

Higher-order symmetry energy parameters and neutron star properties

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in collaboration with

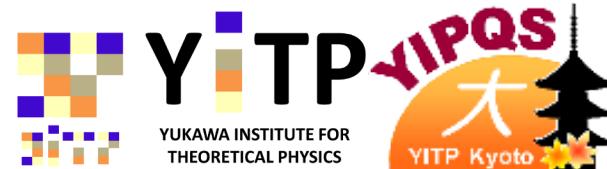
E. E. Kolomeitsev (Matej Bel U.), James M. Lattimer

(Stony Brook), Ingo Tews (LANL), Xuhao Wu (Nankai U./YITP)

*5th Joint Meeting of the NP Divisions
of APS and JPS*

Oct. 23-27, 2018, Waikoloa, Hawaii, USA

- *I. Tews, J. M. Lattimer, AO, E.E.Kolomeitsev, ApJ 848('17) 105 [arXiv:1611.07133]*
- *AO, Kolomeitsev, Lattimer, Tews, X.Wu, in prog.*



First, I would like to thank the audience ...

Saturday Afternoon, October 27

Start Time	Kona 4 Session MA: Quantum Computing and Machine Learning for Nuclear Physics Chair: Phiala Shanahan	Kohala 1 Session MB: Quiescent Stellar Burning Chair: Frank Strieder	Kohala 2 Session MC: Nuclear Reactions 2 Chair: Takashi Nakamura	Kohala 3 Session MD: Light Elements in Nuclear Astrophysics Chair: Carl Brune	King's 1 Session ME QCD Theory II Chair: Atsushi Hosaka	King's 2 Session MF: Nuclear Matter and Nuclear Astrophysics Chair: Jeremy Holt	King's 3 Session MG: Nuclear Theory 4 Chair: Yutaka Utsuno
14:00	Natalie Kloco	Alexander Laminack	Oleg B. Tarasov	Yudong Luo	• Yuto Mori	Alis Rodriguez Manso	B. Alex Brown
14:15		Christopher J. Prokop	Hiroshi Suzuki	Seiya Hayakawa	Takuya Sugiura	Kyle Wayne Brown	David Kekejian
14:30		Daniel Robertson	Mitsunori Fukuda	Ingo Wiedenhoever	Masayuki Wakayama	Masanori Kaneko	Konstantinos Kravvaris
14:45	Christine Muschik	Bryce Frenzt	Midori Miwa	Azusa Inoue	Kadir Utku Can	Matthew E. Caplan	Caroline Robin
15:00		Richard J. DeBoer	Jesus F. Perello	Nabin Rijal	✗ Akio Tomiya	Shigehiro Yasui	Shalom Shlomo
15:15		Chad C. Ummel	Shota Y. Matsumoto	Maria Gatu Johnson	Wayne Nicholas Polyzoou	Chinatsu Watanabe	Matteo Vorabbi
15:30	Akinori Tanaka	Michael T. Febraro	Alan McIntosh	Alex Zylstra	George Rosensteel	Toshiki Maruyama	William Ormand
15:45		Devin Connolly	Sharon Stephenson	Bryant Vande Kolk	Mamiya Kawaguichi	Tsuyoshi Miyatsu	
16:00		Devin Connolly	Lauren Heilborn	Gary Grim	Makito Oi	Kaoru Shoji	
16:15	Shinji Takeda		Young Jin Kim			• Sanjay Reddy	
16:30						• Hajime Togashi	
17:00						• Akira Ohnishi	

Last talks



Symmetry Energy Parameters & Neutron Star Radius

- Nuclear Matter Symmetry Energy parameters (S_0, L) are closely related to Neutron Star Properties, e.g. $R_{1.4} = R_{NS}(M = 1.4M_{\odot})$

c.f. Workshop (Tue, 2WAA & 2WAB), this session, ...

- How can we constrain (S_0, L) ?
→ Nuclear Exp't. & Theory, Astro. Obs., **Unitary gas**
- Conjecture: UG gives the lower bound of neutron matter energy.

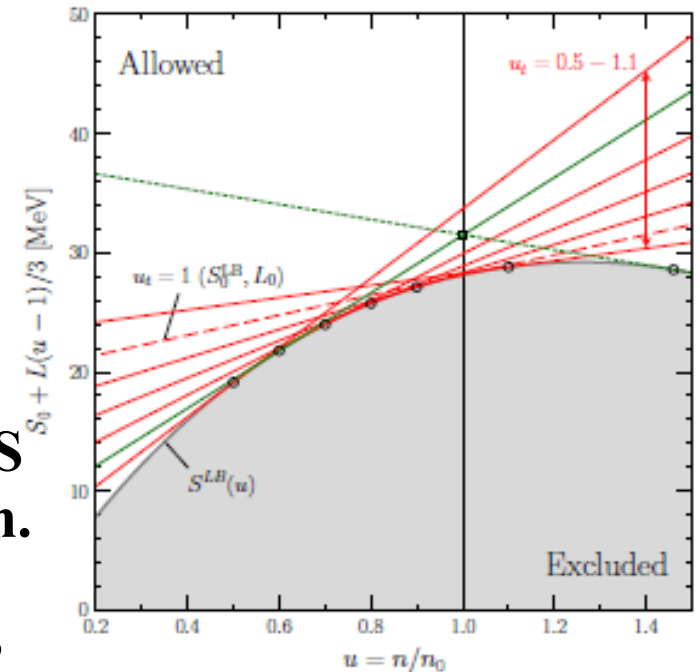
Tews, Lattimer, AO, Kolomeitsev (TLOK), ApJ ('17)

