
Lambda-Lambda interaction from two-particle intensity correlation in relativistic heavy-ion collisions

Akira Ohnishi¹

in collaboration with

Kenji Morita¹²³, Takenori Furumoto⁴

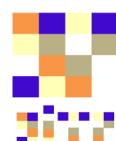
1. YITP, Kyoto Univ., 2. FIAS, 3. U. Wroclaw, 4. Ichinoseki CT

*Molecule-type workshop on "Selected topics in the physics
of the Quark-Gluon Plasma and Ultrarelativistic Heavy Ion Collisions"*
2015/09/07 --- 2015/09/26

based on K. Morita, T. Furumoto, AO,

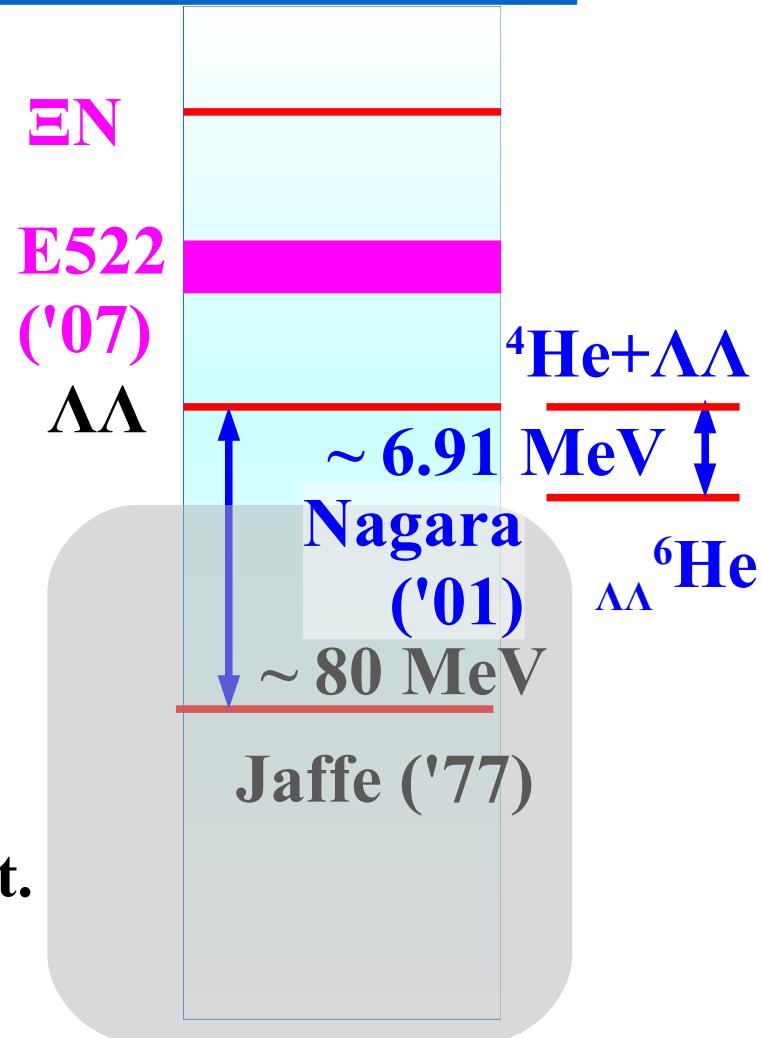
Phys. Rev. C91 (2015), 024916 [arXiv:1408.6682]

Same talk as HYP2015 talk by AO



$\Lambda\Lambda$ interaction & Exotic Hadron (H particle)

- Deeply bound H ? *Jaffe ('77)*
 - Strong Attraction from Color Mag. Int.
→ 80 MeV below $\Lambda\Lambda$
- No observation for a long time.
- Why there is no bound state ?
 - Repulsive Instanton Ind. Int.
Oka, Takeuchi ('91)
- Nagara event $_{\Lambda\Lambda}^6\text{He}$ *Takahashi et al. ('01)*
 - No deeply bound “H”, Weakly Att. $\Lambda\Lambda$ int.
- Resonance or Bound “H” ?
 - 2 σ “bump” at $E_{\Lambda\Lambda} \sim 15$ MeV
Yoon et al. (KEK-E522) ('07)
 - Bound H at large ud quark masses
HAL QCD & NPLQCD ('11), Haidenbauer, Meissner ('11)



$\Lambda\Lambda$ interaction models

■ Boson exchange potential

Nijmegen potentials (ND, NF, NSC89, NSC97, ESC08)

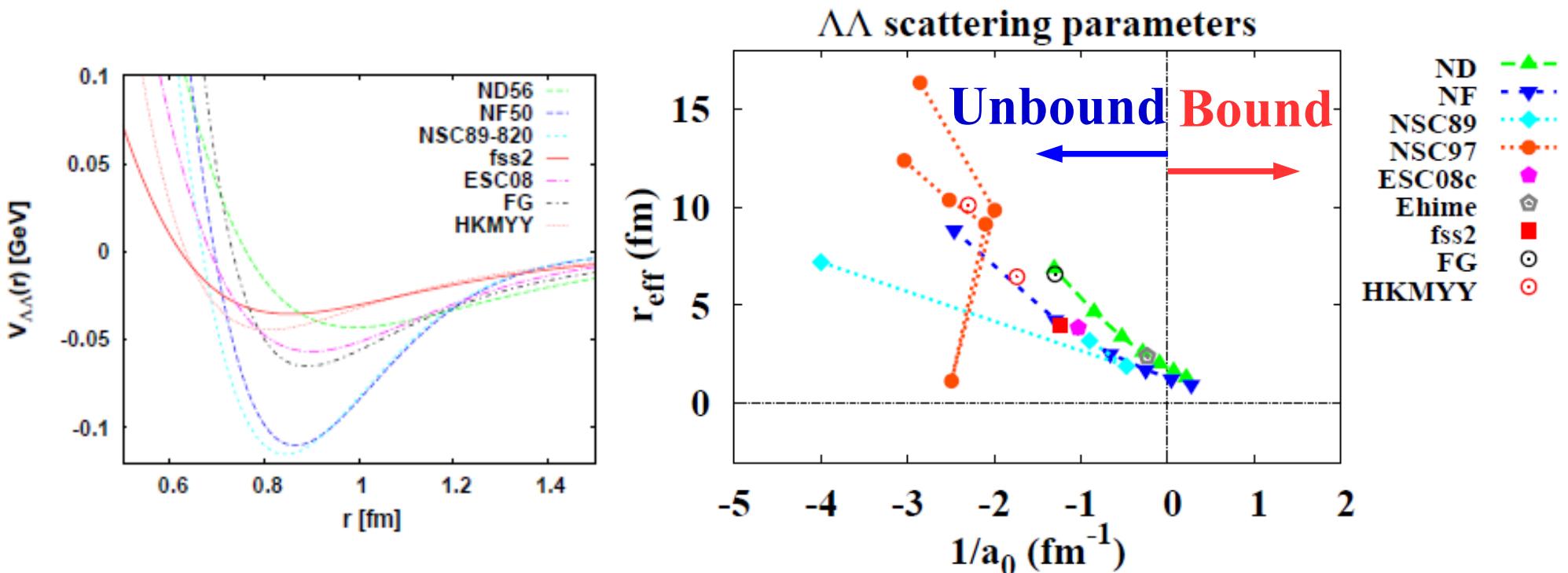
Nagels+ ('77, '79), Maessen+ ('89), Rijken+ ('99, '10)

Ehime *Ueda et al.* ('98)

■ Quark model interaction: fss2 *Fujiwara et al.* ('07)

■ Tuned potential to Nagara

Filikhin, Gal ('02) (FG), Hiyama et al. ('02, '10)(HKMYY)



$\Lambda\Lambda$ interaction and Neutron Star

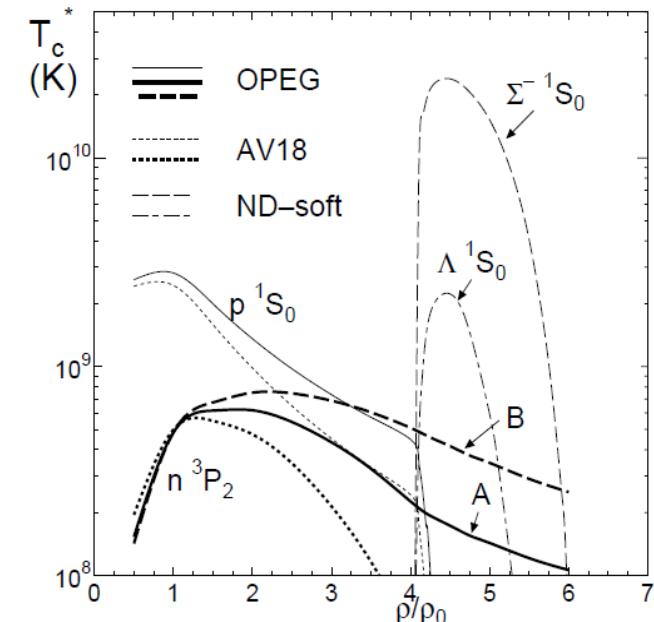
■ Hyperon superfluidity puzzle

- “Exotic” components are necessary to explain fast cooling of some of NSs.
- “Exotic” superfluidity is necessary to forbid too fast cooling.
- Nagara-fit $\Lambda\Lambda$ interaction is too weak for Λ superfluidity.

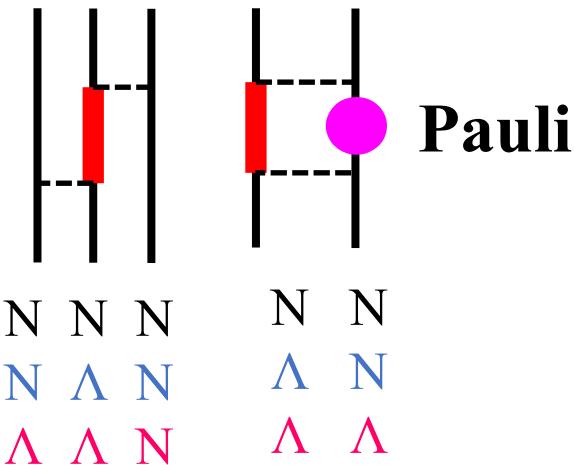
■ Massive Neutron Star Puzzle

- Hyperons should appear in dense matter.
- EOSs w/ Y are too soft to support $2 M_\odot$ NS
- Three-body repulsion involving Y may help to stiffen EOS at high density.

NS session (Gandolfi, Bombaci, Yamamoto, Takatsuka, Togashi, Lonardoni, Muto, Miyatsu, Nakamoto) Kohno ('13) / Myint, Shinmura, Akaishi ('03) / Machleidt.



T. Takatsuka, R. Tamagaki,
PTP 112('04)37



Impact of $\Lambda\Lambda$ interaction

- $\Lambda\Lambda$ interaction is relevant to the existence of H dibaryon.
 - One of the long-standing problem in hadron physics.
- $\Lambda\Lambda$ interaction is important to understand BB interaction models.
 - No pion exchange → middle and core range int. are visible.
 - No quark pauli blocking in the flavor singlet (H) channel.
- $\Lambda\Lambda$ and $\Lambda\Lambda N$ interaction is important to neutron stars.
 - Hypern superfluidity is preferred to explain NS cooling.
 - 3-body force including Y is necessary to support $2 M_{\odot}$ NS.

$V_{\Lambda\Lambda}$ in vacuum & medium $\rightarrow V_{\Lambda\Lambda N}$

*Other observables than double hypernuclei ?
→ $\Lambda\Lambda$ intensity correlation from heavy-ion collisions
Recent data by STAR Collab.
Talk by N. Shah (Adamczyk et al., PRL 114 ('15) 022301.)*

Contents

- **Introduction**
- **$\Lambda\Lambda$ correlation and interaction from heavy-ion collisions**
K. Morita, T. Furumoto, AO, PRC
 - Comparison of correlation function obtained with various $\Lambda\Lambda$ model interactions with STAR data.
 - Expansion effects, Feeddown effects, and Residual source effects
 - Favored $\Lambda\Lambda$ interaction
- **$\Lambda\Lambda$ scattering length: Negative or Positive ?**
 - STAR analysis: $a_0 > 0$, MFO analysis: $a_0 < 0$
 - Re-analysis based on Lednicky-Lyuboshits '81 model
- **Summary**

$\Lambda\Lambda$ correlation and interaction from heavy-ion collisions

*K. Morita, T. Furumoto, AO,
Phys. Rev. C91 (2015), 024916 [arXiv:1408.6682]*

$\Lambda\Lambda$ correlation in HIC

■ Merit of HIC to measure $\Lambda\Lambda$ correlation

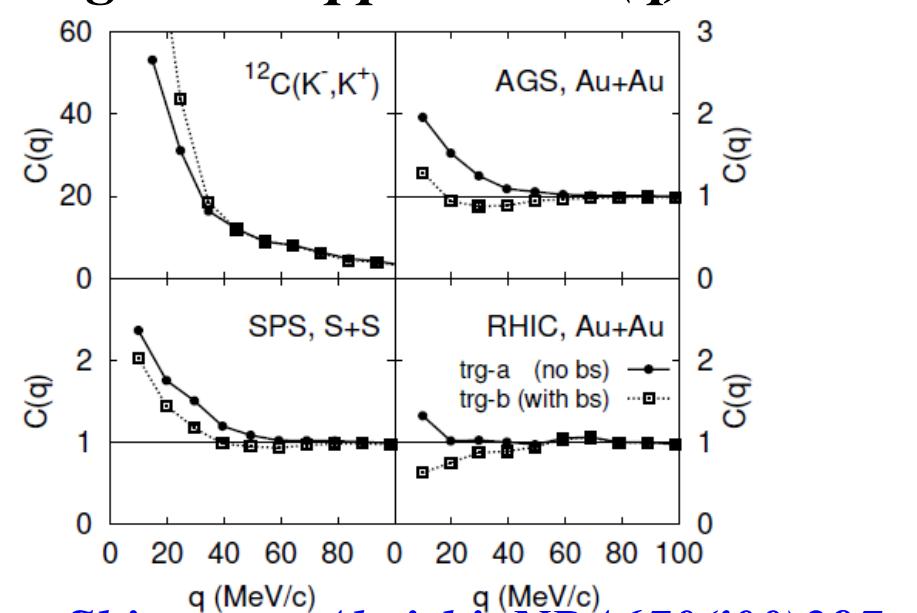
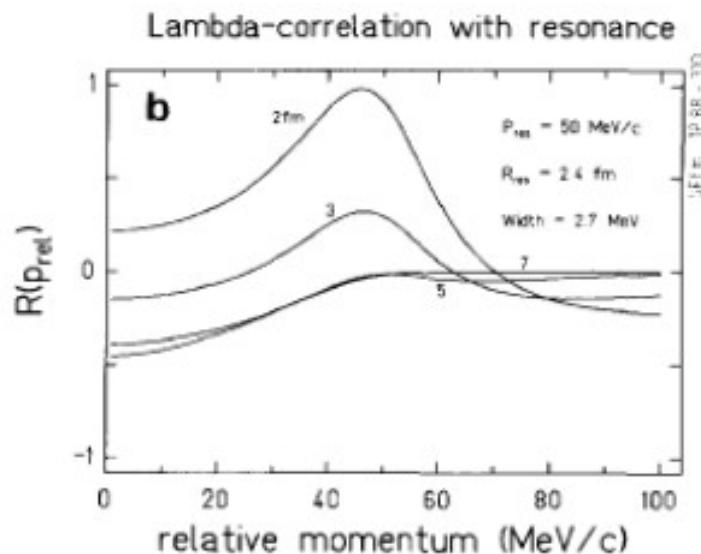
- Source is “Simple and Clean” !
 $T, \mu, \text{flow, size, ...}$ are well-analyzed.

- Nearly Stat. prod.
→ Many exotics will be produced.

Schaffner-Bielich, Mattiello, Sorge ('00), Cho et al.(ExHIC Collab.) ('11)

- Discovery of “H” and/or Constraint on $\Lambda\Lambda$ int.

Bound state exhaust the low q strength → suppressed $C(q)$.



C. Greiner, B. Muller, PLB219('89)199.

AO, Hirata, Nara, Shinmura, Akaishi, NPA670('00)297c

$\Lambda\Lambda$ correlation in HIC and $\Lambda\Lambda$ interaction

- Two particle correlation from chaotic source *Bauer, Gelbke, Pratt ('92)*

$$C_{\Lambda\Lambda}(\mathbf{q}) = \frac{\int dx_1 dx_2 S(x_1, p+q) S(x_2, p-q) |\psi^{(-)}(x_{12}, q)|^2}{\int dx_1 dx_2 S(x_1, p+q) S(x_2, p-q)}$$

- Static spherical source, s-wave only s-wave w.f. enh.

$$C_{\Lambda\Lambda}(\mathbf{q}) \simeq 1 - \frac{1}{2} \exp(-4q^2 R^2) + \frac{1}{2} \left[\int dr S_{12}(r) \left(|\chi_0(r)|^2 - |j_0(qr)|^2 \right) \right]$$

HBT

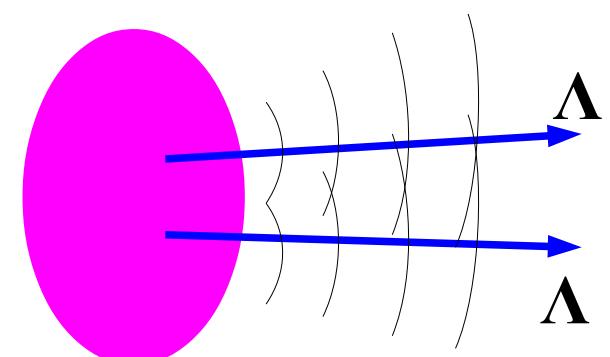
$\mathbf{q} = (\mathbf{p}_1 - \mathbf{p}_2)/2$, χ_0 : s-wave wf, $S_{12}(x) = r^2 \exp(-r^2/4R^2)/2R^3\sqrt{\pi}$

- HBT term: Suppression due to anti-symmetrization for Fermions

$$0 \times 3/4 \text{ (triplet)} + 2 \times 1/4 \text{ (singlet)} = 1/2$$

*Hanbury Brown, Twiss ('56),
Goldhaber, Goldhaber, Lee, Pais ('60)*

- Enhancement of wf by $\Lambda\Lambda$ attraction



Correlation function and scattering length

- Low E scatt. is mainly determined by the scattering length (a_0)

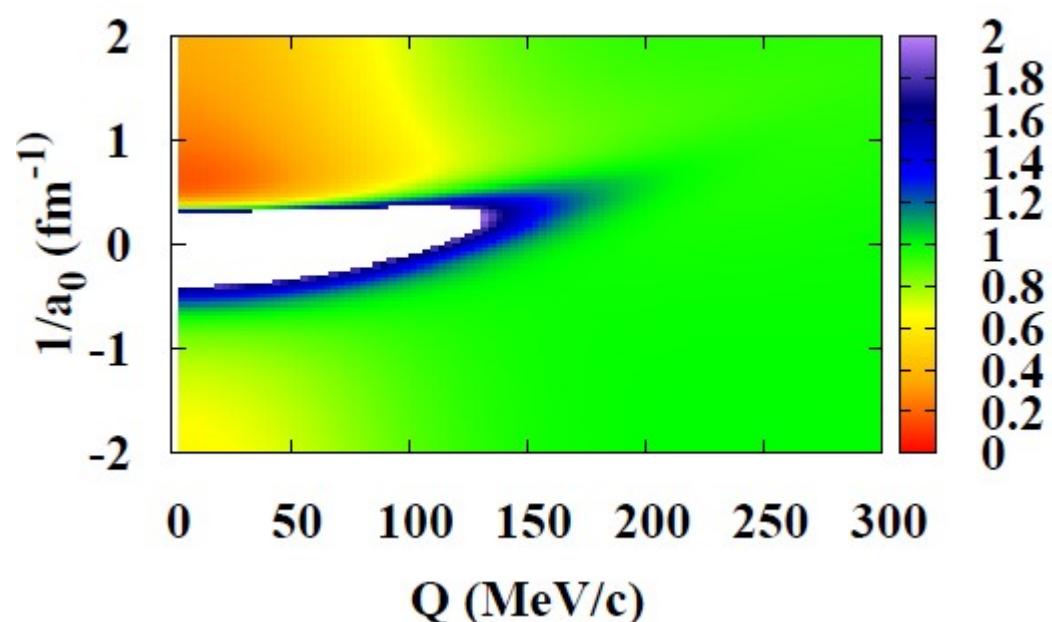
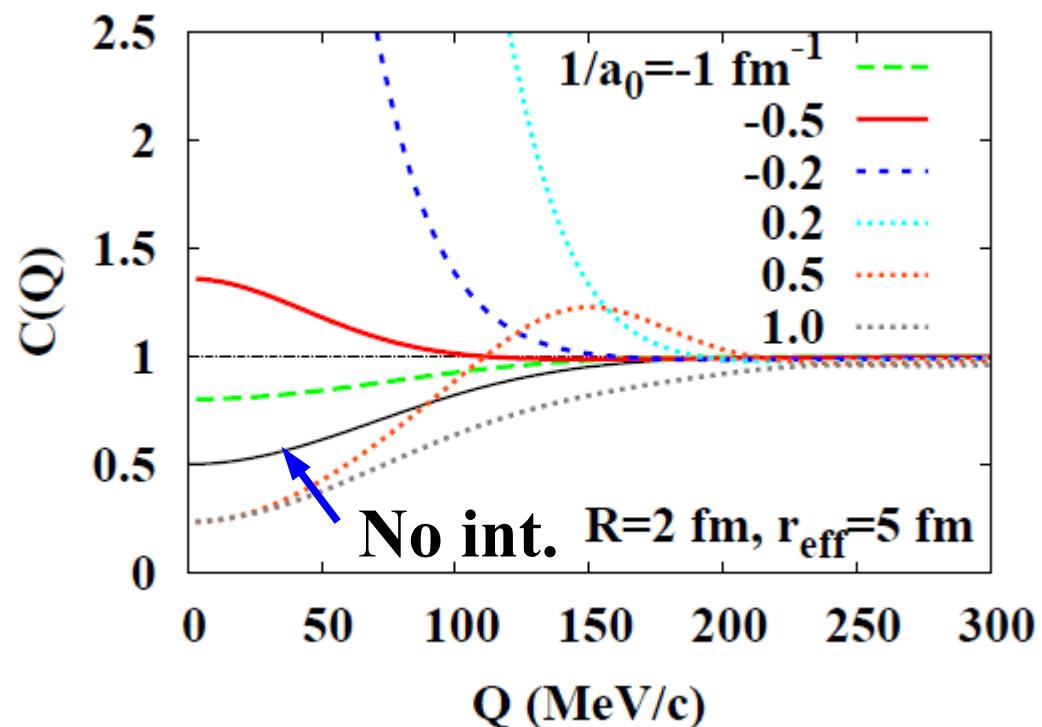
$$q \cot \delta = -\frac{1}{a_0} + r_{\text{eff}} q^2/2 + O(q^4) \quad \text{Nuclear phys. convention}$$

