

Entropy production in the classical Yang-Mills theory as the coherent state dynamics

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in collaboration with

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The Approach to Equilibrium in Strongly Interacting Matter

April 2-4, 2014, BNL

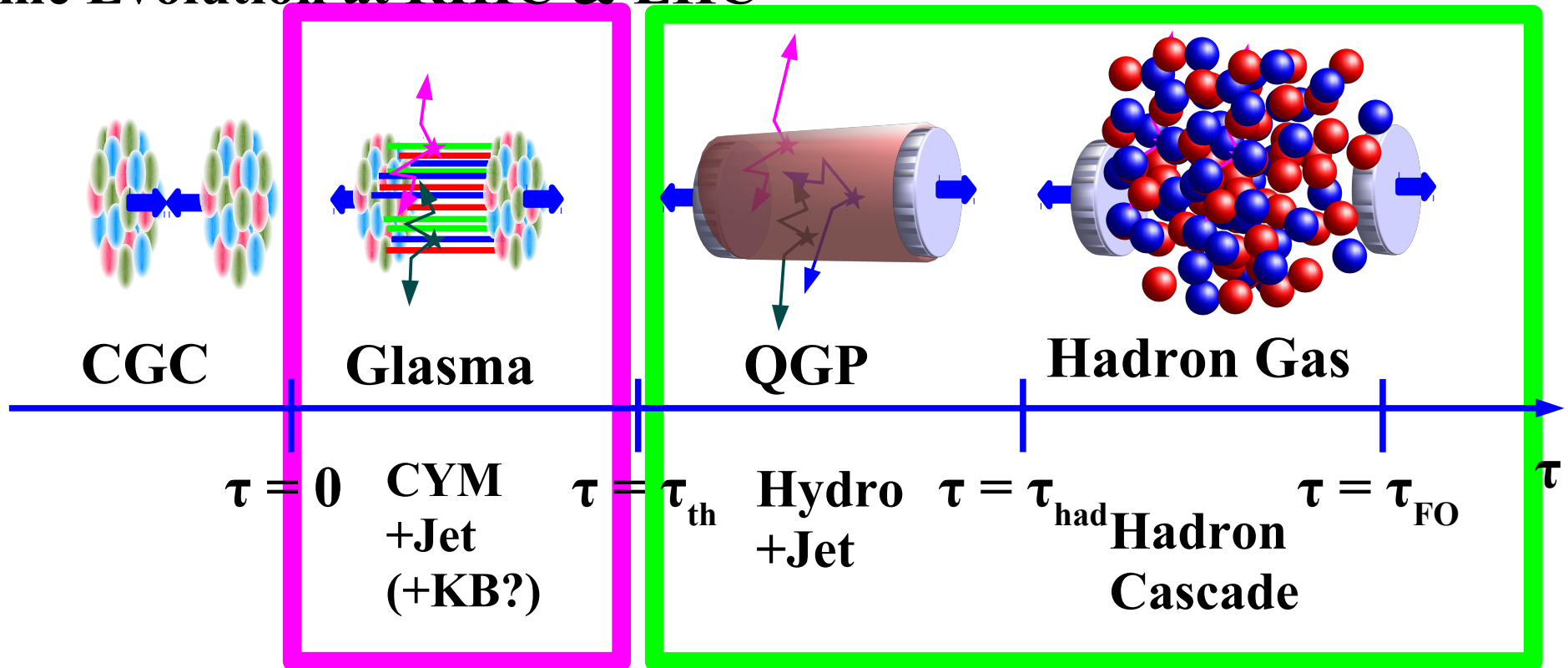
Based on the works,

- *Towards a Theory of Entropy Production in the Little and Big Bang*, T. Kunihiro, B. Müller, A. Ohnishi, A. Schäfer, Prog. Theor. Phys. 121 ('09), 555 [arXiv:0809.4831].
- *Chaotic behavior in classical Yang-Mills dynamics*, T. Kunihiro, B. Müller, A. Ohnishi, A. Schäfer, T. T. Takahashi, A. Yamamoto, Phys. Rev. D 82 (2010), 114015 [arXiv:1008.1156].
- *Entropy production in classical Yang-Mills theory from Glasma initial conditions*, H. Iida, T. Kunihiro, B. Müller, A. Ohnishi, A. Schäfer, T. T. Takahashi, Phys. Rev. D 88 (2013), 094006 [arXiv:1304.1807].
- H. Iida, T. Kunihiro, A. Ohnishi, T.T. Takahashi, arXiv:1404.xxxx



Thermalization in High-Energy Heavy-Ion Collisions

Time Evolution at RHIC & LHC



Theor. Challenges

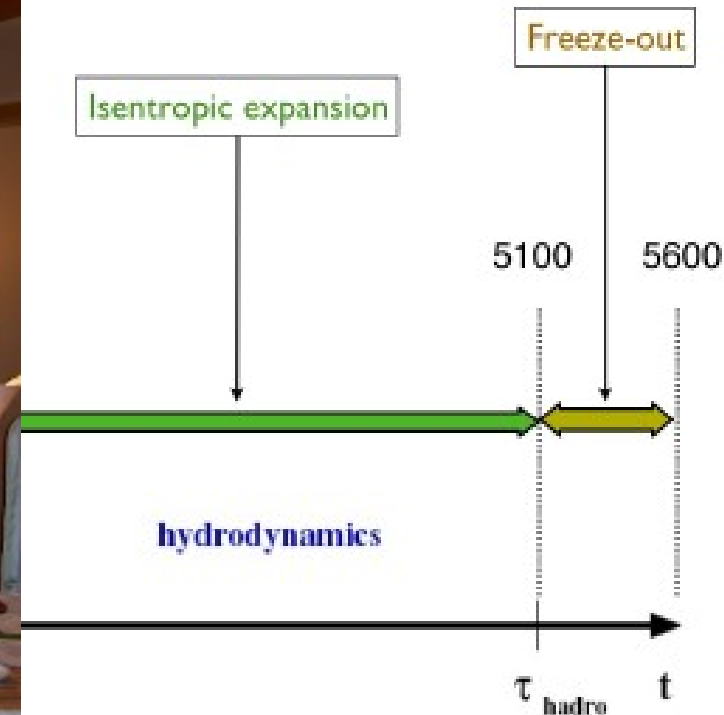
- Thermalization under dynamical classical field
- Theoretically interesting and Phenomenologically important.
dN/ d η , init. cond. of hydro.

Phen. Challenges

- flow, jet, hard probes
→ hydro., transport coef., E-loss, hadron prop., phase diagram, ...

Entropy Production in Heavy-Ion Collisions

- We have been working on entropy production in non-equilibrium stage, since the international Molecule-type workshop on “Entropy Production before QGP” (2008.08.01-28) (A. Schafer, R. Fries, B. Mueller, M. Strickland, T. Schafer, M. Natsuume, Y. Nara, T. Hirano, K. Fukushima, T. Kunihiro, AO)



Muller, Schafer ('11)

Entropy production in Classical Yang-Mills dynamics

- Perturbative estimate of thermalization time is longer than expected from hydrodynamics simulations.
→ Classical Yang-Mills field is expected to play a role of entropy prod.
Baier, A.Mueller, Schiff, Son ('01); Chatterjee, Srivastava ('09), Heinz, Kolb ('02)

- How does CYM field have entropy ? Chaos & Decoherence !

- Entropy from chaoticity

- ◆ (Husimi-)Wehrl entropy *Kunihiro, B. Müller, AO, Schäfer ('09)*

$$S_{Wehrl} = - \int \frac{d^n x d^n p}{(2\pi)^n} H \log H \quad (H = \text{phase space prob. fn.})$$

- ◆ Entropy production rate = Kolmogorov-Sinai entropy

*B. Muller, Trayanov ('92), Biro, Gong, B. Muller ('94), Bolte, B. Muller ('00).
Kunihiro, B. Müller, AO, Schäfer, Takahashi, Yamamoto ('10), Iida,
Kunihiro, B.Müller, AO, Schäfer, Takahashi ('13)*

- Decoherence entropy

*B.Muller, Schafer ('03, '06), Fries, B. Muller, Schafer ('09), Iida, Kunihiro, AO,
Takahashi ('14)*

- Classical Statistical simulation

Berges, Scheffler, Sexty ('08), Epelbaum, Gelis ('13)

Contents

We discuss the CYM entropy and its production rate with emphasis on the decoherence entropy

■ Introduction

■ Entropy production from chaotic nature of CYM

- Chaoticity, Lyapunov exponent, and Kolmogorov-Sinai entropy

- KS entropy from CYM

T. Kunihiro, B. Müller, AO, A. Schäfer, T.T. Takahashi, A. Yamamoto, PRD82('10),114015 [arXiv:1008.1156].

H.Iida, T.Kunihiro, B.Müller, AO, A.Schäfer, T.T.Takahashi, PRD88('13),094006 [arXiv:1304.1807].

