
Symmetry Parameter Constraints from a Lower Bound on the Neutron-Matter Energy (and Its Relevance to Neutron Star Physics)

Akira Ohnishi
(Yukawa Inst. for Theor. Phys., Kyoto U.)

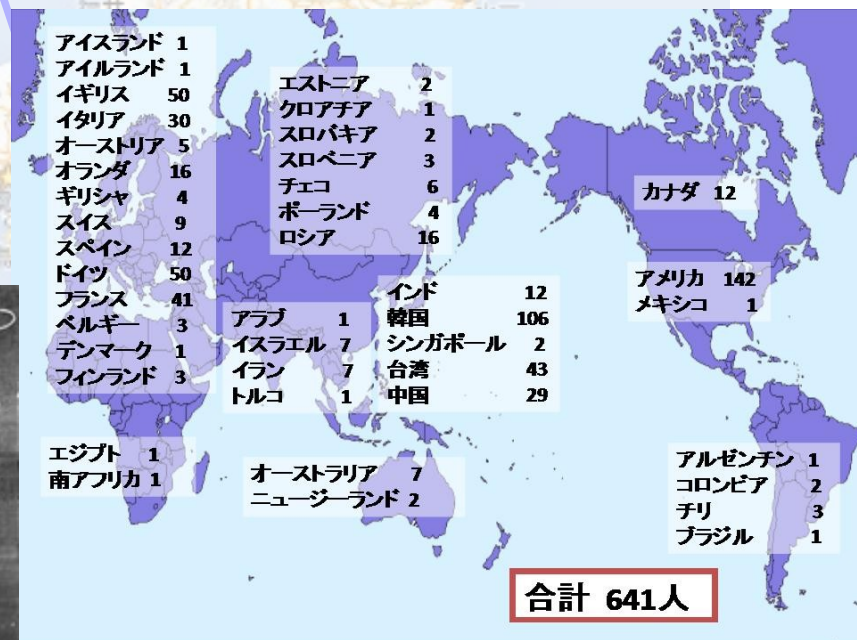
名古屋大学談話会 (2018/12/4)

- *I. Tews, J. M. Lattimer, AO, E.E.Kolomeitsev, ApJ 848('17) 105 [arXiv:1611.07133]*
- *AO, Kolomeitsev, Lattimer, Tews, X.Wu, in prog.*

Do you know Yukawa Institute ?

■ Yukawa Institute for Theoretical Physics, Kyoto University

- Founded in 1953 to memorize Yukawa's Nobel prize (first winner in Japan).
- Domestic & International Collaboration program
20-30 domestic workshops, ~ 10 international workshops,
~ 1000(?) domestic visitors, 600-700 visitors from abroad
- We will have around two long-term workshops annually.
Thank you for participating the long-term workshop NFQCD 2018.



■ Introduction

- **Gravitational Wave from Binary Neutron Star Merger**
- **Neutron Star Matter Mass-Radius Curve and EOS**

■ Nuclear Matter EOS constraints

- **Saturation and Symmetry Energy Parameters**
- **Unitary gas constraint**
- **Effects of higher-order Sym. E. parameters on Neutron Star Radius**

■ Implications to quark-hadron physics in cold dense matter

- **Neutron chemical potential, QCD phase transition**
- **EOS softening signal from heavy-ion collisions**

■ Summary

Introduction

Impact of GW from binary neutron star merger

- **GW170817 from NS-NS → Multi messenger astrophysics**

B.P.Abbott+ [LIGO Sci. and Virgo Collab.], PRL119('17)161101

- **Neutron Star Maximum Mass**

- **No prompt collapse, No GW signal from Hyper Massive NS**
→ $M_{\max}(T=0, \omega=0) < M_{\max}(T=0, \omega) < M < M_{\max}(T, \omega)$

- **Neutron Star Radius**

- **Inspiral region → Tidal deformability (Λ) → NS radius (e.g. R1.4)**

- **Nucleosynthesis site of r-process nuclei**

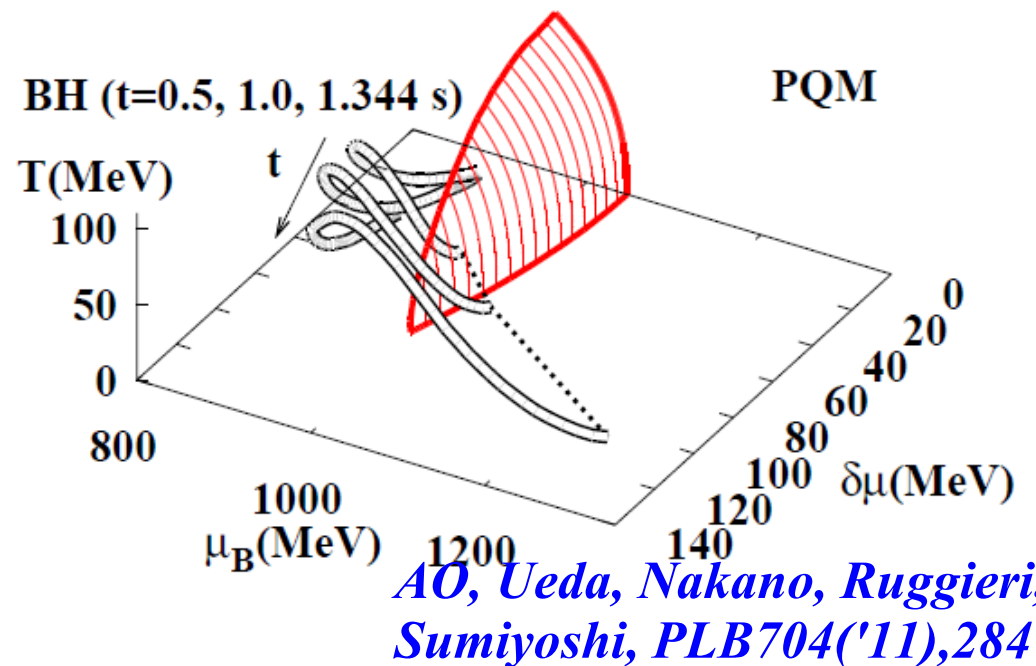
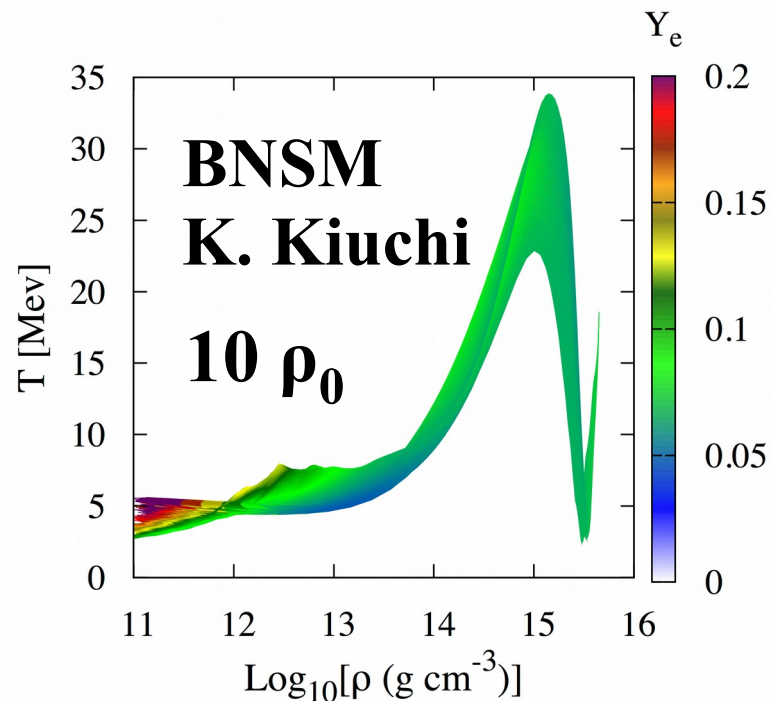
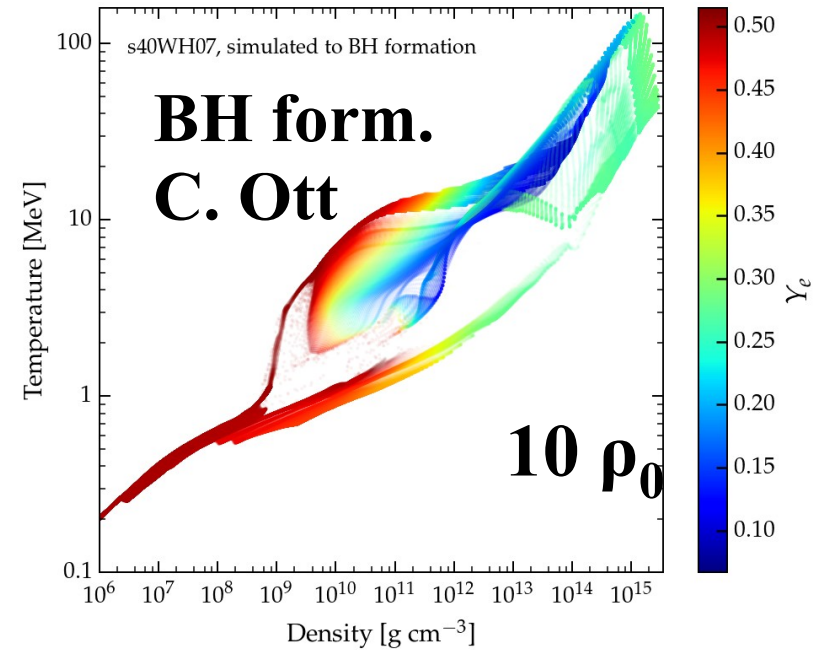
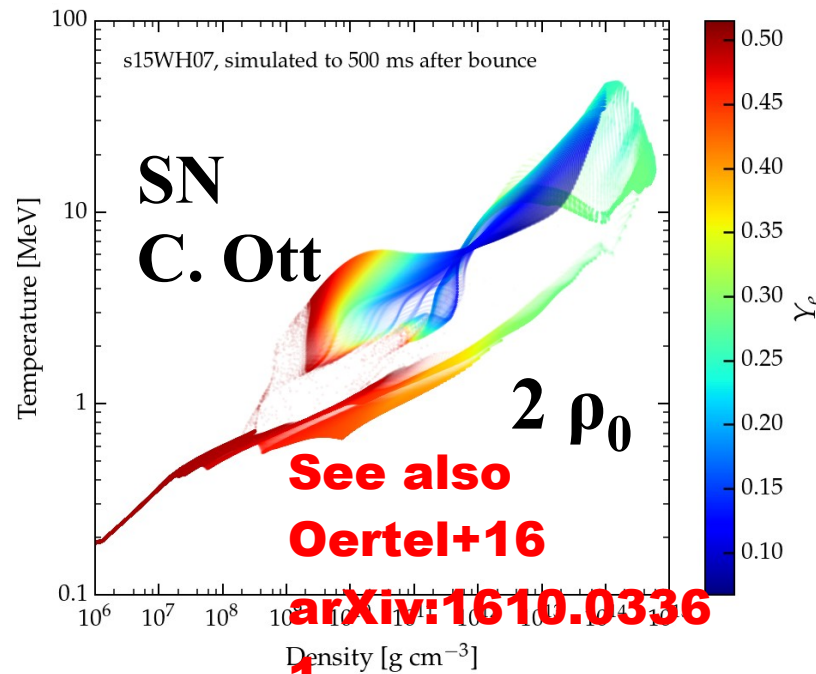
- **kilonova/macronova from decay energy of the synthesized elements**
- **r-process nucleosynthesis seems to occur in BNSM !**

- **Central Engine of (Short) Gamma-Ray Bursts**

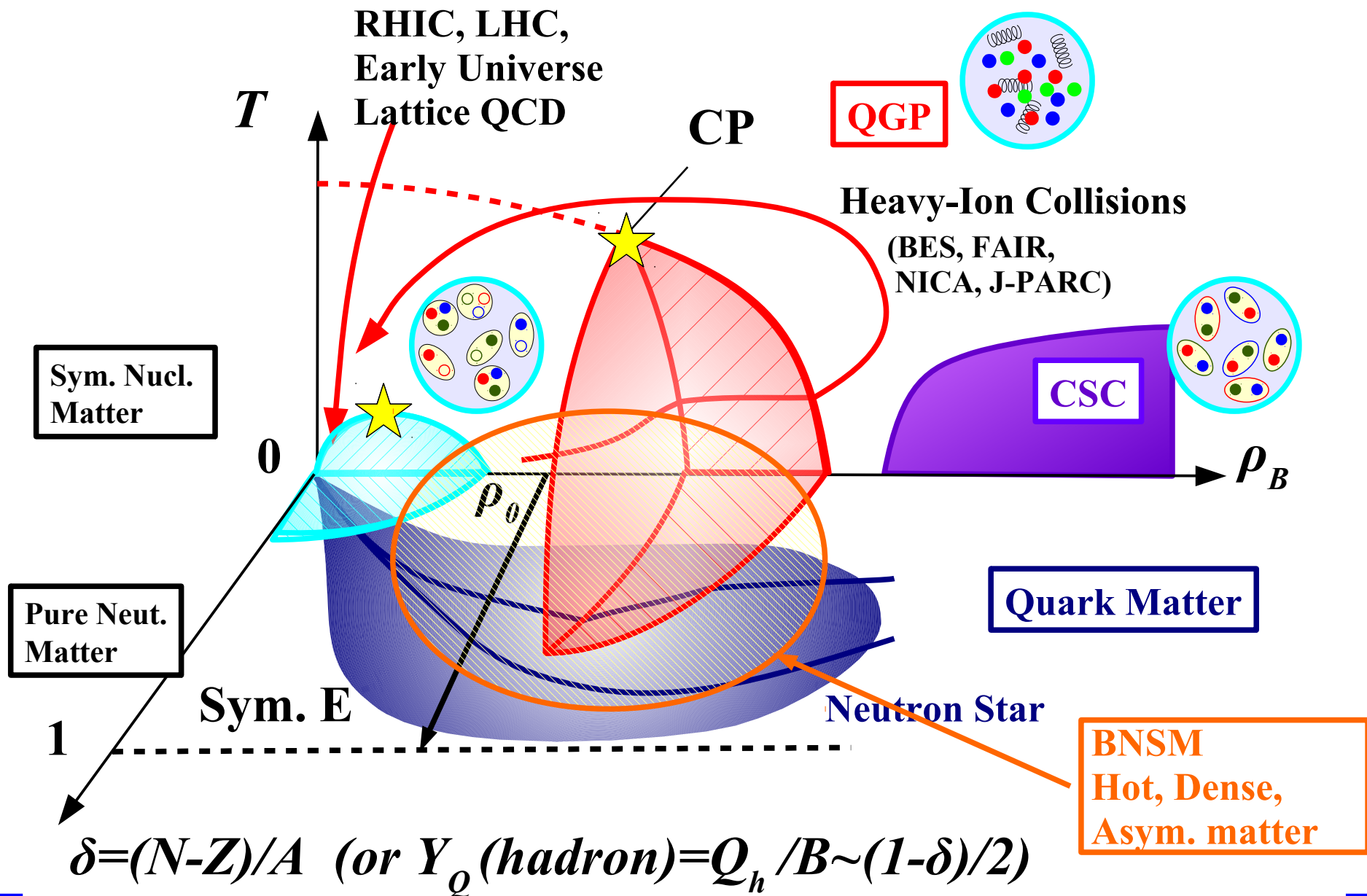
- **GW as standard siren (Hubble constant)**

Courtesy of Y. Sekiguchi @ YKIS2018b

(ρ, T, Y_e) during SN, BH formation, BNSM



QCD Phase Diagram



AO, JPS Conf. Proc. 20 (2018), 011035

A. Ohnishi @ Nagoya U., Dec.4-6, 2018

連星中性子星合体

■ GW170817

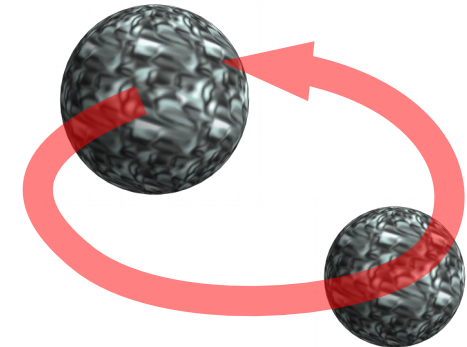
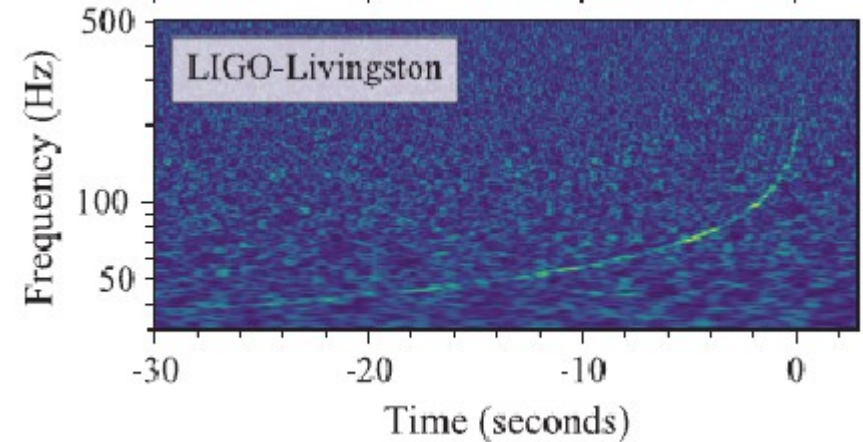
B. P. Abbott et al. (LIGO and Virgo)
PRL 119, 161101 (2017)

- 質量和 $M = 2.74^{+0.04}_{-0.01} M_{\odot}$
- それぞれの質量 $1.17\text{-}1.60 M_{\odot}$
→ 連星中性子星合体
(Binary Neutron Star Merger)
- Gamma Ray Burst (GRB170717A) が
1.7 s 後に起こる。
→ GRB の起源 (の一つ) を特定

- inspiral (徐々に近づいていく段階) における振動数変化を観測
→ 中性子星半径を制限

- 放出された物質の速度から中性子星の最大質量を制限
M. Shibata et al., 1710.07579

$$M_{\text{max}} = (2.15\text{-}2.25) M_{\odot}$$



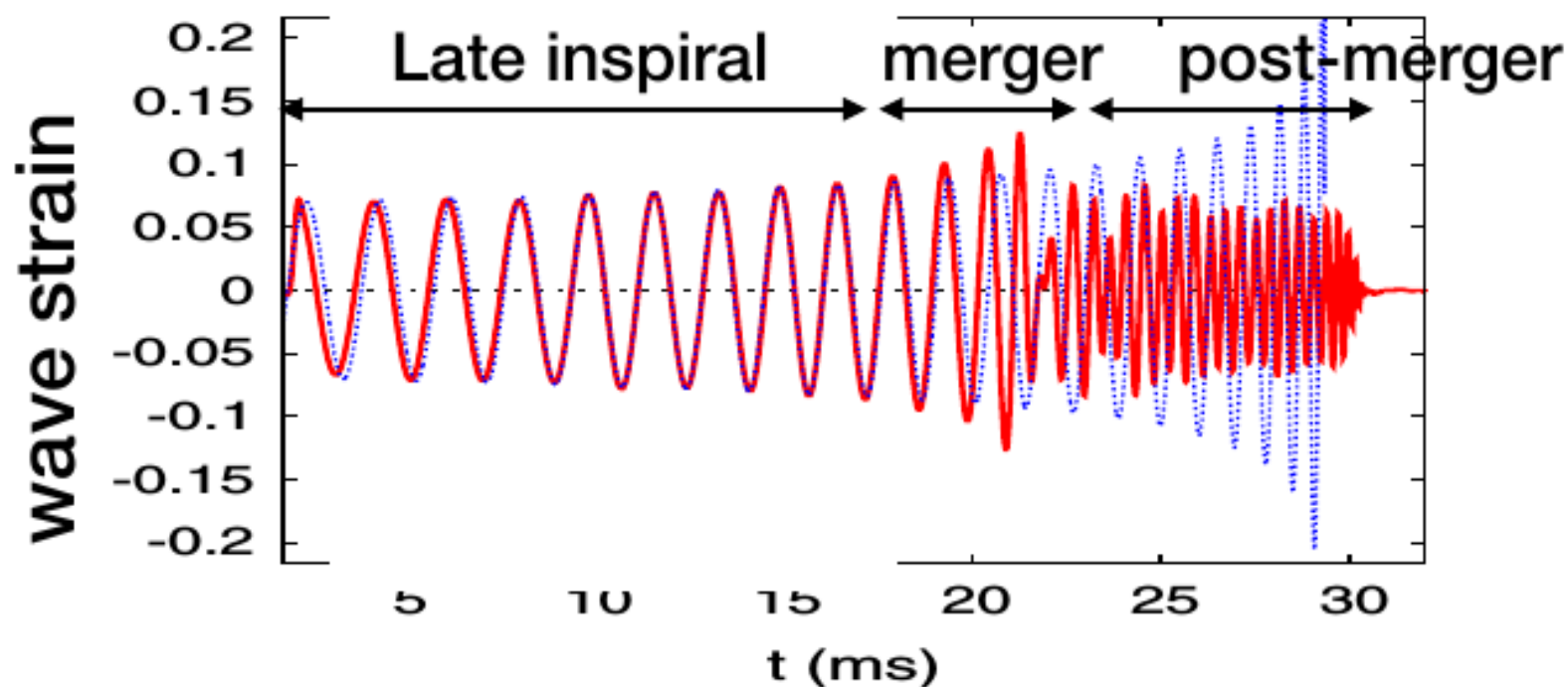
Gravitational Wave in the Inspiral Phase

■ 重力波の波形

- Inspiral, merger, post-merger (hypermassive neutron star)

■ Inspiral phase での重力波の波形

- 質点からの重力波にくらべて、位相が進む
→ Tidal Deformation (潮汐変形) 効果



Hotokezaka

Tidal Deformability (潮汐変形度)

- Tidal deformability λ (無次元化 Λ)
= 潮汐力 E_{ij} に対する変形 Q_{ij} の起こりやすさ
$$Q_{ij} = -\lambda E_{ij}, \quad \lambda = \frac{\Lambda}{G} (GM/c^2)^5 = \frac{\Lambda}{G} (CR)^5$$
$$C = GM/c^2 R \text{ (Compactness)}$$

