

Path optimization method with use of Neural Network for the Sign Problem in Field theories

Akira Ohnishi¹, Yuto Mori², Kouji Kashiwa³

1. *Yukawa Inst. for Theoretical Physics, Kyoto U.,*
2. *Dept. Phys., Kyoto U.,* 3. *Fukuoka Inst. Tech.*

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Collaborators

Akira Ohnishi¹, Yuto Mori², Kouji Kashiwa³

1. Yukawa Inst. for Theoretical Physics, Kyoto U.,

2. Dept. Phys., Kyoto U., 3. Fukuoka Inst. Tech.



Y. Mori
(grad. stu.)



K. Kashiwa



AO (10 yrs ago)

1D integral: Y. Mori, K. Kashiwa, AO, PRD 96 ('17), 111501(R) [arXiv:1705.05605]
 ϕ 4 w/ NN: Y. Mori, K. Kashiwa, AO, PTEP 2018 ('18), 023B04 [arXiv:1709.03208]
Lat 2017: AO, Y. Mori, K. Kashiwa, EPJ Web Conf. 175 ('18), 07043 [arXiv:1712.01088]
NJL thimble: Y. Mori, K. Kashiwa, AO, PLB 781('18),698 [arXiv:1705.03646]
PNJL w/ NN: K. Kashiwa, Y. Mori, AO, arXiv:1805.08940.
0+1D QCD: Y. Mori, K. Kashiwa, AO, in prep.

The Sign Problem

- When the action is complex, strong cancellation occurs in the Boltzmann weight at large volume. = **The Sign Problem**

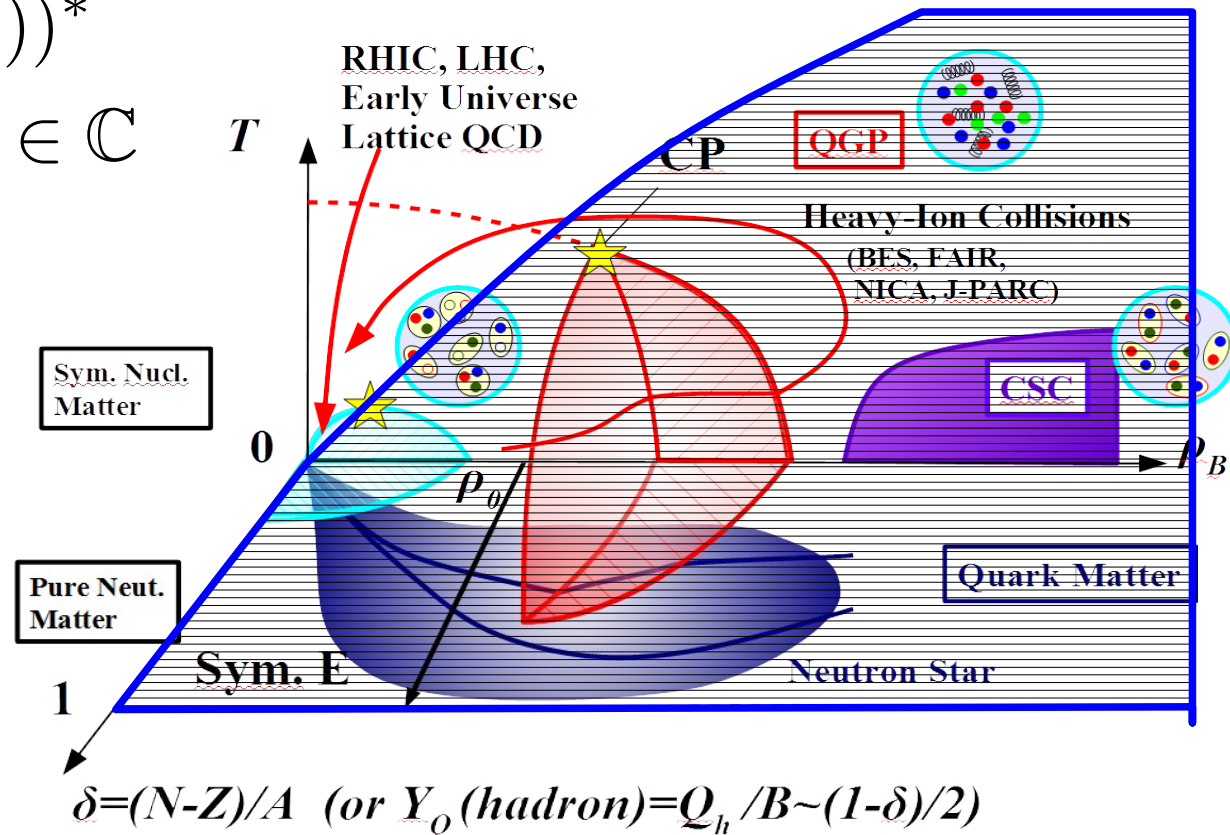
$$\mathcal{Z} = \int \mathcal{D}x e^{-S(x)} (S(x) \in \mathbb{C}) \rightarrow 0 (V \rightarrow \infty)$$

- Fermion det. is complex at finite density

$$\det D(\mu) = (\det D(-\mu^*))^*$$

$$\rightarrow S_{\text{eff}} = S - \log \det D \in \mathbb{C}$$

- Difficulty in studying finite density in LQCD
 \rightarrow Heavy-Ion Collisions, Neutron Star, Binary Neutron Star Mergers, Nuclei, ...



Approaches to the Sign Problem in Lattice 2018

■ Standard approaches

- Taylor expansion [*Ratti(Mon), Mukerjee(Tue), Steinbrecher(Wed)*]
- Imaginary μ (Analytic cont. / Canonical) [*Guenther, Goswami (Wed)*]
- Strong coupling [*Unger, Klegrewe (Fri)*]

→ *Mature, Practically useful, but cannot reach cold dense matter*

■ Integral in Complexified variable space

- Lefschetz thimble method [*Zambello (Mon)*]
- Complex Langevin method
[*Sinclair, Tsutsui, Attanasio, Ito, Josef (Mon), Wosiek (Fri)*]
- Path optimization method [*Lawrence, Warrington, Lamm (Mon), AO (Sat)*]
- Action modification (*e.g. Tsutsui, Doi ('16)*)

→ *Premature, but Developing !*

■ Other Approaches [*Ogilvie (Mon), Jaeger(Fri)*]

Integral in Complexified Variable Space

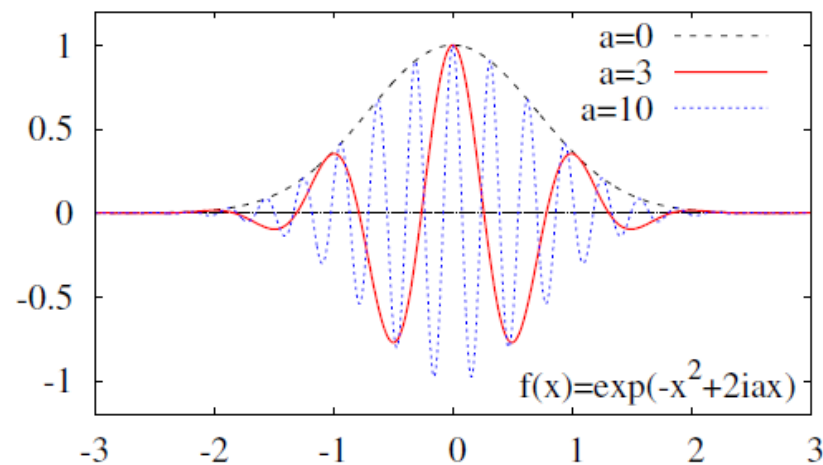
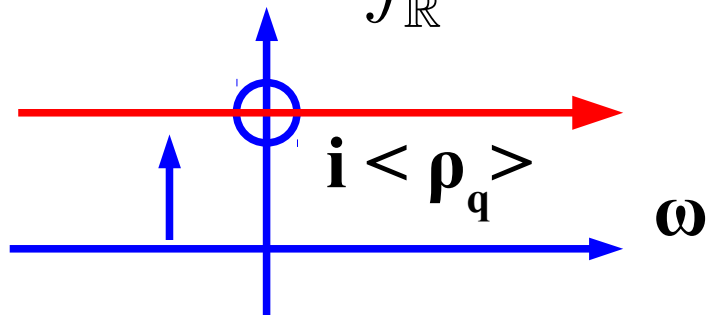
- Phase fluctuations can be suppressed by shifting the integration path in the complex plain.

$$\mathcal{Z} = \int_{\mathcal{C}_{\mathbb{R}}} \mathcal{D}x e^{-S(x)} = \int_{\mathcal{C}} \mathcal{D}z e^{-S(z)} = \int_{\mathcal{C}_{\mathbb{R}}} \mathcal{D}x J e^{-S(z(x))}$$

- Simple Example: Gaussian integral (bosonized repulsive int.)
Mori, Kashiwa, AO ('18b)

$$\int_{\mathbb{R}} d\omega e^{-\omega^2/2 + i\omega\rho_q} = \int_{\mathbb{R} + i\rho_q} d\omega e^{-\underbrace{(\omega - i\rho_q)^2}_{z^2} / 2 - \rho_q^2/2}$$

$$= \exp(-\rho_q^2/2) \int_{\mathbb{R}} dz e^{-z^2/2}$$

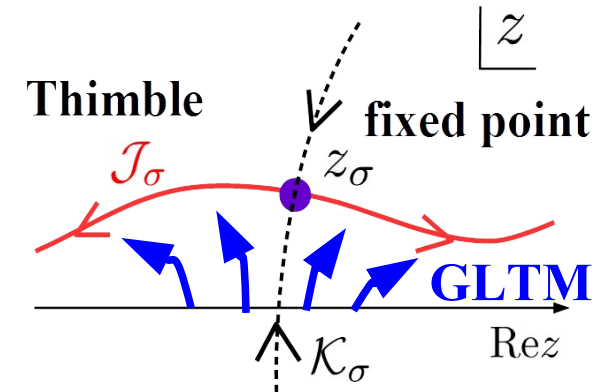


- Lefschetz thimble / Complex Langevin / Path Optimization

Lefschetz thimble method

*E. Witten ('10), Cristoforetti et al. (Aurora) ('12),
Fujii et al. ('13), Alexandru et al. ('16); [Zambello (Mon)]*

