

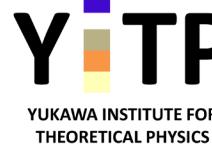
Hadron-Hadron Correlation and Interaction from Heavy-Ion Collisions

Akira Ohnishi (YITP, Kyoto U.)



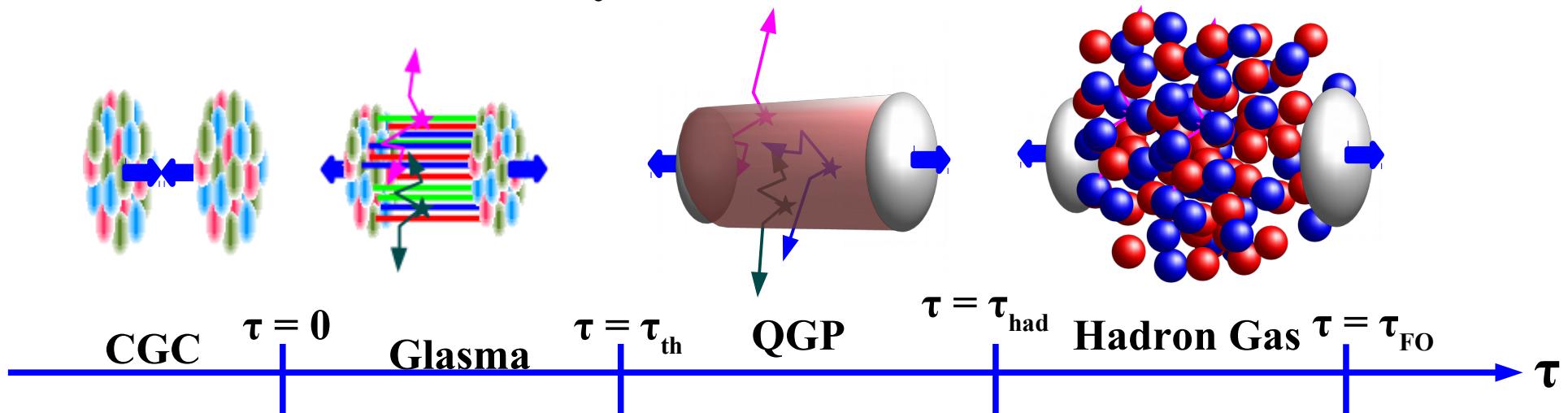
International Nuclear Physics Conference
Adelaide Convention Centre, Australia
11-16 September 2016

- *K. Morita, T. Furumoto, AO, PRC91('15)024916. AA*
- *AO, K. Morita, K. Miyahara, T. Hyodo,
NPA954('16), 294 [arXiv:1603.05761] AA, K-p*
- *K. Morita, AO, F. Etminan, T. Hatsuda,
PRC94('16)031901(R) [arXiv:1605.06765 [hep-ph]] Ω-p*



Heavy-Ion Collision as a Hadron Factory

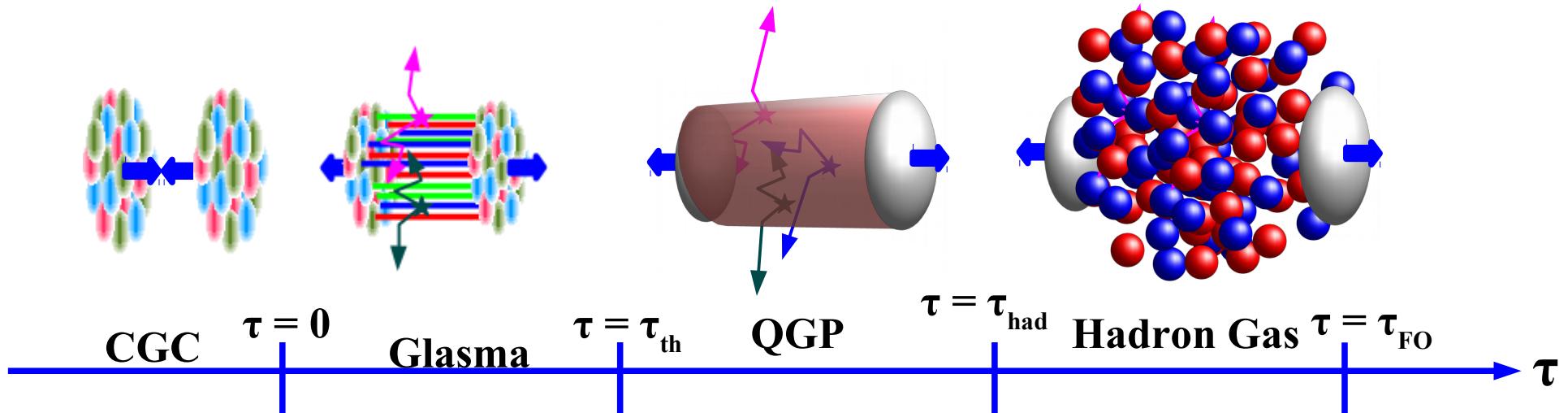
- Many hadrons are produced simultaneously in high-energy HIC
→ HIC = Hadron Factory



- Physics of Hot and Dense Matter Esumi, Mironov, Meyer
 - Yield Ratio, η - $\Delta\Phi$ correlation, R_{AA} , Flow, Fluctuations, ...
→ (T, μ) , Ridge, Jet quenching (E-loss), parton collectivity, CP, ...

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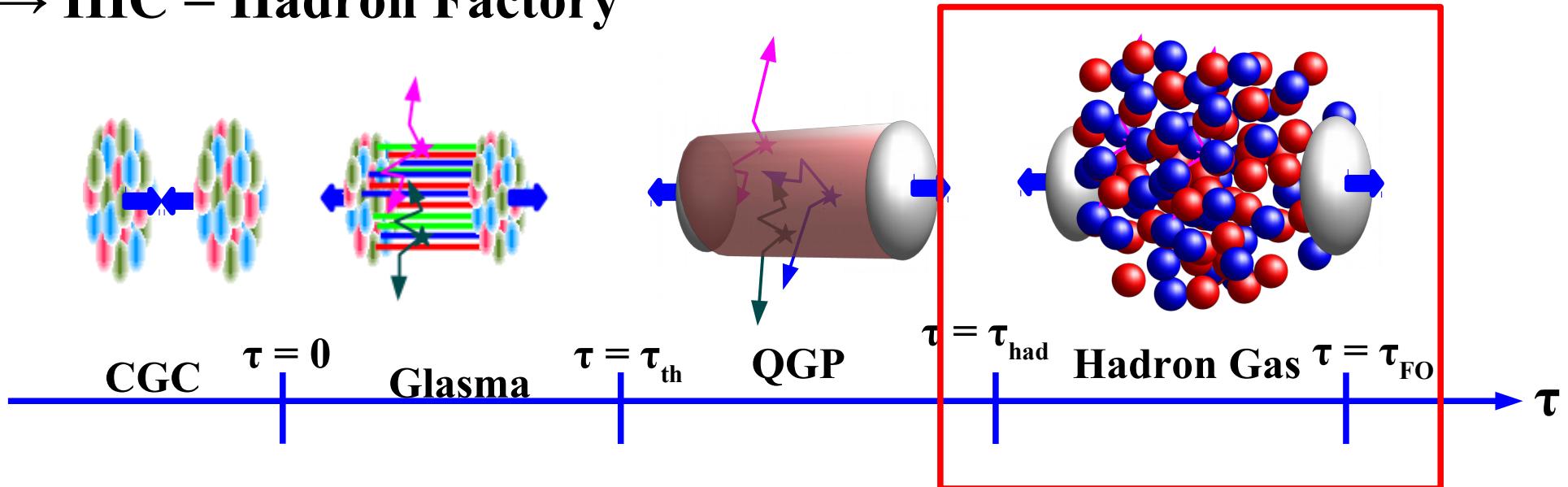
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- Hadron Physics Robbe (LHCb)
 - Medium effects, Formation mechanism, Exotic Hadrons, Interaction between hadrons

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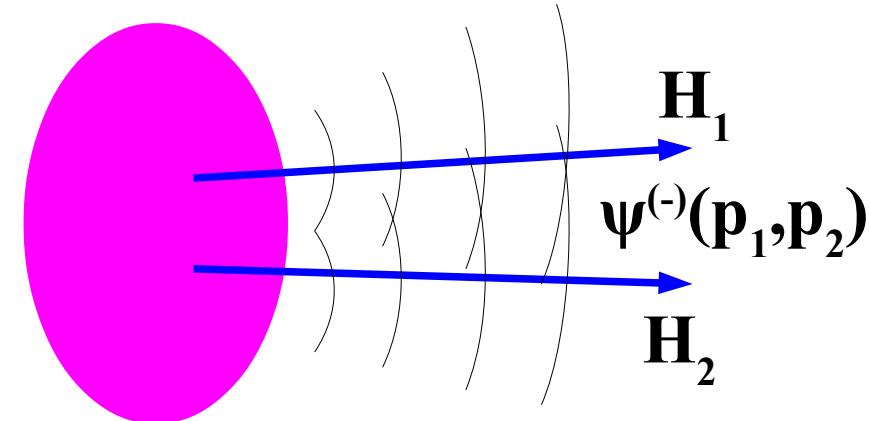


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Hadron-Hadron Correlation in HIC

■ Hadron-Hadron Correlation Func. (Koonin-Pratt (KP) formula)

*S. E. Koonin, PLB 70 ('77) 43; S. Pratt,
T. Csorgo and J. Zimanyi, PRC42 ('90) 2646;
W. Bauer, C.-K. Gelbke, S. Pratt, Annu. Rev.
Nucl. Part. Sci. 42 ('92) 77; R. Lednicky,
V. L. Lyuboshits, Sov. J. Nucl. Phys. 35 ('82) 770.*



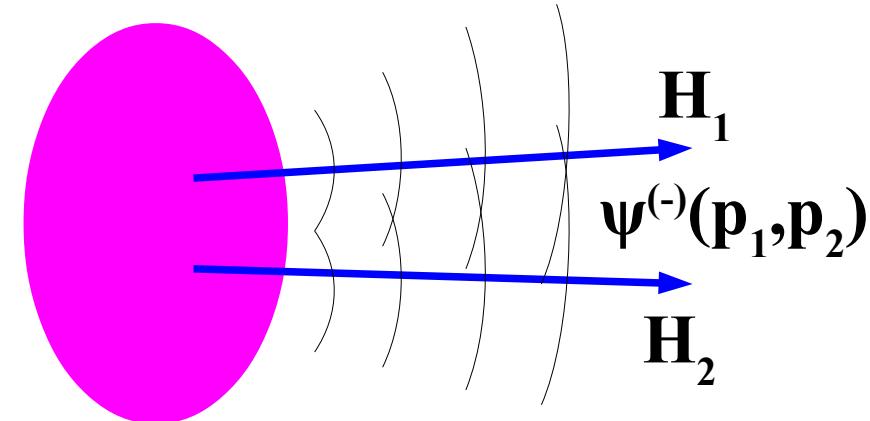
$$\begin{aligned} C(\mathbf{P}, \mathbf{q}) &= \frac{N(\mathbf{p}_1, \mathbf{p}_2)}{N(\mathbf{p}_1)N(\mathbf{p}_2)} = \frac{\int d^4x_1 \int d^4x_2 S(\mathbf{p}_1, x_1)S(\mathbf{p}_2, x_2)|\psi^{(-)}|^2}{\int d^4x_1 S(x_1, \mathbf{p}_1) \int d^4x_2 S(x_2, \mathbf{p}_2)} \\ &= \int d^4r S_{12}(\mathbf{q}, r) \left| \psi_{12}^{(-)}(\mathbf{r}'; \mathbf{q}) \right|^2 \\ &\simeq 1 - \frac{1}{2} \exp(-4q^2 R^2) + \frac{1}{2} \int d^3r S_{12}(\mathbf{r}) \left[|\chi_0(r)|^2 - |j_0(qr)|^2 \right] \end{aligned}$$

P: Total mom., q: rel. mom., Spherical source, Identical Fermion

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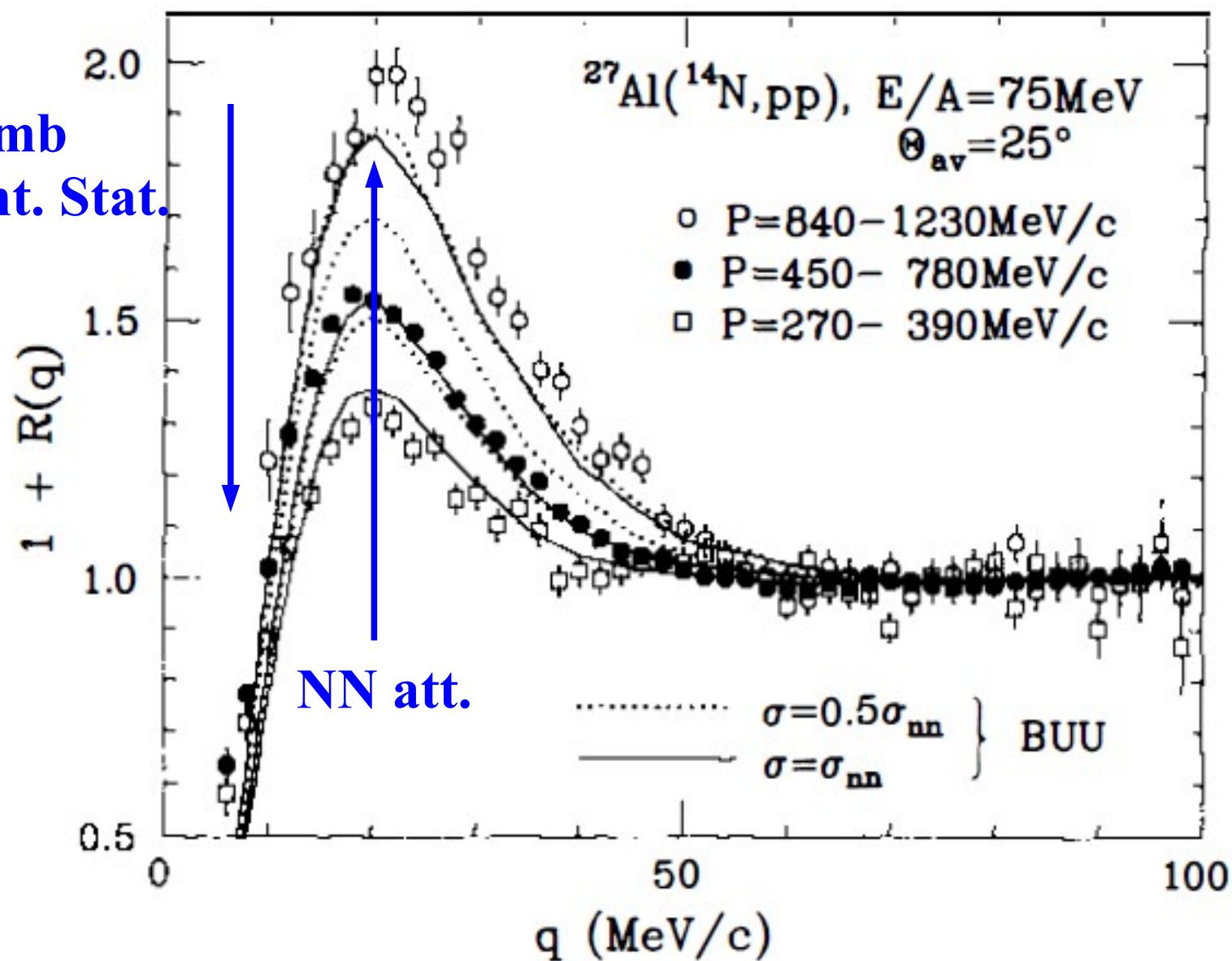


$$\begin{aligned}
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 &= \int d^4r S_{12}(\mathbf{q}, r) \left| \underline{\psi_{12}^{(-)}(\mathbf{r}'; \mathbf{q})} \right|^2 \quad \text{int.} \rightarrow \text{relative w.f.} \\
 &\simeq 1 - \frac{1}{2} \exp(-4q^2R^2) + \frac{1}{2} \int d^3r \underline{S_{12}(\mathbf{r})} \left[\underline{|\chi_0(r)|^2} - \underline{|j_0(qr)|^2} \right] \\
 &\qquad\qquad\qquad \text{Fermion} \qquad\qquad\qquad \text{Source w.f.} \qquad\qquad\qquad \text{free}
 \end{aligned}$$

P: Total mom., q: rel. mom., Spherical source, Identical Fermion

How does interaction modifies correlation ?

Coulomb
+Quant. Stat.



Bauer et al. ('92)

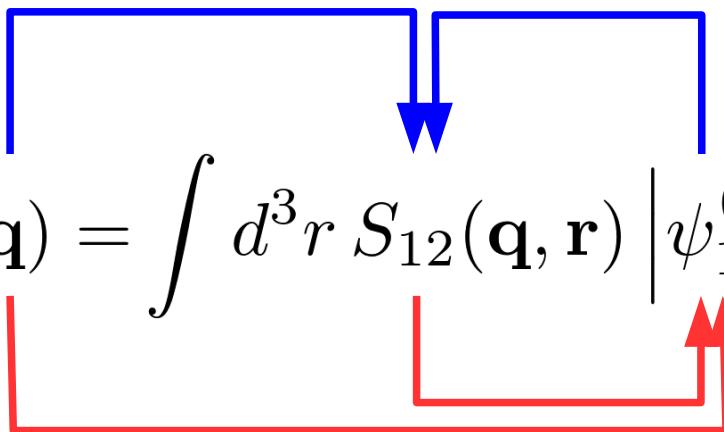
Standard Usage

Known Int. + Corr. Fn. → Source Size

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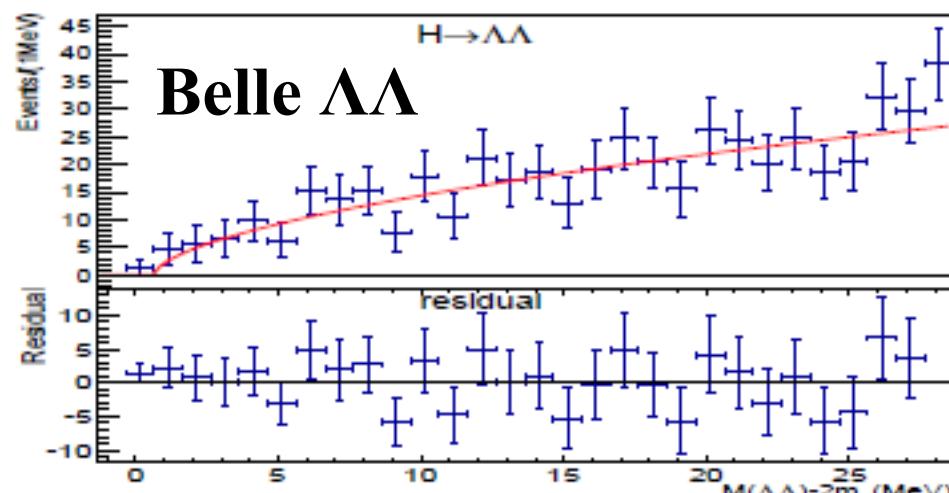
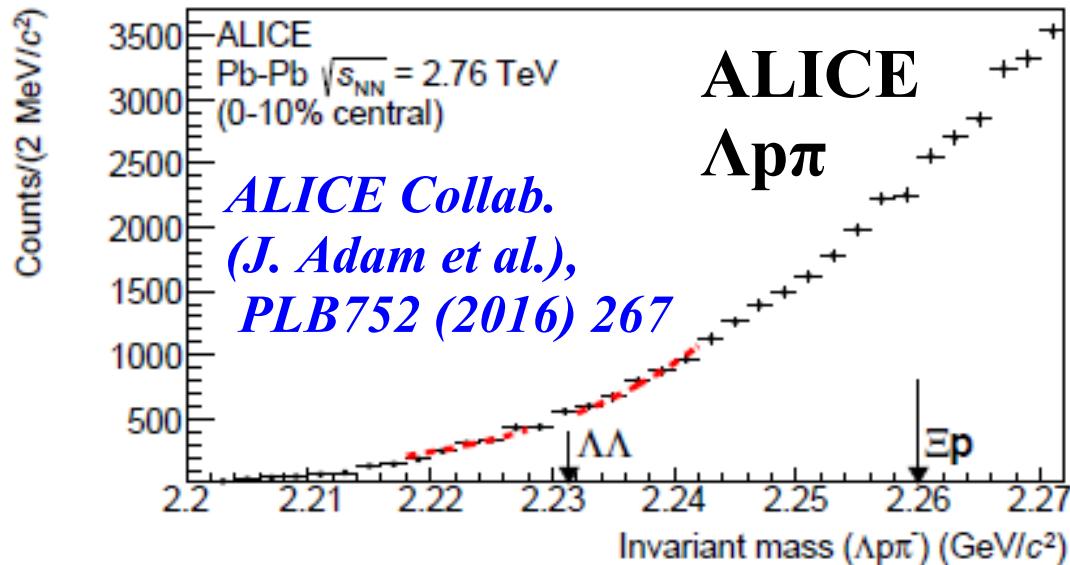
$$C(\mathbf{q}) = \int d^3r S_{12}(\mathbf{q}, \mathbf{r}) \left| \psi_{12}^{(-)}(\mathbf{r}; \mathbf{q}) \right|^2$$


