

---

# ***Quark matter phase diagram in the strong coupling region of lattice QCD***

***Akira Ohnishi, Noboru Kawamoto, Koutaro Miura  
Hokkaido University, Sapporo, Japan***

***Phase diagram at finite temperature and quark density in the strong coupling limit of  
lattice QCD for color SU(3)***

***N. Kawamoto, K. Miura, A. Ohnishi, T. Ohnuma,  
Phys. Rev. D 75(2007), 014502 [arXiv: hep-lat/0512023]***

***Strong Coupling Limit/Region of Lattice QCD***

***A. Ohnishi, N. Kawamoto, K. Miura, K. Tsubakihara, H. Maekawa,  
Prog. Theor. Phys. Suppl. (2007), to appear [arXiv:0704.2823]  
(Proc. of YKIS06.)***

***Phase diagram at finite temperature and quark density in the strong coupling region  
of lattice QCD for color SU(3)***

***A. Ohnishi, N. Kawamoto, K. Miura,  
J. Phys. G (2007), to appear [arXiv:hep-lat/0701024] (Proc. of QM2006).***



# Quark and Hadronic Matter Phase Diagram

- Dense quark & hadronic matter contains rich physics, but Lattice QCD simulation is not yet reliable.  
→ *Model/Approximate approaches are necessary!*

- Monte-Carlo calc. of Lattice QCD:

Improved ReWeighting Method (Fodor-Katz)

Taylor Expansion in  $\mu$

(Bielefeld-Swansea)

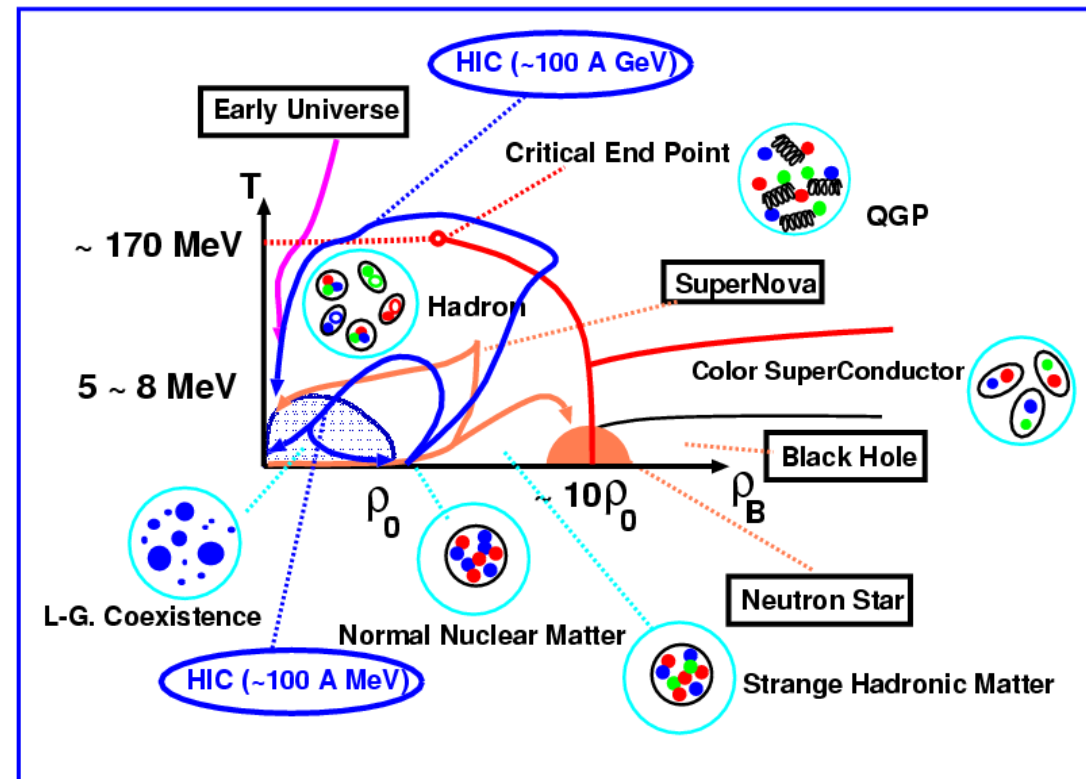
Analytic Continuation

(de Forcrand-Philipssen)

- Model / Phen. Approaches:

(P)NJL, QMC, RMF, ...

- *Strong Coupling Limit of Lattice QCD*



# Previous Works in Strong Coupling Limit LQCD

- Strong Coupling Limit Lattice QCD re-attracts interests  
*c.f. Nakamura @ JHF Symp. for high density matter (2001)*

Ref	$T$	$\mu$	$N_c$	Baryon	CSC	$N_f$
Damgaard-Kawamoto-Shigemoto('84)	Finite	0	$U(N_c)$	X	X	1
Damgaard-Hochberg-Kawamoto('85)	0	Finite	3	Yes	X	1
Bilic-Karsch-Redlich('92)	Finite	Finite	3	X	X	1 ~ 3
Azcoiti-Di Carlo-Galante-Laliena('03)	0	Finite	3	Yes	Yes	1
Nishida-Fukushima-Hatsuda('04)	Finite	Finite	2	Yes (*)	Yes (*)	1
Nishida('04)	Finite	Finite	3	X	X	1~2
<b>Kawamoto-Miura-AO-Ohnuma('07)</b>	<b>Finite</b>	<b>Finite</b>	<b>3</b>	<b>Yes</b>	<b>Yes (+)</b>	<b>1</b>

\*: bosonic baryon=diquark in  $SU(2)$

+: analytically included, but ignored in numerical calc.

