

Heavy-ion and dense matter physics in Hokkaido University

initiated

Akira Ohnishi (YITP, Kyoto U.)

*Workshop on Nuclear Cluster Physics 2017
Oct. 25-27, 2017, Sapporo, Japan.*

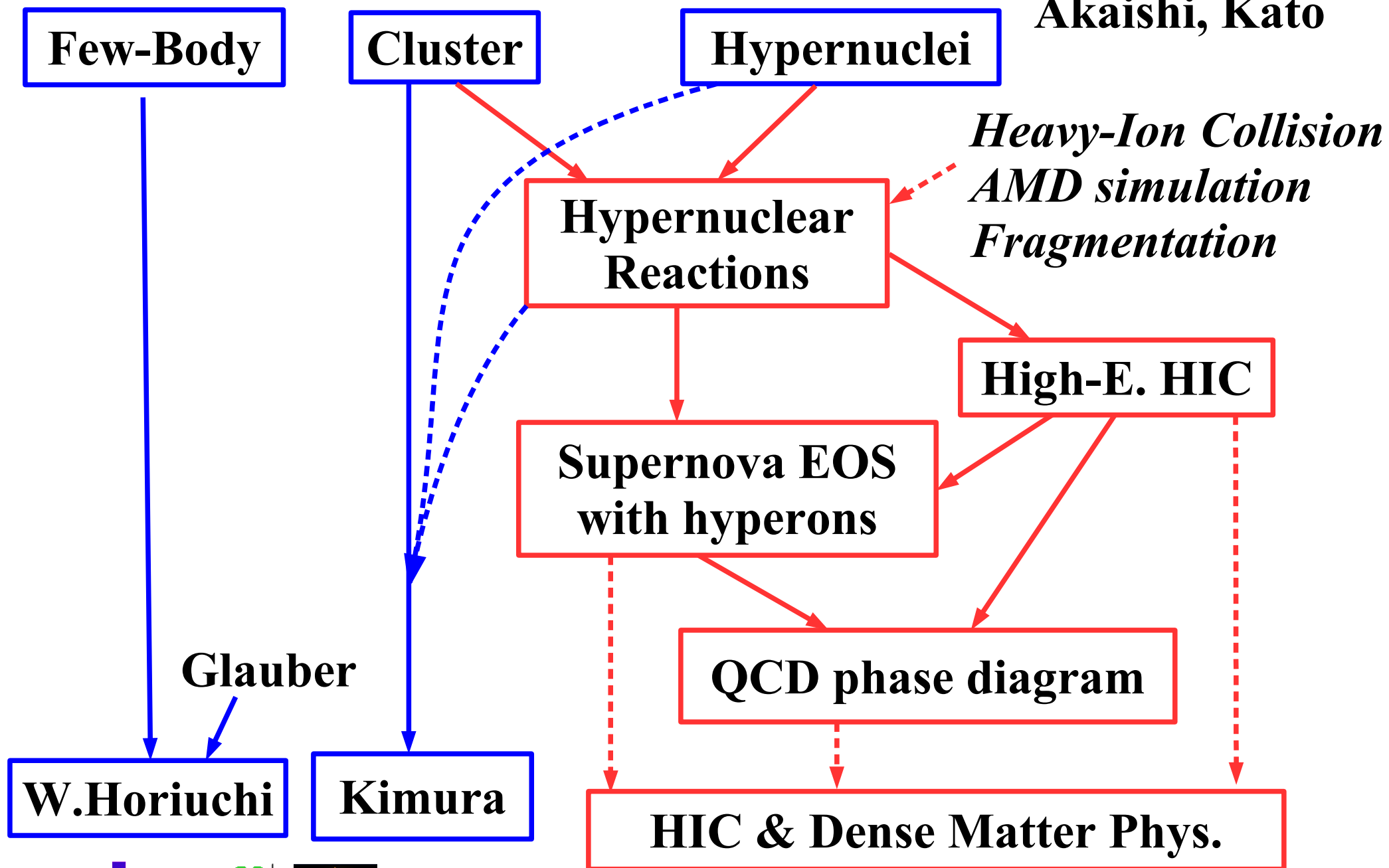


My Research Subjects in Hokkaido University

- I belonged to Hokkaido University from 1993 to 2008.
- Subjects I worked on
 - Strangeness Nuclear Physics
(Nara, Hirata, Maekawa, Tsubakihara, Matsumiya, [Isaka])
 - Heavy-Ion Collisions
(Nara, Otuka, Isse, Yoshino, Mizukawa)
 - Dense Matter Physics
(Okuda, Ishizuka, Ohnuma, Tsubakihara, [Miura])
 - Nuclear Structure (Itagaki, [Myo], Isshiki)
 - Nuclear Reactions (Uchida, Hirata, Maekawa, Yamaguchi)

My Research Subjects in Hokkaido University

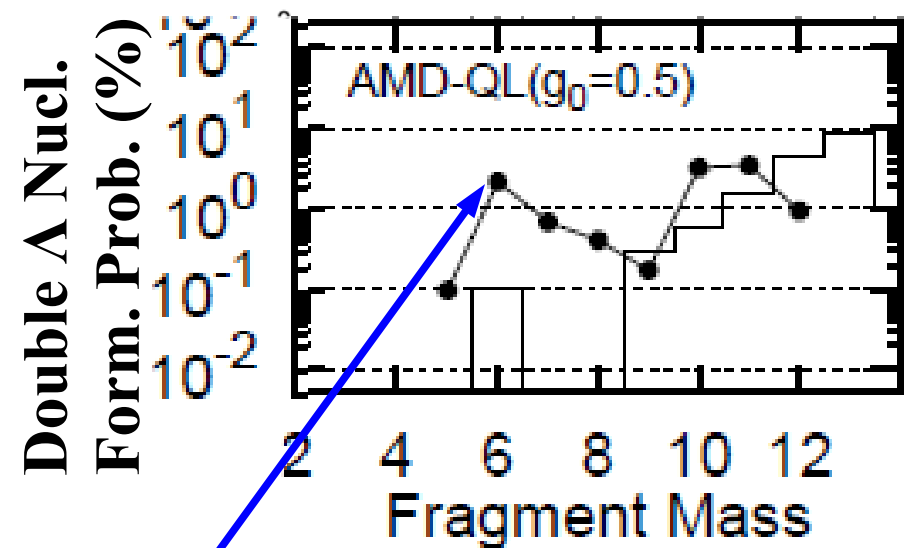
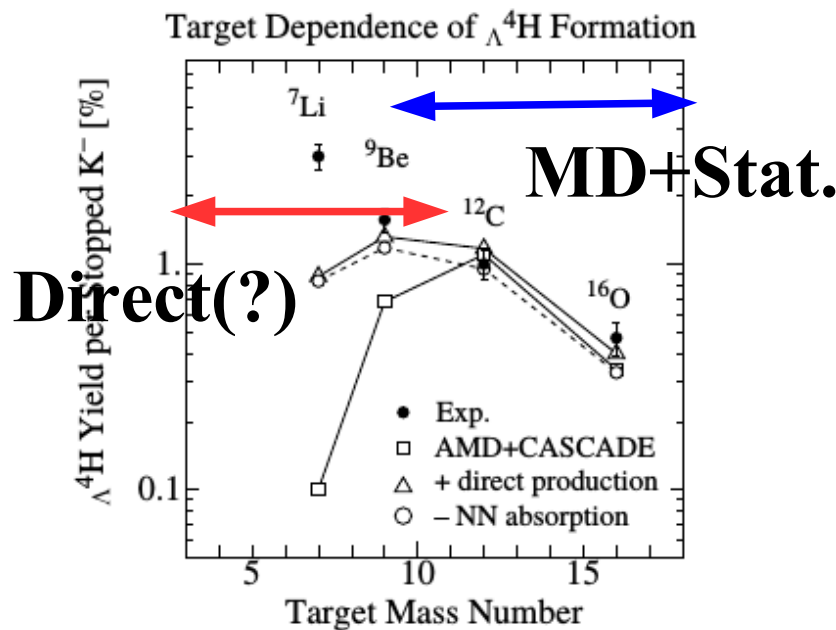
Akaishi, Kato



Hyperfragment Formation

■ Hypernuclear physics

- **Standard: Formation, Structure, Decay** (c.f. Motoba's talk)
Binding energies, Excitation spectra, Direct reactions, Decays ...
- **Non-standard aspects: Hyperfrag. formation, YY correlation, ...**
→ **Simulation calculation is useful!** (c.f. Yamada's work)



$\Lambda^4\text{H}$ form. prob. from stopped K^-

${}^6_{\Lambda\Lambda}\text{He}$ Hyperfrag. form. from stopped Ξ^-

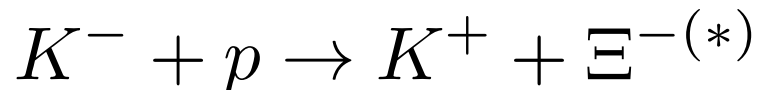
Nara, AO, Harada, PLB346 ('95)217

Hirata, Nara, AO, Randrup, PTP102('99)89

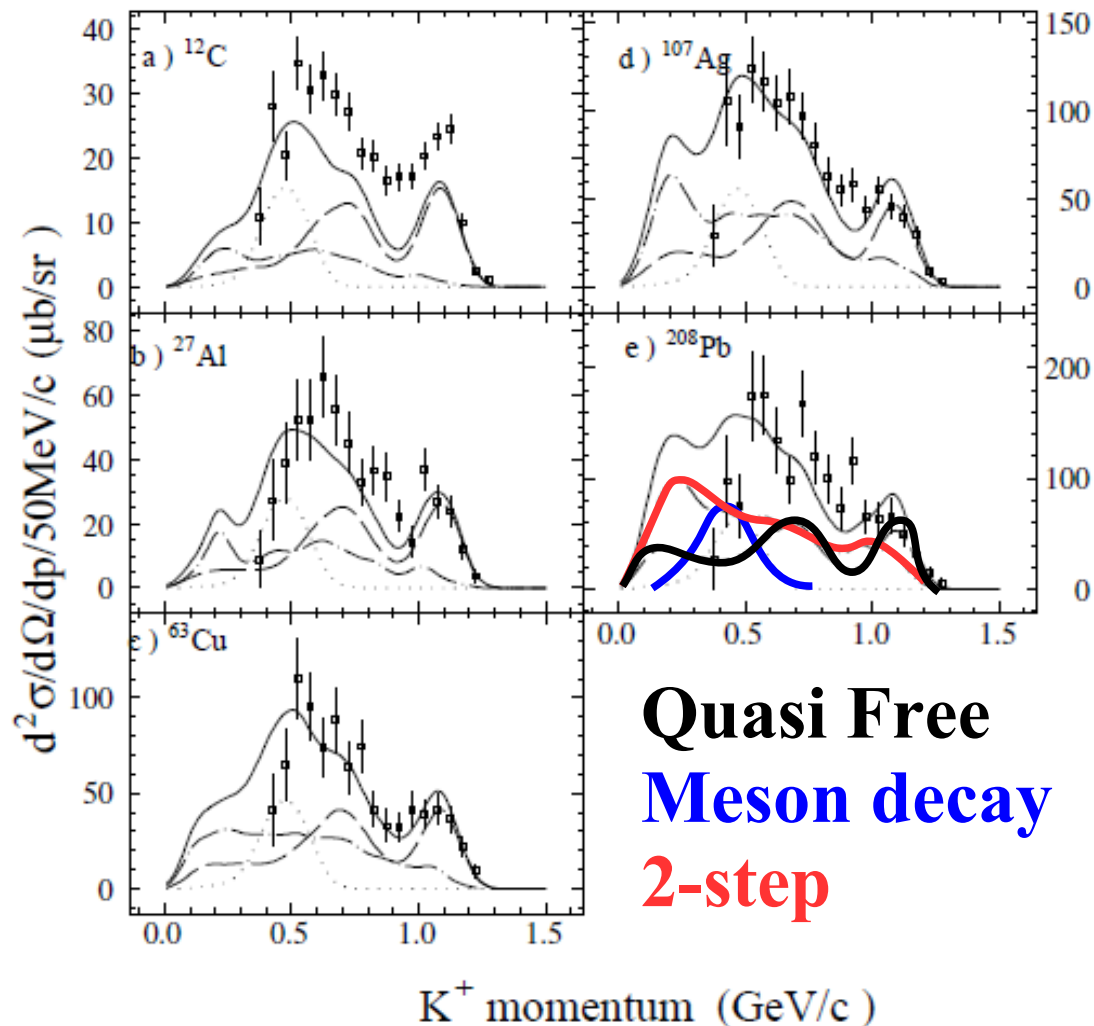
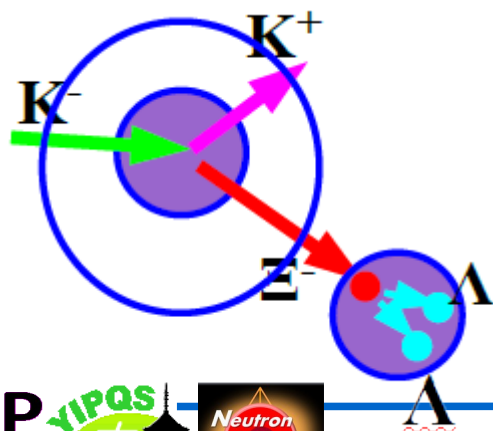
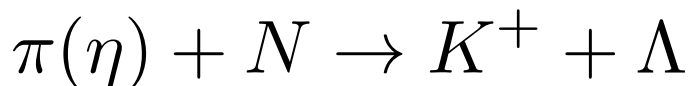
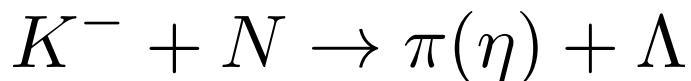
$(K^-, K^+) \text{ Reaction}$

- The primary reaction to produce Ξ^- , which is absorbed to form double Λ hypernuclei.

- Direct (Quasi-Free) prod. of Ξ and Ξ^* is not enough to explain the spectrum.



- Various 2-step processes may contribute !



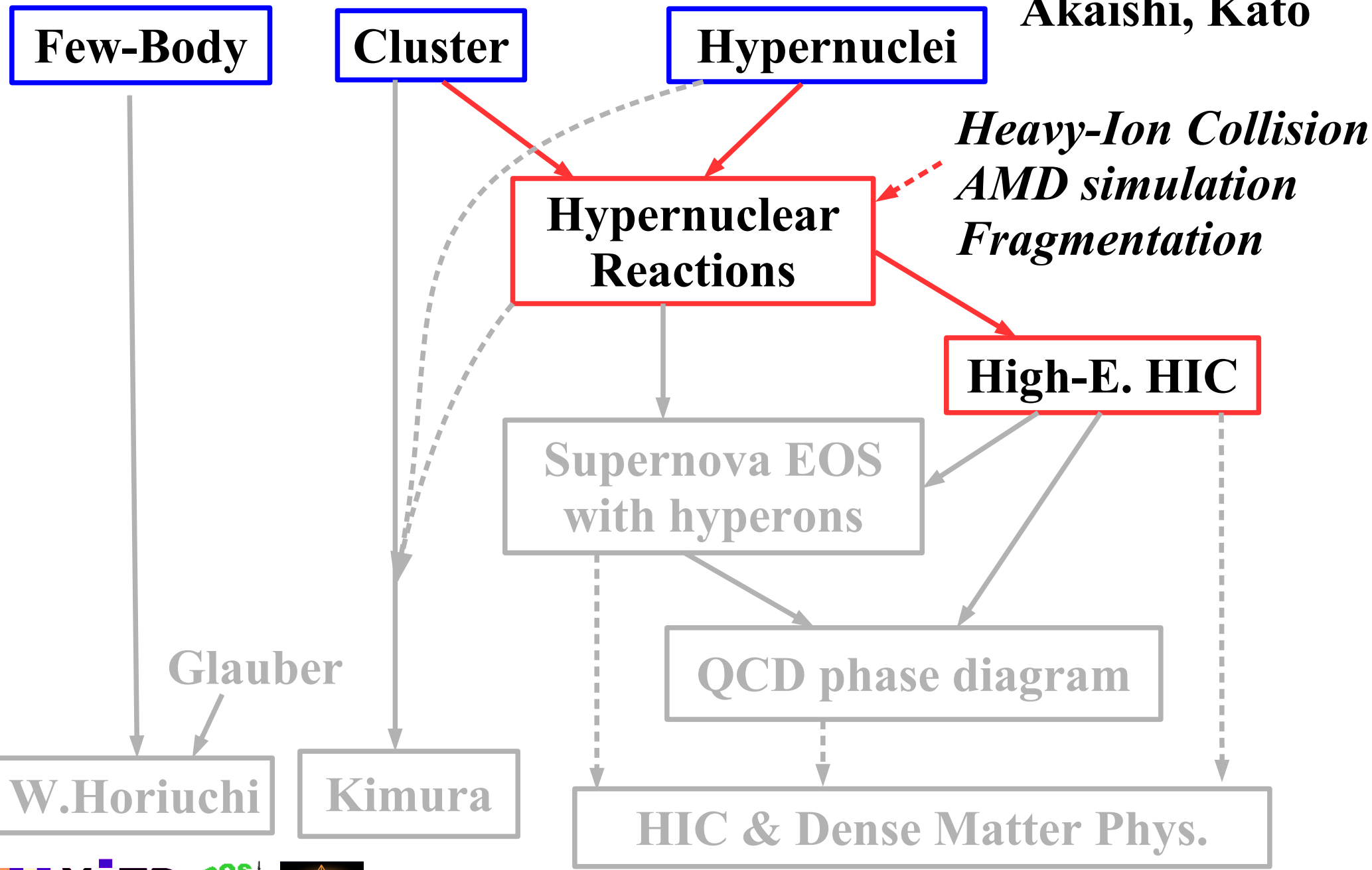
Quasi Free
Meson decay
2-step

Y. Nara, AO, T. Harada, A. Engel, NPA614 ('7), 433

High-Energy Heavy Ion Collisions

My Research Subjects in Hokkaido University

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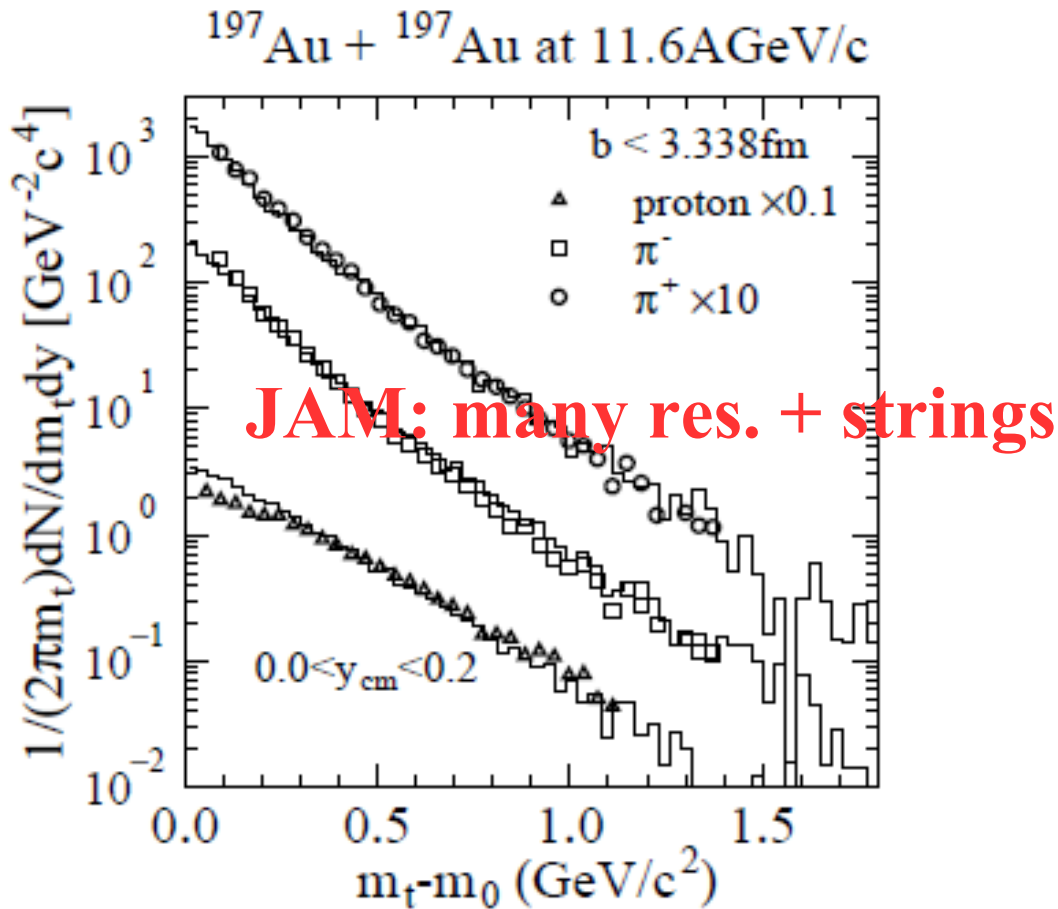


High-Energy Heavy-Ion Collisions

- **Lessons from Hypernuclear Reactions**
= Sum of small cross sections can be significant !
- **Hadron-String transport model in Heavy-Ion Collisions**
[Jet AA Microscopic transport model (JAM)]
 - Include as many processes as possible.
(Cross section book, 300 kyen !)
 - Include as many degrees of freedom as possible.
→ Ground state hadrons, Resonances, Strings, Jets, ...
 - Include as many ingredients as possible
→ Two-body collisions, Mean field potential, Fluctuations, ...
- **JAM is now one of the STANDARD transport models.**
 - Describes AA and pA collisions at $E/A = (1-160)$ GeV
 - Adopted in PHITS (nuclear engineering code)
 - High score also at low energies ($E/A \sim 300$ MeV)

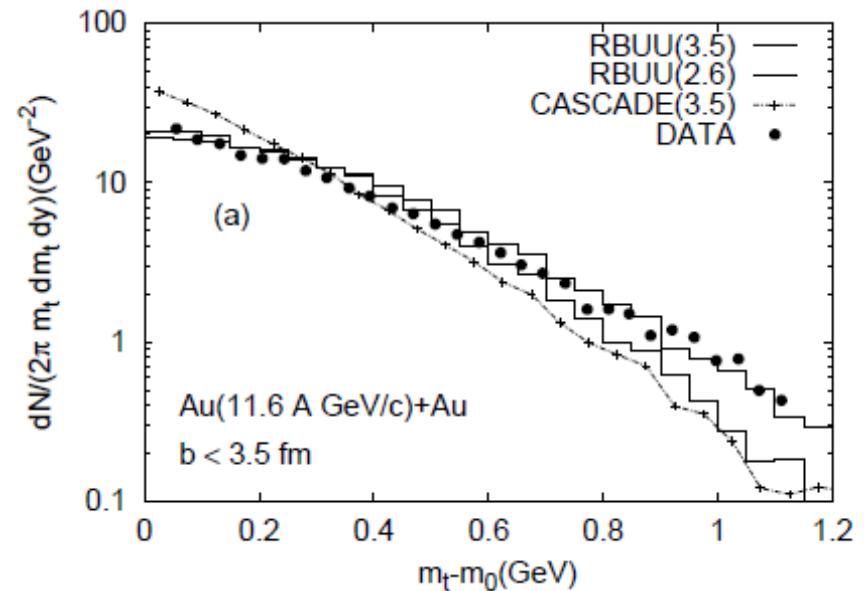
AGS energy ($E/A = 10.6 \text{ GeV}$) HIC

- Hadronic DOF matters.
- Winners in Hadron-String Cascade include
Res. ($M_B < 3 \text{ GeV}$, $M_M < 2 \text{ GeV}$) + String (continuum) (+ MF)



Nara, Otuka, AO, Niita, Chiba,
PRC('00), 024901

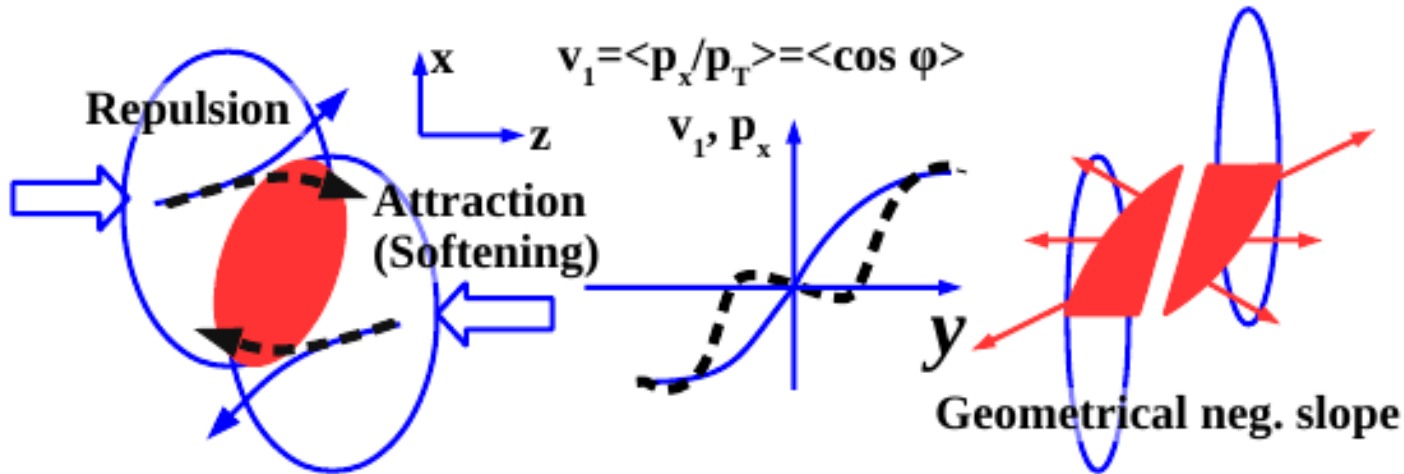
**RBUU(HSD):
low mass res. + string**



Sahu, Cassing, Mosel, AO, NPA672('00)376

Collective Flow

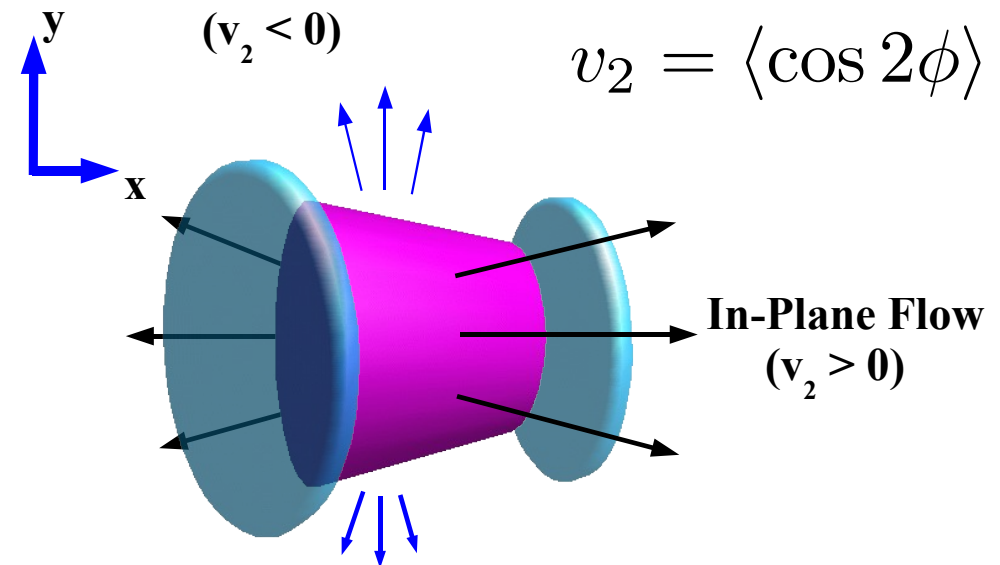
- Directed flow ($v_1, \langle p_x \rangle$), Elliptic flow (v_2)
 → Generated in the Early stage, sensitive to dense matter EOS



Out-of-Plane Flow

($v_2 < 0$)

$$v_2 = \langle \cos 2\phi \rangle$$



*Can we understand
the flows in JAM (+MF) ?*

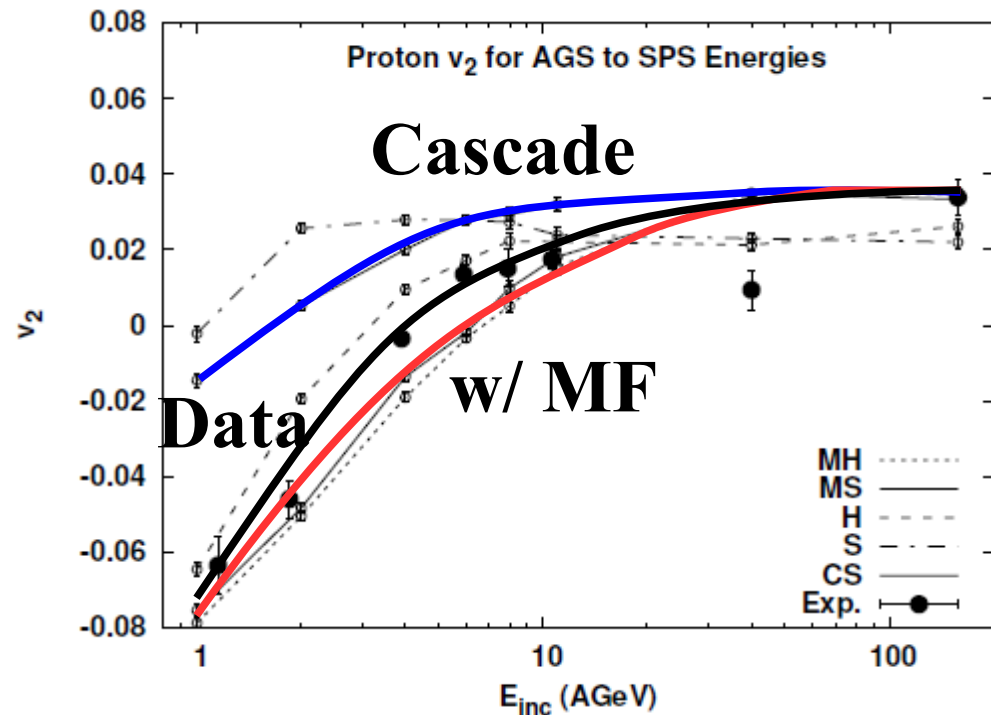
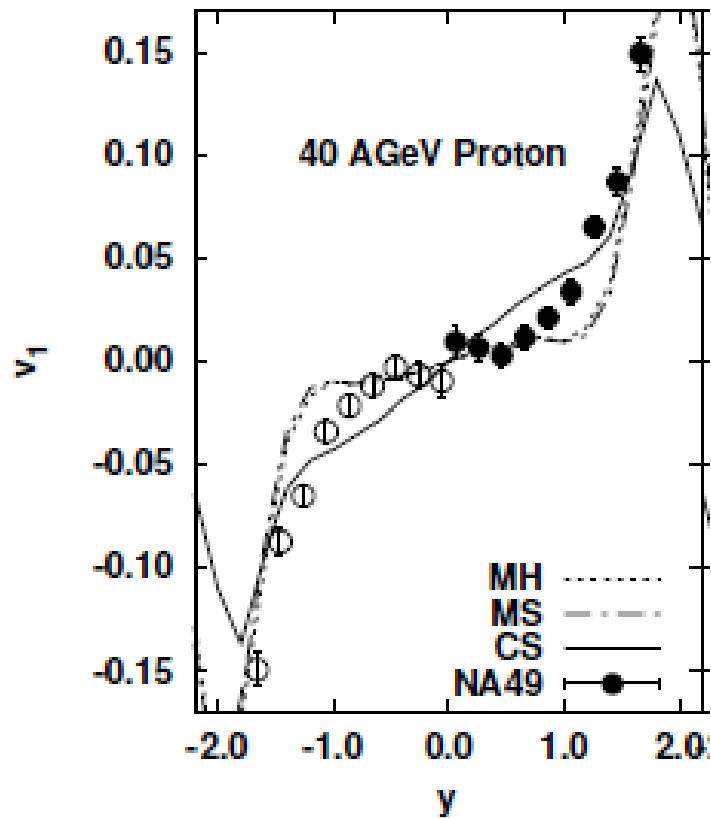
Hadron-String Cascade with Mean Field Potential

■ Hadron-String Cascade + Mean Field (JAM-MF)

● Data are in theoretical (MF) uncertainties.

