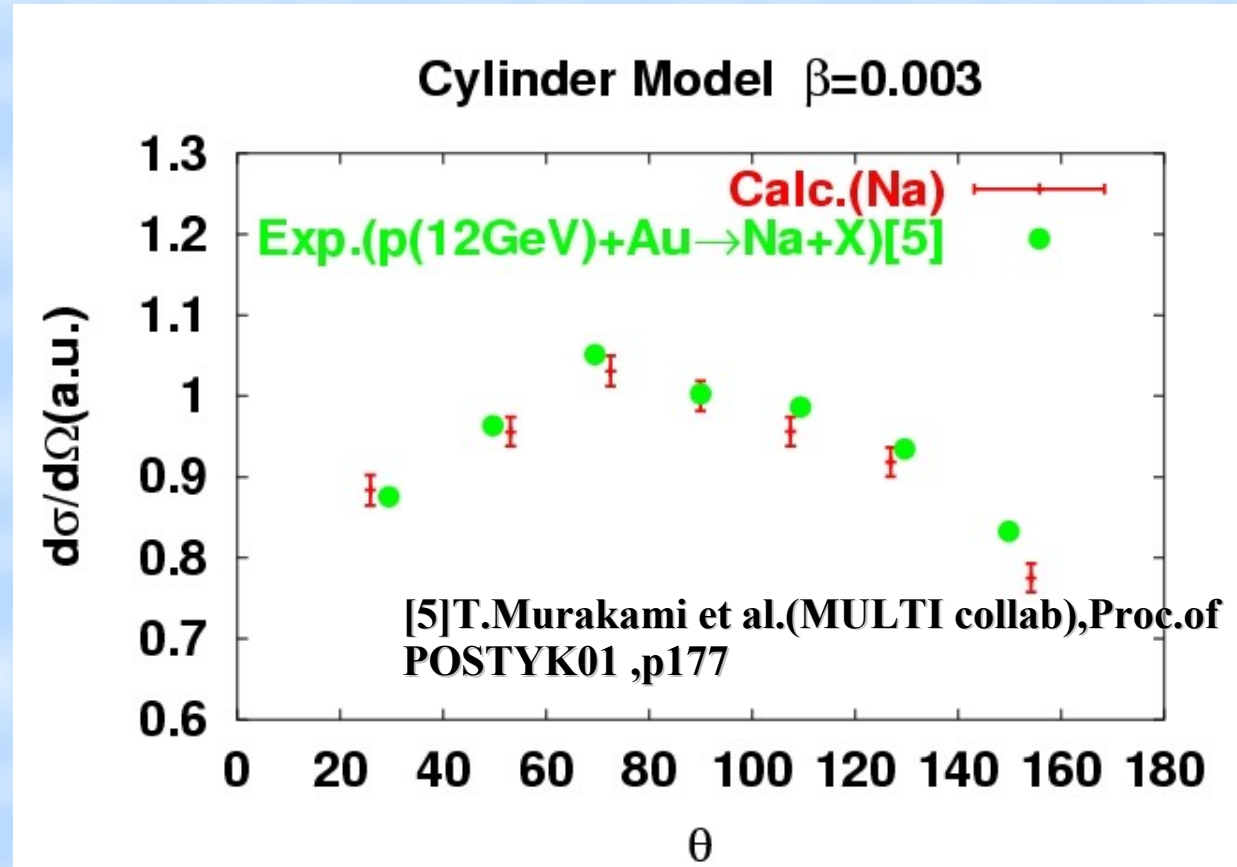
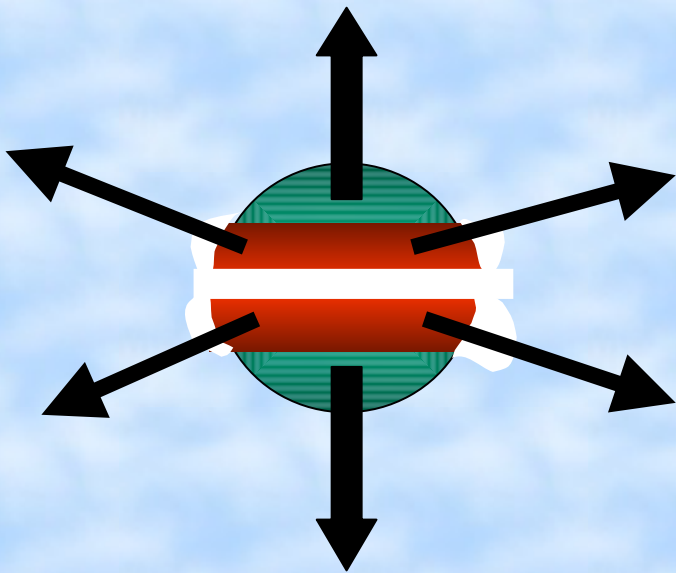
Cylinder Model

A_f Distribution



Better description

\square_f Distribution

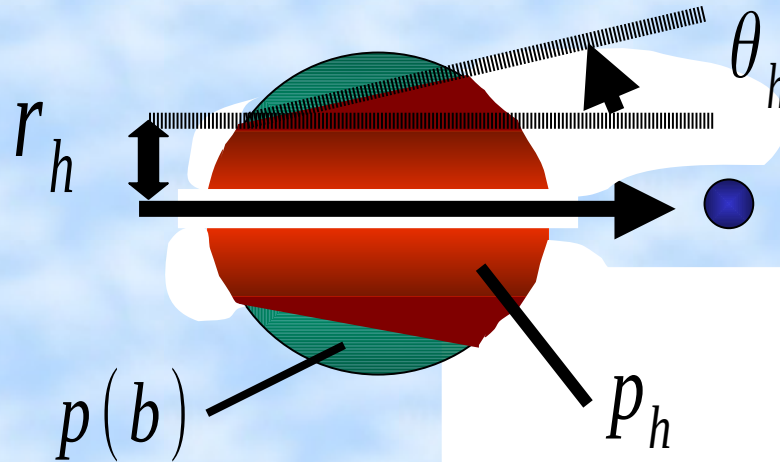
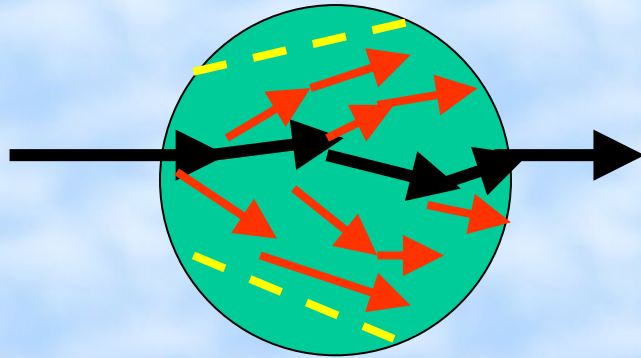


Sideward peak

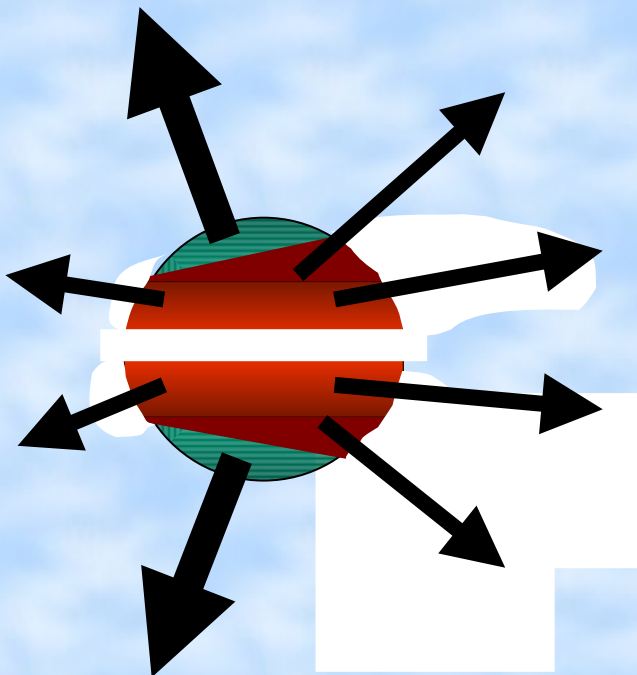
However backward enhancement cannot be explained in this Model.

2.2. Cone Model + SDM

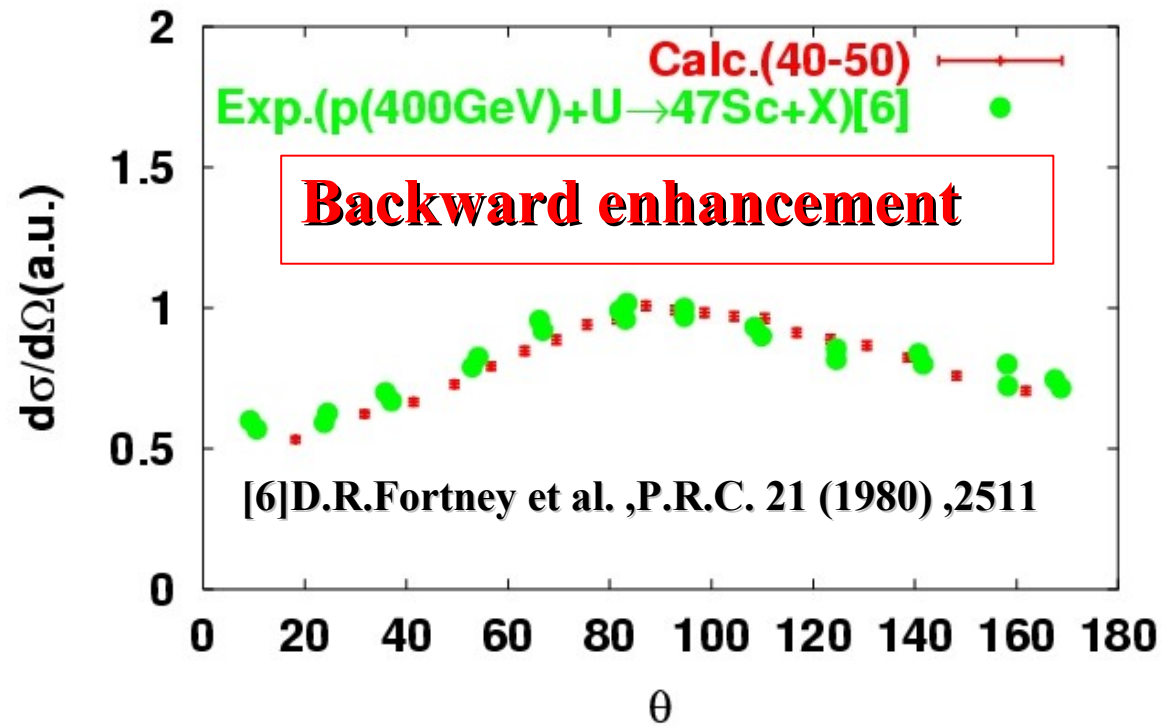
High T conic region



θ_f Distribution



Cone Model $\theta_h = 20^\circ, \beta = 0.005$



3. Summary

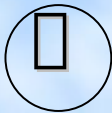
We investigate high-energy proton induced nuclear fragmentation by using the Percolation Model with non-equilibrium shape effects.

For Fragment Mass Distribution

We can reproduce fragment mass distribution well.

For Angular Distribution

The sideward peak and backward enhancement can be explained quantitatively.



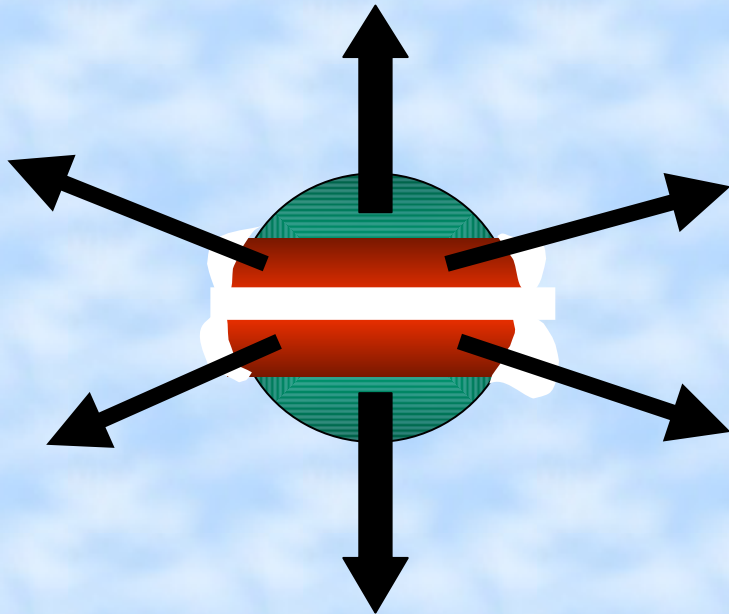
Coulomb repulsion from light fragments

This work implies the necessity to modify the step(2) in the standard picture of high-energy proton induced nuclear fragmentation.

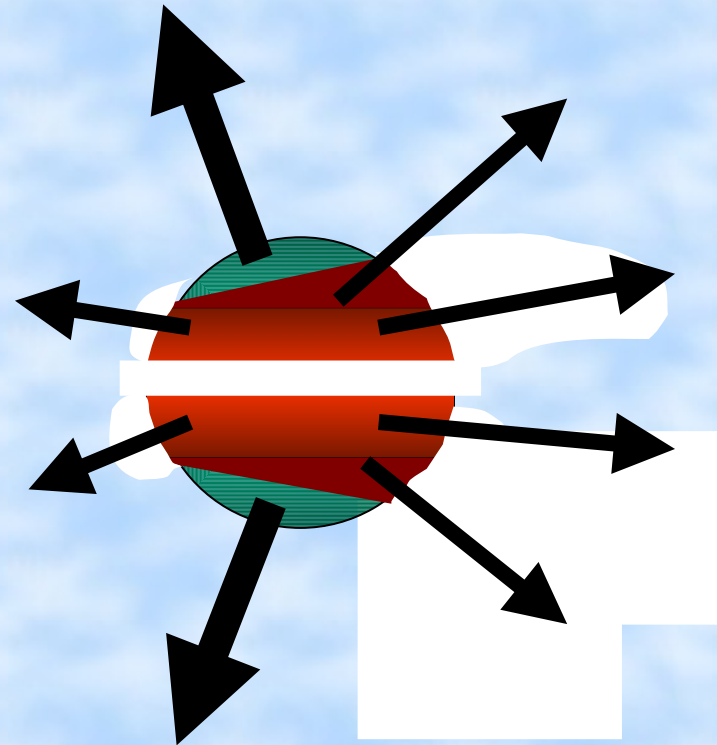
(2)Residual excited nucleus reaches equilibrium.

→ (2)' The excited nucleus breaks up into fragments from non-spherical shaped region before reaching equilibrium.