Inspiral phase での位相の進みは
tidal deformability によって
ほぼ決まる。

- 大きな半径 (硬い EOS)
→ 変形しやすい → 大きな Λ

T.Hinderer, B.D.Lackey, R.N.Lang, J.S.Read,
PRD81 ('10)123016

(A) B.D. Lackey, L.Wade, PRC91('15)043002

(R-A) E. Annala+, PRL120('18)172703

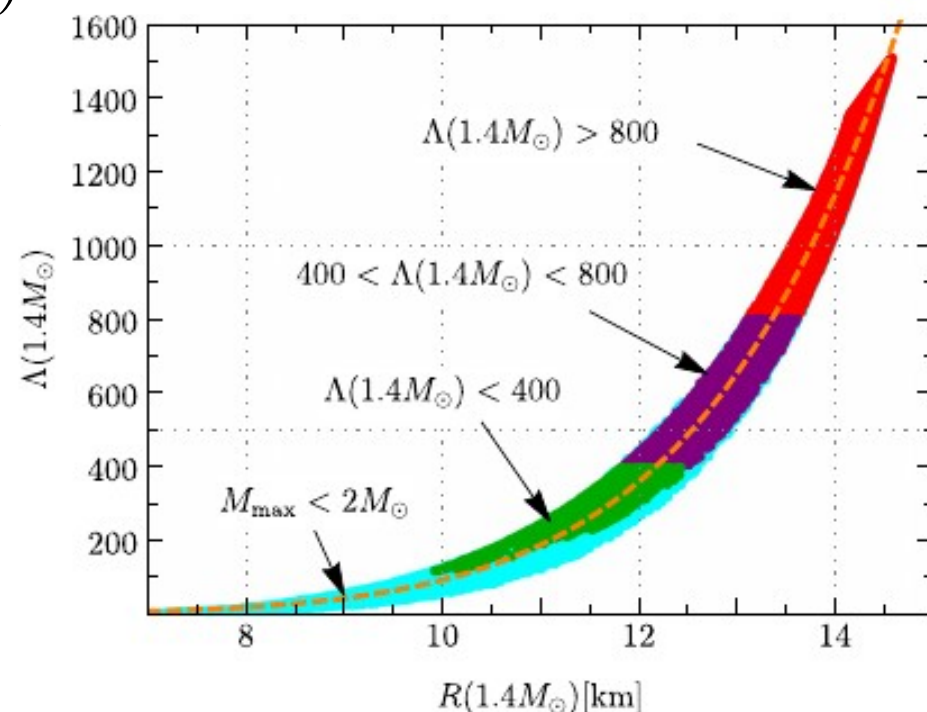
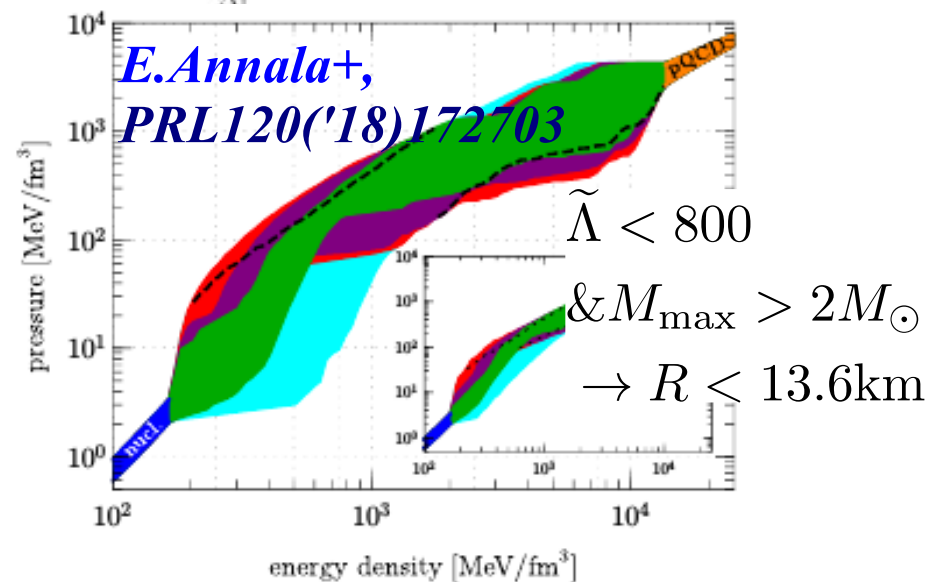
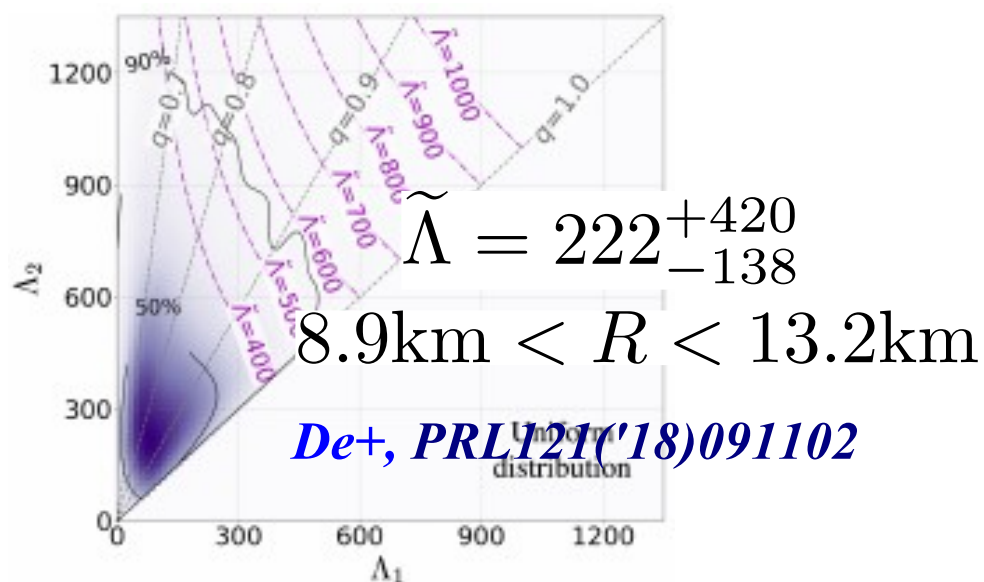
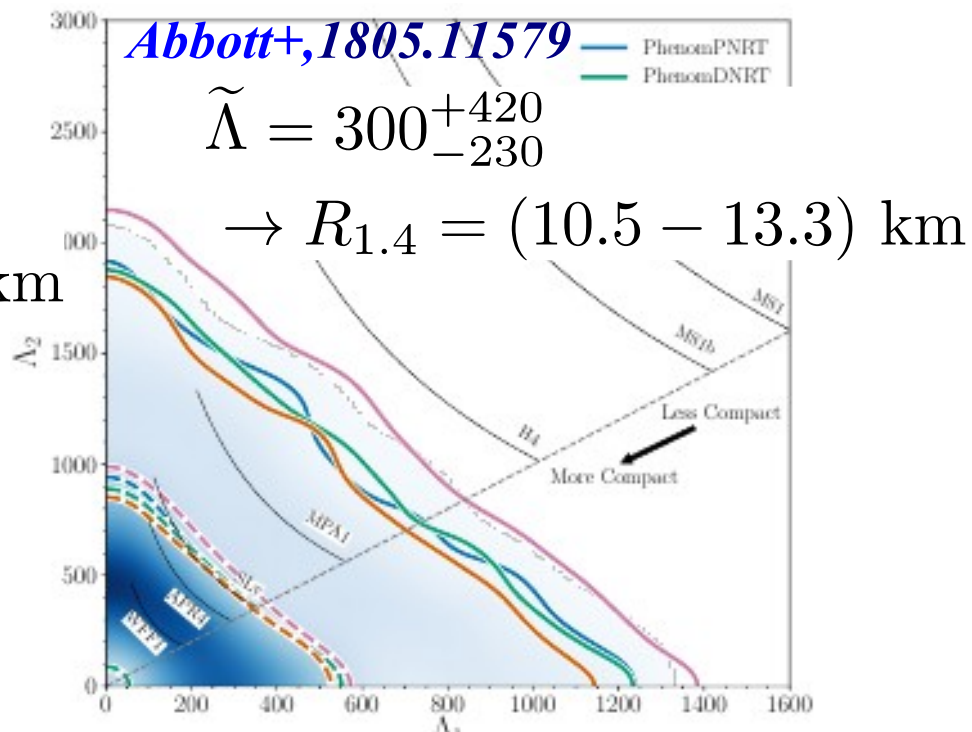
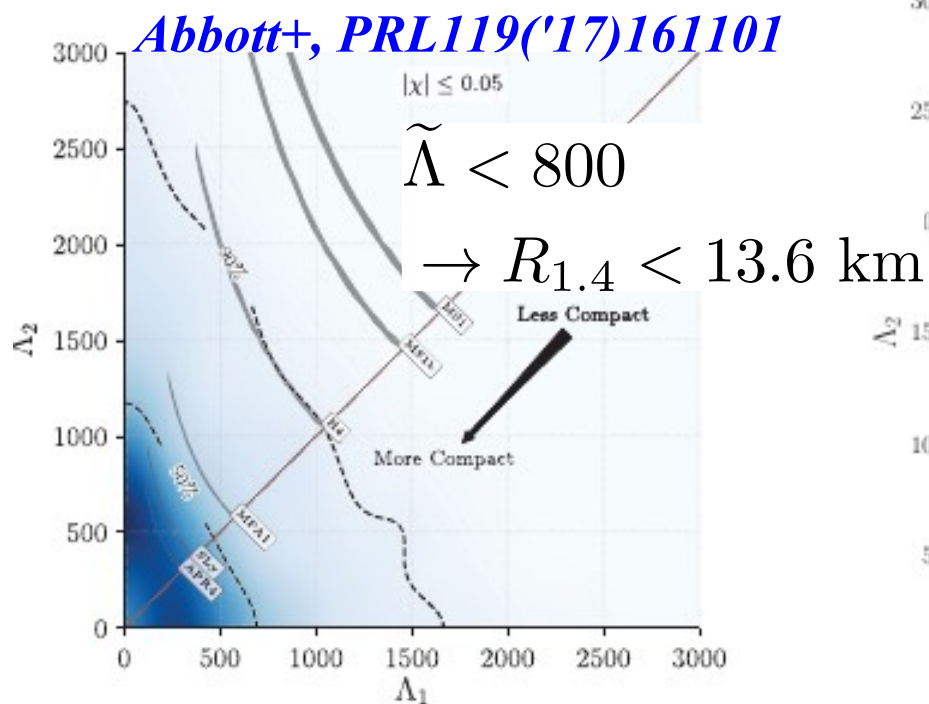
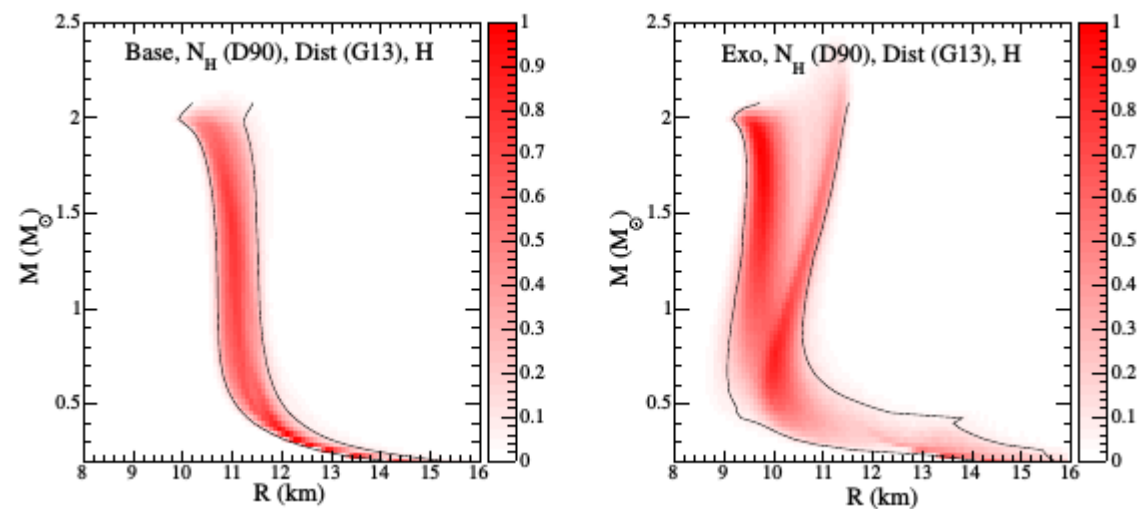
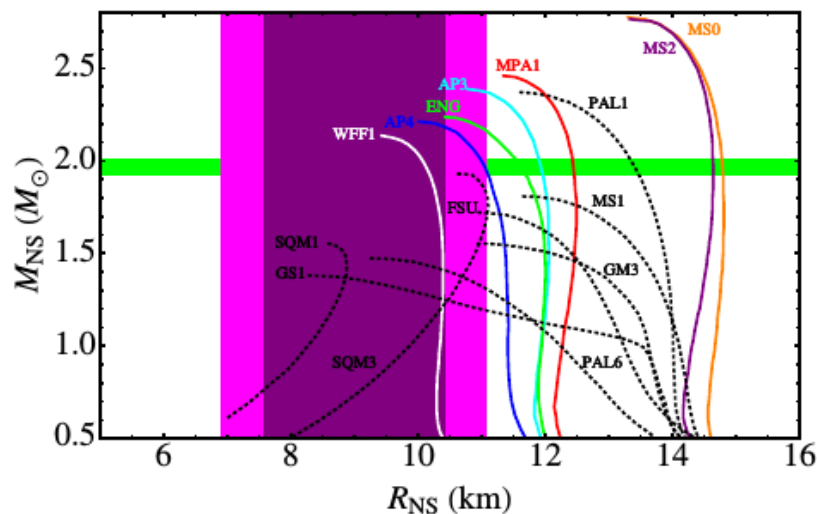
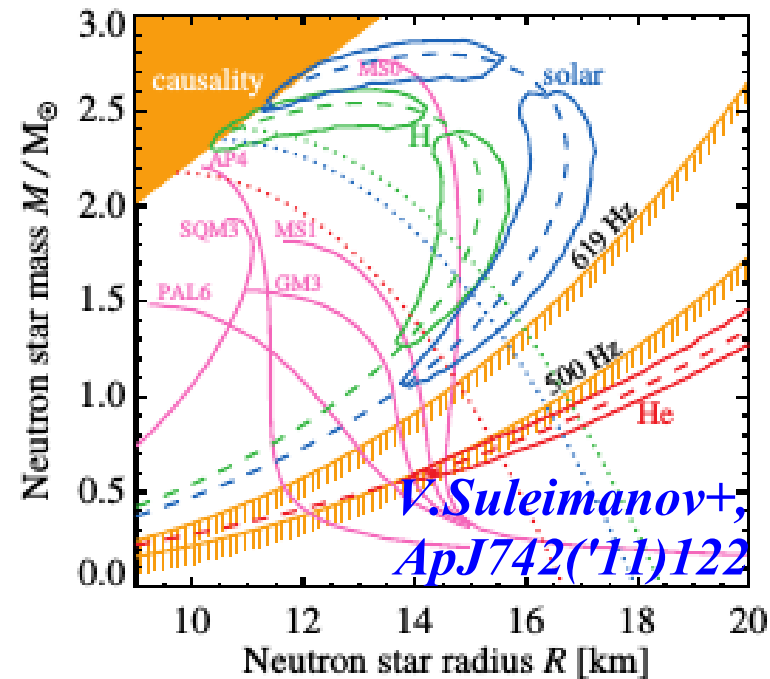
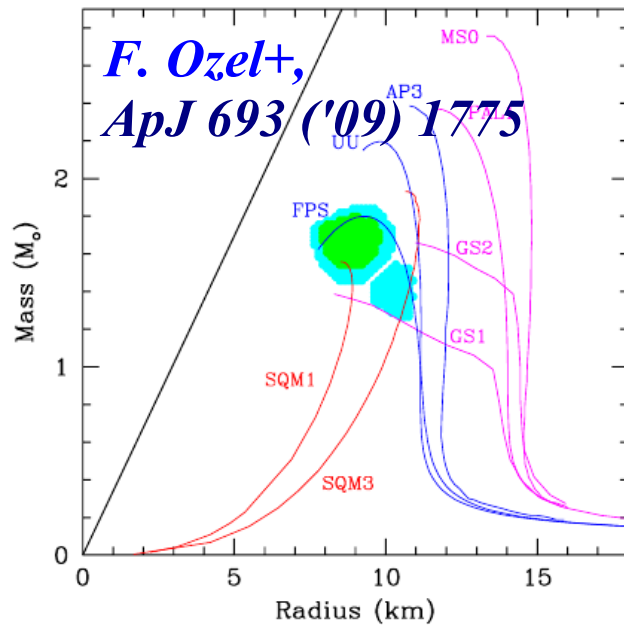


FIG. 2. The Λ values for stars with $M = 1.4 M_\odot$ as functions of the corresponding radius. The color coding follows Fig. 1, while the orange dashed line $\Lambda = 2.88 \times 10^{-6} (R/\text{km})^{7.5}$ has been included just to guide the eye.

Constraints on EOS from GW170817



MR curve from X-ray burst



Neutron Star Radius from Astronomical Observations

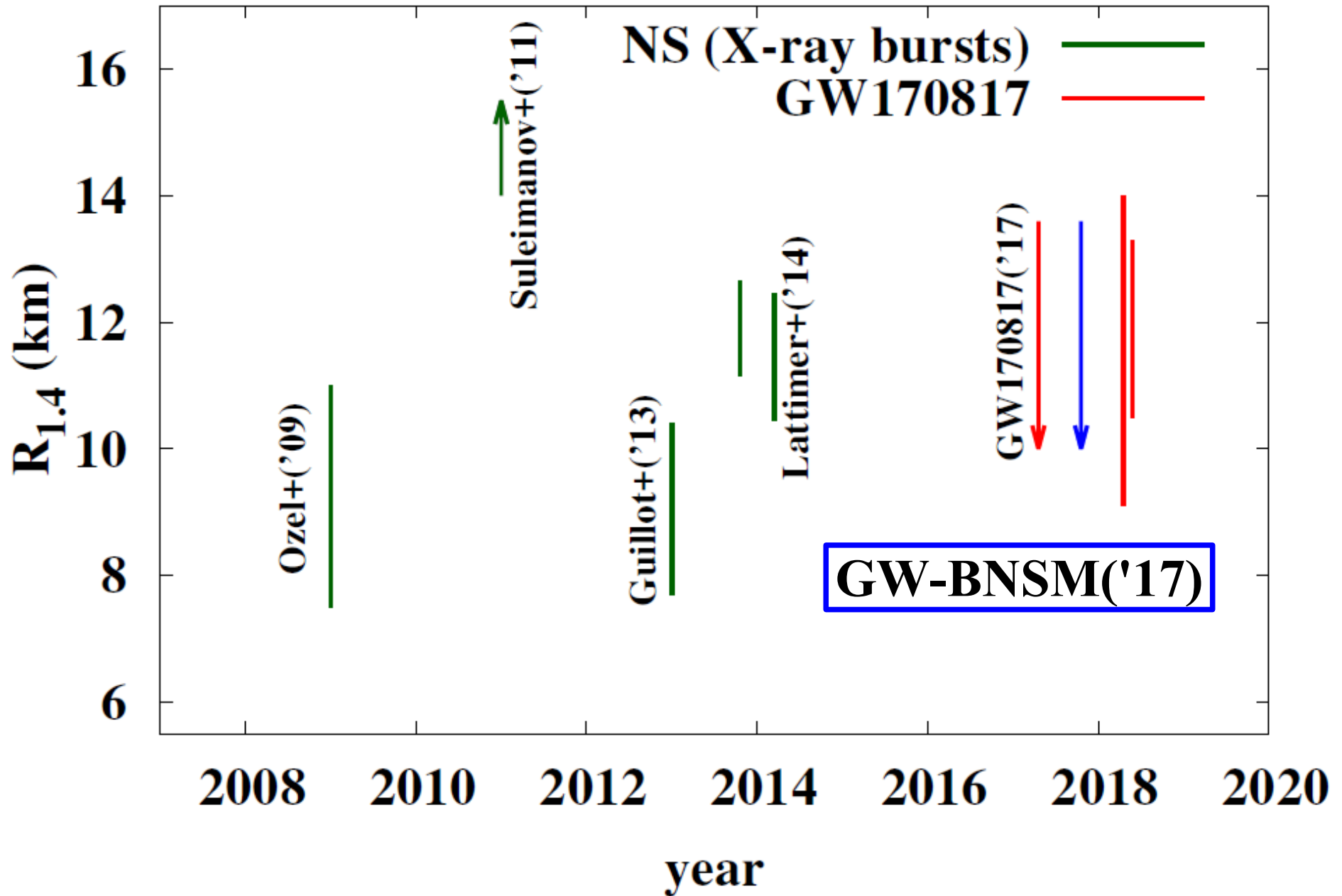
■ From X-ray bursts

- $R=9$ or 11 km (7.5-11 km) *F. Özel+, ApJ693('09)1775*
[touch down = Eddington limit?]
- $R > 14$ km *V.Suleimanov+, ApJ742('11)122*
[color correction factor ?]
- $R=9.1^{+1.3}_{-1.4}$ km *S.Guillot+, ApJ772('13)7*
[Common R ?; Denied later by the authors in ApJ 796('14)L3 (1409.4306)]
- $R=10.4-12.7$ km *J.M.Lattimer, A.W.Steiner, ApJ784('14)123*
($R=(11.15-12.66)$ km (normal EOS), $R=(10.45-12.45)$ km (Exo EOS))

■ Gravitational Waves

- $\bar{\Lambda} < 800$ *B.P.Abbott+, PRL119('17)161101* $\rightarrow R < 13.6$ km
- $\bar{\Lambda}=300^{+420}_{-230} \rightarrow R=(10.5-13.3)$ km *B.P.Abbott+ (LIGO-Virgo), 1805.11579*
- $\bar{\Lambda}=222^{+420}_{-138} \rightarrow R=(9.1-14.0)$ km *S.De+, PRL121('18)091102*
- $\bar{\Lambda} < 800$ & $M_{\max} > 2 M_{\odot} \rightarrow R < 13.6$ km *E.Annala+, PRL120('18)172703*

Time dependence of Neutron Star Radius ($R_{1.4}$)

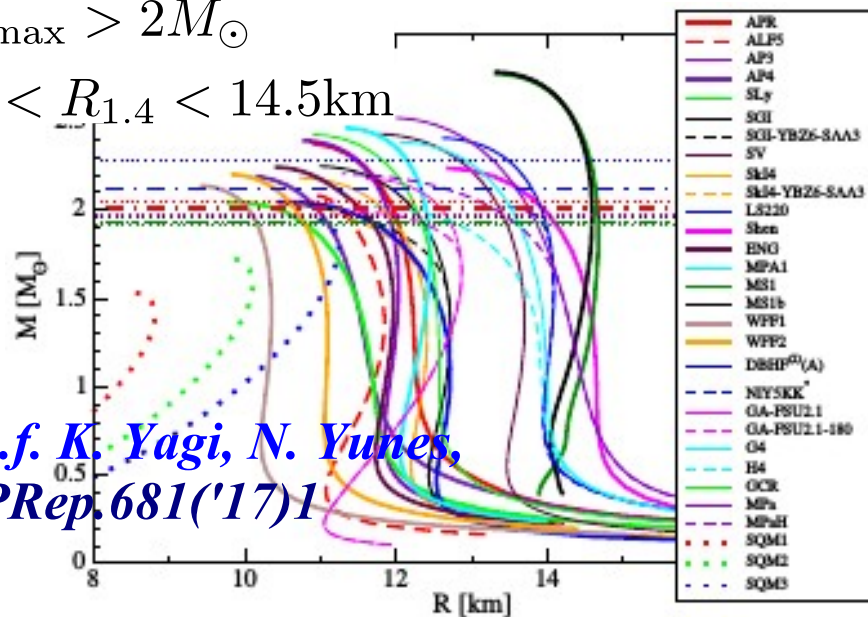


Constraints from Nuclear Physics (+ α)