→ There is no bulk QGP formation at $E/A < 160$ GeV

($\sqrt{s_{NN}} < 20$ GeV)



Isse, AO, Otuka, Sahu, Nara, PRC72('05)064908

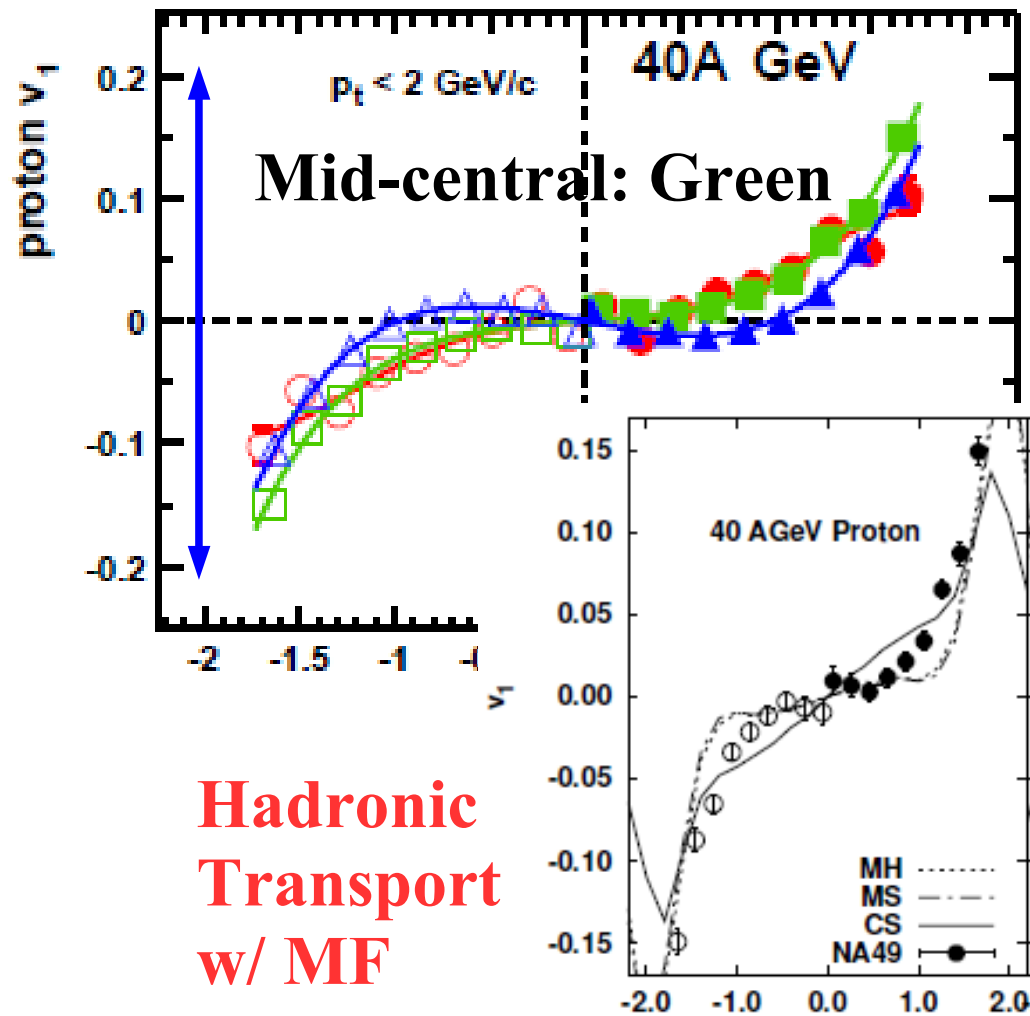
No QCD phase transition below $E/A=160$ GeV !

SPS(NA49) vs RHIC(STAR)

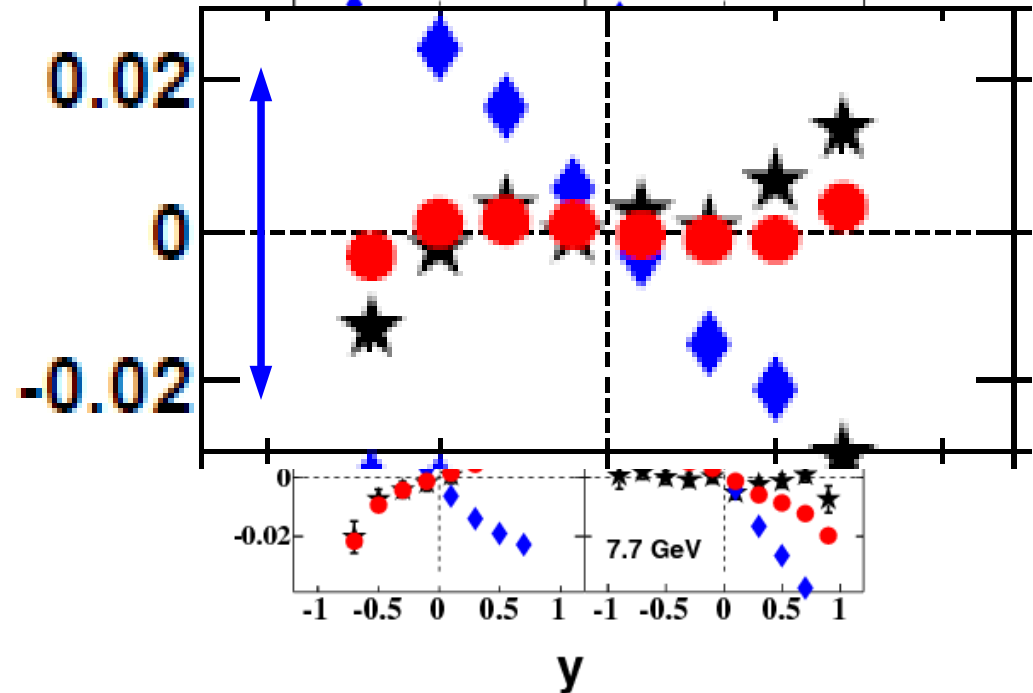
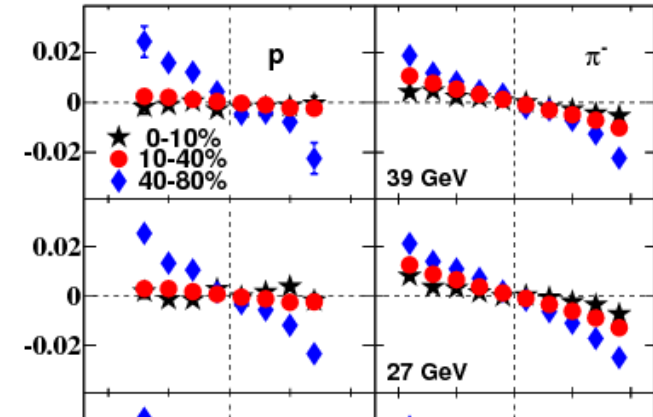
■ SPS (NA49), $\sqrt{s_{NN}} = 8.9$ GeV

■ RHIC(STAR), 7.7-39 GeV

C. Alt et al. (NA49), PRC68 ('03) 034903



M.Isse, A.O., N.Otuka, P.K.Sahu, Y.Nara,
PRC72 ('05)064908



L. Adamczyk et al. (STAR),
PRL 112(2014)162301

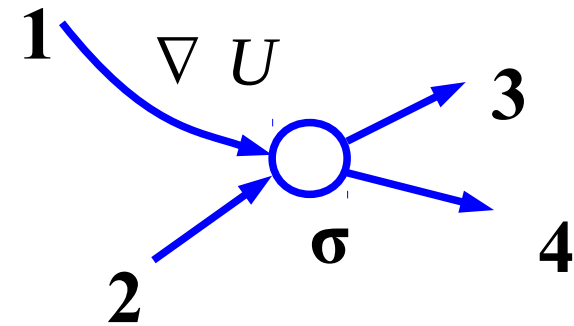
Transport model w/ EOS softening

■ Transport model (Boltzmann+MF)

E.g. Bertsch, Das Gupta, Phys. Rept. 160(88), 190

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \nabla f - \nabla U \cdot \nabla_p f = I_{\text{coll}}$$

$$I_{\text{coll}}(\mathbf{r}, \mathbf{p}) = -\frac{1}{2} \int \frac{d\mathbf{p}_2}{(2\pi)^3} d\Omega v_{12} \frac{d\sigma}{d\Omega} [f f_2 (1 - f_3)(1 - f_4) - (12 \leftrightarrow 34)]$$

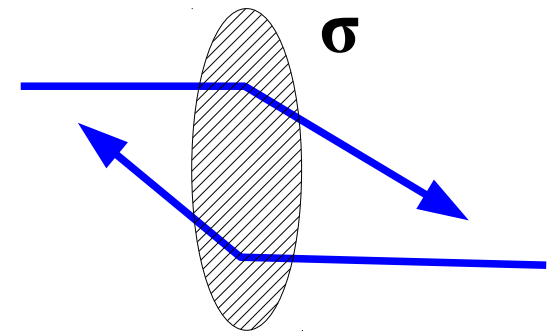


■ Simulating EOS softening in the collision term

Danielewicz, Pratt ('96); Sorge ('99); Nara, Niemi, AO, Stoecker ('16)

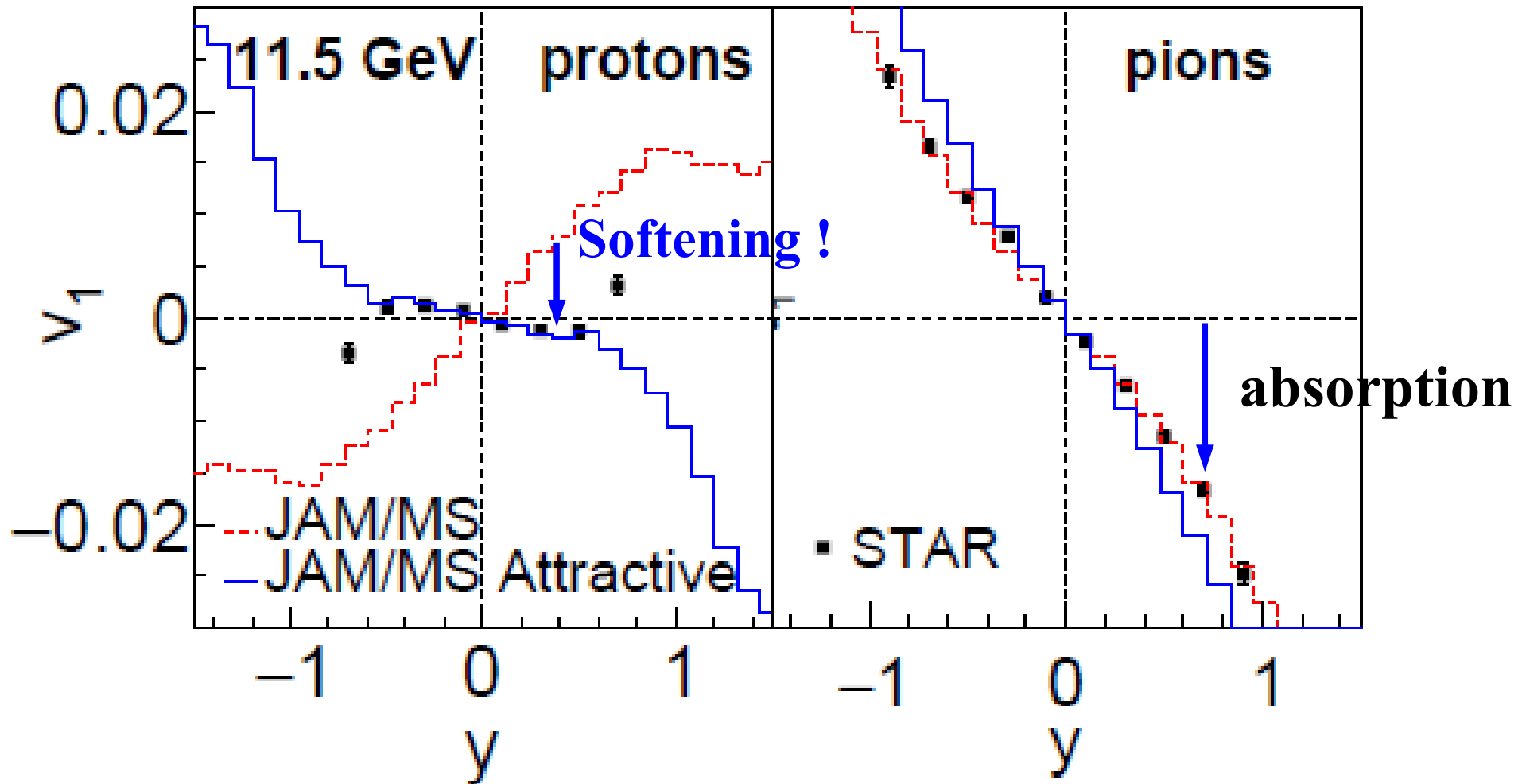
- Attractive orbit scattering can simulate EOS softening (Virial theorem)

$$P = P_f + \frac{1}{3V\Delta t} \sum_{(i,j)} \mathbf{q}_{ij} \cdot (\mathbf{r}_i - \mathbf{r}_j)$$



Mean Field + Attractive Orbit

Nara, Niemi, AO, Stöcker ('16)

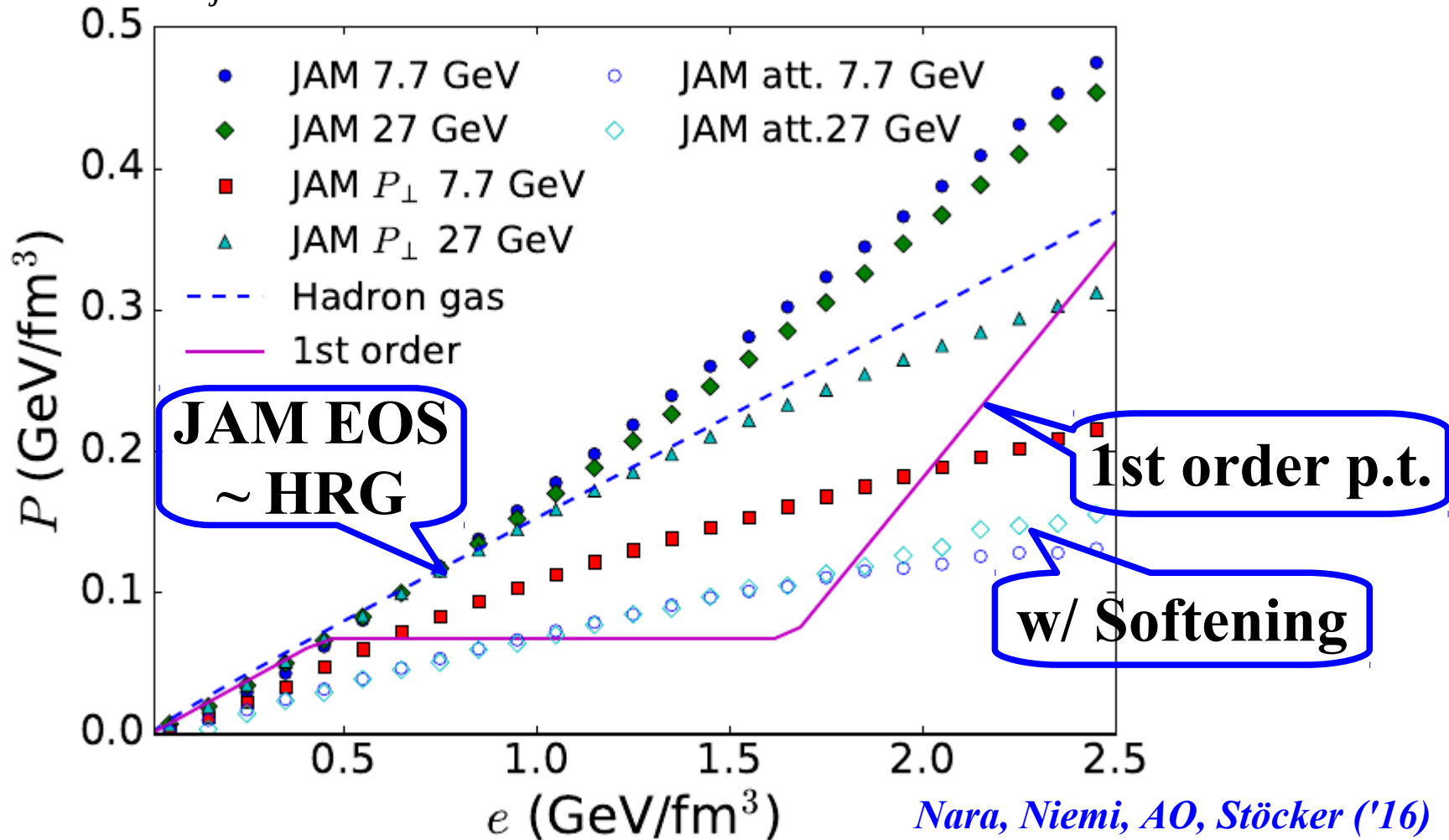


MF+Attractive Orbit make dv_1/dy negative at $\sqrt{s_{NN}} \sim 10 \text{ GeV}$

Softening of EOS by Attractive Orbits

$$\Delta P = - \frac{\rho}{3(\delta\tau_i + \delta\tau_j)} (p_i' - p_i)^\mu (x_i - x_j)_\mu$$

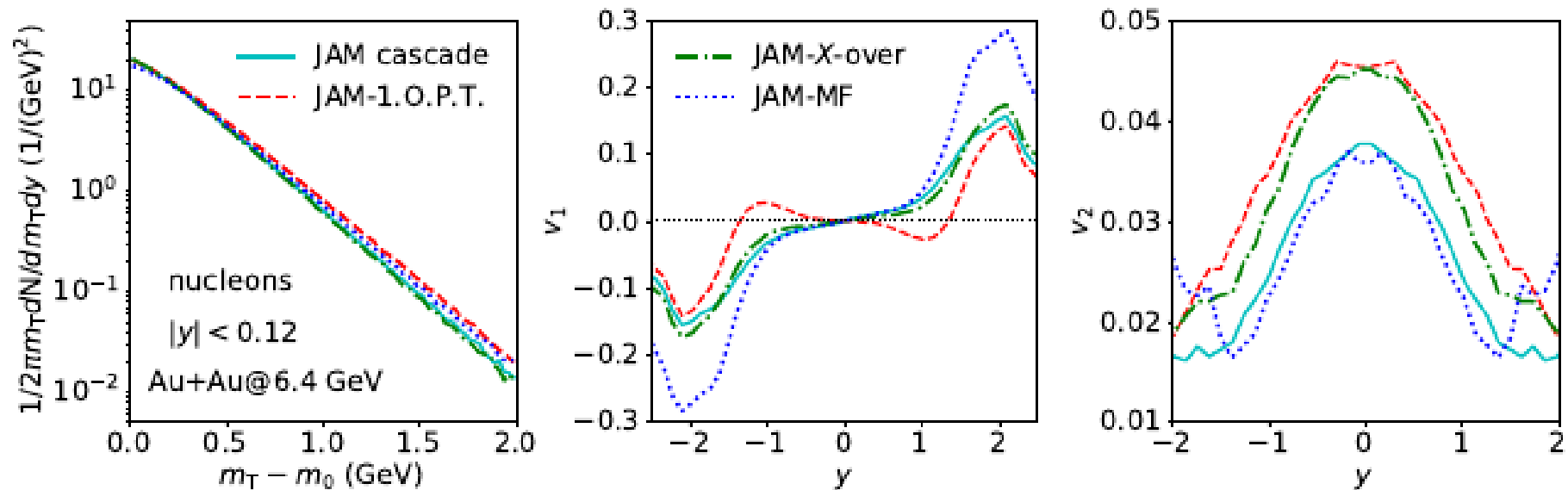
H. Sorge, PRL82('99)2048.



Pressure in simulated EOS ~ EOS-Q (e.g. Song, Heinz ('08))

Can we distinguish Crossover and 1st order ?

- **First ord.:** $T_{\text{eff}} \uparrow$, $dv_1/dy < 0$, $v_2 \uparrow$
- **Crossover:** $T_{\text{eff}} \rightarrow$, $dv_1/dy > 0$, $v_2 \uparrow$
- **Hadronic:** $T_{\text{eff}} \uparrow$, $dv_1/dy > 0$, $v_2 \rightarrow$

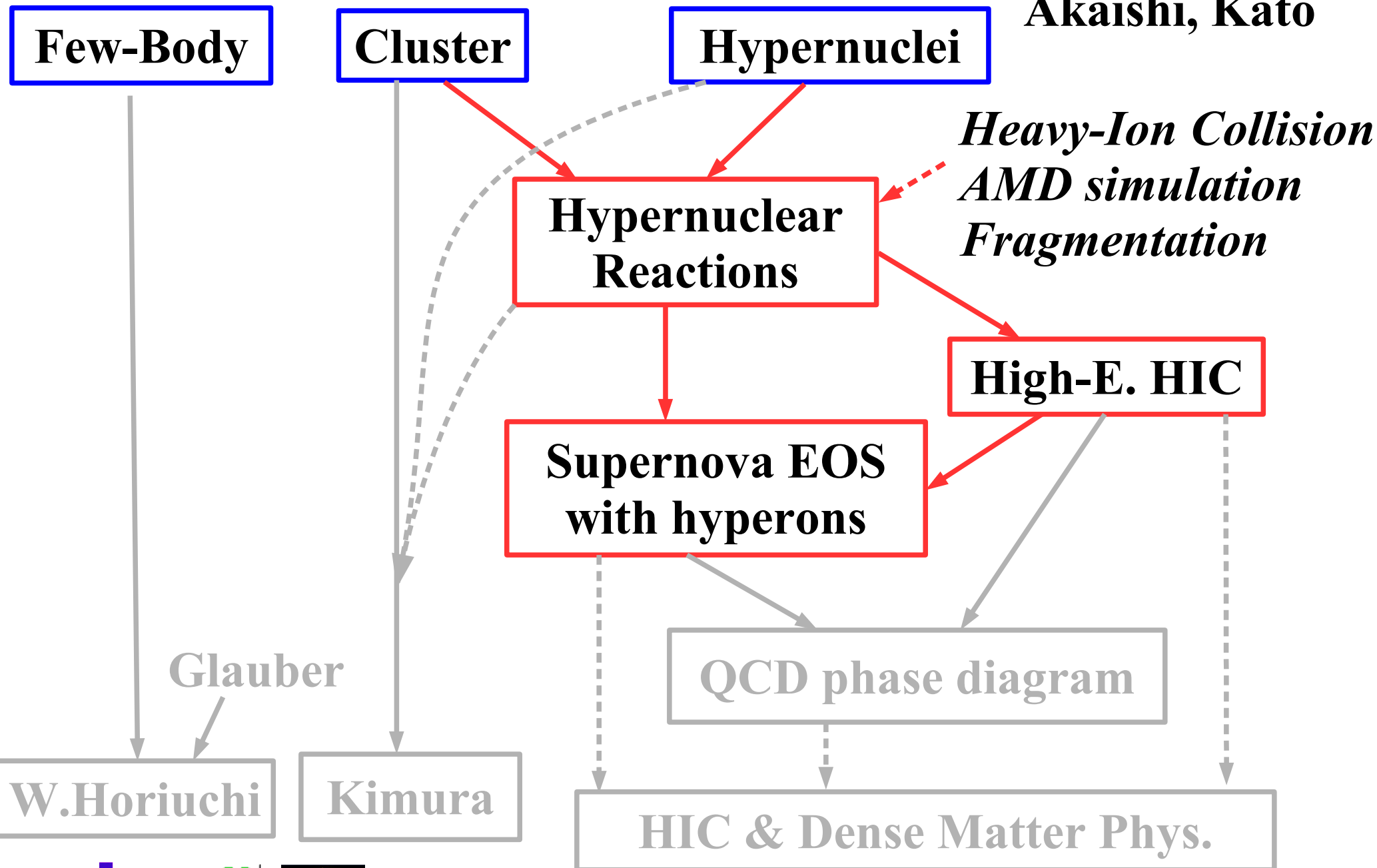


Nara, Niemi, AO, Steinheimer, Luo, Stoecker, EPJA, in press; arXiv:1708.05617

Dense Matter Physics

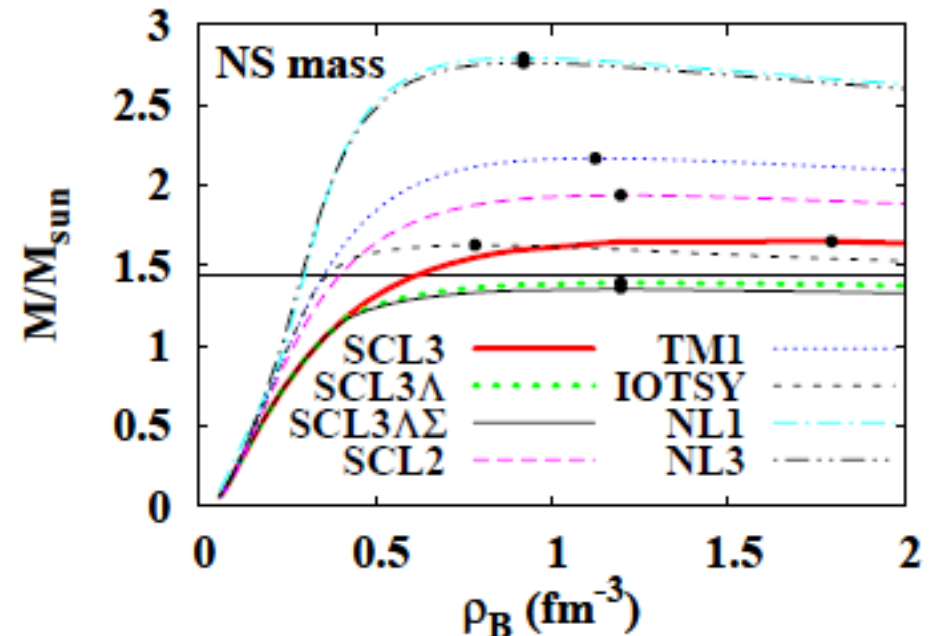
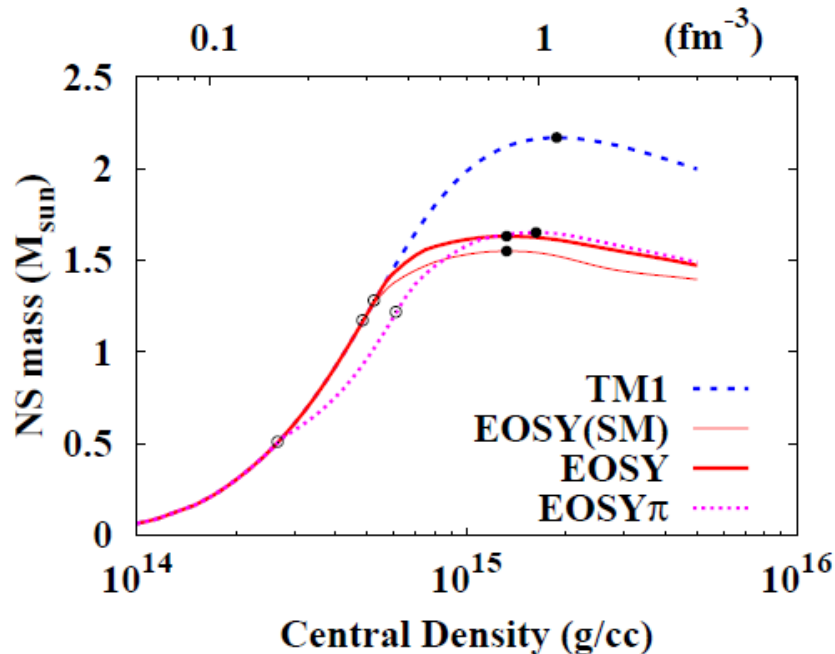
My Research Subjects in Hokkaido University

Akaishi, Kato



Supernova Matter EOS w/ Strangeness

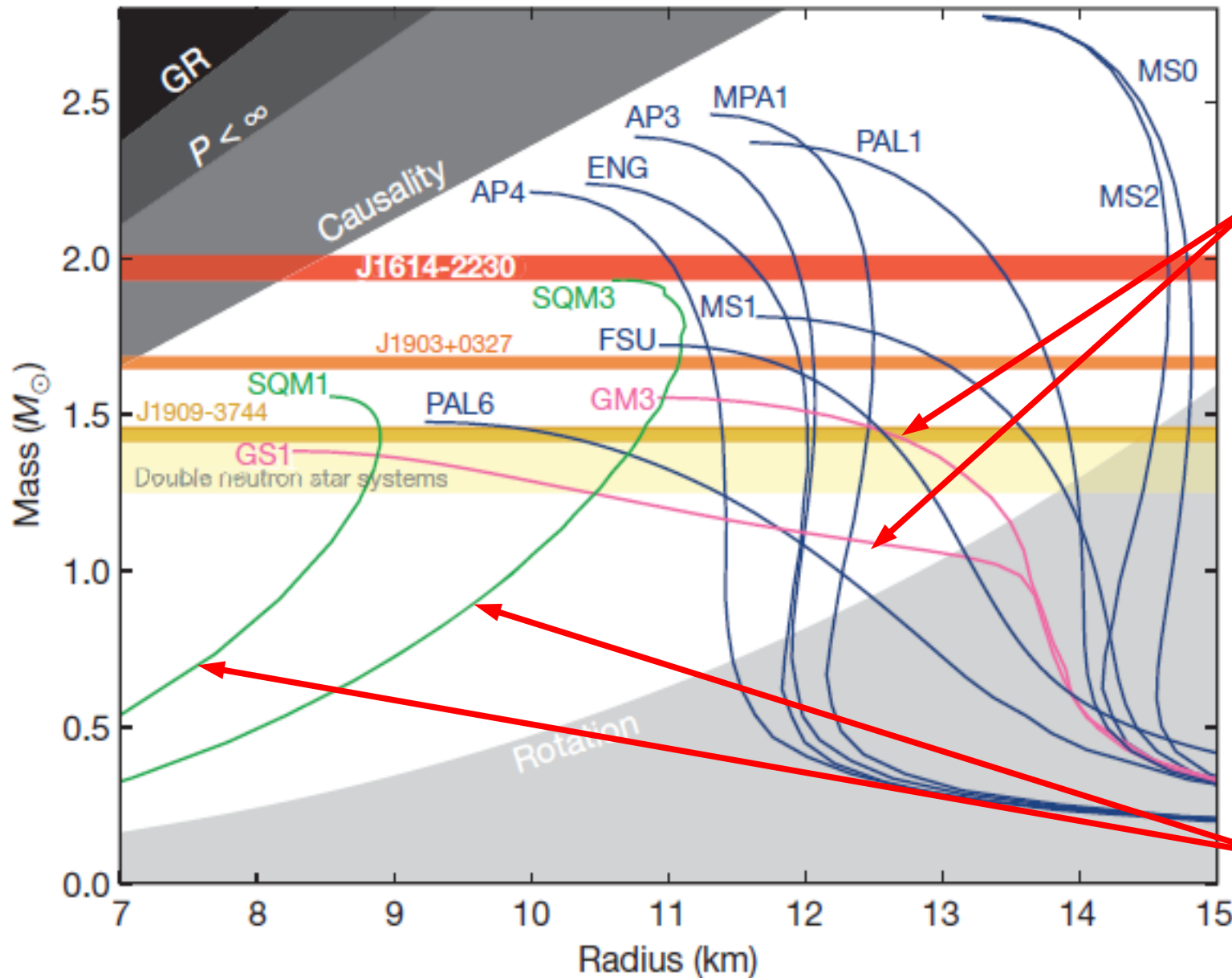
- A conclusion of hypernuclear physics
= Hyperon should appear in Neutron Star Matter at $(2-4) \rho_0$
- Let's try to make hyperonic matter EOS for Supernovae
 - Need EOS in 3D (T, ρ_B, Y_e) in a wide range.
 $0 < T < 150 \text{ MeV}, 10^{11} < \rho_B < 10^{15} \text{ g/cc}$
 - Hyperon potentials; $U_\Lambda = -30 \text{ MeV}, U_\Sigma = +30 \text{ MeV}, U_\Xi = -15 \text{ MeV}$



Ishizuka, AO, Tsubakihara, Sumiyoshi, Yamada ('08)

Hyperon Puzzle

Demorest et al., *Nature* 467 (2010) 1081 (Oct.28, 2010).