$$S(n) = E_{\text{PNM}} - E_{\text{SNM}} \geq E_{\text{UG}} - E_{\text{SNM}}$$

$$E_{\text{UG}} = \xi E_{\text{FG}} \quad (\xi \simeq 0.38)$$

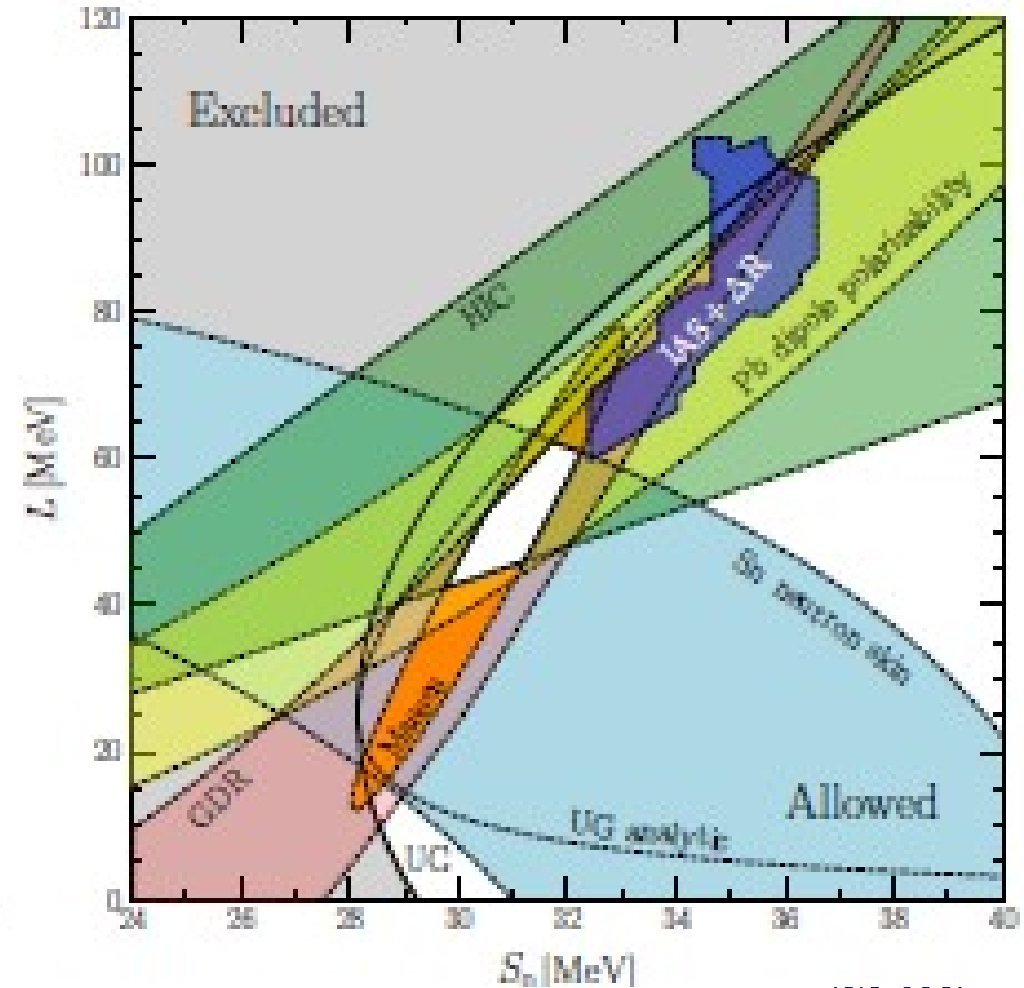
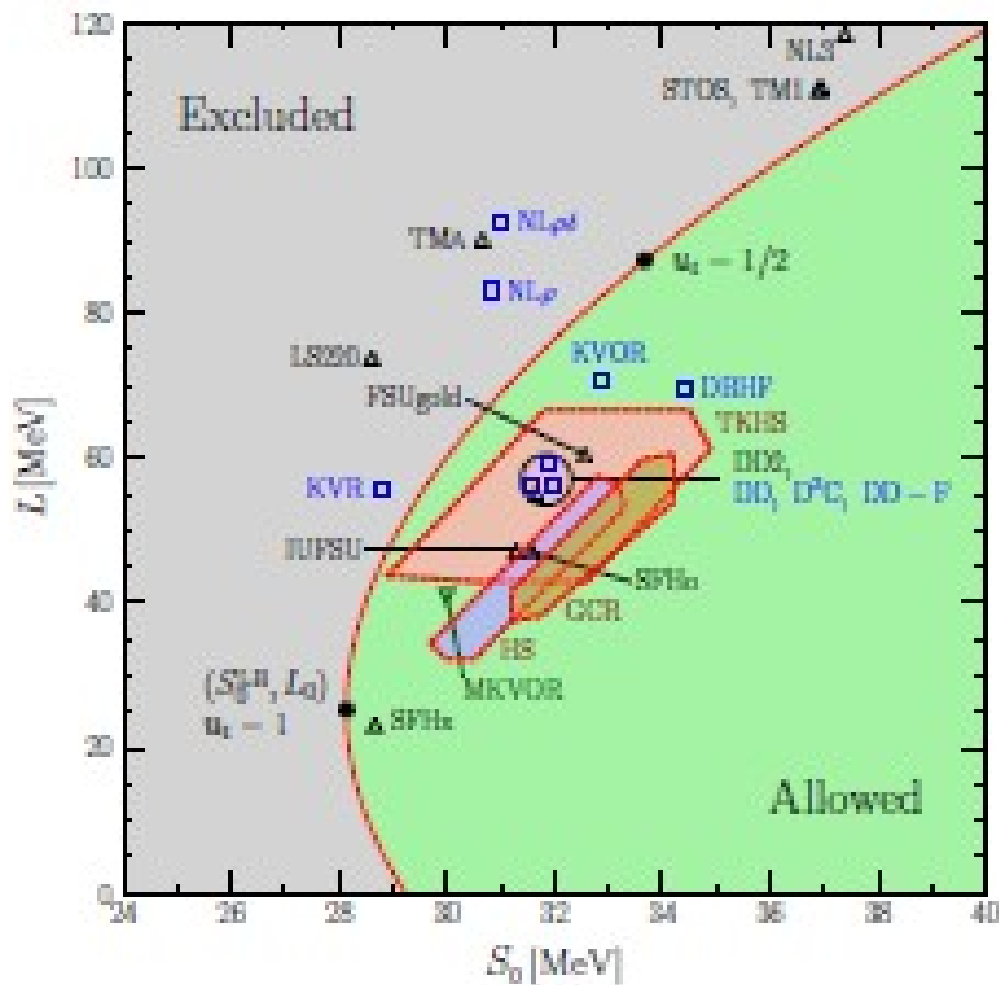
↑
Sym. Nucl. Matter EOS
is relatively well known.

→ For a given L , lower bound of S_0 exists



Constraint on (S_0, L) from Lower Bound of PNM Energy

- Unitary gas + $2 M_{\odot}$ constraints rule out 5 EOSs out of 10 numerically tabulated and frequently used in astrophys. calc.



