

J-PARC エネルギーにおける集団フロー *(Collective Flows at J-PARC Energies)*

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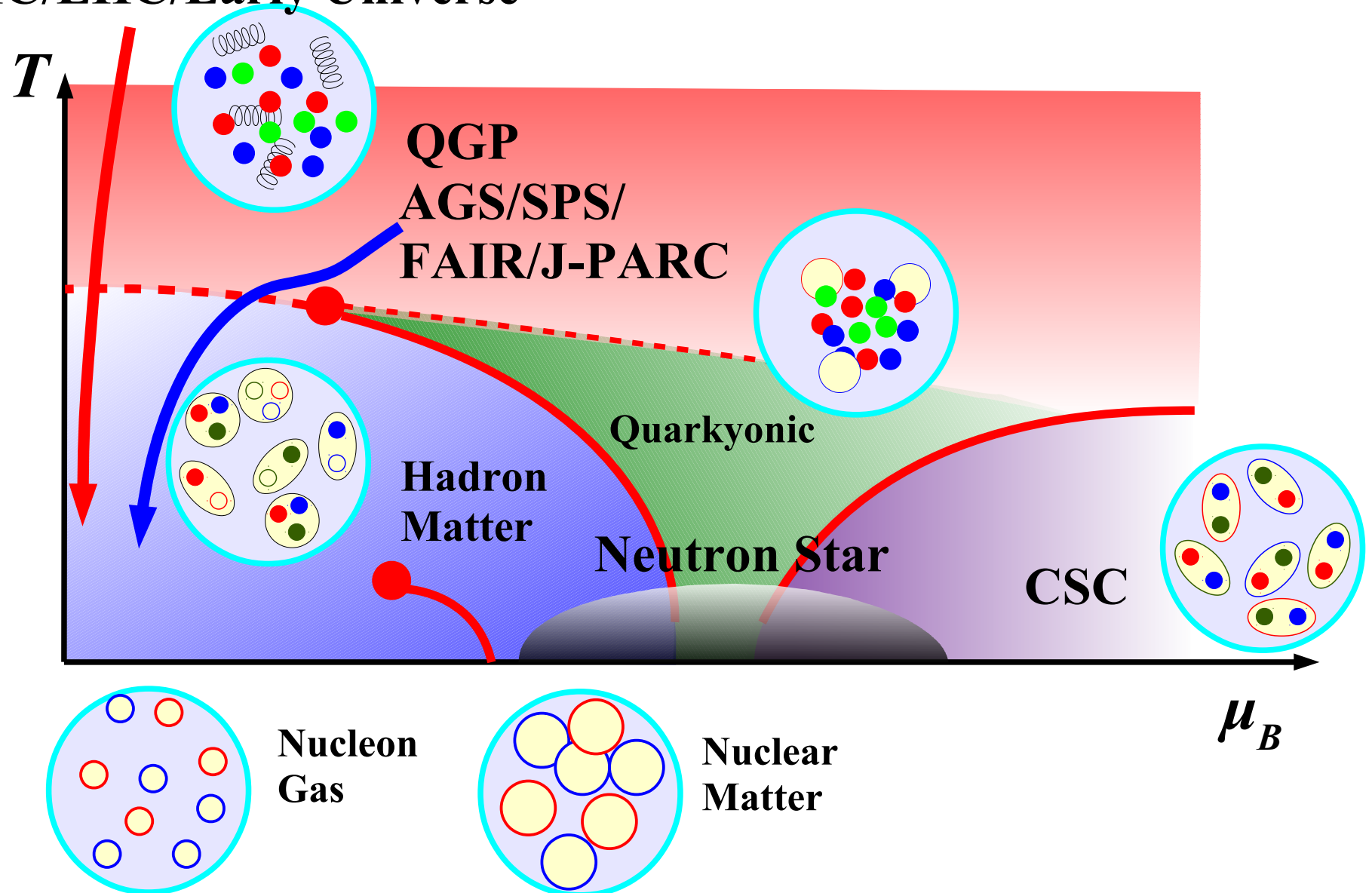
J-PARC における重イオン衝突実験が拓く新しい物理
11/26-27, 2014, KEK

- Introduction
- A Hadronic Transport Model: JAM
- Collective Flows at J-PARC Energy
- Summary



QCD Phase Diagram

RHIC/LHC/Early Universe



Signals of QGP formation & QCD phase transition

■ Signals of QGP formation at $\sqrt{s_{NN}} = 200$ GeV and above (RHIC, LHC)

- Jet quenching in AA collisions (not in dA)
- Large elliptic flow (success of hydrodynamics)
- Quark number scaling (coalescence of quarks)

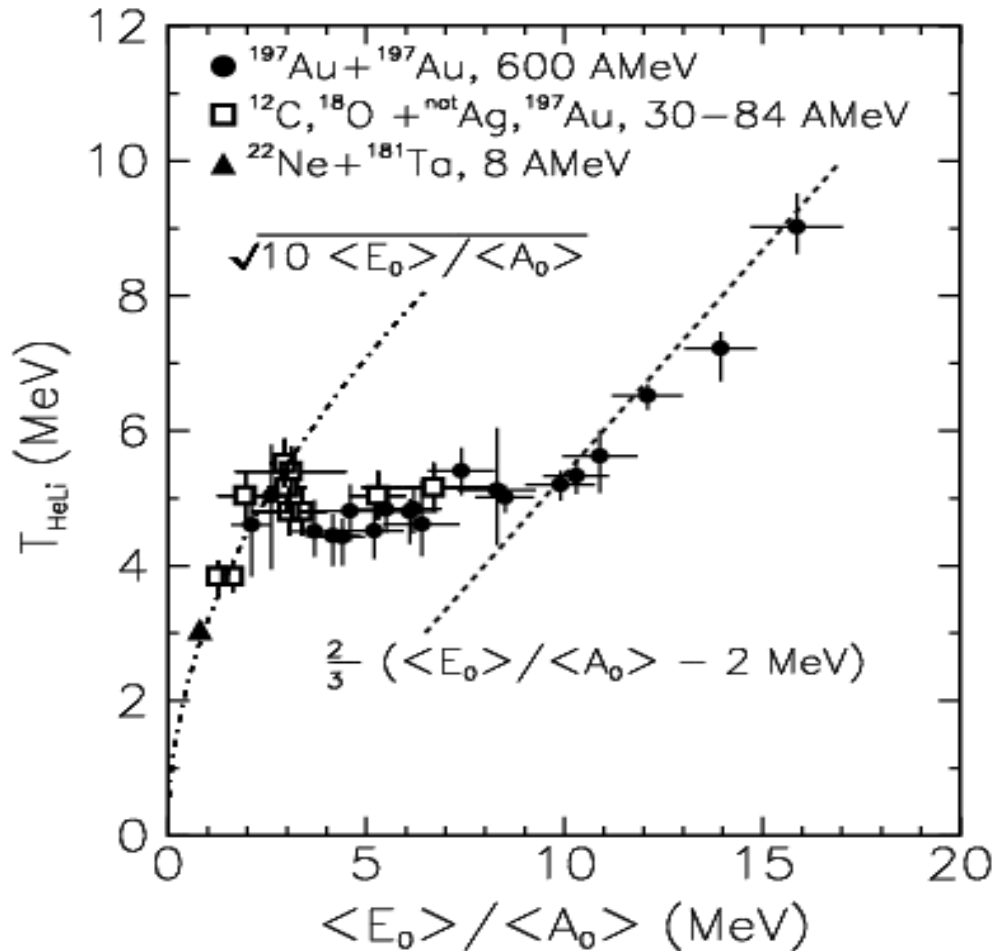
Similar signals are observed at $\sqrt{s_{NN}} > 39$ GeV

■ Signals of QCD phase transition at $\sqrt{s_{NN}} = 4-40$ GeV (?)

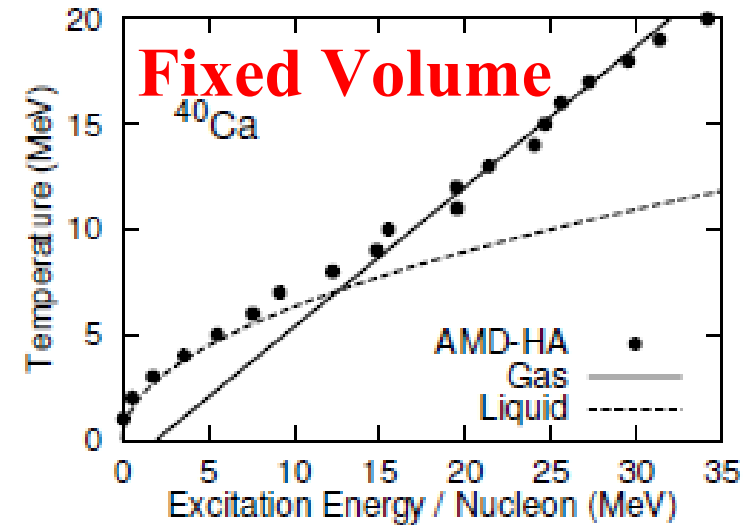
- Hint from Liquid-Gas phase transition: caloric curve
- Horn, Step (Re-Hardening), Dale
- Non-monotonic behavior of proton number moment ($\kappa\sigma^2$) and collective flow (dv_1/dy)

Nuclear Liquid-Gas Phase Transition

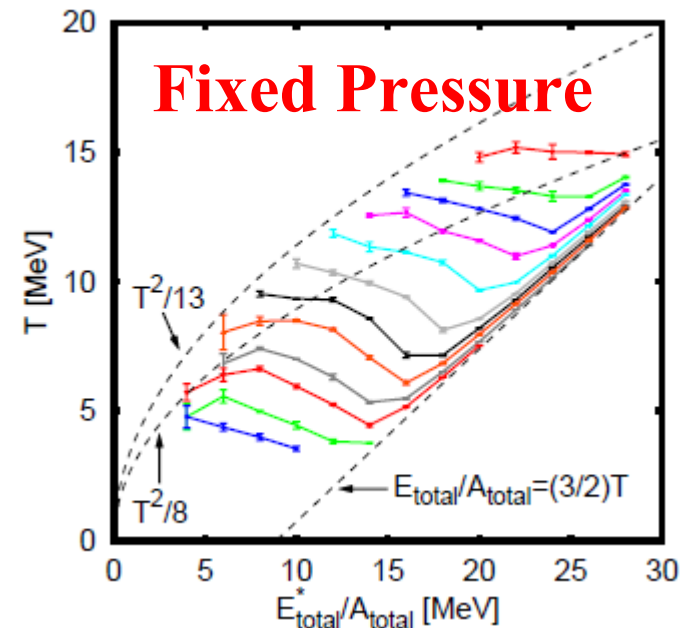
- Caloric curve \rightarrow LG phase transition (Smoking gun)



J. Pochadzalla et al. (GSI-ALLADIN collab.), PRL 75 (1995) 1040.



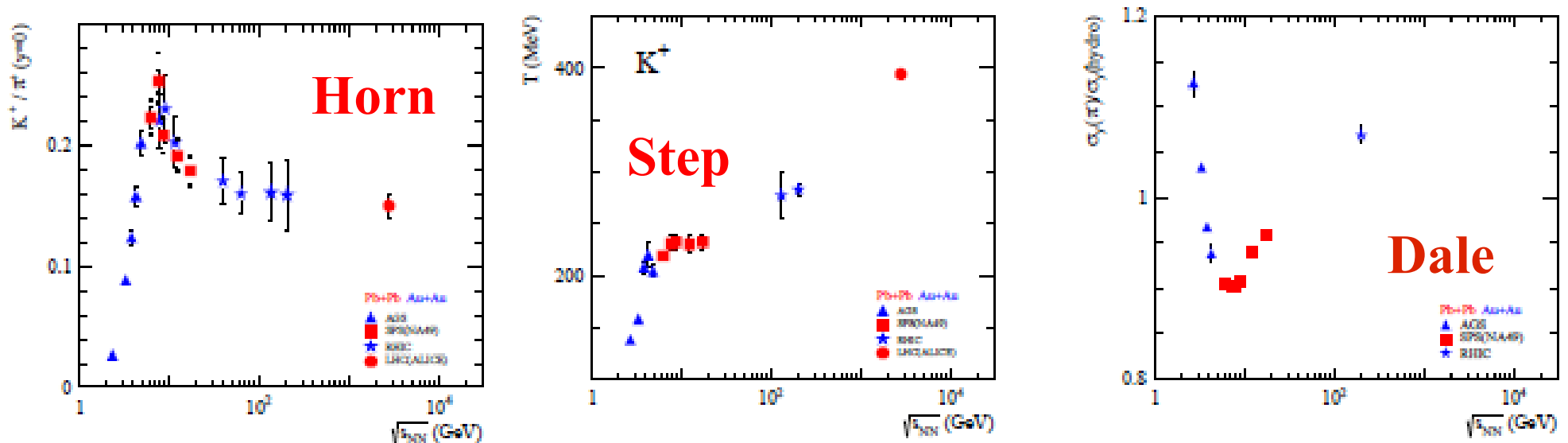
AO, Randrup ('98)



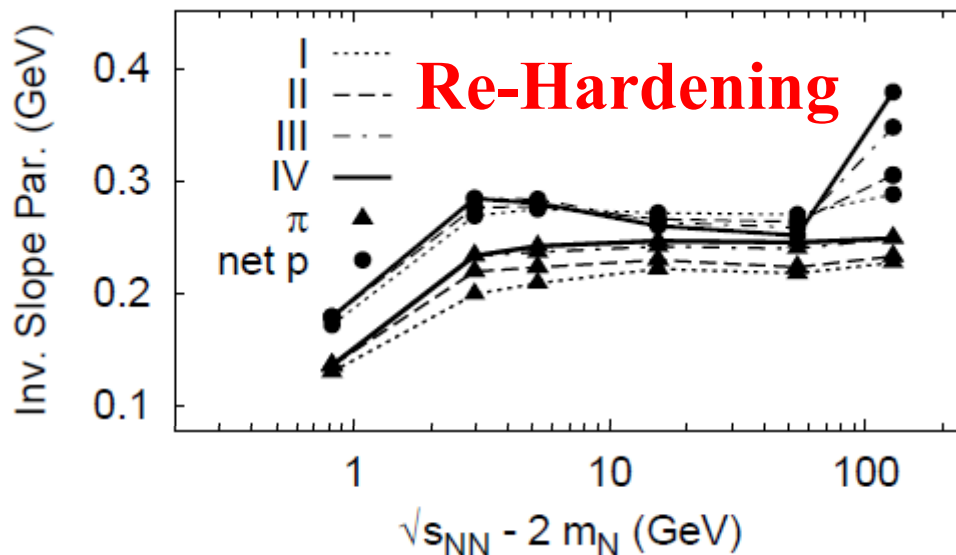
T. Furuta, A. Ono ('09)

Horn, Step and Dale

- Non-monotonic behavior in K^+/π^+ ratio (Horn), m_T slope par. (Step or re-hardening), rapidity dist. width of π (Dale)



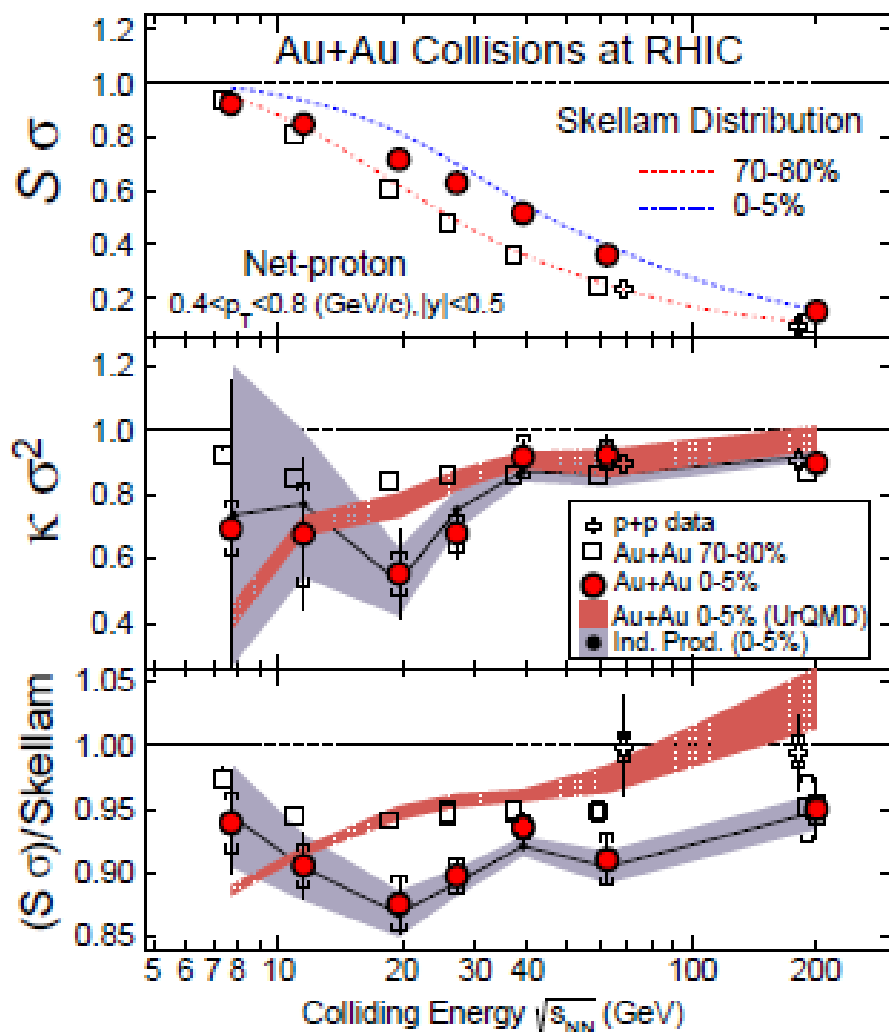
E.g. A. Rustamov (2012)