*Idea of Reversal:
Can we determine hh interaction
from hh correlation ?*

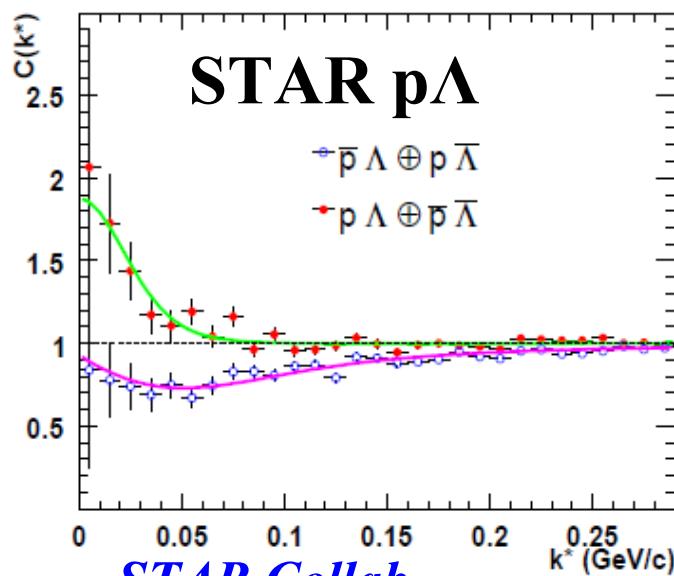
Let us try ! Examples: $\Lambda\Lambda$, $\Omega^- p$, and $K^- p$

*Previous works ($\Lambda\Lambda$): C. Greiner, B. Muller, PLB219('89)199.;
AO, Y. Hirata, Y. Nara, S. Shinmura, Y. Akaishi , NPA 670 ('00), 297c*

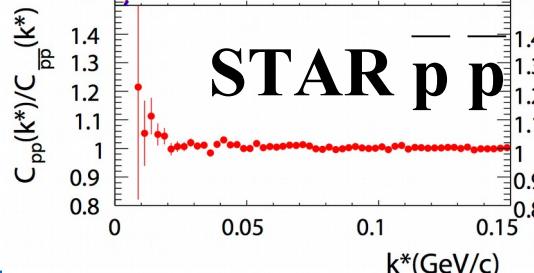
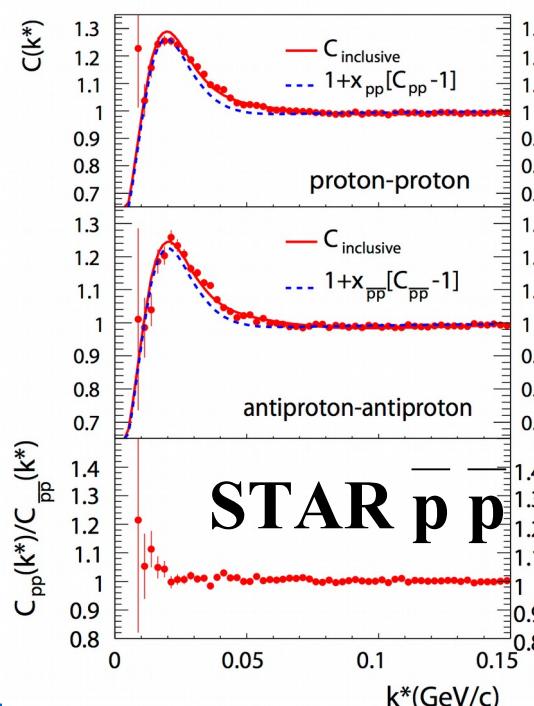
$\Lambda\bar{\Lambda}$ invariant mass / BB correlation function



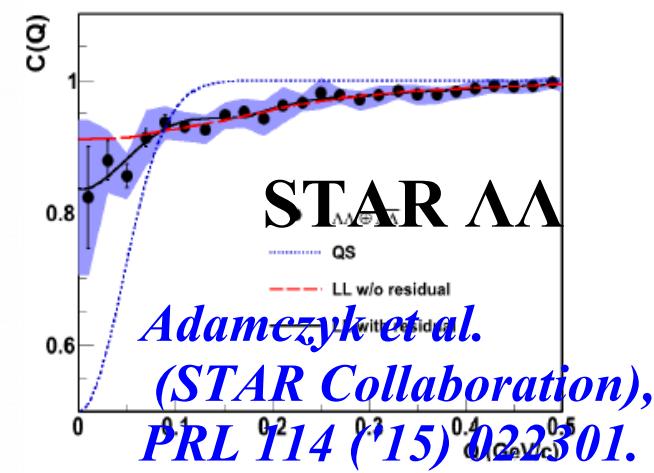
Belle Collaboration (Kim, B.H. et al.),
PRL110('13)222002.



STAR Collab.
(J. Adams et al.),
PRC74('06)064906.



STAR Collab.
(J. Adams et al.)
Nature 527('15)345



Contents

- **Introduction**
- **Interaction dependence of two-particle momentum correlation**
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 - **Lednicky-Lyuboshits (LL) model**
 - **Interaction dependence**
- **Hadron-hadron correlation & interaction**
 - **$\Lambda\Lambda$ correlation and $\Lambda\Lambda$ interaction**
 - **$\Omega^- p$ correlation and Coulomb potential effects**
 - **$K^- p$ correlation and correlation in couple channels**
- **Summary**

Interaction dependence of two particle momentum correlation

Correlation function formula (for identical Fermions)

■ Koonin-Pratt formula (Spherical Gaussian Source)

$$C(\mathbf{q}) = \int d^3r S_{12}(\mathbf{q}, \mathbf{r}) \left| \underline{\psi_{12}^{(-)}(\mathbf{r}; \mathbf{q})} \right|^2 \text{ int.} \rightarrow \text{relative w.f.}$$

$$\simeq 1 - \frac{1}{2} \exp(-4q^2 R^2) + \frac{1}{2} \int d^3r S_{12}(\mathbf{r}) \left[\underline{|\chi_0(r)|^2} - \underline{|j_0(qr)|^2} \right]$$

Fermion	$= C_{\text{HBT}}(q) + \frac{1}{2} \Delta C(q)$	Source	w.f.
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Fermion	$= C_{\text{HBT}}(q) + \frac{1}{2} \boxed{\Delta C(q)}$	Source w.f.	free
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■ Lednicky-Lyuboshits analytic model

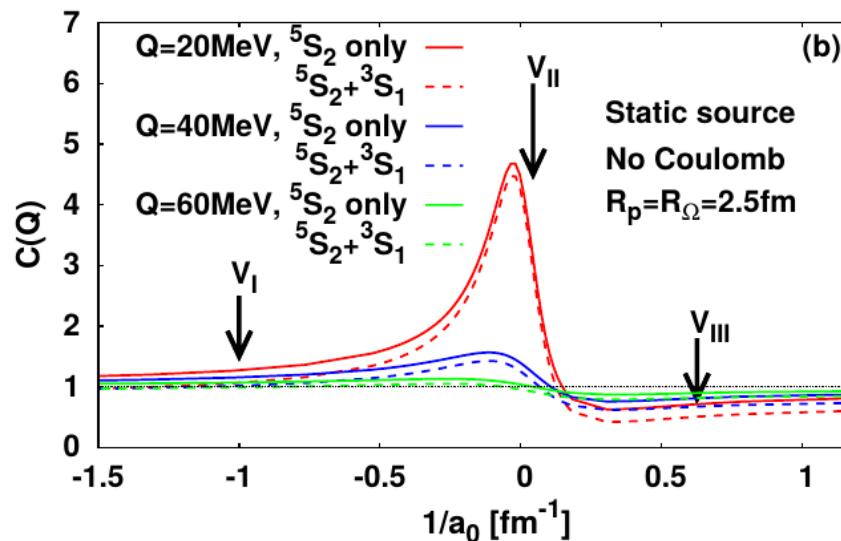
- Asymp. w.f. + Eff. range corr. + $\psi^{(-)} = [\psi^{(+)}]^*$

$$\psi_{\text{asy}}(r) = \frac{e^{-i\delta}}{qr} \sin(qr + \delta) = \mathcal{S}^{-1} \left[\frac{\sin qr}{qr} + f(q) \frac{e^{iqr}}{r} \right]$$

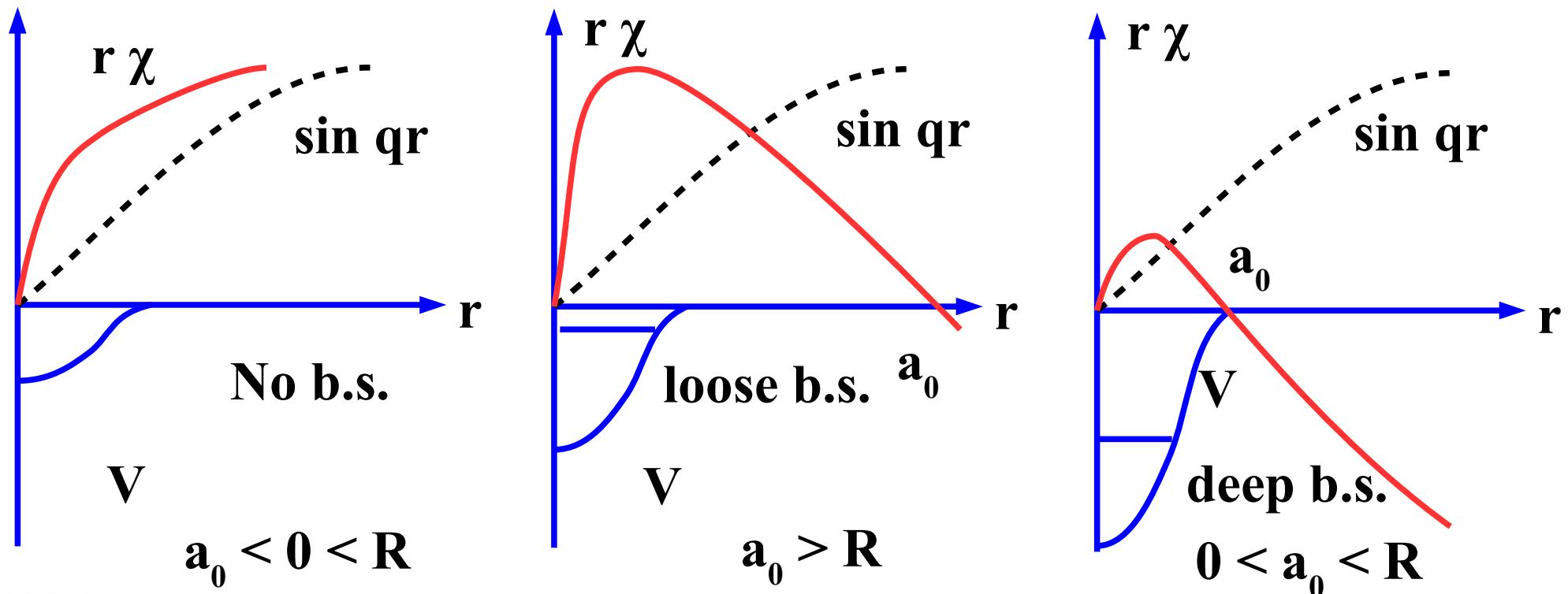
$$\boxed{\Delta C_{\text{LL}}(q)} = \frac{|f(q)|^2}{2R^2} F_3 \left(\frac{r_{\text{eff}}}{R} \right) + \frac{2\text{Re}f(q)}{\sqrt{\pi}R} F_1(x) - \frac{\text{Im}f(q)}{R} F_2(x)$$

$(x = 2qR, R = \text{Gaussian size}, F_1, F_2, F_3 : \text{Known function})$

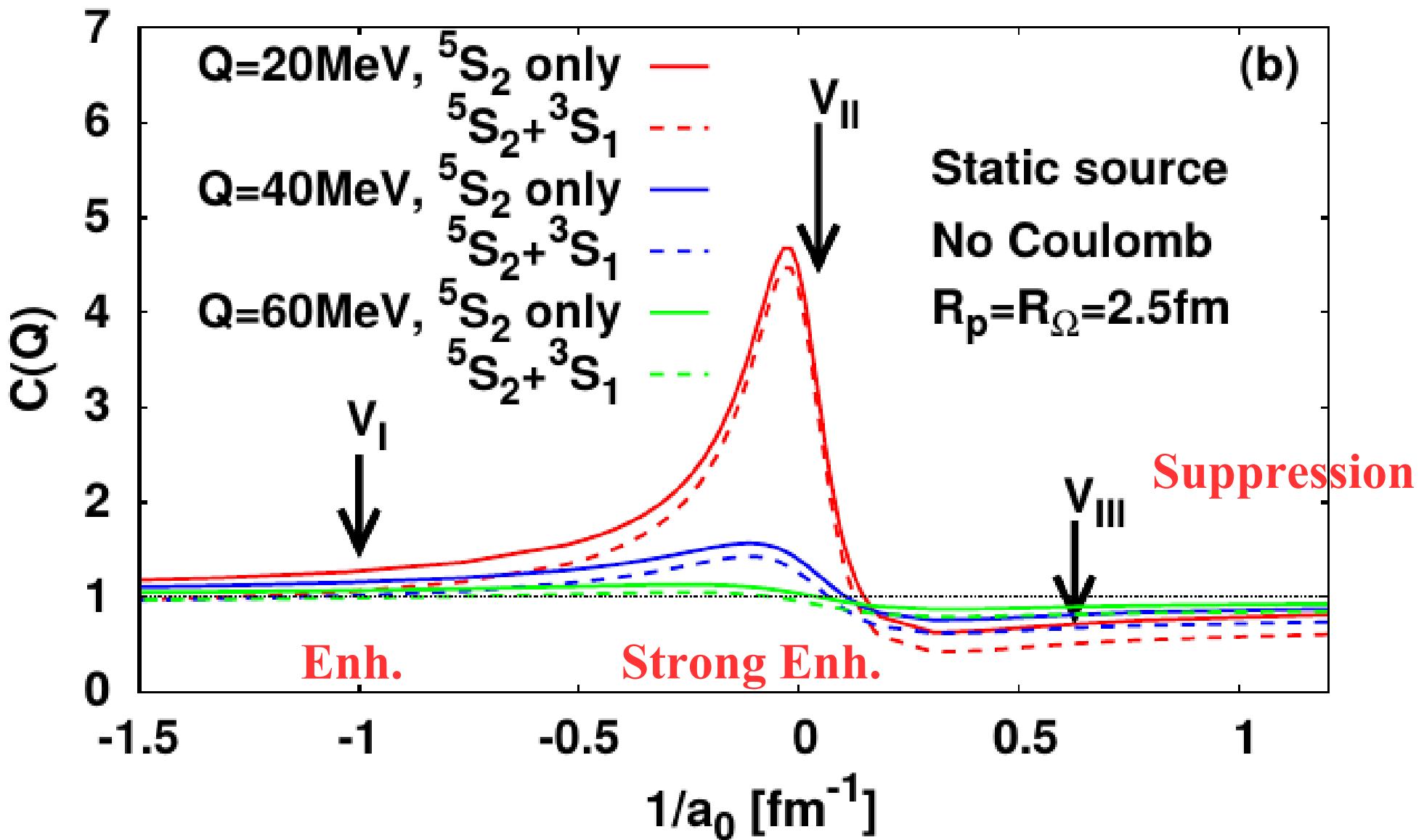
Interaction Dependence of Correlation Function



$$q \cot \delta = -1/a_0 + r_{\text{eff}} q^2/2 + O(q^4)$$



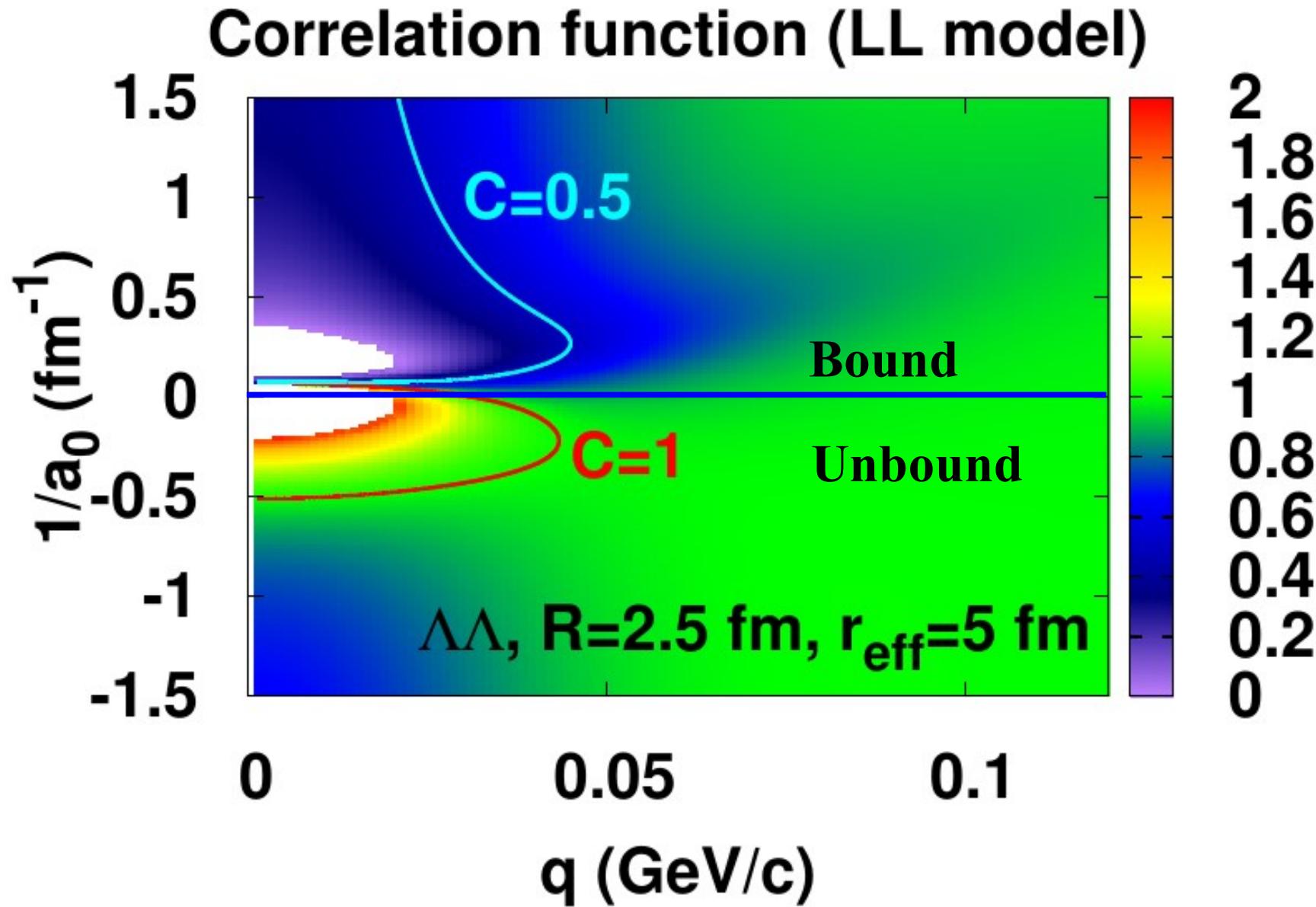
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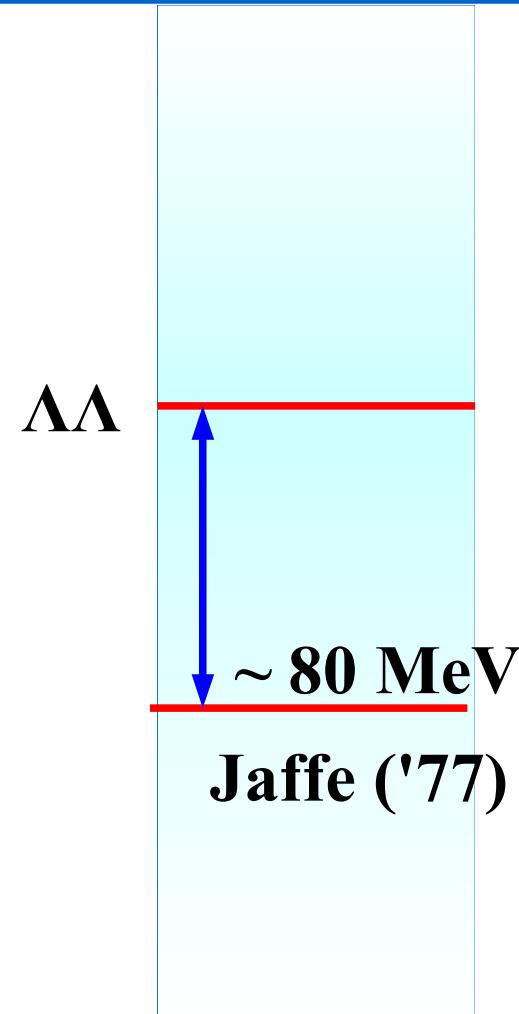
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AO, K. Morita, K. Miyahara, T. Hyodo, NPA954 ('16), 294.