- Negative a_0 : enhanced correlation

- Positive a_0 : suppressed correlation

w/ bound state \rightarrow scatt. wf has a node

w/o bound state \rightarrow repulsion suppress wf



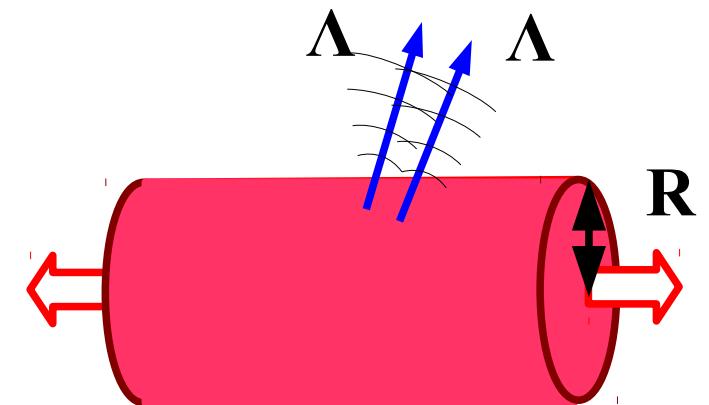
Source models

■ Static & Spherical source

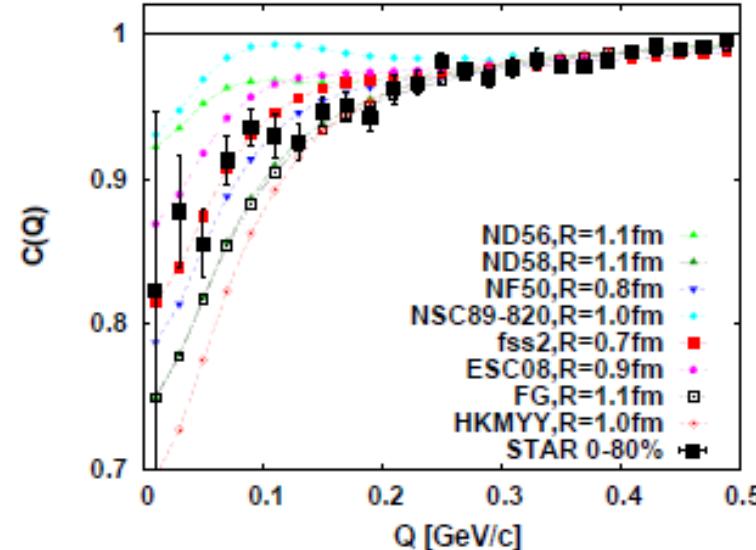
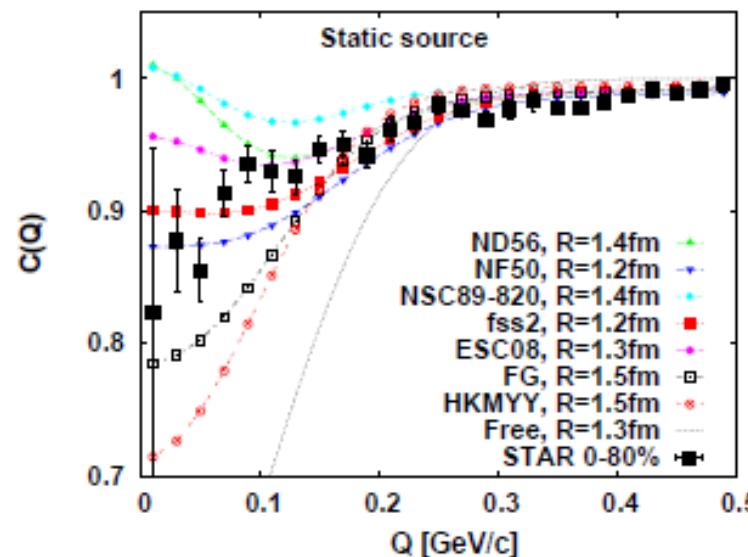
- Source size $R = (1-1.5)$ fm, $\chi^2/\text{DOF} \sim 2$

■ Cylindrical source

- Bjorken expansion in long. direction
- Transverse flow from pT spectrum
- Source size $R = (0.7-1.1)$ fm, $\chi^2/\text{DOF} \sim 1.5$



$$S_{\text{cyl}}(x, k) = \frac{m_T \cosh(y - Y_L)}{(2\pi)^3 \sqrt{2\pi(\Delta\tau)^2}} n_f(u \cdot k, T) \exp \left[-\frac{(\tau - \tau_0)}{2(\Delta\tau)^2} - \frac{x^2 + y^2}{2R^2} \right]$$



Additional Source

■ Feed down effects

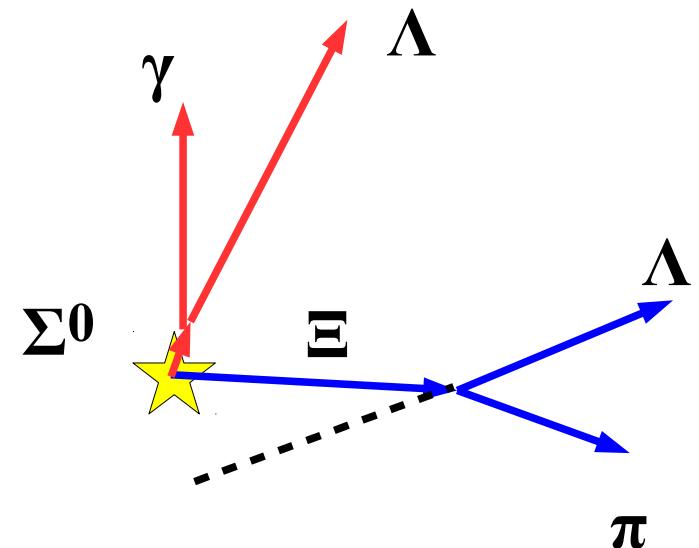
$$C_{\text{corr}}(Q) = 1 + \lambda(C_{\text{bare}}(Q) - 1)$$

λ = Purity of $\Lambda\Lambda$ pair

- Short-lived Y^* \rightarrow mod. of source fn.
- $\Xi \rightarrow \Lambda\pi$ can be excluded ($c\tau=8.71$ cm)
- $\Sigma^0 \rightarrow \Lambda\gamma$ is difficult to reject
- Data based purity $\lambda=(0.67)^2$
 $\Sigma^0/\Lambda=0.278$ (p+Be, 28.5 GeV/c) *Sullivan et al. ('87)*
 $\Xi/\Lambda = 15\%$ (RHIC)

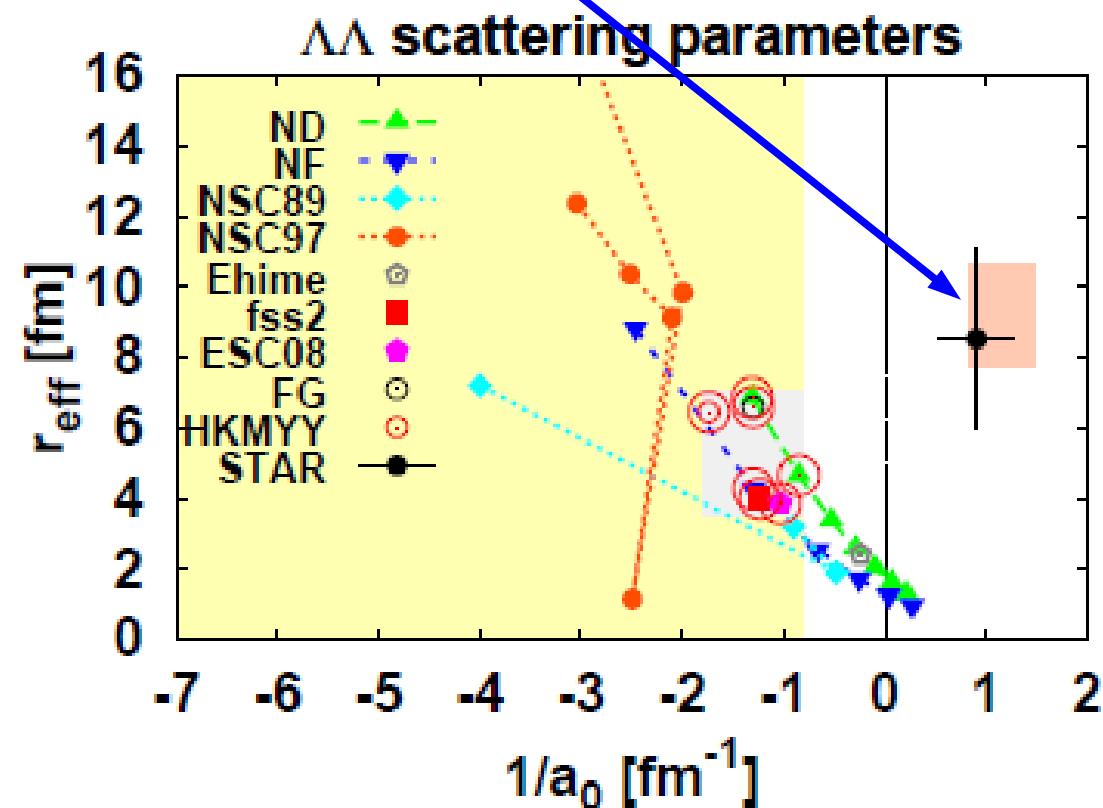
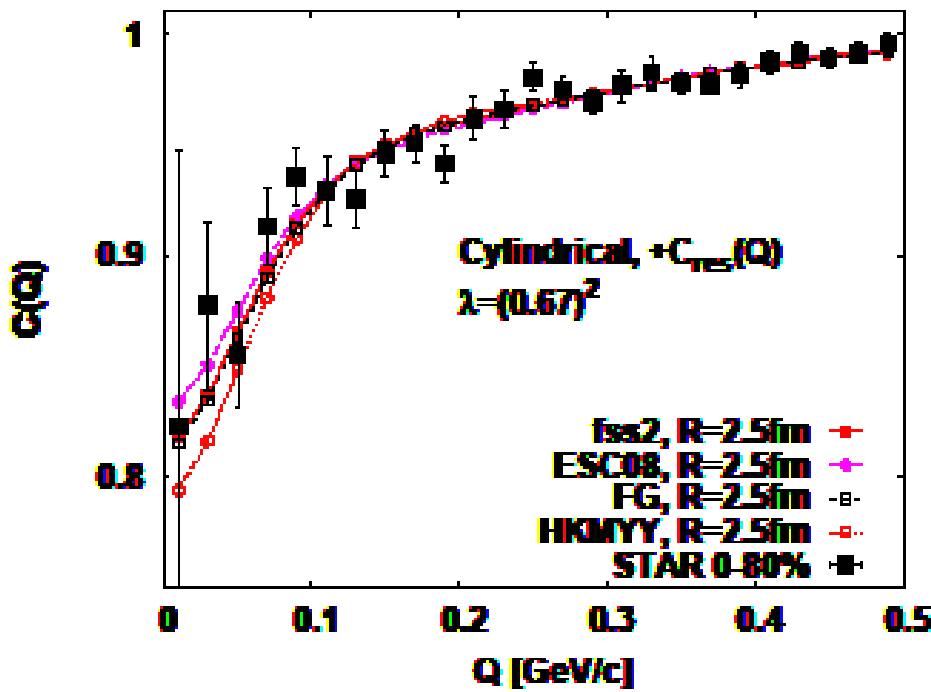
■ “Residual” source

- High-momentum tail $\rightarrow R_{\text{res}} \sim 0.5$ fm (STAR collab.)



$\Lambda\bar{\Lambda}$ interaction from $\Lambda\bar{\Lambda}$ correlation at RHIC

- HBT + $\Lambda\bar{\Lambda}$ int. + Expansion + Feed down + “Residual” source
 - $1/a_0 < -0.8 \text{ fm}^{-1}$
 - Source size dependence is small, $\chi^2/\text{DOF} \sim 1$
- Question: STAR analysis $a_0 = 1.10 \pm 0.37 \text{ fm} > 0$



*K.Morita, T.Furumoto, AO, PRC91('15)024916 [arXiv:1408.6682]
Data: Adamczyk et al. (STAR Collaboration), PRL 114 ('15) 022301.*

$\Lambda\Lambda$ scattering length Positive or Negative ?

STAR collab. analysis

■ Lednicky-Lyuboshits (LL '81) model

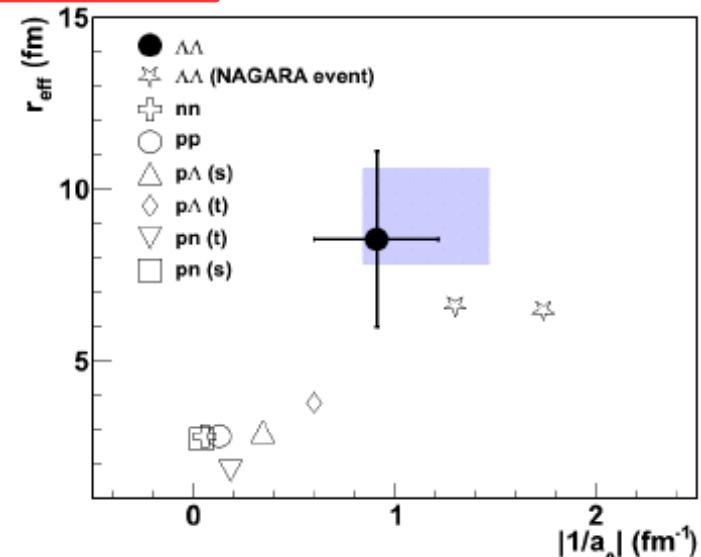
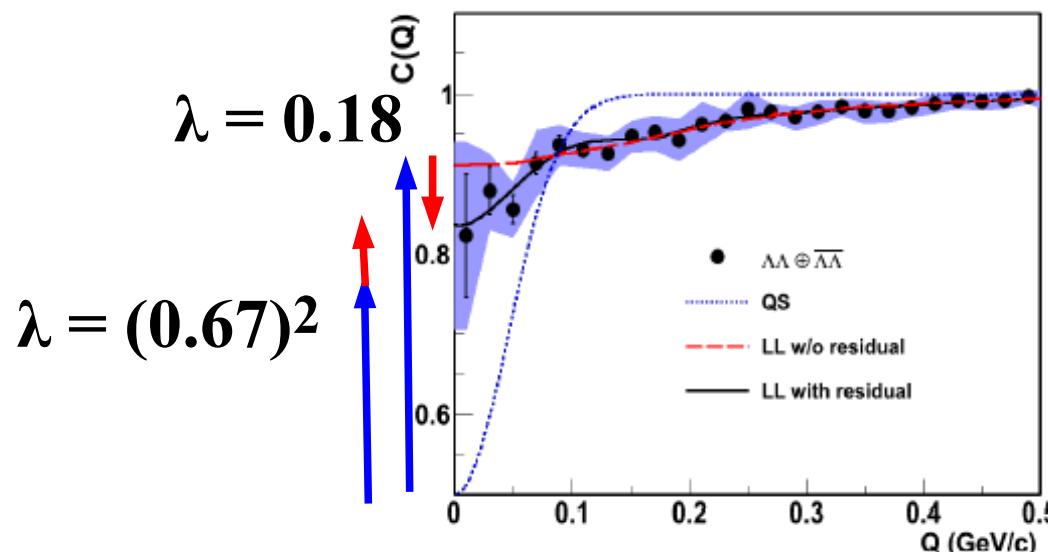
R. Lednicky, V. L. Lyuboshits, Sov.J.Nucl.Phys.35('82)770; Yad.Fiz.35 ('81) 1316.

- Analytic model including a_0 and r_{eff} .

- Intercept (pair purity) parameter is treated as a free parameter.

- STAR Collab. analysis implies positive scatt. length.

$$a_0 = 1.10 \text{ fm} , \quad r_{\text{eff}} = 8.52 \text{ fm} , \quad \lambda = 0.18$$



Adamczyk et al. (STAR Collaboration), PRL 114 ('15) 022301.

Which is correct ? Positive a_0 or Negative a_0 ?

Lednicky-Lyuboshits (LL '81) model

■ Lednicky-Lyuboshits (LL '81) model

R. Lednicky, V. L. Lyuboshits, Sov.J.Nucl.Phys.35('82)770; Yad.Fiz.35 ('81) 1316.

● Analytic model including FSI effects

$$C(Q) = N \left[1 + \lambda \left(-\exp(-R^2 Q^2)/2 + C_{\text{int}}(Q) - 1 \right) + a_{\text{res}} \exp(-R_{\text{res}}^2 Q^2) \right]$$

$$C_{\text{int}}(Q) = 1 + \frac{|f(k)|^2}{4R^2} \left(1 - \frac{r_{\text{eff}}}{2\sqrt{\pi}R} \right) + \frac{\text{Re}f(k)}{\sqrt{\pi}R} F_1(QR) - \frac{\text{Im}f(k)}{2R} F_2(QR)$$

$$f(k) = (-1/a_0 + r_{\text{eff}}k^2/2 - ik)^{-1} \quad (k = Q/2) \quad \text{scattering amp.}$$

$$F_1(z) = \int 0^z e^{x^2 - z^2}/z, \quad F_2(z) = (1 - e^{-z^2})/z$$

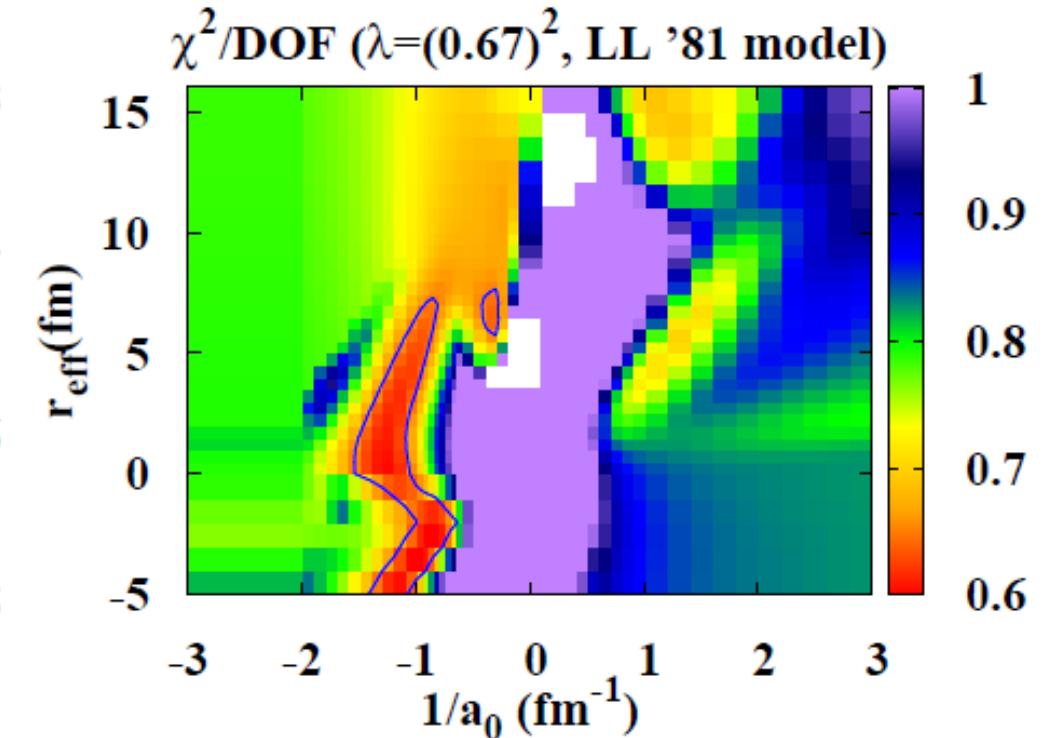
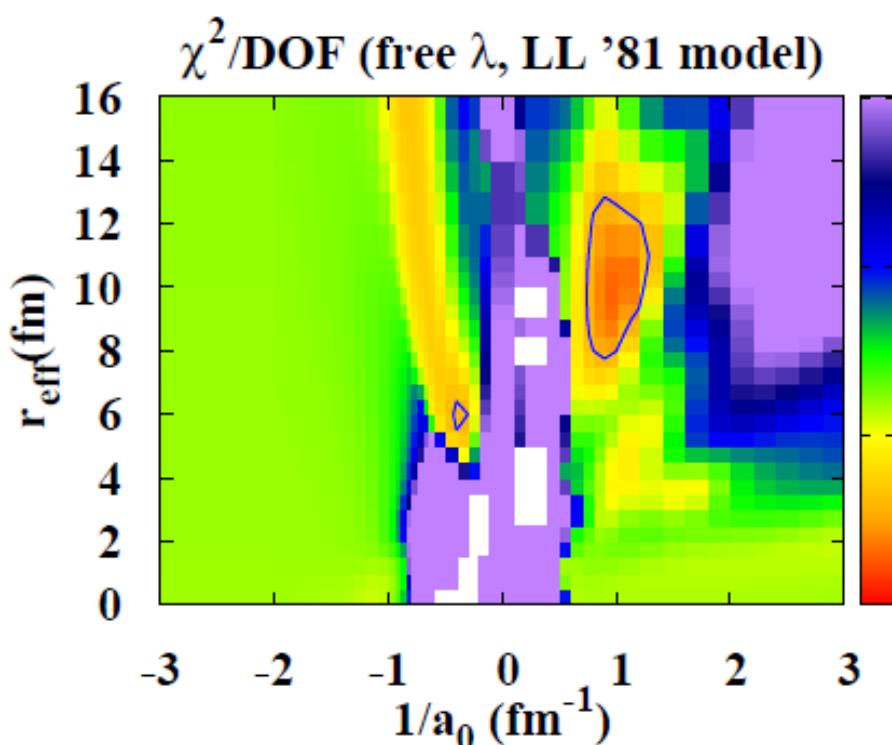
$$k \cot \delta = -\frac{1}{a_0} + \frac{1}{2}r_{\text{eff}}k^2 + \mathcal{O}(k^4) \quad \text{phase shift}$$

Let us examine LL '81 model results with fixed λ !