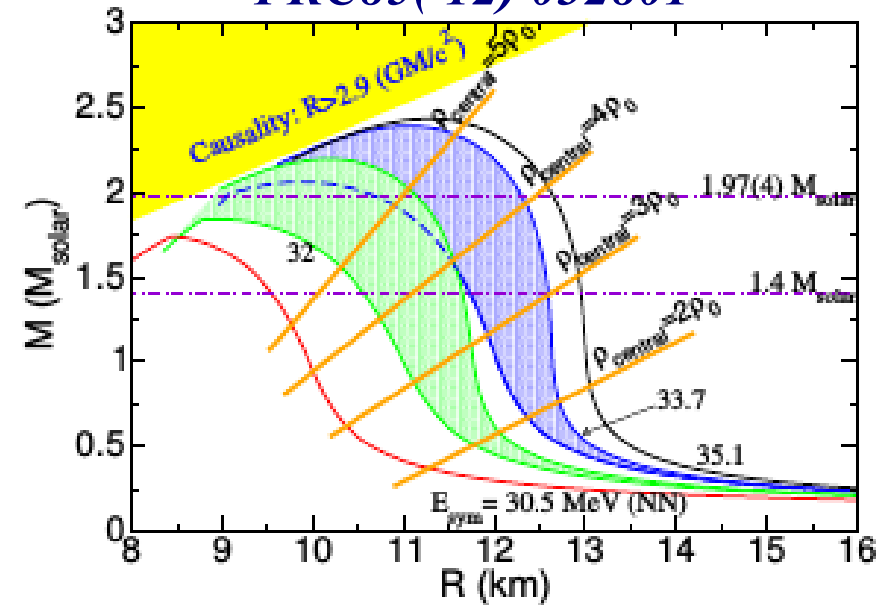
$$M_{\max} > 2M_{\odot}$$

$$10 < R_{1.4} < 14.5 \text{ km}$$

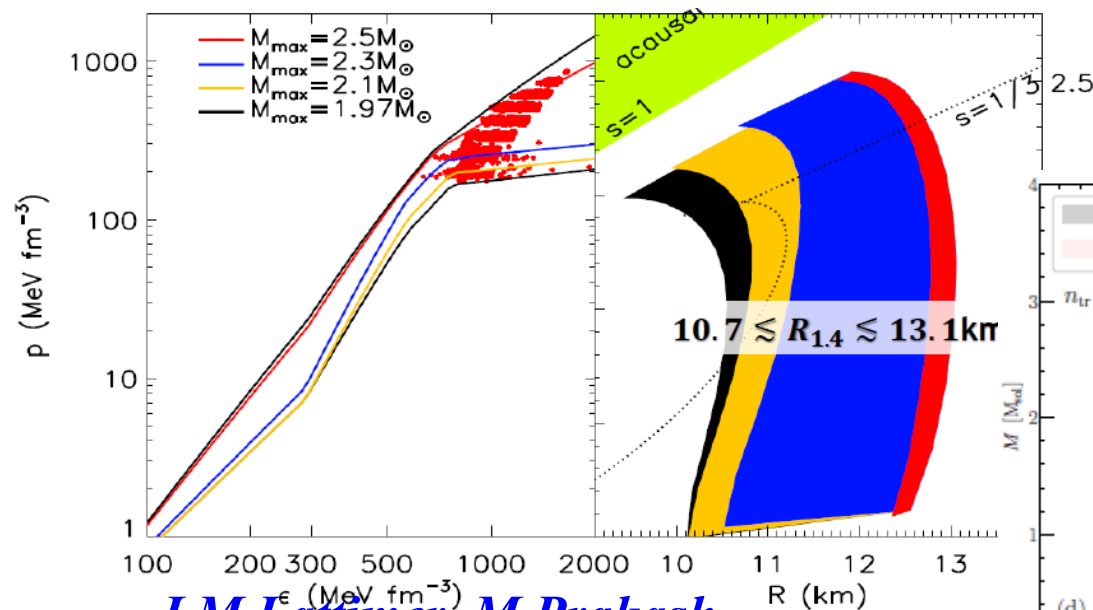
*c.f. K. Yagi, N. Yunes,
PRep.681('17)1*



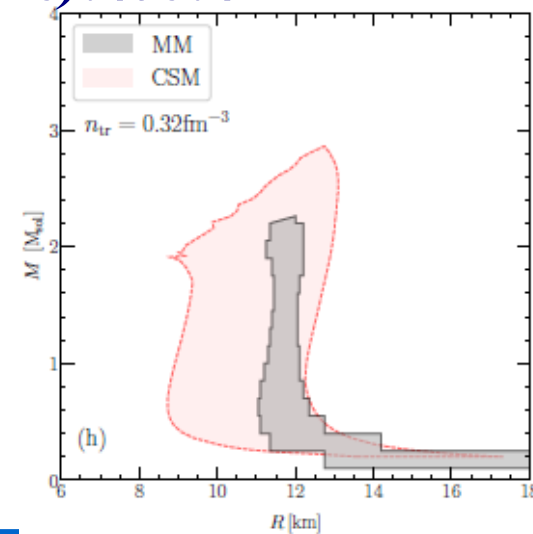
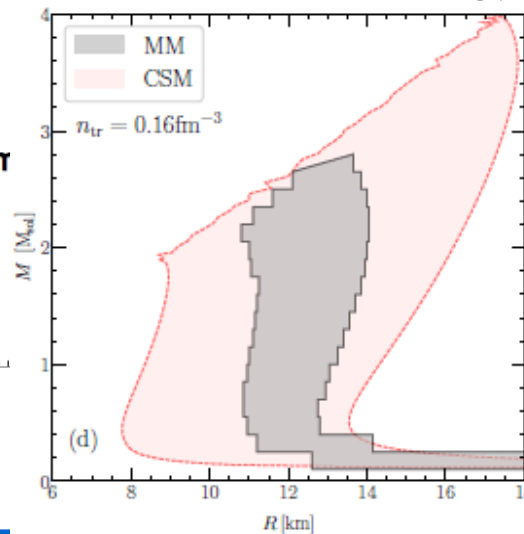
*S. Gandolfi, J. Carlson, S. Reddy,
PRC85('12) 032801*



*I. Tews, J. Margueron, S. Reddy,
PRC98 ('18)045804*



*J.M. Lattimer, M. Prakash,
PRep.621('16)127*



A. Ohnishi @ Nagoya U., Dec.4-6, 2018

Neutron Star Radius from Nuclear Physics (+ α)

■ Impact of $2 M_{\odot}$ neutron star

P.Demorest+, Nature 467('10)1081 (1010.5788),

J.Antoniadis+, Science 340('13)6131(1304.6875)

- $R=(6-15)$ km (before 2010) $\rightarrow R=(10-15)$ km (with $2 M_{\odot}$)

■ Impact of symmetry energy parameters

- $S_0 \rightarrow R=(10.5-13)$ km *S.Gandolfi, J.Carlson, S.Reddy, PRC85('12)032801*

- $(S_0, L) \rightarrow R=(10.7-13.1)$ km *J.M.Lattimer, M.Prakash, PRep.621('16)127*

■ Chiral Effective Field Theory (Chiral EFT)

- $\chi\text{EFT}+\text{pQCD}+\text{GW} \rightarrow R=(10.0-13.6)$ km

E.Annala, T.Gorda, A.Kurkela, A.Vuorinen, PRL120('18)172703

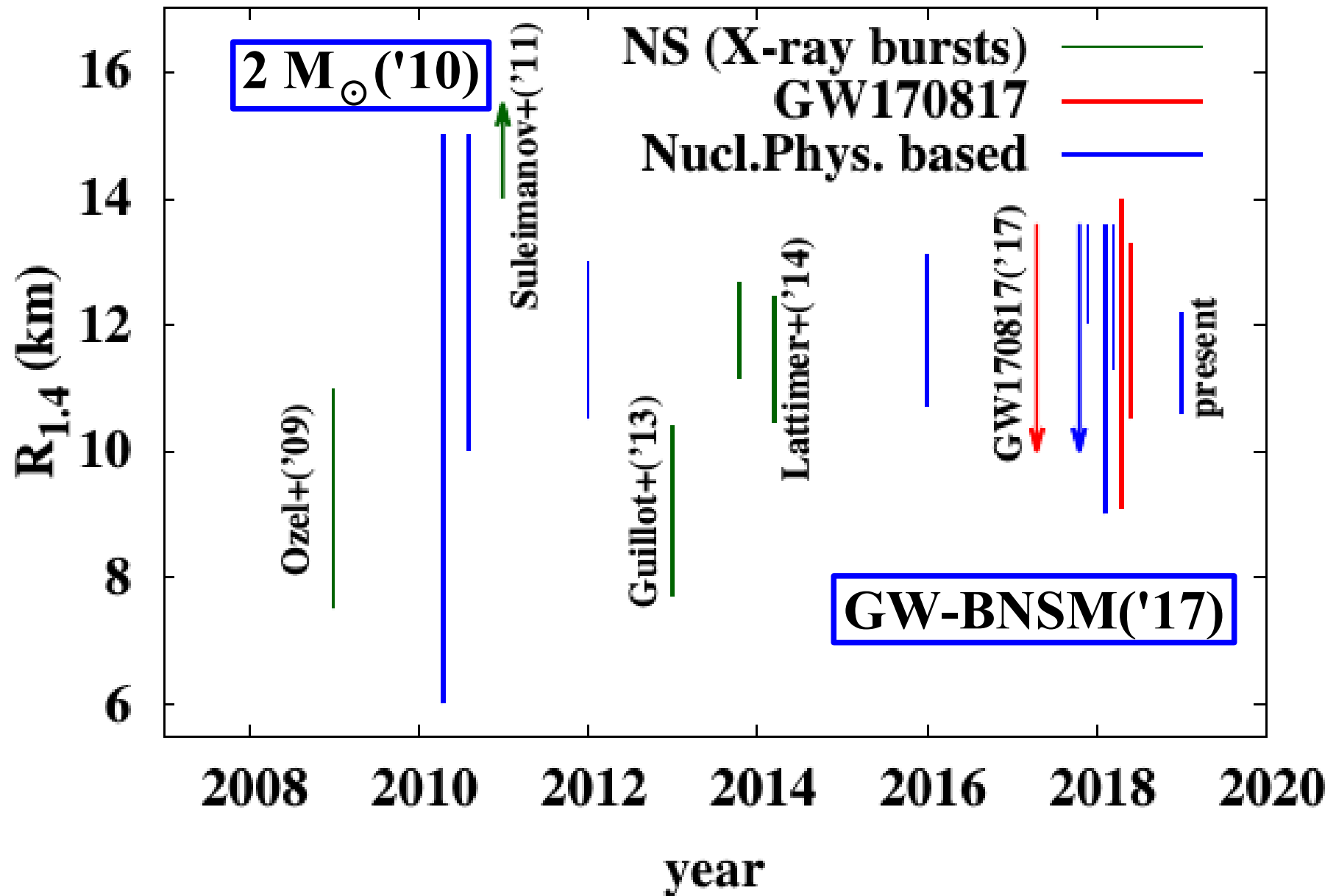
- $\chi\text{EFT}+c_s^2 \rightarrow R=(9.0-13.6)$ km [min. model $R=(11.3-13.6)$ km]

I. Tews, J. Margueron, S. Reddy, PRC98 ('18)045804

- Neutron skin thickness from ν scatt. (PREX) $\rightarrow R=(12.0-13.6)$ km

F.J. Fattoyev, J. Piekarewicz, C.J. Horowitz, PRL120 ('18)172702

Time dependence of Neutron Star Radius ($R_{1.4}$)



*Symmetry Energy Parameters (S_0 , L) affect
Neutron Star Radius.*

How can we constrain (S_0 , L) ?

How about higher-order parameters ?

Unitary Gas Constraint on Symmetry Energy Parameters

現実的な中性子物質 EOS 構築へ向けて

■ 中性子過剰物質物質のキーワード

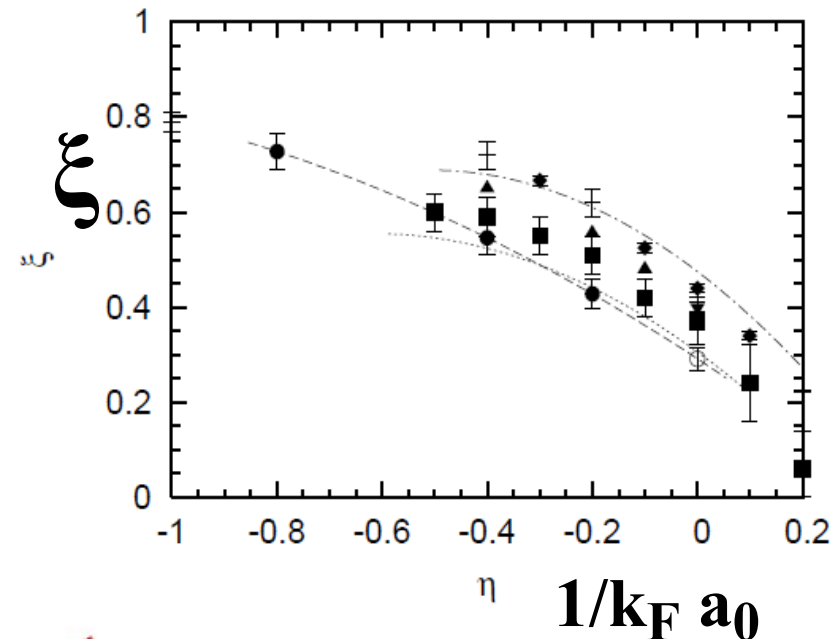
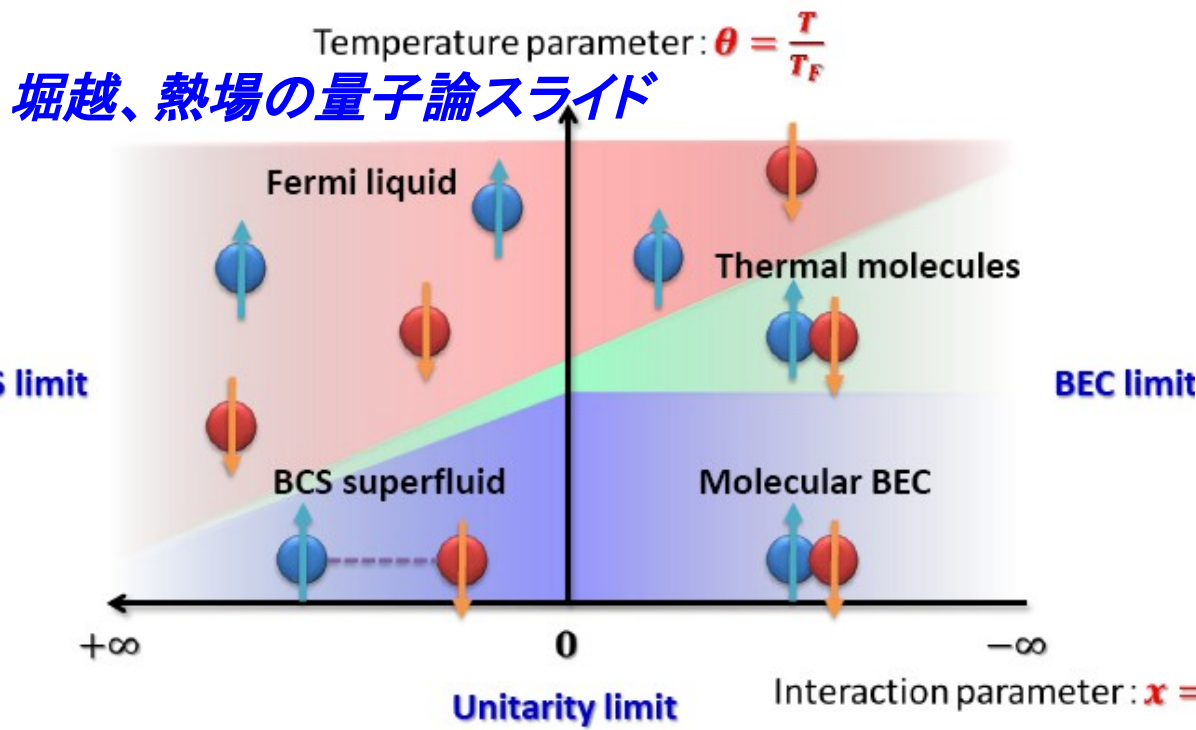
= 対称エネルギー・微視的相互作用・普遍性

● 低密度純中性子物質：第一原理計算が可能 (e.g. Abe, Seki ('09); Takano)

● 大きな散乱長 ($a_0 \sim -16$ fm) $\rightarrow$ 中性子物質 $\sim$ ユニタリーガス

$$E^{\text{Unitary}} = \xi E^{\text{Free}} \quad \xi \approx 0.4 \text{ (Bertsch parameter)}$$

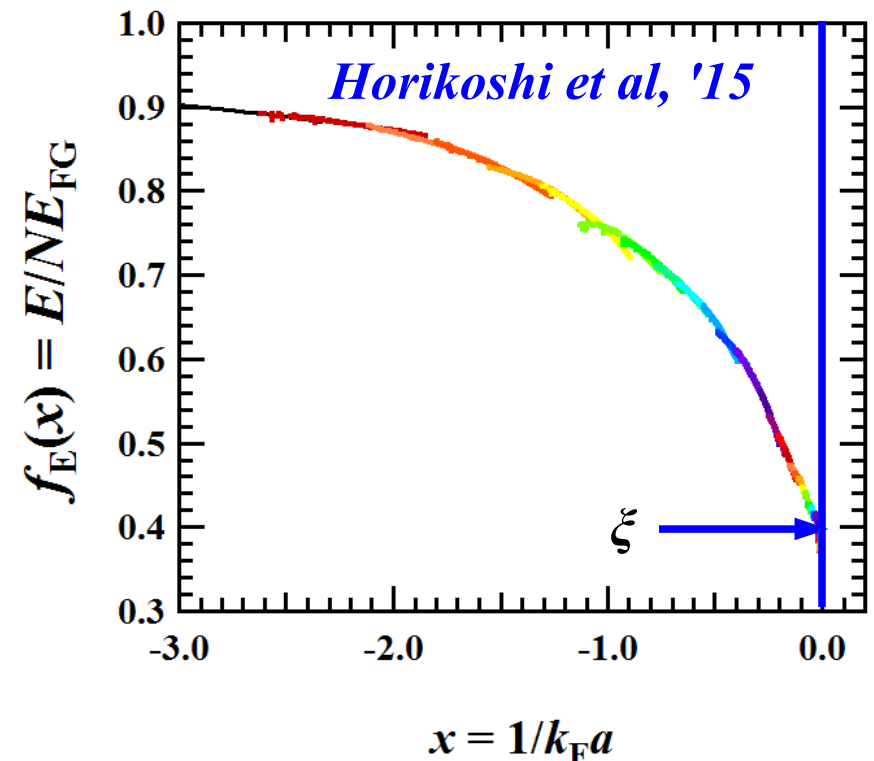
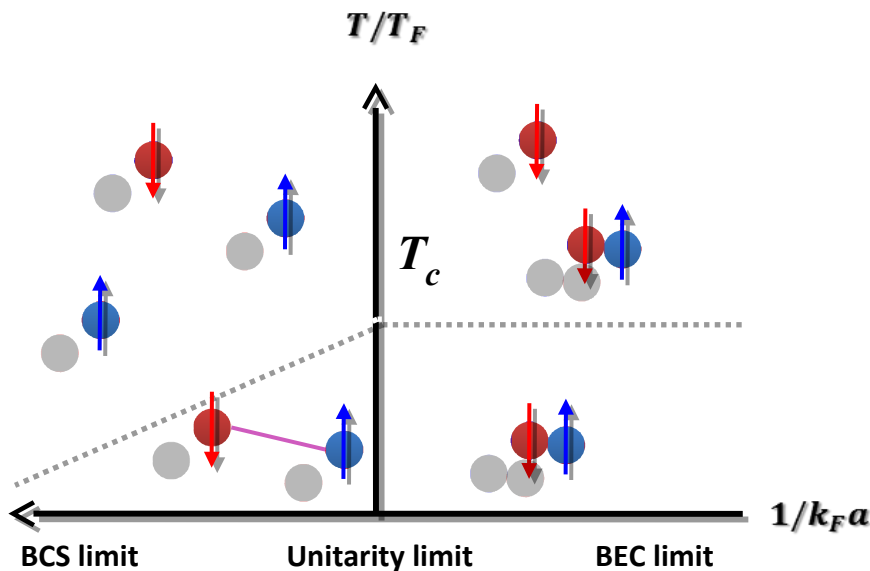
$\rightarrow$ 理論計算手法のテスト、*finite* a_0 , r_{eff} を含む普遍的関数の探索



Abe, Seki ('09)

Cold Atom EOS

- $E = \xi E_{\text{FG}}$ ($\xi \sim 0.4$) in the unitary limit ($1/k_F a_0 \rightarrow 0$, $k_F r_{\text{eff}} = 0$)
 - ξ : Bertsch parameter (INT workshop)
- EOS measurement off unitary limit
Horikoshi et al., 2016
- How can we use the cold atom EOS in the context of NS matter ?



