

高密度物質と中性子星の物理：講義の内容

1. 中性子星の基本的性質

- 質量・半径測定の概論など

2. 状態方程式を記述する理論模型

- 平均場理論、第一原理計算手法、場の理論によるアプローチ

3. 対称エネルギーと非対称核物質の状態方程式

- 対称エネルギーを決める実験手法、現在の制限

4. QCD 有効模型と高密度核物質の性質

- 有限温度・密度の場の理論入門（松原和・摂動論など）
- NJL 模型による相転移と状態方程式の記述

5. ハイパー核物理と中性子星でのハイペロンパズル

- ハイパー核実験の現状、ハイペロンパズルの解決に向けて

■ 談話会

Symmetry Parameter Constraints

from a Lower Bound on the Neutron-Matter Energy

Nuclear Symmetry Energy *– Overview –*

M-R Relation and EOS

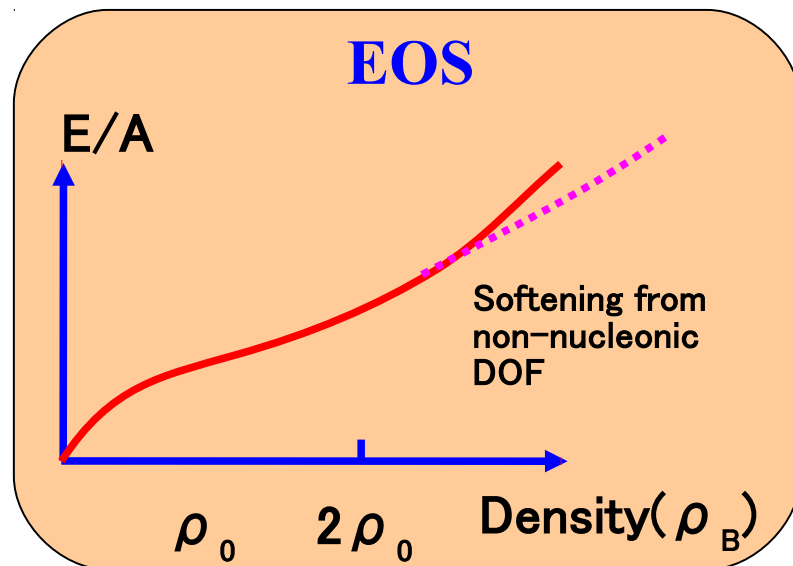
■ Solving TOV eq.

starting from the “initial” condition, $\varepsilon(r=0) = \varepsilon_c = \text{given}$
until the “boundary” condition $P(r)=0$ is satisfied.

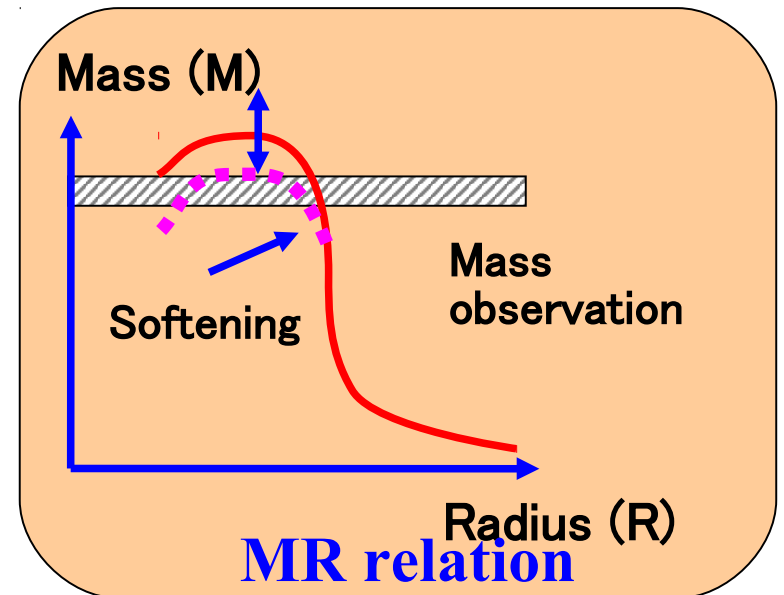
→ M and R are the functions of $\varepsilon(r=0)$
and functionals of EOS, $P=P(\varepsilon)$.

$$M = M(\varepsilon_c)[P(\varepsilon)], \quad R = R(\varepsilon_c)[P(\varepsilon)]$$

→ **M-R curve and NS matter EOS : 1 to 1 correspondence**



TOV Eq.



Nuclear Mass

■ Bethe-Weizsacker mass formula

Nuclear binding energy is roughly given by Liquid drop.

Nuclear size measurement $\rightarrow R = r_0 A^{1/3}$

$$B(A, Z) = \underbrace{a_v A}_{\text{Volume}} - \underbrace{a_s A^{2/3}}_{\text{Surface}} - \underbrace{a_c \frac{Z^2}{A^{1/3}}}_{\text{Coulomb}} - \underbrace{a_a \frac{(N - Z)^2}{A}}_{\text{Symmetry}} + \underbrace{a_p \frac{\delta_p}{A^{1/2}}}_{\text{Paring}}$$

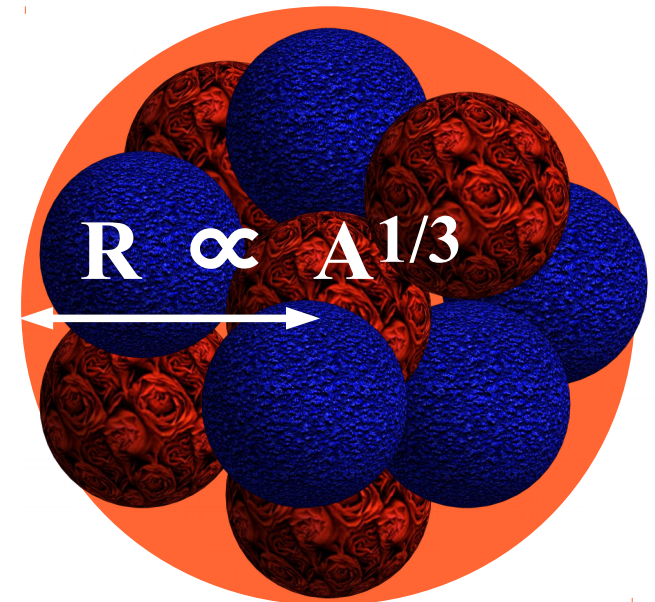
Volume	Surface	Coulomb	Symmetry	Paring
$A \propto \frac{4\pi}{3} R^3$	$A^{2/3} \propto 4\pi R^2$	$\propto \frac{Q^2}{R}$		

■ Ignore Coulomb, consider $A \rightarrow \infty$,

$$B/A = a_v(\rho) - a_a(\rho)\alpha^2 \quad (\alpha = (N - Z)/A)$$

$$a_v \simeq 16 \text{ MeV}, a_a \simeq 30 \text{ MeV}$$

Coef. may depend on the number density ρ
 $\rightarrow$ Nuclear Matter EOS



Nuclear Matter EOS

■ Energy per nucleon in nuclear matter

$$E_{\text{NM}}(\rho, \alpha) = E_{\text{SNM}}(\rho) + S(\rho)\alpha^2 \quad (\alpha = (N - Z)/A)$$

■ Saturation point (ρ_0, E_0)

$$\rho_0 \simeq 0.16 \text{ fm}^{-3}$$

$$E_0 = -a_v \simeq -16 \text{ MeV}$$

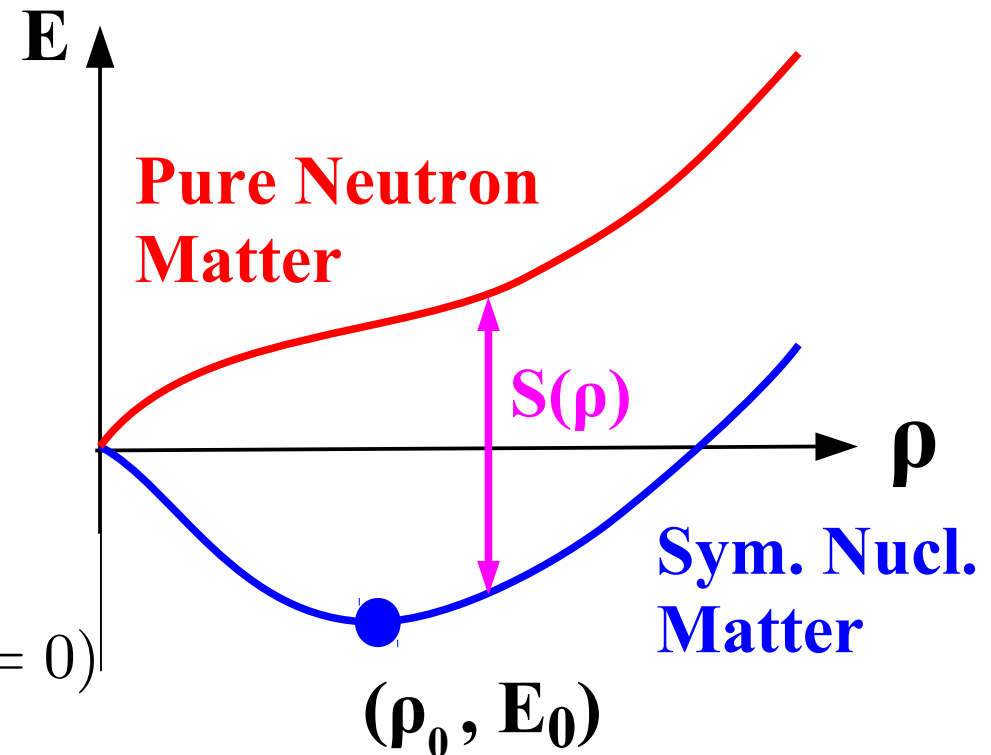
(nuclear radius and mass)

■ Symmetry energy

$$\begin{aligned} S(\rho) &= E_{\text{PNM}}(\rho) - E_{\text{SNM}}(\rho) \\ &= E_{\text{NM}}(\rho, \alpha = 1) - E_{\text{NM}}(\rho, \alpha = 0) \end{aligned}$$

$$S_0 = S(\rho_0) \simeq 30 \text{ MeV}$$

(Mass formula, volume term)



Nuclear Matter EOS can be, in principle, determined by terrestrial (laboratory) nuclear physics experiments !

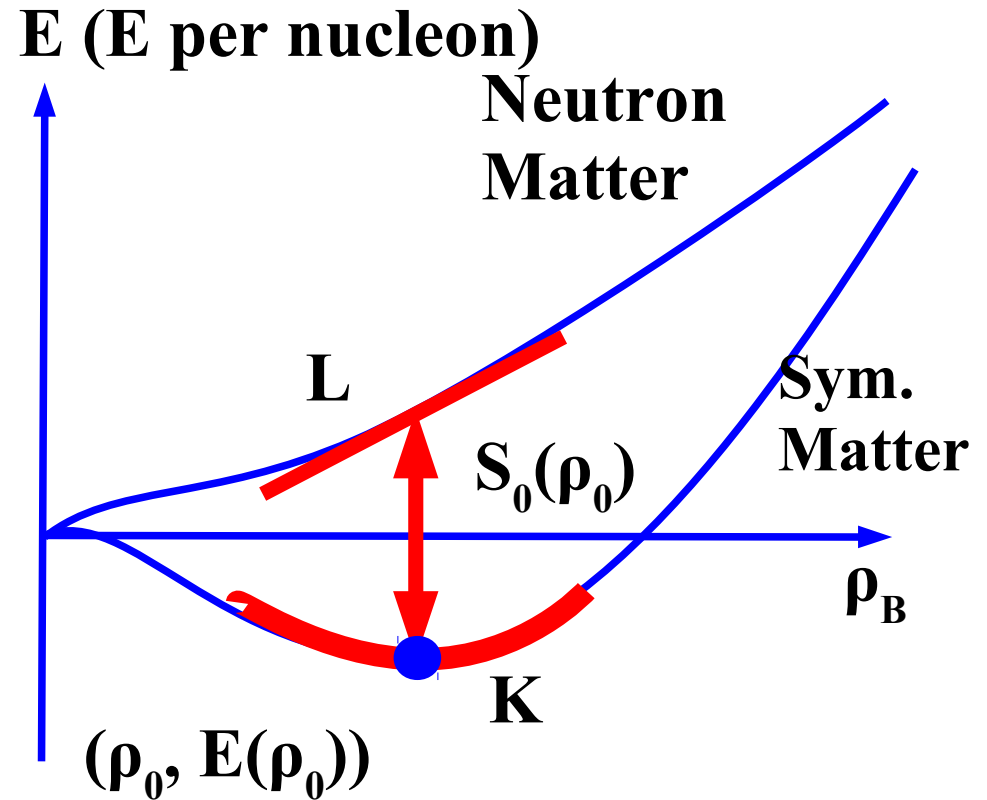
Nuclear Matter EOS

- Additional two important parameters: **K** and **L**
- Pressure is given by the derivative of **E** via ρ

$$P = \rho^2 (\partial E / \partial \rho)$$

At ρ_0 , **L** determines **P**

$$P = \rho_0 L / 3 \text{ (at } \rho = \rho_0 \text{)}$$



Neutron Star Matter EOS

■ What happens in low-density uniform neutron star matter ?

- Constituents = proton, neutron and electron
- Charge neutrality $\rightarrow$ # of electrons = # of protons ($\rho_e = \rho_p = \rho(1 - \delta)/2$)

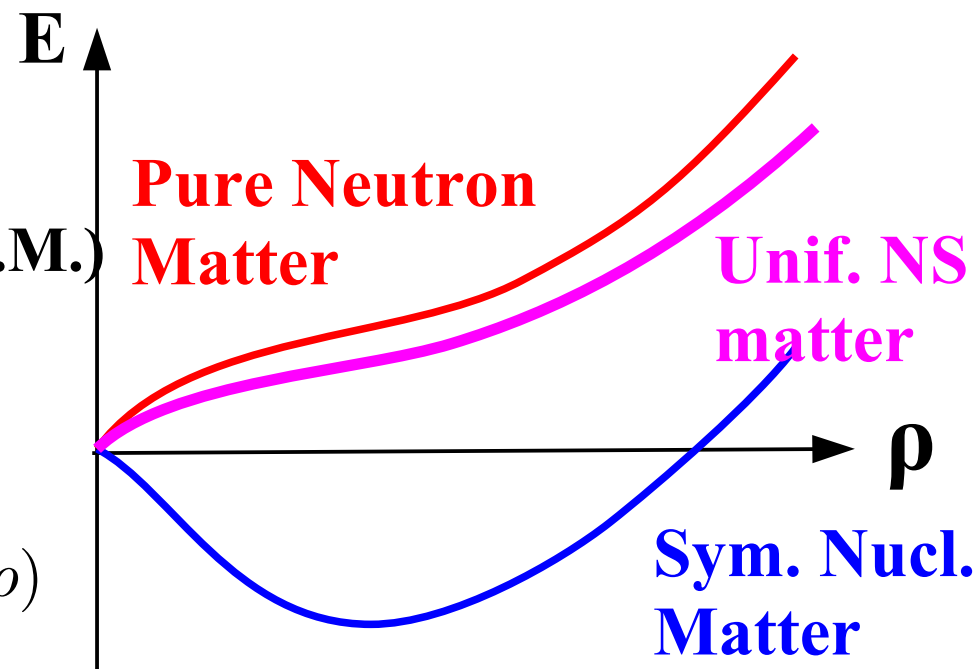
$$E_{\text{NSM}}(\rho) = E_{\text{NM}}(\rho, \alpha) + E_e(\rho_e = \rho_p)$$

$$= E_{\text{SNM}}(\rho) + \alpha^2 S(\rho) + \frac{\Delta M}{2} \alpha + \frac{3}{8} \hbar k_F (1 - \alpha)^{4/3}$$

(electron mass neglected,
neutron-proton mass diff. incl.
 k_F = Fermi wave num. in Sym. N.M.)

- δ is optimized to minimize energy per nucleon

$$E_{\text{NSM}}(\rho) \leq E_{\text{NM}}(\rho, \alpha = 1) = E_{\text{PNM}}(\rho)$$



対称エネルギーの起源

■ Fermi Gas model での核子あたりの運動エネルギー

$$E_{\text{sym}, K} = \frac{Z}{A} \frac{3}{5} \frac{\hbar^2 k_{\text{F}p}^2}{2m} + \frac{N}{A} \frac{3}{5} \frac{\hbar^2 k_{\text{F}n}^2}{2m} = \frac{3}{5} E_F \frac{1}{2} \left[(1 - \delta)^{5/3} + (1 + \delta)^{5/3} \right]$$

$$\simeq \frac{3}{5} E_F + \frac{1}{3} E_F \delta^2 + \mathcal{O}(\delta^4)$$

$a_{\text{sym}}(\text{FG}) = E_F/3 \sim 11 \text{ MeV}$ となり、質量公式の $a_{\text{sym}} \sim 23 \text{ MeV}$ (surface を考えると $a_{\text{sym}}(\text{vol}) \sim 30 \text{ MeV}$) と比べて 1/3-1/2 程度。残りは相互作用。

■ 残りの半分の対称エネルギーを RMF で評価してみましょう。

$$\Delta E_{\text{sym}, \rho} = \frac{1}{4} \frac{m_\rho^2 R^2}{\rho_B} = \frac{1}{4} \frac{g_\rho^2}{m_\rho^2} \rho_B \delta^2 = \Delta a_{\text{sym}} \delta^2 \quad \left(R = \frac{g_\rho(\rho_n - \rho_p)}{m_\rho^2} = \frac{g_\rho \rho_B \delta}{m_\rho^2} \right)$$

$$g_\rho^2 = \frac{2 m_\rho^2 \Delta a_{\text{sym}}}{\rho_B} \simeq (4.3)^2 \quad (a_{\text{sym}} = 30 \text{ MeV}) \quad \leftarrow \text{RMF par. より少し小さめ}$$

$$L \simeq E_F + 3 \Delta a_{\text{sym}} \simeq 90 \text{ MeV}$$

← Optimal value より少し大きい

Simple parametrized EOS

■ Skyrme int. motivated parameterization

$$E_{\text{SNM}} = \frac{3}{5} E_F(\rho) + \frac{\alpha}{2} \left(\frac{\rho}{\rho_0} \right) + \frac{\beta}{2 + \gamma} \left(\frac{\rho}{\rho_0} \right)^{1+\gamma}$$

$$\alpha = \frac{2}{\gamma} \left(E_0(1 + \gamma) - \frac{E_F(\rho_0)(1 + 3\gamma)}{5} \right), \quad \beta = \frac{2 + \gamma}{\gamma} \left[-E_0 + \frac{1}{5} E_F(\rho_0) \right].$$

