

Physics of Neutron Star Matter

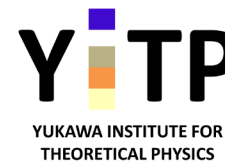
– current status and challenges –

Akira Ohnishi (YITP, Kyoto U.)

Grant-in-Aid for Scientific Research on Innovative Areas

**Nuclear matter in neutron stars
investigated by experiments
and astronomical observations**

JFY 2012-2016



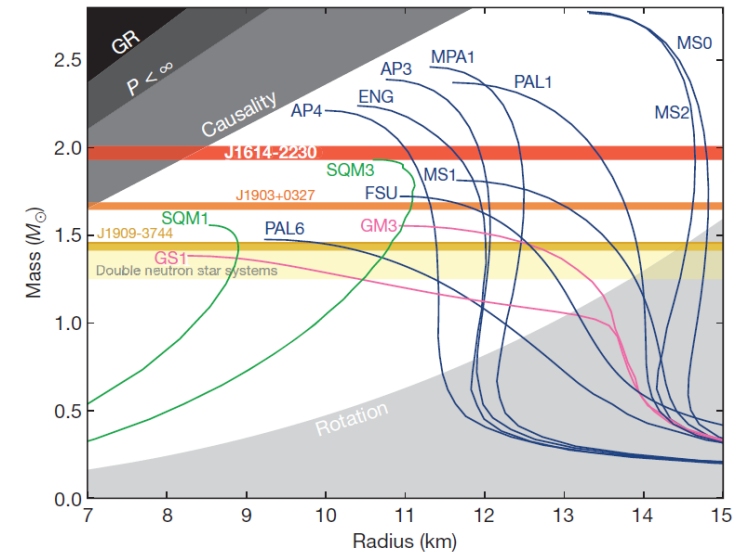
Contents and Summary

■ Neutron star (NS)

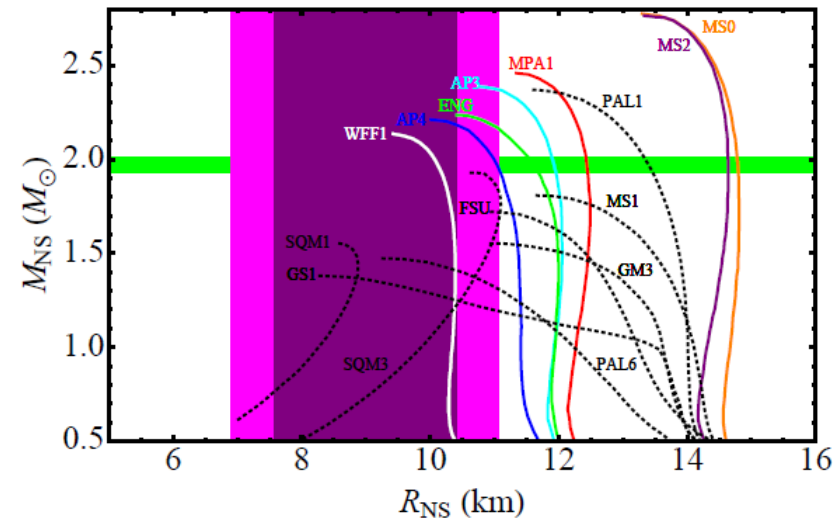
- Giant neutron-rich hypernuclei & Almost causal limit object
- Massive ($\sim 2 M_{\odot}$) & Compact (~ 9 km) NS puzzle.
- Grant-in-aid study on NS matter
J-PARC + RIBF + ASTRO-H + Theory

■ Challenges in neutron star matter physics

- How can we solve massive NS puzzle ?
→ Three-body force as a candidate
- What determines M-R relation of low-mass neutron stars ?
→ New parameter $\eta=(KL^2)^{1/3}$



Demorest et al. (2010)



Guillot et al. (2013)

*Physics of Neutron Stars
– Basics and Puzzles –*

Basic properties of neutron stars

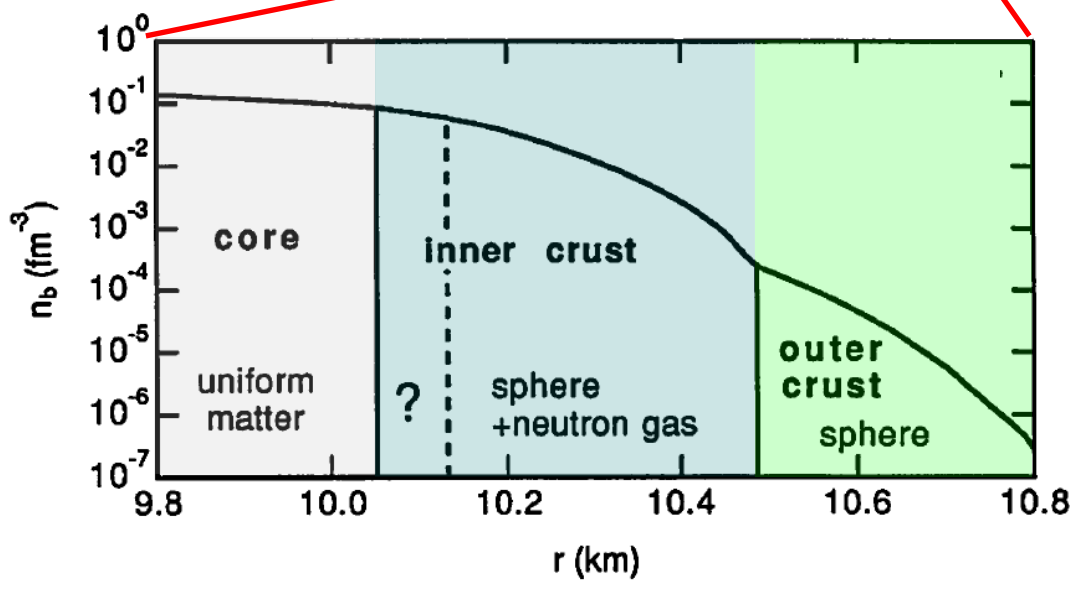
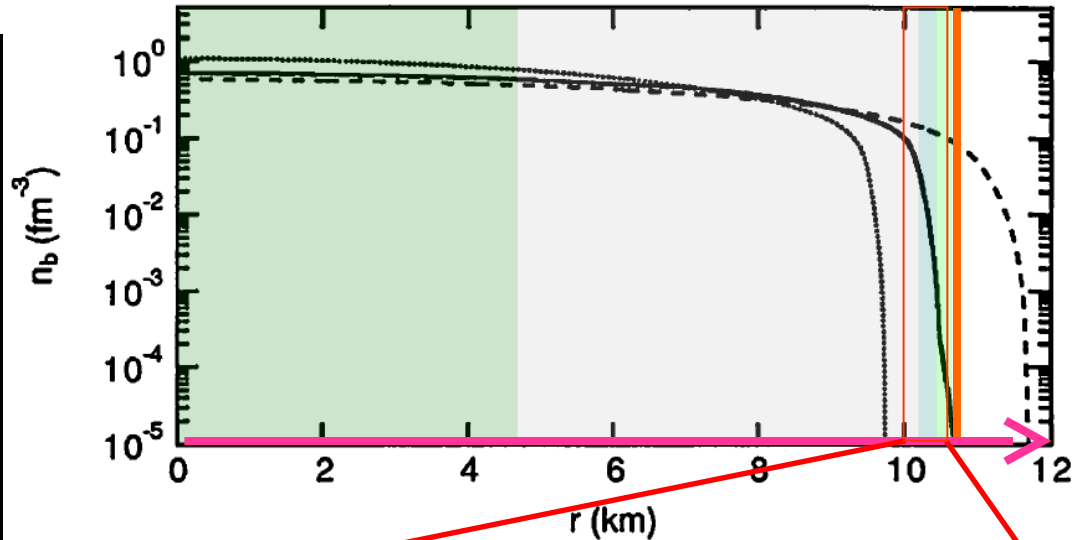
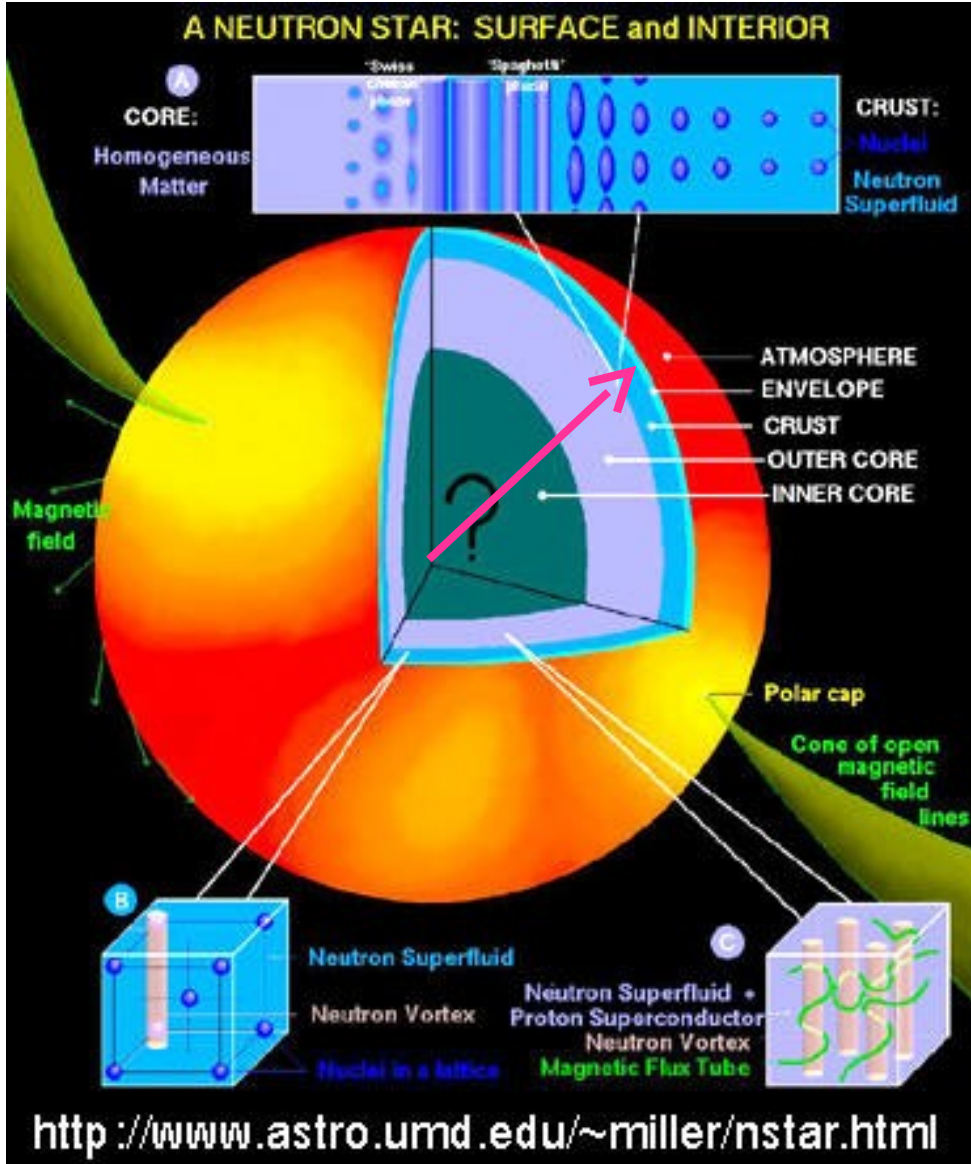
- **Mass:** $M = (1-2) M_{\odot}$ ($M \sim 1.4 M_{\odot}$)
- **Radius:** $5 \text{ km} < R < 20 \text{ km}$ ($R \sim 10 \text{ km}$)
- **Supported by Nuclear Pressure**
c.f. Electron pressure for white dwarfs
- **Cold enough** ($T \sim 10^6 \text{ K} \sim 100 \text{ eV}$) compared to neutron Fermi energy.



google & zenrin

Neutron Star Structure

Dense core + Thin Crust



by Nakazato

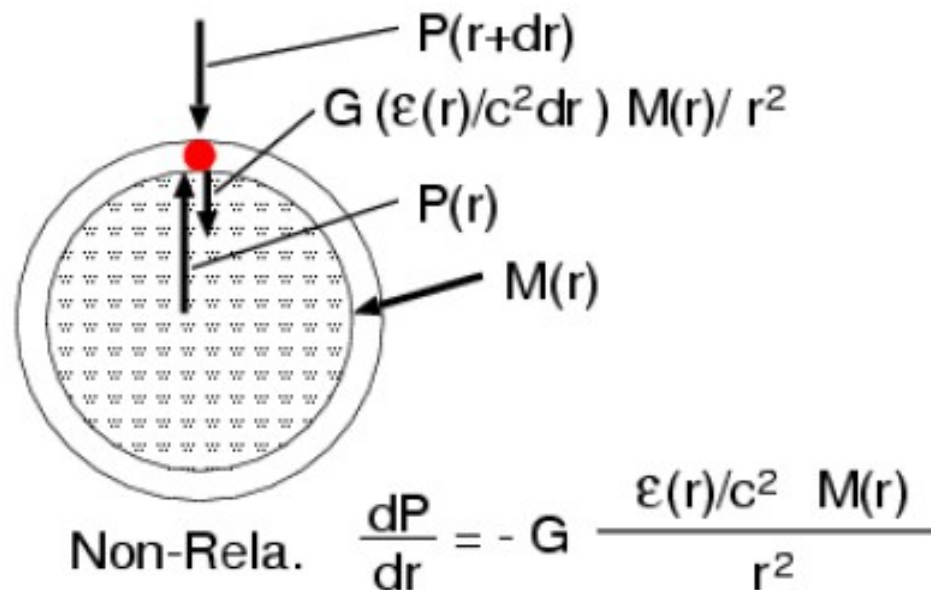
TOV equation

- General Relativistic Hydrostatic Equation
= TOV(Tolman-Oppenheimer-Volkoff) equation

$$\frac{dP}{dr} = -G \frac{(\epsilon/c^2 + P/c^2)(M + 4\pi r^3 P/c^2)}{r^2(1 - 2GM/rc^2)}$$

$$\frac{dM}{dr} = 4\pi r^2 \epsilon/c^2, \quad P = P(\epsilon) \quad (\text{EOS})$$

- Spherical and non-rotating.
- 3 Variables ($\epsilon(r)$, $P(r)$, $M(r)$),
3 Equations.



M-R Relation and EOS

■ Solving TOV eq.

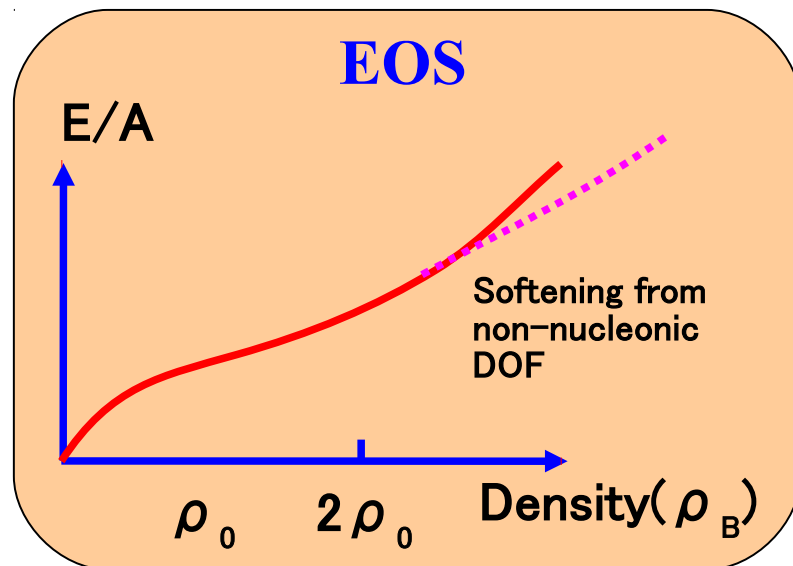
starting from the “initial” condition, $\varepsilon(r=0) = \varepsilon_c = \text{given}$

until the “boundary” condition $P(r)=0$ is satisfied.