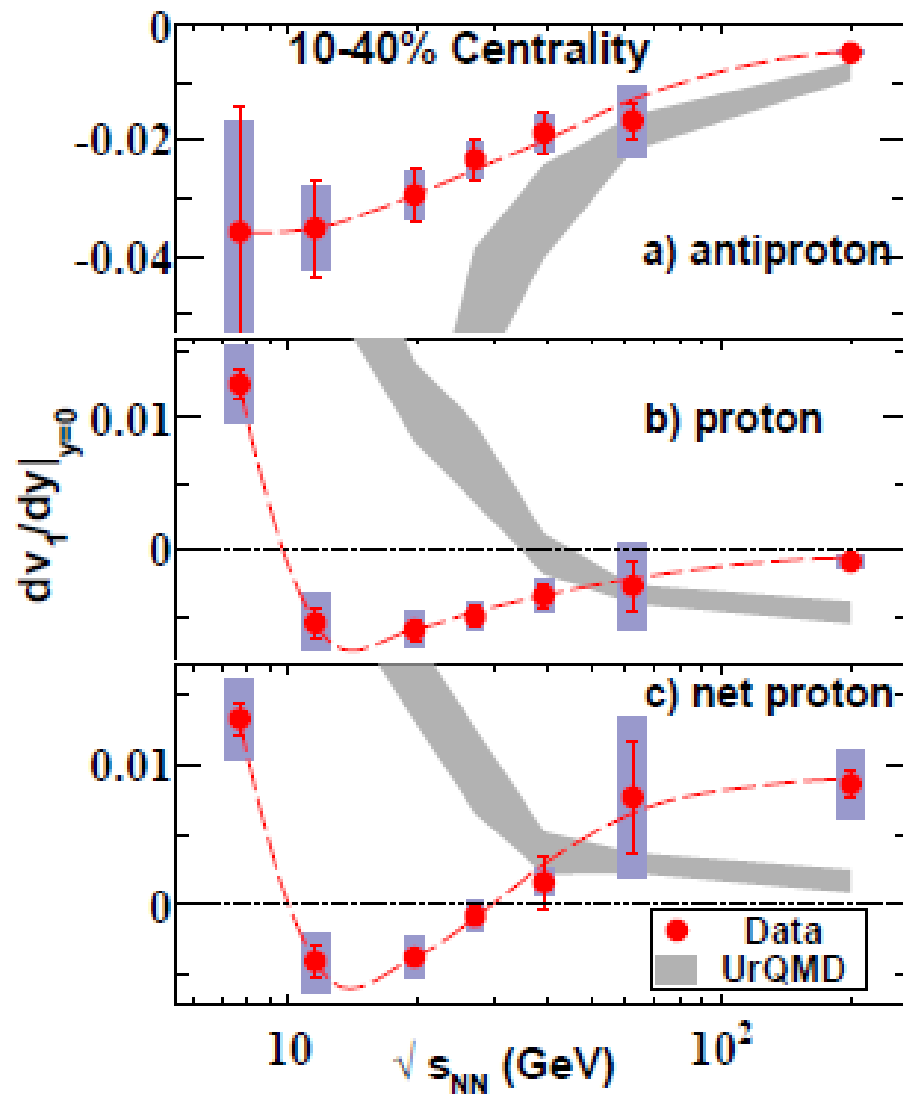
*N. Otuka, P.K.Sahu, M. Isse,
Y. Nara, AO, nucl-th/010205*

Net-Proton Number Moments & Directed Flow

- Non-monotonic behavior of $\kappa\sigma^2$ and dv_1/dy . CP signal ?



STAR Collab. (PRL 112('14)032302)

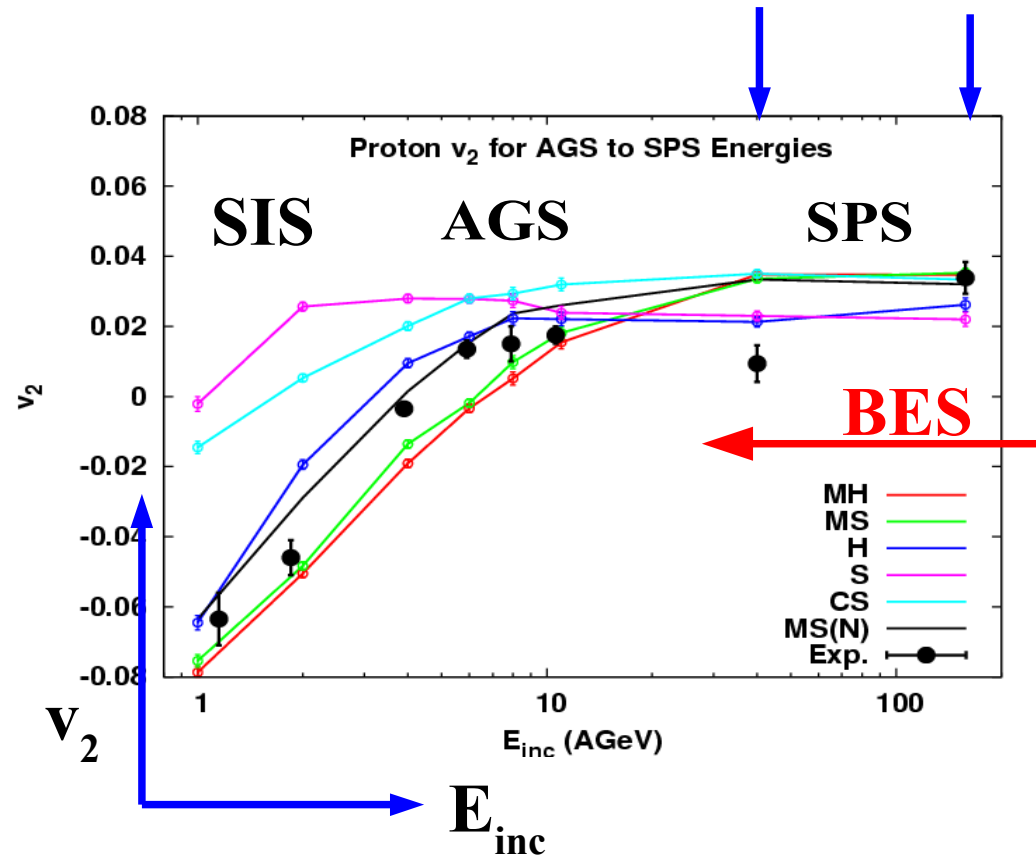
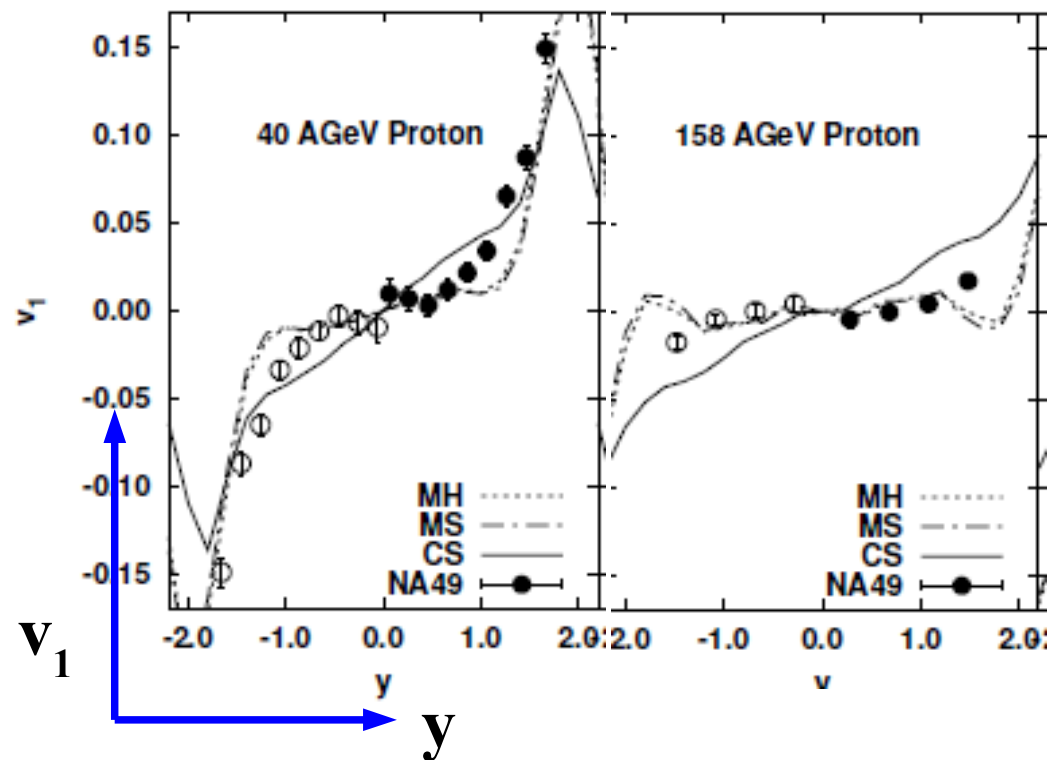


STAR Collab., PRL 112('14)162301.

JAM results at AGS and SPS Energies

- JAM w/ Mean-Field effects roughly explains v_1 and v_2 at AGS & SPS (1-158 A GeV $\rightarrow \sqrt{s_{NN}} = 2.5-20$ GeV)

$$\sqrt{s_{NN}} = 8.9 \text{ GeV} \quad \sqrt{s_{NN}} = 17.3 \text{ GeV}$$

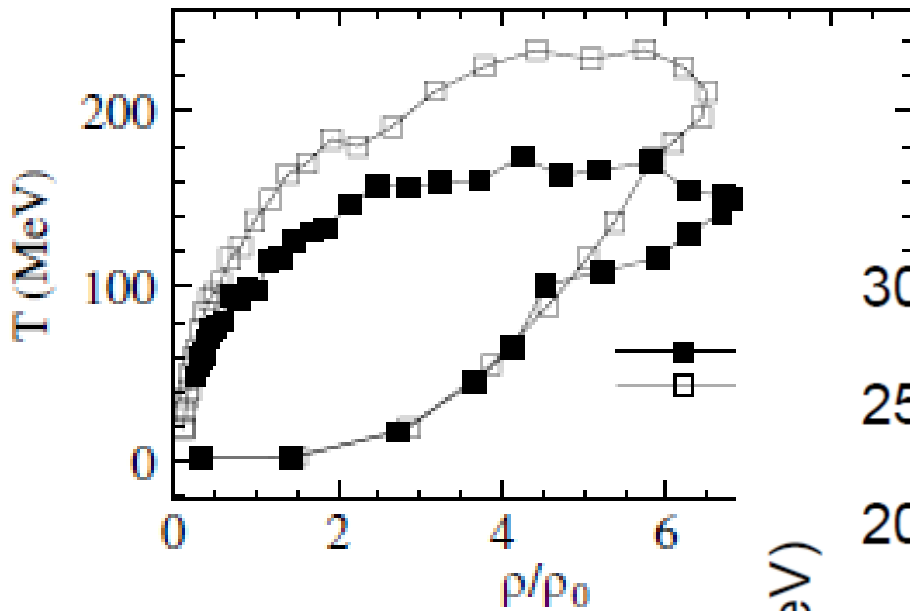


M. Isse, AO, N. Otuka, P. K. Sahu, Y. Nara, PRC72('05)064908

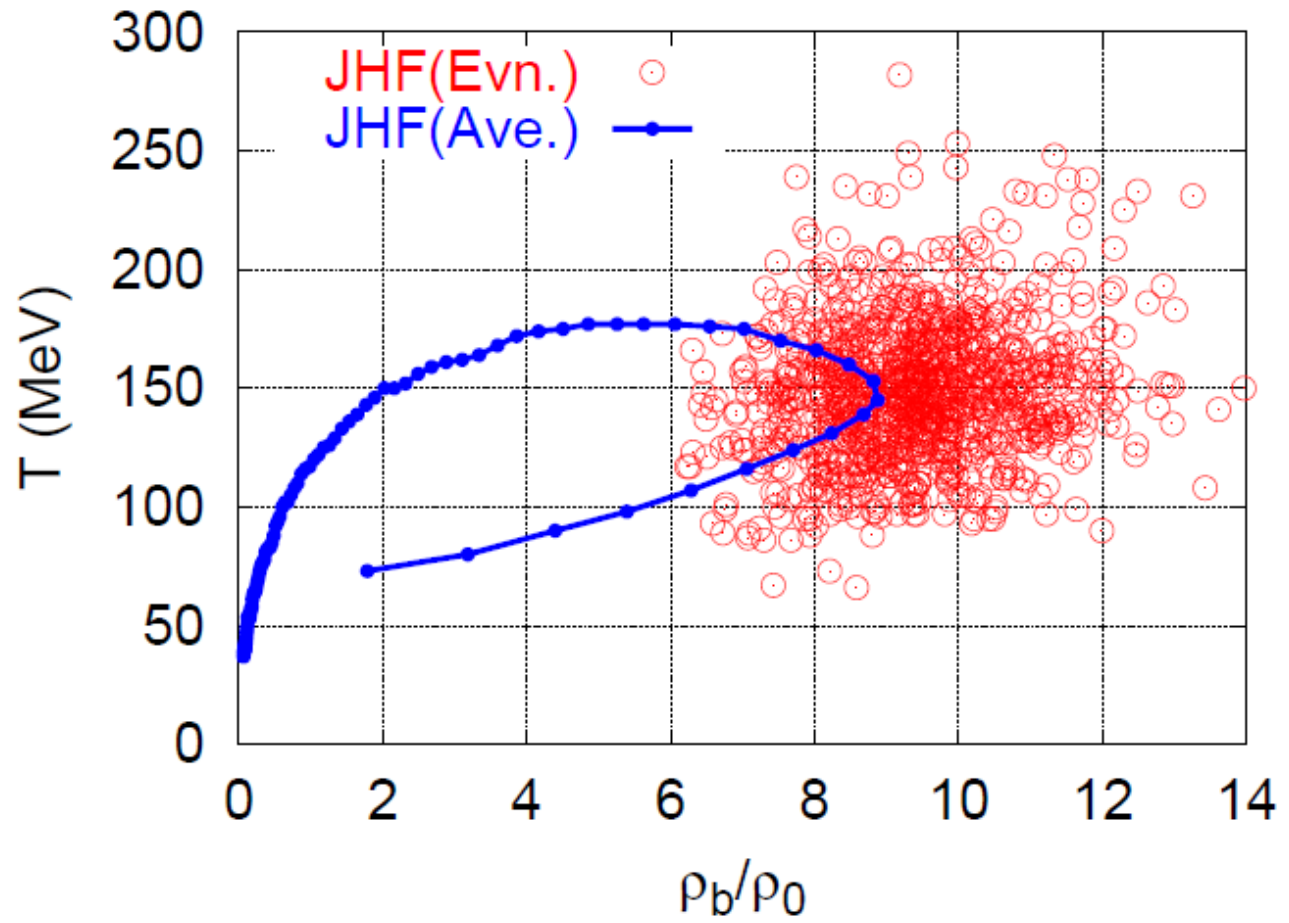
QCD phase transition at J-PARC Energies ?

- **J-PARC Energies:** $\sqrt{s_{NN}} = 4-40 \text{ GeV}$ (or $\sqrt{s_{NN}} = 1.9-6.2 \text{ GeV}$)
 - $E(p)=30 \text{ GeV} \rightarrow E(\text{Au}) \sim 12 \text{ AGeV}$ (full strip, $\sqrt{s_{NN}} = 5.1 \text{ GeV}$ for Au+Au)
 - $E(p)=50 \text{ GeV} \rightarrow E(\text{Au}) \sim 20 \text{ AGeV}$ ($\sqrt{s_{NN}} = 6.4 \text{ GeV}$)
 - $E(p)=30 \text{ GeV}$ (50 GeV) Collider $\rightarrow \sqrt{s_{NN}} = 26 \text{ GeV}$ (42 GeV)
- **Two Aspects of J-PARC energies**
 - Formation of highest baryon density matter
 - Various non-monotonic behaviors \rightarrow Onset of deconfinement

Highest Density Matter at J-PARC ?