- Baryon effects have been ignored in finite  $T$  treatments !***  
***→ This work: Baryonic effects at Finite  $T$  (and  $\mu$ ) for  $SU_c(3)$***

# Strong Coupling Limit Lattice QCD

## QCD Lattice Action

$$Z \simeq \int D[\chi, \bar{\chi}, U] \exp \left[ - \left( S_G + S_F^{(s)} + S_F^{(t)} + m_0 M \right) \right]$$

$$S_G = \frac{1}{g^2} \sum_{x\mu\nu} \left[ \text{Tr} U_{\mu\nu} + \text{Tr} U_{\mu\nu}^+ \right]$$

$$S_F^{(s)} = \frac{1}{2} \sum_{x,j} \eta_j(x) \left( \bar{\chi}_x U_j(x) \chi_{x+\hat{j}} - \bar{\chi}_{x+\hat{j}} U_j^+(x) \chi_x \right)$$

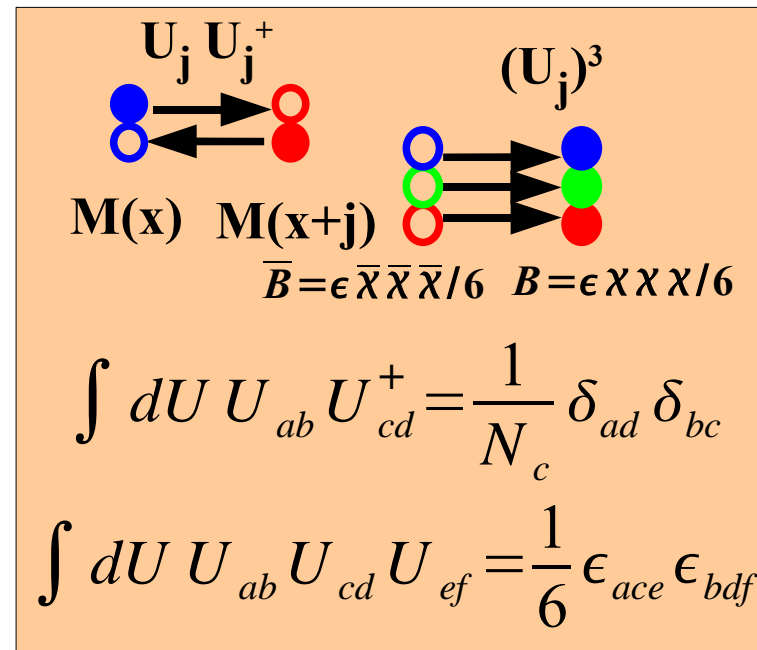
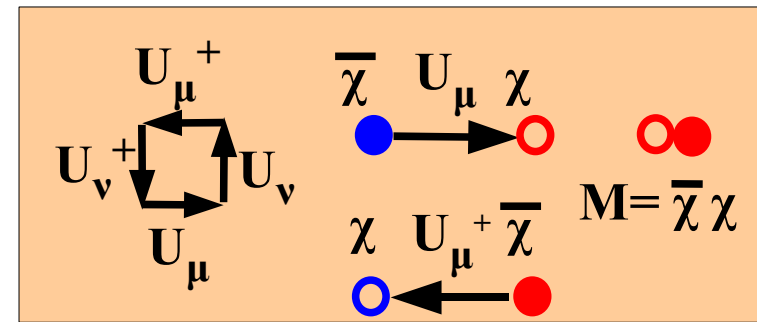
$$S_F^{(t)} = \frac{1}{2} \sum_x \left( e^\mu \bar{\chi}_x U_0(x) \chi_{x+\hat{0}} - e^{-\mu} \bar{\chi}_{x+\hat{0}} U_0^+(x) \chi_x \right)$$

## Strong Coupling Limit: $g \rightarrow \infty$

- We can ignore  $S_G$  and perform one-link integral after 1/d expansion.

$$S_F^{(s)} \rightarrow -\frac{1}{2} (M V_M M) - (\bar{B} V_B B)$$

$$= -\frac{1}{4 N_c} \sum_{x,j>0} M_x M_{x+\hat{j}} + \sum_{x,j>0} \frac{\eta_j}{8} \left[ \bar{B}_x B_{x+\hat{j}} - \bar{B}_{x+\hat{j}} B_x \right]$$



$$\int dU U_{ab} U_{cd}^+ = \frac{1}{N_c} \delta_{ad} \delta_{bc}$$

$$\int dU U_{ab} U_{cd} U_{ef} = \frac{1}{6} \epsilon_{ace} \epsilon_{bdf}$$

# SCL-LQCD w/o Baryons

*Damgaard-Kawamoto-Shigemoto 1984, Faldt-Petersson 1986, Bilic-Karsch-Redlich 1992, Nishida 2004, .....*