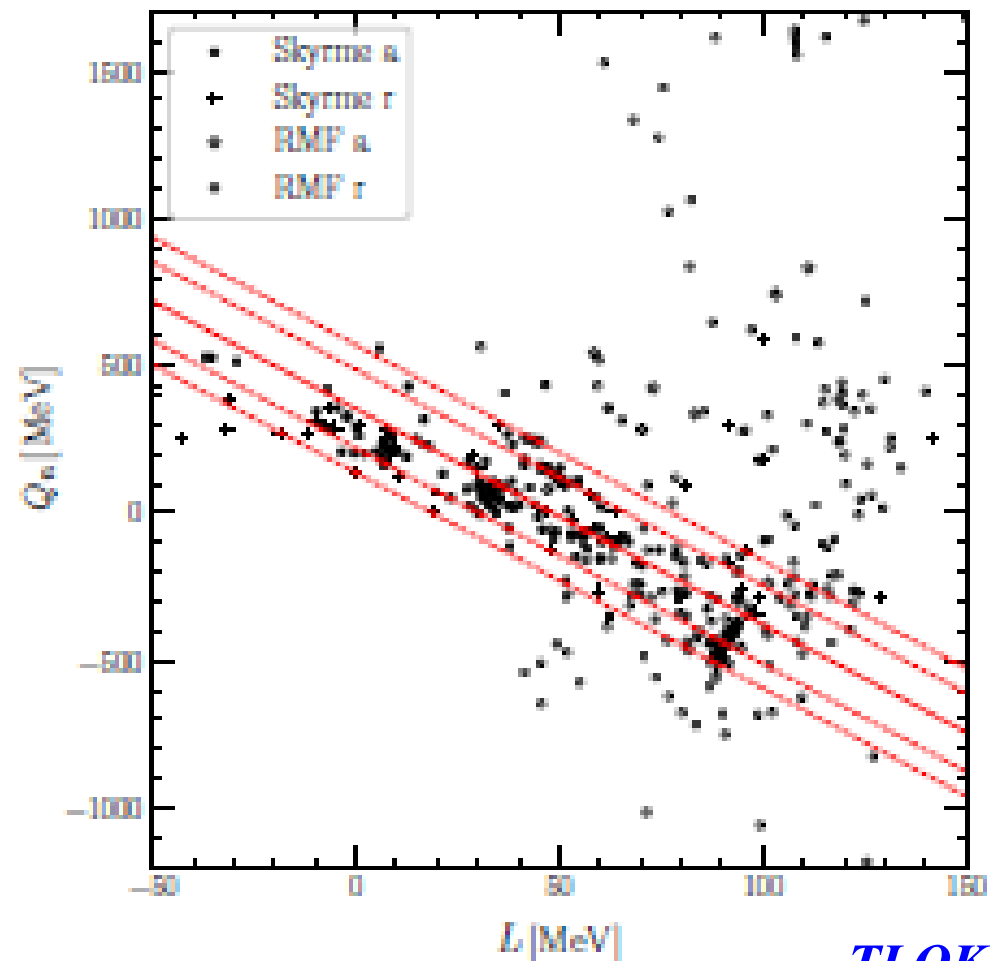
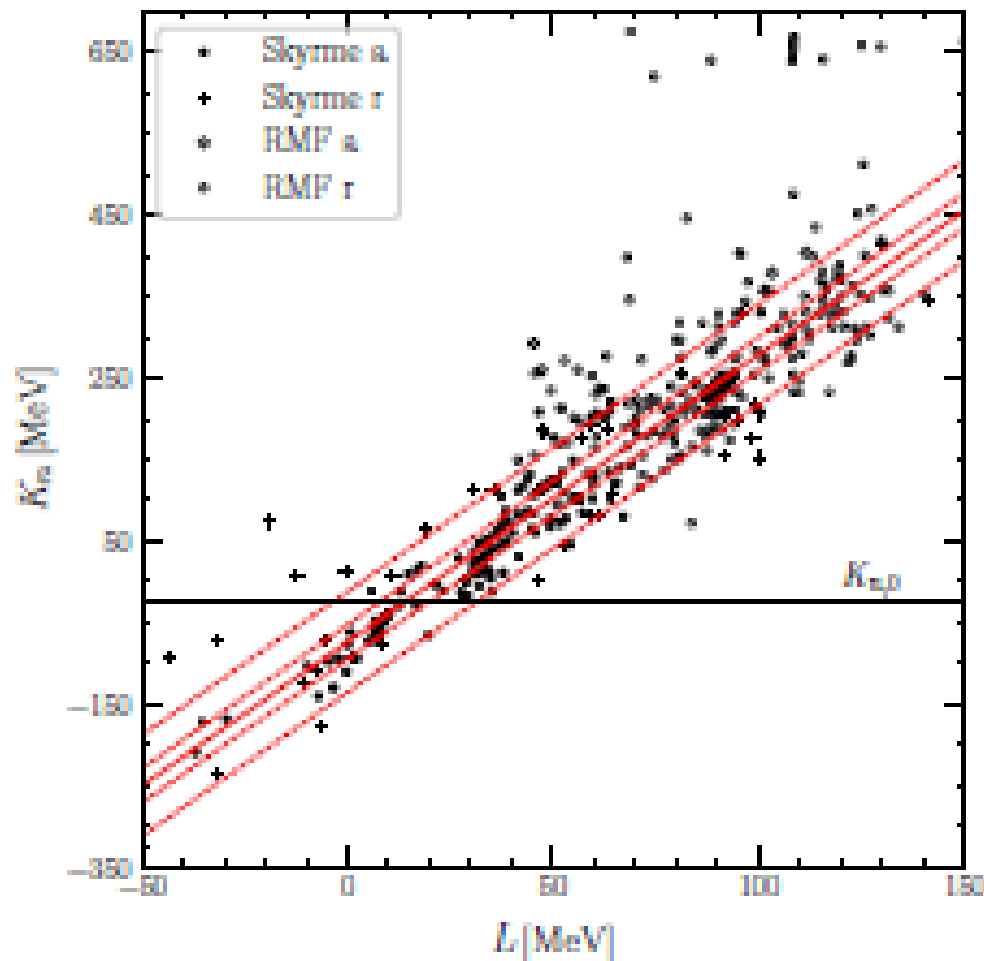
TLOK

Further Constraints on Higher-Order Sym. E. parameters

- K_n and Q_n are correlated with L in “Good” theoretical models.

$$K_n = 3.534L - (74.02 \pm 21.17)\text{MeV}$$

$$Q_n = -7.313L + (354.03 \pm 133.16)\text{MeV}$$



TLOK

Purpose & Contents

- **Question:**

What are the effects of these higher-order sym. E. parameters on MR curve of NS ?

- **This work:**

TLOK + $2 M_{\odot}$ constraints + k_F expansion $\rightarrow R_{1.4}$

- **Contents**

- Introduction
- Fermi momentum (k_F) expansion of EOS
- Neutron Star MR curve
- (Limitations & Prospect)
- Summary

Fermi momentum (k_F) expansion

■ Saturation & Symmetry Energy Parameters

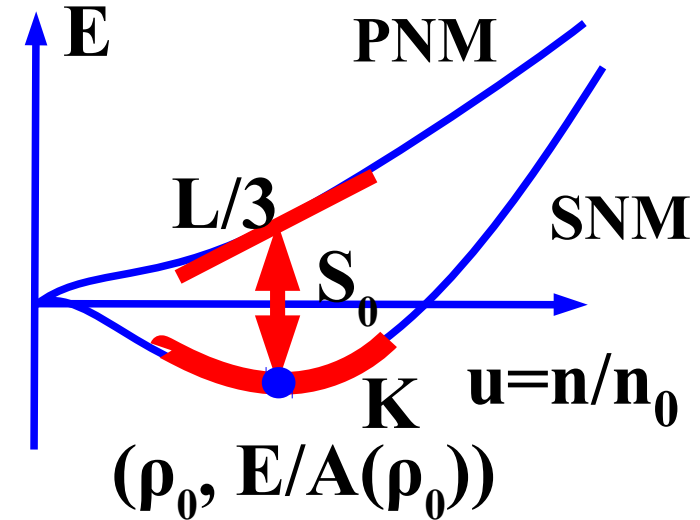
TLOK

$$E_{\text{NM}}(u, \alpha) = E_{\text{SNM}}(u) + \alpha^2 S(u)$$

$$E_{\text{SNM}}(u) \simeq E_0 + \frac{K_0}{18}(u-1)^2 + \frac{Q_0}{162}(u-1)^3$$

$$S(u) \simeq S_0 + \frac{L}{3}(u-1) + \frac{K_s}{18}(u-1)^2 + \frac{Q_s}{162}(u-1)^3$$

$$(u = n/n_0, \alpha = (n_n - n_p)/n)$$



Energy does not approach zero at $n \rightarrow 0$.

■ Fermi momentum expansion (\sim Skyrme type EDF)

- Generated many-body force is given by $k_F \propto u^{1/3}$ m^*

$$E_{\text{SNM}}(u) \simeq T_0 u^{2/3} + \underline{a_0 u} + \underline{b_0 u^{4/3}} + \underline{c_0 u^{5/3}} + \underline{d_0 u^2}$$

$$S(u) \simeq T_s u^{2/3} + \underline{a_s u} + \underline{b_s u^{4/3}} + \underline{c_s u^{5/3}} + \underline{d_s u^2}$$

Kin. E. Two-body Density-dep. pot.

