

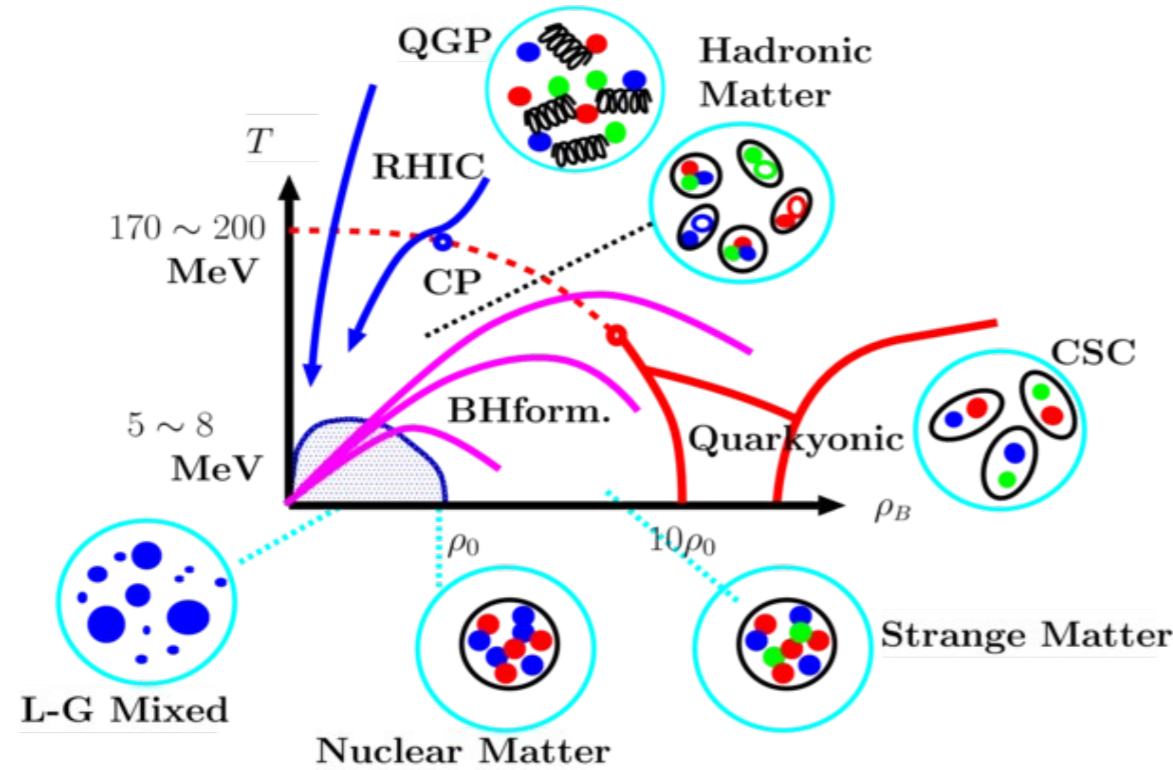
# Chiral and deconfinement transitions in strong coupling lattice QCD with finite coupling and Polyakov loop effects \*

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## Introduction

- QCD phase diagram is of great interest !
  - Transition at high T → Lattice MC, RHIC, LHC
  - High  $\mu$  transition has rich physics
    - Various phases, CEP, Astrophysical applications, ...