- **Lattice Action (staggered fermion) in SCL**

$$Z \simeq \int D[\chi, \bar{\chi}, U] \exp \left[ -S_F^{(s)} - S_F^{(t)} - m_0 \bar{\chi} \chi - S_G \right]$$

- **Spatial Link Integral**

$$\simeq \int D[\chi, \bar{\chi}, U_0] \exp \left[ \frac{1}{2} (M, V_M M) + (\bar{B}, V_B B) - (\bar{\chi} G_0 \chi) \right]$$

*Strong Coupling*

- **Bosonization (Hubburd-Stratonovich transformation)**

$$\simeq \int D[\chi, \bar{\chi}, U_0, \sigma] \exp \left[ -\frac{1}{2} (\sigma, V_M \sigma) - (\sigma, V_M M) - (\bar{\chi} G_0 \chi) \right]$$

*1/d Expansion (1/√d)*

- **Quark and U<sub>0</sub> Integral**

$$\simeq \exp \left( -N_s^3 N_\tau \left[ \frac{1}{2} a_\sigma \sigma^2 - T \log G_U(\sigma) \right] \right) = \exp \left( -N_s^3 F_{\text{eff}}/T \right)$$

$(\bar{\chi} G(\sigma) \chi)$

*Local Bi-linear action in quarks → Effective Free Energy*

# SCL-LQCD with Baryons

## Effective Action up to $O(1/\sqrt{d})$

$$M = \bar{\chi}_a \chi^a$$

$$B = \epsilon_{abc} \chi^a \chi^b \chi^c / 6$$

$$Z \simeq \int D[\chi, \bar{\chi}, U_0] \exp \left[ \frac{1}{2} (M, V_M M) + (\bar{B}, V_B B) - (\bar{\chi} G_0 \chi) \right]$$

$$= \int D[\chi, \bar{\chi}, U_0, b, \bar{b}] \exp \left[ \frac{1}{2} (M V_M M) - (\bar{b} V_B^{-1} b) + (\bar{b}, B) + (\bar{B}, b) - (\bar{\chi} G_0 \chi) \right]$$

## Decomposition of $bB$ by using diquark condensate (Azcoiti et al., 2004)

$$\exp[(\bar{b}, B) + (\bar{B}, b)] = \exp \left[ \frac{1}{6} (\bar{b}, \epsilon \chi \chi \chi) + \frac{1}{6} (\epsilon \bar{\chi} \bar{\chi} \bar{\chi}, b) \right]$$

$$= \int D[\phi_a, \phi_a^*] \exp \left[ -\phi^* \phi + \phi^* \left( \frac{\gamma}{2} \epsilon \chi \chi + \frac{\bar{\chi} b}{3\gamma} \right) + \phi \left( \frac{\gamma}{2} \epsilon \bar{\chi} \bar{\chi} + \frac{\bar{b} \chi}{3\gamma} \right) \right]$$

$$\times \exp(-\gamma M^2 / 2 + M \bar{b} b / 9 \gamma^2)$$

## Decomposition of $Mbb$ using baryon potential field $\omega$

$$\exp(M \bar{b} b / 9 \gamma^2) = \int D[\omega] \exp \left[ \frac{1}{2} \omega^2 - \omega \left( \alpha M + \frac{\bar{b} b}{9 \alpha \gamma^2} \right) - \frac{\alpha^2}{2} M^2 \right]$$

note:  $(\bar{b} b)^2 = 0$  with one species of staggered fermion !



# Effective Free Energy with Baryon Effects

## Effective Action in local bilinear form of quarks

$$S_F = -\frac{1}{2} (M \tilde{V}_M M) + \frac{1}{2} (\omega, \omega) + (\bar{b}, \tilde{V}_B^{-1} (g_\omega \omega) b) + \alpha (\omega, M) + (\bar{\chi} G_0 \chi)$$

*Bosonization + MFA*

*+No diquark cond.*

$$+(\phi^* \phi) + (\phi^* D) + (D^+ \phi)$$

$$= \frac{N_s^3 N_\tau}{2} (a_\sigma \sigma^2 + \omega^2) + (a_\sigma \sigma + \alpha \omega, M) + (\bar{\chi} G_0 \chi) + (\bar{b}, \tilde{V}_B^{-1} (g_\omega \omega) b)$$

*quark & gluon int.*

*b int.*

$$F_{\text{eff}}(\sigma, \omega) = \frac{1}{2} a_\sigma \sigma^2 + \frac{1}{2} \omega^2 + F_{\text{eff}}^{(q)}(a_\sigma \sigma + \alpha \omega) + F_{\text{eff}}^{(b)}(g_\omega \omega)$$

$$= \frac{1}{2} a_\sigma \sigma^2 + \frac{1}{2} a_\omega \omega^2 + F_{\text{eff}}^{(q)}(a_\sigma \sigma + \alpha \omega) + \Delta F_{\text{eff}}^{(b)}(g_\omega \omega)$$

*Linear Approx. ( $\omega \sim \alpha \sigma / a_\omega$ )*

$$F_{\text{eff}}(\sigma) = \frac{1}{2} b_\sigma \sigma^2 + F_{\text{eff}}^{(q)}(b_\sigma \sigma) + \Delta F_{\text{eff}}^{(b)}(g_\sigma \sigma)$$