- Solving the flow eq. from a fixed point σ
→ Integration path (thimble)
Note: $\text{Im}(S)$ is constant on one thimble



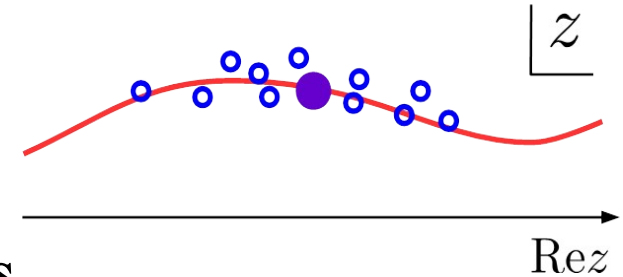
$$\mathcal{J}_\sigma : \frac{dz_i(t)}{dt} = \overline{\left(\frac{\partial S}{\partial z_i} \right)} \rightarrow \frac{dS}{dt} = \sum_i \left| \frac{\partial S}{\partial z_i} \right|^2 \in \mathbb{R}, \quad \mathcal{C} = \sum_\sigma n_\sigma \mathcal{J}_\sigma$$

■ Problem:

- Phase from the Jacobian (residual. sign pr.),
- Different Phases of Multi-thimbles (global sign pr.),
- Stokes phenomena, ...

Complex Langevin method

*Parisi ('83), Klauder ('83), Aarts et al. ('11),
Nagata et al. ('16); Seiler et al. ('13), Ito et al. ('16);
[Sinclair, Tsutsui, Attanasio, Ito, Joseph (Mon)]*



- Solving the complex Langevin eq. → Configs.

$$\frac{dz_i}{dt} = -\frac{\partial S}{\partial z_i} + \eta_i(t) (\eta_i : \text{White noise}), \quad \langle \mathcal{O}(x) \rangle = \langle \mathcal{O}(z) \rangle$$

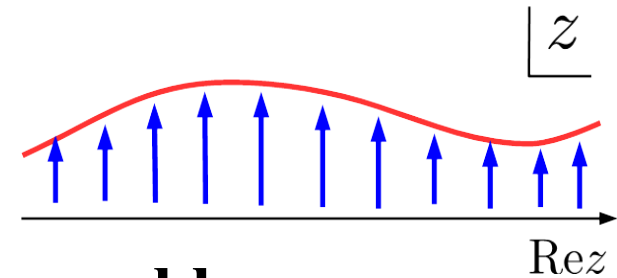
- No sign problem.

- Problem:

- CLM can give converged but wrong results,
and we cannot know if it works or not in advance.

Path optimization method

*Mori et al. ('17), AO, Mori, Kashiwa (Lattice 2017),
Mori et al. ('18), Kashiwa et al. ('18);
Alexandru et al. ('17 (Learnifold), '18 (SOMMe), '18),
Bursa, Kroyter ('18), [Lawrence, Warrington, Lamm (Mon)]*



- Integration **path** is **optimized** to evade the sign problem, i.e. to enhance the average phase factor.

$$\text{APF} = \langle e^{i\theta} \rangle_{\text{pq}} = \int dx e^{-S} / \int dx |e^{-S}| = \mathcal{Z} / \mathcal{Z}_{\text{pq}}$$

Sign Problem → Optimization Problem

- Cauchy(-Poincare) theorem: the partition fn. is invariant if
 - the Boltzmann weight $W = \exp(-S)$ is holomorphic (analytic),
 - and the path does not go across the poles and cuts of W .
- At Fermion det.=0, S is singular but W is not singular
 - Problem: quarter/square root of Fermion det.

Cost Function and Optimization

- **Cost function:** a measure of the seriousness of the sign problem.

$$\begin{aligned}\mathcal{F}[z(x)] &= \frac{1}{2} \int dx \left| e^{i\theta(x)} - e^{i\theta_0} \right|^2 \left| J(x) e^{-S(z(x))} \right| \\ &= |\mathcal{Z}| \left(\left| \langle e^{i\theta} \rangle_{\text{pq}} \right|^{-1} - 1 \right) = \mathcal{Z}_{\text{pq}} - |\mathcal{Z}| \\ &\quad [\theta = \arg(Je^{-S}), \theta_0 = \arg(\mathcal{Z})]\end{aligned}$$

- **Optimization:** the integration path is optimized to minimize the Cost Function.
(via Gradient Descent or Machine Learning)

- **Example: One-dim. integral \rightarrow Complete set**

$$z(x) = x + iy(x), \quad y(x) = \sum_n c_n H_n(x)$$

$$\mathcal{Z} = \int dx J(x) e^{-S(z(x))}, \quad J(x) = \frac{dz(t)}{dx}$$

Benchmark test: 1 dim. integral

- A toy model with a serious sign problem

J. Nishimura, S. Shimasaki ('15)

$$\mathcal{Z} = \int dx (x + i\alpha)^p \exp(-x^2/2)$$

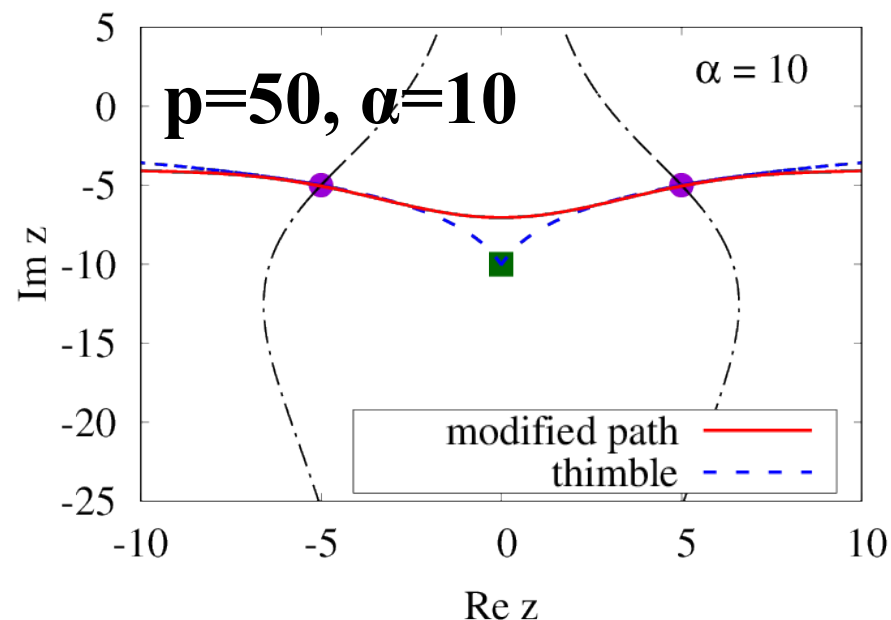
- Sign prob. is serious with large p and small $\alpha \rightarrow$ CLM fails

- Path optimization

$$y(x) = c_1 \exp(-c_2^2 x^2/2) + c_3, \quad J = 1 + i dy/dx$$

- Gradient Descent optimization

- Optimized path ~ Thimble around Fixed Points

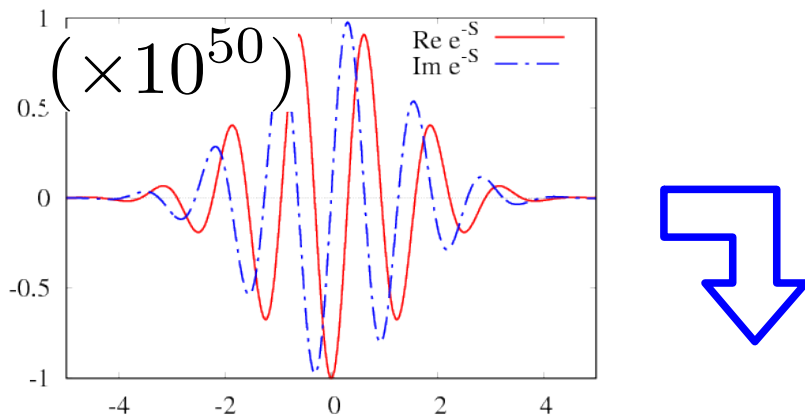


Mori, Kashiwa, AO ('17); AO, Mori, Kashiwa (Lat 2017)