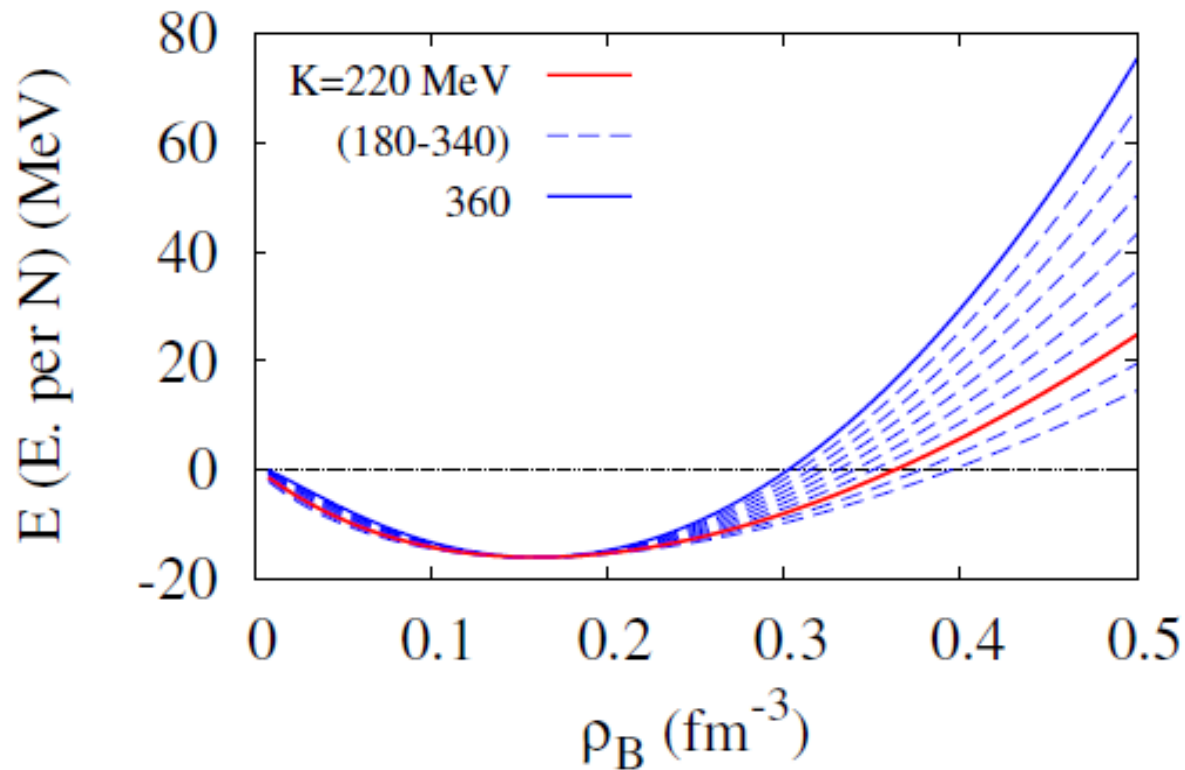
$$K = \frac{1 + 3\gamma}{5} E_F(\rho_0) - 3E_0(1 + \gamma).$$

■ Symmetry energy parameterization

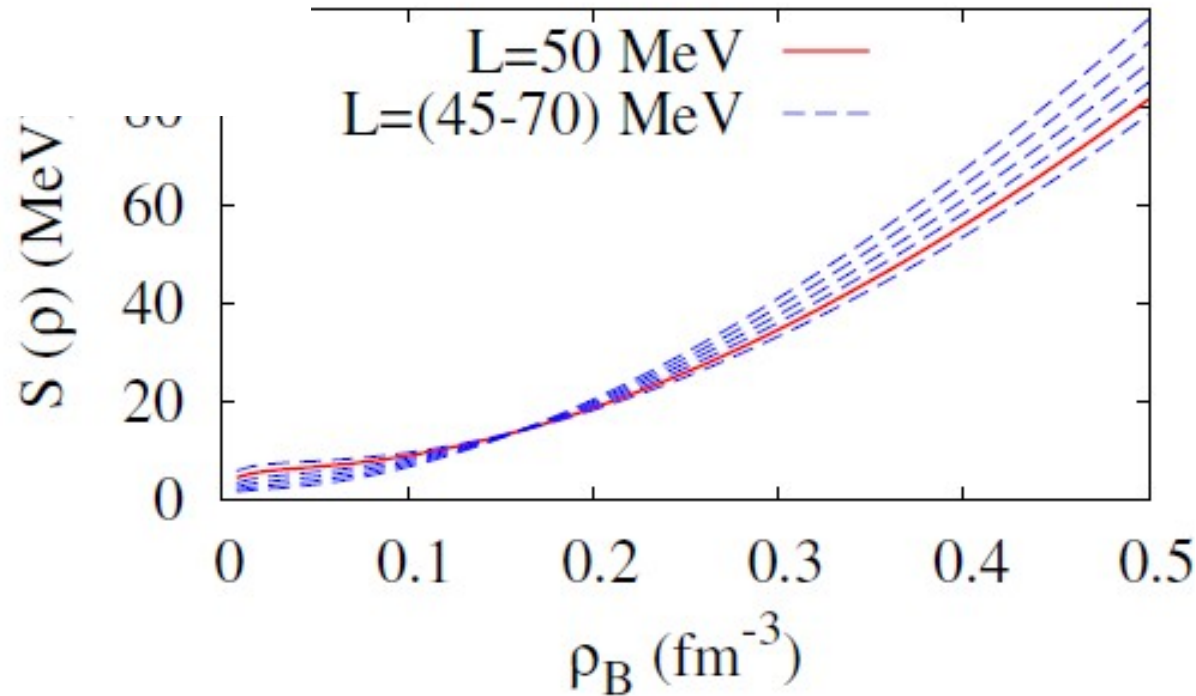
$$S(\rho) = \frac{1}{3} E_F(\rho) + \left[S_0 - \frac{1}{3} E_F(\rho_0) \right] \left(\frac{\rho}{\rho_0} \right)^{\gamma_{\text{sym}}}$$

$$\gamma_{\text{sym}} = \frac{L - \frac{2}{3} E_F(\rho_0)}{3S_0 - E_F(\rho_0)}$$

Simple parametrized EOS

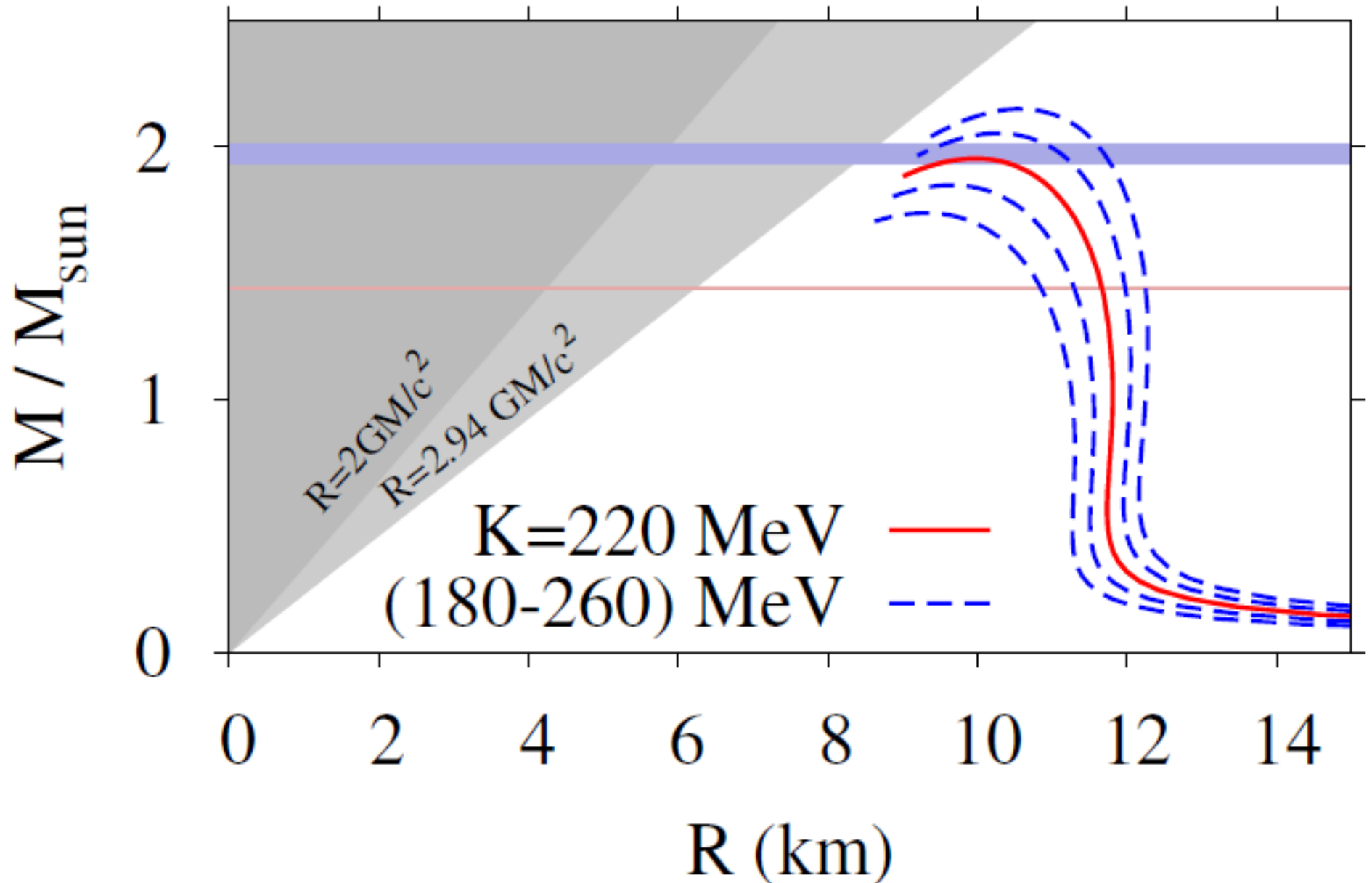


$K=220$ MeV, $S_0=30$ MeV



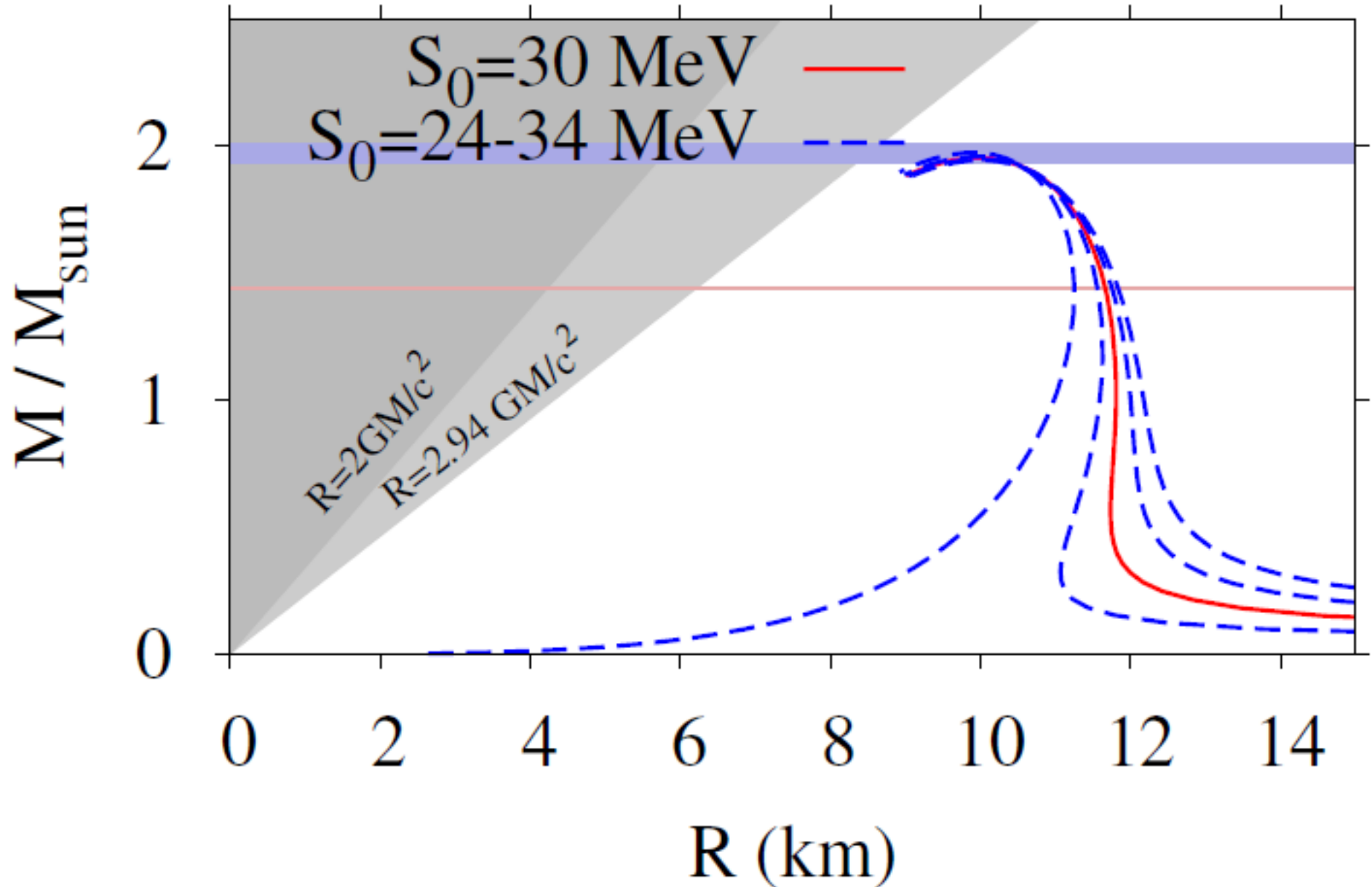
Simple parametrized EOS

$$(S_0, L) = (30 \text{ MeV}, 50 \text{ MeV})$$



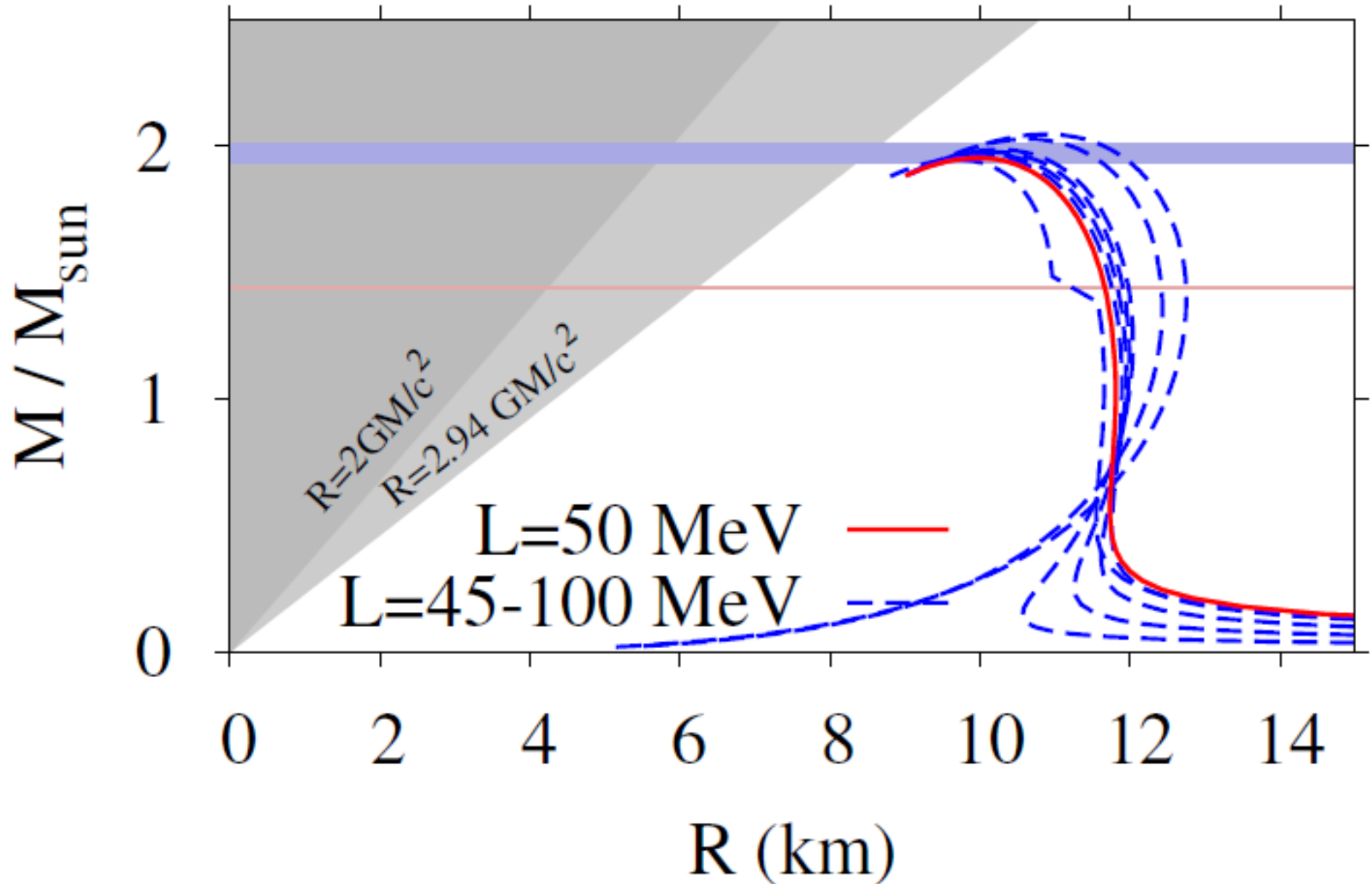
Simple parametrized EOS

$(K, L)=(220 \text{ MeV}, 50 \text{ MeV})$



Simple parametrized EOS

$(S_0, K)=(30 \text{ MeV}, 220 \text{ MeV})$



Symmetry Energy

■ Summary of Nuclear Symmetry Energy workshop

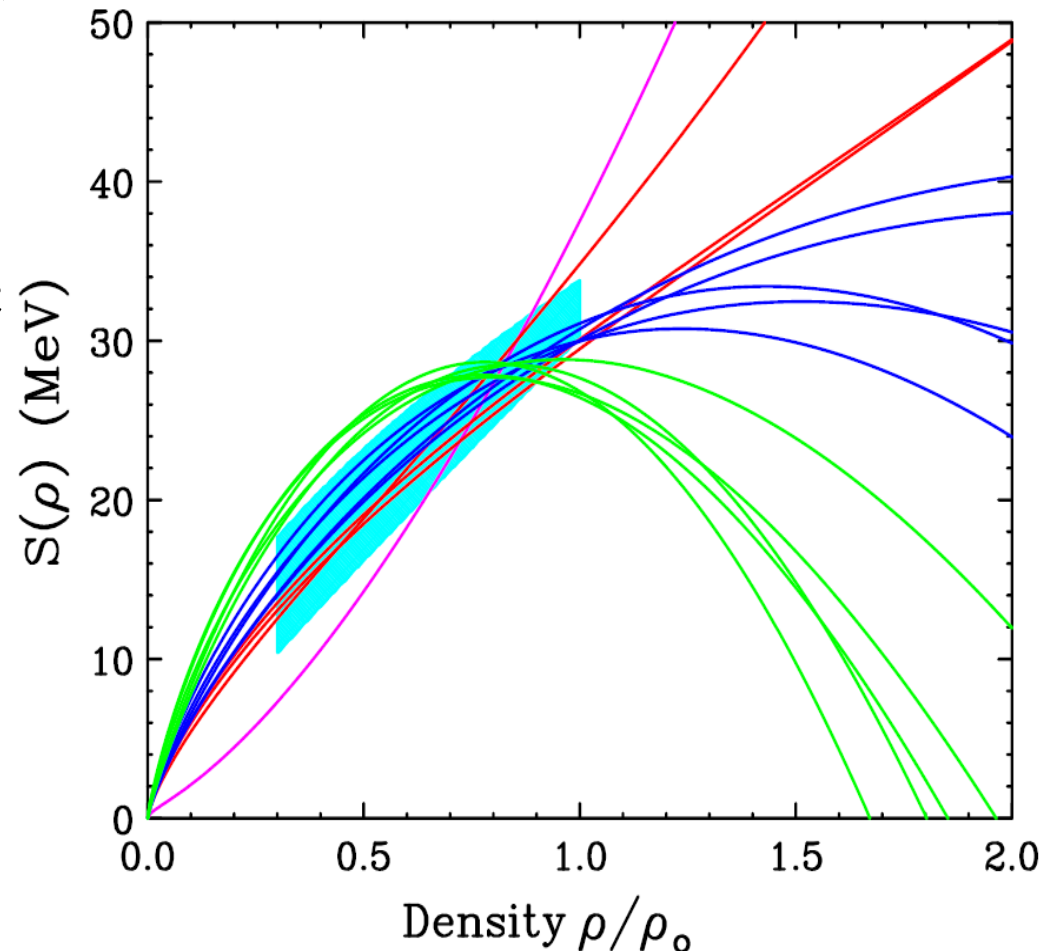
NuSym11 <http://www.smith.edu/nusym11>

$$E_{\text{sym}}(\rho_0) = 31\text{-}34 \text{ MeV}, L = 50\text{-}110 \text{ MeV}$$

extracted from various observations.