# Effective Free Energy with Baryon Effects

(Kawamoto-Miura-AO-Ohnuma, hep-lat/0512023)

$$F_{\text{eff}}(\sigma) = \frac{1}{2} b_{\sigma} \sigma^2 + F_{\text{eff}}^{(q)}(b_{\sigma} \sigma; T, \mu) + \Delta F_{\text{eff}}^{(b)}(g_{\sigma} \sigma)$$

is analytically derived based on many previous works, including

- **Strong Coupling Limit** (Kawamoto-Smit, 1981)
- **1/d expansion** (Kluberg-Stern-Morel-Petersson, 1983)
- **Lattice chemical potential** (Hasenfratz-Karsch, 1983)
- **Quark and time-like gluon analytic integral**  
(Damgaard-Kawamoto-Shigemoto, 1984, Faldt-Petersson, 1986)

$$F_{\text{eff}}^{(q)}(\sigma; T, \mu) = -T \log \left( C_{\sigma}^3 - \frac{1}{2} C_{\sigma} + \frac{1}{4} C_{3\mu} \right) \quad C_{\sigma} = \cosh(\sinh^{-1} \sigma/T) \quad C_{3\mu} = \cosh(3\mu/T)$$

- **Decomposition of baryon-3 quark coupling**  
(Azcoiti-Di Carlo-Galante-Laliena, 2003)

and auxiliary baryon potential and baryon integral



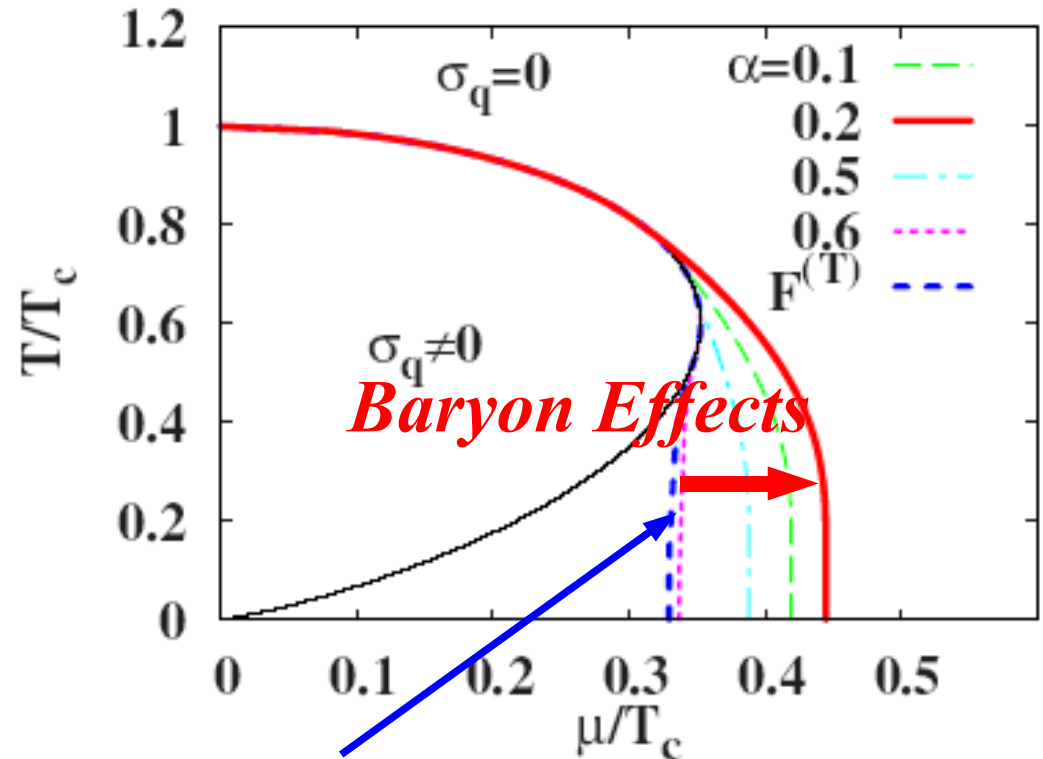
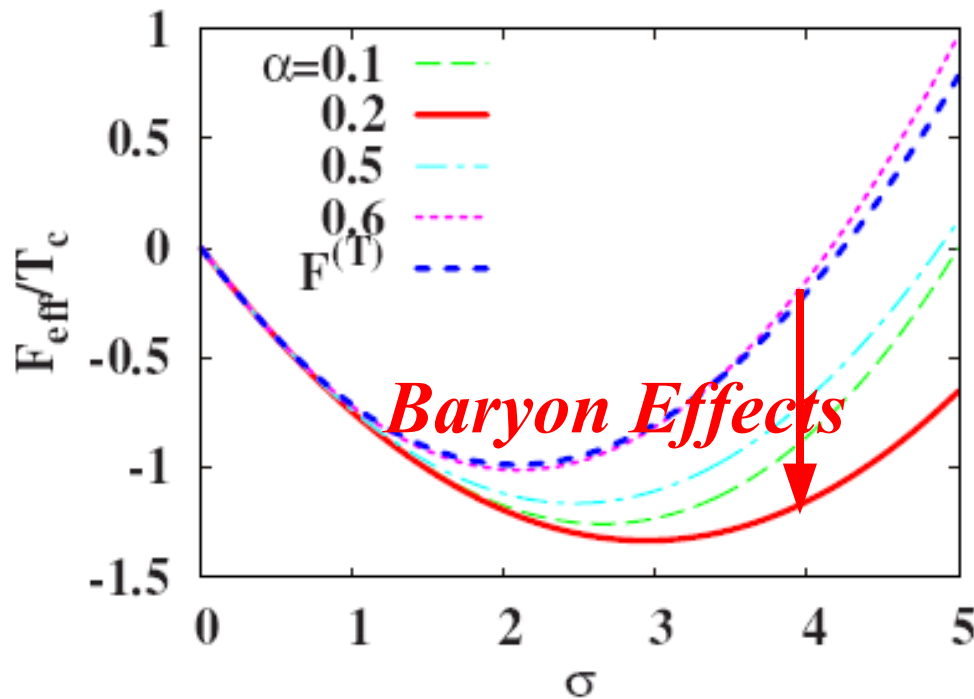
# Phase diagram in *SCL-LQCD* with Baryons