■ CYM as a coherent state and decoherence entropy

- Decoherence entropy

- CYM as a coherent state

- Decoherence entropy from CYM dynamics

H.Iida, T.Kunihiro, AO, T.T.Takahashi, arXiv:1404.xxxx (to be submitted soon).

■ Summary

*Entropy production
from chaotic nature of CYM*

Chaoticity, Lyapunov exponent, and KS entropy

- Entropy in classical dynamics = Wehrl entropy

$$S = - \int d\Gamma H \log H$$

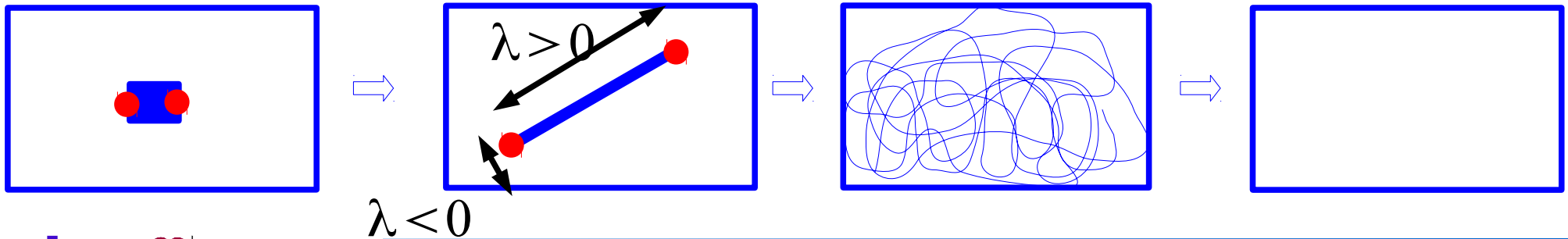
($d\Gamma = dx dp$ = phase space, H = phase space dist. fn., e.g. Husimi fn.)

- Lyapunov exponent and Kolmogorov-Sinai entropy

$$\delta X_i(t) = \delta X_i(t_0) \exp[\lambda_i(t-t_0)] \quad (X = (x, p)),$$

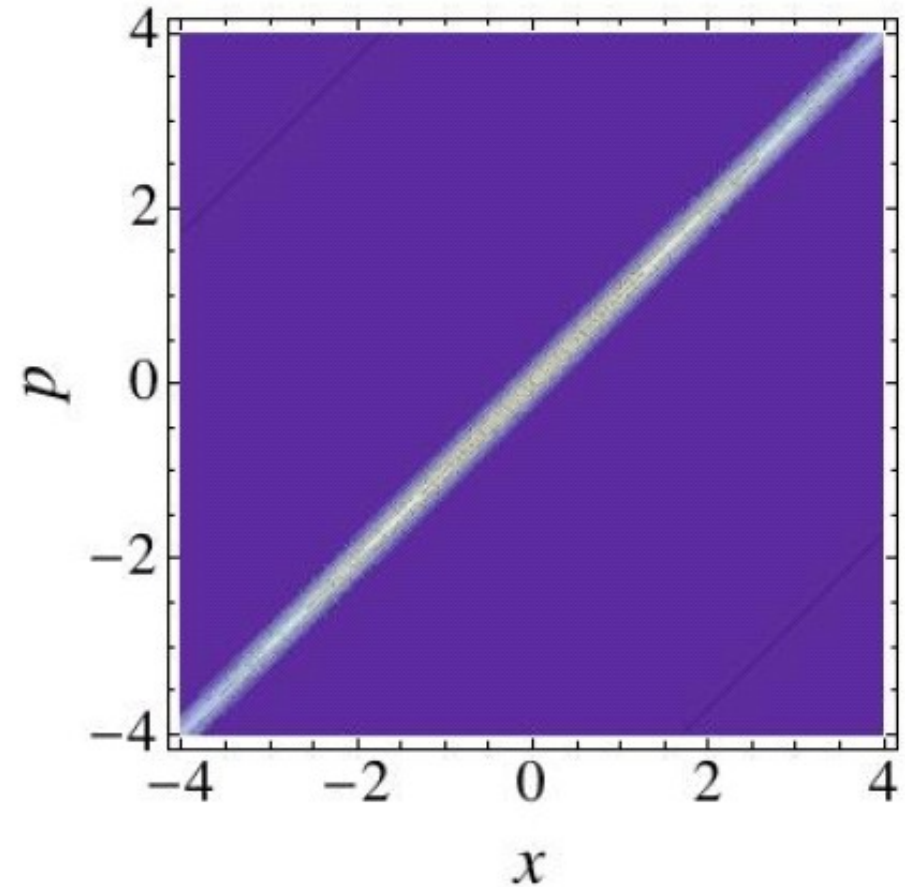
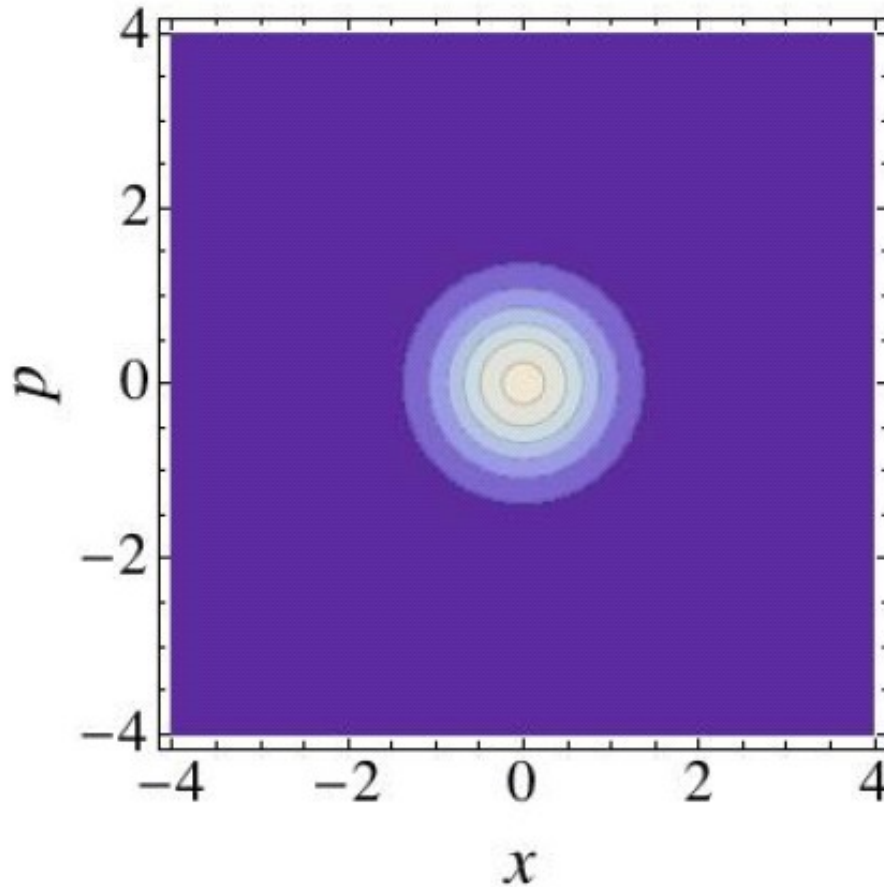
$$dS/dt = S_{\text{KS}} \equiv \sum_{i, \lambda_i > 0} \lambda_i$$

- δX = difference of two trajectories from adjacent initial conditions
 λ = initial state sensitivity (Lyapunov exponent, measure of chaoticity)
- When $\lambda > 0$, exponentially growing number of phase space cells are visited
→ phase space dist. fn. becomes smooth after proper coarse graining
→ entropy production (Kolmogorov-Sinai entropy)



Evolution of the Wigner Function

- Liouville theorem → conservation of the phase space volume
 - Exponential growth in $(x+p/\lambda)$, Exponential narrowing in $(x-p/\lambda)$