**EOS with
hyperons
or Kaons**

**Quark matter
EOS**

Hyperon Puzzle

- When we include hyperons with potentials consistent with data, EOS cannot support $2 M_{\odot}$ neutron stars.
- Proposed solutions: 3-body force, quark matter, modified gravity
- Three-Body Force including hyperons
 - Universal Three-Body Repulsion (NNN, YNN,)
Takatsuka, Nishizaki ('17), Yamamoto, Furumoto, Yasutake, Rijken ('17)
 - Pauli blocking in 2π attraction via Σ exch. (chiral EFT)
Kohno ('17), Petschauer, Haidenbauer, Meissner, Kaiser, Weise ('16)
 - Quark cluster model 3BF
Nakamoto, Suzuki ('16), AO, Kashiwa, Morita ('17)
 - RMF with multi-body coupling (no hyperons)
S. Typel et al. ('99), Steiner et al. ('13)

Relativistic Mean Field with Multi-body couplings

- $\sigma\omega\rho$ model +std. non-linear terms + multi-body couplings

Scalar polarizability (A. Thomas)

$$\mathcal{L} = \bar{\psi}(i\gamma^\mu \partial_\mu - M_N - U_s - \gamma^\mu U_\mu)\psi + \mathcal{L}_{\sigma\omega\rho},$$

ω^2 scalar (Typel)

$$\mathcal{L}_{\sigma\omega\rho} = \frac{1}{2}\partial_\mu\sigma\partial^\mu\sigma - \frac{1}{4}\omega_{\mu\nu}\omega^{\mu\nu} - \frac{1}{4}R_{\mu\nu} \cdot R^{\mu\nu} - \mathcal{V}_{\sigma\omega\rho},$$

$$U_s = -g_\sigma\sigma \left[1 - r_{\sigma\sigma}\sigma/f_\pi \right] + g_\sigma\omega^\mu\omega_\mu \left[r_{\omega\omega} - r_{\sigma\omega\omega}\sigma/f_\pi \right],$$

DD coupling (Ring)

$$U_\mu = g_\omega\omega_\mu \left[1 - r_{\sigma\omega}\sigma/f_\pi + r_{\omega 3}\omega^\nu\omega_\nu/f_\pi^2 \right] + g_\rho\tau \cdot R_\mu \left[1 - r_{\sigma\rho}\sigma/f_\pi + r_{\omega\rho}\omega^\nu\omega_\nu/f_\pi^2 \right],$$

$$\mathcal{V}_{\sigma\omega\rho} = \frac{1}{2}m_\sigma^2\sigma^2 - a_\sigma f \log(\sigma/f_\pi) + \frac{1}{4}c_{\sigma 4}\sigma^4 + \frac{1}{3}c_{\sigma 3}f_\pi\sigma^3$$

$$- \frac{1}{2}m_\omega^2\omega^\mu\omega_\mu \left[1 - c_{\sigma\omega}\sigma/f_\pi \right] - \frac{1}{4}c_{\omega 4}(\omega^\mu\omega_\mu)^2$$

ρ^4 term

$$- \frac{1}{2}m_\rho^2 R^\mu \cdot R_\mu \left[1 - c_{\sigma\rho}\sigma/f_\pi + c_{\omega\rho}\omega^\mu\omega_\mu/f_\pi^2 \right] - \frac{1}{4}c_{\rho 4}(R^\mu \cdot R_\mu)^2,$$

DD meson mass (e.g. Steiner, Fischer, Hempel)

Relativistic Mean Field with Multi-body couplings

■ Phen. Approach: RMF w/ Multi-body coupling

● Naive dimensional analysis (NDA) and naturalness

Manohar, Georgi ('84)

The vertex is called “natural” if $C \sim 1$ (consistent with pQCD).

$$L_{\text{int}} \sim (f_{\pi} \Lambda)^2 \sum_{l,m,n,p} \frac{C_{lmnp}}{m! n! p!} \left(\frac{\bar{\psi} \Gamma \psi}{f_{\pi}^2 \Lambda} \right)^l \left(\frac{\sigma}{f_{\pi}} \right)^m \left(\frac{\omega}{f_{\pi}} \right)^n \left(\frac{R}{f_{\pi}} \right)^p$$

● FST truncation

*R. J. Furnstahl, B. D. Serot, H. B. Tang,
NPA615 ('97)441.*

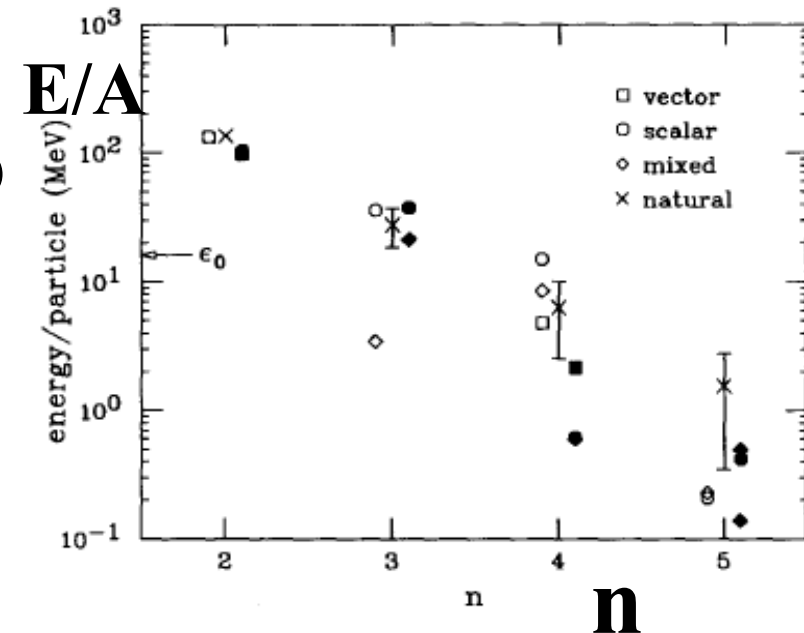
Truncation the index

$$n = B/2 + M + D$$

(B: baryon, M: Non NG boson, D: derivatives)

Natural $\rightarrow V \sim \rho^n/n!$

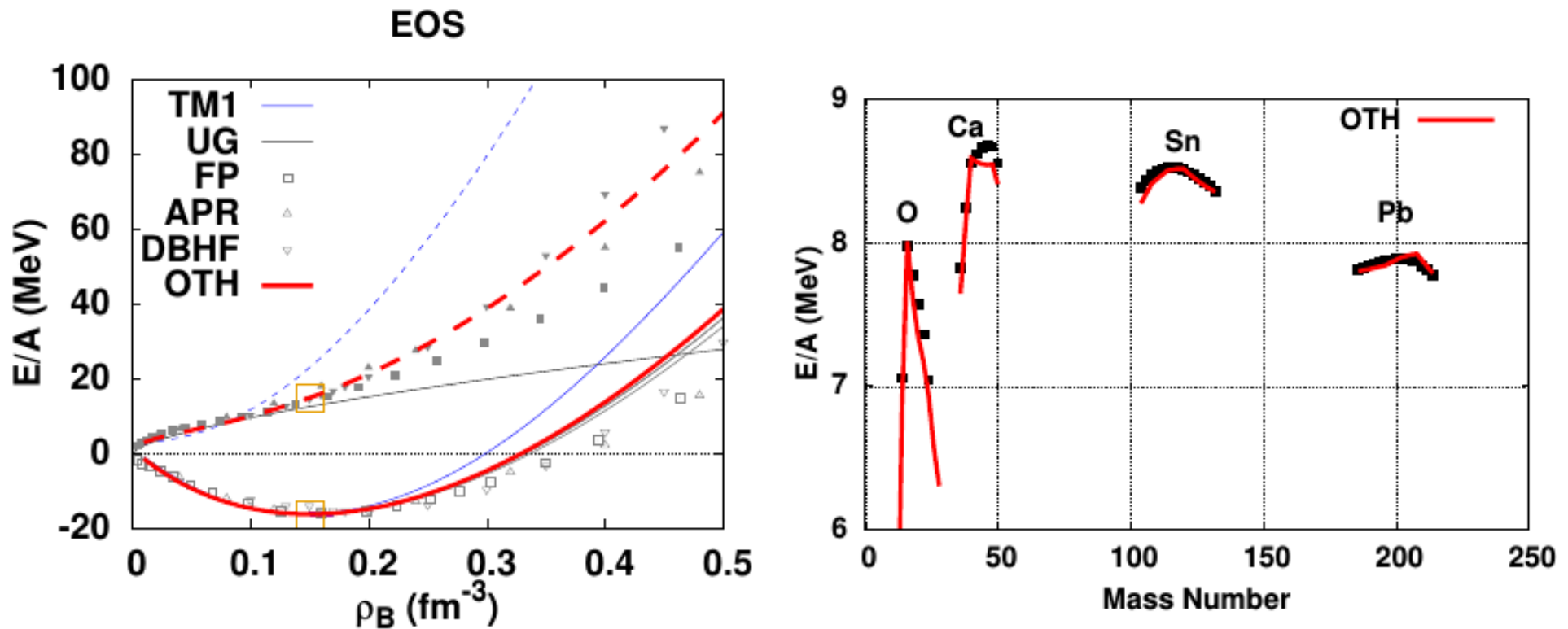
\rightarrow small for large n



Simultaneous Fit to EOS and Finite Nuclei

■ Fitting procedure

= Fit finite nuclear binding energies and charge rms radius under the constraint of given (ρ_0, E_0, K, S_0, L) .

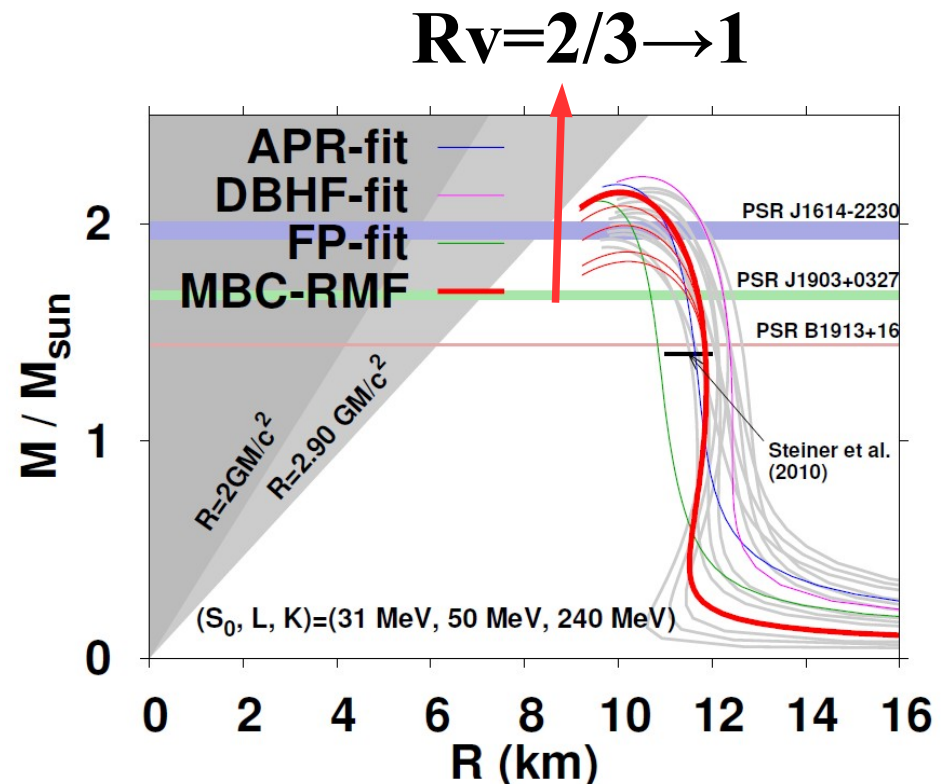
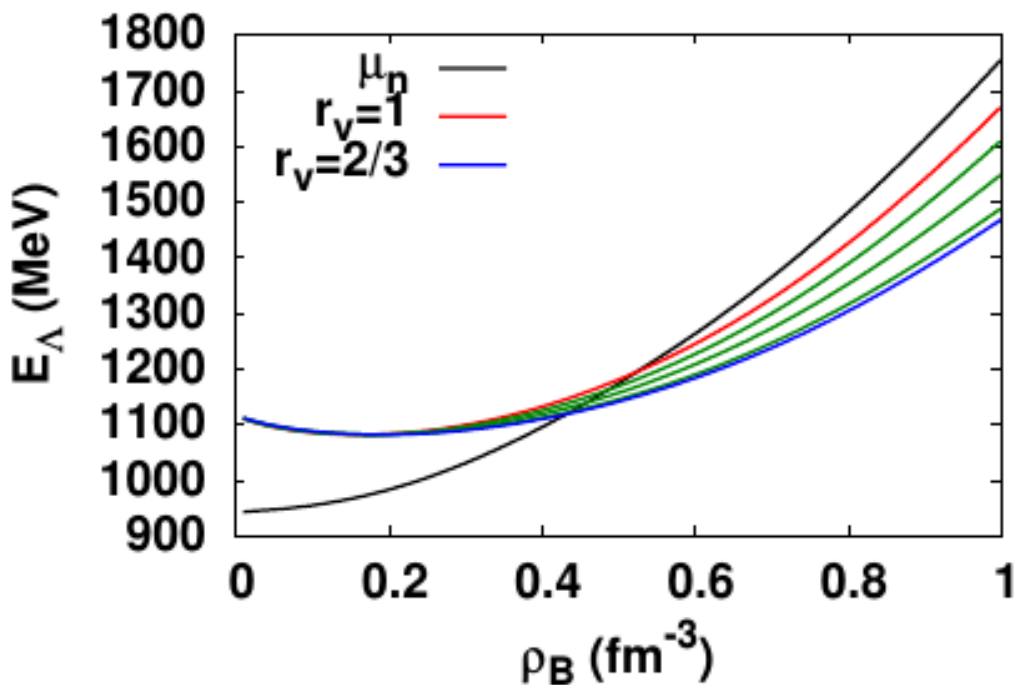


$(S_0, L, K) = (31 \text{ MeV}, 50 \text{ MeV}, 240 \text{ MeV})$

AO, Tsub=akihara, Harada (in prep.)

Hypernuclei and Neutron Star MR

- $R_v = g_{\omega\Lambda} / g_{\omega N} = 2/3 - 1$ is chosen, and $g_{\sigma\Lambda} / g_{\sigma N}$ is fitted to data.
 (Other parameters are assumed to be the same.)
 → Λ emerges at $\rho = 0.4 - 0.5 \text{ fm}^{-3}$
 2 M_{\square} neutron stars may be supported with $R_v > 0.8$
 (Depends on nuclear matter EOS)



AO, Tsubakihara, Harada (in prep.)

Can we distinguish ?

■ Density dependence of U_{Λ}

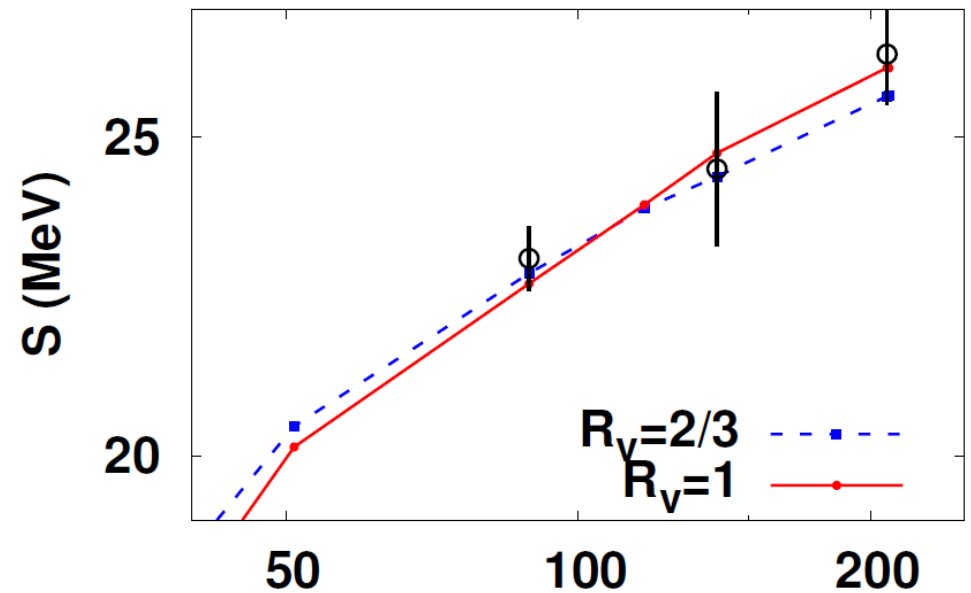
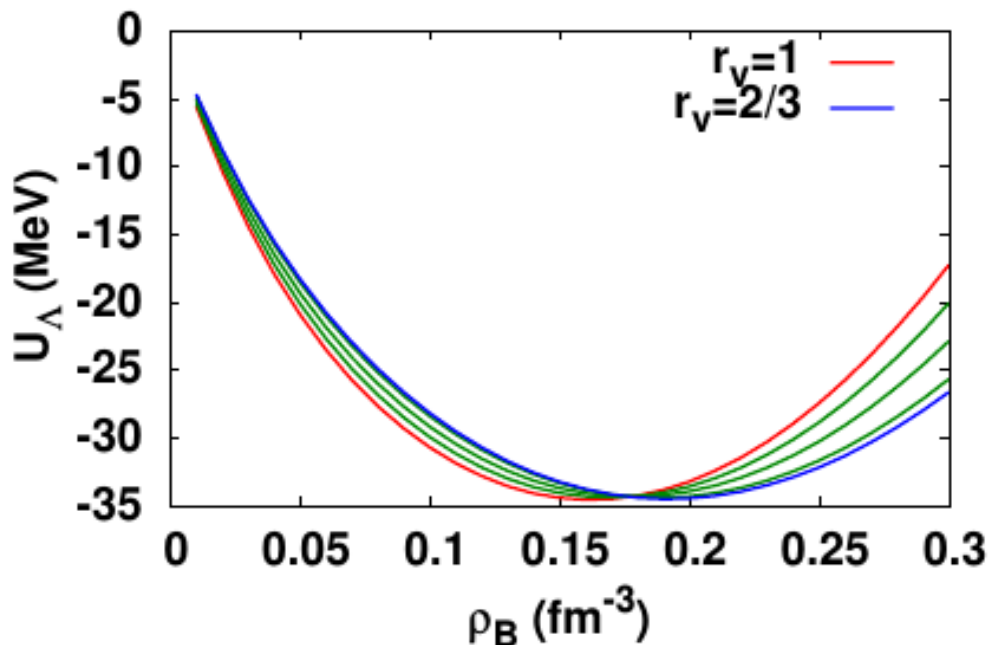
- $dU_{\Lambda}/d\rho$ turns to be positive at around ρ_0

Kohno ('17), Petschauer, Haidenbauer, Meissner, Kaiser, Weise ('16)

- $R_V=2/3$ and 1 leads to the difference of S_{Λ} of a few 100 keV

→ sub MeV hypernuclear spectroscopy is necessary

Isaka, Yamamoto, Rijken('17); Yamamoto, Furumoto, Yasutake, Rijken('17)



AO, Tsubakihara, Harada (in prep.)

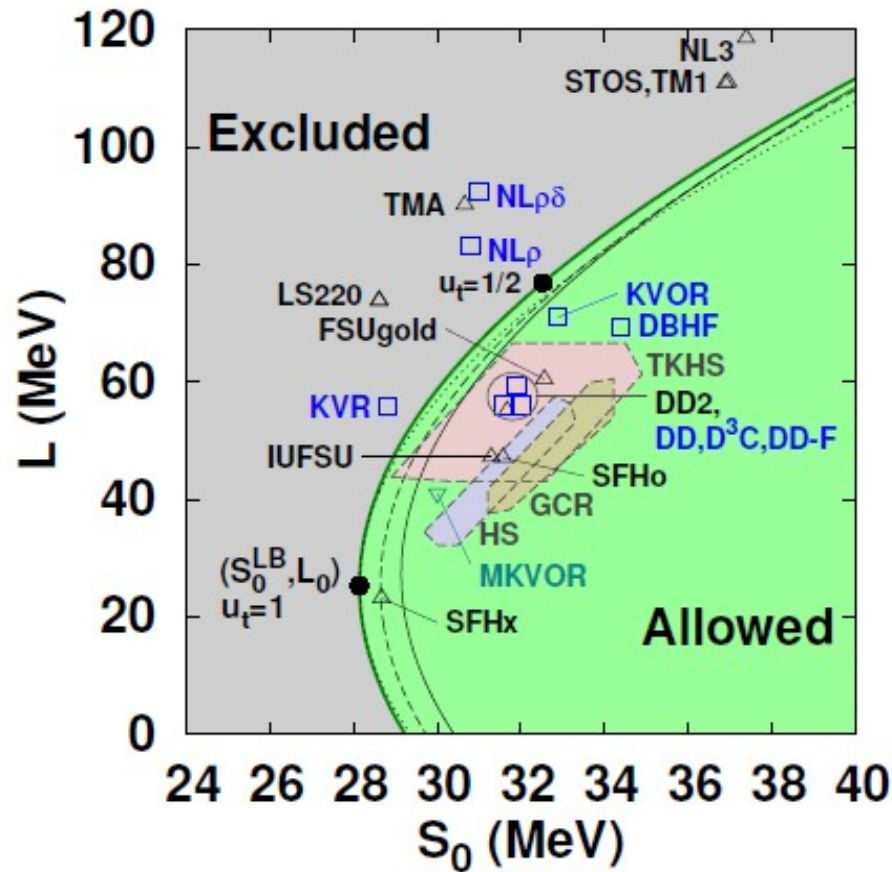
Summary

- We may have seen QCD phase transition (1st or 2nd) signals at BES (or J-PARC) energies in baryon number cumulants and v_1 slope. The transport model (JAM) is utilized to elucidate the EOS softening.
- In order to solve the hyperon puzzle based on data, we need models which describes normal nuclei, hypernuclei, and nuclear matter in a consistent manner. RMF with multi-body coupling may be a handy framework. Turn over density ($dU_{\Lambda}/d\rho = 0$) is found to be around ρ_0 , then massive NS may be supported even with hyperons.
- I enjoyed the time in Hokkaido University. Prof. Kato allowed me to work independently, and the high activity in Hokkaido U. required us to make works with originality and ambition.

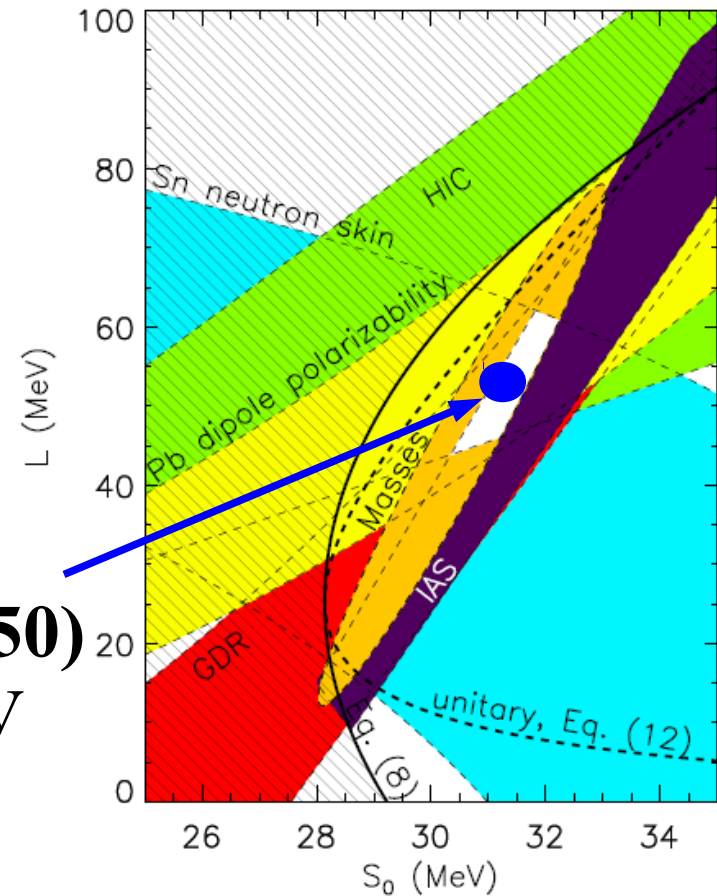
*Kato-san, Congratulations
for your age Seventy,
and thank you very much
for your encouragements !*



Symmetry Energy Constraints



(31,50)
MeV

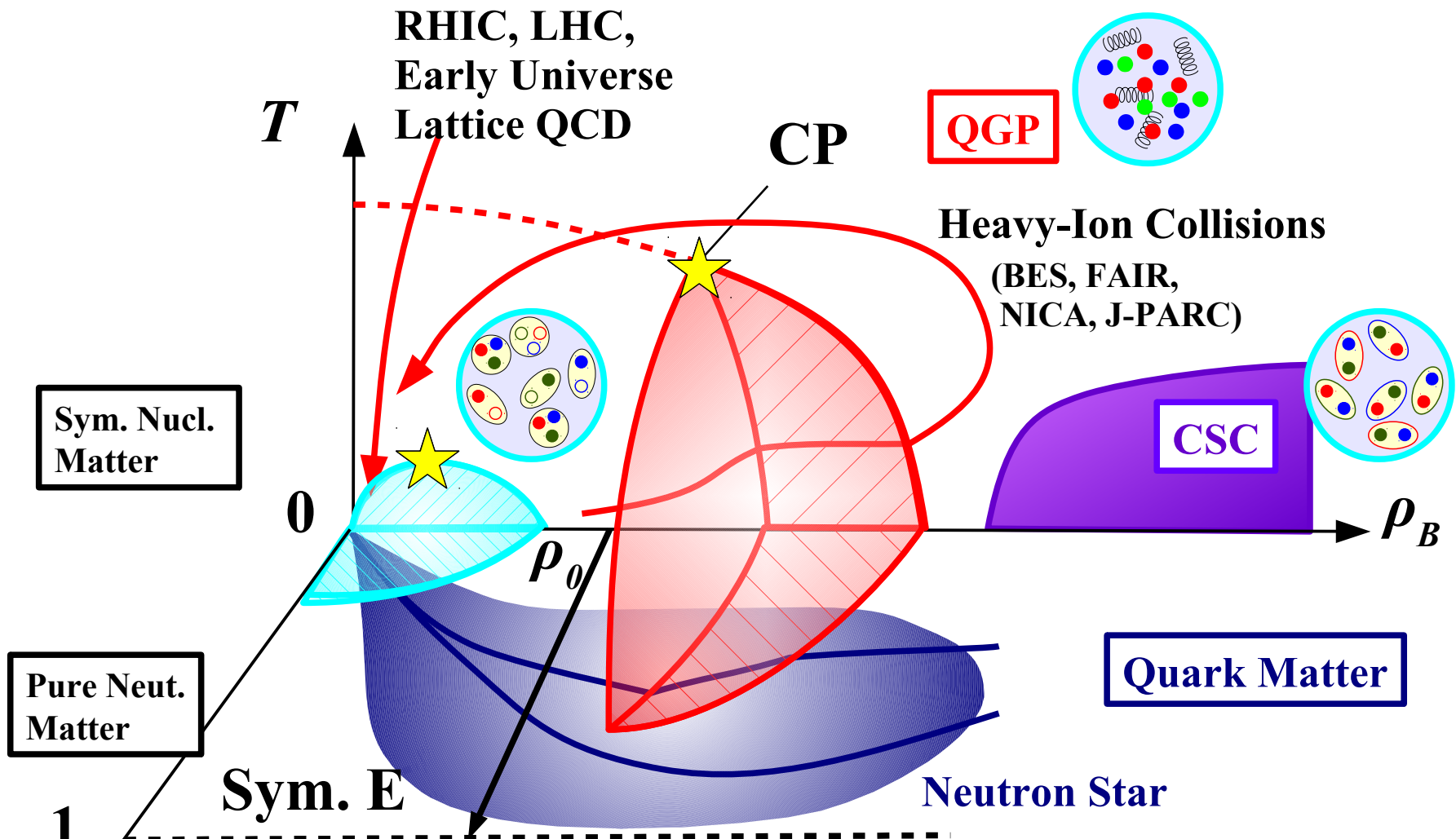


Many of EOSs in active astrophysical use do not satisfy recent symmetry energy constraint or $2 M_{\odot}$ constraint.
 → *SFHo, SHFx, DD2*

Tews, Lattimer, AO, Kolomeitsev ('17)

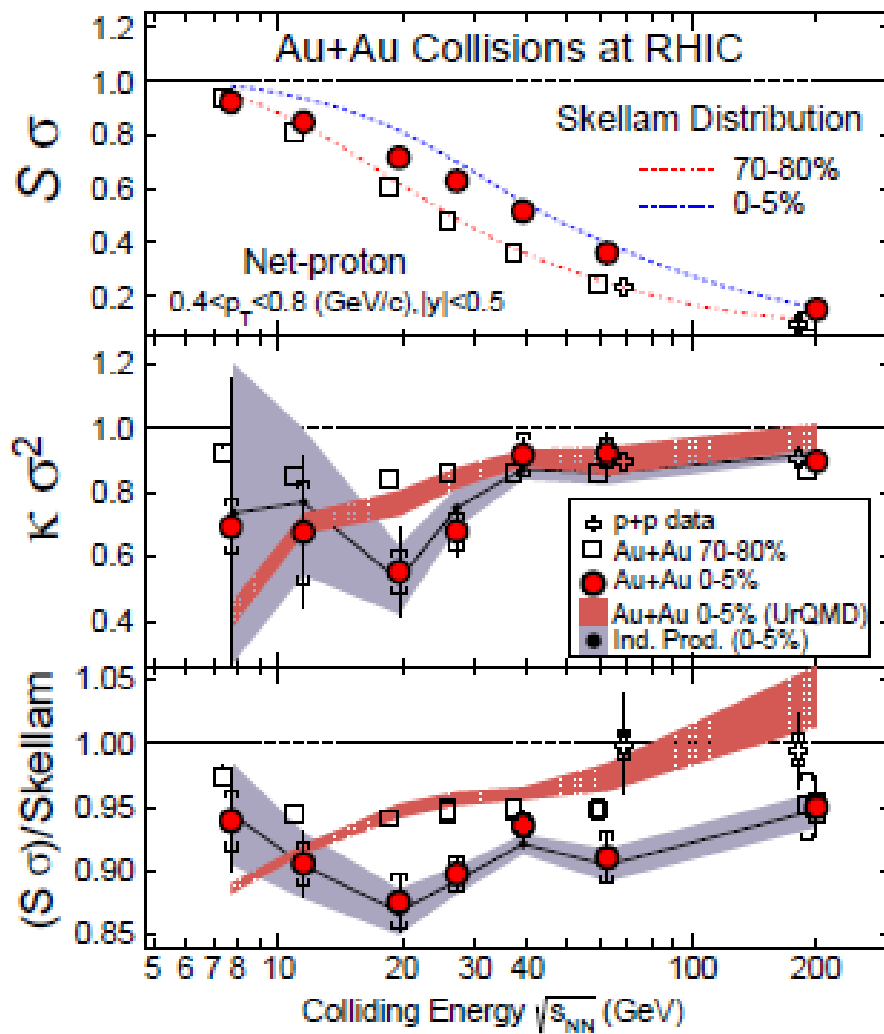
Ohnishi @ WCNP2017, Oct. 27, 2017 29

QCD Phase Diagram

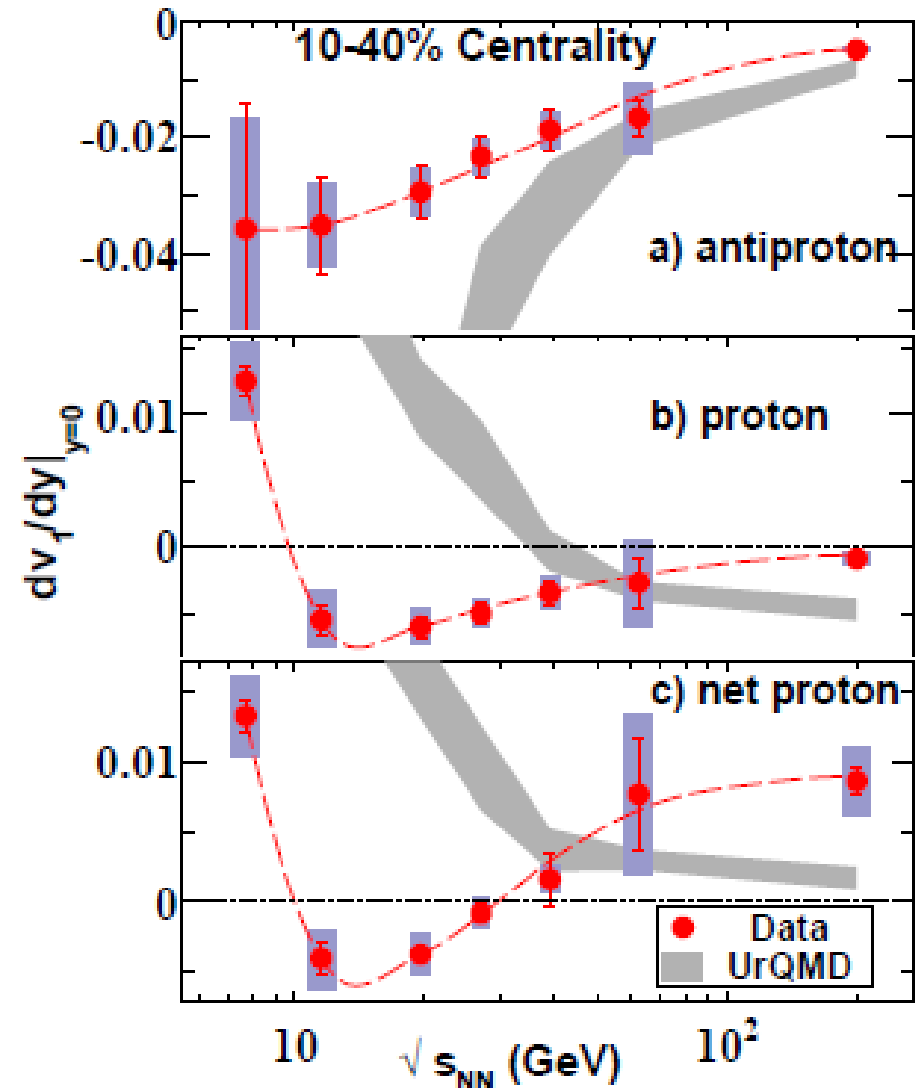


$$\delta = (N-Z)/A \quad (\text{or } Y_Q(\text{hadron}) = Q_h/B \sim (1-\delta)/2)$$