(Kawamoto-Miura-AO-Ohnuma, hep-lat/0512023)

## ■ Baryon effects on phase diagram

### ● Energy gain in larger condensates

→ Extension of hadron phase to larger  $\mu$  by around 30 %.



*Nishida 2004(No B)*



# Discussions

- Present phase diagram  $\leftrightarrow$  real phase diagram
  - One species of staggered fermion  $\sim N_f=4$ . Should be 1st order !
  - $T_c$  seems to be too high.  $\mu_c/T_c(\text{present}) \sim 0.45 \leftrightarrow \mu_c/T_c(\text{real}) \sim (2-3)$
  - No stable CSC phase (*Azcoiti et al., 2003*)  
 $\leftrightarrow$  Stable CSC phase at large  $\mu$  (*Alford, Hands, Stephanov*)
- Two parameters are introduced through identities (HS transf.)
  - The results should be independent from parameter choice !  
 $\rightarrow$  MFA may break the identity...
  - How should we fix these parameters ?
- Is SCL-LQCD useful ?  $\rightarrow$  We would like to answer “Yes” !
  - Chiral RMF derived in SCL-LQCD works well in Nuclear Physics (Tsubakihara, AO, nucl-th/0607046  
Tsubakihara, Maekawa, AO, Proc. of HYP06, to appear)
  - $1/g^2$  expansion may connect SCL-LQCD and real world.

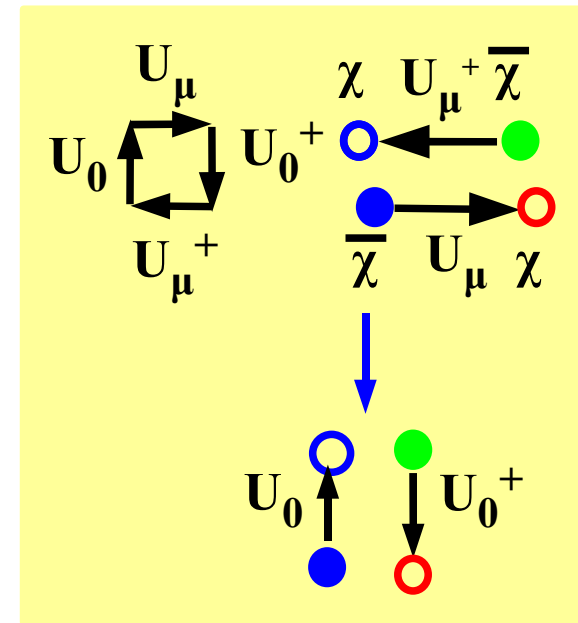


# *1/g<sup>2</sup> expansion (w/o Baryon Effects)*

## ■ T<sub>c</sub> (μ=0) and μ<sub>c</sub> (T=0): Which is worse ?

- 1/g<sup>2</sup> correction reduces T<sub>c</sub>. (*Bilic-Cleymans 1995*)
- Hadron masses are well explained in SCL. (*Kawamoto-Smit 1981, Kawamoto-Shigemoto 1982*)

→ We expect T<sub>c</sub> reduction with 1/g<sup>2</sup> correction !