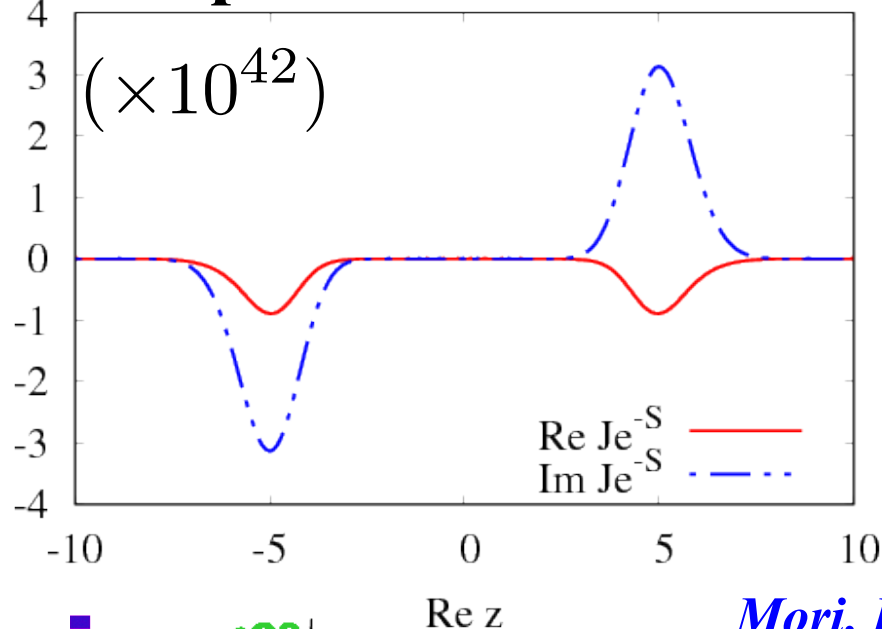
Benchmark test: 1 dim. integral

Stat. Weight $J e^{-S}$

On Real Axis

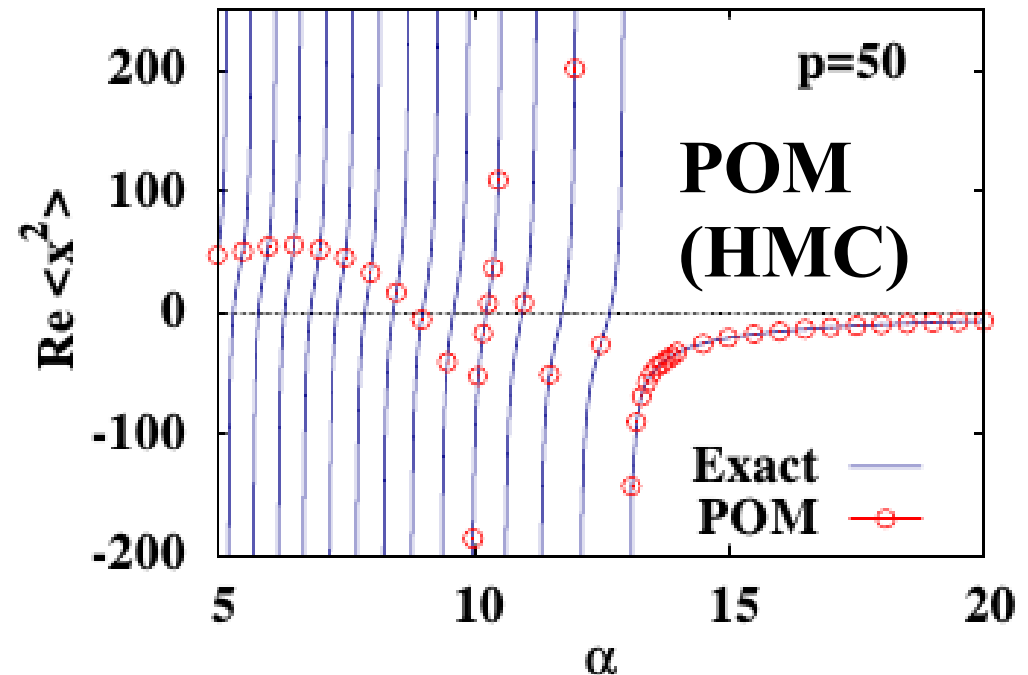
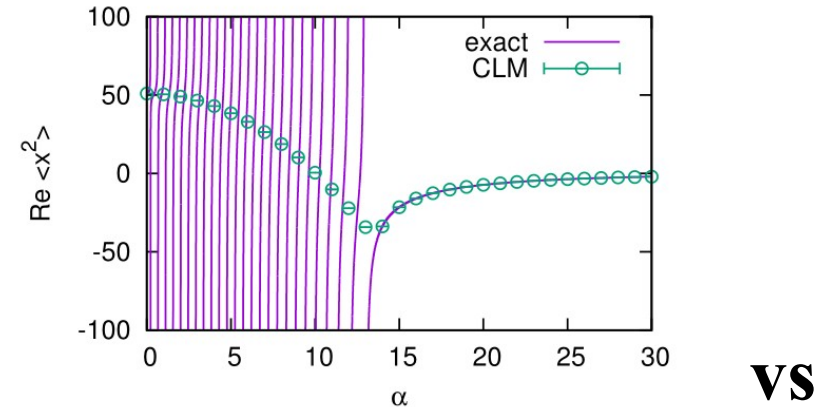


On Optimized Path



Observable

CLM *Nishimura, Shimasaki ('15)*



Mori, Kashiwa, AO ('17); AO, Mori, Kashiwa (Lat 2017)

Ohnishi @ Lattice 2018, July 28, 2018

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*Now it's the time to apply POM
to field theories !*

Lattice 2017 (Granada) → Lattice 2018 (MSU)

- **Introduction to Path Optimization Method**

Y. Mori, K. Kashiwa, AO, PRD 96 ('17), 111501(R) [arXiv:1705.05605]

*AO, Y. Mori, K. Kashiwa, EPJ Web Conf. 175 ('18), 07043 [arXiv:1712.01088]
(Lattice 2017 proceedings)*

- **Application to complex ϕ^4 theory using neural network**

Y. Mori, K. Kashiwa, AO, PTEP 2018 ('18), 023B04 [arXiv:1709.03208]

- **Application to gauge theory: 1-dimensional QCD**

Y. Mori, K Kashiwa, AO, in prep.

- **Discussions**

- **Summary**

*Application to complex ϕ^4 theory
using neural network*

Y. Mori, K. Kashiwa, AO, PTEP 2018 ('18), 023B04 [arXiv:1709.03208]

Application of POM to Field Theory

- Preparation & variation of trial fn. is tedious in multi-D systems

$$z_i(x) = x_i + i \sum_{n_1, n_2, \dots} c_i(n_1 n_2 \dots) H_{n_1}(x_1) H_{n_2}(x_2) H_{n_3}(x_3) \dots$$

- Neural network

- Combination of linear and non-linear transformation.

$$a_i = g(\underline{W}_{ij}^{(1)} x_j + \underline{b}_i^{(1)}) \quad \text{parameters}$$

$$f_i = g(\underline{W}_{ij}^{(2)} a_j + \underline{b}_i^{(2)})$$

$$z_i = x_i + i(\underline{\alpha}_i f_i(x) + \underline{\beta}_i)$$

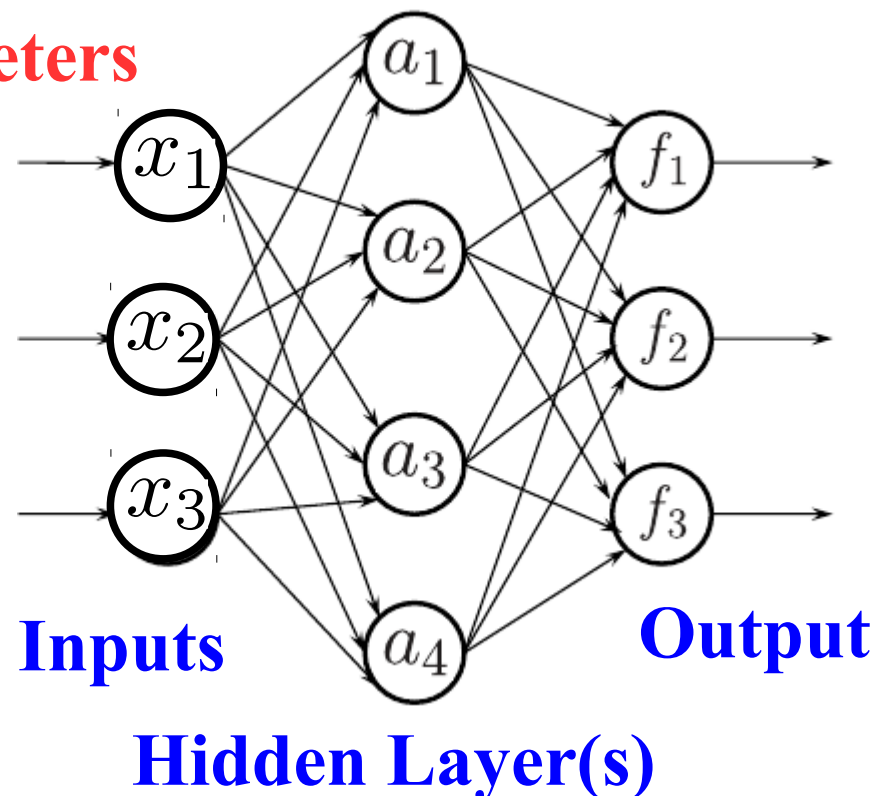
$$g(x) = \tanh x \quad (\text{activation fn.})$$

- Universal approximation theorem

Any fn. can be reproduced
at (hidden layer unit #) $\rightarrow \infty$

G. Cybenko, MCSS 2 ('89) 303

K. Hornik, Neural networks 4('91) 251



Optimization of many parameters

- Stochastic Gradient Descent method, E.g. ADADELTA algorithm
M. D. Zeiler, arXiv:1212.5701

Grad. Desc. :
 $dc_i/dt = -\partial\mathcal{F}/\partial c_i$

par. in (j+1)th step

Learning rate

$$c_i^{(j+1)} = c_i^{(j)} - \eta v_i^{(j+1)}$$

mean sq. ave. of v

$$v_i^{(j+1)} = \frac{\sqrt{s_i^{(j)} + \epsilon}}{\sqrt{r_i^{(j+1)} + \epsilon}} F_i^{(j)}$$

decay rate

mean sq. ave. of F

$$r_i^{(j+1)} = \gamma r_i^{(j)} + (1 - \gamma)(F_i^{(j)})^2$$

$$s_i^{(j+1)} = \gamma s_i^{(j)} + (1 - \gamma)(v_i^{(j+1)})^2$$

gradient
evaluated
in MC

(batch training)

$$F_i = \partial\mathcal{F}/\partial c_i$$

Cost fn.

