

核物質の状態方程式とコンパクト天体現象

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東京理科大セミナー 2011年12月9日

■Introduction

- Neutron Stars, Supernovae, and Black Hole formation
- Equation of State for Dense Baryonic Matter
- Relativistic Mean Field based EOS
- Recent development in dense matter EOS and QCD phase diagram
- Three-body force in RMF and $1.97 M_{\odot}$ neutron star
- Critical point sweep during black hole formation

■Summary



核物質の状態方程式とコンパクト天体現象

Abstract:核物質の状態方程式は、原子核のバルクな性質を説明する基本的な概念であるとともに、コンパクト天体現象を記述する上で欠かせないものである。例えば、高密度での斥力がなければ精度よく測定されている中性子星の質量(1.44Msun)を支えることは不可能である。コアの重力崩壊から始まる爆発天体現象においても状態方程式は重要である。近年の多次元シミュレーションでは超新星爆発を起こすことに成功するという大きな成果が報告されているが、こうした計算で用いられている状態方程式は一般にソフトなものであり、最近発見された重い中性子星(1.97Msun)とのconsistencyが問われる。

このセミナーでは、これまでに超新星爆発シミュレーションで用いられている核物質の状態方程式と高密度物質におけるエキゾチックな自由度(主としてハイペロン)について概観したのち、ブラックホール形成時のQCD臨界点探索、相対論的平均場理論(RMF)における3体力等、講演者らの最近の研究について報告する。

石塚さんからのメール:

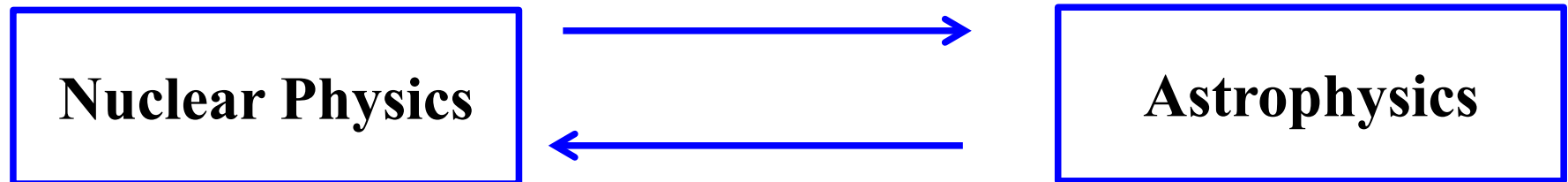
鈴木先生からのリクエストでセミナーの内容は天体核(原子核?)&宇宙物理の境界領域のバックグラウンド的なことをざっくり話しつつ(若い人向け

に)最新のお仕事の話題を組み込んでいただけると助かるよね。とのこと
Ohnishi @ IUS seminar, Dec. 9, 2011, Tokyo Univ. of Science, Noda, Japan 2

原子核物理と宇宙物理

- Nuclear physicists think that
“Astrophysics is a laboratory of nuclear physics”.
- Astrophysicists think that
“Nuclear physics gives us inputs of Astrophysics”.
(~ Kajino's closing at OMEG.)
- Both of them are true, and hopefully we can keep win-win relation.

Microphysics Inputs
(EOS, reaction rate,
 νA , e capture,)



Macrophysics Environment
(T , ρ , Element abundance,)

Compact Astrophysical Phenomena

コンパクト天体現象

■中性子星

- Cold -dense ($\sim 5 \rho_0$) matter (static, v-less)
- Many new forms of matter have been proposed !

■超新星

- Warm ($T \sim 20$ MeV), dense ($\sim 1.6 \rho_0$), dynamical, non-eq. v
- Important site of nucleosynthesis

■ブラックホール形成過程

- Hot ($T \sim 90$ MeV), dense ($\sim 5 \rho_0$), dynamical, non-eq. v
- QCD critical point may be reached

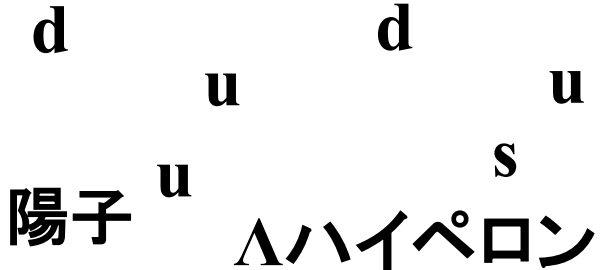
■BH-BH NS-BH NS-NS 融合 → 数値相対論

Nuclear matter at various densities and temperatures is realized in nature !

中性子星コアの状態

■コア領域では様々な可能性

●ストレンジクォークを含むバリオン(ハイペロン)を含む物質

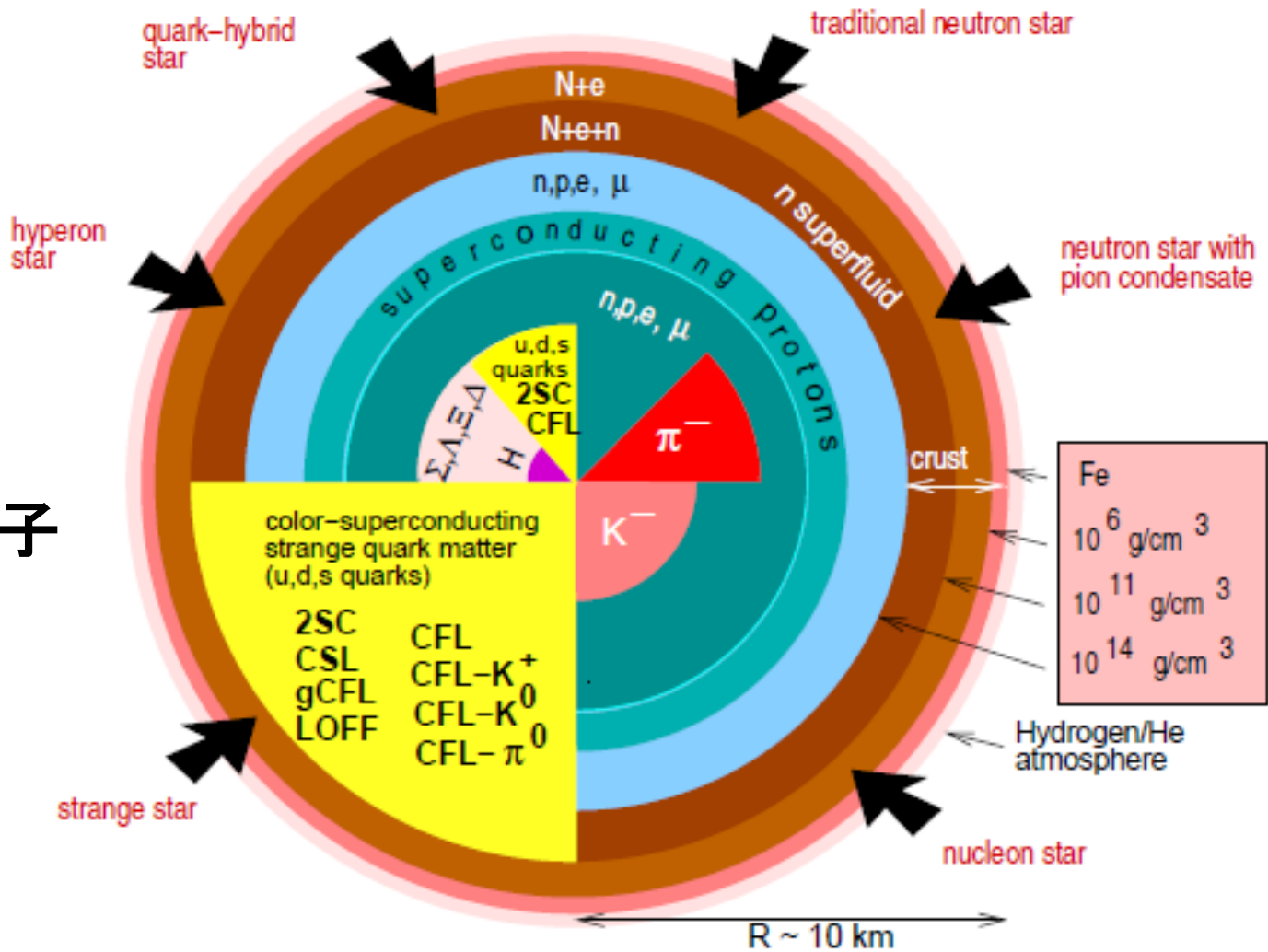


●中間子凝縮(K, π)

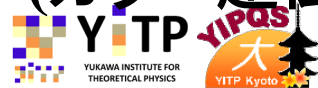
π 中間子 (反) K 中間子

●クォーク物質

●クォーク対凝縮状態 (カラー超伝導)



F. Weber, Prog. Part. Nucl. Phys. 54 (2005) 193



重い中性子星ショック....

■2010年のビッグニュース

「 $1.97 \pm 0.04 M_{\odot}$ の質量をもつ中性子星が発見された」
Demorest et al., Nature 467 (2010) 1081 (Oct.28, 2010).

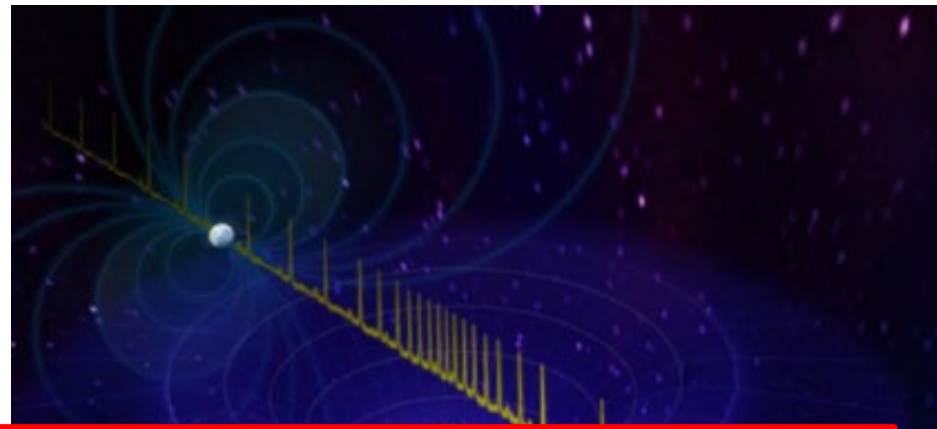
●一般相対論に基づく観測

「パルサー(中性子星)からくる光が伴星(白色矮星)の近くを通り、
時間が遅れる(Shapiro delay)。」

論文での主張

$(1.97 \pm 0.04) M_{\odot}$ の中性子星は、
ハイペロン、中間子凝縮を含む
状態方程式では支えられない。
クォーク物質でも強い相互作用が

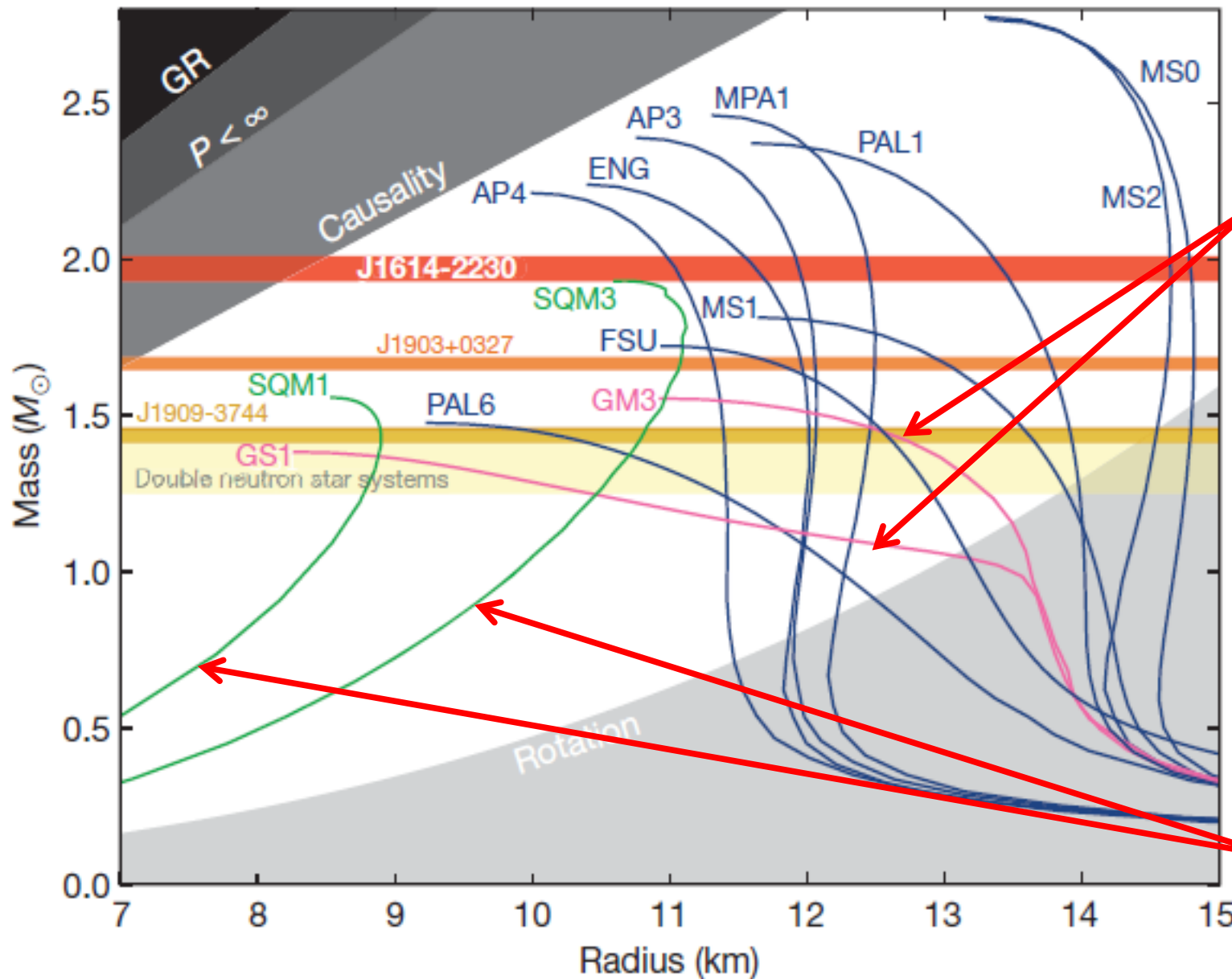
必要である



signature. We calculate the pulsar mass to be $(1.97 \pm 0.04)M_{\odot}$, which rules out almost all currently proposed²⁻⁵ hyperon or boson condensate equations of state (M_{\odot} , solar mass). Quark matter can support a star this massive only if the quarks are strongly interacting and are therefore not ‘free’ quarks¹².

$1.97 \pm 0.04 M_{\odot}$ Neutron Star

Demorest et al., Nature 467 (2010) 1081 (Oct.28, 2010).

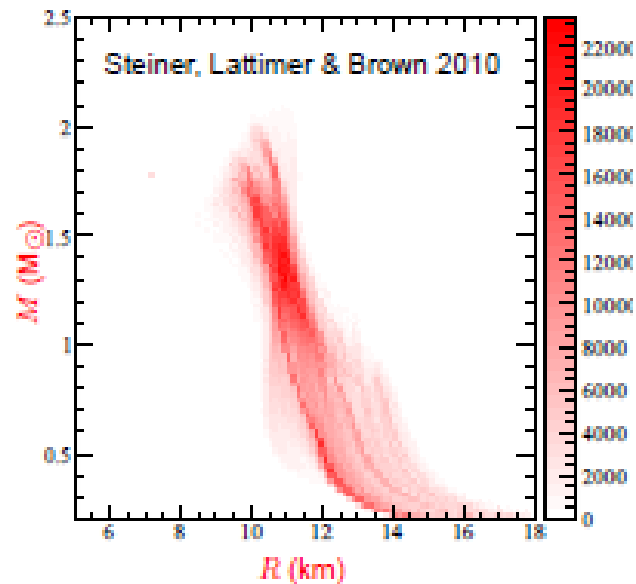
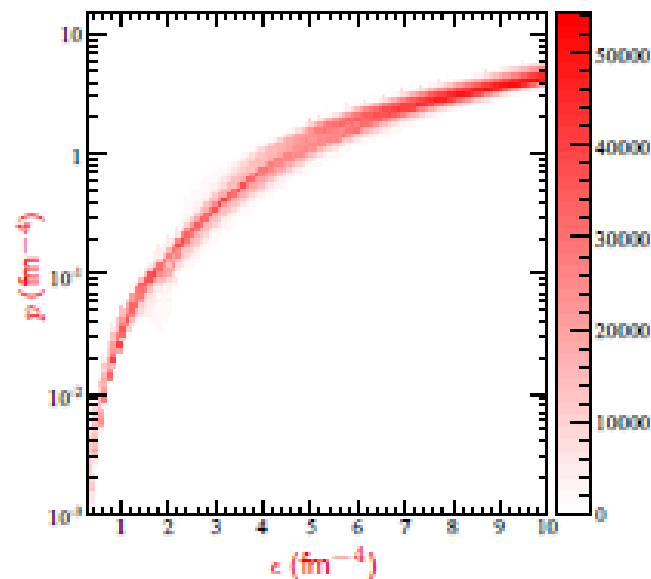
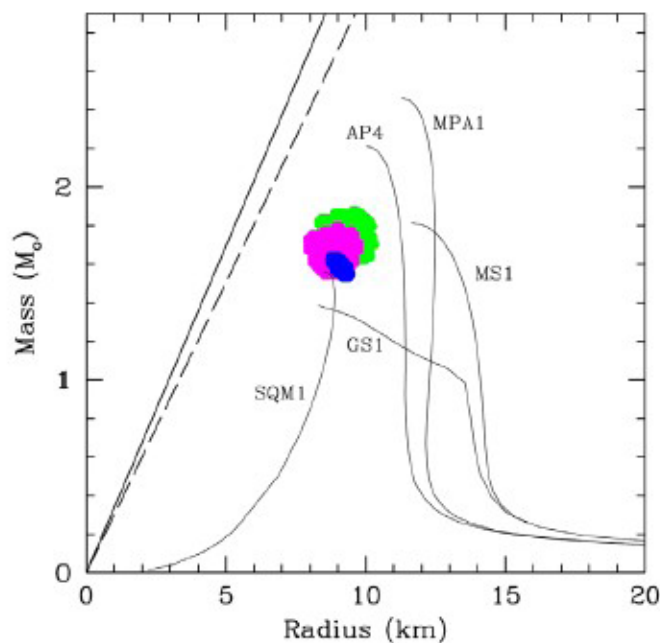


ハイペロンを
含むEOS

クォーク物質
のみのEOS

Neutron Star Radius

- 中性子星の質量・半径同時測定
- TOV方程式を使うと M (質量)- R (半径)関係式と EOS は1対1対応
- M , R が同時に決まると、EOS に非常に強い制限
[観測された(M , R) の”点”を通る必要がある！]
- X線バースト観測 → 半径(+質量)の情報



Ozel, Baym & Guver,
PRD82('10)101301 [arXiv: 1002.3153]

Steiner, Lattimer, Brown, *ApJ 722 (2010) 33*
[arXiv:1005.0811]

Neutron Star Cooling

■表面温度測定と冷却曲線

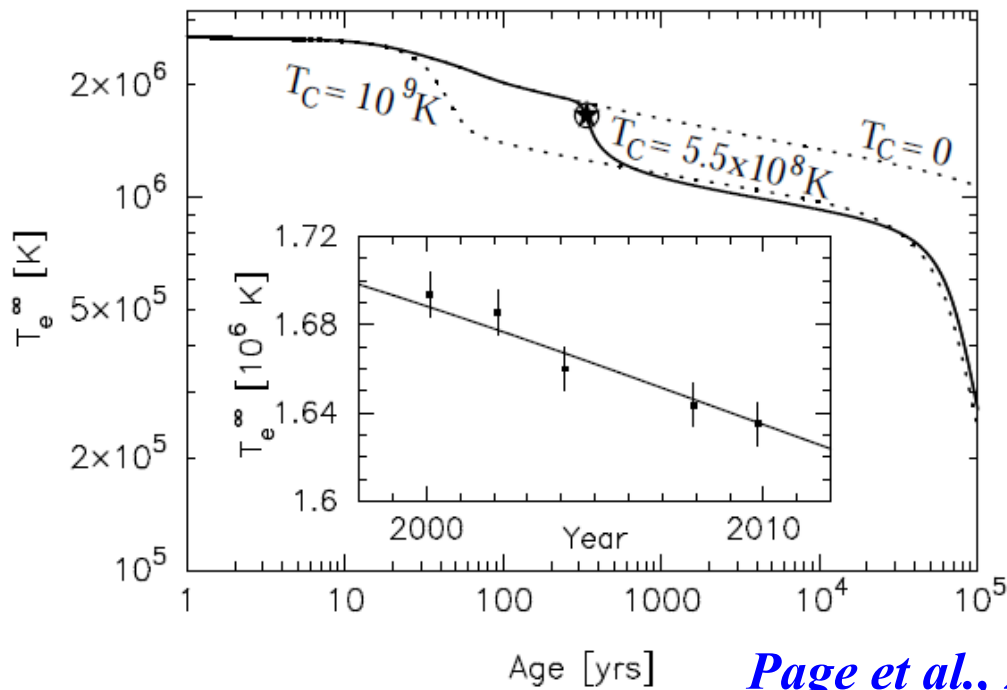
●Cas A の正確な温度測定と冷却率の測定

Heinke, Ho, ApJ 719('10) L167 [arXiv:1007.4719]

Page et al., PRL 106 ('11) 081101 [arXiv:1011.6142]

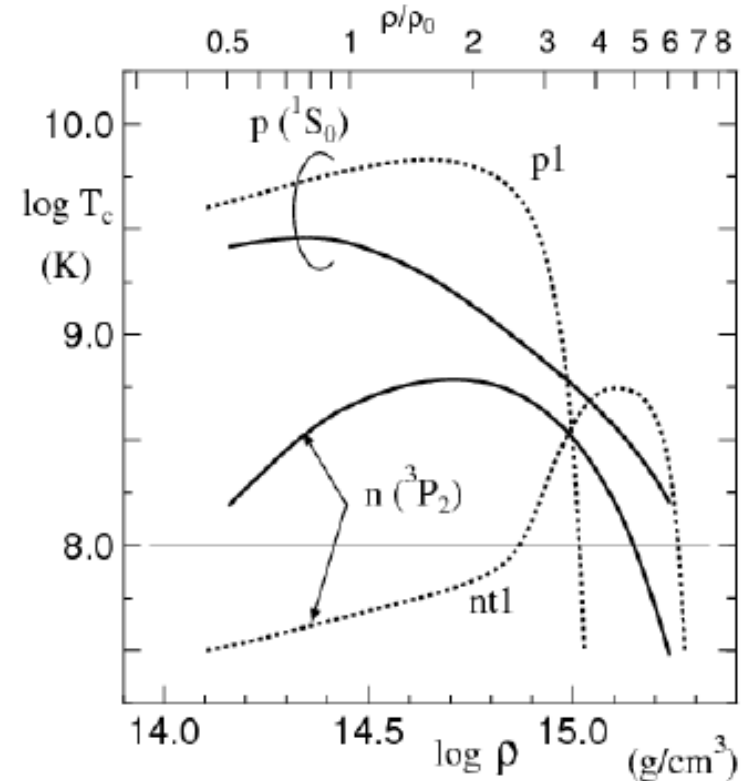
●neutron pair のbreaking & formation

■核物理への宿題: $5\rho_0$ 程度までのギャップを正確に測定・計算できるか?



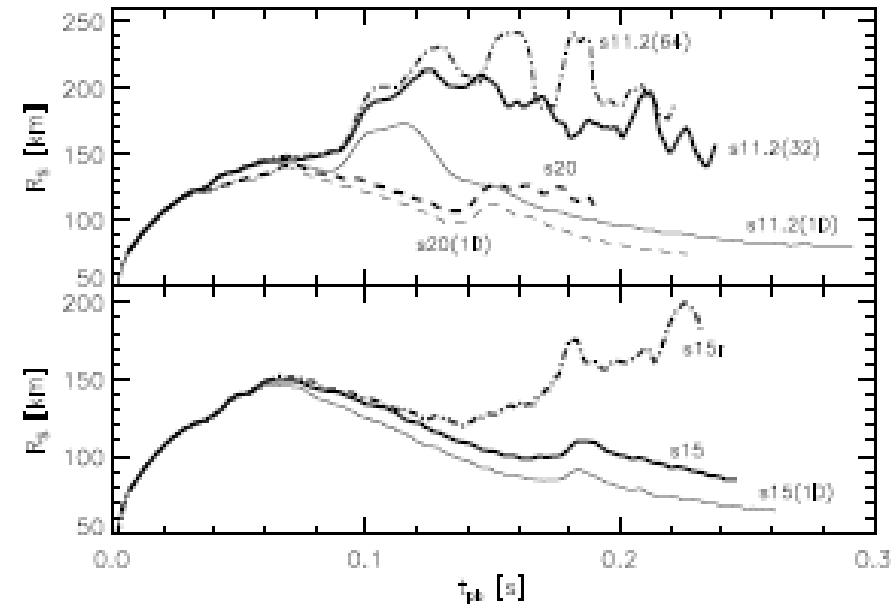
Page et al., 2011

Takatsuka

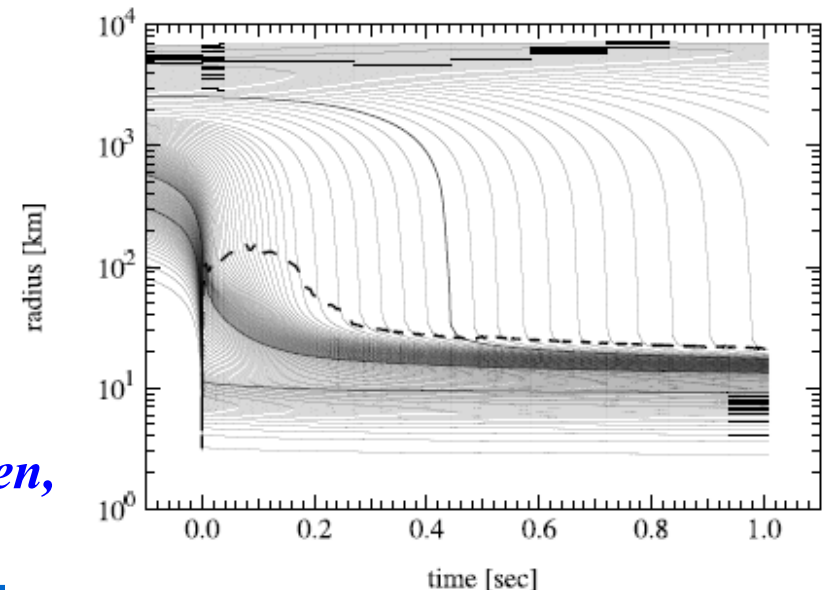


Numerical Simulation of Supernova Explosion

- ν radiation hydrodynamics
- Baryons, Electrons, Photons (Hydro)
- + neutrinos (Boltzmann)
- 1-dim. (Spherical Sym.)
→ Exact ν transport leads to failed supernova explosion failure. (Sumiyoshi et al., 2005)
- 2-dim. Hydrodynamics → marginal (Janka et al., 2002)



R. Buras, M. Rampp, H.-Th. Janka, K. Kifonidis, PRL90(03)241101



K. Sumiyoshi, S. Yamada, H. Suzuki, H. Shen, S. Chiba, H. Toki, ApJ629(05)922

Numerical Simulation of Supernova Explosion

- Recent developments (approximate v transport)

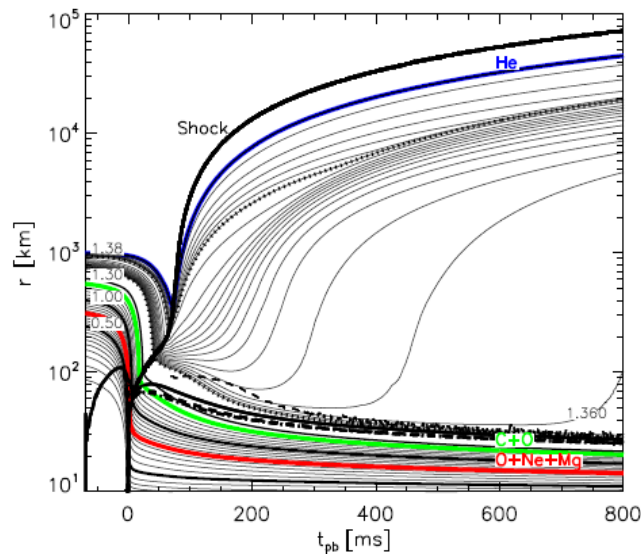
- Light progenitor (8-10 Msun, 1D)

- Successful explosion with simultaneous calc. of nucleosynthesis

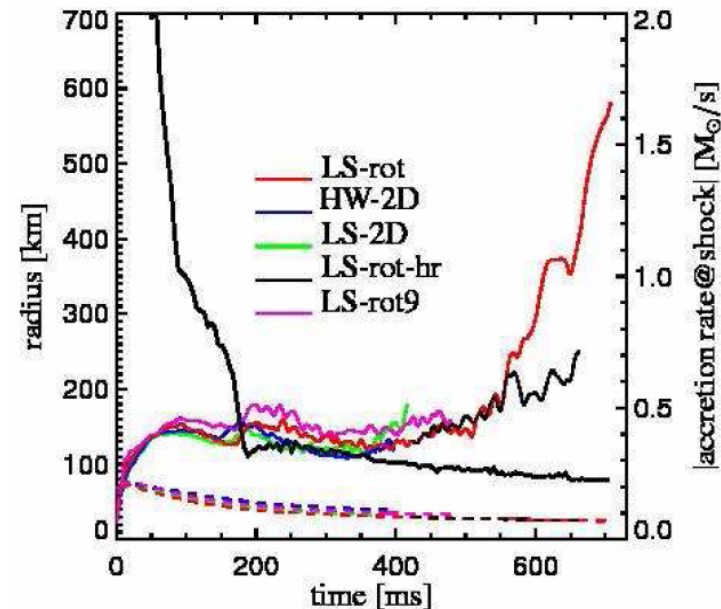
- Heavy progenitor (15 Msun, 2D)

- Standing accretion shock instability (SASI) causes late expl.

There are some successful examples, but not conclusive yet.