Expansion Coefficients

- Coefficients (a,b,c,d) are represented by Saturation and Symmetry Energy Parameters

TLOK

$$\begin{array}{llll}
 a_0 = -4T_0 & +20E_0 & +K_0 & -Q_0/6 \\
 b_0 = 6T_0 & -45E_0 & -5K_0/2 & +Q_0/2 \\
 c_0 = -4T_0 & +36E_0 & +2K_0 & -Q_0/2 \\
 d_0 = T_0 & -10E_0 & -K_0/2 & +Q_0/6 \\
 \\
 a_s = -4T_s & +20S_0 - 19L/3 & +K_s & -Q_s/6 \\
 b_s = 6T_s & -45S_0 + 15L & -5K_s/2 & +Q_s/2 \\
 c_s = -4T_s & +36S_0 - 12L & +2K_s & -Q_s/2 \\
 d_s = T_s & -10S_0 + 10L/3 & -K_s/2 & +Q_s/6
 \end{array}$$

$$\left(T_0 = \frac{3 \hbar^2 k_F (n_0)^2}{5 \cdot 2m}, \quad T_s = T_0 (2^{1/3} - 1) \right)$$

Tedious but straightforward calc.

TLOK+2M_⊙ constraints

■ TLOK constraints

- (S₀, L) is in Pentagon.

- (K_n, Q_n) are from TLOK constraint.

- K₀=(190-270) MeV

- (n₀,E₀) is fixed

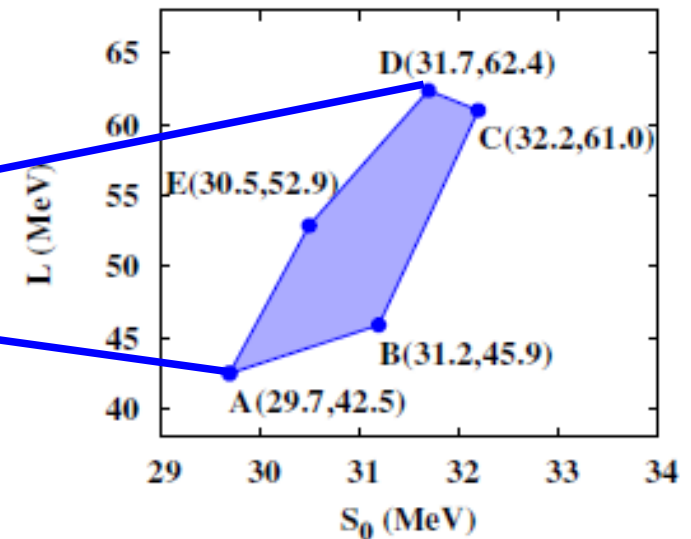
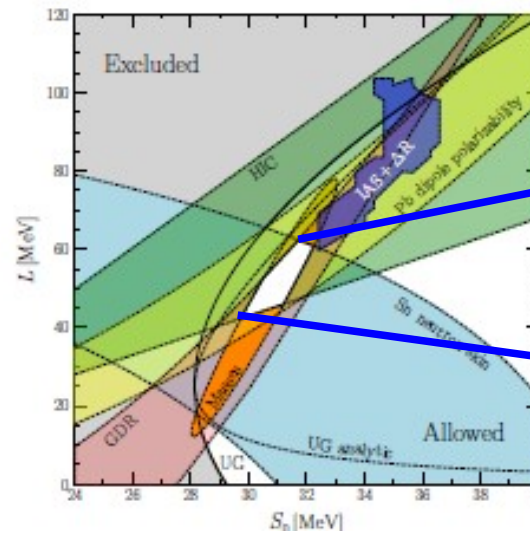
n₀=0.164 fm⁻³, E₀=-15.9 MeV (small uncertainties)

- Q₀ is taken to kill d₀ parameter

(Coef. of u². Sym. N. M. is not very stiff at high-density)

■ 2 M_⊙ constraint

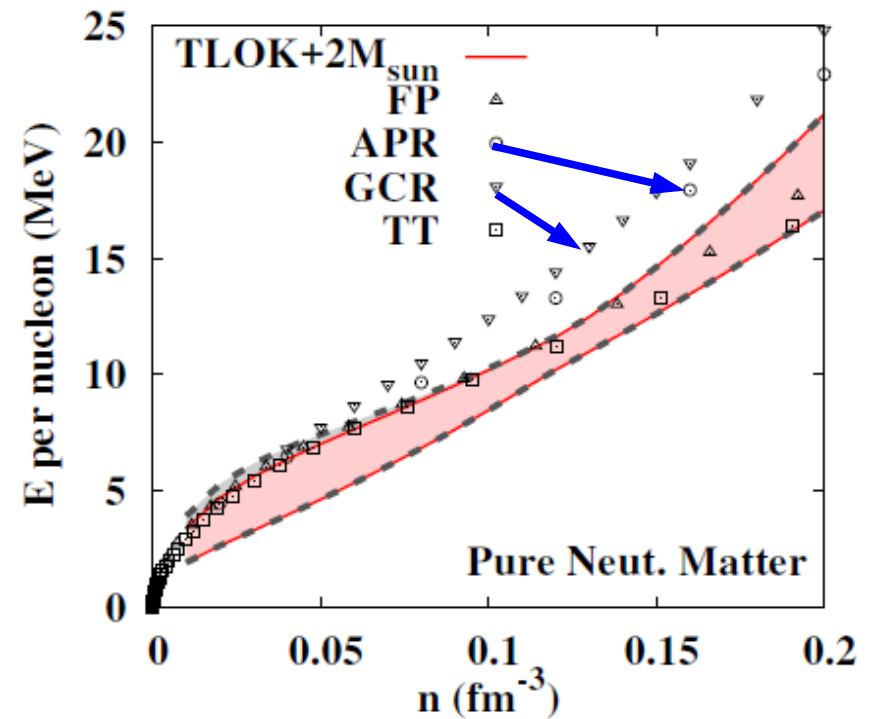
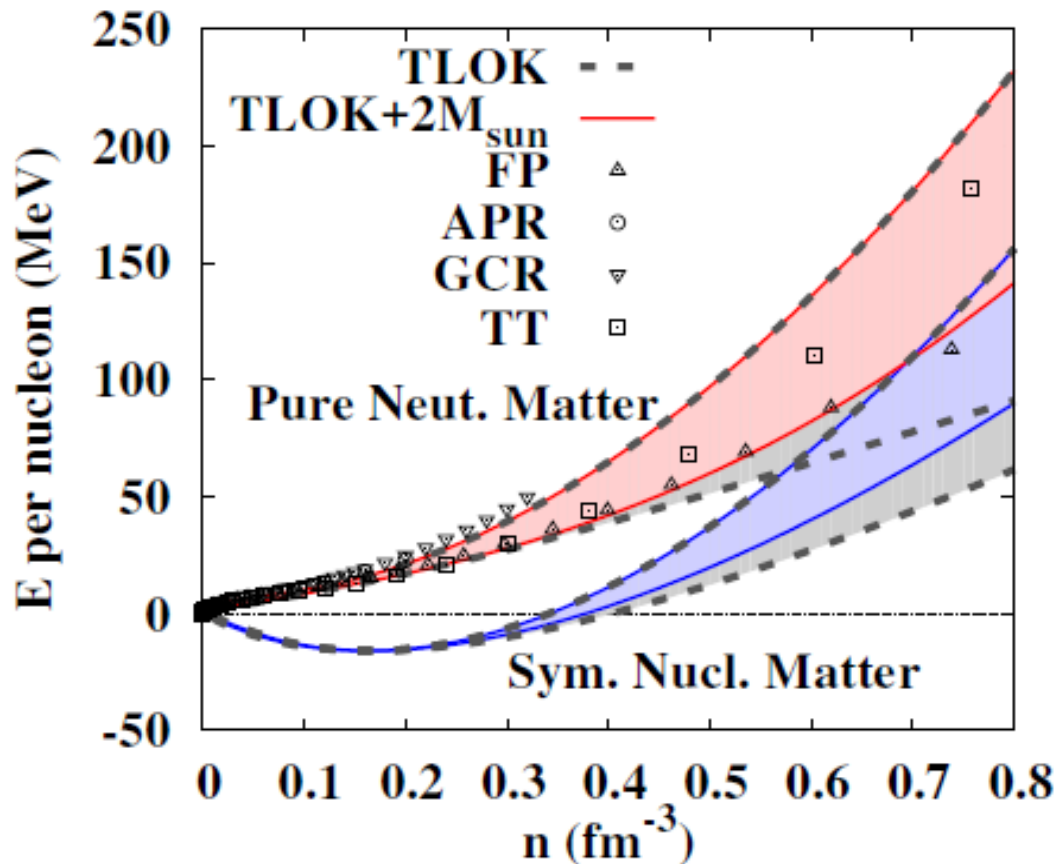
- EOS should support 2 M_⊙ neutron stars.