## ■ 1/d expansion of plaquettes (*Falgt-Petersson 1986*)

### ● Space-like plaquette

$$\exp \left[ \frac{1}{g^2} \sum_{x, i > j > 0} \text{Tr} U_{ij}(x) \right] \rightarrow \exp \left[ \frac{1}{8 N_c^4 g^2} \sum_{x, k > j > 0} M_x M_{x+\hat{j}} M_{x+\hat{k}} M_{x+\hat{k}+\hat{j}} \right]$$

### ● Time-like plaquette

$$\exp \left[ \frac{1}{g^2} \sum_{x, j > 0} \text{Tr} U_{0j}(x) \right] \rightarrow \exp \left[ -\frac{1}{4 N_c^2 g^2} \sum_{x, j > 0} \left( V_x V_{x+\hat{j}}^+ + V_x^+ V_{x+\hat{j}} \right) \right]$$

$$(V_x = \bar{\chi}_x U_0(x) \chi_{x+\hat{0}})$$



# Plaquette Bosonization

- **Bosonization of Plaquettes** ( $O(1/d, 1/g^4)$  and  $\text{Im}(V)$  are ignored) + **MFA**

$$\begin{aligned}
 \exp(-S_F - S_g) &\rightarrow \exp \left[ \frac{1}{2} \sum_x (e^\mu V_x - e^{-\mu} V_x^+) + \frac{1}{4N_c} \sum_{x, j>0} M_x M_{x+\hat{j}} - m_0 \sum_x M_x \right] \\
 &\times \exp \left[ \frac{\beta_t}{2} \varphi_t \sum_x (V_x - V_x^+) + \beta_s \varphi_s \sum_{x, j>0} M_x M_{x+\hat{j}} \right] \\
 &\times \exp \left[ -L^3 N_\tau \left( \frac{\beta_t}{4} \varphi_t^2 + \frac{\beta_s d}{4} \varphi_s^2 \right) \right] \left( \beta_t = \frac{d}{2N_c^2 g^2}, \quad \beta_s = \frac{d-1}{8N_c^4 g^2} \right) \\
 &= \exp \left[ -\frac{L^3}{T} F_\varphi \left[ \frac{\alpha}{2} \sum_x (e^{\tilde{\mu}} V_x - e^{-\tilde{\mu}} V_x^+) + \frac{1}{2} \sum_{x, y} M_x \tilde{V}_M(x, y) M_y \right] \right] \\
 &\alpha = 1 + \beta_t \varphi_t \cosh \mu, \quad \tilde{\mu} = \mu - \beta_t \varphi_t \sinh \mu \\
 &\langle \varphi_t \rangle = \langle V^+ - V \rangle, \quad \langle \varphi_s \rangle = 2 \langle M_x M_{x+\hat{j}} \rangle
 \end{aligned}$$

*Time-like plaquettes modifies effective chemical potential*



# Effective Free Energy with $1/g^2$ Correction (w/o $B$ )

- After Quark and Time-like Link integral, we get  $F$  as

$$F = \frac{d}{4N_c} \sigma^2 (1 + 4N_c \beta_s \varphi_s) + \frac{\beta_t}{4} \varphi_t^2 + \frac{\beta_s d}{4} \varphi_s^2 - N_c \beta_t \varphi_t \cosh \mu + F_q(m_q; \tilde{\mu})$$

$$= \frac{d}{4N_c} \sigma^2 + 3d \beta_s \sigma^4 + \frac{\beta_t}{4} \tilde{\varphi}_t^2 + F_q(m_q; \tilde{\mu})$$

$$\varphi_s = 2\sigma^2, \quad \varphi_t = \tilde{\varphi}_t + 2N_c \cosh \mu \quad \leftarrow \text{Time-like plaquettes remains finite at large } \mu$$

$$m_q = \frac{d}{2N_c} \sigma (1 + 4N_c \beta_s \varphi_s - \beta_t \varphi_t \cosh \mu)$$