$a_0(\Lambda)$: Positive or Negative ?

■ Key parameter = intercept parameter λ

- λ = free parameter $\sim 0.18 \rightarrow$ positive a_0 at χ^2 min.
($\chi^2/\text{DOF} \sim 0.56$)
- Measured Σ^0 / Λ ratio
 $\lambda = (\Lambda / (\Lambda + \Sigma^0))^2 \sim (0.67)^2 \rightarrow$ negative a_0 at χ^2 min.
($\chi^2/\text{DOF} \sim 0.65$)

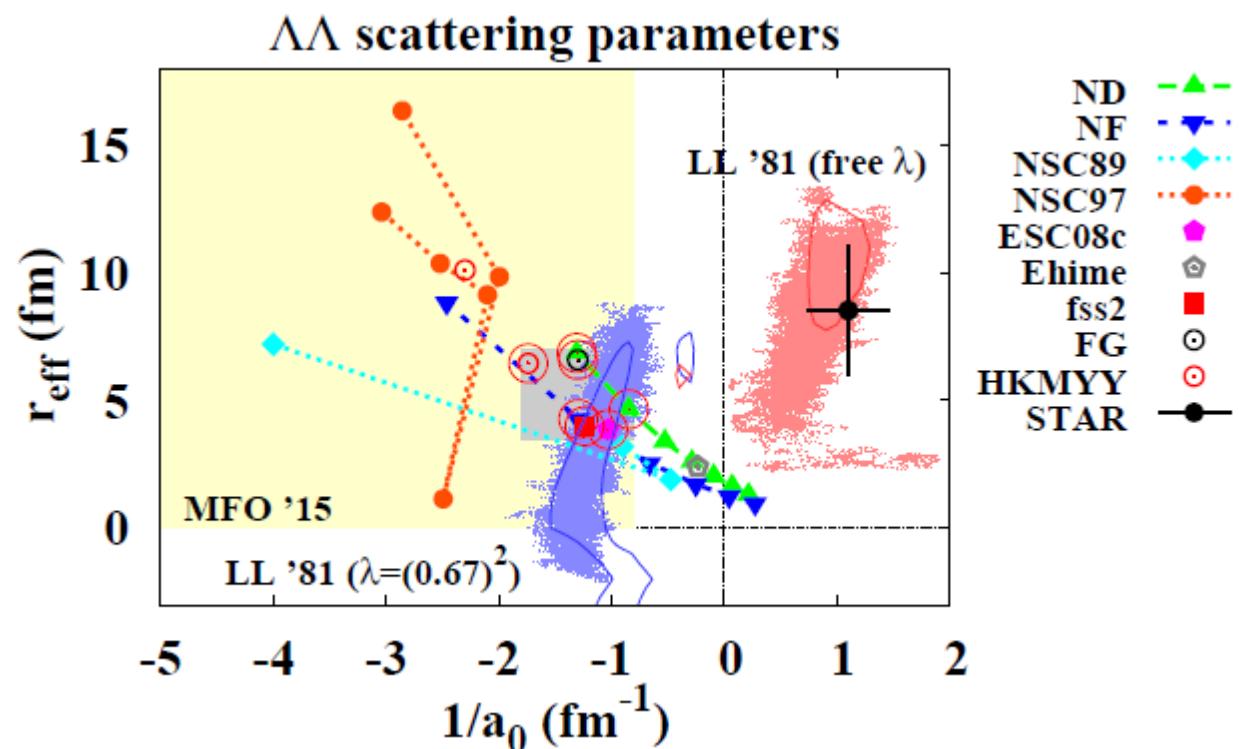


$a_0 (\Lambda\Lambda)$: Positive or Negative ?

■ χ^2 analysis

- Gradient method (gnuplot) & Markov Chain MC (MCMC) in LL '81 model
- Free λ result ~ STAR ('15) analysis : $a_0 > 0$
- Fixed λ result ~ Our analysis : $a_0 < 0$

*Difference btw
STAR collab. analysis
and our analysis
comes from λ*



What is the origin of the long tail ?

- Do we have a physical origin ?
- Two source model + LL '81 model

$$S_{12}(\mathbf{x}) = \frac{w}{(2R_1\sqrt{\pi})^3} \exp\left(-\frac{x^2}{4R_1^2}\right) + \frac{1-w}{(2R_2\sqrt{\pi})^3} \exp\left(-\frac{x^2}{4R_2^2}\right)$$

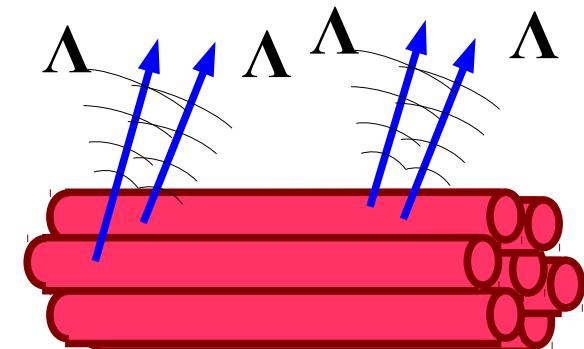
- Fix Λ Λ interaction, and obtain R_1 and R_2 .

$R_1 = (0.65-1.30)$ fm

$R_2 = (0.33-0.54)$ fm

for fss2, ESC08, FG, HKMYY interactions.

$\chi^2/\text{DOF} = (0.6-0.65)$



- Λ might be produced from small tubes.
 - FSI with hot medium ($\sim \pi$ gas) is small for Λ

Summary

- **$\Lambda\Lambda$ intensity correlation in high-energy heavy-ion collisions has sensitivity to $\Lambda\Lambda$ interaction.**
- Our analysis of the STAR data implies that the favored $\Lambda\Lambda$ interaction has negative scattering length ($1/a_0 < -0.8 \text{ fm}^{-1}$, $\delta > 0$), which is consistent with the Nagara event analysis.
 - Anti-symm. + $\Lambda\Lambda$ interaction + Expansion + Feeddown + Residual
- Difference between the STAR collab. analysis and ours lies in the assumption on the pair purity (chaoticity, intercept) parameter.
 - $\Sigma^0/\Lambda \sim 0.67$: p+Be, consistent with stat. model ($T \sim 170 \text{ MeV}$)
- Further studies are necessary to pin-down $\Lambda\Lambda$ interaction.
 - Higher precision data are expected, Σ^0 detection is desired, Comparison with $A(K^-, K^+)\Lambda\Lambda$ reaction **C.J.Yoon et al.('07)/J-PARC**
- We can access other hh interactions using correlations.
 - ΩN interaction from ΩN correlation **K. Morita's talk**

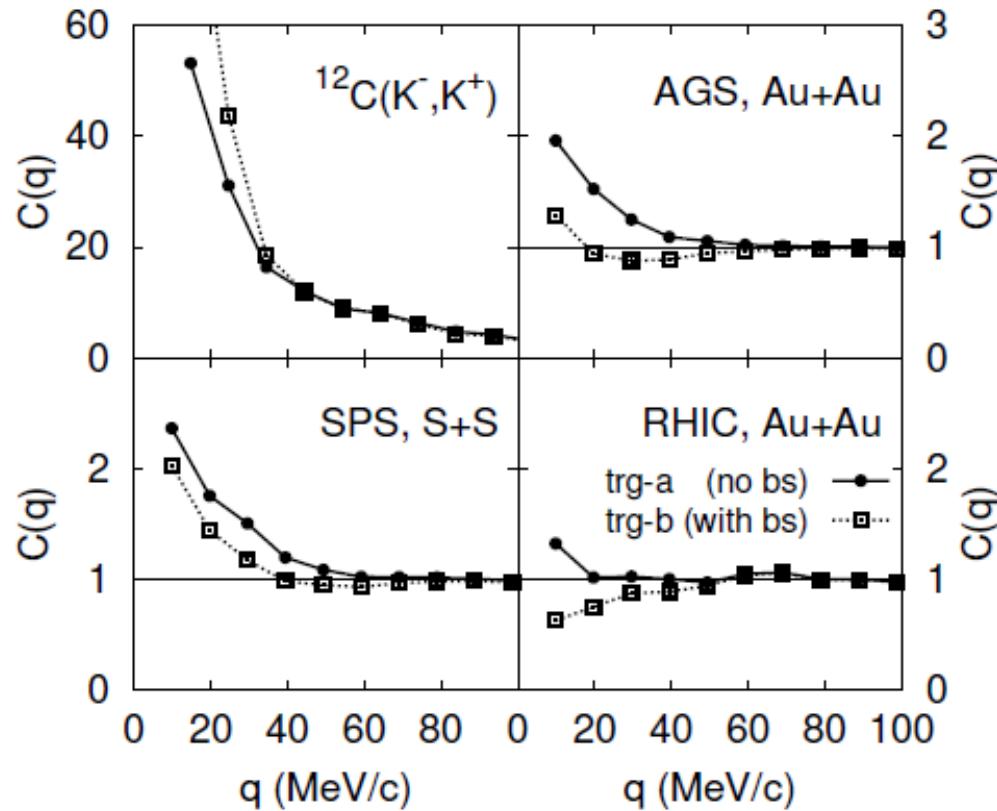
Thank you !

Nuclear Physics, Compact Stars, and Compact-star Mergers 2016 (NPCSM 2016)

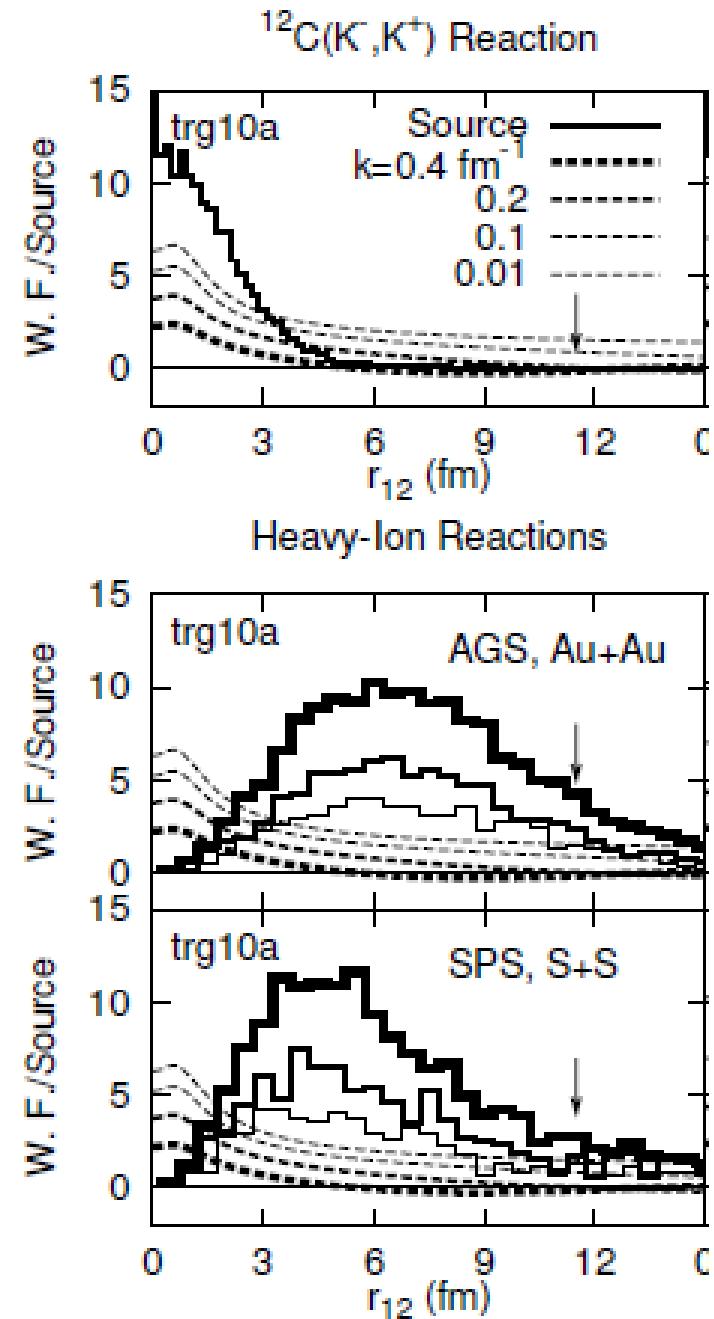
- Dates :
Oct.17 (Mon.)-Nov.18 (Fri.), 2016 (2016 Fall, 5 weeks).
- Organizers:
**Ohnishi (YITP, co-chair), M. Shibata(YITP, co-chair),
Y. Suwa(YITP, LOC), K. Morita(YITP, LOC), H.Tamura (Tohoku
U.), T.Harada (Osaka EC), H.Nakada(Chiba U.), K. Iida (Kochi U.),
Sekiguchi (Toho U.), K. Sumiyoshi (Numazu), Ioka(KEK),
Kotake(Fukuoka U.), Sotani(NAOJ), Wanajo (RIKEN)**

Previous Work (before RHIC & Nagaoka)

- Hadronic transport (JAM)
 - + Two Range Gaussian $V_{\Lambda\Lambda}$
 - w/ bound state \rightarrow w.f. node suppresses $C(q)$

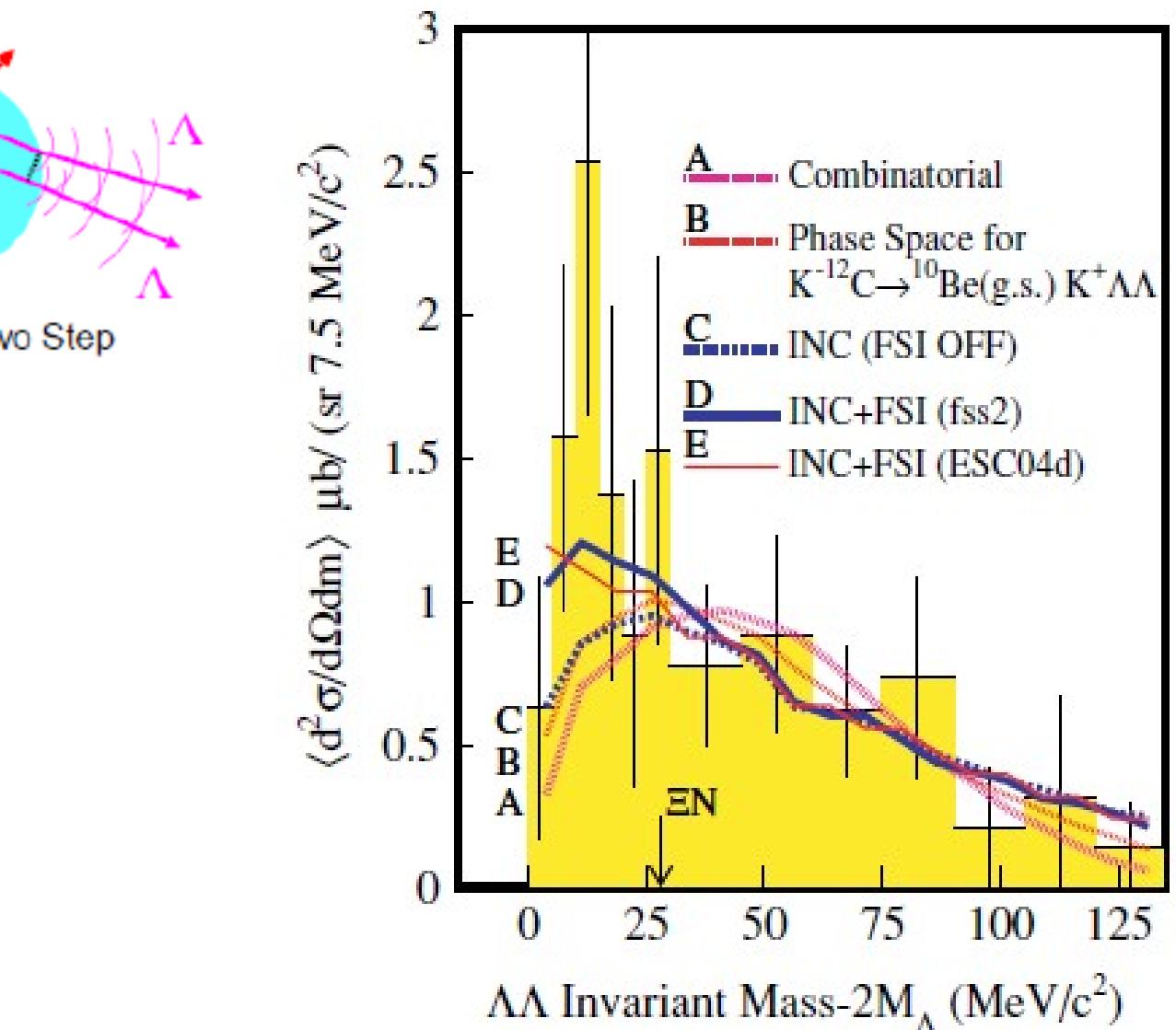
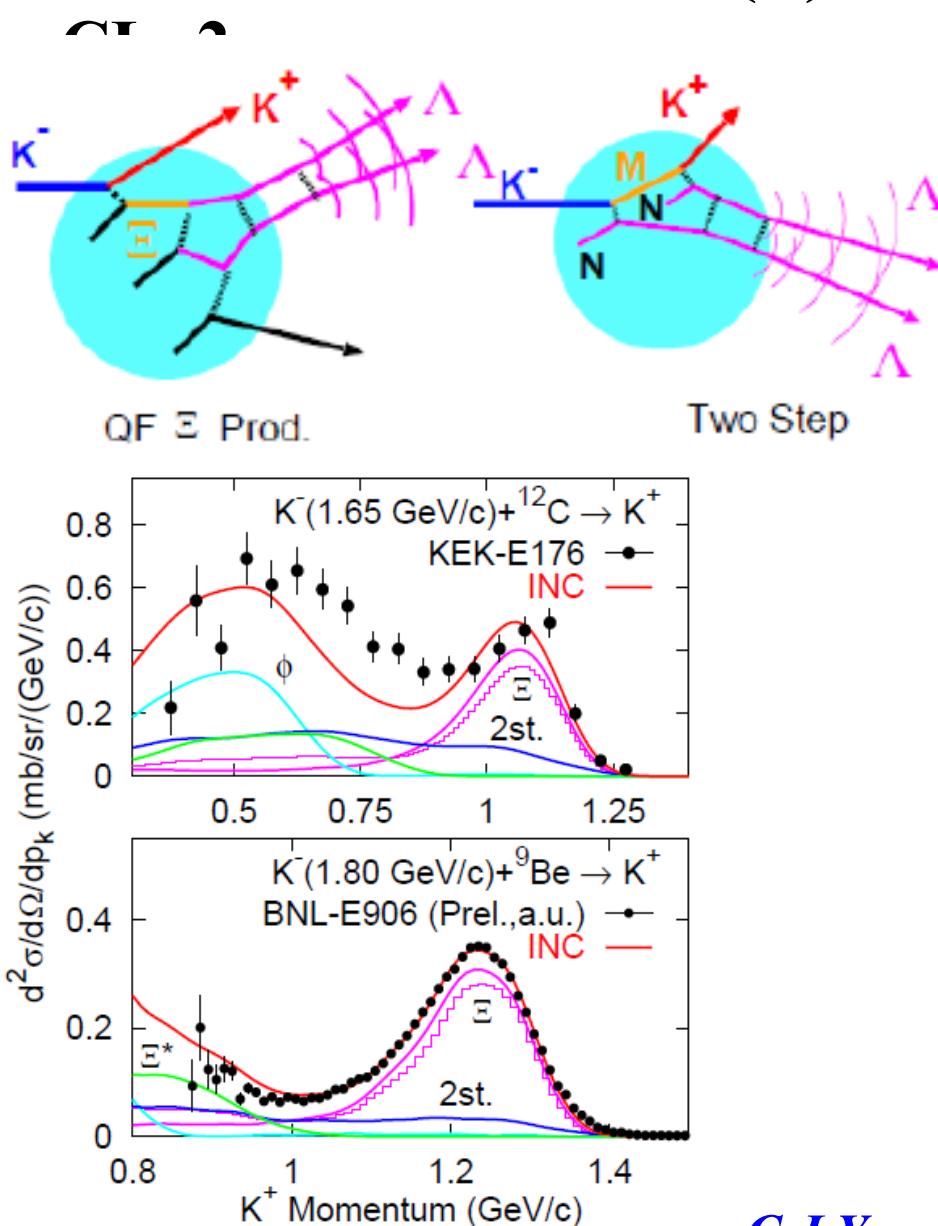