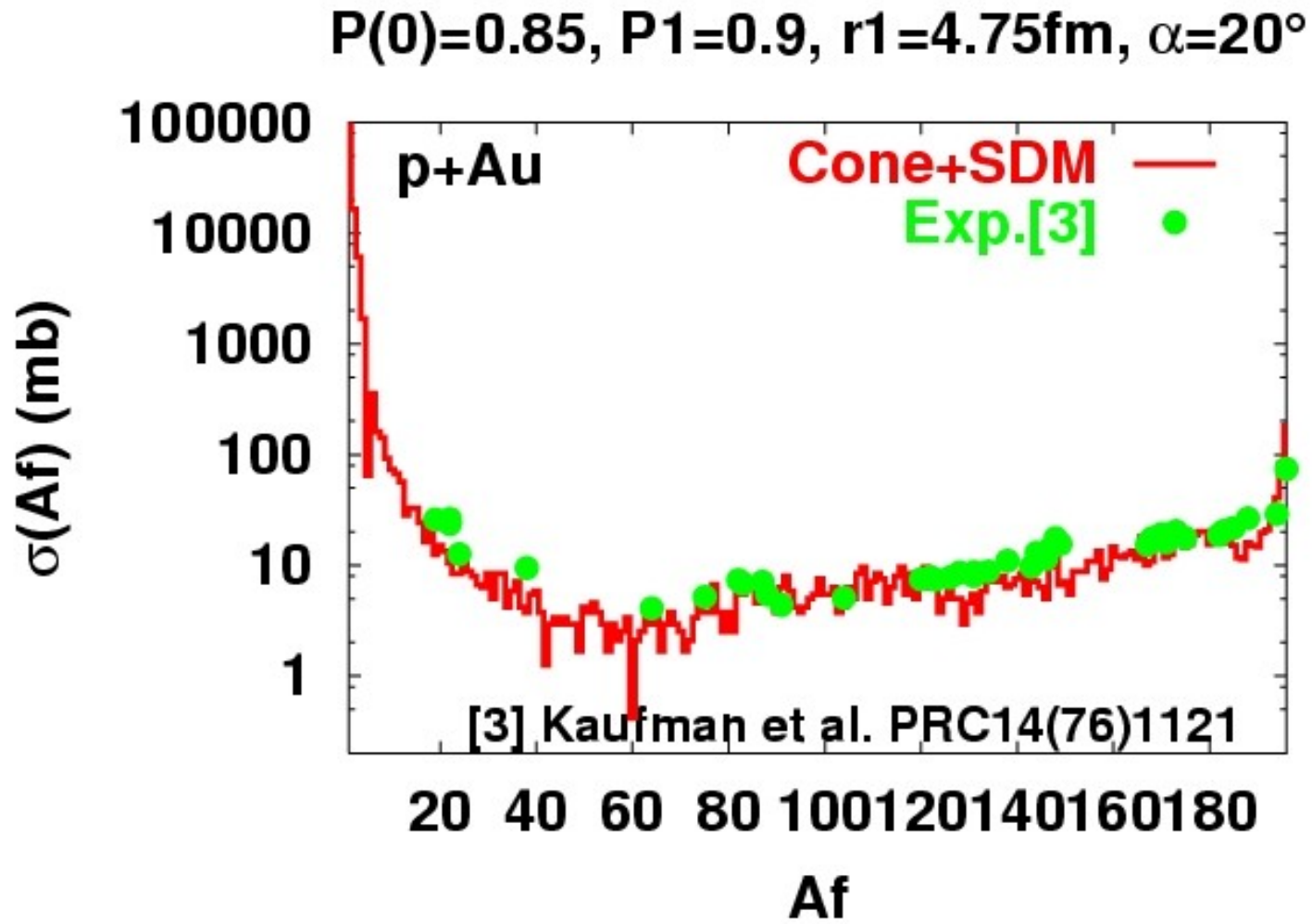
Cylinder model



Cone model

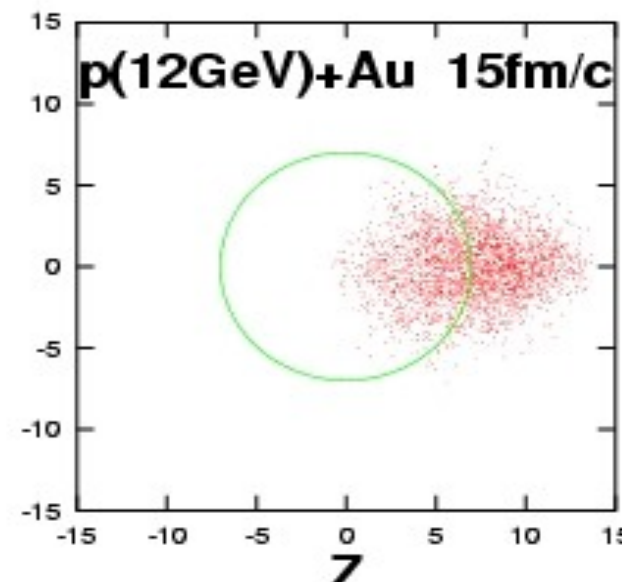
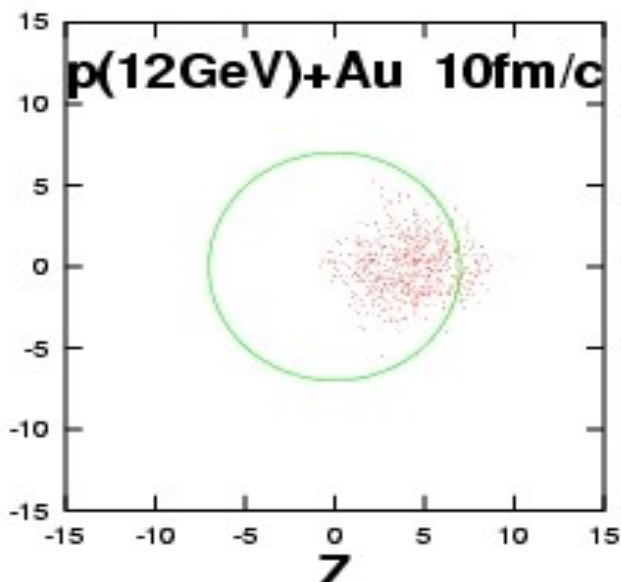
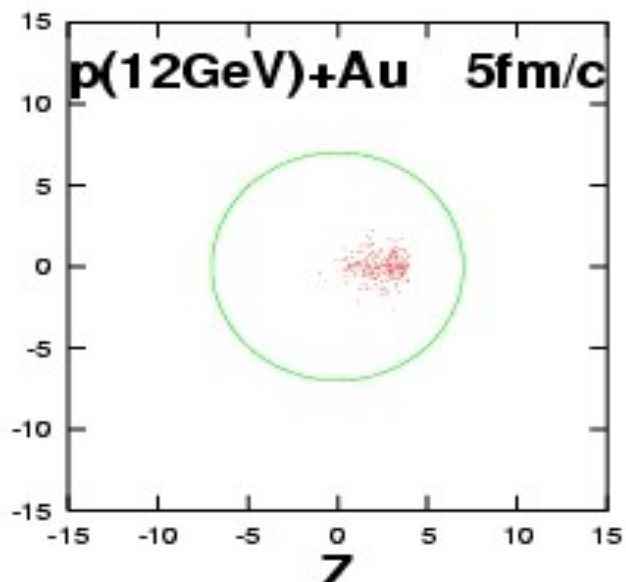
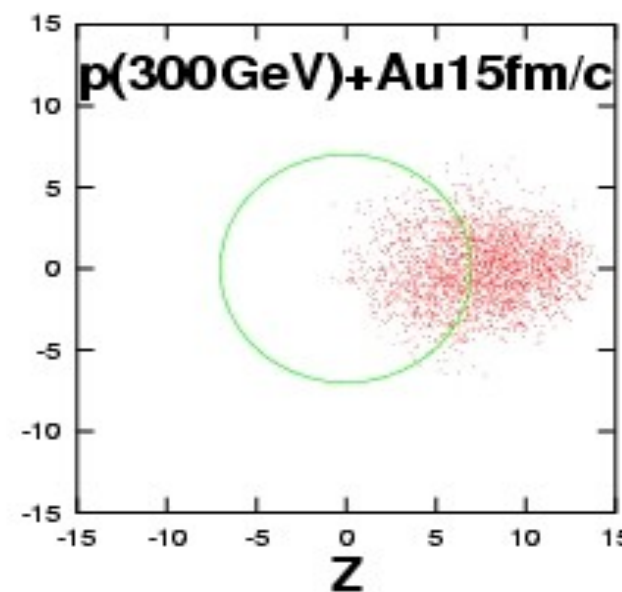
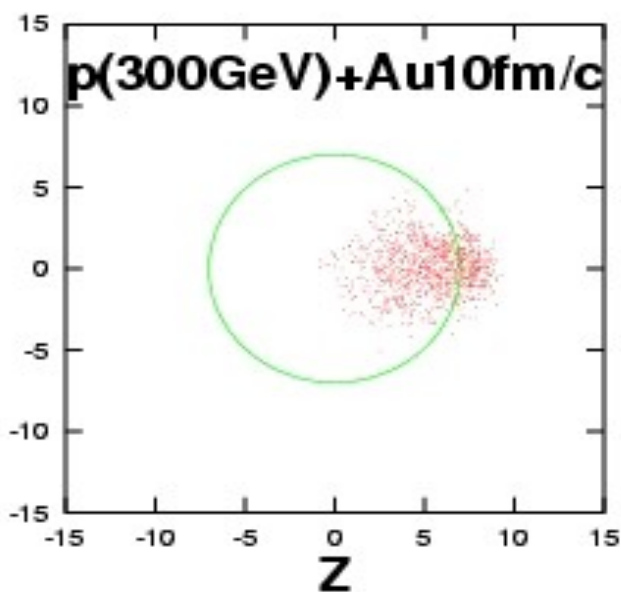
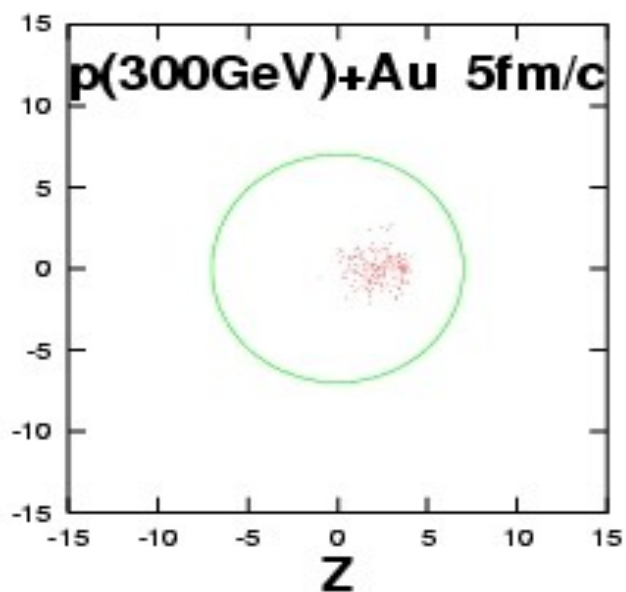
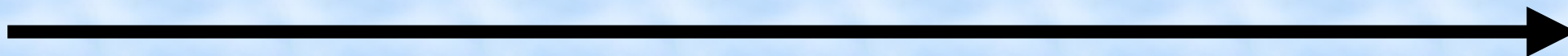


Fragment Mass Distribution of Cone Model



JAM Calculation

time



$$\xi \leq 1 - \exp(-T(b))$$

➔ High temperature region is created.

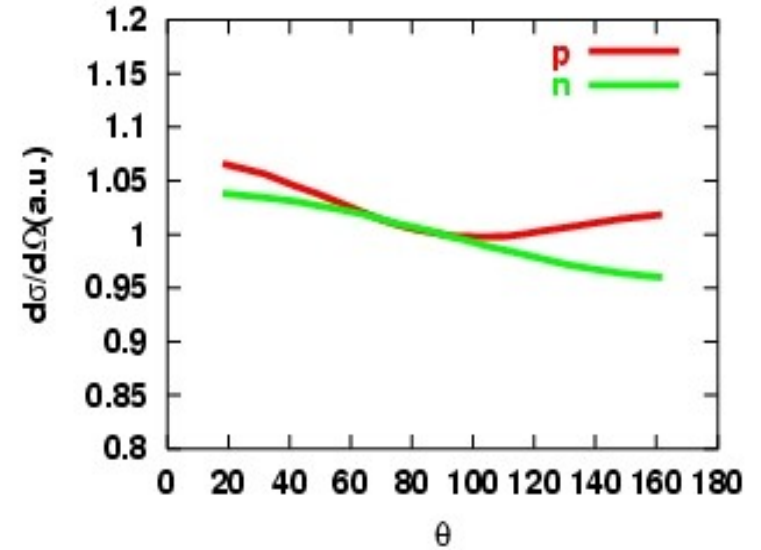
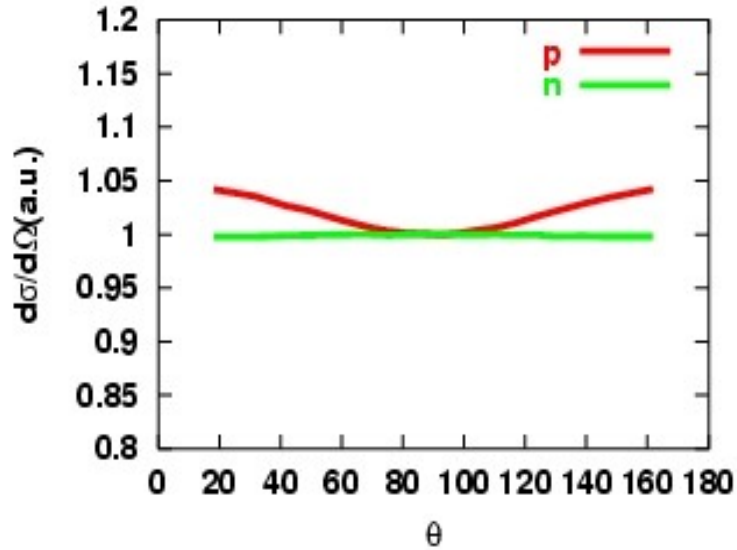
$$\xi \geq 1 - \exp(-T(b))$$

➔ High temperature region is not created.

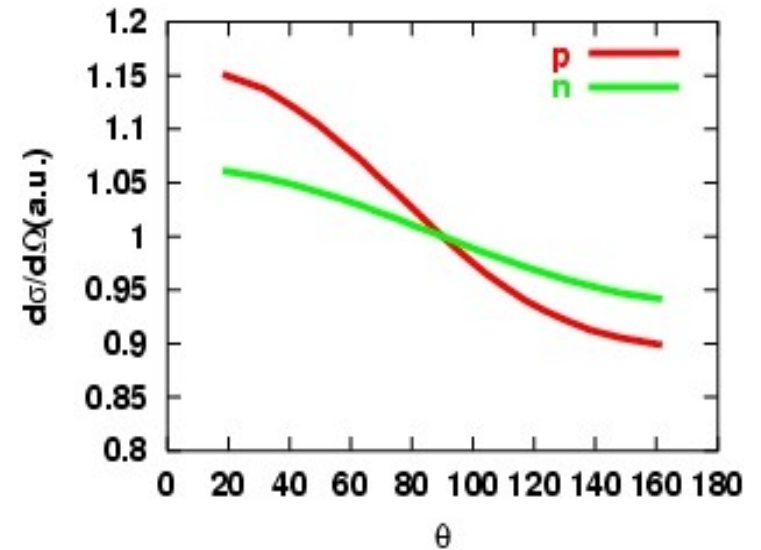
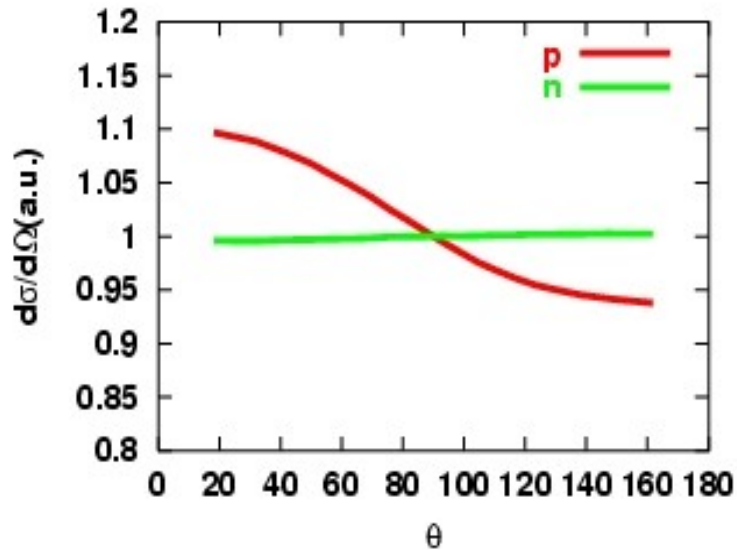
C.M. system