AO, Kolomeitsev, Lattimer, Tews, Wu (OKLTW), in prog.

TLOK+2M_⊙ constraints on EOS

- 2M_⊙ constraint narrows the range of EOS.
- Consistent with FP and TT(Togashi-Takano) EOSs.
- APR and GCR(Gandolfi-Carlson-Reddy) EOSs seems to have larger S₀ values.



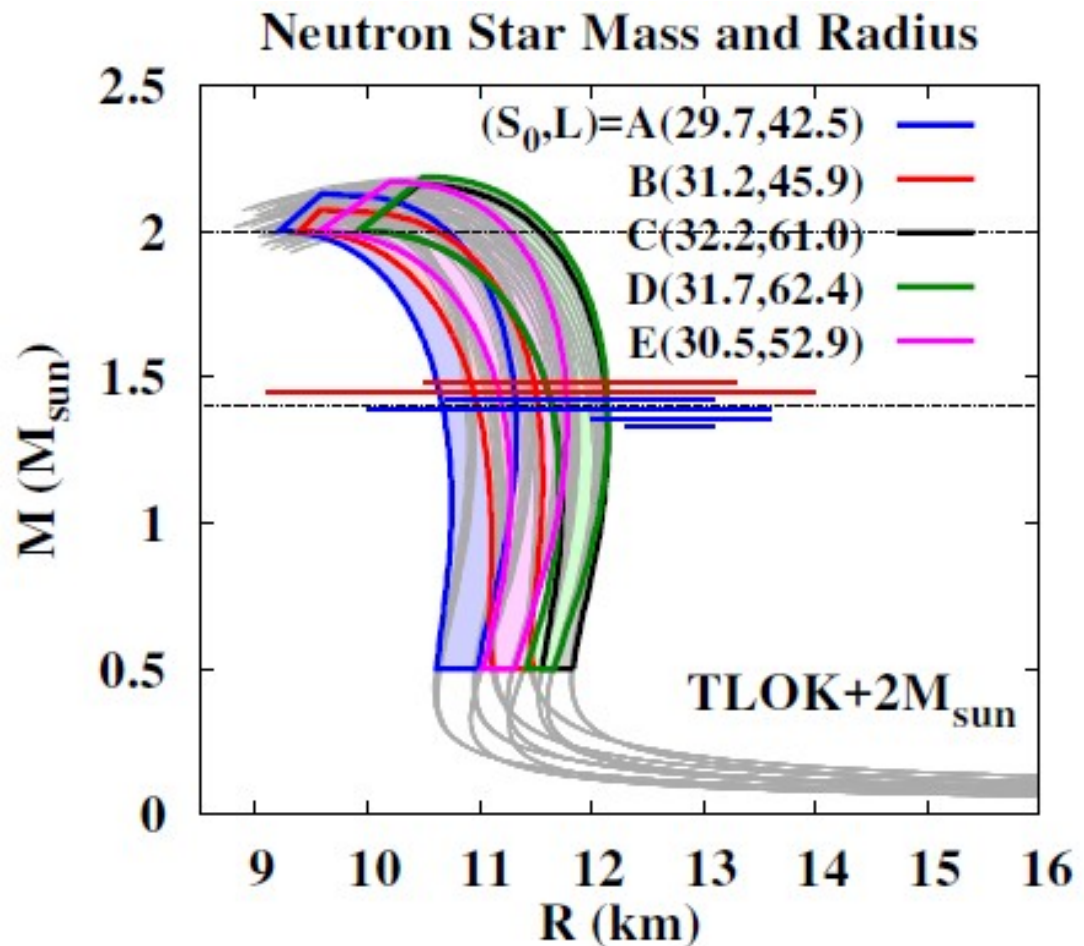
OKLTW, in prog.

Neutron Star MR curve

■ TLOK + $2 M_{\odot}$ constraints $\rightarrow R_{1.4} = (10.6-12.2)$ km

OKLTW, in prog.

- E and P are linear fn. of Sat. & Sym. E. parameters
 \rightarrow Min./Max. appears at the corners of pentagon (ABCDE).
- For a given (S_0, L) ,
unc. of $R_{1.4} \sim 0.5$ km
= unc. from higher-order parameters
- Unc. from $(S_0, L) \sim 1.1$ km
 \rightarrow We still need to fix (S_0, L) more precisely.



Neutron Star MR curve

- Our constraint is consistent with many of previous ones.
 - $R_{1.4}=(10.6-12.2)$ km *Present work (TLOK + 2 M_{\odot}) OKLTW, in prog.*

- LIGO-Virgo (Tidal deformability Λ from BNSM)
 - (10.5-13.3) km *Abbott+('18b)*
 - (9.1-14.0) km *De+('18) (Λ)*

- Theoretical Estimates
 - (10.7-13.1) km

Lattimer, Prakash('16)

(10.0-13.6) km

Annala+('18) (χ EFT+pQCD)