*Machine learning
 ~ Educated algorithm
 to generic problems*

Hybrid Monte-Carlo with Neural Network

Initial Config. on Real Axis

$$\text{HMC } H(x, p) = \frac{p^2}{2} + \text{Re}S(z(x))$$

Jacobian

→ via Metropolis judge

Do k = 1, Nepoch

Do j = 1, Nconf/Nbatch

Mini-batch training of Neural Network

$$c_i^{(j+1)} = c_i^{(j)} - \eta v_i^{(j+1)}$$

Grad. wrt parameters (Nbatch configs.)

$$F_i = \frac{1}{N_{\text{batch}}} \sum_n \partial \mathcal{F}(n) / \partial c_i$$

New Nbatch configs. by HMC

$$H(x, p) = \frac{p^2}{2} + \text{Re}S(z(x))$$

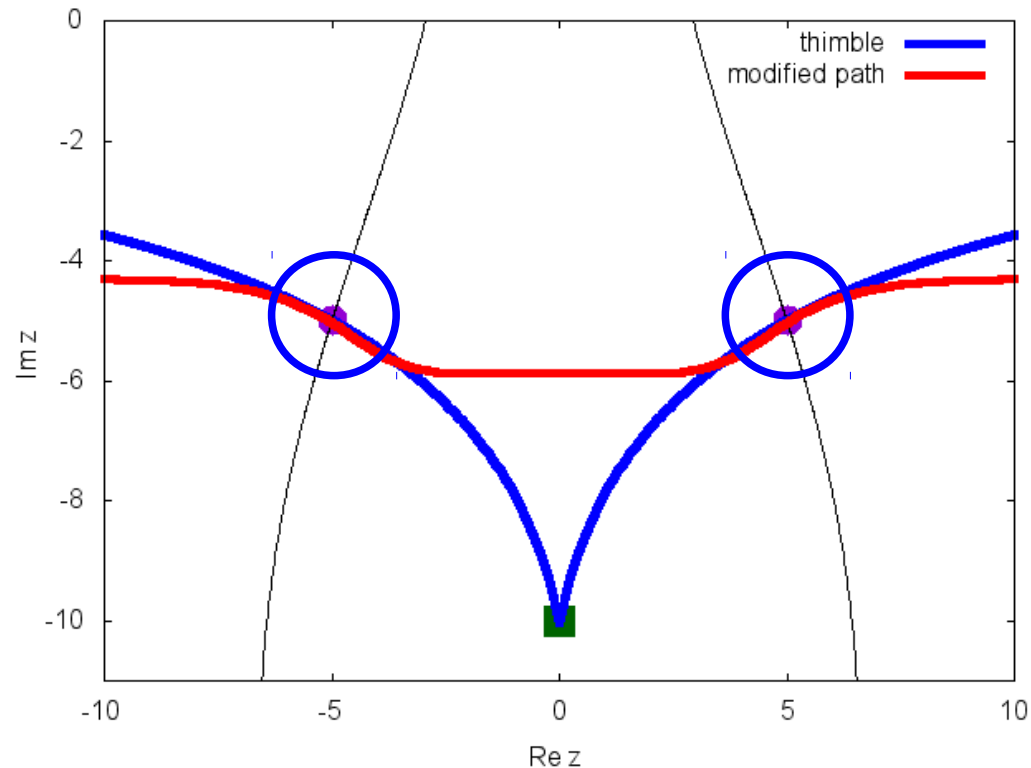
Enddo

Enddo

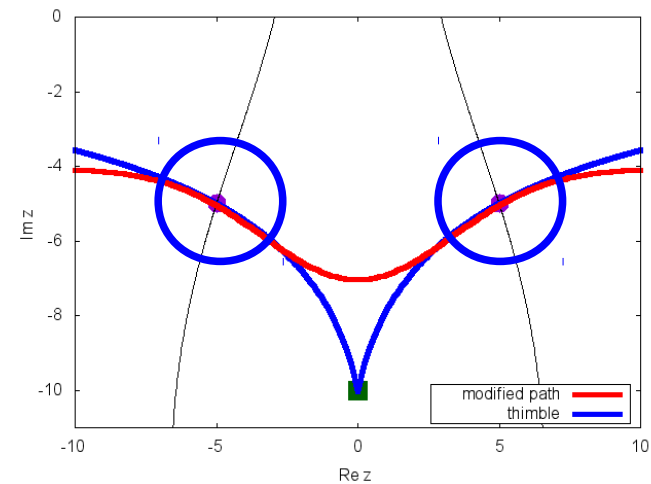
Nbatch ~ 10, Nconfig ~ 10,000, Nepoch ~ (10-20)

Optimized Path by Neural Network

Neural Network



Gaussian +Gradient Descent



*Optimized paths are different,
but both reproduce thimbles around the fixed points !*

AO, Mori, Kashiwa (Lat 2017)

Ohnishi @ Lattice 2018, July 28, 2018

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Complex ϕ^4 theory at finite μ

Complex ϕ^4 theory

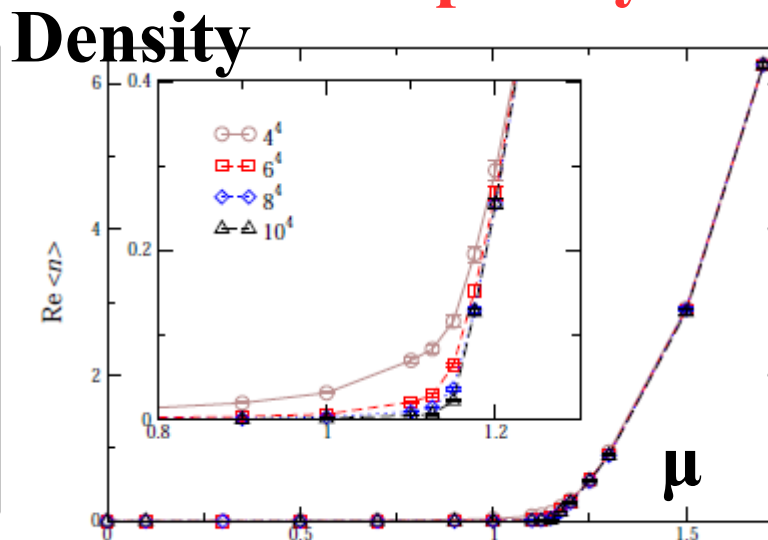
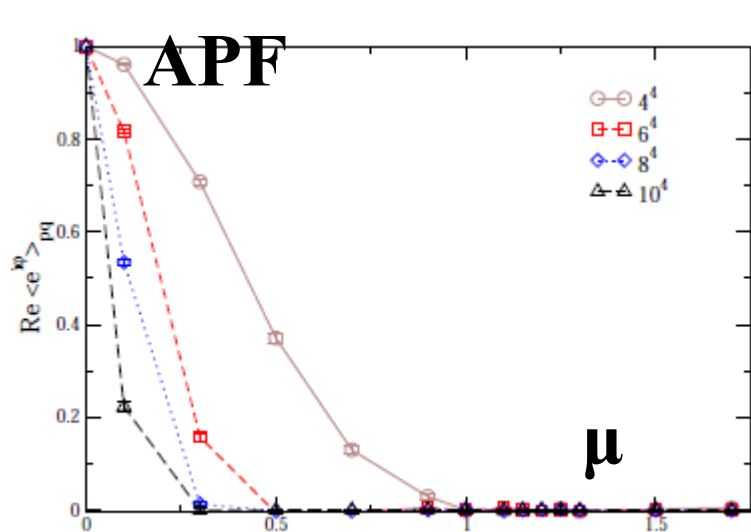
$$\mathcal{L} = \partial_\mu \phi^* \partial^\mu \phi - m^2 \phi^* \phi - \lambda (\phi^* \phi)^2$$

Action on Euclidean lattice at finite μ .

$$S = \sum_x \left[\frac{(4 + m^2)}{2} \phi_{a,x} \phi_{a,x} + \frac{\lambda}{4} (\phi_{a,x} \phi_{a,x})^2 - \phi_{a,x} \phi_{a,x+\hat{1}} - \cosh \mu \phi_{a,x} \phi_{a,x+\hat{0}} + i \epsilon_{ab} \sinh \mu \phi_{a,x} \phi_{b,x+\hat{0}} \right] \left(\phi = \frac{1}{\sqrt{2}} (\phi_1 + i \phi_2) \right)$$

complex

Complexify



Complex
Langevin
& Lefschetz
thimble
work.

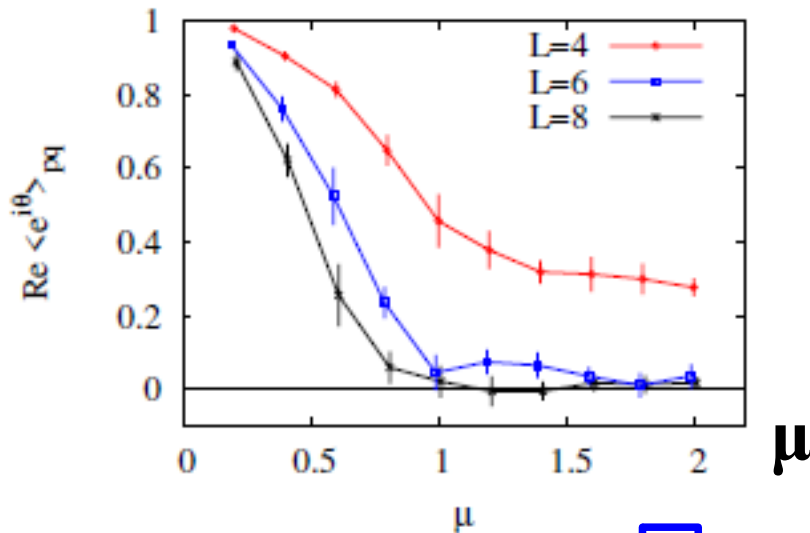
μ G. Aarts, PRL102('09)131601; H. Fujii, et al., JHEP 1310 (2013) 147

POM result (1): Average phase factor

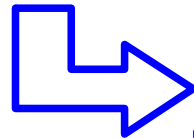
■ POM for 1+1D ϕ^4 theory

- $4^2, 6^2, 8^2$ lattices, $\lambda=m=1$
- $\mu_c \sim 0.96$ in the mean field approximation
- Enhancement of the average phase factor after optimization.