Lab. system

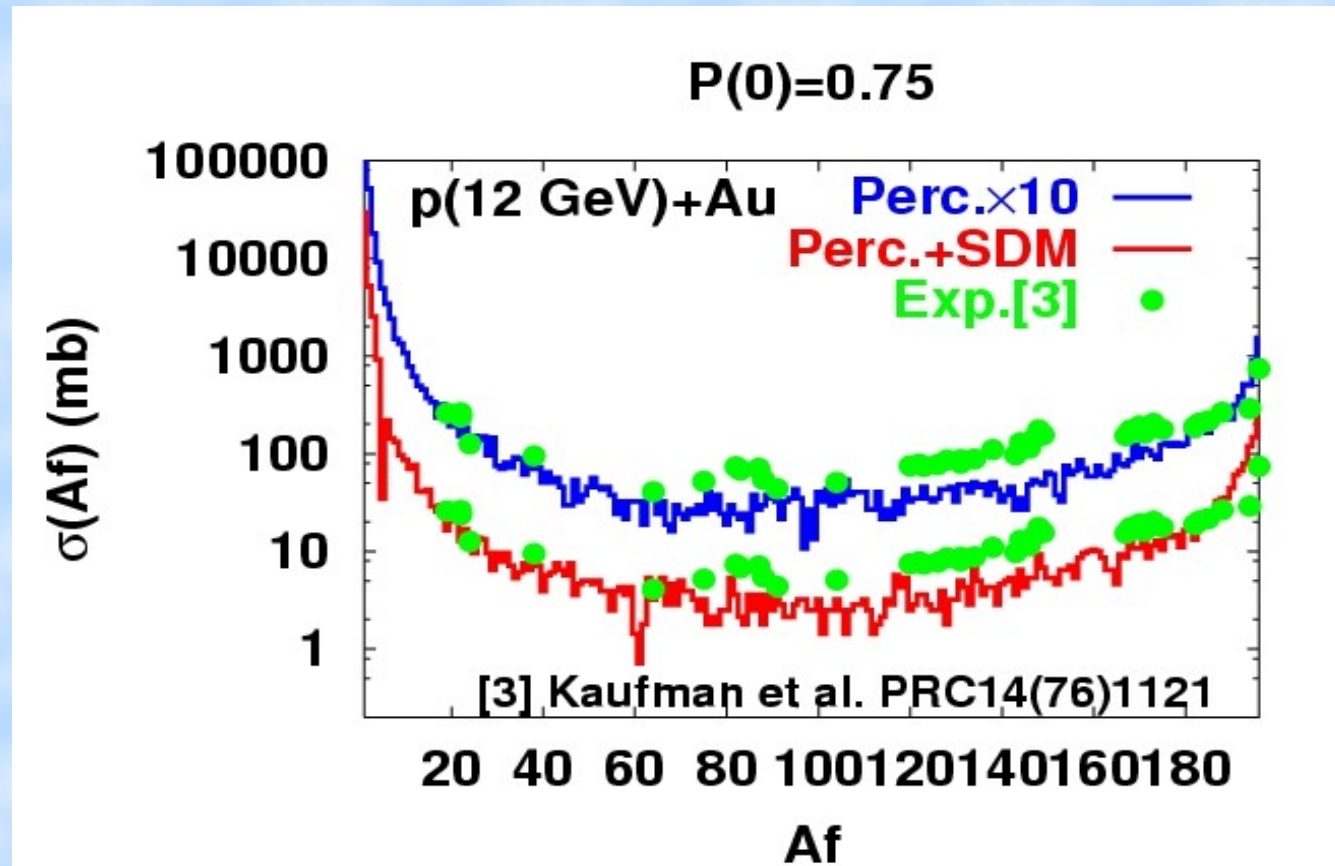
Cylinder



Cone

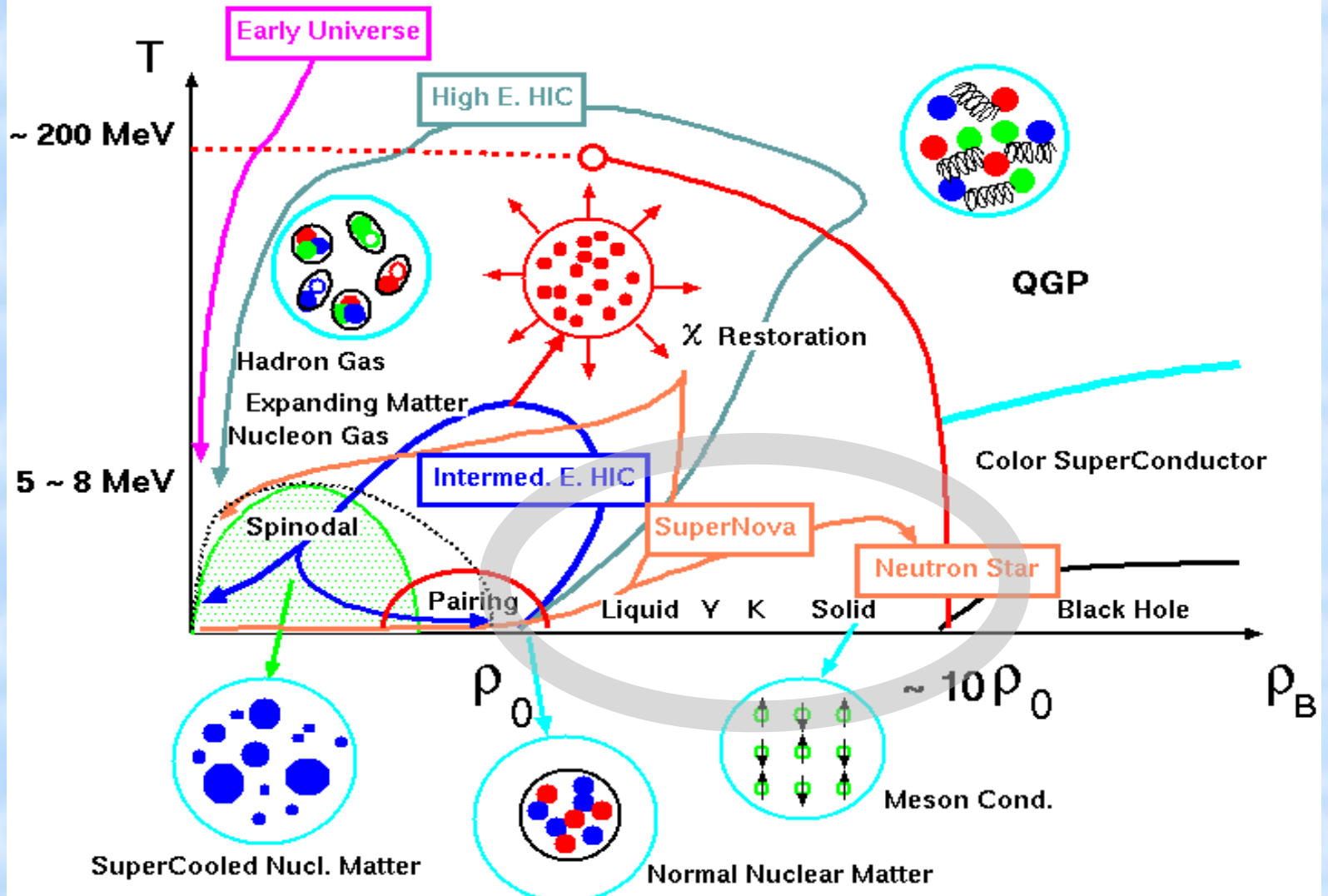


Perc. in comparison with Perc. + SDM



Part II:
Roles of Strangeness in Dense Matter

Hadronic Matter Phase Diagram



Contents

**Introduction — Roles of Hyperons in Dense Matter
What is Known, and What is Unknown ?
Several Recent Topics Related to Hyperons in Dense Matter
Summary**

What is Expected in the Neutron Star Core ?

Nucleon Superfluid ($^1S_0, ^3P_2$)

Pion Condensation

Hyperon Matter ← Strangeness

Tsuruta-Cameron (66), Langer-Rosen (70), Pandharipande (71), Itoh(75), Glendenning, Weber-Weigel, Sugahara-Toki, Schaffner-Mishustin, Balberg-Gal, Baldo et al., Vidana et al., Nishizaki-Yamamoto-Takatsuka, Kohno-Fujiwara et al., ...

Kaon Condensation ← Strangeness

Kaplan-Nelson(88), Forkel-Rho et al.(SUNY), Davidson-Miller, Claymans et al., Politzer-Wise, Miller et al., Muto-Tatsumi, Brown-Thorsson-Lee-Rho-Min, Fujii et al., Yabu et al, Maruyama et al., Ellis-Knorren-Prakashi (with Y), Li-Ning, Li-Brown, Tiwari-Prasad-Singh, Glendenning-Schaffner,

Quark-Gluon Plasma

We cannot understand Highly Dense Hadronic Matter without the Knowledges of Strangeness Nuclear Physics

Why is Strangeness important in Dense Matter ?

Constituents:

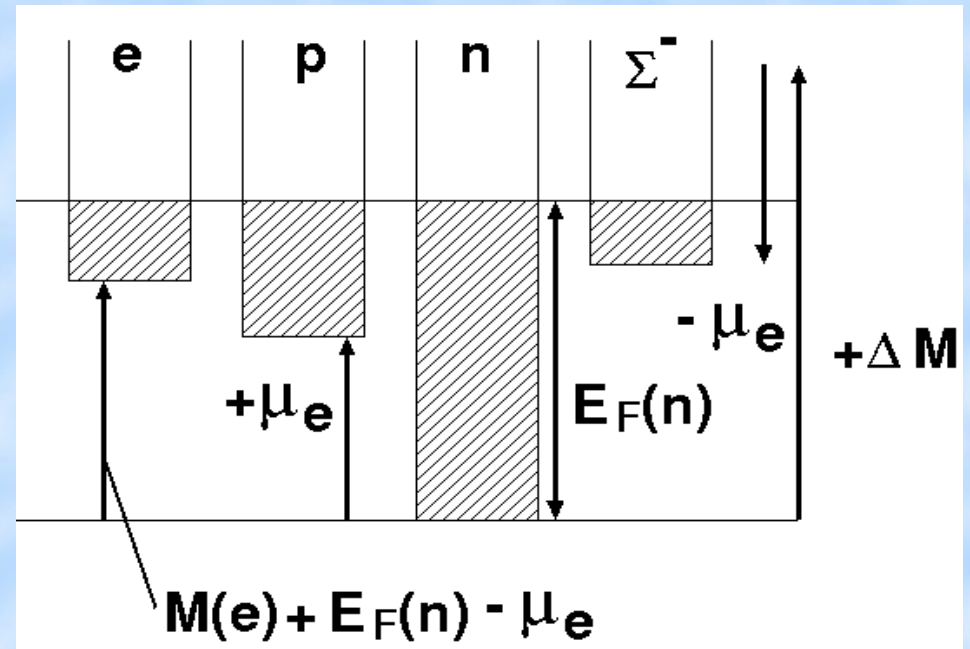
$$p, n, e^{\pm}, \mu^{\pm}, \Lambda, \Sigma^{\pm,0}, \dots$$

Chemical Equilibrium:

× Strangeness (Weak)

× Lepton (ν Emission)

$$\mu_i = B_i \mu_B + Q_i \mu_Q$$

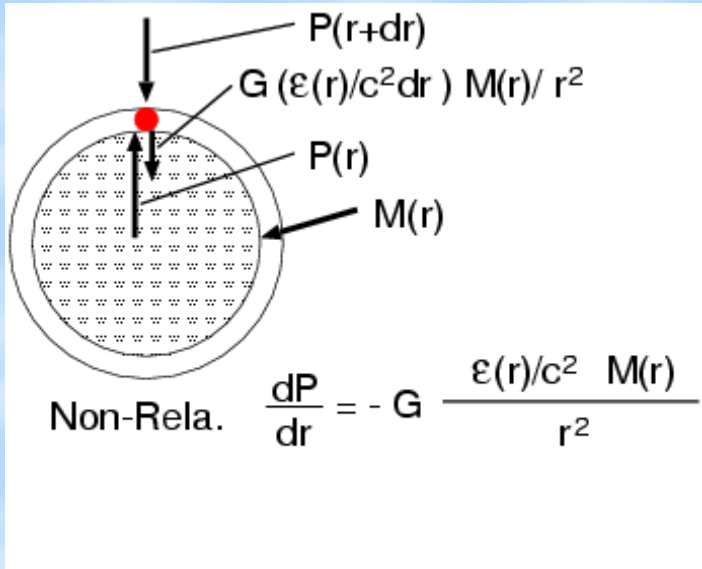


Negatively Charged or Neutral Baryons are Favored

$$E_F^*(n) + U(n) + \mu_e = M^*(\Sigma^-) + U(\Sigma^-) \quad \Sigma \text{ appears}$$

$$E_F^*(n) + U(n) = M^*(\Lambda) + U(\Lambda) \quad \Lambda \text{ appears}$$

TOV Equation: Balance of Pressure and Gravitation



$$\frac{dP}{dr} = -G \frac{(\epsilon/c^2 + P/c^2)(M + 4\pi r^3 P/c^2)}{r^2(1 - 2GM/rc^2)}$$

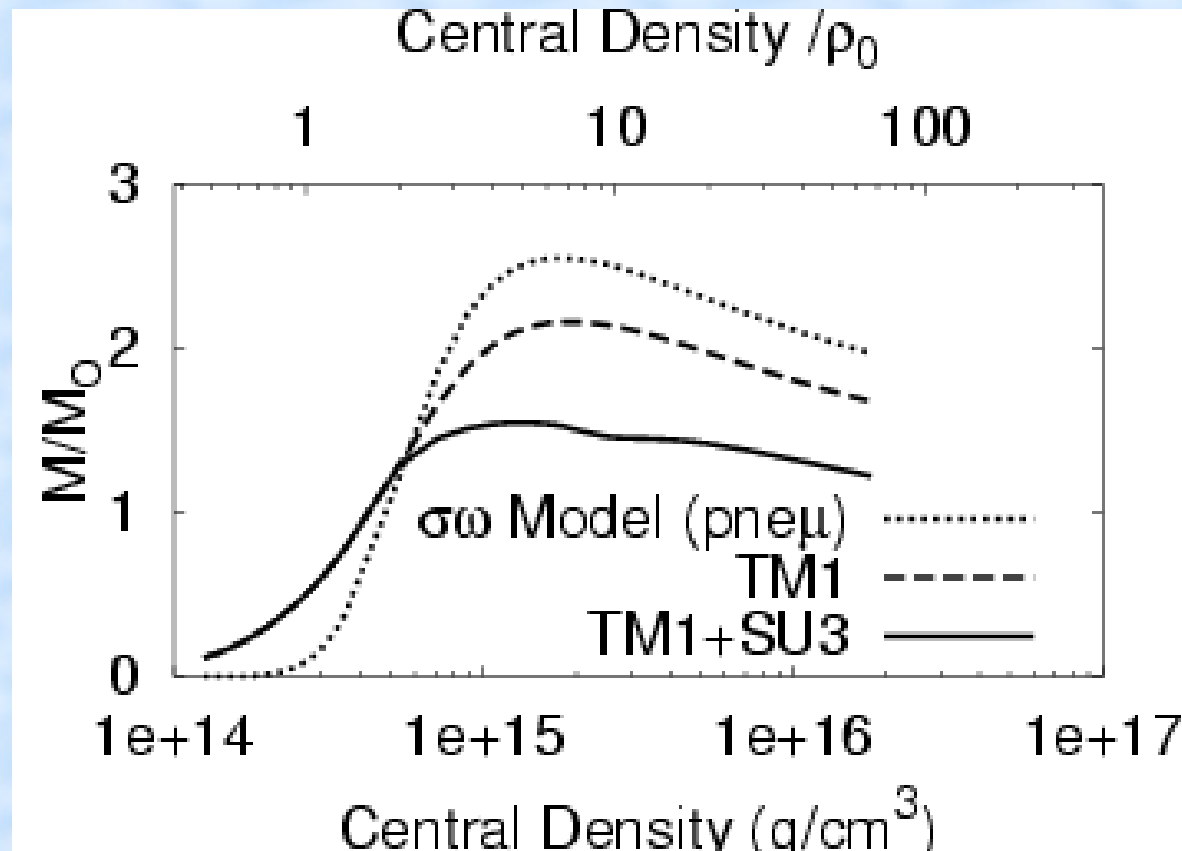
$$\frac{dM}{dr} = 4\pi r^2 \epsilon/c^2, \quad \frac{dP}{dr} = \frac{dP}{d\epsilon} \frac{d\epsilon}{dr}$$

$$P = P(\epsilon), \quad \frac{dP}{d\epsilon} = \frac{dP}{d\epsilon}(\epsilon) \quad (\text{EOS})$$

Neutron Star Mass = $M(R)$, where $P(R) = 0$

When You Make a New EOS, Please Check Neutron Star Mass !

Neutron Star Max. Mass



A. Isshiki, AO, JPS @ Akita; Serot-Walecka ($\sigma\omega$);
Sugahara-Toki (TM1); Schaffner-Mishustin (TM1+SU3); Glendenning, ...

Maximum Mass Reduction $\sim 0.5-1.0 M_{sun}$

Σ Potential Effects on Neutron Star Matter

Potential for Λ ; Relatively Well Known

$$U(\Lambda) \sim -30 \text{ MeV} \quad (\text{Many Single Hypernuclei})$$

Potential for Ξ ; Recently Suggested from (K^-, K^+) Experiments

$$U(\Xi) \sim -(14 - 16) \text{ MeV}$$

(KEK-E224, BNL-E885, BNL-E906)

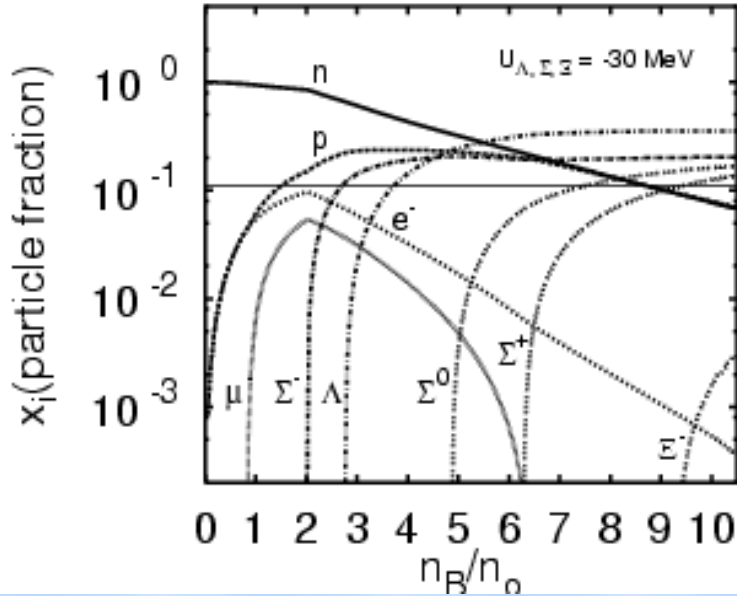
\rightarrow Potential Depth \propto Number of ud Quarks ?

Potential for Σ : Contradicting Conjectures

$$U(\Sigma) \sim -(24 - 30) \text{ MeV} \quad (\text{Old Conjectures})$$

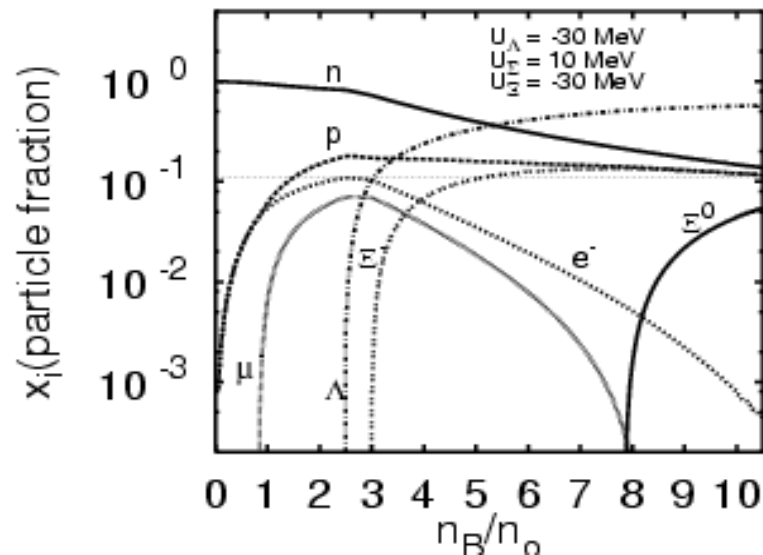
$$U(\Sigma) > 0$$

(Dabrowski, Yamamoto et al., Kohno-Fujiwara et al.)



Attractive Potential for Σ
 $\rightarrow \Sigma$ appears at around

$$\rho \approx 2\rho_0$$



Repulsive Potential for Σ
 $\rightarrow \Sigma$ does not appear

(RMF: Sahu, Ohnishi Nucl. Phys. A691 (2001), 439.)

What is Already Known ?

Light Single Λ Hypernuclear Shell/Cluster Structure

Bare Λ N Interaction

Germanium γ -ray Detector(Tamura et al.)

+Precise Few-Body Calculation (Hiyama et al., Nemura et al.)

Structure of $^4_{\Lambda}\text{He}$:

Coherent $\Lambda\Sigma$ Coupling

(Harada-Akaishi-Shinmura-Myint, Hiyama et al.)

$\Lambda\Lambda$ Interaction in Nuclei = Weakly Attractive

Recent Experiment KEK-E373 (Nagara Event)

What is Still Unknown ?