Net-Proton Number Cumulants & Directed Flow



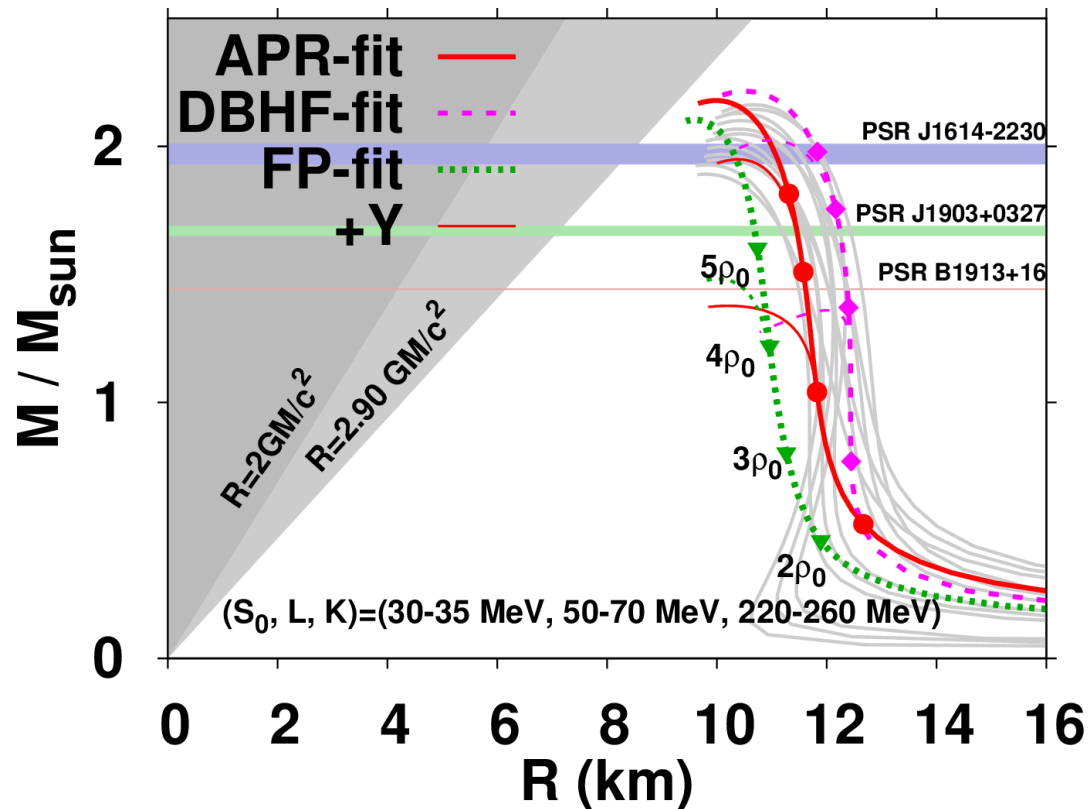
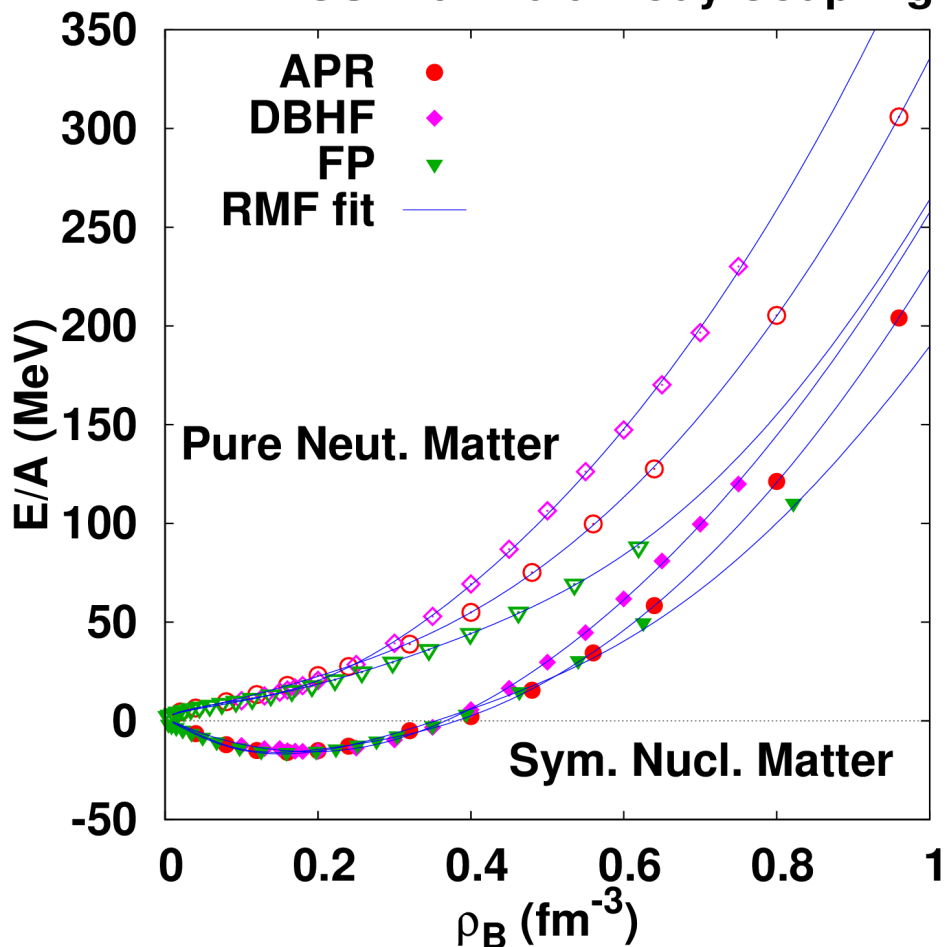
STAR Collab. PRL 112('14)032302



STAR Collab., PRL 112('14)162301.

Fitting “Ab initio” EOS via RMF

RMF EOS with Multi-Body Coupling



RMF fitting EOS does not necessarily describe finite nuclei...

AO, Tsubakihara, Harada ('16, NIC proc.)

My Research Subjects in Hokkaido University

Akaishi, Kato

Few-Body

Cluster

Hypernuclei

Hypernuclear Reactions

*Heavy-Ion Collision
AMD simulation
Fragmentation*

High-E. HIC

Supernova EOS
with hyperons

QCD phase diagram

Glauber

W.Horiuchi

Kimura

HIC & Dense Matter Phys.

Virial Theorem

■ Virial

$$G = \sum_i \mathbf{p}_i \cdot \mathbf{r}_i$$

$$\rightarrow \frac{dG}{dt} = \underbrace{\sum_i \mathbf{p}_i \cdot \mathbf{v}_i}_{\text{Kinetic}} - \underbrace{\sum_i \nabla_i U \cdot \mathbf{r}_i}_{\text{Potential}} + \underbrace{\frac{1}{\Delta t} \sum_{\text{collision}} \mathbf{q}_i \cdot (\mathbf{r}_i - \mathbf{r}_j)}_{\text{Pressure from Collisions}} = 3VP$$

■ Attractive / Repulsive Orbit Scatterings

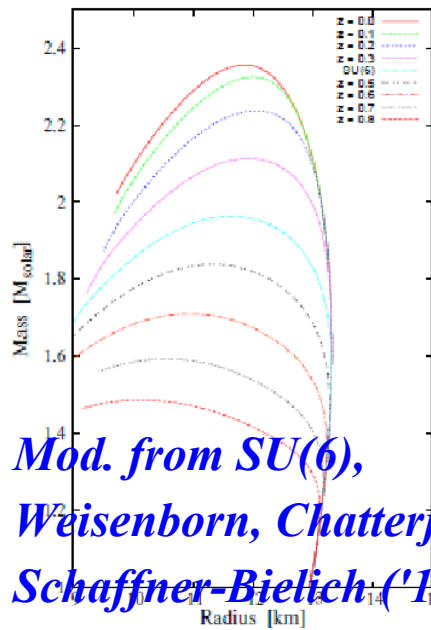
- Random choice of scatt. angle \rightarrow No effect on pressure
- Attractive orbits $\rightarrow \Delta P < 0$ (softening)
- Repulsive orbits $\rightarrow \Delta P > 0$ (hardening)

■ Boltzmann Eq. simulating a given EOS

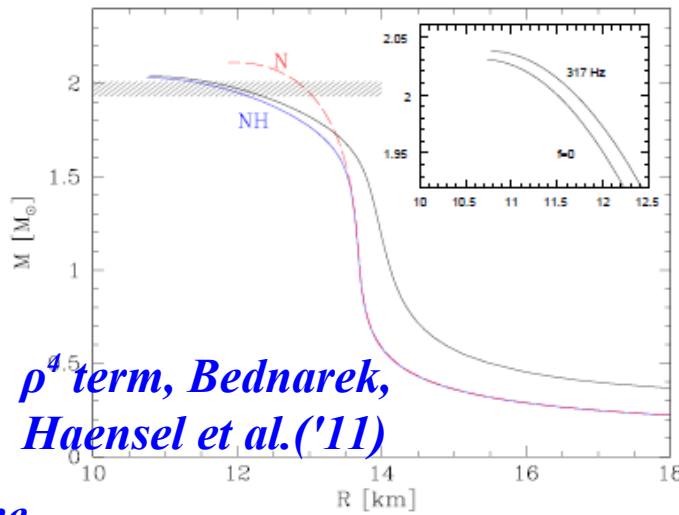
$P > P(\varepsilon) \rightarrow$ Attractive orbit, $P < P(\varepsilon) \rightarrow$ Repulsive orbit

If collisions are frequent enough, we can simulate MF effects in Boltzmann equation.

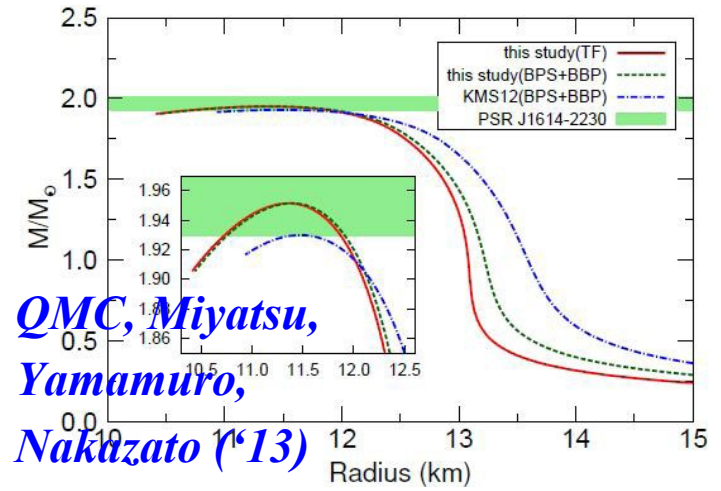
Massive Neutron Stars with Hyperons



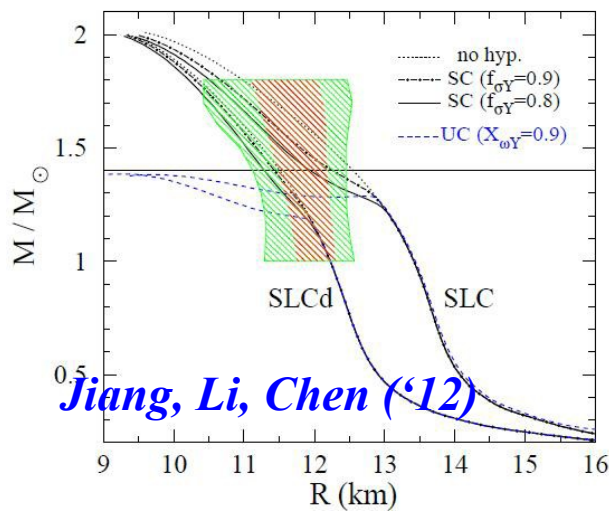
Mod. from SU(6),
Weisenborn, Chatterjee,
Schaffner-Bielich ('11)



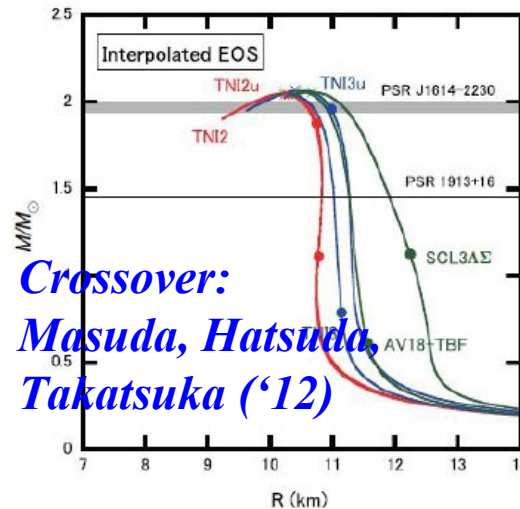
ρ^4 term, Bednarek,
Haensel et al. ('11)



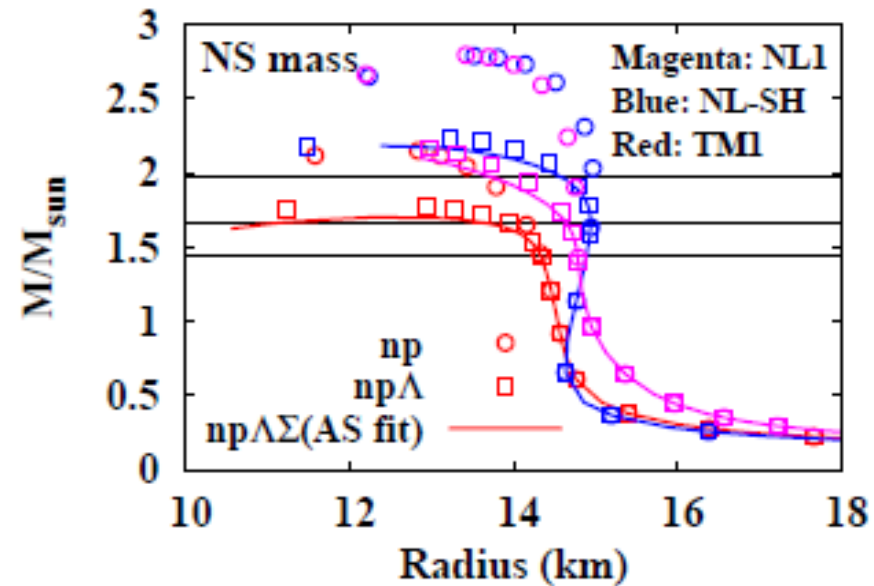
QMC, Miyatsu,
Yamamuro,
Nakazato ('13)



Jiang, Li, Chen ('12)



Crossover:
Masuda, Hatsuda,
Takatsuka ('12)



Tsubakihara, Harada, AO, arXiv:1402.0979

Relativistic Mean Field with Multi-body couplings

■ $\sigma\omega\rho$ model +std. non-linear terms + multi-body couplings

$$\mathcal{L} = \bar{\psi}(i\gamma^\mu \partial_\mu - M_N - U_s - \gamma^\mu U_\mu)\psi + \mathcal{L}_{\sigma\omega\rho} ,$$

$$\mathcal{L}_{\sigma\omega\rho} = \frac{1}{2}\partial_\mu\sigma\partial^\mu\sigma - \frac{1}{4}\omega_{\mu\nu}\omega^{\mu\nu} - \frac{1}{4}R_{\mu\nu} \cdot R^{\mu\nu} - \mathcal{V}_{\sigma\omega\rho} ,$$

$$U_s = -g_\sigma\sigma \left[1 - r_{\sigma\sigma}\sigma/f_\pi \right] + g_\sigma\omega^\mu\omega_\mu \left[r_{\omega\omega} - r_{\sigma\omega\omega}\sigma/f_\pi \right] ,$$

$$U_\mu = g_\omega\omega_\mu \left[1 - r_{\sigma\omega}\sigma/f_\pi + r_{\omega 3}\omega^\nu\omega_\nu/f_\pi^2 \right] + g_\rho\tau \cdot R_\mu \left[1 - r_{\sigma\rho}\sigma/f_\pi + r_{\omega\rho}\omega^\nu\omega_\nu/f_\pi^2 \right] ,$$

$$\begin{aligned} \mathcal{V}_{\sigma\omega\rho} = & \frac{1}{2}m_\sigma^2\sigma^2 \left[-a_\sigma f \log(\sigma/f_\pi) + \frac{1}{4}c_{\sigma 4}\sigma^4 + \frac{1}{3}c_{\sigma 3}f_\pi\sigma^3 \right] \\ & - \frac{1}{2}m_\omega^2\omega^\mu\omega_\mu \left[1 - c_{\sigma\omega}\sigma/f_\pi \right] - \frac{1}{4}c_{\omega 4}(\omega^\mu\omega_\mu)^2 \\ & - \frac{1}{2}m_\rho^2 R^\mu \cdot R_\mu \left[1 - c_{\sigma\rho}\sigma/f_\pi + c_{\omega\rho}\omega^\mu\omega_\mu/f_\pi^2 \right] - \frac{1}{4}c_{\rho 4}(R^\mu \cdot R_\mu)^2 , \end{aligned}$$

Relativistic Mean Field with Multi-body couplings

- $\sigma\omega\rho$ model +std. non-linear terms + multi-body couplings

Scalar polarizability (A. Thomas)

$$\mathcal{L} = \bar{\psi}(i\gamma^\mu \partial_\mu - M_N - U_s - \gamma^\mu U_\mu)\psi + \mathcal{L}_{\sigma\omega\rho},$$

ω^2 scalar (Typel)

$$\mathcal{L}_{\sigma\omega\rho} = \frac{1}{2}\partial_\mu\sigma\partial^\mu\sigma - \frac{1}{4}\omega_{\mu\nu}\omega^{\mu\nu} - \frac{1}{4}R_{\mu\nu} \cdot R^{\mu\nu} - \mathcal{V}_{\sigma\omega\rho},$$

$$U_s = -g_\sigma\sigma \left[1 - r_{\sigma\sigma}\sigma/f_\pi \right] + g_\sigma\omega^\mu\omega_\mu \left[r_{\omega\omega} - r_{\sigma\omega\omega}\sigma/f_\pi \right],$$

DD coupling (Ring)

$$U_\mu = g_\omega\omega_\mu \left[1 - r_{\sigma\omega}\sigma/f_\pi + r_{\omega 3}\omega^\nu\omega_\nu/f_\pi^2 \right] + g_\rho\tau \cdot R_\mu \left[1 - r_{\sigma\rho}\sigma/f_\pi + r_{\omega\rho}\omega^\nu\omega_\nu/f_\pi^2 \right],$$

$$\mathcal{V}_{\sigma\omega\rho} = \frac{1}{2}m_\sigma^2\sigma^2 - a_\sigma f \log(\sigma/f_\pi) + \frac{1}{4}c_{\sigma 4}\sigma^4 + \frac{1}{3}c_{\sigma 3}f_\pi\sigma^3$$

$$- \frac{1}{2}m_\omega^2\omega^\mu\omega_\mu \left[1 - c_{\sigma\omega}\sigma/f_\pi \right] - \frac{1}{4}c_{\omega 4}(\omega^\mu\omega_\mu)^2$$

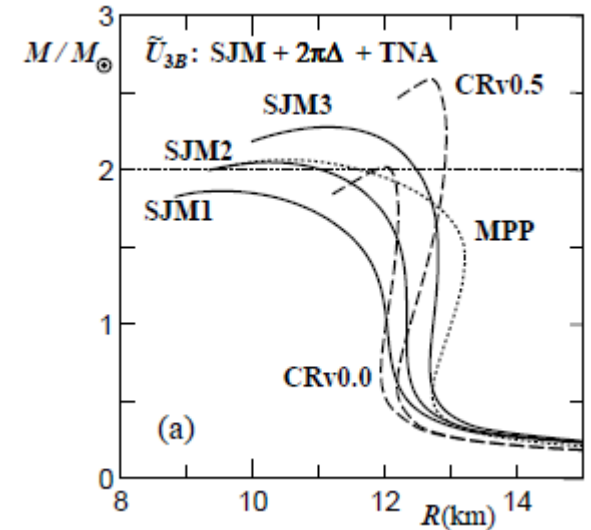
ρ^4 term

$$- \frac{1}{2}m_\rho^2 R^\mu \cdot R_\mu \left[1 - c_{\sigma\rho}\sigma/f_\pi + c_{\omega\rho}\omega^\mu\omega_\mu/f_\pi^2 \right] - \frac{1}{4}c_{\rho 4}(R^\mu \cdot R_\mu)^2,$$

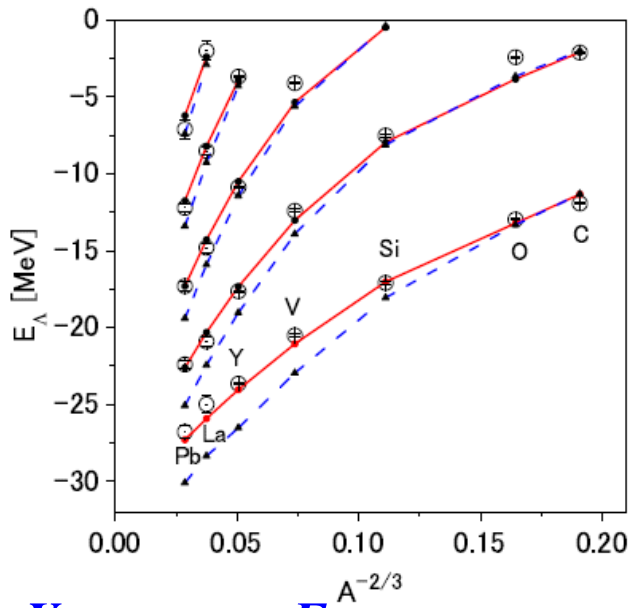
DD meson mass (e.g. Steiner, Fischer, Hempel)

Hyperon Puzzle

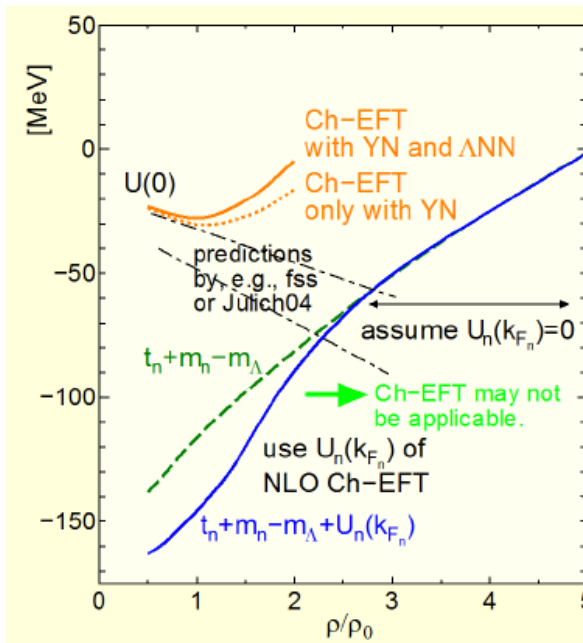
- When we include hyperons with potentials consistent with data, EOS cannot support $2 M_{\odot}$ neutron stars.
- Proposed solutions: 3-body force, quark matter, modified gravity



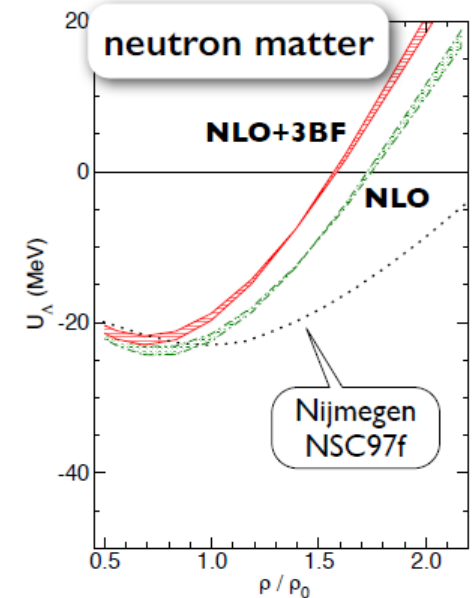
Takatsuka, Nishizaki ('17)



Yamamoto, Furumoto, Yasutake, Rijken ('17)



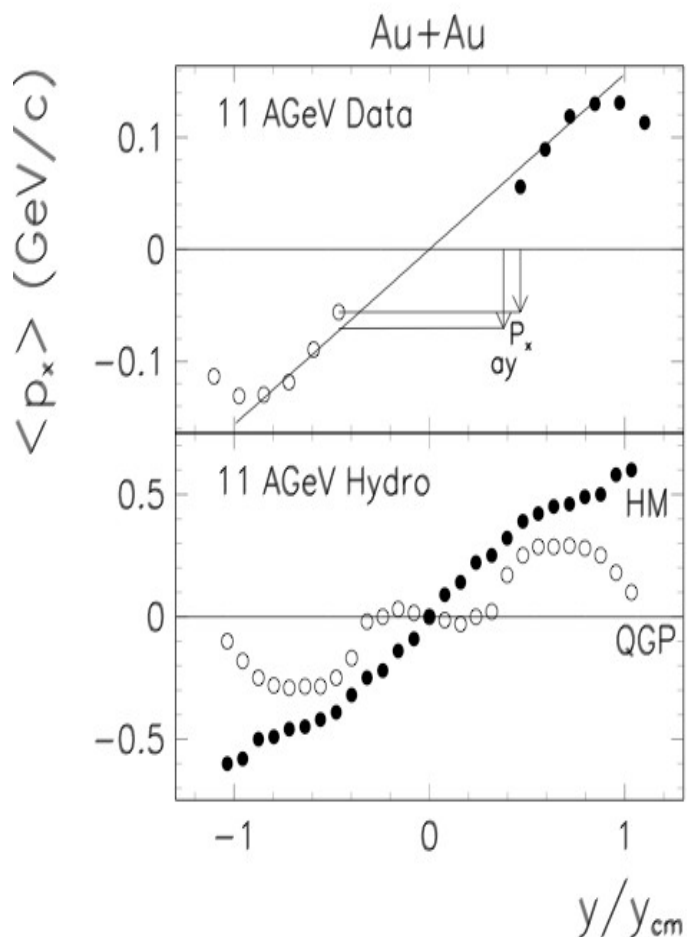
Kohno (SCHDM2017)



Haidenbauer, Meissner, Kaiser, Weise ('17)

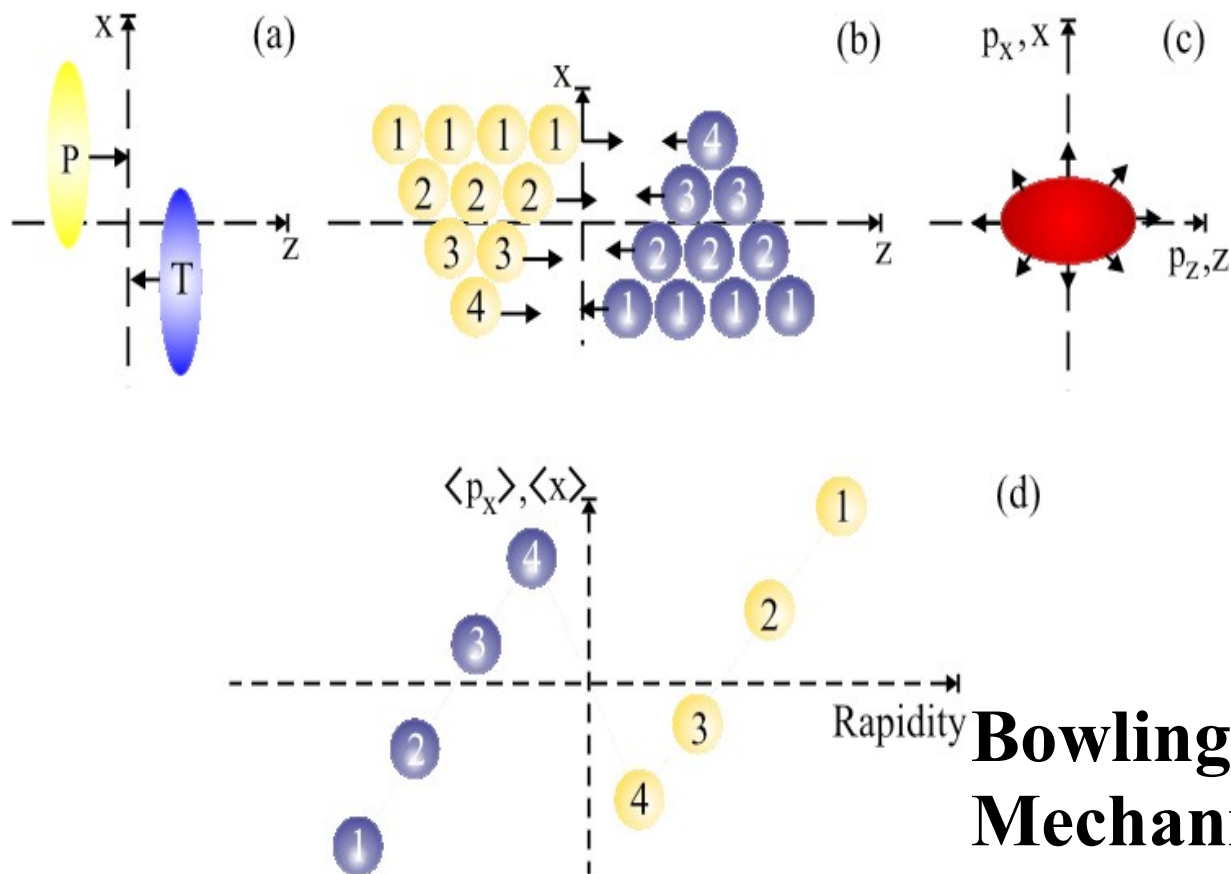
Does the “Wiggle” signal the QGP ?

- Hydro predicts wiggle with QGP EOS.