(10-13.6) km

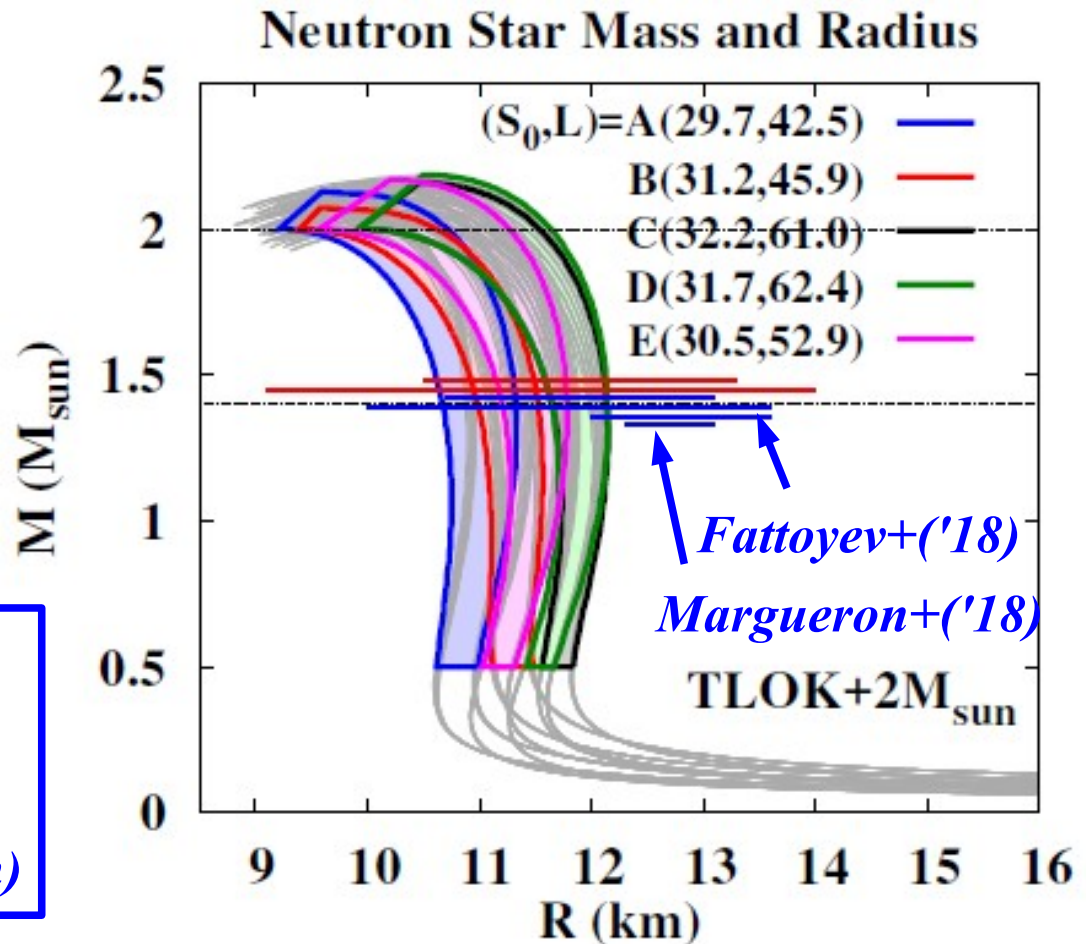
Tews+('18)(χ EFT+ c_{ν})

(12.0-13.6) km

Fattoyev+('18) (PREX)

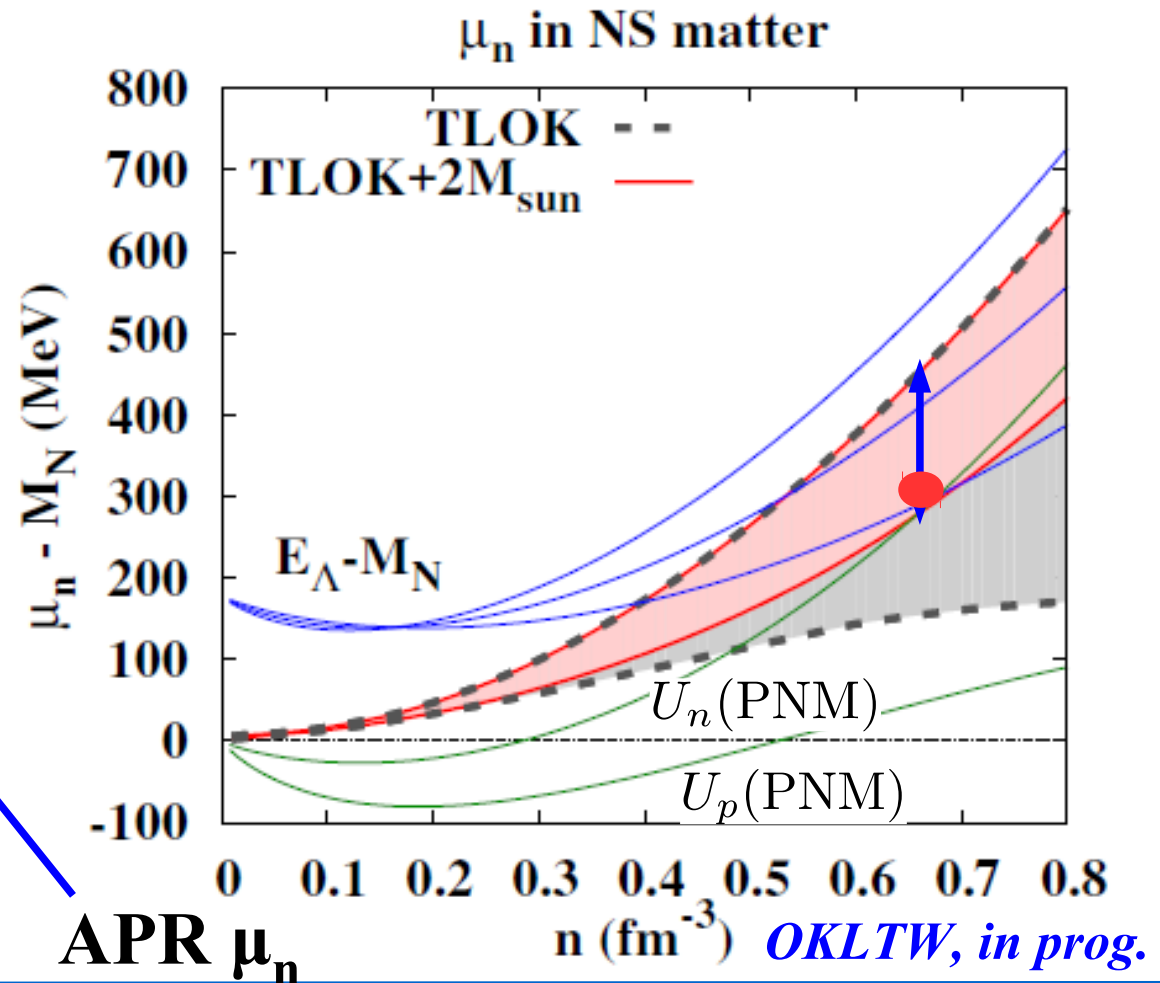
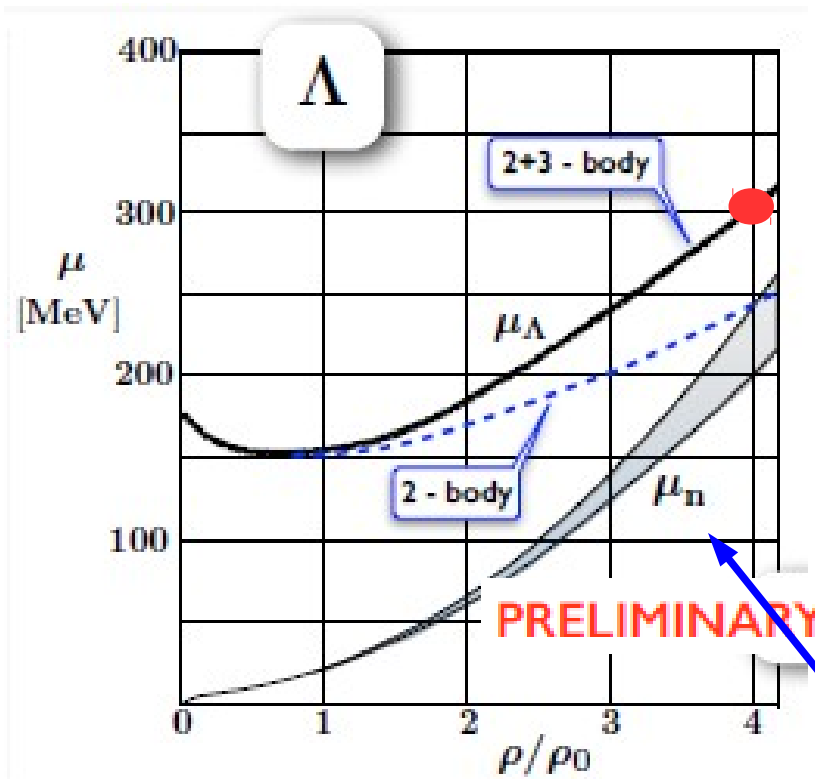
12.7 ± 0.4 km