Nara, Otuka, AO,
Maruyama ('97)



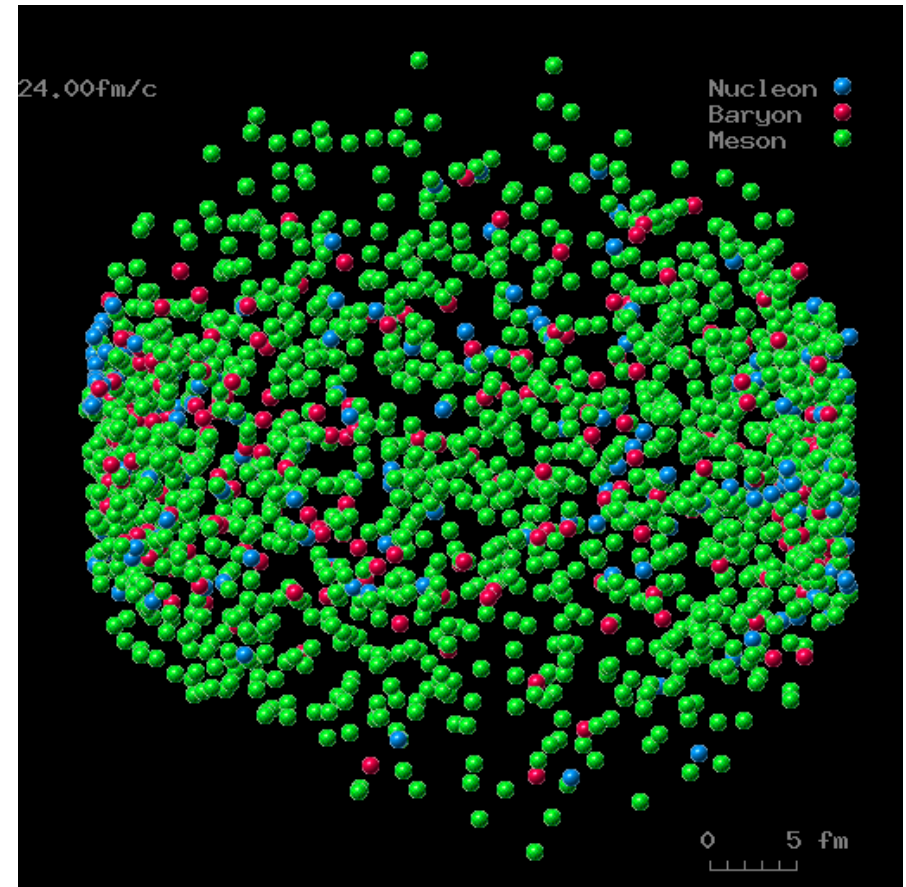
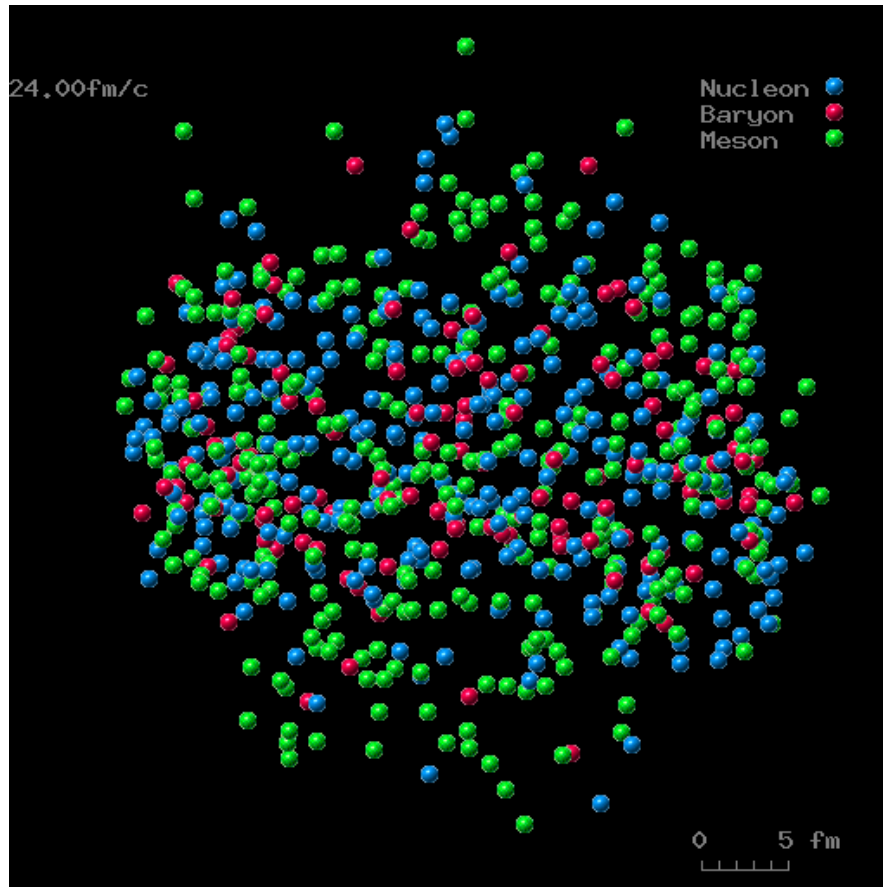
Central 1 fm^3 cube.

大西、JHF workshop (2002)

How do heavy-ion collisions look like ?

Au+Au, 10.6 A GeV

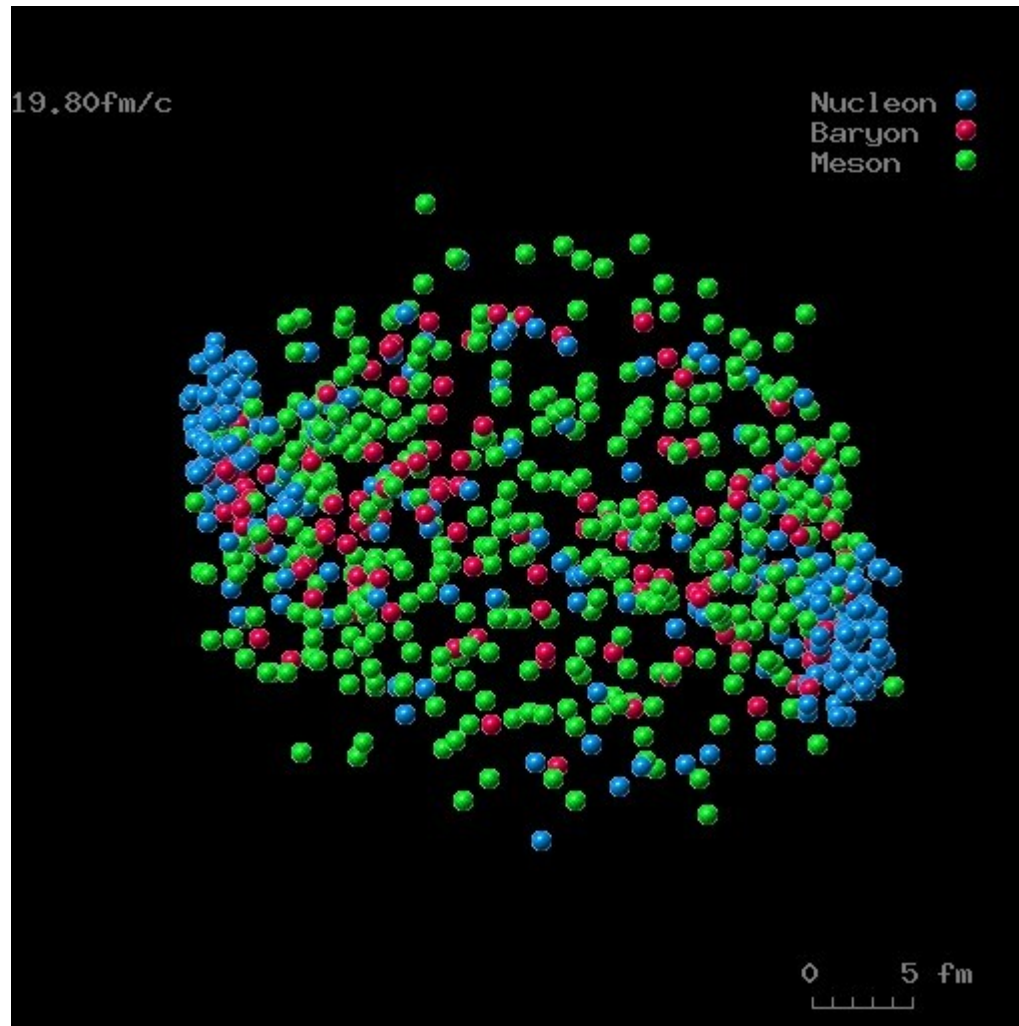
Pb+Pb, 158 A GeV



JAMming on the Web <http://www.jcprg.org/jow/>

Ohnishi @ KEK-HIC workshop, Nov.26-27, 2014

J-PARC energy



Au+Au, 25 AGeV, b=5 fm (JOW)

QCD phase transition at J-PARC Energies ?

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- **Two Aspects of J-PARC energies**
 - Formation of highest baryon density matter
 - Various non-monotonic behaviors \rightarrow Onset of deconfinement

Question

*Do these Non-mono. behaviors signal the onset of QCD phase transition and/or QCD critical point ?
or Do they show some properties of hadronic matter ?
 \rightarrow Let's examine in hadronic transport models !*

A Hadronic Transport Model: JAM

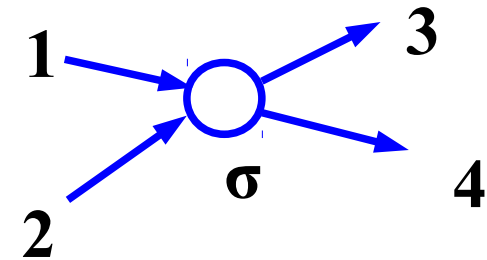
Transport Equation

■ Boltzmann equation with potential effects (BUU Equation)

Bertsch, Das Gupta, Phys. Rept. 160(88), 190

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \nabla f - \nabla U \cdot \nabla_p f = I_{\text{coll}}$$

$$I_{\text{coll}}(\mathbf{r}, \mathbf{p}) = -\frac{1}{2} \int \frac{d\mathbf{p}_2}{(2\pi)^3} d\Omega v_{12} \frac{d\sigma}{d\Omega} [f f_2 (1 - f_3)(1 - f_4)] - (12 \leftrightarrow 34)]$$



(NN elastic scattering case)

■ Inputs of Transport models

● Cross section σ (c.f. Nara's talk)

- ◆ Elementary cross section should be taken from data, if possible.
- ◆ Unknown cross sections need to be given by a model.

● Potential U

- ◆ Hopefully, given by ab initio calc. Usually, given by phen. models.
- ◆ U depends on ρ and p .

JAM (Jet AA Microscopic transport model)

Nara, Otuka, AO, Niita, Chiba, Phys. Rev. C61 (2000), 024901.
Isse, AO, Otuka, Sahu, Nara, Phys.Rev. C 72 (2005), 064908.

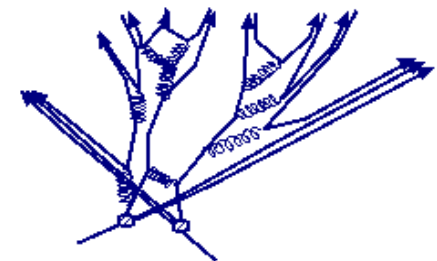
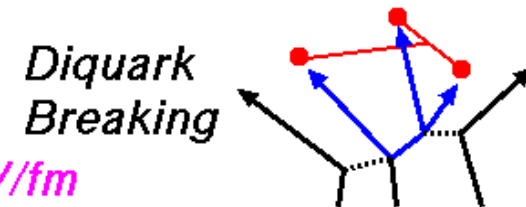
■ Hadron-String Cascade with Jet production

- Hadron Res. up to $m < 2$ GeV
- String & Jet production and decay using Lund string model
T. Sjostrand et al., Comput. Phys. Commun. 135 (2001), 238 (PYTHIA).
- String-Hadron collisions are simulated by hh collisions in the formation time (\sim RQMD) *H. Sorge, PRC52 ('95)3291.*
- Mean field effects are included with density and momentum dependences.
Isse et al., ('05)
- **NO Secondary partonic int.**
- Open to public



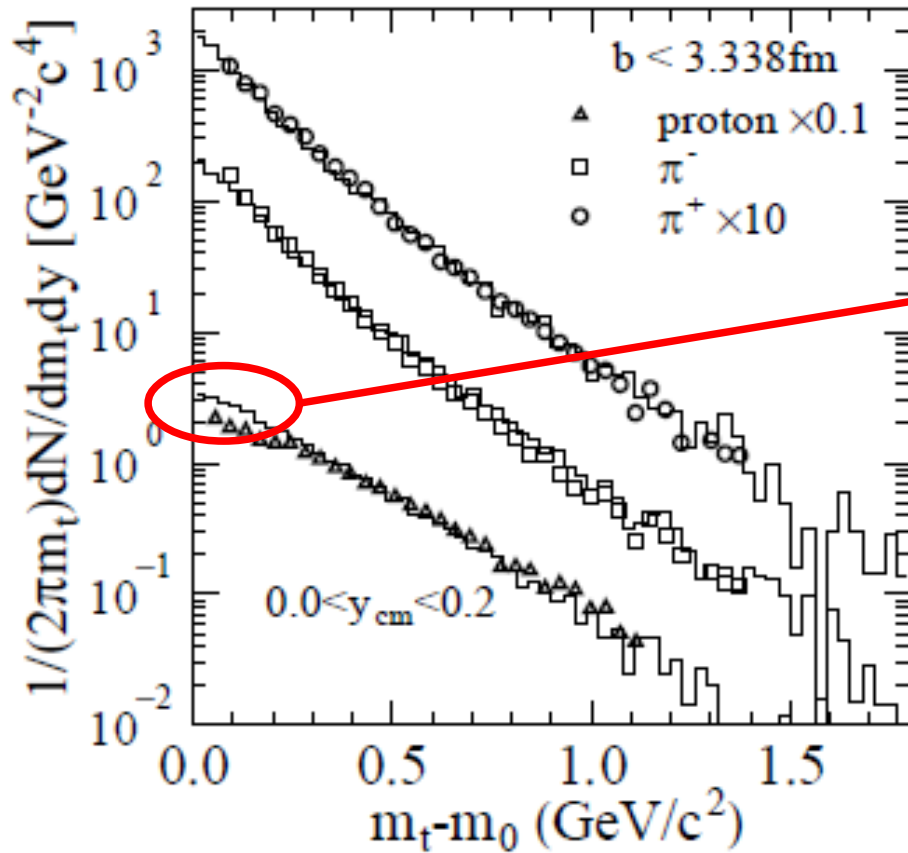
$\tau \sim 1$ fm/c
for $\kappa \sim 1$ GeV/fm

Resonance
+ String
+ Jet

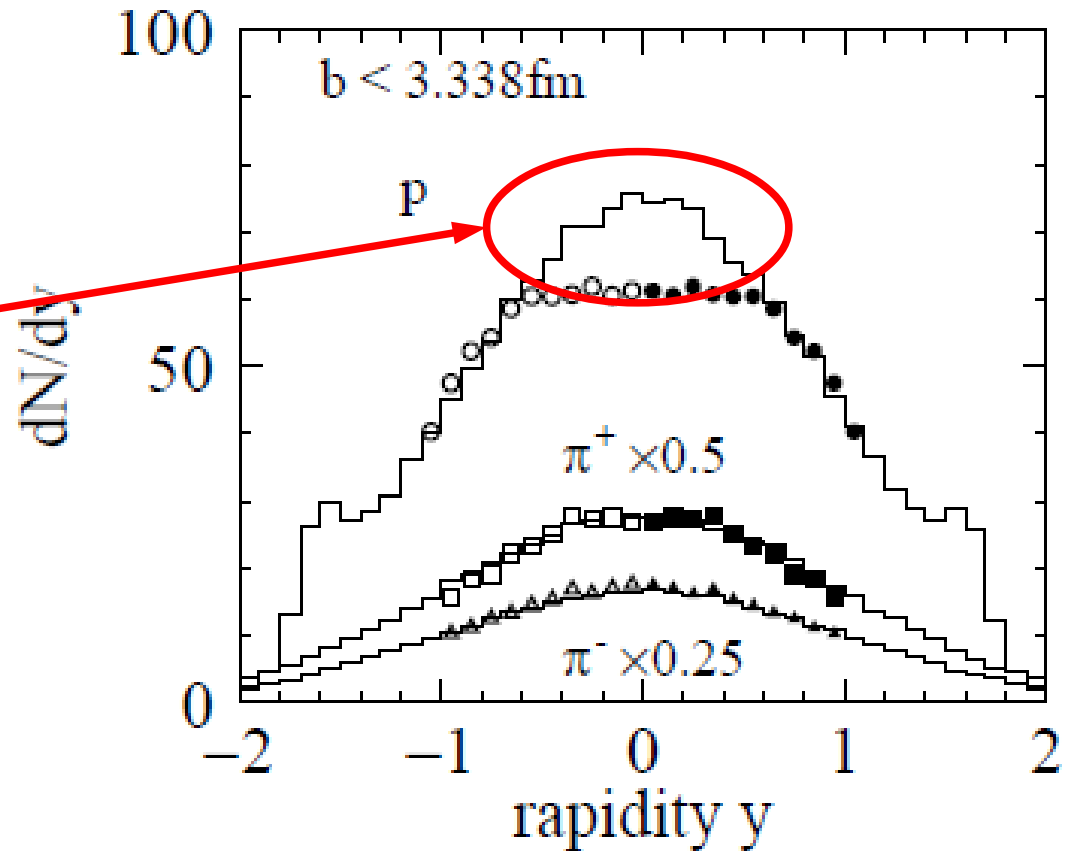


Hadron spectra in Au+Au at AGS

$^{197}\text{Au} + ^{197}\text{Au}$ at 11.6 A GeV/c



$^{197}\text{Au} + ^{197}\text{Au}$ at 11.6 A GeV/c



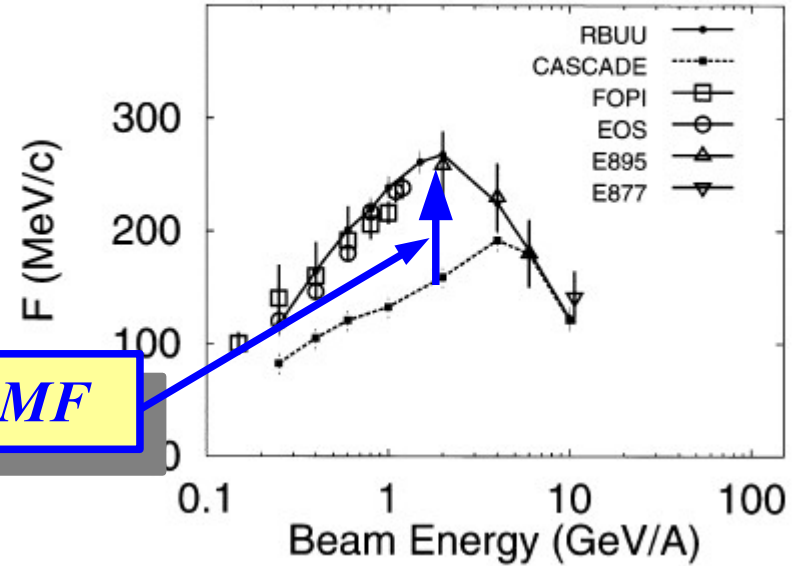
Hadron p_T spectra at AGS are well described, except for low p_T protons (\rightarrow Mean Field Effects).