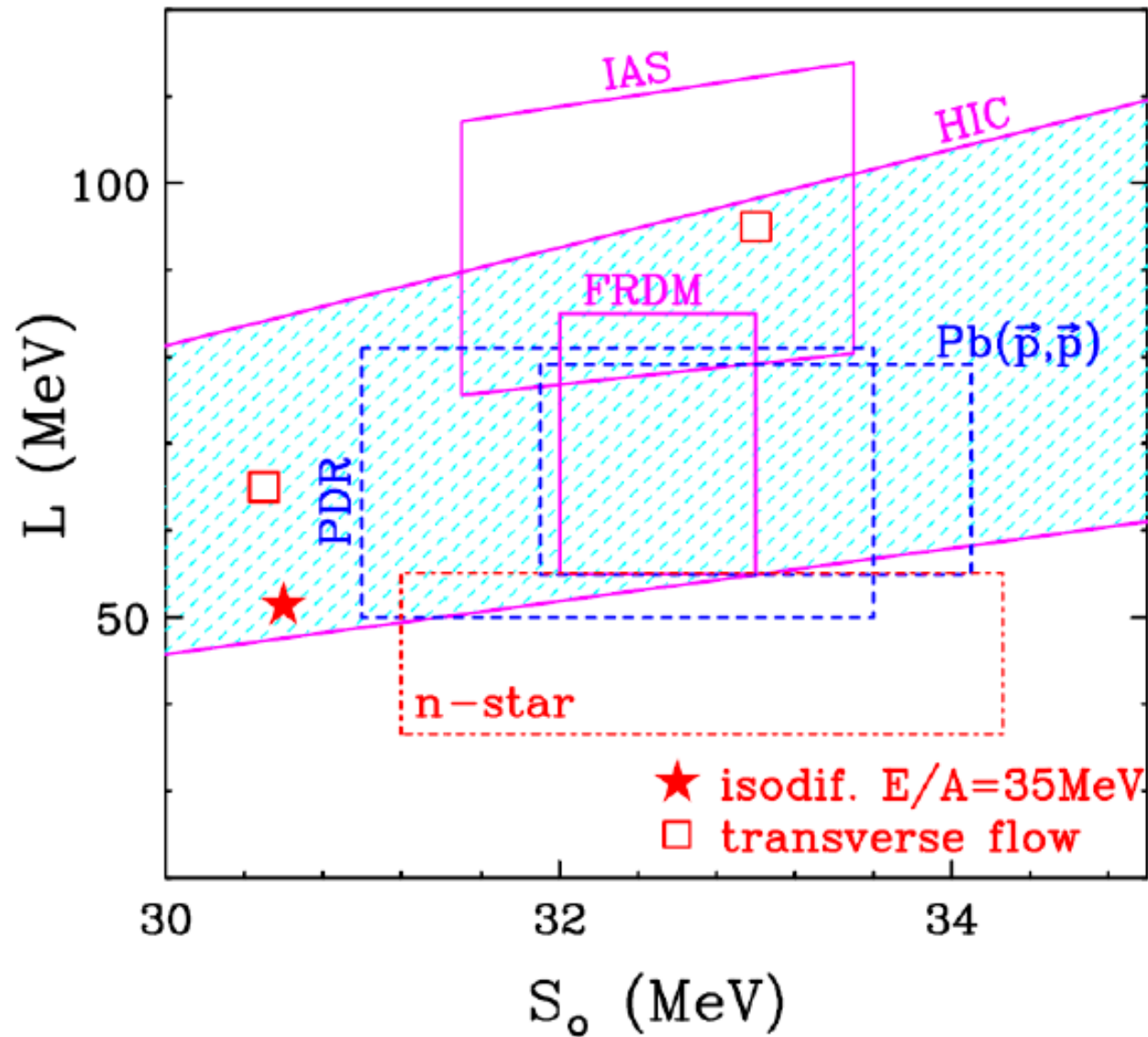
- Mass formula *Moller ('10)*
- Isobaric Analog State
Danielewicz, Lee ('11)
- Pygmy Dipole Resonance
Carbone+ ('10)
- Isospin Diffusion
Tsang et al. ('04)
- Neutron Skin thickness
J.Zenihiro+ ('10)

- これらの多くは ρ_0 以下の密度での E_{sym} に敏感。



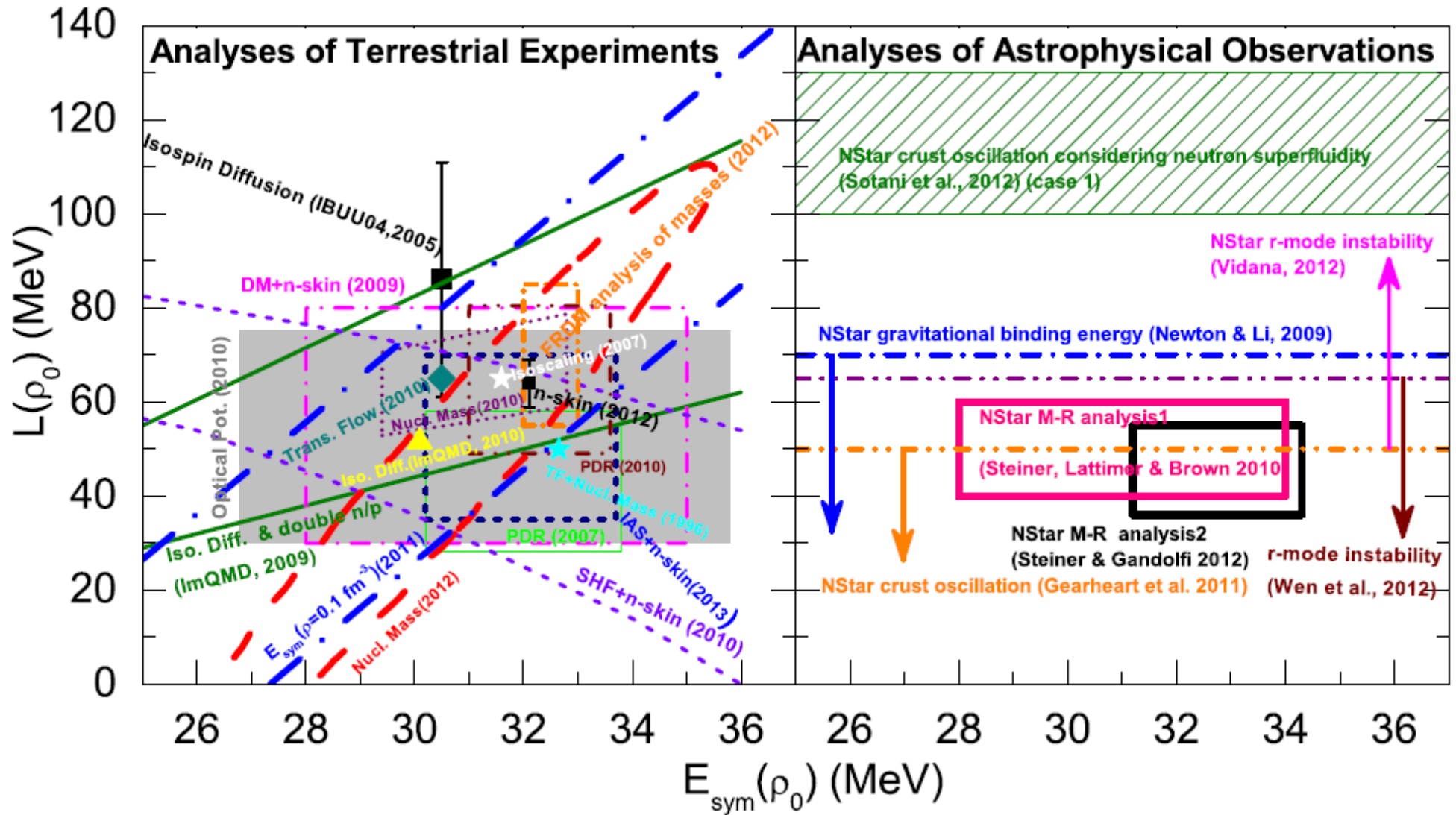
M. B. Tsang et al., Phys. Rev. C 86 (2012) 015803.

Nuclear Symmetry Energy (NuSYM 2011)



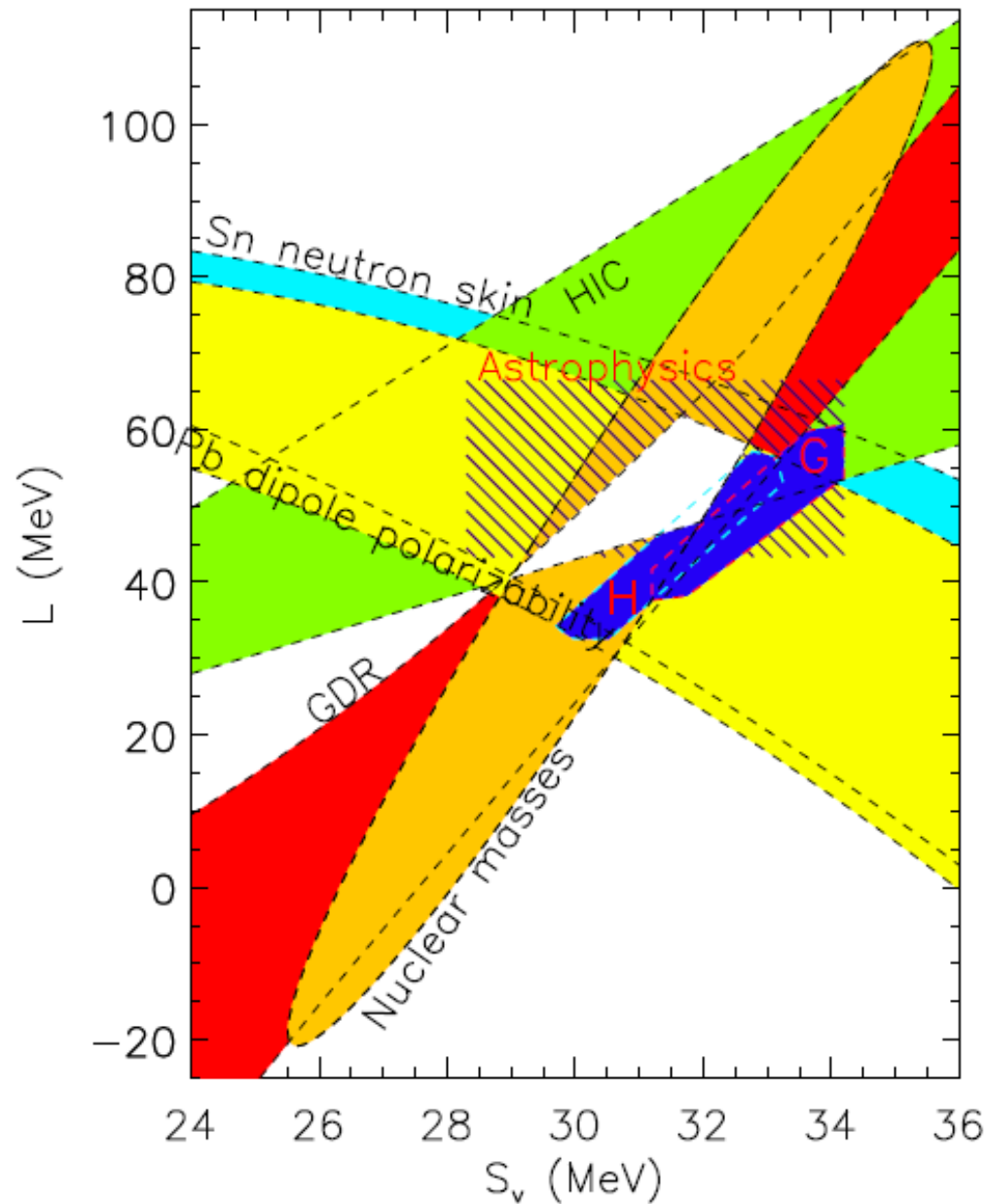
Tsang et al. ('12): NuSYM 2011

Nuclear Symmetry Energy (NuSYM 2013)



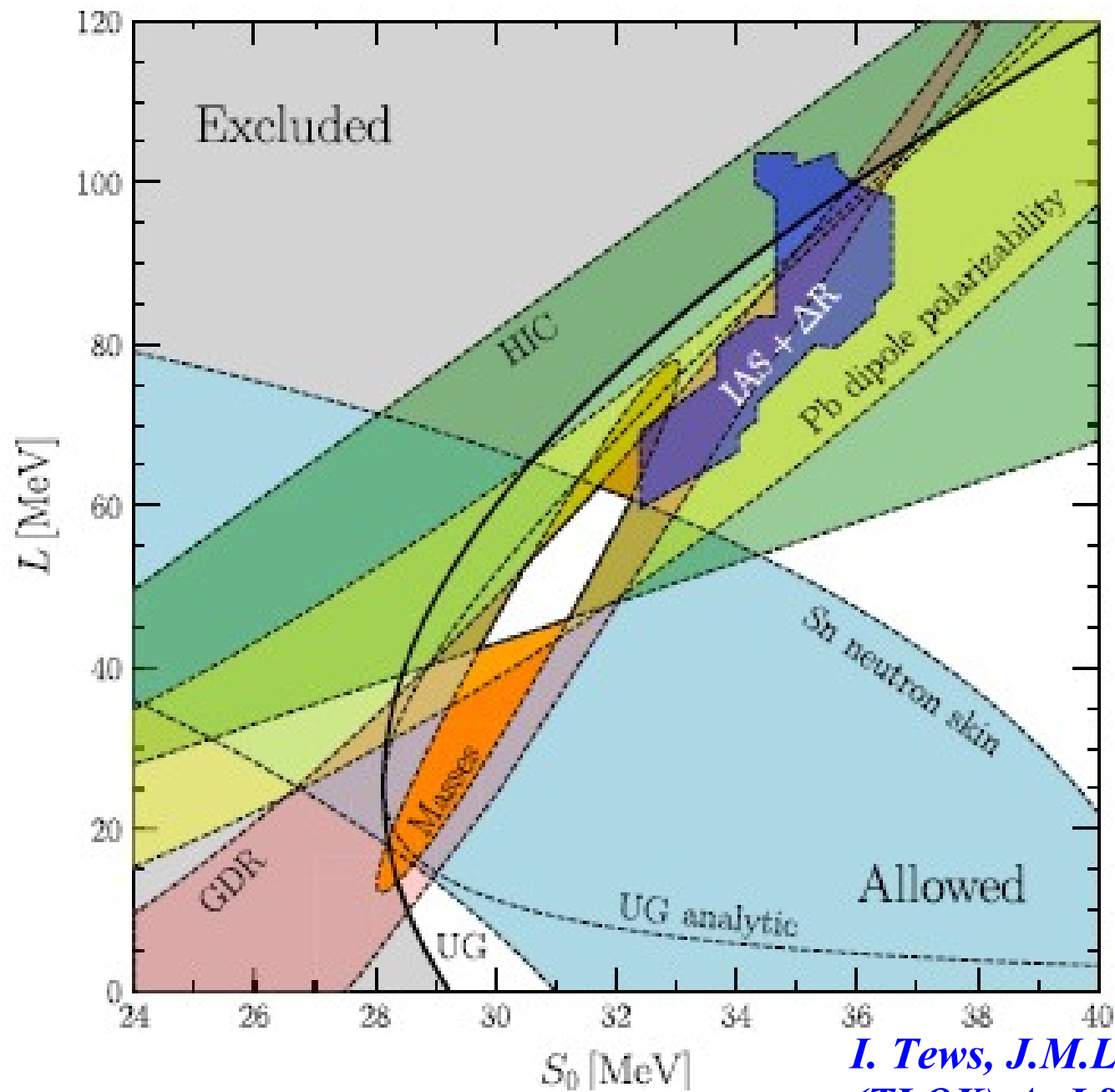
B. A. Li et al. ('13)

Nuclear Symmetry Energy (Lattimer-Lim, 2013)



Lattimer, Lim ('13)

Nuclear Symmetry Energy (Tews+, 2017)

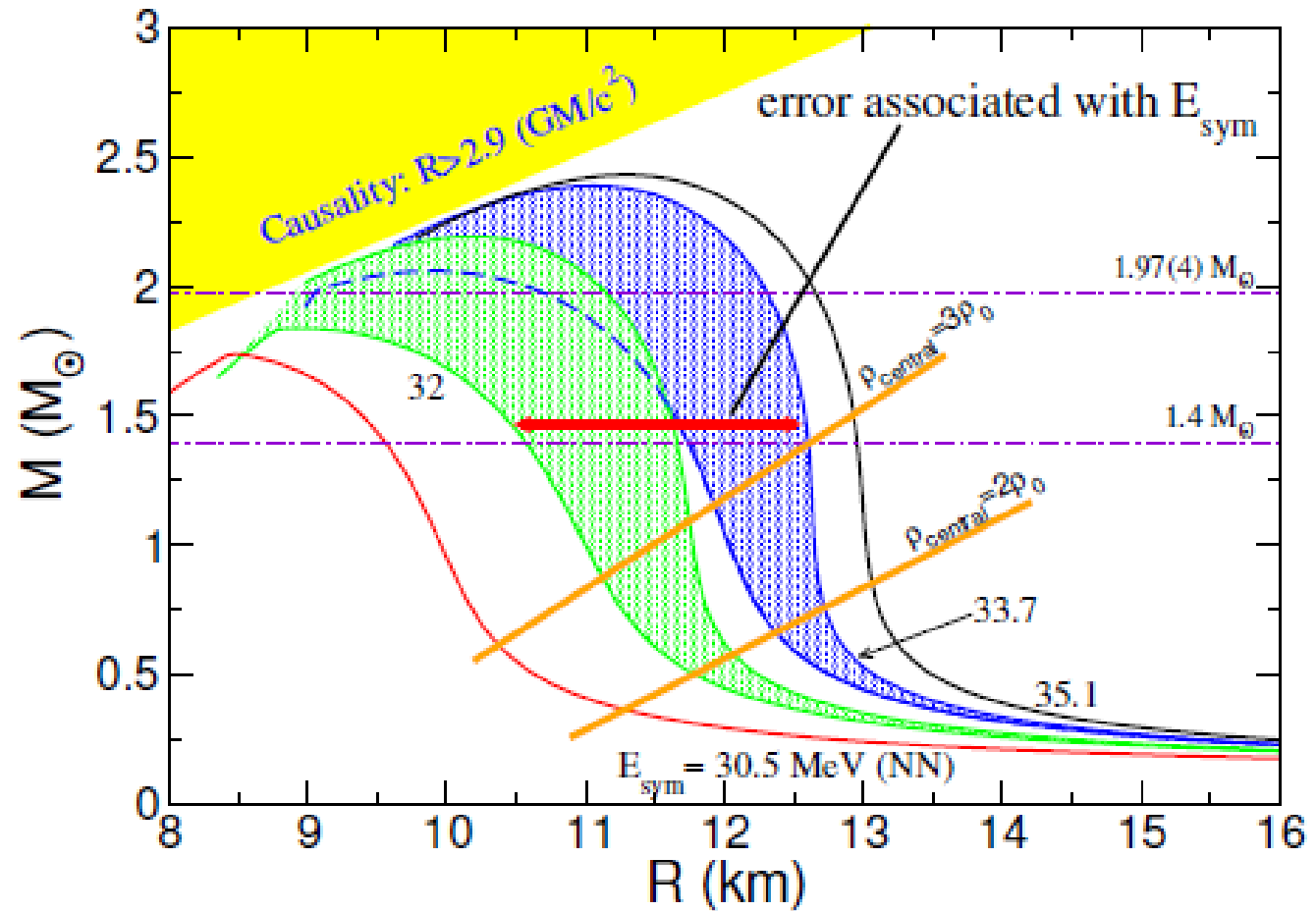


*I. Tews, J.M.Lattimer, A.O., E.E.Kolomeitsev
(TLOK), ApJ 848 ('17)105*

A. Ohnishi @ Nagoya U., Dec.4-6, 2018

Symmetry Energy affects MR Relation of NS

- Nuclear pressure at ρ_0 comes **ONLY** from E_{sym} , then E_{sym} dominates pressure around ρ_0 !
- 5 MeV Difference in E_{sym} results in (3-4) km difference in R_{NS} prediction.