*AO, Hirata, Nara, Shinmura, Akaishi, NPA670('00)297c
[arXiv:nucl-th/9903021]; SNP2000 proc. p175.
JAM: Nara, Otuka, AO, Niita, Chiba, PRC61 ('00), 024901.*



$\Lambda\bar{\Lambda}$ correlation from ($K^-, K^+ \Lambda\bar{\Lambda}$) reaction

- Enhancement at $\sim 2 M(\Lambda) + 10$ MeV,



C.J.Yoon, ..., (KEK-E522), AO, PRC75 (2007) 022201(R)

Exotics from Heavy Ion Collisions

■ High-Energy Heavy-Ion Collisions

- Too complex → Statistical → Simple and Clean !
- High T & Large volume → Abundant hadrons
- RHIC & LHC → Nearly 4π detector / Vertex detector

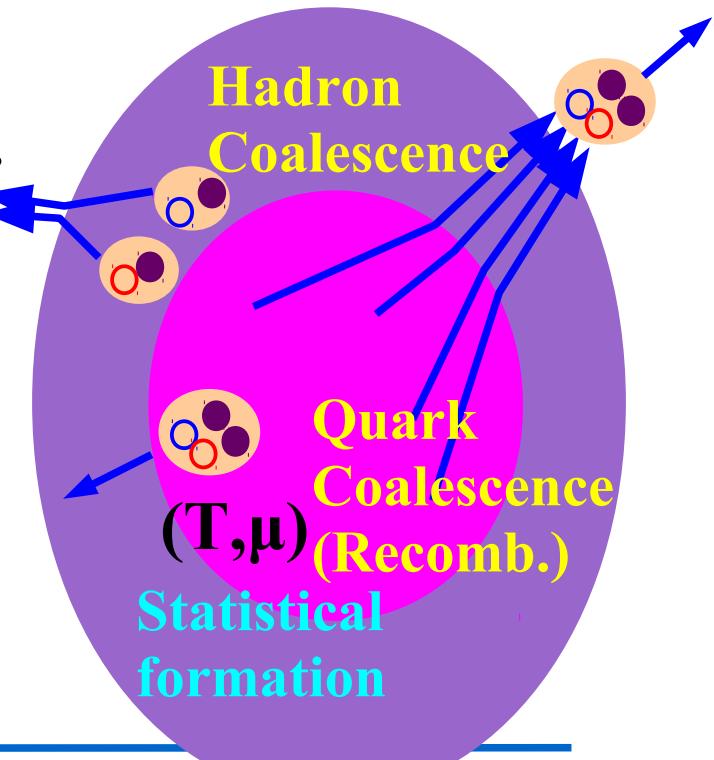
→ Let's regard RHC & LHC as Exotic Hadron Factories

■ What can we learn ? Existence, Mass, Width, Size, Interaction, ..

- Formation mechanism of hadrons
= Statistical , Coalescence, Fragmentation.
→ Yields are sensitive to hadron size
in Coal.

Chen, Greco, Ko, Lee, Liu, ('04)

- Correlation func. $\sim \int \text{Source} \times |\text{w.f.}|^2$
→ Once we know source,
Corr. Func. is sensitive to w.f.
and pairwise interaction.



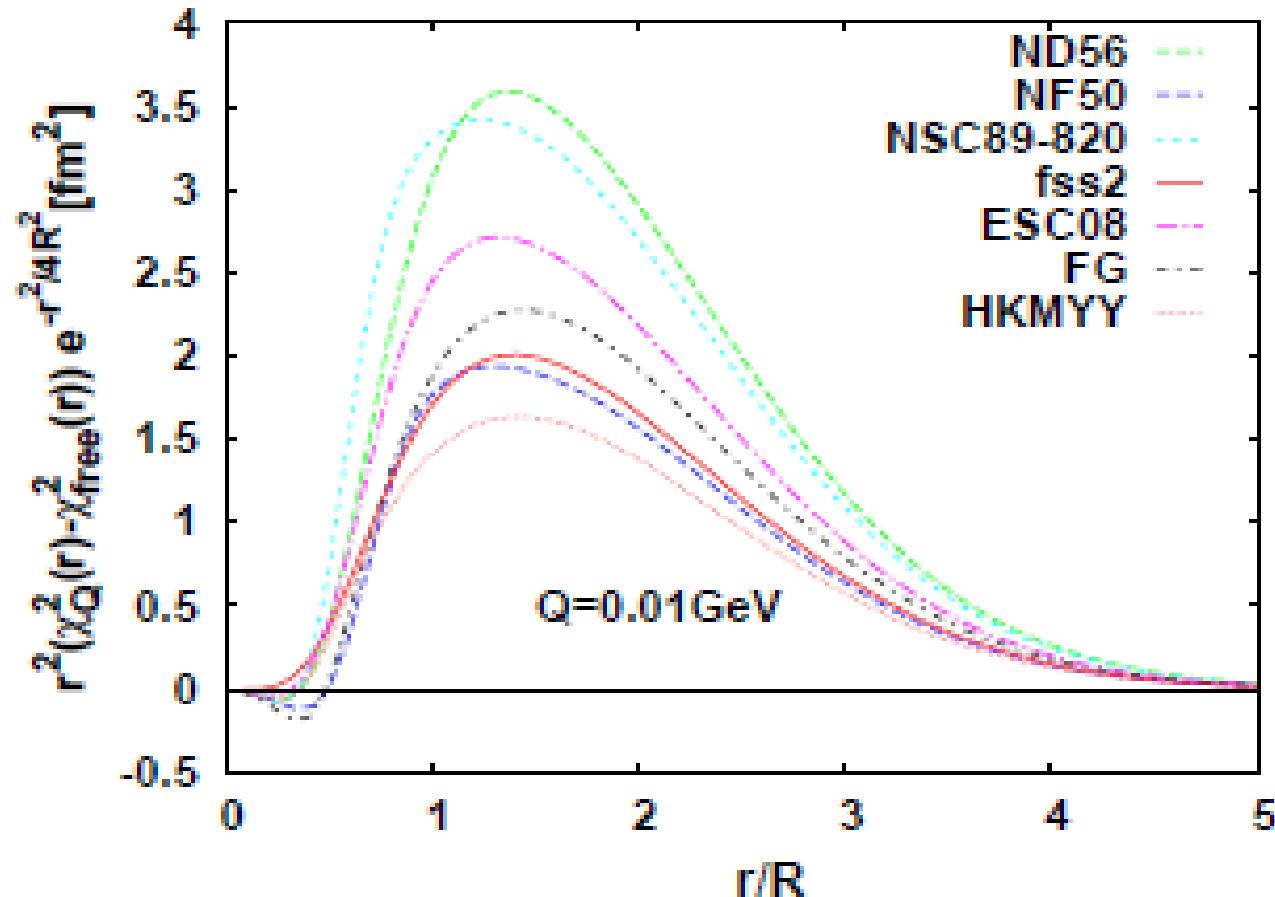
Bauer, Gelbke, Pratt ('92); Lednicky ('09).

A. Ohnishi @ Mol. 2015.09, Sep. 21, 2015, Kyoto, Japan

$\Lambda\Lambda$ wave function

■ Correlation function

$$C_{\Lambda\Lambda}(q) \approx 1 - \frac{1}{2} \exp(-4q^2 R^2) + \frac{1}{2} \int dr S_{12}(r) (|\chi_0(r)|^2 - |j_0(qr)|^2)$$



$\Lambda\Lambda$ interaction models

■ Boson exchange potential

Nijmegen potentials (ND, NF, NSC89, NSC97, ESC08)
Nagels+('77, '79), Maessen+('89), Rijken+('99,'10)
Ehime Ueda et al. ('98)

■ Quark model interaction: fss2 Fujiwara et al.('07)

■ Tuned potential to Nagara

Filikhin, Gal ('02)	$a_0 = -0.77$ fm, $r_{eff} = 6.59$ fm
Hiyama et al. ('02)(HKMYY)	$a_0 = -0.575$ fm, $r_{eff} = 6.45$ fm
Hiyama et al. ('10)(HKMYY)	$a_0 = -0.44$ fm, $r_{eff} = 10.1$ fm

■ Tuned potential to $A(K^-, K^+)\Lambda\Lambda$

Gasparayan et al. ('12)(GHH) $a_0 = -1.2$ fm
A.M.Gasparyan, J. Haidenbauer, C. Hanhart. PRC85('12)015204.

Nagara event

■ $_{\Lambda\Lambda}^6\text{He}$ hypernuclei

Takahashi et al., PRL87('01)212502

(KEK-E373 experiment)

Lambpha

$$m(_{\Lambda\Lambda}^6\text{He}) = 5951.82 \pm 0.54 \text{ MeV}$$

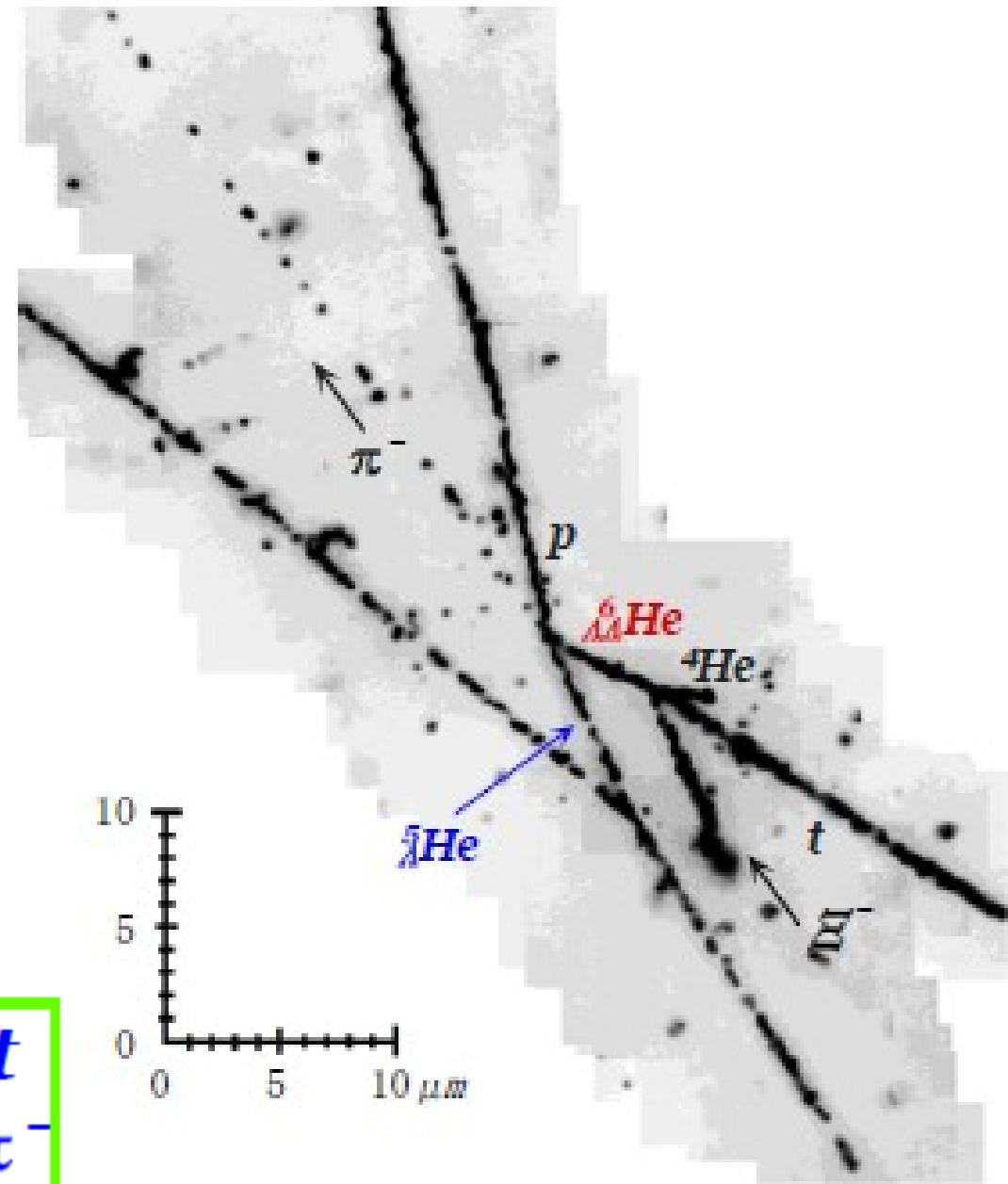
$$B_{\Lambda\Lambda} = 7.25 \pm 0.19^{+0.18}_{-0.11} \text{ MeV}$$

$$\Delta B_{\Lambda\Lambda} = 1.01 \pm 0.20^{+0.18}_{-0.11} \text{ MeV}$$

$$(\text{assumed } B_{\Xi^-} = 0.13 \text{ MeV})$$

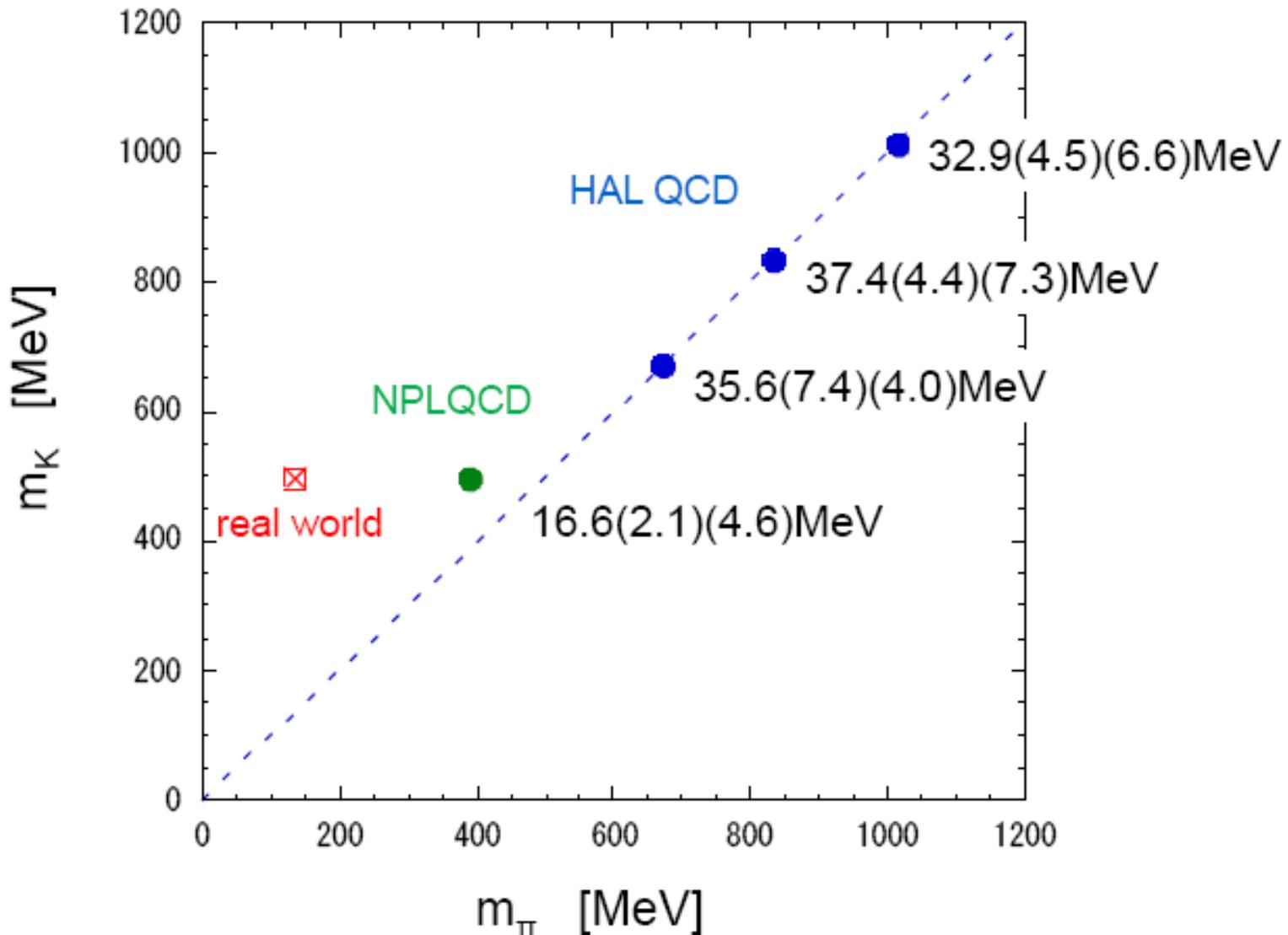
$$\rightarrow B_{\Lambda\Lambda} = 6.91 \text{ MeV}$$

(PDG modified(updated)
 Ξ^- mass)



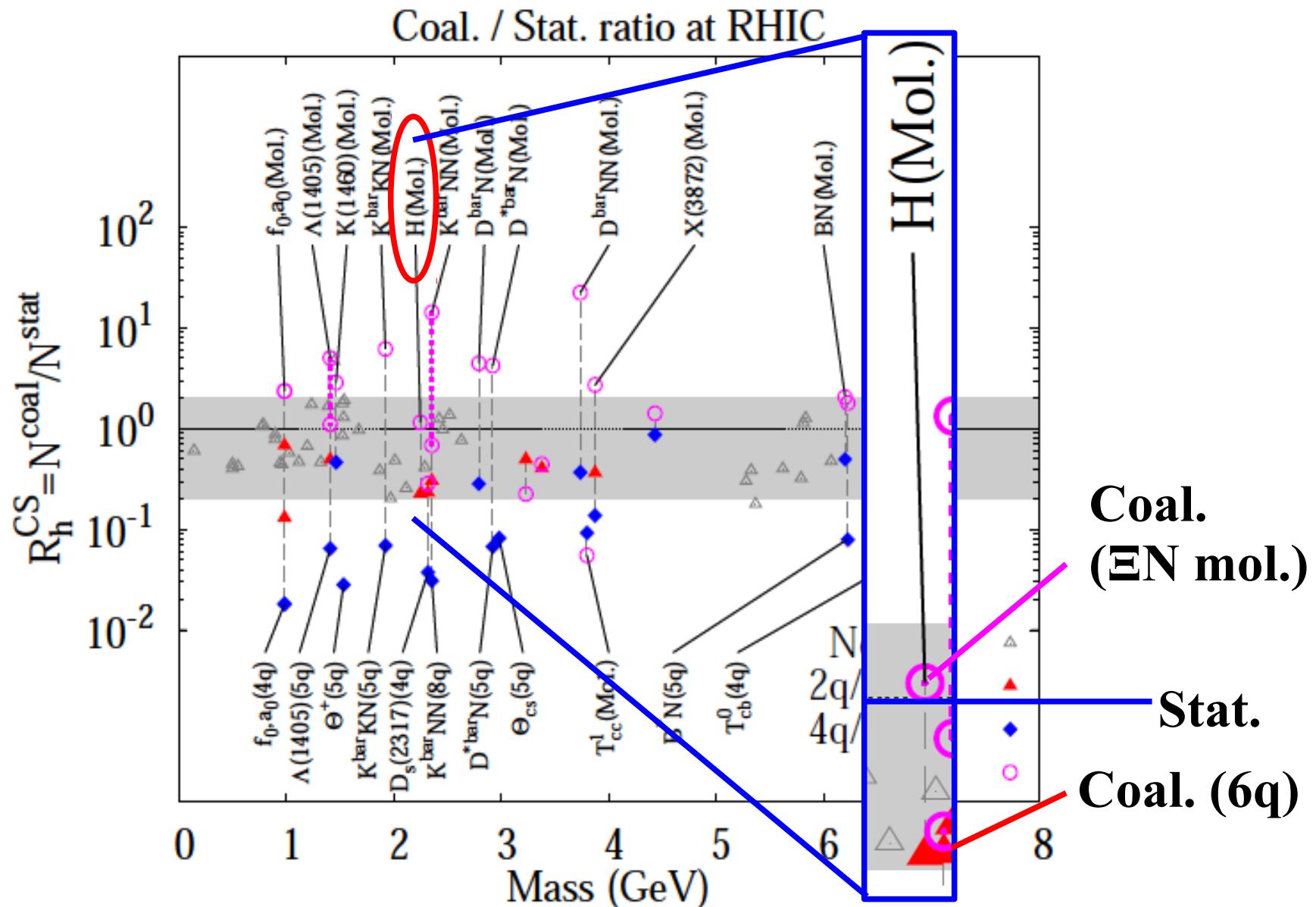
Lattice QCD predicts bound “H”

- “H” bounds with heavy π ($M_\pi > 400$ MeV)