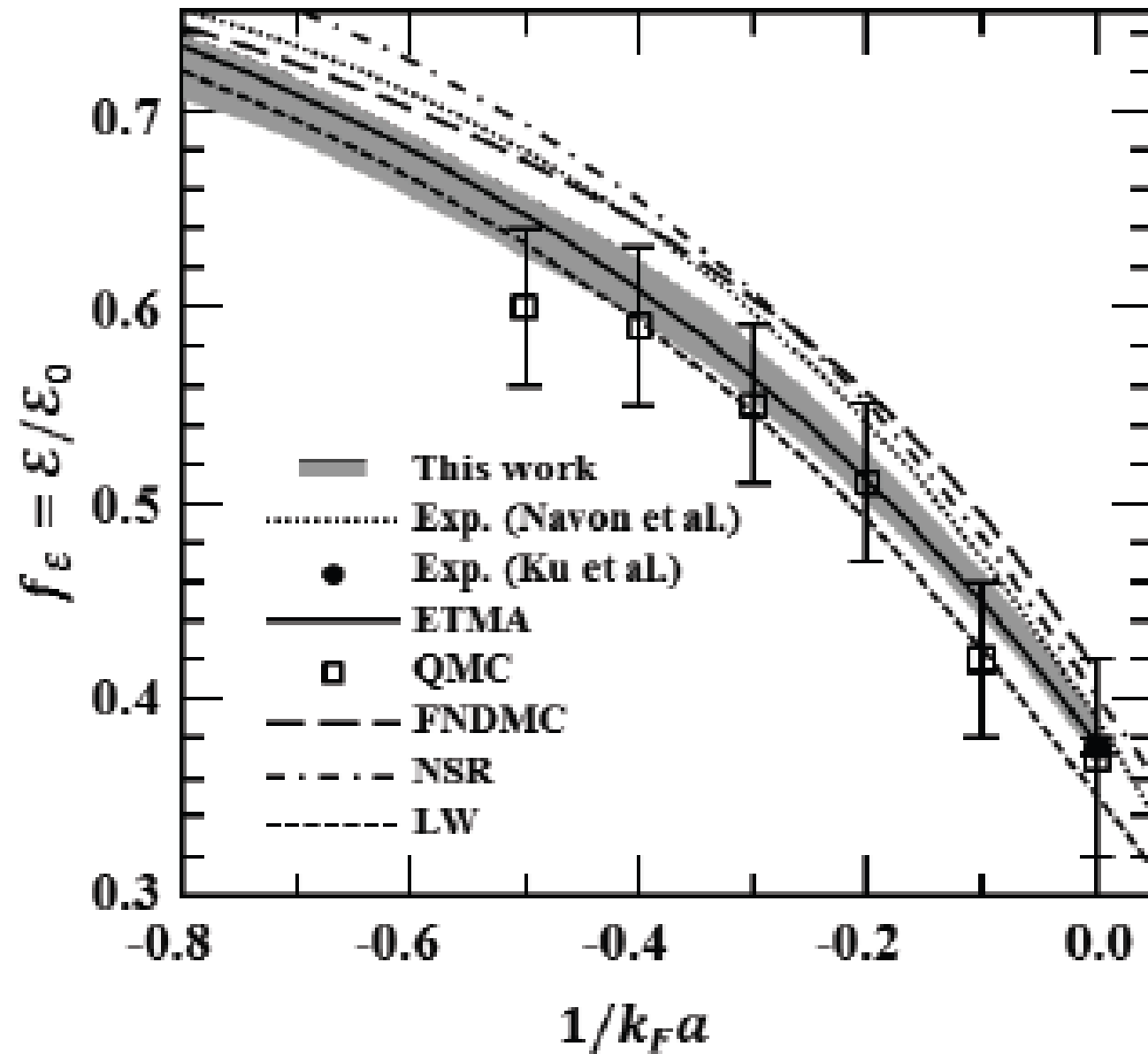
Comparison with Pure Neutron Matter EOS

表 1: 冷却原子と希薄中性子物質

	冷却原子 (${}^6\text{Li}$)	希薄中性子物質
		
粒子系	スピン 1/2 フェルミ粒子	スピン 1/2 フェルミ粒子
質量 (mc^2)	5600MeV	940MeV
力の種類	電磁気力	核力
ポテンシャルサイズ (r_0)	$\sim 2 \text{ nm}$	$\sim 2 \text{ fm}$
平均粒子間距離 ($d = n^{-1/3}$)	$\sim 250 \text{ nm}$	$> 2.3 \text{ fm}$ ($n = n_0/2$)
温度 (T)	$\sim 100 \text{ nK}$	$10^5 \sim 10^7 \text{ K}$
熱的波長 ($\lambda_T = \frac{\hbar}{\sqrt{2\pi m k_B T}}$)	$\sim 350 \text{ nm}$	$100 \sim 900 \text{ fm}$
散乱長 (a)	$-\infty \sim +\infty$	-18.5 fm
有効長 (r_e)	4.7 nm	2.75 fm
ポテンシャル形状	短距離型	短距離型
温度パラメータ (T/T_F)	< 0.1	~ 0
相互作用パラメータ ($1/k_F a$)	$-\infty \sim +\infty$	$-\infty \sim -0.04$ ($0 \sim n_0/2$)
有効長パラメータ ($k_F r_e$)	~ 0.05	< 3.7

※飽和核密度: $n_0 = 0.16 \text{ fm}^{-3}$

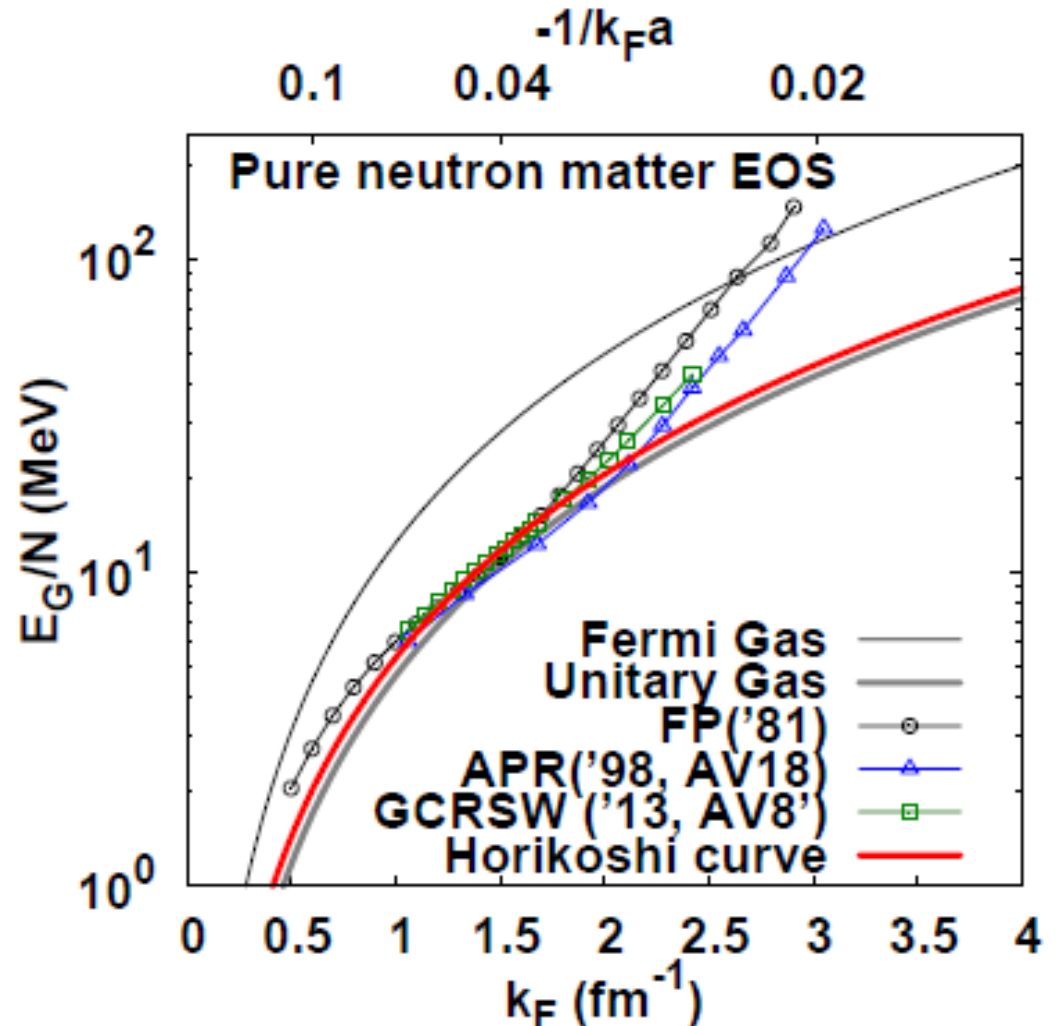
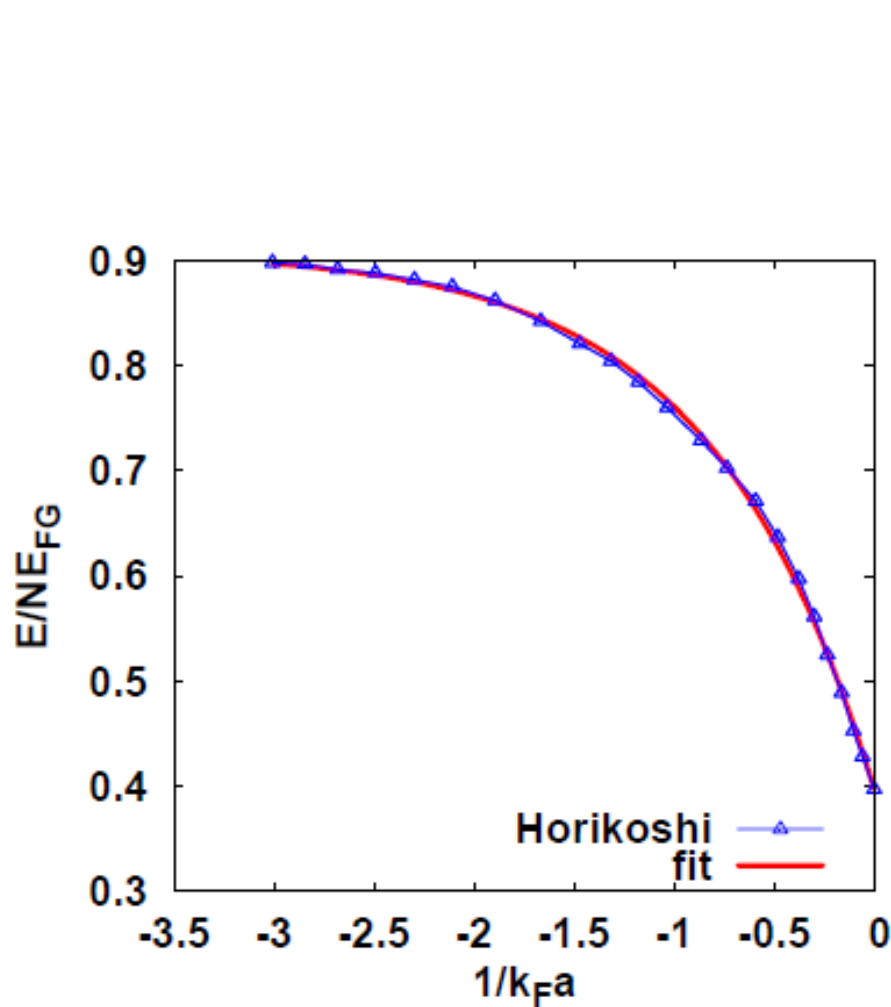
大橋 *et al.* ('16)[原子核研究]



大橋 *et al.* ('16)[原子核研究]

Comparison with Pure Neutron Matter EOS

- Cold Atom EOS is consistent with MC result.
- How can we evaluate the effective range correction ?
- Does it have other impacts in neutron star physics ?



Low Density Neutron Star Matter EOS

■ Nuclei & electron ($\rho < 4 \times 10^{11}$ g/cc)

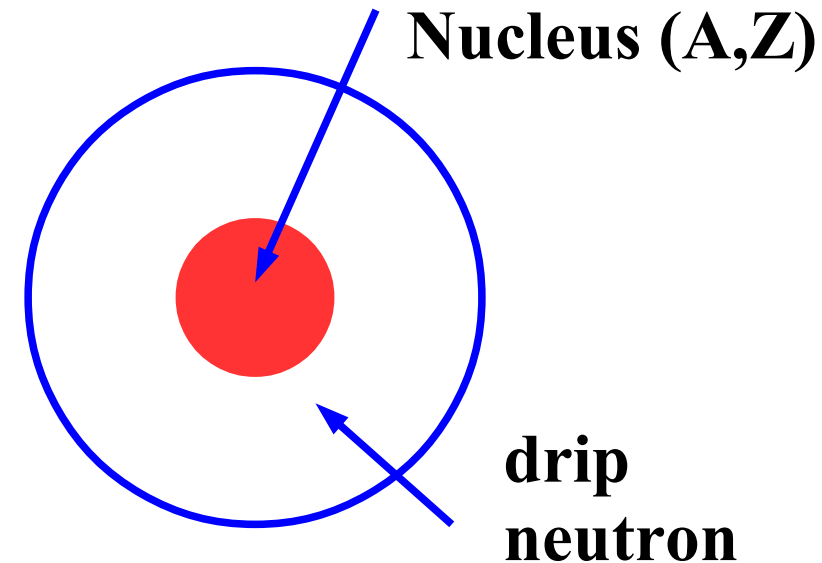
■ Nuclei, drip neutron, electron
(4×10^{11} g/cc $< \rho < 2.5 \times 10^{14}$ g/cc)
→ pure neutron matter EOS

■ Uniform nuclear matter
($\rho > 2.5 \times 10^{14}$ g/cc)
→ nuclear matter EOS at $Z \ll N$

■ “Standard” (low density) NS matter EOS

G. Baym, H.A. Bethe, C. J. Pethick, NPA175('71),225.

Bulk pure neutron matter EOS



$$W(k, 0) \approx \underbrace{19.74k^2}_{\text{kin. E}} - \underbrace{k^3}_{\propto \rho} \frac{(40.4 - 1.088k^3)}{(1 + 2.545k)}$$

↑
↑

Fermi mom.

Another way to use Cold Atom EOS

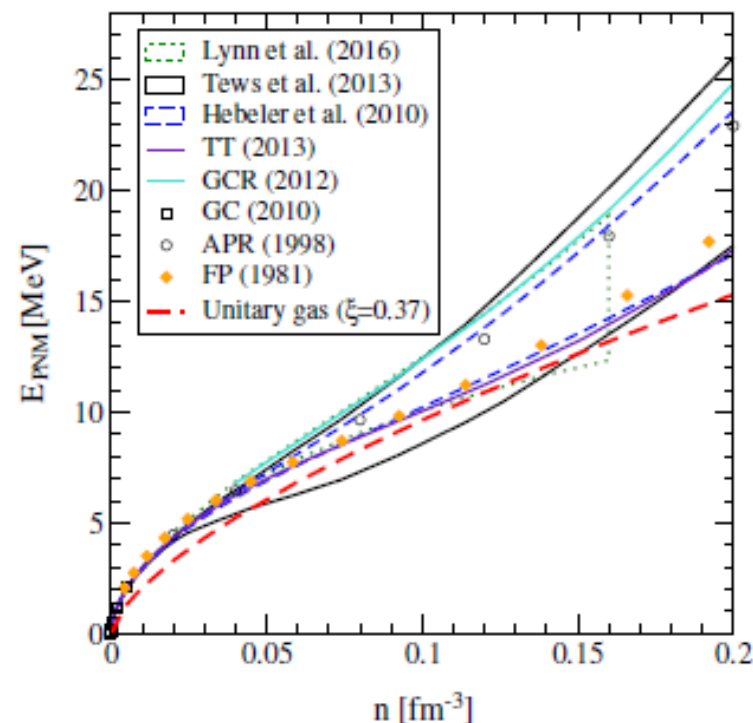
- Unitary gas EOS は ab initio pure neutron matter EOS の lower bound を与えているように見える。(Conjecture)

$$E_{\text{PNM}}(n) \geq E_{\text{UG}}(n)$$

$$E_{\text{UG}}(u) = \frac{3}{5} E_{\text{F}}(n) \xi = \frac{3\hbar^2 k_{\text{F}}^2}{10m_n} \xi = E_{\text{UG}}^0 u^{2/3}$$

$$(\xi \sim 0.37, u=n/n_0)$$

- これが正しいなら、対称エネルギーへの制限が与えられるのでは？



*I. Tews, J.M.Lattimer, AO, E.E.Kolomeitsev
(TLOK), ApJ 848 ('17)105*

Unitary Gas Constraint on Symmetry Energy

- 対称エネルギー =
純中性子物質エネルギー - 対称核物質エネルギー

$$S(u) = E_{\text{PNM}}(u) - E_{\text{SNM}}(u)$$

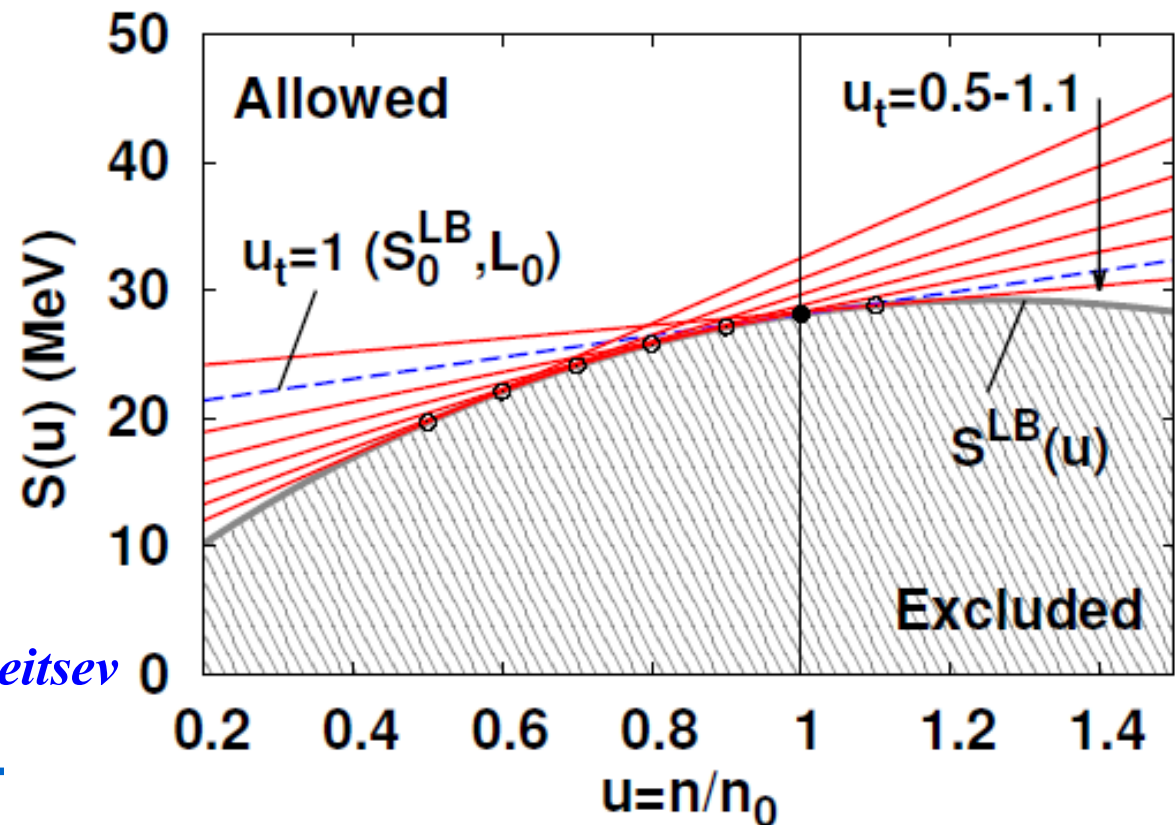
$$E_{\text{SNM}}(u) = E_0 + K(u-1)^2/18 + \mathcal{O}[(u-1)^3]$$

$$S(u) \geq E_{\text{UG}}^0 u^{2/3} - \left[E_0 + \frac{K}{18} (u-1)^2 \right] \equiv S^{\text{LB}}(u)$$

- S_0 への制限

$$S_0 \geq E_{\text{UG}}^0 - E_0 \equiv S_0^{\text{LB}}$$

- $S_0 > S_0^{\text{LB}}$ では、
L への制限



I. Tews, J.M.Lattimer, A.O., E.E.Kolomeitsev
(TLOK), *ApJ* 848 ('17)105

Unitary Gas Constraint on Symmetry Energy Parameters

■ Symmetry Energy Parameters (S_0, L) への制限

$$S_0 + \frac{L}{3}(u - 1) \geq E_{\text{UG}}^0 u^{2/3} - \left[E_0 + \frac{K_n}{18}(u - 1)^2 \right]$$

■ 接線条件 $\rightarrow L$ の領域

$$S_0 = \frac{E_{\text{UG}}^0}{3u_t^{1/3}}(u_t + 2) + \frac{K_n}{18}(u_t - 1)^2 - E_0$$

$$L = \frac{2E_{\text{UG}}^0}{u_t^{1/3}} - \frac{K_n}{3}(u_t - 1)$$

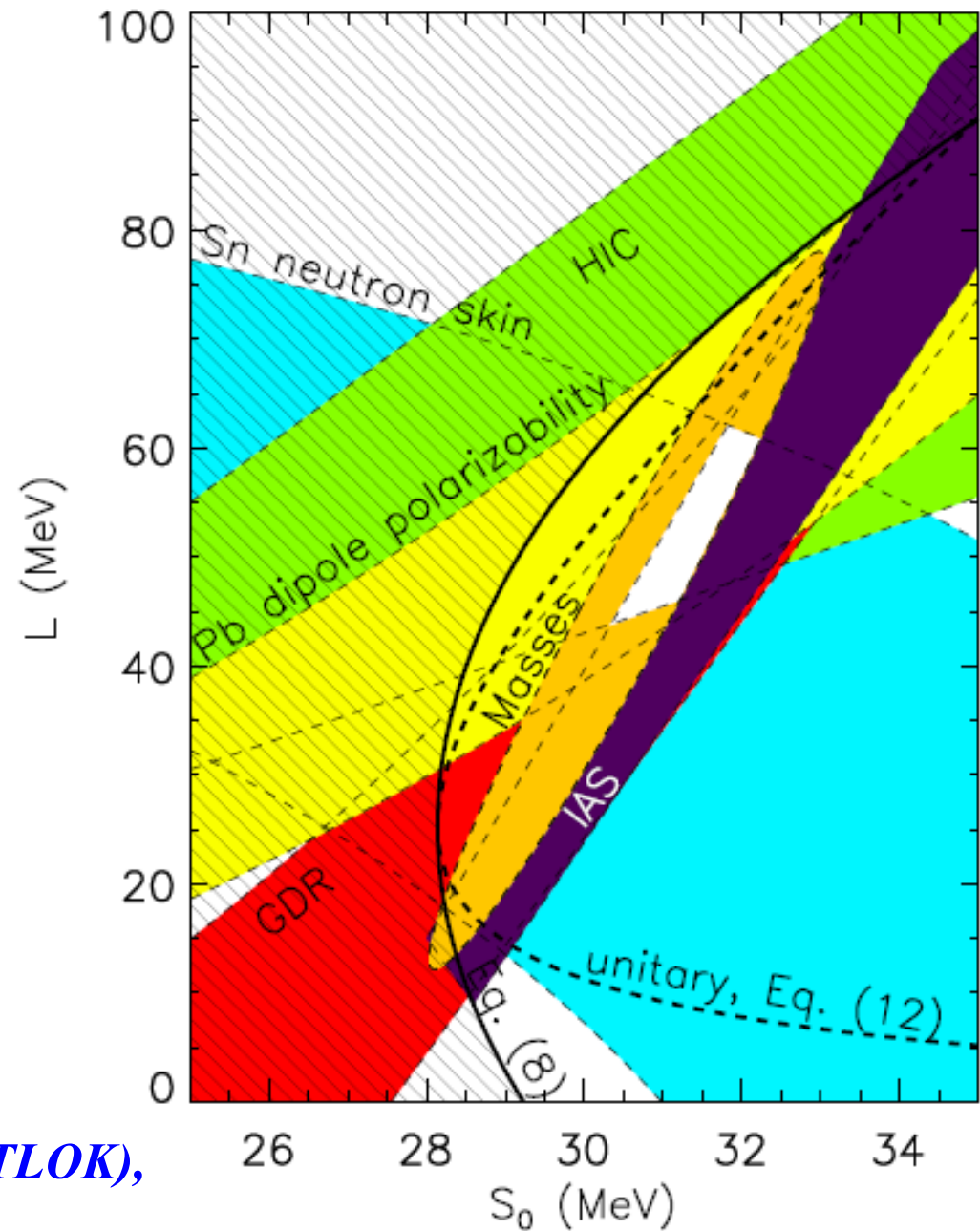
■ $(u-1)^3$ を無視しているので、不等式が成り立つには E_{SNM} を大きく、 E_{UG} を小さく見積もる必要あり。