Gandolfi, Carlson, Reddy, PRC 032801, 85 (2012).

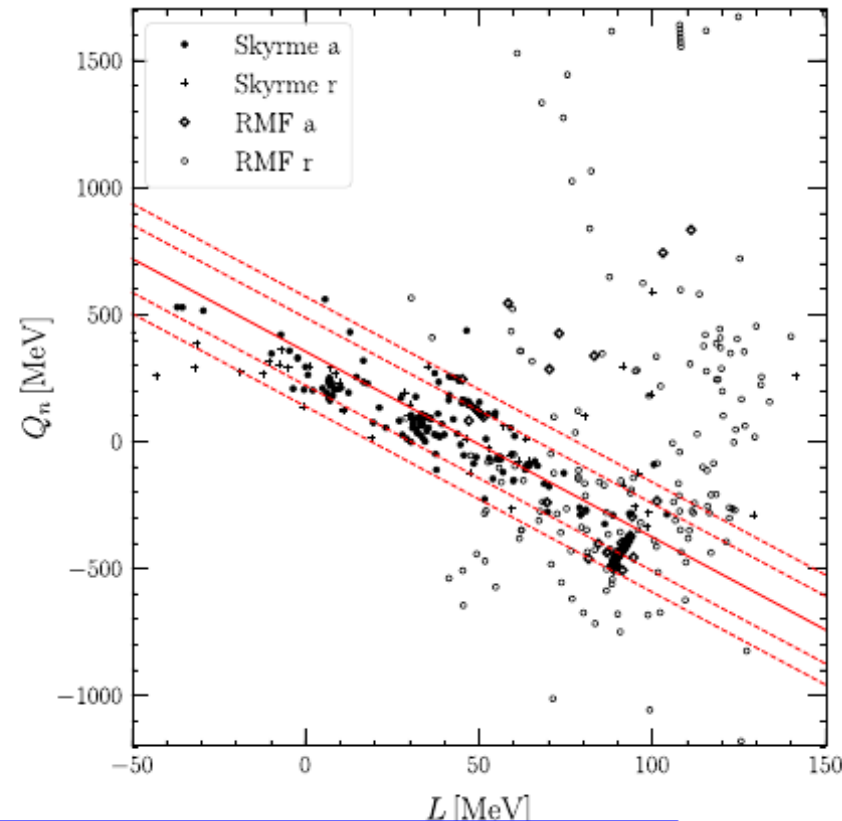
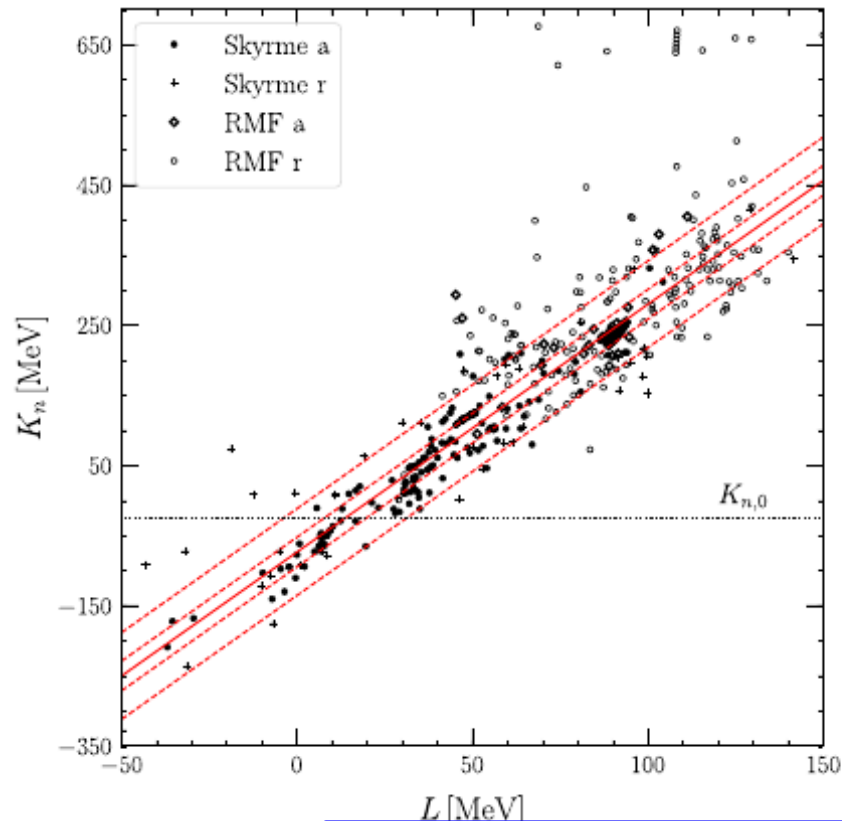
*Higher-Order
Symmetry Energy Parameter Effects
on Nuclear Matter EOS,
and Neutron Star Radius*

Further Constraints on Higher-Order Sym. E. parameters

- K_n and Q_n are correlated with L in “Good” theoretical models.

$$K_n = 3.534L - (74.02 \pm 21.17)\text{MeV}$$

$$Q_n = -7.313L + (354.03 \pm 133.16)\text{MeV}$$



Regard theoretical models as data !

I. Tews, J.M.Lattimer, AO, E.E.Kolomeitsev (TLOK), ApJ 848 ('17)105

Fermi momentum (k_F) expansion

■ Saturation & Symmetry Energy Parameters

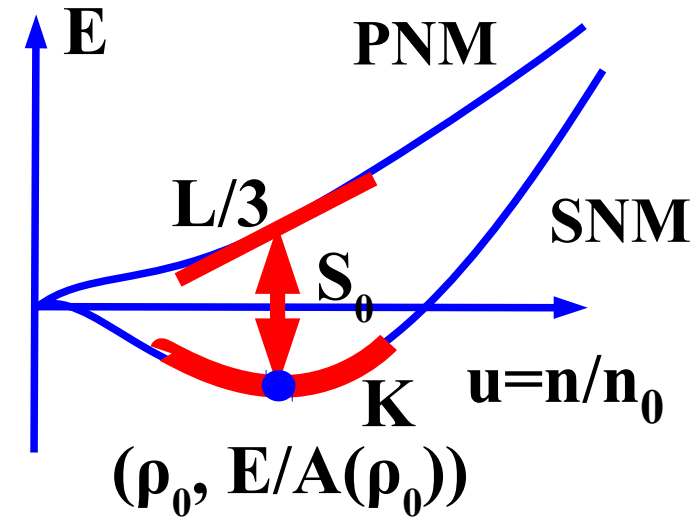
TLOK

$$E_{\text{NM}}(u, \alpha) = E_{\text{SNM}}(u) + \alpha^2 S(u)$$

$$E_{\text{SNM}}(u) \simeq E_0 + \frac{K_0}{18}(u-1)^2 + \frac{Q_0}{162}(u-1)^3$$

$$S(u) \simeq S_0 + \frac{L}{3}(u-1) + \frac{K_s}{18}(u-1)^2 + \frac{Q_s}{162}(u-1)^3$$

$$(u = n/n_0, \alpha = (n_n - n_p)/n)$$



Energy does not approach zero at $n \rightarrow 0$.

■ Fermi momentum expansion ($\sim$ Skyrme type EDF)

- Generated many-body force is given by $k_F \propto u^{1/3}$

$$E_{\text{SNM}}(u) \simeq \boxed{T_0 u^{2/3}} + \underline{a_0 u} + \underline{b_0 u^{4/3}} + \underline{c_0 u^{5/3}} + \underline{d_0 u^2}$$

$$S(u) \simeq \boxed{T_s u^{2/3}} + \underline{a_s u} + \underline{b_s u^{4/3}} + \underline{c_s u^{5/3}} + \underline{d_s u^2}$$

Kin. E. Two-body Density-dep. pot.

m^*

Expansion Coefficients

- Coefficients (a,b,c,d) are represented by Saturation and Symmetry Energy Parameters

TLOK

$a_0 = -4T_0$	$+20E_0$	$+K_0$	$-Q_0/6$
$b_0 = 6T_0$	$-45E_0$	$-5K_0/2$	$+Q_0/2$
$c_0 = -4T_0$	$+36E_0$	$+2K_0$	$-Q_0/2$
$d_0 = T_0$	$-10E_0$	$-K_0/2$	$+Q_0/6$
$a_s = -4T_s$	$+20S_0 - 19L/3$	$+K_s$	$-Q_s/6$
$b_s = 6T_s$	$-45S_0 + 15L$	$-5K_s/2$	$+Q_s/2$
$c_s = -4T_s$	$+36S_0 - 12L$	$+2K_s$	$-Q_s/2$
$d_s = T_s$	$-10S_0 + 10L/3$	$-K_s/2$	$+Q_s/6$

$$\left(T_0 = \frac{3}{5} \frac{\hbar^2 k_F (n_0)^2}{2m}, \quad T_s = T_0 (2^{1/3} - 1) \right)$$

Tedious but straightforward calc.

$TLOK+2M_{\odot}$ constraints

■ TLOK constraints

- (S_0, L) is in Pentagon.

- (K_n, Q_n) are from TLOK constraint.

- $K_0=(190-270)$ MeV

- (n_0, E_0) is fixed

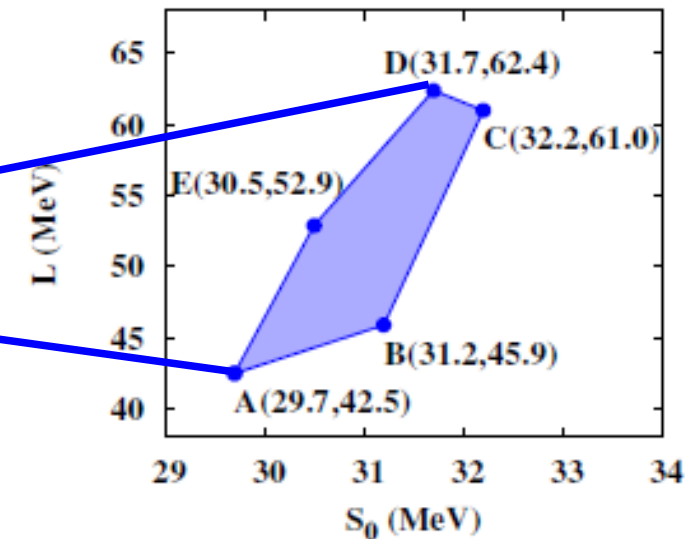
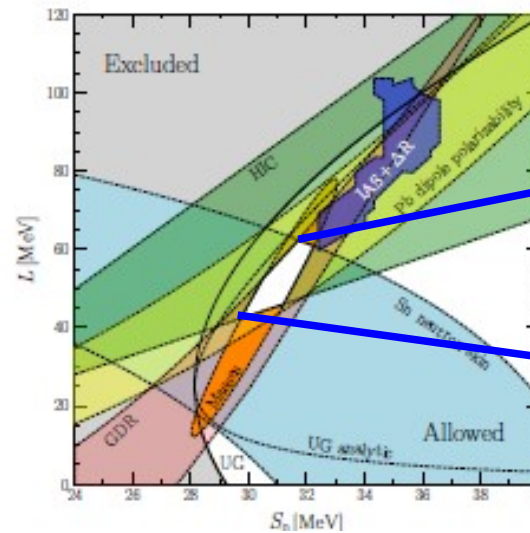
$n_0=0.164$ fm⁻³, $E_0=-15.9$ MeV (small uncertainties)

- Q_0 is taken to kill d_0 parameter

(Coef. of u^2 . Sym. N. M. is not very stiff at high-density)

■ $2 M_{\odot}$ constraint

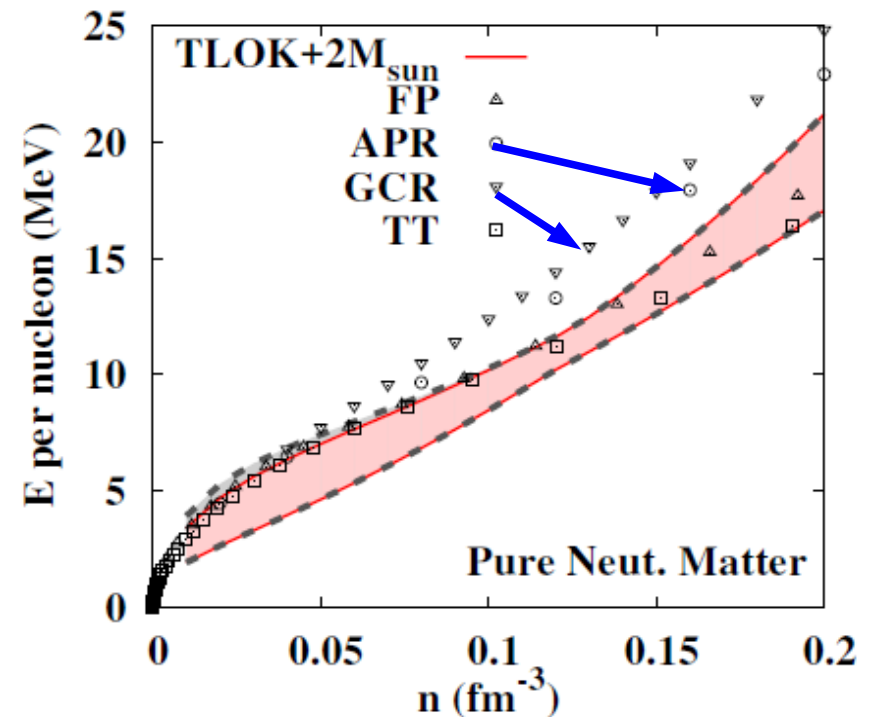
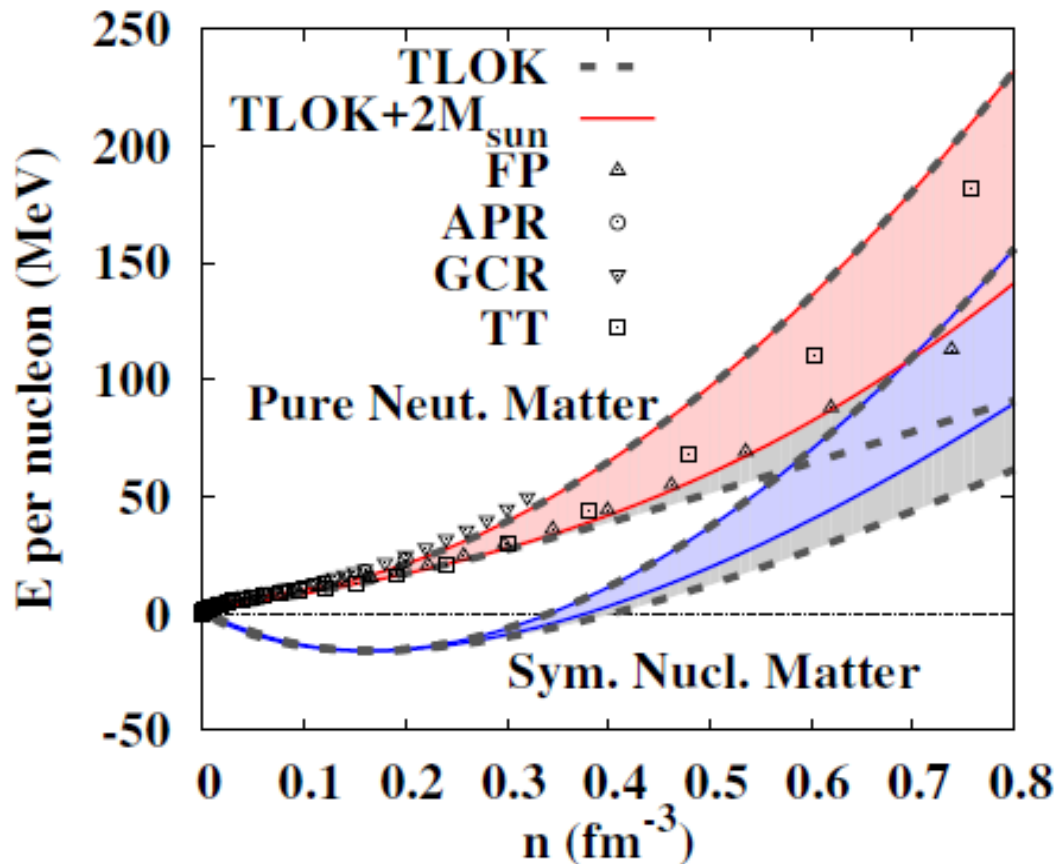
- EOS should support $2 M_{\odot}$ neutron stars.



AO, Kolomeitsev, Lattimer, Tews, Wu (OKLTW), in prog.

$TLOK+2M_{\odot}$ constraints on EOS

- $2M_{\odot}$ constraint narrows the range of EOS.
- Consistent with FP and TT(Togashi-Takano) EOSs.
- APR and GCR(Gandolfi-Carlson-Reddy) EOSs seems to have larger S_0 values.



OKLTW, in prog.