NPLQCD Collab., PRL 106 (2011) 162001; HAL QCD Collab., PRL 106 (2011) 162002

Exotics from Heavy-Ion Collisions



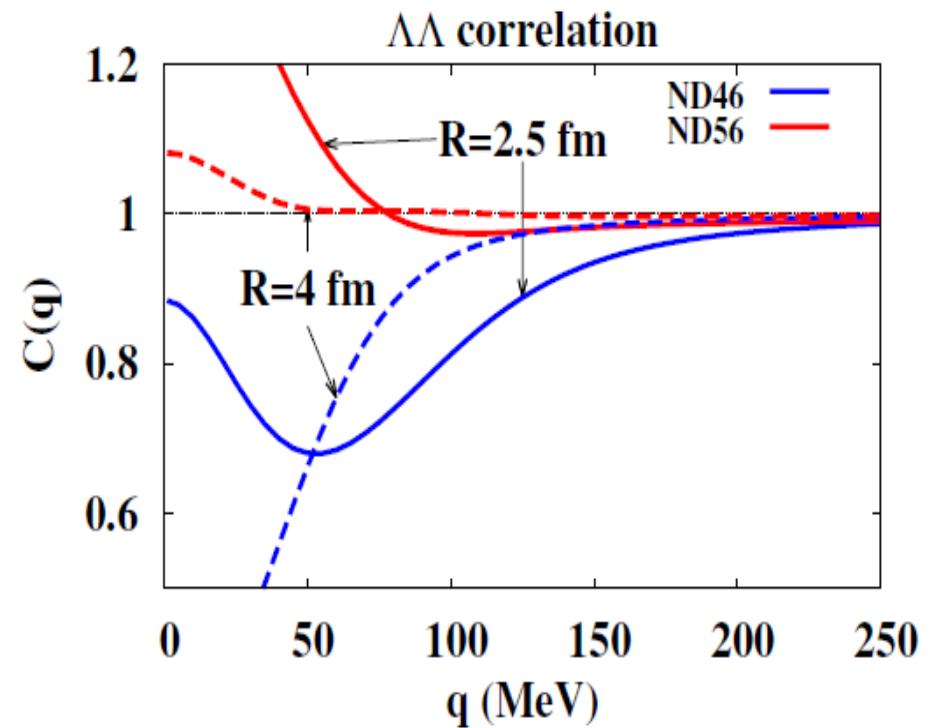
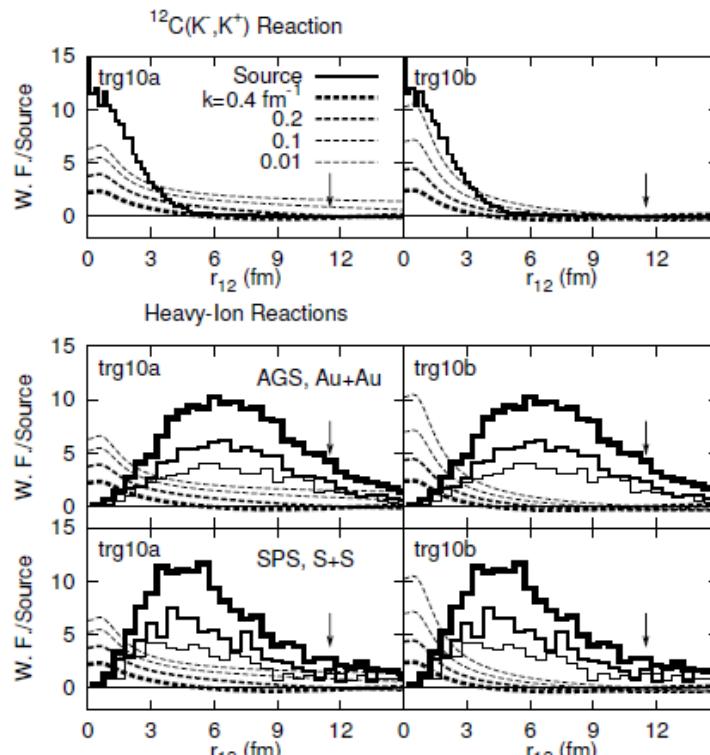
*Cho, Furumoto, Hyodo, Jido, Ko, Lee, Nielsen, AO, Sekihara, Yasui, Yazaki
(ExHIC Collab.), PRL('11)212001; arXiv:t:1107.1302*

Fate of the prediction

Conjecture in 2000

Suppressed $\Lambda\Lambda$ correlation may suggest the existence of bound H

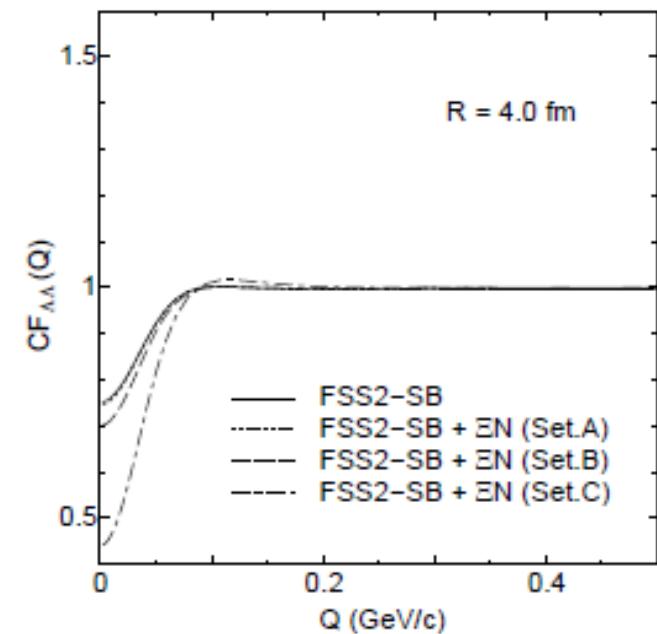
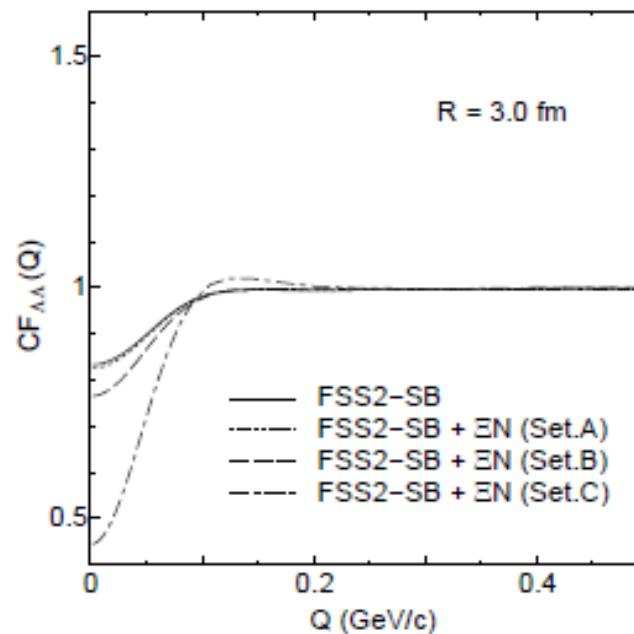
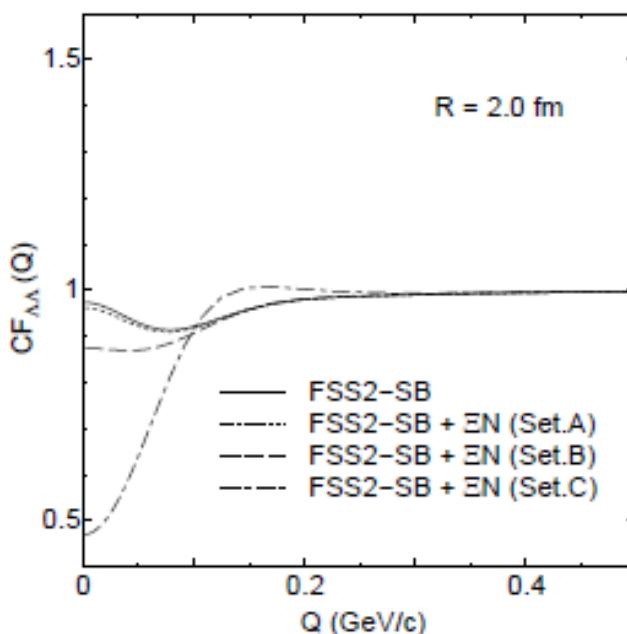
- Bound H → Node in scattering $\Lambda\Lambda$ wf → suppressed correlation
*AO, Hirata, Nara, Shinmura, Akaishi, NPA670('00)297c
[arXiv:nucl-th/9903021]; SNP2000 proc. p175.*
- When the source (homogeneity) size is small, we find a dip with/without bound state.



Source size dependence

- Larger size → Smaller Q region
- No dip structure for larger size.
(Anti-symmetrization effects > Interaction effects)
→ Sensitive only to the scattering length.

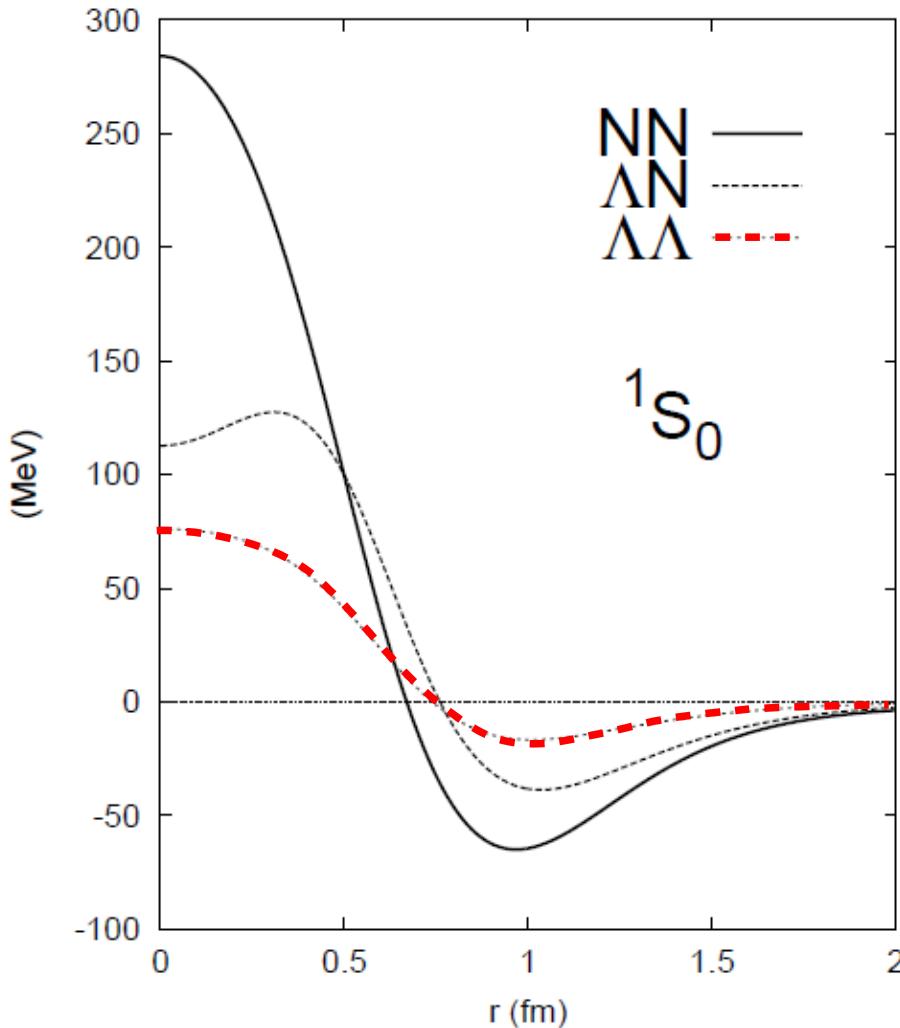
$$C(Q \rightarrow 0) \approx \frac{1}{2} - \frac{2}{\sqrt{\pi}} \frac{a_0}{R} + \left(\frac{a_0}{R} \right)^2 \quad (\text{if "Interaction Range" } \ll R)$$



AO, Furumoto, in prep.

$\Lambda\Lambda$ potential

fss2 Phase shift equivalent potential



fss2

- $a_0 = -0.82$ fm, $r_{\text{eff}} = 4.1$ fm

Nagara fit

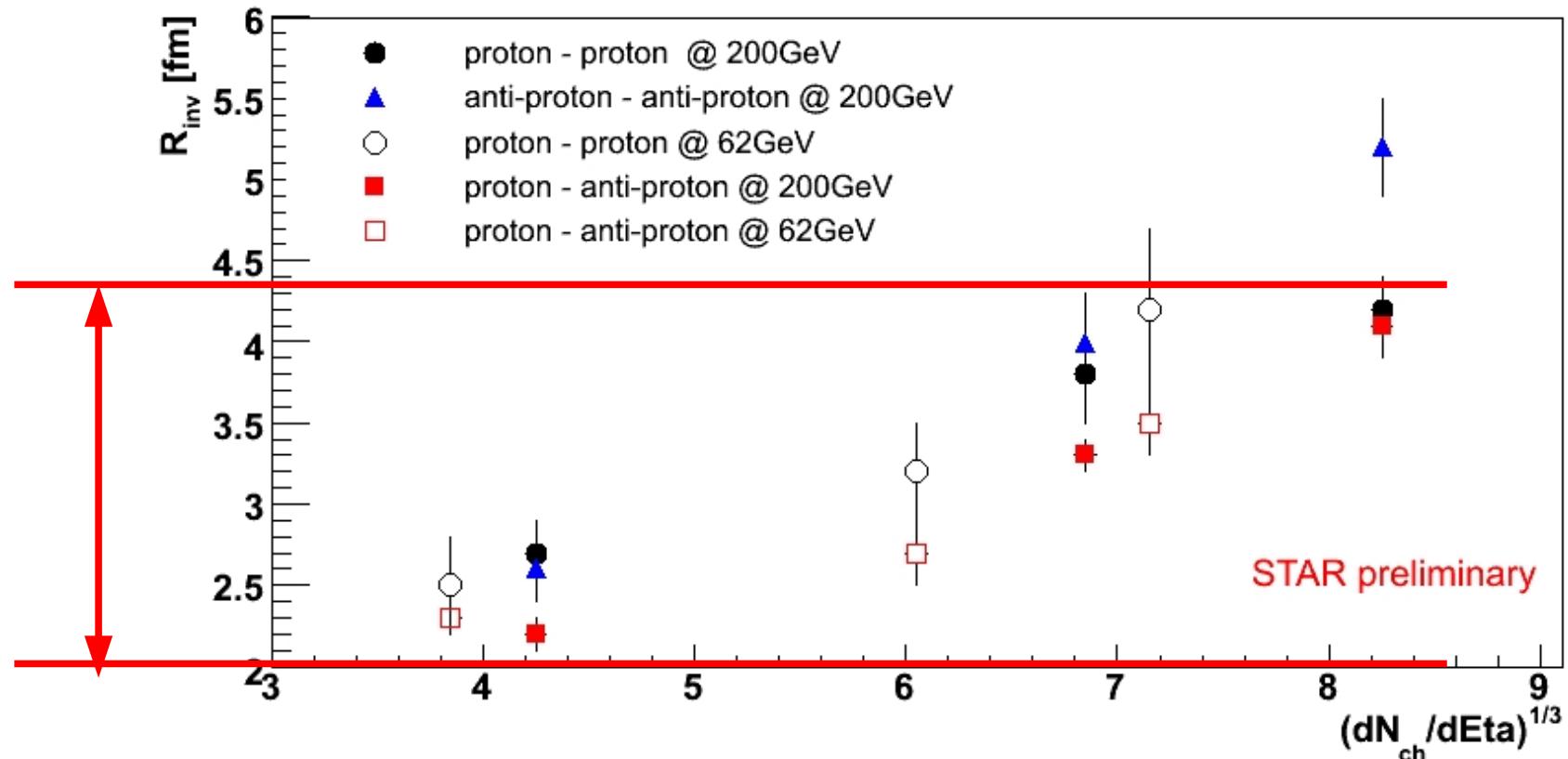
- Filikhin, Gal ('02)
 $a_0 = -0.77$ fm, $r_{\text{eff}} = 6.59$ fm
- Hiyama et al. ('02)(HKMYY)
 $a_0 = -0.575$ fm, $r_{\text{eff}} = 6.45$ fm
- Hiyama et al. ('10)(HKMYY)
 $a_0 = -0.44$ fm, $r_{\text{eff}} = 10.1$ fm

*Y. Fujiwara, Y. Suzuki, C. Nakamoto,
Prog.Part.Nucl.Phys. 58 (2007) 439-520*

Toward $\Lambda\Lambda$ correlation at RHIC: Source Size

■ Source size : $R = (2-4.5)$ fm

- Smaller than last collision point dist. results in hadron cascade (JAM)
→ Interaction in the early stage at RHIC
- Smaller than π , K homogeneity length → Further smaller for Λ ?



A. Kisiel (H. P. Zbroszczyk) (STAR)

Toward $\Lambda\Lambda$ correlation at RHIC: $\Lambda\Lambda$ interaction

■ $\Lambda\Lambda$ interaction

After Nagara, “plausible” $\Lambda\Lambda$ interaction becomes weaker.

Bond energy $\Delta B_{\Lambda\Lambda} = 0.7$ MeV (old guess=(3-6) MeV)

● fss2 (quark model interaction): No bound state

Y. Fujiwara, M. Kohno, C. Nakamoto, Y. Suzuki, PRC64('01)054001

Bond energy $\Delta B_{\Lambda\Lambda} = (1.2-1.9)$ MeV (depending on ΛN int.)

● Nijmegen model D (boson exch., $R_c=0.46$ fm): with bound state

M.M. Nagels, T.A. Rijken, J.J. de Swart, PRD15('77)2547

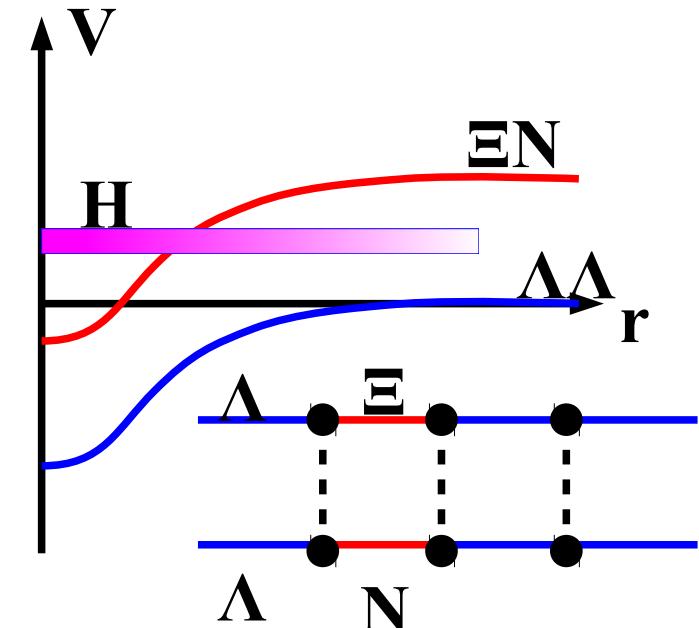
B.E.(H) ~ 1.6 MeV

■ Resonance “H” btw $\Lambda\Lambda$ - ΞN threshold

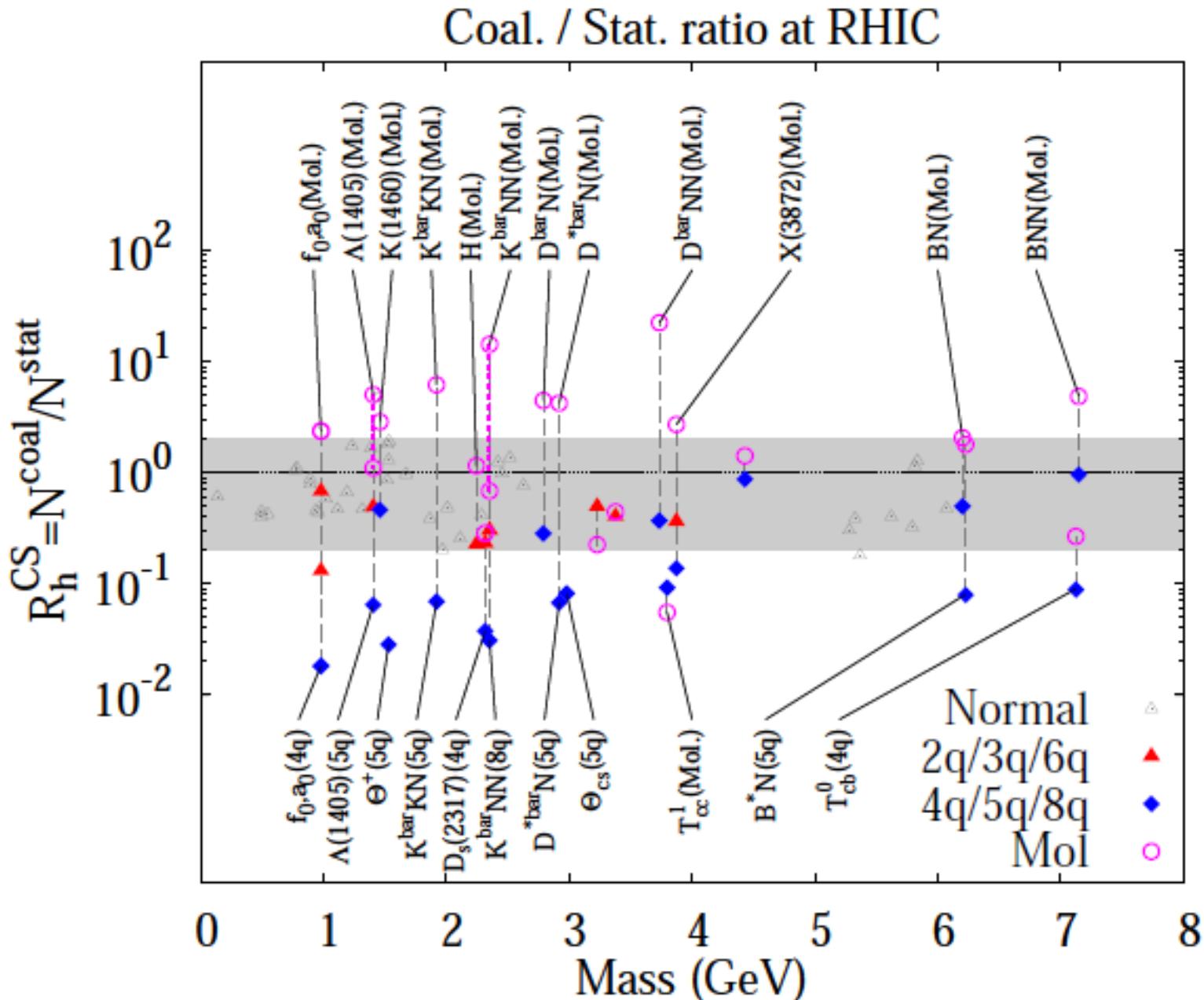
→ Couple channel calc. is required

● One range gaussian coupling potential is assumed.