$\Lambda\Lambda$ correlation and $\Lambda\Lambda$ interaction

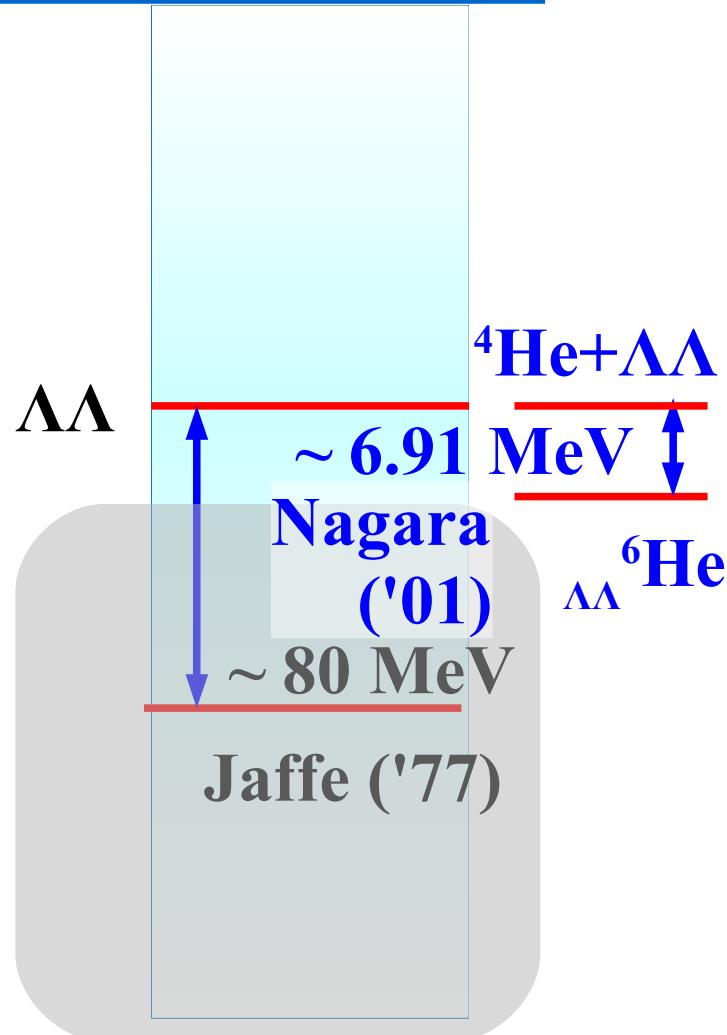
Relevance of $\Lambda\Lambda$ interaction to physics

- H-particle: 6-quark state (uuddss)
 - Prediction: *R.L.Jaffe, PRL38(1977)195*



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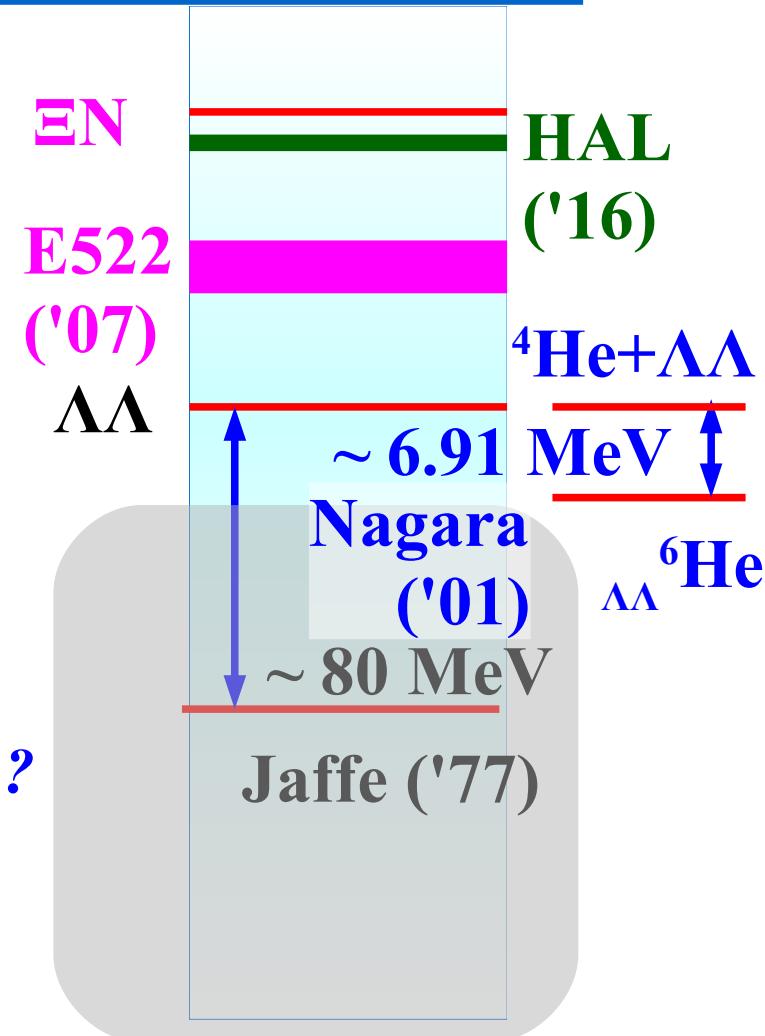
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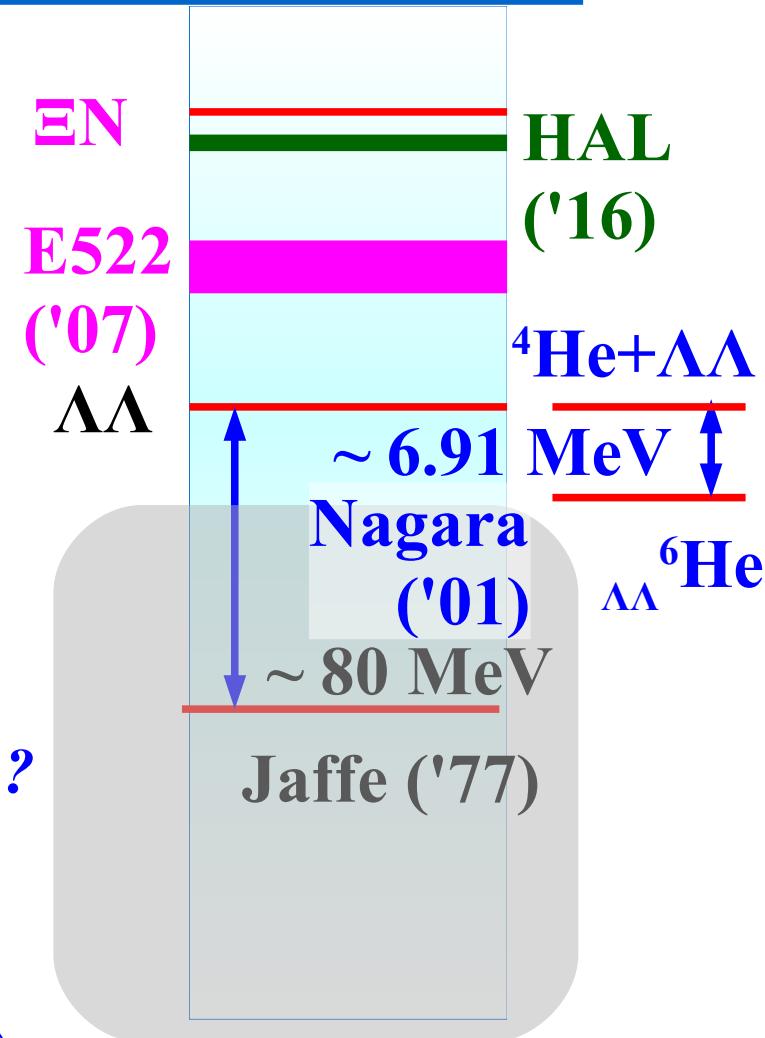
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- Resonance or Bound “H” ?
Yoon et al.(KEK-E522)+AO ('07)
- Lattice QCD
HAL QCD & NPLQCD ('11)
HAL QCD ('16): H as a loosely bound ΞN ?



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■ Neutron Star Matter EOS

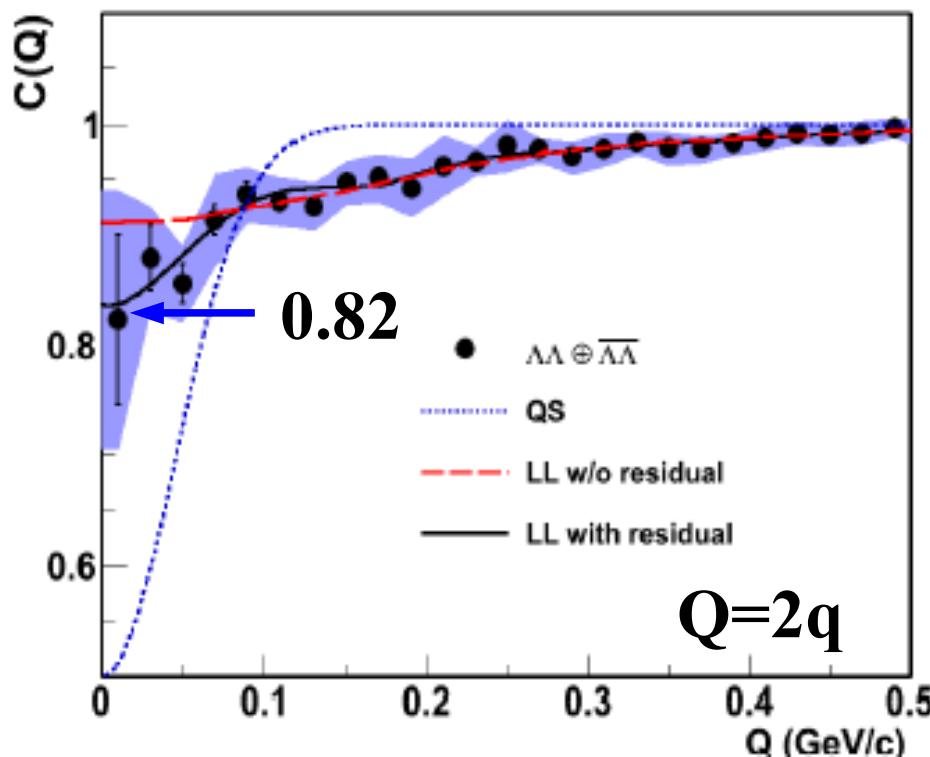
- Hyperon Puzzle
Demorest et al. ('10), Antoniadis et al. ('13)
- Cooling Puzzle ($\Lambda\Lambda$ superfluidity)
T. Takatsuka, R. Tamagaki, PTP 112('04)37

$\Lambda\bar{\Lambda}$ correlation at RHIC

- STAR collaboration at RHIC measured $\Lambda\bar{\Lambda}$ correlation !

Adamczyk et al. (STAR Collaboration), PRL 114 ('15) 022301.

- RHIC, Au+Au ($\sqrt{s_{NN}}=200$ GeV), Weak decay vertex analysis.



$\Lambda\bar{\Lambda}$ correlation at RHIC

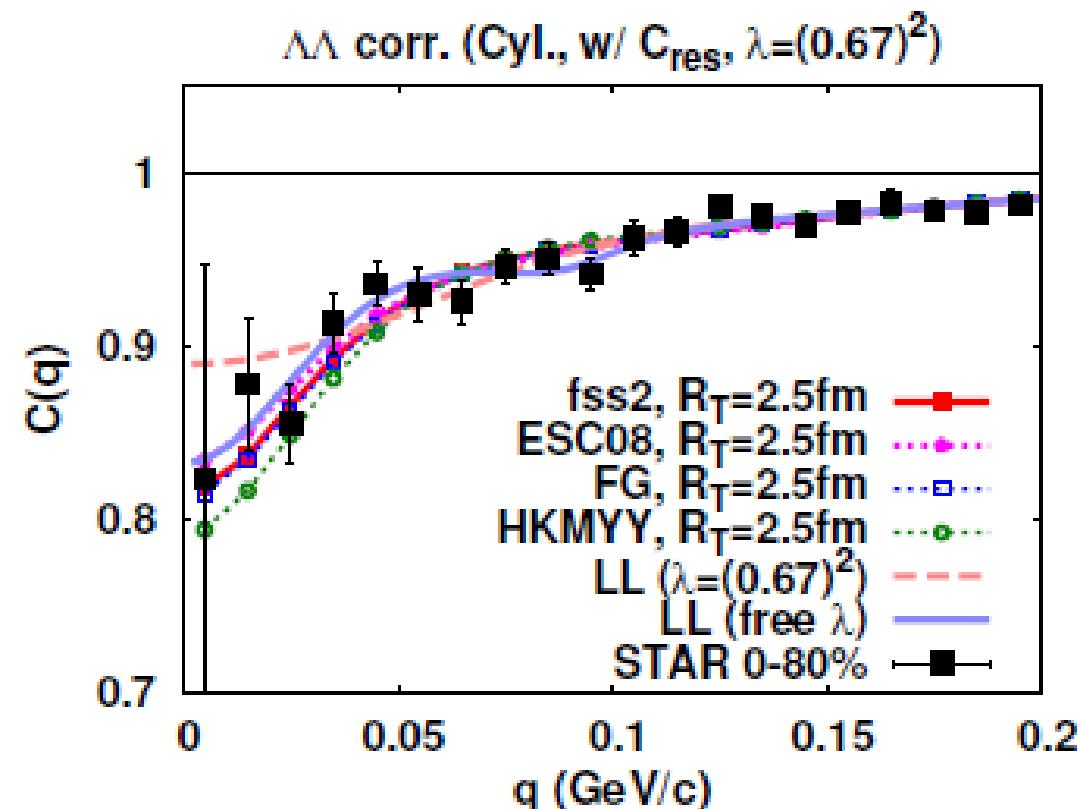
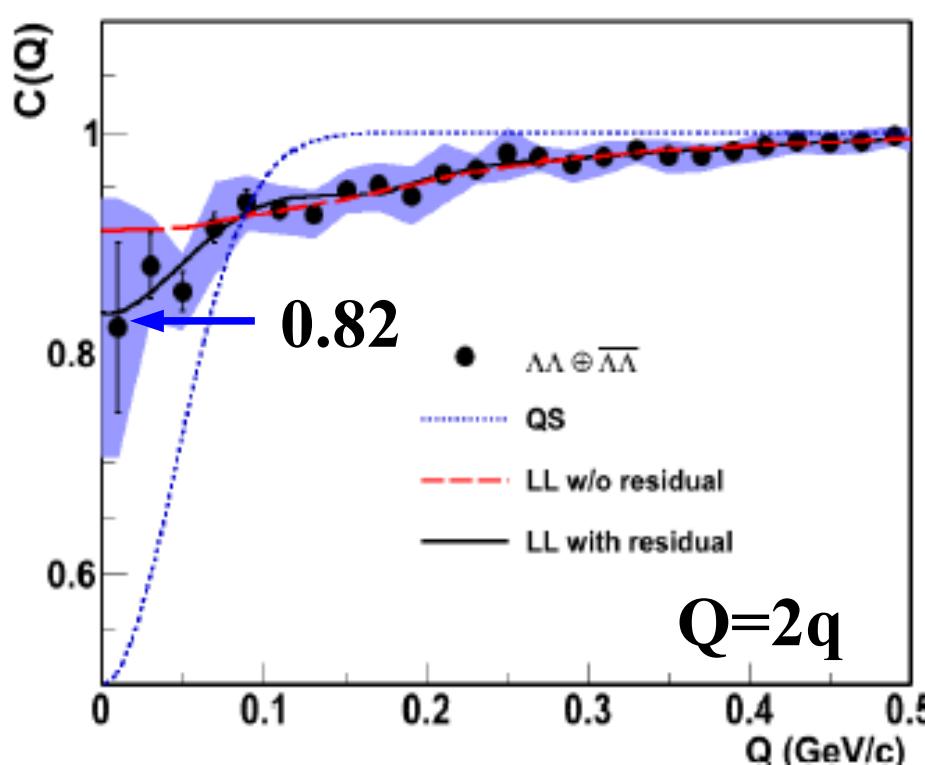
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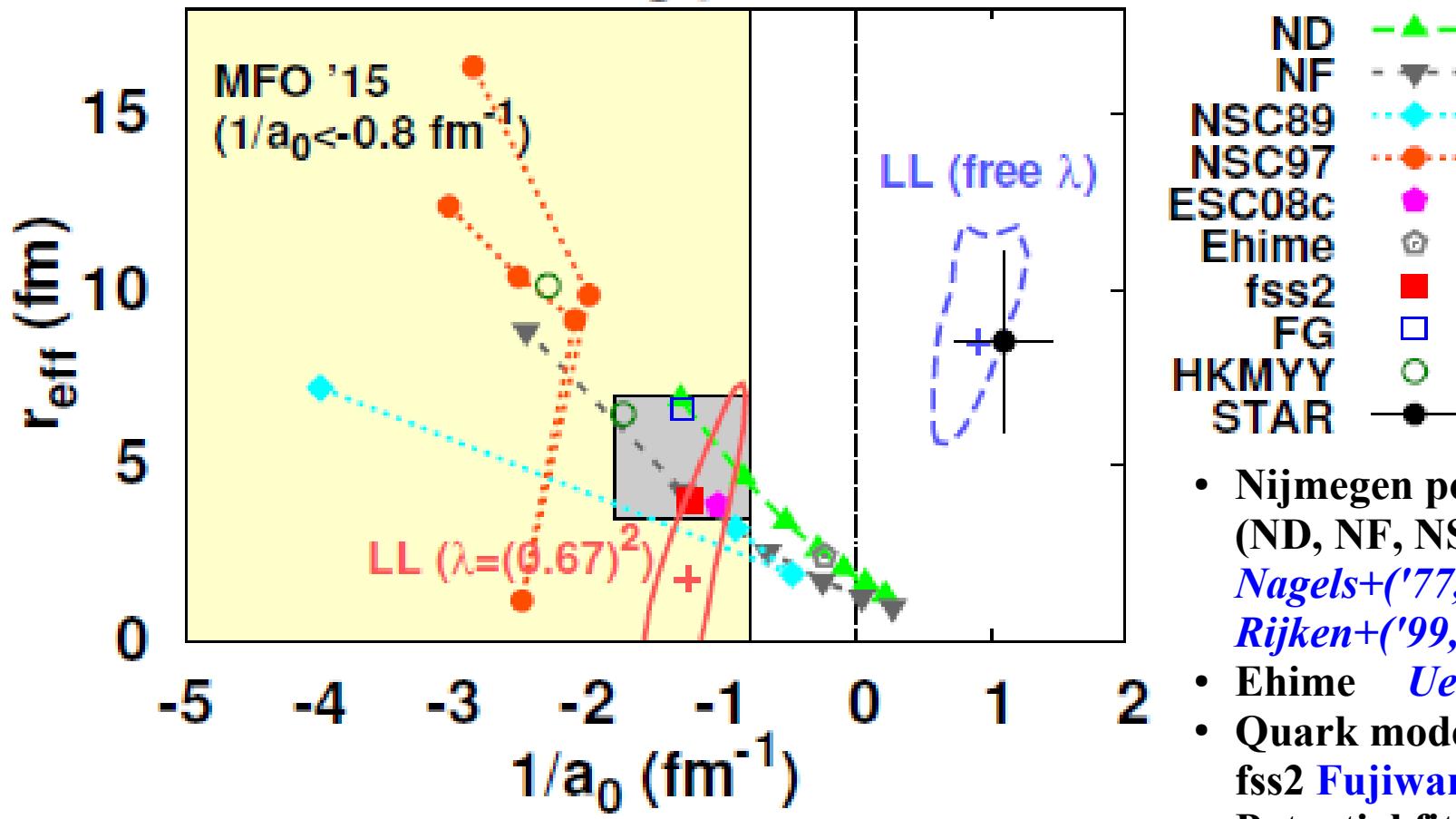
- Theoretical Analysis well explains the data