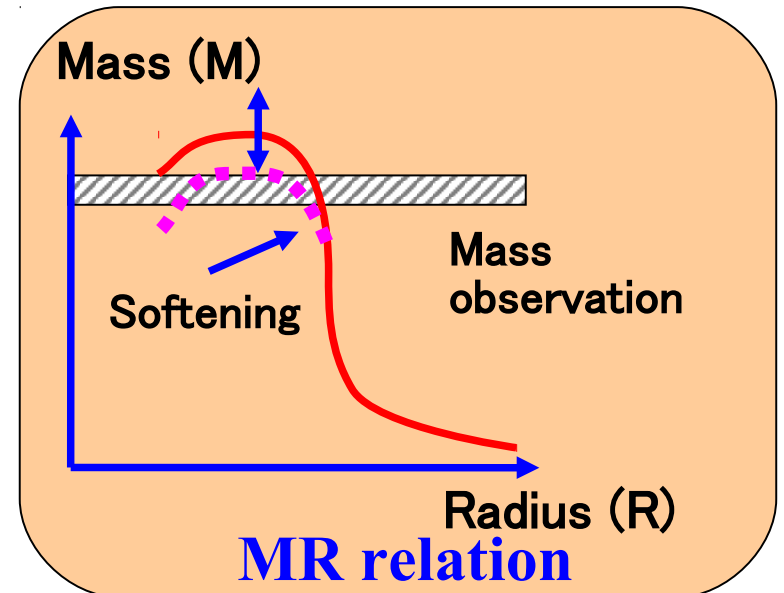
→ M and R are the functions of $\varepsilon(r=0)$ and functionals of EOS, $P=P(\varepsilon)$.

$$M = M(\varepsilon_c)[P(\varepsilon)] \quad , \quad R = R(\varepsilon_c)[P(\varepsilon)]$$

→ M-R curve and NS matter EOS : 1 to 1 correspondence



TOV Eq.



Nuclear Mass

■ Bethe-Weizsacker mass formula

Nuclear binding energy is roughly given by Liquid drop.

Nuclear size measurement $\rightarrow R = r_0 A^{1/3}$

$$B(A, Z) = \underbrace{a_v A}_{\text{Volume}} - \underbrace{a_s A^{2/3}}_{\text{Surface}} - \underbrace{a_C \frac{Z^2}{A^{1/3}}}_{\text{Coulomb}} - \underbrace{a_a \frac{(N - Z)^2}{A}}_{\text{Symmetry}} + \underbrace{a_p \frac{\delta_p}{A^\gamma}}_{\text{Paring}}$$

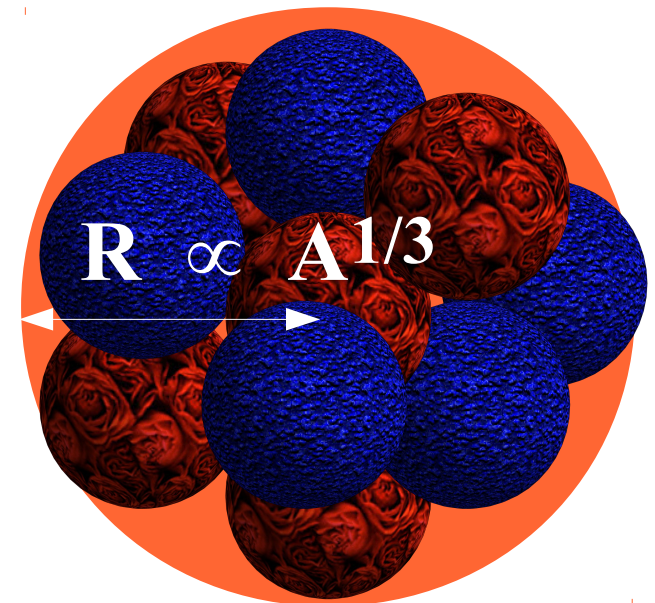
Volume Surface Coulomb Symmetry Paring

$$A \propto \frac{4\pi}{3} R^3 \quad A^{2/3} \propto 4\pi R^2 \quad \propto \frac{Q^2}{R}$$

■ Ignore Coulomb, consider $A \rightarrow \infty$,

$$B/A = a_v(\rho) - a_a(\rho) \delta^2, \quad \delta = (N - Z)/A$$

$$a_v \simeq 16 \text{ MeV}, \quad a_a \simeq 30 \text{ MeV}$$



Coef. may depend on the number density ρ

\rightarrow Nuclear Matter EOS

Nuclear Matter EOS

■ Energy per nucleon in nuclear matter

$$E(\rho, \delta) = E_{\text{SNM}}(\rho) + E_{\text{Sym}}(\rho) \delta^2, \quad \delta = (N - Z) / A$$

■ Saturation point (ρ_0, E_0)

$$\rho_0 \sim 0.15 \text{ fm}^{-3}$$

$$E_0 = -a_v \sim -16 \text{ MeV}$$

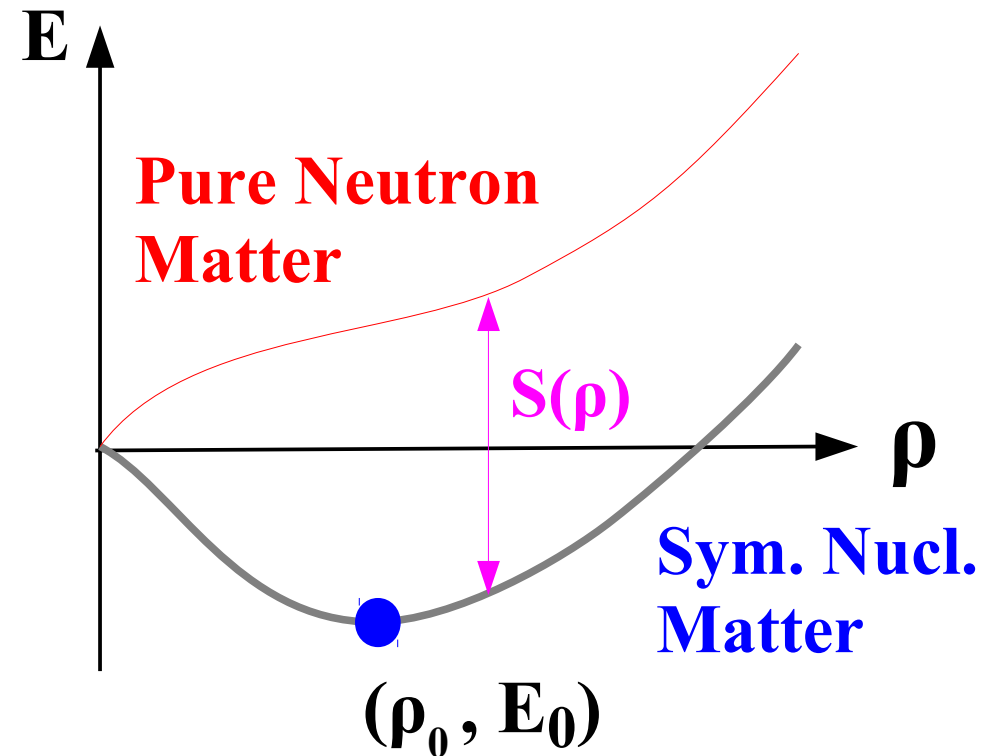
(nuclear radius and mass)

■ Symmetry energy

$$\begin{aligned} S(\rho) &= E_{\text{PNM}}(\rho) - E_{\text{SNM}}(\rho) \\ &= E(\rho, \delta=1) - E(\rho, \delta=0) \end{aligned}$$

$$S_0 = S(\rho_0) \sim 30 \text{ MeV}$$

(mass formula)



Nuclear Matter EOS can be, in principle, determined by terrestrial (laboratory) nuclear physics experiments !

Nuclear Matter EOS

- Additional two important parameters: **K** and **L**
- Pressure is given by the derivative of **E** via ρ

$$P = \rho^2 (\partial E / \partial \rho)$$

At ρ_0 , **L** determines **P**

$$P = \rho_0 L / 3 \quad (\text{at } \rho = \rho_0)$$

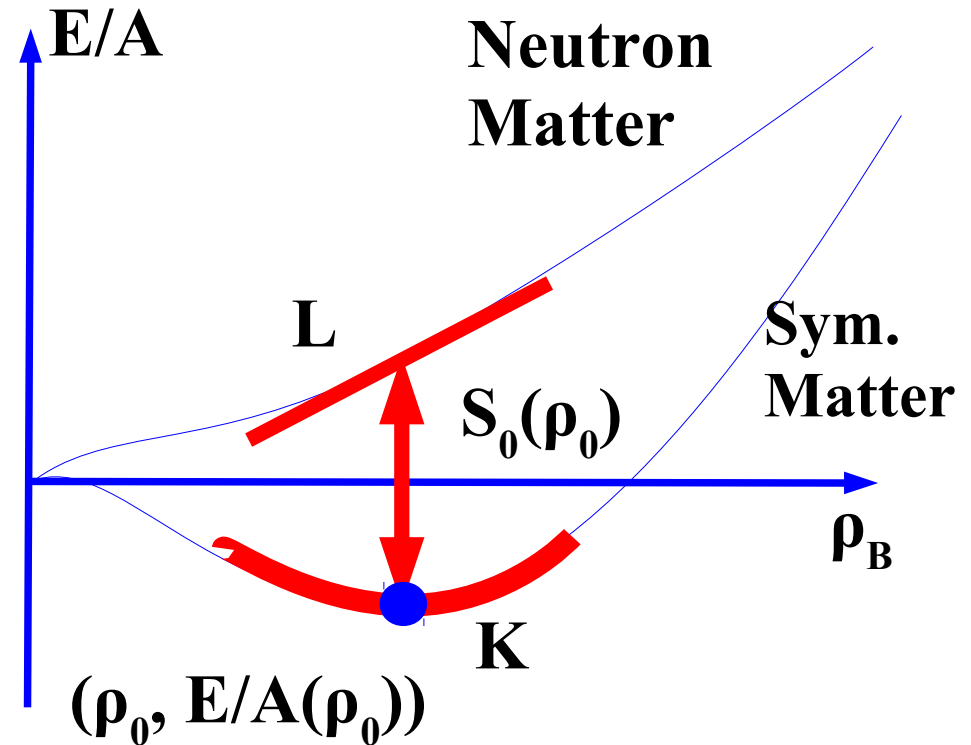
$$E/A(\rho, \delta) = \varepsilon(\rho) + E_{\text{sym}}(\rho) \delta^2 + O(\delta^4)$$

Symmetric Nuclear Matter

$$\varepsilon(\rho) = \varepsilon(\rho_0) + \frac{K(\rho - \rho_0)^2}{18\rho_0^2} + O((\rho - \rho_0)^3)$$

Symmetry Energy ($\delta = (N - Z)/A = 1 - 2Y_p$)

$$E_{\text{sym}}(\rho) = S_0 + \frac{L(\rho - \rho_0)}{3\rho_0} + \frac{K_{\text{sym}}(\rho - \rho_0)^2}{18\rho_0^2} + O((\rho - \rho_0)^3)$$



Neutron Star Matter EOS

■ What happens in low-density uniform neutron star matter ?

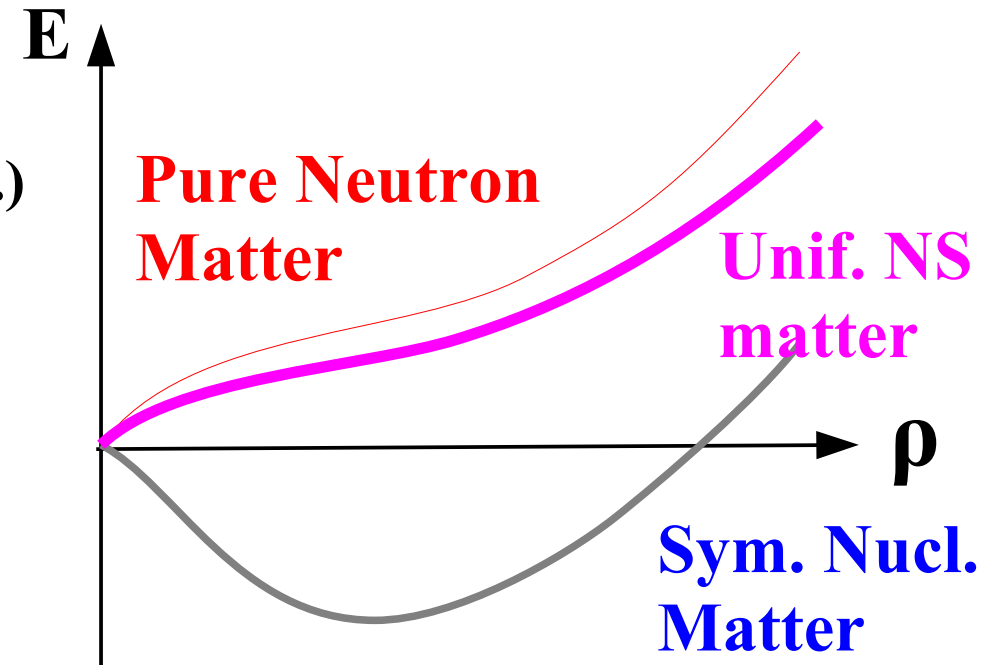
- Constituents = proton, neutron and electron
- Charge neutrality \rightarrow # of electrons = # of protons ($\rho_e = \rho_p = \rho (1 - \delta)/2$)

$$\begin{aligned}
 E_{\text{NSM}}(\rho) &= E_{\text{NM}}(\rho, \delta) + E_e(\rho_e = \rho_p) \\
 &= E_{\text{SNM}}(\rho) + S(\rho)\delta^2 + \frac{\Delta M}{2}\delta + \frac{3}{8}\hbar k_F(1 - \delta)^{4/3}
 \end{aligned}$$

(electron mass neglected,
neutron-proton mass diff. incl.
 k_F = Fermi wave num. in Sym. N.M.)

- δ is optimized to minimize energy per nucleon

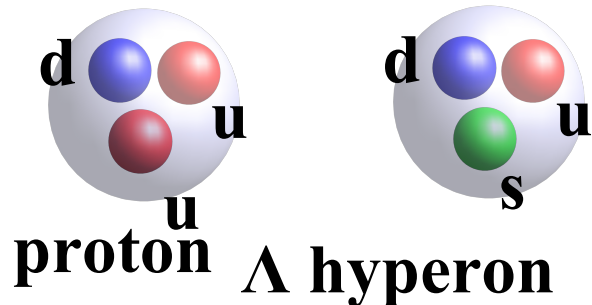
$$E_{\text{NSM}}(\rho) \leq E_{\text{NM}}(\rho, \delta = 1) = E_{\text{PNM}}(\rho)$$



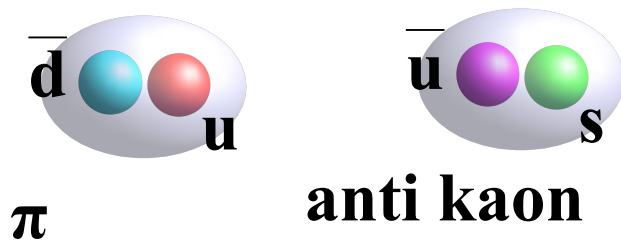
Neutron star – Is it made of neutrons ?