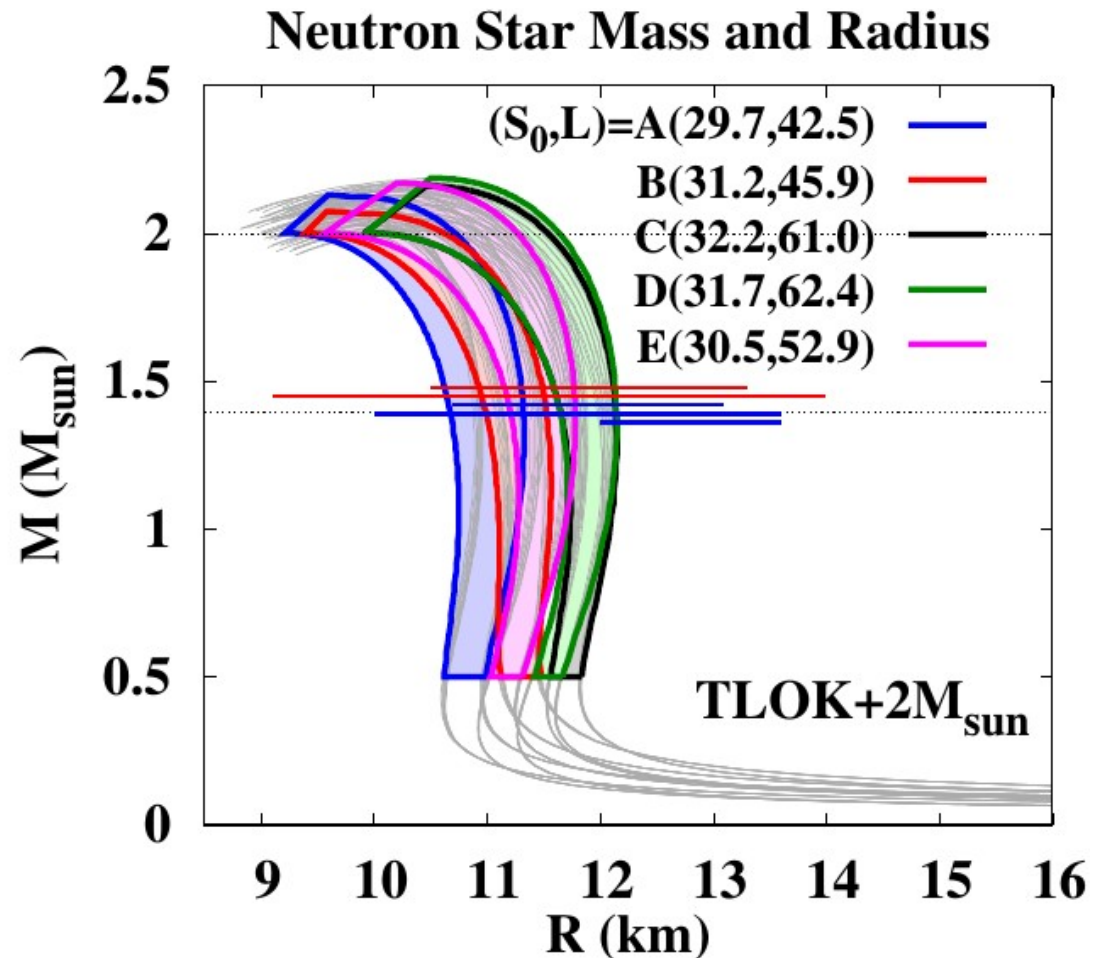
Neutron Star MR curve

■ TLOK + $2 M_{\odot}$ constraints $\rightarrow R_{1.4}=(10.6-12.2)$ km *OKLTW, in prog.*

- E and P are linear fn. of Sat. & Sym. E. parameters
 $\rightarrow$ Min./Max. appears at the corners of pentagon (ABCDE).

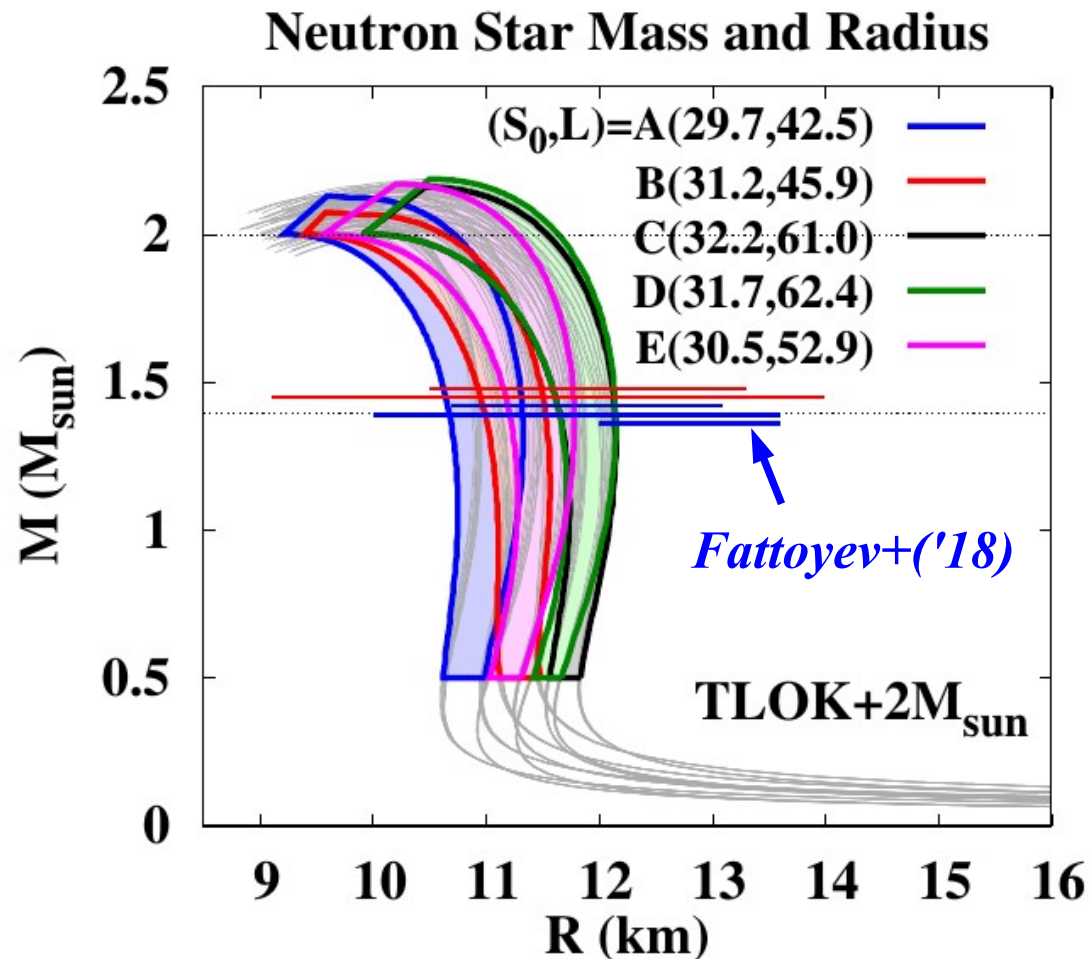
- For a given (S_0, L) ,
unc. of $R_{1.4} \sim 0.5$ km
= unc. from higher-order parameters

- Unc. from $(S_0, L) \sim 1.1$ km
 $\rightarrow$ We still need to fix
 (S_0, L) more precisely.

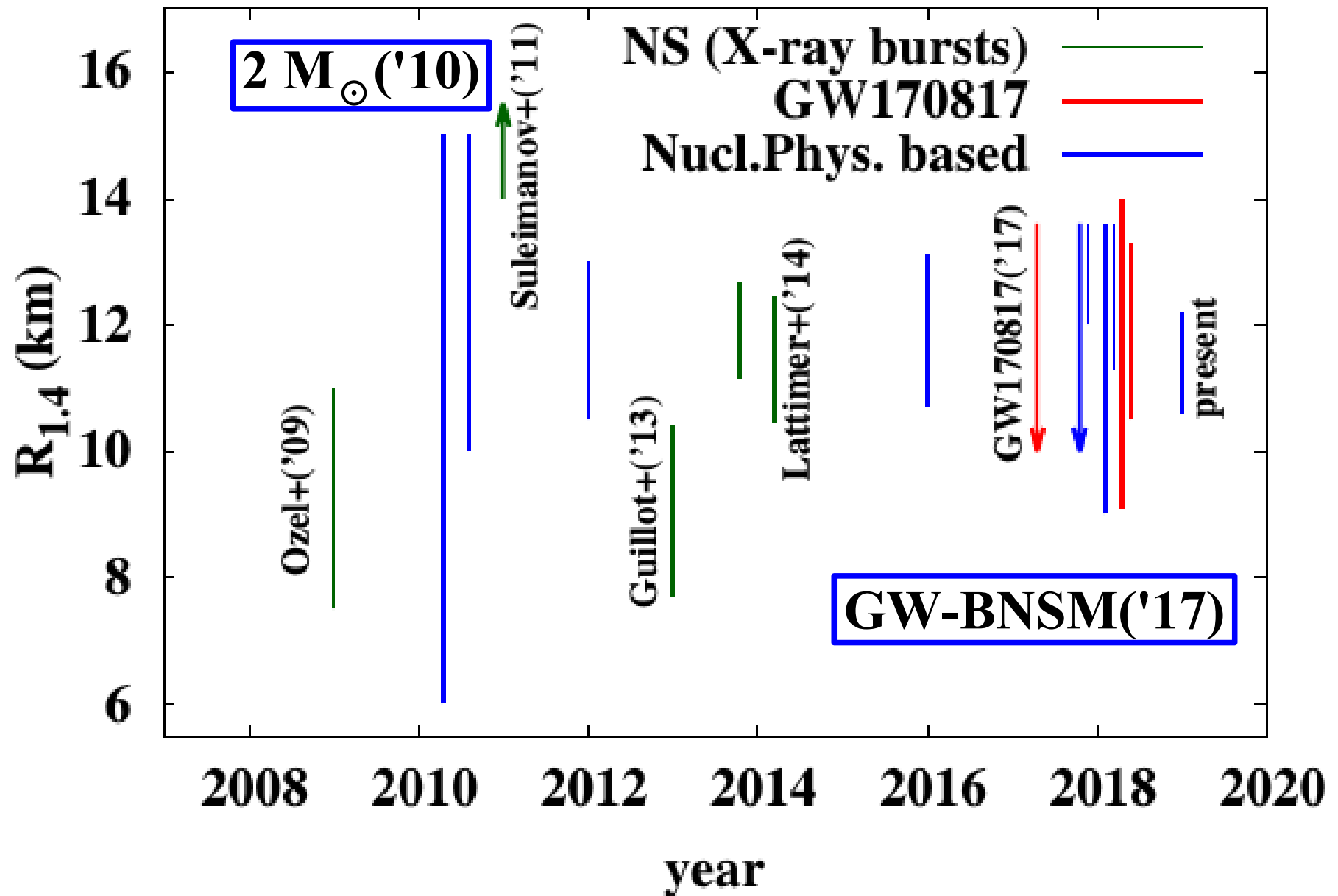


Neutron Star MR curve

- Our constraint is consistent with many of previous ones.
 - $R_{1.4}=(10.6-12.2)$ km *Present work (TLOK + 2 $M_{\odot}$) OKLTW, in prog.*
 - LIGO-Virgo (Tidal deformability Λ from BNSM)
 - (10.5-13.3) km *Abbott+('18b)*
 - (9.1-14.0) km *De+('18) (Λ)*
 - Theoretical Estimates
 - (10.7-13.1) km
Lattimer+, PRep.621('16)127
 - (10.0-13.6) km
Annala+, PRL120('18)172703
 - (9-13.6) km
Tews+, PRC98 ('18)045804
 - (12.0-13.6) km
F.J.Fattoyev+(PREX), PRL120 ('18)172702



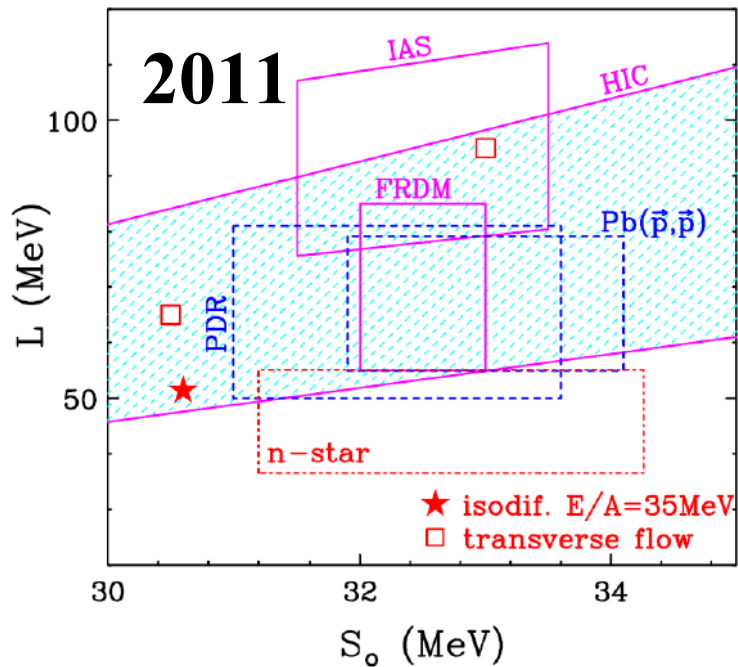
Time dependence of Neutron Star Radius ($R_{1.4}$)



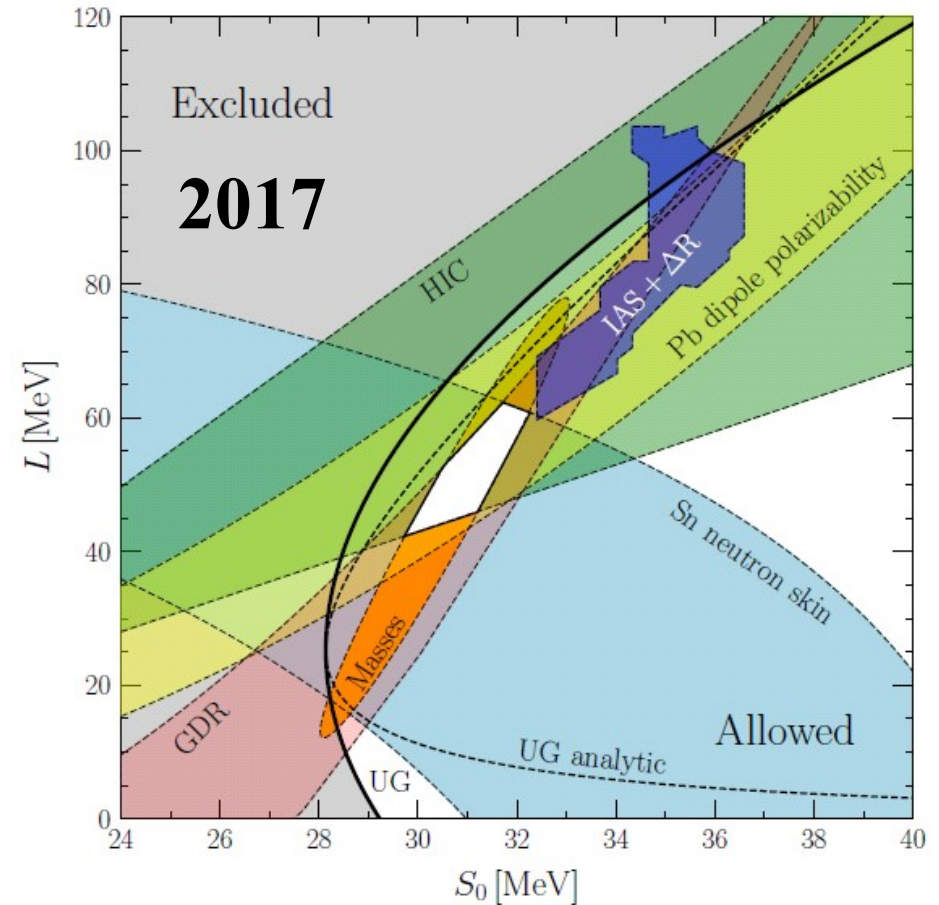
対称エネルギーパラメータ

■ Symmetry Energy Parameters

$$S(u) = S_0 + \frac{L}{3} (u-1) + \frac{K_{\text{sym}}}{18} (u-1)^2 + \mathcal{O}[(u-1)^3] \quad (u=n/n_0)$$

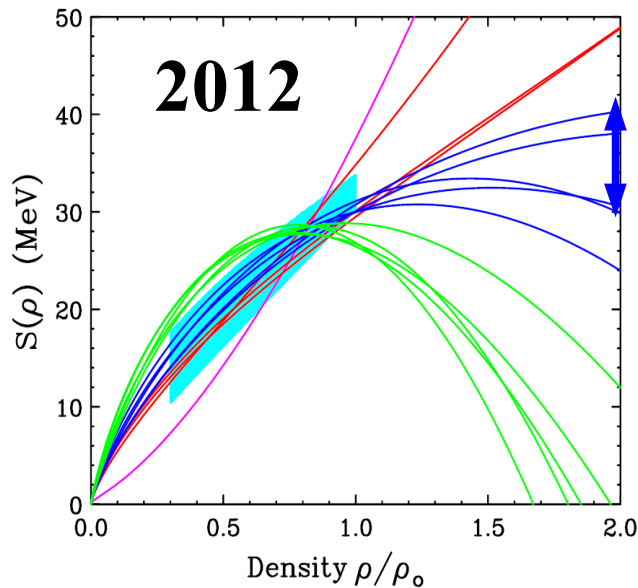


Tsang et al. ('12): NuSYM 2011

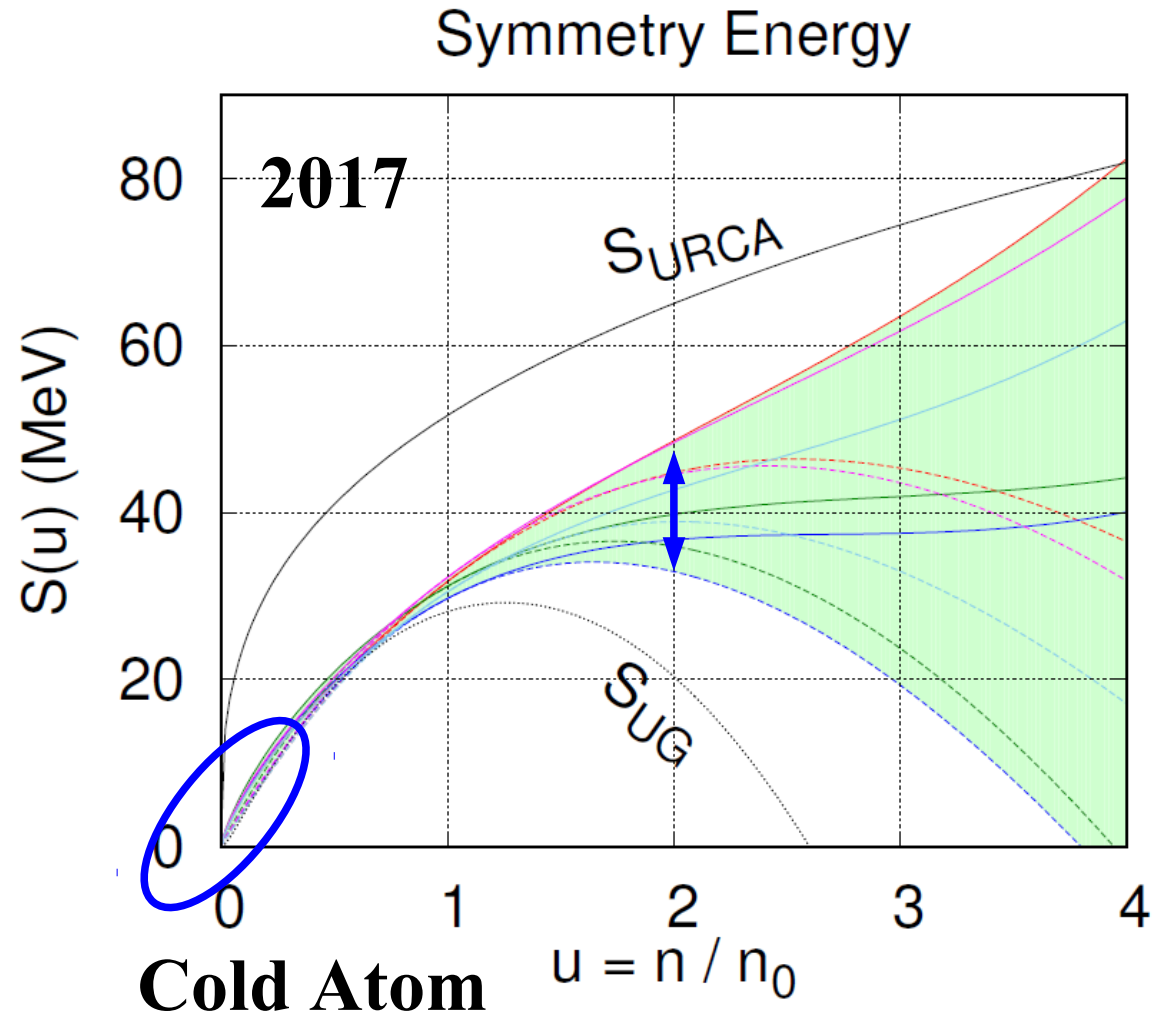


Lattimer, Lim ('13), Lattimer, Steiner ('14)
Tews, Lattimer, AO, Kolomeitsev ('16)

対称エネルギー



*M. B. Tsang et al.,
PRC86 ('12) 015803.*



Based on Tews et al. ('17)

OKLTW, in prog.