Margueron+('18) (n expansion)



Neutron Chemical Potential in NS

- Λ appears in neutron stars if $E_{\Lambda}(p=0) = M_{\Lambda} + U_{\Lambda} < \mu_n$
- W. Weise's conjecture: U_{Λ} in χ EFT (2+3 body) is stiff enough.
- But μ_n is larger with TLOK+2 M_{\odot} constraints



W. Weise, NFQCD2018 (2018.06);
Gerstung, Kaiser, Weise, in prog.

Neutron Chemical Potential in NS

Neutron Chemical Potential

$$\mu_n + M_N = \frac{\partial(nE)}{\partial n_n} = E + u \frac{\partial E}{\partial u} + 2\alpha(1 - \alpha)S(u)$$

Single particle potential

$$U_\Lambda(u) = \frac{\partial(nV)}{\partial n_\Lambda}$$

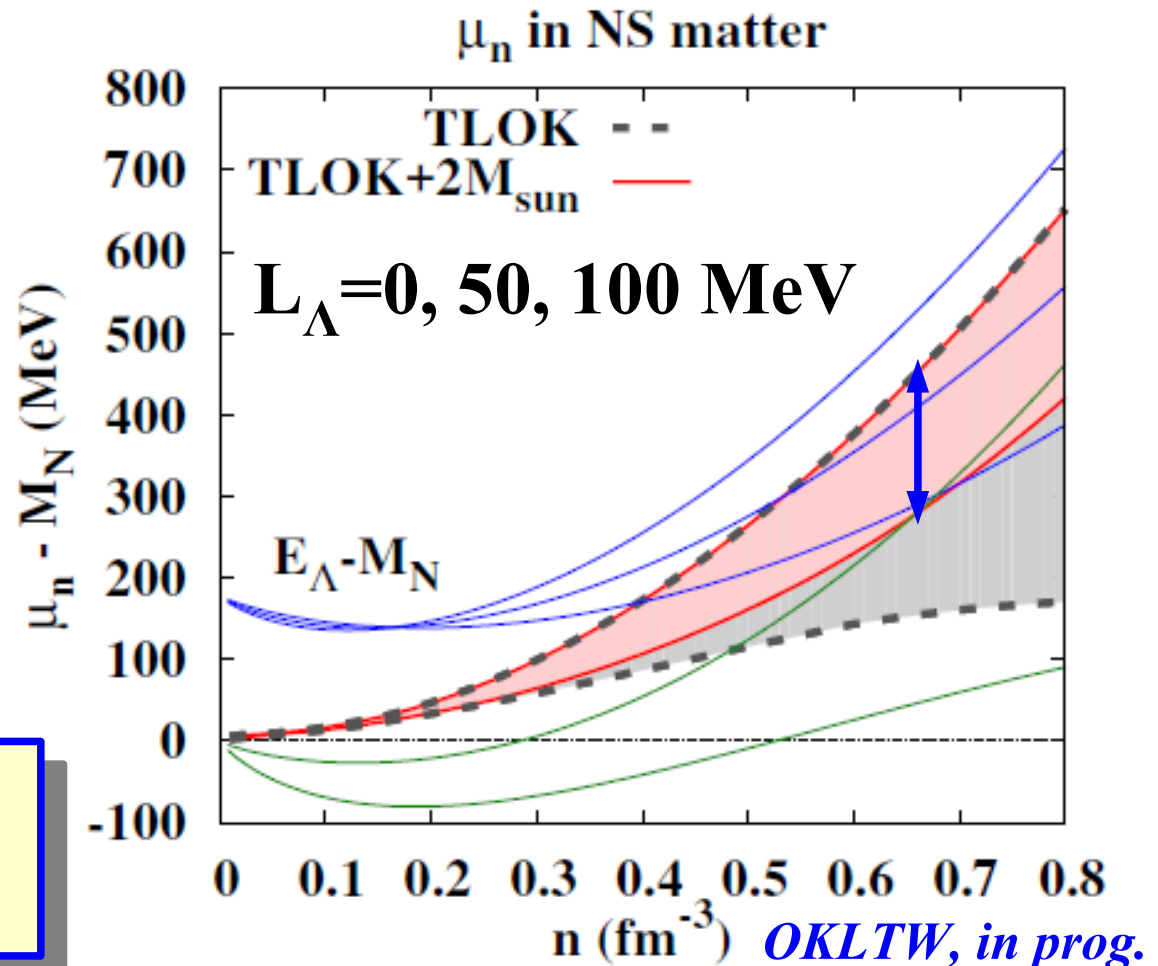
$$\simeq U_{0\Lambda} + \frac{L_\Lambda}{3}(u - 1)$$

$$U_{0\Lambda} \simeq -30 \text{ MeV}$$

$$L_\Lambda = ???$$

($L_\Lambda < 0$ in most of RMF before 2010)

Sym. E. and L_Λ determine the emerging density of Λ .



Reservations

- Only massless electrons are considered and Crust EOS is ignored.
 - With μ , chemical potential may be reduced a little.
- Non-relativistic kinetic energy is used.
 - With rel. K.E., E per nucleon is modified by 0.03 MeV @ $10 n_0$ as long as Sat. and Sym. E parameters are fixed.
- Function form is limited to k_F expansion with $u^{k/3}$ (k=2-6).
 - $R_{1,4}$ range becomes narrower with k=2-5.
 - Density expansion gives EOSs very sensitive to parameters.
- Smooth $E(u)$ (= No phase transition) is assumed.
 - We expect QCD phase transition at (5-10) n_0 from recent BES data of directed flow *Nara, Niemi, AO, Stoecker ('16)*
 - Transition to quark matter may not soften EOS drastically.
- Causality is violated at high densities, $n > (4-6) n_0$.

To Do (or Prospect)

- Baryons other than nucleons Λ , Δ , Ξ , Σ , ...
- Connecting to Hadron Resonance Gas (HRG) EOS
 - HRG EOS

mass and kinetic E of hadrons with $M < 2$ GeV + simple potential E

$$\varepsilon_{\text{HRG}} = \mathcal{T} + cn^2$$

or Lattice EOS in HIC (No saturation, No constraint from NS).

- We need to guess the potential energy density more seriously for consistent understanding of HIC, Nuclear, and NS physics.

$$\varepsilon = \mathcal{T} + \mathcal{V} \leftarrow \text{Nuclear and NS physics}$$

- Connecting to Quark(-Gluon) matter EOS
 - Embed model-H singularities *E.g. Nonaka, Asakawa ('04)*
 - “Interpolation” of nuclear and quark matter EOS

Summary

- Tews-Lattimer-AO-Kolomeitsev ('17) constraints (S_0 , L , K_n , Q_n) and $2 M_\odot$ constraint with the aid of Fermi momentum (k_F) expansion lead to the constraint on $1.4 M_\odot$ neutron star radius of (10.6-12.2) km.
 - Consistent with many of other constraint.
- Appearance of hyperons and Deltas may be sensitive to the symmetry energy in addition to potential parameters, (U_{0B} , L_B).
 - We need to know the slope of potential in addition to the depth.
- Global EOS (HIC and Nuclear/NS matter) needs to be given in a way where HIC physicists and NS physicists admit.
E.g. “Hadron Resonance Gas (HRG)+Potential from NS”

Thank you for your attention and for staying until the end of scientific program of the meeting.