● ΞN potential (diagonal) effects on $C(q)$ is almost negligible.



Coalescence / Statistical Ratio



S. Cho et al. (ExHIC Collab.), PRL106('11)212001; PRC 84 ('11) 064910

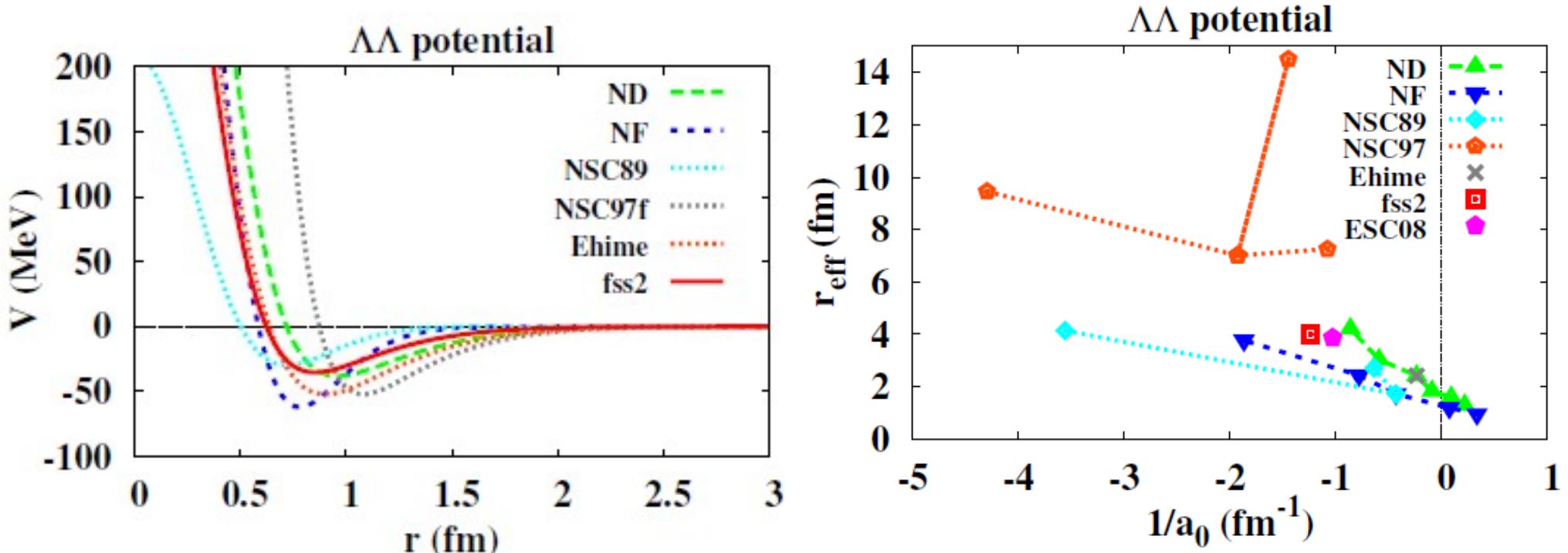
A. Ohnishi @ Mol. 2015.09, Sep. 21, 2015, Kyoto, Japan

*Exotic Interactions
from Heavy Ion Collisions
--- H particle and $\Lambda\Lambda$ interaction ---*

$\Lambda\Lambda$ interaction

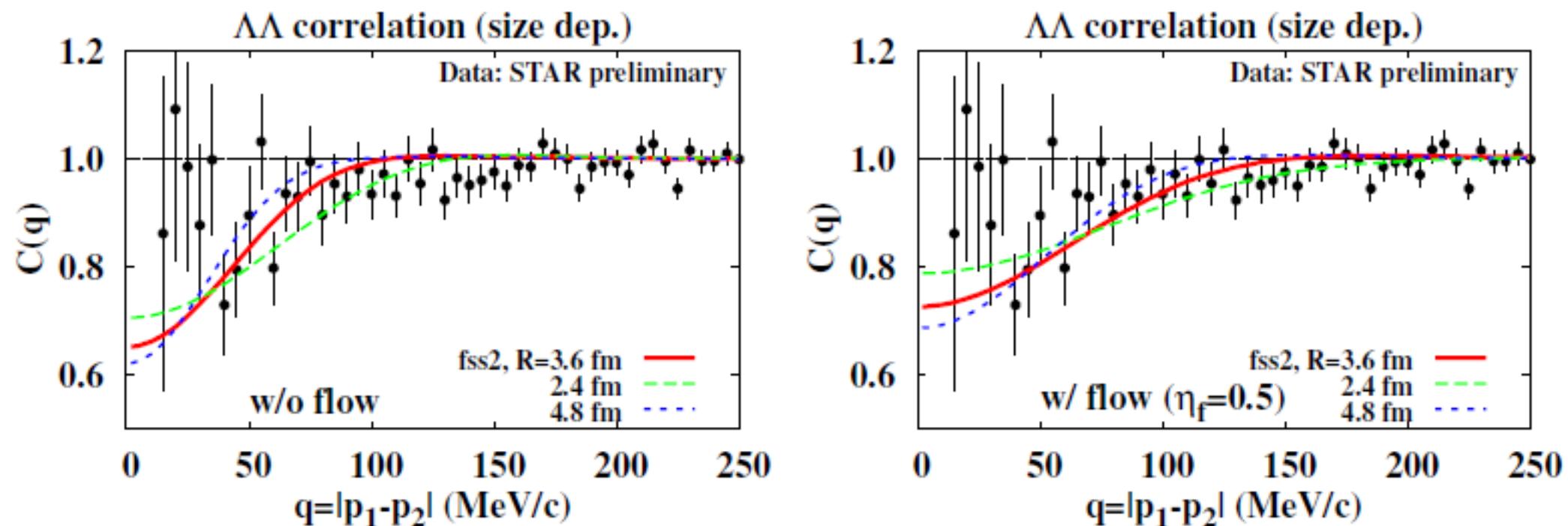
Type of $\Lambda\Lambda$ interactoin

- Meson exchange models: Nijmegen model D, F, Soft Core (89, 97), ESC08
Nagels, Rijken, de Swart ('77, '79), Maessen, Rijken, de Swart ('89), Rijken, Stoks, Yamamoto ('99); Rijken, Nagels, Yamamoto ('10).
- Quark cluster model interaction: fss2
Fujiwara, Fujita, Kohno, Nakamoto, Suzuki ('00)
- Phenomenological model: Ehime T. Ueda et al. ('99).



Flow Effects

- Too small source size ~ 1.7 fm with Σ^0 feed down effects ?
- Flow effects make the “apparent” size smaller.
 - Relative momentum is enhanced by the flow.
→ Actual size ~ 3 fm (consistent with proton source size)



Transverse rapidity, $Y_T = \eta_f r_T / R$ Morita

Summary

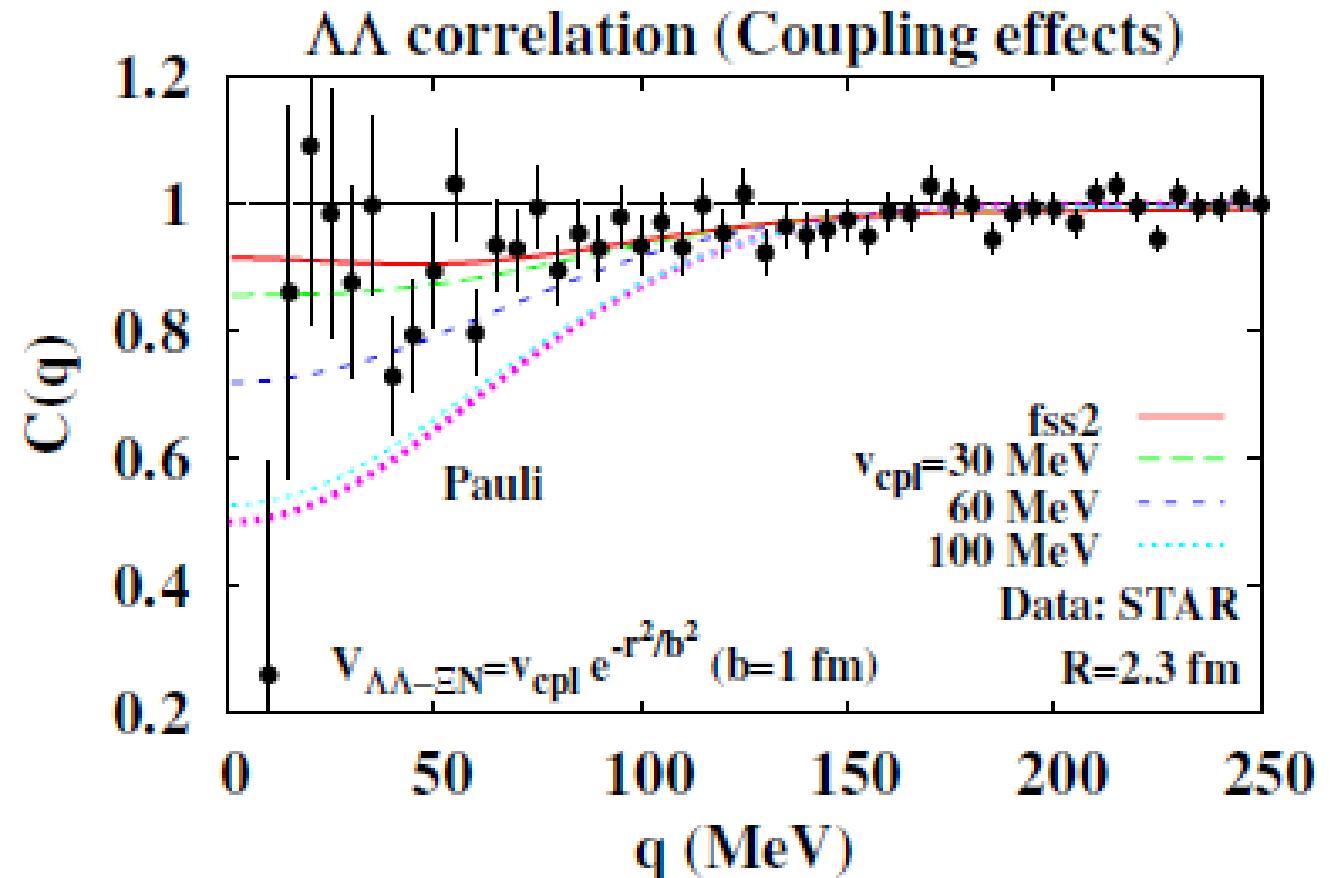
- Various exotic hadrons will be formed in heavy ion collisions, and it is the right time to search them.
 - Extended hadron molecule production is enhanced in HICs.
 - Compact multiquark states may be suppressed in HICs.
 - Systematic study of exotics from e^+e^- reaction may be promising.
- $\Lambda\Lambda$ correlation observed at RHIC is useful to distinguish proposed $\Lambda\Lambda$ potentials.

N.Shah et al. (STAR Collab.), Acta Phys. Pol. Suppl. 5 ('12) 593.

- Preferred $\Lambda\Lambda$ interactions have $1/a_0 < -0.8 \text{ fm}^{-1}$, $r_{\text{eff}} > 3 \text{ fm}$.
Weakly attractive. Consistent with Nagara event ($a_0 = -(0.7-1.3) \text{ fm}$)
E. Hiyama et al. PRC66('02)024007; A.M.Gasparyan et al. PRC85('12)015204.
- Source size is $R \sim 3 \text{ fm}$, if we take account the flow effects.
Consistent with proton source size (2.5-4) fm (*STAR preliminary*).
- Existence of resonance “H” requires higher statistics.

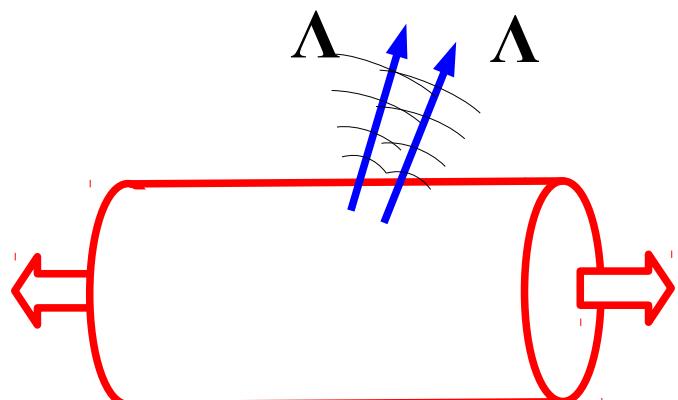
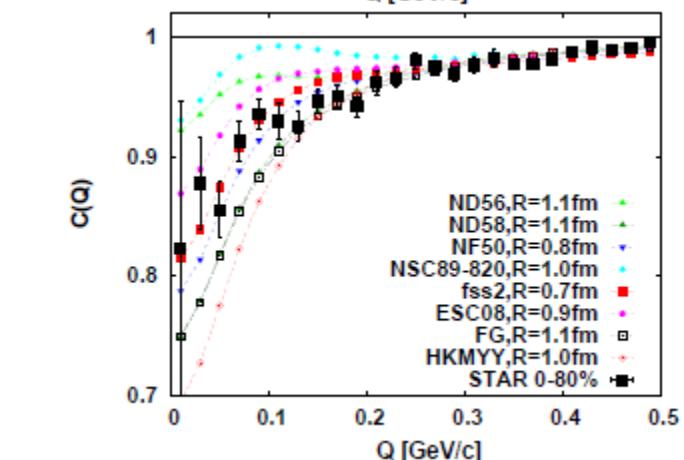
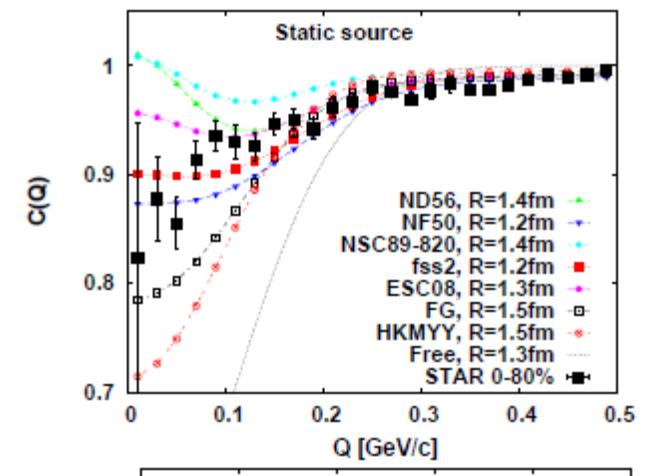
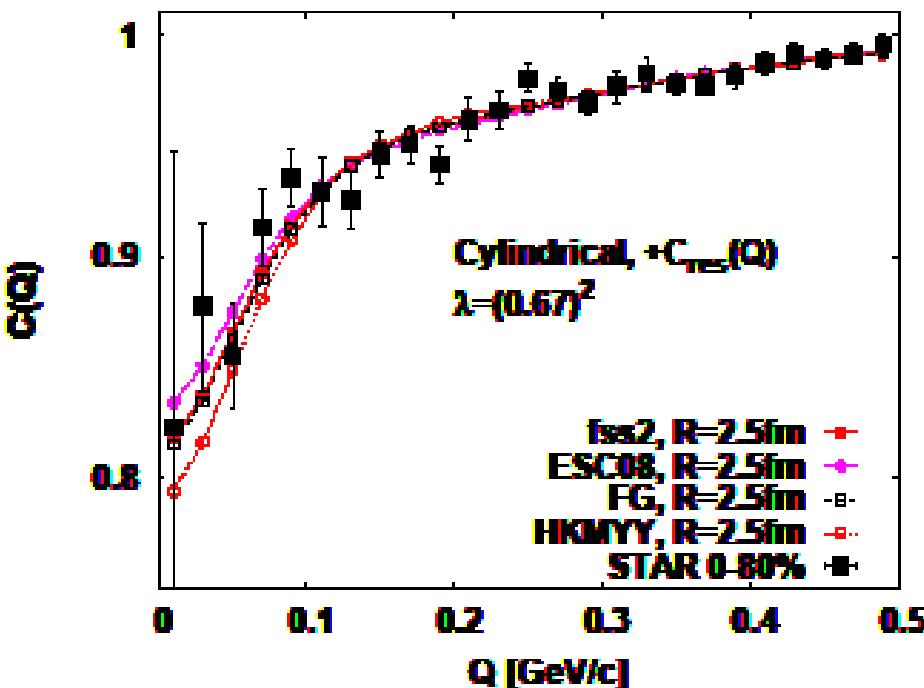
Coupling Effects

- Coupled channels effects with ΞN channel is considered.
 - Coupling with ΞN channel suppresses $C(q)$ at low q . (\sim Imag. pot.)
 - Unreasonably large coupling would meaningfully modify $C(q)$.



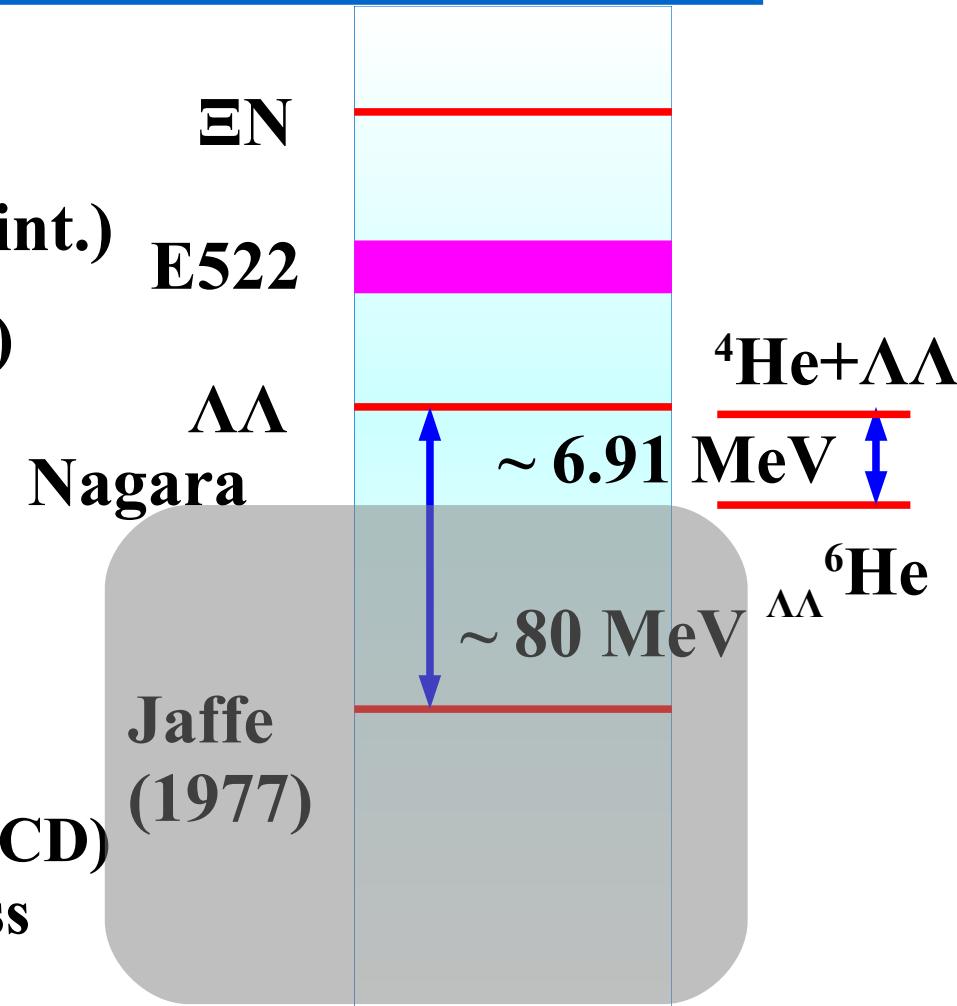
Source models

- Static & Spherical source
- Cylindrical model
 - Bjorken expansion in long. direction
 - Transverse flow from pT spectrum
- Cylindrical + Feed down ($\Sigma 0$) + Residual



Where is the $S=-2$ dibaryon ($uuddss$) “H” ?