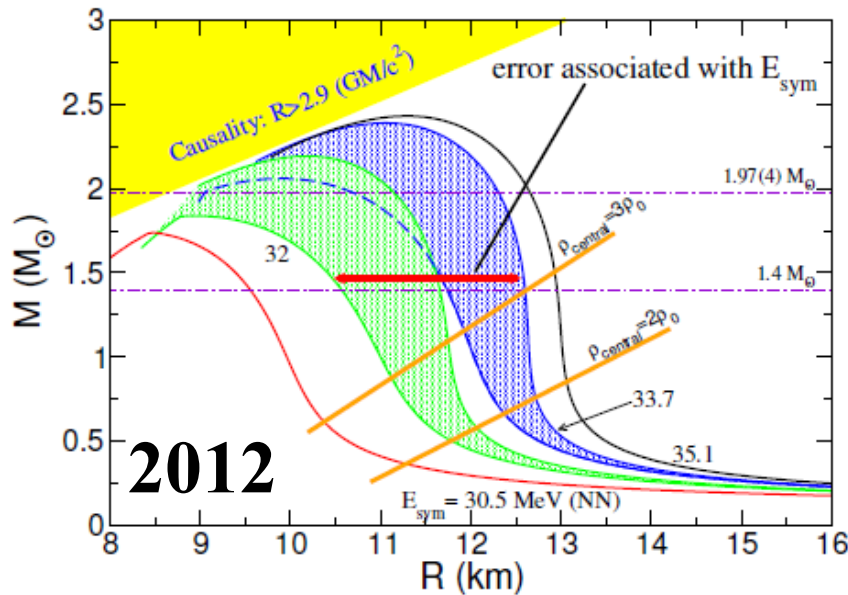
仮定

- (S_0, L) が5 角形の中
- k_F^n ($n=2,3,4,5,6$) で展開
- 2, 3 次の係数 ($K_{\text{sym}}, Q_{\text{sym}}$) が L と相関 (模型からの推定)

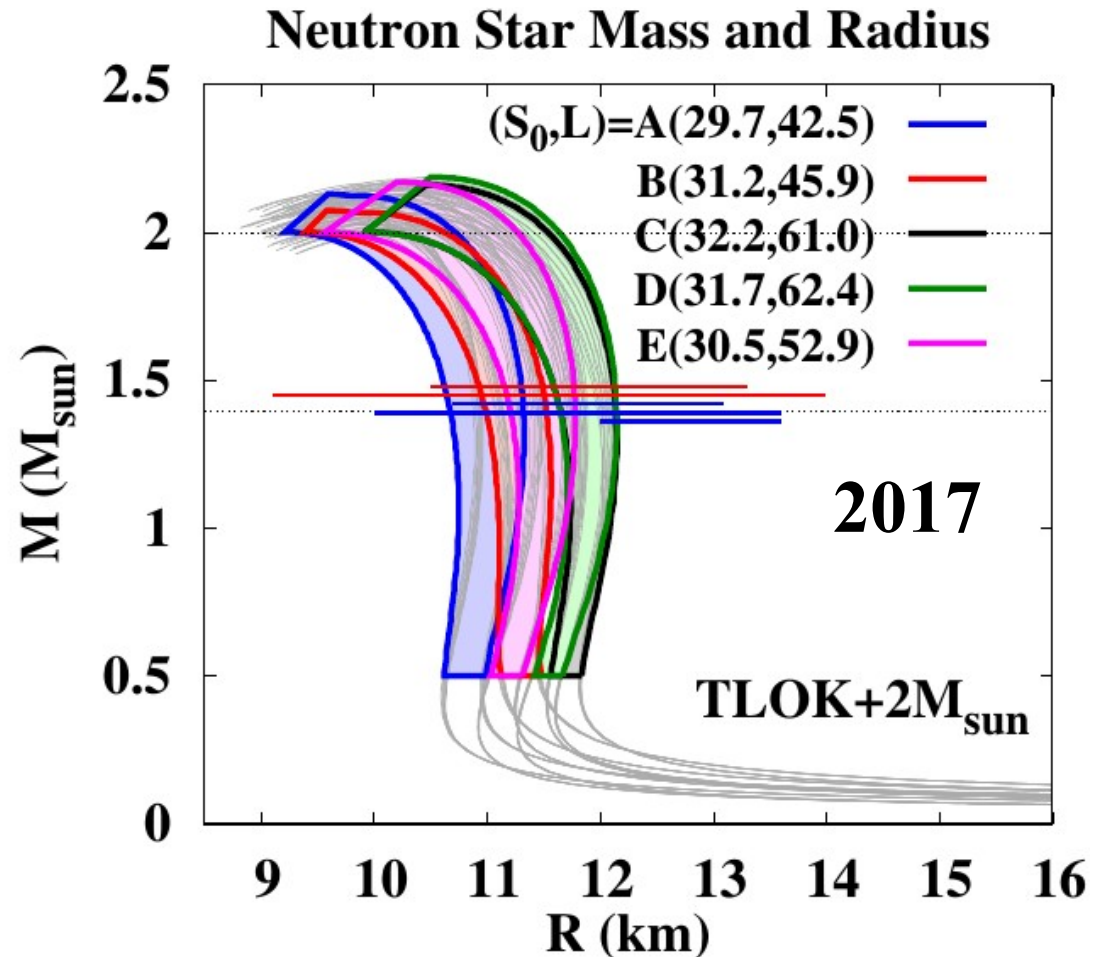
中性子星のMR 曲線

Based on Tews et al. ('17)
OKLTW, in prog.

**$R(1.4) = (10.6 - 12.2)$ km
(unif. matter)**



Gandolfi, Carlson, Reddy,
PRC85('12) 032801.



Experiments & Theories towards Nuclear Symmetry Energy

Nuclear Mass

- Larger symmetry energy $\rightarrow$ B.E. of n-rich nuclei become smaller.

- Volume and surface symmetry energy

$$S(A) = a_a(A) = S_{\text{vol}} - S_{\text{surf}} A^{-1/3}$$

$$B(A, Z) = a_v A - a_s A^{2/3} - a_c \frac{Z^2}{A^{1/3}} - a_a \frac{(N - Z)^2}{A} + a_p \frac{\delta_p}{A^\gamma}$$

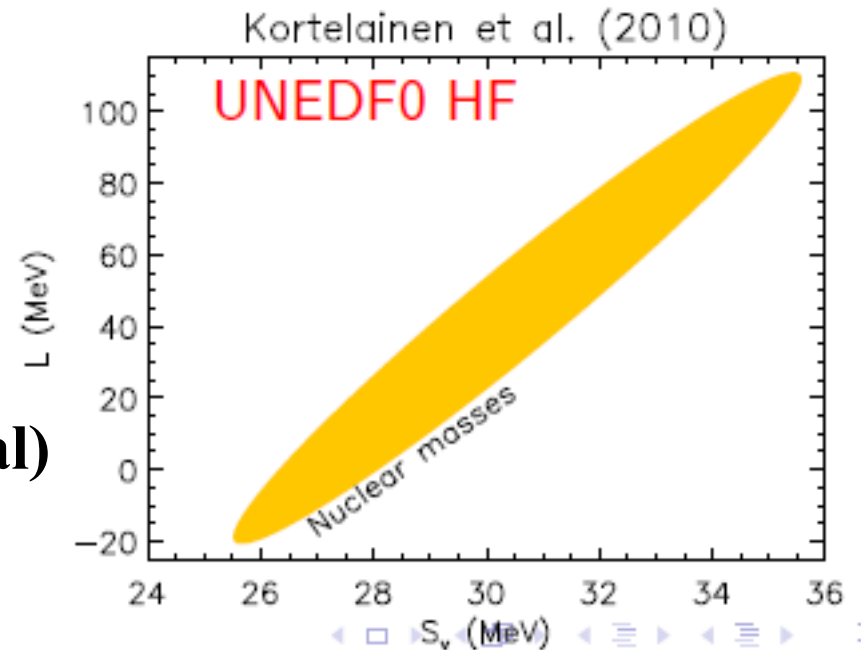
- Finite Range Droplet Model

*P. Moller, W.D. Myers, H. Sagawa, S. Yoshida,
Phys. Rev. Lett. 108, 052501 (2012)*

$$S_v = 32.5 \pm 0.5 \text{ MeV}, L = 70 \pm 15 \text{ MeV}$$

- Density Functional
(Universal Nuclear Density Functional)

Kortelainen et al. (2010)



Lattimer, Lim ('13)

Well, relax (rough idea)

- When L and S_v are linearly correlated as

$$L = a S_v + b,$$

Symmetry energy is given as

$$E_{\text{sym}}(x) = S_v + L(x-1)/3 = S_v (1 + a(x-1)/3) + \text{const. (indep. of } S_v)$$
$$(x = \rho / \rho_0)$$

→ That observable determines symmetry energy most effectively at $x = 1 - 3/a$.

- Nuclear mass: $a \sim 14$

原子核質量は $x \sim 1 - 3/14 = 0.78$ 近辺の対称エネルギーをよく決める

Pigmy Dipole Resonance

■ E1 response of nuclei

- **Giant Resonance:** p and n oscillates collectively.

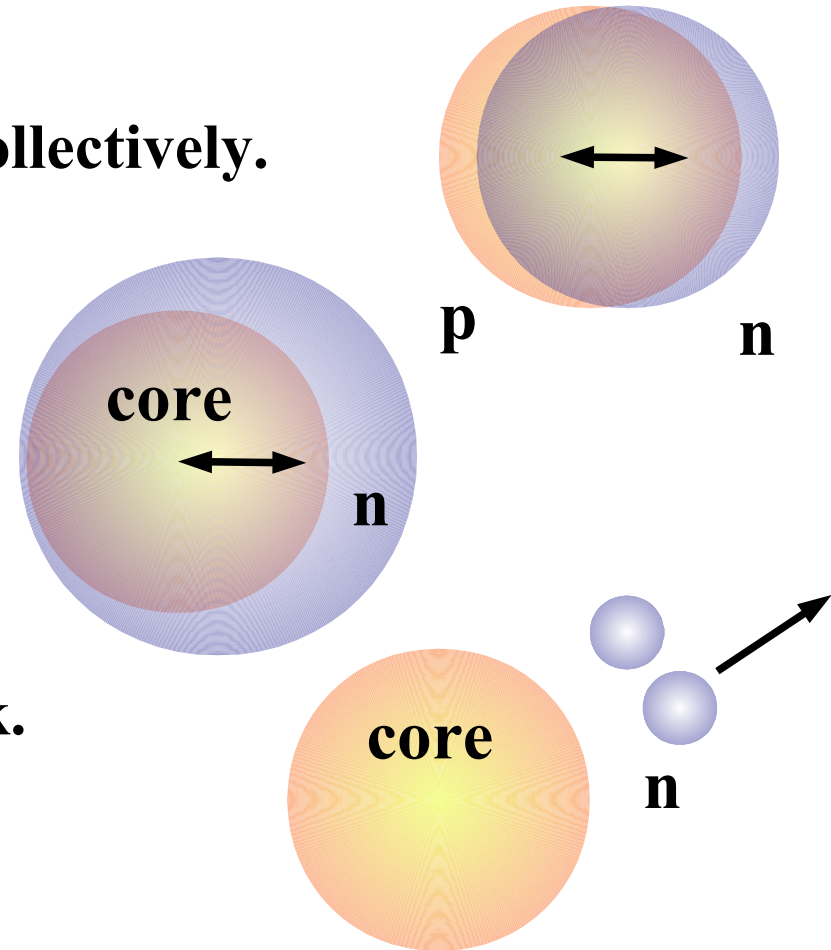
$$E^* \sim 80 A^{-1/3} \text{ MeV}$$

- **Pigmy Dipole Resonance:**
Core oscillates in neutron skin/halo

$$E^* \sim (5-10) \text{ MeV}$$

- **Soft E1 excitation**

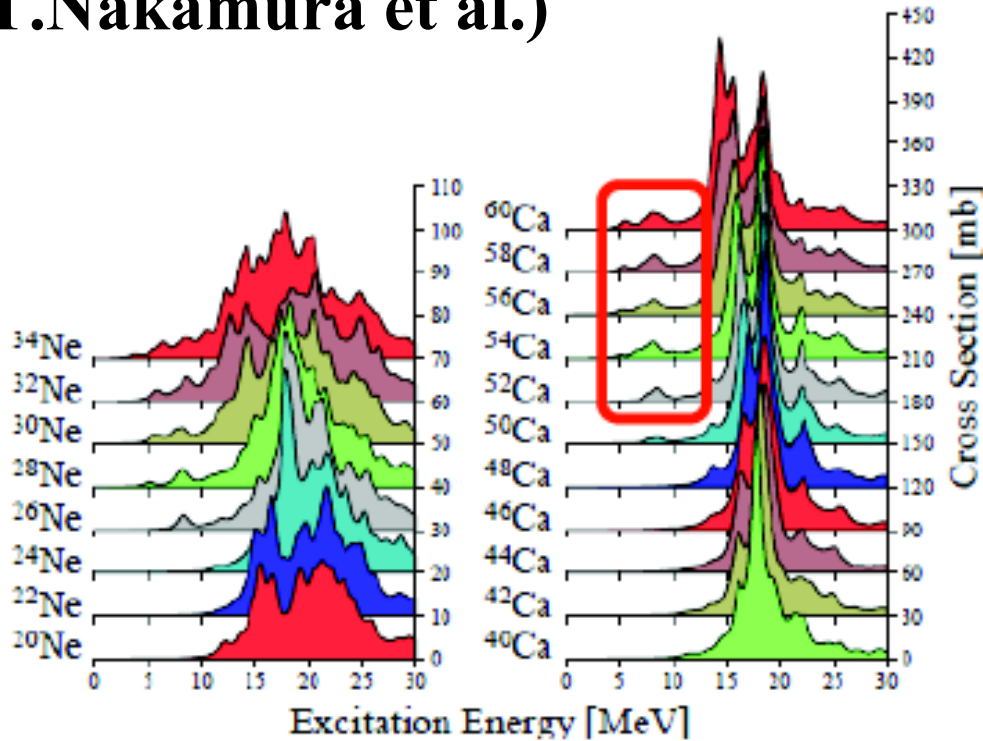
When n wf is extended,
direct dissociation σ also shows a peak.



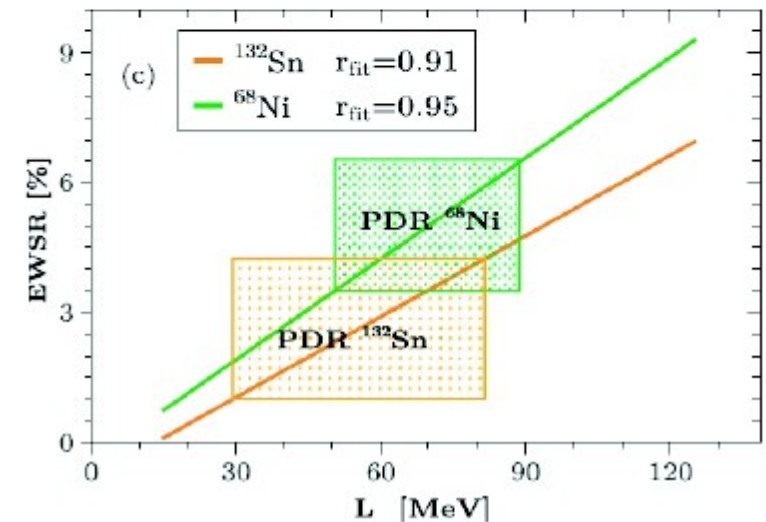
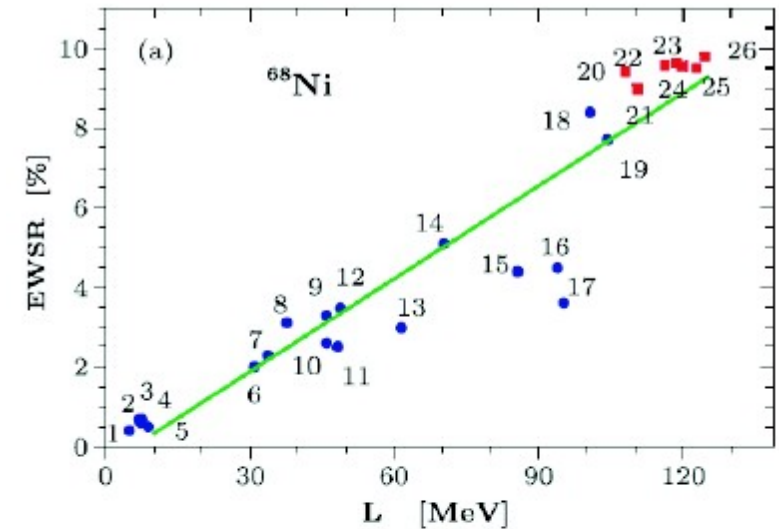
PDR would be sensitive to E_{sym}

PDR of neutron skin nuclei

- Energy Weighted Sum Rule value of PDR would have linear dep. on L
- PDR of very neutron rich nuclei will be measured at RIBF-SAMURAI (T.Nakamura et al.)