$$E_0 = -15.9 \pm 0.4 \text{ MeV}, n_0 = 0.164 \pm 0.007 \text{ fm}^{-3},$$

$$K = 230 \pm 40 \text{ MeV}, \xi = 0.370 \pm 0.005$$

Unitary Gas Constraint vs Experimental Constraints

- Unitary gas constraint と
実験による制限は consistent



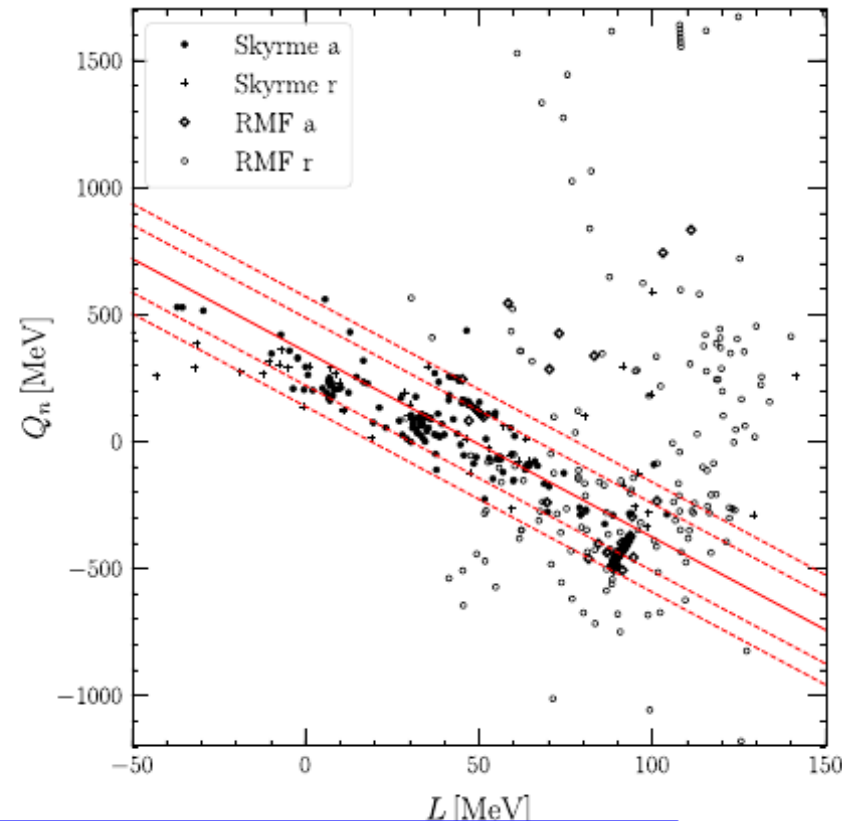
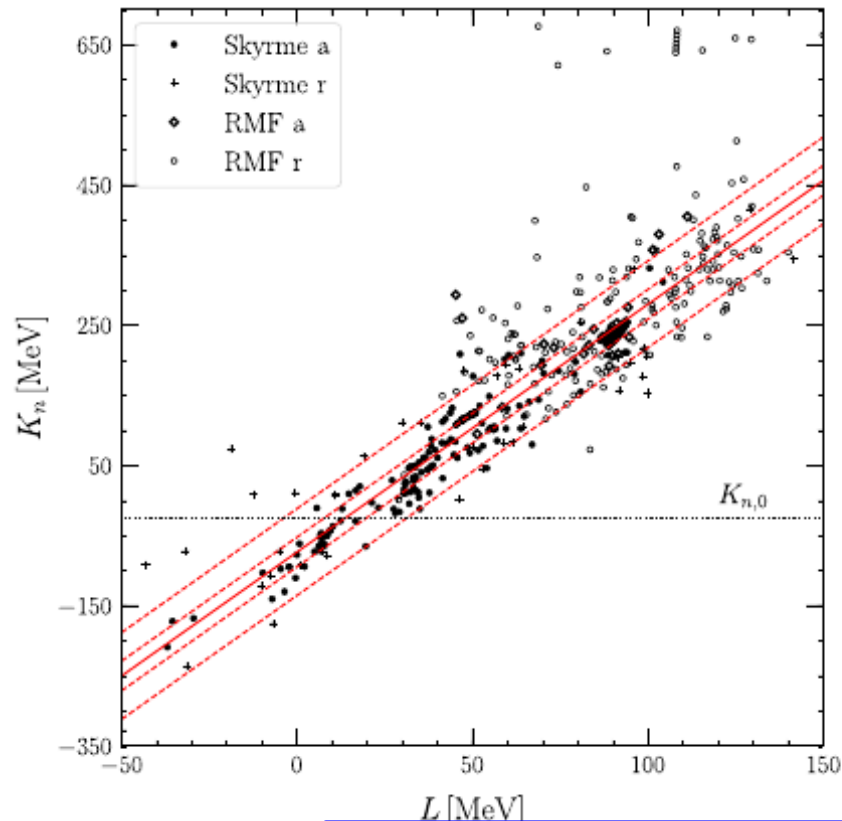
*I. Tews, J.M.Lattimer, AO, E.E.Kolomeitsev (TLOK),
ApJ 848 ('17)105*

Further Constraints on Higher-Order Sym. E. parameters

- K_n and Q_n are correlated with L in “Good” theoretical models.

$$K_n = 3.534L - (74.02 \pm 21.17)\text{MeV}$$

$$Q_n = -7.313L + (354.03 \pm 133.16)\text{MeV}$$



Regard theoretical models as data !

I. Tews, J.M.Lattimer, AO, E.E.Kolomeitsev (TLOK), ApJ 848 ('17)105

*Symmetry Energy Parameters,
Nuclear Matter EOS,
and Neutron Star Radius*

Fermi momentum (k_F) expansion

■ Saturation & Symmetry Energy Parameters

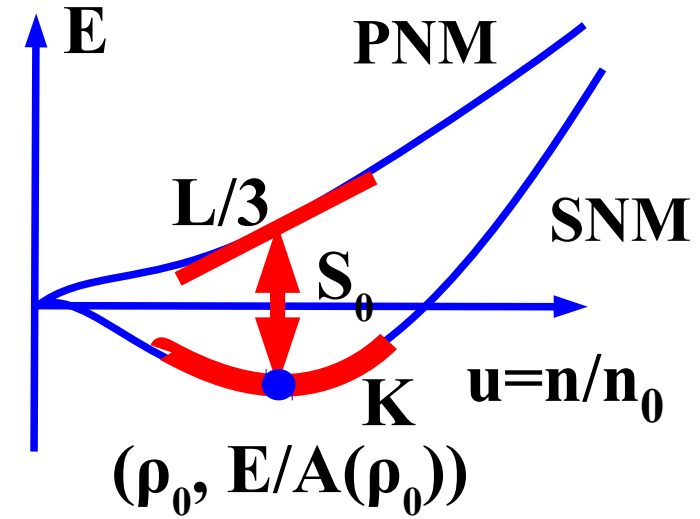
TLOK

$$E_{\text{NM}}(u, \alpha) = E_{\text{SNM}}(u) + \alpha^2 S(u)$$

$$E_{\text{SNM}}(u) \simeq E_0 + \frac{K_0}{18}(u-1)^2 + \frac{Q_0}{162}(u-1)^3$$

$$S(u) \simeq S_0 + \frac{L}{3}(u-1) + \frac{K_s}{18}(u-1)^2 + \frac{Q_s}{162}(u-1)^3$$

$$(u = n/n_0, \alpha = (n_n - n_p)/n)$$



Energy does not approach zero at $n \rightarrow 0$.

■ Fermi momentum expansion ($\sim$ Skyrme type EDF)

- Generated many-body force is given by $k_F \propto u^{1/3}$

$$E_{\text{SNM}}(u) \simeq \boxed{T_0 u^{2/3}} + \underline{a_0 u} + \underline{b_0 u^{4/3}} + \underline{c_0 u^{5/3}} + \underline{d_0 u^2}$$

$$S(u) \simeq \boxed{T_s u^{2/3}} + \underline{a_s u} + \underline{b_s u^{4/3}} + \underline{c_s u^{5/3}} + \underline{d_s u^2}$$

Kin. E. Two-body Density-dep. pot.

m^*

Expansion Coefficients

- Coefficients (a,b,c,d) are represented by Saturation and Symmetry Energy Parameters

TLOK

$a_0 = -4T_0$	$+20E_0$	$+K_0$	$-Q_0/6$
$b_0 = 6T_0$	$-45E_0$	$-5K_0/2$	$+Q_0/2$
$c_0 = -4T_0$	$+36E_0$	$+2K_0$	$-Q_0/2$
$d_0 = T_0$	$-10E_0$	$-K_0/2$	$+Q_0/6$
$a_s = -4T_s$	$+20S_0 - 19L/3$	$+K_s$	$-Q_s/6$
$b_s = 6T_s$	$-45S_0 + 15L$	$-5K_s/2$	$+Q_s/2$
$c_s = -4T_s$	$+36S_0 - 12L$	$+2K_s$	$-Q_s/2$
$d_s = T_s$	$-10S_0 + 10L/3$	$-K_s/2$	$+Q_s/6$

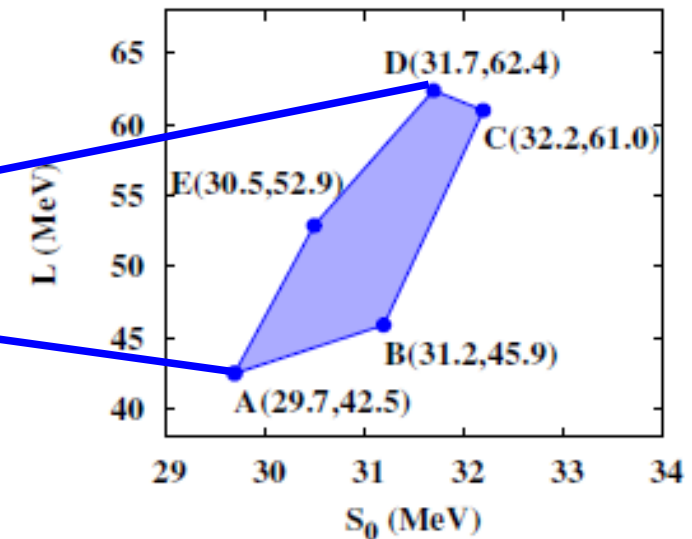
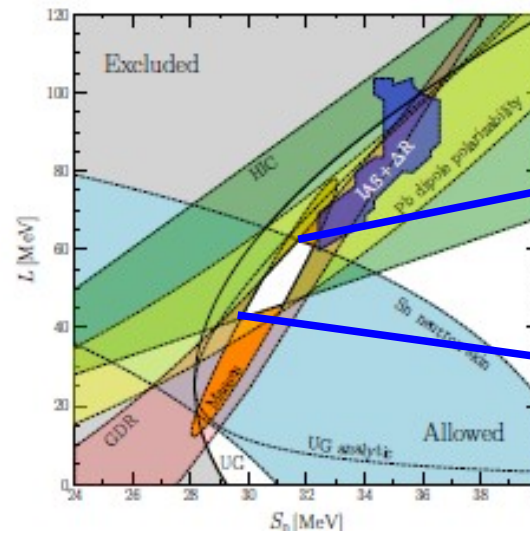
$$\left(T_0 = \frac{3}{5} \frac{\hbar^2 k_F (n_0)^2}{2m}, \quad T_s = T_0 (2^{1/3} - 1) \right)$$

Tedious but straightforward calc.

$TLOK+2M_{\odot}$ constraints

■ TLOK constraints

- (S_0, L) is in Pentagon.
- (K_n, Q_n) are from TLOK constraint.
- $K_0=(190-270)$ MeV
- (n_0, E_0) is fixed
 $n_0=0.164$ fm⁻³, $E_0=-15.9$ MeV (small uncertainties)
- Q_0 is taken to kill d_0 parameter
(Coef. of u^2 . Sym. N. M. is not very stiff at high-density)



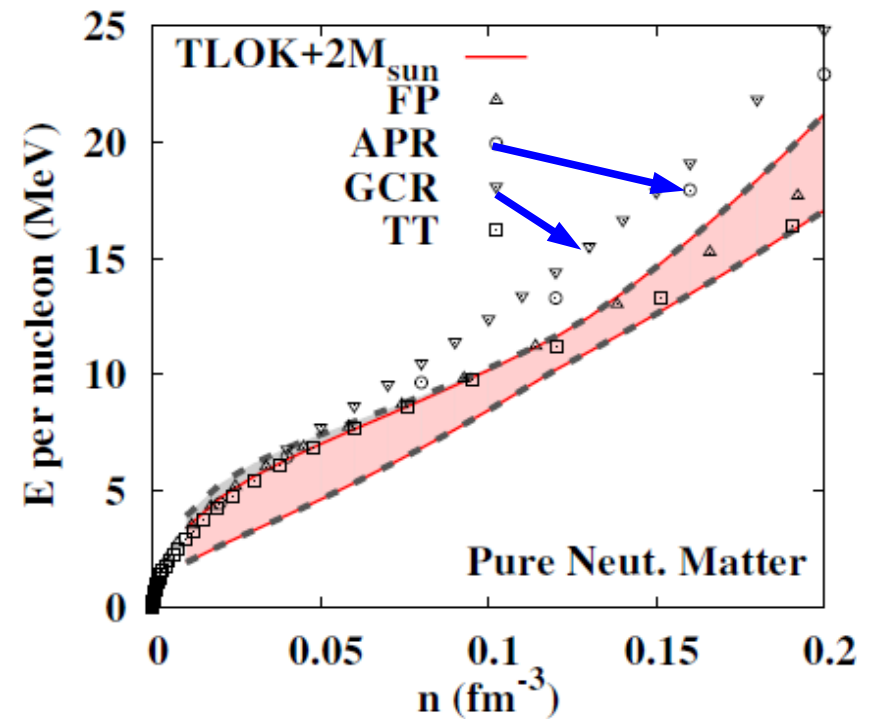
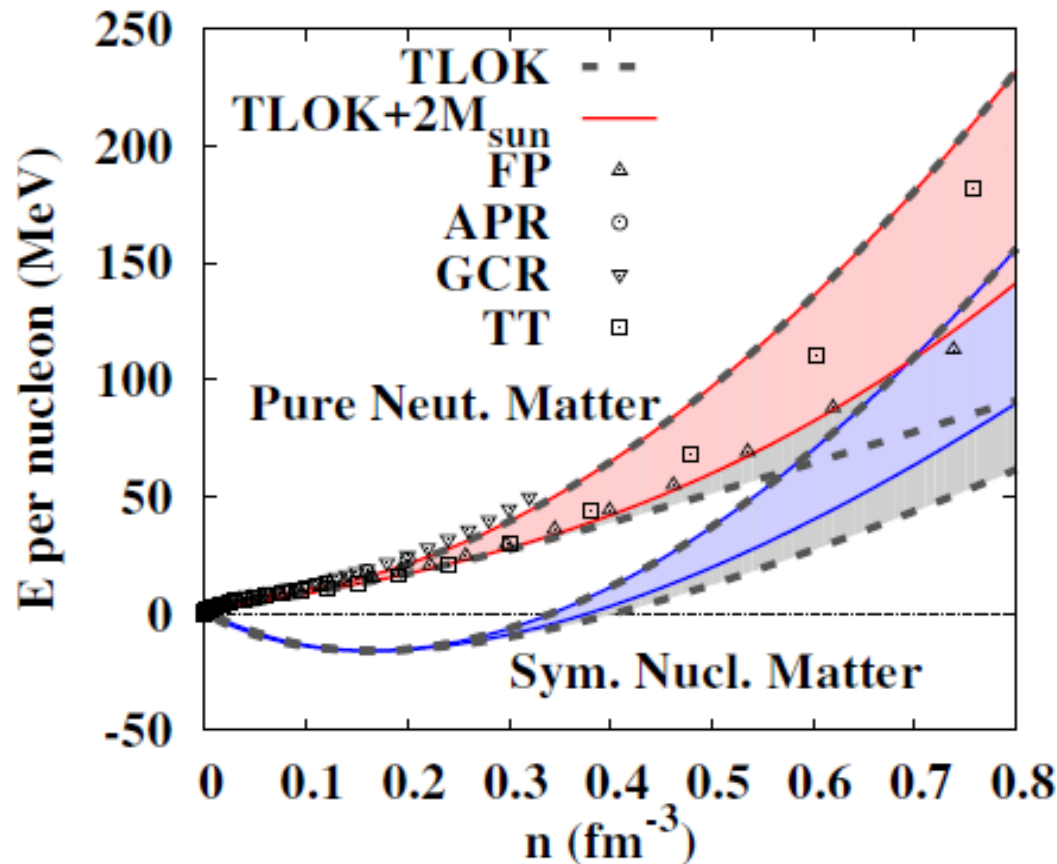
■ $2 M_{\odot}$ constraint

- EOS should support $2 M_{\odot}$ neutron stars.

AO, Kolomeitsev, Lattimer, Tews, Wu (OKLTW), in prog.

$TLOK+2M_{\odot}$ constraints on EOS

- $2M_{\odot}$ constraint narrows the range of EOS.
- Consistent with FP and TT(Togashi-Takano) EOSs.
- APR and GCR(Gandolfi-Carlson-Reddy) EOSs seems to have larger S_0 values.



OKLTW, in prog.

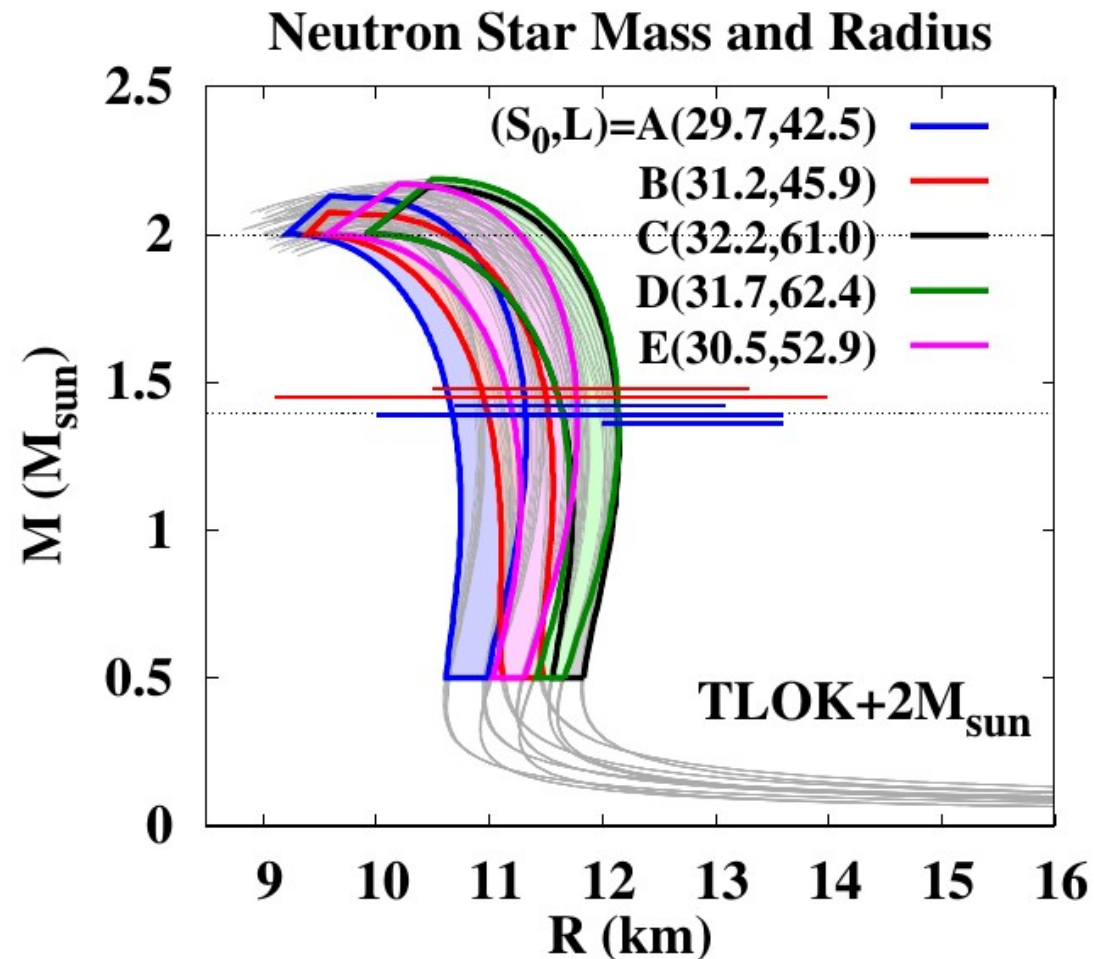
Neutron Star MR curve

■ TLOK + $2 M_{\odot}$ constraints $\rightarrow R_{1.4}=(10.6-12.2)$ km *OKLTW, in prog.*

- E and P are linear fn. of Sat. & Sym. E. parameters
 $\rightarrow$ Min./Max. appears at the corners of pentagon (ABCDE).

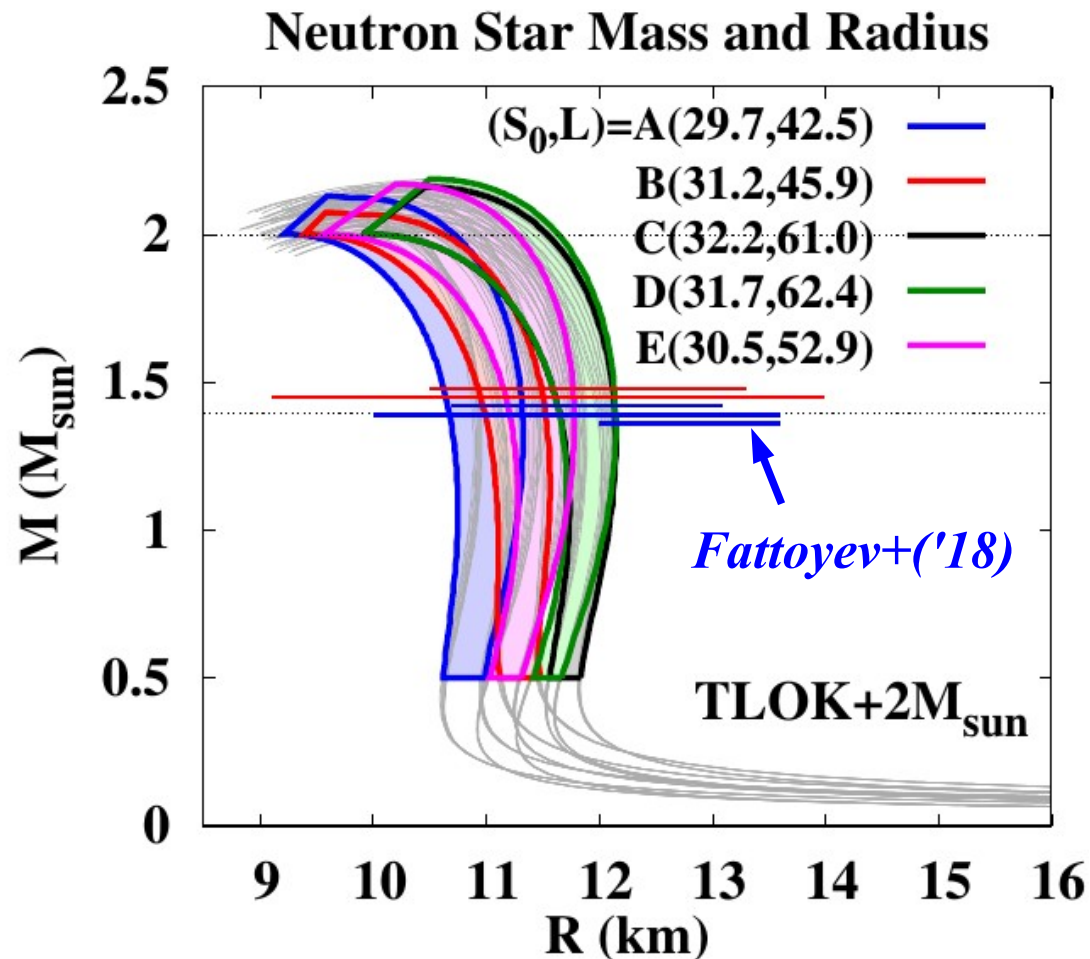
- For a given (S_0, L) ,
 unc. of $R_{1.4} \sim 0.5$ km
 = unc. from higher-order parameters

- Unc. from $(S_0, L) \sim 1.1$ km
 $\rightarrow$ We still need to fix
 (S_0, L) more precisely.