Mean Field and Particle DOF Effects @ AGS

- Mean Field Effects at AGS
 - Visible but small for p_T spectrum
 - Essential for Flow
- Particle DOF Effects
 - Seen at high p_T

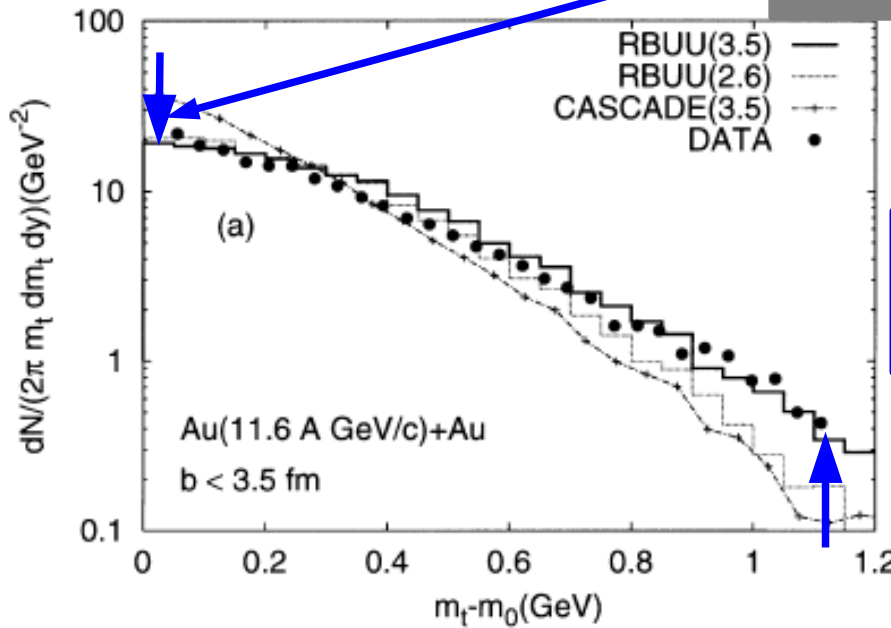
Sahu, Cassing, Mosel, Ohnishi (2000)

P.K. Sahu et al. / Nuclear Physics A 672 (2000) 376–386



P.K. Sahu et al. / Nuclear Physics A 672 (2000) 376–386

Repulsive MF



Switching $\sqrt{s} = 3.5$ GeV (JAM fit)

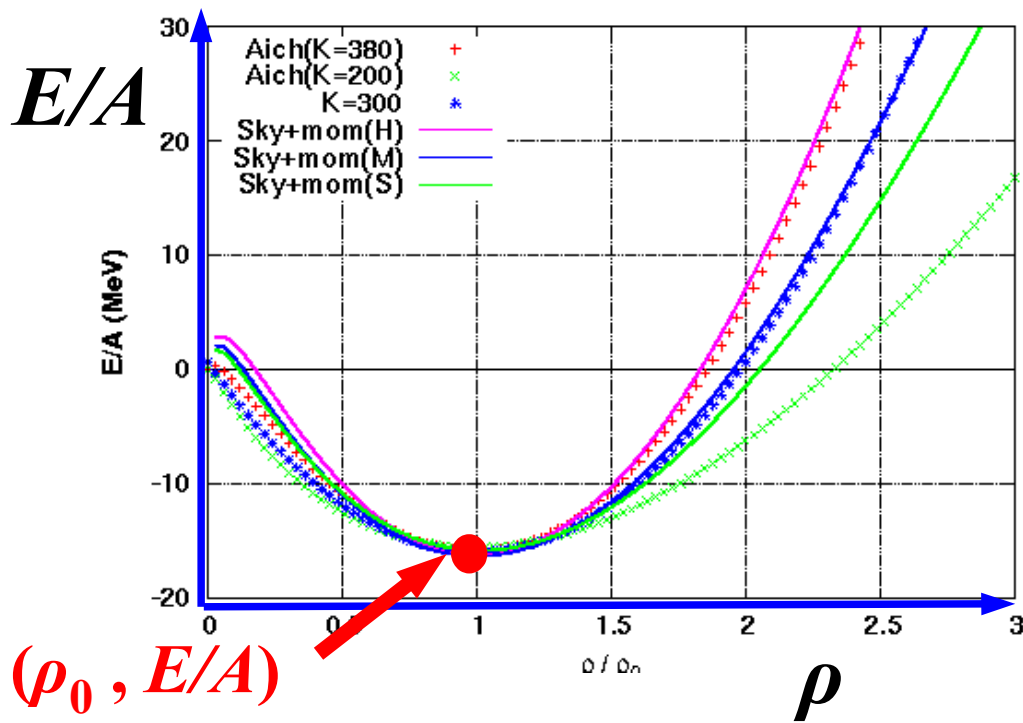
Switching $\sqrt{s} = 2.6$ GeV (HSD default)

Nuclear Mean Field

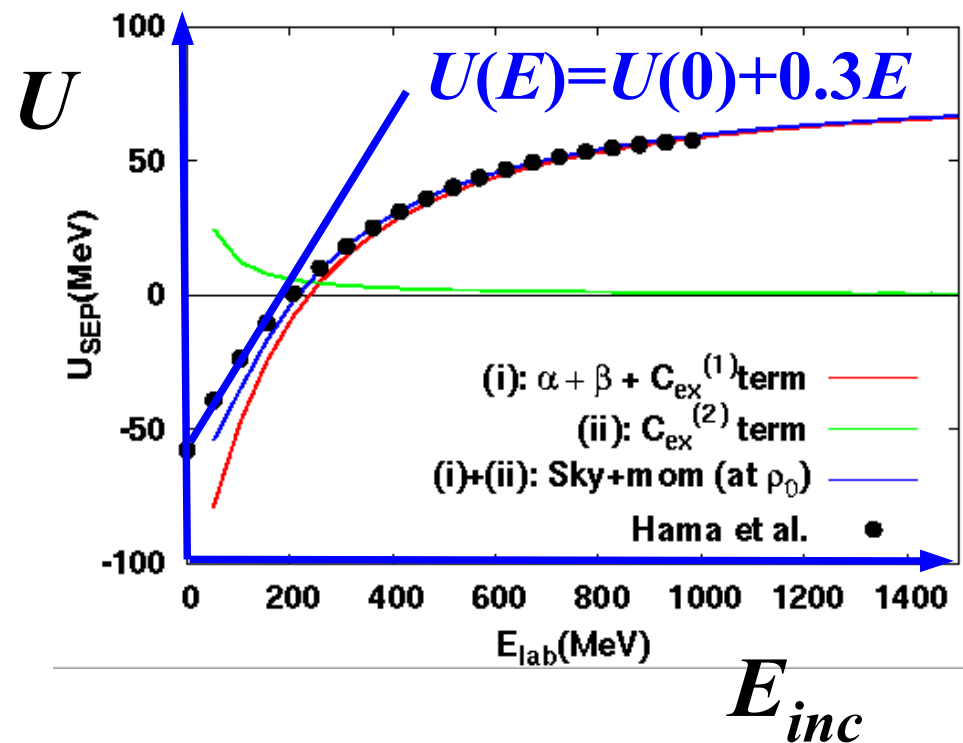
MF has ρ and p -deps.

- ρ dep.: Saturation point $(\rho_0, E/A) = (0.15 \text{ fm}^{-3}, -16.3 \text{ MeV})$
- p dep.: Global potential up to $E=1 \text{ GeV}$ is known from pA scattering

$$U(\rho_0, E) = U(\rho_0, E=0) + 0.3 E$$



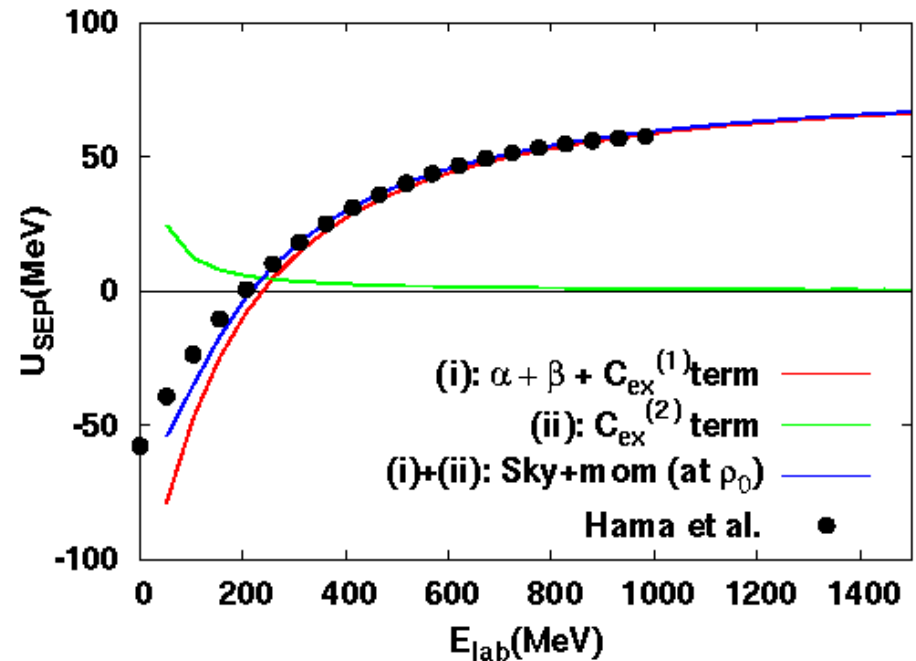
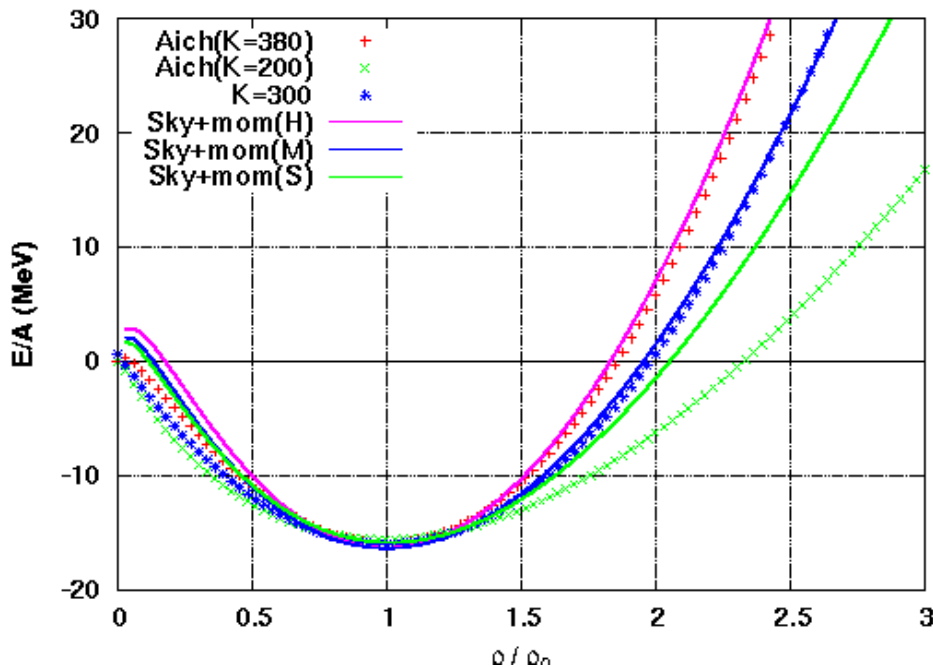
$(\rho_0, E/A)$
 $= (0.15 \text{ fm}^{-3}, -16.3 \text{ MeV})$



Phenomenological Mean Field

■ Skyrme type ρ -Dep. + Lorentzian p -Dep. Potential

$$V = \sum_i V_i = \int d^3 r \left[\frac{\alpha}{2} \left(\frac{\rho}{\rho_0} \right)^2 + \frac{\beta}{\gamma+1} \left(\frac{\rho}{\rho_0} \right)^{\gamma+1} \right] \\ + \sum_k \int d^3 r d^3 p d^3 p' \frac{C_{ex}^{(k)}}{2\rho_0} \frac{f(\mathbf{r}, \mathbf{p}) f(\mathbf{r}, \mathbf{p}')}{1 + (\mathbf{p} - \mathbf{p}')^2 / \mu_k^2}$$



Simplified RQMD treatment of p - and ρ -dep. mean field in JAM
Isse, AO, Otuka, Sahu, Nara, Phys.Rev. C 72 (2005), 064908

Relativistic QMD/Simplified (RQMD/S)

■ RQMD = Constraint Hamiltonian Dynamics

(Sorge, Stocker, Greiner, Ann. of Phys. 192 (1989), 266.)

■ Constraints: $\varphi \approx 0$ (Satisfied on the realized trajectory, by Dirac)

● Variables in Covariant Dynamics = $8N$ phase space: (q_μ, p_μ)

● Variables in EOM = $6N$ phase space

→ We need $2N$ constraints to get EOM

■ On Mass-Shell Constraints

$$H_i \equiv p_i^2 - m_i^2 - 2m_i V_i \approx 0$$

■ Time-Fixation in RQMD/S

$$\chi_i \equiv \hat{a} \cdot (q_i - q_N) \approx 0 \quad (i = 1, \sim N - 1) \quad , \quad \chi_N \equiv \hat{a} \cdot q_N - \tau \approx 0$$

\hat{a} = Time-like unit vector in the Calculation Frame

(Tomoyuki Maruyama et al., Prog. Theor. Phys. 96(1996), 263.)

RQMD/S (cont.)

- Hamiltonian is made of constraints

$$H = \sum_i u_i \phi_i \quad (\phi_i = H_i (i=1 \sim N), \chi_{i-N} (i=N+1 \sim 2N))$$

- Time Development $\frac{d f}{d \tau} = \frac{\partial f}{\partial \tau} + \{f, H\}$, $\{q_\mu, p_\nu\} = g_{\mu\nu}$

- Lagrange multipliers are determined to keep constraints
 → *We can obtain the multipliers analytically in RQMD/S*

$$\frac{d \phi_i}{d \tau} \approx 0 \rightarrow \delta_{i,2N} + \sum_j u_j \{ \phi_i, \phi_j \} \approx 0$$

- Equations of Motion

$$H = \sum_i (p_i^2 - m_i^2 - 2m_i V_i) / 2p_i^0, \quad p_i^0 = E_i = \sqrt{\vec{p}_i^2 + m_i^2 + 2m_i V_i}$$

$$\frac{d \vec{r}_i}{d \tau} \approx -\frac{\partial H}{\partial \vec{p}_i} = \frac{\vec{p}}{p_i^0} + \sum_j \frac{m_j}{p_j^0} \frac{\partial V_j}{\partial \vec{p}_i}, \quad \frac{d \vec{p}_i}{d \tau} \approx \frac{\partial H}{\partial \vec{r}_i} = -\sum_j \frac{m_j}{p_j^0} \frac{\partial V_j}{\partial \vec{r}_i}$$

We can include MF in an almost covariant way in molecular dynamics

Particle “DISTANCE”

$$r_{Tij}^2 \equiv r_\mu r^\mu - \left(r_\mu P_{ij}^\mu \right)^2 / P_{ij}^2 = \vec{r}^2 \quad (\text{in } CM)$$

$$P_{ij} \equiv p_i + p_j \quad , \quad r \equiv r_i - r_j$$

Particle “Momentum Difference”

$$p_{Tij}^2 \equiv p_\mu p^\mu - \left(p_\mu P_{ij}^\mu \right)^2 / P_{ij}^2 = \vec{p}^2 \quad (\text{in } CM)$$

$$p \equiv p_i - p_j$$

Lorentz Invariant, and Becomes Normal Distance in CM !

Collective Flows at J-PARC Energies

What is Collective Flow ?

(Directed) Flow (dP_x/dY)

Stiffness (Low E)
+ Time Scale (High E)

Elliptic Flow (V_2)

Thermalization
& Pressure Gradient

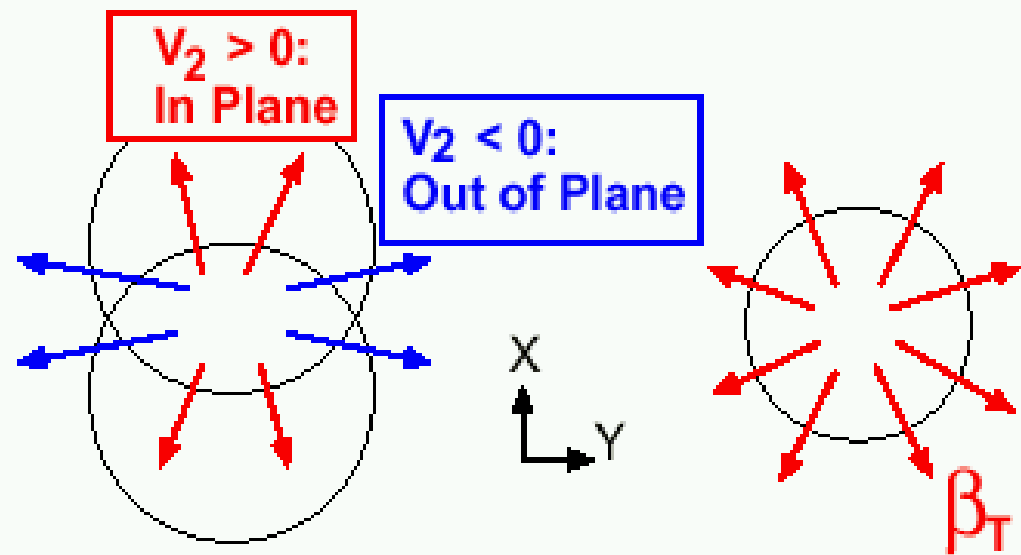
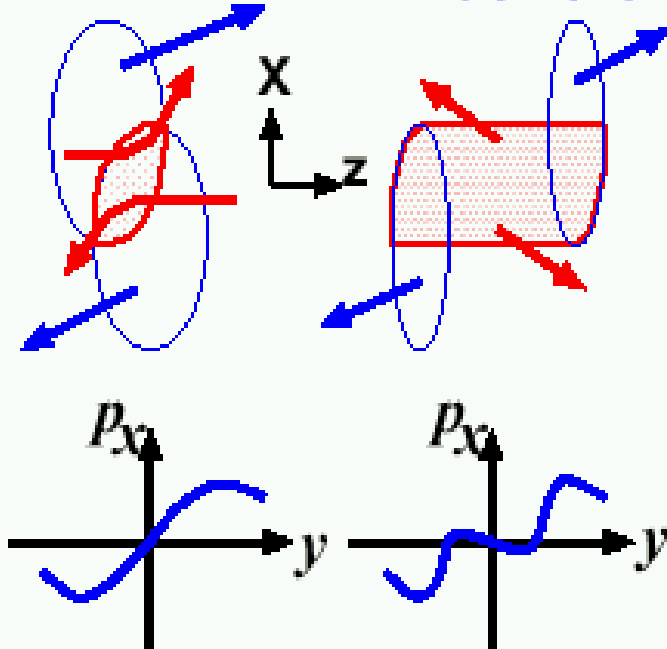
Radial Flow (β_T)