*K.Morita et al., T.Furumoto, AO, PRC91('15)024916;
AO, K.Morita, K.Miyahara, T.Hyodo, NPA954 ('16), 294.*



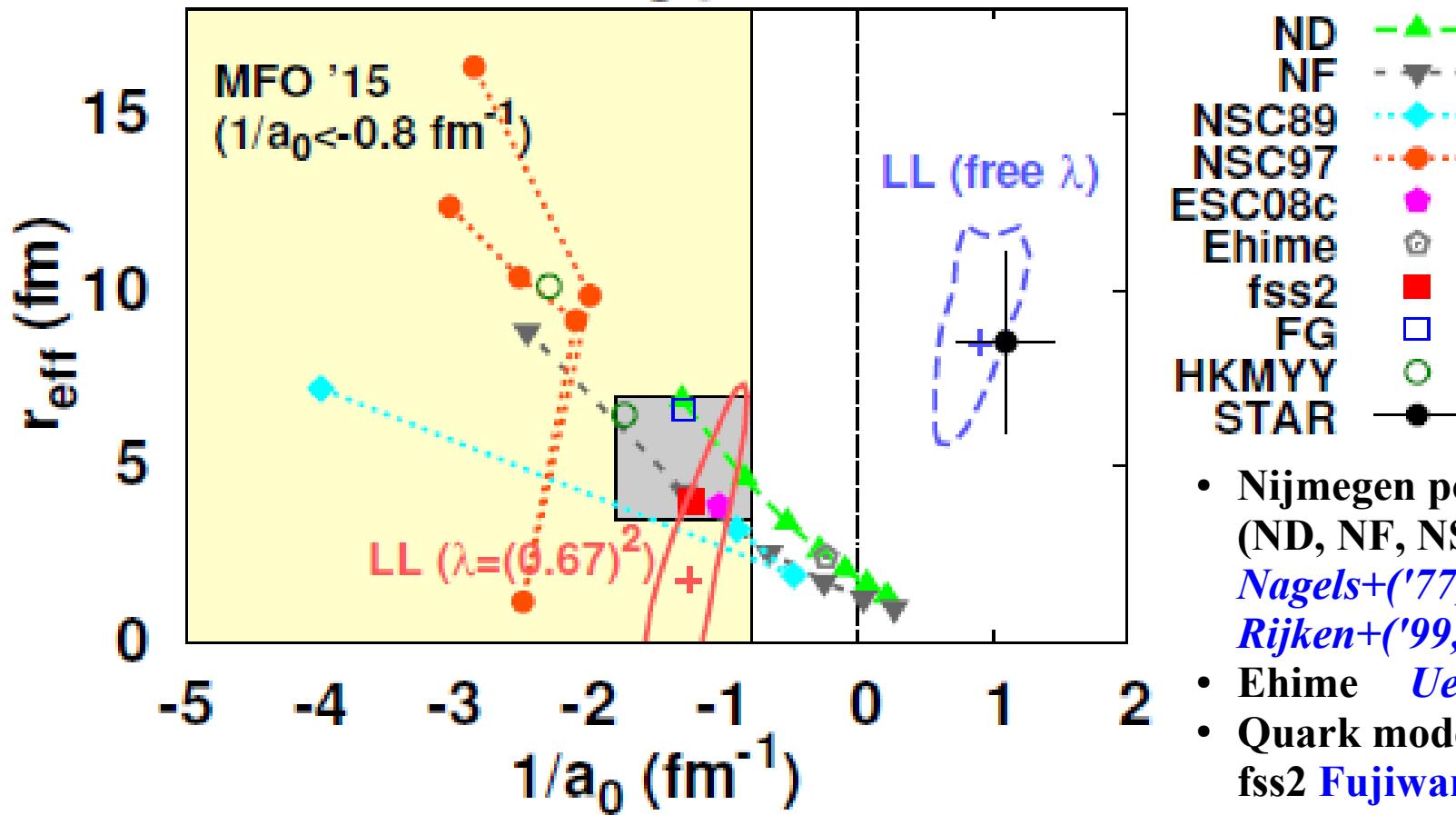
$\Lambda\Lambda$ interaction from $\Lambda\Lambda$ correlation

$\Lambda\Lambda$ scattering parameters



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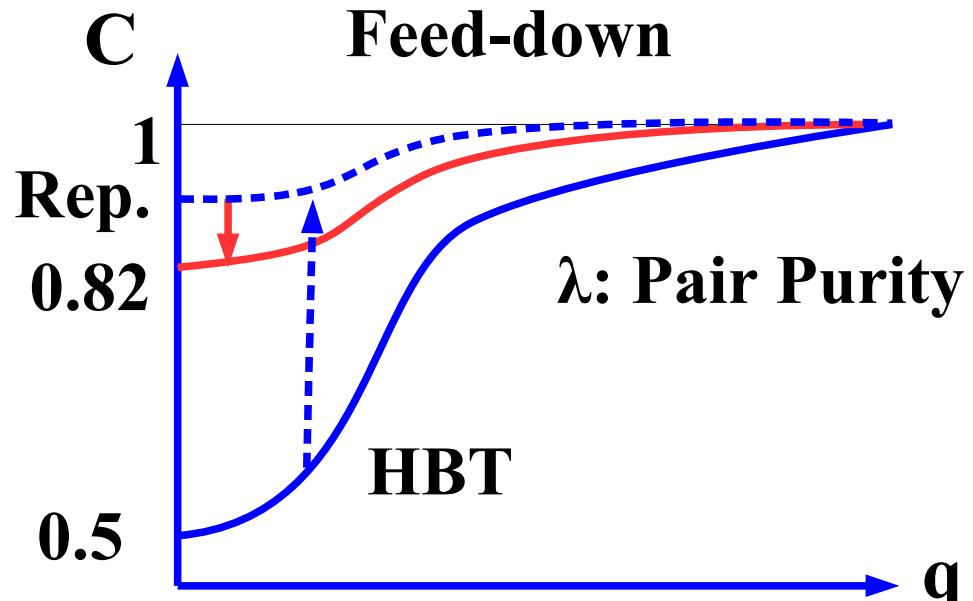
- Nijmegen potentials
(ND, NF, NSC89, NSC97, ESC08)
Nagels+ ('77, '79), Maessen+ ('89), Rijken+ ('99, '10)
- Ehime *Ueda et al. ('98)*
- Quark model interaction:
fss2 *Fujiwara et al. ('07)*
- Potential fitted to Nagara
Filiakin, Gal ('02) (FG), Hiyama et al. ('02, '10) (HKMYY)

$\text{Positive } a_0 \text{ (STAR)} \longleftrightarrow \text{Negative } a_0 \text{ (MFO'15)}$

Feed-Down Effects & Residual Source

- Correlation Fn. w/ Feed-down & Residual source effects.

$$C_{\text{corr}}(q) = 1 + \lambda(C_{\text{bare}}(q) - 1) + a_{\text{res}} \exp(-4r_{\text{res}}^2 q^2)$$



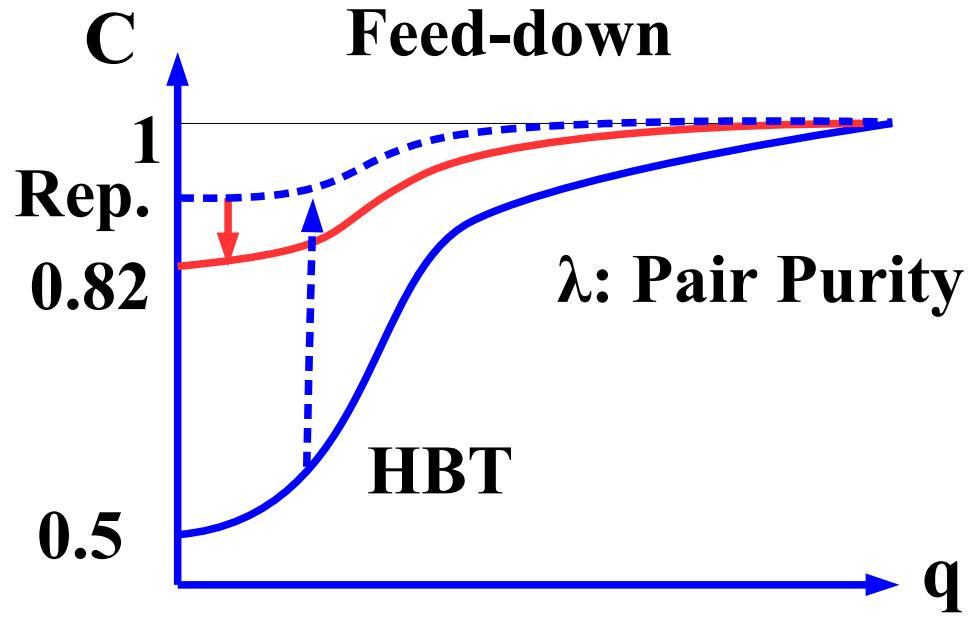
STAR:
 $\lambda \sim 0.18$ (free para.)

AO, Morita, Mihayara, Hyodo ('16)

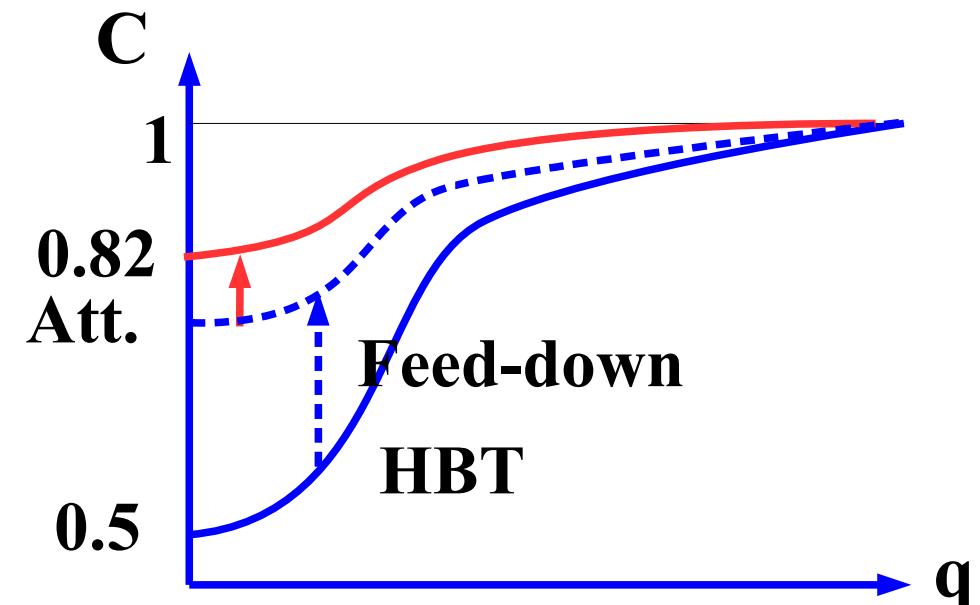
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Morita et al. (MFO15):
 $\lambda \sim 0.45$

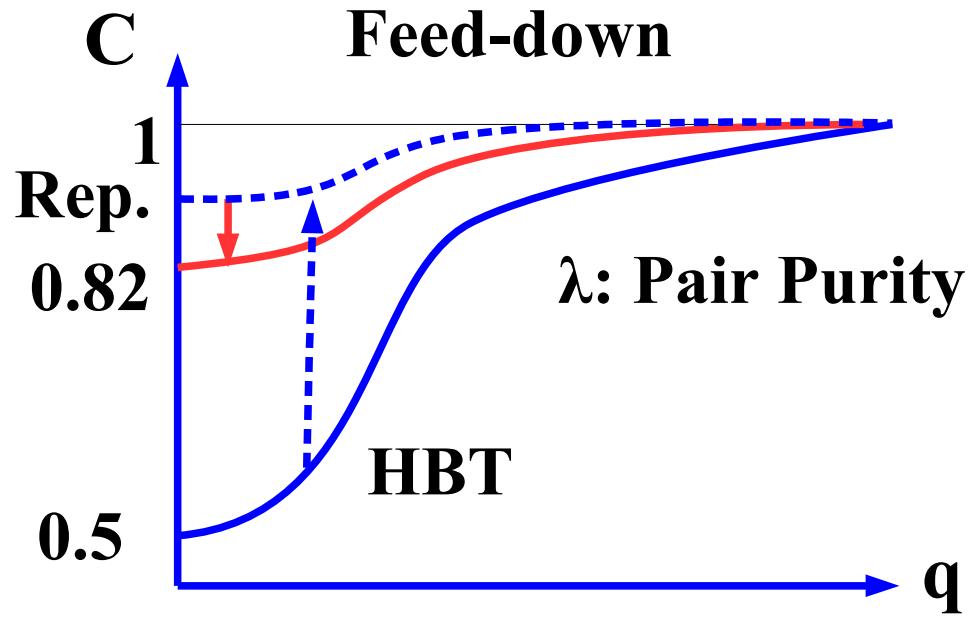
$\Sigma^0/\Lambda = 0.278$ (p+Be, 28.5 GeV/c)
Sullivan et al. ('87)
 $\Xi/\Lambda = 15\%$ (RHIC)

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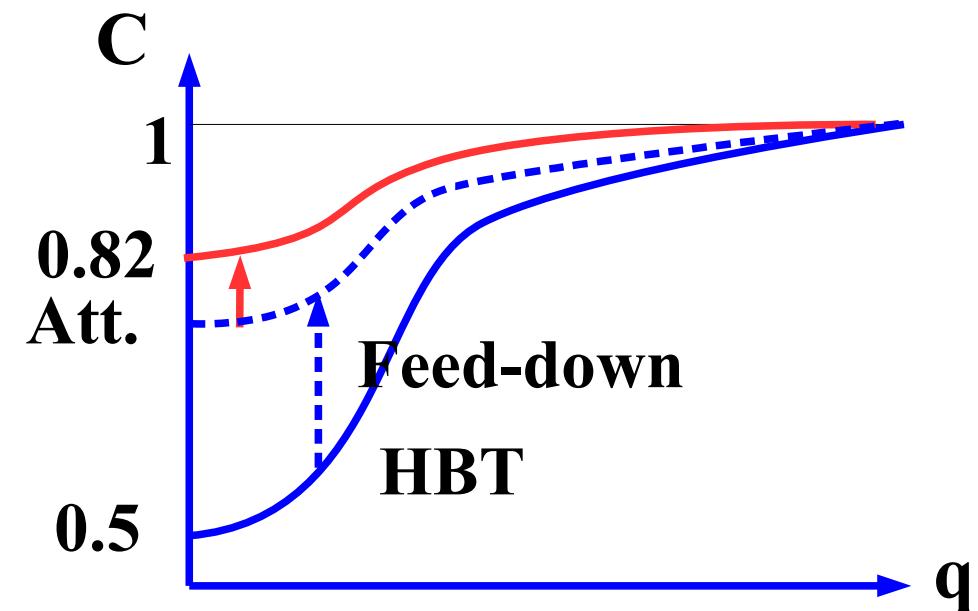
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*Pair purity (λ) should be determined experimentally !
Puzzle: Residual source*

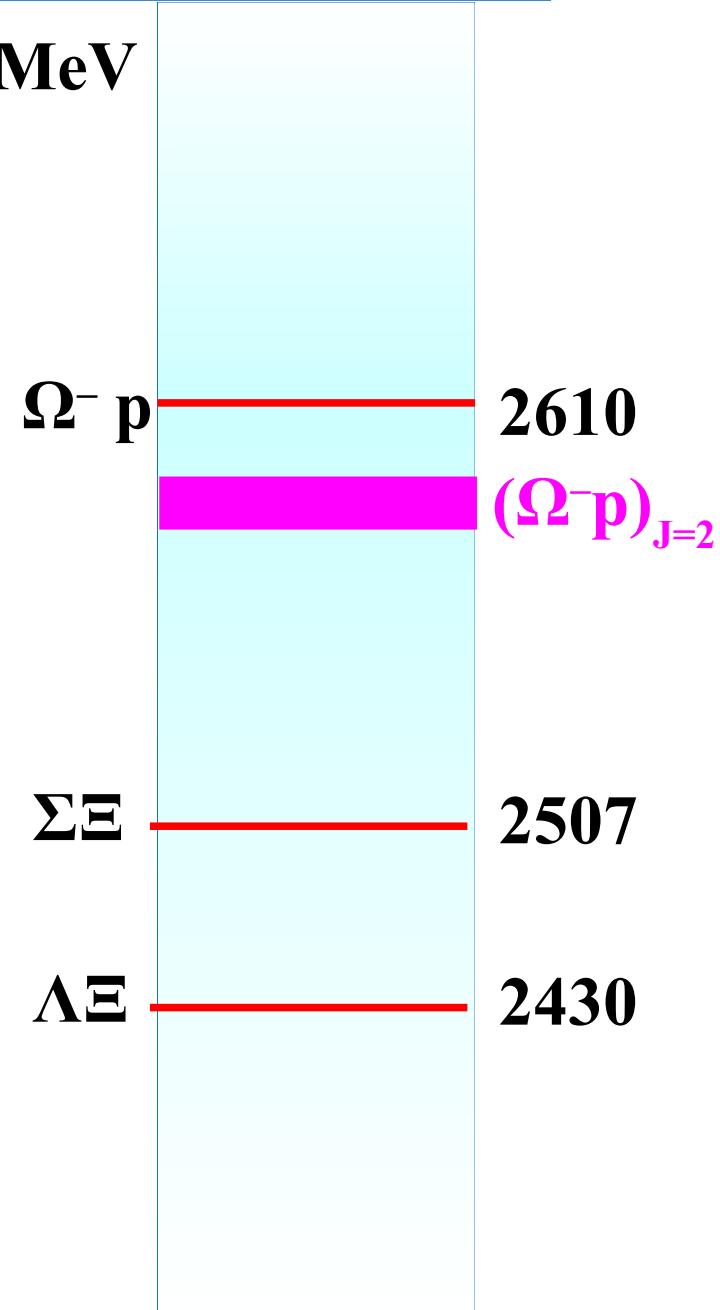
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$\Omega^- p$ correlation

$\Omega^- p$ interaction

- Ω^- : quark content=sss, $J\pi=3/2+$, $M=1672$ MeV
- $\Omega^- p$ bound state as a $S=-3$ dibaryon ?
F.Etminan et al. (HAL QCD Collab.), NPA928('14)89.