■ Possibilities of various constituents in neutron star core

● Strange Hadrons

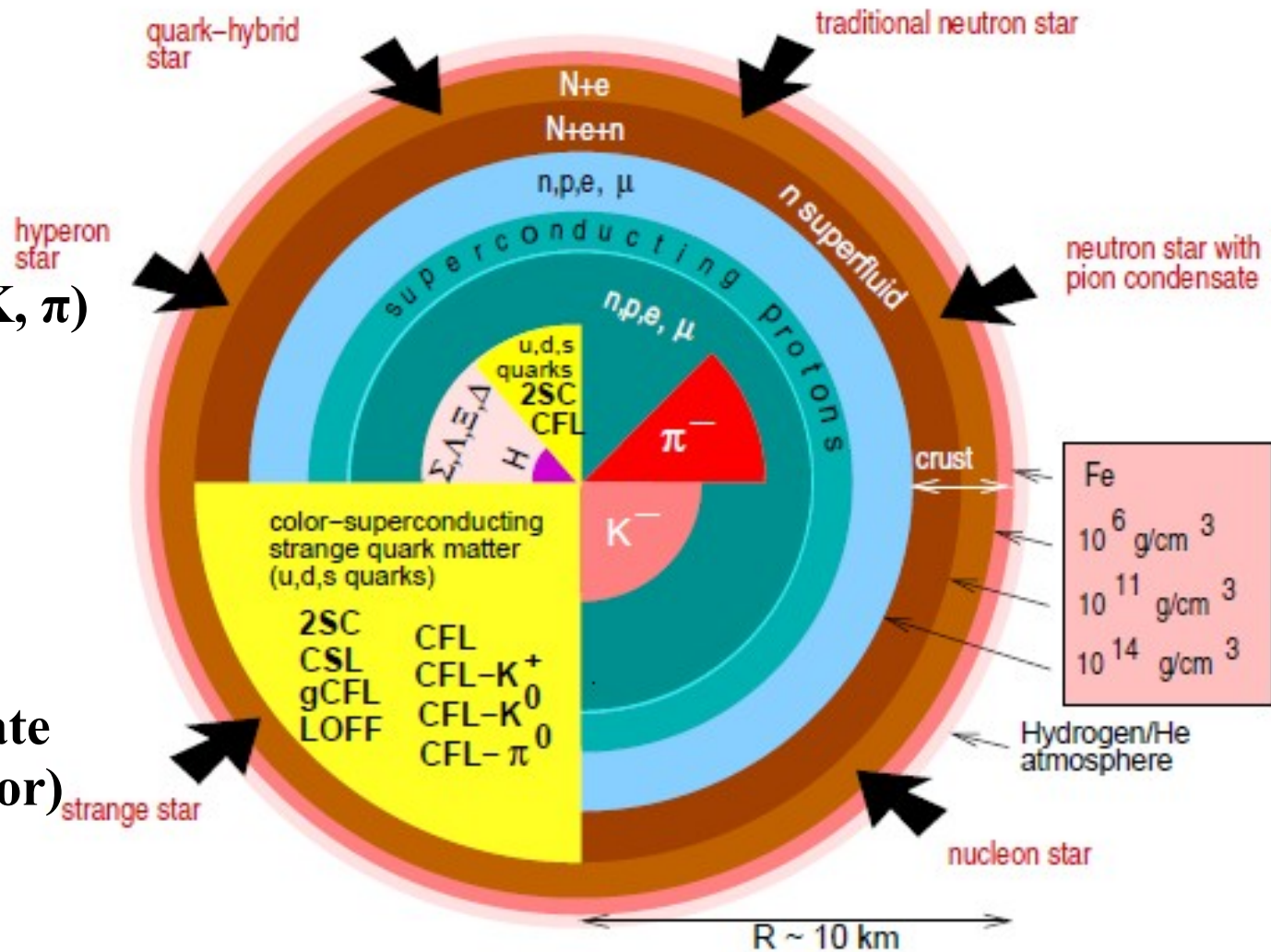


● Meson condensate (K , π)



● Quark matter

● Quark pair condensate (Color superconductor)



F. Weber, Prog. Part. Nucl. Phys. 54 (2005) 193

High Density Neutron Star Matter

- **Hadrons other than nucleons can admix at high densities.**

- **Proposed constituents = N, e, μ , π , Y, \bar{K} , q, qq**
- **Conserved charge = Elec. Charge & Baryon number**

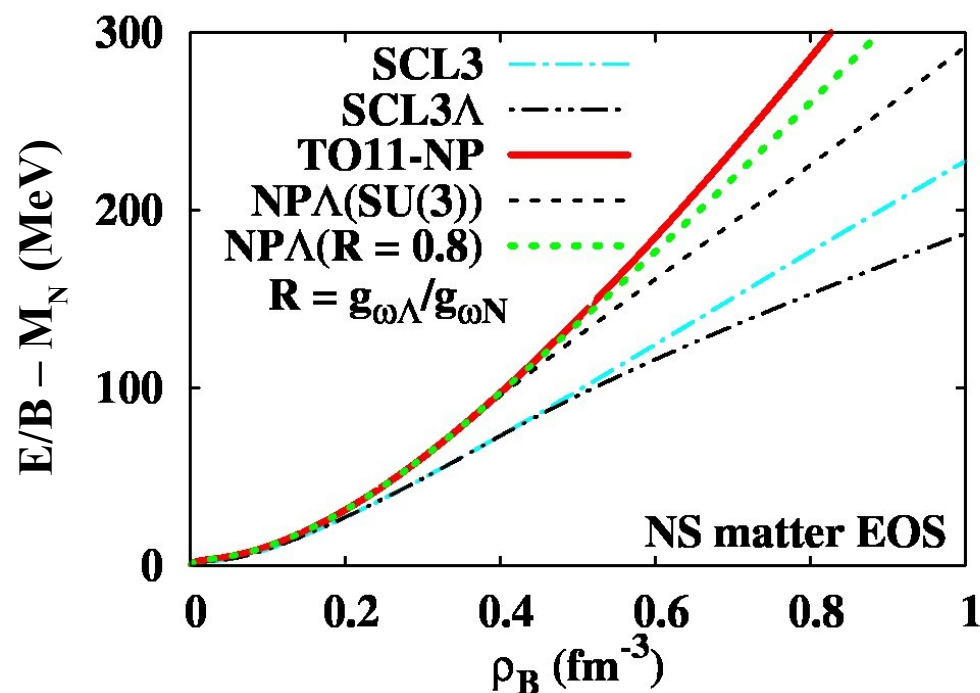
$$E_{\text{NSM}}(\rho) = E_{\text{HM}}(\rho_n, \rho_p, \rho_\Lambda, \dots) + E_e(\rho_e) + E_\mu(\rho_\mu)$$

$$\rho = \rho_B = \sum_{i \in B} \rho_i, \quad \rho_Q = \sum_i Q_i \rho_i$$

$$\rightarrow \mu_i = B_i \mu_B - Q_i \mu_e$$

Negative charged baryons are favored.

- **Each particle fraction is determined to minimize Energy per Baryon**
→ Softening of NS matter EOS



Common Understanding of Neutron Star Matter (<2010)

■ Neutron star crust ($\rho < \rho_0/2 \sim 10^{14}$ g/cc, 1-1.5 km)

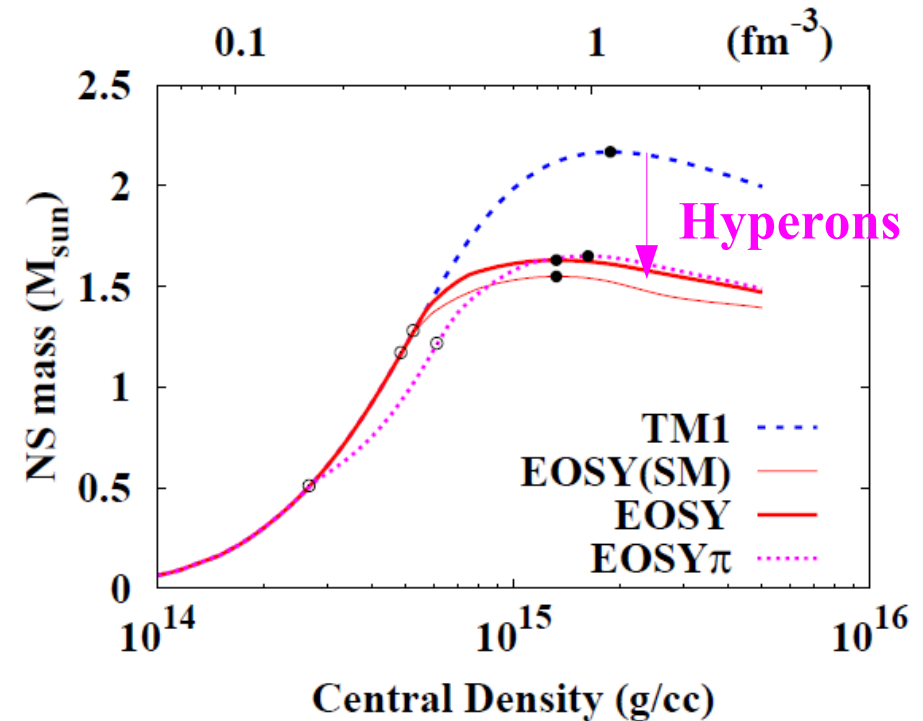
- Made of nuclei, free nucleons and electrons.
We may have pasta nuclei in the inner crust.

■ Outer core ($\rho_0/2 < \rho < 2\rho_0$)

- Made of nucleons and leptons. Nucleons are in superfluid.

■ Inner core ($\rho > 2\rho_0$)

- Constituents are unknown.
- Many calculations suggest hyperon admixture starting at $\rho = (2-4) \rho_0$, and EOS is softened.
- \bar{K} or quarks may appear at $\rho > 3 \rho_0$.

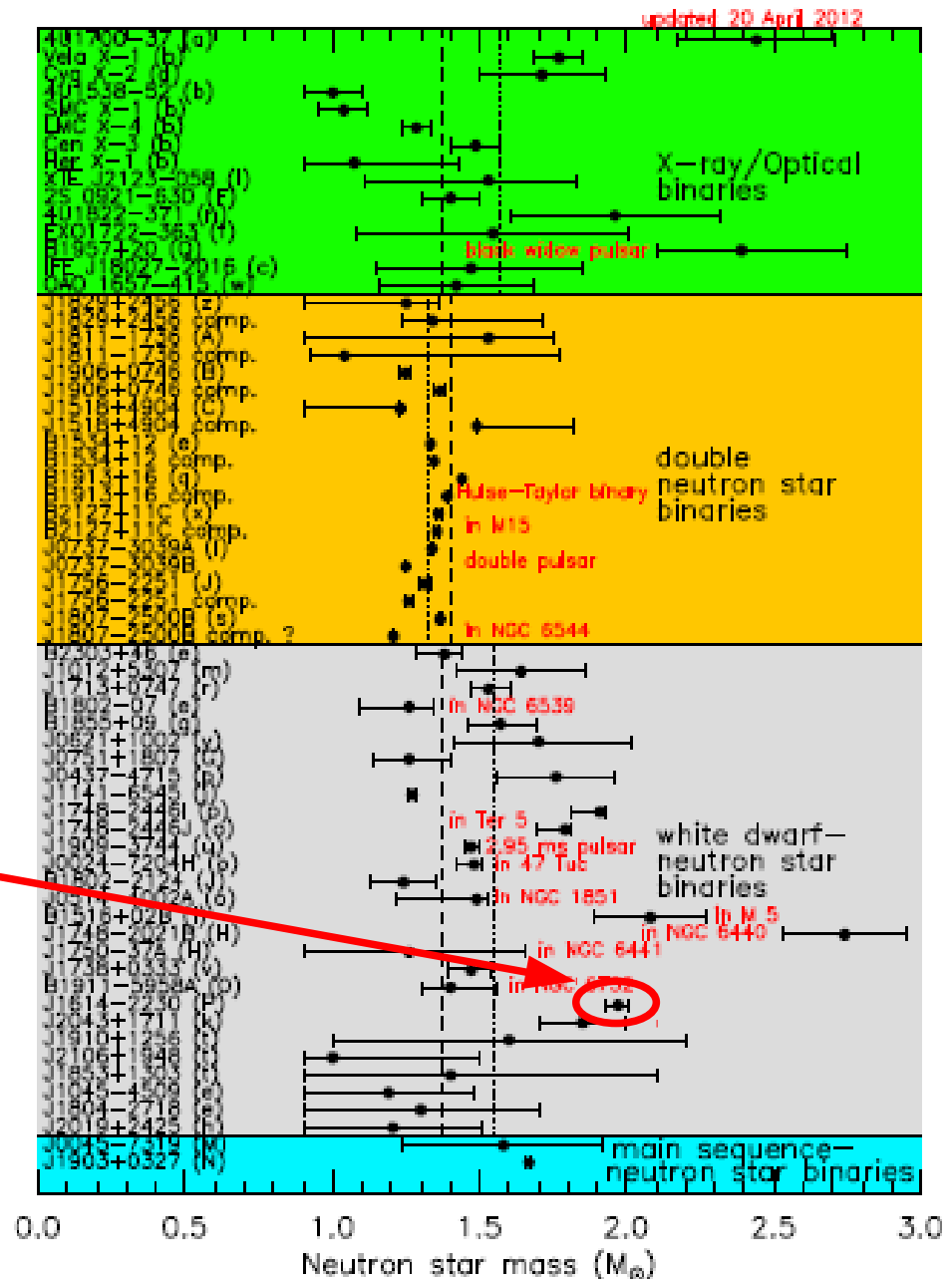


Ishizuka et al. (2008)

Neutron Star Masses

- Many NSs have masses $\sim 1.4 M_{\odot}$.
- Massive NS masses had large error bars, and were considered not to be conclusive (before 2010).

This is it !



Lattimer (2013)

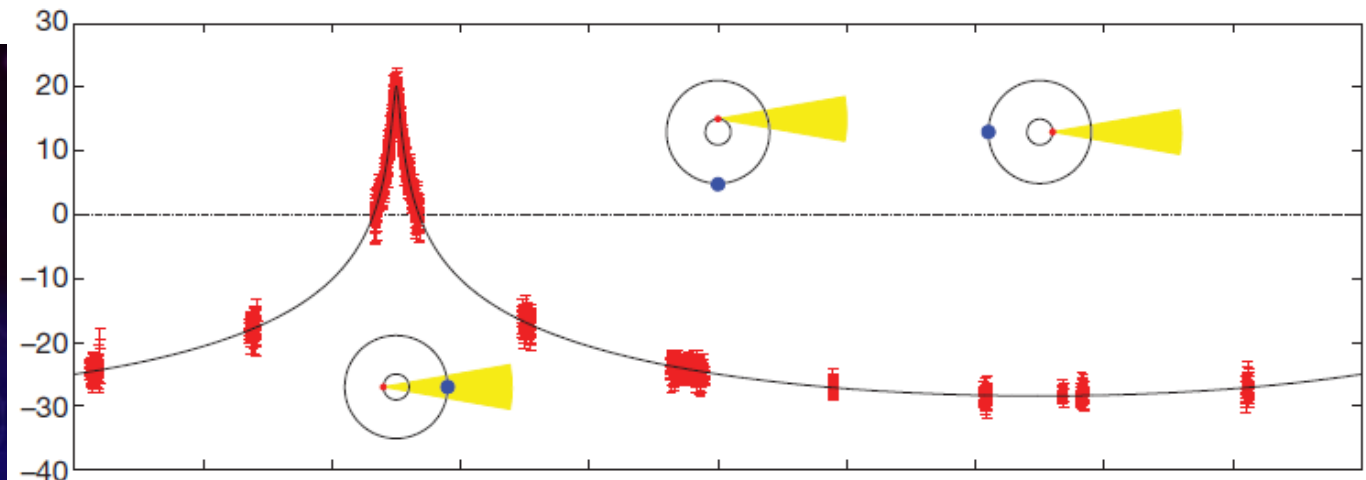
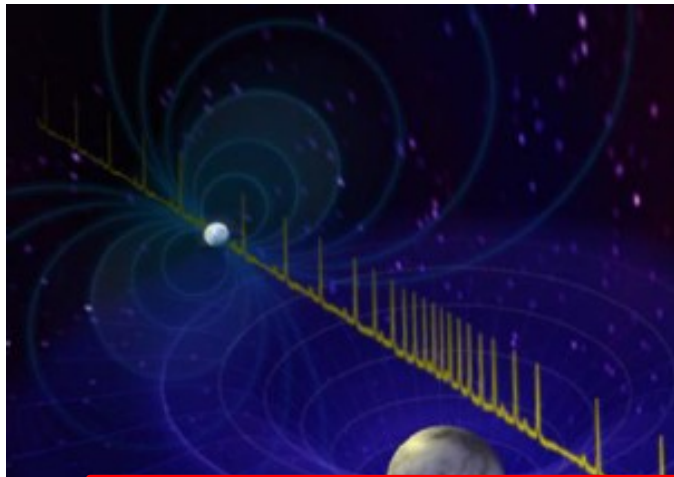
Massive Neutron Star “Shock” (2010)

■ Big news in 2010 autumn

- A neutron star is found to have the mass of $1.97 \pm 0.04 M_{\odot}$ ”
Demorest et al., Nature 467 (2010) 1081 (Oct.28, 2010).