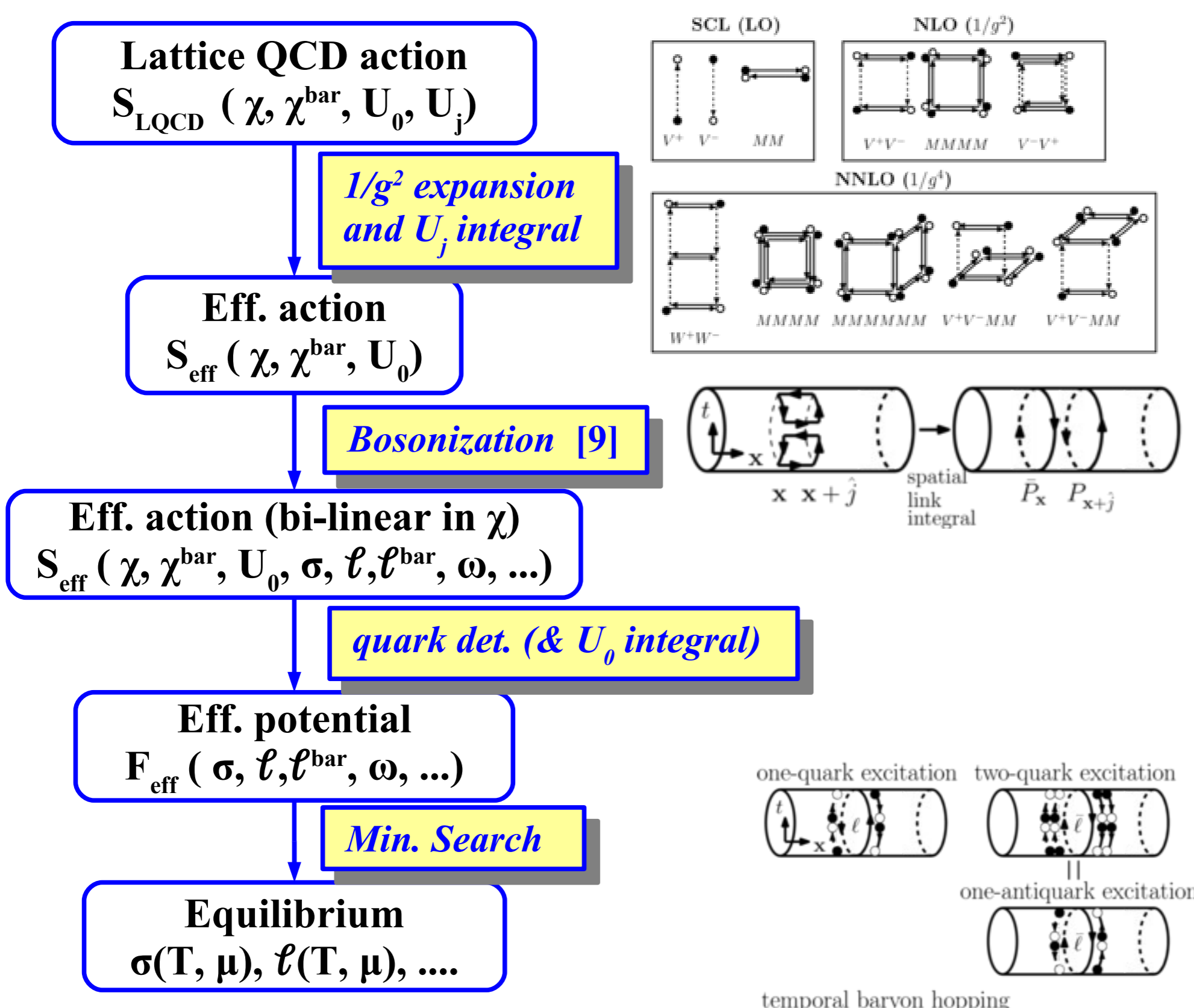


- Approaches to Dense QCD Matter
  - Perturbative QCD → Applicable only to high T or  $\mu$
  - Lattice QCD → Sign Problem
  - Eff. Models (NJL [1], PNJL [2], PQM [3], ...)
    - Model & Par. deps.
  - Strong Coupling Lattice QCD**
    - Strong coupling expansion &  $U_j$  integral
      - Sign problem is weak and applicable to large  $\mu$  region !
      - Successes in pure Yang-Mills [4], Chiral transitions [5], Deconf. transitions [6,7].
- No work on self-consistent description of chiral and deconf. transitions with finite coupling effects based on SC-LQCD.

*We develop SC-LQCD framework including both Finite Coupling & Polyakov loop effects Polyakov loop extended SC-LQCD (P-SC-LQCD), and discuss chiral and deconf. transitions.*

## Polyakov loop extended Strong Coupling Lattice QCD

- Lattice action → Eff. potential
  - Strong coupling ( $1/g^2$ ) expansion (NNLO in with quarks [8], LO in Polyakov loop [6,7])
  - Large dimensional ( $1/d$ ) expansion (LO)
  - Mean Field Approx. (static and constant aux. fields)



- Eff. Potential (in Haar measure method)
 
$$F_{\text{eff}} = F_q(\Phi; \mu, T) + U_g(\ell, \bar{\ell}; T) + F_{\text{eff}}^{(X)}(\Phi)$$