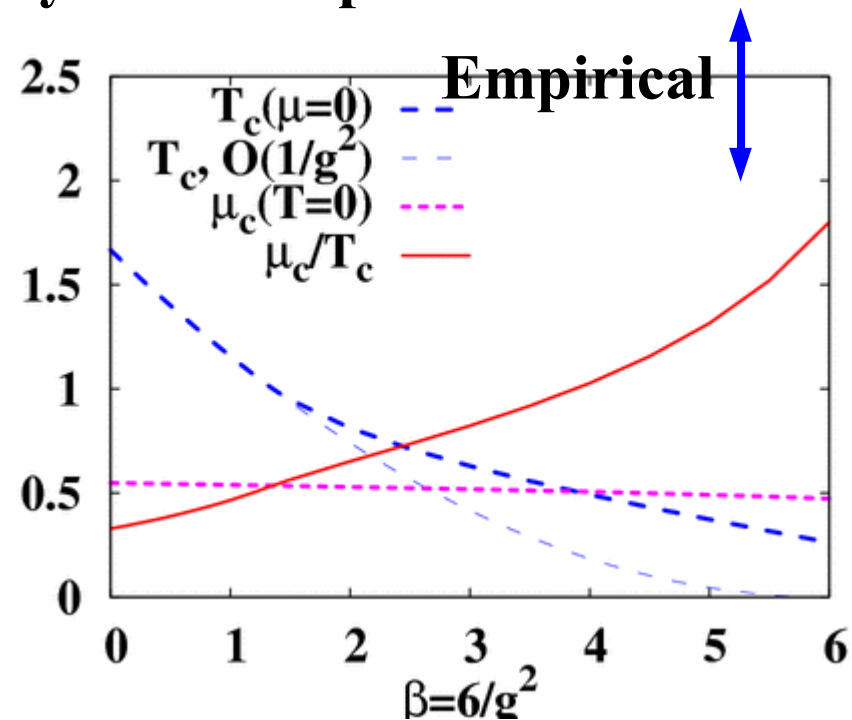
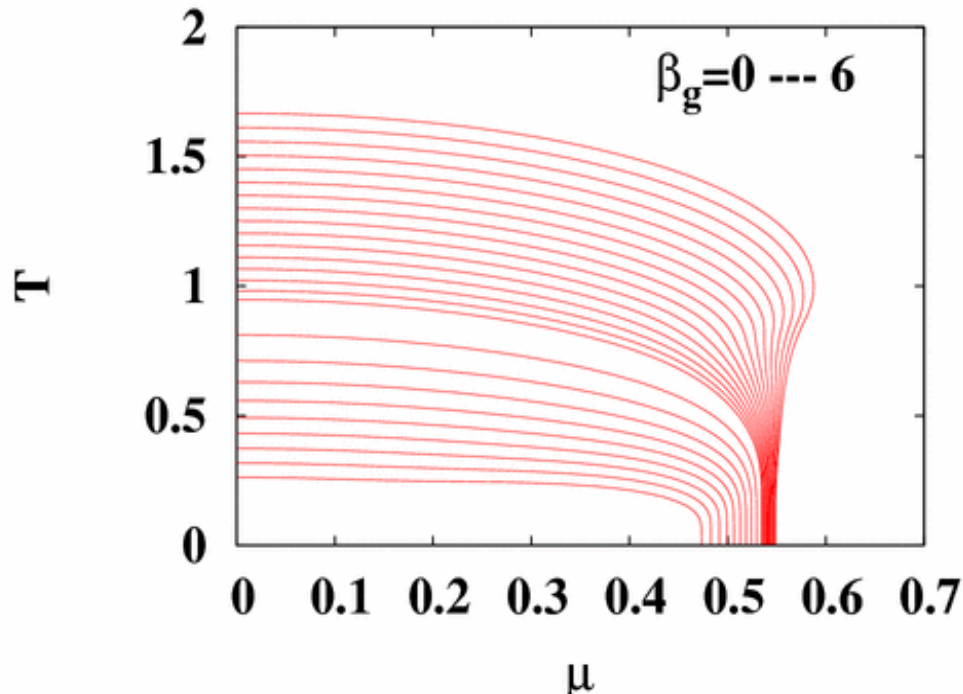
$$= \frac{d}{2N_c} \sigma (1 - 2N_c \beta_t \cosh^2 \mu + 8N_c \beta_s \sigma^4 - \beta_t \tilde{\varphi}_t \cosh \mu)$$

$$\tilde{\mu} = \mu - \beta_t \varphi_t \sinh \mu = \mu - 2N_c \beta_t \cosh \mu \sinh \mu - \beta_t \tilde{\varphi}_t \sinh \mu$$

- Space-like plaquette  $\rightarrow$  Repulsive pot.  $\propto \sigma^4$ , Enh.  $\sigma$ -quark coupling
- Time-like plaquette  $\rightarrow$  Reduces  $\mu$  and  $\sigma$ -quark coupling  
( $\varphi_t$  has to be determined to minimize  $F_{\text{eff}}$ )