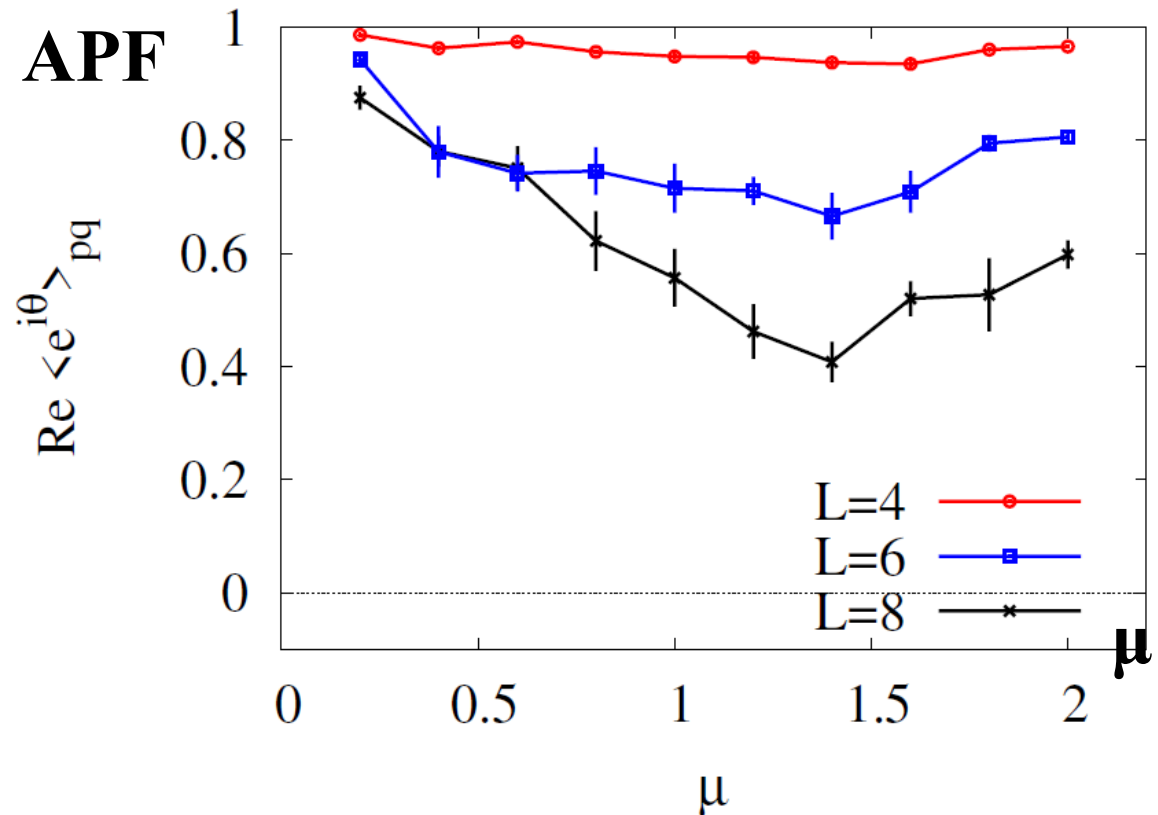
APF



Optimization



APF



Y. Mori, K. Kashiwa, AO, PTEP 2018 ('18), 023B04 [arXiv:1709.03208]

Ohnishi @ Lattice 2018, July 28, 2018

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POM result (2): Density

- Results on the real axis
Small average phase factor, Large errors of density
- On the optimized path
Finite average phase factor, Small errors

Mean Field App.

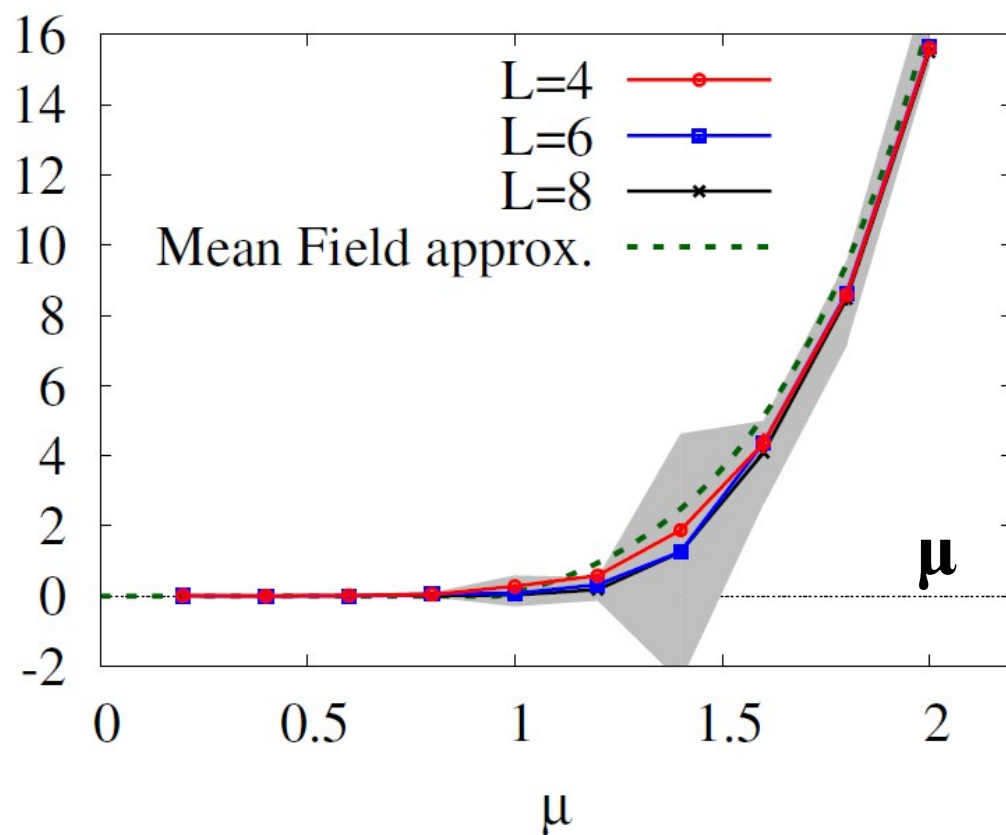
$$\frac{S}{V} = \left(1 + \frac{m^2}{2} - \cosh \mu \right) \phi^2 + \frac{\lambda}{4} \phi^4 ,$$

$$n = \phi^2 \sinh \mu ,$$

$$\phi_{\text{stat.}}^2 = \begin{cases} 0 & (|\mu| < \mu_c) , \\ \frac{2}{\lambda} (\cosh \mu - 1 - \frac{m^2}{2}) & (|\mu| \geq \mu_c) , \end{cases}$$

Density

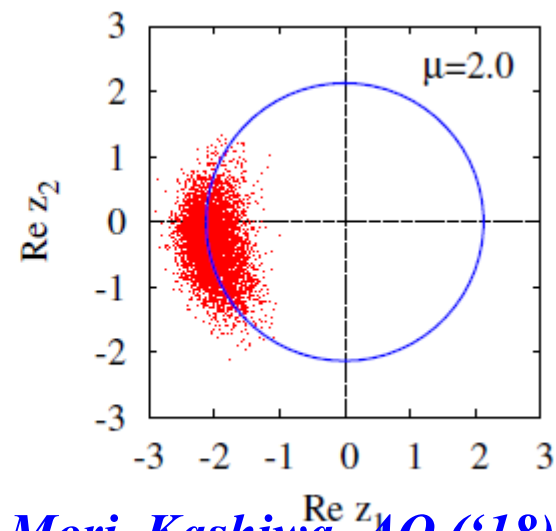
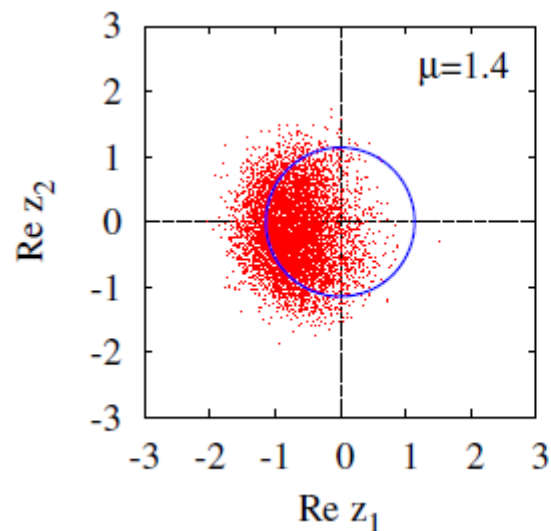
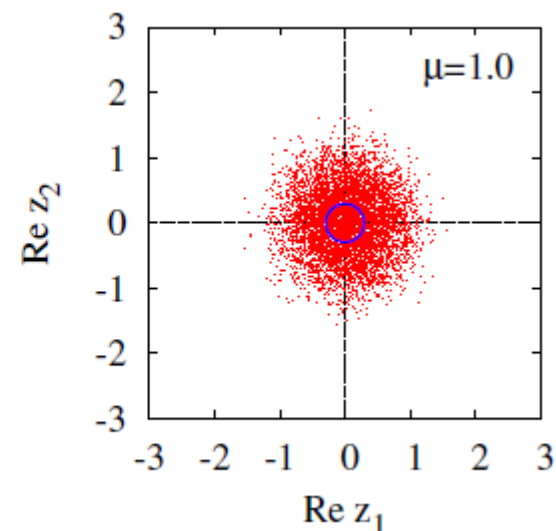
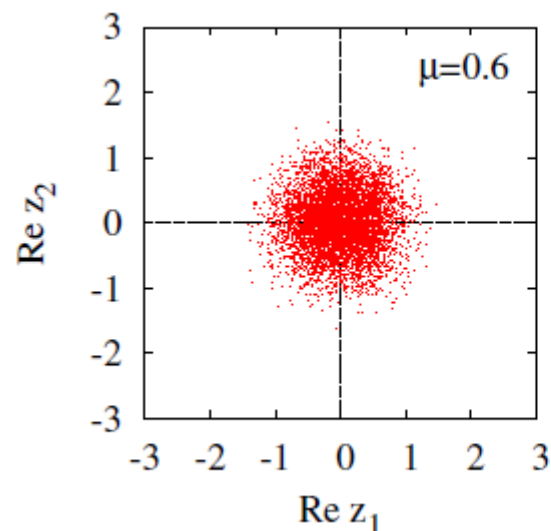
Re <n>



Mori, Kashiwa, AO ('18)

POM result (3): Configurations

- Updated configurations after optimization
→ sampled around the mean field results
- Global U(1) symmetry in (φ_1, φ_2) is broken(*)
by the optimization
or by the sampling.