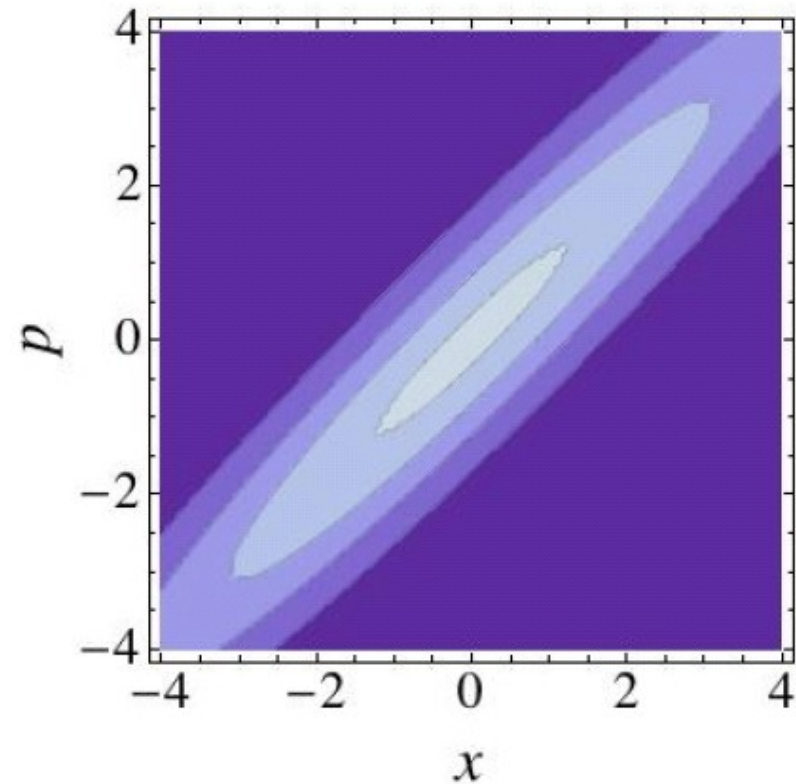
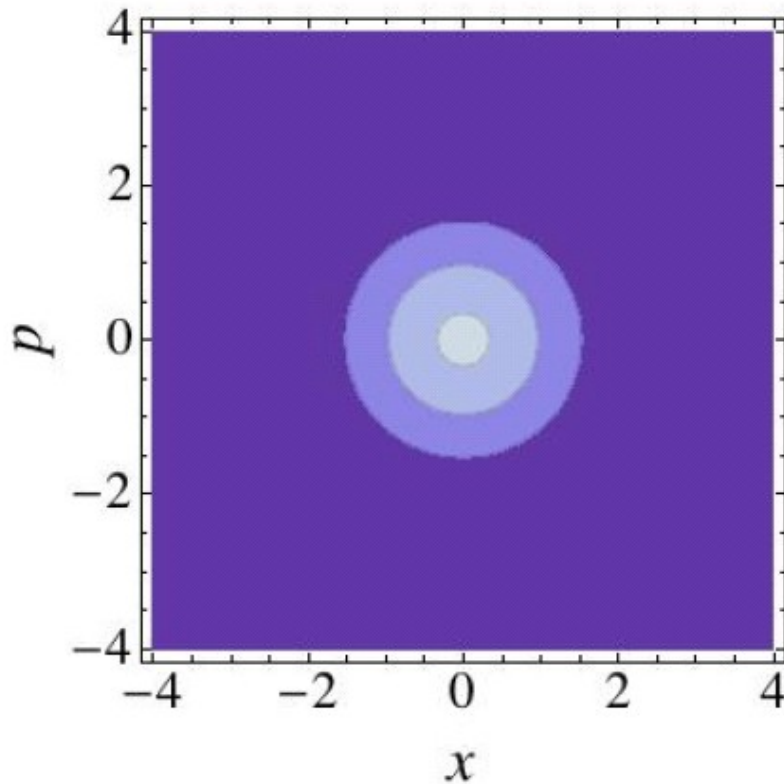


Kunihiro, Müller, AO, Schäfer ('09)

$$\hat{\mathcal{H}} = \frac{1}{2}\hat{p}^2 - \frac{1}{2}\lambda^2\hat{x}^2 \quad \lambda=1, \lambda t=0, 2$$

Evolution of the Husimi Function

- Coherent state broadening of phase space
 - Minimum width in $(x-p/\lambda) \rightarrow$ phase space dist. func. is smeared !



Kunihiro, Müller, AO, Schäfer ('09)

$$\hat{\mathcal{H}} = \frac{1}{2}\hat{p}^2 - \frac{1}{2}\lambda^2\hat{x}^2 \quad \lambda=1, \lambda t=0, 2$$

Classical Yang-Mills dynamics on the lattice

- Lattice CYM Hamiltonian in temporal gauge ($A_0=0$) in the lattice unit

$$H = \frac{1}{2} \sum_{x, a, i} \left[E_i^a(x)^2 + B_i^a(x)^2 \right]$$

$$F_{ij}^a(x) = \partial_i A_j^a(x) - \partial_j A_i^a(x) + \sum_{b, c} f^{abc} \boxed{A_i^b(x) A_j^c(x)} = \varepsilon_{ijk} B_k^a(x)$$

Non-linear & coupling

- Non-compact (A, E) form !

- Demerit: Gauge invariance is not fully satisfied at finite lattice spacing.
- Merit: Easy to consider the coherent state, and conformality is manifest.

- Initial conditions ($E_i^a(x)=0$ is assumed here.)

- Random initial condition: $A_i^a(x) = \text{random in } [-\eta, \eta]$,

- Modulated init. cond.: $A_i^a(\vec{r}) = \delta_{i2} \left[\epsilon_1 \sin\left(\frac{2x\pi}{N_x}\right) + \epsilon_2 \sin\left(\frac{2nx\pi}{N_x}\right) \sin\left(\frac{2mz\pi}{N_z}\right) \right]$

- Constant-A init. cond. $A_i^a(x) = \sqrt{B/g} (\delta_{i2} \delta^{a3} + \delta_{i3} \delta^{a2})$ *Berges et al.('12)*

magnetic field ~ z direction ($\epsilon_1 \gg \epsilon_2$), w and w/o fluc.

How to obtain Lyapunov exponents

■ EOM of $\delta X \rightarrow$ Integral

$$\dot{X} = \begin{pmatrix} \dot{x} \\ \dot{p} \end{pmatrix} = \begin{pmatrix} H_p \\ -H_x \end{pmatrix} \rightarrow \delta \dot{X} = \begin{pmatrix} H_{px} & H_{pp} \\ -H_{xx} & -H_{xp} \end{pmatrix} \delta X \equiv \tilde{H} \delta X$$

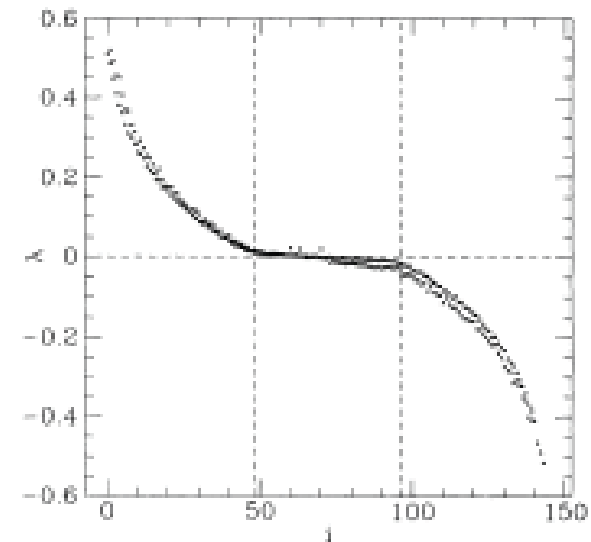
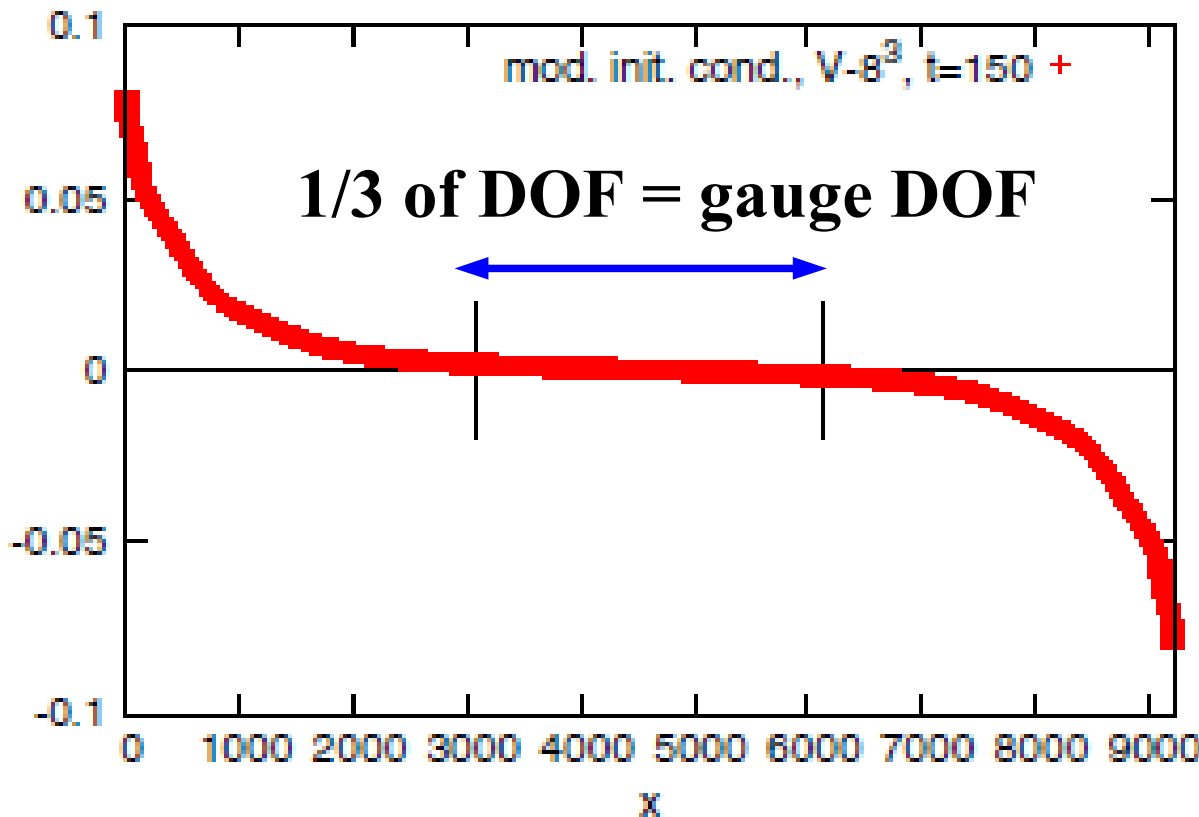
$(H_{px} \equiv \partial^2 H / \partial p \partial x \text{ etc})$