Kitaura, Janka, Hillebrandt, 2006



Marek, Janka, 2008

Numerical Simulation of Supernova Explosion

More recent example

Soft EOS + 2D + rotation

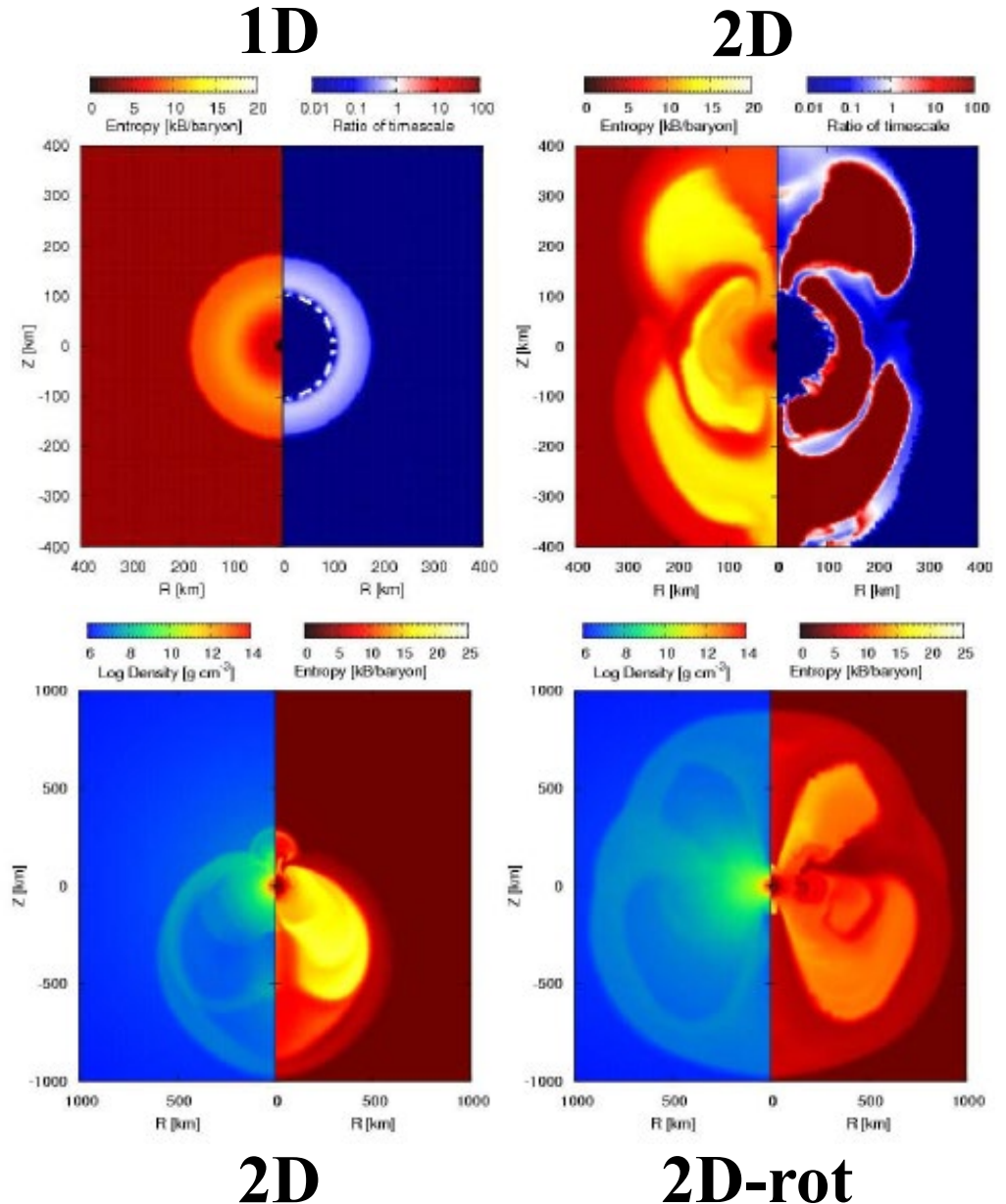
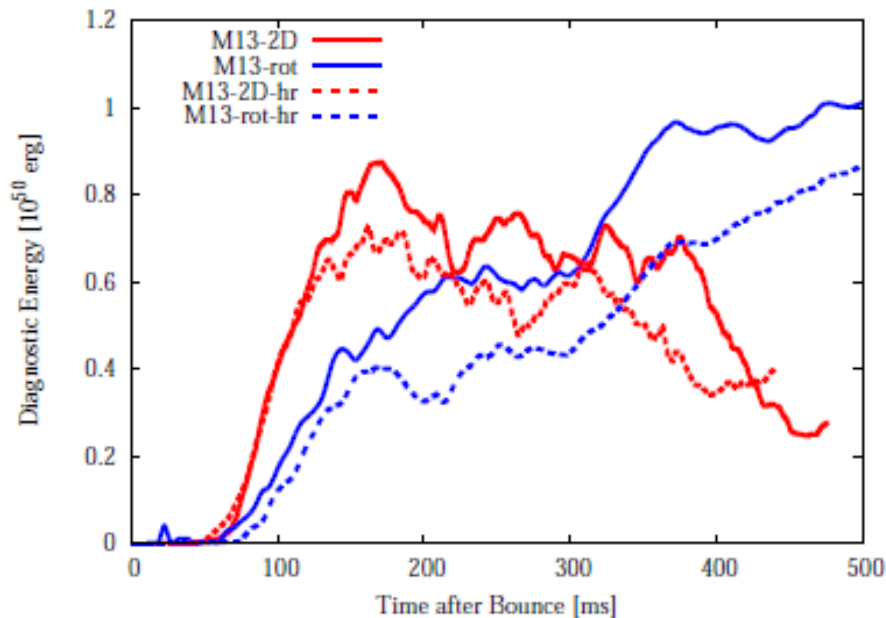
leads stronger explosion

($\sim 10^{50}$ erg $\sim 10\%$ of observed E.)

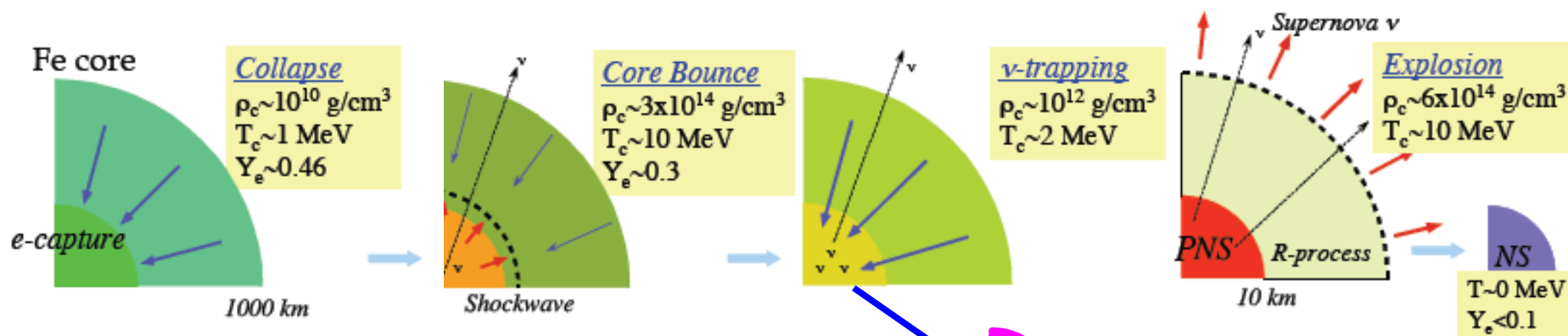
Y.Suwa, K.Kotake, T.Takiwaki, S.C.

Whitehouse, M. Liebendoerfer, K.Sato,

Publ.Astron.Soc.Jap. 62 (2010) L49.



Dynamical Black Hole Formation



■コアの重力崩壊の後、超新星爆発が起こらない(失敗する)と...

→ ブラックホール形成

$M > (20-25) M_{\odot}$: *S.J.Smartt, J.J. Eldridge, R.M. Crockett, J.R. Maund, MNRAS 395('09), 1409; K.Nomoto, N.Tominaga, H.Umeda, C.Kobayashi, K.Maeda, NPA777('06)42.*

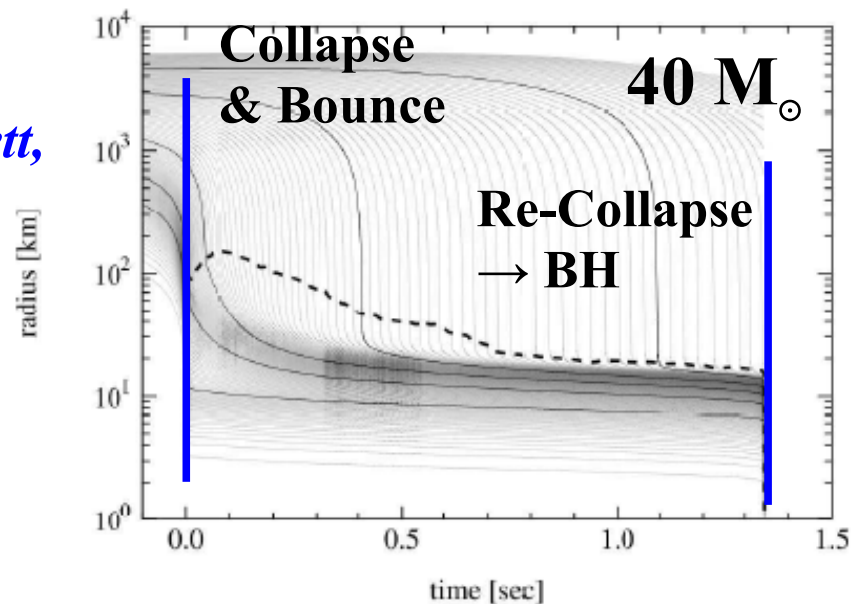
■ ν radiation hydro. シミュレーション

→ $t \sim 1$ sec. で horizon 形成

→ 急速な ν 減少

Sumiyoshi, Yamada, Suzuki, Chiba, PRL 97('06) 091101.

BH $T \sim (70-90) \text{ MeV}$
 $Y_e \sim (0.1-0.3)$



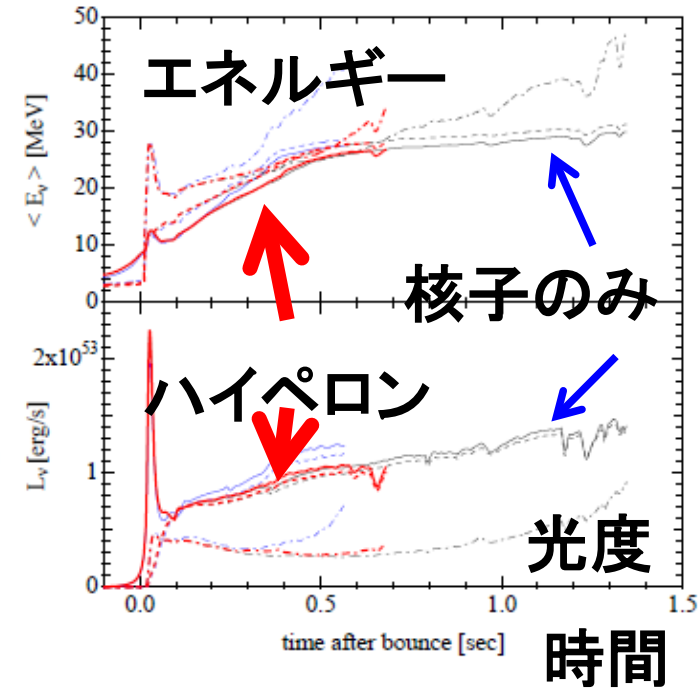
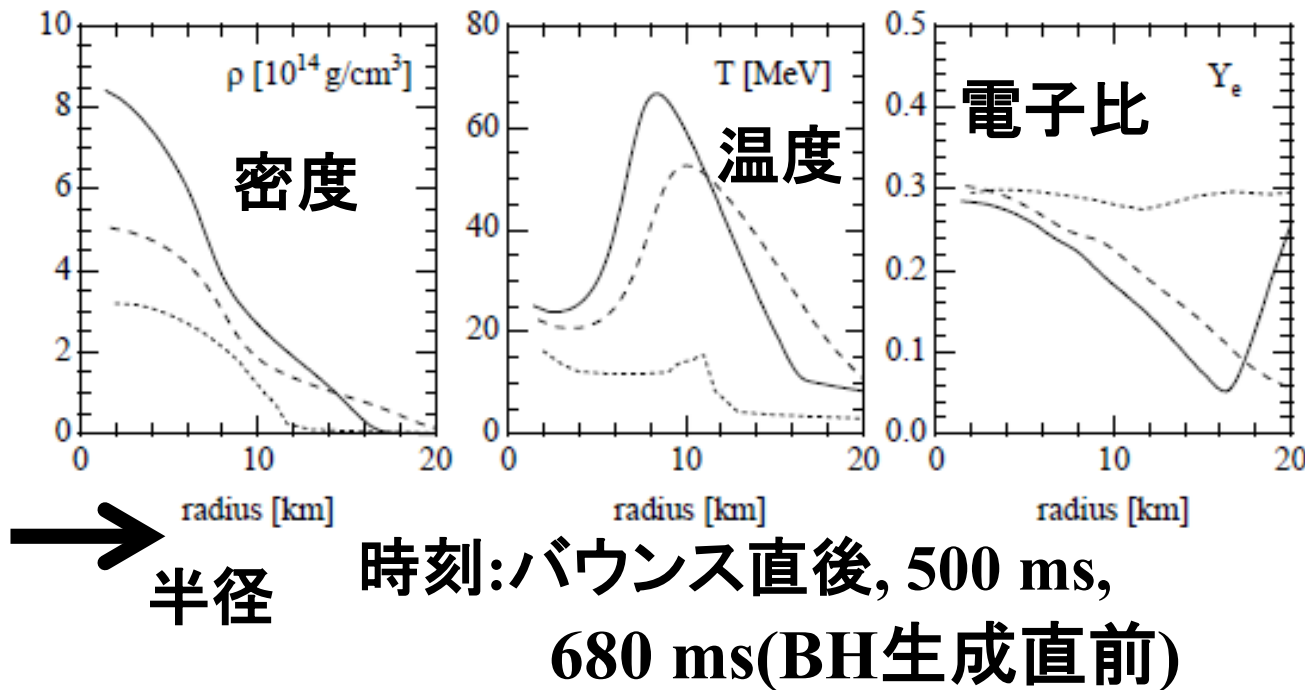
Black Hole Formation (Failed Supernova)

■ “Hot” and “Dense” matter in BH formation process !

● $T_{\max} \sim 90$ MeV (Nucleon), 70 MeV (w/ Hyperon)

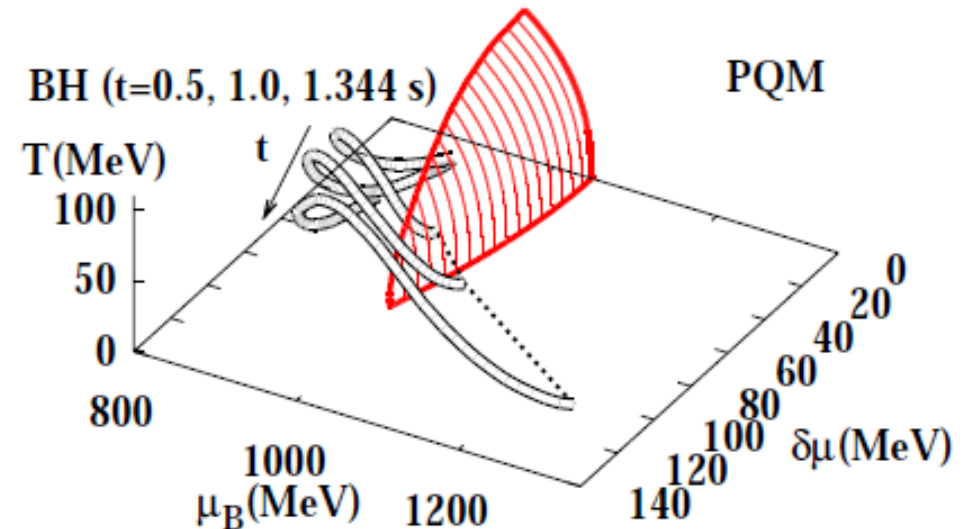
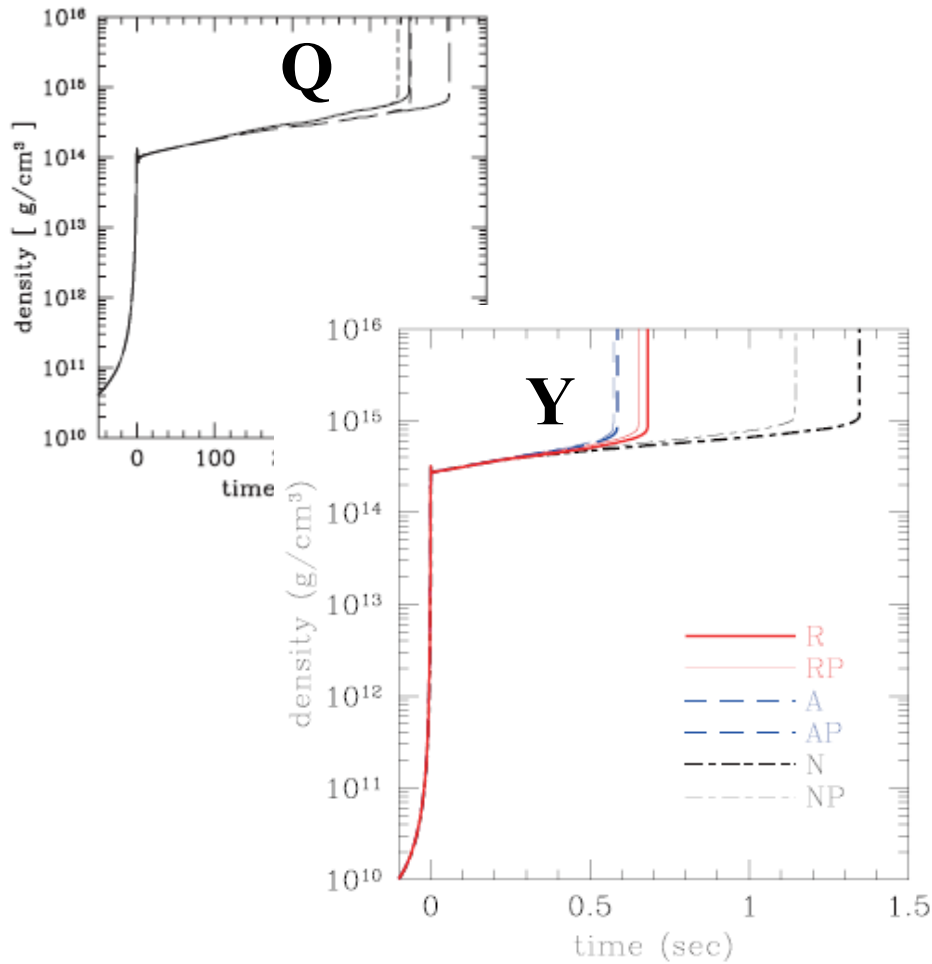
● $\rho_{\text{Bmax}} \sim 4 \rho_0$

■ ν spectrum is sensitive to dense matter EOS (stiffness, hyperon, quark, ...)



Sumiyoshi, Ishizuka, AO, Yamada, Suzuki, 2009

Black Hole Formation (Failed Supernova)



AO, Ueda, Nakano, Ruggieri, Sumiyoshi, PLB 704 ('11), 284 [arXiv:1102.3753]

*Nakazato, Furusawa, Sumiyoshi, AO, Yamada, Suzuki, ApJ, to appear
K.Nakazato, K.Sumiyoshi, S. Yamada, PRD77 (2008) 103006*

Short Summary of Compact Astrophysical Phenomena

■ Neutron Stars

- μ_B (center) ~ 1650 MeV (TM1, $M_{\max} = 2.17 M_{\odot}$)

→ Much larger than Λ mass and hyperons are expected to admix.

- Challenge: How to find the mechanism to suppress hyperons
OR stiffen to support $1.97 M_{\odot}$ with hyperons (very stiff quark matter ?)

■ Supernovae

- Some 2D v radiation hydrodynamics results show explosions.
(Explosion energy is too small / soft ($K \sim 180$ MeV) EOS is used.)

- No consensus on the explosion mechanism.

■ Black hole formation

- Hot and dense nuclear matter is formed.

- v sp

In all of these phenomena, nuclear matter EOS from low to high density is necessary !

*Nuclear Matter EOS
for Compact Astrophysical Phenomena*

コンパクト天体現象に用いる核物質状態方程式

- 超新星爆発計算 = v 輸送を取り入れた流体模型
- 時間スケール ~ 数100 msec = v 以外は熱・化学平衡
- 状態方程式: 核子、電子、光子、**原子核**、**ハイペロン**、 π , K, クォーク、...
- 輸送方程式(Boltzmann) : v -A 断面積、 e -捕獲率
- 状態方程式 → 有限温度効果、広い密度・ Y_p 範囲、公開
- 第一原理計算 (LQCD, GFMC, Variational, DBHF, G-matrix)
→ 飽和性の説明には現象論的3体力などが必要
- Lattimer-Swesty (LS) EOS (*J. M. Lattimer, F.D. Swesty, NPA535('91)331*)
 - ◆ 一様物質 → スキルム力(密度依存ゼロレンジ力)での平均場
 - ◆ 非一様効果 → 圧縮性液滴
- Relativistic EOS (Shen EOS)
(*H. Shen, H. Toki, K. Oyamatsu, K. Sumiyoshi, NPA637(1008)435*)
 - ◆ 一様物質 → Relativistic Mean Field (RMF, TM1)

コンパクト天体現象に用いる核物質状態方程式

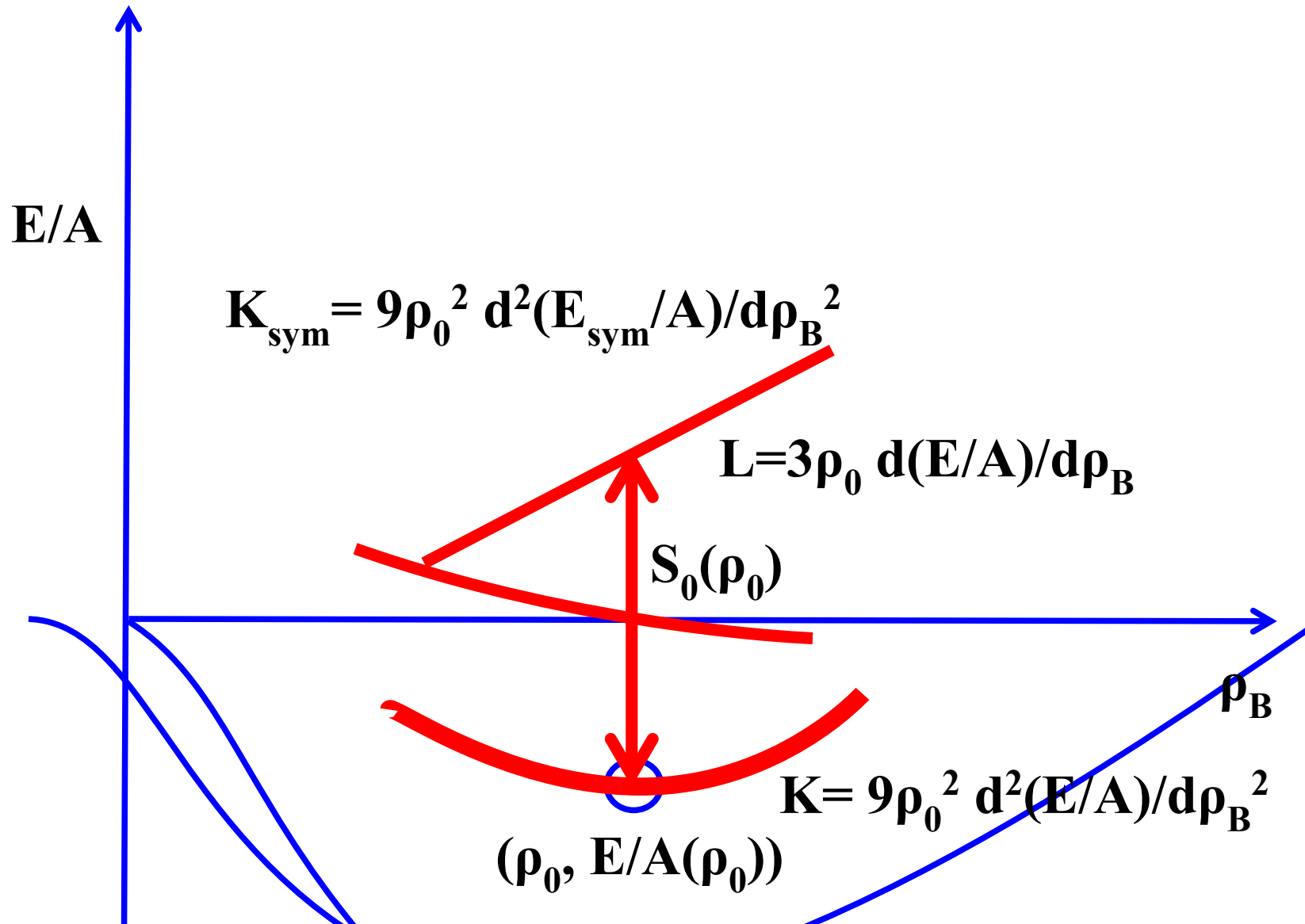
- 超新星爆発計算 = v 輸送を取り入れた流体模型
- 時間スケール ~ 数100 msec = v 以外は熱・化学平衡
- 状態方程式: 核子、電子、光子、**原子核**、**ハイペロン**、 π , K, クォーク、...
- 輸送方程式(Boltzmann) : v -A 断面積、e-捕獲率
- 状態方程式 \rightarrow 有限温度効果、広い密度・ Y_p 範囲、公開

● 第一原理計算 (LQCD, GFMC, Variational, DBHF, G-matrix)

\rightarrow 飽和性の説明には現象論的2体力 V_{ij} with operators $(\sigma_i \cdot \sigma_j)_{ij}$, L_{ij}^2 , $\sigma_i \cdot \sigma_j L_{ij}^2$ and $(L \cdot S)_{ij}^2$. The UIX model of V_{ijk} contains two static terms; the two-pion exchange Fujita-Miyazawa interaction, $V_{ijk}^{2\pi}$, and a phenomenological, intermediate range repulsion V_{ijk}^R . The strength of the $V_{ijk}^{2\pi}$ interaction was determined by reproducing the binding energy of the triton via Green's-function Monte Carlo (GFMC) calculations [20], while that of V_{ijk}^R was adjusted to reproduce the saturation density of SNM.

APR paper

Key quantities in Nuclear Matter EOS



Notations for symmetry energy

$$E(\rho, \delta) = E(\rho, 0) + E_{sym}(\rho)\delta^2 + o(\delta^4)$$

$$E(\rho, 0) = E(\rho_0, 0) + \frac{K_0}{2}\varepsilon^2 + o(\varepsilon^3)$$

$$E_{sym}(\rho) = E_{sym}(\rho_0) + L\varepsilon + \frac{K_{sym}}{2}\varepsilon^2 + o(\varepsilon^3)$$

$$K_0 = 9\rho_0^2 \left. \frac{\partial^2 E(\rho, 0)}{\partial \rho^2} \right|_{\rho=\rho_0}$$

$$\delta = (\rho_n - \rho_p) / \rho$$

$$\varepsilon = (\rho - \rho_0) / 3\rho_0$$

$$S_0 = E_{sym}(\rho_0)$$

$$L = 3\rho_0 \left. \frac{\partial E_{sym}(\rho)}{\partial \rho} \right|_{\rho=\rho_0} = (3/\rho_0)P_0$$

$$K_{sym} = 9\rho_0^2 \left. \frac{\partial^2 E_{sym}(\rho)}{\partial \rho^2} \right|_{\rho=\rho_0}$$

$$K_\tau \approx K_{sym} - 6L$$

DBHF and Dirac Phenomenology

Dirac Bruckner-Hartree-Fock

R. Brockmann, R. Machleidt, PRC42('90),1965

Non Rel. Brueckner calculations do not reproduce saturation point (Coester Line).

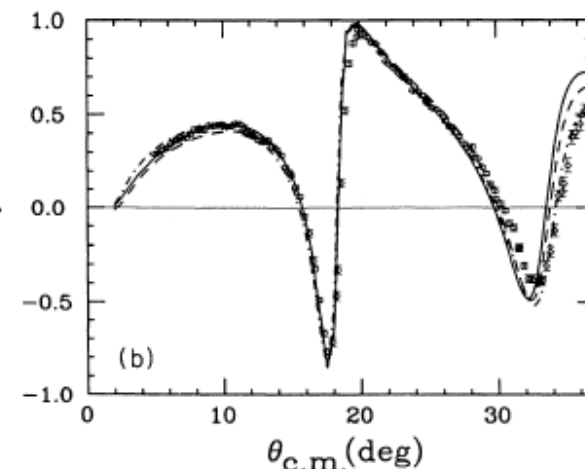
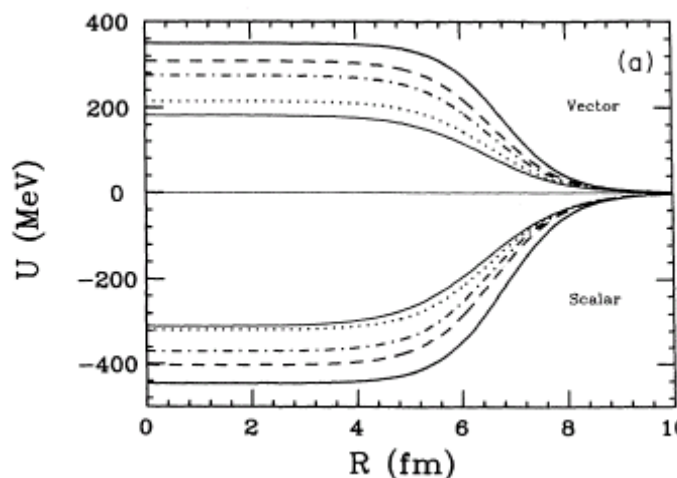
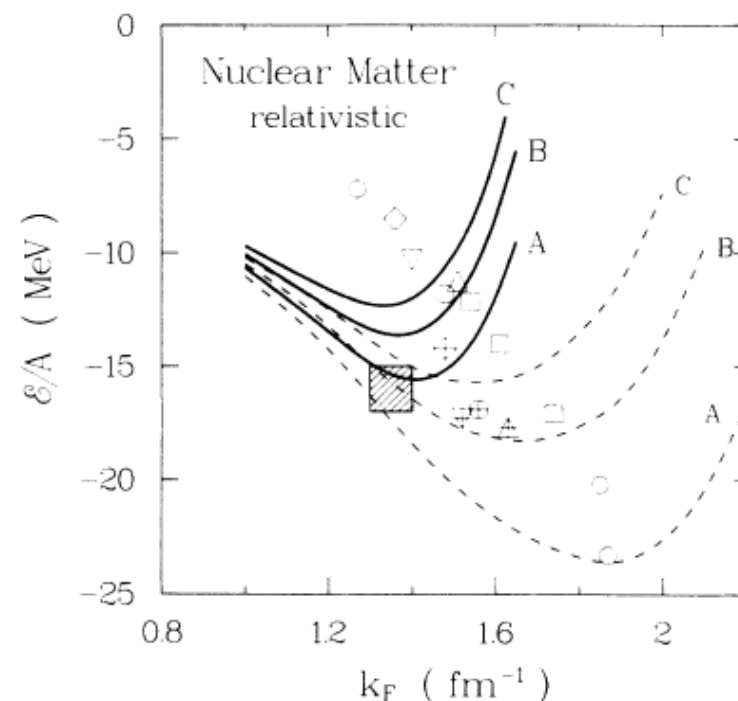
Relativity gives additional repulsion,
→ Saturation point !

Dirac phenomenology

E.D. Cooper, S. Hama, B.C. Clark, R.L. Mercer, PRC47('93),297

Scalar + Vector pA potent
(-400 MeV + 350 MeV)

→ Cross Section,
& Spin Observables



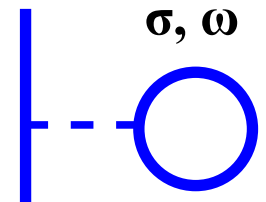
$\sigma\omega$ Model

Serot, Walecka, *Adv.Nucl.Phys.16 (1986),1*

■ Consider only σ and ω mesons

■ Lagrangian

$$L = \bar{\psi}(i\gamma^\mu \partial_\mu - M + g_s \sigma - g_v \gamma^\mu \omega_\mu) \psi + \frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma - \frac{1}{2} m_s^2 \sigma^2 - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} m_v^2 \omega_\mu \omega^\mu$$
$$(F_{\mu\nu} = \partial_\mu \omega_\nu - \partial_\nu \omega_\mu)$$



■ Equation of Motion

$$\frac{\partial}{\partial x^\mu} \left[\frac{\partial L}{\partial (\partial_\mu \phi_i)} \right] - \frac{\partial L}{\partial \phi_i} = 0$$

● Euler-Lagrange Equation

$$\sigma: [\partial_\mu \partial^\mu + m_s^2] \sigma = g_s \bar{\psi} \psi$$

$$\omega: \partial_\mu F^{\mu\nu} + m_v^2 \omega^\nu = g_v \bar{\psi} \gamma^\nu \psi \rightarrow [\partial_\mu \partial^\mu + m_v^2] \omega^\nu = g_v \bar{\psi} \gamma^\nu \psi$$

$$\psi: [\gamma^\mu \partial_\mu - g_v \gamma^\mu \omega_\mu - (M - g_s \sigma)] \psi = 0$$

EOM of ω (for beginners)

■ Euler-Lagrange Eq.

$$\partial_\mu F^{\mu\nu} + m_\nu^2 \omega^\nu = g_\nu \bar{\psi} \gamma^\nu \psi$$

■ Divergence of LHS and RHS

$$\partial_\nu \partial_\mu F^{\mu\nu} + m_\nu^2 (\partial_\nu \omega^\nu) = m_\nu^2 (\partial_\nu \omega^\nu) = g_\nu (\partial_\nu \bar{\psi} \gamma^\nu \psi) = 0$$

LHS: derivatives are sym. and $F_{\mu\nu}$ is anti-sym.