* This does not contradict the Elitzur's theorem.

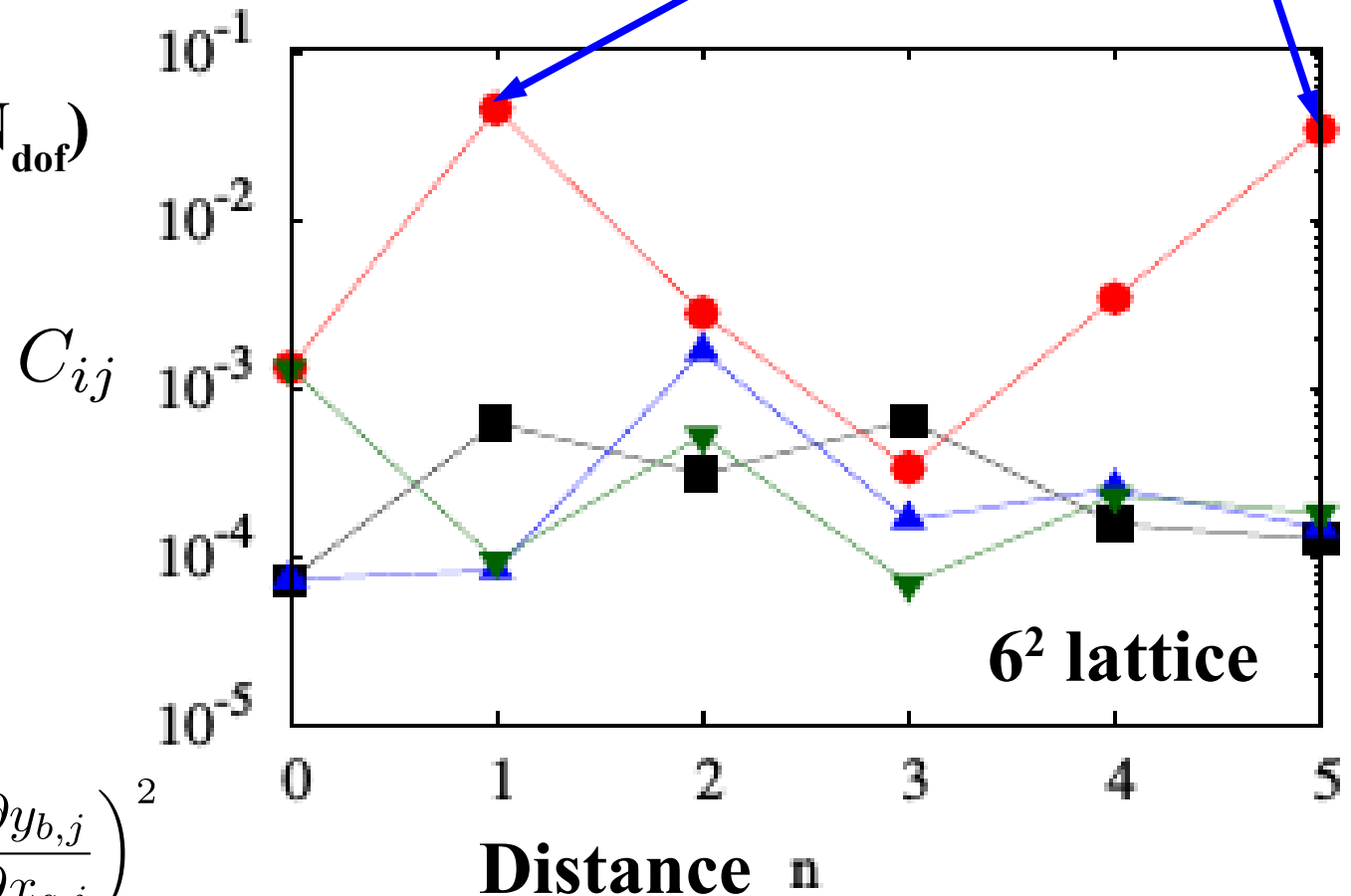
Mori, Kashiwa, AO ('18)

Which y 's should be optimized ?

- Correlation btw (z_1, z_2) of temporal nearest neighbor sites are strong. Other correlations $\sim 10^{-2}$ times smaller

$$\text{Im}(S) = \sum_x \epsilon_{ab} \sinh \mu \phi_{a,x} \phi_{b,x+\hat{0}}$$

- Hope to reduce the cost to be $O(N_{\text{dof}})$



$$C_{ij} \equiv \left(\frac{\partial y_{a,i}}{\partial x_{b,j^*}} \right)^2 + \left(\frac{\partial y_{b,j}}{\partial x_{a,i}} \right)^2$$

Y. Mori, Master thesis

*Application to Gauge Theory:
1 dimensional QCD*

0+1 dimensional QCD

0+1 dimensional QCD (1 dim. QCD)

with one species of staggered fermion on a 1xN lattice

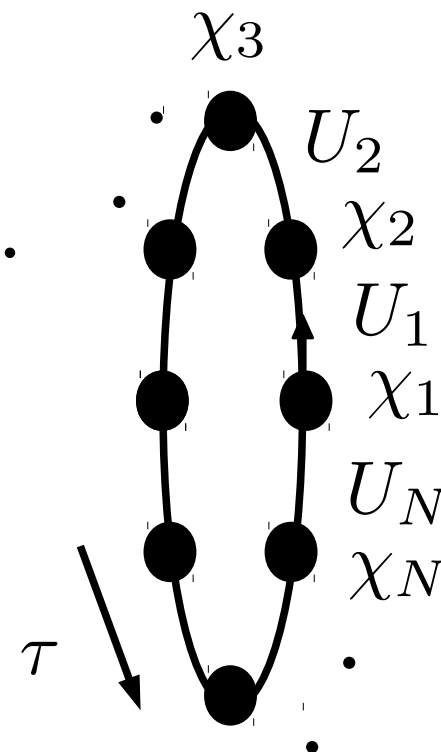
$$S = \frac{1}{2} \sum_{\tau} (\bar{\chi}_{\tau} e^{\mu} U_{\tau} \chi_{\tau+\hat{0}} - \bar{\chi}_{\tau+\hat{0}} e^{-\mu} U_{\tau}^{-1} \chi_{\tau}) + m \sum_{\tau} \bar{\chi}_{\tau} \chi_{\tau} = \frac{1}{2} \bar{\chi} D \chi$$

$$\mathcal{Z} = \int \mathcal{D}U \det D[U] = \int dU \det \left[X_N + (-1)^N e^{\mu/T} U + e^{-\mu/T} U^{-1} \right]$$

$$X_N = 2 \cosh(E/T), \quad E = \operatorname{arcsinh} m, \quad U = U_1 U_2 \cdots U_N, \quad T = 1/N$$

*Bilic('88), Ravagli('07), Aarts('10, CLM), Bloch('13, subset),
Schmidt('16, LTM), Di Renzo('17, LTM)*

- A toy model, but the actual source of QCD sign prob.
- Studied well in the context of strong coupling LQCD
*E.g. Miura, Nakano, AO, Kawamoto('09, '09, '17),
de Forcrand, Langelage, Philipsen, Unger ('14)*



1 dim. QCD in diagonal gauge

■ Diagonal gauge

$$U = (e^{iz_1}, e^{iz_2}, e^{iz_3}) \quad (z_1 + z_2 + z_3 = 0)$$

$$\mathcal{Z} = \int dU e^{-S} = \int dx_1 dx_2 J H e^{-S}$$

$$= \int dx_1 dx_2 \det \left(\frac{\partial z_a}{\partial x_b} \right) \left[\frac{8}{3\pi^2} \prod_{a < b} \sin^2 \left(\frac{z_a - z_b}{2} \right) \right] \left[\prod_a (X_N + 2 \cos(z_a - i\mu)) \right]$$

Jacobian

Haar measure

exp(-S)

■ Path optimization (t: fictitious time)

→ $y(x_1, x_2)$ itself is the parameter on the (x_1, x_2) mesh point

$$z_i = x_i + iy_1, \quad y_i = y_i(x_1, x_2)$$

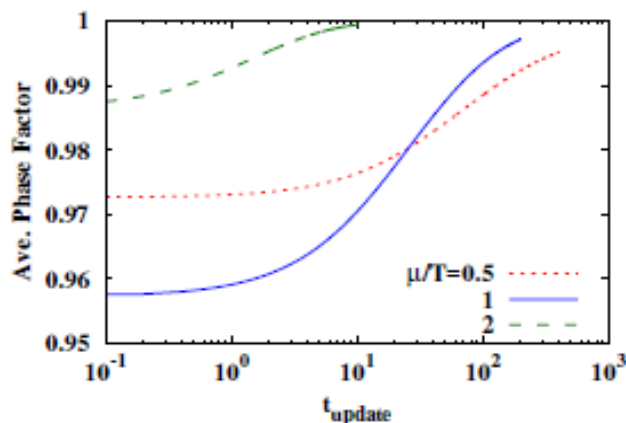
$$\frac{dy_i}{dt} = -\frac{\partial \mathcal{Z}_{pq}}{\partial y_i}, \quad \mathcal{Z}_{pq} = \int dx_1 dx_2 |J H e^{-S}|$$