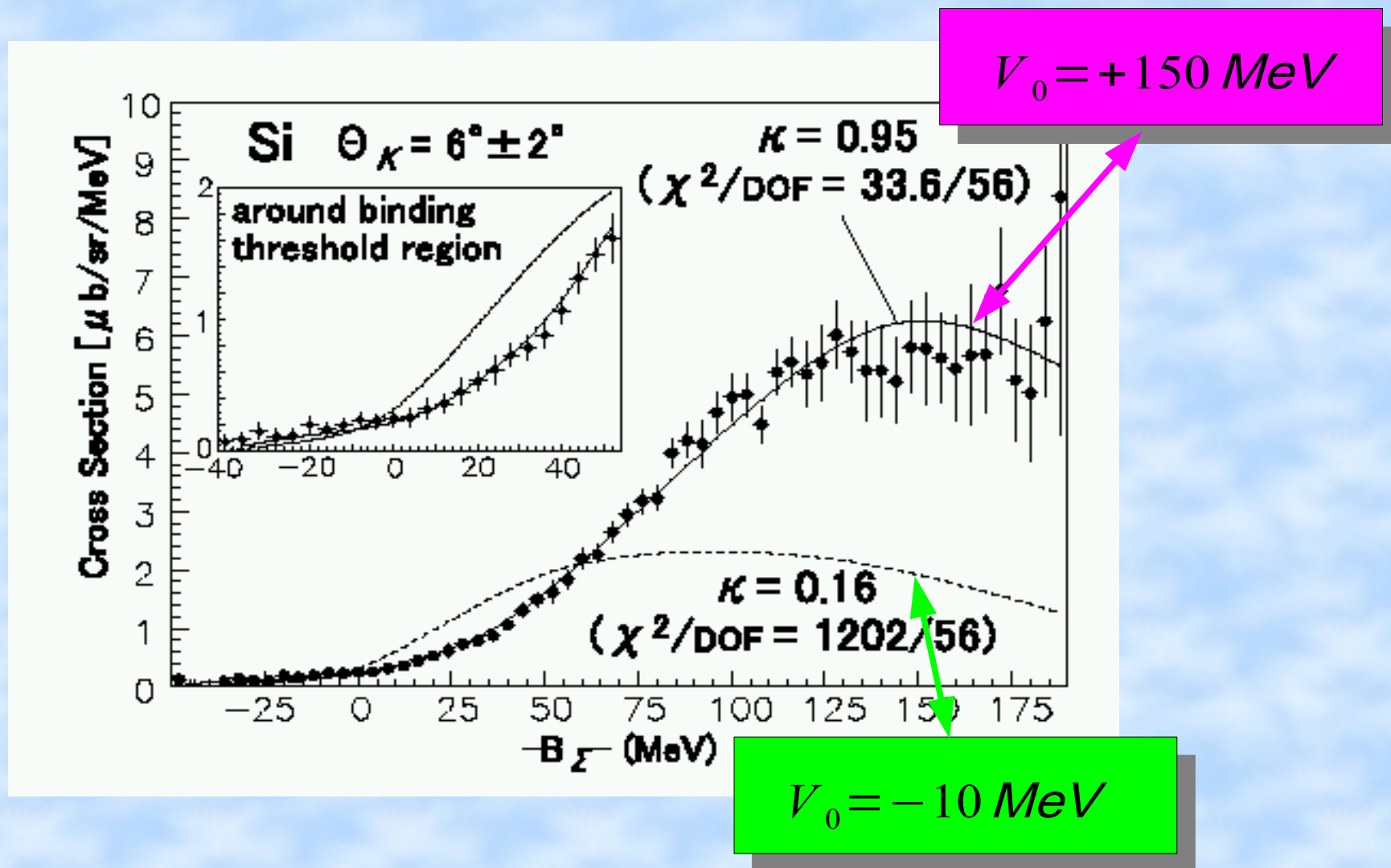
- ◆ **Properties of Hyperons (All) at Higher Densities.**
- ◆ **Σ Potential at ρ_0 and Higher Densities**
- ◆ **$\Lambda\Lambda$ Interaction in “Free” Space**

➡ Very Recent Experiments !

Direct Quasi-Free Production of Σ (Noumi et al.)
Strangeness Enhancement in HIC at SPS (NA49)

Does Σ Feel +150 MeV (Repulsive) in Nuclei ?

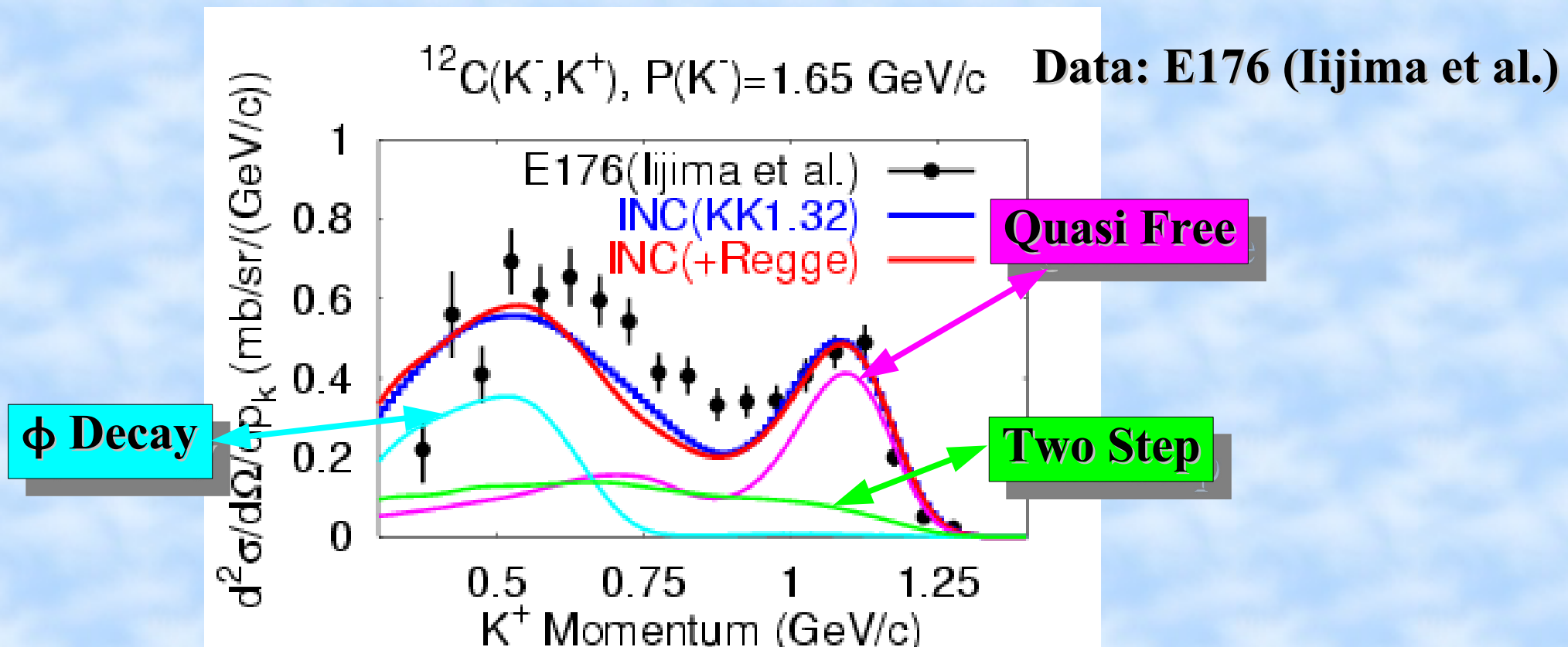
Noumi et al., Phys. Rev. Lett. 89 (2002) , 072301.



No Theoretical Model Support $V_0 = +150 \text{ MeV}$! \rightarrow Big Puzzle !!

Multi-Step Effects ?

In the case of (K^-, K^+) Reaction,
Multi-Step Effects was Important to Understand
Highly Excited Region.

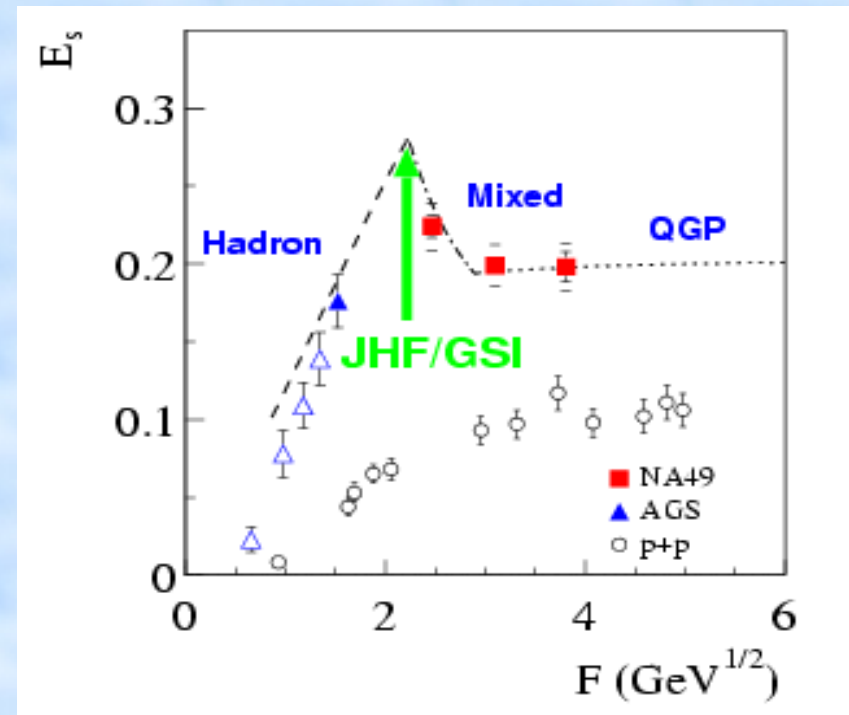
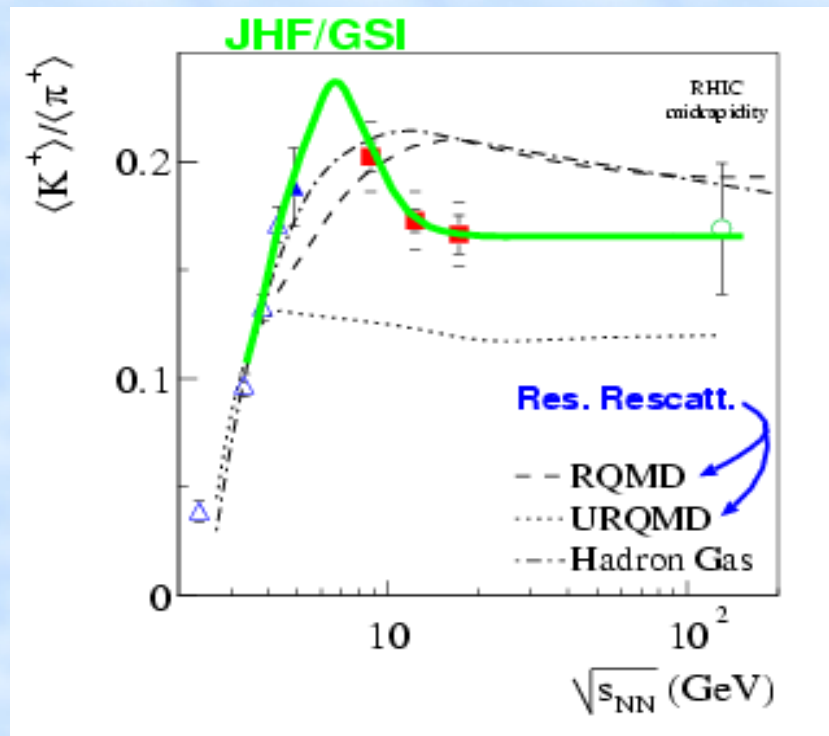


Is it also Important in (π^-, K^+) Reactions ?

Strangeness Enhancement: Rescattering, Potential, or Phase Transition ?

Strangeness is Enhanced Sharply at $E_{inc} = 10 \text{ } \square \text{ } 40 \text{ GeV/A !}$

NA49 (nucl-ex/0205002)



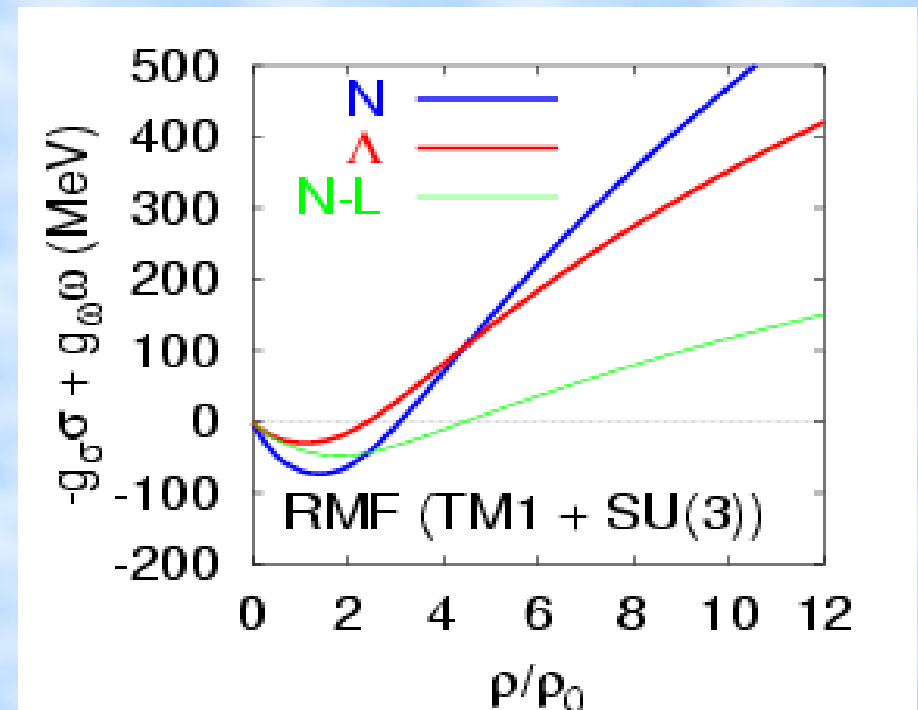
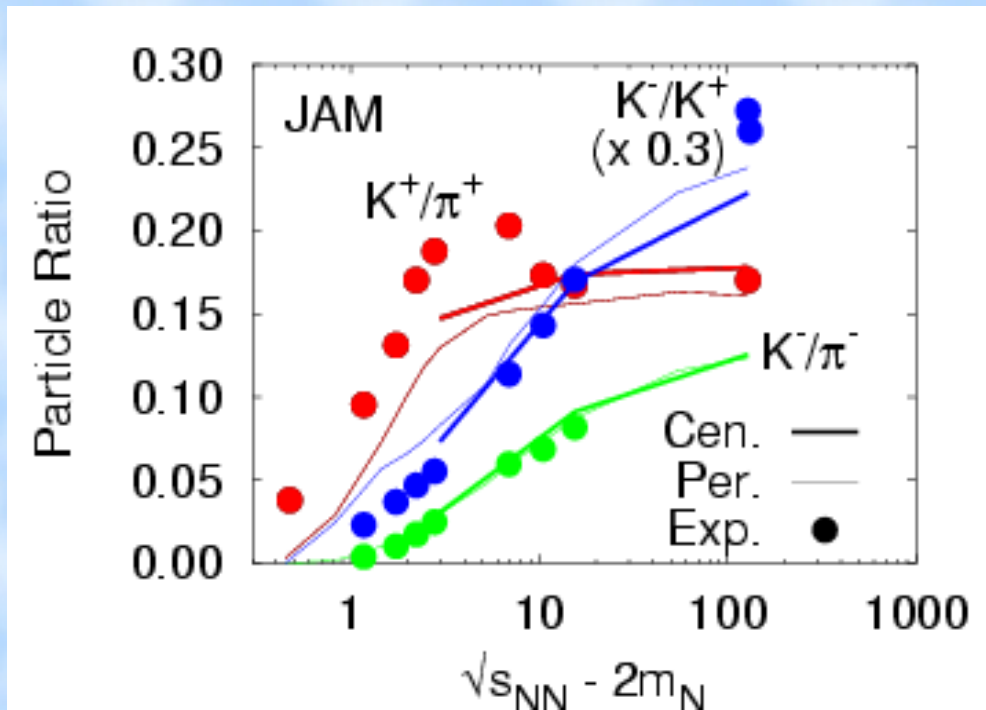
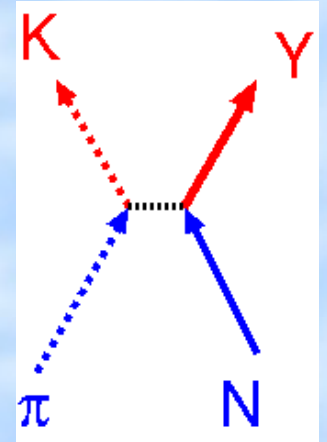
JHF Energy: \sim Maximum K/π ratio

Does Hyperon Potential Help It ?

Rescattering of Resonances/Strings (RQMD)

Baryon Rich QGP Formation

High Baryon Density Effect (Associated Prod. of Y)



At $\rho > 5 \rho_0$ Hyperon Feels More Attractive Potential than N

Is Lambda-Lambda Interaction Really Weak ?

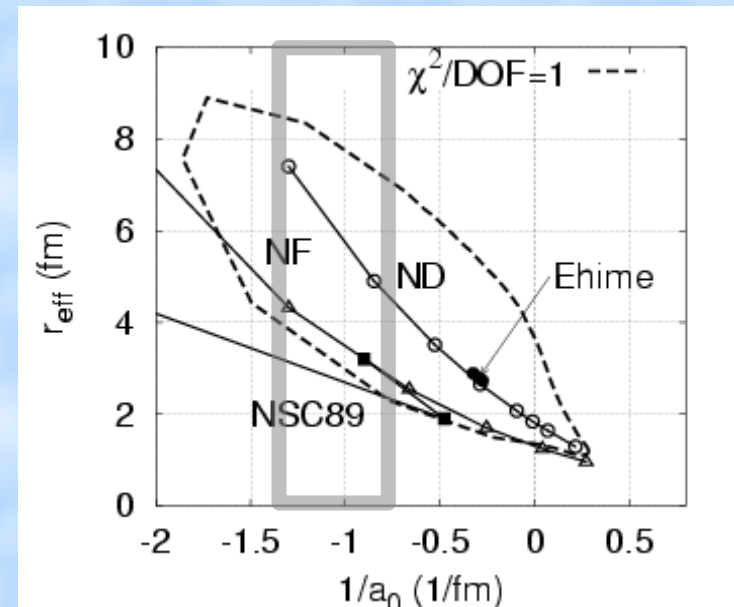
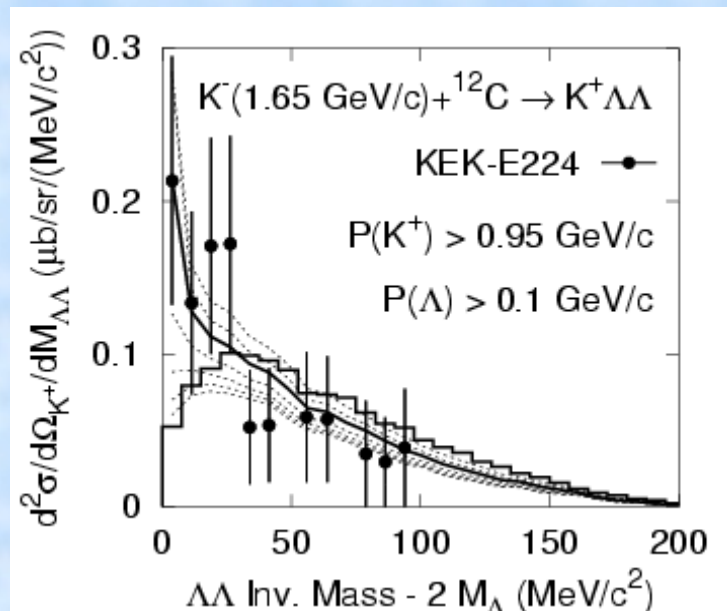
Khin Swe Myint, Shinmura, Akaishi, Euro. Phys. J. A16 (2003) 21.