RHS: Baryon Current = Conserved Current

■ Put it in the Euler-Lagrange Eq.

$$\partial_\mu F^{\mu\nu} = \partial_\mu (\partial^\mu \omega^\nu - \partial^\nu \omega^\mu) = \partial_\mu \partial^\mu \omega^\nu - \partial^\nu (\partial_\mu \omega^\mu) = \partial_\mu \partial^\mu \omega^\nu$$

Nuclear Matter in $\sigma\omega$ Model

Serot, Walecka, *Adv.Nucl.Phys.16 (1986),1*

Uniform Nuclear Matter

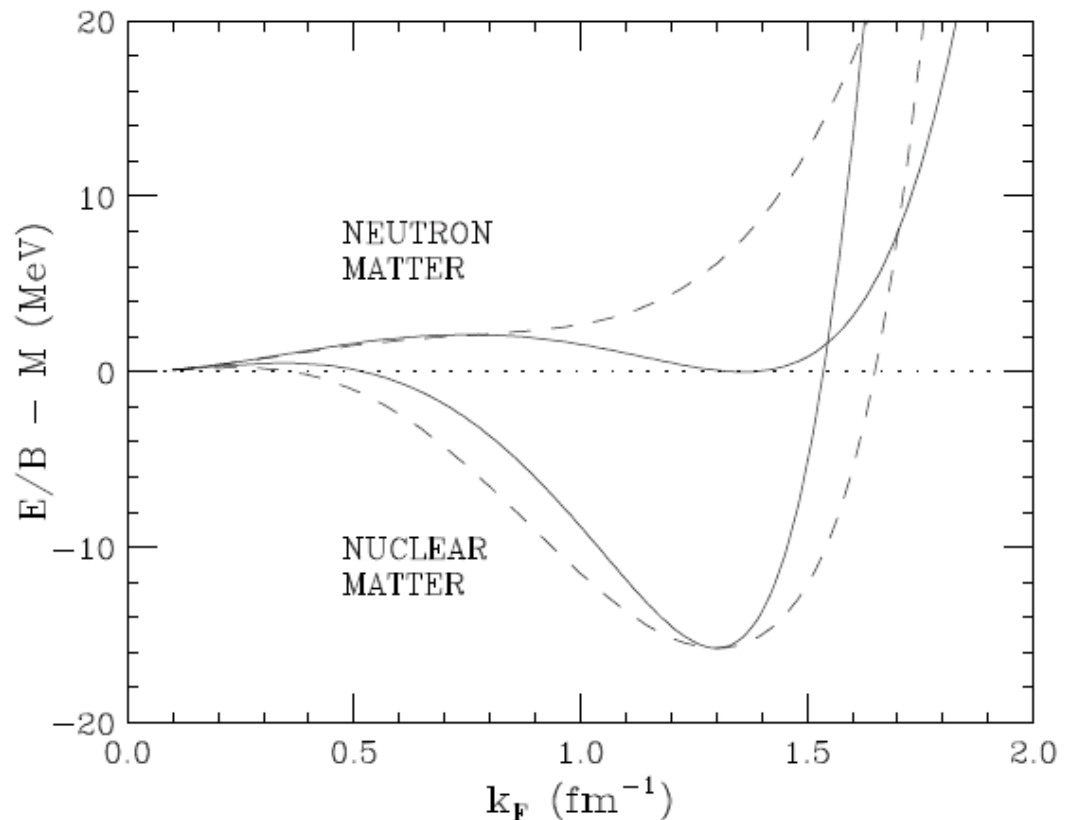
$$E/V = \gamma_N \int^{P_F} \frac{d^3p}{(2\pi)^2} E^* + \frac{1}{2} m_s^2 \sigma^2 - \frac{1}{2} m_v^2 \omega^2 + g_v \rho_B \omega$$

$$\sigma = \frac{g_s}{m_s^2} \rho_s = \gamma_N \frac{g_s}{m_s^2} \int^{P_F} \frac{d^3p}{(2\pi)^2} \frac{M^*}{E^*} \quad \left(M^* = M + U_s = M - g_s \sigma, E^* = \sqrt{p^2 + M^{*2}} \right)$$

$$\omega = \frac{g_v}{m_v^2} \rho_B = \gamma_N \frac{g_v}{m_v^2} \int^{P_F} \frac{d^3p}{(2\pi)^3}$$

$\gamma_N =$ Nucleon degeneracy
(=4 in sym. nuclear matter)

Problem: EOS is too stiff
 $K \sim (500-600) \text{ MeV}!$
 \rightarrow How can we solve ?



σ ω model --- pros and cons

■ Pros (merit)

- Foundation is clear: based on the success of Dirac phen. and DBHF.
- Simple description of scalar and vector potential in σ and ω mesons.
- Saturation is well described in two parameters.
- Natural explanation of large LS potential in nuclei.

■ Cons (shortcomings)

- Relation with the bare NN interaction is not clear.
- Especially, pion effects are not included.
- Symmetry energy is too small.
- Incompressibility is too large ($K \sim 600\text{-}700$ MeV)
(c.f. Empirical value $K \sim (200\text{-}300)$ MeV)
- Chiral symmetry is not respected.

High Quality RMF models

■ Variety of the RMF models

→ MB couplings, meson masses, meson self-energies

● σN , ωN , ρN couplings are well determined

→ almost no model deps. in Sym. N.M. at low ρ

● ω^4 term is introduced to simulate DBHF results of vector pot.

TM: Y. Sugahara, H. Toki, NPA579('94)557;

R. Brockmann, H. Toki, PRL68('92)3408.

● σ^3 and σ^4 terms are introduced to soften EOS at ρ_0 .

J. Boguta, A.R. Bodmer NPA292('77)413,

NL1: P.-G. Reinhardt, M. Rufa, J. Maruhn,

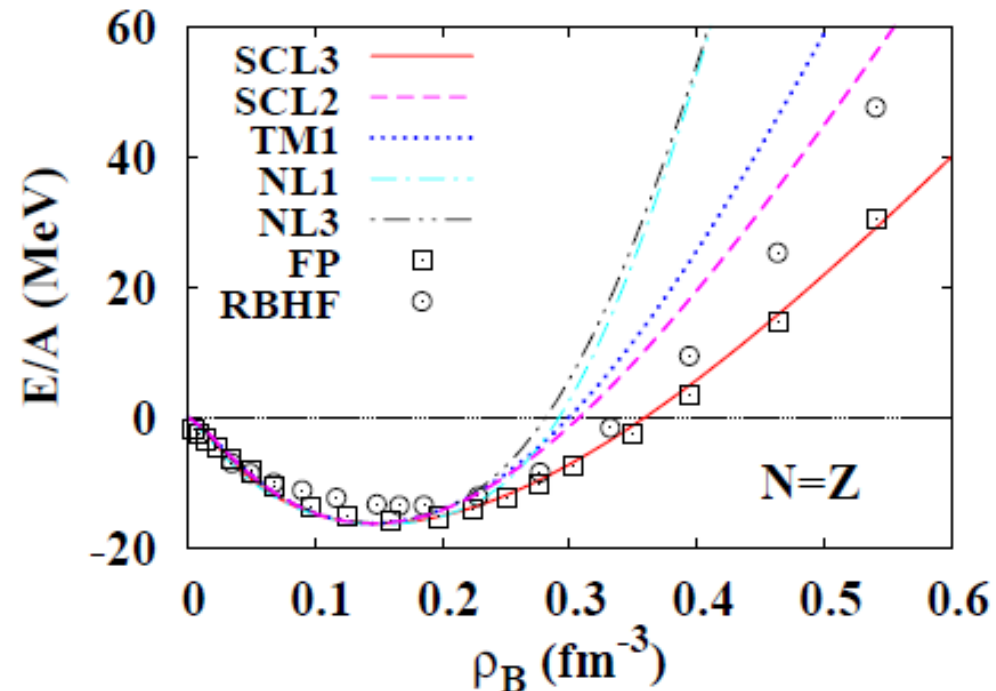
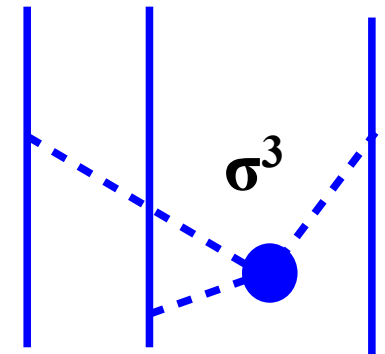
W. Greiner, J. Friedrich, ZPA323('86)13.

NL3: G.A. Lalazissis, J. Konig, P. Ring,

PRC55('97)540.

■ → Large differences are found at high ρ

K. Tsubakihara, H. Maekawa, H. Matsumiya, AO, PRC81('10)065206.



How to determine higher order terms ?

■ Nucleon-meson coupling can be well determined from data !

■ Higher order terms are not well determined, *and give EOS uncertainties at high density !*

■ We need some guiding principle to obtain hadronic Lagrangian including higher order (higher mass dimensional) coupling.

c.f. Naive dim. analysis,

Chiral effective field,

Quark Meson Coupling

(Miyatsu, Saito),

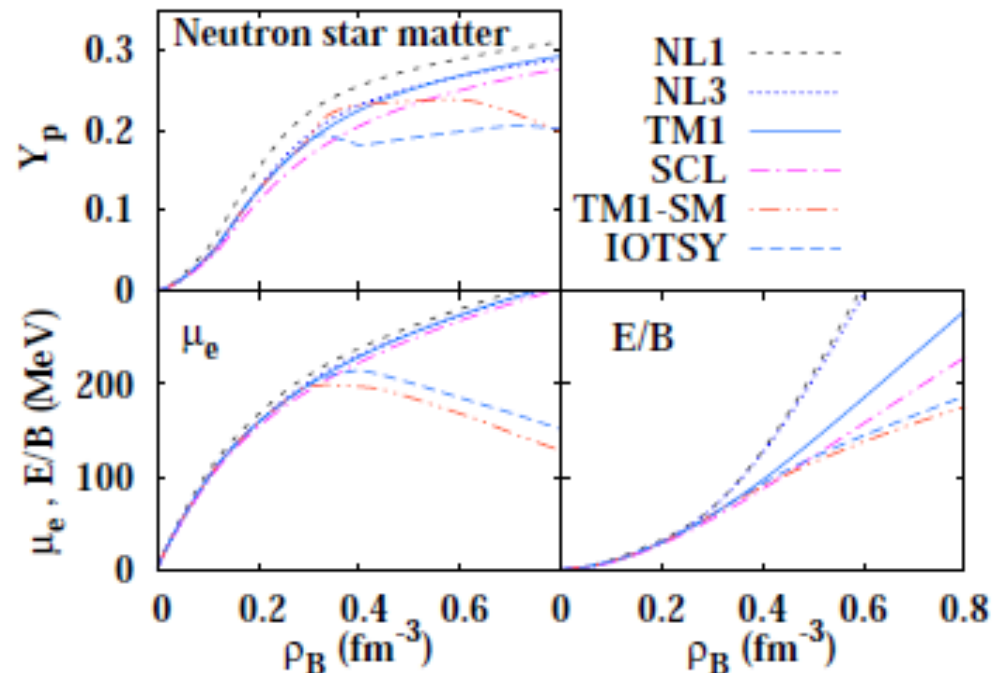


TABLE II. RMF parameters. In SCL, g_3 and g_4 are from the expansion of f_{SCL} .

	$g_{\sigma N}$	$g_{\omega N}$	$g_{\rho N}$	$g_3(\text{MeV})$	g_4	c_ω	$m_\sigma(\text{MeV})$	$m_\omega(\text{MeV})$	$m_\rho(\text{MeV})$
NL1[26]	10.138	13.285	4.976	2401.9	-36.265	0	492.25	795.359	763
NL3[27]	10.217	12.868	4.474	2058.35	-28.885	0	508.194	782.501	763
TM1[28]	10.0289	12.6139	4.6322	1426.466	0.6183	71.3075	511.198	783	770
SCL[29]	10.08	13.02	4.40	1255.88	13.504	200	502.63	783	770

高密度核物質の状態方程式

■構成粒子

●超流動核子、 π 、K、ハイペロン、クォーク、クォーク対、...

■非圧縮率(K): 決まっていない

●GMR (原子核の圧縮振動) $\rightarrow K = 210 \pm 30$ MeV (非相対論的平均場)

●重イオン反応 \rightarrow 平均場の運動量依存性がK依存性を隠す

(Sahu, Cassing, Mosel, AO, 2000; Danielewicz et al., 2002; Isse et al., 2005)

■対称エネルギーの密度依存性: 分りつつある。

●不安定核半径の精密測定で推定可能 (Oyamatsu, Iida, 2007)

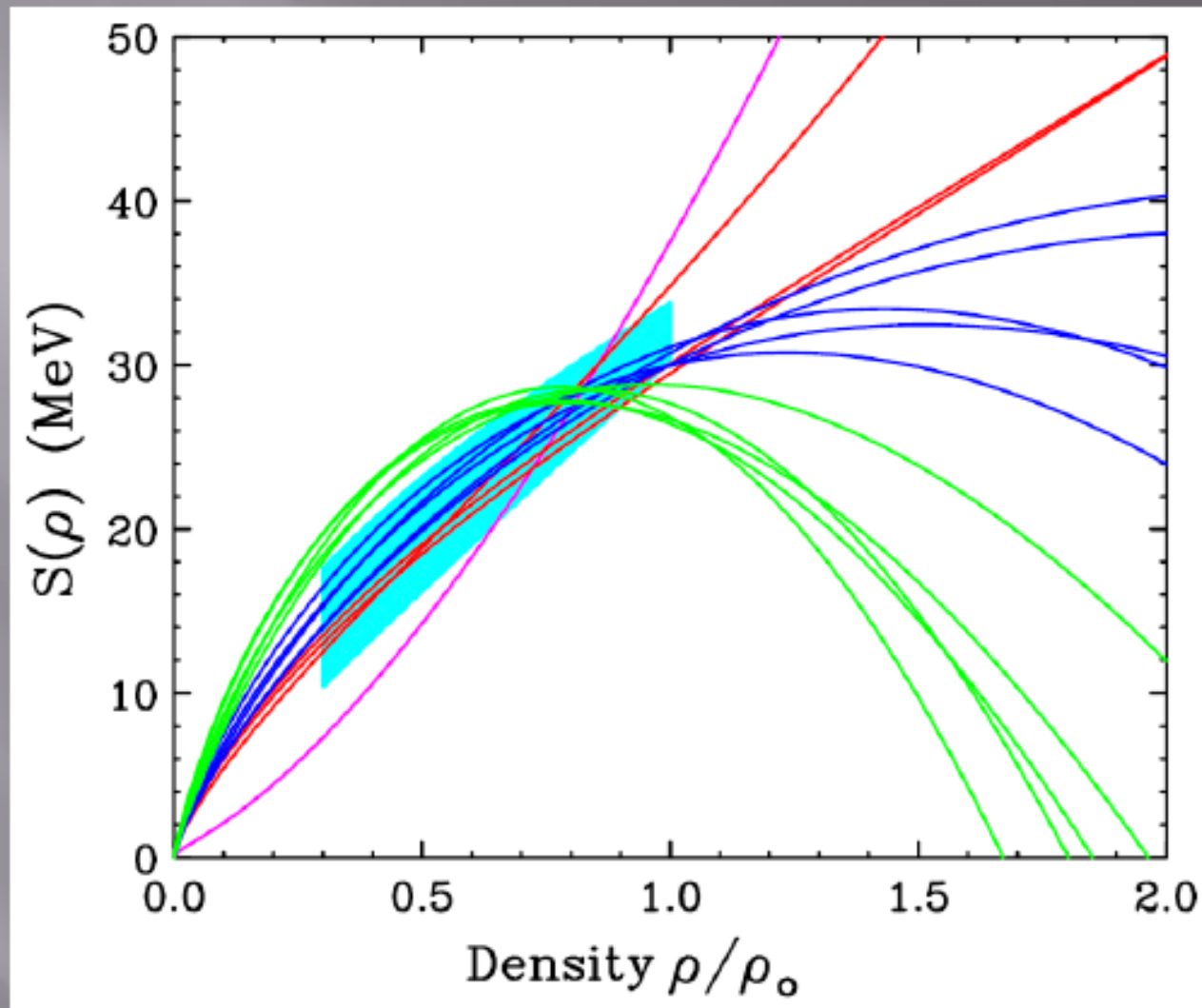
●実験からの制限 (Murakami, 2011)

■核物質中のハドロン・ポテンシャル: 進んでいる

●ハイペロン $U_{\Lambda}(\rho_0) = -30$ MeV, $U_{\Sigma}(\rho_0) = +(15-90)$ MeV, $U_{\Xi}(\rho_0) \sim -15$ MeV

●pionic atom $\rightarrow U_{\pi}(2\rho_0, Y_p \sim 0.2) > +50$ MeV

(AO, Jido, Sekihara, Tshubakihara)



Our consensus is $S_0=31-34$ MeV and $L=50-110$ MeV

Now preparing a summary article on outcomes of NuSYM11.

核物質中でのハイペロン・ポテンシャル

1粒子ポテンシャル

$$U_Y(r) \simeq g_{\sigma Y}\sigma + g_{\omega Y}\omega + g_{\rho Y}R$$

- 核物質中では σ , ω , R ($=\rho$ 中間子期待値)は、与えられている
→ ハイペロン・中間子結合定数によりハイペロン・ポテンシャルが決まる。

結合定数の与え方(hyperon – meson)

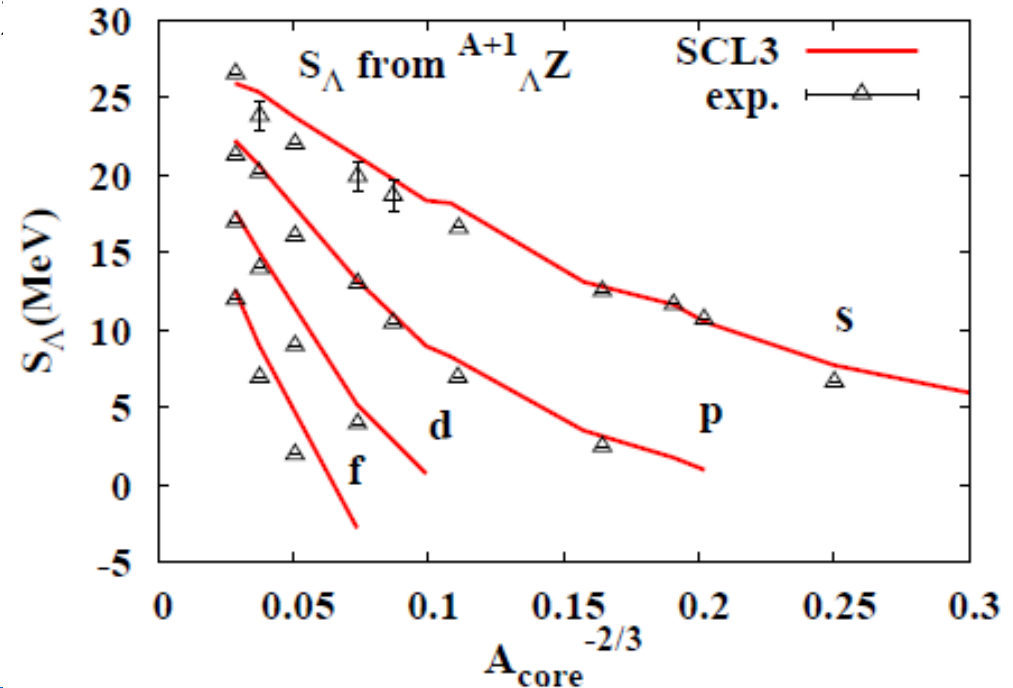
- ベクトル結合: クォーク数カウンティング or フレーバーSU(3)

Quark number counting $\sigma_{\omega Y} =$

$$\text{Flavor } g_{\omega\Lambda} = \frac{5}{6}g_{\omega N} - \frac{1}{2}g_{\rho N}$$

- スカラー結合:

ハイペロン・ポテンシャル
orハイパー核束縛エネルギー
を fit

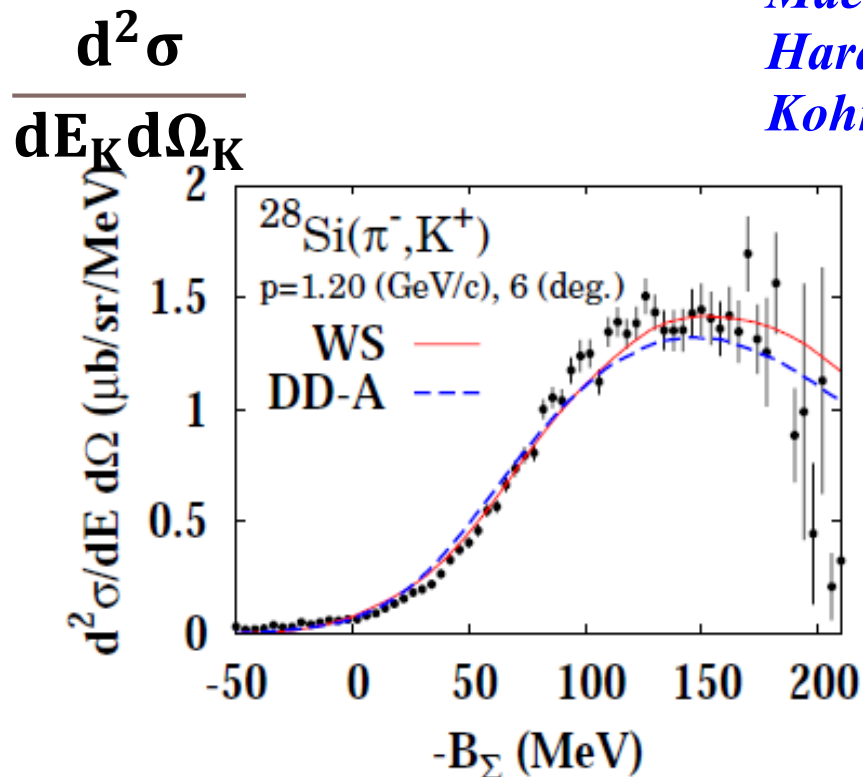


Σ Potential in Nuclear Matter

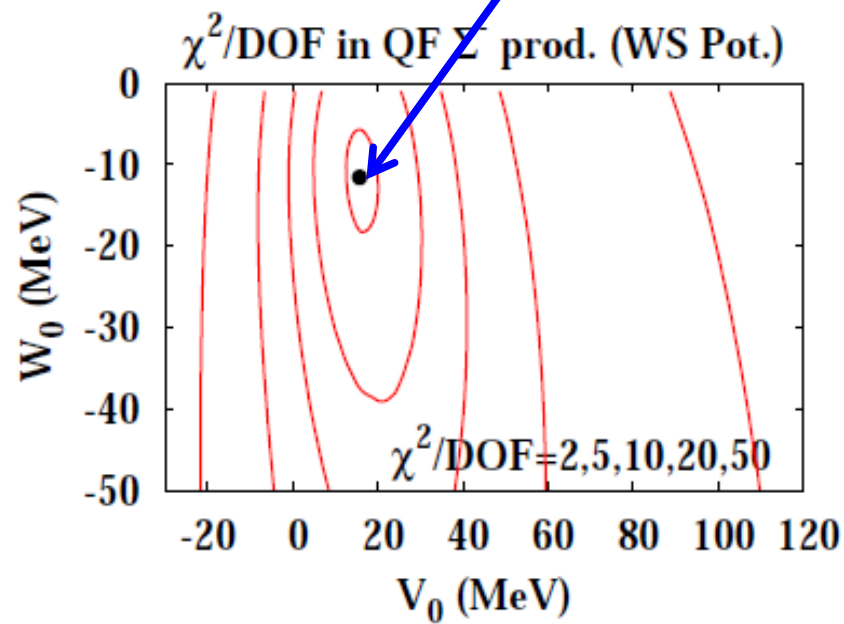
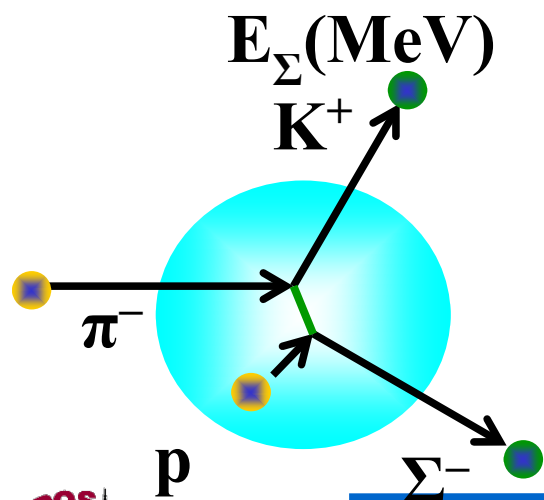
Maekawa, Tsubakihara, AO, EPJA 33(2007),269.

Harada, Hirabayashi, NPA744('04),323.

Kohno, Fujiwara, Kawai, et al. PTP112('04)895



$U_\Sigma(\rho_\rho) \sim +15 \text{ MeV} - i 10 \text{ MeV}$
 with Woods-Saxon potential,
 no Atomic shift fit



Σ feels repulsive potential in nuclei

Relativistic EOS of Supernova Matter with Hyperons

- Extention of the Relativistic (Shen) EOS to $SU_f(3)$ with updated Hyperon Potentials in Nuclear Matter

(Ishizuka, AO, Tsubakihara, Sumiyoshi, Yamada, J. Phys. G35(08),085201)

- Relativistic (Shen) EOS *(Shen, Toki, Oyamatsu, Sumiyoshi, PTP 100('98), 1013)*
Rel. Mean Field (RMF) + Local Density Approx. (Nuclear Formation)

- $SU_f(3)$ Extention of RMF *(Schaffner, Mishustin, PRC53 (1996), 1416)*

$$\mathcal{L} = \mathcal{L}_{Free}(B, \sigma, \omega_\mu, \vec{R}_\mu, \zeta, \phi_\mu) - U_\sigma(\sigma) + \frac{1}{4} c_\omega (\omega^\mu \omega_\mu)^2 - \sum_B \bar{\Psi}_B \left(g_{\sigma B} \sigma + g_{\omega B} \omega + g_{\rho B} \vec{R} \cdot \vec{\tau}_B + g_{\zeta B} \zeta - g_{\phi B} \gamma^\mu \phi_\mu \right) \Psi_B$$

Coupling ~ Quark Number Counting

- $g_{\sigma Y}$ is tuned to fit Hyperon Potential in Nuclear Matter
 $U_\Lambda = -30 \text{ MeV}, U_\Sigma = +30 \text{ MeV}, U_\Xi = -15 \text{ MeV}$
- Nuclear Formation is included using Shen EOS table

Hyperon Composition in Dense Matter

■ Hyperon start to emerge at $(2-3)\rho_0$ in Neutron Star Matter !

■ Hyperon composition in NS is sensitive to Hyperon potential.

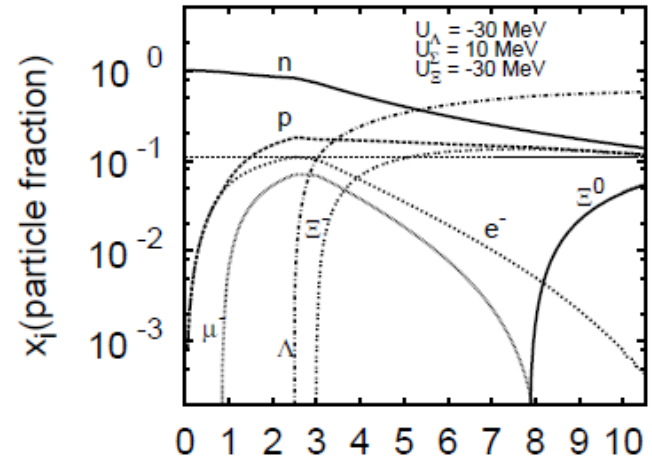
● $U_\Lambda \sim -30$ MeV: Well-known

● $U_\Sigma, U_\Xi \sim -30$ MeV (Old conjecture)
 → Σ - appears prior to Λ

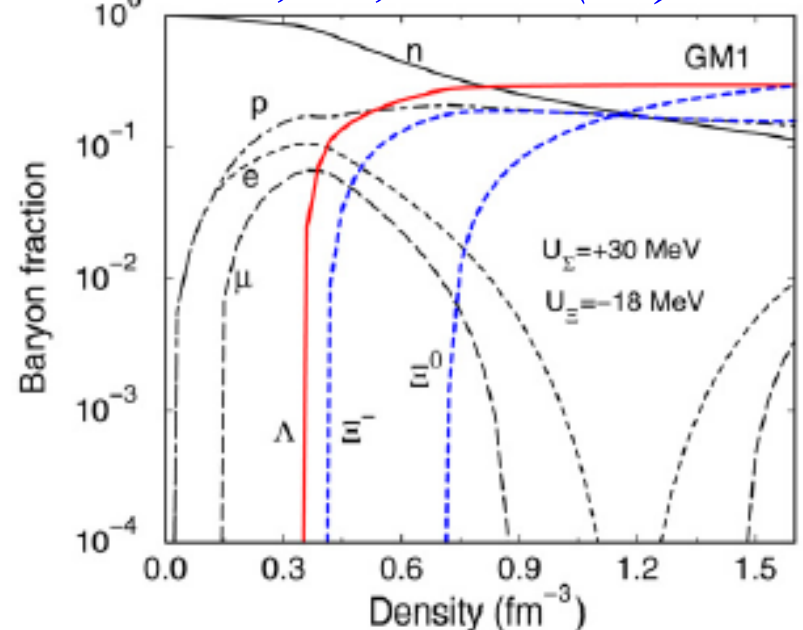
● $U_\Sigma > 0$ (repulsive) → No Σ in NS
 Σ atom (phen. fit), QF prod.

S. Balberg, A. Gal, NPA625('97)435;
H. Noumi et al., PRL89('02)072301;
T. Harada, Y. Hirabayashi, NPA759('05)143;
M. Kohno et al. PRC74('06)064613.

● $U_\Xi \sim -(12-15)$ MeV
 (K^-, K^+) reaction, twin hypernuclei
P. Khaustov et al. (E885), PRC61('00)054603;
S. Aoki et al., PLB355('95)45.