Further Constraint on Q_n

- $2 M_{\odot}$ requirement constrains Q_n further.

$$Q_n > -9.3L + 480 \text{ MeV}$$

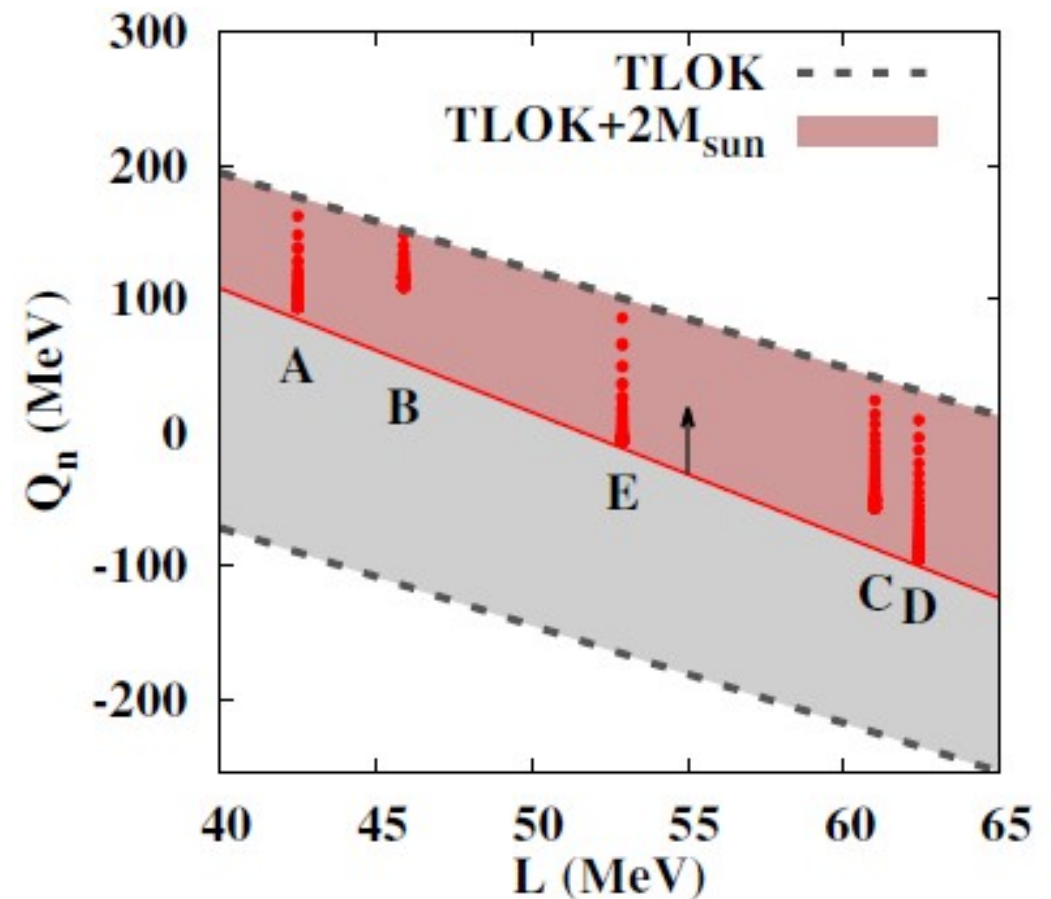


FIG. 4. Constraint on Q_n

AO, Kolomeitsev, Lattimer, Tews, Wu (OKLTW), in prog.

Unitary Gas Constraint

Tews, Lattimer, AO, Kolomeitsev (TLOK), ApJ ('17)

■ Conjecture:

Unitary gas gives the lower bound of neutron matter energy.

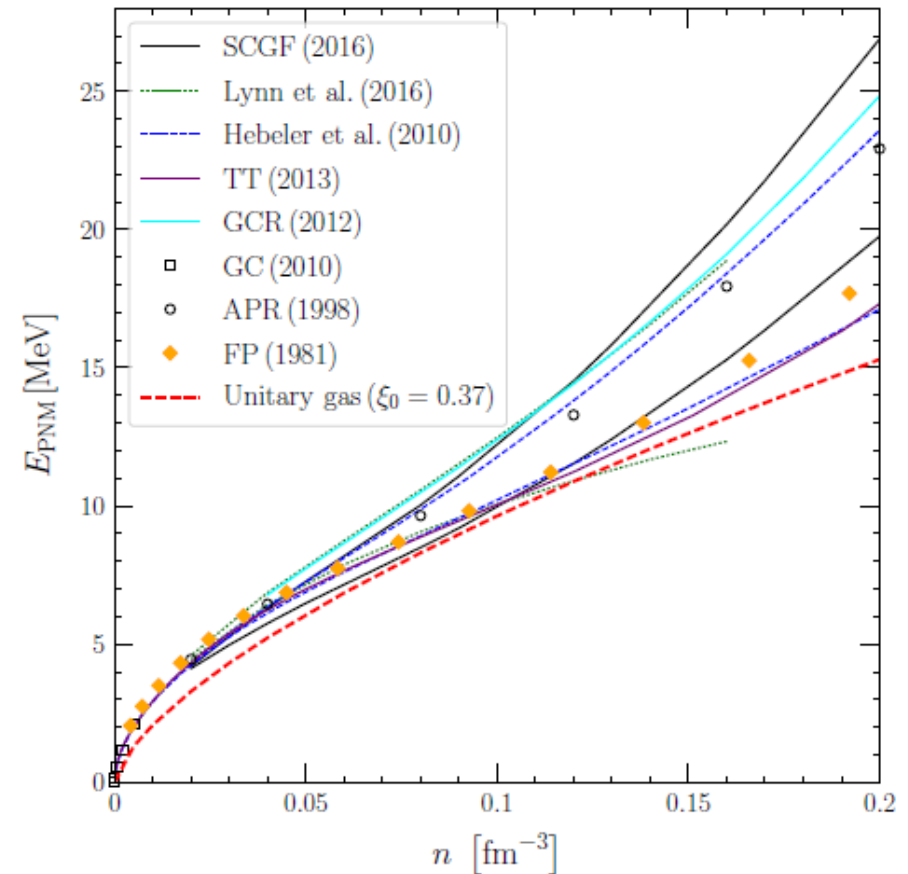
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$$E_{\text{UG}} = \xi E_{\text{FG}} \quad (\xi \simeq 0.38)$$

■ $a_0 = \infty$ in unitary gas

→ lower bound energy of $a_0 < 0$ systems (w/o two-body b.s.) ?

■ Supported by (most of) ab initio calc.



Potential Energy Density

- Potential Energy Density in the Fermi momentum expansion

$$\mathcal{V} = nV = \sum_{i,j \in B} n_i n_j v_{ij}(n)$$

Density-dependent NN interactions v_{ij} ($i, j = p$ or n) are known.

- Single particle potential

$$\begin{aligned} U_i &= \frac{\partial \mathcal{V}}{\partial n_i} = \sum_j n_j v_{ij}(n) + \sum_{jk} n_j n_k \frac{\partial v_{jk}(n)}{\partial n_i} \\ &= U_{0i} + \frac{L_i}{3} (u - 1) + \mathcal{O}((u - 1)^2) \\ &\simeq au + bu^{4/3} \end{aligned}$$

rearrangement
term

Again, a and b are given as a linear function of U_{0i} and L_i .