From Nagara Event,

$a_{\Lambda\Lambda} = 0.7 \text{ fm}$ (Weak $\Lambda\Lambda - \Xi N$ Coupling)

$\sim 1.3 \text{ fm}$ (Strong Coupling, Pauli Suppressed in Nuclei)

Momentum Correlation of $\Lambda\Lambda$



We have the range of overlap !

Summary

1. **Strangeness is important in dense matter such as in neutron star core.**

Strangeness changes the max. mass of neutron star, modifies the order of QCD phase transition, probes deeply inside the nucleus, mixes elementary particles in nuclei.

2. **Hypernuclear spectroscopy have developed a lot in these years, but we need more data for the understanding of dense matter.**

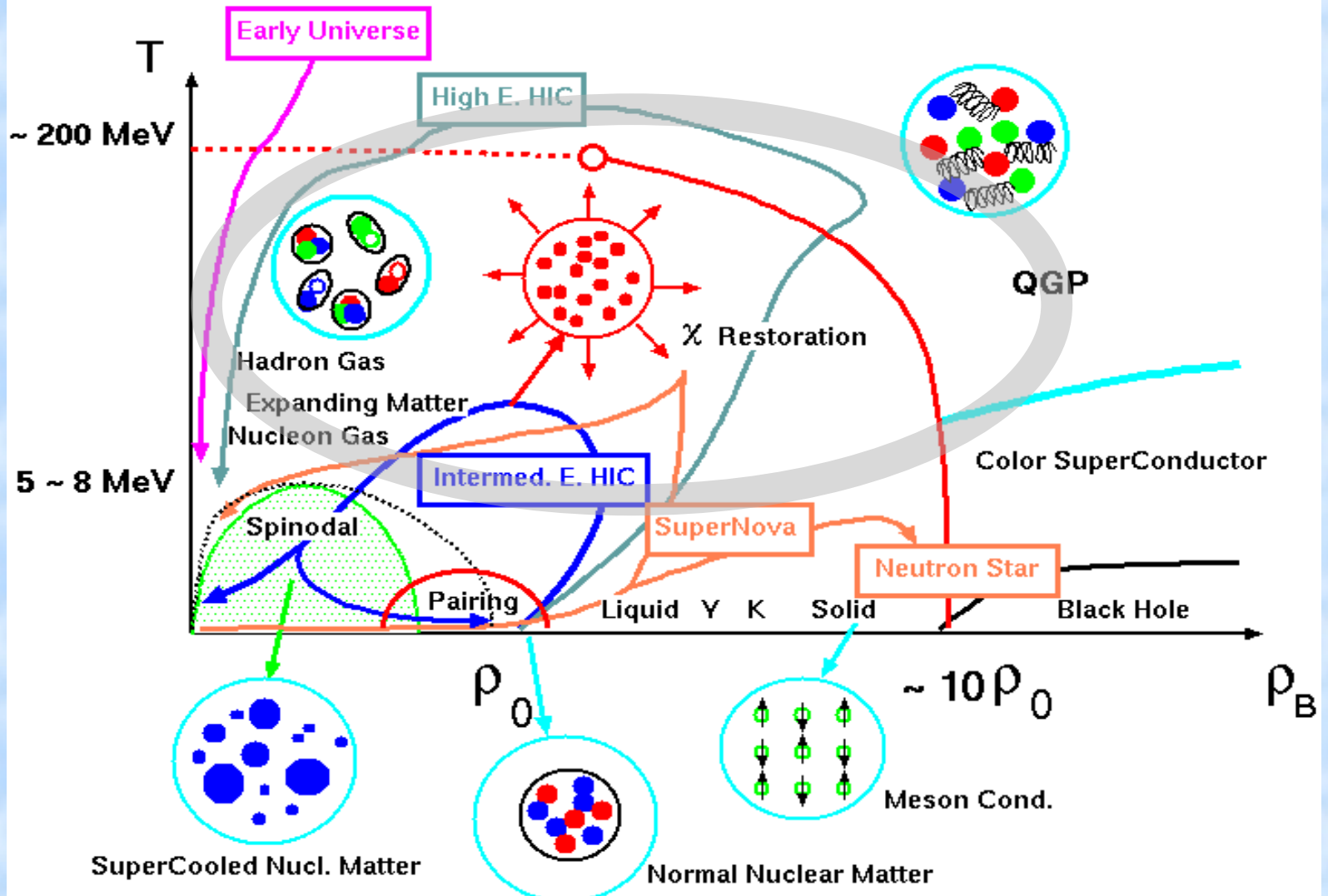
Σ Potential, $\Lambda\Lambda$ Interaction, ΛN - ΣN and $\Lambda\Lambda$ - ΞN Coupling, Hyperon Potential in Dense Matter,

3. **Recent Data would be Helpful to Understand Hyperons in Dense Matter based on *Real Data***

Quasi Free Σ Production, Kaon Enhancement, $\Lambda\Lambda$ Nuclei, $\Lambda\Lambda$ Correlation,

Part III:
**High Energy Heavy-Ion Collisions,
and Equation of State of Hot / Dense Matter**

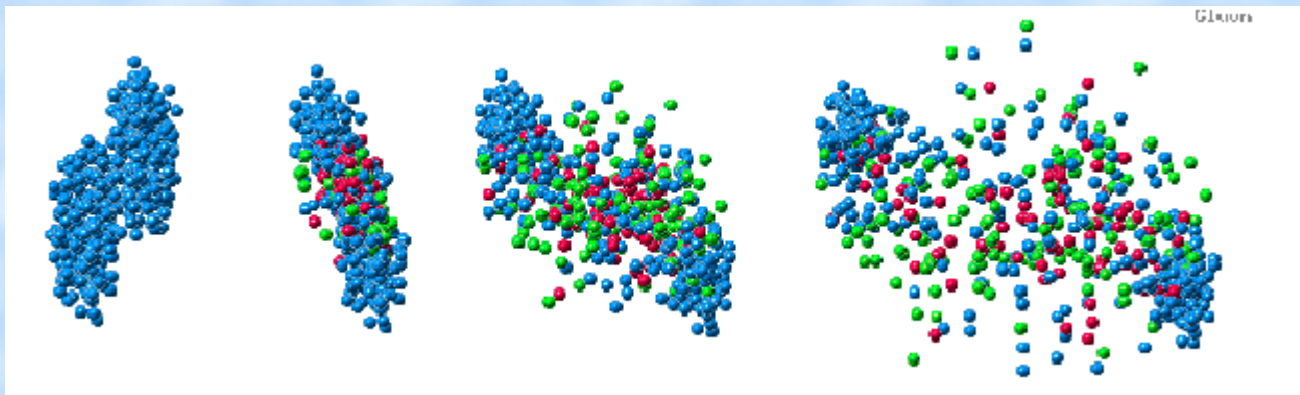
Hadronic Matter Phase Diagram



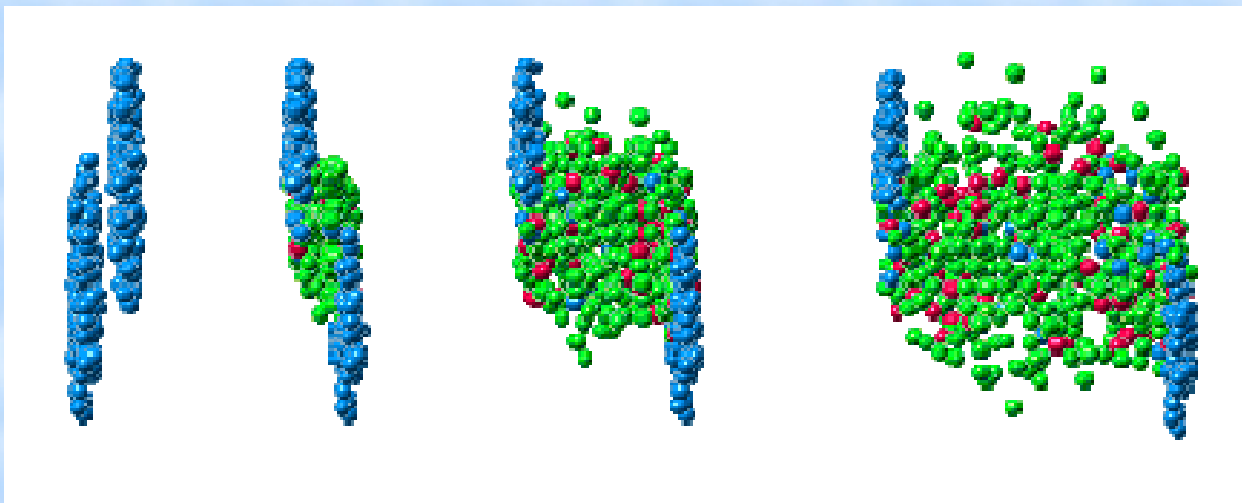
JAMming on the Web

<http://nova.sci.hokudai.ac.jp/~ohtsuka/>

AGS

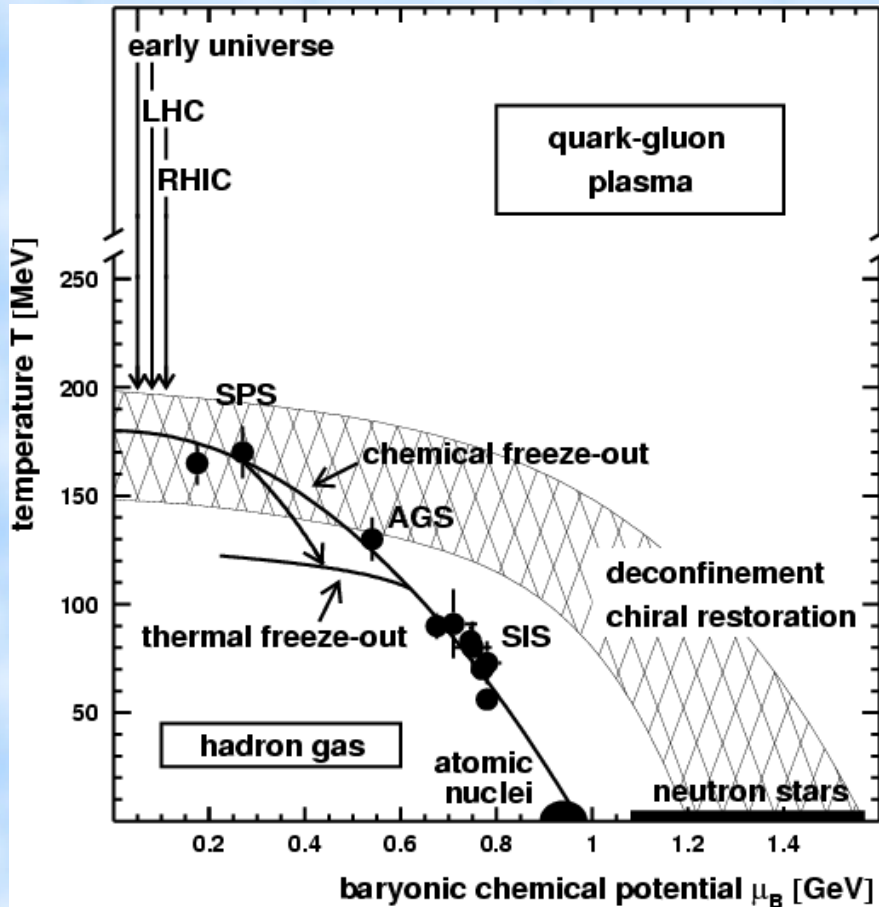


SPS

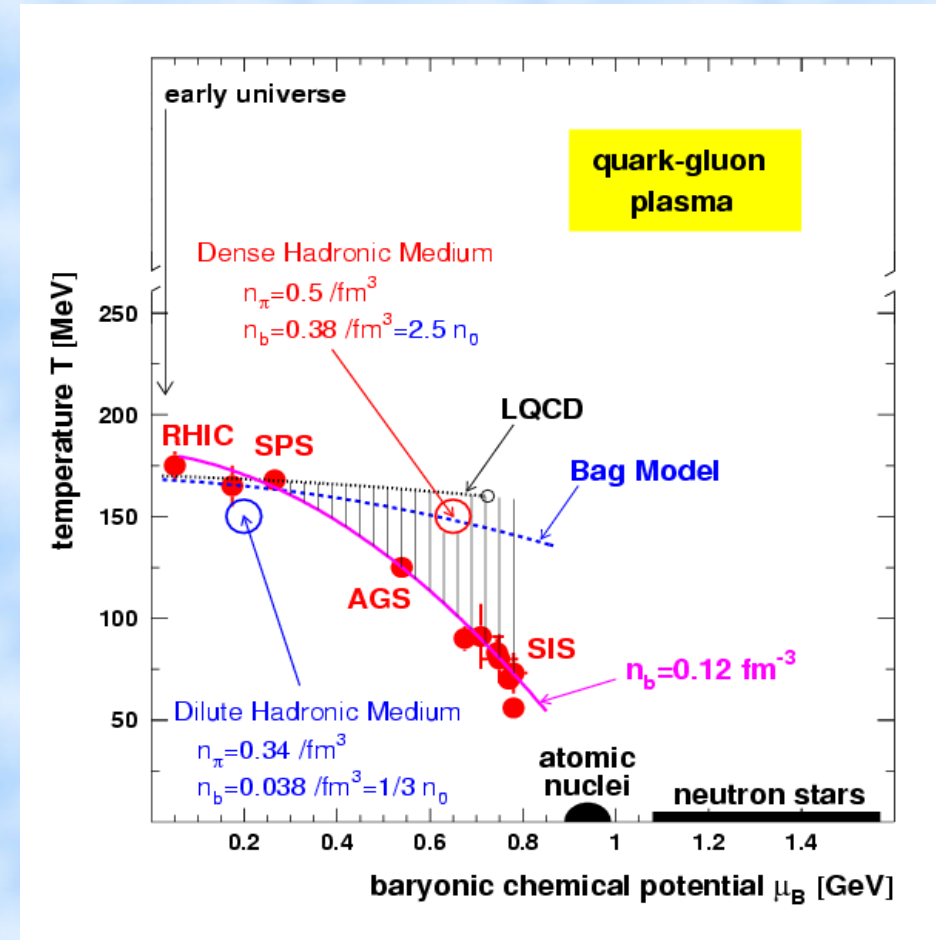


Experimentally Estimated Phase Diagram

Chemical Freeze-Out Points in High-Energy Heavy-Ion Collisions



1998 (J. Stachel et al.)

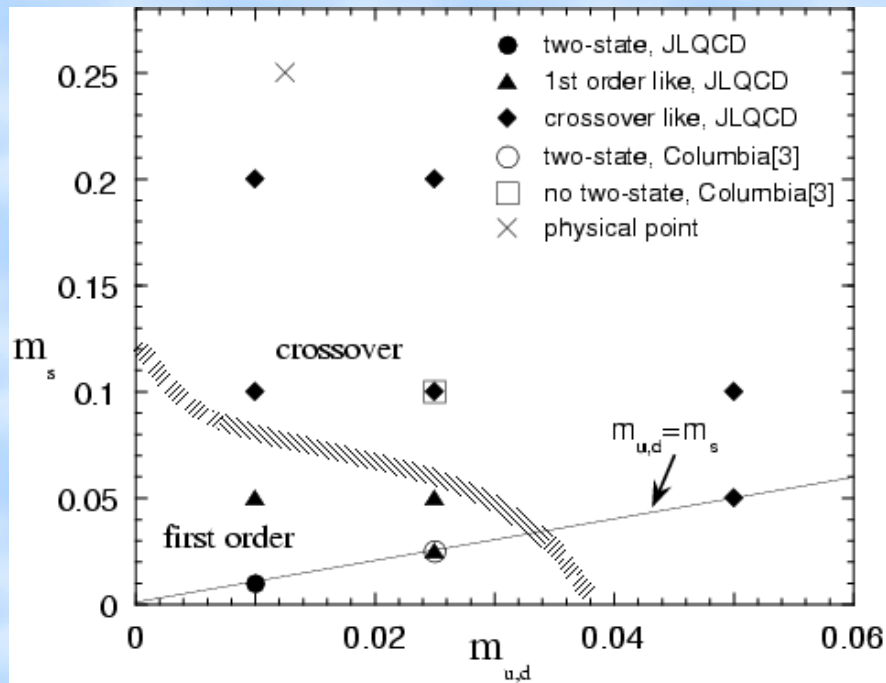


2002 (Braun-Munzinger et al.
J. Phys. G28 (2002) 1971.)