P.K.Sahu, AO, NPA691('01)439c

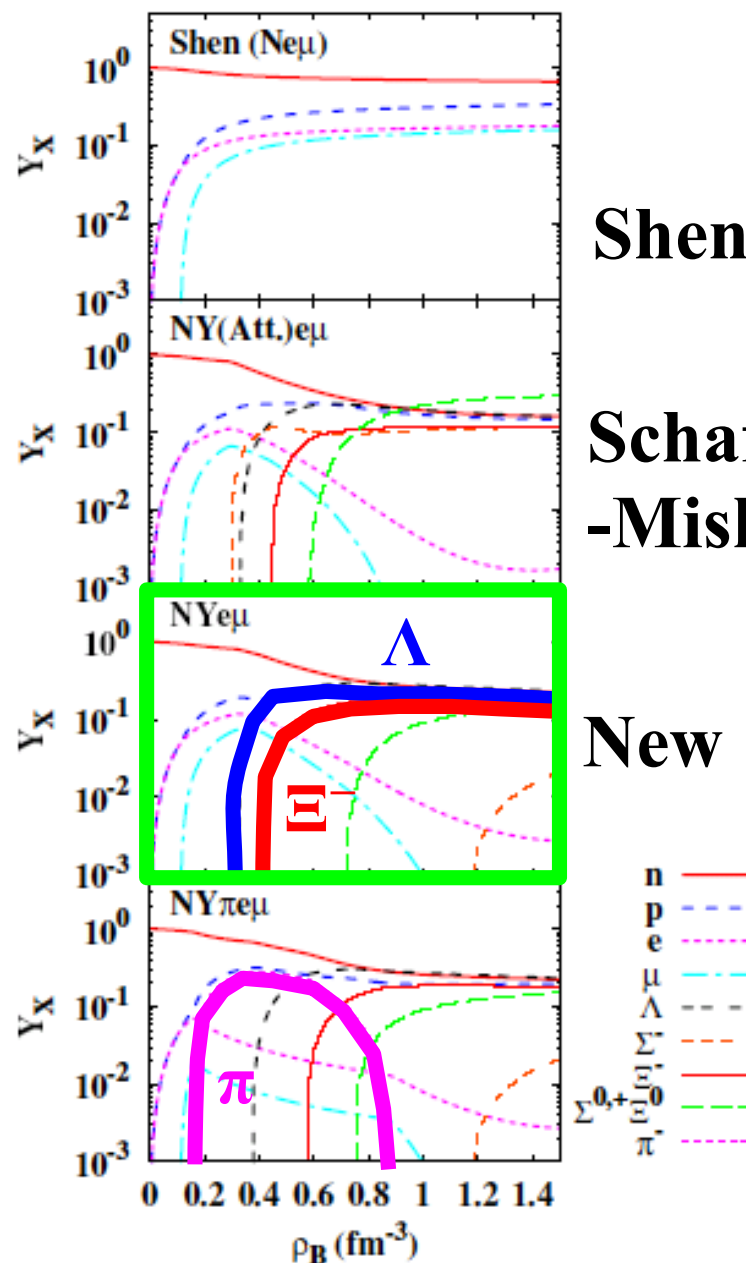
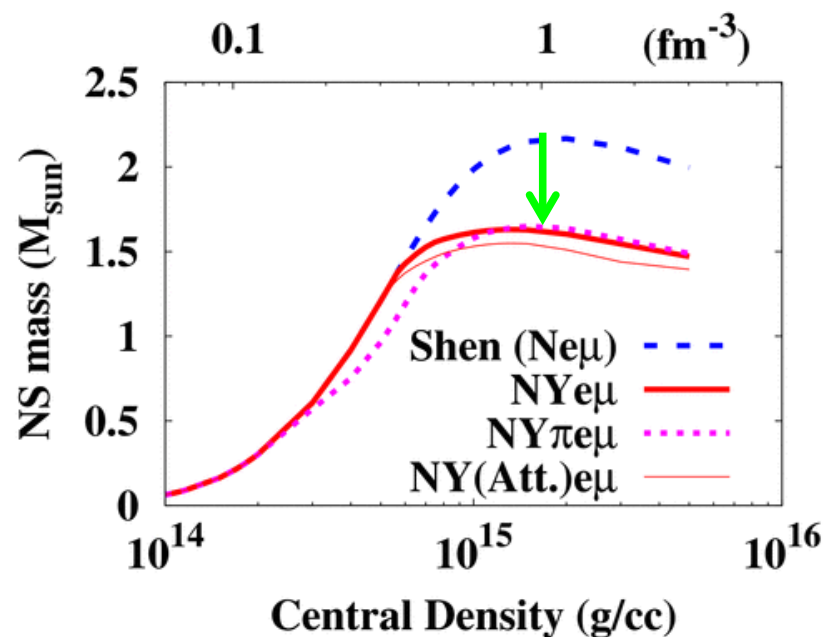


J. Schaffner-Bielich, NPA804('08)309.

Neutron Star

Ishizuka, AO, Tsubakihara, Sumiyoshi, Yamada, *J. Phys. G35(08),085201*

- Hyperon Effect is **DRASTIC**
- $M_{\text{max}} = 2.1 M_{\text{sun}} \rightarrow 1.56 M_{\text{sun}}$
- Composition $Y_{\Lambda} \sim Y_n$
- Large fraction of Ξ
- Thermal (free) pions can admix at $\rho > 1.5 \rho_0$



Shen

Schaffner
-Mishustin

New

Finite Temperature and Supernova

Ishizuka, AO, Tsubakihara, Sumiyoshi, Yamada, *J. Phys. G35(08),085201*

■ Example: $T=10$ MeV, $Y_e = 0.4$

● Λ starts to increase at $\rho \sim 2\rho_0$, becomes significant at $\rho \sim 3\rho_0$.

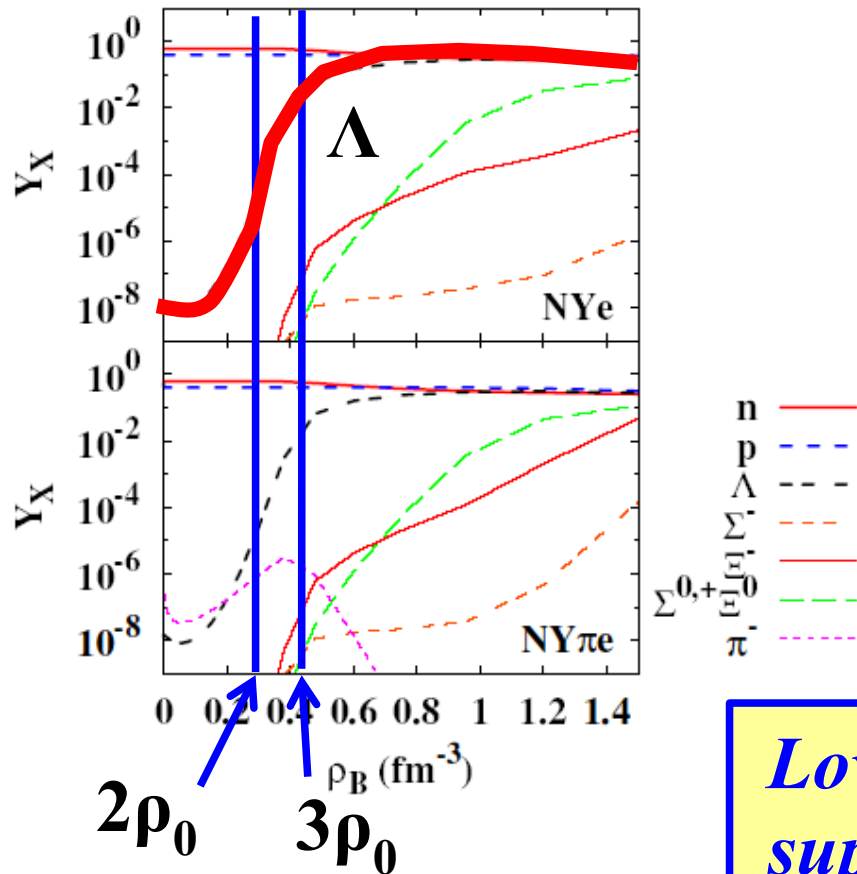
■ Prompt explosion

(without ν transport)

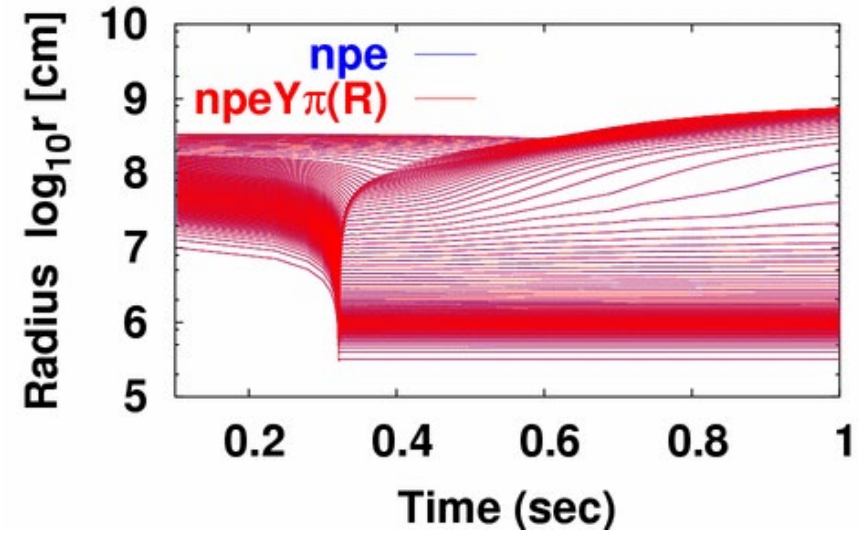
→ Almost no change

(Expl. E. increase $\sim (0.1-0.5\%)$)

$T=10$ MeV, $Y_C=0.4$



$15 M_{\text{solar}}$



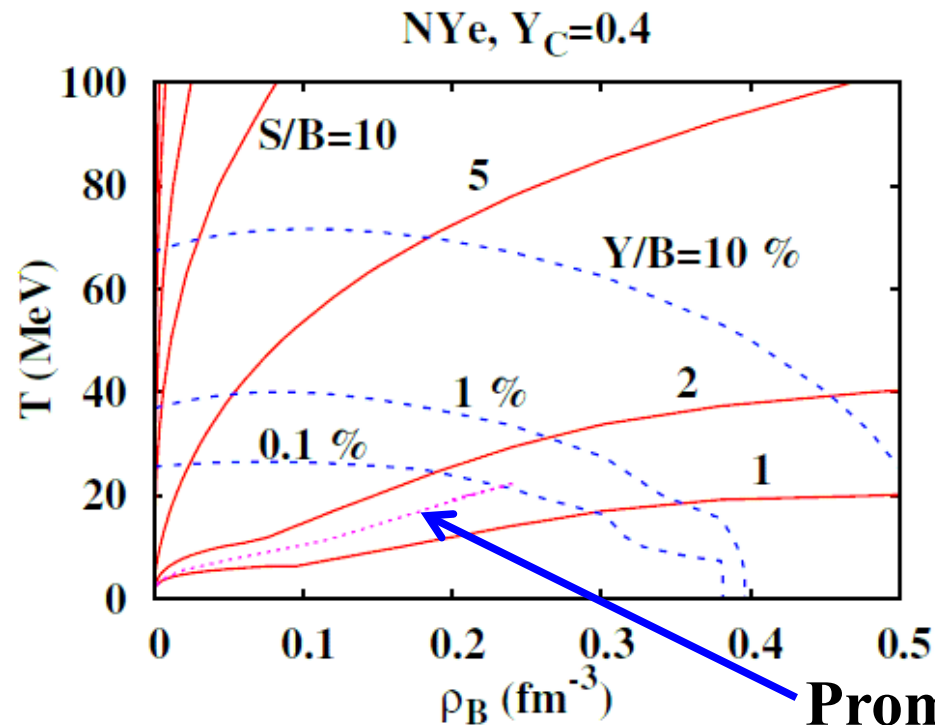
WW95 + 1 Dim. Hydro. (Sumiyoshi, Yamada)

Low density and High Y_e suppresses Hyperons in the Early Stage

Where do we see Hyperons ?

- Hyperon Fraction is sensitive to Y_e , T , and ρ_B .
- $Y_v \sim 0$ (Neutron Star) $\rightarrow \rho_B > 2 \rho_0$
- $Y_e \sim 0.4$ (Supernova, early stage) $\rightarrow T > 40$ MeV or $\rho_B > 3 \rho_0$

Hyperons would be important in Late Stage(Nstar cooling), BH formation, and Heavy-Ion Collisions



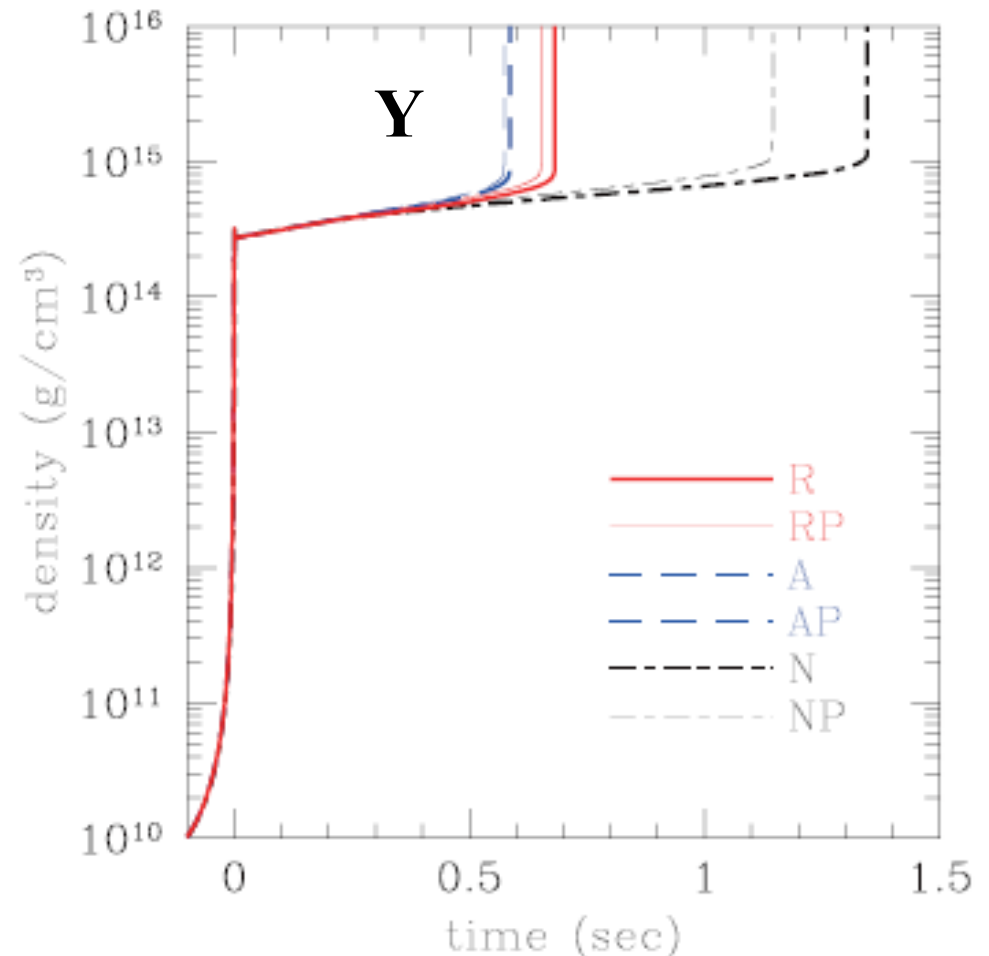
Prompt Expl. (15 Msun)

Black Hole Formation (Failed Supernova)

■ Can we obtain information on Σ repulsion strength from astrophysics ?

● ν duration time during black hole formation is more sensitive to the Σ potential depth in nuclear matter !

Nakazato, Furusawa, Sumiyoshi, AO, Yamada, Suzuki, ApJ, to appear



RMF is a phenomenological MODEL !

■ Baryon one-loop approximation (Hartree approximation) makes RMF a phenomenological model.

→ We need DATA and AB INITIO results.

● Saturation point (ρ_0 and $E/A(\rho_0)$) from mass formula

● Nuclear binding energies

● U_v and U_s from DBHF results

● $P(\rho_B)$ from heavy-ion data

● Λ separation energy from single Λ hypernuclear data

● $\Lambda\Lambda$ bond energy from double Λ hypernuclear data

● Σ atomic shift

● Σ and Ξ potential depth from quasi-free production data

● P_{Δ} pure neutron matter EOS from ab initio calculations (not used here)

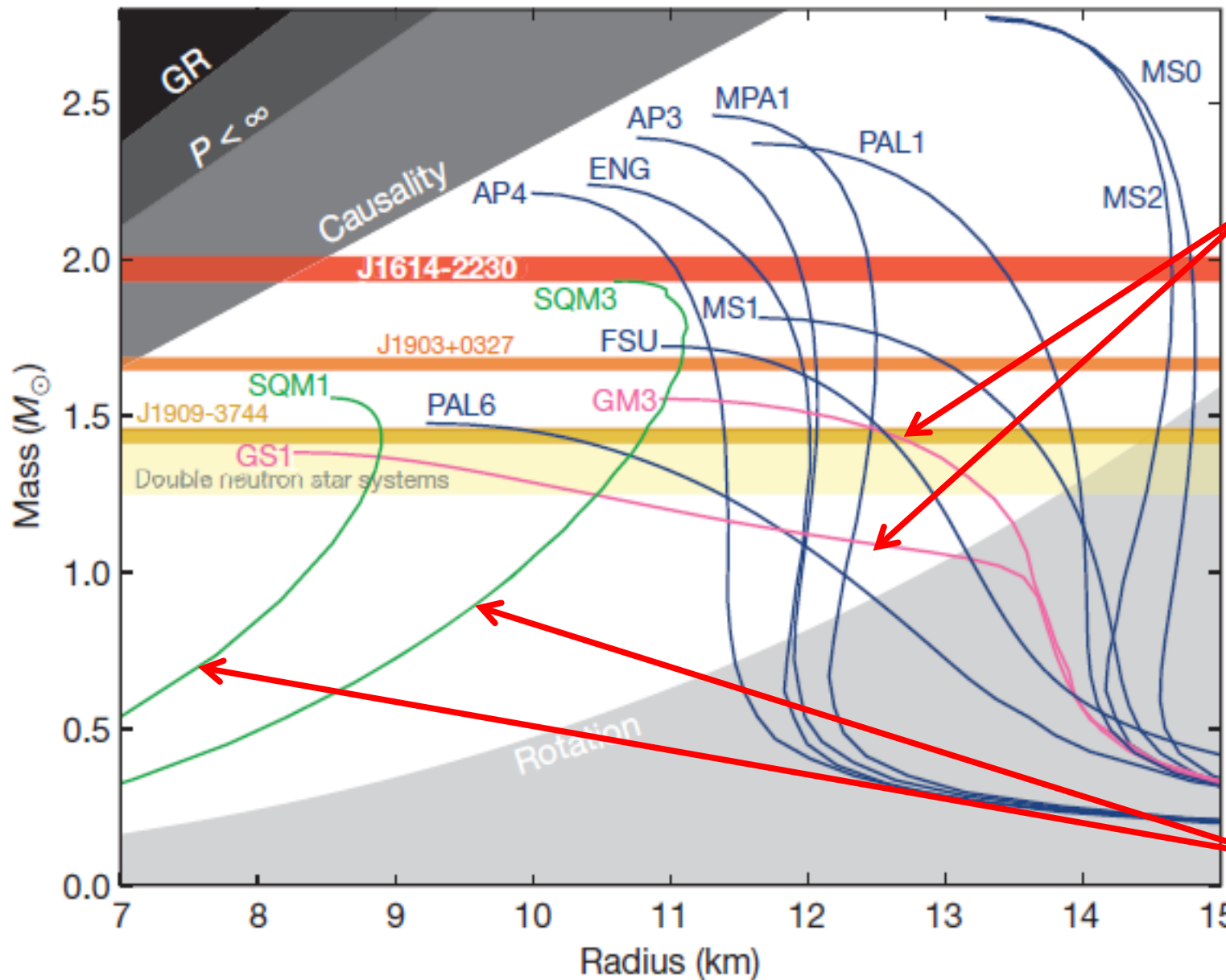
● Neutron Star Max. Mass ~ 1.4

The Judgement Day, Oct. 28, 2010.

*Can hyperon survive in $1.97 M_{\odot}$
neutron star ?*

$1.97 \pm 0.04 M_{\odot}$ Neutron Star

Demorest et al., Nature 467 (2010) 1081 (Oct.28, 2010).



ハイペロンを含むEOS

クォーク物質のみのEOS

Which type of EOSs are rejected ?

Rejected Hyperonic Matter EOS

Relativistic Mean Field model

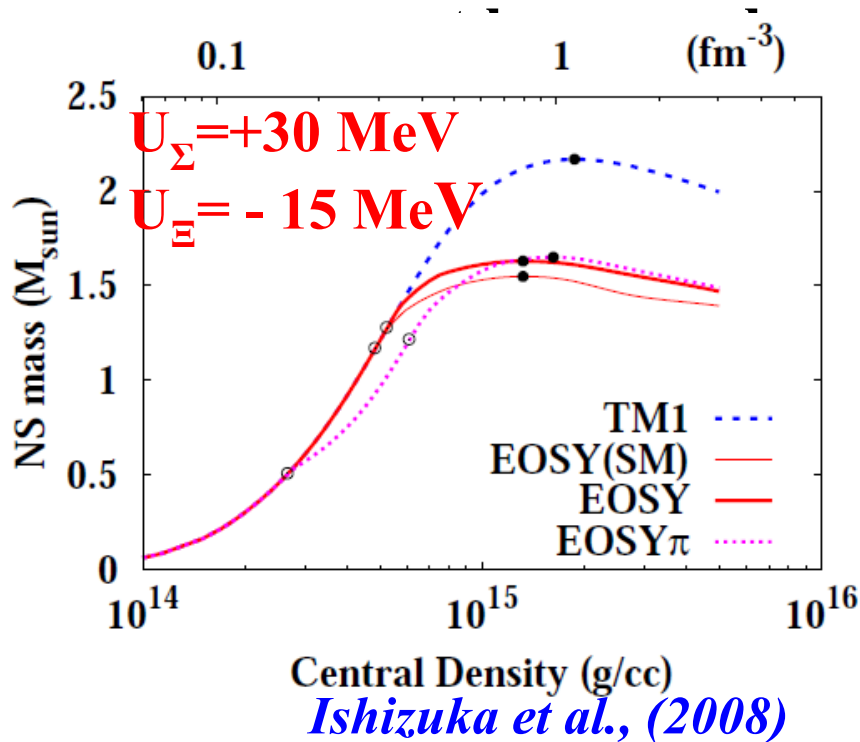
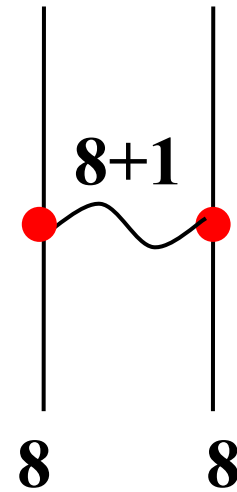
GM3: Glendenning & Moszkowski (1991)(npY)

GS1: Glendenning & Schaffner-Bielich (1999)(npK)

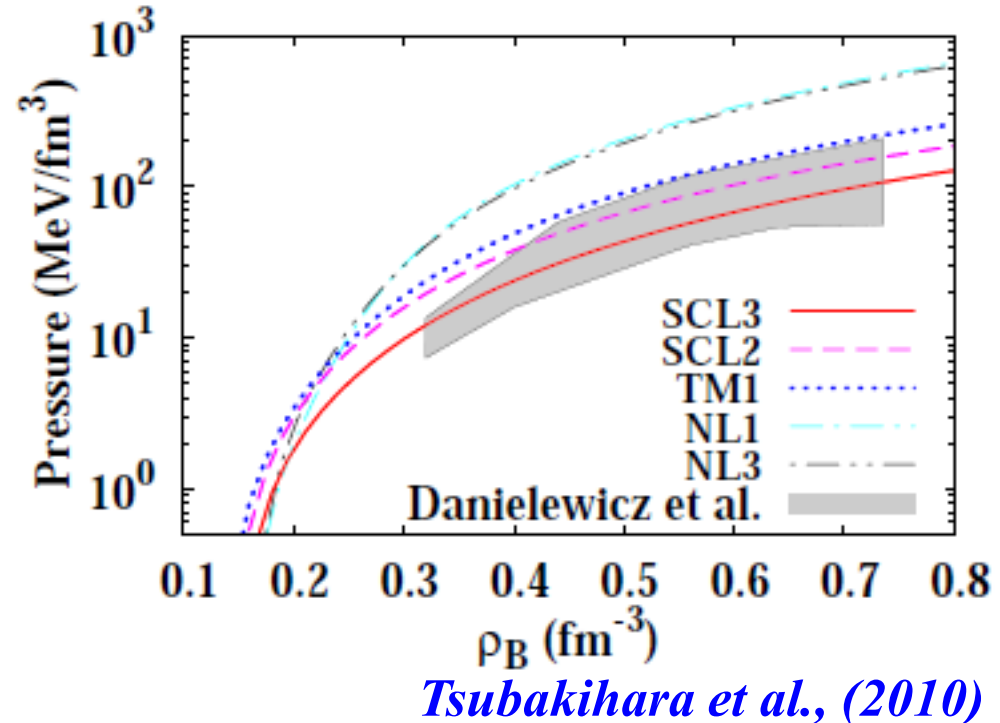
Coupling \sim Quark Counting ($g_{\omega Y}/g_{\omega N} \sim 2/3$)

Even with rel. effects, we cannot support $1.97M_{\odot}$

as I



& HIC data



Glendenning & Moszkowski (1991)

RMF with hyperons

n, p, Y, σ , ω , ρ / σ^3 , σ^4

Give $x_\sigma = g_{\sigma Y} / g_{\sigma N}$ and fix $x_\omega = g_{\omega Y} / g_{\omega N}$ to fit Λ separation energy.

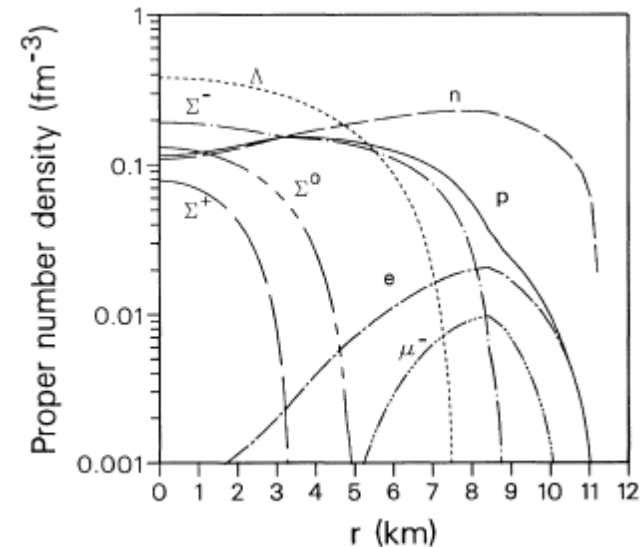
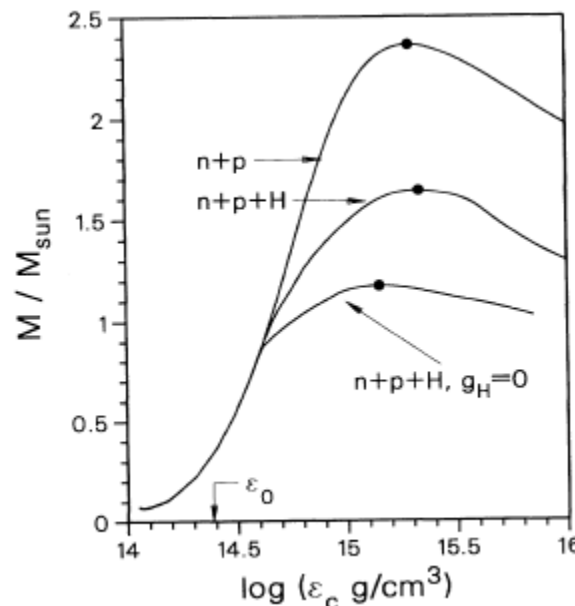
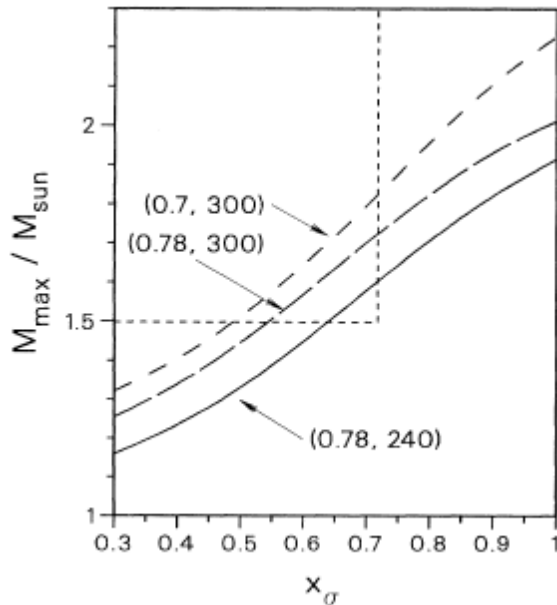
$x_\sigma = 0.6 \rightarrow m^*/m = 0.7, x_\omega = 0.653$

(similar to quark number counting result, $x = 2/3$)

TABLE I. Values of the hyperon-to-nucleon scalar and vector coupling that are compatible with the binding of -28 MeV for Λ hyperons in nuclear matter for two values of the nucleon (Dirac) effective mass at saturation density.

x_σ	$m^*/m = 0.7$	x_ω	$m^*/m = 0.78$
0.2	0.131		0.091
0.3	0.261		0.233
0.4	0.392		0.375
0.5	0.522		0.517
0.6	0.653		0.568
0.7	0.783		0.800
0.8	0.913		0.942
0.9	1.04		1.08
1	1.17		1.23

$M_{\text{max}} \sim 1.6 M_{\text{sun}}$



N.K.Glendenning, S.A.Moszkowski, PRL67('91)2414

How can we solve it ?