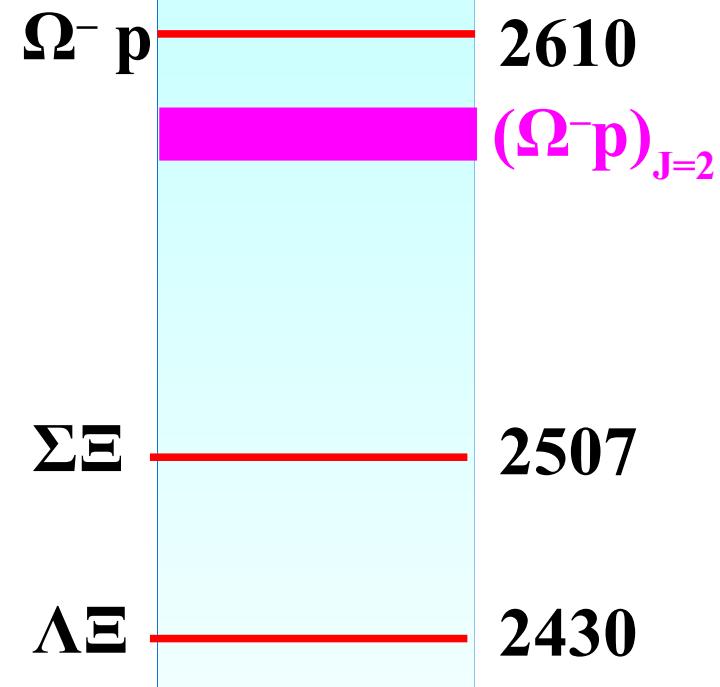
- No quark Pauli blocking in ΩN and $H=uuddss$ channels.
Oka ('88), Gal ('16)
- $J=2$ state (5S_2) couples to Octet-Octet baryon pair only with $L \geq 2$
→ Small width is expected.
Etminan et al. (HAL QCD)('14)
- Correlation is measurable at RHIC !
Neha Shah (STAR), private commun.



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*Let us try to discover the first(?) dibaryon (after deuteron) !
(First dibaryon with $S<0$!)*

$\Omega^- p$ potential from lattice QCD

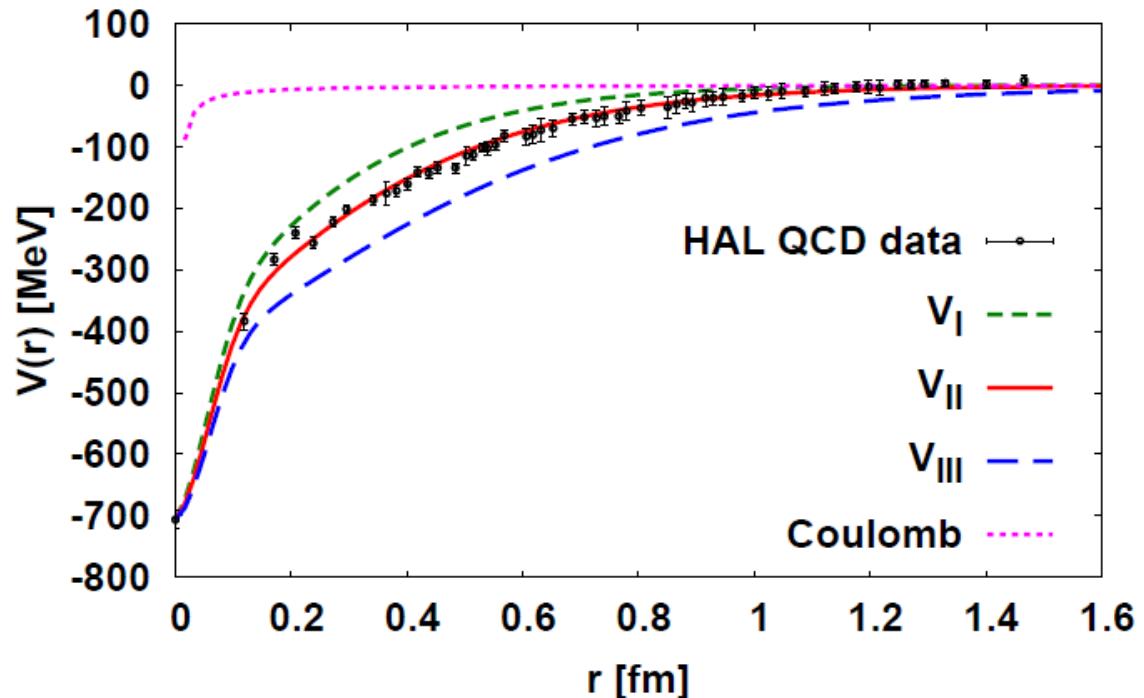
- Lattice QCD predicts $\Omega^- p$ bound state at large quark mass, $m\pi=875$ MeV (B.E.~ ~ 19 MeV) in 5S_2 channel.

F.Etminan et al. (HAL QCD Collab.), NPA928('14)89.

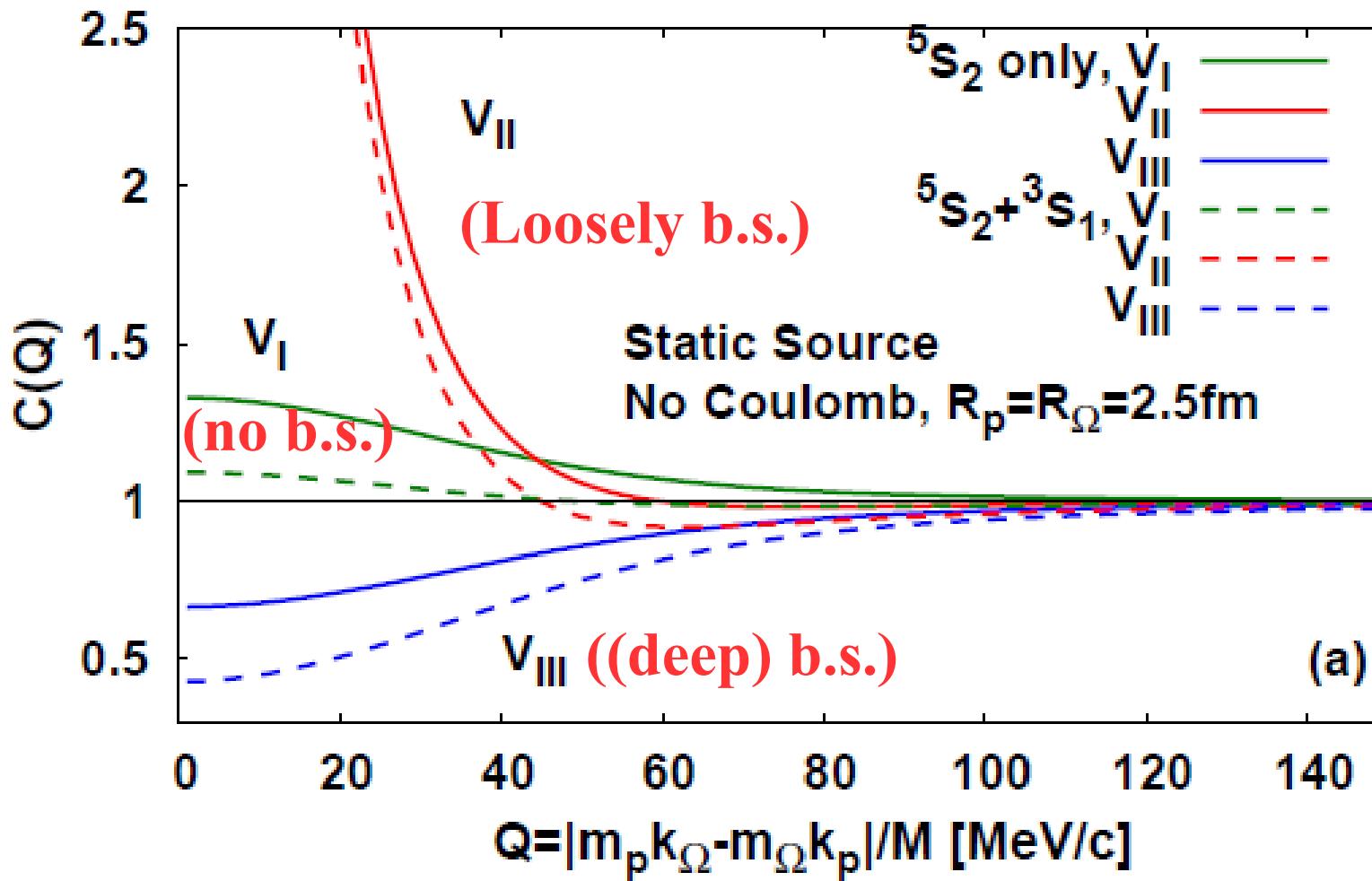
- Extrapolation to physical quark mass

- VI → Weaker potential (no b.s.)
- VII → Same potential (shallow b.s.)
- VIII → Stronger potential (deep b.s.)

Spin-2 $N\Omega$ Potentials		V_I	V_{II}	V_{III}
	E_B [MeV]	—	0.05	24.8
without Coulomb	a_0 [fm]	-1.0	23.1	1.60
	r_{eff} [fm]	1.15	0.95	0.65
	E_B [MeV]	—	6.3	26.9
with Coulomb	a_0 [fm]	-1.12	5.79	1.29
	r_{eff} [fm]	1.16	0.96	0.65



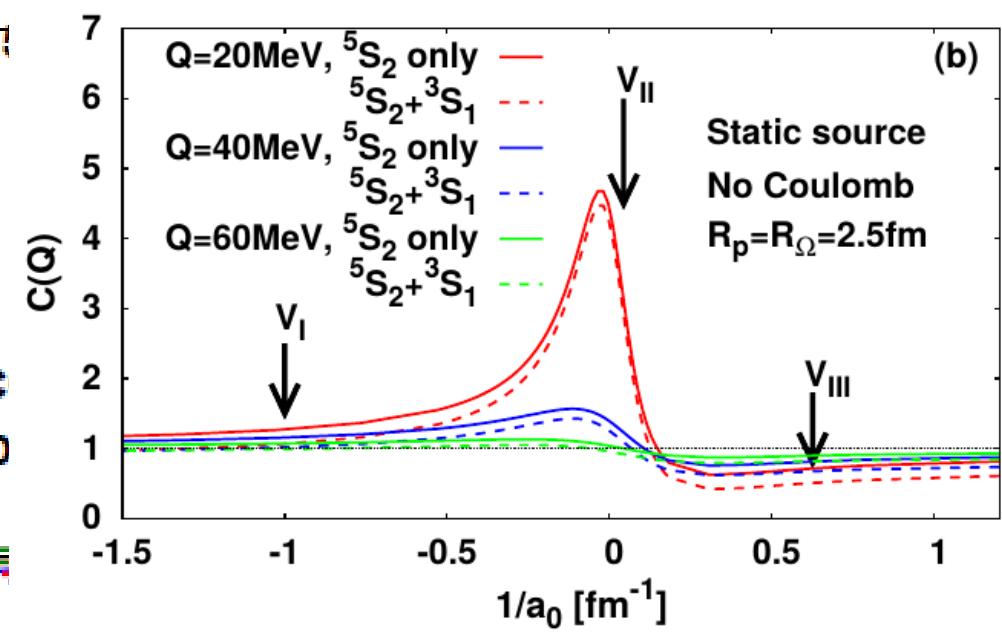
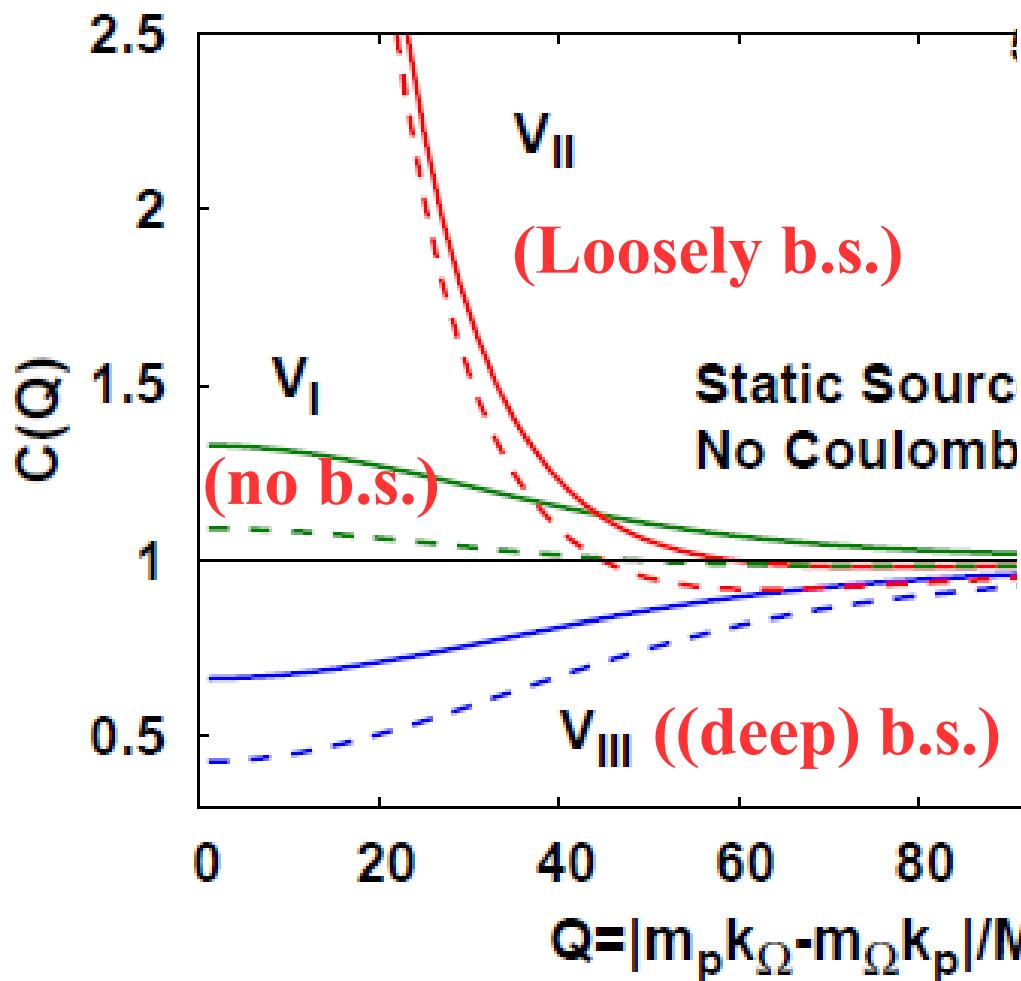
$\Omega^- p$ correlation



(w/o Coulomb, Strong absorption at $r < 2$ fm in 3S_1 (decay to 8-8 in S-wave))

K. Morita, AO, F. Etminan, T. Hatsuda, PRC94('16)031901(R) [arXiv:1605.06765 [hep-ph]]

$\Omega^- p$ correlation

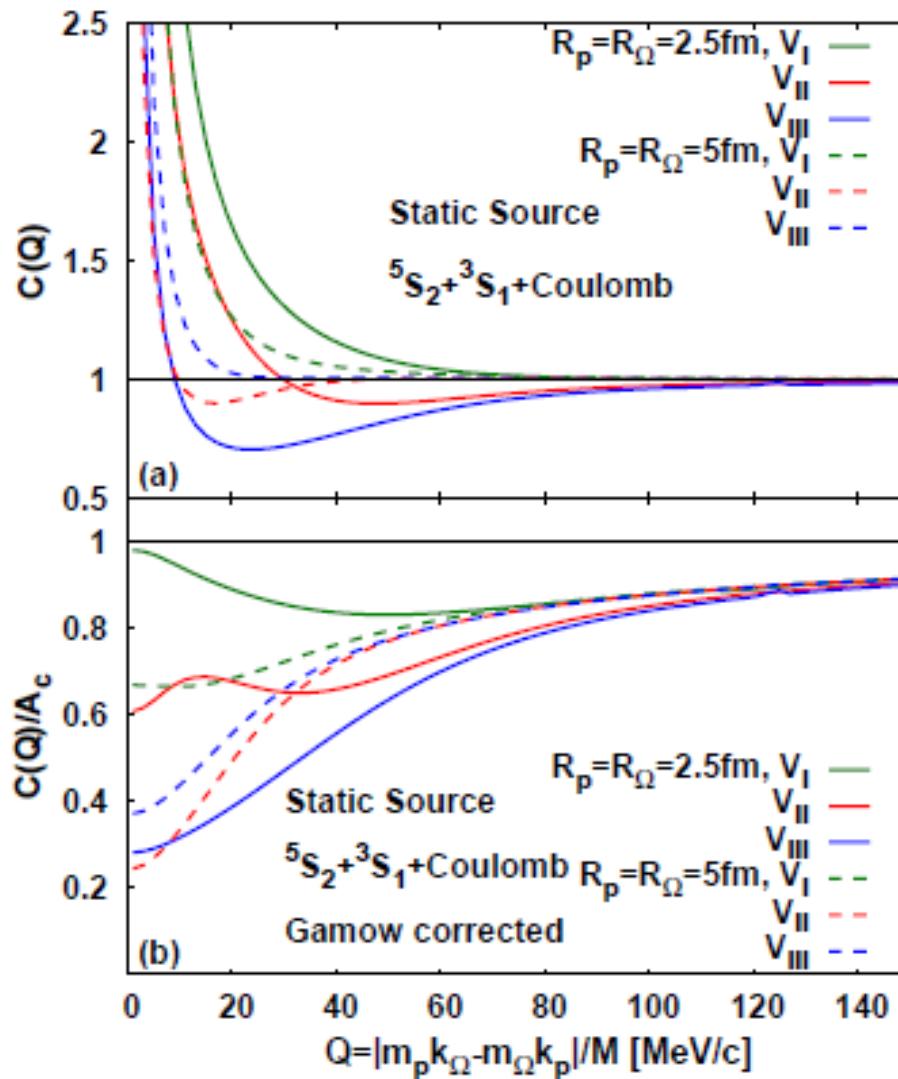


(a)

(w/o Coulomb, Strong absorption at $r < 2$ fm in ${}^3\text{S}_1$ (decay to 8-8 in S-wave))