L. P. Csernai, D. Röhrich, PLB 45 (1999), 454.

- Baryon stopping + Positive space-momentum correlation leads wiggle (w/o QGP)

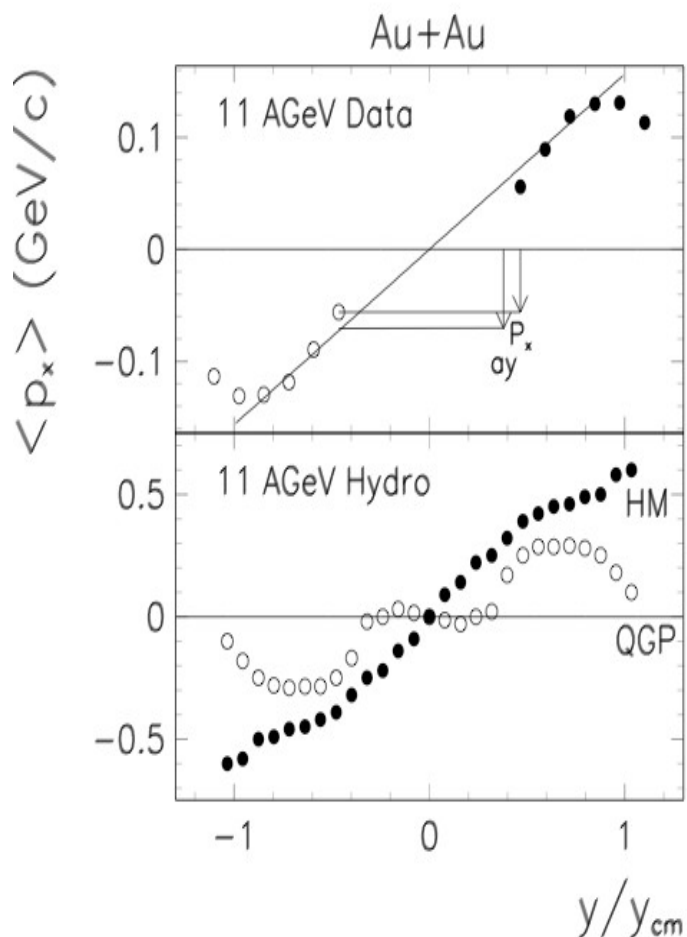


Bowling Pin Mechanism

R. Snellings, H. Sorge, S. Voloshin, F. Wang, N. Xu, PRL (84) 2803(2000)

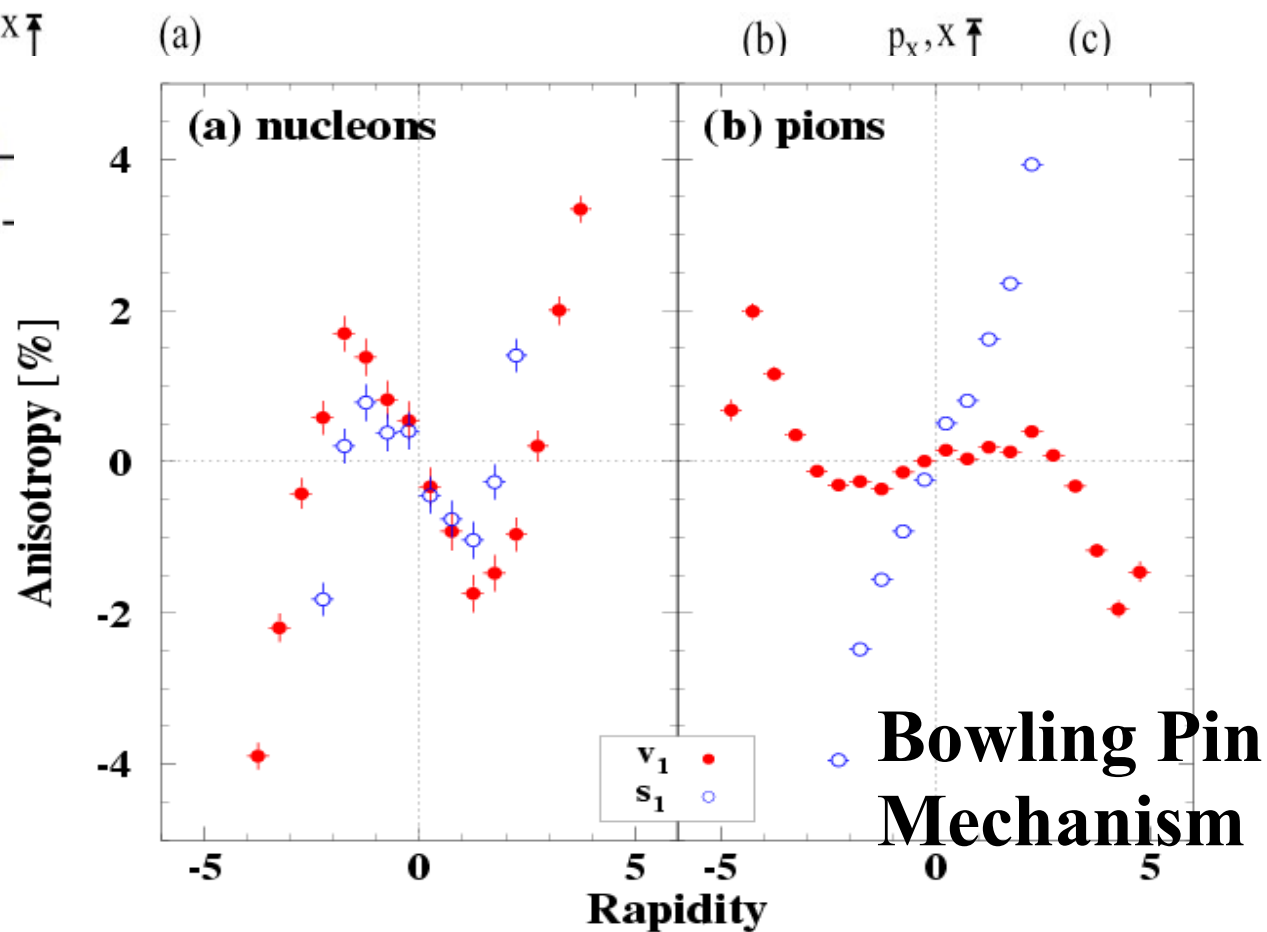
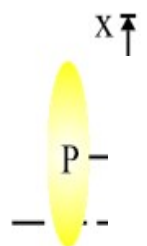
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L. P. Csernai, D. Röhrich, PLB 45 (1999), 454.

- Baryon stopping + Positive space-momentum correlation leads wiggle (w/o QGP)



Bowling Pin Mechanism

R. Snellings, H. Sorge, S. Voloshin, F. Wang, N. Xu, PRL (84) 2803(2000)

Does Directed Flow Collapse Signal Phase Tr. ?

- **Negative dv_1/dy at high-energy ($\sqrt{s_{NN}} > 20$ GeV)**
 - **Geometric origin (bowling pin mechanism), not related to FOPT**
R.Snellings, H.Sorge, S.Voloshin, F.Wang, N. Xu, PRL84,2803('00)
- **Negative dv_1/dy at $\sqrt{s_{NN}} \sim 10$ GeV**
 - **Yes, in three-fluid simulations. → Thermalization ?**
Y. B. Ivanov and A. A. Soldatov, PRC91('15)024915
 - **No, in transport models incl. hybrid.**
E.g. J. Steinheimer, J. Auvinen, H. Petersen, M. Bleicher, H. Stoecker, PRC89('14)054913.
Exception: *B.A.Li, C.M.Ko ('98) with FOPT EOS*

*We investigate the directed flow at BES energies
in hadronic transport model
with / without mean field effects
with / without softening effects via attractive orbit.*

Transport Model

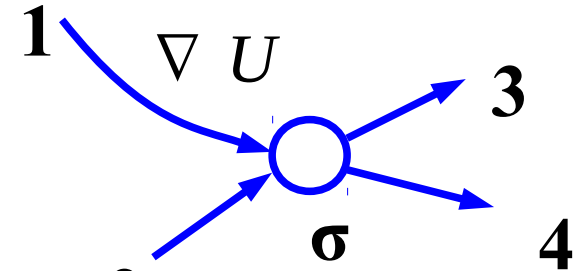
■ Boltzmann equation with (optional) potential effects

E.g. Bertsch, Das Gupta, Phys. Rept. 160(88), 190

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \nabla f - \nabla U \cdot \nabla_p f = I_{\text{coll}}$$

$$I_{\text{coll}}(\mathbf{r}, \mathbf{p}) = -\frac{1}{2} \int \frac{d\mathbf{p}_2}{(2\pi)^3} d\Omega v_{12} \frac{d\sigma}{d\Omega} [f f_2 (1 - f_3)(1 - f_4) - (12 \leftrightarrow 34)]$$

(NN elastic scattering case)



■ Hadron-string transport model JAM

- Collision term → Hadronic cascade with resonance and string excitation

Nara, Otuka, AO, Niita, Chiba, Phys. Rev. C61 (2000), 024901.

- Potential term → Mean field effects in the framework of RQMD/S

Sorge, Stocker, Greiner, Ann. of Phys. 192 (1989), 266.

Tomoyuki Maruyama et al., Prog. Theor. Phys. 96(1996), 263.

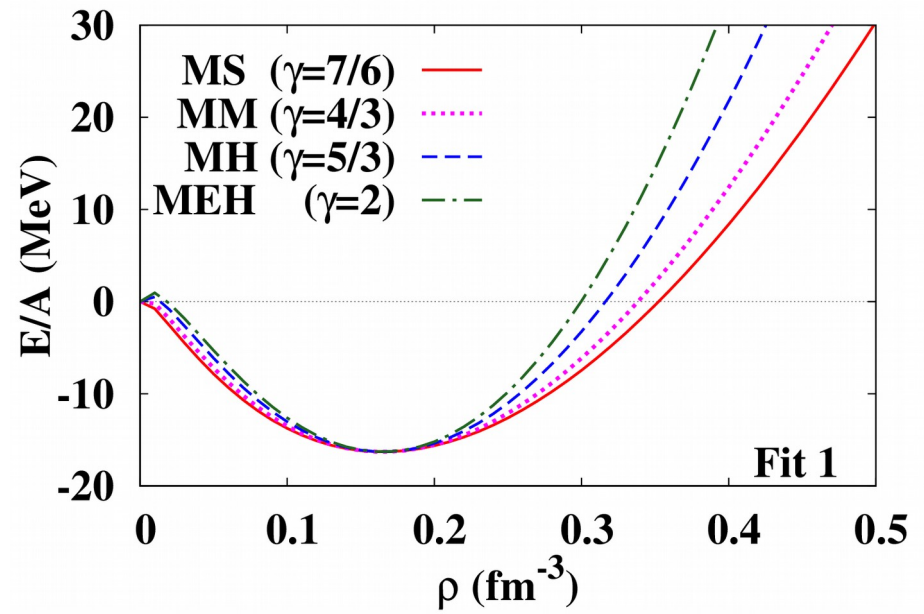
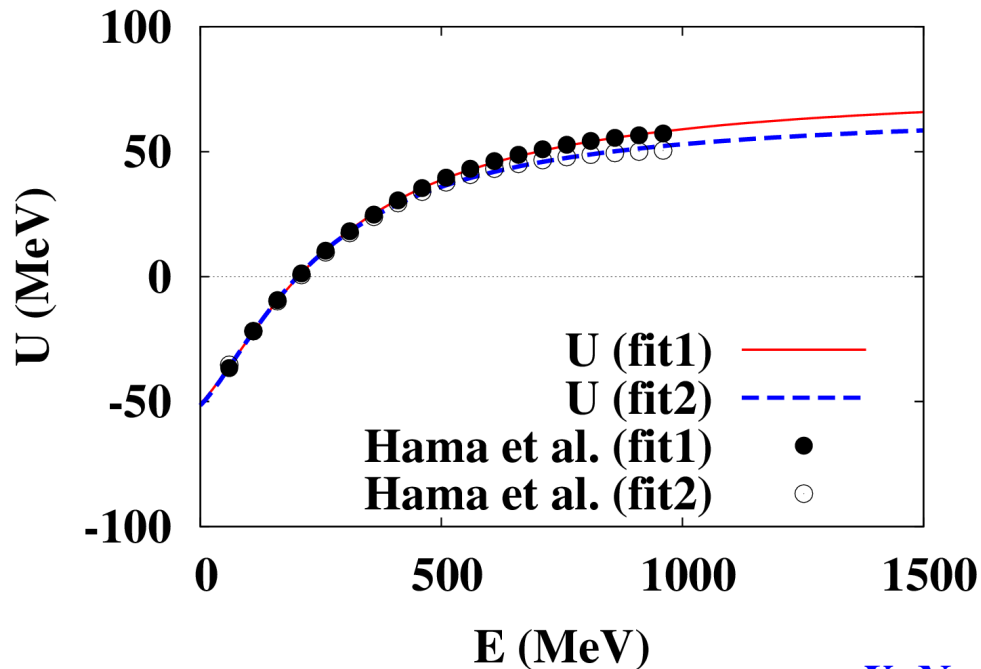
Isse, AO, Otuka, Sahu, Nara, Phys.Rev. C 72 (2005), 064908.

Mean Field Potential

■ Skyrme type density dependent + momentum dependent potential

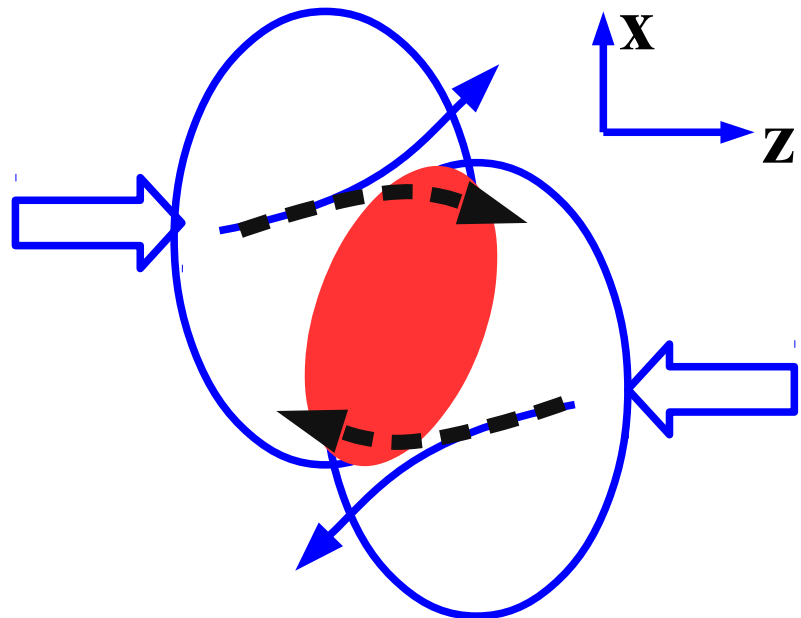
$$V = \sum_i V_i = \int d^3r \left[\frac{\alpha}{2} \left(\frac{\rho}{\rho_0} \right)^2 + \frac{\beta}{\gamma+1} \left(\frac{\rho}{\rho_0} \right)^{\gamma+1} \right] + \sum_k \int d^3r d^3p d^3p' \frac{C_{ex}^{(k)}}{2\rho_0} \frac{f(\mathbf{r}, \mathbf{p})f(\mathbf{r}, \mathbf{p}')}{1 + (\mathbf{p} - \mathbf{p}')^2 / \mu_k^2}$$

Type	α (MeV)	β (MeV)	γ	$C_{ex}^{(1)}$ (MeV)	$C_{ex}^{(2)}$ (MeV)	μ_1 (fm ⁻¹)	μ_2 (fm ⁻¹)	K (MeV)
MH1	-12.25	87.40	5/3	-383.14	337.41	2.02	1.0	371.92
MS1	-208.89	284.04	7/6	-383.14	337.41	2.02	1.0	272.6



*Y. Nara, AO, arXiv:1512.06299 [nucl-th] (QM2015 proc.)
Isse, AO, Otuka, Sahu, Nara, PRC 72 (2005), 064908.*

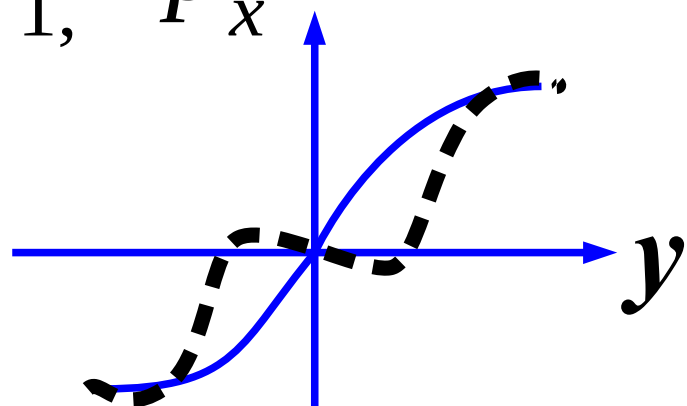
What is directed flow ?



- v_1 or $\langle p_x \rangle$ as a function of y is called directed flow.
- Created in the overlapping stage of two nuclei
→ Sensitive to the EOS in the early stage.
- Becomes smaller at higher energies.

Attraction
(Softening)

$v_1, \langle p_x \rangle$



$$v_1 = \langle p_x / p \rangle = \langle \cos \phi \rangle$$

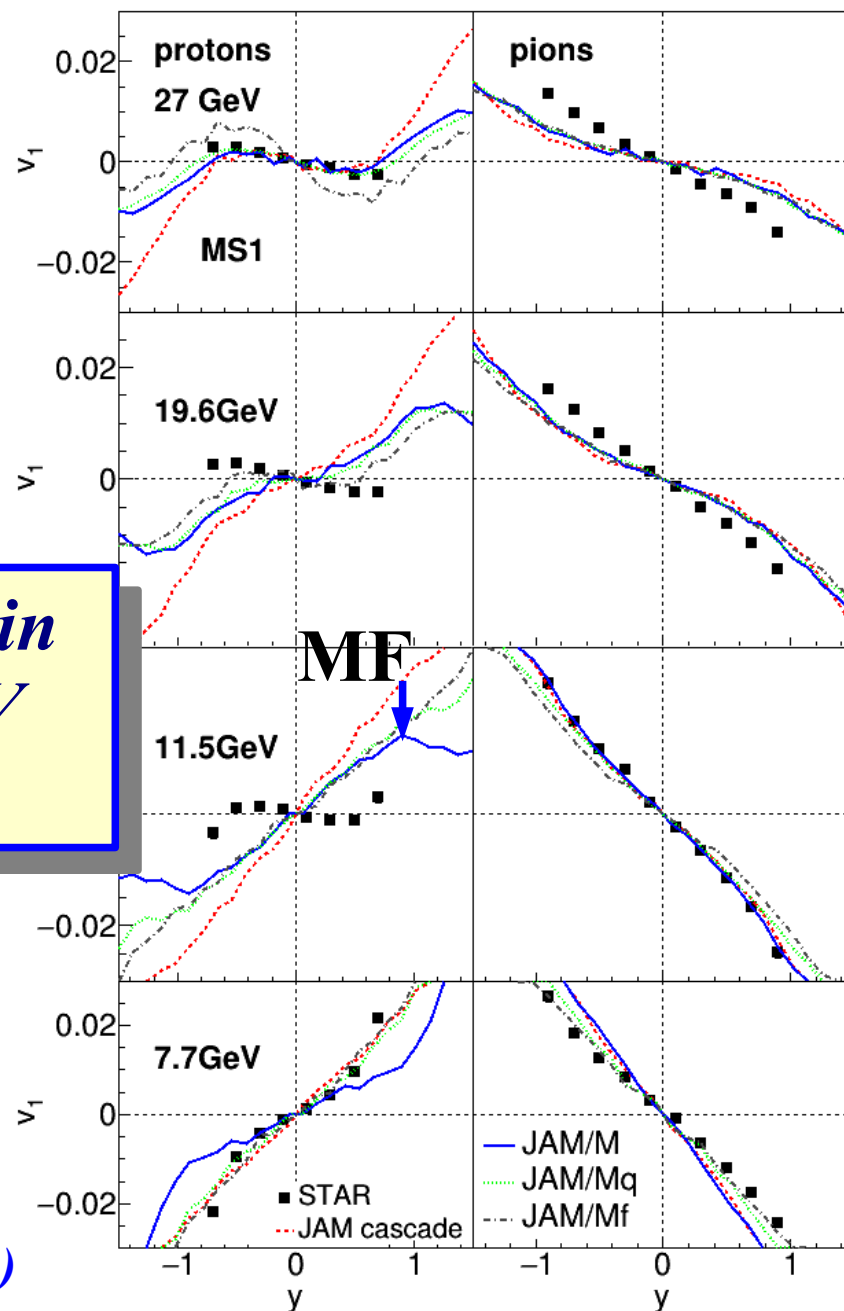
How can we explain non-monotonic dependence of dv_1/dy ?
→ *Softening or Geometry*

Comparison with RHIC data on v_1

- Pot. Eff. on the v_1 is significant, but dv_1/dy becomes negative only at $\sqrt{s_{NN}} > 20$ GeV.

Hadronic approach does not explain directed flow collapse at 10-20 GeV even with potential effects.

- JAM/M: only formed baryons feel potential forces
- JAM/Mq: pre-formed hadron feel potential with factor 2/3 for diquark, and 1/3 for quark
- JAM/Mf: both formed and pre-formed hadrons feel potential forces.



Y. Nara, AO, arXiv:1512.06299 [nucl-th] (QM2015 proc.)

Softening Effects via Attractive Orbit Scattering

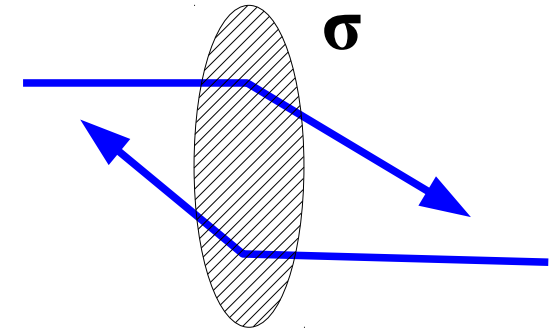
- Attractive orbit scattering simulates softening of EOS

P. Danielewicz, S. Pratt, PRC 53, 249 (1996)

H. Sorge, PRL 82, 2048 (1999).

$$P = P_f + \frac{1}{3TV} \sum_{(i,j)} (\mathbf{q}_i \cdot \mathbf{r}_i + \mathbf{q}_j \cdot \mathbf{r}_j)$$

(Virial theorem)



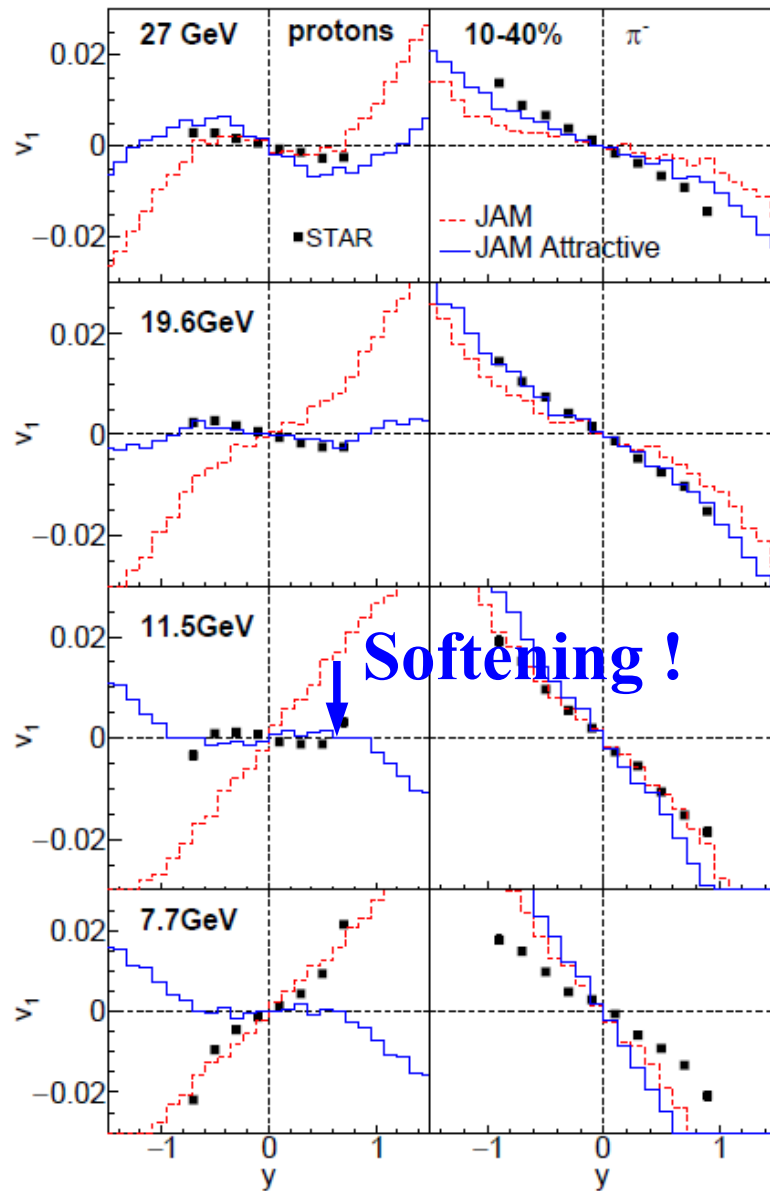
- Attractive orbit → particle trajectory are bended in denser region

*Let us examine the EOS softening effects,
which cannot be explained in hadronic mean field potential,
by using attractive orbit scatterings !*

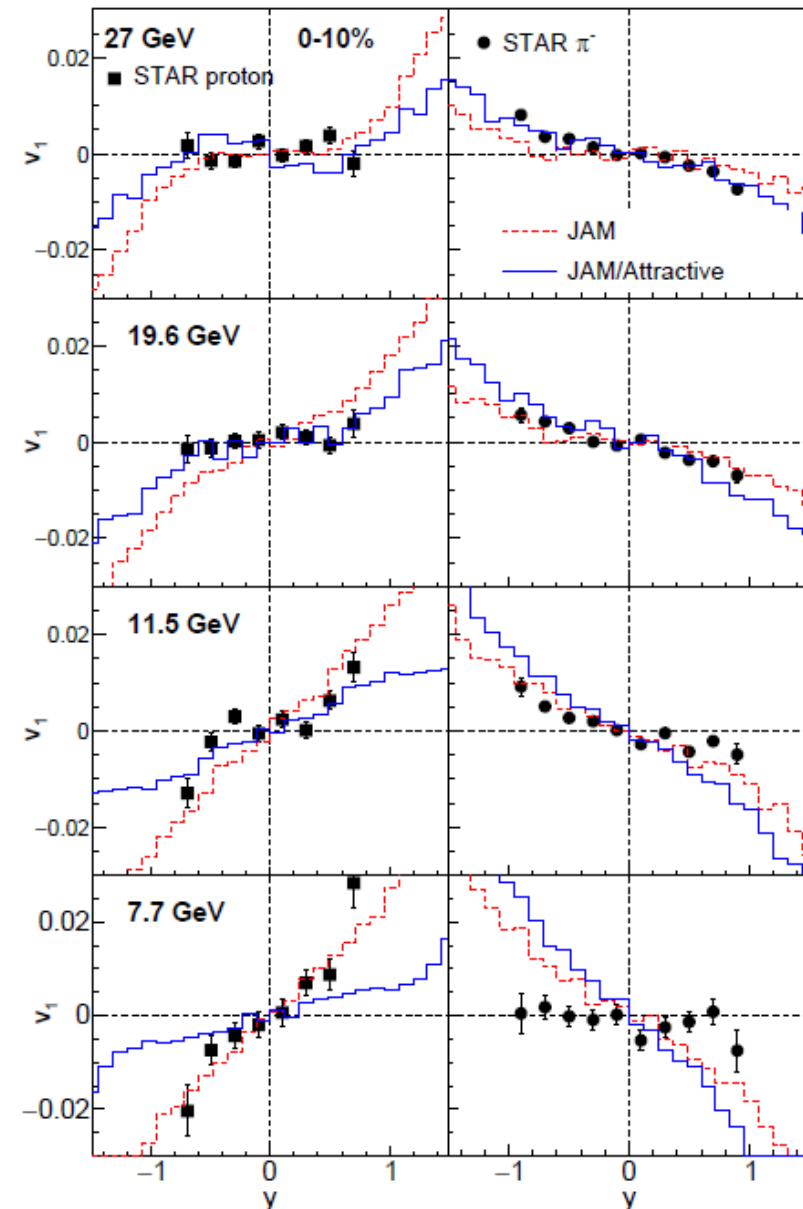
Y. Nara, Niemi, AO, H. Stöcker ('16)

Directed Flow with Attractive Orbits

Nara, Niemi, AO, Stöcker ('16)

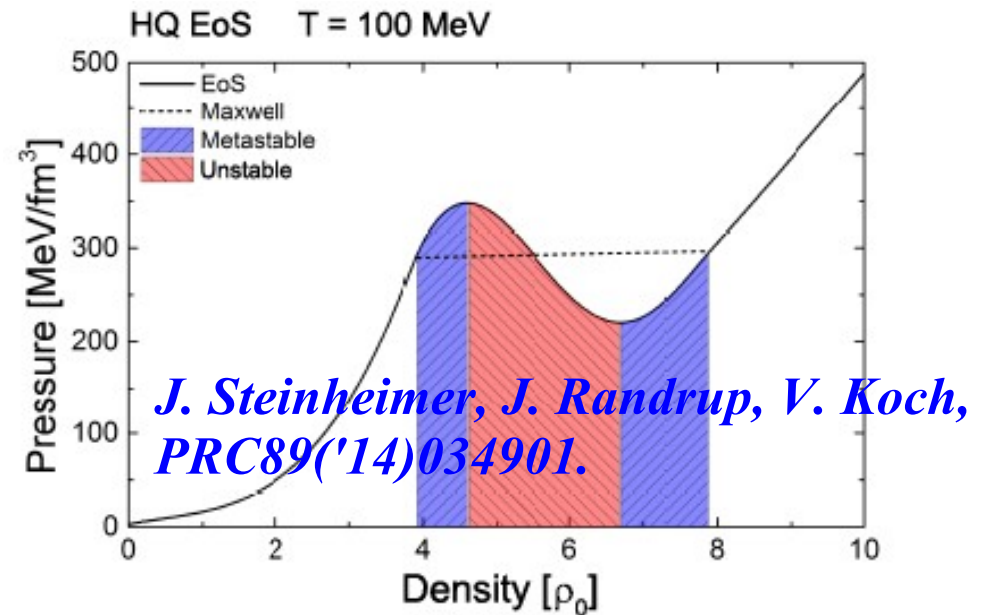
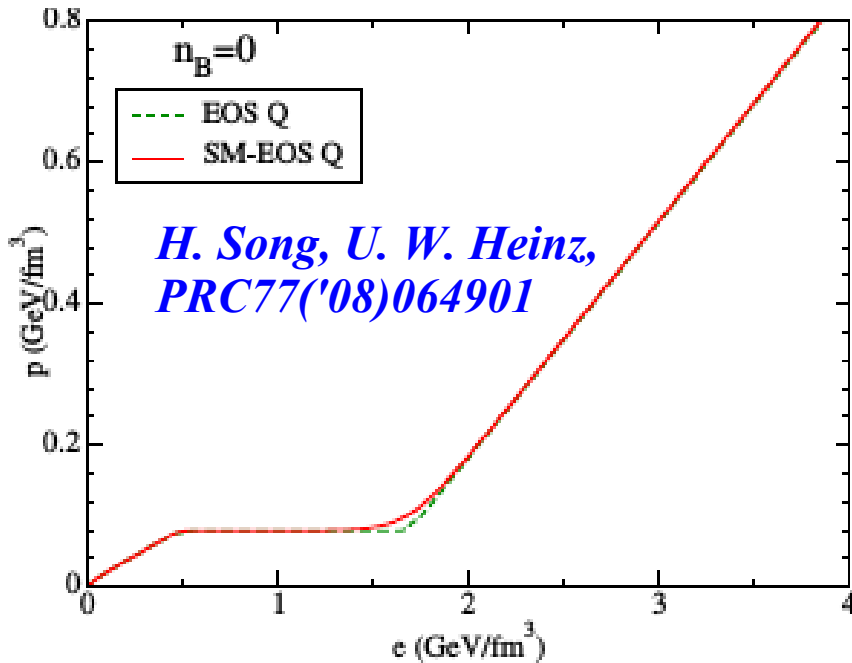
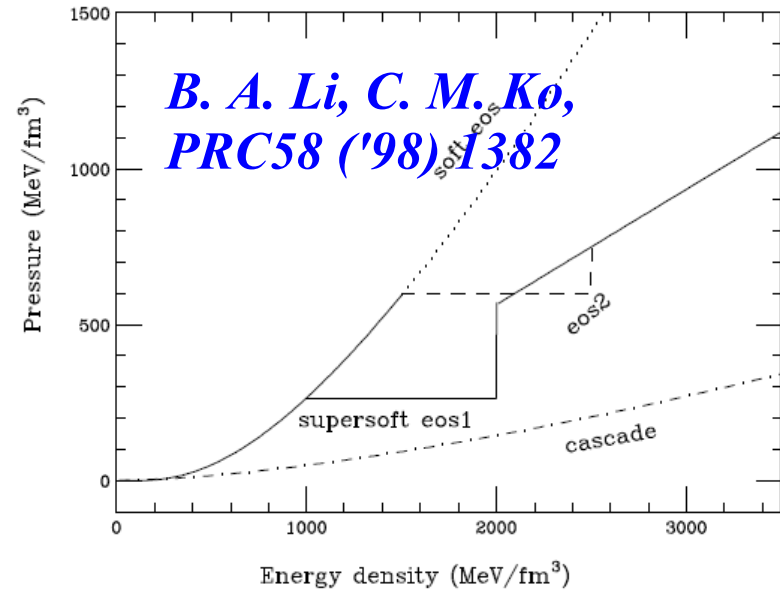
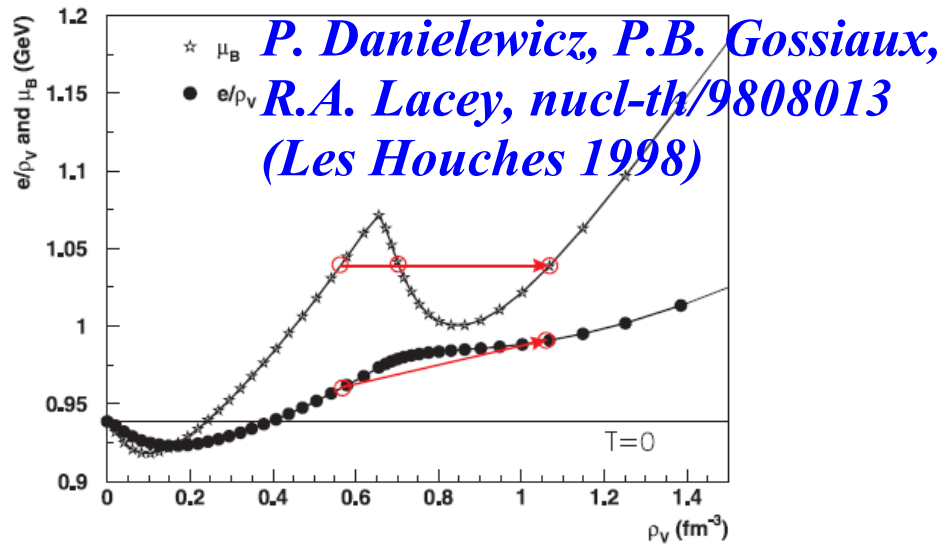


mid-central (10-40 %)