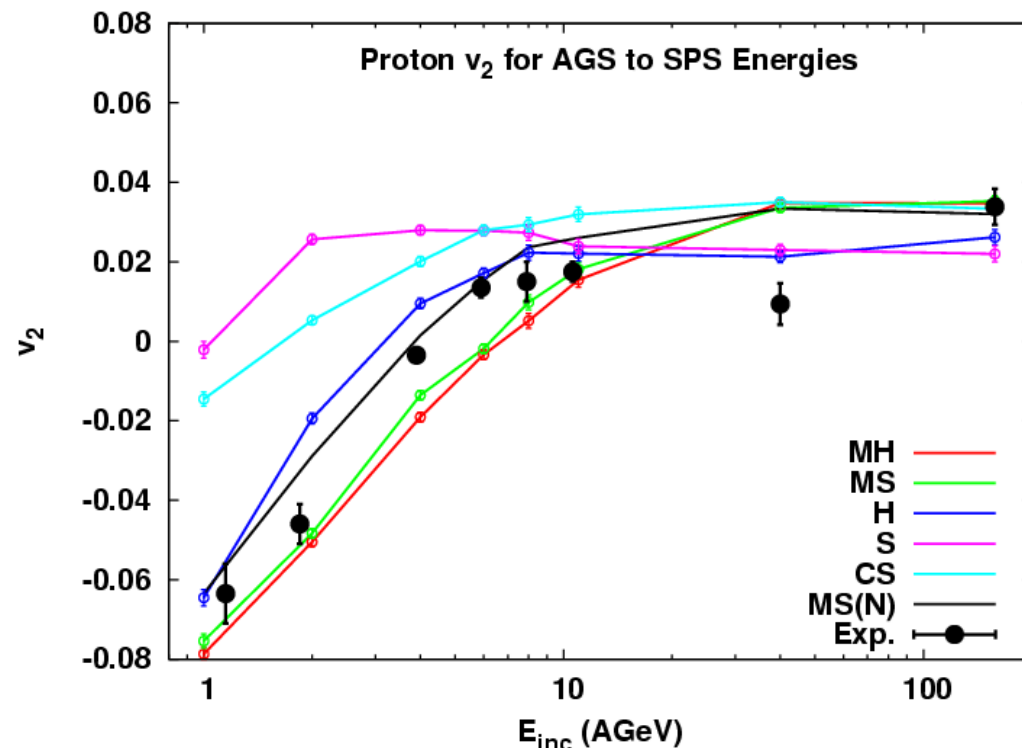
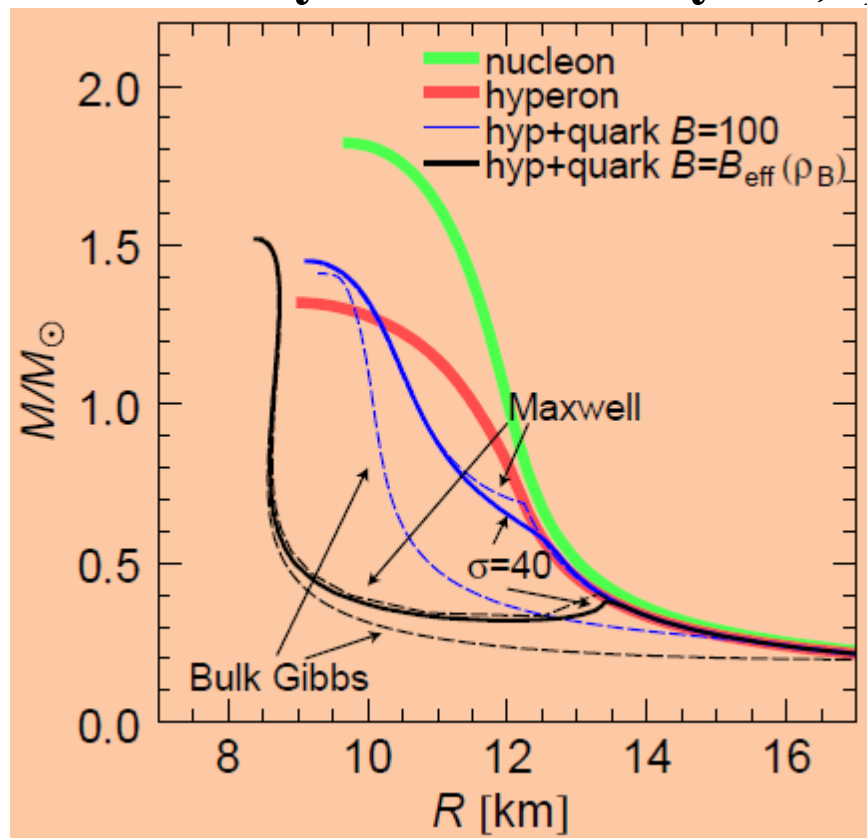
■ No Hyperons, No Kaons

→ How can it be consistent with YN interaction ?

■ Stiff nuclear matter EOS + transition to quark matter at small ρ_B

→ How can it be consistent with HIC data at AGS-SPS energies ?

■ Three-body force for baryons, quarks, ...



M. Isse, AO, N. Otuka, P. K. Sahu, Y. Nara, PRC72 ('05)064908

H.-J. Schulze, NFOCD10

Ohnishi @ TUS seminar, Dec.9, 2011, Tokyo Univ. of Science, Noda, Japan

$SU(3)_f$ “violating” coupling

■ Naïve RMF assumption = BM coupling follows $SU(3)_f$

■ Short range BB interaction comes from quark Pauli blocking + one-gluon exch.

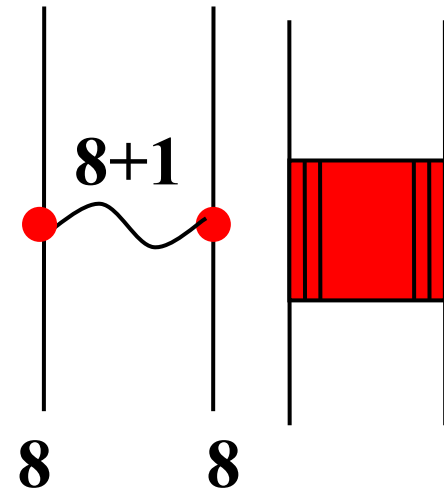
Oka, Yazaki; Faessler et al.; Fujiwara et al.; HAL QCD collab.

■ Short-range BB repulsion is sensitive to (S,T)

in the s-channel. When we include those interactions in the 8

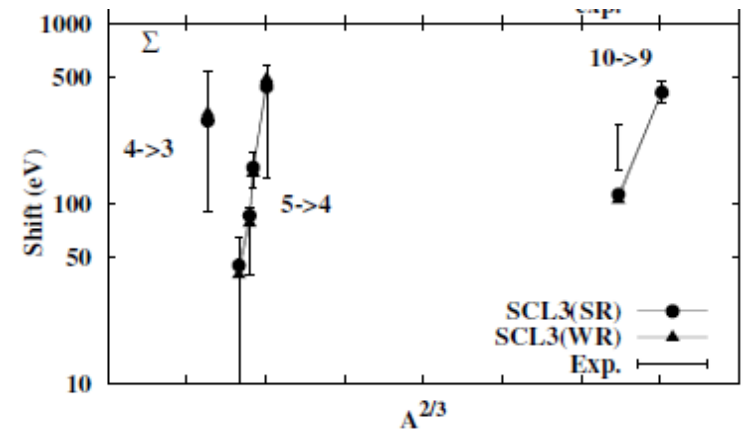
bosonized form, BM coupling violates $SU(3)_f$.

$$V = \sum_{\alpha,\beta} (\bar{\psi}\psi)_{\alpha} \Gamma_{\alpha\beta} (\psi\psi)_{\beta} \rightarrow -\frac{1}{2} \sum_{\alpha} m_{\alpha}^2 \omega_{\alpha}^2 + \sum_{\alpha} g_{\alpha} \omega_{\alpha} (\psi\Gamma\psi)_{\alpha}$$



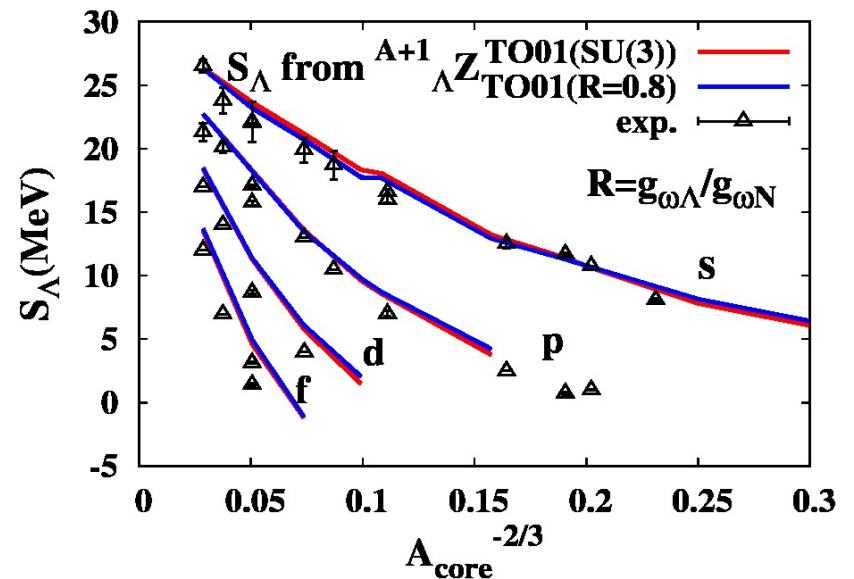
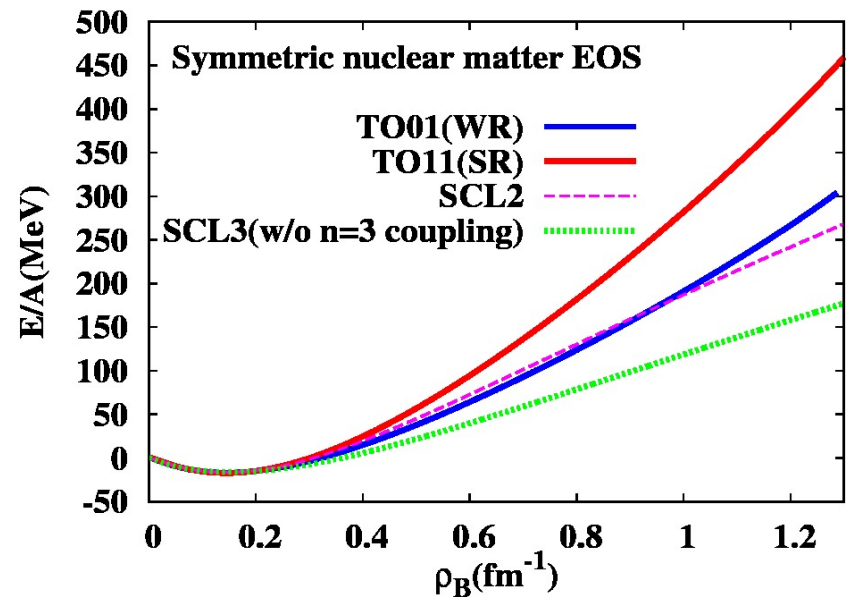
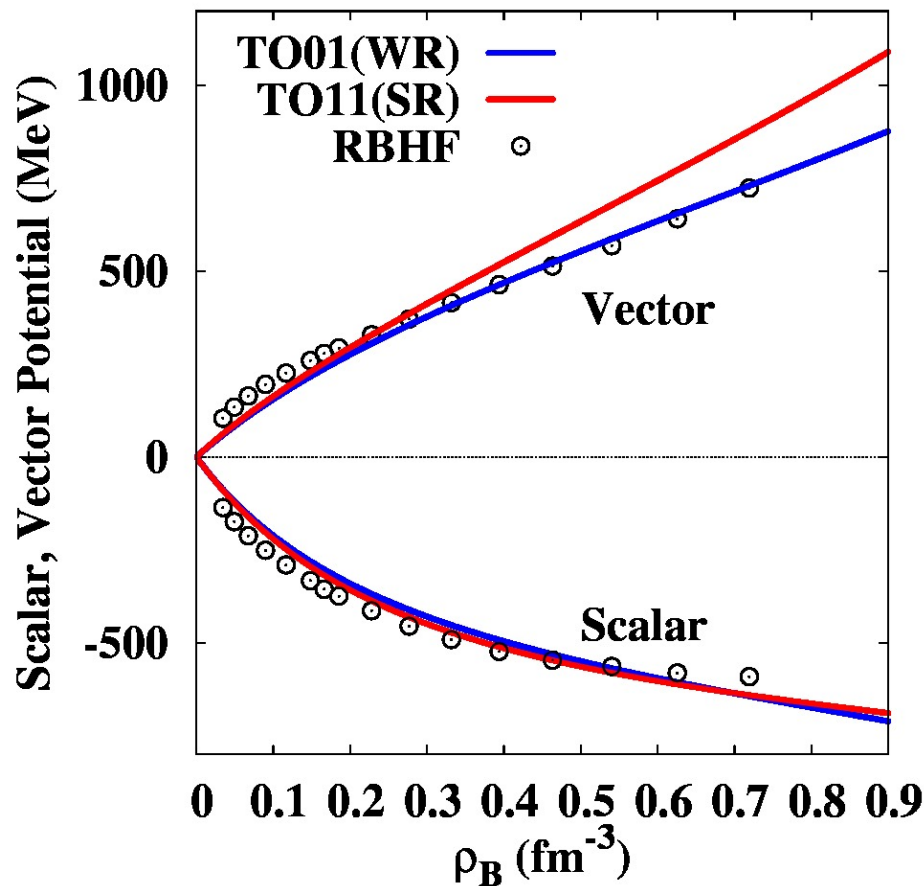
E.g., Σ atomic shift

$\rightarrow g_{\sigma\Sigma} \sim g_{\sigma\Sigma} (SU(3)) \times (0.2-0.3)$



Tsubakihara et al., (2010)

■ Nucleon vector potential $U_v(\rho)$
 in DBHF: Non-linear behavior in ρ_B .
 → EOS becomes gradually stiffer



RMF with 3BF + SU(3)_f “violation”

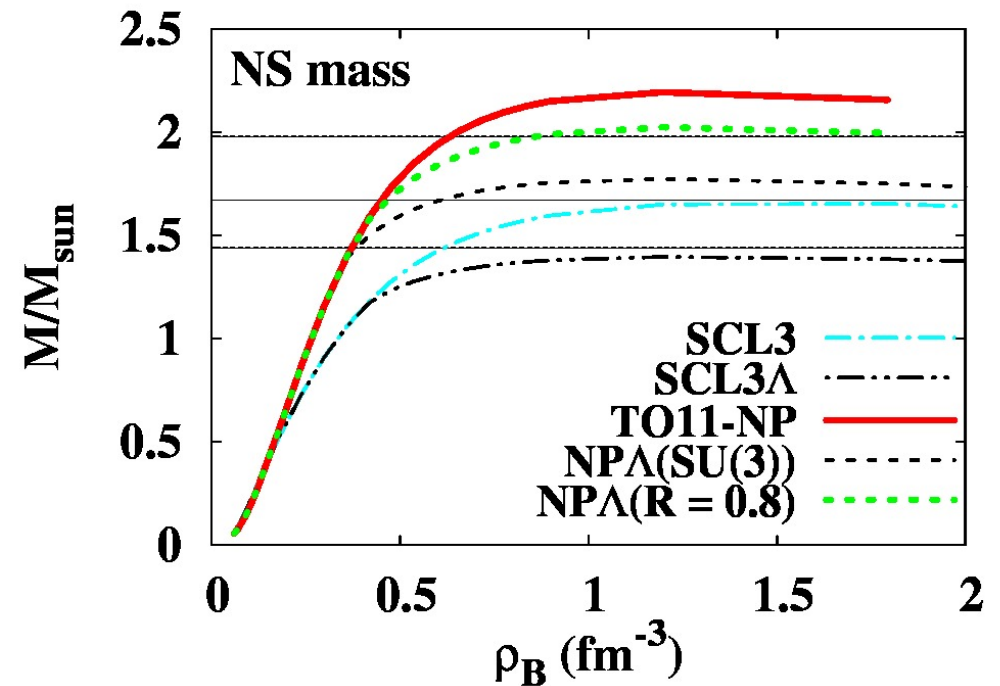
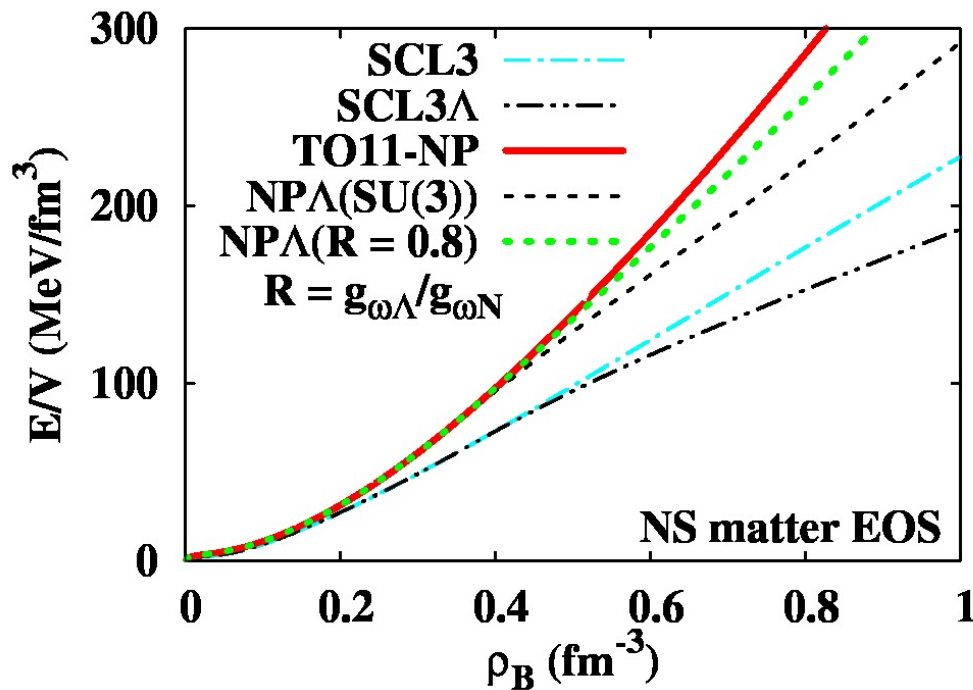
Tsubakihara, AO, in prep.

Two types of modification

3-baryon repulsion → EOS becomes stiff gradually at high density.
(Fitting meson mass (E325) and U_v in RBHF)

$R = g_{\omega\Lambda} / g_{\omega N} \sim 0.8$ ($\sim 2/3$ (SU(3)))

→ $M_{\max} \sim 2.02 M_{\odot}$ with hyperons ($\sim 1.4 M_{\odot}$ w/o 3BF, violation)



*Critical Point Sweep
during Black Hole Formation*

From Supernova Matter EOS to Phase Diagram

■ Supernova matter EOS

● Lattimer-Swesty EOS (Skyrme-type int. + Droplet)

J.M.Lattimer, F.D.Swesty, NPA535('91)331.

● Shen EOS (Relativistic Mean Field + Thomas Fermi)

H.Shen et al., NPA637('98)435;PTP100('98)1013.

● Ishizuka EOS (Shen EOS + Hyperons)

C. Ishizuka, AO, K.Tsubakihara, K.Sumiyoshi, S.Yamada, JPG 35 ('08)085201.

■ Does quark matter exist in compact stars ?

● Suggested in Supernovae: Warm(~ 20 MeV), mildly dense ($\sim 1.8 \rho_0$)

T. Hatsuda, MPLA2('87)805; I. Sagert et al., PRL102 ('09) 081101; Nishimura talk.

● Probable in Neutron Stars: Cold ($T \sim 0$), Dense ($\rho_B \sim 5 \rho_0$)

E.g. N. Glendenning, "Compact Stars"; F. Weber, Prog.Part.Nucl.Phys.54('05)193

● How about Black hole formation ?

M. Liebendorfer et al., ApJS 150('04)263; K. Sumiyoshi et al., PRL97('06) 091101;

K.Sumiyoshi, C.Ishizuka, AO, S.Yamada, H.Suzuki, ApJL690('09),L43;

K.Nakazato et al., ApJ, to appear [arXiv:1111.2900] (Nakazato, Poster)

Purpose and Methods

■ We compare (T, μ_B) during BH formation and QCD phase transition boundary by using

● v radiation Hydrodynamics (1D) for BH formation

Sumiyoshi et al., PRL97('06)091101;

◆ Shen EOS (npe μ) *Shen et al., NPA637('98)435; PTP100('98)1013*

◆ Grav. collapse of $40 M_{\text{sun}}$ star with WW95 initial condition.
S.E. Woosley, T.A. Weaver, ApJS 101 ('95) 181.

● Chiral Effective Models for phase boundary and Critical Point

◆ NJL (Nambu, Jona-Lasinio), PNJL (Polyakov loop extended NJL), PNJL with 8 quark int., PQM (Pol. loop ext. quark meson) models
Nambu, Jona-Lasinio('61); Hatsuda, Kunihiro('94), Fukushima('04); Ratti, Thaler, Weise('06); Roessner et al.('07); Kashiwa, Kouno, Matsuzaki, Yahiro('08), Schaefer, Pawłowski, Wambach ('07), Skokov et al. ('10).

◆ Vector coupling: unknown \rightarrow compare results with $G_v/G_s=0, 0.2$

◆ Flavor SU(2) models are considered.

QCD phase diagram in Asymmetric Matter

■ Characteristic features of Compact Star Matter

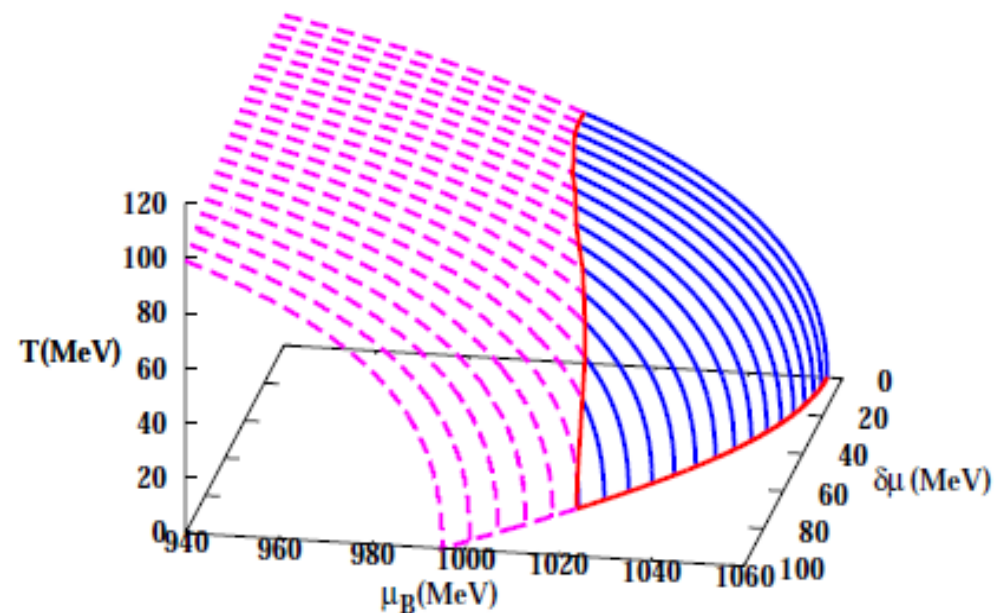
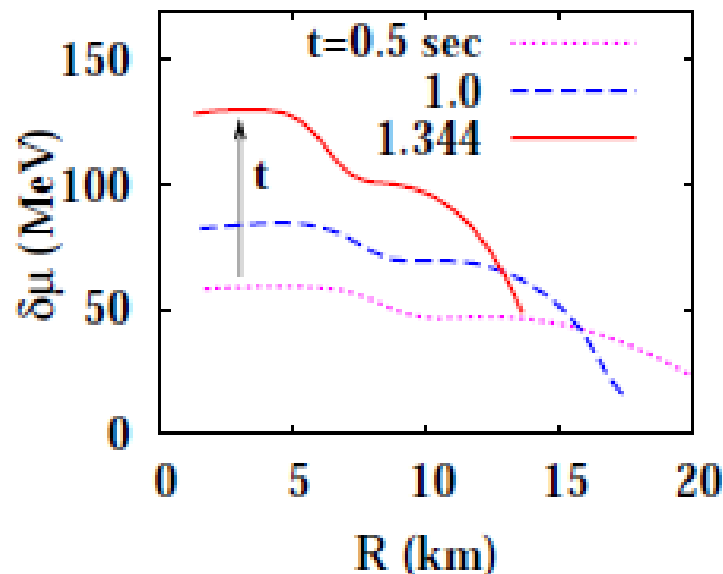
- Hot and/or Dense

- Unbalanced n and p yields (Isospin Asymmetric Matter)

Isospin chemical potential $\delta\mu = (\mu_n - \mu_p)/2 = (\mu_d - \mu_u)/2 > 0$

- T_{CP} (critical point T) decreases at finite $\delta\mu$

- Decrease of effective number of flavors

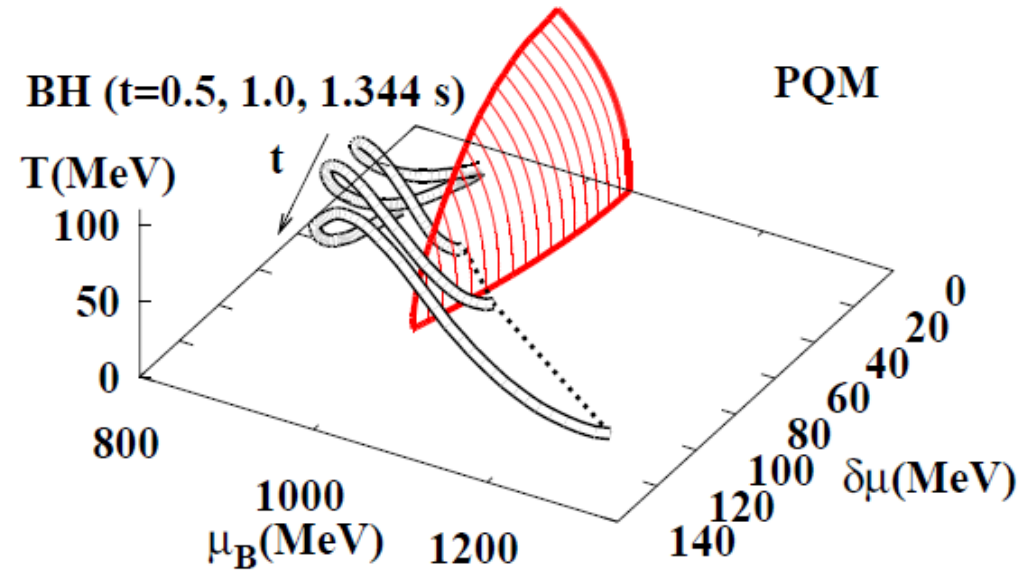
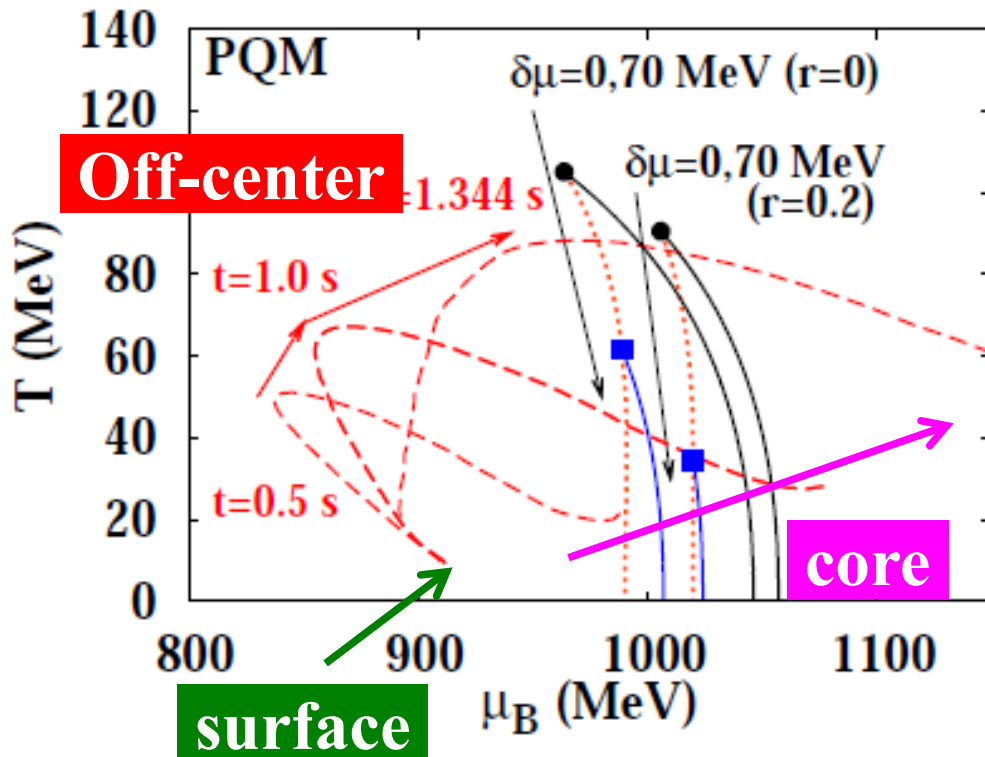


Ueda et al, in preparation

AO, Ueda, Nakano, Ruggieri, Sumiyoshi ('11)

How is quark matter formed during BH formation ?

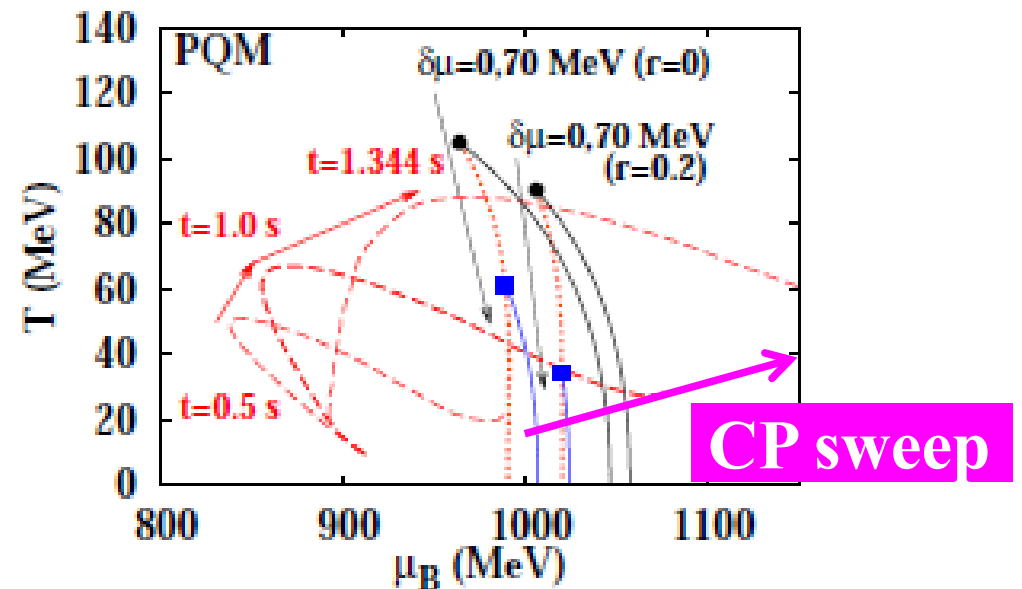
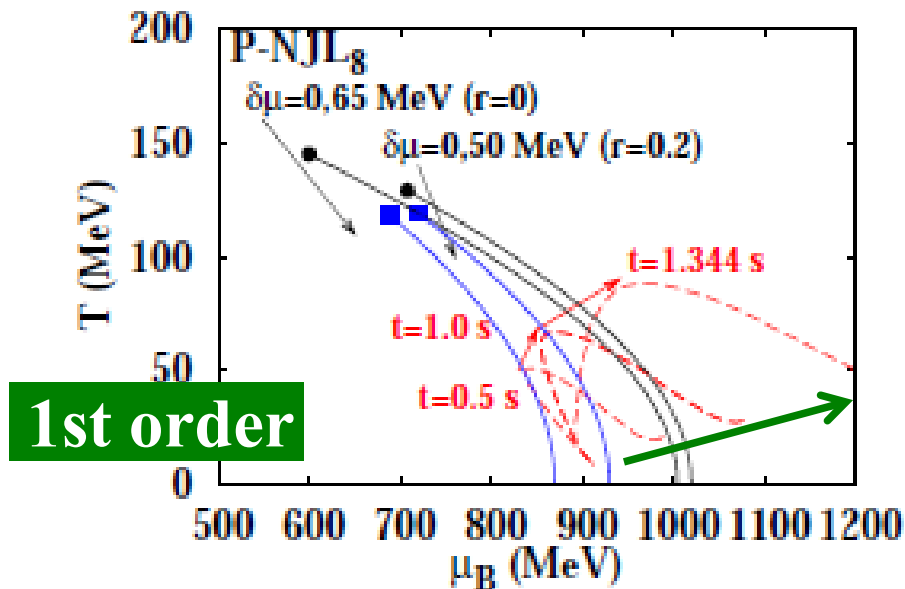
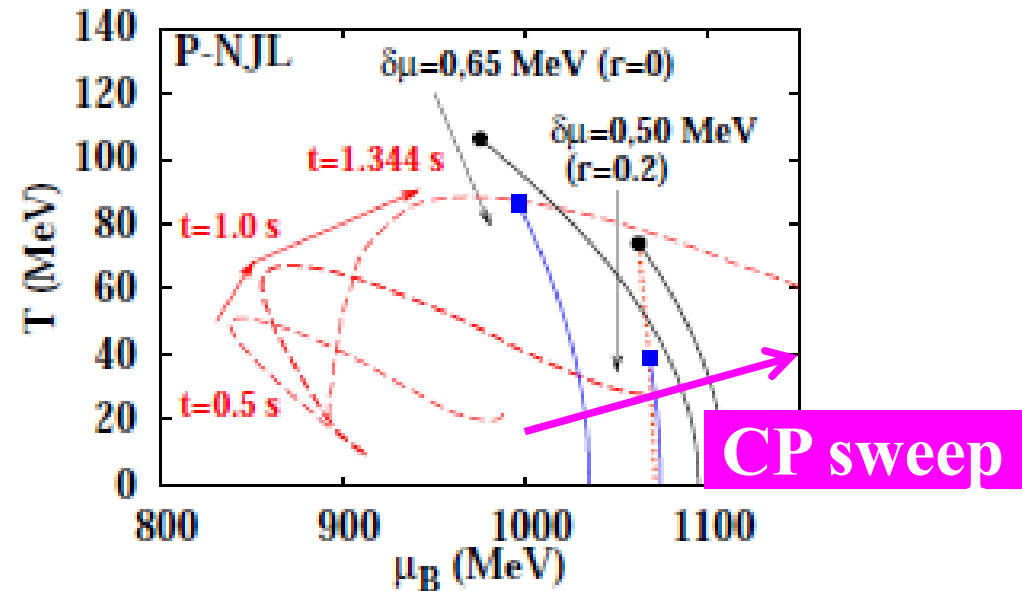
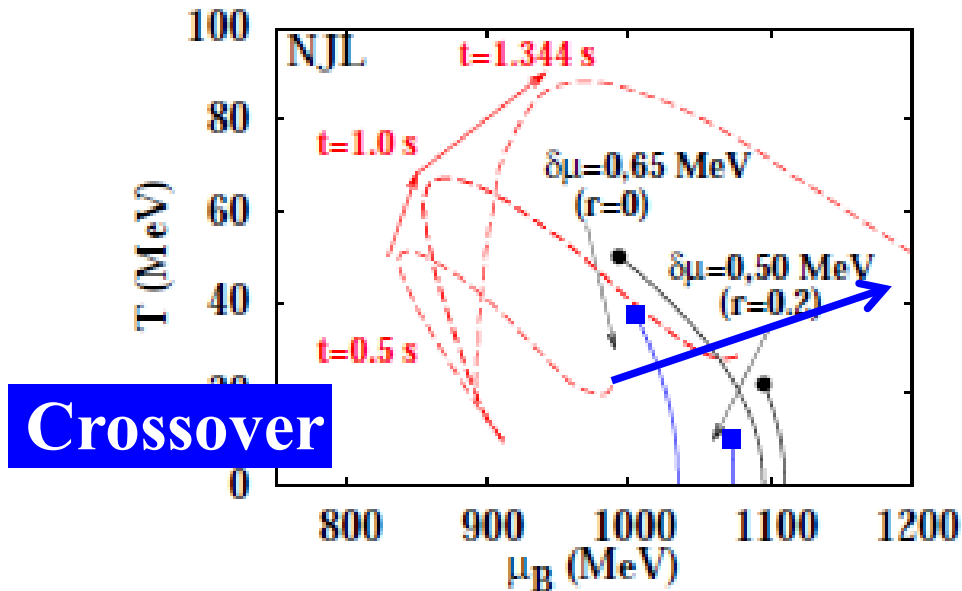
- Highest $\mu_B \sim 1300 \text{ MeV} > \mu_c$ (1000-1100 MeV in eff. models)
 → *Quark matter is formed before BH formation*
- Highest $T \sim 90 \text{ MeV} > T_{CP}$ (at $\delta\mu \sim 50 \text{ MeV}$)
 Core evolves below CP, Off-center goes above CP → *CP sweep*
- Convenient to consider 3D phase diagram ($T, \mu_B, \delta\mu$)



1.344 s = Just before BH formation

How is quark matter formed during BH formation ?