● Based on “kinetic” observable

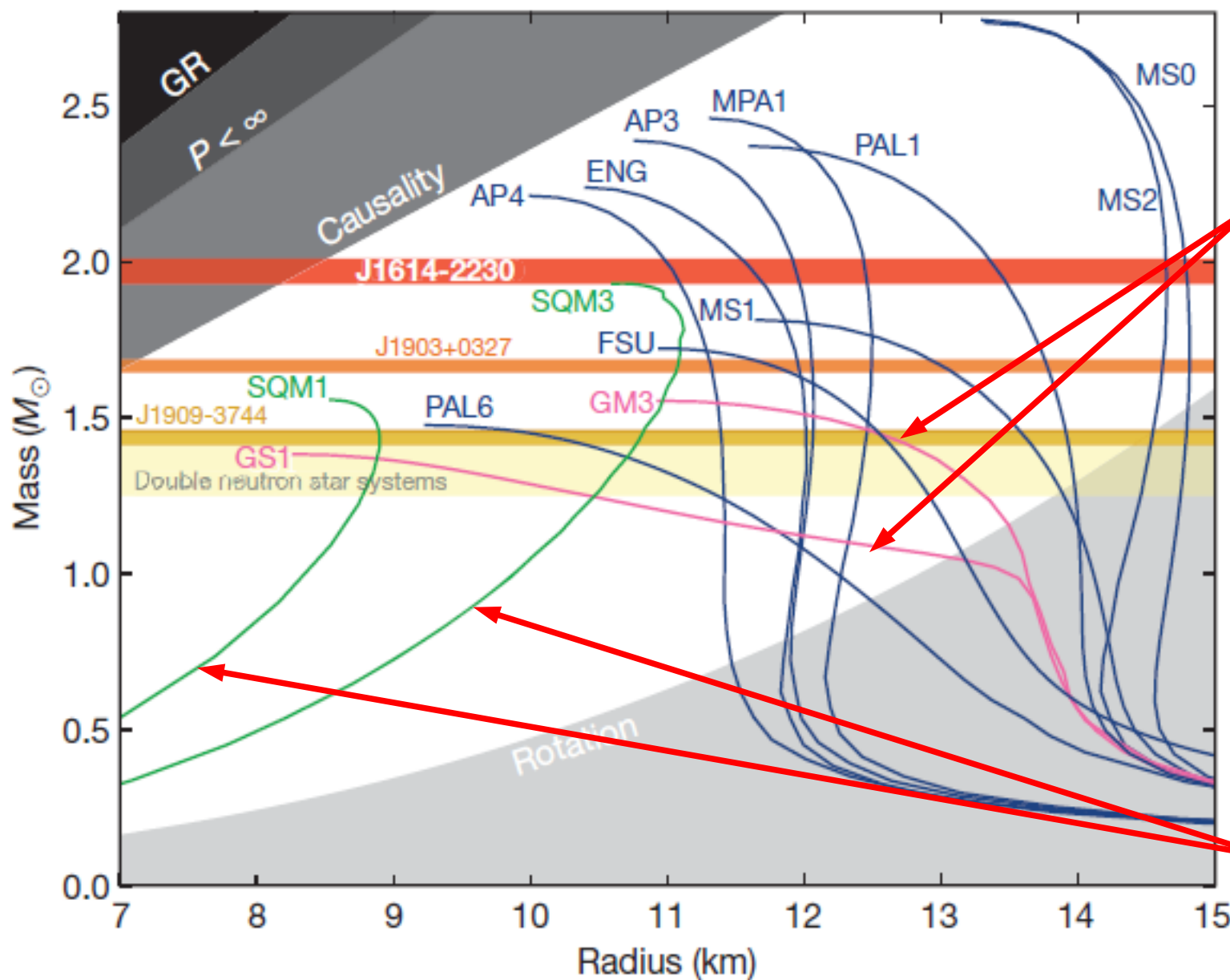
- Photon from the pulsar passes near the companion white dwarf, and Shapiro delay (GR effect) was clearly observed.



signature. We calculate the pulsar mass to be $(1.97 \pm 0.04)M_{\odot}$, which rules out almost all currently proposed²⁻⁵ hyperon or boson condensate equations of state (M_{\odot} , solar mass). Quark matter can support a star this massive only if the quarks are strongly interacting and are therefore not ‘free’ quarks¹².

$1.97 \pm 0.04 M_{\odot}$ Neutron Star

Demorest et al., Nature 467 (2010) 1081 (Oct.28, 2010).



**EOS with
hyperons
or Kaons**

**Quark matter
EOS**

Compact Neutron Star “Shock” (2013)

■ Compact NS “Puzzle”

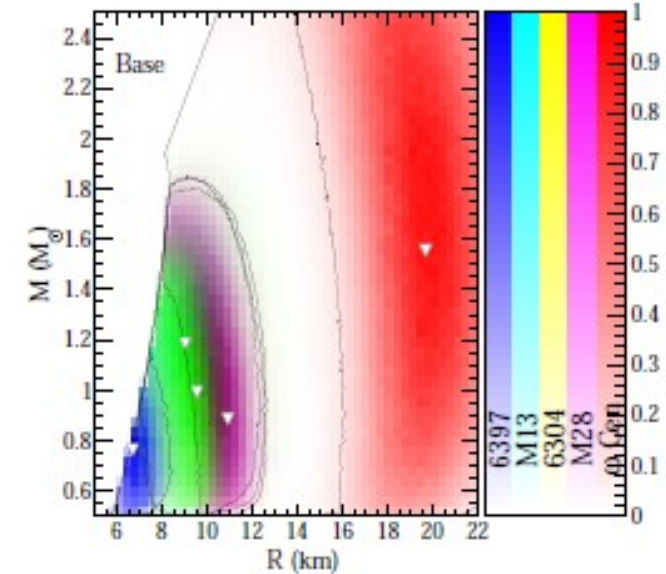
Guillot et al. (2013); Lattimer, Steiner (2013).

- Markov-Chain MC fit of X-ray spectrum
Bayesian analysis
from quiescent Low Mass X-ray Binaries
(qLMXB)

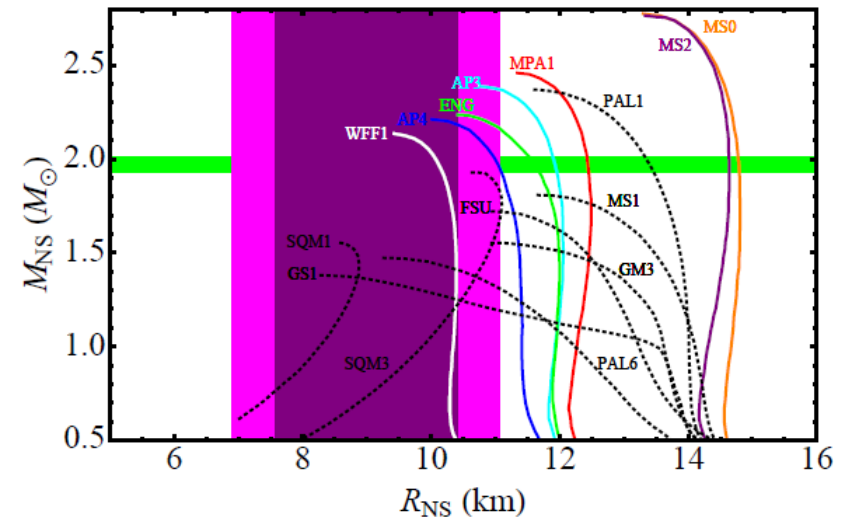
- Guillot+ (2013) results

$$R_{\text{NS}} = 9.1^{+1.3}_{-1.4} \text{ km}$$

→ Most of proposed EOSs are ruled out.
WFF: Wiringa, Fiks, Fabrocini (1988)



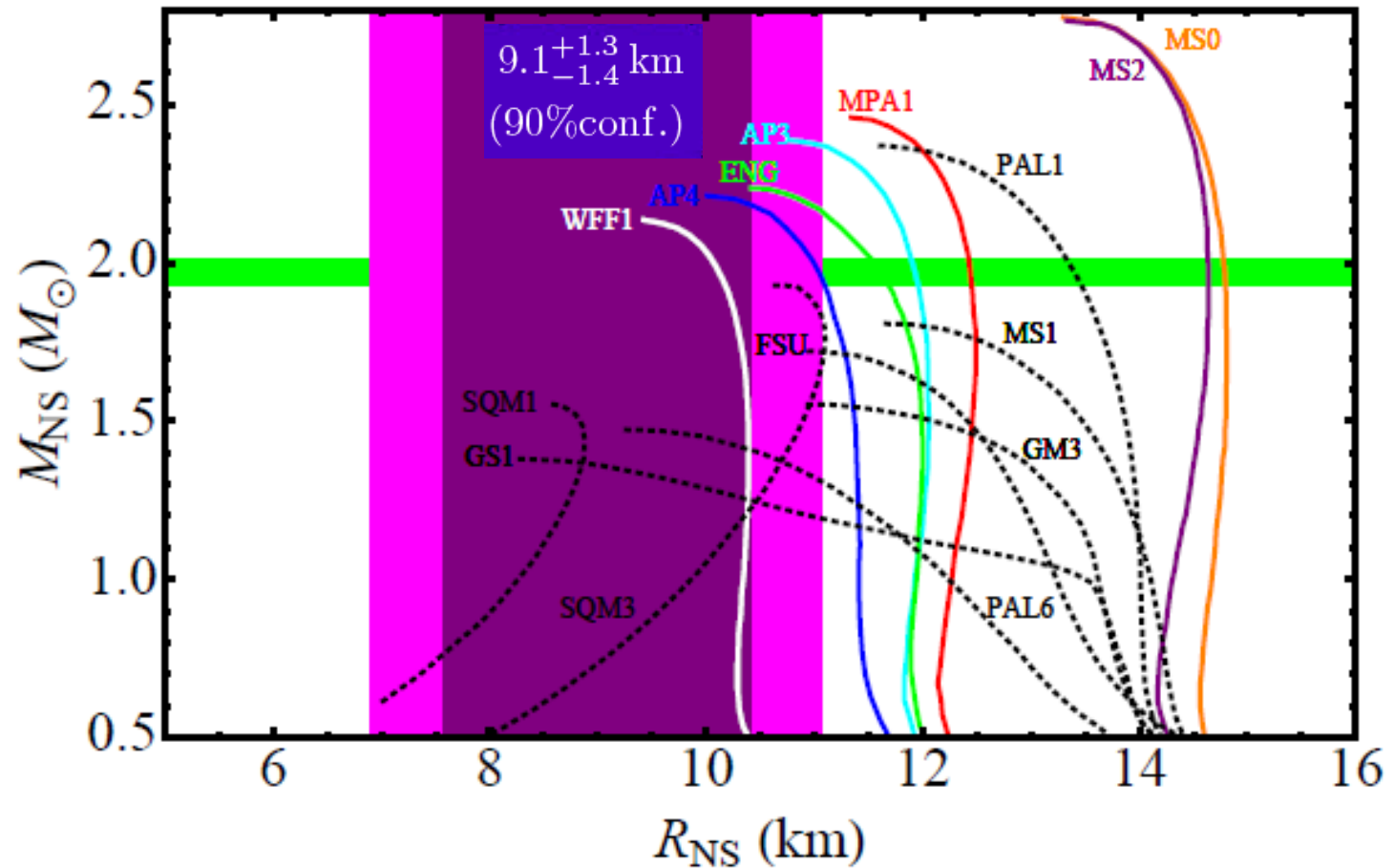
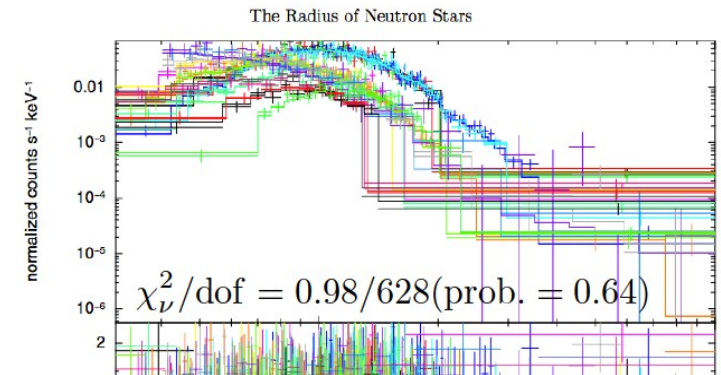
Lattimer, Steiner (2013).



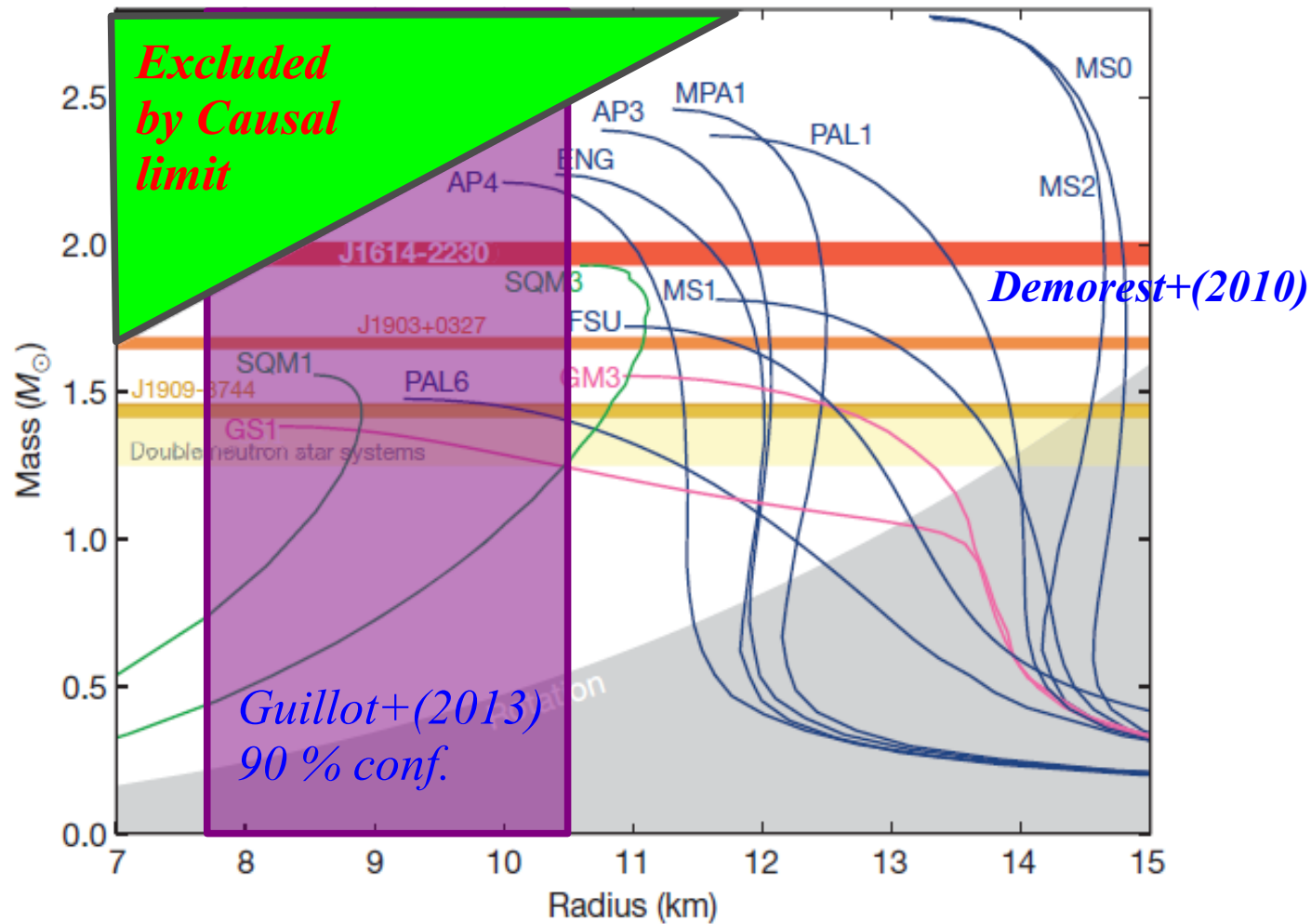
Guillot et al. (2013)

Assumptions

- H atmosphere neutron stars.
- Low B-field ($<10^{10}$ G) neutron stars.
- Emitting isotropically.