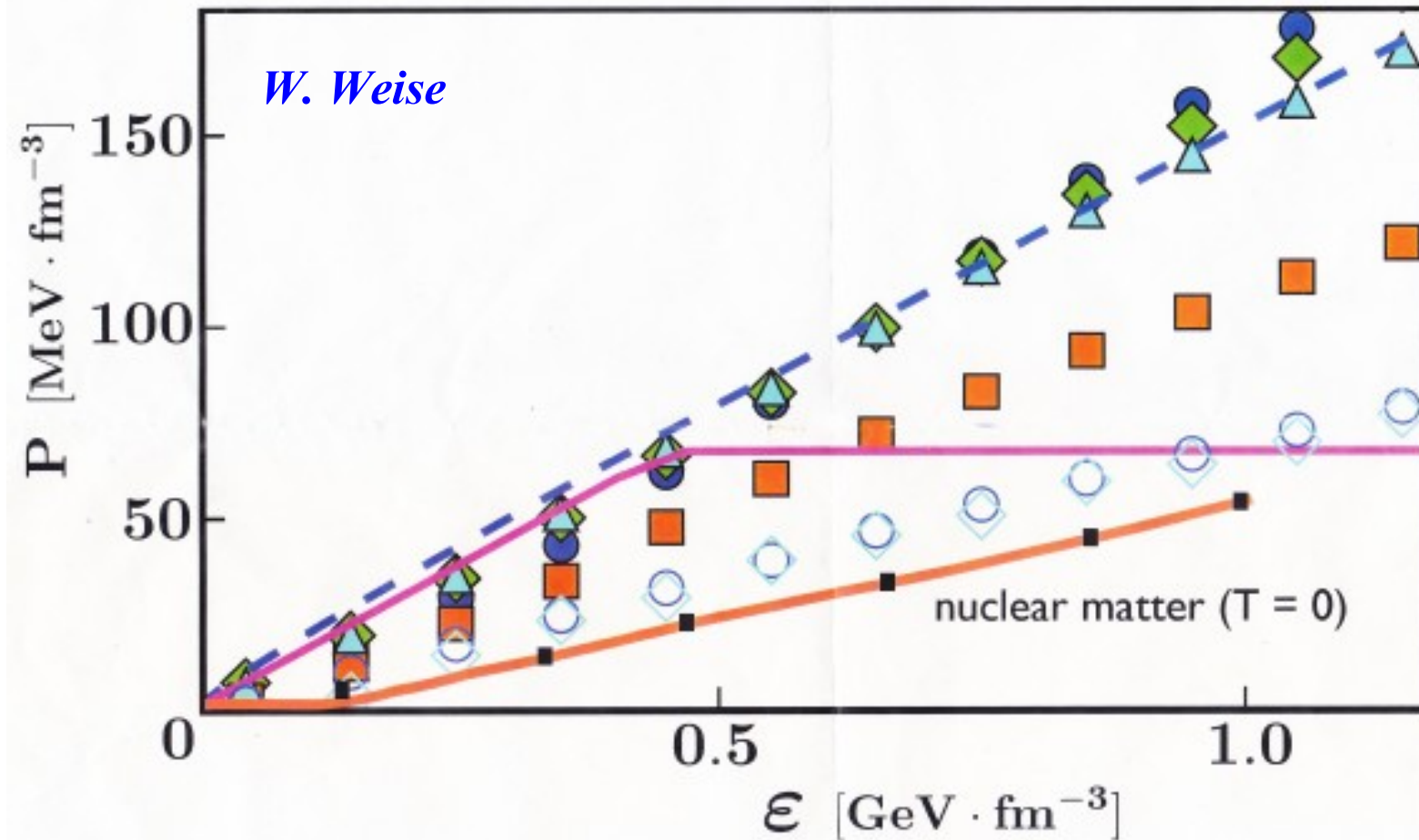
central (0-10 %)

Softening: Where and How much ?



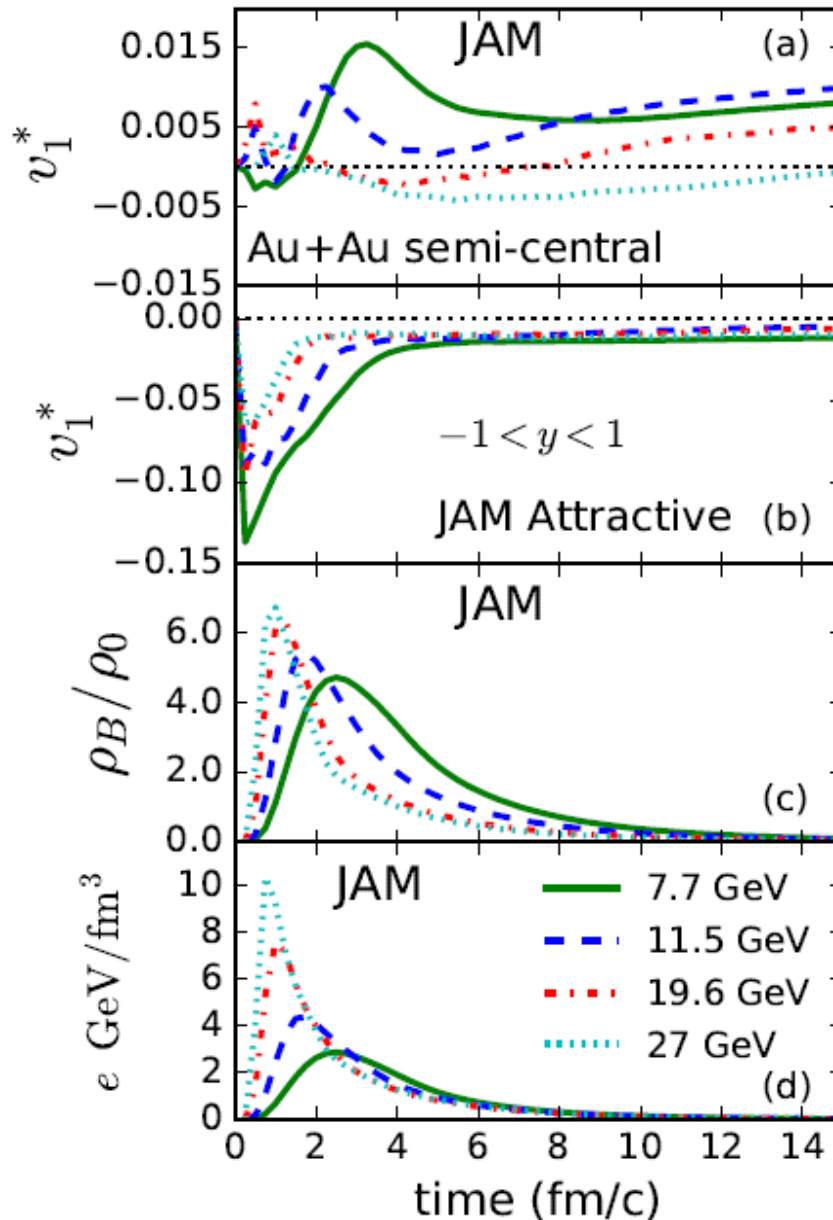
Previous analyses: $\rho_B = (3-10) \rho_0$, $P = (80-700) \text{ MeV/fm}^3$

Comparison with Cold Matter EOS



*FRG EOS does not reach $P \sim 70 \text{ MeV}/\text{fm}^3$ at $\epsilon < 1 \text{ GeV}/\text{fm}^3$
→ Consistent with no FOPT at $\epsilon < 1 \text{ GeV}/\text{fm}^3$*

At which density is the softening required ?



**Softening is
required at
 $\rho > 5\rho_0$**

Short Summary of the 1st part

- We may have seen **QCD phase transition (1st or 2nd) signals at BES (or J-PARC) energies in baryon number cumulants and v_1 slope.**
- **Hadronic transport models cannot explain negative v_1 slope below $\sqrt{s_{NN}} = 20$ GeV.**
 - Geometric mechanism becomes manifest at higher energies.
- **Hadronic transport with EOS softening can describe negative v_1 slope below $\sqrt{s_{NN}} = 20$ GeV.**

Y. Nara, H. Niemi, A. Ohnishi, H. Stoecker, PRC94 ('16), 034906.

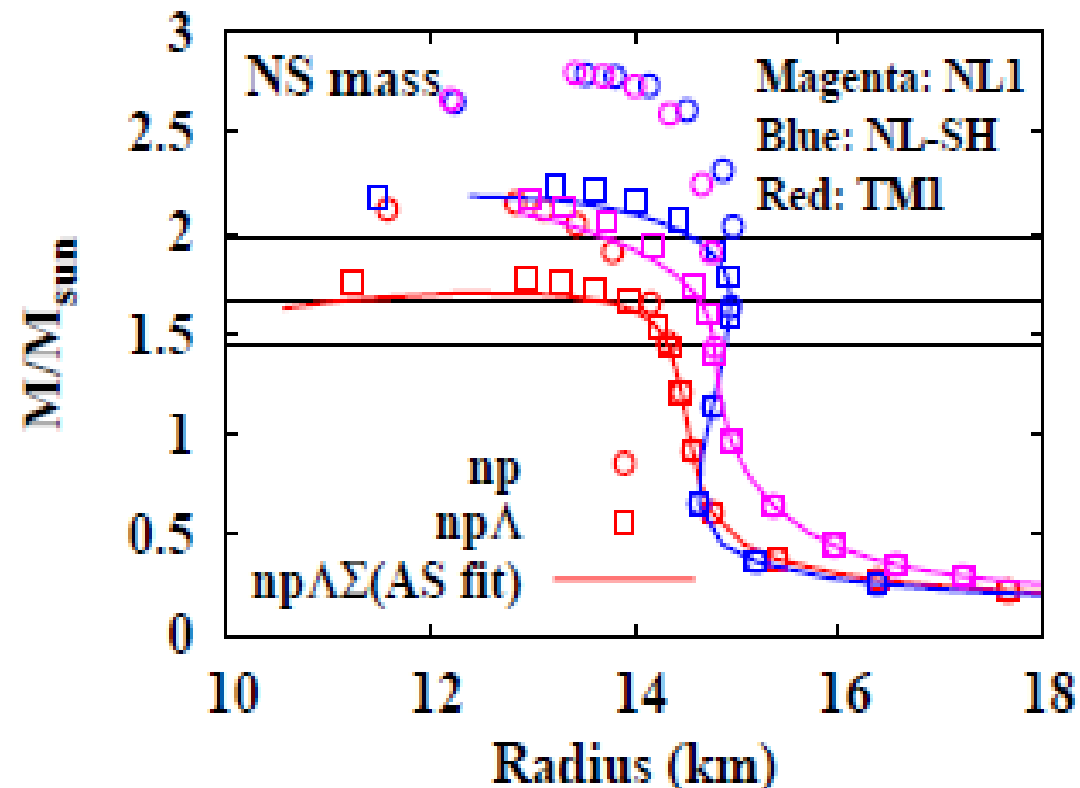
 - **Attractive orbit scattering** simulates EOS softening (virial theorem).
 - We need more studies to confirm its nature.
First-order phase transition ? Crossover ? Forward-backward rapidities ? MF leading to softer EOS ?
- ***We need “re-hardening” at higher energies, e.g. $\sqrt{s_{NN}} = 27$ GeV.***

Dense matter EOS in neutron stars

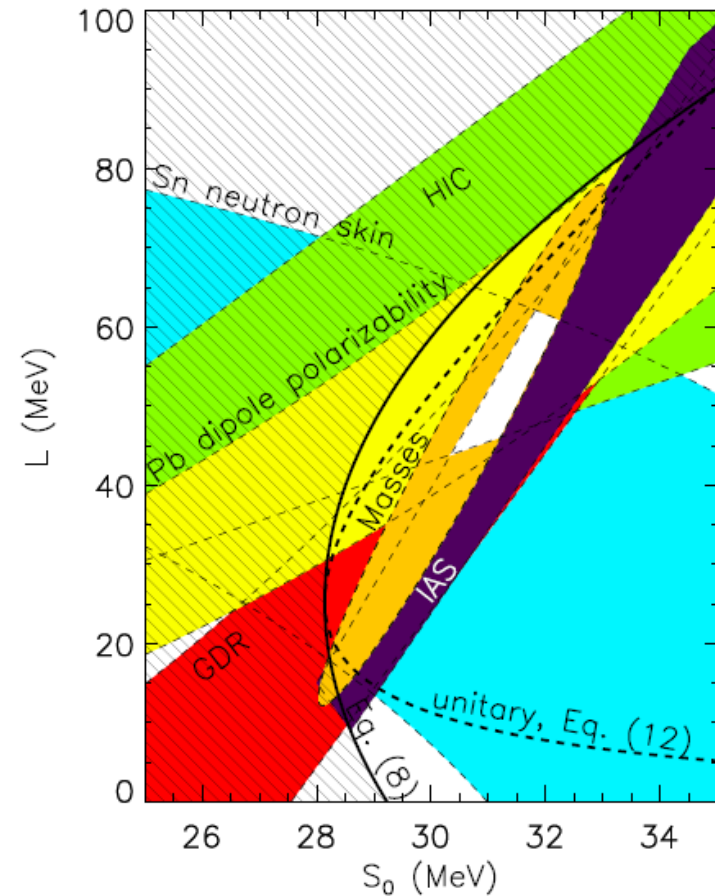
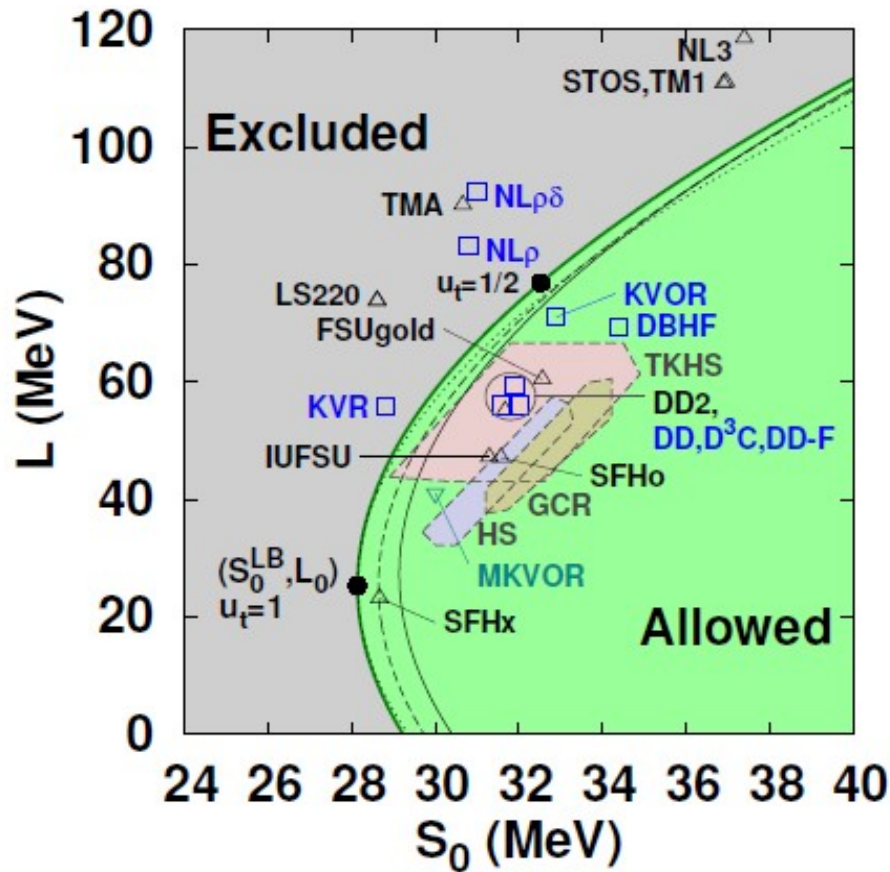
Massive Neutron Stars with Hyperons

Tsubakihara, Harada, AO, arXiv:1402.0979

- Ruled-out EOS with hyperons = GM3
Glendenning & Moszkowski (1991)
- We did NOTHING special and find $2 M_{\odot}$ NS can be supported.
 - “Typical” RMF for nucl. matter
NL1, NL-SH, TM1
Reinhardt et al. ('86); Sharma, Nagarajan, Ring ('93); Sugahara, Toki ('94).
 - $\bar{s}s$ mesons are introduced
 - Hypernuclear data
 Λ , $\Lambda\Lambda$ hypernuclei
 Σ atomic shifts
SU(3) relation to isoscalar
-vector couplings



Symmetry Energy Constraints



Many of EOSs in active astrophysical use do not satisfy recent symmetry energy constraint or $2 M_{\odot}$ constraint.

→ SFHo, SHFx, DD2

Kolomeitsev, Lattimer, AO, Tews ('16)

Ohnishi @ WCNP2017, Oct. 27, 2017 54

What is necessary ?

- Saturation properties (ρ_0 , E_0 , K)
- Symmetry energy parameters (S_0 , L)
- Finite nuclear properties (mass, radius)
- Hypernuclear separation energies (S_Λ)
- Support $2 M_\odot$ neutron stars
- (Neutron star radius at $1.4 M_\odot$ of 12 ± 1 km)
- Hopefully based on microscopic calculations and/or QCD

Relativistic mean field model with multi-body couplings

Thank you !

I review heavy-ion and dense matter physics developed in Hokkaido University.

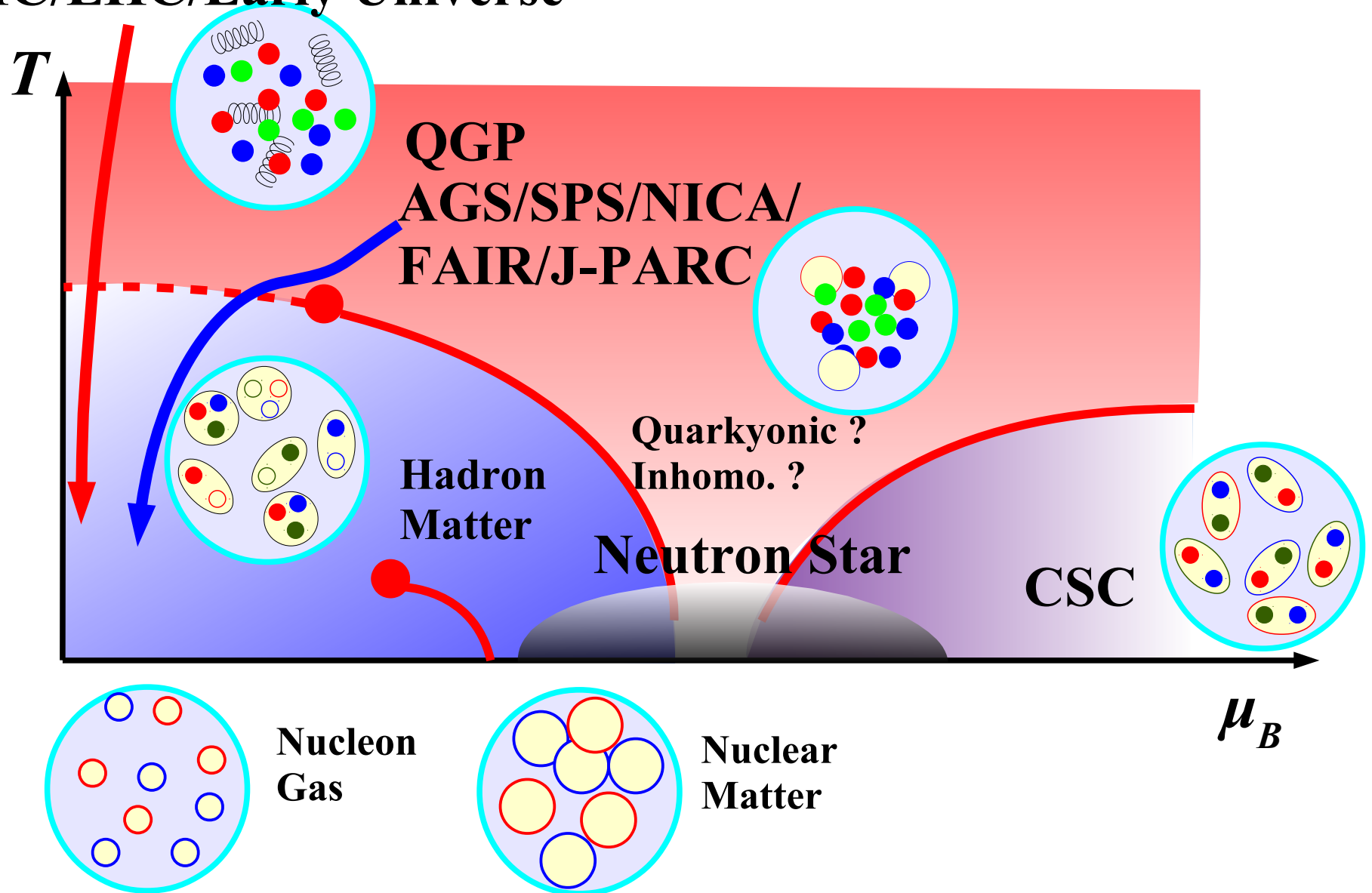
Heavy-ion collision dynamics at 1-158 A GeV were investigated based on hadronic transport models in Hokkaido University. Now its achievements are utilized to understand the QCD phase transition at high densities in heavy-ion collision experiments performed in FAIR, NICA, J-PARC-HI and the Beam Energy Scan program at RHIC.

Dense matter physics investigated in Hokkaido University is now utilized to solve the hyperon puzzle. Hyperon potentials in nuclear matter were used to predict high density baryonic matter EOS, and it is found that we can support neutron star with mass $M < 1.7 M_{\text{sun}}$. In order to explain massive neutron stars with $M \sim 2 M_{\text{sun}}$, we need three-baryon interactions, which may be clarified via precise hypernuclear data.

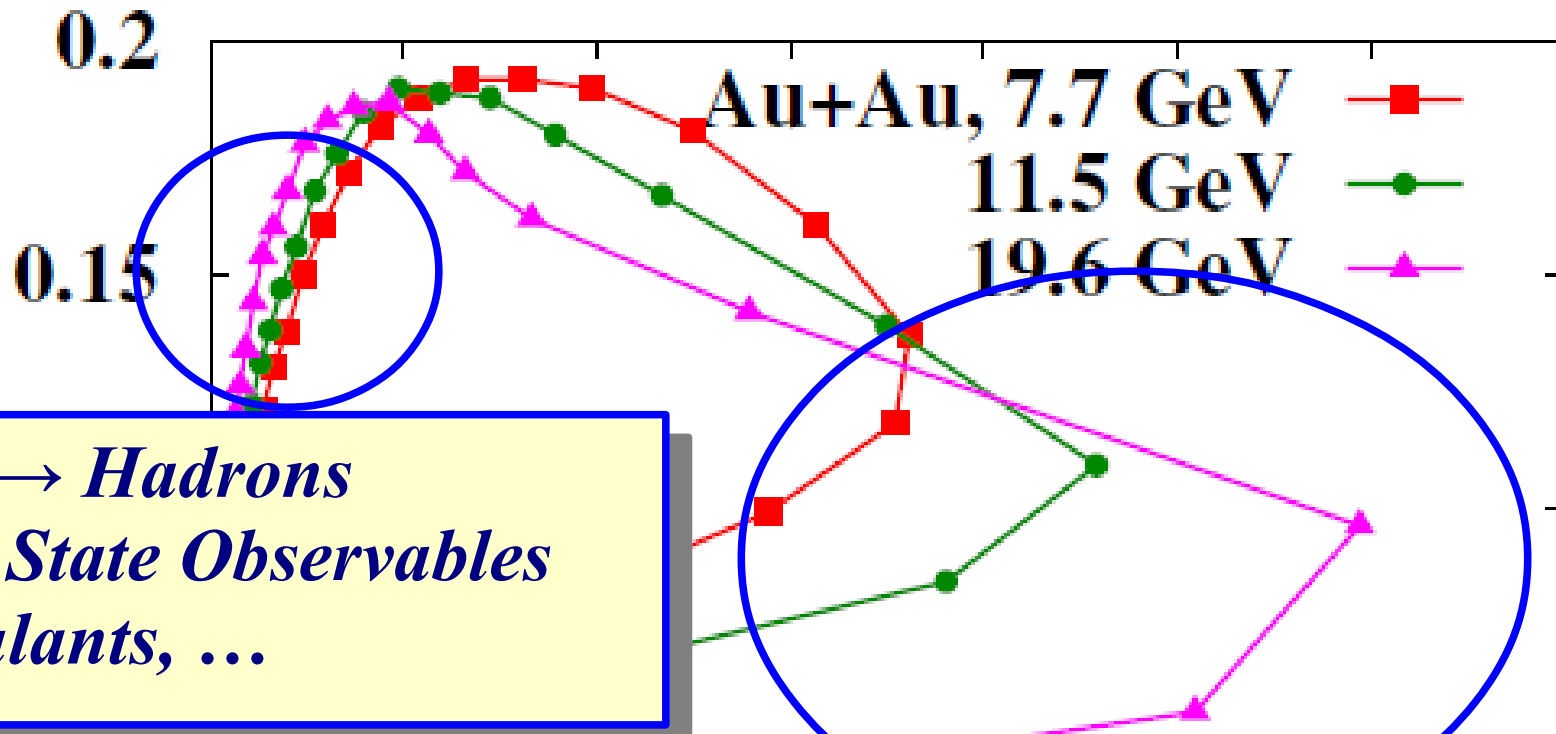
-
- 1993 Nara (PhD 1996.3)
 - 1994 Itagaki (PhD 1999.3), Otsuka (PhD 2001.3)
 - 1995 Uchida (MS 1997.3)
 - 1996 Hirata (PhD 2001.3), [Myo (PhD 2002.3)]
 - 1997 Isshiki (PhD 2005.3?)
 - 1998 Okuda (~1999.3)
 - 1999 [Fukuzaki (B 2000.3)]
 - 2000 Ishizuka (PhD 2005.3)
 - 2001 Isse (PhD 2006.3)
 - 2002 Yamaguchi (MS 2004.3)
 - 2003 Maekawa (PhD 2008.3)
 - 2004 Tsubakihara (PhD 2009.3), Ohnuma (MS 2006.3)

QCD Phase Diagram

RHIC/LHC/Early Universe

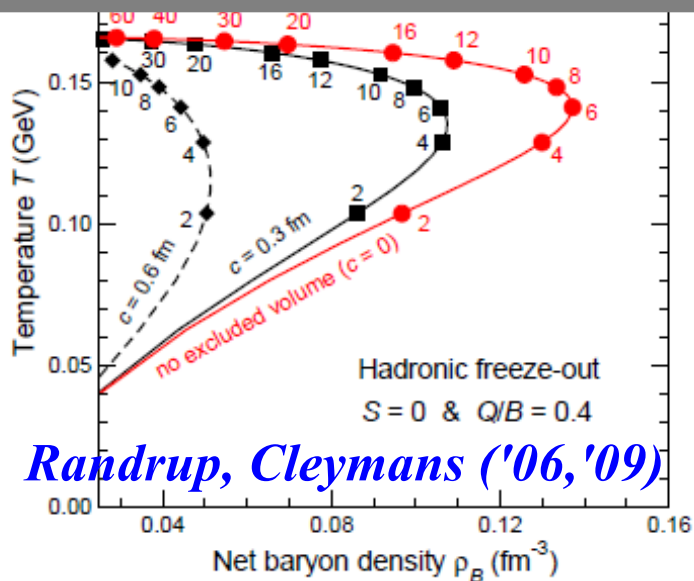


Two ways to probe QCD phase transition



QGP → Hadrons
Final State Observables
Cumulants, ...

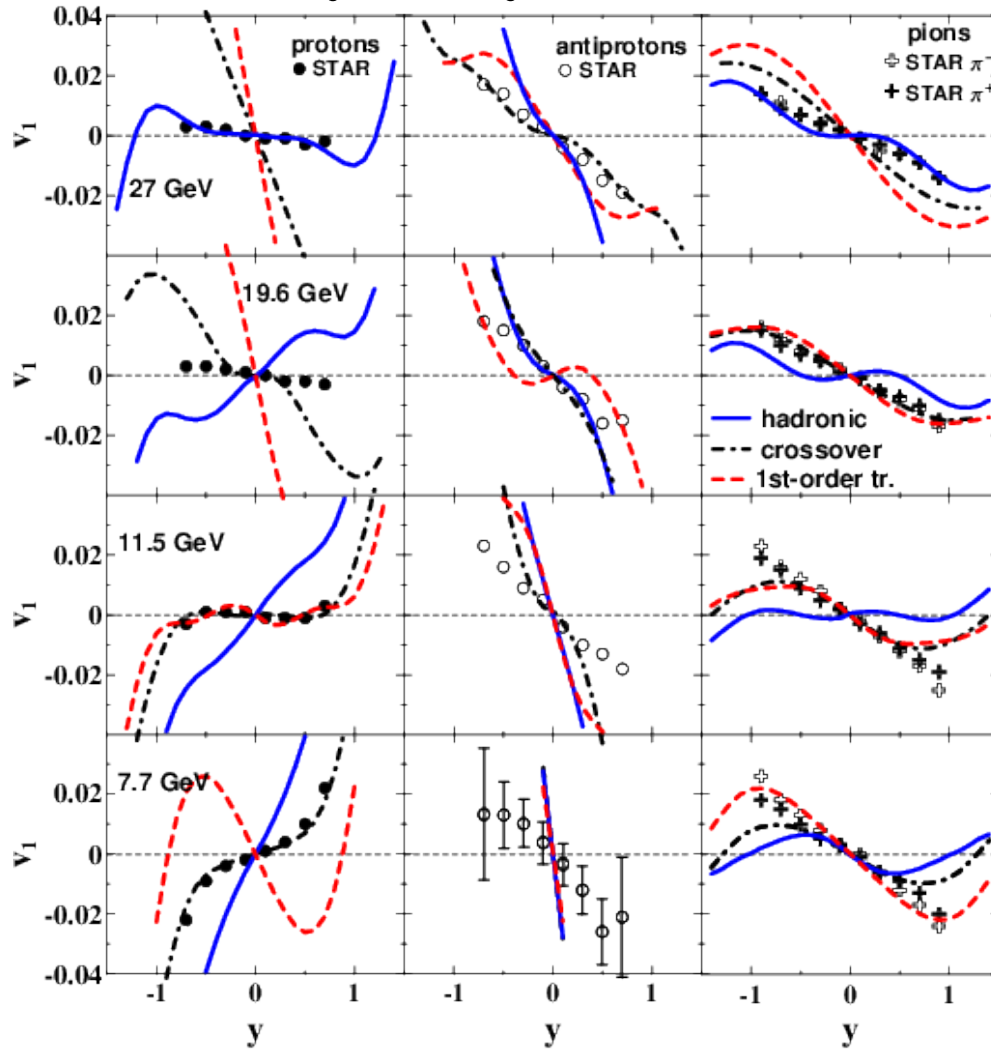
Hadrons → QGP
Early Stage Observables
Caution: (Partial) Equilibration
is necessary !