Model dependence to form quark matter → Three ways



Swept Region of Phase Diagram during BH formation

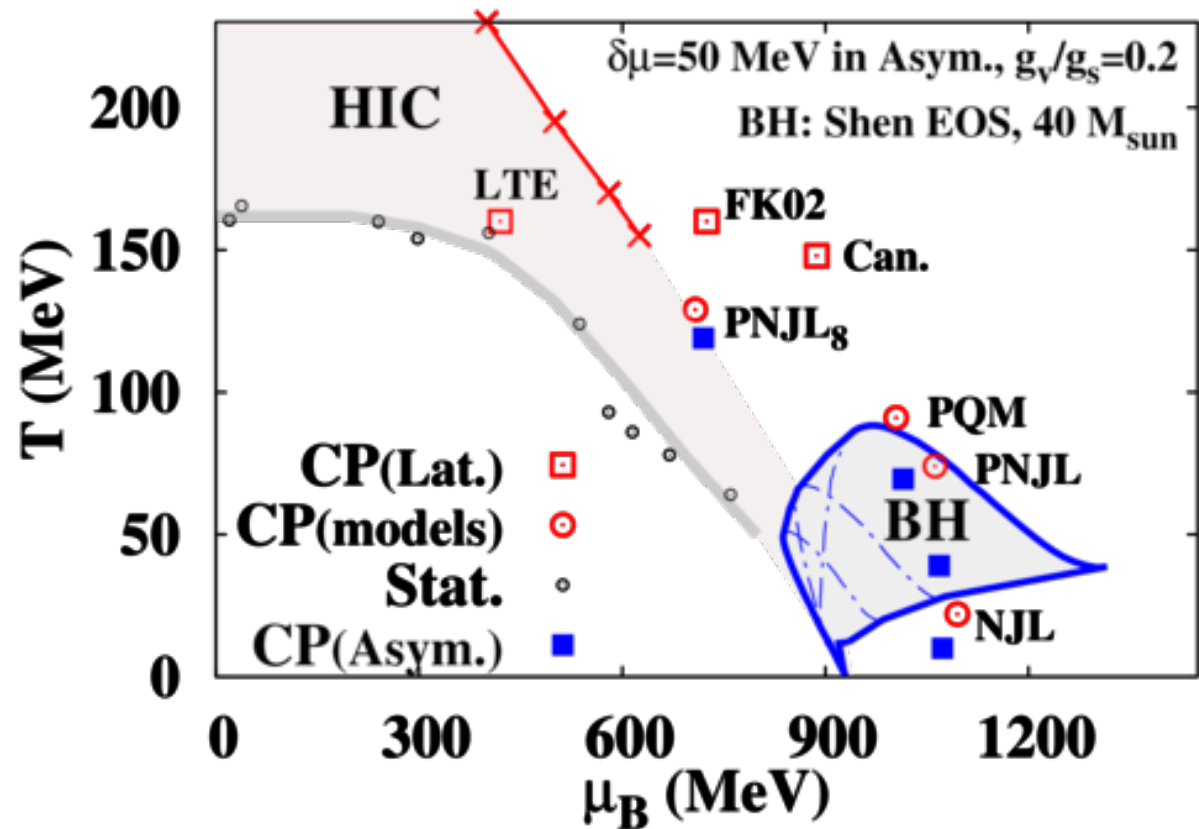
■ CP location
in Symmetric Matter

● Lattice QCD
 $\mu_{CP} = (400-900)$ MeV

● Effective models
 $\mu_{CP} = (700-1050)$ MeV

■ CP in Asymmetric Matter
(E.g. $\delta\mu = 50$ MeV)

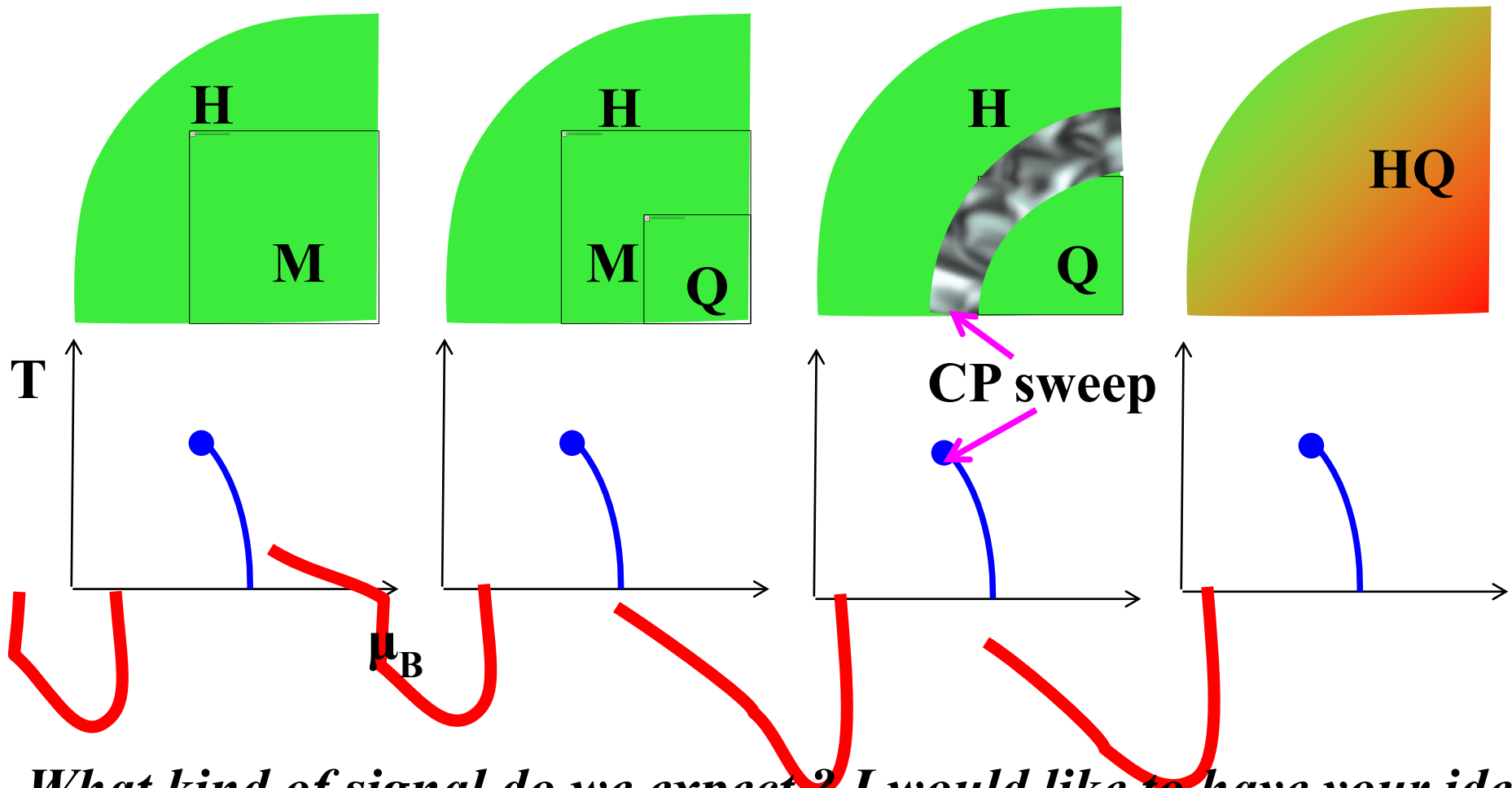
● T_{CP} decreases at finite $\delta\mu$.
→ Accessible (T, μ_B) region
during BH formation



M.A.Stephanov, Prog.Theor.Phys.Suppl.153 ('04)139;
FK02:Z. Fodor, S.D.Katz, JHEP 0203 (2002) 014
LTE:S. Ejiri et al., Prog.Theor.Phys.Suppl. 153 (2004) 118;
Can: S. Ejiri, PRD78 (2008) 074507
Stat.:A. Andronic et al., NPA 772('06)167

What happens at CP sweep ?

- Large density fluctuation is expected around CP.
- Three layers (hadron, mixed, quark) merges to be one at a time.



What kind of signal do we expect? I would like to have your idea

Summary

- **Compact Astrophysical Phenomena**
- **Possibilities of various state of matter, New observations, New calculations**
- **Dense baryonic matter Equation of State from Relativistic Mean Field**
- **Successfully applied to Compact Astrophysical Phenomena.**
- **Needs further studies: Higher order terms, pion contribution, E_{sym}, Hyperons,**
- **Standard RMF with hyperons cannot support 1.97 M_⊙ neutron star.**
- **Various data / DBHF results can be fitted in RMF.**
- **Vector Coupling ~ SU(3)_f, linear BM coupling (BMB)**
- **RMF with 3BF + SU(3)_f “violation” may help to support the heavy NS.
→ Can we support 1.97 M_⊙ neutron star with hyperons ? Open problem.**
- **Quark matter formation and critical point sweep may take place in neutron stars and black hole formation.**

Thank you for your attention !

核物理からみた状態方程式の面白さ・問題点

■なぜ面白い？

- 高温・高密度物質が「平衡状態」で作られている
→ 相転移現象がみえる？

- 物質の構成要素は？
→ 原子核・ハドロン・クォークの3階層の関連

■問題点

●高密度側

- ◆非圧縮率、対称エネルギーの密度依存性 → Unknown
- ◆核物質中のハドロン・ポテンシャル
- ◆クォーク物質への相転移の性質は？

●低密度側

- ◆物質の構成粒子は？ Fermi gas ? Nuclei ? Pasta ?
- ◆「原子核熱的分布」と「一様物質」をスムーズにつないだEOS はない。

Theories/Models for Nuclear Matter EOS

■ Ab initio Approaches to Nuclear Matter

→ LQCD, Variational, GFMC, BHF(G-Matrix), DBHF, ...

● LQCD-MC: Not (yet) applicable to cold dense matter, $A \leq 4$

SC-LQCD: Nuclear matter does not bound

● Variational, BHF: Need phen. 3-body repulsion to reproduce saturation point.

● GFMC: Limited to be $A \leq 12$.

● DBHF: Good, but E/A is not enough. Not yet extensively investigated.

■ → Not easy to handle, Not yet satisfactory for phen. purposes

■ Mean Field Models (~ Nuclear Density Functional approach)

● Skyrme Hartree-Fock(-Bogoliubov)

◆ Nuclear Mass is very well explained (HFB, Total B.E. $\Delta E \sim 0.6$ MeV)

◆ Causality is violated at very high densities.

● Relativistic Mean Field

◆ Relativistic, Meson-Baryon coupling, Meson self-energies

EOS in Dirac-Brueckner-Hartree-Fock

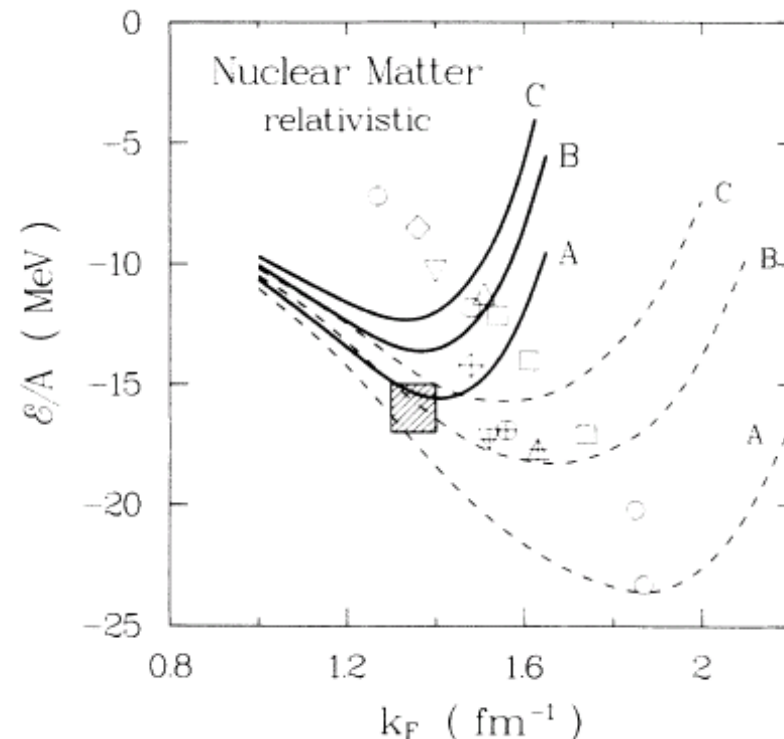
R. Brockmann, R. Machleidt, PRC42('90),1965

■ Non Relativistic Brueckner Calculation

→ Nuclear Saturation Point cannot be reproduced (Coester Line)

■ Relativistic Approach (DBHF)

→ Relativity gives additional repulsion, leading to successful description of the saturation point.



Bruckner-Hartree-Fock

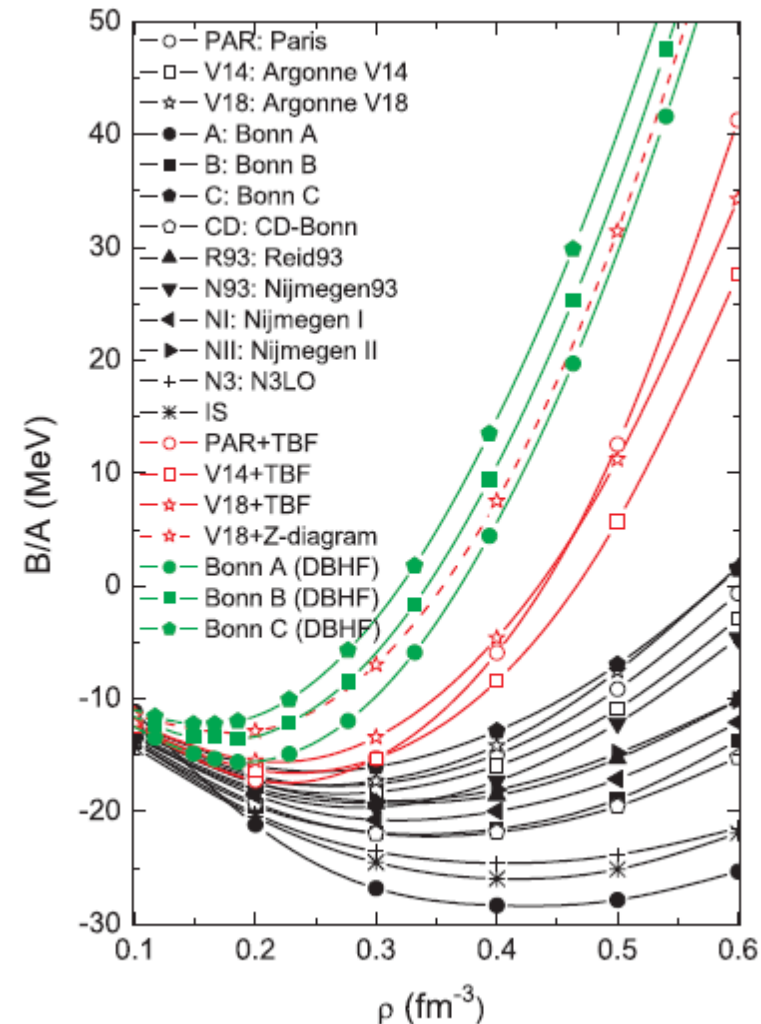
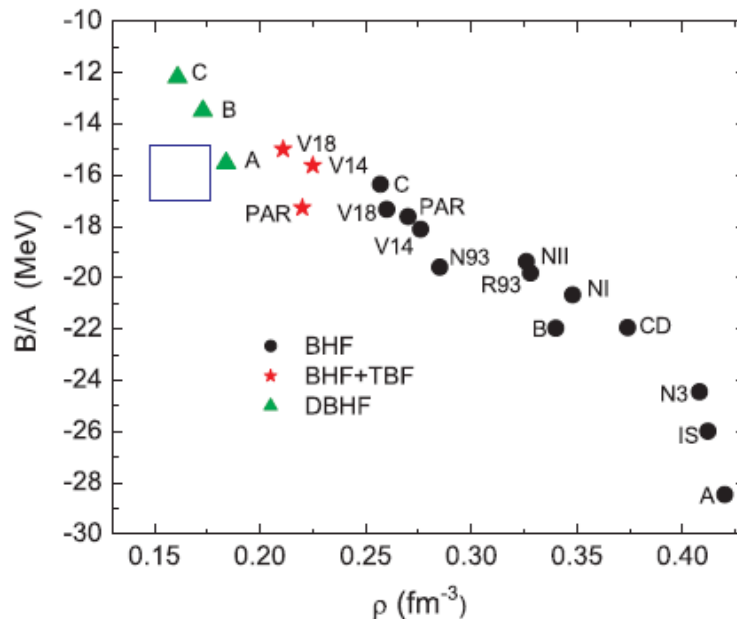
Self-consistent treatment of

Effective interaction (G-matrix) in the Bruckner Theory

and Single particle energy from G-matrix

Need 3-body force to reproduce saturation point.

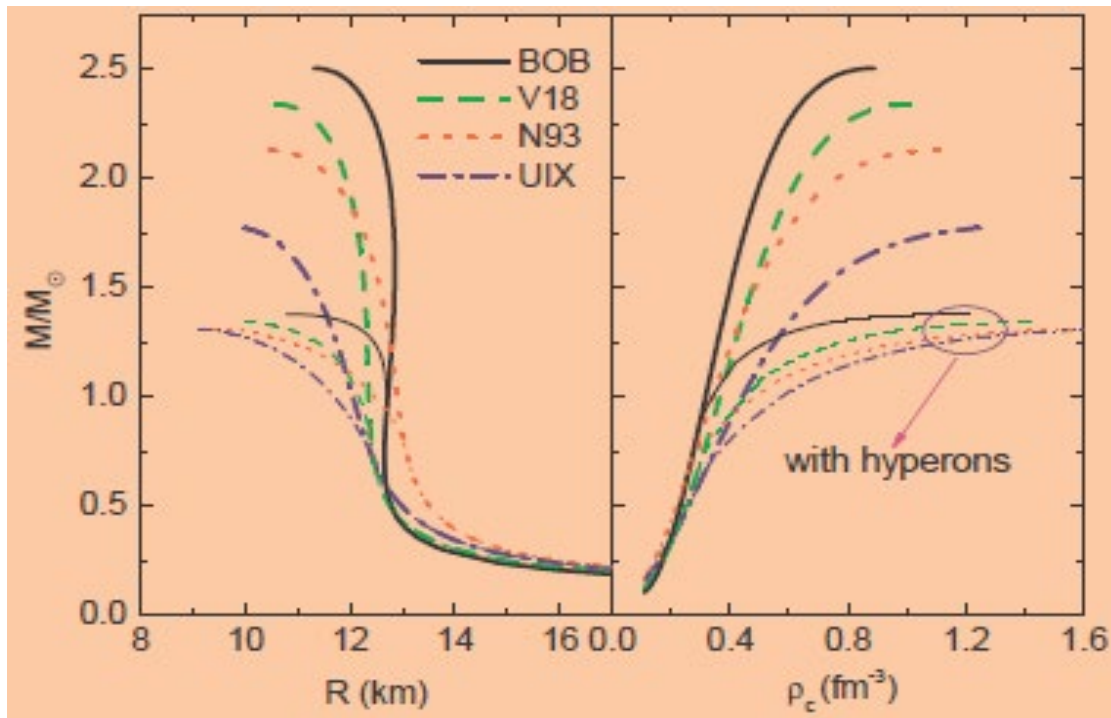
→ FY type 2 π exchange
+ phen. or Z-diagram



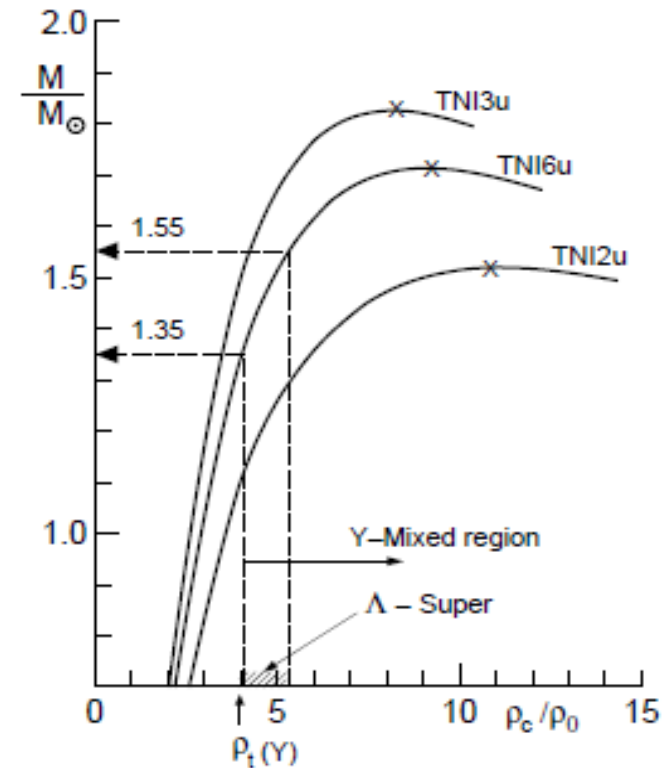
Z.H.Li, U. Lombardo, H.-J. Schulze, W. Zuo, L. W. Chen, H. R. Ma, PRC74('06)047304.

Bruckner-Hartree-Fock theory with Hyperons

- Microscopic G-matrix calculation with realistic NN, YN potential and microscopic (or phen.) 3N force (or 3B force).
- Interaction dep. (V18, N93, ...) is large → Need finite nuclear info.
E.Hiyama, T.Motoba, Y.Yamamoto, M.Kamimura / M.Tamura et al.
- NS collapses with hyperons w/o 3BF.



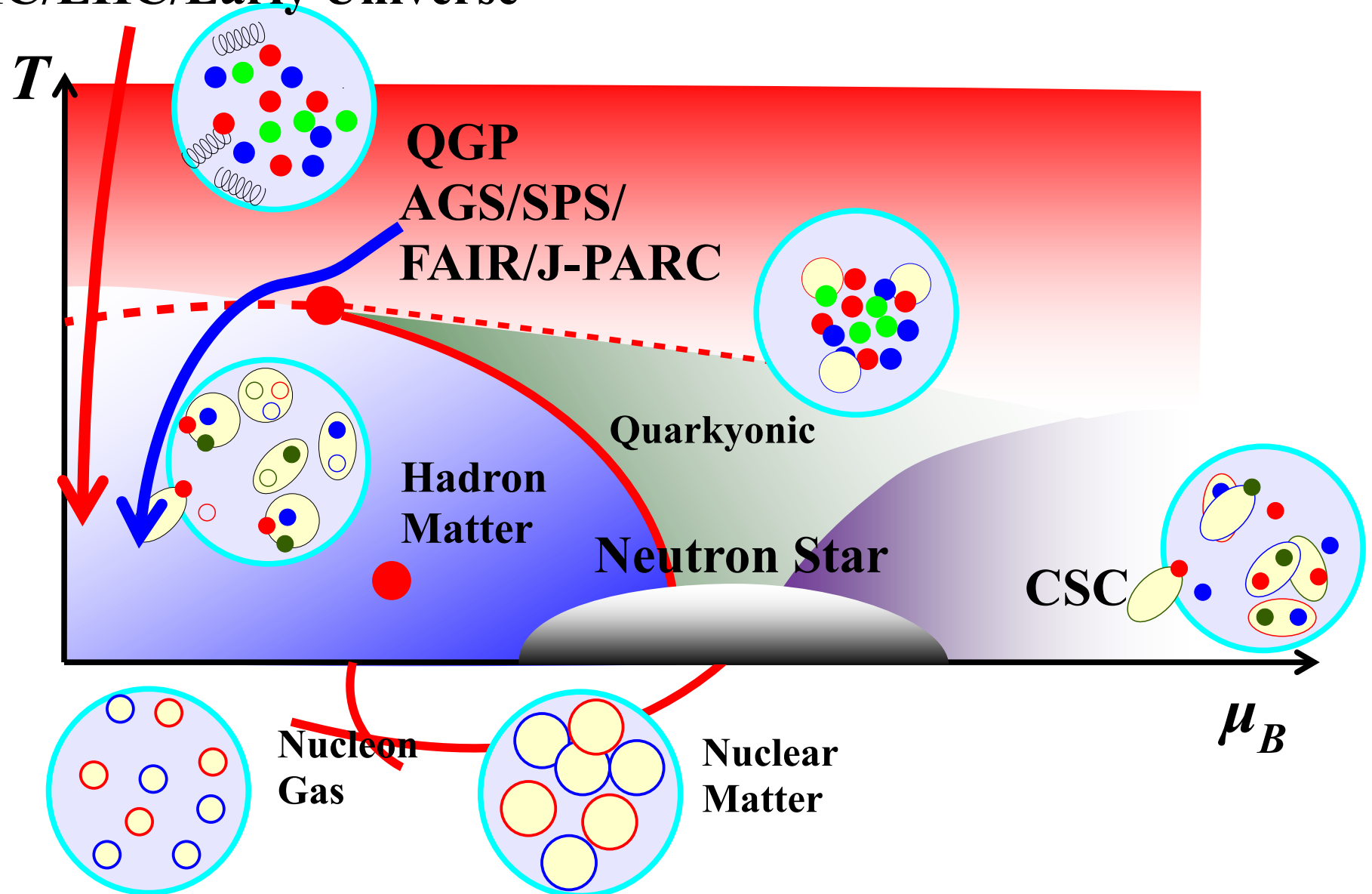
H.J.Schulze, A.Polls, A.Ramos, I.Vidana, PRC73('06),058801.



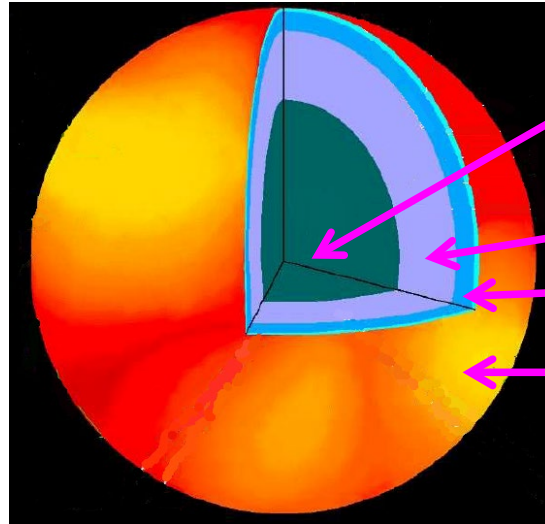
S. Nishizaki, T. Takatsuka, Y. Yamamoto, PTP108('02)703.