$$\delta X(t) = T \exp \left(\int_0^t dt' \tilde{H}(t') \right) \delta X(t=0) \simeq T \prod_{k=1, N} (1 + \tilde{H} \Delta t) \delta X(t=0)$$
$$= U(0, t) \delta X(t=0)$$

- Diagonalizing U and the eigen value becomes λt .
- Matrix size = $3 \text{ (xyz)} \times (N_c^2 - 1) \times L^3 \times 2 \text{ (A,E)}$

Typical Lyapunov spectrum

- Sum of all Lyapunov exponent = 0 (Liouville theorem)
- 1/3 Positive, 1/3 negative, and 1/3 zero (or pure imag.)



cf) Lyapunov spectrum ($V=2^3$)
Gong, Phys.Rev.D49, 2642 (1994).

Iida, Kunihiro, B.Müller, AO, Schäfer, Takahashi ('13)

KS entropy in CYM from random init. cond.

Kunihiro, Müller, AO, Schäfer, Takahashi, Yamamoto ('10)

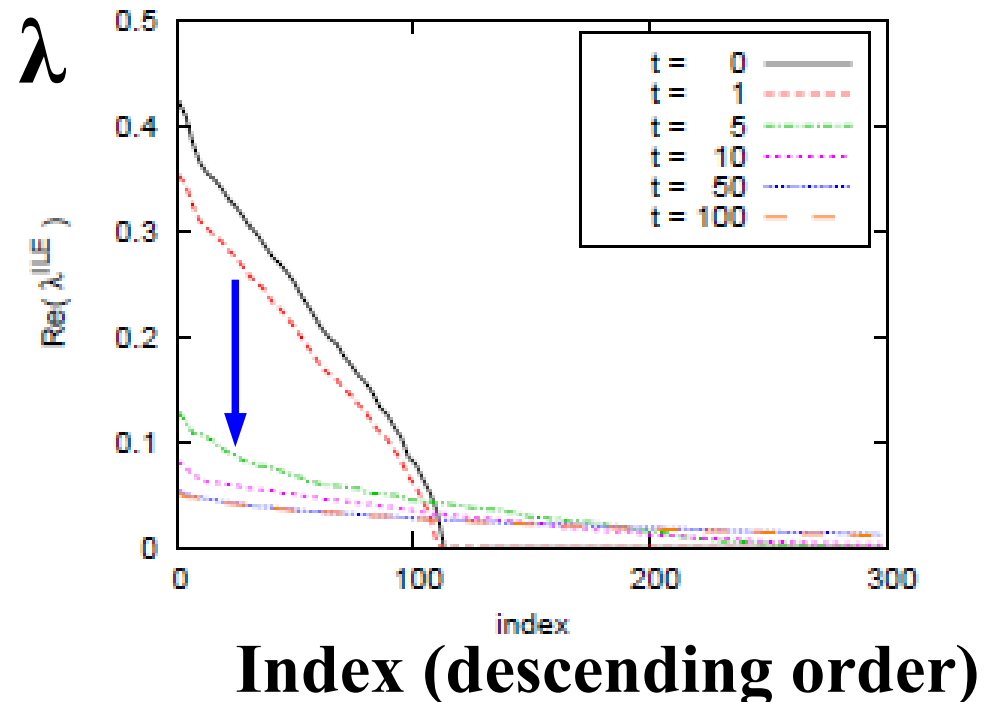
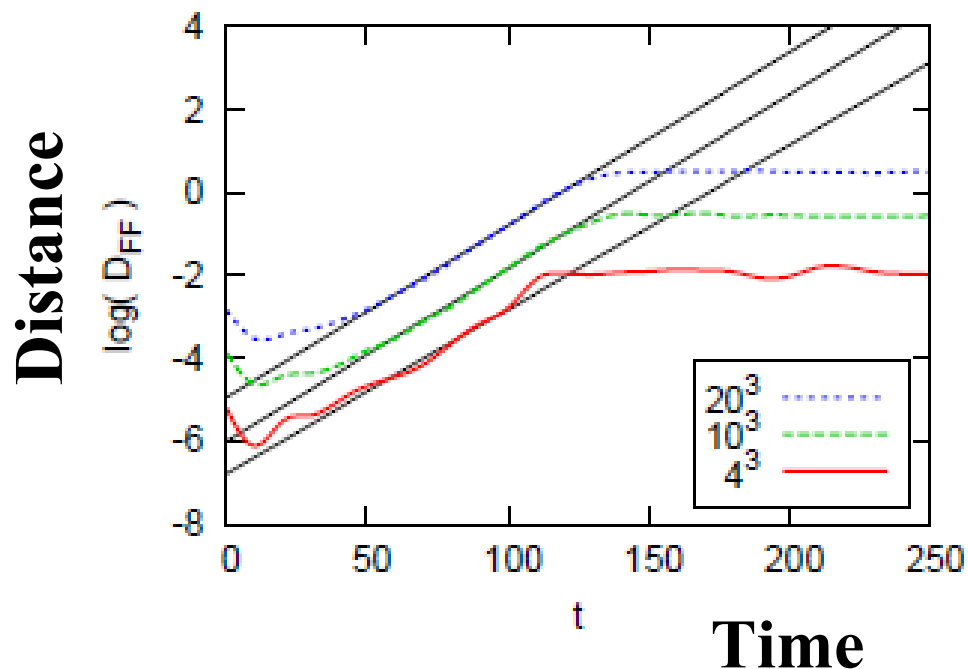
■ Evolution of distance from adjacent init. cond.