Path Opt. of 1 dim. QCD in diagonal temporal gauge

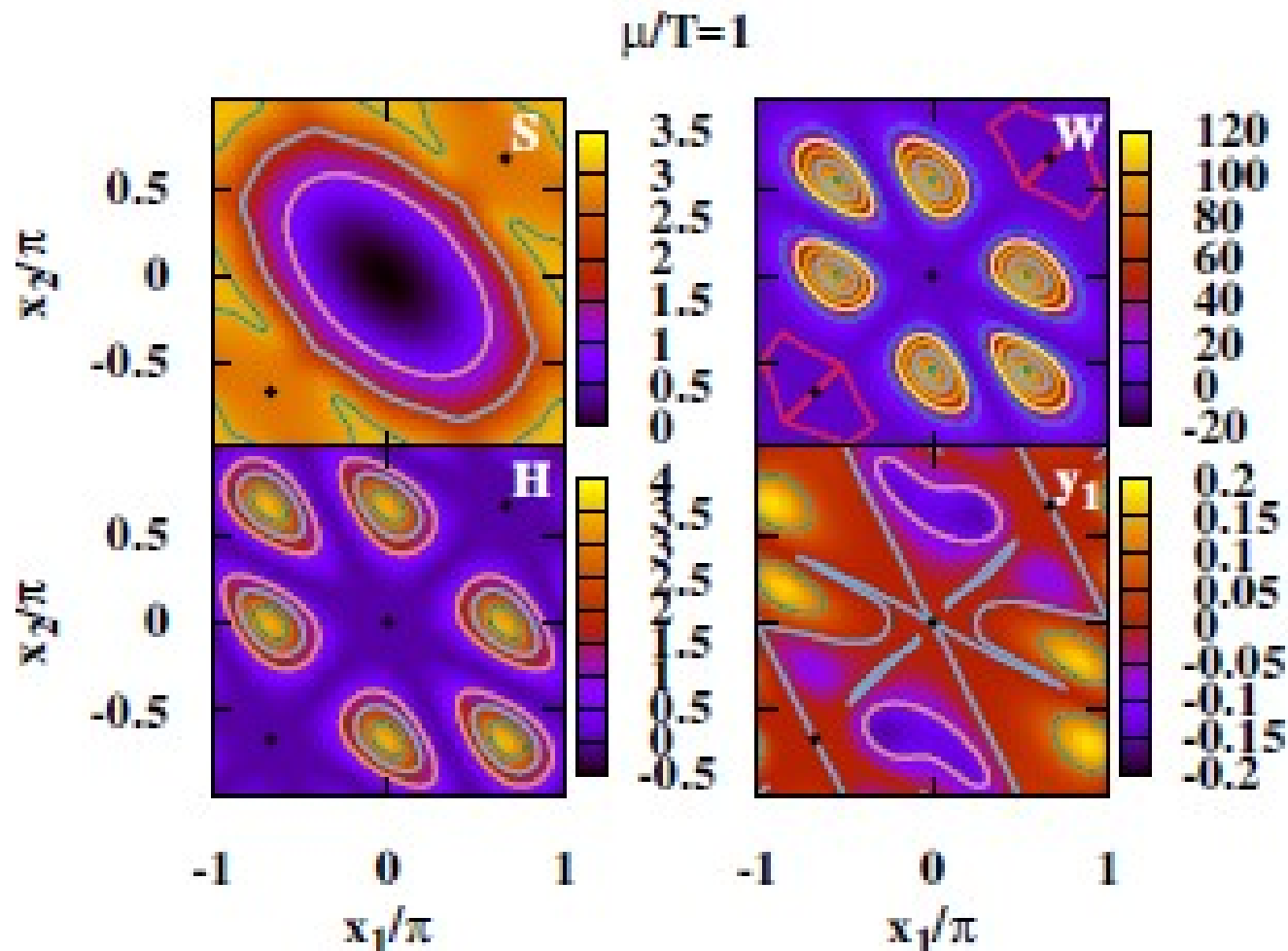
■ Path optimization

- Average phase factor > 0.99 \rightarrow Easily achieved
- $\exp(-S)$ and Haar Measure \rightarrow “six pads” *Schmidt+('16, LTM)*

APF



fictitious time



Mori, Kashiwa, AO, in prep.

1 dim. QCD with Hybrid MC

■ Concern...

- Six pads are separated by the Haar measure barrier.

$$\text{Symmetry : } S(-z) = (S(z^*))^*, z_i \leftrightarrow z_j (i, j = 1, 2, 3)$$

Do we need exchange MC or different tempering ?

E.g. Fukuma, Matsumoto, Umeda ('17)

■ Hybrid Monte-Carlo in 1 dim. QCD

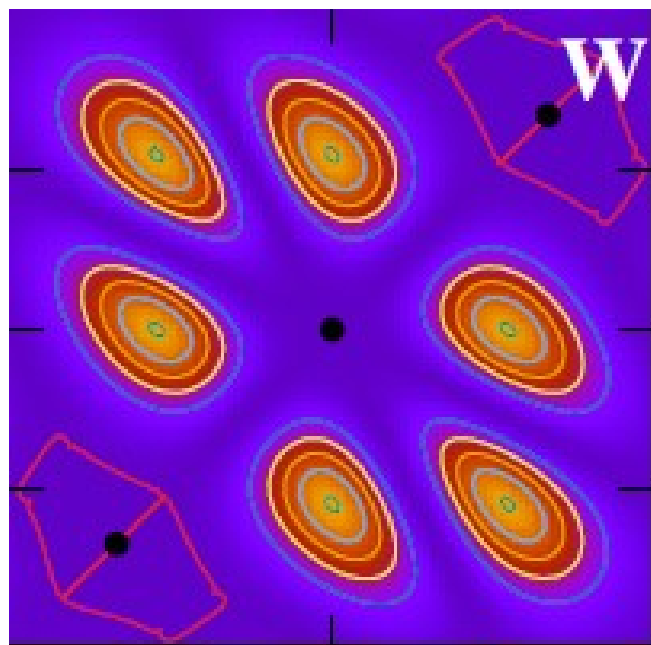
$$U \rightarrow \mathcal{U}(U) = U \prod_{a=1}^{N_c^2-1} e^{-y_i \lambda_i / 2}, \quad H = \frac{P^2}{2} + \text{Re}(S(\mathcal{U}(U)))$$

SL(3)

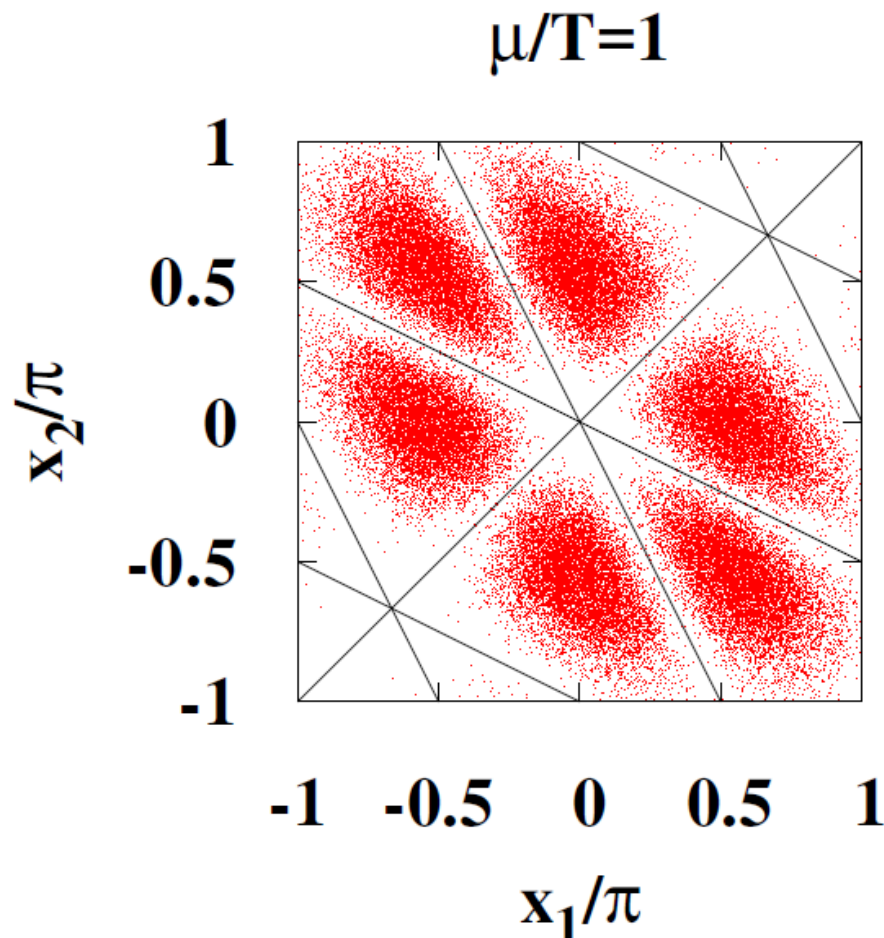
- 8 variables → path optimization using Neural Network

1 dim. QCD with Hybrid MC

- HMC + diagonalization of the link
→ All six pads are visited, and no Ex. MC needed.



Mesh point + Grad. Desc.



HMC + NN

Mori, Kashiwa, AO, in prep.

Discussions

Frequently Asked Questions

- How many parameters do you have ?

→ Many ;) For generic trial function ($V = \#$ of variables)

$$y_i = y_i(x_1, x_2, \dots, x_V)$$

$$N_{\text{par}} = (N_{\text{layer}} + 1) \times V \times (N_{\text{unit}} + 1) + 2V$$

- How about the numerical cost ?

→ A lot ;) Derivative of J with respect to parameters cost most.

$$\frac{\partial J}{\partial c_i} = J \frac{\partial J_{jk}^{-1}}{\partial z_l} \frac{\partial z_l}{\partial c_i} \rightarrow \mathcal{O}(V^3)$$

- It is still polynomial.

Does the sign problem becomes “P” problem?

→ No. The average phase factor is still $\exp(-\# V)$.

If extrapolation is possible from finite V , we have a hope.

- How can we reduce the cost ? → Next page

How can we reduce the numerical cost ?

- Restrict the function form of $y(x)$.

- Imaginary part is a function of its real part.