K. Morita, A.O., F. Etminan, T. Hatsuda, PRC94('16)031901(R) [arXiv:1605.06765 [hep-ph]]

$\Omega^- p$ correlation w/ Coulomb

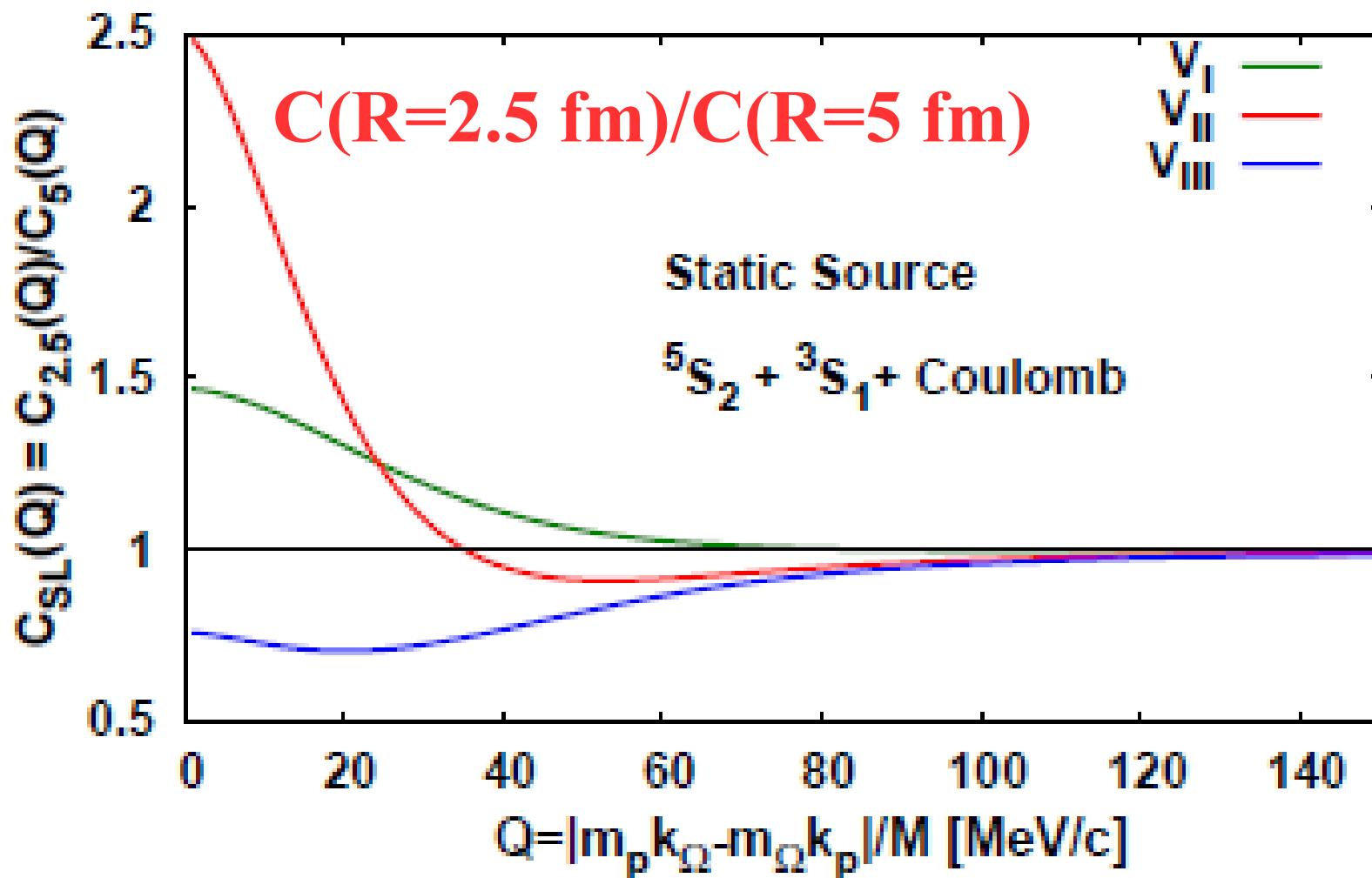


With Coulomb

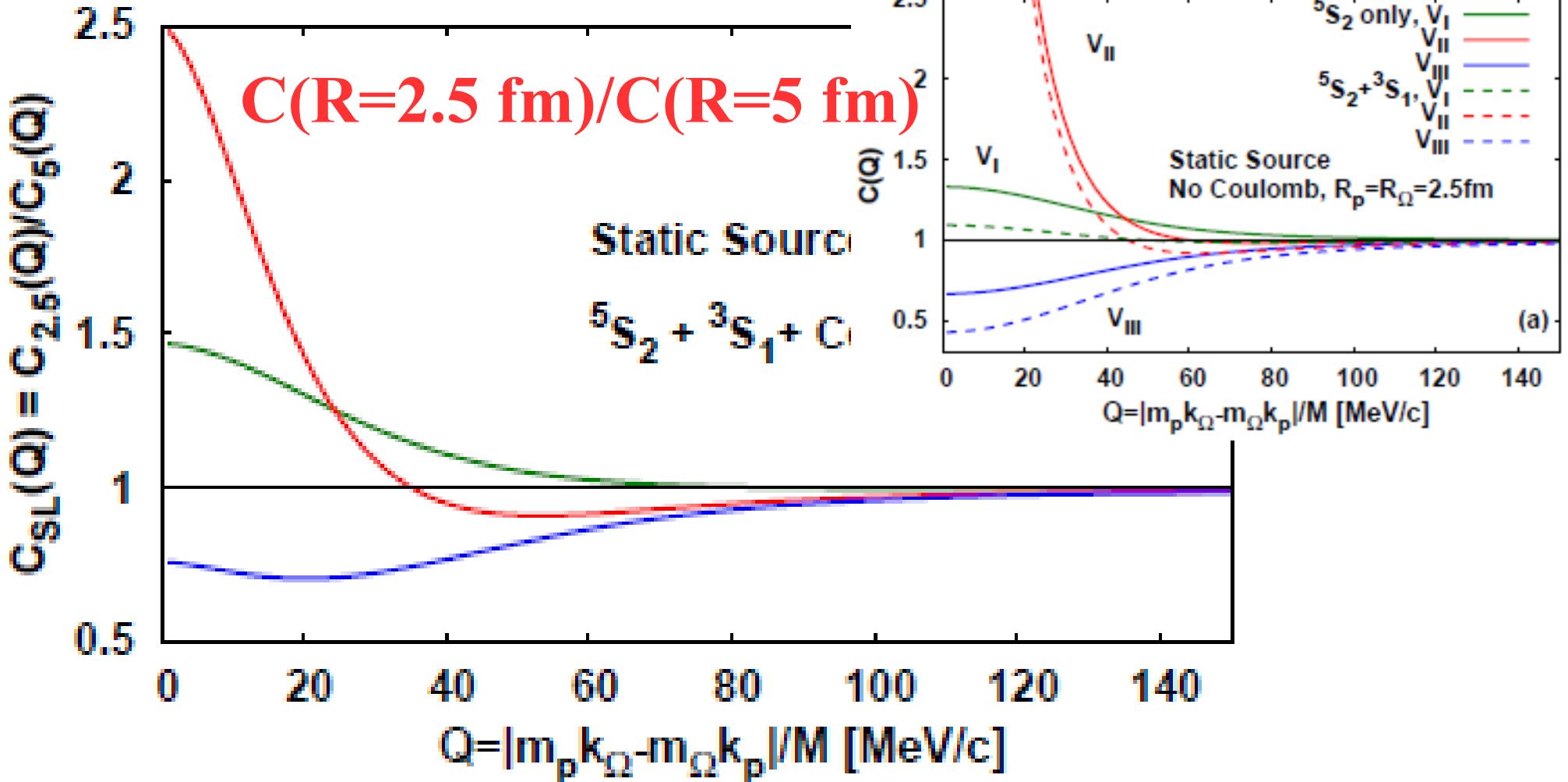
Coulomb + Gamow corr.

Coulomb potential washes out the features of V_I , V_{II} , V_{III} , and Gamow correction is not enough.

$\Omega^- p$ correlation: Small / Large Ratio



$\Omega^- p$ correlation: Small / Large Ratio



*By taking small ($R=2.5$ fm) / large ($R=5$ fm) ratio,
we approximately see the corr. fn. w/o Coulomb !*

K⁻p correlation

K⁻ p interaction

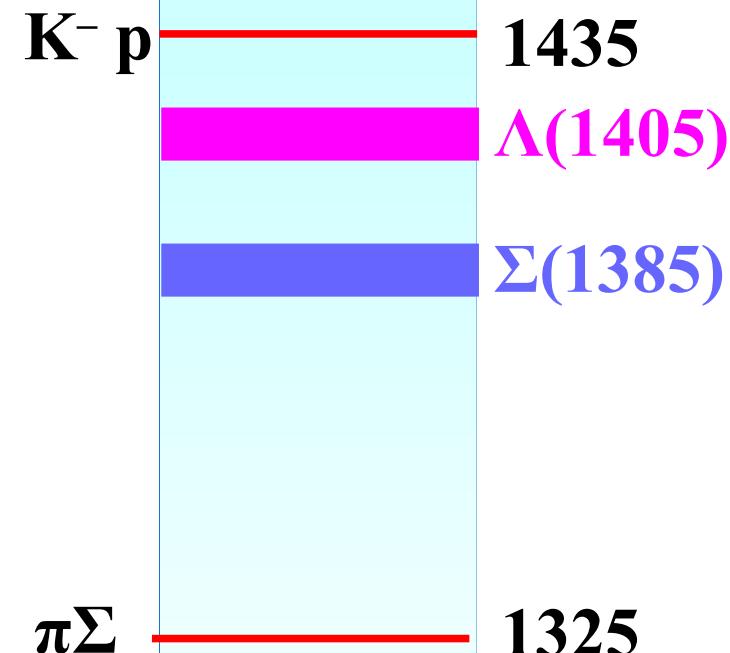
■ $\Lambda(1405)$ $\bar{K}N$ quasi-bound state

Dalitz, Tuan ('60), Koch ('94), Kaiser, Siegel, Weise ('95)

● Positive scattering length in K⁻ atoms

M.Iwasaki et al. PRL78('97)3067;

*M.Bazzi et al. [SIDDHARTA Collab.],
PLB704('11)113.*



K⁻ p interaction

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■ Kaonic nuclei ?

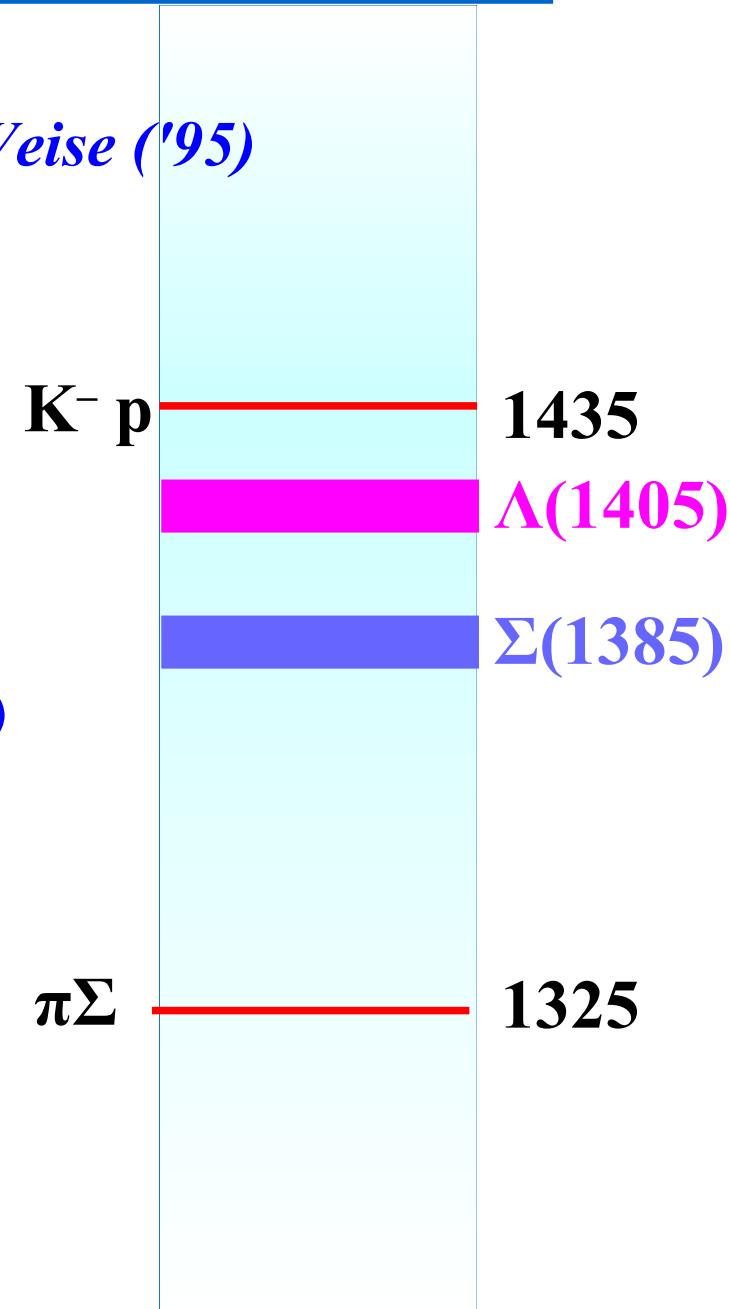
Nogami ('63); Akaishi, Yamazaki ('02);

Shevchenko, Gal, Mares ('07), Ikeda, Sato ('07)

→ Needs precise info. on $\bar{K}N$ int.

- Scattering amplitude and Potential fitting scattering and SIDDARTA data in chiral approach

Ikeda, Hyodo, Weise ('11, '12), Miyahara, Hyodo ('16)



$K^- p$ interaction

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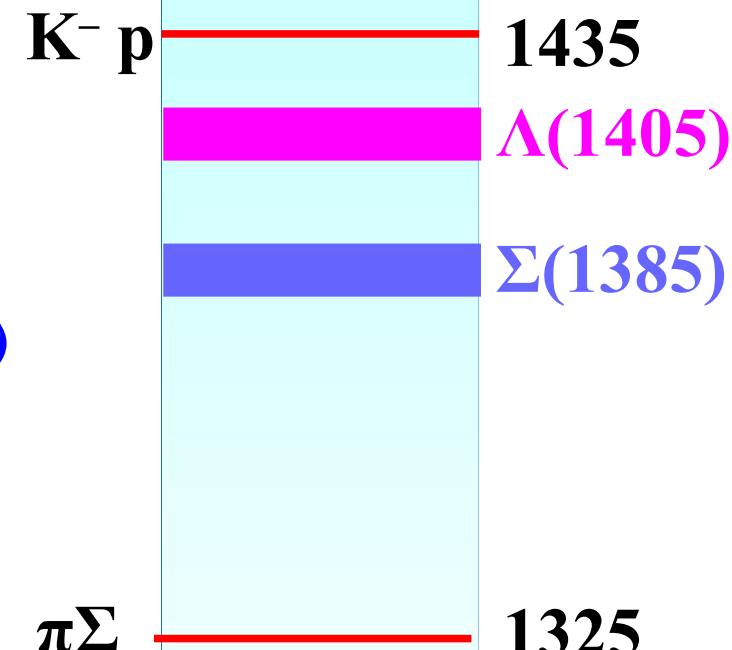
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Miyahara, Hyodo ('16)*



How about $K^- p$ correlation ?

K⁻ p scattering and K⁻ p correlation

■ K⁻ p scattering: Plane wave + Outgoing spherical wave

$$\begin{aligned}\Psi^{(+)}(K^- p) \rightarrow & \exp(i k z) \chi(K^- p) \\ & + \frac{\exp(i k r)}{r} [f(K^- p \rightarrow K^- p) \chi(K^- p) + f(K^- p \rightarrow \bar{K}^0 n) \chi(\bar{K}^0 n)]\end{aligned}$$

K⁻ p scattering and K⁻ p correlation

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■ K⁻ p correlation: Plane wave + Incoming spherical wave

$$\begin{aligned}\Psi^{(-)}(K^- p) \rightarrow & \exp(i k z) \chi(K^- p) \\ & + \frac{\exp(-i k r)}{r} [\tilde{f}(K^- p \rightarrow K^- p) \chi(K^- p) + \tilde{f}(K^- p \rightarrow \bar{K}^0 n) \chi(\bar{K}^0 n)]\end{aligned}$$

$$\tilde{f}(K^- p \rightarrow K^- p) = \frac{\tilde{S}-1}{2iq} \neq f(K^- p \rightarrow K^- p) , \quad \tilde{S} = \left[\frac{S_{I=0}^{-1} + S_{I=1}^{-1}}{2} \right]^{-1}$$

K⁻ p scattering and K⁻ p correlation

■ K⁻ p scattering: Plane wave + Outgoing spherical wave

$$\begin{aligned}\Psi^{(+)}(K^- p) \rightarrow & \exp(ikz)\chi(K^- p) \\ & + \frac{\exp(ikr)}{r} [f(K^- p \rightarrow K^- p)\chi(K^- p) + f(K^- p \rightarrow \bar{K}^0 n)\chi(\bar{K}^0 n)]\end{aligned}$$

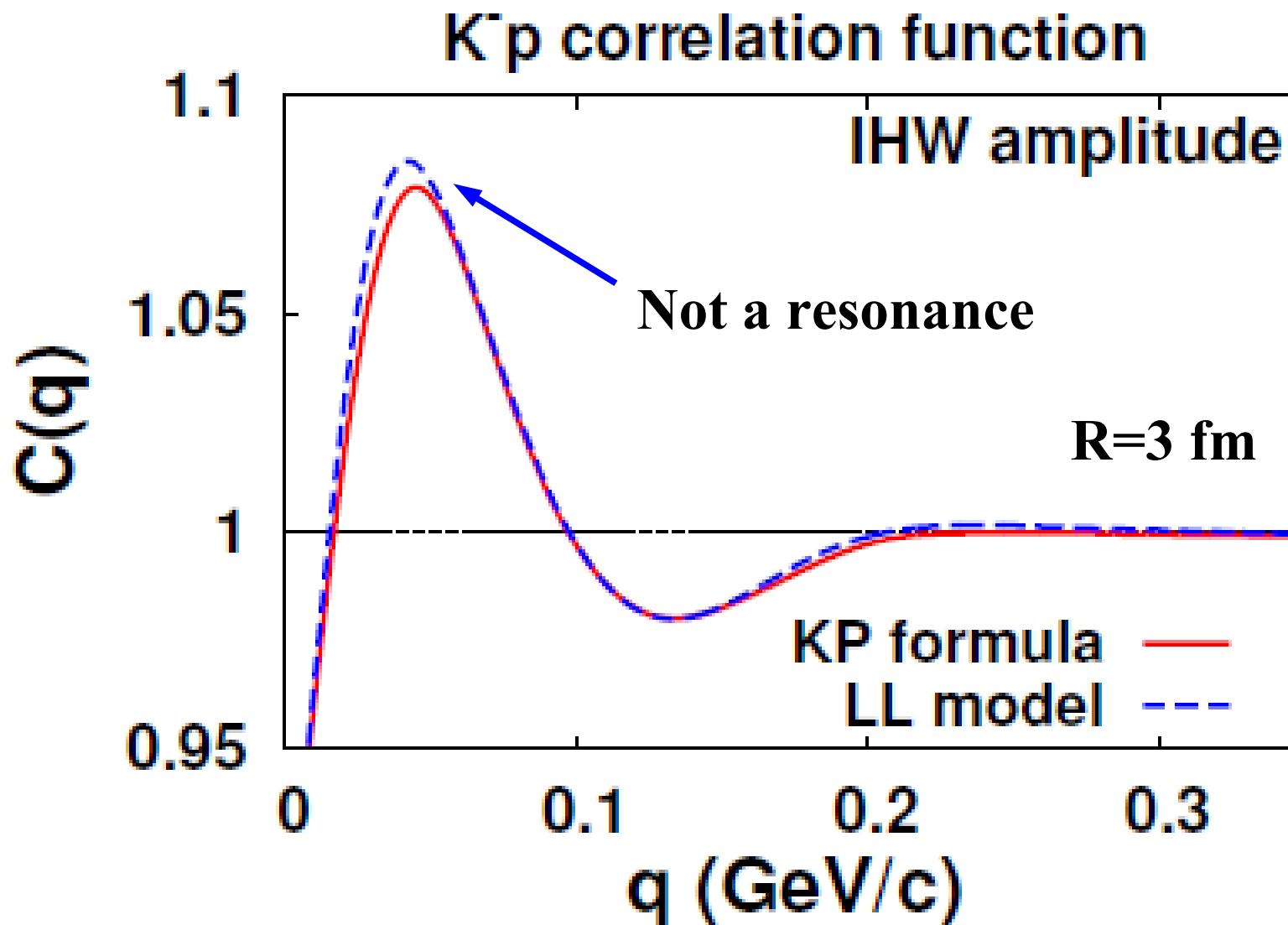
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$$\tilde{f}(K^- p \rightarrow K^- p) = \frac{\tilde{S}-1}{2iq} \neq f(K^- p \rightarrow K^- p) , \quad \tilde{S} = \left[\frac{S_{I=0}^{-1} + S_{I=1}^{-1}}{2} \right]^{-1}$$

$\Psi^{(-)} \neq (\Psi^{(+)})^* \rightarrow K^- p$ correlation probes
different combination of I=0, 1 from K⁻ p scattering

K⁻p correlation



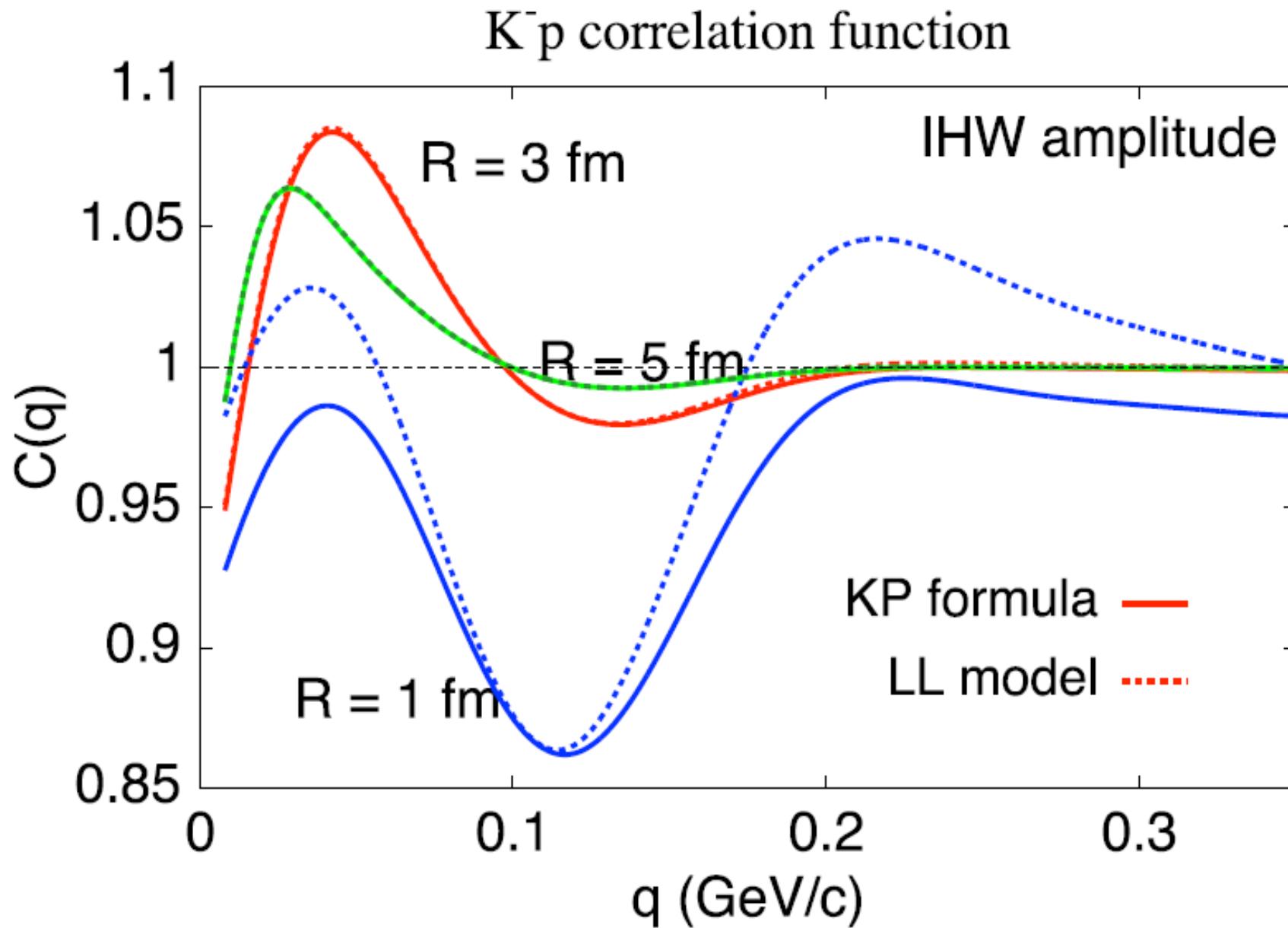
AO, K.Morita, K.Miyahara, T.Hyodo, arXiv:1603.05761

Potential: K. Miyahara, T. Hyodo, PRC93 ('16) 015201.