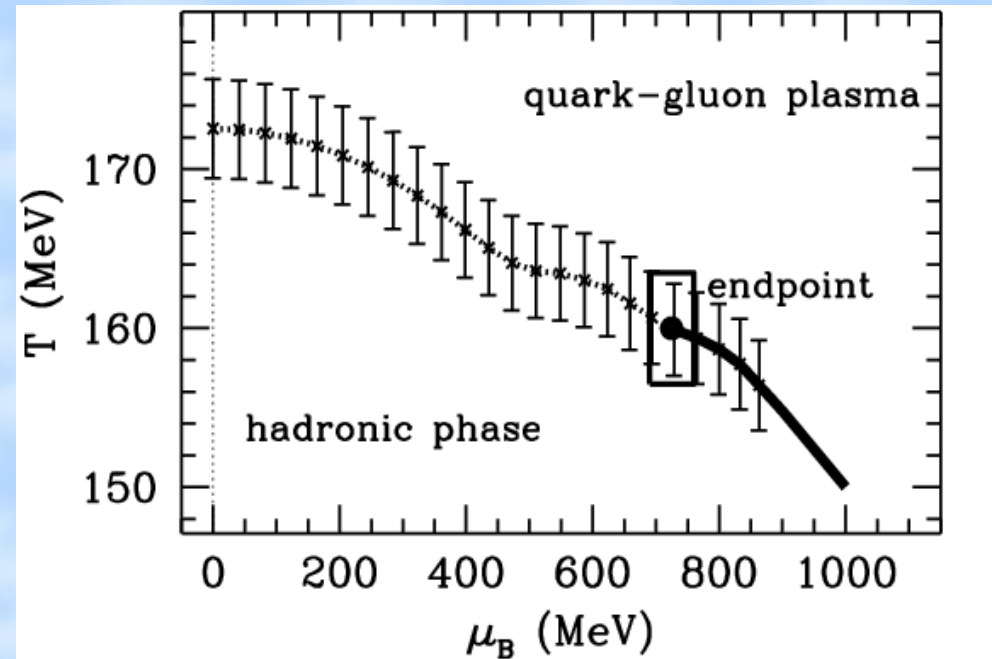
Chem. Freeze-Out Points are very Close to
Expected QCD Phase Transition Boundary

Theoretically Expected QCD Phase Diagram

Zero Chem. Pot.



Finite Chem. Pot.



JLQCD Collab. (S. Aoki et al.),
Nucl. Phys. Proc. Suppl. 73 (1999) 459.

Finite \square : Fodor & Katz,
JHEP 0203 (2002), 014.

Zero Chem. Pot. : *Cross Over*
Finite Chem. Pot.: *Critical End Point*

Is QGP Formed at AGS, SPS and/or RHIC ?

Proposed and/or Measured Signals

Collective Flow (AGS, SPS, RHIC)

EOS modification / Thermalization Degree

Low-Mass Lepton Pair (Yes @ SPS, Not Yet @ RHIC)

Partial Restoration at High Temperature/Density

High-Mass Lepton Pair (Yes @ SPS, Preliminary @ RHIC)

J/Ψ Suppression at High Temperature

Jet Energy Loss (@ RHIC)

Parton Dynamics at High (Freed) Gluon Density

Strangeness Enhancement (Yes @ AGS, Lower E. SPS, No @ RHIC)

Rescattering or Potential at High Density or QGP

Later on, I mainly Discuss *Collective Flows*

What is Collective Flow ?

(Directed) Flow (dP_x/dY)

Stiffness (Low E)
+ Time Scale (High E)

Elliptic Flow (V_2)

Thermalization
& Pressure Gradient

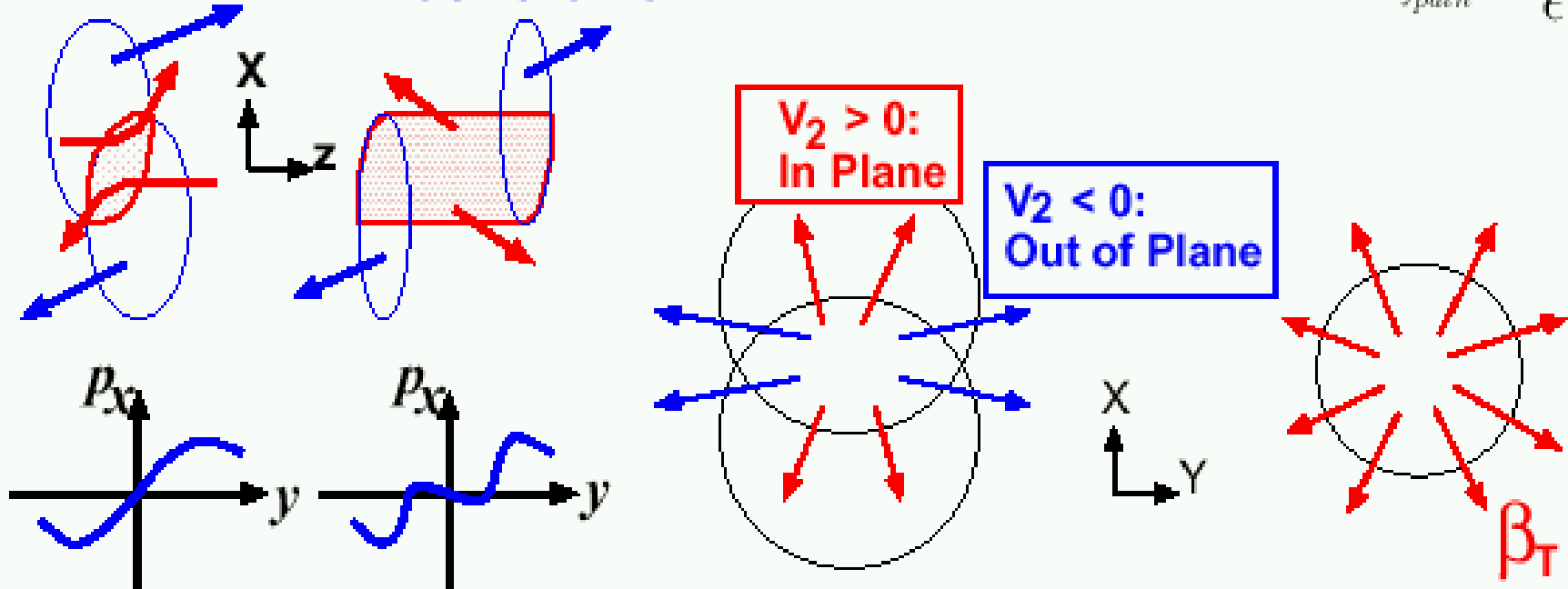
Radial Flow (β_T)

Pressure History

$$\epsilon \frac{DV}{Dt} = -\nabla P$$

$$\rightarrow V = \int_{\text{path}} \frac{-\nabla P dt}{\epsilon}$$

Until AGS **Above SPS**

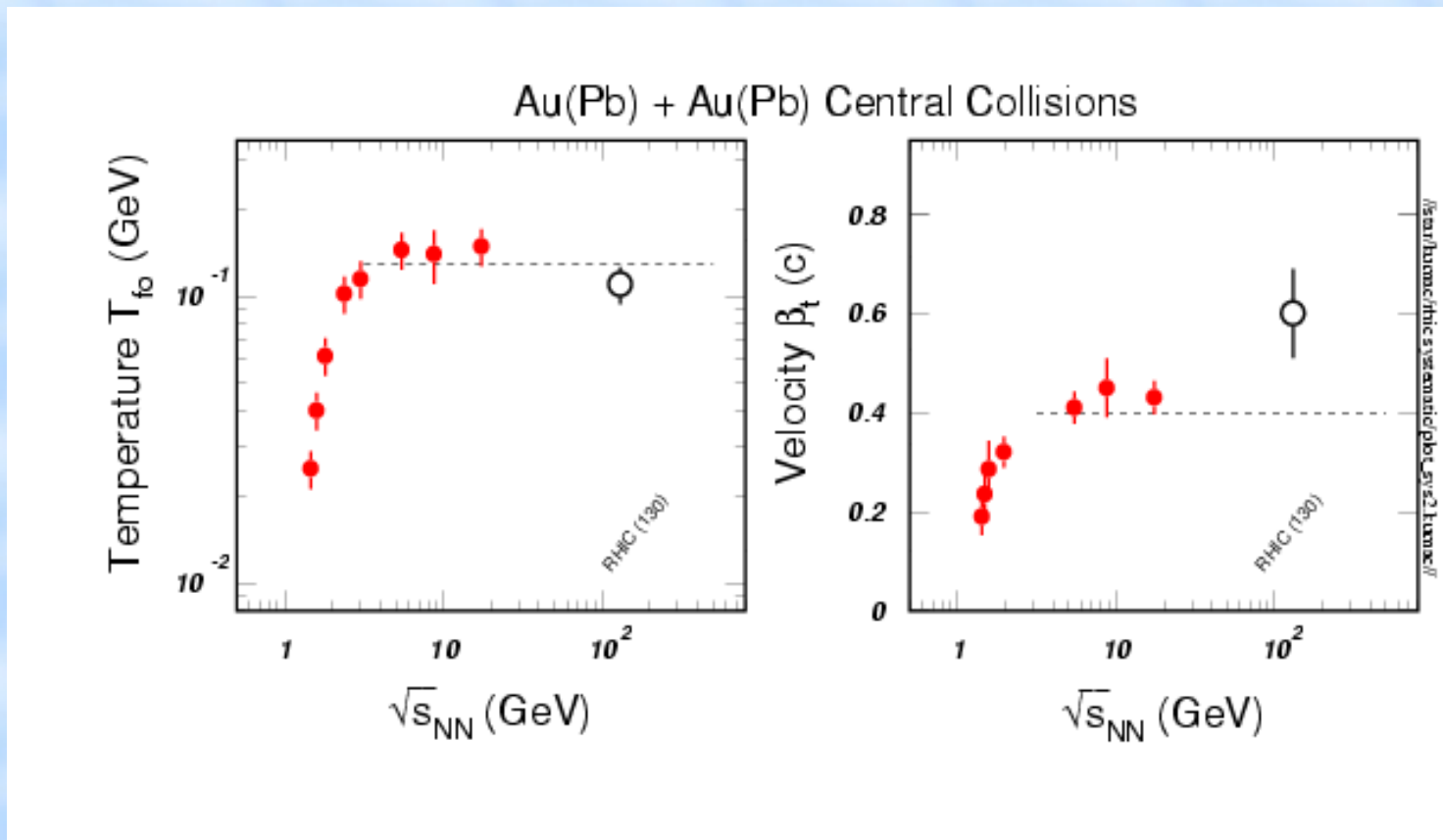


Complex Observables, but Closely Related to EOS

Can we see ANOMALIES in Collective Flows at RHIC ?

Answer = Yes !

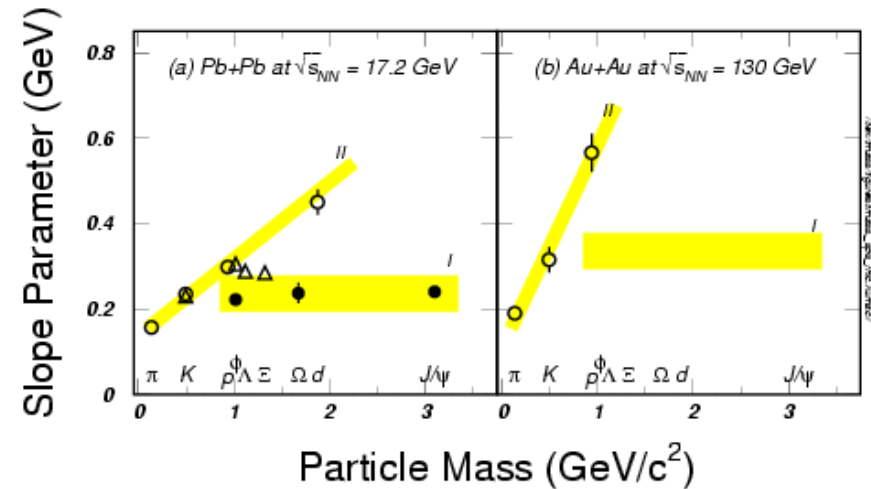
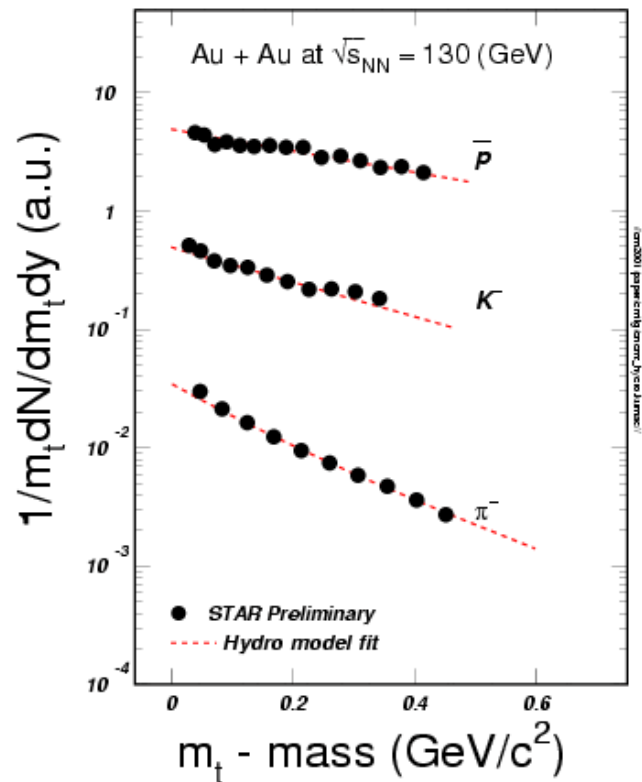
Radial Flow Nu Xu and M. Kaneta (STAR)



Formed matter seems to become STIFF Quickly at RHIC

How can we estimate Collective Radial Flow ?

.... Mt Spectra of Several Particles



$$\frac{d^2 N}{M_t dM_t dY d\phi} \propto \exp(-M_t/T'), \quad T'(M) = T + \frac{1}{2} M \beta^2$$

Simple Model of QCD Phase Transition

Massless Particles at Zero Chem. Pot.

$$\Omega = -\frac{\pi^2 V}{90} \left(\sum_B g_B + \frac{7}{8} \sum_F g_F \right) T^4,$$

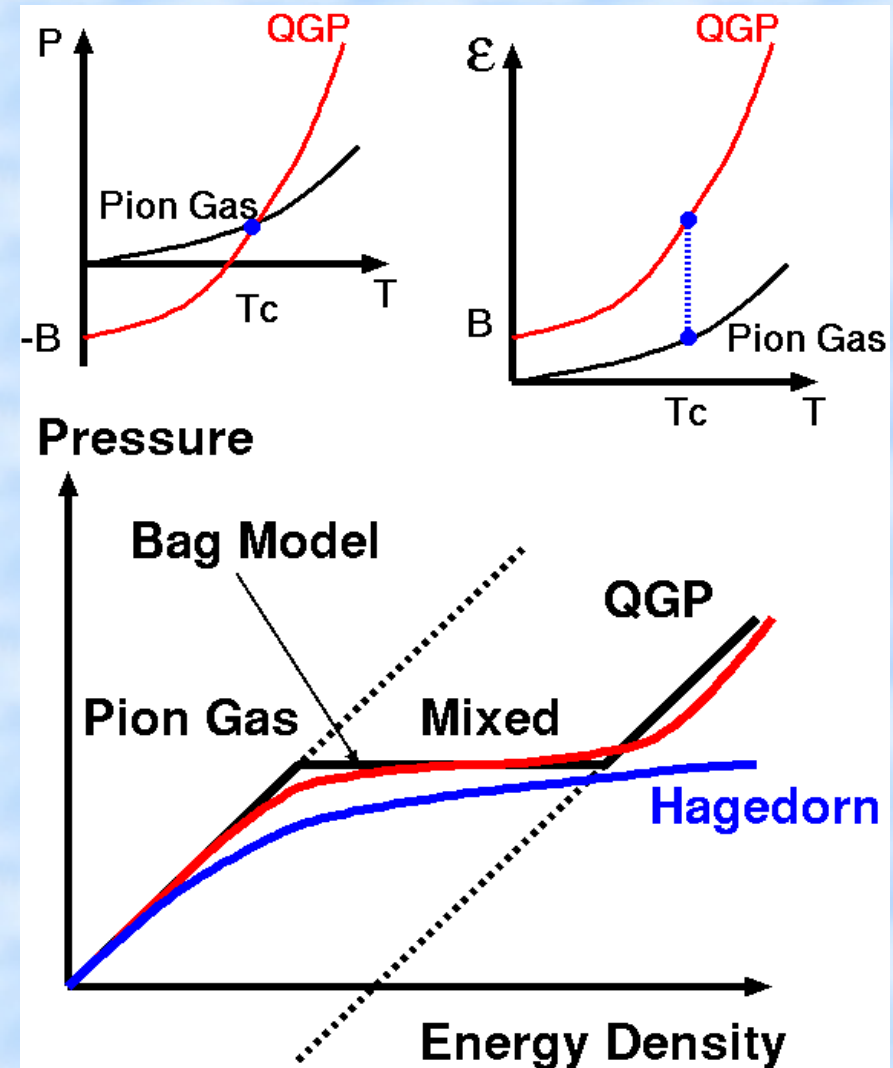
Massless (Free) Pion Gas

$$P_\pi = \frac{\pi^2}{30} T^4, \quad \epsilon_\pi = \frac{\pi^2}{10} T^4$$

QGP with Finite Bag Constant

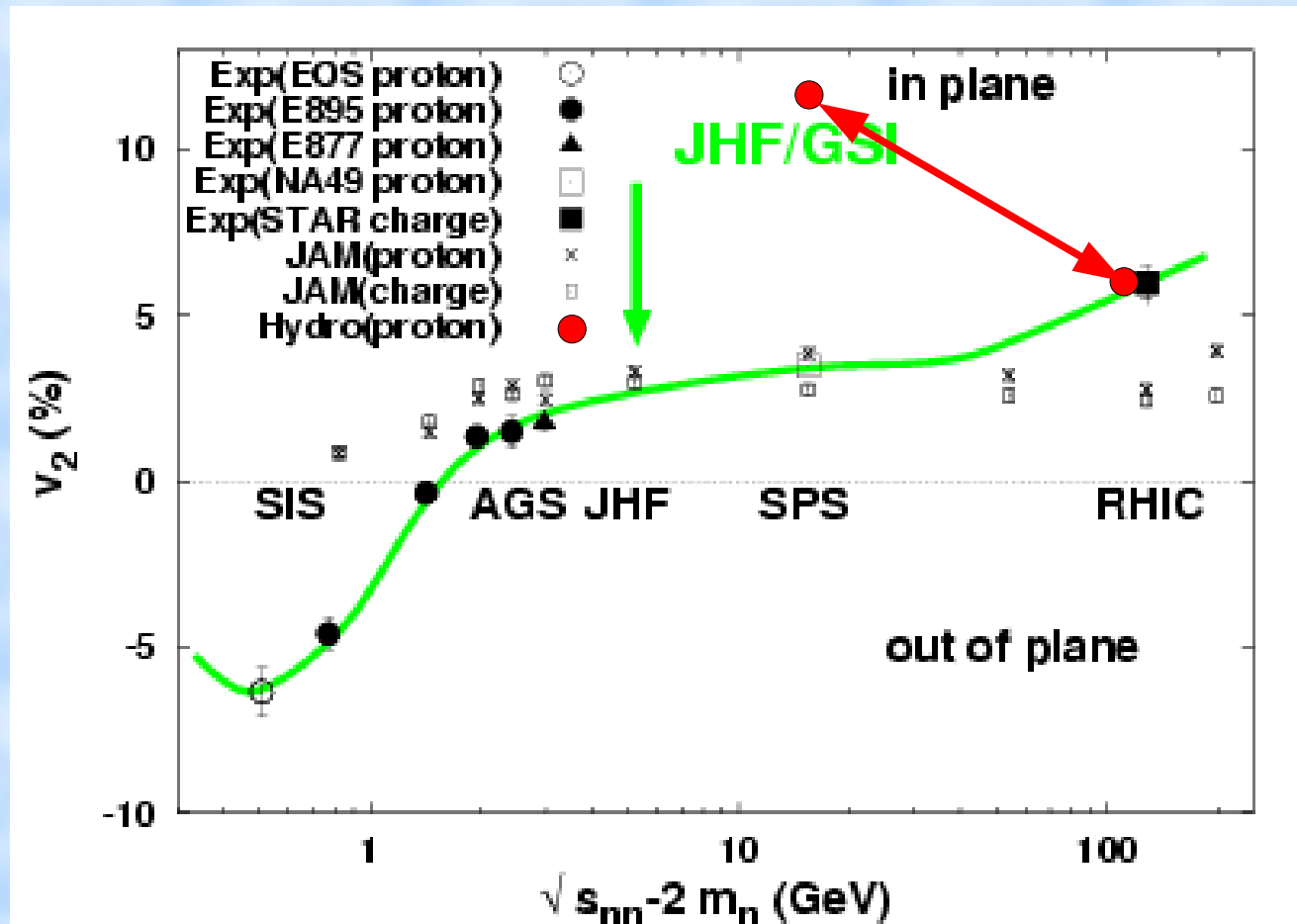
$$P_{QGP} = \frac{37 \pi^2}{90} T^4 - B$$