Massive and Compact Neutron Star



NS ~ Almost causal limit object ?

Massive and Compact Neutron Star Puzzles

■ Puzzle 1:

Massive ($\sim 2 M_{\odot}$) NS cannot be supported by “standard” EOS with hyperons or kaons, while theoretical model calculations based on laboratory hypernuclear physics experiments predict hyperon appearance.

→ Something is wrong !

■ Puzzle 2:

X-ray spectra from qLMXB seems to imply neutron star masses around 9 km, while most of theoretical model calculations based on laboratory normal and neutron-rich nuclear physics experiments predict neutron star radius above 11 km.

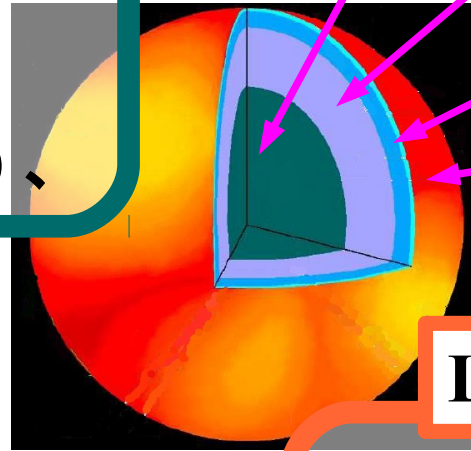
→ Something is wrong !

We need studies of neutron star matter EOS cooperated by Strangeness Nuclear Physics, Neutron-rich Nuclear Physics, Astronomical Observations, and Theories.

Grant-in-Aid Study on Neutron Star Matter

High ρ (Group A)

Hypernuclei, Kaonic nuclei
 YN & YY int.,
 Eff. Interaction
 (Heavy-ion collisions)



Hyperons, mesons, quarks

Asym. nuclear matter
 +elec.+ μ

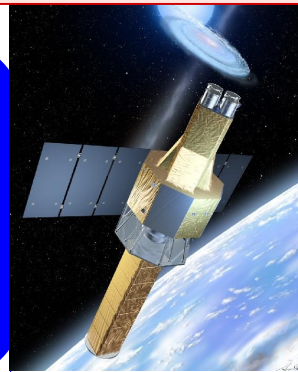
Nuclei+neutron gas+elec.

Nuclei + elec.

Low ρ (Group B)

Sym. E, Pairing gap,
 BEC-BEC cross over,
 Cold atom, Unitary gas

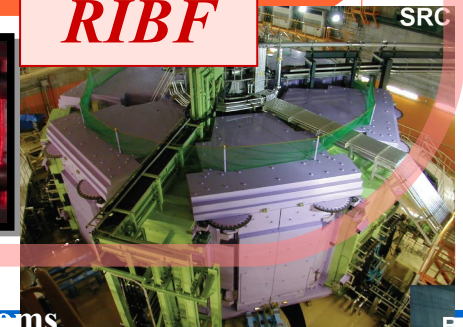
ASTRO-H



NS Obs. (Group C)

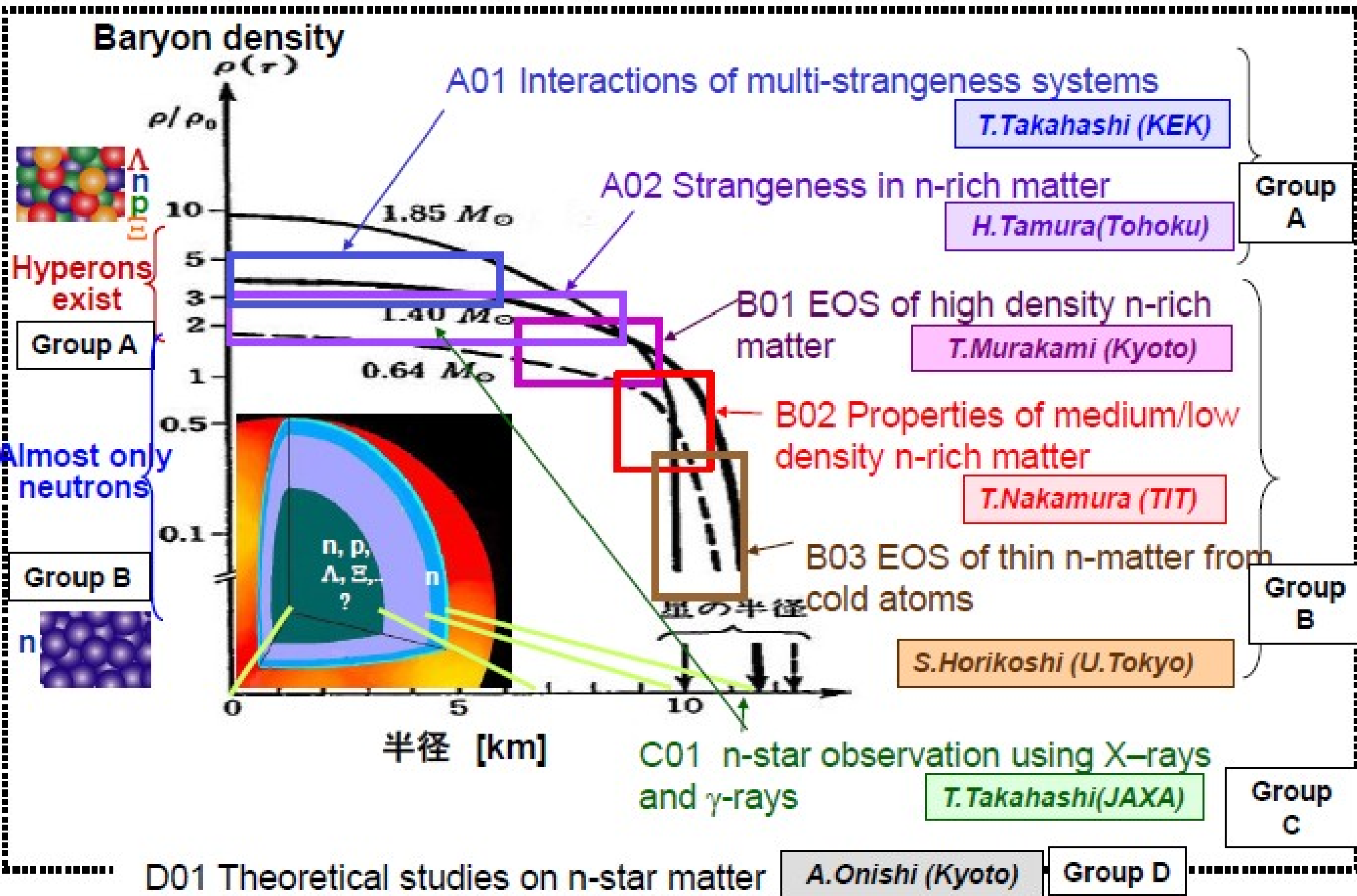
Radius, Mass,
 Temp. (Cooling),
 Star quake, Pasta

RIBF



Theory (Group D)

Groups and research subjects

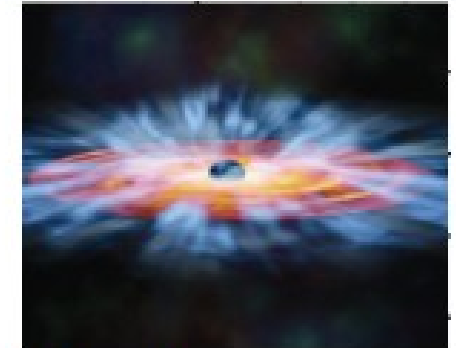




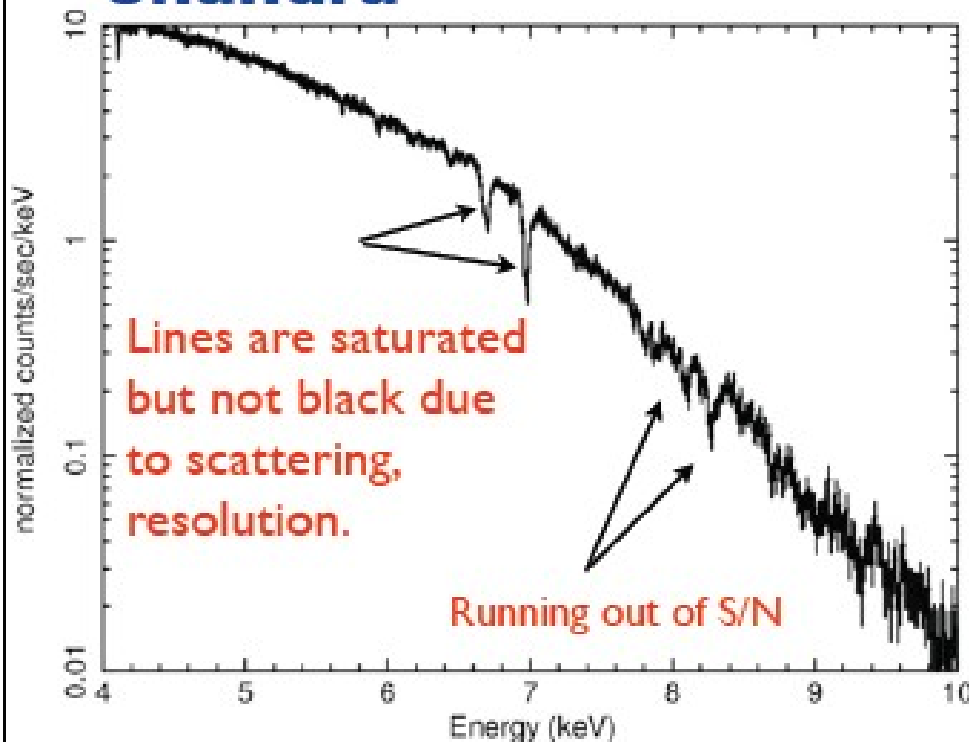
6. ASTRO-H Features --- High Resolution Spectroscopy ---

GRO J1655-40

The superior resolution of SXS in the Fe K band enables the unambiguous detection of weak and narrow lines from a wind.

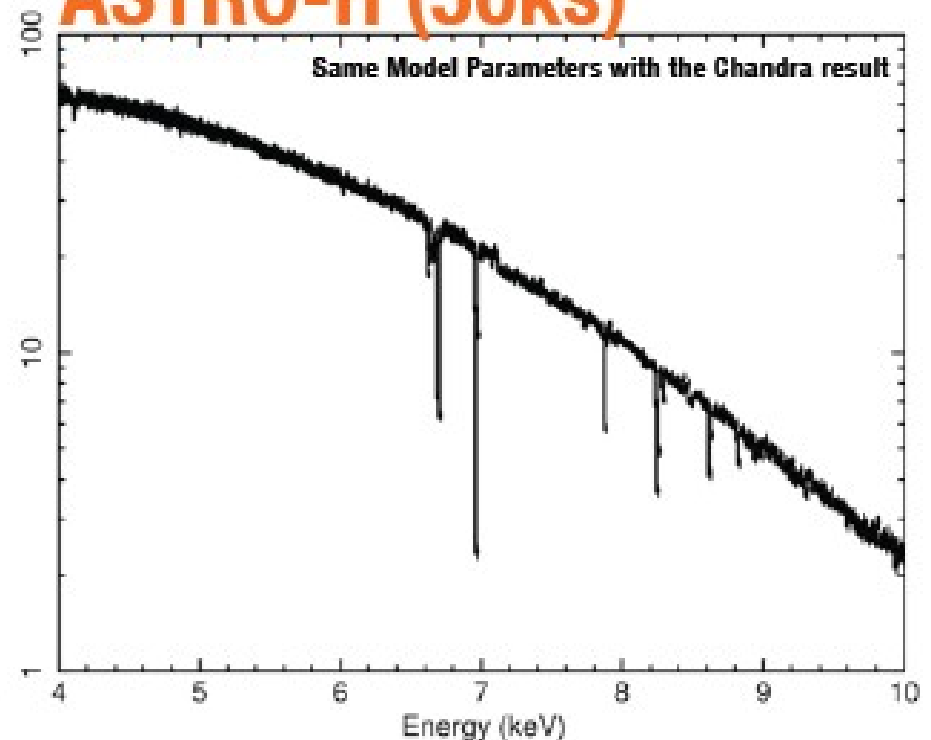


Chandra



(Simulation by J. Miller)

ASTRO-H (50ks)