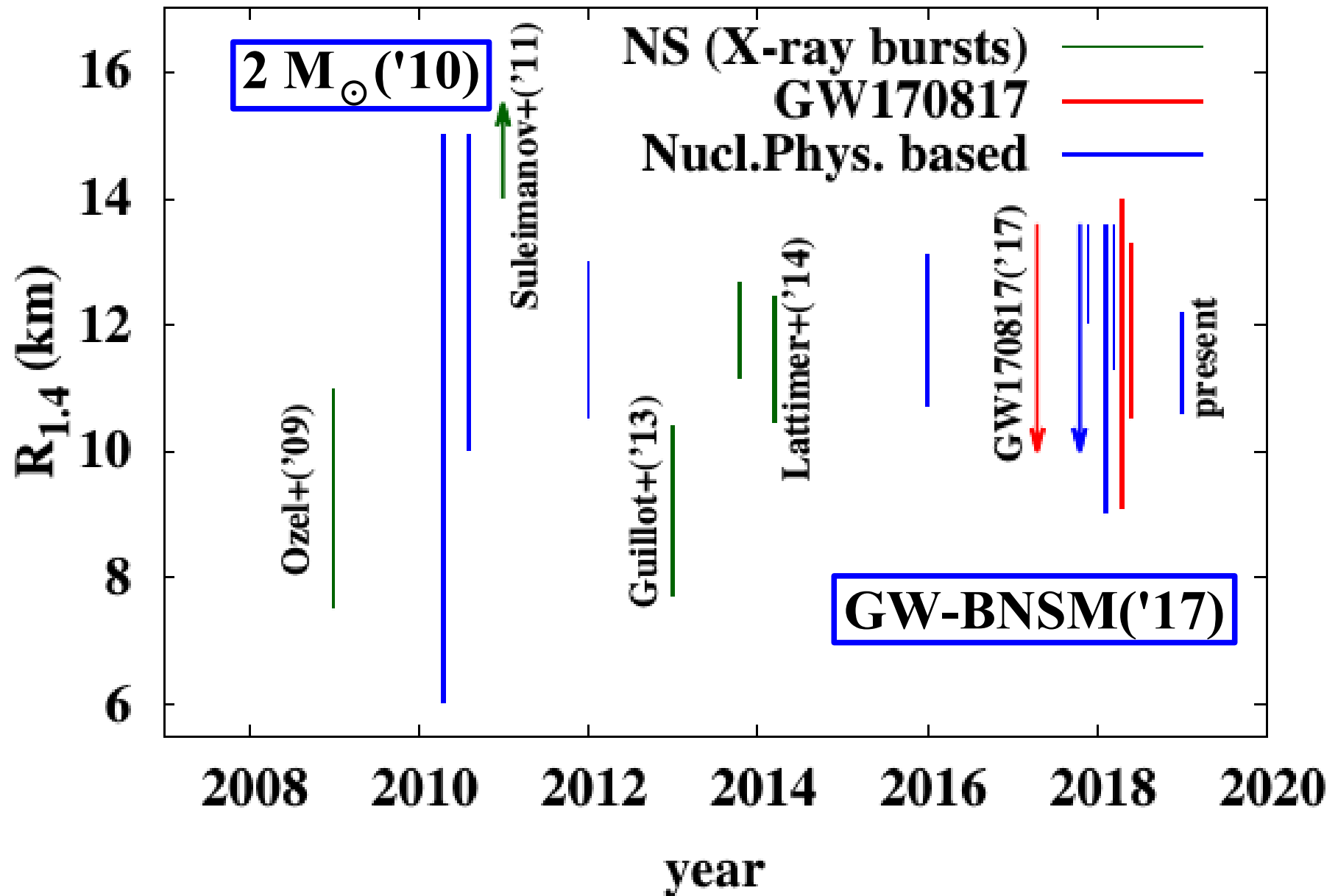


Neutron Star MR curve

- Our constraint is consistent with many of previous ones.
 - $R_{1.4}=(10.6-12.2)$ km *Present work (TLOK + 2 $M_{\odot}$) OKLTW, in prog.*
 - LIGO-Virgo (Tidal deformability Λ from BNSM)
 - (10.5-13.3) km *Abbott+('18b)*
 - (9.1-14.0) km *De+('18) (Λ)*
 - Theoretical Estimates
 - (10.7-13.1) km
Lattimer+, PRep.621('16)127
 - (10.0-13.6) km
Annala+, PRL120('18)172703
 - (9-13.6) km
Tews+, PRC98 ('18)045804
 - (12.0-13.6) km
F.J.Fattoyev+(PREX), PRL120 ('18)172702



Time dependence of Neutron Star Radius ($R_{1.4}$)



*Implications to quark-hadron physics
in cold dense matter*

Questions !

■ Hyperon puzzle

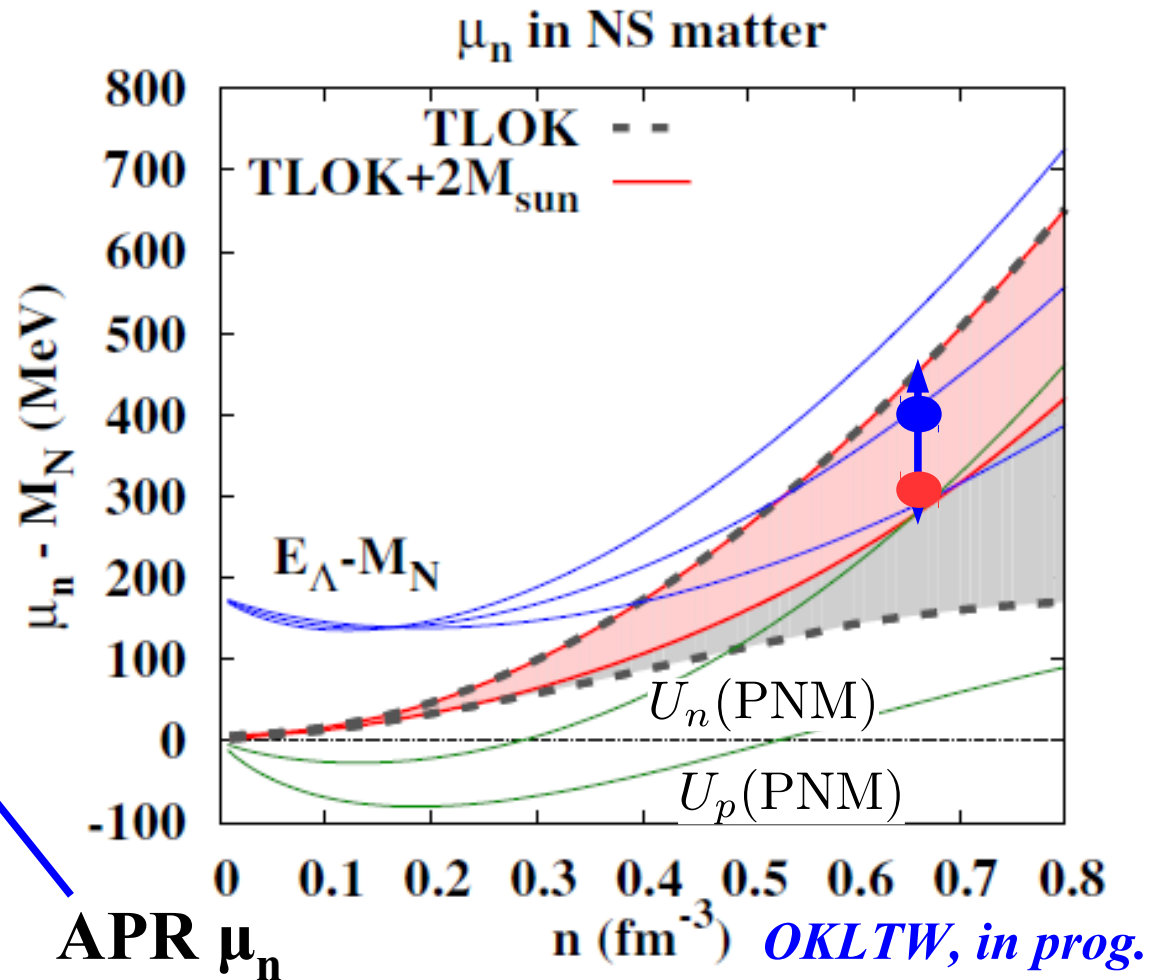
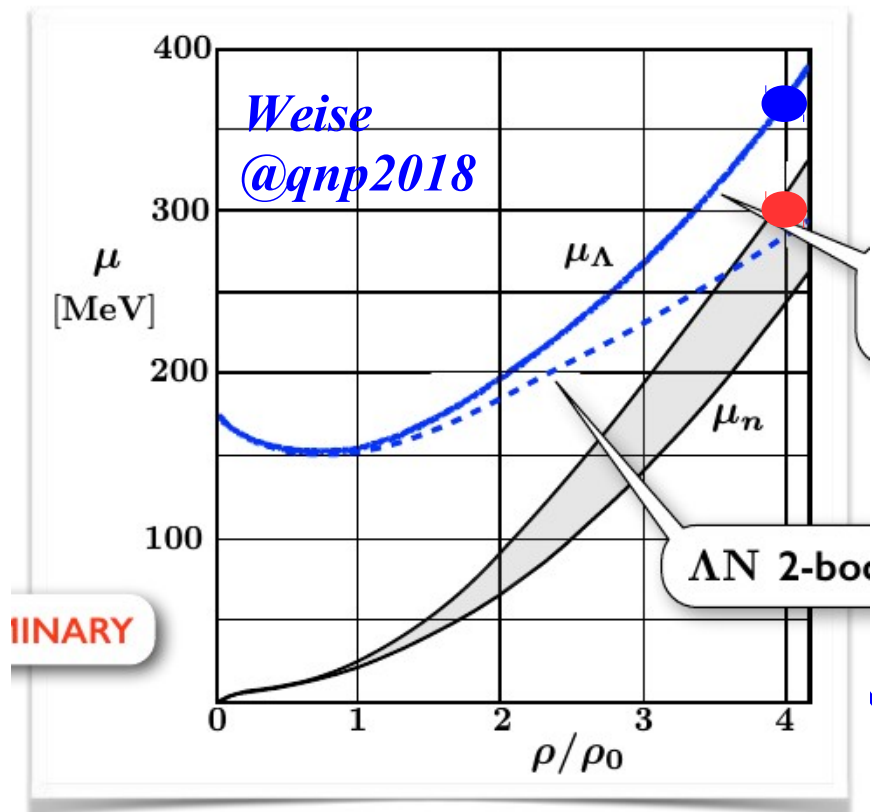
- At what density do hyperons appear ? $\rightarrow U_{\Lambda} = \mu_n$
- In STANDARD EOS with hyperons with $U_{\Lambda}(n_0) = -30$ MeV, Λ appears at $n = (2-3)n_0$
- Density dep. of U_{Λ} is essential.
- Neutron chemical potential strongly depends on sym. E.

■ QCD phase transition in cold dense matter

- Do we have the first order phase transition in cold dense matter ?
If yes, at which density ?
- Recent high-energy heavy-ion collision data suggest strong softening of EOS at $n = (5-10) n_0$.
- With hadronic matter EOS with $L = 50$ MeV and NJL model, mixed phase would appear at $n = (5-10) n_0$ in neutron stars.

Neutron Chemical Potential in NS

- Λ appears in neutron stars if $E_\Lambda(p=0) = M_\Lambda + U_\Lambda < \mu_n$
- U_Λ in χ EFT (2+3 body) is stiff.
- But μ_n is larger with TLOK+2M $_\odot$ constraints



*W. Weise, NFQCD2018 (2018.06);
Gerstung, Kaiser, Weise, in prog.*

Neutron Chemical Potential in NS

■ Neutron Chemical Potential

$$\mu_n + M_N = \frac{\partial(nE)}{\partial n_n} = E + u \frac{\partial E}{\partial u} + 2\alpha(1 - \alpha)S(u)$$

■ Single particle potential

$$U_\Lambda(u) = \frac{\partial(nV)}{\partial n_\Lambda}$$

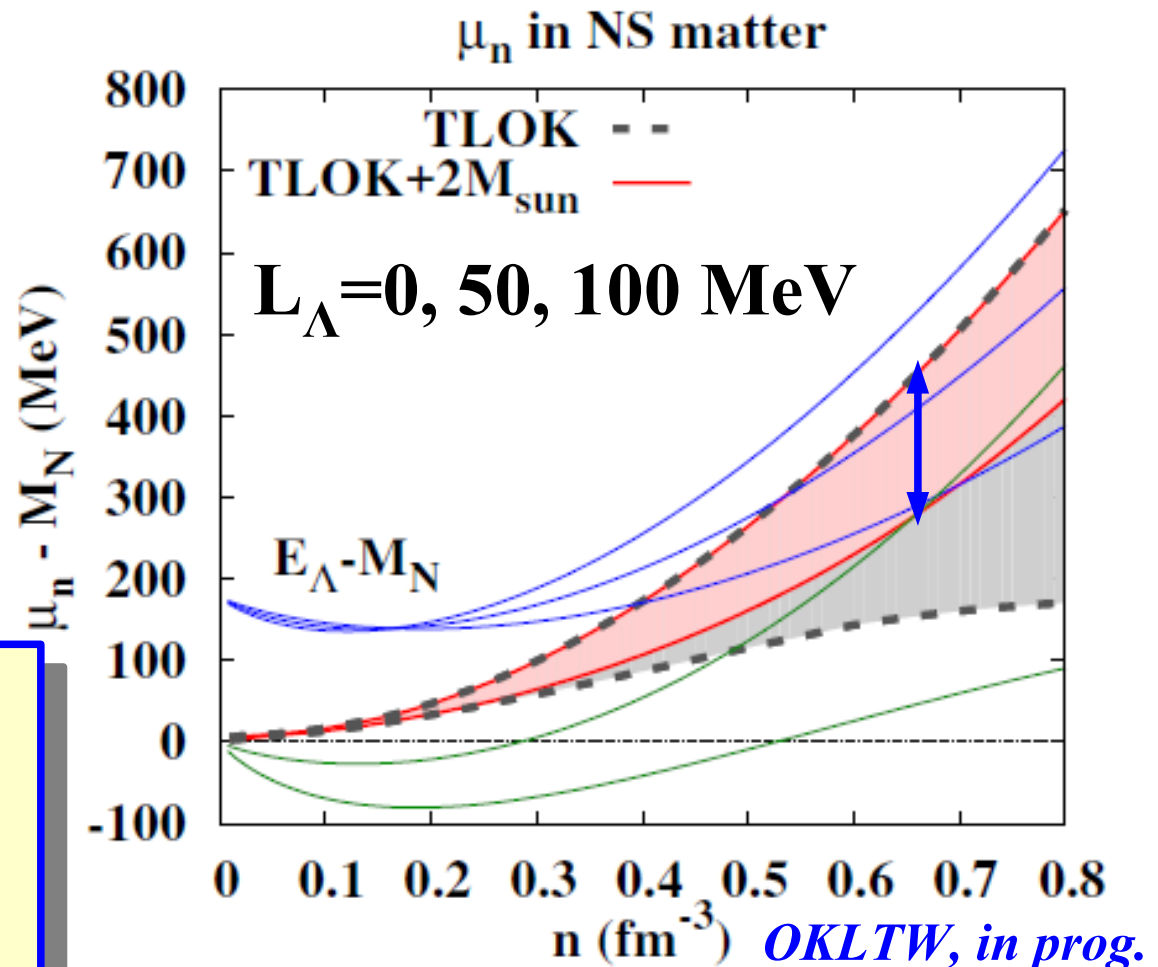
$$\simeq U_{0\Lambda} + \frac{L_\Lambda}{3}(u - 1)$$

$$U_{0\Lambda} \simeq -30 \text{ MeV}$$

$$L_\Lambda = ???$$

($L_\Lambda < 0$ in most of RMF before 2010)

Sym. E and L_Λ determine the onset density of Λ . (Already mentioned in Millener, Dover, Gal paper)

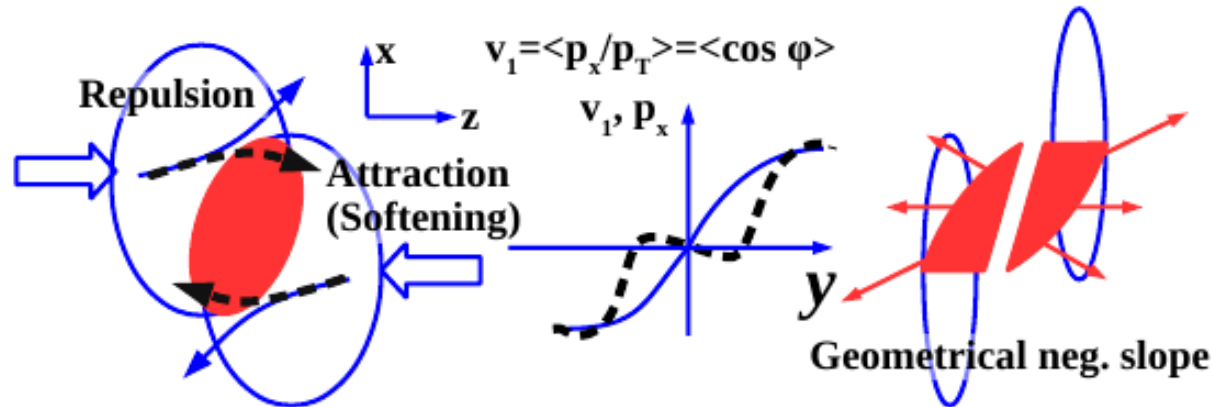


QCD phase transition in cold dense matter

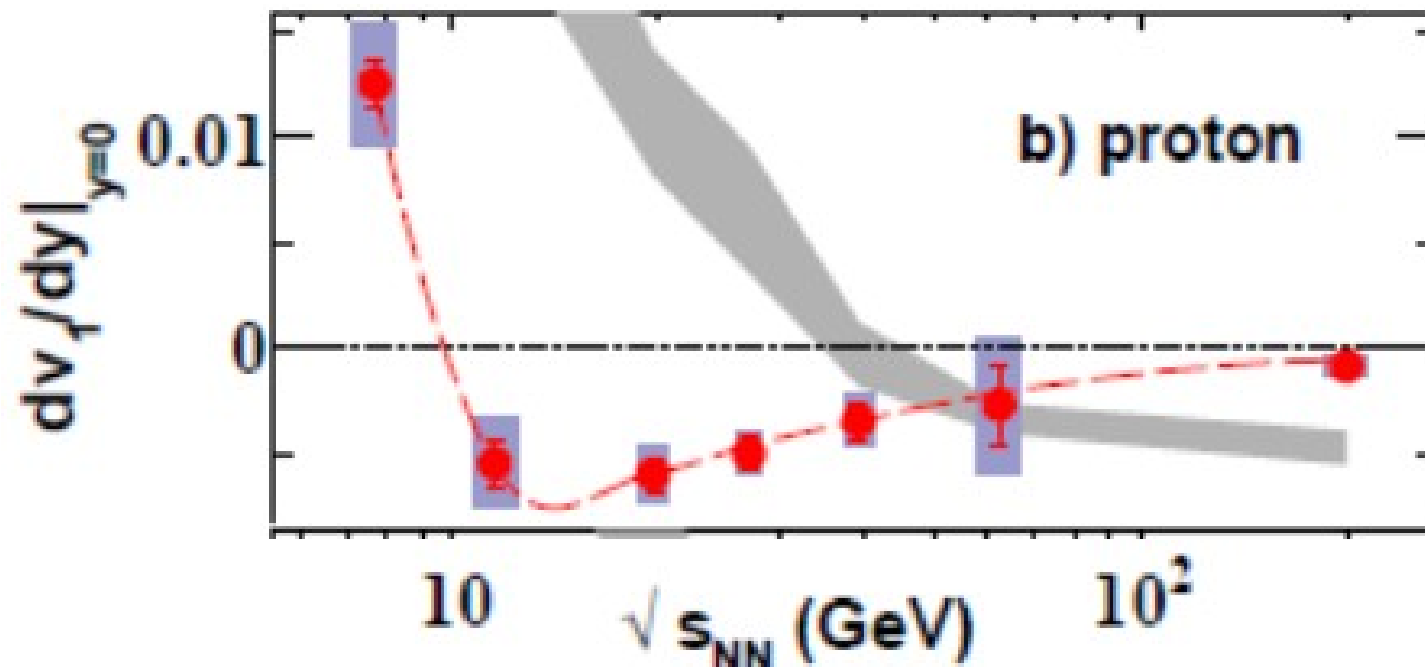
- Transition to quark matter in cold-dense matter
1st order or crossover ?
- Crossover: Masuda, Hatsuda, Takatsuka, Kojo, Baym, ...
- 1st order p.t.
 - Many effective models predict, e.g. Asakawa-Yazaki CP
 - Recent phenomenological support: Negative Directed Flow in HIC
Y.Nara, H.Niemi, AO, H.Stoecker, PRC94('16)034906.
Y. Nara, H. Niemi, AO, J. Steinheimer, X.-F. Luo, H. Stoecker, EPJA 54 ('18)18
 - The phase transition density may be above NS central density
X.Wu, AO, H.Shen, PRC to appear (arXiv:1806.03760)

Negative Directed Flow

- Directed Flow $v_1 = \langle \cos \phi \rangle = \langle p_x / p_T \rangle$, Slope = dv_1/dy