E.g. Alexandru, Bedaque, Lamm, Lawrence, PRD97('18)094510

[Lawrence, Warrington, Lamm (Mon)]

Thirring model, 1+1D QED

$$y_i = f(x_i), f(x) = \lambda_0 + \lambda_1 \cos x$$

- Nearest neighbor site

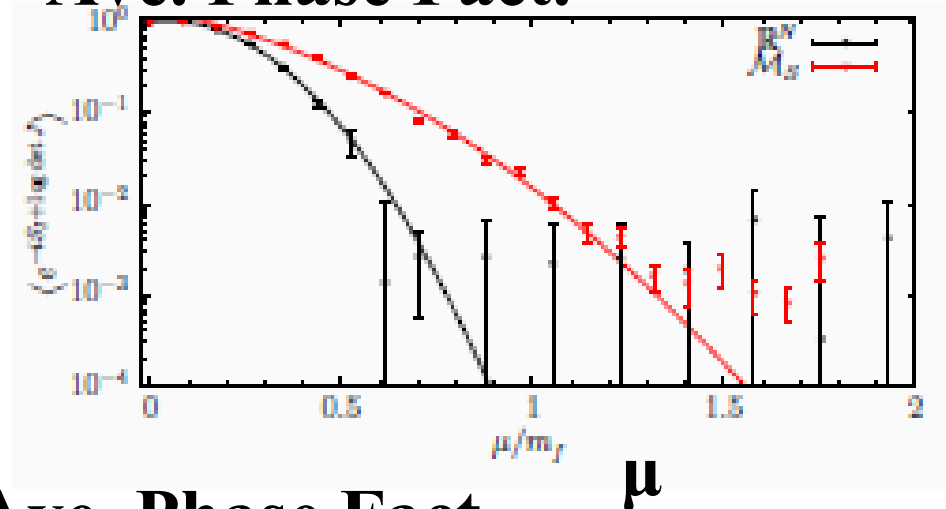
F. Bursa, M. Kroyter, arXiv:1805.04941

0+1 D ϕ^4 theory

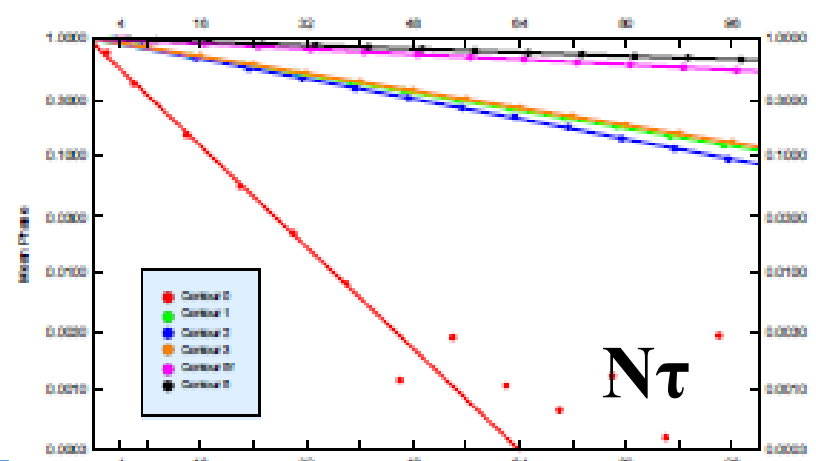
Translational inv. + U(1) sym.

$$y_{a,i} = \frac{\epsilon_{ab} x_{a,i+1}}{1 + x_{1,i}^2 + x_{2,i}^2}$$

Ave. Phase Fact.

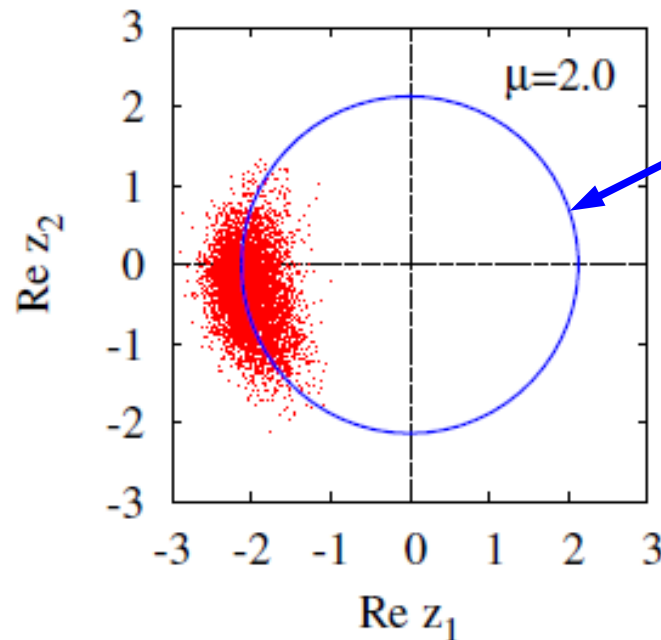


Ave. Phase Fact.



Frequently Asked Questions (cont.)

- What happens when we have 10^{10} fixed points ?
 - In that case we should give up. (My answer @ Lattice 2017)
 - If those fixed points are connected by the symmetry, we may be able to perform path optimization.



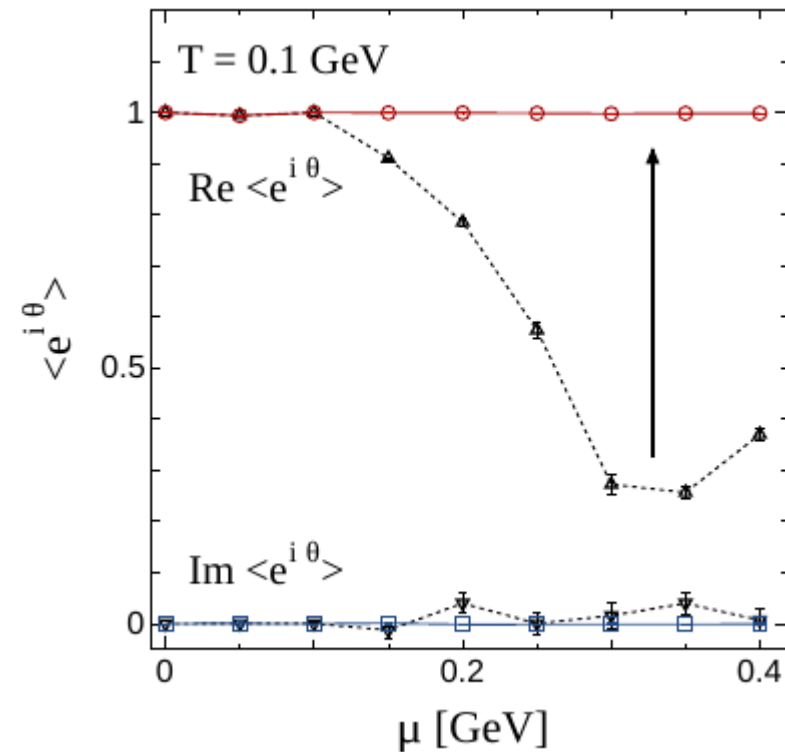
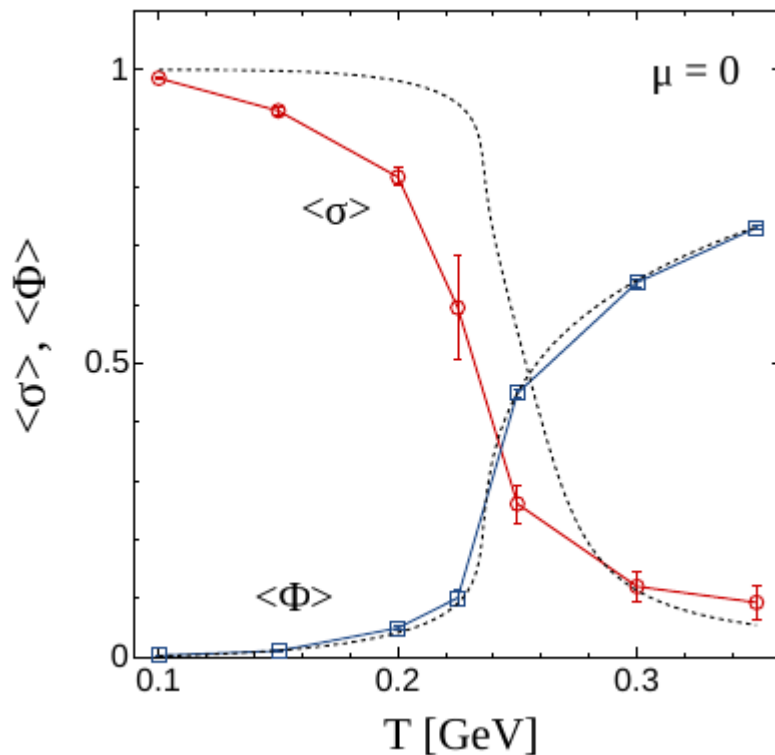
Mean field results
= Degenerate fixed points
(All have the same θ .)

If they have different complex phases, the global sign problem emerges and the partition function would be almost zero.

E.g. H. Fujii, S. Kamata, Y. Kikukawa, arXiv:1710.08524

Application to PNJL

- PNJL model with homogeneous condensates, $(\sigma, \pi, \Phi, \bar{\Phi})$.
 - Has Sign problem in finite volume
 - Converges to mean field results in the large volume limit



K. Kashiwa, Y. Mori, AO, arXiv:1805.08940.

Summary

- The sign problem is a grand challenge in theoretical physics, and appears in many fields of physics,
 - finite density QCD, real time evolution, Hubbard model off half-filling, other quantum MC with fermions, ...and complexified variable methods (LTM, CLM, POM) would be promising to evade the sign problem.
- **Path optimization** with the use of the **neural network** is demonstrated to work in **field theories** having many variables.
 - 1+1D ϕ^4 theory at finite μ (neural network)
 - 0+1D QCD w/ fermions (grad. descent, neural network)
 - 3+1D homogeneous PNJL (neural network)
- Neural network (single hidden layer) is the simplest device of machine learning, and it helps us to **generate and optimize generic multi-variable functions**, $y_i = y_i(\{\mathbf{x}\})$.

Prospect

- **Path optimization in 3+1 D field theories would require reduction of numerical cost.**
Imaginary part
= f (real parts of same point and nearest neighbor points)
may be a good guess.
- **Deep learning (# of hidden layers > 3) may be helpful to explore complex path, which human beings (~ 7 layers) cannot imagine, while “Understanding” the results of machine learning need to be done by human beings (at present).**

Thank you for your attention !