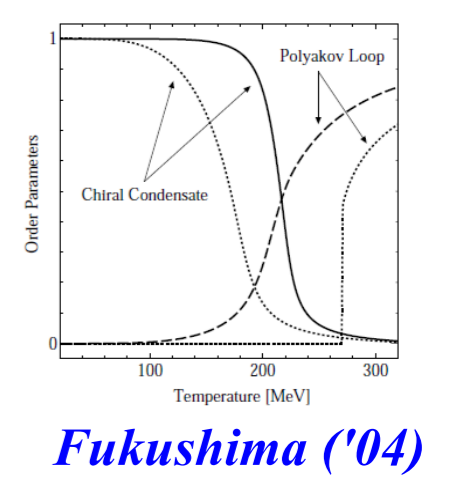
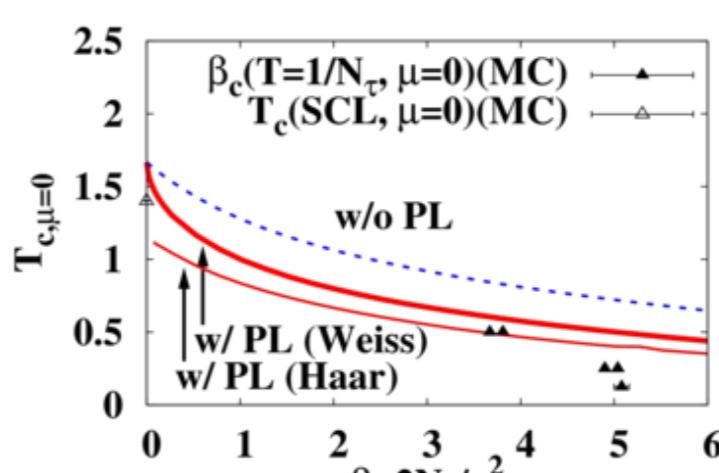
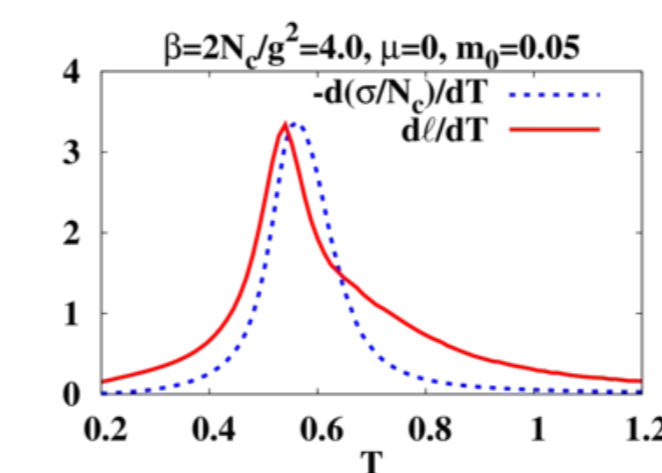
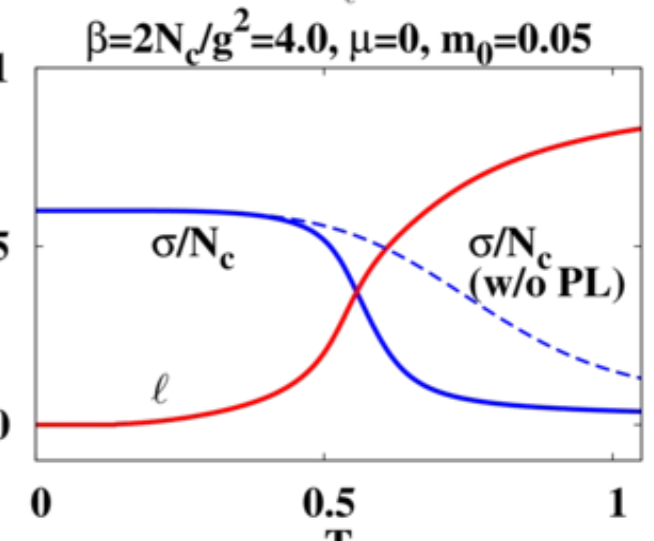
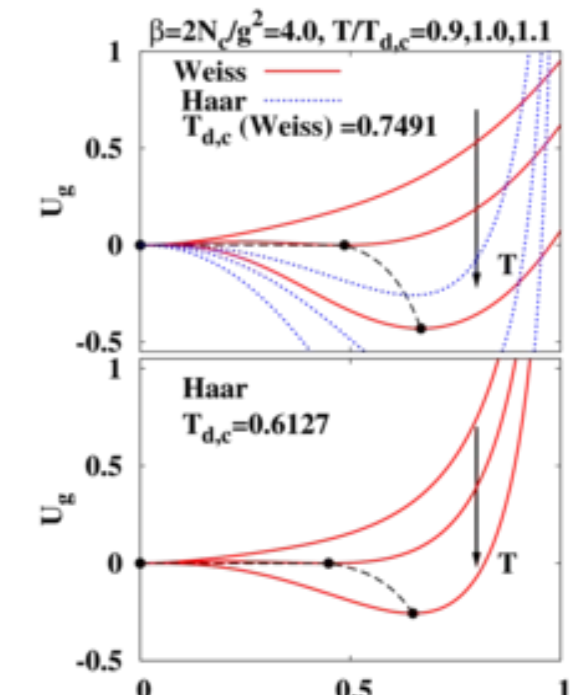
$$F_q = -N_c E_q - T \log R(E_q - \tilde{\mu}, N_c \bar{\ell}, N_c \bar{\ell}) - T \log R(E_q + \tilde{\mu}, N_c \bar{\ell}, N_c \bar{\ell})$$