Pressure History

$$\epsilon \frac{DV}{Dt} = -\nabla P$$

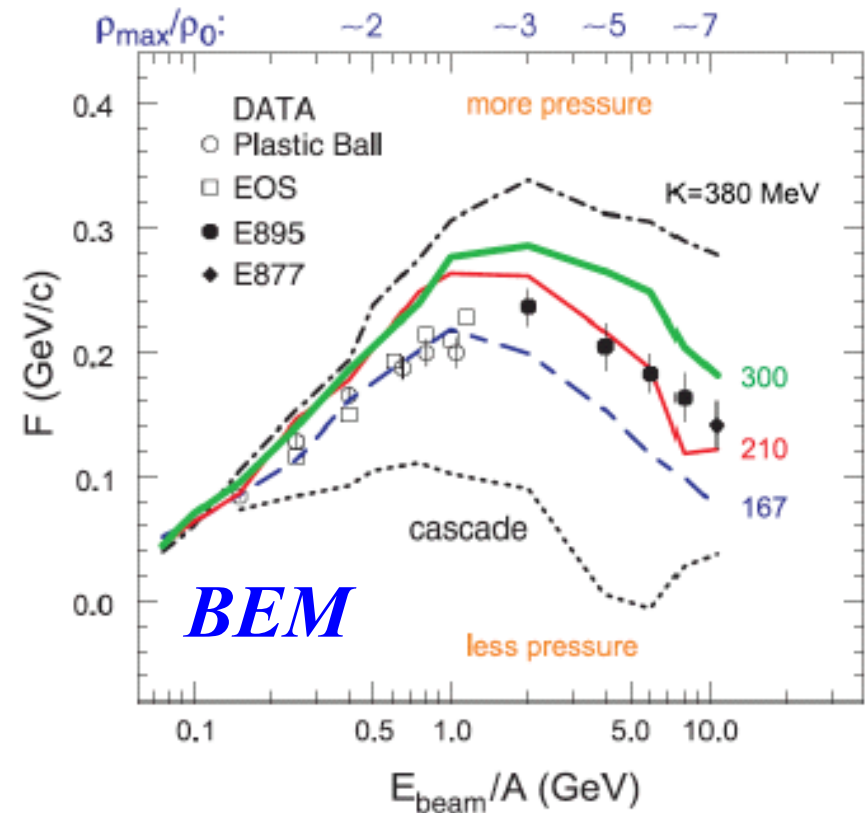
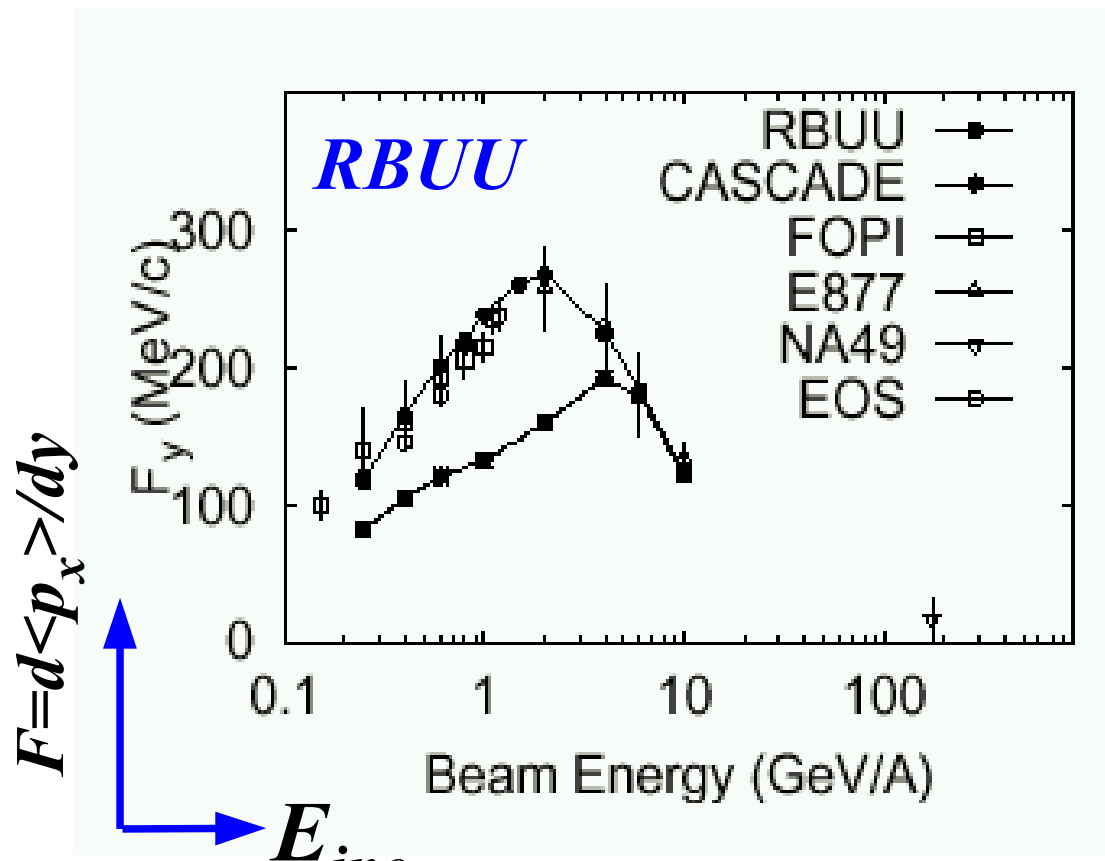
$$\rightarrow V = \int_{\text{path}} \frac{-\nabla P dt}{\epsilon}$$

Until AGS Above SPS



Side Flow at AGS Energies

- Relativistic BUU (RBUU) model: $K \sim 300 \text{ MeV}$
(Sahu, Cassing, Mosel, AO, Nucl. Phys. A672 (2000), 376.)
- Boltzmann Equation Model (BEM): $K=167\sim 210 \text{ MeV}$
(P. Danielewicz, R. Lacey, W.G. Lynch, Science 298(2002), 1592.)

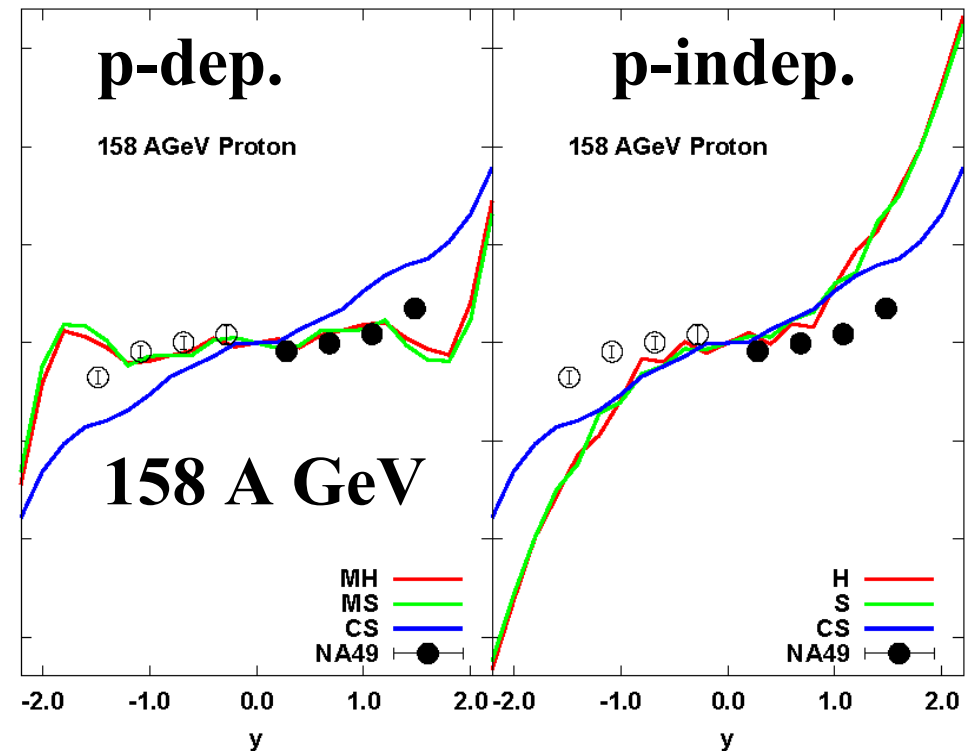
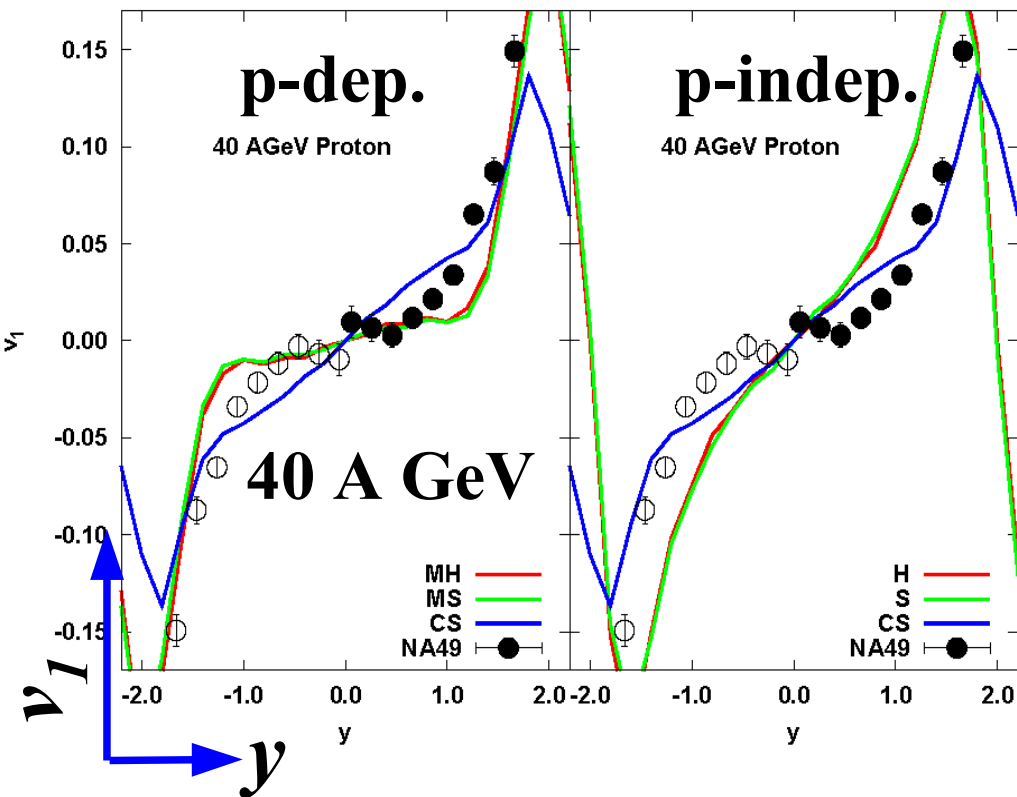


Directed flow v_1 at SPS

Isse, AO, Otuka, Sahu, Nara, PRC 72 (2005), 064908

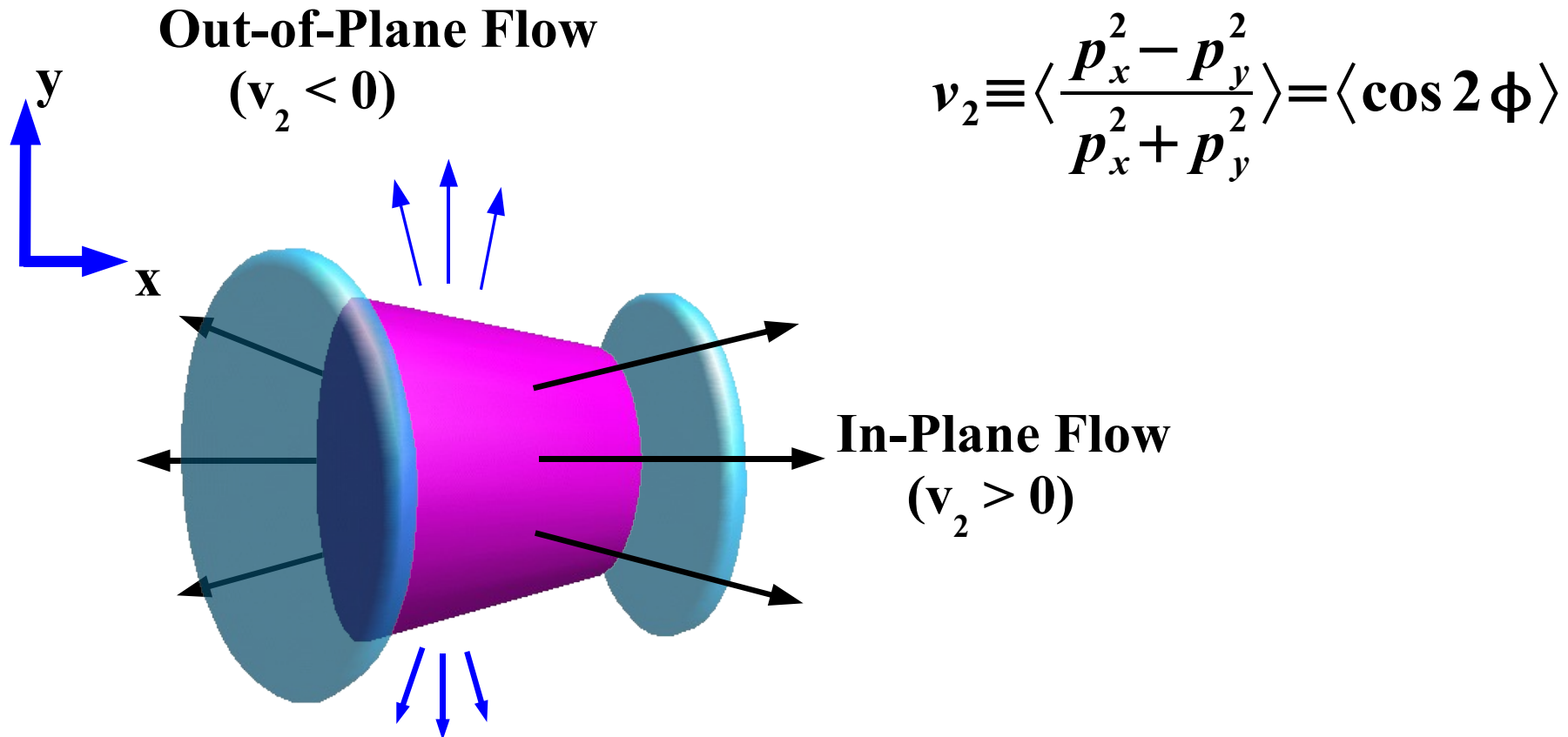
JAM-RQMD/S

- p-dep. (indep.) MF suppresses (enhances) v_1 . $v_1 = \langle \cos \phi \rangle = \langle p_x / p_T \rangle$
- “Wiggle” behavior appears with p-dep. MF at 158 A GeV.



Elliptic Flow

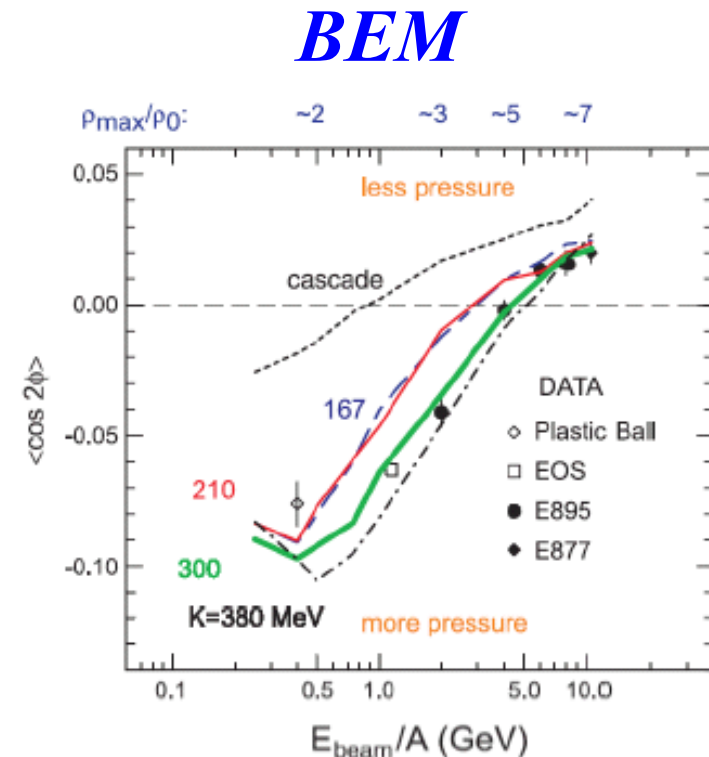
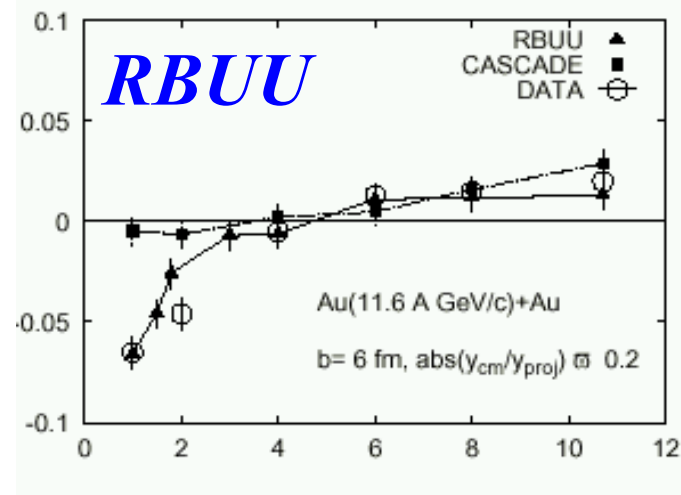
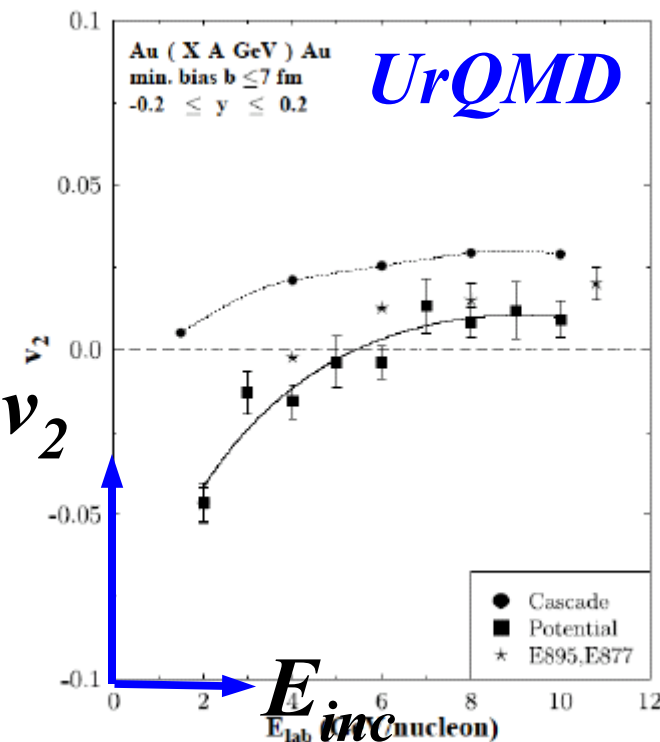
- What is Elliptic Flow ? → Anisotropy in P space
- Hydrodynamical Picture
 - Sensitive to the Pressure Anisotropy in the Early Stage
 - Early Thermalization is Required for Large v_2