- Exponential growth of distance → Instability or Chaoticity

$$D_{FF} = \sqrt{\sum_x \left\{ \sum_{a,i,j} F_{ij}^a(x)^2 - \sum_{a,i,j} F'^a_{ij}(x)^2 \right\}^2}$$

■ Lyapunov exponent distribution

- Rapid spread of positive LEs and macroscopic # of positive LEs → Chaoticity



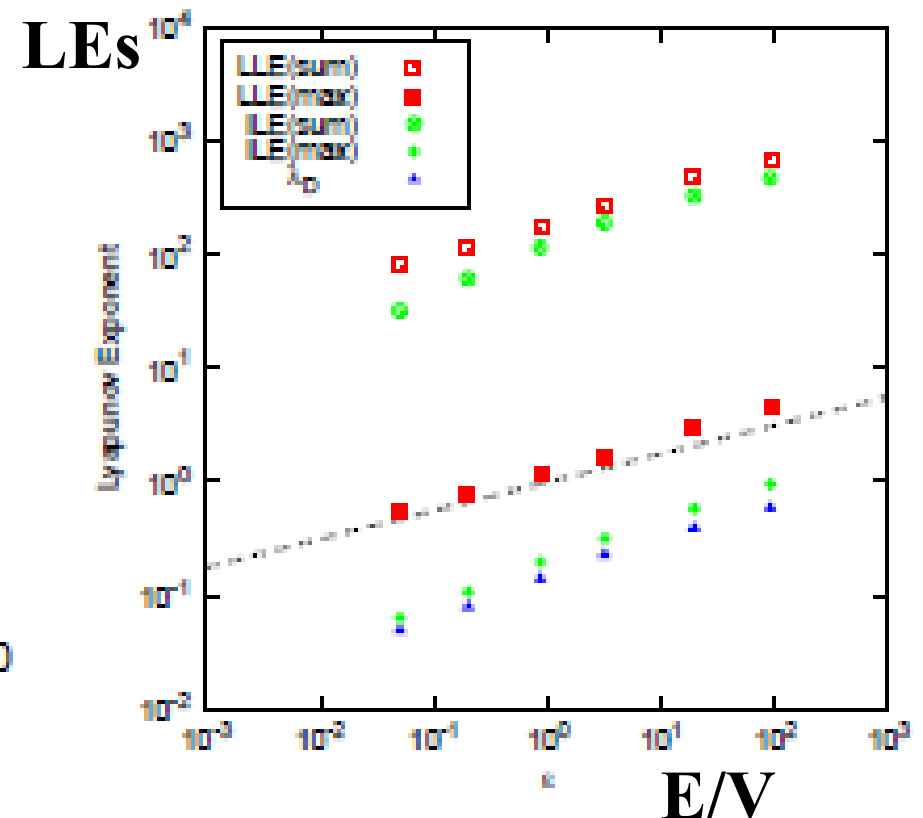
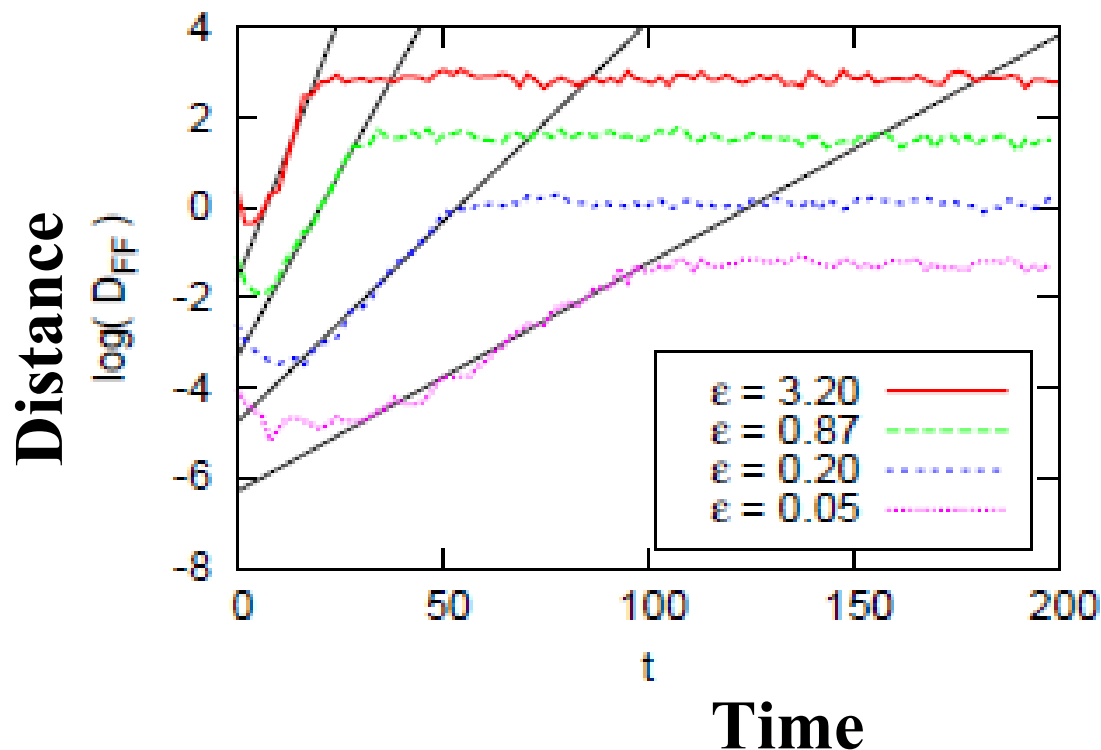
KS entropy in CYM from random init. cond.

Kunihiro, Müller, AO, Schäfer, Takahashi, Yamamoto ('10)

■ Energy density dependence

- Larger energy \rightarrow Larger Lyapunov exp.
- CYM is conformal $\rightarrow \lambda \propto \varepsilon^{1/4}$ ($\varepsilon = E/V$)

$$\frac{dS}{dt} = S_{KS} \sim c_{KS} \varepsilon^{1/4}, \quad c_{KS} \sim 2 \quad (\text{Lattice unit})$$

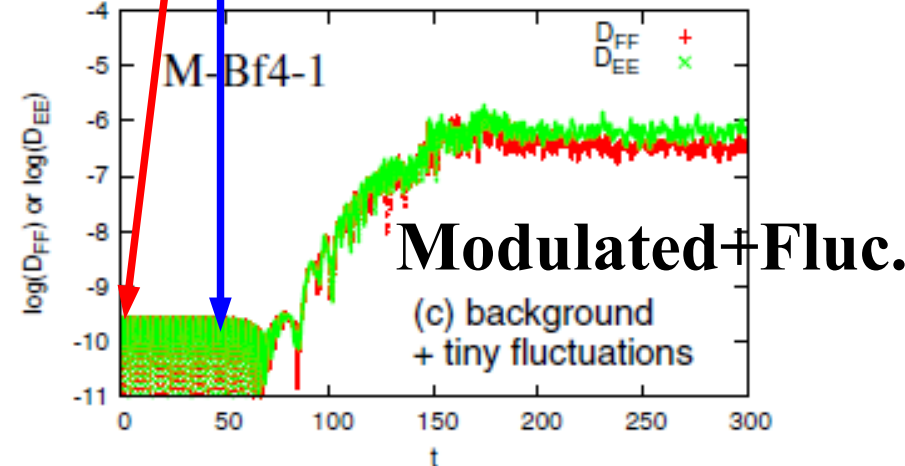
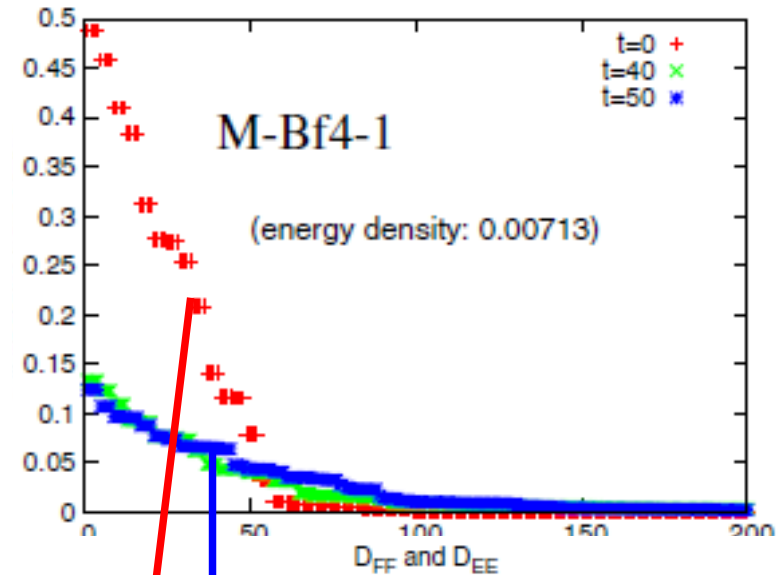
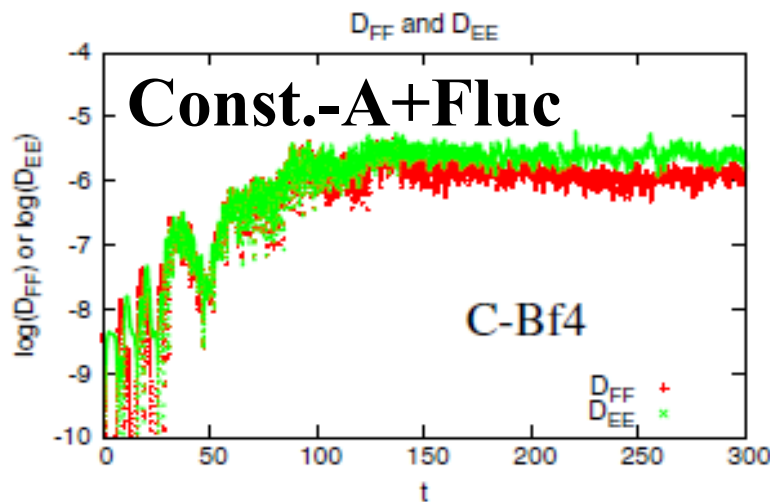


KS entropy in CYM from glasma-like init. cond.