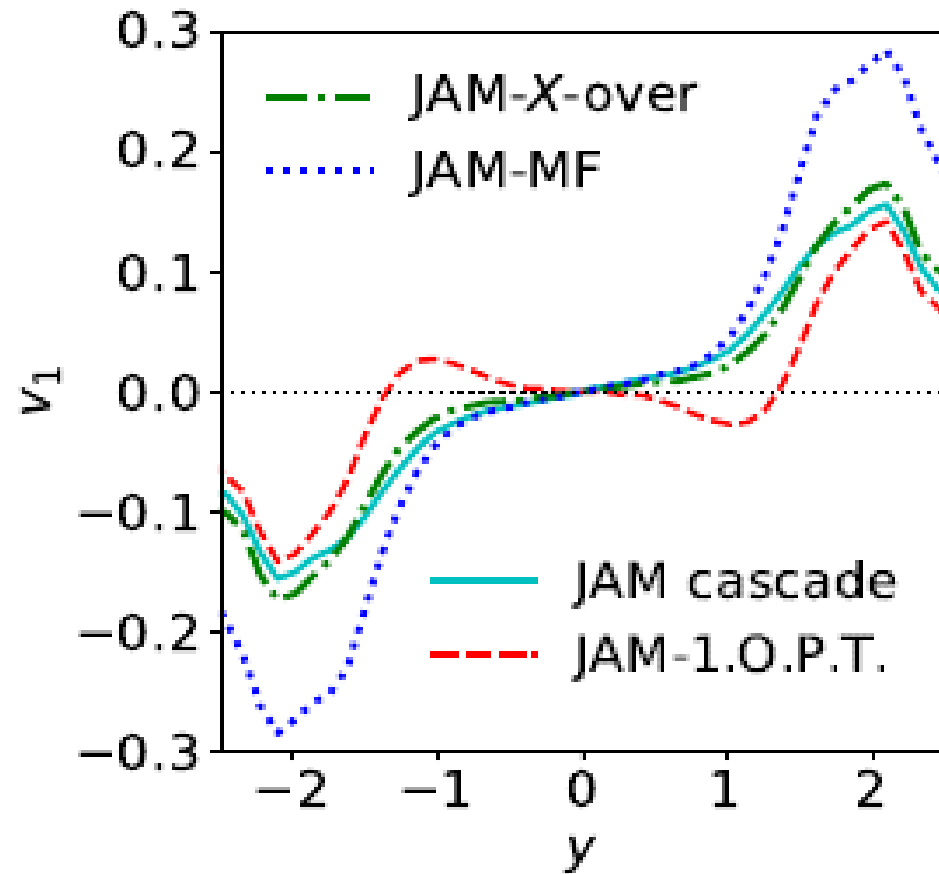
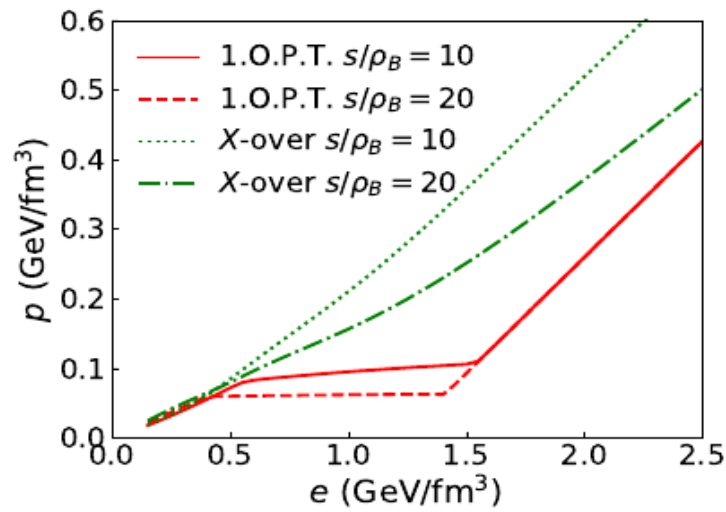
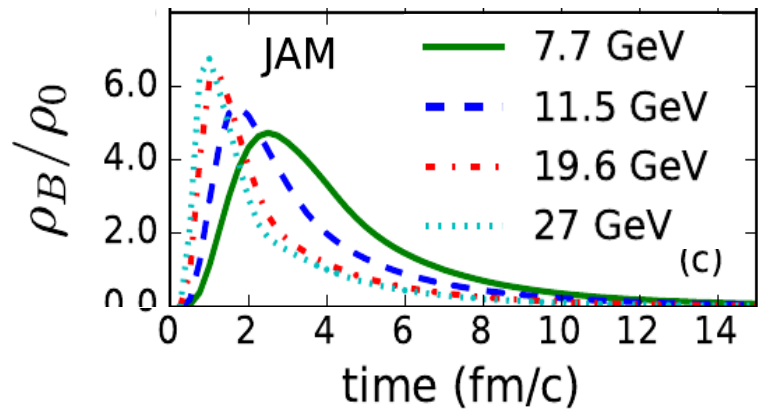
- Negative Flow in Heavy-Ion Collisions



STAR Collab. (L. Adamczyk et al.), Phys.Rev.Lett. 112 ('14), 162301

Negative Directed Flow

- Negative Directed Flow slope at $\sqrt{s_{NN}} = 11.5$ GeV (STAR ('14))
 → Strong softening of EOS is necessary at $n > (5-10) n_0$



*Y.Nara, H.Niemi, AO, H.Stoecker, PRC94('16)034906.
 Y. Nara, H. Niemi, AO, J. Steinheimer, X.-F. Luo,
 H. Stoecker, EPJA 54 ('18)18*

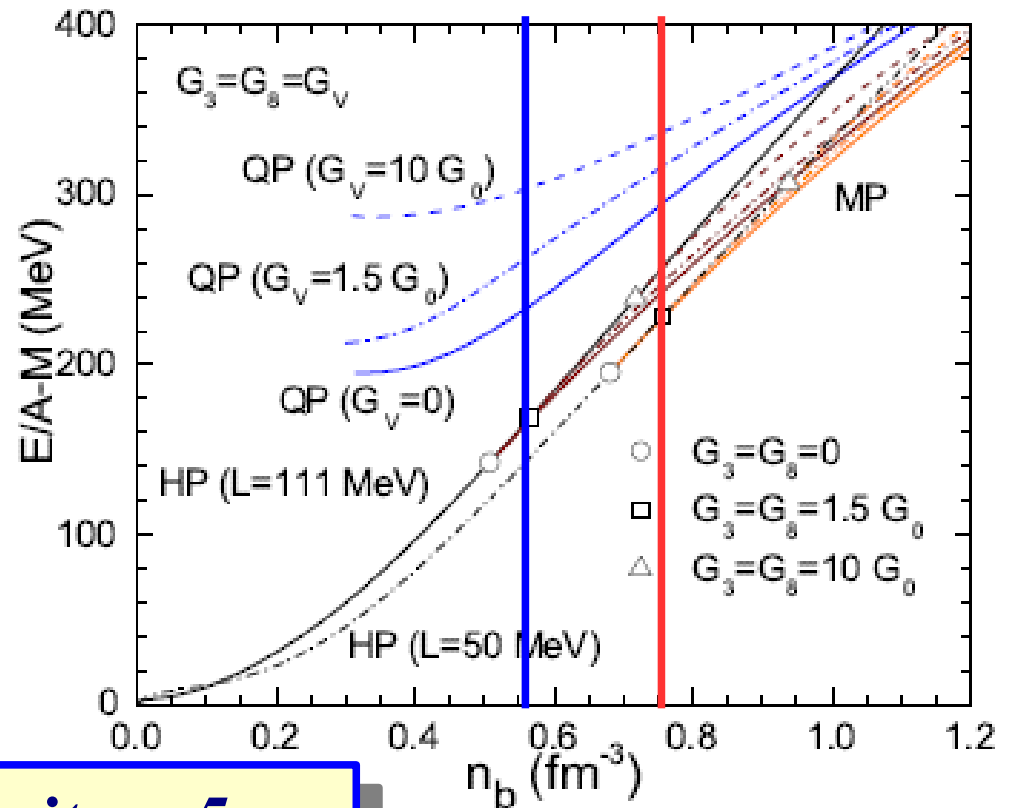
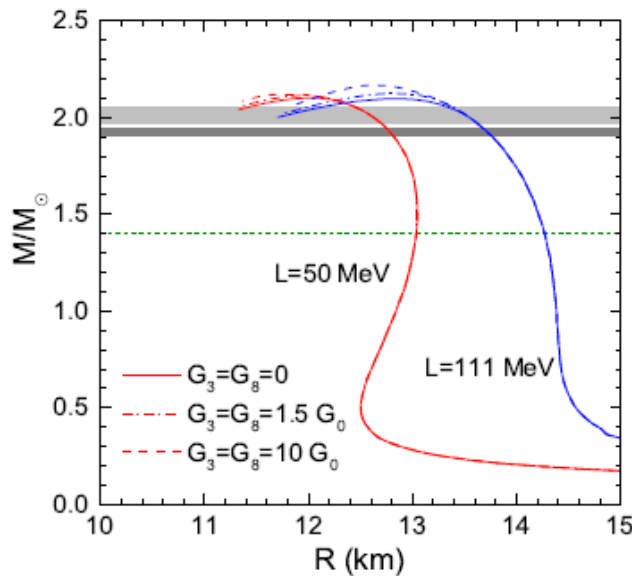
Isospin & Hypercharge Sym. E in quark matter

- Two types of vector int. in NJL $\rightarrow$ Isospin & Hypercharge Sym. E

X.Wu, AO, H.Shen, PRC to appear (arXiv:1806.03760)

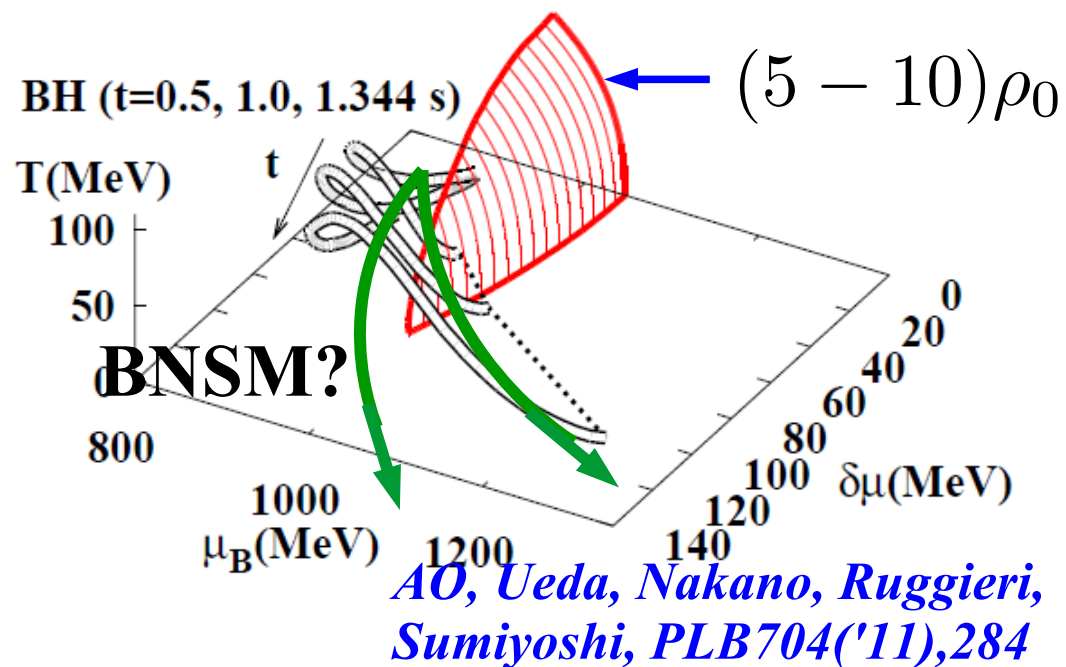
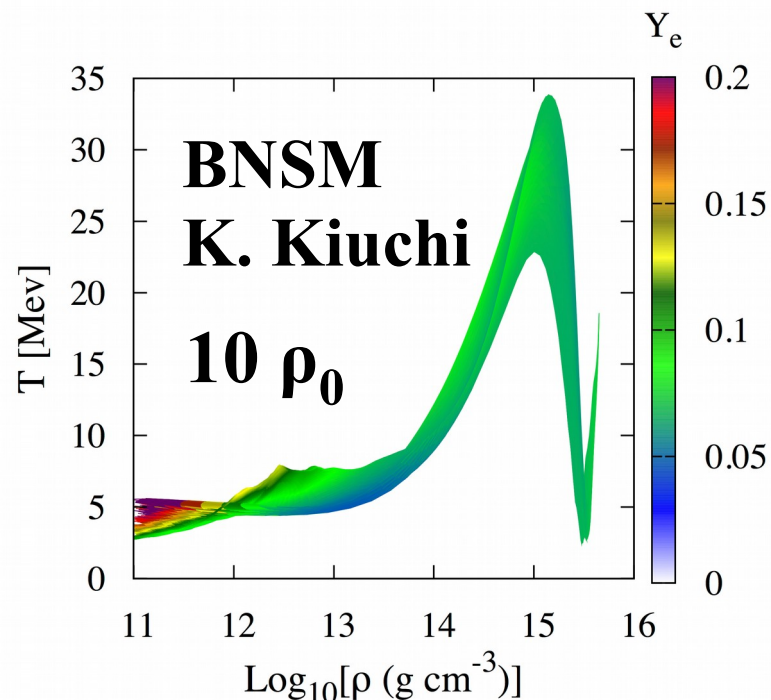
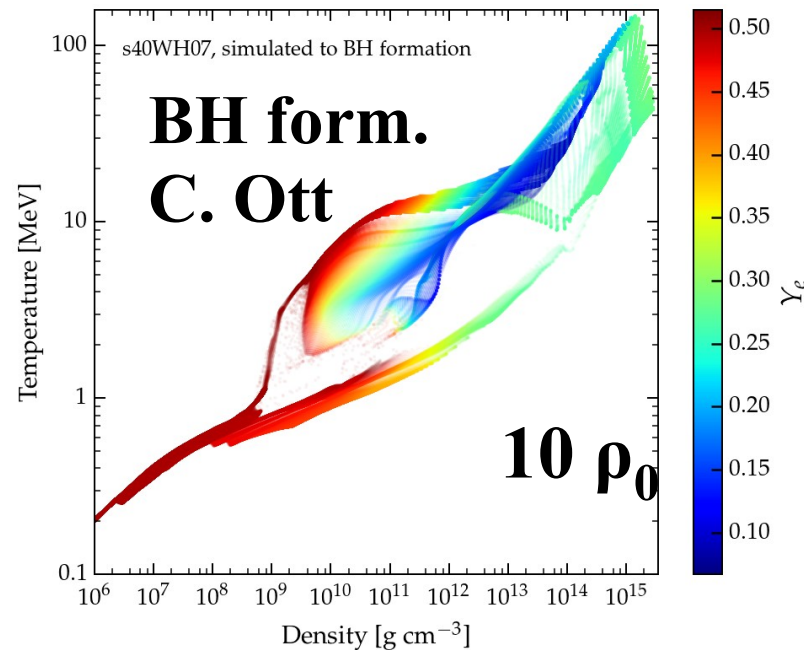
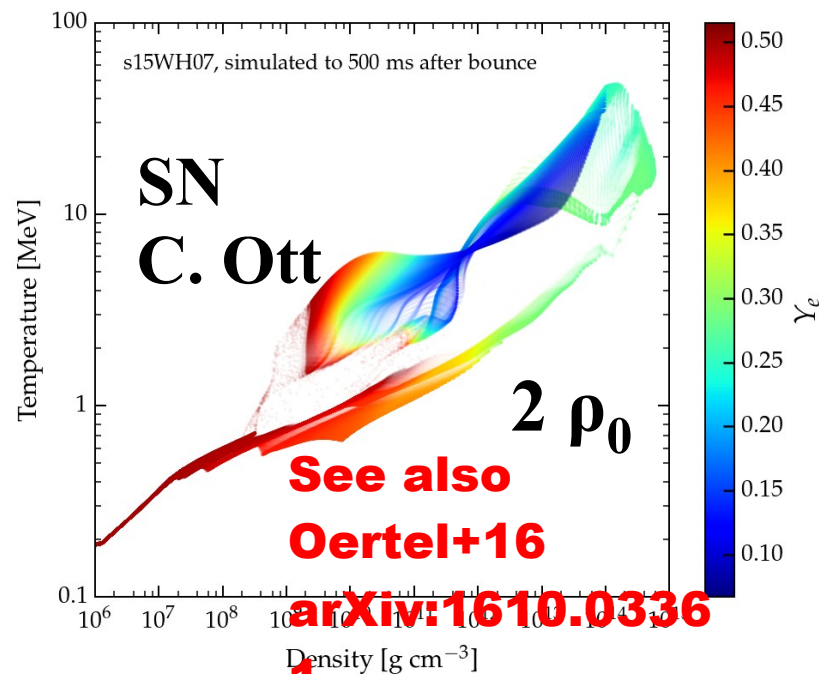
$$\mathcal{L}_v = -G_0(\bar{q}\gamma_\mu q)^2 - G_v \sum_i [(\bar{q}\gamma_\mu \lambda_i q)^2 + (\bar{q}i\gamma_5 \gamma_\mu \lambda_i q)^2]$$

$$E = \alpha^2 S(n) + \alpha_Y^2 S_Y(n), \quad \alpha = -2\langle T_z \rangle / B, \quad \alpha_Y = \langle B + S \rangle / B$$



$L=50$ MeV $\rightarrow$ transition density $\sim 5 n_0$

(ρ, T, Y_e) during SN, BH formation, BNSM



Summary

- **Constraint on symmetry energy parameters (S_0 , L , K_n , Q_n) together with $2 M_\odot$ constraint gives the $1.4 M_\odot$ neutron star radius in the range of (10.6-12.2) km.**
 - Consistent with constraint from GW, e.g. (10.5-13.3) km by LIGO
 - Fermi momentum (k_F) expansion is invoked.
Smooth extrapolation to (2-3) n_0 seems to work.
 - Let's wait for the NICER data and next NS-NS merger event.
- **Onset density of hyperons may be sensitive to the symmetry energy in addition to potential parameters, (U_{0B} , L_B).**
 - We need to know the slope of potential in addition to the depth.
- **QCD phase transition with strong EOS softening is expected at $n=(5-10)n_0$ in almost sym. n.m. from heavy-ion data.**
 - GW data from HMNS would clarify 3D phase diagram structure.

Thank you for your attention !

新学術領域「中性子星核物質」(2012-2017)

High ρ (Group A)
head: Tamura, Takahashi

Hypernuclei, Kaonic nuclei
YN & YY int.,
Eff. Interaction
(Heavy-ion collisions)

J-PARC



PI: H. Tamura

Hyperons, mesons, quarks

Asym. nuclear matter
+elec.+ μ

Nuclei+neutron gas+elec.

Nuclei + elec.

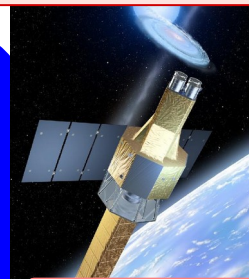
Low ρ (Group B)
head: Murakami,
Nakamura, Horikoshi

Sym. E, Pairing gap,
BEC-BEC cross over,
Cold atom, Unitary gas

NS Obs. (Group C)
head: Takahashi

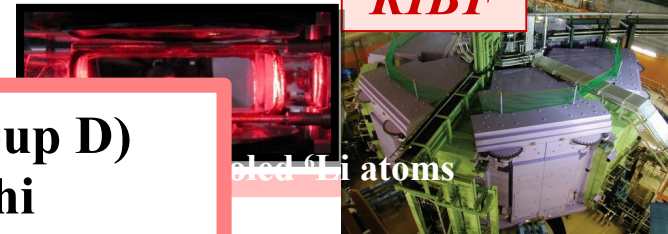
Radius, Mass,
Temp. (Cooling),
Star quake, Pasta

ASTRO-H



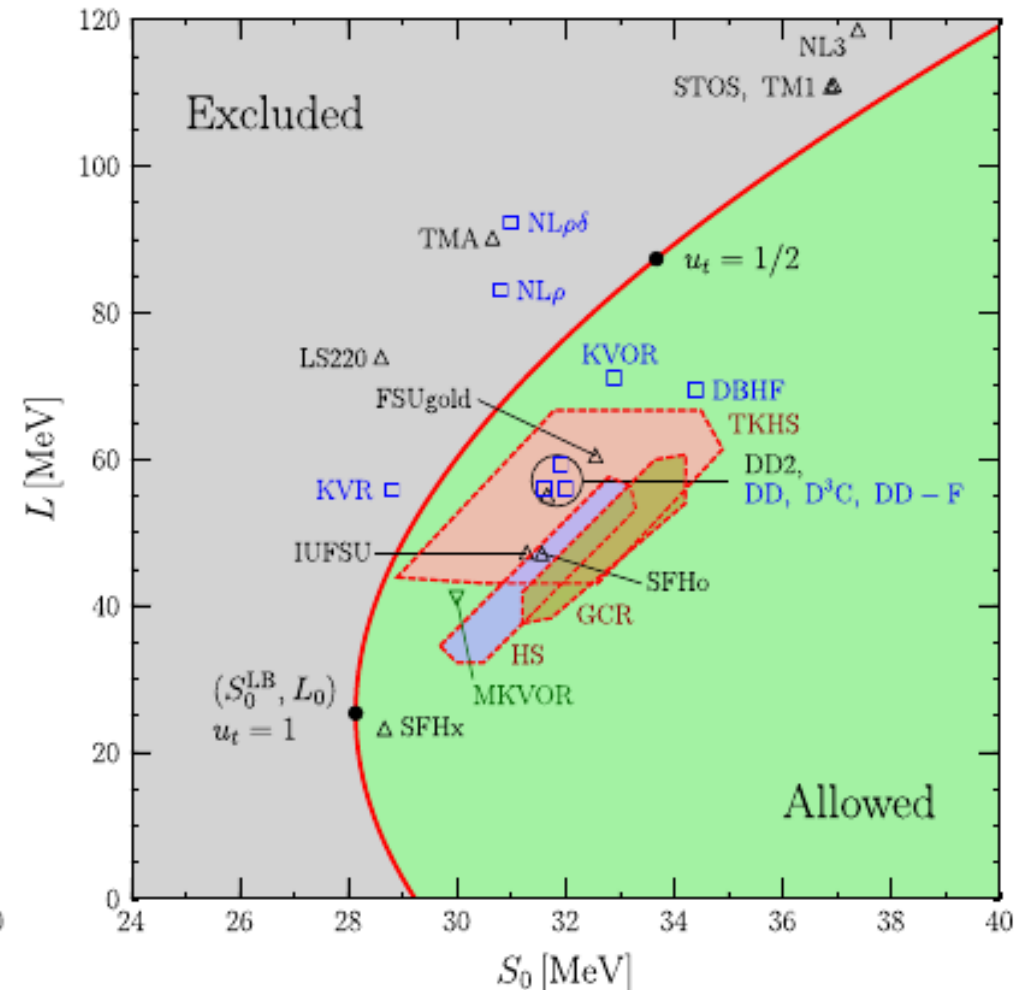
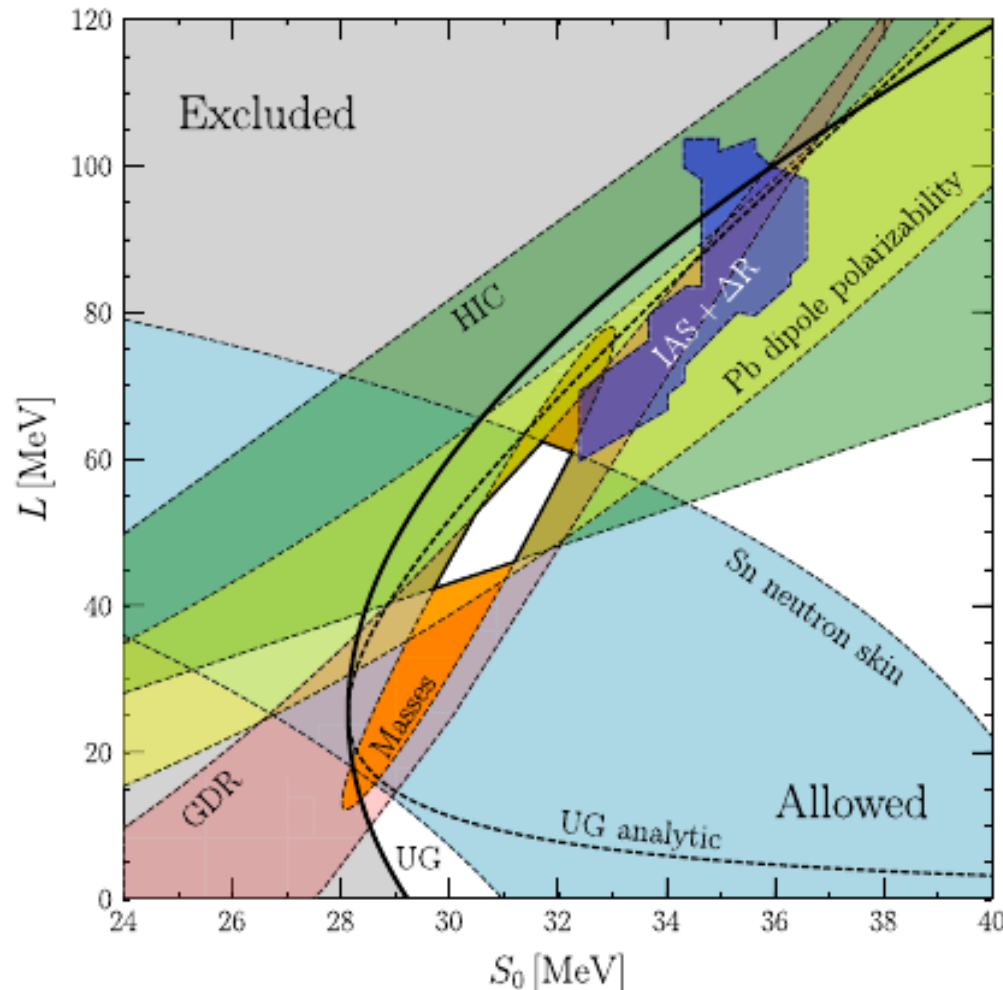
Theory (Group D)
head: Ohnishi

RIBF



Constraint on (S_0, L) from Lower Bound of PNM Energy

- Unitary gas ($E_{\text{PNM}} > E_{\text{UG}}$) + $2 M_\odot$ constraints rule out 5 EOSs (incl. LS220, Shen) out of 10 numerically tabulated ones.