$$\epsilon_{QGP} = \frac{37 \pi^2}{30} T^4 + B$$



$$DOF = 2(\text{spin}) \times 2(q, \bar{q}) \times 3(\text{color}) \times 2(\text{flavor}) \times \frac{7}{8}(\text{Fermion}) + 2(\text{spin}) \times 8(\text{color})$$

Elliptic Flow

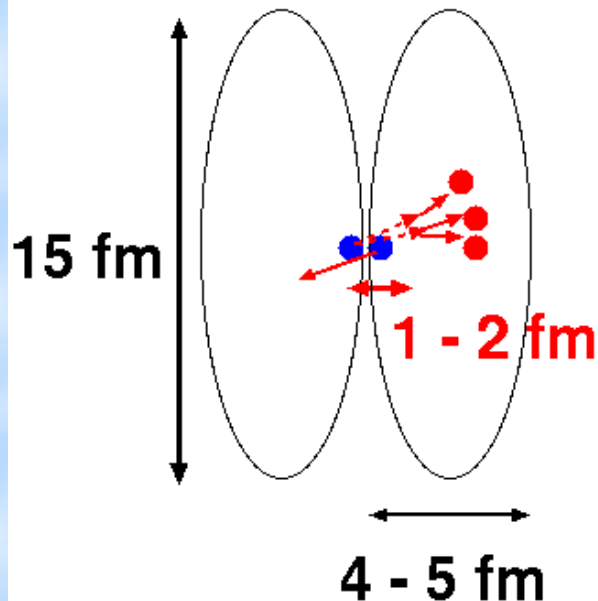


**Anisotropic Pressure is close to Hydrodynamical Values @ RHIC
→ Particles are interacting before Almond Shape is obscured.**

Hadron Formation Time

JHF Energies

$$\gamma_{\text{cm}} \simeq 3.5, \quad \tau \simeq 0.5 - 1 \text{ fm/c}$$

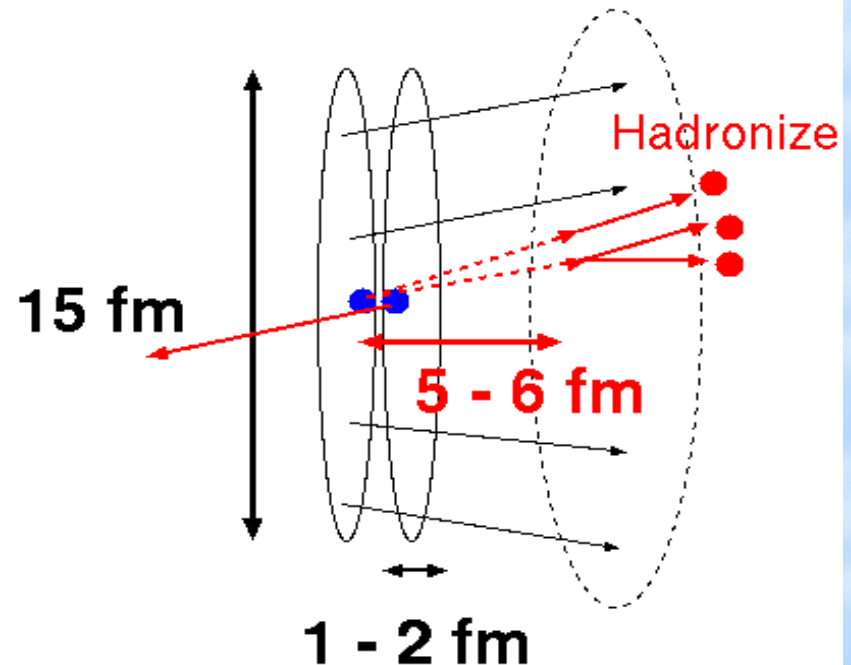


Multiple Hadron-Hadron Collisions

 (Approx.) Thermalized Hadron Gas

SPS Energies

$$\gamma_{\text{cm}} \simeq 10, \quad \tau \simeq 0.5 - 1 \text{ fm/c}$$



String-String, String-Hadron Int.
+ Int. within Co-Movers

**It takes ≈ 1 fm for hadrons to be formed (and thus to interact)
→ *Pre-Hadronic* Interactions are necessary at SPS & RHIC**

What is Suggested from Collective Flows

Radial Flow

→ **Re-Hardening Behavior**

Elliptic Flow

→ **Pre-Hadronic Interaction**

Jet Observation (\square Correlation, Energy Loss)

→ **Partons are Propagating**

Do these really require QGP formation ?

→ *Verification by Hadron-String Cascade Model is Necessary*

JAM (Jet AA Microscopic transport model)

Y. Nara et al., Phys. Rev. C61 (2000), 024901.

DOF

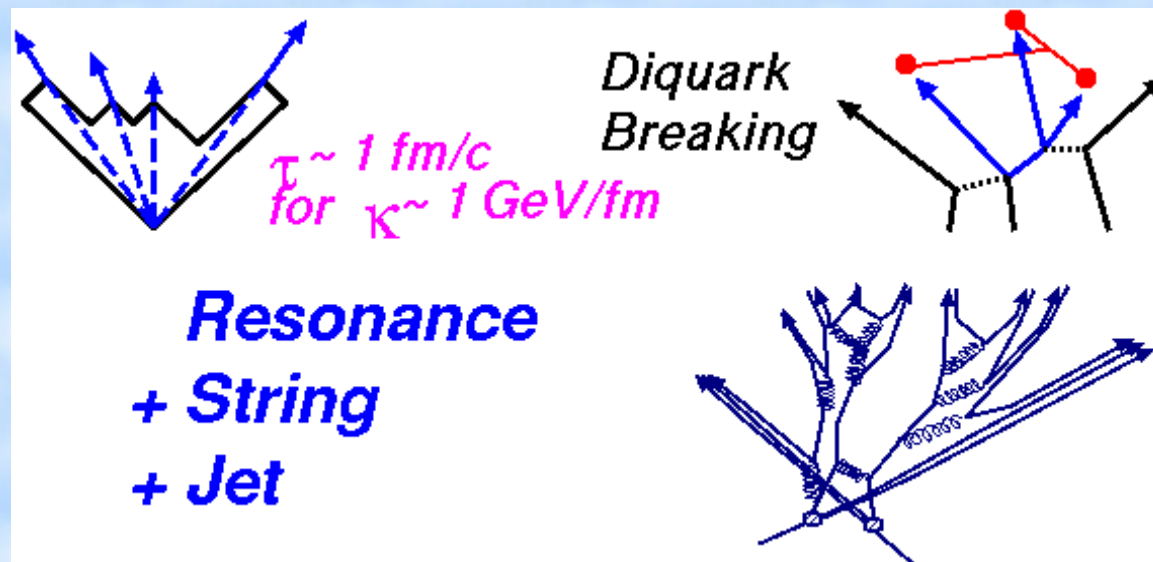
Hadrons (h , $m < 2 \text{ GeV}$) + Strings (s) + Partons (in Jet)

Cross Sections

Hadronic ($hh \rightarrow hh$, $hh \rightarrow h$, $h \rightarrow hh$)

+ Soft ($hh \rightarrow s$, $hh \rightarrow ss$, $s \rightarrow hh$, $hh \rightarrow hs$ [1],
 $sh \rightarrow s'h$,[2])

+ Hard (Jet Production)



[1] "DPM + Lund" (\rightarrow HIJING) + Phase Space

[2] Constituent Rescattering (\rightarrow RQMD)

Followings are NOT included in JAM

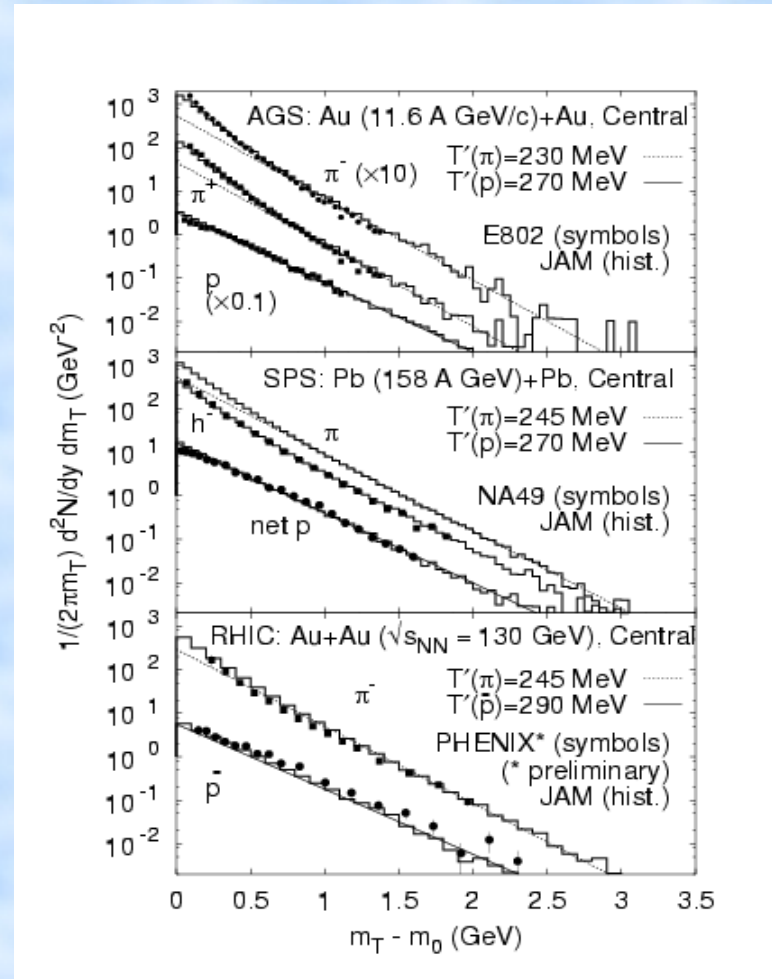
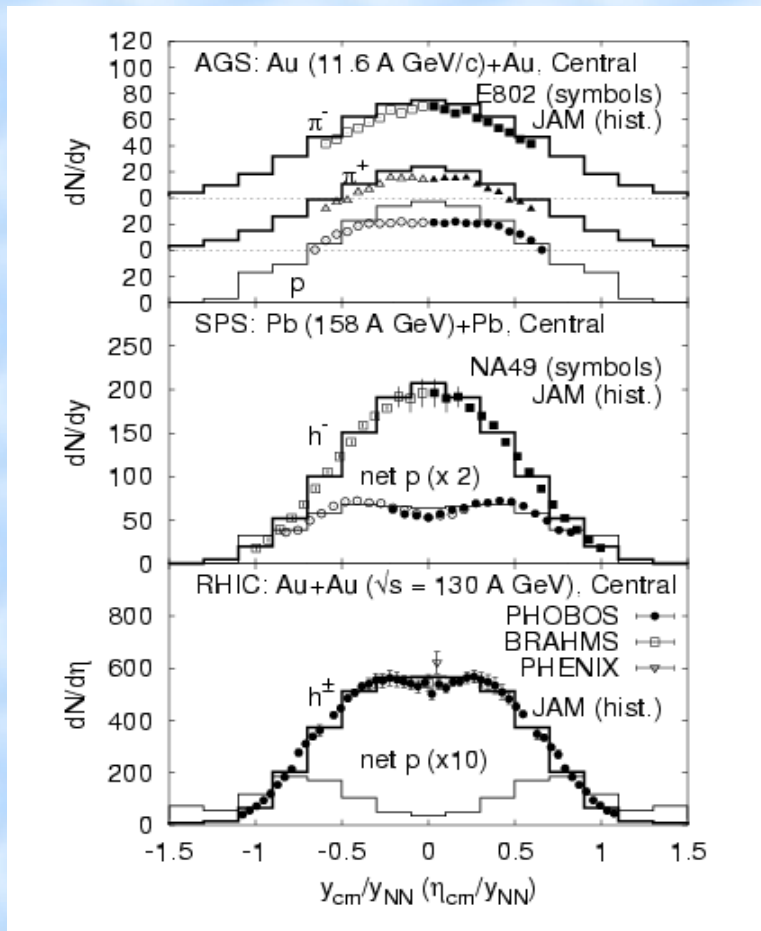
Mean Field (in progress)

Medium Modification

Secondary Interaction of Partons

with Other Hadrons, String and Partons from Other Jets

Hadron Spectra at AGS-SPS-RHIC

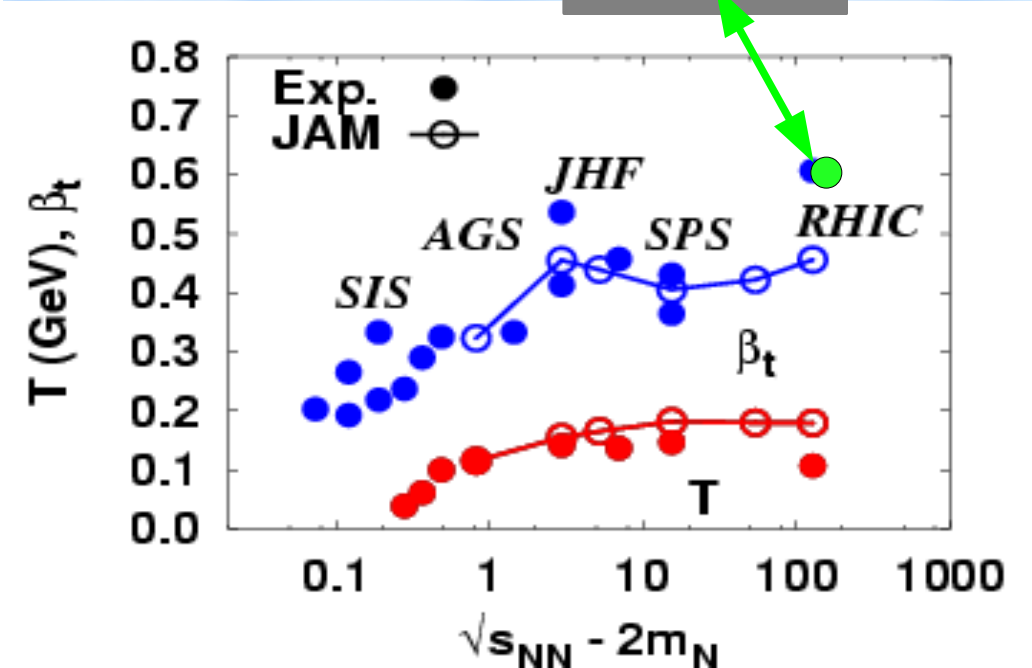
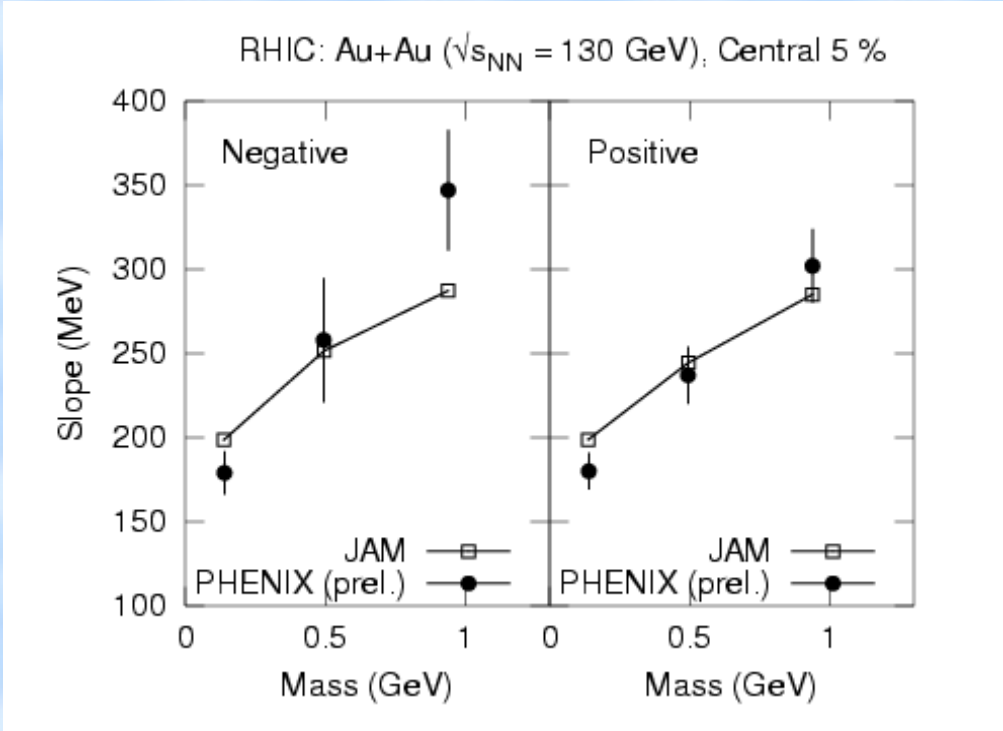


Hadron Spectra @ RHIC is too soft in JAM.