Amplitude: Y. Ikeda, T. Hyodo, W. Weise, PLB 706 ('11) 63; NPA 881 ('12) 98.

Fitting SIDDHARTA data: M. Bazzi et al. [SIDDHARTA Collab.], PLB 704 ('11) 113.

K⁻p correlation



$\psi^{(-)}$: Another source of fake peak ?

K. Miyahara

Summary

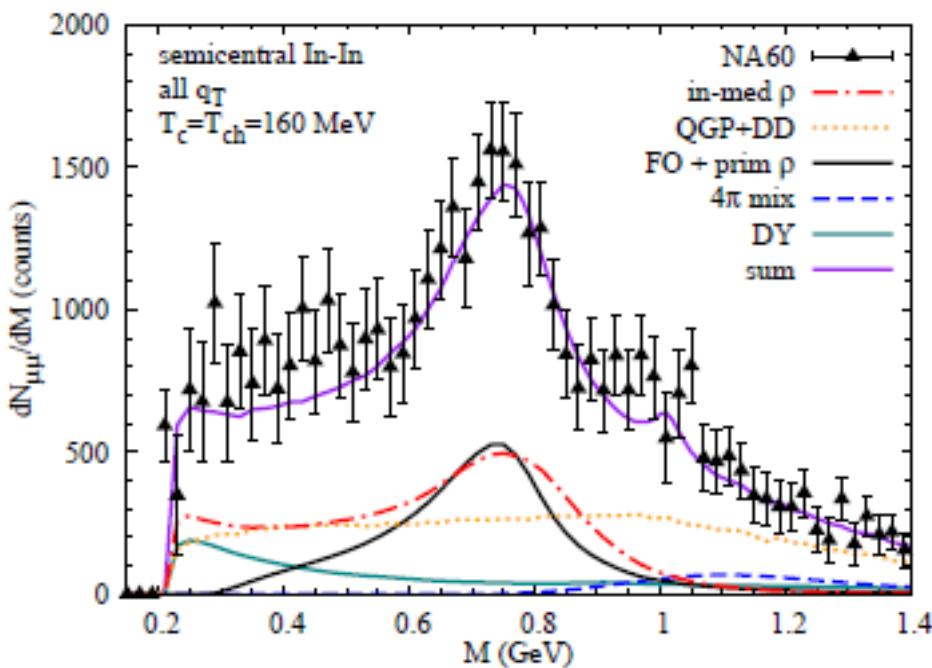
- Heavy-ion collisions are hadron factories, and can be utilized to perform hadron physics as well.
- Hadron-Hadron correlation contains information on interactions.
 - Correlations in various pairs have been measured:
 $\pi\pi$, KK , pp , nn , $p\bar{p}$, $\Lambda\Lambda$, Λp , $(K^- p, \Omega^- p)$...
 - When the pair purity is large enough, corr. fn. has sensitivity to hh interaction.
- Some of hh correlations have been discussed.
 - $\Lambda\Lambda$ correlation data is consistent with $\Lambda\Lambda$ interaction with $1/a_0 < -0.8 \text{ fm}^{-1}$, provided that pair purity is large enough, $\lambda > 0.35$.
 - $\Omega^- p$ correlation may tell us the existence of a S=-3 dibaryon.
 - $K^- p$ correlation is found to probe other combination of scattering amplitudes of I=0 and I=1.
- Many other type of pairs are waiting for us.

Thank you !

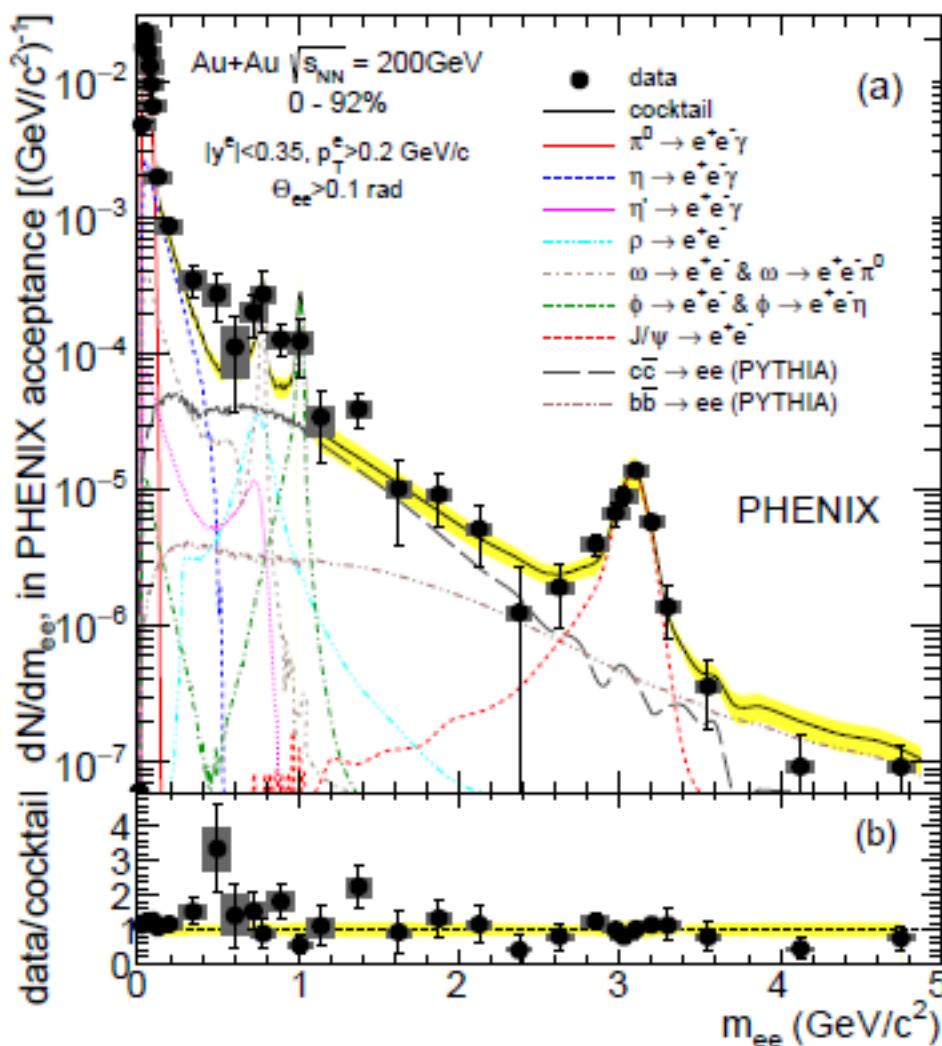
Thermal Modification of Hadron Spectral Function

Low mass dilepton

- Enhanced dilepton invariant mass spectrum may be explained with broadening of ρ meson and other mechanisms.



NA60 Collab. (R. Arnaldi *et al.*),
PRL96('06)162302.
H. van Hees, R. Rapp, *NPA806*('08)339



PHENIX Collab. (A. Adare *et al.*),
PRC93('16)014904

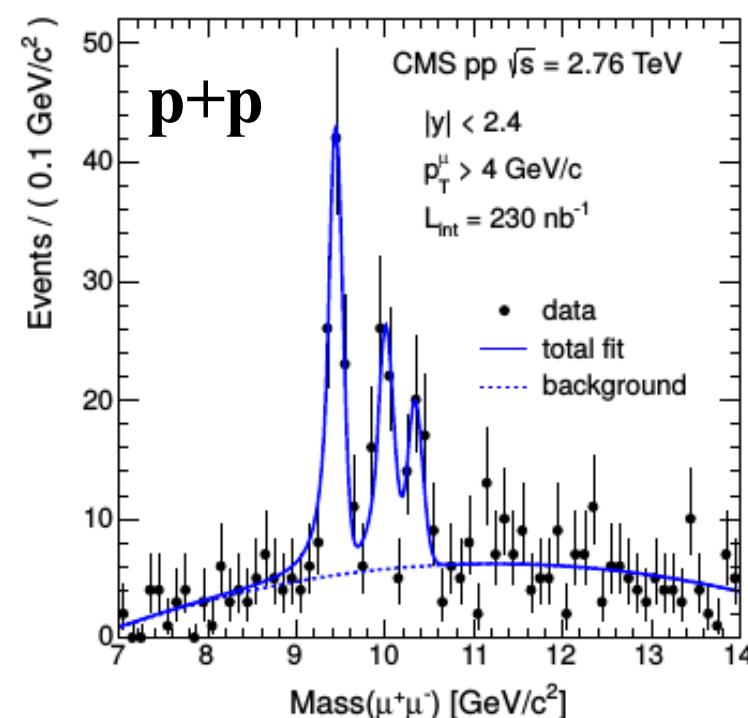
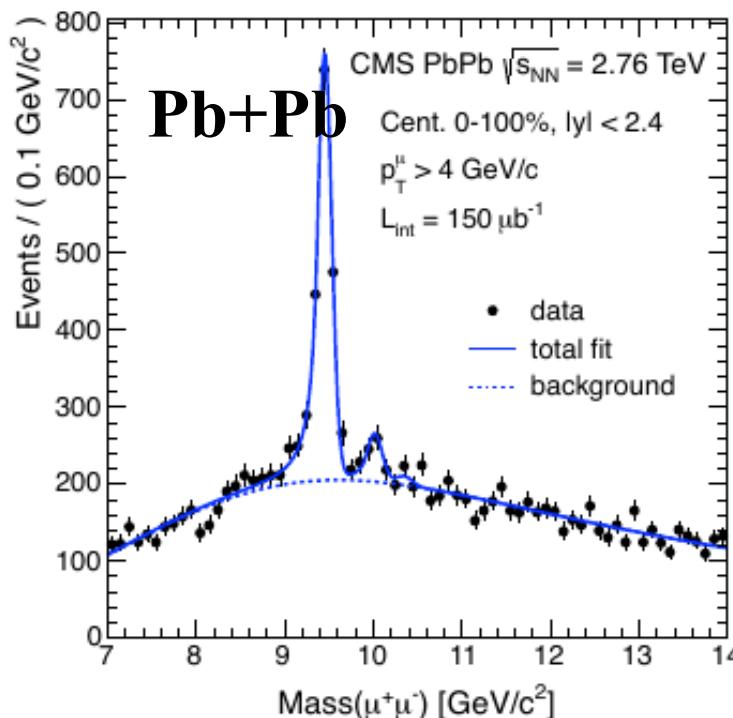
STAR Collab. (L. Adamczyk *et al.*),
PRL113 ('14)022301

Thermal Modification of Hadron Spectral Function

- Excited Upsilon (2S, 3S) are suppressed in Pb+Pb

CMS Collab. (S. Chatrchyan et al.), PRL107('11)052302; PRL 109 ('12) 222301

→ Excited quarkonium may melt before g.s. does by hot-medium effects



- η' mass reduction ? T.Csorgo, R.Vertesi, J.Sziklai, PRL105('10)182301

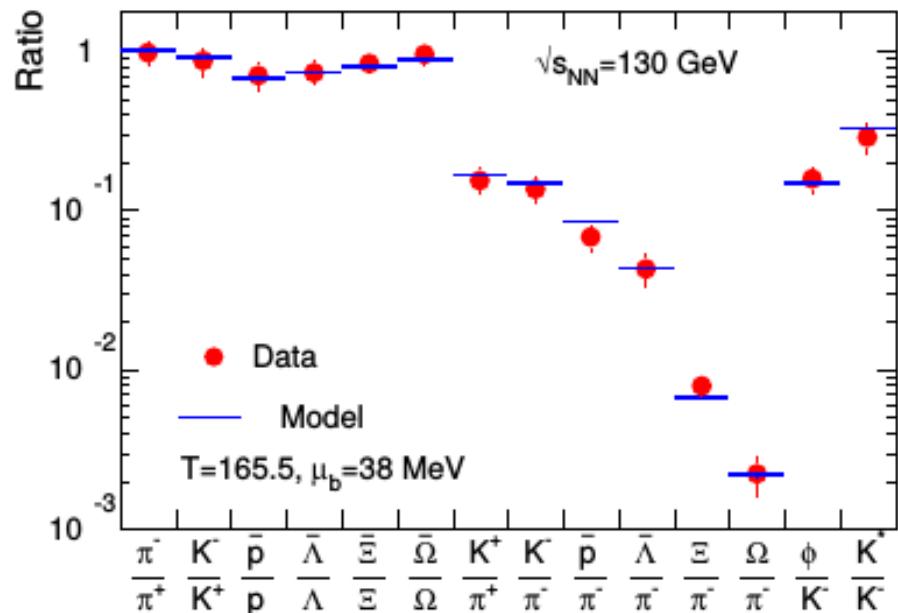
Modification by T , complementary to pA reactions.

Hadron Formation Mechanism

■ Statistical Model

*A. Andronic, P. Braun-Munzinger, J. Stachel,
NPA772('06)167; J. Cleymans, H. Oeschler,
K. Redlich, S. Wheaton, PRC 73 ('06), 034905*

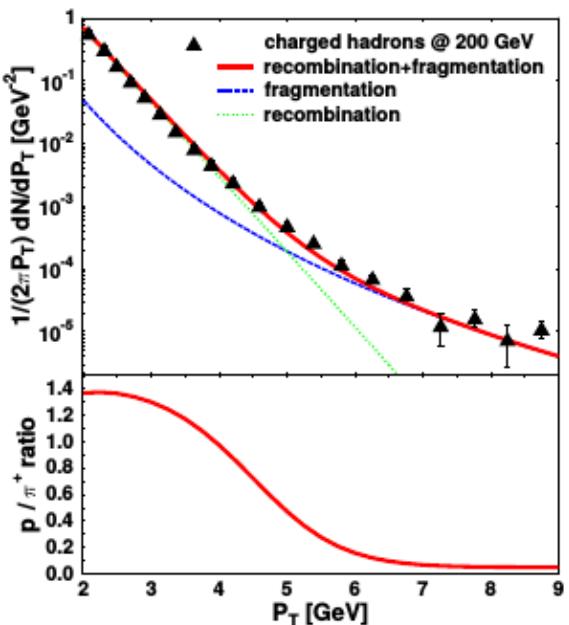
**Hadron yields are well explained
by the statistical model.**



■ Coalescence of quarks

*R.J.Fries, B. Muller, C. Nonaka, S.A.Bass,
PRL90('03)202303;
V. Greco, C.M.Ko, P. Levai, PRL90('03)202302;
R.C.Hwa, C.B.Yang, PRC67('03)034902*

**pT spectra and v2 are explained
by coalescence of quarks.**



■ Coalescence of hadrons

H.Sato, K. Yazaki, PLB98('81)153

**Deuteron spectrum is explained
by the coalescence of pn.**

Exotic Hadron Production in HIC

- How do exotic hadrons are produced ?

Statistical ?

Quark Coalescence ?

Hadron Coalescence ?

- Statistical model overestimates

$\Lambda(1520, 3/2^-)$ yield.

L=2 effects: Kanada-En'yo, Muller ('06)

- Statistical model underestimates $f_0(980)$ yield

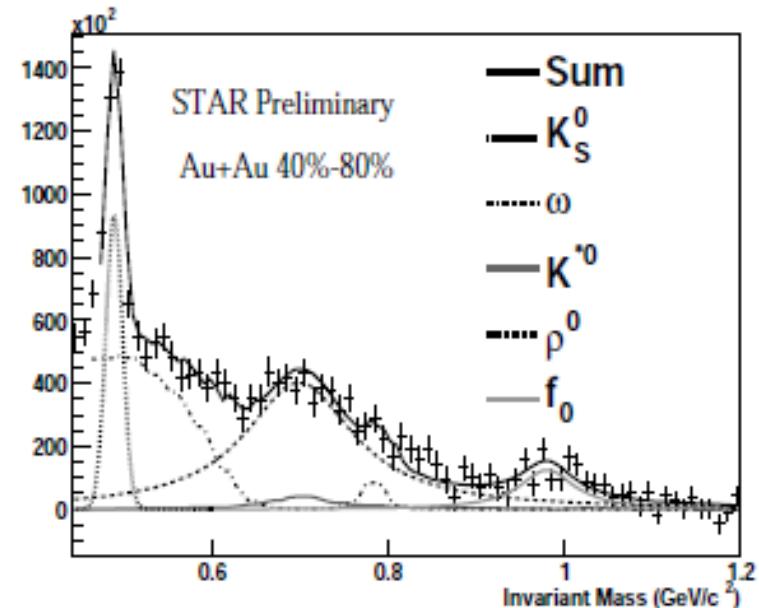
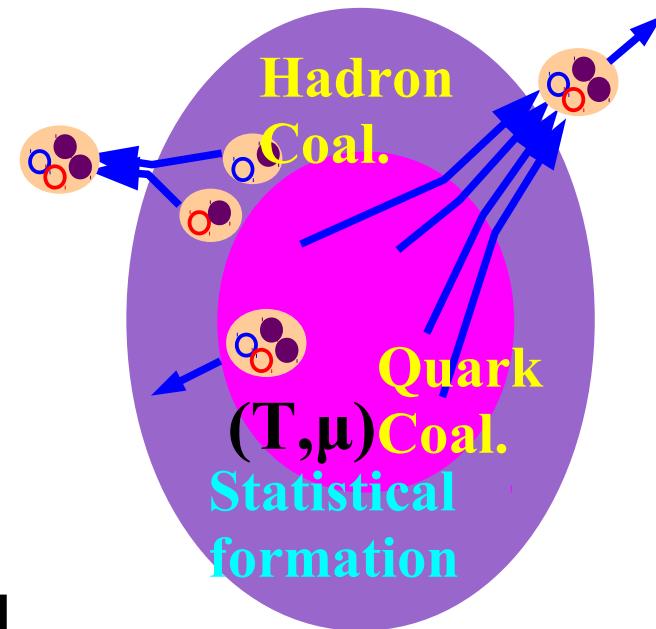
P. Fachini [STAR Collab.J, NPA715('03)462.