■ Instability under strong color-magnetic field

Nielsen, Olesen ('78), Fujii, Itakura ('08), Berges, Scheffler, Schlichting, Sexty ('12)

- No chaotic behavior is observed with sine waves and constant-A w/o fluctuations.
- Small fluctuations activate instability and chaoticity.
- Chaoticity emerges after instability spreads to many modes.



*CYM as a coherent state
and
Decoherence entropy*

Decoherence Entropy

Muller, Schafer ('03,'06), Fries, Muller, Schafer ('09), Iida, Kunihiro, AO, Takahashi ('14)

■ Coherent State

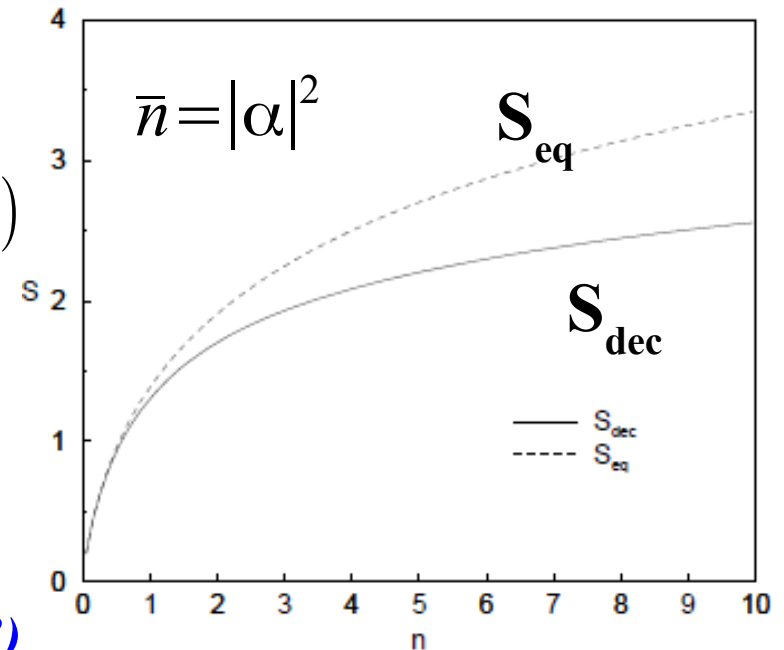
$$\hat{a} |\alpha\rangle = \alpha |\alpha\rangle$$

$$|\alpha\rangle = N \exp(\alpha \hat{a}^+) |0\rangle = \exp(-|\alpha|^2/2) \sum_n \frac{\alpha^n}{\sqrt{n!}} |n\rangle$$

- n-quanta states are coherently superposed in a coherent state.
- When this coherence is broken, entropy is generated (decoherence entropy)

$$P_n = \frac{|\alpha|^{2n}}{n!} \exp(-|\alpha|^2) \quad (\text{Poisson dist.})$$

$$\rightarrow S_{\text{dec}} = - \sum_{n=0}^{\infty} P_n \log P_n > 0$$



Muller, Schafer ('03)

CYM as a Coherent State

Muller, Schafer ('03,'06), Fries, Muller, Schafer ('09), Iida, Kunihiro, AO, Takahashi ('14)

- What kind of state does the CYM correspond to ?
→ Natural guess = Coherent State

$$| \text{CYM} \rangle \simeq \prod_{k, a, i} | \alpha_{k ai} \rangle$$

- Decoherence entropy from CYM

$$S_{\text{dec}} = - \sum_{k, a, i} \sum_n P_n(\alpha_{k ai}) \log P_n(\alpha_{k ai})$$

$$\alpha_{k ai} = \frac{1}{\sqrt{2} \omega_k} \left[\omega_k A_{ai}(\mathbf{k}, t) + i E_{ai}(\mathbf{k}, t) \right], \quad \omega_k = \sqrt{\sin^2 k_x + \sin^2 k_y + \sin^2 k_z}$$

- Is the above assignment unique ?

- Coherent state in each “coherent domain” *Fries, Muller, Schafer ('09)*
- Deviation from Poisson dist. with coupled oscillator
Glauber ('66), Gelis, Venugopalan ('06)

Initial Condition and Time Evolution

■ “Glasma-like” init. cond.

- MV model (boost inv.)
+ Longitudinal fluctuations

→ $B_{x,y}$, $E_{x,y}$, B_η , E_η

McLerran, Venugopalan ('94), Romatschke, Venugopalan ('06), Fukushima, Gelis ('12)

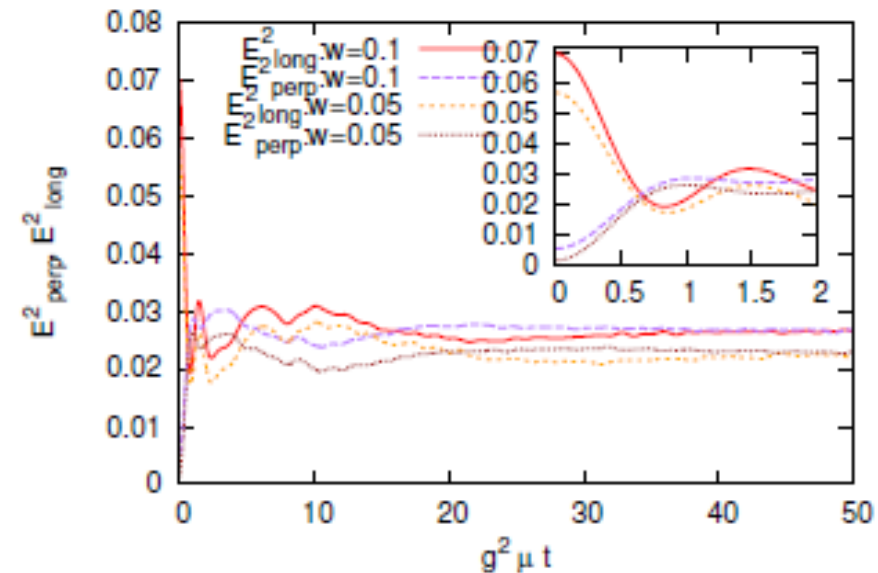
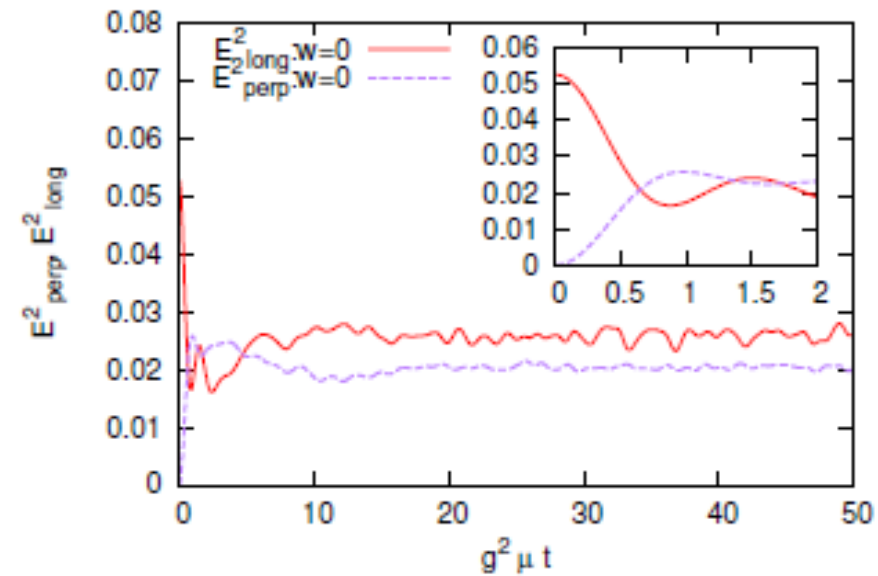
- Non-expanding geometry is assumed,
Substitute B_η and E_η in MV model
into B_z and E_z at $t=0$.