** Mean Field Effects are included for AGS and SPS energies*

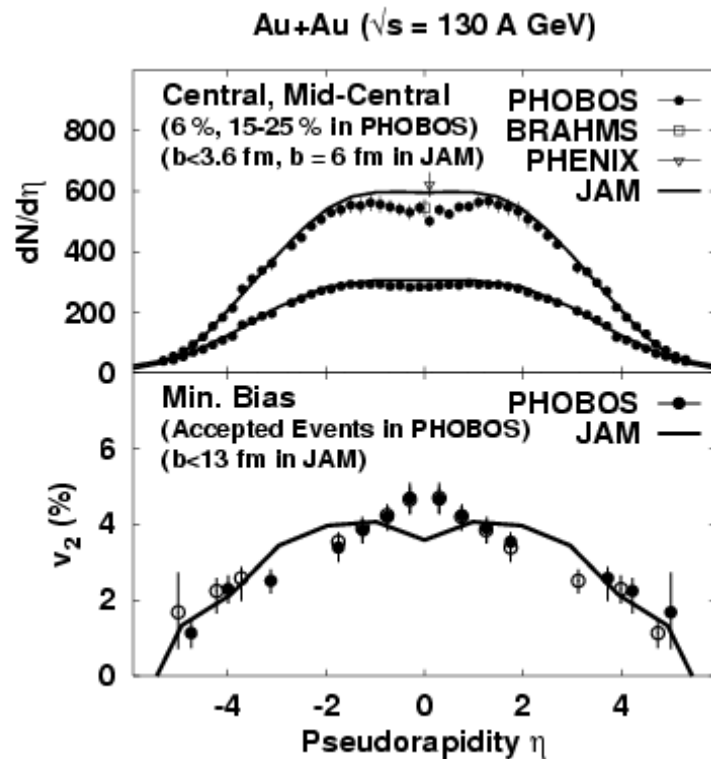
Radial Flow and Temperature in JAM

RHIC

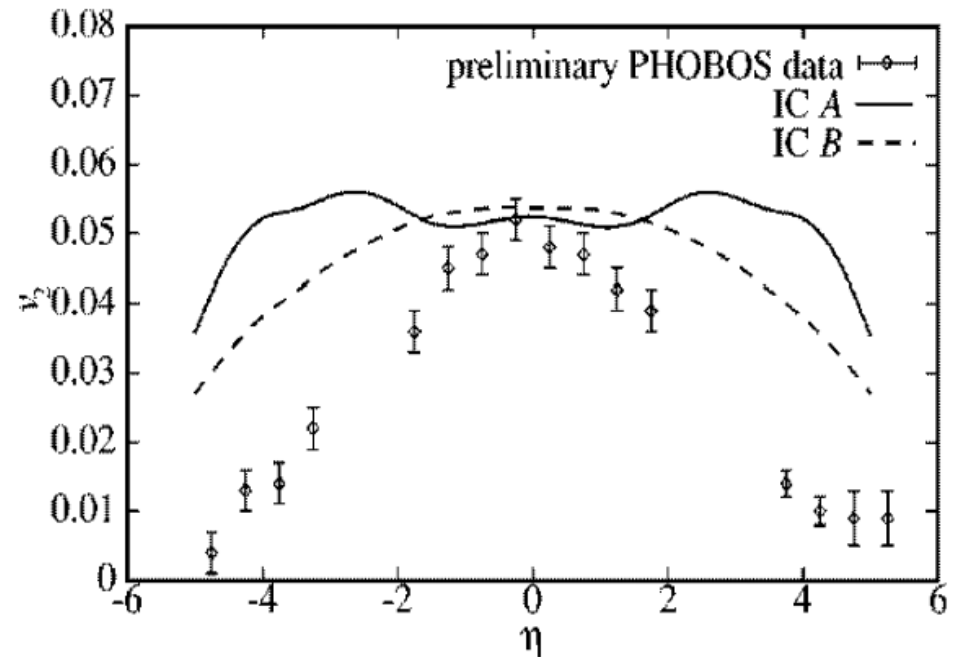


Re-Hardening Behavior Cannot Be Explained !

Pseudo Rapidity Dep. of Elliptic Flow



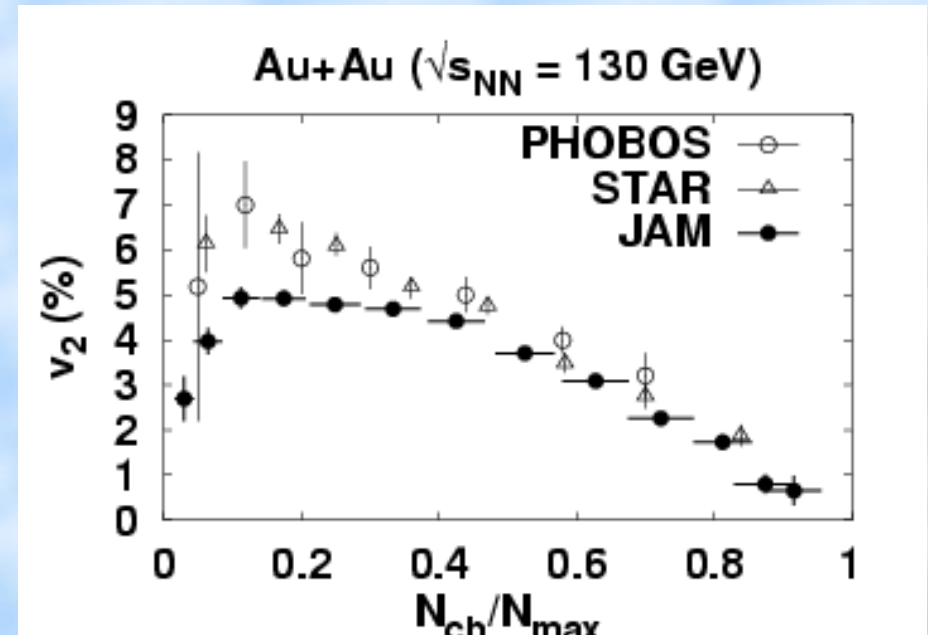
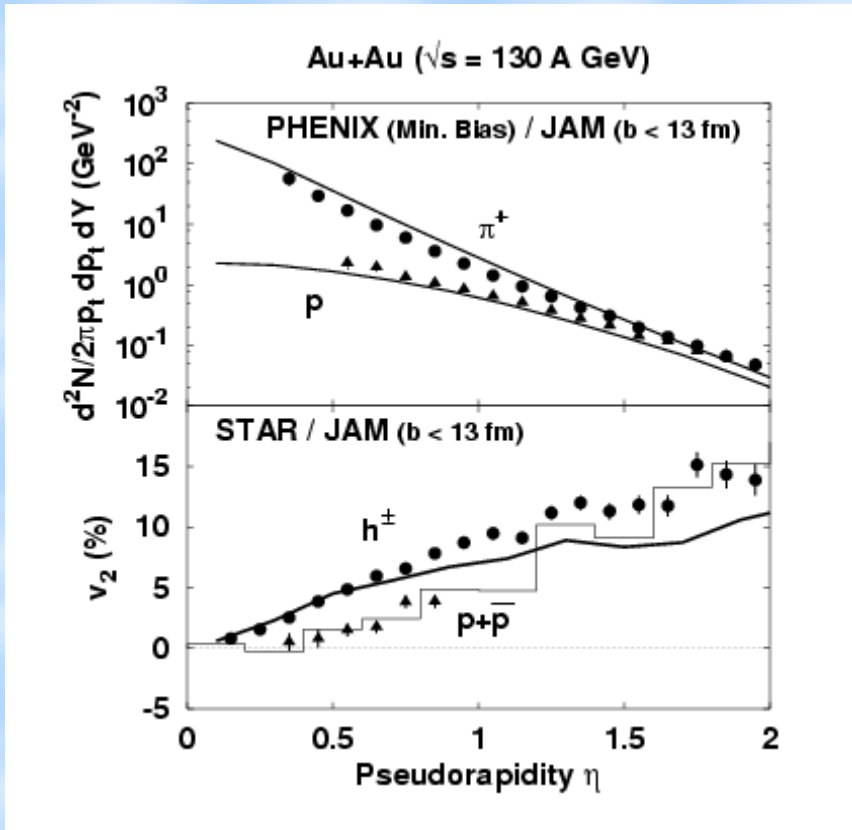
Hydro Results (Hirano, 2001)



Mid-Rapidity : Underestimate of v_2 (QGP ?)
Large Rapidities : JAM Explains the Data (Hadronic ?)

Pt and Impact Par. Dep. of Elliptic Flow

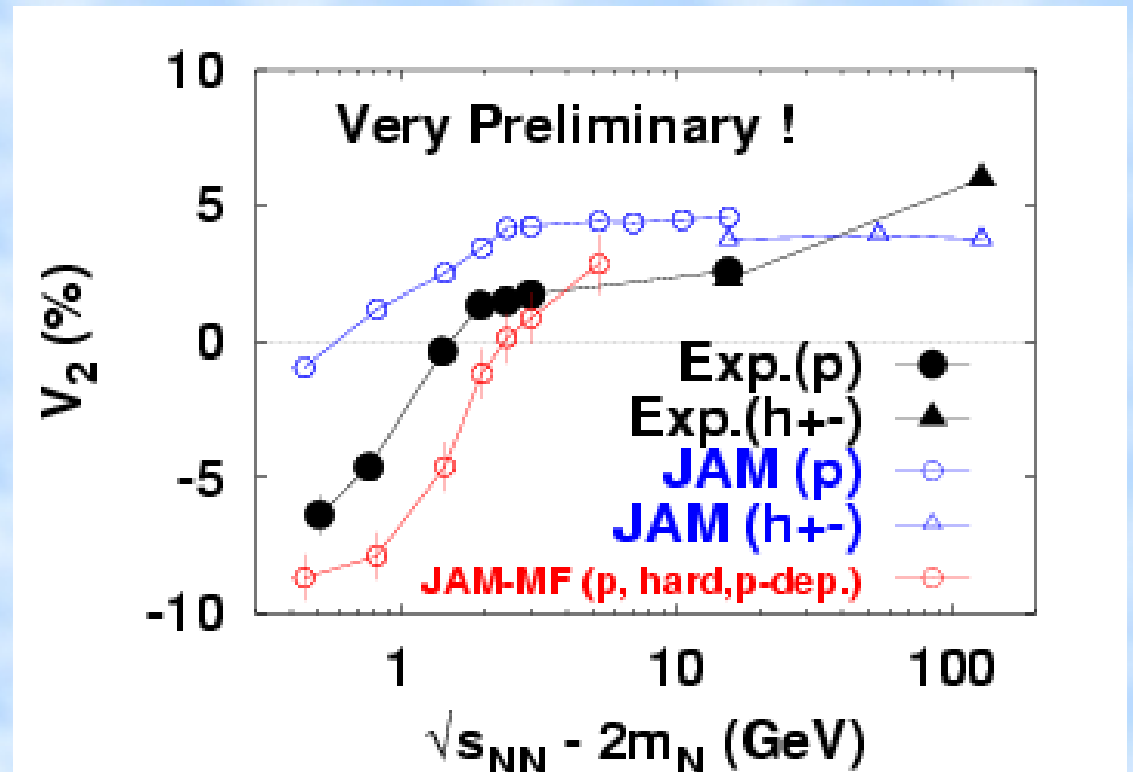
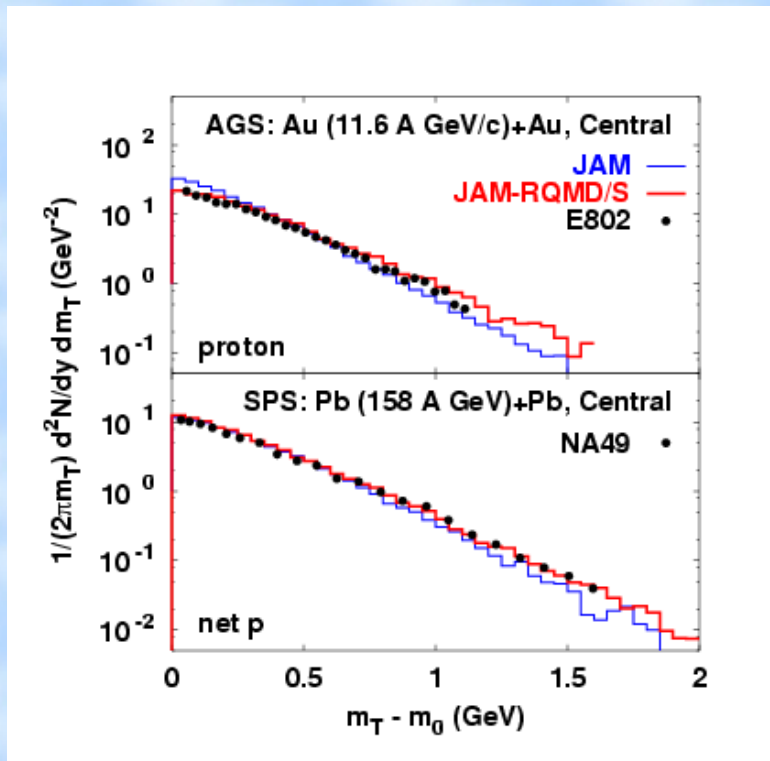
Where Do We Underestimate ?



Answer = High Pt Regions !

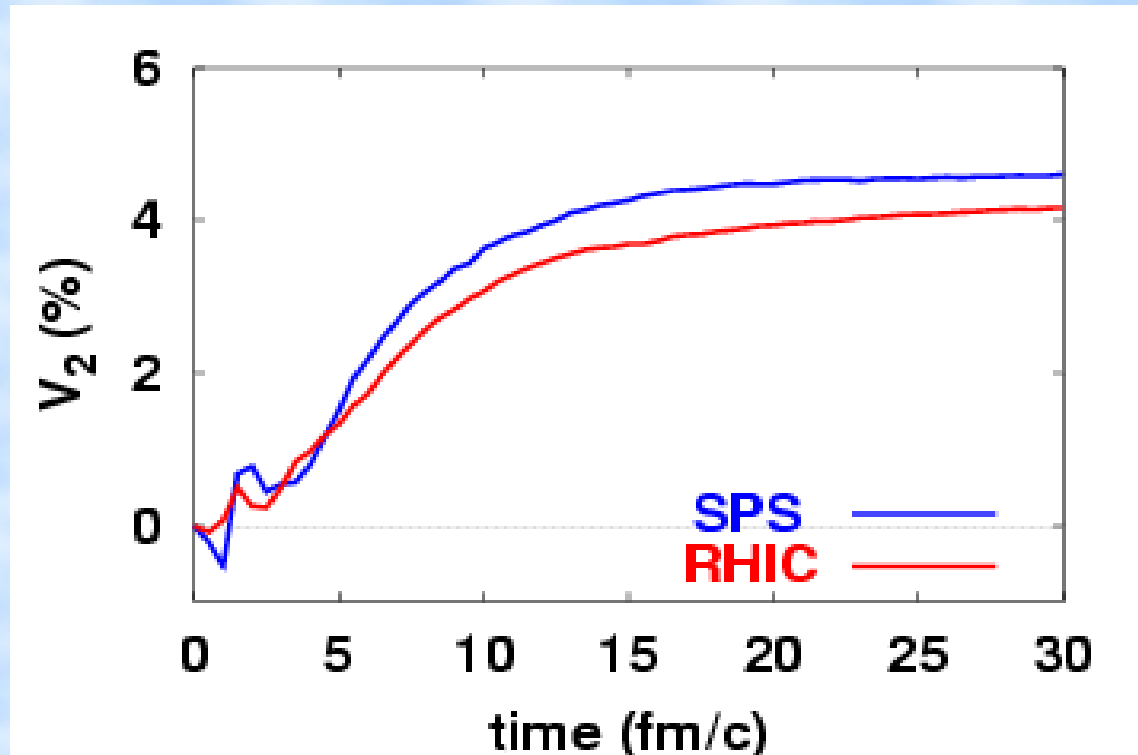
Incident Energy Deps. of V_2

Mean Fields Effects (JAM-RQMD/S)



Mean Field Effects → Downward Shift in V_2

***When are Collective Flows Generated ?
Why Do We Underestimate ?***



(b = 6 fm)

**V_2 is Generated at a long time scale in Hadron-String Cascade.
→ Almond Shape is already obscured.**

Summary

Collective Flow Data at RHIC seems to suggest QGP formation.

Large V_2 at Mid-Rapidity : Early Thermalization

Strong Radial Flow : Re-Hardening

Jet Quenching : Partonic Interaction

JAM (Hadron-String Cascade with Jet Prod.) cannot explain RHIC Data, Especially in High Pt Region at Mid Rapidity.

Slow Growth of Elliptic Flow

No Secondary Partonic Interaction

Mean Field Does NOT Help much at RHIC

There are many things to do, especially at RHIC and LHC.

Elementary Cross Sections

Coherence & Incoherence in Collision

Side and Out Radii Puzzle

Parton-Hadron, Parton-String, Parton-Parton Interactions

Modification of the “Vacuum”