T.Inakura, T.Nakatsukasa, K.Yabana
Phys.Rev. C 84, 021302(R) (2011)

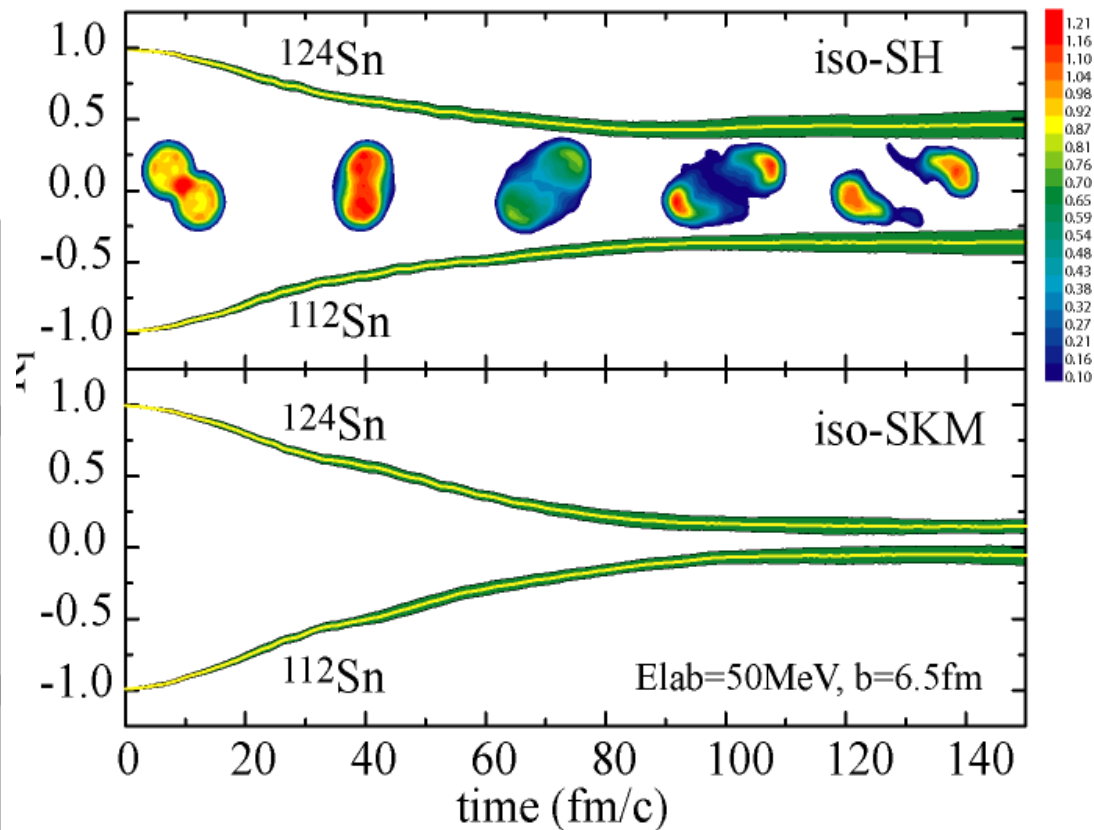
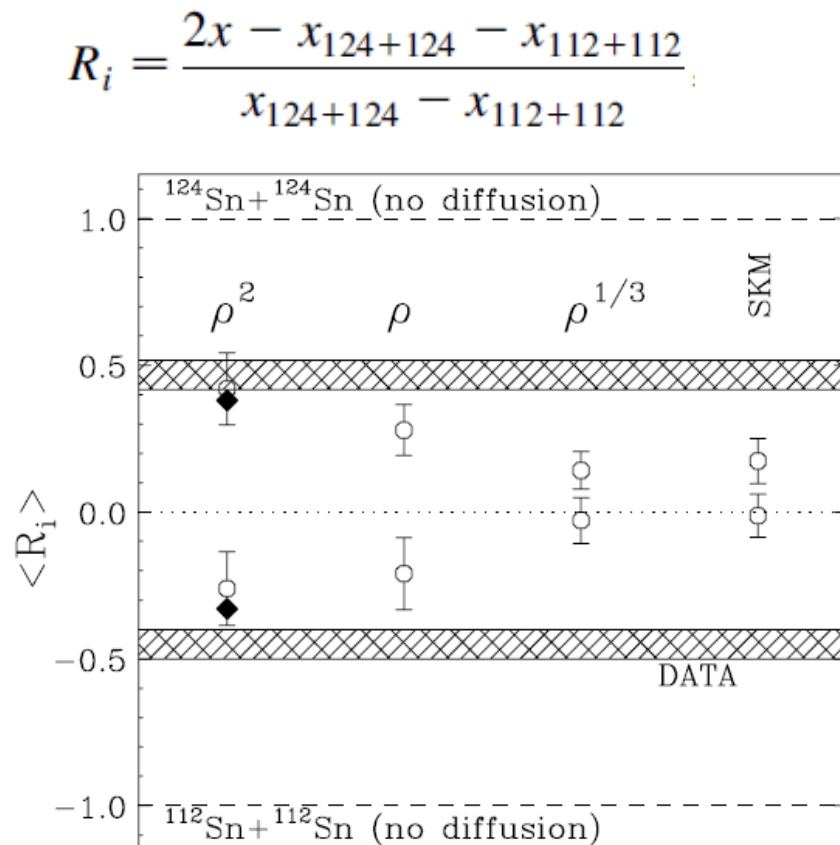


P. Adrich et al., *PRL* 95, 132501
 (2005). (GSI) $^{130,132}\text{Sn}$

O.Wieland et al., *PRL* 102, 092502
 (2009). (GSI) ^{68}Ni

HIC (Isospin Diffusion)

- Collision of nuclei having different n/p ratio
→ fragments with medium n/p ratio will be formed
(Isospin diffusion)
- Driving force of isospin diffusion
= Symmetry energy



Tsang et al. ('04)

Skin Thickness & Dipole Polarizability

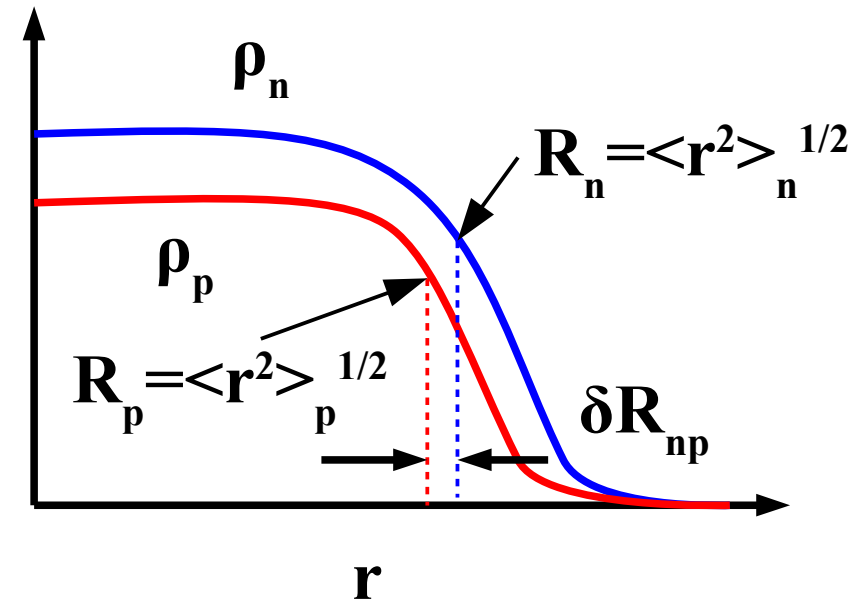
■ Skin Thickness δR_{np}

● Larger L

→ Small E_{sym} at low ρ

Large E_{sym} at high ρ

→ Larger δR_{np}



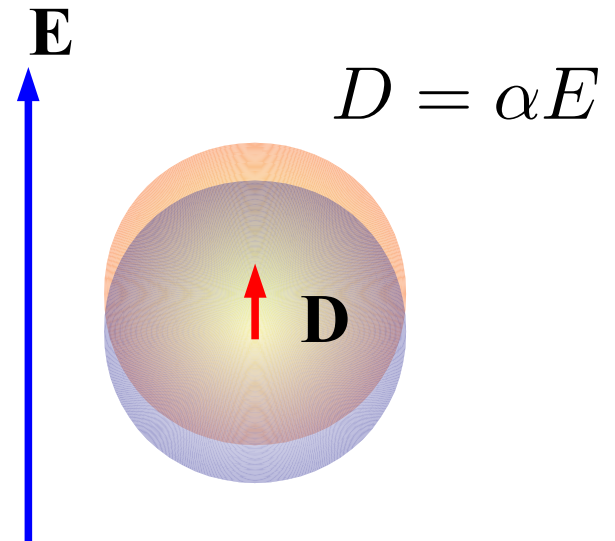
■ Electric Dipole Polarizability α

$$H = H_0 - eE \sum_{i \in p} x_i = H_0 - E \hat{D}$$

$$|\psi\rangle = |0\rangle - \sum_{n>0} \frac{|n\rangle \langle n | V | 0\rangle}{E_n - E_0} + O(E^2)$$

$$D = \langle \psi | \hat{D} | \psi \rangle = E \times 2 \sum_{n>0} \frac{\langle 0 | \hat{D} | n \rangle \langle n | \hat{D} | 0 \rangle}{E_n - E_0}$$

$$\alpha = \frac{8\pi}{9} \int \frac{dB(E1)}{\omega}$$

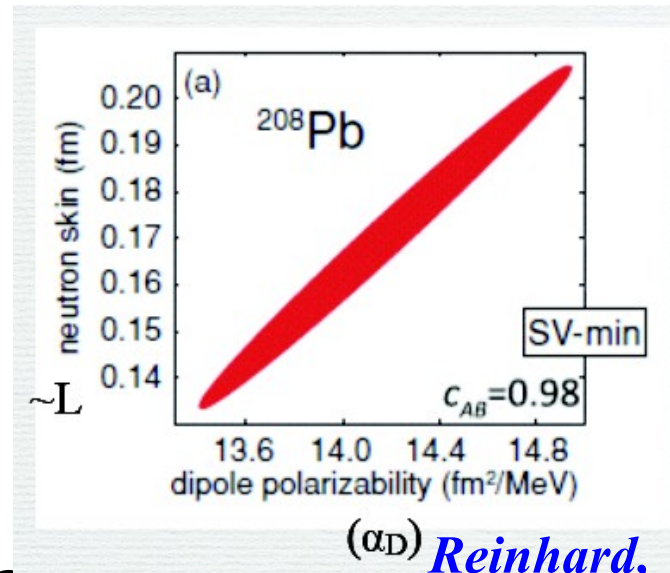


Skin Thickness & Dipole Polarizability

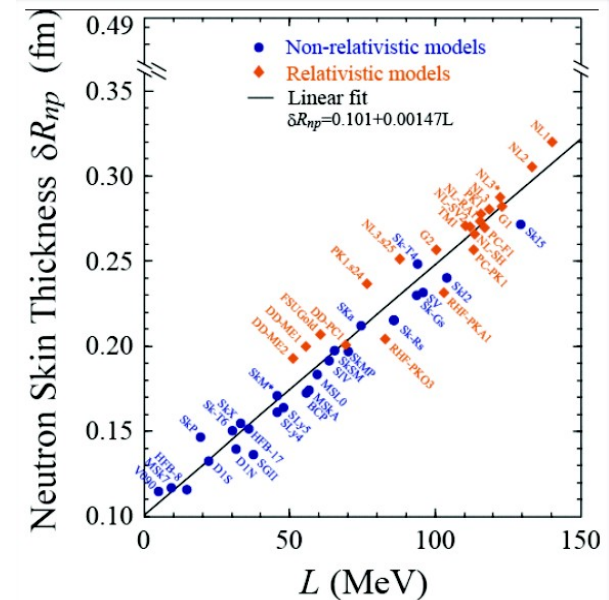
- Strong corr. btw α and skin thickness (smaller restoring force $\rightarrow$ soft)

- Skin thickness is also correlated with L .

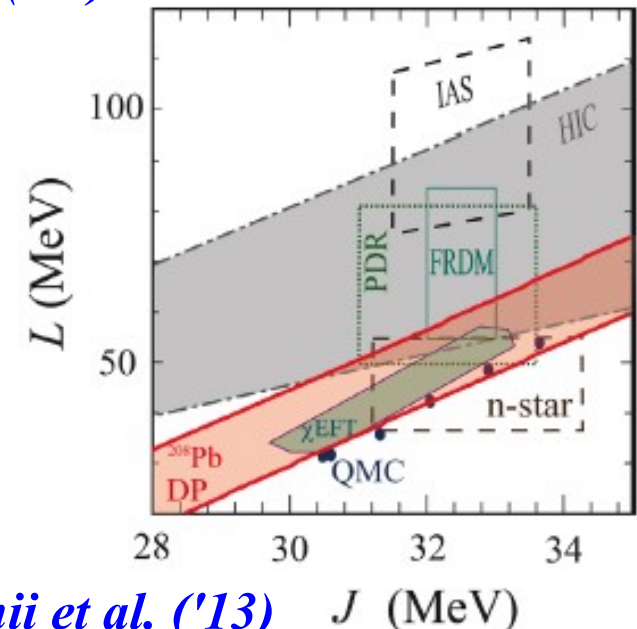
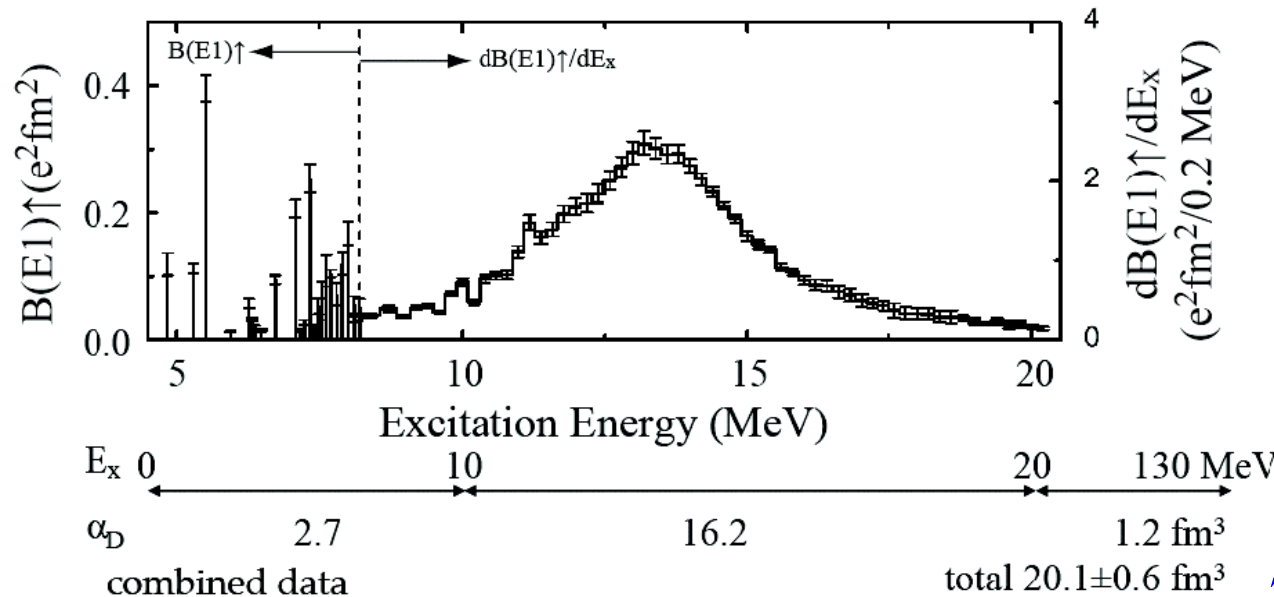
- Precise data from RCNP



Reinhard,
Nazarewicz ('10)



Roca-Maza et al. ('11)



Tamii et al. ('13)

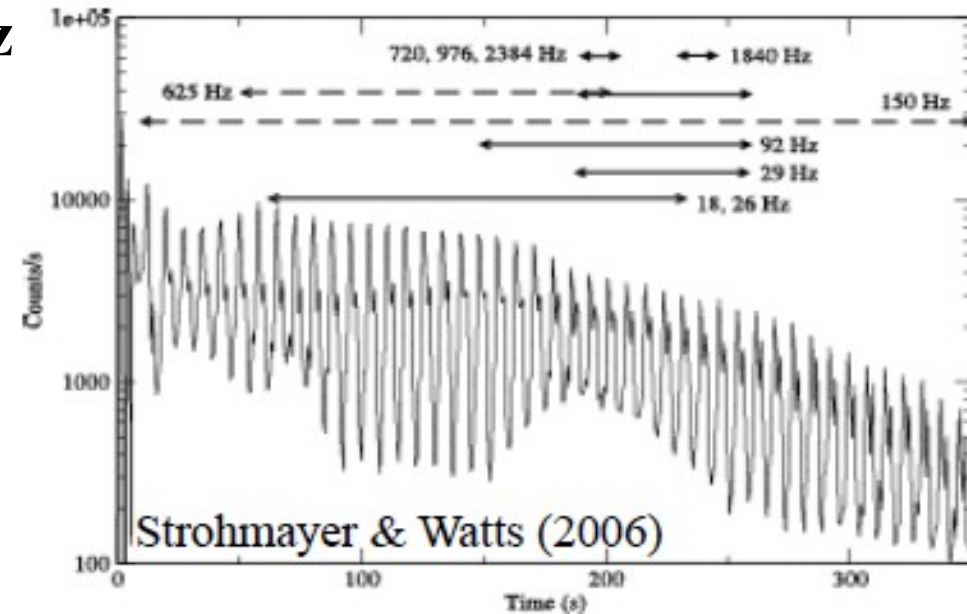
Quasi Periodic Oscillation of Neutron Stars

■ QPOs in afterglow of giant flares from soft-gamma repeaters (SGRs) (*Barat+ 83, Israel+ 05, Strohmayer & Watts 05, Watts & Strohmayer 06*)

- SGR 0526-66 (5th/3/1979) : 43 Hz

- SGR 1900+14 (27th/8/1998) :
28, 54, 84, 155 Hz

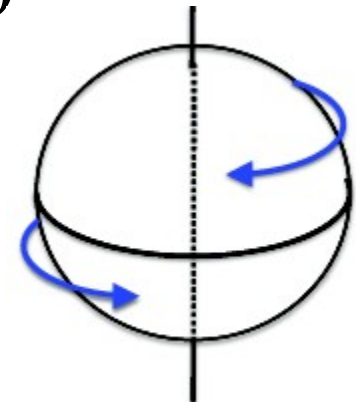
- SGR 1806-20 (27th/12/2004) :
18, 26, 30, 92.5, 150,
626.5, 1837 Hz



■ Asteroseismology

- From star quake to stellar properties (M, R, B, EOS ...)