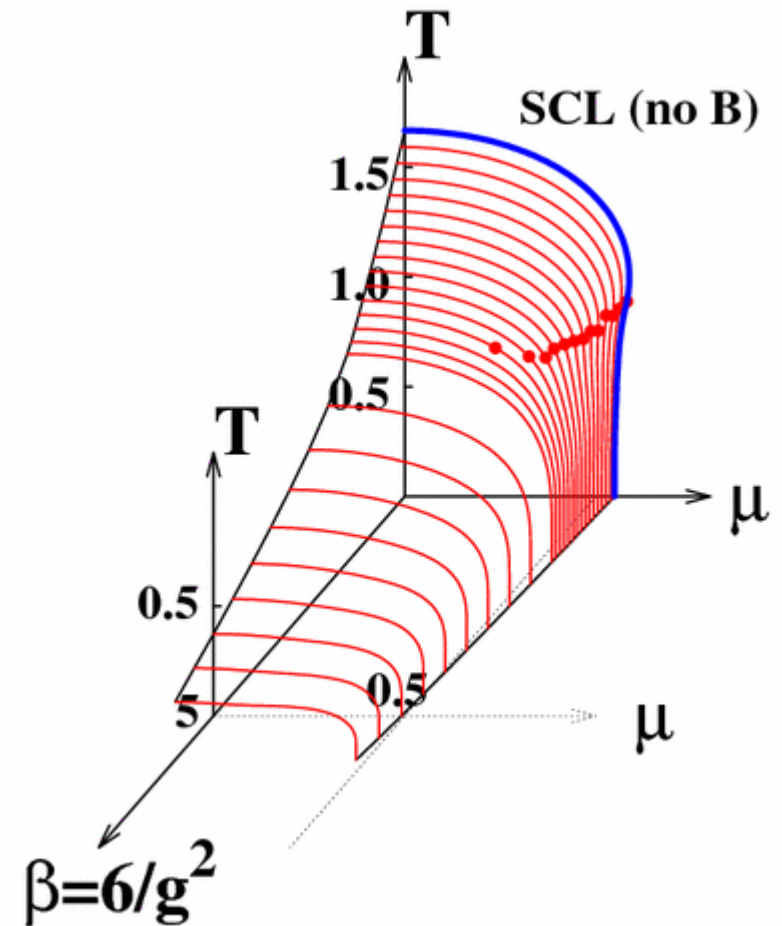
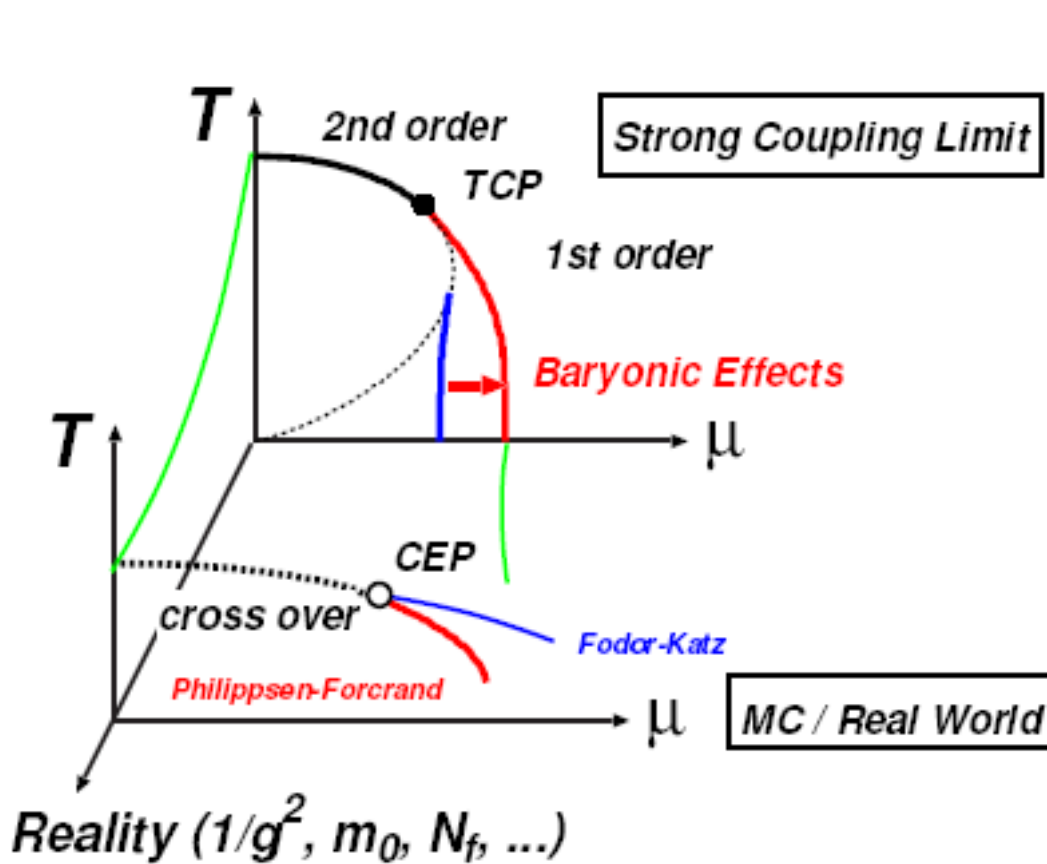
# Phase Boundary with $1/g^2$ correction

- Rapid decrease of  $T_c(\mu=0)$ , and slow decrease of  $\mu_c(T=0)$ .
  - Similar reduction of  $\sigma$ -quark coupling and effective  $\mu$  at small condensate  $\rightarrow$  can be mimicked by the scaling of  $T$  (c.f. Bilic-Claymans 1995 ( $T_c$  goes down), Arai-Yoshinaga (Poster, goes up)).
- Ratio  $\mu_c/T_c \sim 1.8$  @  $g=1$ .
  - with baryonic effects ( $\sim 30\%$ ), it may reach empirical value.



# Evolution of Phase Diagram

- “Reality” Axis:  $1/g^2$ ,  $n_f$ ,  $m_0$ , .... would enhance  $\mu_c/T_c$  ratio
- Example:  $1/g^2$  correction enhances  $\mu_c/T_c$  by a factor  $\sim(2-3)$ .



# Summary

- We obtain an analytical expression of effective free energy *at finite  $T$  and finite  $\mu$*  with *baryonic composite* effects in the strong coupling limit of lattice QCD for color SU(3).
  - *MFA, QG integral,  $1/d$  expansion (NLO,  $O(1/\sqrt{d})$ ), bosonization with diquarks and baryon potential field using  $(\bar{b}b)^2=0$ , Linear approx., zero diquark cond.(Color Angle Average), variational parameter choice*
- Baryonic action is found to result in *Free Energy Gain* and *Extension of Hadron Phase to Larger  $\mu$*  by around 30 %.
  - *Problem: Too small  $\mu_c/T_c$  in the Strong Coupling Limit.*
- Strong Coupling Limit is useful to understand Dense Matter
  - *SCL gives a qualitative insight.*
  - *$1/g^2$  correction seems to work well (Do not believe us yet ...)*
  - *Application to chiral RMF (K. Tsubakihara, AO, nucl-th/0607046)*  
→ *3rd week Poster by Tsubakihara*