ASTRO-H SXS can handle 250 cts/s

Recent Challenges

中性子星物質における
対称エネルギーのもたらす影響について

椿原 康介

Osaka Electro-Communication University

in collaboration with

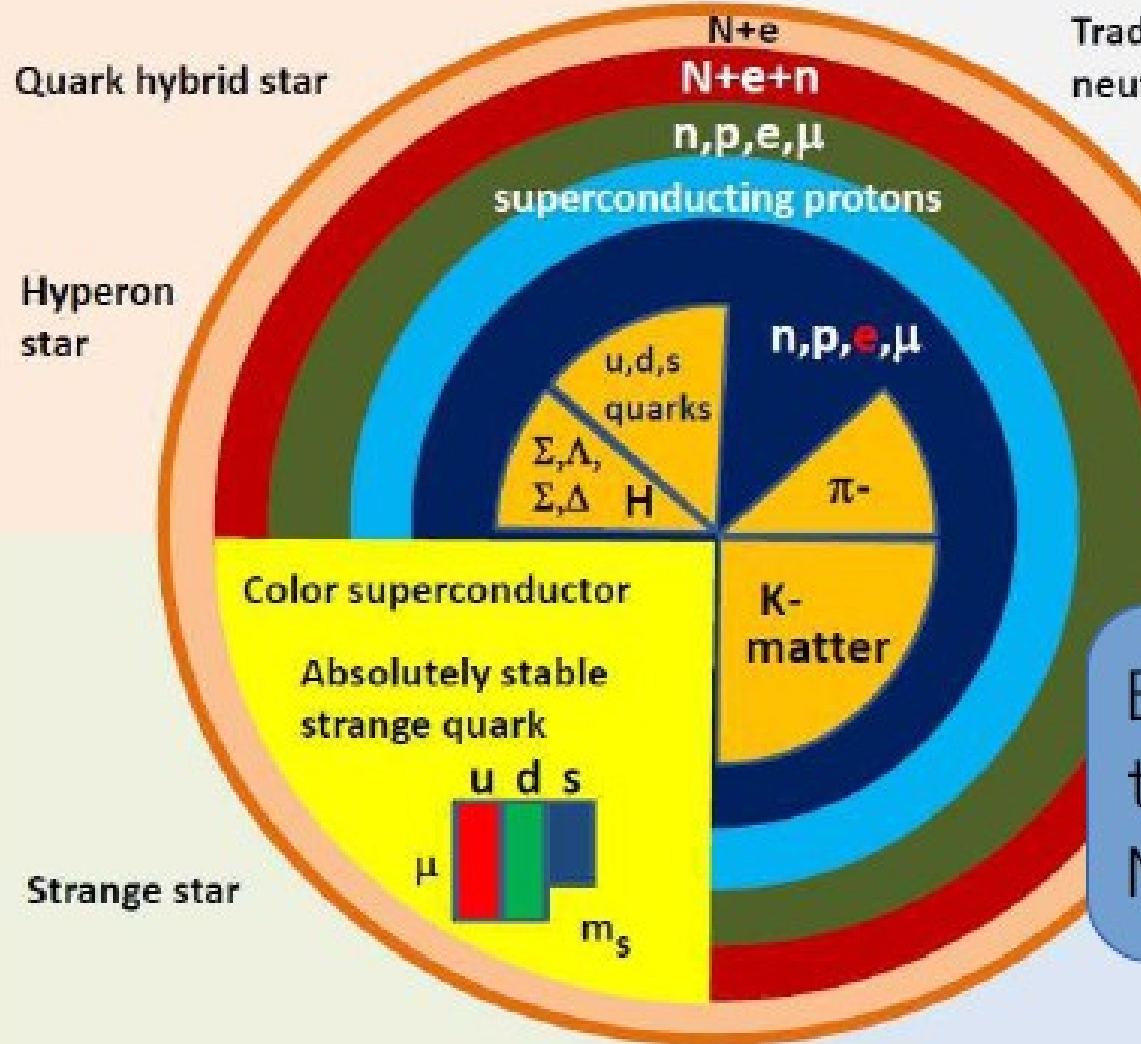
大西 明

YITP, Kyoto University

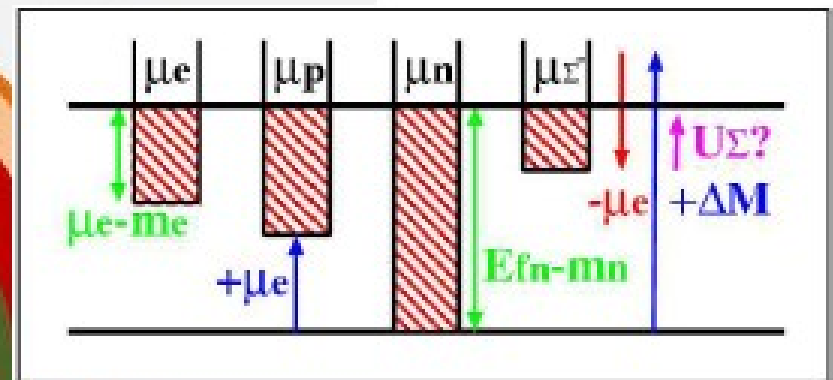
原田 融

Osaka Electro-Communication University

Introduction: Composition of NS



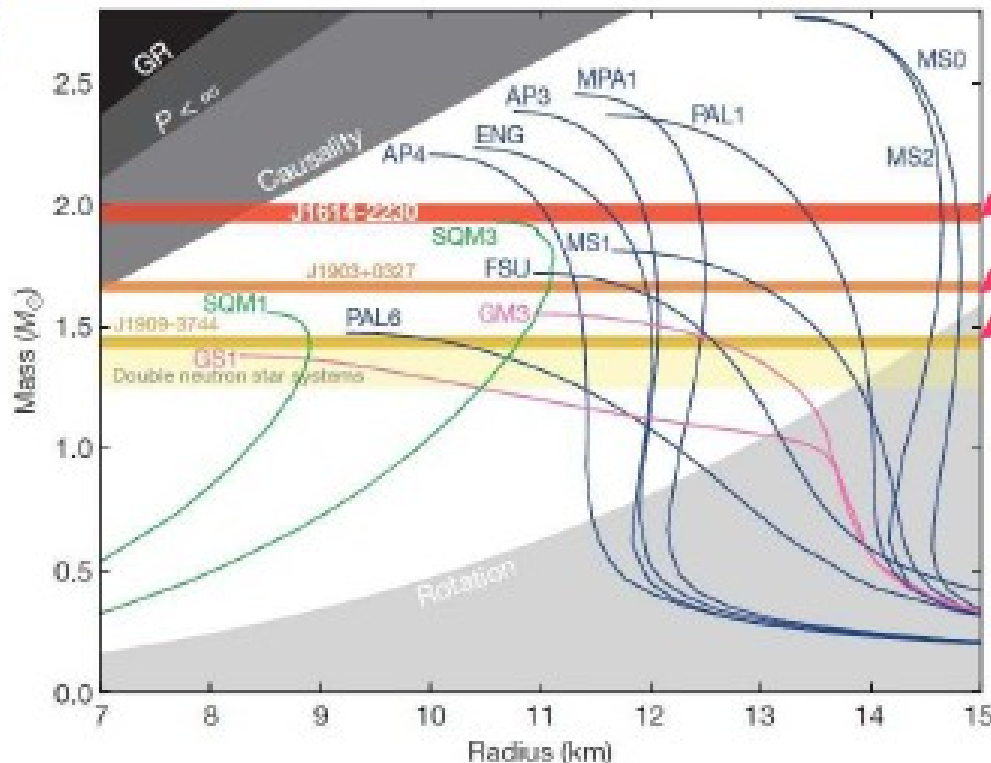
Traditional neutron star



Exotic particle: expected to emerge in the core of NS naturally.

Introduction: Massive NS

- Maximum mass of NS



1.97M_⊙ Obsvd. 2009

1.67M_⊙

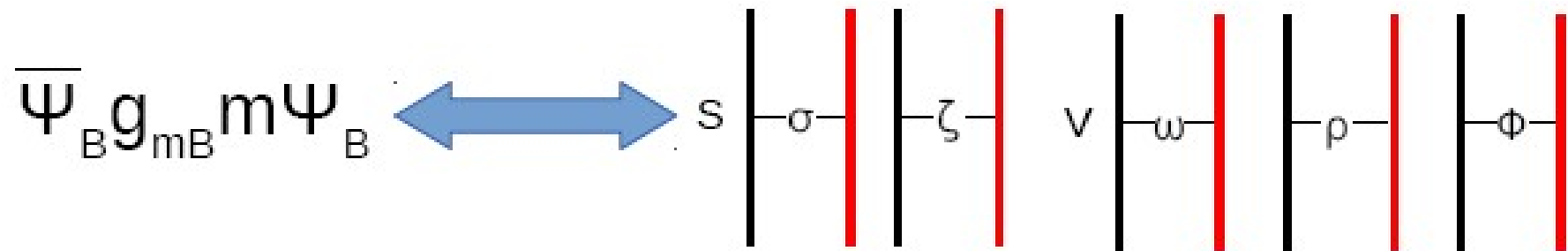
1.44M_⊙ Obsvd. 1974

“Most EoS curves involving exotic matter, such as kaon condensates or hyperons tend to predict maximum masses well below 2.0 M and are therefore ruled out”

Strong restriction to high-dense matter EoS

Introduction: RMF model

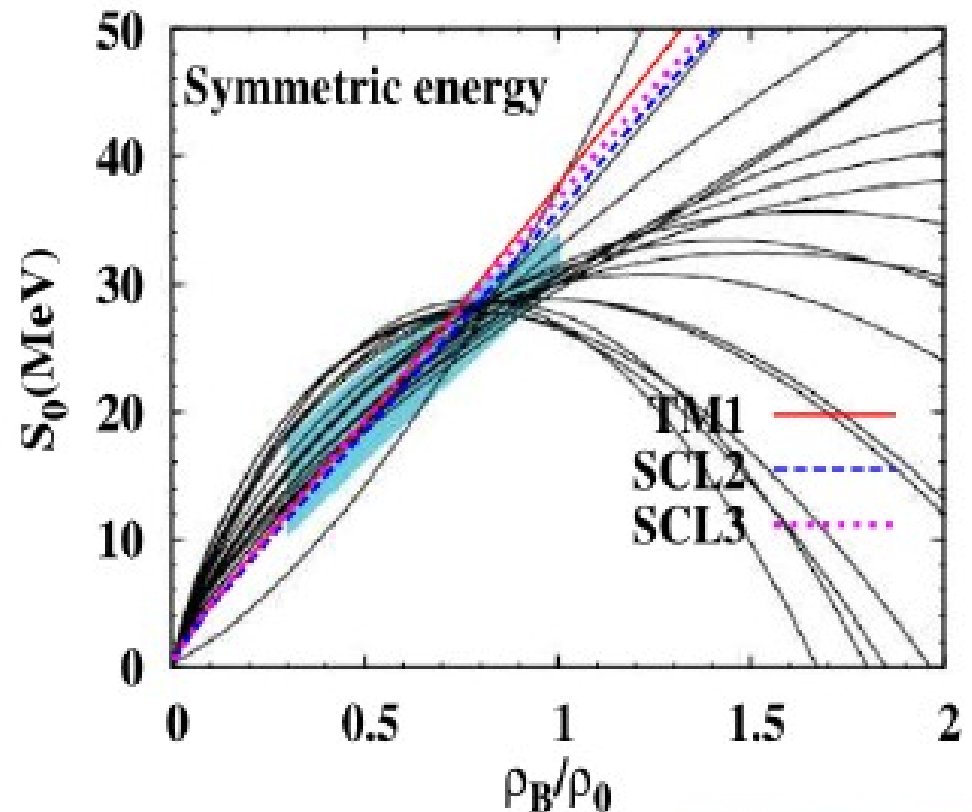
- RMF model: reasonable description to both nuclear matter and finite nuclei
- $n=B/2+M+D=2$ RMF model (+ effective pot.)
 → 2-body interaction (and rel. 3-body corr.)



$$\begin{aligned}
 \mathcal{L} = & \sum_i \bar{\psi} [i\partial - M_i^* - \gamma_\mu V^\mu] - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\
 & + \frac{1}{2} \partial_\mu \varphi_\sigma \partial^\mu \varphi_\sigma - \frac{m_\sigma^2}{2} \varphi_\sigma^2 + \frac{1}{2} \partial_\mu \varphi_\zeta \partial^\mu \varphi_\zeta - \frac{m_\zeta^2}{2} \varphi_\zeta^2 - V_{\sigma\zeta} \\
 & - \frac{1}{4} W_{\mu\nu} W^{\mu\nu} + \frac{m_\omega^2}{2} \omega_\mu \omega^\mu - \frac{1}{4} R_{\mu\nu} R^{\mu\nu} + \frac{m_\rho^2}{2} \rho_\mu \rho^\mu - \frac{1}{4} \phi_{\mu\nu} \phi^{\mu\nu} + \frac{m_\phi^2}{2} \phi_\mu \phi^\mu + \frac{C_\omega}{4} (\omega_\nu \omega^\nu)^2
 \end{aligned}$$

Symmetry Energy

- Nuclear symmetry energy: The size of NS and/or thickness of neutron skin.
- Symmetry energy in RMF
→ Larger L than the suggestion from Exp.
- Density dependence: almost linear since E/A is proportional to ρ field if we adopt $n=2$ RMF Lagrangian.

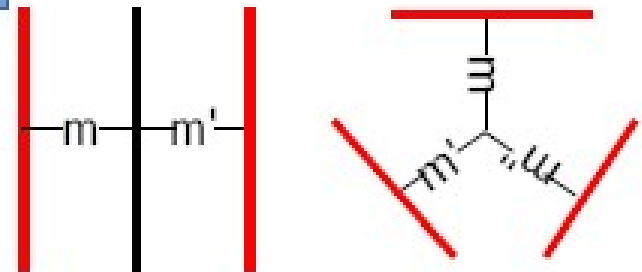


Motivation

- Symmetry energy and its effect to NS-EoS:
Can we modify them by including IV type $n=3$ ρ couplings to an RMF model?

$n=3$ coupl. \rightarrow 3-body int. among baryons

$$\bar{\Psi}_{Bmm'}\Psi_B, \quad m m' m''$$



- Other systematic features:
Can we reasonably reproduce bulk properties of finite/infinite nuclear systems?
- How to justify determined $n=3$ couplings:
Can neutron skin thickness be candidate?

Model description

- $n=B/2+M+D=2$ RMF model + effective pot.

$$\mathcal{L} = \sum_i \bar{\psi} [i\partial - M_i^* - \gamma_\mu V^\mu] - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} \partial_\mu \varphi_\sigma \partial^\mu \varphi_\sigma - \frac{m_\sigma^2}{2} \varphi_\sigma^2 + \frac{1}{2} \partial_\mu \varphi_\zeta \partial^\mu \varphi_\zeta - \frac{m_\zeta^2}{2} \varphi_\zeta^2 - V_{\sigma\zeta} - \frac{1}{4} W_{\mu\nu} W^{\mu\nu} + \frac{m_\omega^2}{2} \omega_\mu \omega^\mu - \frac{1}{4} R_{\mu\nu} R^{\mu\nu} + \frac{m_\rho^2}{2} \rho_\mu \rho^\mu - \frac{1}{4} \phi_{\mu\nu} \phi^{\mu\nu} + \frac{m_\phi^2}{2} \phi_\mu \phi^\mu + \frac{C_{\omega 4}}{4} (\omega_\nu \omega^\nu)^2$$

N=3 Coupl.