I. Tews, J.M.Lattimer, AO, E.E.Kolomeitsev (TLOK), ApJ 848 ('17)105

Reservations and Prospects

Reservations

- Only massless electrons are considered and Crust EOS is ignored.
 - With μ , chemical potential may be reduced a little.
- Non-relativistic kinetic energy is used.
 - With rel. K.E., E per nucleon is modified by 0.03 MeV @ $10 n_0$ as long as Sat. and Sym. E parameters are fixed.
- Function form is limited to k_F expansion with $u^{k/3}$ ($k=2-6$).
 - $R_{1.4}$ range becomes narrower with $k=2-5$.
 - Density expansion gives EOSs very sensitive to parameters.
- Smooth $E(u)$ (= No phase transition) is assumed.
 - We expect QCD phase transition at $(5-10) n_0$ from recent BES data of directed flow *Nara, Niemi, AO, Stoecker ('16)*
 - Transition to quark matter may not soften EOS drastically.
- Causality is violated at high densities, $n > (4-6) n_0$.

To Do (or Prospect)

- Baryons other than nucleons Λ , Δ , Ξ , Σ , ...
 - Connecting to Hadron Resonance Gas (HRG) EOS
 - HRG EOS
- mass and kinetic E of hadrons with $M < 2$ GeV + simple potential E

$$\varepsilon_{\text{HRG}} = \mathcal{T} + cn^2$$

or Lattice EOS in HIC (No saturation, No constraint from NS).

- We need to guess the potential energy density more seriously for consistent understanding of HIC, Nuclear, and NS physics.

$$\varepsilon = \mathcal{T} + \mathcal{V} \longleftarrow \text{Nuclear and NS physics}$$

- Connecting to Quark(-Gluon) matter EOS
 - Embed model-H singularities *E.g. Nonaka, Asakawa ('04)*
 - “Interpolation” of nuclear and quark matter EOS

Further Constraint on Q_n

- $2 M_{\odot}$ requirement constrains Q_n further.

$$Q_n > -9.3L + 480 \text{ MeV}$$

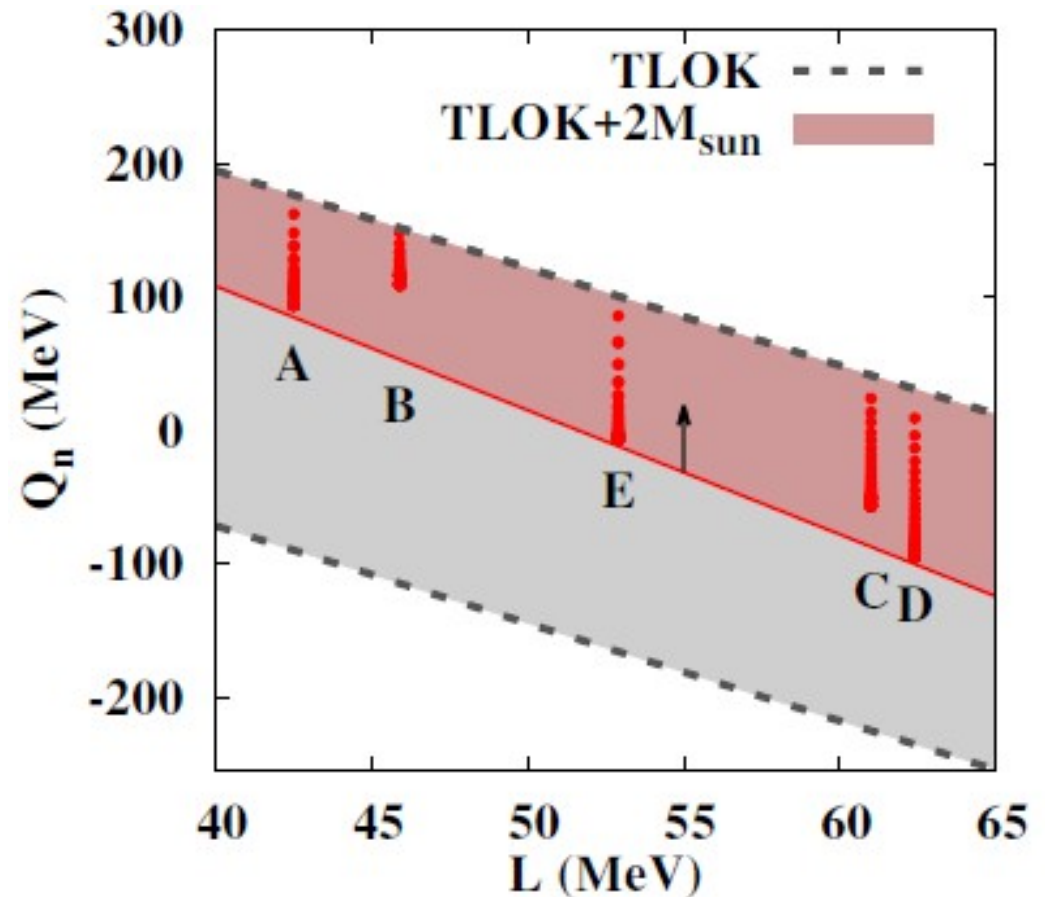


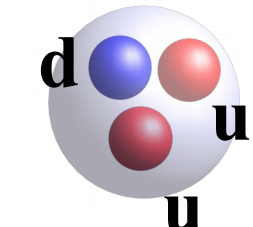
FIG. 4. Constraint on Q_n

AO, Kolomeitsev, Lattimer, Tews, Wu (OKLTW), in prog.

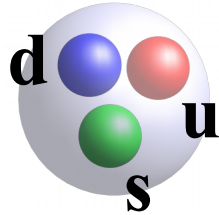
Neutron star – Is it made of neutrons ?

- Possibilities of various constituents in neutron star core

- Strange Hadrons

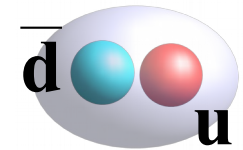


proton

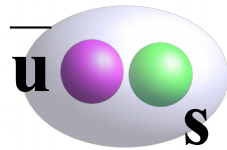


Λ hyperon

- Meson condensate (K, π)



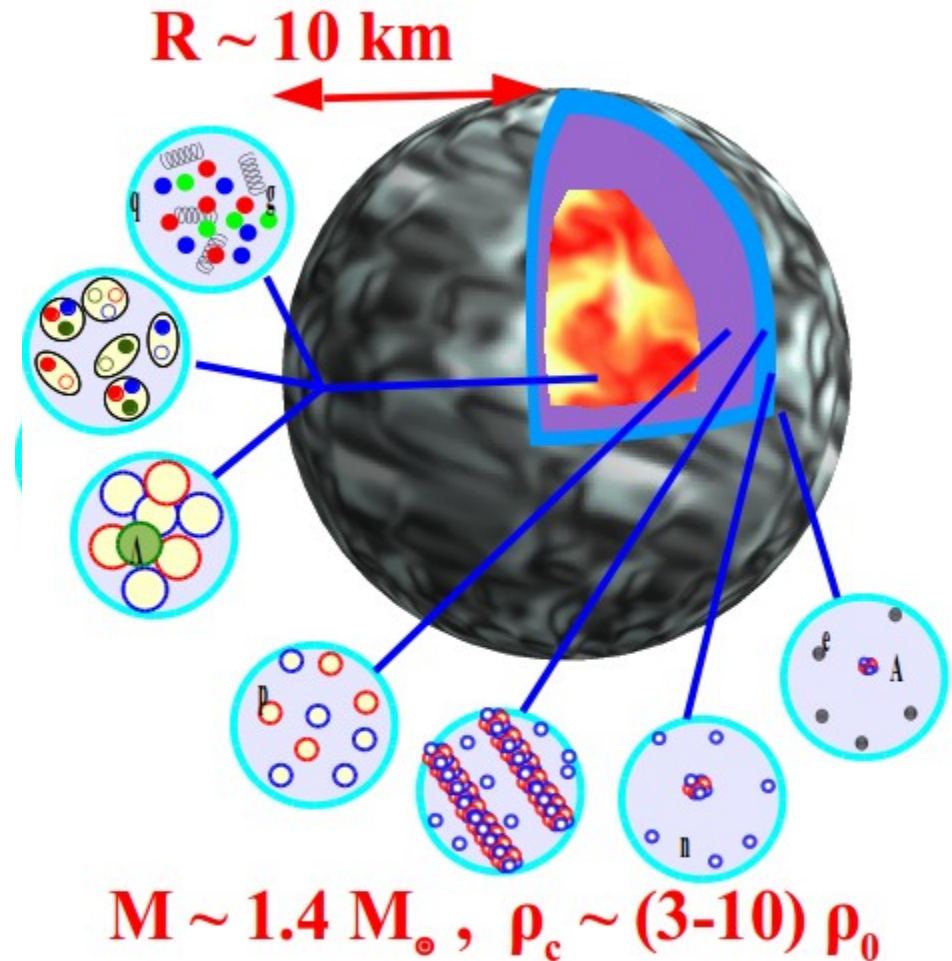
π



anti kaon

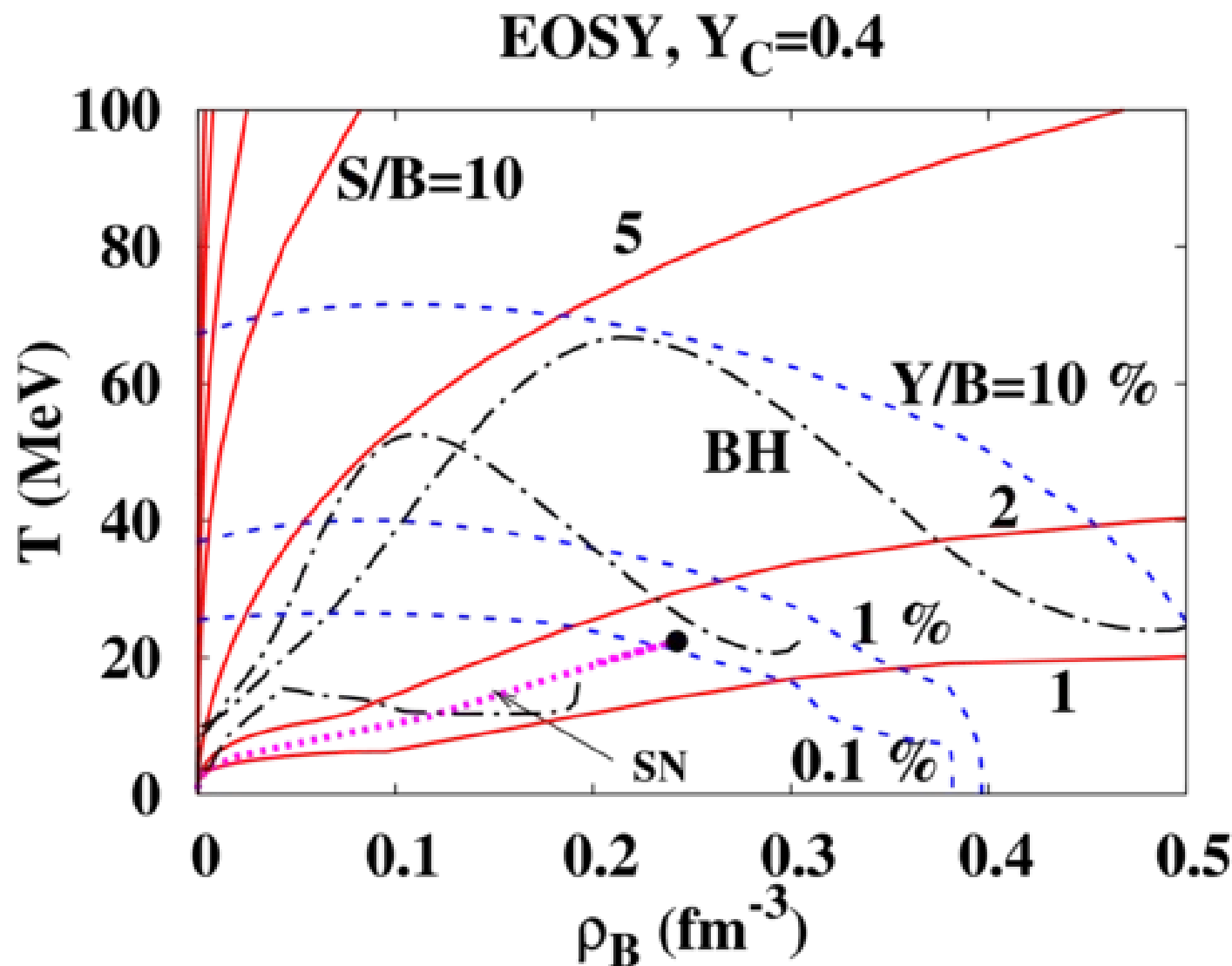
- Quark matter

- Quark pair condensate
(Color superconductor)



*NS core = Densest stable matter
existing in our universe.*

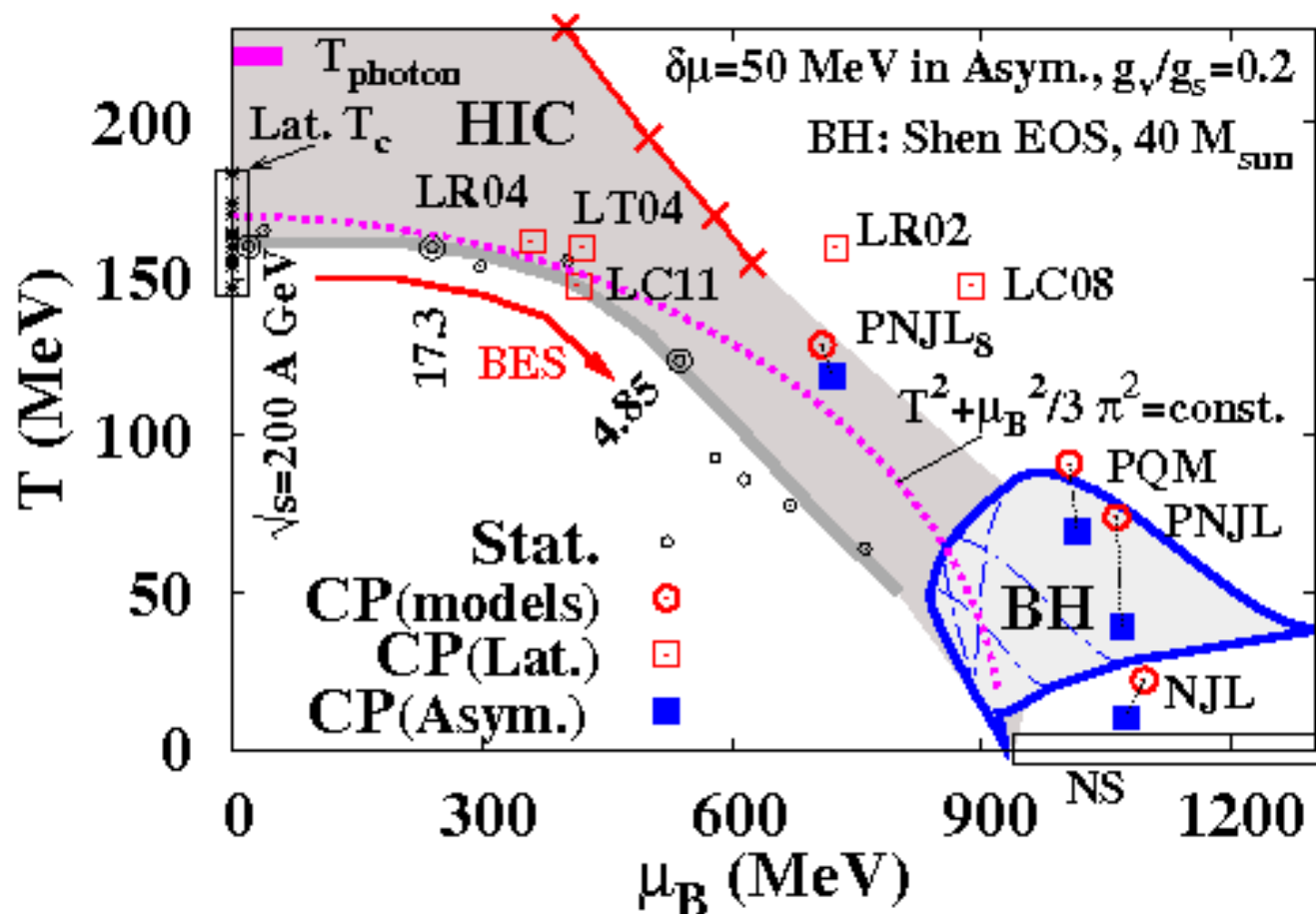
(ρ, T) during SN & BH formation



**Shen EOS
+ hyperons**

*Ishizuka, AO, Tsubakihara, Sumiyoshi, Yamada, JPG 35('08) 085201;
AO et al., NPA 835('10) 374.*

QCD phase diagram (Exp. & Theor. Studies)



QCD phase transition is not only an academic problem, but also a subject which would be measured in HIC or Compact Stars

Unitary Gas Constraint

Tews, Lattimer, AO, Kolomeitsev (TLOK), ApJ ('17)

■ Conjecture:

Unitary gas gives the lower bound of neutron matter energy.

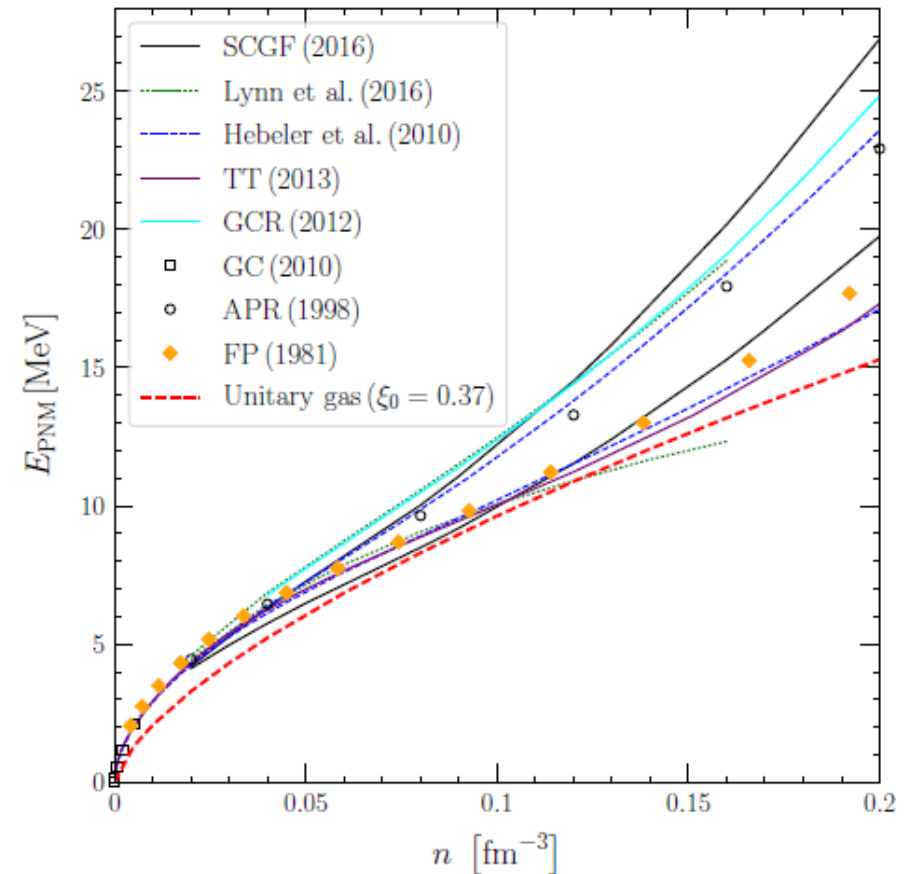
$$S(n) = E_{\text{PNM}} - E_{\text{SNM}} \geq E_{\text{UG}} - E_{\text{SNM}} \leftarrow \text{Sym. Nucl. Matter EOS is relatively well known.}$$

$$E_{\text{UG}} = \xi E_{\text{FG}} \quad (\xi \simeq 0.38)$$

■ $a_0 = \infty$ in unitary gas

→ lower bound energy
of $a_0 < 0$ systems
(w/o two-body b.s.) ?

■ Supported by (most of) ab initio calc.



Potential Energy Density

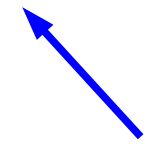
■ Potential Energy Density in the Fermi momentum expansion

$$\mathcal{V} = nV = \sum_{i,j \in B} n_i n_j v_{ij}(n)$$

Density-dependent NN interactions v_{ij} ($i, j=p$ or n) are known.

■ Single particle potential

$$\begin{aligned} U_i &= \frac{\partial \mathcal{V}}{\partial n_i} = \sum_j n_j v_{ij}(n) + \sum_{jk} n_j n_k \frac{\partial v_{jk}(n)}{\partial n_i} \\ &= U_{0i} + \frac{L_i}{3} (u - 1) + \mathcal{O}((u - 1)^2) \\ &\simeq au + bu^{4/3} \end{aligned}$$

 **rearrangement
term**

Again, a and b are given as a linear function of U_{0i} and L_i .