$$R(x, L, \bar{L}) = 1 + L \exp(-x/T) + \bar{L} \exp(-2x/T) + \exp(-3x/T)$$

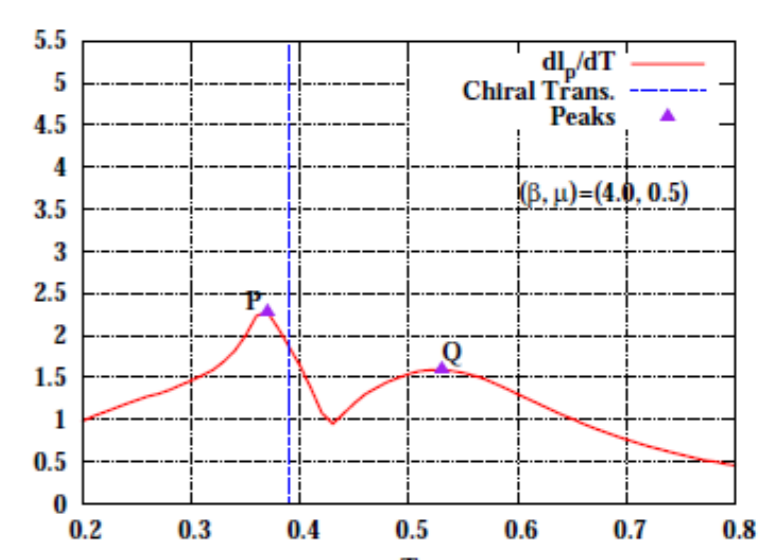
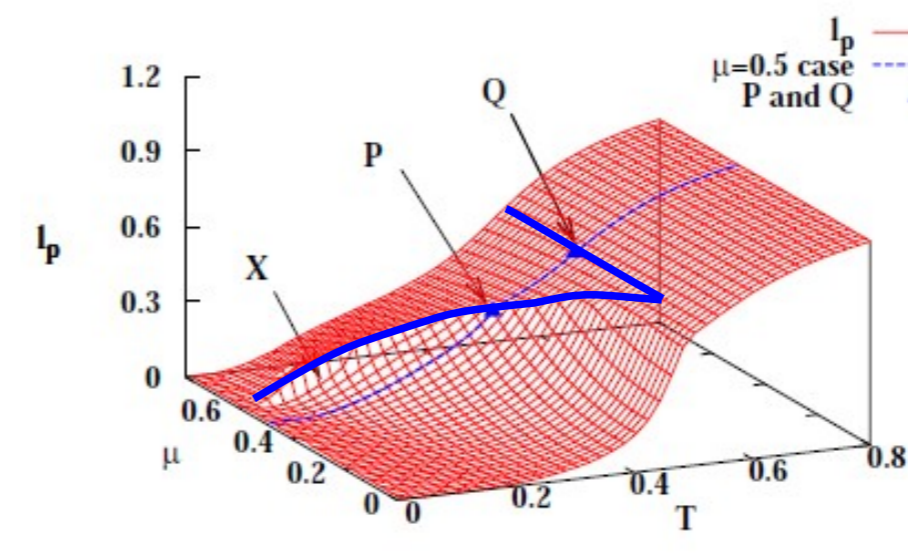
$$U_g = -2T \beta_p \bar{\ell} l - T \log [1 - 6 \bar{\ell} l + 4(l^3 + \bar{\ell}^3) - 3(\bar{\ell} l)^3]$$