QCD Phase Diagram

RHIC/LHC/Early Universe



Neutron Star Composition

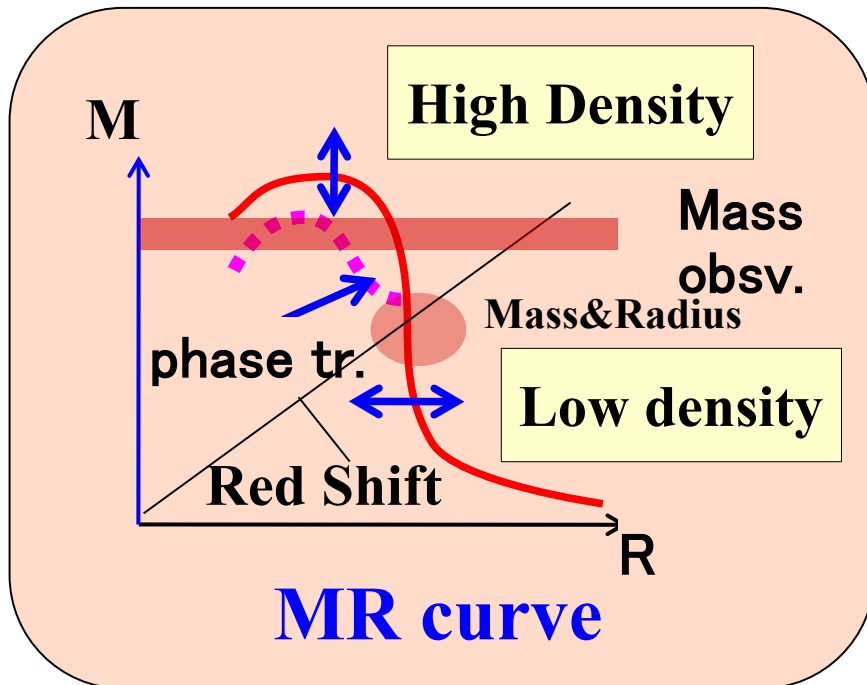


Hyperons, mesons, quarks

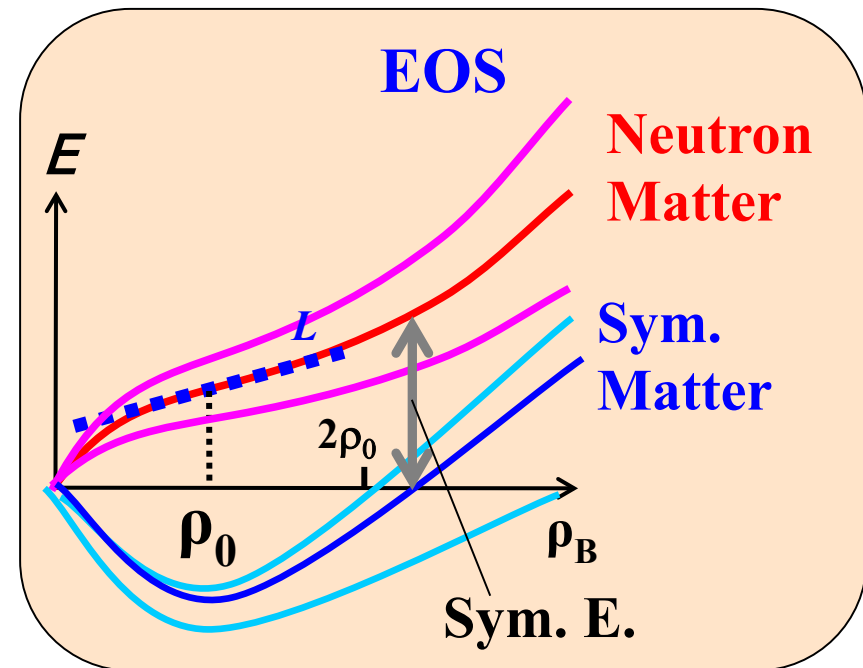
Asym. nuclear matter+elec.+ μ

Nuclei+neutron gas+elec.

Nuclei + elec.



TOV eq.



Hyperons in Dense Matter

- Hyperons are HOT now !
- What makes NS matter core ? Nucleons ? Quarks ? Hyperons ?
- How can we suppress hyperon appearance in NS ?
or How can hyperonic matter be so stiff ?
or Which inter-quark interaction supports $1.97 M_{\odot}$ NS ?

*We stick to hyperonic matter (rather than quark matter),
and discuss possible mechanism*

to stiffen the EOS at high density

$1.97 \pm 0.04 M_{\odot}$ Neutron Star

■重い中性子星(2倍の太陽質量)の観測

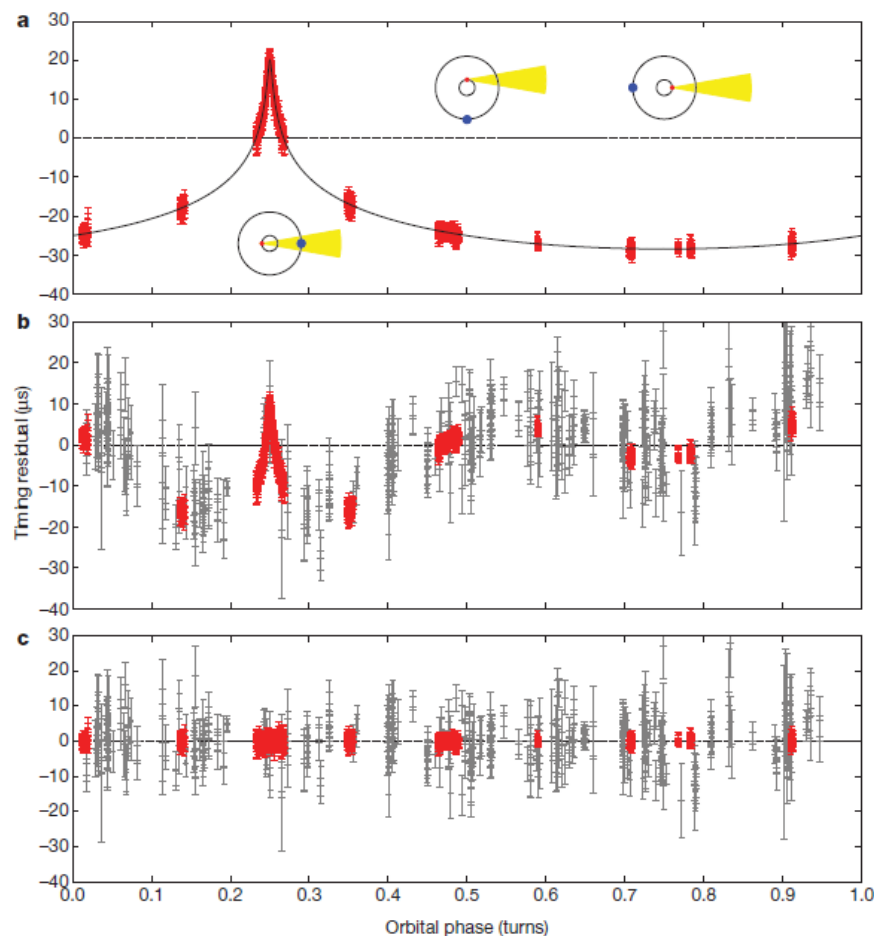
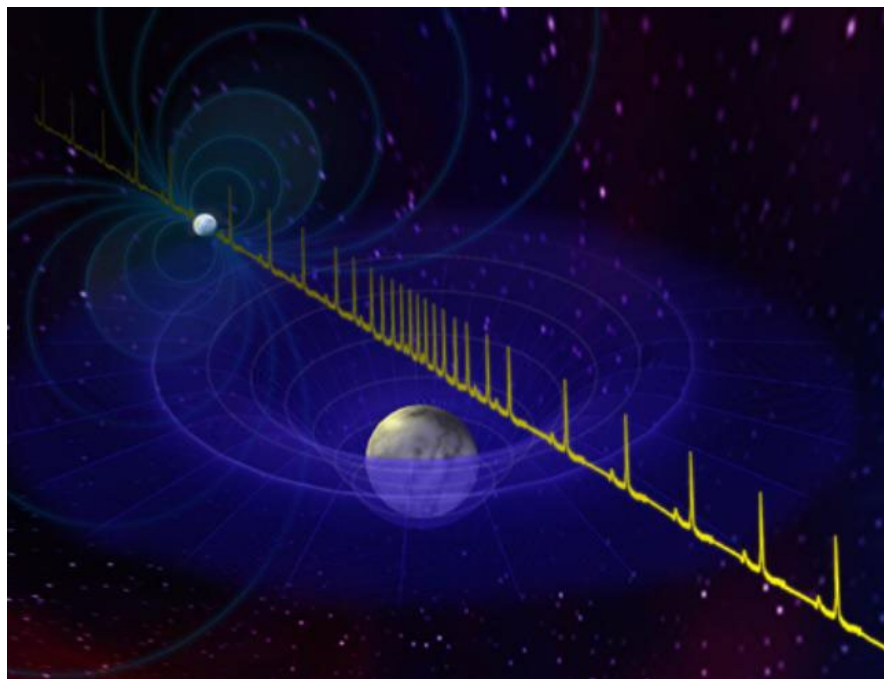
Demorest et al., Nature 467 (2010) 1081 (Oct.28, 2010).

PSR J1614-2230 (NS-WD binary), $1.97 \pm 0.04 M_{\text{sun}}$

●一般相対性理論に基づく時間の遅れ(Shapiro delay)による質量決定(模型によらない)

●幸運な公転面の向き

●美しい観測結果



From Supernova Matter EOS to Phase Diagram

■ Supernova matter EOS

● Lattimer-Swesty EOS (Skyrme-type int. + Droplet)

J.M.Lattimer, F.D.Swesty, NPA535('91)331.

● Shen EOS (Relativistic Mean Field + Thomas Fermi)

H.Shen et al., NPA637('98)435;PTP100('98)1013.

● Ishizuka EOS (Shen EOS + Hyperons)

C. Ishizuka, AO, K.Tsubakihara, K.Sumiyoshi, S.Yamada, JPG 35 ('08)085201.

■ Does quark matter exist in compact stars ?

● Suggested in Supernovae: Warm(~ 20 MeV), mildly dense ($\sim 1.8 \rho_0$)

T. Hatsuda, MPLA2('87)805; I. Sagert et al., PRL102 ('09) 081101; Nishimura talk.

● Probable in Neutron Stars: Cold ($T \sim 0$), Dense ($\rho_B \sim 5 \rho_0$)

E.g. N. Glendenning, "Compact Stars"; F. Weber, Prog.Part.Nucl.Phys.54('05)193

● How about Black hole formation ?

M. Liebendorfer et al., ApJS 150('04)263; K. Sumiyoshi et al., PRL97('06) 091101;

K.Sumiyoshi, C.Ishizuka, AO, S.Yamada, H.Suzuki, ApJL690('09),L43;

K.Nakazato et al., ApJ, to appear [arXiv:1111.2900] (Nakazato, Poster)

Chiral Symmetry

■ Fundamental symmetry of massless QCD, and its spontaneous breaking generates hadron masses.

Nambu, Jona-Lasinio ('61)

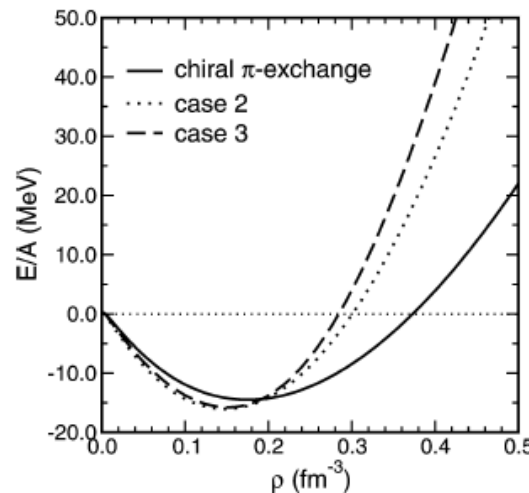
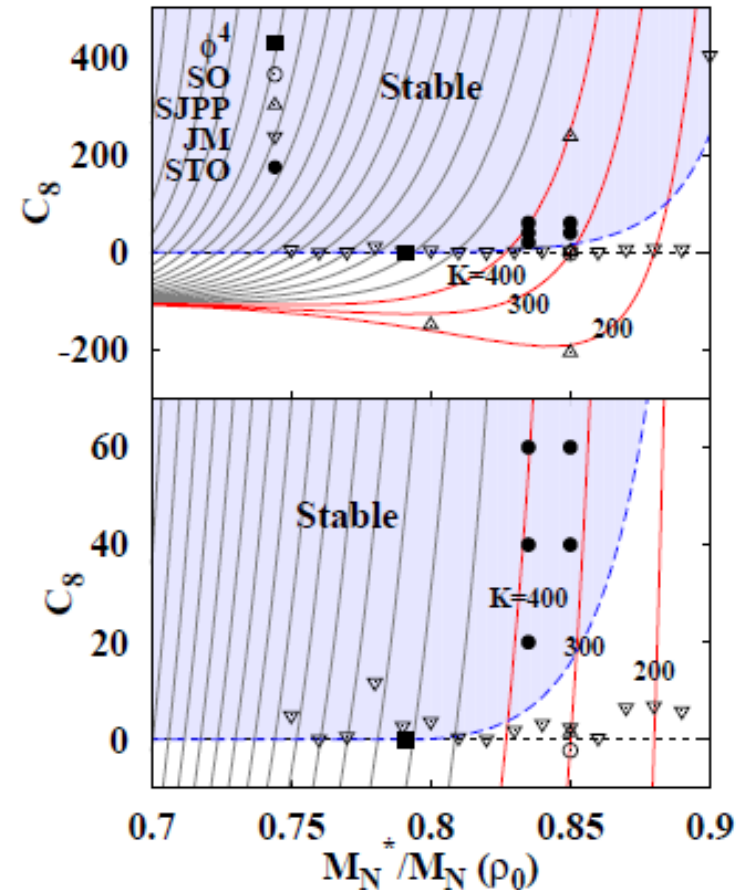
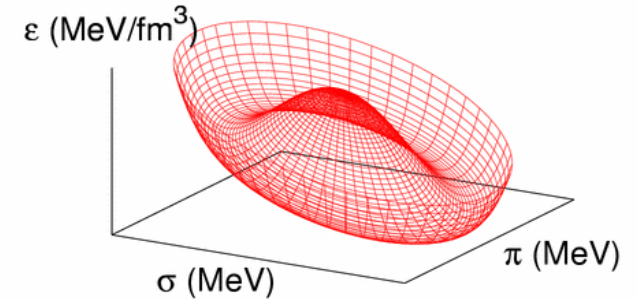
■ Many of the linear σ models are unstable against finite density (chiral collapse).

→ Log type chiral potential

Sahu, Tsubakihara, AO('10), Tsubakihara, AO('07)

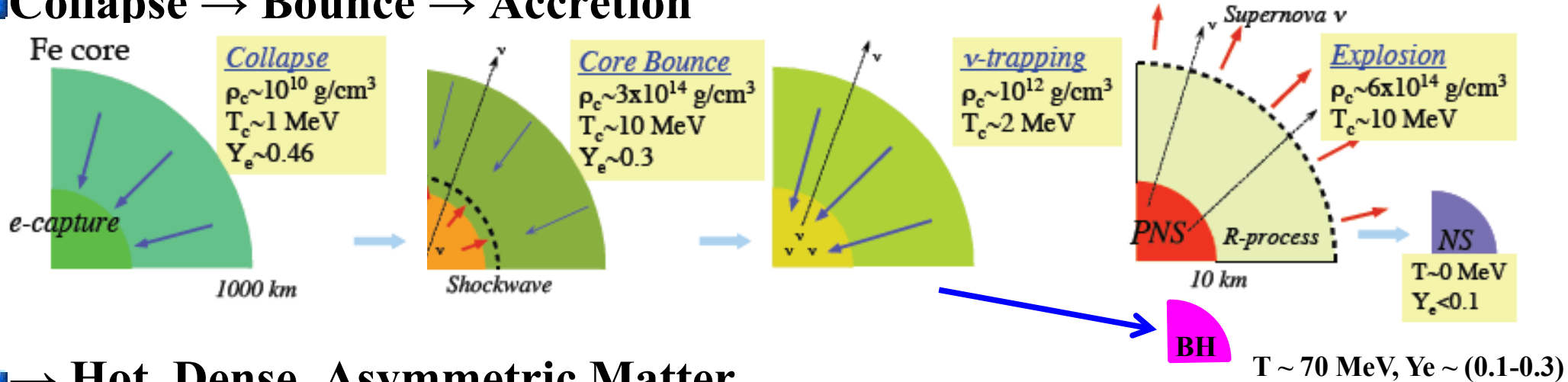
■ Non-linear representation (chiral pert.) leads to density dependent coupling from pion loops.

*Kaiser, Fritsch, Weise ('02),
Finelli, Kaiser, Vretener
Weise ('04)*



Dynamical Black Hole Formation

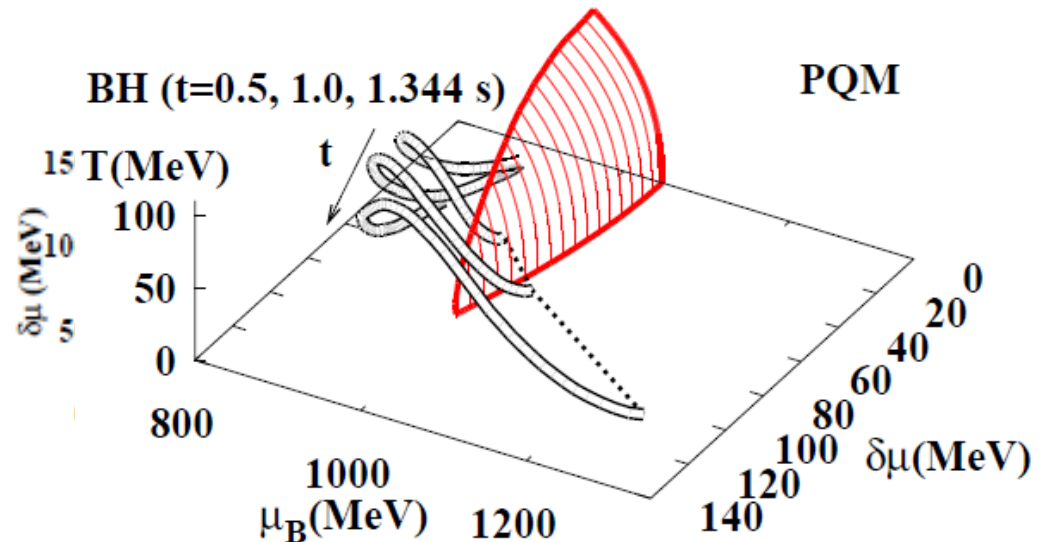
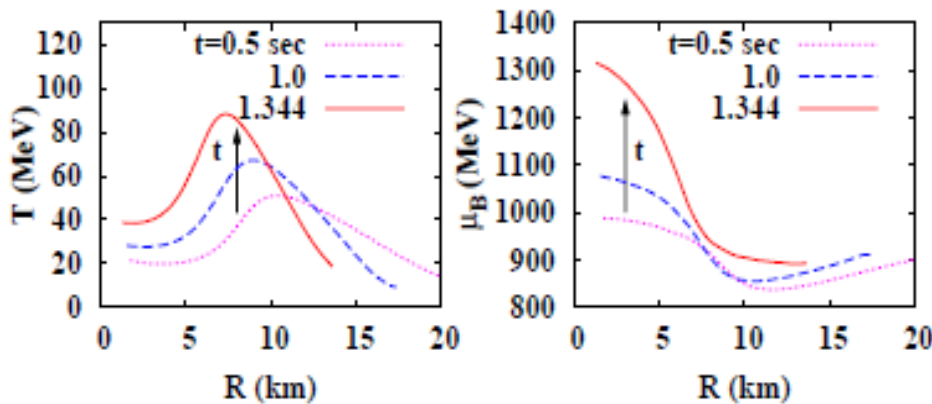
■ Collapse → Bounce → Accretion



■ → Hot, Dense, Asymmetric Matter

$$T \sim 70 \text{ MeV}, \mu_B \sim 1300 \text{ MeV}, \delta\mu = \mu_e/2 \sim 130 \text{ MeV}$$

■ → CP may be reachable



*K. Sumiyoshi, et al., ('06); K. Sumiyoshi, C. Ishizuka, A.O., S. Yamada, H. Suzuki ('09)
A.O., H. Ueda, T.Z. Nakano, M. Ruggieri, K. Sumiyoshi, PLB in press.*

RMF is a phenomenological MODEL !

- **Baryon one-loop approximation (Hartree approximation) makes RMF a phenomenological model.**
→ **We need DATA and AB INITIO results.**
- **Saturation point (ρ_0 and $E/A(\rho_0)$) from mass formula**
- **Nuclear binding energies**
- **U_v and U_s from DBHF results**
- **$P(\rho_B)$ from heavy-ion data**
- **Λ separation energy from single Λ hypernuclear data**
- **$\Lambda\Lambda$ bond energy from double Λ hypernuclear data**
- **Σ atomic shift**
- **Σ and Ξ potential depth from quasi-free production data**
- ***Pure neutron matter EOS from ab initio calculations (not used here)***

RMF models

■ Variety of the RMF models

→ MB couplings, meson masses, meson self-energies

● σN , ωN , ρN couplings are well determined

→ almost no model deps. in Sym. N.M. at low ρ

● ω^4 term is introduced to simulate DBHF results of vector pot.

TM1&2: Y. Sugahara, H. Toki, NPA579('94)557;

R. Brockmann, H. Toki, PRL68('92)3408.

● σ^3 and σ^4 terms are introduced to soften EOS at ρ_0 .

J. Boguta, A.R. Bodmer NPA292('77)413,

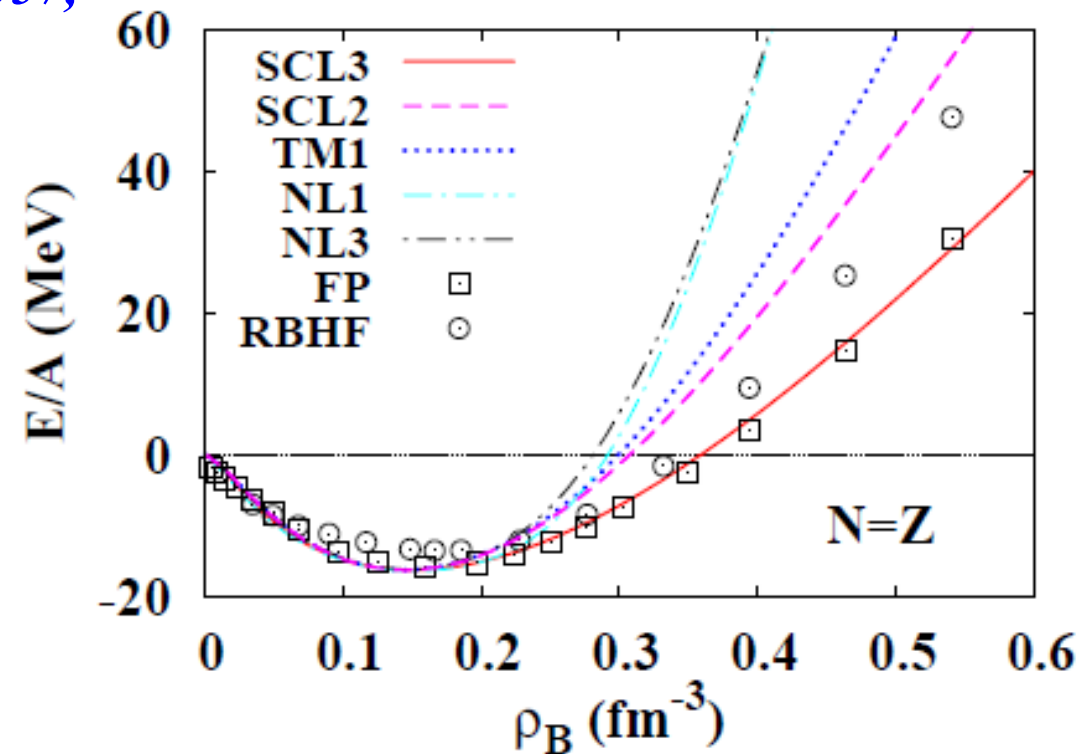
NL1: P.-G. Reinhardt, M. Rufa, J. Maruhn,

W. Greiner, J. Friedrich, ZPA323('86)13.

NL3: G.A. Lalazissis, J. Konig, P. Ring,

PRC55('97)540.

● → Large differences are found at high ρ



K. Tsubakihara, H. Maekawa, H. Matsumiya, AO, PRC81('10)065206.

Vector potential in RMF

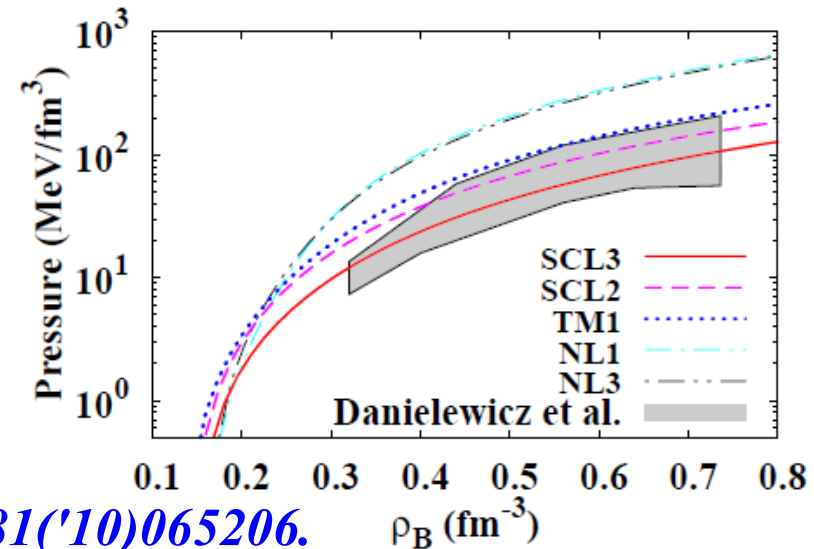
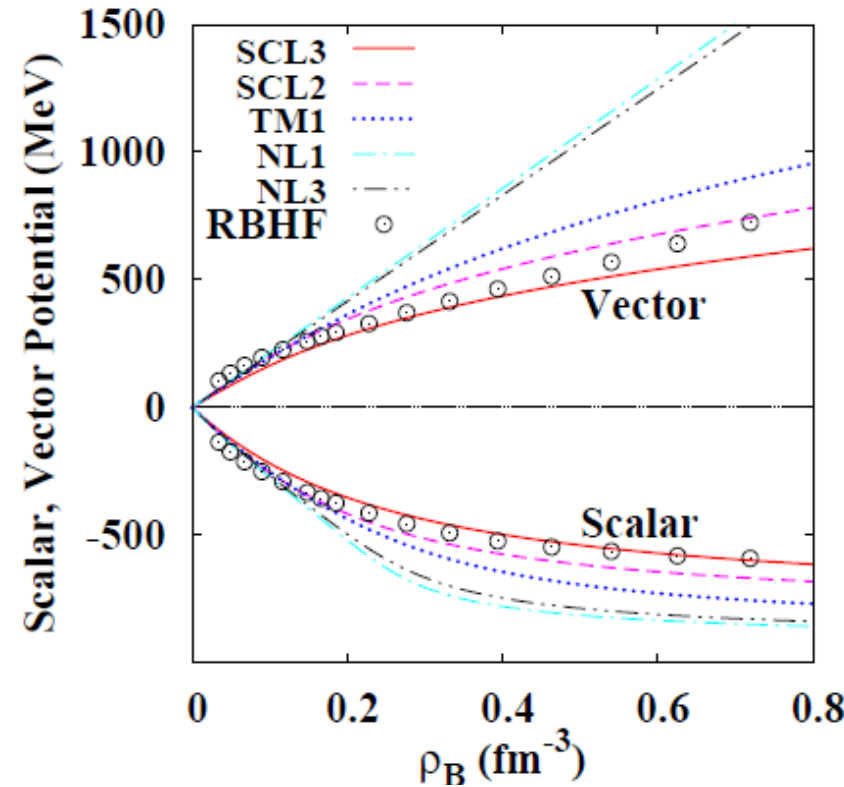
- Vector potential from ω dominates at high density !

$$U_v(\rho_B) = g_\omega \omega \sim \frac{g_\omega^2}{m_\omega^2} \rho_B$$

- Dirac-Bruckner-Hartree-Fock shows suppressed vector potential at high ρ_B .
R. Brockmann, R. Machleidt, PRC42('90)1965.

- Collective flow in heavy-ion collisions suggests pressure at high ρ_B .
P. Danielewicz, R. Lacey, W. G. Lynch, Science298('02)1592.

- Self-interaction of $\omega \sim c_\omega (\omega_\mu \omega^\mu)^2$
→ DBHF results & Heavy-ion data



K. Tsubakihara, H. Maekawa, H. Matsumiya, AO, PRC81('10)065206.

RMF with Hyperons (Single Λ hypernuclei)

RMF for Λ hypernuclei

$x \sim 1/3$: R. Brockmann, W. Weise, PLB69('77)167; J. Boguta and S. Bohrman, PLB102('81)93.

$x \sim 2/3$: N. K. Glendenning, PRC23('81)2757, PLB114('82)392;

Tensor: Y. Sugahara, H. Toki, PTP92('94)803; H. Shen, F. Yang, H. Toki, PTP115('06)325;

J. Mares, B. K. Jennings, PRC49('94)2472.

ρ -dep. coupling: H. Lenske, Lect. Notes Phys. 641('04)147; C. M. Keil, F. Hofmann, H. Lenske, PRC 61('00)064309.

SU(3) or SU(6) (ζ, φ): J. Schaffner, C. B. Dover, A. Gal, C. Greiner, H. Stoecker, PRL71('93)1328;

Schaffner et al., Ann.Phys.235('94)35; J. Schaffner, I. N. Mishustin, PRC 53('96)1416.

Chiral SU(3) RMF: K. Tsubakihara, H. Maekawa, H. Matsumiya, AO, PRC81('10)065206.

Sep. E. of Λ is well fitted

by $U_{\Lambda} \sim -30 \text{ MeV} \sim 2/3 U_N$

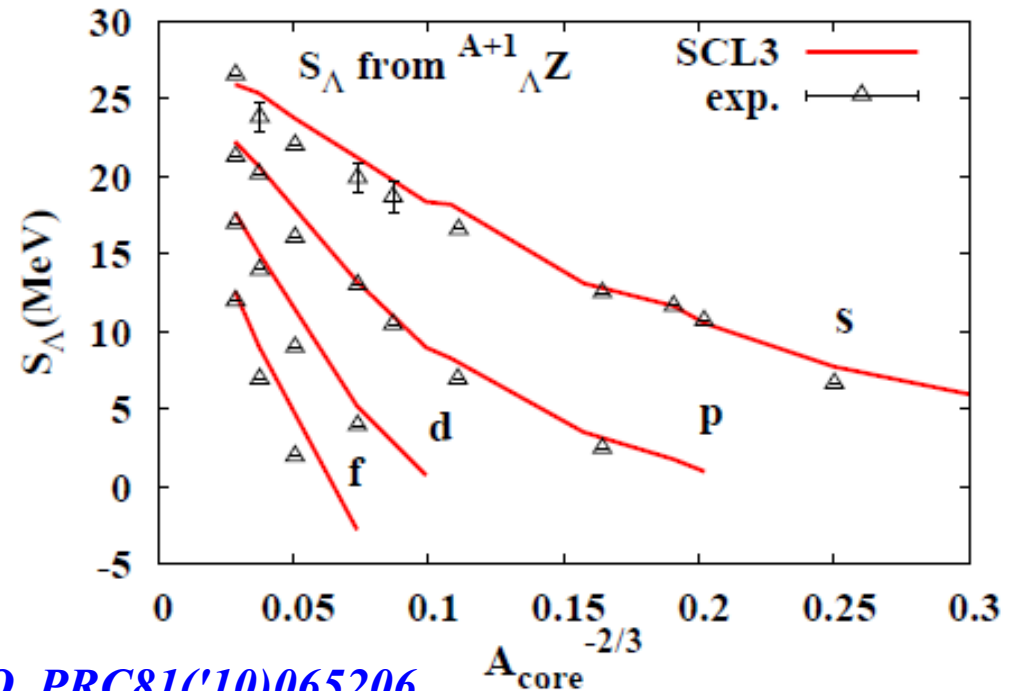
Coupling with mesons

$$x_M = g_{M\Lambda} / g_{MN}$$

quark counting: $x_{\sigma} \sim 2/3$

π exchanges: $x_{\sigma} \sim 1/3$

→ Which is true ?