- Low frequency (e.g. 28 Hz) requires long wave mode.
→ Torsional oscillations of the crust



QPO and Symmetry Energy

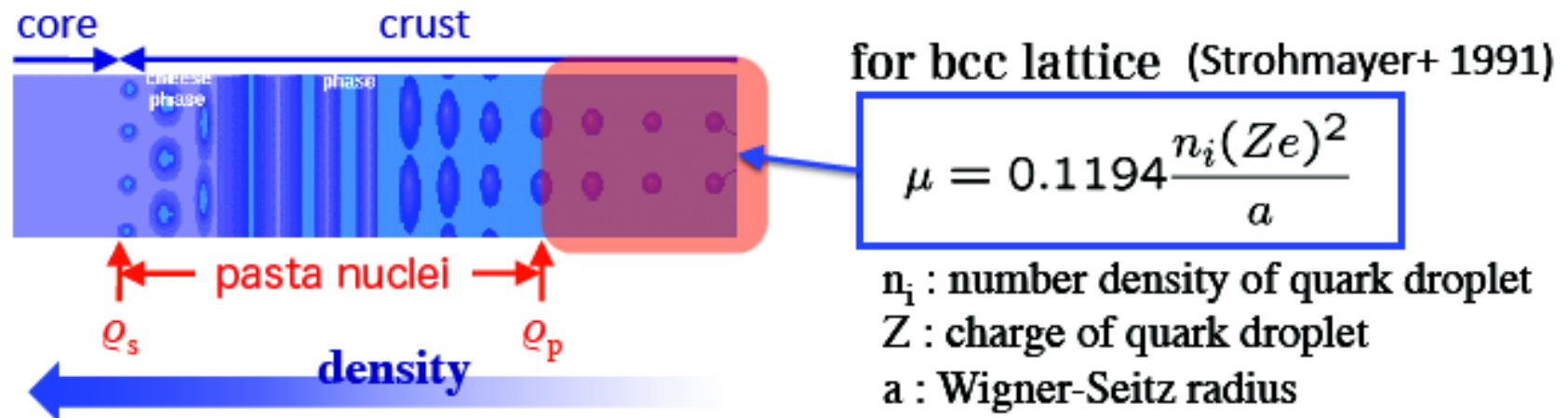
■ Torsional oscillations (ねじれ振動)

- incompressible (no density perturbations)

- Frequency ${}_lt_0 = \sqrt{l(l+1)} \frac{v_s}{2\pi R}, \quad v_s = \sqrt{\mu/\rho}$

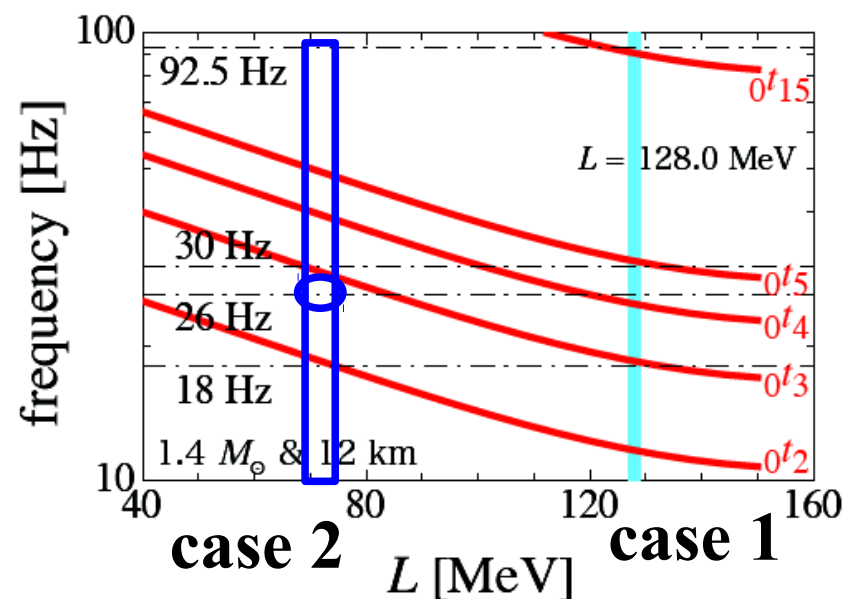
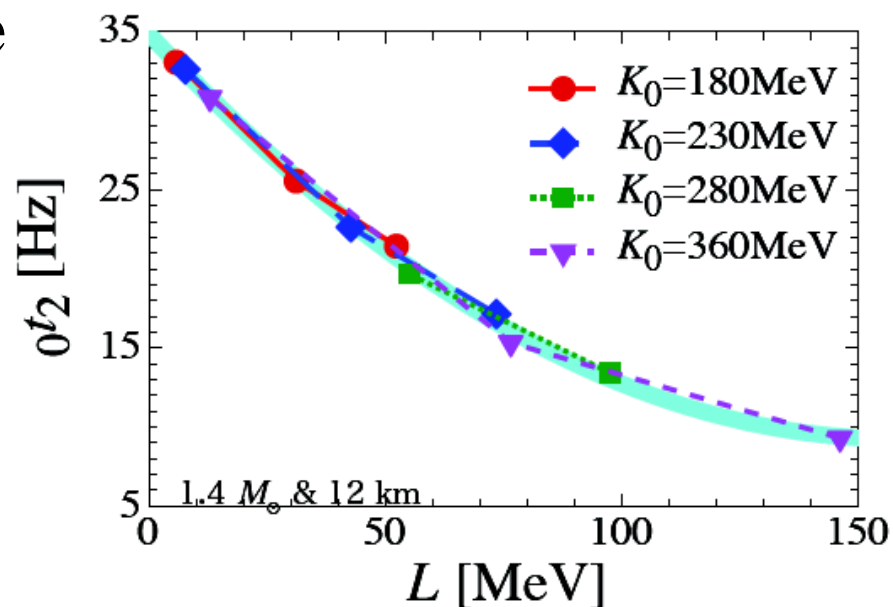
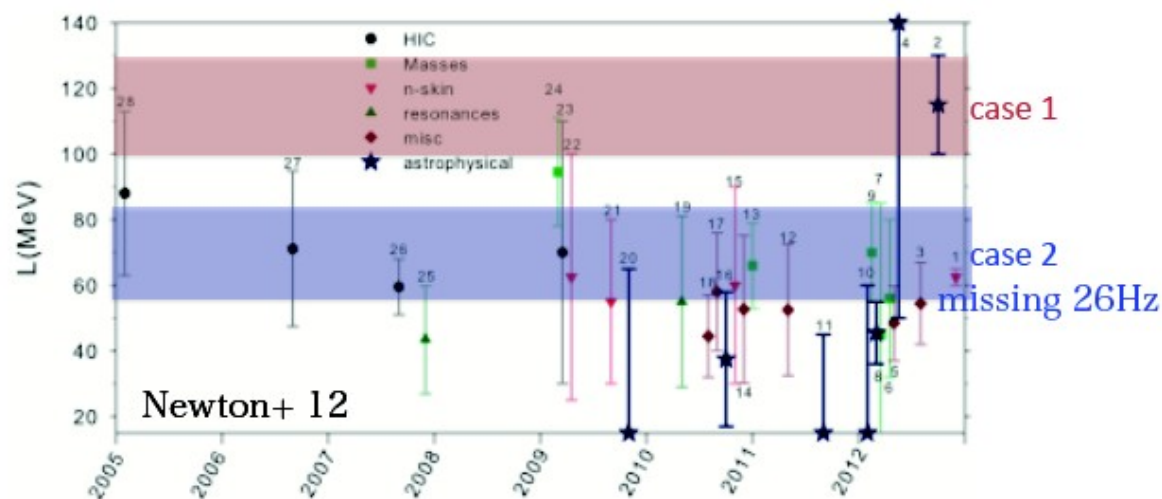
μ : shear modulus, v_s : shear velocity
(Hansen & Cioff 1980)

- Shear modulus of bcc lattice depends on nuclear charge
→ Dependence on the symmetry energy



QPO and Symmetry Energy

- For a given set of (M,R), we can solve TOV equation from the surface.
 - Torsional oscillation frequencies are calculated using EOSs with various (L,K).
- Oyamatsu, Iida ('03,'07)
- Compared with observed freq.
→ constrains on L (small dep. on K)



Large L or missing 26 Hz

Sotani+ '13

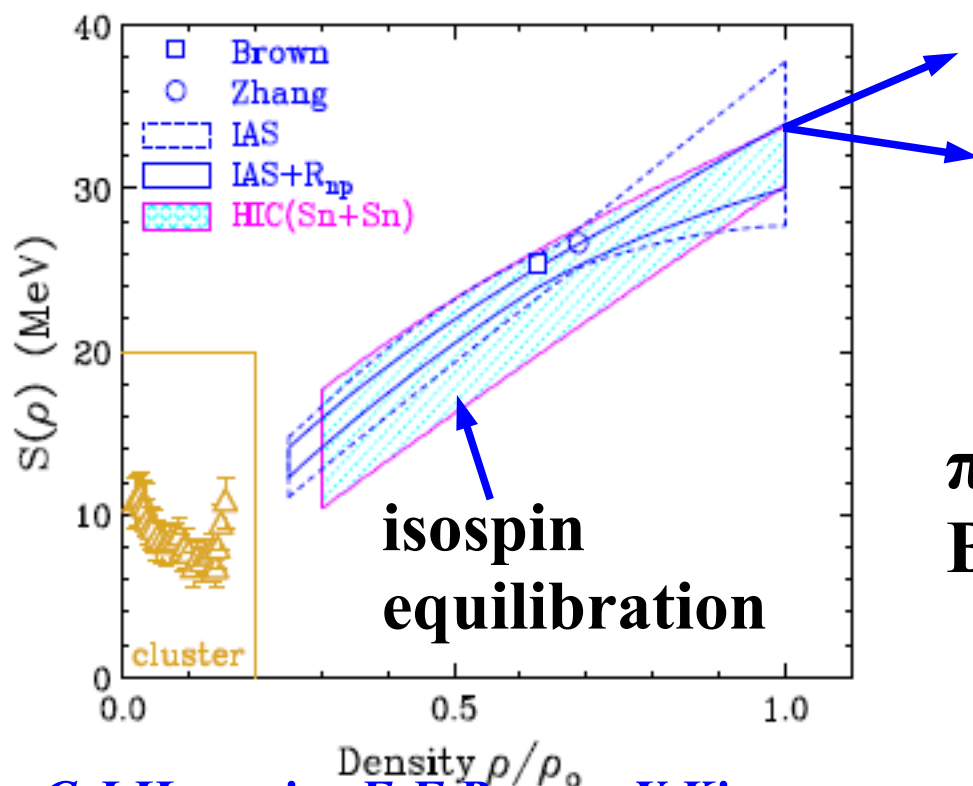
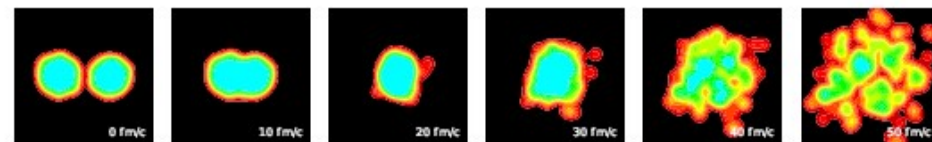
@ Nagoya U., Dec.4-6, 2018

High Density Symmetry Energy

- Symmetry energy at $\rho=(2-4) \rho_0$ dominantly determines NS radius.

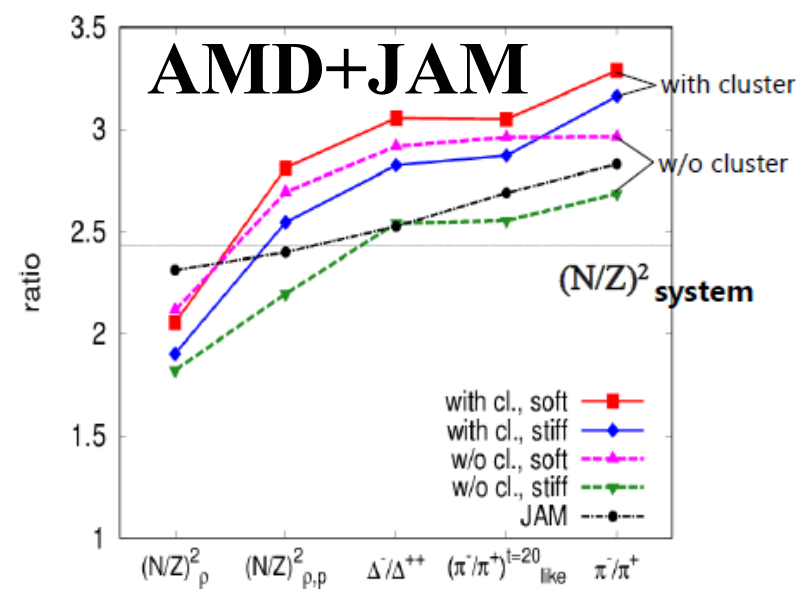
→ Central Heavy-Ion collisions
at a few 100 MeV !

(Li, Trautmann, Murakami, Ono)



*Let's wait for SπRIT
B01 results (Murakami) !
More theor. work needed.*

π^-/π^+
B.A.Li

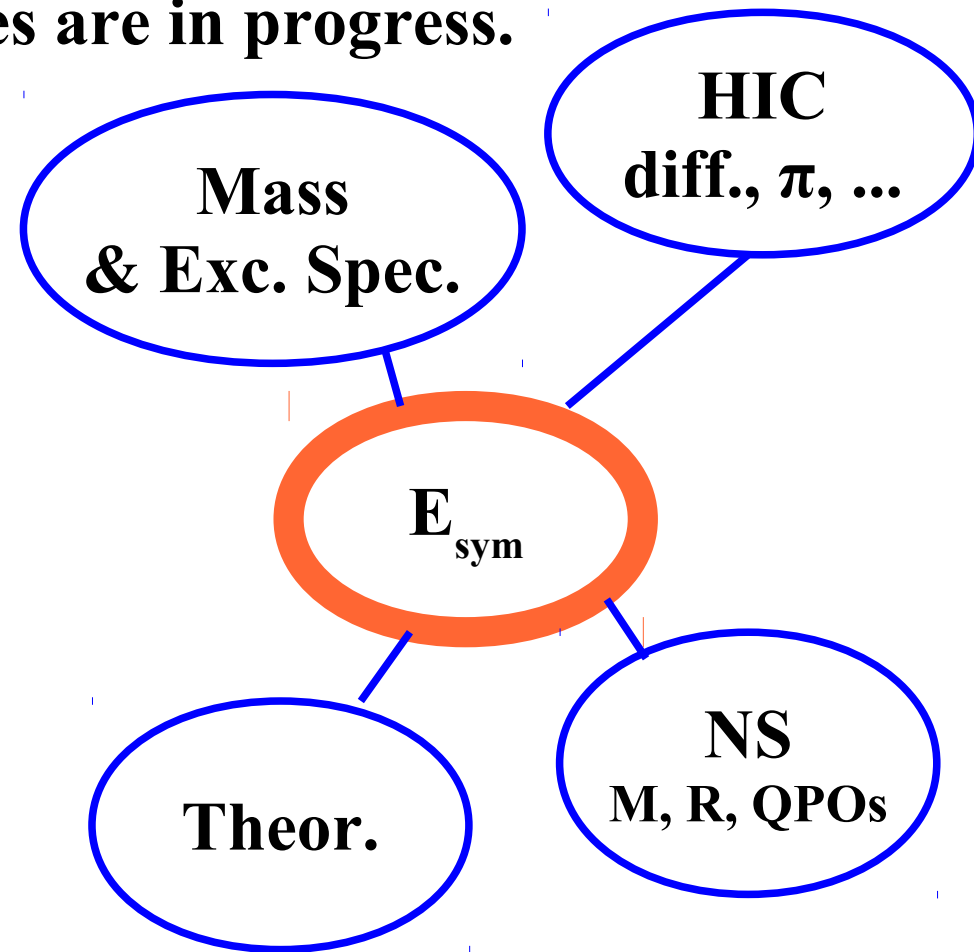
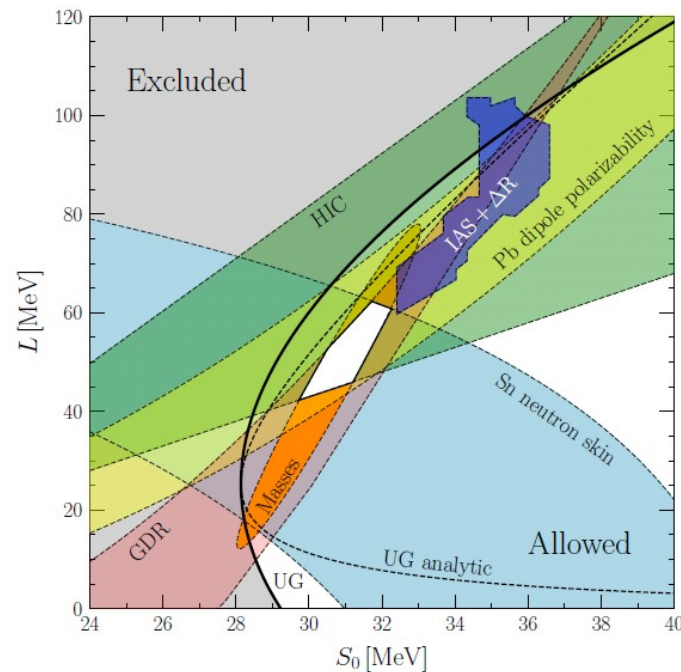


*C.J.Horowitz, E.F.Brown, Y.Kim,
W.G.Lynch, R.Michaels, A. Ono, J.
Piekarewicz, M. B. Tsang, H.H.Wolter
(NuSYM13), JPG41('14) 093001*

Ikeno, Ono, Nara, AO ('16)

Summary

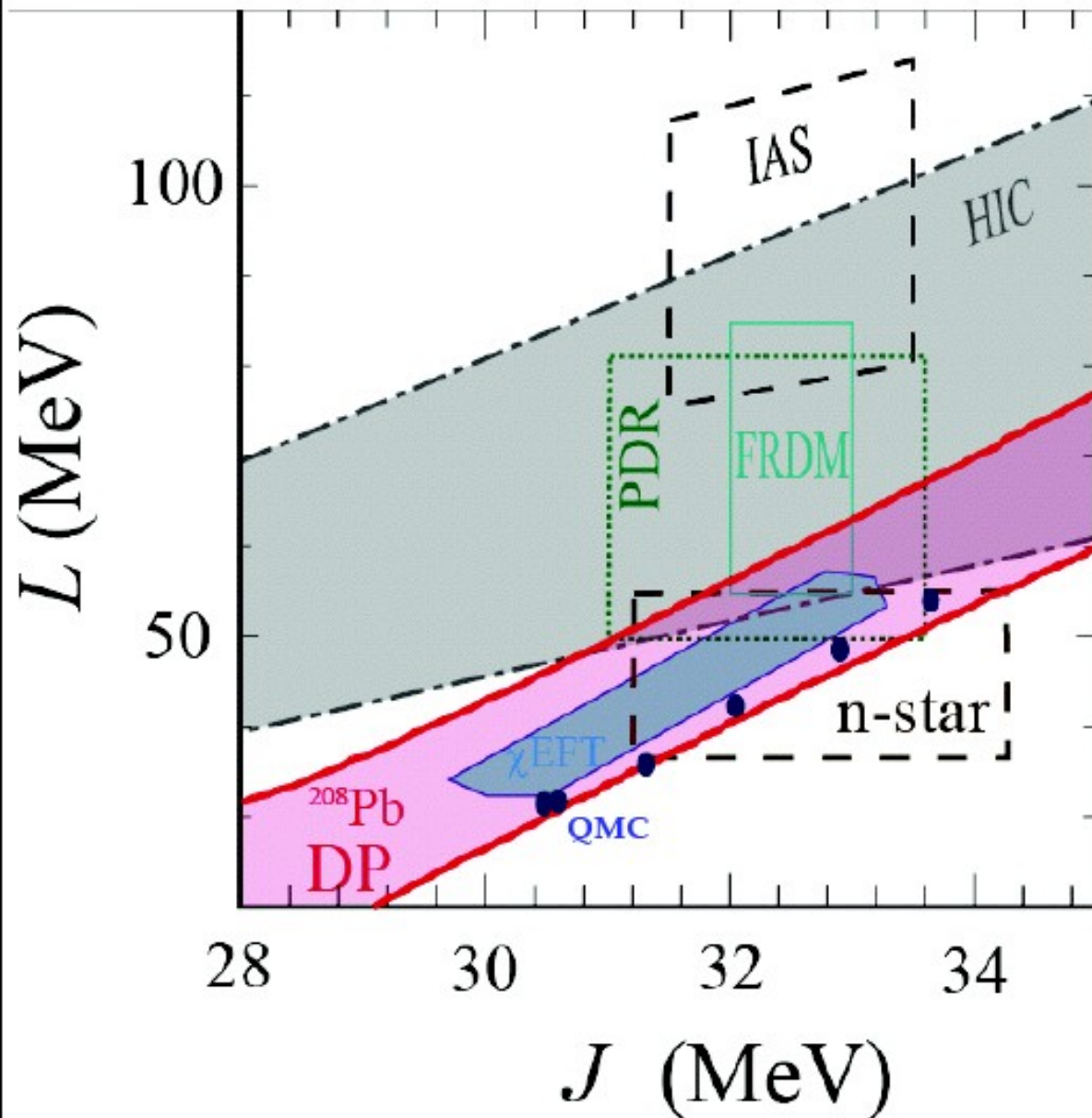
- Symmetry energy is decisive in neutron star matter EOS, and is related to various properties of nuclei.
- Experimental & Theoretical studies are in progress.
- If you have a new idea to determine E_{sym} , please propose as soon as possible.



Thank you !

Constraints on J and L

by Tamii



AT *et al.*, to be published in EPJA.

M.B. Tsang *et al.*, PRC86, 015803 (2012).

I. Tews *et al.*, PRL110, 032504 (2013)

DP: Dipole Polarizability (this work)

HIC: Heavy Ion Collision

PDR: Pygmy Dipole Resonance

IAS: Isobaric Analogue State

FRDM: Finite Range Droplet

Model (nuclear mass analysis)

n-star: Neutron Star Observation

(A.W. Steiner *et al.*,)

cEFT: Chiral Effective Field Theory

QMC: Quantum Monte-Carlo Calc.