## Chiral and Deconf. Transitions in P-SC-LQCD

- Effective Potential of the Polyakov loop
  - Haar measure method [6]
    - $\log(\text{Haar measure})$  as the potential (No  $U_0$  integral)
  - Bosonization method [7]
    - Perform  $U_0$  integral after bosonization (Weiss mean field approximation)
- Chiral cond. and Polyakov loop ( $\mu=0$ )
  - Chiral and Deconf. transition correlate !
  - SC-LQCD w/o PL: quarks are confined.**
    - PL promote quarks to deconfine ! (cf. Quarks are not confined in NJL
      - PL confines quarks in PNJL [2].)
  - $T_c$  is suppressed by PL
- Transition Temperature at  $\mu=0$ 
  - $-\text{d}\sigma/\text{d}T, \text{d}l/\text{d}T$  are peaked at similar T's.
  - Improved from SC-LQCD w/o PL.
  - Quantitatively, not bad for  $\beta < 4$  in  $T_c(\beta_c)$  [10]



- Phase diagram [11]
  - Separated chiral & deconf. boundaries at finite  $\mu$ 
    - Quarkyonic matter [12] would be measurable from two peak structure of  $\text{d}l/\text{d}T$  !



## Summary

- SC-LQCD framework is extended to include finite coupling and Polyakov loop effects (P-SC-LQCD).
  - Two order parameters (chiral cond. & Polyakov loop) are included on the same footing based on QCD.
  - Mixed strong coupling expansion orders.
    - Double expansion in fermionic and pure YM parts.
- Chiral & Deconf. transitions are discussed in P-SC-LQCD.
  - Chiral transition temperature is suppressed by Polyakov loop in P-SC-LQCD.
  - MC results on chiral transition T are roughly reproduced for  $\beta_g = 2N_c/g^2 \leq 4$ .
  - The existence of Quarkyonic matter is supported, and may be measurable via the behavior of the Polyakov loop.
- Future works
  - Improvement:  $1/g^2, 1/d$ , mean field, staggered fermions.
  - Model construction based on P-SC-LQCD.

[1] Y. Nambu, G. Jona-Lasinio, PR122('61)345; T.Hatsuda, T. Kunihiro, PRep247('94) 221. [2] K.Fukushima, PLB591('04)277; S.Roessner, C.Ratti, W.Weise, PRD75('07) 034007. [3] B.J. Schaefer, J.M.Pawlowski, J.Wambach, PRD76('07)074023; V.Skokov, B.Friman, E. Nakano, K.Redlich, B.J. Schaefer, PRD82('10)034029. [4] K.G.Wilson, PRD10('74)2445; G. Munster, NPB180('81),23. [5] N. Kawamoto, J. Smit, NPB192('81) 100; P.H.Damgaard, N. Kawamoto, K. Shigemoto, PRL53('84)2211. [6] E.M.Ilggenfritz, J. Kripfganz, ZPC29('85) 79; A.Gocksch, M.Ogilvie, PRD31('85)877. [7] J.B.Kogut, M.Snow, M.Stone, NPB200('82) 2311. [8] T.Z.Nakano, K.Miura, AO, PTP123('10)825. [9] K.Miura, T.Z.Nakano, AO, N.Kawamoto, PRD80('09)074034; K.Miura, T.Z.Nakano, AO, PTP122('09)1045. [10] Ph. de Forcrand, M. Fromm, PRL104('10)112005; S.A. Gottlieb et al. PRD35('87)3972; M. D'Elia, M.P.Lombardo, PRD67('03)014505; Z.Fodor, S.D.Katz, JHEP03('02)014; R. Gavai et al.[MT(c) Collab.], PLB241('90)567. [11] K. Miura, T.Z.Nakano, AO, N.Kawamoto, PoS(Lat10)202. [12] L.McLerran, R.D.Pisarski, NPA796('07)83; Y.Hidaka, L.McLerran, R.D.Pisarski, NPA808('08)117.

\* T.Z.Nakano, K. Miura, A. Ohnishi, [arXiv:1009.1518]