4

Negative dv_1/dy around $\sqrt{s}_{NN} \sim 10$ GeV

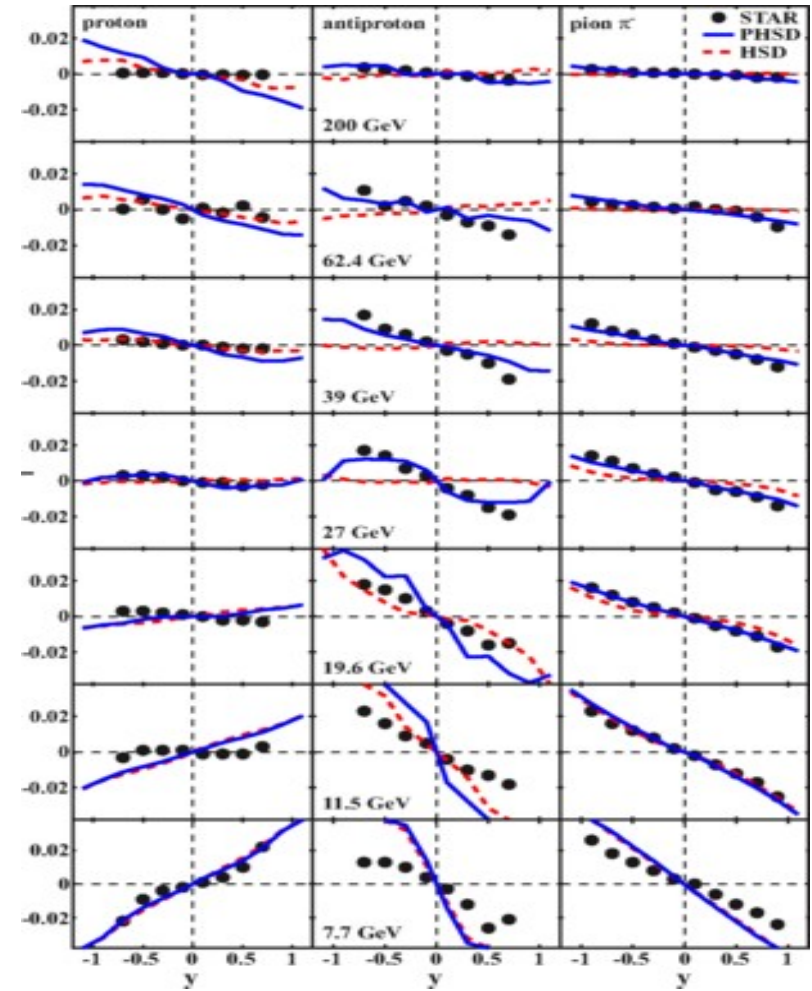
Yes in Hydrodynamics



Black: Crossover, Red: 1st

Y. B. Ivanov and A. A. Soldatov,
PRC91 (2015)024915

No at around $\sqrt{s}_{NN} \sim 10$ GeV in transport models.



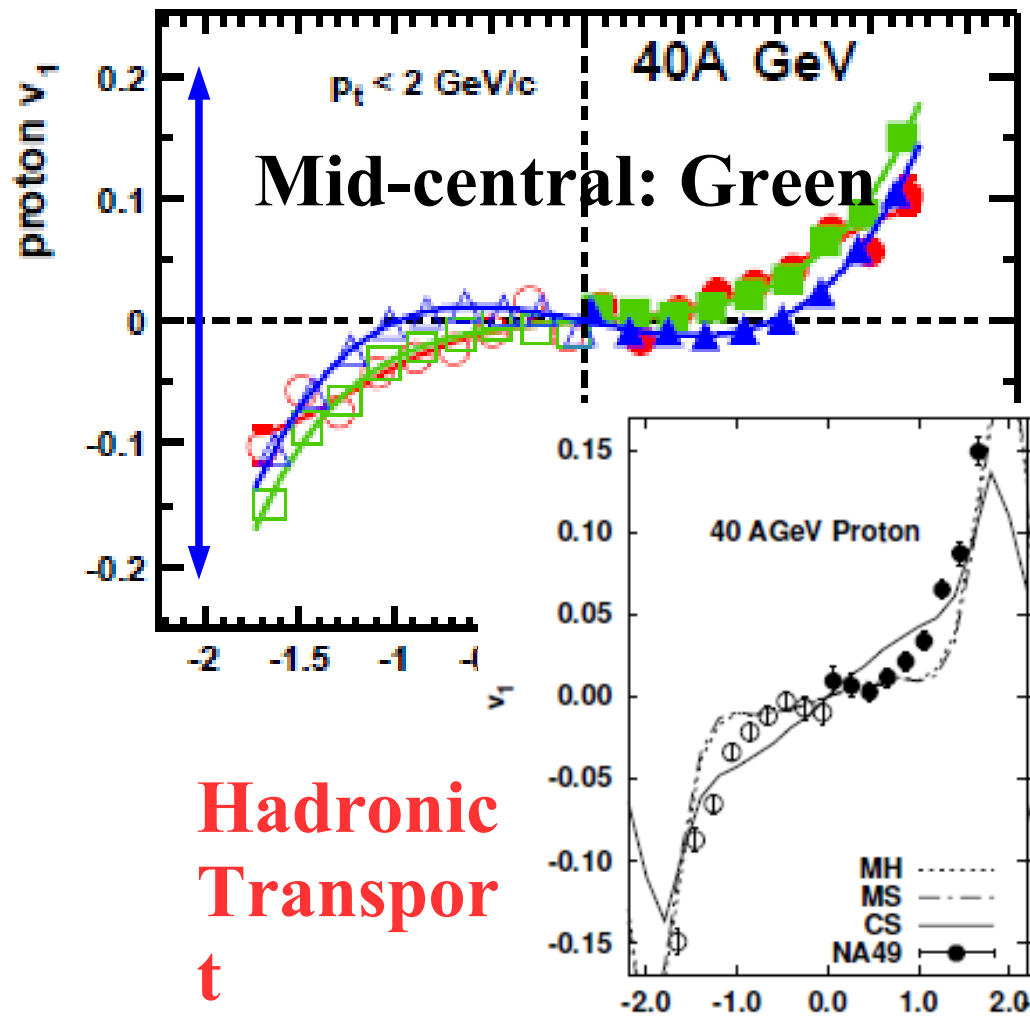
V. P. Konchakovski, W. Cassing, Y. B. Ivanov,
V. D. Toneev, PRC90('14)014903

SPS(NA49) vs RHIC(STAR)

■ SPS (NA49), $\sqrt{s_{NN}} = 8.9$ GeV

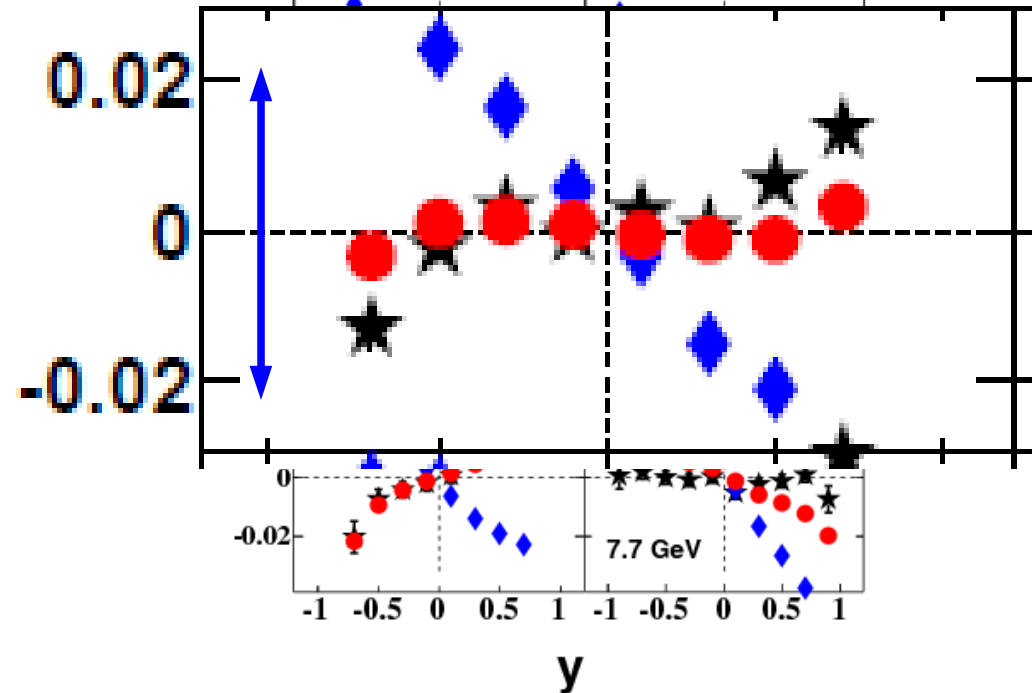
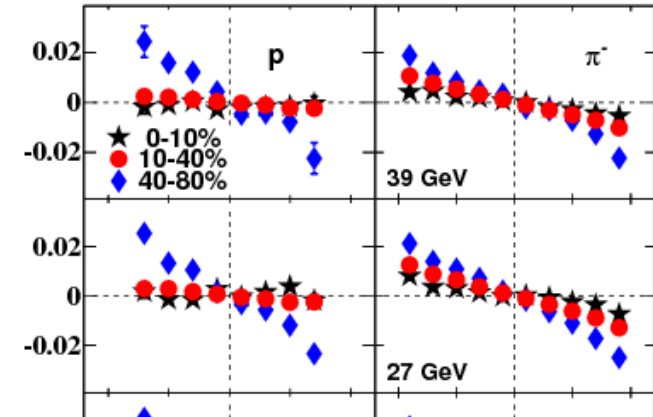
■ RHIC(STAR), 7.7-39 GeV

C. Alt et al. (NA49), PRC68 ('03) 034903



Hadronic
Transpor
t

M. Ise, A. O. N. Otuka, P. K. Sahu, Y. Nara,
PRC72 ('05) 064908



L. Adamczyk et al. (STAR),
PRL 112(2014)162301

- **Introduction**
 - **Two ways to probe QCD phase transition**
 - **Collapse of Directed Flow at $\sqrt{s_{\text{NN}}} \sim 10$ GeV**
- **Hadronic Transport Model Approaches**
 - **Boltzmann equation with potential effects**
 - **Jet AA Microscopic transport model (JAM)**
- **Additional Softening Effects**
 - **Attractive Orbit Scattering**
 - **Transition Density and Pressure (conjecture)**
- **Summary**

Hadronic Transport Models

■ Transport models: Boltzmann + (optional) Potential Effects

$$\frac{\partial f}{\partial t} + \vec{v} \cdot \vec{\nabla} f - \vec{\nabla} U \cdot \vec{\nabla}_p f = C[f]$$

$$C[f] = \int \frac{d\vec{p}_2 d\vec{p}_3 d\vec{p}_4}{(2\pi)^3} v_{12} \sigma_{12 \rightarrow 34} \delta^4(p_1 + p_2 - p_3 - p_4) \\ \times [f_1 f_2 (1 \pm f_3)(1 \pm f_4) - f_3 f_4 (1 \pm f_1)(1 \pm f_2)]$$

■ Commonly used transport models

- **UrQMD 3.4** Frankfurt **public**
- **PHSD** Giessen (Cassing) **upon request**
- **GiBUU 1.6** Giessen (Mosel) **public**
- **AMPT** **public**
- **JAM** (Y. Nara) **public**

Relativistic QMD/Simplified (RQMD/S)

- RQMD is developed based on constraint Hamiltonian dynamics
H. Sorge, H. Stoecker, W. Greiner, Ann. Phys. 192 (1989), 266.

- 8N dof \rightarrow 2N constraints \rightarrow 6N (phase space)
- Constraints = on-mass-shell constraints + time fixation

- RQMD/S uses simplified time-fixation

Tomoyuki Maruyama, et al. Prog. Theor. Phys. 96(1996),263.

- Single particle energy (on-mass-shell constraint)

$$p_i^0 = \sqrt{\mathbf{p}_i^2 + m_i^2 + 2m_i V_i}$$

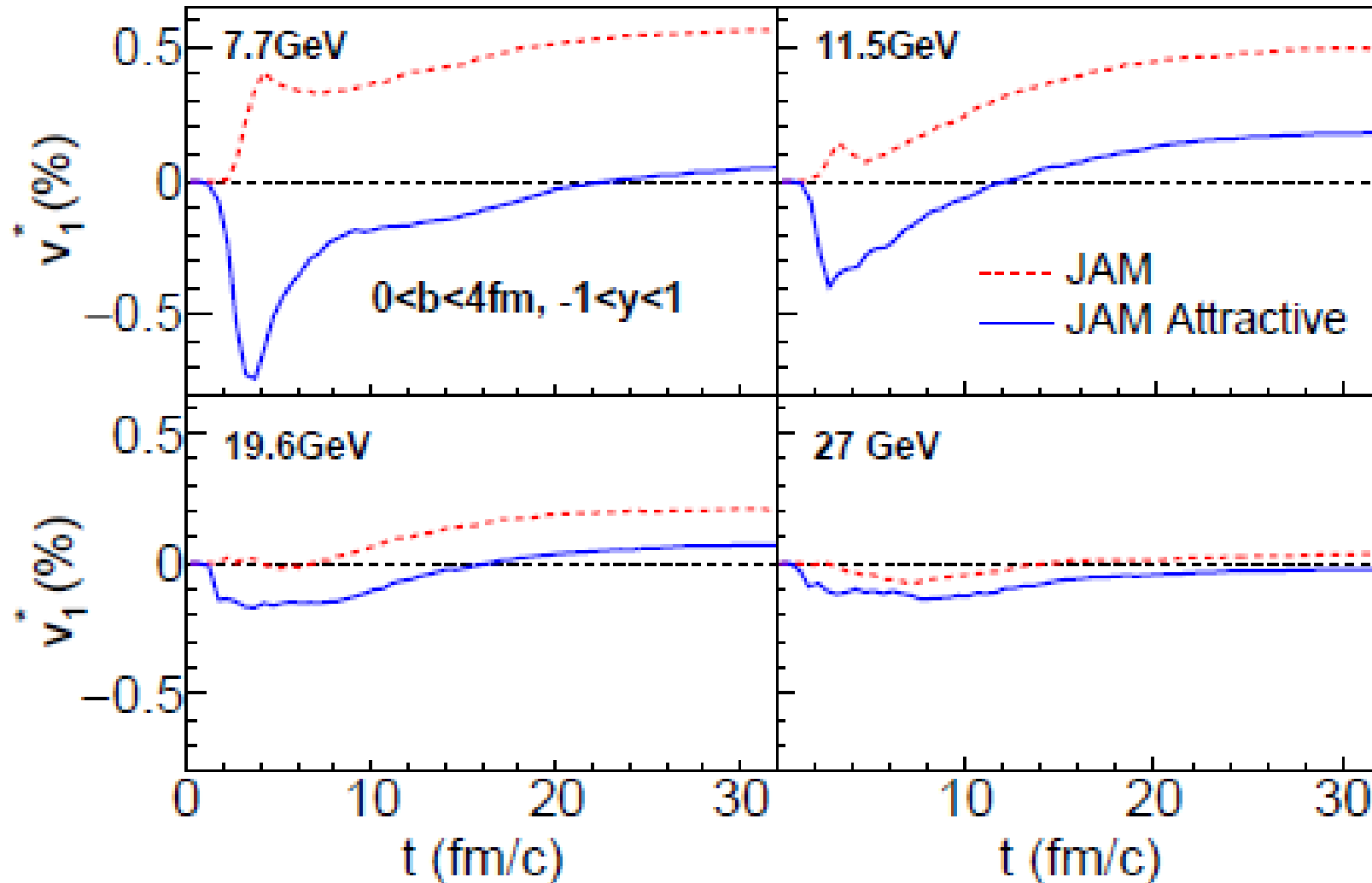
- EOM after solving constraints

$$\dot{\mathbf{r}}_i = \frac{\mathbf{p}_i}{p_i^0} + \sum_j \frac{m_j}{p_j^0} \frac{\partial V_j}{\partial \mathbf{p}_i} \quad \dot{\mathbf{p}}_i = - \sum_j \frac{m_j}{p_j^0} \frac{\partial V_j}{\partial \mathbf{r}_i}$$

- Relative distances $(\mathbf{r}_i - \mathbf{r}_j)^2$ are replaced with those in the two-body c.m.
 \rightarrow Potential becomes Lorentz scalar

When is negative v_1 slope generated ?

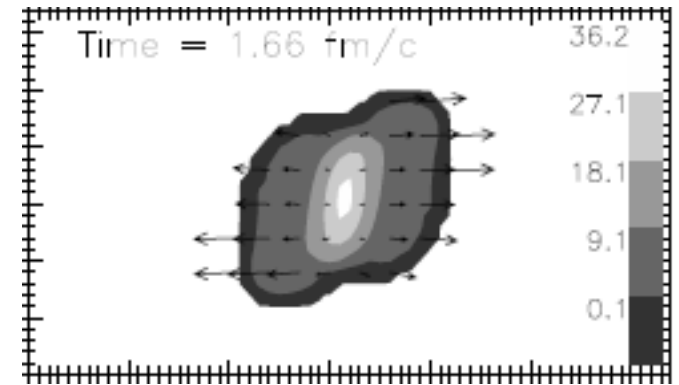
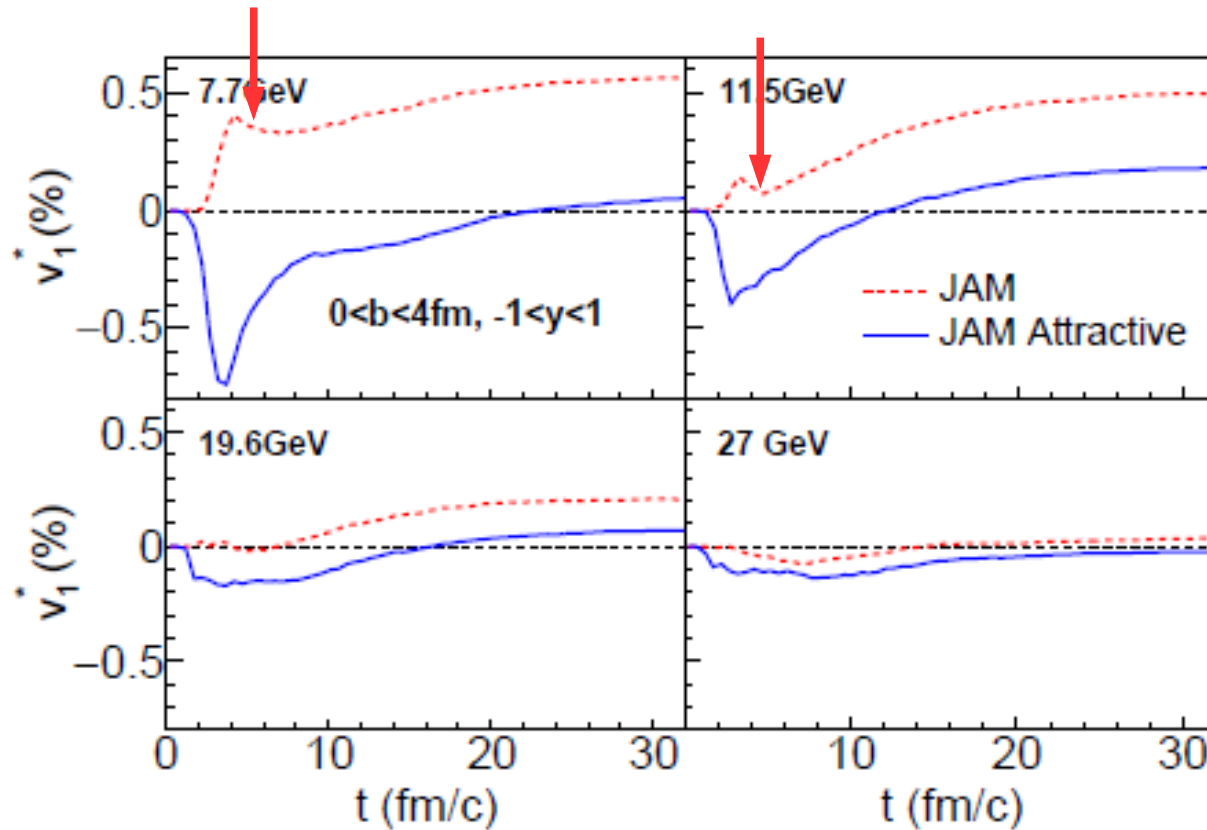
Nara, AO, Stöcker ('16)



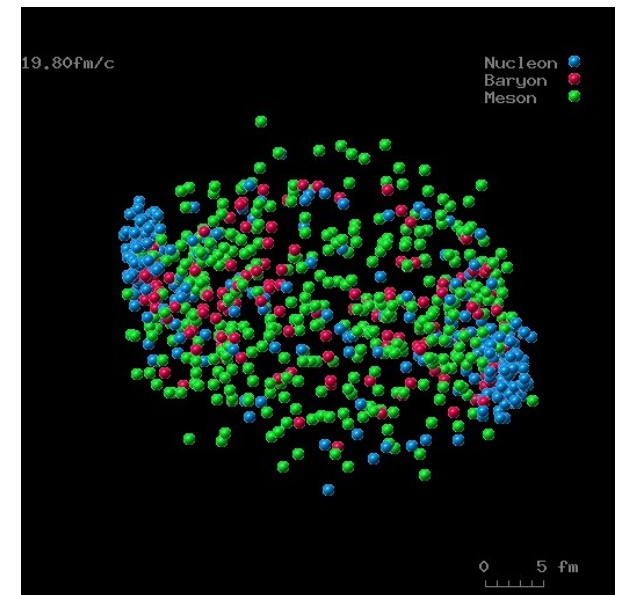
We need to make v_1 slope negative in the compressing stage.

Tilted Ellipsoid ?

Nara, AO, Stöcker ('16)



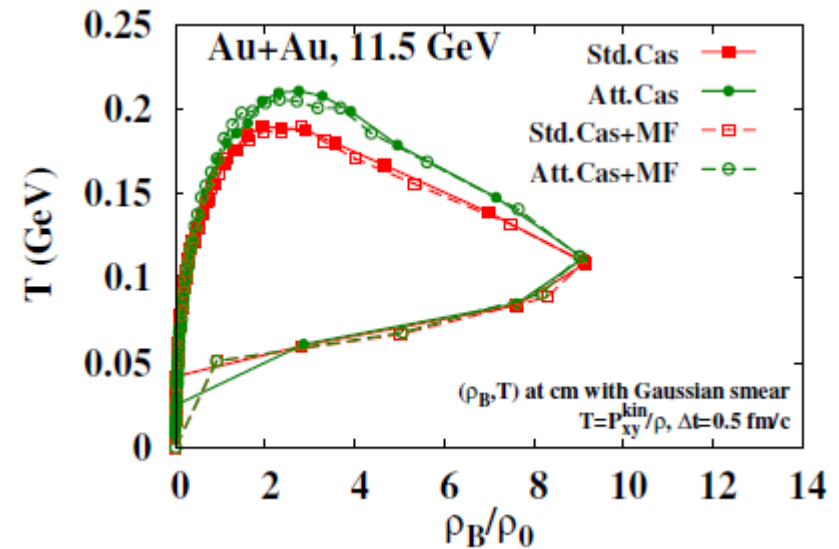
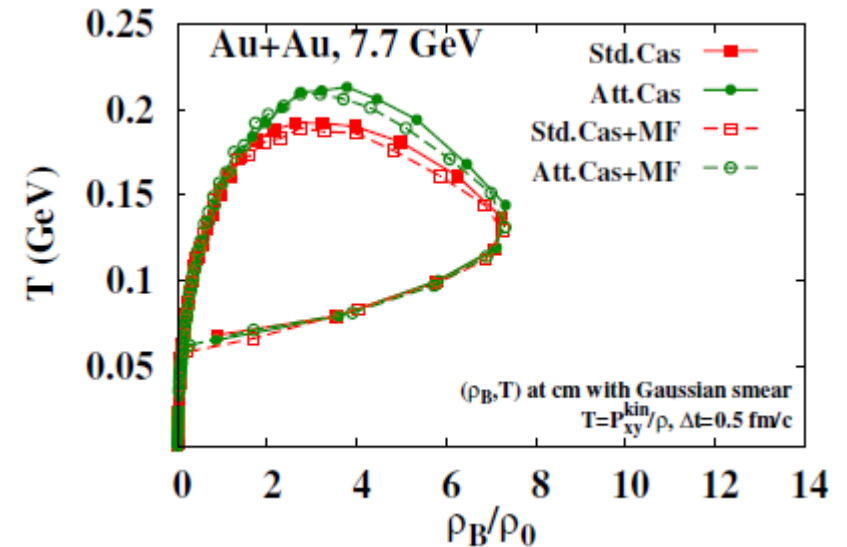
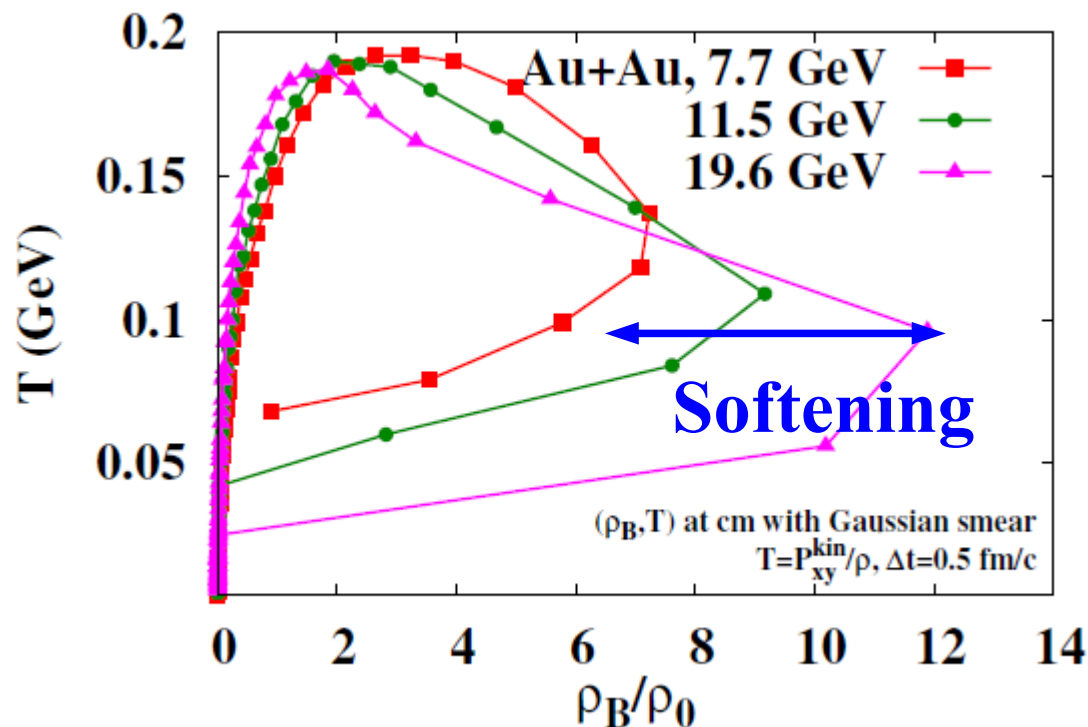
18 GeV, 3-fluid
Toneev et al. ('03)



Transport model results also show tilted-ellipsoid-like behavior, but it is not enough.

Softening of EOS: Where and How much ?

- “Softening” should take place at $\sqrt{s_{NN}}=11.5$ GeV $\rightarrow \rho/\rho_B \sim (6-10)$
- Attractive orbit
 \rightarrow Larger interactions
 & Higher T at later times



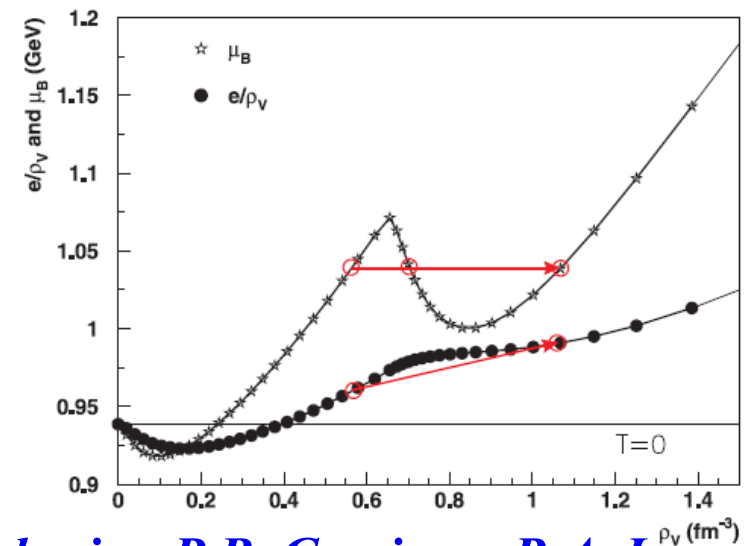
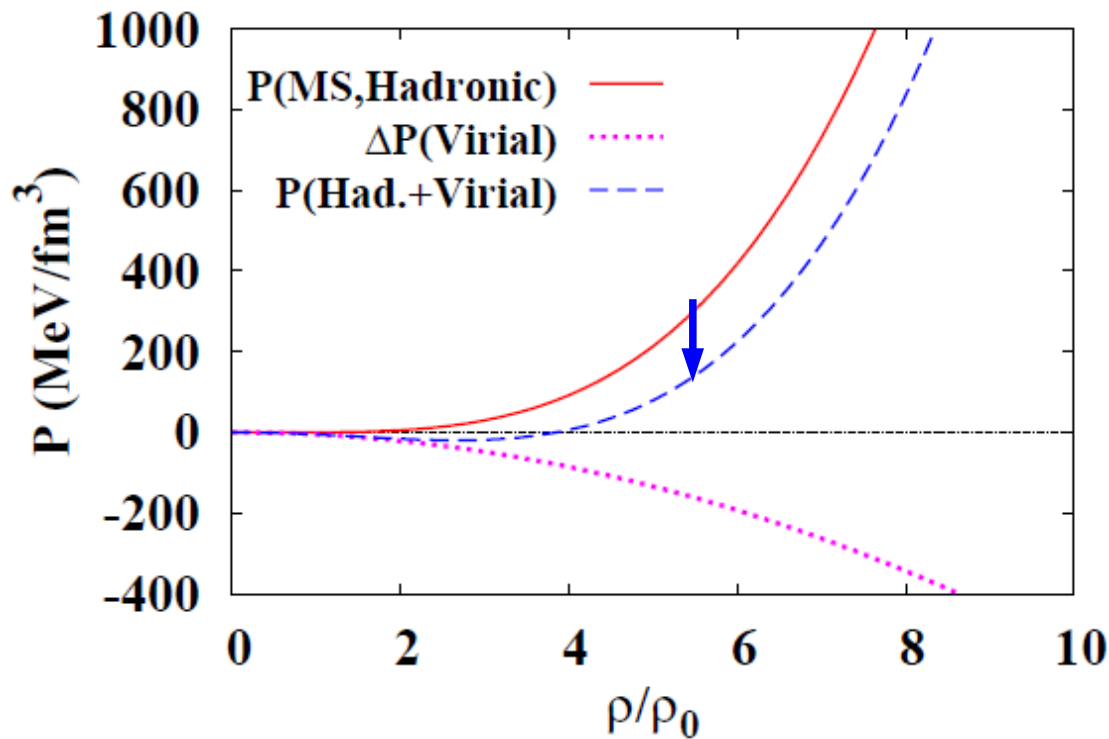
How much softening do we need ?