■ Time-evolution

- Short time behavior of E^2 does not depend on the fluctuation strength.
(and similar to expanding geo. results.)
E.g. Lappi, McLerran ('06)

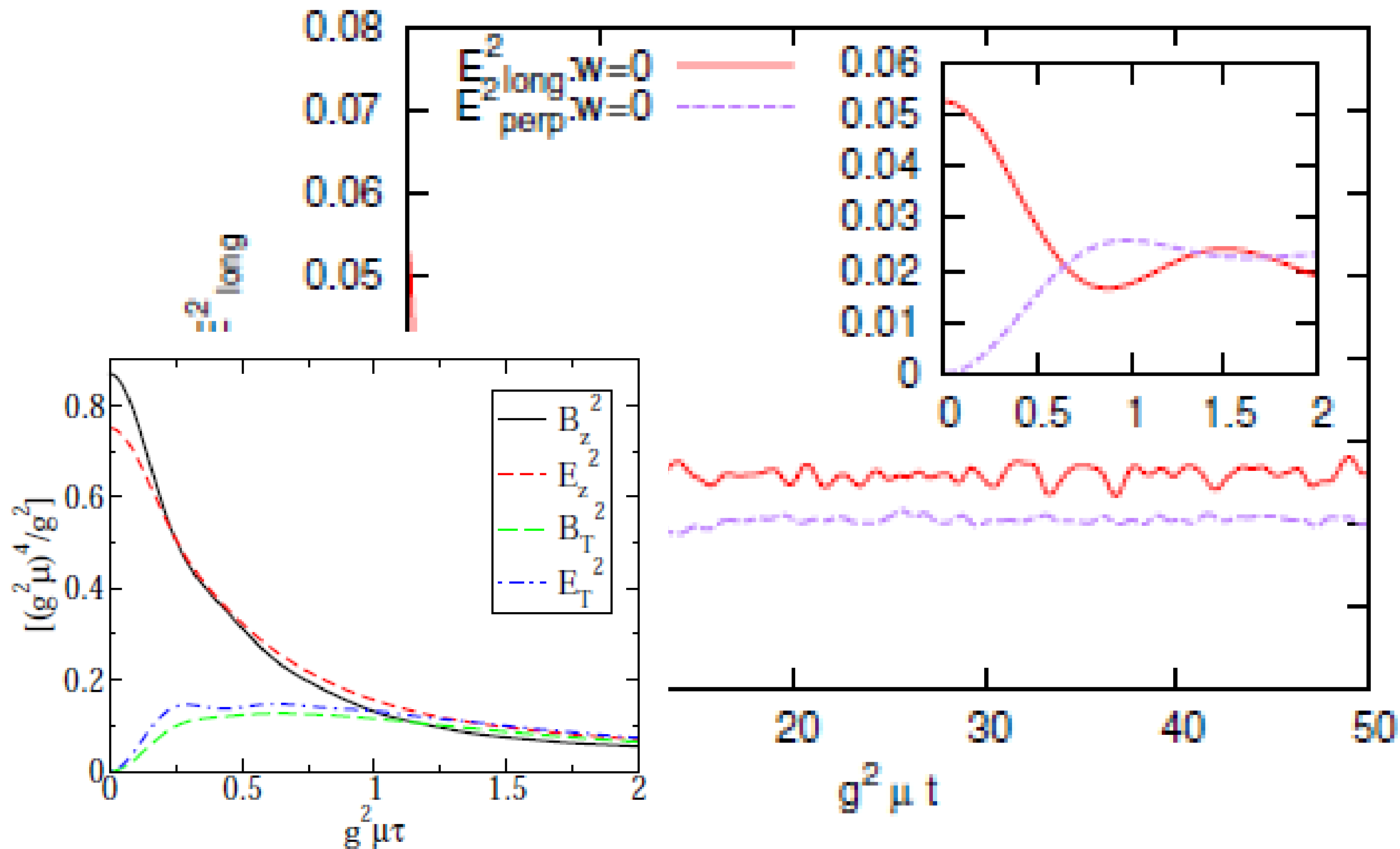
- Long-time behavior:
Earlier “isotropization” in perp. and long. directions of E^2 .

20³ lattice



Iida, Kunihiro, AO, Takahashi ('14)

Initial Condition and Time Evolution



Lappi, McLerran ('06)

Iida, Kunihiro, AO, Takahashi ('14)

Physical Scale

- Rough estimate: $L^2 \sim \pi R_{\text{Au}}^2 \rightarrow 1/a = g^2 \mu \sim 0.32 \text{ GeV}$ ($a \sim 0.63 \text{ fm}$)
- Thermal energy estimate *Kunihiro et al.('10), Muller, Schafer('11)*

$$\varepsilon_{\text{CYM}} = 2(N_c^2 - 1) \frac{T}{a^3}, \quad \varepsilon_{\text{SB}} = 2(N_c^2 - 1) \frac{\pi^2}{30} T^4$$

- **CYM energy density should not exceed Stefan-Boltzmann energy density in equilibrium.**

$$\varepsilon_{\text{CYM}} < \varepsilon_{\text{SB}} \\ \rightarrow a > \frac{1}{T} \left(\frac{30}{\pi^2} \right)^{1/3} \sim \frac{1.4}{T}$$

This estimate gives

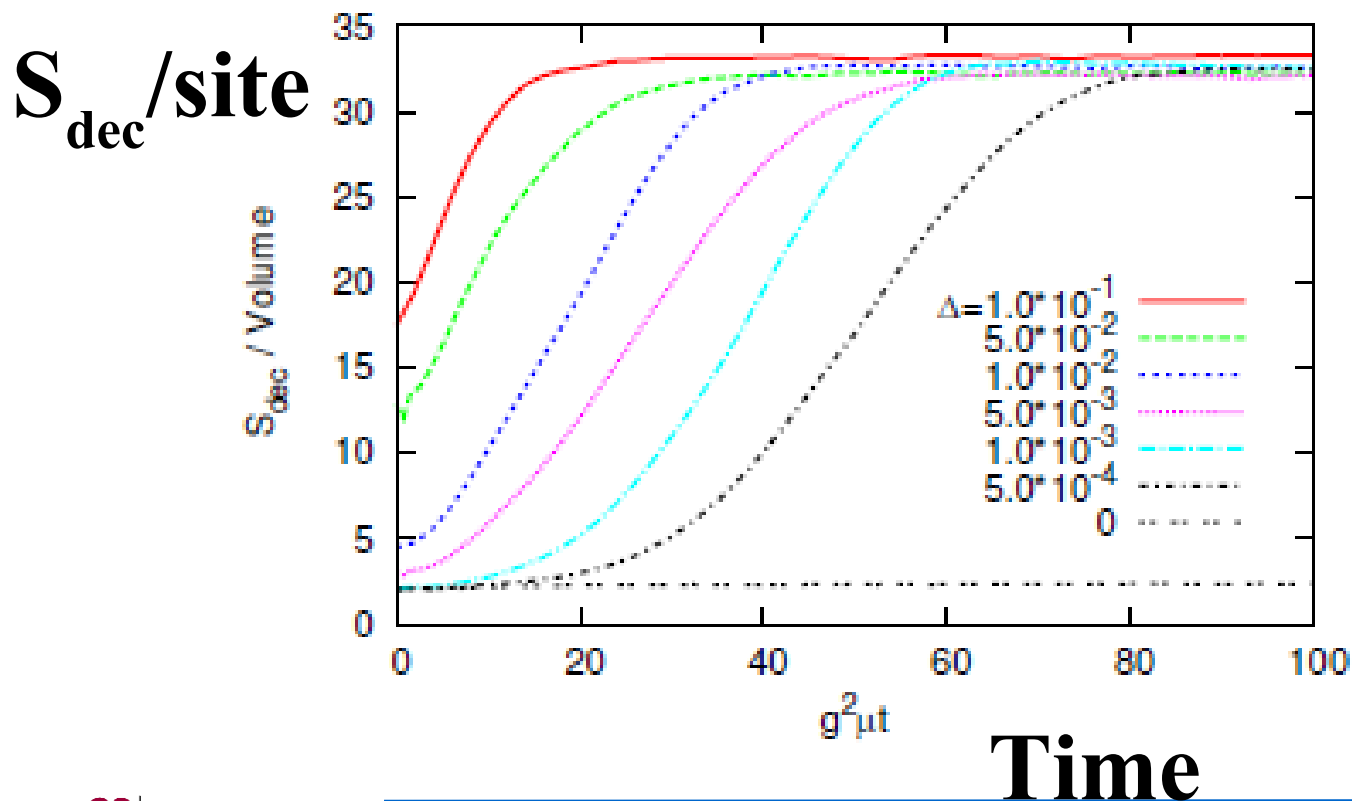
$a > 0.8 \text{ fm}$ ($T=350 \text{ MeV}$),

$a > 0.6 \text{ fm}$ ($T=500 \text{ MeV}$).

Decoherence Entropy of CYM

■ How about the decoherence entropy ?