Isoscalar \rightarrow two-different-type repulsive int. (TBC5 and 6)

$$\bar{\Psi}_B (g_{\sigma\sigma B} \varphi_\sigma^2 / f_\pi) \Psi_B, \bar{\Psi}_B (g_{\sigma\omega B} \varphi_\sigma \omega / f_\pi) \Psi_B, \bar{\Psi}_B (g_{\omega\omega B} \omega^2 / f_\pi) \Psi_B, \frac{C_{\sigma\omega 2}}{2} f_\pi \varphi_\sigma \omega^2$$

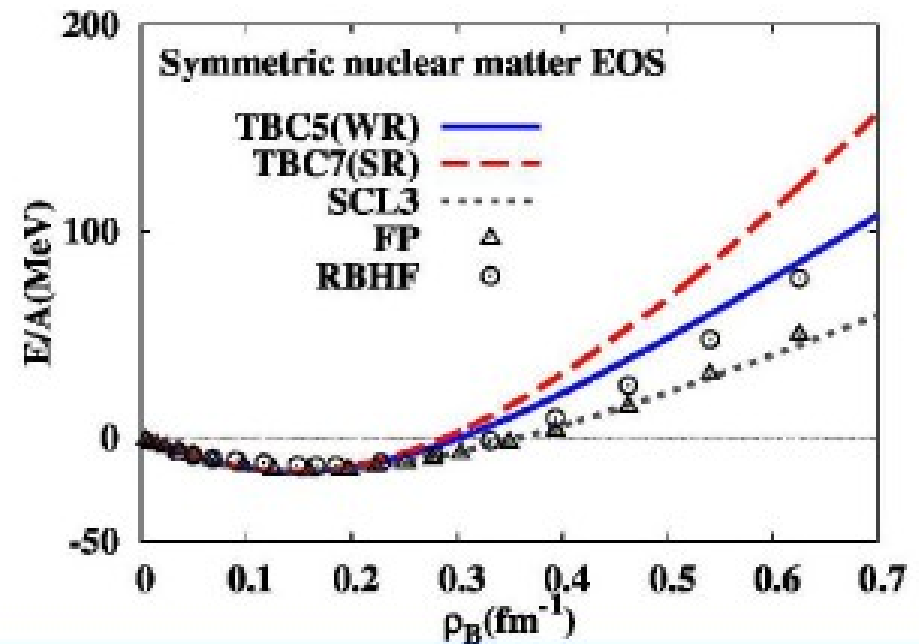
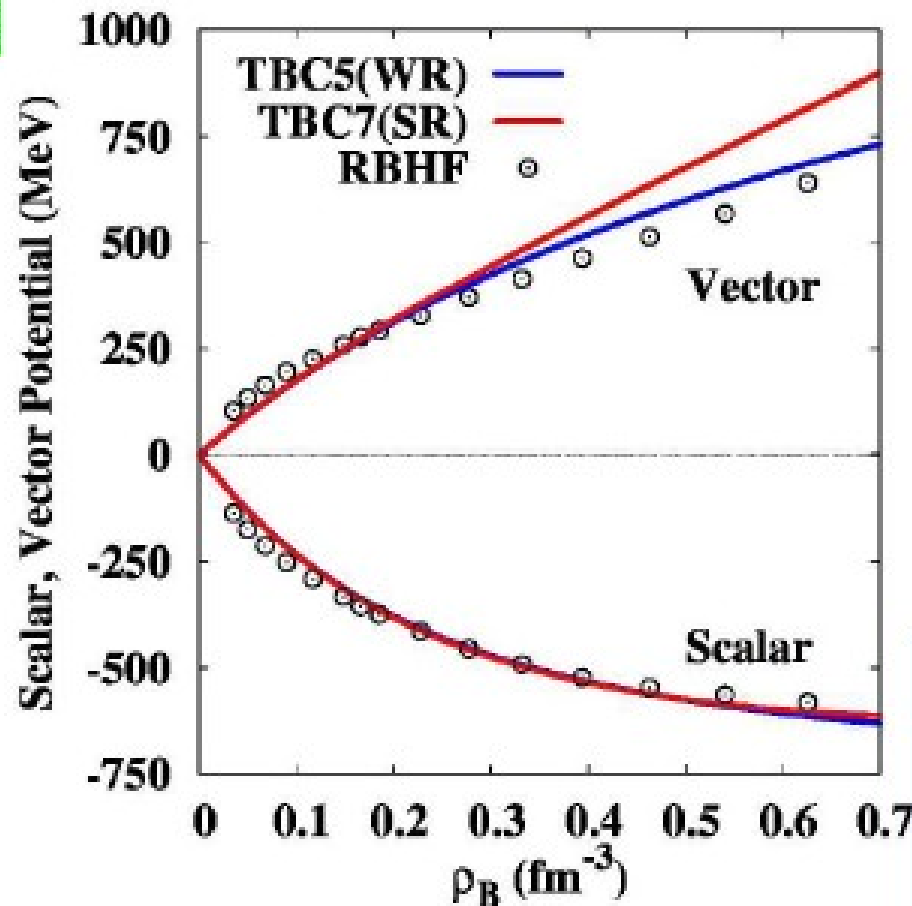
Isovector \rightarrow Newly introduced so as to control IV part of EoS
(Systematics of B/A and symmetry energy at ρ_0)

$$\bar{\Psi}_B (g_{\rho\sigma B} \rho \varphi_\sigma / f_\pi) \Psi_B, \bar{\Psi}_B (g_{\omega\rho B} \omega \rho / f_\pi) \Psi_B, \bar{\Psi}_B (g_{\rho\rho B} \rho^2 / f_\pi) \Psi_B, \frac{C_{\sigma\rho 2}}{2} f_\pi \varphi_\sigma \rho^2$$

Red: Vector, Blue: Scalar type n=3 couplings

Results (1): Nuclear matter

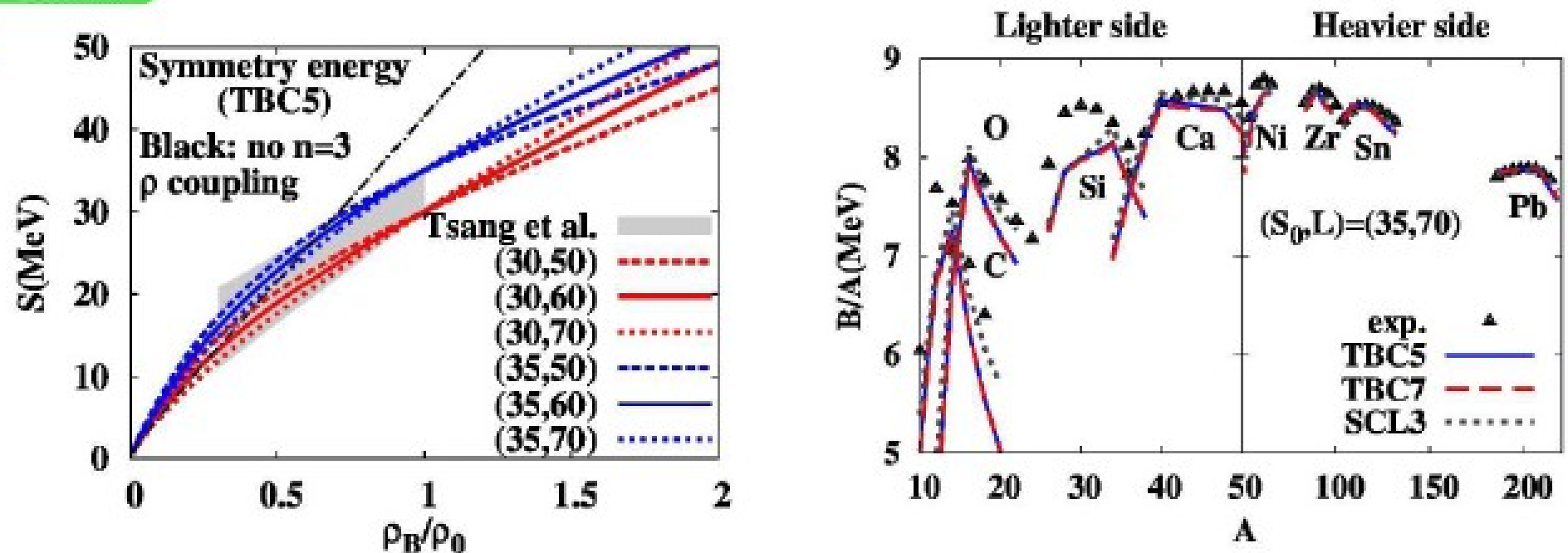
- Symmetric nuclear matter EoS



TBC5: determined so as to reproduce RBHF calc.
TBC7: More repulsive parameter set than TBC5

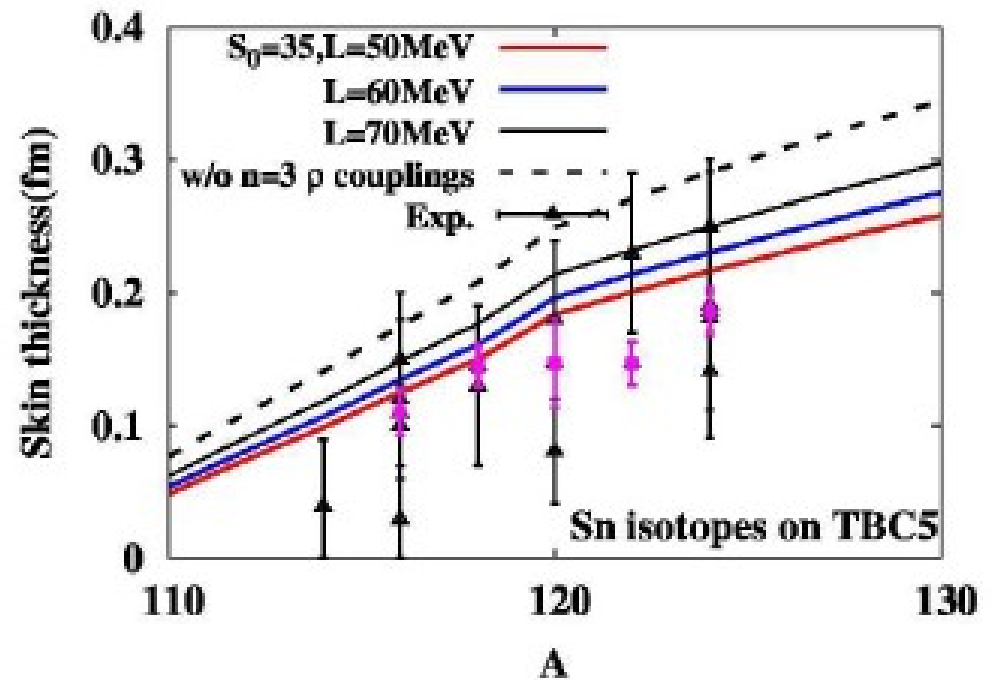
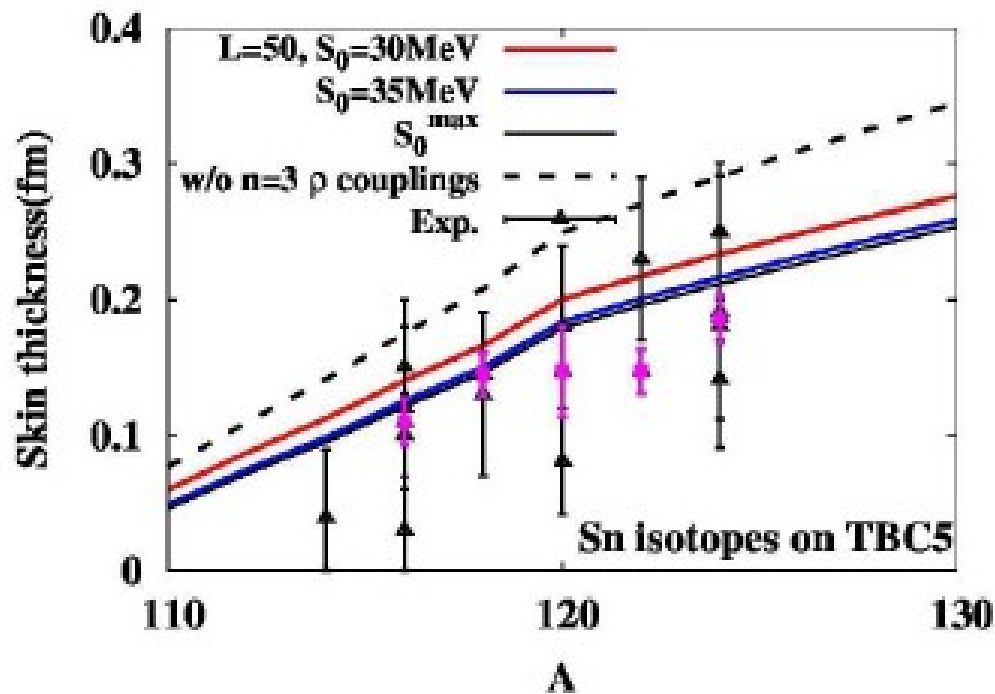
Results(2): Symmetry energy

- Neutron matter w or w/o n=3 ρ coupling



- n=3 ρ couplings: reasonably constrained by (S_0, L)
- Neutron matter: Linear behavior
→ good agreement with HIC suggestion
- B/A of finite nuclei: well-reproduced even if we include n=3 IV type couplings

Results (3): Neutron skin in Sn

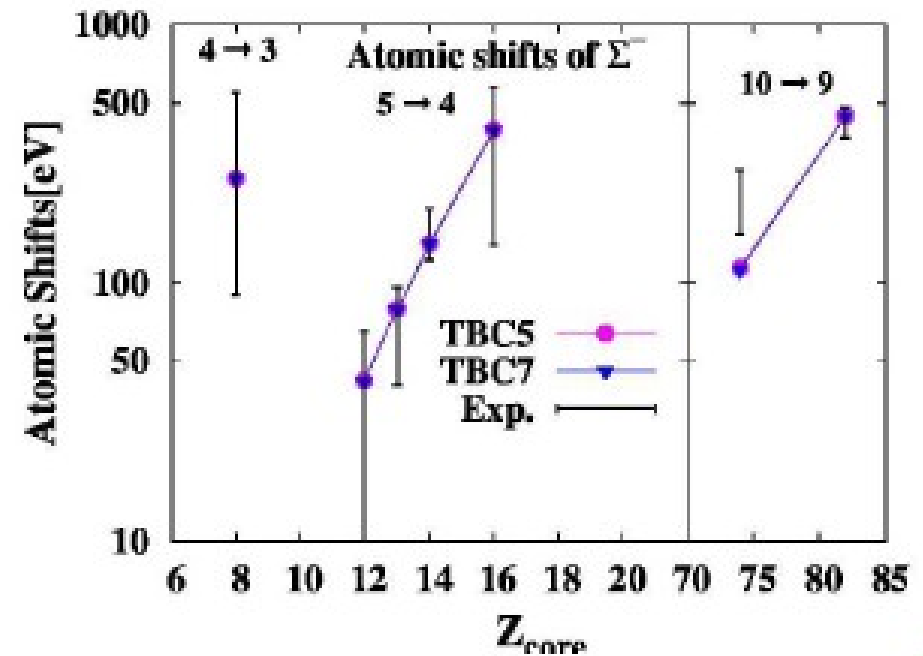
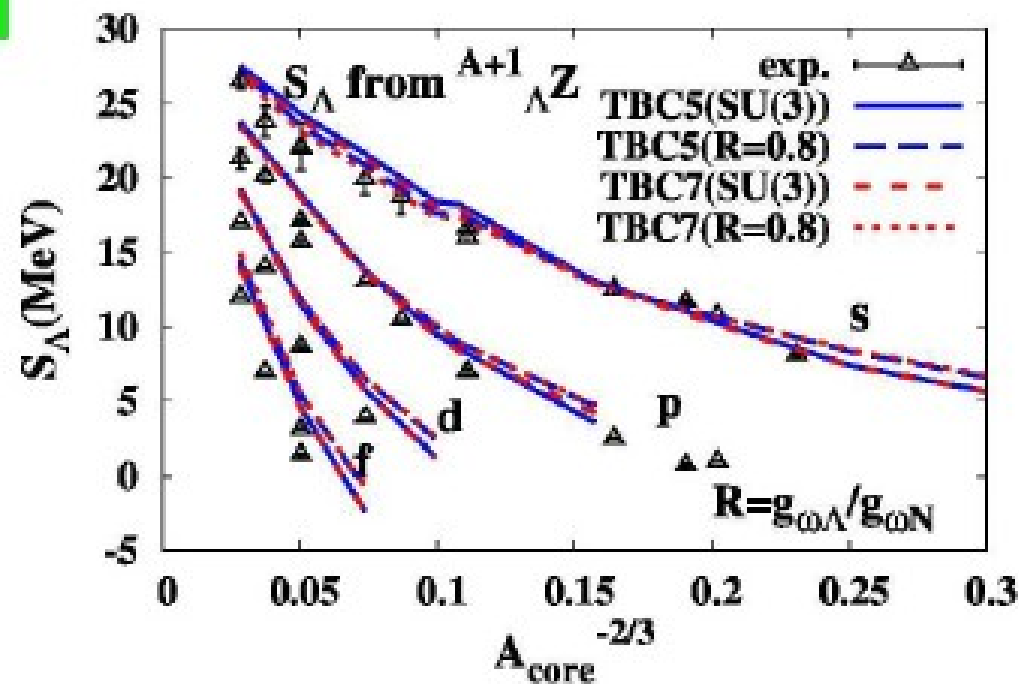


- Symmetry energy: controllable by introducing IV type $n=3$ couplings (TBC5: $S_0 \sim 41.5\text{MeV} \gg 30\text{-}40\text{MeV}$, $L \sim 120\text{MeV} \gg 50\text{-}70\text{MeV}$)
- $(S_0, L)=(35, 70)$: good agreement with the data taken at RCNP