K. Tsubakihara, H. Maekawa, H. Matsumiya, AO, PRC81('10)065206.

RMF with Hyperons (Double Λ hypernuclei)

■ Nagara event $\Delta B_{\Lambda\Lambda} \sim 1.0$ MeV (weakly attractive)

● TM & NL-SH based RMF

H. Shen, F. Yang, H. Toki, PTP115('06)325.

Model 1: $x_\sigma = 0.621$, $x_\omega = 2/3$ (no ζ , φ)

Model 2: $R_\zeta = g_{\zeta\Lambda} / g_{\zeta N} = 0.56-0.57$, $R_\varphi = g_{\varphi\Lambda} / g_{\varphi N} = -\sqrt{2/3}$

● Chiral SU(3) RMF

K. Tsubakihara, H. Maekawa, H. Matsumiya, AO, PRC81('10)065206.

SU(3)f for vector coupling

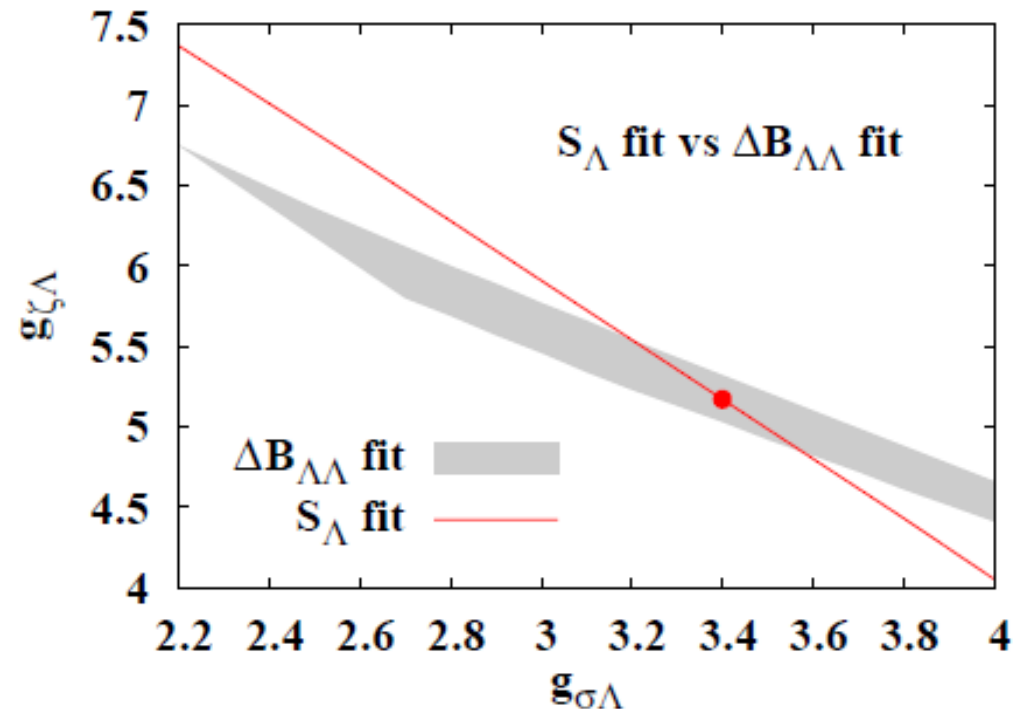
$x_\omega = 0.64$, $R_\varphi = 0.504$

Det. (KMT) int. mixes σ and ζ

M. Kobayashi, T. Maskawa, PTP44('70)1422;

G. 't Hooft, PRD14('76)3432.

→ $x_\sigma = 0.335$, $R_\zeta = 0.509$



Hyperon Composition in Dense Matter

■ Hyperon start to emerge at $(2-3)\rho_0$ in Neutron Star Matter !

■ Hyperon composition in NS is sensitive to Hyperon potential.

● $U_\Lambda \sim -30$ MeV: Well-known

● $U_\Xi \sim -(12-15)$ MeV

(K^-, K^+) reaction, twin hypernuclei

P. Khaustov et al. (E885), PRC61('00)054603;
S. Aoki et al., PLB355('95)45.

● $U_\Sigma \sim -30$ MeV (Old conjecture)

→ Σ^- appears prior to Λ

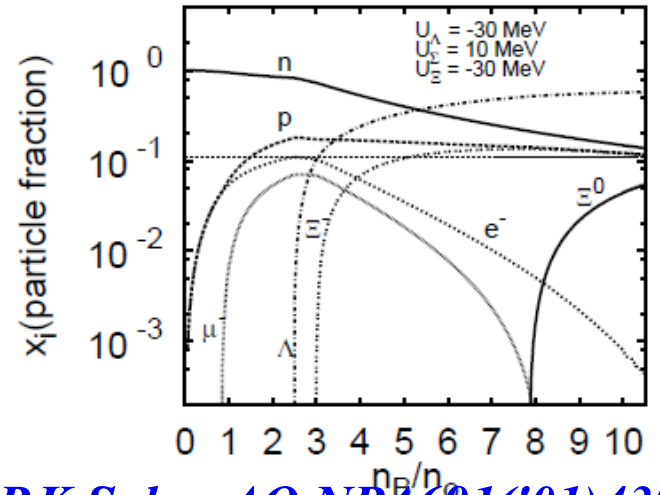
● $U_\Sigma > 0$ (repulsive) → No Σ in NS
 Σ atom (phen. fit), QF prod.

S. Balberg, A. Gal, NPA625('97)435;

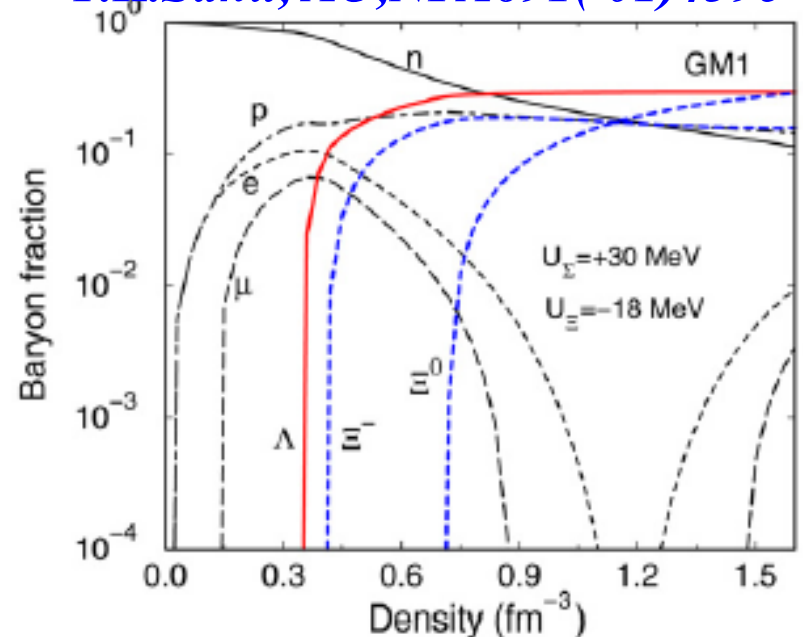
H. Noumi et al., PRL89('02)072301;

T. Harada, Y. Hirabayashi, NPA759('05)143;

M. Kohno et al. PRC74('06)064613.



P.K.Sahu, AO, NPA691('01)439c



J. Schaffner-Bielich, NPA804('08)309.

Hyperon Composition in Dense Matter

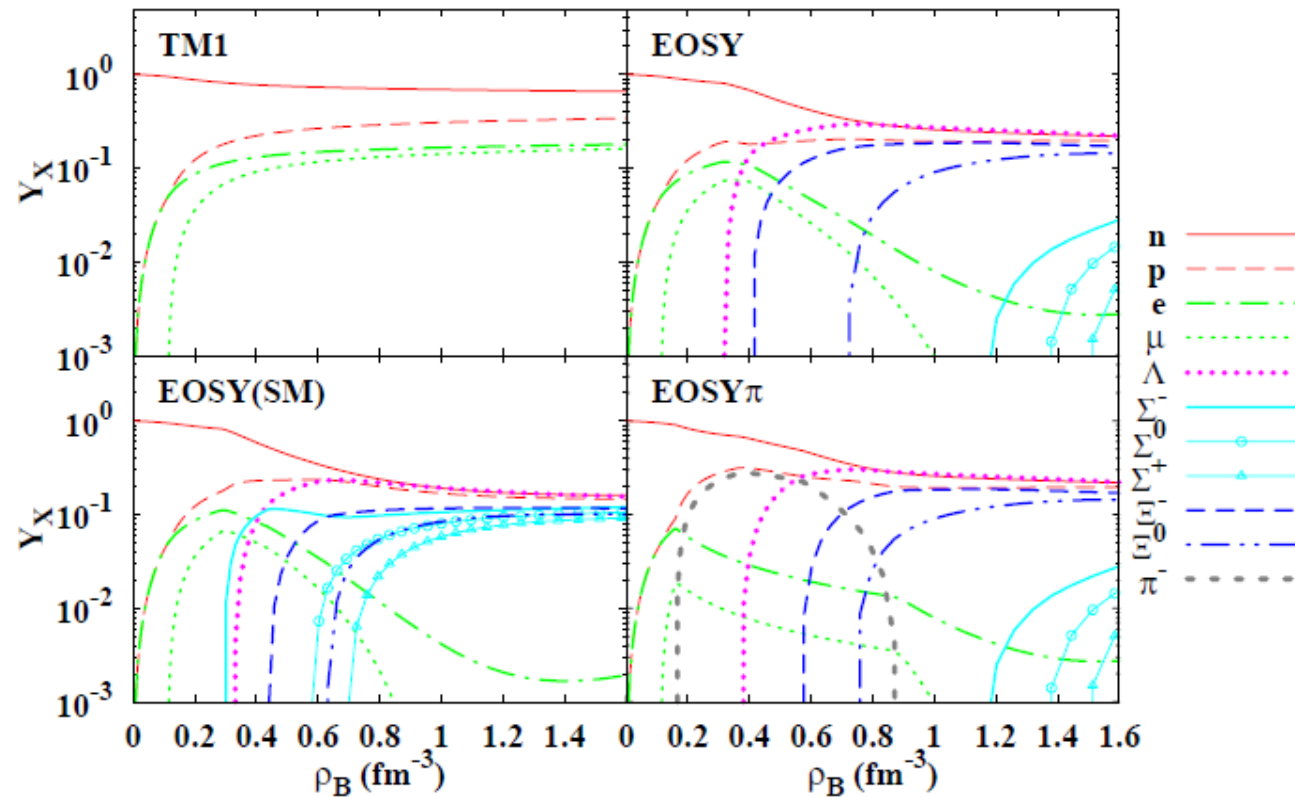
Comparison of Hyperon Composition

$U_{\Sigma} = -30 \text{ MeV}$, $U_{\Xi} = -28 \text{ MeV} \rightarrow \text{SU}(3) \text{ sym. matter at } \rho_B \sim 10 \rho_0$
 Schaffner, Mishustin ('94)

$U_{\Sigma} = +30 \text{ MeV}$, $U_{\Xi} = -15 \text{ MeV} \rightarrow \Sigma \text{ baryons are strongly suppressed.}$
C. Ishizuka, AO, K. Tsubakihara, K. Sumiyoshi, S. Yamada, JPG35('08)085201.

→ Does Σ play no role
 in NS ?

Neutron Star Matter



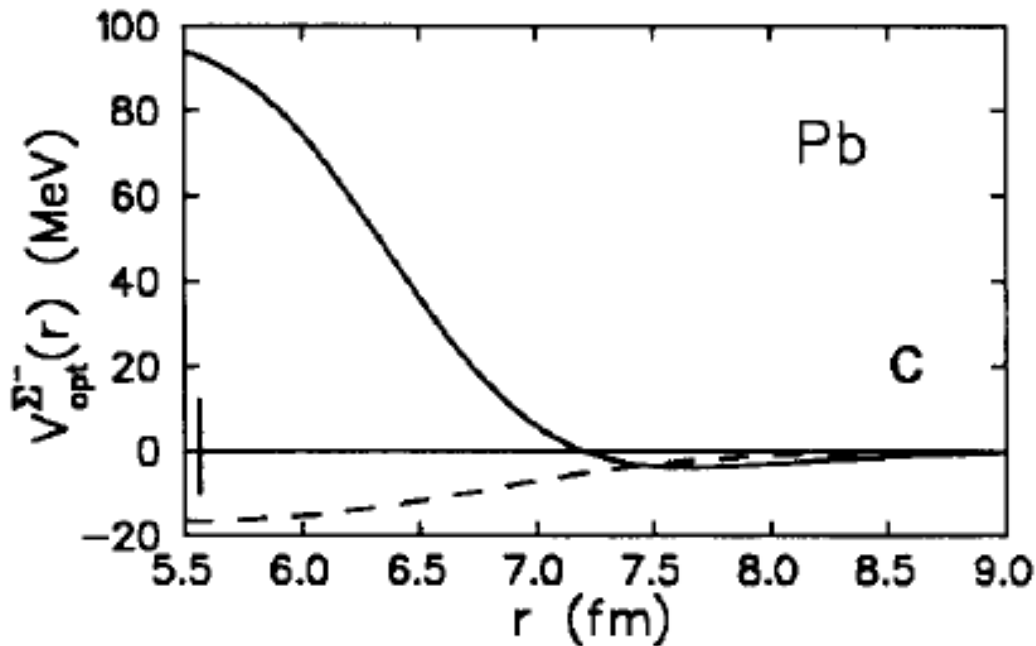
Σ^- atom data

■ Σ^- atom data suggested repulsion in the interior of nuclei !

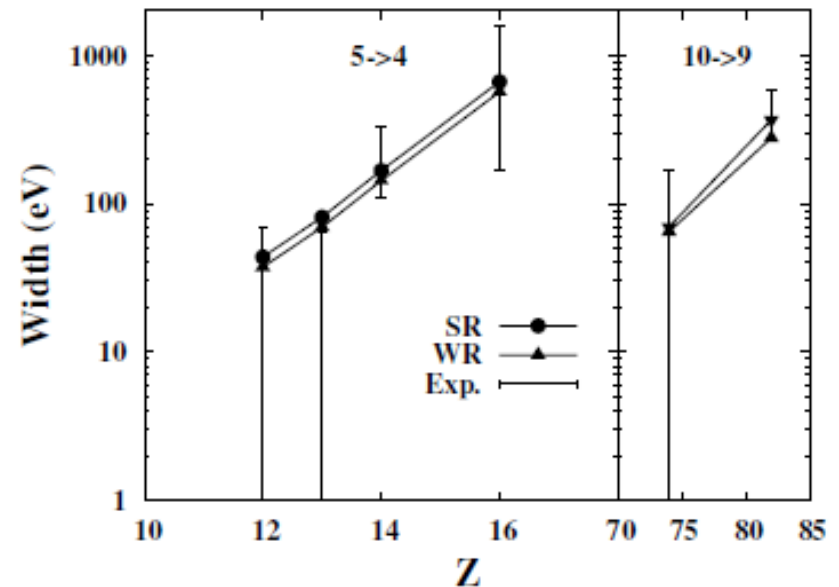
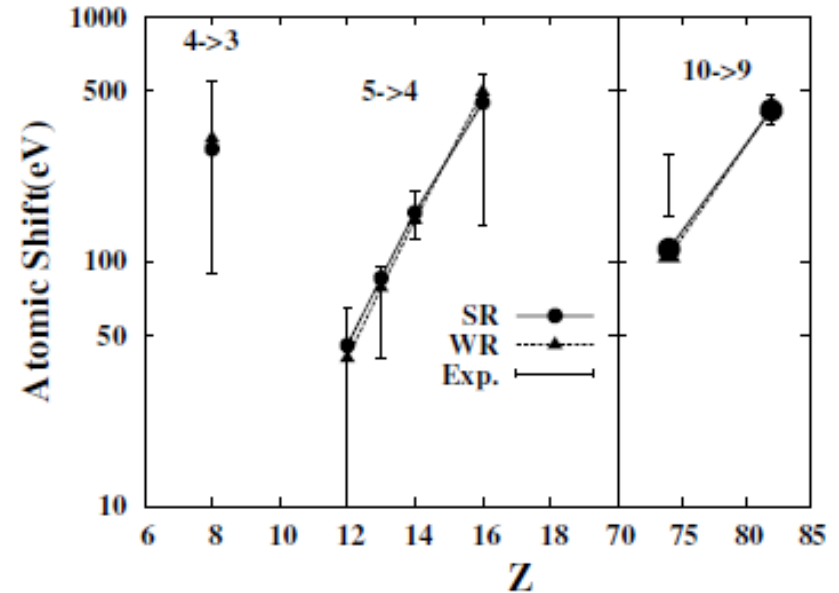
C.J.Batty, E.Friedman, A.Gal, PLB335('94)273

Batty's DD potential is very repulsive inside nuclei.

→ No Σ baryon in dense matter.



J.Mares, E.Friedman, A.Gal, B.K.Jennings, NPA594('95)311.



K.Tsubakihara, H.Maekawa, AO, EPJA33('07)295.

Σ^- atom in RMF

RMF fit of Si and Pb Σ^- atom

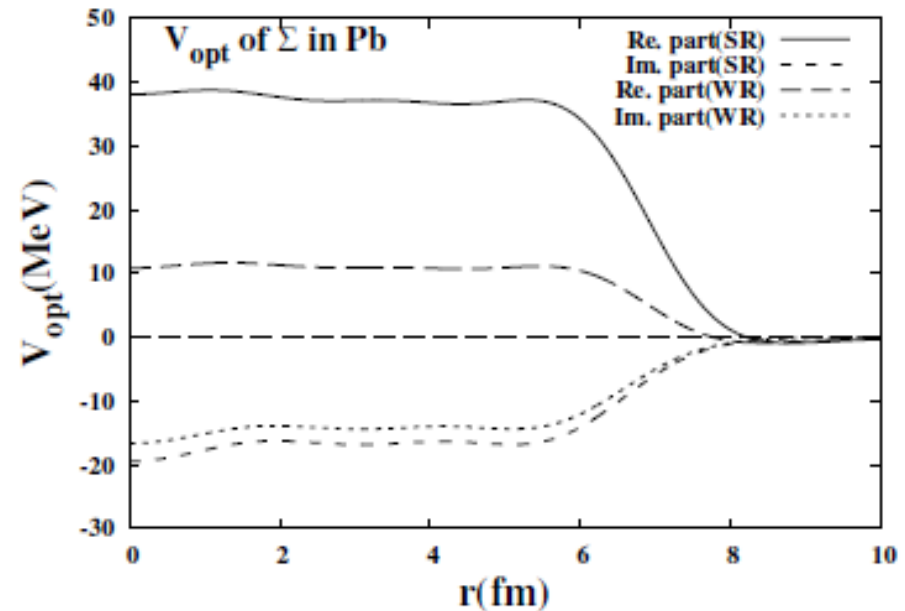
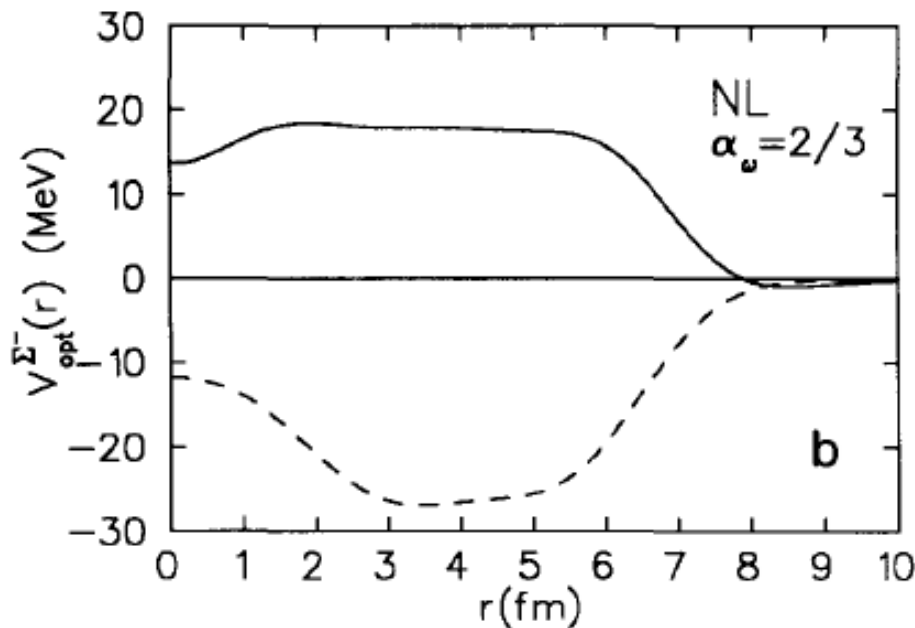
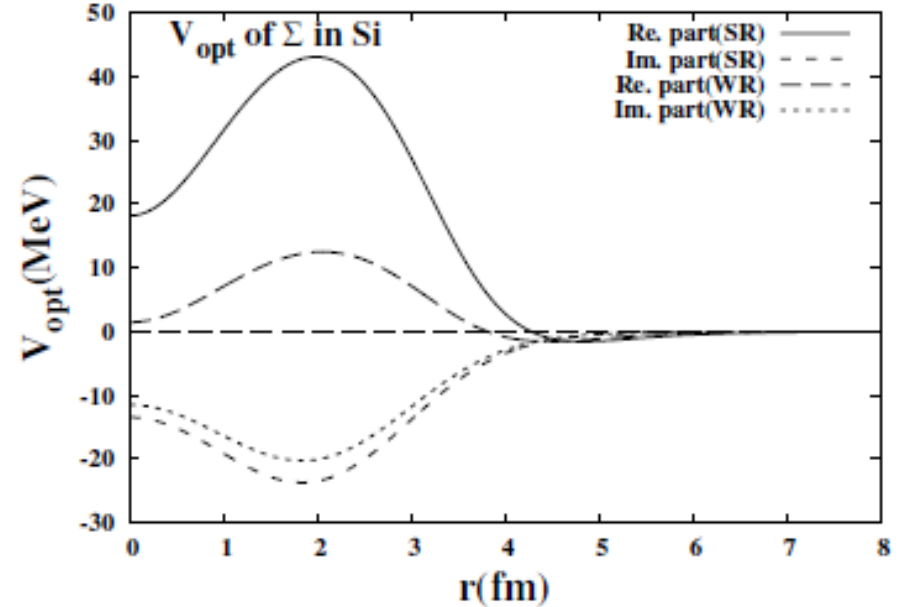
$$\alpha_\omega = g_{\omega\Sigma} / g_{\omega N} \sim 2/3(\text{M}), 0.69(\text{T})$$

$$\alpha_\rho = g_{\rho\Sigma} / g_{\rho N} \sim 2/3(\text{M}), 0.434(\text{T})$$

J.Mares, E.Friedman, A.Gal, B.K.Jennings,

NPA594('95)311; Tsubakihara et al.('10)

● Much smaller $g_{\rho\Sigma}$ than naïve SU(3) ($g_{\rho\Sigma} / g_{\rho N} = 2$), which has been applied in some of previous works.



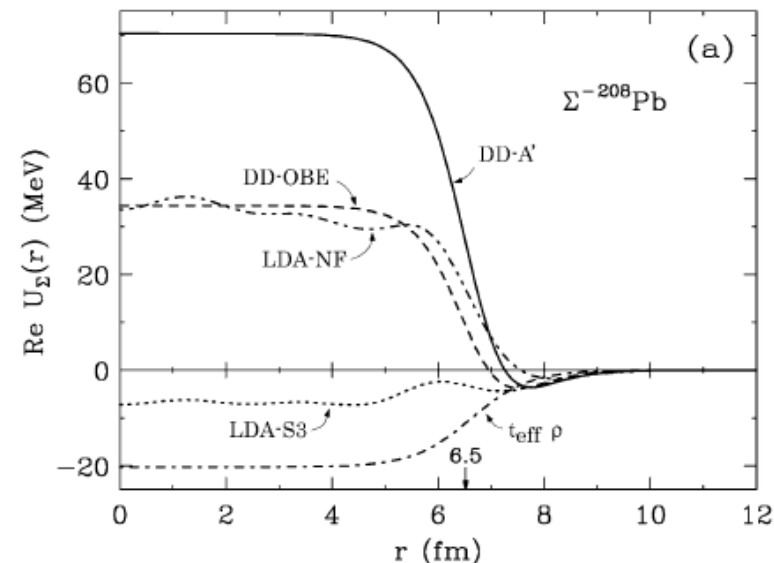
Σ atom and Neutron Star

■ Σ may not feel *very* repulsive potential in neutron star....

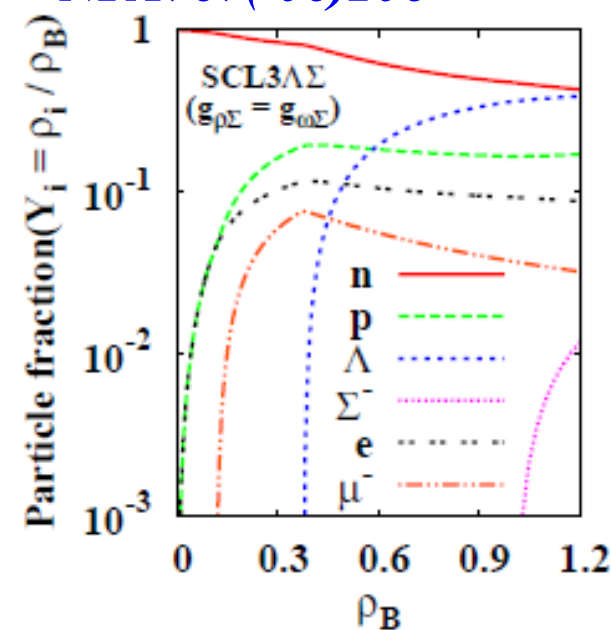
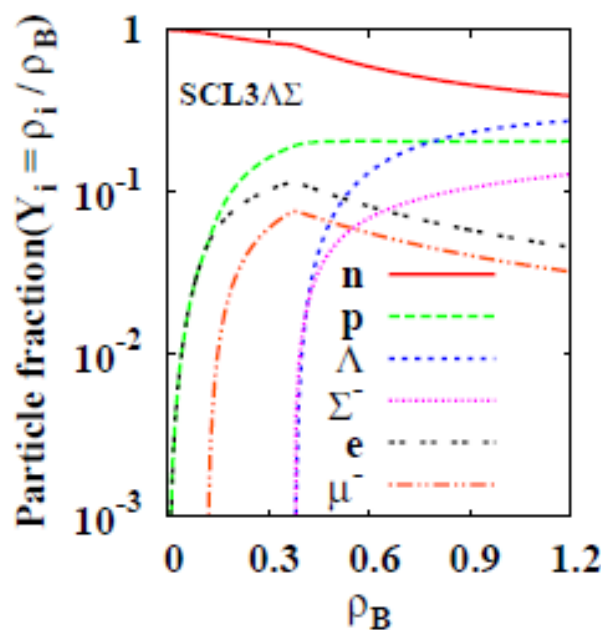
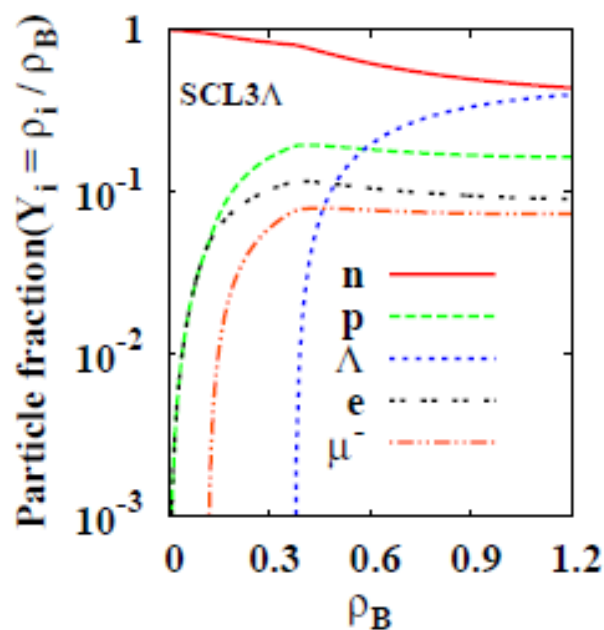
● ρ^γ -type fit \rightarrow very repulsive

● RMF fit \rightarrow small isovector potential

■ \rightarrow QF prod. may support the latter.
 Σ^- would appear in NS.



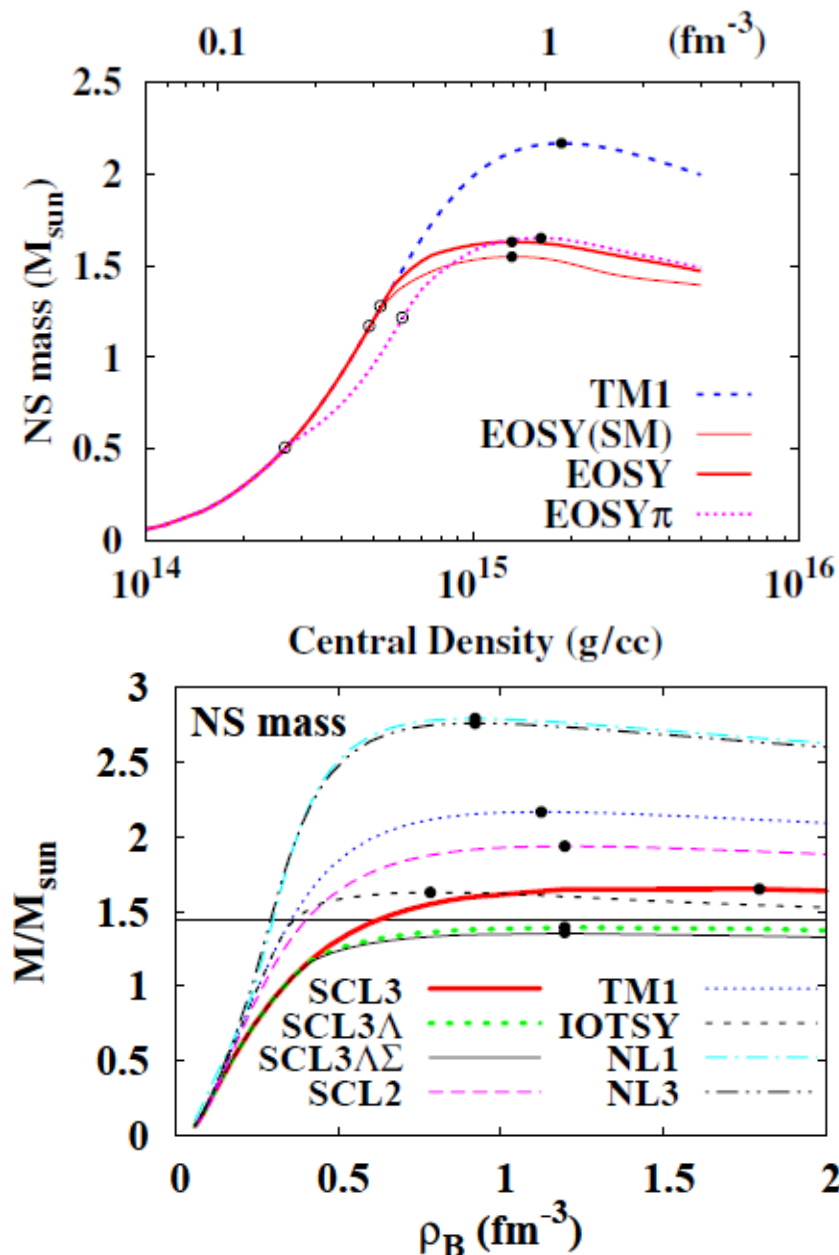
*T. Harada, Y. Hirabayashi,
 NPA767('06)206*



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Neutron Star Mass

- Large fraction of hyperons softens EOS at $\rho_B > (0.3-0.4) \text{ fm}^{-3}$
- NS star max. mass red. $\sim 1 M_{\text{sun}}$.
- RMF generally predicts stiff EOS at high density.
(Scalar attraction saturation, or Z-graph in NR view.)
- Some of RMF with Y do not support $1.44 M_{\text{sun}}$.
- Additional Repulsion at high ρ ?
- Vector mass mod.
→ stronger repulsion at high ρ .
- Another term such as $\text{NN}\omega\sigma$.



C. Ishizuka, AO, K. Tsubakihara, K. Sumiyoshi, S. Yamada, JPG35('08)085201.

K. Tsubakihara, H. Maekawa, H. Matsumiya, AO, PRC81('10)065206.