Elliptic Flow at AGS

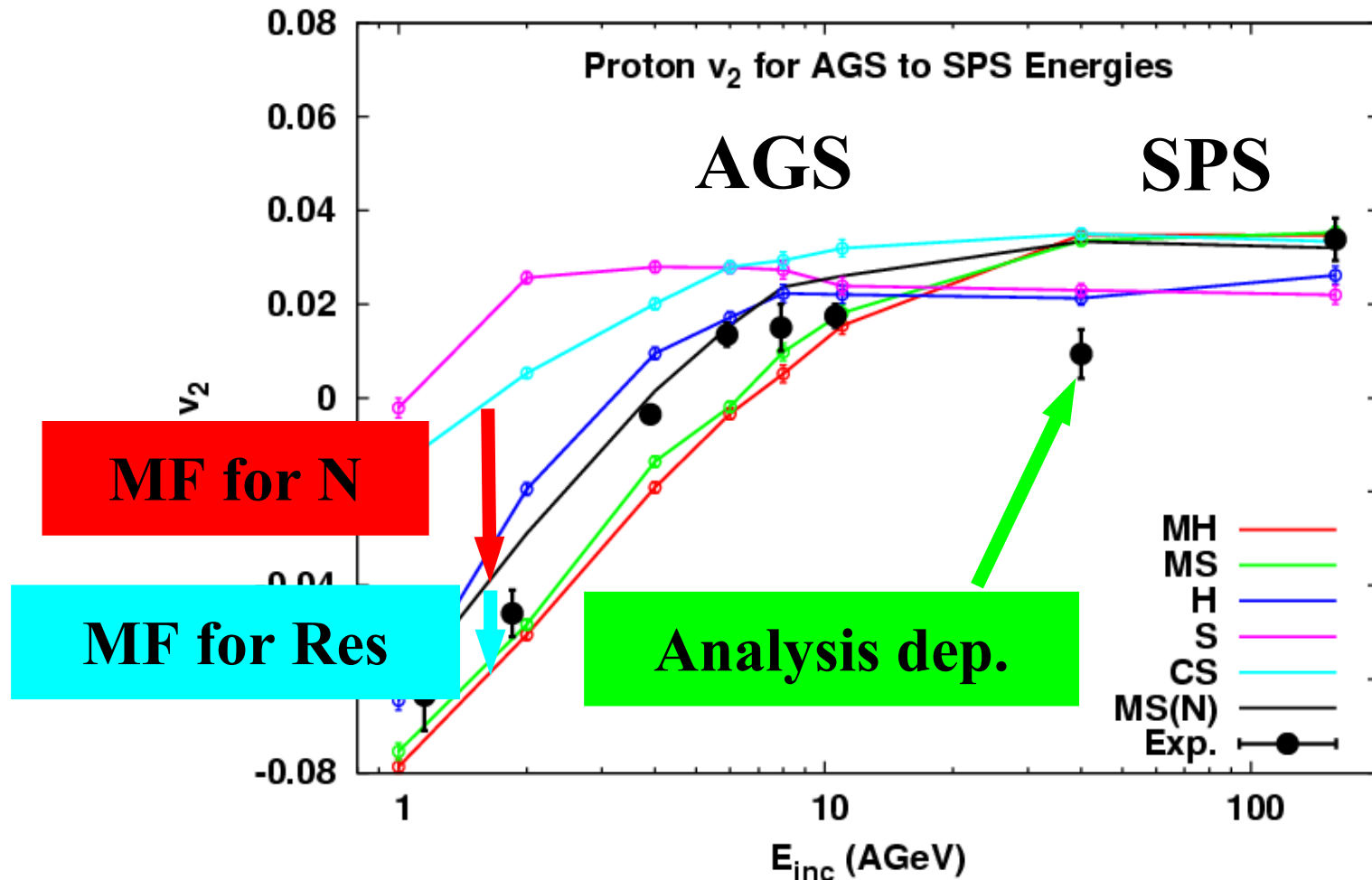
Strong Squeezing Effects at low E (2-4 A GeV)

- UrQMD: Hard EOS (S.Soff et al., nucl-th/9903061)
- RBUU (Sahu-Cassing-Mosel-AO, 2000): $K \sim 300$ MeV
- BEM(Danielewicz2002): $K = 167 \rightarrow 300$ MeV

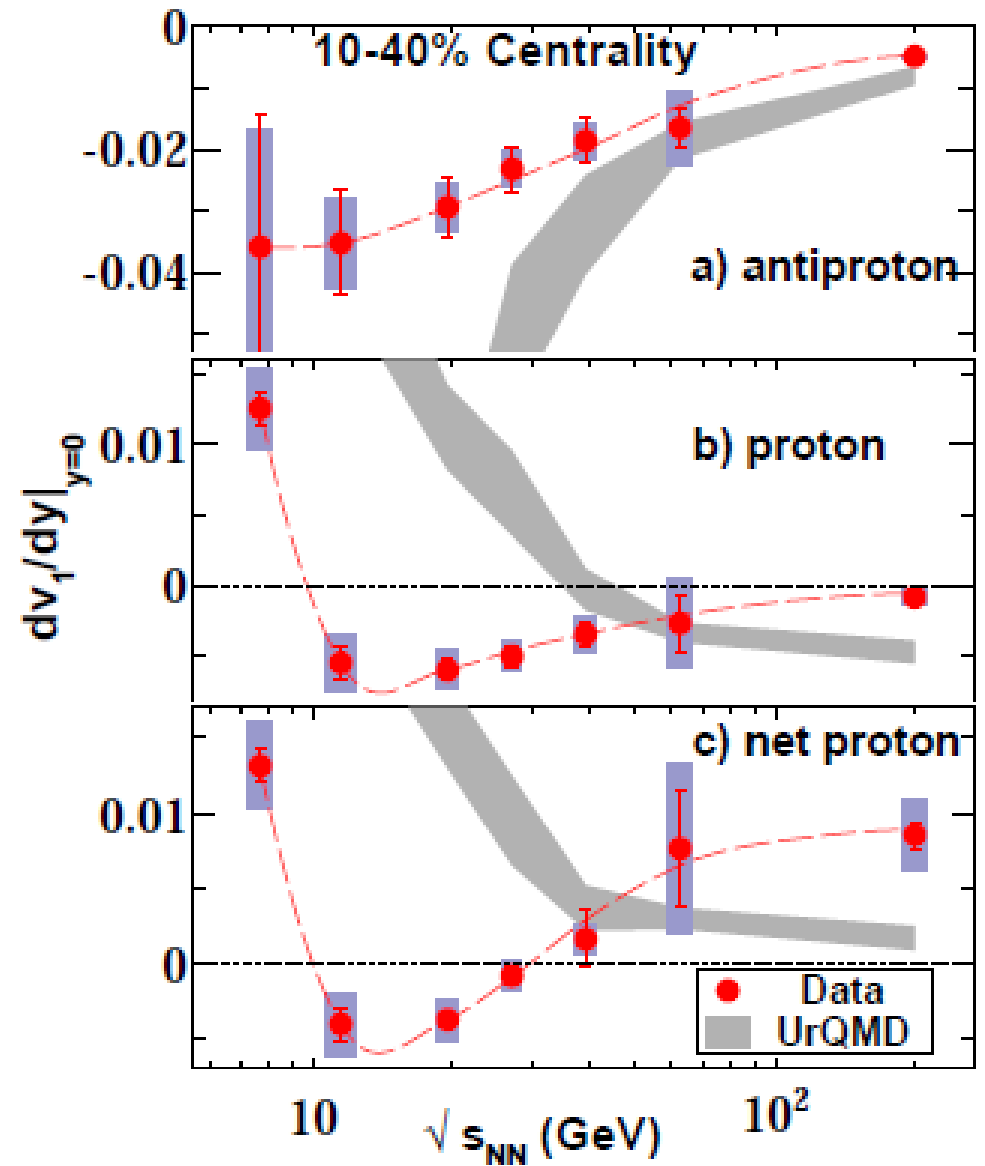
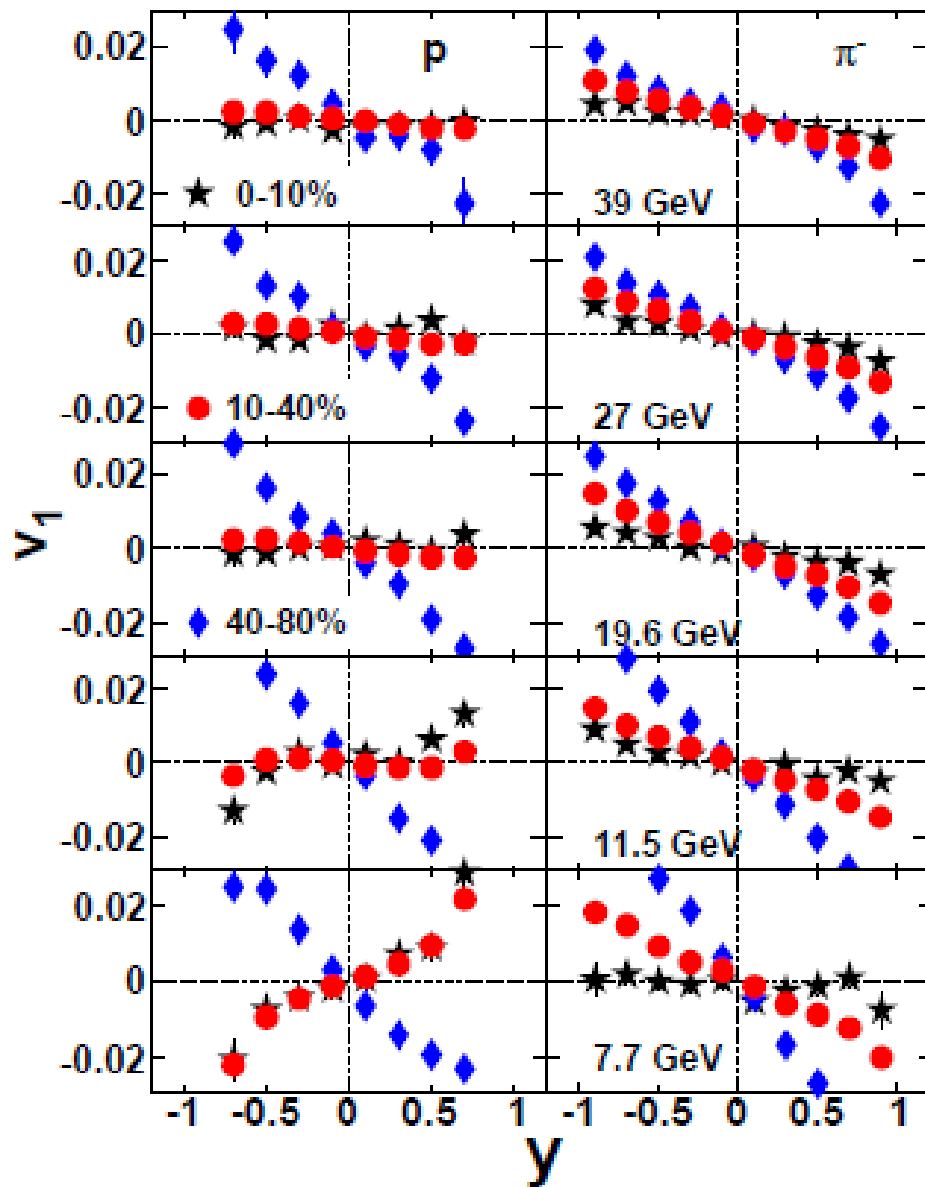


Elliptic Flow from AGS to SPS

- JAM-MF with p dep. MF explains proton v_2 at 1-158 A GeV
 - v_2 is not very sensitive to K (incompressibility)
 - Data lies between MS(B) and MS(N)

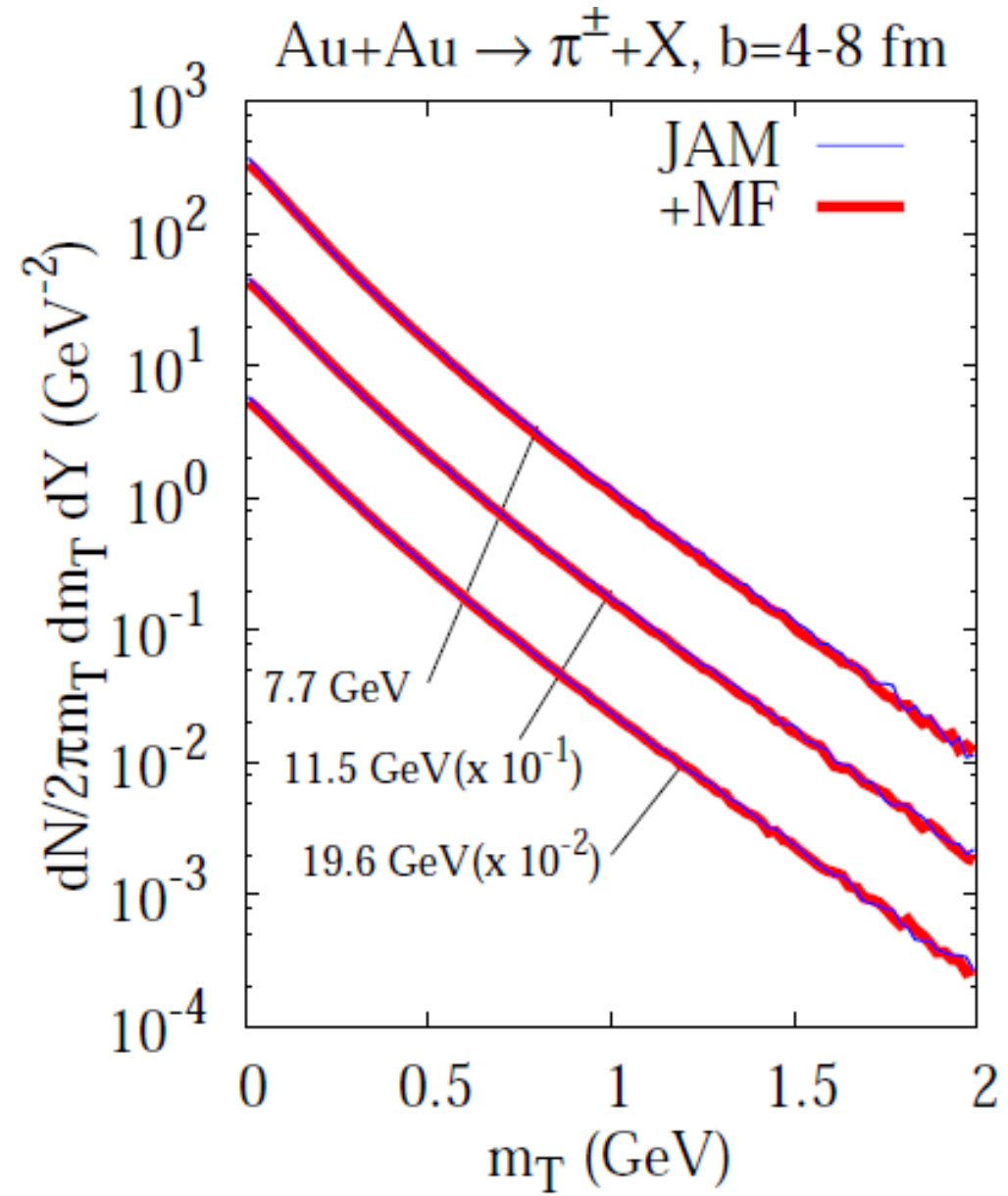
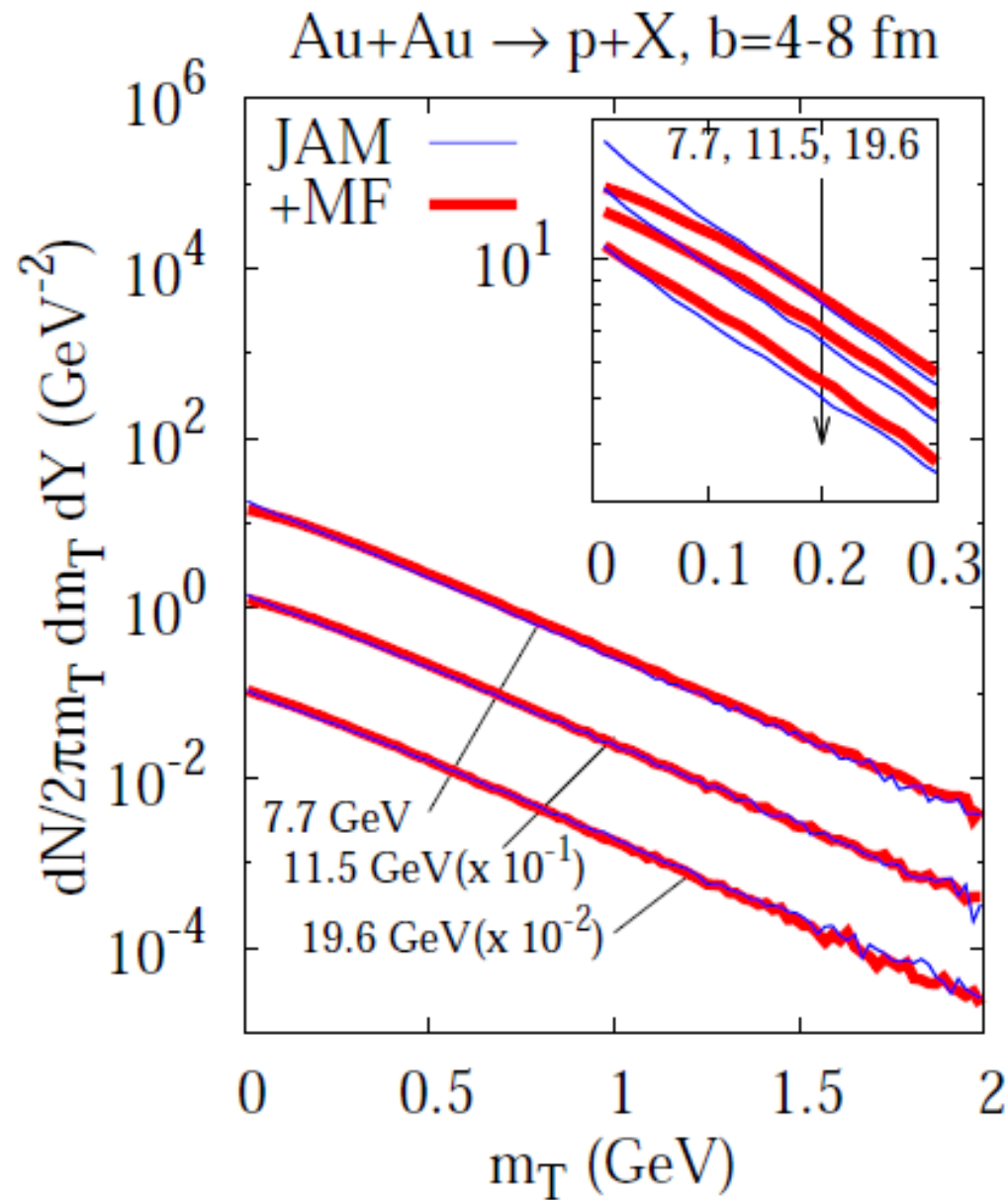