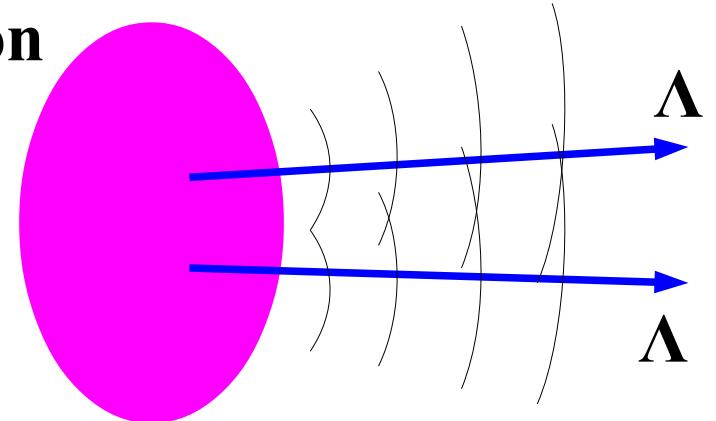
- Jaffe's prediction (1977)
→ 80 MeV below $\Lambda\Lambda$
(strong attraction from color mag. int.)
- Double hypernuclei $_{\Lambda\Lambda}^6\text{He}$ (Nagara)
→ No deeply bound “H”
- Resonance or Bound “H” ?
 - KEK-E522 (Yoon et al., ('07))
→ “bump” at $E_{\Lambda\Lambda} \sim 15$ MeV
 - Lattice QCD (HAL QCD & NPLQCD)
→ bound H at large ud quark mass
- How about HIC ?
 - RHIC & LHC = Hadron Factory including Exotics
 - “H” would be formed as frequently as stat. model predicts.
*Cho, Furumoto, Hyodo, Jido, Ko, Lee, Nielsen, AO, Sekihara, Yasui, Yazaki
(ExHIC Collab.), PRL('11)212001; arXiv:t:1107.1302*



$\Lambda\Lambda$ correlation in HIC

■ Merit of HIC to measure $\Lambda\Lambda$ correlation

- Source is “simple and clean” !
T, μ , flow, size, ... are well-analyzed.
- Source size is big and probes w.f. tail.
- Discovery of “H” and/or Constraint
on $\Lambda\Lambda$ int.



■ Gaussian Source + s-wave int.

c.f. P. Danielewicz; Bauer, Gelbke, Pratt, Annu. Rev. Nucl. Part. Sci. 42('92)77.

$$C_{\Lambda\Lambda}(q) = \frac{\int dx_1 dx_2 S(x_1, p+q) S(x_2, p-q) |\psi^{(-)}(x_{12}, q)|^2}{\int dx_1 dx_2 S(x_1, p+q) S(x_2, p-q)}$$
$$\approx 1 - \frac{1}{2} \exp(-q^2 R^2) + \frac{1}{2} \int dr S_{12}(r) (|\chi_0(r)|^2 - |j_0(qr)|^2)$$

(χ_0 : s-wave wave func., $S_{12}(x) = (R\sqrt{\pi})^{-3} \exp(-r^2/R^2)$)

Question:

*Can we constrain $\Lambda\Lambda$ interaction
from correlation data at RHIC ?*

*Is the constraint consistent
with $\Lambda\Lambda$ corr. in (K^-, K^+) reaction ?*

Does H exist as a bound state or a resonance ?

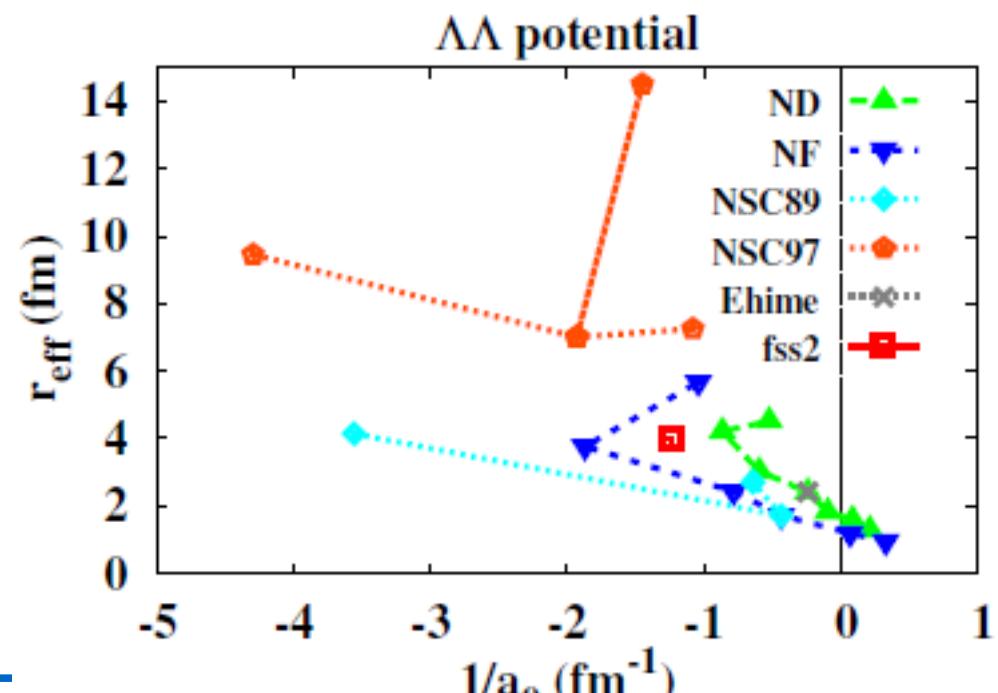
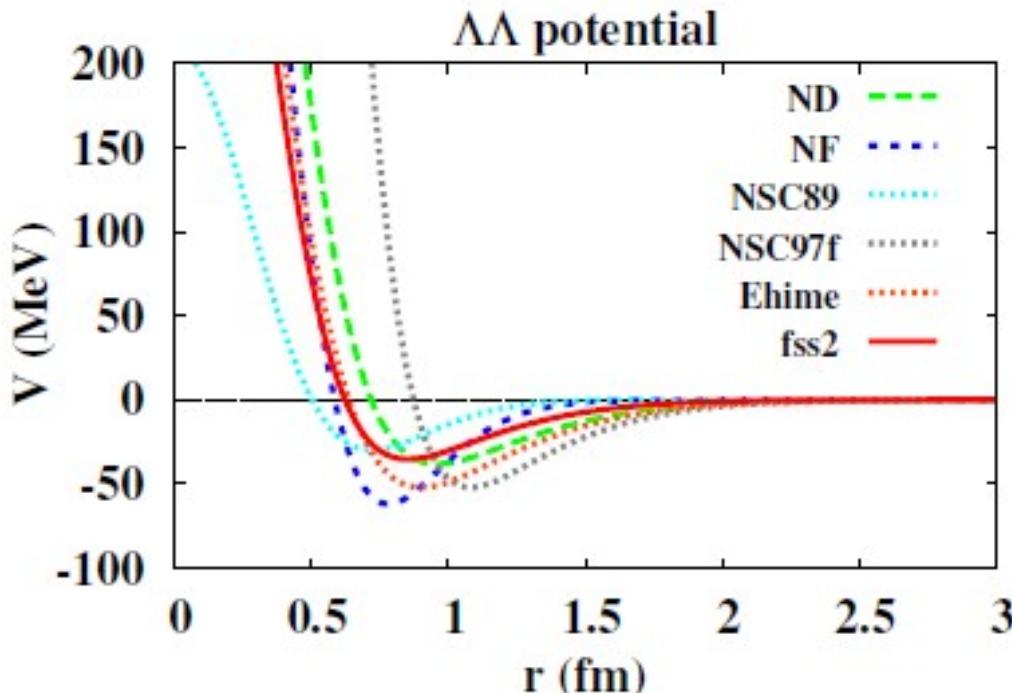
- Source: Gaussian (HIC), Cascade model results (K^-, K^+)
- Interaction: ND, NF, NSC89, NSC97, Ehime, fss2
(two or three range gaussian fit)
- Other channel: $\Lambda\Lambda$ - ΞN couple channels, Σ^0 decay feed effects

$\Lambda\Lambda$ interaction

Type of $\Lambda\Lambda$ interactoin

- Meson exchange models: Nijmegen model D, F, Soft Core (89, 97)
Nagels, Rijken, de Swart ('77, '79), Maessen, Rijken, de Swart ('89), Rijken, Stoks, Yamamoto ('99)
- Quark cluster model interaction: fss2
Fujiwara, Fujita, Kohno, Nakamoto, Suzuki ('00)
- Phenomenological model: Ehime

Two (or three) range gaussian fit results are used in the analysis.



$\Lambda\bar{\Lambda}$ Correlation in (K^- , K^+) Reaction

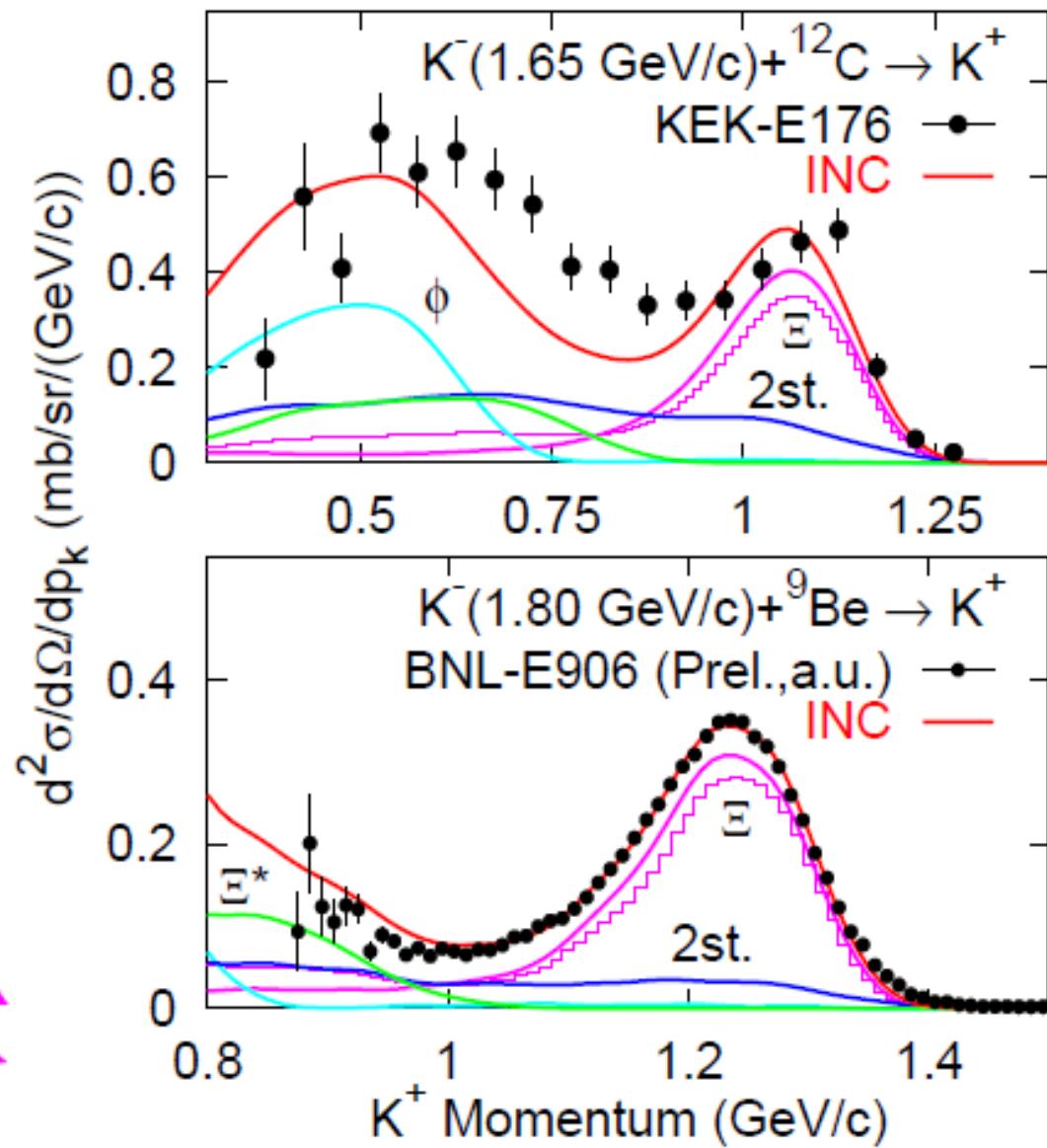
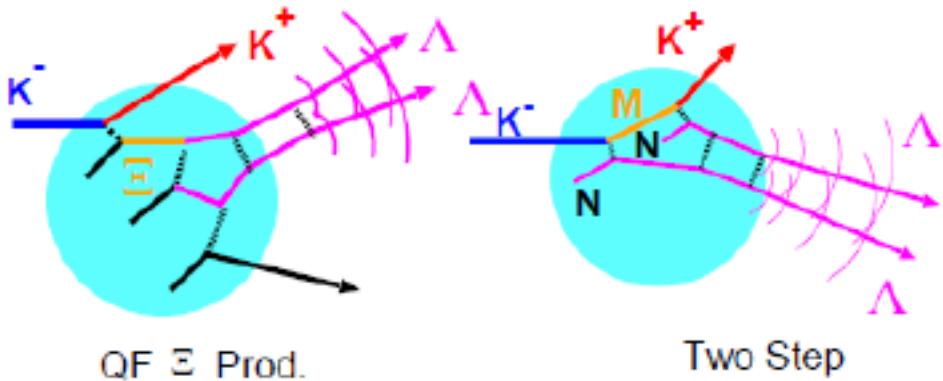
$\Lambda\bar{\Lambda}$ Correlation in (K^-, K^+) Reaction (1)

■ K^+ production mechanism

- QF Ξ production
- Heavy meson production and Decay

*Gobbi, Dover, Gal,
PRC50 (1994) 1594.*

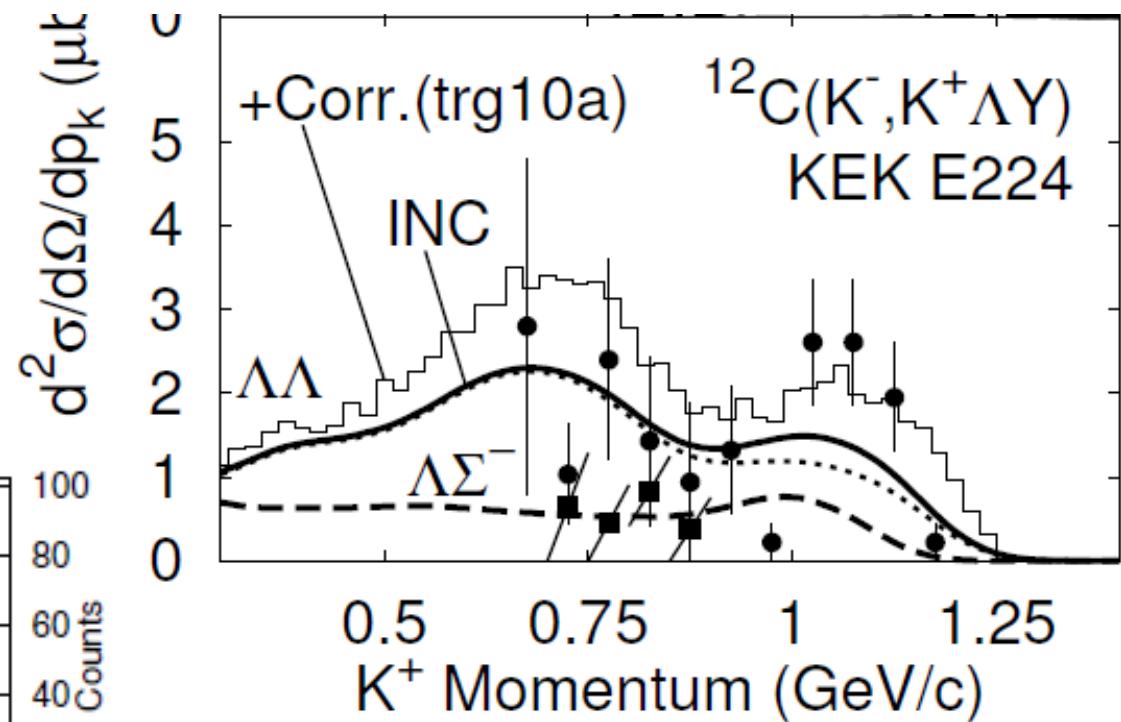
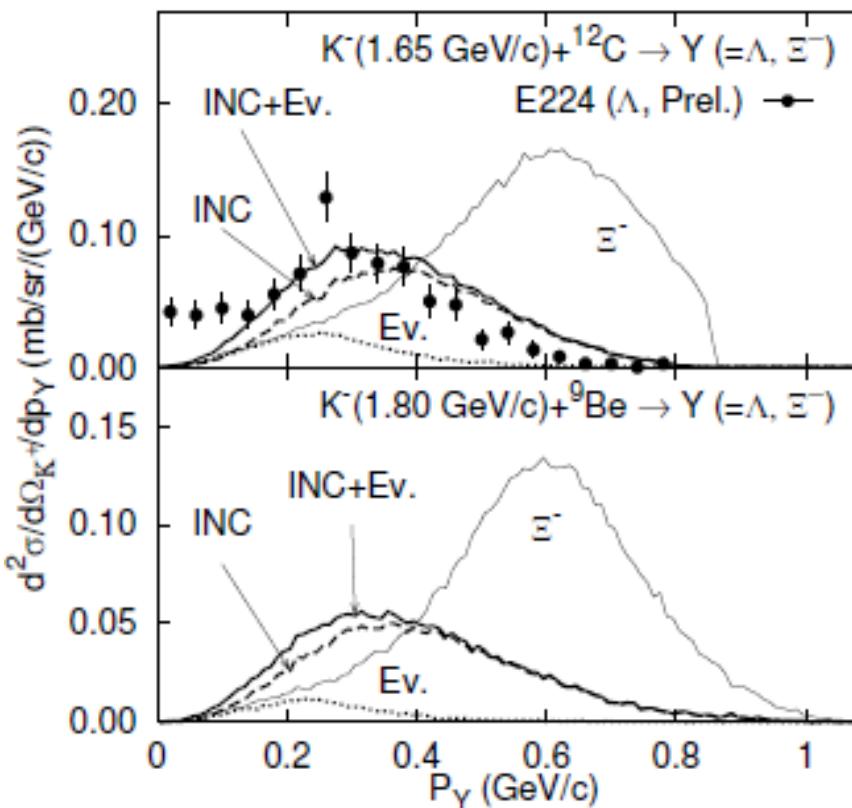
- Two step processes
*Nara, AO, Harada, Engel,
NPA614 (1997) 433*



*AO, Hirata, Nara, Shinmura, Akaishi,
Few-Body Syst. Suppl. 12 (2000), 367*

$\Lambda\bar{\Lambda}$ Correlation in (K^-, K^+) Reaction (2)

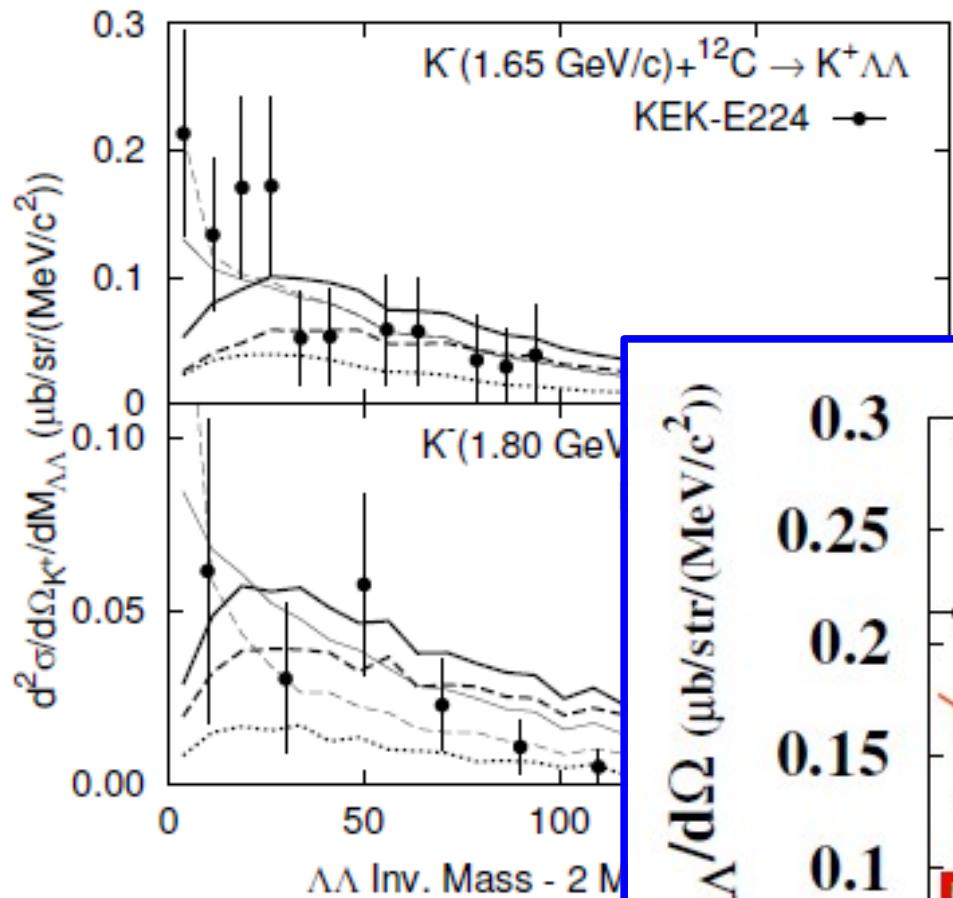
- Λ production mechanism
 - Cascade processes
 - Evaporation from hyper compound nuclei