- $\langle \delta E^2 \rangle / \langle E^2 \rangle \sim 0.1$ ($\Delta=0.05$) and 0.3 ($\Delta=0.1$)
- $S_{\text{dec}} \sim 2.3$ ($\Delta=0$) and 33 ($\Delta=0.05, 0.1$)
- Entropy from initial state fluc. and chaoticity
- No long. fluc. results in 2D (pz=0 mode) entropy, while 3D entropy is realized with finite long. fluc. (non-zero Δ).



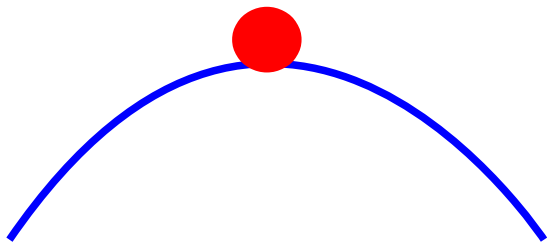
Decoherence Entropy Production Rate

Decoherence entropy growth rate should be compared with KS entropy

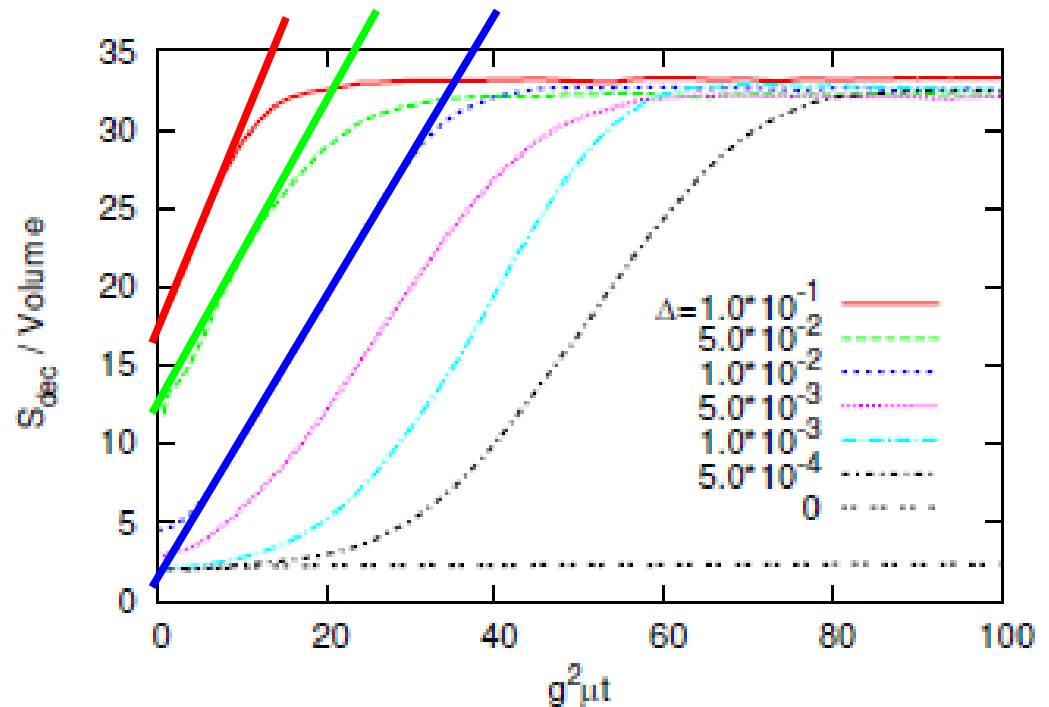
- $dS_{\text{dec}}/dt \sim 0.88$ ($\Delta=0.01$), 1.05 ($\Delta=0.05$), 1.36 ($\Delta=0.1$)
- KS entropy estimate: $S_{\text{KS}} \sim c_{\text{KS}} \varepsilon^{1/4}$, $c_{\text{KS}} \sim 2$ (conformal chaotic value)
- Energy density: $\varepsilon = 0.17$ ($\Delta=0.01$), 0.18 ($\Delta=0.05$), 0.21 ($\Delta=0.1$)
 $\rightarrow c_{\text{KS}} = dS^{\text{dec}}/dt/\varepsilon^{1/4} = 1.4$ ($\Delta=0.01$), 1.6 ($\Delta=0.05$), 2.0 ($\Delta=0.1$)

$$\frac{1}{S_{\text{KS}}} \frac{dS_{\text{dec}}}{dt} \sim (0.7 - 1.0)$$

- KS entropy
= Potentially realized
growth rate



$\Delta=0$: unstable
but stationary



Physical Scale

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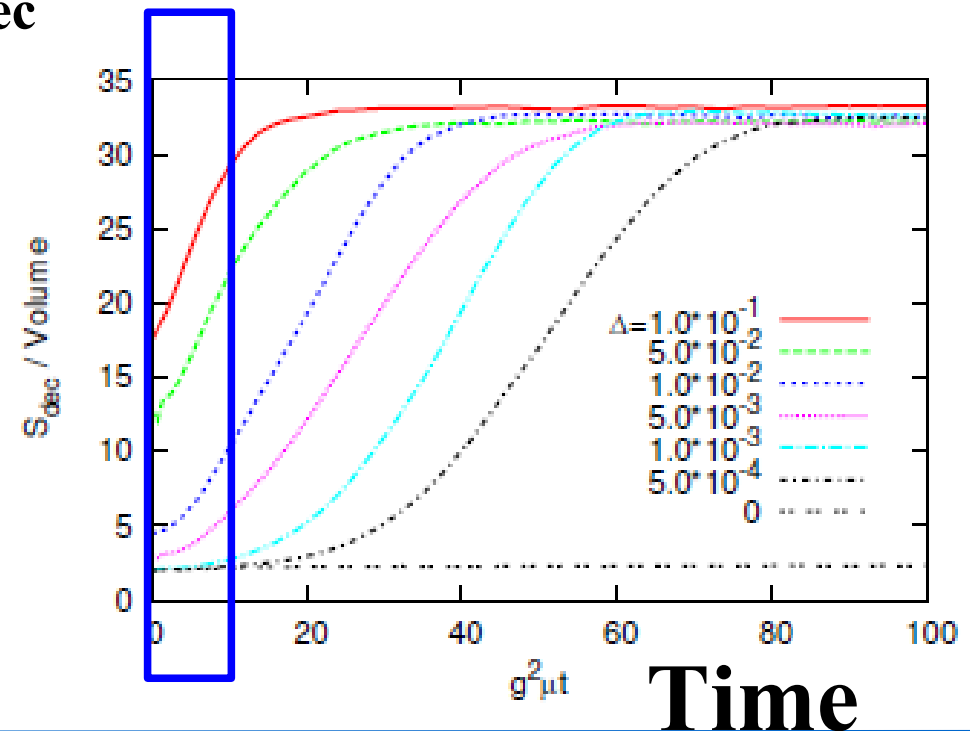
This estimate gives

$a > 0.8 \text{ fm}$ ($T=350 \text{ MeV}$),

$a > 0.6 \text{ fm}$ ($T=500 \text{ MeV}$).

- First 6 fm/c corresponds to lattice time $g^2\mu t < 10$!

$S_{\text{dec}}/\text{site}$ relevant to HIC



Summary

- We have evaluated the entropy from classical Yang-Mills field using
 - Kolmogorov-Sinai entropy (as a growth rate),
 - Decoherence entropy.
- Entropy could be produced even before classical Yang-Mills field decays into particles.

- Suggested scenario: Fluctuation
→ Realization of instability & Spread to many modes → Chaoticity
- Rough estimate of entropy production rate in non-expanding CYM

$$\frac{dS}{dt} = S_{KS} \sim c_{KS} \epsilon^{1/4}, \quad c_{KS} \sim 2 \quad (\text{Lattice unit})$$

- Decoherence entropy grows at about the $c_{KS} = (1-2)$ rate, and saturates, and it is sensitive to longitudinal fluctuations.
- For the initial stage entropy, both the time-evolution and the initial entropy value would be important.

Future works

- **Decoherence entropy in expanding glasma with realistic fluctuation strength**
E.g. Epelbaum, Gelis ('13)
- **Estimate of decoherence time during CYM evolution and Defining coherent domain in CYM**
c.f. Fries, Muller, Schafer ('09)
- **Coupling to particle DOF**
 - **Fluctuation in classical statistical simulation ~ particles ?**
 - **CYM + gluon test particles**
Dumitru, Nara, Strickland ('07)
 - **2PI formalism of CYM and gluon propagator**
c.f. Nishiyama, AO ('10, w/o CYM), Hatta, Nishiyama ('11), Cassing ('09)

Thank you for your attention !