New Data from RHIC-BES



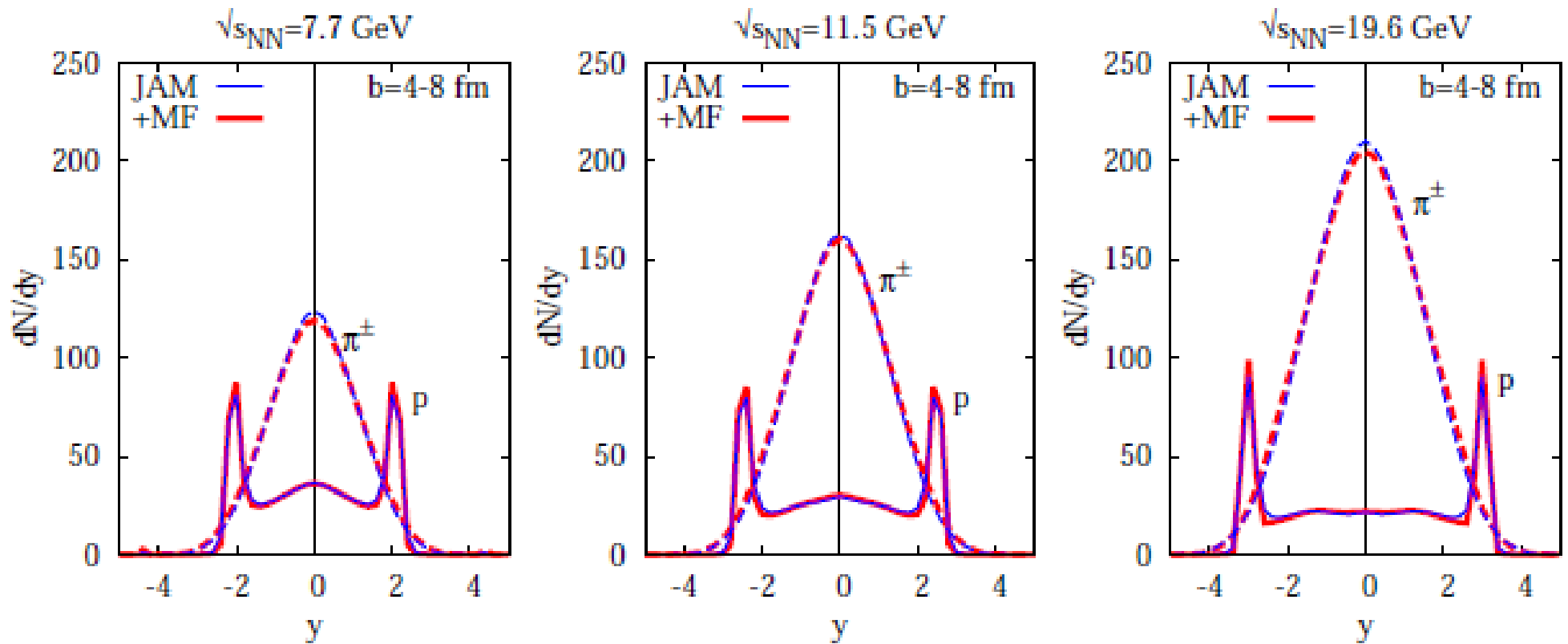
L. Adamczyk et al. (STAR Collab.), PRL 112('14)162301.

JAM Results at $\sqrt{s_{NN}}=7.7, 11.5$ and 19.6 GeV



preliminary, Nara, AO, work in progress

JAM Results at $\sqrt{s_{NN}}=7.7, 11.5$ and 19.6 GeV

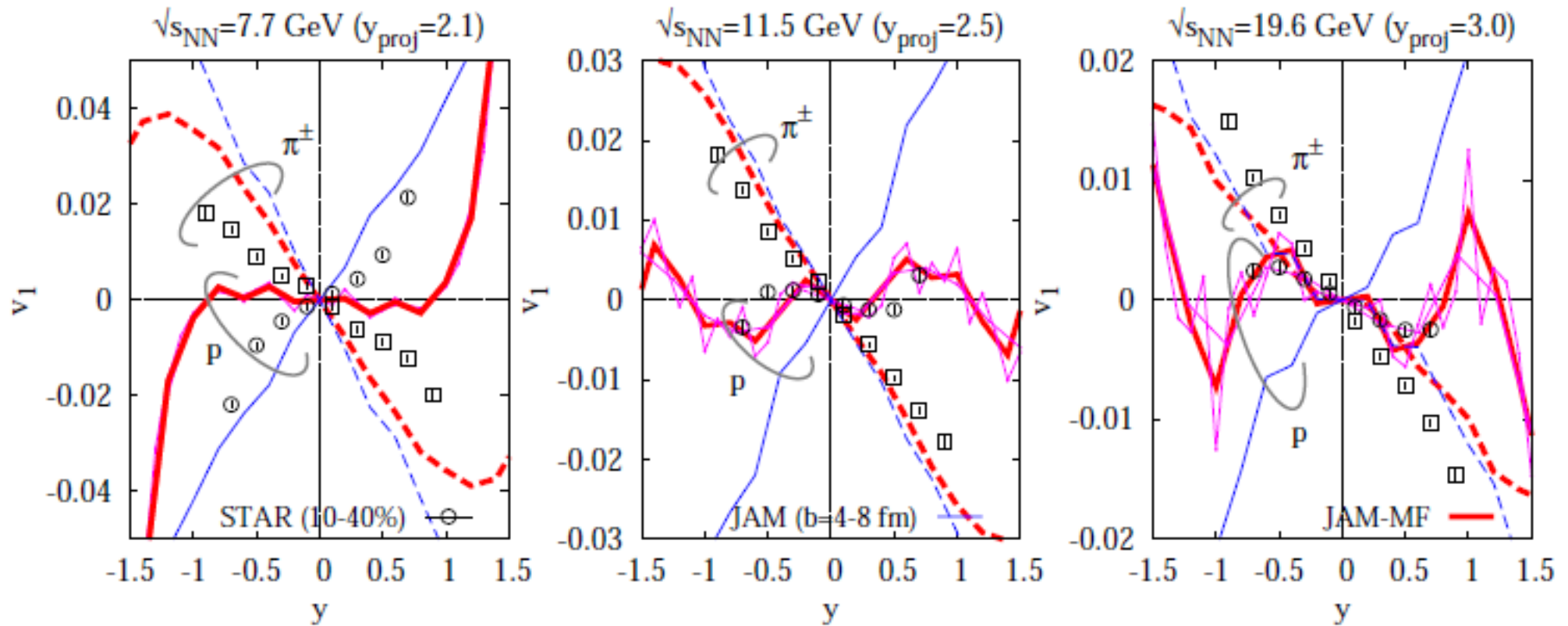


preliminary, Nara, AO, work in progress

Ohnishi @ KEK-HIC workshop, Nov.26-27, 2014

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JAM Results at $\sqrt{s_{NN}}=7.7, 11.5$ and 19.6 GeV



preliminary, Nara, AO, work in progress

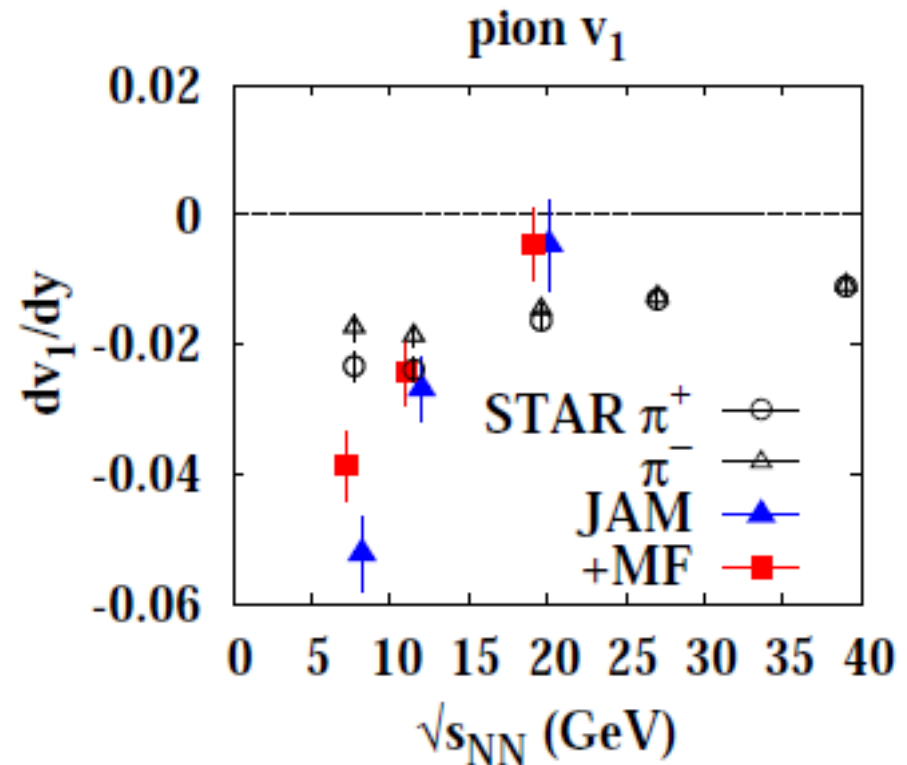
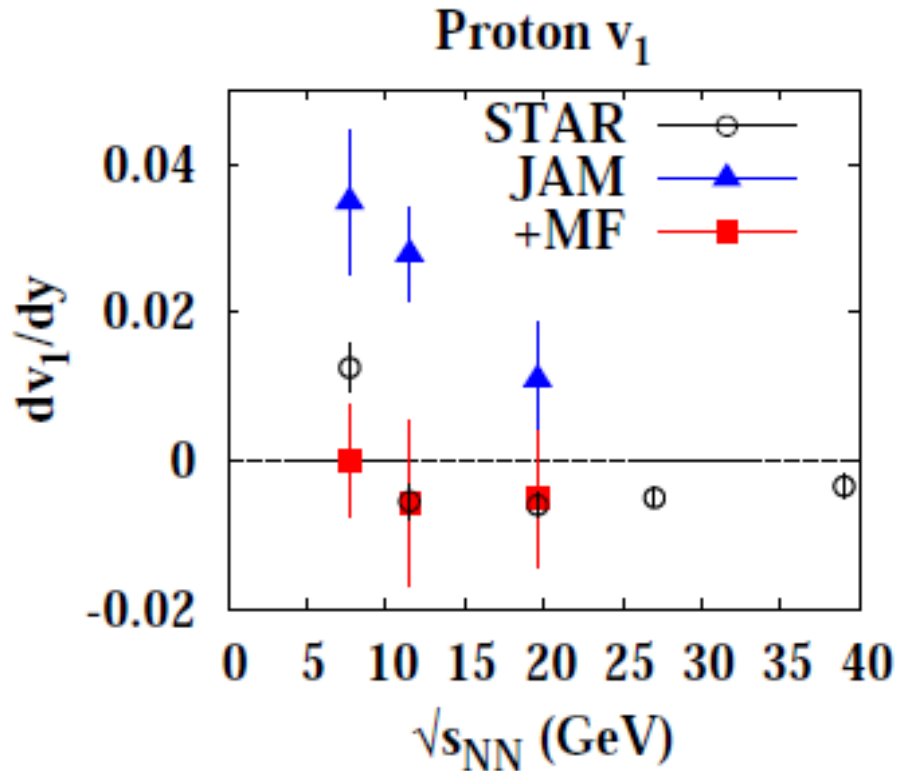
Ohnishi @ KEK-HIC workshop, Nov.26-27, 2014

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JAM Results at $\sqrt{s_{NN}}=7.7, 11.5$ and 19.6 GeV

■ Incident E. dependence of dv_1/dy

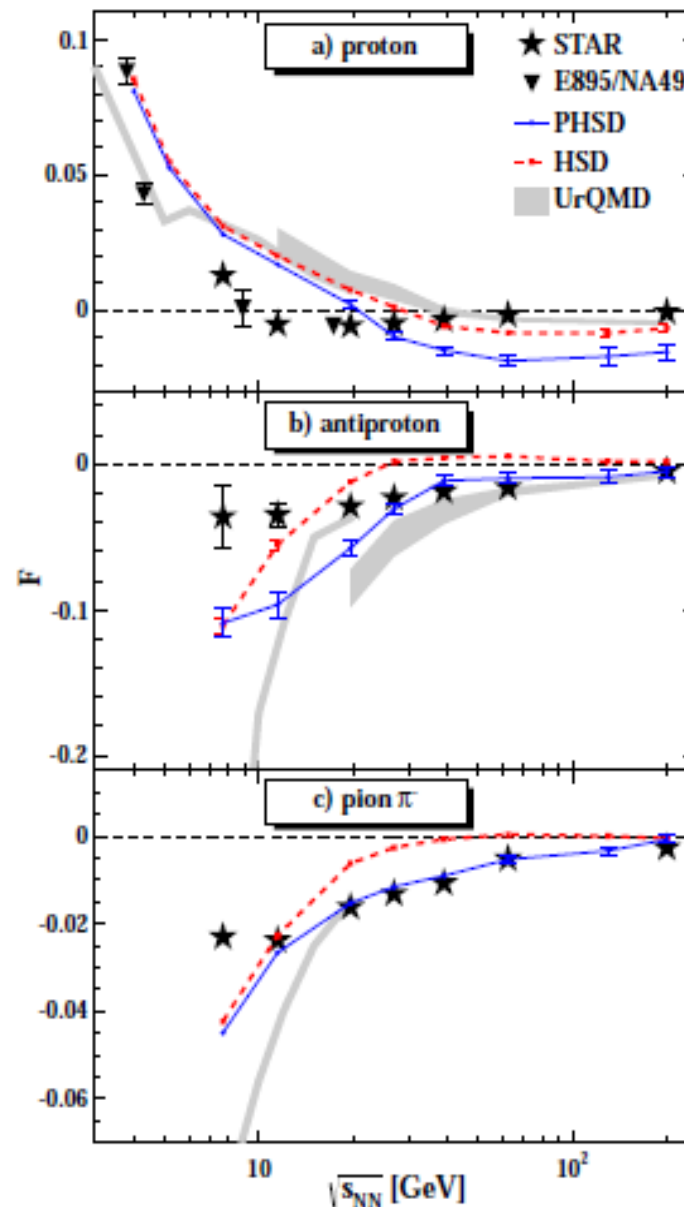
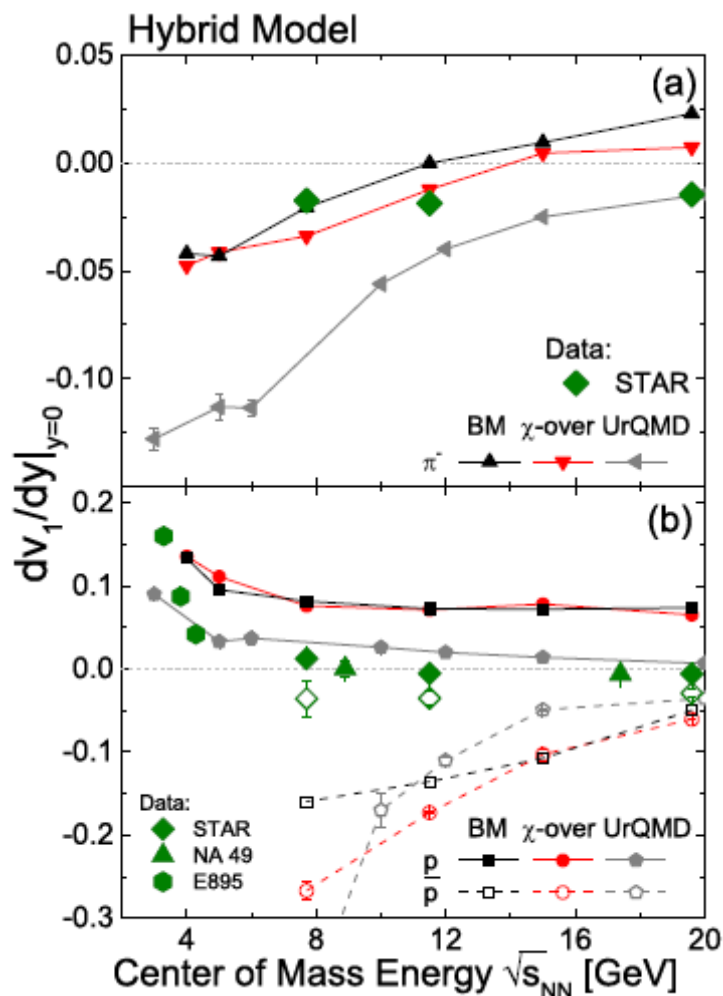
- proton: JAM (w/o MF) \sim UrQMD,
JAM (w/ MF) \sim STAR (underestimate at 7.7 GeV)
- pion: energy dep. is too much (Note: No direct potential effects on pions.)



preliminary, Nara, AO, work in progress

Comparison with other approaches

- Both Hybrid model (Frankfurt), PHSD (Giessen) show higher balance E.