Highlighted RCNP data: Terashima et al. PRC 77 024317

Results (4): NS-EoS

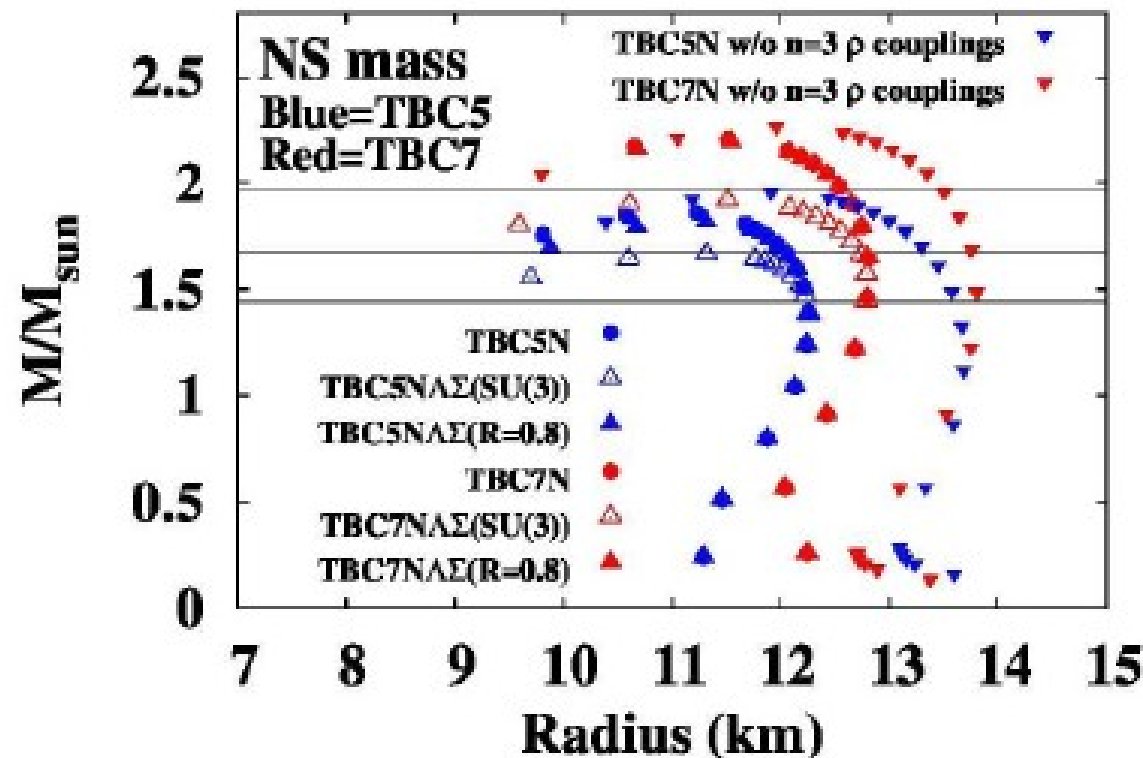
- M-R curve on TBC parameter sets with hyperon



- Hypernuclear data: well reproduced even if we introduce SU(3) violation
- Introduction of $n=3$ ρ coupling: Large modification to the M-R relation \leftrightarrow not so large modification to maximum mass of NS

Results (4): NS-EoS

- M-R curve on TBC parameter sets with hyperon



- Hypernuclear data: well reproduced even if we introduce SU(3) violation
- Introduction of n=3 ρ coupling: Large modification to the M-R relation
↔ not so large modification to maximum mass of NS

Summary

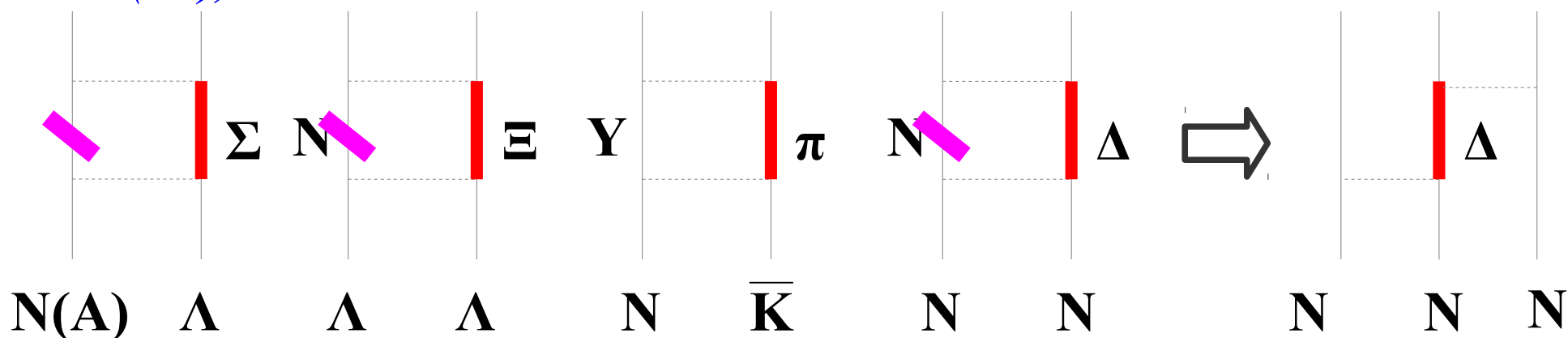
- IV-type $n=3$ TBC in RMF model:
Well-constrained by the data of finite and infinite nuclear systems
- Symmetry energy: Controllable by introducing the $n=3$ TBC and good agreement in HIC results
- Neutron skin: good constraint to (S_0, L)
>> (35MeV, 50MeV) case seems to be in agreement with RCNP data
- If combined with hypernuclear data, TBC7(SR) parameter set with SU(3)-violation in m_ν -B coupling can explain massive 2-times-solar mass.

Do we have strangeness hadrons in NSs ?

- A way to answer massive neutron star puzzle = 3-body repulsion
Nishizaki, Takatsuka, Yamamoto ('03); Doi et al.(HALQCD)

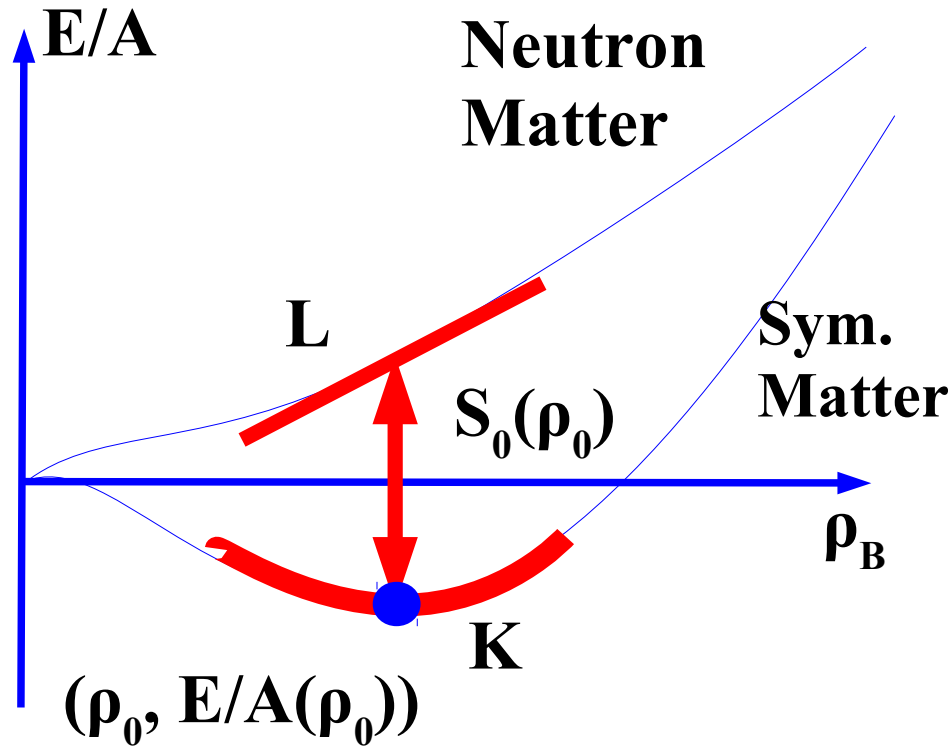
- One of the possible origin of 3-body repulsion
→ Pauli blocking in the intermediate state with resonances

Kohno ('13), Otsuka et al.



- Can we determine these couplings ?
 - Detailed study of hypernuclear properties
incl. n-rich hypernuclei and S=2 hypernuclei.

Nuclear Matter EOS



$$E/A(\rho, \delta) = \varepsilon(\rho) + E_{\text{sym}}(\rho)\delta^2 + O(\delta^4)$$

Symmetric Nuclear Matter

$$\varepsilon(\rho) = \varepsilon(\rho_0) + \frac{K(\rho - \rho_0)^2}{18\rho_0^2} + O((\rho - \rho_0)^3)$$

Symmetry Energy ($\delta = (N - Z)/A = 1 - 2Y_p$)

$$E_{\text{sym}}(\rho) = S_0 + \frac{L(\rho - \rho_0)}{3\rho_0} + \frac{K_{\text{sym}}(\rho - \rho_0)^2}{18\rho_0^2} + O((\rho - \rho_0)^3)$$

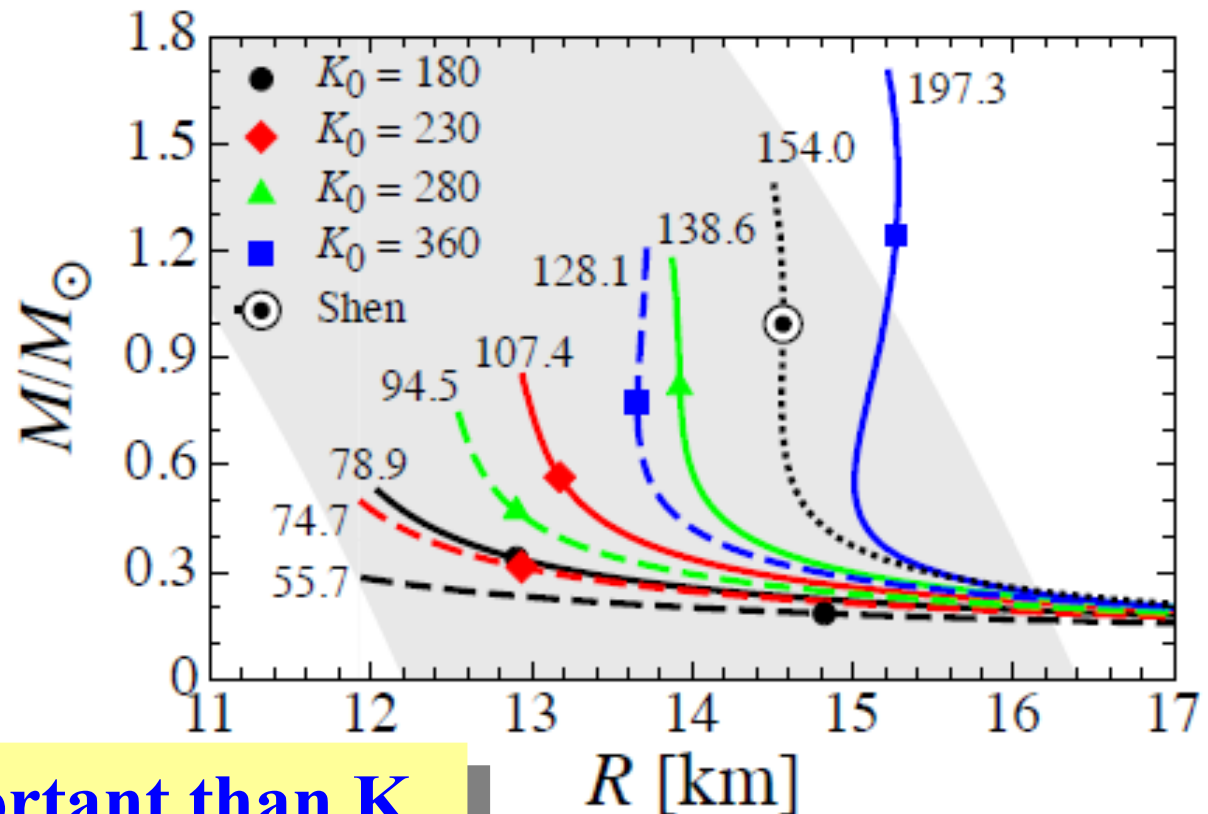
K or L

- We have two EOS parameters, which are the derivatives of E, K and L. Which mainly determines the M-R relation ?

→ Answer

One parameter $\eta = (KL^2)^{1/3}$ characterizes M-R curve (for low mass neutron stars) !

Sotani, Iida, Oyamatsu, AO (to be submitted)



L is twice more important than K.

Contents and Summary

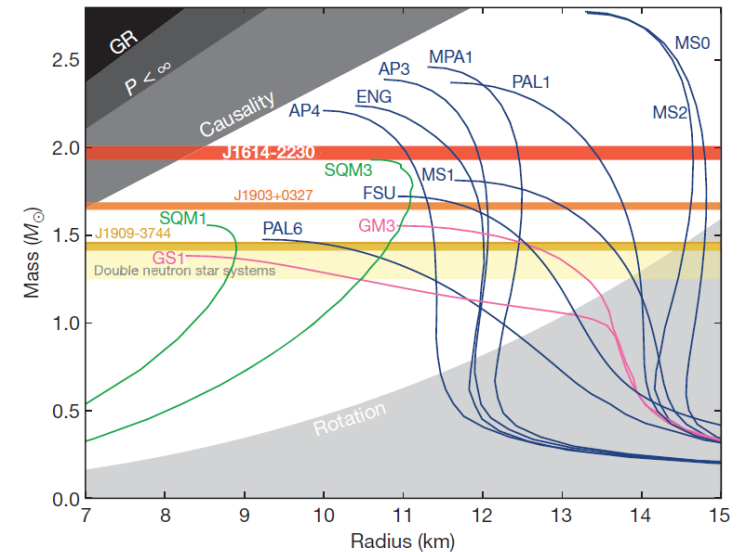
■ Neutron star (NS)

- Giant neutron-rich hypernuclei & Almost causal limit object
- Massive ($\sim 2 M_{\odot}$) & Compact (~ 9 km) NS puzzle.
- Grant-in-aid study on NS matter
J-PARC + RIBF + ASTRO-H + Theory

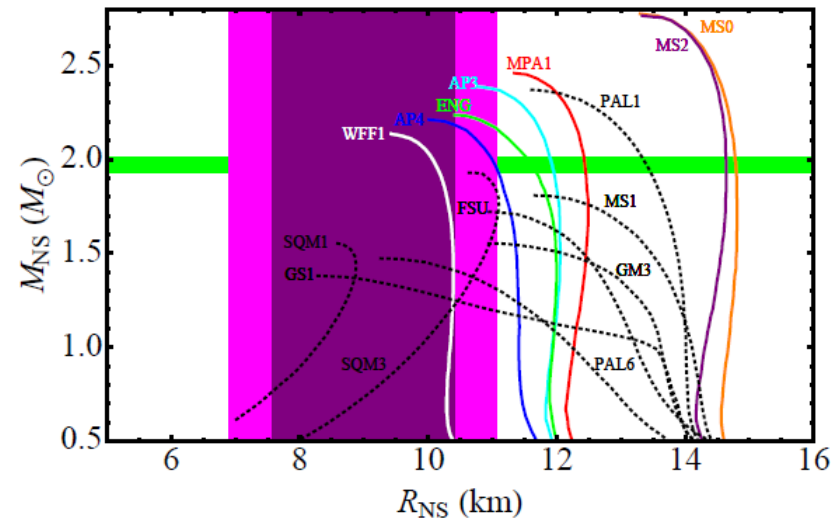
■ Challenges in neutron star matter physics

- How can we solve massive NS puzzle ?
→ Three-body force as a candidate
- What determines M-R relation of low-mass neutron stars ?
→ New parameter $\eta = (KL^2)^{1/3}$

We have big challenges in neutron star physics !



Demorest et al. (2010)



Guillot et al. (2013)

Thank you for your attention.