*AO, Hirata, Nara, Shinmura, Akaishi,
 NPA670(2000), 297c*

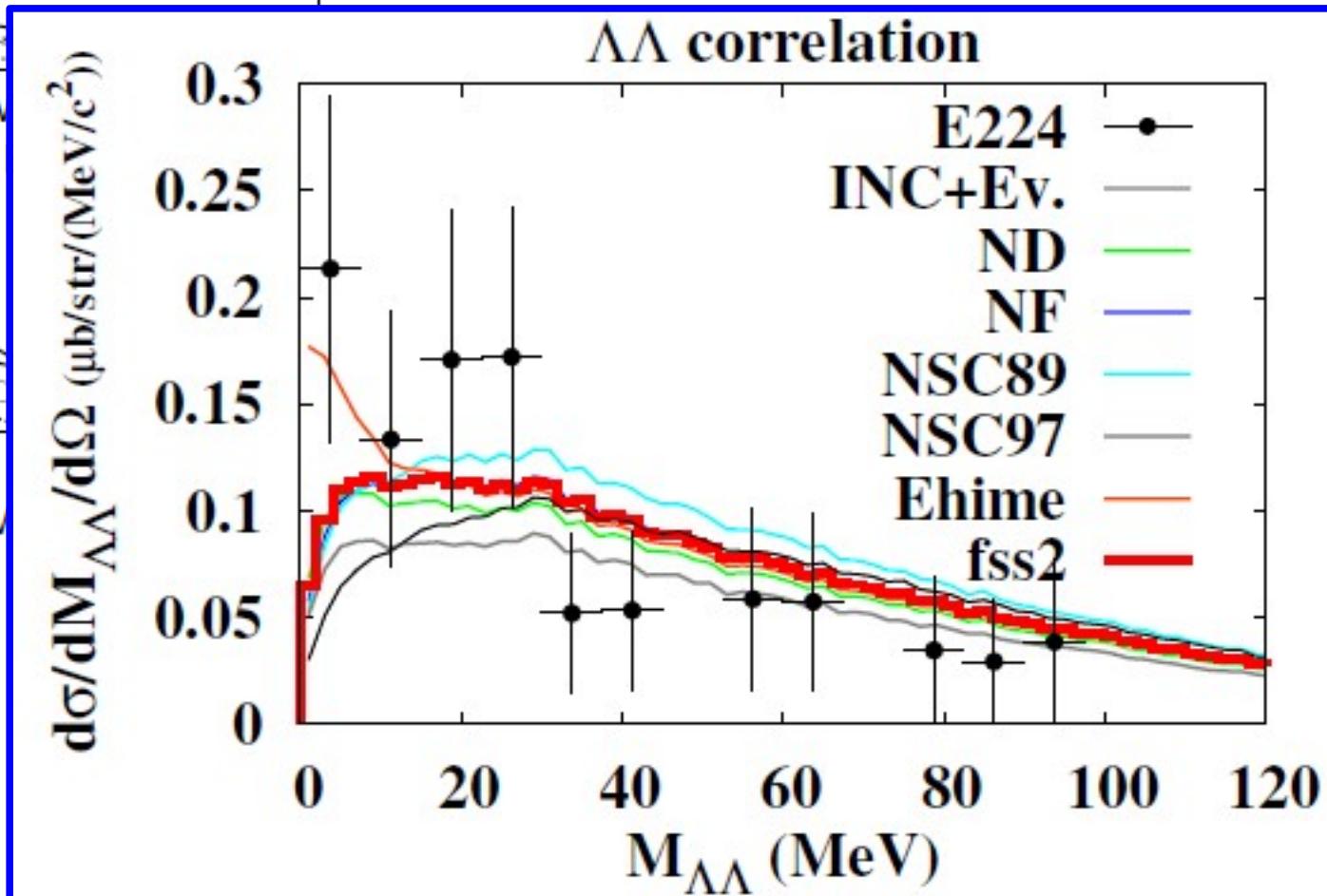
*AO, Hirata, Nara, Shinmura, Akaishi,
 NPA691(2001), 242c*

$\Lambda\Lambda$ Invariant Mass Spectrum



*AO, Hirata, Nara,
Shinmura, Akaishi,
NPA684(2001), 595c*

$\Lambda\Lambda$ int. constrained in HIC
are consistent with $\Lambda\Lambda$ inv.
mass spectrum in (K , $K^+\Lambda\Lambda$)



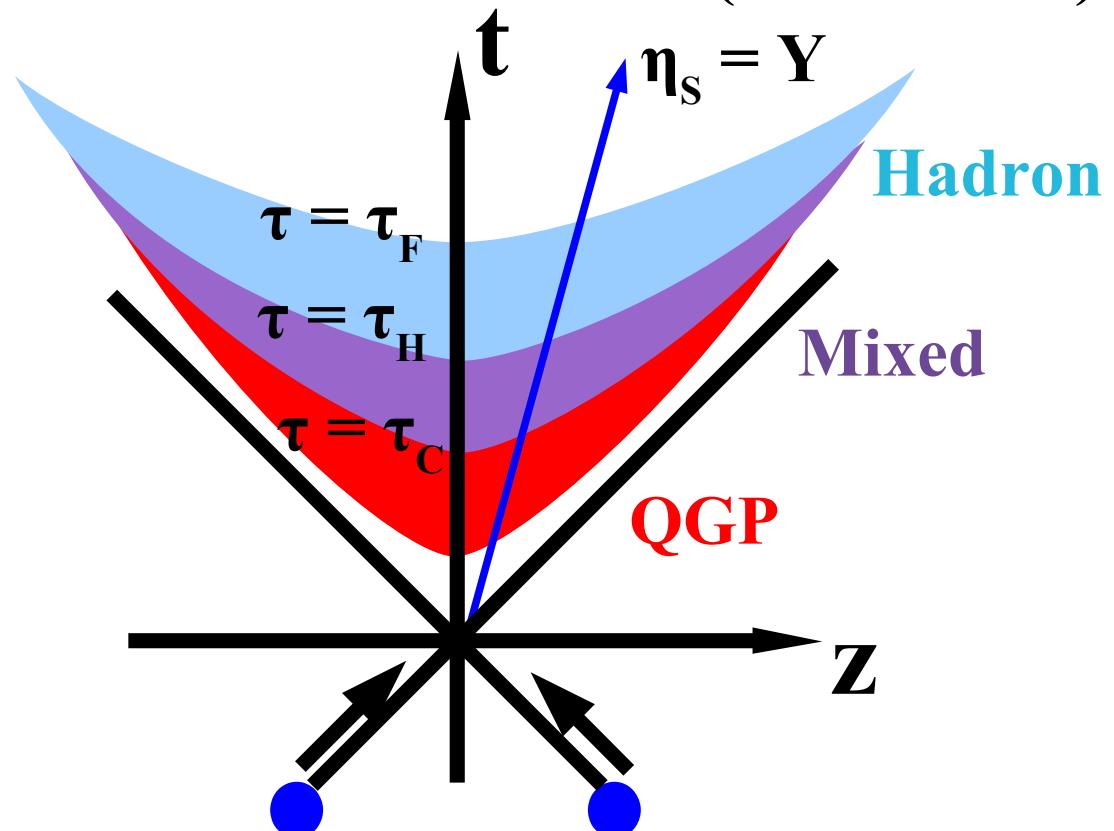
Exotic Hadron Yields in Heavy Ion Collisions

Schematic picture of HIC

HIC picture based on the first order phase transition

- $\tau = \tau_c$, $T = T_c$, $V = V_c \rightarrow$ QGP start to hadronize (quark coal.)
- $\tau = \tau_h$, $T = T_h = T_c$, $V = V_h \rightarrow$ Hadronization is over (stat. model)
- $\tau = \tau_f$, $T = T_f$, $V = V_f \rightarrow$ Hadronic Freeze-out (hadron coal.)

	RHIC	LHC
$N_u = N_d$	245	662
$N_s = N_{\bar{s}}$	150	405
$N_c = N_{\bar{c}}$	3	20
$N_b = N_{\bar{b}}$	0.02	0.8
V_c	1000 fm^3	2700 fm^3
$T_c = T_h$	175 MeV	175 MeV
V_h	1908 fm^3	5152 fm^3
μ_B	20 MeV	20 MeV
μ_s	10 MeV	10 MeV
V_f	11322 fm^3	30569 fm^3
T_f	125 MeV	125 MeV



L.W.Chen, V.Greco, C.M.Ko, S.H.Lee, W.Liu, PLB 601('04)34.

Statistical Model

■ Statistical model

$$N_h^{\text{stat}} = V_H \frac{g_h}{2\pi^2} \int_0^\infty \frac{p^2 dp}{\gamma_h^{-1} e^{E_h/T_H} \pm 1}$$

$(N_h = dN_h/dy \text{ (y=rapidity)}, V_H = \text{Chem. freeze-out vol.})$

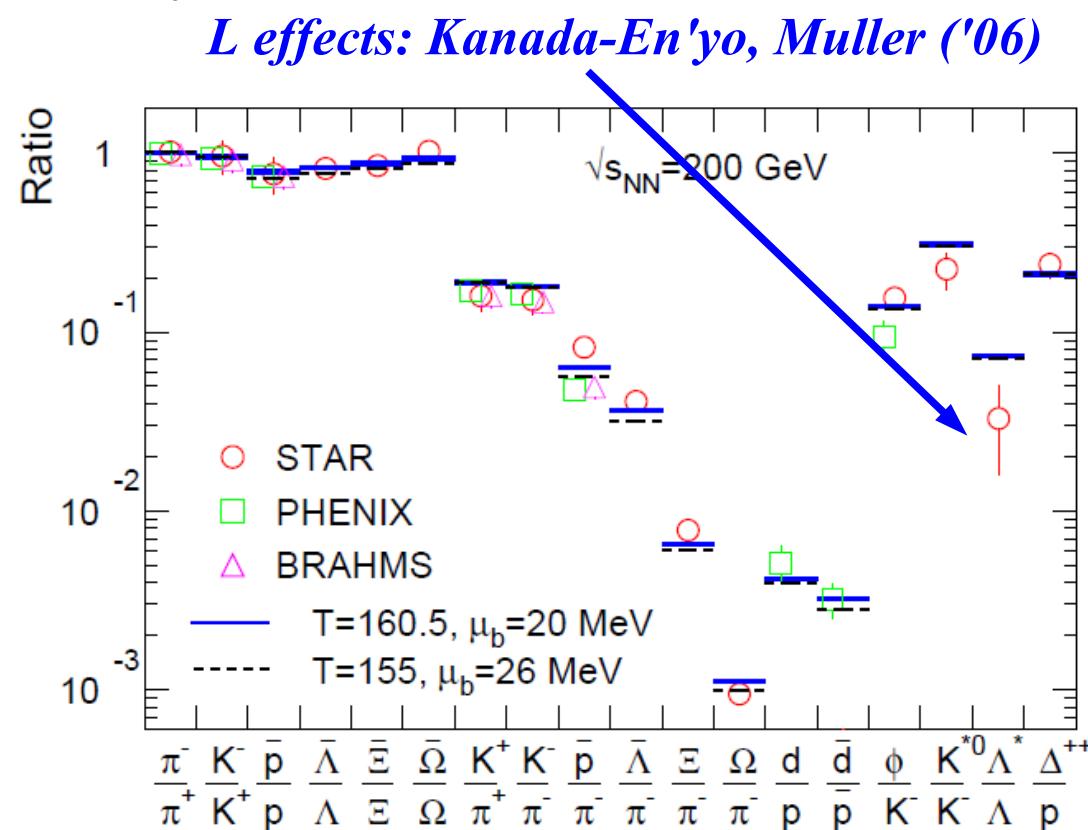
→ Successful to predict the hadron yield ratio at RHIC

■ Fugacity factor γ

- u,d,s: chem. equil.
- c,b: enhanced by initial hard processes

Fugacities of c and b quarks are set to reproduce expected c and b quark numbers.

$$\gamma_h = \gamma_c^{n_c + n_{\bar{c}}} \gamma_b^{n_b + n_{\bar{b}}} e^{(\mu_B B + \mu_s S)/T_H}$$



A. Andronic, P. Braun-Munzinger, J. Stachel, NPA772('06)167.

Coalescence model

- Yield = Overlap of const. dist. & Hadron intrinsic Wigner func.
(Sudden approximation)

Sato, Yazaki (1984), Hwa, Yang (2003), Greco, Ko, Levai (2003), Fries, Muller, Nonaka, Bass (2003), Chen, Ko, Lee (2003)

$$N_h^{\text{coal}} = \frac{g_h \int \left[\prod_{i=1}^n \frac{1}{g_i} \frac{p_i \cdot d\sigma_i}{(2\pi)^3} \frac{d^3 p_i}{E_i} f(x_i, p_i) \right] \times f^W(x_1, \dots, x_n : p_1, \dots, p_n)}{\text{Dist. of constituents} \quad \text{Intrinsic Wigner func.}}$$

- Yield in HIC

- Quark & hadron dist. = Transverse Boltzmann + Bjorken
Chen, Ko, Liu, Nielsen (2007)
- Hadron intr. Wigner func. = s-wave and p-wave HO w.f.
Kanada-En'yo, Muller (2006)

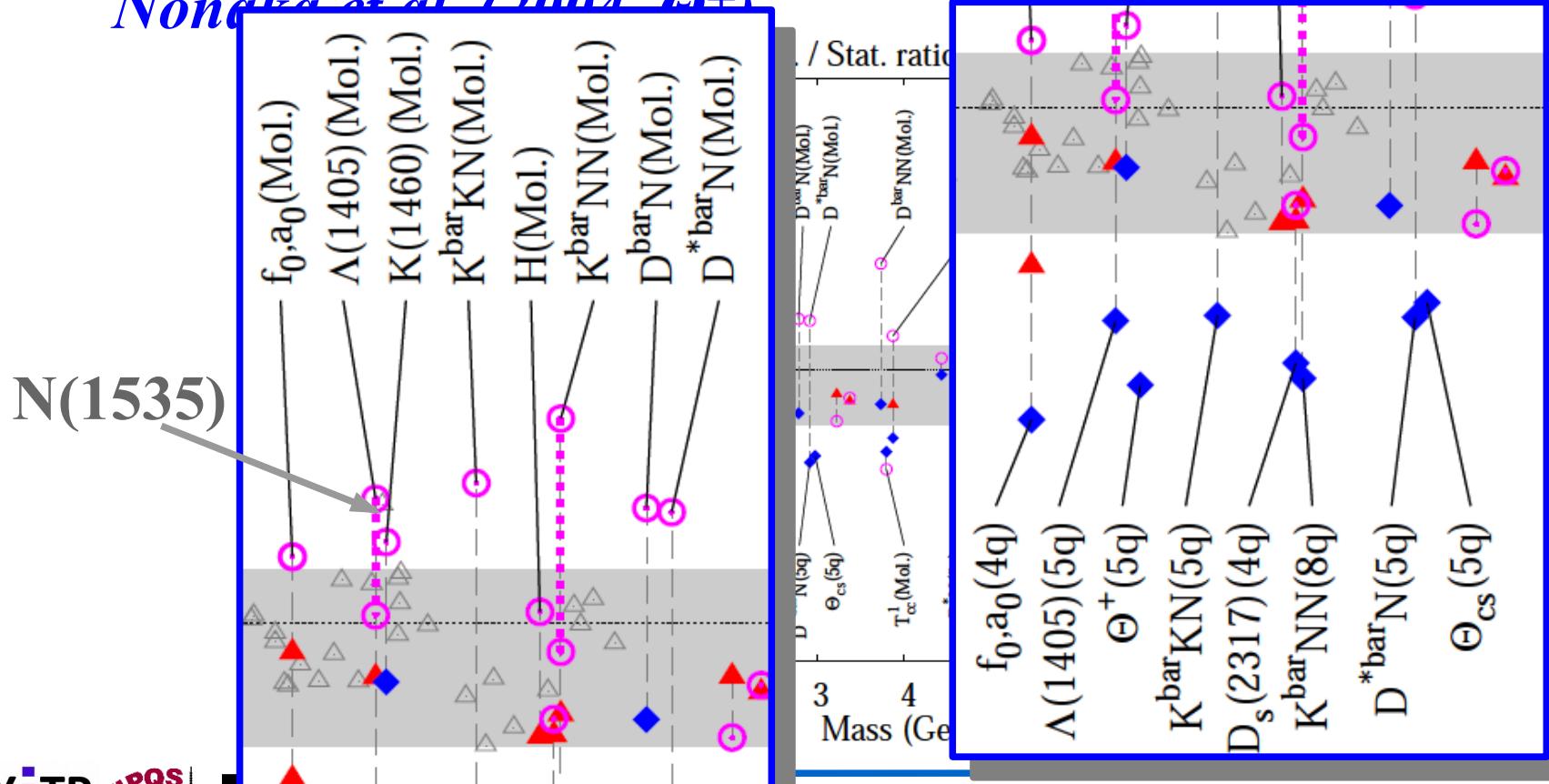
$$N_h^{\text{coal}} \simeq g_h \prod_{j=1}^n \frac{N_j}{g_j} \prod_{i=1}^{n-1} \frac{(4\pi\sigma_i^2)^{3/2}}{V(1 + 2\mu_i T\sigma_i^2)} \left[\frac{4\mu_i T\sigma_i^2}{3(1 + 2\mu_i T\sigma_i^2)} \right]^{l_i}$$

σ = Gaussian width, μ =reduced mass, N = constituent yield

Coalescence / Statistical Ratio

- Normal hadrons: $R_{CS} = (0.2-2)$ *Normal band*
- Extended hadronic molecule: Large yield is expected, $R_{CS} > 2$.
 $\Lambda(1405) = \bar{K}N, \bar{K}NN, \bar{K}KN, \bar{D}NN, \dots$ ($\hbar\omega = (6-50)$ MeV)
- Compact Multiquark states will be suppressed in HICs, $R_{CS} < 0.2$
 $f_0/a_0(q\bar{q}\bar{q}\bar{q}), \Theta^+(uudd\bar{s}), H(uudd\bar{s}\bar{s}), \Theta_{cs}(uud\bar{s}\bar{c}), \dots$

Nomura et al. (2004, Ω^+)



Ohnishi @ Mol. 2015.09, Sep. 21, 2015, Kyoto, Japan

Enhancement of Hadronic Molecules: Why ?

- Simple estimate: 2-body, Gaussian w.f. + Thermal dist. of constituents

$$N_h \propto \int \frac{d^D x d^D p}{(2\pi\hbar)^D} \frac{f_w(x, p)}{f_{th}(x, p)} = \left[\left(\frac{4}{\hbar^2} \right) ((\Delta p)^2 + \mu T) ((\Delta x)^2 + 2R^2) \right]^{-D/2}$$

Intrinsic

Constituents (thermal)

$$f_w(x, p) = \left(\frac{\hbar}{\Delta x \Delta p} \right)^D \exp \left(-\frac{x^2}{2(\Delta x)^2} - \frac{p^2}{2(\Delta p)^2} \right),$$

$$f_{th}(x, p) = \left(\frac{\hbar^2}{2\mu T R^2} \right)^{D/2} \exp \left(-\frac{x^2}{4R^2} - \frac{p^2}{2\mu T} \right)$$

- N_h is large when f_w shape is similar to f_{th} in phase space.

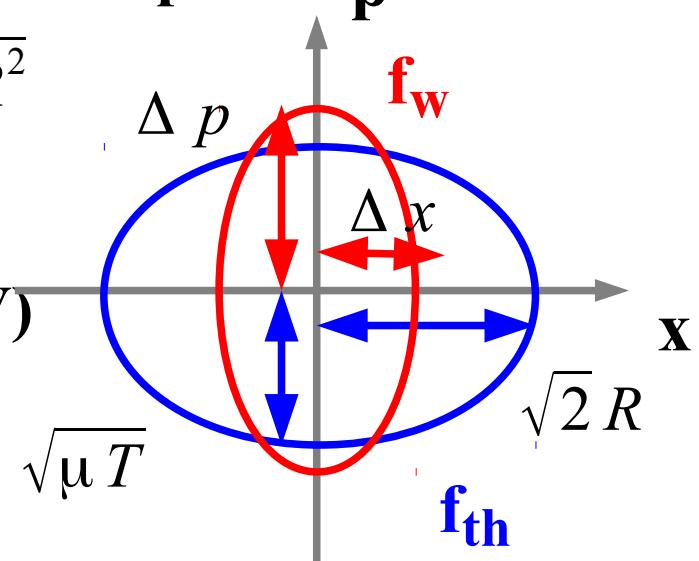
$$\Delta x : \Delta p = \sqrt{2} R : \sqrt{\mu T} \rightarrow \hbar \omega = \sqrt{\hbar^2 T / 2 \mu R^2}$$

- Example: $T=170$ MeV, $\mu=500$ MeV

(red. mass), $R = 5$ fm (source size)

→ optimal $\hbar \omega \sim 16$ MeV ($\ll 300$ -500 MeV)

*Large source size & Moderate T
prefer extended hadrons*



Why are Multi-quark Configs. Suppressed ?

- Hadron yield is sensitive to the structure in coal.

- Additional q penalty factor

s-wave $\frac{1}{g_i} \frac{N_i}{V} \frac{(4\pi\sigma_i^2)^{3/2}}{(1 + 2\mu_i T \sigma_i^2)} \sim 0.36$

p-wave $\frac{1}{g_i} \frac{N_i}{V} \frac{2}{3} \frac{(4\pi\sigma_i^2)^{3/2} 2\mu_i T \sigma_i^2}{(1 + 2\mu_i T \sigma_i^2)^2} \sim 0.09$

Nonaka et al. (2004, Θ^+)

Kanada-En'yo, B. Muller (2006, $\Lambda(1520)$)

Large V disfavors multi-quarks !

- STAR data (2003): $N(f_0(980)) \sim 8.4$

[$f_0(980)/\rho^0 \sim 0.2$, stat. $N(\rho^0) \sim 42$]

Stat: 5.6, 2q:0.76-3.8, 4q:0.1, Mol: 13

→ Tetra-quark picture

underestimate the measured yield of f_0 .

*P. Fachini [STAR Collaboration],
Nucl. Phys. A 715, 462 (2003).*