J. Steinheimer, J. Auvinen, H. Petersen, M. Bleicher, H. Stöcker, PRC89 ('14) 054913

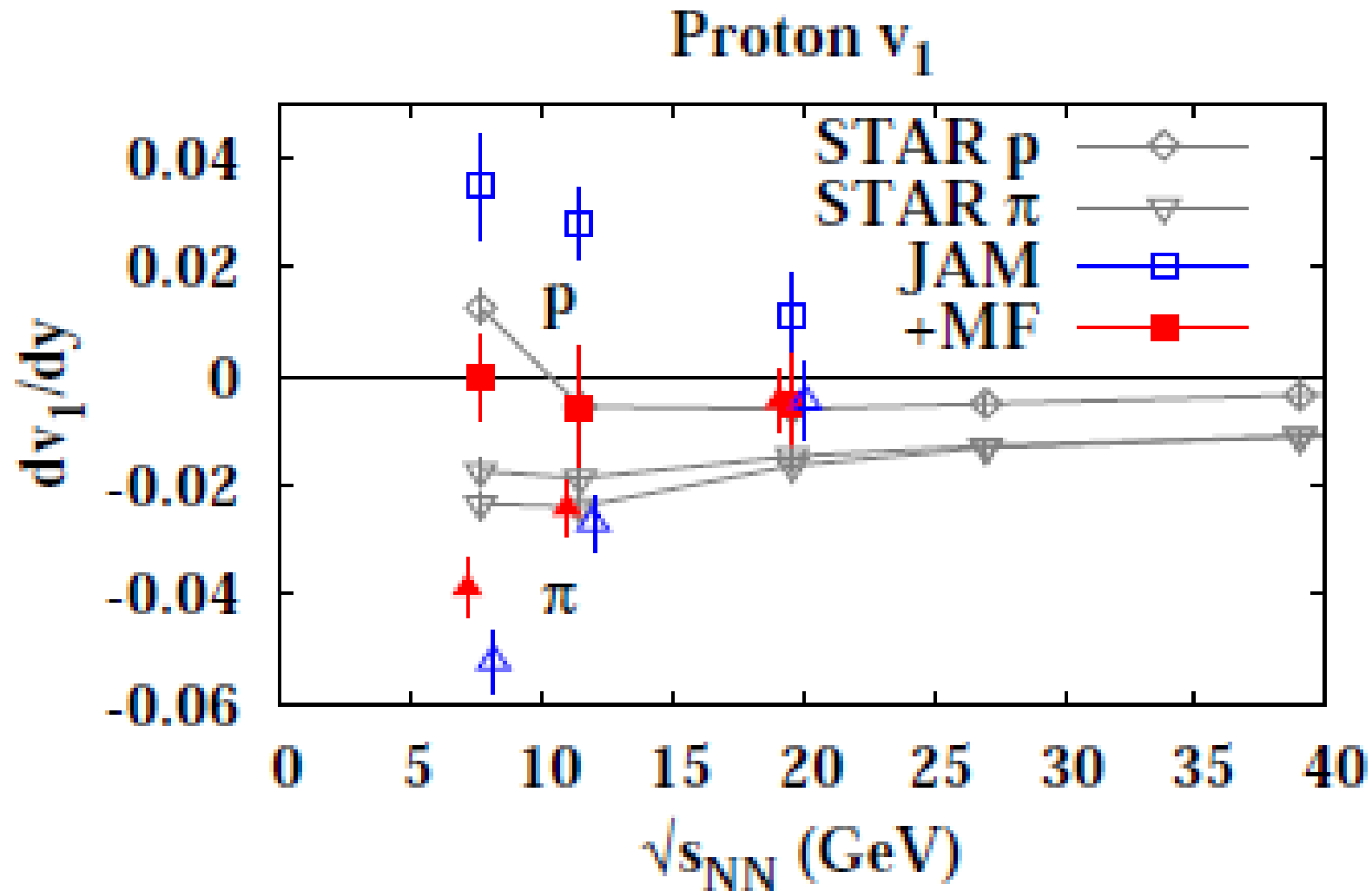
V. P. Konchakovski, W. Cassing, Yu. B. Ivanov, V. D. Toneev, PRC90('14)014903

Summary

- J-PARC energy での重イオン衝突では超高密度の物質が作られ、かつ解かれていない問題が多く残されている。
 - Horn, Step, Dale, proton # fluc., Collective flows, ...
- ハドロン輸送モデルにより、 $E/A=1-158$ GeV ($\sqrt{s_{NN}}=2.5-20$ GeV) における集団フローは定性的に説明される。
 - カスケードモデルにより、ハドロンスペクトルはほぼ説明可能。
 - v_1, v_2 などの集団フローは平均場に敏感である。
 $E/A > 300$ MeV では低密度でも核子・核ポテンシャルは斥力なので、運動量依存ポテンシャルを用いることが決定的。
 - RHIC-BES で見られる非単調な dv_1/dy は定性的 (半定量的) に説明される。(Isse et al. (2005) のプログラムをそのまま利用)
- To do
 - 素過程断面積の改善、 $\pi, \Delta, N^*, K, Y...$ への平均場効果、quark-gluon 自由度の導入、(viscous)hydro との結合、...
 - 興味のある方はやってみませんか？

Thank you for your attention.

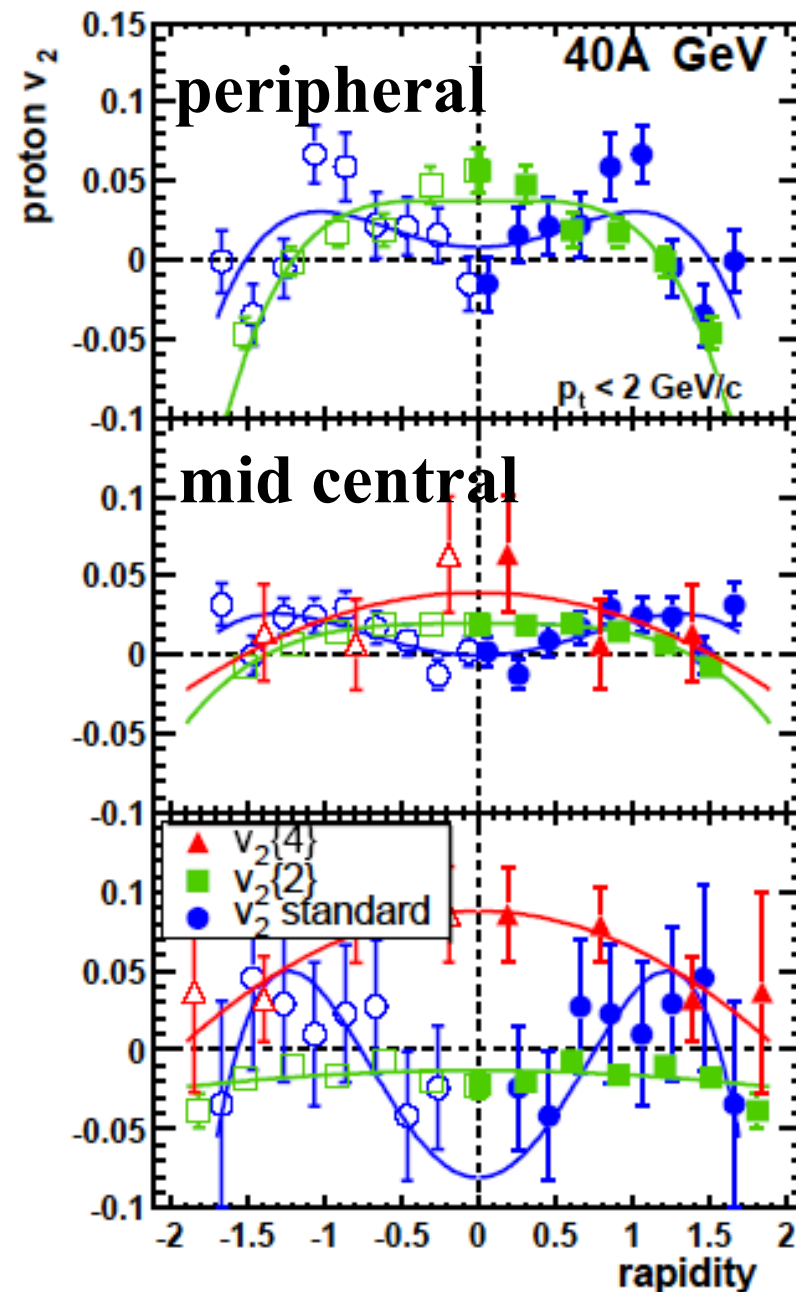
JAM Results at $\sqrt{s_{NN}}=7.7, 11.5$ and 19.6 GeV



Nara, AO, work in progress

Dip of V_2 at 40 A GeV: Phase Transition ?

- Dip of V_2 at 40 A GeV may be a signal of QCD phase transition at high baryon density.
(Cassing et al.)
- However, the data is too sensitive to the way of the analysis (reaction plane/two particle correlation).
 - We have to wait for better data.



Flow and EOS; to be continued

- In addition to the ambiguities in in-medium cross sections, Res.-Res. cross sections, we have model dependence.

- RBUU (*e.g. Sahu, Cassing, Mosel, AO, Nucl. Phys. A672 (2000), 376.*)

- ◆ In RMF, Strong cut-off for meson-N coupling in RMF
→ Smaller EOS dep.

- Scalar potential interpretation in BUU

Larionov, Cassing, Greiner, Mosel, PRC62,064611('00), Danielewicz, NPA673,375('00)

$$\varepsilon(\mathbf{p}, \rho) = \sqrt{[m + U_s(\mathbf{p}, \rho)]^2 + \mathbf{p}^2} = \sqrt{m^2 + \mathbf{p}^2} + U(\mathbf{p}, \rho)$$

- ◆ Due to the Scalar potential nature, EOS dependence is smaller.

- Scalar/Vector Combination *Danielewicz, Lacey, Lynch, Science 298('02), 1592*

$$\varepsilon(p, \rho) = m + \int_0^p dp' v^*(p', \rho) + \tilde{U}(\rho), \quad v^*(p, \rho) = \frac{p}{\sqrt{p^2 + [m^*(p, \rho)]^2}}$$

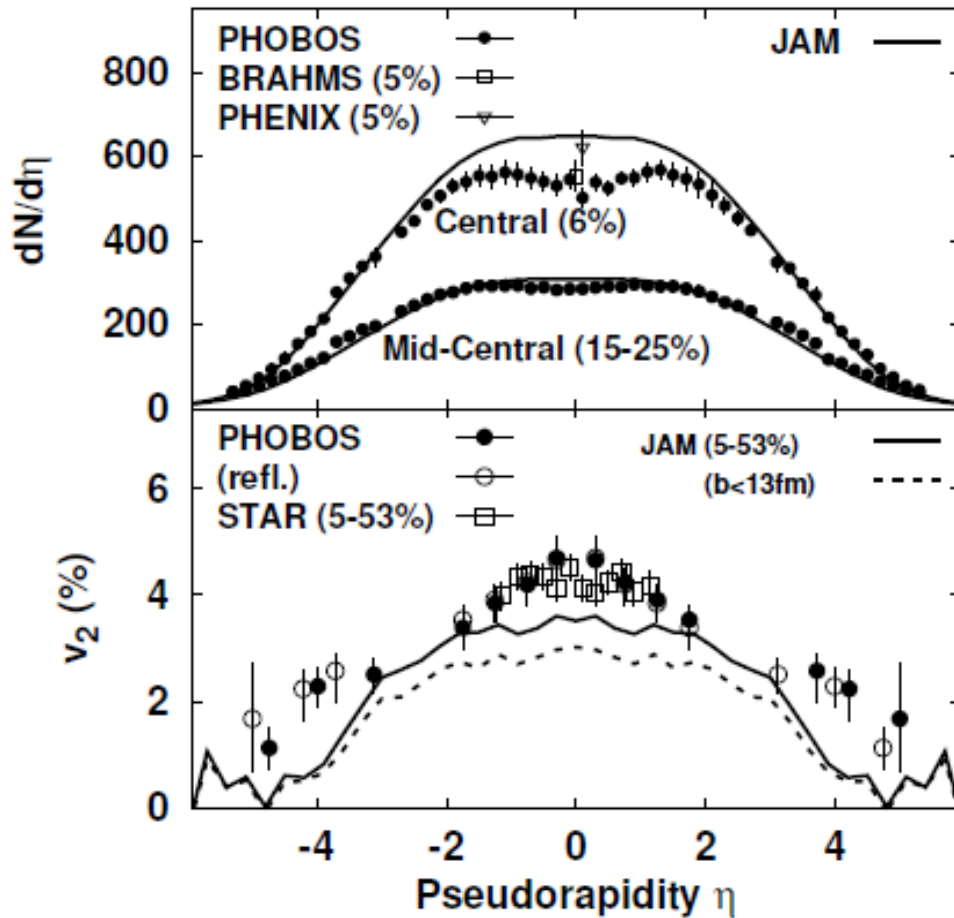
- ◆ Relatively Strong EOS dependence even at high energy

- JAM-RQMD/S *Isse, AO, Otuka, Sahu, Nara, PRC 72 (2005), 064908*

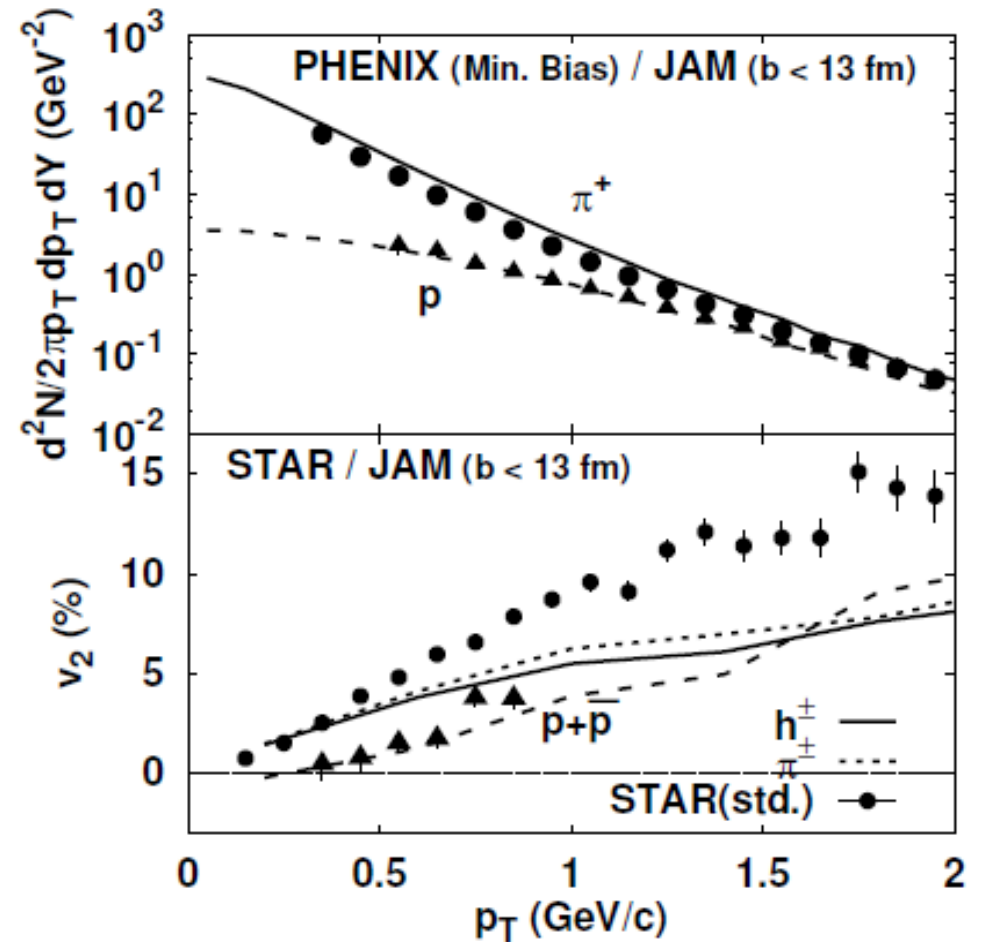
- ◆ Similar to the Scalar model BUU

JAM at RHIC

Au+Au ($\sqrt{s} = 130$ A GeV)



Au+Au ($\sqrt{s} = 130$ A GeV)



Sahu, AO, Isse, Otuka, Phatak (2006)