STAR: $N(f_0(980)) \sim 8.4$, Stat: 5.6,

Coal: 2q:0.76-3.8, 4q:0.1, Mol: 13

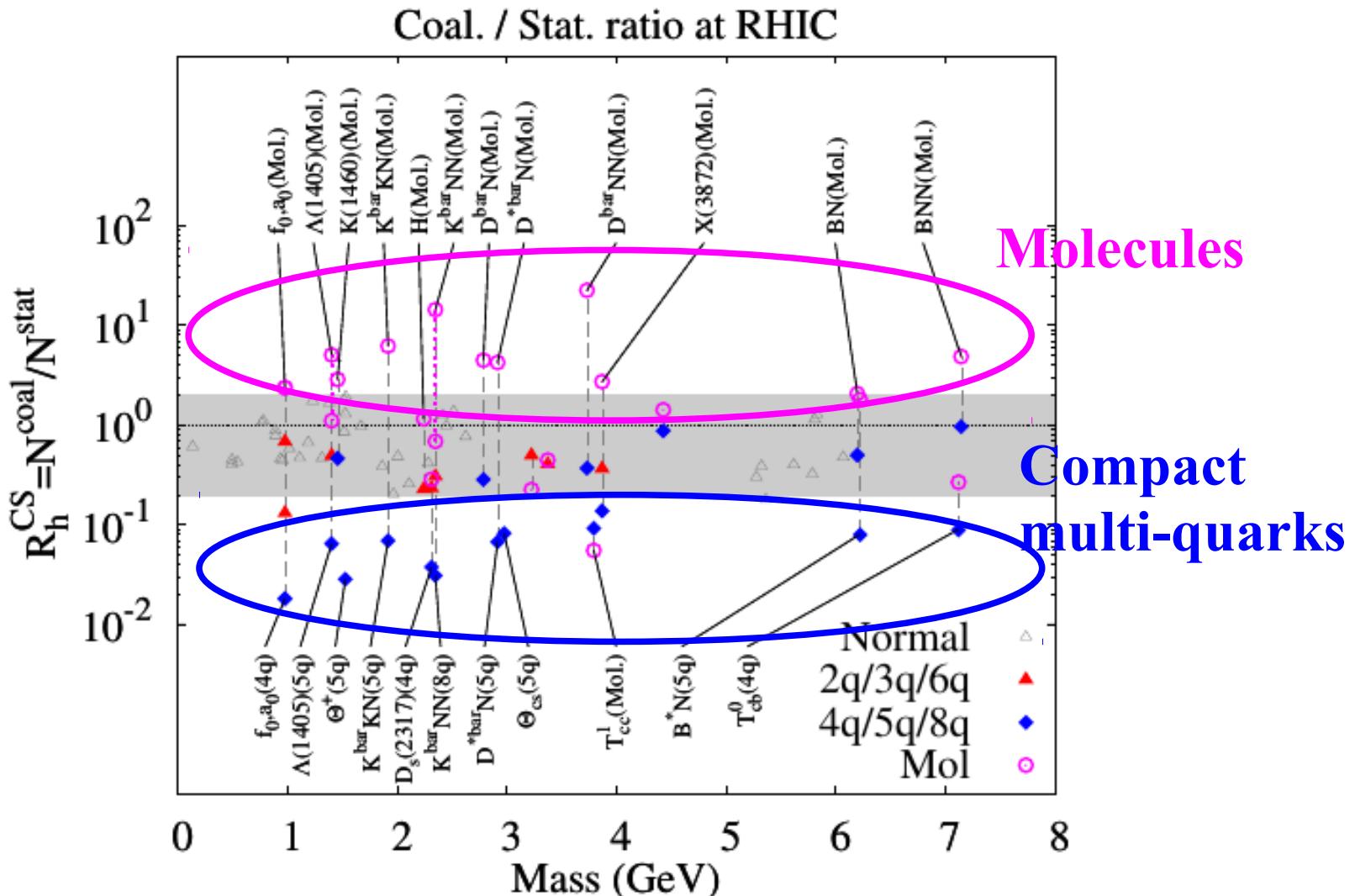
S. Cho et al. [ExHIC Collab.J, PRL106 ('11)212001.

*Production yield in HIC may tell us
the nature of exotic hadrons !*



Exotic Hadron Production in HIC

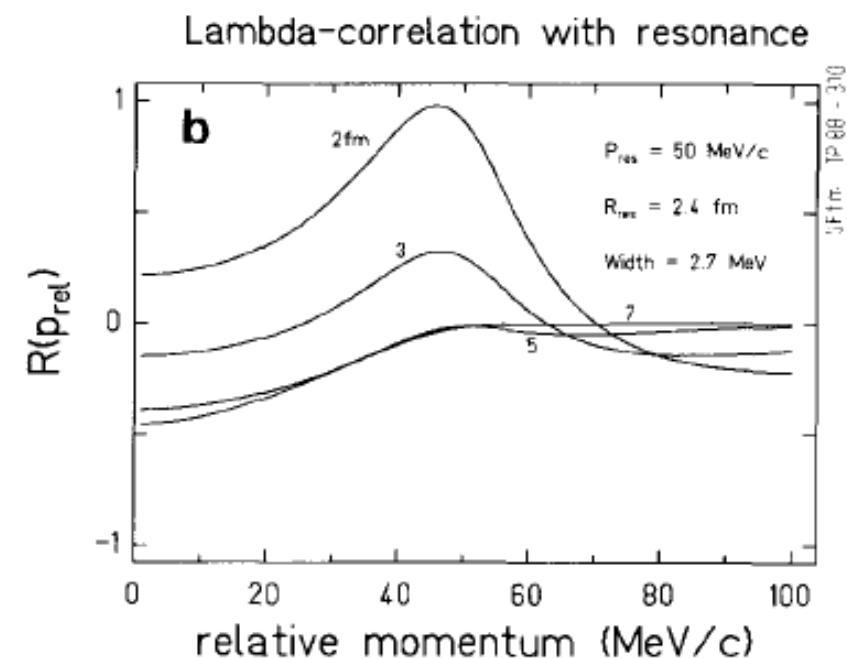
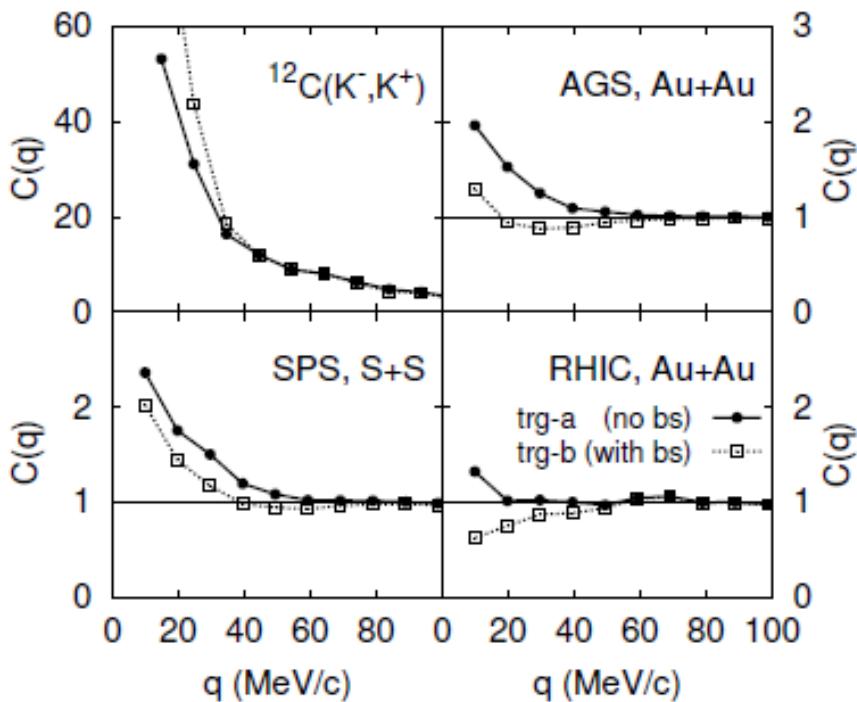
S. Cho et al. [ExHIC Collab.], PRL106 ('11)212001; PRC84('11)064910.



Hadronic molecules will be made more abundantly in HIC than stat. model predictions.

Correlation → Interaction

- If we know correlation and source, it should be possible to get knowledge of hadron interactions !
- How about $\Lambda\Lambda$ interaction ?
 - Λ particle is too short-lived to perform scattering experiments.
Mass=1115.6 MeV, $\tau = 2.6 \times 10^{-10}$ s, $c\tau = 7.89$ cm
quark content = uds (p=uud, n=udd)



AO, Y. Hirata, Y. Nara, S. Shinmura, Y. Akaishi C. Greiner, B. Muller, PLB219(1989)199.
Nucl. Phys. A 670 (2000), 297c

Feed-Down Effects & Residual Source

■ Correlation Fn. w/ Feed-down & Residual source effects.

$$C_{\text{corr}}(q) = 1 + \lambda(C_{\text{bare}}(q) - 1) + a_{\text{res}} \exp(-4r_{\text{res}}^2 q^2)$$

■ Purity of $\Lambda\Lambda$ pair λ

- Not easy to reject Σ^0 decay in exp't.

$$\rightarrow \lambda \sim (1 - \Sigma^0/\Lambda)^2$$

- Morita, Furumoto, AO ('15)

$$\rightarrow \lambda = (0.67)^2 = 0.45$$

$\Sigma^0/\Lambda = 0.278$ (p+Be, 28.5 GeV/c)

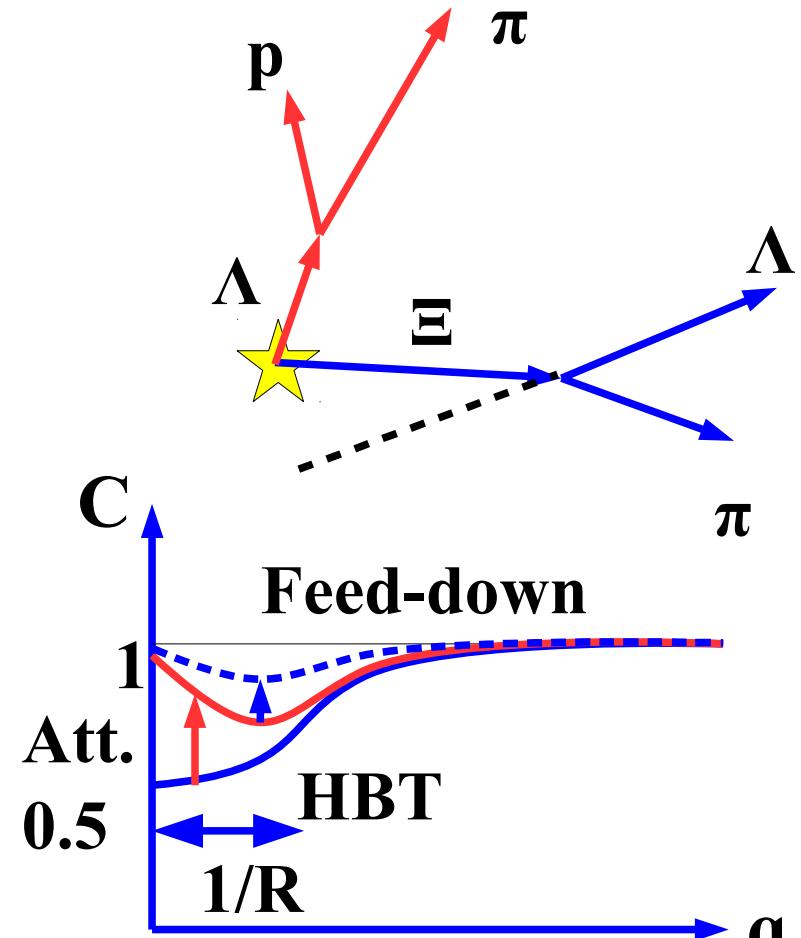
Sullivan et al. ('87)

$\Xi/\Lambda = 15\%$ (RHIC)

- STAR ('15)

$\rightarrow \lambda = \text{Free parameter} \sim 0.18$

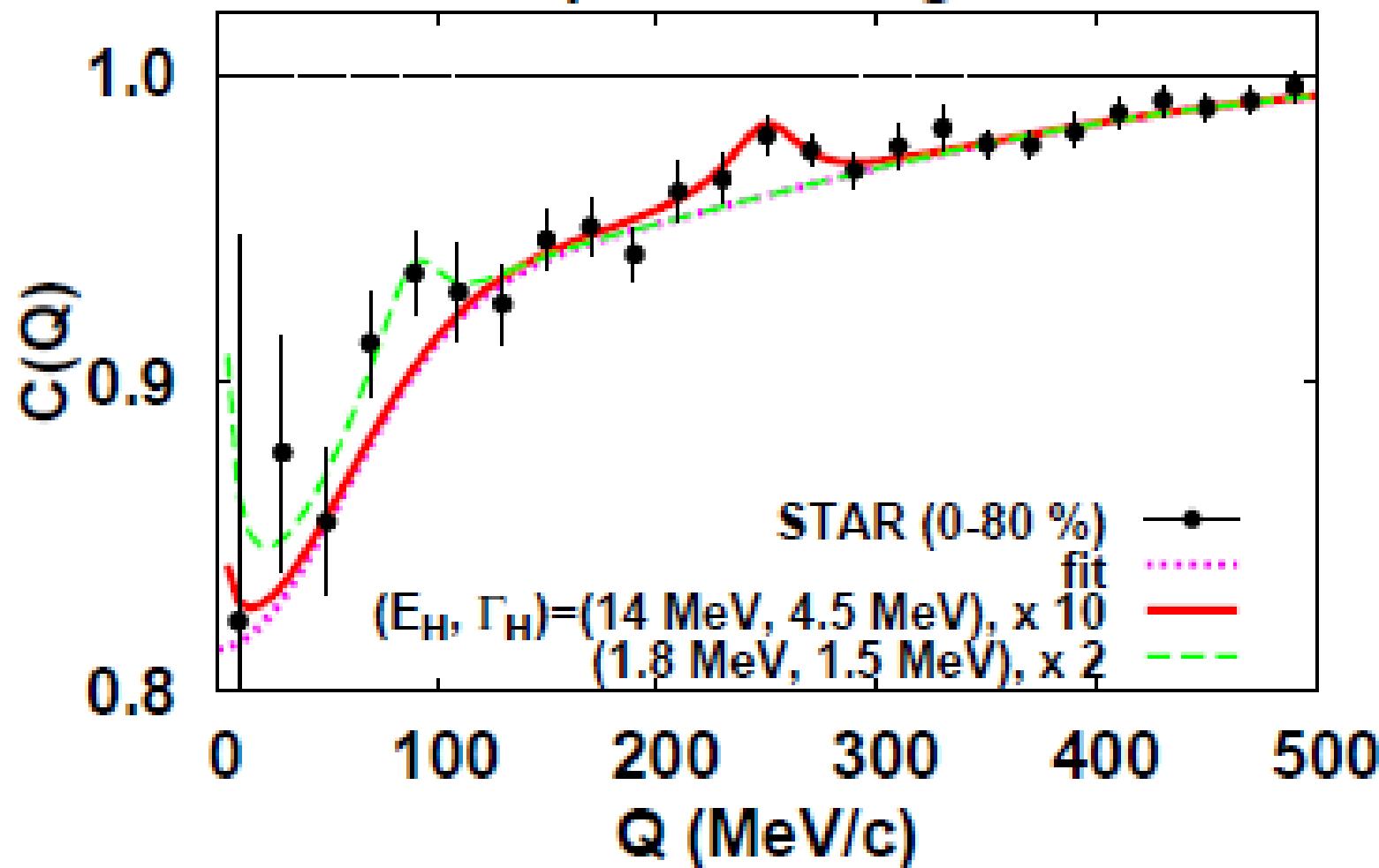
C(q~0)~0.8 suggests attraction with large λ ($\lambda > 0.35$) repulsion with small λ ($\lambda < 0.35$)



AO, Morita, Mihayara, Hyodo ('16)

Detecting H Resonance

Expected H signal

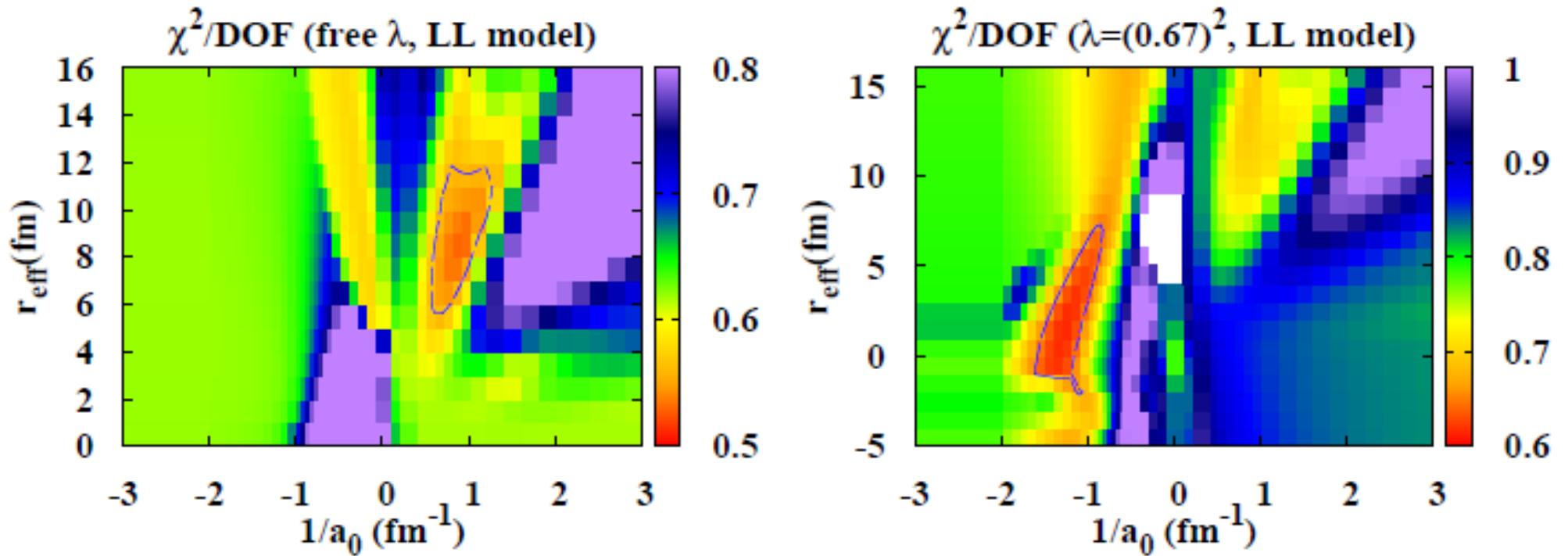


*When the resonance energy is much above the threshold,
detecting a resonance is not easy because of huge background.*

3S_1 contribution

- $\Omega(J=3/2)$ and p ($J=1/2$) can take spin quintet (5S_2) and triplet (3S_1) states for $L=0$.
 - 5S_2 ($J=2$) couples with octet-octet (e.g. $\Lambda\Xi$) via $L=2$.
 - 3S_1 ($J=1$) couples with octet-octet (e.g. $\Lambda\Xi$) via $L=0$.
→ strong absorption is expected at short distances.
- Our assumption: Strong absorption at $r < 2$ fm.
 - Strong absorption limit = Hard core boundary condition.
 - Intermediate absorption gives similar results.

$a_0(\Lambda\Lambda)$: Positive or Negative



free $\lambda \rightarrow$ positive a_0 is preferred by the STAR data.

fixed λ ($\lambda > 0.35$) \rightarrow negative a_0 is preferred by the STAR data.