Virial theorem

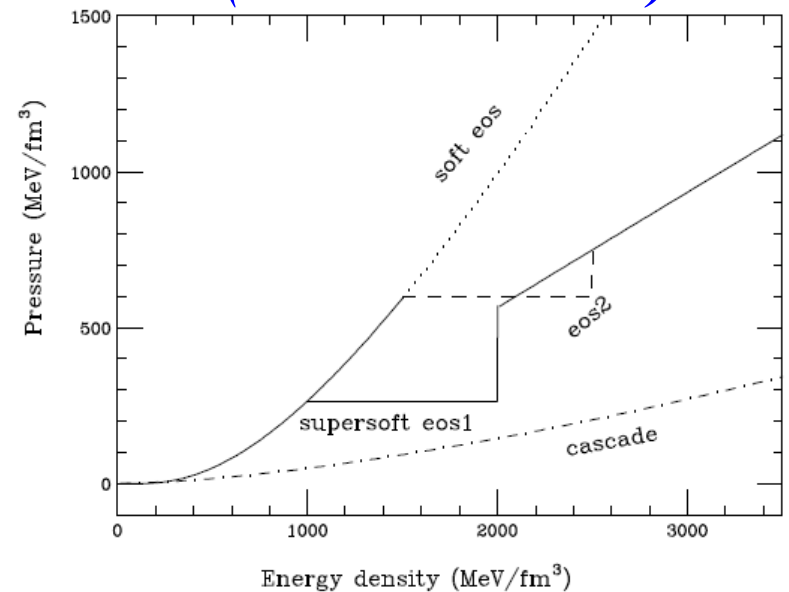
$$\Delta P = \frac{1}{3} \langle v \rho^2 \sigma \mathbf{q} \cdot \Delta \mathbf{r} \rangle$$

Simple estimate:

$$\sigma = 30 \text{ mb}, \langle \mathbf{q} \cdot \Delta \mathbf{R} \rangle \sim -1$$



Danielewicz, P.B. Gossiaux, R.A. Lacey, nucl-th/9808013 (Les Houches 1998)

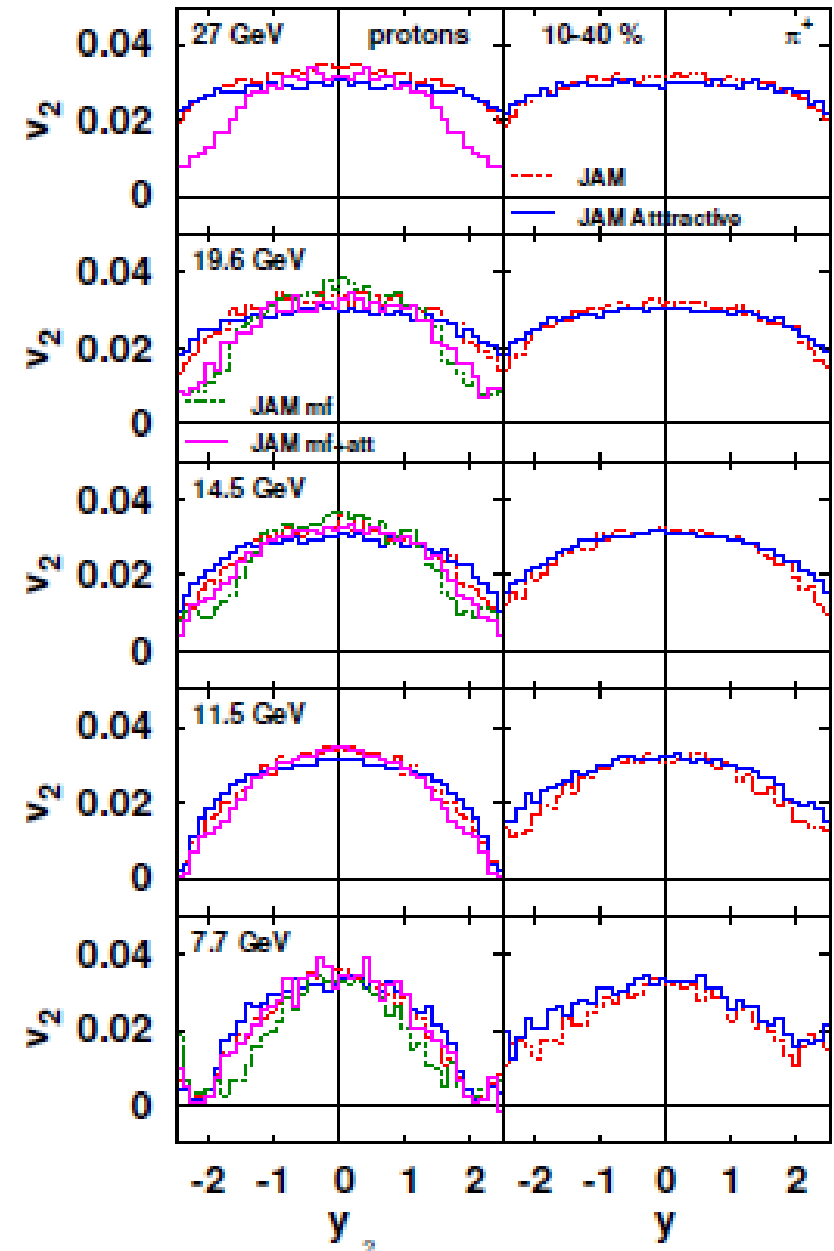


B. A. Li, C. M. Ko, PRC58 ('98) 1382

Ohnishi @ WCNP2017, Oct. 27, 2017 69

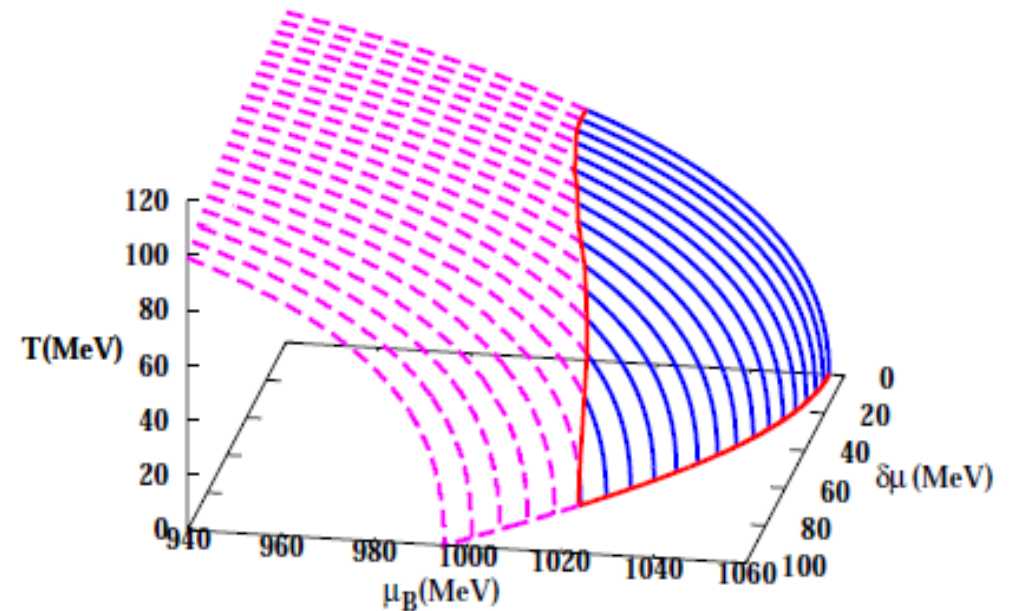
How about v_2 ?

- Do we see softening effects in other observables, e.g. v_2 ?
- Yes, attractive orbits reduces proton v_2 by ~ 0.2 %.
(but there is no qualitative change.)



Relation to Neutron Star Matter

- We may need early transition ($2-5 \rho_0$) to quark matter to solve the hyperon puzzle. Contradicting ?
 - Temperature effects ($T \sim 0 \text{ MeV} \text{ \& } 100 \text{ MeV}$)
 - Isospin chem. pot. (Weaker transition with finite $\delta\mu$)**
 - Hyperon repulsion may push up the transition density.

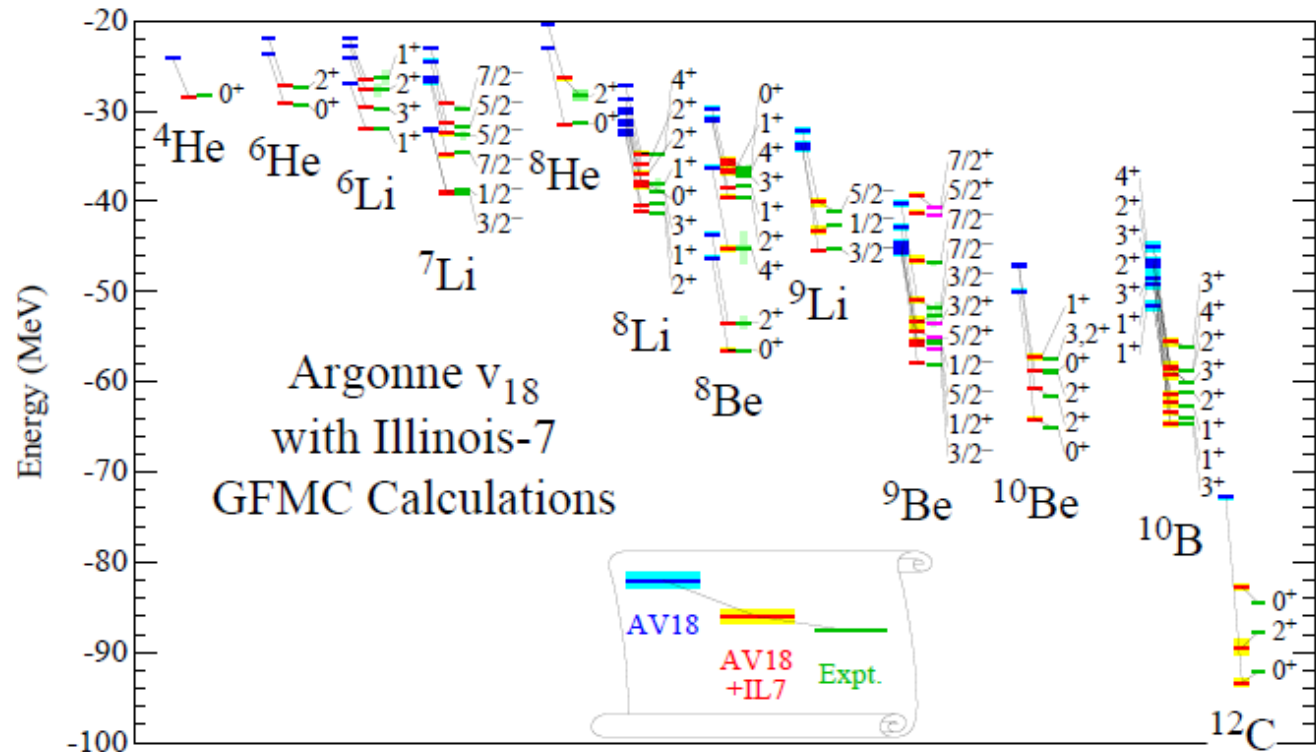


AO, Ueda, Nakano, Ruggieri, Sumiyoshi, PLB704('11),284

H. Ueda, T. Z. Nakano, AO, M. Ruggieri, K. Sumiyoshi, PRD88('13),074006

What is necessary to solve the massive NS puzzle ?

- There are many “model” solutions.
- Ab initio calculation including three-baryon force (3BF)
 - Bare 2NF+Phen. 3NF(UIX, IL2-7) + many-body theory (verified in light nuclei).
 - Chiral EFT (2NF+3NF) + many-body theory
 - Dirac-Bruckner-HF (no 3NF)



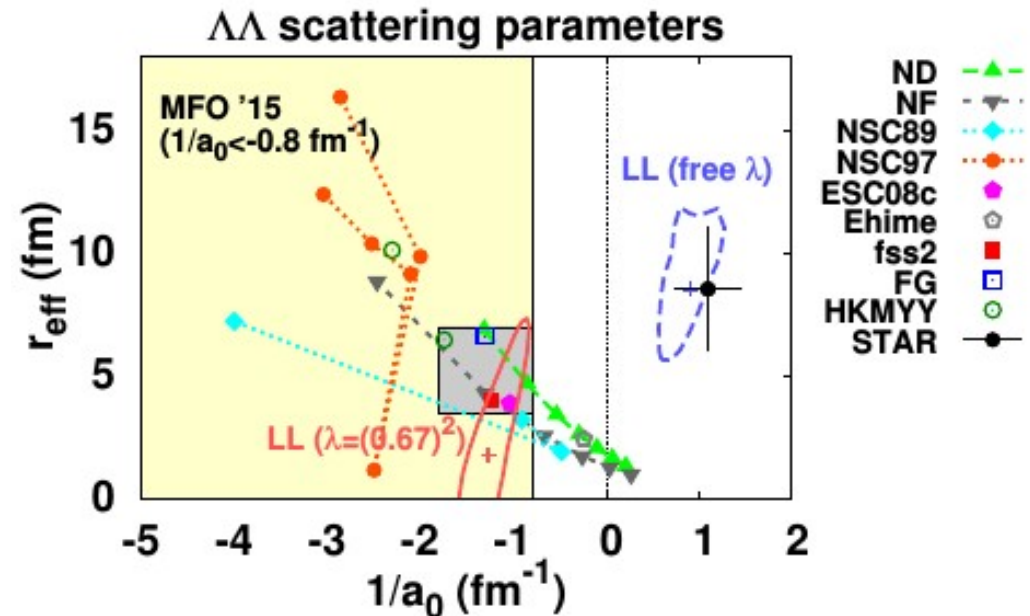
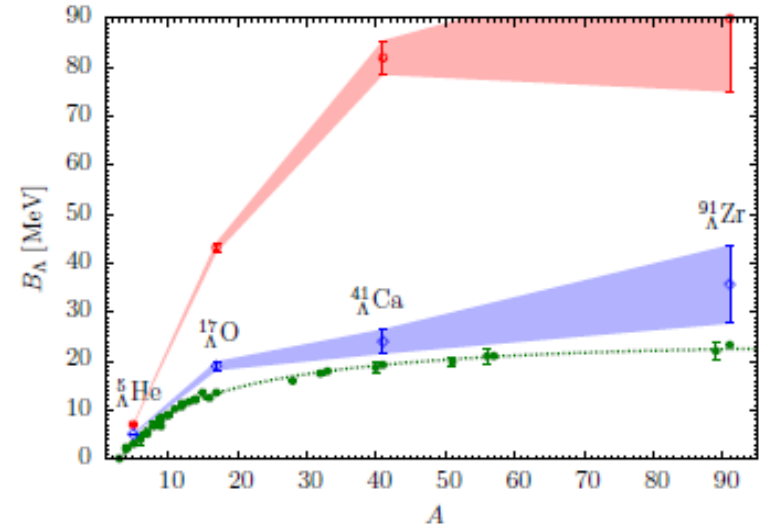
J. Carlson et al. ('14)

3BF including Hyperons

- 3BF incl. YNN, YYN and YYY should exist and contribute to EOS.

Nishizaki, Takatsuka, Yamamoto ('02)

- Chiral EFT, Multi-Pomeron exch., Quark Pauli, Lattice 3BF, SJ, ..
Kohno('10); Heidenbauer+('13); Yamamoto+('14); Nakamoto, Suzuki; Doi+(HALQCD,'12); Tamagaki('08); ...
- Quant. MC study *Lonardon et al.('14)*
- Quark Meson Coupling
Miyatsu et al.; Thomas (HHIQCD)
- $\Lambda\Lambda$ *K. Morita, T. Furumoto, AO, PRC91('15)024916*



Caveat: Missing data

Relativistic Mean Field with Multi-body couplings

$\sigma\omega\rho$ model +std. non-linear terms + multi-body couplings

$$\mathcal{L}_N = \bar{\psi} (i\gamma^\mu \partial_\mu - M_N - U_s - \gamma^\mu U_\mu) \psi + \mathcal{L}_{\sigma\omega\rho}$$

$$\mathcal{L}_{\sigma\omega\rho} = \frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma - \frac{1}{4} \omega_{\mu\nu} \omega^{\mu\nu} - \frac{1}{4} R_{\mu\nu} \cdot R^{\mu\nu} - \mathcal{V}_{\sigma\omega\rho}$$

$$U_s = -g_\sigma \sigma [1 + r_{\sigma\sigma}(1 - \sigma/f_\pi)] + g_\sigma \omega^\mu \omega_\mu / f_\pi [r_{\omega\omega} + r_{\sigma\omega\omega}(1 - \sigma/f_\pi)]$$

$$U_\mu = g_\omega \omega_\mu [1 - r_{\sigma\omega}\sigma/f_\pi + r_{\omega 3}\omega^\nu \omega_\nu / f_\pi^2]$$

$$+ g_\rho \tau \cdot R_\mu [1 - r_{\sigma\rho}\sigma/f_\pi + r_{\omega\rho}\omega^\nu \omega_\nu / f_\pi^2]$$

$$\mathcal{V}_{\sigma\omega\rho} = \frac{1}{2} m_\sigma^2 \sigma^2 - a_\sigma f \log(\sigma/f_\pi) + \frac{1}{4} c_{\sigma 4} (\sigma^4 - 4f_\pi \sigma^3)$$

$$- \frac{1}{2} m_\omega^2 \omega^\mu \omega_\mu [1 - c_{\sigma\omega}\sigma/f_\pi] - \frac{1}{4} c_{\omega 4} (\omega^\mu \omega_\mu)^2$$

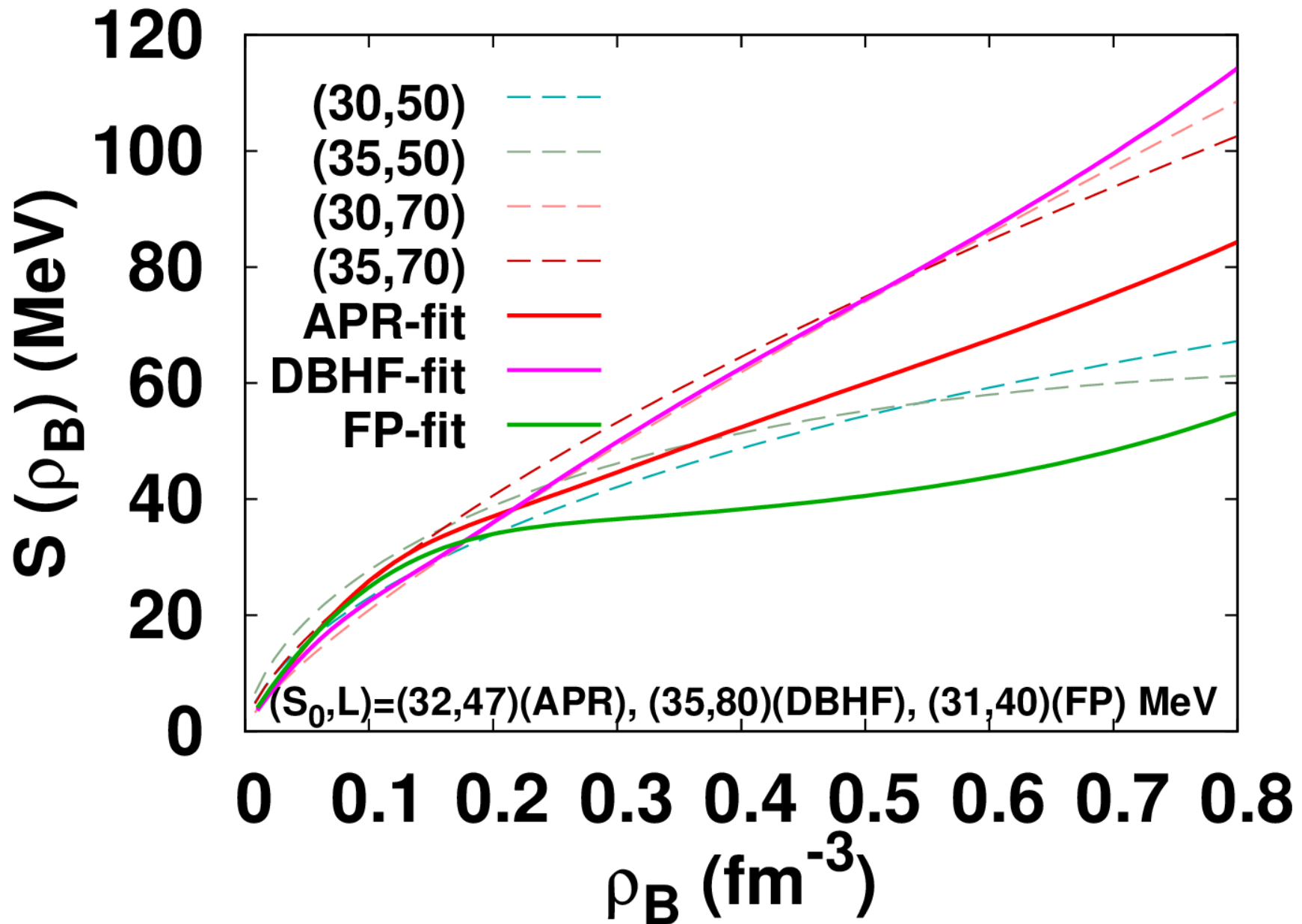
$$- \frac{1}{2} m_\rho^2 R^\mu \cdot R_\mu [1 - c_{\sigma\rho}\sigma/f_\pi + c_{\omega\rho}\omega^\mu \omega_\mu / f_\pi^2] - \frac{1}{4} c_{\rho 4} (R^\mu \cdot R_\mu)^2$$

$$f \log(x) = \log(1-x) + x + \frac{1}{2} x^2 \quad a_\sigma = f_\pi^2 (m_\sigma^2 - m_\pi^2) / 2 - f_\pi^4 c_{\sigma 4}$$

Fitting “Ab initio” EOS via RMF

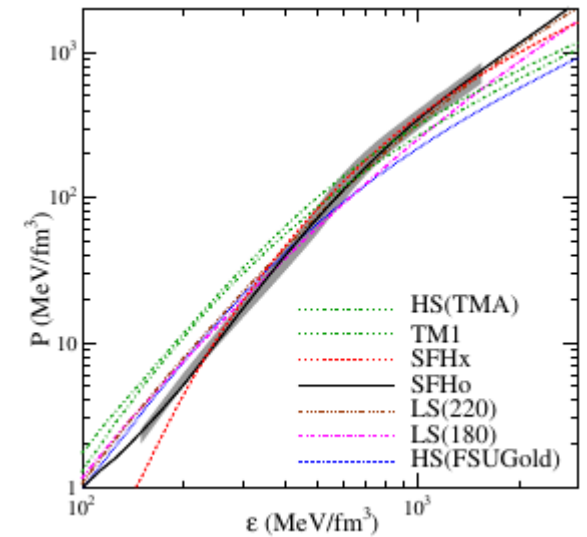
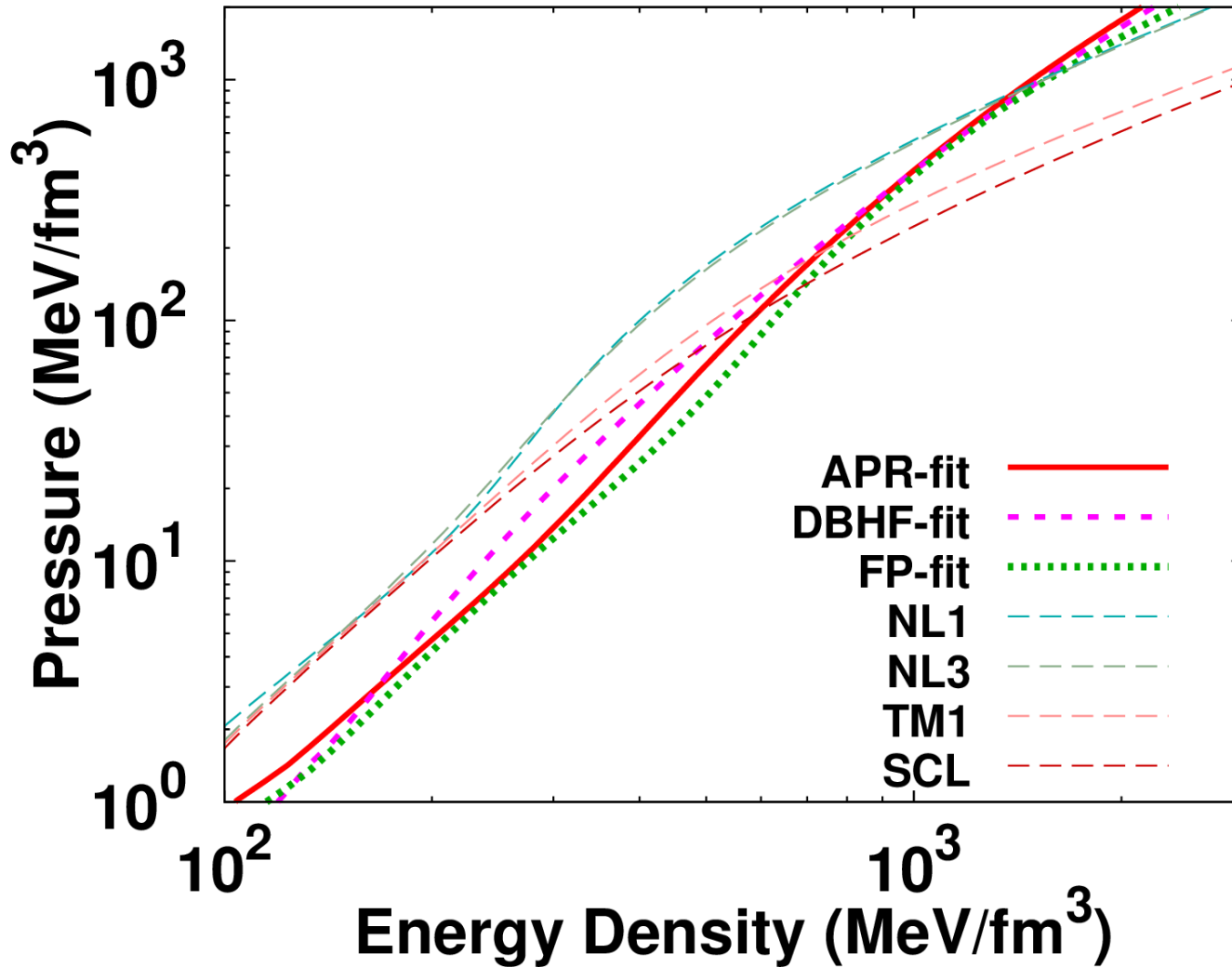
- “Ab initio” EOS under consideration
 - FP: Variational calc. (Av14+3NF(att.+repl.))
B. Friedman, V.R. Pandharipande, NPA361('81)502.
 - APR: Variational chain summation (Av18+ rel. corr.+3NF)
A. Akmal, V.R.Pandharipande, D.G. Ravenhall, PRC58('98)1804.
 - DBHF: Dirac Bruckner approach (Bonn A)
G. Q. Li, R. Machleidt, R. Brockmann, PRC45('92)2782
- RMF with multi-body couplings: 16 parameters
 - n=3 *Tsubakihara, AO, NPA914 ('13), 438.*
 - Working hypothesis: σ self-energy: SCL2 model *Tsubakihara, AO ('07)*
 $M_N \rightarrow 0 @ \sigma \rightarrow f_\pi$
- Markov Chain Monte-Carlo (MCMC)-like parameter search
 - Langevin type shift+Metropolis judge
 - Simultaneous fit of SNM and PNM \rightarrow std. dev=0.5-1.0 MeV

Symmetry Energy



Neutron Star Matter EOS

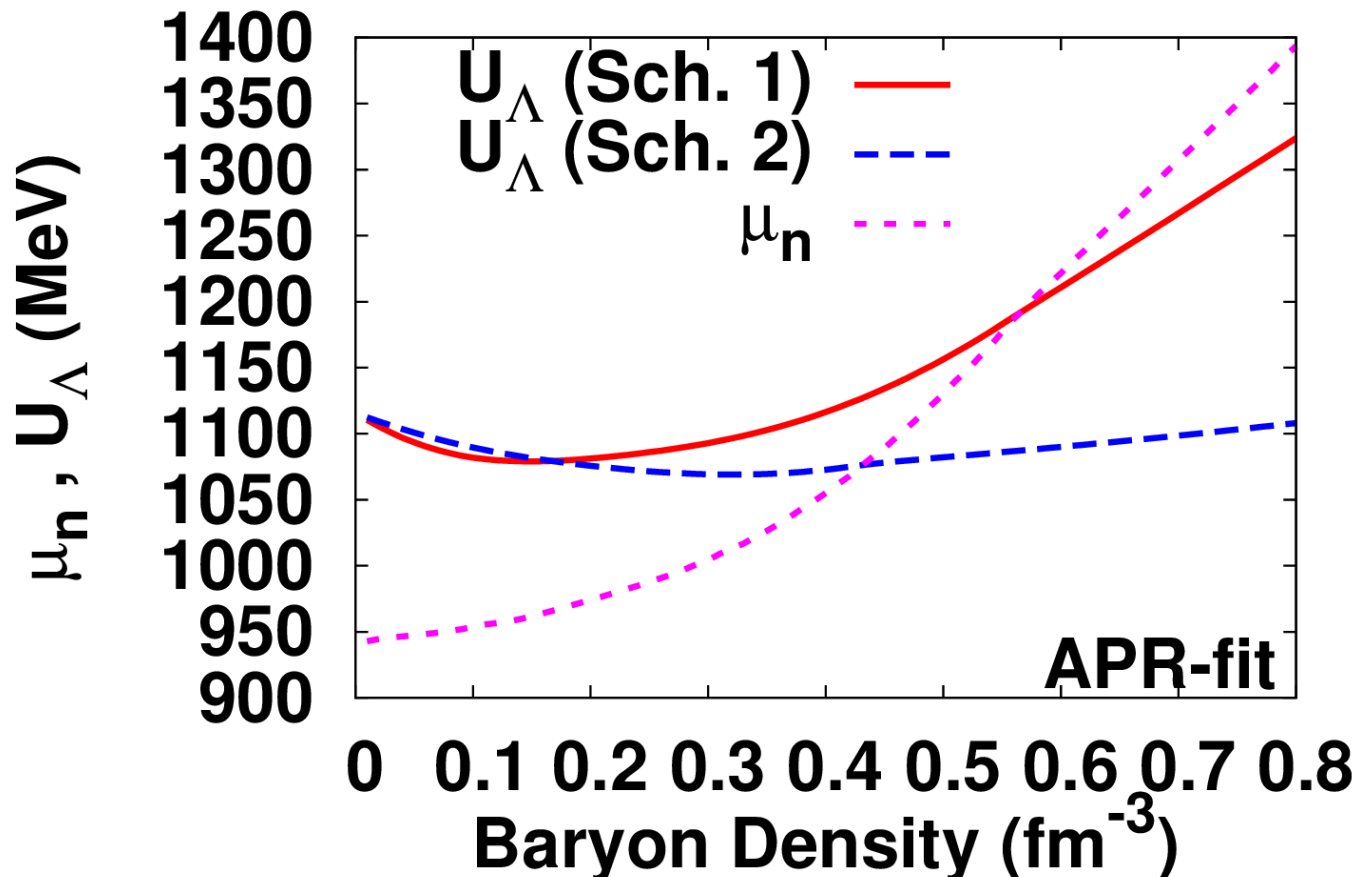
Neutron Star Matter EOS



A. W. Steiner, M. Hempel, T. Fischer, ApJ 774 (2013) 17 (TMA+NSE w/ excl. vol.)

NS matter in “ab initio”-fit + Λ

- Λ potential in nuclear matter at $\rho_0 \sim -30$ MeV
 - Scheme 1: $U_\Lambda(\rho) = \alpha U_N(\rho)$
 - Scheme 2: $U_\Lambda(\rho) = 2/3 U_{N^{n=2}}(\rho) + \beta U_{N^{n>2}}(\rho)$



M-R curve of Neutron Stars

